



# ABRACADABRA

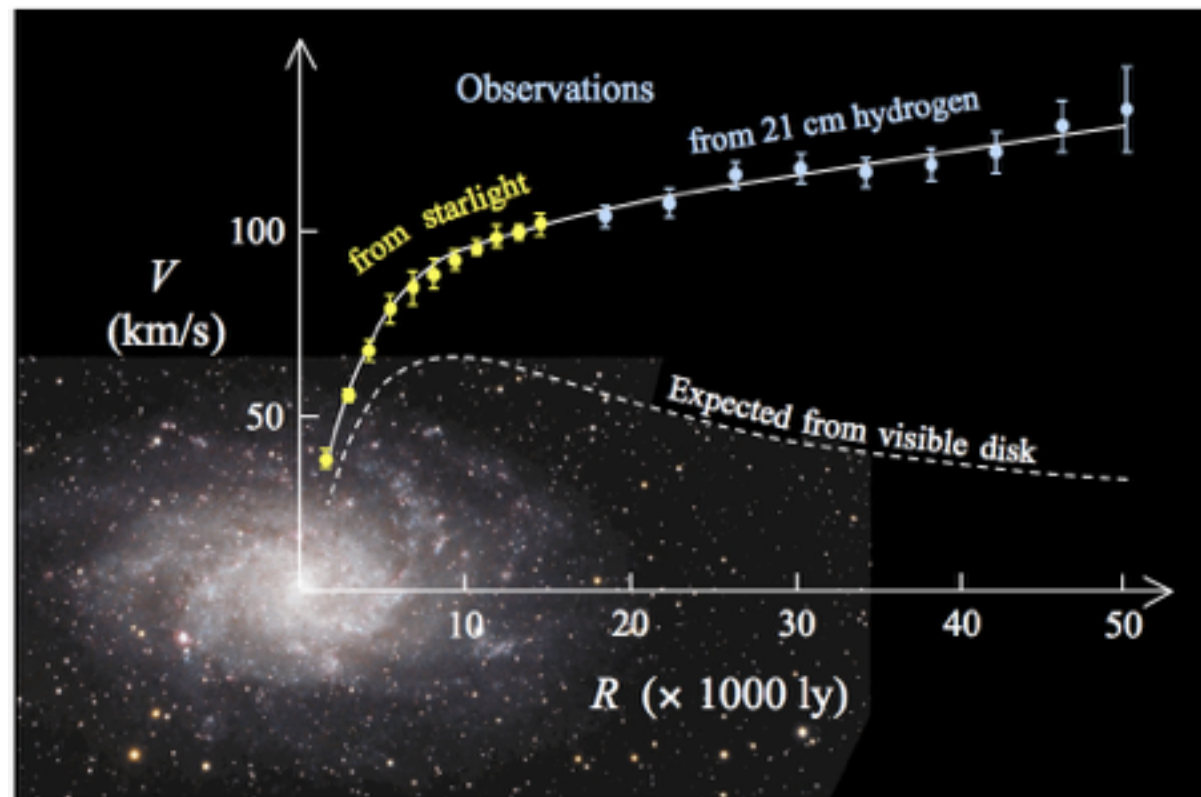
## A Broadband Search for Axion Dark Matter

Jonathan Ouellet  
*Massachusetts Institute of Technology*

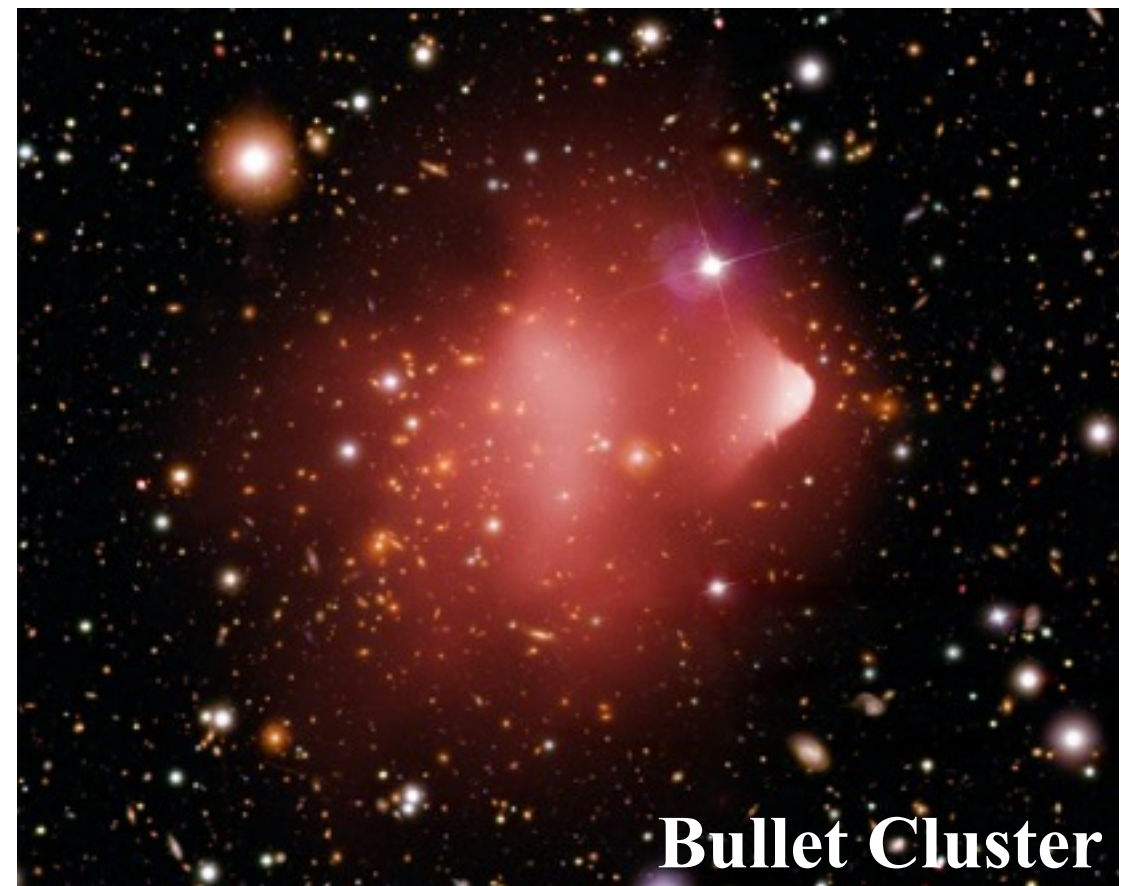
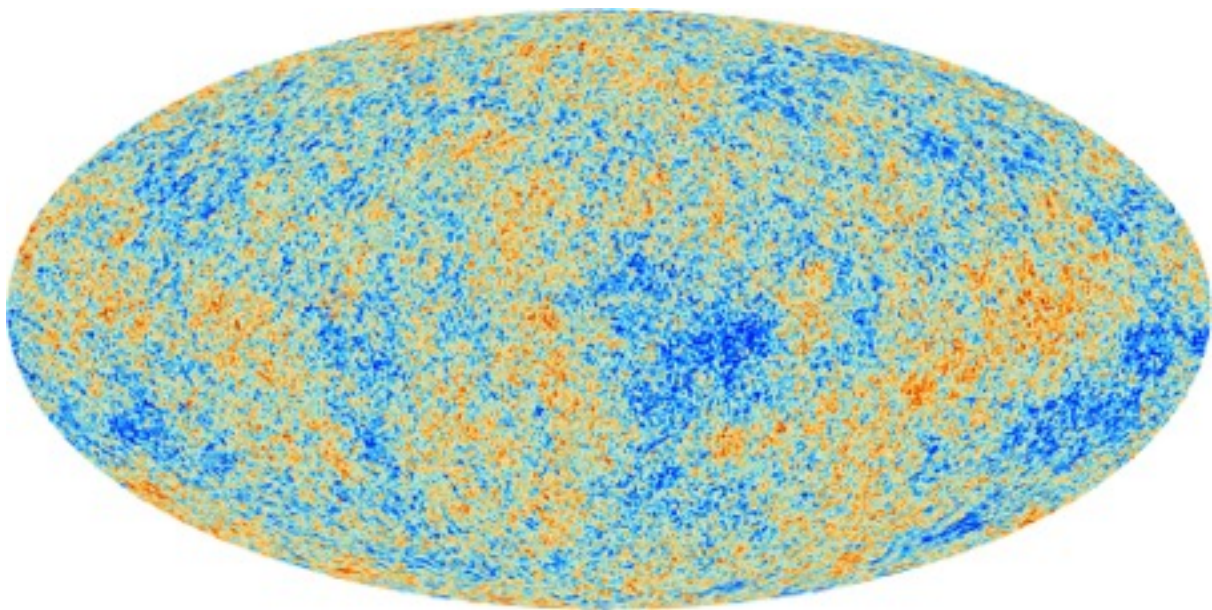
March 23, 2017



# We can see there's something we can't see!



- ▶ Galaxies are spinning too fast, but not flying apart!
- ▶ There's a footprint in the CMB!
- ▶ Galaxy mergers



**Bullet Cluster**



# So, what is Dark Matter?

## ► We know what it's not:

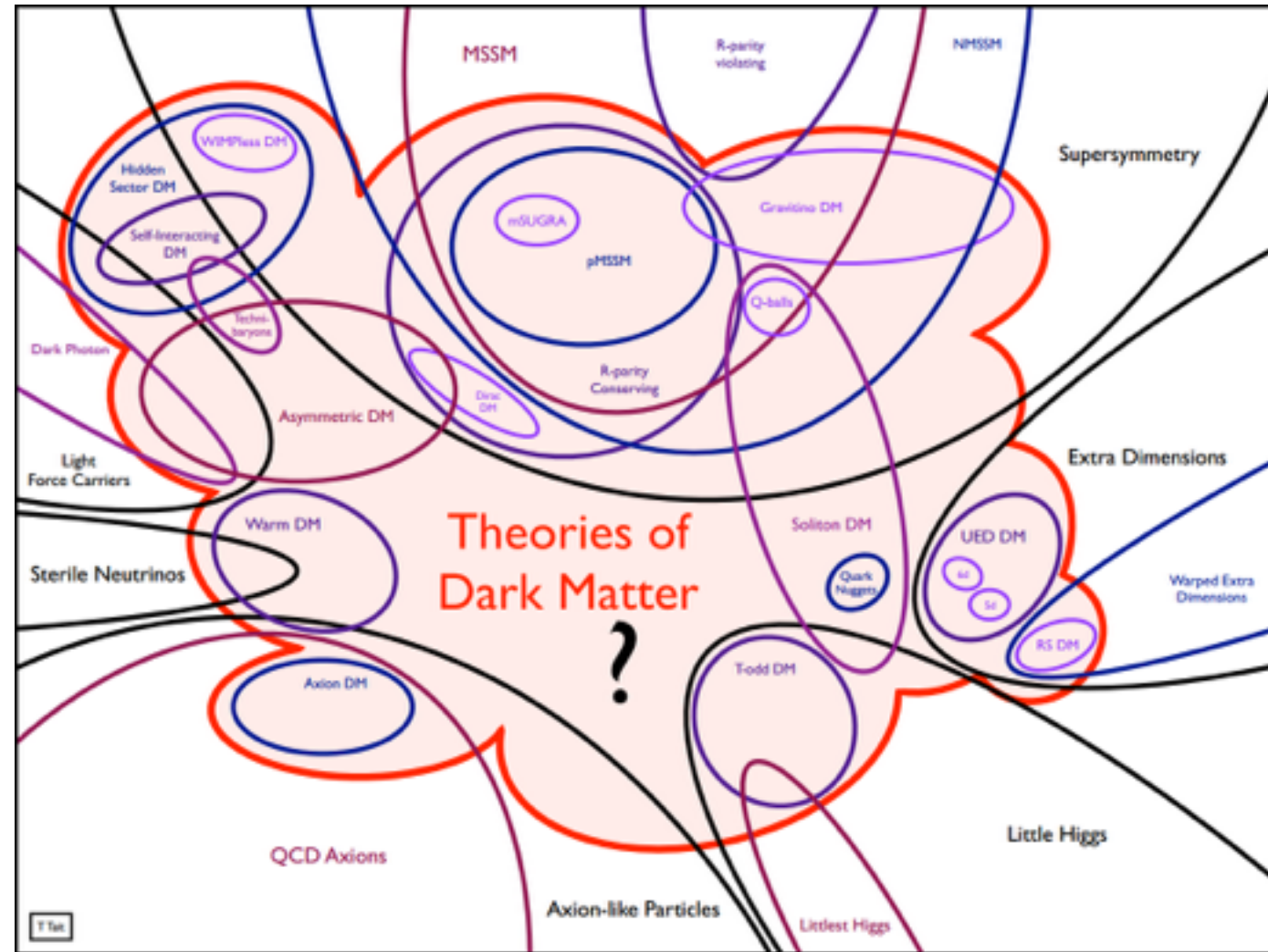
- It is not normal matter

## ► But we have no evidence about what it could be

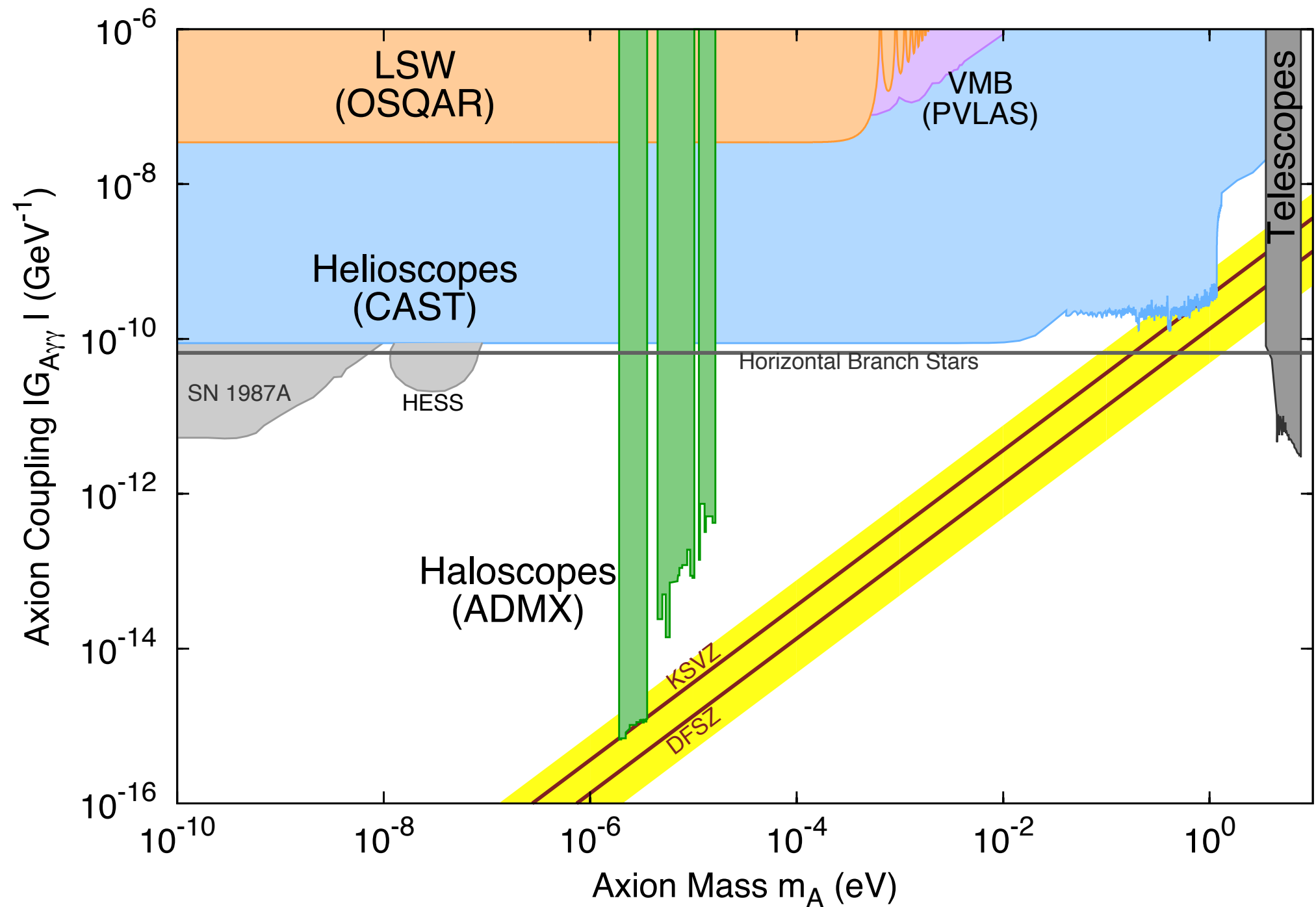
- ~50 orders of magnitude in viable mass range
- Little restriction on how it interacts (just feebly and gravitationally)

## ► The best motivated theories simultaneously solve two problems:

- DM + Supersymmetry (WIMPs)
- DM + BAU (~GeV Scale WIMPs)
- DM + Strong CP problem (Axions)

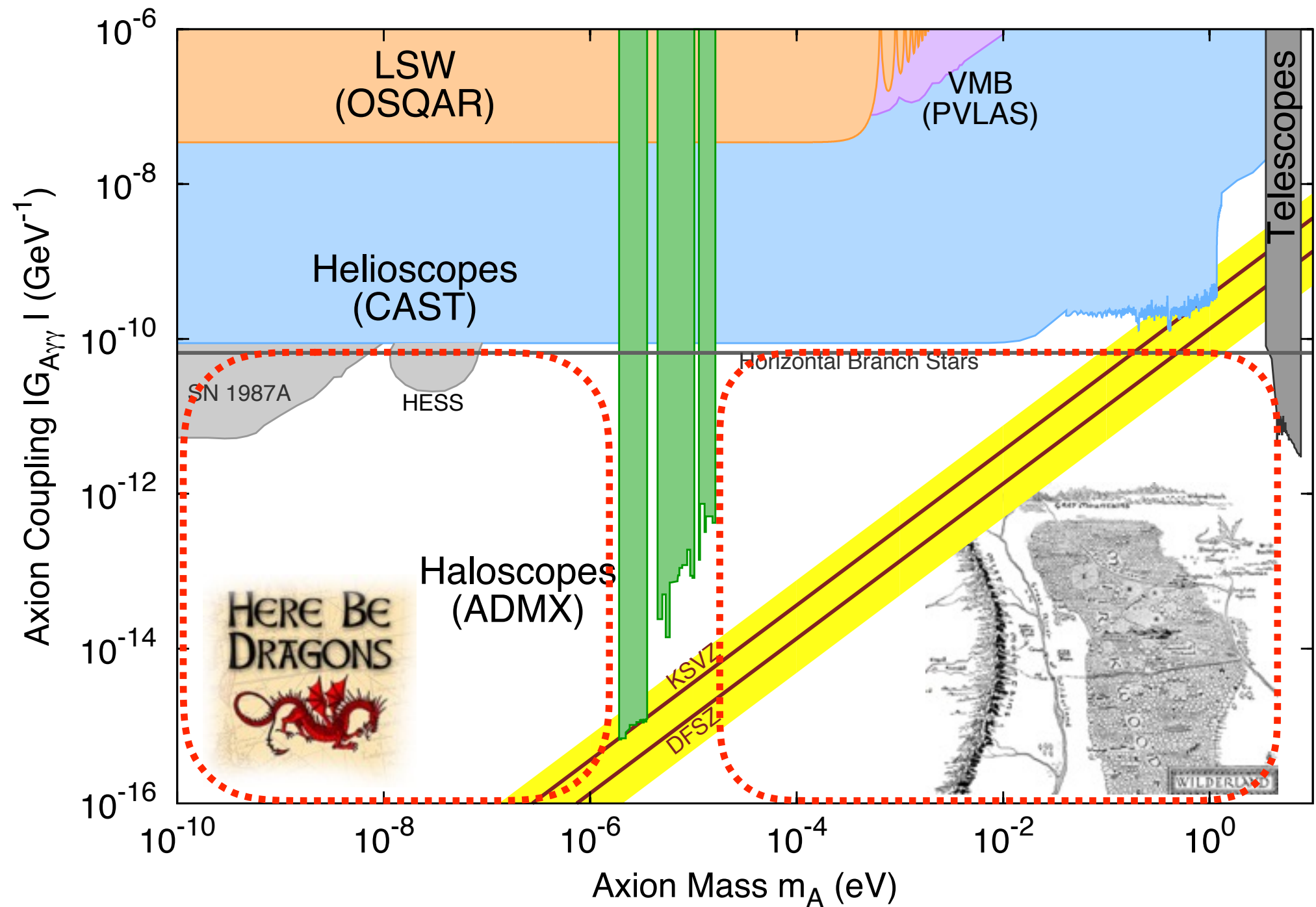


# Axion and ALP Parameter Landscape



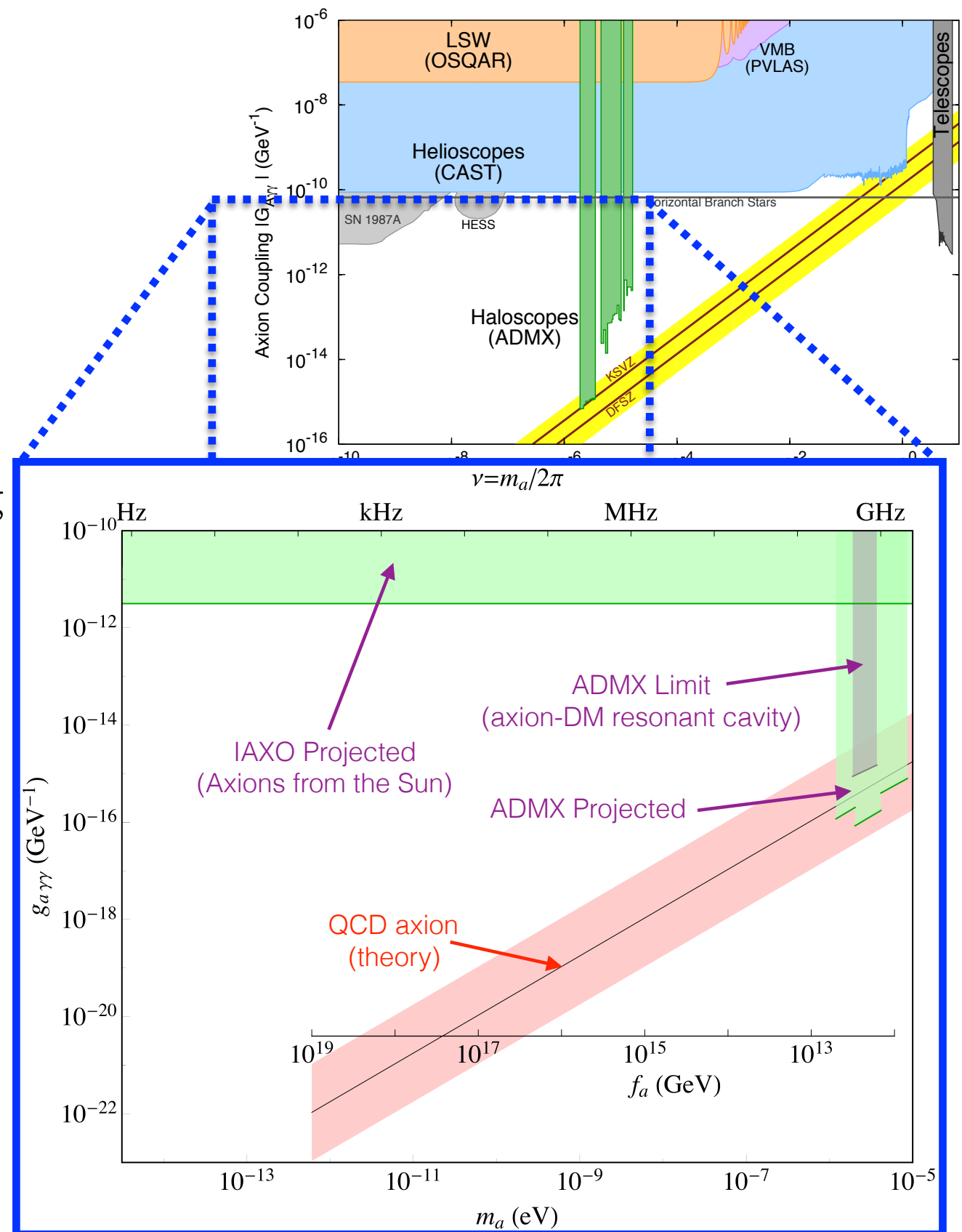
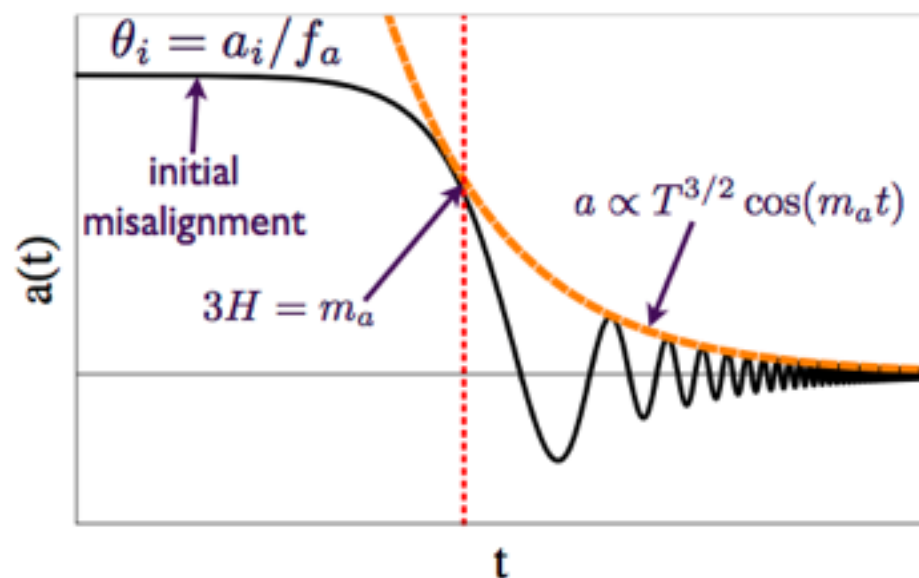


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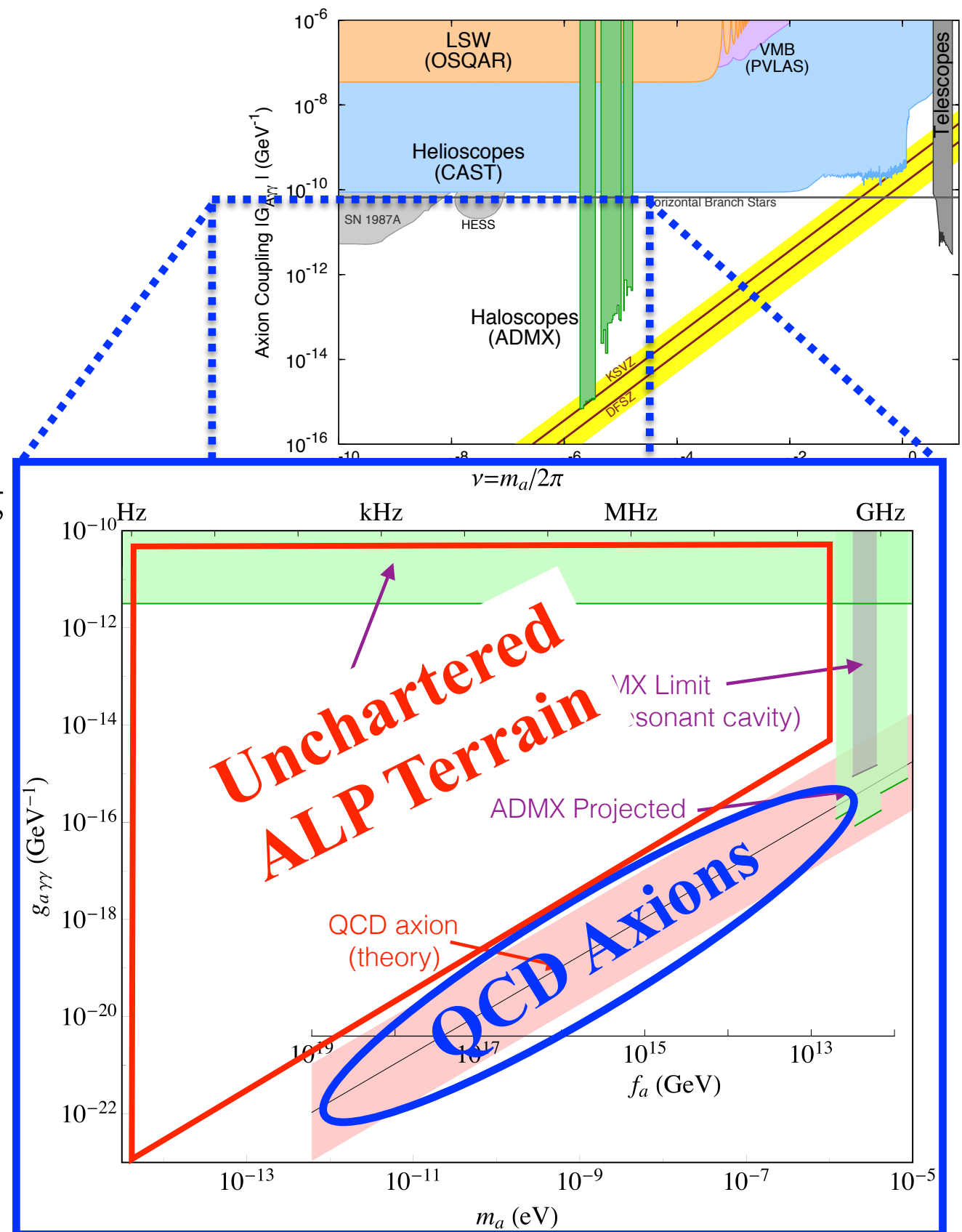
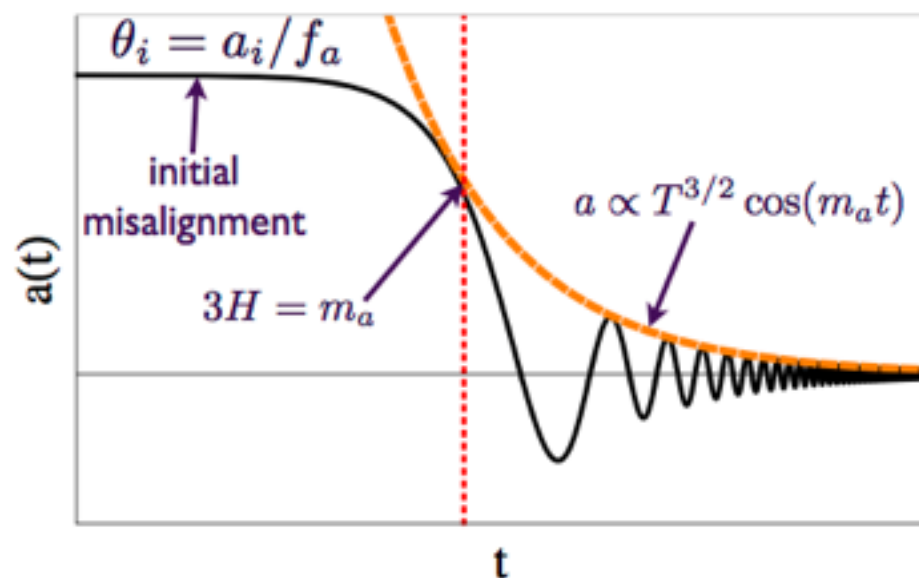
# Axion and ALP Parameter Landscape

- Low mass ( $m_a < 10^{-3}$  eV) Axion DM created through the realignment mechanism
- This parameter space was initially discounted on theoretical grounds
- Models require a pre-inflationary PQ transition, combined with a  $O(1\%)$  tuning on the initial alignment angle.



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# Axion Modifications to EM Lagrangian

New terms in the Lagrangian:

$$\mathcal{L}_{\text{aEM}} = -\frac{1}{4}g_{a\gamma\gamma}aF_{\mu\nu}\tilde{F}^{\mu\nu}$$

Leads to modifications to Maxwell's Equations:

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} - g_{a\gamma\gamma} \left( \mathbf{E} \times \nabla a - \mathbf{B} \frac{\partial a}{\partial t} \right)$$

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 $\lambda \gg \text{Detector Scale}$

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$\lambda \gg \text{Detector Scale}$

In the magneto-quasistatic limit, this behaves like an effective current parallel to the magnetic field

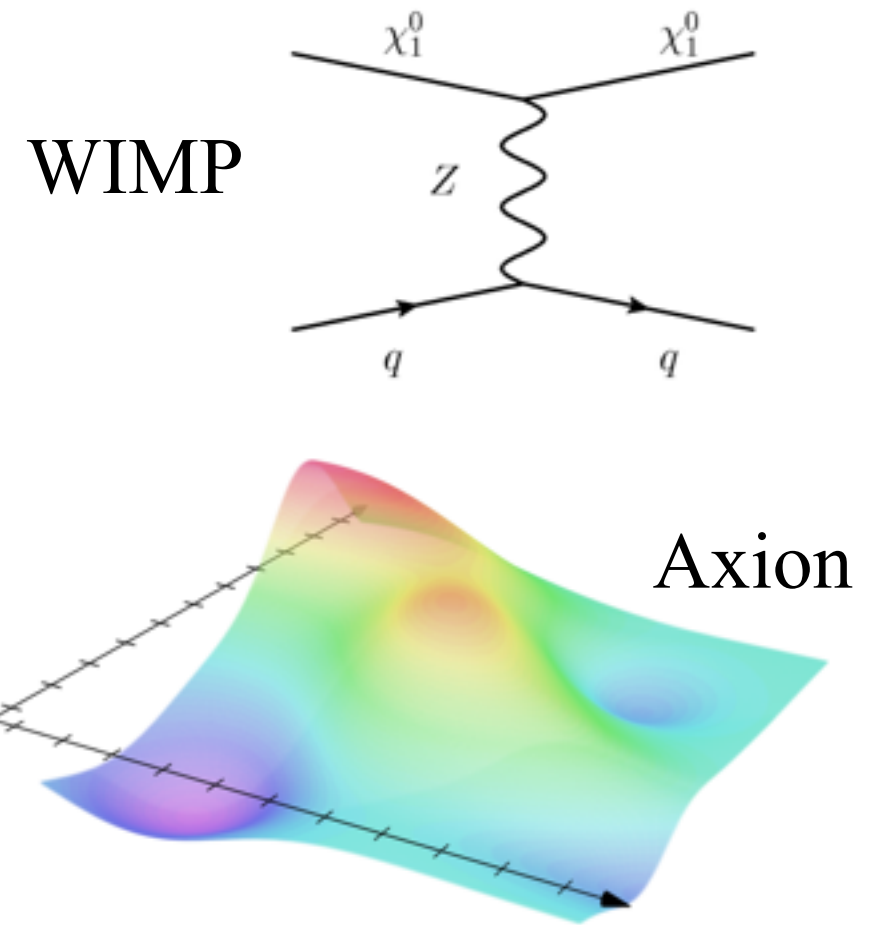
$$\nabla \times \mathbf{B} = g_{a\gamma\gamma} \mathbf{B} \frac{\partial a}{\partial t} \quad \Rightarrow \quad \mathbf{J}_{\text{eff}} \equiv g_{a\gamma\gamma} \mathbf{B} \frac{\partial a}{\partial t}$$



# Axion Modifications to EM

- Rather than being particles (e.g. WIMPs) floating around bumping into detector, they act as a coherent axion field (similar to an E&M field)
- If the local DM density ( $\sim 0.3 \text{ GeV/cm}^3$ ) is dominated by axions, the axion field oscillates in time

$$a(t) = \frac{\sqrt{2\rho_{\text{DM}}}}{m_a} \sin(m_a t)$$



➡ **Constant magnetic field  $\Rightarrow$  effective AC current at a frequency of  $f = m_a/2\pi$ :**

$$\mathbf{J}_{\text{eff}}(t) = g_{a\gamma\gamma} \sqrt{2\rho_{\text{DM}}} \cos(m_a t) \mathbf{B}_0$$

➡ **Signal is very coherent**

$$\frac{\Delta f}{f} \sim v^2 \sim 10^{-6}$$

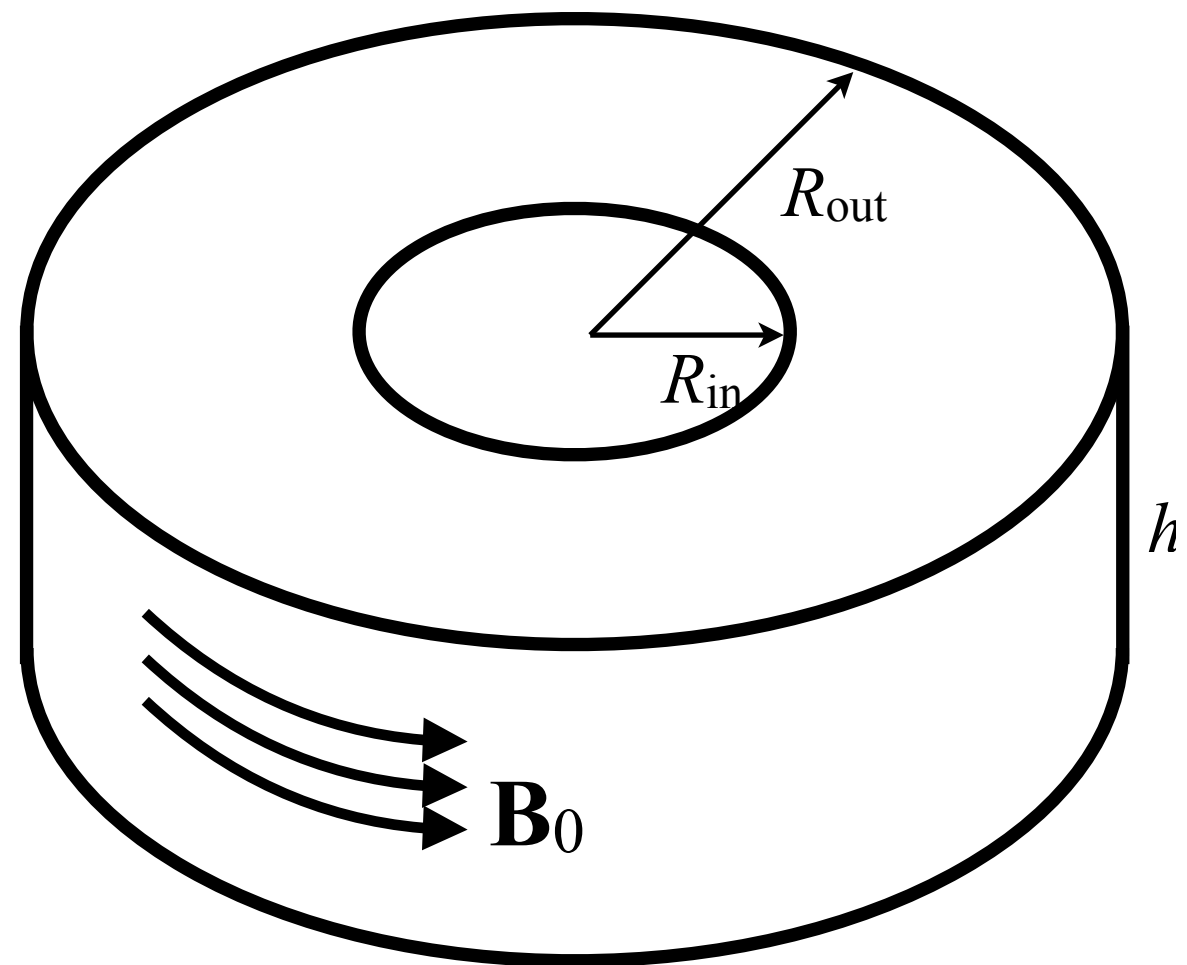
# ABRACADABRA Search Principle

**A** **B**roadband/**R**esonant **A**pproach to **C**osmic **A**xion **D**etection with an  
**A**mplifying **B**-field **R**ing **A**pparatus

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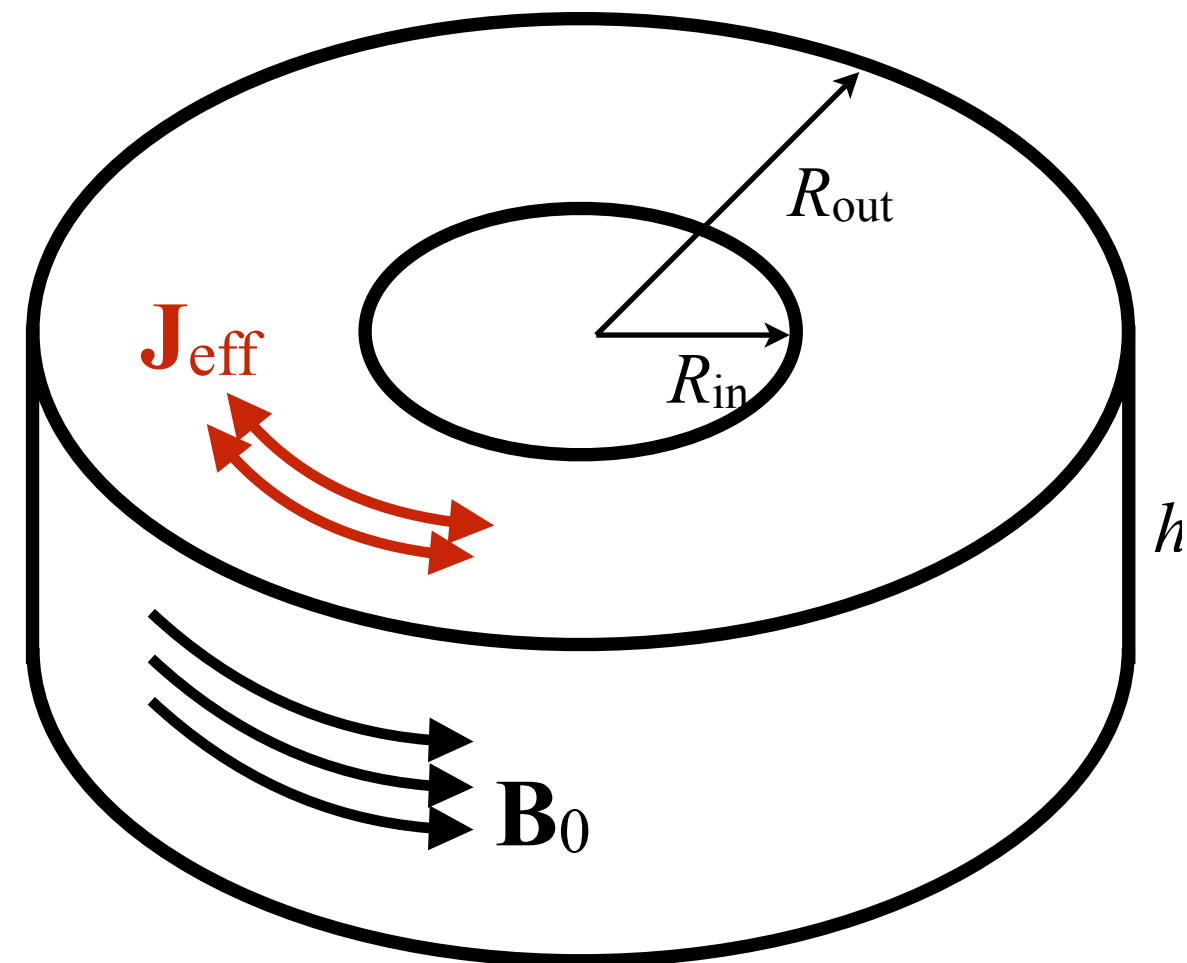




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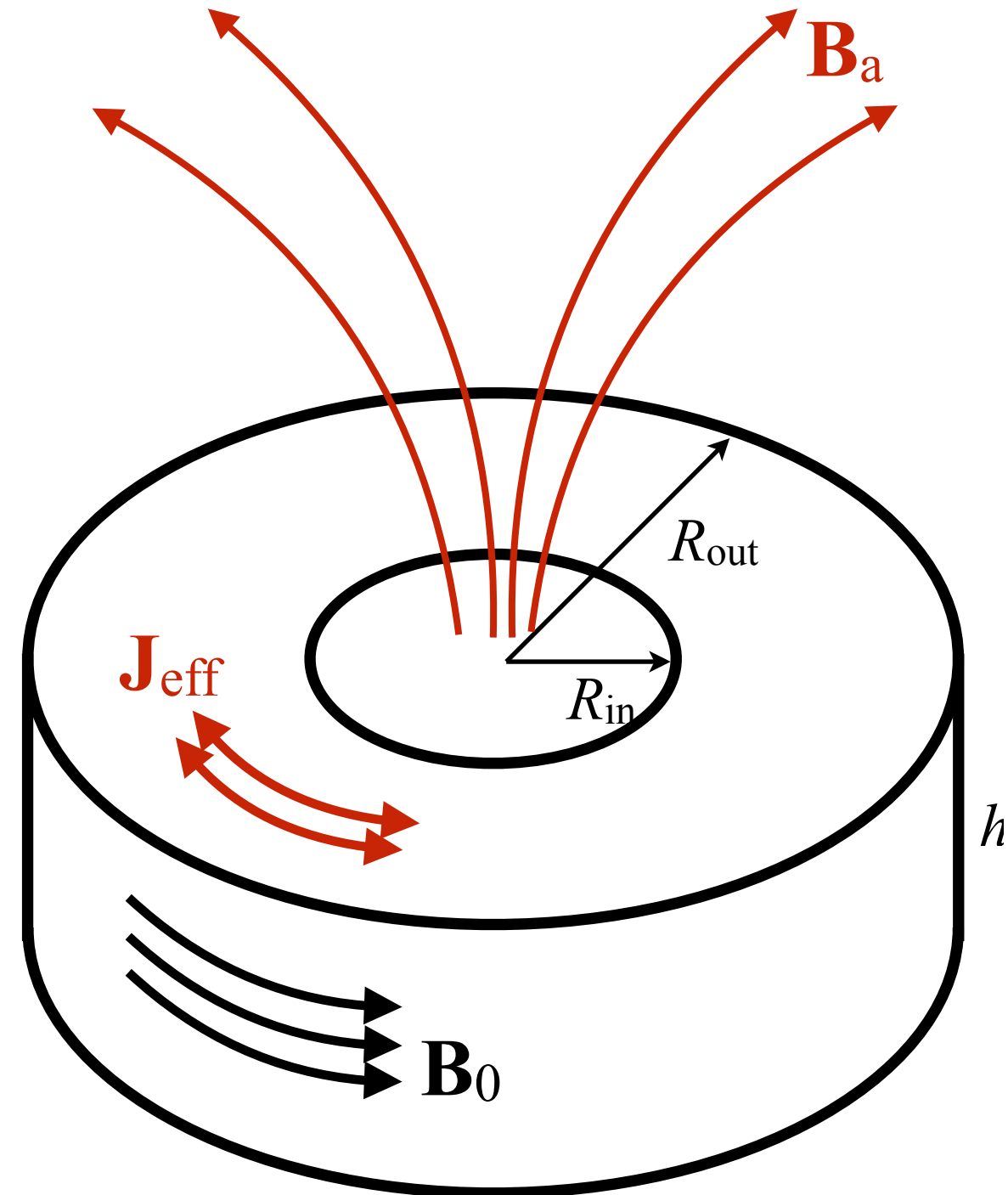
- Start with a toroidal magnet, with a fixed magnetic field,  $\mathbf{B}_0$
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- ... this generates an oscillating magnetic field through the center



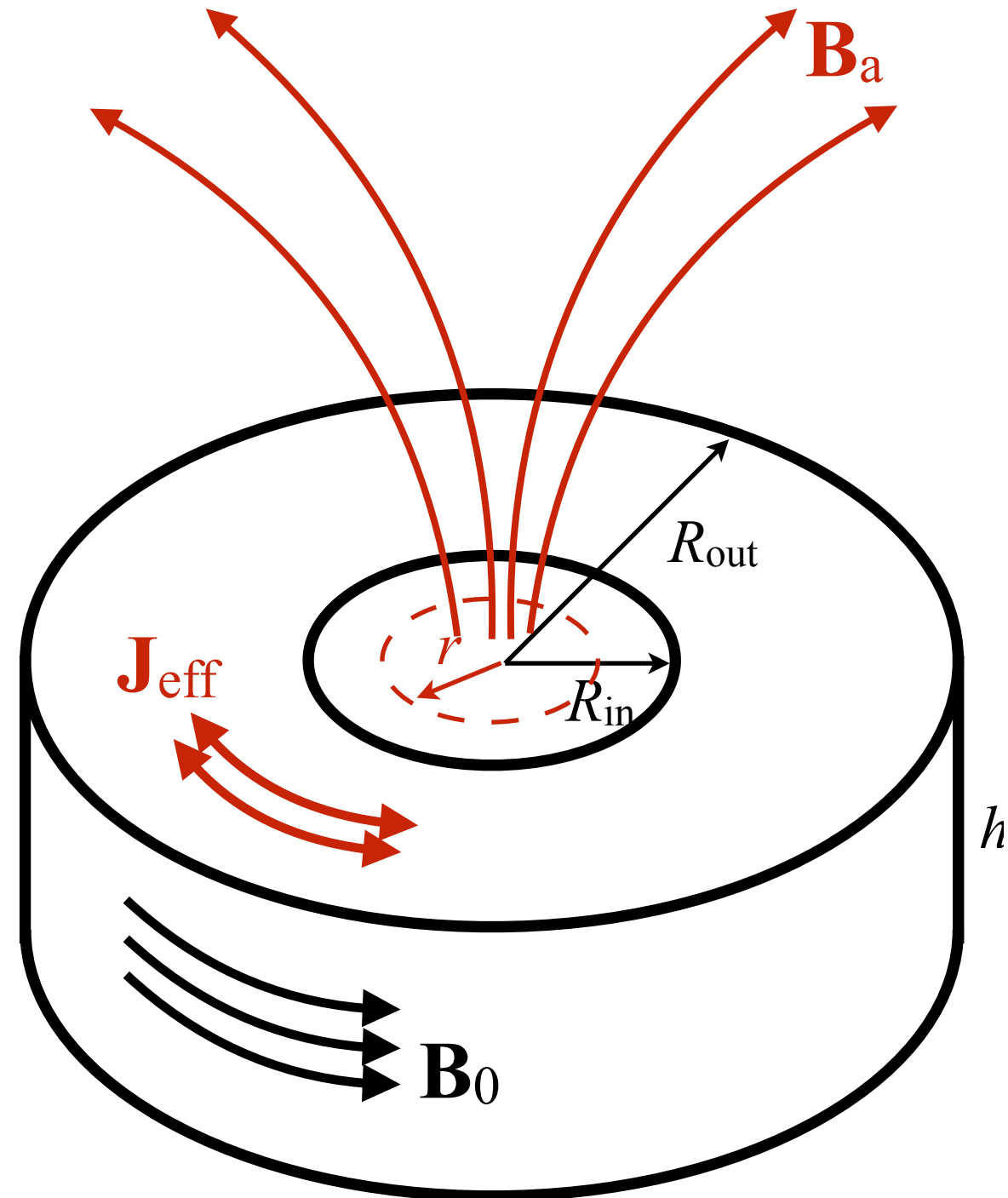
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(MQS approx:  $2\pi/m_a \gg R_{\text{in}}, R_{\text{out}}, h$ )
- ... this generates an oscillating magnetic field through the center
- Insert a pickup loop in the center of the toroid to detect the oscillating magnetic field

$$\Phi_a(t) = g_{a\gamma\gamma} B_{\text{max}} \sqrt{2\rho_{\text{DM}}} \cos(m_a t) \mathcal{G}_V V$$

$\mathcal{G}_V$ : Geometric factor, depends on magnet geometry

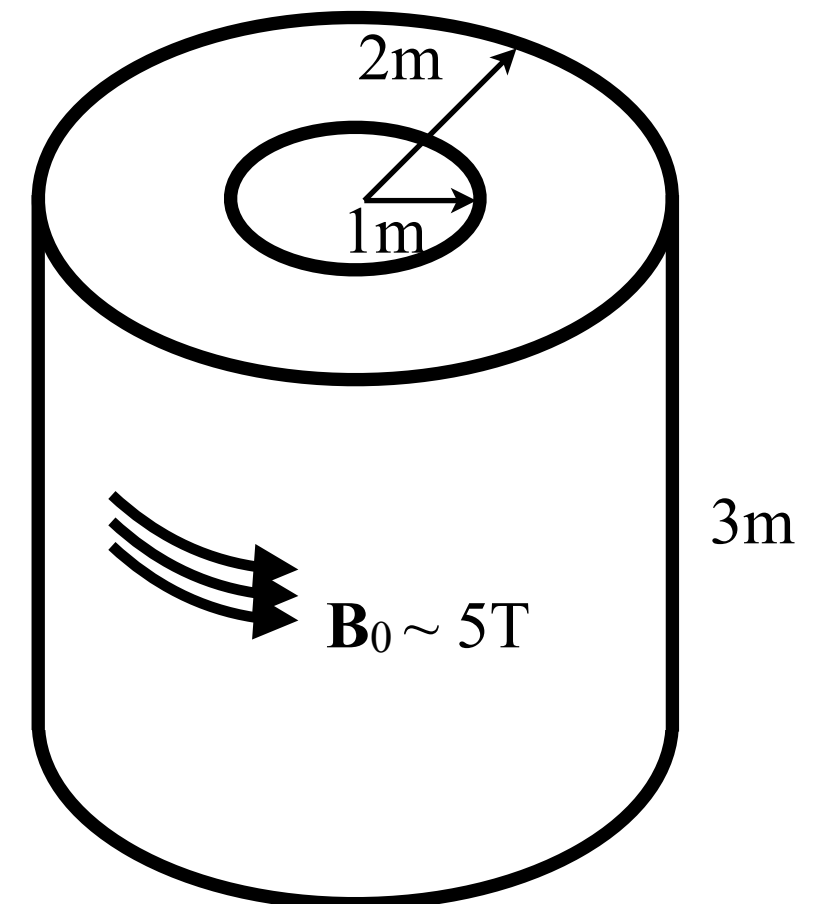
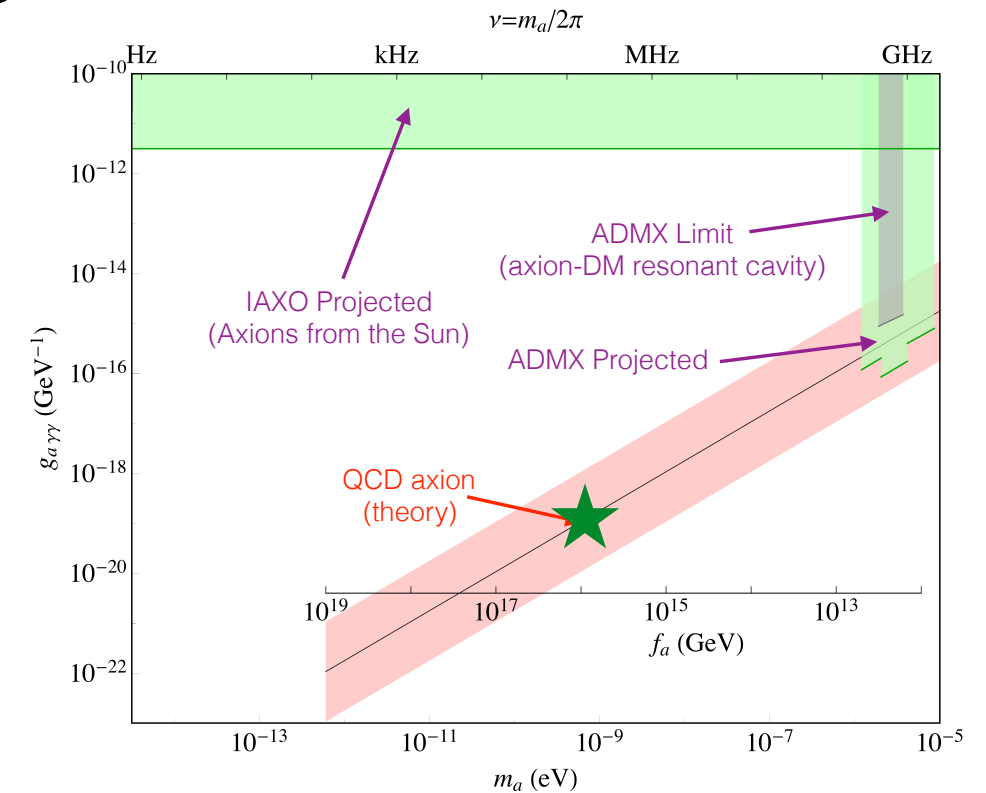




# The Challenge

To get an idea of the size of the effect:

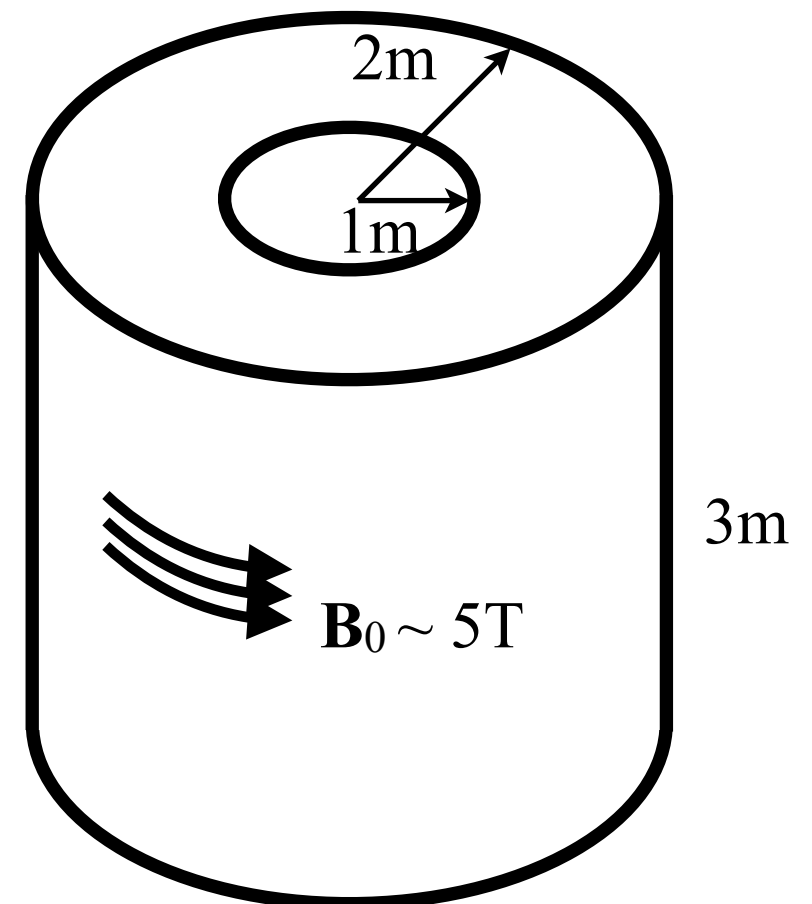
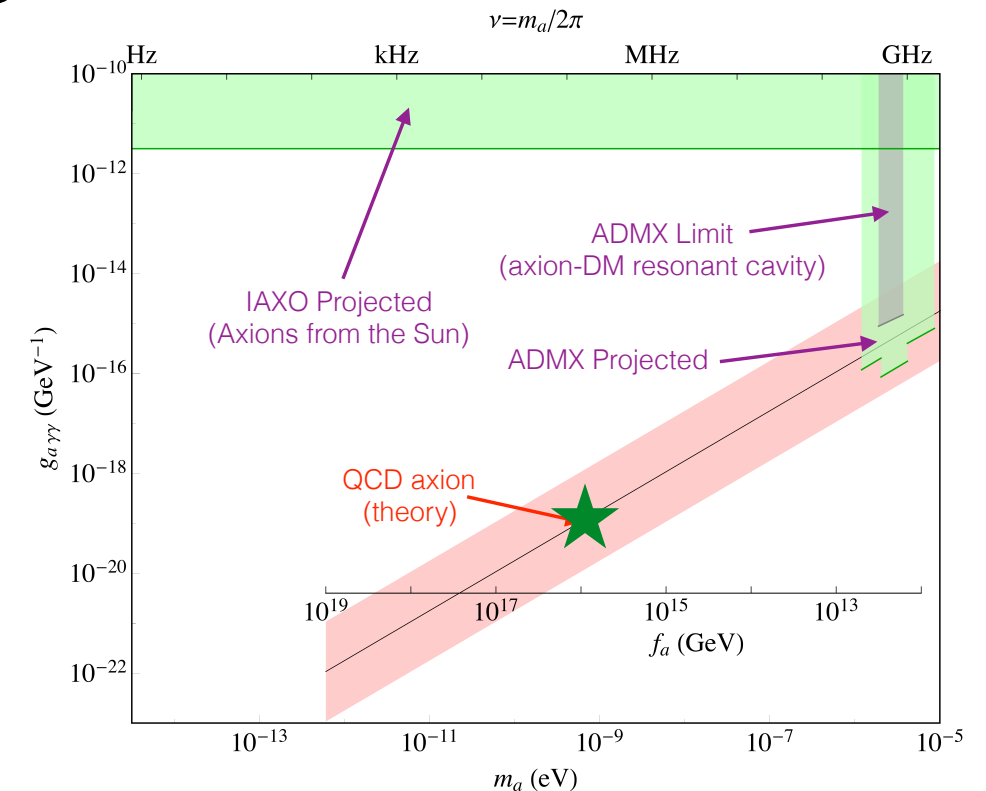
- Take as an example, the geometry  
 $r = R_{\text{in}} = R_{\text{out}}/2 = h/3 = 1 \text{ m}$ ,  $B_{\text{max}} = 5\text{T}$
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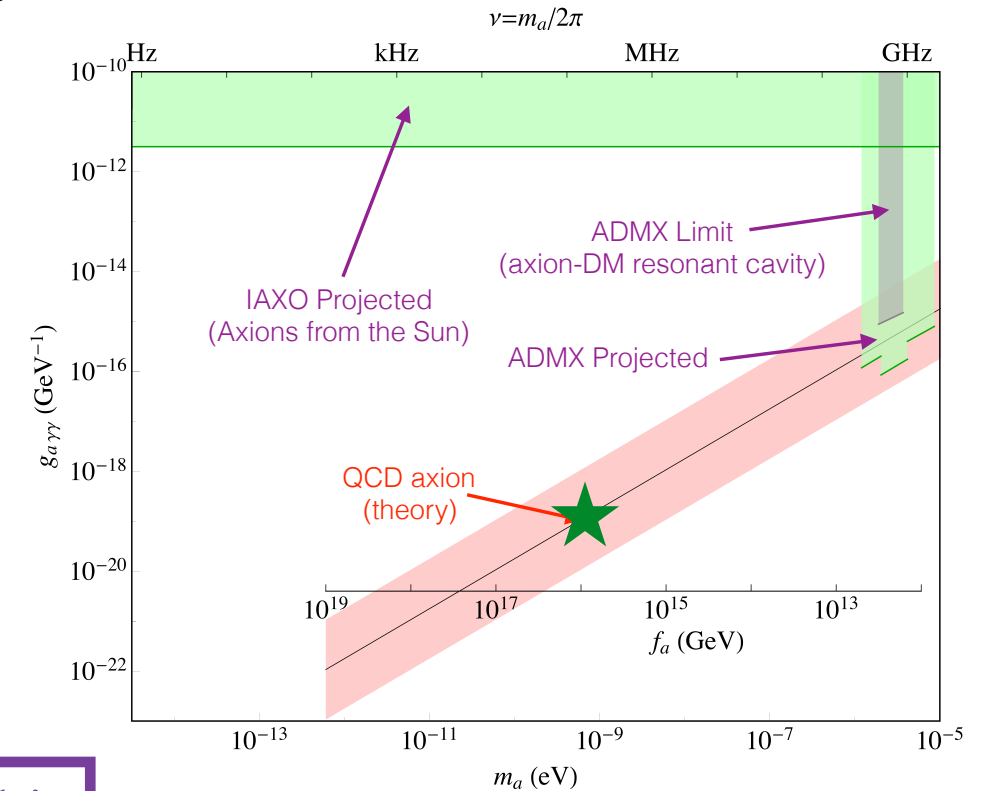
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 ➔  $B \sim 5 \times 10^{-22} \text{ T}$  @  $f = 240 \text{ kHz}$ ,  
 $\Delta f \sim 240 \text{ mHz}$



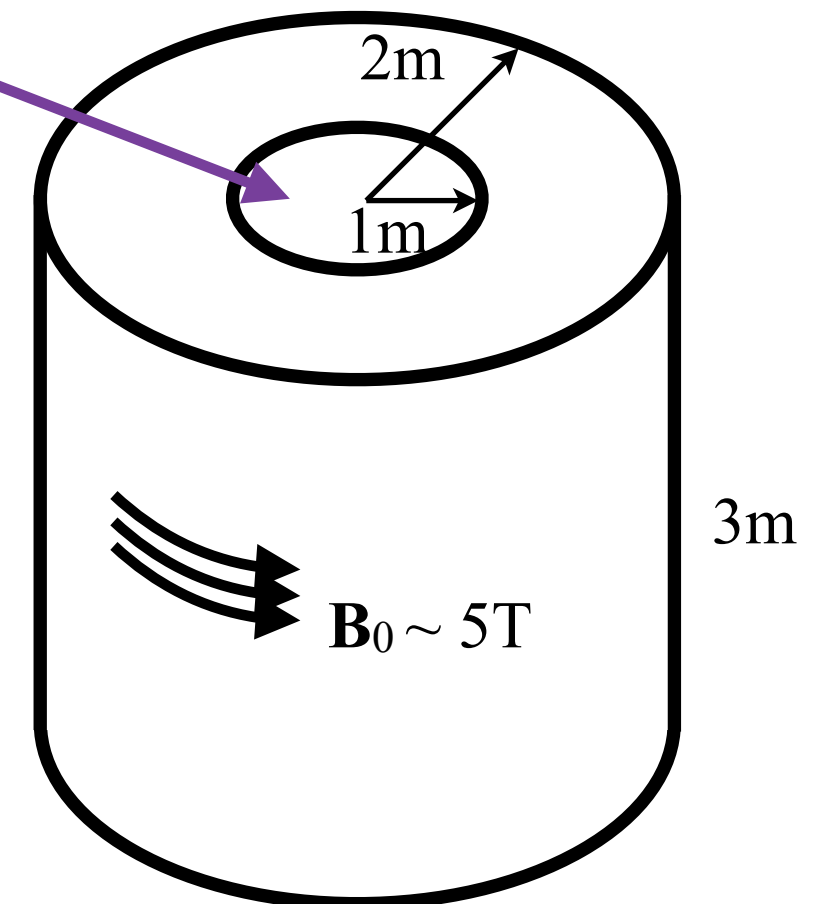
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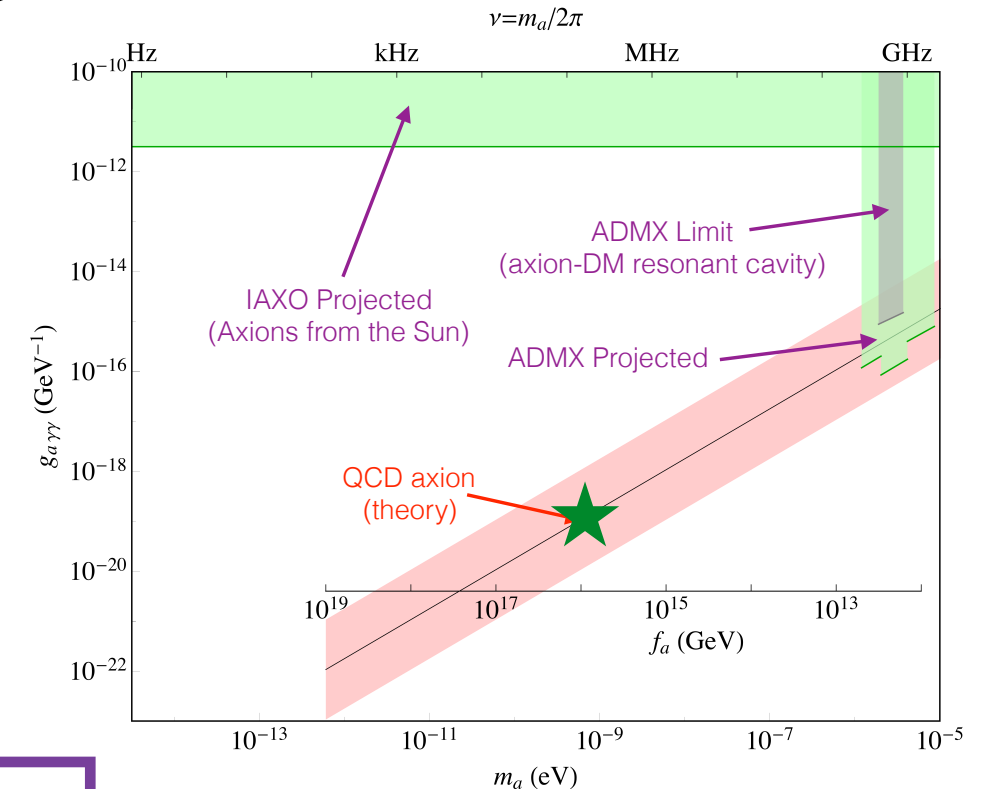
In the absence of  $a(t)$ , this  
is a zero field region!



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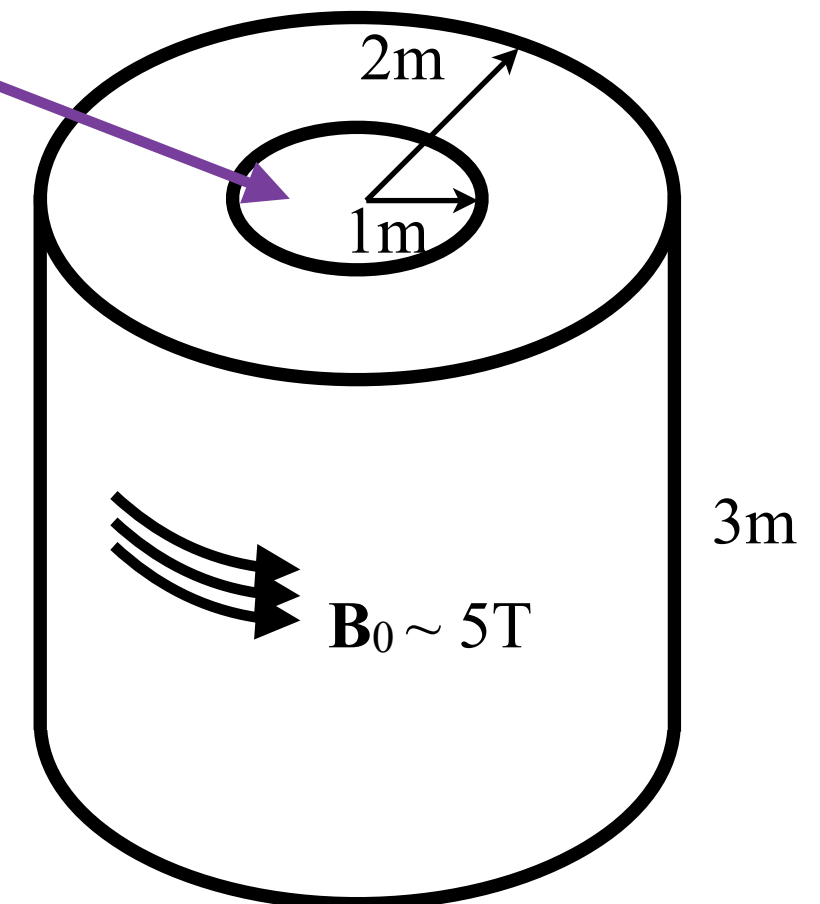
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 $\Delta f \sim 240 \text{ mHz}$



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Example from MRI:

- Demonstrated sensitivity  $S_B \sim 10^{-17} \text{ T/Hz}^{1/2}$  in SQUID  $R \sim 3.3 \text{ cm}$
- Scaling to  $R = 1 \text{ m}$ ,  $S_B \sim 10^{-20} \text{ T/Hz}^{1/2}$
- Access to QCD axion scale after  $\sim 2$  year of scan (when integration time  $\gg 1/\Delta f$ ):  
 $g_{a\gamma\gamma} \sim S_B(t/\Delta f)^{1/4}$

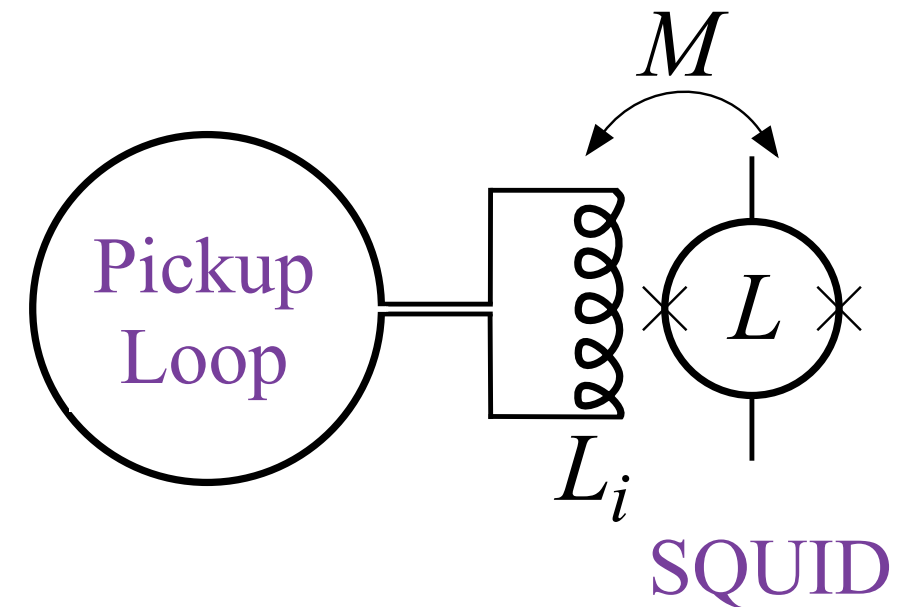


# Two Readout Approaches

Exploring two SQUID based readout approaches:

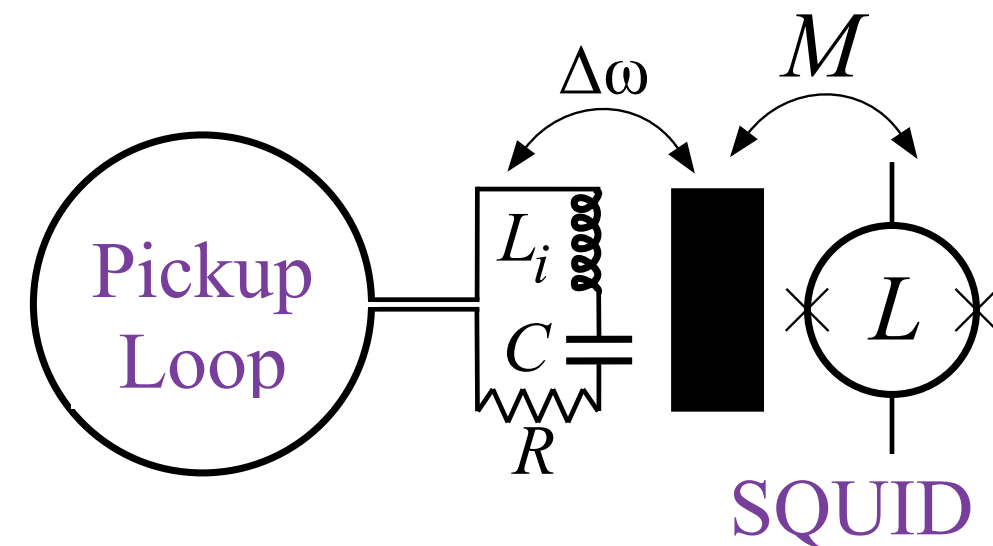
► **A broadband approach**

- Scan the full frequency range in “one” measurement
- No resonance amplification



► **A resonance amplification approach**

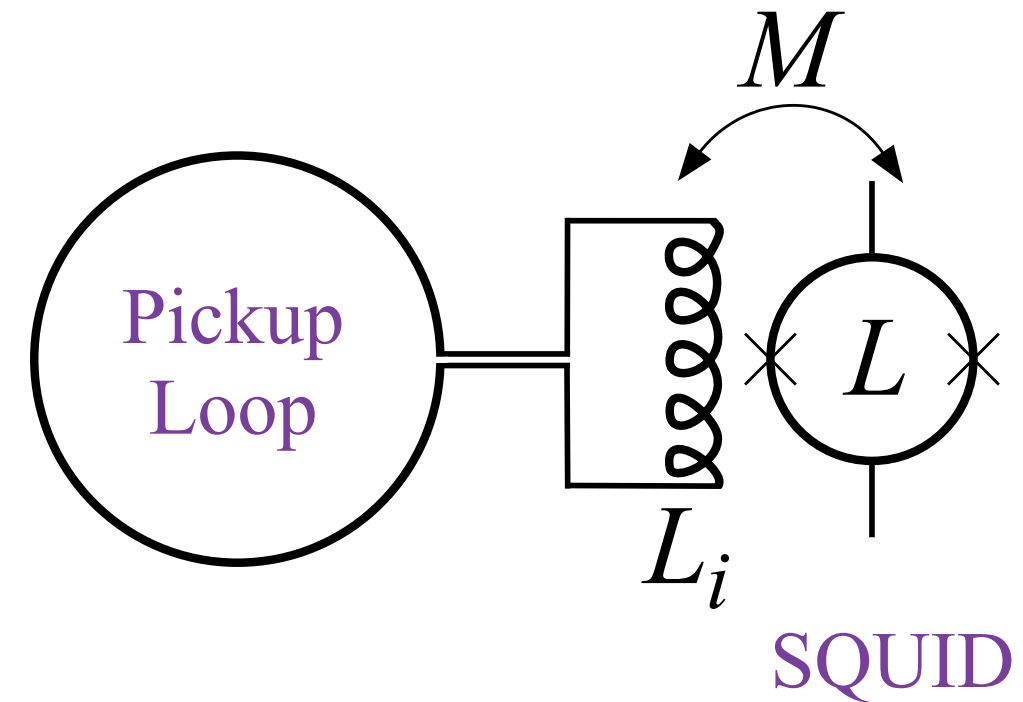
- Resonance circuit can enhance signal by  $Q_0 \approx 10^6$
- “Narrow” but tunable frequency line slowly scans through the full frequency range





# Broadband Axion Search

- SQUID readout systems can measure very small AC currents coming out of the pickup loop
- Over  $\sim 5 - 50$  Hz, the primary noise source is flux noise from the SQUID
- Typical noise level is
 
$$S_{\Phi,0}^{1/2} \sim 10^{-6} \Phi_0 / \sqrt{\text{Hz}}$$
- Below  $\sim 5 - 50$  Hz,  $1/f$  noise takes over and becomes dominant



- Cannot resolve thermal noise floor, so somewhat insensitive to temperature

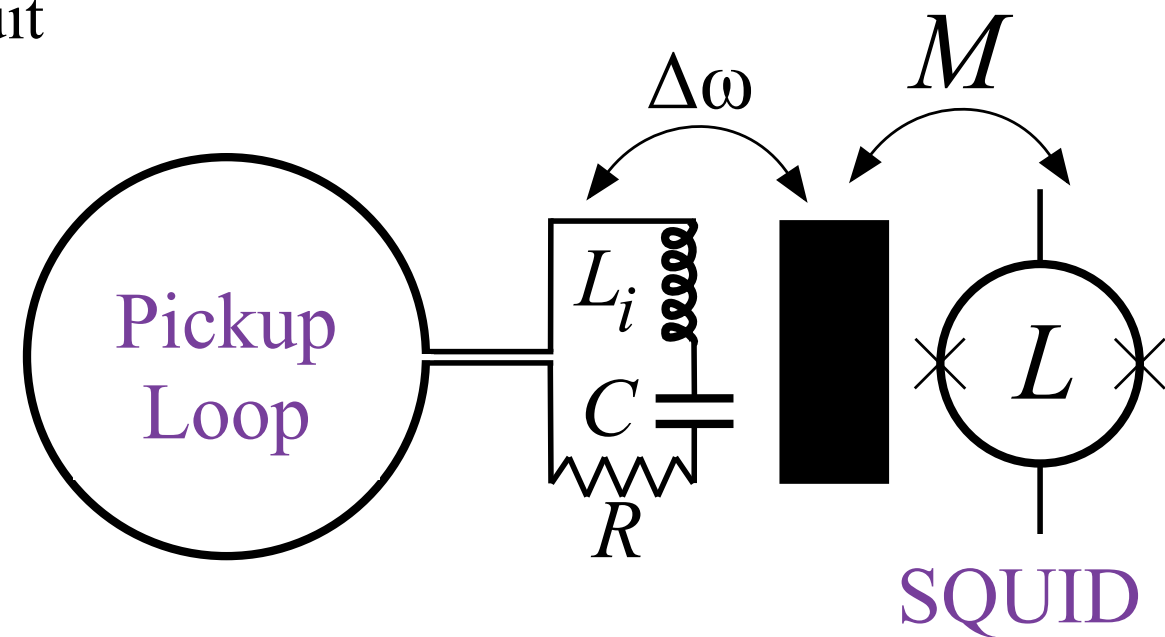
$$t \gg \frac{1}{\Delta\nu}$$

- Broadband only sensitivity:

$$g_{a\gamma\gamma} \propto \left( \frac{m_a}{t} \right)^{\frac{1}{4}} \frac{1}{B_{\text{max}}} \frac{1}{\mathcal{G}_V V} \frac{1}{\sqrt{\rho_{\text{DM}}}} S_{\Phi,0}^{1/2}$$

# A Resonance Enhanced Search

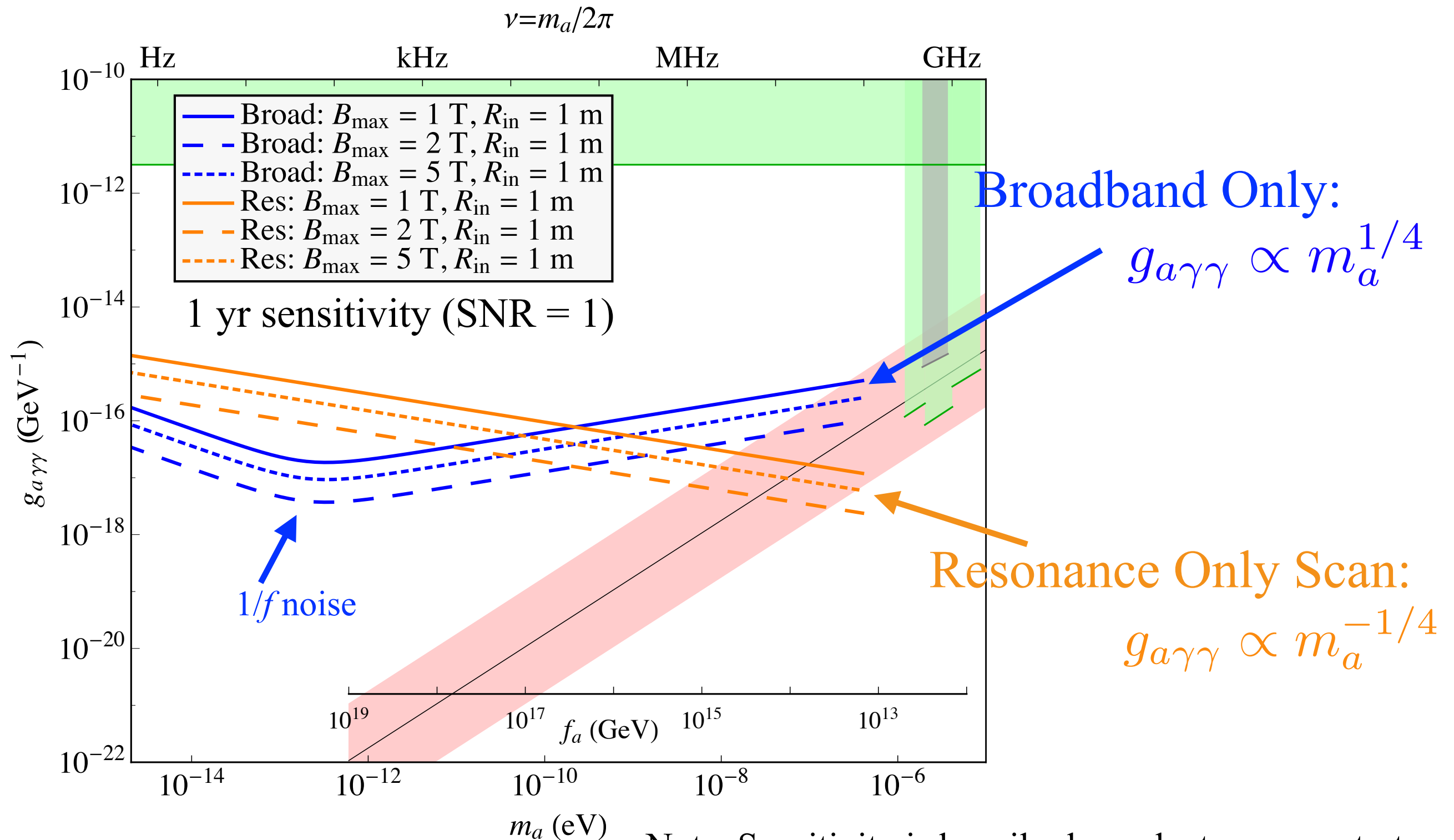
- We can enhance the signal in a narrow frequency window by adding a capacitor to create a RLC circuit
- This can amplify the signal by  $Q_0$ , in a narrow bandwidth,  $\Delta\omega = \omega/Q_0$
- Resonators with  $Q_0 \sim 10^6$  are achievable
- Dominant noise source becomes dissipation in the oscillator circuit  $\Rightarrow$  temperature dependent
- Benchmark at operating temperature of 100mK
- Naive resonance scan sensitivity:



$$t \gg \frac{1}{\Delta\nu}$$

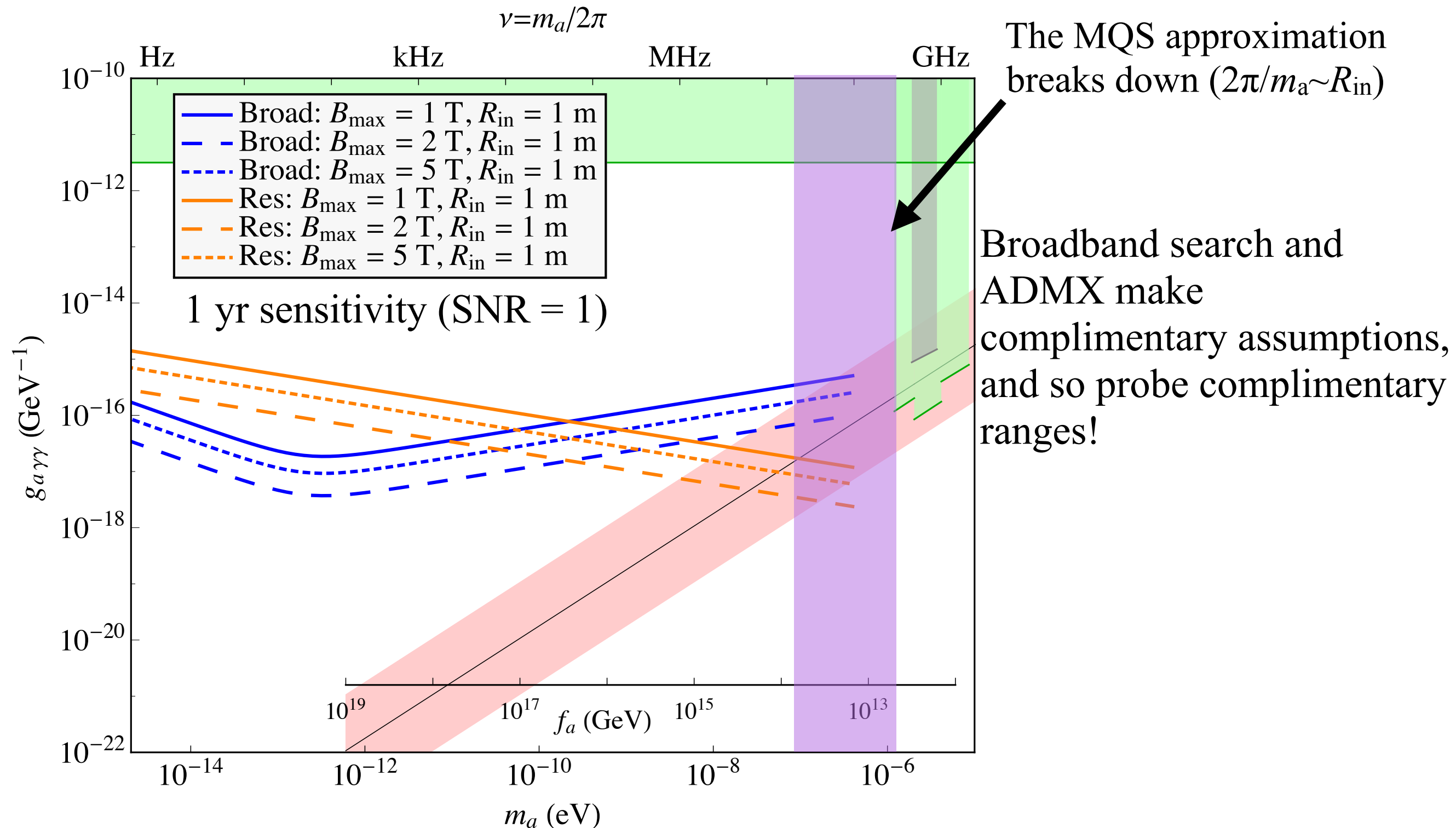
$$g_{a\gamma\gamma} \propto \sqrt{L_T} \left( \frac{1}{m_a t_{\text{scan}}} \right)^{\frac{1}{4}} \frac{1}{B_{\text{max}}} \frac{1}{\mathcal{G}_V V} \sqrt{\frac{1}{\rho_{\text{DM}}} \frac{k_B T}{Q_0}}$$

# Broadband Axion Search Sensitivity



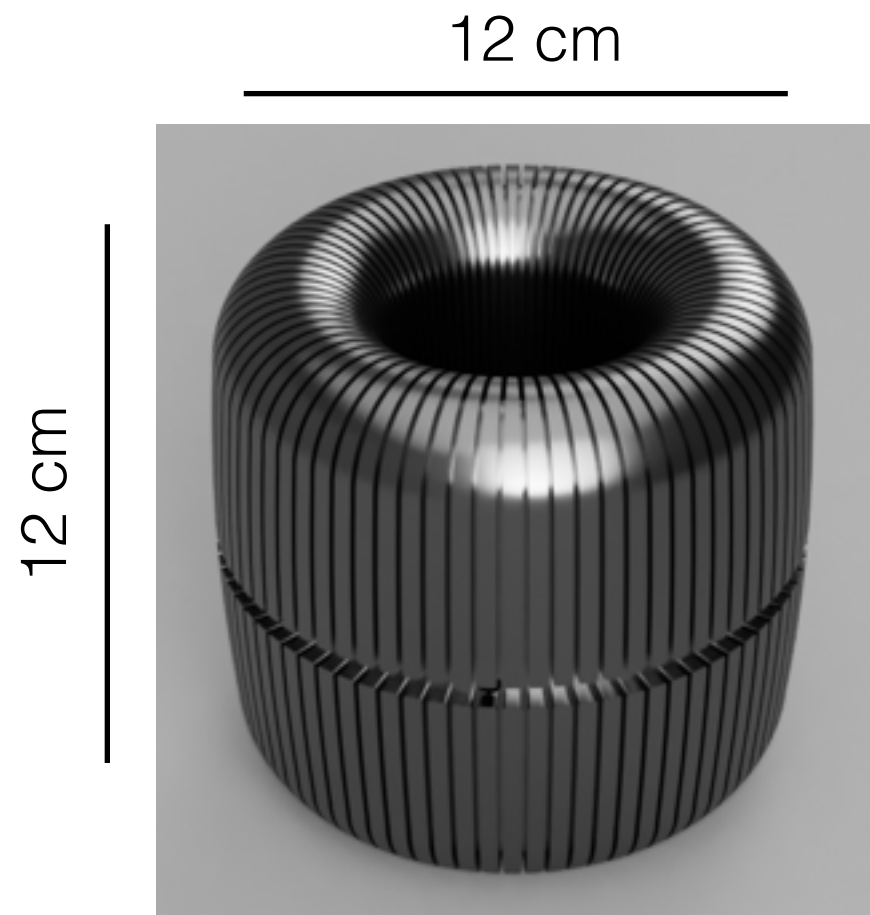
Note: Sensitivity is heavily dependent on scan strategy!  
 These curves are not optimal. In discussions with K. Irwin's group about a more optimal scan strategy.

# Broadband Axion Search Sensitivity

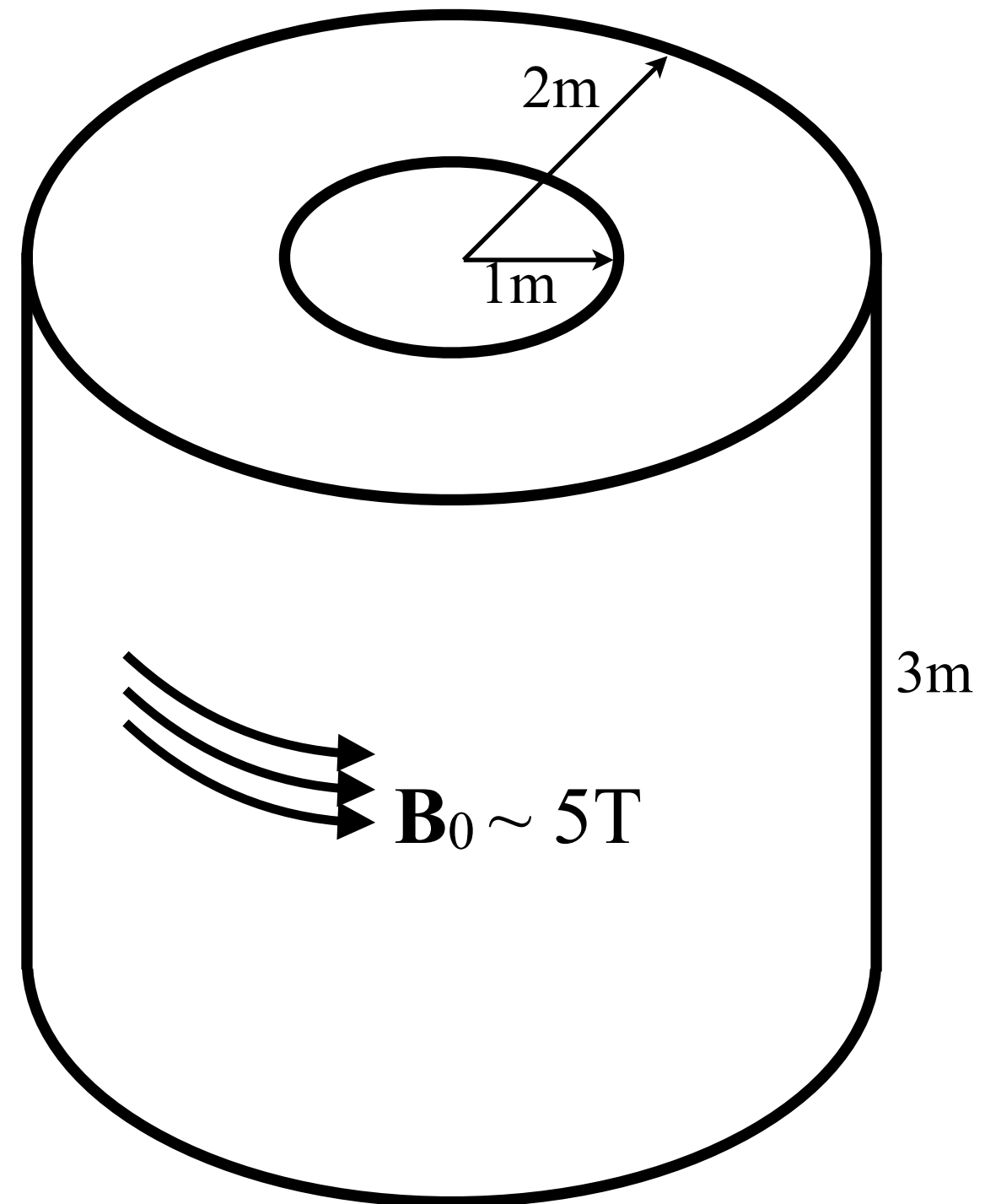
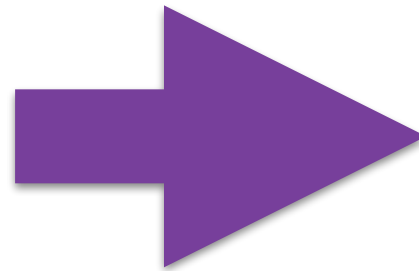


We have some interesting calculations to do to understand what happens in this regime

# ABRACADABRA Program



ABRACADABRA 10cm



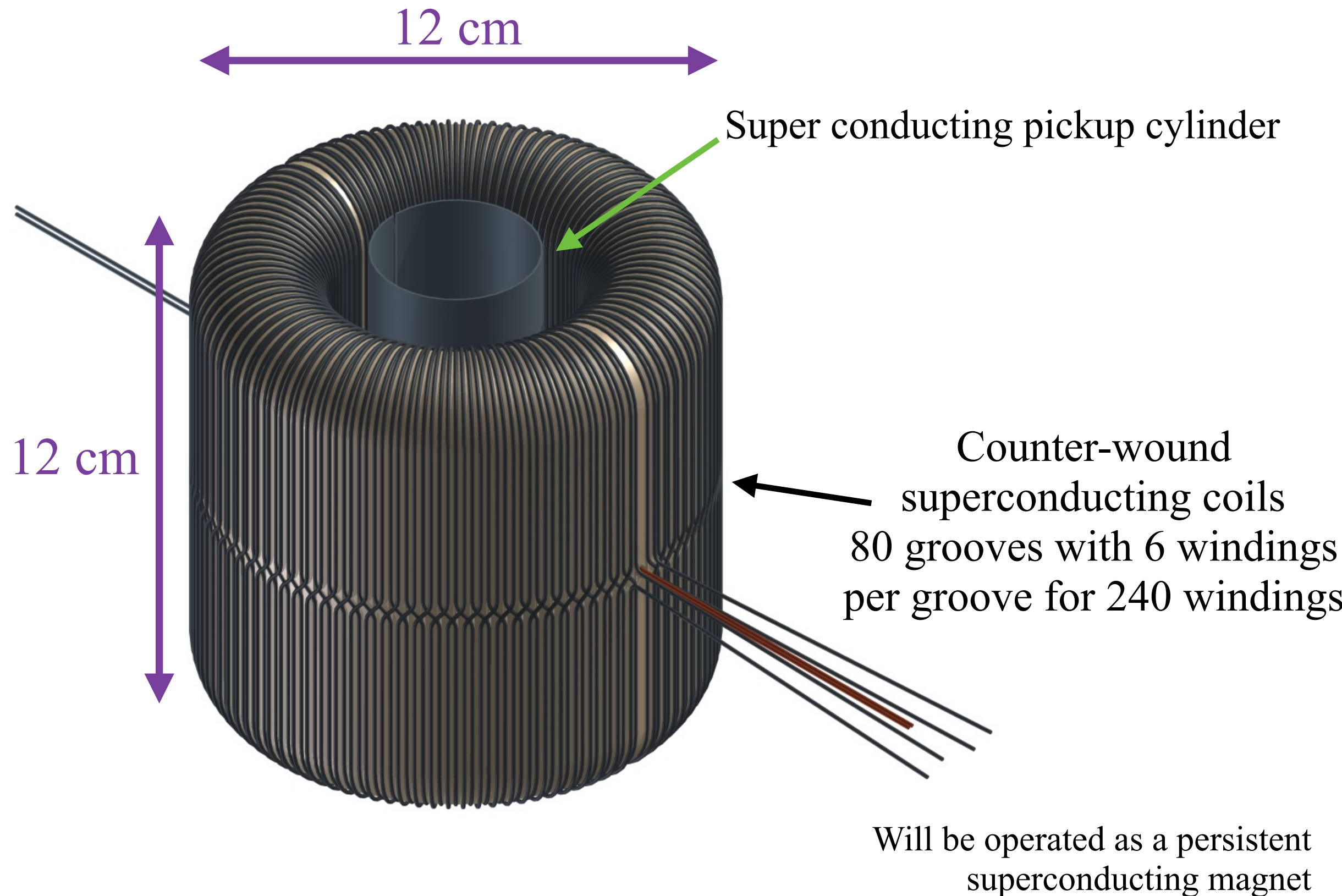
ABRACADABRA



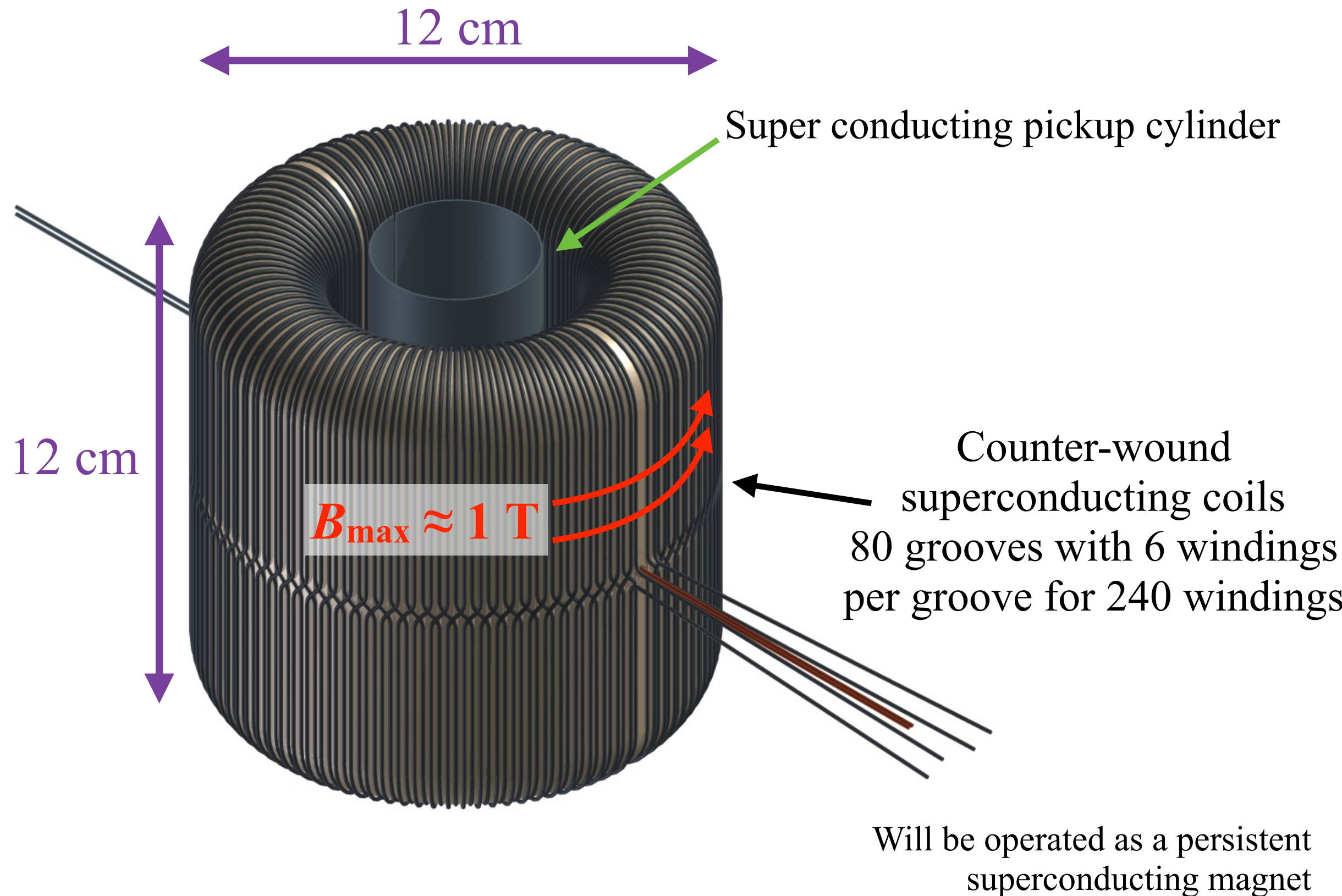
# ABRACADABRA-10 cm

- We are building a 10 cm scale prototype version at MIT
- **Dimensions:**  $R_{\text{in}} = 3 \text{ cm}$ ,  $R_{\text{out}} = 6 \text{ cm}$ ,  $h = 12 \text{ cm}$ .  $G_V \sim 5\%$ .  $B_{\text{max}} = 1 \text{ T}$ .
- **People:** Janet Conrad, Joe Formaggio, Sarah Heine, Reyco Henning (UNC), Yoni Kahn (Princeton), Joe Minervini, Jonathan Ouellet, Kerstin Perez, Alexey Radovinsky, Ben Safdi, Jesse Thaler, Daniel Winklehner, Lindley Winslow
- Funded by an \$80k NSF grant
  - Expected magnet delivery in May
  - Hope to have first results before the end of 2017, though we will spend quite a bit of time investigating different data taking configurations.

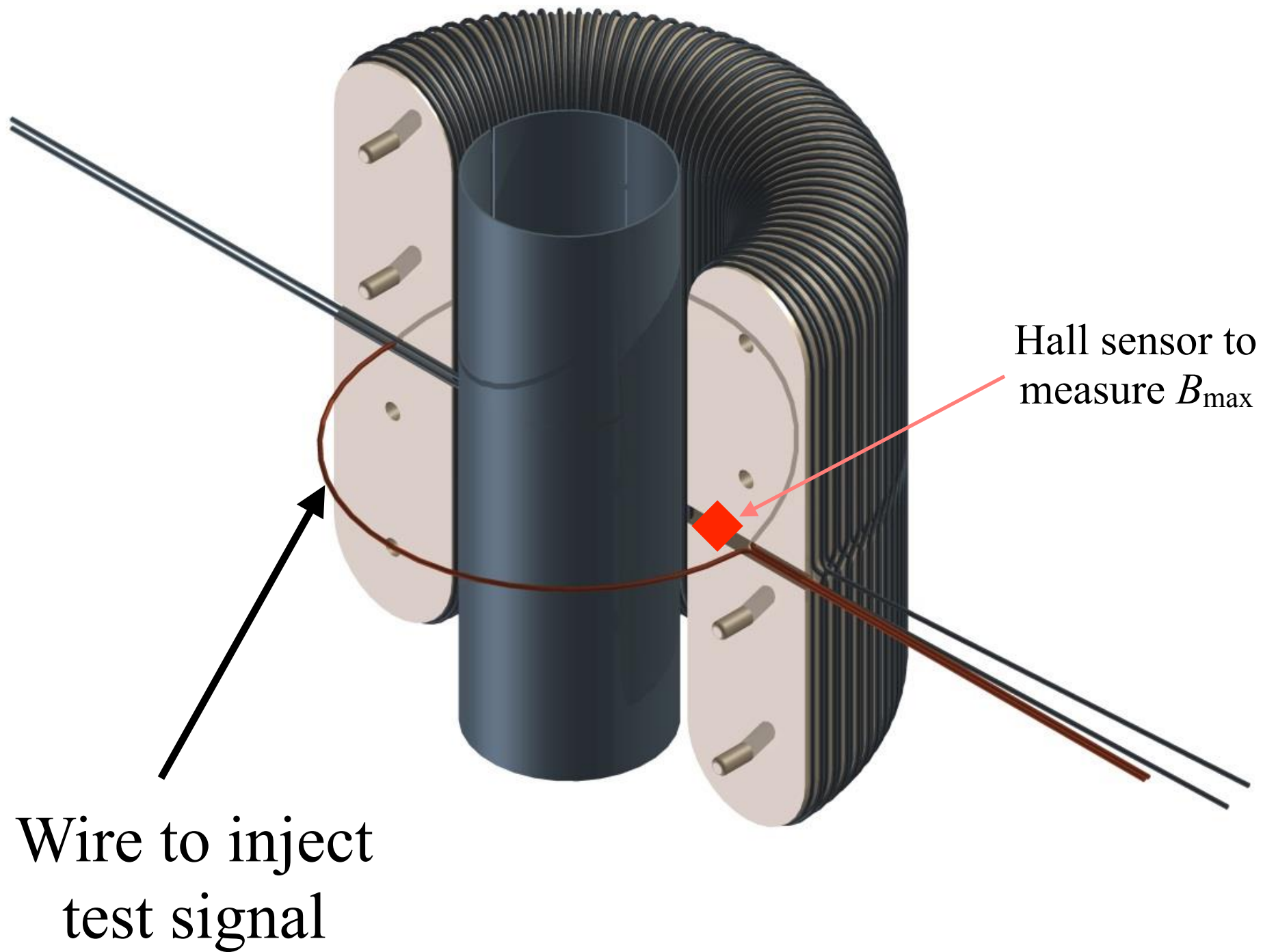
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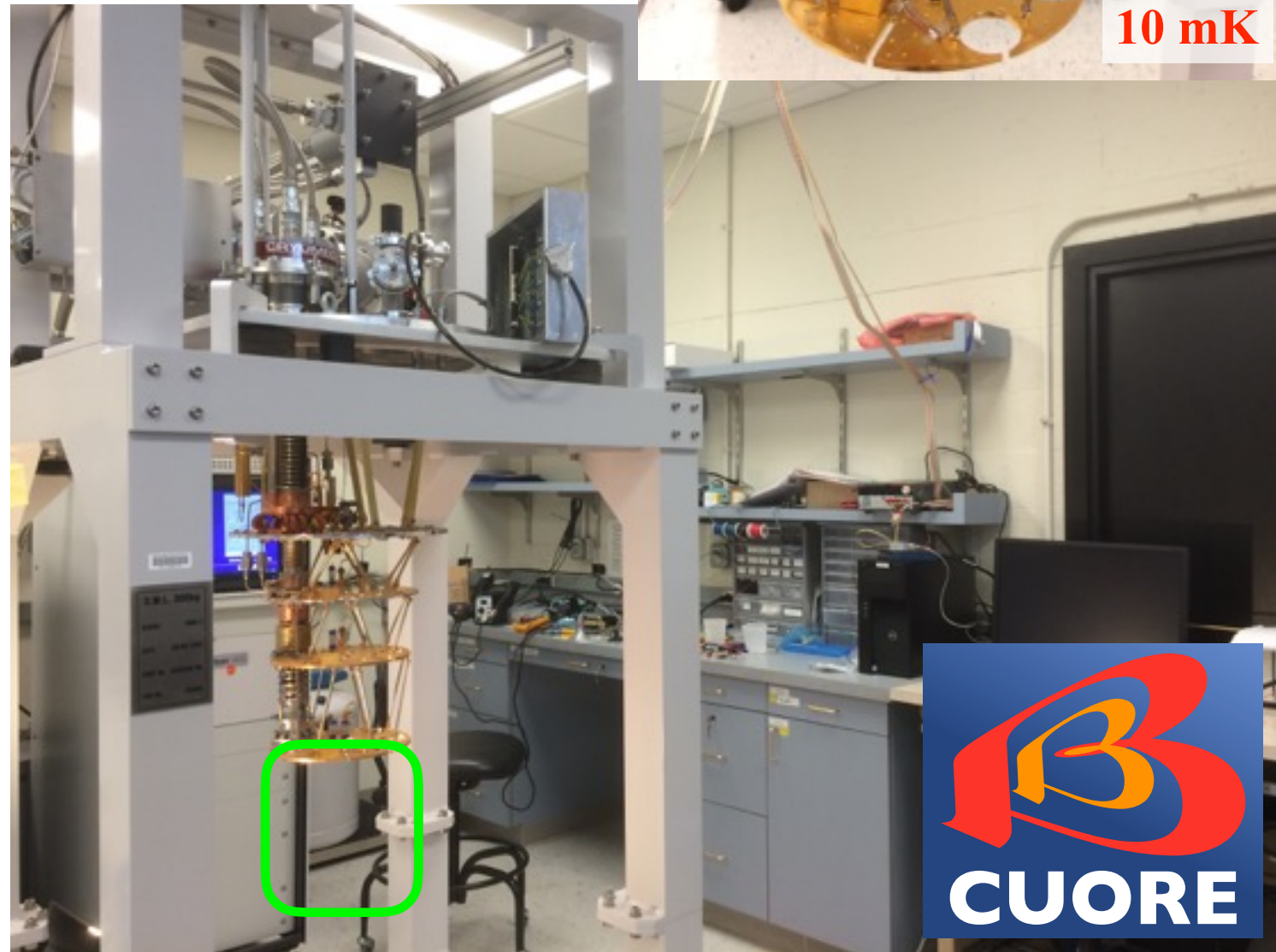
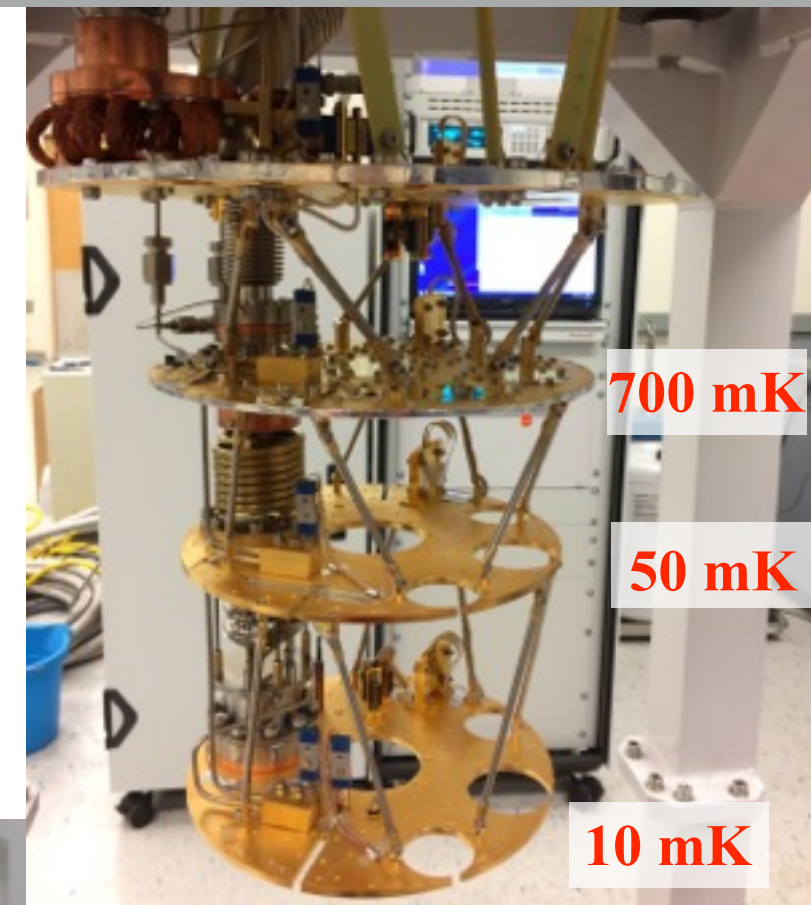
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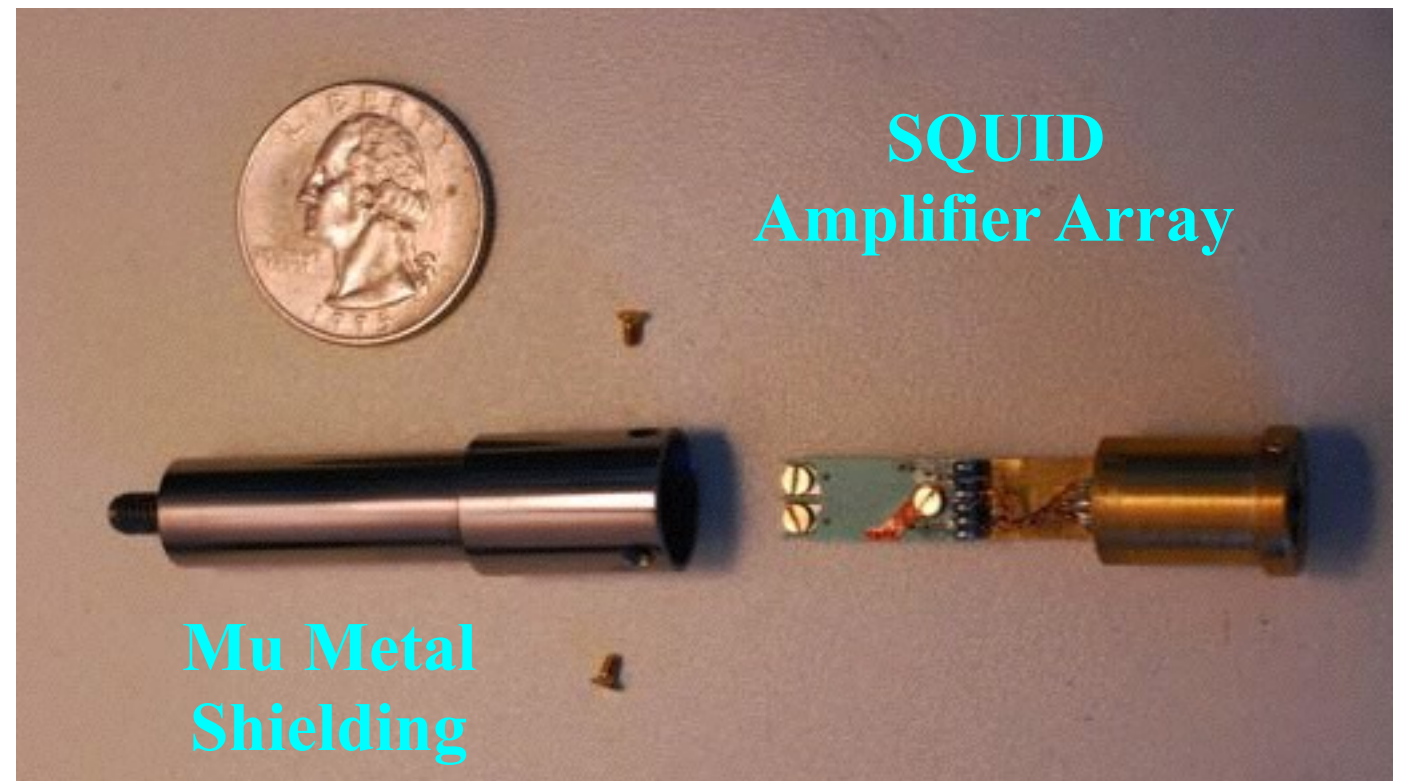
# Lab At MIT

- We have an Oxford Instruments Triton 400 dilution refrigerator
- Normally used for R&D for CUORE/CUPID  $0\nu\beta\beta$
- Capable of a base temperature of  $<10$  mK
- Working volume of  $\sim 12$  L
  - $\sim 25$  cm diameter by  $\sim 24$  cm height
- Cryogen free, so does not require He refilling
- Can run  $\sim 2$  weeks unattended



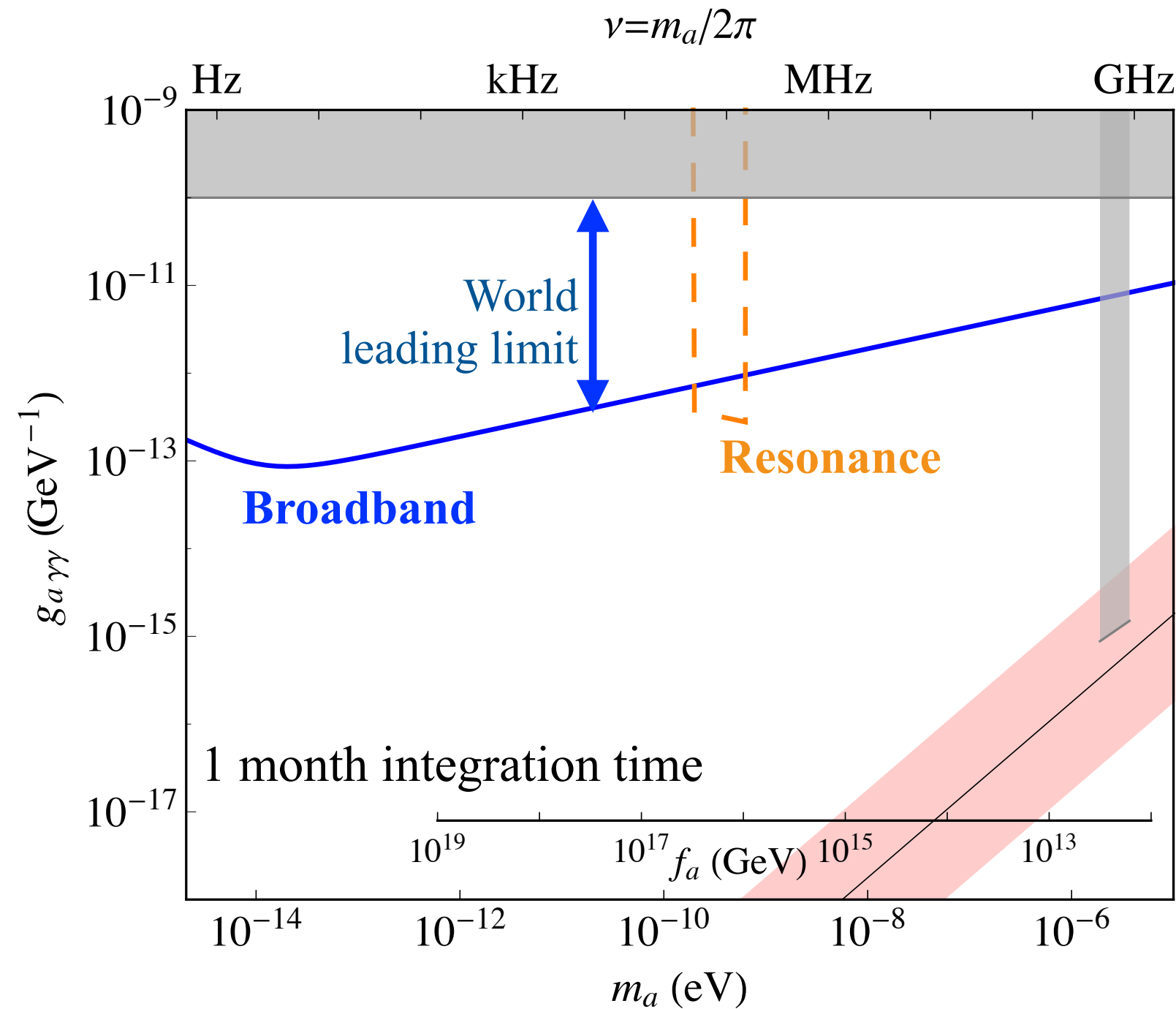
# SQUIDs and SQUID Arrays

- We currently have a set of Magnicon SQUID amplifier arrays
- These will act as a second stage amplification to a first stage SQUID
- We are looking to purchase a Magnicon SQUID current sensor
- Typical noise @ 4K:  $1.2 \times 10^{-6} \Phi_0/(\text{Hz})^{1/2}$ , with  $1/f$  corner at 3 Hz
- Bandwidth of 6MHz
- Might be able to do slightly better with SQUID amplifiers or at lower temperatures.





# ABRACADABRA-10 cm



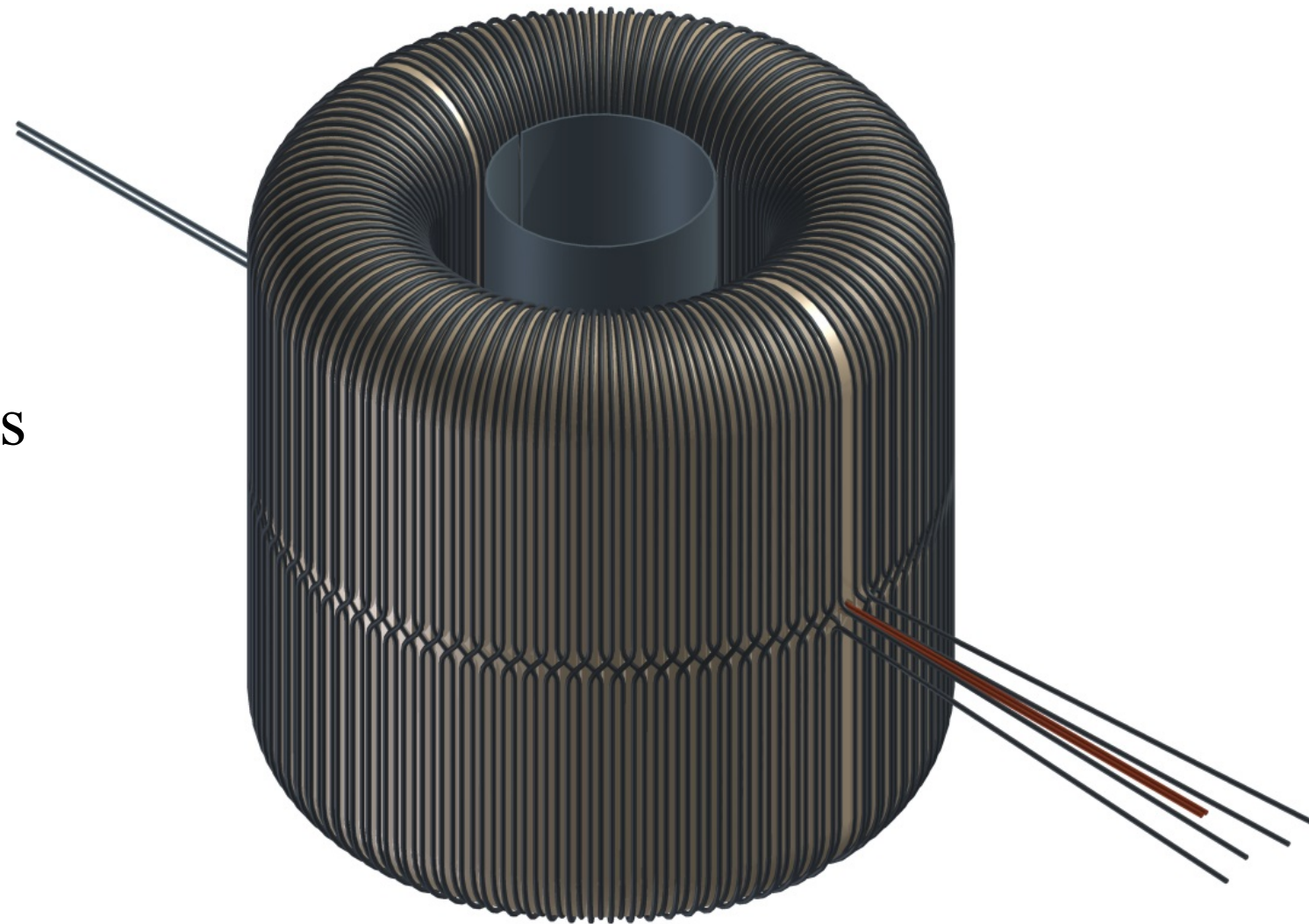
*World leading limit after 1 month of data!*

- ▶ SNR=1
- ▶ Assumes flux limited background
- ▶  $1/f$  corner between  $\sim 3\text{Hz}$  —  $50\text{Hz}$
- ▶ Bandwidth of  $\sim 6\text{MHz}$

# ABRACADABRA 10cm Prototype Goals

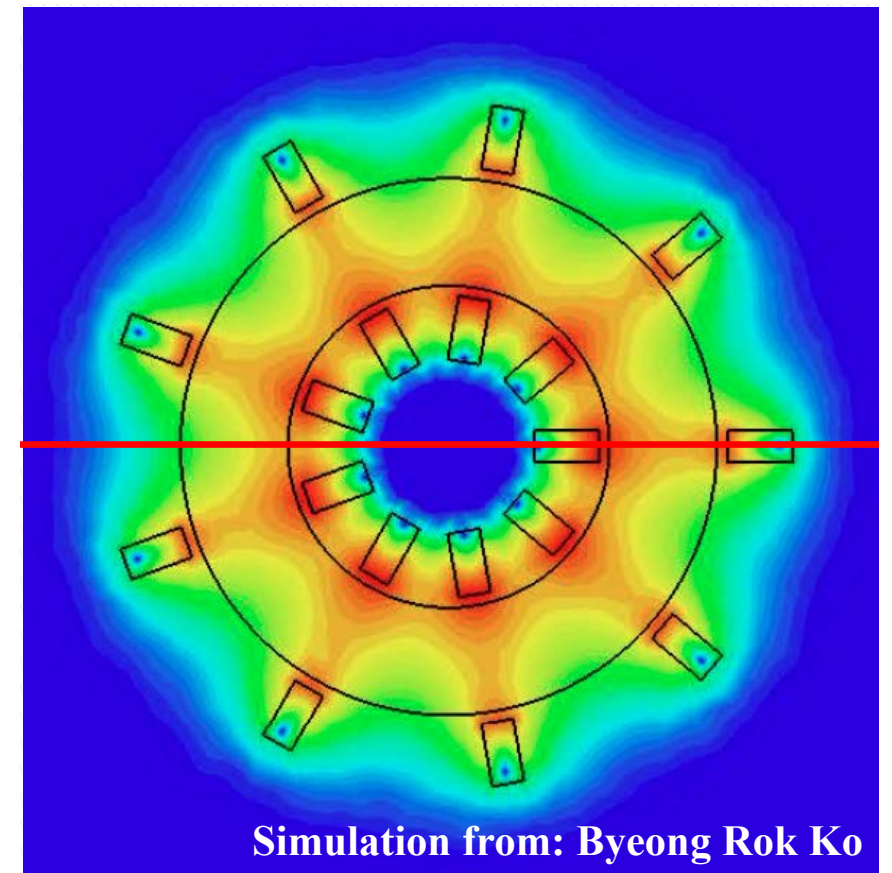
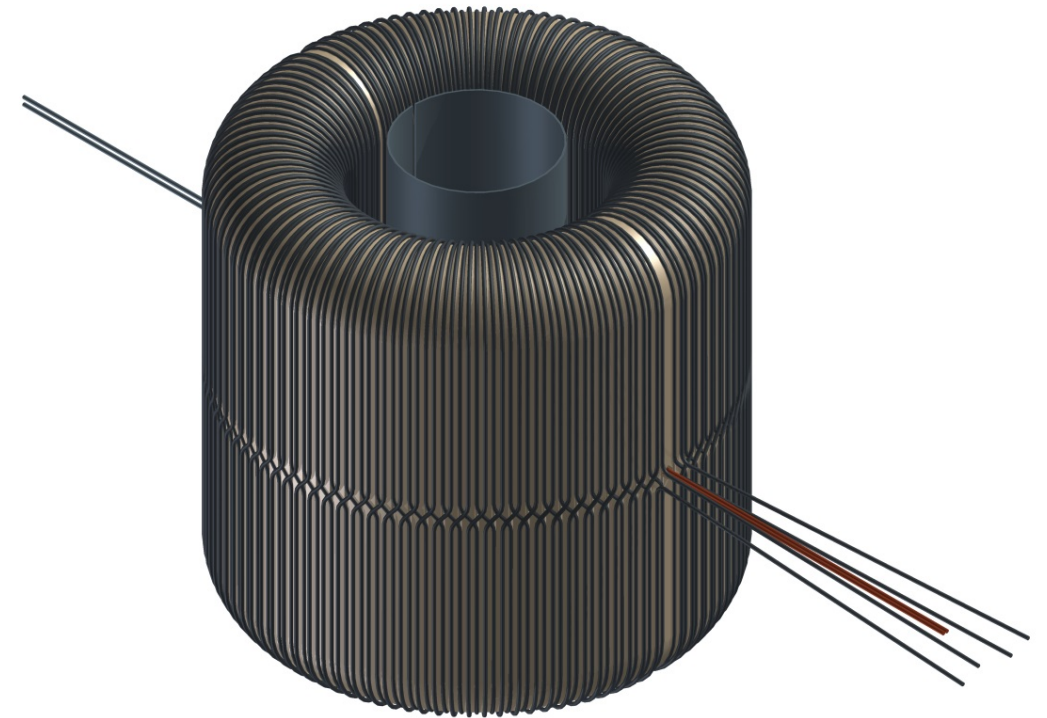
Questions we hope to address with the prototype:

1. Stray fields
2. Shielding
3. Pickup loop geometry
4. Temperature
5. Vibration
6. Decoherence and Scan Strategies



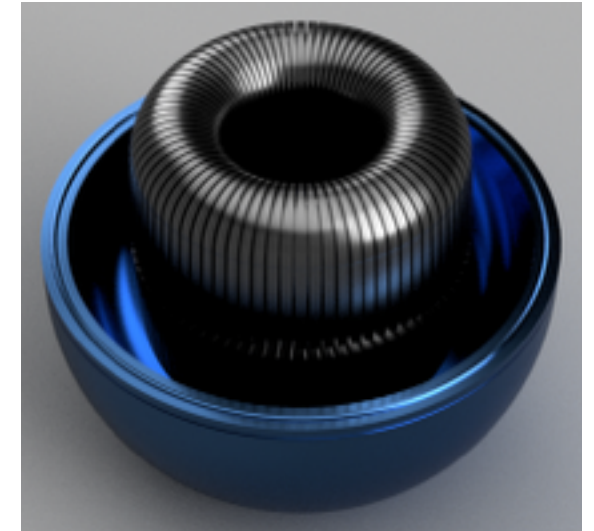
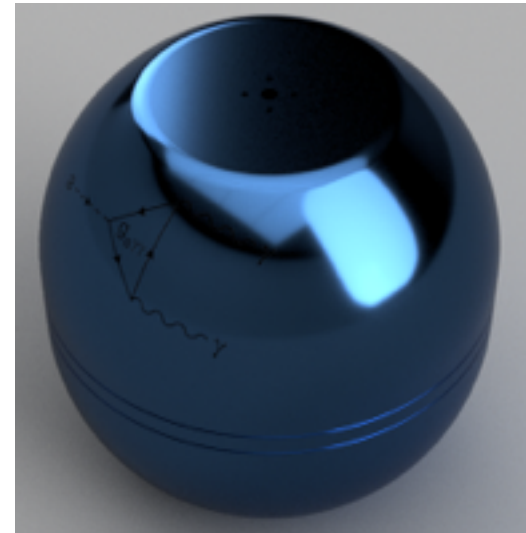
# Stray Fields And Toroid Construction

- ▶ The advantage of this approach is that we are searching for a signal in a **zero field** region
  - 4 orders of magnitude for free!
  - Need to minimize the stray fields in the center of the toroid
  - Want to keep wiring out of the center of the toroid
  - Persistent superconducting magnet to minimize noise in the coils
- ▶ Possibly investigate a segmented toroidal magnet
  - Easier to scale up in size
  - Possibly more stray fields
  - Worse geometric factor,  $G_V$



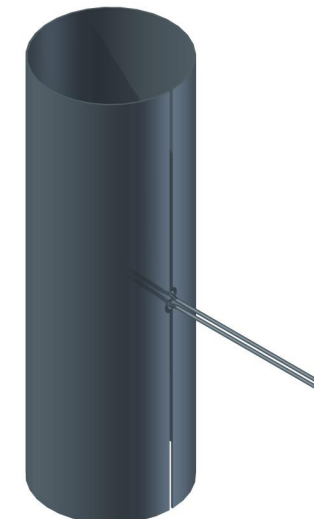
# Magnetic Shielding

- ▶ Also need to have significant shielding of environmental magnetic fields
  - Looking into encapsulating toroid in a superconducting shield
- ▶ Johnson noise in normal conducting metals generates magnetic field noise
  - All materials must be superconducting or insulators (or outside)
- ▶ Exploring setting up Helmholtz coils to cancel Earth field during cool down.



# Pickup Cylinder

- ▶ A pickup cylinder (rather than a simple wire) reduces the inductance
- ▶ Need to design detector to be disassembled and reassembled with different pickup configurations





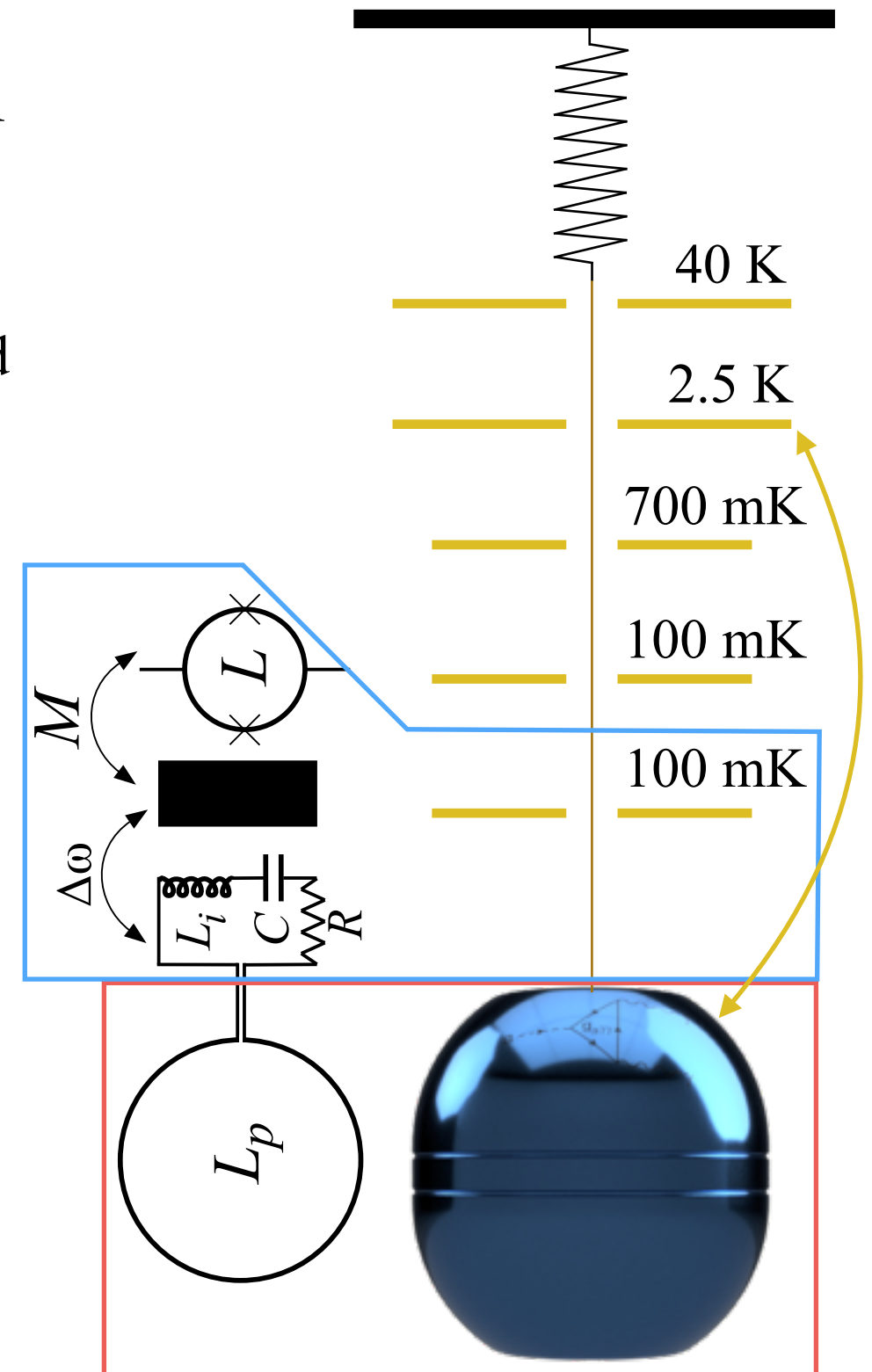
# Operating Temperature and Vibration Isolation

## The operating temperature is a big question

- ▶ The toroid and shield need to be superconducting,  $\approx 1 - 4$  K (depending on material, field, etc)
- ▶ The SQUIDs can operate between  $\sim 100$  mK and 4 K, but do not need to be the same temperature as the toroid and shield
- ▶ The only temperature dependence is the resonant circuit
- ▶ Lower temperature reduces Johnson noise, but raises  $1/f$  corner
- ▶ Does black body radiation induce noise?

## Of course vibration can induce incoherent noise

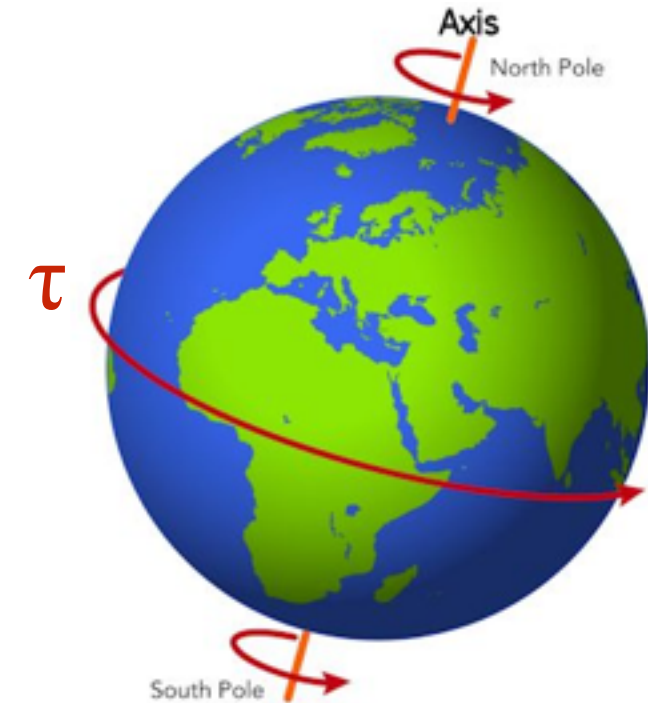
- ▶ Microphonic noise pickup creating a noise background
- ▶ Vibration of pickup relative to toroid  $\Rightarrow$  stray DC B-fields into AC B-fields



# Data Taking and Scan Strategies

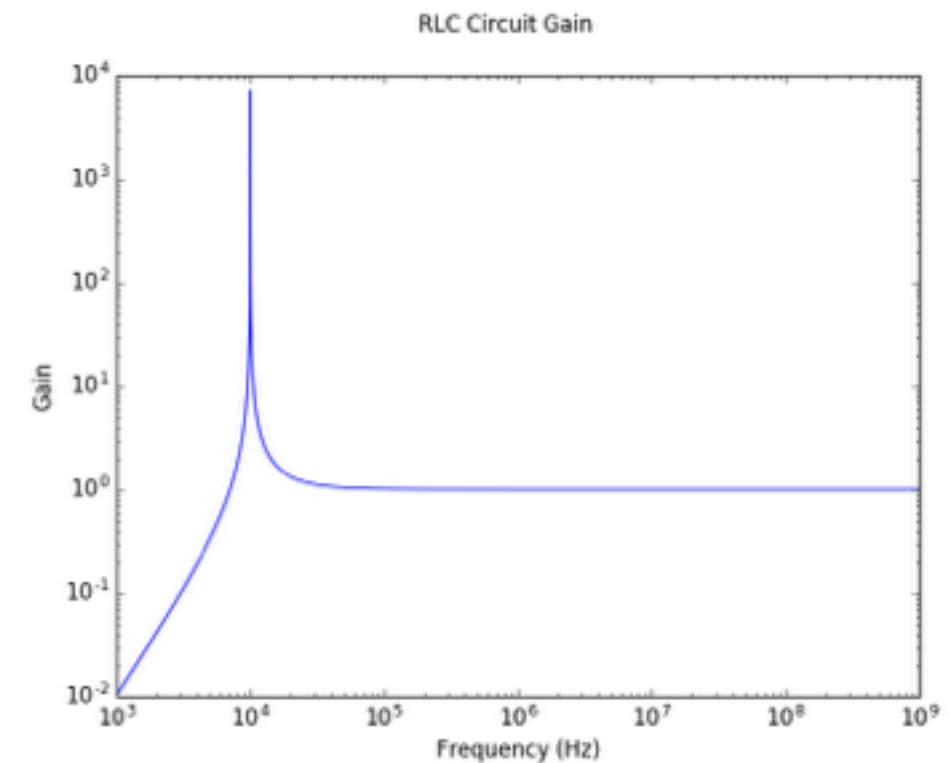
## Decoherence can broaden a signal line and reduce sensitivity

- On board timing resolution of the digitizer
  - High precision clock?
- How high of a data rate can we handle?
- At very low frequency, correcting for Earth's motion



## An optimal scan strategy can significantly improve sensitivity

- K. Irwin has pointed out that broadband data can be collected simultaneously with resonance data
- Currently studying optimal scan approaches





# ABRACADABRA Budget Estimates

## ABRACADABRA-10cm

Dilution Refrigerator	✓
SQUID Readout System	✓
Magnet	\$80k
Shielding	

## ABRACADABRA 1m

Cooling System:	
Dilution refrigerator systems	O( $\leq$ \$1M)
Larger system to cool the toroid	
SQUID Readout Systems	
Custom system with larger bandwidth and resonator	O( $\leq$ \$1M)
Shielding	
To be determined	
<b>Magnet</b>	
Typical scaling number (cost driver)	\$250k/MJ
$R_{\min}=1\text{m}, R_{\max}=2\text{m}, h=3\text{m}, B_{\max}=1\text{T}$	\$1.2M
$R_{\min}=2.2\text{m}, R_{\max}=4.5\text{m}, h=6.7\text{m}, B_{\max}=1\text{T}$	\$6M
$R_{\min}=1\text{m}, R_{\max}=2\text{m}, h=3\text{m}, B_{\max}=5\text{T}$	\$30M

\*All numbers are ballpark estimates

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Same Sensitivity

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# Summary

- ▶ Axions are interesting candidates for both Dark Matter and explaining the strong CP problem
- ▶ A broadband search using a toroidal magnet geometry has the potential to quickly probe into previously untested regions of parameter space
- ▶ Long term, this type of search could hope to probe down into the QCD axion regime
- ▶ At MIT, we are building a prototype called ABRACADABRA-10 cm with the goal of scaling this up to a 1 m scale experiment
- ▶ We are aiming to have early results by the end of 2017

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*Thank you for you attention!*

Backup Slides

# Broadband Sensitivity

Our benchmark geometry  $R_{\text{in}} = R_{\text{out}}/2 = r = h/3$ :

- The S/N ratio depends on the coherence time,  $\tau$

$$S/N \sim |\Phi_a| (t\tau)^{1/4} / S_{\Phi,0}^{1/2}$$

- The coherence time is given by

$$\tau \sim \frac{2\pi}{m_a v^2} \approx 3 \times 10^4 \text{ s} \left( \frac{10^{-12} \text{ eV}}{m_a} \right)$$

- The sensitivity goes as (depends strongly on scan strategy):

$$g_{a\gamma\gamma} > 6.3 \times 10^{-18} \text{ GeV}^{-1} \left( \frac{m_a}{10^{-12} \text{ eV}} \frac{1 \text{ year}}{t} \right)^{1/4} \frac{5 \text{ T}}{B_{\text{max}}} \\ \times \left( \frac{0.85 \text{ m}}{R} \right)^{5/2} \sqrt{\frac{0.3 \text{ GeV/cm}^3}{\rho_{\text{DM}}} \frac{S_{\Phi,0}^{1/2}}{10^{-6} \Phi_0 / \sqrt{\text{Hz}}}}$$



# Resonance Sensitivity

Our benchmark geometry  $R_{\text{in}} = R_{\text{out}}/2 = r = h/3$ :

- The signal and noise power goes as

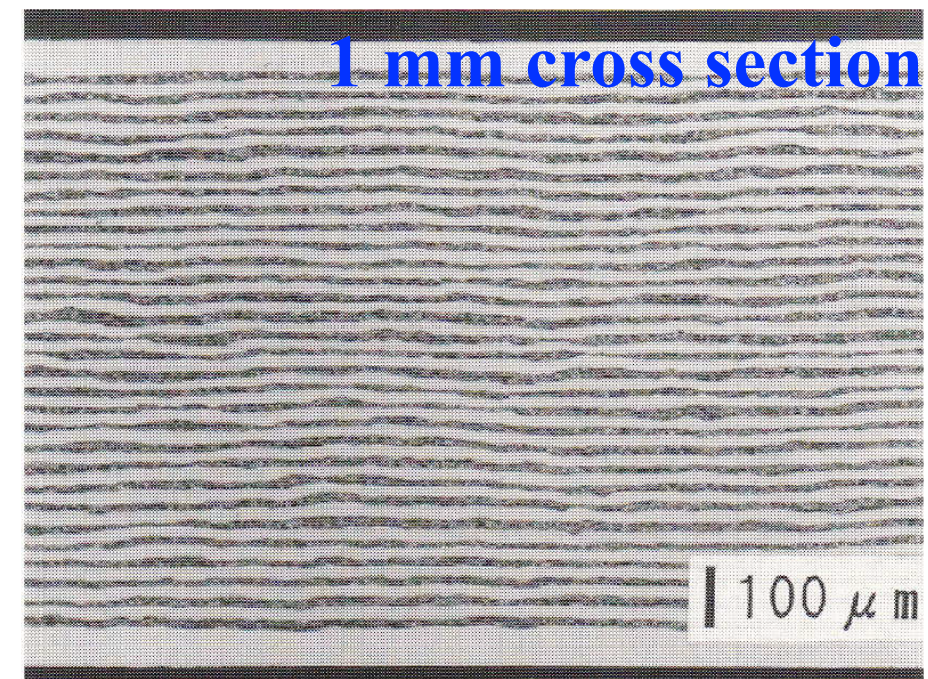
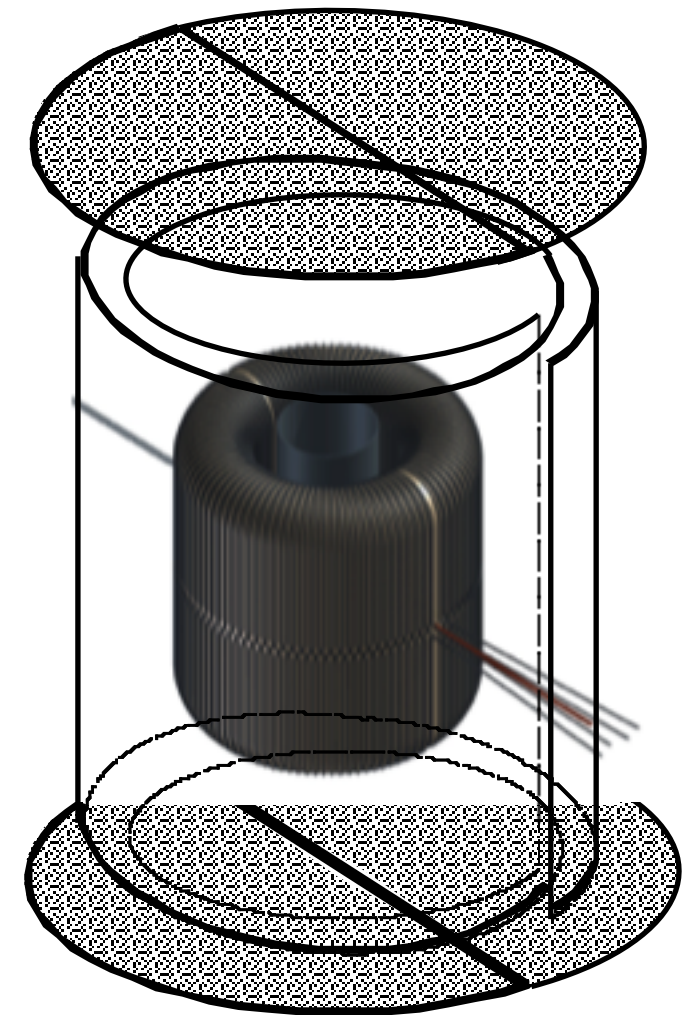
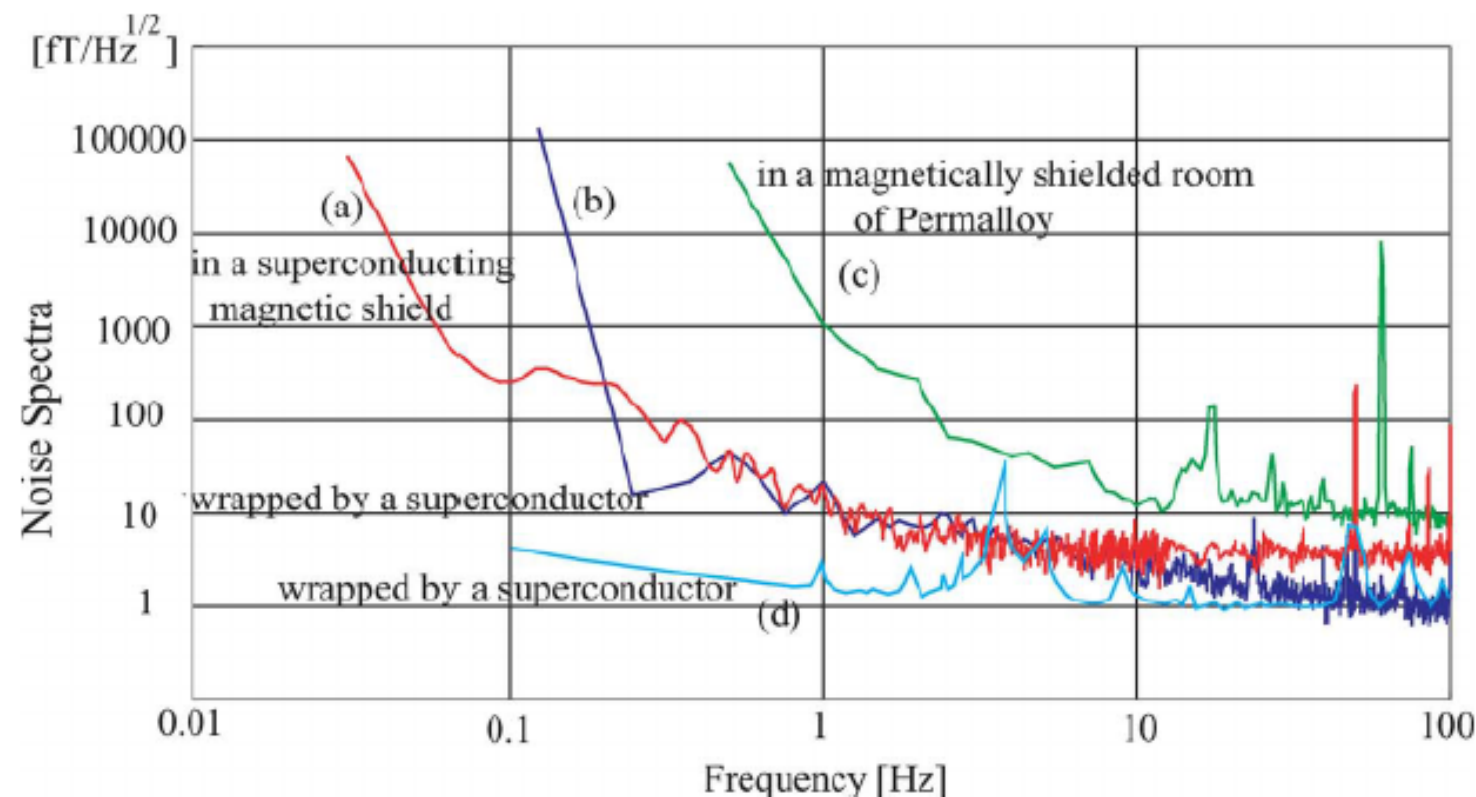
$$P_S = Q_0 \frac{m_a \Phi_a^2}{2L_T} \qquad P_N = k_B T \sqrt{\frac{m_a}{2\pi t_{\text{e-fold}}}}$$

- The sensitivity goes as

$$g_{a\gamma\gamma} > 9.0 \times 10^{-17} \text{GeV}^{-1} \left( \frac{10^{-12} \text{eV}}{m_a} \frac{20 \text{ days}}{t_{\text{e-fold}}} \right)^{1/4} \\ \times \frac{5 T}{B_{\text{max}}} \left( \frac{0.85 m}{R} \right)^{5/2} \sqrt{\frac{0.3 \text{ GeV/cm}^3}{\rho_{\text{DM}}} \frac{10^6}{Q_0} \frac{T}{0.1 K}}$$

# Magnetic Shielding

- We will need to shield the toroid from environmental magnetic noise
- Ideally, use a superconducting shield
- Looking into this NbTi/Nb/Cu multilayer sheet
  - Achieve  $\sim \text{fT}/(\text{Hz})^{1/2}$  noise levels in SQUID magnetometers



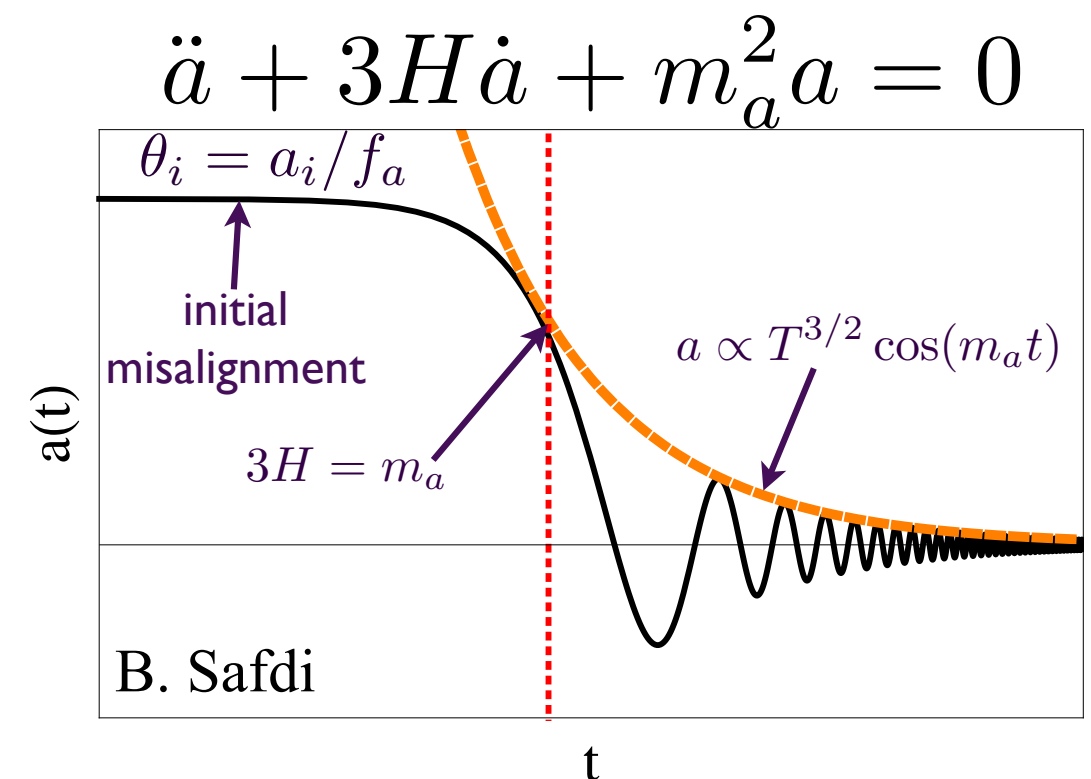
# Axion Dark Matter Density

- The axion energy matter density today is given by

$$\Omega_a h^2 \sim 0.1 \left( \frac{f_a}{10^{12} \text{ GeV}} \right)^{7/6} \theta_i^2$$

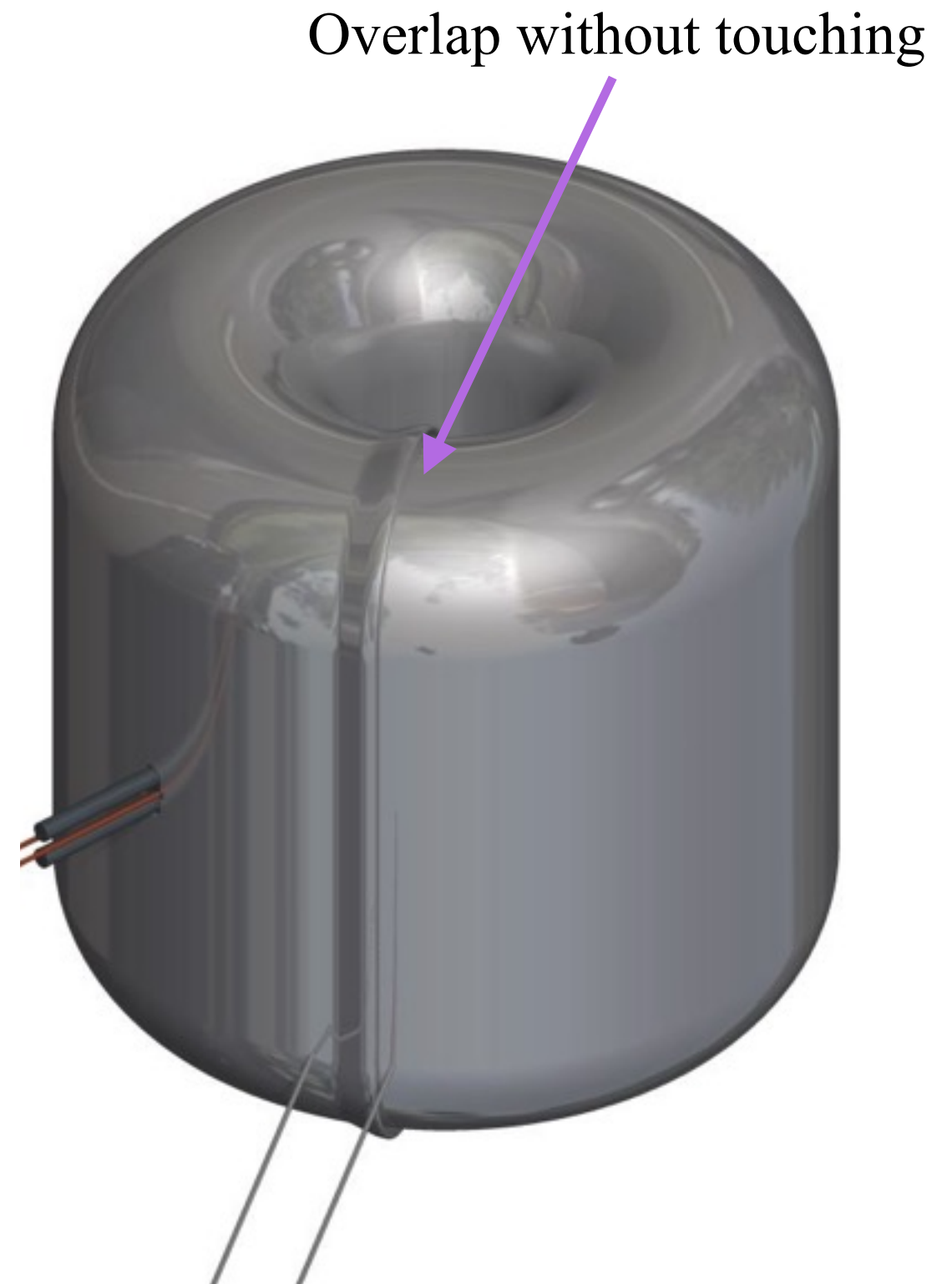
- For  $f_a = 10^{16} \text{ GeV} \Rightarrow |\theta_i| \lesssim 10^{-3} - 10^{-2}$

- Alternatively, we can also begin with a larger misalignment angle, and suppress the energy density by dumping entropy into the universe after  $3H=m_a(T)$  but before BBN.



# Future Directions?

- We have also considered surrounding the toroid in a non-overlapping superconducting shield
- The axion induced B-field is trapped inside, and the FULL axion current returns through the SC shield!



# Large Cold Detectors

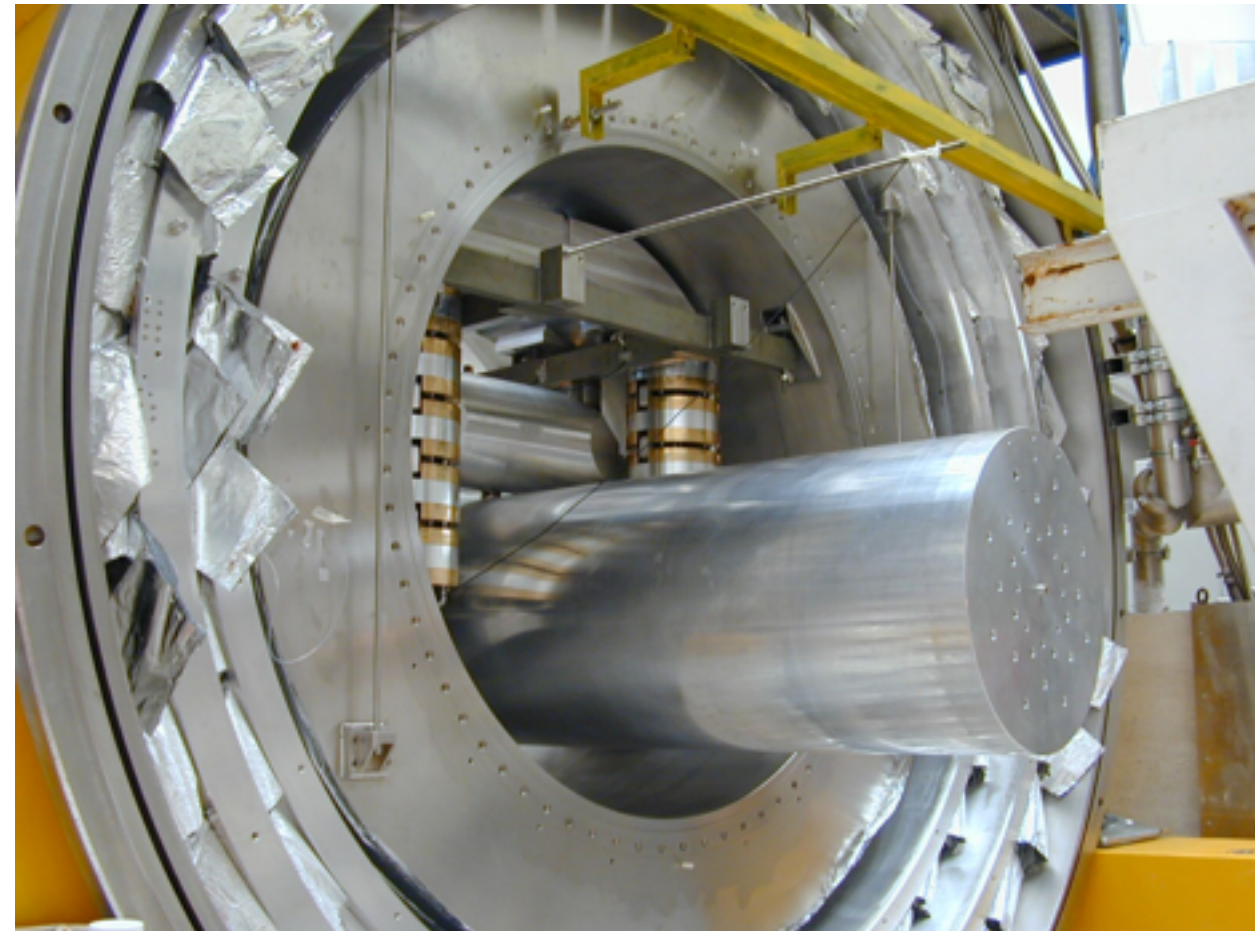
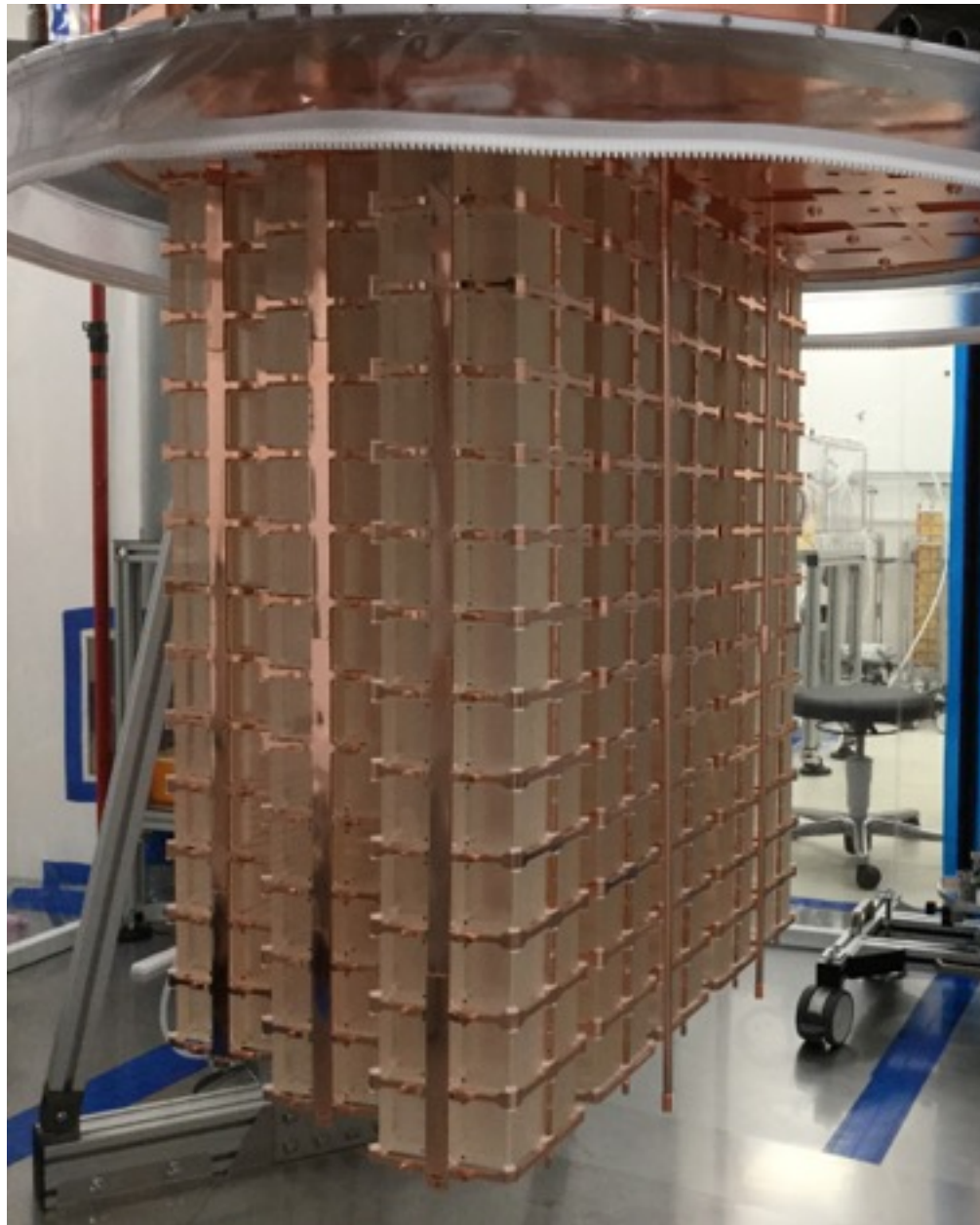
- Can we really cool down something that big??



# Large Cold Detectors

- Can we really cool down something that big??

➔ *Yes!*



AURIGA: 60 cm x 3 m  
Cryogenic gravitational  
wave antenna

CUORE: 1.5 t of material at 10 mK  
Coldest Cubic Meter in the Universe!