

```
In[16]= mp = 938. × 106; qe = 1.6 × 10-19; Co = 2.997 × 108; ω = 2 Pi 805 × 106;
Eo = 10 × 106; d = .104; epso = 8.85 × 10-12; "MKS units everywhere";
```

## Heavy charged particles drift in RF field depending on initial phase

Consider transition of non relativistic particles across a cavity, ie a charged droplet of copper. The induced dipole moment from the electric field does not contribute to a net force as its force = 0 for a uniform field. However if the droplet has a net charge  $q_d$  and mass  $m_d$  then there can be a net drift, depending on the phase of when the droplet is launched. The motion is calculated below. We normalize the electric field to 10 MV/m and a cavity gap of  $d = 0.104$  m. The phase when the particle starts is  $\phi$ .

$$dz^2/dt^2 = q_d E_o / m_d \cos[\omega t + \phi]$$

$$\text{set } k = q_d / m_d E_o$$

Set  $k_p = q_e / m_p E_o$  for a proton

$k_p = C_o^2 E_o / m_p / \omega$   $k_p$  is in m/sec. We will come back later for charged droplets. For example a single charged Cu atom would have  $k_{cu} = k_p / 64$  and its scaled velocity would be 1/64 that for a proton.

```
In[3]= kp = Co2 Eo / mp / ω
```

```
Out[3]= 189 319.
```

```
In[4]= v[t_, φ_, kk_] := kk (Sin[ω t + φ] - Sin[φ]); "kk is kp for a proton";
```

```
In[17]= z[tt_, phi_, kkk_] := Integrate[v[t, phi, kkk], {t, 0, tt}];
```

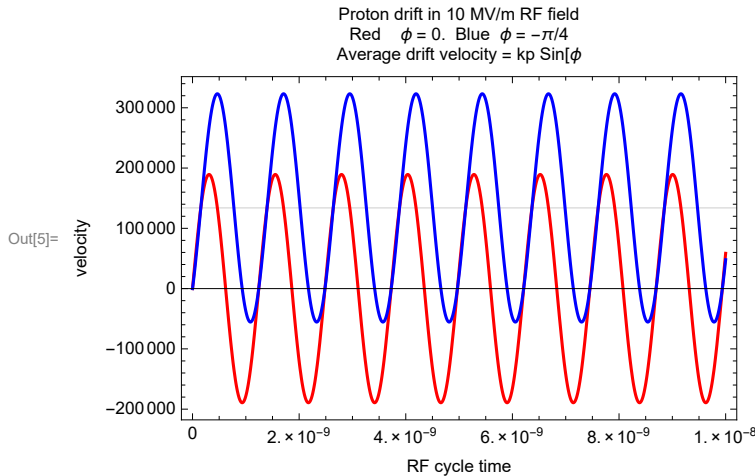
```
In[18]= z[tt, Pi / 4, kp]
```

```
Out[18]= 0.000026467 - 133 869. tt - 0.0000374299 Cos[ $\frac{\pi}{4} + 1 610 000 000 \pi tt$ ]
```

```

In[5]= Plot[{v[t, 0, kp], v[t, -Pi / 4, kp]}, {t, 0, 10^-8},
  PlotStyle -> {Red, Blue}, GridLines -> {None, {kp Sin[pi / 4]}},
  Frame -> True, FrameLabel -> {"RF cycle time", "velocity"},
  "Proton drift in 10 MV/m RF field\nRed    phi = 0. Blue
    phi = -pi/4\nAverage drift velocity = kp Sin[phi", ""}]

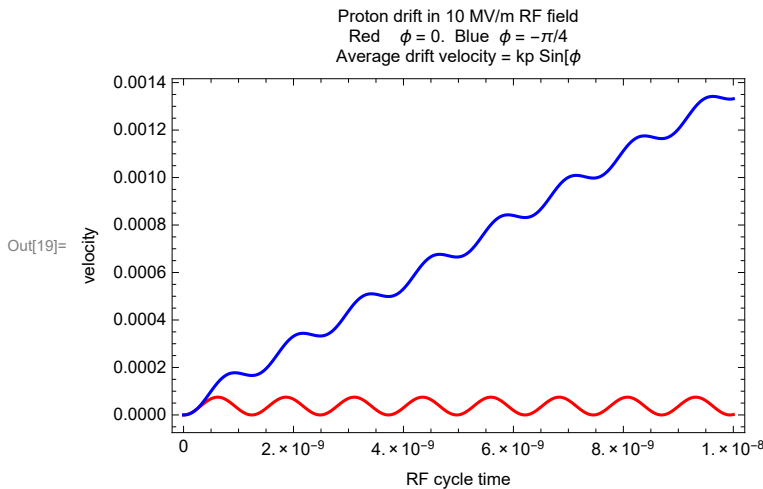
```



```

In[19]= Plot[{z[tt, 0, kp], z[tt, -Pi / 4, kp]}, {tt, 0, 10^-8},
  PlotStyle -> {Red, Blue}, GridLines -> {None, {kp Sin[pi / 4]}},
  Frame -> True, FrameLabel -> {"RF cycle time", "velocity"},
  "Proton drift in 10 MV/m RF field\nRed    phi = 0. Blue
    phi = -pi/4\nAverage drift velocity = kp Sin[phi", ""}]

```



From the above we see the average velocity of a proton is  $k \sin[\phi] = 189319 \sin[\phi]$  meters/sec or 0.23 mm / rf cycle. It would take a proton at max  $\sin[\phi]$  400 cycles to drift across the cavity. Its peak energy is about  $1/2 mp c^2 \beta^2$ .  $\beta^2$  is about  $10^{-6}$  (see above plot). The case for a droplet is modified as its charge/mass ratio will be much smaller for two reasons: the mass is higher by 64 x Number of copper atoms in the droplet and the charge is much smaller than the number of copper atoms.

## Droplets

Lets consider a droplet.. Lets estimate the charge possible on a droplet. A sphere in a uniform field has an induced dipole moment and the max field on its surface is  $3 E_0$  at the pole. he droplet has an induced dipole, the net charge is zero. Suppose we place an additional charge of electrons on the droplet so that the surface charge is negative everywhere. Since this will be uniform, the net number of electrons will be  $4 \pi a^2 \epsilon_0 (3 E_0) / q_e$ . The number of copper atoms is  $4/3 \pi a^3 / (4/3 \pi 1.25 \cdot 10^{-10})^3$  where  $a$  is the radius of the droplet and  $1.25 \cdot 10^{-10}$  is the radius of a copper atom in meters. So can compute  $q_d / m_d$  in terms of the  $q_e / m_p$  by calculating the number of electrons and number of copper atoms.

So the  $q_d / m_d = 4 \pi a^2 (3 E_0 \epsilon_0 / q_e) / (a / 1.25 \cdot 10^{-10})^3 q_e / m_p$

$$\text{fac}[a\_]:= 4 \pi a^2 (3 E_0 \epsilon_0 / q_e) / (a / (1.25 \times 10^{-10}))^3$$

$\text{fac}[a]$  is used to multiply  $k_p$  (the drift velocity of a proton) and obtain the droplet velocity as a fraction of what the proton drift velocity would be. Note that the surface charge we have used is already very high! and I would expect the charge would be much lower. In any case the above charge means the field at every spot on the droplet is repulsive.

```
In[6]:= fac[a_] := 4 π a^2 (3 E0 eps0 / qe) / (a / (1.25 × 10^-10))^3
```

```
In[7]:= fac[10^-5] kp
```

```
Out[7]= 0.000771044
```

Note:  $a$  is in meters.. So we see the velocity of drift for a  $10 \mu$  radius droplet is 0.7 mm/sec. It takes over 100 sec to cross cavity!

## Time constant for sphere to solidify with radiation cooling

```
In[8]:= Hcu = 209; ρ = 8.9; ε = 0.2; σ = 5.67 × 10^-4; Tm = 1358;
```

Units above mks except length in cm, weight in grms. We calculate how many joules are stored in the latent heat of melting. If the droplet is just starting to solidify, then this heat must be removed before it becomes solid. If the temperature is above 1358 K then it must be first be cooled to this temperature. The cooling is by radiation. The emittance of molten copper is about 0.2. We calculate the cooling in two stages:  $\tau_c$  From the initial temperature  $T_i$  to 1358 and then  $\tau$  the time to get the latent heat of melting out.

$$\tau = 1/3 a \rho H_{cu} / (\epsilon \sigma (T_m^4 - 300^4)).$$

Above the melting point, the specific heat of copper is constant and equal 0.625 joules/grm up to about 2000 K. So for instance the additional heat from from 1358 to 2000 is about 400 joutes. Integrating, we get

$$\tau_c[a_, T_i_] := -3 / 5 a \rho .625 (T_i^{-3} - 1358^{-3}) / (\epsilon \sigma)$$

The total time to solidify is the sum of these two times.

Note: a is in cm! sorry

```
In[9]:=  $\tau[a_] := 1/3 a \rho H_{cu} / (\epsilon \sigma (T_m^4 - 300^4))$ 
```

```
In[10]:=  $\tau[a]$ 
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```
Out[10]=  $1.61153 \times 10^{-6} a$ 
```

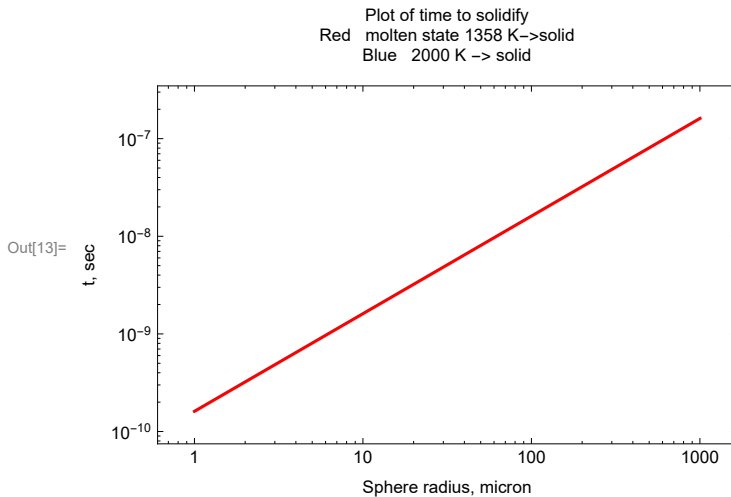
```
In[11]:=  $4/3 \text{Pi} a^3 \rho H_{cu}$  "Joules stored in latent heat of fusion times"
```

```
Out[11]= 7791.57 Joules stored in latent heat of fusion times a3
```

```
In[12]:=  $4 \text{Pi} a^2 \epsilon \sigma (T_m^4 - 300^4)$  "Joules Radiated/sec times"
```

```
Out[12]=  $4.83489 \times 10^9$  Joules Radiated/sec times a2
```

```
In[13]:= LogLogPlot[{ $\tau[a 10^{-4}]$ ,  $\tau[a 10^{-4}] + \tau_c[10^{-4} a, 2000]$ },
  {a, 1, 103}, PlotStyle -> {Red, Blue}, Frame -> True,
  FrameLabel -> {"Sphere radius, micron", "t, sec", "Plot of time to solidify\nRed
  molten state 1358 K->solid\nBlue 2000 K -> solid", ""}]
```



```
In[14]:=  $\tau_c[a_, T_i_] := -a \rho .625 (T_i^{-3} - 1358^{-3}) / (9 \epsilon \sigma)$ 
```

```
In[15]:=  $\tau_c[1, 1678]$ 
```

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Out[15]=  $1.02273 \times 10^{-6}$ 
```

## Conclusion: Mystery

The cavity has a gap of 0.1 meter. A 100 $\mu$  drop has a cooling time constant of few times 10<sup>-8</sup> sec. To get across the cavity before solidifying its velocity would have to be 10<sup>7</sup> m/sec!!!!

So droplet must come from same side. If so it must be nearby. The splash marks are remarkably symmetrical...doesn't look like a side wise hit.

There isn't molten copper till late in the discharge. How does one eject a charged droplet? I would guess that an expert could estimate how fast the droplet hit from the splash mark. That would give us

an estimate of the charge.

Notice the little tiny spheres around the splash. I would suggest that they are so small that their cooling time is  $< ns$  and they are solid when they hit.