# Robust Collider Limits on Heavy-Mediator DM <br> Andrea Wulzer <br> with D.Racco and F.Zwirner 



## Introduction

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Model-independence, i.e. broad exploration of the parameter space is mandatory here!

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$\left\{\begin{array}{l}\text { 1) thermal relic calculation } \\ \text { 2) direct search limits } \\ \text { 3) indirect search limits }\end{array}\right.$

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what instead about ...
4) collider limits ??

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## "The cutoff is physical !!"

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Cutoff is part of the EFT definition, one of its free parameters. In any specific microscopic model, we might read its true value

$$
M_{\mathrm{cut}} \sim M_{\mathrm{Med}}
$$

mass of the specific "mediators", or scale of strong UV theory

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LHC might carry us above the cutoff:

however restricting the signal to the predictable region sets lower bound on the "true" signal, which holds for any mediator model

$$
\left.\sigma_{E F T}^{S}\right|_{E_{\mathrm{cm}}<M_{\mathrm{cut}}} \leq \sigma_{\text {true }}^{S}<\sigma_{\mathrm{exc}}
$$

compared with exclusion upper bound, model indep. limit is set

## ATLAS mono-jet recast

chosen operator: $\quad \mathcal{L}_{\text {int }}=-\frac{1}{M_{*}^{2}}\left(\bar{X} \gamma^{\mu} \gamma^{5} X\right)\left(\sum_{q} \bar{q} \gamma_{\mu} \gamma^{5} q\right)$
counting in four SR

| signal region | SR1 | SR2 | SR3 | SR4 |
| :--- | :---: | :---: | :---: | :---: |
| $p_{T}^{\text {jet }}$ and MET | $>120$ | $>220$ | $>350$ | $>500$ |
| $\sigma_{\mathrm{exc}}[\mathrm{pb}]$ | 2.7 | 0.15 | $4.810^{-2}$ | $1.510^{-2}$ |

restricted signal definition:

$$
\sigma_{\mathrm{SR} i}\left(M_{*}, m_{D M}, M_{\mathrm{cut}}\right)=\sigma\left(M_{*}, m_{D M}, M_{\mathrm{cut}}\right) \times A_{i}\left(m_{D M}, M_{\mathrm{cut}}\right) \times \epsilon
$$

NOTE: the EFT has three parameters

1) $m_{\mathrm{DM}}$
2) $M_{*}$
3) $M_{\text {cut }}$ (as physical as the other two)

## ATLAS mono-jet recast

colored lines: fixed $M_{\text {cut }}$



Hard signal regions are favored at high cutoff (naive EFT)
But rapidly lose sensitivity: the cut makes distributions softer

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Theoretical connection among $M_{*}$ and $M_{\text {cut }}$ :

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Two justifications:

1) from examples: $\frac{1}{M_{*}^{2}}=\frac{g_{*}^{2}}{M_{\text {med }}^{2}}$

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Useful redefinition: $\left\{\begin{array}{lc}\text { We know for sure that: } & g_{*}<4 \pi \\ \text { Expected for a WIMP: } & g_{*} \sim 1\end{array}\right.$

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Fixed $g_{*}$ limits: (from all the SR )


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Similar plot in De Simone et. al 1402.1275. Comparison in backup.

## Conclusions and Outlook

- Model-independent test of H-M DM is possible
- Parameter space currently far from fully tested progress needed in the soft region


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1) define "hard scale" to be cut on. Using MLM matching 2) find optimal statistics for limits (shape an. with $>0$ th. errors?)

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- Model-independent test of H-M DM is possible
- Parameter space currently far from fully tested progress needed in the soft region
- Towards the concrete implementation of the method [with M.Zanetti (CMS) and F.Pobbe]

1) define "hard scale" to be cut on. Using MLM matching
2) find optimal statistics for limits (shape an. with $>0$ th. errors?)
-Beyond EFT's, the improvement is from mediator prod.:
3) turn to mediator search, appropriate interpretation is $\sigma \times B R$
4) other search channels for the mediator (e.g., model $B$ is squark )
5) final goal is cover all models by patches (EFT + mediator search)

## Backup

From De Simone et. al 1402.1275:


Conceptual difference: (to me...)
Their aim was show up to when naive EFT limits coincide with UV theory ones. Our aim is set limits that hold for any UV.
Practical differences:
Model-dependent cut variable $Q_{\mathrm{tr}}$.
The contours are open! Issue due to Naive EFT limit rescaling.

## Other variables

By further specifying mediator dynamics (s- or t-channel)
$Q_{\mathrm{tr}}=\max$ virtuality of mediator propagator

Model A: Z' coup. to q and DM


In all cases (kinematical bound):

Model B: squark-DM-quark coup.

$Q_{\mathrm{tr}}=\ldots$
$Q_{\mathrm{tr}}<E_{\mathrm{cm}}$

## Other variables

Worth dedicated s-and t-channel analyses for a better bound?



We consider the improvement not sufficient

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## Simplified models reinterpretation

Properly set EFT limits hold in any microscopic theory. They are correct, but conservative.

Model A: Z' coup. to q and DM

$M_{*}=\frac{m_{Z^{\prime}}}{\sqrt{g_{q} g_{X}}}$

$g_{*}=\sqrt{g_{q} g_{X}}$

Model B: squark-DM-quark coup.

$M_{*}=\frac{2 \widetilde{m}}{g_{\mathrm{DM}}}$
$g_{*}=\frac{g_{\mathrm{DM}}}{2}$

Compute parameters, use EFT limits, obtain bounds.
Compare with direct recasting of mono-jet.

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Model A: 95\% CL limit on $M_{*}$


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Lines for $\Gamma_{Z^{\prime}} / m_{Z^{\prime}}=1 / 8 \pi$ and $1 / 3$
Caution remark: most of these lines are inconsistent!

$$
\frac{\Gamma_{Z^{\prime}}}{m_{Z^{\prime}}}=\alpha g_{q}^{2}+\beta g_{X}^{2} \geq g_{q} g_{X} \sqrt{4 \alpha \beta}=\frac{m_{Z^{\prime}}^{2}}{M_{*}^{2}} \sqrt{4 \alpha \beta}
$$

## Simplified models reinterpretation



Lines for $\Gamma_{Z^{\prime}} / m_{Z^{\prime}}=1 / 8 \pi$ and $1 / 3$
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## Simplified models reinterpretation



Lines for $\Gamma_{Z^{\prime}} / m_{Z^{\prime}}=1 / 8 \pi$ and $1 / 3$
Notice: improvement due to resonant mediator production

