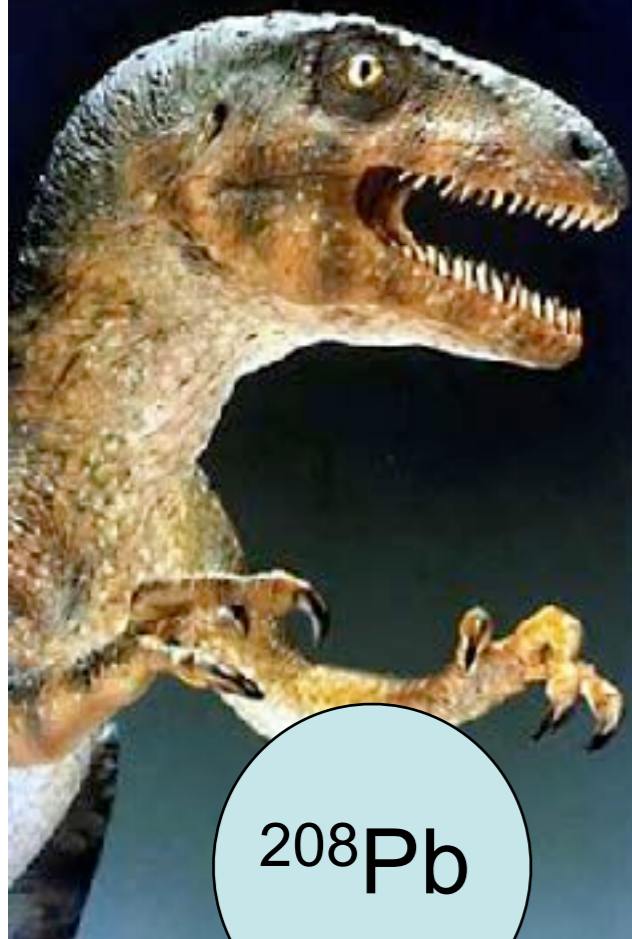


PREX at JLAB



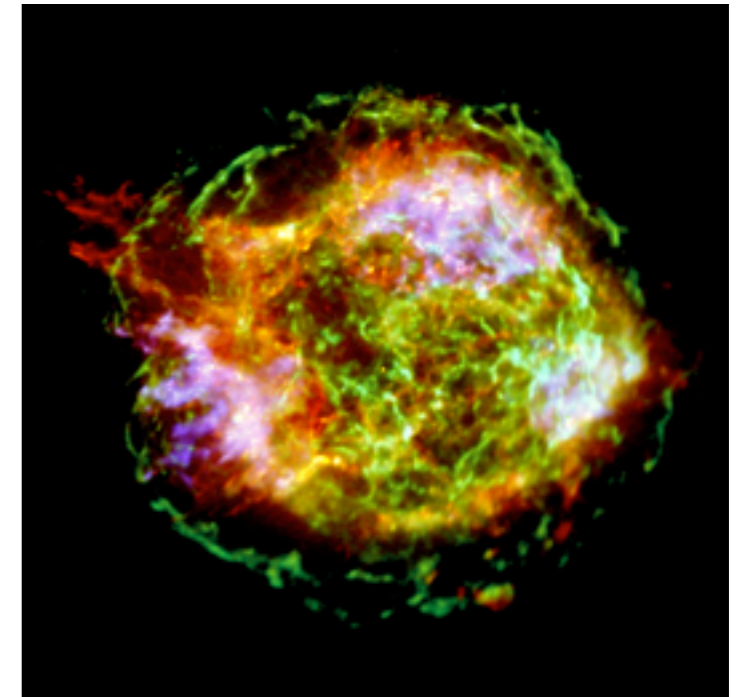
$^{208}\text{Pb}$

# Laboratory and Astronomical Observations of Neutron Rich Matter

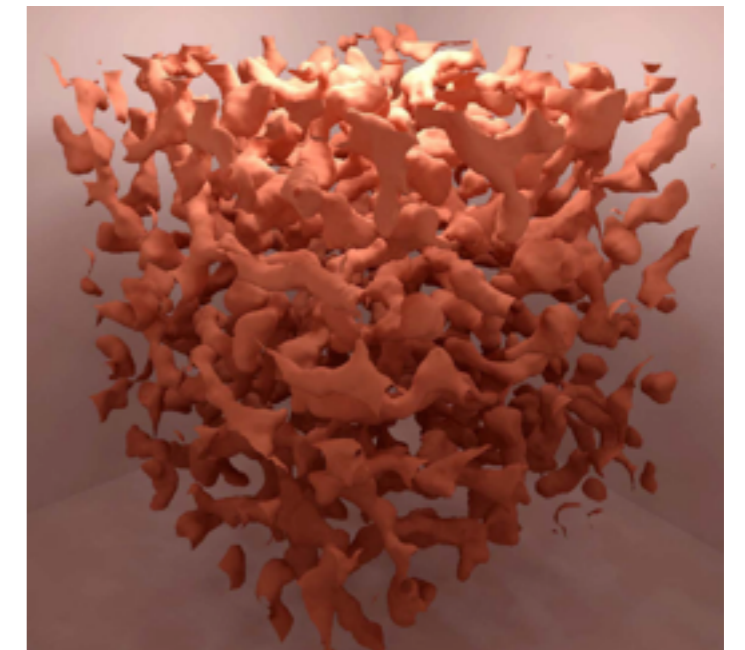
C. J. Horowitz, Indiana University  
Joint CNA/JINA-CEE Winter School on Nuclear  
Astrophysics,  
Shanghai, China, Dec. 2016

# Neutron Rich Matter

- Compress almost anything to  $10^{11}+$  g/cm<sup>3</sup> and electrons react with protons to make neutron rich matter. This material is at the heart of many fundamental questions in nuclear physics and astrophysics.
  - What are the high density phases of QCD?
  - Where did the chemical elements come from?
  - What is the structure of many compact and energetic objects in the heavens, and what determines their electromagnetic, neutrino, and gravitational-wave radiations?
- Interested in neutron rich matter over a tremendous range of density and temperature were it can be a *gas, liquid, solid, plasma, liquid crystal (nuclear pasta), superconductor, superfluid, color superconductor...*



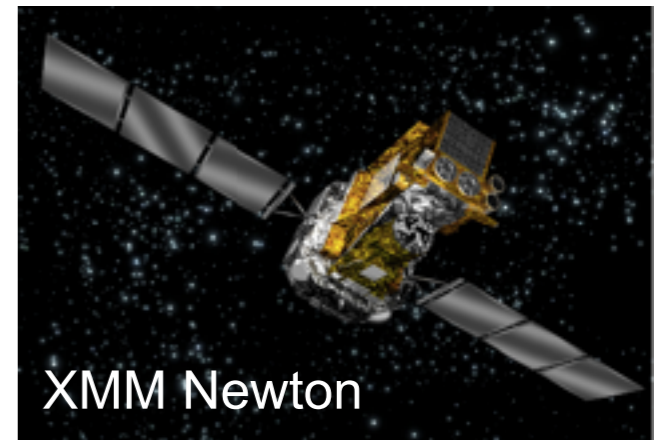
Supernova remanent  
Cassiopea A in X-rays



MD simulation of Nuclear  
Pasta with 100,000 nucleons

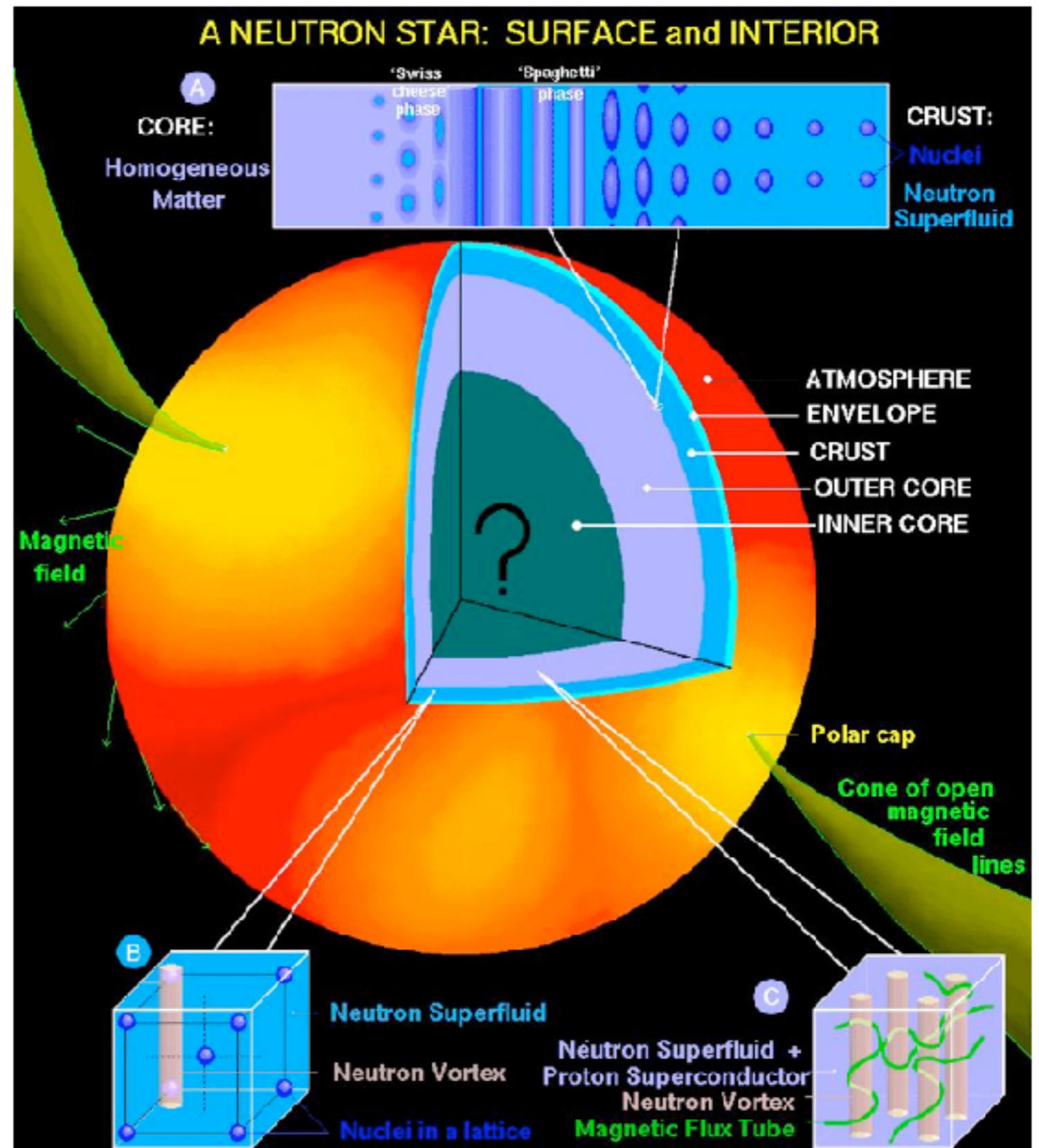
# Probes of Neutron Rich Matter

- **Multi-Messenger Astronomy:** “seeing” the same event with very different probes should lead to fundamental advances. Often **photons** from *solid* neutron star crust, supernova **neutrinos** from low density *gas*, and **gravitational waves** from energetic motions of *liquid* interior of neutron stars.
- **Laboratory:** Nuclei are liquid drops so most experiments probe liquid n rich matter. However one can also study vapor phase by evaporating nucleons.
  - Electroweak measurements, Heavy ion collisions, Radioactive beams of neutron rich nuclei...
- **Computational:**
  - Chiral effective field theory (and MC calc. with phenomenological NN and NNN forces) depends on important and poorly known *three neutron forces*.
  - **Chiral expansion does not converge at high densities. This strongly limits first principle calculations.**
- Increases importance of laboratory experiments and astrophysical observations.



# Neutron stars

- Neutron stars are formed from the collapse of a massive star in a supernova explosion.
- Mass  $\sim 1.4 M_{\text{sun}}$ , Radius  $\sim 10$  km
- Solid crust  $\sim 1$  km thick ( $\sim 10^{12}$  g/cm<sup>3</sup>) over liquid outer core of neutron rich matter ( $\sim 10^{15}$  g/cm<sup>3</sup>).
- Possible exotic phase in center: de-confined quark matter, strange matter, meson condensates, color superconductor...
- Structure determined by Equation of State (pressure vs density) of n rich matter.
- Figure: **Dany Page**, UNAM



# Equation of state

- Neutron star structure is determined by the equation of state: pressure as a function of energy density.
- $P = -dE/dV = -d(E/A)/d(V/A) = n^2 d(E/A)/dn$ ,  
 $V$ =volume.
- Here the **baryon density**  $n=(A/V)$ . [particles/fm<sup>3</sup>]
- The **energy density** is  $\epsilon=E/V=(E/A)n$ . [MeV/fm<sup>3</sup>]
- The **mass density** is  $\rho=\epsilon/c^2$ . [g/cm<sup>3</sup>] Note that  $E$  includes rest mass  $mc^2$ . (In nonrel. limit  $\rho=mn$ )
- Calculate both  $\epsilon(n)$ ,  $P(n)$  and tabulate  **$P=P(\epsilon)$**   
**which is the equation of state.**

# NS hydrostatic equilibrium

- As you go deeper into a NS pressure  $p$  rises because of weight of material above you.

$$\frac{dp}{dr} = -\frac{G\rho(r)\mathcal{M}(r)}{r^2} = -\frac{G\epsilon(r)\mathcal{M}(r)}{c^2 r^2}$$

$$\frac{d\mathcal{M}}{dr} = 4\pi r^2 \rho(r) = \frac{4\pi r^2 \epsilon(r)}{c^2}$$

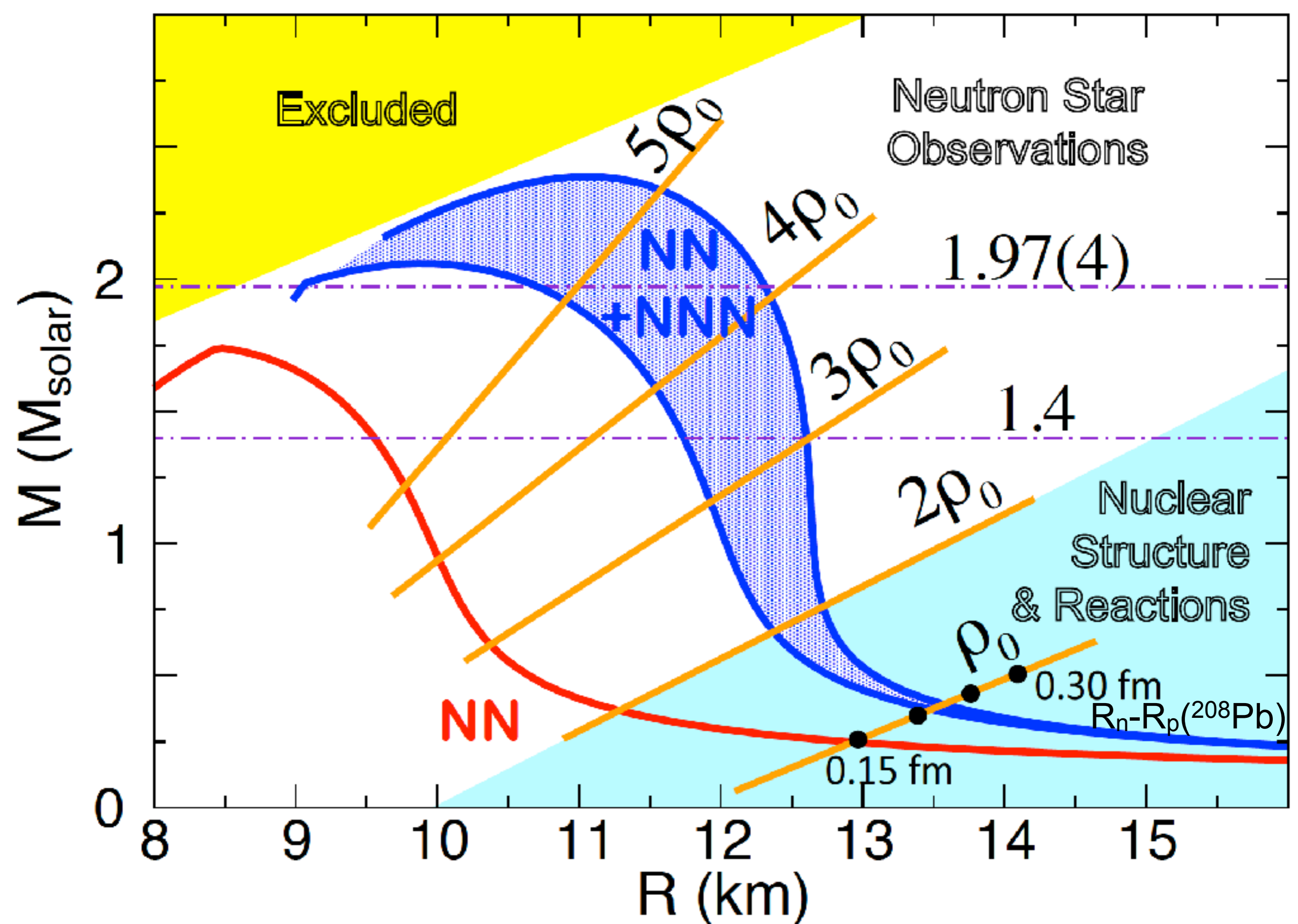
$$\mathcal{M}(r) = 4\pi \int_0^r r'^2 dr' \rho(r') = 4\pi \int_0^r r'^2 dr' \epsilon(r') / c^2$$

- Above assumes Newtonian gravity. See NS for undergraduates R. R. Silbar, S. Reddy, Am. J. Phys. **72** (2004) 892

# TOV Equation

$$\frac{dp}{dr} = -\frac{G\epsilon(r)\mathcal{M}(r)}{c^2r^2} \left[1 + \frac{p(r)}{\epsilon(r)}\right] \left[1 + \frac{4\pi r^3 p(r)}{\mathcal{M}(r)c^2}\right] \left[1 - \frac{2G\mathcal{M}(r)}{c^2r}\right]^{-1}$$

- First two terms in brackets are special relativistic corrections while last term is a GR correction from the curvature of space.
- All three terms act to increase  $p$  in star. Effects are nonlinear, eventually gravity wins and star will collapse to black hole  $\rightarrow$  maximum mass for NS.
- For given EOS  $p=p(\epsilon)$  guess central  $p$  and  $\epsilon(p)$ . Integrate  $p$  outward from  $r=0$  until  $p=0$ . This gives radius and mass of star. Repeat for different central  $p$  to determine mass of NS as a function of radius.

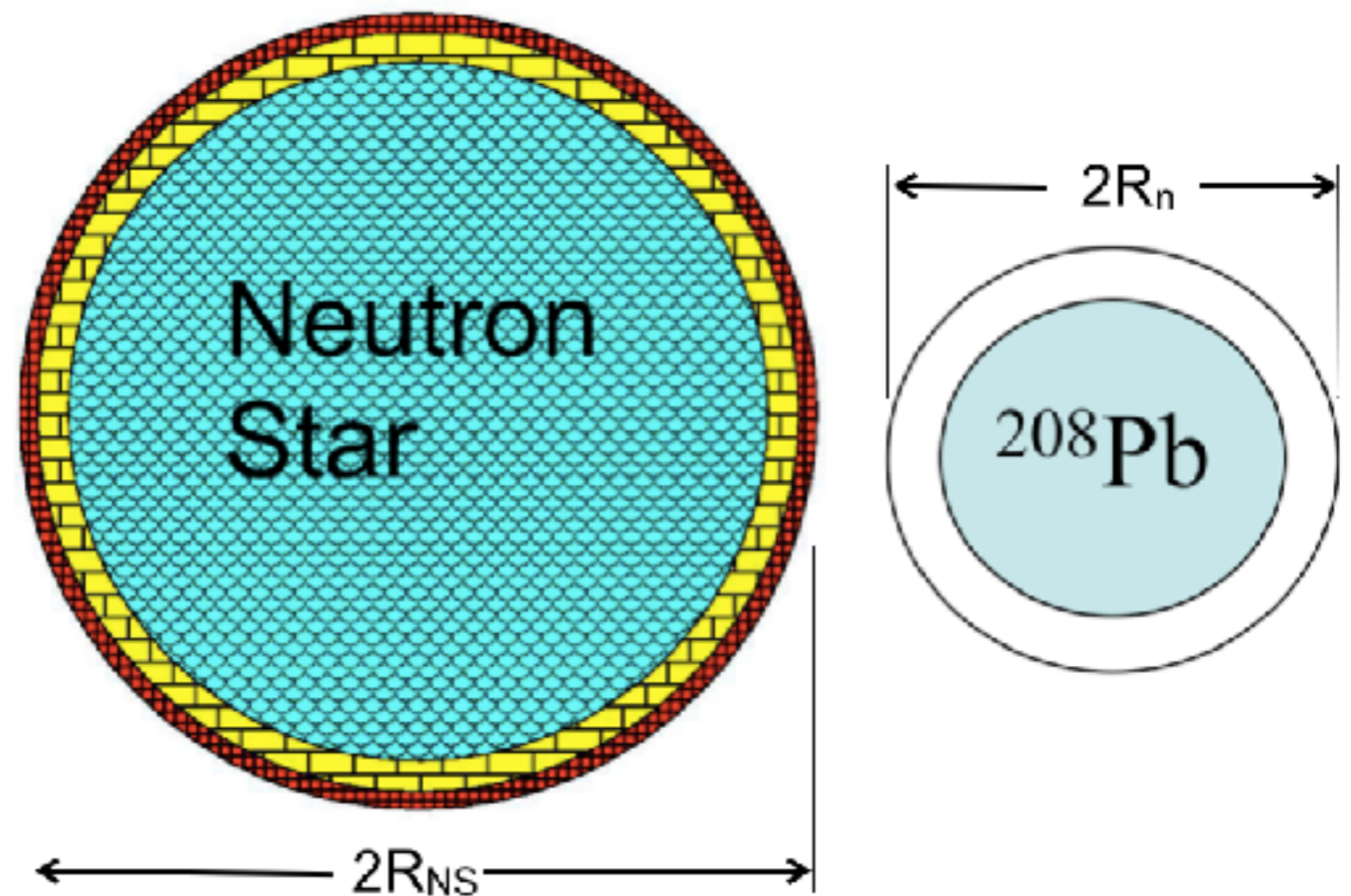




# Equation of state and Radii

- Neutron radius of heavy nucleus determines:
  - Equation of state (pressure vs density) of neutron matter.
  - Density dependence of symmetry energy.
- Radius of a neutron star:
  - Also determined by equation of state of neutron matter.
  - Can be measured with X-ray observations.

## Neutron Star radius versus $^{208}\text{Pb}$ Radius

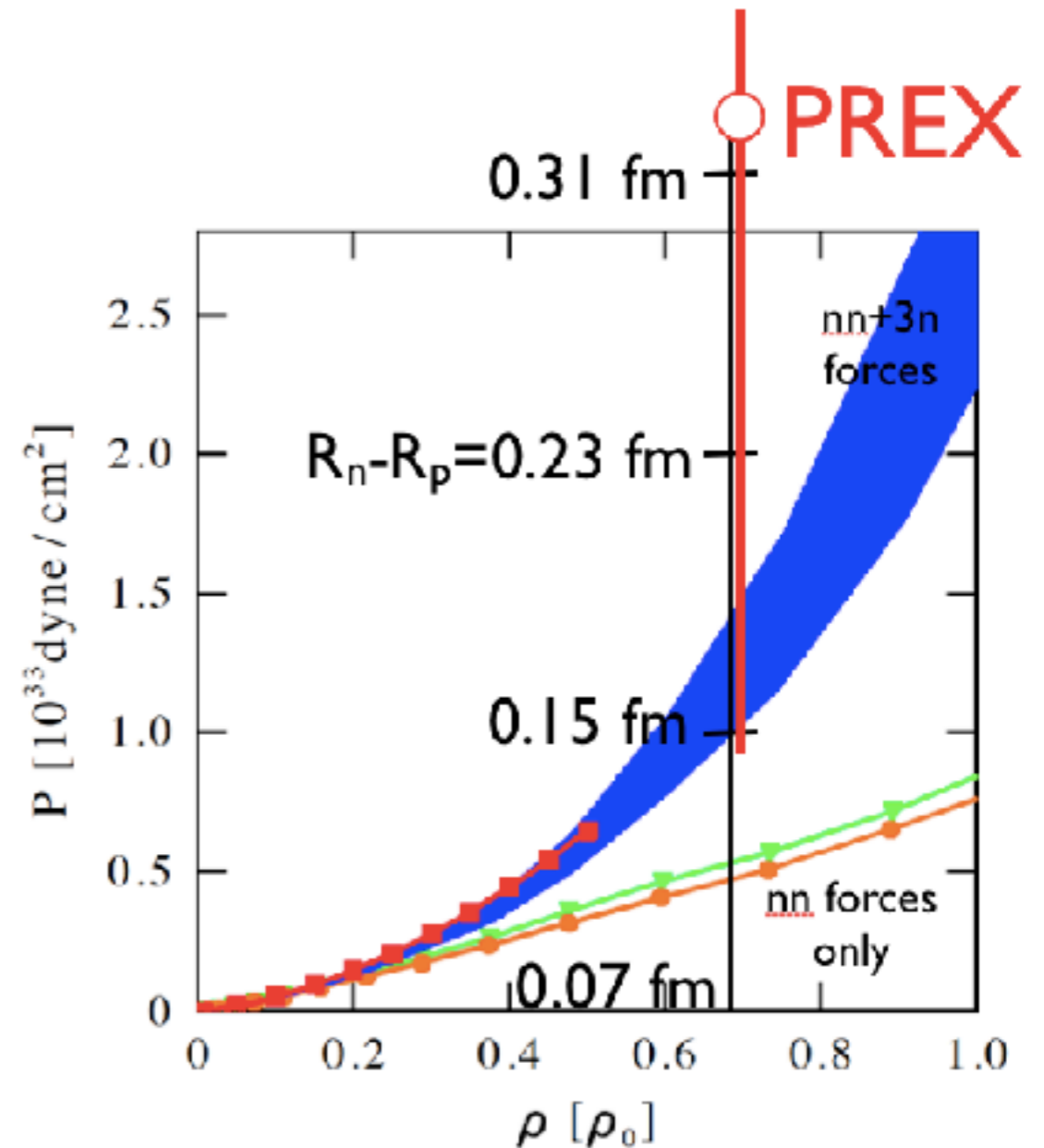


Neutron star is 18 orders of magnitude larger than  $^{208}\text{Pb}$  nucleus but both radii determined by same strong interaction physics.

# Pb Radius vs Neutron Star Radius

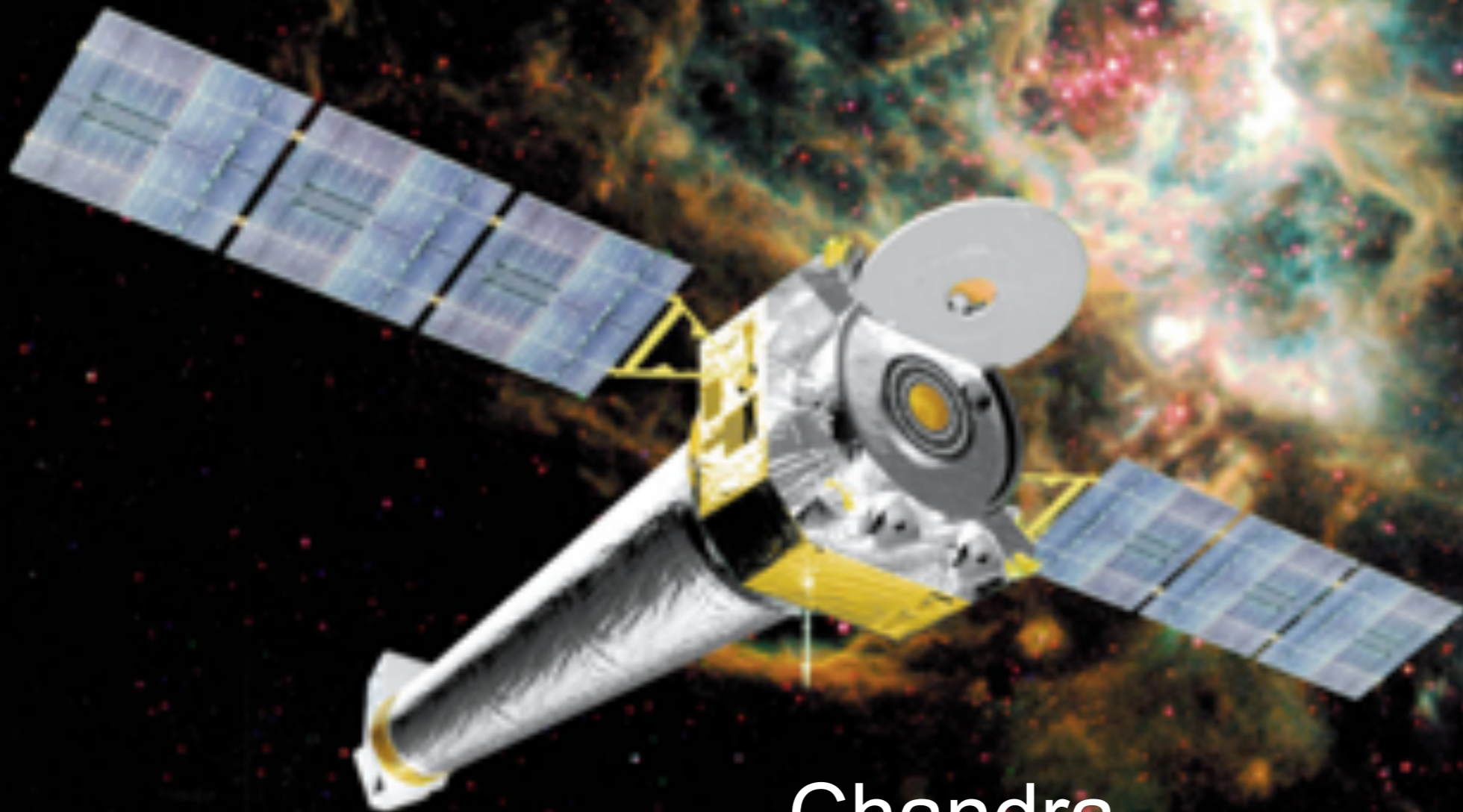
- The  $^{208}\text{Pb}$  radius constrains the pressure of neutron matter at subnuclear densities. Typel + Brown find sharp correlation between  $P$  at  $2/3 \rho_0$  and  $R_n$ .
- The NS radius depends on the pressure at nuclear density and above. Central density of NS few to  $10 \times$  nuclear density.
- *Pb radius probes low density, NS radius medium density, and maximum NS mass probes high density equation of state.*
- An observed softening of EOS with density (smaller increase in pressure) could strongly suggest a transition to an exotic high density phase such as quark matter, strange matter, or a color superconductor...

J. Piekarewicz, CJH



Chiral EFT calc. of pressure  $P$  of neutron matter by Hebeler et al. including three neutron forces (blue band) agree with PREX results but two nucleon only calculations yield smaller  $P$ .

# Electromagnetic Messengers



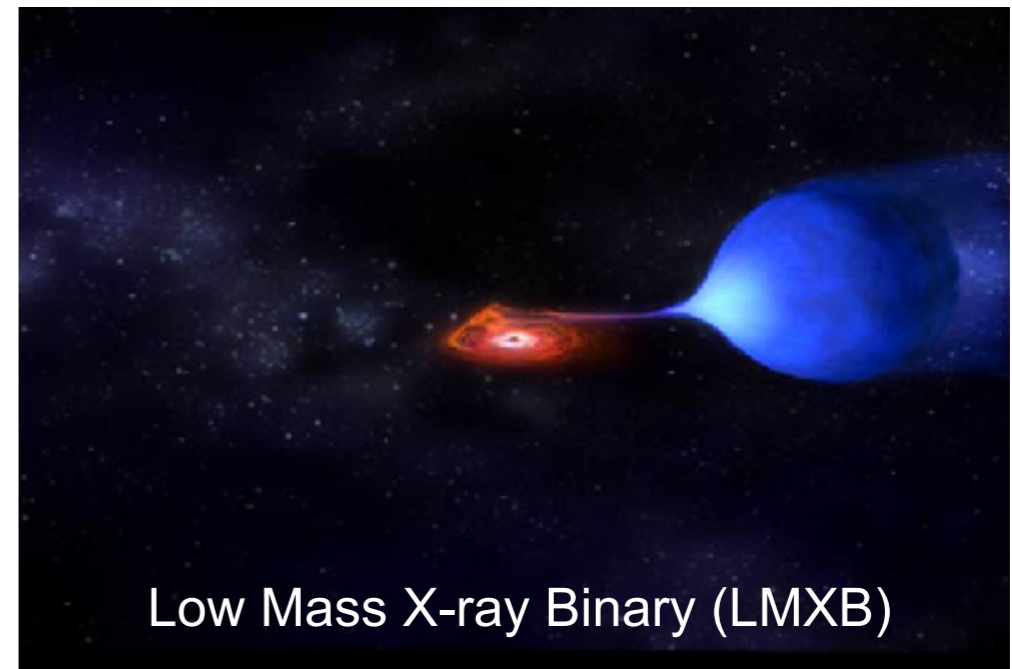
Chandra

# X-ray observations of NS radii, masses

- Deduce surface area from luminosity, temperature from X-ray spectrum.

$$L_{\gamma} = 4\pi R^2 \sigma_{\text{SB}} T^4$$

- Complications:
  - Need distance (parallax for nearby isolated NS...)
  - Non-blackbody corrections from atmosphere models can depend on composition and B field.
  - Curvature of space: measure combination of radius and mass.
- **NS in globular clusters:** expect simple nonmagnetic hydrogen atmospheres and know distance.
- **X-ray bursts:** NS accretes material from companion that ignites a runaway thermonuclear burst.

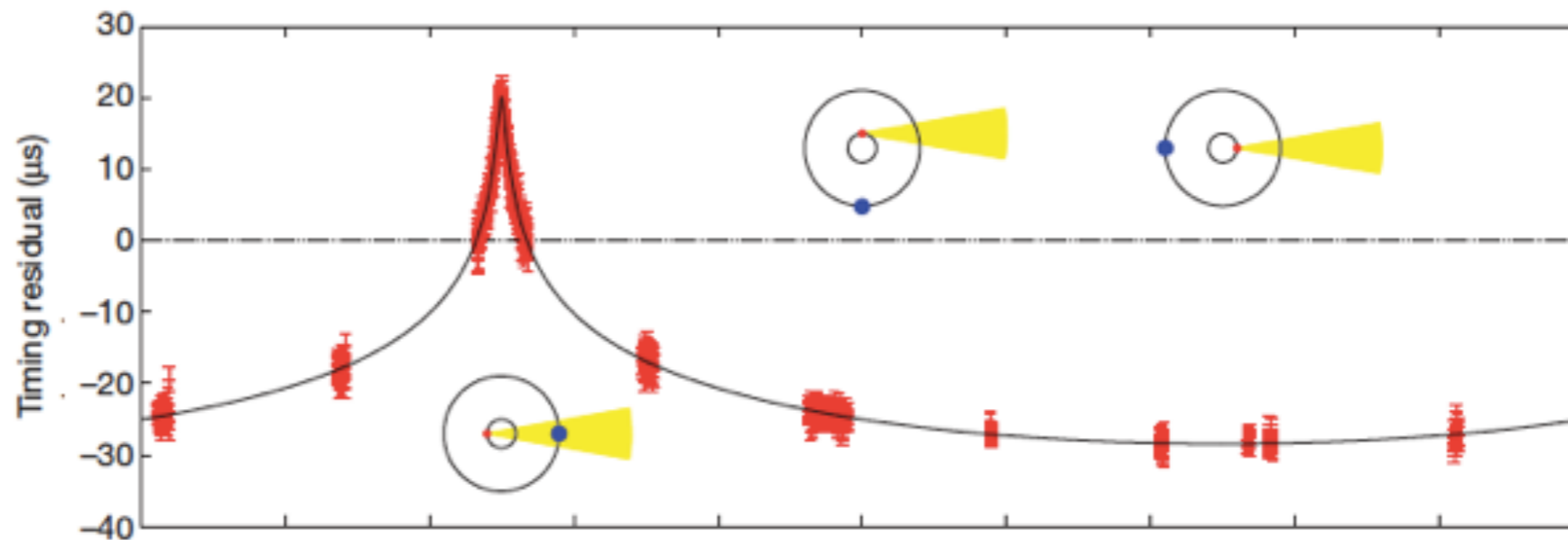


- **Eddington luminosity:** when radiation pressure balances gravity --> gives both M and R!
- Steiner, Lattimer, Brown combine X-ray observations --> 1.4  $M_{\text{sun}}$  star has  $\sim 12$  km radius.
- However important uncertainties may remain in extracted radii. Suleimanov and Poutanen use more sophisticated atmosphere models and find larger radii.

# Discovery of $2M_{\text{sun}}$ Neutron Star

Demorest et al: PSR J1614-2230 has  $1.97 \pm 0.04 M_{\text{sun}}$ .

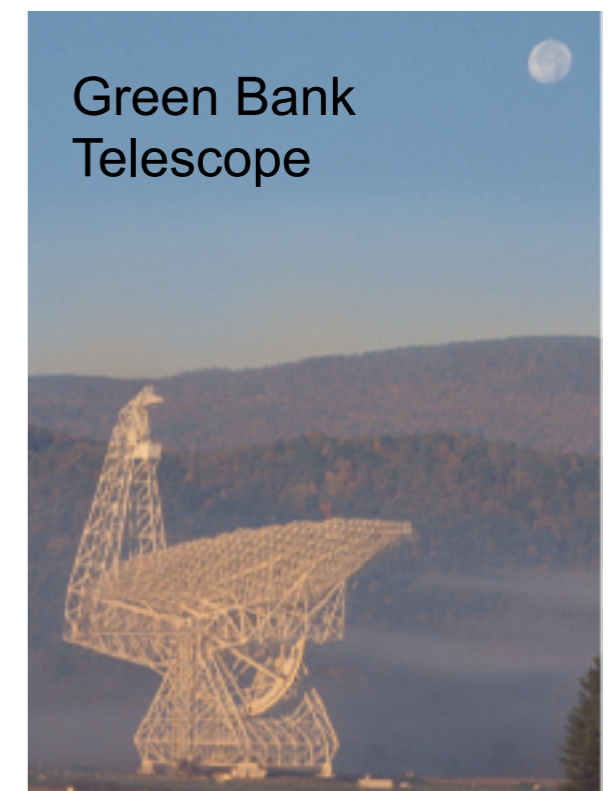
Delay  
in  
pulse  
arrival



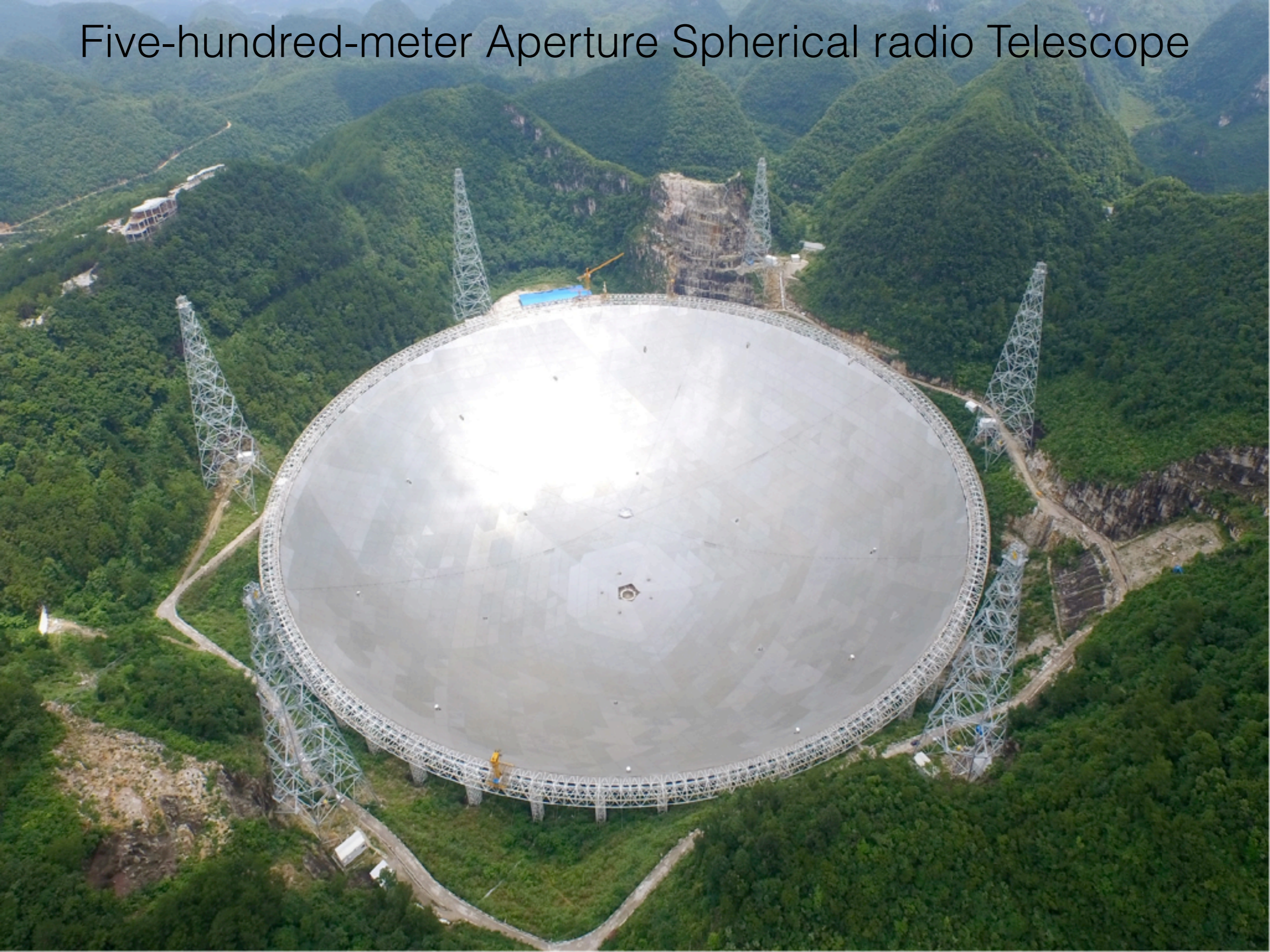
NS +  
White  
Dwarf  
Binary

Orbital phase

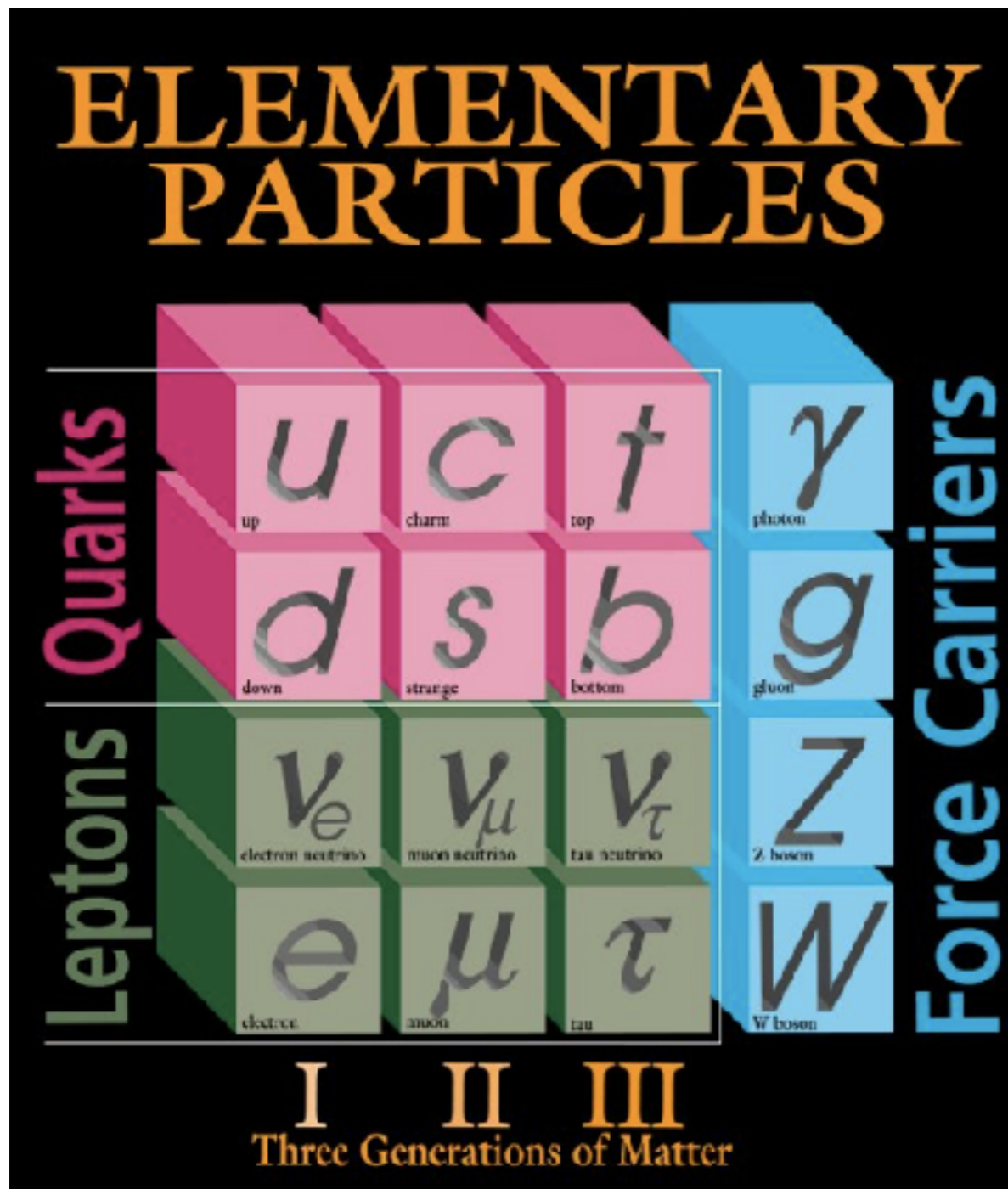
- The equation of state of neutron rich matter (pressure vs density) at high densities must be stiff enough (have a high enough  $p$ ) to support this mass against collapse to a black hole. *All soft EOS are immediately ruled out!*
- *However this does not tell composition of dense matter be it neutron/ proton, quark, hyperon...*
- *NS cooling (by neutrinos) sensitive to composition.*



# Five-hundred-meter Aperture Spherical radio Telescope



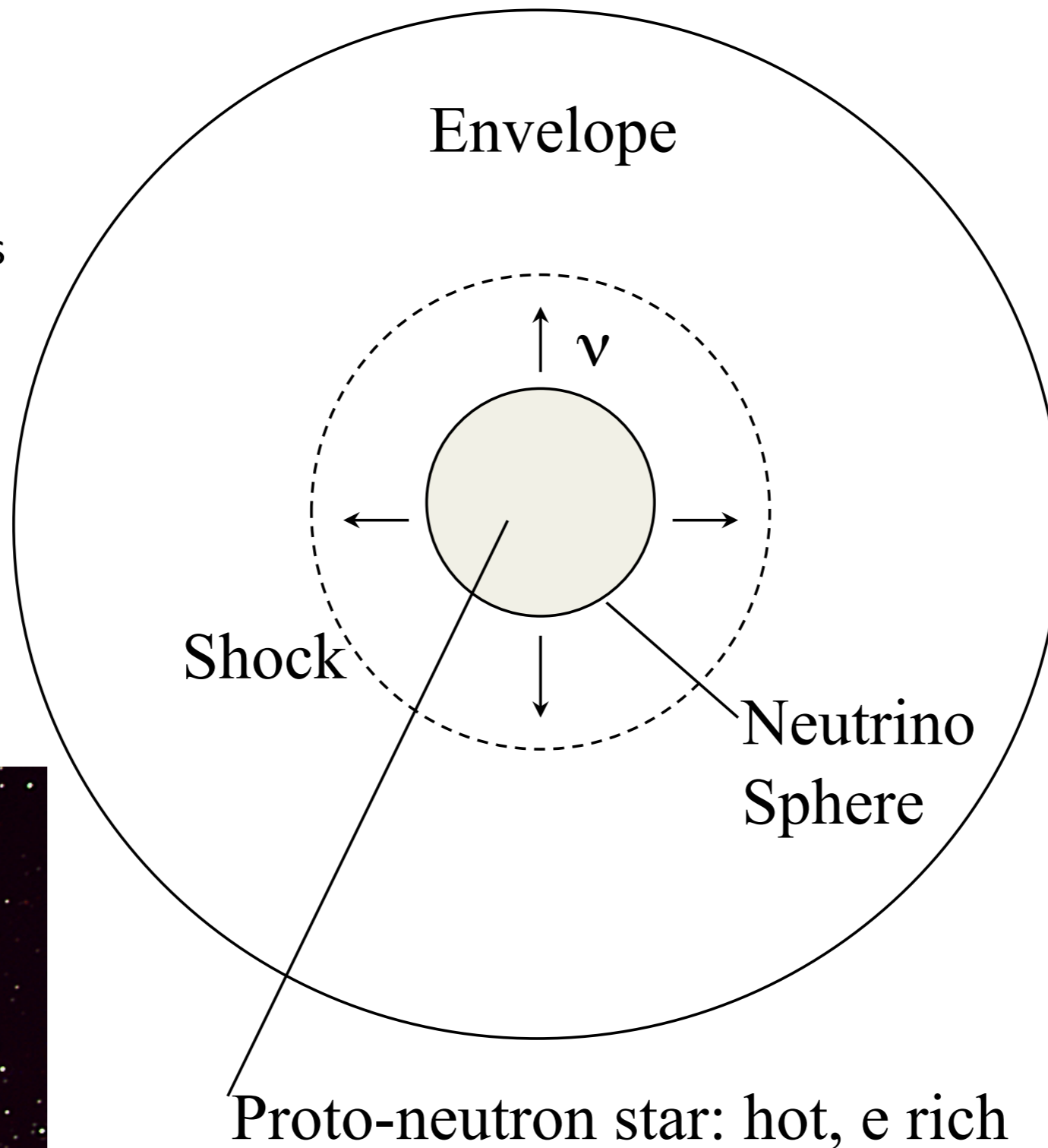
# Neutrino messengers



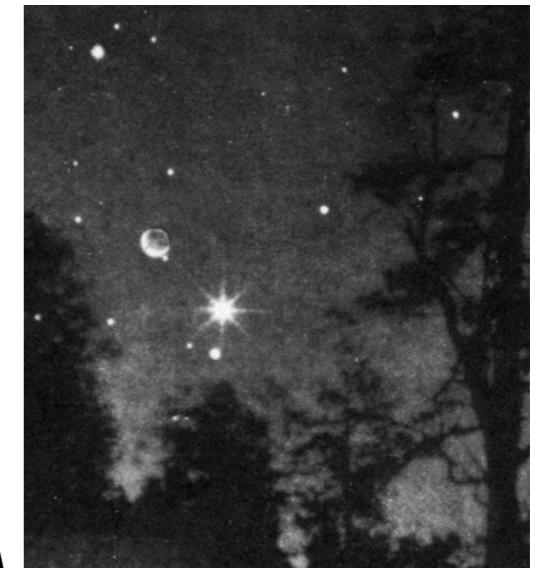
- All three flavors of neutrinos and their antineutrinos (electron, mu, tau) are radiated in core collapse supernovae.
- **Neutrinos carry unique flavor information all the way to earth.**
- Note, neutrinos are somewhat forgetful messengers because of oscillations.

# NS Born in Core Collapse Supernovae

Core of massive star collapses to form proto-neutron star.  $\nu$ s form neutron star energizes shock that ejects outer 90% of star.

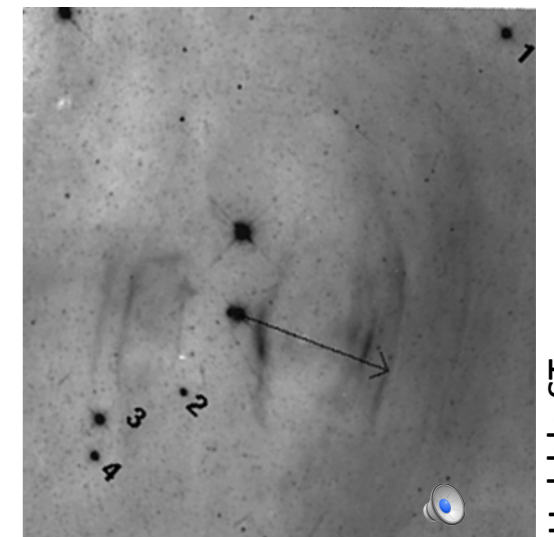


Crab nebula



July 5, 1054

Crab Pulsar

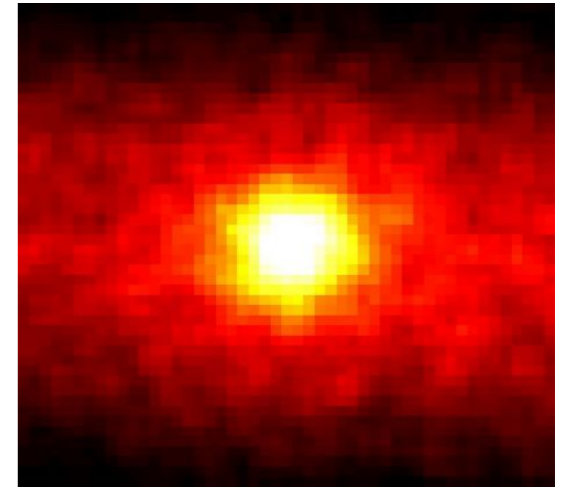


Hubble ST



# Astrophysical Neutrinos

- Much progress with **solar neutrinos** over last decade documenting neutrino oscillations and confirming solar model.
- Core collapse supernovae radiate gravitational binding energy of NS,  $\sim 3/5 GM^2/R = 0.1$  to  $0.2 M_{\text{sun}} c^2$  in  $10^{58}$  neutrinos of  $\sim 10$  MeV each.
- Historic first detection of  $\sim 20$  **supernova neutrinos** from SN1987A (detection efficiency  $10^{-57}$ ).
- New underground dark matter, solar nu,... experiments will be sensitive to nu from the next galactic supernova (SN).
- Example: ton scale dark matter detectors sensitive to SN neutrinos via nu-nucleus elastic scattering. [CJH, K. Coakley, D. McKinsey, PRD**68**(2003)023005]
- Expect ten thousand or more events from next galactic SN!



Sun in neutrinos

# Neutrino Spectra

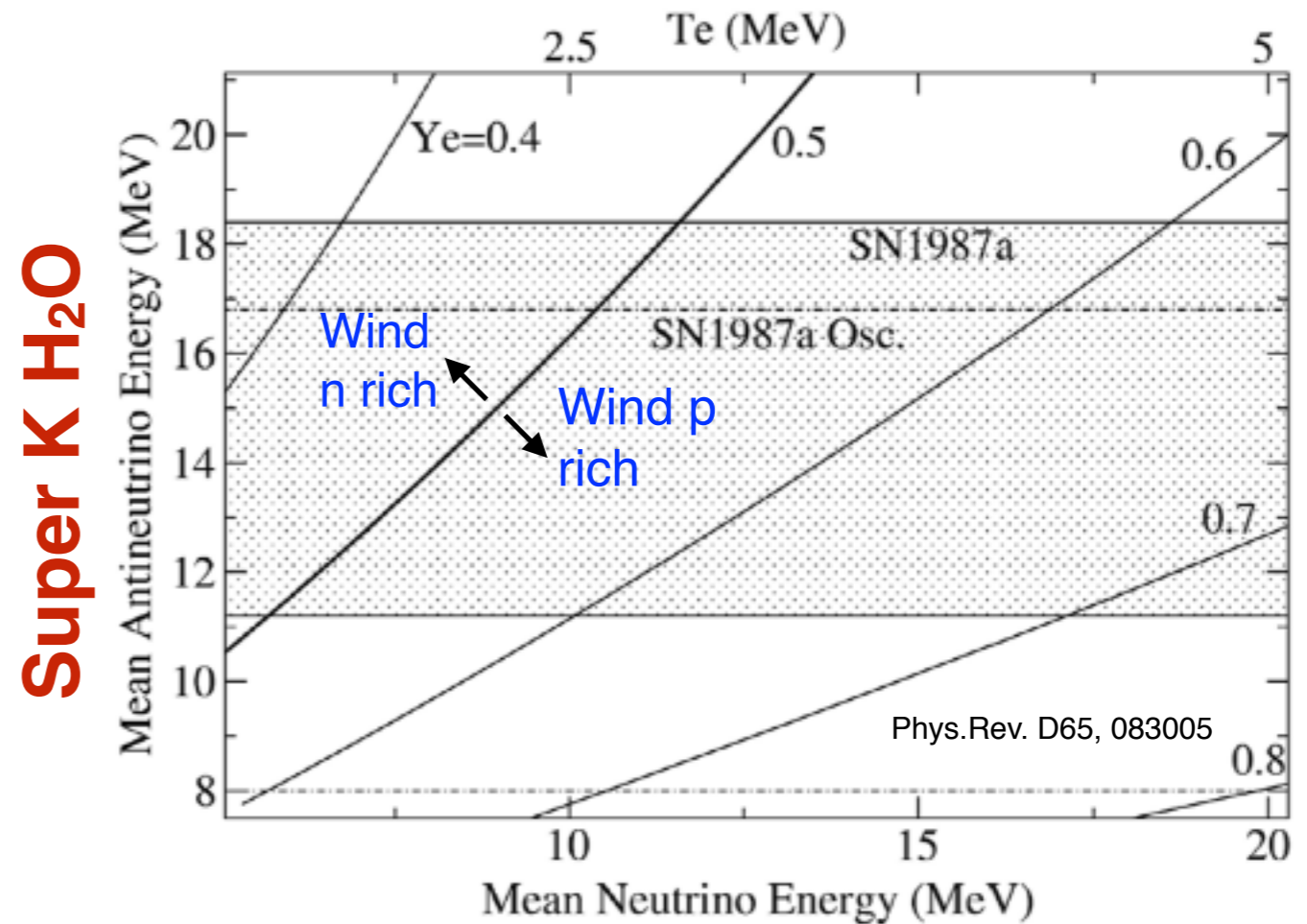
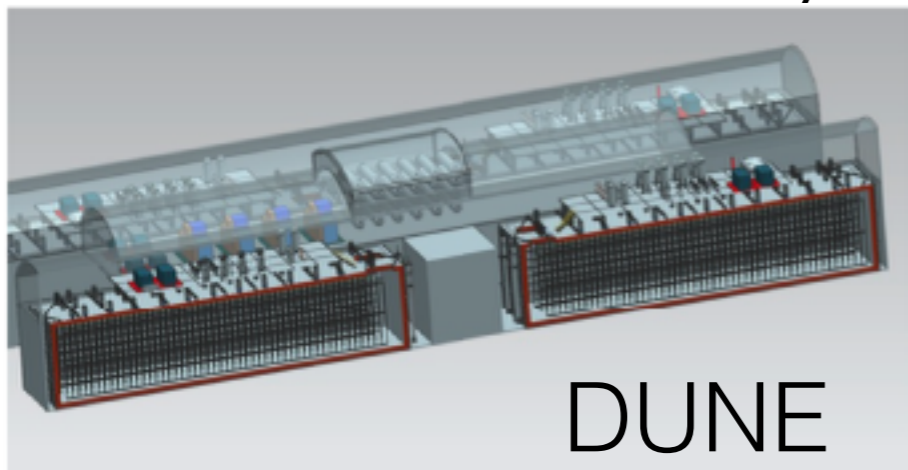
- The stronger the neutrino interactions, the longer a neutrino stays in thermal equilibrium with matter to lower densities and temperatures, and the lower is the emitted neutrino energy.
- Mu and tau neutrinos have only neutral current reactions (not enough energy to make muons) and so decouple at highest energies.
- Electron antineutrinos capture on protons while neutrinos capture on neutrons. Matter is neutron rich so neutrinos have large opacity. Therefore electron neutrinos are emitted with lowest energy.
- Expect order  $E(\mu, \tau) > E(\text{anti-}\nu_e) > E(\nu_e)$ .

# Deep Underground Neutrino Experiment

- Neutrino driven wind in a Supernova is an important nucleosynthesis site.
- Ratio of neutrons to protons in wind set by capture rates that depend on neutrino and antineutrino energies.



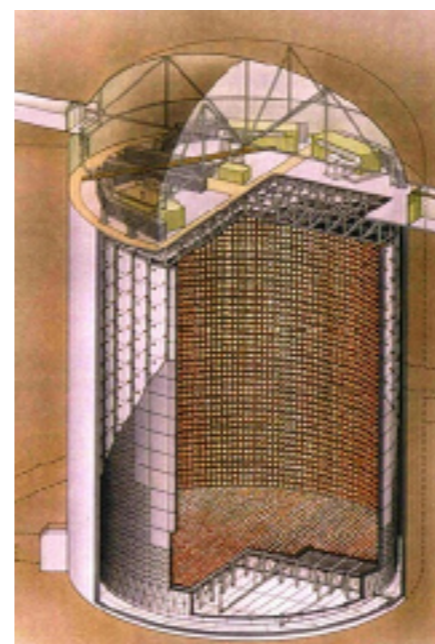
- Measure spectrum of both **antineutrinos** and **neutrinos** from next galactic SN.
- Super Kamiokande is large H<sub>2</sub>O detector good for **antineutrinos**.
- DUNE will be large liquid Ar detector that should measure **neutrino** spectrum well. If neutrinos are not much colder than antineutrinos then wind not very n rich.



## DUNE liquid Ar

Present SN simulations find too few neutrons for (main or 3rd peak: Au, U) r-process.

Likely that main r-process occurs in neutron star mergers. Expect gravitational wave observations.



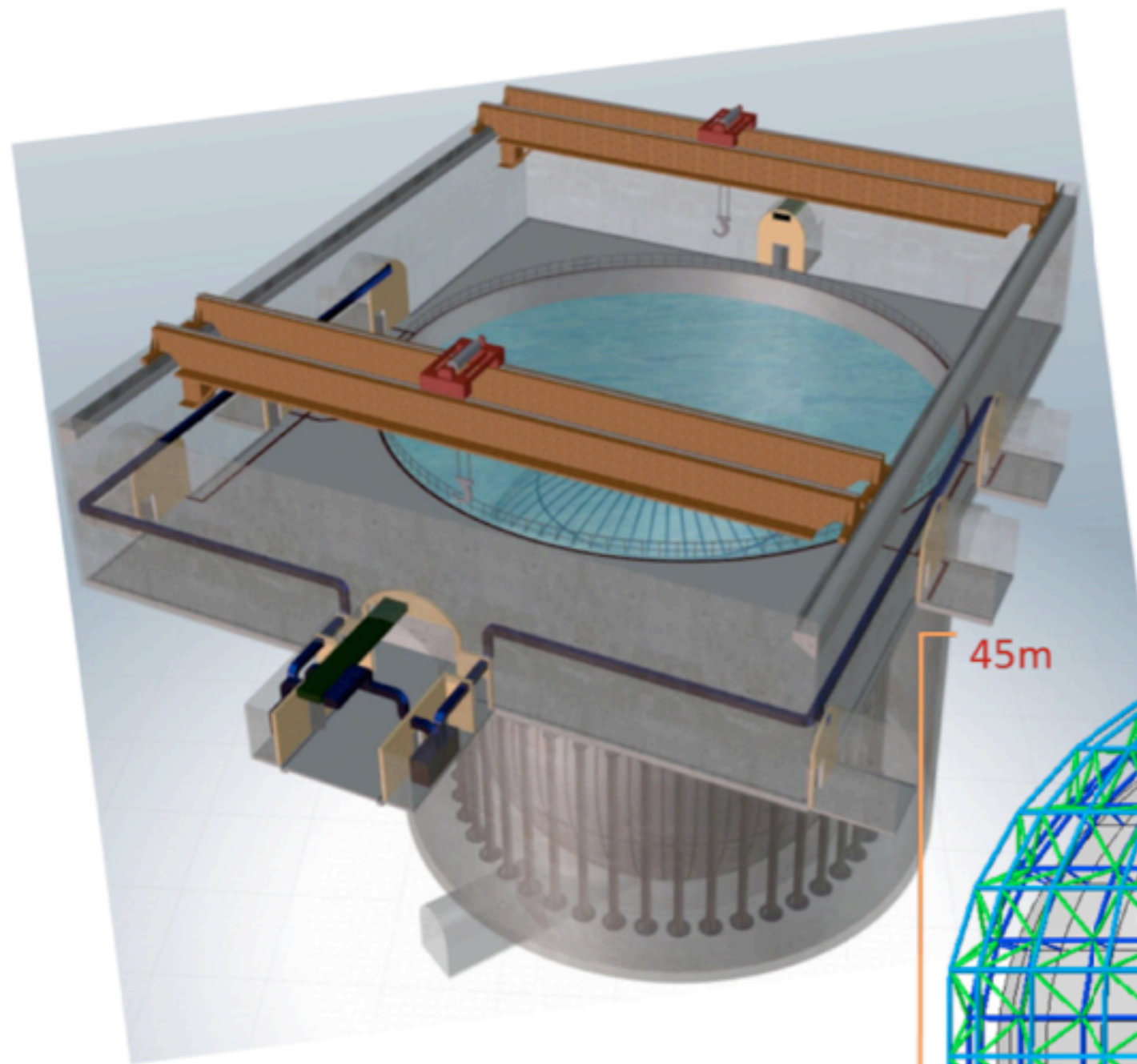
Super K

# *JUNO detector*

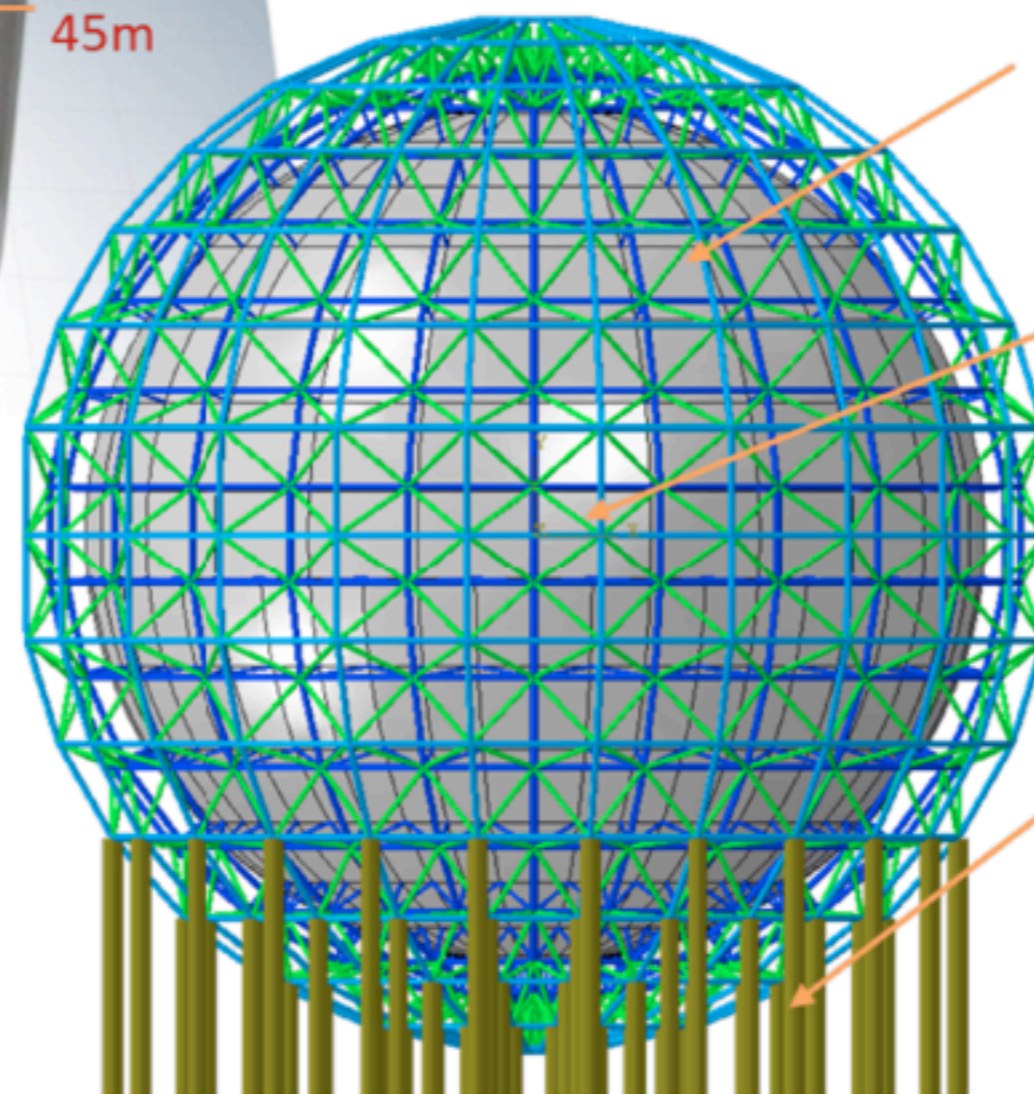
*35m diameter*

*20kt of liquid scintillator*

*15,000 PMTs*



45m



vessel made  
from acrylic

truss made from  
stainless steel

column made from  
stainless steel

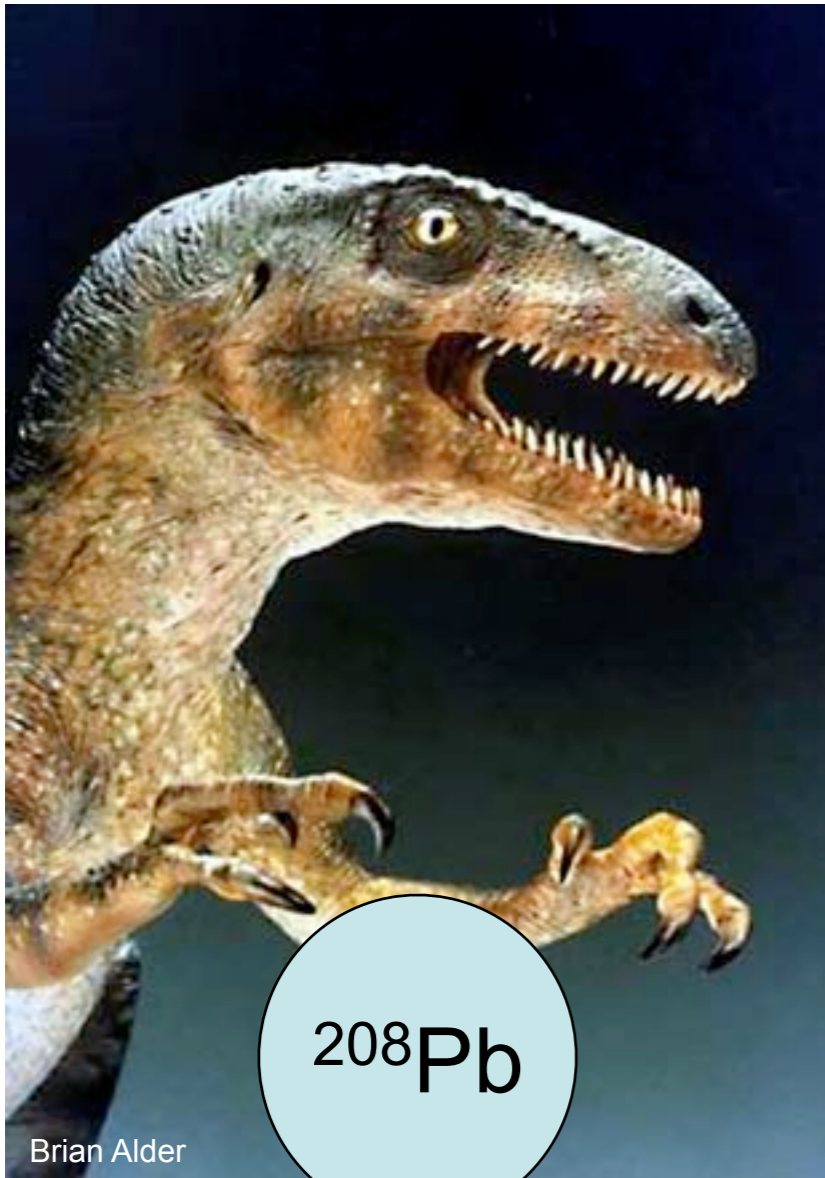
# Supernova Neutrinos

- Many thousands of events for a galactic SN.
- Electron antineutrinos in Super-K, JUNO, ...
- Electron neutrino spectra in DUNE (liquid Ar).
- Mu and tau neutrinos and antineutrinos via  $\nu_x + p \rightarrow \nu_x + p$  elastic scattering in JUNO.

# Site of main r-process

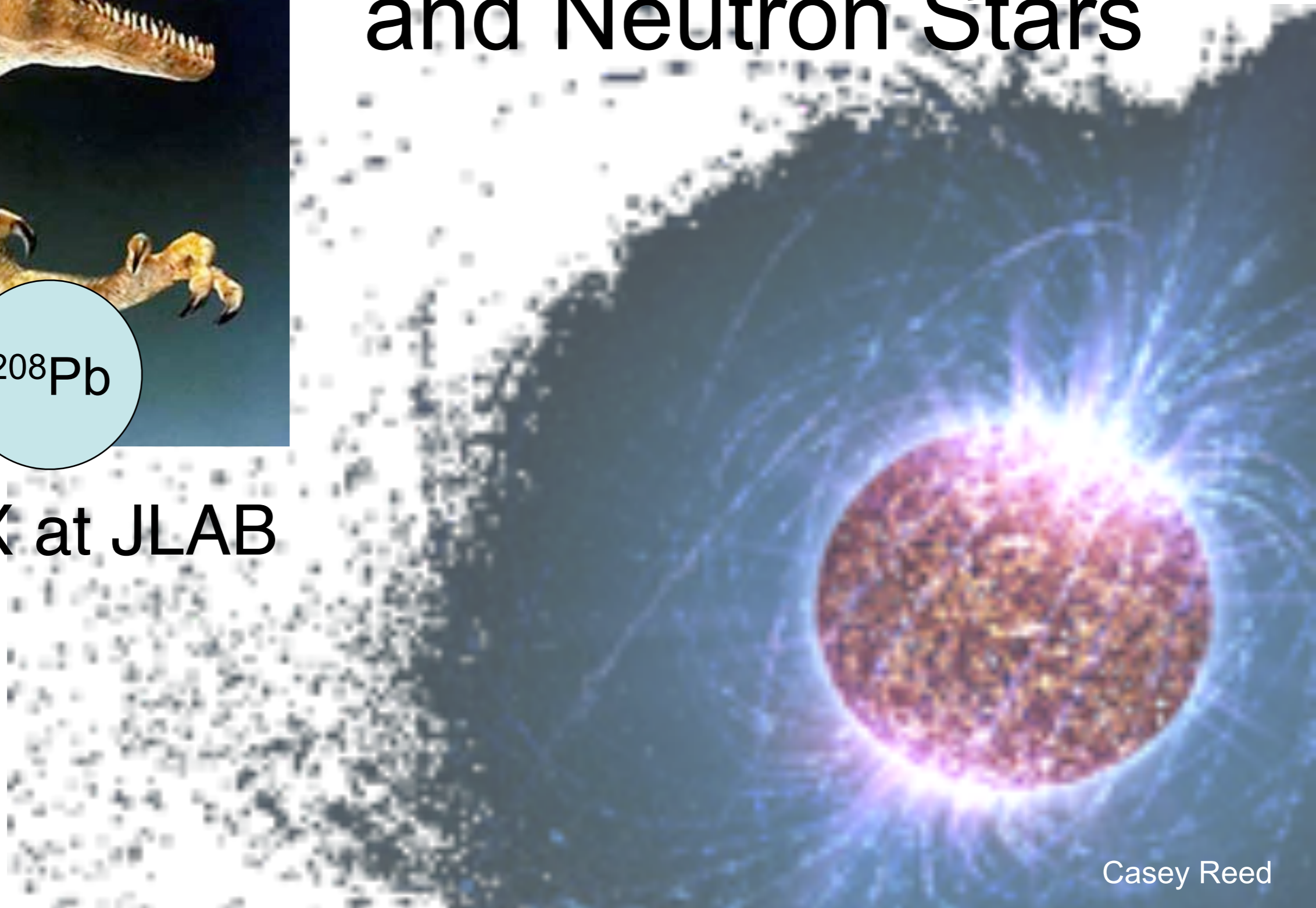
- Could be core collapse supernovae →  
Neutrino observations from next galactic SN will provide fundamental information.
- Could be neutron star mergers →  
Gravitational wave observations from LIGO are directly determining the merger rate.

# Laboratory Measurements and Neutron Stars



$^{208}\text{Pb}$

PREX at JLAB



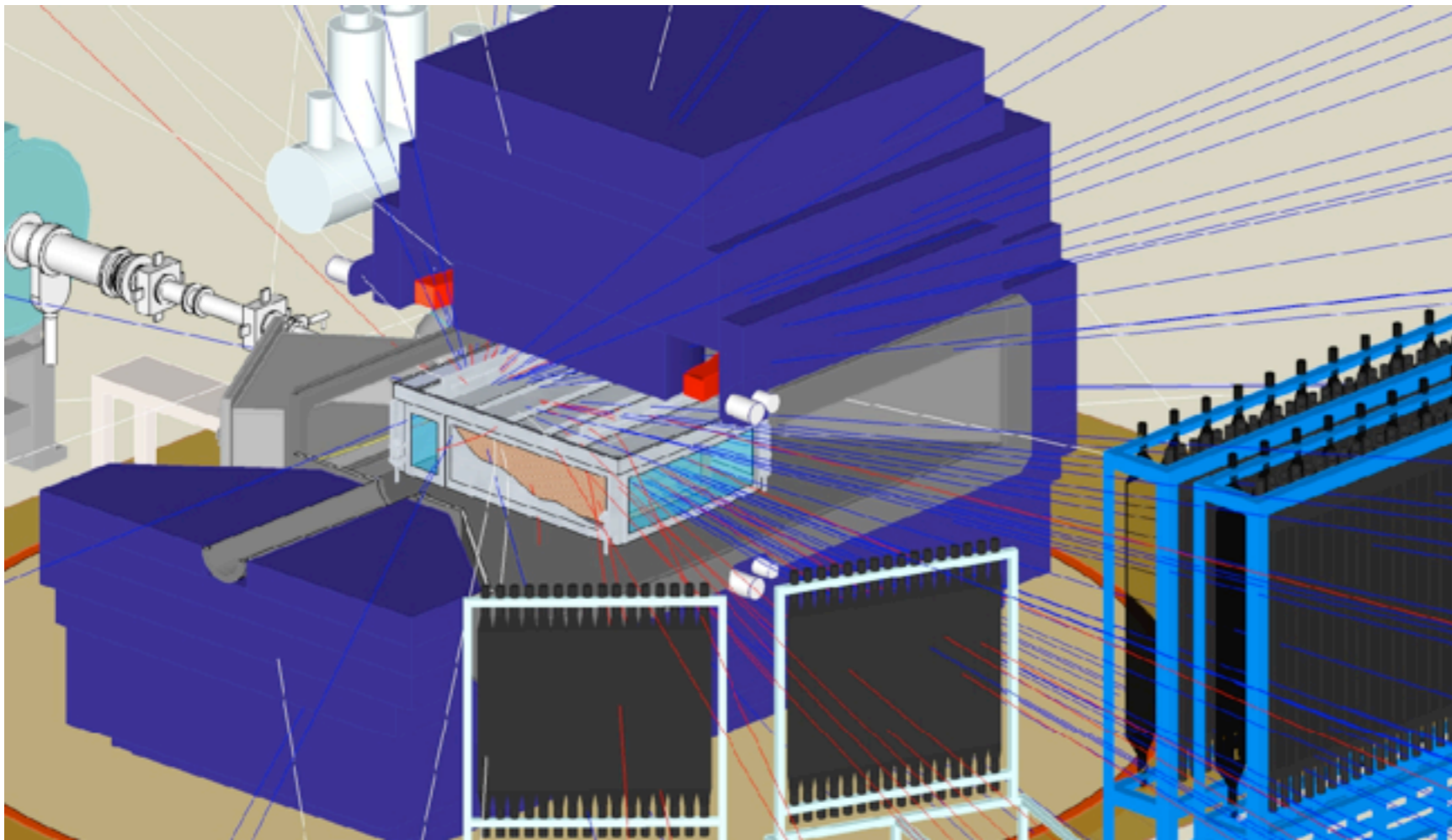
# Symmetry Energy $S(\rho)$

- Describes how energy of nuclear matter rises with increasing neutron excess.
- Important for extrapolating laboratory measurements to very neutron rich systems in astrophysics.
- *$S(\rho)$  at high densities ( $\rho > \rho_0$ ) is the single laboratory observable most closely related to the structure of neutron stars.*
- Heavy ion collisions, with radioactive beams, can produce high density n rich matter in the laboratory.



# Samurai TPC and $S(\rho)$ at $\rho > \rho_0$

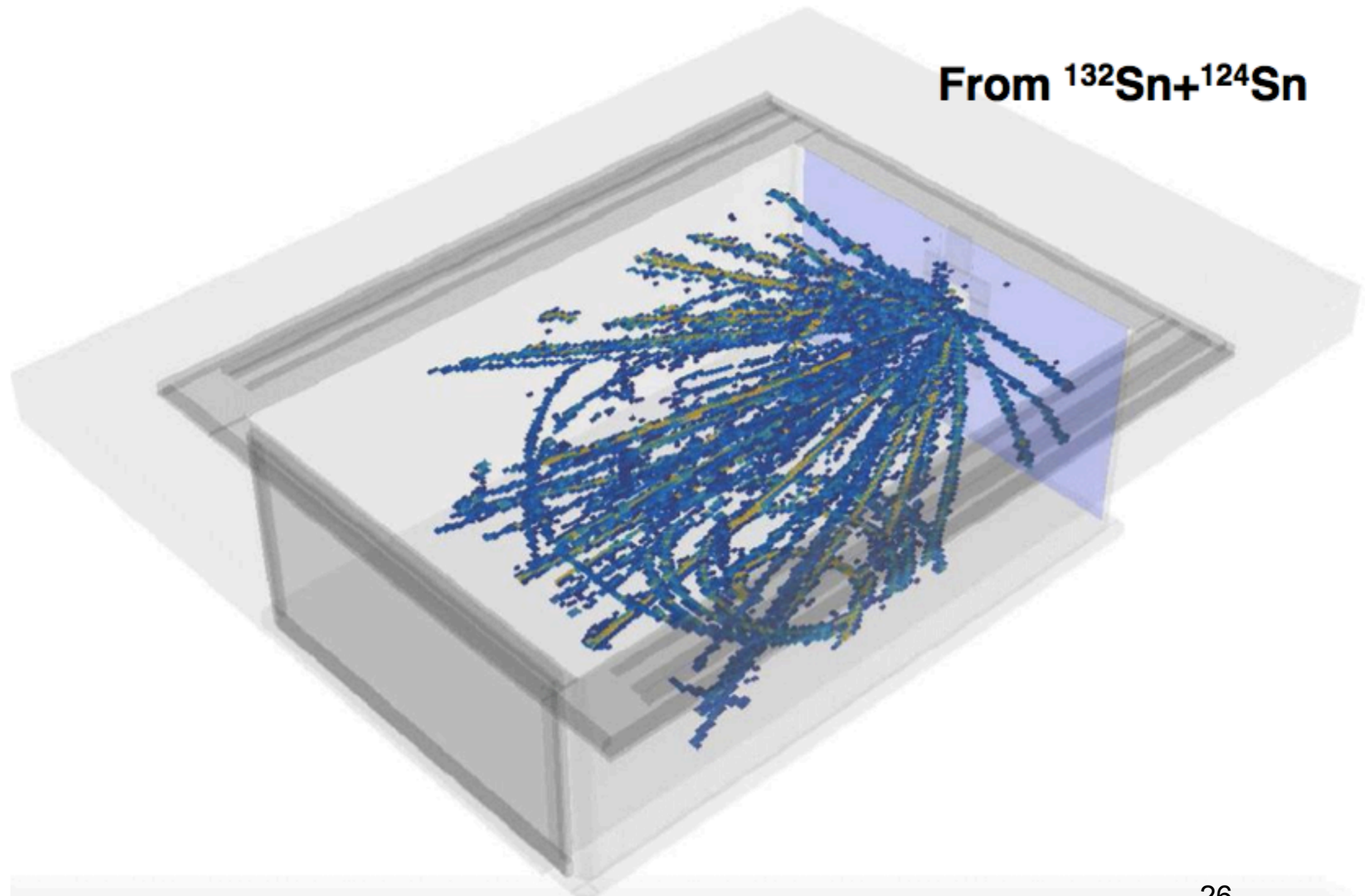
- Determining  $S(\rho)$  from HI exp. may depend on transport models. Look at pion production and  $\pi^+/\pi^-$  ratios, n/p flow...
- Experimental program underway at RIKEN RIBF using SAMURAI magnet and



time projection chamber (TPC).

Exp. with  $^{108}\text{Sn}$ ,  $^{132}\text{Sn}$  beams 2016

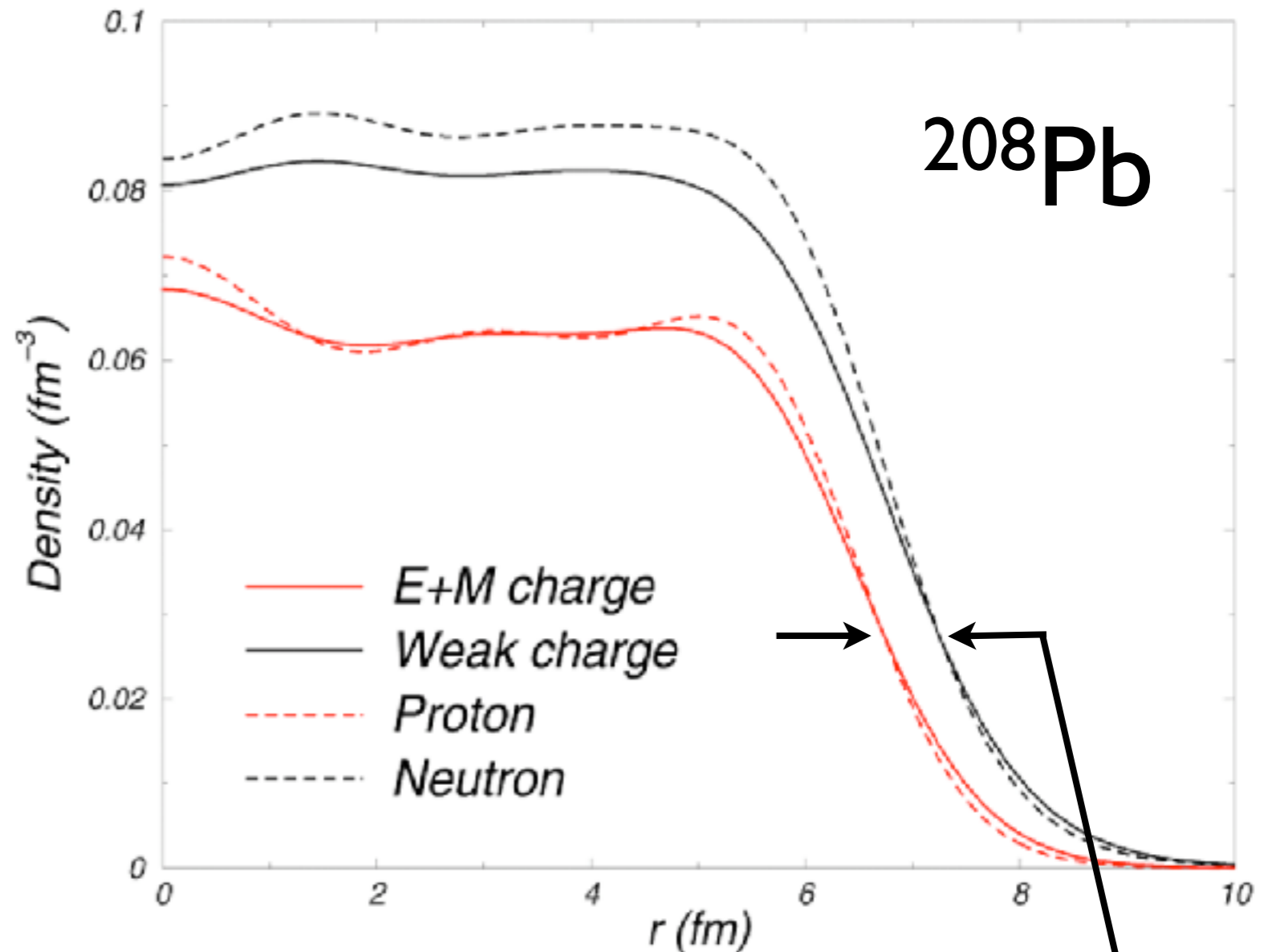
# First results in a year



Event from Tetsuya MURAKAMI talk

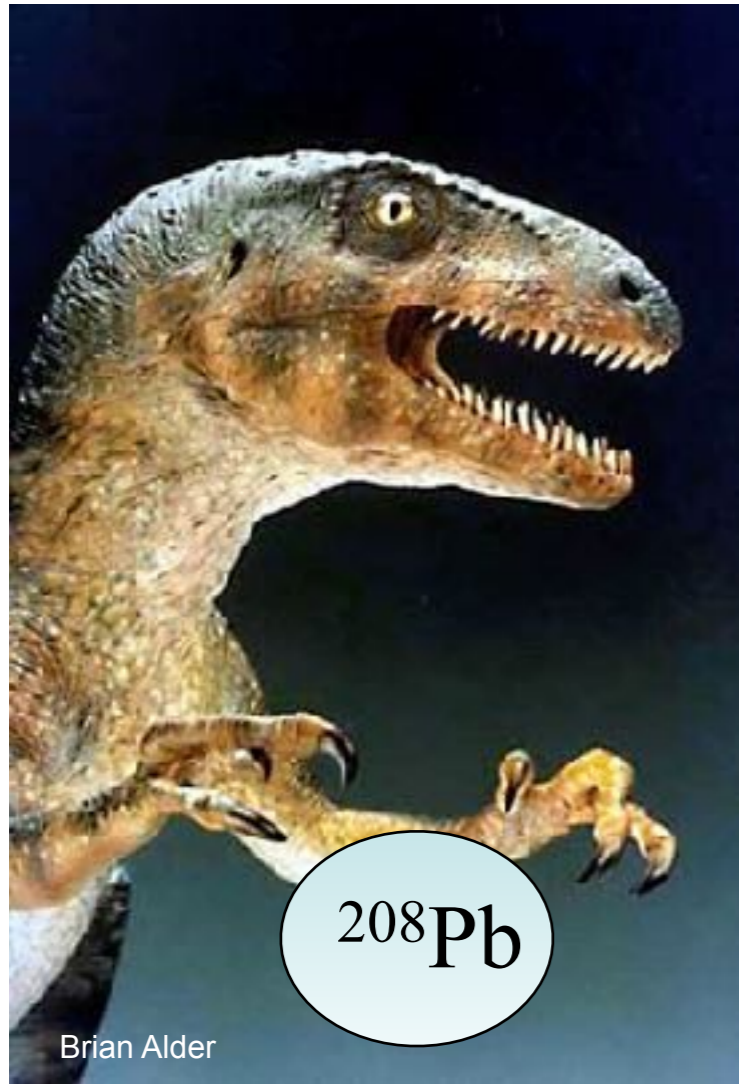
# Neutron skins and $dS/d\rho$

- Cleanest way to get  $dS/d\rho$  is to measure neutron skin thickness.
- $^{208}\text{Pb}$  has 44 more n than p. If extra n are in the center they will cost  $S(\rho)$  at relatively high  $\rho$ . But if extra n are in the surface they will only cost  $S(\rho)$  at low surface densities.
- The density dependence of  $S$  ( $dS/d\rho$ ) will push n out to the surface and give a thick n skin.



Measure how much neutrons stick out past protons

# Laboratory probe of neutron rich matter



**PREX** uses parity violating electron scattering to accurately measure the neutron radius of  $^{208}\text{Pb}$ .

This has important implications for neutron rich matter and astrophysics.

# Parity Violation Isolates Neutrons

- In Standard Model  $Z^0$  boson couples to the weak charge.

- Proton weak charge is small:

$$Q_W^p = 1 - 4\sin^2\Theta_W \approx 0.05$$

- Neutron weak charge is big:

$$Q_W^n = -1$$

- **Weak interactions, at low  $Q^2$ , probe neutrons.**

- Parity violating asymmetry  $A_{pv}$  is cross section difference for positive and negative helicity electrons

$$A_{pv} = \frac{d\sigma/d\Omega_+ - d\sigma/d\Omega_-}{d\sigma/d\Omega_+ + d\sigma/d\Omega_-}$$

- $A_{pv}$  from interference of photon and  $Z^0$  exchange. In Born approximation

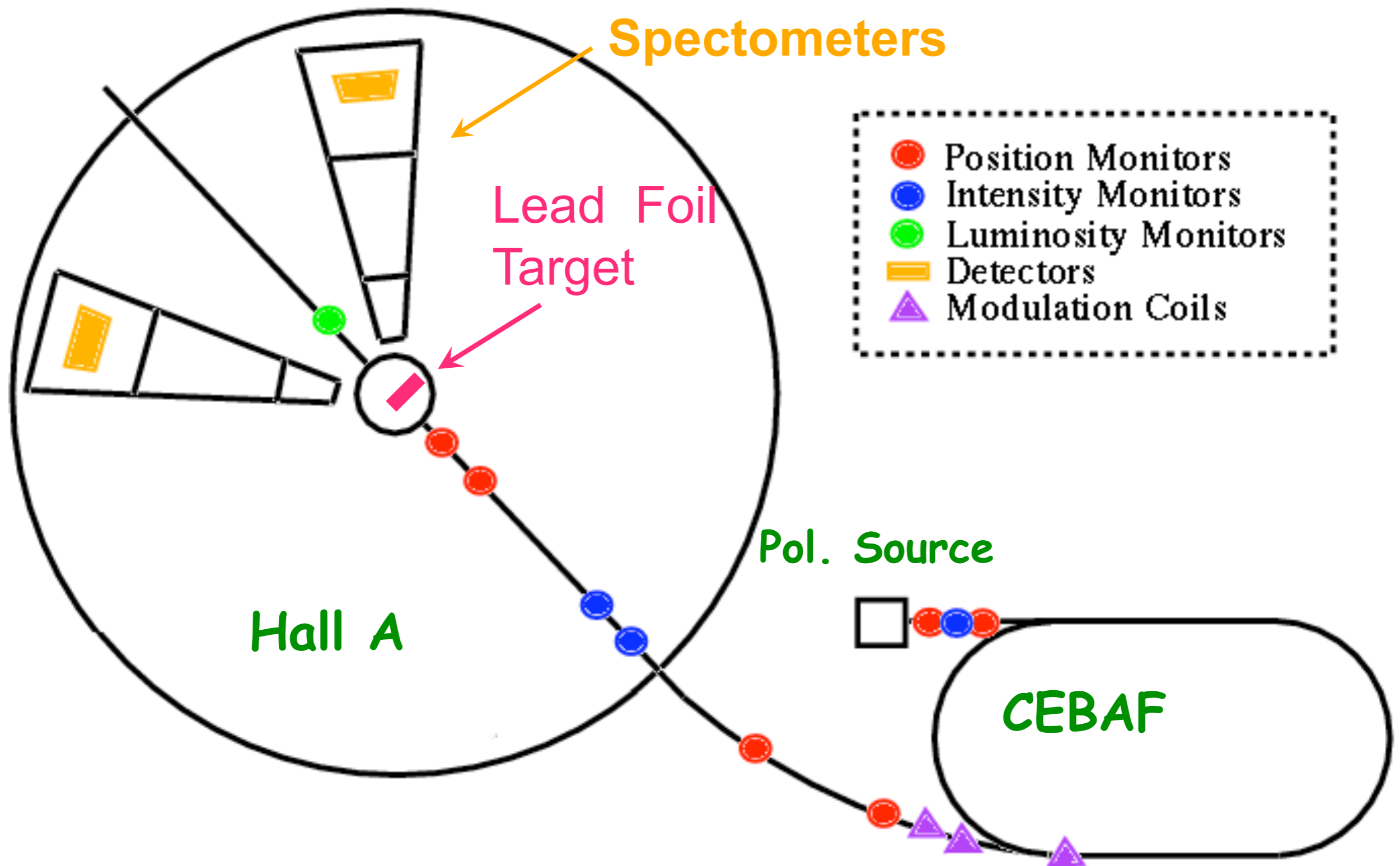
$$A_{pv} = \frac{G_F Q^2}{2\pi\alpha\sqrt{2}} \frac{F_W(Q^2)}{F_{ch}(Q^2)}$$

$$F_W(Q^2) = \int d^3r \frac{\sin(Qr)}{Qr} \rho_W(r)$$

- Model independently map out distribution of weak charge in a nucleus.
- **Electroweak reaction free from most strong interaction uncertainties.**

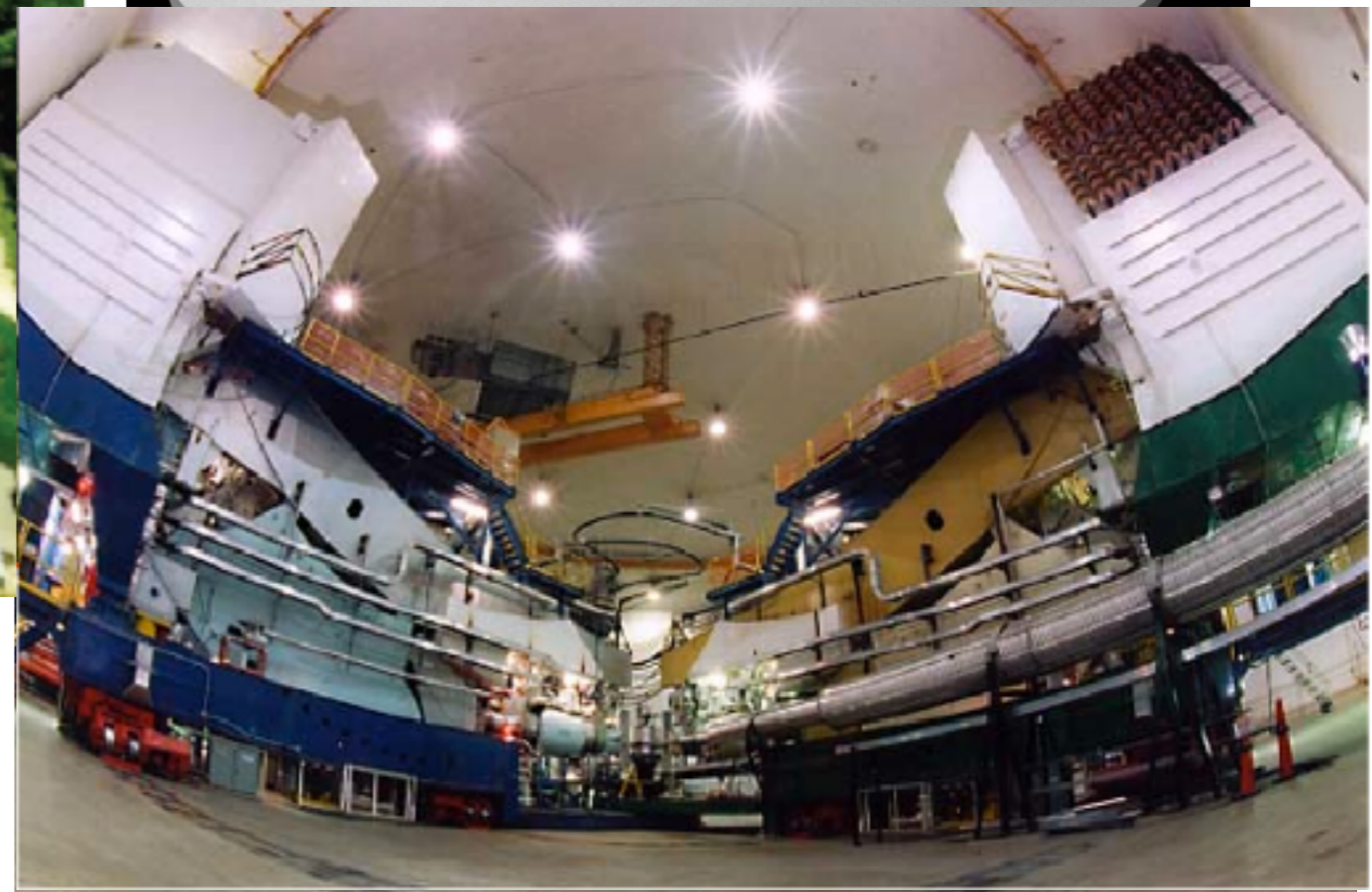
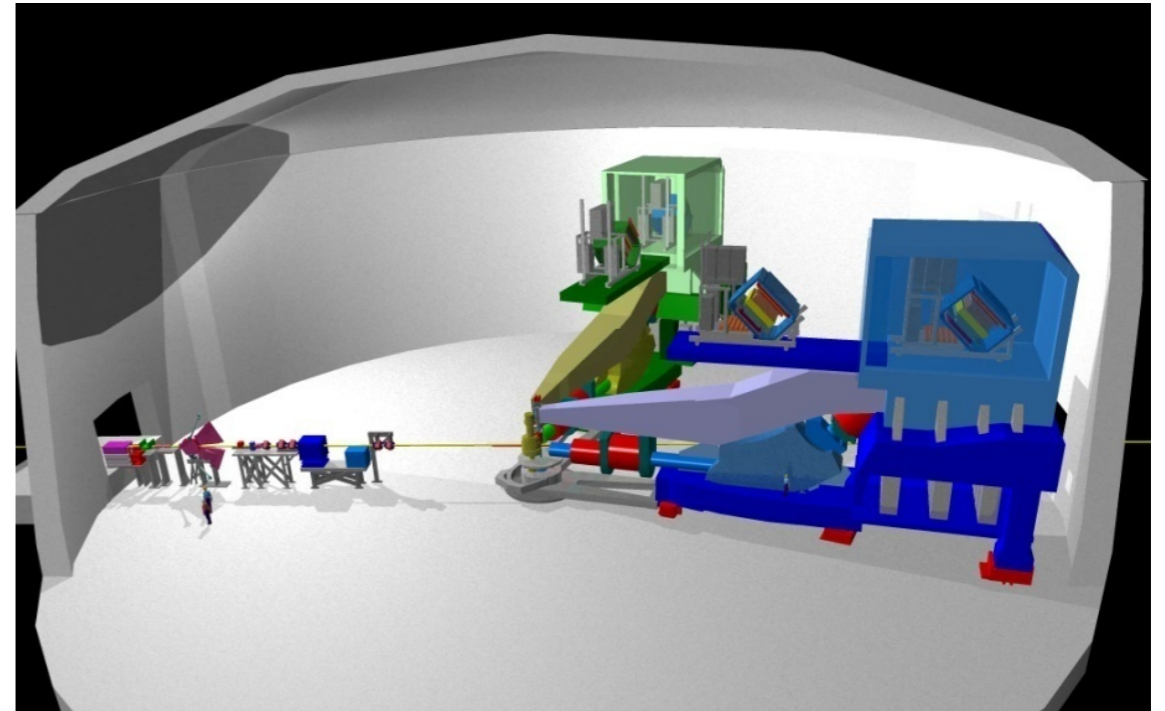
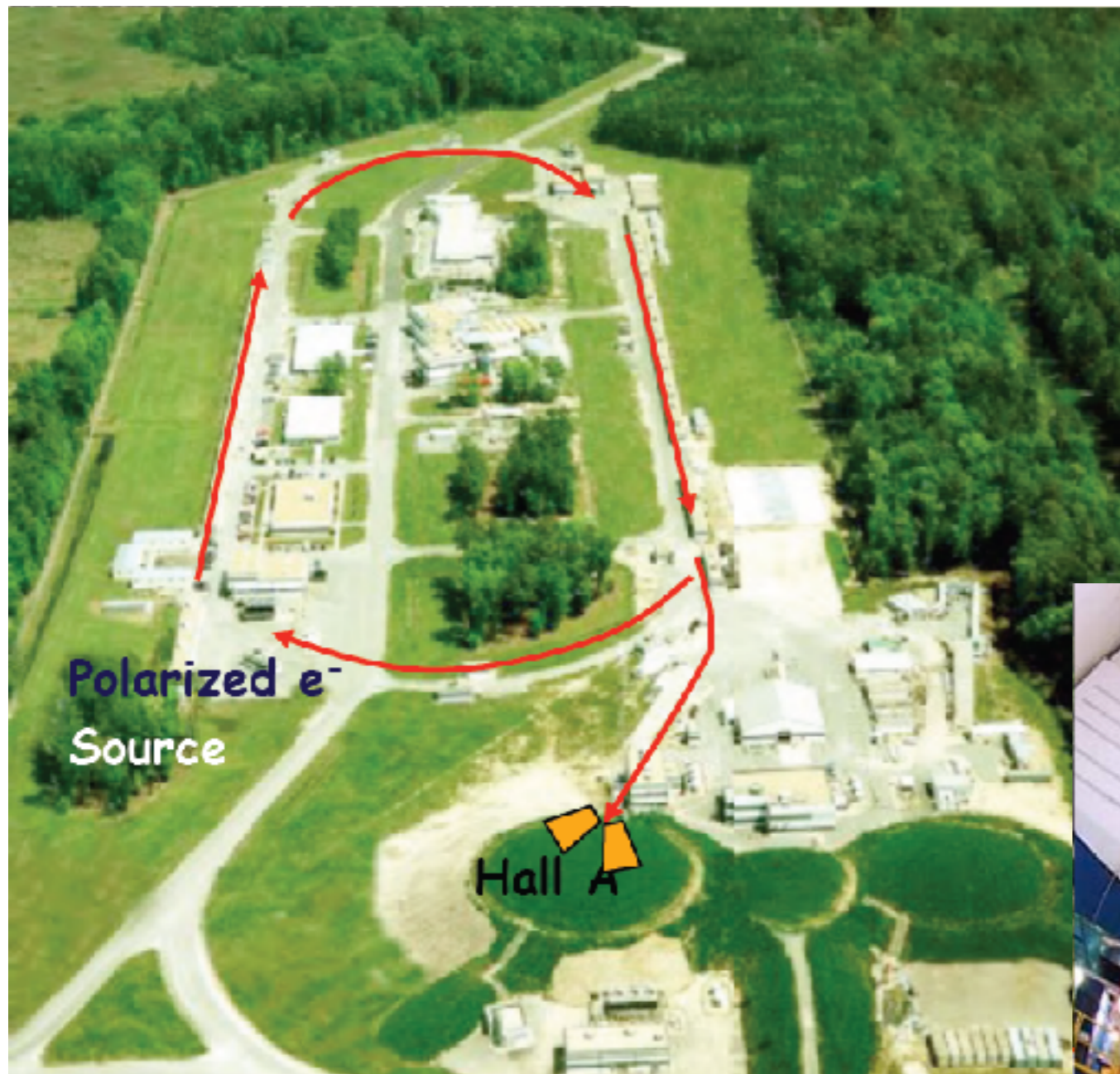
–Donnelly, Dubach, Sick

# PREX in Hall A at JLab



R. Michaels

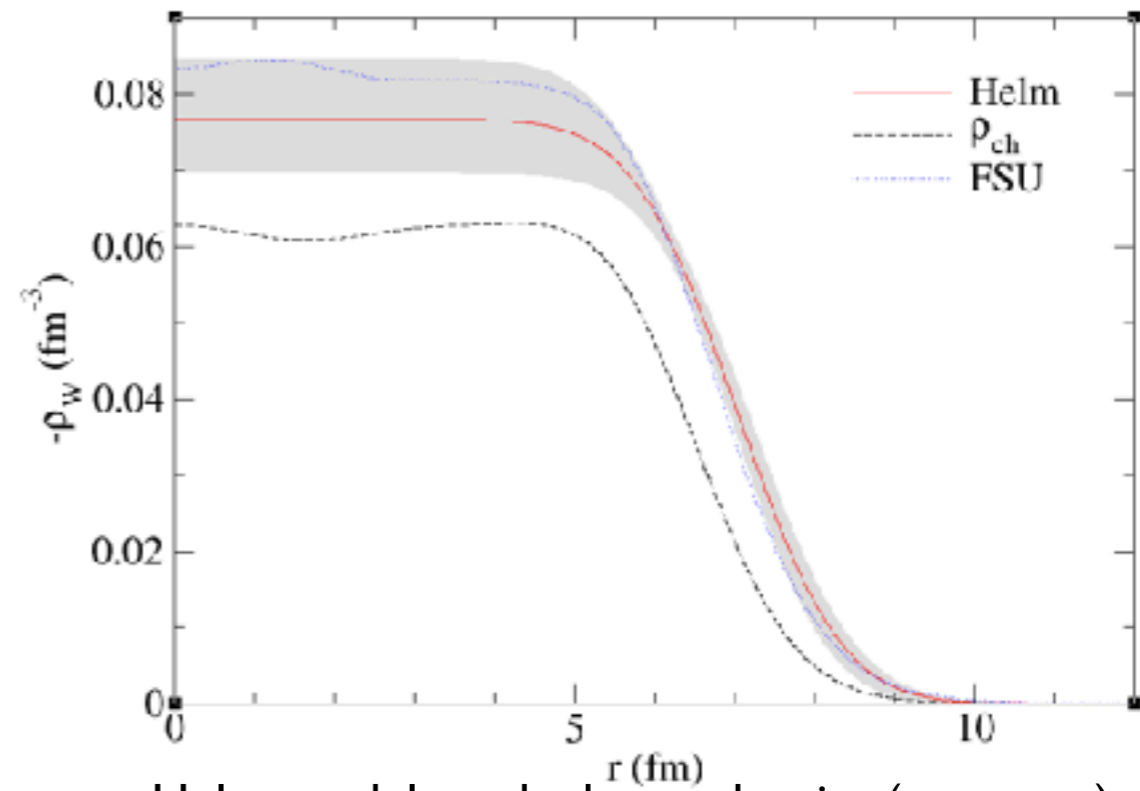
# Hall A at Jefferson Lab



R. Michaels, JLAB

# First PREX results

- 1.05 GeV electrons elastically scattering at  $\sim 5$  deg. from  $^{208}\text{Pb}$
- $A_{PV} = 0.657 \pm 0.060(\text{stat}) \pm 0.014(\text{sym})$  ppm
- Weak form factor at  $q=0.475 \text{ fm}^{-1}$ :  
 $F_W(q) = 0.204 \pm 0.028$
- Radius of weak charge distr.  
 $R_W = 5.83 \pm 0.18 \pm 0.03 \text{ fm}$
- Compare to charge radius  
 $R_{ch}=5.503 \text{ fm}$  --> Electroweak skin:  
 $R_W - R_{ch} = 0.32 \pm 0.18 \text{ fm}$
- First observation that weak charge density more extended than (E+M) charge density --> weak skin.
- Unfold nucleon ff--> neutron skin:  
 $R_n - R_p = 0.33^{+0.16}_{-0.18} \text{ fm}$
- Phys Rev Let. **108**, 112502 (2012),  
Phys. Rev. C **85**, 032501(R) (2012)



Helm model weak charge density (gray area) consistent with PREX results.

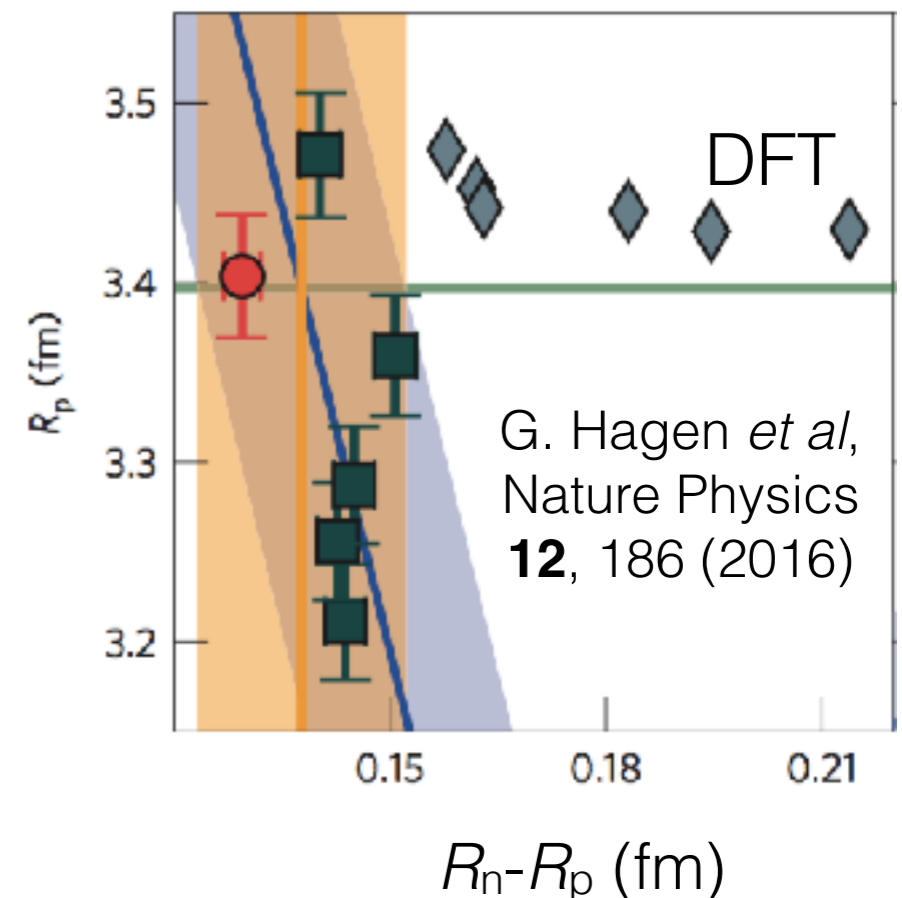
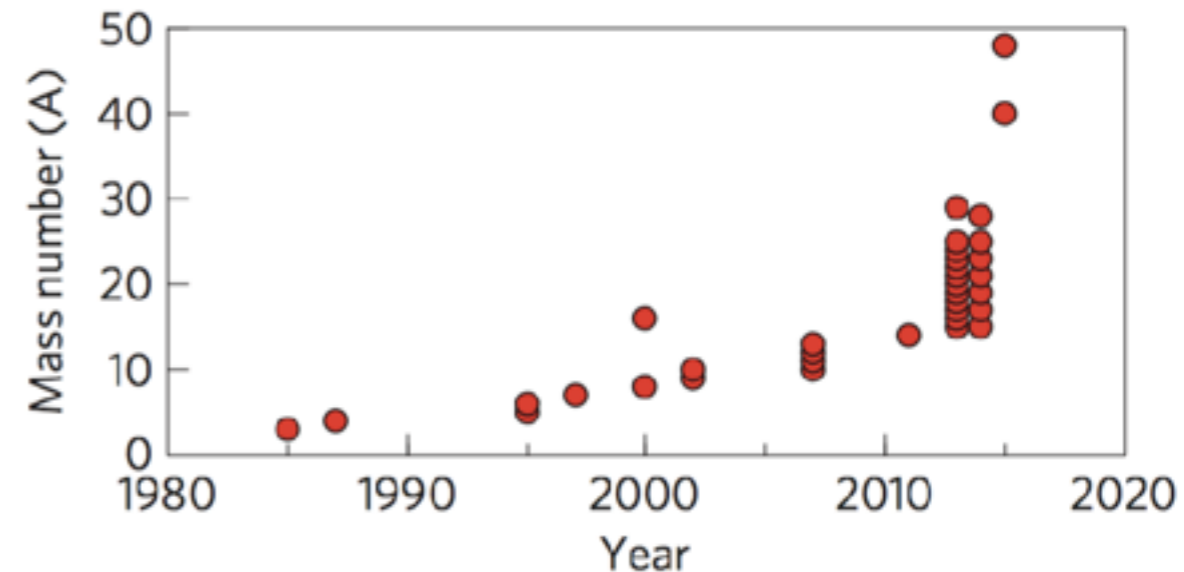
## Next Steps

- PREX-II:  $^{208}\text{Pb}$  with more statistics.  
Goal:  $R_n$  to  $\pm 0.06 \text{ fm}$ . Will large  $R_n - R_p$  be confirmed?
- CREX: Measure  $R_n$  of  $^{48}\text{Ca}$  to  $\pm 0.02 \text{ fm}$ .  
Microscopic calculations feasible for light n rich  $^{48}\text{Ca}$  (but not  $^{208}\text{Pb}$ ) to relate  $R_n$  to three neutron forces.



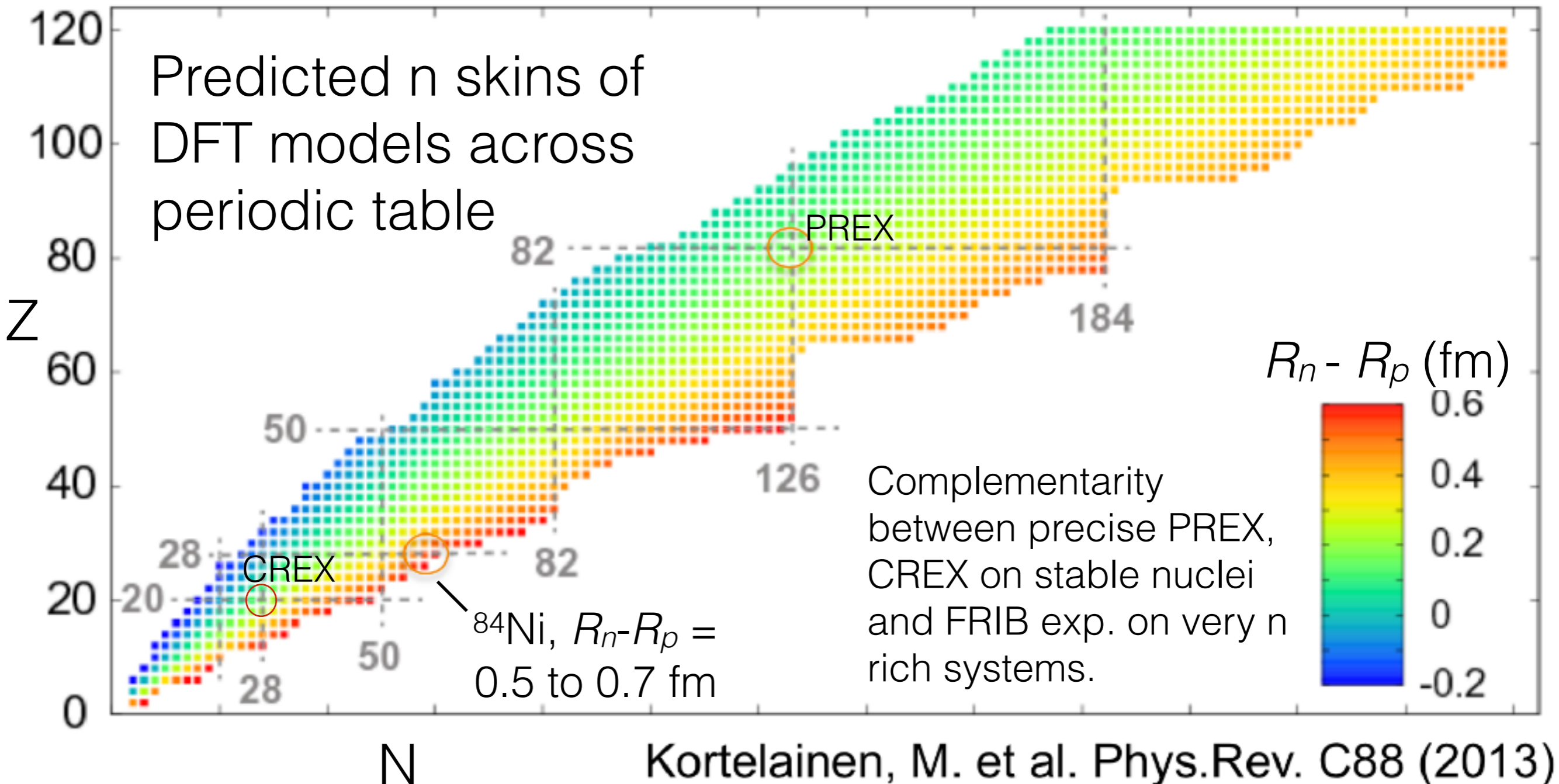
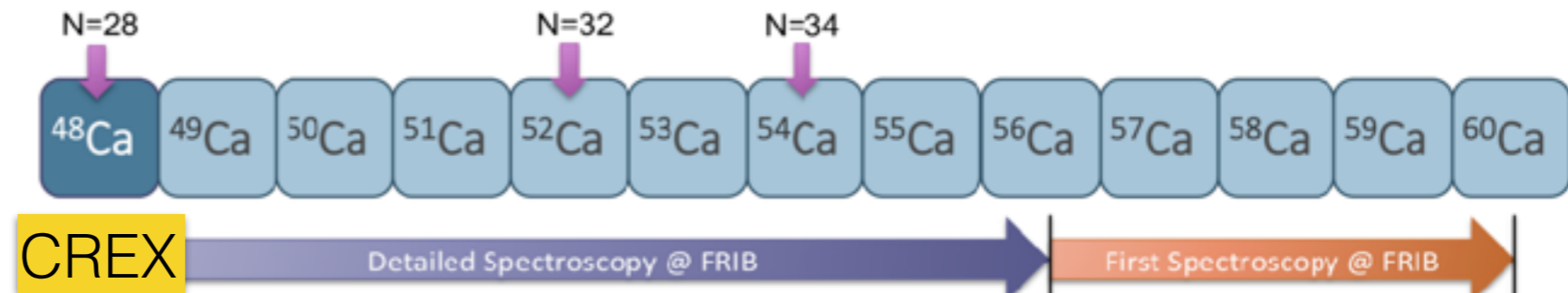
# Study 3 neutron forces in $^{48}\text{Ca}$

- Large computational advances allow microscopic calculations of structure of medium mass ( $A=48$ ) nuclei using realistic two nucleon and three nucleon forces from Chiral EFT.
- Coupled cluster calculations by G. Hagen *et al* make sharp prediction  $R_n - R_p(^{48}\text{Ca}) = 0.135 \pm 0.015$  fm.
- Three neutron forces play an important role. Many DFT models predict larger neutron skin.
- Prediction will be directly tested by CREX with goal of  $R_n$  to  $\pm 0.02$  fm.



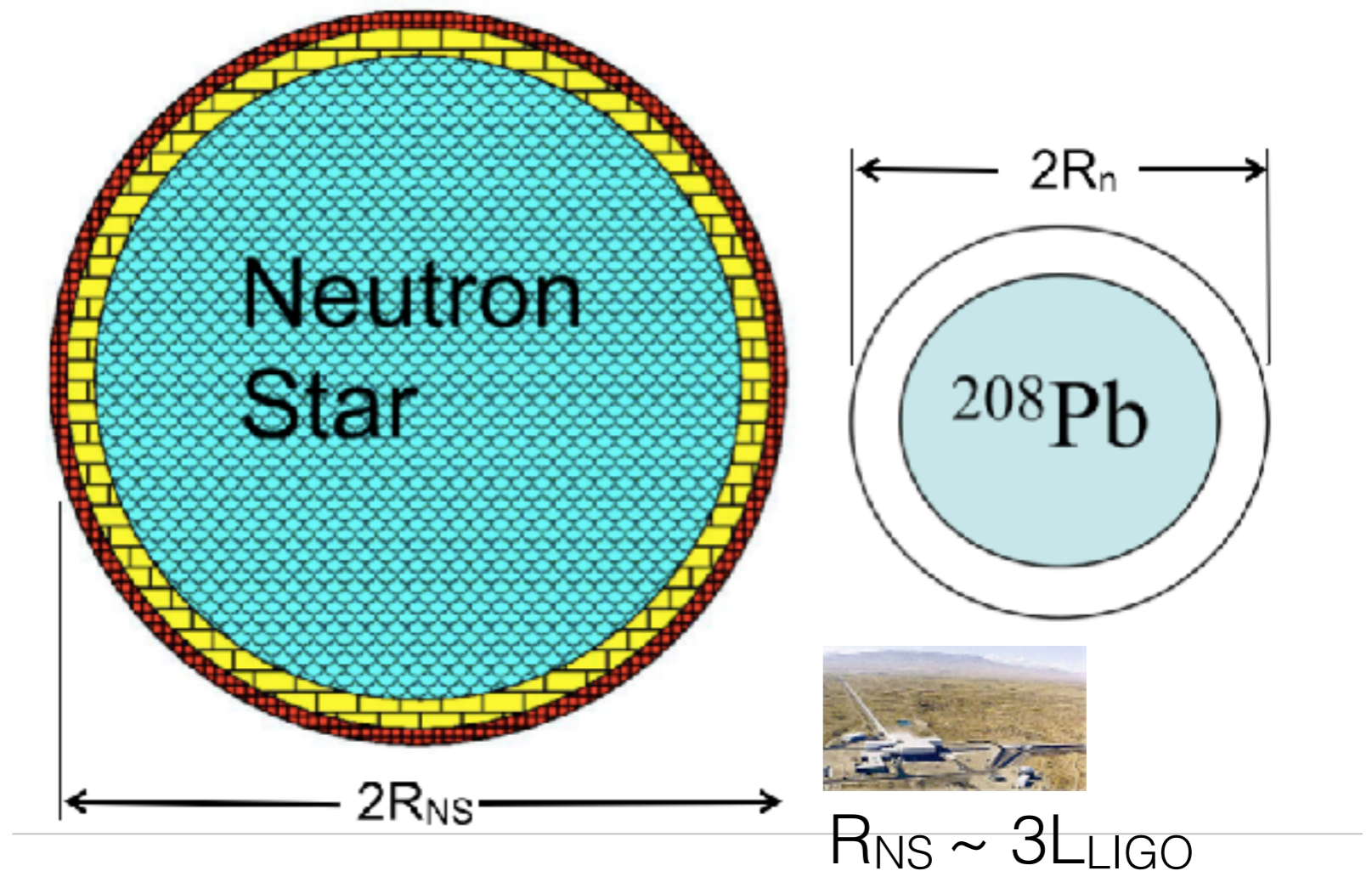
# Study more n rich nuclei at FRIB

- $^{40}\text{Ca}$  ( $Z=N=20$ ) is stable. FRIB can make  $^{60}\text{Ca}$  ( $N=40$ )



# Density Dependence of EOS

- Pressure of neutron matter pushes neutrons out against surface tension  $\implies R_n - R_p$  of  $^{208}\text{Pb}$  determines  $P$  at low densities near  $\rho_0$
- Radius of ( $\sim 1.4M_{\text{sun}}$ ) NS depends on  $P$  at medium densities  $> \rho_0$ .
- Maximum mass of NS depends on  $P$  at high densities.
- These three measurements constrain density dependence of EOS.



Neutron star is 18 orders of magnitude larger than Pb nucleus but has same neutrons, strong interactions, and equation of state.

PREX II:  $R_n(^{208}\text{Pb})$  to  $\pm 0.06$  fm

CREX:  $R_n(^{48}\text{Ca})$  to  $\pm 0.02$  fm or  $\sim 5\Delta L_{LIGO}$

# Laboratory and Astronomical Observations of n Rich Matter

- Parity violating PREX and CREX measure neutron skin of  $^{208}\text{Pb}$ ,  $^{48}\text{Ca}$  and constrain pressure of neutron rich matter near nuclear density, density dependence of the symmetry energy, and three neutron forces.
- Neutron star radius constrains pressure of neutron rich matter at medium densities  $\sim 2$  nuclear density.
- Maximum mass of a neutron star constrains pressure at high densities.
- Supported in part by DOE

# TALENT Classes

- Training in Advanced Low Energy Nuclear Theory (TALENT) teaches three week advanced classes in nuclear theory / astrophysics in North America and Europe.
- Example in 2015 “Nuclear and neutrino physics and astrophysics of neutron stars and supernovae” was taught at the INT in Seattle.
- Check web site for this years classes.
- Should TALENT be expanded to teach classes in China?