Lepton Universality (Violation, and its consequences)

Diego Guadagnoli LAPTh Annecy (France)

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- Disagreement is rather in muons, that are among the most reliable objects within LHCb

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$$P + 2 + 3 \quad \Rightarrow \quad There seems to be BSM LFNU and the effect is in \mu\mu, not ee$$



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 ${\pmb B} \to {\pmb K}^*\, {\pmb \mu} {\pmb \mu}$ angular analysis: The ${P'}_{_5}$ anomaly

LHCb can perform a fully angular analysis of the decay products in $B \rightarrow K^* \mu \mu$ One can then construct observables with limited sensitivity to form factors. One of such "clean" observables is called P'₅ $B \rightarrow K^* \mu \mu$ angular analysis: The P'₅ anomaly

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$$\frac{BR(B_s \rightarrow \mu \mu)_{exp}}{BR(B_s \rightarrow \mu \mu)_{SM}} = 0.77 \pm 0.20$$

 $BR(B_{s} \rightarrow \mu \mu)_{exp} = (2.8^{+0.7}_{-0.6}) \times 10^{-9}$ [LHCb&CMS full-Run I combination]

 $BR(B_s \rightarrow \mu \mu)_{SM} = (3.65 \pm 0.23) \times 10^{-9}$ [C. Bobeth et al., PRL 14]







More discrepancies: b → c decays

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There are long-standing discrepancies in $b \rightarrow c$ transitions as well.

$$R(D^{(*)}) = \frac{BR(B \rightarrow D^{(*)}\tau\nu)}{BR(B \rightarrow D^{(*)}\ell\nu)} (\text{with } \ell = e,\mu)$$

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- Each of the mentioned effects needs confirmation from Run II to be taken seriously
- Yet, focusing (for the moment) on the $b \rightarrow s$ discrepancies
 - **Q1:** Can we (easily) make theoretical sense of data?
 - **Q2:** What are the most immediate signatures to expect ?

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In fact:

Consider a new, LFNU interaction above the EWSB scale, e.g. with •

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 - (This basis doesn't yet even exist. We are above the EWSB scale.)
- Rotating q and l to the mass eigenbasis generates LFV interactions.

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Or more generally, take the SM plus a minimal mechanism for v masses.

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So, BSM LFNU \implies BSM LFV (i.e. not suppressed by m_{y})



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• Yes we can. Consider the following Hamiltonian

$$H_{\rm SM+NP}(\bar{b} \rightarrow \bar{s} \,\mu \,\mu) = -\frac{4 \,G_F}{\sqrt{2}} \,V_{tb}^* V_{ts} \,\frac{\alpha_{\rm em}}{4 \,\pi} \left[\bar{b}_L \,\gamma^\lambda s_L \cdot \left(C_9^{(\mu)} \,\bar{\mu} \,\gamma_\lambda \mu + C_{10}^{(\mu)} \,\bar{\mu} \,\gamma_\lambda \gamma_5 \mu \right) \right]$$

Concerning Q1: can we easily make theoretical sense of these data?
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- Advocating the same $(V A) \times (V A)$ structure also for the corrections to $C_{9,10}^{SM}$ (in the $\mu\mu$ -channel only!) would account for:
 - R_{κ} lower than 1
 - $B \rightarrow K \mu \mu \& B_s \rightarrow \mu \mu$ BR data below predictions
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A fully quantitative test requires a global fit.

new physics contributions to the Wilson coefficients. We find that the by far largest decrease in the χ^2 can be obtained either by a negative new physics contribution to C_9 (with $C_9^{\text{NP}} \sim -30\% \times C_9^{\text{SM}}$), or by new physics in the $SU(2)_L$ invariant direction $C_9^{\text{NP}} = -C_{10}^{\text{NP}}$, (with $C_9^{\text{NP}} \sim -12\% \times C_9^{\text{SM}}$). A positive NP contribution to C_{10} alone would also improve the fit, although to a lesser extent. [Altmannshofer, Straub, EPJC '15]

For analogous conclusions, see also [Ghosh, Nardecchia, Renner, JHEP '14]

Model example: Glashow, DG, Lane, PRL 2015 4......

As we saw before, all $b \rightarrow s$ data • are explained at one stroke if:

- $C_{9}^{(t)} \approx -C_{10}^{(t)}$ (V A structure) $|C_{9,\text{NP}}^{(\mu)}| \gg |C_{9,\text{NP}}^{(e)}|$ (LFNU)

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- This pattern can be generated from a purely 3rd-generation interaction of the kind
 - $H_{\rm NP} = G \bar{b}'_{L} \gamma^{\lambda} b'_{L} \bar{\tau}'_{L} \gamma_{\lambda} \tau'_{L}$ with $G = 1/\Lambda_{\rm NP}^{2} \ll G_{F}$ expected e.g. in partial-compositeness frameworks

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 - Fields are in the "gauge" basis (= primed)
 - They need to be rotated to the mass eigenbasis

$$b'_{L} \equiv (d'_{L})_{3} = (U_{L}^{d})_{3i} (d_{L})_{i}$$
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• Recalling our full Hamiltonian

$$H_{\rm SM+NP}(\bar{b} \rightarrow \bar{s}\mu\mu) = -\frac{4G_F}{\sqrt{2}} V_{tb}^* V_{ts} \frac{\alpha_{\rm em}}{4\pi} \left[\bar{b}_L \gamma^\lambda s_L \cdot \left(C_9^{(\mu)} \bar{\mu} \gamma_\lambda \mu + C_{10}^{(\mu)} \bar{\mu} \gamma_\lambda \gamma_5 \mu \right) \right]$$



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the shift to the $C_{_9}$ Wilson coeff. in the $\mu\mu$ -channel becomes

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The NP contribution has opposite sign than the SM one if

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$$k_{\rm SM} C_9^{(e)} = k_{\rm SM} C_{9,\rm SM} + \frac{G}{2} (U_L^d)_{33}^* (U_L^d)_{32} |(U_L^\ell)_{31}|^2$$

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the shift to the C_g Wilson coeff. in the $\mu\mu$ -channel becomes

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• On the other hand, in the ee-channel

$$k_{SM} C_9^{(e)} = k_{SM} C_{9,SM} + \frac{G}{2} \left(U_{L,33}^d (U_{L,32}^d) | U_{L,33}^f \right)^2$$
The NP contrib. In the ee-channel is negligible, as

$$||U_{L,31}^f|^2 \ll ||U_{L,32}^f|^2$$



• So, in the above setup

$$R_{K} \approx \frac{|C_{9}^{(\mu)}|^{2} + |C_{10}^{(\mu)}|^{2}}{|C_{9}^{(e)}|^{2} + |C_{10}^{(e)}|^{2}} \simeq \frac{2|C_{10}^{\text{SM}} + \delta C_{10}|^{2}}{2|C_{10}^{\text{SM}}|^{2}}$$







D. Guadagnoli, Lepton universality




























More on LFV model signatures

- Bottom line: we can reasonably expect one of the $B \rightarrow K \ell \ell'$ decays in the 10⁻⁸ ballpark and one of the $B \rightarrow \ell \ell'$ decays in the 10⁻¹⁰ one, namely ~ 5% of $BR(B_s \rightarrow \mu\mu)$
- The most suppressed of the above modes is most likely $B_s \rightarrow \mu e$. . (The lepton combination is the farthest from the 3^{rd} generation, and it's chirally suppressed.)
- What about $B_s \rightarrow \mu e \gamma$? •
 - γ = "hard" photon

(hard = outside of the di-lepton Invariant-mass signal window)

Chiral-suppression factor, of $O(m_{\mu}^{}/m_{Bs}^{})^{2}$ replaced by $\alpha_{em}^{}/\pi$ suppression

DG, Melikhov, Reboud, 2016



D. Guadagnoli, Lepton universality



The interaction advocated in Glashow et al.

$$H_{\rm NP} = G \, \bar{b}'_L \gamma^{\lambda} b'_L \, \bar{\tau}'_L \gamma_{\lambda} \tau'_L$$

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- $K_L^0 \rightarrow e^{\pm} \mu^{\mp}$ $K^+ \rightarrow \pi^+ e^{\pm} \mu^{\mp}$

Exp limits

$$\frac{\Gamma(K_{L}^{0} \rightarrow e^{\pm}\mu^{\mp})}{\Gamma(K^{+} \rightarrow \mu^{+}\nu_{\mu})} < 1.7 \times 10^{-12}$$

$$BNL E871 Collab., PRL 1998$$

$$\frac{\Gamma(K^{+} \rightarrow \pi^{+}\mu^{+}e^{-})}{\Gamma(K^{+} \rightarrow \pi^{0}\mu^{+}\nu_{\mu})} < 3.9 \times 10^{-10}$$

$$BNL E865 Collab., PRD 2005$$





$$\beta^{(K)} = \frac{G(U_L^d)_{32}^* (U_L^d)_{31} (U_L^t)_{31}^* (U_L^t)_{32}}{\frac{4G_F}{\sqrt{2}} V_{us}^*}$$

$$|\beta^{(K)}|^2 = 2.15 \times 10^{-14}$$

(within "model A" of DG, Lane, PLB 2015)

I obtain

$$\frac{\Gamma(K_L^0 \rightarrow e^{\pm} \mu^{\mp})}{\Gamma(K^+ \rightarrow \mu^+ \nu_{\mu})} = \left|\beta^{(K)}\right|^2$$



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(within "model A" of DG, Lane, PLB 2015)

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$$\frac{\Gamma(K_{L}^{0} \rightarrow e^{\pm}\mu^{\mp})}{\Gamma(K^{+} \rightarrow \mu^{+}\nu_{\mu})} = |\beta^{(K)}|^{2} \qquad \qquad \frac{\Gamma(K^{+} \rightarrow \pi^{+}\mu^{\pm}e^{\mp})}{\Gamma(K^{+} \rightarrow \pi^{0}\mu^{+}\nu_{\mu})} = 4 |\beta^{(K)}|^{2}$$

$$\mathbb{BR}(K_{L}^{0} \rightarrow e^{\pm}\mu^{\mp}) \approx 6 \times 10^{-14} \qquad \qquad \mathbb{BR}(K^{+} \rightarrow \pi^{+}\mu^{\pm}e^{\mp}) \approx 3 \times 10^{-15}$$
with
$$\mathbb{BR}(K^{+} \rightarrow \mu^{+}\nu_{\mu}) \approx 64\%$$

$$\Gamma(K^{+})/\Gamma(K_{L}^{0}) \approx 4.2$$

$$\mathbb{BR}(K^{+} \rightarrow \pi^{0}\mu^{+}\nu_{\mu}) \approx 3\%$$











Properly taking into account RG	E running from the NP scale to the scale(s) of the low-energy
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 B → K vv 	See also: Calibbi, Crivellin, Ota, PRL 2015

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Properly taking into account processes, one finds non-triv	RGE running from the NP scale to the vertical constraints from:	he scale(s) of the low-energy
– B → K vv	See also: Calibbi, Crivellin,	Ota, PRL 2015
 Modifications to LEP-m 	easured $Z \rightarrow \ell \ell$ couplings	
 LFU-breaking effects in 	$ au au o \ell' extsf{v} extsf{v}$ (tested at per mil accu	uracy)







Introduce one single leptoquark scalar, transforming as (3, 1, -1/3) under SU(3)_c x SU(2)_L x U(1)_Y



Some models explaining R_{κ} and $R(D^*)$

- Introduce one single leptoquark scalar, transforming as (3, 1, -1/3)• under $SU(3)_c \times SU(2)_L \times U(1)_Y$
- One coupling does all the job: $\ \bar{Q^c}_{Li} \ \lambda_{ij} \ i \ au_2 \ L_{Lj} \ \phi$ •

Bauer-Neubert,

Summer and the second second

PRL 2016









• New non-Abelian strongly interacting sector with $N_{\tau c}$ new "techni-fermions" (TC fermions).



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The basic idea can easily be understood in analogy to QCD:

 The TC-fermion condensate breaks spontaneously a large global symmetry G to a smaller group H, at a scale of about 1 TeV Buttazzo, Greljo, Isidori, Marzocca 1604.03940

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One of the pNGB is the 750-GeV state seen by Atlas & CMS It couples to 2 gluons and decays to 2γ via the anomaly

1604

Buttazzo, Greljo, lsidori, Marzocca

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- To explain the flavor deviations, the mixing needs be hierarchical across generations (largest for the 3rd one, as in partial compositeness)
- Integrating out the vector mesons then yields automatically (among the others) the effective operator

$$H_{\rm NP} = G \, \bar{b}'_{L} \gamma^{\lambda} b'_{L} \, \bar{\tau}'_{L} \gamma_{\lambda} \tau'_{L}$$

proposed in [Glashow, DG, Lane, PRL 15]

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- Early to draw conclusions. But Run II will provide a definite answer
- Timely to propose further tests. One promising direction is that of LFV. Plenty of channels, many of which largely untested.