## Lepton Universality (Violation, and its consequences)

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(1) $+\mathbf{2}+\boldsymbol{3} \quad \Rightarrow \quad$ There seems to be BSM LFNU and the effect is in $\mu \mu$, not ee

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The measured branching fraction is compatible with the previous measurement [3] and lies below SM expectations. For the $q^{2}$ region $1.0<q^{2}<6.0 \mathrm{GeV}^{2} / c^{4}$ the differential branching fraction of $\left(2.58_{-0.31}^{+0.33} \pm 0.08 \pm 0.19\right) \times 10^{-8} \mathrm{GeV}^{-2} c^{4}$ is more than $3 \sigma$ below the SM prediction of $(4.81 \pm 0.56) \times 10^{-8} \mathrm{GeV}^{-2} c^{4}[4,5,32]$.

## $B \rightarrow K^{*} \mu \mu$ angular analysis:

The $P_{5}^{\prime}$ anomaly

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this obs needs be taken cum grano salis
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- What cancels is the dependence on the large- $\mathrm{m}_{\mathrm{b}}$ form factors.
- Debate on the role of
- Subleading terms in $1 / m_{b}$
- cc̄ loops and their resummation


## See:

Jäger \& Martin-Camalich, PRD 2016
Ciuchini et al., 1512.07157

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- Conclusion:

If it's new physics, it is expected to show up elsewhere in the $B \rightarrow K^{*} \mu \mu$ angular analysis.

Run II will tell for sure

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- "large- $\Delta \Gamma_{\mathrm{s}}$ " effect [K. De Bruyn et al., PRL 12]
- soft-photon corr's [Buras, Girrbach, DG, Isidori, EPJC 12]
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- Exp error will go to: $\sim 10 \%$ by end of Run II
$\sim 5 \%$ w/ LHCb upgrade


## More discrepancies: <br> b $\rightarrow$ c decays

There are long-standing discrepancies in $b \rightarrow c$ transitions as well.

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2016: Belle also starts to See an R(D*) excess (semi-lep. tau)

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- Yet, focusing (for the moment) on the $b \rightarrow s$ discrepancies
- Q1: Can we (easily) make theoretical sense of data?
- Q2: What are the most immediate signatures to expect ?
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## Concerning Q2: most immediate signatures to expect

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- Rotating $q$ and $\ell$ to the mass eigenbasis generates LFV interactions.
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- Advocating the same $(V-A) \times(V-A)$ structure also for the corrections to $C_{9,10}{ }^{\text {SM }}$ (in the $\mu \mu$-channel only!) would account for:
- $\quad R_{K}$ lower than 1
- $B \rightarrow K \mu \mu \& B_{s} \rightarrow \mu \mu \quad B R$ data below predictions
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- A fully quantitative test requires a global fit.
new physics contributions to the Wilson coefficients. We find that the by far largest decrease in the $\chi^{2}$ can be obtained either by a negative new physics contribution to $C_{9}$ (with $\left.C_{9}^{\mathrm{NP}} \sim-30 \% \times C_{9}^{\mathrm{SM}}\right)$, or by new physics in the $S U(2)_{L}$ invariant direction $C_{9}^{\mathrm{NP}}=-C_{10}^{\mathrm{NP}}$, (with $C_{9}^{\mathrm{NP}} \sim-12 \% \times C_{9}^{\mathrm{SM}}$ ). A positive NP contribution to $C_{10}$ alone would also improve the fit, although to a lesser extent.
[Altmannshofer, Straub, EPJC '15]
For analogous conclusions, see also [Ghosh, Nardecchia, Renner, JHEP '14]
D. Guadagnoli, Lepton universality


## Model example:

Glashow, DG, Lane, PRL 2015

- As we saw before, all $b \rightarrow s$ data are explained at one stroke if:

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& -C_{9}^{(e)} \approx-C_{10}^{(e)} \quad \text { (V - A structure) } \\
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expected e.g. in partial-compositeness frameworks

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## Explaining $\mathbf{b} \rightarrow \mathbf{s}$ data

- Recalling our full Hamiltonian

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H_{\mathrm{SM}+\mathrm{NP}}(\bar{b} \rightarrow \overline{\mathrm{~s}} \mu \mu)=-\frac{4 G_{F}}{\sqrt{2}} V_{t b}^{*} V_{t s} \frac{\alpha_{\mathrm{em}}}{4 \pi}\left[\bar{b}_{L} \gamma^{\lambda} s_{L} \cdot\left(C_{9}^{(\mu)} \bar{\mu} \gamma_{\lambda} \mu+C_{10}^{(\mu)} \bar{\mu} \gamma_{\lambda} \gamma_{5} \mu\right)\right]
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R_{K} \approx \frac{\left|C_{9}^{(u)}\right|^{2}+\left|C_{10}^{(u)}\right|^{2}}{\left|C_{9}^{(e)}\right|^{2}+\left|C_{10}^{(e)}\right|^{2}} \simeq \frac{2\left|C_{10}^{\mathrm{SM}}+\delta C_{10}\right|^{2}}{2\left|C_{10}^{S M}\right|^{2}}
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- Note as well

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0.77 \pm 0.20=\frac{B R\left(B_{s} \rightarrow \mu \mu\right)_{\exp }}{B R\left(B_{s} \rightarrow \mu \mu\right)_{\mathrm{SM}}}=\frac{B R\left(B_{s} \rightarrow \mu \mu\right)_{\mathrm{SM}+\mathrm{NP}}}{B R\left(B_{s} \rightarrow \mu \mu\right)_{\mathrm{SM}}}=\frac{\left|C_{10}^{\mathrm{SM}}+\delta C_{10}\right|^{2}}{\left|C_{10}^{\mathrm{SM}}\right|^{2}}
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implying (within our model) the correlations

$$
\frac{B R\left(B_{s} \rightarrow \mu \mu\right)_{\text {exp }}}{B R\left(B_{s} \rightarrow \mu \mu\right)_{\mathrm{SM}}} \simeq R_{K} \simeq \frac{B R\left(B^{+} \rightarrow K^{+} \mu \mu\right)_{\mathrm{exp}}}{B R\left(B^{+} \rightarrow K^{+} \mu \mu\right)_{\mathrm{SM}}}
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## LFV model signatures

As mentioned: if $R_{K}$ is signaling BSM LFNU, then expect BSM LFV as well

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\nabla \quad \frac{B R\left(B^{+} \rightarrow K^{+} \mu e\right)}{B R\left(B^{+} \rightarrow K^{+} \mu \mu\right)}=\frac{\left|\delta C_{10}\right|^{2}}{\left|C_{10}^{S M}+\delta C_{10}\right|^{2}} \cdot \frac{\left|\left(U_{L}^{\ell}\right)_{31}\right|^{2}}{\left|\left(U_{L}^{\ell}\right)_{32}\right|^{2}} \cdot 2
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The current $B R(B+\rightarrow K+\mu e)$ limit yields the weak bound

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## Parenthesis: More quantitative LFV predictions

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Reminder:

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the flavor-SU(3) rotations are not all independent. Choosing 3 to be the independent ones allows to predict one SM Yukawa in terms of the other two.


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- Bottom line: we can reasonably expect one of the $B \rightarrow K \ell \ell^{\prime}$ decays in the $10^{-8}$ ballpark and one of the $B \rightarrow \ell \ell^{\prime}$ decays in the $10^{-10}$ one, namely $\sim 5 \%$ of $B R\left(B_{s} \rightarrow \mu \mu\right)$


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## LFV in K decays

- The interaction advocated in Glashow et al.

$$
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can also manifest itself in $K \rightarrow$ (п) $\ell \ell^{\prime}$, for example
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- Exp limits

$$
\begin{aligned}
& \frac{\Gamma\left(K_{L}^{0} \rightarrow e^{ \pm} \mu^{\mp}\right)}{\Gamma\left(K^{+} \rightarrow \mu^{+} v_{\mu}\right)}<1.7 \times 10^{-12} \\
& \frac{\Gamma\left(K^{+} \rightarrow \pi^{+} \mu^{+} e^{-}\right)}{\Gamma\left(K^{+} \rightarrow \pi^{0} \mu^{+} v_{u}\right)}<3.9 \times 10^{-10}
\end{aligned}
$$


D. Guadagnoli, Lepton universality

## LFV in K decays

- Defining the basic quantity

$$
\beta^{(K)}=\frac{G\left(U_{L}^{d}\right)_{32}^{*}\left(U_{L}^{d}\right)_{31}\left(U_{L}^{\ell}\right)_{31}^{*}\left(U_{L}^{\ell}\right)_{32}}{\frac{4 G_{F}}{\sqrt{2}} V_{u s}^{*}}
$$

$$
\left|\beta^{(K)}\right|^{2}=2.15 \times 10^{-14}
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(within "model A" of DG, Lane, PLB 2015)

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\operatorname{BR}\left(K_{L}^{0} \rightarrow e^{ \pm} \mu^{\mp}\right) \approx 6 \times 10^{-14}
\end{gathered}
$$

$$
\begin{aligned}
& \text { with } \\
& \mathrm{BR}\left(K^{+} \rightarrow \mu^{+} v_{u}\right) \approx 64 \% \\
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\beta^{(K)}=\frac{G\left(U_{L}^{d}\right)_{32}^{*}\left(U_{L}^{d}\right)_{31}\left(U_{L}^{\ell}\right)_{31}^{*}\left(U_{L}^{\ell}\right)_{32}}{4 G_{F \mathrm{~V}^{*}}} \quad \square \quad\left|\beta^{(K)}\right|^{2}=2.15 \times 10^{-14}
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I obtain
$\frac{\Gamma\left(K_{L}^{0} \rightarrow e^{ \pm} \mu^{\mp}\right)}{\Gamma\left(K^{+} \rightarrow \mu^{+} v_{\mu}\right)}=\left|\beta^{(K)}\right|^{2}$

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$\Gamma\left(K^{+}\right) / \Gamma\left(K_{L}^{0}\right) \approx 4.2$

$$
\left.\frac{\Gamma\left(K^{+} \rightarrow \pi^{+} \mu^{ \pm} e^{\mp}\right)}{\Gamma\left(K^{+} \rightarrow \pi^{0} \mu^{+} v_{\mu}\right)}=4 \right\rvert\, \beta^{\left.(K)\right|^{2}}
$$



$$
\operatorname{BR}\left(K^{+} \rightarrow \pi^{+} \mu^{ \pm} e^{\mp}\right) \approx 3 \times 10^{-15}
$$

with
$\operatorname{BR}\left(K^{+} \rightarrow \pi^{0} \mu^{+} v_{\mu}\right) \approx 3 \%$

## More signatures



- Being defined above the EWSB scale, our assumed operator $G \bar{b}^{\prime}{ }_{L} \gamma^{\lambda} b^{\prime}{ }_{L} \bar{\tau}^{\prime}{ }_{L} \gamma_{\lambda} \tau^{\prime}{ }_{L}$ must actually be made invariant under $S U(3)_{c} \times S U(2)_{L} \times U(1)_{Y}$


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See: 
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- Thus, the generated structures are all of:
$t^{\prime} t^{\prime} v_{\tau}^{\prime} v_{\tau}^{\prime}$,
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- Thus, the generated structures are all of:
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$b^{\prime} b^{\prime} \tau^{\prime} \tau^{\prime}$,
$t^{\prime} b^{\prime} \tau^{\prime} v_{\tau}^{\prime}$
- After rotation to the mass basis (unprimed), the last structure contributes to $\Gamma\left(b \rightarrow c \tau \bar{v}_{i}\right)$

$$
\square \text { Can explain BaBar }+ \text { Belle }+L H C b \text { deviations on } R\left(D^{(*)}\right)=\frac{B R\left(\bar{B} \rightarrow D^{(*)+} \tau^{-} \bar{v}_{\tau}\right)}{B R\left(\bar{B} \rightarrow D^{(*)+} \ell^{-} \bar{v}_{\ell}\right)}
$$

## But this coin has a flip side

- Properly taking into account RGE running from the NP scale to the scale(s) of the low-energy processes, one finds non-trivial constraints from:
- $B \rightarrow K v v$

See also:
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- Also LFV decays of leptons are generated, and they provide sensitive probes.
E.g.:

$$
\operatorname{BR}(\tau \rightarrow 3 \mu) \& \operatorname{BR}(\tau \rightarrow \mu \rho) \sim 5 \times 10^{-8}
$$

## Some models explaining $R_{K}$ and $R\left(D^{*}\right)$

- Introduce one single leptoquark scalar, transforming as (3, 1, $-1 / 3$ ) under $\operatorname{SU}(3)_{c} \times S U(2)_{L} x U(1)_{Y}$


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- Two insertions (making a tree diag.) contribute to $B \rightarrow D$ Tv
- Four insertions (making a box) contribute to $B \rightarrow K$ e

- With $M_{\phi} \sim 1 \mathrm{TeV}$ and $O(1)$ generation-diagonal couplings, contributions are just the right size


## One model explaining all flavor

 anomalies and the diphoton resonance- New non-Abelian strongly interacting sector with $N_{T C}$ new "techni-fermions" (TC fermions).


## Buttazzo, Greljo Isidori 1604.03940

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The basic idea can easily be understood in analogy to QCD:

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One of the pNGB is the $750-\mathrm{GeV}$ state seen by Atlas \& CMS It couples to 2 gluons and decays to $2 \gamma$ via the anomaly

One model explaining all flavor anomalies and the diphoton resonance:
continued

- There are also vector mesons, like QCD's rho.


## Buttazzo, Greljo, Isidori, Marzeljo,

Their coupling to quarks and leptons explains the flavor anomalies.

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- To explain the flavor deviations, the mixing needs be hierarchical across generations (largest for the $3^{\text {rd }}$ one, as in partial compositeness)
- Integrating out the vector mesons then yields automatically (among the others) the effective operator

$$
H_{\mathrm{NP}}=G \bar{b}_{L}^{\prime} \gamma^{\lambda} b_{L}^{\prime} \bar{\tau}_{L}^{\prime} \gamma_{\lambda} \tau_{L}^{\prime}
$$

proposed in [Glashow, DG, Lane, PRL 15]

## Conclusions

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- Early to draw conclusions. But Run II will provide a definite answer
- Timely to propose further tests. One promising direction is that of LFV. Plenty of channels, many of which largely untested.

