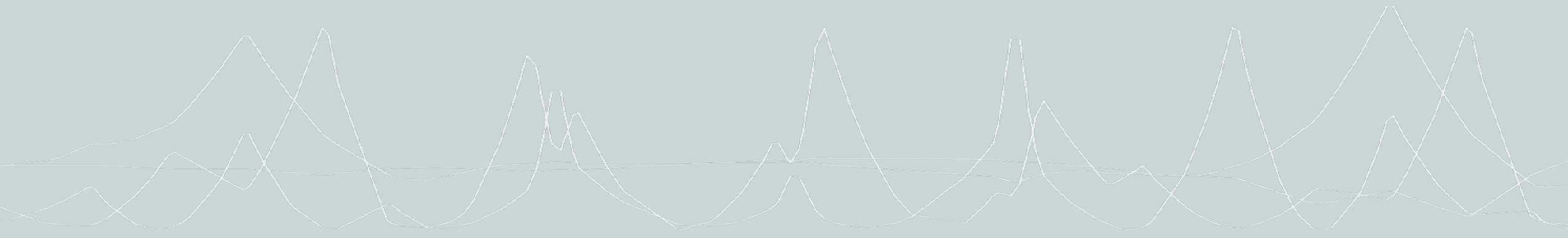


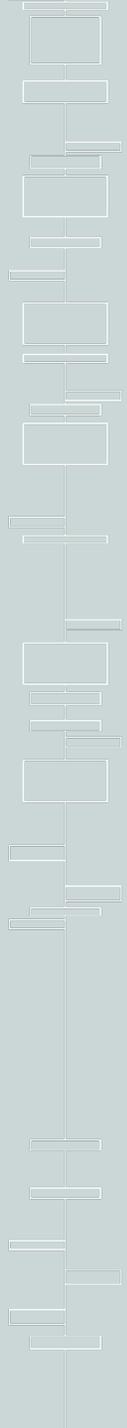


Investigation of Emittance Analysis Methods Using the MuCool Test Area

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Supervisors: Dr. Carol Johnstone & Adam Watts





Outline

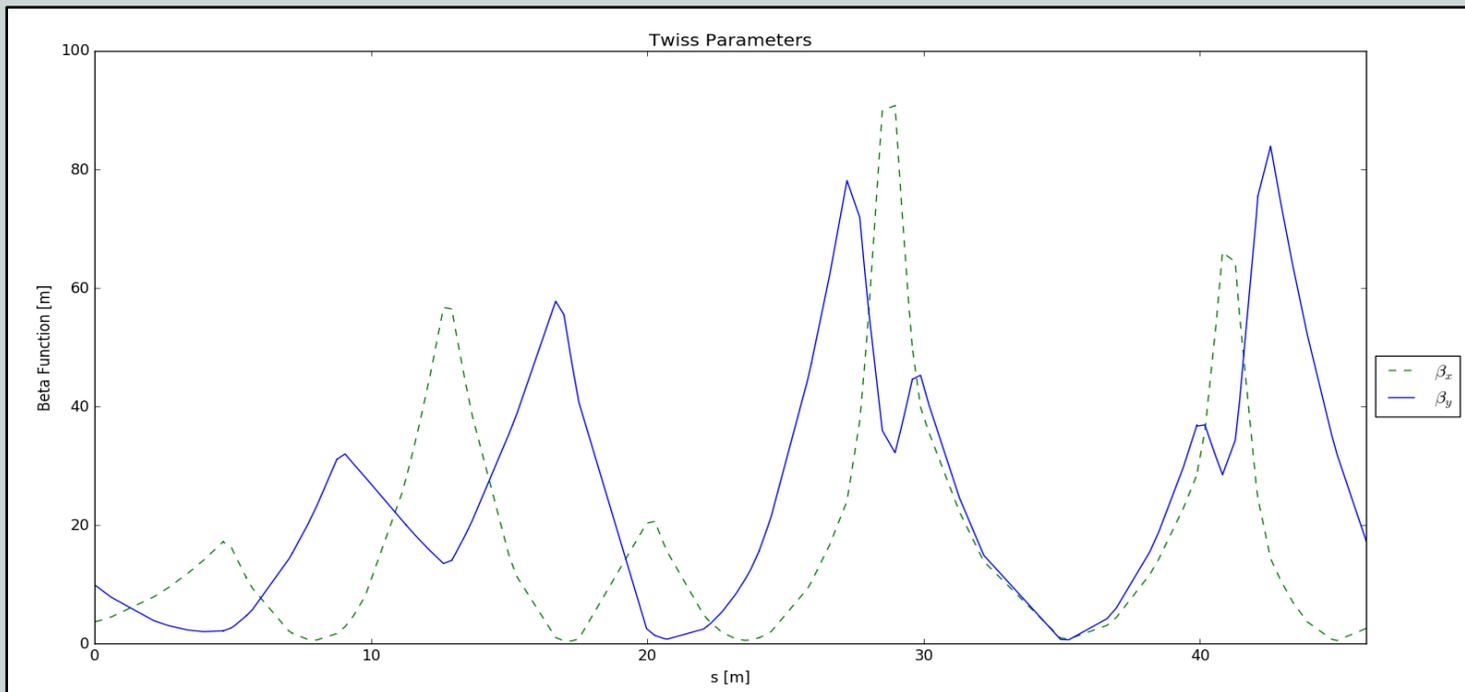
- I. Introduction to Beamline Physics
 - I. Emittance
- II. MTA Beamline
- III. Multiwires
 - I. Instrumentation
 - II. Data Analysis
- IV. Emittance Calculation
 - I. MAD Simulation
 - II. 3-Wire Method
- V. Conclusions & Future Directions

Purpose

- ⇒ Use the MuCool Test Area (MTA) beamline as a diagnostic line to measure a parameter known as emittance coming out of Linac
- ⇒ Knowledge of the transverse properties of the Linac beam can help improve Booster injection efficiency

Courant-Snyder Parameters

- ⇒ Model the beam and understand how the envelope propagates
- ⇒ Utilize transfer matrices derived from magnet optics to calculate beam trajectory and properties



$$x(s) = A\sqrt{\beta_x(s)}\cos(\psi_x(s) + \phi)$$

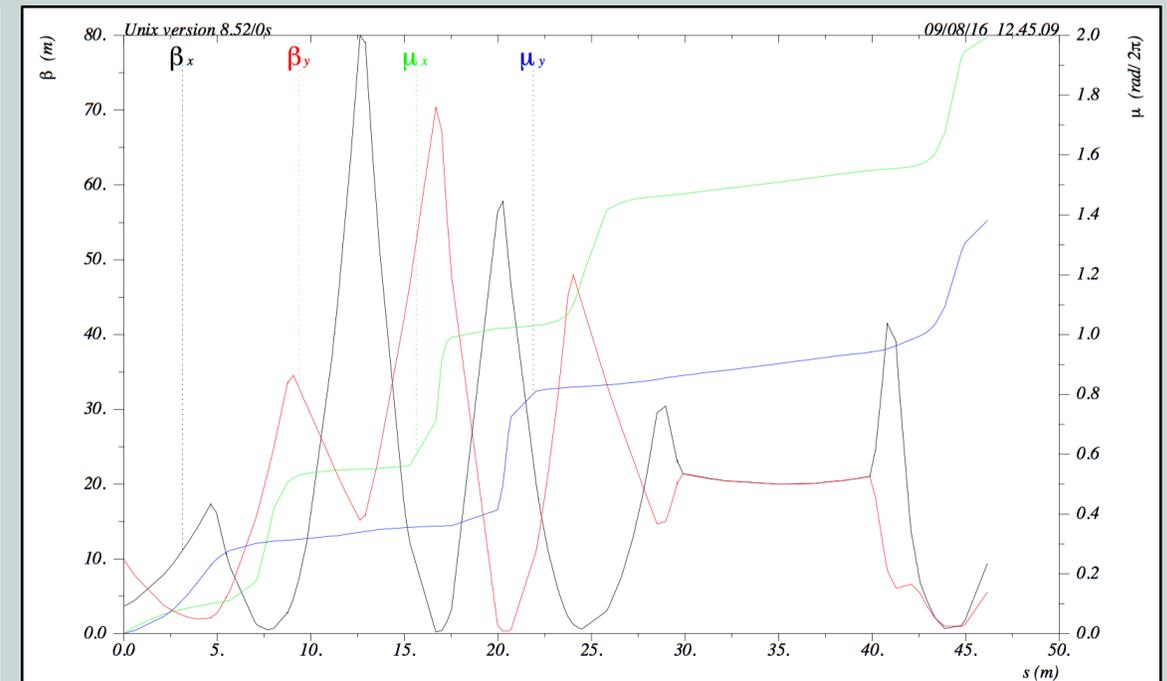
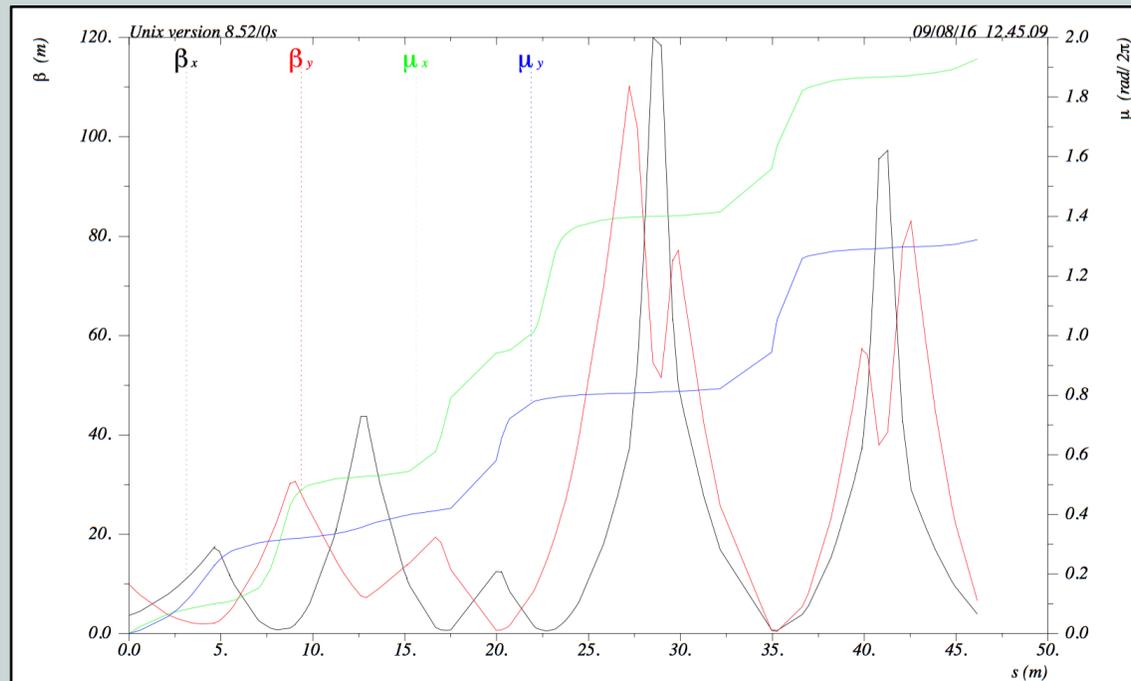
$$\alpha(s) \stackrel{\text{def}}{=} -\frac{1}{2} \frac{d\beta(s)}{ds}$$

$$\gamma(s) \stackrel{\text{def}}{=} \frac{1 + \alpha(s)^2}{\beta(s)}$$

Phase Advance

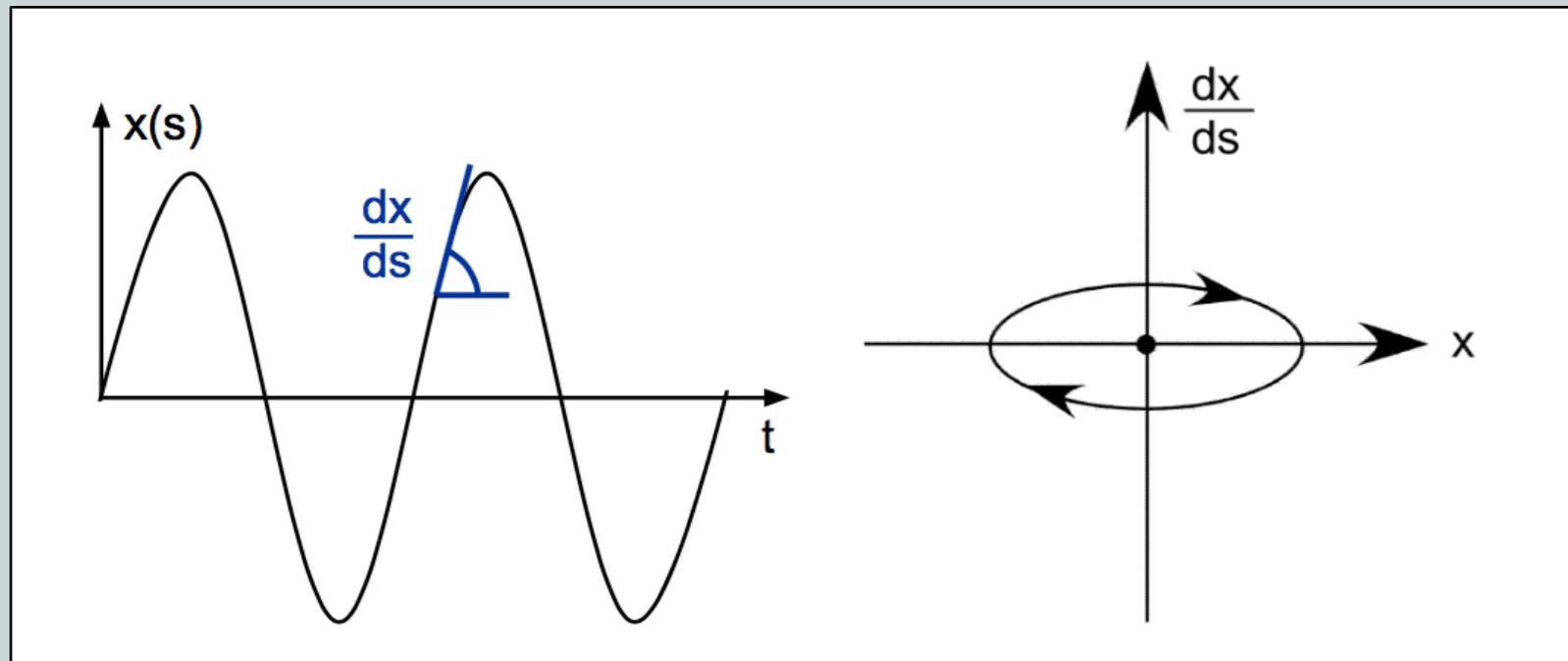
⇒ Amount of oscillation period a particle has gone through in a certain arc – will be used to parameterize data sets

$$\psi_{s_1 s_2} = \int_{s_1}^{s_2} \frac{ds}{\beta(s)}$$



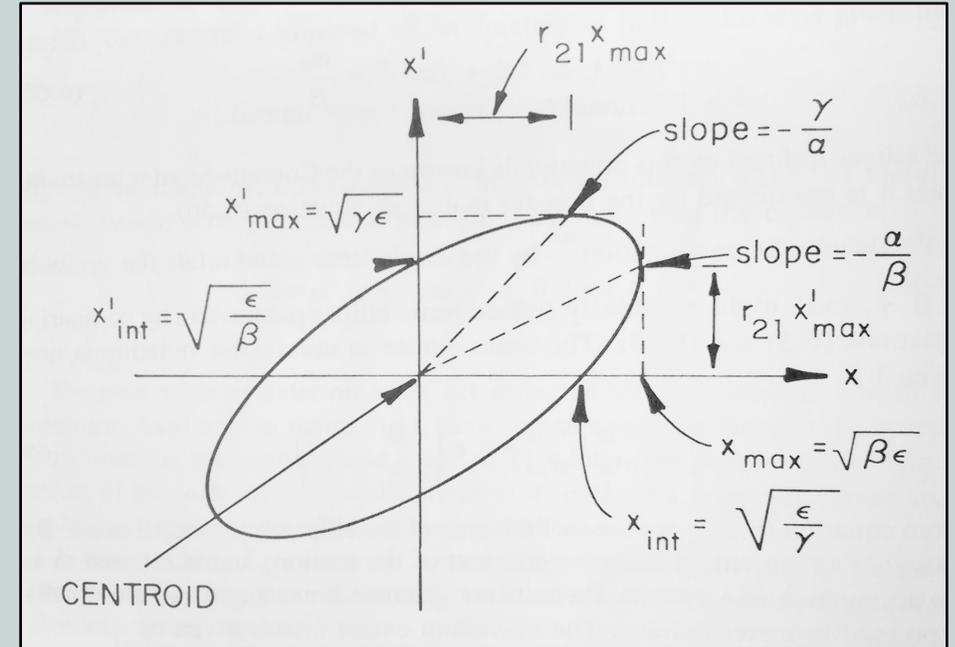
Phase Space

- ⇒ Betatron oscillation as a function of distance along the beam path
- ⇒ Parameter space that is used for emittance calculations



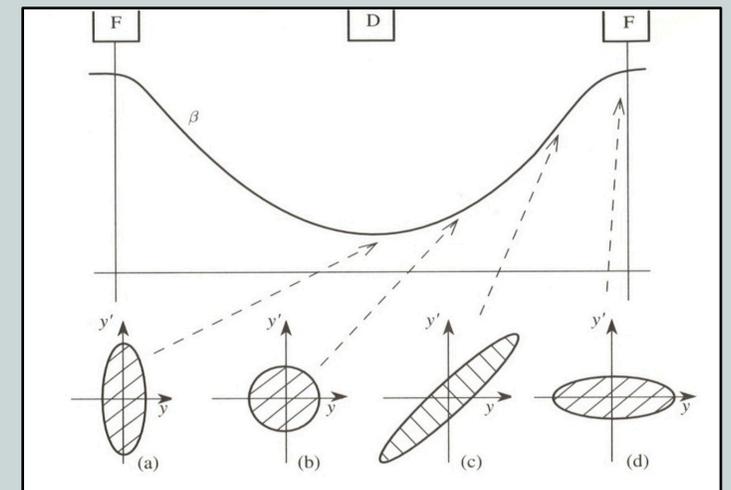
Emittance

- ⇒ Envelope of particles in phase space is elliptical when beam is exposed to linear forces
- ⇒ Emittance is equal to the area of this envelope – invariant quantity
- ⇒ Important metric for understanding beam properties & optics models
- ⇒ Need to have a method of calculating the emittance from beam position measurements



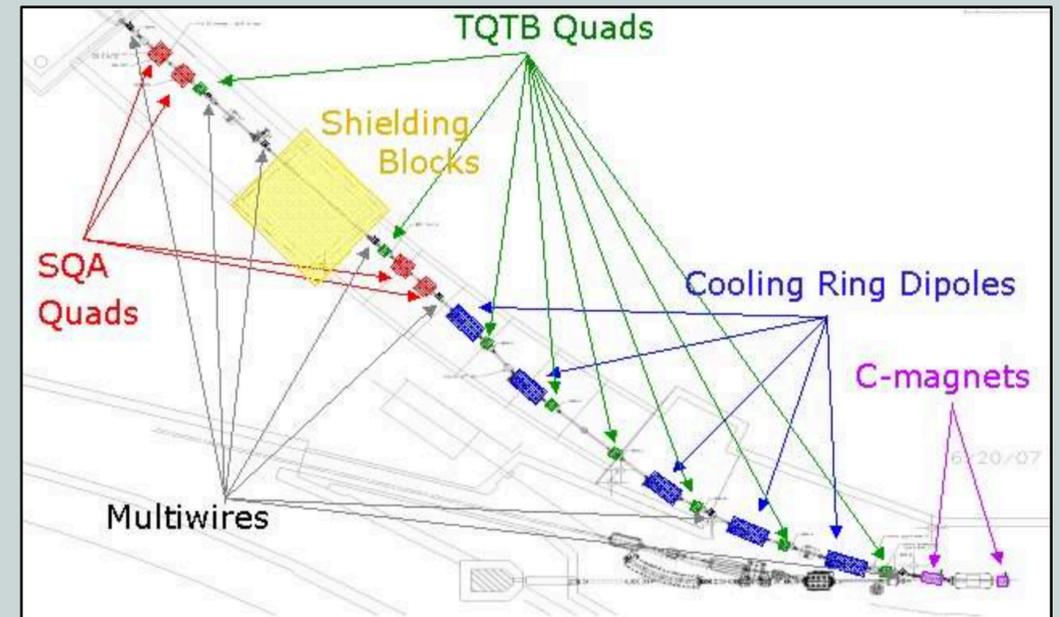
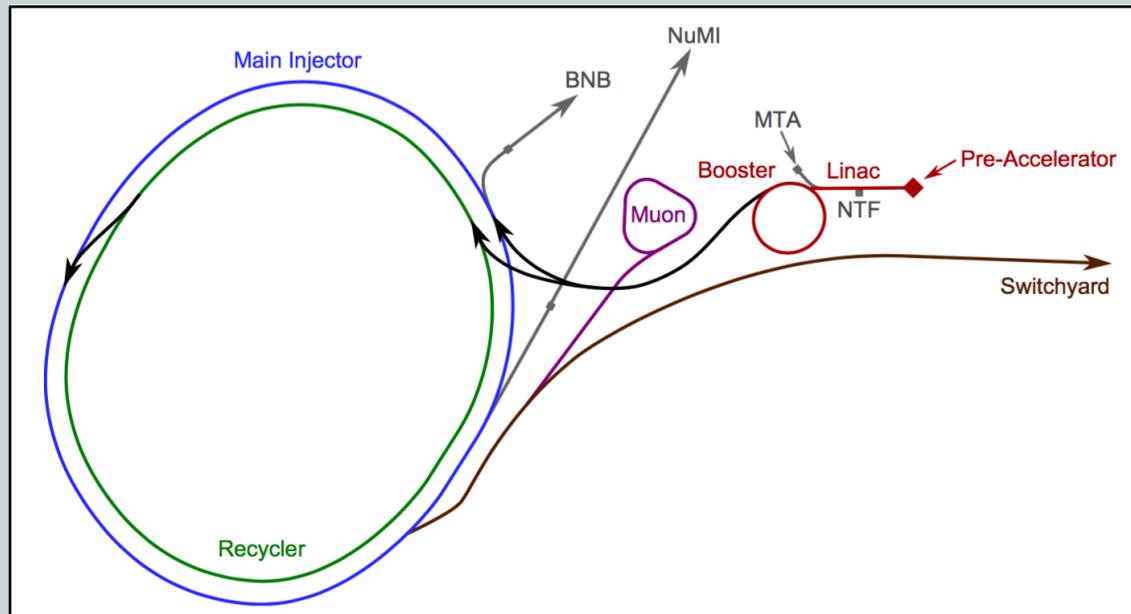
$$\frac{\epsilon}{\pi} = \gamma x^2 + 2\alpha x x' + \beta x'^2$$

$$\epsilon = \frac{\sigma^2}{\beta}$$



Muon Test Area (MTA)

- ⇒ 400 MeV line off of Fermi's linear accelerator
- ⇒ Useful for beam diagnostics and testing
- ⇒ Data used for beam profile reconstruction



Secondary Emission Multiwires

- ⇒ 48 evenly spaced parallel wires in both planes
- ⇒ Protons strike the wires and liberate electrons, producing a net charge proportional to the beam intensity
- ⇒ 0.5, 1, or 2mm spacing
- ⇒ Profile monitor can be rotated in and out of beam to minimize losses



Analyzing Multiwire Data

⇒ Gaussian Curve Fit

⇒ *Non-linear least squares method done using Python to determine function parameters*

⇒ χ^2 goodness of fit test

$$g(x) = \frac{A}{\sigma\sqrt{2\pi}} e^{-\frac{1}{2}\left(\frac{x-\mu}{\sigma}\right)^2}$$

⇒ Root Mean Square (RMS) Method

⇒ *Highly sensitive to tails*

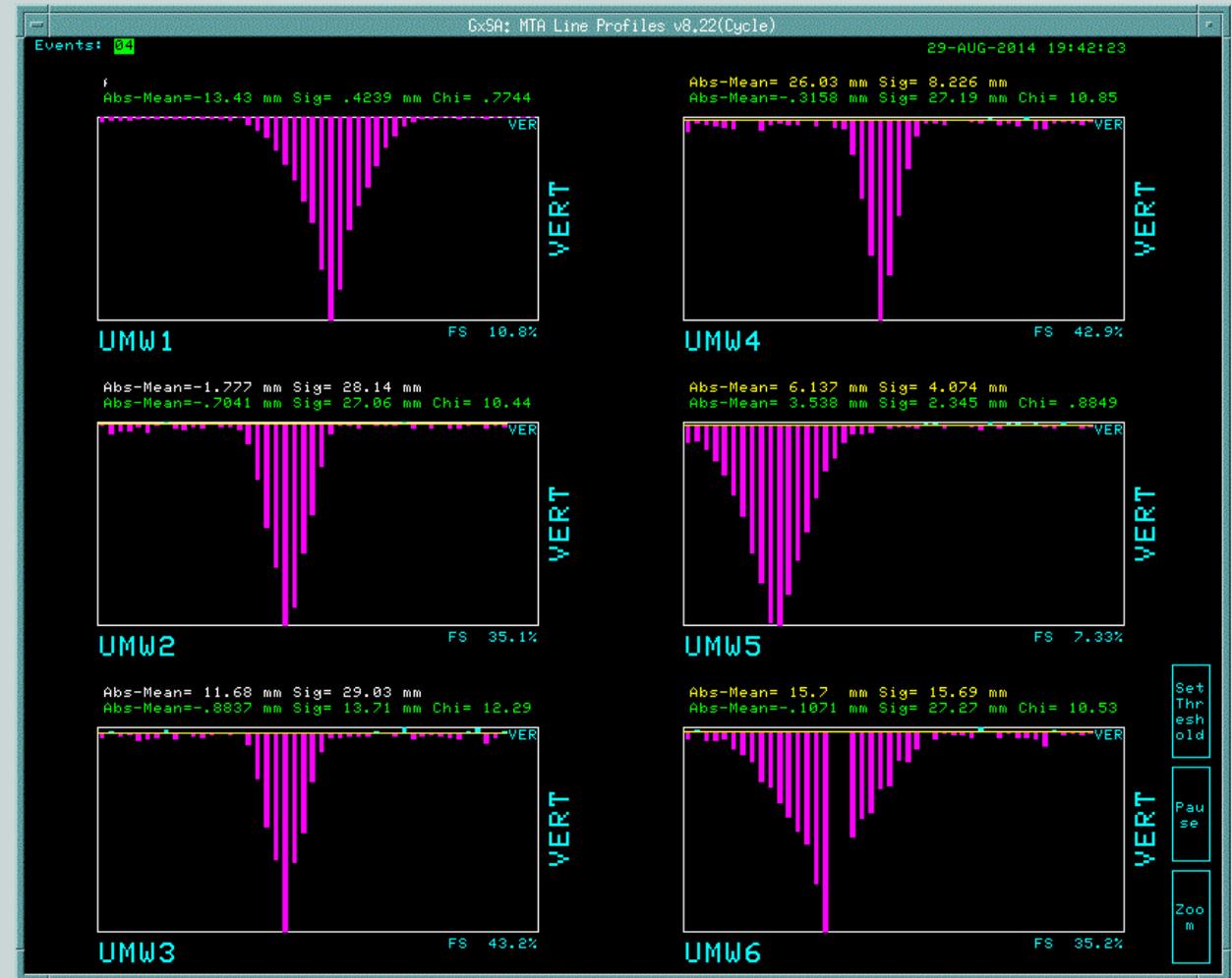
⇒ *Where x represents the wire position and $P(n)$ is the signal at that wire*

$$\mu = \frac{\sum_n n |P(n)|}{\sum_n |P(n)|}$$

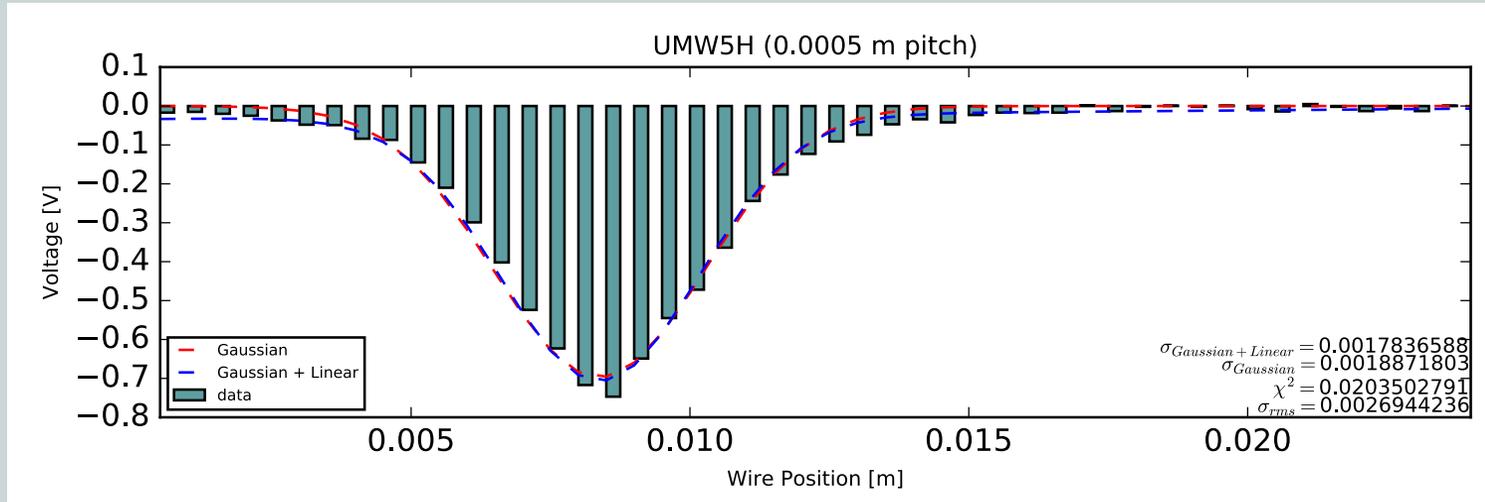
$$\sigma_{rms}^2 = \frac{\sum_n (x - \mu)^2 |P(n)|}{\sum_n |P(n)|}$$

Secondary Emission Multiwires

- ⇒ Assume Gaussian distribution, compute beam radius in both planes & compare to RMS
- ⇒ Specifically looking at 10 m drift in MTA (wires 4, 5, & 6 with 2mm, 0.5mm, & 1mm pitches respectively)
- ⇒ August 2014: 5 data sets taken at each multiwire for all of the 5 phase advances, also referred to as the tune



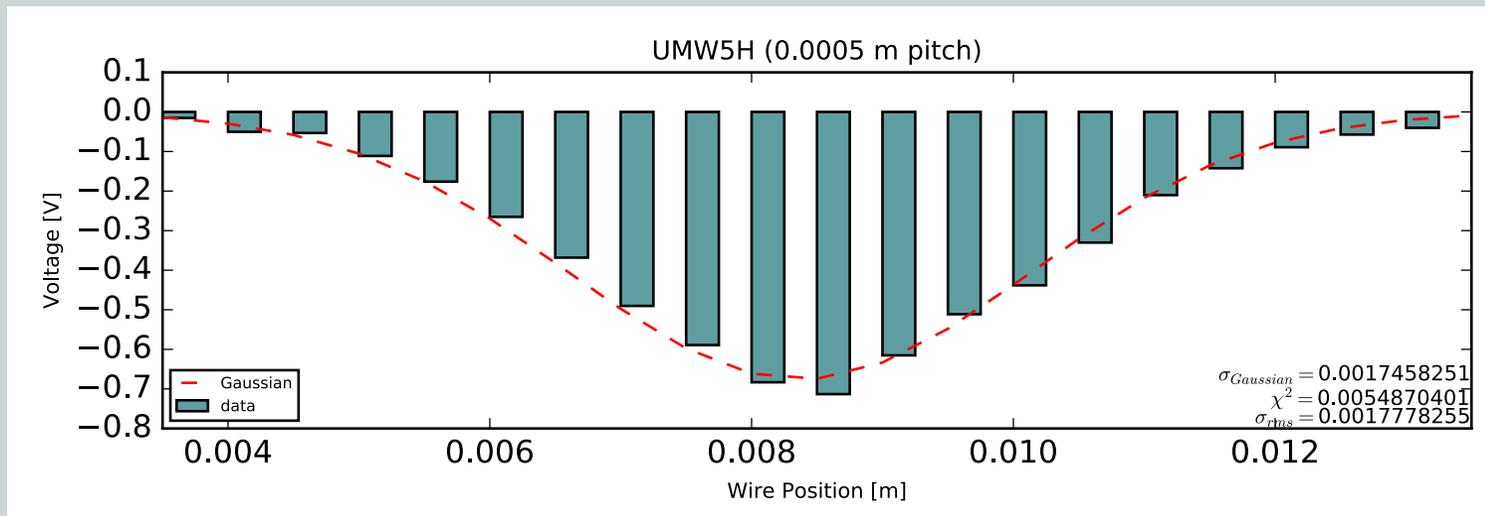
Cleaning up Data



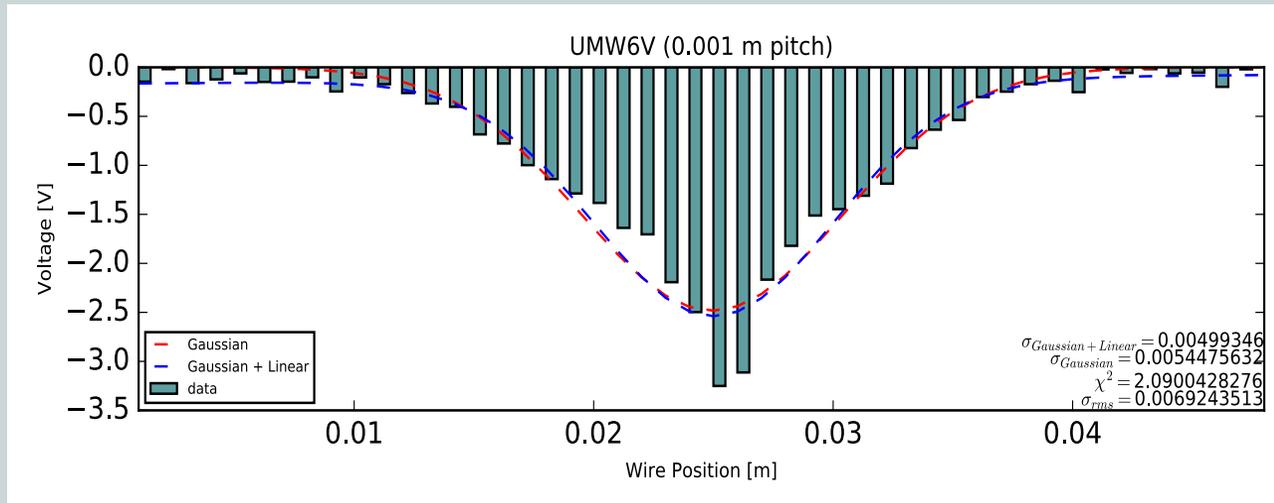
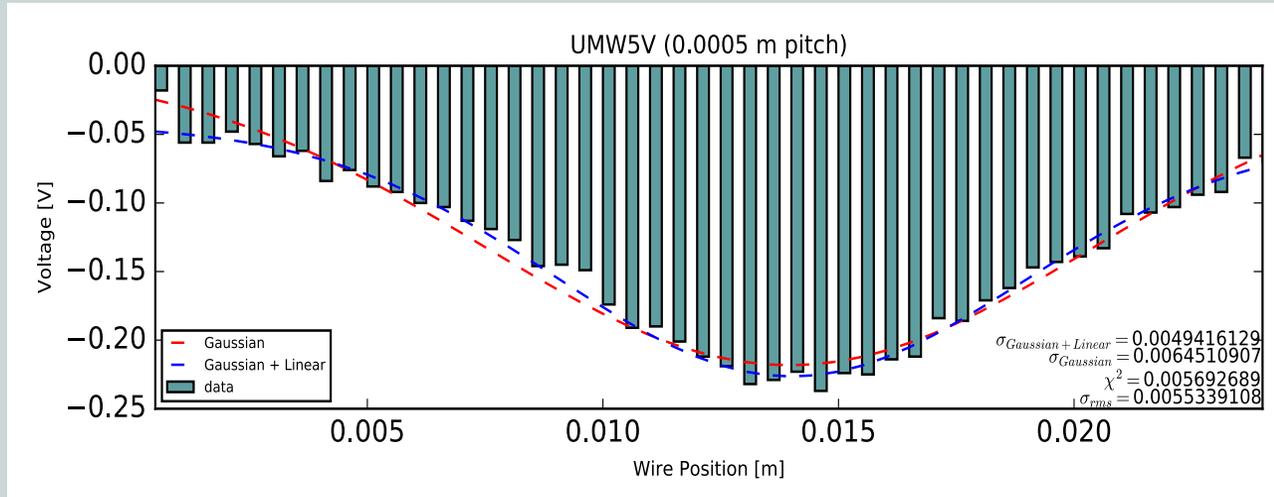
⇒ Gaussian distribution + linear background component fit

⇒ Linear intercept subtracted as background

⇒ Gaussian tails outside $3\sigma/4\sigma$ from peak cut for more accurate RMS calculations



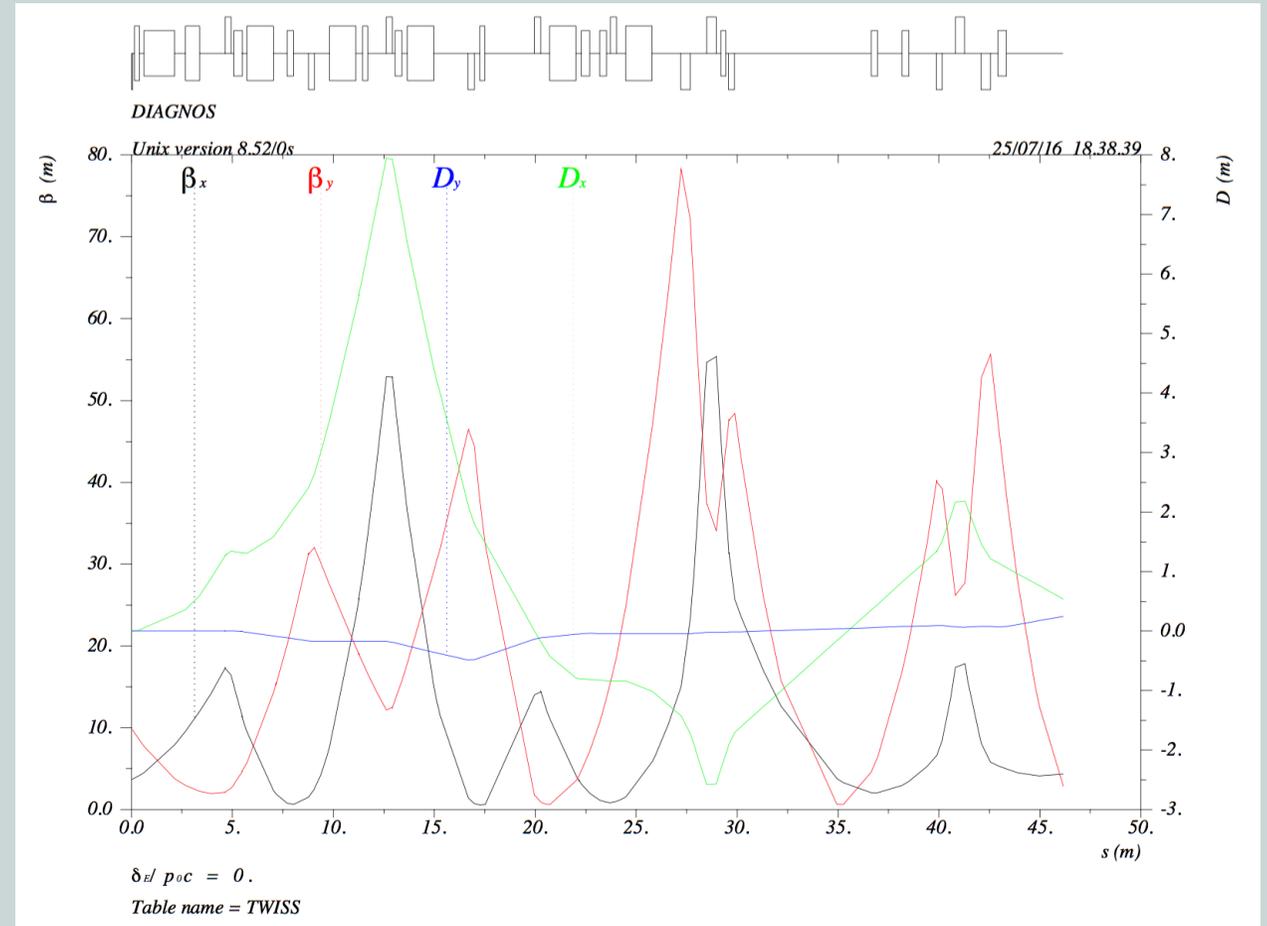
Issues with Gaussian Fit and RMS Calculation



- ⇒ Beam scraping leads to fractured beam profiles
- ⇒ Distribution is not always well defined with a Gaussian function
- ⇒ RMS method highly sensitive to cutoff point because of tails

Methodical Accelerator Design (MAD)

- ⇒ Simulation program developed at CERN
- ⇒ Models the beam using Courant-Snyder parameters
- ⇒ Input lattice of all magnets & respective strengths at each tune



Current Conversion & Simulations

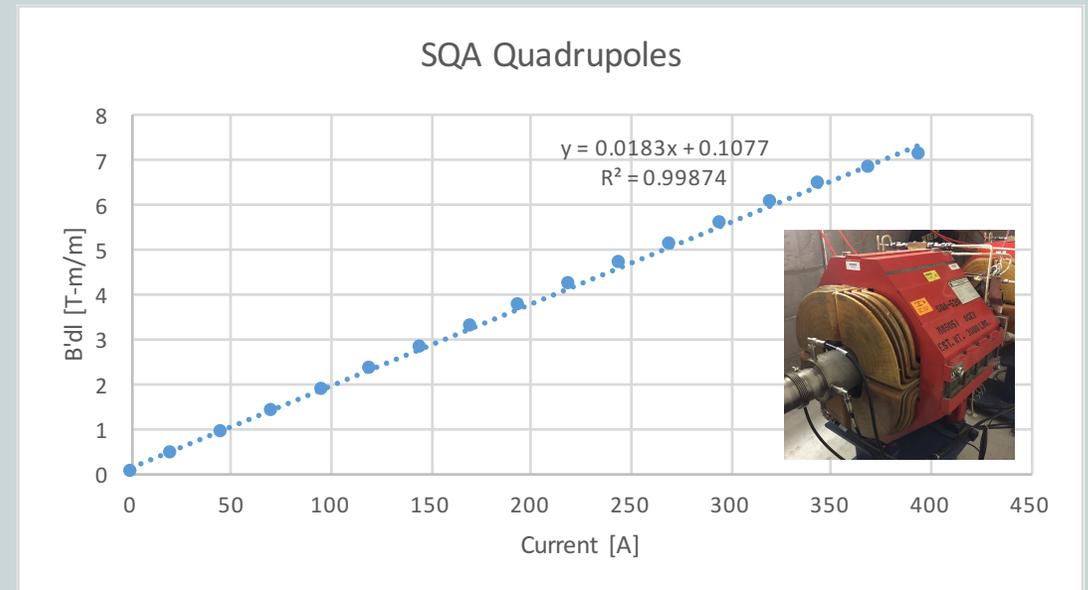
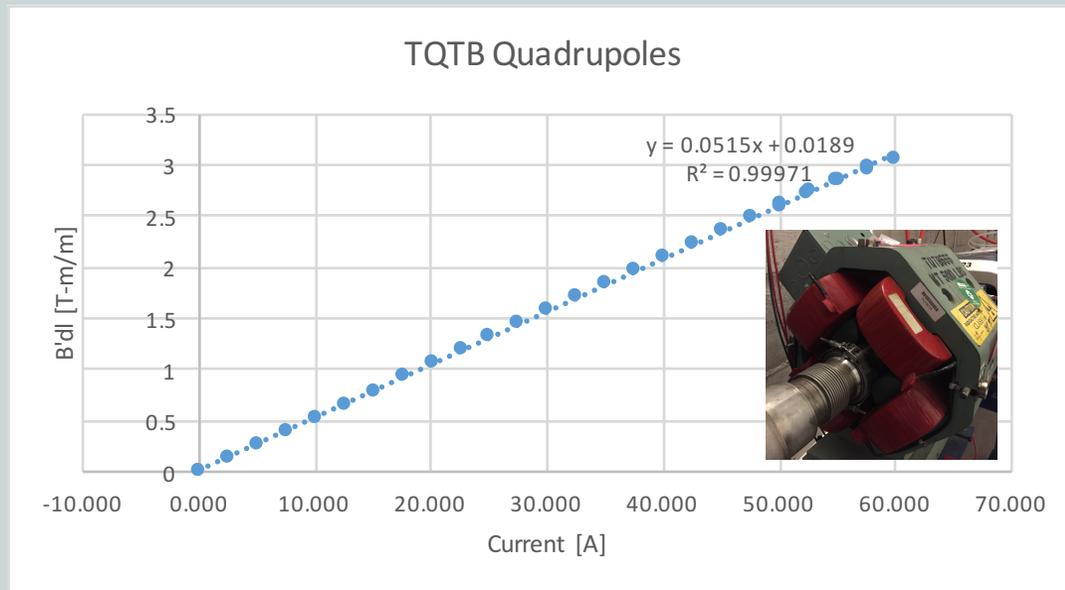
- ⇒ Use currents from operator log to determine quadrupole k-values
- ⇒ Run simulation with different lattices based on calculations

```

!MAJOR BENDS AND QUAD POWER SUPPLIES-----
-E:UHB07  L BEND DS Q06      426   424   420.1  Amps ...
-E:UQ07   SQA Quad UQ07 Curre 21.54  26.34  26.08  Amps ...
-E:UQ08   SQA Quad UQ08 Curre  88     111    109.6  Amps ...
-E:UQ11   SQA Quad UQ11 Current 141     * 140  Amps ...
-E:UQ12   SQA Quad UQ12 Current  89.7   * 88.34 Amps ...
!QUAD TRIMS----- (50 AMP CORRECTORS)-----
-E:UQ01   Current Monitor  26.5   28.7   28.67  A  ..
-E:UQ02   Current Monitor  18.72  15.8   15.78  A  ..
-E:UQ03   Current Monitor  27.34  28.1   28.09  A  ..
-E:UQ04   Current Monitor  23.3   24.5   24.47  A  ..
-E:UQ05   Current Monitor  35.1   31.8   31.77  A  ..
-E:UQ06   Current Monitor -23.1  -23.9  -23.89 A  .T
-E:UQ09   Current Monitor  29.86  19.66  19.65  A  ..
    
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$$k = \frac{B'}{B\rho}$$

$$B\rho = 3.1823 \text{ GeV}/c$$



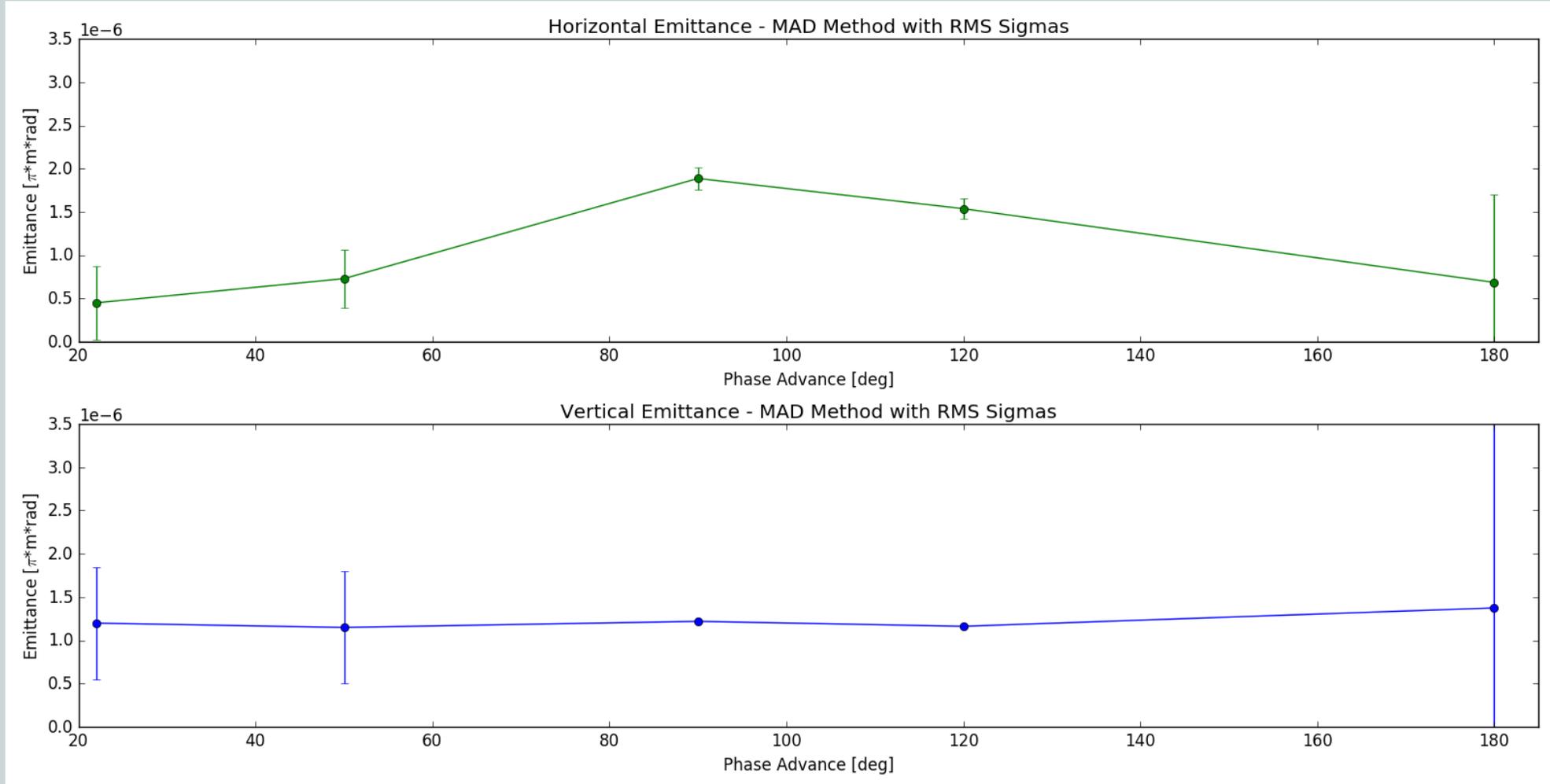
MAD Method

- ⇒ Use calculated quadrupole settings to recreate conditions when beam profiles were recorded
- ⇒ 'Match' method used to fit emittance values
 - ⇒ *Introduce constraint using computed beam width values at each wire*

$$\epsilon = \frac{\sigma^2}{\beta}$$

- ⇒ MAD returns emittance values with penalty function

MAD Method: Results



Issues with MAD Method

- ⇒ This method relies on initial conditions coming out of Linac
 - ⇒ *The measurements we have were taken before recent upgrades*
- ⇒ Have minimal control over convergence of MAD's matching algorithm
 - ⇒ *Our system is underconstrained and the algorithm will not always iterate enough times to find a solution with high confidence values*
- ⇒ Heavily reliant on a lattice that requires correct component spacing data, magnet current values, & consistent fields and strengths from magnet to magnet

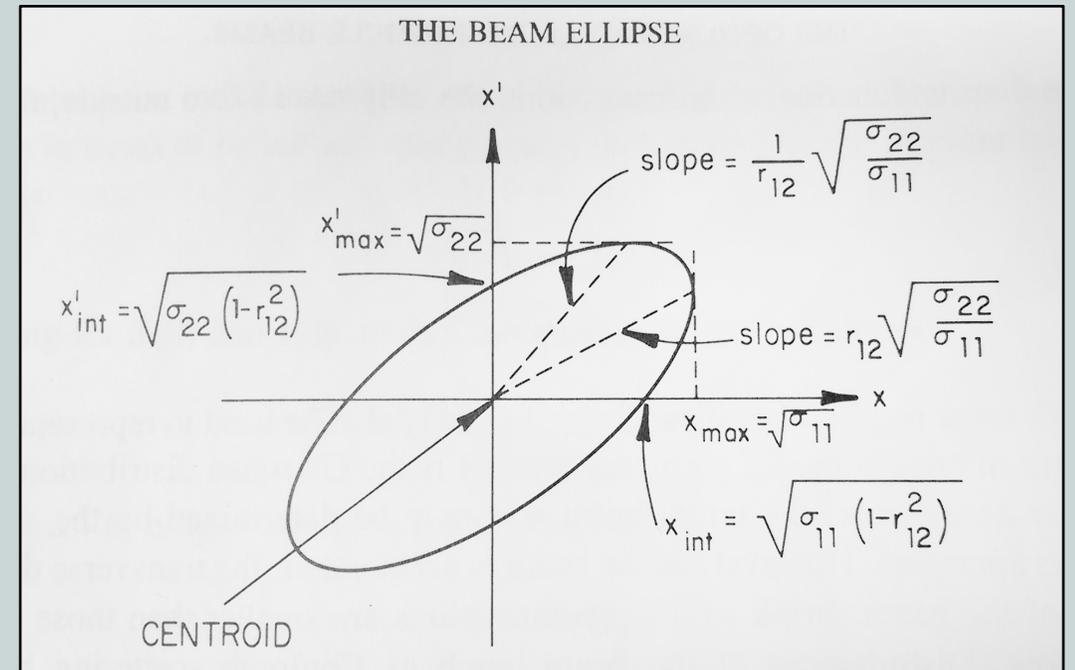
3-Wire Method

- ⇒ Beam width calculated at 3 multiwires in a drift
- ⇒ Solve linear system derived from linear optics transfer matrices

$$R_2^2 = R_1^2 + 2L_1\sigma_{12} + L_1^2\sigma_{22}$$

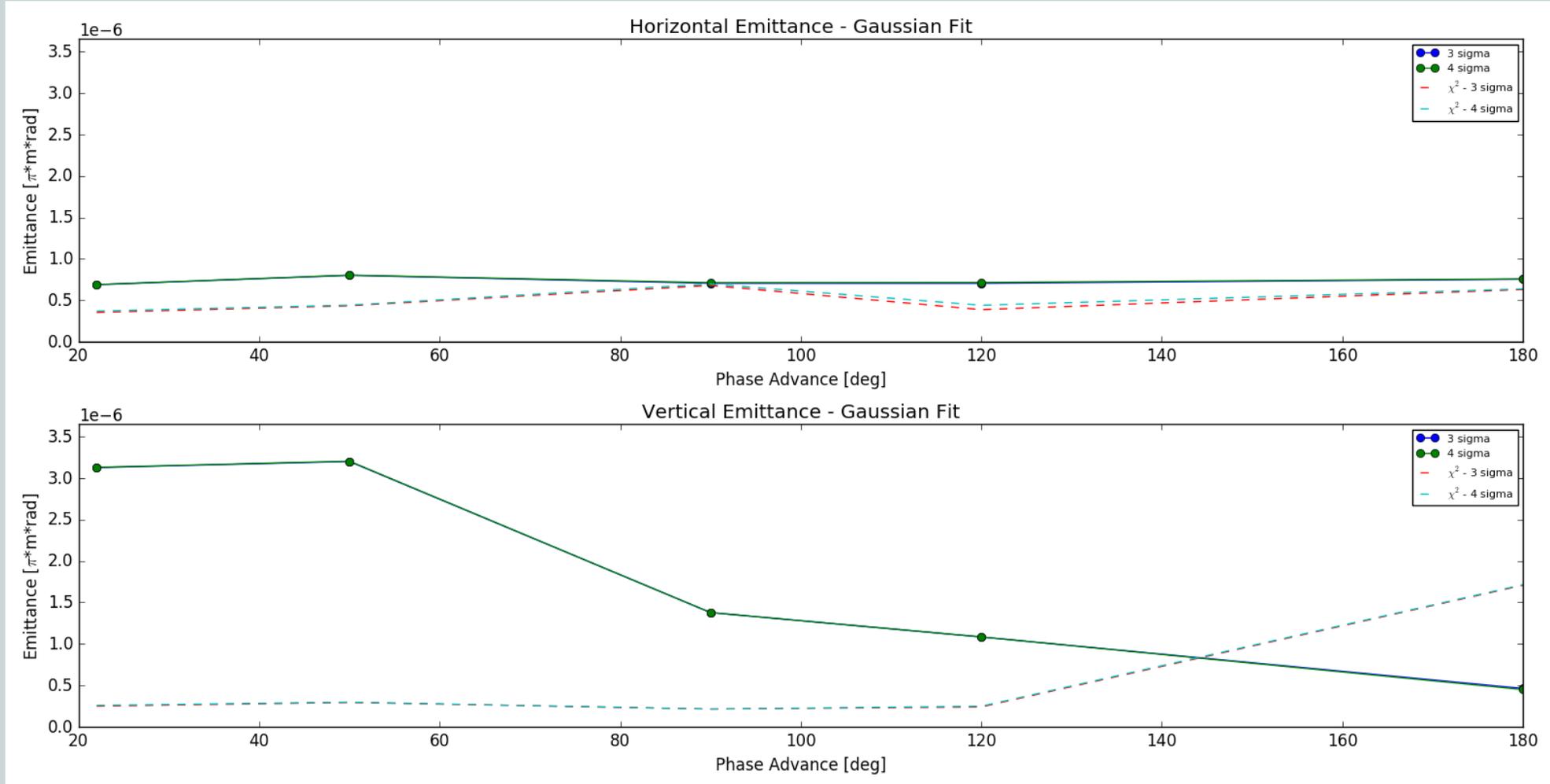
$$R_3^2 = R_1^2 + 2(L_1+L_2)\sigma_{12} + (L_1 + L_2)^2\sigma_{22}$$

where L_1 & L_2 are distances along nominal trajectory and σ values are beam radii ($\sigma = R^2$)

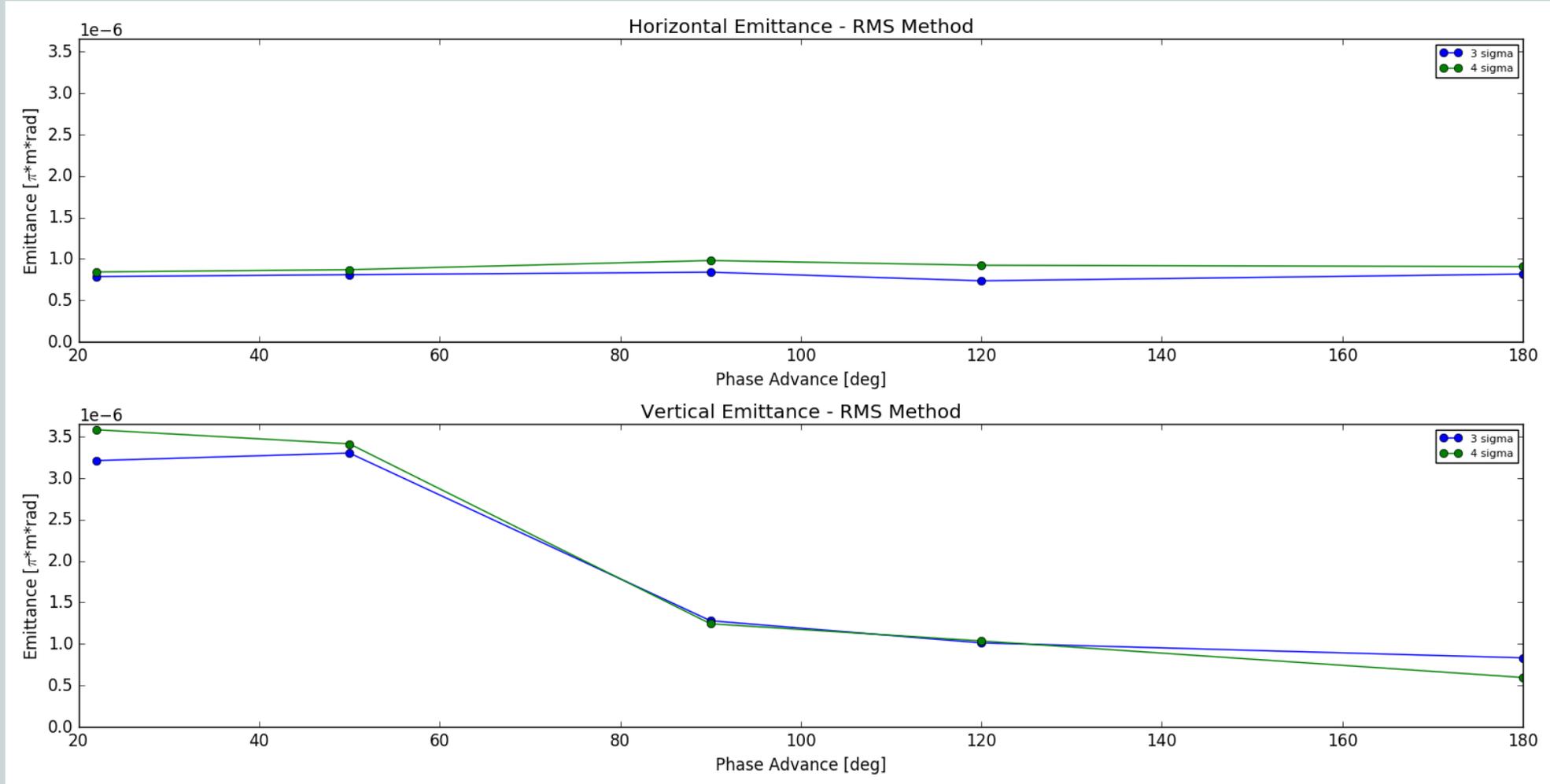


$$\epsilon_{rms} = \sqrt{\sigma_{11}\sigma_{22} - \sigma_{12}^2}$$

3-Wire Method Using Gaussian σ



3-Wire Method Using RMS σ



Conclusions & Future Directions

- ⇒ We appear to see violation of emittance conservation in the vertical but not in the horizontal
- ⇒ Most likely source of this violation is non-elliptical beam
- ⇒ Using tomographic techniques to reconstruct the beam profile and test this assertion
 - ⇒ *Move away from reliance on Courant-Snyder parameters which assume elliptical properties*
- ⇒ Once beam is running again, we can take more data for tomographic reconstruction
- ⇒ Investigation into causes of non-elliptical envelope
- ⇒ Better understanding of beam in transverse phase space will help beam losses into Booster

Closing Remarks

I would like to extend my deepest gratitude to my supervisors and mentors, Dr. Carol Johnstone and Adam Watts. I would also like to thank Michael Backfish and Dr. John Johnstone for their assistance and insight throughout the duration of this project.

Finally, I would like to thank to the SIST committee, Fermilab, and North Central College for this opportunity.

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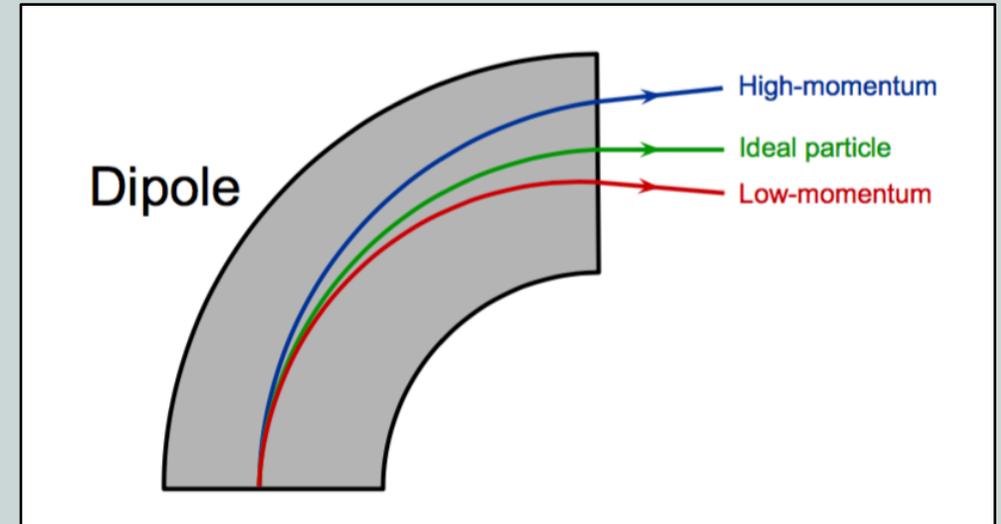
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Dispersion

- ⇒ Force on particles is velocity dependent, so particle trajectories are dependent on momentum
- ⇒ Analogous to a prism effect in optics
- ⇒ Needs to be compensated for in measured beam widths, but we do not have a confident measure of $\frac{\sigma p}{p}$ at this time



$$\sigma_{measured} = \sqrt{\beta\epsilon - \left(D \frac{\sigma p}{p}\right)^2}$$

