

Soft Gluon Resummation for associated $t\bar{t}H$ Production at the LHC

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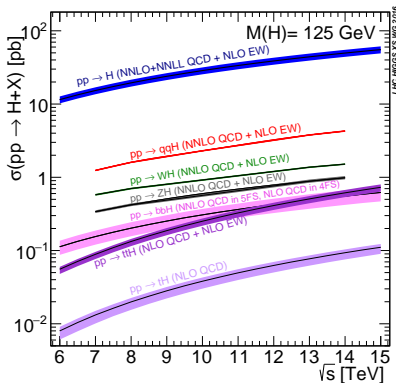
In Collaboration with: Anna Kulesza, Leszek Motyka, Tomasz Stebel

arXiv:1509.02780 and in preparation

Loopfest, 08-15-2016

Importance of $pp \rightarrow t\bar{t}H$

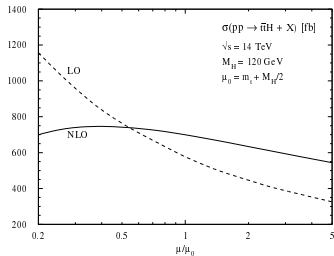
- A Higgs boson with a mass close to 125 GeV
- Precision study needed to determine if it is SM Higgs
- $t\bar{t}H$ direct way to access Yukawa coupling



[LHCHXSWG]

Current status of $pp \rightarrow t\bar{t}H$

- QCD Corrections up to NLO [*Beenakker et al. , '02*] [*Dawson et al. , '02*]
- Matched to parton showers by: aMC@NLO [*Frederix et al. , '11*], PowHel [*Garzelli et al. , '11*], Sherpa [*Hoeche et al., '12*], POWHEG-BOX [*Hartanto et al. , '14*]
- Electroweak correction [*Frixione et al. , '14,'15*][*Zhang, '14*]
- Including top decays [*Denner, Feger, '15*]
- Absolute threshold at NLL [*Kulesza, Motyka, Stebel, VT, '15*]
- Expansion of NNLL in SCET [*Broggio et al., '15*]



[*Beenakker et al. , '02*]

Why resummation for $t\bar{t}H$?

Gains

- NNLO corrections out of reach
- Resummation can help reduce scale uncertainty
- Good process to start:
 - Simple color structure
 - Massive particles \rightarrow no final state collinear divergences

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Pitfalls

- $2 \rightarrow 3$ phase space suppressed near threshold ($\sigma \propto \beta^4$)
- Small corrections from near absolute threshold

Definition of Threshold

Q-approach

Threshold variable $\hat{\tau}_Q = \frac{Q^2}{\hat{s}}$

Q^2 : the invariant mass final state particles

$$1 - \hat{\tau}_Q = 1 - \frac{Q^2}{\hat{s}}$$

$$\sim \frac{\text{energy of the emitted gluons}}{\text{total available energy}}$$

M-approach (absolute threshold)

Threshold variable $\hat{\tau}_M = \frac{M^2}{\hat{s}}$

M : the sum of final state masses

$$1 - \hat{\tau}_M = 1 - \frac{M^2}{\hat{s}}$$

$$\sim \frac{\text{maximum energy of the emitted gluons}}{\text{total available energy}}$$

$\sqrt{\hat{s}}$: the partonic center of mass energy

Logarithms

Q-approach

The IR divergences lead to logarithms:

$$(1 - \hat{\tau}_Q)^{-1-2\epsilon} = -\frac{1}{2\epsilon} \delta(1 - \hat{\tau}_Q) + \left(\frac{1}{1 - \hat{\tau}_Q} \right)_+ - 2\epsilon \left(\frac{\log(1 - \hat{\tau}_Q)}{1 - \hat{\tau}_Q} \right)_+$$

$$\alpha_s^n \left(\frac{\log^m(1 - \hat{\tau}_Q)}{1 - \hat{\tau}_Q} \right)_+$$

In general logarithms of $1 - \hat{\tau}_Q$

M-approach

After integration over $\hat{\tau}_Q$: logarithms of $1 - \hat{\tau}_M$:

$$\alpha_s^n \log^m(1 - \hat{\tau}_M)$$

Logarithms become large in threshold limit: $\hat{\tau} \rightarrow 1$

Mellin Transform

Mellin transform is used with respect to τ (needed for factorization of phase space):

$$\begin{aligned}\tilde{\sigma}_{pp \rightarrow t\bar{t}H}(N) &\equiv \int_0^1 d\tau \tau^{N-1} \sigma_{pp \rightarrow t\bar{t}H}(\tau, \mu_R, \mu_F) \\ &= \sum_{i,j} \tilde{f}_{i/p}(N+1, \mu_F) \tilde{f}_{j/p}(N+1, \mu_F) \tilde{\sigma}_{ij \rightarrow t\bar{t}H}(N, \mu_R, \mu_F)\end{aligned}$$

- $\tilde{f}_{i/p}(N+1, \mu_F)$: Mellin transform with respect to x
- $\tilde{\sigma}_{ij \rightarrow t\bar{t}H}(N, \mu_R, \mu_F)$: Mellin transform with respect to $\hat{\tau}$

$\log^n(1 - \hat{\tau}) \Rightarrow \log^n N$ and threshold $\hat{\tau} \rightarrow 1 \sim N \rightarrow \infty$

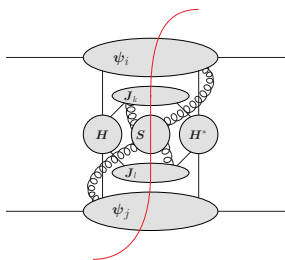
First application for $2 \rightarrow 3$ in Mellin space

General factorization

In general cross section factorizes into:

$$\hat{\sigma}_{ij \rightarrow kl \dots} = H_{ij \rightarrow kl \dots, IJ} \otimes \psi_i \otimes \psi_j \otimes S_{JI} \otimes J_k \otimes J_l \dots$$

- $H_{ij \rightarrow kl, IJ}$ Hard function
- $\psi_{i,j}$ Initial state collinear emission
- $J_{k,l, \dots}$ Final state collinear emission
- S_{JI} Soft emission



Each of these functions is computed through renormalization group equations

Differential formalism

General formalism developed for $2 \rightarrow 2$. [*Kidonakis, Oderda, Serman'98*],
 [*Laenen, Oderda, Serman'98*]

Using a infra-red safe weight (ω) to describe the soft limit of the emission.

For $2 \rightarrow 2$:

- Pair invariant mass $\omega_{\text{PIM}} = z = 1 - Q^2/\hat{s}$
- One particle inclusive $\omega_{1\text{PI}} = s_4 = ((p_1 + p_2 - p_3)^2 - m_4^2)/\hat{s}$

Resummation for differential distributions

$$W = \omega_1 c_1 + \omega_2 c_2 + \omega_s + \omega_k + \omega_l + \dots$$

$$\begin{aligned} & \int dW e^{-NW} \frac{d\sigma_{ij \rightarrow kl\dots}}{d\hat{\Pi}_n} \\ &= \int dW e^{-NW} H_{IJ}(\hat{\Pi}_n) \int d\omega_1 d\omega_2 d\omega_s d\omega_k d\omega_l \dots \\ & \delta(W - \omega_1 c_1 - \omega_2 c_2 - \omega_s - \omega_k - \omega_l - \dots) \\ & \psi_{i/i}(\omega_1) \psi_{j/j}(\omega_2) S_{JI}(\omega_s) J_k(\omega_k) J_l(\omega_l) \dots \\ &= H_{IJ}(\hat{\Pi}_n) \tilde{\psi}_{i/i}(Nc_1) \tilde{\psi}_{j/j}(Nc_2) \tilde{S}_{JI}(N) \tilde{J}_k(N) \tilde{J}_l(N) \dots \end{aligned}$$

- $\omega_{1/2}$ Initial state collinear weights
- ω_s Soft-wide angle weight
- $\omega_{k/l\dots}$ Final state collinear weights

$2 \rightarrow 3$ weights for differential notation

$$pp \rightarrow Q\bar{Q}B + X \quad \Rightarrow Q^2$$

$$z_5 = \frac{\hat{s} - s_{345}}{\hat{s}}$$

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$$pp \rightarrow B + X[Q\bar{Q}] \quad \Rightarrow P_{T,5}, y_5$$

$$v_5 = \frac{(p_1 + p_2 - p_5)^2 - s_{34}}{\hat{s}} = \frac{\hat{s} + \tilde{t}_{15} + \tilde{t}_{25} + m_5^2 - s_{34}}{\hat{s}}$$

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$$pp \rightarrow Q\bar{Q} + X[B] \quad \Rightarrow P_{T,34}, y_{34}, Q_{34}^2$$

$$s_5 = \frac{(p_1 + p_2 - p_3 - p_4)^2 - m_5^2}{\hat{s}} = \frac{\hat{s} + \tilde{t}_{13} + \tilde{t}_{14} + \tilde{t}_{23} + \tilde{t}_{24} + s_{34} - m_5^2}{\hat{s}}$$

Orders of Resummation

Large logarithms $\log N \equiv L$ for $N \rightarrow \infty$

Perturbation needs to be reordered in α_s and L :

$$\tilde{\sigma} \sim \tilde{\sigma}_{LO} \times \mathcal{C}(\alpha_s) \exp [Lg_1(\alpha_s L) + g_2(\alpha_s L) + \alpha_s g_3(\alpha_s L) + \dots]$$

With orders of precision: \Downarrow \Downarrow \Downarrow
LL NLL NNLL
 \Downarrow \Downarrow \Downarrow
 $\alpha_s^n \log^{n+1}(N)$ $\alpha_s^n \log^n(N)$ $\alpha_s^{n+1} \log^n(N)$

Exponential functions are universal for initial state emission

[Kodaira, Trentadue, '82][Sterman, '87][Catani, d'Emilio, Trentadue, '88][Catani, Trentadue, '89]

Color space

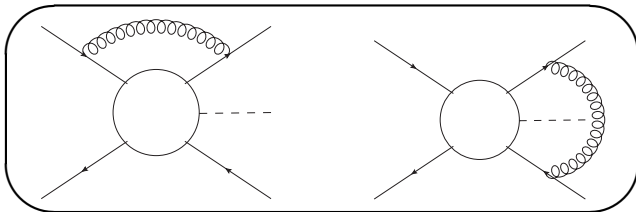
Need to project the matrix element onto a color basis.
Use the s-channel color basis:

$$\begin{array}{ll}
 q\bar{q} & gg \\
 \mathbf{1} : & \delta_{a_2 a_1} \delta_{a_3 a_4} \\
 \mathbf{8} : & T_{a_2 a_1}^D T_{a_3 a_4}^D \\
 & \mathbf{1} : \delta^{A_1 A_2} \delta_{a_3 a_4} \\
 & \mathbf{8}_S : T_{a_3 a_4}^D d^{DA_1 A_2} \\
 & \mathbf{8}_A : iT_{a_3 a_4}^D f^{DA_1 A_2}
 \end{array}$$

Basis for diagonalization of soft anomalous dimension in absolute threshold limit.

Same as for $t\bar{t}$ production

Soft wide-angle



[Kidonakis et al., '97-'01]

$$\begin{aligned} \tilde{S}_{ij \rightarrow kl} \left(\frac{Q}{\mu N} \right) &= \bar{P} \exp \left[\int_{\mu}^{Q/N} \frac{dq}{q} \Gamma_{ij \rightarrow kl}^{\dagger} (\alpha_s (q^2)) \right] \tilde{S}_{ij \rightarrow kl} \\ &\times P \exp \left[\int_{\mu}^{Q/N} \frac{dq}{q} \Gamma_{ij \rightarrow kl} (\alpha_s (q^2)) \right] \end{aligned}$$

Soft anomalous dimension matrix

$$\Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,11}^{(1)} = -\frac{\alpha_s}{2\pi} 2C_F (L\beta_{34} + 1)$$

$$\Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,12}^{(1)} = \frac{\alpha_s}{2\pi} \frac{2C_F}{N_C} \Omega_3 = \frac{C_F}{2N_C} \Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,21}^{(1)}$$

$$\Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,22}^{(1)} = \frac{\alpha_s}{2\pi} [(N_C - 2C_F)(L\beta_{34} + 1) + N_C \Lambda_3 + (8C_F - 3N_C) \Omega_3]$$

$$\Omega_3 = (T_{13} + T_{24} - T_{14} - T_{23}) / 2$$

$$\Lambda_3 = (T_{13} + T_{24} + T_{14} + T_{23}) / 2$$

$$T_{ij} = \log \left(\frac{m_j^2 - t_{ij}}{m_j \sqrt{s}} \right) + \frac{i\pi - 1}{2}$$

with $t_{ij} = (p_i - p_j)^2$ and $\beta_{34}^2 = 1 - (m_3 + m_4)^2 / s_{34}$

Soft anomalous dimension matrix

$$\Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,11}^{(1)} = -\frac{\alpha_s}{2\pi} 2C_F (L\beta_{34} + 1) \quad \text{Agrees with } q\bar{q} \rightarrow Q\bar{Q}$$

$$\Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,12}^{(1)} = \frac{\alpha_s}{2\pi} \frac{2C_F}{N_C} \Omega_3 = \frac{C_F}{2N_C} \Gamma_{q\bar{q} \rightarrow Q\bar{Q}X,21}^{(1)} \quad \text{for } p_5 \rightarrow 0 \text{ and } m_5 \rightarrow 0$$

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with $t_{ij} = (p_i - p_j)^2$ and $\beta_{34}^2 = 1 - (m_3 + m_4)^2 / s_{34}$

For 2 → 2:

$$\begin{array}{cc} \Downarrow & \Downarrow \\ t_{13} = t_{24} & s_{34} = \hat{s} \\ t_{23} = t_{14} & \beta_{34} = \beta \end{array}$$

Soft wide-angle emission absolute threshold

Absolute threshold limit for soft anomalous dimension:

$$2\mathcal{R}e\Gamma_{q\bar{q}\rightarrow Q\bar{Q}B,IJ,\text{thr.}}^{(1)} = \frac{\alpha_s}{\pi} \text{diag}(0, -N_C)$$

$$2\mathcal{R}e\Gamma_{gg\rightarrow Q\bar{Q}B,IJ,\text{thr.}}^{(1)} = \frac{\alpha_s}{\pi} \text{diag}(0, -N_C, -N_C)$$

results in soft wide-angle contribution:

$$\Delta_I^{\text{NLL}} = \exp[h_2(\alpha_s L, -C_I)]$$

With C_I the quadratic Casimir invariant

Hard Matching Coefficient

$$\mathcal{C}(\alpha_s) = 1 + \frac{\alpha_s}{\pi} \mathcal{C}^{(1)} + \dots$$

- Massive final state dipoles [Catani, Dittmaier, Seymour, Trócsányi, '02]
- Virtual contribution from PowHeg-Box [Hartanto, Jäger, Reina, Wackerroth, '15]
confirmed by aMC@NLO [Hirschi, Frederix, Frixione, Garzelli, Maltoni, Pittau, '11]
- Include Coulomb correction $\frac{1}{\beta_{34}}$

Matching to Fixed Order

Resummed Cross Section

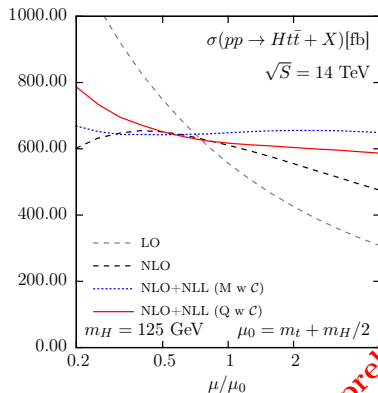
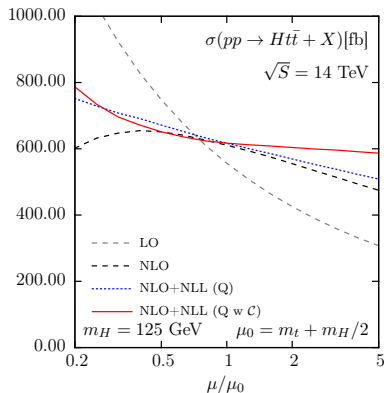
$$\begin{aligned} \sigma^{(\text{NLO+NLL})}(\tau) &= \sigma^{(\text{NLO})}(\tau) \\ &+ \int_{\text{CT}} \frac{dN}{2\pi i} \tau^{-N} \tilde{f}_{g/p}(N+1) \tilde{f}_{g/p}(N+1) \\ &\times \left[\tilde{\sigma}^{(\text{NLL})}(N) - \tilde{\sigma}^{(\text{NLL})}(N)|_{(\text{NLO})} \right] \end{aligned}$$

Matching to fixed order required to avoid double counting.

Results

[Kulesza, Motyka, Stebel, VT, '15 and in preparation]

PDFs used: MMHT2014NLO

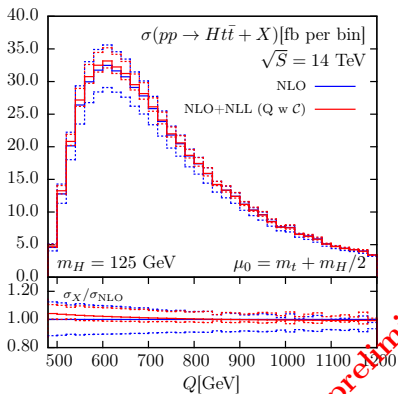


preliminary

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preliminary

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\sqrt{S} [TeV]	NLO [fb]	NLO+NLL M with \mathcal{C}		NLO+NLL Q with \mathcal{C}	
		Value [fb]	K-factor	Value [fb]	K-factor
13	$506^{+5.9\%}_{-9.4\%}$	$537^{+8.2\%}_{-5.5\%}$	1.06	$512^{+5.1\%}_{-6.2\%}$	1.01
14	$613^{+6.2\%}_{-9.4\%}$	$650^{+7.9\%}_{-5.7\%}$	1.06	$619^{+5.2\%}_{-6.4\%}$	1.01

Using 7-point method:

$$(\mu_F/\mu_0, \mu_R/\mu_0) = \{(0.5, 0.5), (0.5, 1), (1, 0.5), (1, 1), (1, 2), (2, 1), (2, 2)\}$$

preliminary

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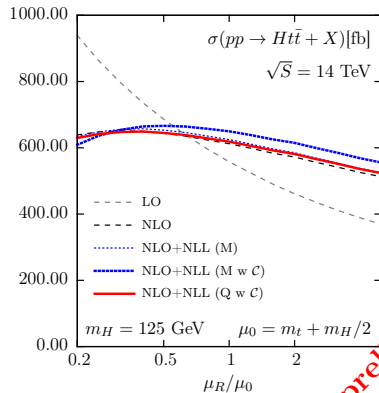
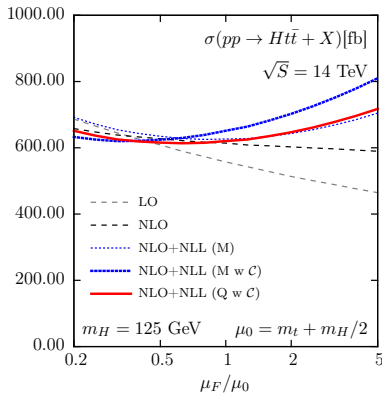
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Thank you for your attention

Backup

[Kulesza, Motyka, Stebel, VT, 'in preparation]

PDFs used: MMHT2014NLO



preliminary

Backup

[Kulesza, Motyka, Stebel, VT, in preparation]

PDFs used: MSTW2008NNLO

