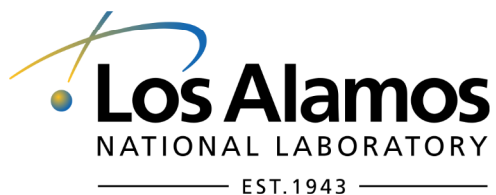
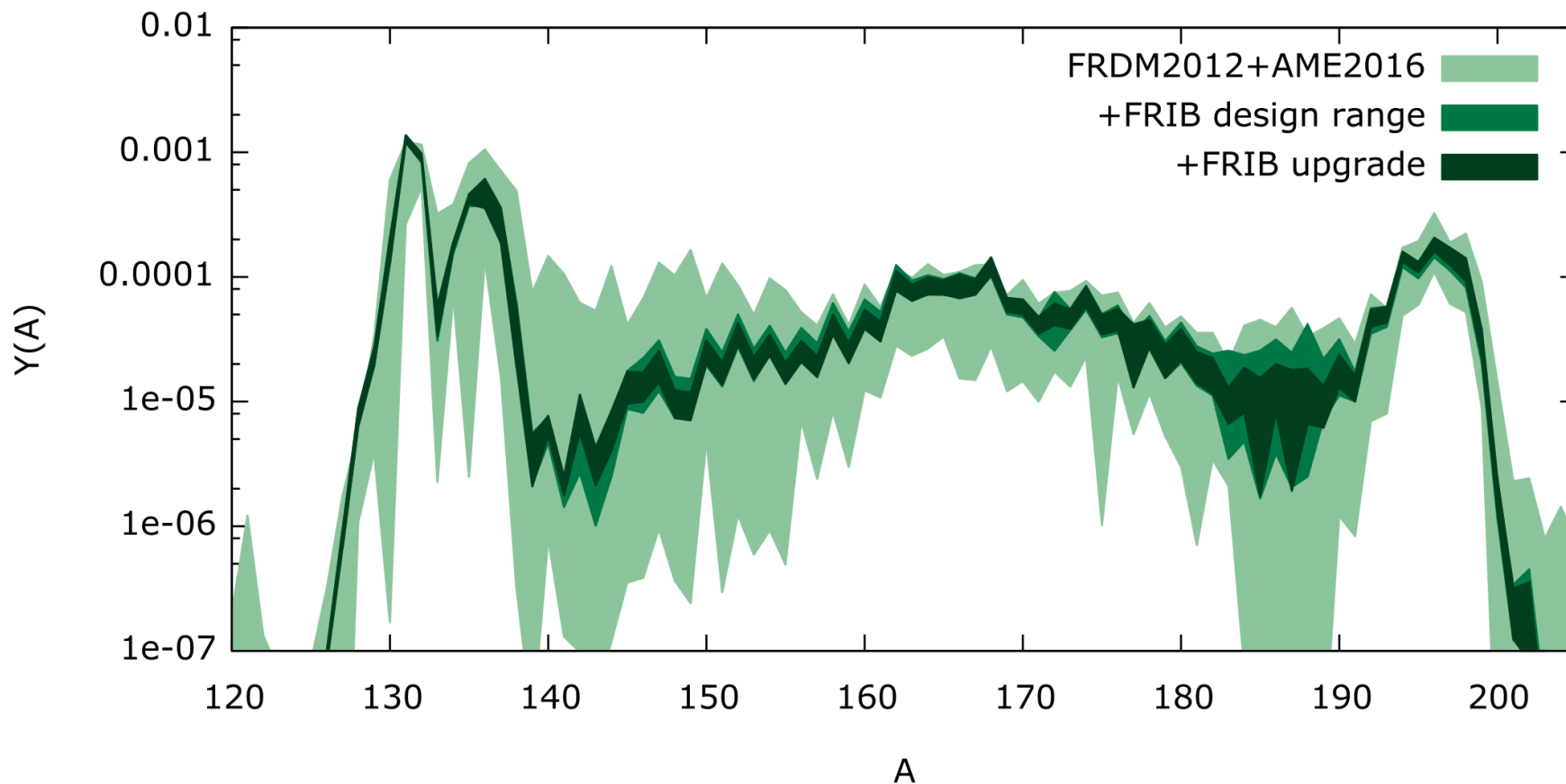


FRIB UPGRADE: ADVANTAGES FOR r-PROCESS STUDIES



LA-UR-18-27611

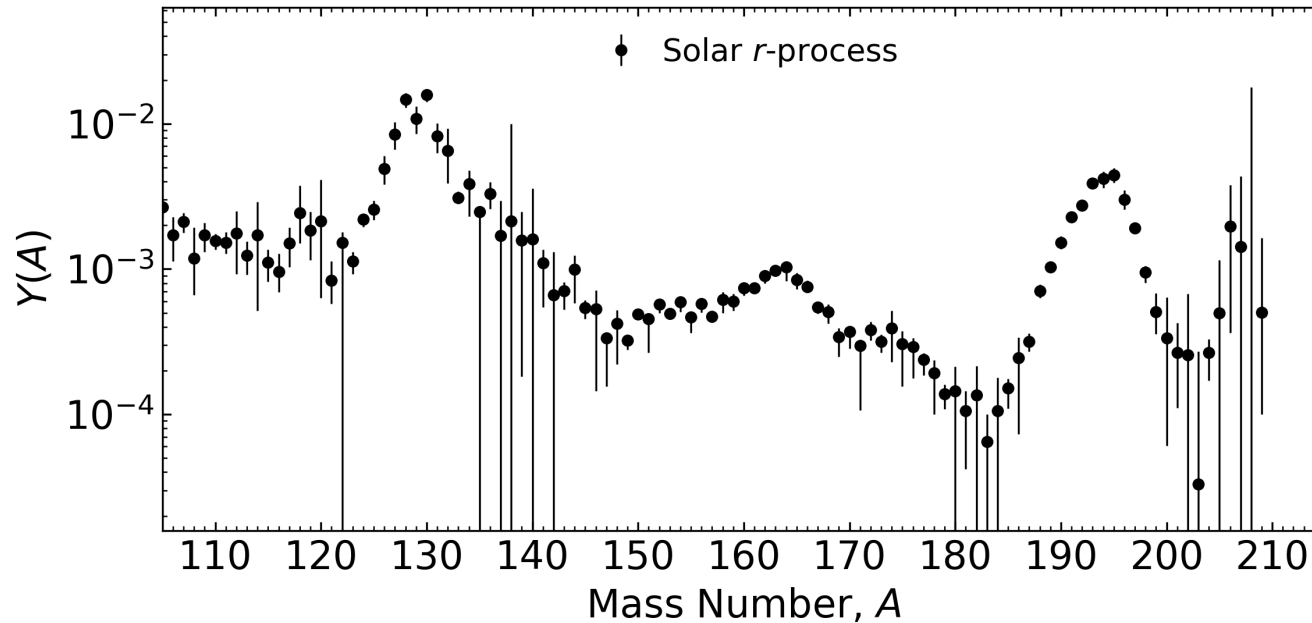
MATTHEW MUMPOWER
Los Alamos National Lab

FRIB Upgrade Workshop
Thurs. Aug. 9th 2018



THE PROBLEM

We want to describe the abundances observed in nature



But there is uncertainty in:

The astrophysical conditions

The nuclear physics inputs

Both are required to model the nucleosynthesis

INPUTS FROM NUCLEAR PHYSICS

1st order: masses, β -decay rates, reaction rates & branching ratios



See review paper: **Mumpower et al.** PPNP 86 (2016)

WHAT DO WE KNOW?

The chart of nuclides

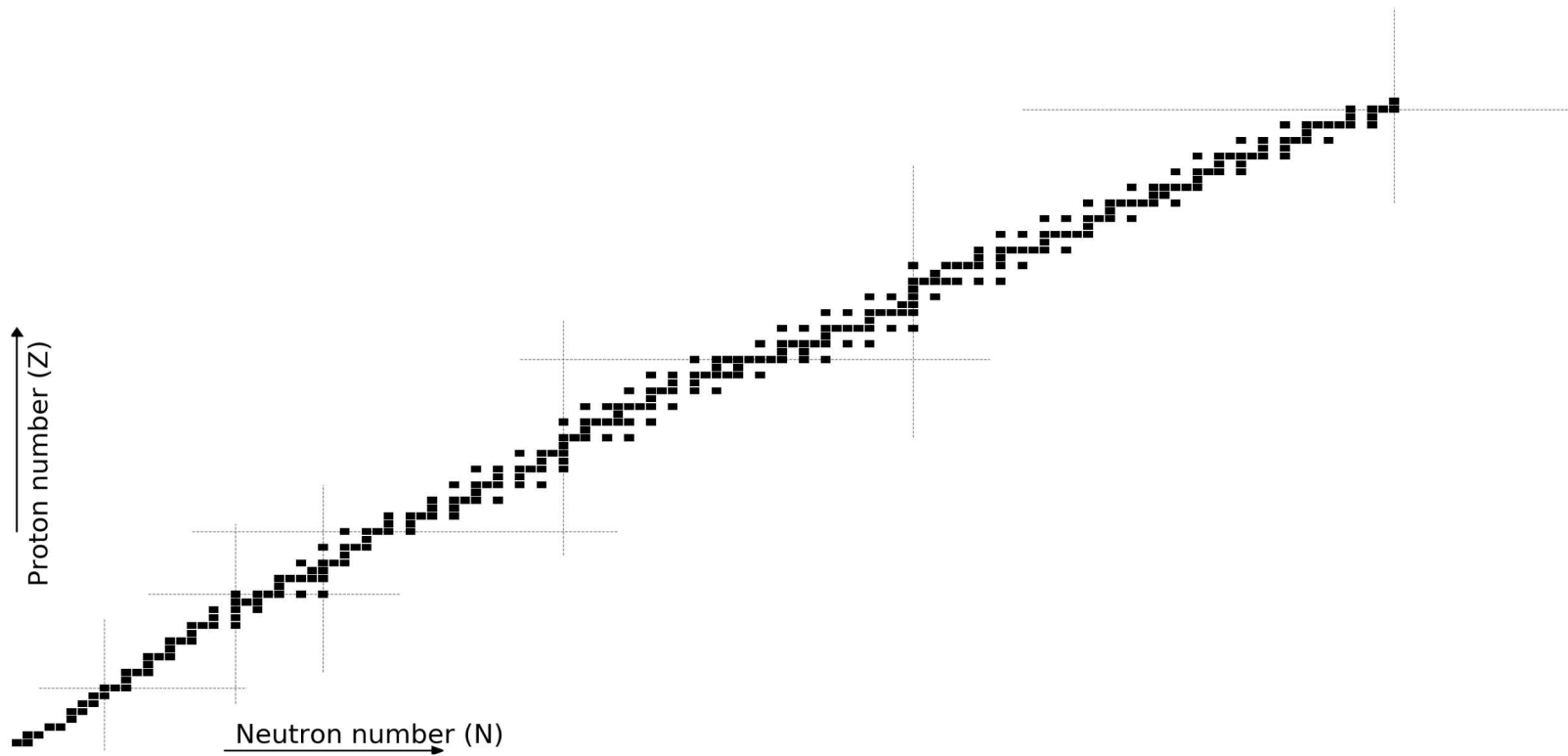


Figure by Mumpower

WHAT DO WE KNOW?

All half-lives

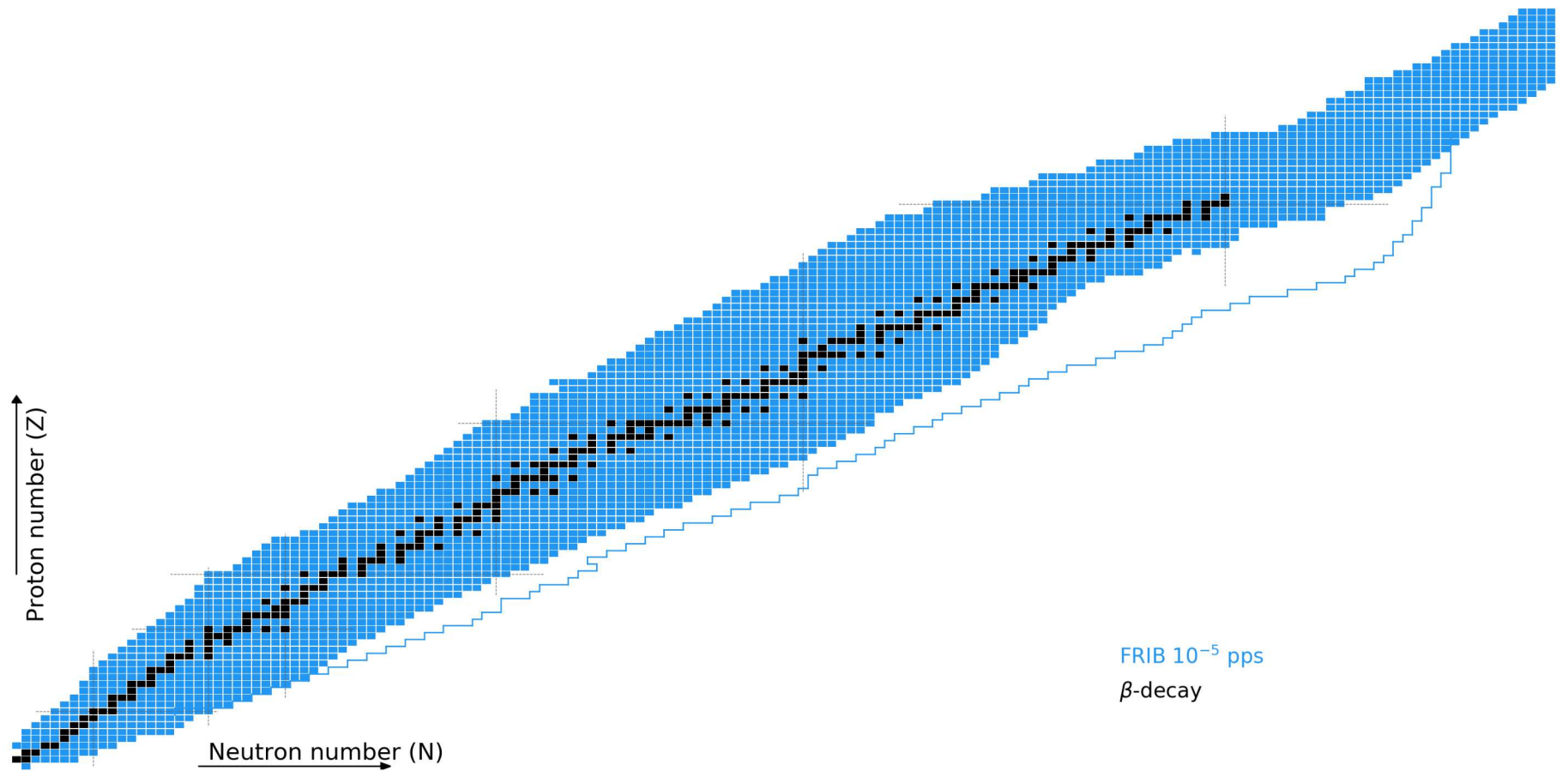


Figure by Mumpower

WHAT DO WE KNOW?

Neutron capture rates

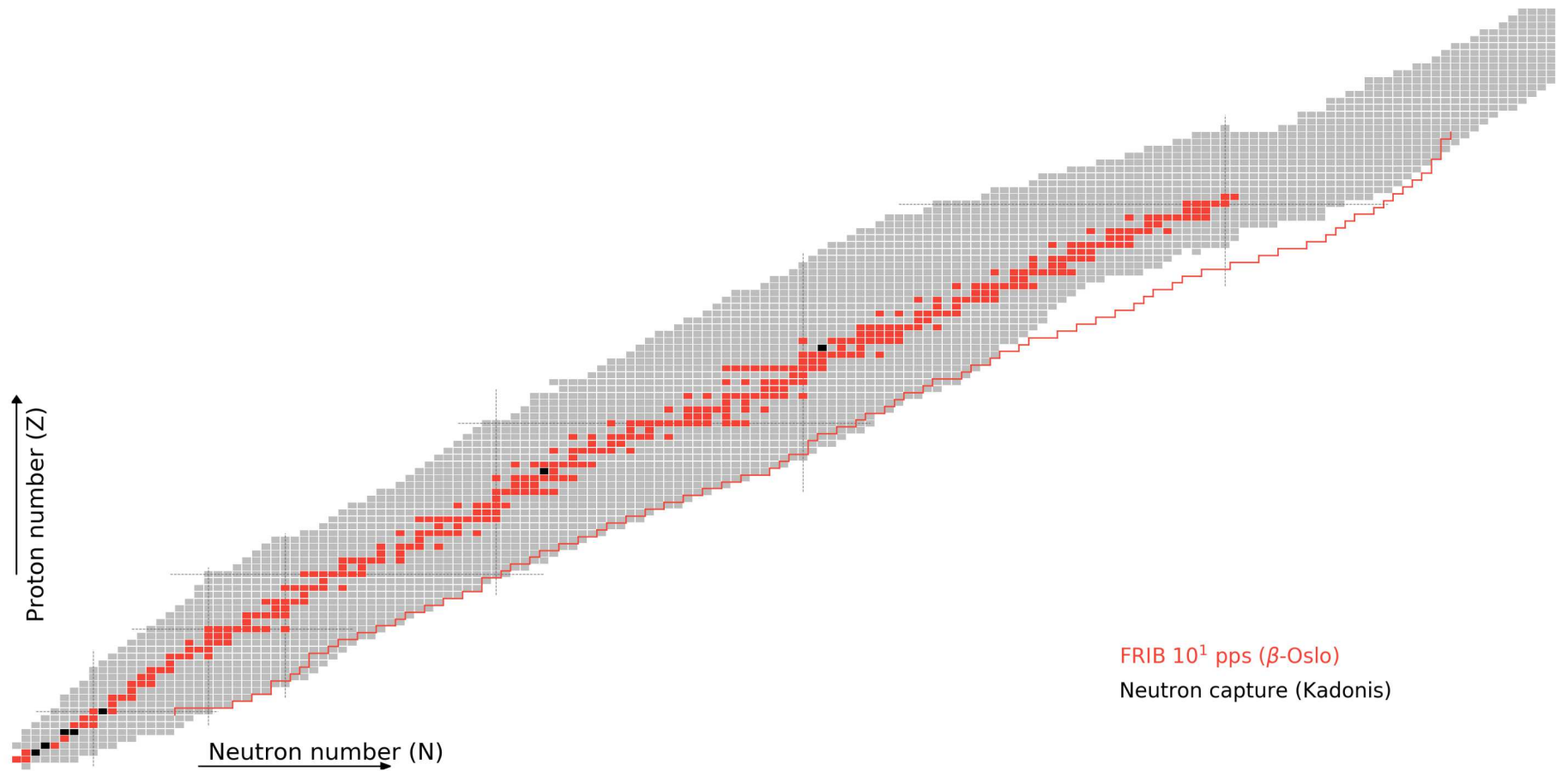
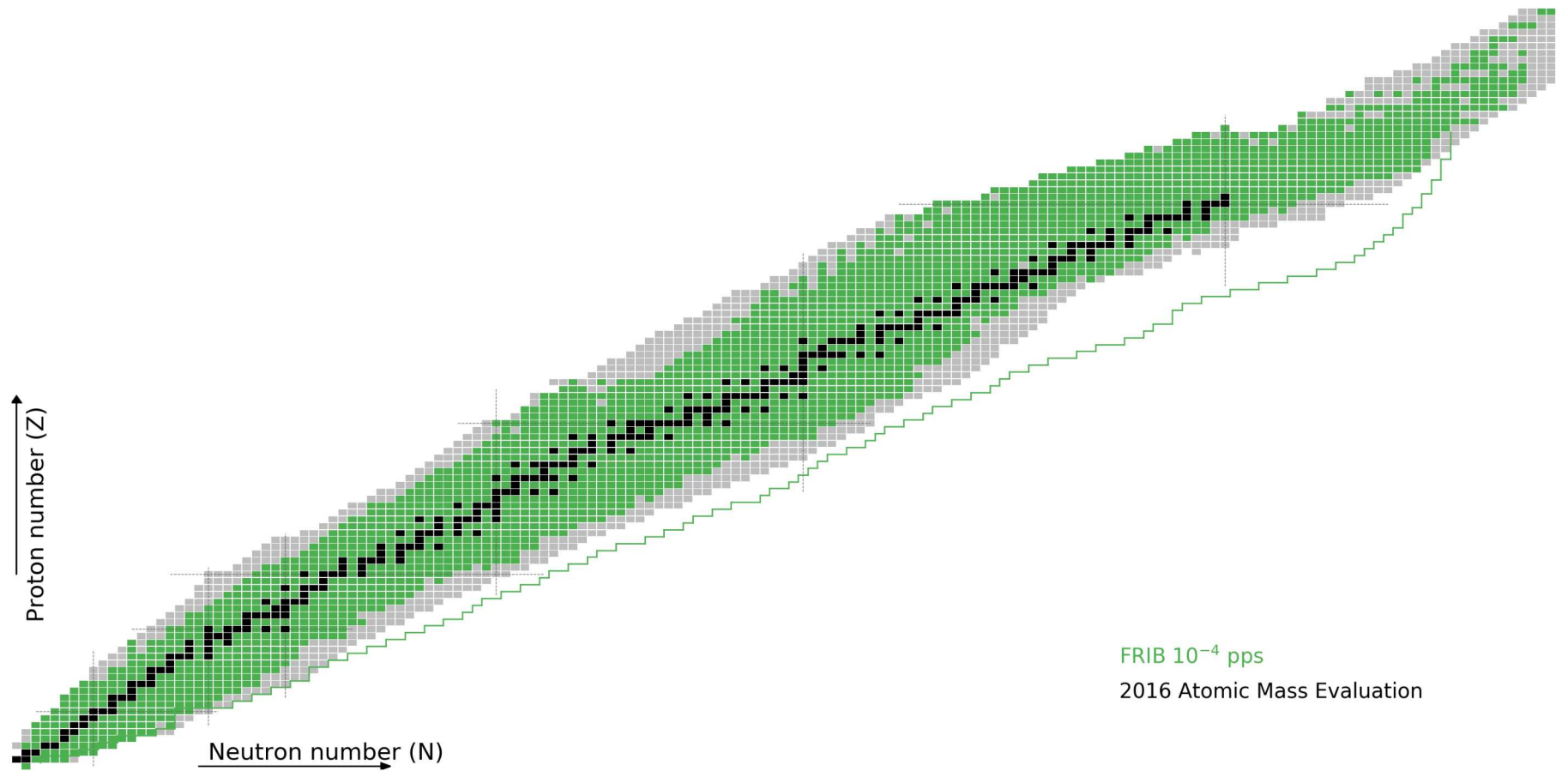


Figure by Mumpower

WHAT DO WE KNOW?

Nuclear masses



FRIB 10^{-4} pps
2016 Atomic Mass Evaluation

Figure by Mumpower

WHAT DO WE KNOW?

As of today, to varying degrees of accuracy

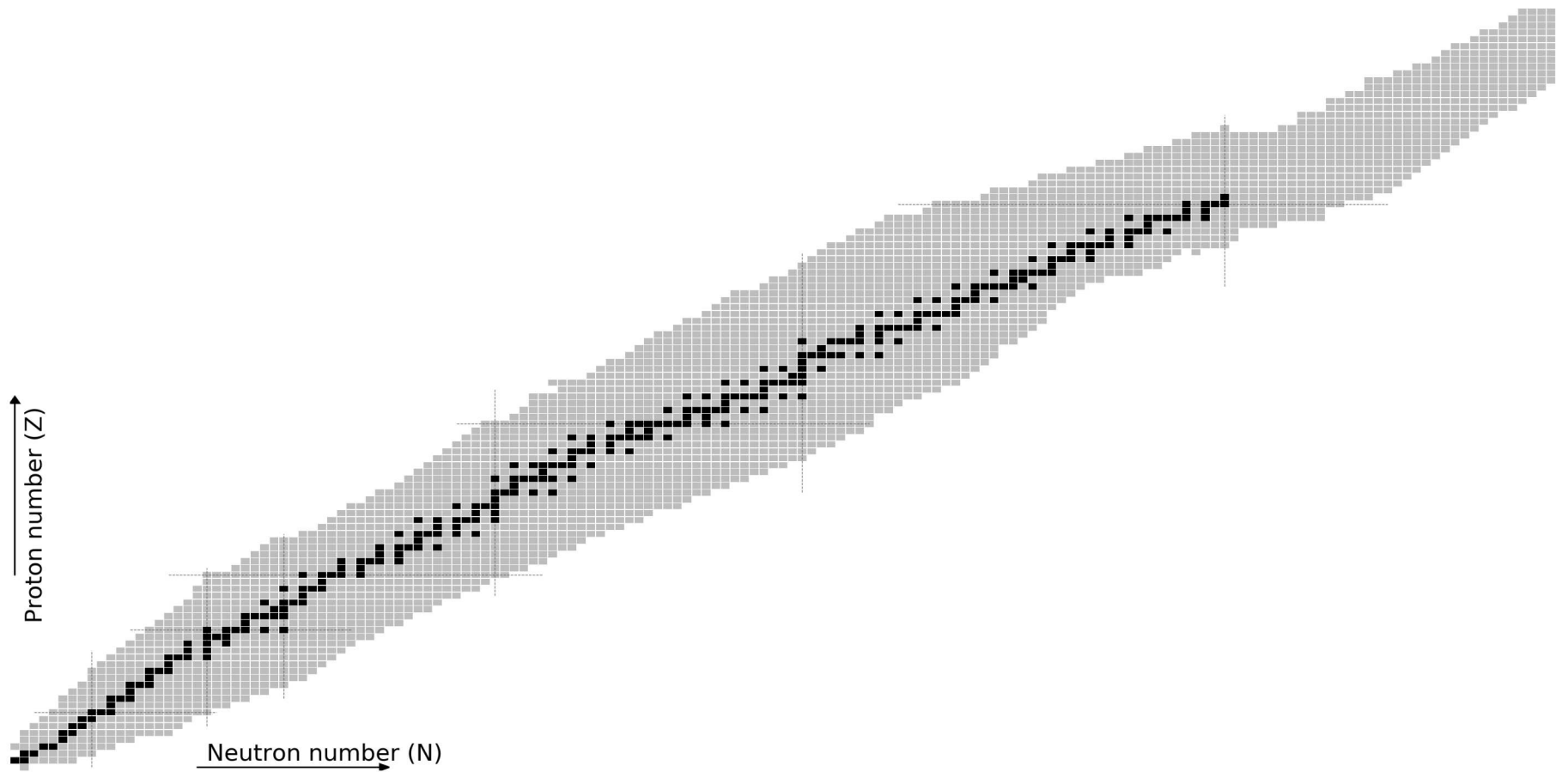


Figure by Mumpower

WHERE WE'RE GOING

FRIB as the **r-process** machine

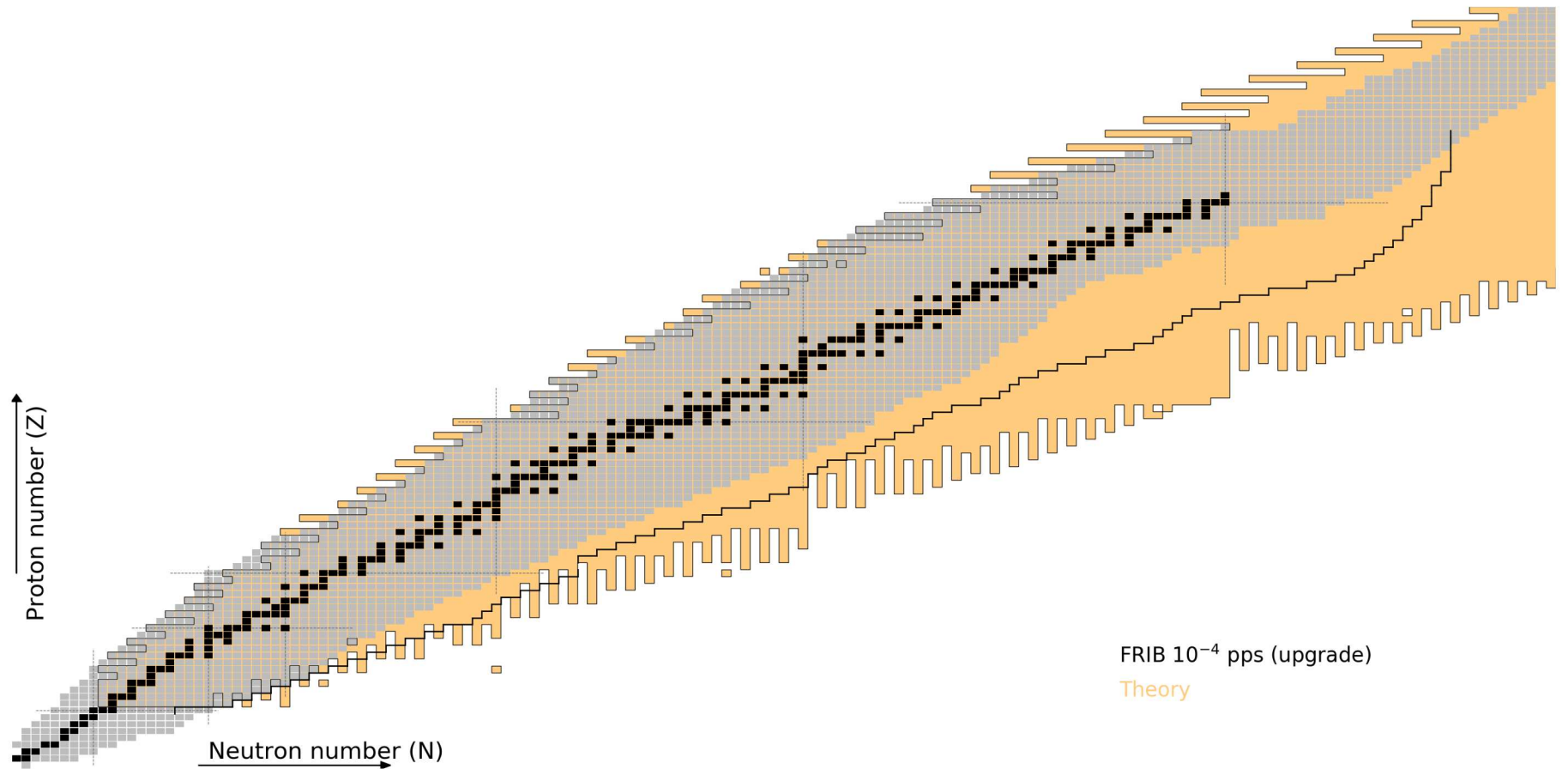


Figure by Mumpower

r-PROCESS CALCULATION

nuclear physics inputs
(S_n , β -rates, n-capture rates, ...)

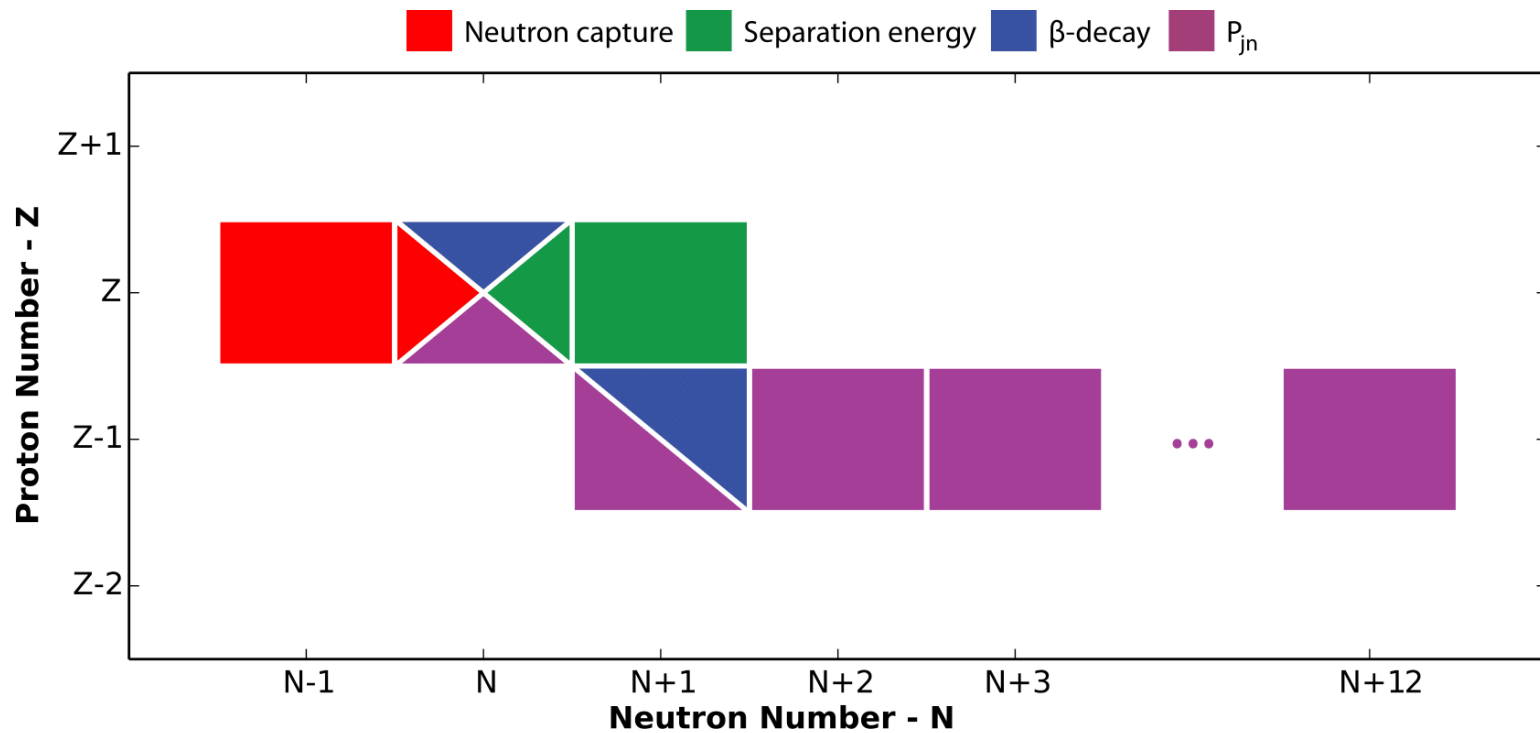


thermodynamic conditions
(temperature, density, ...)

Sprouse & Mumpower in prep (2018)

PRISM: Portable Routines for Integrated nucleoSynthesis Modeling

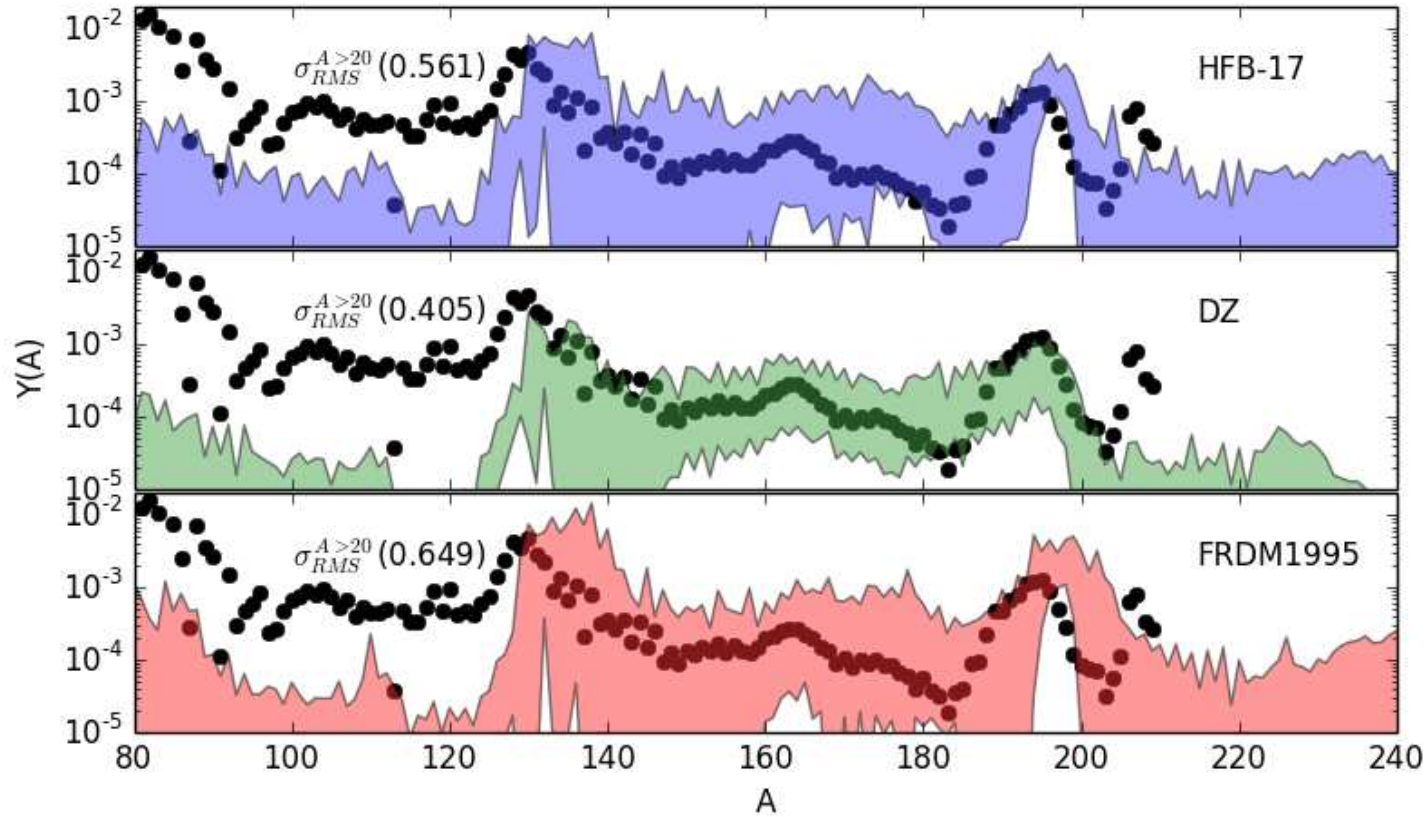
WHY FOCUS ON MASSES?



Masses go into the calculation of all other relevant quantities... Mumpower et al. PRC 92.035807 (2015)

UNCERTAINTIES FROM MASSES

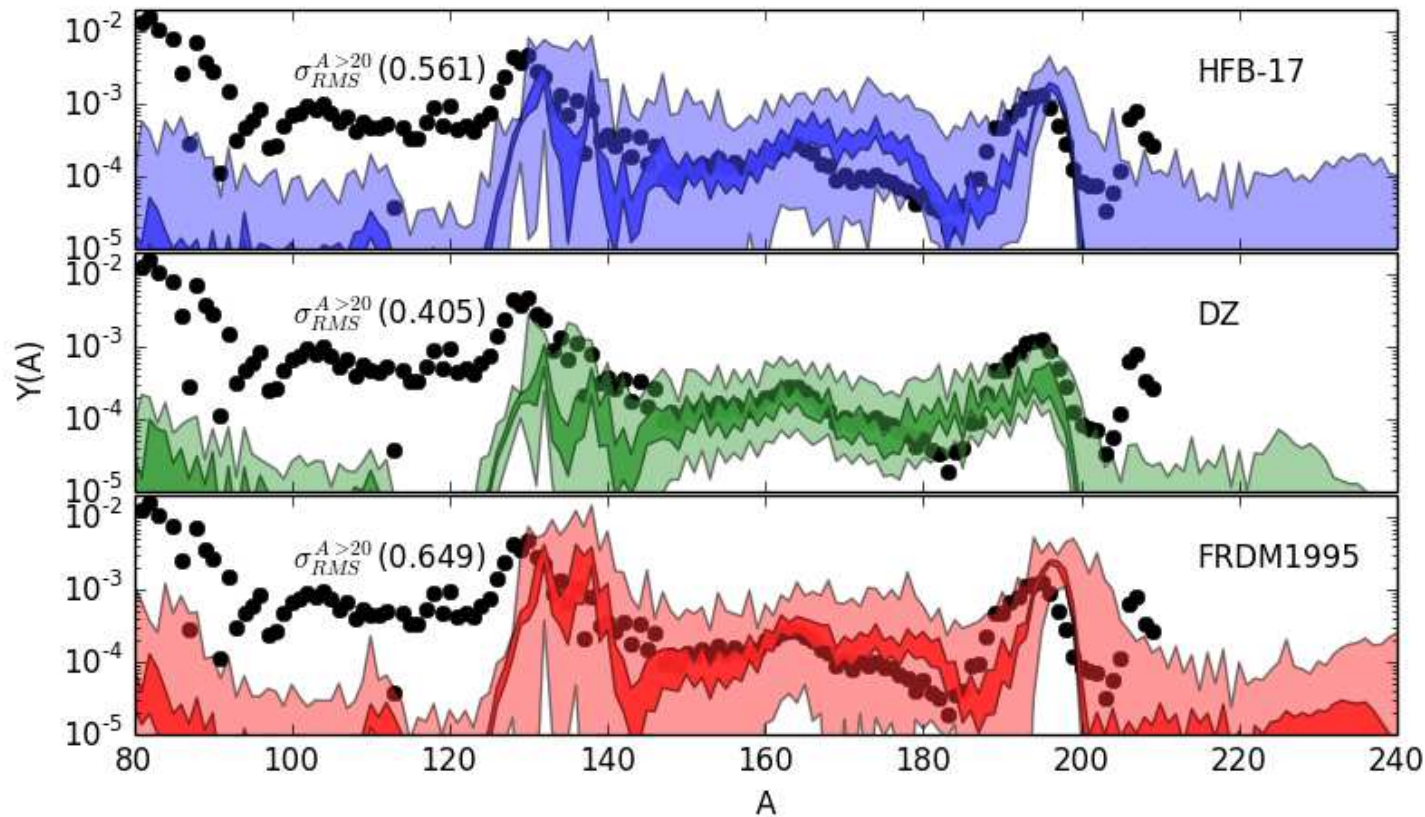
Hot wind: $S \sim 200$, $\tau = 80$ ms, $Y_p = 0.3$



Uncorrelated mass Monte Carlo study; *not full propagation* Mumpower et al., PPNP 86 86-126 (2016)

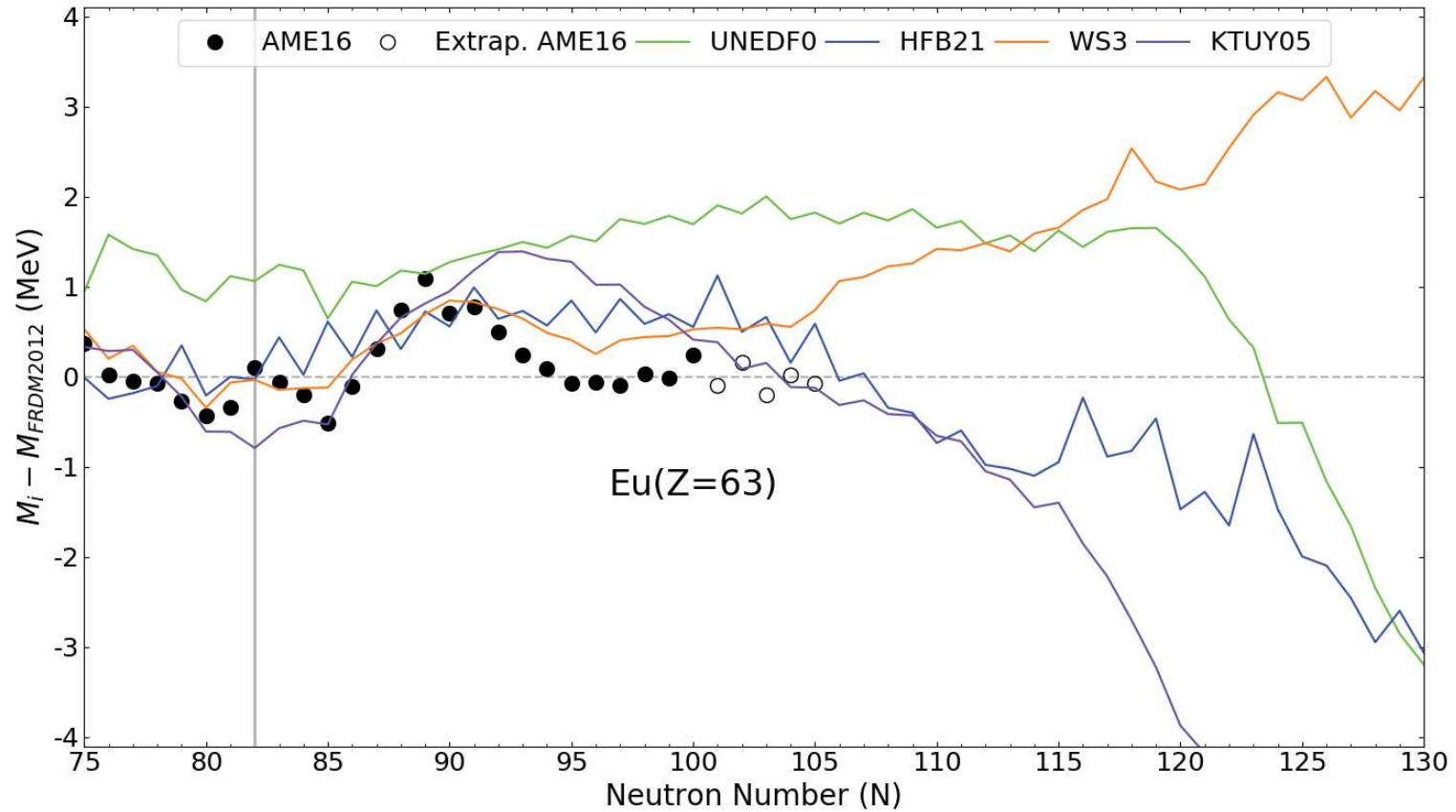
UNCERTAINTIES FROM MASSES

Hot wind: $S \sim 200$, $\tau = 80$ ms, $Y_p = 0.3$



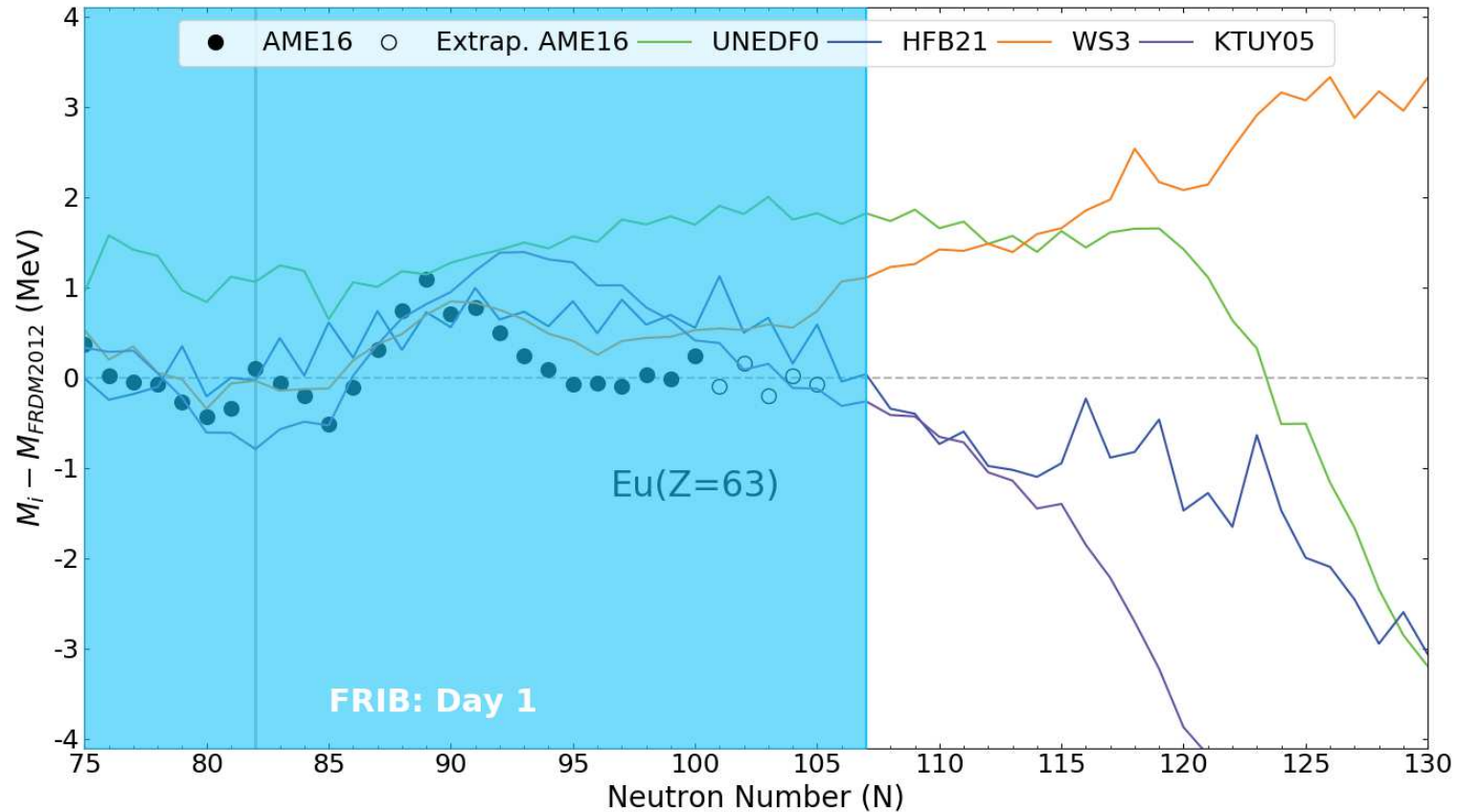
Rule of thumb: $\Delta_{\text{mass}} \sim 500$ keV $\Rightarrow \Delta_Y \sim 2 - 3$ orders of magnitude (2016)

VARIATION IN MASSES



Large variation in mass model predictions further from stability Modified from Mumpower et al. PRNP 16 06 126 (2016)

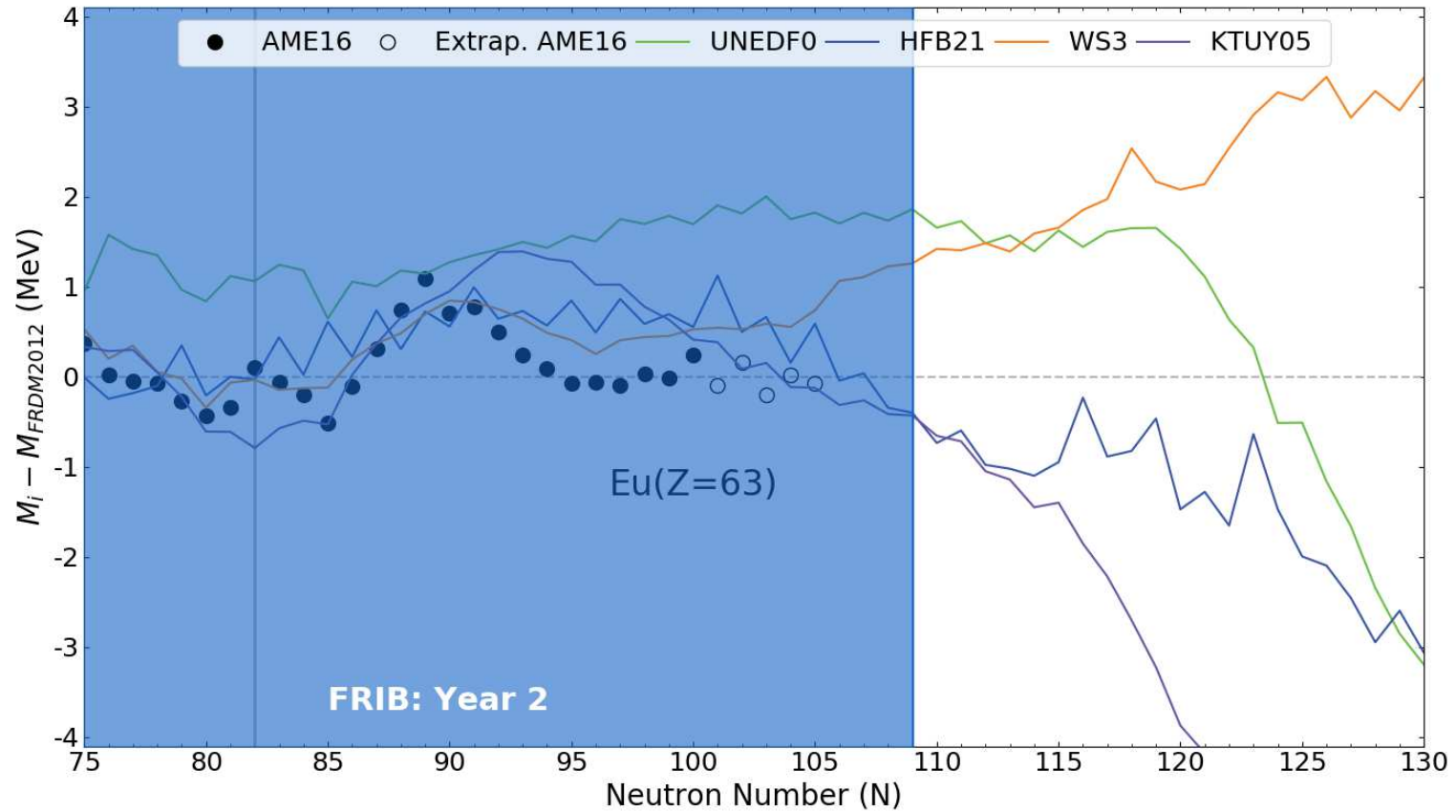
DISCRIMINATORY POWER



FRIB: Day 1 already allows probing of the $N = 104$ subshell closure

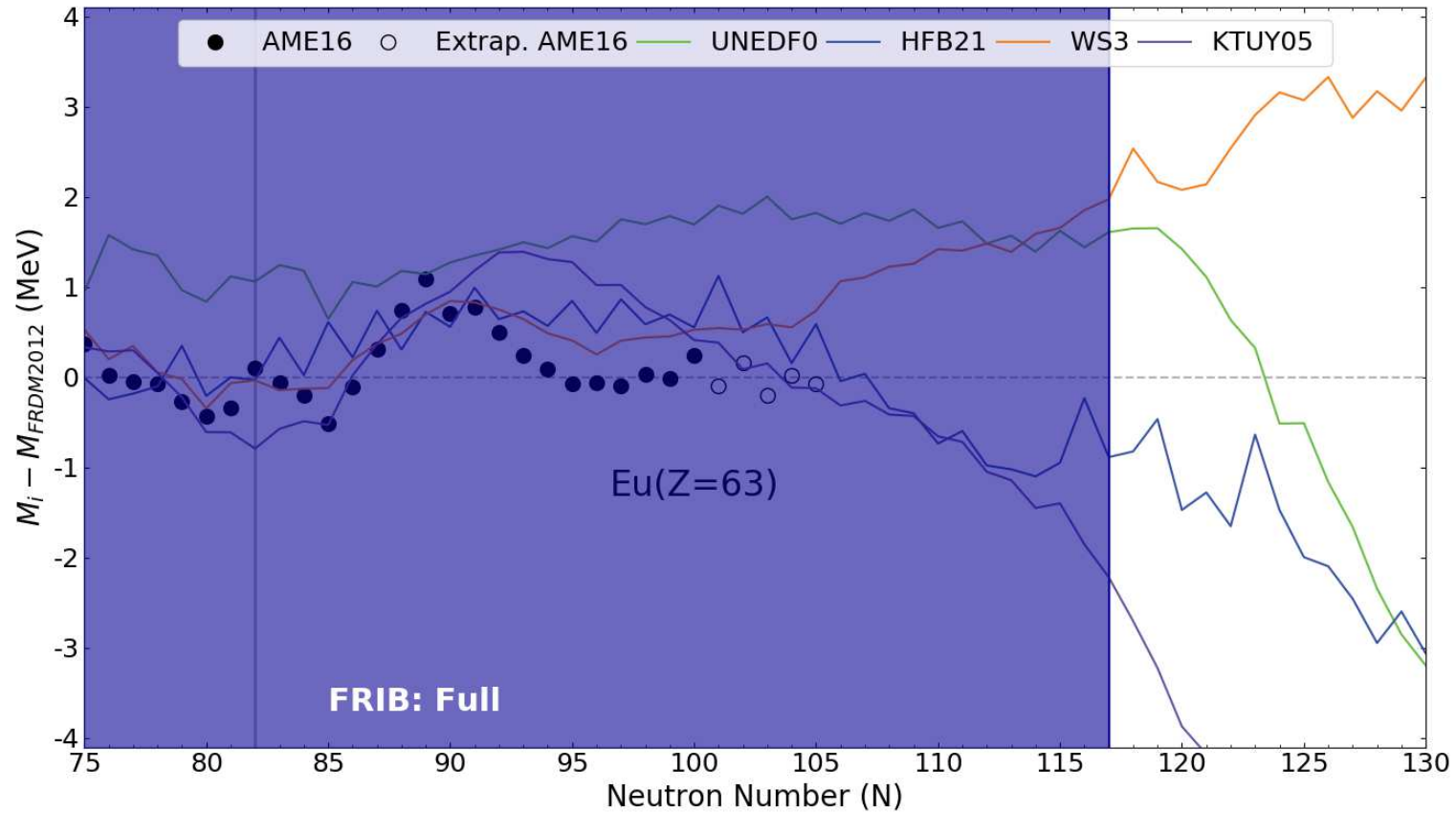
Modified from Mumpower et al., PPNP 86:84-126 (2016)

DISCRIMINATORY POWER



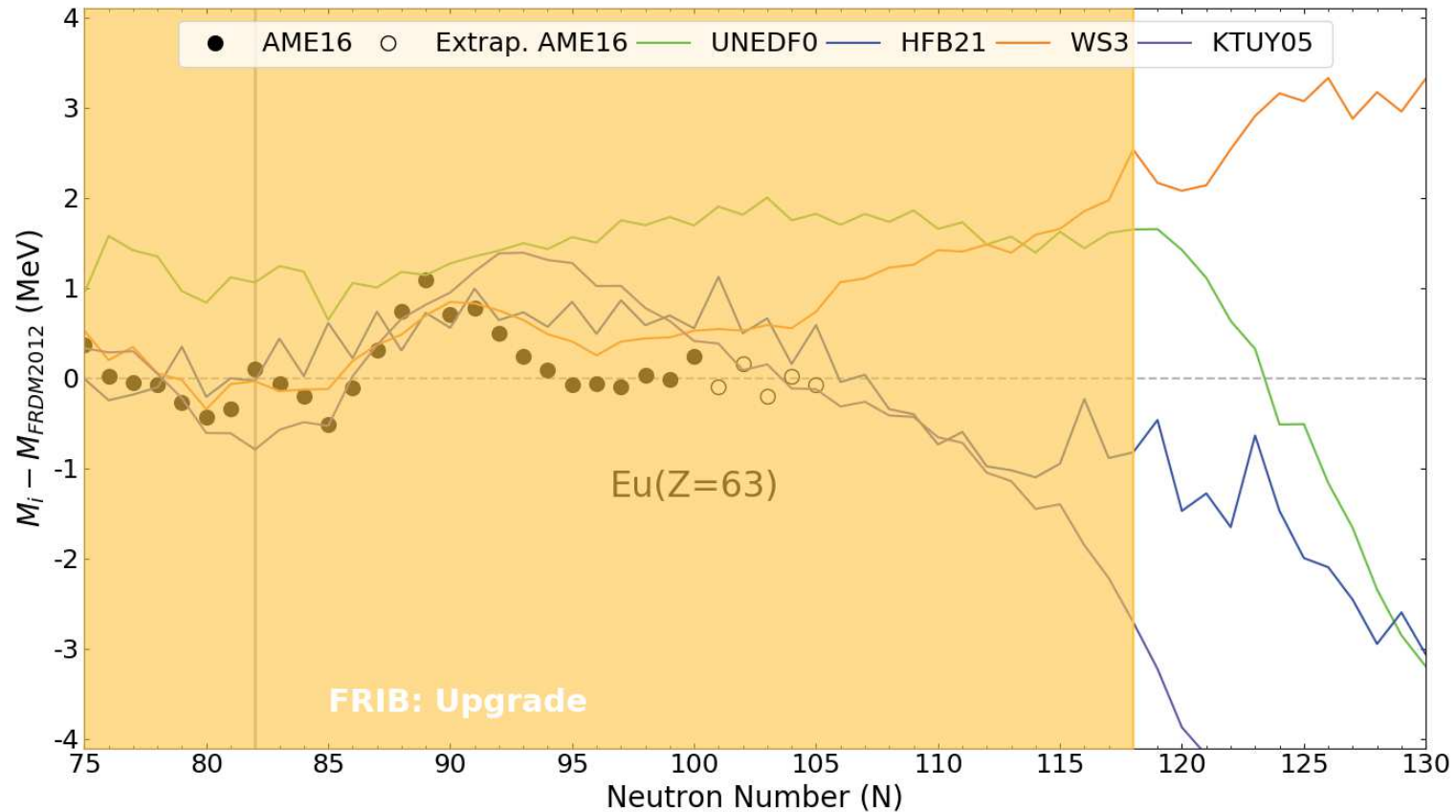
FRIB: Year 2 provides nuclei with several more neutrons Modified from Mumpower et al. PPNP 86-86-126 (2016)

DISCRIMINATORY POWER



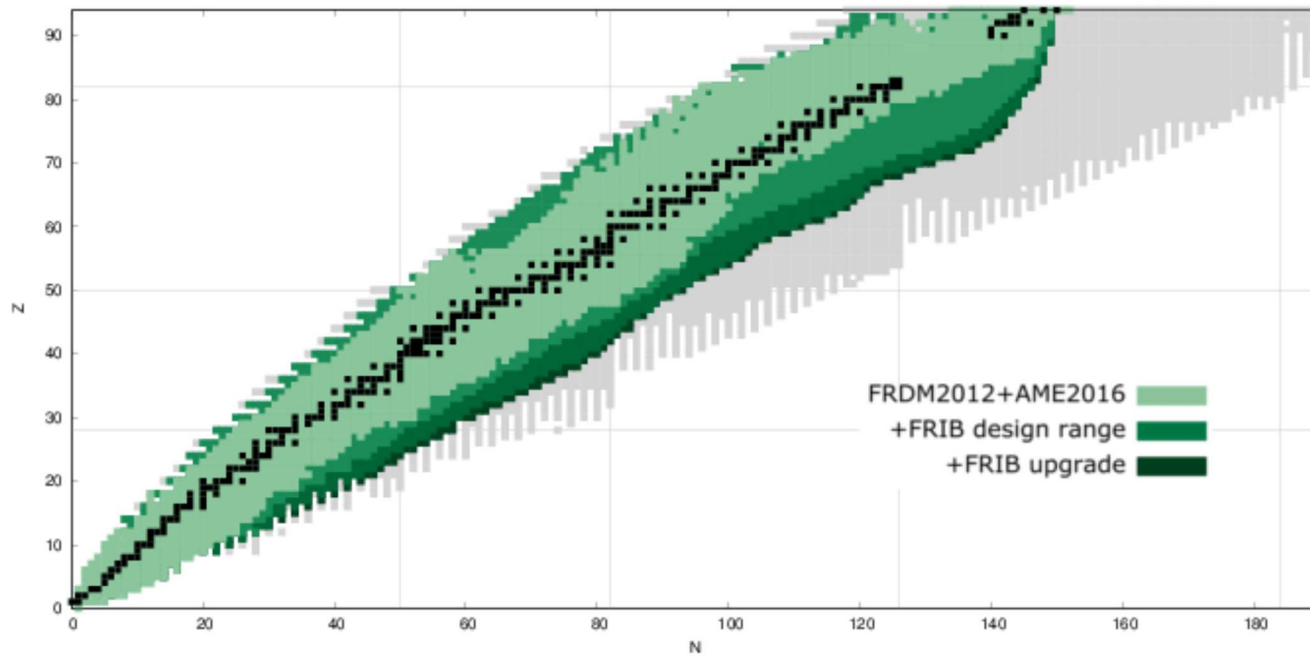
FRIB: Full design spec provides information out to N ~ 117 Modified from Mumpower et al., PNP 868-126 (2016)

DISCRIMINATORY POWER



Upgrade provides 1 or 2 more nuclei per isotopic chain... so what? Modified from Mumpower et al., PPNR 86-86, 126 (2016)

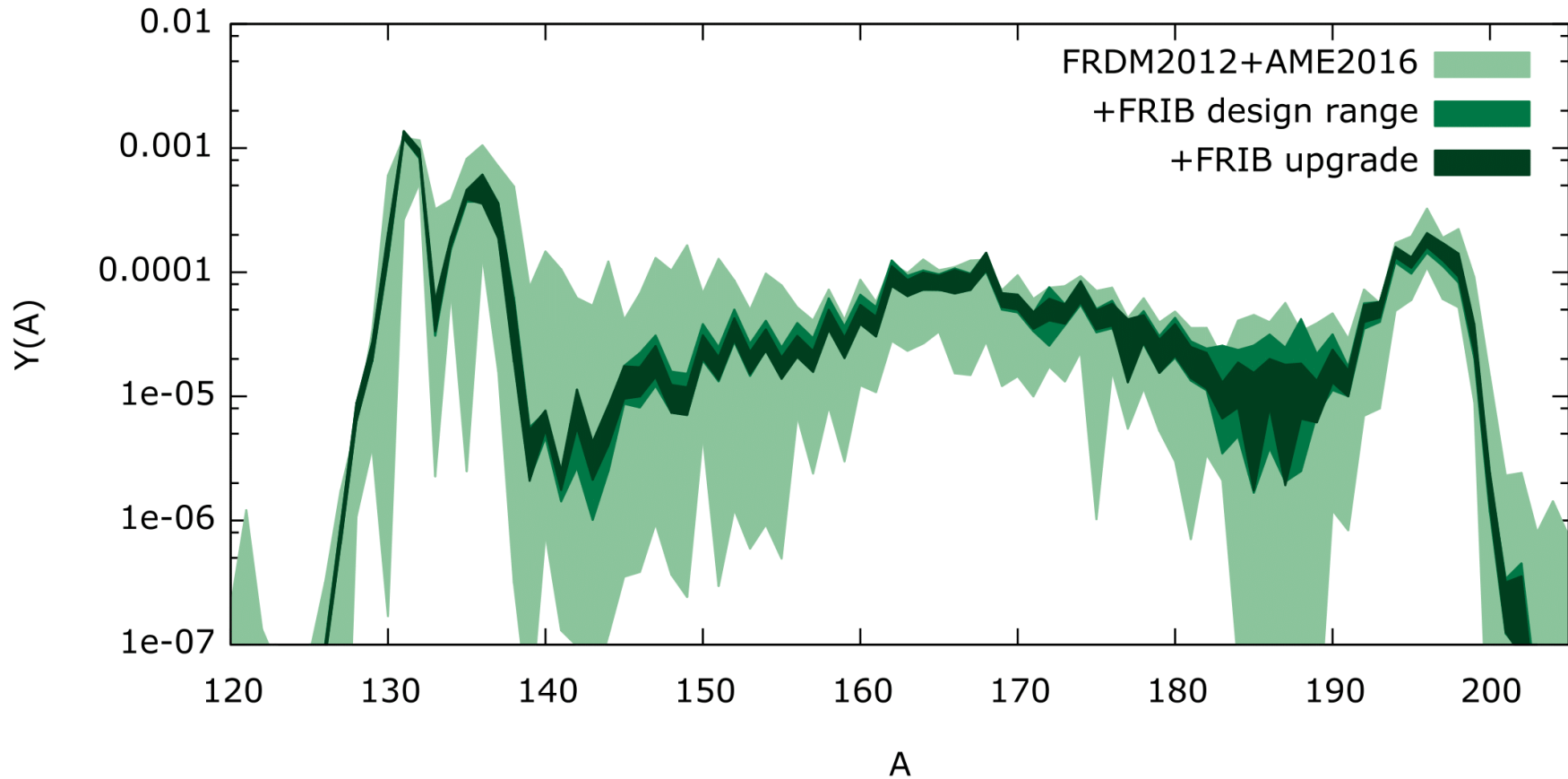
MONTE CARLO RESULTS



Uncorrelated mass Monte Carlo study; *not full propagation*

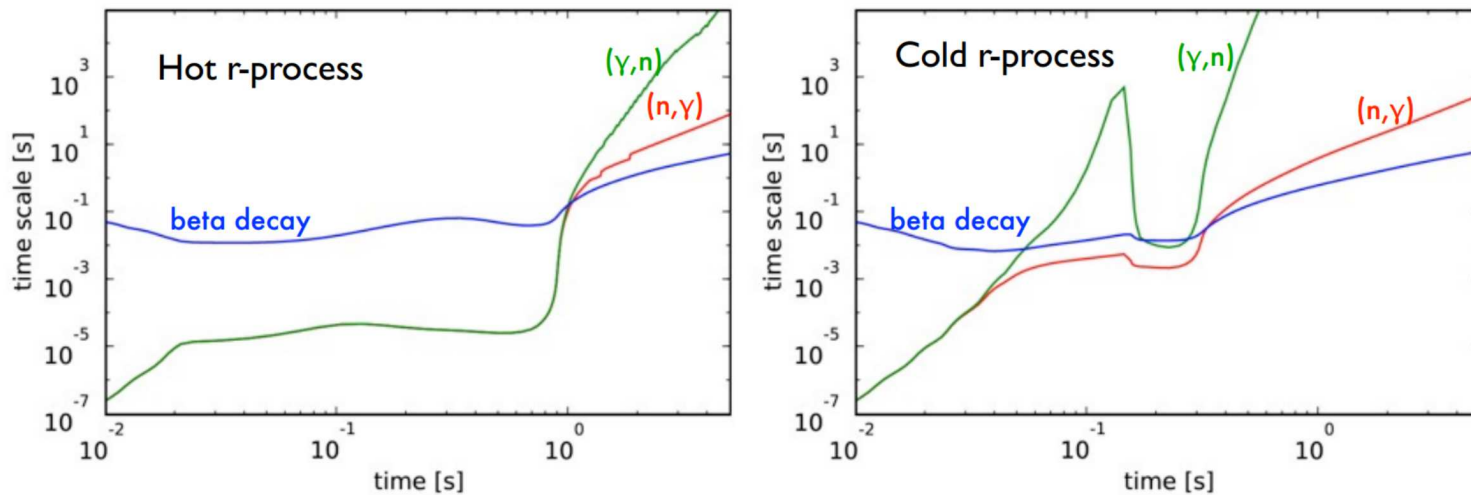
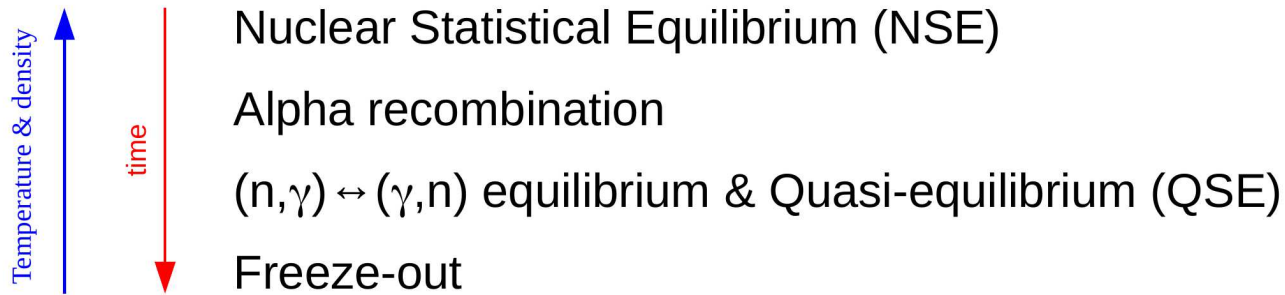
What is the impact of the FRIB upgrade? Surman & Mumpower CGS16 in press (2018)

MONTE CARLO RESULTS



Reduction in uncertainties in Rare Earth and transition regions Surman & Mumpower, CGS16 in press (2018)

A TYPICAL r-PROCESS CALCULATION



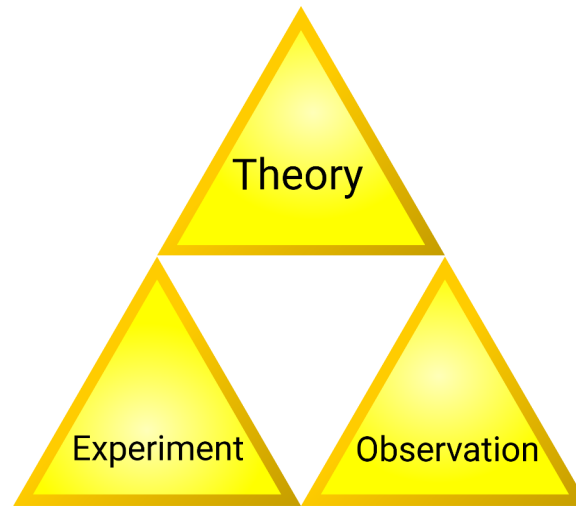
nuclear physics + astrophysics \mapsto abundances

Figure by A. Arcones (2011)

REVERSE ENGINEERING

What if we take a different approach?

Constrain **nuclear physics** with experiment and *additionally* observation using **feedback** from our calculated **abundances**



If we try to fit a particular part of the pattern we can ask what nuclear properties are responsible for its formation and learn how they are required to evolve with neutron excess

REVERSE ENGINEERING

Our pursuit must satisfy several constraints:

We must be able to make measurements on these nuclei

Limits us to nuclei closer to stability

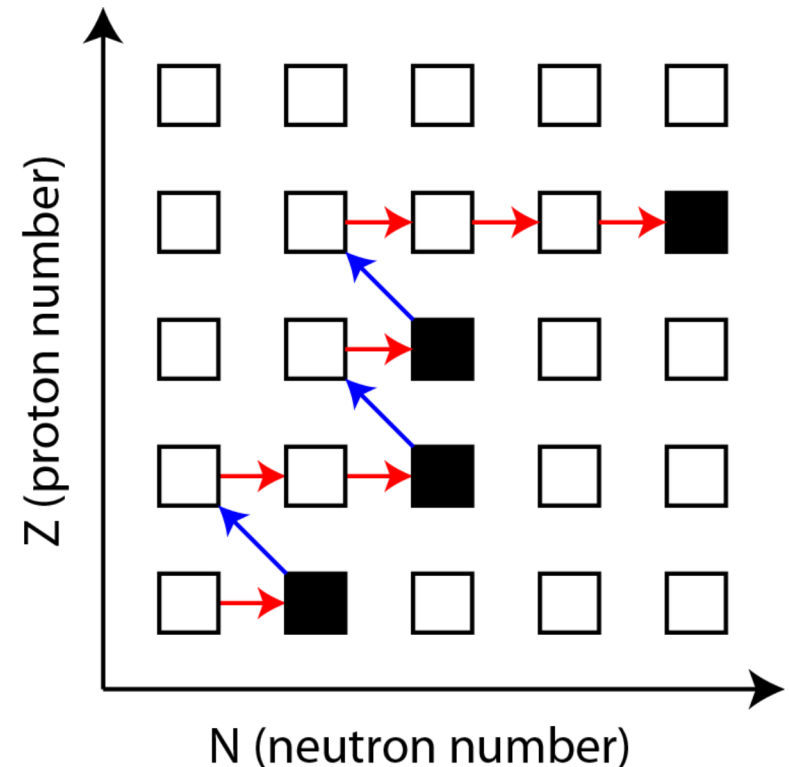
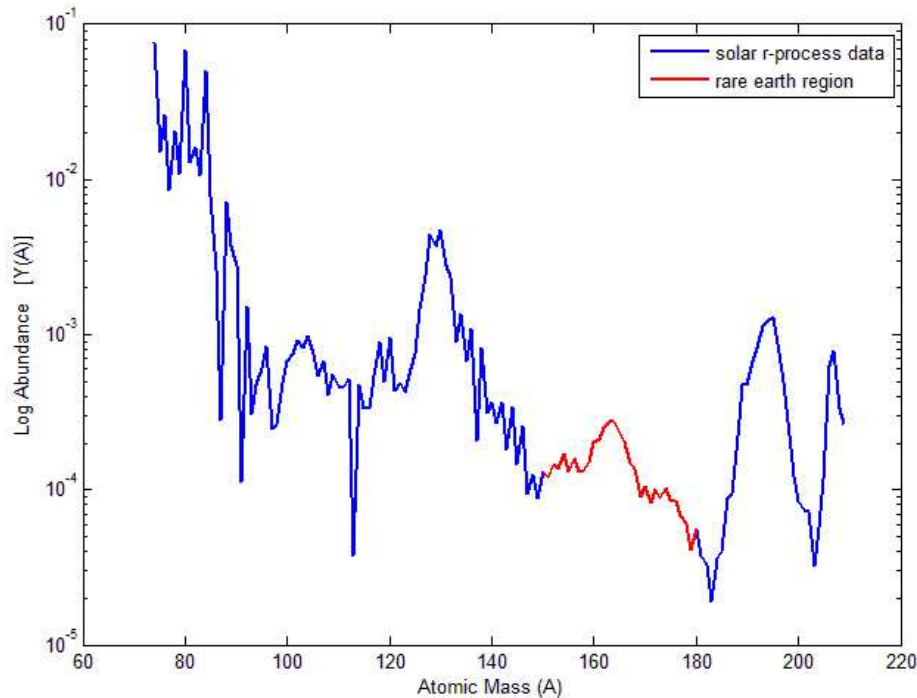
So, we must explore the freeze-out phase of the r process

We must be able to use a recognizable signature in the abundances

The rare earth peak

FORMATION OF THE RARE EARTH PEAK

PROPOSED WAYS TO FORM THE REP



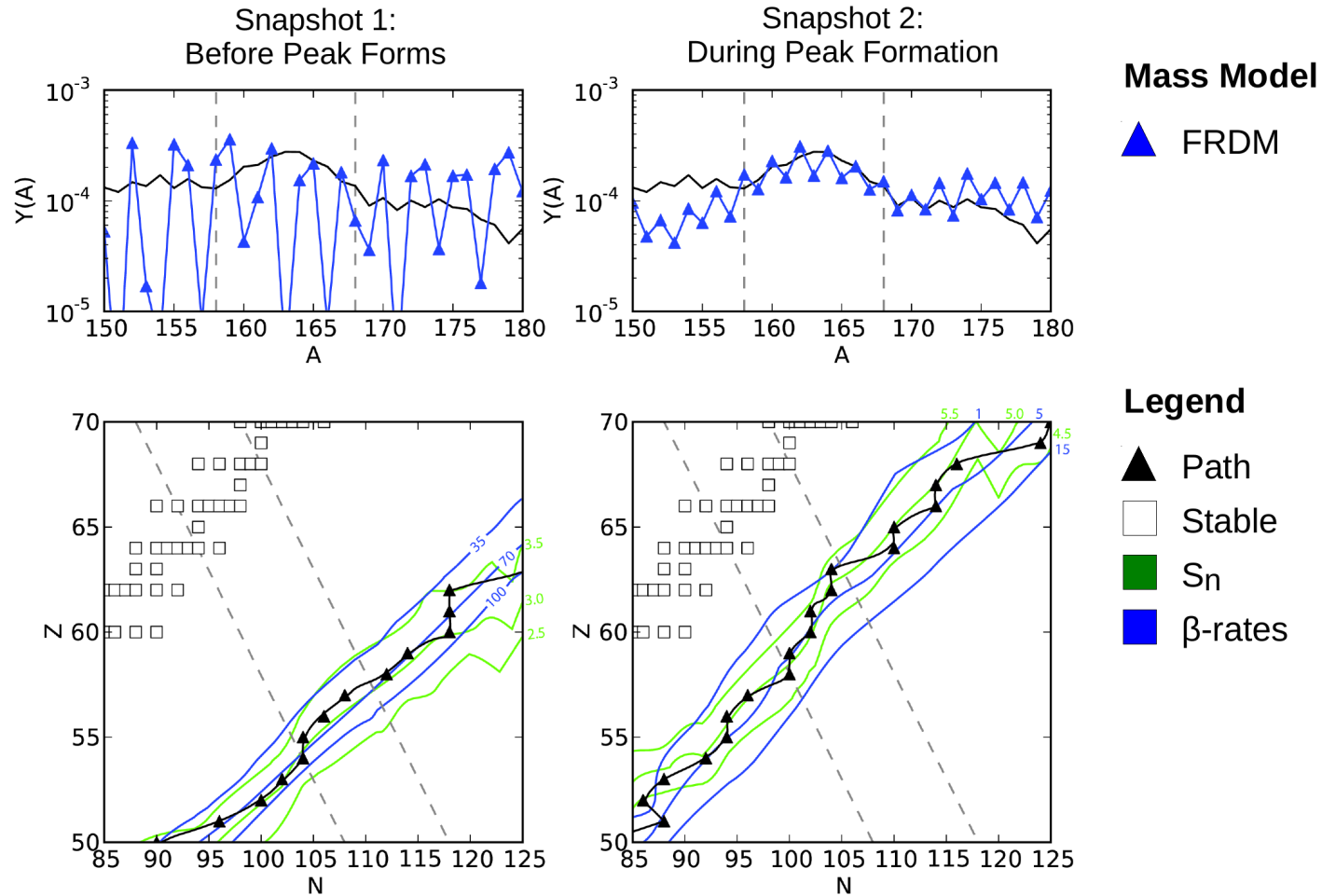
1. Dynamical formation during freeze-out ($R \lesssim 1$)
Requires a localized nuclear structure effect (kink)
2. Via fission fragment yields

Surman & Engel, PRL 79, 1809 (1997); Mumpower et al., PRC 85, 045801 (2012); Mumpower et al., ApJ 752, 117 (2012)

Requires dumping heavy products in exactly the right spot

FORMATION OF REP

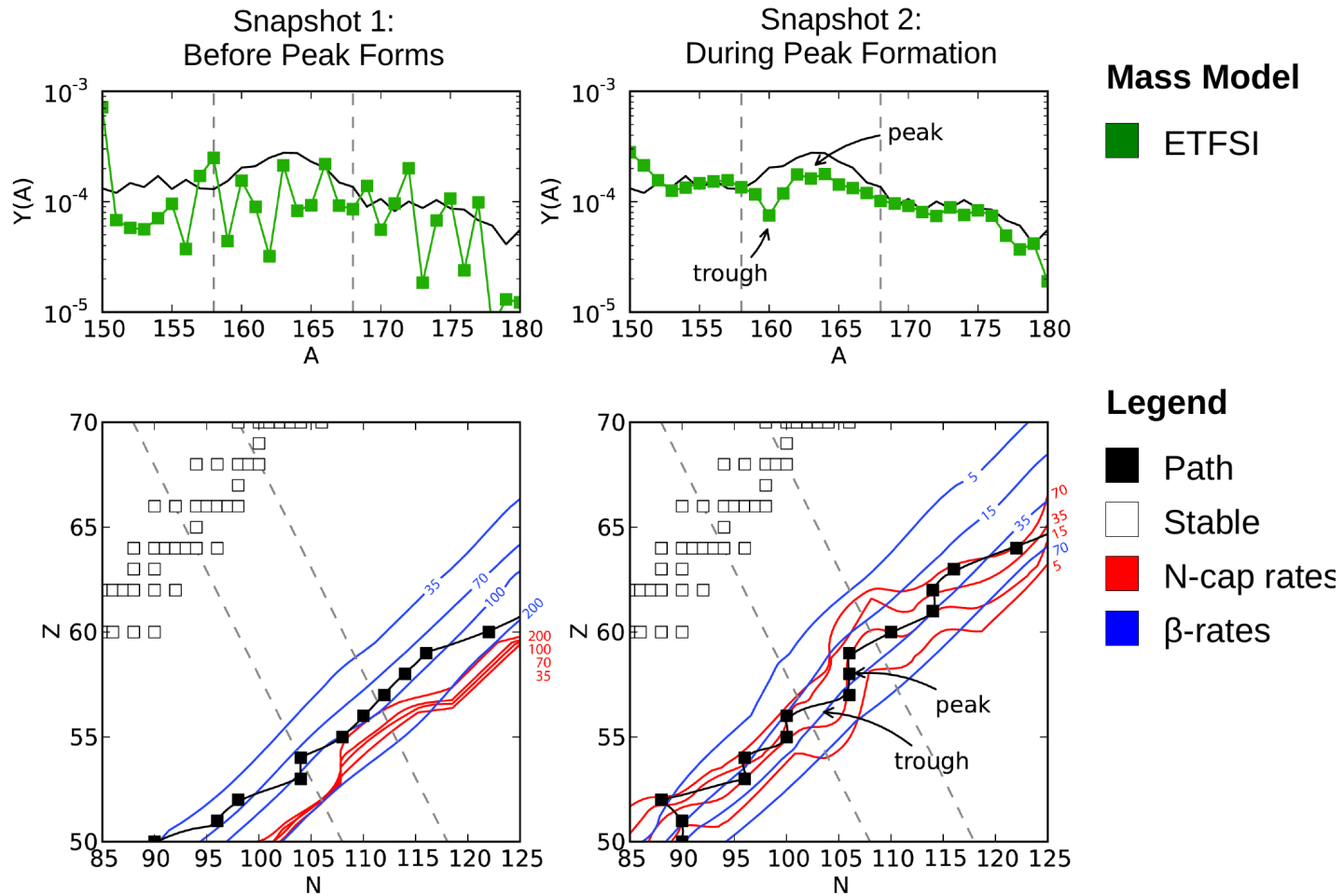
Hot wind: $S \sim 200$, $\tau = 80$ ms, $Y_e = 0.3$



Kink in separation energies forms peak under **hot** freeze-out conditions Mumpower et al. PRC 85, 045801 (2012)

FORMATION OF REP

Cold wind: $S \sim 300$, $\tau = 80$ ms, $Y_e = 0.4$



Kink in neutron capture rates forms peak under cold freeze-out conditions [Mumpower et al. PRC 85,045801 \(2012\)](#)

REP FORMATION: IDEAL CANDIDATE

We choose to study method 1 for reverse engineering

1. Dynamical formation during freeze-out ($R \lesssim 1$)

Requires a localized nuclear structure (kink)

Relatively few nuclei to measure, close to stability

Hints from Jin Wu's $T_{1/2}$ measurements

Very close to making necessary mass measurements

2. Via fission fragment yields

Requires dumping heavy products in exactly the right spot

Extreme r-process conditions necessary

Need to make measurements on hundreds of the heaviest nuclei

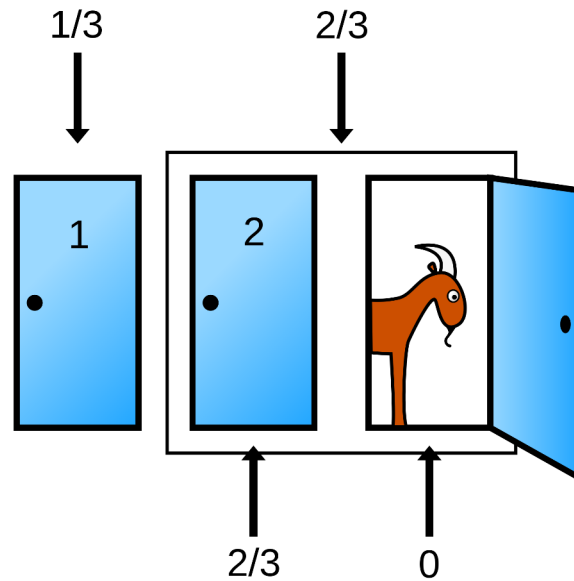
Problem: We can't reach these nuclei, even with FRIB

THE REVERSE ENGINEERING FRAMEWORK

THE BAYESIAN APPROACH

An example... The Monty Hall problem

A new car is hidden behind one of the doors



The optimal strategy is to switch the initial pick - twice the chance of winning the new car

We **update** our probabilities based off new information

APPLY IDEA TO REP FORMATION

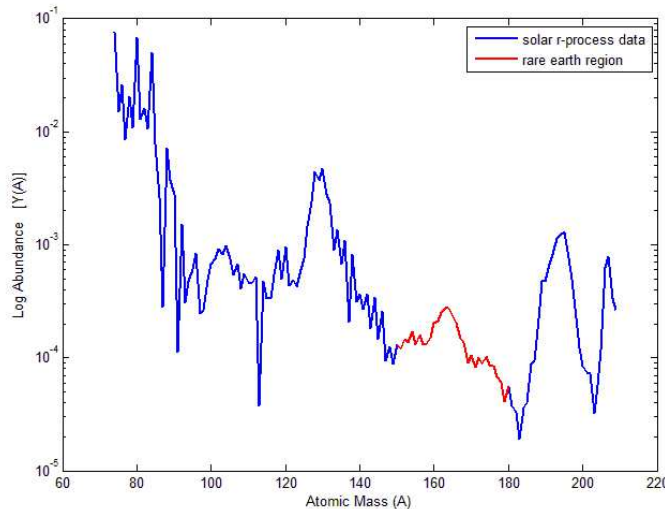
For fixed astrophysical conditions (**hot**, **cold** or **merger**)...

Let's allow the nuclear masses to vary

We have to update all relevant nuclear physics self-consistently

The rare earth abundances provide **feedback** to the change in masses

Use the Metropolis algorithm to traverse the parameter space

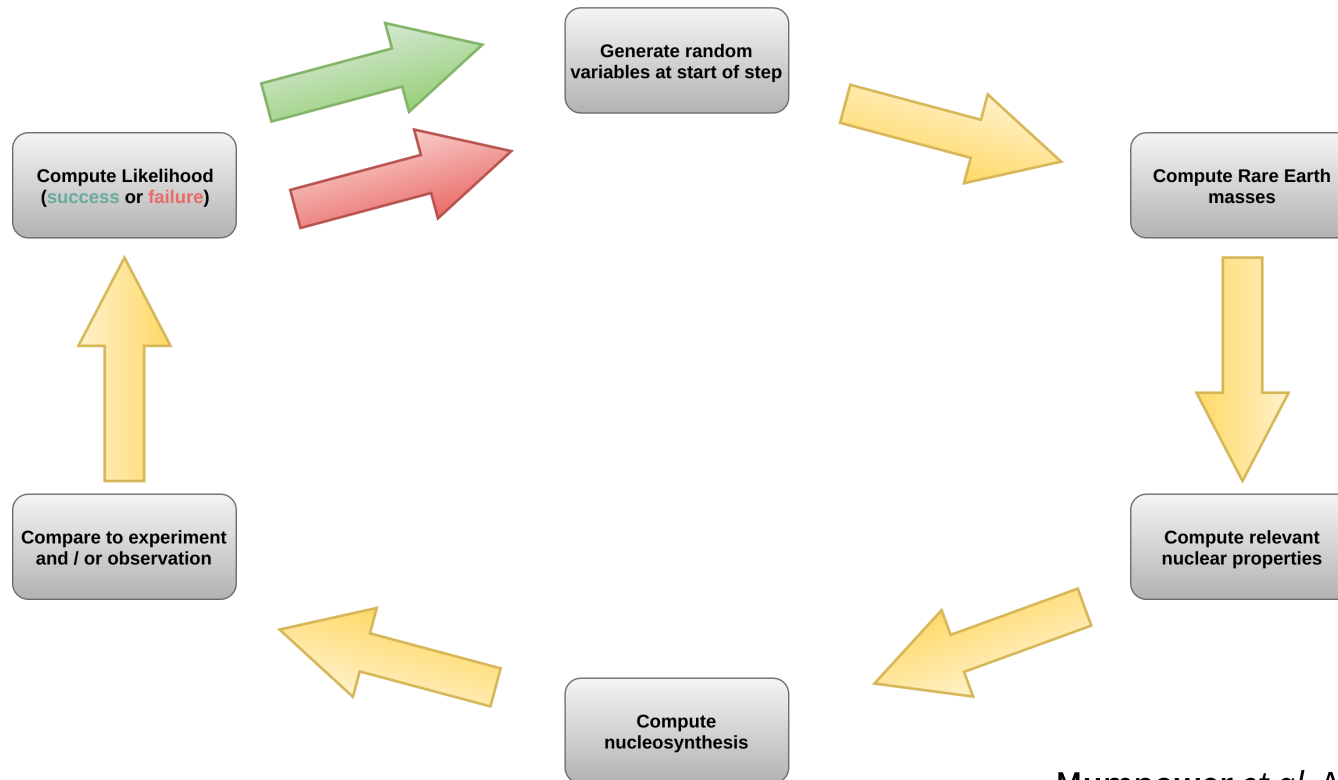


Compute likelihood
 $L \sim$ **match red abundances**

Mumpower *et al.* ApJ 833 282 (2016)

REVERSE ENGINEERING PROCEDURE

How it works in a nutshell



Mumpower *et al.* ApJ 833 282 (2016)

UPDATING NUCLEAR PROPERTIES

Every time the masses change we recalculate...

Relevant Q-values

β -decay properties ($T_{1/2}$ and branching ratios)

Neutron capture rates

For hundreds of nuclei...

This is computationally expensive but necessary!

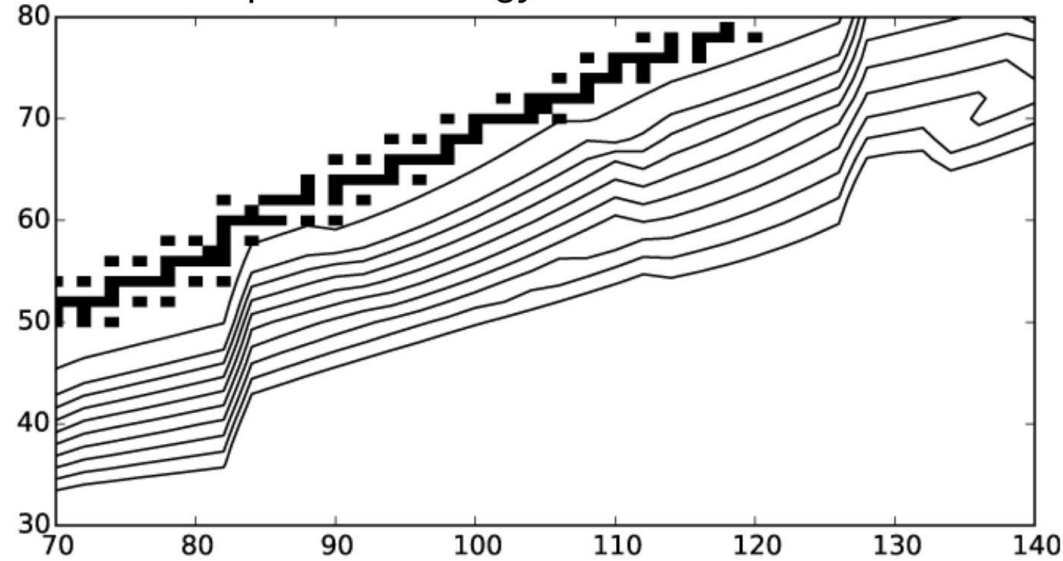
Mumpower *et al.* PRC 92 035807 (2015)

RESULTS OF REVERSE ENGINEERING

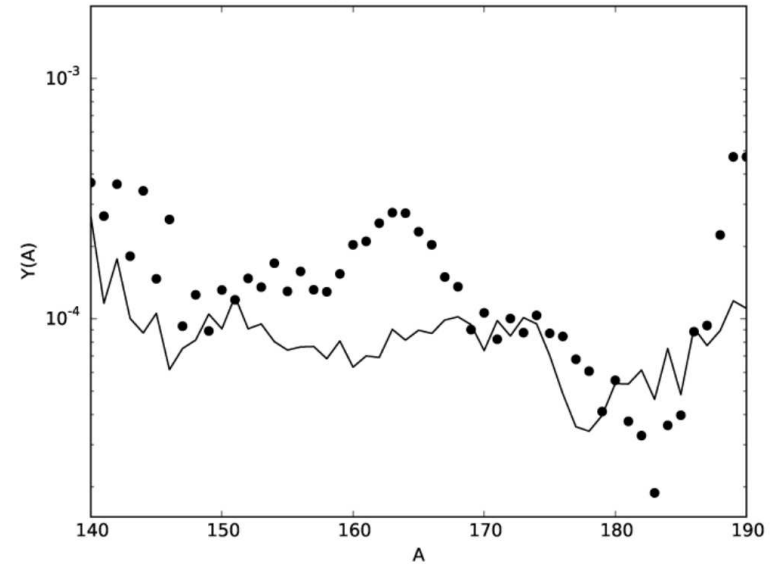
FIRST ATTEMPT

Hot wind r-process with default DZ parameters

1 neutron separation energy contours for even-N nuclei



Predicted r-process abundances



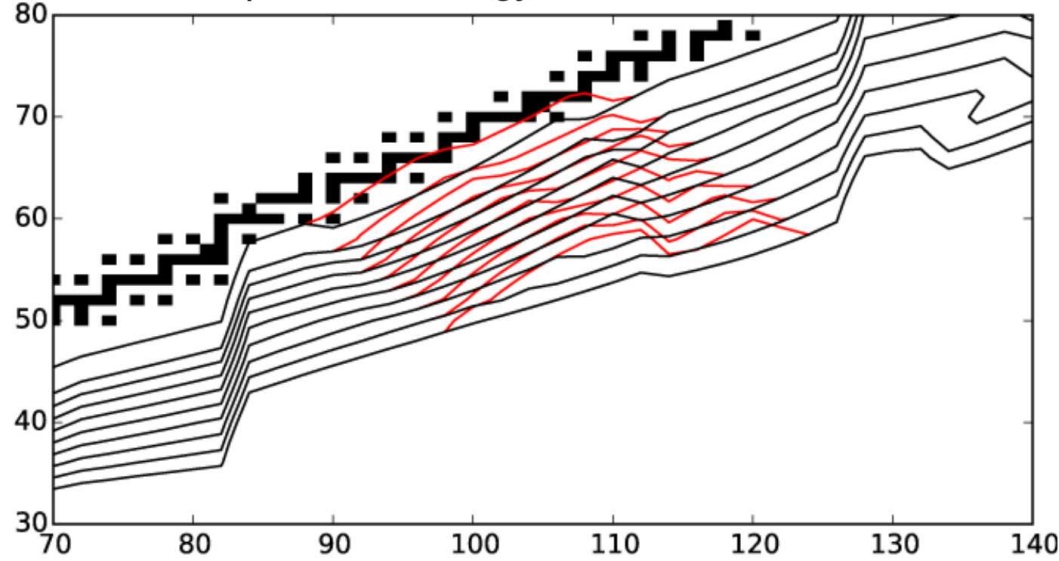
Success?! ... We found a peak! But there's a problem!

L ~ match red abundances Mumpower et al. J. Phys. G 44 3 034003 (2017)

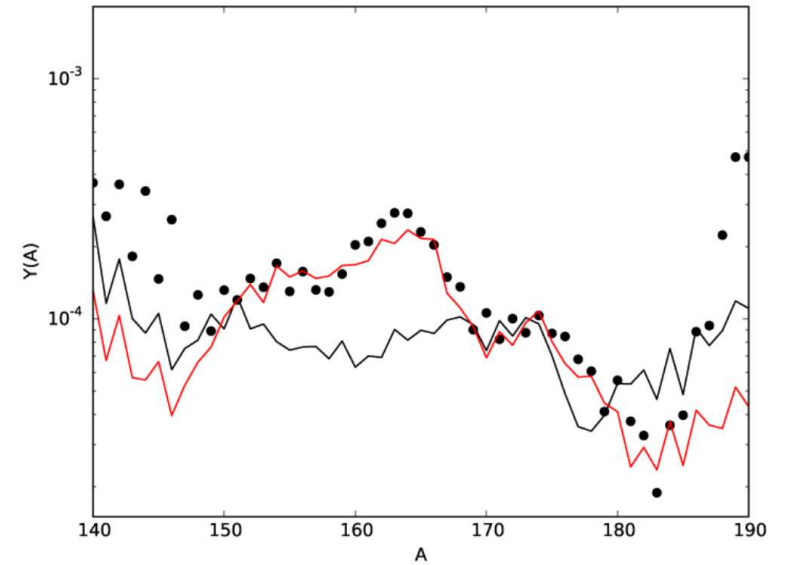
FIRST ATTEMPT

Hot wind r-process with **new** DZ parameters

1 neutron separation energy contours for even-N nuclei



Predicted r-process abundances



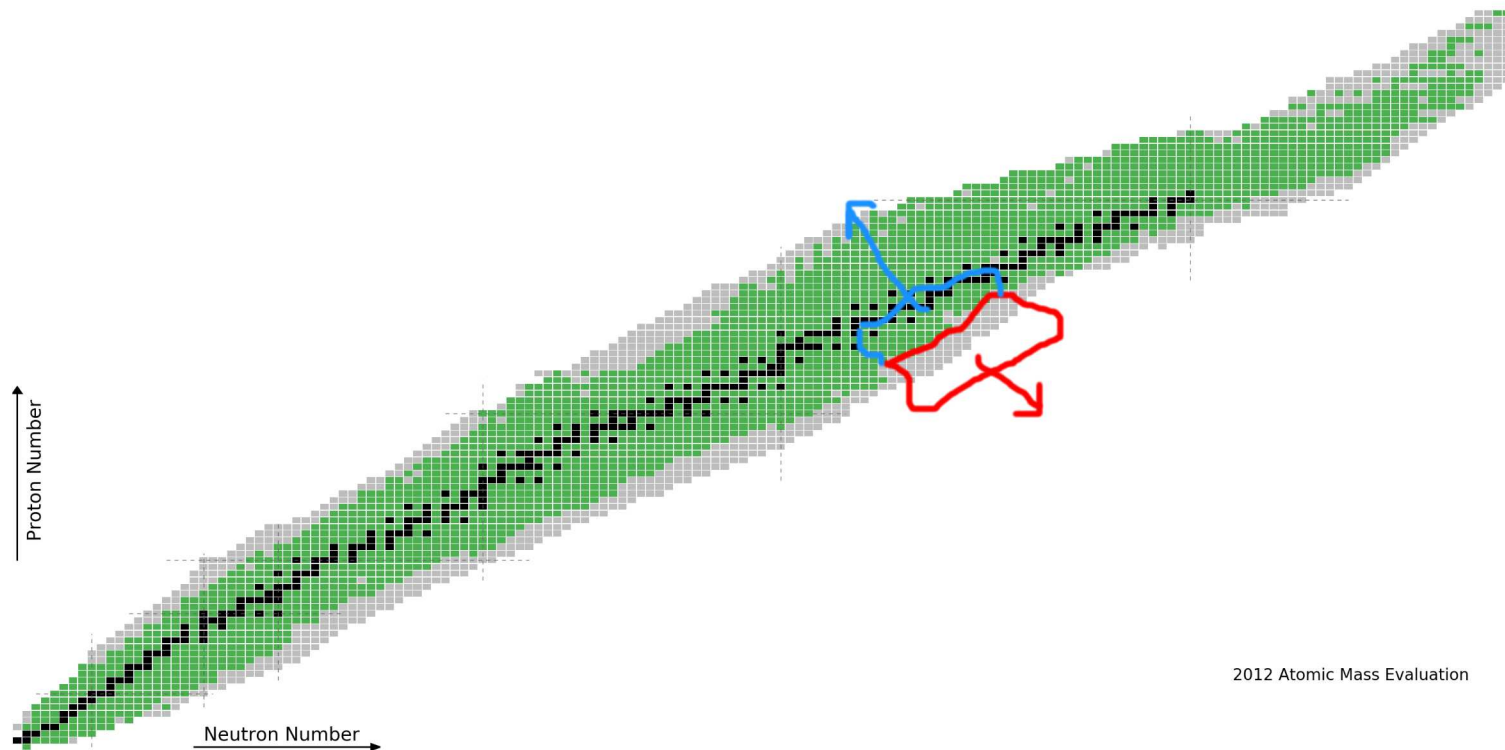
Success?! ... We found a peak! **But there's a problem!**

L ~ **match red abundances** Mumpower et al. J. Phys. G 44 3 034003 (2017)

DIDN'T MATCH KNOWN MASSES

We need to tell the Metropolis algorithm to match both
Update Likelihood function:

$$L \sim \text{match abundances} + \text{match known masses}$$



Mumpower *et al.* J. Phys. G 44 3 034003 (2017)

RESULTS WITH DZ ALONE

No combination of DZ parameters can **simultaneously** reproduce the rare earth peak and match the known masses at the same time

The **nuclear structure** information responsible for the rare earth peak is *missing from the model*

We could move to a nuclear model, but these are more complicated to analyze, with many coupled parameters.

The benefit to DZ is that the abundances are flat to start with.

Let's try to add the missing physics!

Mumpower et al J. Phys. G 44 3 034003 (2017)

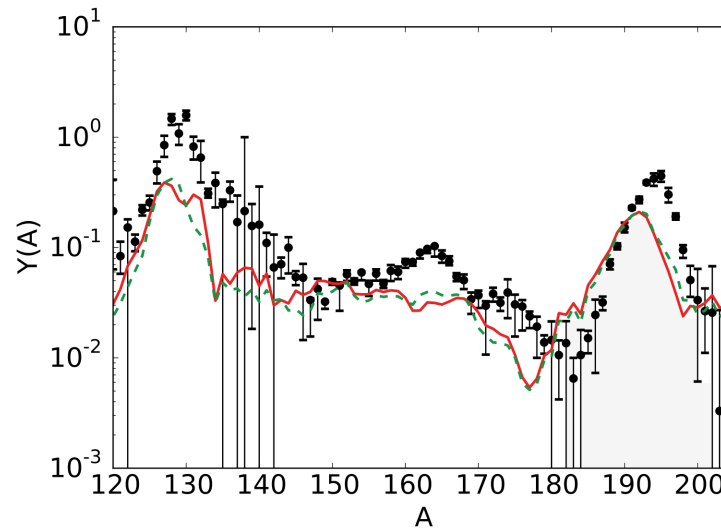
PARAMETERIZE MISSING PHYSICS

$$M(Z, N) = M_{DZ}(Z, N) + a_N \exp[-(Z-C)^2/2f]$$

a_N - Strength of change for given neutron number in MeV

C - Center of the distribution in proton number

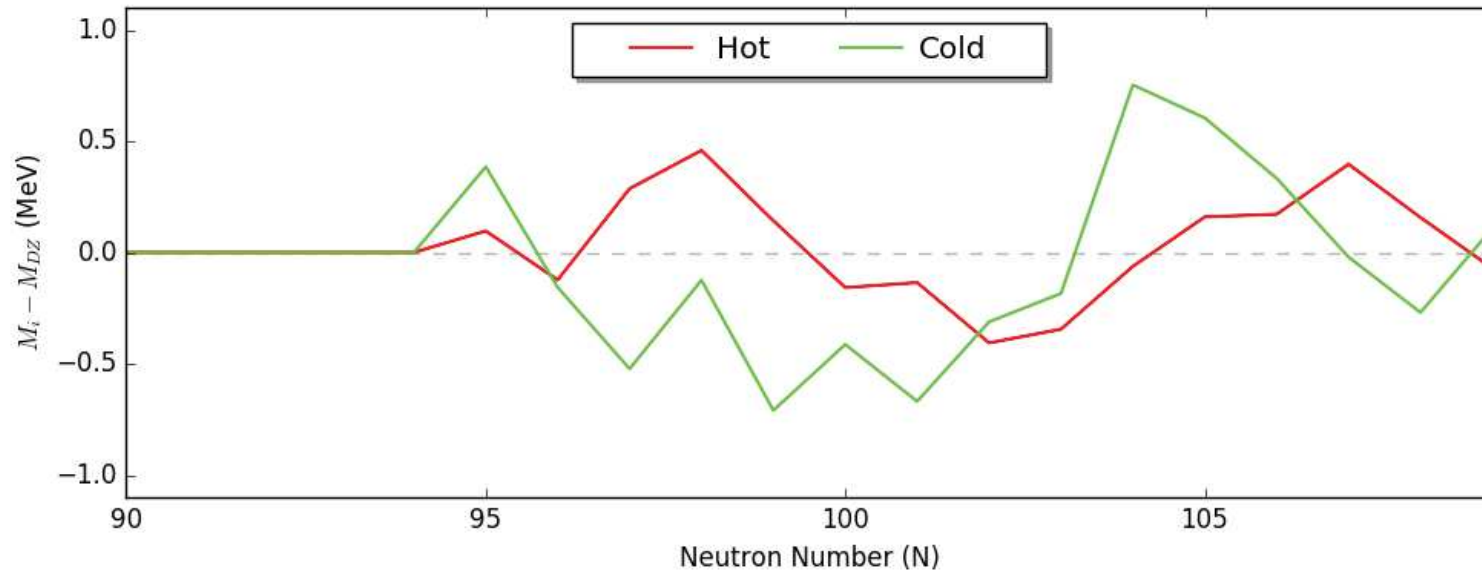
f - Rate of fall off back to stability



Now we repeat the Monte Carlo calculations, letting these parameters vary Mumpower et al. ApJ 833, 282 (2016)

RESULTS WITH NEW PARAMETERS

The predicted masses for $Z = 60$ (Nd)



Distinguishable predictions given different astrophysical conditions

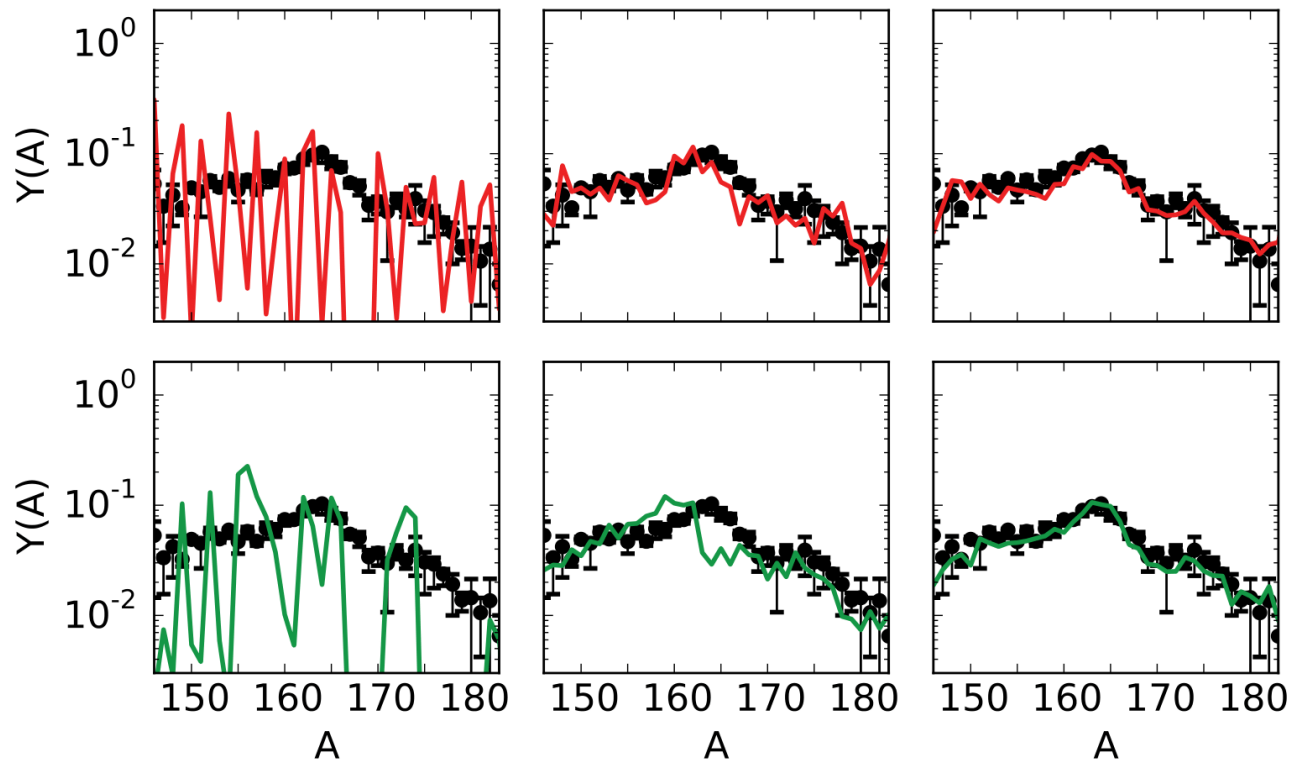
Hot: local min even-N • wider in N • smaller change to masses

Cold: local min odd-N • tighter in N • larger change to masses

Mumpower et al. ApJ 833, 282 (2016)

EVOLUTION OF ABUNDANCES

Before • During • After peak formation

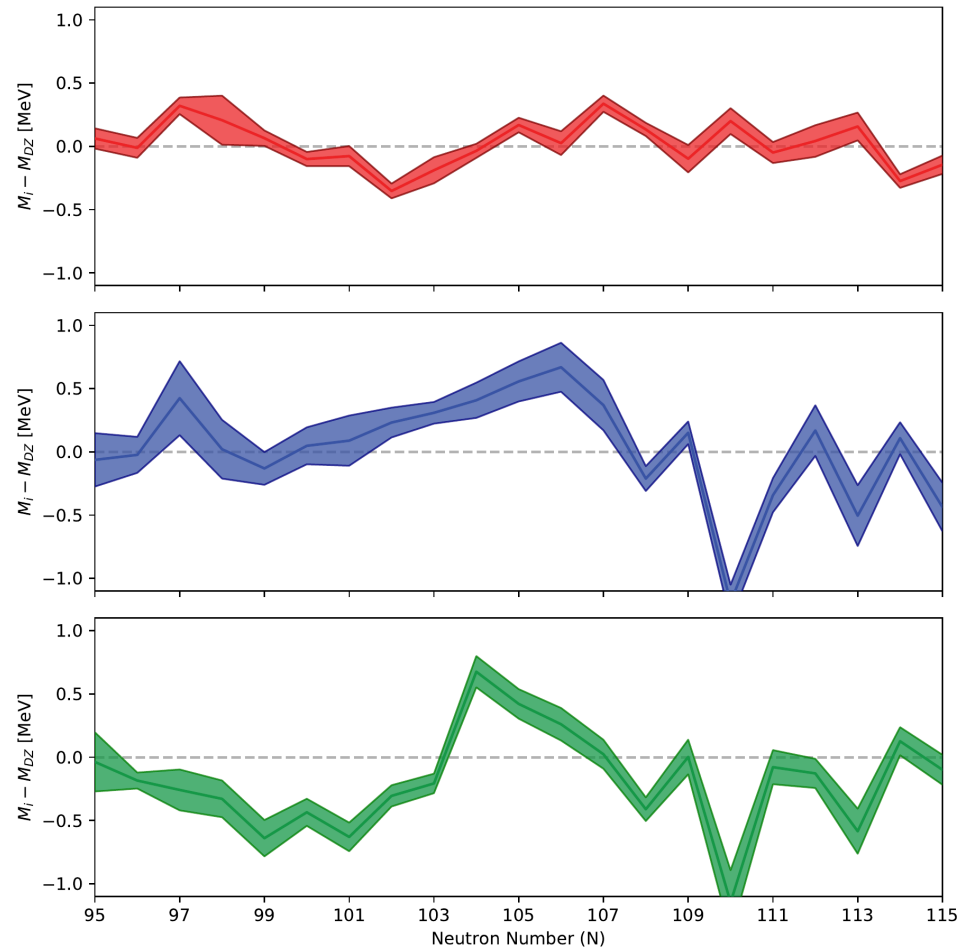


Great success!

Difference is encoded in the astrophysical conditions Mumpower et al. ApJ 833, 282 (2016)

PREDICTED MASSES

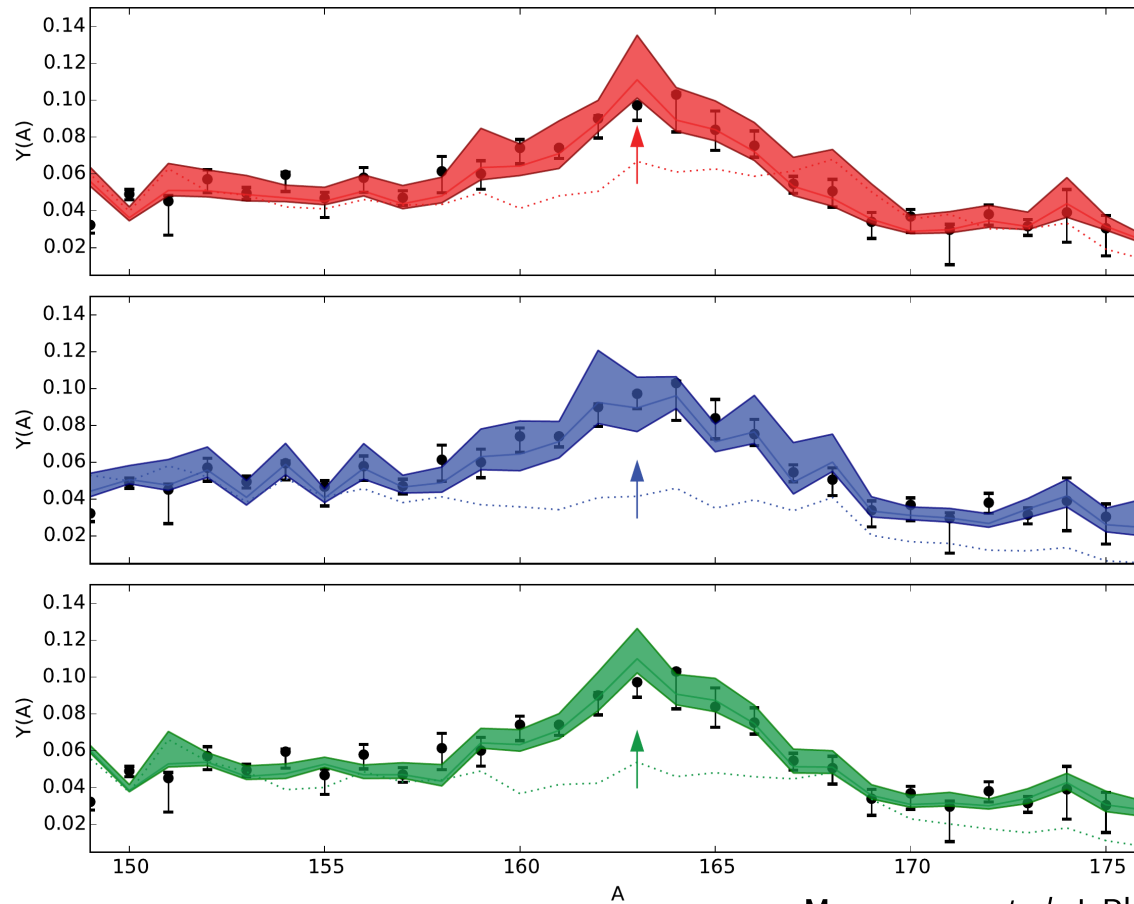
For three astrophysical evolutions: hot, cold or merger



The trend in the masses is important for forming the REP Mumpower et al., J. Phys. G 44:3 034003 (2017)

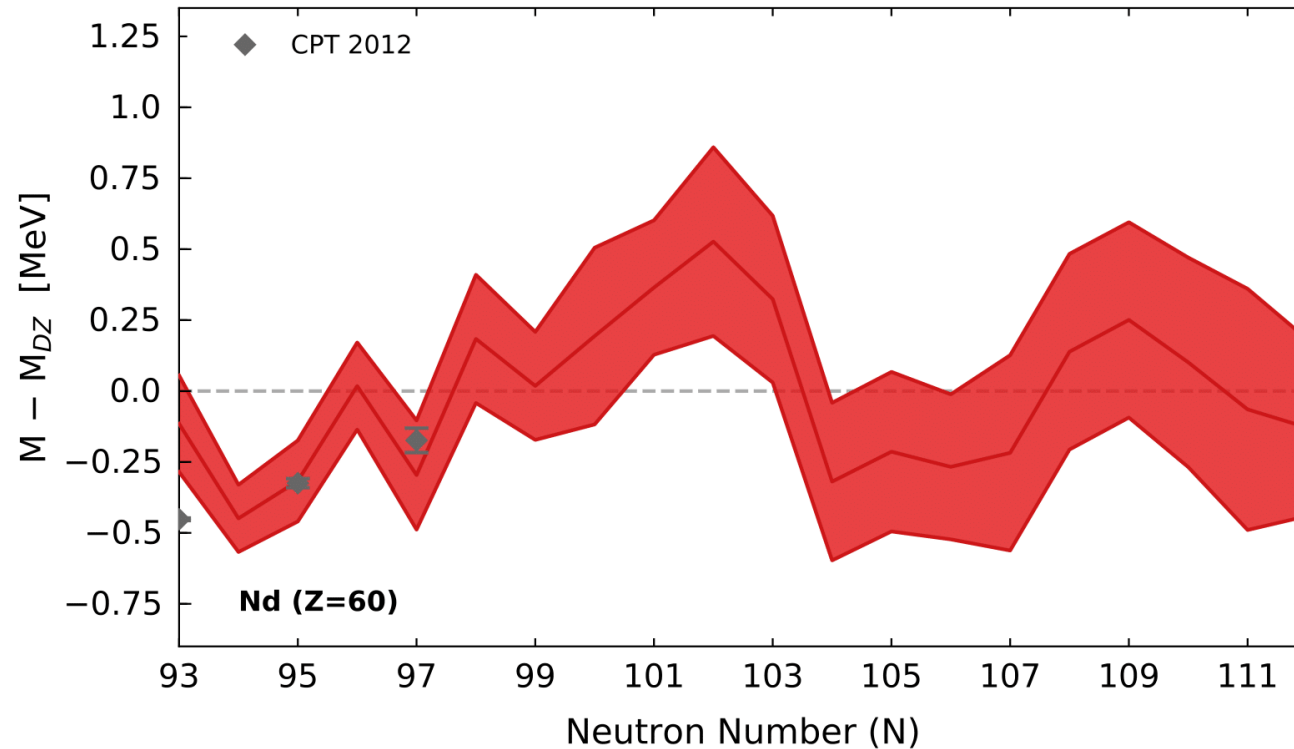
FIT TO ABUNDANCES

For three astrophysical evolutions: **hot**, **cold** or **merger**



Mumpower et al. J. Phys. G 44 3 034003 (2017)

LATEST RESULTS

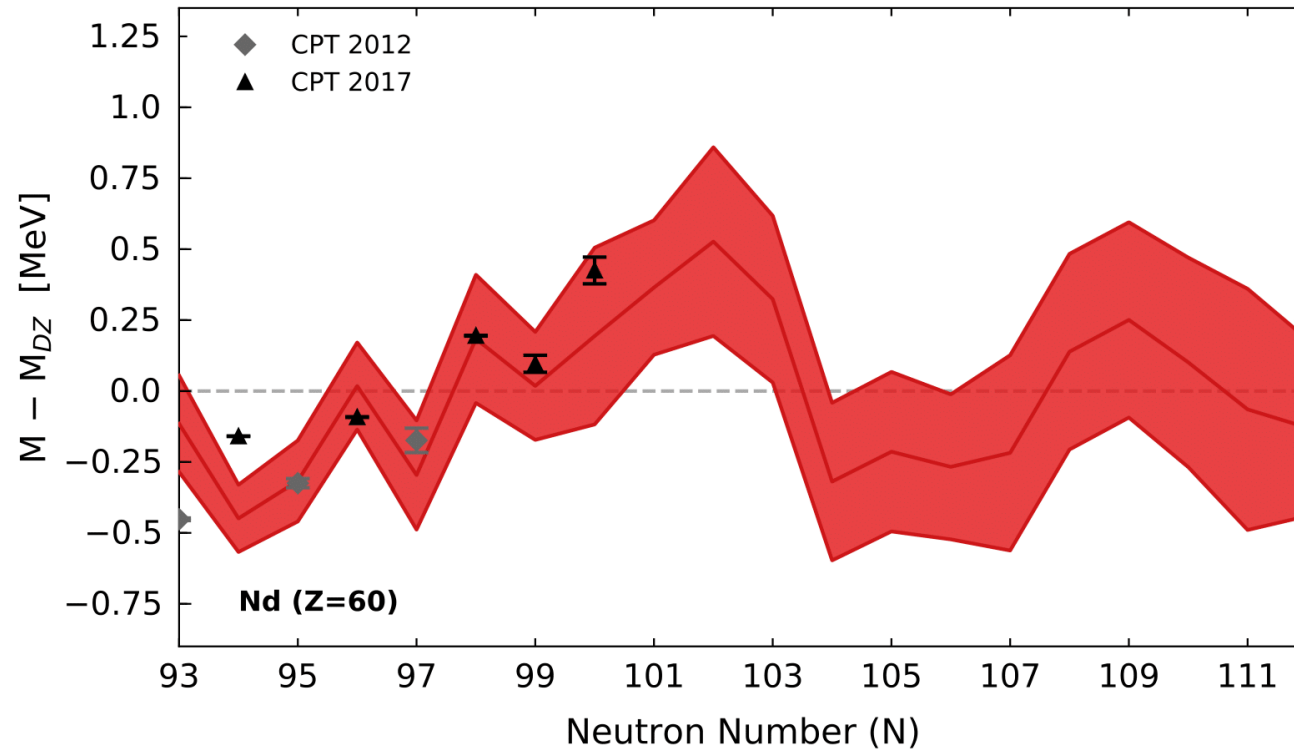


Nicole Vassh has made major upgrades to the algorithm

We now use solar data & uncertainties and also parallel chains Monte Carlo

Vassh *et al.*, in prep (2018), O'ford *et al.*, PRL 120, 262702 (2018)

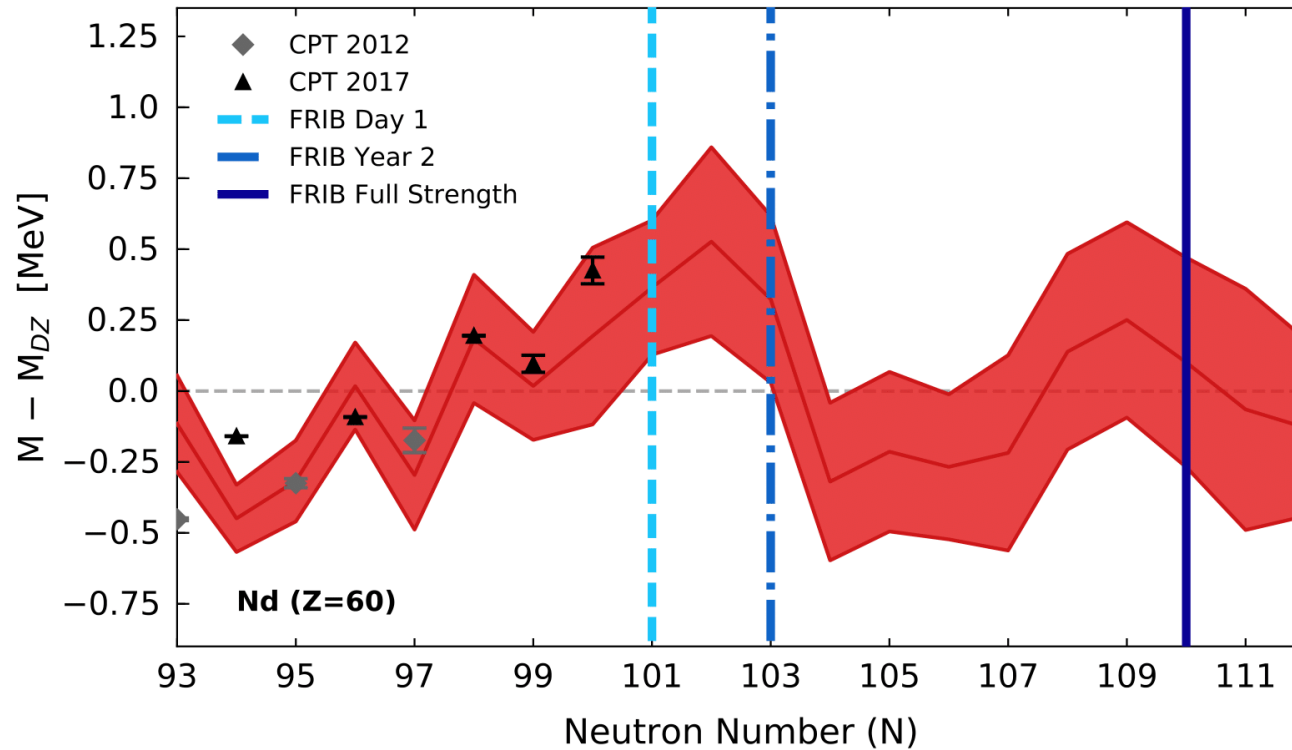
LATEST RESULTS



The algorithm is now ready to be used in matching data

We compared to CARIBU results & Yasli et al. in prep (2018), Orford et al. PRL 120,262702 (2018) independently match this new data

LATEST RESULTS



FRIB will be able to probe most of relevant regions for REP formation

Only the upgrade will allow the reach out to the potential N = 110 structure

Vassh *et al.* in prep (2018), Orford *et al.* PRL 120,262702 (2018)

REVERSE ENGINEERING MASSES

What are the consequences of the future measurements?

Either we **find the structure**... or **we don't**

If we **do**: we favor precise conditions for the main r-process

If we **don't**: we favor extreme conditions that REQUIRE fission recycling... only option is the **tidal tails of mergers**

Perhaps nature is more complicated than we think... and we learn something even more profound

Make the measurements to find out! Munro et al. J. Phys. G 44 3 034003 (2017)

SUMMARY

Theoretical calculations for the r-process are currently **data starved**

Current models result in **orders of magnitude uncertainty** in calculated abundances

FRIB will provide a plethora of **new** nuclear data

This will be invaluable in **benchmarking** theoretical models

(in particular, for **nuclear masses**)

As well as constraining the astrophysical conditions of the r-process

Results at MatthewMumpower.com