

Perspectives on Nuclear Physics Input into High-Energy Cosmic Ray Interactions

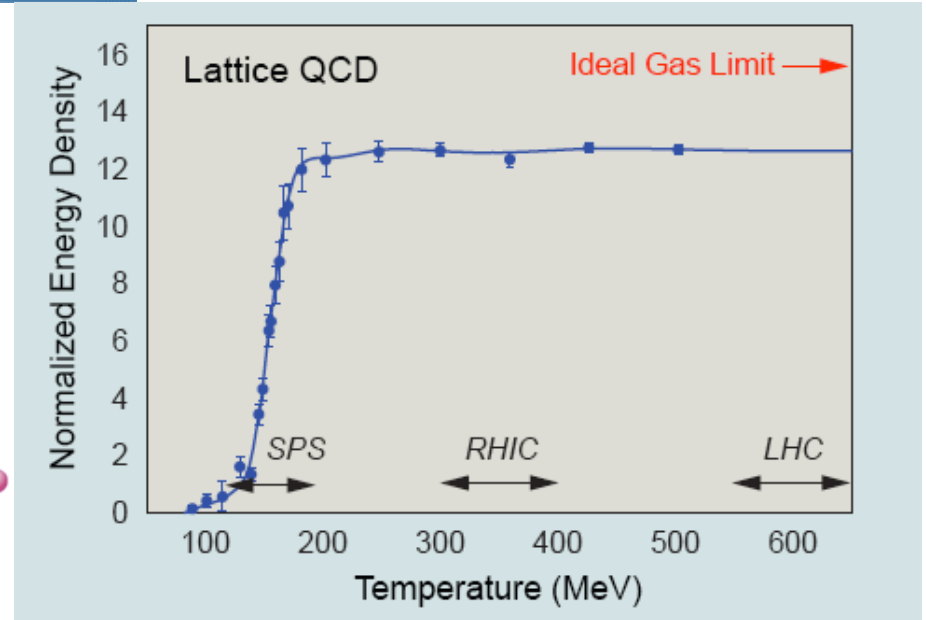
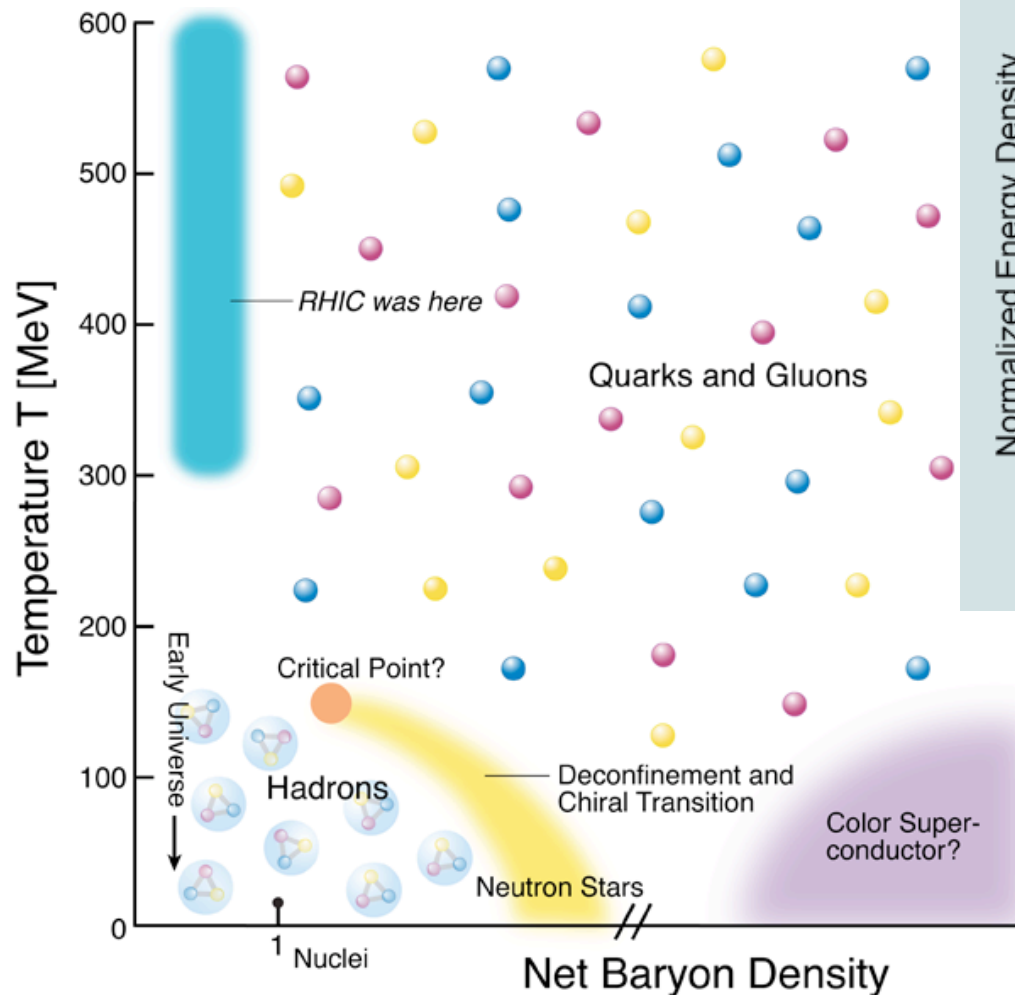
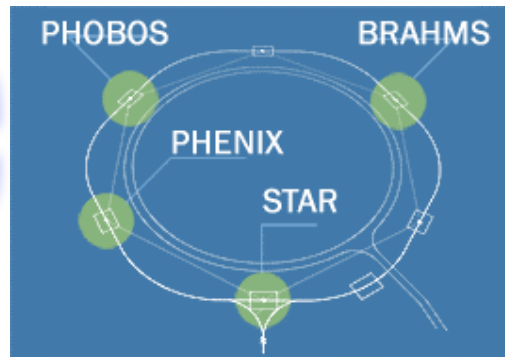
A.B. Balantekin

University of Wisconsin-Madison

XVI International Symposium on Very High Energy
Cosmic Ray Interactions, Fermilab, June 2010

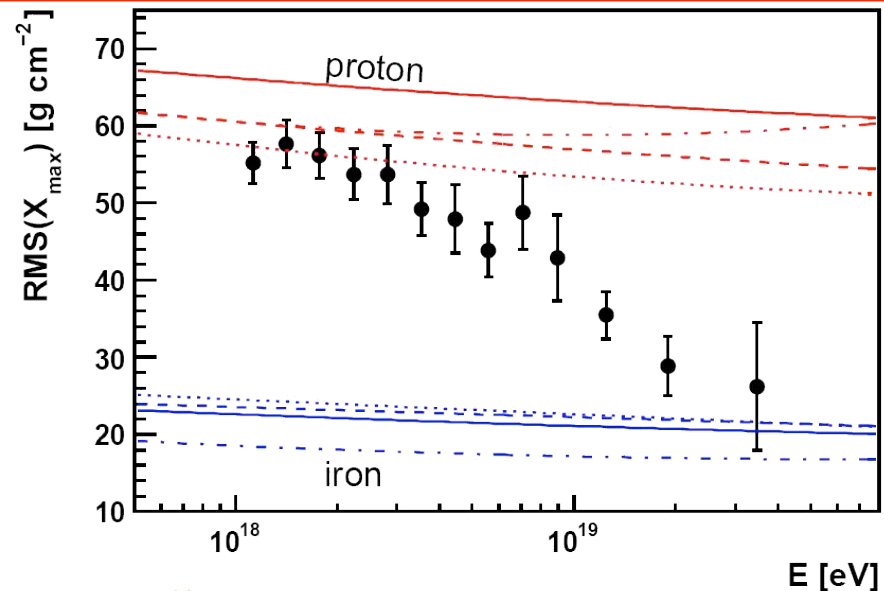
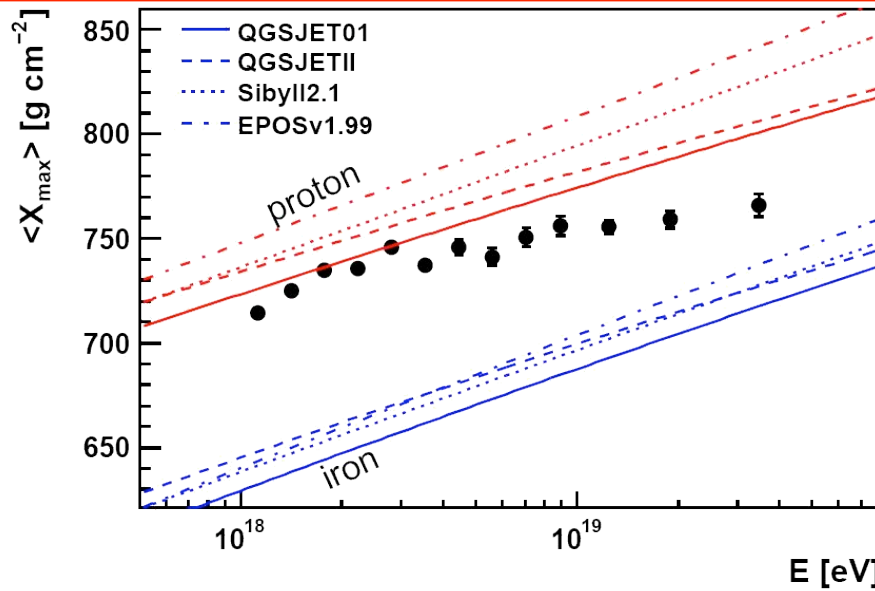
Disclaimer: My expertise in nuclear collisions is mostly at low energies; my expertise in high energies is mostly with neutrinos. So this is essentially an outsider's perspective!

RHIC

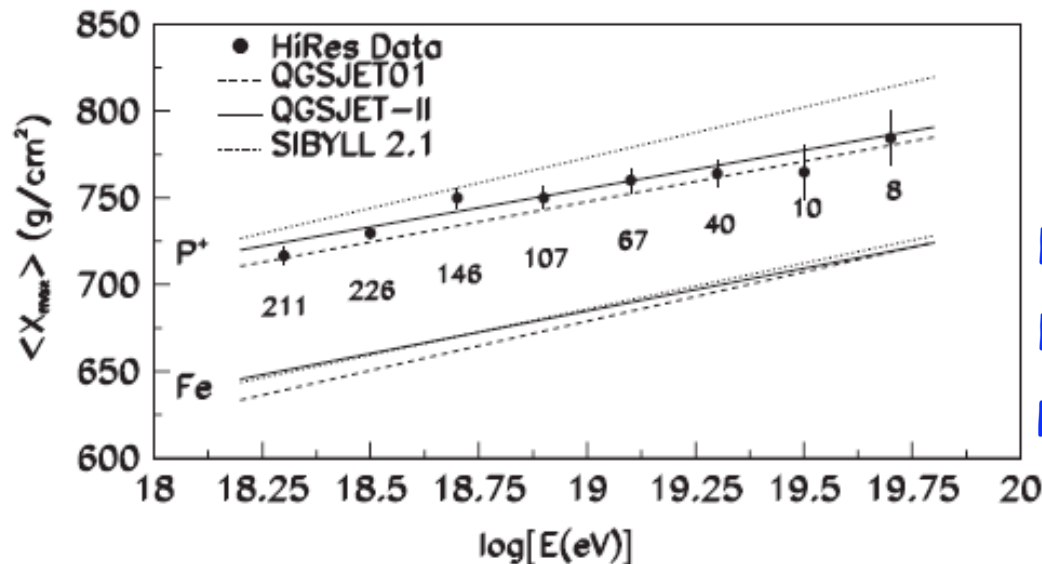


Why are laboratory nuclear experiments relevant to the cosmic-ray physics?

Recent results suggest presence of a significant nuclear component in the higher-energy cosmic-ray flux from the measurements of the depth of the shower maximum

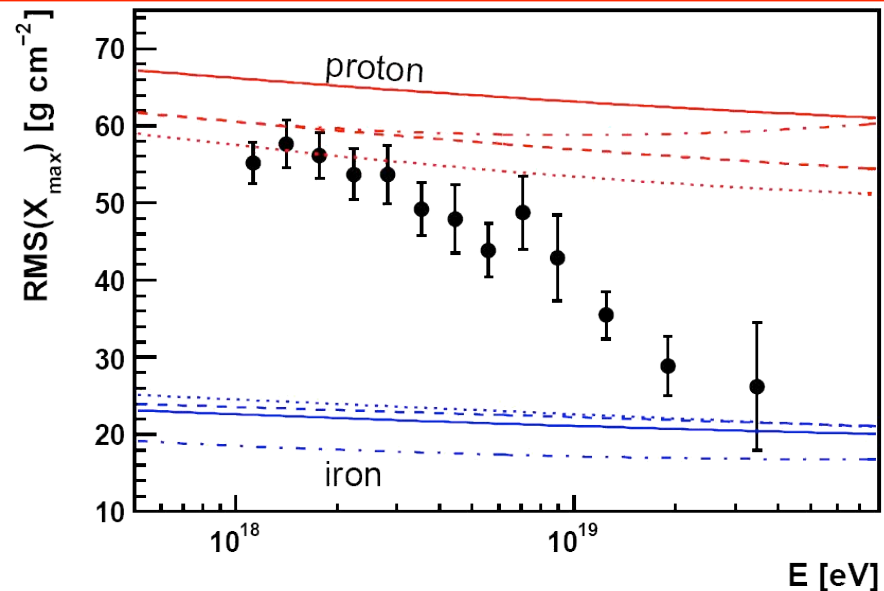
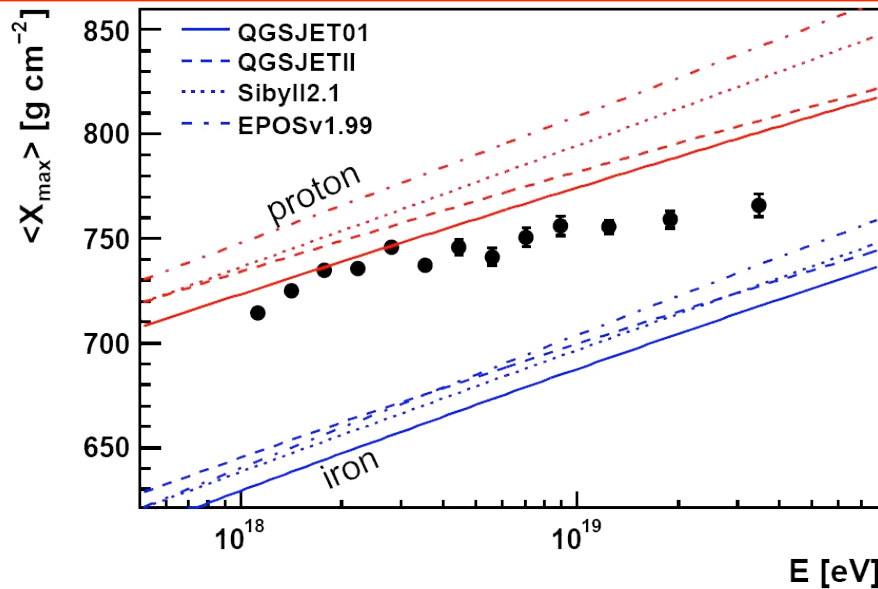


AUGER Collaboration, PRL 104, 091101 (2010)



However, see
HIRES Collaboration,
PRL 104, 161101 (2010)

Recent results suggest presence of a significant nuclear component in the higher-energy cosmic-ray flux from the measurements of the depth of the shower maximum



AUGER Collaboration, PRL **104**, 091101 (2010)

CAUTION: Heitler's original formula:

$$\langle X_{\max} \rangle = \alpha (\ln E - \langle \ln A \rangle) + \beta$$

assumes that heavier nuclei are basically superposition of the nucleons

(see however Ulrich et al., arXiv:0906.0418)

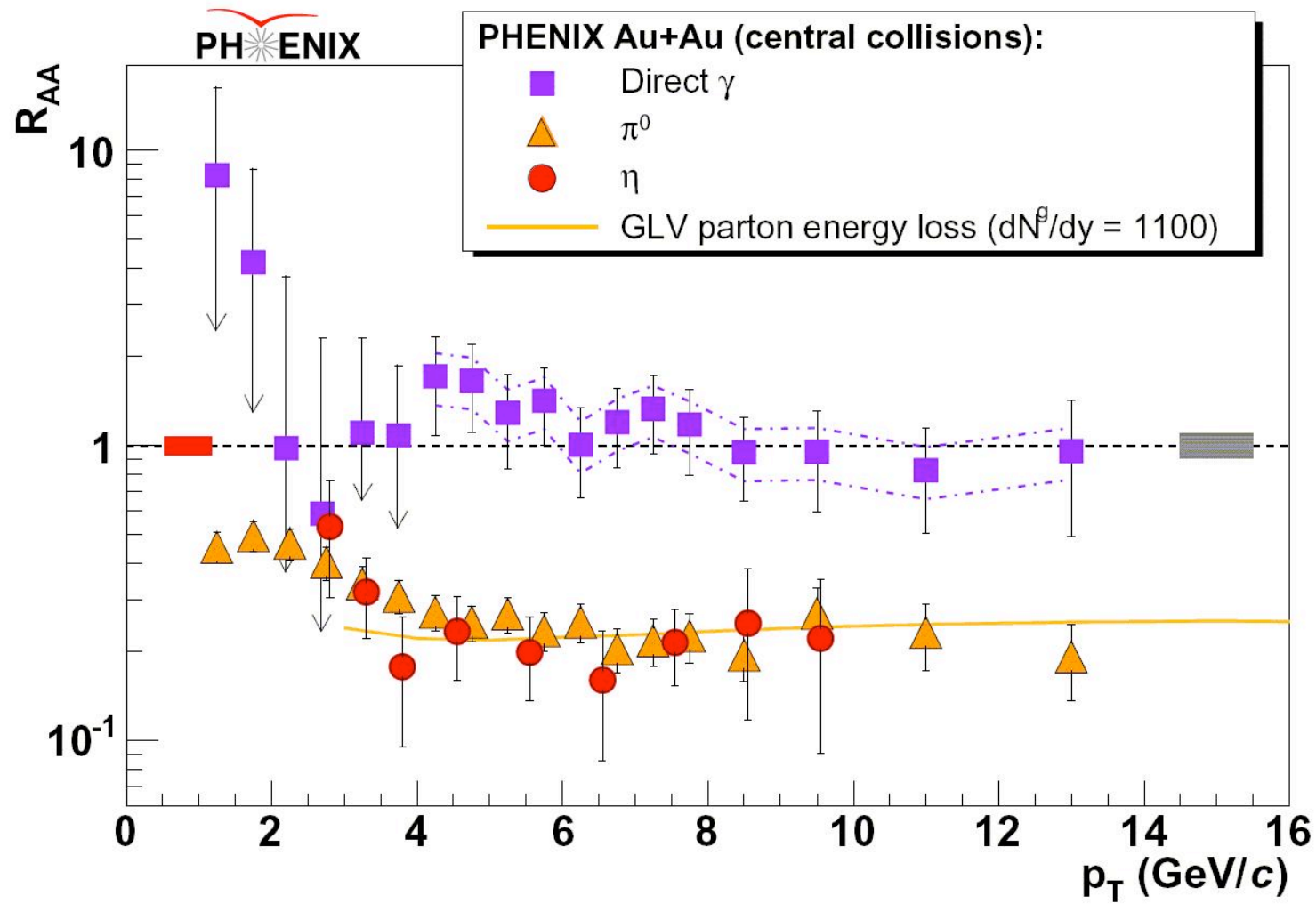
Are nuclear collisions simply a superposition of nucleon collisions?

$$R_{AA} = \frac{dN_{AA}/dp_T}{N_{\text{coll}} (dN_{pp}/dp_T)}$$

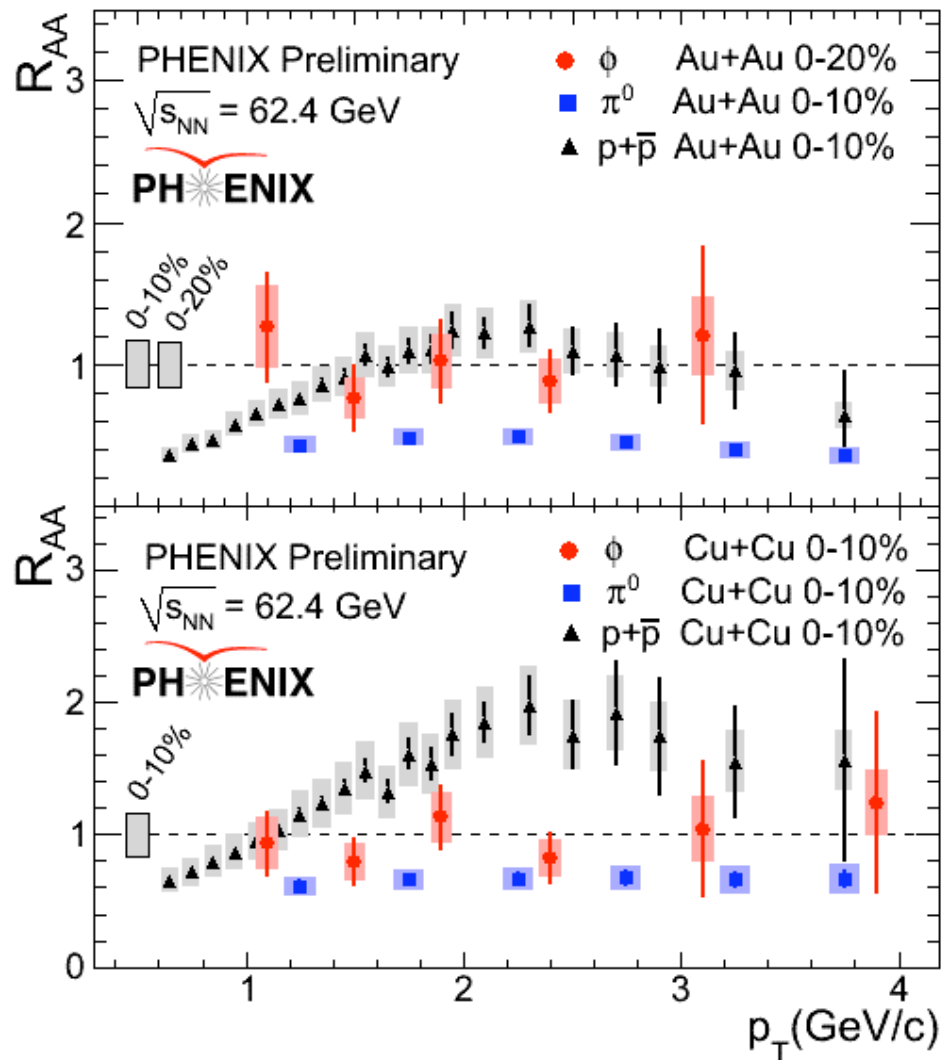
N_{coll} = number of binary collisions

Absence of nuclear medium interactions (i.e. γ 's) $\Rightarrow R_{AA} \approx 1$

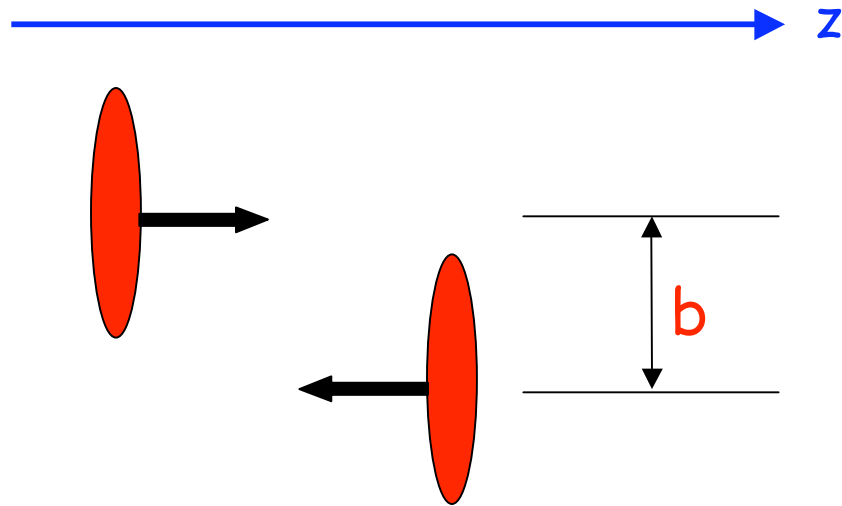
Energy loss in the medium \Rightarrow reduction of p_T



PRL 96, 202301 (2006)



QCD jets are
quenched by the
nuclear medium.
Nuclear collisions
are NOT simply a
superposition of
pp collisions!

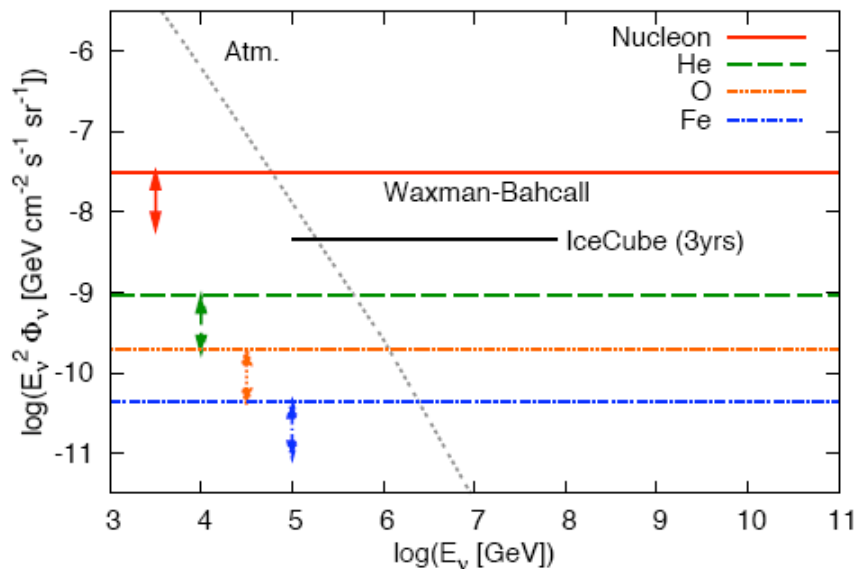
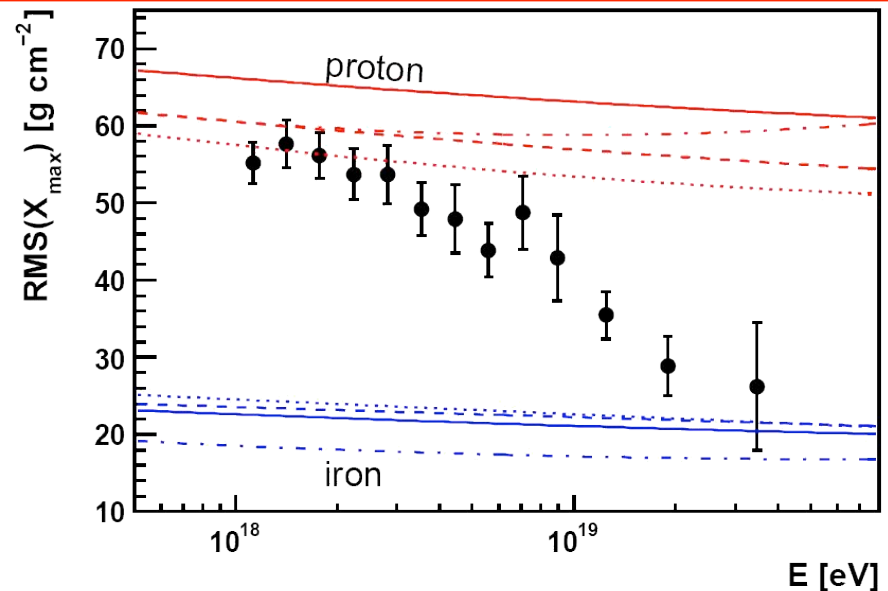
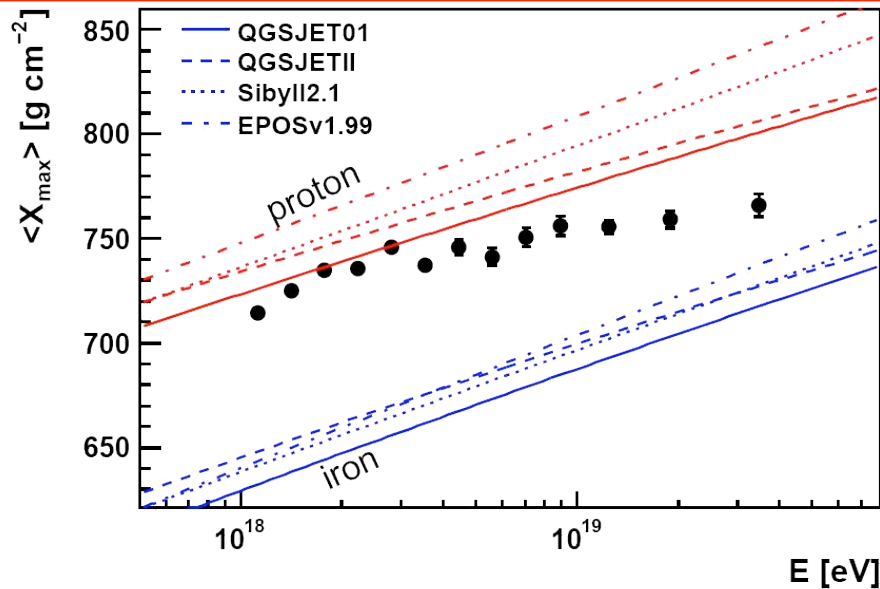


Glauber formula

$$\sigma_{AA} = \int d^2b \left(1 - \exp \left[-\sigma_{pA} \int \rho_A(b) dz \right] \right)$$

Glauber formula and its extensions represent multiple scatterings in the target, but do not take into account the emergent properties of the quark-gluon system for which there are strong experimental hints.

Recent results suggest presence of a significant nuclear component in the higher-energy cosmic-ray flux from the measurements of the depth of the shower maximum



If there are sources of ultra-high energy cosmic-ray nuclei, these sources should also produce neutrinos!

Murase & Beacom,
PRD **81**, 123001 (2010).

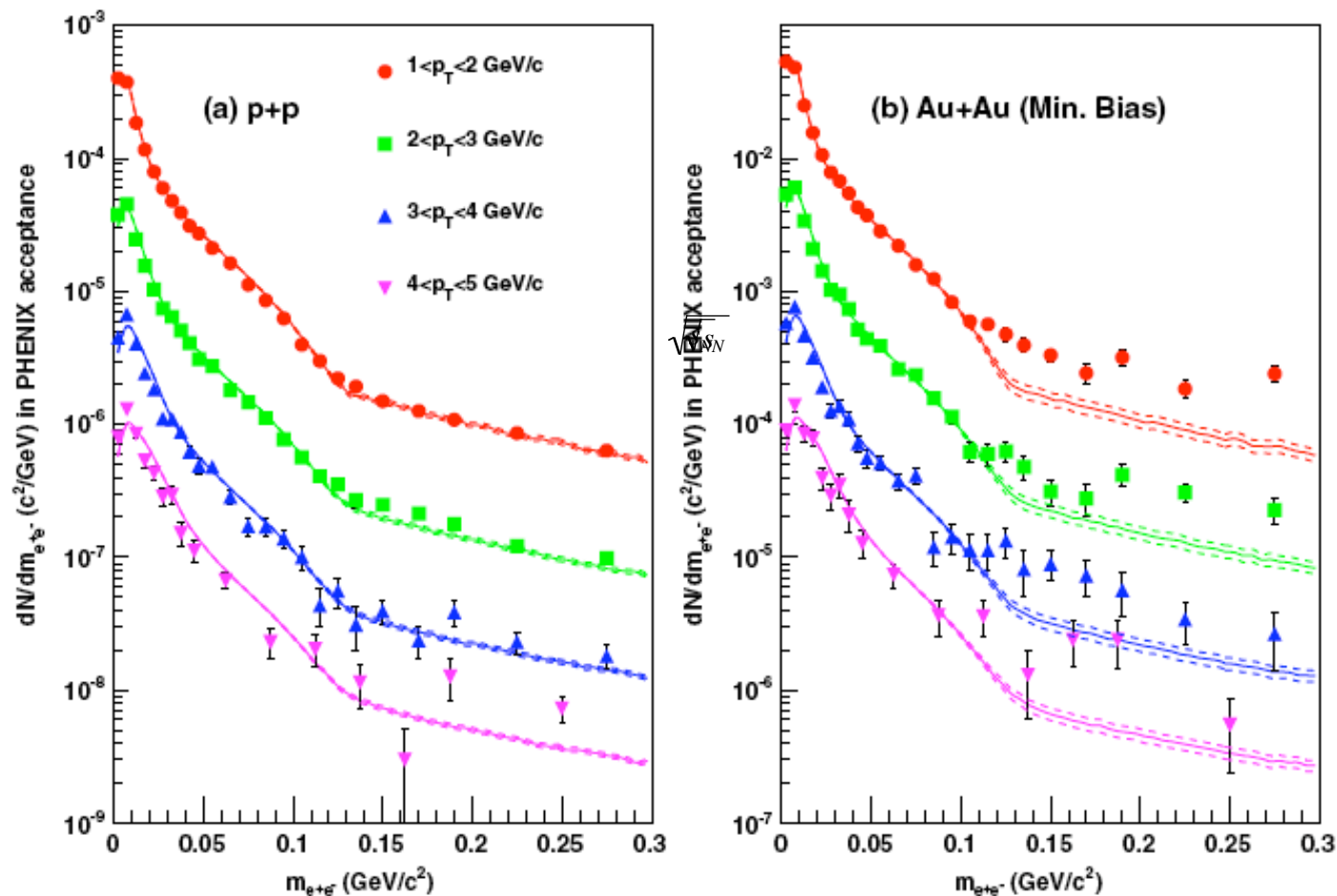
What have we recently learned from relativistic heavy-ion experiments?

- ① An effective “temperature” in 200 GeV Au-Au collisions has been measured. Result is not exactly what we expected. 😊
- ② Negative Binomial Distributions continue to fit the data well. 😐
- ③ There are strong experimental indications that the quark-gluon system formed in relativistic heavy-ion collisions is not a gas, but *almost* a perfect liquid. 😄

1

Measuring the "temperature" at $\sqrt{s_{NN}} \sim 200 \text{ GeV}$ Au-Au collisions

First measure opposite-charge lepton pairs



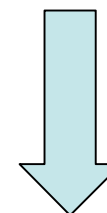
PRL 104, 132301 (2010)

...then convert to real photons by going to zero invariant mass

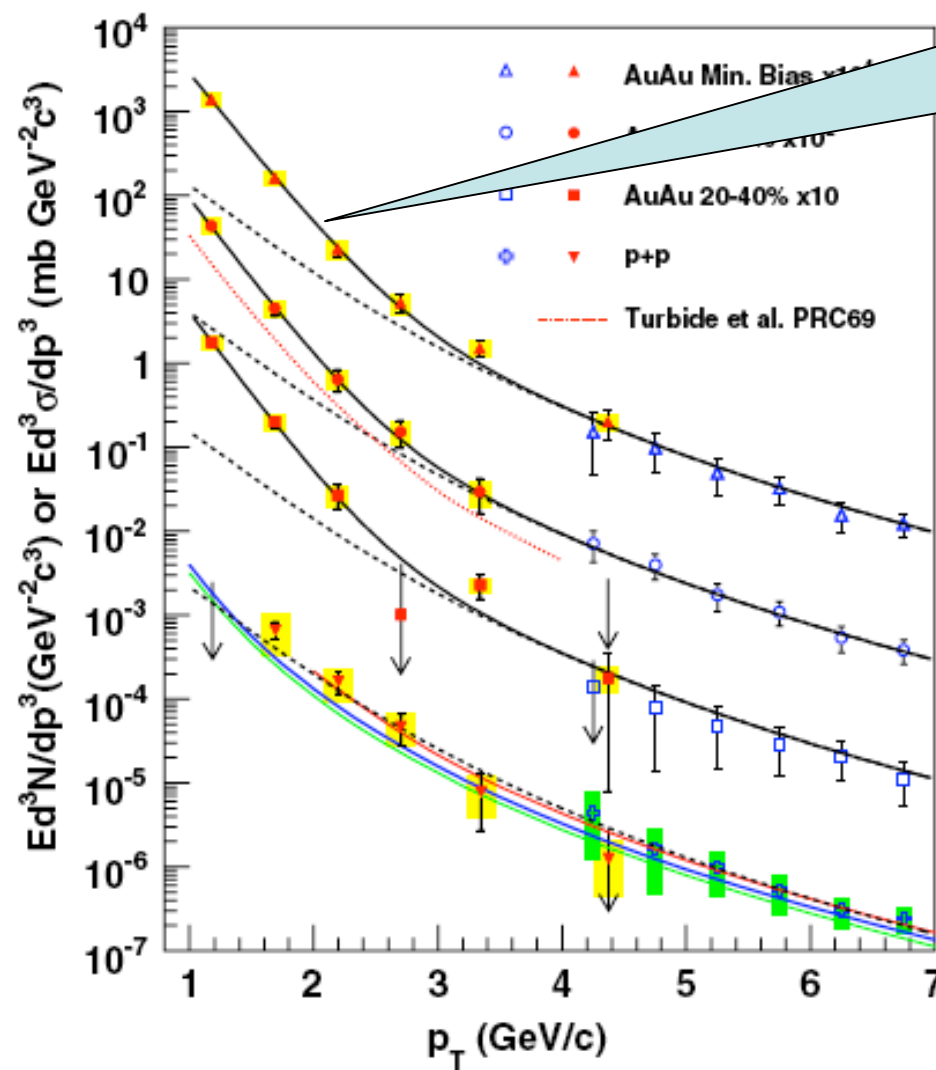
$T_{\text{eff}} = 221 \pm 19 \pm 19 \text{ MeV}$
(effective because γ 's are emitted as the temperature evolves)

+

theoretical input

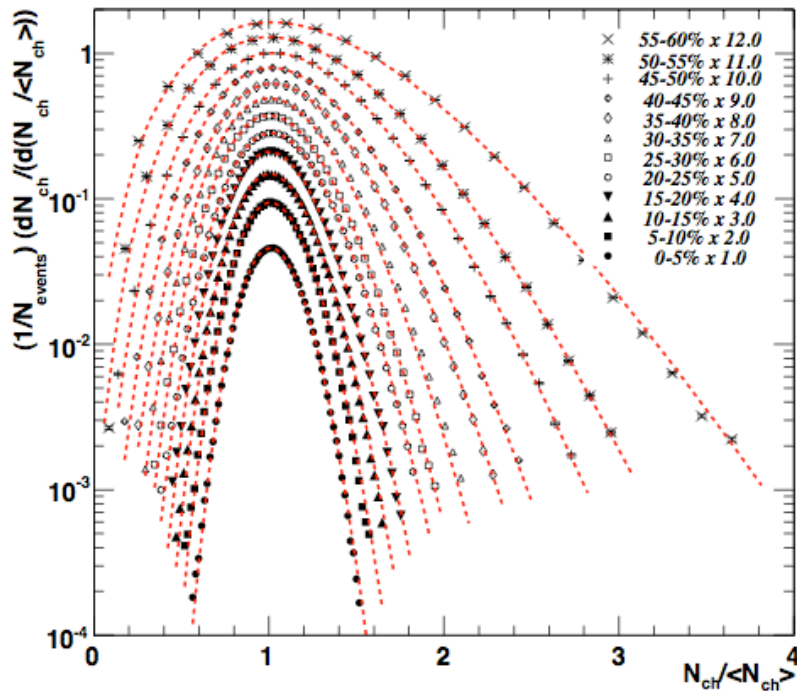


$300 \text{ MeV} < T_{\text{initial}} < 600 \text{ MeV}$
as opposed to the QCD prediction of $\sim 170 \text{ MeV}$!

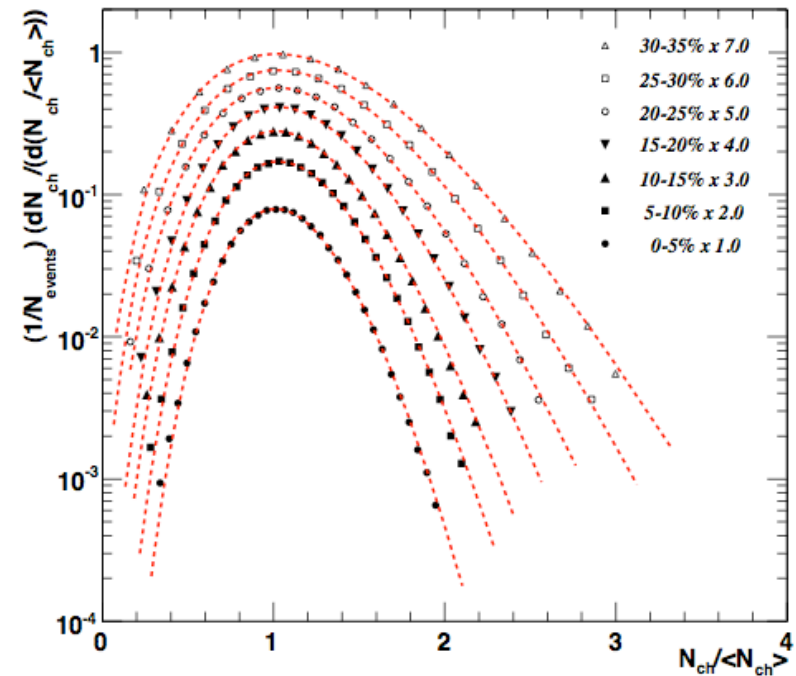


2

Negative Binomial Distribution continues to fit multiplicity fluctuations well



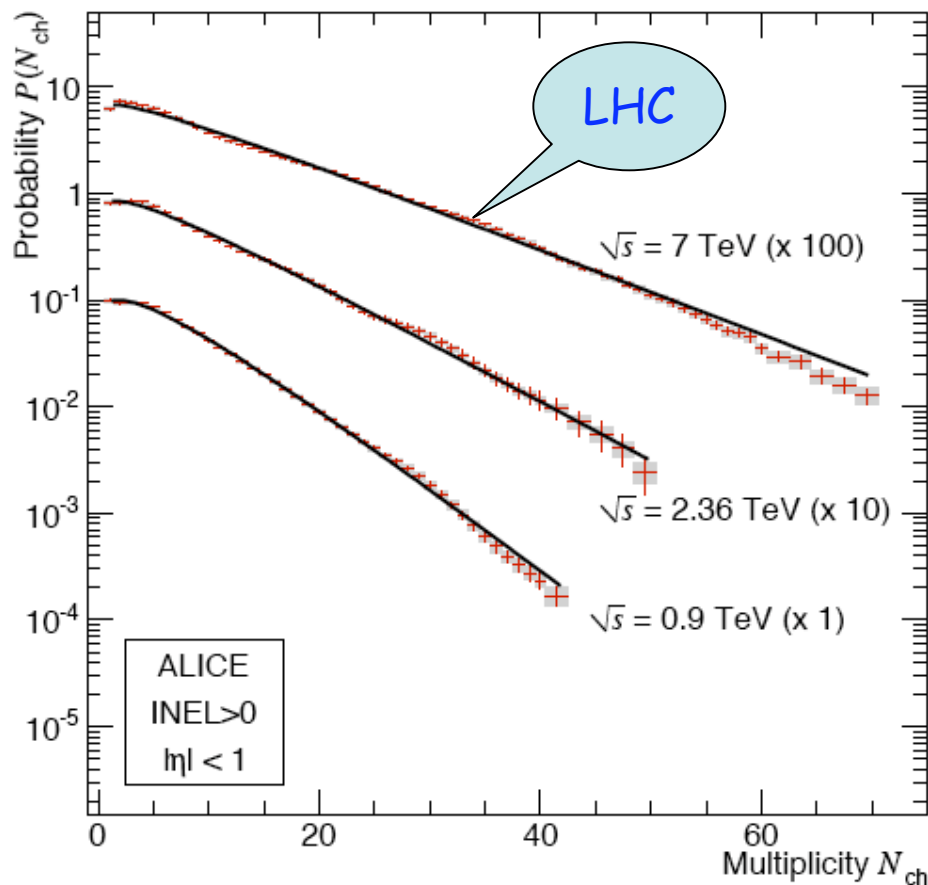
200 GeV Au+Au



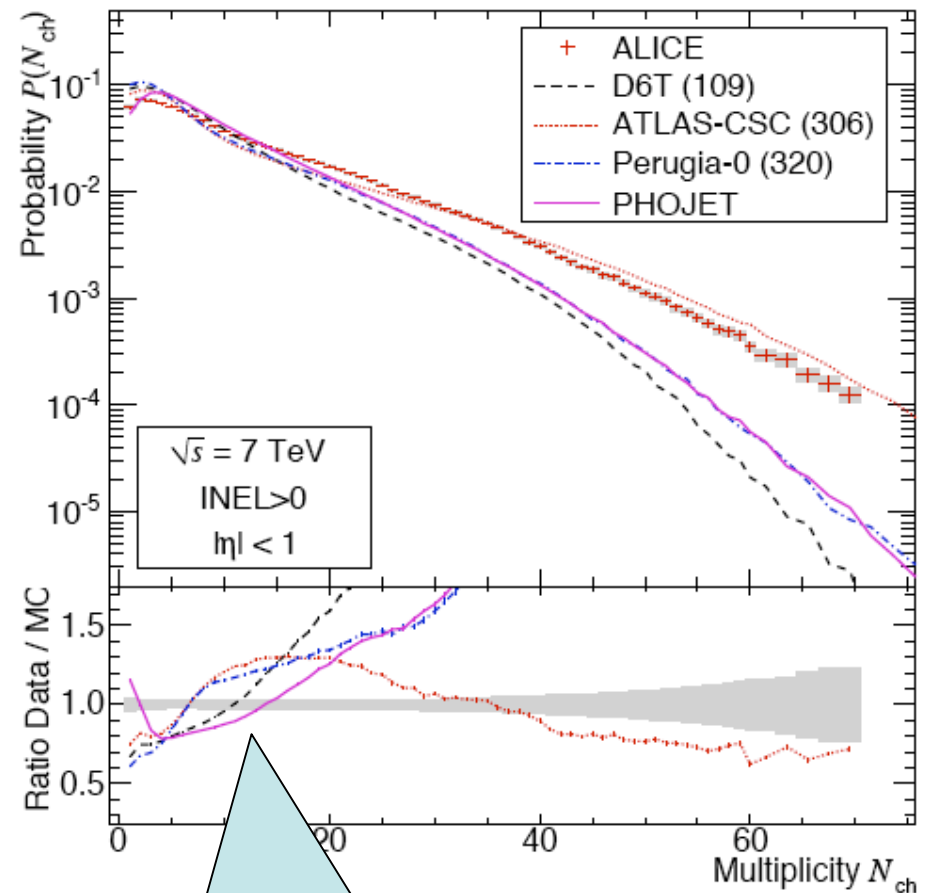
200 GeV Cu+Cu

PHENIX Collaboration, PRC 78, 044902 (2008)

Negative Binomial Distribution continues to fit multiplicity fluctuations well



ALICE Collaboration,
arXiv:1004.3514



The charged-particle
density is higher than
theoretical expectations!

Negative Binomial Distribution

Basic assumption: Probability of emitting n particles by the i th source is b_i^n .

$$\Rightarrow F(\lambda) = \sum_{n=0}^{\infty} P_n \lambda^n = \prod_i^{k_{NBD}} \frac{1 - b_i}{1 - \lambda b_i}$$

All b_i 's are the same \Rightarrow Negative Binomial Distribution:

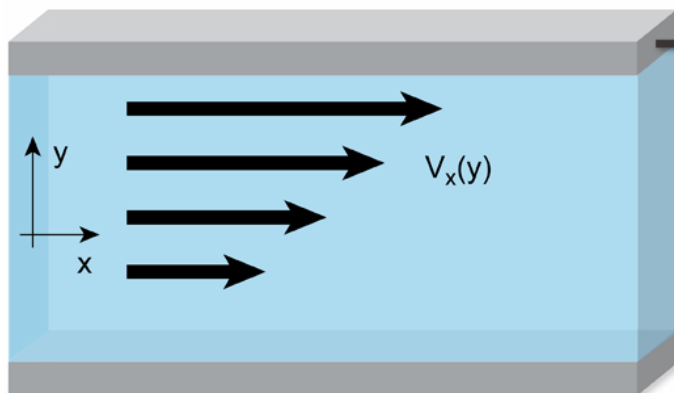
$$\sum_{n=0}^{\infty} P_n \lambda^n = \left(\frac{1 - b}{1 - \lambda b} \right)^{k_{NBD}}$$

Note: In fits k_{NBD} is not constrained to be an integer.

Note: P_n is the complete symmetric function of degree n in the arguments b_i . The ubiquity of negative binomial distribution is likely to be statistical.

3

What is a perfect fluid?



$$\frac{F_x}{A} = \eta \frac{\partial v_x}{\partial y}$$

"good" fluid \Rightarrow
low viscosity, η

Is there a lower bound on shear viscosity, η ? - Heuristic argument

$$F_x = \eta A \frac{\partial v_x}{\partial y}$$

For a dilute (weakly-interacting) gas of quasi-particles:

$$\eta = \frac{1}{3} n \langle p \rangle L_{\text{mean}},$$

n : density, $\langle p \rangle$: ave. momentum, $L_{\text{mean}} = 1/n\sigma$: mean free path

Uncertainty principle:

$$\langle p \rangle L_{\text{mean}} \geq \hbar \quad \Rightarrow \quad \frac{\eta}{n} \geq \hbar$$

Danielewicz & Gyulassy, 1985

Relativistic fluids

For relativistic fluids it is more appropriate to normalize the viscosity to the entropy density

Entropy: $S \propto k_B N$; Entropy density: $s = \frac{S}{V} \propto k_B n$

Is there a lower limit to $\frac{\eta}{s} k_B$?

Kovtun, Son, Starinets, 2005

AdS/CFT duality

Large N conformal field theory in 4D \iff String theory on 5D Anti deSitter space $\times S_5$

Thermal CFT \iff *AdS₅ Black Hole*

CFT entropy \iff Hawking entropy of the black hole

Shear viscosity \iff $\sigma_{\text{graviton absorption}}$

Then one can calculate the viscosity in the strong coupling limit:

$$\frac{\eta}{s} = \frac{1}{4\pi} \frac{\hbar}{k_B}$$

Policastro, Son, Starinets, 2001

Hydrodynamic behavior of the quark-gluon system

Hydrodynamic evolution:

Initial state pressure gradient \Rightarrow Final state velocity gradient

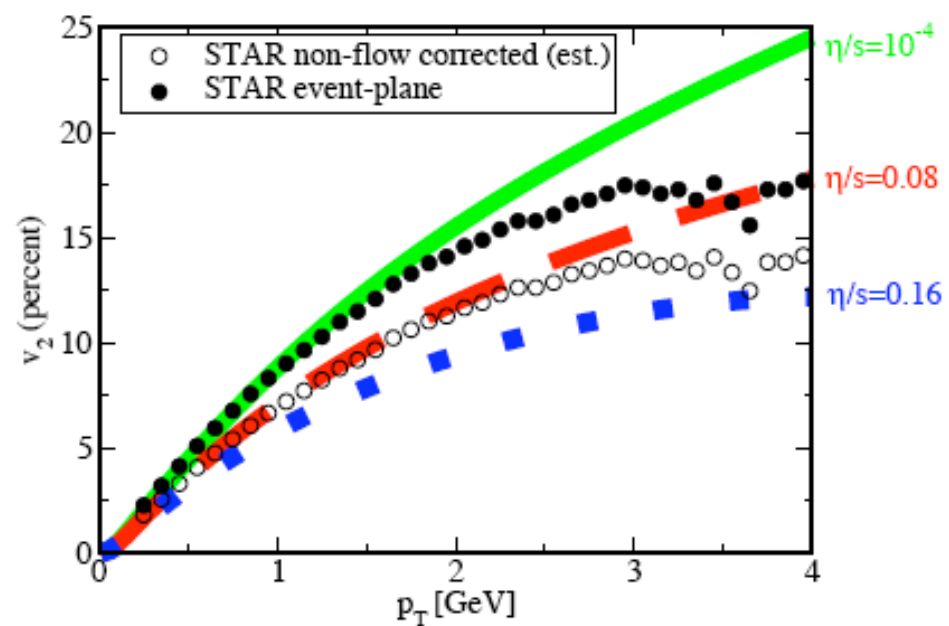
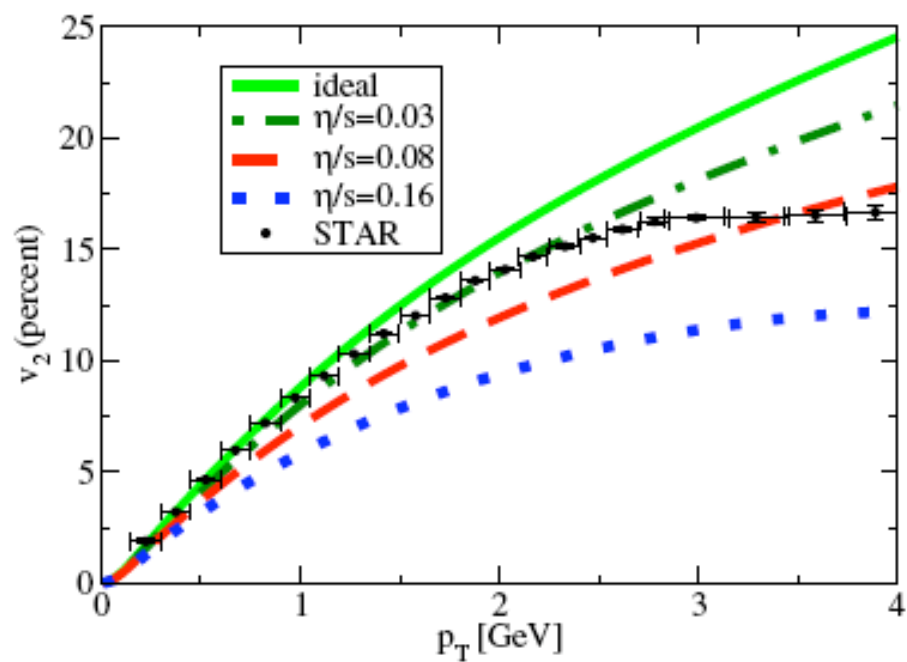
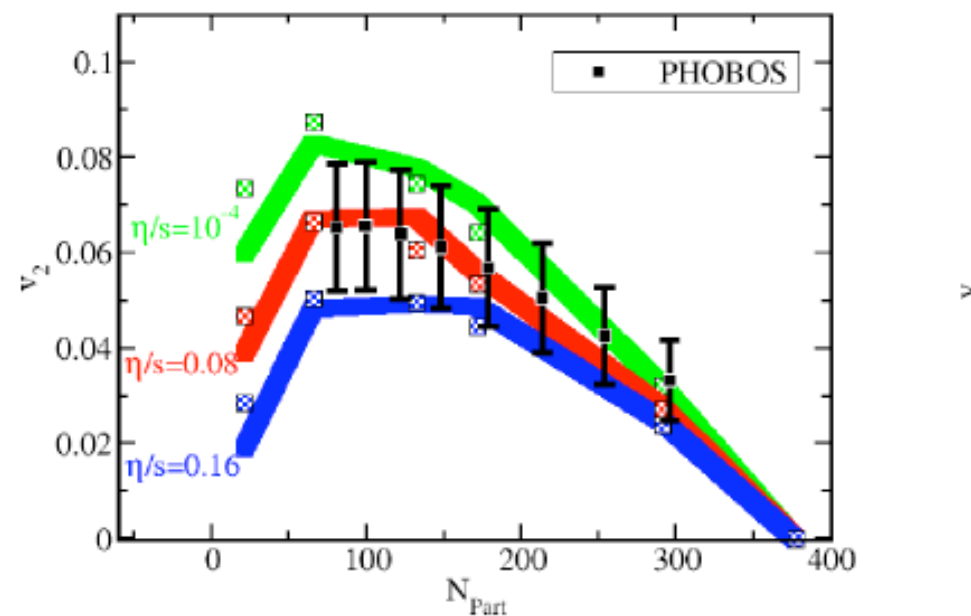
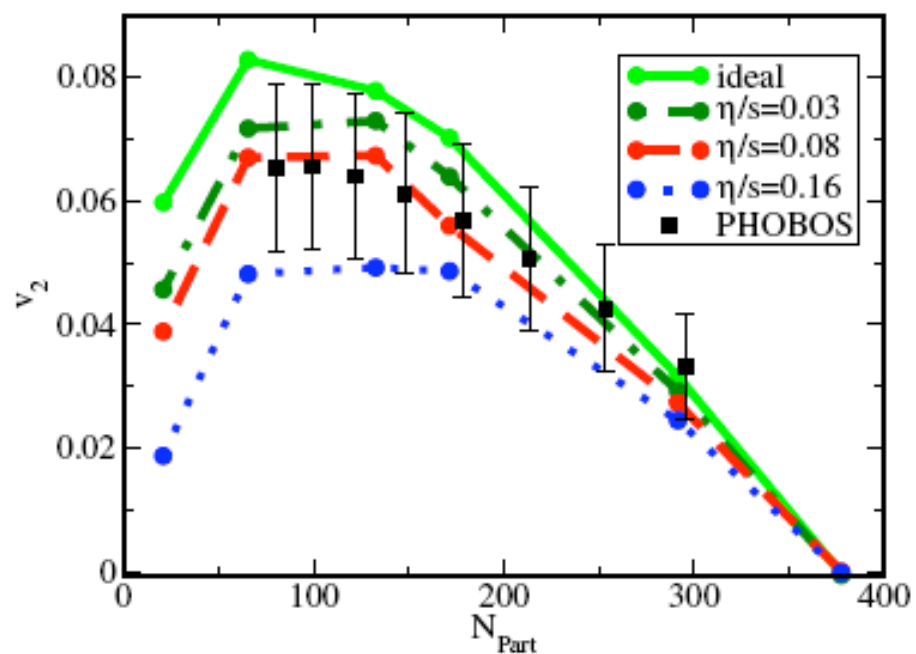
$$\frac{dN_i}{p_T dp_T dy d\phi_p}(b) = \frac{1}{2\pi} \frac{dN_i}{p_T dp_T dy}(b) \left(1 + 2 \sum_i v_n^i \cos(n\phi_p) \right)$$

ϕ_p : azimuthal emission angle relative to the reaction plane

$$\boxed{\text{midrapidity, } y=0} \Rightarrow \boxed{v_{n=\text{odd}} = 0}$$

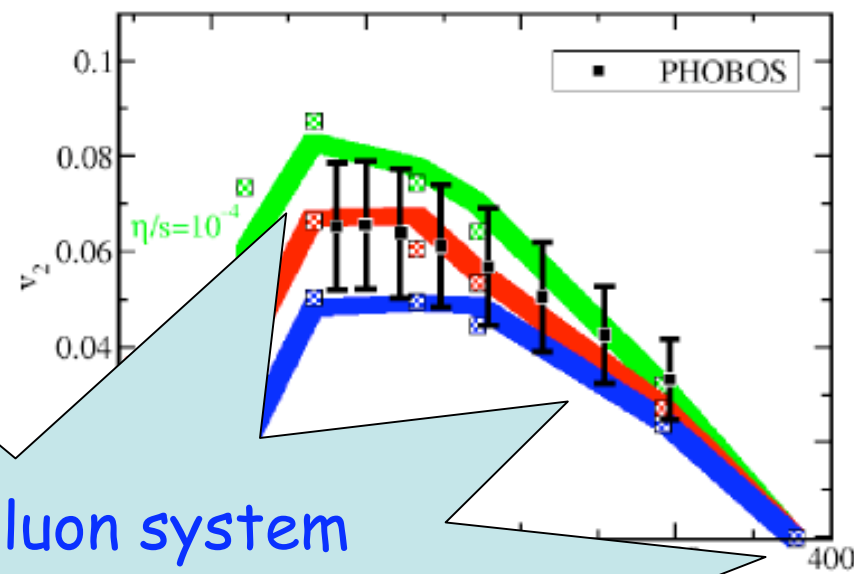
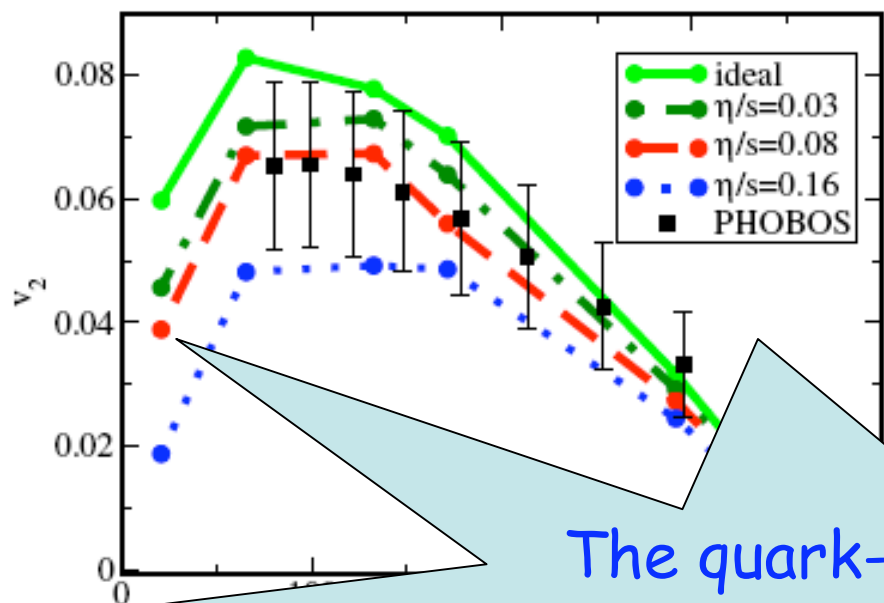
$$\boxed{v_2 > 0} \Rightarrow \boxed{\text{elliptic flow}}$$

The larger the shear velocity, η , is the more transverse expansion of the system slows down.

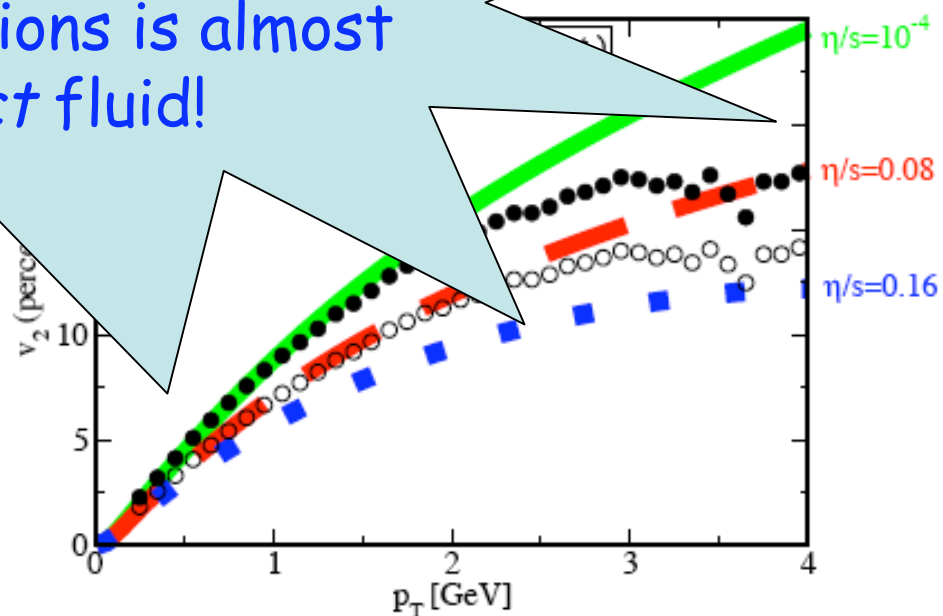
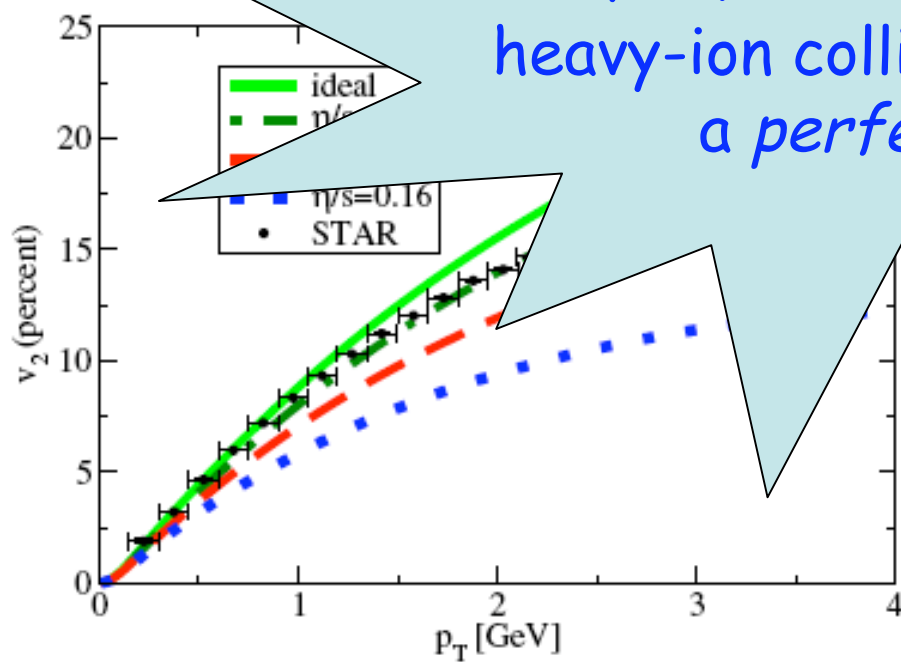


Romatschke & Romatschke, PRL **99**, 172301 (2207)

Heinz, arXiv:0901.4355



The quark-gluon system
formed in relativistic
heavy-ion collisions is almost
a perfect fluid!



Concluding remarks

- At higher energies nuclei are not simply a “collection” of nucleons. Much interesting physics comes into play!
- Recent relativistic heavy-ion experiments found a broad spectrum of interesting phenomena, ranging from the observation of the quark-gluon system as a “perfect fluid” to measuring its temperature.
- Some of the recent cosmic ray experiments suggest an increase in the nuclear component of the cosmic-ray flux at higher energies. Insight gained from recent relativistic heavy-ion experiments could help to understand this nuclear component.