

Healthy degenerate theories with arbitrary higher-order derivatives

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Scalar-tensor theories

Ostrogradsky Ghost

Healthy theories with arbitrary higher-order derivatives

Dark energy

Inflation

Healthy theories with 2nd-order derivatives

DHOST / EST

Extended Galileon

GLPV

Horndeski theory

$$f(\phi, X)R$$

Brans-Dicke

$$G^{\mu\nu}\nabla_\mu\nabla_\nu\phi$$

$$X\Box\phi$$

$$f(R)$$

$$K(\phi, X)$$

Ostrogradsky ghost for $L = \frac{1}{2} \ddot{\phi}^2$ $\phi = \phi(t)$

- 4th order Euler-Lagrange eq

$$\ddot{\phi} = 0$$

1 DOF = 2 ini conds

requires 4 initial conditions = 2 DOF.

- Hamiltonian is unbounded

$$\Phi \equiv \frac{\phi - \psi}{\sqrt{2}}, \Psi \equiv \frac{\phi + \psi}{\sqrt{2}}$$

$$L = \ddot{\phi}\psi - \frac{1}{2}\psi^2 = -\dot{\phi}\dot{\psi} - \frac{1}{2}\psi^2$$

$$= \frac{1}{2}\dot{\Phi}^2 - \frac{1}{2}\dot{\Psi}^2 - \frac{1}{4}(\Phi - \Psi)^2$$

$$H = \frac{1}{2}p_{\Phi}^2 - \frac{1}{2}p_{\Psi}^2 + \frac{1}{4}(\Phi - \Psi)^2$$

2 DOF = 1 healthy + 1 ghost

Problematic when coupled
to a normal system

Ostrogradsky theorem for $L(\ddot{\phi}, \dot{\phi}, \phi)$

For Lagrangian $L(\ddot{\phi}, \dot{\phi}, \phi)$ with $\phi = \phi(t)$,

$K \equiv \partial^2 L / \partial \ddot{\phi}^2 \neq 0 \implies H$ is unbounded.

Woodard, 1506.02210

↖ 'L is nondegenerate'.

- EL eq

$$K\ddot{\ddot{\phi}} + \dot{K}\ddot{\phi} = (\text{terms up to } \ddot{\phi})$$

$K \neq 0 \implies$ 4th order system = 2 DOF

- Hamiltonian is unbounded

Ostrogradsky theorem for $L(\ddot{\phi}, \dot{\phi}, \phi)$

For Lagrangian $L(\ddot{\phi}, \dot{\phi}, \phi)$ with $\phi = \phi(t)$,

$K \equiv \partial^2 L / \partial \ddot{\phi}^2 \neq 0 \Rightarrow H$ is unbounded.

Let us impose $K = 0$

← Degeneracy condition

✓ Remove $\ddot{\phi}$ and $\dot{\phi}$
from EL eq

✓ EL eq

$$K\ddot{\phi} + \dot{K}\dot{\phi} = (\text{terms up to } \dot{\phi})$$

$K = 0 \Rightarrow$ 2nd order system = 1 DOF

✓ Hamiltonian is bounded

✓ The most general ghost-free Lagrangian is

$$L = \ddot{\phi}f(\dot{\phi}, \phi) + g(\dot{\phi}, \phi) \sim G(\dot{\phi}, \phi)$$

\Rightarrow Detailed model-dependent analysis

Eliminating Ostrogradsky ghost

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi) \Rightarrow L(\dot{\phi}, \phi)$$

Eliminating Ostrogradsky ghost

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi) \Rightarrow L(\dot{\phi}, \phi)$$

- $L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a)$

HM, Suyama, 1411.3721

- $L(\phi^{a(d)}, \phi^{a(d-1)}, \dots, \phi^a)$

Ostrogradsky theorem for $L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a)$

For $L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a)$ with $\phi^a = \phi^a(t)$ and $a = 1, \dots, n$,

$\det K \neq 0 \Rightarrow H$ is unbounded,

where $K_{ab} \equiv \frac{\partial^2 L}{\partial \ddot{\phi}^a \partial \ddot{\phi}^b}$.

← 'kinetic matrix'

Woodard, 1506.02210

- EL eq
$$K_{ab} \ddot{\phi}^b + (\dot{K}_{ab} + M_{ab}) \ddot{\phi}^b = (\text{terms up to } \ddot{\phi}^a)$$
$$M_{ab} \equiv \frac{\partial^2 L}{\partial \ddot{\phi}^a \partial \dot{\phi}^b} - \frac{\partial^2 L}{\partial \dot{\phi}^b \partial \ddot{\phi}^a}$$

$\det K \neq 0 \Rightarrow$ 4th order system.

- H is unbounded.
 n healthy + n ghost DOFs.

Eliminating Ostrogradsky ghost

For $L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a)$ with $\phi^a = \phi^a(t)$ and $a = 1, \dots, n$,

$\det K \neq 0 \Rightarrow H$ is unbounded,

where $K_{ab} \equiv \frac{\partial^2 L}{\partial \ddot{\phi}^a \partial \ddot{\phi}^b}$.

$\Rightarrow K_{ab} = 0 \longrightarrow \checkmark$ Remove $\ddot{\phi}^a$ (highest)

• EL eq
$$K_{ab} \ddot{\phi}^b + (\dot{K}_{ab} + M_{ab}) \dot{\phi}^b = (\text{terms up to } \dot{\phi}^a)$$

$$M_{ab} \equiv \frac{\partial^2 L}{\partial \dot{\phi}^a \partial \dot{\phi}^b} - \frac{\partial^2 L}{\partial \dot{\phi}^b \partial \dot{\phi}^a}$$

• Still remain ghosts.

$K_{ab} = 0$, $\det M \neq 0$ (n : even) \Rightarrow 3rd order system

• H is unbounded. $3n$ ini. conds. are required.

• Need to update the theorem.

HM, Suyama, 1411.3721

Eliminating Ostrogradsky ghost

For $L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a)$ with $\phi^a = \phi^a(t)$ and $a = 1, \dots, n$,

$\det K \neq 0$ or $\det M \neq 0 \Leftrightarrow H$ is unbounded,

where $K_{ab} \equiv \frac{\partial^2 L}{\partial \ddot{\phi}^a \partial \ddot{\phi}^b}$ and $M_{ab} \equiv \frac{\partial^2 L}{\partial \ddot{\phi}^a \partial \dot{\phi}^b} - \frac{\partial^2 L}{\partial \ddot{\phi}^b \partial \dot{\phi}^a}$.

$\Rightarrow K_{ab} = 0 \longrightarrow \checkmark$ Remove $\ddot{\phi}^a$ (highest)

$M_{ab} = 0 \longrightarrow \checkmark$ Remove $\dot{\phi}^a$ (next-highest)

\checkmark EL eq is 2nd order system.

\checkmark H is bounded. **All ghost DOFs are removed.**

\checkmark The most general ghost-free Lagrangian: $L \sim G(\dot{\phi}^a, \phi^a)$

Eliminating Ostrogradsky ghost

For $L(\phi^{a(d)}, \phi^{a(d-1)}, \dots, \phi^a)$,

$\det K \neq 0$ or $\det M \neq 0 \Rightarrow H$ is unbounded,

where $K_{ab} \equiv \frac{\partial^2 L}{\partial \phi^{a(d)} \partial \phi^{b(d)}}$ and

$$M_{ab} \equiv \frac{\partial^2 L}{\partial \phi^{a(d)} \partial \phi^{b(d-1)}} - \frac{\partial^2 L}{\partial \phi^{b(d)} \partial \phi^{a(d-1)}}.$$

$\Rightarrow K_{ab} = 0 \longrightarrow \checkmark$ Remove $\phi^{a(2d)}$ (highest)

$M_{ab} = 0 \longrightarrow \checkmark$ Remove $\phi^{a(2d-1)}$ (next-highest)

- Still remain ghosts from lower (> 2) derivatives.

Eliminating Ostrogradsky ghost

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi) \Rightarrow L(\dot{\phi}, \phi)$$

$$\checkmark L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a) \Rightarrow L(\dot{\phi}^a, \phi^a)$$

$$\bullet L(\phi^{a(d)}, \phi^{a(d-1)}, \dots, \phi^a)$$

HM, Suyama, 1411.3721

Eliminating Ostrogradsky ghost

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi) \Rightarrow L(\dot{\phi}, \phi)$$

$$\checkmark L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a) \Rightarrow L(\dot{\phi}^a, \phi^a)$$

HM, Suyama, 1411.3721

$$\bullet L(\phi^{a(d)}, \phi^{a(d-1)}, \dots, \phi^a)$$

- Ostrogradsky theorem \leftrightarrow Highest deriv. in EL eq.
- To eliminate all ghost DOFs, we need to update the theorem for each L (Degeneracy condition).
- So far, higher-order derivatives in L and/or EL eq are not allowed.
- However, for L with mixed order of derivatives (\supset ST theory), the situation is crucially different.

Horndeski theory

3 DOF = 2 for $g_{\mu\nu}$ + 1 for ϕ

Most general theory for 2nd order EL eq is Horndeski theory or generalized Galileon

Horndeski, 1974

$$L_2 = G_2(\phi, X) \quad X \equiv \nabla_\mu \phi \nabla^\mu \phi$$

Deffayet, Gao, Steer, Zahariade, 1103.3260

$$L_3 = G_3(\phi, X) \square \phi$$

Kobayashi, Yamaguchi, Yokoyama, 1105.5723

$$L_4 = G_4(\phi, X) R - 2G_{4X}(\phi, X) \left[(\square \phi)^2 - (\nabla_\mu \nabla_\nu \phi)^2 \right]$$

$$L_5 = G_5(\phi, X) G^{\mu\nu} \nabla_\mu \nabla_\nu \phi + \frac{1}{3} G_{5X}(\phi, X) \left[(\square \phi)^3 - 3 \square \phi (\nabla_\mu \nabla_\nu \phi)^2 + (\nabla_\mu \nabla_\nu \phi)^3 \right]$$

✓ 2nd order EL eq \Rightarrow Ostrogradsky ghost-free

Transformation

One can consider invertible disformal transformation

$$g_{\mu\nu} \rightarrow A(\phi, X)g_{\mu\nu} + B(\phi, X)\partial_\mu\phi\partial_\nu\phi$$

of Horndeski.

Zumalacárregui, García-Bellido, 1308.4685

While EL eq becomes a priori higher order, it is still healthy second order system. \Rightarrow beyond Horndeski

- ✓ Disformal invariance of ζ HM, White, 1504.00846
- ✓ For general field theories, Lagrangians related through general invertible transformation are physically equivalent.

Takahashi, HM, Suyama, Kobayashi, 1702.01849

Example: $L = \frac{1}{2} (\dot{x} - \ddot{y})^2 + \frac{1}{2} \dot{y}^2$ \longleftrightarrow beyond Horndeski

- EL eqs are a priori 4th order

$$\frac{d}{dt}(\dot{x} - \ddot{y}) = 0, \quad \frac{d^2}{dt^2}(\dot{x} - \ddot{y}) + \ddot{y} = 0$$

but can be rearranged to 2nd order system

$$\ddot{x} = 0, \quad \ddot{y} = 0$$

- Invertible transformation

$$\begin{aligned} X &= x - \dot{y}, & Y &= y \\ (x &= X + \dot{Y}, & y &= Y) \end{aligned} \quad \longleftrightarrow \quad \text{Disformal transformation}$$

leads Lagrangian to

$$L = \frac{1}{2} \dot{X}^2 + \frac{1}{2} \dot{Y}^2 \quad \longleftrightarrow \quad \text{Horndeski}$$

Example: $L = \frac{1}{2} \frac{\dot{x}^2}{1+\ddot{y}} + \frac{1}{2} \dot{y}^2$

Gabadadze, Hinterbichler, Khoury,
Pirtskhalava, Trodden, 1208.5773

- EL eqs are a priori 4th order

$$\frac{d}{dt} \left(\frac{\dot{x}}{1+\ddot{y}} \right) = 0, \quad \frac{d^2}{dt^2} \left(\frac{\dot{x}^2}{(1+\ddot{y})^2} \right) + \dot{y} = 0$$

but can be rearranged to 2nd order system

$$\ddot{x} = 0, \quad \ddot{y} = 0$$

- Is it equivalent to $L = \frac{1}{2} \dot{X}^2 + \frac{1}{2} \dot{Y}^2$?

At least, $X = x + a\dot{y}$, $Y = y$ cannot transform it.

- Is there some other transformation?

Not clear.

GLPV theory

$$L_2 = G_2(\phi, X)$$

$$X \equiv \phi_{;\mu} \phi^{;\mu}$$

Gleyzes, Langlois, Piazza, Vernizzi,
1404.6495, 1408.1952

$$L_3 = G_3(\phi, X) \square \phi$$

Domenech et al, 1507.05390

$$L_4 = G_4(\phi, X) R$$

Deffayet, Esposito-Farese, Steer, 1506.01974

$$-2G_{4X}(\phi, X) [(\square \phi)^2 - \phi^{;\mu\nu} \phi_{;\mu\nu}]$$

$$+ F_4(\phi, X) \epsilon^{\mu\nu\rho\sigma} \epsilon^{\tilde{\mu}\tilde{\nu}\tilde{\rho}\tilde{\sigma}} \phi_{;\mu} \phi_{;\tilde{\mu}} \phi_{;\nu\tilde{\nu}} \phi_{;\rho\tilde{\rho}}$$

$$L_5 = G_5(\phi, X) G^{\mu\nu} \phi_{;\mu\nu}$$

$$+ \frac{1}{3} G_{5X}(\phi, X) [(\square \phi)^3 - 3\phi \phi^{;\mu\nu} \phi_{;\mu\nu} + \phi_{;\mu\nu} \phi^{;\mu\sigma} \phi_{;\sigma}^{;\nu}]$$

$$+ F_5(\phi, X) \epsilon^{\mu\nu\rho\sigma} \epsilon^{\tilde{\mu}\tilde{\nu}\tilde{\rho}\tilde{\sigma}} \phi_{;\mu} \phi_{;\tilde{\mu}} \phi_{;\nu\tilde{\nu}} \phi_{;\rho\tilde{\rho}} \phi_{;\sigma\tilde{\sigma}}$$

- ✓ 3 DOF = 2 $g_{\mu\nu}$ + 1 ϕ / Ostrogradsky ghost-free
- ✓ Higher-order EL eqs are reducible to 2nd order system
- ✓ Includes subclass related to Horndeski through disformal transformation $g_{\mu\nu} \rightarrow A(\phi)g_{\mu\nu} + B(\phi, X)\partial_\mu\phi\partial_\nu\phi$

Healthy scalar-tensor theories

quadratic DHOST / EST

Langlois, Noui, 1510.06930, 1512.06820

Crisostomi, Koyama, Tasinato, 1602.03119

Achour, Langlois, Noui, 1602.08398

- Write down all possible derivative terms with coefficient functions at $(\nabla\nabla\phi)^2$, $(\partial\partial g)^2$ order

$$S = \int d^4x \sqrt{-g} \left[F_2(\phi, X) R + \sum_{i=1}^5 A_i(\phi, X) L_i^{(2)} \right]$$

$$L_1^{(2)} = (\phi_{;\mu\nu})^2, L_2^{(2)} = (\square\phi)^2, L_3^{(2)} = (\square\phi)\phi^{;\mu}\phi_{;\mu\nu}\phi^{;\nu},$$
$$L_4^{(2)} = \phi^{;\mu}\phi_{;\mu\nu}\phi^{;\nu\rho}\phi_{;\rho}, L_5^{(2)} = (\phi^{;\mu}\phi_{;\mu\nu}\phi^{;\nu})^2.$$

- Impose degeneracy condition on F_2 and A_i .

Healthy scalar-tensor theories

quadratic DHOST / EST

- Needs to derive degeneracy condition for L with mixed orders of derivatives.

- Toy model was first studied [Langlois, Noui, 1510.06930](#)

$$L = \frac{1}{2} a \ddot{\phi}^2 + \frac{1}{2} k_0 \dot{\phi}^2 + \frac{1}{2} k_{ij} \dot{q}^i \dot{q}^j + b_i \ddot{\phi} \dot{q}^i + c_i \dot{\phi} \dot{q}^i - V(\phi, q)$$

which applies to only $(\nabla\nabla\phi)^2$ order theory.

- More general analysis is required beyond $(\nabla\nabla\phi)^2$.

[HM, Noui, Suyama, Yamaguchi, Langlois, 1603.09355](#)

[See also Klein, Roest, 1604.01719; Crisostomi, Klein, Roest, 1703.01623](#)

Eliminating Ostrogradsky ghost

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi) \Rightarrow L(\dot{\phi}, \phi)$$

$$\checkmark L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a) \Rightarrow L(\dot{\phi}^a, \phi^a)$$

HM, Suyama, 1411.3721

$$\bullet L(\phi^{a(d)}, \phi^{a(d-1)}, \dots, \phi^a)$$

HM, Noui, Suyama, Yamaguchi, Langlois, 1603.09355

$$\bullet L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}, q)$$

$$\bullet L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}^i, q^i) \sim \phi + g_{\mu\nu}$$

$$\bullet L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a; \dot{q}^i, q^i) \sim \phi^a + g_{\mu\nu}$$

$$L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}, q)$$

$\phi(t), q(t)$ have different orders of derivatives.

Equivalent form

$$L(\dot{Q}, Q, \phi; \dot{q}, q) + \lambda(\dot{\phi} - Q).$$

Canonical momenta for ϕ and λ

\Rightarrow 2 primary constraints

8 ini. conds.

From $P = L_{\dot{Q}}$ and $p = L_{\dot{q}}$ ($L_X \equiv \partial L / \partial X$)

$$\begin{pmatrix} \delta P \\ \delta p \end{pmatrix} = \begin{pmatrix} L_{\dot{Q}\dot{Q}} & L_{\dot{q}\dot{Q}} \\ L_{\dot{Q}\dot{q}} & L_{\dot{q}\dot{q}} \end{pmatrix} \begin{pmatrix} \delta \dot{Q} \\ \delta \dot{q} \end{pmatrix}.$$

kinetic matrix K

If $\det K \neq 0 \Rightarrow$ No further primary constraints.

We thus impose $\det K = 0$.

$$L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}, q)$$

The degeneracy condition: $\det K = 0$

In particular, we consider the case

$$L_{\dot{q}\dot{q}} \neq 0, \quad \text{and} \quad L_{\dot{Q}\dot{Q}} - L_{\dot{q}\dot{Q}}^2/L_{\dot{q}\dot{q}} = 0,$$

✓ q -sector is "normal"

✓ Remove $\ddot{\phi}$ & $\dot{\phi}$

$$\begin{pmatrix} \delta P \\ \delta p \end{pmatrix} = \begin{pmatrix} L_{\dot{Q}\dot{Q}} & L_{\dot{q}\dot{Q}} \\ L_{\dot{q}\dot{Q}} & L_{\dot{q}\dot{q}} \end{pmatrix} \begin{pmatrix} \delta \dot{Q} \\ \delta \dot{q} \end{pmatrix}$$

$$\Rightarrow \delta P = \frac{L_{\dot{q}\dot{Q}}}{L_{\dot{q}\dot{q}}} \delta p, \quad \delta p = L_{\dot{q}\dot{Q}} \delta \dot{Q} + L_{\dot{q}\dot{q}} \delta \dot{q}$$

$$\Rightarrow P = F(p, q, Q, \phi)$$

$$L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}, q)$$

The degeneracy condition: $\det K = 0$

In particular, we consider the case

$$L_{\dot{q}\dot{q}} \neq 0, \quad \text{and} \quad L_{\dot{Q}\dot{Q}} - L_{\dot{q}\dot{Q}}^2/L_{\dot{q}\dot{q}} = 0,$$

✓ q -sector is "normal"

✓ Remove $\ddot{\phi}$ & \ddot{Q}

Additional 1 primary constraint $P = F(p, q, Q, \phi)$

⇒ Additional 1 secondary constraint

⇒ $(8 - 2 - 2)/2 = 2$ healthy DOF.

✓ EL eqs can be reducible to 2nd order system.

$$L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}^i, q^i) \quad i = 1, \dots, m$$

The degeneracy condition: $\det K = 0$

In particular, we consider the case

$$\det L_{\dot{q}^i \dot{q}^j} \neq 0, \quad \text{and} \quad L_{\dot{Q} \dot{Q}} - L_{\dot{q}^i \dot{Q}} L_{\dot{q}^i \dot{q}^j}^{-1} L_{\dot{q}^j \dot{Q}} = 0,$$

✓ q^i -sector is "normal" ✓ Remove $\ddot{\phi}$ & $\dot{\phi}$

Additional 1 primary constraint $P = F(p_i, q^i, Q, \phi)$

⇒ Additional 1 secondary constraint

⇒ $(2m + 6 - 2 - 2)/2 = m + 1$ healthy DOF.

✓ EL eqs can be reducible to 2nd order system.

$$L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a; \dot{q}^i, q^i) \quad a = 1, \dots, n \quad i = 1, \dots, m$$

The degeneracy condition: $\det K = 0$

In particular, we consider the case

$$\det L_{\dot{q}^i \dot{q}^j} \neq 0, \quad \text{and} \quad L_{\dot{q}^a \dot{q}^b} - L_{\dot{q}^i \dot{q}^a} L_{\dot{q}^i \dot{q}^j}^{-1} L_{\dot{q}^j \dot{q}^b} = 0,$$

✓ q^i -sector is "normal" ✓ Remove $\ddot{\phi}^a$ but not $\dot{\phi}^a$

Additional n primary constraint

$$\Xi_a \equiv P_a - F_a(p_i, q^i, Q^b, \phi^b) = 0$$

⇒ Impose $\{\Xi_a, \Xi_b\} = 0$ ✓ Remove $\ddot{\phi}^a$

⇒ Additional n secondary constraint

⇒ $(2m + 6n - 2n - 2n)/2 = m + n$ healthy DOF.

✓ EL eqs can be reducible to 2nd order system.

Eliminating Ostrogradsky ghost

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi) \Rightarrow L(\dot{\phi}, \phi)$$

$$\checkmark L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a) \Rightarrow L(\dot{\phi}^a, \phi^a)$$

HM, Suyama, 1411.3721

$$\bullet L(\phi^{a(d)}, \phi^{a(d-1)}, \dots, \phi^a)$$

HM, Noui, Suyama, Yamaguchi, Langlois, 1603.09355

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}, q)$$

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}^i, q^i) \sim \phi + g_{\mu\nu}$$

$$\checkmark L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a; \dot{q}^i, q^i) \sim \phi^a + g_{\mu\nu}$$

Healthy scalar-tensor theories

cubic DHOST / EST

Achour, Crisostomi, Koyama, Langlois, Noui, Tasinato, 1608.08135

- At $(\nabla\nabla\phi)^3, (\partial\partial g)^2\nabla\nabla\phi$ order

$$S = \int d^4x \sqrt{-g} \left[F_3(\phi, X)R + \sum_{i=1}^{10} B_i(\phi, X)L_i^{(3)} \right]$$

$$L_1^{(3)} = (\square\phi)^3, L_2^{(3)} = \square\phi(\phi_{;\mu\nu})^2, L_3^{(3)} = \phi_{;\mu\nu}\phi^{;\nu\rho}\phi_{;\rho}^{;\mu},$$
$$L_4^{(3)} = (\square\phi)^2\phi^{;\mu}\phi_{;\mu\nu}\phi^{;\nu}, L_5^{(3)} = \dots.$$

- Impose degeneracy condition on F_3 and B_i .

Healthy scalar-tensor theories

quadratic & cubic DHOST / EST

✓ Ghost-free Lagrangian for $(\nabla\nabla\phi)^2$ and $(\nabla\nabla\phi)^3$

Open questions

- Can we check whether \exists redefinition of fields for given higher-order theory?

- Higher powers? Multi-field?

HM, Noui, Suyama, Yamaguchi, Langlois, 1603.09355

- Even higher order derivatives?

HM, Suyama, Yamaguchi, in prep.

Eliminating Ostrogradsky ghost

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$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}, q)$$

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}^i, q^i) \sim \phi + g_{\mu\nu}$$

$$\checkmark L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a; \dot{q}^i, q^i) \sim \phi^a + g_{\mu\nu}$$

HM, Suyama, Yamaguchi, in prep.

$$\bullet L(\ddot{\psi}, \dot{\psi}, \psi, \psi; \dot{q}^i, q^i)$$

$$\bullet L(\ddot{\psi}^n, \dot{\psi}^n, \psi^n, \psi^n; \ddot{\phi}^a, \dot{\phi}^a, \phi^a; \dot{q}^i, q^i)$$

$$\bullet L(\phi^{i_d(d+1)}, \dots; \phi^{i_{d-1}(d)}, \dots; \dots; \dot{\phi}^{i_0}, \phi^{i_0})$$

$$L(\ddot{\psi}, \dot{\psi}, \psi, \psi; \dot{q}^i, q^i) \quad i = 1, \dots, I$$

Equivalent form

$$L(\dot{Q}, Q, R, \psi; \dot{q}^i, q^i) + \xi(\dot{\psi} - R) + \lambda(\dot{R} - Q).$$

Canonical momenta for $(Q, R, \psi, q^i, \xi, \lambda)$

\Rightarrow 4 primary constraints

10 + 2I ini. conds.

Need more constraints to remove Ostrogradsky ghost.

Impose $\det K = 0$ or

$$\det L_{\dot{q}^i \dot{q}^j} \neq 0, \quad \text{and} \quad L_{\dot{Q} \dot{Q}} - L_{\dot{q}^i \dot{Q}} L_{\dot{q}^i \dot{q}^j}^{-1} L_{\dot{q}^j \dot{Q}} = 0.$$

$$L(\ddot{\psi}, \dot{\psi}, \psi, \psi; \dot{q}^i, q^i) \quad i = 1, \dots, I$$

The degeneracy condition: $\det K = 0$ or

$$\det L_{\dot{q}^i \dot{q}^j} \neq 0, \quad \text{and} \quad L_{\dot{Q} \dot{Q}} - L_{\dot{q}^i \dot{Q}} L_{\dot{q}^i \dot{q}^j}^{-1} L_{\dot{q}^j \dot{Q}} = 0,$$

→ ✓ Remove $\psi^{(6)}$ & $\psi^{(5)}$ ↓

Additional 1 primary constraint $\Psi \equiv P_Q - F(p_i, x) = 0$

⇒ 1 secondary constraint $\Upsilon \equiv P_R - G(p_i, x) = 0$

Impose $\{\Upsilon, \Psi\} = 0$ → ✓ Remove $\psi^{(4)}$ & $\psi^{(3)}$

⇒ 1 tertiary constraint $\Lambda \equiv \pi - I(p_i, x) = 0$

⇒ 1 quaternary constraint $\Omega \equiv -\{\Lambda, H\} = 0$

Impose $Z \equiv \{\Omega, \Psi\} \neq 0$ → ✓ Stop Dirac algorithm

⇒ $(10 + 2I - 4 - 4)/2 = 1 + I$ healthy DOF.

✓ EL eqs can be reducible to 2nd order system.

Eliminating Ostrogradsky ghost

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi) \Rightarrow L(\dot{\phi}, \phi)$$

$$\checkmark L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a) \Rightarrow L(\dot{\phi}^a, \phi^a)$$

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HM, Noui, Suyama, Yamaguchi, Langlois, 1603.09355

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}, q)$$

$$\checkmark L(\ddot{\phi}, \dot{\phi}, \phi; \dot{q}^i, q^i) \sim \phi + g_{\mu\nu}$$

$$\checkmark L(\ddot{\phi}^a, \dot{\phi}^a, \phi^a; \dot{q}^i, q^i) \sim \phi^a + g_{\mu\nu}$$

HM, Suyama, Yamaguchi, in prep.

$$\checkmark L(\ddot{\psi}, \dot{\psi}, \psi; \dot{q}^i, q^i)$$

$$\checkmark L(\ddot{\psi}^n, \dot{\psi}^n, \psi^n; \ddot{\phi}^a, \dot{\phi}^a, \phi^a; \dot{q}^i, q^i)$$

$$\checkmark L(\phi^{i_d(d+1)}, \dots; \phi^{i_{d-1}(d)}, \dots; \dots; \dot{\phi}^{i_0}, \phi^{i_0})$$

Summary

- We derived the degeneracy condition for various types of Lagrangians to eliminate Ostrogradsky ghost, which allow us to construct healthy scalar-tensor theories, and proceed to detailed model-dependent analysis.
- We clarified that it is possible to construct healthy Lagrangian involving arbitrary higher-order derivatives.
- Application to scalar-tensor theories and their classifications are work in progress.

Scalar-tensor theories

Ostrogradsky Ghost

Healthy theories with arbitrary higher-order derivatives

Dark energy

Inflation

Healthy theories with 2nd-order derivatives

DHOST / EST

Extended Galileon

GLPV

Horndeski theory

$$f(\phi, X)R$$

Brans-Dicke

$$G^{\mu\nu}\nabla_\mu\nabla_\nu\phi$$

$$X\Box\phi$$

$$f(R)$$

$$K(\phi, X)$$