

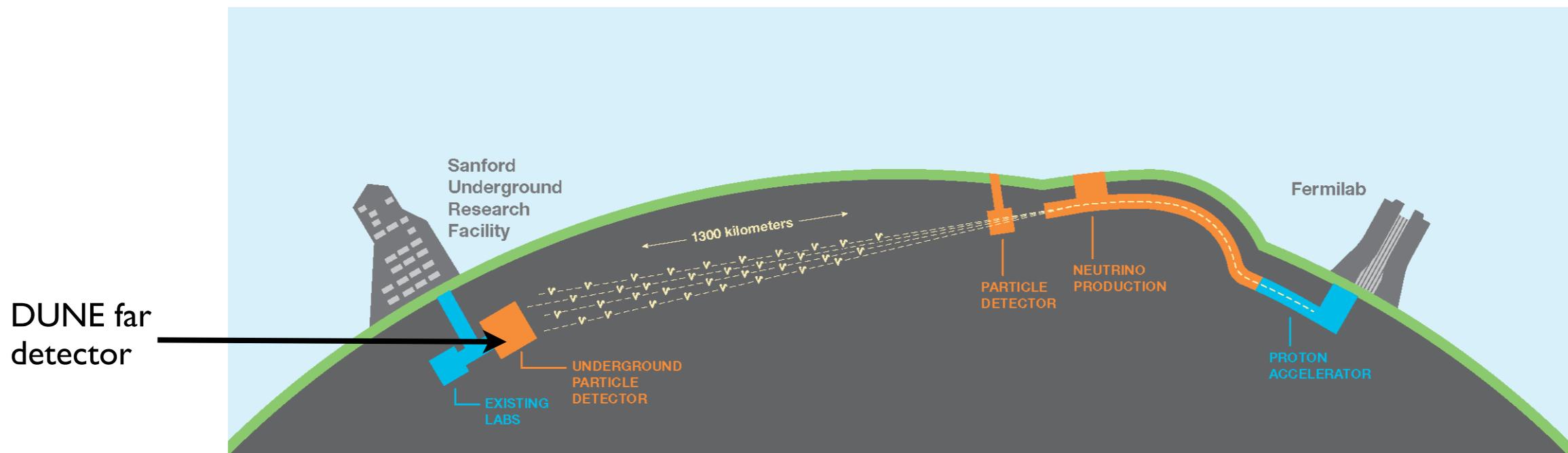


# Neutrino energy reconstruction in the DUNE far detector

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for the DUNE collaboration

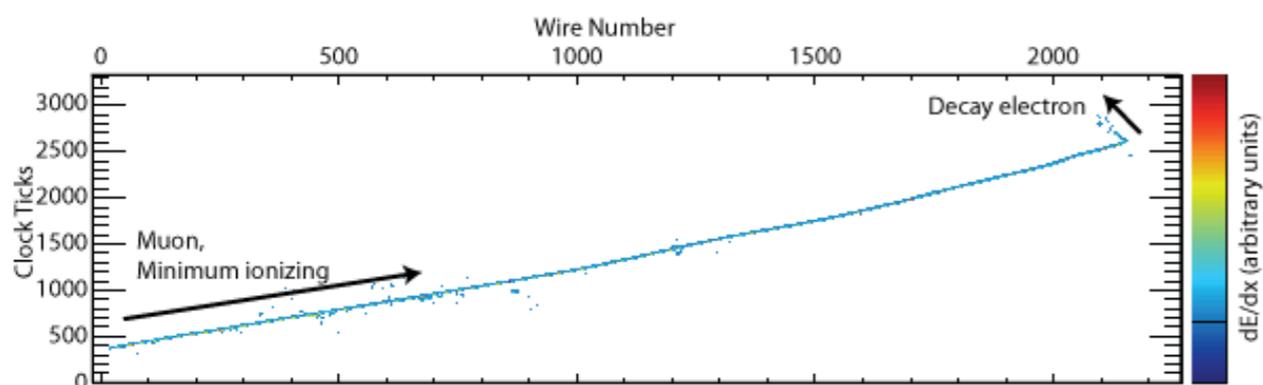
Meeting of APS Division of Particles and Fields

31 July - 4 August 2017

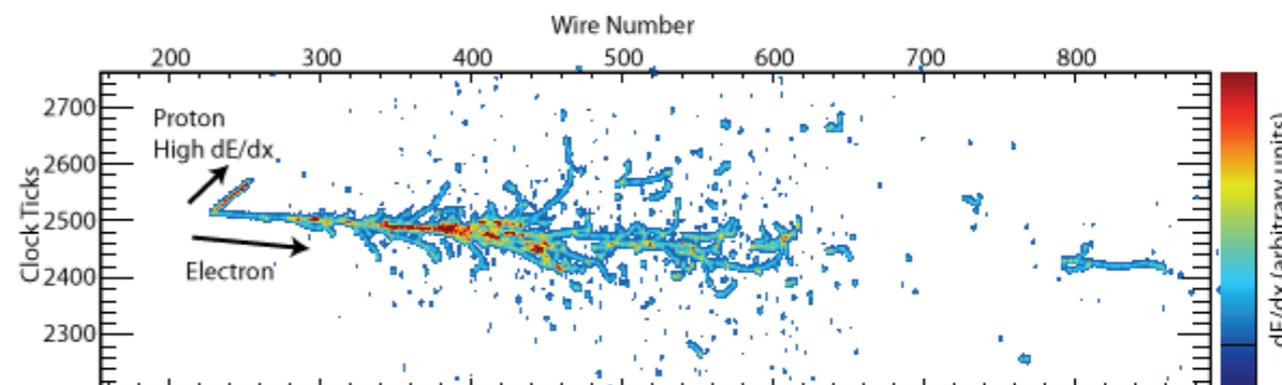


We are working on reconstructing neutrino energy in events selected as  $\nu_\mu$  CC and  $\nu_e$  CC in the DUNE far detector, which is a liquid argon time projection chamber (LArTPC).

$\nu_\mu$  CC event in LArTPC

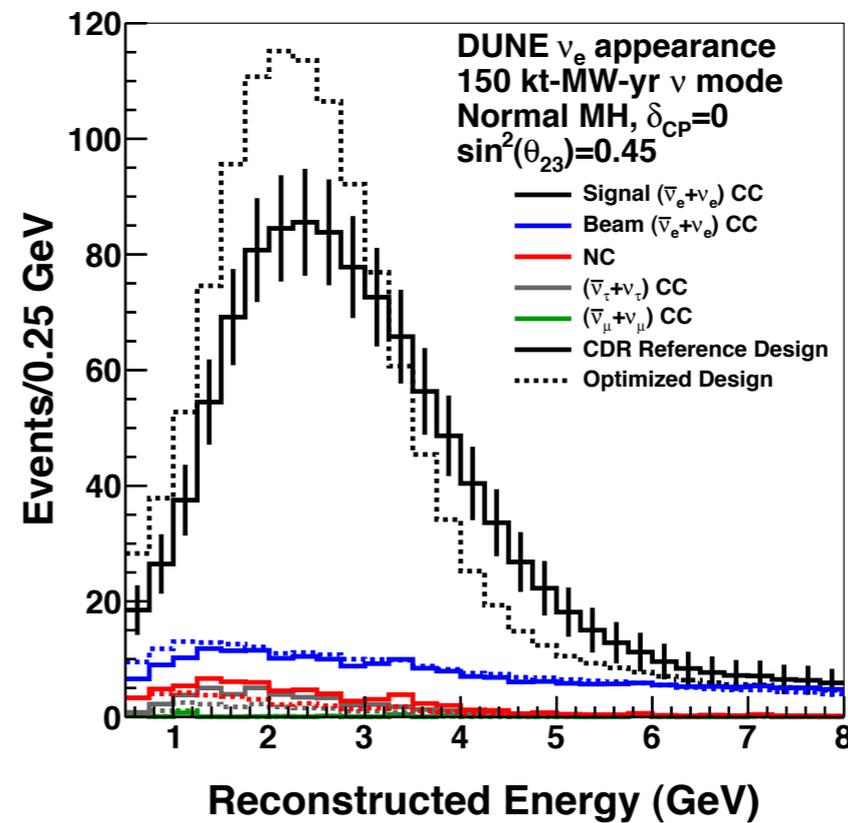
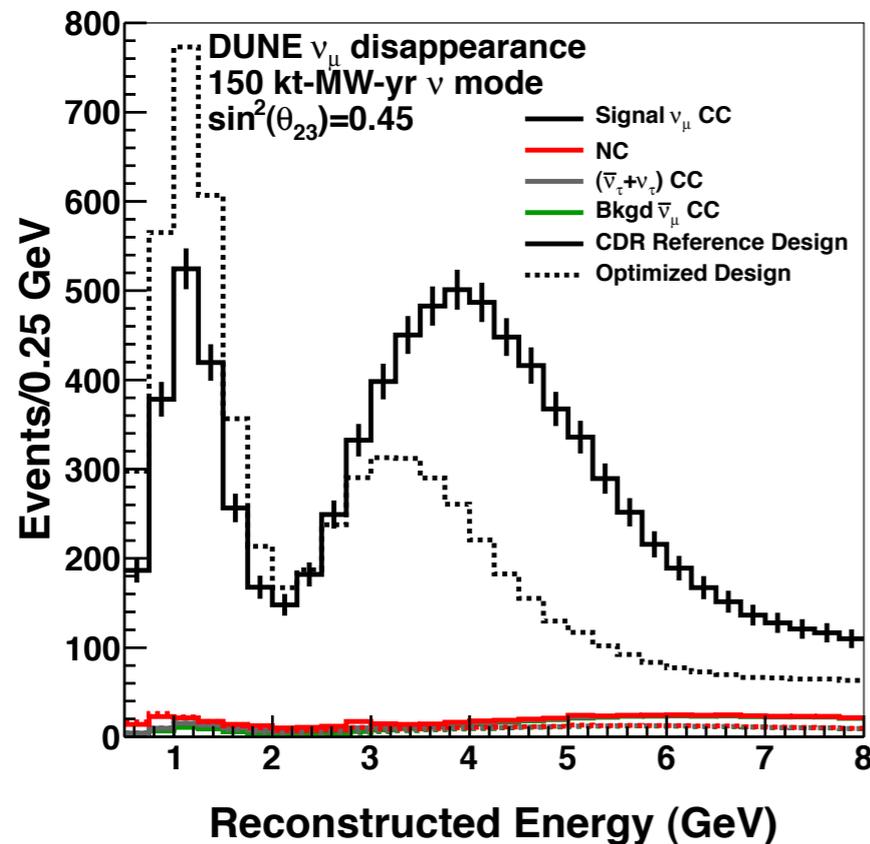


$\nu_e$  CC event in LArTPC



Ref: <http://docs.dunescience.org/cgi-bin/ShowDocument?docid=183&asof=2016-1-30>

Neutrino energy reconstruction is extremely important for DUNE since it will make measurements of neutrino oscillation parameters including the possible CP-violating phase  $\delta_{CP}$ . Neutrino oscillation depends on true energy, but the energies of the neutrinos interacting in the far detector are unknown, and must be reconstructed as accurately as possible.





# Introduction



We use realistic simulations of the far detector and real reconstruction.

For events selected as  $\nu_{\mu}$  CC, divide event into longest reconstructed track and hadronic energy.

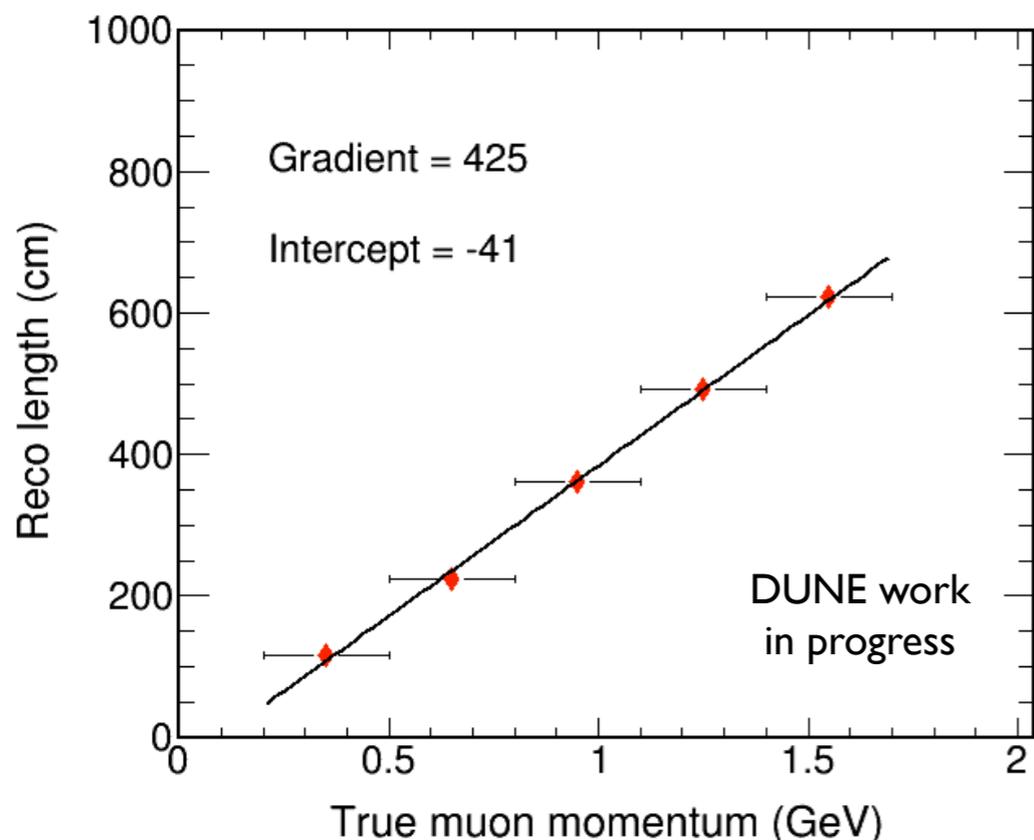
For events selected as  $\nu_e$  CC, divide event into reconstructed shower with highest charge and hadronic energy.

Estimate the hadronic energy using calorimetric information from hits not in track or shower, converting ADC counts to energy after making corrections for electron lifetime and recombination.

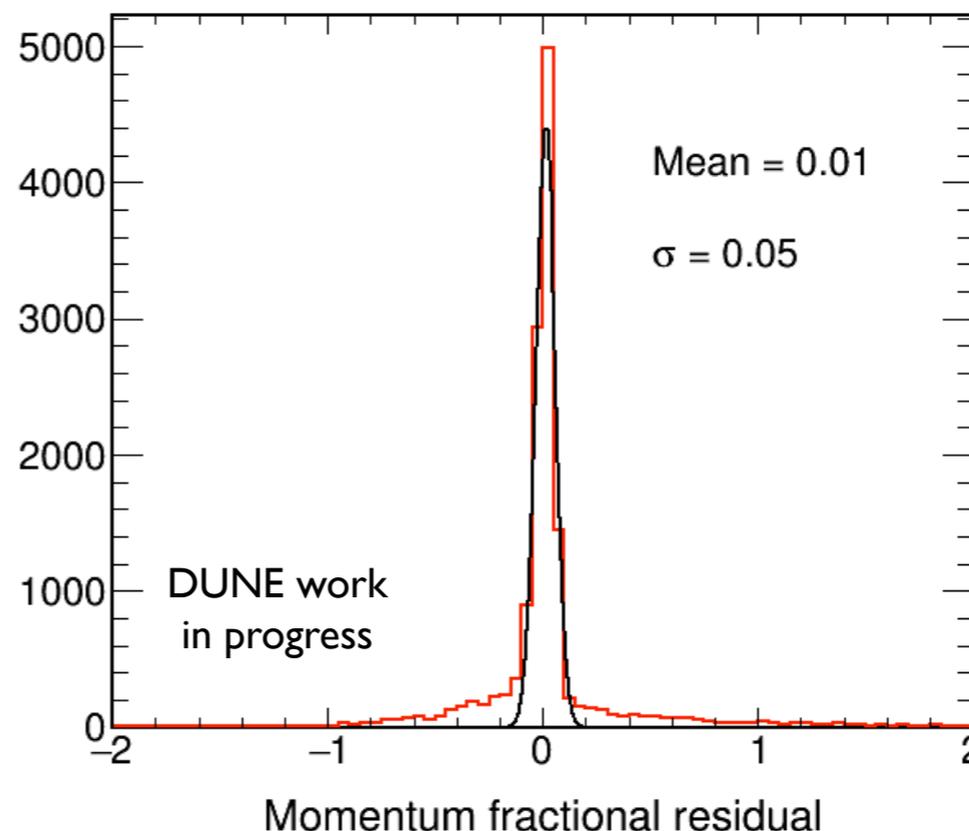
If the longest track is contained within the detector, estimate its momentum from its range. Calibrate track momentum by range using Monte Carlo.

Estimate reco track momentum as  $(\text{range} - \text{intercept}) / \text{gradient}$ .

Range momentum calibration



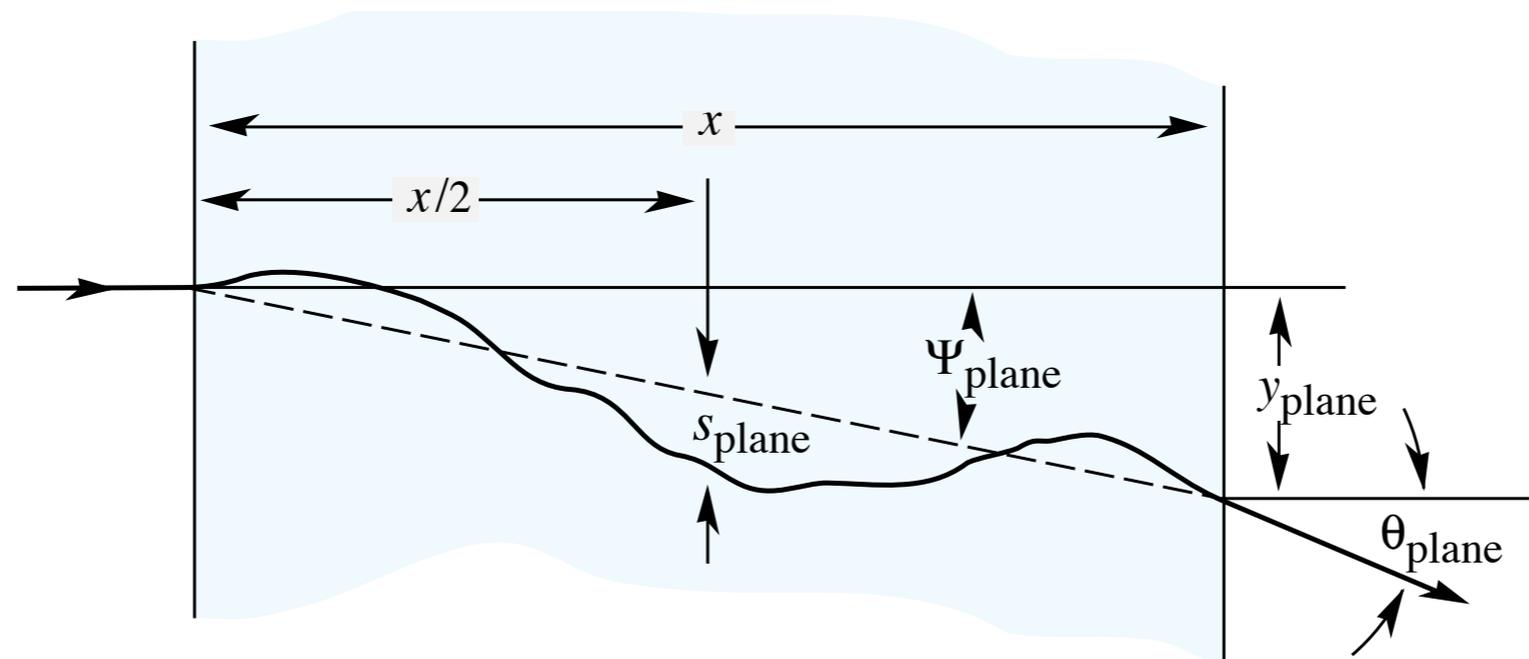
Range momentum fractional residual



True  $\nu_\mu$  CC events with contained track

If the longest reco track exits the detector, estimate its momentum from multi-Coulomb scattering (MCS).

Divide the track into segments of equal length. Fit a straight line to each segment. Scattering angle is angle between successive segments.



Ref: <http://pdg.lbl.gov/2014/reviews/rpp2014-rev-passage-particles-matter.pdf>

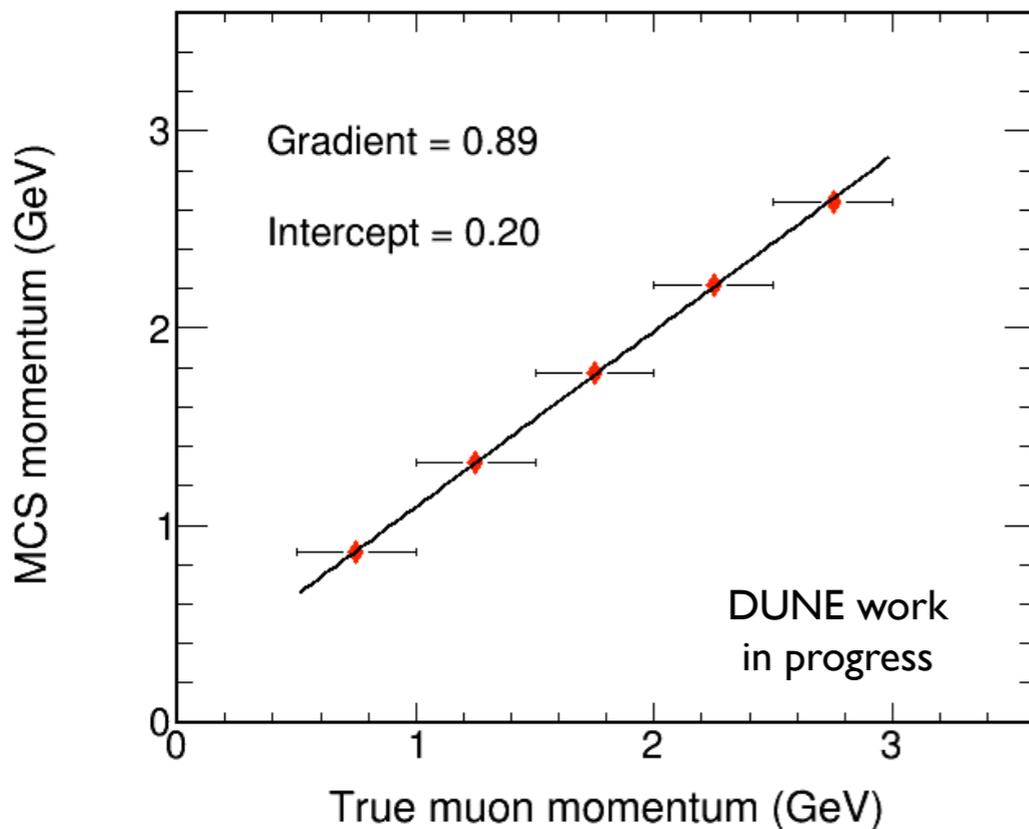
At high momentum, the scattering angle is small, while it is larger at low momentum.

Refs: <https://arxiv.org/abs/1703.06187>, <https://arxiv.org/pdf/hep-ex/0606006v1.pdf>

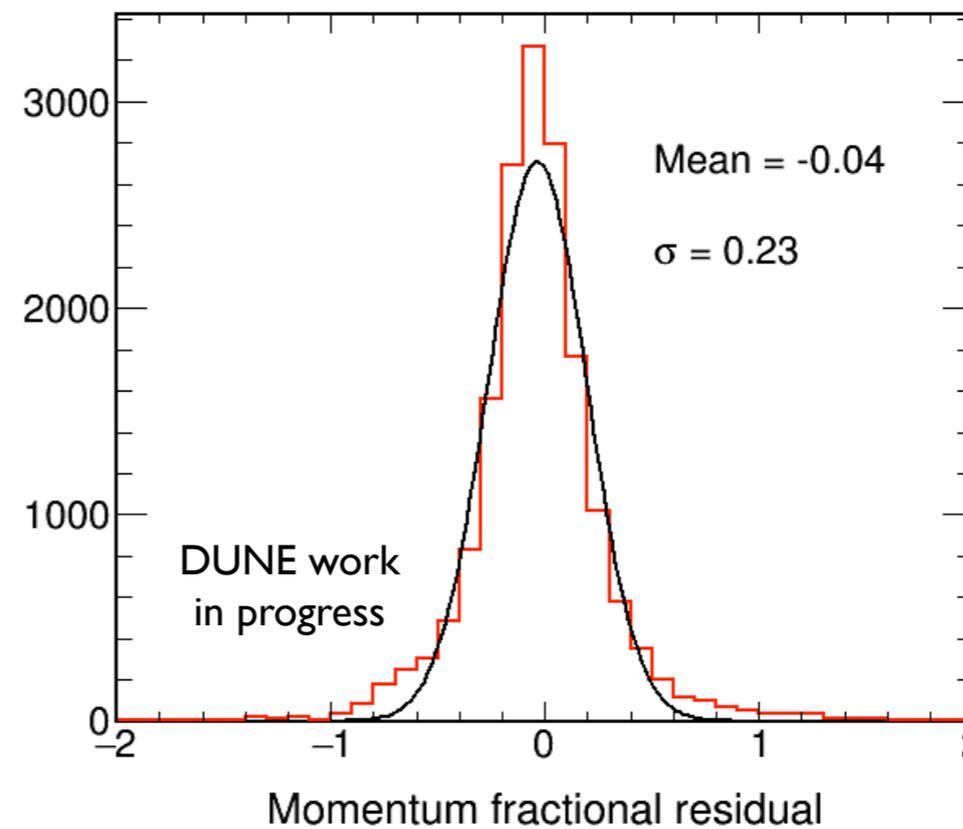
Calibrate MCS momentum using Monte Carlo.

Estimate reco track momentum as  $(\text{MCS} - \text{intercept}) / \text{gradient}$

MCS momentum calibration



MCS momentum fractional residual

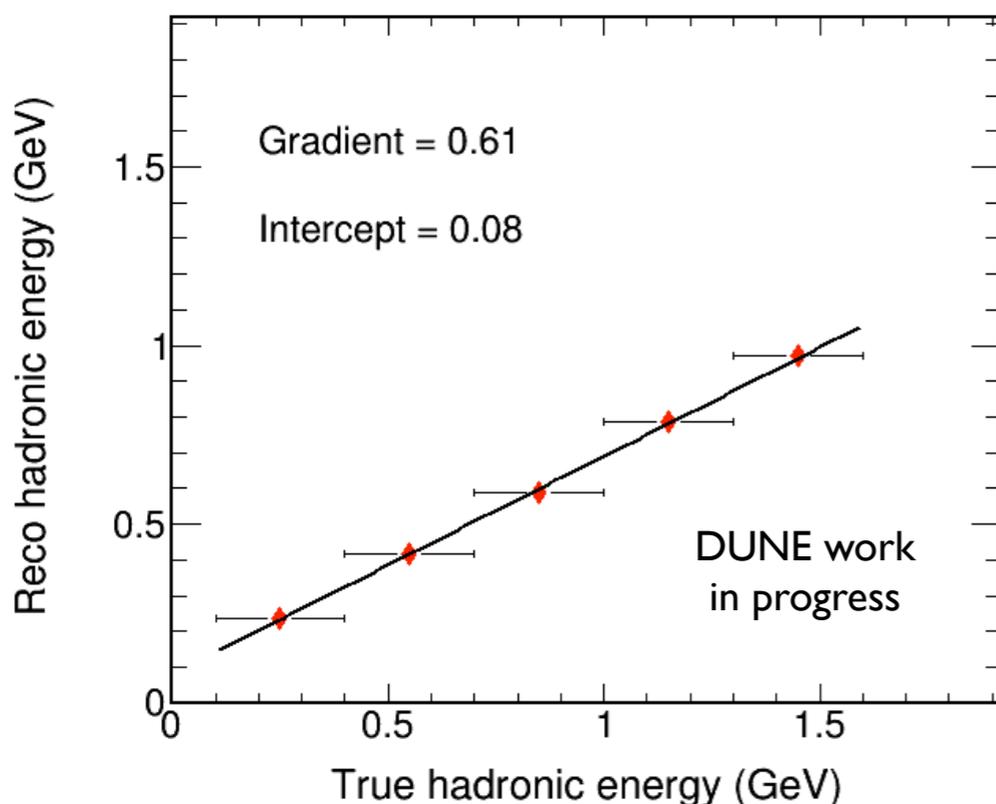


True  $\nu_\mu$  CC events with exiting track

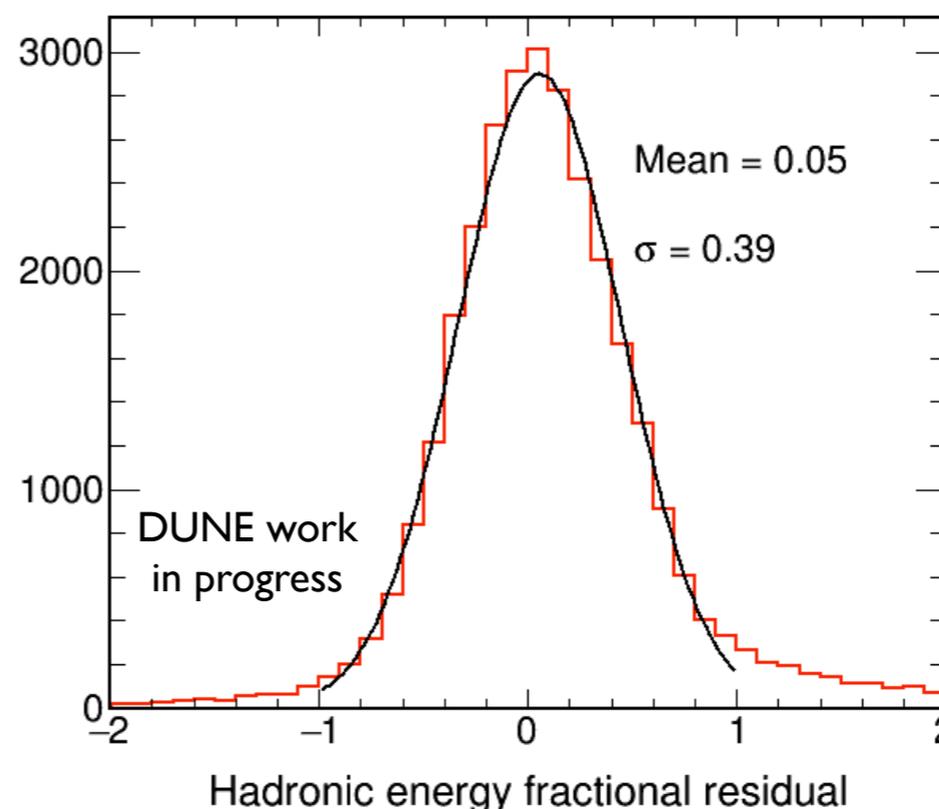
Estimate the hadronic energy from reconstructed hits that are not in the longest track. Make calibration using Monte Carlo.

Reco hadronic energy tends to be too low since neutral particles are not reconstructed in the DUNE far detector. On average 40% of the energy is missing, but there are fluctuations in this from event to event, and this limits energy resolution.

Hadronic energy calibration



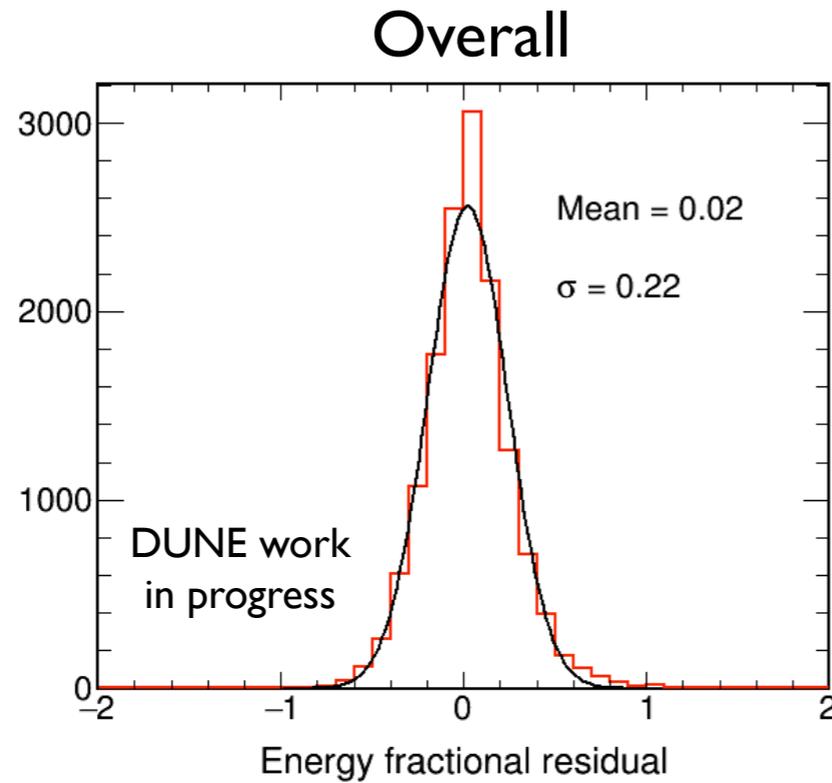
Hadronic energy fractional residual



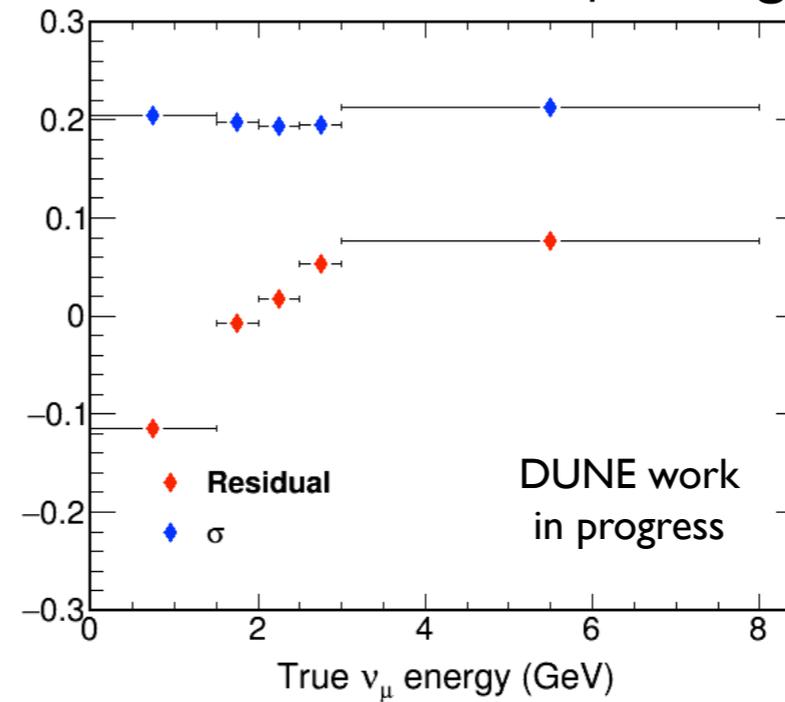
True  $\nu_\mu$  CC events.

Tail  $< -1$  in resolution plot is due to subtraction of intercept (which is  $> 0$ ) when making correction.

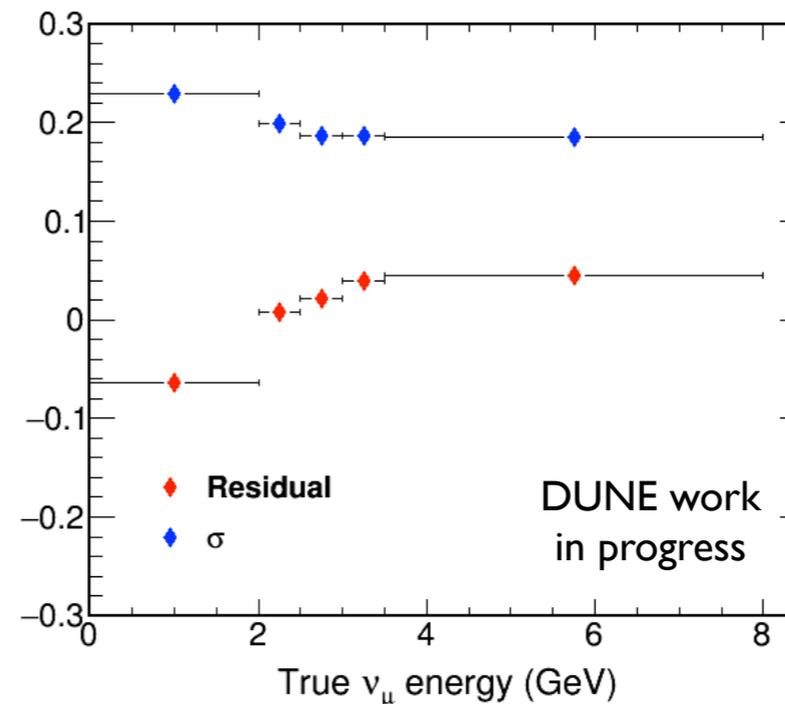
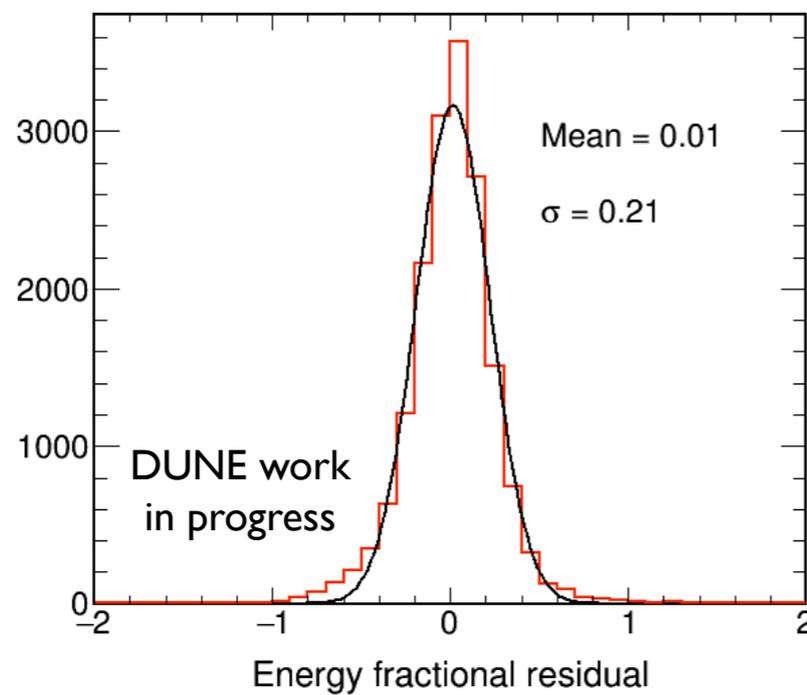
True  $\nu_\mu$  CC events with contained track



As function of true  $\nu_\mu$  energy

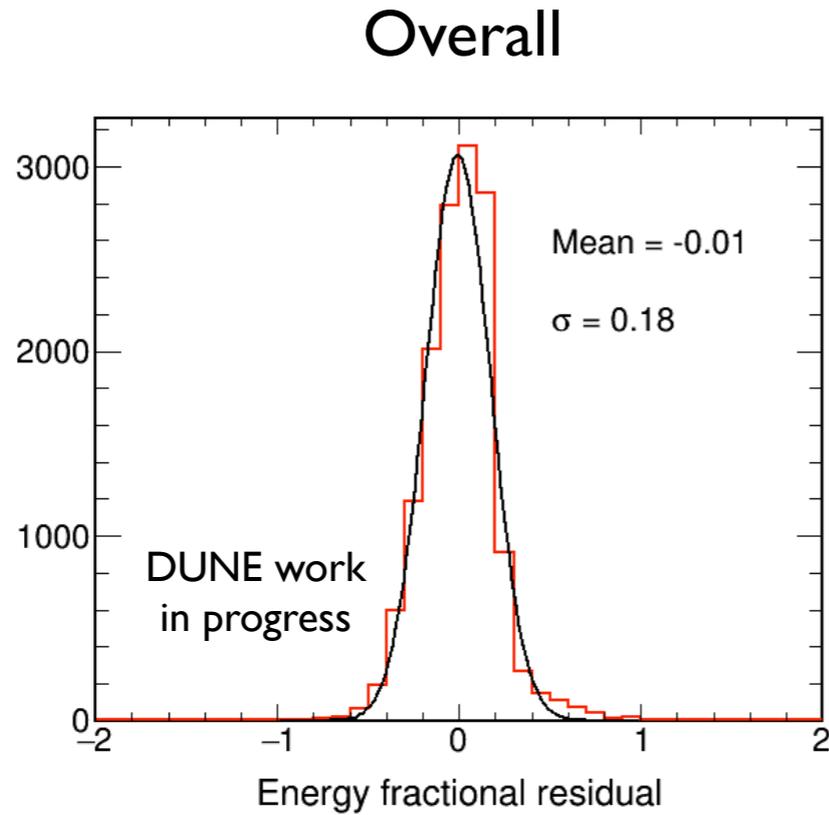


True  $\nu_\mu$  CC events with exiting track

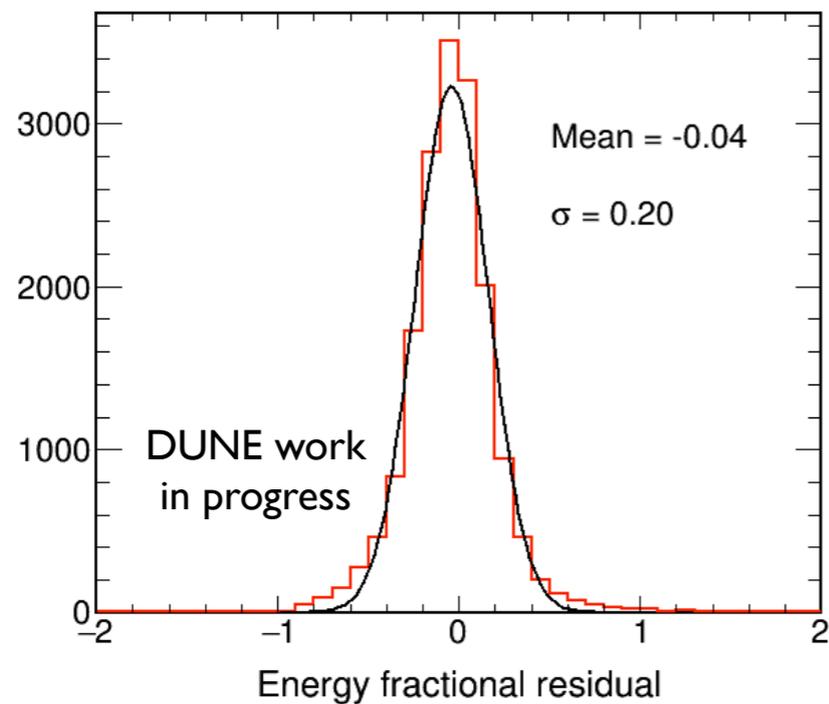


There is a clear bias as a function of true  $\nu_\mu$  energy.

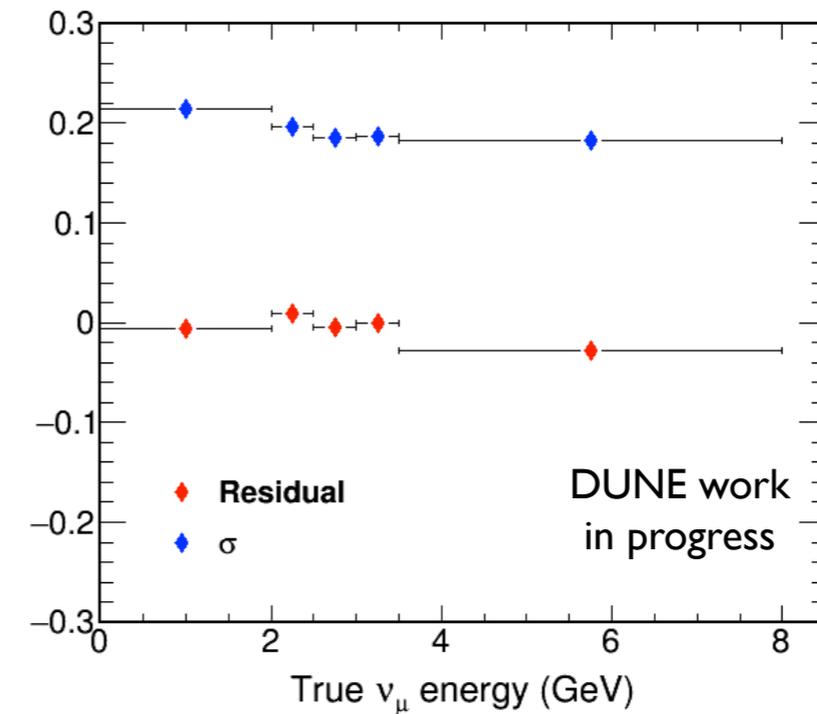
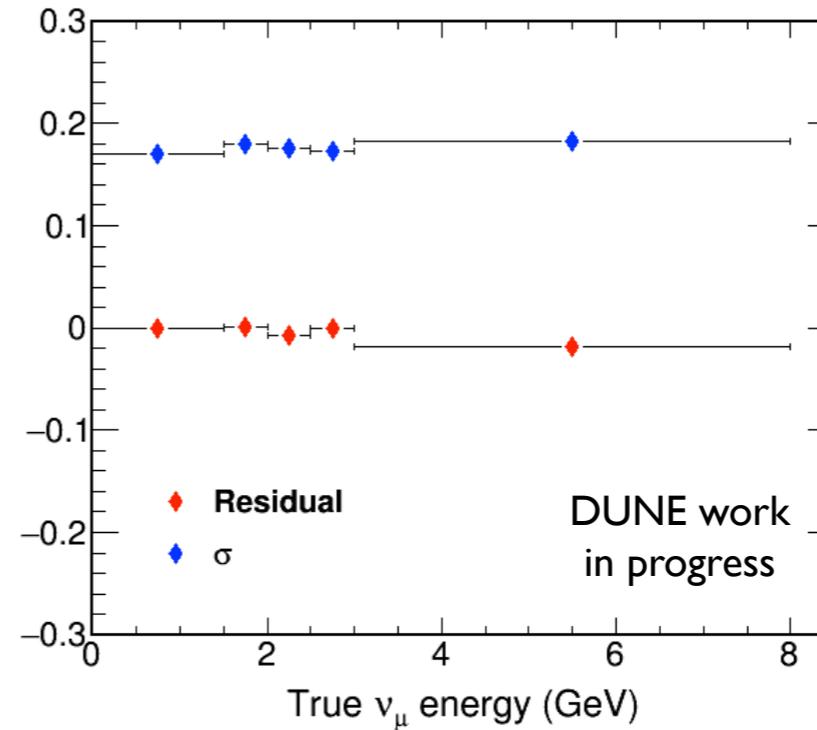
True  $\nu_\mu$  CC events with contained track



True  $\nu_\mu$  CC events with exiting track



As function of true  $\nu_\mu$  energy

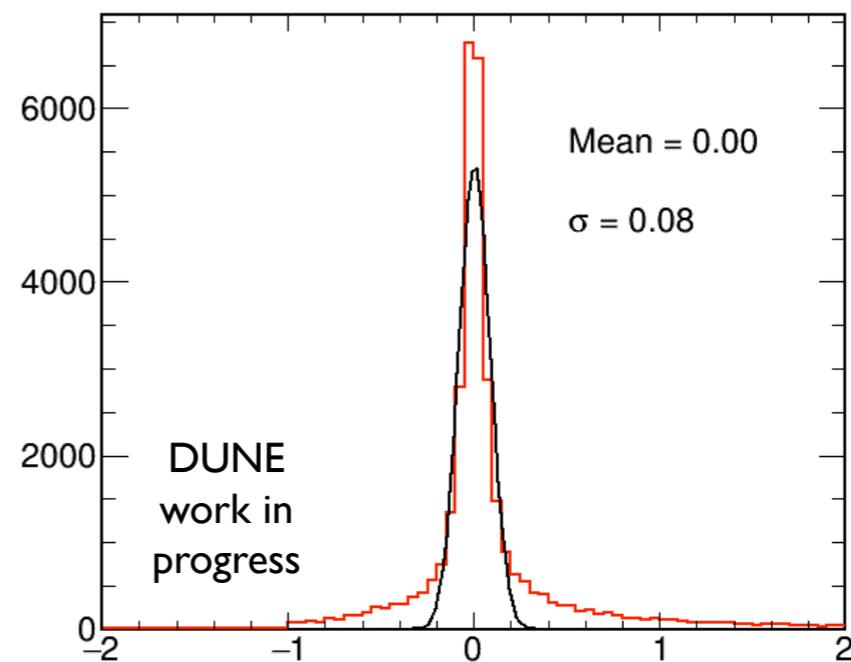


Make ad hoc tweaks of gradient and intercept of correction of hadronic energy.

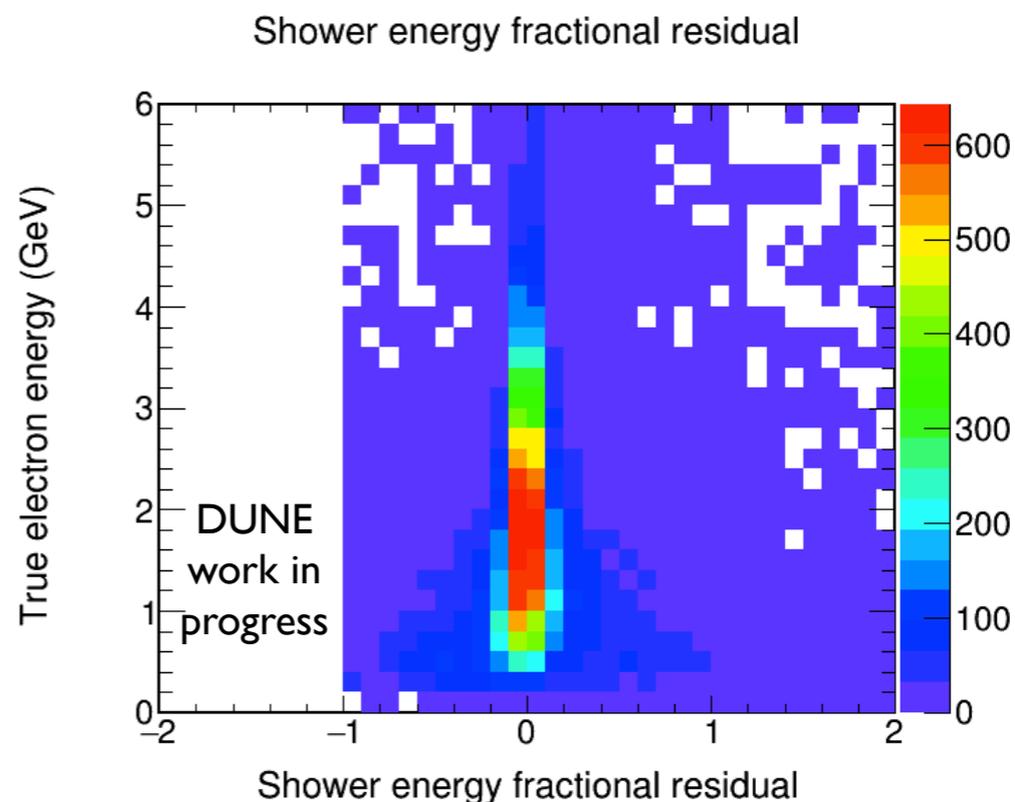
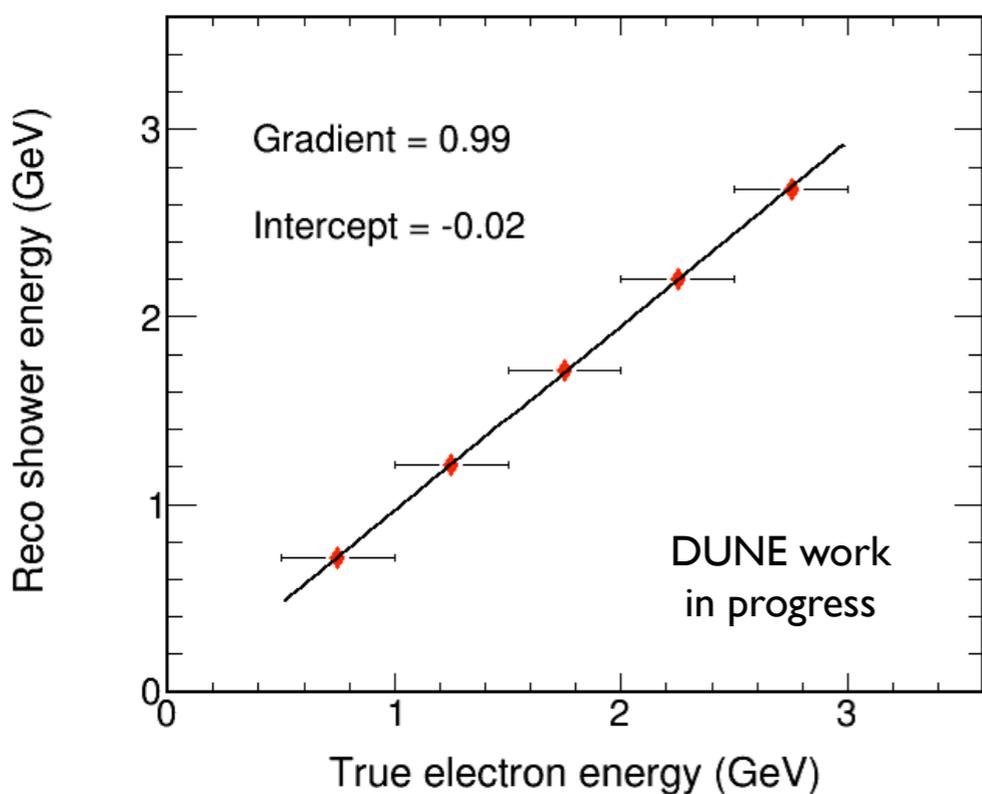
This reduces bias as a function of  $\nu_\mu$  energy.

Make calibration of energy of reconstructed shower with highest charge using Monte Carlo.

Shower energy fractional residual

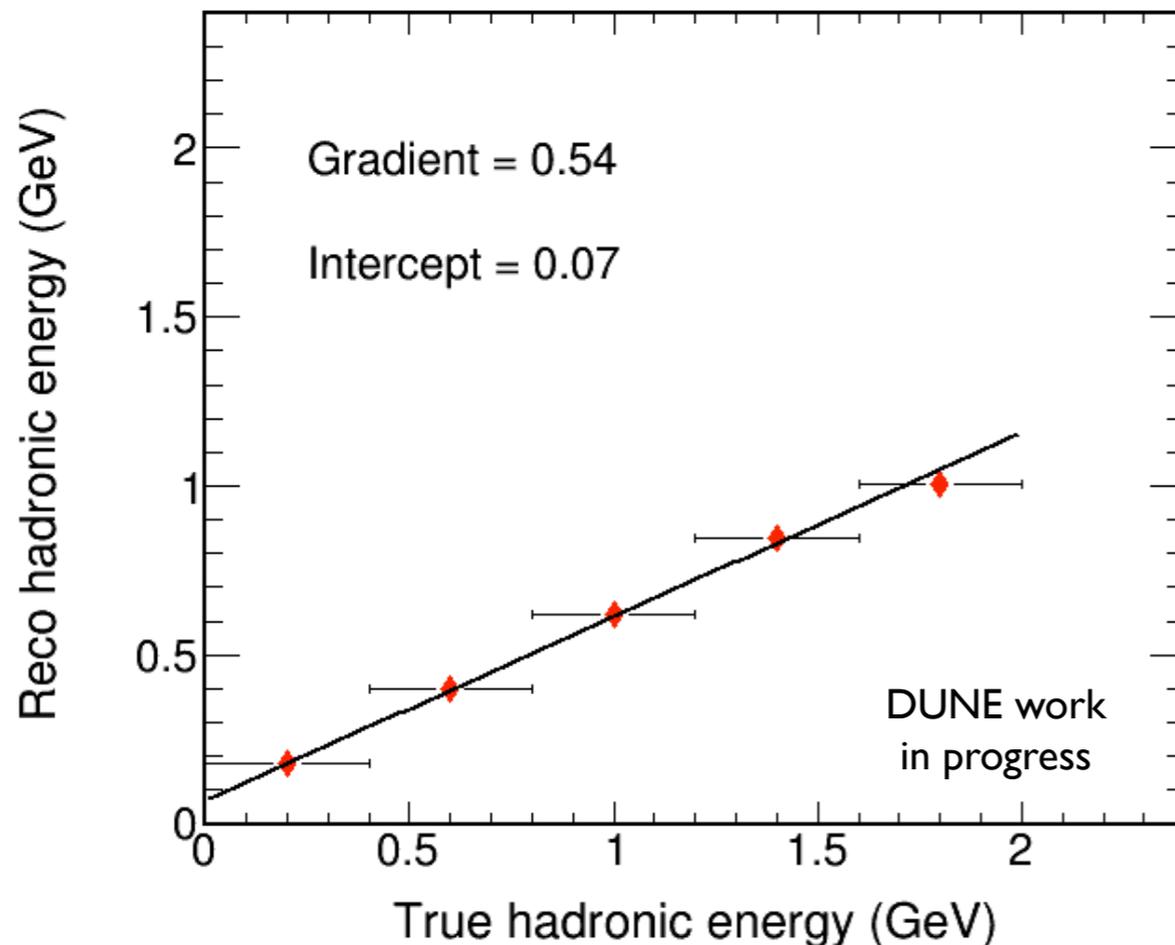


Shower energy calibration

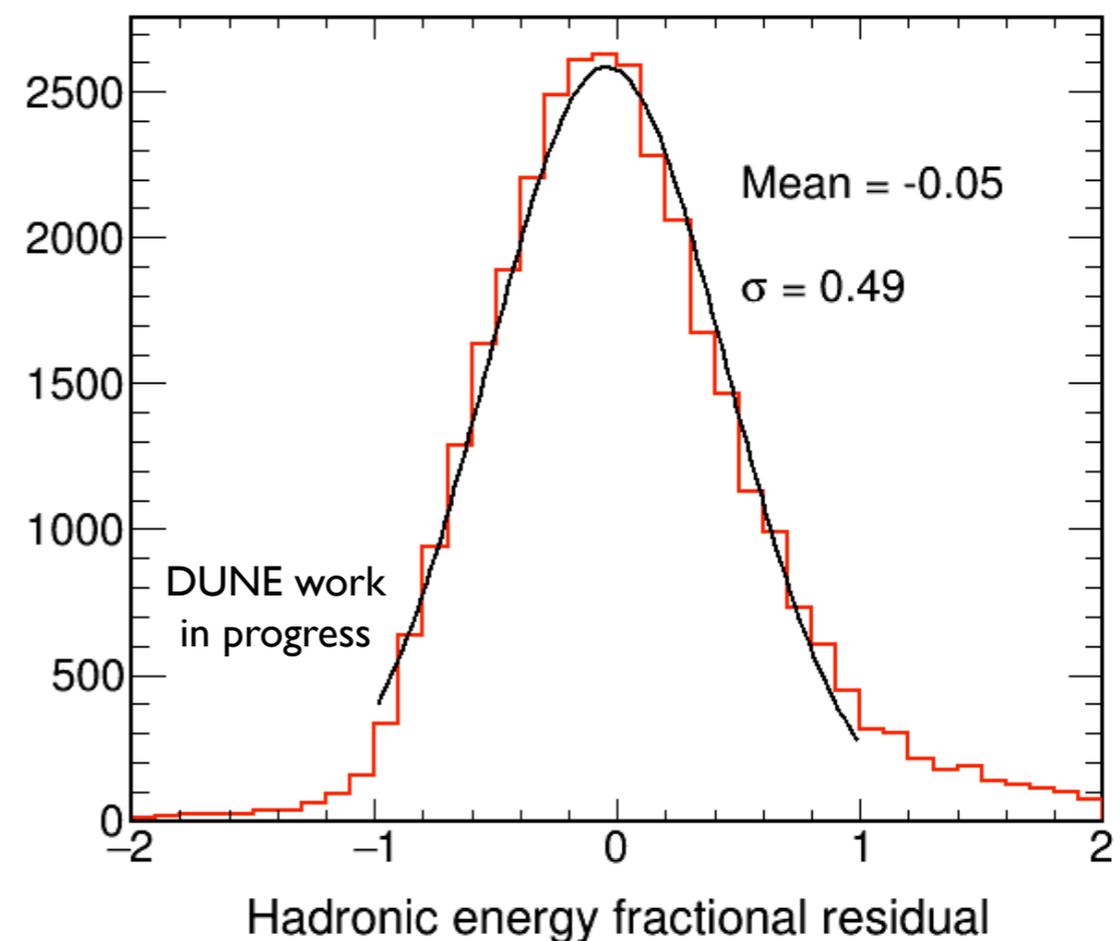


Also make a calibration of  $\nu_e$  CC reco hadronic energy using Monte Carlo. Again the missing energy from neutral particles fluctuates from event to event, which limits the overall energy resolution.

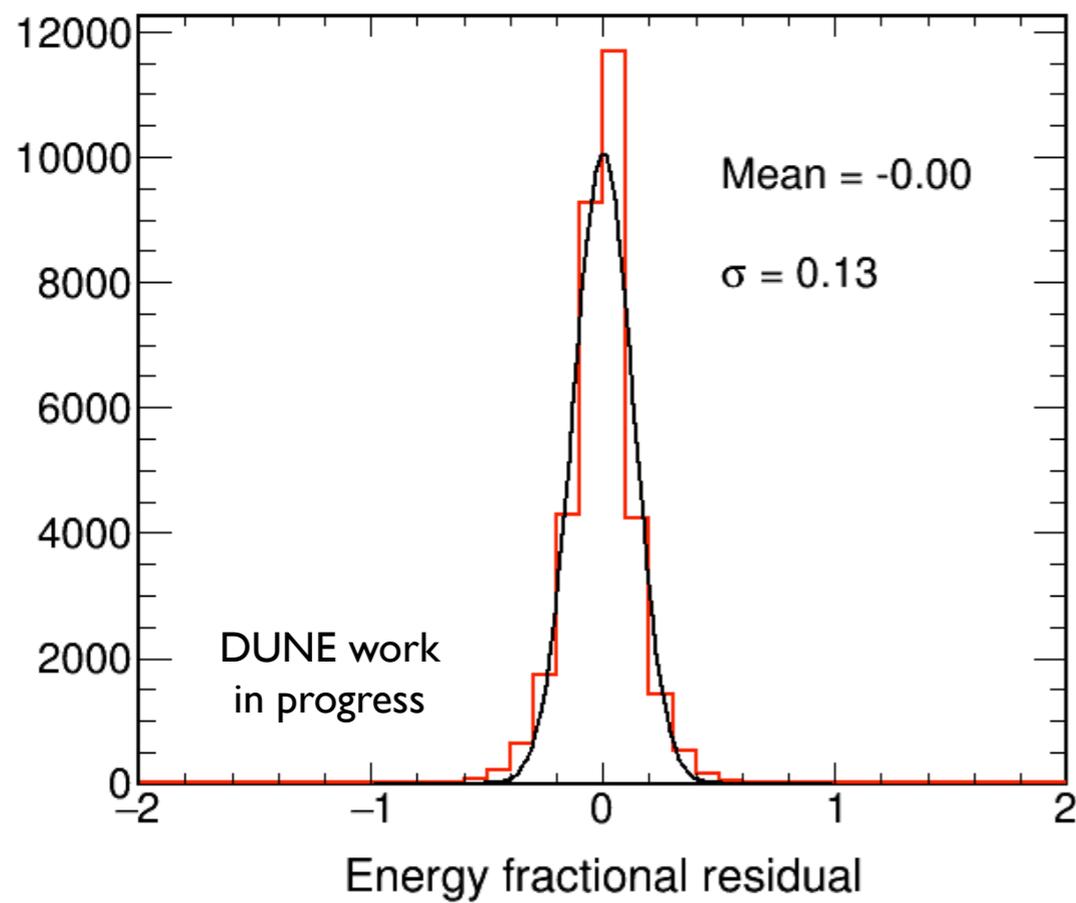
### Hadronic energy calibration



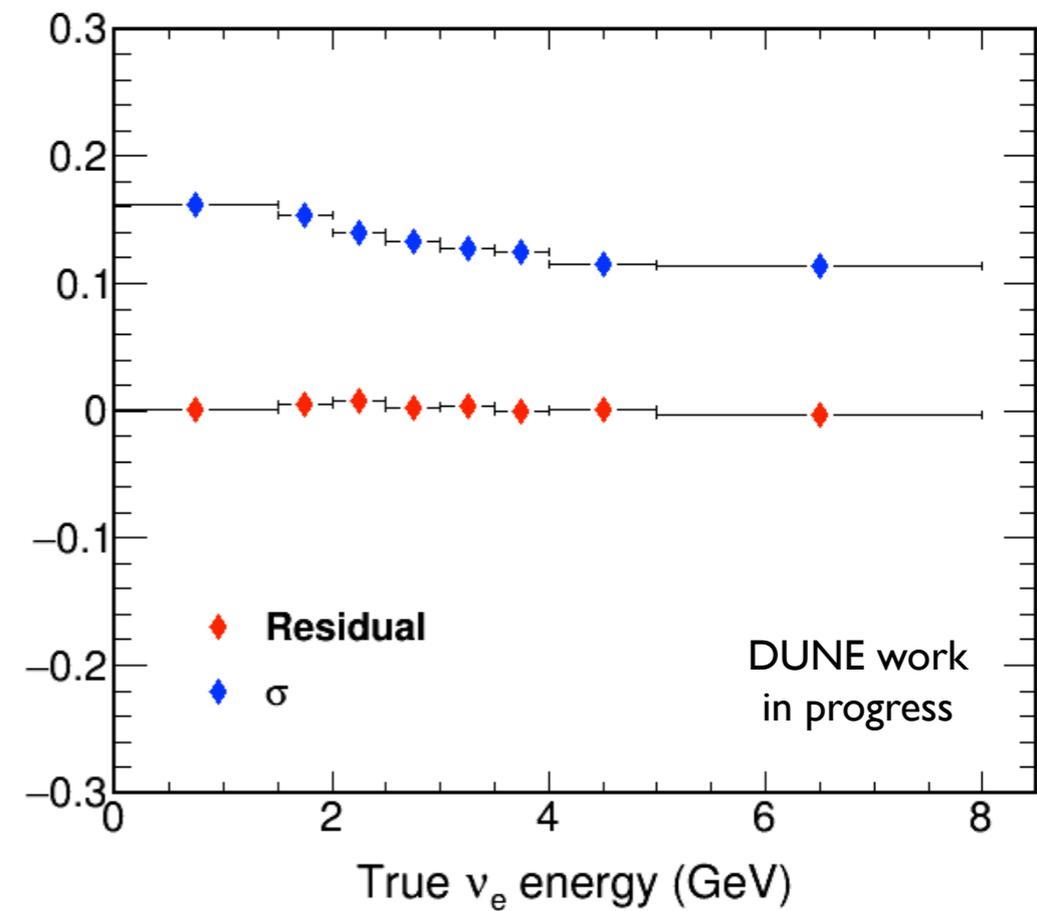
### Hadronic energy fractional residual



## Overall $\nu_e$ energy resolution



## $\nu_e$ energy resolution as function of true energy





## Conclusion



We have implemented a complete first version of neutrino energy reconstruction in the DUNE far detector. This reconstruction is done by making separate estimates of the lepton and hadronic energies in CC events. We achieve overall energy resolutions of 20% for  $\nu_\mu$  CC and 13% for  $\nu_e$  CC events.

The resolutions of each component of the neutrino energy reconstruction in % are summarised below:

	$\nu_\mu$ CC	$\nu_e$ CC
Longest reco track (contained)	5	-
Longest reco track (exiting)	20	-
Reco shower with highest charge	-	8
Hadronic energy	39	49

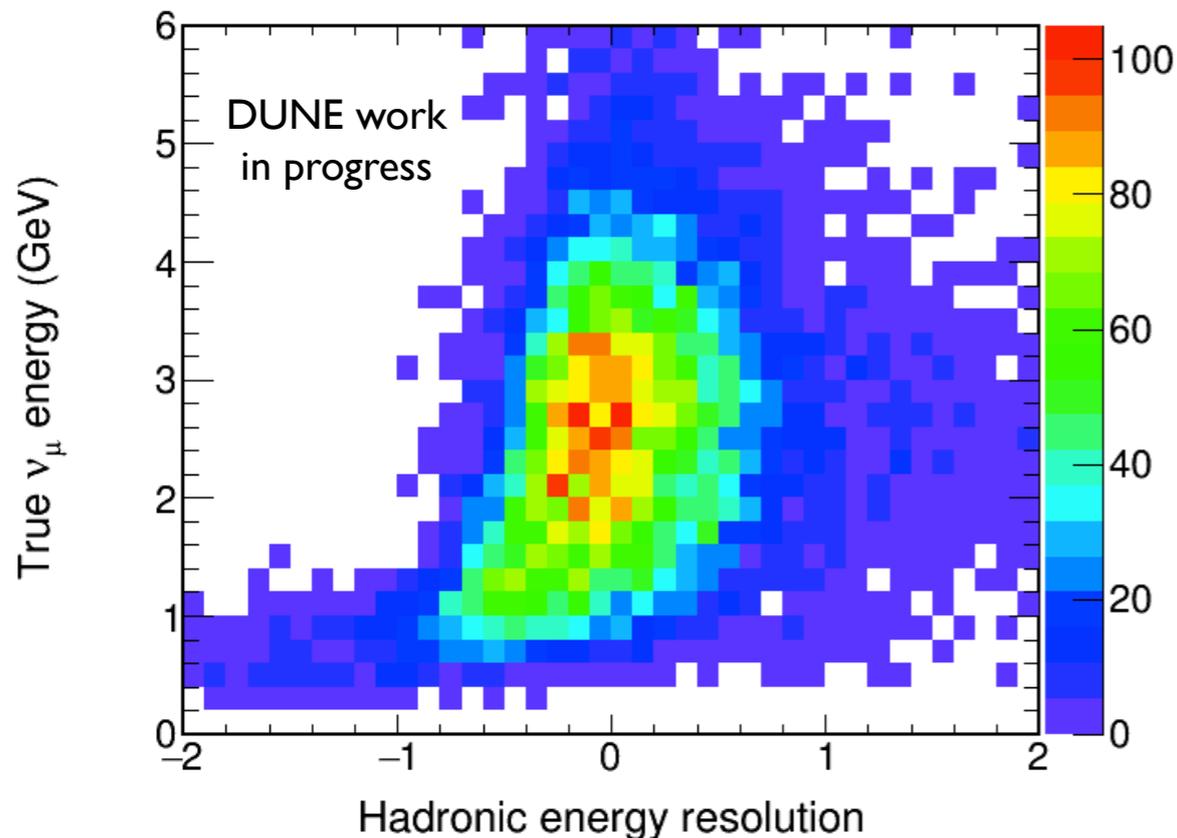


# BACKUP SLIDES

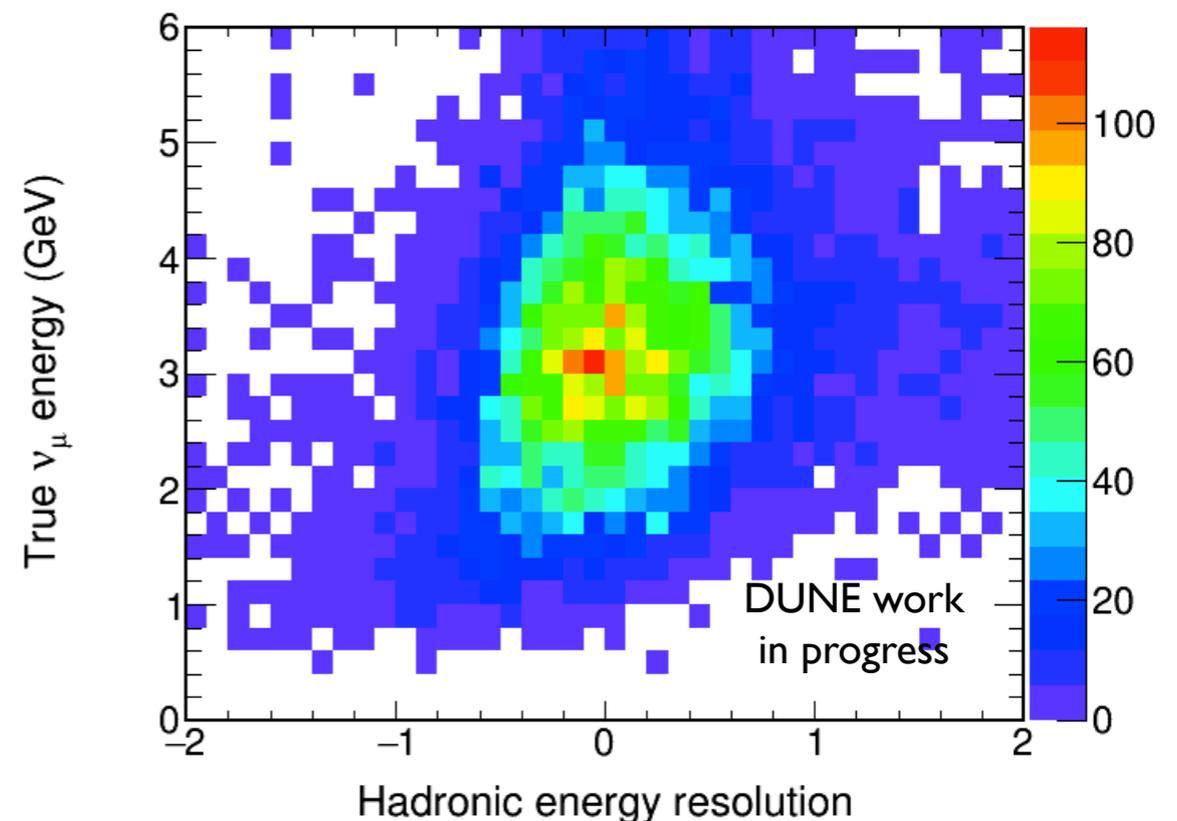
What is reason for bias as a function of true  $\nu_\mu$  energy ?

There is no bias in reco momentum by range or MCS momentum as a function of either true muon momentum or true  $\nu_\mu$  energy. Nor is there a bias in reco hadronic energy as a function of true hadronic energy (please see backup). The culprit is hadronic energy fractional residual as a function of true  $\nu_\mu$  energy:

True  $\nu_\mu$  CC events  
with contained track



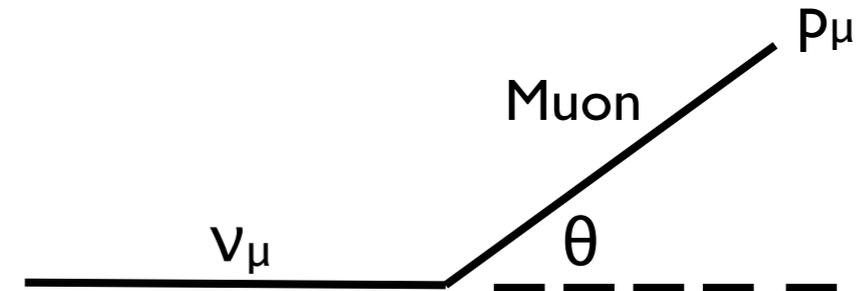
True  $\nu_\mu$  CC events  
with exiting track



Bias as a function of true energy is worse for events with contained tracks and worse at low true energy. It is also worse for  $\nu_\mu$  CCQE events than for CC res or CC DIS events.

Try kinematic reconstruction of  $\nu_\mu$  energy to see whether it can help for some or all  $\nu_\mu$  CCQE events:

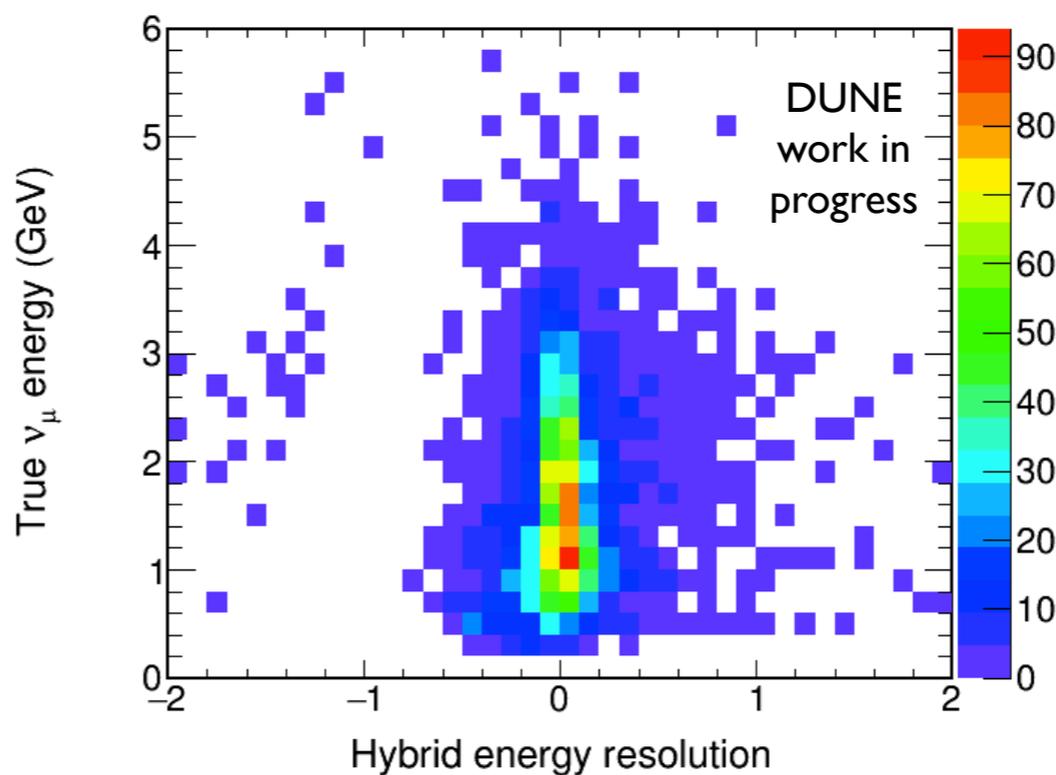
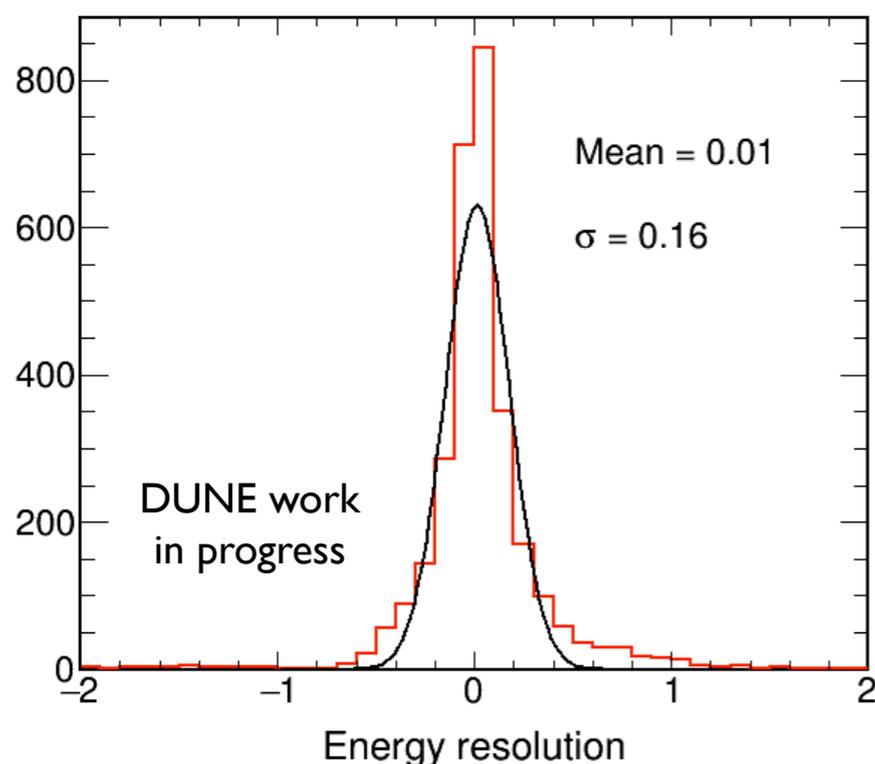
$$E_\nu = \frac{m_p^2 - (m_n - E_b)^2 - m_\mu^2 + 2(m_n - E_b)E_\mu}{2(m_n - E_b - E_\mu + p_\mu \cos\theta_\mu)}$$



$E_b$  = binding energy. This formula is only valid for CCQE events. It neglects the Fermi momentum of the struck nucleon and any possible final-state interactions.

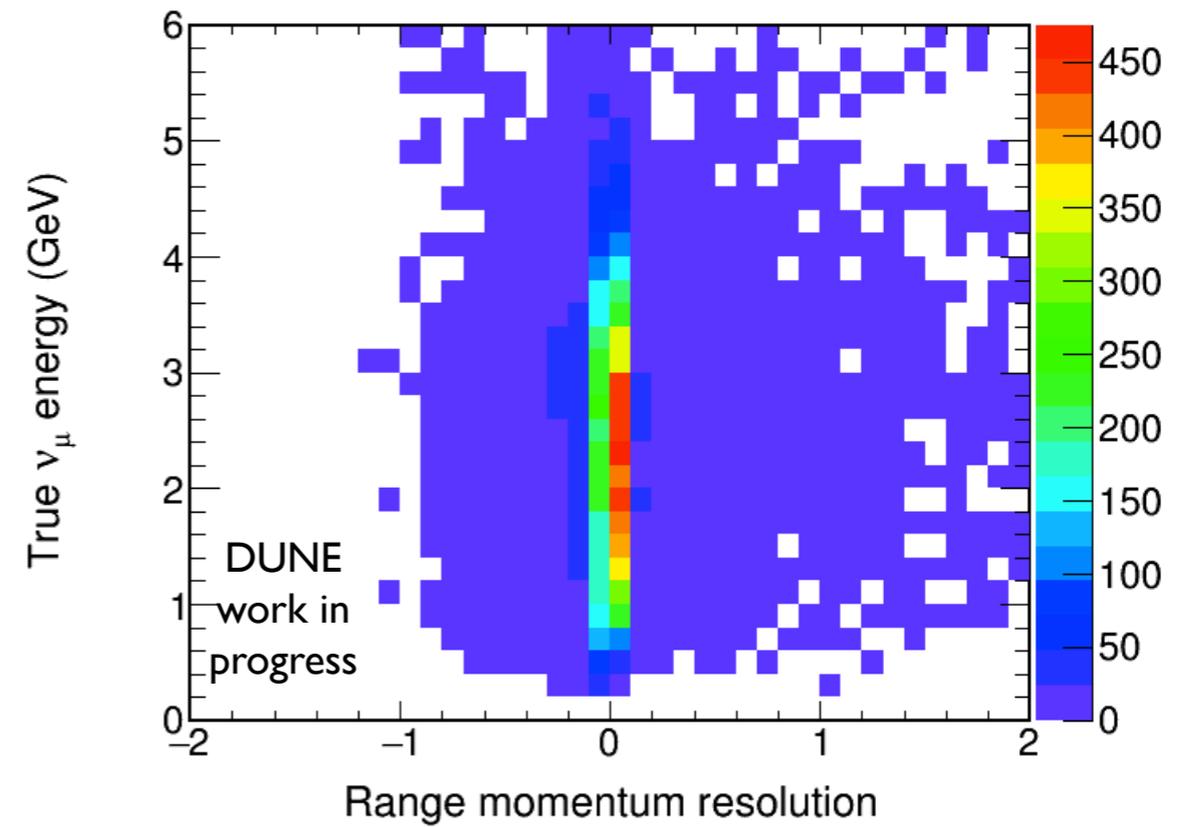
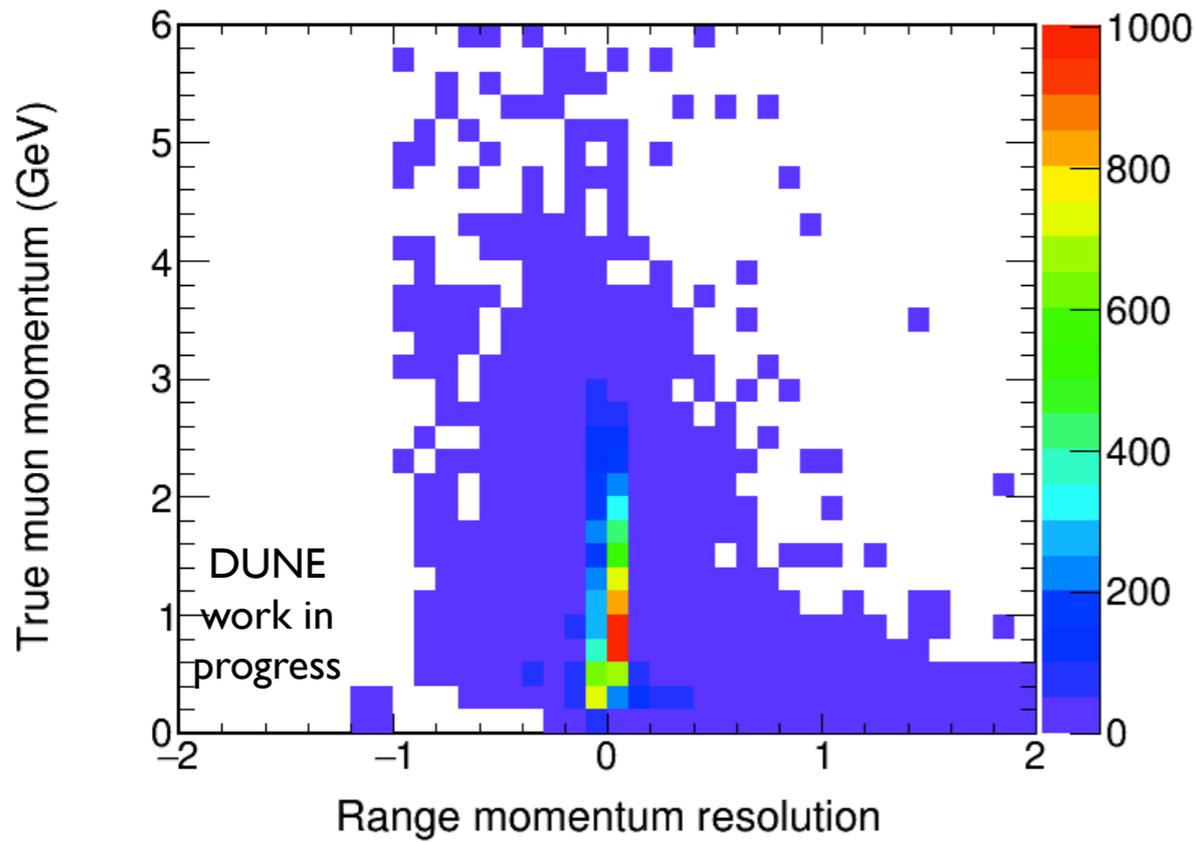
Use hybrid reconstruction for  $\nu_\mu$  CCQE events: if ratio of track + hadronic reco energy / kinematic reco energy  $> 1.5$ , use track + hadronic, otherwise use kinematic reconstruction.

True  $\nu_\mu$  CCQE events with contained track

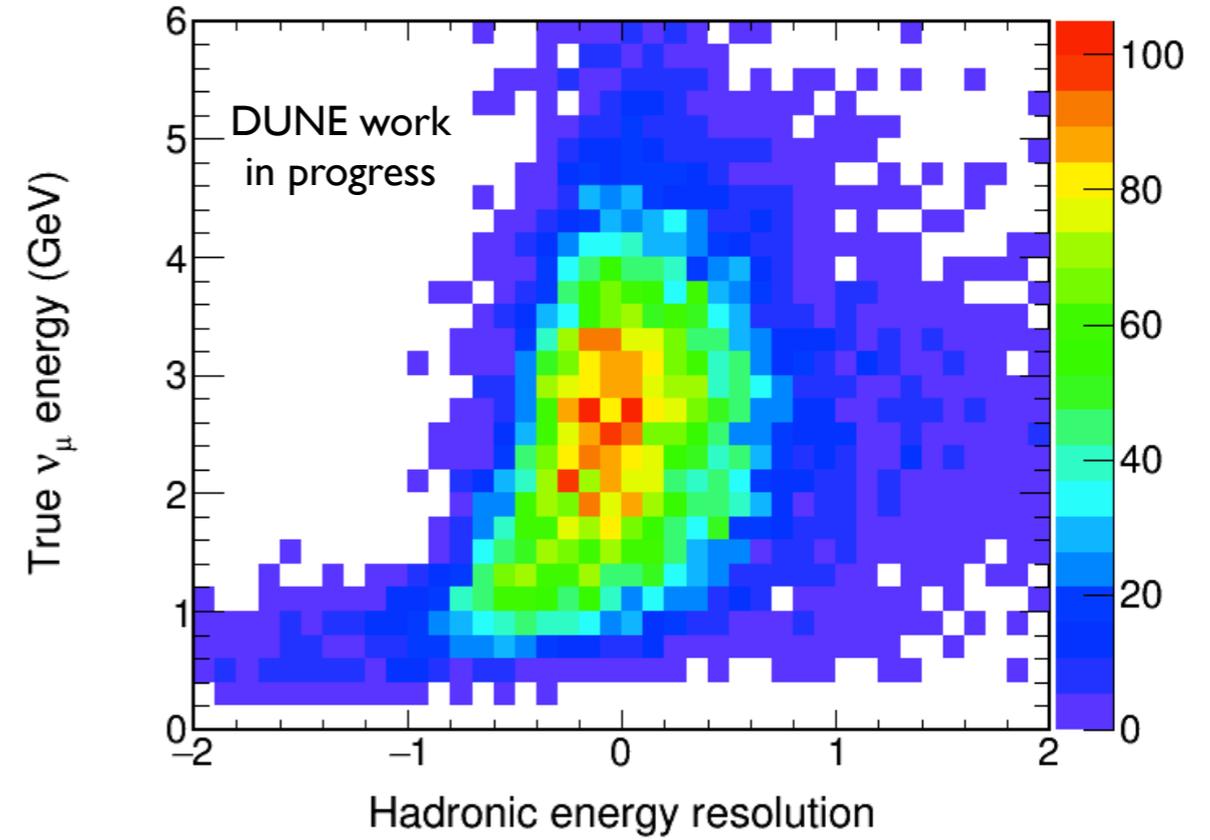
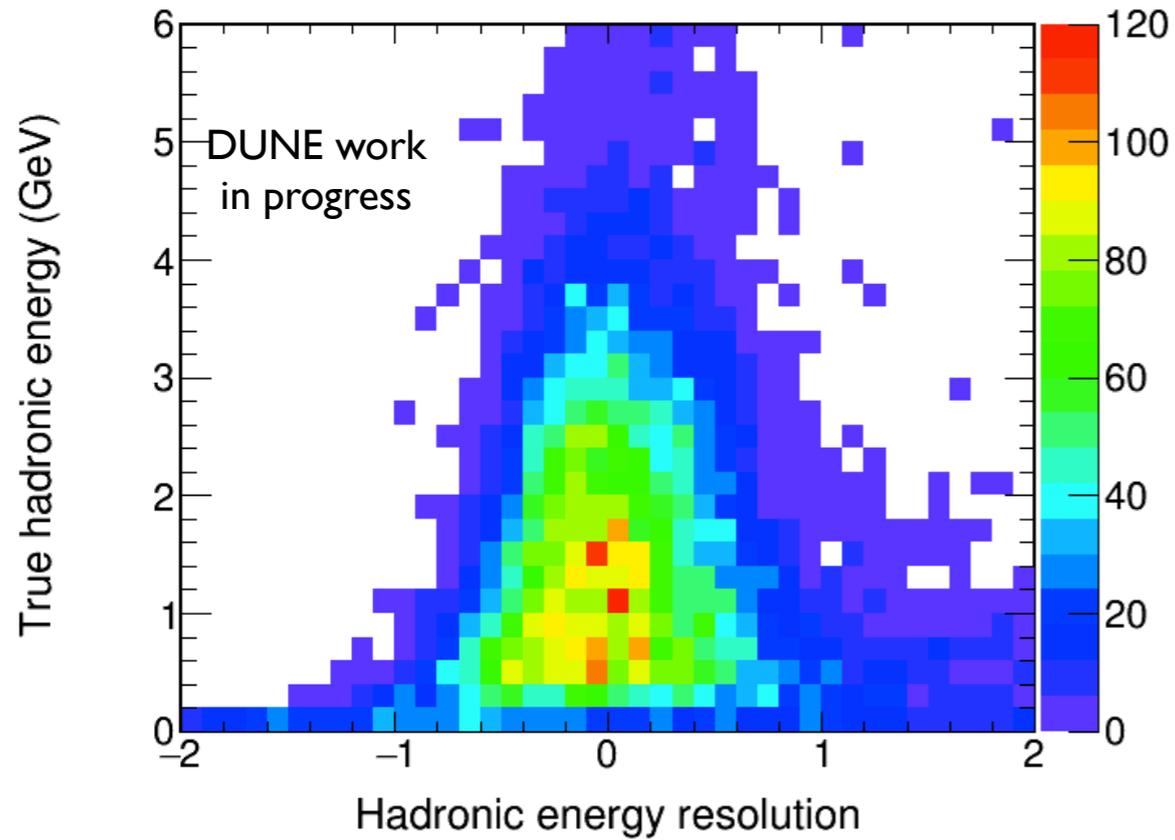


Very small dependence on true energy.

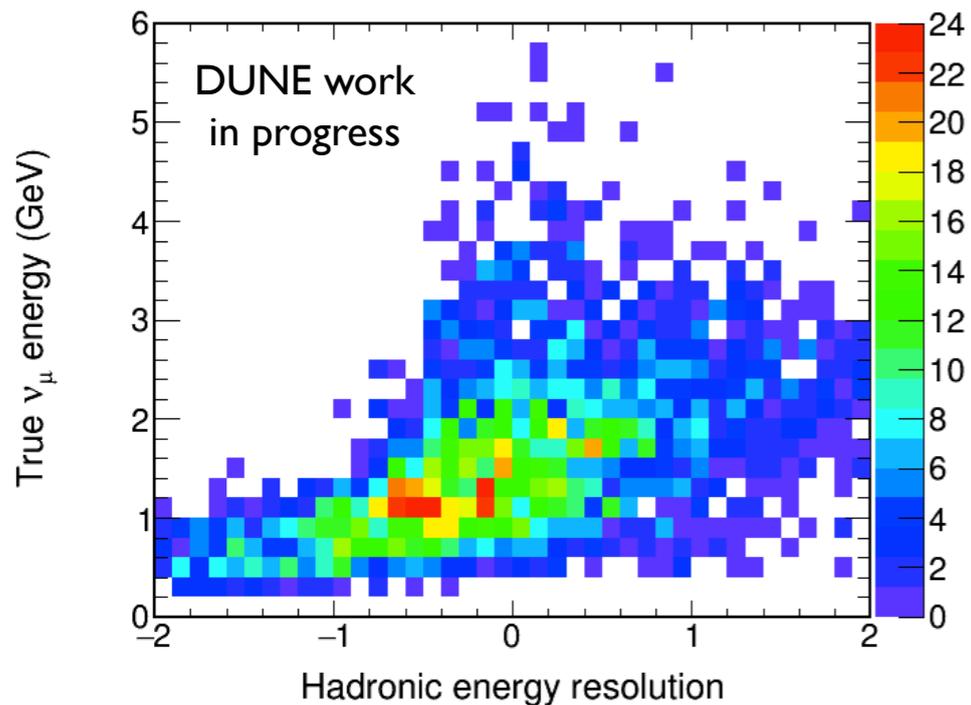
## True $\nu_\mu$ CC events with contained track



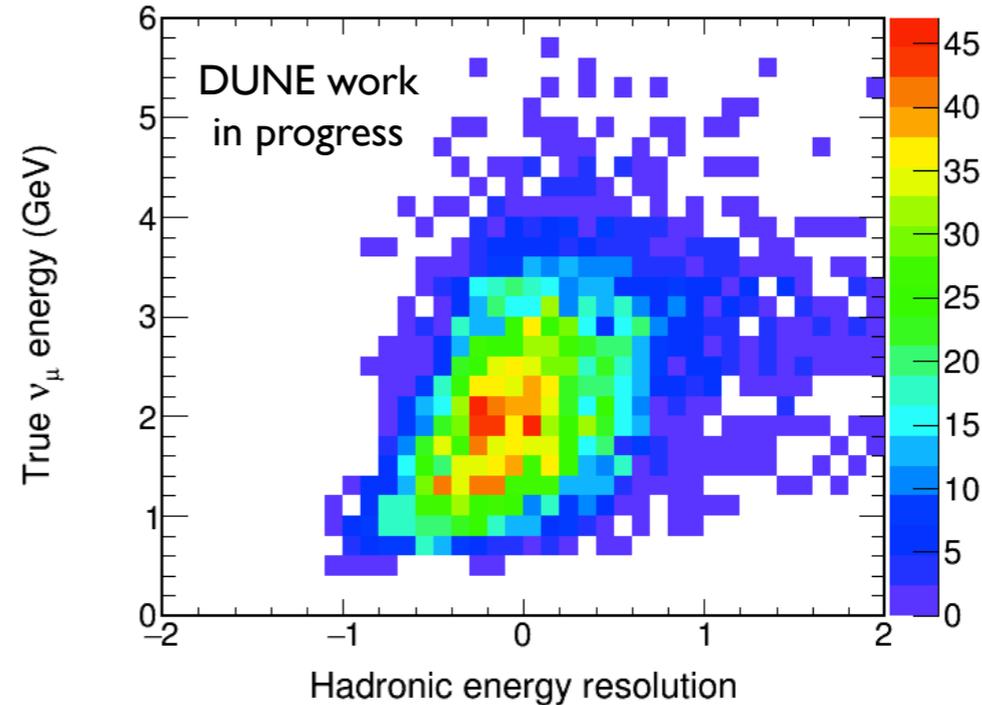
## True $\nu_\mu$ CC events with contained track



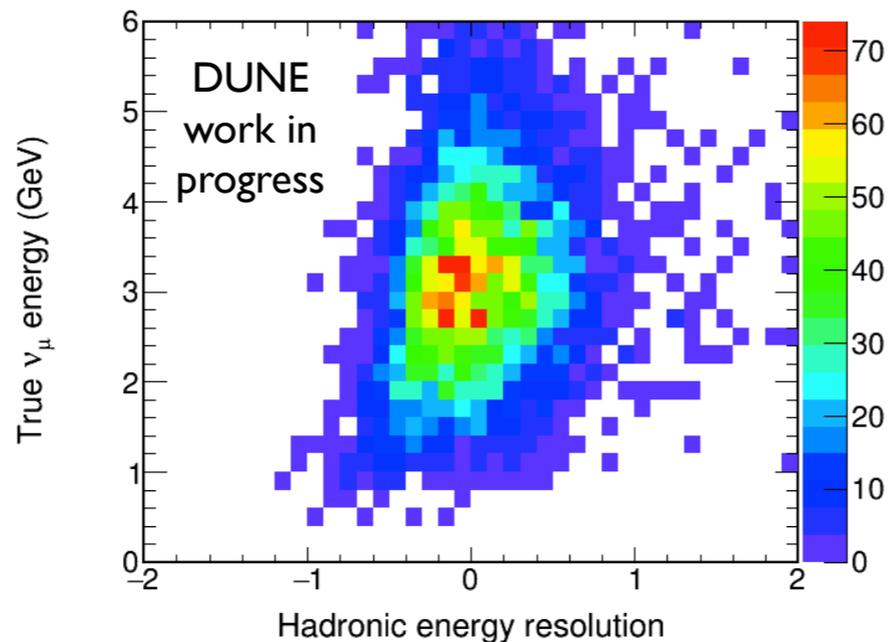
CC QE



CC Res

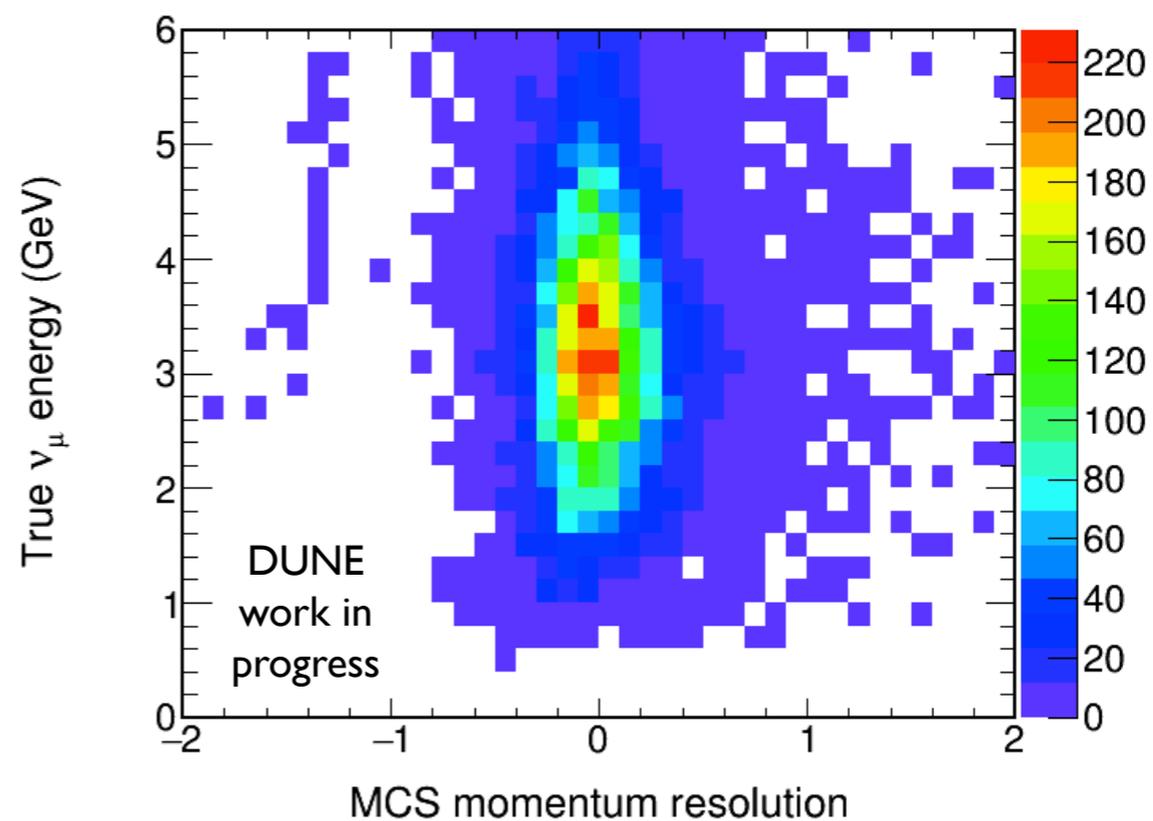
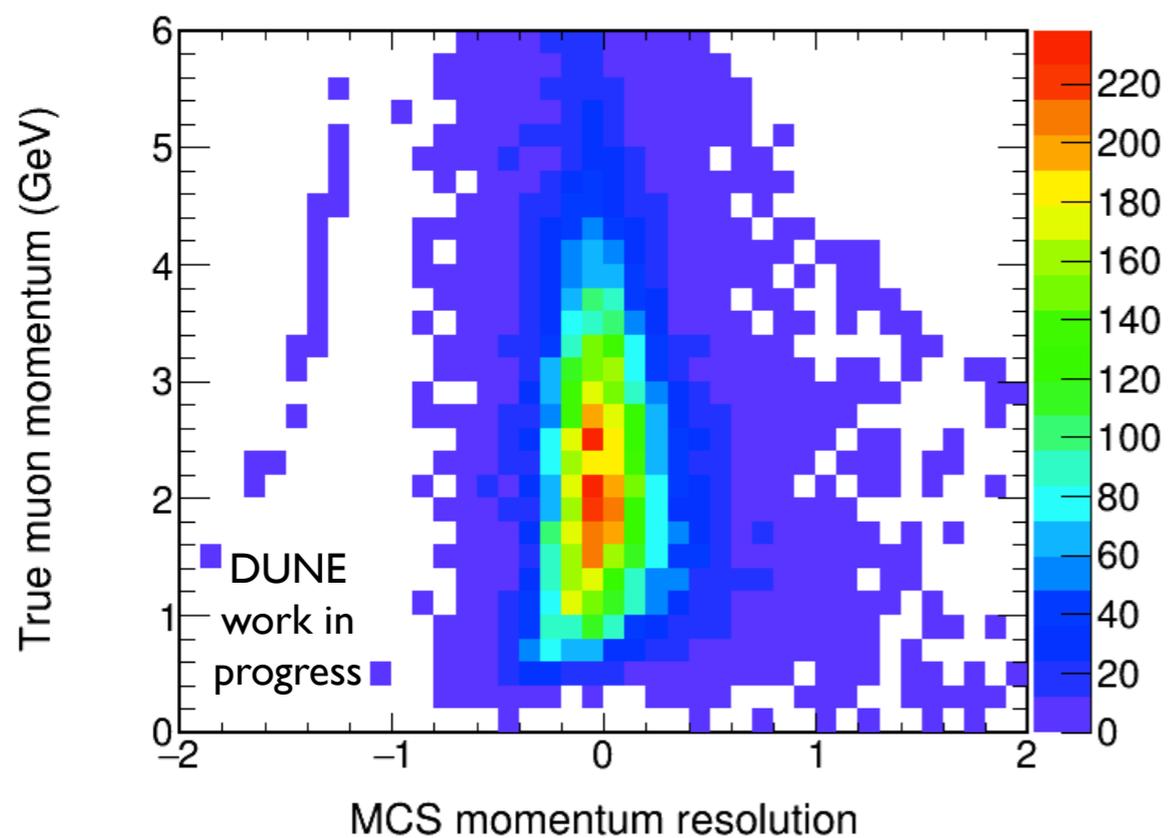


CC DIS

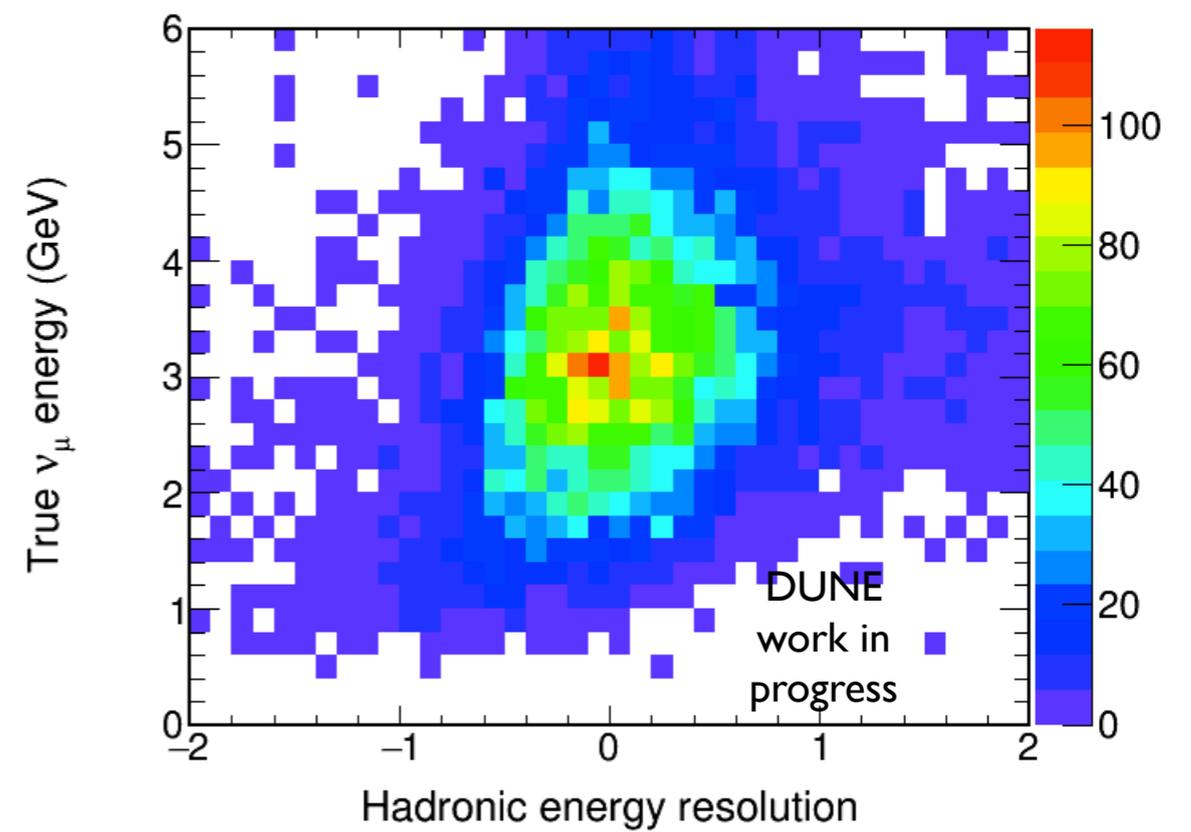
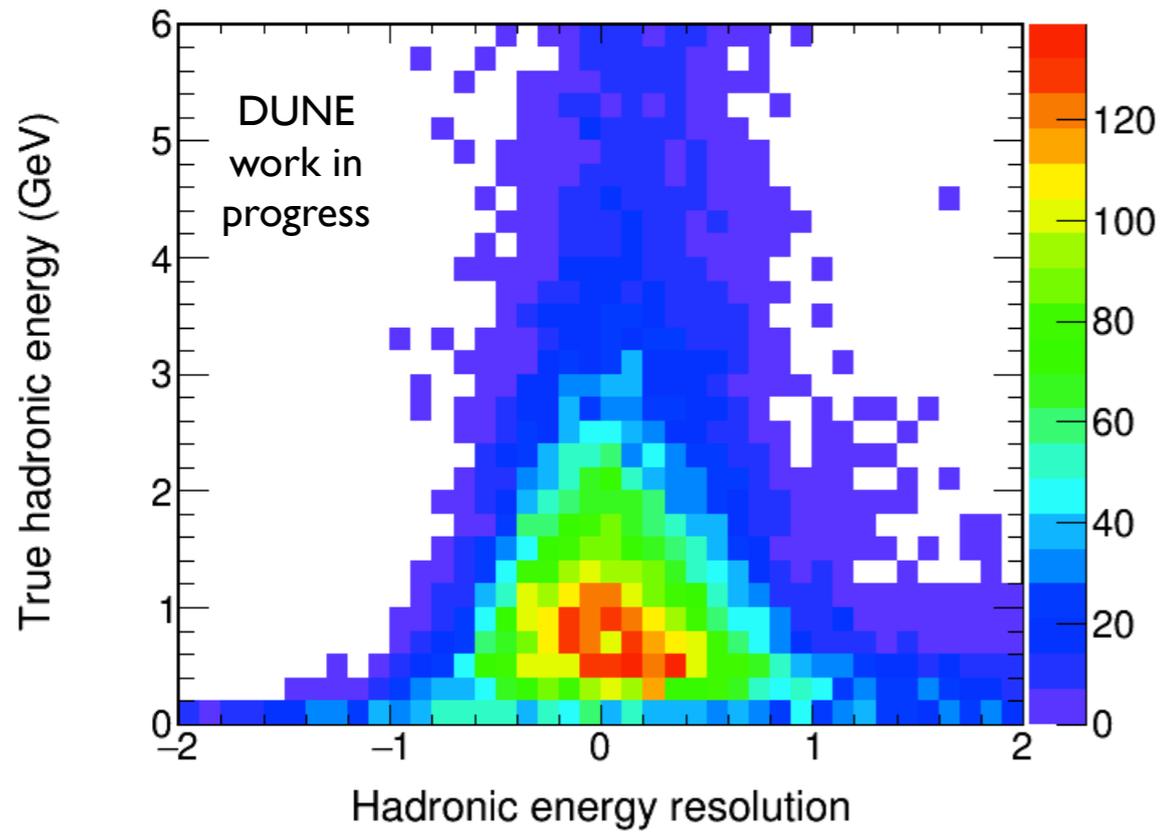


True  $\nu_\mu$  CC events  
with contained track

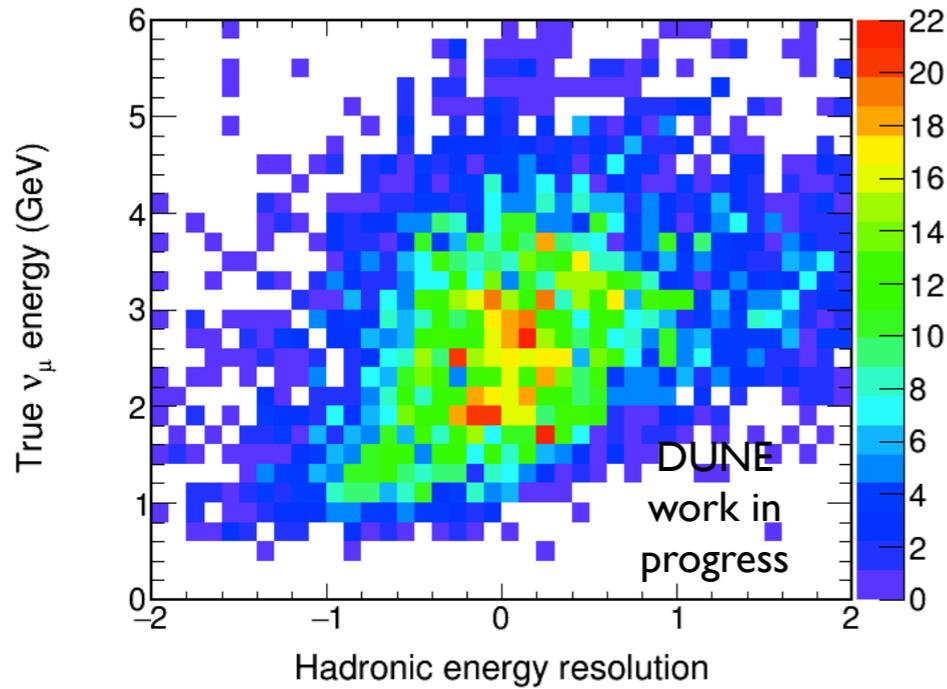
## True $\nu_\mu$ CC events with exiting track



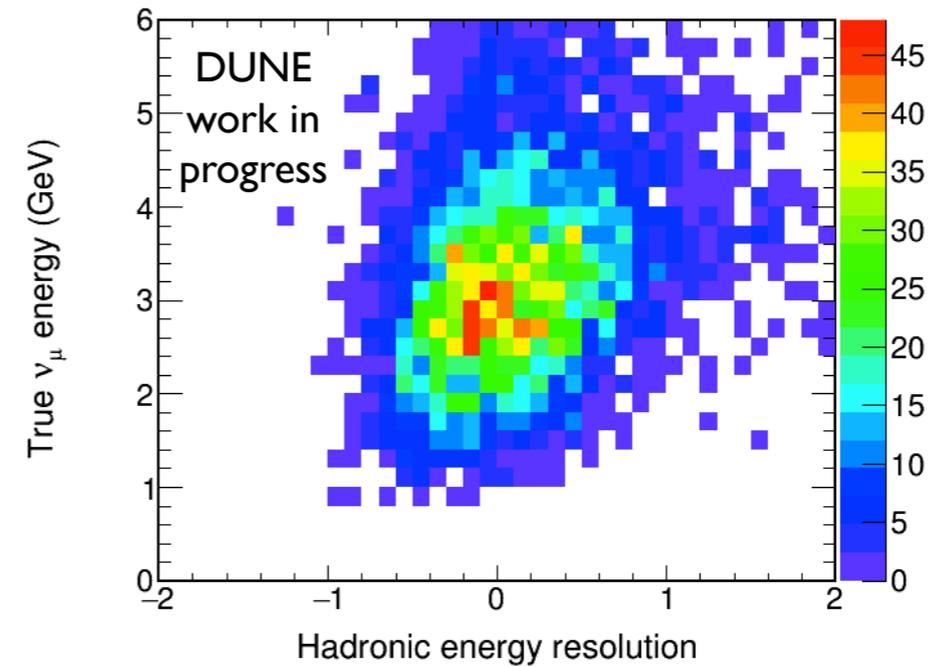
## True $\nu_\mu$ CC events with exiting track



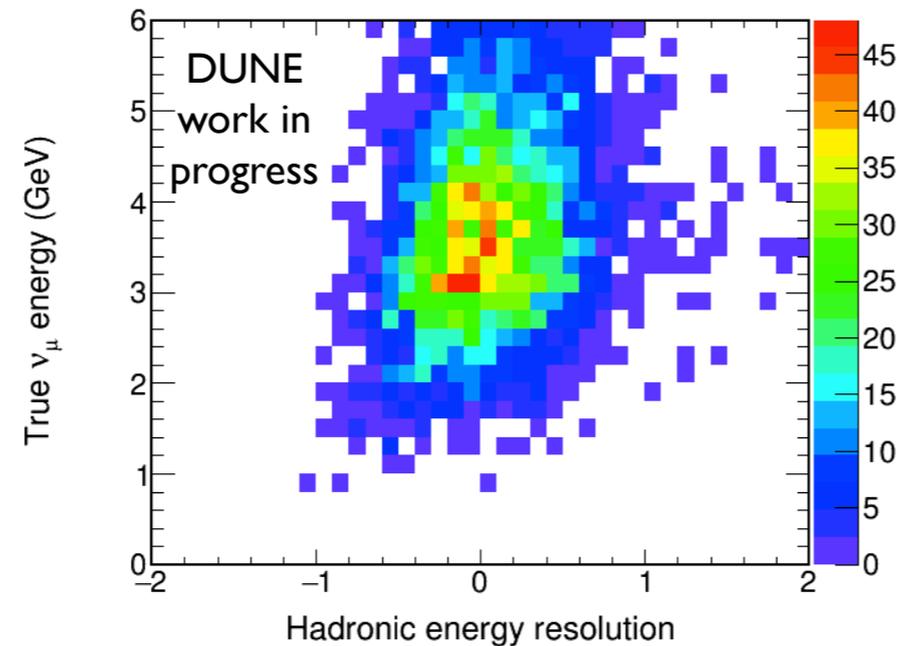
CC QE



CC Res

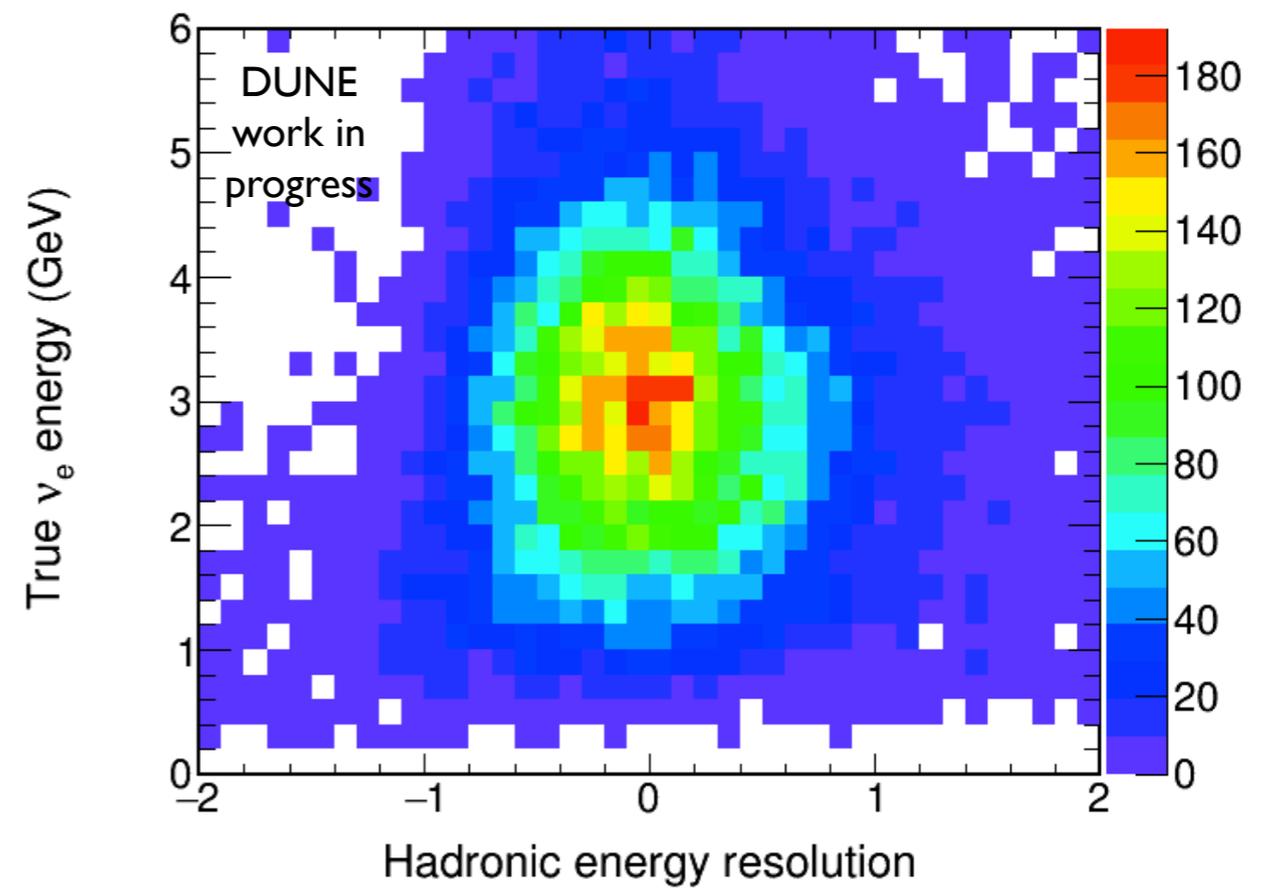
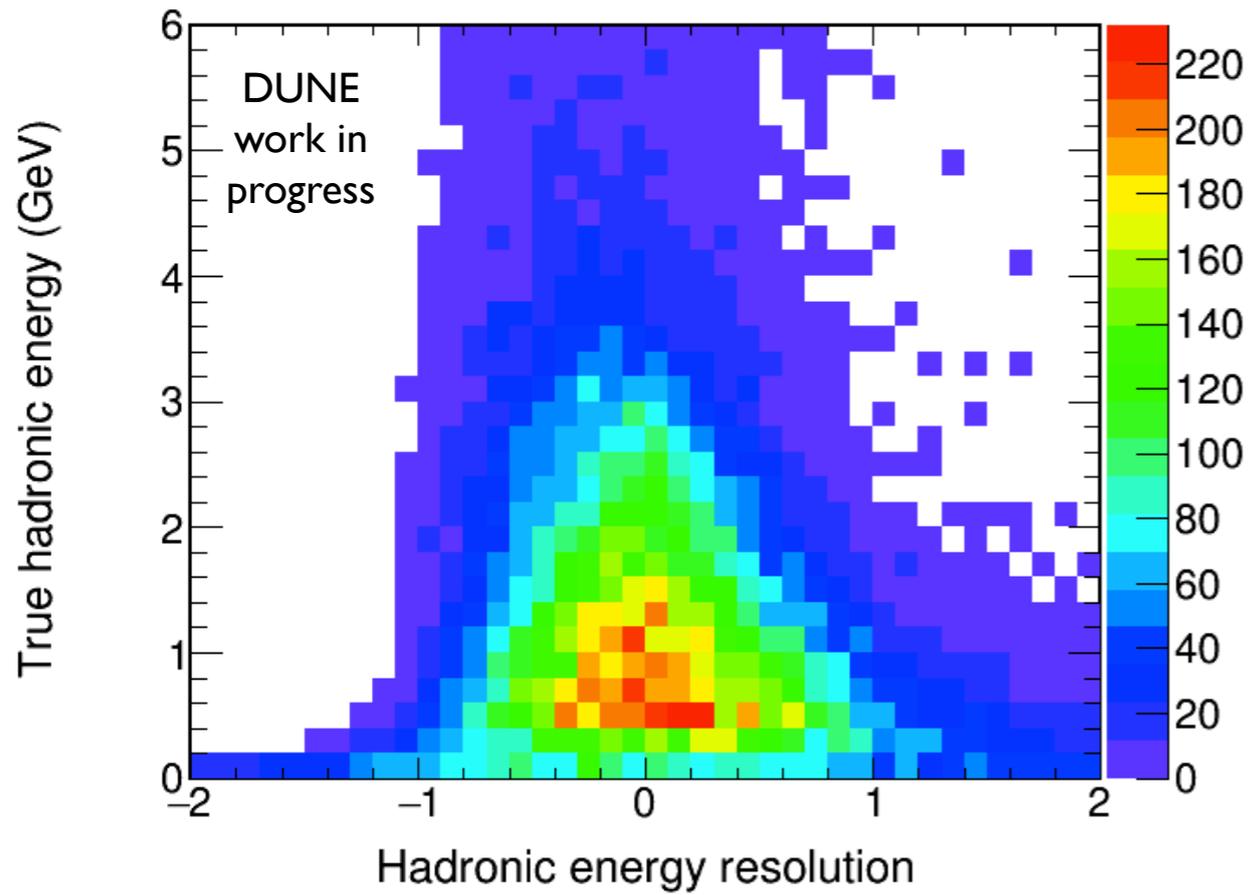


CC DIS



True  $\nu_\mu$  CC events  
with exiting track

## True $\nu_e$ CC events





## Multiple Coulomb scattering



This description is taken from an ICARUS paper at

<https://arxiv.org/pdf/hep-ex/0606006v1.pdf>

A charged particle traversing a medium is deflected through many small angle scatterings.

Divide the track into segments of equal length. Fit a straight line to each track segment.

For each consecutive pair of segments, the scattering angle is the difference between the segment angles.

Compute the RMS of the scattering angle distribution

Repeat for different segment lengths, and fit for  $p$  and  $C$  in

$$\theta_{meas}^{rms} = \sqrt{\left( \frac{13.6 \text{ MeV}}{\beta c p} z \sqrt{\frac{l}{X_0}} \cdot \left[ 1 + 0.038 \cdot \ln \left( \frac{l}{X_0} \right) \right] \right)^2 + (C \cdot l^{-3/2})^2}$$

where  $\beta$  is velocity,  $p$  momentum and  $z$  charge of particle  
 $X_0$  is radiation length,  $l$  is segment length,  $C l^{-3/2}$  is noise

One advantage of this method is that it can be used equally well for stopping and exiting particles.



## Hit charge correction



Lifetime correction to correct for recombination with ions in “bulk” of liquid argon, i.e. not in a column or box around the charged particle track being considered (but they can be from another track).

Ionisation charge = charge collected  $\times$   $\exp(\text{drift time} / \text{lifetime})$ .