Lepton Universality (Violation, and its consequences)

Diego Guadagnoli LAPTh Annecy (France)

$$R_K = \frac{BR(B^+ \to K^+ \mu \mu)_{[1,6]}}{BR(B^+ \to K^+ e e)_{[1,6]}} = 0.745 \cdot (1 \pm 13\%)$$

LHCb and B factories measured several key $b \rightarrow s$ and $b \rightarrow c$ modes. Agreement with the SM is less than perfect.

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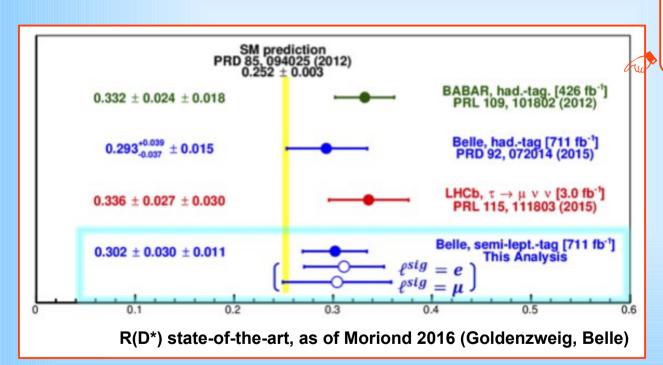
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There are long-standing discrepancies in b \rightarrow c transitions as well.

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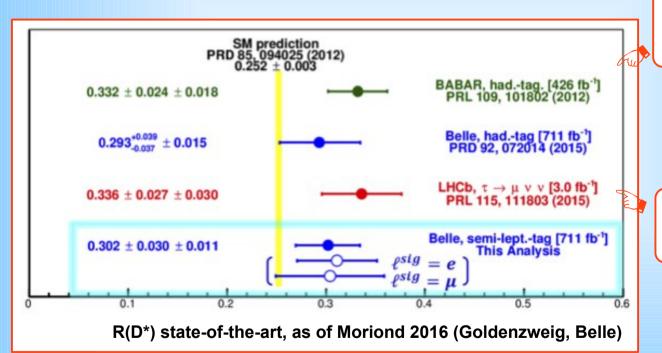
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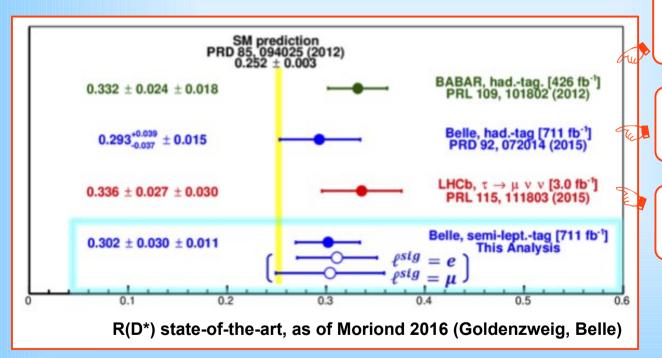


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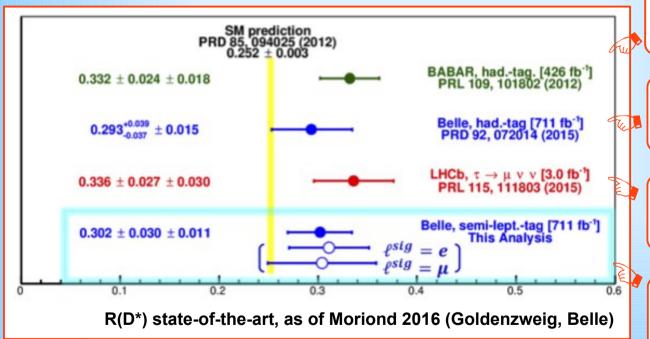
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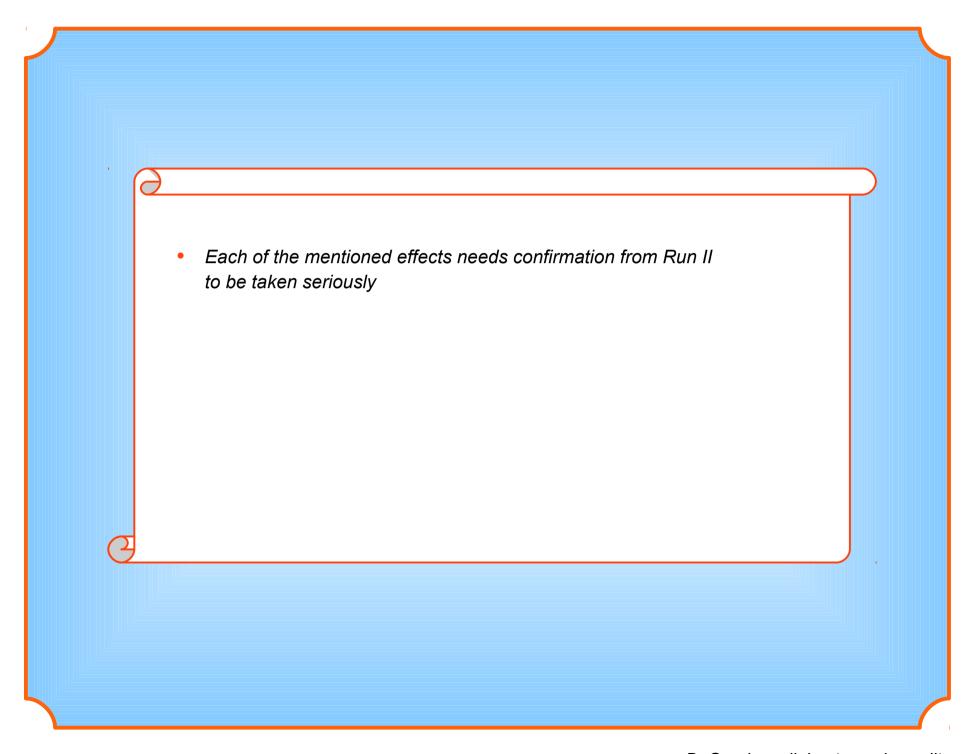


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2016: Belle also starts to see an R(D*) excess (semi-lep. tau's)



Each of the mentioned effects needs confirmation from Run II to be taken seriously Yet, focusing (for the moment) on the $b \rightarrow s$ discrepancies **Q1:** Can we (easily) make theoretical sense of data? **Q2:** What are the most immediate signatures to expect?

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In fact:

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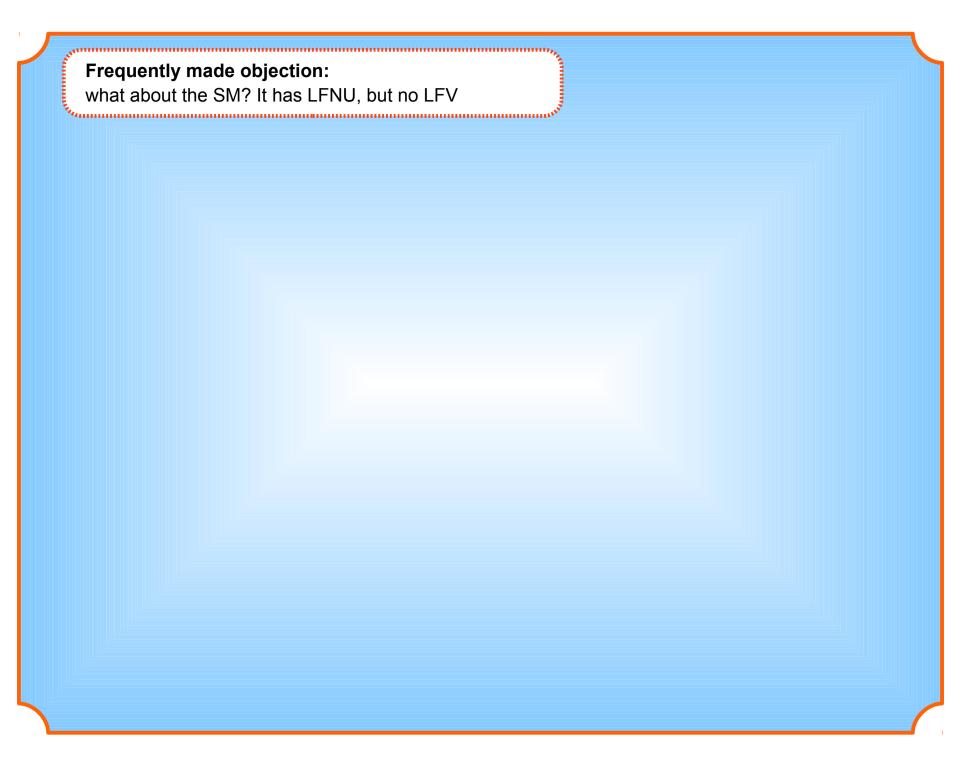
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 Generically, it's not the mass eigenbasis.
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- Rotating q and ℓ to the mass eigenbasis generates LFV interactions.



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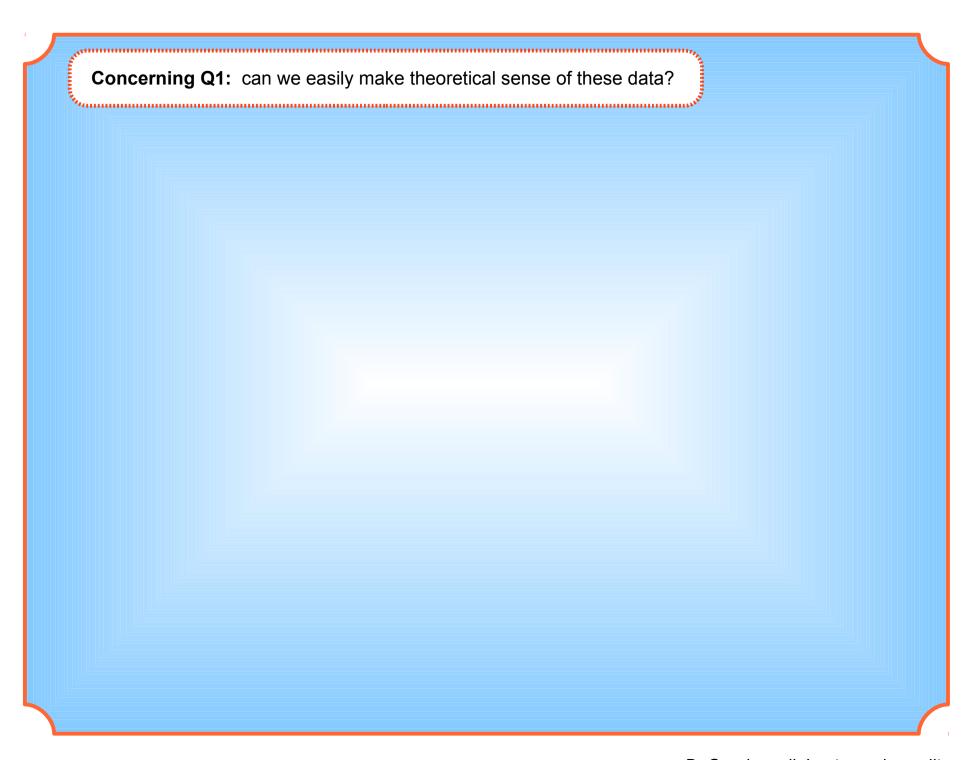
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So, BSM LFNU \Rightarrow BSM LFV (i.e. not suppressed by $m_{_{V}}$)



Yes we can. Consider the following Hamiltonian

$$H_{\text{SM+NP}}(\bar{b} \rightarrow \bar{s} \mu \mu) = -\frac{4 G_F}{\sqrt{2}} V_{tb}^* V_{ts} \frac{\alpha_{\text{em}}}{4 \pi} \left[\bar{b}_L \gamma^{\lambda} s_L \cdot \left(C_9^{(\mu)} \bar{\mu} \gamma_{\lambda} \mu + C_{10}^{(\mu)} \bar{\mu} \gamma_{\lambda} \gamma_5 \mu \right) \right]$$

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About equal size & opposite sign in the SM (at the m_b scale)

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- Advocating the same $(V A) \times (V A)$ structure also for the corrections to $C_{9,10}^{SM}$ (in the $\mu\mu$ -channel only!) would account for:
 - R_{κ} lower than 1
 - **B** \rightarrow K $\mu\mu$ & B_s \rightarrow $\mu\mu$ BR data below predictions
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- A fully quantitative test requires a global fit.

new physics contributions to the Wilson coefficients. We find that the by far largest decrease in the χ^2 can be obtained either by a negative new physics contribution to C_9 (with $C_9^{\rm NP} \sim -30\% \times C_9^{\rm SM}$), or by new physics in the $SU(2)_L$ invariant direction $C_9^{\rm NP} = -C_{10}^{\rm NP}$, (with $C_9^{\rm NP} \sim -12\% \times C_9^{\rm SM}$). A positive NP contribution to C_{10} alone would also improve the fit, although to a lesser extent. [Altmannshofer, Straub, EPJC '15]

For analogous conclusions, see also [Ghosh, Nardecchia, Renner, JHEP '14]

Glashow, DG, Lane, PRL 2015

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This rotation induces <u>LFNU and LFV</u> effects

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Explaining $b \to s$ data

Recalling our full Hamiltonian

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the shift to the C_α Wilson coeff. in the μμ-channel becomes

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The NP contrib. in the eechannel is negligible, as

$$\left|\left(\boldsymbol{U}_{L}^{t}\right)_{31}\right|^{2} \ll \left|\left(\boldsymbol{U}_{L}^{t}\right)_{32}\right|^{2}$$

So, in the above setup

$$R_K \approx \frac{|C_9^{(\mu)}|^2 + |C_{10}^{(\mu)}|^2}{|C_9^{(e)}|^2 + |C_{10}^{(e)}|^2} \simeq \frac{2|C_{10}^{SM} + \delta C_{10}|^2}{2|C_{10}^{SM}|^2}$$

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$$0.77 \pm 0.20 = \frac{BR(B_s \to \mu \mu)_{\text{exp}}}{BR(B_s \to \mu \mu)_{\text{SM}}} = \frac{BR(B_s \to \mu \mu)_{\text{SM+NP}}}{BR(B_s \to \mu \mu)_{\text{SM}}} = \frac{|C_{10}^{\text{SM}} + \delta C_{10}|^2}{|C_{10}^{\text{SM}}|^2}$$

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implying (within our model) the correlations

$$\frac{BR(B_s \to \mu \mu)_{\text{exp}}}{BR(B_s \to \mu \mu)_{\text{SM}}} \simeq R_K \simeq \frac{BR(B^+ \to K^+ \mu \mu)_{\text{exp}}}{BR(B^+ \to K^+ \mu \mu)_{\text{SM}}}$$

Another good reason to pursue accuracy in $B_s \rightarrow \mu\mu$ measurements

See also Hiller, Schmaltz, PRD 14

As mentioned: if R_{κ} is signaling BSM LFNU, then expect BSM LFV as well

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- $oxed{BR}(B^+
 ightarrow K^+ \mu \, au)$ would be even more promising, as it scales with $|(U_L^t)_{33} / (U_L^t)_{32}|^2$
- ✓ An analogous argument holds for purely leptonic modes

6.....

• More quantitative LFV predictions require knowledge of the U_L^{ℓ}

Reminder:

$$(U_L^{\ell})^{\dagger} Y_{\ell} U_R^{\ell} = \hat{Y}_{\ell}$$

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Reminder:

One approach:

DG, Lane, PLB 2015

Appelquist-Bai-Piai ansatz:
 the flavor-SU(3) rotations are not all independent. Choosing 3 to be the independent ones
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4......

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Chiral-suppression factor, of $O(m_{\mu}/m_{Bs})^2$ replaced by α_{em}/π suppression

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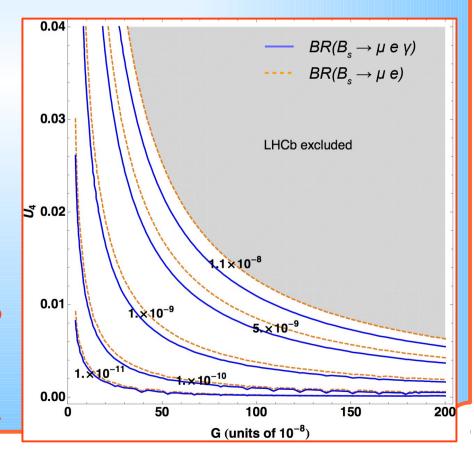
DG, Melikhov, Reboud, 2016

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Enhancement by ~ 30%

Inclusion of the radiative mode more-thandoubles statistics of the non-radiative



D. Guadagnoli, Lepton universality

The interaction advocated in Glashow et al.

$$H_{\rm NP} = G \, \bar{b}'_{L} \gamma^{\lambda} b'_{L} \, \bar{\tau}'_{L} \gamma_{\lambda} \tau'_{L}$$

can also manifest itself in $K \to (\pi) \ \ell'$, for example

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Exp limits

$$\frac{\Gamma(K_L^0 \rightarrow e^{\pm}\mu^{\mp})}{\Gamma(K^+ \rightarrow \mu^+\nu_{\mu})} < 1.7 \times 10^{-12}$$

$$\frac{\Gamma(K^{+} \rightarrow \pi^{+}\mu^{+}e^{-})}{\Gamma(K^{+} \rightarrow \pi^{0}\mu^{+}\nu_{\mu})} < 3.9 \times 10^{-10}$$

BNL E871 Collab., PRL 1998

BNL E865 Collab., PRD 2005

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$$\left|\beta^{(K)}\right|^2 = 2.15 \times 10^{-14}$$

within "model A" of DG, Lane, PLB 2015

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$$BR(K^+ \rightarrow \pi^+ \mu^{\pm} e^{\mp}) \approx 3 \times 10^{-15}$$

with
$$BR(K^+ \rightarrow \pi^0 \mu^+ \nu_{\mu}) \approx 3\%$$

For a recent discussion: Alonso, Grinstein, Martin-Camalich,

• Being defined above the EWSB scale, our assumed operator $G\ \bar{b}'_L \gamma^{\lambda} b'_L \ \bar{\tau}'_L \gamma_{\lambda} \tau'_L$ must actually be made invariant under $SU(3)_c x SU(2)_L x U(1)_{\gamma}$

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Bhattacharya, Datta, London, Shivashankara, PLB 15

$$\bar{b}'_L \gamma^{\lambda} b'_L \bar{\tau}'_L \gamma_{\lambda} \tau'_L$$



$$\begin{array}{c} \text{SU(2)}_{\text{\tiny L}} \\ \\ \text{inv.} \end{array} \qquad \left\{ \begin{array}{c} \bullet \quad \bar{Q}\,{'}_L\, \gamma^\lambda Q\,{'}_L \; \bar{L}\,{'}_L \gamma_\lambda L\,{'}_L \\ \\ \bullet \quad \bar{Q}\,{'}_L^i \, \gamma^\lambda Q\,{'}_L^j \; \bar{L}\,{'}_L^j \gamma_\lambda L\,{'}_L^i \end{array} \right.$$

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[also charged-current int's]

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•
$$\bar{Q}'_L \gamma^{\lambda} Q'_L \bar{L}'_L \gamma_{\lambda} L'_L$$

$$ar{Q}^{\prime i}_{L} \gamma^{\lambda} Q^{\prime j}_{L} ar{L}^{\prime j}_{L} \gamma_{\lambda} L^{\prime j}_{L}$$

Thus, the generated structures are all of:

$$t't'v'_{\tau}v'_{\tau}$$
, $t't'\tau'\tau'$, $b'b'v'_{\tau}v'_{\tau}$, $b'b'\tau'\tau'$, $t'b'\tau'v'_{\tau}$

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 $\begin{array}{c} \text{SU(2)}_{\text{L}} \\ & \bar{Q}^{\,\prime}_{\,L} \, \gamma^{\lambda} Q^{\,\prime}_{\,L} \, \bar{L}^{\,\prime}_{\,L} \gamma_{\lambda} L^{\,\prime}_{\,L} \\ & \bar{Q}^{\,\prime}_{\,L} \, \gamma^{\lambda} Q^{\,\prime}_{\,L} \, \bar{L}^{\,\prime}_{\,L} \gamma_{\lambda} L^{\,\prime}_{\,L} \end{array} \quad \text{[neutral-current int's only]}$

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, $t't'\tau'\tau'$, $b'b'v'_{\tau}v'_{\tau}$, $b'b'\tau'\tau'$, $t'b'\tau'v'_{\tau}$

After rotation to the mass basis (unprimed), the last structure contributes to $\Gamma(b \rightarrow c \tau \bar{\nu}_i)$



Can explain BaBar + Belle + LHCb deviations on $R(D^{(*)}) = \frac{BR(\bar{B} \rightarrow D^{(*)^+} \tau^- \bar{\nu}_{\tau})}{BR(\bar{B} \rightarrow D^{(*)^+} \ell^- \bar{\nu}_{\iota})}$

- Properly taking into account RGE running from the NP scale to the scale(s) of the low-energy processes, one finds non-trivial constraints from:
 - $-B \rightarrow K vv$

See also:

Calibbi, Crivellin, Ota, PRL 2015

Feruglio, Paradisi, Pattori, 2016

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Feruglio, Paradisi, Pattori, 2016

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The latter are the most dangerous.

They "strongly disfavour an explanation of the R(D(*)) anomaly model-independently"

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Feruglio, Paradisi, Pattori, 2016

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Also LFV decays of leptons are generated, and they provide sensitive probes.
 E.g.:

$$BR(\tau \rightarrow 3\mu) \& BR(\tau \rightarrow \mu\rho) \sim 5 \times 10^{-8}$$

Some models explaining R_{κ} and $R(D^*)$

• Introduce one single leptoquark scalar, transforming as (3, 1, -1/3) under $SU(3)_c \times SU(2)_L \times U(1)_Y$

Bauer-Neubert, PRL 2016

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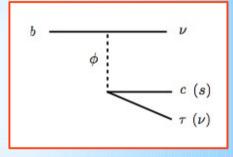
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Bauer-Neubert, PRL 2016

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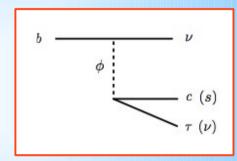
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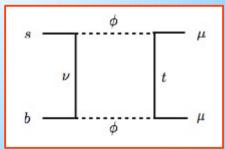
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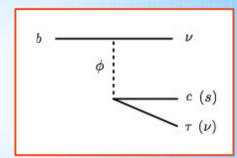
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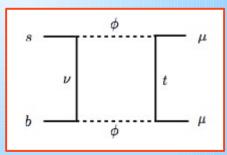
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With M_φ ~ 1 TeV and O(1) generation-diagonal couplings, contributions are just the right size

• New non-Abelian strongly interacting sector with N_{TC} new "techni-fermions" (TC fermions).

Buttazzo, Greljo, Isidori, Marzocca 1604.03940

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The basic idea can easily be understood in analogy to QCD:

The TC-fermion condensate breaks spontaneously a large global symmetry G to a smaller group H, at a scale of about 1 TeV

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The TC-fermion condensate breaks spontaneously a large global symmetry G to a smaller group H, at a scale of about 1 TeV

The broken G/H symmetry gives rise to (pseudo) Goldstone bosons. "Pseudo" because G/H is also broken explicitly by the TC-fermion masses

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The basic idea can easily be understood in analogy to QCD:

- The TC-fermion condensate breaks spontaneously a large global symmetry G to a smaller group H, at a scale of about 1 TeV
- The broken G/H symmetry gives rise to (pseudo) Goldstone bosons. "Pseudo" because G/H is also broken explicitly by the TC-fermion masses



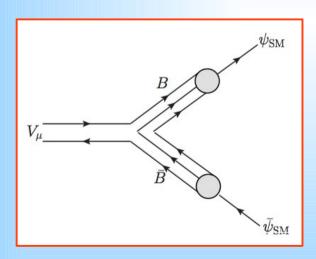
One of the pNGB is the 750-GeV state seen by Atlas & CMS It couples to 2 gluons and decays to 2γ via the anomaly

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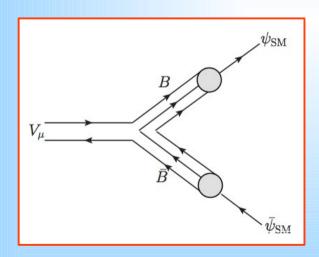
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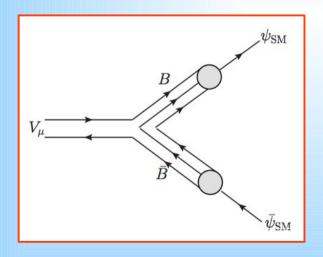


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- To explain the flavor deviations, the mixing needs be hierarchical across generations (largest for the 3rd one, as in partial compositeness)
- Integrating out the vector mesons then yields automatically (among the others) the effective operator

$$H_{\rm NP} = G \, \bar{b}'_{L} \gamma^{\lambda} b'_{L} \, \bar{\tau}'_{L} \gamma_{\lambda} \tau'_{L}$$

proposed in [Glashow, DG, Lane, PRL 15]

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- Timely to propose further tests. One promising direction is that of LFV.
 Plenty of channels, many of which largely untested.