Extended electroweak precision fits and their implications

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# LoopFest XV

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<u>Thanks</u> to the extended **HEPfit** collaboration.

# Electroweak precision physics in the LHC era

- LHC, Run II: after the Higgs-boson discovery
  - $\hookrightarrow M_H$  becomes a precision electroweak (EW) parameter;
  - $\hookrightarrow$  precision fits now probe consistency of the Standard Model (SM) and can provide indirect evidence of new physics;
  - $\hookrightarrow$  fits can be extended to include Higgs-boson results: rates and distributions  $\longrightarrow$  constraints on anomalous Higgs-boson couplings;
  - $\hookrightarrow$  actual sensitivity depends on experimental precision and theoretical accuracy.
- HEPfit: a global fit of existing electroweak precision data (EWPD) and Higgs-boson observables
  - $\rightarrow\,$  General framework, new features
  - $\rightarrow\,$  Results of EWPD fit, constraints on new physics
- Main results for Higgs-boson couplings and effective interactions
  - $\rightarrow$  In terms of  $\kappa_i$  rescaling factors.
  - $\rightarrow$  In terms of  $C_i$  coefficients of EFT operators.
- Outlook

# LHC Run I has discovered the Higgs, measured its mass and spin . . .



ATLAS+CMS, Phys. Rev. Lett. 114, 191803

 $M_H$  is now among the EW precision observables!

Effects of New Physics can now be more clearly disentangled in both **EW observables** and **Higgs-boson couplings** 

Moreover, from decays  $(H \to VV \text{ and } H \to ff)$ 

- $\rightarrow$  Spin: highly constrained to be s = 0
- $\rightarrow$  Parity: scalar vs pseudoscalar, exploring the tensor structure of decay amplitudes

# Fits of electroweak precision data

- Set of input parameters
  - $\rightarrow$  fixed:  $G_F$ ,  $\alpha$  (best measured)
  - $\rightarrow$  floating:  $M_Z, M_H, m_t, \alpha_s(M_Z), \Delta \alpha_{had}^{(5)}$
- Compute EW precision observables (EWPO), including all known higher-order corrections (in a given renormalization scheme):
   M<sub>W</sub>, Γ<sub>W</sub> (LEP2/Tevatron), Z-pole observables: Γ<sub>Z</sub>, A<sub>f</sub>, ... (LEP/SLD)
- **Perform best fit** and compare with experimental measurement: tension might signal new physics.
- Parametrize new physics effects (ex: *S*, *T*, *U* parameters) and constrain deviations in terms of chosen parameters.
- Several groups:
  - $\rightarrow$  GAPP [Erler]
  - $\rightarrow$  ZFITTER: [Akhundov, Arbuzov, S.Riemann & T.Riemann]
  - $\rightarrow$  Gfitter: [Baak, Cúth, Haller, Hoecker, Kogler, Mönig, Schott, Stelzer]
  - → Now also part of  $| HEPfit | \longrightarrow HEPfit$  Collaboration. For this study: [de Blas, Ciuchini, Franco, Mishima, Pierini, L.R., Silvestrini]

HEPfit developer repository: https://github.com/silvest/HEPfit HEPfit webpage: http://hepfit.roma1.infn.it



In this talk: EW precision physics and Higgs-boson physics [de Blas, Ciuchini, Franco, Mishima, Pierini, L.R., Silvestrini, arXiv:1608.01509]

# The fitting procedure $\rightarrow \texttt{HEPfit}$

- Both electroweak and Higgs observables are calculated as a SM core plus corrections:
  - $\hookrightarrow$  the SM cores include all existing higher order corrections [ $\rightarrow$  loops!]
  - $\hookrightarrow$  the NP corrections are at the lowest order in all SM couplings.
- Experimental results are taken from the most recent published analyses
- The fit procedure uses BAT (Bayesan Analysis Toolkit) with flat priors for all input parameters, and posteriors calculated using a Markov Chain Monte Carlo.

(Caldwell, Kollar, Kröninger, arXiv:0808.2552+ Beaujean, Greenwald, Schulz)

- Stand-alone or library mode to compute observables in a given model:
  - $\hookrightarrow$  Implemented models:
    - $\hookrightarrow$  SM,
    - $\hookrightarrow$  Oblique parameters (S,T,U),  $\varepsilon_i$  parameters, Modified  $Zb\bar{b}$  couplings,
    - $\hookrightarrow$  Modified Higgs couplings ( $\kappa_i$ ), SMEFT (d=6),
    - $\hookrightarrow$  2HDM.
  - $\hookrightarrow$  Implemented observables: EWPO, Flavor ( $\Delta F = 2$ , UT, B-decays).

# Results of SM fit to EW precision data

	Measurement	Result	Prediction	1D Pull	nD Pull
$\alpha_s(M_Z)$	$0.1179 \pm 0.0012$	$0.1180 \pm 0.0011$	$0.1185 \pm 0.0028$	-0.2	
$\delta \alpha_5^{\rm had}(M_Z)$	$0.02750 \pm 0.00033$	$0.02747 \pm 0.00025$	$0.02743 \pm 0.00038$	0.04	
$M_Z$ [GeV]	$91.1875 \pm 0.0021$	$91.1879 \pm 0.0020$	$91.199\pm0.011$	-1.0	
$m_t$ [GeV]	$173.34\pm0.76$	$173.61\pm0.73$	$176.6\pm2.5$	-1.3	
$m_H$ [GeV]	$125.09\pm0.24$	$125.09\pm0.24$	$102.8\pm26.3$	0.8	
$M_W$ [GeV]	$80.385 \pm 0.015$	$80.3644 \pm 0.0061$	$80.3604 \pm 0.0066$	1.5	
$\Gamma_W [\text{GeV}]$	$2.085\pm0.042$	$2.08872 \pm 0.00064$	$2.08873 \pm 0.00064$	-0.2	
$\sin^2 \theta_{\rm eff}^{\rm lept}(Q_{\rm FB}^{\rm had})$	$0.2324 \pm 0.0012$	$0.231464 \pm 0.000087$	$0.231435 \pm 0.000090$	0.8	
$P_{\tau}^{\rm pol} = \mathcal{A}_{\ell}$	$0.1465 \pm 0.0033$	$0.14748 \pm 0.00068$	$0.14752 \pm 0.00069$	-0.4	
$\Gamma_Z [{\rm GeV}]$	$2.4952 \pm 0.0023$	$2.49420 \pm 0.00063$	$2.49405 \pm 0.00068$	0.5	
$\sigma_h^0$ [nb]	$41.540 \pm 0.037$	$41.4903 \pm 0.0058$	$41.4912 \pm 0.0062$	1.3	0.7
$R_\ell^0$	$20.767\pm0.025$	$20.7485 \pm 0.0070$	$20.7472 \pm 0.0076$	0.8	0.7
$A_{ m FB}^{0,\ell}$	$0.0171 \pm 0.0010$	$0.01631 \pm 0.00015$	$0.01628 \pm 0.00015$	0.8	
$\mathcal{A}_{\ell}$ (SLD)	$0.1513 \pm 0.0021$	$0.14748 \pm 0.00068$	$0.14765 \pm 0.00076$	1.7	
$\mathcal{A}_{c}$	$0.670 \pm 0.027$	$0.66810 \pm 0.00030$	$0.66817 \pm 0.00033$	0.02	
$\mathcal{A}_b$	$0.923 \pm 0.020$	$0.934650 \pm 0.000058$	$0.934663 \pm 0.000064$	-0.6	
$A_{ m FB}^{0,c}$	$0.0707 \pm 0.0035$	$0.07390 \pm 0.00037$	$0.07399 \pm 0.00042$	-0.9	1.5
$A_{ m FB}^{ar 0,ar b}$	$0.0992 \pm 0.0016$	$0.10338 \pm 0.00048$	$0.10350 \pm 0.00054$	-2.6	
$R_c^{0}$	$0.1721 \pm 0.0030$	$0.172228 \pm 0.000023$	$0.172229 \pm 0.000023$	-0.05	
$R_b^{0}$	$0.21629 \pm 0.00066$	$0.215790 \pm 0.000028$	$0.215788 \pm 0.000028$	0.7	
$\sin^2 \theta_{\rm eff}^{ee}$	$0.23248 \pm 0.00052$			2.1	
$\sin^2 \theta_{\rm eff}^{\mu\mu}$	$0.2315 \pm 0.0010$			0.07	
$\sin^2 \theta_{\rm eff}^{\ddot{e}e}$	$0.23146 \pm 0.00047$	$0.231464 \pm 0.000087$	$0.231/35 \pm 0.00000$	0.1	
$\sin^2  heta_{ ext{eff}}^{ec{ee},\mu\mu}$	$0.2308 \pm 0.0012$	$0.201404 \pm 0.000007$	$0.201400 \pm 0.000090$	-0.5	
$\sin^2 \theta_{\rm eff}^{\mu\mu}$	$0.2287 \pm 0.0032$			-0.8	
$\sin^2 \theta_{\text{eff}}^{\mu\mu}$	$0.2314 \pm 0.0011$			-0.1	

→ New 2016 world average for  $\alpha_s(M_Z)$  (previously:  $\alpha_s(M_Z) = 0.1185 \pm 0.0005$ ) → Successful comparison with both ZFITTER and Gfitter.

# Good agreement between direct and indirect determination of the values of the input parameters



# Good agreement between direct and indirect determination of the values of the input parameters



## EW precision, example of future projections

#### Present

Observable	Exp. Error	Theor. Error
$M_W  [{\rm MeV}]$	15	4
$\sin^2 \theta_{\rm eff}^l \ [10^{-5}]$	16	4.5
$\Gamma_Z[{ m MeV}]$	2.3	0.5
$R_b \ [10^{-5}]$	66	15

#### Future

Observable	ILC	FCC-ee	CEPC	Theor. Error
$M_W [{ m MeV}]$	3-4	1	3	1
$\sin^2 \theta_{\rm eff}^l \ [10^{-5}]$	1	0.6	2.3	1.5
$\Gamma_Z[MeV]$	0.8	0.1	0.5	0.2
$R_b \ [10^{-5}]$	14	6	17	5-10

[A. Freitas, arXiv:1604.00406]

 $(\delta m_t = 50 \text{ MeV}, \delta \alpha_s = 0.001, \delta M_Z = 2.1 \text{ MeV}, \delta(\Delta \alpha) \simeq 5 \cdot 10^{-5})$ 

ILC  $[e^+e^-, \sqrt{s}=90-500 \text{ GeV}] \rightarrow \text{hep-ph/0106315}, \text{arXiv:1306.6352}$ FCC-ee  $[e^+e^-, \sqrt{s}=90-400 \text{ GeV}] \rightarrow \text{arXiv:1308.6176}$ CEPC  $[e^-p, \sqrt{s}=90-250 \text{ GeV}] \rightarrow \text{IHEP-CEPC-DR-2015-01}$ 

 $\rightarrow$  Theoretical errors may become leading source of error

Limits on beyond SM physics from EW precision data and Higgs-boson data

Parametrizing indirect evidence of new physics beyond the SM (BSM) in a **model-independent** way via

- Oblique corrections (ex.: S,T,U parameters)
- Non-standard  $Zb\bar{b}$  couplings
- Non-standard Higgs couplings
- SM effective field theory (SMEFT)

Oblique parameters, S, T, U [Peskin and Takeuchi, Phys. Rev. D46 (1992) 381] Dominant effects of NP in gauge-boson vacuum polarization corrections,

$$\alpha S = 4e^2 \left[ \Pi_{33}^{NP'}(0) - \Pi_{3Q}^{NP'}(0) \right]$$
  
$$\alpha T = \frac{e^2}{s_W^2 c_W^2 M_Z^2} \left[ \Pi_{11}^{NP}(0) - \Pi_{33}^{NP}(0) \right]$$
  
$$\alpha U = 4e^2 \left[ \Pi_{11}^{NP'}(0) - \Pi_{33}^{NP'}(0) \right]$$

NP contributions to given EWPO (linearized in terms of S, T, U)

$$O = O_{\rm SM} + O_{\rm NP}(S, T, U)$$

 $U \to NP$  contributions to  $M_W$  and  $\Gamma_W$ 

 $U \ll S, T$  in many NP models (linearly realized EWSB)  $\rightarrow U = 0$ 

Equivalently: use  $\varepsilon_{1,2,3,b}$  parameters [Altarelli, Barbieri, Phys. Lett. B253 (1991) 161]





blue shaded areas  $\rightarrow 68\%, 95\%, 99\%$ 

blue shaded areas  $\rightarrow 68\%, 95\%$ 

Projected sensitivity to EW oblique parameters at a glance:



FCCee: several projected runs

	Z pole	WW threshold	HZ threshold	$t\bar{t}$ threshold	above $t\bar{t}$ threshold
$\sqrt{s}  [\text{GeV}]$	90	160	240	350	> 350
$\mathcal{L} [ab^{-1}/yr]$	86	15	3.5	1.0	1.0
Years of run	0.3/2.5	1	3	0.5	3
Events	$10^{12}/10^{13}$	$6 \times 10^7$	$2 \times 10^6$	$2 \times 10^5$	$7.5  imes 10^4$

## Non-standard $Zb\bar{b}$ couplings

Tension in  $A_{\rm FB}^{0,b}$  (pull of EWPO fit  $\rightarrow 2.8\sigma$ )

$$A_{\rm FB}^{0,b} = \frac{3}{4} A_e A_b \ , \ A_f = \frac{2 \operatorname{Re} \frac{g_V^f}{g_A^f}}{1 + (\operatorname{Re} \frac{g_V^f}{g_A^f})^2} \longrightarrow g_i^b = g_{i,\rm SM}^b + \delta g_i^b \ {}_{(i=\rm V,A,L,R)}$$



$$\begin{split} \delta A_{\rm FB}^{0,b}, \delta A_b &\propto g_{L,\rm SM}^b \delta g_R^b - g_{R,\rm SM}^b \delta g_L^b \\ \delta R_b &\propto g_{R,\rm SM}^b \delta g_R^b + g_{L,\rm SM}^b \delta g_L^b \end{split}$$



# Higgs couplings analysis



ATLAS: arXiv:1507.04548

$$\begin{split} \mu &= \sum_{i} w_{i} r_{i} \text{ where} \\ w_{i} &= \frac{[\sigma \times \mathrm{Br}]_{i}}{[\sigma_{\mathrm{SM}} \times \mathrm{Br}_{\mathrm{SM}}]_{i}} \\ r_{i} &= \frac{\epsilon_{i} [\sigma_{\mathrm{SM}} \times \mathrm{Br}_{\mathrm{SM}}]_{i}}{\sum_{j} \epsilon_{j}^{\mathrm{SM}} [\sigma_{\mathrm{SM}} \times \mathrm{Br}_{\mathrm{SM}}]_{j}} \\ \sigma_{i} &= \sigma_{i}^{\mathrm{SM}} + \delta \sigma_{i} \\ \Gamma_{j} &= \Gamma_{j}^{\mathrm{SM}} + \delta \Gamma_{j} \end{split}$$

 $\sigma_i^{\rm SM}, \Gamma_j^{\rm SM} \to {\rm YR} \text{ of HXSWG}$  $\delta \sigma_i \to {\rm FR+Madgraph+Kfactors}$  $\delta \Gamma_j \to {\rm eHdecay}$ 

 $h\gamma\gamma$ : ATLAS(1408.7084), CMS(1407.0558)  $h\tau\tau$ : ATLAS(1501.04943), CMS(1401.5041) hZZ: ATLAS(1408.5191), CMS(1412.8662) hWW: ATLAS(1412.2641,1506.06641), CMS(1312.1129) hbb: ATLAS(1409.6212, 1503.05066), CMS(1310.3687, 1408.1682), CDF (1301.6668), D0 (1303.0823)

### Non-standard Higgs-boson couplings

Minimal assumptions (inspired by strong-dynamics EWSB models):

- $\hookrightarrow$  only one Higgs boson below the cutoff  $\Lambda$ ;
- $\hookrightarrow$  custodial symmetry approximately realized;
- $\hookrightarrow$  corrections from new physics flavor-diagonal and universal;
- $\hookrightarrow$  no NP corrections in Hgg,  $H\gamma\gamma$ ,  $HZ\gamma$  loop-induced couplings.

Ex.: Contino, Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089

$$\mathcal{L}_{\text{eff}} = \frac{v^2}{4} \text{tr} \left( D_{\mu} \Sigma^{\dagger} D^{\mu} \Sigma \right) \left( 1 + 2\kappa_V \frac{h}{v} + \cdots \right) - m_i \bar{f}_L^i \left( 1 + 2\kappa_f \frac{h}{v} + \cdots \right) f_R^i$$

where  $\Sigma(x) = \exp i\sigma^a \chi^a(x)/v \rightarrow \text{longitudinal } W/Z \text{ polarizations.}$ 

Defining:  $\kappa_X = g_X/g_X^{\text{SM}} \text{ (SM} \to \kappa_X = 1),$ 

 $\kappa_V \rightarrow$  rescaling of all hVV couplings  $\kappa_f \rightarrow$  rescaling of all  $hf\bar{f}$  couplings

# Considering both $\kappa_V$ and $\kappa_f$

#### Higgs only

	68%	95%	correl	lation
$\kappa_V$	$1.01 \pm 0.04$	[0.93, 1.10]	1.00	
$\kappa_{f}$	$1.03\pm0.10$	[0.83,  1.23]	0.31	1.00



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#### Higgs+EWPO

	68%	95%	correlation
$\kappa_V$	$1.02\pm0.02$	[0.99, 1.06]	1.00
$\kappa_f$	$1.03\pm0.10$	[0.85,  1.23]	0.14 1.00

## Zooming into $\kappa_V$ and $\kappa_f$ ...



#### Custodial symmetry

 $(\kappa_V \to \kappa_W, \kappa_Z)$ 

#### Higgs only

	68%	95%	correlation		L
$\kappa_W$	$1.00\pm0.05$	[0.89, 1.10]	1.00		
$\kappa_Z$	$1.07\pm0.11$	[0.85, 1.27]	-0.17	1.00	
$\kappa_f$	$1.01\pm0.11$	[0.80, 1.22]	0.41	-0.14	1.00





$$(\kappa_f \to \kappa_u, \kappa_d, \kappa_l)$$

#### Higgs only

	68%	95%		correl	lation	
$\kappa_V$	$0.97\pm0.08$	[0.80, 1.13]	1.00			
$\kappa_l$	$1.01\pm0.14$	[0.73,  1.30]	0.54	1.00		
$\kappa_u$	$0.97\pm0.13$	[0.73, 1.25]	0.42	0.41	1.00	
$\kappa_d$	$0.91\pm0.21$	[0.48, 1.35]	0.81	0.61	0.77	1.00





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#### Higgs+EWPO

	68%	95%		corre	lation	
$\kappa_V$	$1.02\pm0.02$	[0.99, 1.06]	1.00			
$\kappa_l$	$1.07\pm0.12$	[0.82, 1.32]	0.15	1.00		
$\kappa_u$	$1.01\pm0.12$	[0.79, 1.27]	0.10	0.24	1.00	
$\kappa_d$	$1.01\pm0.13$	[0.76,  1.30]	0.31	0.38	0.78	1.00

#### **Projected sensitivity to** $\kappa_i$ **parameters at a glance**:



### SM Effective Field Theories

Systematic extension of the SM Lagrangian by d > 4 operators,

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_{d>4} \frac{1}{\Lambda^{d-4}} \mathcal{L}_d, \text{ with } \mathcal{L}_d = \sum_i C_i \mathcal{O}_i, \quad [\mathcal{O}_i] = d,$$

including effects in NP on EWPO and SM Higgs-boson coupling, but also allowing for new structures.

Consider:

- $\rightarrow~d=6$  operators only, obeying SM gauge symmetry, L and B conservation
- $\rightarrow$  one Higgs doublet of  $SU(2)_L$ , linearly realized SSB
- → assuming flavor universality: 59 operators [basis by Grzadkowski et al., JHEP 1010 (2010) 085]
- $\rightarrow$  CP even and with at least one Higgs: 27 operators
- $\rightarrow\,$  contributing to the observables considered: 17 operators
- $\rightarrow$  with a specific model in mind: running  $C_i(\Lambda) \rightarrow C_i(\Lambda_{\rm EW})$  more meaningful
- $\rightarrow$  otherwise take  $C_i = C_i(\Lambda_{\rm EW})$ , no running included

$$\mathcal{O}_{HG} = (H^{\dagger}H) G^{A}_{\mu\nu} G^{A\mu\nu}$$
  

$$\mathcal{O}_{HW} = (H^{\dagger}H) W^{I}_{\mu\nu} W^{I\mu\nu}$$
  

$$\mathcal{O}_{HB} = (H^{\dagger}H) B_{\mu\nu} B^{\mu\nu}$$
  

$$\mathcal{O}_{HWB} = (H^{\dagger}\tau^{I}H) W^{I}_{\mu\nu} B^{\mu\nu}$$
  

$$\mathcal{O}_{HD} = (H^{\dagger}D^{\mu}H)^{*} (H^{\dagger}D_{\mu}H)$$
  

$$\mathcal{O}_{H\Box} = (H^{\dagger}H)\Box(H^{\dagger}H)$$

bosonic operators

corrections to:  $\rightarrow$ 

- oblique parameters
- *hVV*
- WWZ and  $WW\gamma$

$$\begin{aligned} \mathcal{O}_{HL}^{(1)} &= (H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\overline{L}\gamma^{\mu}L) \\ \mathcal{O}_{HL}^{(3)} &= (H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\overline{L}\tau^{I}\gamma^{\mu}L) \\ \mathcal{O}_{He} &= (H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\overline{e}_{R}\gamma^{\mu}e_{R}) \\ \mathcal{O}_{HQ}^{(1)} &= (H^{\dagger}i\overleftarrow{D}_{\mu}H)(\overline{Q}\gamma^{\mu}Q) \\ \mathcal{O}_{HQ}^{(3)} &= (H^{\dagger}i\overleftarrow{D}_{\mu}H)(\overline{Q}\tau^{I}\gamma^{\mu}Q) \\ \mathcal{O}_{Hu} &= (H^{\dagger}i\overleftarrow{D}_{\mu}H)(\overline{u}_{R}\gamma^{\mu}u_{R}) \\ \mathcal{O}_{Hd} &= (H^{\dagger}i\overleftarrow{D}_{\mu}H)(\overline{d}_{R}\gamma^{\mu}d_{R}) \\ \mathcal{O}_{Hud} &= i(\widetilde{H}^{\dagger}D_{\mu}H)(\overline{u}_{R}\gamma^{\mu}d_{R}) \end{aligned}$$

single-fermionic-vector-currentoperators

- corrections to:  $\rightarrow$ 
  - *hff Vff*

$$\mathcal{O}_{eH} = (H^{\dagger}H)(\bar{L}\,e_R H)$$

$$\mathcal{O}_{uH} = (H^{\dagger}H)(\bar{Q}\,u_R \tilde{H})$$

$$\mathcal{O}_{dH} = (H^{\dagger}H)(\bar{Q}\,d_R H)$$

$$\mathcal{O}_{dH} = (H^{\dagger}H)(\bar{Q}\,d_R H)$$

$$\mathcal{O}_{LL} = (\bar{L}\gamma^{\mu}L)(\bar{L}\gamma^{\mu}L)$$

<u>Notice</u>:  $Vf\bar{f}$  and indirect effects (e.g.  $G_F$ ) strongly constrained by EW precision observables.

Upon SSB, direct effect on Higgs-boson couplings

$$\mathcal{L}_{\text{Higgs}} = \mathcal{L}_{hVV} + \mathcal{L}_{hff} + \mathcal{L}_{hVff} + \mathcal{L}_{hTff}$$

each term contains the interactions to either vector bosons or fermions.

<u>Ex.1</u>:  $\mathcal{L}_{hVV}$  contains all the non-fermionic interactions with the SM vector bosons,

$$\mathcal{L}_{hVV} = h \left( g_{hZZ}^{(1)} Z_{\mu\nu} Z^{\mu\nu} + g_{hZZ}^{(2)} Z_{\nu} \partial_{\mu} Z^{\mu\nu} + g_{hZZ}^{(3)} Z_{\mu} Z^{\mu} - g_{hAA} A_{\mu\nu} A^{\mu\nu} + g_{hZA}^{(1)} Z_{\mu\nu} A^{\mu\nu} + g_{hZA}^{(2)} Z_{\nu} \partial_{\mu} A^{\mu\nu} - g_{hWW}^{(1)} W_{\mu\nu}^{+} W^{-\mu\nu} + \left( g_{hWW}^{(2)} W_{\nu}^{+} D_{\mu} W^{-\mu\nu} + (g_{hWW}^{(2)})^{*} W_{\nu}^{-} D_{\mu} W^{+\mu\nu} \right) + g_{hWW}^{(3)} W_{\mu}^{+} W^{-\mu} + g_{hGG} \operatorname{Tr} \left[ G_{\mu\nu} G^{\mu\nu} \right] \right)$$

where (both directly and indirectly  $\rightarrow G_F$ , field renormalization, ...),

$$C_{HG} \longrightarrow g_{hGG}$$

$$C_{HW} \longrightarrow g_{hWW}^{(1)}$$

$$C_{HW}, C_{HB}, C_{HWB} \longrightarrow g_{hZZ}^{(1)}, g_{hZA}^{(1)}, g_{hAA}^{(1)}$$

$$C_{HD} \longrightarrow g_{hZZ}^{(3)}$$

while Ex. 2:  $\mathcal{L}_{hff}$  contains the interactions with the fermions only:

$$\mathcal{L}_{hff} = h \sum_{f} g_{hff} \overline{f_L} f_R + \text{h.c.}$$

where,

 $\begin{array}{l} C_{eH} \longrightarrow g_{h\tau\tau} \\ C_{uH} \longrightarrow g_{hcc}, g_{htt} \\ C_{dH} \longrightarrow g_{hbb} \end{array}$ 

The corresponding rescaling factors  $\kappa_V = \frac{g_{hVV}}{g_{hVV}^{SM}}$  and  $\kappa_f = \frac{g_{hff}}{g_{hff}^{SM}}$ , are

$$\kappa_{Z} = 1 + \delta_{h} + \frac{1}{2} \frac{v^{2}}{\Lambda^{2}} C_{HD} - \frac{1}{2} \delta_{G_{F}}$$
  

$$\kappa_{W} = 1 + \delta_{h} - \frac{1}{2} (c_{W}^{2} - s_{W}^{2}) (4s_{W}c_{W} \frac{v^{2}}{\Lambda^{2}} C_{HWB} + c_{W}^{2} \frac{v^{2}}{\Lambda^{2}} C_{HD} + \delta_{G_{F}})$$
  

$$\kappa_{f} = 1 + \delta_{h} - \frac{1}{2} \delta_{G_{F}} - \frac{v}{m_{f}} \frac{v^{2}}{\Lambda^{2}} \frac{C_{fH}}{\sqrt{2}}$$

where

 $\delta_h \to \text{NP}$  corrections to h wave-function renormalization  $\delta_{G_F} \to \text{NP}$  corrections to  $G_F$ 

## 95% bounds on coefficients of d=6 interactions

#### $\rightarrow$ Fitting one operator at a time

	Only EW	Only Higgs	EW + Higgs
		1	1
	$O(\Lambda^2)$ (m. r. 2)	$C / \Lambda^2$ (m to 2)	$\alpha$ / $\Lambda^2$ (m ev. 3)
	$C_i/\Lambda^2$ [TeV <sup>-2</sup> ]	$C_i/\Lambda^2$ [TeV <sup>-2</sup> ]	$C_i/\Lambda^2$ [TeV <sup>-2</sup> ]
Operator $(O_i)$	at $95\%$	at $95\%$	at $95\%$
$O_{HG} = (H^{\dagger}H)  G^A_{\mu\nu} G^{A\mu\nu}$		[-0.005,  0.009]	[-0.005,  0.009]
$O_{HW} = (H^{\dagger}H) W^{I}_{\mu\nu} W^{I\mu\nu}$		[-0.033, 0.015]	[-0.033,  0.015]
$O_{HB} = (H^{\dagger}H) B_{\mu\nu} B^{\mu\nu}$		[-0.009, 0.004]	[-0.009, 0.004]
$O_{HWB} = (H^{\dagger}\tau^{I}H) W^{I}_{\mu\nu} B^{\mu\nu}$	[-0.010, 0.004]	[-0.008, 0.017]	[-0.007,  0.005]
$O_{HD} = (H^{\dagger}D^{\mu}H)^* \left(\dot{H^{\dagger}}D_{\mu}H\right)$	[-0.032,  0.006]	[-1.38, 1.35]	[-0.032,  0.005]
$O_{H\square} = (H^{\dagger}H)\square(H^{\dagger}H)$		[-1.12, 1.72]	[-1.12, 1.72]
$\overline{O_{HL}^{(1)} = (H^{\dagger} i \overleftrightarrow{D}_{\mu} H)(\overline{L} \gamma^{\mu} L)}$	[-0.006, 0.011]		[-0.006, 0.011]
$O_{HL}^{(3)} = (H^{\dagger} i \overleftrightarrow{D}_{\mu}^{I} H) (\overline{L} \tau^{I} \gamma^{\mu} L)$	[-0.013,  0.006]	[-0.64,  0.49]	[-0.013,  0.006]
$O_{He} = (H^{\dagger}i\overleftarrow{D}_{\mu}H)(\overline{e}_{R}\gamma^{\mu}e_{R})$	[-0.017,  0.006]		[-0.017,  0.006]
$O_{HQ}^{(1)} = (H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\overline{Q}\gamma^{\mu}Q)$	[-0.025,  0.046]	[-4.3,  1.3]	[-0.025,  0.046]
$O_{HQ}^{(3)} = (H^{\dagger}i\overleftarrow{D}_{\mu}^{I}H)(\overline{Q}\tau^{I}\gamma^{\mu}Q)$	[-0.011,  0.016]	[-0.35,  0.18]	[-0.011,  0.016]
$O_{Hu} = (H^{\dagger} i \overleftrightarrow{D}_{\mu} H) (\overline{u}_R \gamma^{\mu} u_R)$	[-0.069,  0.088]	[-1.9,  2.2]	[-0.069,  0.088]
$O_{Hd} = (H^{\dagger}i\overleftarrow{D}_{\mu}H)(\overline{d}_{R}\gamma^{\mu}d_{R})$	[-0.160,  0.058]	[-6.2, 7]	[-0.160,  0.058]
$O_{eH} = (H^{\dagger}H)(\bar{L}e_RH)$		[-0.053, 0.027]	[-0.053, 0.027]
$O_{uH} = (H^{\dagger}H)(\bar{Q}u_R\tilde{H})$		[-0.350, 0.510]	[-0.350,  0.510]
$O_{dH} = (H^{\dagger}H)(\bar{Q}d_RH)$		[-0.036,  0.086]	[-0.036,  0.086]
$O_{LL} = (\bar{L}\gamma^{\mu}L)(\bar{L}\gamma^{\mu}L)$	[-0.010, 0.023]	[-1.970, 1.260]]	[-0.010, 0.023]

 $\hookrightarrow$ see also Corbett, Eboli, Gonçalves, Gonzales-Fraile, Plehn, Rauch, arXiv:1505.0551

95% bounds on coefficients of d=6 interactions

 $\rightarrow$  Fitting all EW operators at the same time

	One at a time	Combined
Operator $(O_i)$	$C_i/\Lambda^2 \ [{ m TeV^{-2}}]$ at 95%	$C_i/\Lambda^2 \; [{ m TeV^{-2}}]$ at 95%
$\begin{aligned} O_{HL}^{(1)} &= (H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\overline{L}\gamma^{\mu}L)\\ O_{HL}^{(3)} &= (H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\overline{L}\tau^{I}\gamma^{\mu}L)\\ O_{He} &= (H^{\dagger}i\overleftarrow{D}_{\mu}H)(\overline{e}_{R}\gamma^{\mu}e_{R})\\ O_{HQ}^{(1)} &= (H^{\dagger}i\overleftarrow{D}_{\mu}H)(\overline{Q}\gamma^{\mu}Q)\\ O_{HQ}^{(3)} &= (H^{\dagger}i\overleftarrow{D}_{\mu}^{I}H)(\overline{Q}\tau^{I}\gamma^{\mu}Q)\\ O_{Hu} &= (H^{\dagger}i\overleftarrow{D}_{\mu}H)(\overline{u}_{R}\gamma^{\mu}u_{R})\\ O_{Hd} &= (H^{\dagger}i\overleftarrow{D}_{\mu}H)(\overline{d}_{R}\gamma^{\mu}d_{R}) \end{aligned}$	$\begin{bmatrix} -0.006, \ 0.011 \end{bmatrix} \\ \begin{bmatrix} -0.013, \ 0.006 \end{bmatrix} \\ \begin{bmatrix} -0.017, \ 0.006 \end{bmatrix} \\ \begin{bmatrix} -0.025, \ 0.046 \end{bmatrix} \\ \begin{bmatrix} -0.011, \ 0.016 \end{bmatrix} \\ \begin{bmatrix} -0.069, \ 0.088 \end{bmatrix} \\ \begin{bmatrix} -0.160, \ 0.058 \end{bmatrix}$	$\begin{bmatrix} -0.012, 0.036 \\ [-0.064, 0.009] \\ [-0.026, 0.014] \\ [-0.106, 0.070] \\ [-0.189, -0.001] \\ [-0.220, 0.420] \\ [-1.180, -0.150] \end{bmatrix}$
$O_{LL} = (\bar{L}\gamma^{\mu}L)(\bar{L}\gamma^{\mu}L)$	[-0.010, 0.023]	[-0.084, 0.030]

Only 8 combinations of EW operators can be fitted at the same time: drop  $O_{HWB}$  and  $O_{HD}$  [ $\hookrightarrow$  e.g. Falkowski, Riva, arXiv:1411.0669]

## 95% bounds on scale of new physics $\Lambda$

	Only EW	Only Higgs	EW + Higgs
Operator $(O_i)$	$\begin{array}{l} \Lambda \ [\text{TeV}] \\  C_i  = 1 \end{array}$	$\begin{array}{l} \Lambda \ [\text{TeV}] \\  C_i  = 1 \end{array}$	$\Lambda \text{ [TeV]} \\  C_i  = 1$
$O_{} (\mathbf{u}^{\dagger} \mathbf{u}) C^{A} C^{A \mu \nu}$		19	19
$O_{HG} = (\Pi^{\dagger}\Pi) G_{\mu\nu} G^{\dagger}$ $O_{HW} = (H^{\dagger}H) W^{I}_{\mu\nu} W^{I\mu\nu}$		5.9	5.9
$O_{HB} = (H^{\dagger}H) B^{\mu\nu}_{\mu\nu} B^{\mu\nu}$		12	12
$O_{HWB} = (H^{\dagger}\tau^{I}H)W^{I}_{\mu\nu}B^{\mu\nu}$	11	8.2	12
$O_{HD} = (H^{\dagger}D^{\mu}H)^* \left(H^{\dagger}D_{\mu}H\right)$	5.9	0.9	6
$O_{H\Box} = (H^{\dagger}H)\Box(H^{\dagger}H)$		0.8	0.8
$O_{HL}^{(1)} = (H^{\dagger} i \overleftrightarrow{D}_{\mu} H) (\overline{L} \gamma^{\mu} L)$	10	_	10
$O_{HL}^{(3)} = (H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\overline{L}\tau^{I}\gamma^{\mu}L)$	9.4	1.3	9.7
$O_{He} = (H^{\dagger}i \overleftarrow{D}_{\mu} H)(\overline{e}_R \gamma^{\mu} e_R)$	8.2	—	8.2
$O_{HQ}^{(1)} = (H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\overline{Q}\gamma^{\mu}Q)$	5.0	0.5	5.0
$O_{HQ}^{(3)} = (H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\overline{Q}\tau^{I}\gamma^{\mu}Q)$	8.6	1.8	8.7
$O_{Hu} = (H^{\dagger}i D_{\mu}H)(\overline{u}_R \gamma^{\mu} u_R)$	3.5	0.7	3.5
$O_{Hd} = (H^{\dagger}i\overleftarrow{D}_{\mu}H)(\overline{d}_{R}\gamma^{\mu}d_{R})$	2.7	0.4	2.7
$O_{eH} = (H^{\dagger}H)(\bar{L}e_RH)$		4.7	4.7
$O_{uH} = (H^{\dagger}H)(\bar{Q}u_R\tilde{H})$	1.5	1.5	
$O_{dH} = (H^{\dagger}H)(\bar{Q}d_RH)$		3.7	3.7
$O_{LL} = (\bar{L}\gamma^{\mu}L)(\bar{L}\gamma^{\mu}L)$	7.9	0.9	7.9

 $\rightarrow$  For  $|C_i| \simeq 1$  NP is beyond LHC reach, need perturbative  $C_i$ .

# 95% bounds on scale of new physics $\Lambda$ - Present vs Future

	95% present bound on		95% future bound on	
Coefficient	$\frac{C_i}{\Lambda^2} \; [\text{TeV}^{-2}]$	$\begin{array}{l} \Lambda \ [\text{TeV}] \\ (C_i = \pm 1) \end{array}$	$\frac{C_i}{\Lambda^2}  \left[ \text{TeV}^{-2} \right]$	$\begin{array}{l} \Lambda \ [\text{TeV}] \\ (C_i = \pm 1) \end{array}$
$\begin{array}{c} C_{HWB} \\ C_{HD} \end{array}$	$\begin{bmatrix} 0.009,  0.003 \\ 0.027,  0.004 \end{bmatrix}$	$\begin{array}{c} 12 \\ 6.6 \end{array}$	[ 0.0001, 0.0001 ] [ 0.0005, 0.0005 ]	$93 \\ 45$
$C^{(1)}_{HL} \ C^{(3)}_{HL} \ C_{He} \ C^{(1)}_{HQ} \ C^{(3)}_{HQ} \ C_{HQ} \ C_{Hu} \ C_{Hu} \ C_{Hu}$	$\begin{bmatrix} 0.005, 0.012 \\ 0.011, 0.005 \\ 0.015, 0.007 \end{bmatrix}$ $\begin{bmatrix} 0.027, 0.043 \\ 0.011, 0.015 \\ 0.071, 0.081 \end{bmatrix}$ $\begin{bmatrix} 0.14, 0.070 \end{bmatrix}$	$9.9 \\ 10 \\ 8.6 \\ 5.3 \\ 9.1 \\ 3.7 \\ 2.9$	$\begin{bmatrix} 0.0003, 0.0003 \\ 0.0002, 0.0002 \end{bmatrix}$ $\begin{bmatrix} 0.0003, 0.0003 \\ 0.0018, 0.0018 \end{bmatrix}$ $\begin{bmatrix} 0.0018, 0.0018 \\ 0.0005, 0.0005 \end{bmatrix}$ $\begin{bmatrix} 0.0035, 0.0035 \\ 0.0046, 0.0046 \end{bmatrix}$	$56 \\ 70 \\ 58 \\ 24 \\ 44 \\ 17 \\ 15$
$C_{LL}$	[0.0096, 0.023]	7.3	[0.0003, 0.0003]	61

 $\rightarrow$  Precision (×10)  $\rightarrow$  reach  $\Lambda \simeq 100$  TeV

 $\hookrightarrow$  Controlling the theoretical uncertainty will be crucial  $\rightarrow$  parameters, NLO HEFT, ...

# Outlook

- The SM offers a incredibly solid theoretical framework that we can use to extract indications of new physics.
- Indirect evidence of new physics from Higgs-boson and EW precision measurements can come from the synergy between
  - $\rightarrow\,$  accurate theoretical prediction,
  - $\rightarrow\,$  a systematic approach to the study of new effective interactions,
  - $\rightarrow\,$  the intuition and experience of many years of Beyond SM searches!
- Increasing the precision of input parameters could allow to test higher scales of new physics: a factor of 10 in precision could give access to scales as high as 100 TeV.
- Identifying and controlling the main sources of theoretical error become very important for future developments.
- Direct evidence of new physics can boost this process, as the discovery of a Higgs-boson has prompted and guided us in this new era of LHC physics.