

# Jets at an EIC: An Experimental Perspective

Brian Page

Brookhaven National Laboratory

Santa Fe Jets and Heavy Flavor Workshop

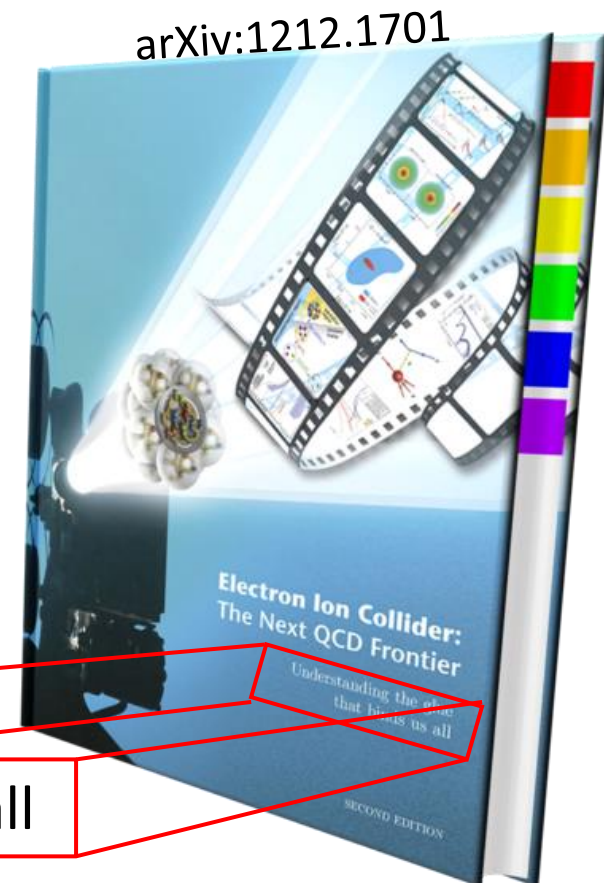
# Outline

- Brief Introduction to the Electron Ion Collider (EIC)
- Underlying Event Characteristics
- Accessing Photon Structure and Gluon Spin with Dijets
- Quark – Gluon Discrimination
- Detector Smearing

# EIC Goals in a Nutshell

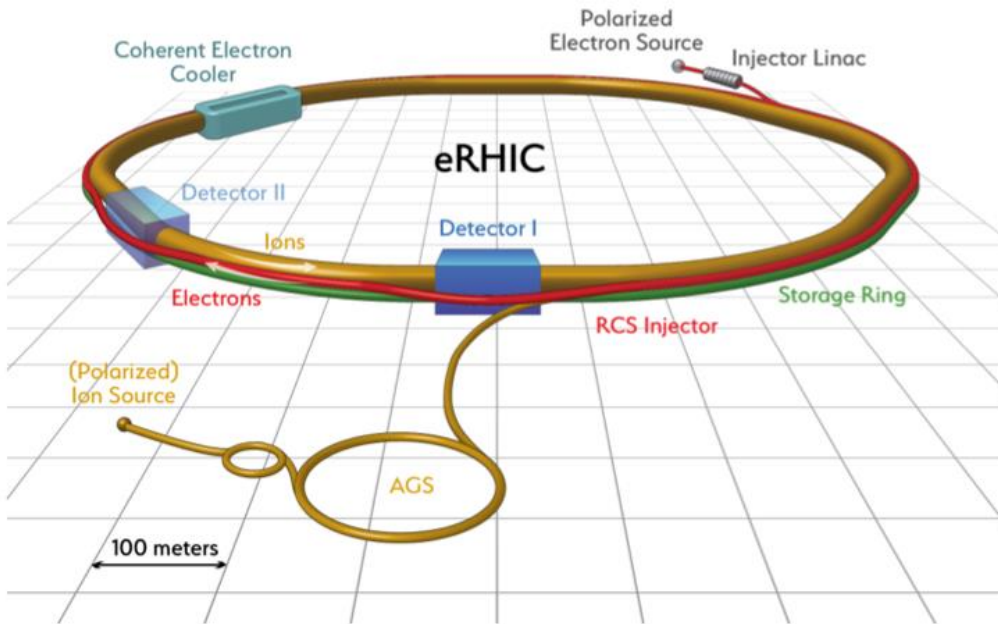
## Gain a Better Understanding of QCD via Precision Measurements of the Bound States of the Theory

- How are the sea quarks and gluons, and their spins, distributed in space and momentum inside the nucleon?
- Where does the saturation of gluon densities set in?
- How does the nuclear environment affect the distribution of quarks and gluons and their interactions in nuclei?

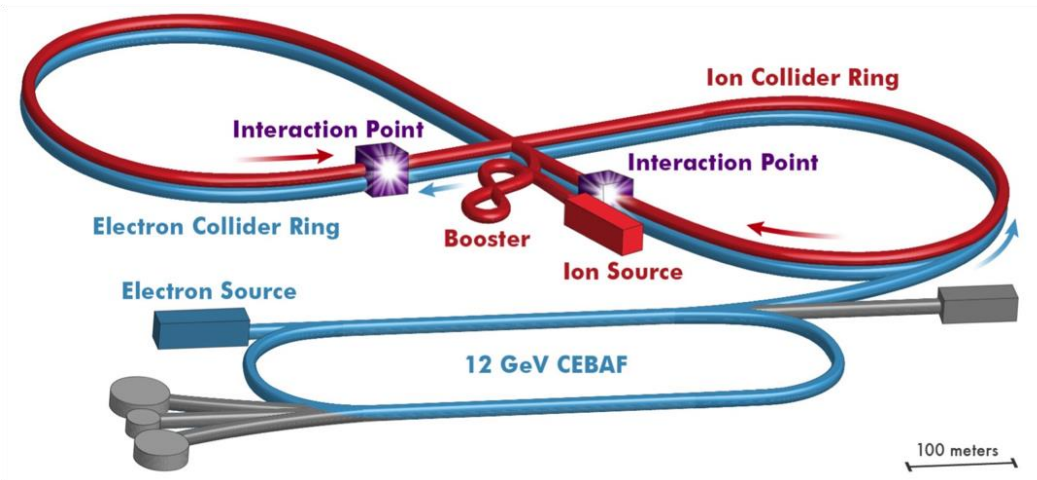


Understanding the glue that binds us all

# Potential EIC Realizations

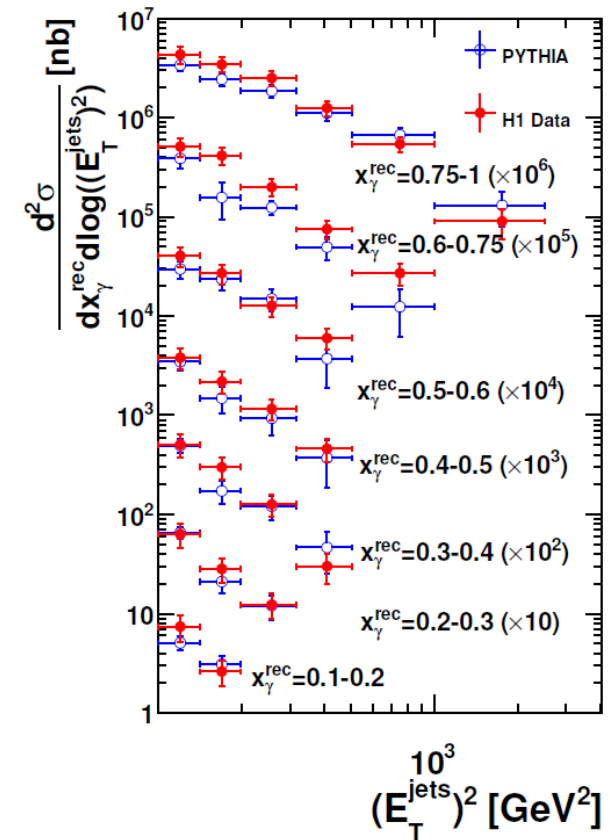


- Two designs are in active development:
  - eRHIC (BNL)
  - JLEIC (JLab)
- eRHIC utilizes the existing RHIC hadron facility and adds an electron ring and injector
- JLEIC utilizes CEBAF as an electron accelerator and adds a hadron source / booster and collider rings
- Broad tradeoff: eRHIC will start with lower luminosities but have larger center of mass energies while JLEIC will prioritize luminosity but with smaller collision energies



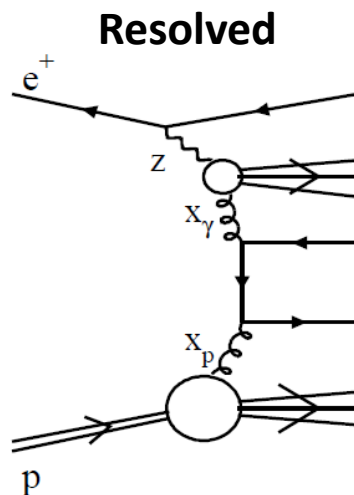
# Simulation Details / Particle Cuts

- Electron – Proton events generated at  $\sqrt{s} = 141$  GeV using PYTHIA (Full energy eRHIC design 20x250 GeV electron x proton)
- Cut on inelasticity:  $0.01 \leq y \leq 0.95$
- Jet Algorithm: Anti\_ $k_T$  ( $R = 1.0$ )
- Jets found in Breit frame
- Particles used in jet finding:
  - Stable
  - $p_T \geq 250$  MeV
  - $\eta \leq 4.5$
  - Parent cannot originate from scattered electron

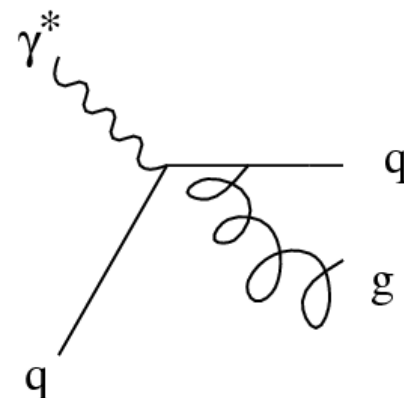


PRD 96, 074035 (2017)

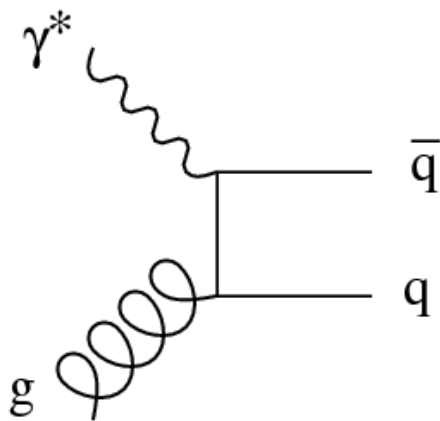
# Relevant Subprocesses



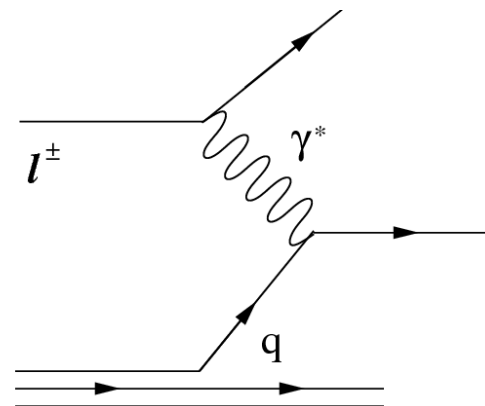
**QCD-Compton (QCDC)**



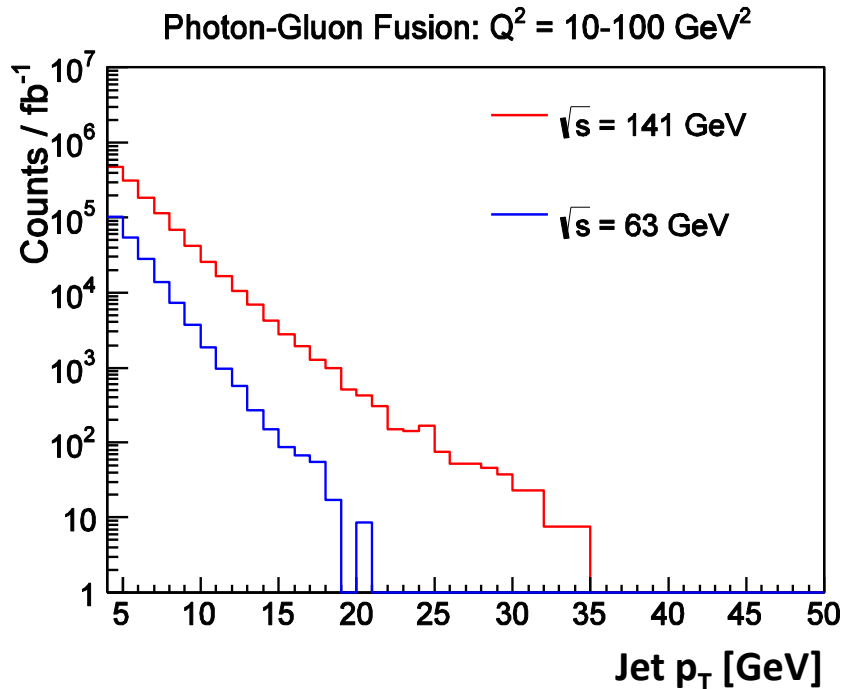
**Photon-Gluon Fusion (PGF)**



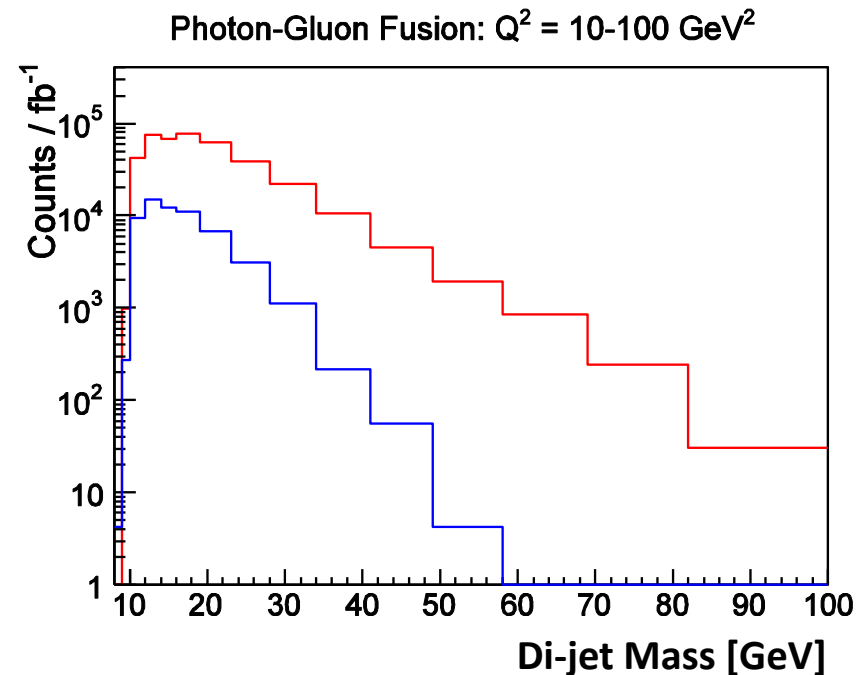
**DIS**



# Jets at an EIC: Points to Remember

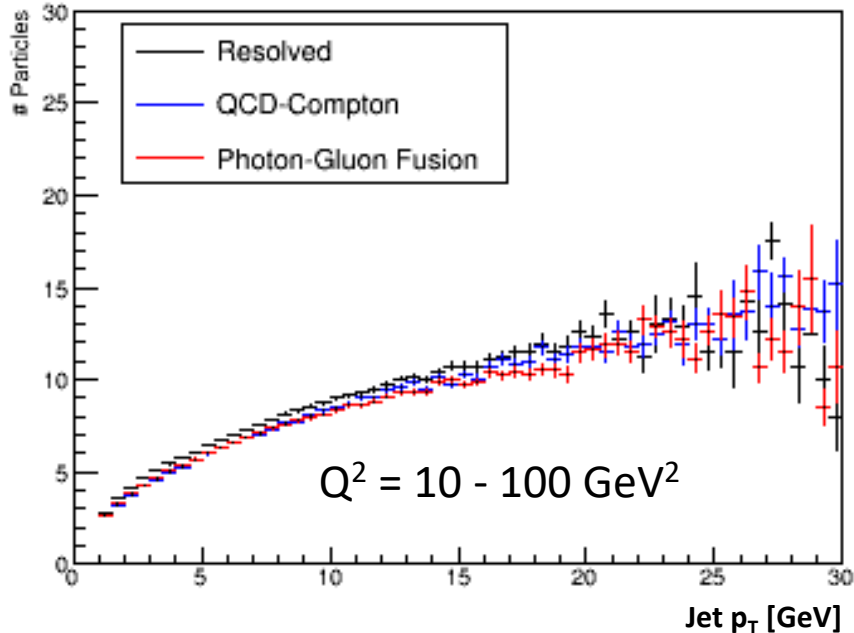


- Lower center of mass energies will lead to lower jet / di-jet yields and more limited  $p_T$  / mass reach
- Will need largest available energies and high luminosity to accumulate reasonable statistics at high  $p_T$  / mass – use  $\sqrt{s} = 141 \text{ GeV}$  for all that follows



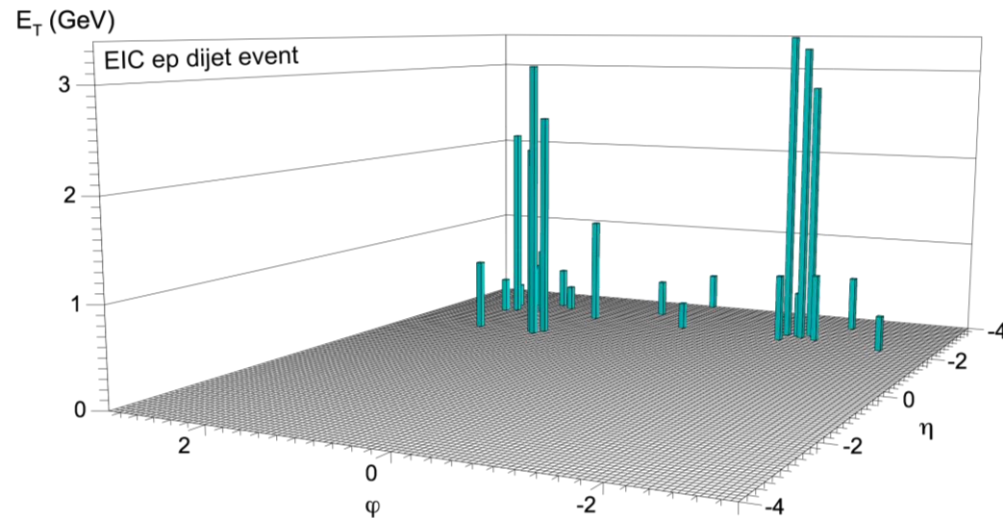
# Jets at an EIC: Points to Remember

Number of Particles in Jet Vs Jet Pt



- Jets contain relatively few particles overall
- Events should be relatively clean with moderate underlying event
- Typical particle  $p_T$  is small  $\rightarrow$  precision tracking important for reducing jet energy scale uncertainties

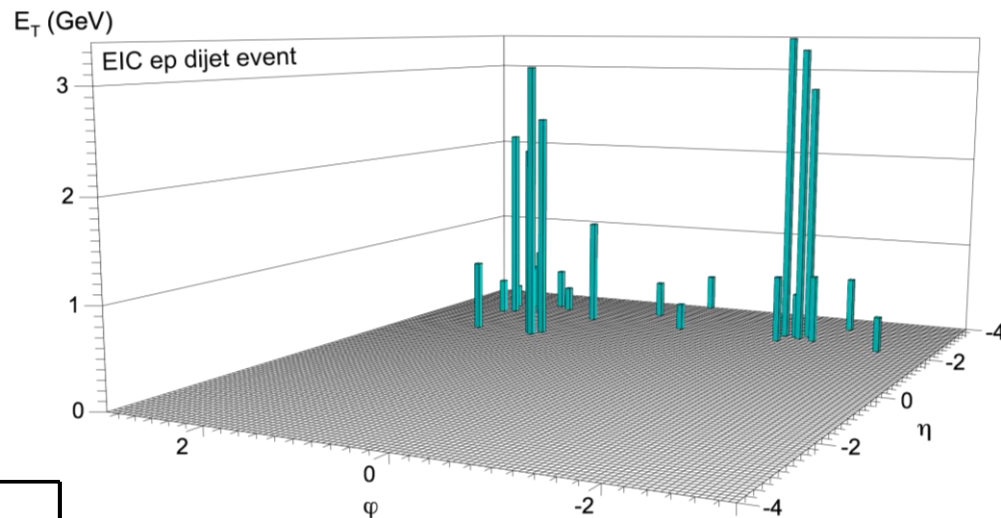
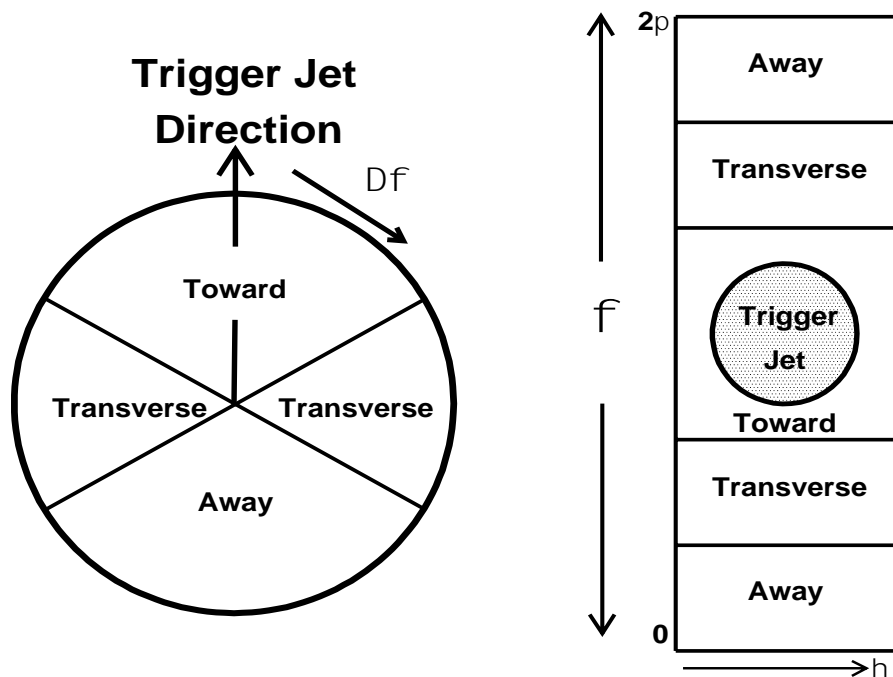
- Lower center of mass energies will lead to lower jet / di-jet yields and more limited  $p_T$  / mass reach
- Will need largest available energies and high luminosity to accumulate reasonable statistics at high  $p_T$  / mass – use  $\sqrt{s} = 141$  GeV for all that follows





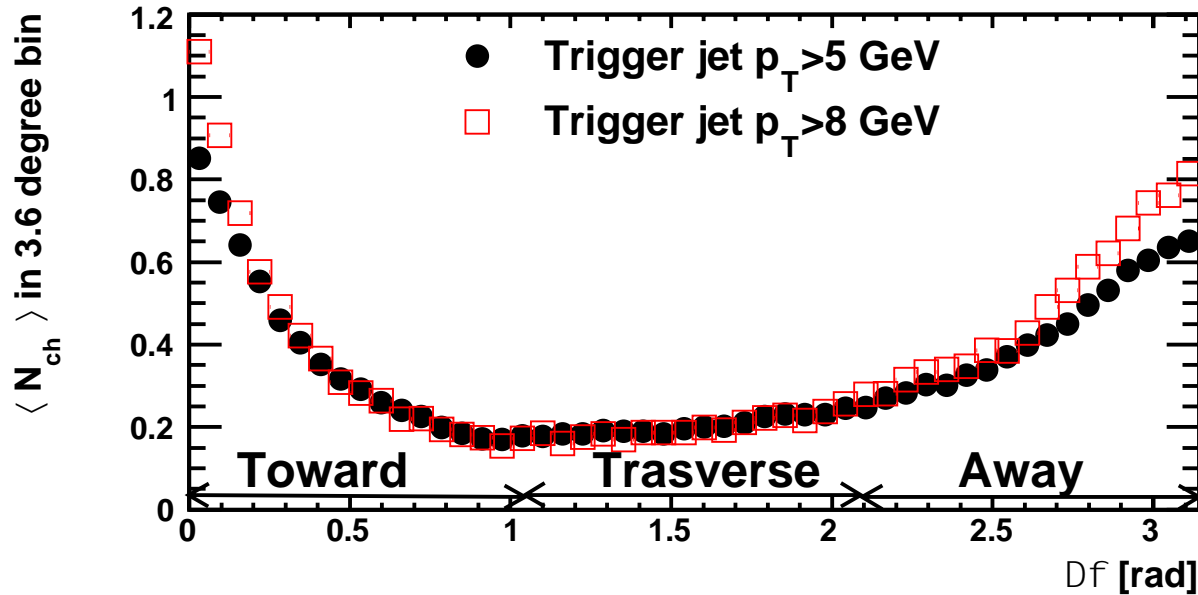
# Underlying Event Study

- ep events are expected to be relatively clean, with moderate underlying event activity
- Want to systematically quantify the amount of underlying event present in a typical event



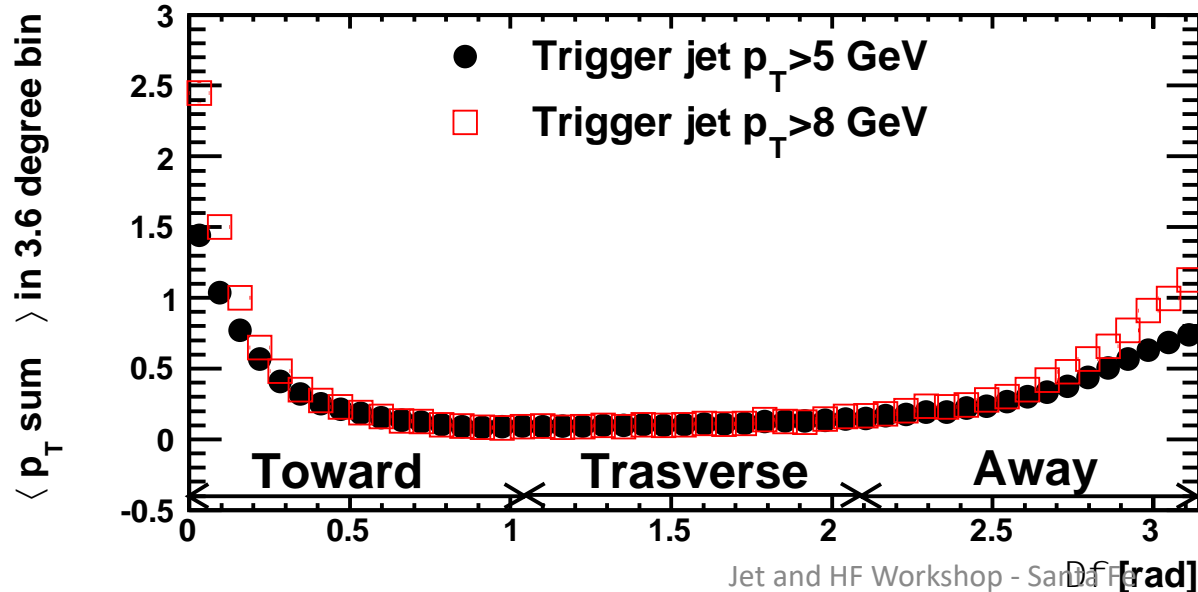
- Divide event into regions based on position of a trigger jet
- Transverse regions sensitive to underlying event contribution
- For this study: Dijet events from Resolved, QCDC, and PGF subprocesses;  $Q^2 < 1 \text{ GeV}^2$ ;  $p_{T1} > 5$ ,  $p_{T2} > 4.5 \text{ GeV}/c$

# Underlying Event Characteristics



- Plot average number of charged particles per event as a function of azimuthal angle from trigger jet

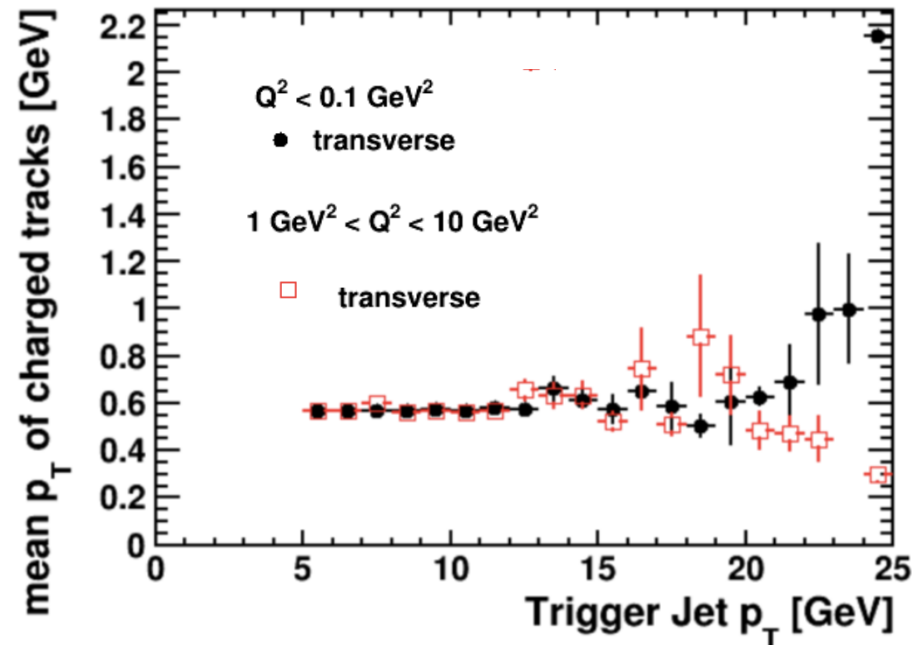
- Also plot the average summed particle  $p_T$



- See little dependence on trigger jet  $p_T$

- The number of charged particles and  $p_T$  sum in transverse region is small

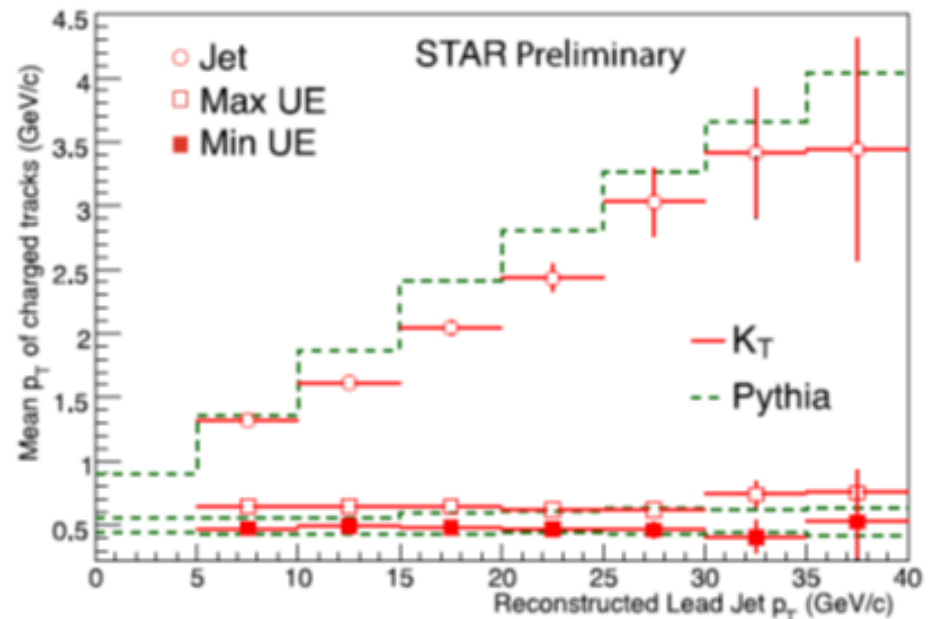
# Comparison with STAR



- See similar behavior in 200 GeV pp events at STAR
- Can we use STAR data to study certain EIC jet observables?

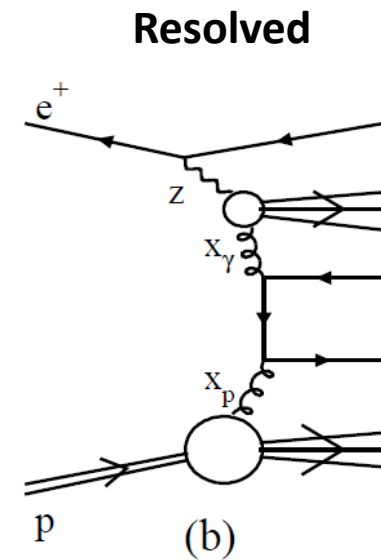
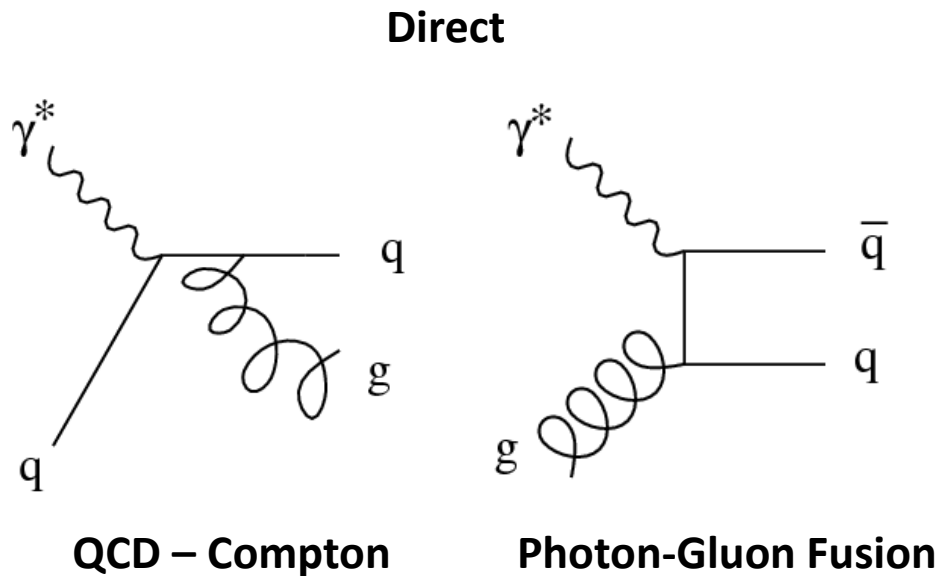
- Plot the average  $p_T$  for charged tracks as a function of trigger jet  $p_T$
- See that these quantities are independent of the trigger jet  $p_T$  in transverse region as well as  $Q^2$

arXiv:1107.4891



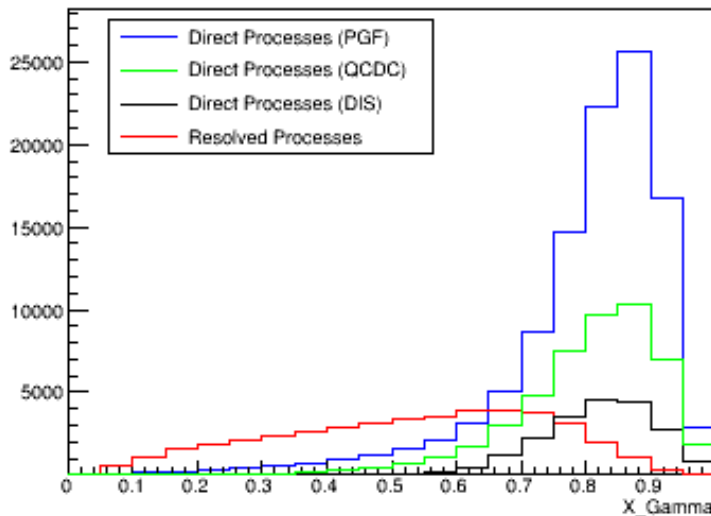
# Jets as Parton Surrogates

- Jets should approximate the energy and momentum of the partons from which they arise allowing the reconstruction of event kinematics such as  $x_\gamma$  (photon momentum fraction) and  $x_p$  (parton momentum fraction) among many other applications
- $x_\gamma$  will allow tagging of direct vs resolved subprocesses which will be important for studies of photon structure (Phys. Rev. D 96, 074035) as well as alternative methods for accessing  $\Delta G$

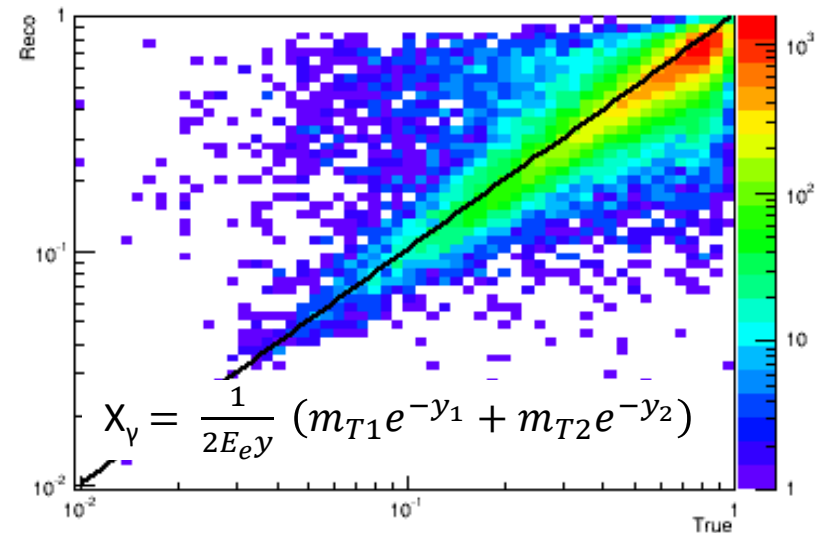


# Subprocess Tagging and Kinematics

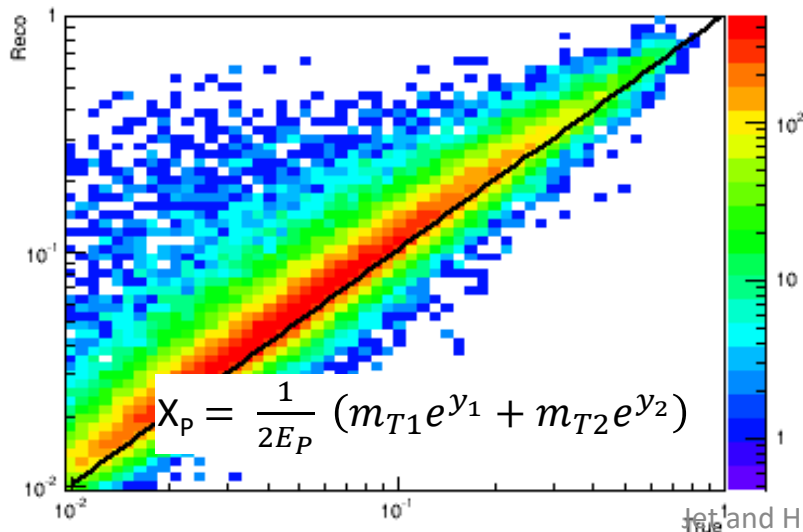
Reconstructed X\_Gamma: Q2 = 10-100 GeV^2



Reco Vs True X\_Gamma: hQCD: Q2 = 10-100 GeV^2

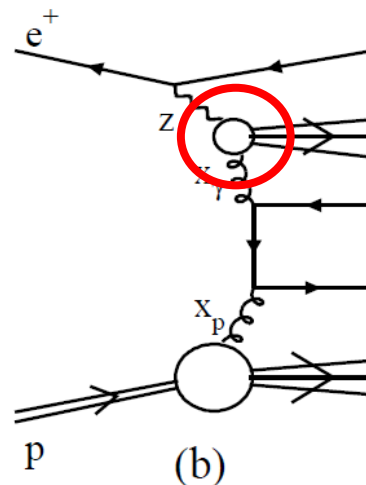
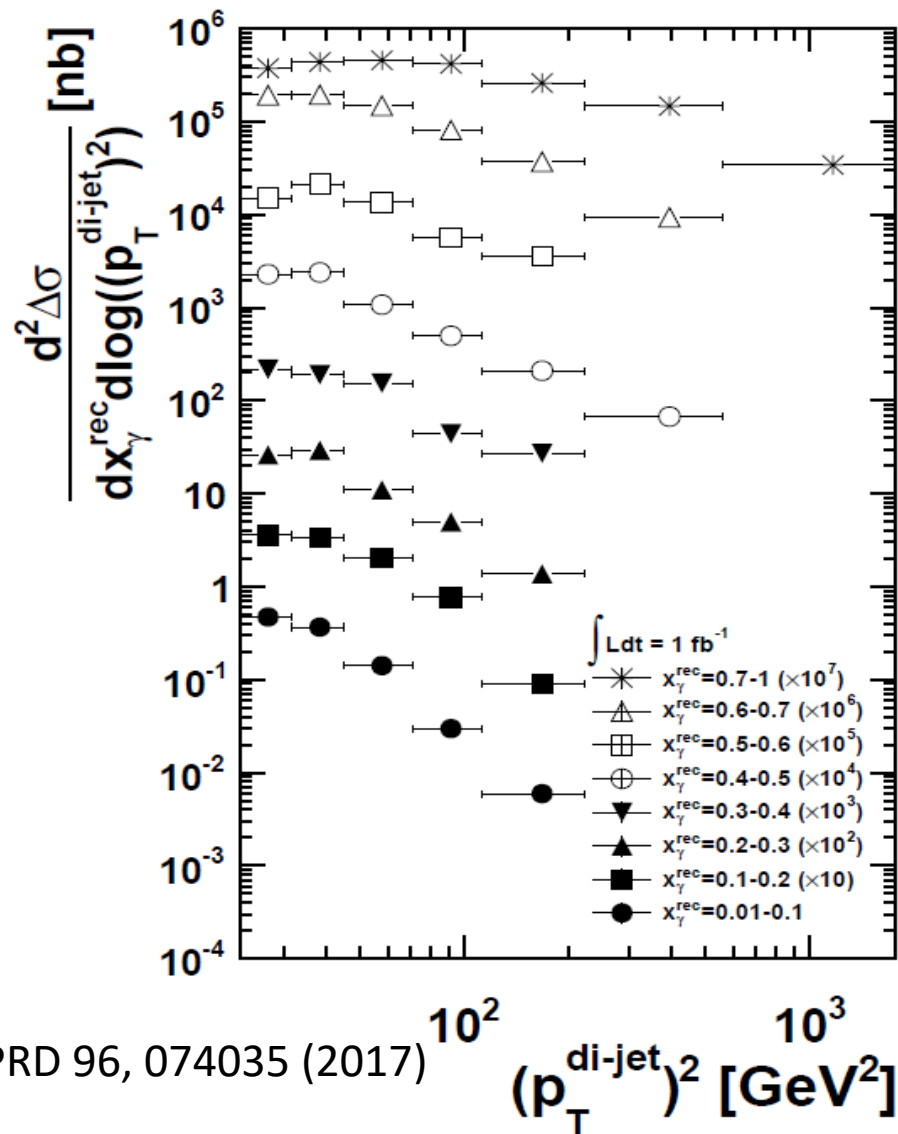


Reco Vs True X\_Proton (X\_Gamma > 0.8): PGF: Q2 = 10-100 GeV^2



- Use dijet energy and momentum to reconstruct  $x_y$  and  $x_p$
- Cutting on  $x_y$  can enhance or reduce resolved contribution (which becomes more prominent at low  $Q^2$ ) depending on the analysis needs
- Both  $x_y$  and  $x_p$  accurately reconstructed

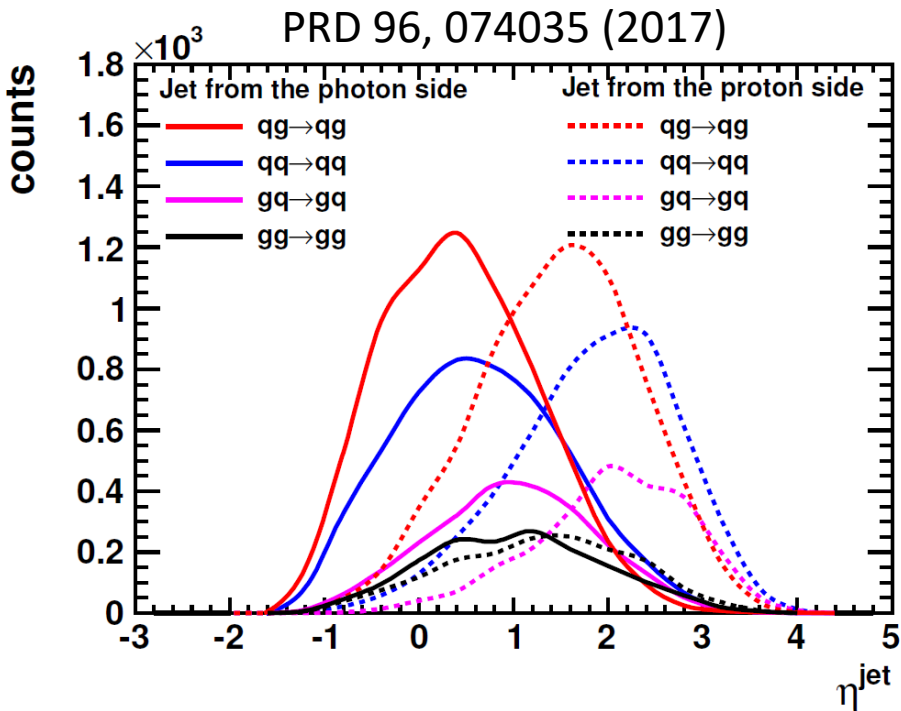
# Example: Photon Structure



Study the polarized and unpolarized hadronic structure of the photon

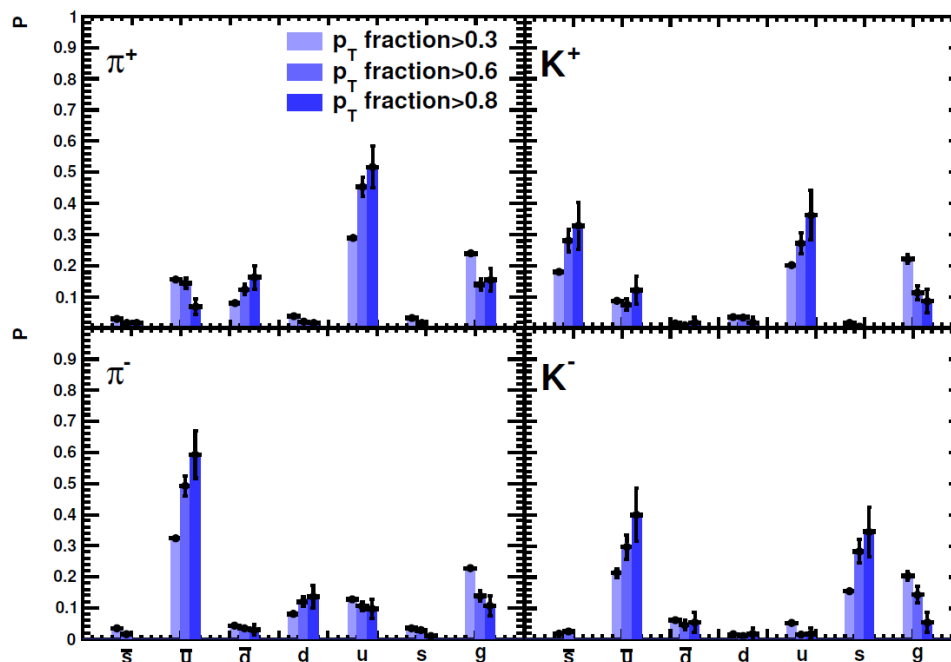
- In QCD, the photon can be considered a superposition of a bare photon state and a hadronic state
- Want to characterize the polarized and unpolarized structure of this hadronic state (photon PDFs)
- EIC cross section data will allow very precise extractions of these PDFs and give access to the polarized structure for the first time

# Flavor Tagging



- Can tag the highest  $p_T$  hadron inside the jet associated with the photon to enhance certain flavors
- See  $\pi^+$  and  $\pi^-$  enhance  $u$  and  $u$ -bar fractions while kaons enhance  $u/u$ -bar and  $s/s$ -bar
- Take advantage of the excellent PID capabilities of the planned EIC detectors

- Would also like to look more differentially and constrain photon PDFs for different parton flavors
- See that the jet associated with the photon preferentially goes to lower pseudorapidities

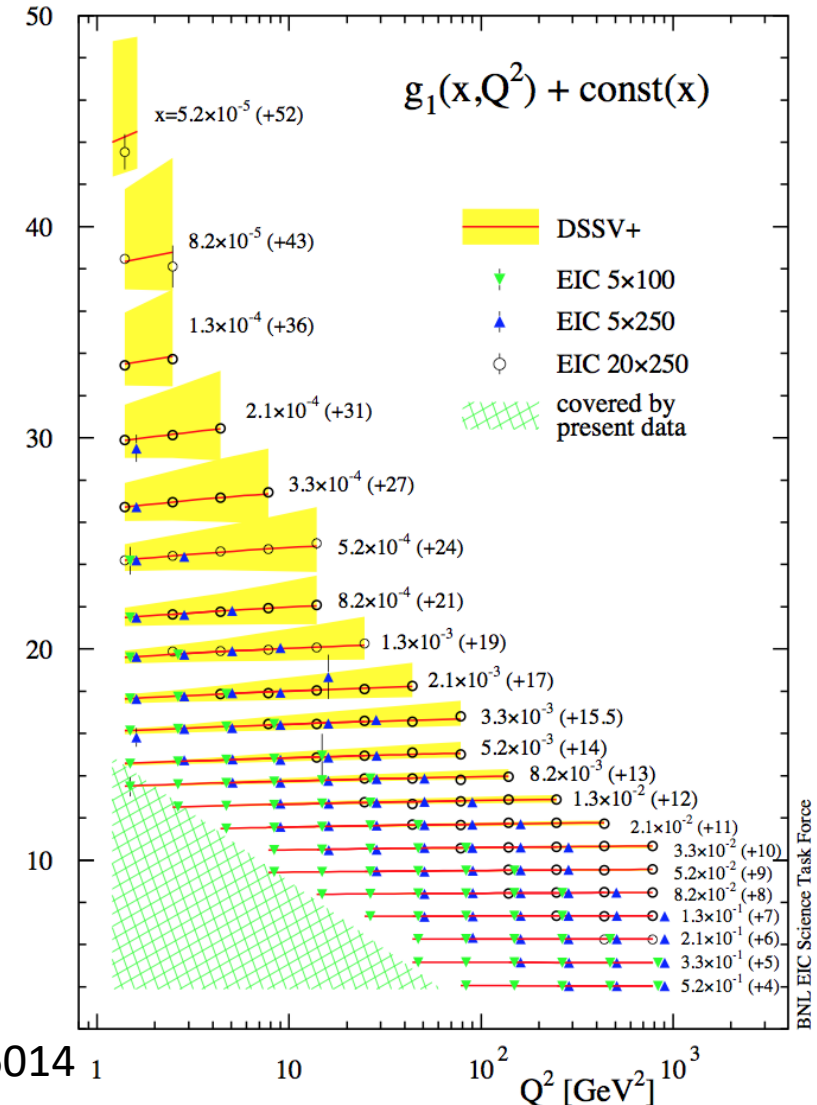


# Example: Accessing $\Delta G$ with Dijets

- Several observables are sensitive to  $\Delta G$  in DIS but golden measurement at an EIC would be scaling violation of  $g_1(x, Q^2)$

$$\frac{dg_1(x, Q^2)}{d\ln(Q^2)} \approx -\Delta g(x, Q^2)$$

- Can also get access to  $\Delta G$  by using dijets to tag the photon-gluon fusion process, providing a cross-check and allowing studies of the evolution of  $\Delta G$  with respect to  $Q^2$
- Reconstruction of  $x_\gamma$  will facilitate rejection of resolved events  $x_p$  will help isolate PGF from the quark-induced QCD-Compton process

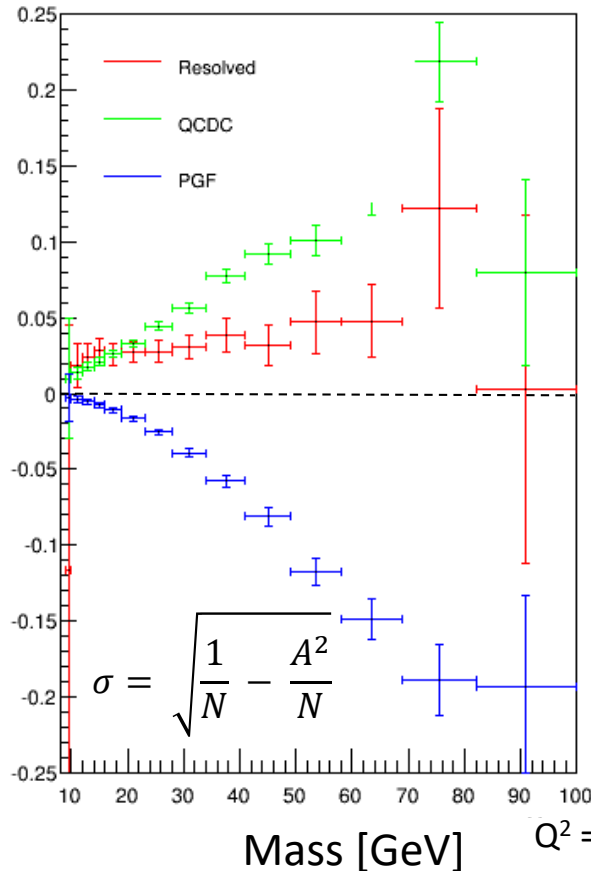


arXiv:1206.6014

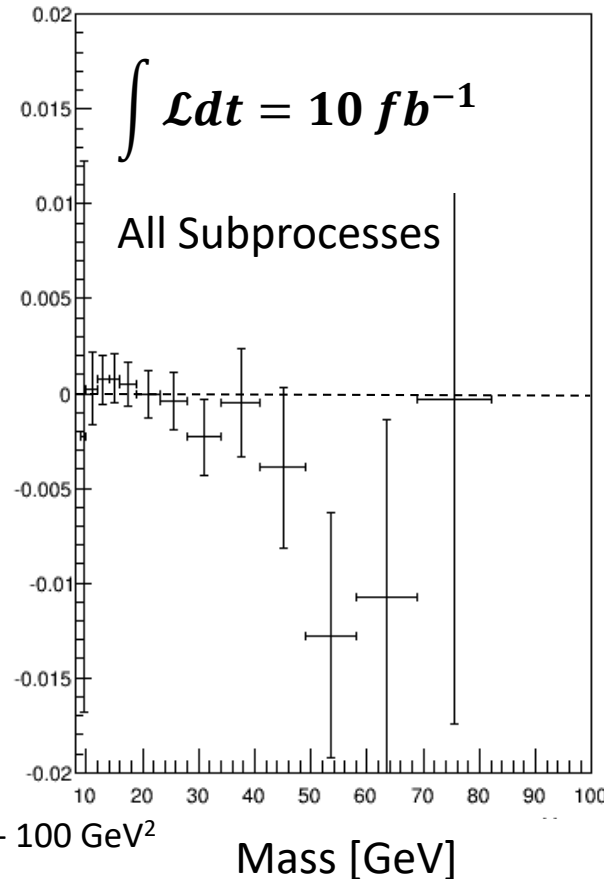


# $A_{LL}$ Vs Di-jet Mass

$A_{LL}$  Vs Di-jet Mass



$A_{LL}$  Vs Di-jet Mass

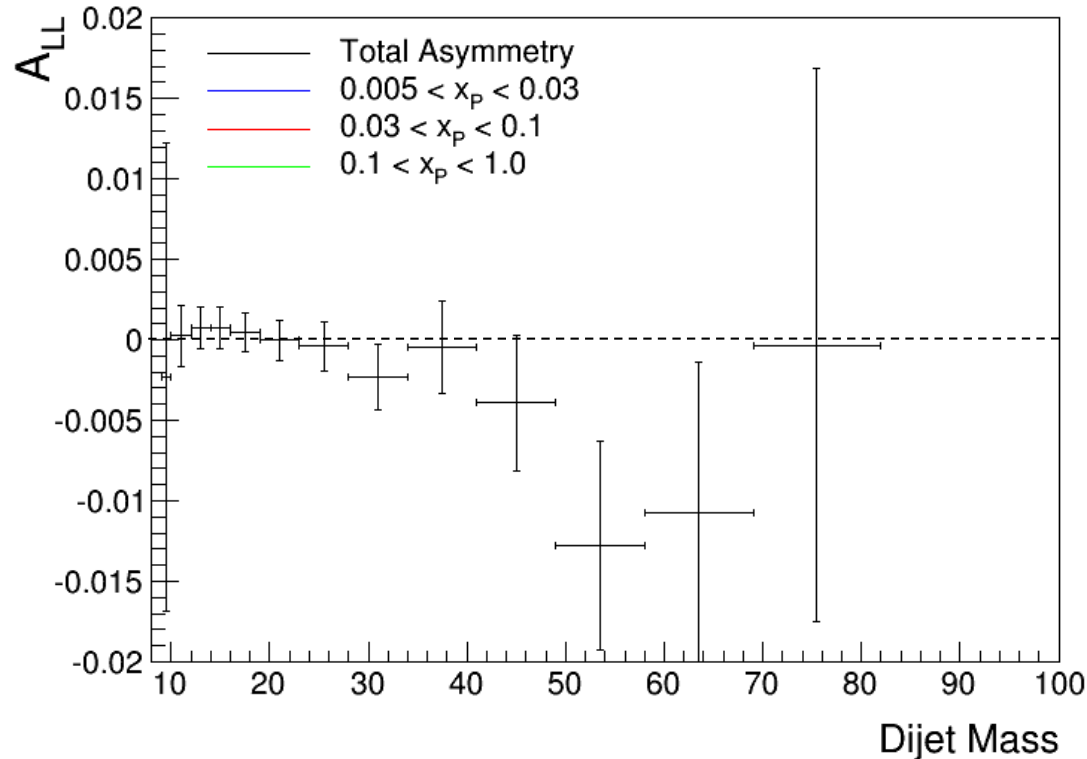


- Weight simulated events by product of the partonic asymmetry and the ratios of the polarized over unpolarized photon and proton PDFs to obtain realistic estimate of  $A_{LL}$
- Plot the expected  $A_{LL}$  as a function of di-jet invariant mass for each sub-process separately as well as the combined sample
- PGF asymmetry is nearly canceled out by QCDC asymmetry with opposite sign – would like to reduce QCDC contribution

$$w = \hat{a}(s, t, \mu^2, Q^2) \cdot \frac{\Delta f_a^{\gamma^*}(x_a, \mu^2)}{f_a^{\gamma^*}(x_a, \mu^2)} \cdot \frac{\Delta f_b^N(x_b, \mu^2)}{f_b^N(x_b, \mu^2)}$$

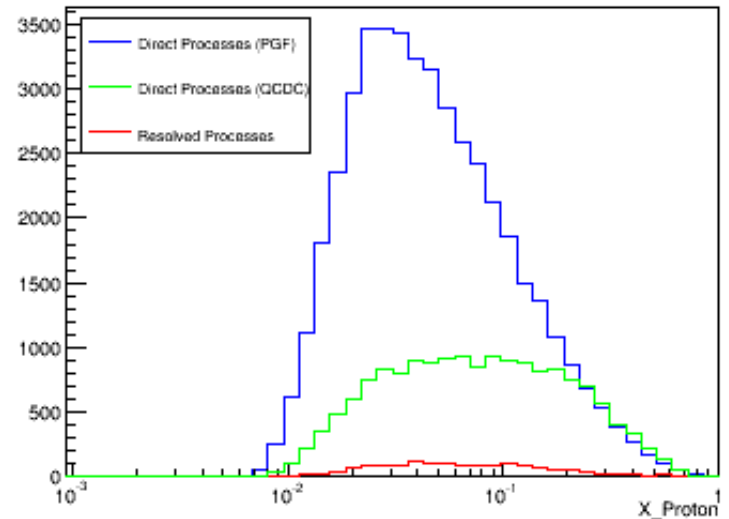
# $A_{LL}$ Vs Di-jet Mass: $x_p$ Cuts

Total  $A_{LL}$  Vs Dijet Mass:  $Q^2 = 10-100 \text{ GeV}^2$

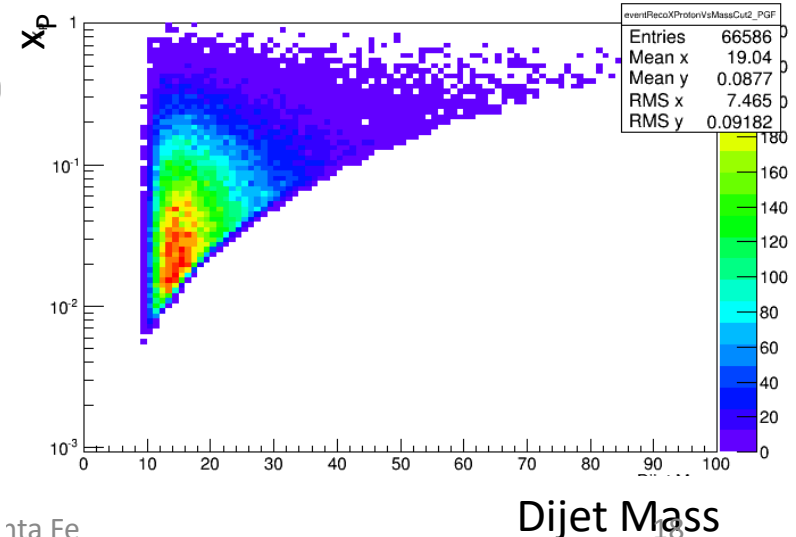


- QCDC and PGF asymmetries largely cancel out making overall asymmetry small
- Want to enhance PGF subprocess w.r.t. QCDC
- PGF events peaked to lower  $x_p$  values

Reco X Proton ( $X_{\text{Gamma}} \geq 0.8$ ):  $Q^2 = 10-100$

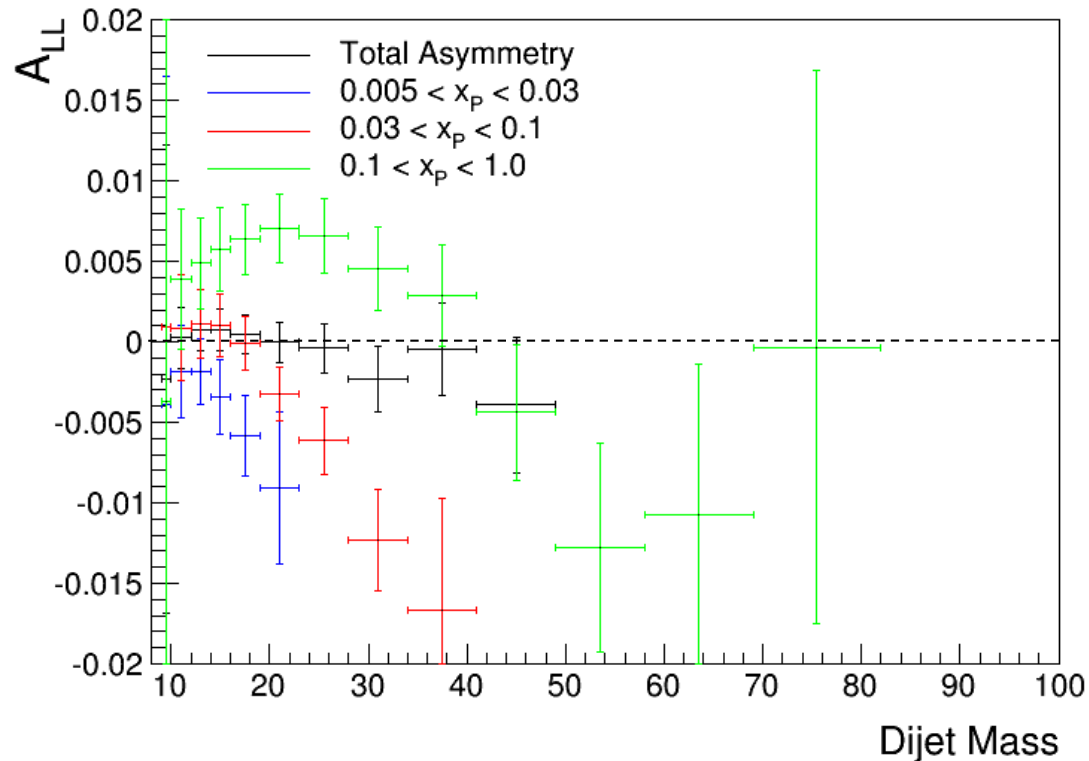


Reco X Proton Vs Dijet Mass  $Q^2 = 10-100 \text{ GeV}^2$ : PGF



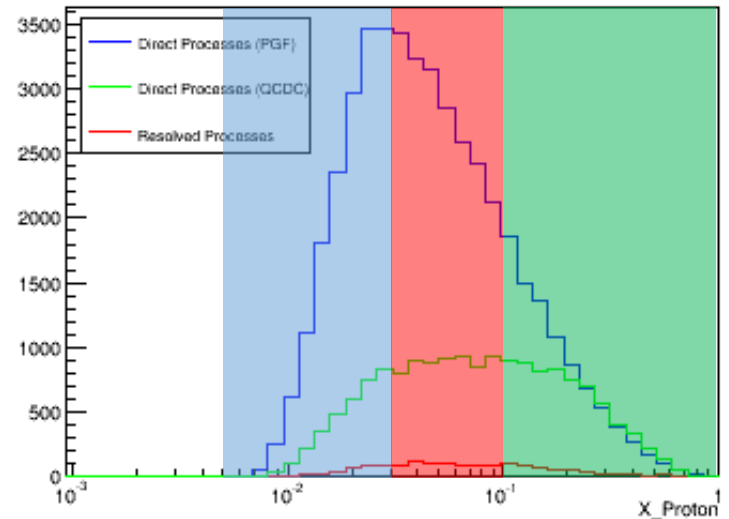
# $A_{LL}$ Vs Di-jet Mass: $x_p$ Cuts

Total  $A_{LL}$  Vs Dijet Mass:  $Q^2 = 10-100 \text{ GeV}^2$

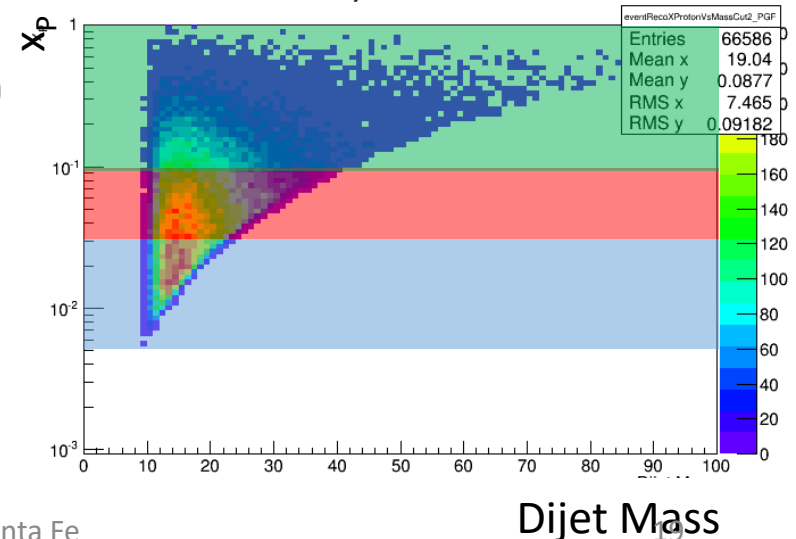


- Selecting events with  $0.005 < x_p < 0.03$  enhances PGF asymmetry but restricts mass range
- Intermediate  $x_p$  values get more QCDC contribution
- Largest  $x_p$  values have roughly equal amounts of PGF and QCDC

Reco X Proton ( $X_{\text{Gamma}} \geq 0.8$ ):  $Q^2 = 10-100$



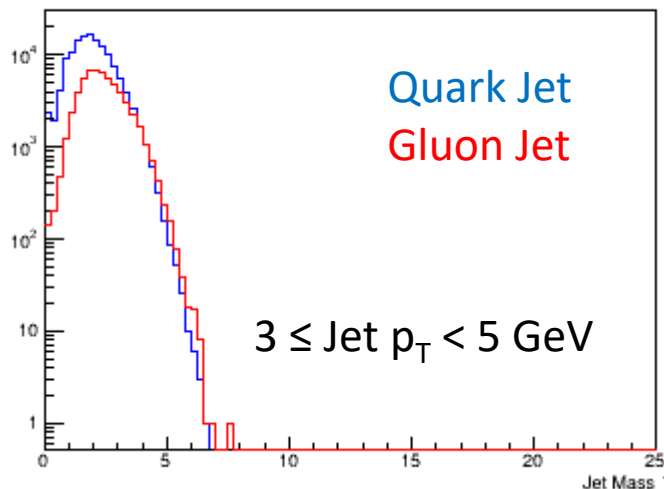
Reco X Proton Vs Dijet Mass  $Q^2 = 10-100 \text{ GeV}^2$ : PGF



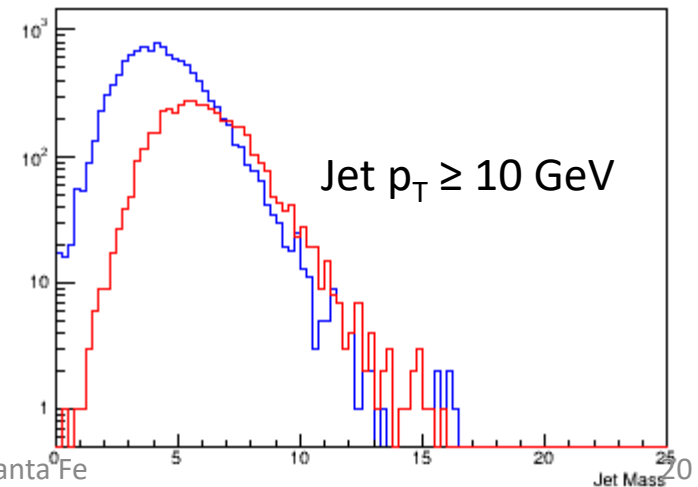
# Quark – Gluon Discrimination

- Can we use the distribution of energy within a jet to determine if that jet arose from a quark or a gluon? Possibility to tag QCD-Compton process via detection of a gluon
- This is a preliminary look at jet substructure at eRHIC; eventually want to explore the utility of substructure for studying how partons lose energy and hadronize in the cold nuclear medium
- For this study, look at jets with  $p_T \geq 10$  GeV as this is where separation between quark and gluon jets is seen. Only consider light quarks: u, d, and s

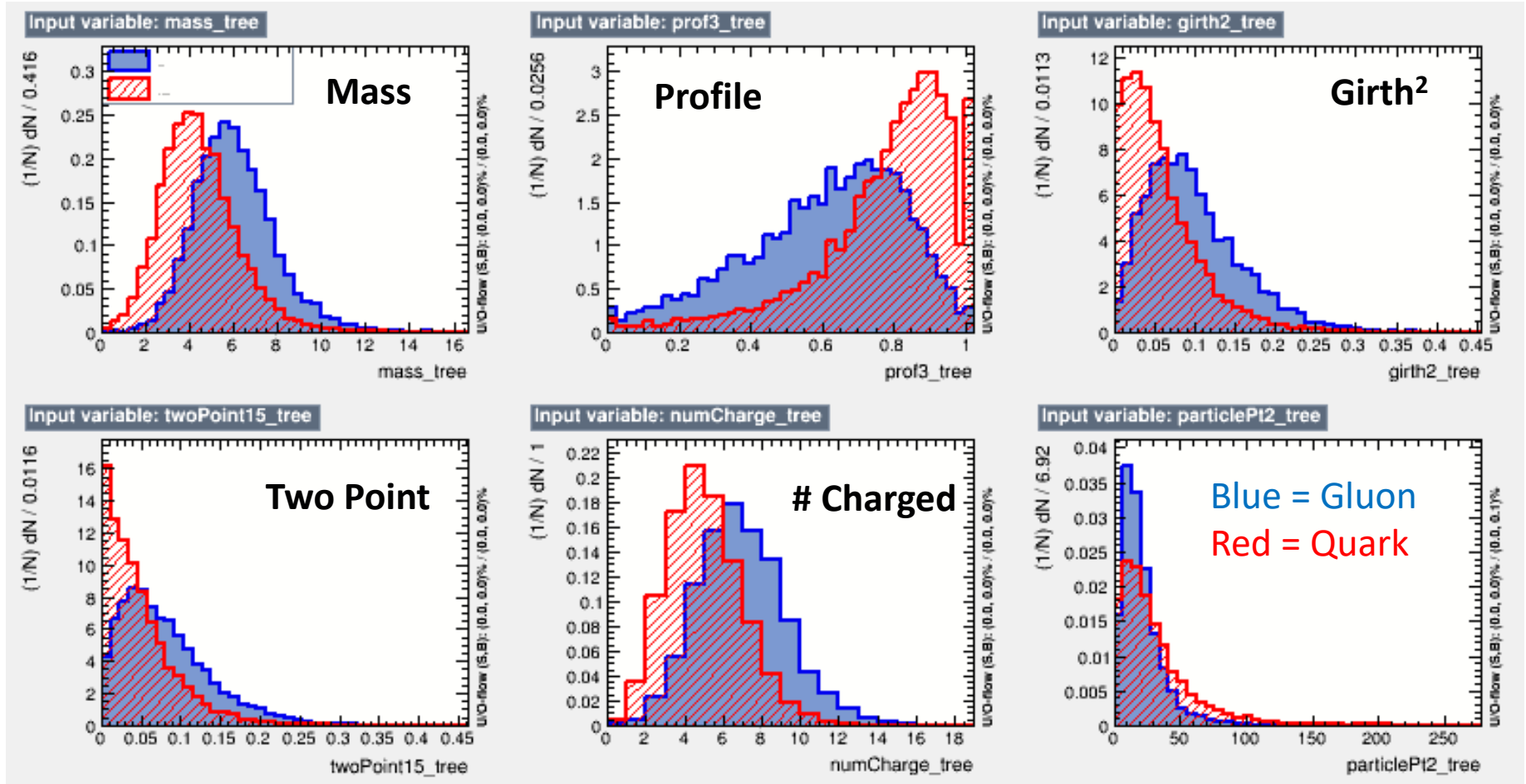
Jet Mass: Low Jet  $p_T$



Jet Mass: High Jet  $p_T$



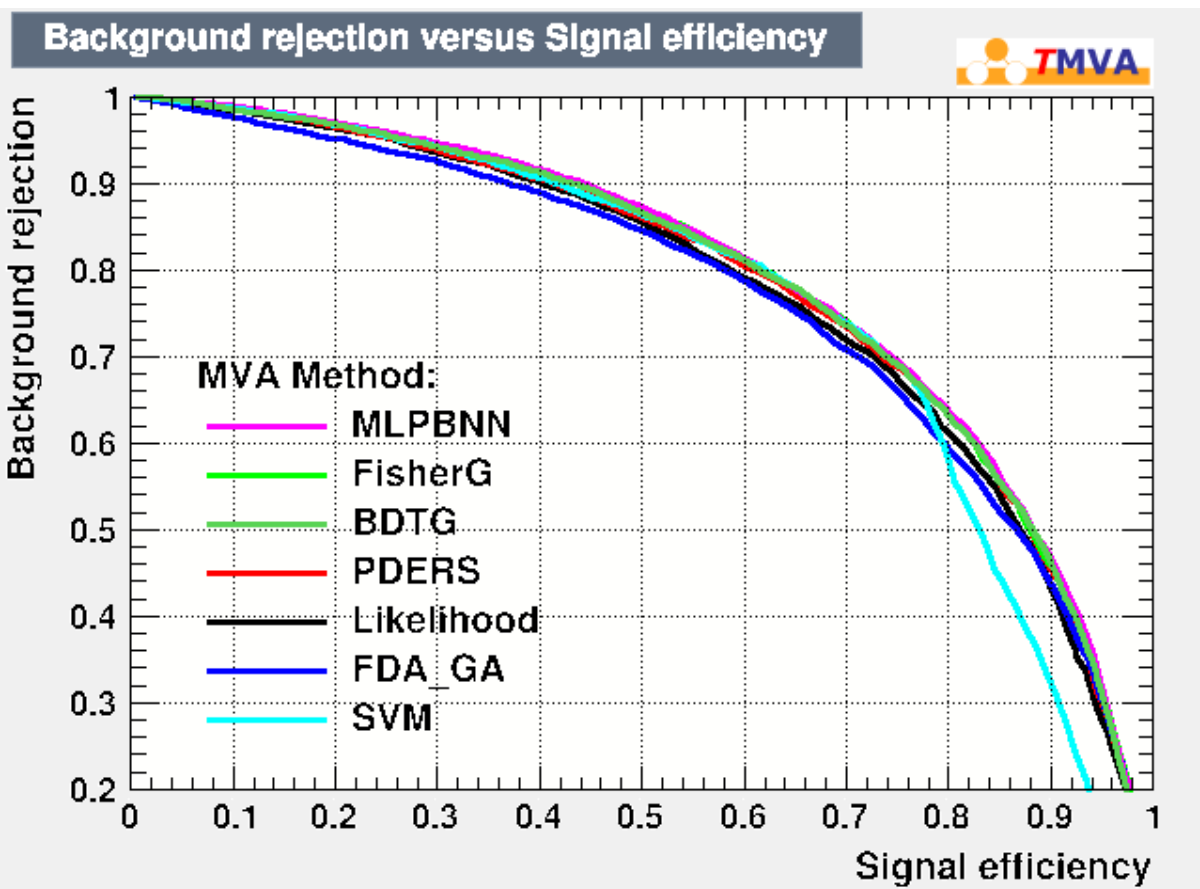
# Input Variables



$$\text{Girth}^2 = \sum_i \frac{p_{Ti}}{p_{Tjet}} |r_i|^2$$

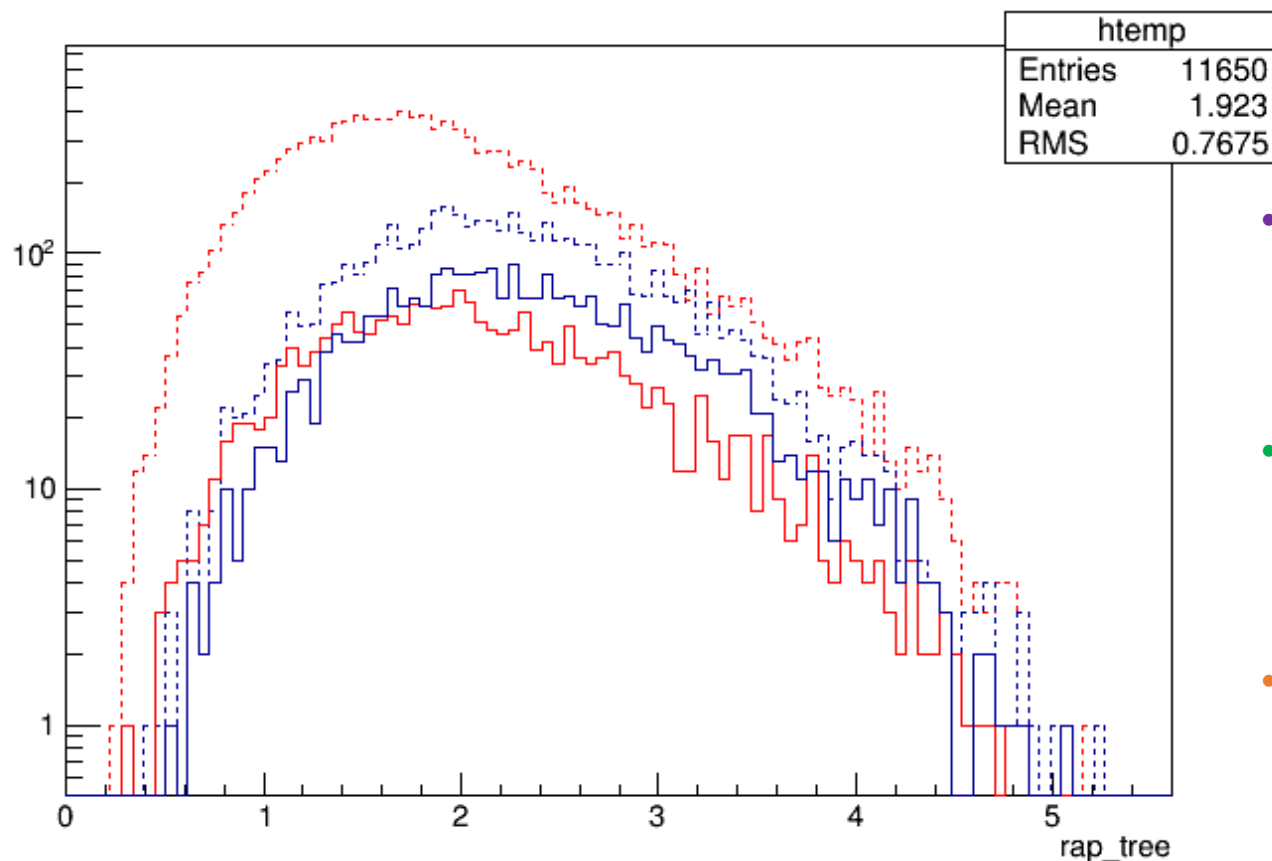
$$\text{2 Point} = \frac{1}{p_{Tjet}^2} \sum_{i \neq j} p_{Ti} * p_{Tj} * |r_{ij}|^\beta$$

# Discrimination Performance



- Characterize a number of multivariate methods by percentage of background (quarks) rejected vs signal (gluons) retained
- All methods performed roughly the same
- For the following, use MLPBNN which is a neural network implementation

# Jet Rapidity Spectra



- After cut is applied, can plot quark and gluon jets vs any relevant variable
- Here we see that gluons dominate at higher rapidity
- Look at jets with rapidity  $> 1.8$  to further enhance gluon fraction

Dotted Red = All Quarks (11650)  
Dotted Blue = All Gluons (4511)  
Solid Red = Quarks After Cut (1964)  
Solid Blue = Gluons After Cut (2568)

G/Q Before Cut = 0.39  
G/Q After Cut = 1.31  
G/(G+Q) Before = 28%  
G/(G+Q) After = 57%

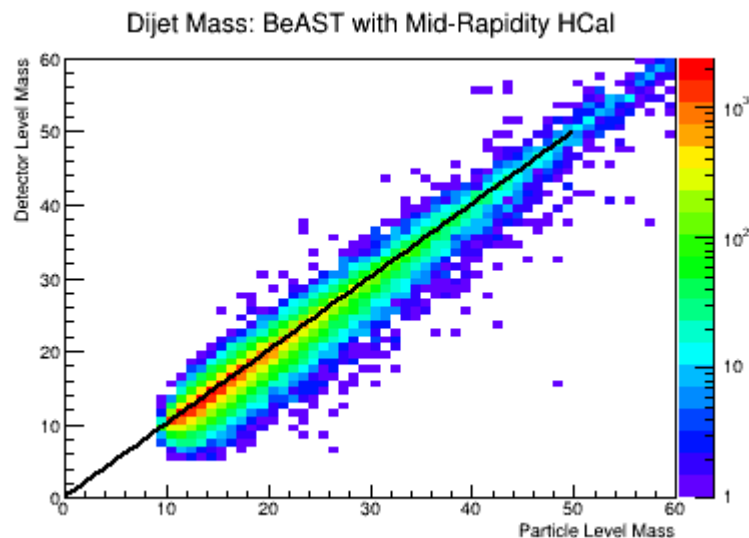
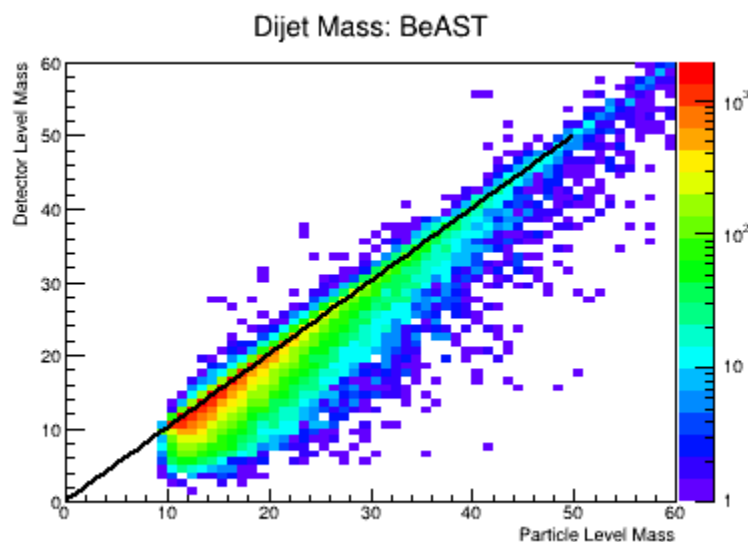
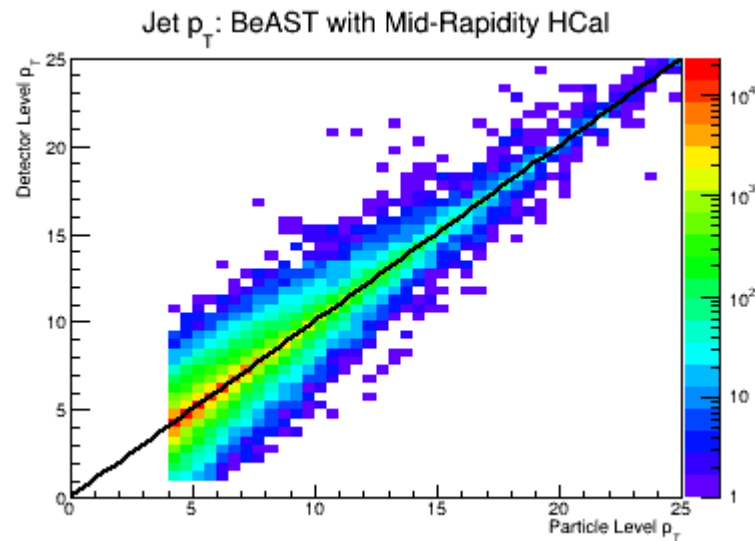
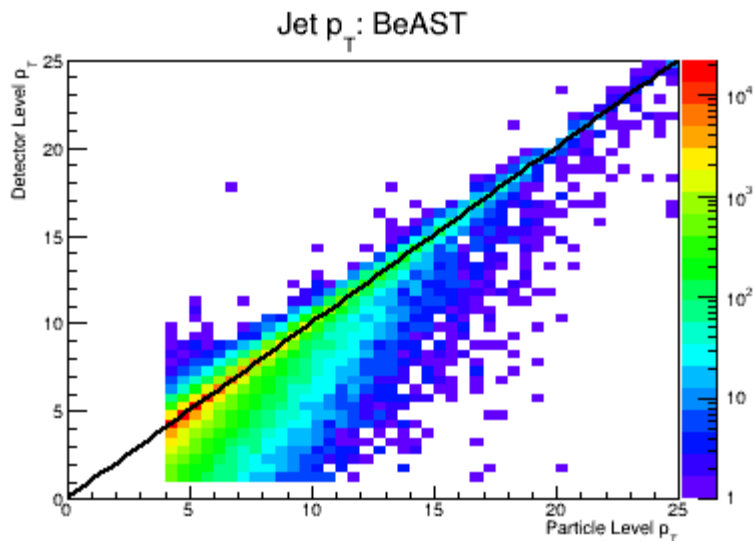
# Simulating Detector Response

- Need to study the effect that detector smearing will have on jet observables and kinematic quantities reconstructed from jets
- Do this via a fast smearing program based on the BeAST detector design
- Electromagnetic calorimeter coverage spans range of  $\pm 4.5$  in pseudorapidity with resolution between 7 and 15%/VE
  - EM clusters assumed to be massless
- Tracking coverage spans range of  $\pm 3.5$  in pseudorapidity
  - Tracks assumed to have pion mass
  - Tracking inefficiencies of 5 and 10% also considered
- Hadronic calorimetry in forward / backward region ( $1 < |\eta| < 4.5$ ) with resolution of 1.5% $\oplus$ 50%/VE
  - Optional mid-rapidity calorimeter with 7% $\oplus$ 85%/VE also considered



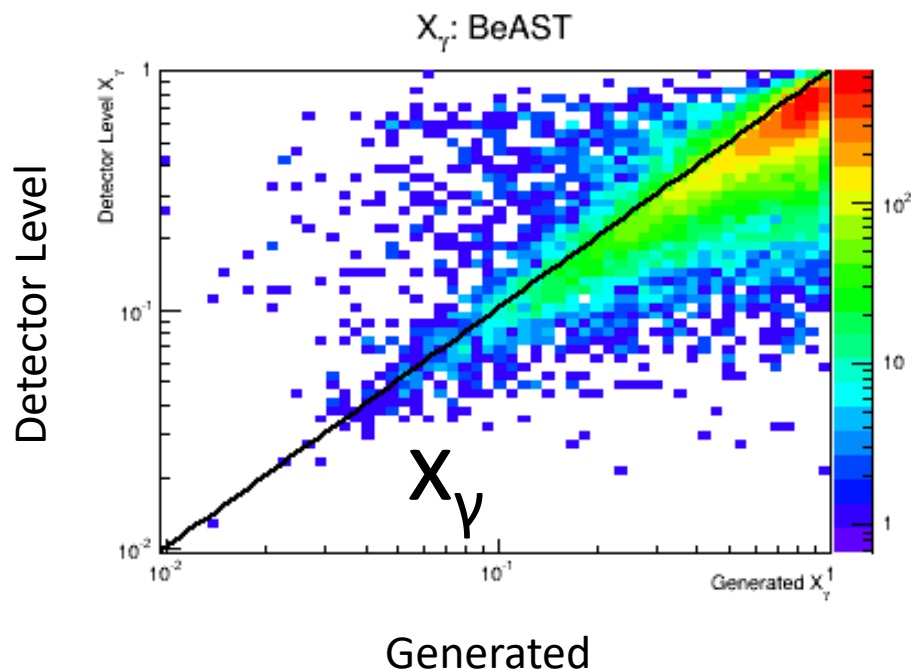
# Particle – Detector Comparison

Detector Level



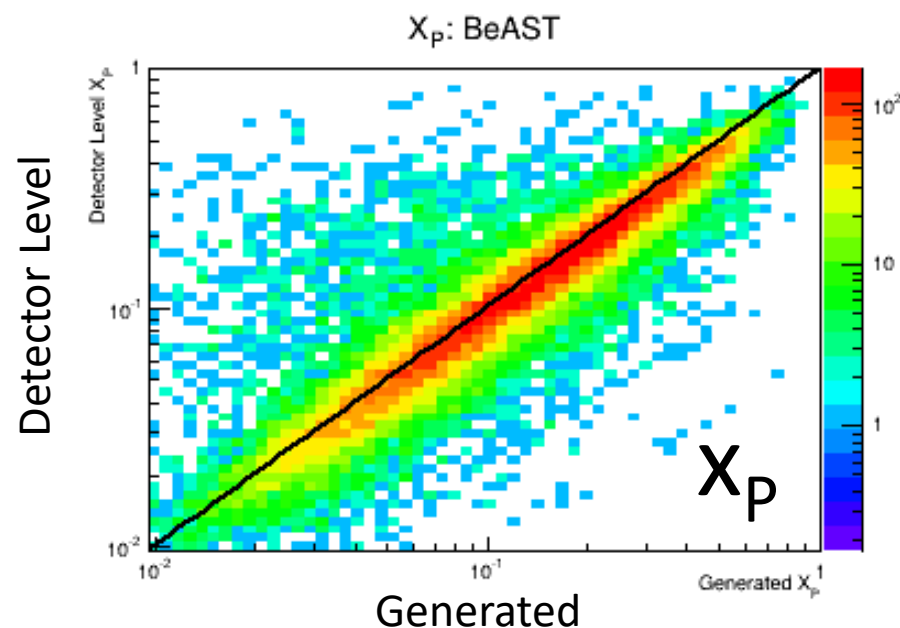
Particle Level

# Particle – Detector Comparison



- Both  $x_\gamma$  and  $x_p$  show good agreement between the generated and smeared values
- Level of smearing is similar to that seen between particle level and generated values

- Have seen that the individual jet quantities are well reproduced after smearing
- What about quantities derived from jet properties such as  $x_\gamma$  and  $x_p$ ?

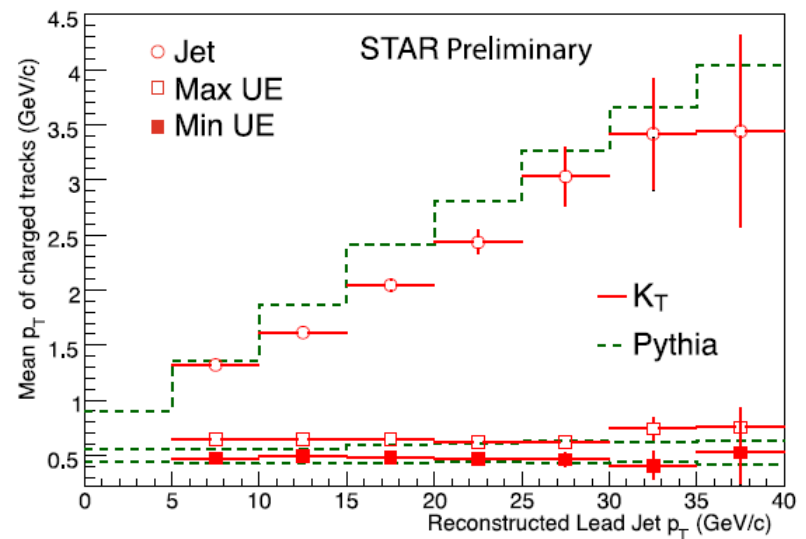
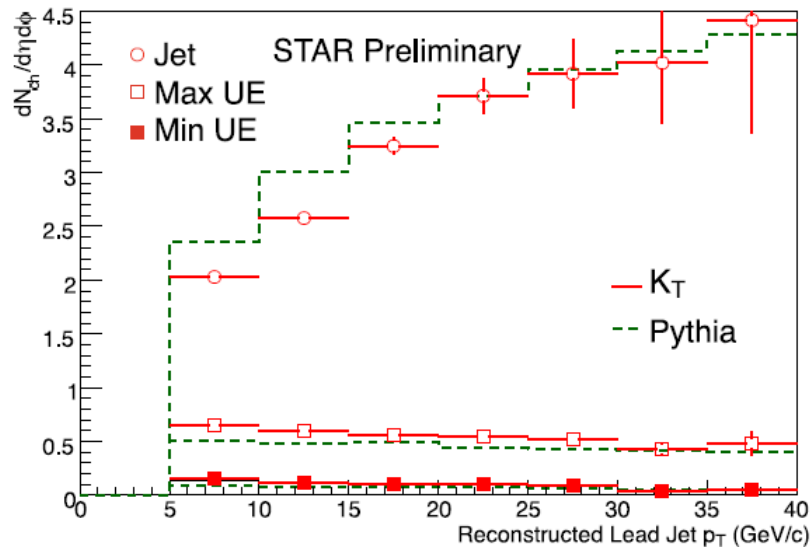
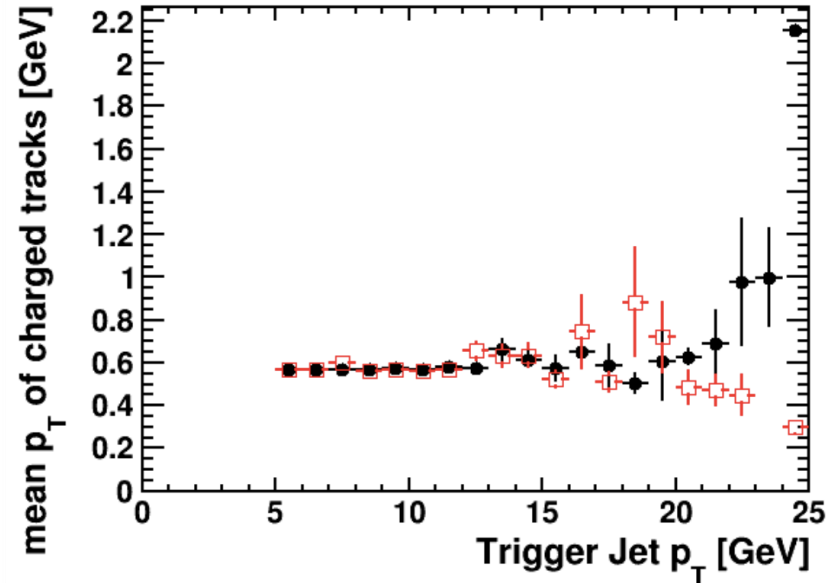
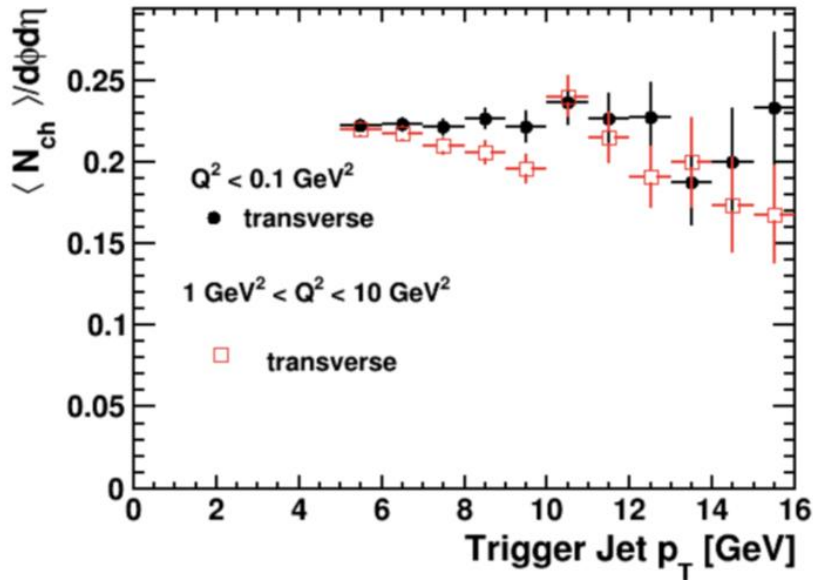


# Summary

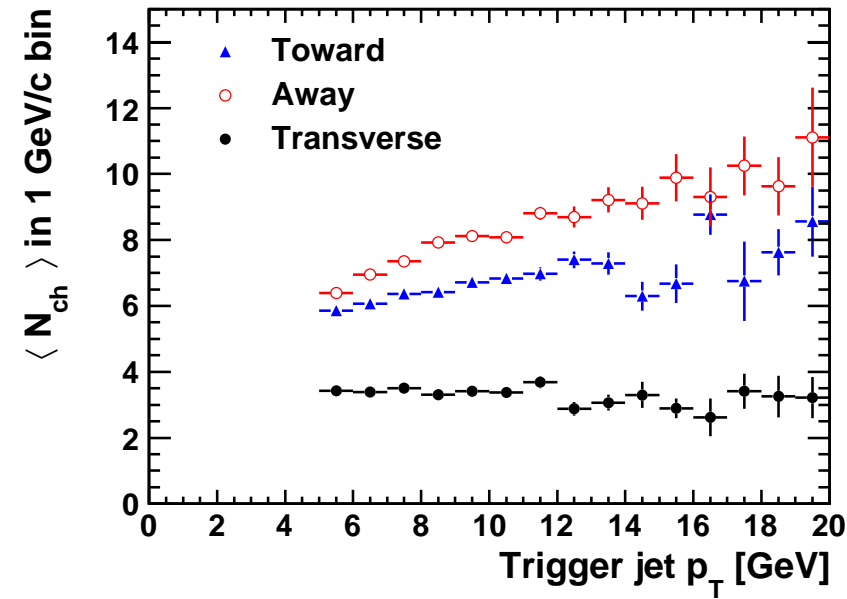
- Characterization of the underlying event was made and the contribution was found to be rather small
- Methods for studying photon structure and gluon polarization using dijets were detailed
- Studies of jet substructure are beginning and a preliminary look at their utility for discriminating between quark and gluon initiated jets has been done
- Impacts of detector effects on jet observables are being investigated using a fast smearing package and BeAST detector parameters

# Backup

# Comparison with STAR

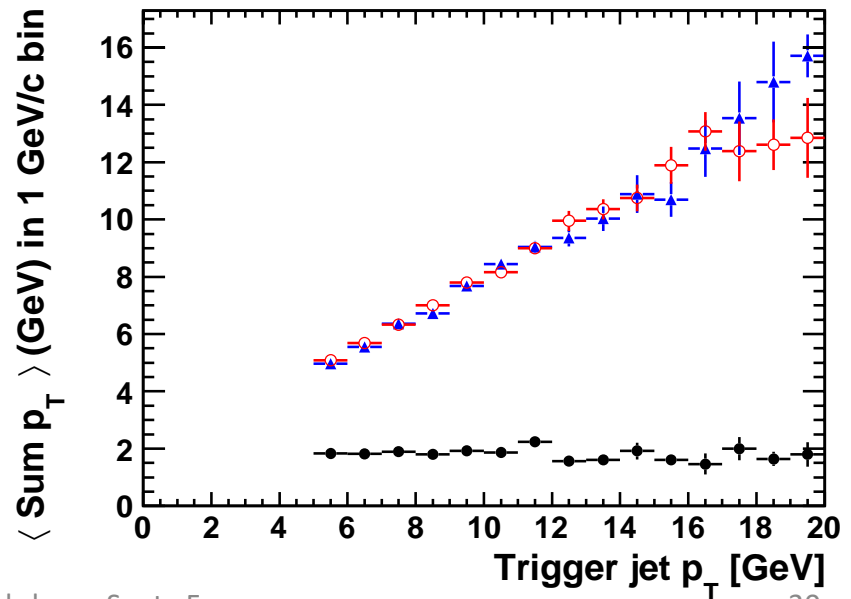


# Underlying Event Characteristics

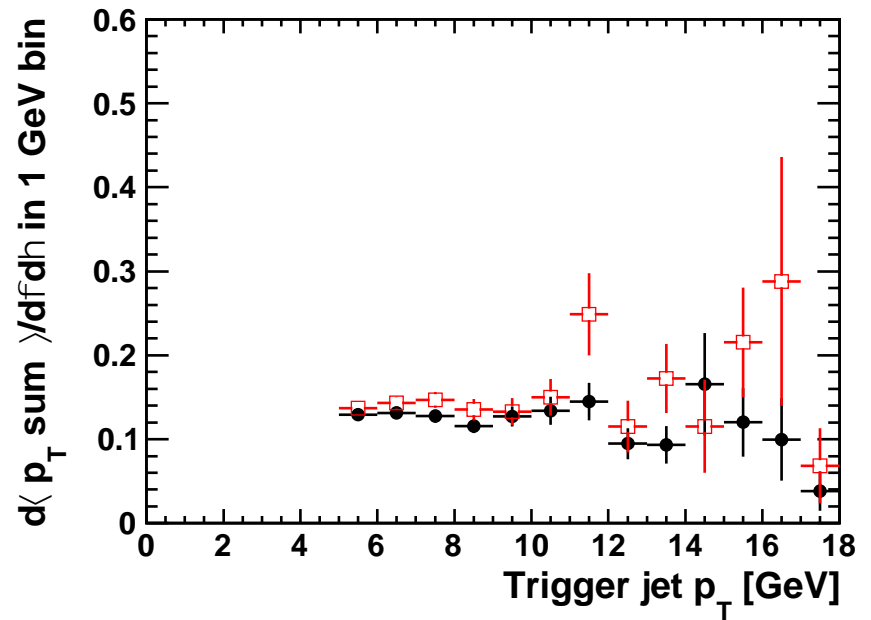
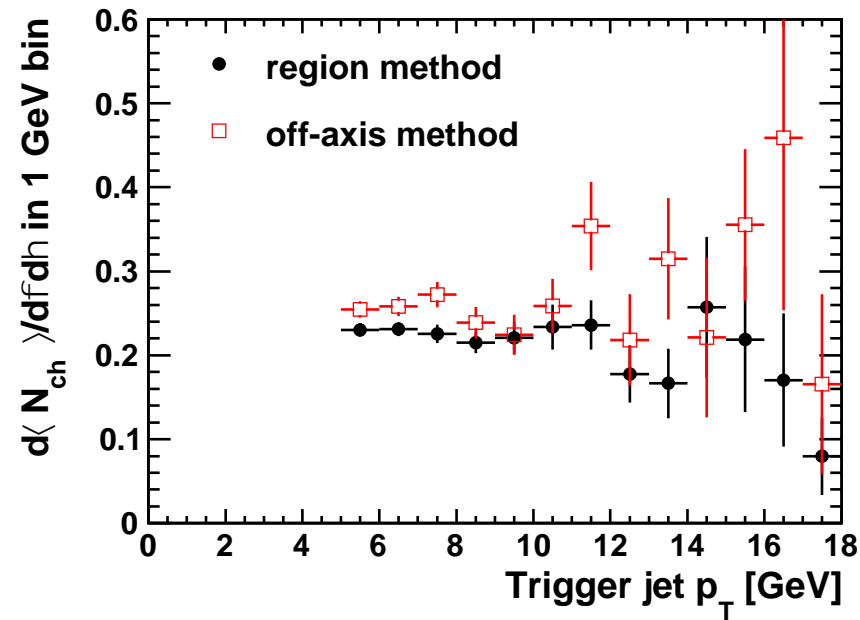


- Underlying event activity similar to that seen in pp collisions at STAR at  $\sqrt{s} = 200$  GeV

- Plot the average number of charged particles and average summed  $p_T$  for the three regions as a function of trigger jet  $p_T$
- See that these quantities are independent of the trigger jet  $p_T$  in transverse region

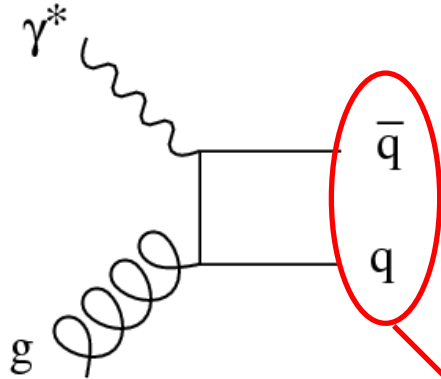


# UE Method Comparison

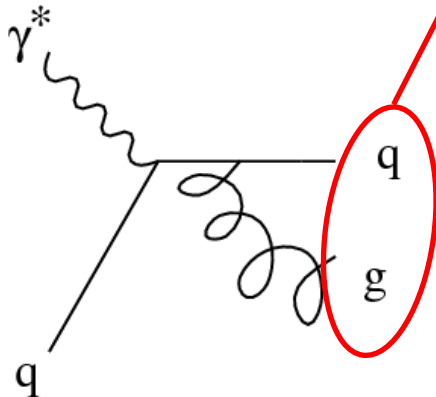


# Gluon Polarization with Di-jets

## Photon-Gluon Fusion



## QCD – Compton

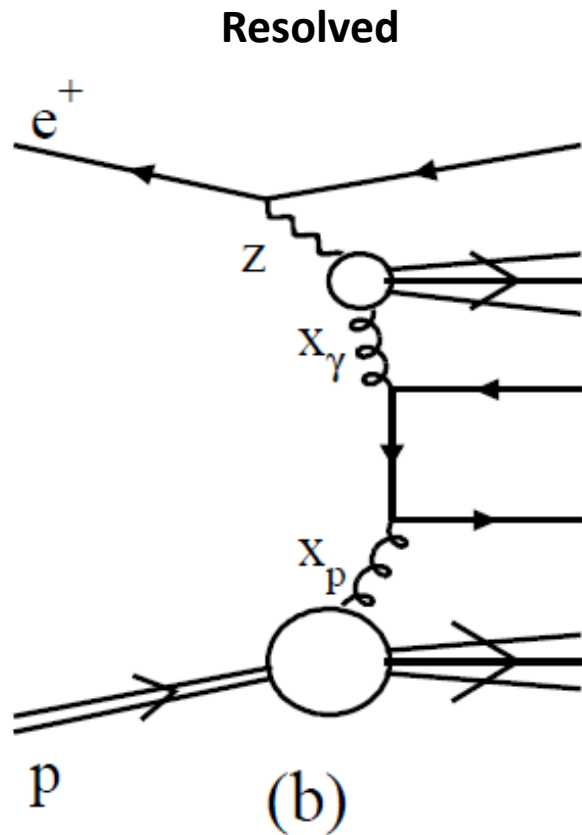


**Di-Jet**

- Gluons can also be probed in DIS via the higher-order photon gluon fusion process
- Also have the QCD – Compton process which probes quarks at the same order
- Both processes produce 2 angularly separated hard partons -> Di-jet



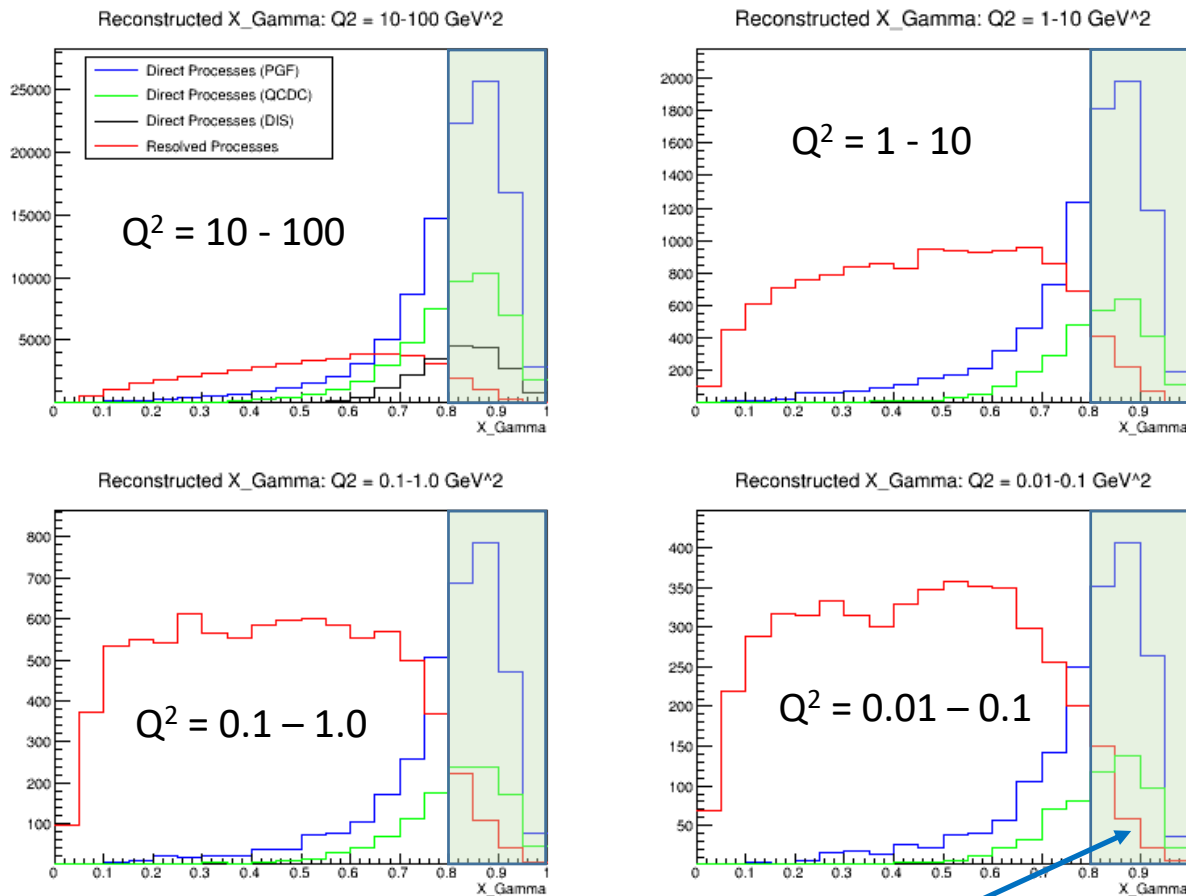
# Gluon Polarization with Di-jets



- Gluons can also be probed in DIS via the higher-order photon gluon fusion process
- Also have the QCD – Compton process which probes quarks at the same order
- Both processes produce 2 angularly separated hard partons  $\rightarrow$  Di-jet
- At lower  $Q^2$ , resolved processes in which the photon assumes a hadronic structure begin to dominate
- Asymmetry is a convolution of polarized PDF from the proton and polarized photon structure – which is completely unconstrained
- Would like to suppress the resolved component

# Direct Vs Resolved Processes

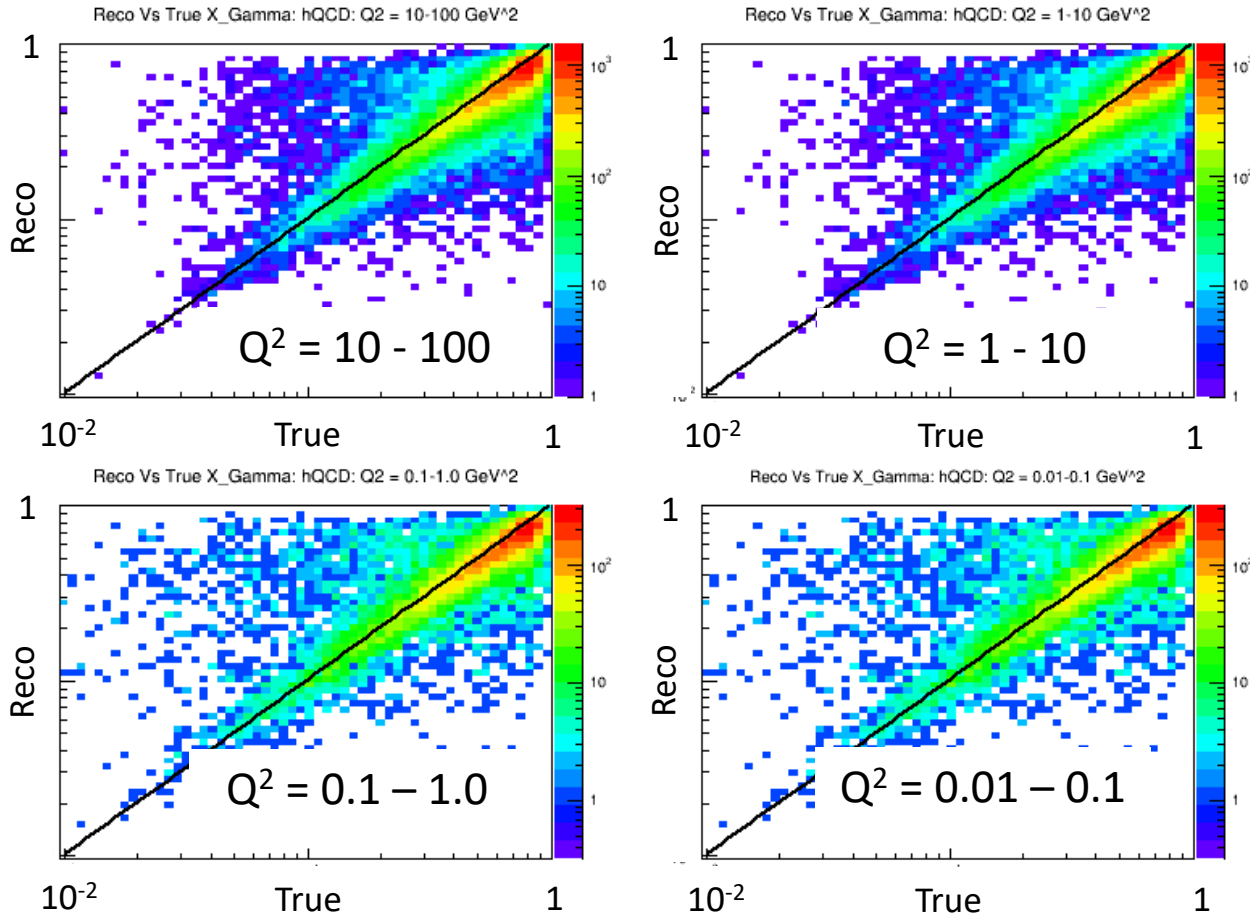
$$X_\gamma = \frac{1}{2E_e y} (m_{T1} e^{-y_1} + m_{T2} e^{-y_2})$$



Accepted Region

- Plot reconstructed  $X_\gamma$  for direct and resolved processes
- Direct processes should concentrate toward 1 while resolved processes are at lower values
- Direct processes dominate at higher  $Q^2$  while resolved are more prevalent at low  $Q^2$
- Cut of  $X_\gamma > 0.8$  enhances the direct fraction at all  $Q^2$

# $X_\gamma$ : Reconstructed Vs True



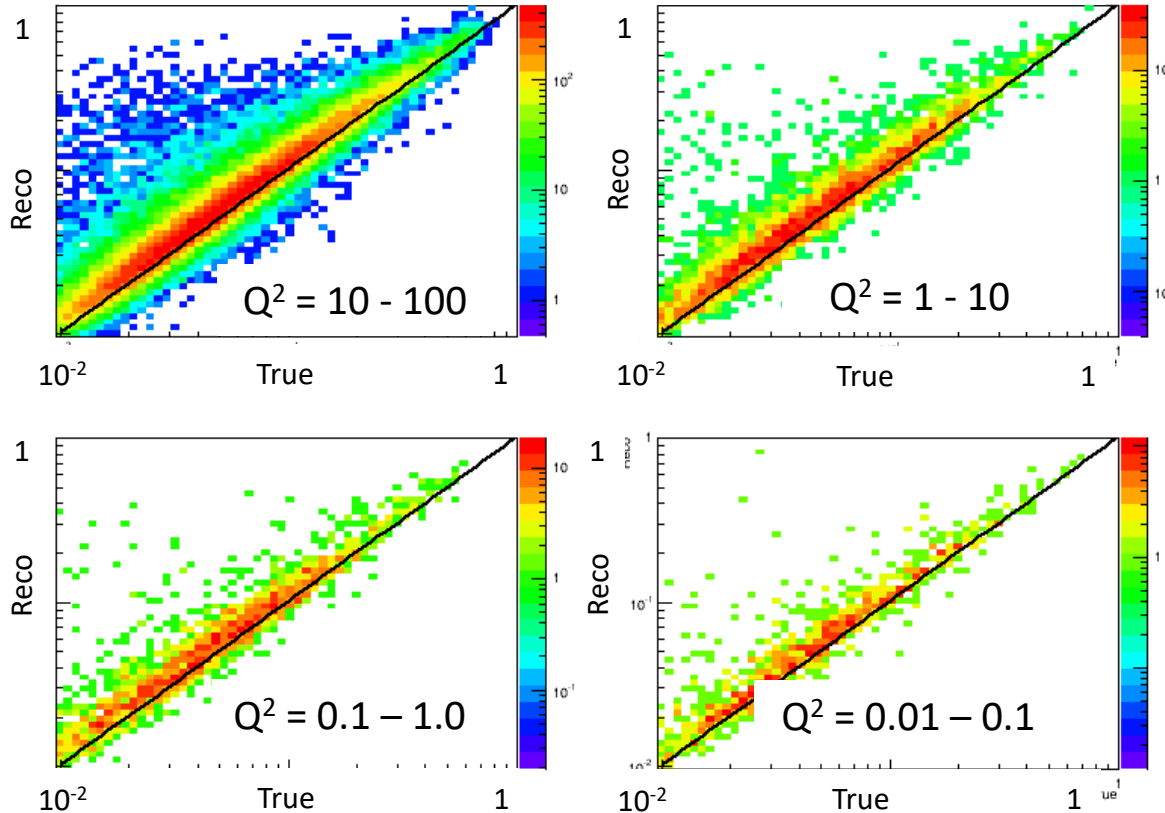
- Will use virtual photon momentum fraction to discriminate between resolved and direct processes
- See good agreement between reconstructed and true  $X_\gamma$  for all  $Q^2$  ranges
- Di-jets found in Breit frame and required one jet with  $p_T \geq 5 \text{ GeV}$  and the other with  $p_T \geq 4 \text{ GeV}$

$$X_\gamma = \frac{1}{2E_{ey}} (m_{T1} e^{-y_1} + m_{T2} e^{-y_2})$$

# Proton Partonic Kinematics

Parton Momentum Fraction: Photon Gluon Fusion

- To measure  $\Delta G$ , need to probe the parton coming from the proton
- Momentum fraction of the parton from proton is well reconstructed




$$x_p = \frac{1}{2E_P} (m_{T1}e^{y_1} + m_{T2}e^{y_2})$$

# Proton Partonic Kinematics

$$X_P = x_B \left( 1 + \frac{M^2}{Q^2} \right)$$

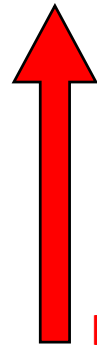
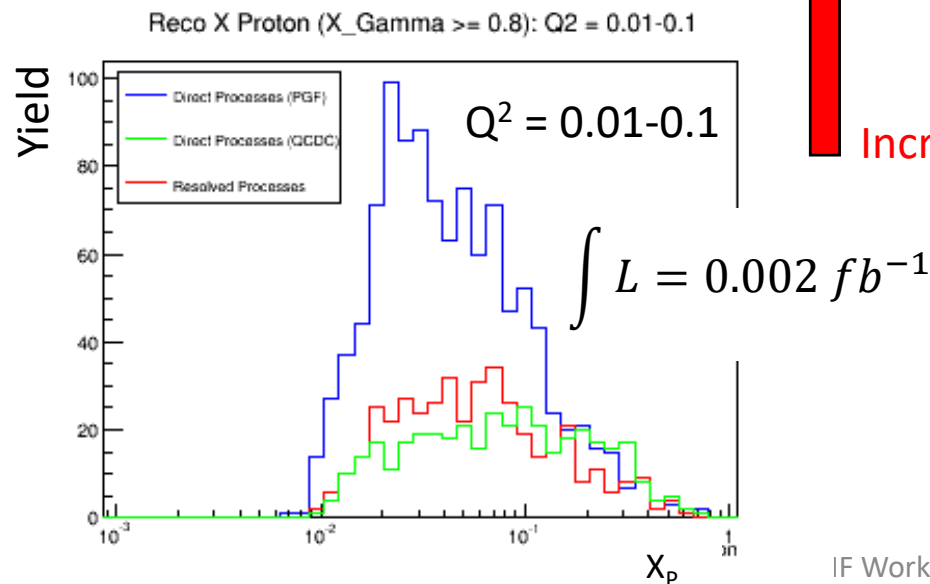
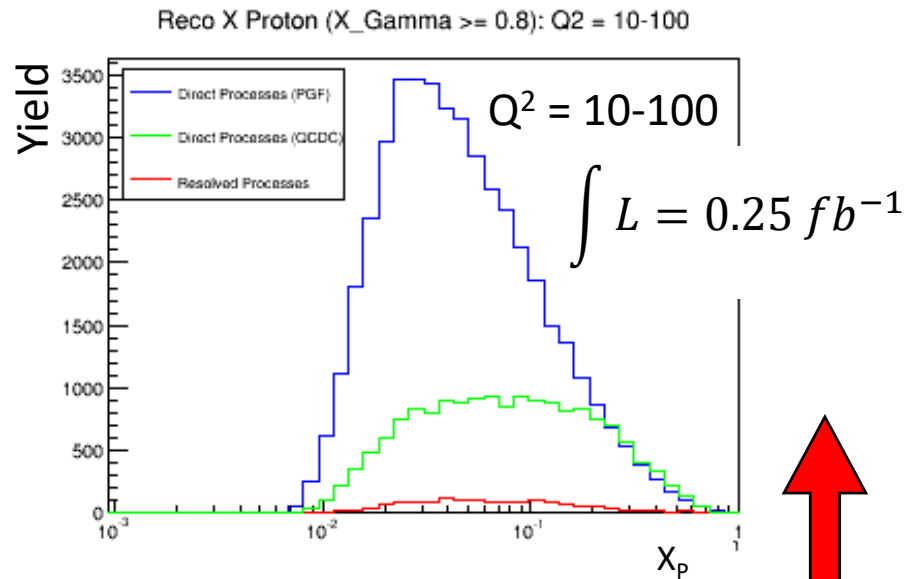
$$Q^2 = syx_B$$

$$X_P = x_B + \frac{M^2}{sy}$$


$$\approx \frac{100}{(20000 \times 0.95)} \approx 0.005$$

- To measure  $\Delta G$ , need to probe the parton coming from the proton
- Momentum fraction of the parton from proton is well reconstructed
- $X_p$  is related to Bjorken- $x$  and  $Q^2$  at leading order
- $Q^2$  and Bjorken- $x$  are also related via the collision energy and inelasticity
- Accessible  $X_p$  range basically determined by beam energies
- Lowest  $X_p$  we can probe is about 0.005

# $X_p$ For Different $Q^2$



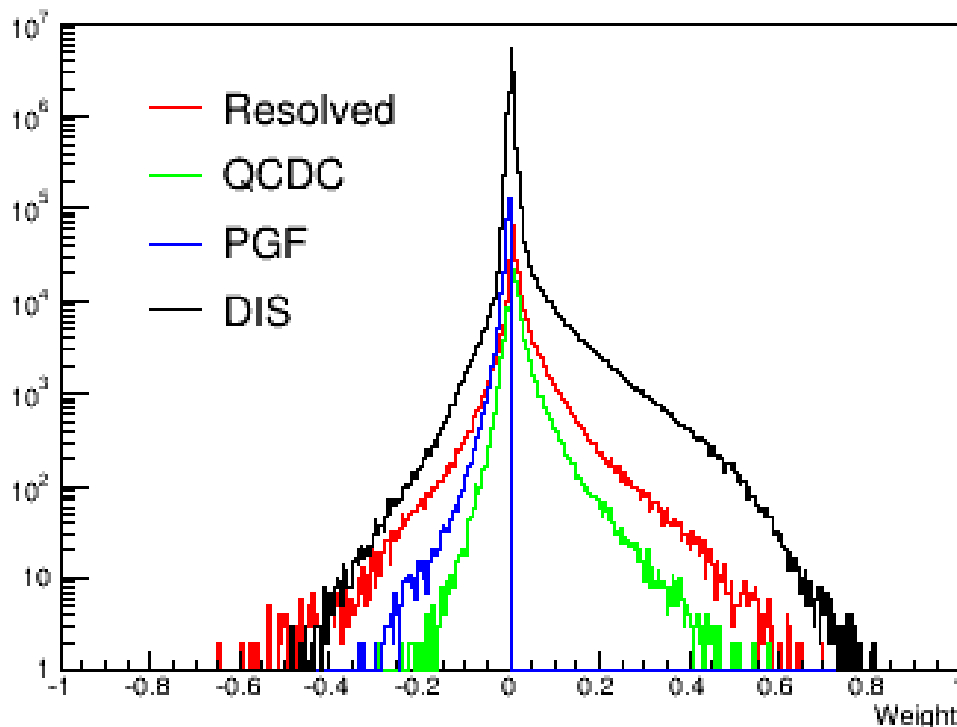
Increasing  $Q^2$

- At lower  $Q^2$ , contribution from resolved process increases while QCD Compton contribution decreases
- For a given di-jet mass range (10 – 20 GeV in this case), same  $X_p$  can be reconstructed event-by-event and probed over large range of  $Q^2$
- This will allow for robust tests of the evolution of  $\Delta G$

# Weighting PYTHIA

$$w = \hat{a}(\hat{s}, \hat{t}, \hat{\mu}^2, Q^2) \cdot \frac{\Delta f_a^{\gamma^*}(x_a, \mu^2)}{f_a^{\gamma^*}(x_a, \mu^2)} \cdot \frac{\Delta f_b^N(x_b, \mu^2)}{f_b^N(x_b, \mu^2)}$$

Total Weight (DSSV14):  $Q^2 = 10\text{-}100 \text{ GeV}$

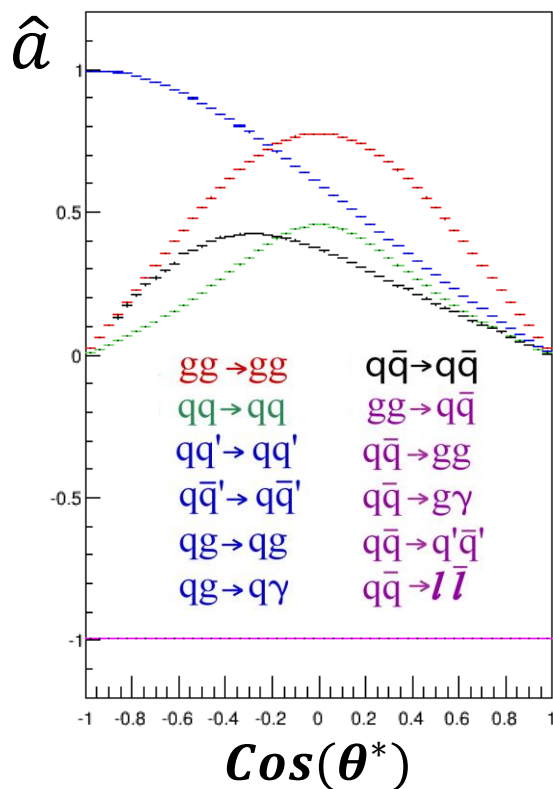


- PYTHIA does not include parton polarization effects, but an asymmetry can be formed by assigning each event a weight depending on the hard-scattering asymmetry and (un)polarized photon and proton PDFs
- Expected asymmetry is then the average over weights
- Weights are sharply spiked near zero -> expect small asymmetries

# Weighting PYTHIA

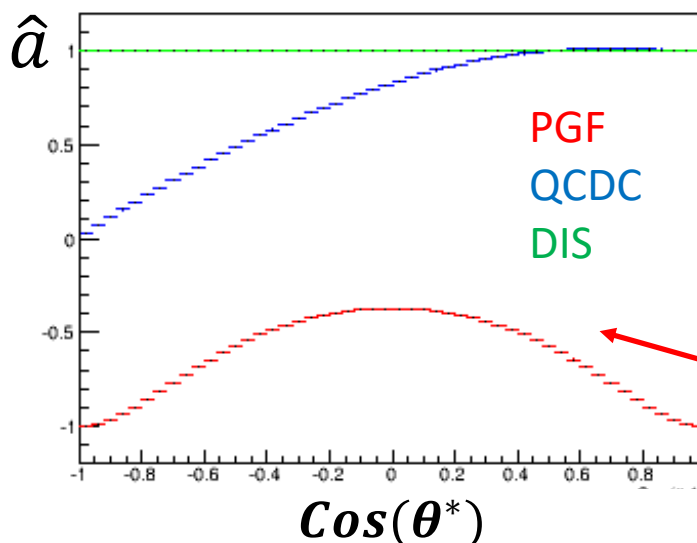
$$w = \hat{a}(\hat{s}, \hat{t}, \hat{\mu}^2, Q^2) \bullet \frac{\Delta f_a^{\gamma^*}(x_a, \mu^2)}{f_a^{\gamma^*}(x_a, \mu^2)} \bullet \frac{\Delta f_b^N(x_b, \mu^2)}{f_b^N(x_b, \mu^2)}$$

**Resolved**



$$\hat{a}(\hat{s}, \hat{t}, \hat{\mu}^2) = \Delta \hat{\sigma} / (2 \hat{\sigma})$$

**Direct**



- Process-dependent hard scattering asymmetry is a function of Mandelstam variables ( $\cos(\theta^*)$ )

- The direct process distributions will be smeared by the additional depolarization term

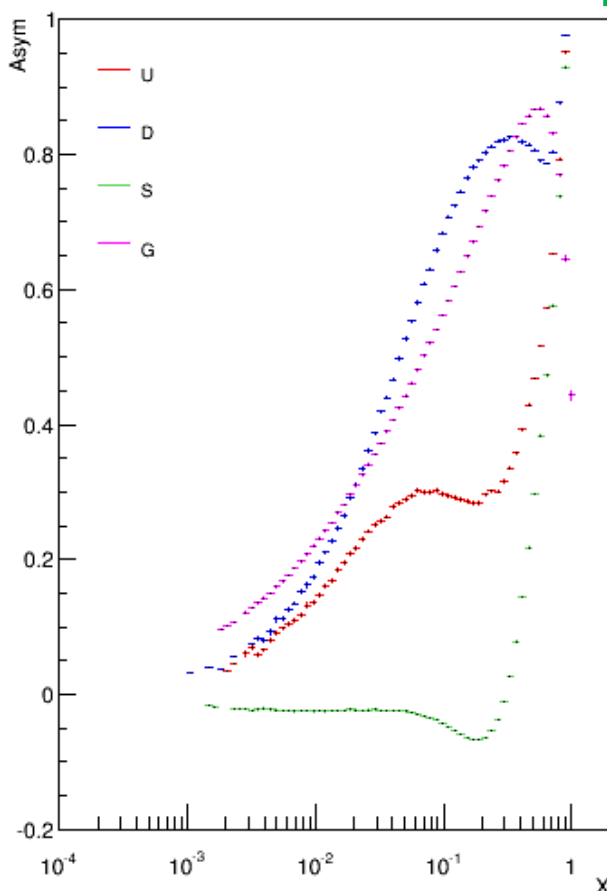
- Note that the asymmetry for PGF is negative



# Weighting PYTHIA

$$w = \hat{a}(\hat{s}, \hat{t}, \hat{\mu}^2, Q^2) \bullet \frac{\Delta f_a^{\gamma^*}(x_a, \mu^2)}{f_a^{\gamma^*}(x_a, \mu^2)} \bullet \frac{\Delta f_b^N(x_b, \mu^2)}{f_b^N(x_b, \mu^2)}$$

Photon Delta\_f/f: Profile: Q2 = 10-100 GeV

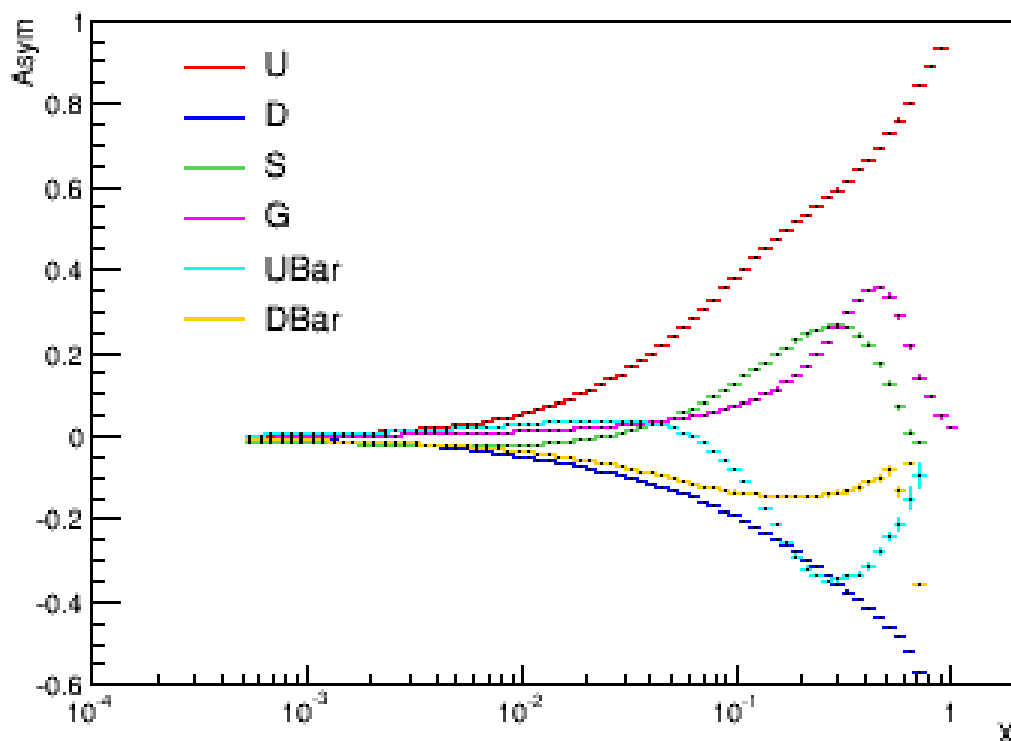


- Second term is the ratio of the polarized to unpolarized photon PDFs
- Use maximal scheme for polarized and GRV-G for unpolarized
- For direct processes such as Photon-Gluon Fusion, this term is identically unity

# Weighting PYTHIA

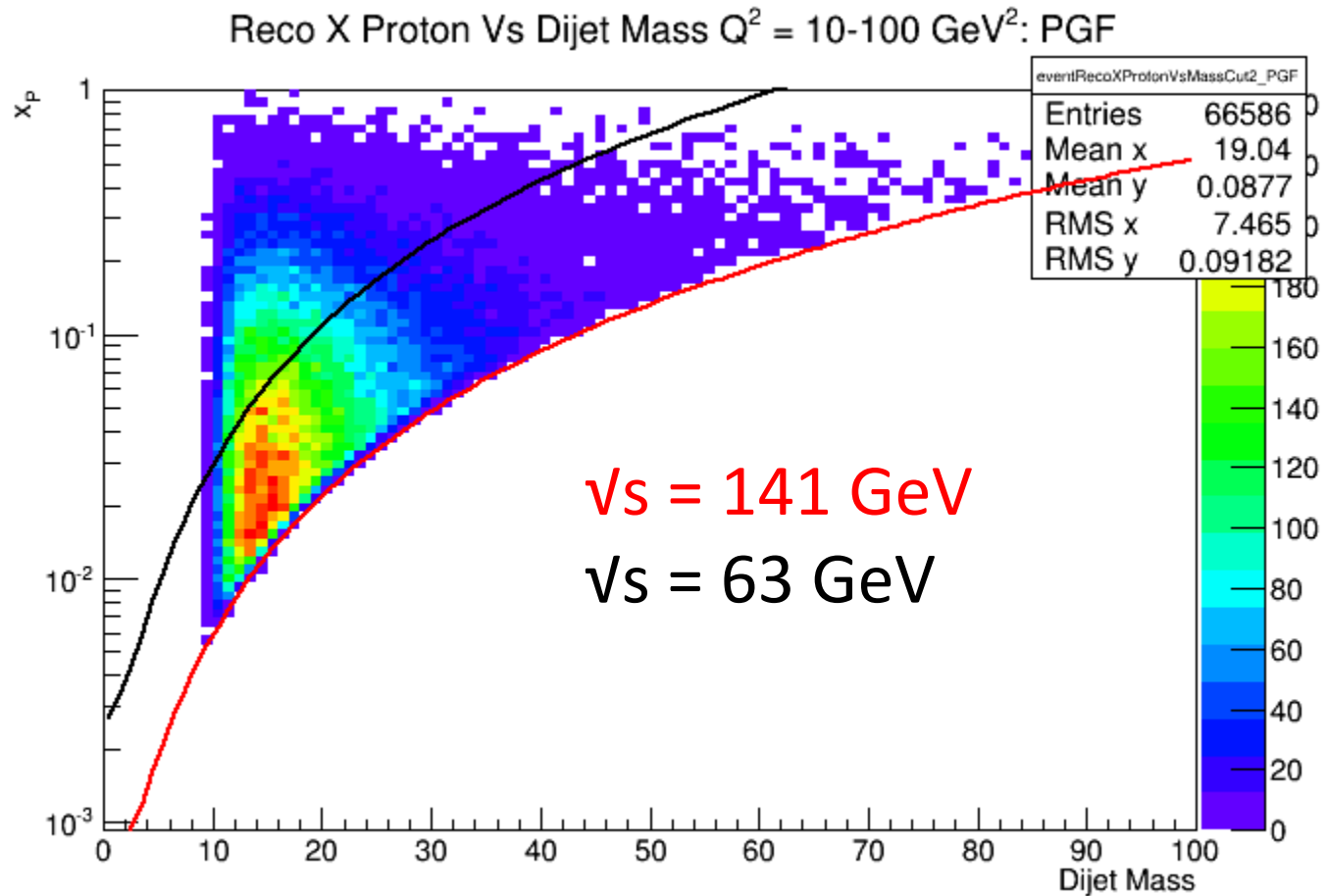
$$w = \hat{a}(\hat{s}, \hat{t}, \hat{\mu}^2, Q^2) \bullet \frac{\Delta f_a^{\gamma^*}(x_a, \mu^2)}{f_a^{\gamma^*}(x_a, \mu^2)} \bullet \frac{\Delta f_b^N(x_b, \mu^2)}{f_b^N(x_b, \mu^2)}$$

Proton Delta\_f/f: Profile (DSSV14): Q2 = 10-100 GeV



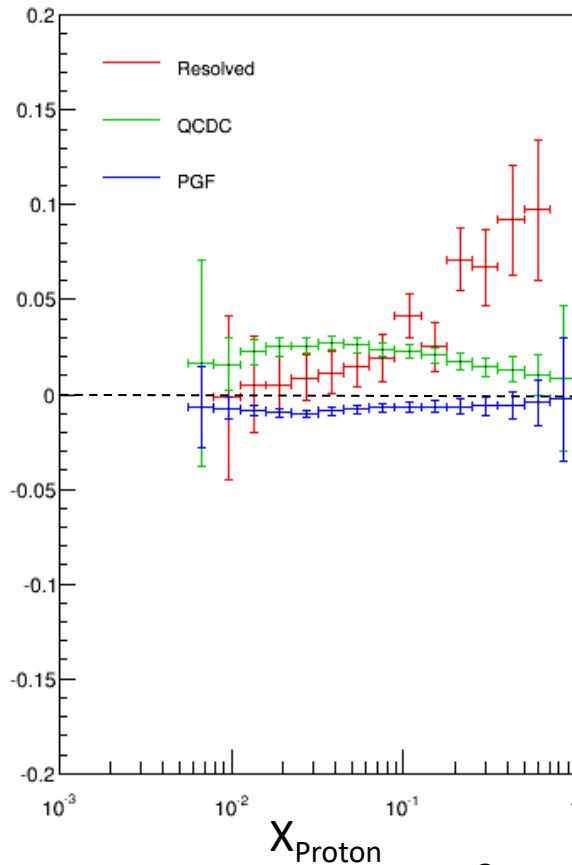
- Last term is the ratio of the polarized to unpolarized proton PDFs
- Use DSSV14 for polarized and CTEQ5M for unpolarized

# Dijet Phase Space

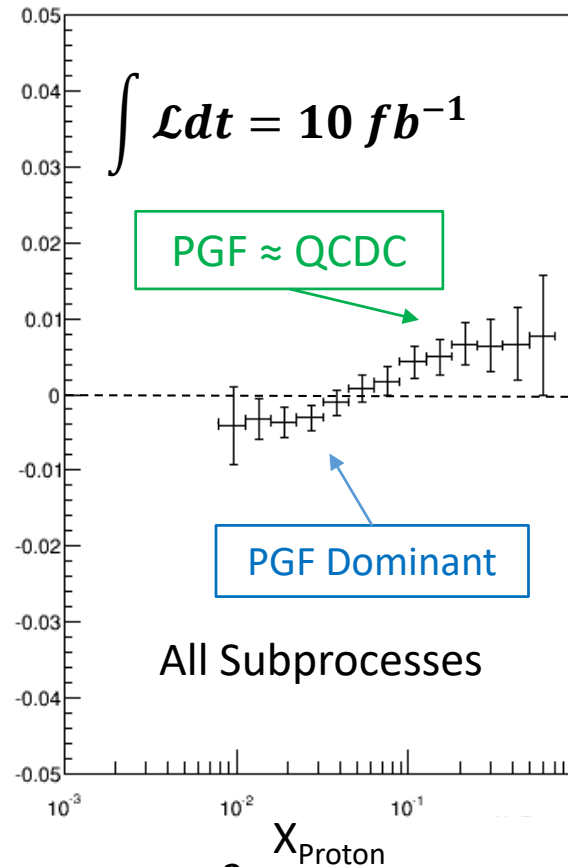


# $A_{LL}$ Vs $X_{\text{Proton}}$

$A_{LL}$  Vs Proton X



$A_{LL}$  Vs Proton X

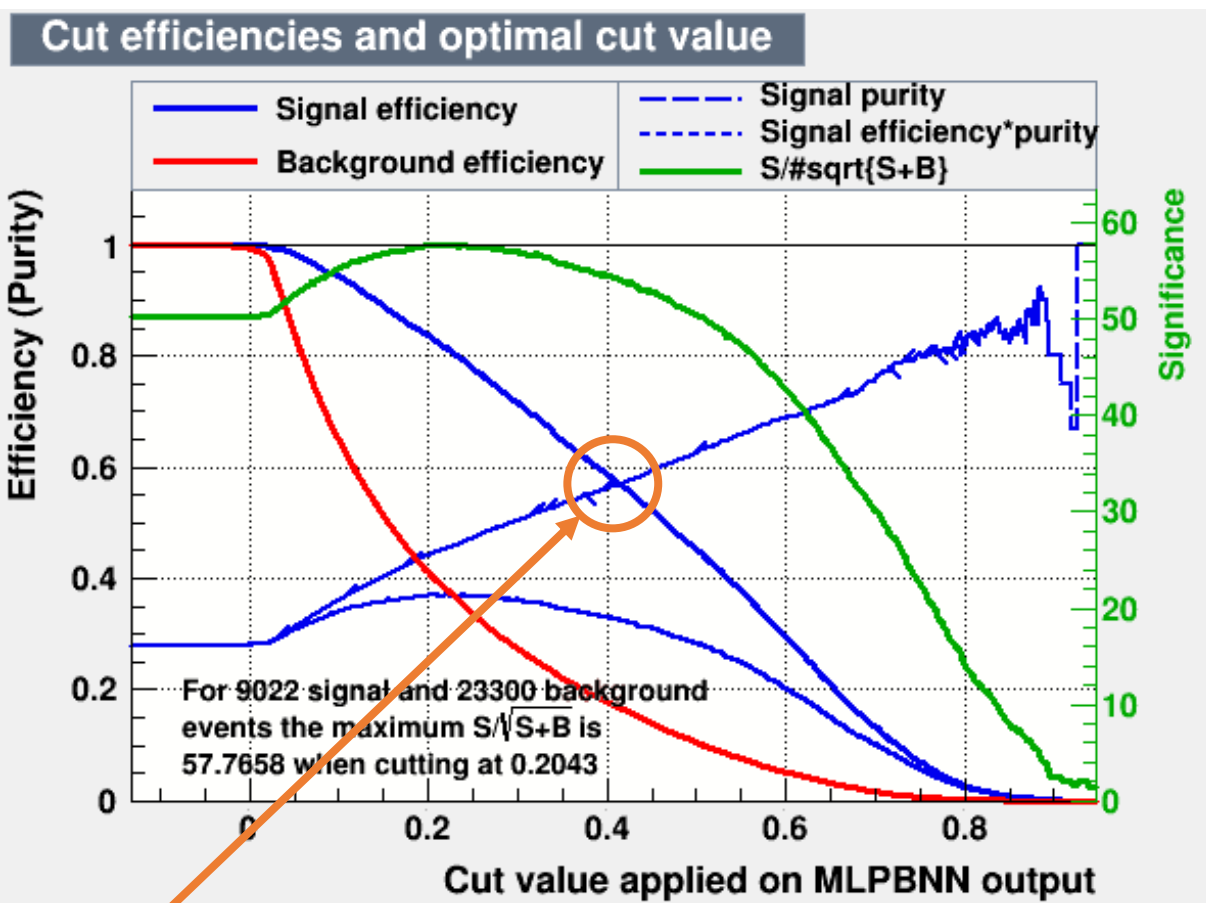


$$Q^2 = 10 - 100 \text{ GeV}^2$$

$$\sigma = \sqrt{\frac{1}{N} - \frac{A^2}{N}}$$

- Asymmetry is plotted as a function of the momentum fraction of the parton from the proton
- Asymmetry shown for di-jet invariant masses between 10 and 20  $\text{GeV}/c^2$
- Error bars are statistical and scaled to the given integrated luminosity
- Different mass ranges will emphasize different momentum fraction ranges and subprocess mixes

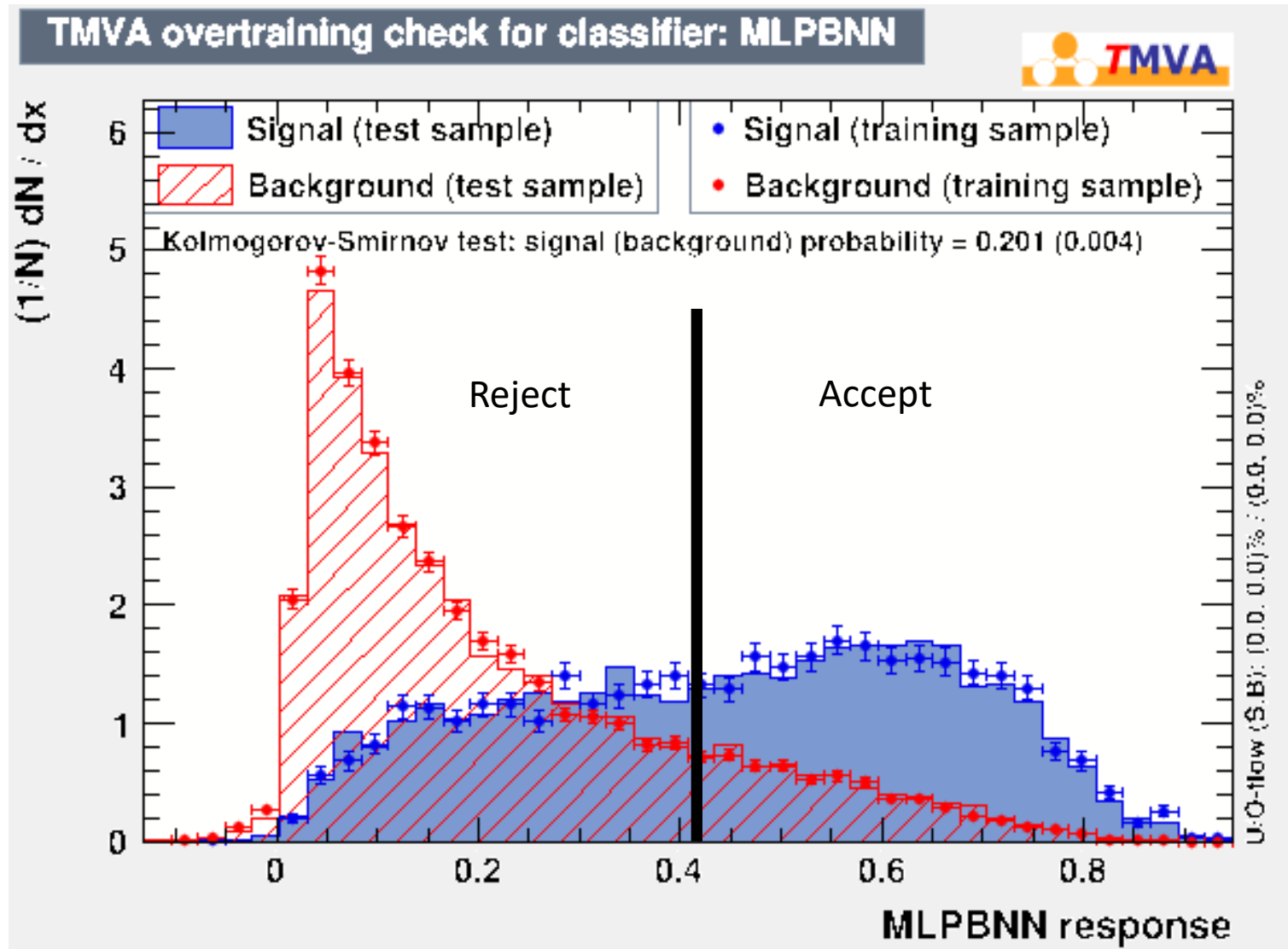
# Cut Optimization



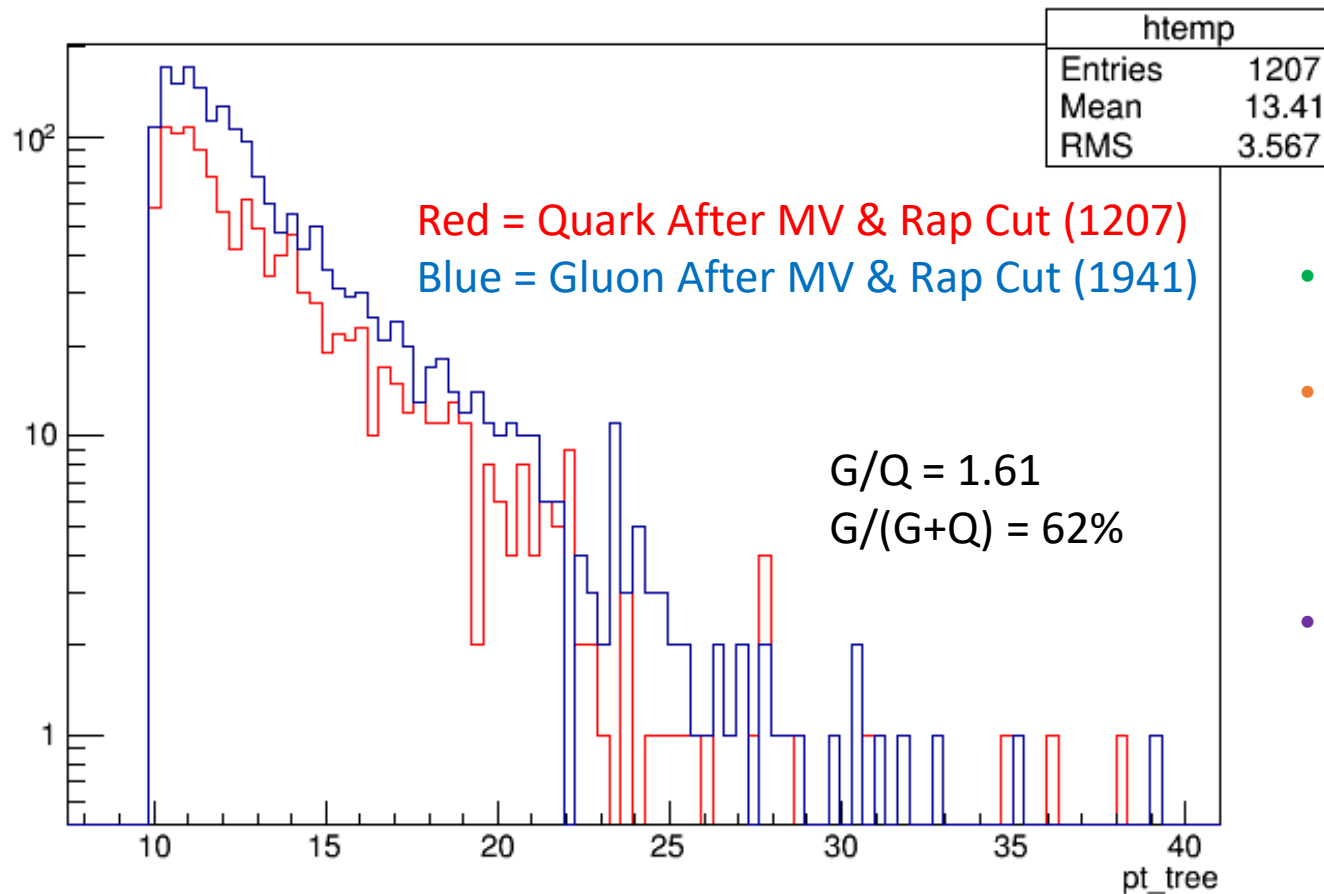
- For current study, place cut where signal purity = signal efficiency

- TMVA evaluates all input and maps them to a single variable with more signal-like events having a higher value
- Plot signal & background efficiency, signal purity, significance, etc as a function of this cut value
- This plot shows where to place cut in order to maximize purity, efficiency, or whatever an analysis requires

# MLPBNN Response

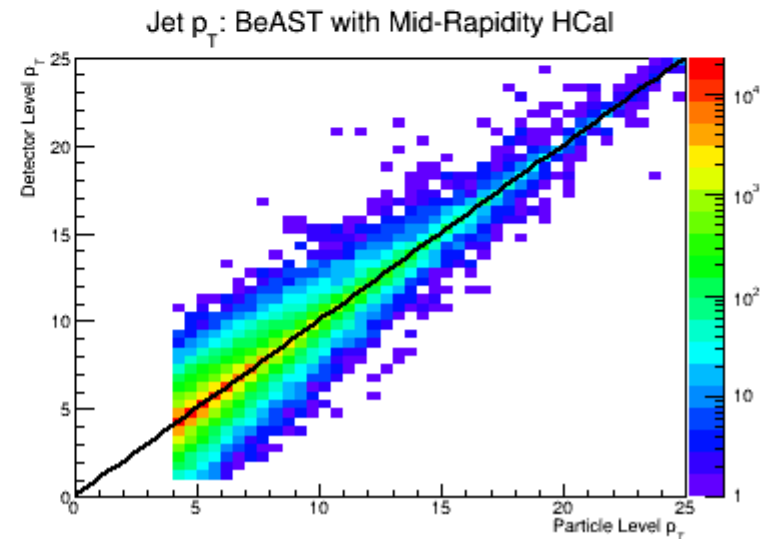
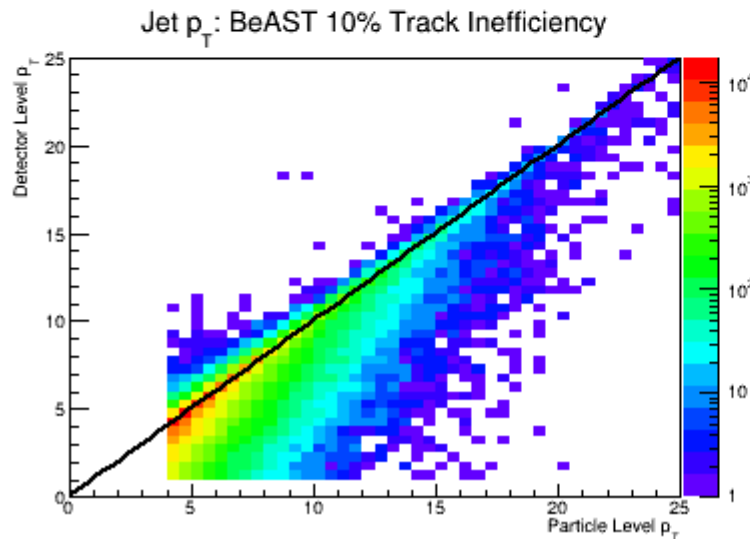
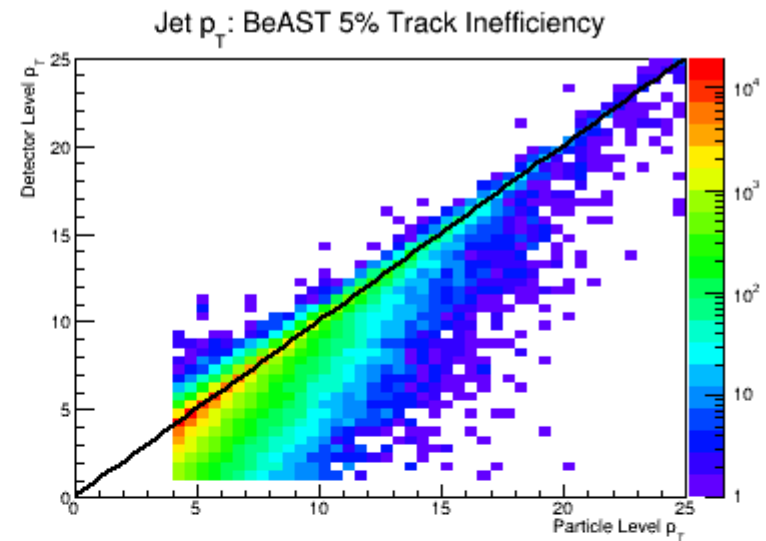
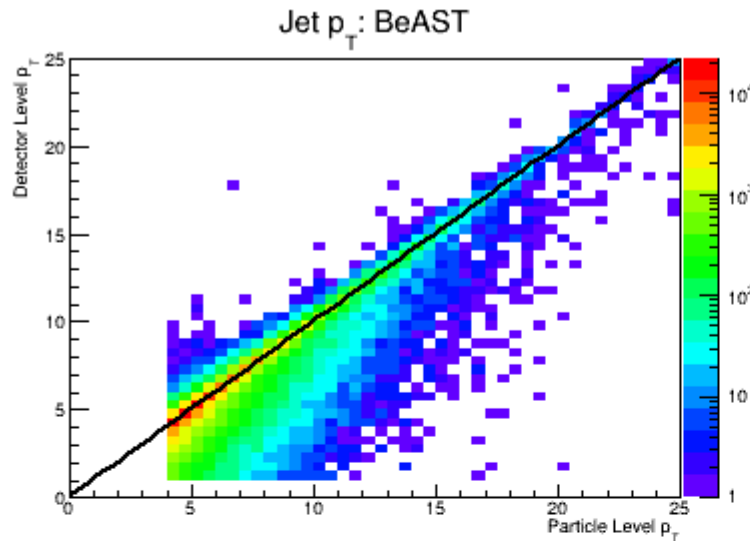


# Jet $p_T$ Spectra With Rapidity Cut



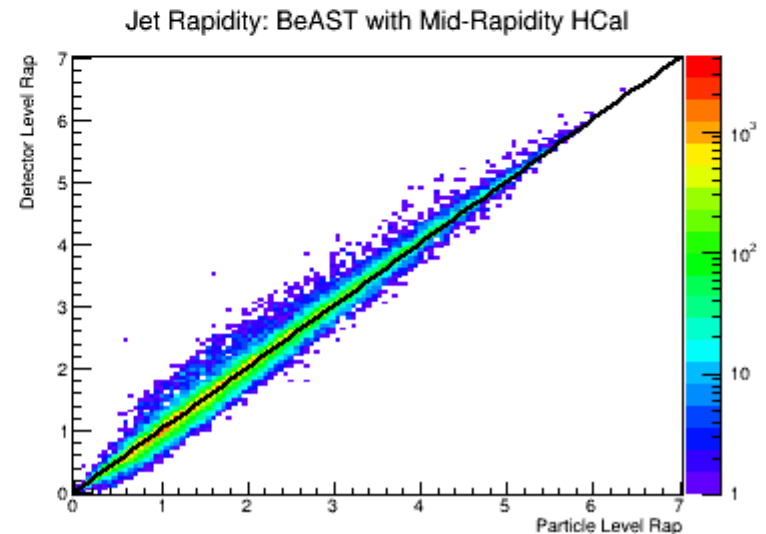
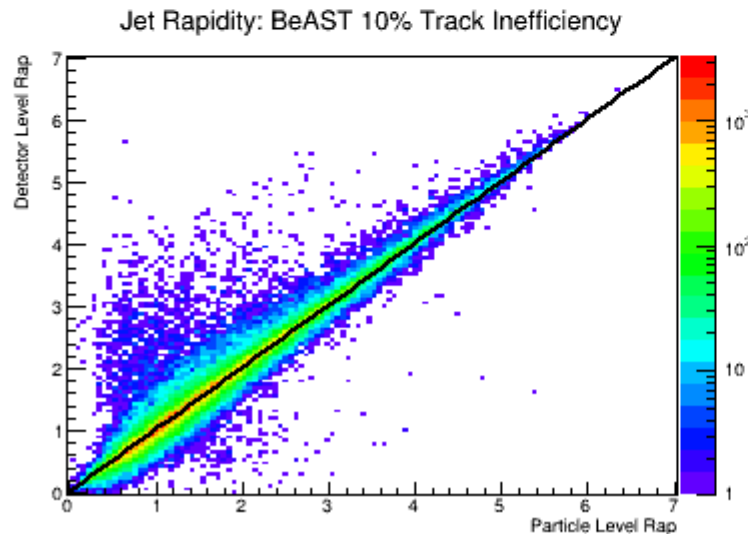
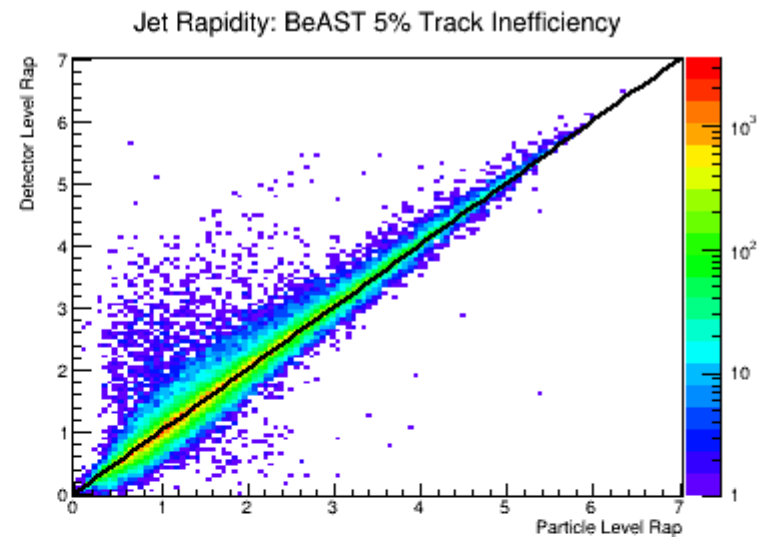
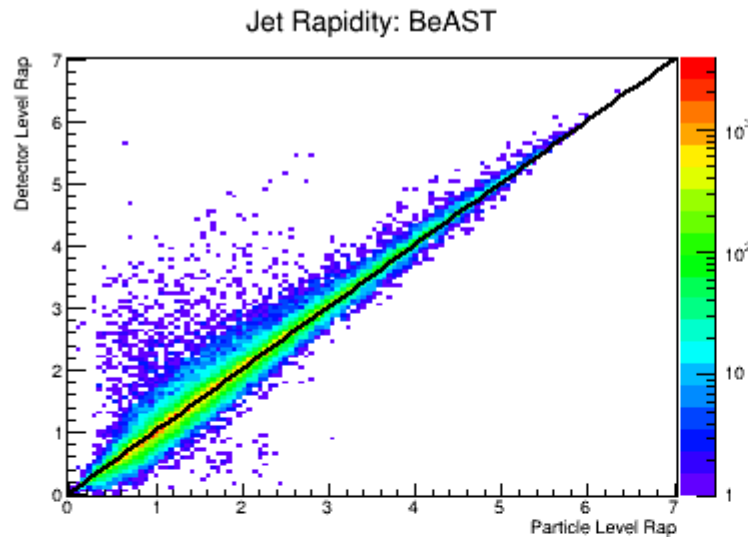
- Plot jet  $p_T$  after all cuts
- See reasonable enhancement of gluon jets over  $p_T$  range
- Should be able to get relatively pure quark sample and enhanced gluon sample for applications which require identification

# Particle – Detector Jet $p_T$



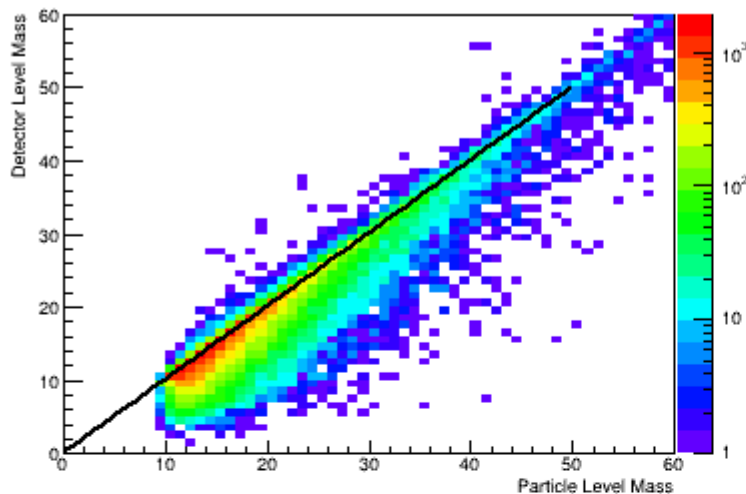


# Particle – Detector Jet Rapidity

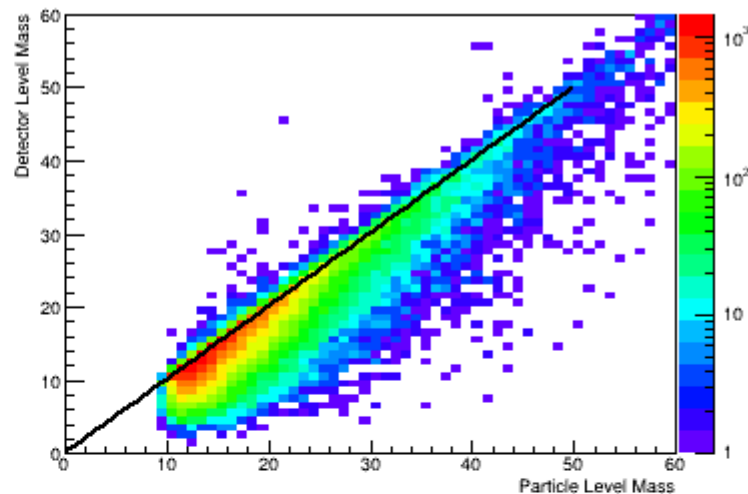


# Particle – Detector Dijet Mass

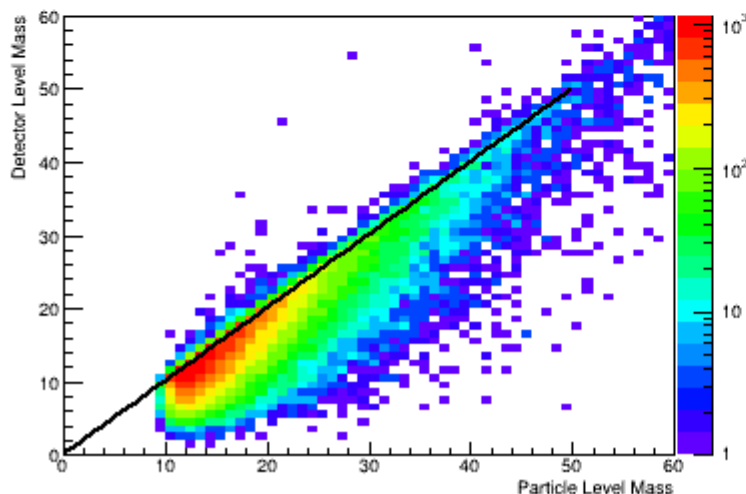
Dijet Mass: BeAST



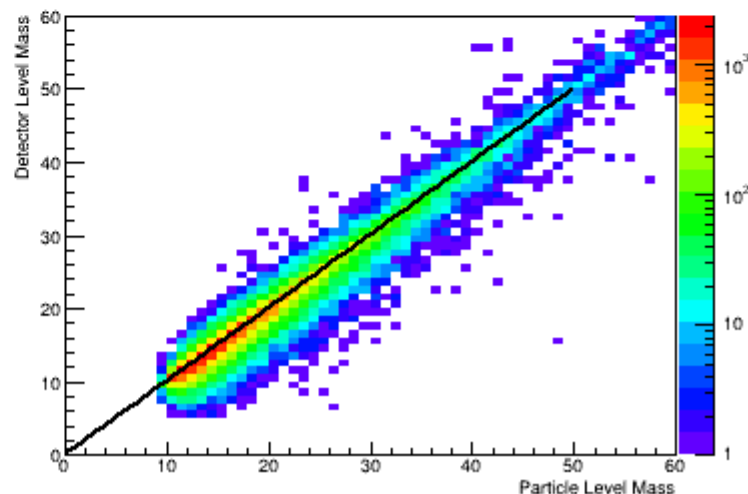
Dijet Mass: BeAST 5% Track Inefficiency



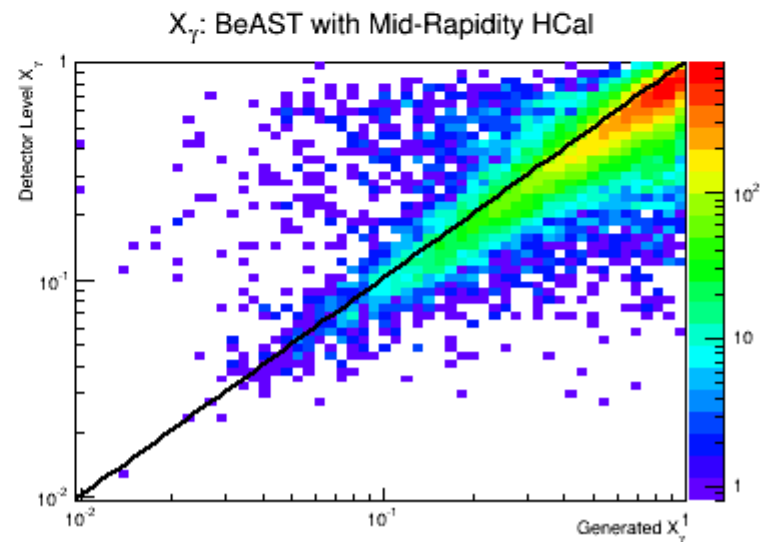
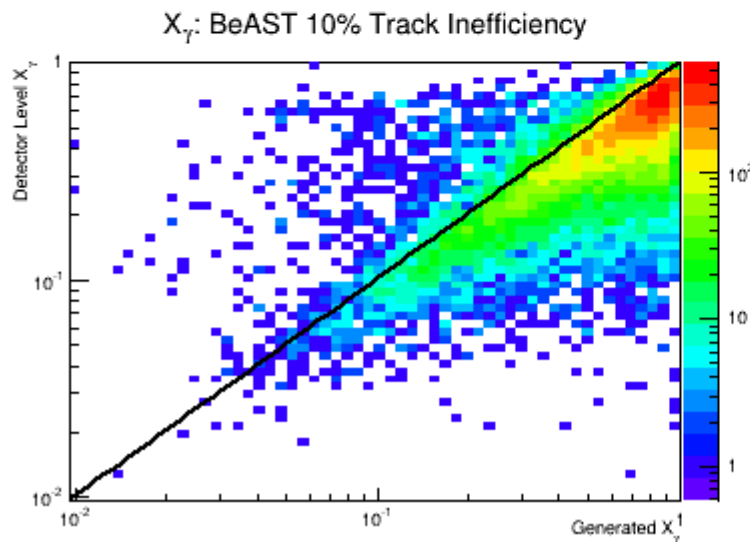
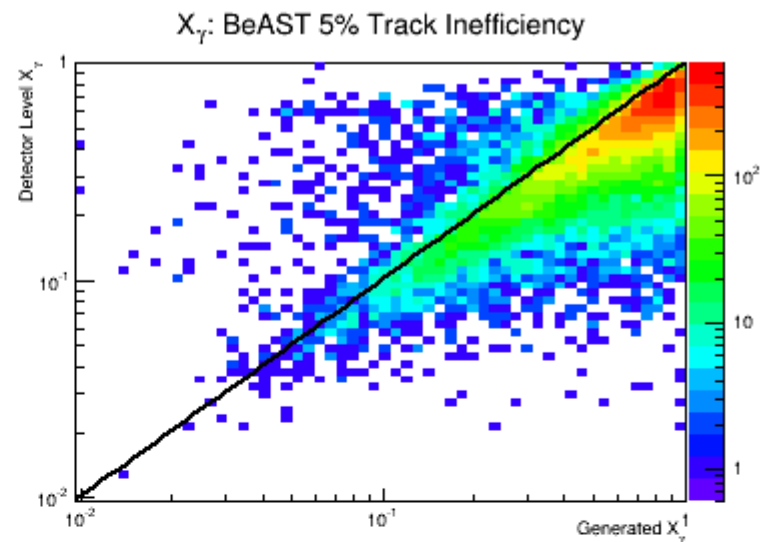
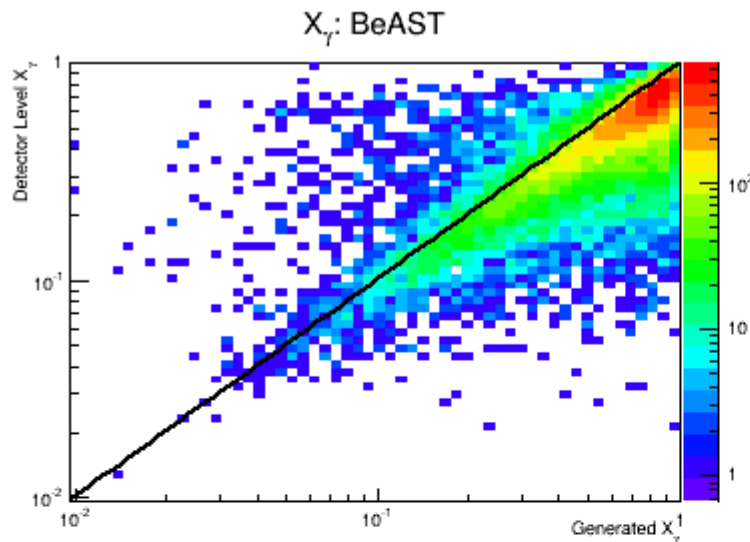
Dijet Mass: BeAST 10% Track Inefficiency



Dijet Mass: BeAST with Mid-Rapidity HCal



# Particle – Detector X\_Gamma



# Particle – Detector $X_p$ Proton

