# **Overview of ATLAS Heavy Flavor Measurements**

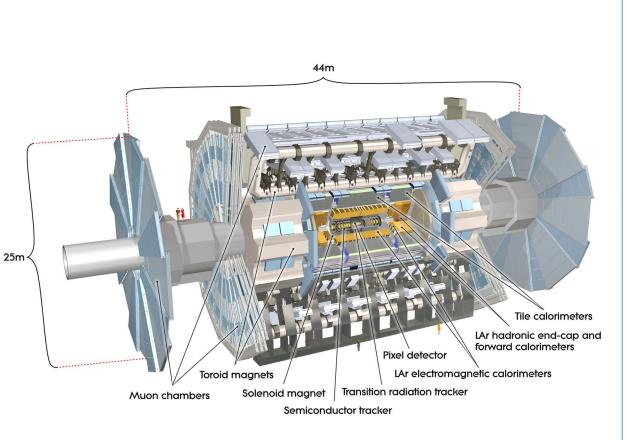
Sally Seidel University of New Mexico

Santa Fe Jets and Heavy Flavor Workshop 31 January 2018

- I. Introduction to ATLAS
- II. b-Hadron Pair Production Cross-section
- III. Prompt J/ $\psi$  Pair Production Cross-section
- IV.  $\psi(2S) \rightarrow J/\psi\pi^+\pi^-$  and  $X(3872) \rightarrow J/\psi\pi^+\pi^-$  Production
- V. Angular Analysis of  $B_d^0 \to K^* \mu^+ \mu^-$  Decays

### Introduction

The 4 most recent public results in B-Physics from ATLAS, all using LHC pp data collected at  $\sqrt{s} = 8$  TeV, and released in 2017.



ATLAS from inside to out:

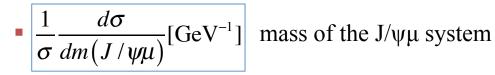
- Inner detector (pixel, silicon microstrips, strawtube TRT)  $|\eta| < 2.5$ , surrounded by a 2T axial B field from the solenoid
- Sampling calorimeters (LAr EM  $|\eta| < 3.2$ ; Scint tile HAD  $|\eta| < 3.2$ ; LAr HAD  $1.5 < |\eta| < 4.9$ )
- Air core toroids provide B field for Muon drift tubes + cathode strip chambers (muon tracking to  $|\eta| < 2.7$ ) and resistive plate + thin gap chambers (triggering to  $|\eta| < 2.4$ ) 3

# Measurement of b-hadron Pair Production Cross-section\*

*Message*: This total cross section is measured:  $\sigma (B(\rightarrow J/\psi [\rightarrow \mu^+ \mu^-] + X)B(\rightarrow \mu + X))$ 

Using it, 8 differential cross sections are obtained:

separation between the J/ $\psi$  and the third  $\mu$  in the azimuth-rapidity plane



 $\frac{1}{\sigma} \frac{d\sigma}{d\Delta R(J/\psi\mu)}$ 

 $\frac{1}{\sigma} \frac{d\sigma}{dp_{\tau} (J/\psi u)}$ 

 $\frac{1}{\sigma} \frac{d\sigma}{d\Delta\phi (J/\psi\mu)}$  [rad<sup>-1</sup>] azimuthal separation  $\Delta\phi$  between the J/ $\psi$  and the third  $\mu$ 

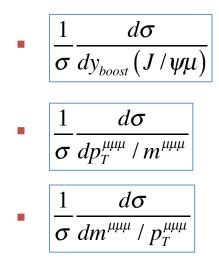
transverse momentum  $p_T$  of the 3-muon system

$$\frac{1}{\sigma} \frac{d\sigma}{d\Delta y (J/\psi\mu)}$$

rapidity separation  $\Delta y$  between the J/ $\psi$  and the third  $\mu$ 

the list continues.....

\*JHEP 11 (2017) 062.



magnitude  $y_{\text{boost}}$  of the avg. rapidity of the J/ $\psi$  and the third  $\mu$ 

ratio of the  $p_T$  to the invariant mass of the 3-muon system,

and its inverse

These differential cross sections are compared to predictions from several event generators.

### **Motivation:**

- Factorization of QCD calculations into parton distribution functions, hard matrix elements, and soft parton shower components allows the heavy (b) quark mass to be introduced at *different stages*.
- *Several schemes are possible* for inclusion of the heavy quark masses
- Previous analyses of heavy flavor production highlighted disagreements *among* theoretical predictions and *between* predictions and data. *This analysis constrains the options*.
- The region of small-angle  $b\overline{b}$  production is *especially sensitive* to details of the calculations but has previously been *only loosely constrained* by data.
- Searches for Higgs produced in association with a vector boson (VH) and decaying to  $b\overline{b}$ rely on the modeling of the background  $b\overline{b} + V$

# **Details of the analysis (1)**

- Trigger: 2 oppositely charged muons with a common vertex,  $p_T(\mu) > 4$  GeV,  $|\eta(\mu)| < 2.4$ , 2.5  $< m(\mu\mu) < 4.3$  GeV
- Integrated luminosity = 11.4 fb<sup>-1</sup>
- Primary vertex:  $\geq 2$  tracks, each with  $p_T > 400$  MeV, with largest summed  $p_T^2$ .
- Form the muon candidates:
  - use combined inner detector and muon spectrometer tracks
  - $p_{\rm T}(\mu) > 6 \text{ GeV}, |\eta(\mu)| < 2.5$
- J/ψ candidates:
  - opposite-sign muon pairs with  $|\eta(\mu)| < 2.3$  and directional correspondence with the trigger-level candidate
  - $2.6 < m(\mu\mu) < 3.5 \text{ GeV}$
  - If multiple candidates per event, choose the one with mass closest to  $J/\psi_{PDG}$ .
- Third muon: choose the highest- $p_T$  one not included in the J/ $\psi$  reconstruction.
- The  $J/\psi$  and the third  $\mu$  may come from feed-down or cascade.
- The data are first compared to these simulations:
  - Inclusive b-hadron pairs from PYTHIA8.186 (2->2 matrix element with parton shower); CTEQ6L1 pdf, AU2 tune; b quarks are massless in the pdf but the mass is reinstated during the shower; pile-up included with PYTHIA8 + MSTW2008 pdf + A2 tune.
  - $pp \rightarrow b\overline{b}$  simulated with HERWIG++, CTEQ6L1, UE-EE5 tune; b-quarks are massive in the matrix element and in the parton shower.
- 4-momenta of photons near muon ( $\Delta R_{\eta}(\mu,\gamma) < 0.1$ ) added to muon

### Analysis details (2)

Corrections:

- for trigger efficiency including vertex recon and spatial overlap of muons
- for muon reconstruction efficiency
- To collect the  $J/\psi$ 's produced in decays of b-hadrons:
  - Define  $L_{xy}$ : transverse distance between primary vertex (PV) and dimuon vertex, signed positively for momentum pointing away from primary vertex.
  - Define pseudo-proper decay time:

$$\tau \equiv \frac{L_{xy} \cdot m \left( J / \psi_{PDG} \right)}{p_T \left( \mu^+ \mu^- \right)}$$

- J/ $\psi$ 's from most b decays are non-prompt, so to optimize for signal events, require  $\tau > 0.25$  mm/c.
- simultaneous maximum likelihood fit to the distributions of dimuon mass and  $\tau$ .
- Extract # non-prompt  $J/\psi$ 's.

### Analysis details (3)

- To select the third muon, reject bkgs: prompt muons, muons from charged π/K decay, fake muons from decay in flight and hadron shower leakage, muons combined with continuum (false) J/ψ, and muons in pile-up.
- Discriminate third-muon signal from bkg with a simultaneous fit on 2 observables:
  - transverse impact parameter significance

$$S_{d_0} \equiv d_0 \, / \, \boldsymbol{\sigma}_{d_0}$$

( $d_0$  is distance of closest approach of the muon track to the PV in the r- $\phi$  projection, with sign given by the sign of the angular momentum of the track around the beam at point of closest approach)

- Output of a boosted decision tree using kinematic variables related to track deflection significance, momentum balance, and  $|\eta|$ .
- Subtract 3 remaining irreducible bkgs from fitted yields:
  - $B_c \rightarrow J/\psi + \mu + X$  (very small, taken from simulation)
  - Semileptonic decays of c-hadrons not resulting from b-hadron feed-down
  - "Sail through" charged π/K: traverses the detector to the muon spectrometer without interacting or decaying (mimics a muon, taken from simulation)

### Analysis details (4)

Corrections:

- for the τ requirement: extrapolate to full range
- for detector resolution on momentum and η of muons. Issue: migration between bins and in/out of fiducial volume.

Repeat for every kinematic bin for each differential cross section.

Systematic uncertainties:

- Muon efficiency corrections to data
- $J/\psi$  model
- Background components in the fits

# Statistical uncertainties:

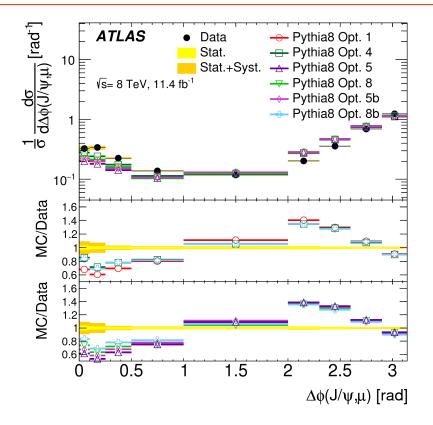
- On the data statistics
- On the third-muon templates taken from simulation

Luminosity uncertainty: 1.9%

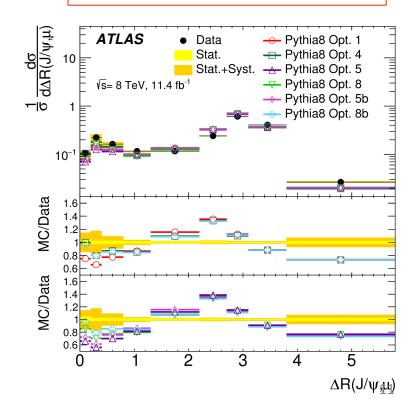
**Result 1:** 
$$\sigma \left( B \left( \rightarrow J / \psi \left[ \rightarrow \mu^+ \mu^- \right] + X \right) B \left( \rightarrow \mu + X \right) \right) = 17.7 \pm 0.1 \text{(stat)} \pm 2.0 \text{(syst) nb.}$$

**Result 2**: Is the scale of  $\alpha_s$  during splitting set by *relative*  $p_T$  or by *mass*? Compare differential cross sections using 6 options in PYTHIA8 for the  $g \rightarrow b\overline{b}$  splitting kernel (dominates small angle b-hadron production).

PYTHIA8 does not reproduce the shape of the angular distributions for any of the 6 options.



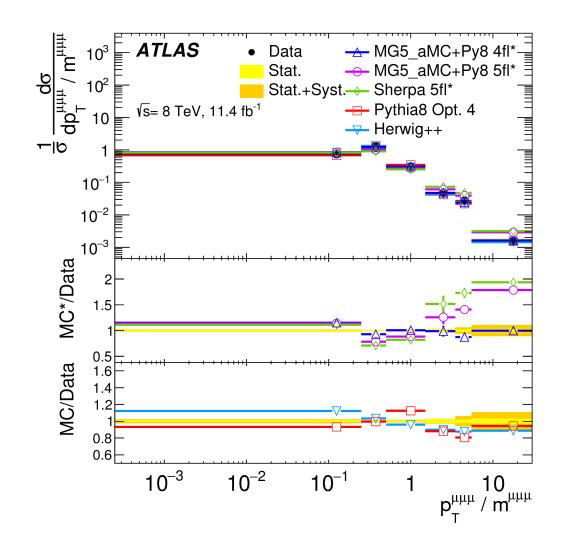
Some of the options of splitting function form and  $\alpha_s$  scale fit better to the mass or  $\Delta R$  distributions. Example:



### **Result 3:**

Extend the comparison of data to HERWIG++, SHERPA, and MADGraph5\_AMC@NLOv2.2.2 + PYTHIA8.186 parton shower model. These cover a range of matrix element calculations and parton shower models. Consider options with 4 or 5 massless flavors. Compare all of these to PYTHIA8.

- HERWIG++ reproduces the ΔR and Δφ graphs best.
- 4-massless flavors models ΔR and Δφ better than 5.
- Δy spectrum is well modeled by MadGraph and SHERPA
- All models reproduce y<sub>boost</sub> well.
- 5-massless flavor MadGraph models low mass distribution better than 4,
- but 4-massless flavor MadGraph models high p<sub>T</sub>/m best.



### **Conclusions:**

- Considering all distributions, the 4-massless flavor prediction from MadGraph5\_AMC@NLO+PYTHIA8 best describes the data.
- Predictions of PYTHIA8 and HERWIG++ are comparable.
- Among PYTHIA8 options studied, the p<sub>T</sub>-based splitting kernel is best.

# Measurement of Prompt J/ $\psi$ Pair Production Cross-section<sup>\*</sup>

*Message:* The cross section for production of 2 prompt centrally-produced J/ $\psi$  mesons is measured. *"Prompt" means: produced at a point consistent with the primary vertex, not as a product of the decay of a long-lived hadron.* 

Differential cross sections are produced as a function of:

- $p_T$  of the lower- $p_T$  meson (called "J/ $\psi_2$ ") Measurements use subleading meson J/ $\psi_2$  to access full kinematic region.
- $di-J/\psi p_T$
- di-J/ψ mass
- Δy between the 2 mesons
- $\Delta \phi$  between the 2 mesons

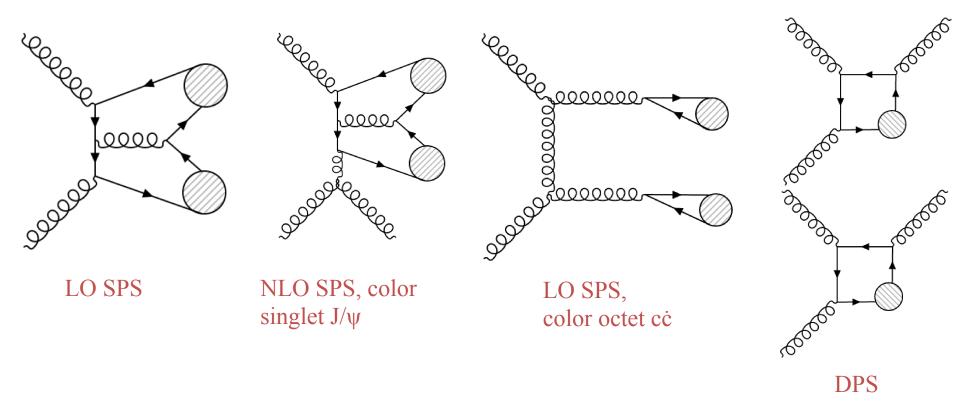
Characterization of kinematic correlations between the 2 J/ $\psi$ 's is used to extract the fraction of prompt pair events arising from double parton scattering.

Total and double parton scattering cross sections are compared with predictions.

The effective cross section of a double parton scattering is measured.

This is the first such measurement at 8 TeV, and it probes a different kinematical range from previous (1.96 TeV and 7 TeV) measurements. \*Eur. Phys. J. C (2017) 77:76.

What the events look like: each shaded circle is a  $J/\psi$  meson. Di-J/ $\psi$ 's can be produced from single parton (g-g) scattering ("SPS") or from double parton scattering ("DPS").



#### **Motivation:**

**Goal #1:** measure the *fraction of events that result from double parton scattering*. The DPS cross section is sensitive to the spatial distribution of gluons in the proton.

**Goal #2:** use the fraction of DPS events  $f_{DPS}$  to measure the effective cross section of DPS. Effective cross section is:

$$\sigma_{eff} = \frac{1}{2} \frac{\sigma_{J/\psi}^2}{f_{DPS} \cdot \sigma_{J/\psi J/\psi}}$$

It relates the production cross section of the 2 individual interactions to the total production cross section. Testing correlations of non-perturbative origin between the partons in a DPS may improve understanding of non-perturbative QCD.

**Goal #3:** DPS can be modeled and subtracted to *provide input to SPS quarkonium production models.* Quarkonium production is a background to new physics searches. Make comparisons between the data and various production models using different techniques to compute di-J/ $\psi$  production at LO, NLO, NLO color singlet NRQCD without loops (NLO\*), and intrinsic parton transverse momentum fractions.

### **Details of the analysis (1)**

- Integrated luminosity = 11.4 fb<sup>-1</sup>
- Accept prompt-prompt mesons produced directly or through feed-down from  $\psi(2S)$  decay
- Dimuon trigger, each muon's  $p_T > 4$  GeV;  $2.5 < m(\mu^+\mu^-) < 4.3$  GeV
- Reconstruction:
  - $\geq$  3 muons in the muon spectrometer data
  - Record  $|d_z|$  of 2 J/ $\psi$  decay vertices projected onto the beam axis
  - $|\eta^{\mu}| < 2.3, p_T^{\mu} > 2.5 \text{ GeV}$
  - $2.8 < m(\mu\mu) < 3.4 \text{ GeV}$
  - $|y^{J/\psi}| < 2.1, p_T^{J/\psi} > 8.5 \text{ GeV}$
- For each  $J/\psi$  candidate, find the signed transverse decay length  $L_{xy}$  (recall from page 8)
- Because J/ $\psi$  mass resolution is worse in forward region, measure cross section separately for 2 rapidity regions:  $|y^{J/\psi}| < 1.05$  ("central") and  $1.05 < |y^{J/\psi}| < 2.1$  ("forward")

### Analysis details (2)

Signal extraction procedure:

- Weight each event by efficiency of trigger, reconstruction, and selection, and by geometrical acceptance.
- First find *all* di-J/ψ events:
  - Build a 2-d distribution of the m(J/ψ<sub>1</sub>) vs. m(J/ψ<sub>2</sub>) from inclusive single J/ψ events. Signal for each J/ψ is modeled by Crystal Ball function, bkg (muons from semileptonic decays of b-hadrons and from continuum) is modeled by polynomial
  - Fit the data to this 2-d probability density function. Subtract this non-J/ $\psi$  bkg. What remains is inclusive di-J/ $\psi$  signal (prompt and non-prompt).
- To extract only prompt-prompt ("PP") events from that inclusive di-J/ $\psi$  sample:
  - Construct 2 L<sub>xy</sub> probability distributions from the inclusive J/ψ sample one for prompt-prompt and one for nonprompt-nonprompt. Prompt events have L<sub>xy</sub> consistent with resolution, non-prompt with an exponential (decay constant τ). Mixed prompt-nonprompt events are negligible.
  - Classify events according to the rapidity bins (central or forward) of the 2 J/ $\psi$ 's and apply bin-specific decay constants  $\tau$  to the exponentials.
  - For each event, plot  $L_{xy}$  of  $J/\psi_1$  versus  $L_{xy}$  of  $J/\psi_2$  and compare it to the PP and NP-NP PDFs, then classify it.
- Divide the PP-weighted PDF by the full PDF to get the likelihood that the event is PP as a function of its mesons' values for L<sub>xy</sub> and rapidity.
- Subtract pileup bkg: remove events with  $|d_z| > 1.2$  mm.

### Analysis details (3)

- Determine DPS fraction:
  - Construct a data-driven DPS template by combining  $J/\psi$  mesons from different random events in the di- $J/\psi$  sample.
  - Construct a data-driven SPS template by subtracting the DPS template from the di-J/ψ samples' Δy vs. Δφ distribution.
  - Normalize the DPS sample to the data in the region  $\Delta y > 1.8$  and  $\Delta \phi < \pi/2$  (where SPS is negligible).

• Construct weights 
$$w_{DPS(SPS)}(\Delta\phi, \Delta y) = \frac{N_{DPS(SPS)}(\Delta\phi, \Delta y)}{N_{Data}(\Delta\phi, \Delta y)}$$

- For every event, apply these weights, apply the PP weight, fit to template of m(J/ψ<sub>1</sub>) vs. m(J/ψ<sub>2</sub>) in bins of the chosen variable, extract PP SPS signal and PP DPS signals, compute f<sub>DPS</sub>.
- For this fixed f<sub>DPS</sub>, compare distribution to LO DPS and NLO\* SPS model distributions.
- Extract effective cross section

## Analysis details (4)

Corrections:

- Dimuon trigger efficiency including (1) correlations between vtx resolution and oppositesign requirement and (2) muons overlapped and unresolvable by the trigger
- Muon recon efficiency
- Kinematic acceptance (from simulation effect of  $p_T$  and  $\eta$  cuts on fiducial region)
- Signal efficiency on  $d_z$  and  $L_{xy}$ .
- $p_T$ -dependence of reconstructed mass and mass resolution of J/ $\psi$

### Systematic uncertainties:

Trigger selection, muon recon, kinematic acceptance, mass model developed from inclusive J/ $\psi$  sample, J/ $\psi$  mass and width bias function of p<sub>T</sub>, prompt-prompt model (from inclusive J/ $\psi$  sample) dependence on p<sub>T</sub>, pile-up, J/ $\psi$  to dimuon branching fraction, luminosity, DPS model and binning.

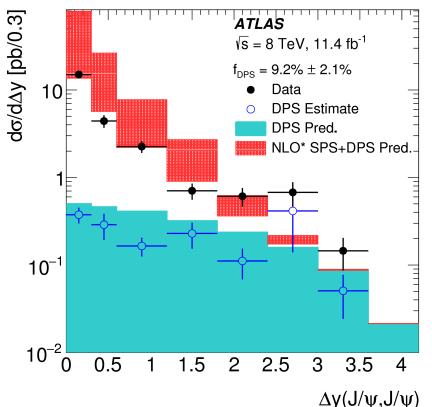
### **Results:**

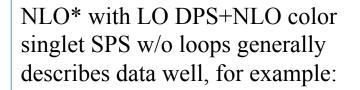
Prompt-prompt cross sections measured:

 $\sigma_{central} (J/\psi J/\psi) = 82.2 \pm 8.3 \text{ (stat.)} \pm 6.3 \text{ (syst.)} \pm 0.9 \text{ (BF)} \pm 1.6 \text{ (lumi) pb}$  $\sigma_{forward} (J/\psi J/\psi) = 78.3 \pm 9.2 \text{ (stat.)} \pm 6.6 \text{ (syst.)} \pm 0.9 \text{ (BF)} \pm 1.5 \text{ (lumi) pb}$ 

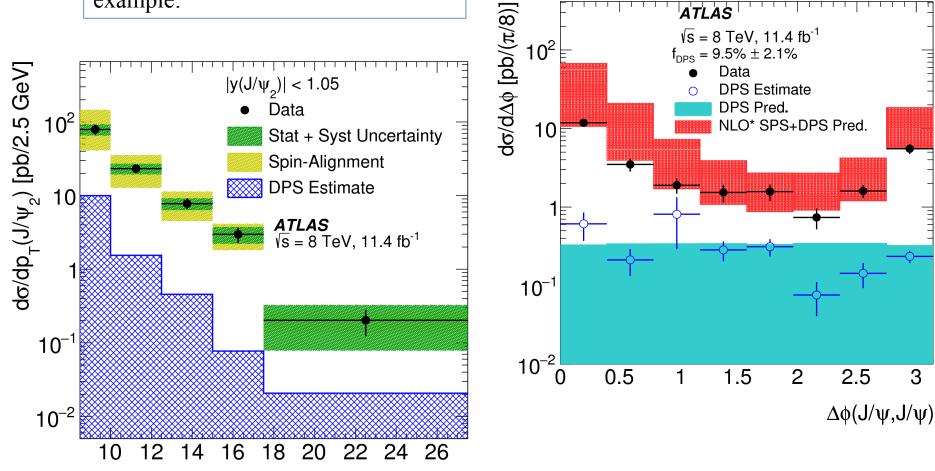
- Data are compared to theoretical distributions.<sup>§</sup> Shapes of DPS distributions are consistent with models. For SPS, the data distributions in Δy, |Δφ|, m(J/ψJ/ψ), and p<sub>T</sub>(J/ψJ/ψ) are wider than predicted by the NLO calculation.
- Data and predictions especially diverge for Δy > 1.8. This may indicate a large effect due to k<sub>T</sub> or contributions via feeddown from color-singlet ψ(2S).

<sup>§</sup>LO DPS: C. Borschensky and A. Kulesza, arXiv: 1610.00666 [hep-ph]; NLO\* SPS: J.P. Lansbert, H.S. Shao, Phys. Lett. B 751, 479 (2015) and PRL 111, 122001 (2013).

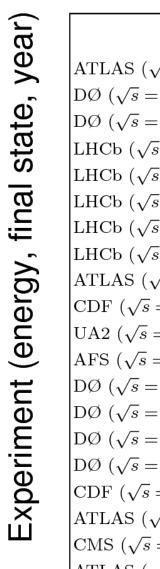




 $f_{DPS} = (9.2 \pm 2.1 \text{ (stat)} \pm 0.5 \text{ (syst)})\%$ , consistent with model predictions, for example:



 $p_{\tau}(J/\psi_{2})$  [GeV]



ATLAS

 ATLAS (
$$\sqrt{s} = 8 \text{ TeV}, J/\psi + J/\psi, 2016$$
)

 DØ ( $\sqrt{s} = 1.96 \text{ TeV}, J/\psi + J/\psi, 2014$ )

 DØ ( $\sqrt{s} = 1.96 \text{ TeV}, J/\psi + \Upsilon, 2016$ )

 LHCb ( $\sqrt{s} = 7.88 \text{ TeV}, \Upsilon(1S) + D^{0,+}, 2015$ )

 LHCb ( $\sqrt{s} = 7 \text{ TeV}, J/\psi + \Lambda_c^+, 2012$ )

 LHCb ( $\sqrt{s} = 7 \text{ TeV}, J/\psi + D^+, 2012$ )

 LHCb ( $\sqrt{s} = 7 \text{ TeV}, J/\psi + D^+, 2012$ )

 LHCb ( $\sqrt{s} = 7 \text{ TeV}, J/\psi + D^0, 2012$ )

 ATLAS ( $\sqrt{s} = 7 \text{ TeV}, 4 \text{ jets}, 1993$ )

 LHCb ( $\sqrt{s} = 7 \text{ TeV}, 4 \text{ jets}, 1993$ )

 LHCb ( $\sqrt{s} = 630 \text{ GeV}, 4 \text{ jets}, 1993$ )

 UA2 ( $\sqrt{s} = 630 \text{ GeV}, 4 \text{ jets}, 1993$ )

 UA2 ( $\sqrt{s} = 630 \text{ GeV}, 4 \text{ jets}, 1986$ )

 DØ ( $\sqrt{s} = 1.96 \text{ TeV}, \gamma + 3 \text{ jets}, 2014$ )

 DØ ( $\sqrt{s} = 1.96 \text{ TeV}, \gamma + 4 \text{ jets}, 2014$ )

 DØ ( $\sqrt{s} = 1.96 \text{ TeV}, \gamma + 3 \text{ jets}, 2014$ )

 DØ ( $\sqrt{s} = 1.96 \text{ TeV}, \gamma + 3 \text{ jets}, 2014$ )

 DØ ( $\sqrt{s} = 1.8 \text{ TeV}, \gamma + 3 \text{ jets}, 1997$ )

 ATLAS ( $\sqrt{s} = 8 \text{ TeV}, Z + J/\psi, 2015$ )

 CMS ( $\sqrt{s} = 7 \text{ TeV}, W + 2 \text{ jets}, 2014$ )

 ATLAS ( $\sqrt{s} = 7 \text{ TeV}, W + 2 \text{ jets}, 2013$ )

 $\sigma_{eff} = (6.3 \pm 1.6 \text{ (stat)} \pm 1.0 \text{ (syst)} \pm 0.1 \text{ (BF)} \pm 0.1 \text{ (lumi)}) \text{ mb}$ 

 $\sigma_{_{eff}}\,[\text{mb}]$ 

# Measurements of $\psi(2S) \rightarrow J / \psi \pi^+ \pi^-$ and $X(3872) \rightarrow J / \psi \pi^+ \pi^-$ Production<sup>\*</sup>

### Message:

Differential cross sections of X(3872) and  $\psi(2S)$  are measured and compared to models, for *prompt and non-prompt* production.

The ratio of production cross sections  $X(3872)/\psi(2S)$  is measured.

*The fraction of non-prompt X(3872) and the fraction of non-prompt \psi(2S) are measured.* 

The non-prompt X(3872) sample requires 2 lifetimes in the fit. The short lifetime component involves X(3872)'s produced in B<sub>c</sub> decays.

The *invariant mass of the dipion* system in the  $J/\psi\pi^+\pi^-$  final state is measured and found to be consistent with the process  $\rho^0 \to \pi^+\pi^-$ .

\*JHEP 01 (2017) 117.

### **Motivation:**

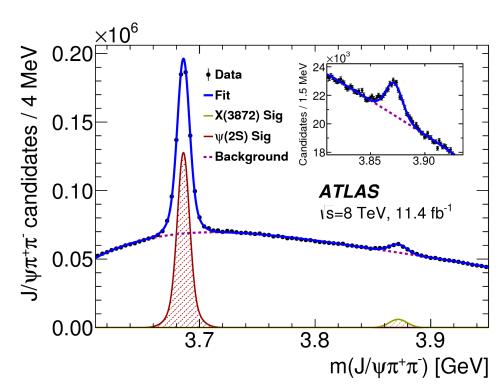
This study examines:

- the production mechanisms for these hidden charm states: direct versus feed-down from one or more heavy hadrons.
- the production mechanism of the dipion in the final state

The current best model for the X(3872) is a mixed  $\chi_{c1}(2P) - D^0 \overline{D}^{*0}$  state.

### **Details of the analysis (1)**

- Integrated luminosity = 11.4 fb<sup>-1</sup>
- Trigger: dimuons fitted to a common vertex
- Reconstruction:
  - Muons well matched to trigger objects
  - $p_T(\mu) > 4$  GeV,  $|\eta(\mu)| < 2.3$ ,  $m(\mu\mu)$  within  $m_{J/\psi} \pm 120$  MeV
- Find  $J/\psi\pi^+\pi^-$  candidates:
  - Constrain m( $\mu\mu$ ) to m(J/ $\psi$ )<sub>PDG</sub>, then assign pion masses to 2 additional oppositely charged non-muon tracks and fitted to a common vertex with the muons. p<sub>T</sub>( $\pi$ ) > 0.6 GeV, | $\eta(\pi)$ | < 2.4.
  - $J/\psi \pi^+ \pi^-$  rapidity |y| < 0.75;10 < p<sub>T</sub> < 70 GeV.
  - ΔR(J/ψ,π<sup>±</sup>) < 5: the angular distance between momenta of the dimuon system and each candidate.</li>
  - Require:  $Q = m(J/\psi\pi^+\pi^-) m(J/\psi) m(\pi\pi) < 0.3$  GeV: suppresses combinatorial bkg while saving 90% of signal.



### Analysis details (2)

- Bin candidates in p<sub>T</sub>.
- Weight each candidate for  $p_T$  and  $\eta$ -dependent selection and recon efficiencies
- Subdivide candidates in each p<sub>T</sub> bin according to pseudo-proper lifetime where

$$\tau = \frac{L_{xy}m}{cp_T}$$
$$L_{xy} = \frac{\vec{L} \cdot \vec{p}_T}{p_T}$$

 $\vec{L}$  is the vector pointing from the PV to the J/ $\psi \pi^+ \pi^-$  vertex.

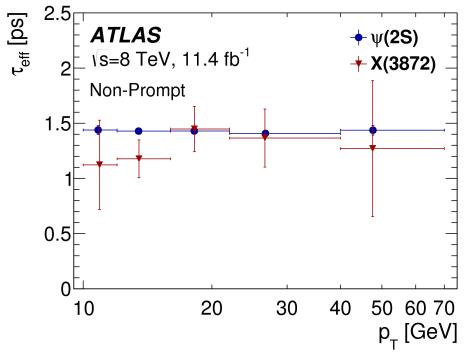
- Define 4 lifetime intervals.
- Fit data distribution in each lifetime interval to a function with 2 double-Gaussian signal functions (the  $\psi(2S)$  and the X(3872), polynomial bkg. Extract signal yields Y.

### Analysis details (3)

Apply yields Y to find double differential cross sections × branching ratios, for *i* = X(3872) or ψ(2S):

$$B(i \to J / \psi \pi^+ \pi^-) B(J / \psi \to \mu^+ \mu^-) \cdot \frac{d^2 \sigma(i)}{dp_T dy} = \frac{Y(i)}{\Delta p_T \Delta y \int L dt}$$

• First fit data in each  $p_T$  bin assuming one prompt component and one non-prompt ( $\tau_{eff}$ ) component. Observe:  $\tau_{eff}$  is different for low- $p_T X(3872)$  decays. Do these proceed by a different mechanism?



#### Analysis details (4) and results:

- Try 2 lifetimes for the non-prompt decays. (Short component: from B<sub>c</sub> decays, long component from all other B<sup>±</sup>, B<sup>0</sup>, B<sub>s</sub>, and b-baryons.)
  - Observe: no short-lived non-prompt component in  $\psi(2S)$  production.
  - For X(3872), short-lived non-prompt fraction is 25%:

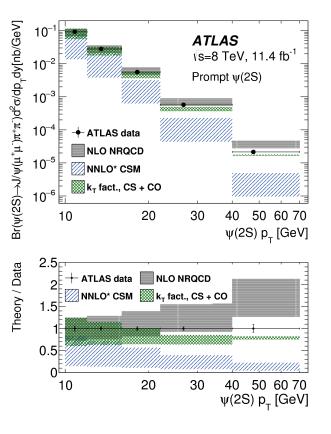
$$\frac{\sigma(pp \to B_c) \cdot B(B_c \to X(3872))}{\sigma(pp \to \text{non-prompt } X(3872))} = (25 \pm 13 \text{ (stat)} \pm 2 \text{ (syst)} \pm 5 \text{ (spin)})\%$$

Measure ratio:

$$R = \frac{B(B \to X(3872) + any) \cdot B(X(3872) \to J / \psi \pi^+ \pi^-)}{B(B \to \psi(2S) + any) \cdot B(\psi(2S) \to J / \psi \pi^+ \pi^-)} = (3.57 \pm 0.33 \pm 0.11)) \times 10^{-2}$$

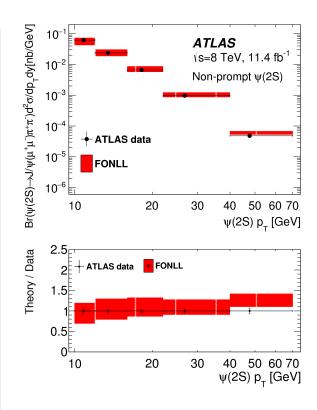
This ratio is below the value inferred from the ratio of Tevatron data (numerator) to the world average of branching fractions (denominator):  $0.18 \pm 0.08$ .

**Results on differential cross sections:** 



### For prompt $\psi(2S)$ :

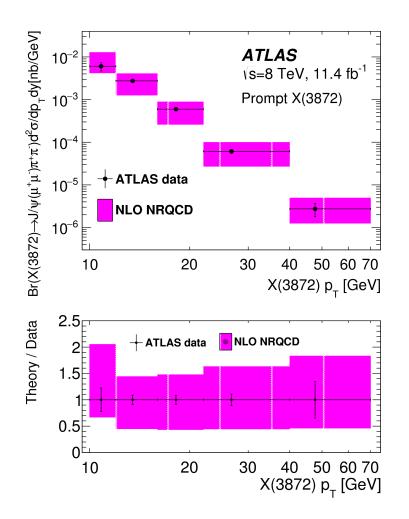
- Generally good agreement between data and NLO NRQCD using long distance matrix elements derived from Tevatron data, below highest p<sub>T</sub>
- k<sub>T</sub> factorization model including color-octet contributions tuned on 7 TeV CMS data + color singlet contributions describes data well but underestimates at highest p<sub>T</sub>.
- NNLO\* color singlet model agrees at low p<sub>T</sub>.



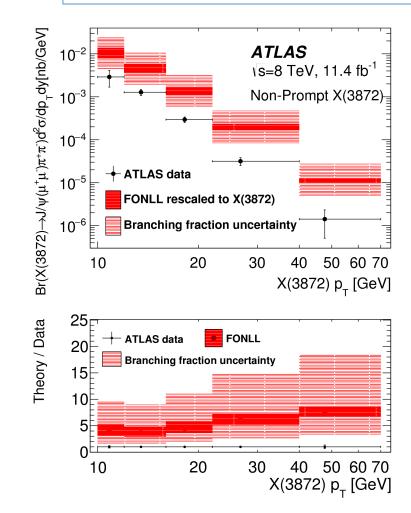
For non-prompt  $\psi(2S)$ : good agreement with FONLL over full  $p_T$  range

### Results on differential cross sections, continued:

*For prompt X(3872):* described adequately by NRQCD as  $\chi_{c1}(2P) - D^0 \overline{D}^{*0}$  mixture

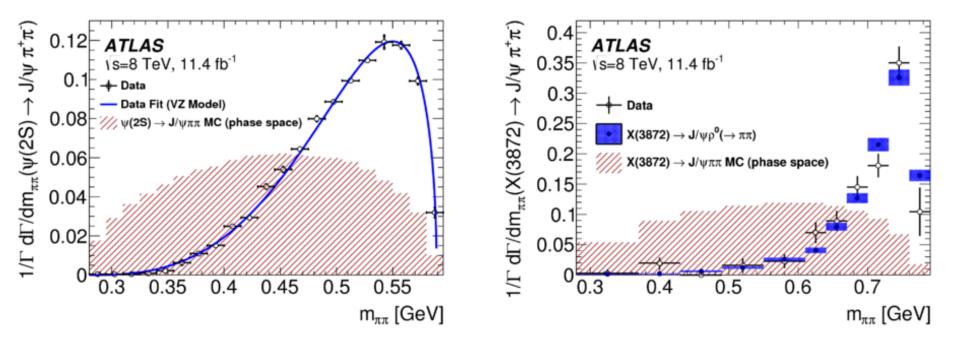


For non-prompt X(3872): FONLL model overestimates data by factor 4-8, increasing with  $p_T$ 



#### **Results on differential cross sections**, continued:

Using the normalized differential decay width in bins of dipion invariant mass, we see that phase space decay is disfavored: the pion production occurs through  $\rho^0 \rightarrow \pi^+ \pi^-$ .



# Angular Analysis of $B_d^0 \to K^* (\to K^+ \pi^-) \mu^+ \mu^-$ Decays<sup>\*</sup>

### The message:

The longitudinal polarization of the  $K^*$  is measured and compared to theoretical predictions.

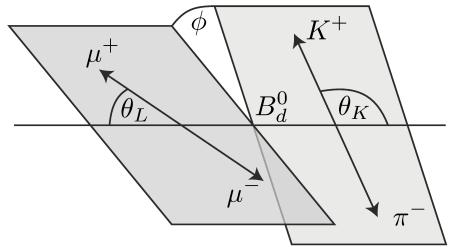
This polarization can be influenced by *penguin diagrams involving new physics*.

Hadron form factors dominate the prediction at leading order. LHCb has adopted *a method<sup>§</sup> for minimizing uncertainties in hadron form factors* in this measurement. LHCb observes<sup>¶</sup> a 3.4 sigma deviation from Standard Model calculations. The LHCb method is used here.

\*ATLAS-CONF-2017-023 (3 April 2017). \$LHCb Collaboration, PRL 111 (2013) 191801. \$LHCb Collaboration, JHEP 02 (2016) 104.

### The method:

- 3 angular variables:
  - θ<sub>K</sub>, between the K<sup>+</sup> and the direction opposite the B<sub>d</sub>, in the K<sup>\*</sup> frame
  - $\theta_L$ , between the  $\mu^+$  and the direction opposite the  $B_d$ , in the dimuon frame
  - φ, between the two decay planes formed by the Kπ and dimuon systems, in the B<sub>d</sub> frame.



• Measure:

$$\frac{1}{d\Gamma/dq^2} \frac{d^4\Gamma}{d\cos\theta_L d\cos\theta_K d\phi dq^2} = \frac{9}{32\pi} \left[ F_L \cos^2\theta_K + \frac{3(1-F_L)}{4} \sin^2\theta_K + \frac{1-F_L}{4} \sin^2\theta_K \cos 2\theta_L \right]$$
$$-F_L \cos^2\theta_K \cos 2\theta_L + S_3 \sin^2\theta_K \sin^2\theta_L \cos 2\phi$$
$$+S_4 \sin 2\theta_K \sin 2\theta_L \cos\phi + S_5 \sin 2\theta_K \sin\theta_L \cos\phi$$
$$+S_6 \sin^2\theta_K \cos\theta_L + S_7 \sin 2\theta_K \sin\theta_L \sin\phi$$
$$+S_8 \sin 2\theta_K \sin 2\theta_L \sin\phi + S_9 \sin^2\theta_K \sin^2\theta_L \sin 2\phi$$

The familiar forward-backward asymmetry is given by  $A_{FB} = 3S_6/4$ .

The method to reduce hadronic form factor dependence is this: the S<sub>i</sub> depend on the form factors and have significant uncertainty at LO. Transform the S<sub>i</sub> using ratios constructed to cancel the form factor dependence at LO:

$$P_{1} = \frac{2S_{3}}{1 - F_{L}}$$

$$P_{2} = \frac{2}{3} \frac{A_{FB}}{1 - F_{L}}$$

$$P_{3} = -\frac{S_{9}}{1 - F_{L}}$$

$$P'_{i=4,5,6,8} = \frac{S_{j=4,5,7,8}}{\sqrt{F_{L}(1 - F_{L})}}$$

All these parameters depend on the invariant mass squared of the dilepton system (q<sup>2</sup>), so analyze data in 6 partially-overlapping q<sup>2</sup> bins.

# **Analysis details:**

- Integrated luminosity =  $20.3 \text{ fb}^{-1}$
- Trigger: 1, 2, or 3 muons
- **Reconstruct muons**:  $p_T > 3.5$  GeV,  $|\eta| < 2.5$
- Require:  $\mu^+\mu^-$  reconstruct to a common vertex
- Candidate kaon, pion tracks:  $p_T > 0.5$  GeV [no dedicated particle ID in ATLAS]
- Select K<sup>\*</sup> mesons:  $p_T(K^*) > 3.0$  GeV,  $m(K\pi)$  within [846,946] MeV.
- Reconstruct B candidate:
  - Flavor assigned from K charge.
  - Vertex  $K^*$  with  $\mu^+\mu^-$
  - Require consistent vectors: vector from PV to  $B_d$  decay vertex, and  $B_d$  momentum vector
- Suppress combinatorial bkg with lifetime significance cut:  $\tau(B_d)/\sigma_{\tau} > 12.5$
- Suppress partially recon decays with tight lower cut around nominal B<sub>d</sub> mass:  $5150 < m(K\pi\mu\mu) < 5700 \text{ MeV}$
- to eliminate extra candidates per event, choose best match to  $m(K^*)_{PDG}$ , and best B vertex fit.
- $q^2$  bin range: [0.04,6.0] excluding [9.8,1.1] (to remove  $\phi$  resonance)
- Compare data to a model using maximum likelihood, for Gaussian signal, with parameters taken from a control region, and 4 bkg components
- To overcome low statistics, a "folding procedure" of transformations is used that exploits trigonometric relations among the angular parameters. 36

### The models:

- Ciuchini et al. (CFFMPSV)<sup>1</sup> QCD factorization framework to perform consistency checks of the LHCb data with theory expectations
- Descotes-Genon et al. (DHMV)<sup>2</sup> QCD factorization
- Jäger and Camalich (JC)<sup>3</sup> QCD factorization, focus on impact of long distance corrections using a helicity amplitude approach

# **Results:**

Good agreement except in 3  $q^2$  bins:  $P'_4$  and  $P'_5$  in  $q^2$  bin [4.0,6.0] and  $P'_8$  in  $q^2$  bin [2.0,4.0].

The  $P'_4$  ( $P'_5$ ) deviations are consistent with the LHCb observation and are 2.5 (2.7) sigma from the DHMV model. All measurements are within 3 sigma of the SM theory band. They are also compatible with the LHCb result.

<sup>1</sup> JHEP 06 (2016)116, arXiv: 1512.07157 [hep-ph]

<sup>2</sup> JHEP 12 (2014) 125, arXiv: 1407.8526 [hep-ph]

<sup>3</sup> JHEP 05 (2013) 043, arXiv: 1212.2263 [hep-ph]; PRD 93 (2016) 014028, arXiv: 1412.3183 [hep-ph]

The 3 noted deviations, for ATLAS and LHCb data and theoretical models

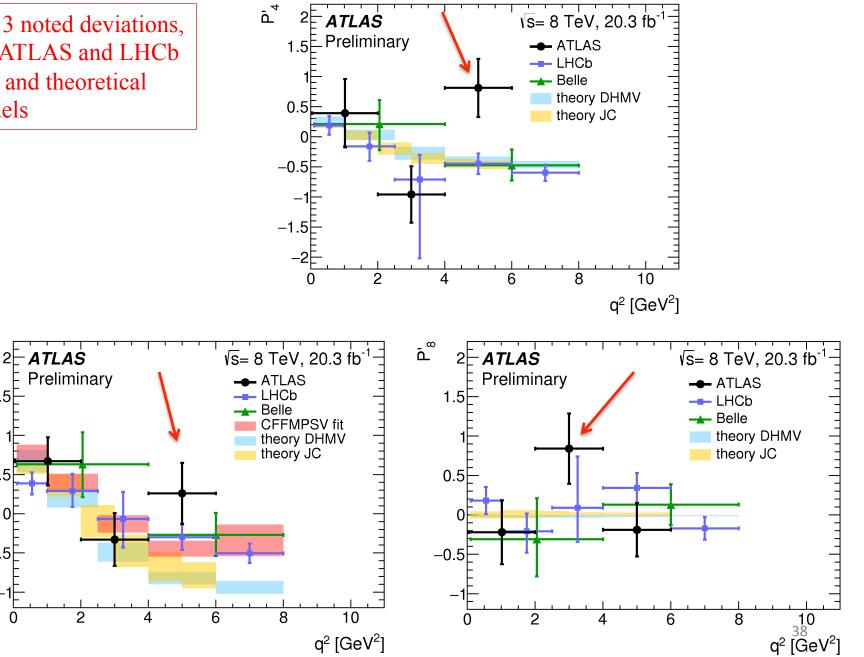
٦

1.5

0.5

-0.5

0



# Summary

ATLAS presents 4 measurements using data recorded at  $\sqrt{s} = 8$  TeV at the LHC. All are compared to contemporary models.

- Differential cross sections for b-hadron pair production to improve the theoretical description of quarkonium production and to facilitate background subtractions in new physics searches.
- Prompt J/ψ pair production differential cross sections to characterize double parton scattering as a probe of the gluon distribution in the proton, and to investigate correlations in the non-perturbative regime.
- Differential production cross sections for  $\psi(2S)$  and X(3872), both observed in decays to  $J / \psi \pi^+ \pi^-$  - a study of production mechanisms through examination of prompt and non-prompt signals.
- An angular analysis of  $B_d^0 \to K^* \mu^+ \mu^-$  decays a potential probe of new physics contributions through penguin diagrams.