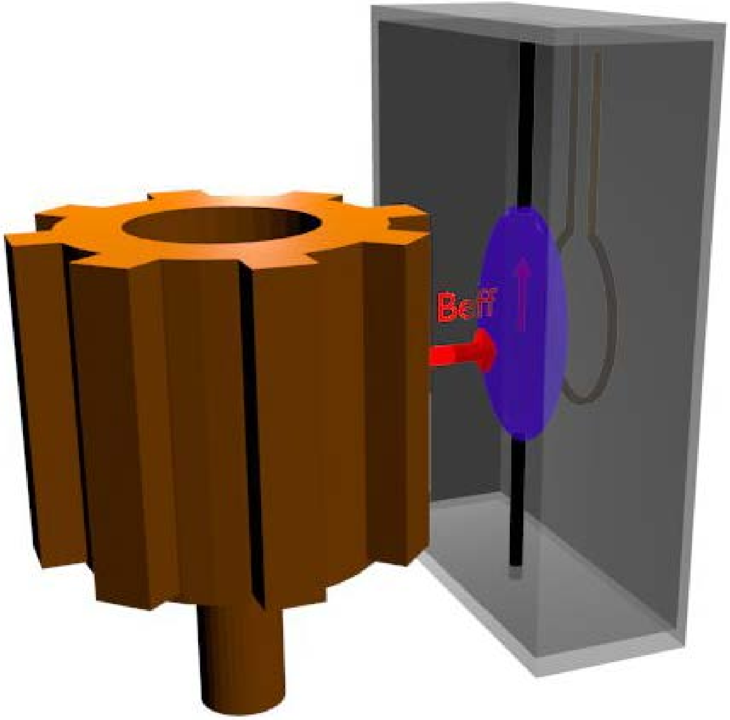
The Axion Resonant InterAction Detection Experiment (ARIADNE)



A. Arvanitaki and AG., *Phys. Rev. Lett*. 113,161801 (2014).

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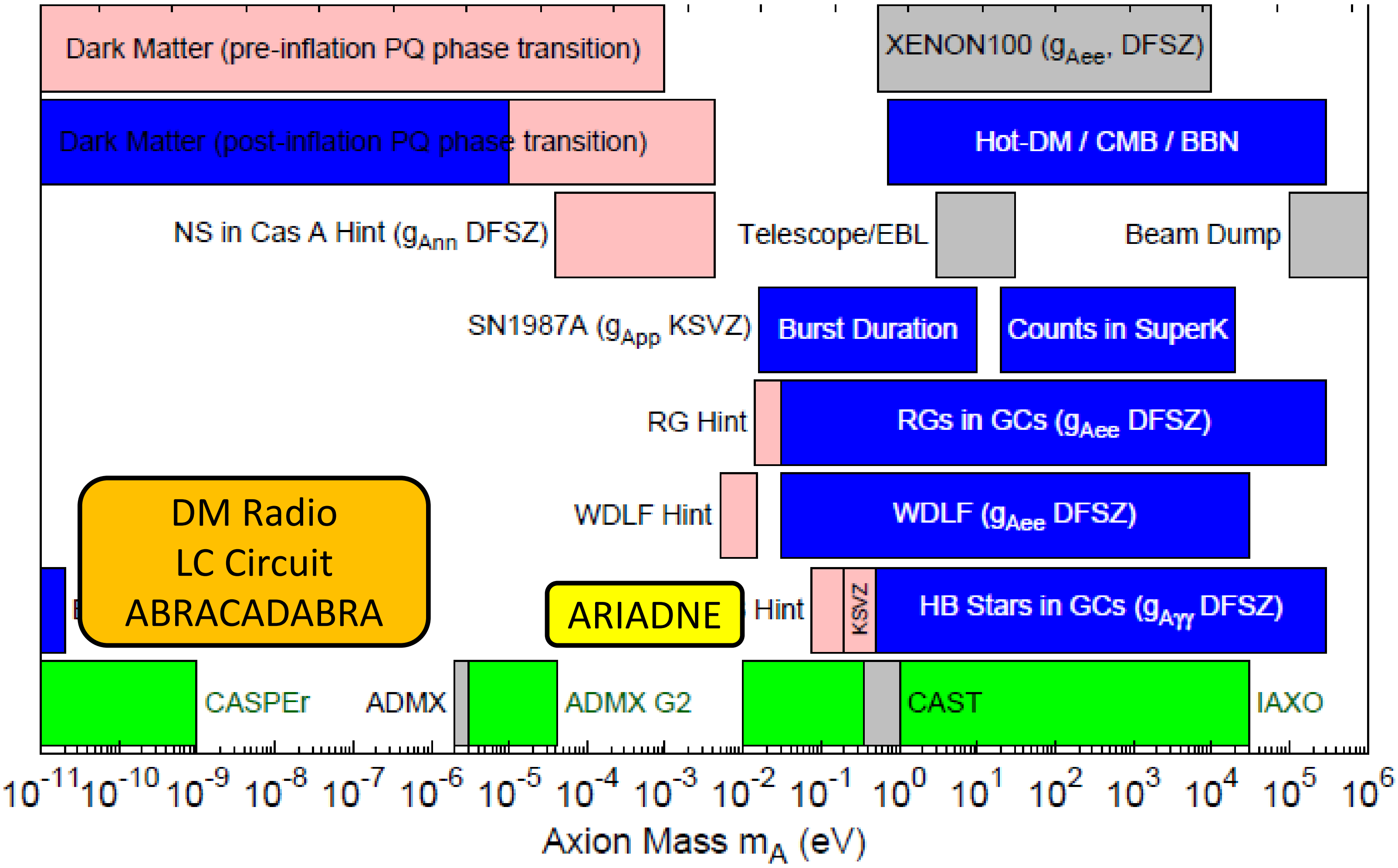
DOE Cosmic Visions Workshop, Mar. 22, 2017

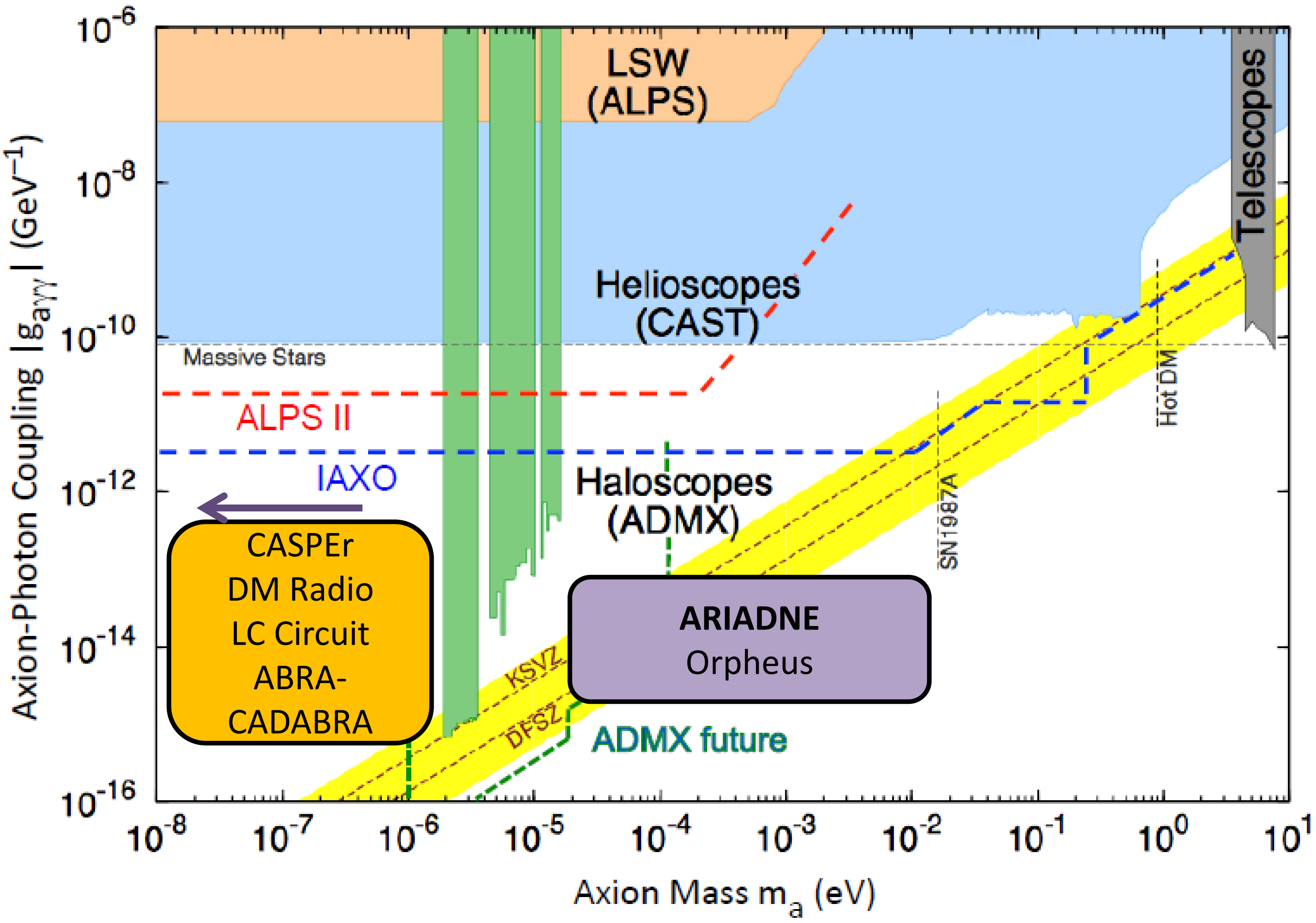
A. Geraci, University of Nevada Reno

QCD Axion parameter space



Adapted from http://pdg.lbl.gov/2015/reviews/rpp2015-rev-axions.pdf

Axion Parameter space



Adapted from Graham, arxiv: 1602.00039

Axion and ALP searches

Source	**Coupling**	
	Photons	**Nucleons**
Dark Matter (Cosmic) axions	**ADMX, ADMX-HF** **DM Radio, ABRA-CADABRA,LC Circuit** **Orpheus, CULTASK**	**CASPEr-Electric** **CASPEr-Wind**
Solar axions	CAST IAXO	
Lab-produced axions	Light-shining-thru-walls (ALPS, ALPS-II)	ARIADNE

Axion-exchange between nucleons

- Scalar coupling ∝ θ_{QCD}

$$\mathcal{L} \supset \frac{\theta_{QCD}}{f_a}\mu\; a\bar{\psi}\psi$$

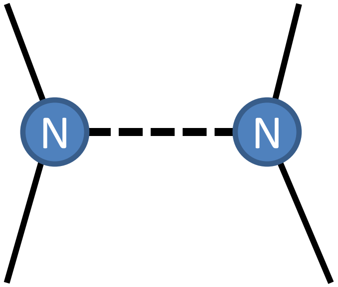
$$\mathcal{L} \supset \frac{\partial_\mu a}{f_a}\bar{\psi}\gamma_\mu\gamma_5\psi$$

- Pseudoscalar coupling

In the non-relativistic limit:

$$\mathcal{L} \supset \frac{\vec{\nabla} a}{f_a}\cdot\vec{\sigma}$$

Axion acts a force mediator between nucleons



$$(g_s^N)^2$$

Monopole-monopole

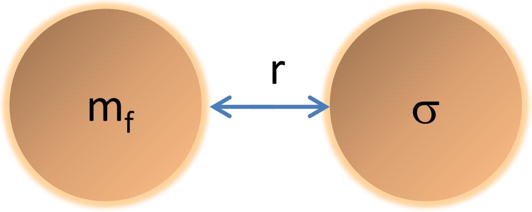
$$g_s^N g_P^N$$

Monopole-dipole

$$(g_p^N)^2$$

dipole-dipole

Spin-dependent forces



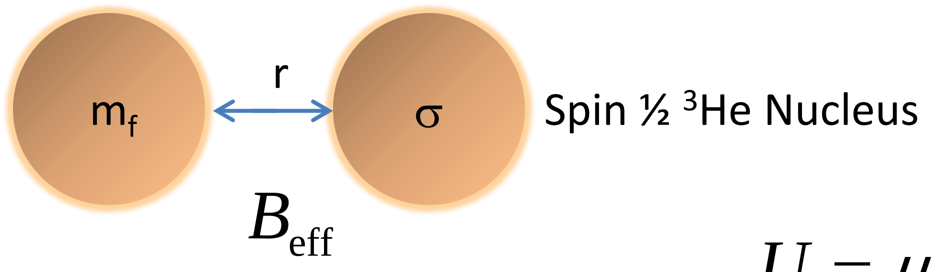
Monopole-Dipole axion exchange

$$U(r) = \frac{\hbar^2 g_s{}^N g_p{}^N}{8\pi m_f}\left(\frac{1}{r\lambda_a} + \frac{1}{r^2}\right) e^{-r/\lambda_a} (\hat{\sigma}\cdot\hat{r}) \qquad \equiv \mu \cdot B_{\text{eff}}$$

Fictitious magnetic field

- Different than ordinary B field
- Does not couple to angular momentum
- Unaffected by magnetic shielding

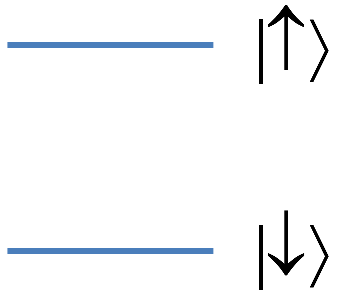
Using NMR for detection



$$U = \mu \cdot B_{\text{ext}}$$

Bloch Equations

$$\frac{d\vec{M}}{dt} = \gamma \vec{M} \times \vec{B}$$

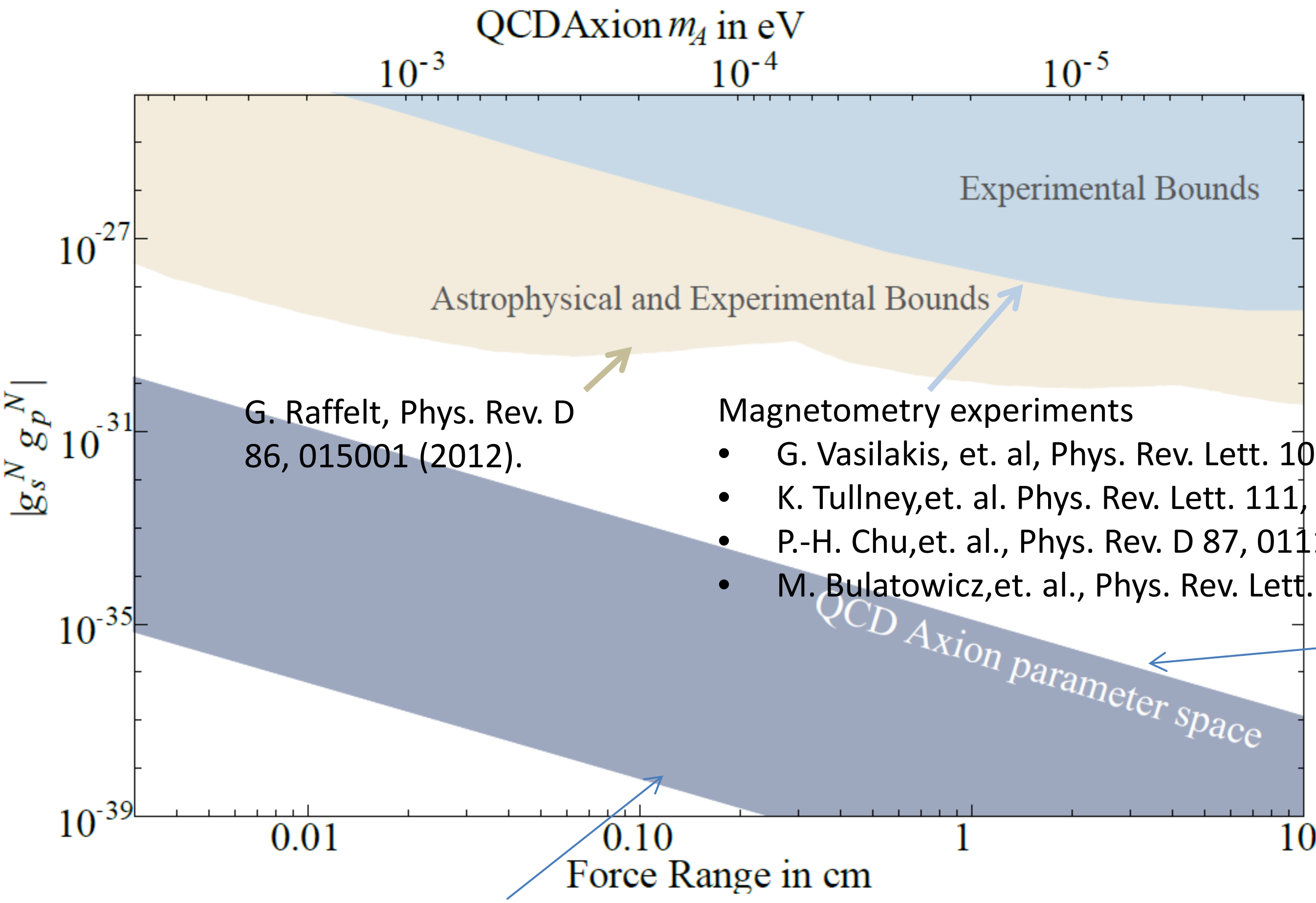


$$\omega = \frac{2\mu_N \cdot B_{\text{ext}}}{\hbar}$$

Spin precesses at nuclear spin Larmor frequency $\omega = \gamma B$

Axion B_{eff} modifies measured Larmor frequency

Constraints on spin dependent forces



G. Raffelt, Phys. Rev. D 86, 015001 (2012).

Magnetometry experiments

- G. Vasilakis, et. al, Phys. Rev. Lett. 103, 261801 (2009).
- K. Tullney,et. al. Phys. Rev. Lett. 111, 100801 (2013)
- P.-H. Chu,et. al., Phys. Rev. D 87, 011105(R) (2013).
- M. Bulatowicz,et. al., Phys. Rev. Lett. 111, 102001 (2013).

nEDM bound

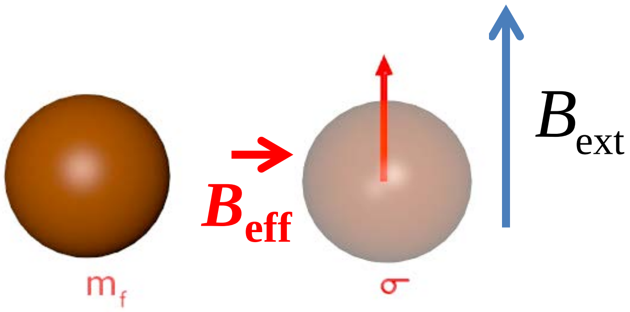
$$\theta_{QCD} < 10^{-10}$$

Standard Model lower bound

$$\theta_{QCD} > 10^{-16}$$

ARIADNE: uses resonant enhancement

Spin ½ ^{3}He Nucleus



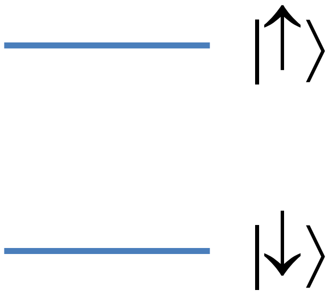
Oscillate the mass at Larmor frequency

$$B_{\text{eff}} = B_{\perp} \cos(\omega t)$$

$$U = \mu \cdot B_{\text{ext}}$$

Bloch Equations

$$\frac{d\vec{M}}{dt} = \gamma \vec{M} \times \vec{B}$$



$$\omega = \frac{2\mu_N \cdot B_{\text{ext}}}{\hbar}$$

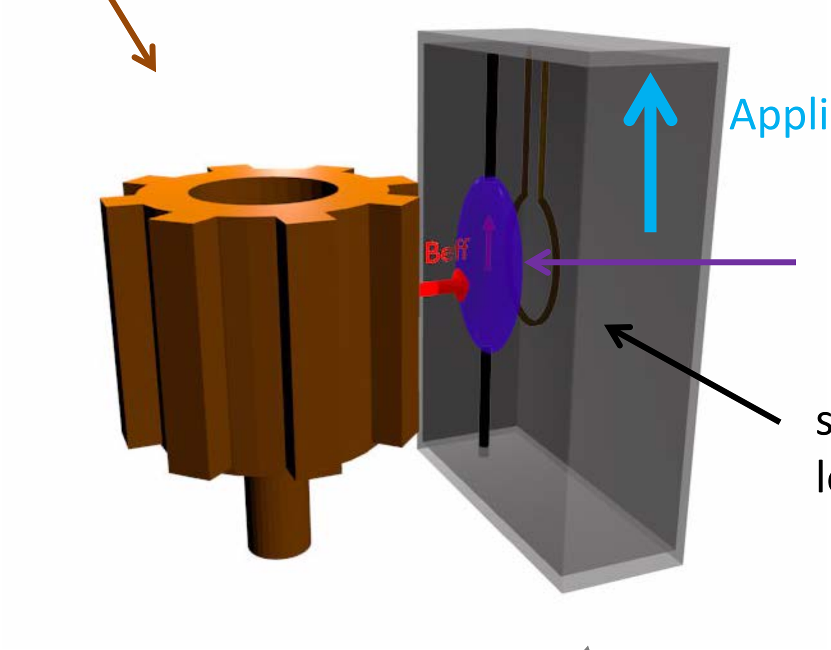
Time varying Axion B_{eff} drives spin precession
→ produces transverse magnetization

Amplitude is resonantly enhanced by Q factor ~ ωT_2.

Can be detected with a SQUID

Concept for ARIADNE

Unpolarized (tungsten) segmented cylinder sources B_{eff}



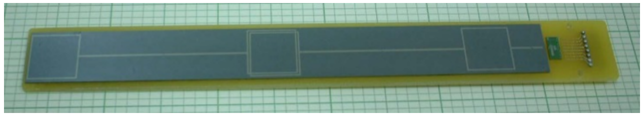
Applied Bias field B_{ext}

$$\omega = \frac{2\mu_N \cdot B_{ext}}{\hbar}$$

Laser Polarized ^{3}He gas senses B_{eff} (Indiana U)

squid pickup loop

Y.-H. Lee (KRISS)



Superconducting shielding (Stanford)

Limit: Transverse spin projection noise

$$B_{\min} \approx p^{-1}\sqrt{\frac{2\hbar}{n_s \mu_{^3\mathrm{He}} \gamma V T_2}} = 10^{-20}\frac{T}{\sqrt{Hz}} \times \left(\frac{1}{p}\right)\left(\frac{1\ \mathrm{cm}^3}{V}\right)^{1/2}\left(\frac{10^{21}\ \mathrm{cm}^{-3}}{n_s}\right)^{1/2}\left(\frac{1000\ \mathrm{sec}}{T_2}\right)^{1/2}$$

A. Arvanitaki and A. Geraci, *Phys. Rev. Lett*. 113, 161801 (2014).

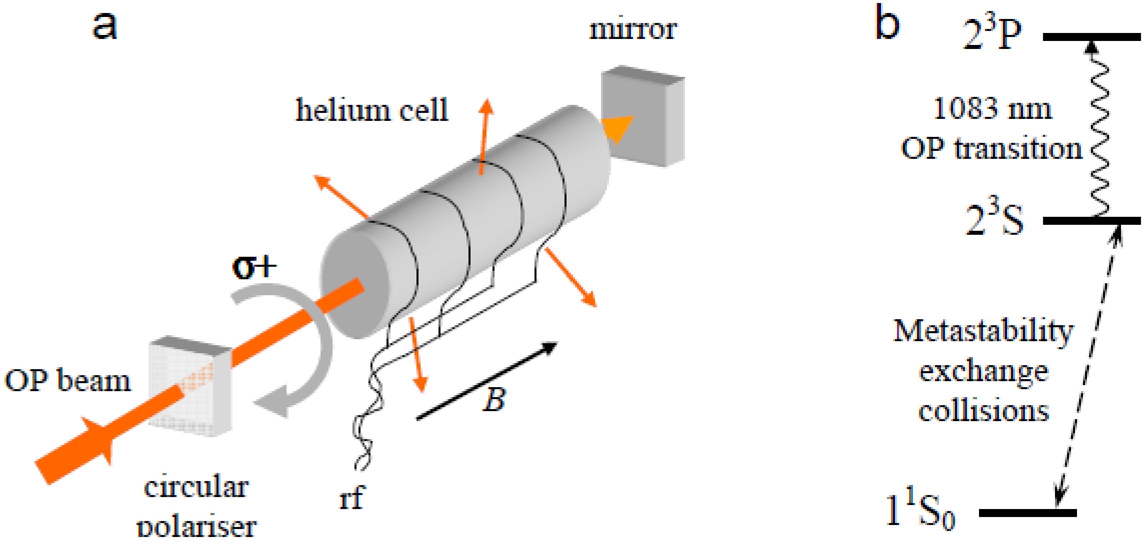
Hyperpolarized ^{3}He

- Ordinary magnetic fields cannot be used to reach near unity polarization

$$\exp[-\mu_N B / k_B T]$$

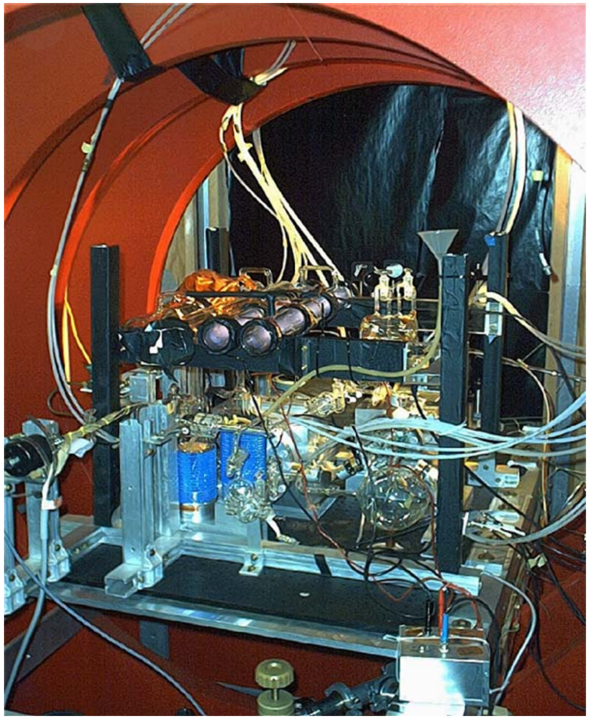
Optical pumping techniques

- Metastability exchange optical pumping



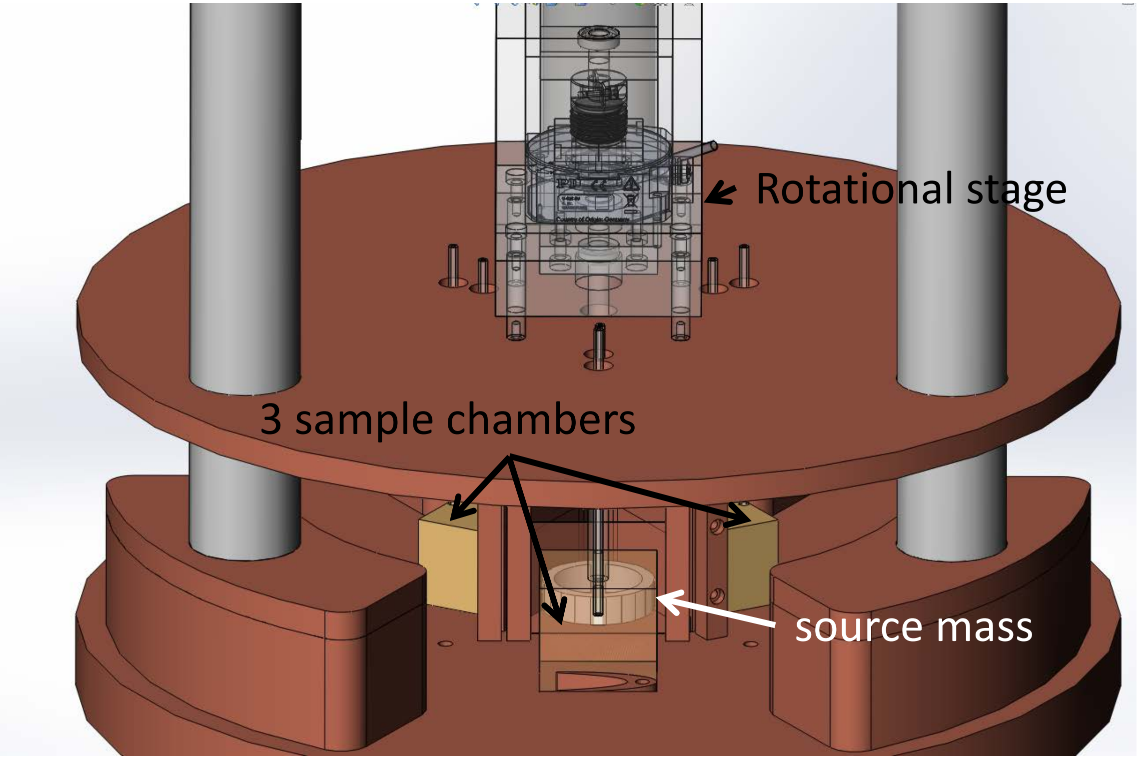
M Batz, P-J Nacher and G Tastevin, Journal of Physics: Conference Series **294** (2011) 012002

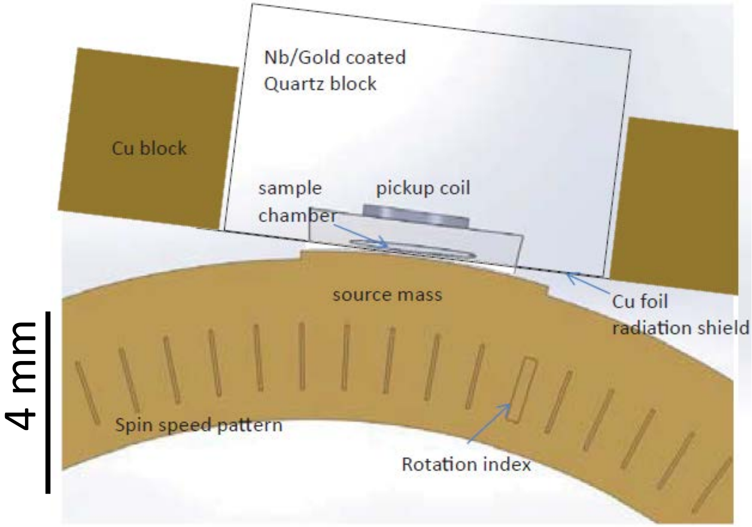
Indiana U. MEOP apparatus



Rev. Sci. Instrum. 76, 053503 (2005)

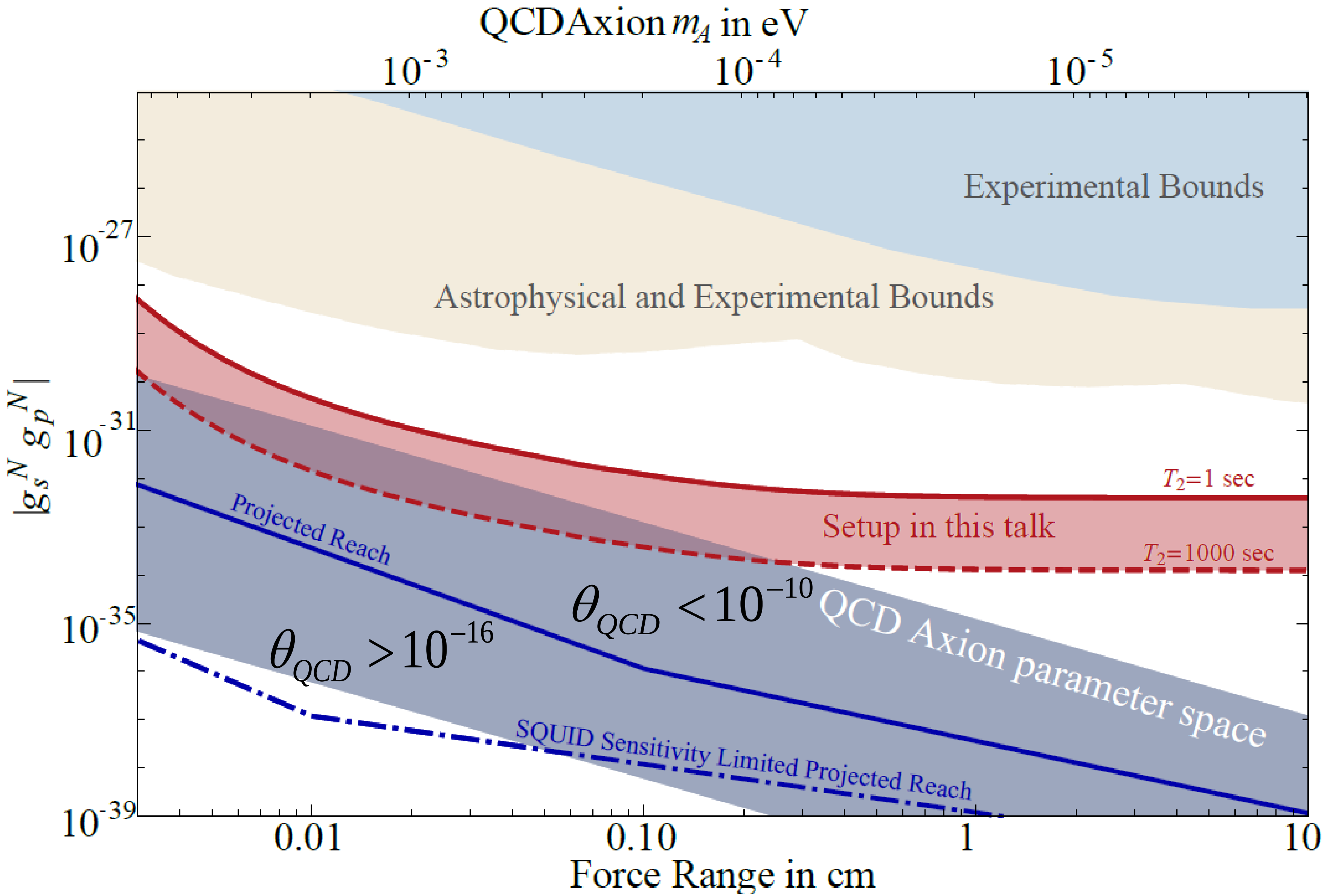
Experimental parameters





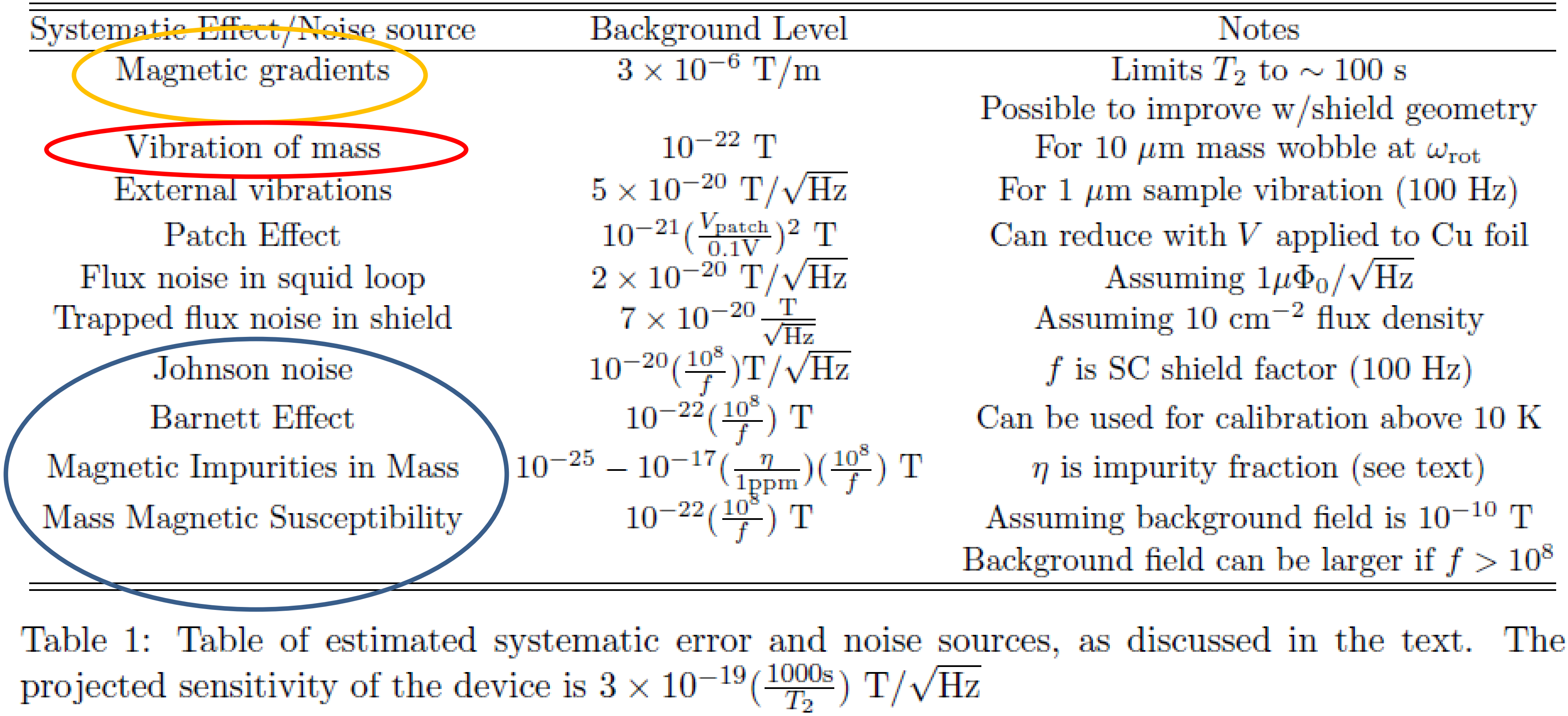
11 segments
100 Hz nuclear spin precession frequency
2 x 10^{21} / cc ^{3}He density
10 mm x 3 mm x 150 µm volume
Separation 200 µm
Tungsten source mass (high nucleon density)

Sensitivity



A. Arvanitaki and AG., *Phys. Rev. Lett*. 113,161801 (2014).

Experimental challenges



Systematic Effect/Noise source	Background Level	Notes
Magnetic gradients	3×10^{-6} T/m	Limits T_2 to ~ 100 s Possible to improve w/shield geometry
Vibration of mass	10^{-22} T	For 10 μm mass wobble at $\omega_{\rm rot}$
External vibrations	5×10^{-20} T/$\sqrt{\rm Hz}$	For 1 μm sample vibration (100 Hz)
Patch Effect	$10^{-21}(\frac{V_{\rm patch}}{0.1{\rm V}})^2$ T	Can reduce with V applied to Cu foil
Flux noise in squid loop	2×10^{-20} T/$\sqrt{\rm Hz}$	Assuming $1\mu\Phi_0/\sqrt{\rm Hz}$
Trapped flux noise in shield	$7 \times 10^{-20} \frac{\rm T}{\sqrt{\rm Hz}}$	Assuming 10 cm^{-2} flux density
Johnson noise	$10^{-20}(\frac{10^8}{f})$T/$\sqrt{\rm Hz}$	f is SC shield factor (100 Hz)
Barnett Effect	$10^{-22}(\frac{10^8}{f})$ T	Can be used for calibration above 10 K
Magnetic Impurities in Mass	$10^{-25} - 10^{-17}(\frac{\eta}{1{\rm ppm}})(\frac{10^8}{f})$ T	η is impurity fraction (see text)
Mass Magnetic Susceptibility	$10^{-22}(\frac{10^8}{f})$ T	Assuming background field is 10^{-10} T Background field can be larger if $f > 10^8$

Table 1: Table of estimated systematic error and noise sources, as discussed in the text. The projected sensitivity of the device is $3 \times 10^{-19}(\frac{1000{\rm s}}{T_2})$ T/$\sqrt{\rm Hz}$

- Design/Simulation Work: Magnetic gradient reduction strategy
- Experimental testing in progress: Vibration tests, Shielding factor *f* test thin-film SC

Superconducting Magnetic Shielding

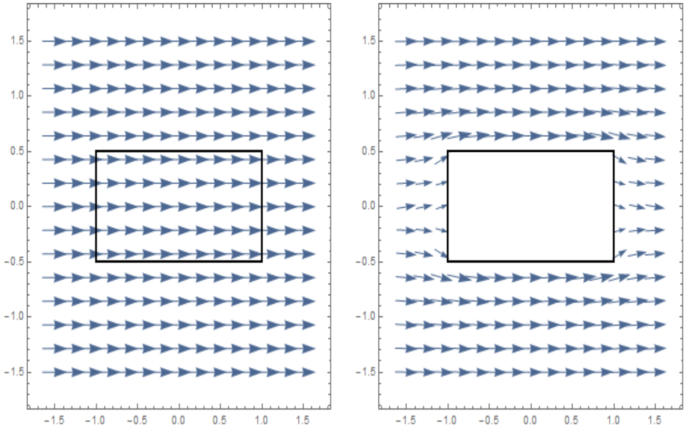
→ Essential to avoid Johnson noise

Meissner Effect

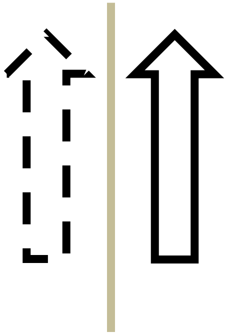
- No magnetic flux across superconducting boundary

Method of Images

- Make “image currents” mirrored across the superconducting boundary



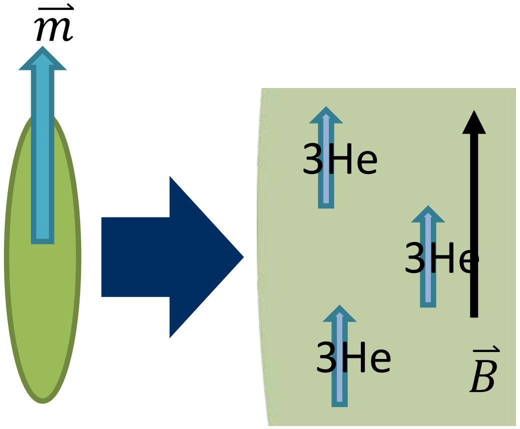
$T > T_c$ $T < T_c$



Dipole with image

The Problem of Unwanted Images

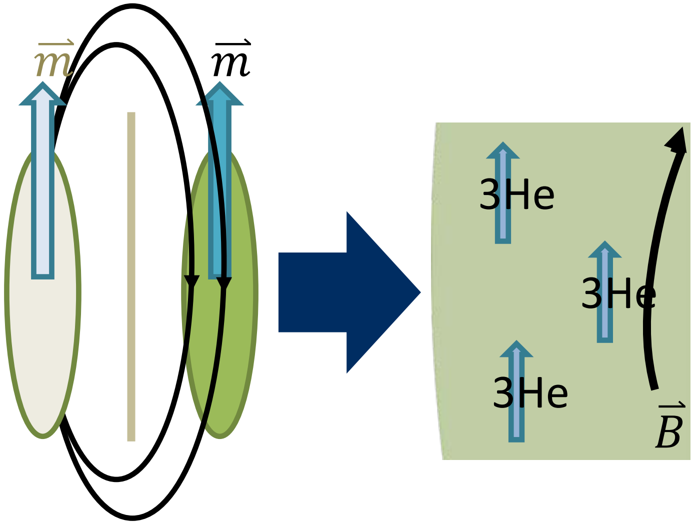
- ARIADNE uses magnetized spheroid
  - Constant interior field



- $B_{in} = \text{const.}$
- $\overrightarrow{B_{in}} \parallel \overrightarrow{m_i}$

- Magnetic shielding introduces “image spheroid”
  - Interior field varies

→ variations in nuclear Larmor frequency!

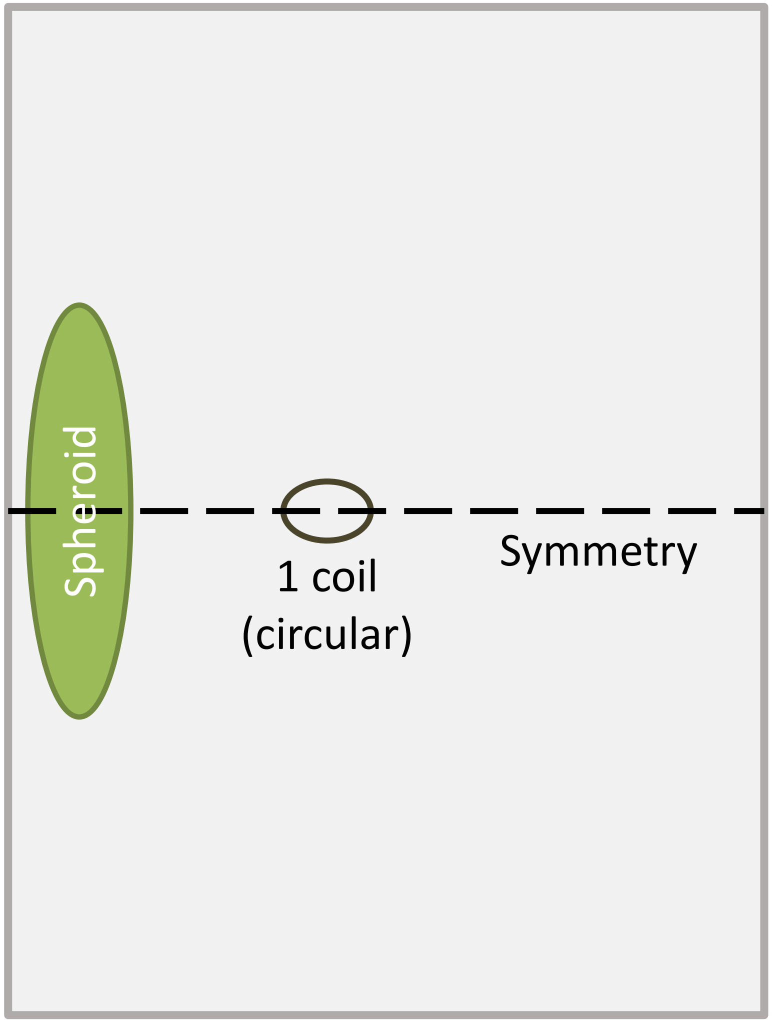


- $B_{in} \neq \text{const.}$
- $\overrightarrow{B_{in}} \nparallel \overrightarrow{m_i}$

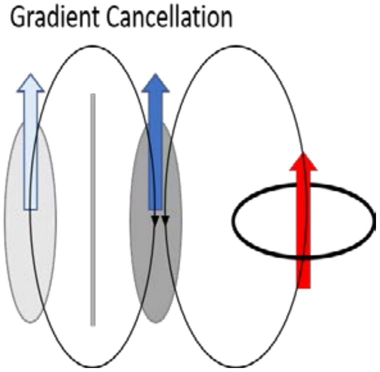
But want to drive entire sample on resonance

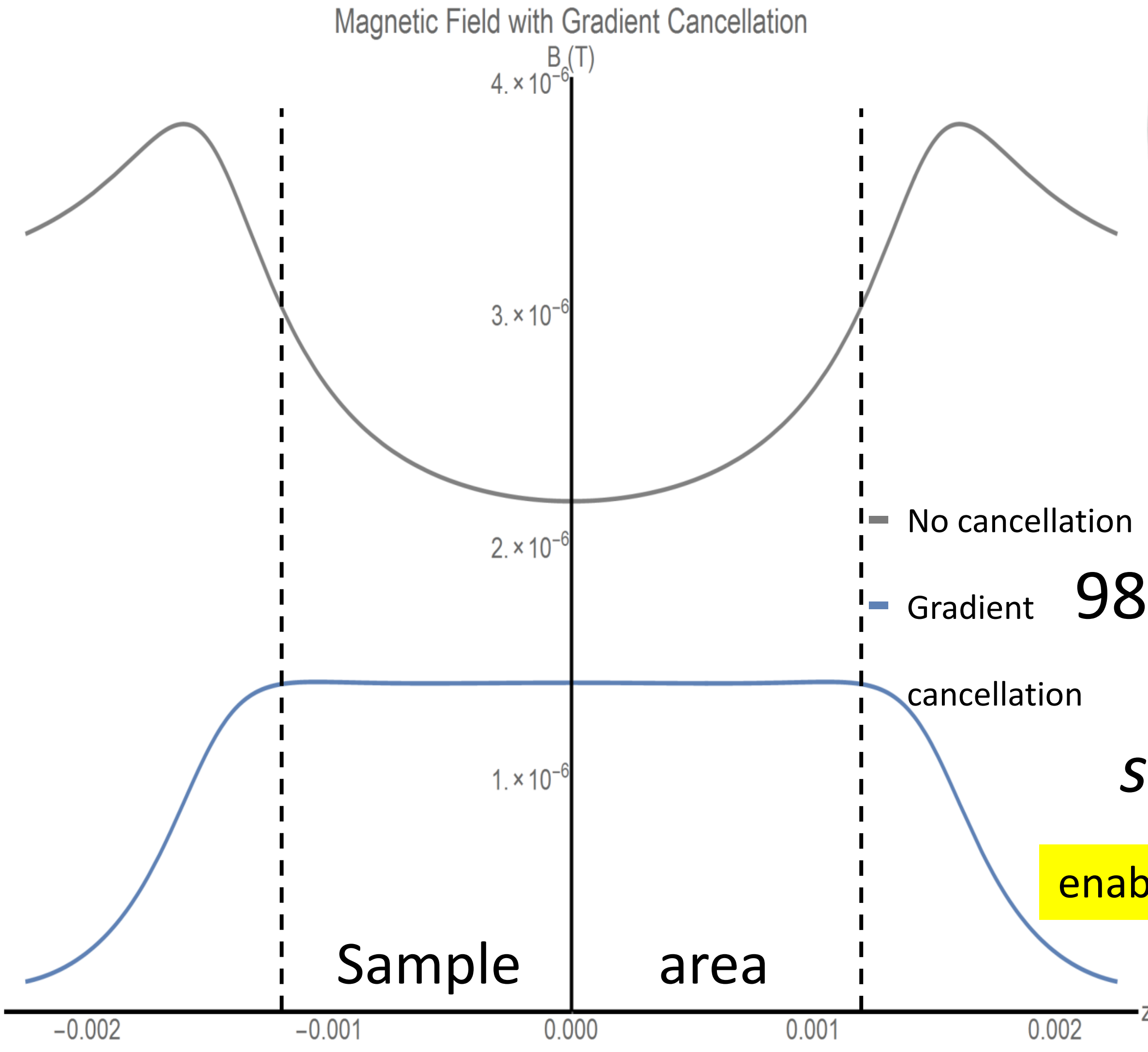
Flattening Solution

- 1 coil – simple configuration
- Expected field from spheroid ~1 μT
  - I on the 0.1 – 1 A range



Gradient Cancellation





98 times flatter

$I = 1.6$ A

$s_{\text{Frac}} = 0.17\%$

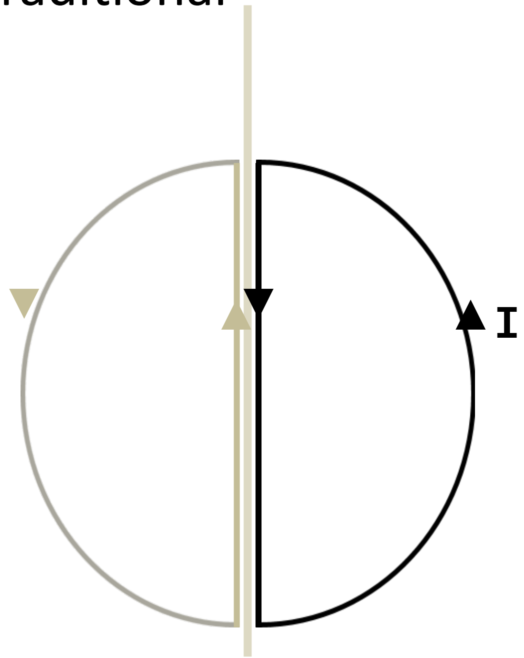
enabling T_2 of ~100 s

Tuning Solution – “D” Coils

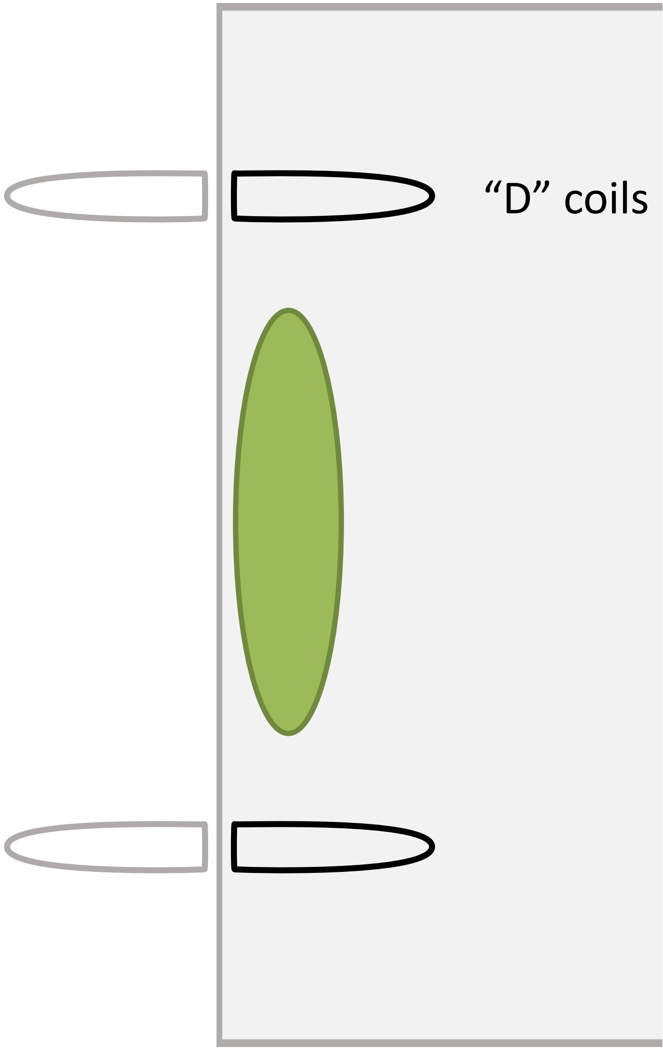
- Tune field with Helmholtz coils
  - Helmholtz field only flat near the center
  - Geometry restrictions prevent the spheroid from being centered in traditional Helmholtz coils

- “D” coils look like Helmholtz coils when their images are included

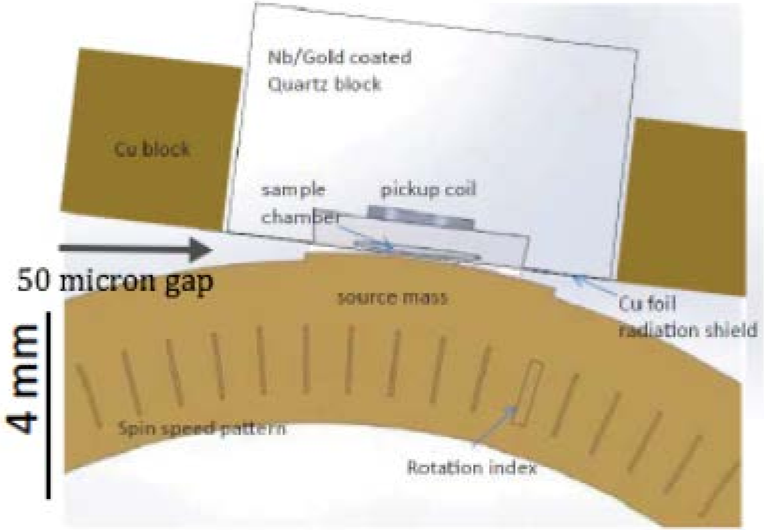
- Inner straight-line currents cancel
- Outer currents do not

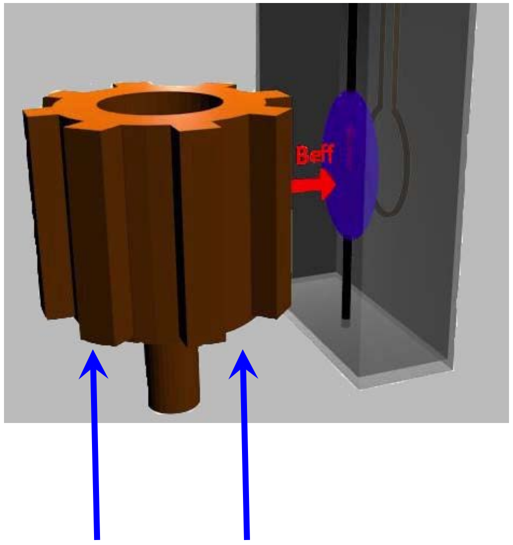


One “D” coil and image (bird’s eye view)

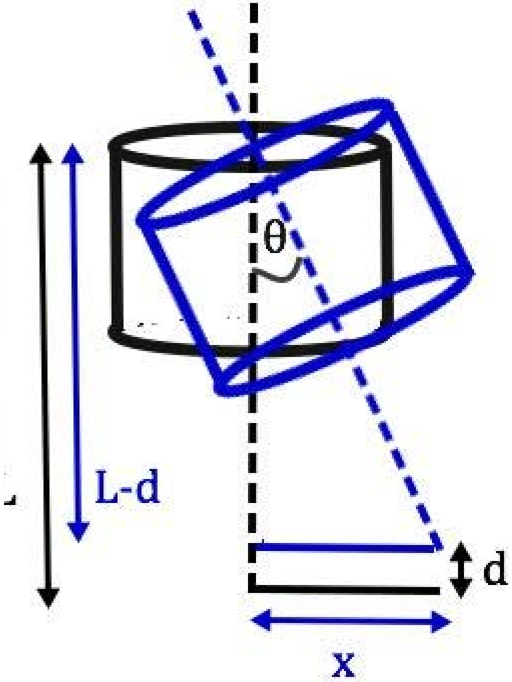


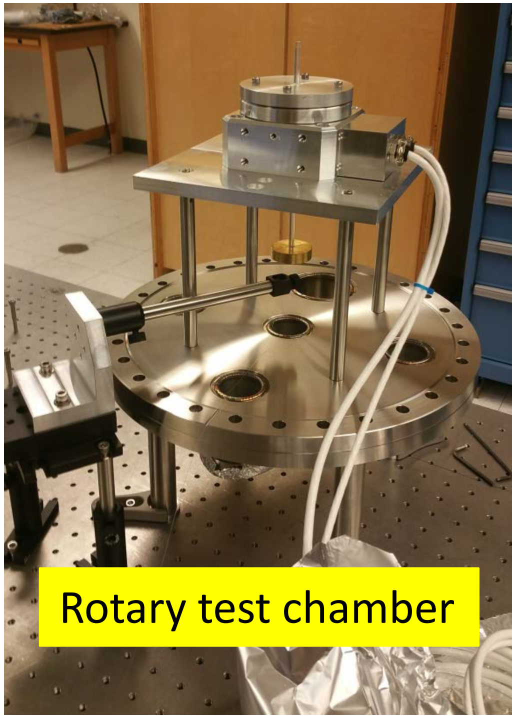
Rotary stage vibration and tilt





Interferometers





Rotary test chamber

- Build an interferometer to measure the change in distance (d).
- We can find theta (Θ) from:

  $$\Theta = \cos^{-1}((L-d)/L)$$

- We can solve for the wobble distance (X) by:

  $$X = L\sin(\Theta)$$

Fiber-coupled laser interferometers

Laser Source

Signal Photodiode

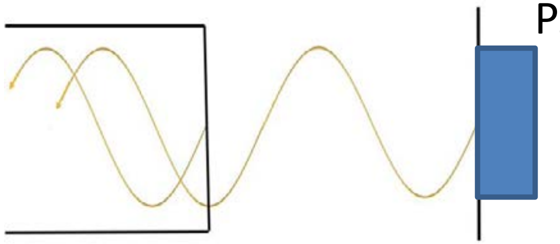
90:10 fiber coupler

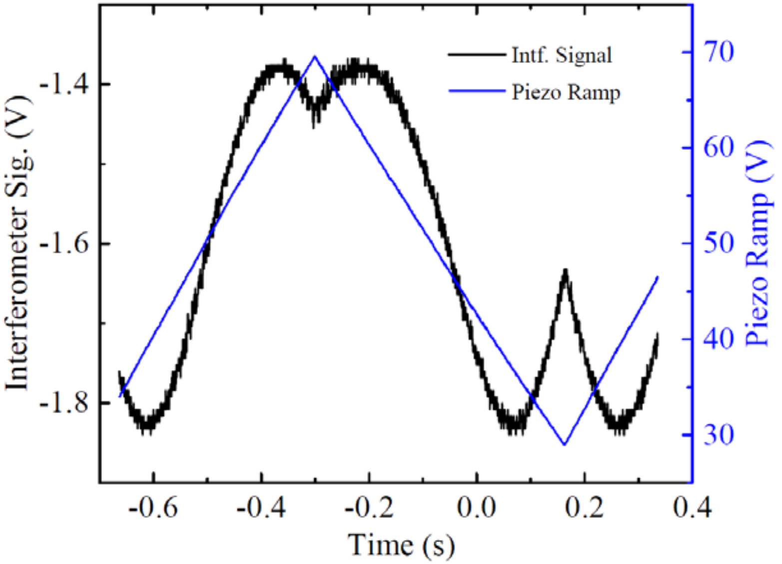
Reference Photodiode

Fiber Tip

Reflector

PZT





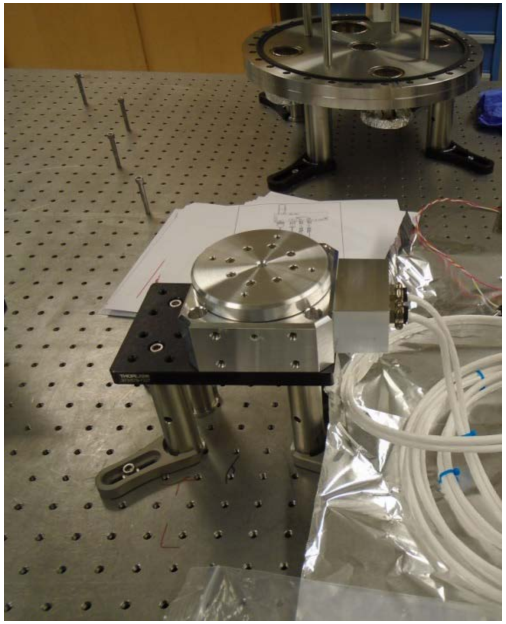
Fringe visibility
~0.13

Sensitivity
~160 nm/V

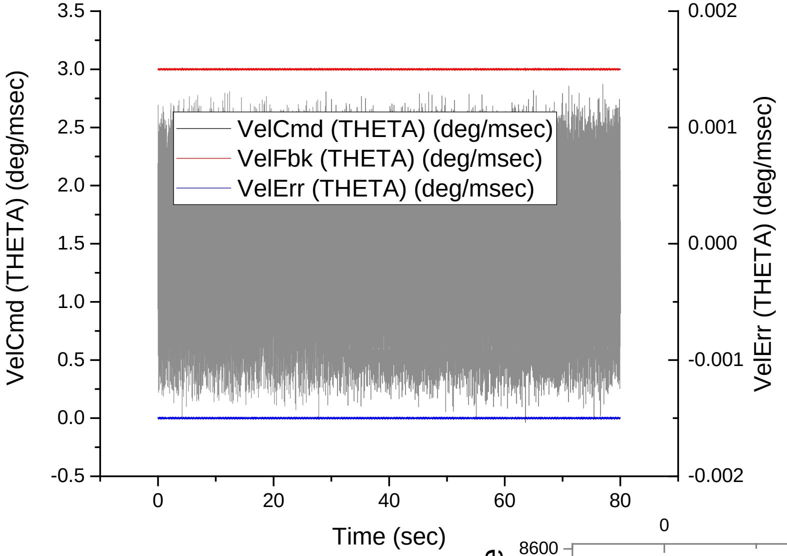
Shot noise limit
~20 pm/Hz$^{1/2}$

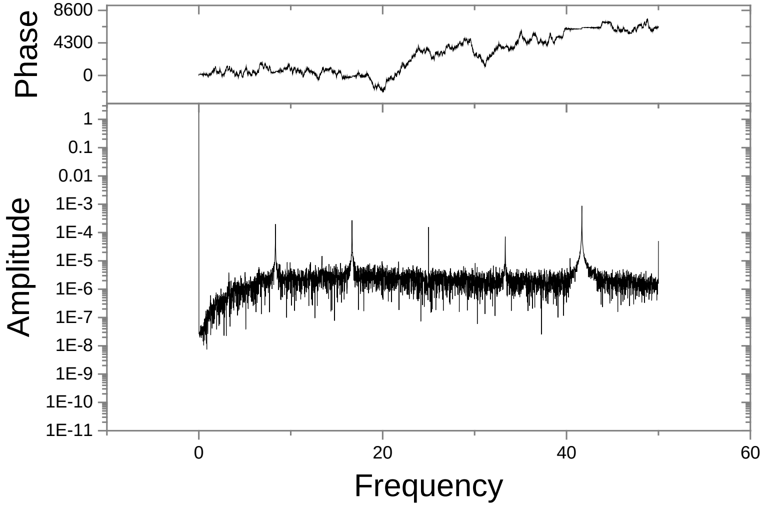
Speed stability test - direct drive stage

- Optical encoder
- Current feedback control



Stage speed stability error – unloaded, in air

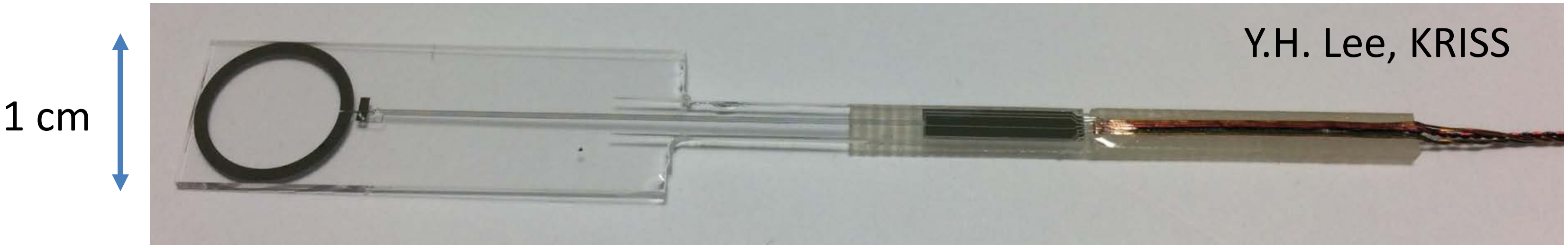




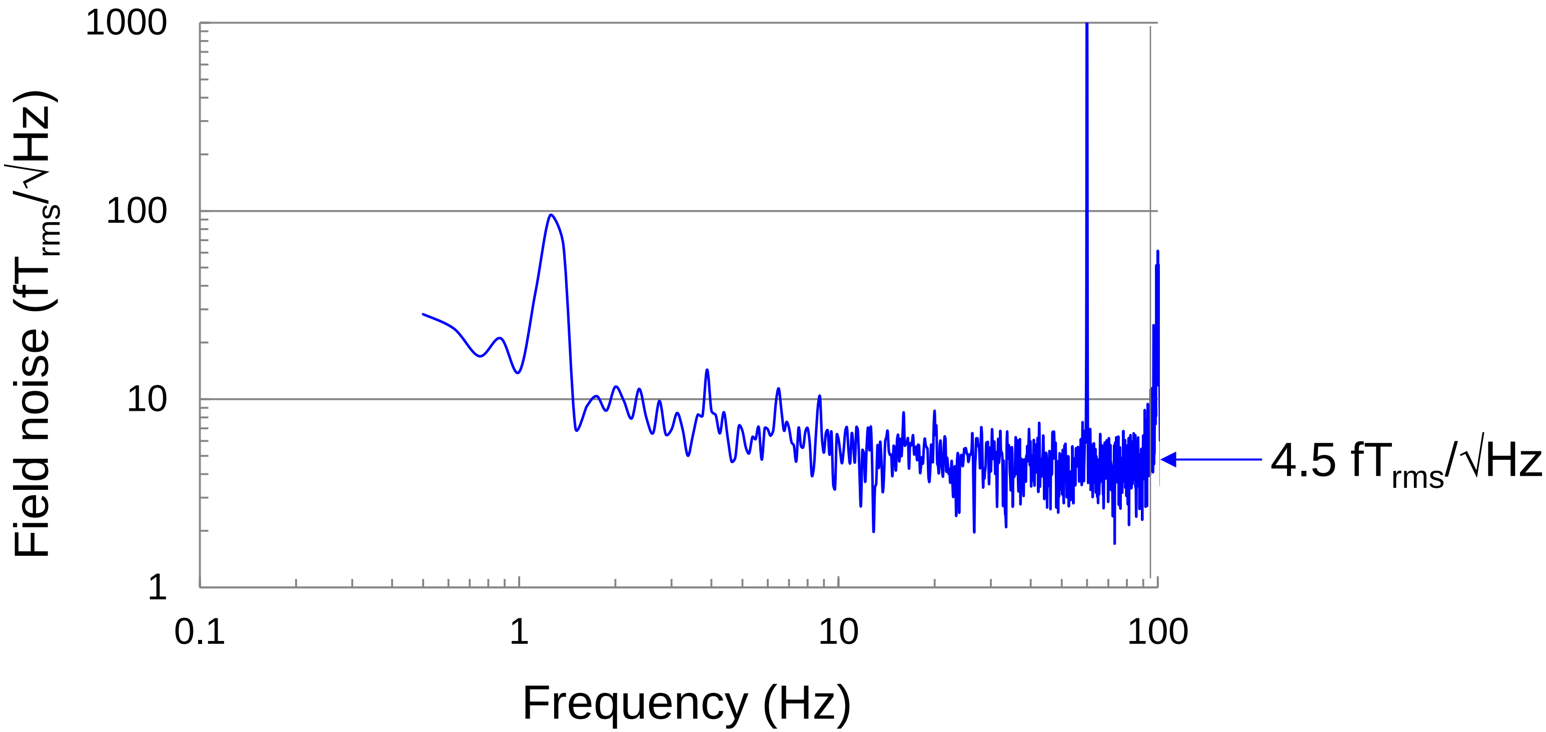
Rotation speed control
8.3 Hz ~ 1 part in 10000
RMS ~ 1 part in 3000

Allows utilization of T2 > 100s

SQUID Magnetometers



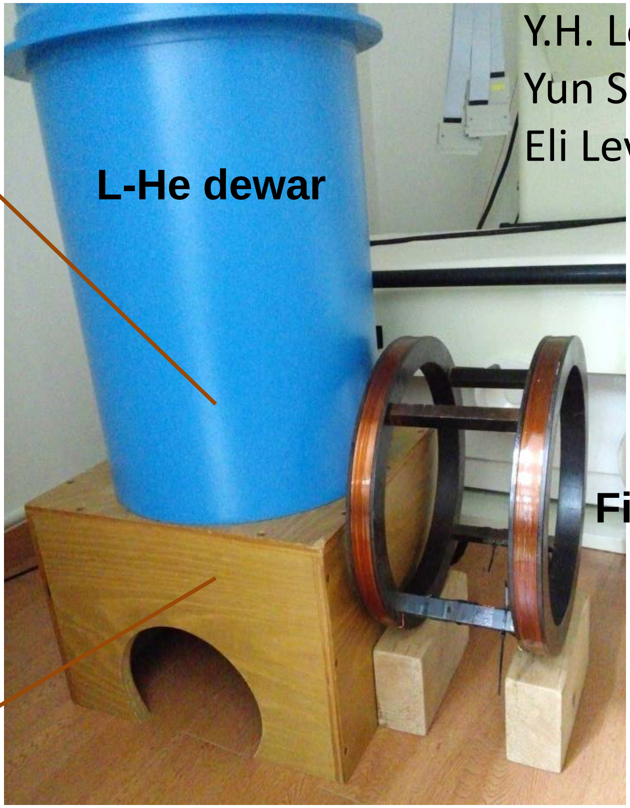
Y.H. Lee, KRISS



Measured inside a magnetically shielded room (without Nb tube)

Preliminary test of superconductive shielding

Nb tube:
23 mm ID
1 mm thick
Length 200 mm



Y.H. Lee, KRISS,
Yun Shin (CAPP)
Eli Levenson-Falk (Stanford)

Applied field: 10-100 μT_{pp} range (at 8 Hz)

SQUID magnetometer: Near the center of Nb tube
Shielding factor: ≒(0.5-3) x10^9 for transverse field

Goal: 10^8 with thin film Nb SC shield – tests planned Apr 2017

R&D and Operating Costs/Timeline

- Current funding for 1yr ($270k) for R&D, initial testing -- thank you!
- Approx. 2-3 years to finish design/construction and bring to commissioning phase ($0.6 M/yr)
- 3 years data taking phase to reach QCD axion / R&D for upgrades ($0.6 M/yr)

Summary

- ARIADNE → New resonant NMR method
- Gap in experimental QCD axion searches

  $$0.1 \text{ meV} < m_a < 10 \text{ meV}$$
- Complementary to cavity-type (e.g. ADMX) experiments
- No need to scan mass, indep. of local DM density
- Next tests – shielding (Stanford/Korea), vibration (UNR), ^{3}He system (Indiana)

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Kirsten Casey (UG)

Bella Rodriguez(UG)

Gambhir Ranjit (PD)

Not pictured: Apryl Witherspoon (UG), Ohidul Mojumder (UG), Hannah Mason (UG)