



# Likelihood-based test statistic in Solar Axion searches

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#### **Solar Axions**



= 0

z = -d

z = -2d





#### **Example Experiments**



#### **MAJORANA DEMONSTRATOR**

#### **EDELWEISS-III**

Array of High Purity Germanium (HPGe) detectors Extremely clean, very low background rate edelweiss.in2p3.fr Ge detector/crystal sanfordlab.org



## **Axion Bragg Diffraction**





#### Motion of the Sun w.r.t. the Earth $\rightarrow$ Changing of attacking angles → Complicated signal patterns **Tables of Sun's trajectory**

Calculated for  $\lambda$ =1 and M<sub>D</sub>=1kg HPGe Current limits:  $\lambda$  < around 1 x 10<sup>-3</sup>

DAMA, Phys. Lett. B 515 (2001) 6, 90% CL; EDELWEISS JCAP 11 (2013) 067, 95% CL

→ Numerical expression of signal rate  $R^{(j)} = R^{(j)}(\phi^{(j)})$  depends on angle of the *j*<sup>th</sup> detector

Calculation of Axion signal closely follows: R. J. Creswick, et al., PLB 427 (1998) 235-240 SOLAX, Phys. Rev. Lett. 81, 5068 (1998)

0.9



### **The Statistical Problem**



For the *j*-th detector (*j*=1, ..., n<sub>D</sub>), the likelihood function isJoint likelihood
$$\mathcal{L}_D^{(j)} = e^{-m^{(j)}} \prod_{i=1}^{N^{(j)}} [b^{(j)} + R^{(j)}(t_i^{(j)}, E_i^{(j)})]$$
 $\mathcal{L} = \prod_{j=1}^{n_D} \mathcal{L}_D^{(j)}$ where *i* is the *i*-th event in this detector $R^{(j)} \propto \lambda M_D^{(j)}$ 

Assuming required angles of  $n_D$  detectors ( $\vec{\phi}$ ) are all precisely known

Parameter of interest: Axion-photon coupling constant ( $\lambda$ , universal)

Axion signal

- Nuisance parameter: background ( $\overline{b}$ , universal)
- Apply the profile likelihood method Ref.:Rolke et al, NIMA 551 (2005) 493 -2ln(profile likelihood) should approximately have a  $\chi^2(1)$  distribution
- $\equiv -2 \ln \mathcal{L}_p(\lambda_{true})$ Monte Carlo simulations generated with  $\lambda_{true}$  $= -2\ln \frac{\sup\{\mathcal{L}(\lambda_{true}, \bar{b}); \bar{b}\}}{\sup\{\mathcal{L}(\lambda, \bar{b}); \lambda, \bar{b}\}}$ 1000 simulations per ensemble Define a test statistic *D*, based on profile likelihood.



## **Effect of Physical Boundary**



- Coupling  $\lambda_{true}$ =0 used in Monte Carlo simulations, no Axion signals
- Physical boundary makes the test statistic distribution narrower than  $\chi^2(1.0)$   $\rightarrow$  Over-coverage with nominal critical value 2.7 for 90% confidence level
  - $\rightarrow$  Proper critical value for 90% C.L. is smaller than 2.7



# Los Alamos Detector Angles: Poorly Known



Laboratory

Comprehensive orientation surveys are required to obtain the absolute angle for every Ge detector. Difficult to measure, and could suffer from large uncertainty ~ 15°, i.e. poorly known



Crystal axis angle measurable by Laue diffraction or charge carrier drift time (*in situ*)



Detector Array

**Experiment Apparatus** 



Figures of **MAJORANA DEMONSTRATOR** courtesy of sanfordlab.org.

Topological contours of Sanford Underground Research Facility from Journal of Physics: Conference Series **606** (2015) 012015 Solar Axions

Figure of a HPGe Detector , Courtesy of James Loach, LBNL Los Alamos The Statistical Problem Revised



Parameter of interest: Axion-photon coupling constant (universal) Nuisance parameters: background (universal) and

angles of  $n_D$  detectors (poorly known to within 15°)

Test statistic  $D \equiv -2 \ln \mathcal{L}_p(\lambda_{true})$  $= -2 \ln \frac{\sup\{\mathcal{L}(\lambda_{true}, \bar{b}, \vec{\phi}); \bar{b}, \vec{\phi}\}}{\sup\{\mathcal{L}(\lambda, \bar{b}, \vec{\phi}); \lambda, \bar{b}, \vec{\phi}\}}$ 

- Supposedly the nuisance parameters can be profiled out from the data.
- Works for background
- Works for the detector angles only if the data itself has sensitivity on them.
- Axion signal depends on the detector angles; the background events don't.
- → None to few axion signals = none to little ability to profile out the angles.

#### Only considering four detectors, $n_D=4$ simulated 1000day kg exposure for each detector

# **Los Alamos Poor Angle Info + Big Coupling**



Signal  $\lambda_{true} = 5 \times 10^{-4}$  , bakground  $\overline{b} = 0.1$  cts/kg/day/keV

- Total signal ~ 0.1 cts/day/kg, total background ~ 0.6 cts/day/kg
- Signal large enough for the detector angles to be accurately profiled out from data. Nominal critical value 2.7 for 90% is OK



# • Los Alamos Poor Angle Info + Small Coupling



Signal  $\lambda_{true} = 1 \times 10^{-4}$  , bakground  $\overline{b} = 0.1$  cts/kg/day/keV

- Total signal ~ 0.02 cts/day/kg, total background ~ 0.6 cts/day/kg
- Signals not large enough for accurate profiling of angles, but still some sensitivity to angles. Nominal critical value 2.7 for 90%CL gives under-coverage



## • Los Alamos Poor Angle Info + Zero Coupling



Signal  $\lambda_{true} = 0$ , bakground  $\overline{b} = 0.1$  cts/kg/day/keV

- No signals in MC data, no sensitivity on angles
- Test statistic distribution approximates  $\chi^2(5.0)$
- 1 unknown coupling + 4 unknown angles of 4 detectors  $\rightarrow$  5 degrees of freedom.





## χ<sup>2</sup> Degree of Freedom



- Very poorly constrained nuisance parameters could affect test statistic to deviate from  $\chi^2(1.0)$ , if they cannot be profiled out in the data.
- Nominal critical values for  $\chi^2(1.0)$  cause under-coverage.

If  $\lambda_{true} = 0$  and large uncertainty on  $n_D$  angles, almost  $\chi^2 (1 + n_D)$ How about tiny but non-zero  $\lambda_{true}$ ? Need extensive Monte Carlo simulations  $\rightarrow$  Numerically obtain critical value

Is there a general solution without extensive Monte Carlo simulations?

Avoid the problem: measure all the angles For one detector, *small enough* angular uncertainty

- $\rightarrow$  angle can be treated as known  $\rightarrow$  Test statistic stays effectively  $x^2$  with
- → Test statistic stays effectively  $\chi^2$  with 1 d.o.f.
- $\rightarrow$  For n<sub>D</sub> detectors, all angles can be treated as known, still  $\chi^2(1)$ .

In this example: 15° uncertainty =17% of the allowed range of 90°, already too large *How small is small enough?* 



## **Maximum Angle Uncertainty**





![](_page_13_Picture_0.jpeg)

## **Angles Unknown**

![](_page_13_Picture_2.jpeg)

Current and future experiments have many detectors. Challenging to measure every detector angle to a 3°-4° precision. *Alternatively*,

![](_page_13_Figure_4.jpeg)

- n<sub>D</sub> detectors with equal mass M<sub>D</sub>
- Joint likelihood.

- One detector with mass  $n_D \times M_D$
- Reduce the dimension of the problem
- At the cost of losing individual detector information
- Single likelihood based on R<sub>Total</sub>

## An Averaging Approximation

Los Alamos

![](_page_14_Picture_1.jpeg)

![](_page_14_Figure_2.jpeg)

![](_page_15_Picture_0.jpeg)

![](_page_15_Picture_2.jpeg)

![](_page_15_Figure_3.jpeg)

- Close to 50 detectors is sufficient to use the averaging approximation in the model
- A factor of 2.5-3 increase of experimental uncertainty on the coupling constant, even for numerous detectors, due to loss of information.

![](_page_16_Picture_0.jpeg)

![](_page_16_Picture_1.jpeg)

![](_page_16_Picture_2.jpeg)

#### **Findings:**

Very poorly constrained nuisance parameters cause likelihood-based test statistic to deviate from  $\chi^2(1.0)$ , if they cannot be profiled out in the data.

Physical boundary also affects the test statistic distribution as well.

#### In solar axion search:

Extensive Monte Carlo simulations + likelihood-based test statistic
→ Numerically obtain correct critical values for confidence intervals

OR measure angles to 3°-4° precision  $\rightarrow$  test statistic ~  $\chi^2(1.0)$ 

OR use an averaging method with largely reduced sensitivity, if 50 or more detectors

![](_page_17_Picture_0.jpeg)

![](_page_17_Picture_1.jpeg)

![](_page_17_Picture_2.jpeg)

![](_page_18_Picture_0.jpeg)

### Backup

**Question:** Both the axion-photon coupling and the angles can be treated as

- parameters of interest
- little to none sensitivity from the data
- somewhat constrained otherwise (physical boundary, subsidiary measurements)
   Statistical methods to treat this situation without extensive simulation?

![](_page_19_Picture_0.jpeg)

![](_page_19_Picture_1.jpeg)

![](_page_19_Picture_2.jpeg)

Axion: originally postulated as the pseudo Nambu–Goldstone boson for the breaking of Peccei-Quinn U(1) symmetry, which is involved to explain the <u>strong CP problem</u> **Candidate for low mass dark matter** 

![](_page_19_Figure_4.jpeg)

Upper limits on Axion-photon coupling constant. Figure adapted CDMS Phys. Rev. Lett. 103, 141802 (2009) Examples of search efforts: Semi-conductor and solid scintillator (crystal) based: SOLAX, CDMS, DAMA, EDELWEISS, TEXONO, MAJORANA DEMONSTRATOR

Man-made magnetic field based: Tokyo Helioscope, CAST, ADMX, IAXO

Xenon based: XMASS, XENON100

Space based:

XMM-Newton: Potential solar axion signatures in earth magnetic field MNRAS **445**, 2146–2168 (2014)

<sup>7</sup> This talk only considers the Axion-photon coupling and experiments using crystals

![](_page_20_Picture_0.jpeg)

### **Statistical Method**

![](_page_20_Picture_2.jpeg)

Construct likelihood function

$$L(\lambda, b, \vec{\phi} | \vec{X}) = \prod_{i=1}^{n} f(X_i | \lambda, b, \vec{\phi})$$

 $\vec{\phi}$  are the azimuthal angles, and b is background level. They are the nuisance parameters

Profile Likelihood (
$$\Lambda$$
) = Maximized L against all nuisance parameters.

$$\Lambda(\lambda_0 | \vec{X}) = \frac{\sup\{L(\lambda_0, b, \vec{\phi} | \vec{X}); b, \vec{\phi}\}}{\sup\{L(\lambda, b, \vec{\phi} | \vec{X}); \lambda, b, \vec{\phi}\}}$$

Rolke et al, NIMA 551 (2005) 493

![](_page_20_Figure_9.jpeg)

Rolke et al, NIMA 551 (2005) 493: "  $-2\log\Lambda$  has an approximated  $\chi^2$ distribution with 1 degree of freedom" i.e.  $\chi^2(1)$ 

If  $\chi^2(1)$  is indeed followed, change  $-2\log \Lambda$  by 2.71 (3.84) for 90% (95%) confidence level.

Physical limits and other factors will change the distribution of profile likelihood.
 Need adjustments to ensure coverage.

![](_page_21_Picture_0.jpeg)

### **In-situ Angle Measurement**

![](_page_21_Picture_2.jpeg)

![](_page_21_Figure_3.jpeg)

The charge carrier velocity has an angular dependence.

drift time v.s. angle

![](_page_21_Figure_6.jpeg)

Alternatively, use a Laue measurement to study the Ge detectors:

- \* Higher angular precision
- \* May introduce more surface time, undesired for underground Low background Ge experiment

![](_page_22_Picture_0.jpeg)

### **Experiment design**

![](_page_22_Picture_2.jpeg)

![](_page_22_Figure_3.jpeg)

Utilize the coincidence between the 356keV gamma and the 81keV gamma of <sup>133</sup>Ba

For <sup>133</sup>Ba radioactive decay, the 356keV gamma is always immediately (~ns) followed by a 81keV gamma.

Use the 356keV gamma signal as a the starting time of an event.

Use the 81keV gamma to probe the Ge detector.

Most of this low energy gammas will not penetrate deep into Ge  $\rightarrow$  similar drift distance

 $\rightarrow$  Avoids smearing in the drift time

![](_page_23_Picture_0.jpeg)

#### **Example efforts**

![](_page_23_Picture_2.jpeg)

![](_page_23_Figure_3.jpeg)

Exact treatment of waveform varies, typically including:

Baseline removal, smoothing (e.g. via averaging) and pole-zero correction.

Typically 50% or 90% rise time of the Ge detector signal are regarded as the end of the drift  $\rightarrow$  Angle dependent

50% or 10% rise time of the scintillation signal are regarded as the start of the drift

 $\rightarrow$  Common for all angles on average

![](_page_24_Picture_0.jpeg)

## **Goodness of Fit with Approx.**

![](_page_24_Picture_2.jpeg)

![](_page_24_Figure_3.jpeg)

- Model's single likelihood function constructed with the *approximation*.
- To describe Monte Carlo data. Does it work? Perform  $\chi^2$  Goodness of Fit test, convert  $\chi^2$  to p-value

If axion signals exist, the model (with the approximation) well describes an experiment with 150 detectors, does NOT describe an experiment with 5 detectors model cannot be used for 5-detector expt.

![](_page_24_Figure_7.jpeg)

If no axion signal at all, the model well describes an experiment regardless of the detectors. No axion signals  $\rightarrow$  expression of  $R_{Total}$  is irrelevant

![](_page_25_Picture_0.jpeg)

#### Alternatively

![](_page_25_Picture_2.jpeg)

Current and future experiments many detectors

Tremendous challenge to measure every detector angle to 3°-4° precision

![](_page_25_Figure_5.jpeg)

 $n_G$  groups (e.g.  $n_G$  = 90), with equally  $n_D/n_G$  detectors in each group, if  $n_D \rightarrow \infty$ . Flat angle prior assumed.

Treat n<sub>D</sub> detectors with equal mass M<sub>D</sub> as ndetector with mass  $n_D \times M_D$ 

$$\frac{\sum_{j=1}^{n_D} R^{(j)}(\phi^{(j)})}{n_D} \to \frac{\int_{-\pi/4}^{\pi/4} R(\phi) d\phi}{\pi/2}, \text{ if } n_D \to \infty$$