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IOTA electron lens: nonlinear optics, cooling, and space-charge compensation

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Introduction

Electron lens layout in IOTA (top view)



Electron-lens layout in IOTA





Roles of the IOTA electron lens

Nonlinear element for integrable optics

- thin McMillan lens
- thick axially symmetric lens
- Electron cooler
 - extend range of proton emittances and lifetimes for experiments
 - new research on electron cooling reach in nonlinear lattice
- Space-charge compensator for rings
 - shaped beam from electron gun
 - trapped electron column from residual gas

["Electron lens" = magnetically confined electrons acting on the circulating beam]

Antipov et al., JINST 12, T03002 (2017)



Typical e-lens parameters

Cathode-anode voltage	0.1 10 kV
Electron beam current	5 mA 5 A
Current density on axis	0.1 12 A/cm ²
Main solenoid length	0.7 m
Main solenoid field	0.1 0.8 T
Gun/collector solenoid fields	0.1 0.4 T
Max. cathode radius	15 mm
Lattice amplitude function	0.5 10 m
Circulating beam size (rms), e-	0.1 0.5 mm
Circulating beam size (rms), p	1 5 mm



Electron lenses for nonlinear integrable optics

Nonlinear integrable optics with electron lenses

Use the electromagnetic field generated by the electron distribution to provide the desired nonlinear field.

Linear focusing strength on axis ~ 1/m: $k_e = 2\pi \frac{j_0 L(1 \pm \beta_e \beta_z)}{(B\rho)\beta_e \beta_z c^2} \left(\frac{1}{4\pi\epsilon_0}\right)$.

1. Axially symmetric thin kick of McMillan type

current density
$$j(r) = \frac{j_0 a^4}{(r^2 + a^2)^2}$$

transverse
kick $\theta(r) = \frac{k_e a^2 r}{r^2 + a^2}$
achievable $\sim \frac{\beta k_e}{\gamma}$

 4π

Larger tune spreads in IOTA More sensitive to kick shape

2. Axially symmetric kick in long solenoid

Any axially-symmetric current distribution

$$\sim \frac{L}{2\pi\beta} = \frac{LB_z}{4\pi(B\rho)}$$

Smaller tune spreads in IOTA More robust



tune spread

Design of beam transport in electron lens



Design of McMillan e-gun

Is it possible to generate the required current-density profile? Contrasting requirements of high yield and peaked distribution

Optimization of the e-gun geometry to match the desired profile

$$j(r) = \frac{j_0 a^4}{\left(r^2 + a^2\right)^2}$$

Space-charge-limited emission determined mostly by E-field at surface =>
optimize E-field first (fast)
then, refine beam profile (slower), iterating calculation of space-charge-

limited emission



Electron gun parametric geometry



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Optimization of the electric field at the emitter



Deviations from target distribution < 5% of peak field Comparable agreement between codes Spiffe and Warp

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Electric field distribution



There is margin to decrease distances and increase perveance

Current density distribution and sensitivity

Space-charge-limited simulation of beam emission



Field optimization provides good starting point for current-density optimization (within 10% of target)

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Fermilab electron-lens test stand

dimensions in mm



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Accuracy of profile measurements in e-lens test stand

Current density measured with beam scans over collector pinhole Very reproducible, but misalignments may introduce systematics No observed effects due to repetition rate (ions, ...) Recent studies done with a hollow beam (larger and more sensitive to distortions)



Can we also measure the velocity distribution vs. radius? Charge density vs. current density



Space-charge compensation

Space-charge compensation in rings

Space-charge compensation routinely used in linacs, rf photoinjectors In rings, it would enable higher intensities A challenging subject: local correction of global effect? Issues: high charge densities, lattice distortions, beam-plasma instabilities Implementation with electron lens has advantage of magnetic confinement for stability

- Two concepts:
 - given profile (transverse/longitudinal?) from electron gun or
 - electrons from residual-gas ionization trapped in Penning-Malmberg configuration ("electron column")

Numerical simulation studies to guide experiments in IOTA



Concept of electron column



In strong field, ionization electrons mirror transverse profile of protons How does the e-column evolve?

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Electron-column simulations in IOTA with Warp

Trap configuration: 1 m length, [0 T, 0.2 T] solenoid, [0 V, -200 V] electrodes Residual gas: H₂, [10⁻³ torr, 10⁻⁵ torr] (to enhance ionization rates) Protons: 2.5 MeV, 8 mA (~5 V on axis), various profiles and time structures New protons injected in interaction volume (no ring circulation yet) Ionization processes: $p + H_2 \rightarrow p + H_2^+ + e^ \sigma = 1.5 \times 10^{-21} \text{ m}^2$

 $e^- + H_2 \rightarrow 2e^- + H_2^+$ $\sigma = 1.3 \times 10^{-20} \text{ m}^2$

No recombination, double ionization, hydrogen clusters, ...



Simulation layout



Electron density buildup vs. electrode voltage



Example of particle distributions

Continuous proton beam B = 0.1 T, V = -5 V, $p = 5 \times 10^{-4}$ torr



Electrons reach equilibrium transverse distribution They may have enough energy to escape longitudinally

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Numerical calculations with pulsed proton beam

How does the electron column evolve after one proton beam traversal?

Protons on for 100 ns, then off for 900 ns (as an example)

Transit time in e-column is 46 ns, revolution time in IOTA is 1.8 us

link to animation



Longitudinal electron density vs. time



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Longitudinal electron density vs. time



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Evolution of electron density



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Numerical calculations with pulsed proton beam

How does the electron column evolve after one proton beam traversal?

Protons on for 100 ns, then off for 900 ns

Electrons are confined transversely and oscillate longitudinally, with little loss

lons are lost both transversely and longitudinally

Oscillations are determined by secondary electron velocity and by plasma frequency

Distributed electrode voltages can help shape the charge distribution



Electron cooling

Electron cooling

1.36-keV electrons match the velocity of 2.5-MeV protonsA wider range of proton lifetimes and brightnesses will be available for experiments

Cooling option determined the co-propagating configuration of the e-lens

Cooling rates of 0.1 s are achievable Emittances can be reduced by a factor 10 Better models of magnetized cooling are needed for predictions

Does nonlinear integrable optics combined with cooling enable higher brightnesses?

Stancari et al., COOL15 Antipov et al., JINST **12**, T03002 (2017)



Proton beam diagnostics through recombination

Spontaneous recombination generates neutral hydrogen with distribution of Rydberg states, some of which are Lorentz-stripped in e-lens toroid and IOTA dipole $m + e^{-} \rightarrow H^0 + hu$

$$p + e^- \to H^0 + h\nu$$

Recombination rate at detector is \sim 50 kHz; good compromise between beam lifetime and measurement time



A critical diagnostic tool for cooling and proton beam evolution Hardware options identified; needs final design



Gun and collector solenoids reused from Tevatron Now awaiting for magnetic measurements



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Several Tevatron power supplies for magnets and HV can be reused







Magnetic system (main and transport solenoids) is behind schedule, limited by resources



Electrode and pickup structure



- beam-position monitors
- •signal pickups for charge oscillations
- •trapping/clearing electrodes
- Antenna for high-frequency plasma and cyclotron osc.

=> estimates of electron density and temperature

Need technical design, hardware, electronics



5 diagonally split cylinders or combined horiz./vert. striplines



New diagnostic stations for e-beam current and profile



Conclusions

Conclusions

The electron lens in IOTA will enable new experiments in nonlinear optics, electron cooling, and space-charge compensation

Design challenges are related to the multiple functions and the limited physical space

The project is closely related to electron-lens applications in other machines, such as beam-beam compensation and hollow electron beam collimation

There are several opportunities for collaborators to make an impact: theory, numerical calculations, diagnostics, hardware, experiments

Thank you for your attention!

