

# Cross Sections and Spin Observables for Forward Jet Production

A talk primarily about...

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Physics Letters B

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Cross sections and transverse single-spin asymmetries in forward jet production from proton collisions at  $\sqrt{s} = 500$  GeV

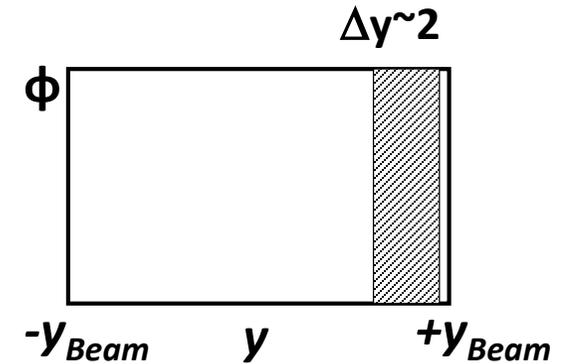
A<sub>N</sub>DY Collaboration \*

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Workshop on Jets and Heavy Flavor  
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# Forward Particle Production

- In this talk, forward means when the observed particle Feynman-x ( $x_F=2p_z/\sqrt{s}$ ) scaling variable is larger than 0.1
- In general, sufficient  $p_T$  is required so that pQCD is applicable. Consequently, forward is further defined to require sufficient  $p_T$  [which looks to be  $\sim 2$  GeV/c for inclusive  $\pi^0$  production]
- RHIC interaction regions have uniquely large length for a collider, when scaled by  $\sqrt{s}$ . This interaction length does permit space for forward instrumentation



	Free Space (m)	$\sqrt{s}$ (GeV)	Ratio ( $L/\sqrt{s}$ )
Tevatron	13	1600	0.0081
LHC	38	13000	0.0029
RHIC	16	500	0.032
	16	200	0.080

Consider the separation in x-y plane ( $d_{\gamma\gamma}$ ) of a pair of photons from the decay  $M \rightarrow \gamma\gamma$ , when the plane is  $L$  from where  $M$  (mass  $m_M$ ) is produced:

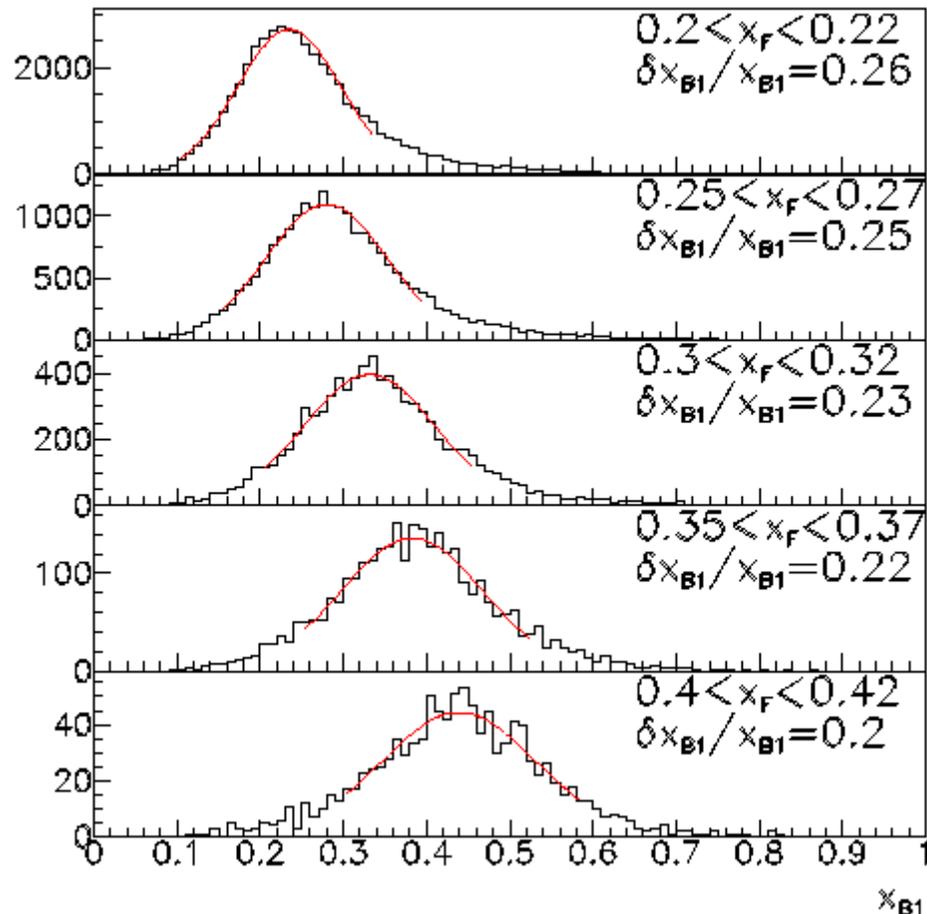
$$d_{\gamma\gamma}^{min} = \frac{L}{\sqrt{s}} \frac{4m_M}{x_F}$$

$\Rightarrow$  Large  $L/\sqrt{s}$  enables reconstruction of light mesons to large  $x_F$  at large  $\sqrt{s}$

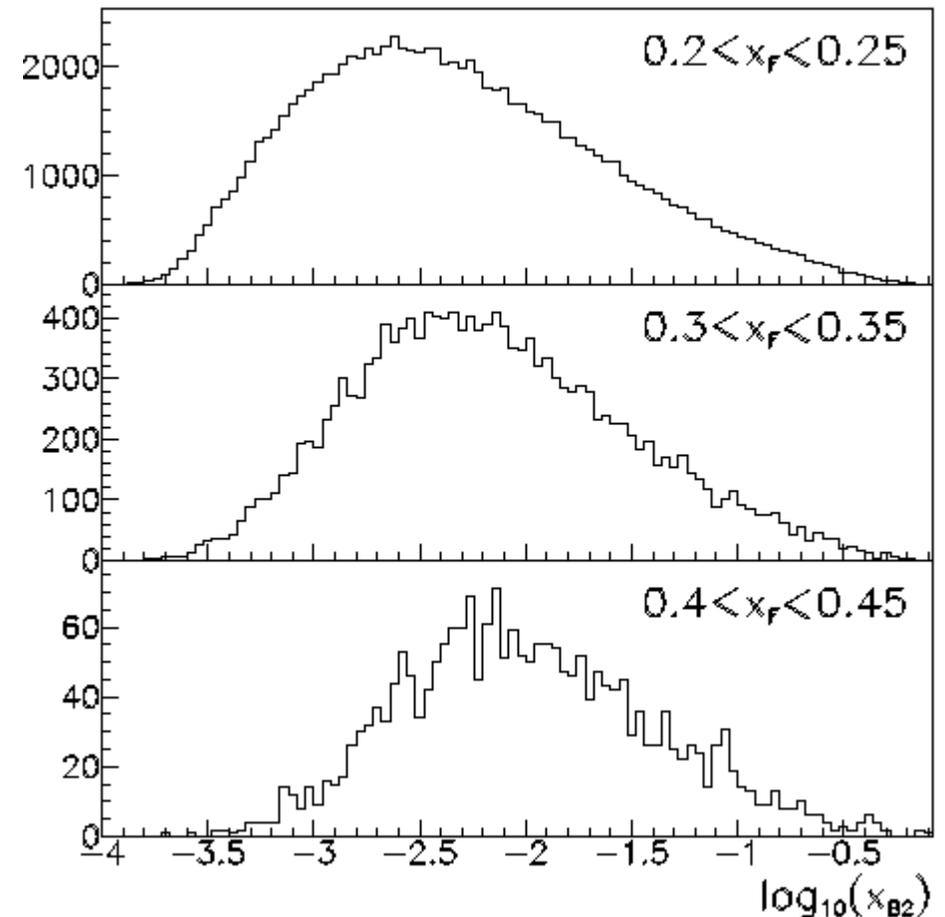
# Why is large $x_F$ useful? - I

For hard scattering ( $2 \rightarrow 2$  processes),  $x_F \sim x_1 - x_2$ , where  $x_1$  is the Bjorken  $x$  of the parton from the hadron heading towards the apparatus and  $x_2$  is the Bjorken  $x$  of the parton from the other colliding hadron. In general, forward particle production probes these  $x$  values at “low scale” (as set by  $p_T$ ). Distributions are for inclusive forward jets.

$p+p$ ,  $\sqrt{s}=510$  GeV, PYTHIA 6.222/GEANT, tower jets



$p+p$ ,  $\sqrt{s}=510$  GeV, PYTHIA 6.222/GEANT, tower jets



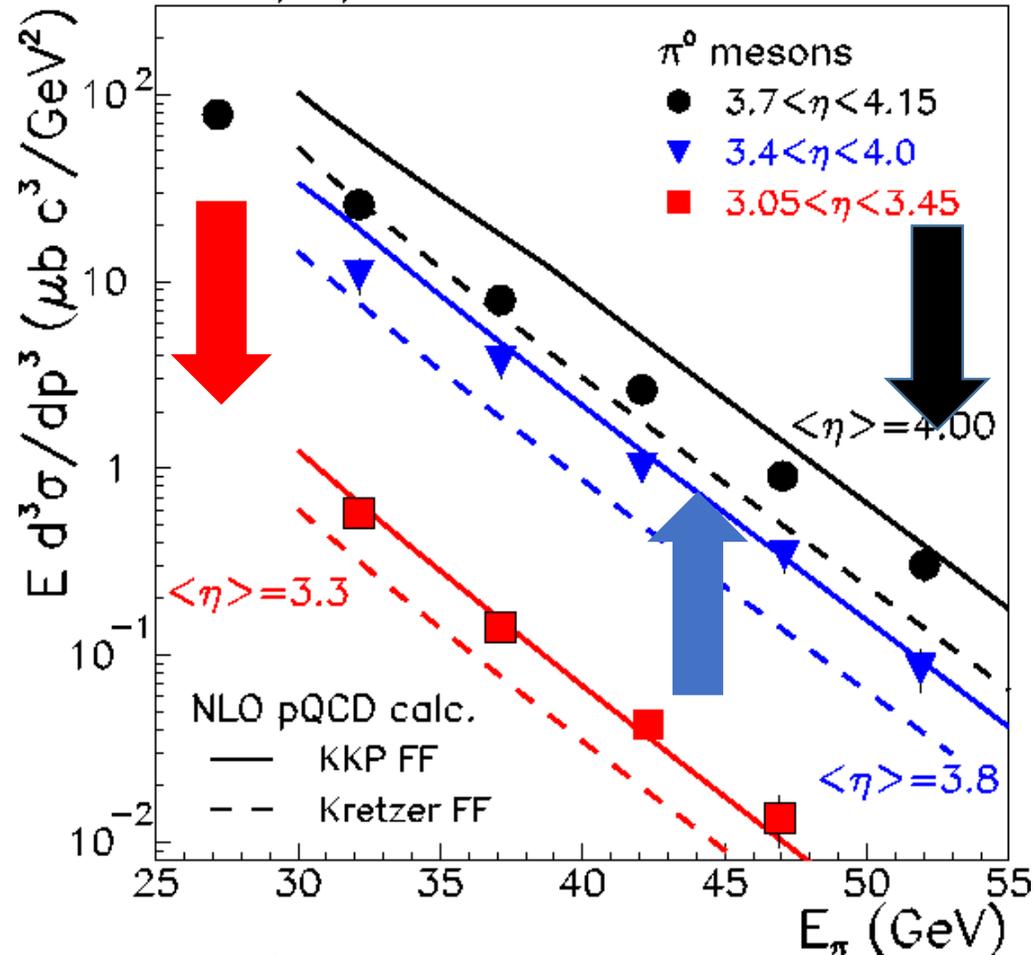
1/13/2016 Valence-like quarks for  $x_F > 0.1$

$x_2$  is broad, but extends to very low  $x$  ( $\sim \text{few} \times 10^{-4}$ ).  
Forward dijets can select low  $x$  (see below)

# Why is large $x_F$ useful? - II

PRL 97 (2006) 192302

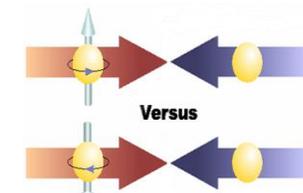
$p+p \rightarrow \pi^0 + X$   $\sqrt{s}=200$  GeV



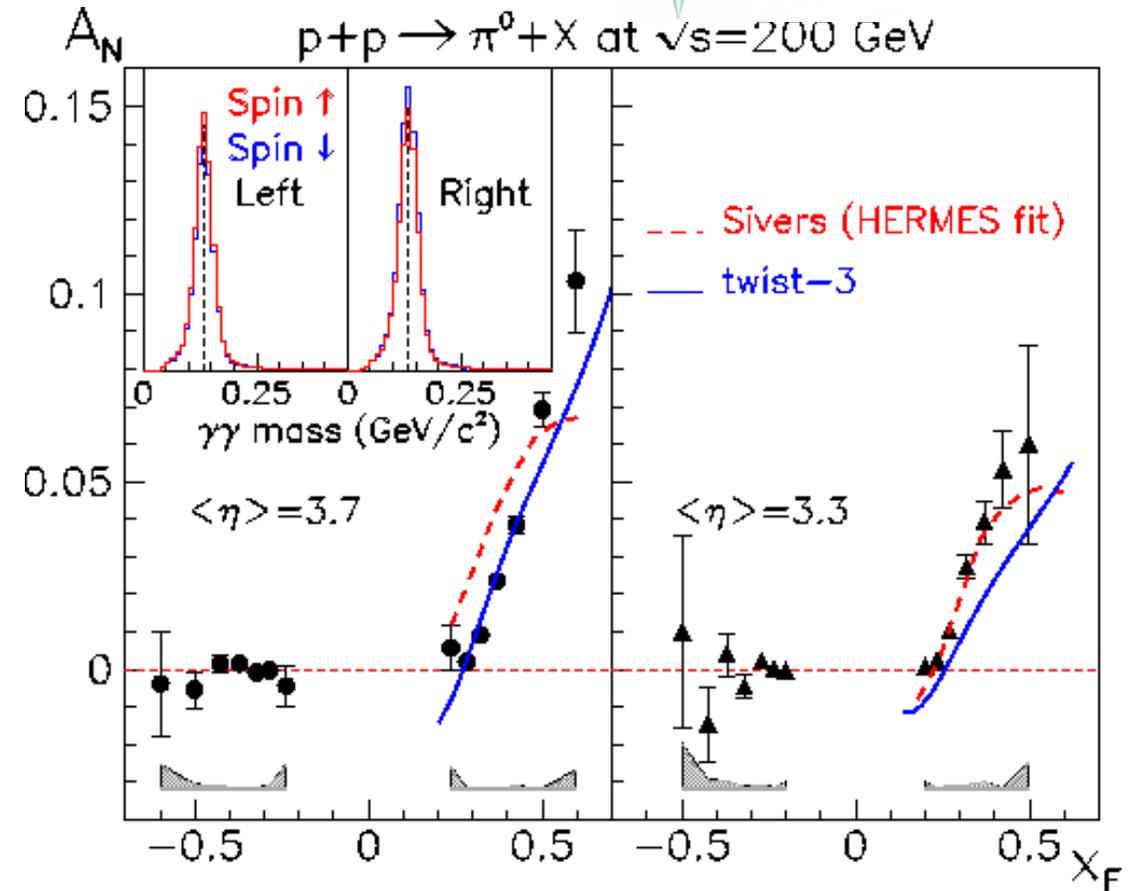
For  $p_T > 2$  GeV/c (arrow positions), measured cross sections are in good agreement with NLO pQCD, albeit with large scale dependence which is smaller for jets (see below)

1/13/2016

PRL 101 (2008) 222001



$$A_N = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}}$$

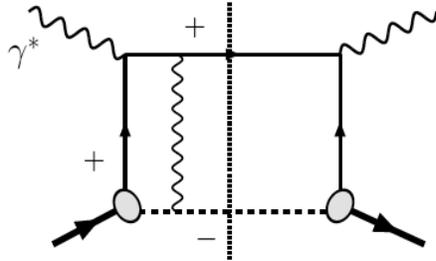


Although cross sections can be described by NLO pQCD, there are still large transverse single-spin asymmetries (SSA), that are expected to be zero in naive pQCD but can arise from spin-correlated  $k_T$

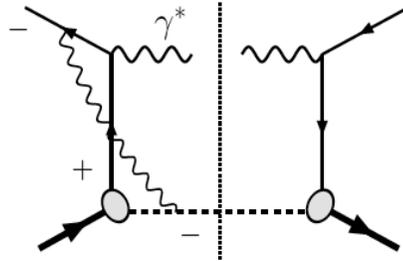
# Attractive vs Repulsive Sivers Effects

## Unique Prediction of Gauge Theory !

Simple QED example:

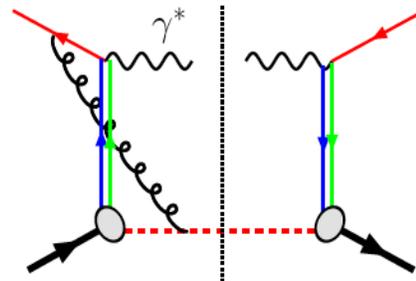
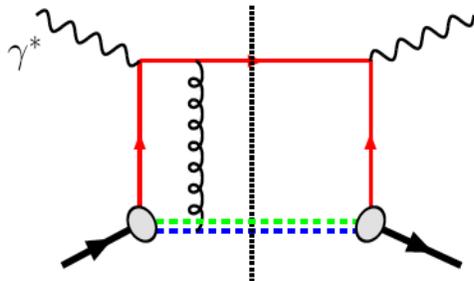


DIS: attractive



Drell-Yan: repulsive

Same in QCD:



$$\text{Sivers}|_{\text{DIS}} = -\text{Sivers}|_{\text{DY}}$$

As a result:

## $A_N$ DY Goal

Measure the transverse single spin asymmetry for forward low-mass dileptons produced via the Drell-Yan process to test theoretical predictions of a sign change for the initial-state spin-correlated  $k_T$ -dependent distribution function (Sivers function).

The objective was to match as closely as experimentally possible kinematics between DY [dilepton mass and  $x_1 \sim x_F$ ] and semi-inclusive deep inelastic scattering ( $Q^2$  and Bjorken  $x$ ).

*Transverse Spin Drell-Yan Physics at RHIC (2007)*

# A<sub>N</sub>DY Setup at IP2 for 2011 RHIC Run

Left/right symmetric HCal

Trigger/DAQ electronics

Left/right symmetric ECal

Left/right symmetric preshower

Blue-facina BBC

Beryllium vacuum pipe

- This was a stage-1 test that could not have worked for forward DY
- The stage-1 test did measure forward jets
- There were not further stages

Absolute Polarimeter ( $H^\uparrow$  jet)

RHIC pC Polarimeters

A<sub>N</sub>DY

Siberian Snakes

Siberian Snakes

PHENIX

STAR

Spin Rotators (longitudinal polarization)

Spin Rotators (longitudinal polarization)

Pol. H<sup>-</sup> Source

LINAC

BOOSTER

AGS

200 MeV Polarimeter

AGS pC Polarimeter

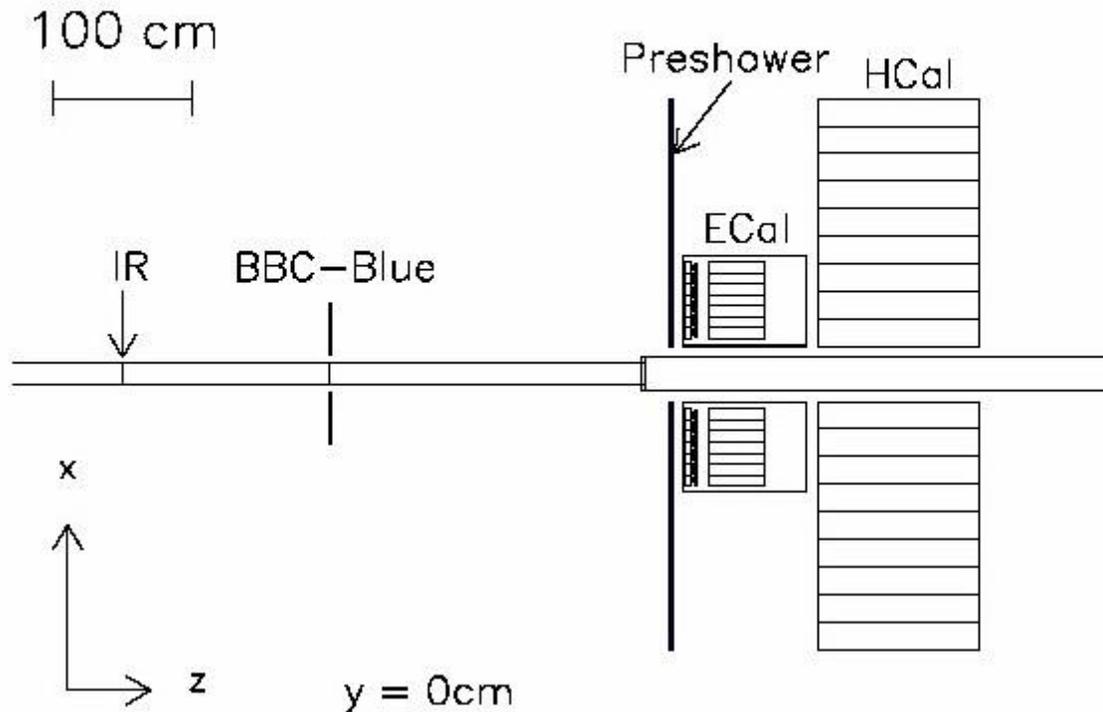
Strong AGS Snake

Helical Partial Siberian Snake

# $A_N$ DY Setup at IP2 for 2011 RHIC Run

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IP2/DY-Run11

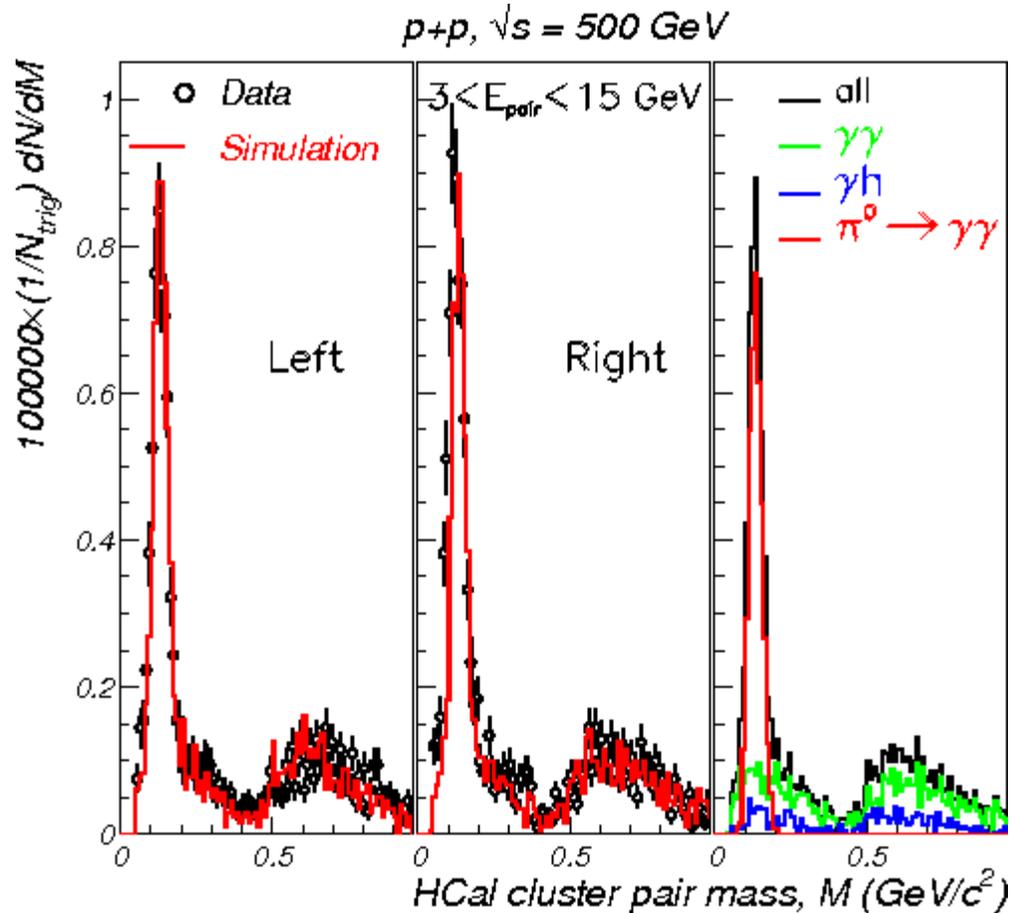


- Beam-beam counter (BBC) for minimum-bias trigger and luminosity measurement (from PHOBOS [NIM A474 (2001) 38])
- Zero-degree calorimeter and shower maximum detector for luminosity measurement and local polarimetry (ZDC/ZDC-SMD, not shown)
- Hadron calorimeter (HCal) are L/R symmetric modules of 9x12 lead-scintillating fiber cells,  $(10\text{cm})^2 \times 117\text{cm}$  (from AGS-E864 [NIM406(1998)227])
- Small ECal - 7x7 matrices of lead glass cells,  $(4\text{cm})^2 \times 40\text{cm}$  (loaned from BigCal at JLab)
- Preshower detector - two planes, 2.5 & 10 cm
- In 2012, modular calorimeters were replaced by an annular calorimeter

# Calibrations-I

## Electromagnetic Response

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- Cosmic-ray muons were used to adjust relative gains in advance of collisions
- The primary determination of the energy scale was from reconstruction of  $\pi^0 \rightarrow \gamma\gamma$  from single-tower cluster pairs. The maximum energy for this calibration was limited by photon merging into the coarse  $(10 \text{ cm})^2$  towers. [See below for pixelization results from this same calorimeter]
- Full PYTHIA/GEANT simulation agrees with data, for both the pair-mass resolution of the calorimeter, as well as the neutral pion reconstruction efficiency.
- Subsequent test-beam studies at FNAL [T1064] are consistent with an excellent response of this calorimeter to incident photons and electrons.

# Calibrations-II

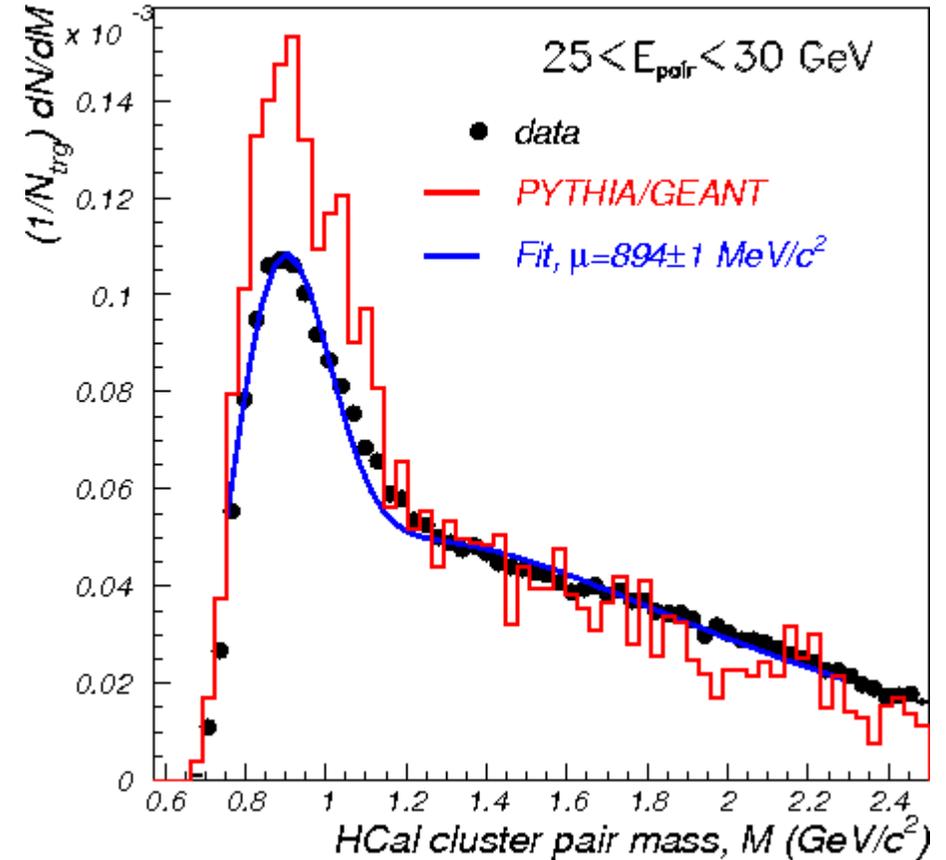
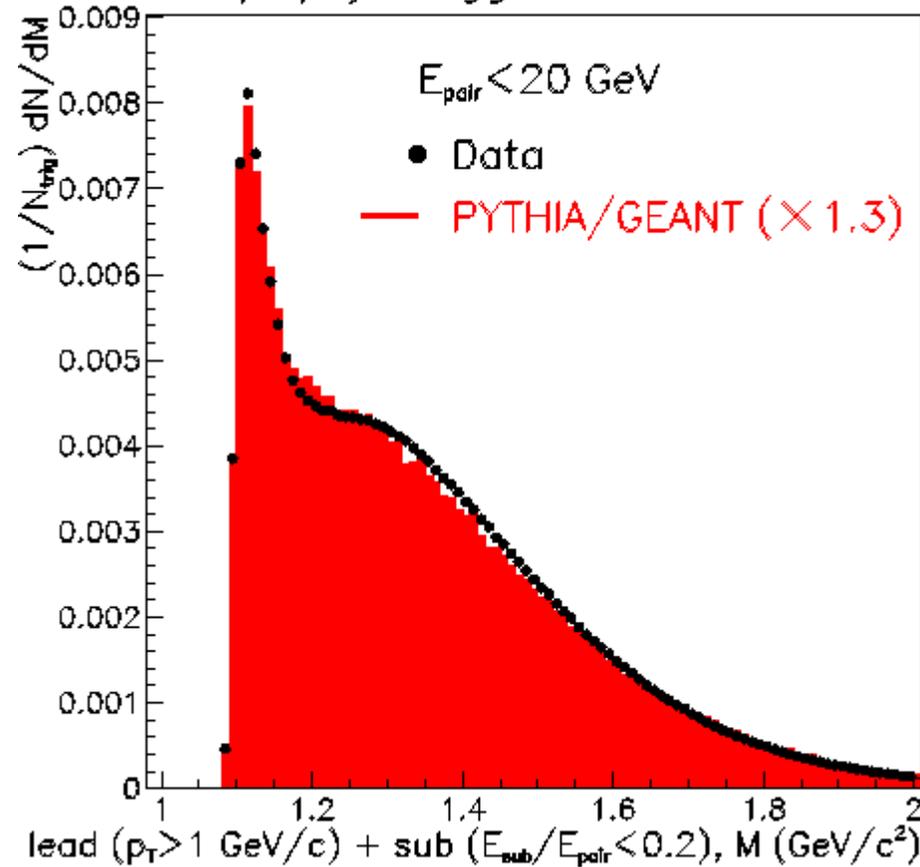
arXiv:1308.4705

Hadronic Reponse

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p+p, jet trigger,  $\sqrt{s}=510$  GeV

p+p,  $\sqrt{s}=510$  GeV

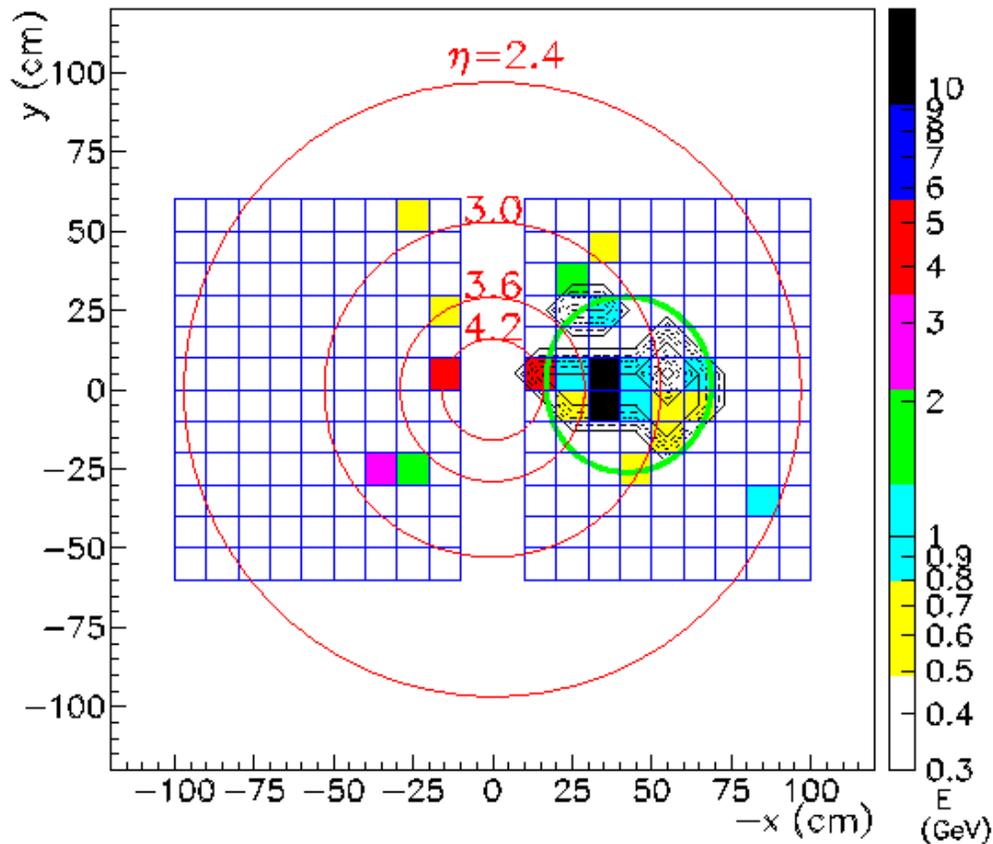


- Use BBC detector to tag HCal clusters made by incident charged hadrons. Mass assignments are then made.
- Tagged cluster-pair mass distributions are consistent with  $\Lambda \rightarrow \pi^- p$  (left) and  $K^*(892) \rightarrow \pi^+ K^-$  (right) and charge conjugates
- Use  $E = 1.12E' - 0.1$  GeV for jet finding from an event list of tower energies that use the photon calibration ( $E'$ )

# Jet Reconstruction – Anti- $k_T$ Jet Finder

Trigger on HCal masked ADC Sum in L/R Modules  
 Display anti- $k_T$  jet clusters satisfying acceptance cuts

Run=12107004.001, trig=Jet, Event=15, mod=2, anti- $k_T$



## • Anti- $k_T$ Jet Finder Procedure :

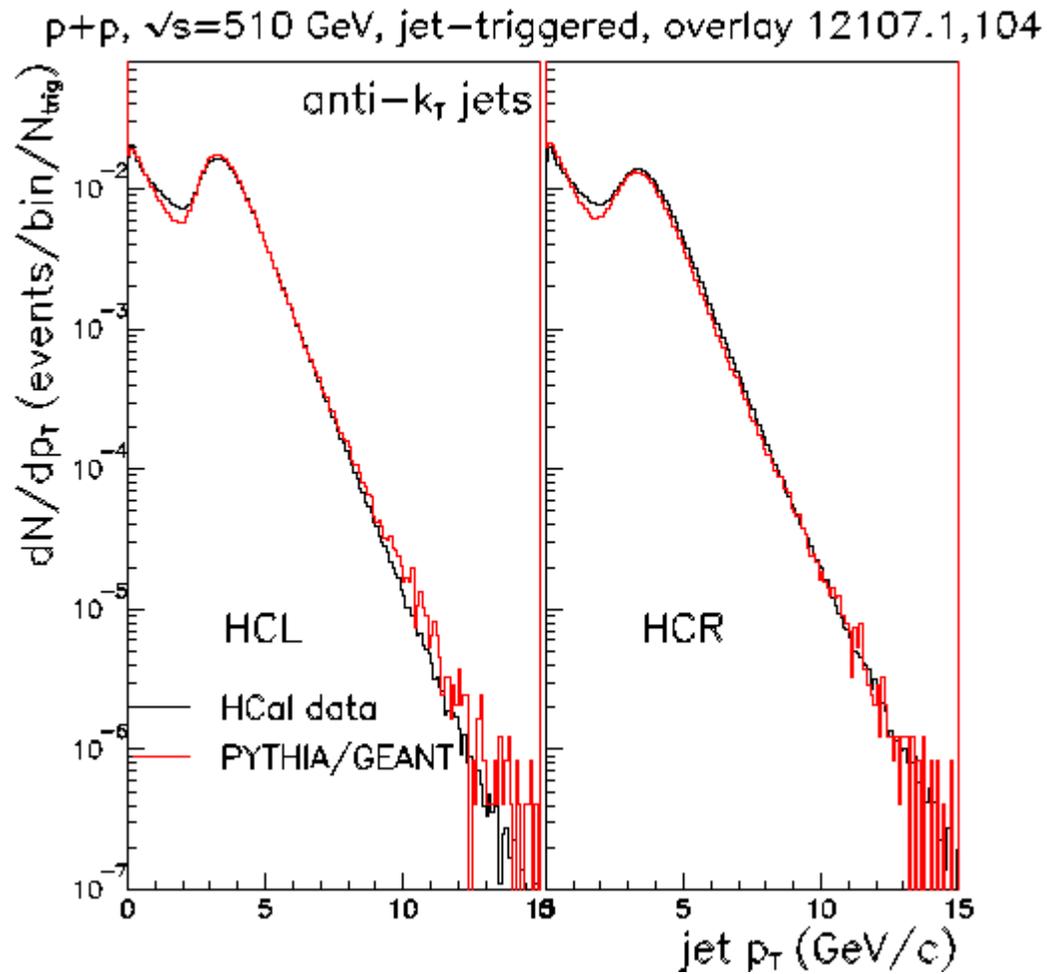
- Iteratively merge pairs of towers until towers cease to satisfy distance criteria
  - No Seed
  - Towers can be outside trigger region
- Distance Criteria (clusters  $j,k$ ) :
  - $d_{jk} = \min(k_{Tj}^{-2}, k_{Tk}^{-2})(R_{jk}^2/R^2)$
  - $R_{jk}^2 = (\eta_j - \eta_k)^2 + (\Phi_j - \Phi_k)^2$
  - If  $d_{jk} < k_{Tj}^{-2}$  then merge clusters  $j,k$
- Use cone with  $R_{jet} = 0.7$  in  $\eta$ - $\Phi$  space but cluster towers can fall outside of cone
- Impose acceptance cuts to accept/reject jet:
 
$$|\eta_j - 3.25| < 0.25$$

$$|\Phi_j - \Phi^{Off}| < 0.50$$

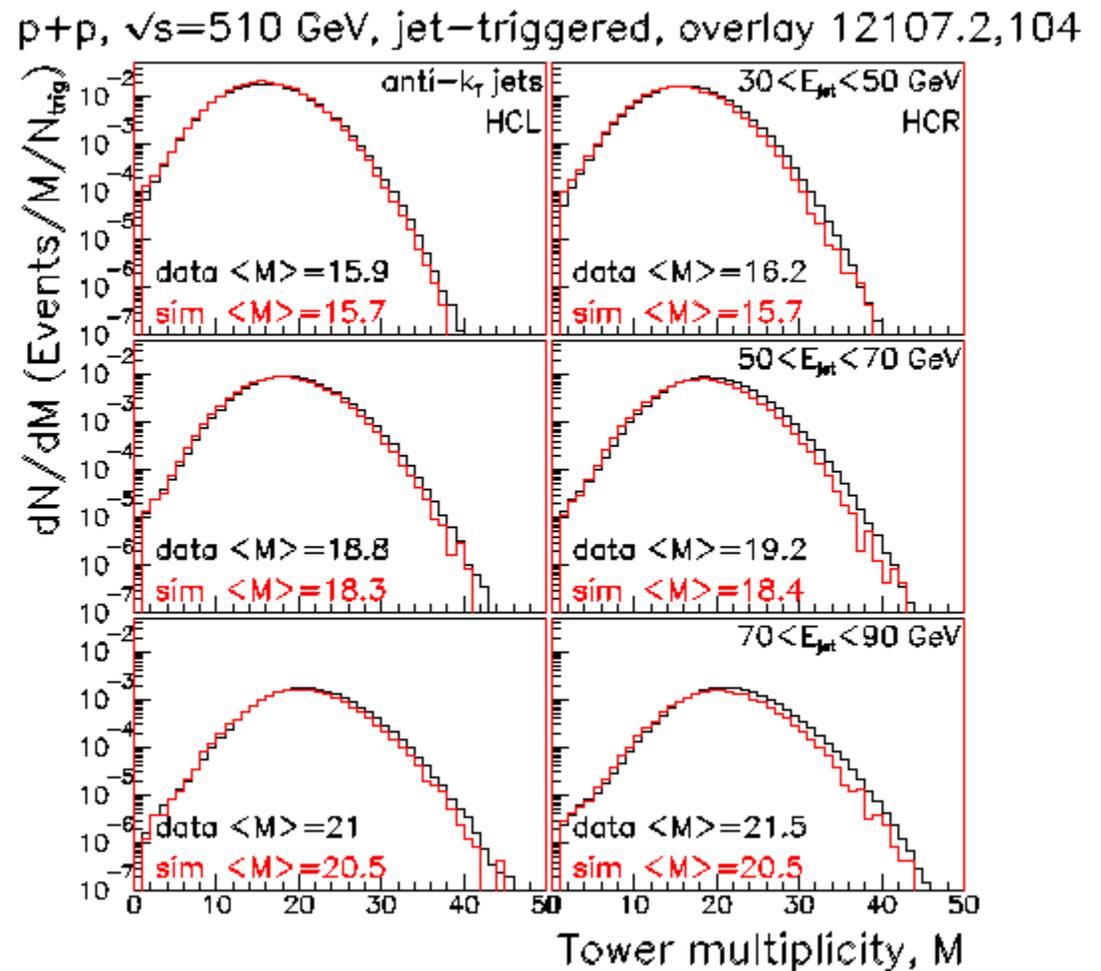
where  $\Phi^{Off} = 0$  for HCL  
 $\Phi^{Off} = \pi$  for HCR
- Energy Cut :  $E_{jet} > 30$  GeV
- Algorithm : arXiv : 0802.1189  
 arxiv : 1209.1785

1/13/2016 Events look “jetty” / Results with anti- $k_T$  algorithm similar to midpoint cone algorithm

# Comparison of Data to PYTHIA 6.222/GEANT Simulation



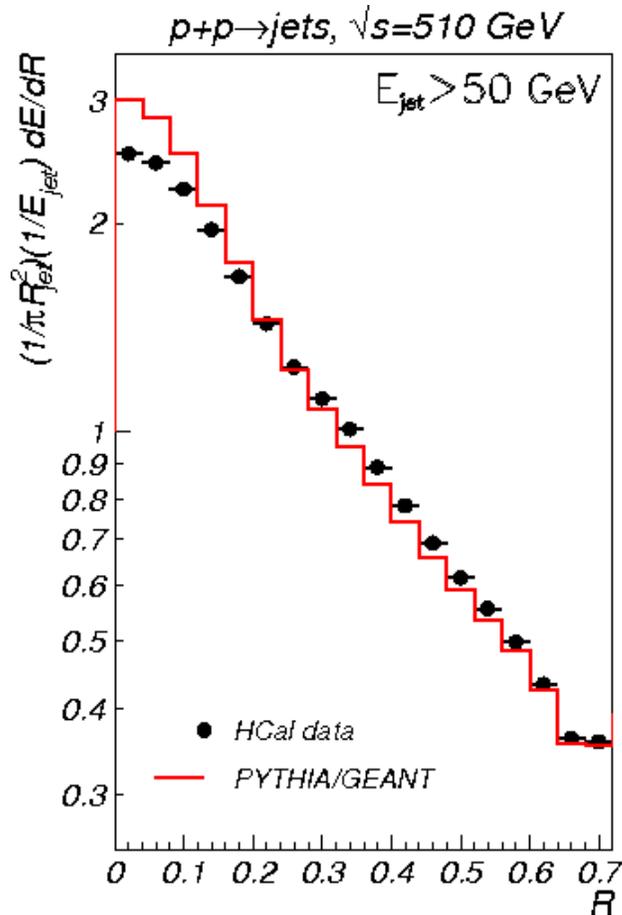
Uncorrected  $p_T$  distribution of anti-kT clusters



Uncorrected multiplicity of towers in anti-kT cluster

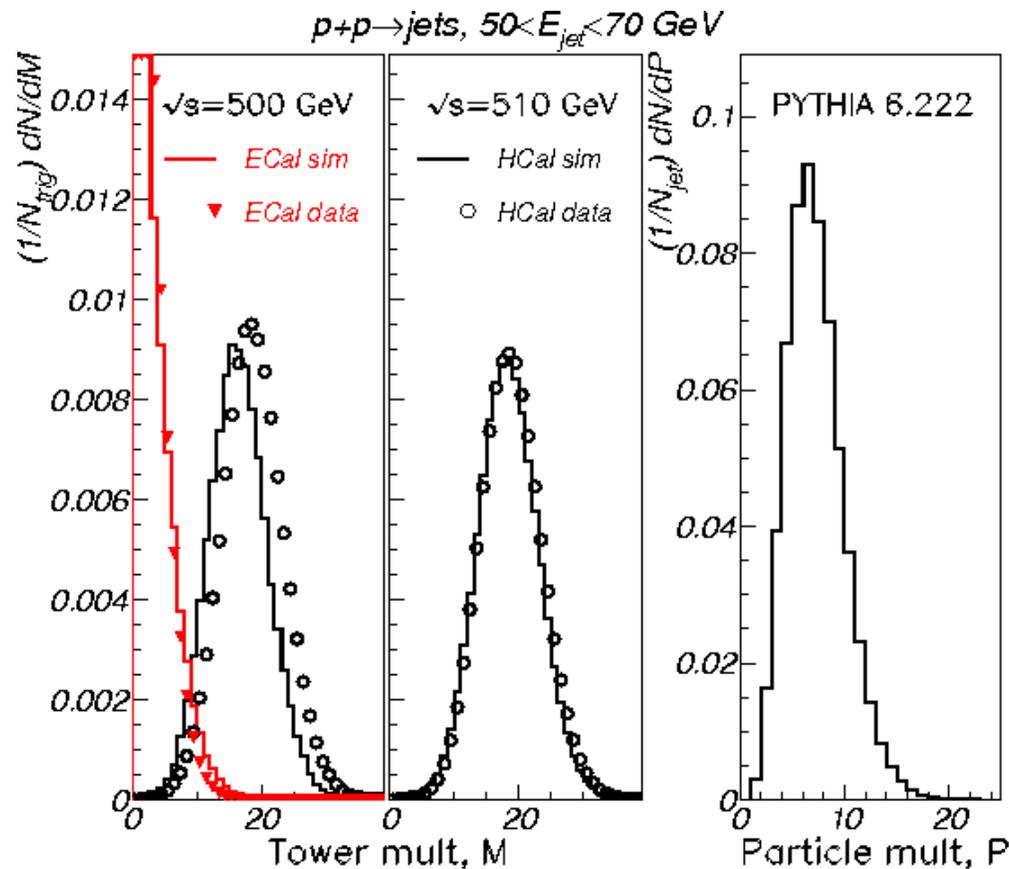
Good description of data by simulation  $\rightarrow$  use simulation for efficiency correction

# What is a forward jet?



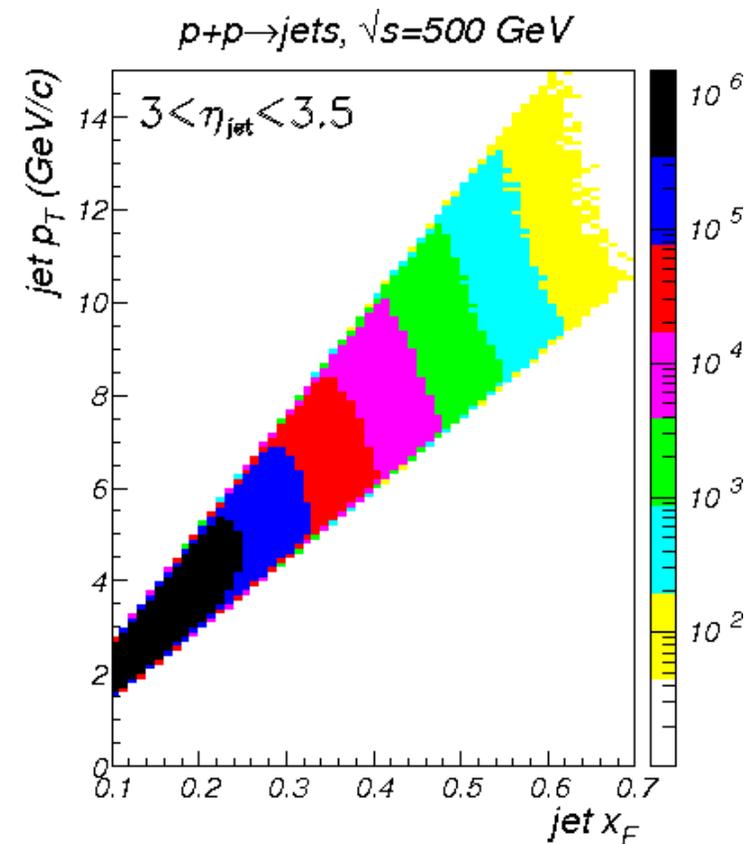
Event averaged jet shape: how the energy is distributed a distance  $R$  in  $\eta, \phi$  from the thrust axis

⇒ the anti-kT clusters have shapes similar to midrapidity jets



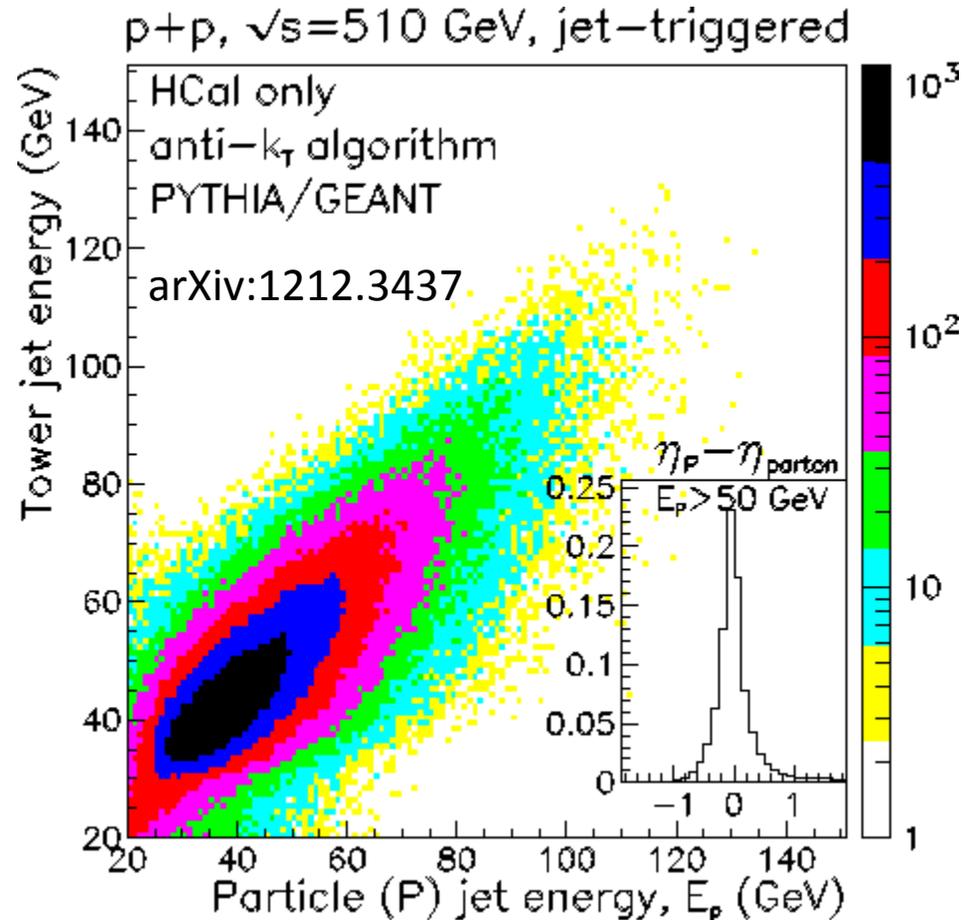
(left) tower multiplicities, as used for  $A_N$ ;  
 (middle) tower multiplicities, as used for  $\sigma$ ;  
 (right) incident particle multiplicity from simulation

⇒ multiplicity similar to jets of comparable scale



Acceptance of contained jets from particles with  $2.4 < \eta < 4.2$  correlates  $x_F$  and  $p_T$  for the jet cluster

# Jet Energy Scale - I

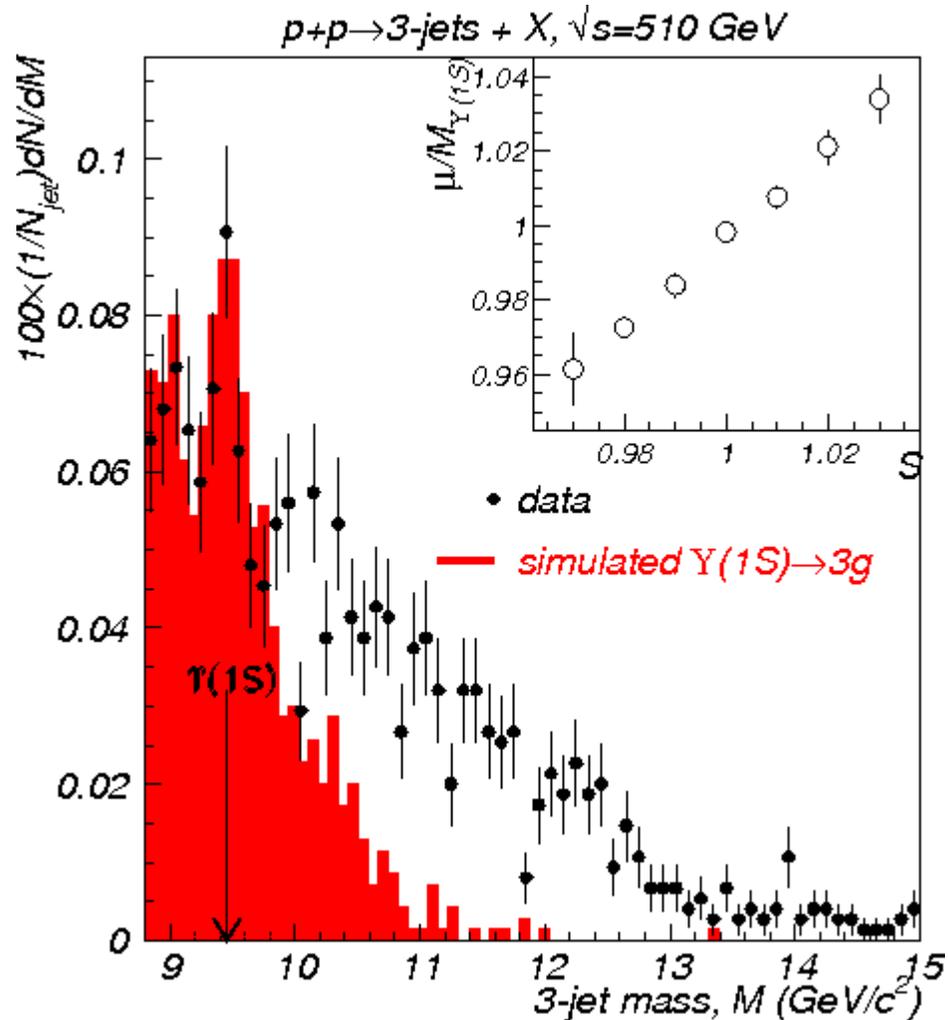


- Simulations confirm energy scale of jets, by comparison of “tower” jets [with full detector response] to “particle” jets [excluding detector response].
- Reconstructed jets are directionally matched to hard-scattered partons as generated by PYTHIA

Correlation between tower jet [from PYTHIA/GEANT] to particle jet [from PYTHIA]. The inset shows the  $\eta$  component of the directional match ( $\Delta\eta$ ) between particle jets and a hard-scattered parton, whose direction is defined by  $\eta_{parton}, \phi_{parton}$ . There is a 82% match requiring  $|\Delta\eta|, |\Delta\phi| < 0.8$

# Jet Energy Scale - II

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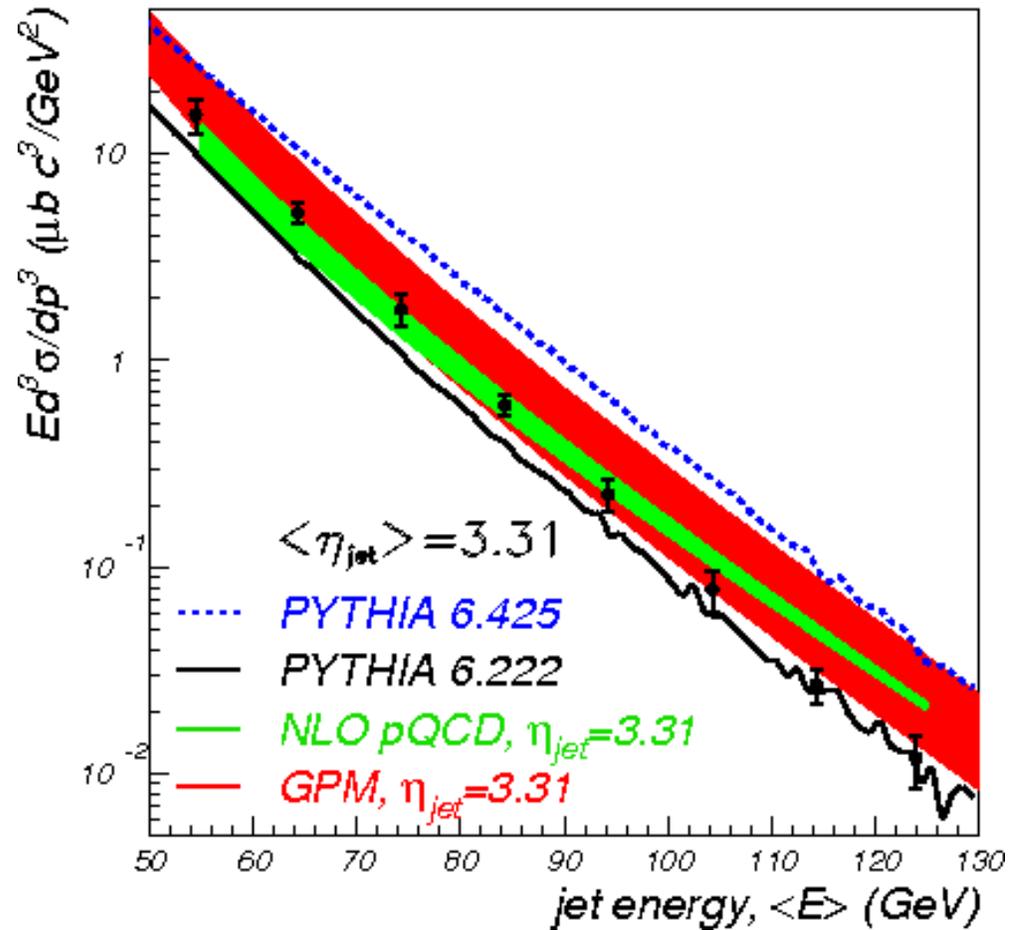


- Test jet energy scale by reconstruction of invariant mass for multi-jet events
- Observe  $3.5\sigma$  statistical significance peak, attributed to  $\Upsilon(1S) \rightarrow 3g$ . The red overlay is a simulation of the signal from the PYONIA generator of  $\Upsilon(1S) \rightarrow 3g$ , run through GEANT, and then reconstructed as done for the data
- For the inset,  $S$  rescales the energy calibrations, so tests the jet-energy scale.
- Peaks are also observed in 2-jet mass attributed to  $\chi_{2b} \rightarrow 2$  gluons

# Forward Jet Cross Sections

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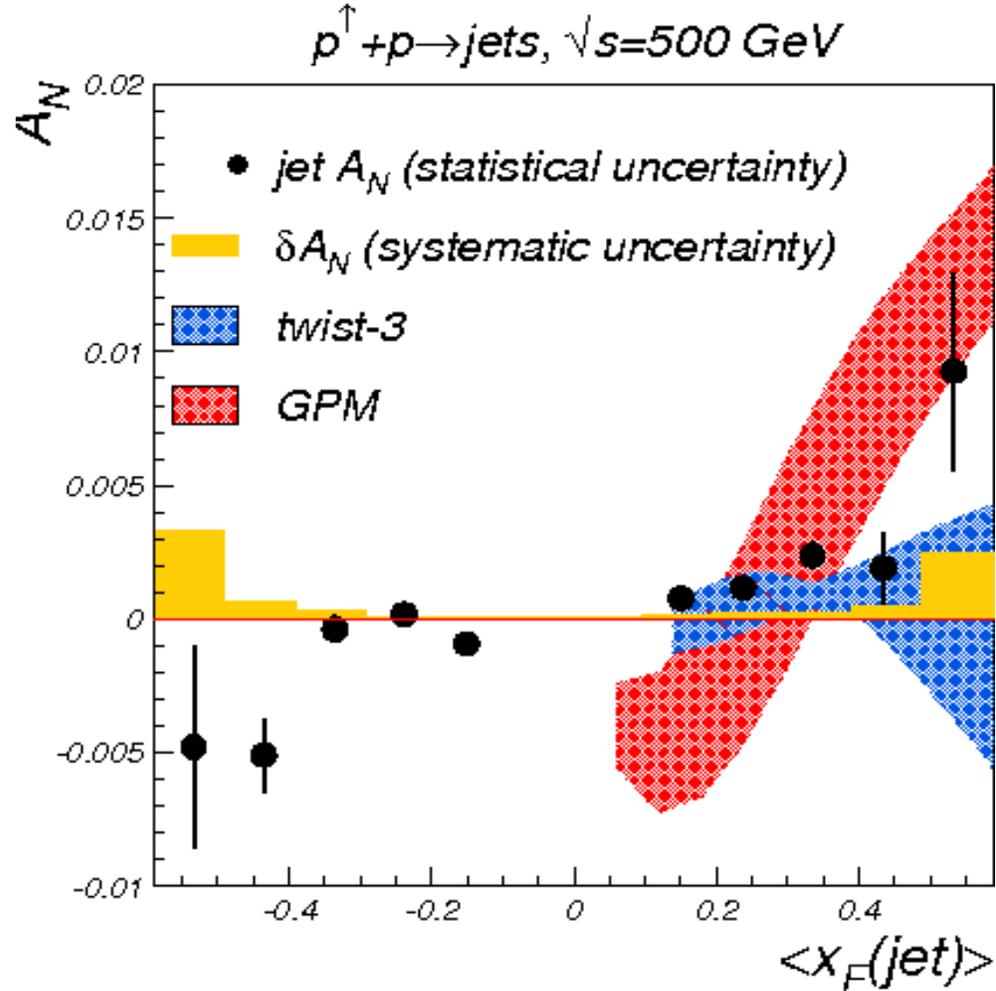
$p+p \rightarrow \text{jets}, \sqrt{s}=510 \text{ GeV}$



- Uncertainties include both statistical and systematic estimates [as described in backup]
- Strong dependence on both  $x_F$  and  $p_T$  requires data/theory comparisons at  $\langle \eta_{\text{jet}} \rangle$
- NLO pQCD [PRD 86 (2012) 094009] calculation provides a good description of the data using CTEQ6.6M PDF. Note the small scale dependence [band represents range of scale from  $\mu=2p_T$  to  $\mu=p_T/2$ ]
- Leading-order pQCD model calculation assuming factorization in the use of  $k_T$  dependent distribution functions [generalized parton model (GPM), PRD 88 (2013) 054023] also describes the data. The larger scale dependence is likely a consequence of a leading-order calculation
- Particle jet reconstructions [no detector effects beyond acceptance] with the anti-kT algorithm with  $R_{\text{jet}}=0.7$  are used to compare default PYTHIA 6.222 [prior to tunings for the LHC] and PYTHIA 6.425 [“Field tune A”] to data. PYTHIA 6.222 was previously found to describe forward  $\pi^0$  production at  $\sqrt{s}=200 \text{ GeV}$  [arXiv:hep-ex/040312].

# Forward Jet Analyzing Power

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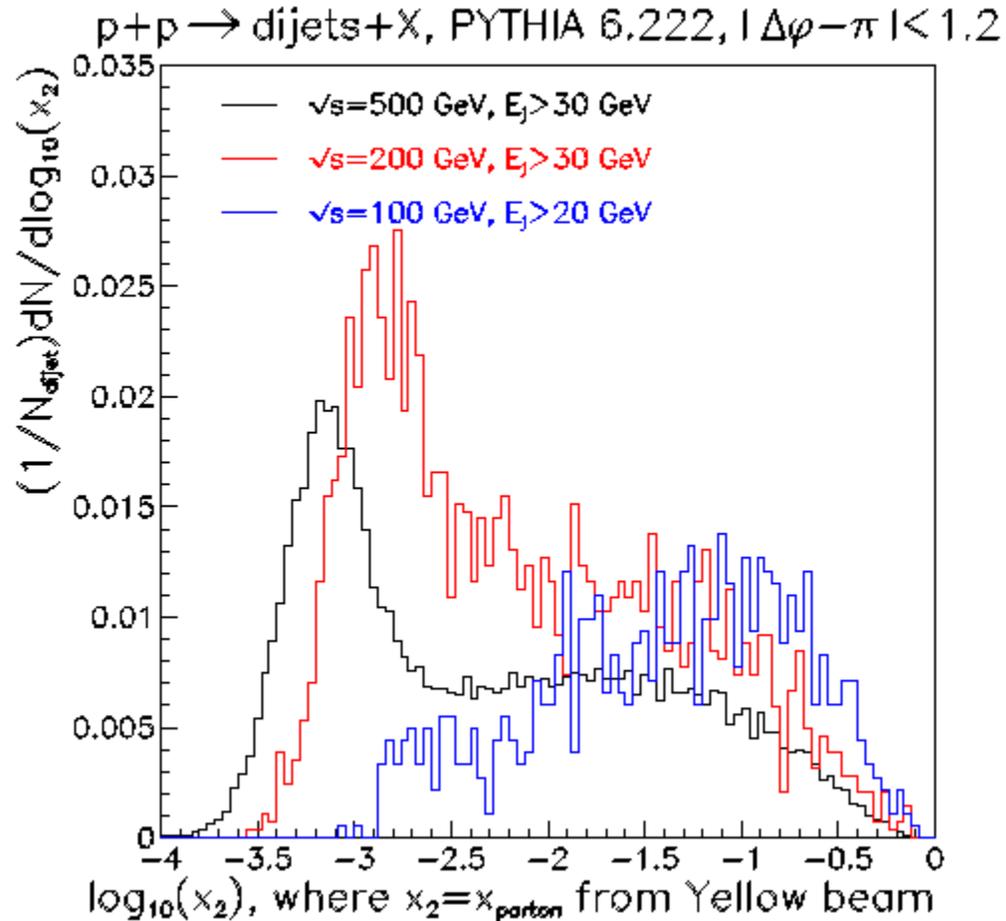


- Analyzing power is computed via the cross-ratio method that exploits the mirror symmetry of the apparatus

$$P_{beam} A_N = \frac{\sqrt{N_L^\uparrow N_R^\downarrow} - \sqrt{N_R^\uparrow N_L^\downarrow}}{\sqrt{N_L^\uparrow N_R^\downarrow} + \sqrt{N_R^\uparrow N_L^\downarrow}}$$

- $P_{beam}$  – beam polarization [ $0.526 \pm 0.027$  for  $x_F > 0$ ]
- $N_{L(R)}^{\uparrow(\downarrow)}$  – number of jets in left or right module for beam-spin up or down in each bin of  $\langle x_F(\text{jet}) \rangle$
- Systematic uncertainties [as described in backup] are quoted separately from statistical uncertainties, and are available in tabulated form in the published paper
- Both theory curves for  $x_F > 0$  fit the Siverts function in semi-inclusive deep inelastic scattering. “twist-3” is a collinear approach [PRL 110 (2013) 232301] with color gauge link effects. “GPM” is a generalized-parton model calculation [PRD 88 (2013) 054023]
- Jets with  $x_F < 0$  would be described by low- $x$  spin effects: e.g., PRD 89 (2014) 074050 and PRD 89 (2014) 034029

# Forward Dijets



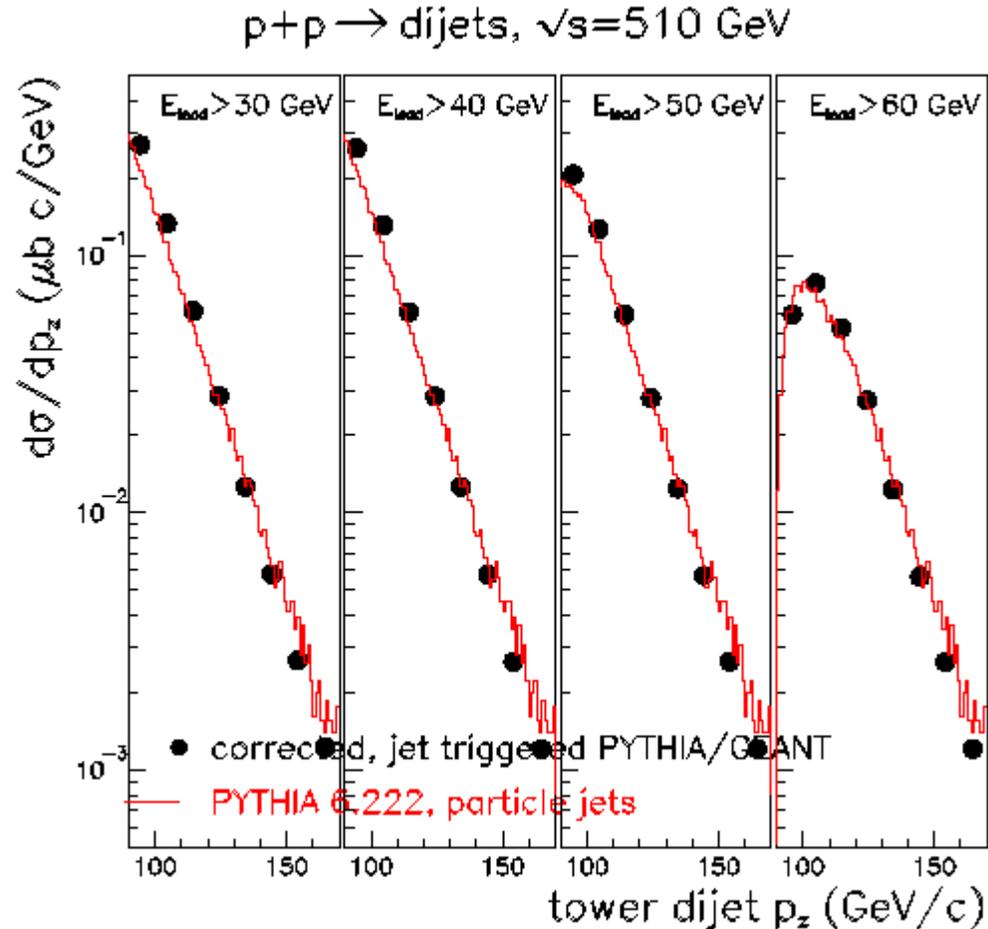
Reconstruction of  $>1$  jet in the forward direction can emphasize hard-scattering contributions from low- $x$  gluons

Examples of why this is important are

- Extending probes of gluon polarization to low- $x$  by measurement of longitudinal double-spin asymmetries
- QCD processes are the reducible background to forward Drell-Yan production of low-mass virtual photons

# Test of Dijet Corrections

Comparison of corrected PYTHIA/GEANT tower dijets to PYTHIA particle dijets

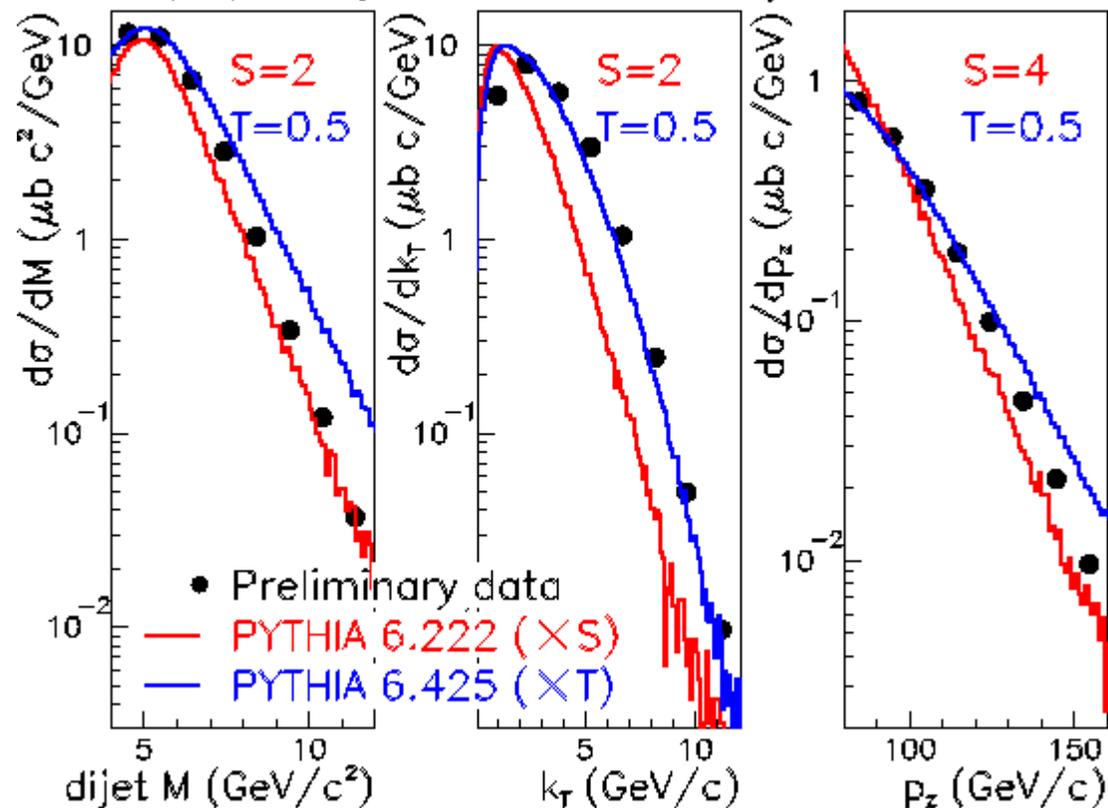


- It is found that the dijet  $\varepsilon_{\text{trig}}(V)$  [for  $V=M, k_T, p_z$ ] is the only correction required; i.e.,  $\varepsilon_{\text{det}}(V)=1$
- The dijet correction procedure when applied to PYTHIA/GEANT tower dijets reproduces the input PYTHIA particle dijets (animate for  $V=k_T$  and  $p_z$  distributions)
- Require  $M > 4$  GeV/ $c^2$  when reporting  $d\sigma/dk_T$  and  $d\sigma/dp_z$ .

# Dijet Results

arXiv:1308.4705

$p+p \rightarrow$  dijets,  $\sqrt{s}=510$  GeV,  $E_{\text{jet}} > 30$  GeV

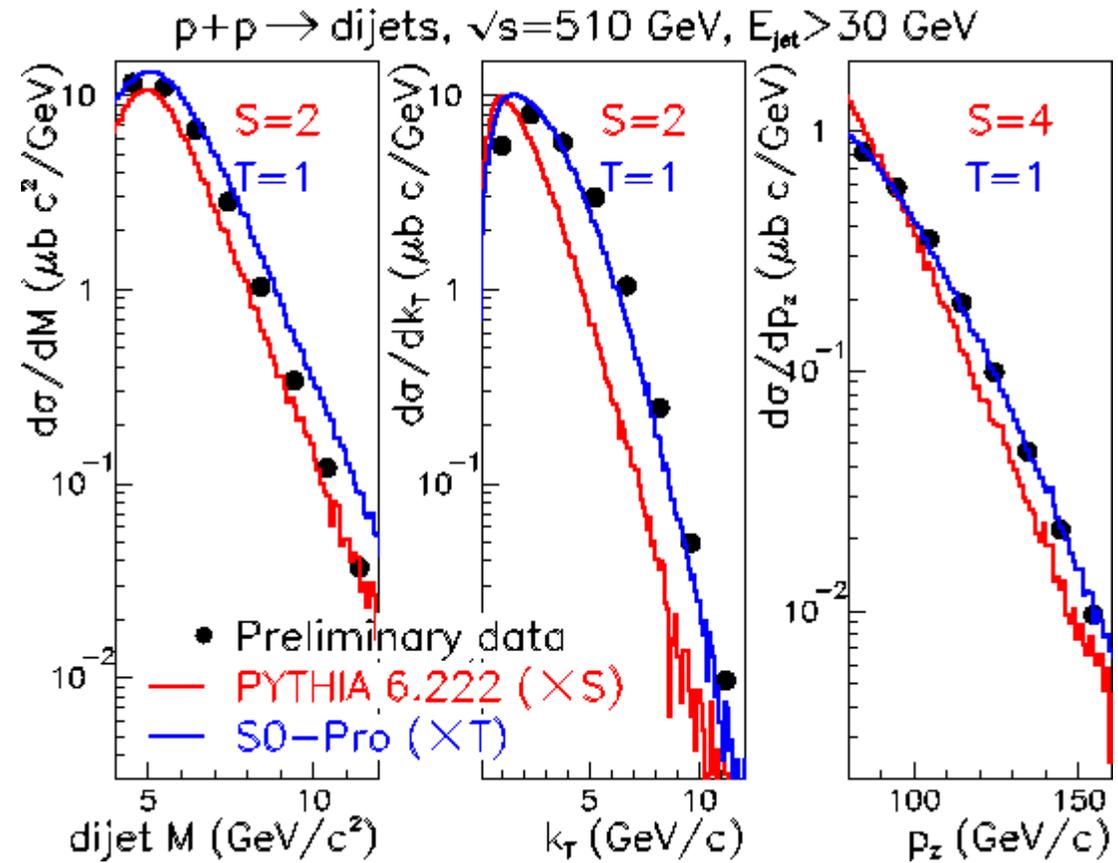
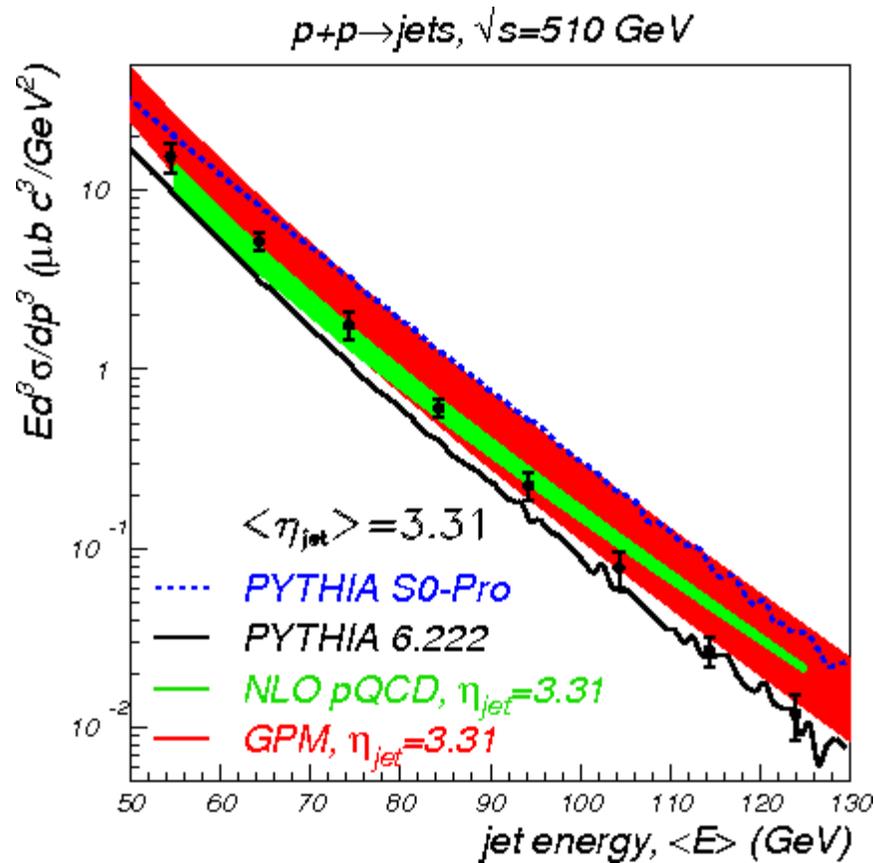


- Use “DY-like” variables, where  $M$  is dijet mass assuming massless partons;  $k_T$  is the net transverse momentum of the dijet; and  $p_z$  is the longitudinal momentum of the dijet
- Preliminary results are efficiency corrected, but at present do not reflect the acceptance imposed on the jet pair [each jet of pair requires  $3.0 < \eta_{\text{jet}} < 3.5$ , where jets are reconstructed from a nearly annular calorimeter spanning  $\sim 2.4 < \eta < 4.2$ ]
- Results are compared to particle dijet results, using default settings for PYTHIA 6.222 and PYTHIA 6.425. Neither version explains the data.

# PYTHIA Tunings

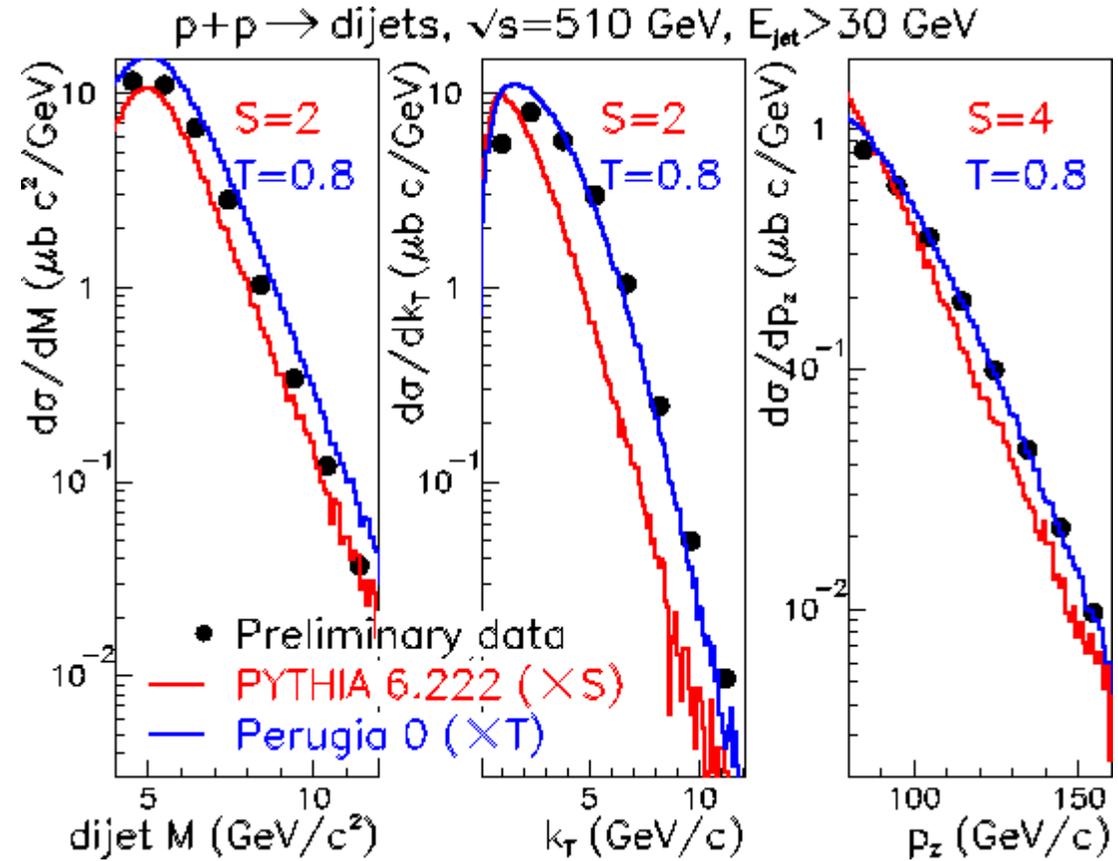
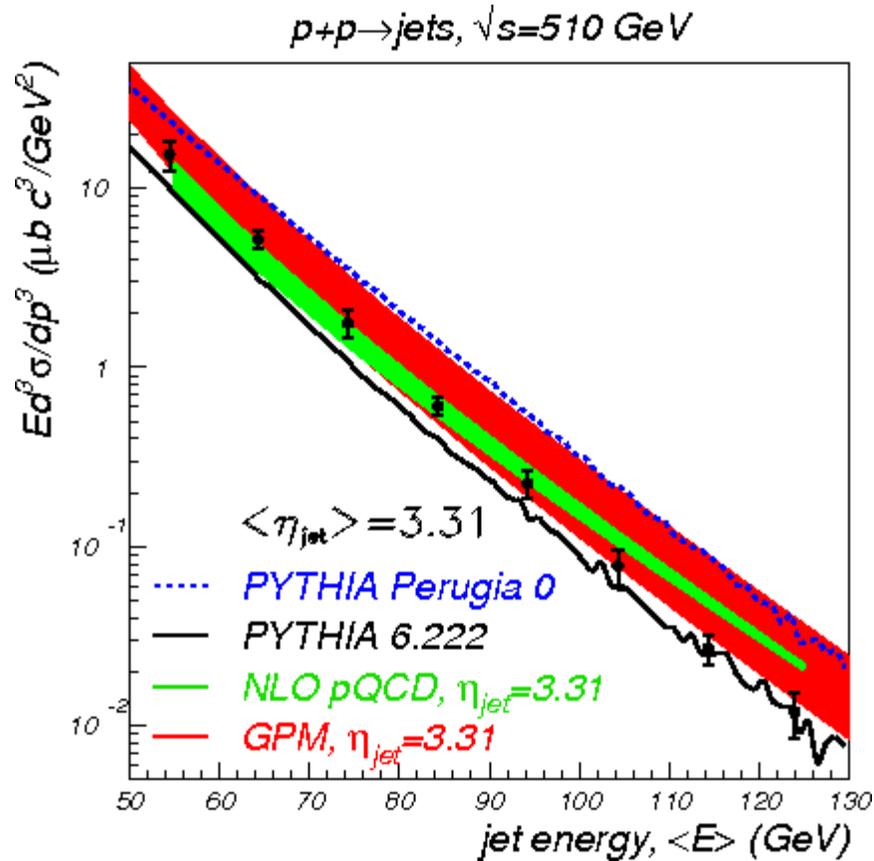
- The LHC high-energy program has prompted many retunings of PYTHIA, so that backgrounds in e.g. dijet mass are well modeled to allow new particle searches. See P.Z. Skands, PRD 82 (2010) 074018 [arXiv:1005.3457]
- PYTHIA tunings most commonly adjust initial-state and final-state showering parameters; multi-parton interaction model parameters; etc. As will be shown, inclusive forward jets and forward dijets from RHIC are sensitive to these tunings [as should be expected, since the rapidities involved for forward dijets at RHIC rival those from midrapidity at the LHC]
- In general, any serious low-x physics study of forward particle production will need to deal with the physics of parton showers and multi-parton interactions. It is not good to attempt to “correct” measurements for these effects. Experimental results should report what’s observed, rather than subtracting model-dependent quantities from what is measured [in my opinion...]

# Data versus S0-Pro Tune



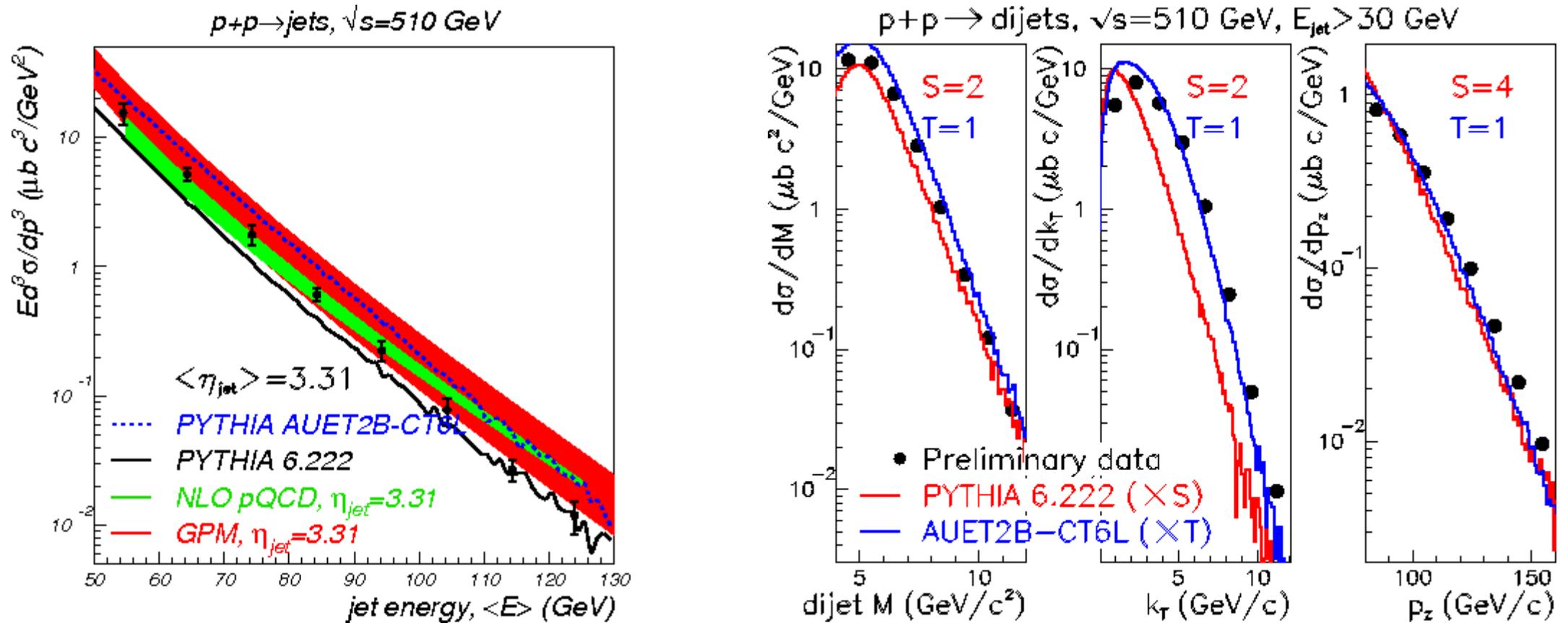
- Good fit to dijets in distribution shape and normalization
- Overpredicts the inclusive jet cross section

# Data versus Perugia 0 Tune



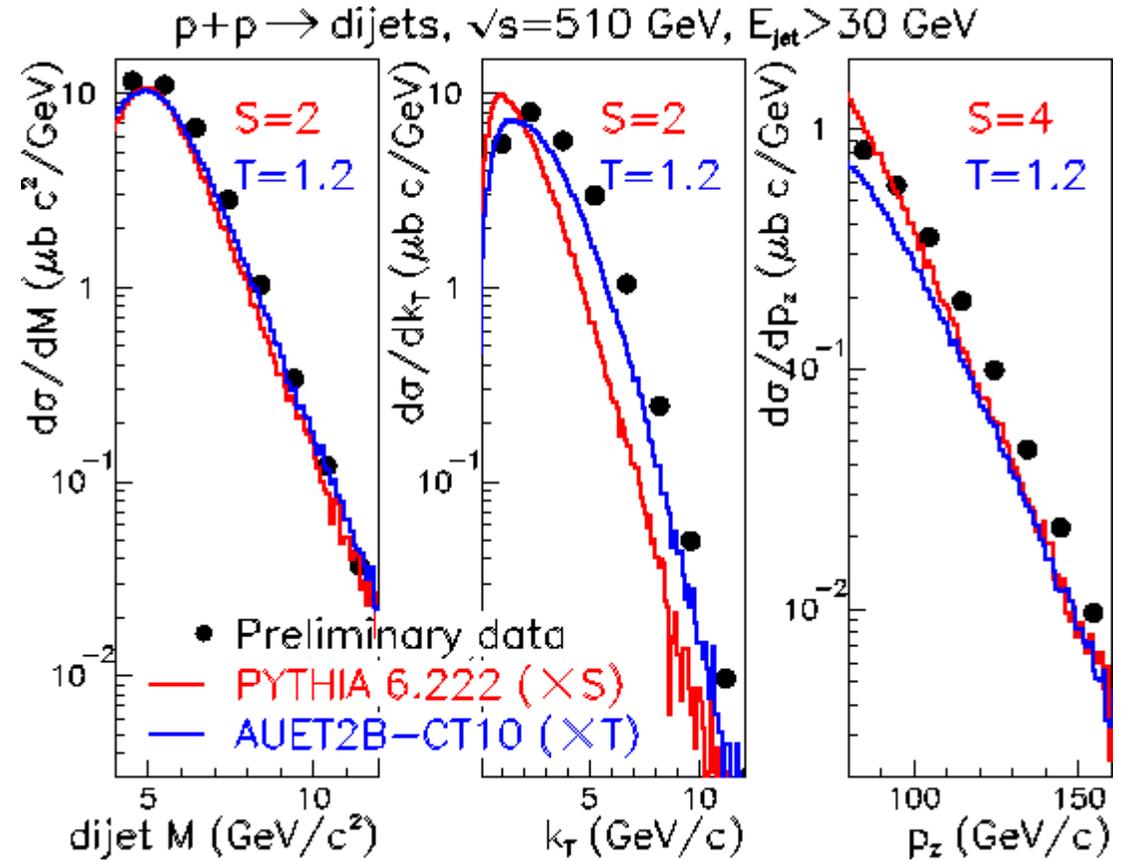
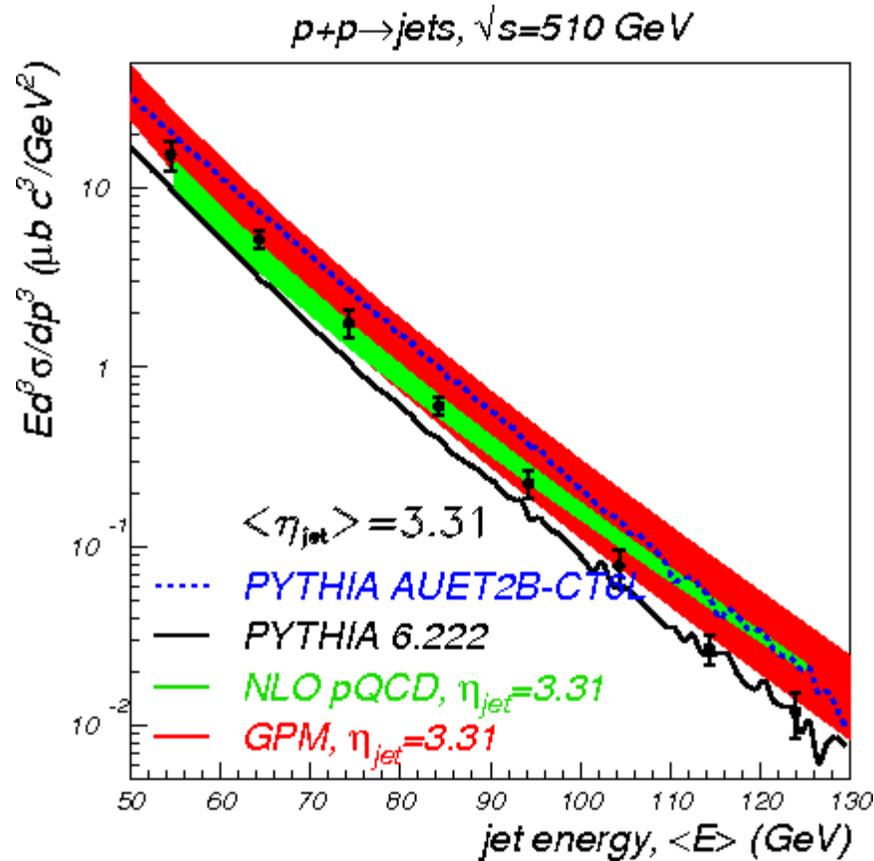
- Perugia 0 is commonly used for RHIC midrapidity data
- Overpredicts dijet cross sections by  $\sim 20\%$
- Overpredicts inclusive jet cross section

# Data versus Atlas [AUET2B-CT6L] Tune



- PYTHIA tune developed by Atlas [arXiv:1512.001917], including LHC 7 TeV data
- Reasonable representation of dijet distributions and normalization
- Better description of inclusive jet result than other tune

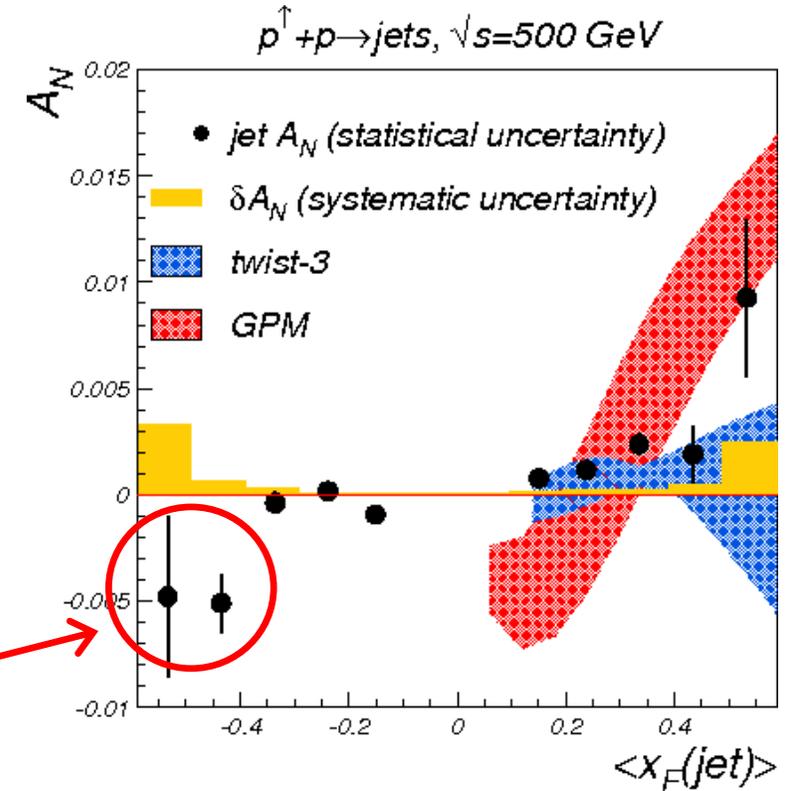
# Data versus Atlas [AUET2B with CTEQ10] Tune



- This is AUET2B-CT6L, replacing the PDF by CTEQ10 [which differs from CTEQ6 for low-x gluons]
- Reasonable description of inclusive jet data
- ~20% overprediction of dijet data

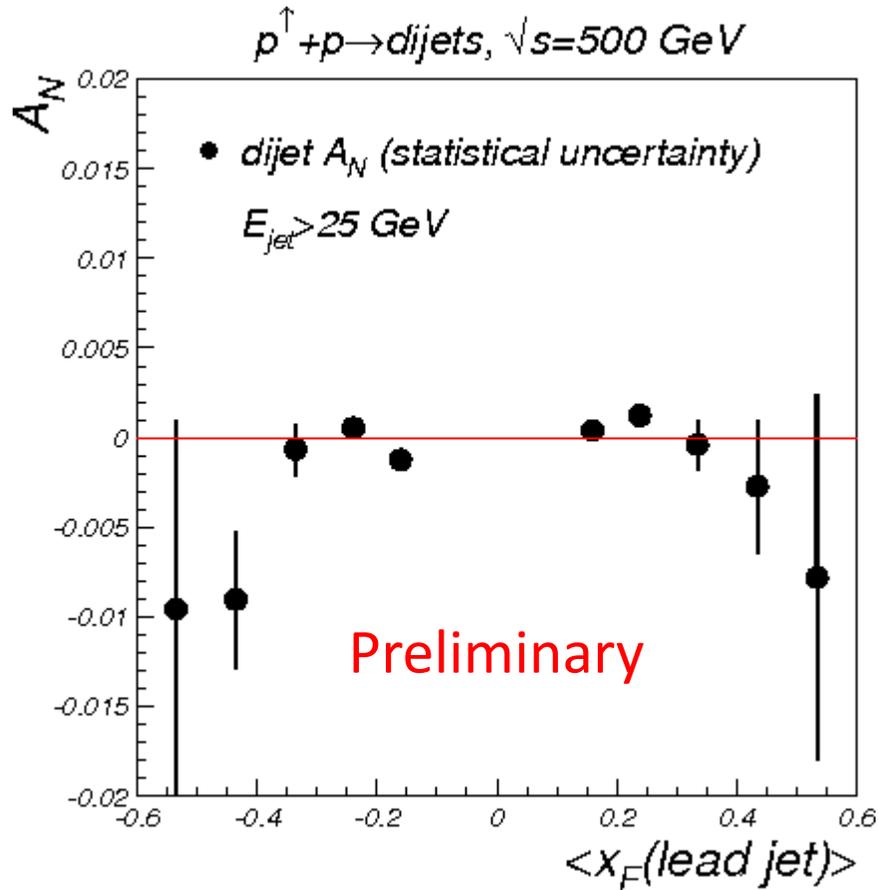
# Are there transverse SSA for dijets?

- Relevant because dijets are the dominant background for forward DY production of dileptons
- Important for low-x physics [e.g. PRD 89 (2014) 074050 and PRD 89 (2014) 034029], in that forward dijets emphasize low-x whereas inclusive jets involve a broad distribution of Bjorken x



Are these hints of non-zero  $A_N$  from low-x physics?

# First Look at Dijet Analyzing Power

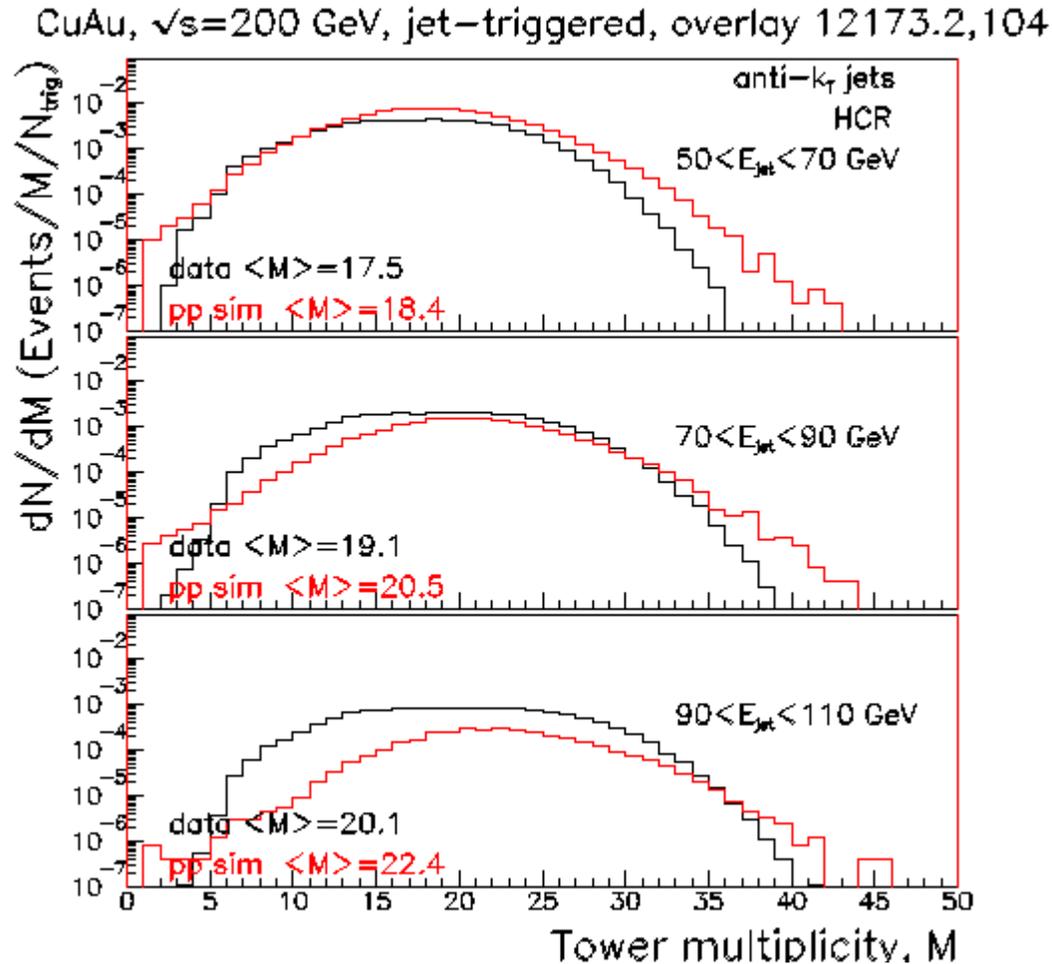


- Identify jet-like clusters with the anti-kT algorithm with  $R_{\text{jet}}=0.7$  [identical to the inclusive jet analysis]. Select events with  $>1$  valid jet. Each valid jet is with  $\eta_{\text{jet}}, \phi_{\text{jet}}$  acceptance selections imposed and has  $E_{\text{jet}} > 25 \text{ GeV}$ .
- The azimuthal angle of the jet pair is used for spin sorting [replacing the azimuthal angle of a single jet, for the inclusive analysis]
- $A_N$  for dijets is computed via the cross-ratio method, so systematic uncertainties are expected to be similar in size to inclusive jet systematic uncertainties
- The results are a function of the lead-jet  $\langle x_F \rangle$ , to facilitate comparison to the inclusive-jet results
- $A_N(\text{dijet})$  is consistent with zero for  $x_F > 0$
- $2.4\sigma$  statistical significance non-zero  $A_N(\text{dijet})$  for  $x_F < 0$

# Conclusions

- Forward jets [ $x_F > 0.1$ ] are consistent with next-to-leading order pQCD for pp collisions at  $\sqrt{s} = 500$  GeV
- Forward jets have a non-zero transverse single spin asymmetry consistent with spin-correlated  $k_T$  [Sivers]
- Forward jet pairs are sensitive to low-x gluons at low scale, with the caveat that multi-parton interactions and partonic showers can be a significant background.
- Forward inclusive jet and dijet cross sections from RHIC can be described by the same PYTHIA tunes as used at the LHC

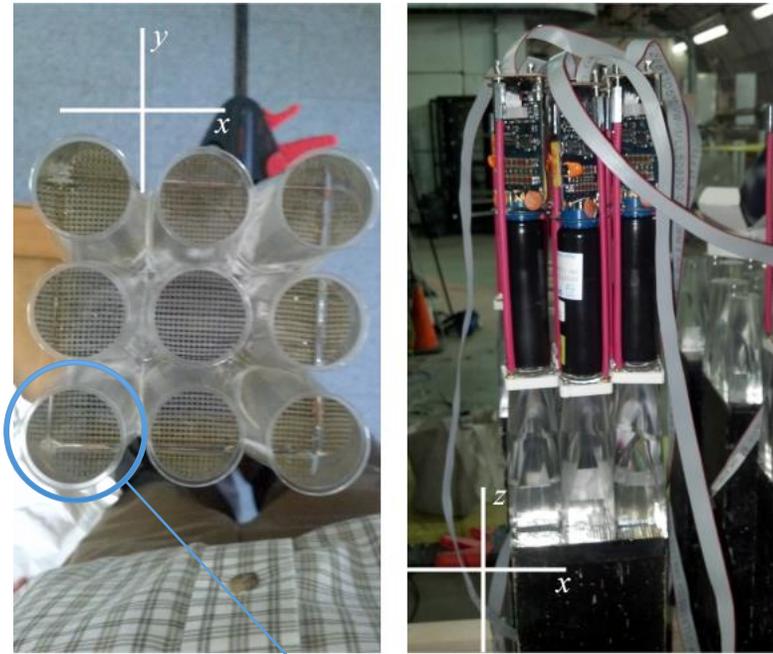
# Outlook – I: Forward Jets in HI Collisions?



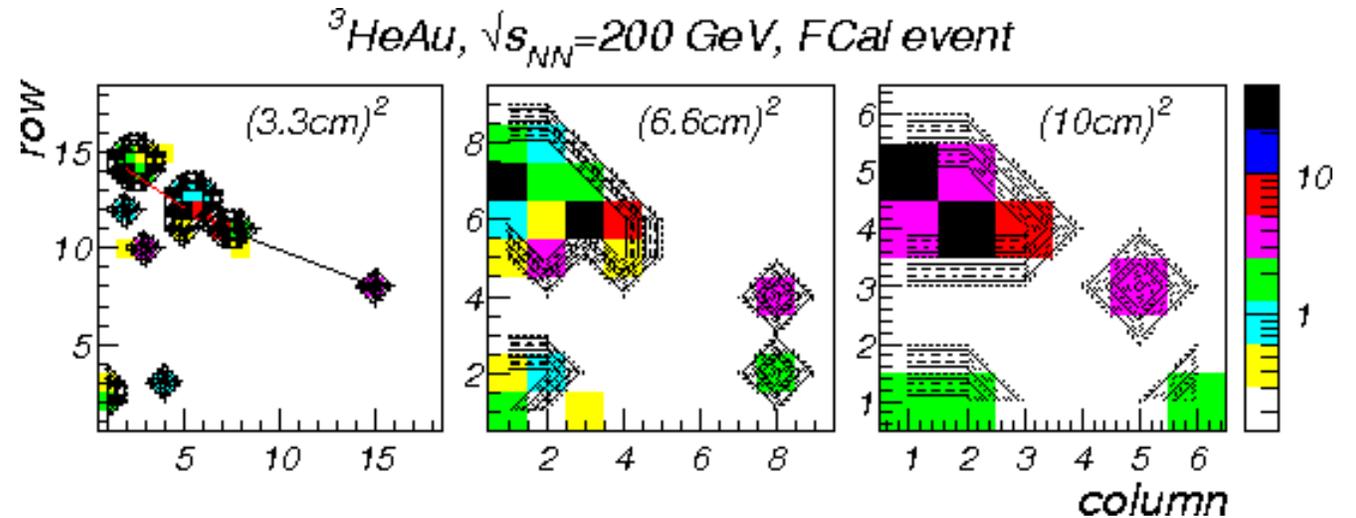
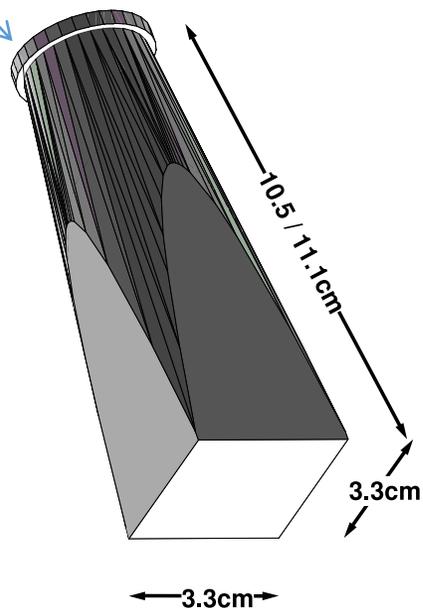
“pp sim” is in reference to PYTHIA 6.222/GEANT for p+p at  $\sqrt{s}=510$  GeV, with jets reconstructed with  $R_{jet}=0.7$ . The comparison of the normalizations to CuAu are irrelevant. The tower multiplicities of the jets in CuAu are comparable to those

- The anti- $k_T$  jet finder developed for  $p+p$  collisions at  $\sqrt{s}=500$  GeV [arXiv:1304.1454] produces reasonable jets for centrality-averaged CuAu, when compensated  $R_{jet}=0.5$  jets [arXiv:1308.4705] are reconstructed.
- A first look was made for the modular HCal, in comparison to p+p PYTHIA/GEANT simulations at  $\sqrt{s}=510$  GeV.
- A modular HCal at STAR could enable rapidity correlation studies in AuAu for at least some of the centrality values.

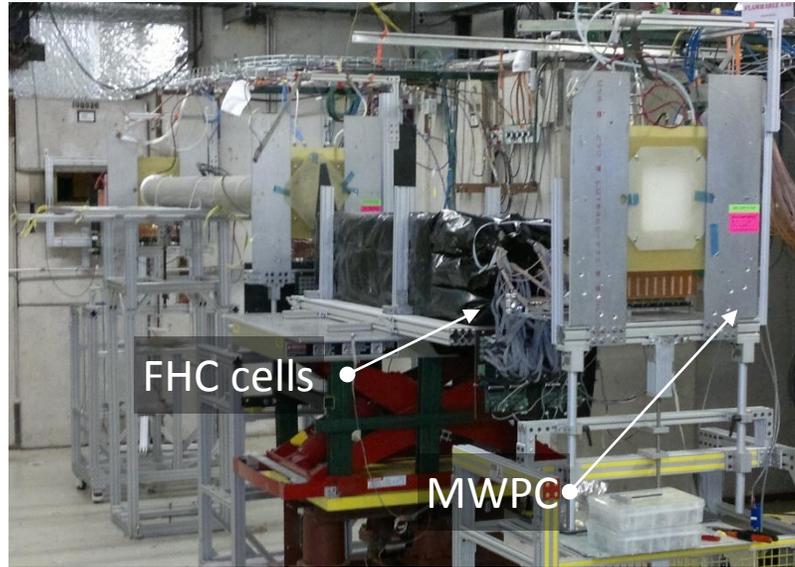
# Outlook-II : Pixelizing AGS-E864 cells



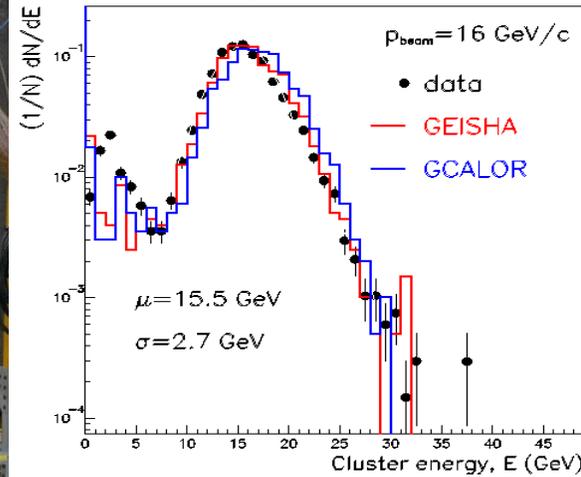
- Transverse resolution is increased by pixelizing
  - $(10\text{ cm})^2$  cells  $\rightarrow$   $3\times 3$  array of  $(3.3\text{ cm})^2$  cells
  - Stack  $(6\text{-cell})^2$  forward calorimeter at STAR in 2014
  - “looking within” cell shows clear structures



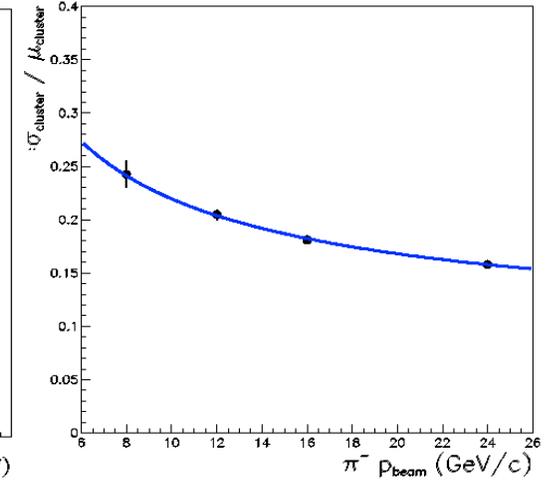
# Outlook III - Results from Fermilab Test Beam Facility-T1064



Cluster energy distribution of center pixel for  $\pi^-$  beam

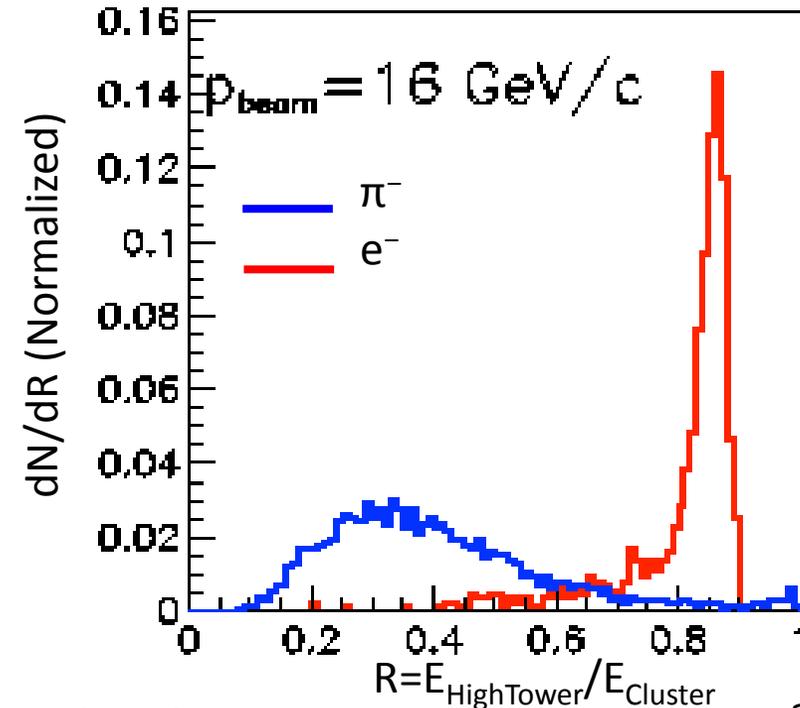


Cluster energy resolution for  $\pi^-$  on central pixel



- 1 GeV ( $\pi, K, p$ ) to 120 GeV  $p$  (resolution  $< 3\%$ )
- Cerenkov Detector (Particle Identification)
- MWPC Tracking System (Beam profile, trigger)

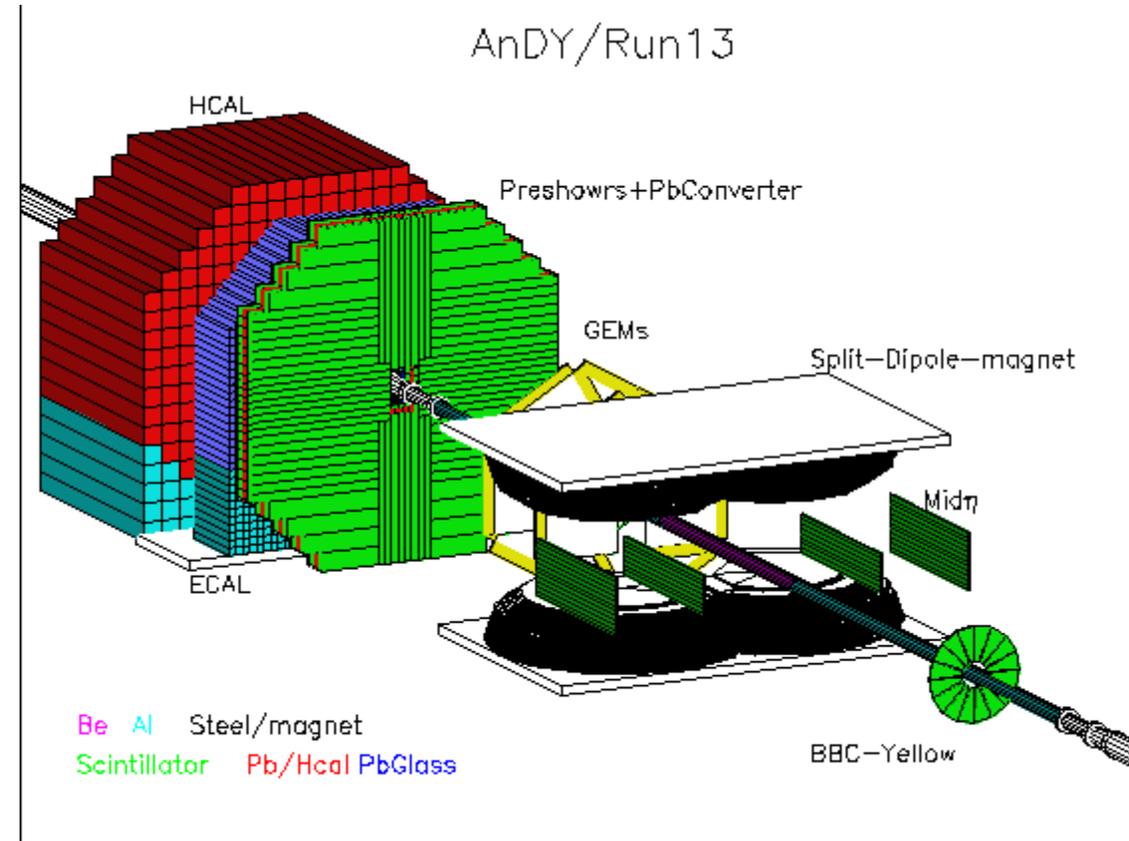
- 3x3 Cells (9x9 pixel) were used
- Studied shower shapes of  $e^-$  &  $\pi^-$  at beam momenta : 8, 12, 16, 24 GeV/c
- Simulations shows **good** agreement with data
- **Shows clean separation between  $e^-$  &  $\pi^-$  shower shapes**



# Backup

# Goal of $A_N$ DY Project

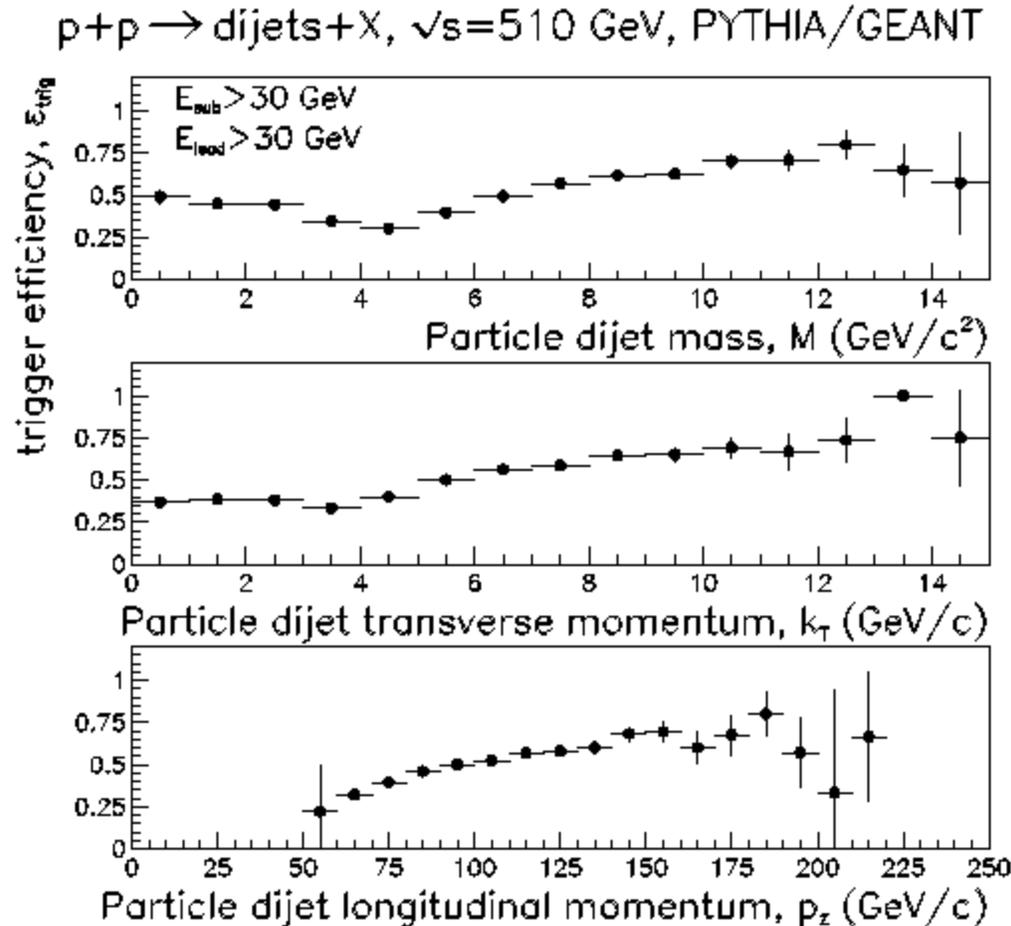
*Measure the analyzing power for forward Drell-Yan production to test the predicted change in sign from semi-inclusive deep inelastic scattering to DY associated with the Sivvers function*



*GEANT model of proposed  $A_N$ DY apparatus (run-13)*

# Dijet Trigger Efficiency

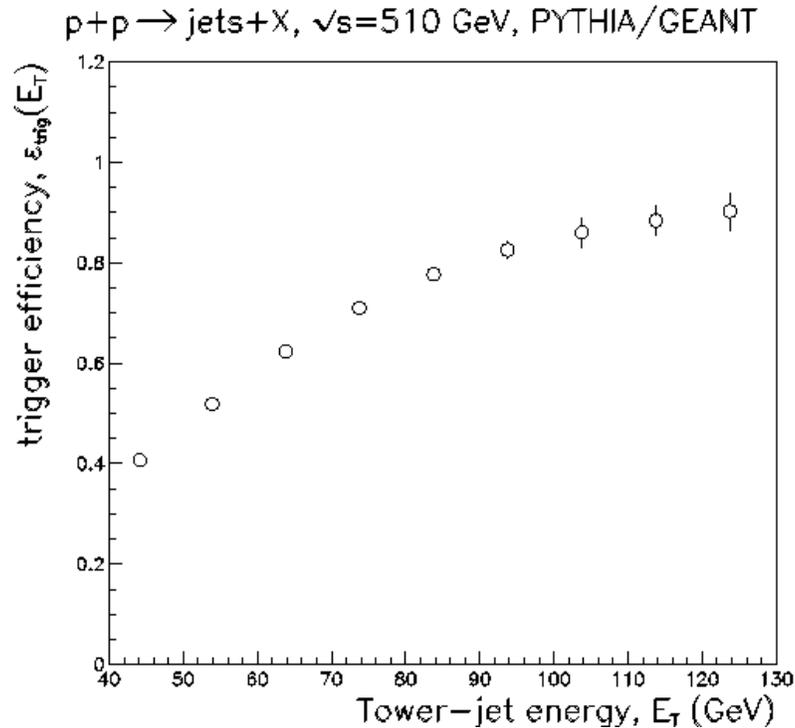
Results prior to convolution integrals



- Analyze the correlation between particle and tower jets from PYTHIA/GEANT simulation to determine dijet finding efficiency
- The dijet efficiencies  $\epsilon_{\text{trig}}(V_{\text{part}})$  [for  $V_{\text{part}}=M, k_T, p_z$ ] behave as expected, becoming larger at larger energies.
- The dijet variable  $M$  is not well measured for  $M < 4$   $\text{GeV}/c^2$ , since small dijet mass corresponds to small dijet opening angle, and the “leading jet” being large leaves little acceptance for the “subleading jet”.
- Consequently, the focus will be on  $M > 4$   $\text{GeV}/c^2$ . [Note: the reducible background for DY with  $M < 4$   $\text{GeV}/c^2$  is dominated by inclusive jet production]

# Jet Cross Section-I

## Definition



This shows an evaluation of the trigger efficiency from PYTHIA/GEANT. Inefficiency results from variation of  $\eta$  for each tower for the extended source for the colliding beams.  $\epsilon_{\text{trig}}$  is checked by extracting cross section from minimum-bias triggers

The jet invariant cross section is:

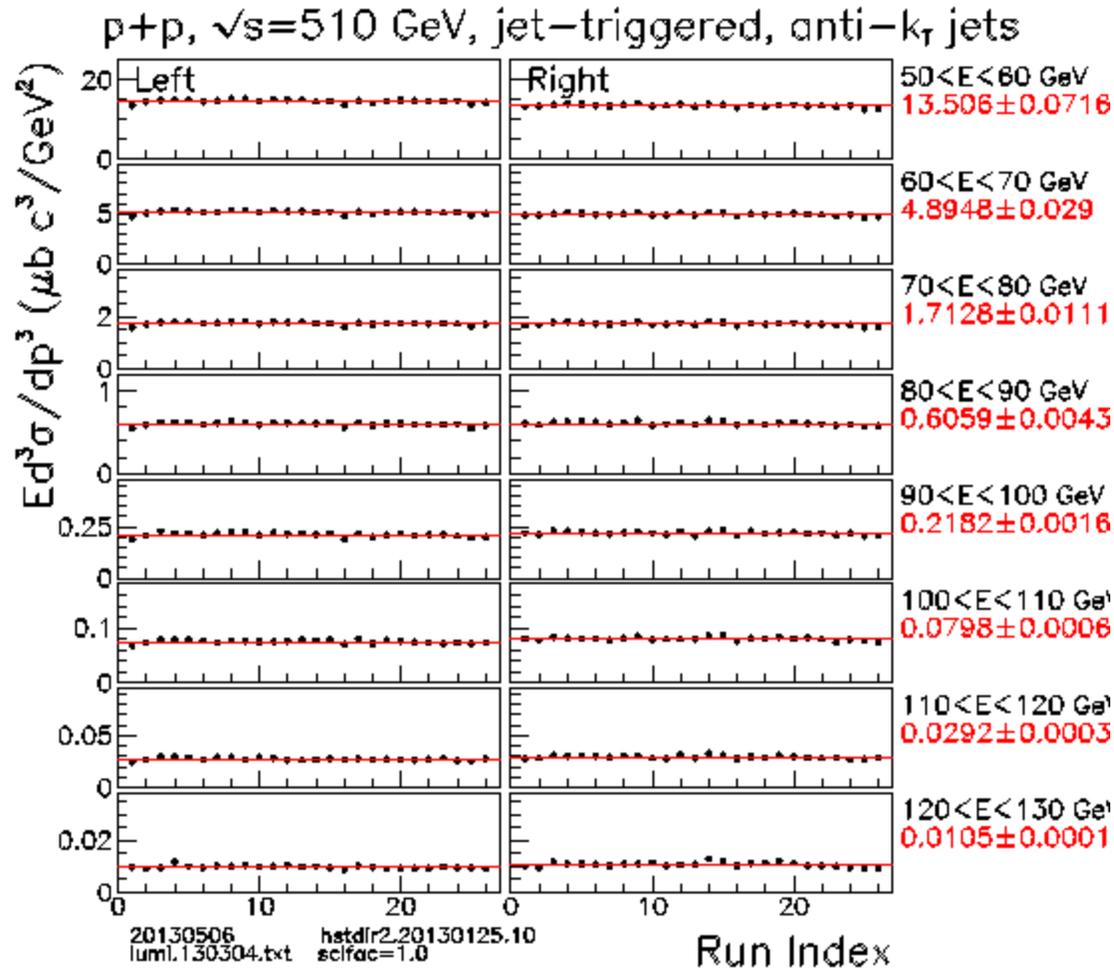
$$E \frac{d^3\sigma}{dp^3} = \frac{N \langle \cosh \eta \rangle}{\epsilon_{\text{trig}} \epsilon_{\text{det}} L_{\text{samp}} \langle p_T \rangle \Delta \eta \Delta \phi \Delta E}$$

where

- N – number of particles detected
- $\epsilon_{\text{trig}}$  – trigger efficiency
- $\epsilon_{\text{det}}$  – detection efficiency
- $L_{\text{samp}}$  – sampled luminosity (time integrated), calibrated by vernier scan
- $\langle \cosh(\eta) \rangle$  - average value of  $\cosh(\eta)$  over the acceptance,  $\Delta \eta \approx \Delta y$
- $\langle p_T \rangle$  - average value of transverse momentum in acceptance
- $\Delta \eta, \Delta \phi$  - specifies the geometry of the acceptance
- $\Delta E$  – width in energy of bin considered

# Jet Cross Section-II

Run Dependence



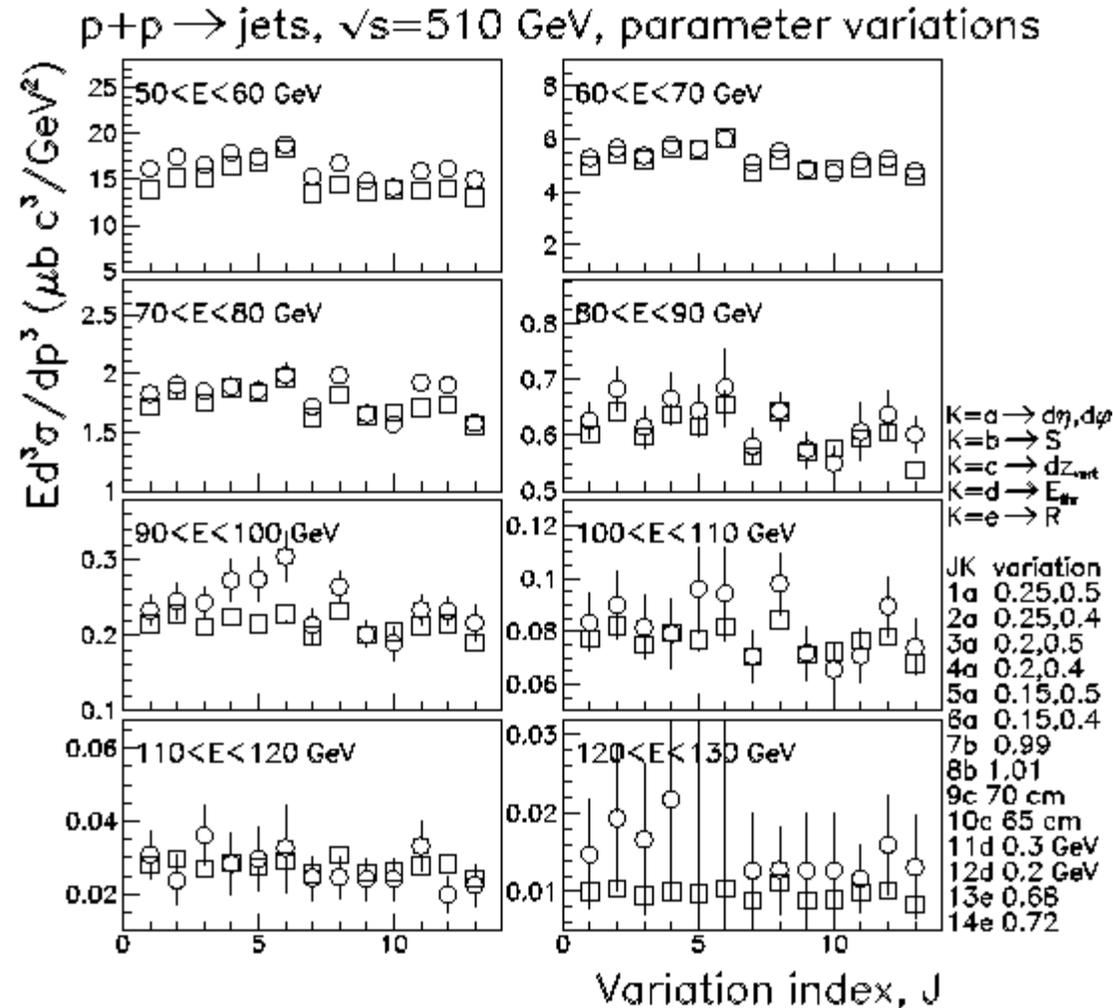
Multiple systematic checks were made for the cross section. This plot shows two:

- Comparison of cross sections from left and right modules
- Stability of cross section with time.

In addition, results were obtained from jet-triggered and minimum-bias triggered samples, to check consistency.

# Jet Cross Section-III

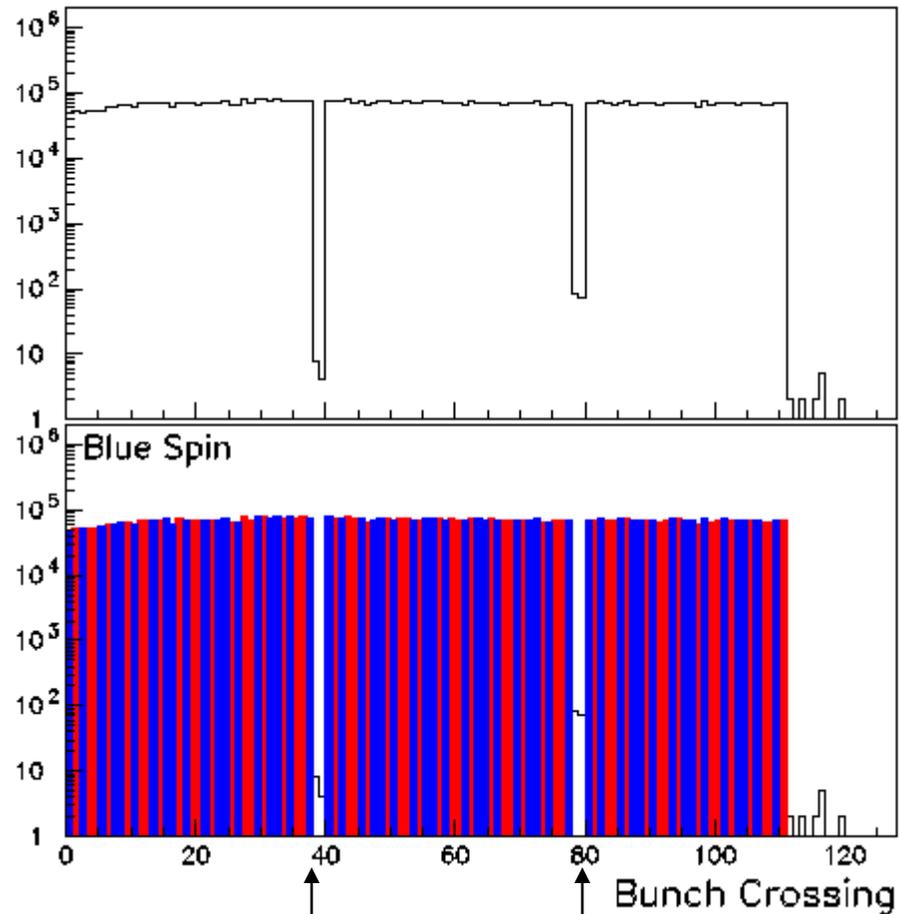
## Systematic Errors



- The stability of the jet cross section was examined as jet-finder (R,  $E_{\text{thr}}$ ); jet acceptance ( $d\eta, d\phi$ ); jet energy scale (S) and vertex selection ( $dz_{\text{vert}}$ ) parameters were varied.
- Results with jet triggered (open squares) and minimum-bias triggered (open circles) events are shown.
- Projections of the resulting cross section on the variation index J result in distributions for each energy bin used to estimate the systematic error for that bin.

# Spin Sorting

p+p,  $\sqrt{s} = 500$  GeV, run=15457, Red=+/Blue=-



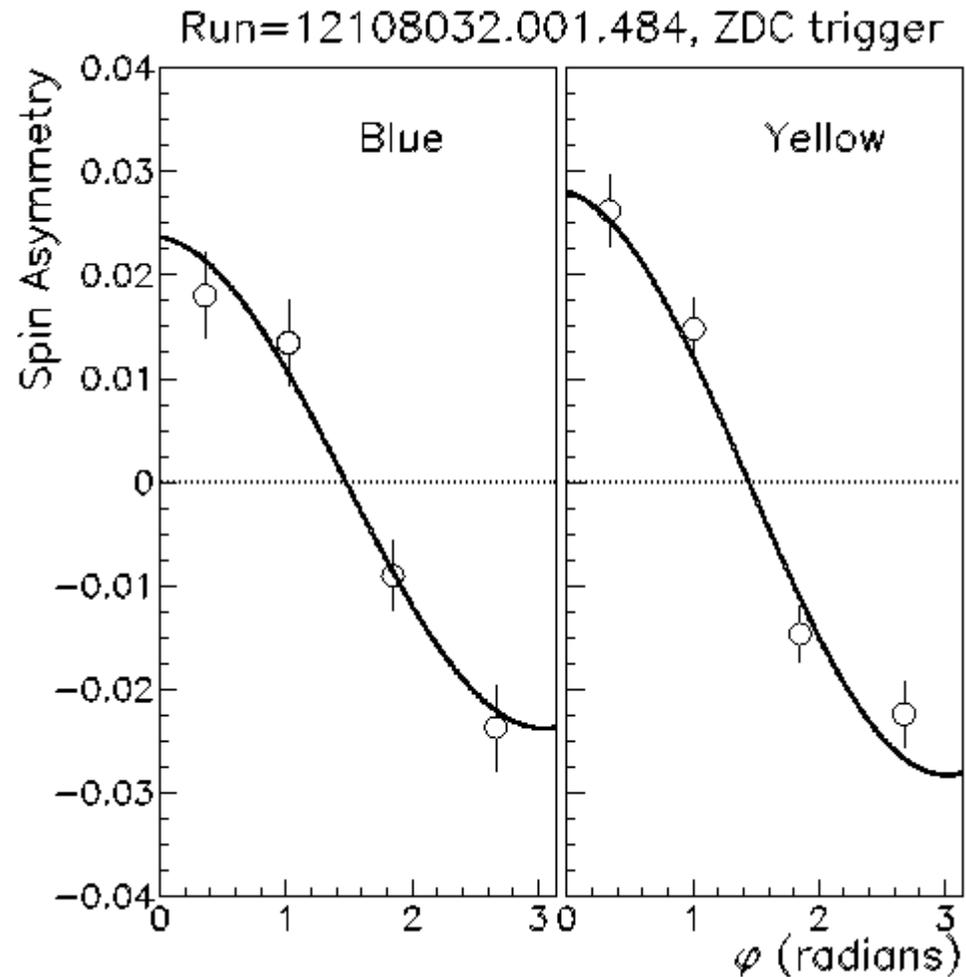
· RHIC has a pattern of polarization directions injected for each fill.

· Polarization for colliding beams is established by counting (C) the 9.38 MHz clock, and identifying specific bunch crossings by  $B = \text{mod}(C, 120)$

· Polarization pattern for a fill is communicated from the accelerator to the experiments.

· Bunch counter distributions also assess single-beam backgrounds

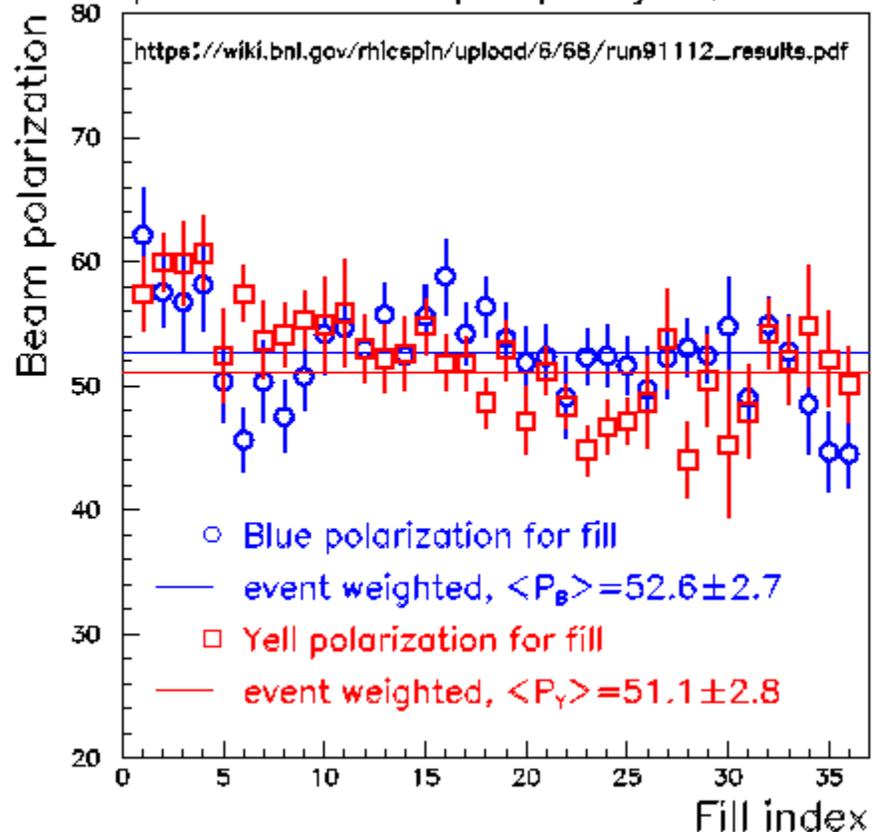
# Spin Direction



- The analyzing power for forward neutron production [ $A_N(n)$ ] has been measured to be positive [PLB **650** (2007) 325]
- $A_N(n)$  is measured with zero-degree calorimeters [NIM A **499** (2003) 433], and provides colliding beam experiments with a local polarimeter.
- Confirm the spin direction used for jet measurements by measuring  $A_N(n)$  concurrent with measuring  $A_N(\text{jet})$ .
- This fixes the sign of  $A_N(\text{jet})$ .

# Beam Polarization

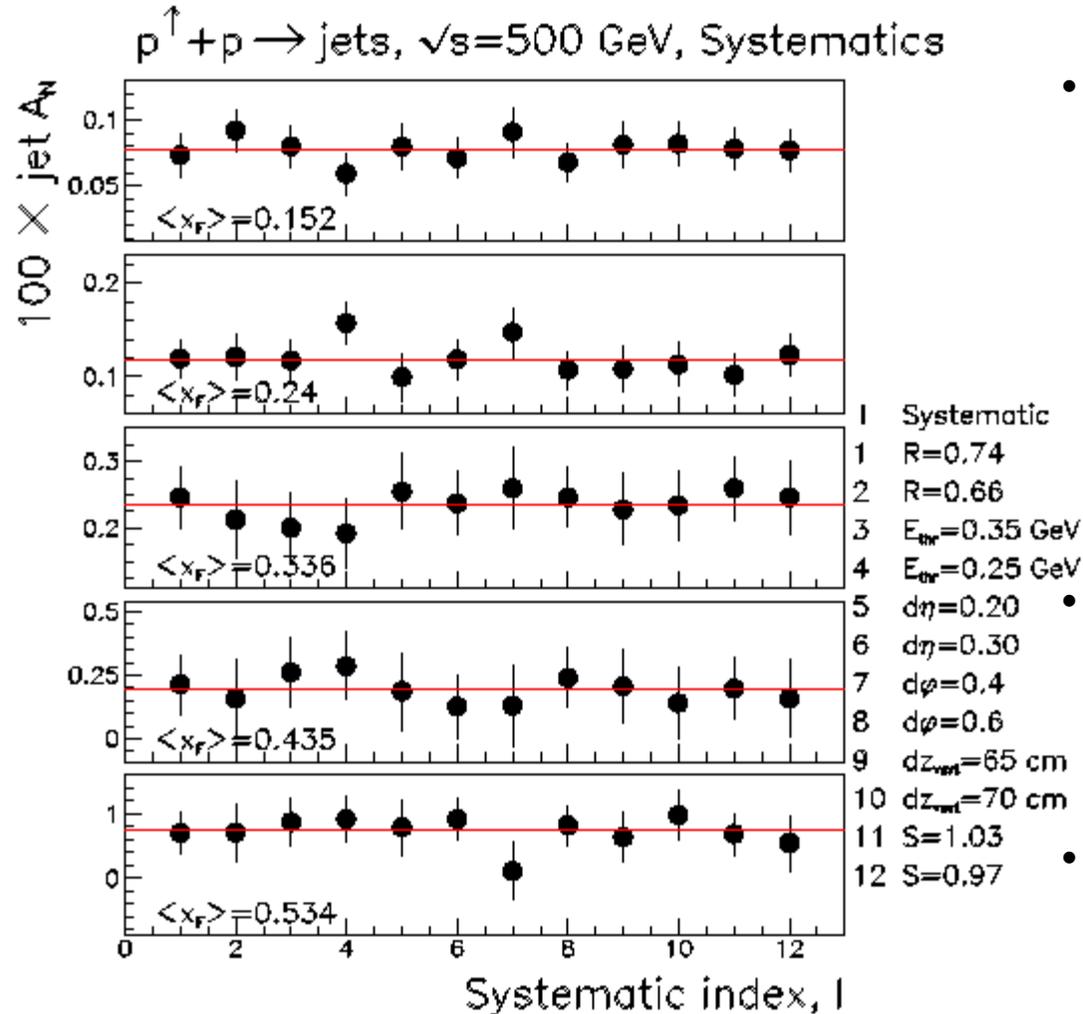
Beam polarization for  $p^\uparrow + p \rightarrow \text{jets}$ ,  $\sqrt{s}=500$  GeV



- Polarization of colliding beams is measured by the polarimeter group [see reference noted in plot].
- Measurements of p+carbon elastic scattering in the Coulomb-nuclear interference region provide a relative polarimeter
- Measurements of p+p elastic scattering in the Coulomb-nuclear interference region from a polarized gas jet target provides an absolute polarimeter

# Jet Analyzing Power

## Definition and Systematics

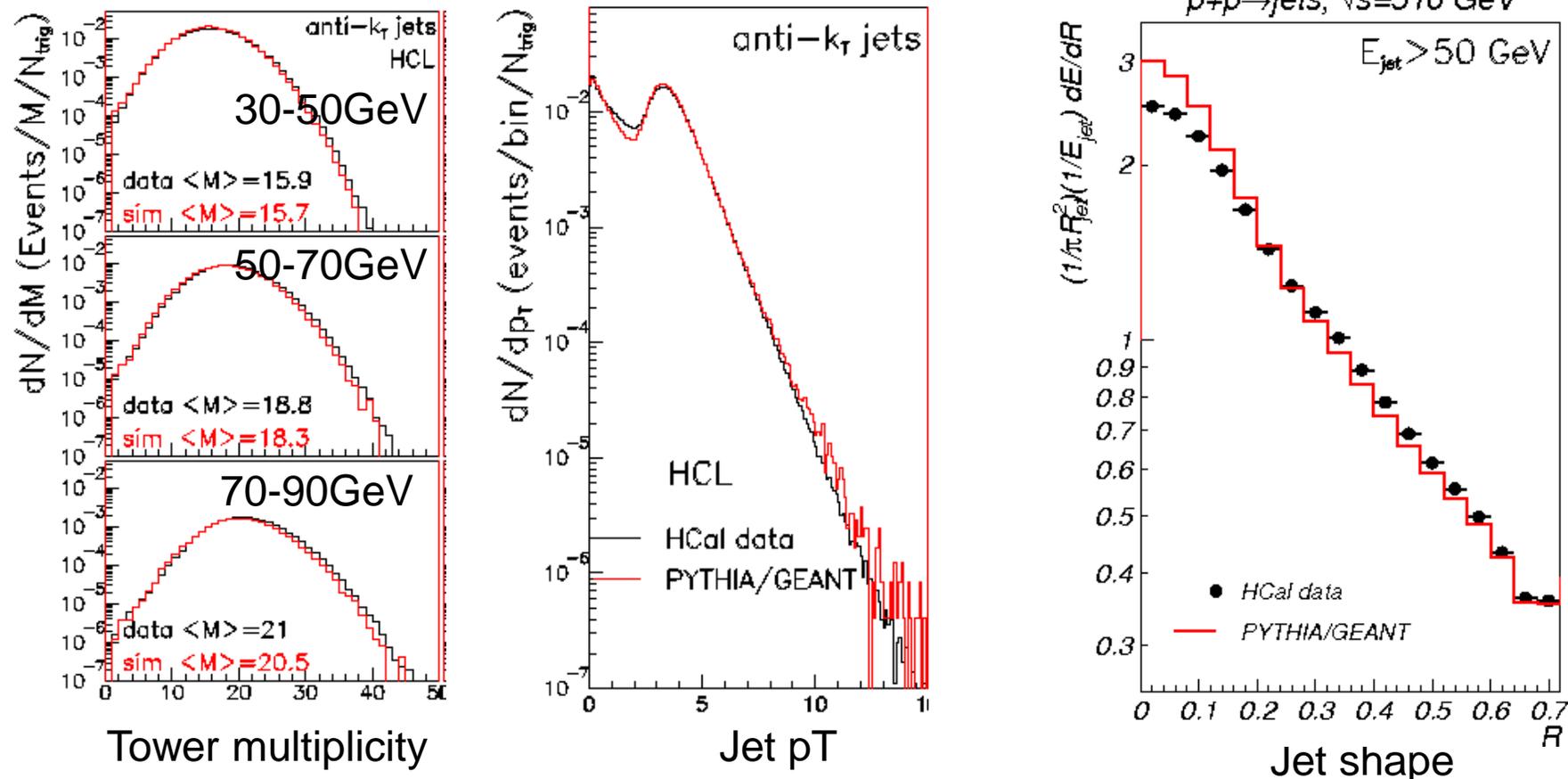


- $A_N(\text{jet})$  exploits mirror (left/right) symmetry of apparatus with spin- $\uparrow$ /spin- $\downarrow$  of colliding beams, via a cross-ratio...

$$A_N = \frac{1}{P} \frac{\sqrt{N_L^\uparrow N_R^\downarrow} - \sqrt{N_L^\downarrow N_R^\uparrow}}{\sqrt{N_L^\uparrow N_R^\downarrow} + \sqrt{N_L^\downarrow N_R^\uparrow}}$$

- Systematic errors for  $A_N(\text{jet})$  are in part computed by varying parameters analogous to manner done for cross section.
- Bottom line is that  $A_N(\text{jet})$  is statistics limited, because of cancellation of systematic errors from symmetry.

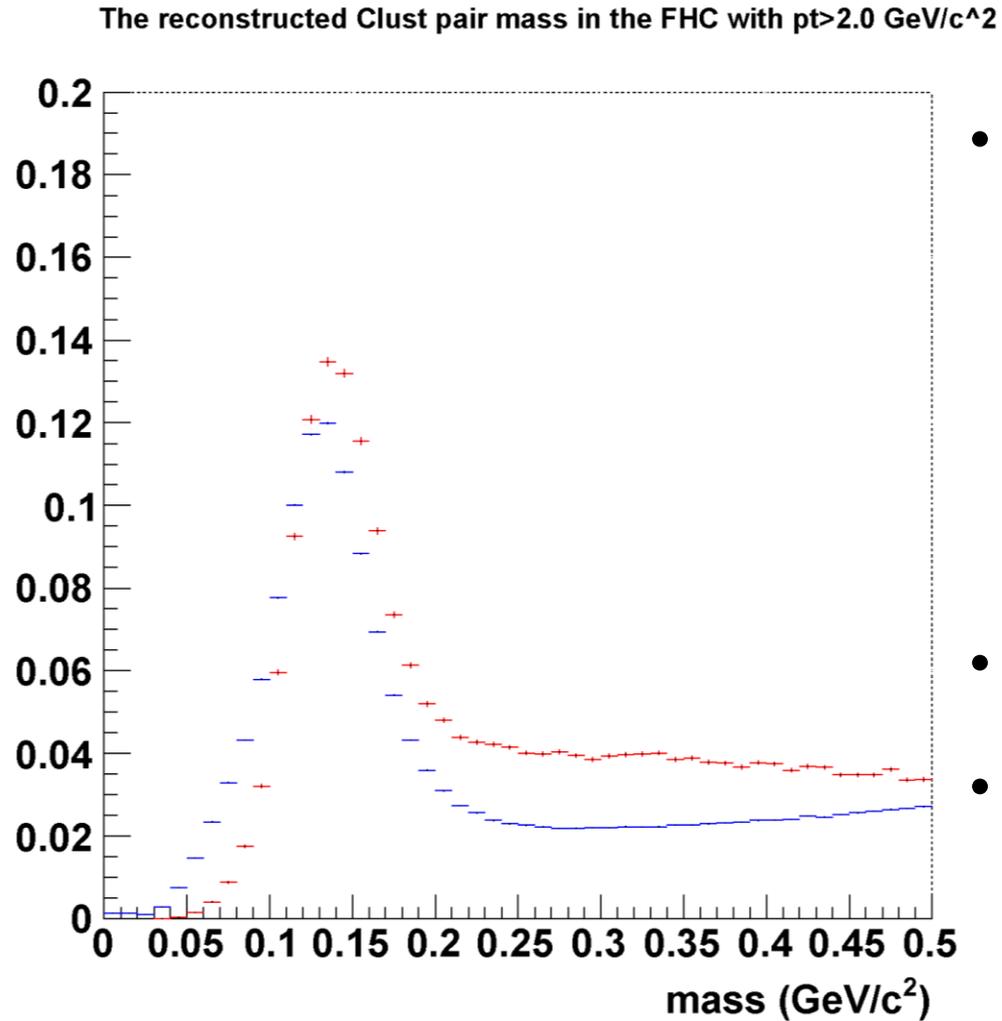
# Comparison of Data to PYTHIA 6.222/GEANT Simulation



- Jet  $p_T$  and  $x_F$  are calculated from anti- $k_T$  cluster with  $R_{\text{jet}} = 0.7$  ignoring mass
- Tower multiplicity agrees with full simulation, meaning particle multiplicity can be deduced
- Given the agreement between data and full simulation, the latter is used for efficiency corrections, e.g.  $\epsilon_{\text{trig}}$  [trigger efficiency (see backup)]

**Jet-triggered data is well described by simulation**

# Resolution



- Compare forward pair mass from
  - (red) <sup>3</sup>He+Au with FHC
  - (blue) d+Au with FMS  
(*Xuan Li - QM12 proceedings*)
- Results are comparable
- Spaghetti calorimetry has demonstrated in electron test beam to give good resolution [Leverington, NIM **A596** (2008) 327]