

# Revisiting the problem of inclusive breakup within the IAV model

Antonio M. Moro

University of Seville (Spain)



Workshop on “Deuteron-induced reactions and beyond: inclusive breakup fragment cross sections, MSU, July 2016”

# Outline

- 1 Motivation
- 2 Formulation
- 3 Numerical implementation
- 4 Applications
- 5 Possible extensions for incomplete fusion
- 6 Perspectives

### In 2012:

- Understanding (and quantification) of large inclusive breakup cross sections observed in reactions with neutron-halo nuclei ( ${}^6\text{He}$ ,  ${}^{11}\text{Li}$ ,  ${}^{11}\text{Be}$ )

### In 2016:

- Understanding of larger inclusive yields for other (non-halo) weakly-bound nuclei ( ${}^6,7\text{Li}$ ,  ${}^7\text{Be}$ , ...).
- Possible extension to incomplete fusion.
- Surrogate reactions.

Ph.D thesis of Jin Lei:

*"Study of inclusive breakup reactions induced by weakly bound nuclei"* (Univ. of Seville, July 2016)

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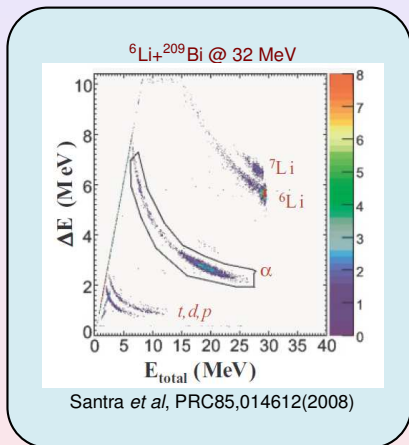
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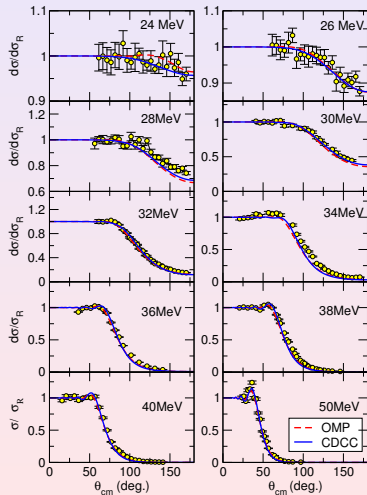
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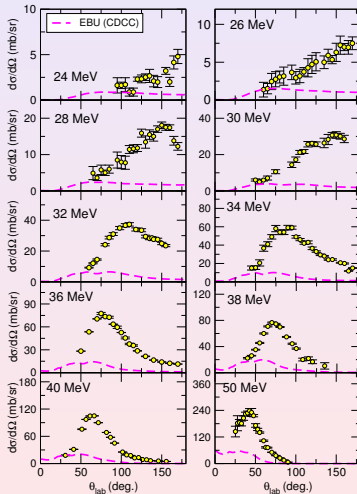


$\alpha$  yields much larger than  $d$  yields  
 $\Rightarrow$  expect other breakup modes  
rather than EBU

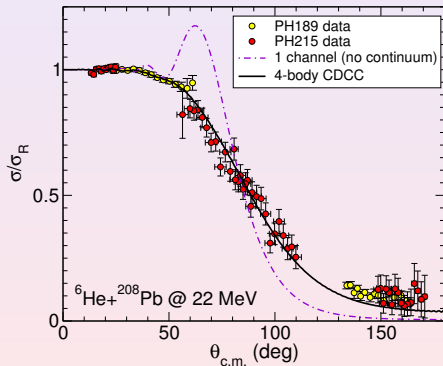
### Elastic scattering



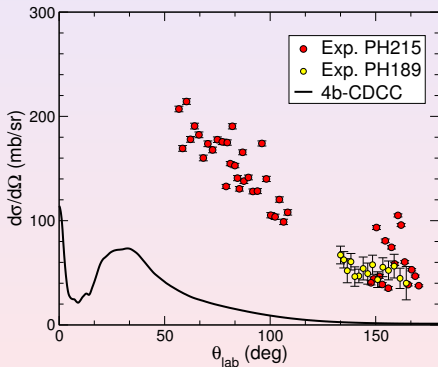
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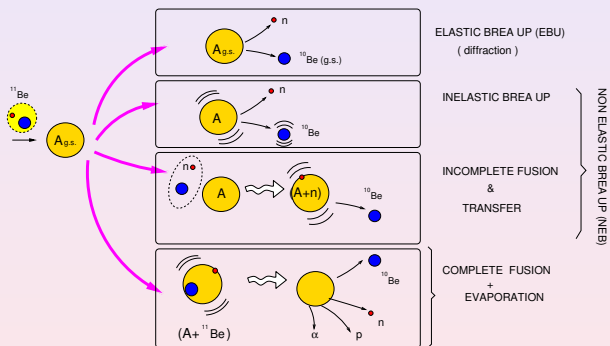


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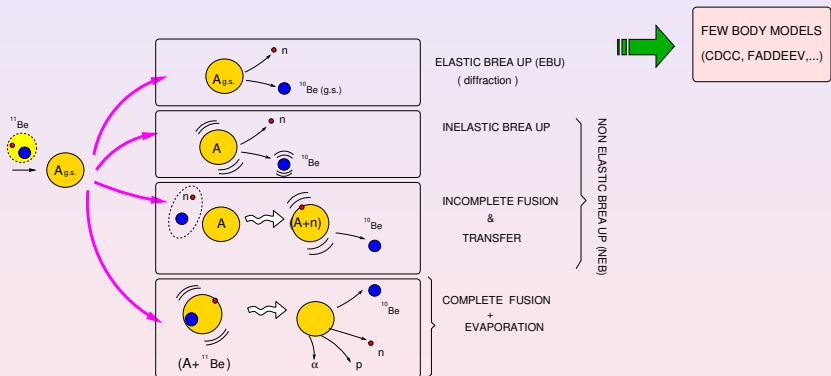




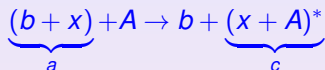
# Breakup modes: $^{11}\text{Be} + A \rightarrow ^{10}\text{Be} + n + A$ example



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- Inclusive breakup:

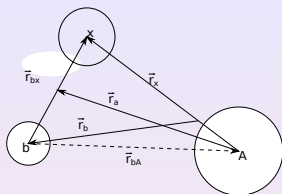


- Inclusive differential cross section:  $\sigma_b^{\text{BU}} = \sigma_b^{\text{EBU}} + \sigma_b^{\text{NEB}}$
- Post-form expression for inclusive breakup:

$$\frac{d^2\sigma}{d\Omega_b E_b} = \frac{2\pi}{\hbar v_a} \rho(E_b) \sum_c |\langle \chi_b^{(-)} \Psi_{xA}^{c,(-)} | V_{bx} | \Psi^{(+)} \rangle|^2 \delta(E - E_b - E^c)$$

- ⇒  $\Psi_{xA}^{c,(-)}$  wavefunctions for  $c \equiv x + A$  states
- ⇒  $\Psi^{(+)}$  exact scattering wavefunction

*Inclusion of all relevant  $c = x + A$  channels is not feasible in general ⇒ use closed-form models*



- Treat  $b$  particle as **spectator**  $\Rightarrow \chi_b^{(-)}(\vec{k}_b, \vec{r}_b)$
- Model Hamiltonian (still many-body for  $x + A$ ):

$$H = K_b + V_{bx} + U_{bA}(\mathbf{r}_{bA}) + \underbrace{K_x + H_A(\xi) + V_{xA}(\xi, \mathbf{r}_x)}_{H_B}$$

- $x - A$  wave function for a given final  $b$  state:

$$Z_x^{(b)}(\xi, \mathbf{r}_x) \equiv (\chi_b^{(-)}(\vec{k}_b) | \Psi) = [E^+ - E_b - H_B]^{-1} (\chi_b^{(-)} | V_{\text{post}} | \Psi) \quad V_{\text{post}} \approx V_{bx}(\mathbf{r}_x)$$

- Define projector onto target g.s. ( $A_{g.s.}$ ):  $\mathcal{P} = |\phi_A^{(0)}\rangle\langle\phi_A^{(0)}|$
- 3-body reduction:  $\mathcal{P}Z_x^{(b)}(\xi, \mathbf{r}_x) = \varphi_x^{(0)}(\mathbf{r}_x)\phi_A^{(0)}(\xi)$

$$[K_x + U_{xA} - E_x]\varphi_x^{(0)}(\mathbf{r}_x) = (\chi_b^{(-)} | V_{bx} | \Psi_{3b}^{(+)})$$

- ⇒  $\Psi_{3b}^{(+)}$  = three-body scattering wave function (DWBA, CDCC, Faddeev...)
  - ⇒  $U_{xA}$  = optical model potential operator.
  - ⇒ The asymptotics of  $\varphi_x^{(0)}(\mathbf{r}_x)$  gives ONLY the EBU part.
- $\sigma_b^{\text{NEB}}$  is the flux loss ("absorption") in the  $x + A_{gs}$  channel:

$$\frac{d\sigma^{\text{NEB}}}{d\Omega_b dE_b} = -\frac{2}{\hbar v_a} \rho_b(E_b) \langle \varphi_x | W_{xA} | \varphi_x \rangle \quad (\text{optical theorem})$$

- Ignore internal spins  $\Rightarrow \vec{l}_a + \vec{l}_{bx} = \vec{l}_b + \vec{l}_x$
- DWBA:  $\Psi_{3b}^{(+)} \approx \chi_d^{(+)}(\mathbf{R})\phi_d(\mathbf{r}_{bx})$
- $\chi_b^{(-)}(\vec{k}_b, \vec{r}_b)$  distorted waves averaged in energy *bins*.

Details in:

J. Lei and A.M.M., PRC 92, 044616 (2015); PRC 92, 061602 (2015)

Two regularization procedures tested and compared:

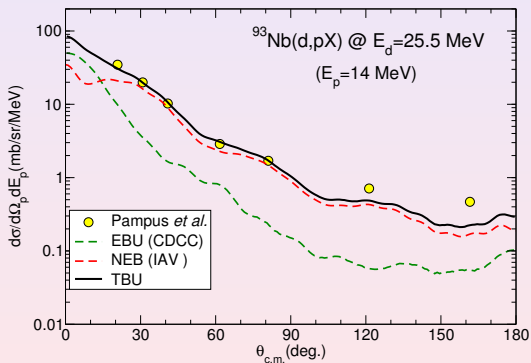
- 1 Damping factor method (Huby and Mines):

$$\rho(\vec{k}_b, \vec{r}_x) \rightarrow \lim_{\alpha \rightarrow 0} e^{-\alpha r_x} \rho(\vec{k}_b, \vec{r}_x)$$

- 2 Binning method (I.J. Thompson): For  $b - B$  distorted waves:

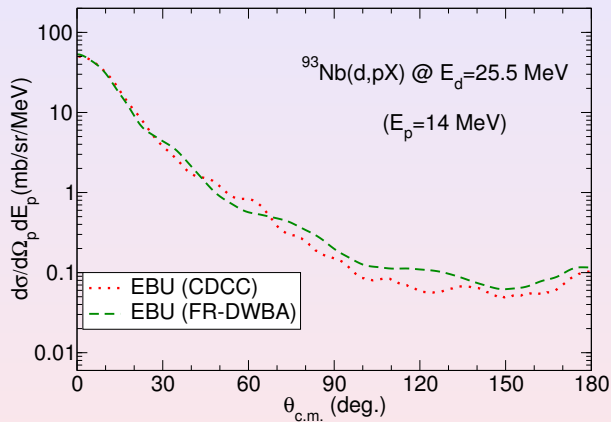
$$\bar{R}_{\ell_b}(r_b, k_b^i) = N \int_{k_b^i - \Delta k_b/2}^{k_b^i + \Delta k_b/2} dk_b R_{\ell_b}(r_b, k_b),$$

- EBU  $\rightarrow$  CDCC (Fresco).
- NBU  $\rightarrow$  post-form IAV model.



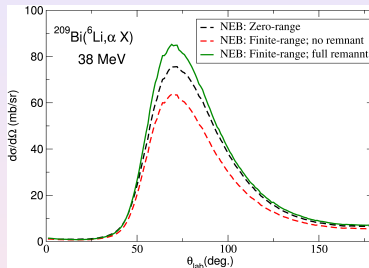
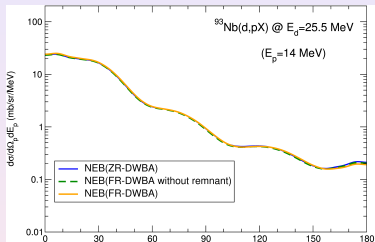
Data: Pampus *et al.*, NPA311 (1978)141





⇒ Good agreement in this case, but needs to be tested for other cases

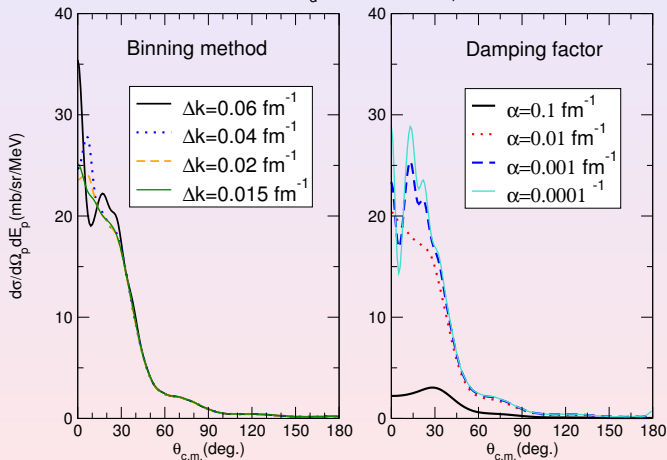
# Importance of finite-range and remnant term



- ⇒ Zero-range accurate for deuterons, but not for other projectiles like  $^6\text{Li}$
- ⇒ Remnant term also important for  $^6\text{Li}$

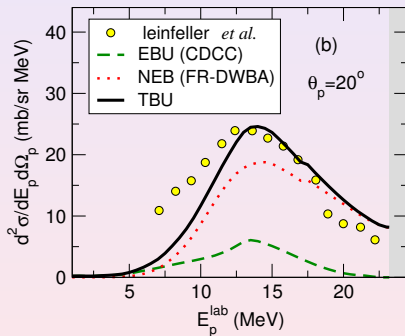
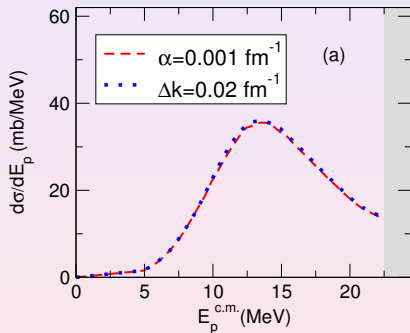
# Comparison of regularization procedures

$^{93}\text{Nb}(d,pX)$  @  $E_d=25.5$  MeV ( $E_p=14$  MeV)

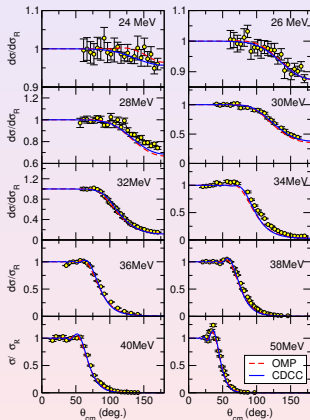
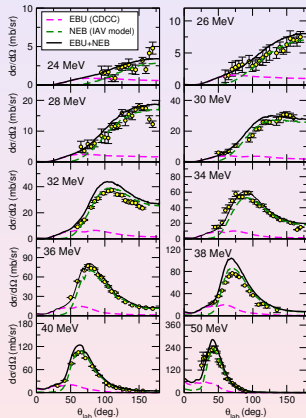


# Comparison of regularization procedures

$^{62}\text{Ni}(d,pX)$  @  $E_d=25.5$  MeV

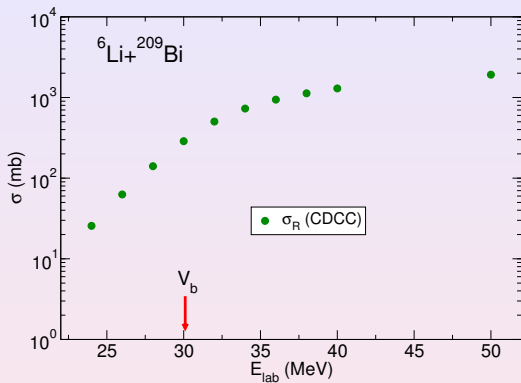


## Elastic scattering

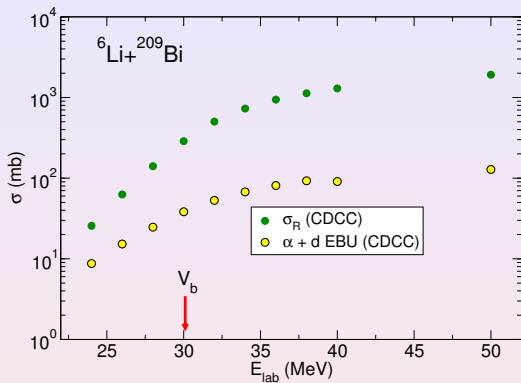
Inclusive  $\alpha$ 's

Good agreement but...  $W_{dA}$  has to be readjusted to give good elastic scattering.

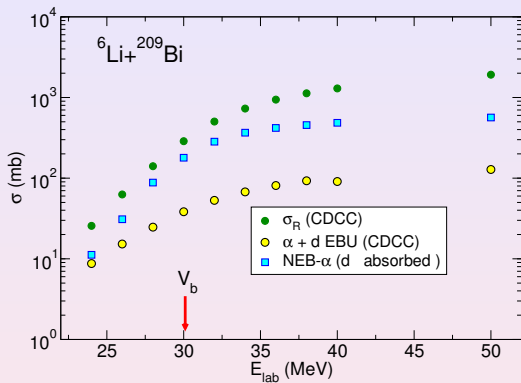
# ${}^6\text{Li}+{}^{209}\text{Bi}$ : incident energy dependence of cross sections



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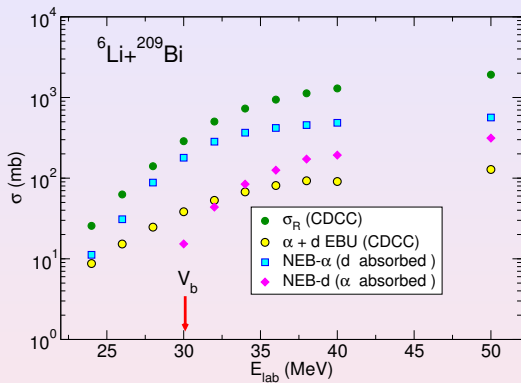


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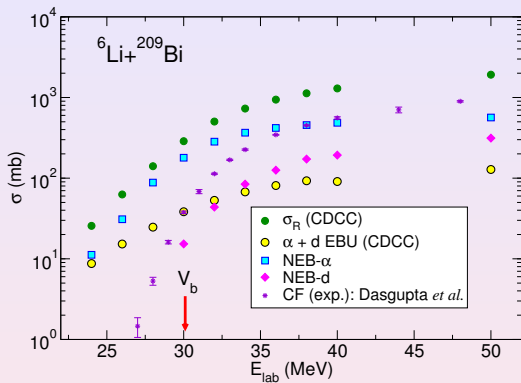


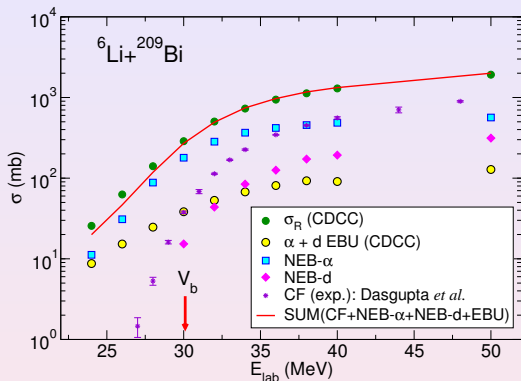


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$$\sigma_{\text{reac}} \approx \sigma_{\alpha+d}(\text{EBU}) + \sigma_{\alpha}(\text{NBU}) + \sigma_d(\text{NBU}) + \sigma(\text{CF})$$

$$\left. \frac{d^2\sigma}{dE_b d\Omega_b} \right|_{\text{NEB}} = -\frac{2}{\hbar v_a} \rho_b(E_b) \langle \psi_x | W_x | \psi_x \rangle,$$

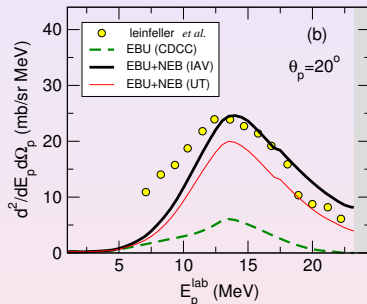
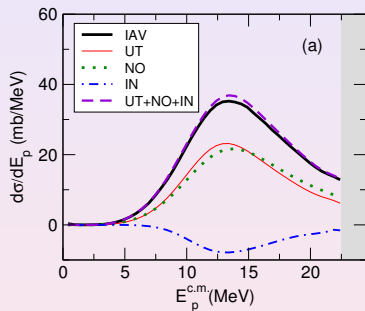
$$(E_x^+ - K_x - U_x) \psi_x^{\text{post}}(\mathbf{r}_x) = (\chi_b^{(-)} | V_{\text{post}} | \chi_a^{(+)} \phi_a); \quad V_{\text{post}} \equiv V_{bx} + U_{bA} - U_{bB}$$

$$(E_x^+ - K_x - U_x) \psi_x^{\text{prior}}(\mathbf{r}_x) = (\chi_b^{(-)} | V_{\text{prior}} | \chi_a^{(+)} \phi_a); \quad V_{\text{prior}} \equiv U_{xA} + U_{bA} - U_{aA}$$

$$\psi_x^{\text{post}} = \psi_x^{\text{prior}} + \psi_x^{\text{NO}} \quad \psi_x^{\text{NO}}(\vec{r}_x) \equiv \langle \chi_b^{(-)} | \chi_a^{(+)} \phi_a \rangle$$

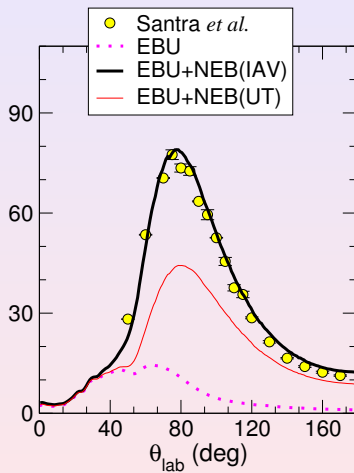
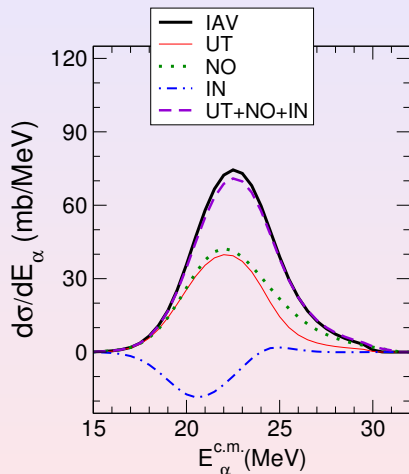
$$\left. \frac{d^2\sigma}{dE_b d\Omega_b} \right|_{\text{NEB}}^{\text{IAV}} = \left. \frac{d^2\sigma}{dE_b d\Omega_b} \right|_{\text{NEB}}^{\text{UT}} + \left. \frac{d^2\sigma}{dE_b d\Omega_b} \right|_{\text{NEB}}^{\text{NO}} + \left. \frac{d^2\sigma}{dE_b d\Omega_b} \right|_{\text{NEB}}^{\text{IN}}$$

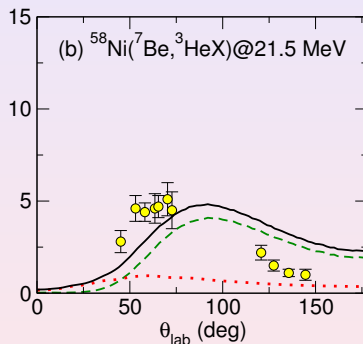
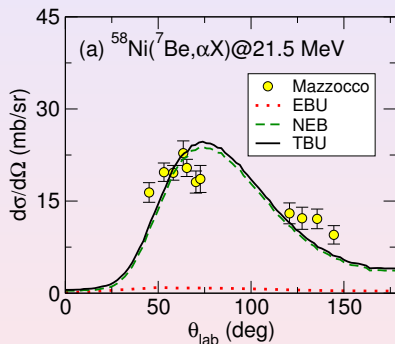
$^{62}\text{Ni}(d,pX) @ E_d=25.5 \text{ MeV}$



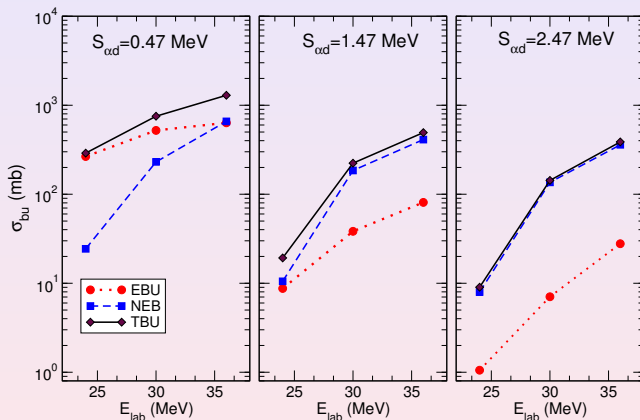
J. Lei, A.M.M., PRC92, 061602(R)(2015)

$^{209}\text{Bi}(^6\text{Li},\alpha X)$  @  $E=36$  MeV



Data: Mazzocco *et al.*:  $\sigma_\alpha \approx 5\sigma_{^3\text{He}}$ 

# Effect of binding energy and incident energy: $^{209}\text{Bi}(^6\text{Li}^+, \alpha)\text{X}$

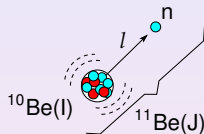




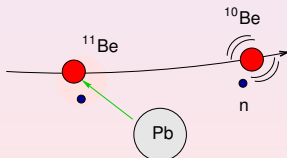
Core excitations will affect:

- 1 the **structure** of the projectile  $\Rightarrow$  core-excited admixtures

$$\Psi_{JM}(\vec{r}, \xi) = \sum_{\ell, j, l} [\varphi_{\ell, j, l}^J(\vec{r}) \otimes \Phi_l(\xi)]_{JM}$$

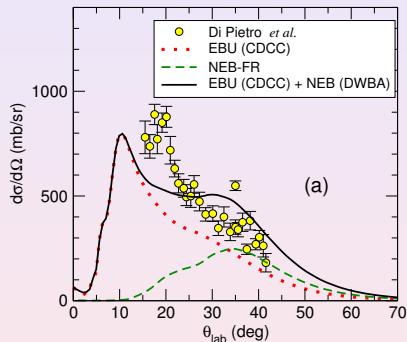


- 2 the **dynamics**  $\Rightarrow$  collective excitations of the  $^{10}\text{Be}$  during the collision compete with halo (single-particle) excitations.

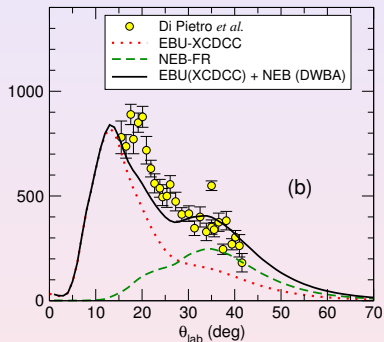


$\Rightarrow$  Both effects have been recently implemented in an extended version of of the CDCC formalism (CDCC): [Summers \*et al\*, PRC74 \(2006\) 014606](#), [R. de Diego \*et al\*, PRC 89, 064609 \(2014\)](#)

### Standard CDCC + NEB



### Extended CDCC + NEB



⇒ Core deformation/excitations very important for EBU. What about for NEB?

The IAV model provides the total NEB cross section, but for some applications  $\sigma^{\text{ICF}}$  is needed:

$$\sigma^{\text{NEB}} = \sigma^{\text{ICF}} + \sigma^{\text{DR}}$$

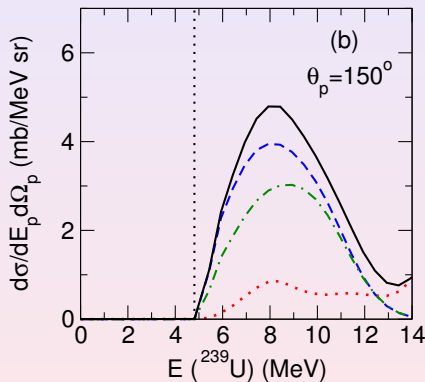
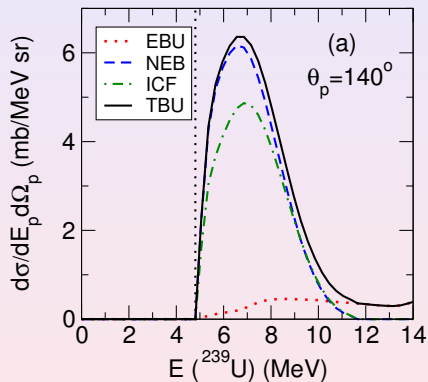
Some ideas to evaluate  $\sigma^{\text{ICF}}$ :

- 1  $W_{xA} = W_{xA}^{\text{DR}} + W_{xA}^{\text{CN}}$ , with  $W_{xA}^{\text{CN}}$  constrained by CC calculations or  $x + A$  fusion.

$$\frac{d\sigma^{\text{ICF}}}{d\Omega_b dE_b} = -\frac{2}{\hbar v_a} \rho_b(E_b) \langle \varphi_x | W_{xA}^{\text{CN}} | \varphi_x \rangle$$

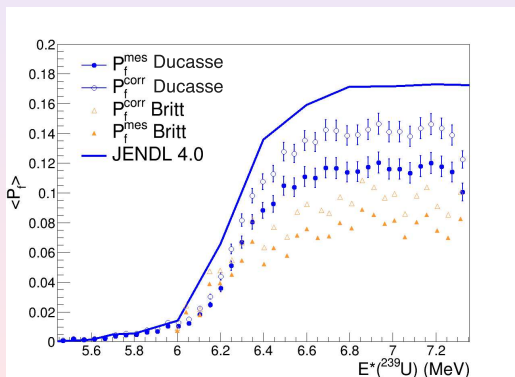
- 2 Extended CDCC with target excitation (should provide EBU + INBU consistently).
- 3 Extended IAV model.

Application to ( $d, p$ ) surrogate reactions: the  $^{238}\text{U}(d, pf)$  case at  $E_d=15$  MeV



$$P_f^{\text{corr}}(E_n) = P_f^{\text{meas}}(E_n) \frac{\sigma_{\text{TBU}}(E_n)}{\sigma_{\text{ICF}}(E_n)},$$

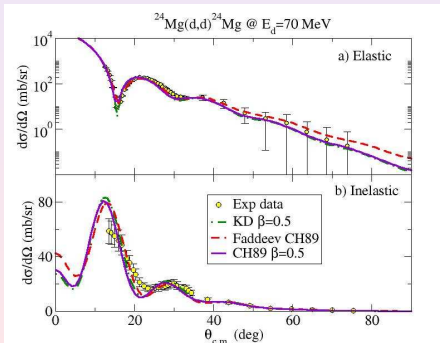
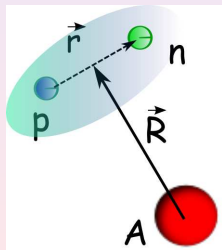
Q. Ducasse et al,  
arXiv:1512.06334



Effective model 3-body Hamiltonian with target excitation:

$$H = H_{\text{proj}}(\mathbf{r}) + H_{\text{tar}}(\xi_t) + \hat{T}_{\mathbf{R}} + U_{CA}(\mathbf{r}_{CA}, \xi_t) + U_{VA}(\mathbf{r}_{VA}, \xi_t),$$

Benchmark with Faddeev formalism:



M. Gomez-Ramos and A.M.M., in preparation

- Inclusion of intrinsic **spins** of clusters.
- Explicit inclusion of **target excitation**, either prior to breakup (in  $a + A$  interaction) and/or in the  $x - A$  channel WF.
- Calculations for “stripping” part in **nucleon removal cross sections (knockout)** at intermediate energies may provide a benchmark for standard semiclassical approaches.
- Extension to **3-body** projectiles ( ${}^9\text{Be}$ ,  ${}^6\text{He}$ ,  ${}^6\text{Li}$ , etc).
- Proper application to more weakly bound (e.g. **halo**) nuclei might require going beyond (DWBA) in the calculation of  $\varphi_x$  (eg. CDCC):

$$[K_x + U_{xA} - E_x]\varphi_x(\mathbf{r}_x) = (\chi_b^{(-)} | V_{bx} | \Psi_{3b}^{(+)} \rangle$$