BSM Theory: The SM and Beyond

UCHITIAR

Patrick "PaDDy" Fox

Fermilab

Bugs and Features

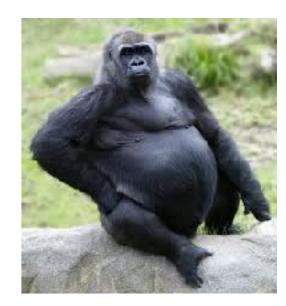
 $\mathcal{L}_2 = \pm \mu^2 |H|^2$

Hierarchy (Naturalness) problem

Why is μ so much smaller than M_{GUT} , M_{Pl} ?

Unlike fermions (and gauge bosons) no symmetry protects scalar mass parameter
1.Nature is fine-tuned (anthropics?)
2.The SM has no high scales (gravity?, unification?)
3.New dynamics/symmetries keeps mass scale low





Bugs and Features

Hierarchy (Naturalness) problem

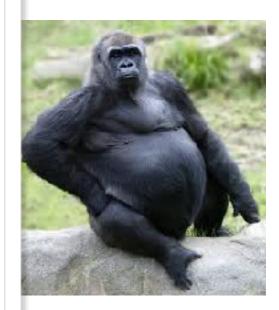


$$\mathcal{L}_2 = \pm \mu^2 |I|$$

Why is μ sc

There is something fascinating about science. One gets such wholesale returns of conjecture out of such a trifling investment of fact.

—Mark Twain



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Hierarchy (Naturalness) problem



Scalars are sensitive to the highest scale in the theory!

Expect new physics (A) at
$$rac{4\pi}{g}m_h$$

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Hierarchy (Naturalness) problem

SM Higgs sensitivity (how low can you go)



$$\delta m_h^2 = \alpha_t \Lambda_t^2 + \alpha_g \Lambda_g^2 + \alpha_h \Lambda_h^2$$

$$\alpha_t = \frac{3m_t^2}{4\pi^2 v^2}, \quad \alpha_g = -\frac{6m_W^2 + 3m_Z^2}{16\pi^2 v^2}, \quad \alpha_h = -\frac{3m_h^2}{16\pi^2 v^2}$$

(One) Measure of fine tuning: $D_i(m_h) \equiv \left| \frac{\partial \log m_h^2}{\partial \log \Lambda_i^2} \right| = \frac{|\alpha_i|\Lambda_i^2}{m_h^2}$

No guaranteed discovery, unlike Higgs mechanism Should not stop us looking!! SM Higgs sensitivity (how low does Nature Joes Nature Joes

$$\delta m_h^2 = \alpha_t \Lambda_t^2 + \alpha_g \Lambda_g^2 + \alpha_h \Lambda_h^2$$

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Supersymmetry...a BSM case study Bosons Fermions

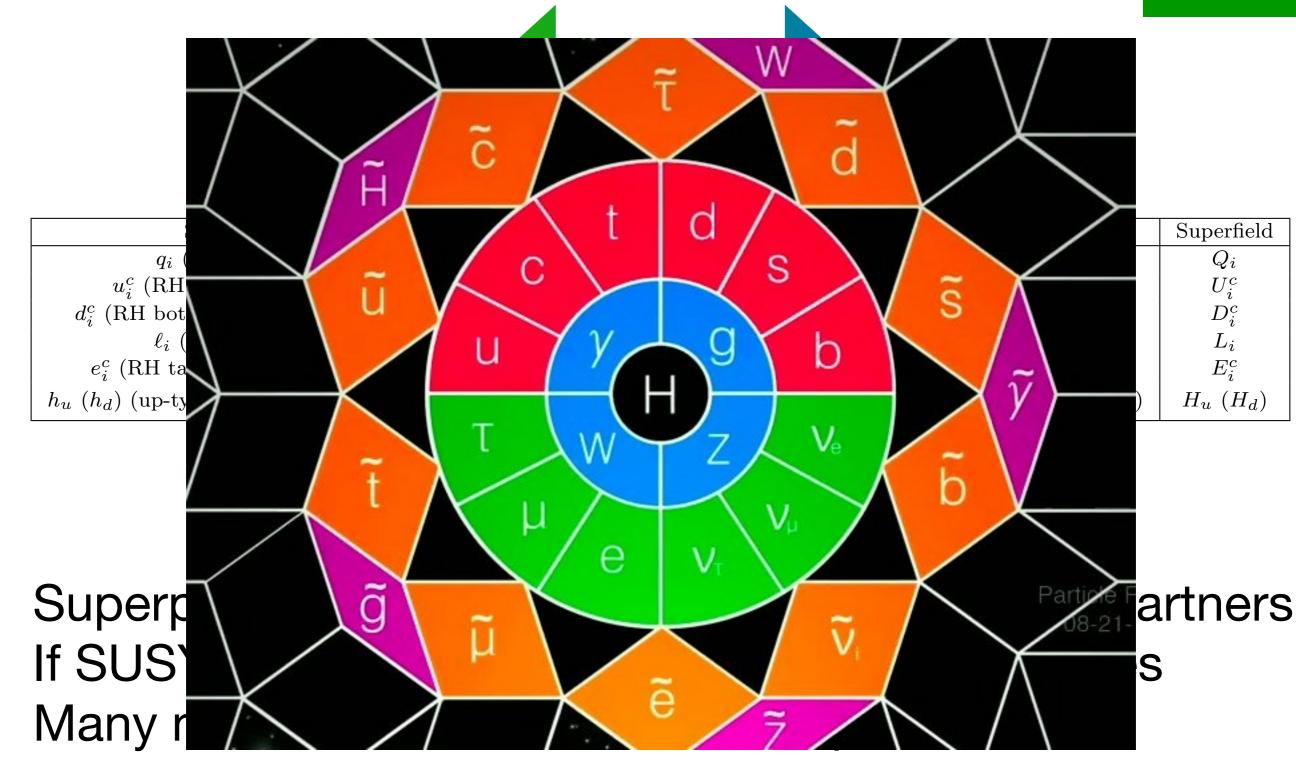
(more than) Doubling of the spectrum

SM Field	SU(3)	SU(2)	U(1)	MSSM partner	Superfield
q_i (LH quarks)	3	2	$\frac{1}{6}$	\tilde{q}_i (LH squarks)	Q_i
u_i^c (RH top, charm, up)	3	1	$\frac{\frac{1}{6}}{-\frac{2}{3}}$	\tilde{u}_i^c (RH stop, scharm, sup)	U_i^c
d_i^c (RH bottom, strange, down)	3	1	$\frac{1}{3}$	\tilde{d}_i^c (RH sbottom, sstrange, sdown)	D_i^c
ℓ_i (LH leptons)	1	2	$-\frac{1}{2}$	$\tilde{\ell}_i$ (LH sleptons)	L_i
e_i^c (RH tau, muon, electron)	1	1	1	\tilde{e}_i^c (RH stau, smuon, selectron)	E_i^c
h_u (h_d) (up-type (down-type) Higgs)	1	2	$\frac{1}{2}\left(-\frac{1}{2}\right)$	$\tilde{h}_u \left(\tilde{h}_d \right)$ (up-type (down-type) higgsino)	H_u (H_d)
gluino	8	1	0	gluino	
W/Z	1	3	0	Wino/Zino	
B/photon	1	1	0	bino/photino	

Superpartners have the same couplings as SM partners If SUSY (softly) broken they have different masses Many new interactions...>100 new parameters!

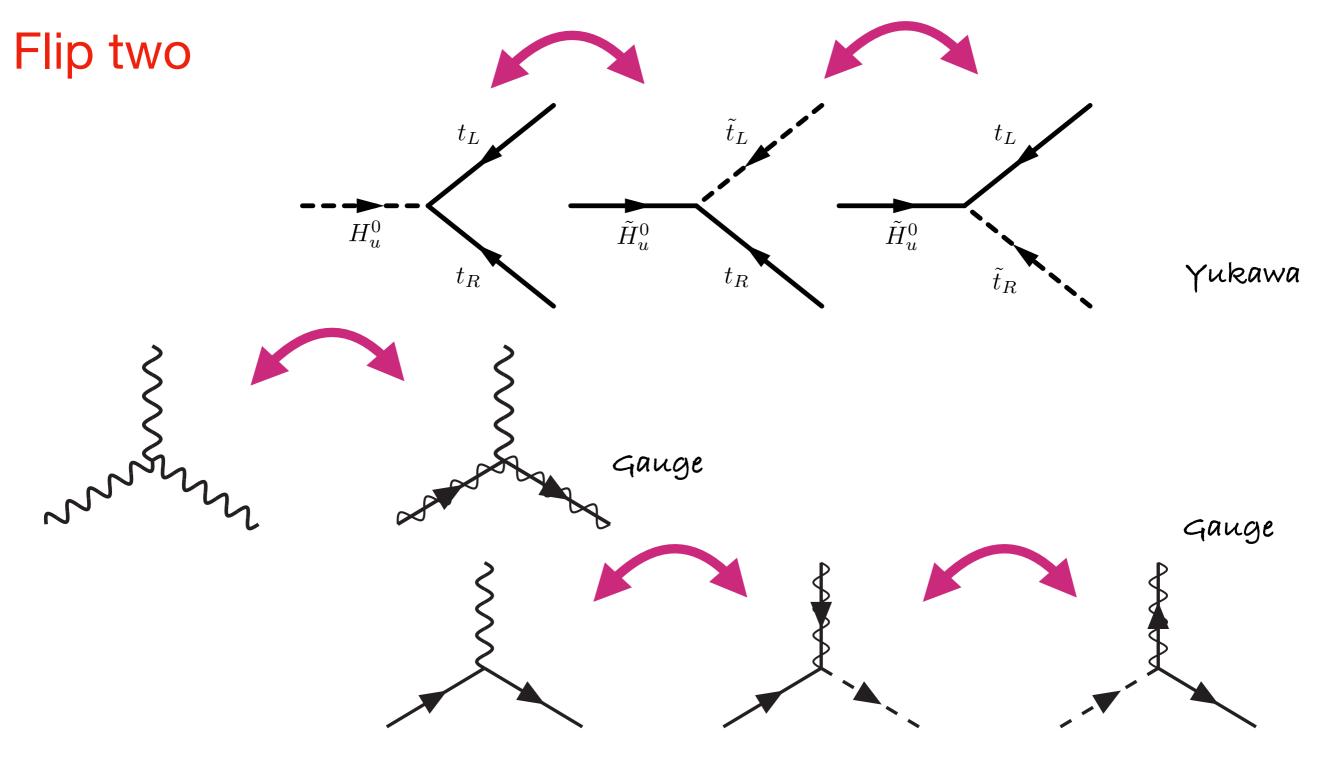
> Many constrained by flavour, CP-violation SUSY breaking models (GMSB, AMSB,...) predict relations

Supersymmetry...a BSM case study



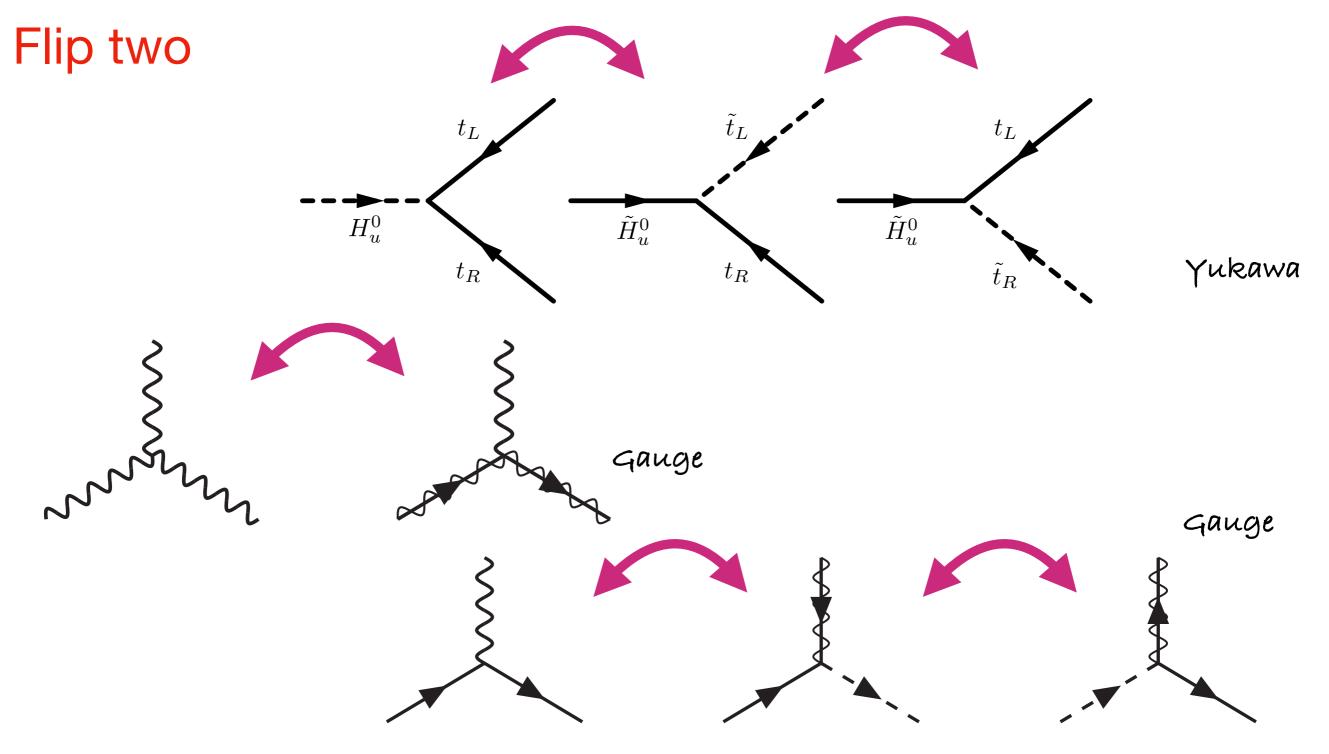
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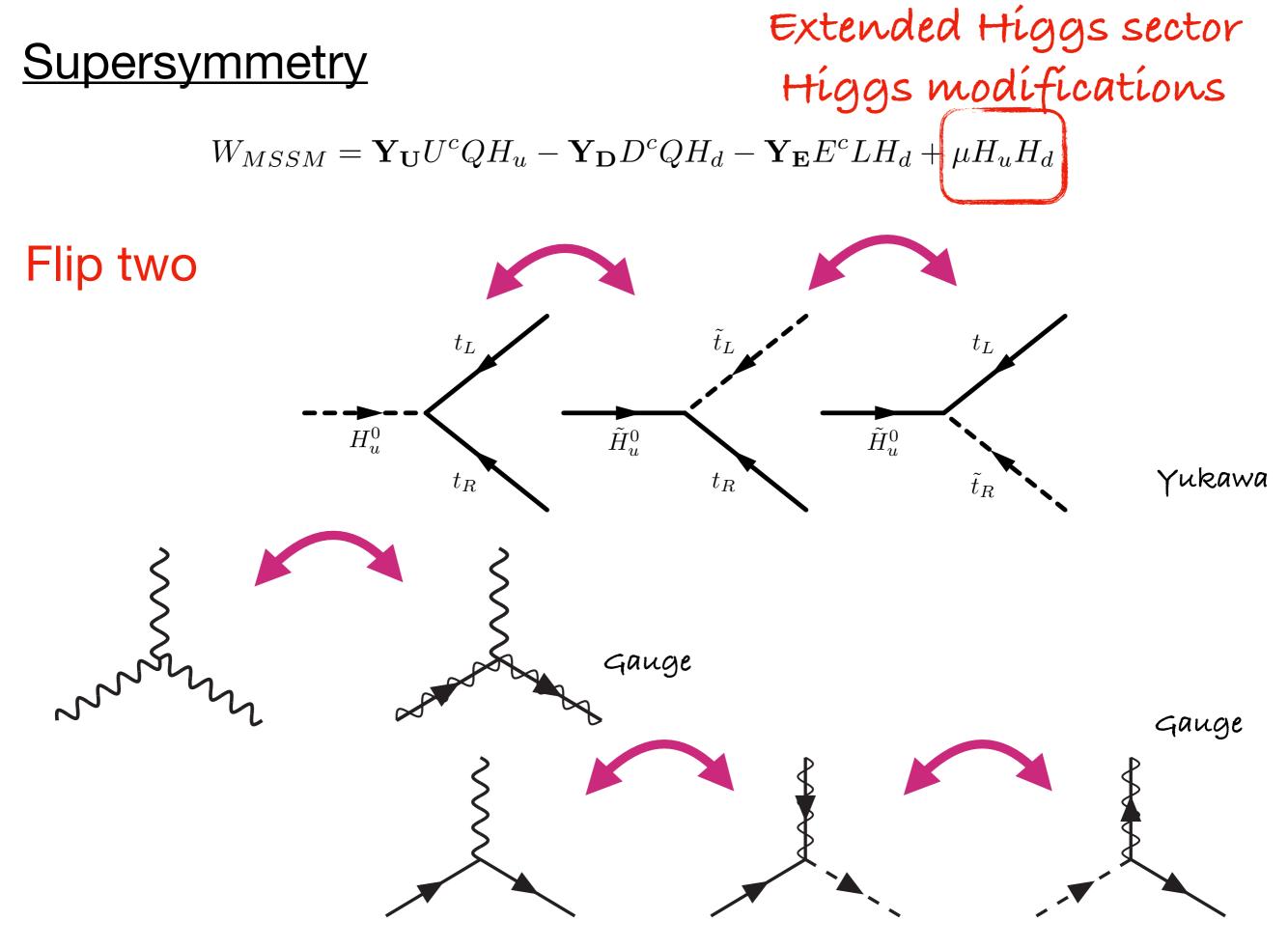
 $W_{MSSM} = \mathbf{Y}_{\mathbf{U}} U^{c} Q H_{u} - \mathbf{Y}_{\mathbf{D}} D^{c} Q H_{d} - \mathbf{Y}_{\mathbf{E}} E^{c} L H_{d} + \mu H_{u} H_{d}$



Extended Híggs sector Híggs modífications

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 t_R

New fields allow for new interactions (EFT philosophy)

 t_R

 \tilde{t}_R

POF

$$\Delta L = 1 \qquad \Delta B = 1 \qquad \Delta B = 1 \qquad W_{\Delta B,L} = \kappa_1^{ijk} Q_i L_j D_k^c + \kappa_2^{ijk} L_i L_j E_k^c + \kappa_3^i L^j H_u + \kappa_4^{ijk} D_i^c D_j^c U_k^c$$

Proton lifetime: $\Gamma \sim \frac{\kappa_1 \kappa_4}{16\pi} \frac{m_p^{\circ}}{m_{\tilde{q}}^4} \qquad \kappa < 10^{-12}!$

Forbid these (RPV) operators with a parity (R-parity)

 $SM \to SM$ $BSM \to -BSM$

1.SM and partners don't mix2.SUSY states pair produced3.Lightest parity odd particle stable (DM?)

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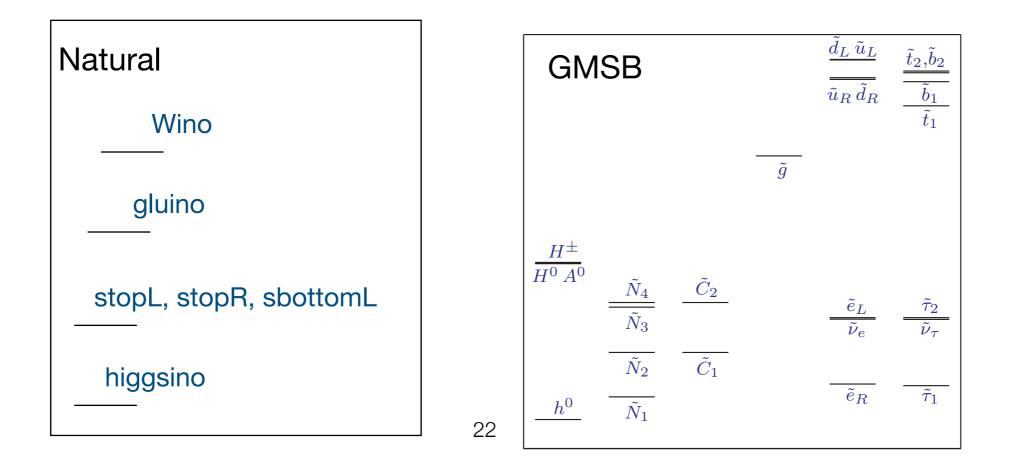
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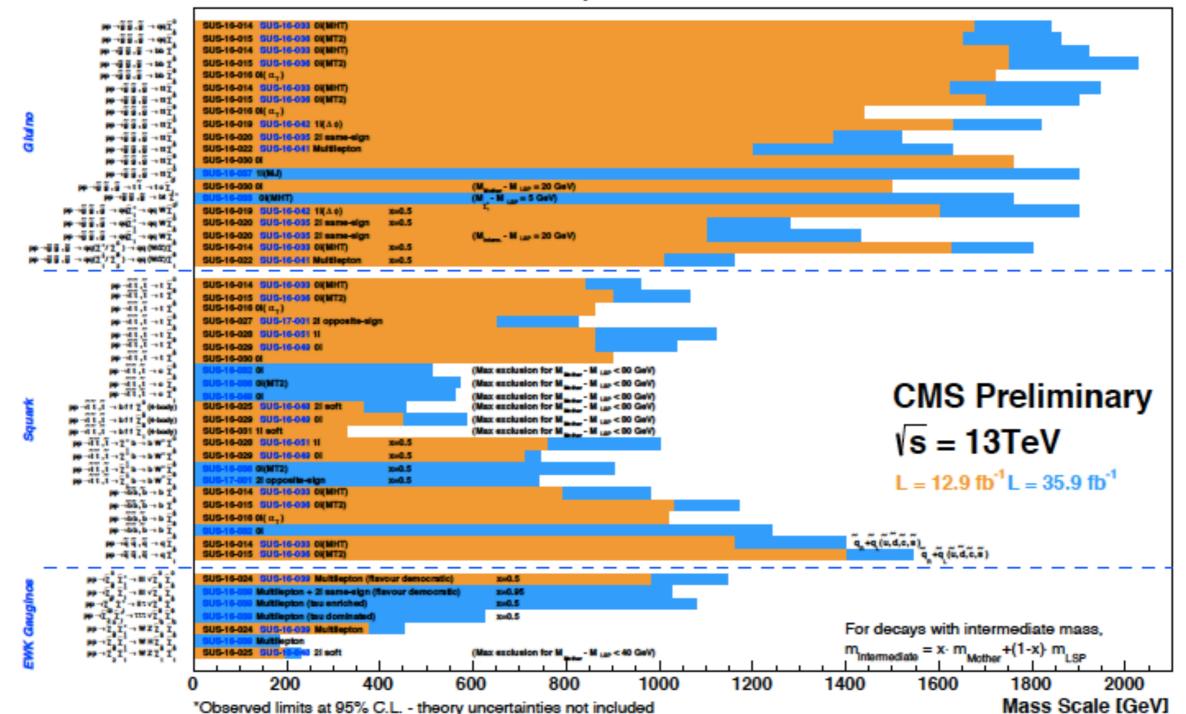
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Complicated spectrum, details depend on model GMSB, Effective/Natural SUSY, Dirac gauginos.... Many, many interesting collider signatures Lightest coloured states made first, decays involve MET Compressed/Stealth spectra can hide SUSY (a little) Electroweakino sector starting to be probed (DM) SUSY is a great signal generator

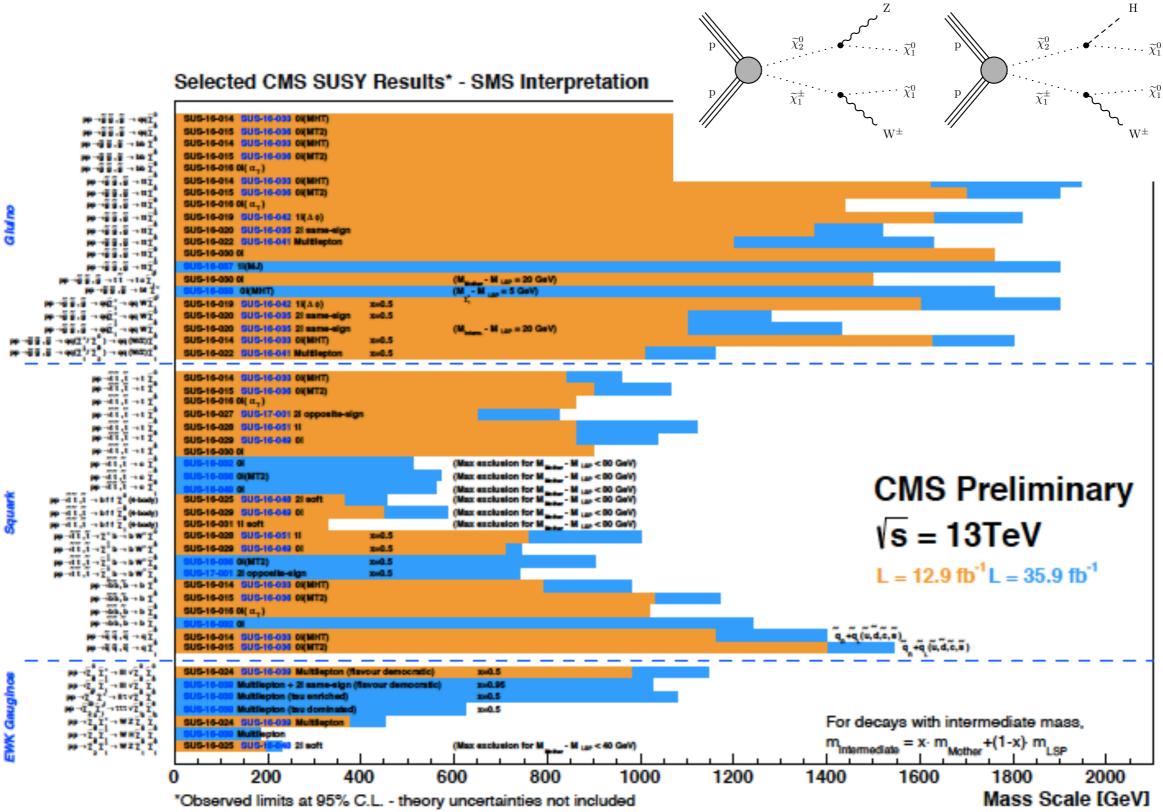




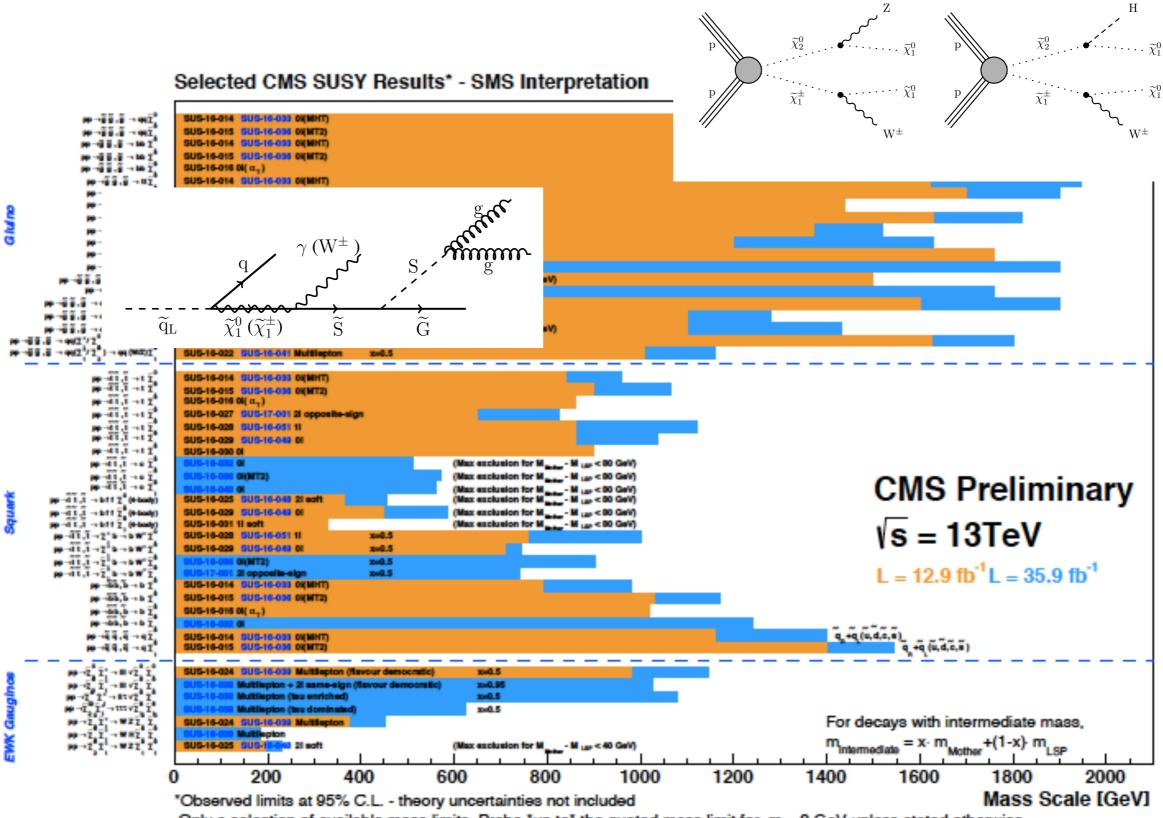
Selected CMS SUSY Results* - SMS Interpretation

Only a selection of available mass limits. Probe *up to* the quoted mass limit for me =0 GeV unless stated otherwise

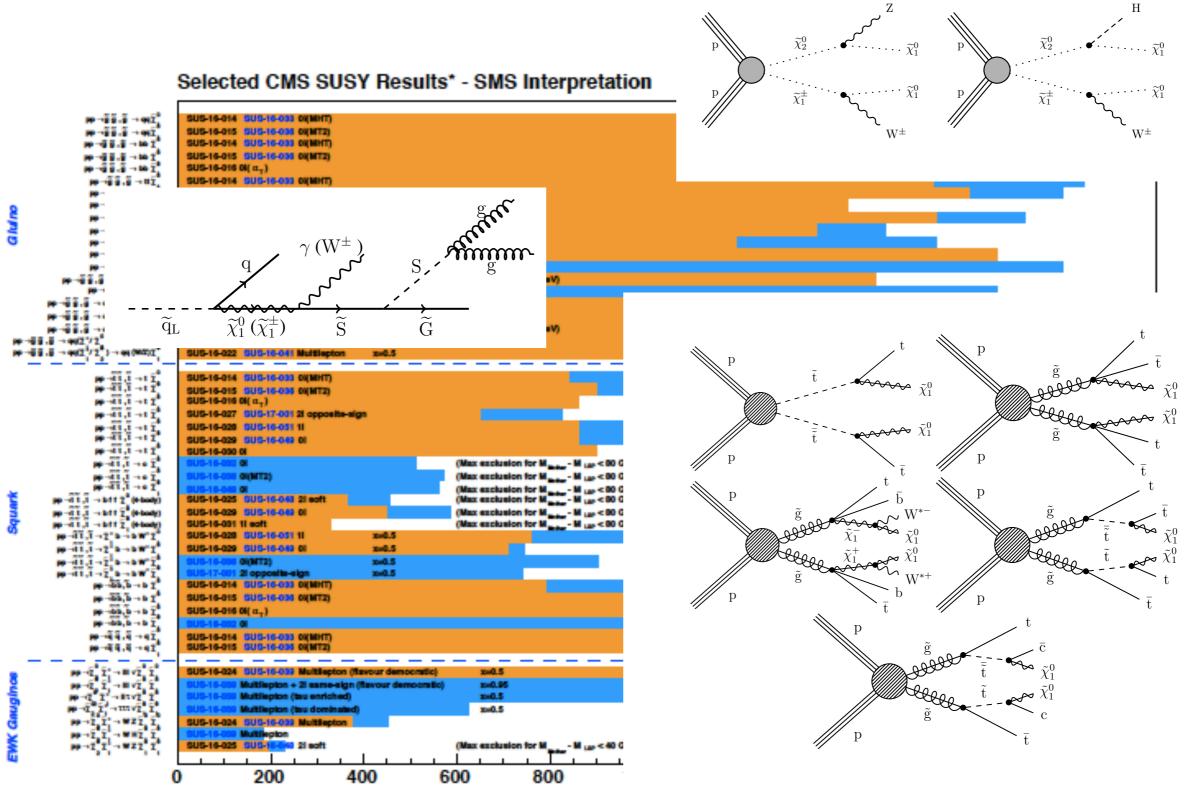
ICHEP '16 - Moriond '17



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*Observed limits at 95% C.L. - theory uncertainties not included Mass Only a selection of available mass limits. Probe *up to* the quoted mass limit for m =0 GeV unless stated otherwise

Mass Scale [GeV]

Top partners (fermions/bosons) Higgs sector modifications LPOPs, parity, DM, MET Extra matter in fundamental and adjoint reps.

New gauge groups? Lighter, more weakly coupled particles? Resonances?





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Top partners (fermions/bosons) **Higgs sector modifications** LPOPs, parity, DM, MET Extra matter in fundamental and adj Randall Sundrum

New gauge groups? Lighter, more weakly coupled particles? Little Higgs **Resonances?**







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[Chacko, Goh, Harnik]

 $\mathrm{SM}_A \times \mathrm{SM}_B \times \mathbb{Z}_2$

$\mathcal{L} \supset y \, Q_A H_A U_A^c + y \, Q_B H_B U_B^c$

Higgs is a PNGB, and Higgs potential is O(8) symmetric

$$V = -m^2 H^{\dagger} H + \lambda \left(H^{\dagger} H \right)^2$$

 $O(8) \rightarrow O(7)$:7 Goldstone bosons, 3 eaten by B gauge bosons

$$H = \begin{pmatrix} H_A \\ H_B \end{pmatrix} = e^{ih^a t^a / f} \begin{pmatrix} 0 \\ 0 \\ 0 \\ f \end{pmatrix}$$



[Chacko, Goh, Harnik]

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\text{SM Higgs} \\
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\end{array}$$

Twin Higgs

Gauge and Yukawa interactions explicitly break the O(8)

Mass for Higgs?

$$\frac{3}{8\pi^2}\Lambda^2 \left(y_A^2 H_A^{\dagger} H_A + y_B^2 H_B^{\dagger} H_B \right)$$

$$Z_2 \Rightarrow y_A = y_B$$

$$\frac{3}{8\pi^2}\Lambda^2 y^2 H^{\dagger} H$$

Loop corrections to mass is O(8) symmetric, does not lead to quadratically divergent Higgs mass

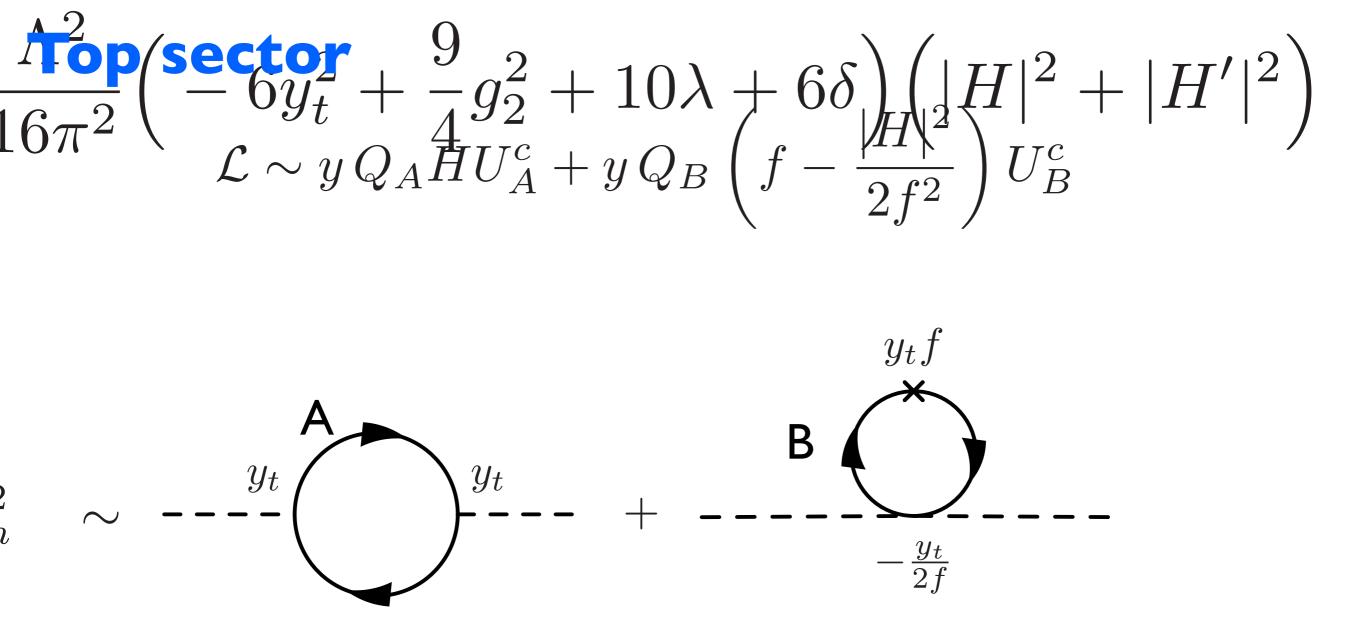
EFT aside

Low energy degrees of freedom, non-linearly realized symm.

$$H = \begin{pmatrix} H_A \\ H_B \end{pmatrix} = e^{ih^a t^a / f} \begin{pmatrix} 0 \\ 0 \\ 0 \\ f \end{pmatrix}$$

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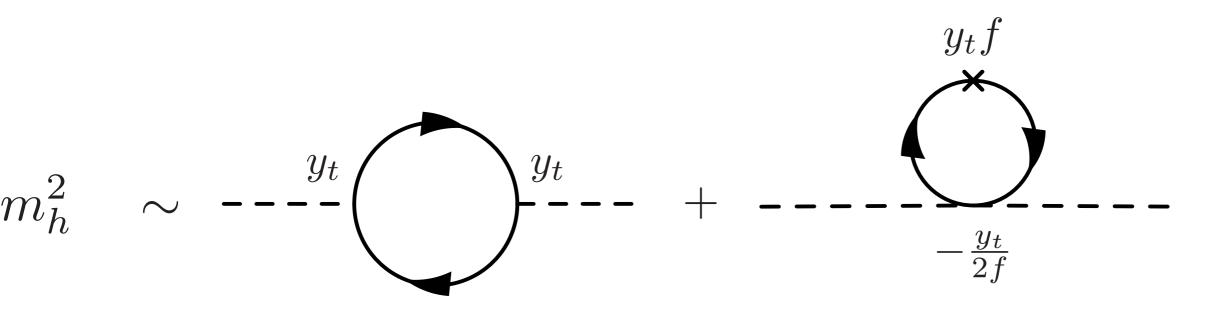
$$H = \begin{pmatrix} \mathbf{h} \frac{if}{\sqrt{\mathbf{h}^{\dagger}\mathbf{h}}} \sin \frac{\sqrt{\mathbf{h}^{\dagger}\mathbf{h}}}{f} \\ 0 \\ f \cos \frac{\sqrt{\mathbf{h}^{\dagger}\mathbf{h}}}{f} \end{pmatrix} = \begin{pmatrix} i\mathbf{h} \\ 0 \\ f - \frac{1}{2f}\sqrt{\mathbf{h}^{\dagger}\mathbf{h}} \end{pmatrix} + \dots$$



Quadratic divergences cancel, states running in loop have <u>no</u> SM charge, <u>same</u> spin (3 is just a number)

Cancelling states not coloured: small production x-sec at LHC $= \frac{\Lambda^2}{10 \text{ separate vand f and make the SM Higgs he mostly in A H'|^2 }$

 $\langle H_A \rangle = v_{SM} \ll \langle H_B \rangle = f \sim 1 \,\mathrm{TeV}$



Tuning grows with f/v

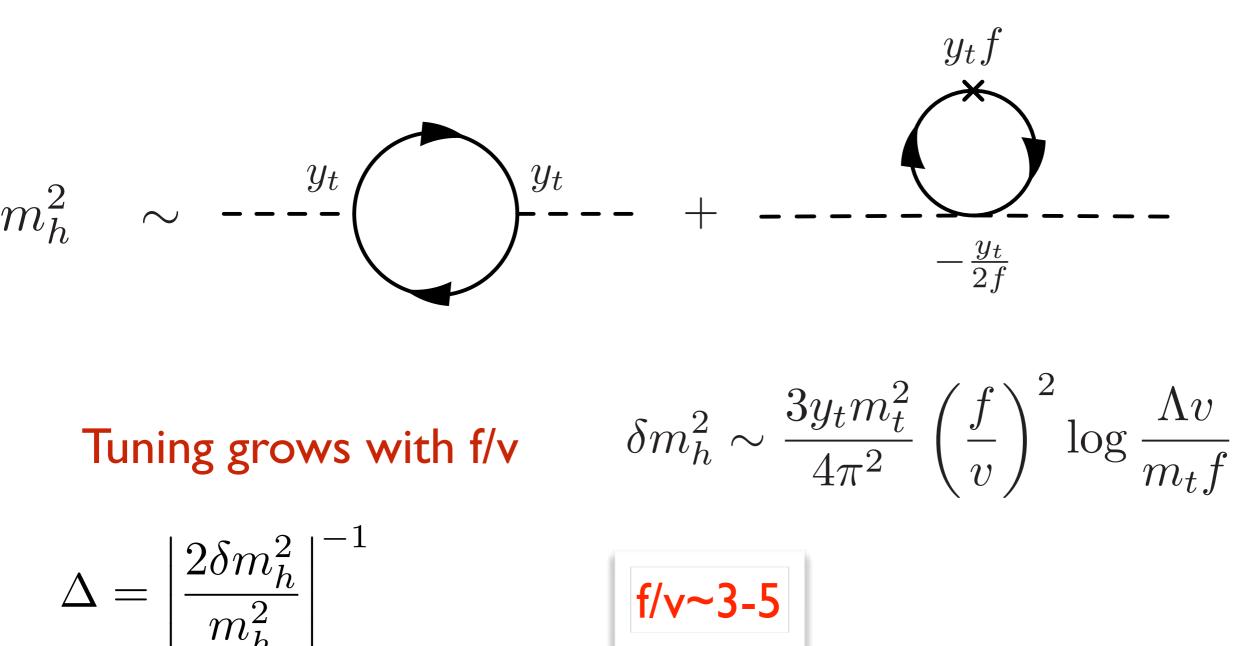
1

$$\delta m_h^2 \sim \frac{3y_t m_t^2}{4\pi^2} \left(\frac{f}{v}\right)^2 \log \frac{\Lambda v}{m_t f}$$

$$\Delta = \left|\frac{2\delta m_h^2}{m_h^2}\right|^-$$

 $= \frac{\Lambda^2}{10 \text{ separate vand f and make the SM Higgs if (mpstly in AH'|^2) in the set of the set$

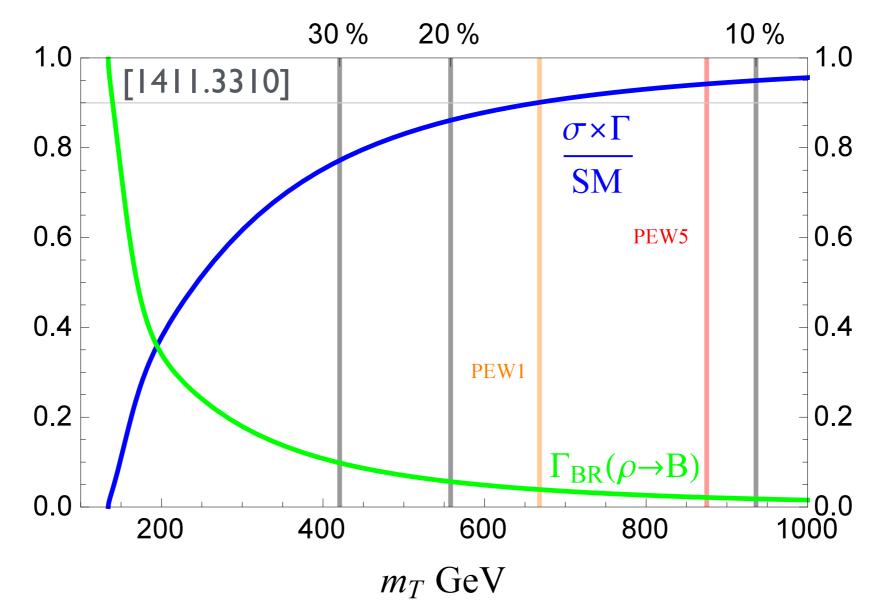
 $\langle H_A \rangle = v_{SM} \ll \langle H_B \rangle = f \sim 1 \,\mathrm{TeV}$



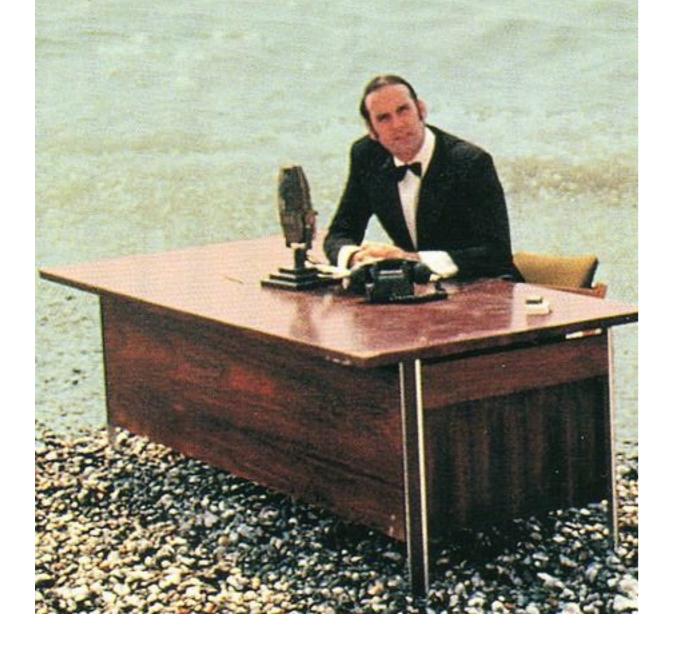
$$V = -m^2 \left(H_A^{\dagger} H_A + H_B^{\dagger} H_B \right) + \lambda \left(H_A^{\dagger} H_A + H_B^{\dagger} H_B \right)^2$$

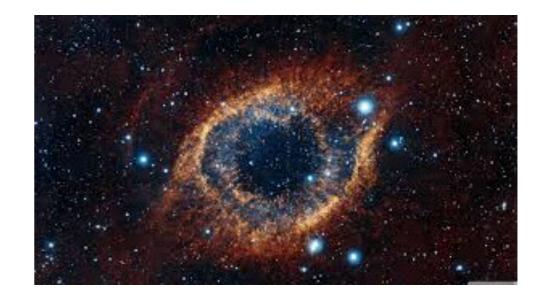
Higgs portal between A and B sectors

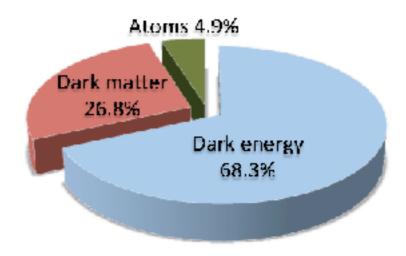
- Higgs mixing and corrections to Higgs pheno at $rac{v^2}{f^2}$
- Higgs invisible decay width tight B sector stuff



And now for something completely different...



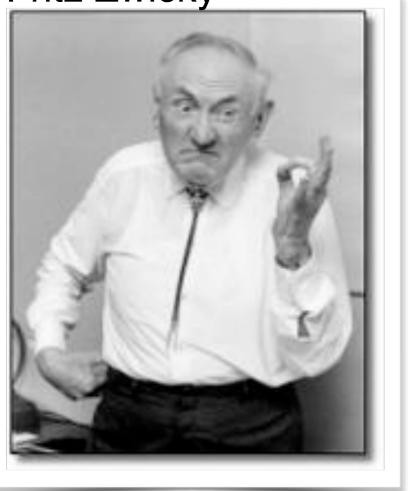






"You spin me right round..."

Fritz Zwicky



Coma Cluster

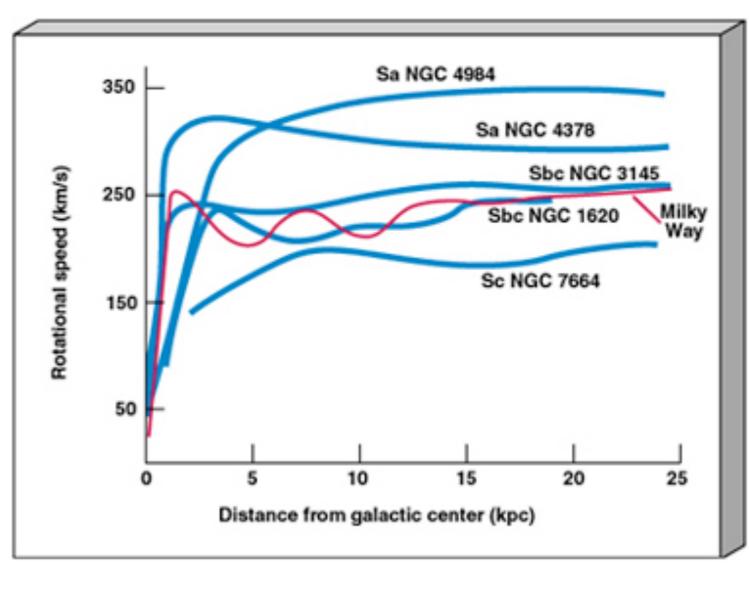


Virial theorem: $2\langle K \rangle = -\langle V \rangle$ $M = \frac{v^2 R}{G_N}$

90% of the matter in the cluster doesn't shine

Vera Rubin



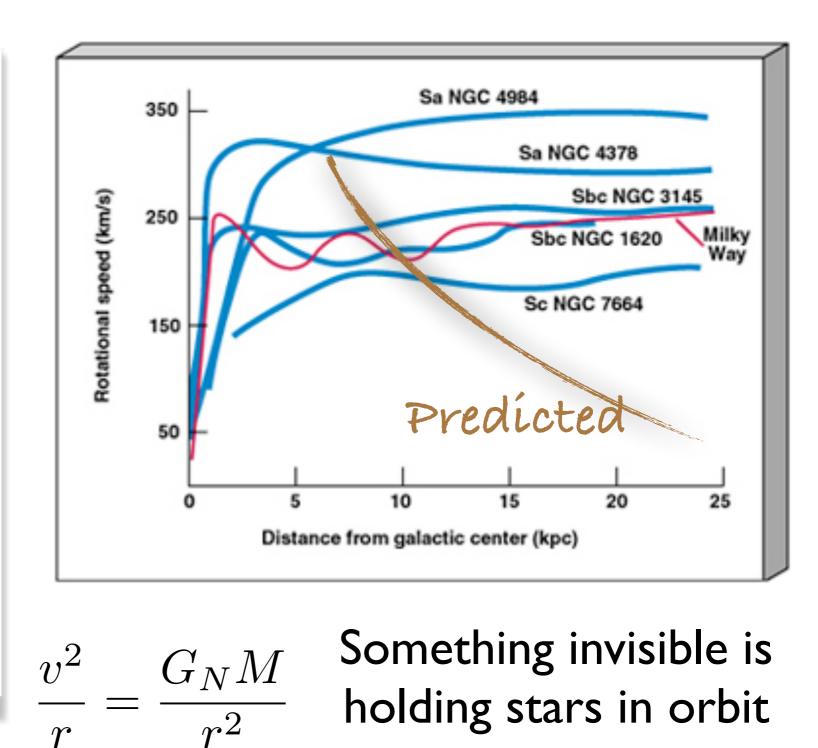


Something invisible is holding stars in orbit

Has been repeated in many systems on many scales. Alway same result: never enough stuff

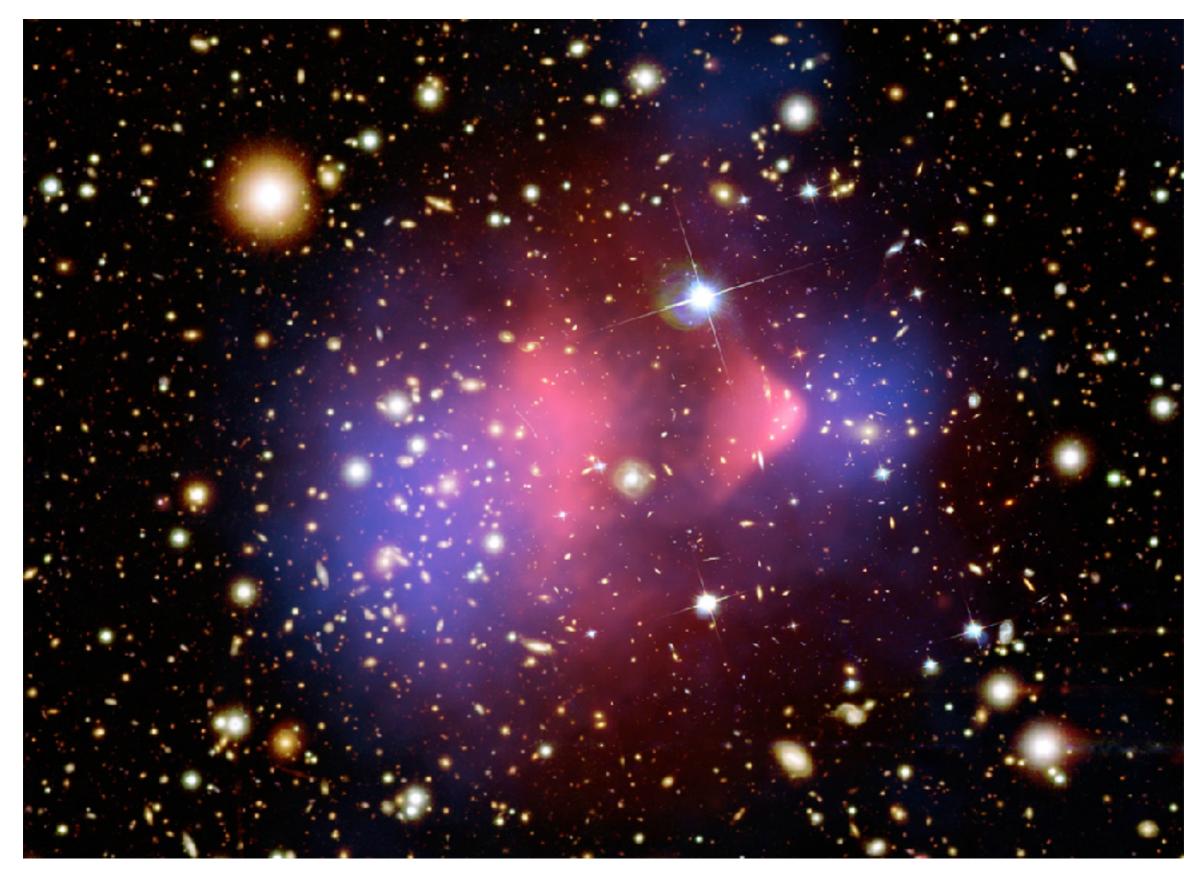
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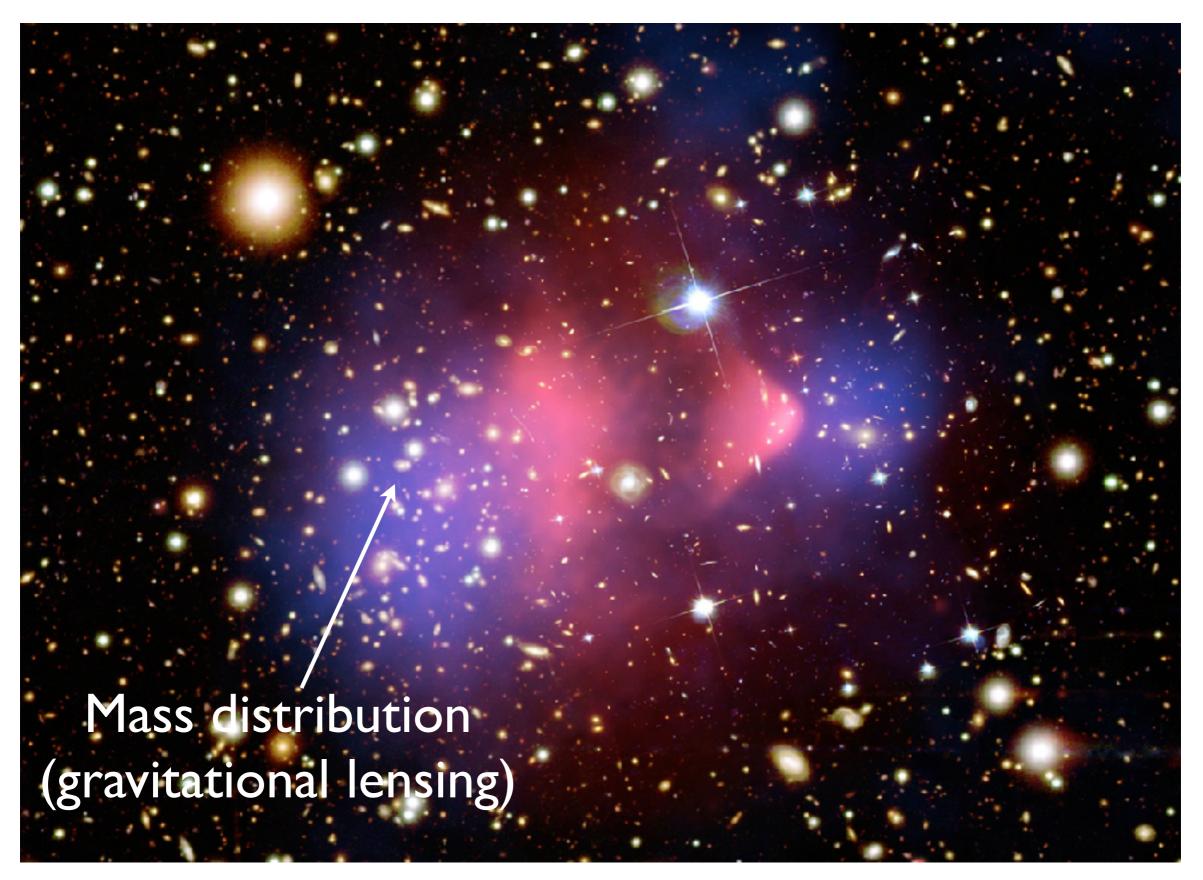


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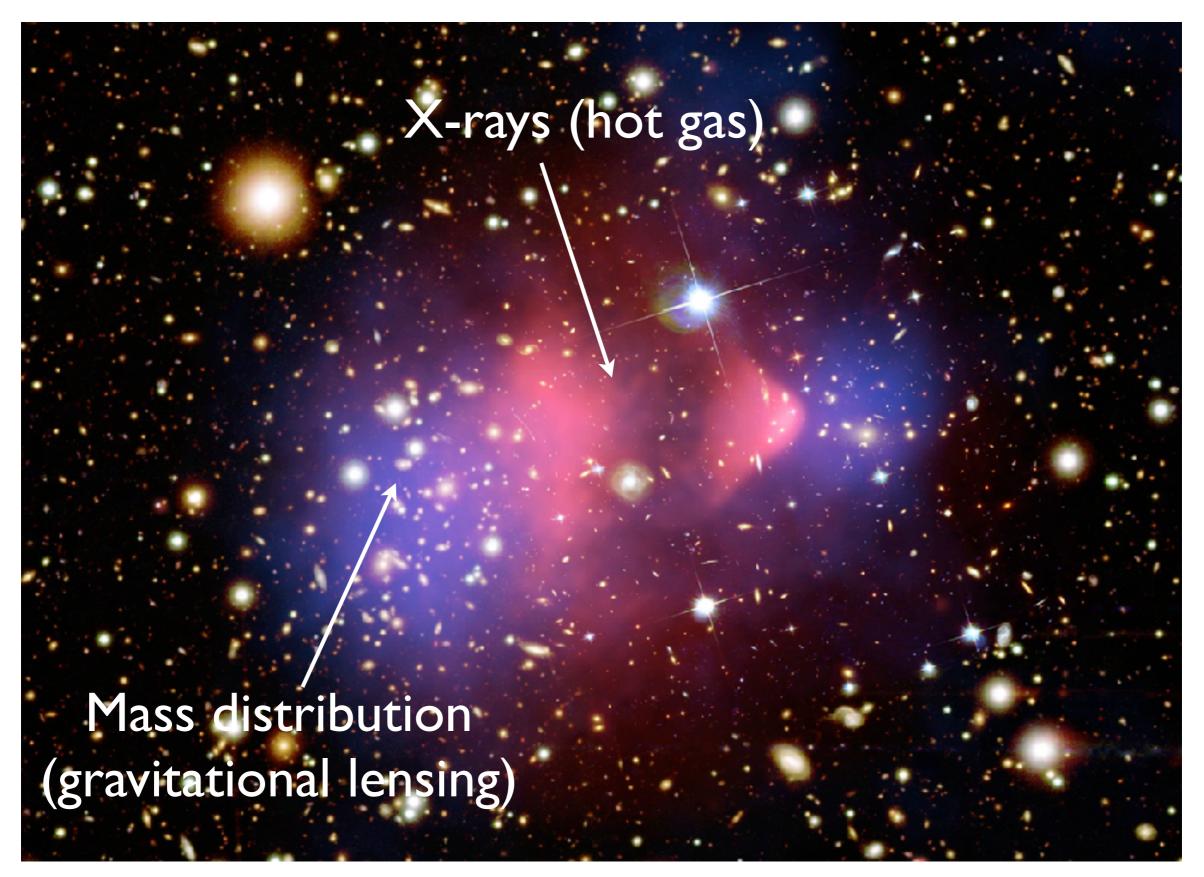
Evidence for Dark Matter



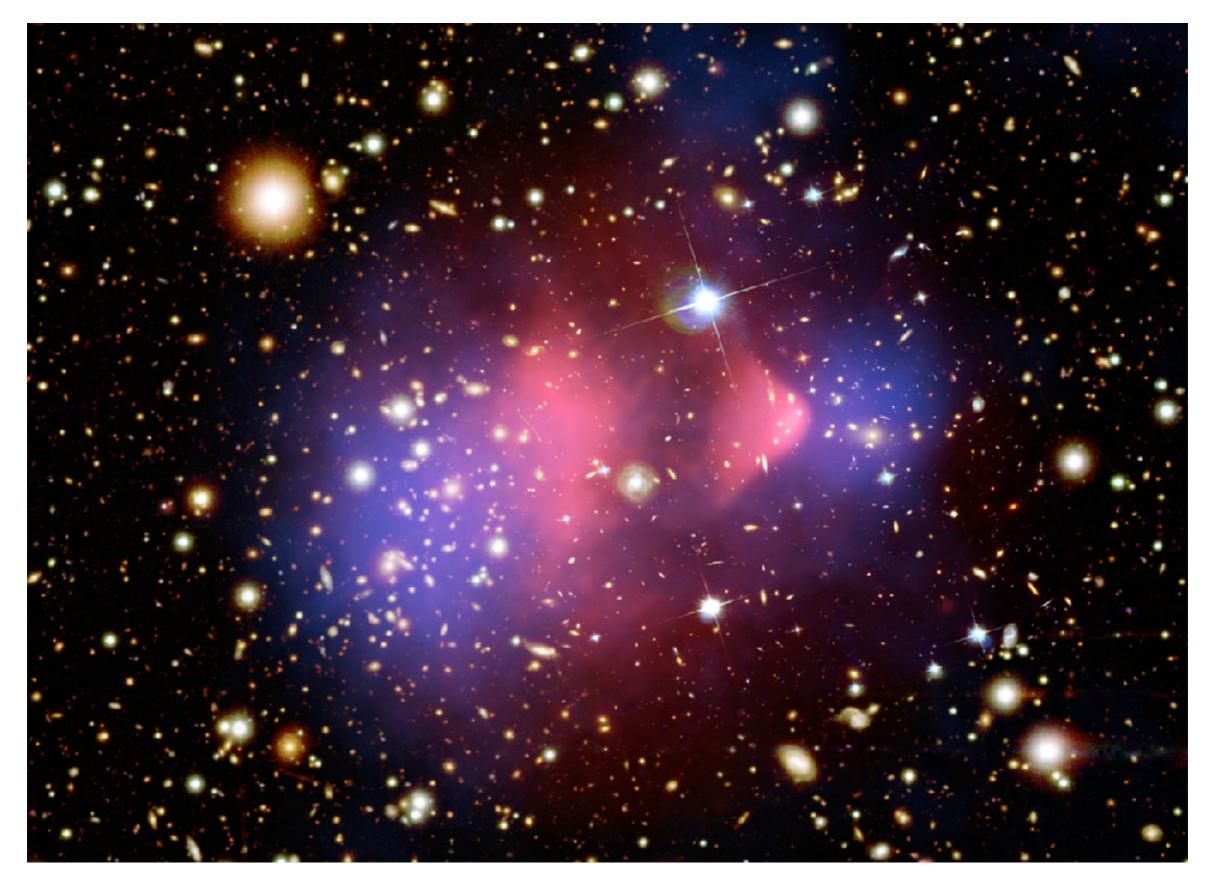
The Bullet Cluster

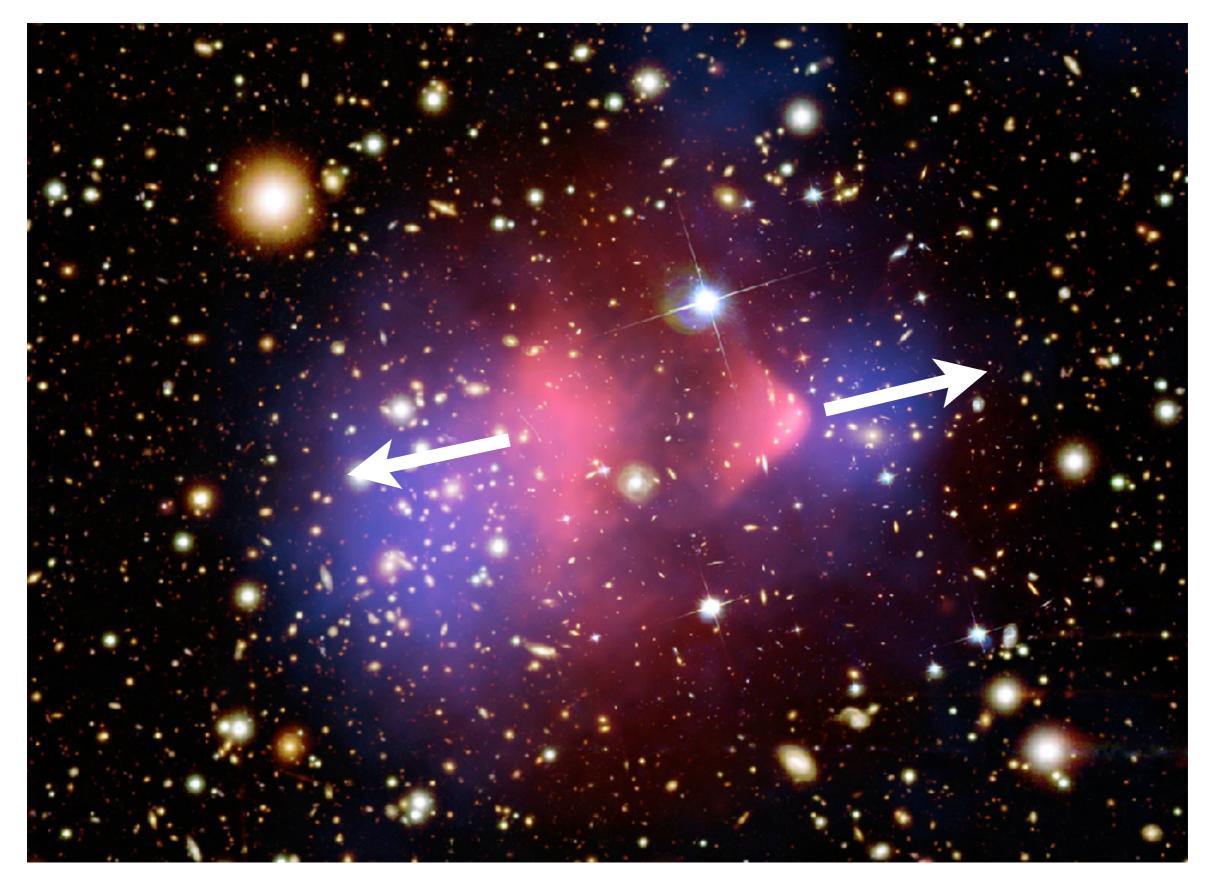


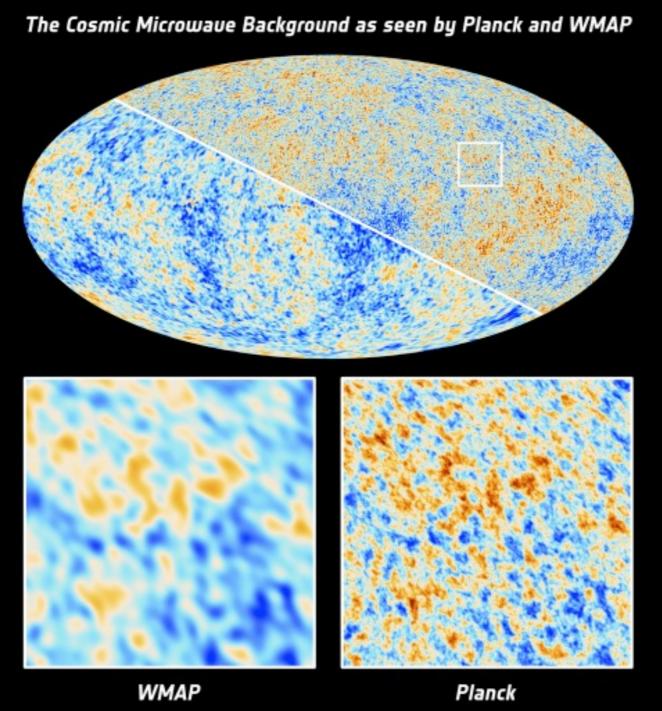
The Bullet Cluster



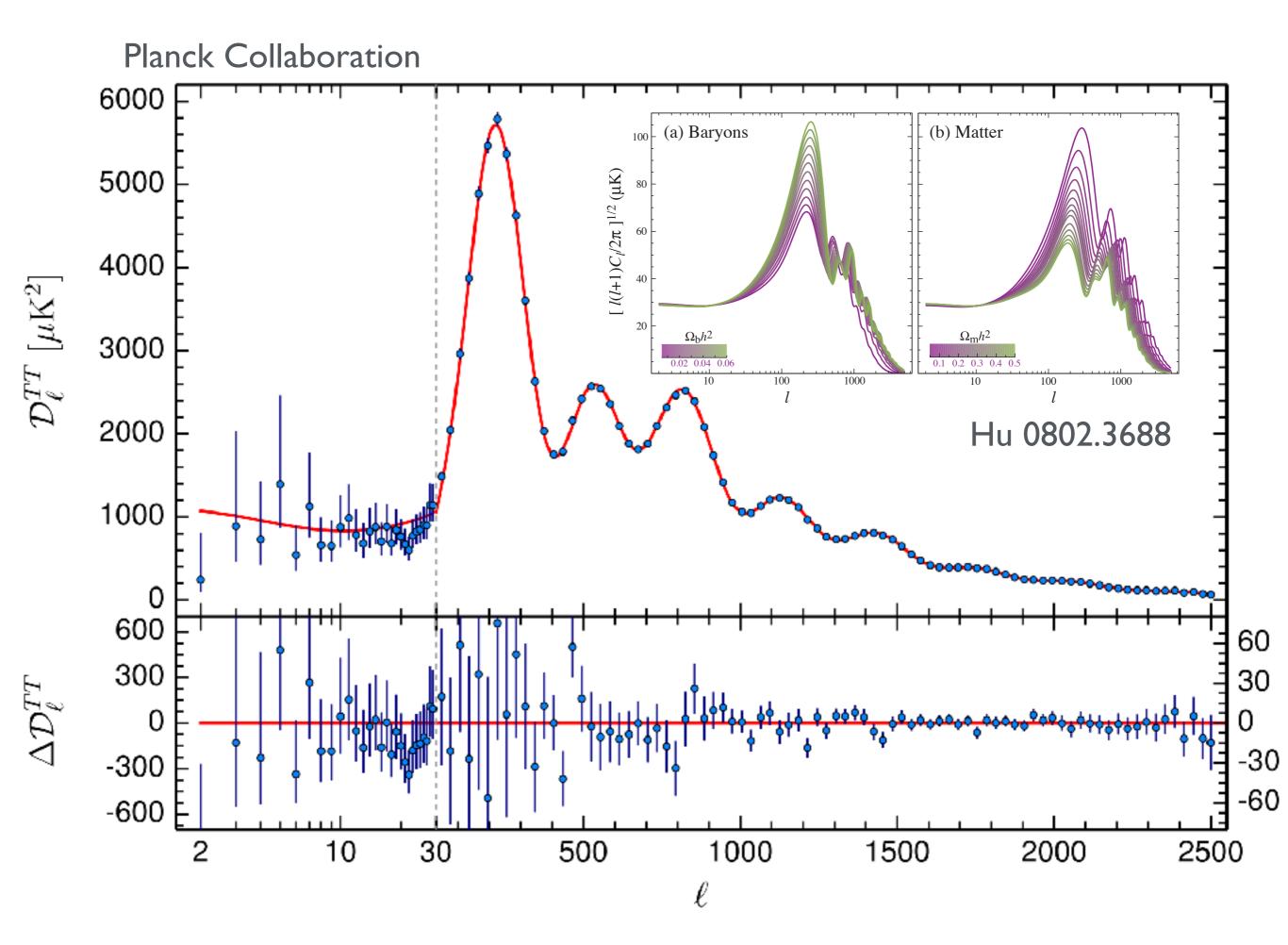
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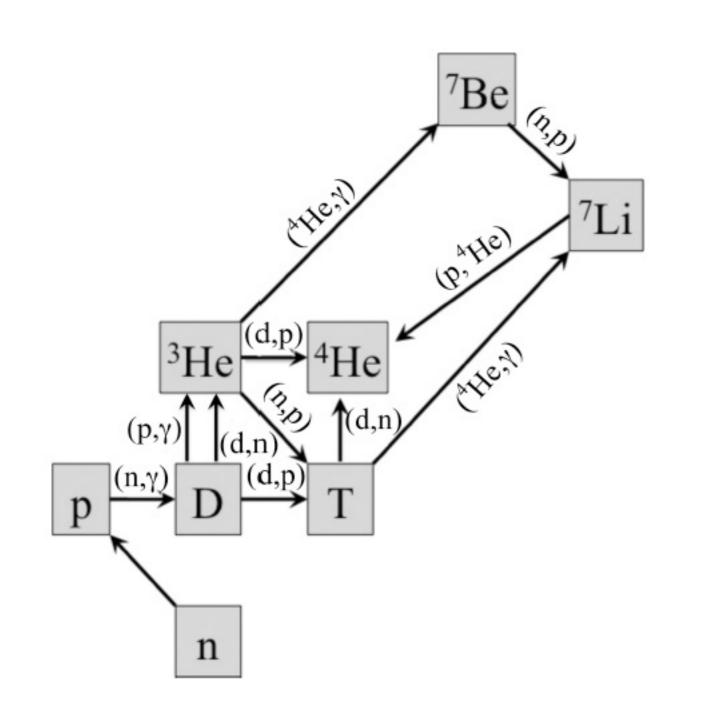


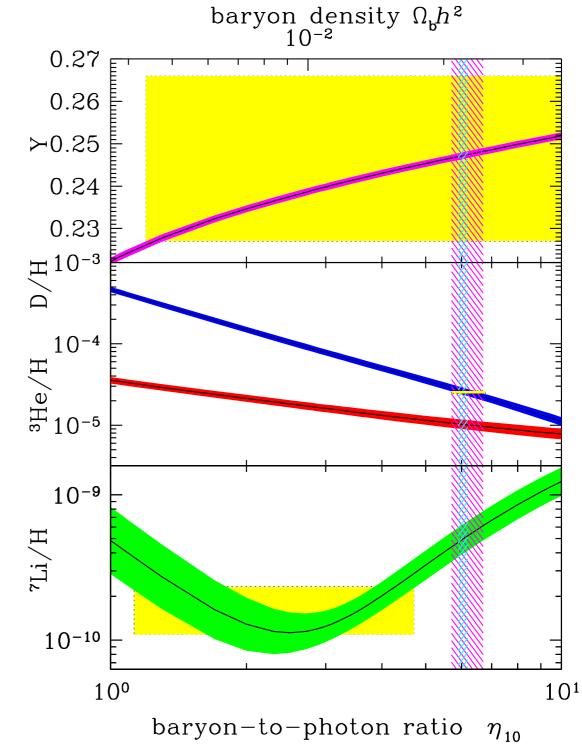


Hot plasma of hydrogen atoms and photons, and DM and cc



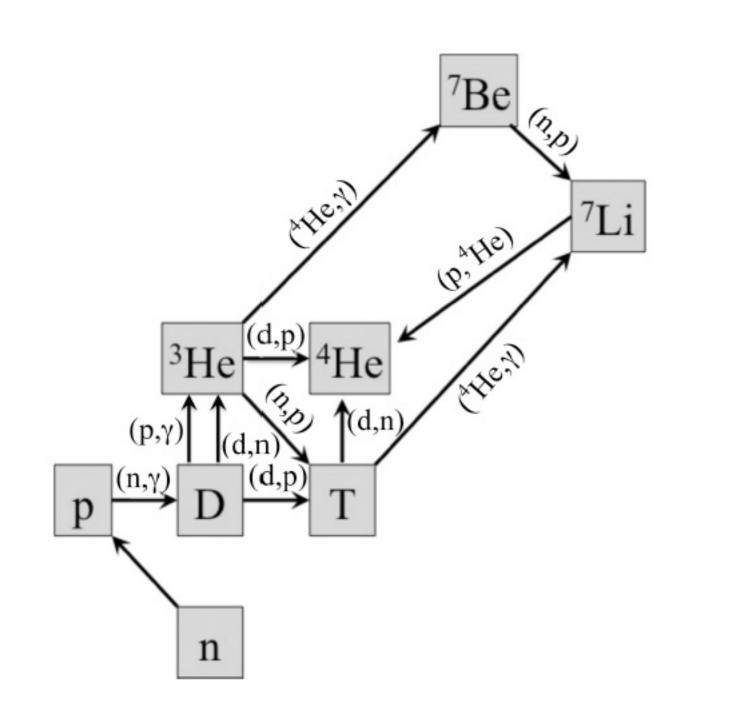
Big Bang Nucleosynthesis

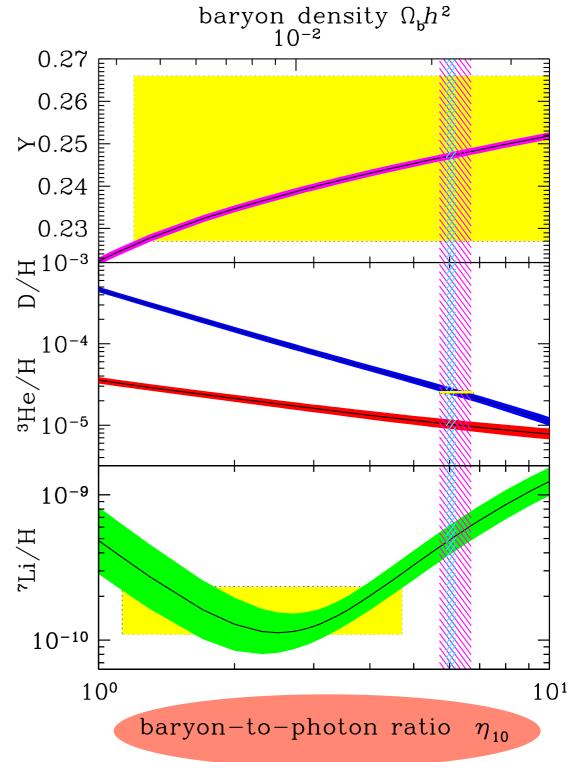




Hot soup of protons and neutrons, can predict light element abundance

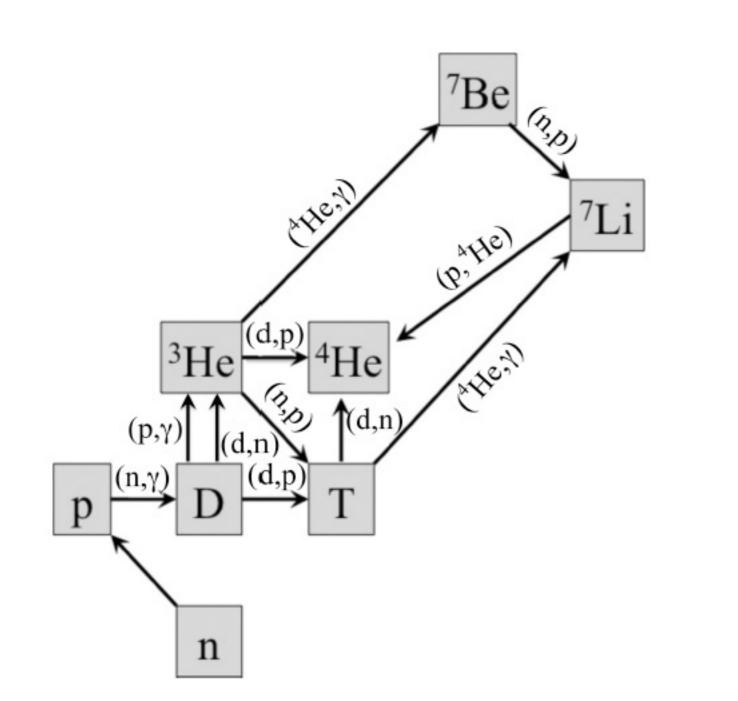
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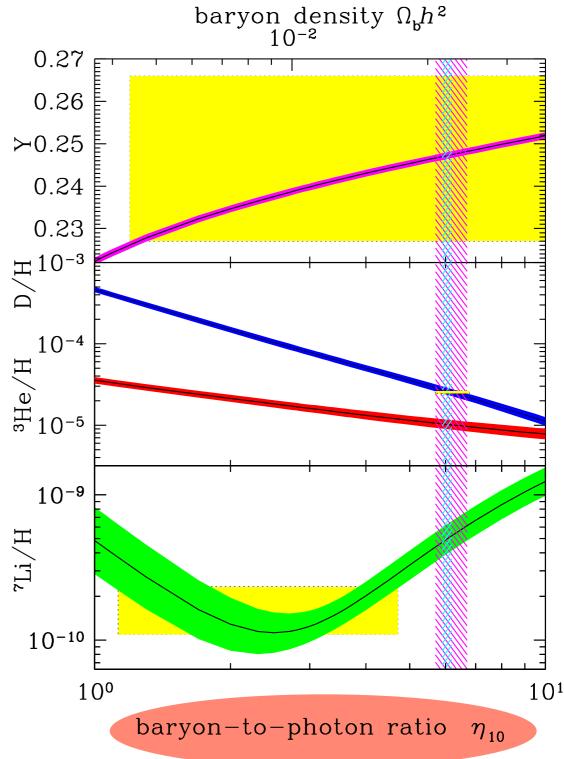




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Big Bang Nucleosynthesis





Hot soup of protons and neutrons, can predict light element abundance $\sim 5\%$ in baryons

So far all probes have been gravitational in nature

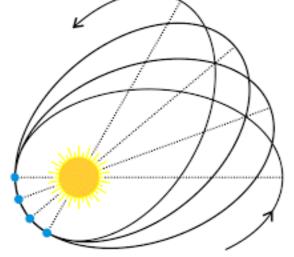
What about other interactions?

HISTORY LESSON

Neptune discovered by wobble in orbit of Uranus —original DM!



Advance in Perihelion of Mercury needed new physics (general relativity) to explain it. (Originally thought to be planet Vulcan!) —MOND??





A weak scale particle (WIMP) freezes out to leave the correct relic abundance - the WIMP "miracle"





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•At high T production and annihilation in equilibrium



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"Freeze out":
$$n\langle \sigma v \rangle \sim H \sim \frac{T^2}{M_{pl}}$$

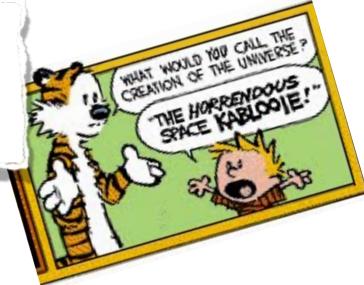


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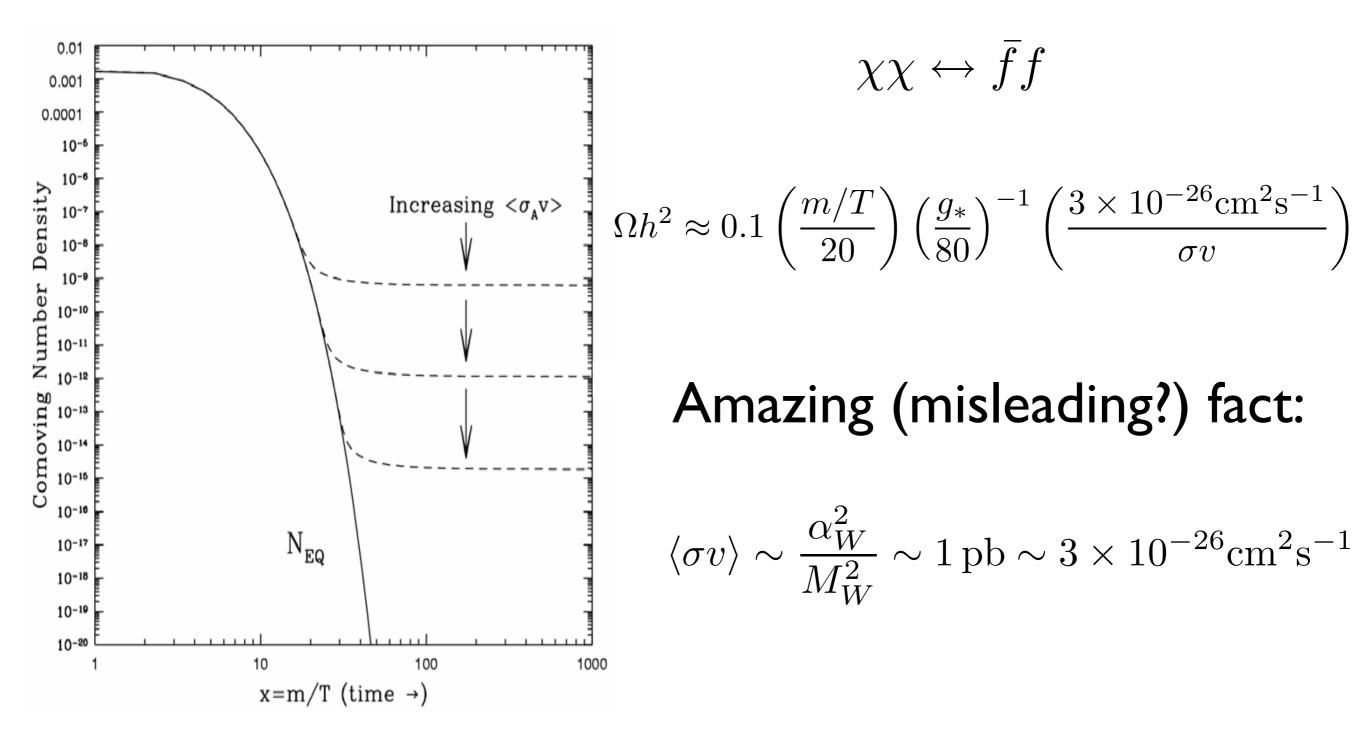
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$$\frac{dn_{\chi}}{dt} + 3Hn_{\chi} = -\langle \sigma v \rangle \left(n_{\chi}^2 - n_{eq}^2 \right)$$



A weak scale particle (WIMP) freezes out to leave the correct relic abundance - the WIMP "miracle"



DM, the story so far

- •DM makes up 23% of the universe
- •Gravitates like ordinary matter, but is non-baryonic
- •Is dark i.e. neutral under SM (not coloured, or charged)
- Does not interact much with itself
- •Does not couple to massless particle
- •Was not relativistic at time of CMB
- •Is long lived
- Is BSM physics
- IF DM is a thermal relic:
- •A weak scale annihilation x-sec gives correct abundance •Mass range is $10 \, \mathrm{keV} \lesssim m_\chi \lesssim 70 \, \mathrm{TeV}$

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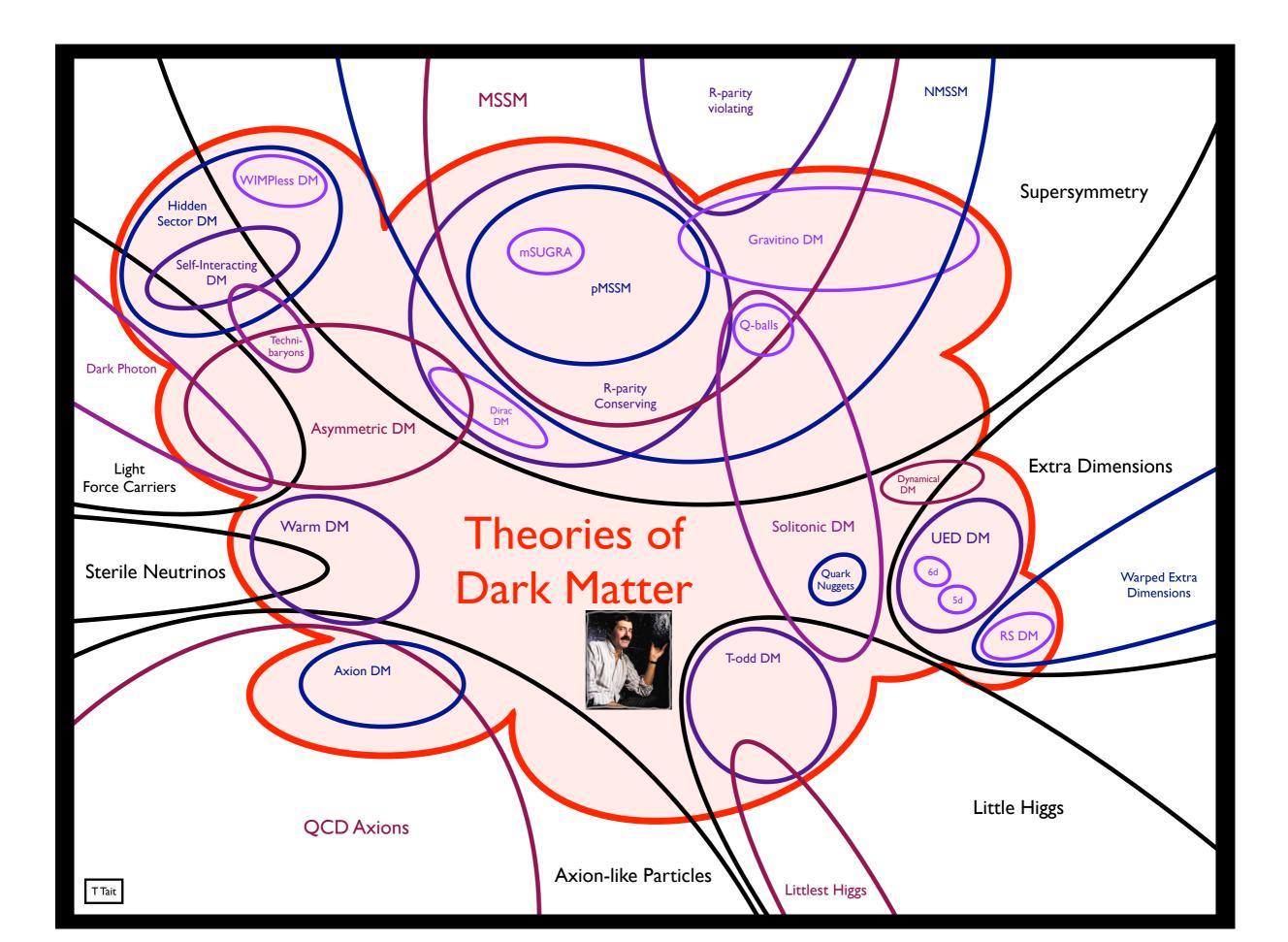
LPOPs

$\begin{array}{ll} \mbox{Many models of BSM physics contain a parity} \\ \mbox{SM} \rightarrow \mbox{SM} & \mbox{BSM} \rightarrow - \mbox{BSM} \end{array}$

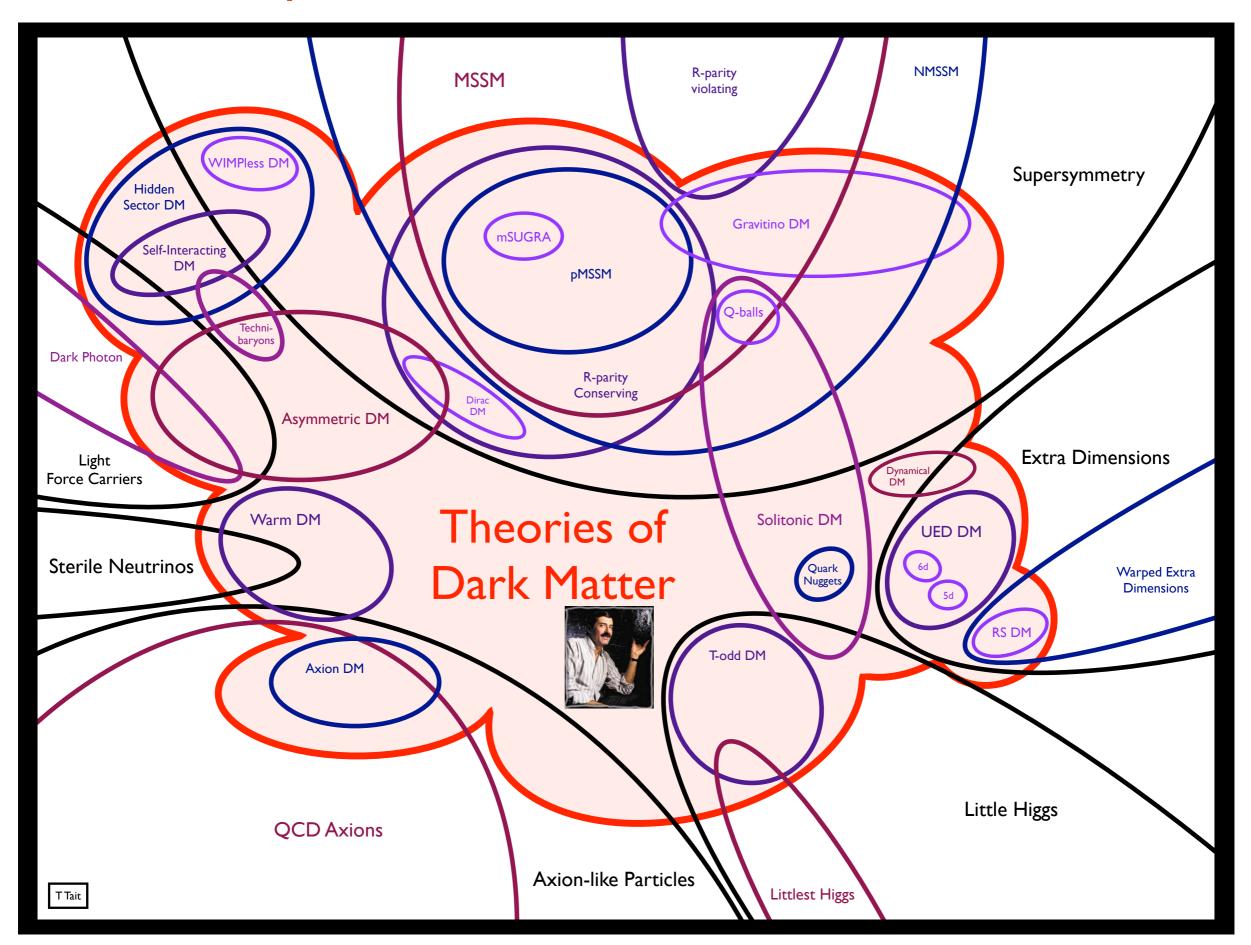
e.g. R-parity in SUSY (proton decay) T-parity in little higgs models (precision EW observables) KK-parity in extra-dimensional models

Lightest Parity Odd Particle is stable, may be a DM candidate

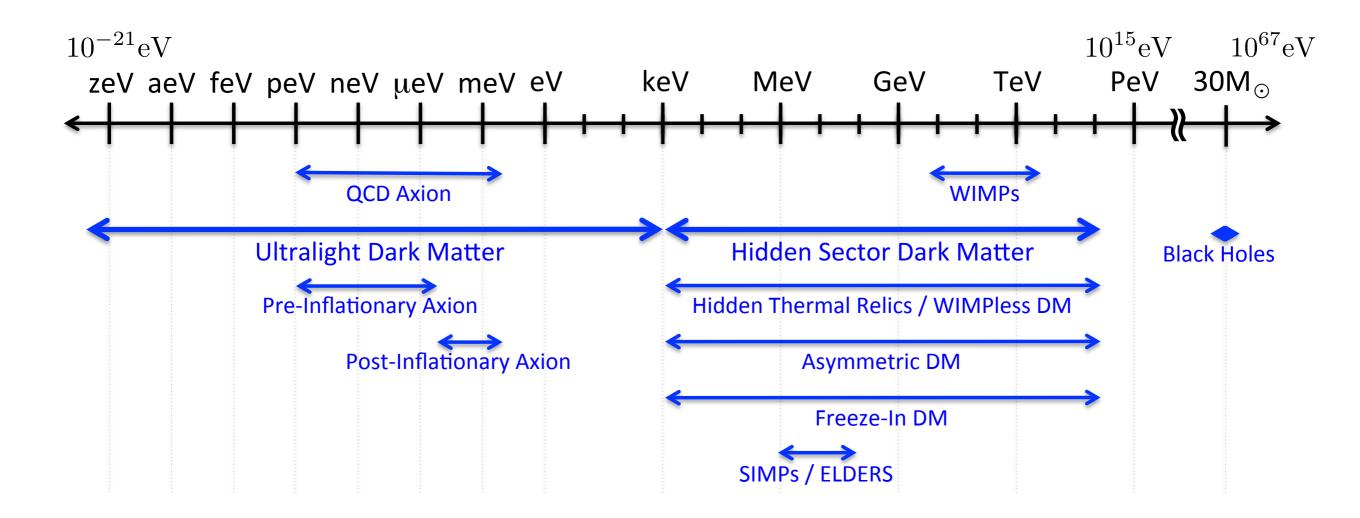
Always produced in pairs and leaves detector as MET



But such particles exist in MANY BSM models

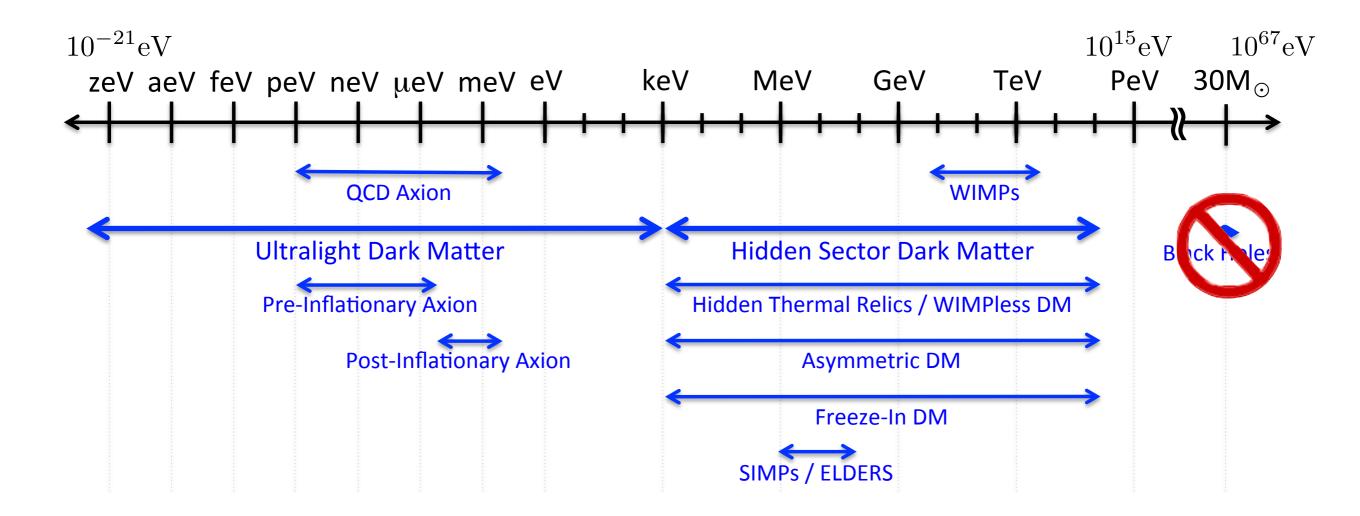


Particle theories

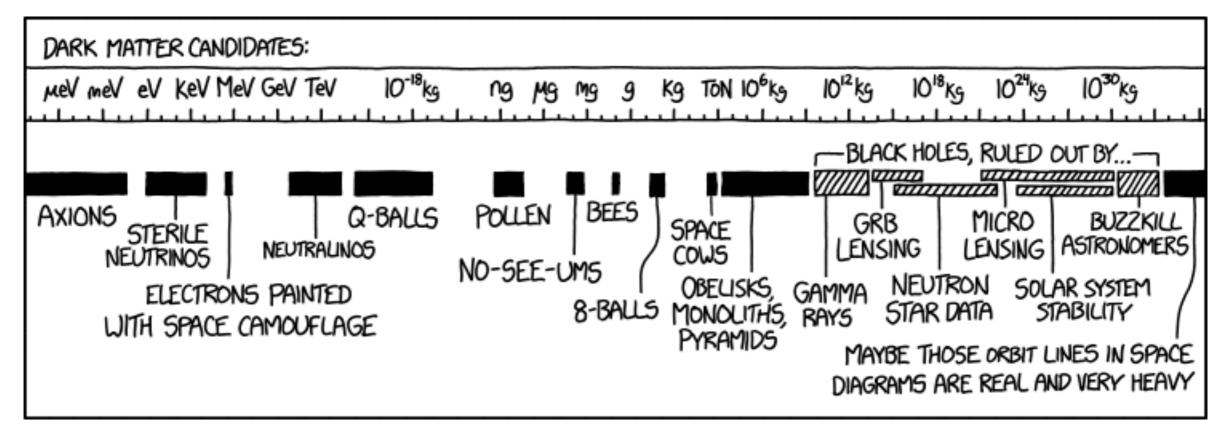


[Feng-US Cosmic Visions White papers]

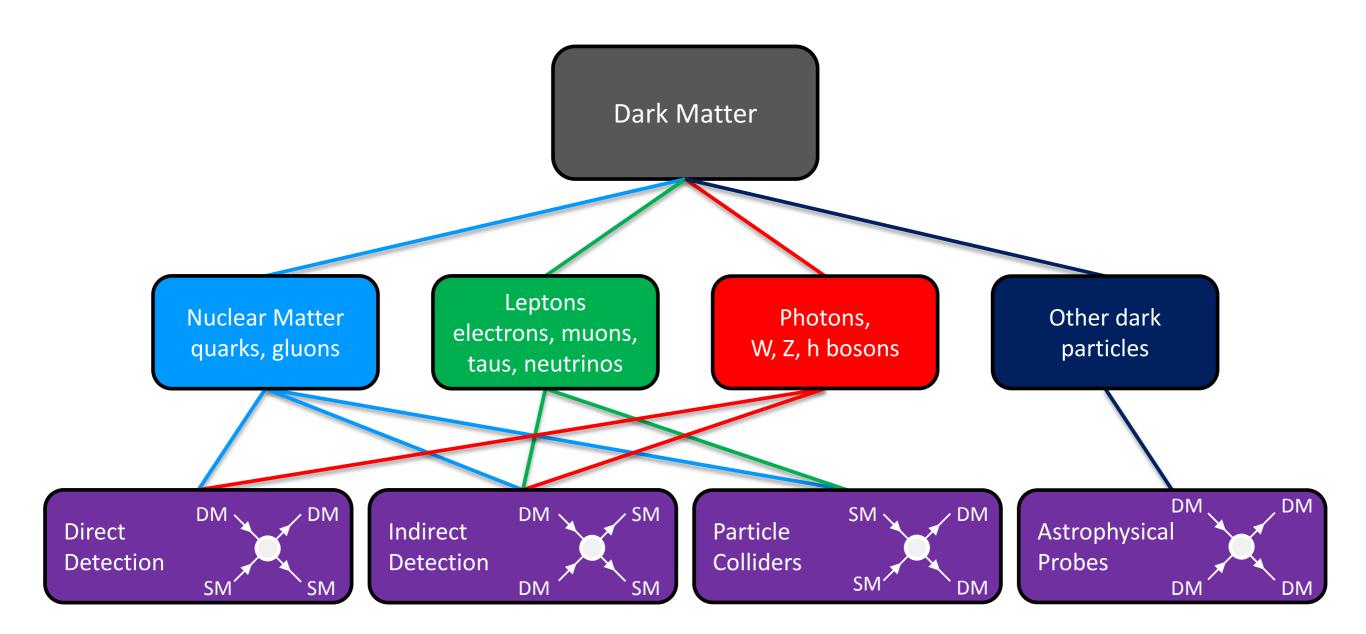
Particle theories



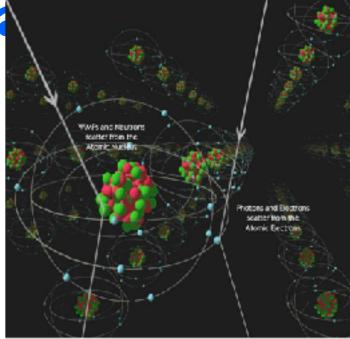
[Feng-US Cosmic Visions White papers]

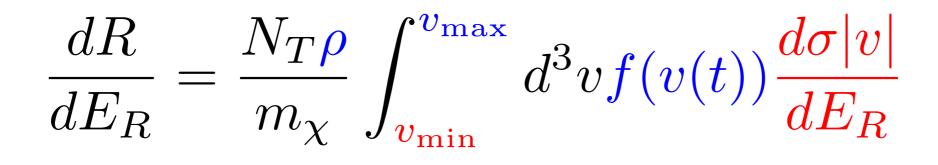




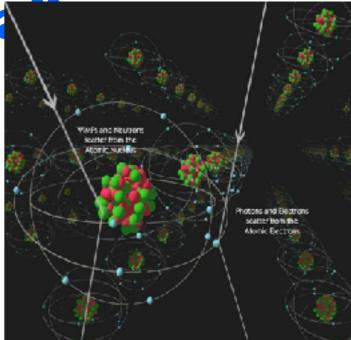


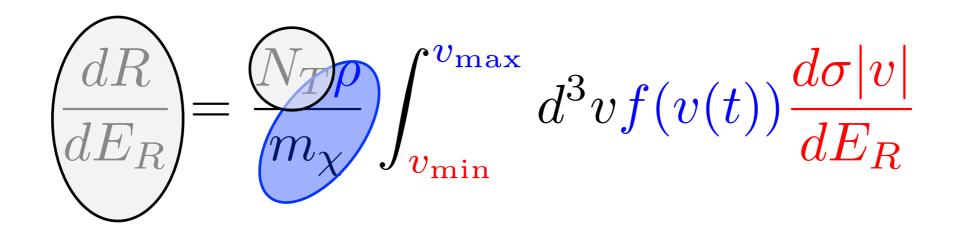
Q:Are these different search strategies separate, redundant, complementary, relatable,....?





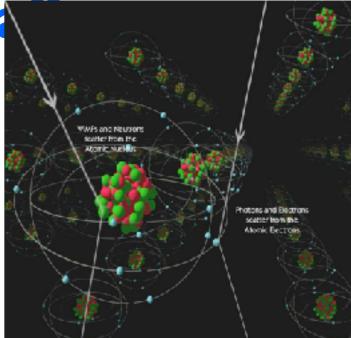
Recoil rate as a function of recoil energy Depends on how much DM is around...

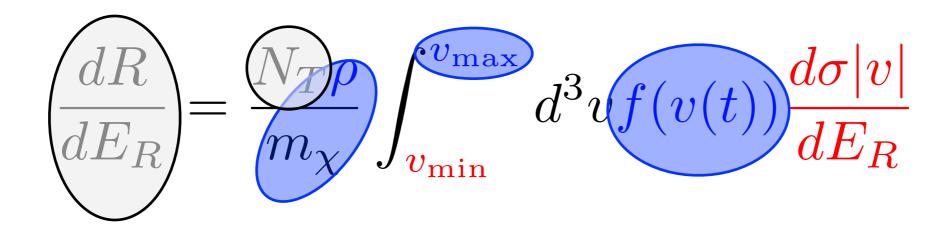




Number of targets in experiment

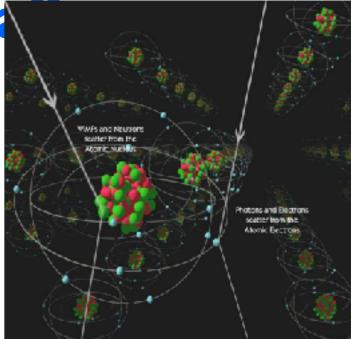
Recoil rate as a function of recoil energy Depends on how much DM is around...

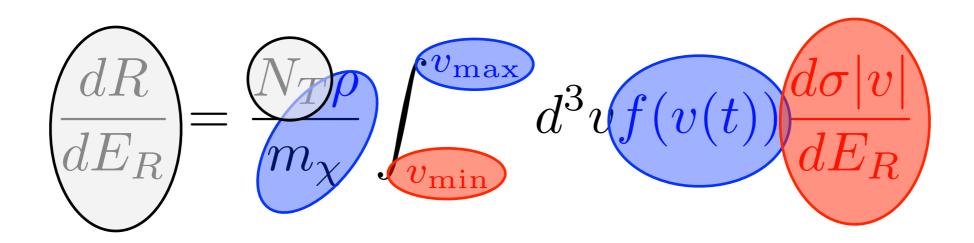




Number of targets in ...and how it's experiment moving...

Recoil rate as a function of recoil energy Depends on how much DM is around...



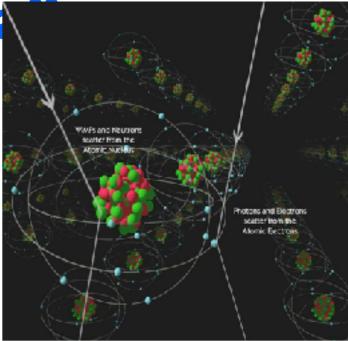


Number of targets in experiment

...and how it's moving...

...and how it interacts with nuclei.

Recoil rate as a function of recoil energy Depends on how much DM is around...

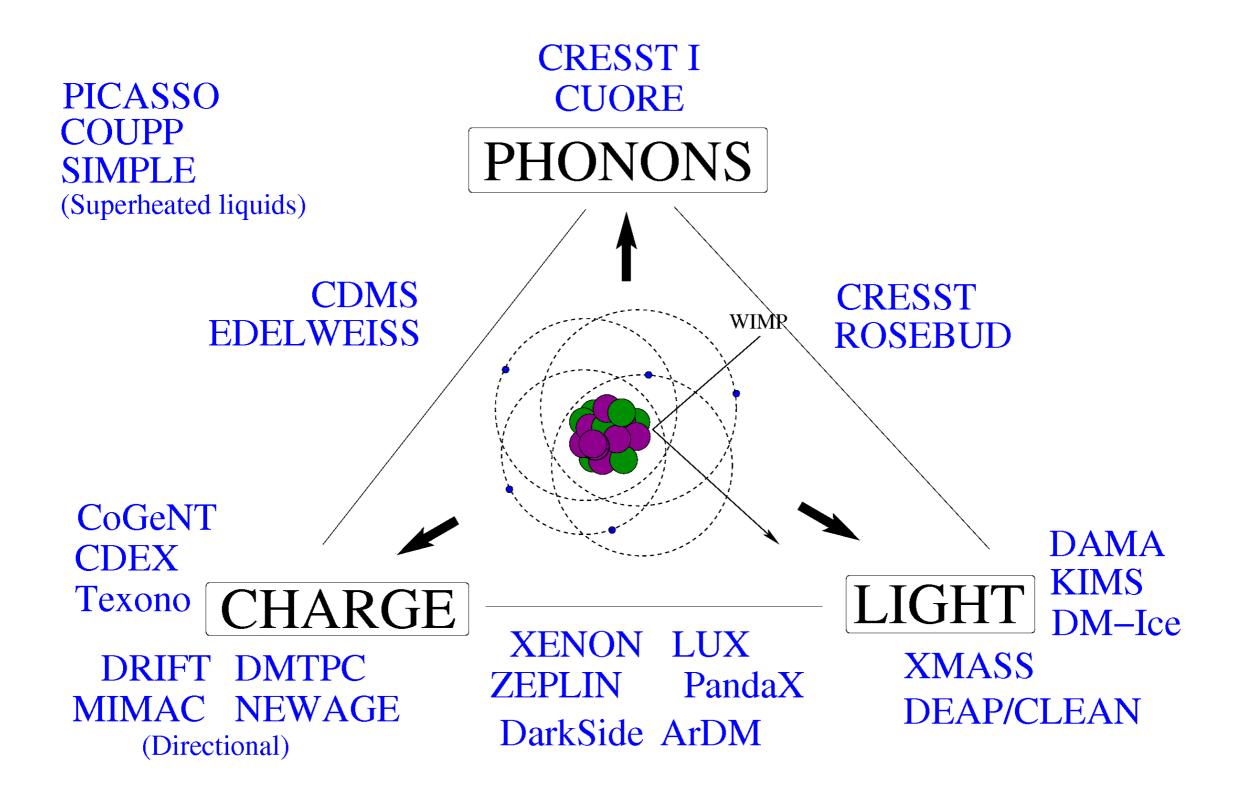




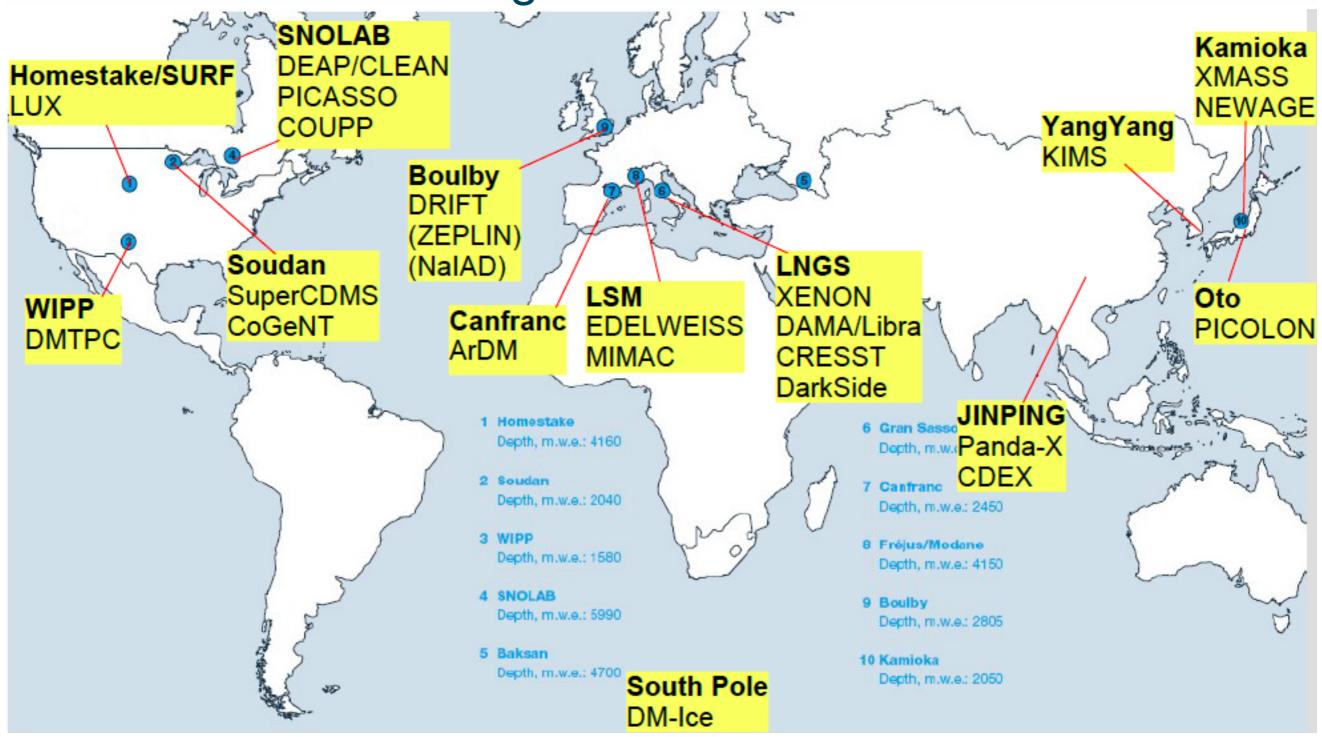
Number of targets in experiment

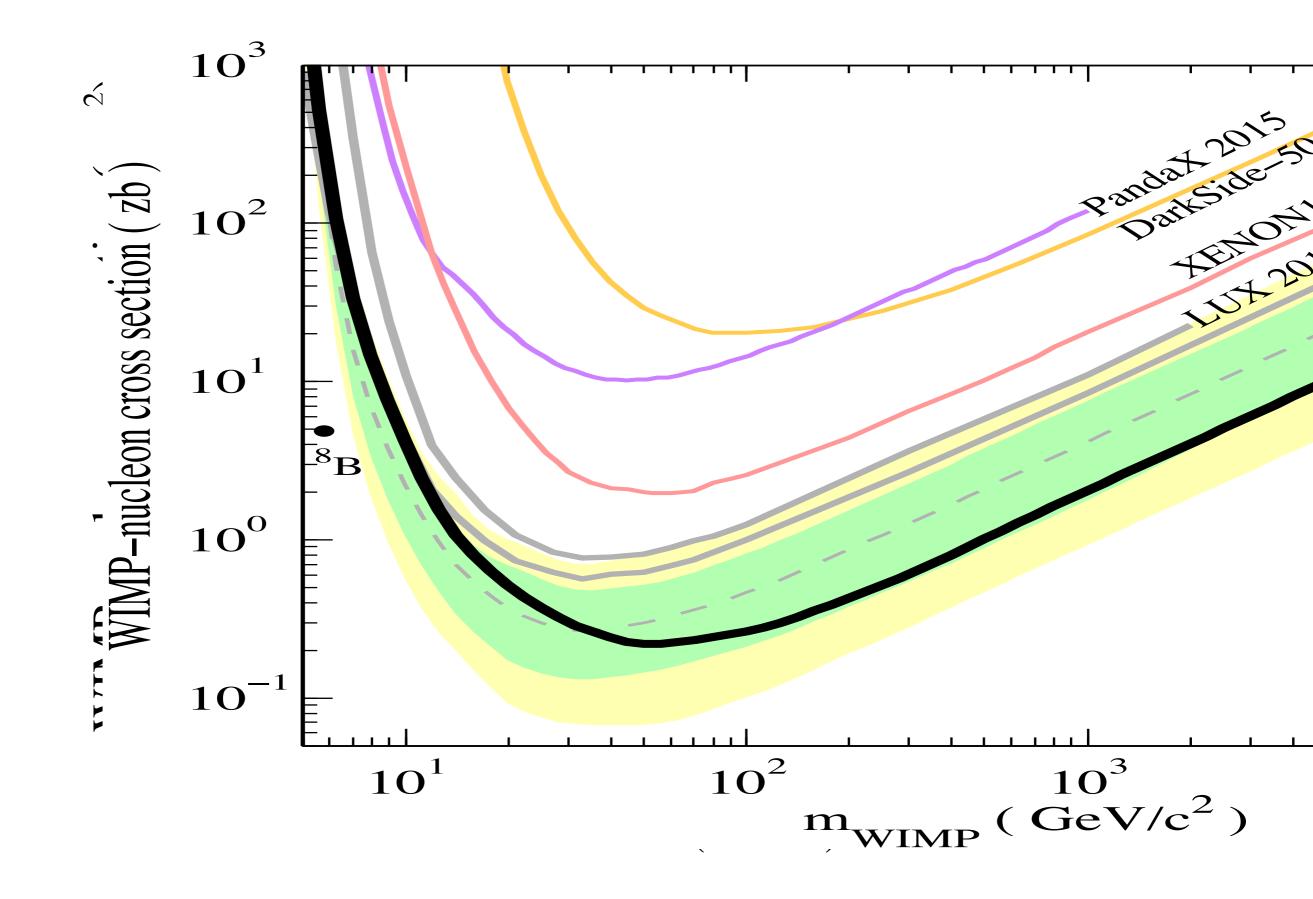
...and how it's moving...

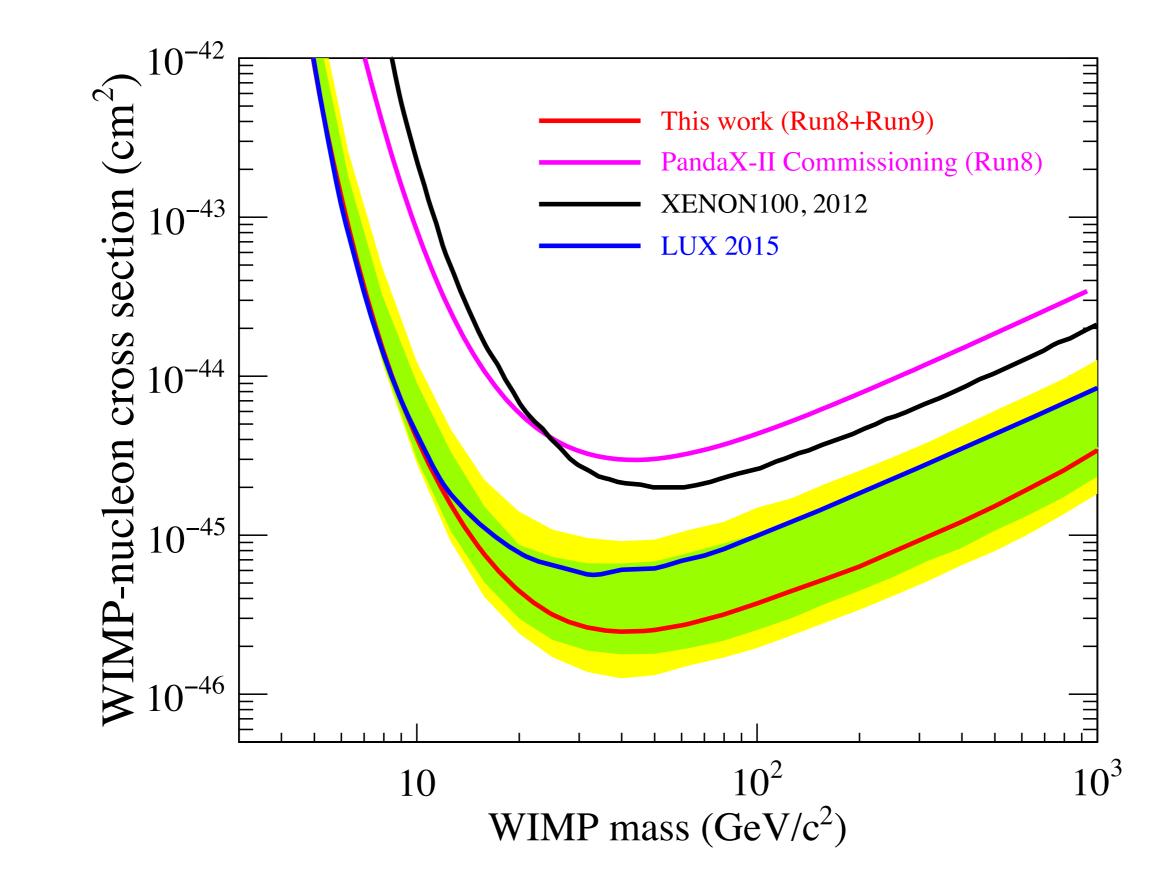
...and how it interacts with nuclei.



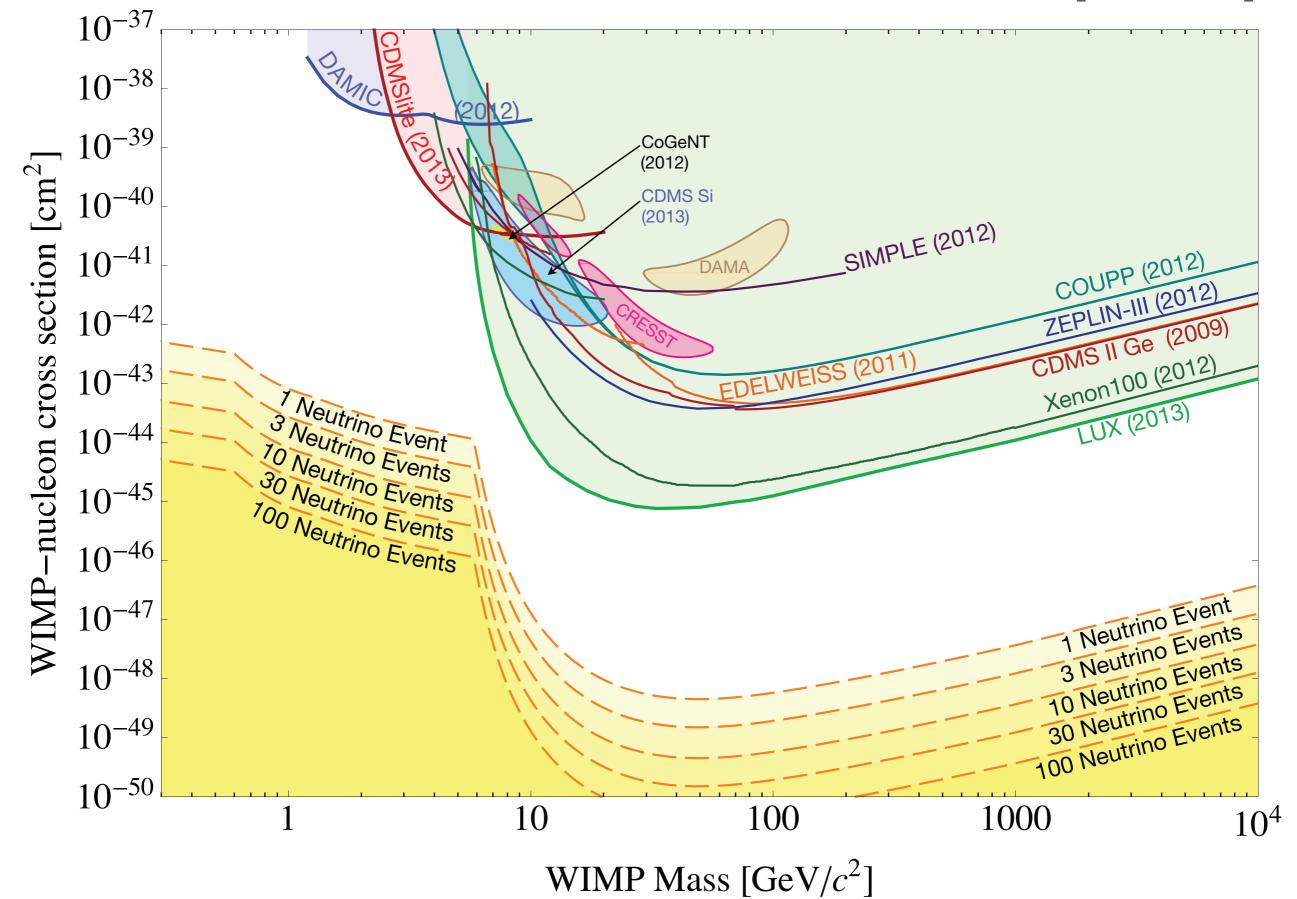
Underground laboratories

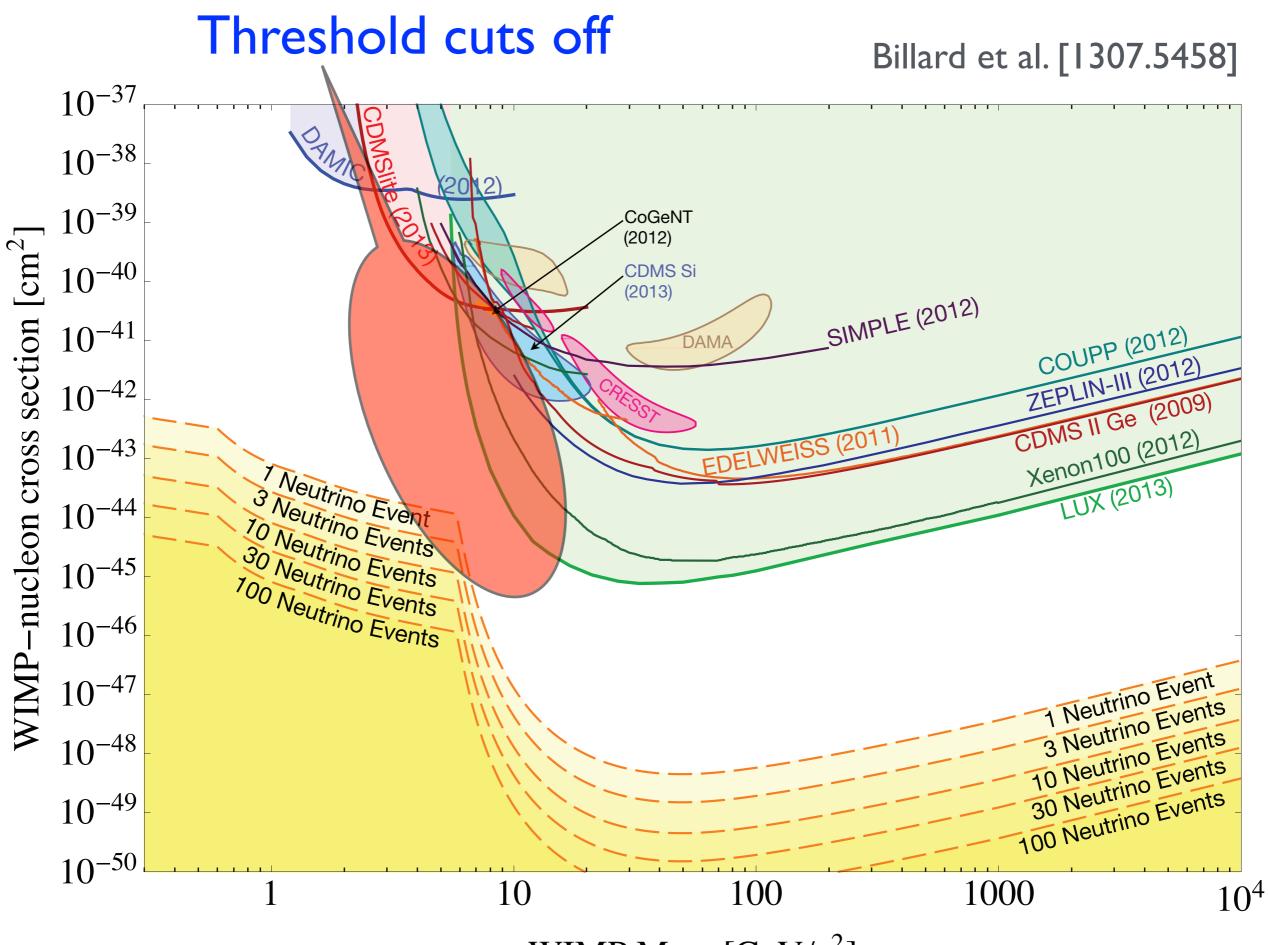




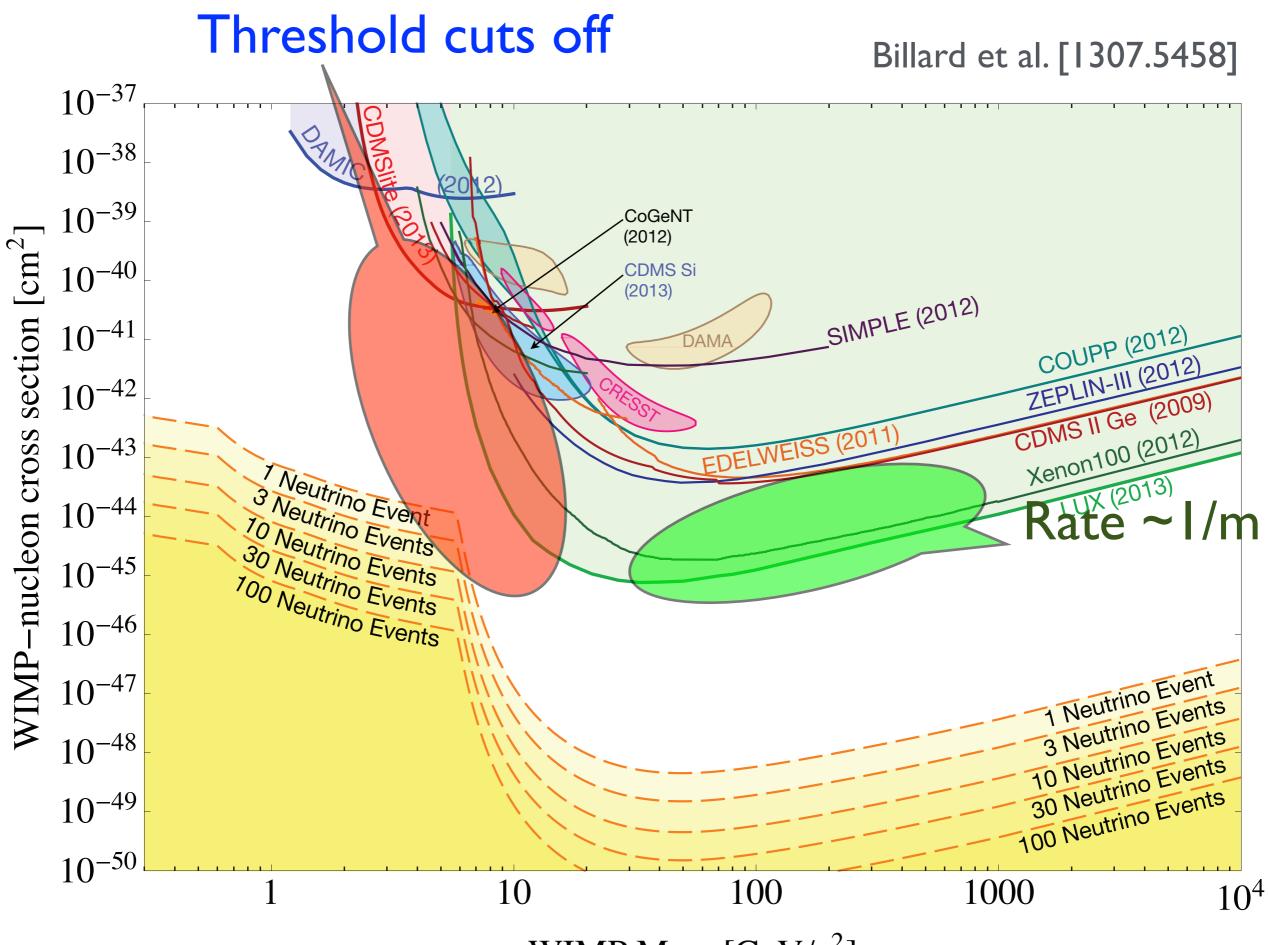


Billard et al. [1307.5458]



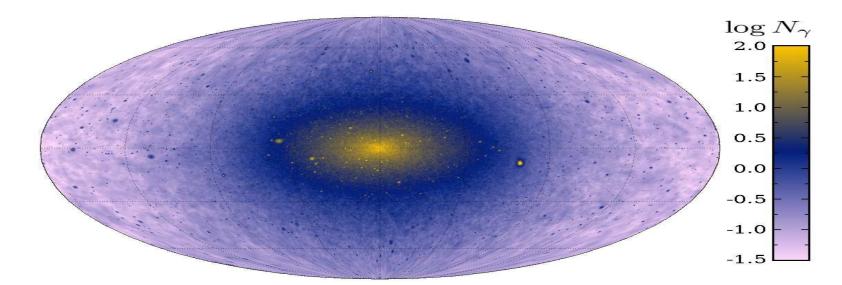


WIMP Mass [GeV/ c^2]



WIMP Mass [GeV/ c^2]

ndirect Detection "Master formula"



$$\frac{dN}{d\Omega dE}(\psi) = \frac{1}{4\pi\eta} \frac{f_{\chi}^2 J(\psi)}{m_{\chi}^2} \sum_i \langle \sigma v \rangle_i \frac{dN^i}{dE_{\gamma}}$$

Spectrum of particles in final state

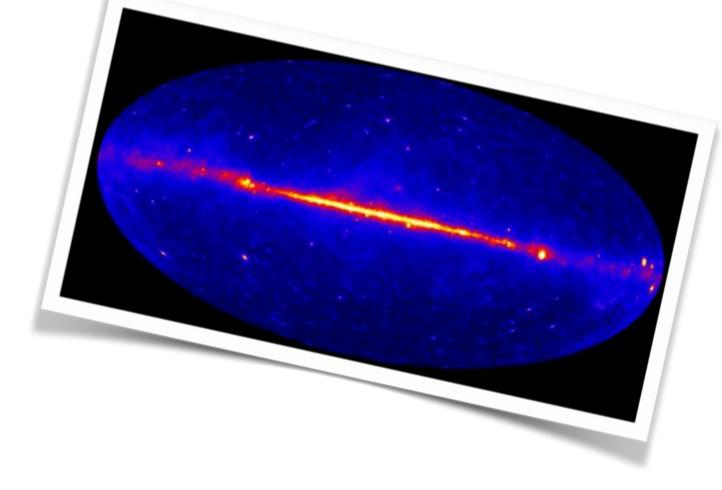
 $J(\psi) = \int_{\text{l.o.s.}} ds \,\rho(r)^2$

Line of sight integral

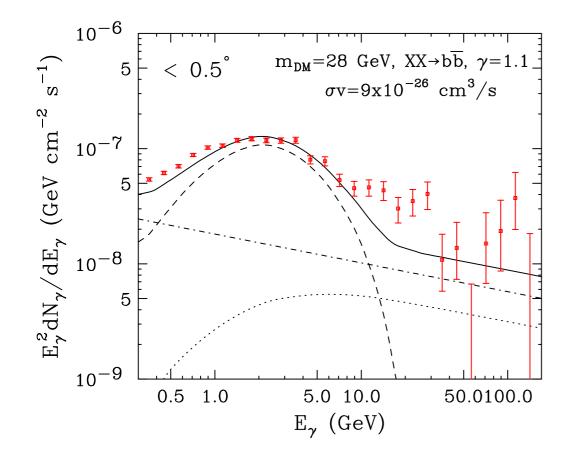
Dark Matter Indirect Detection

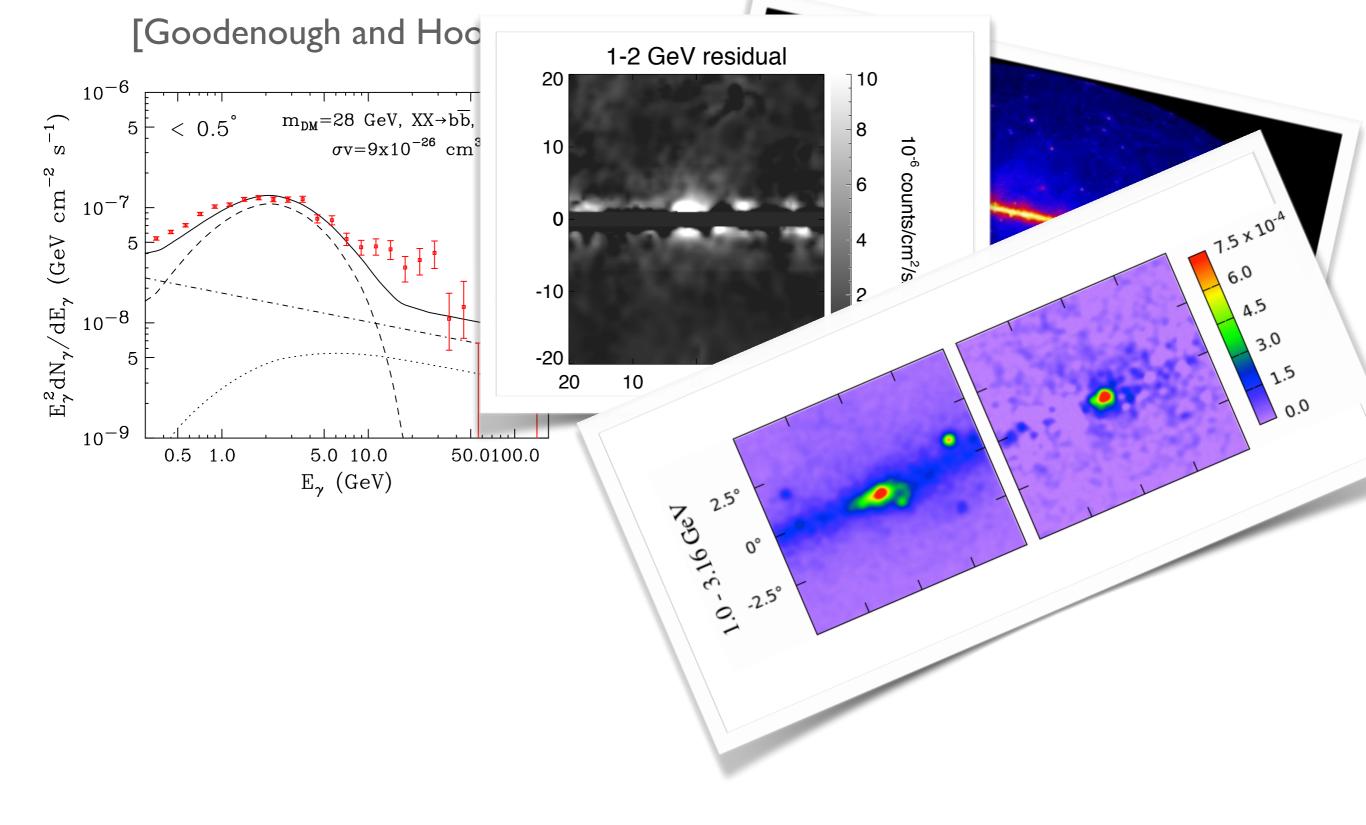
DM annihilates in our galaxy, or nearby dwarf galaxy e.g.

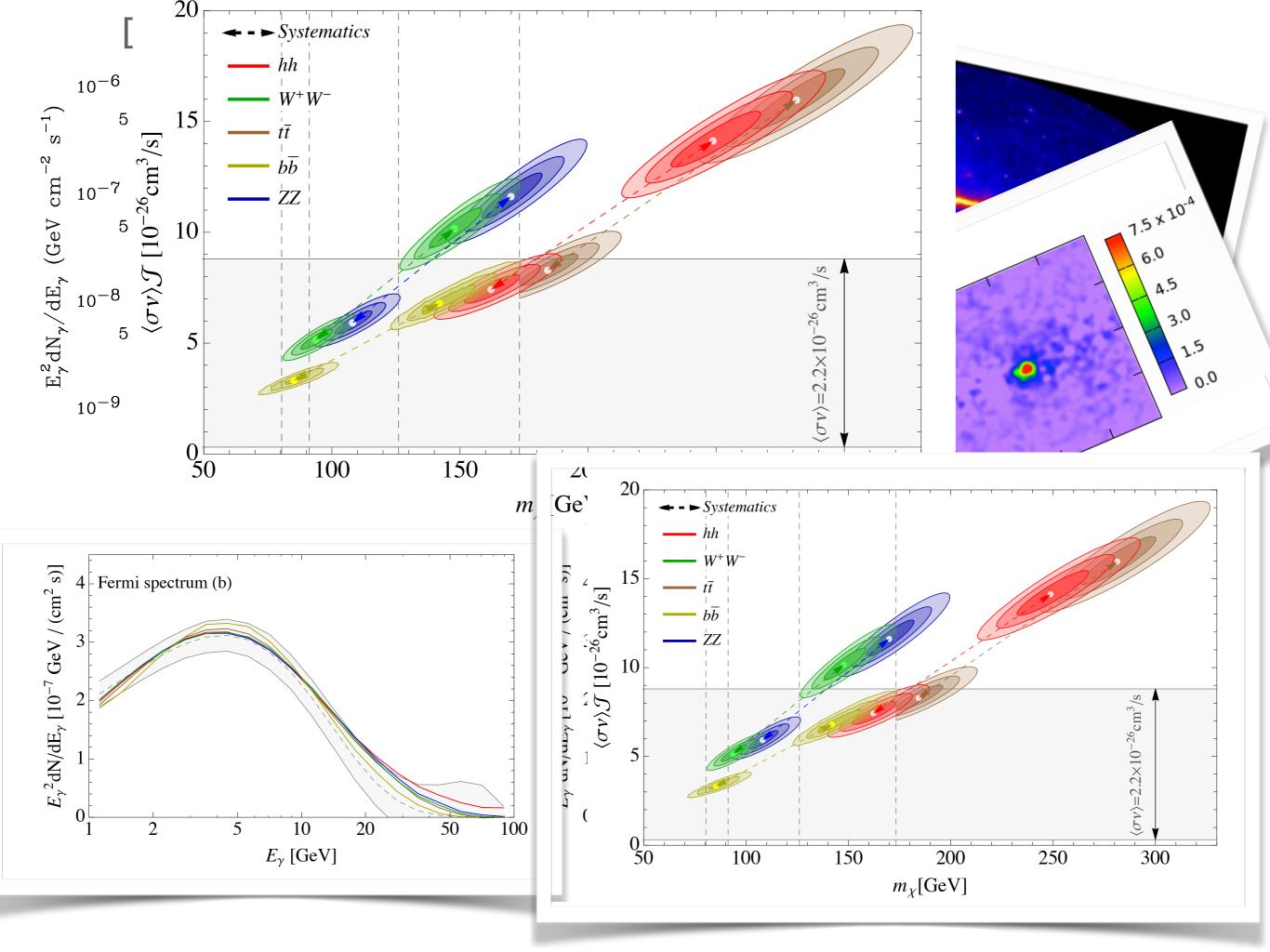
$\chi\chi \to p\bar{p}, e^+e^-$	Look for antimatter in cosmic rays, does not point back to source, limited range. PAMELA, AMS02, Fermi
$\chi\chi \to \nu\bar{\nu}$	Point back to source, low cross section. IceCube, ANTARES, Super-K
$\chi \chi o \gamma \gamma$	Point back to source, spectral line, low rate Fermi, HESS
$\chi \chi \to \mathrm{SM} \ \mathrm{SM}$ $\hookrightarrow \ldots + \gamma \gamma$	Point back to source, continuum with edge, backgrounds Fermi, HESS



[Goodenough and Hooper, 2009]







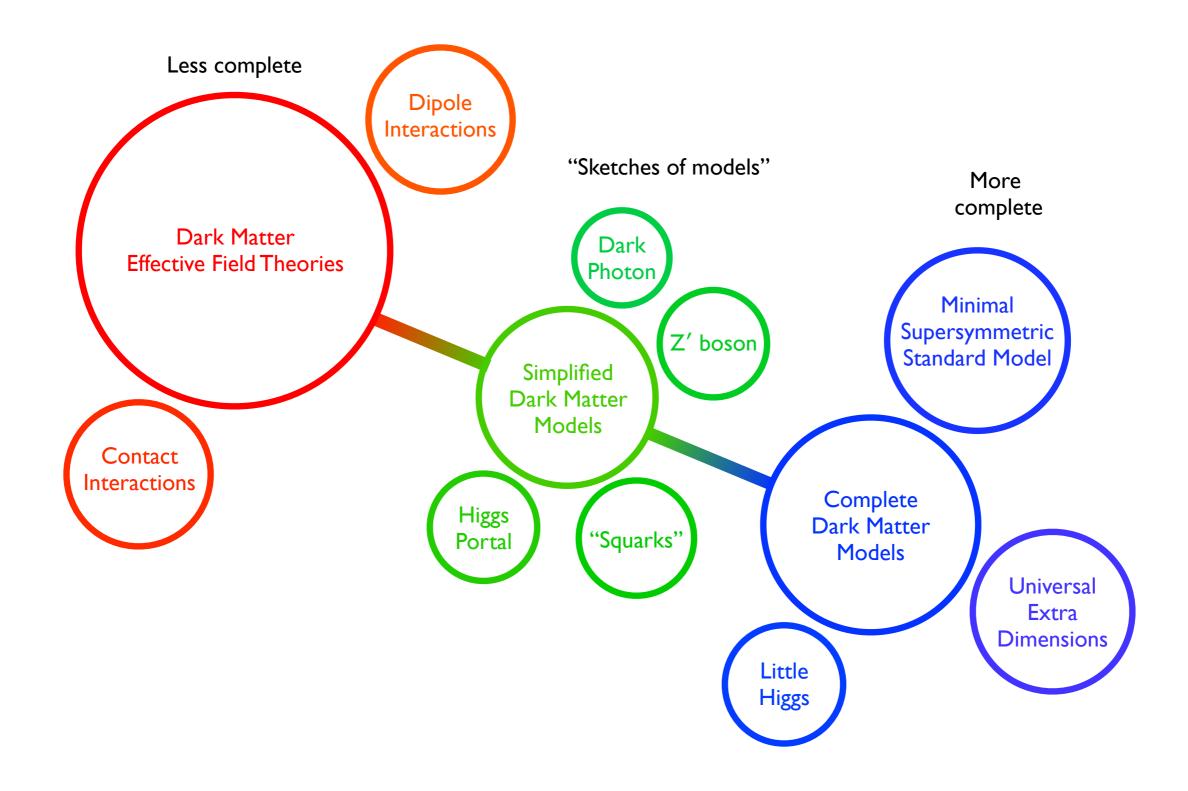
Are the excess photons from the Galactic centre DM?

- •Source is spherical, with the expected radial dependence
- Cross section is close to thermal
- •Centred in the right place



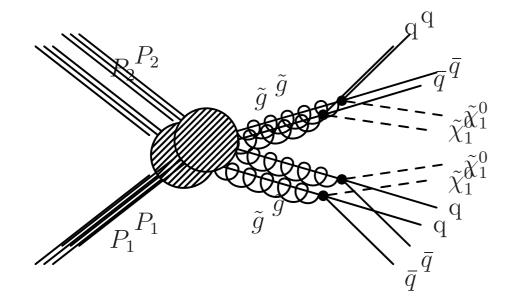
- •Statistical significant, and Fermi-team sees it too
 - •Galactic centre is a confusing place
 - •Not as clear as a spectral line
 - •Milli-second pulsars (but we would have seen more, also spectrum different from those observed)
 - Look at other DM "bright spots"--dwarf galaxies
 - Cosmic ray anti-particles
 - •Correlated signals, LHC, direct detection
 - Interesting times ahead

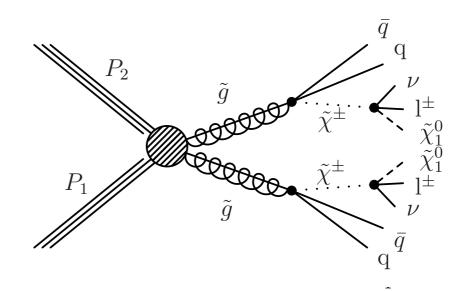
Ways to search for DM at colliders



Ways to search for DM at colliders

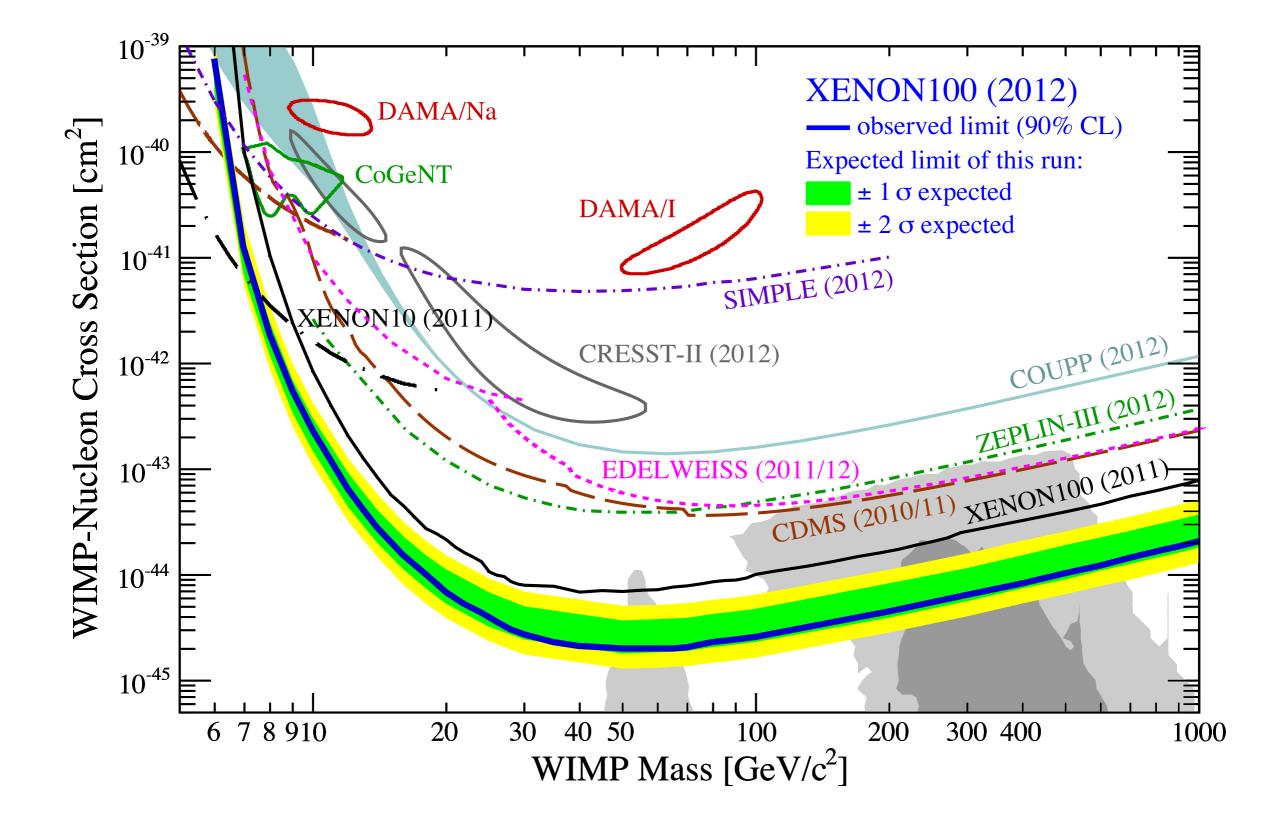
Use a full UV model (e.g. SUSY)

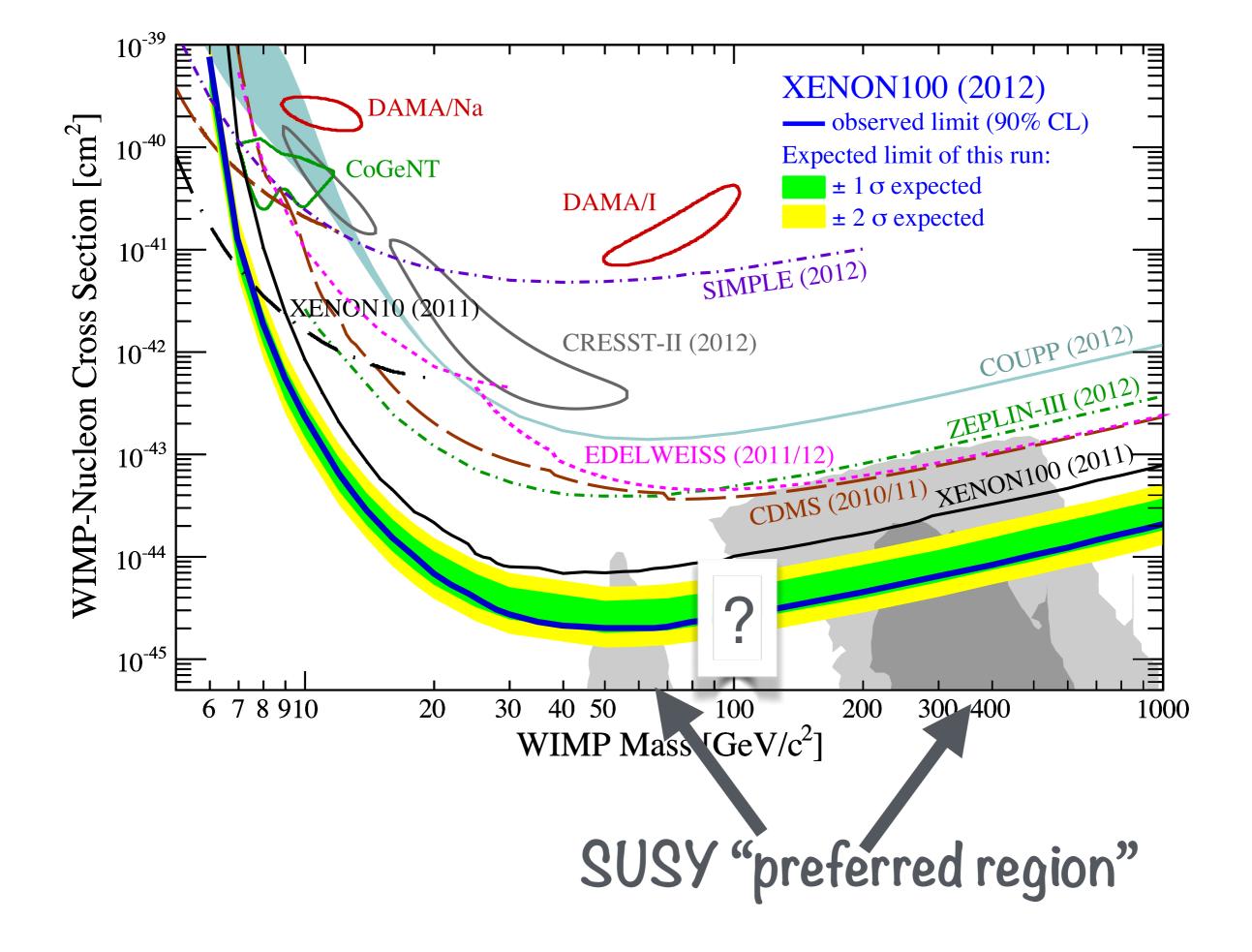




Thursday, 2 August 2012 Thursday, 2 August 2012

Complicated/interesting final state. Tuned analyses No clear relation between different search strategies

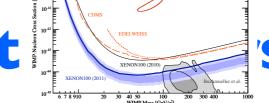




Q:Are these different search strategies separate, redundant, complementary, relatable,....?

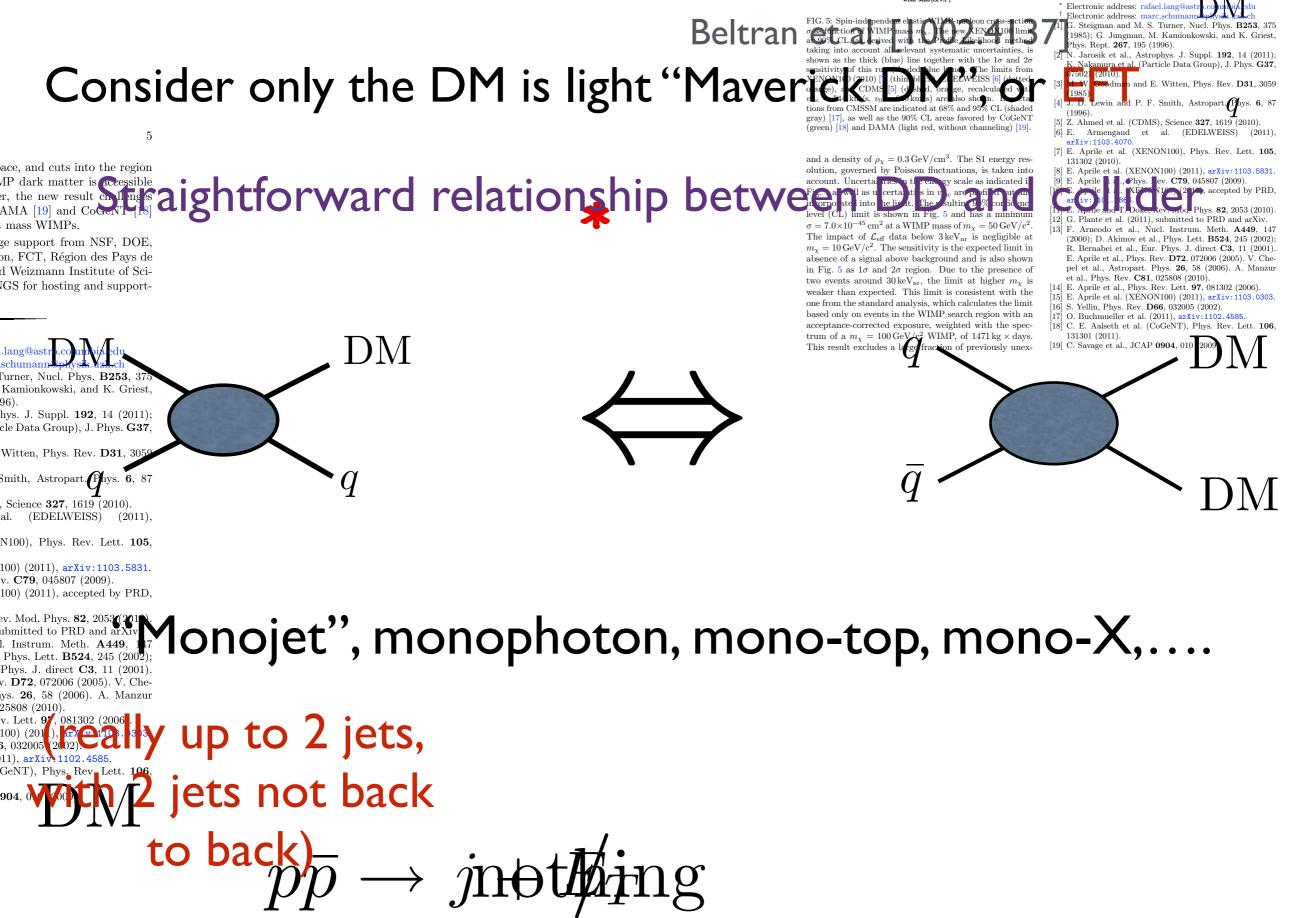
A: traditionally there was no clear way to relate them

Ways to search for DM at



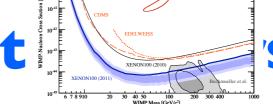
the interpretation of the DAMA [19] and CoGeNT [18] results as being due to light mass WIMPs.

We gratefully acknowledge support from NSF, DOE, SNF, Volkswagen Foundation, FCT, Région des Pays de la Loire, STCSM, DFG, and Weizmann Institute of Science. We are grateful to LNGS for hosting and supporting XENON.



Ways to search for DM at

DM



Consider only the DM is light "Maver" Maver to the busic states of the back of

ace, and cuts into the region AP dark matter is accessible or, the new result collinges aightforward relationship between account. Uncerta AMA [19] and CoGeNT [18]

mass WIMPs.

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Turner, Nucl. Phys. **B253**, 37 Kamionkowski, and K. Griest,

hys. J. Suppl. 192, 14 (2011); cle Data Group), J. Phys. G37,

Witten, Phys. Rev. D31, 30

Smith, Astropart. Phys. 6, 87

Science **327**, 1619 (2010). al. (EDELWEISS) (2011),

N100), Phys. Rev. Lett. 105,

100) (2011), arXiv:1103.5831. v. C79, 045807 (2009). 100) (2011), accepted by PRD,



ev. Mod. Phys. 82, 2053 (2012). ubmitted to PRD and arXiv I. Instrum. Meth. A449, 117 Phys. Lett. B524, 245 (2002); Phys. Lett. B524, 245 (2002);

(green) [18] and DAMA (light red, without channeling) [19].

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* Electronic address: rafael.lang@a Electronic address: marc.schuma

hys. Rept. 267, 195 (1996).

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arXiv:1103.4070.

131302 (2010).

] E. Aprile

Mono-mania at the LHC



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ng@astro.commoia.edu numann@physik.uza.ch

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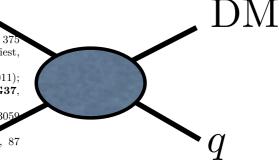
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Mod. Ph/s. 82 2053 (2010). mitted to BB and arXiv. Instrum. Meth. A449, 147 hys. Lett. B524, 245 (2002); nys. J. direct C3, 11 (2001). D72, 072006 (2005). V. Che-. 26, 58 (2006). A. Manzur

808 (2010). Lett. **97**, 087502 (2006). 0) (2011), arXiv:1103.03<u>05.</u> 032005 (2012).), arXiv:1102.4585.

eNT), Phys. Rev. Lett. 106,

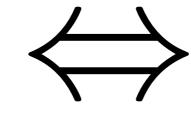
 α_s



 $\frac{(\bar{\chi}\gamma_{\mu}\chi)(\bar{q}\gamma^{\mu}q)}{\Lambda^2},$

 $(\bar{\chi}\gamma_{\mu}\gamma_{5}\chi)(\bar{q}\gamma^{\mu}\gamma_{5}q)$

 $(\bar{\chi}\chi) (G^a_{\mu\nu}G^{a\mu\nu})$



and a density of $\rho_{\chi} = 0.3 \,\text{GeV/cm}^3$. The S1 energy resolution, governed by Poisson fluctuations, is taken into account. Uncertainties in the energy scale as indicated in Fig. 1 as well as uncertainties in v_{esc} are profiled out and incorporated into the limit. The resulting 90% confidence level (CL) limit is shown in Fig. 5 and has a minimum $\sigma = 7.0 \times 10^{-45} \text{ cm}^2$ at a WIMP mass of $m_{\chi} = 50 \text{ GeV/c}^2$. The impact of \mathcal{L}_{eff} data below $3 \, \mathrm{keV}_{nr}$ is negligible at $m_{\chi} = 10 \,\mathrm{GeV/c^2}$. The sensitivity is the expected limit in absence of a signal above background and is also shown in Fig. 5 as 1σ and 2σ region. Due to the presence of two events around 30 keV_{nr} , the limit at higher m_{χ} is weaker than expected. This limit is consistent with the one from the standard analysis, which calculates the limit based only on events in the WIMP search region with an acceptance-corrected exposure, weighted with the spectrum of a $m_\chi=100\,{\rm GeV/c^2}$ WIMP, of 1471 kg × days. This result excludes a large fraction of previously unex-

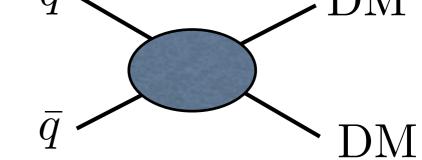
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SI, vector exchange

SD, axial-vector exchange

SI, scalar exchange

SI, scalar exchange

Typically consider each operator separately

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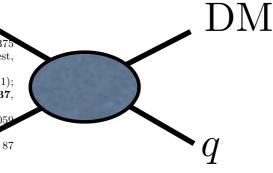
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d∕arXiv. Instrum. Meth. **Å449**, 147 hys. Lett. **B524**, 245 (2002); nys. J. direct **C3**, 11 (2001). **D72**, 072006 (2005). V. Che-

26, 58 (2006). A. Manzur 808 (2010) Lett. **97**. (2006)

), arXiv:1102.458 eNT), Phys. Rev. Lett. 106,

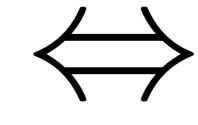
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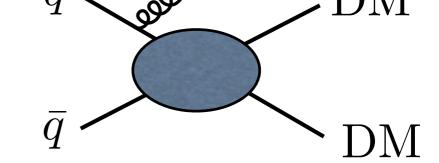
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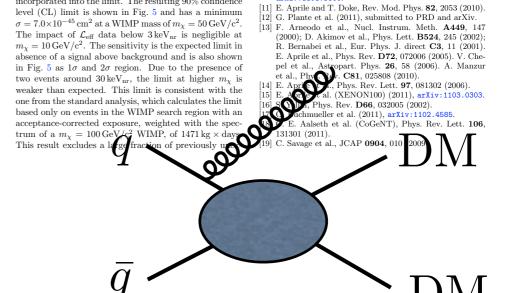
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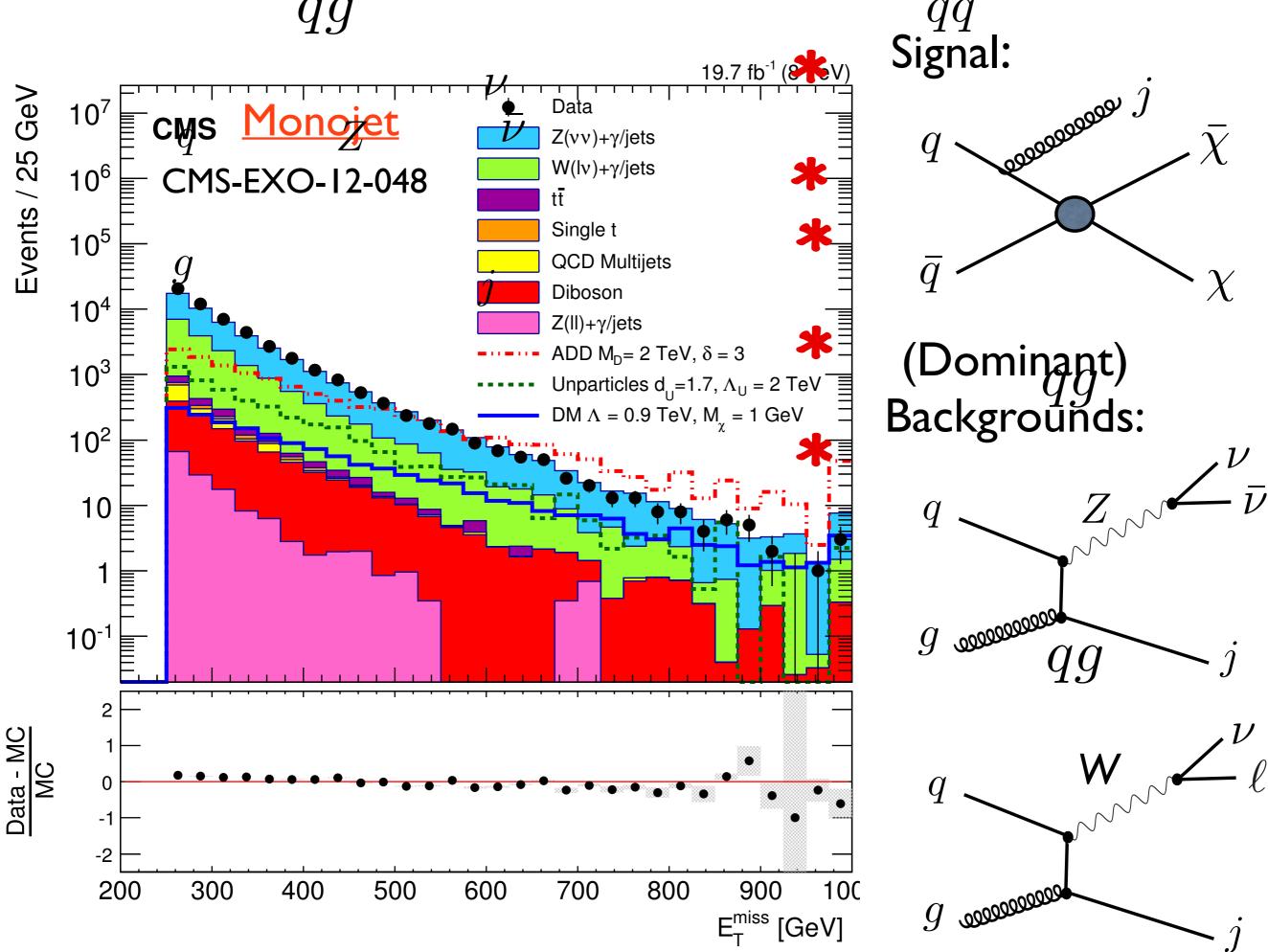
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0) (2011), arXiv:1103.5831. $\frac{(\bar{\chi}\gamma_{\mu}\chi)(\bar{q}\gamma^{\mu}q)}{\Lambda^2}$ C79, 045807 (2009). $(\underline{\bar{q}P_L}, \underline{See Goodman et al. [1008.1783]}_{for more complete list})$ 0) (2011), accepted by PRD, mitted to **RE** nd/arXiv. Instrum. Meth. **Å449**, 147 hys. Lett. **B524**, 245 (2002); nys. J. direct **C3**, 11 (2001). **D72**, 072006 (2005). V. Che- $(\bar{\chi}\gamma_{\mu}\gamma_{5}\chi)(\bar{q}\gamma^{\mu}\gamma_{5}q)$ 26, 58 (2006). A. Manzur 808 (2010) Lett. 97. 032005 (2)), arXiv:1102.458 eNT), Phys. Rev. Lett. 106,

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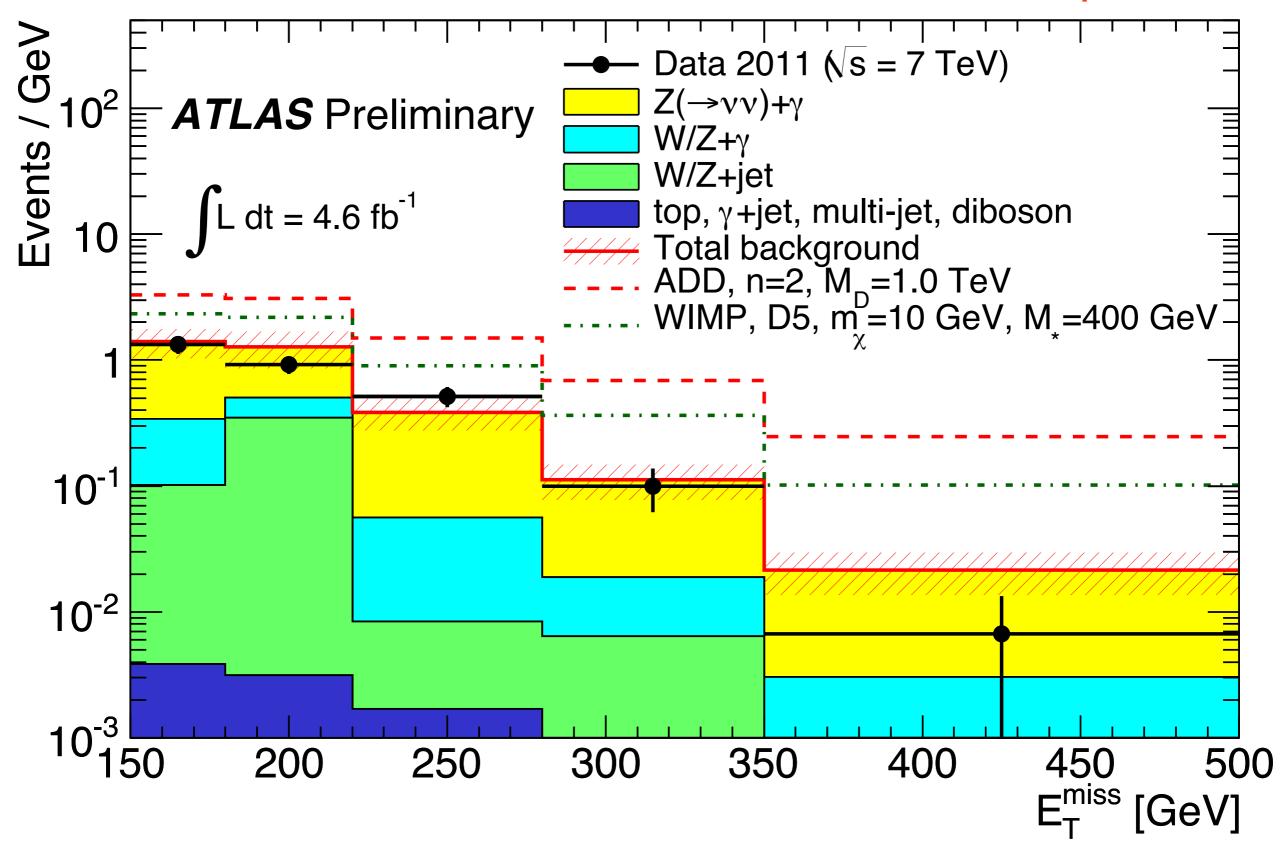
 $(G^a_{\mu\nu}G^{a\mu\nu})$



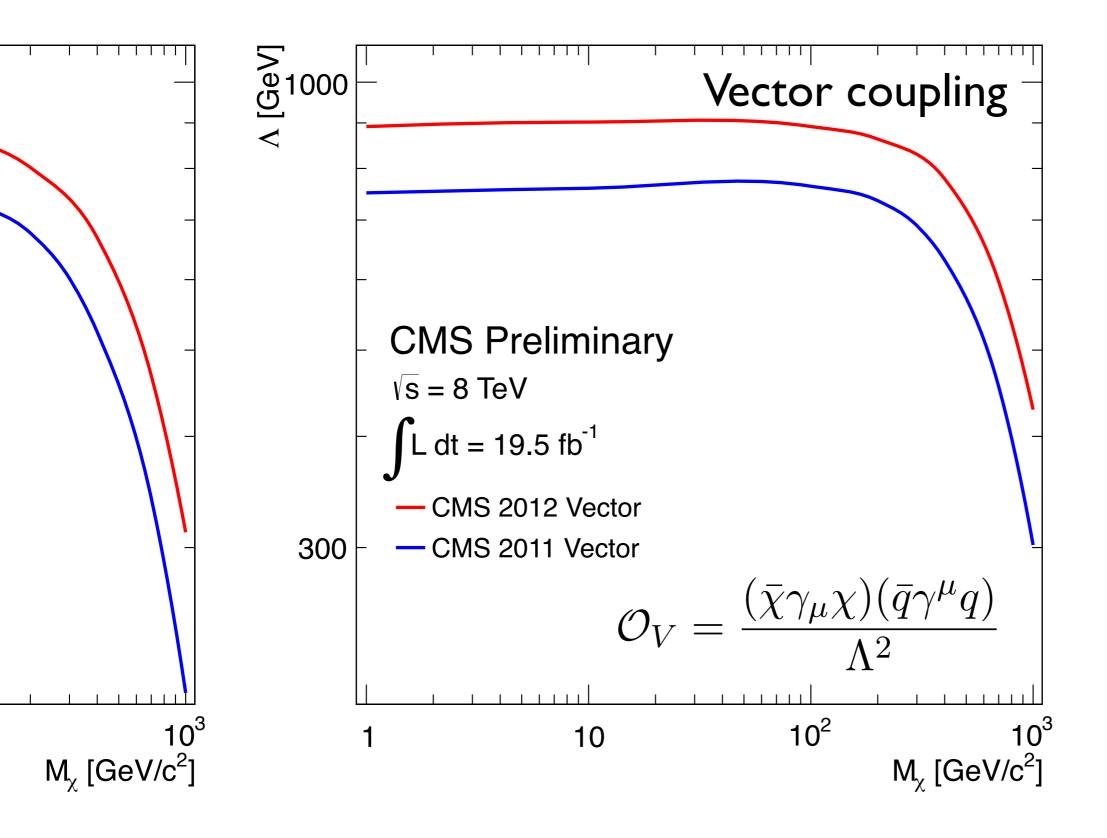


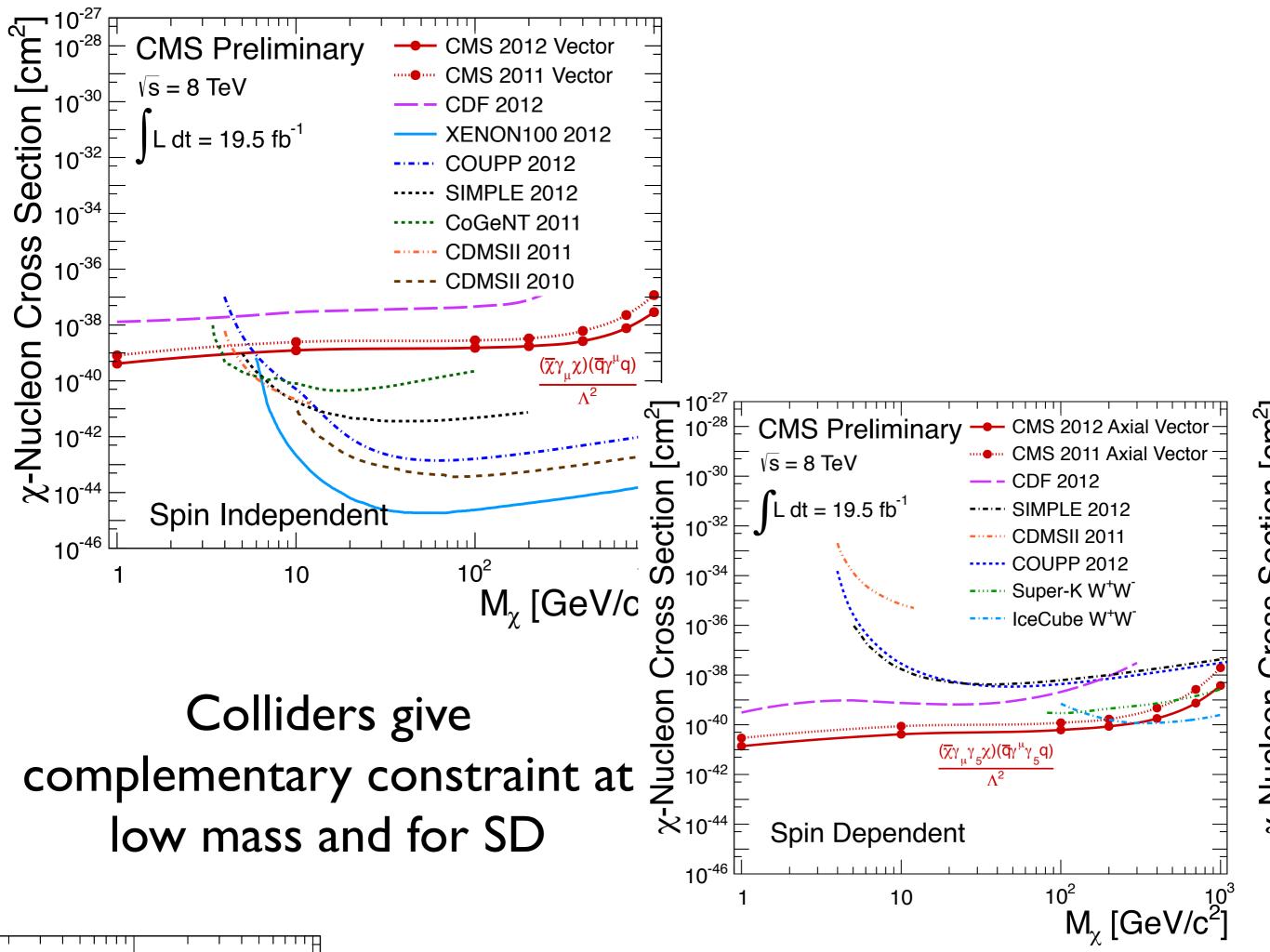
ATLAS-CONF-2012-085

Monophoton



How to quantify nothing?



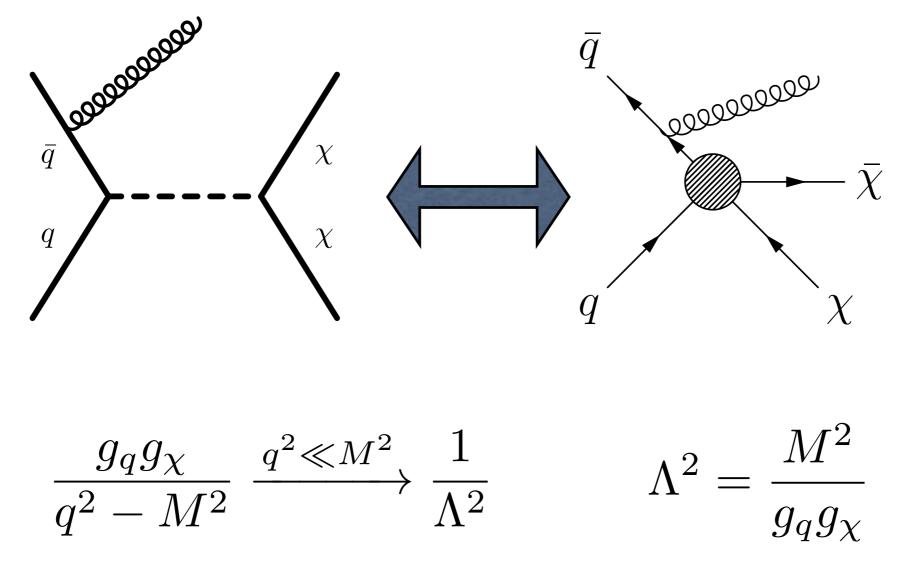


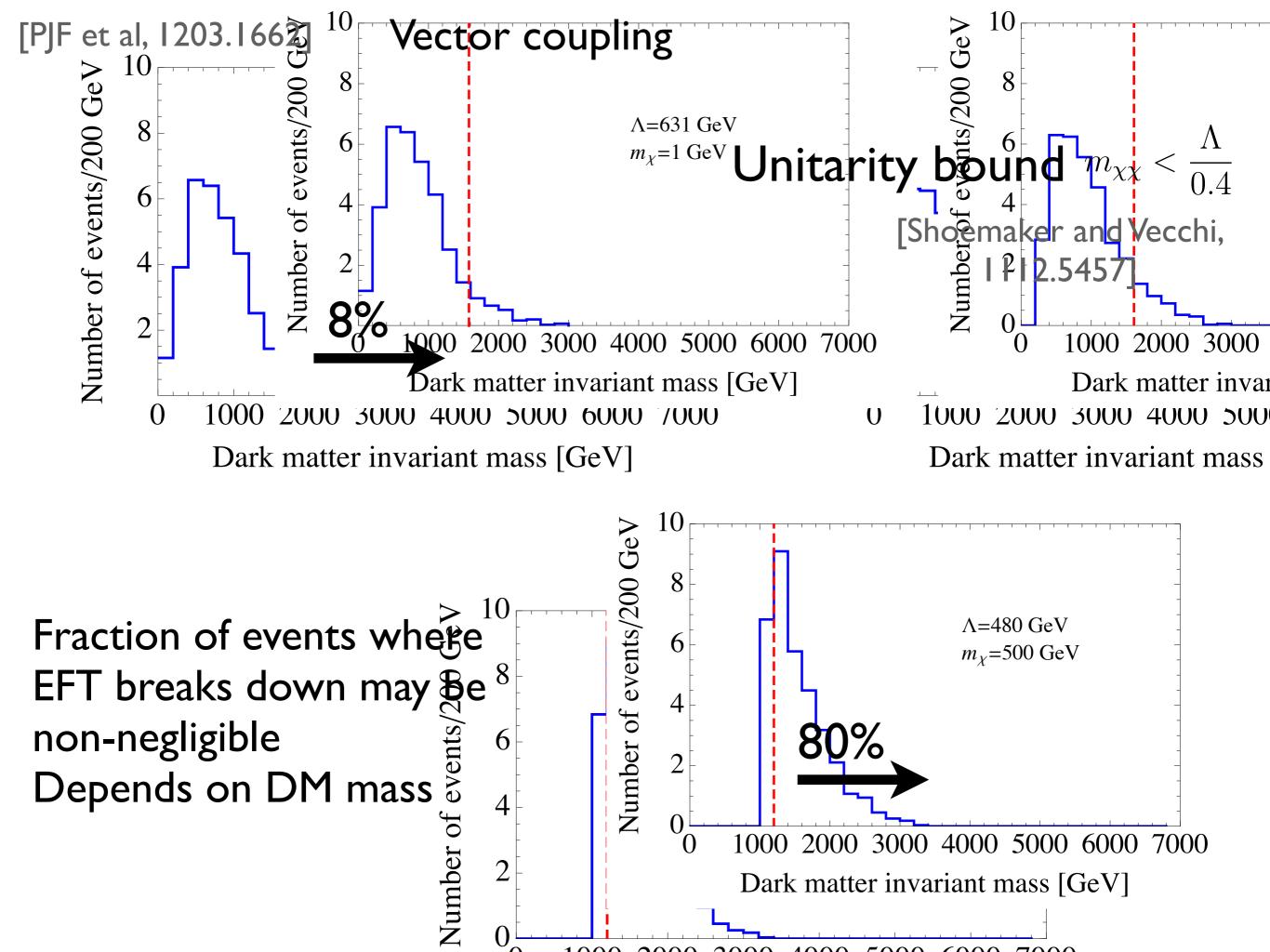
Light Mediators

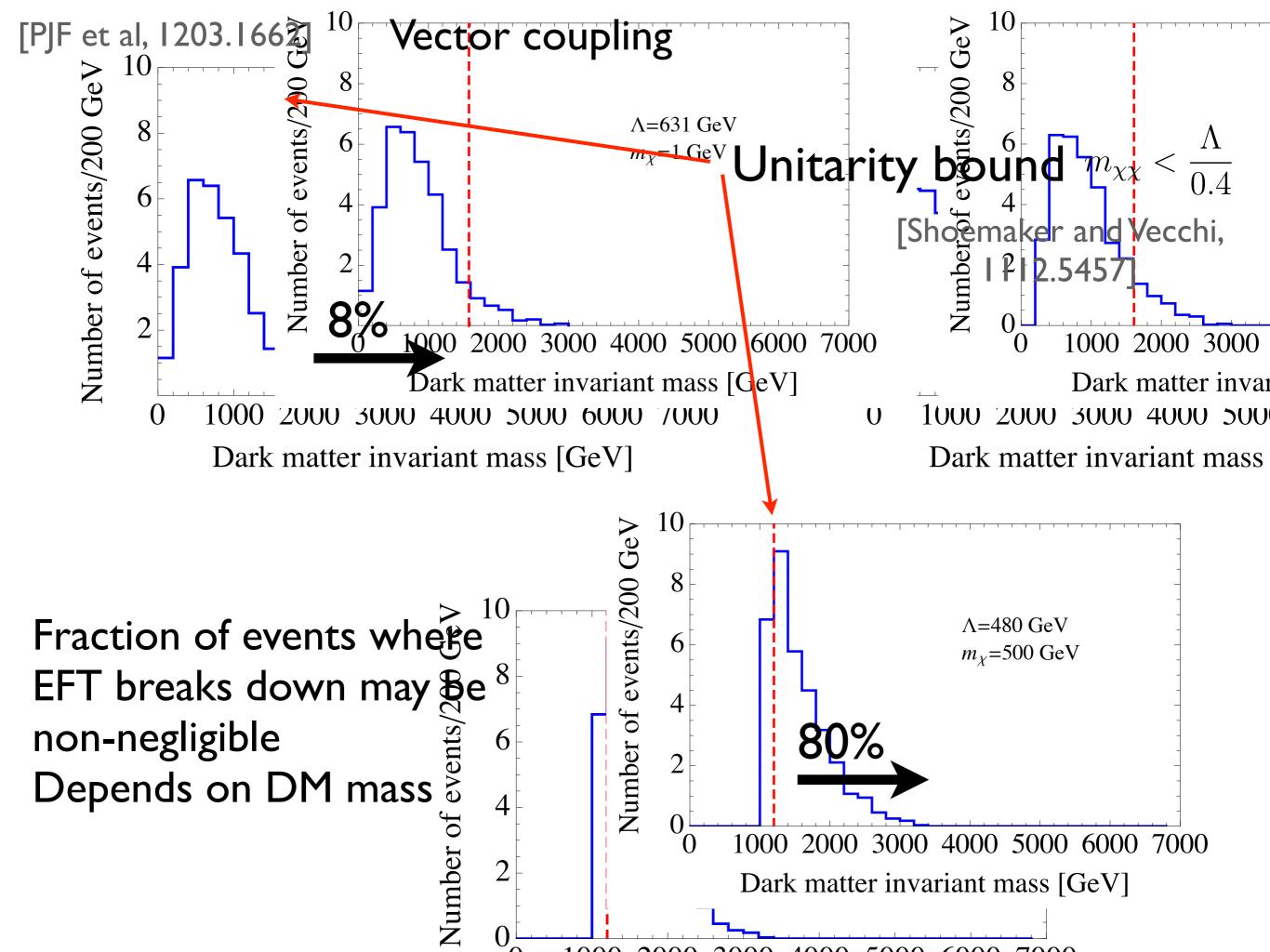
For all but the lightest mediators EFT is good for direct detection

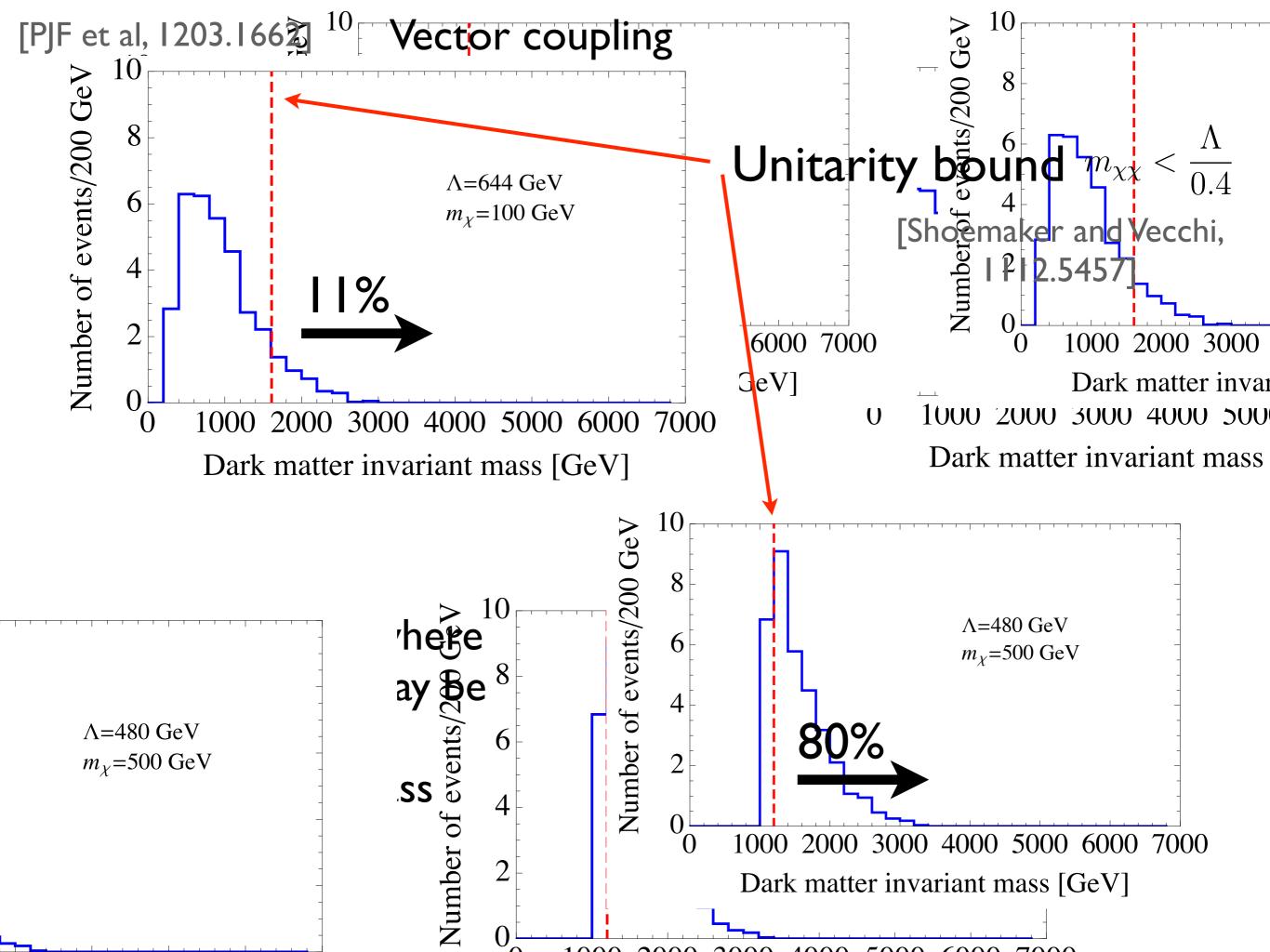
$$\sigma(\chi N \to \chi N) \sim \frac{g_q^2 g_\chi^2}{M^4} \mu_{\chi N}^2$$

What fraction of collider events have momentum transfers sufficient to probe the UV completion?

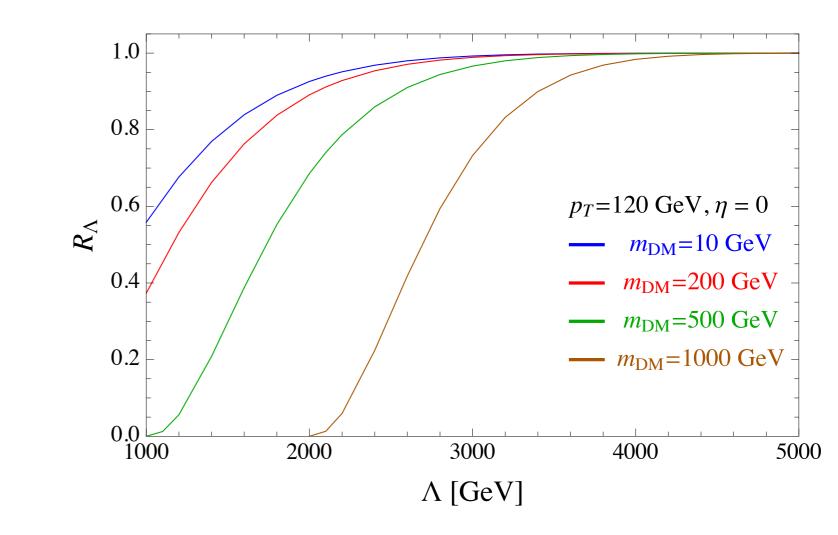


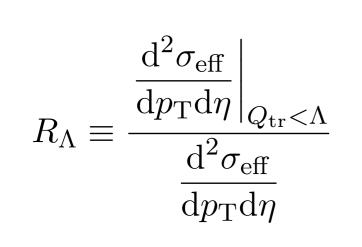


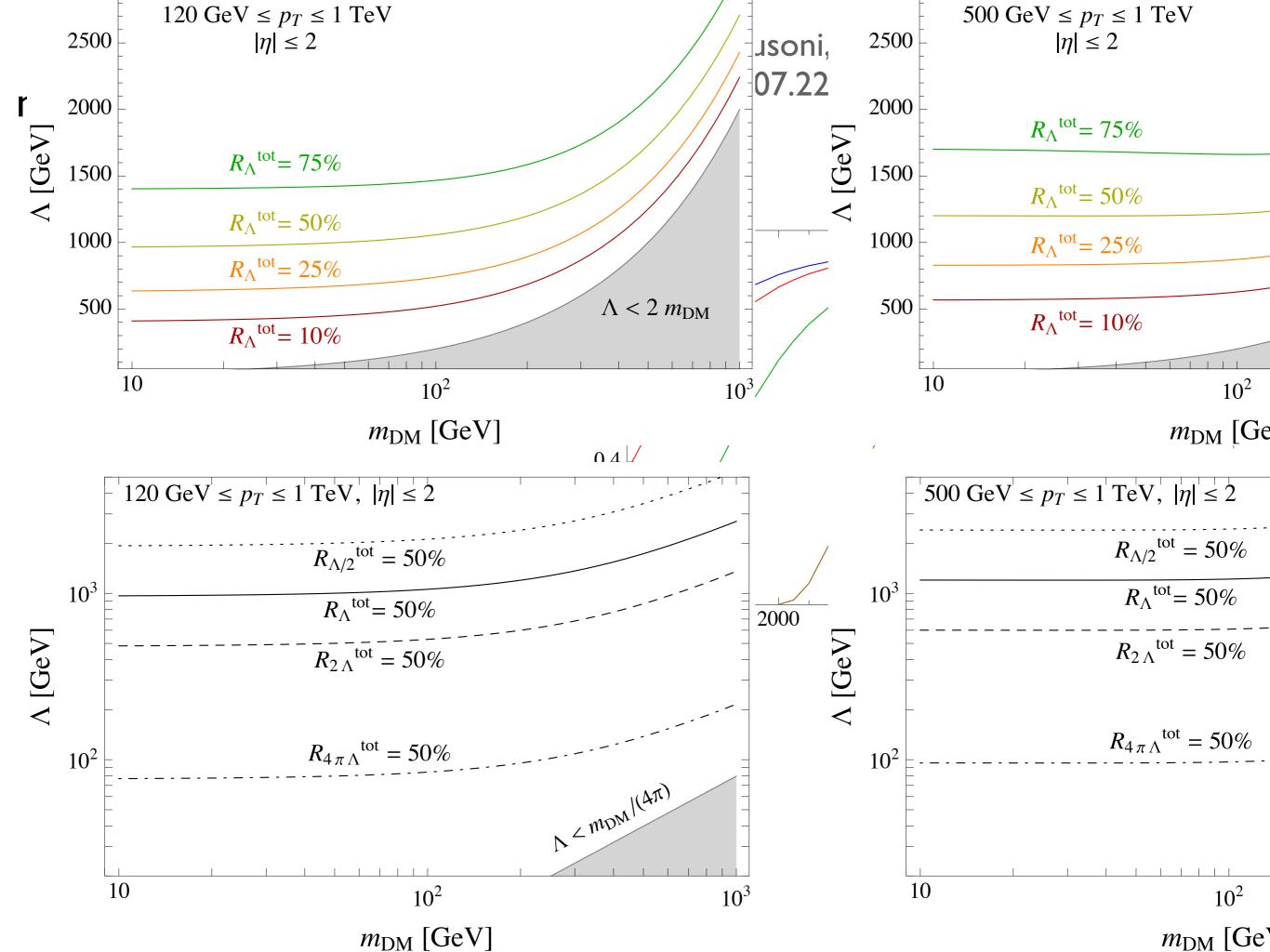


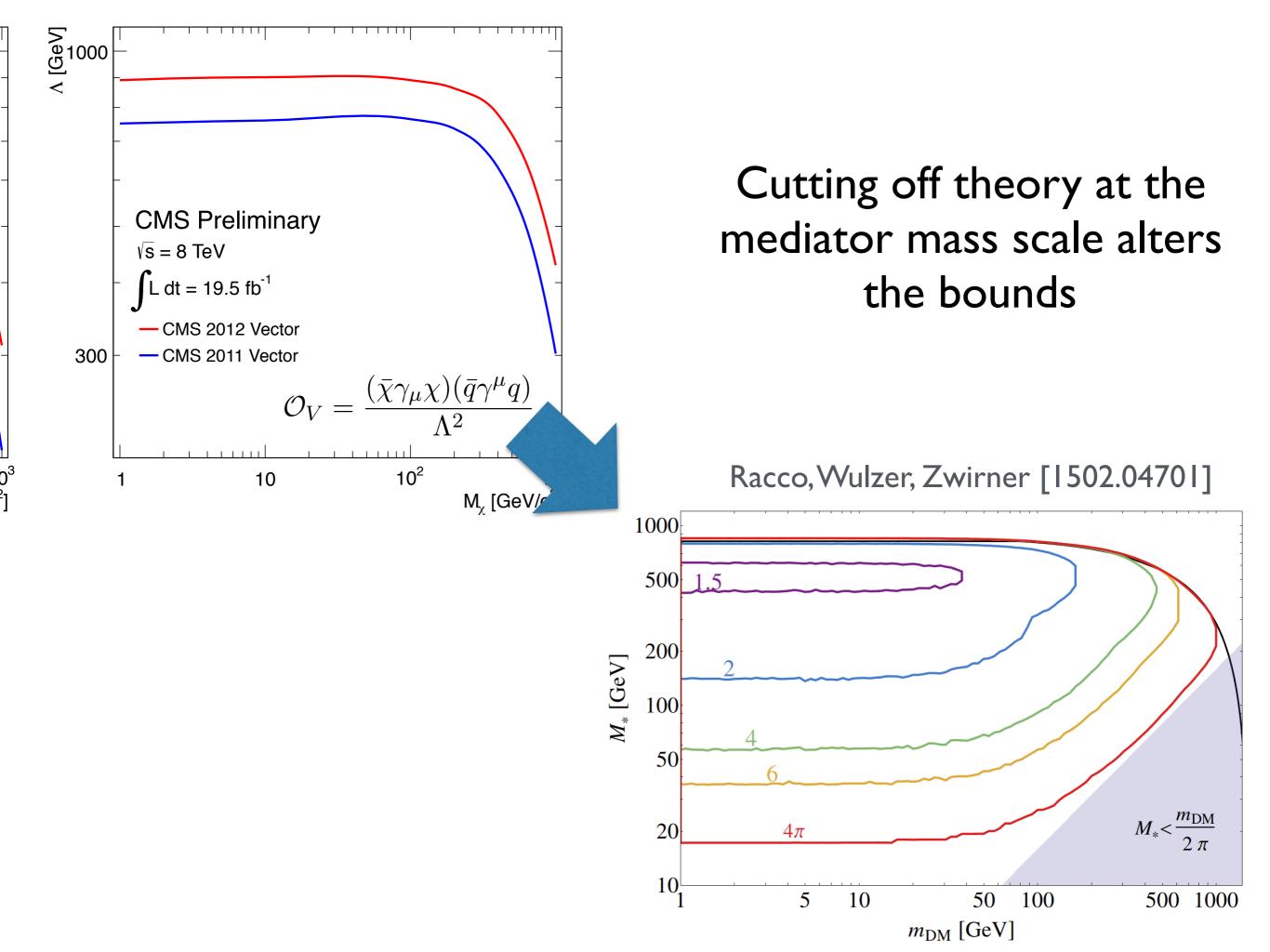


What fraction of events have momentum transfers sufficient to probe the UV completion? [Busoni, De Simone, Morgante, Riotto, 1307.2253, 1402.1275, 1405.3103]

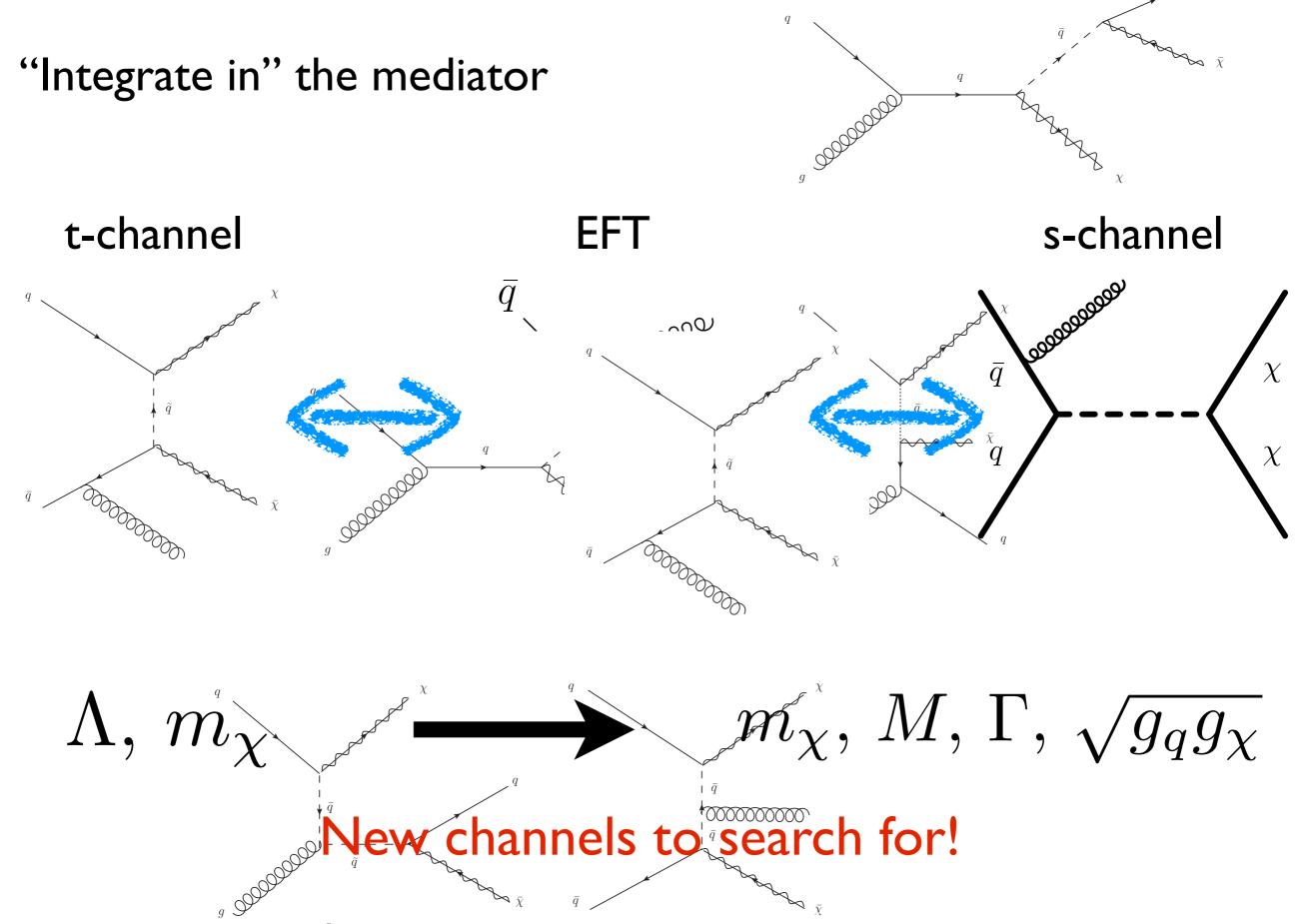






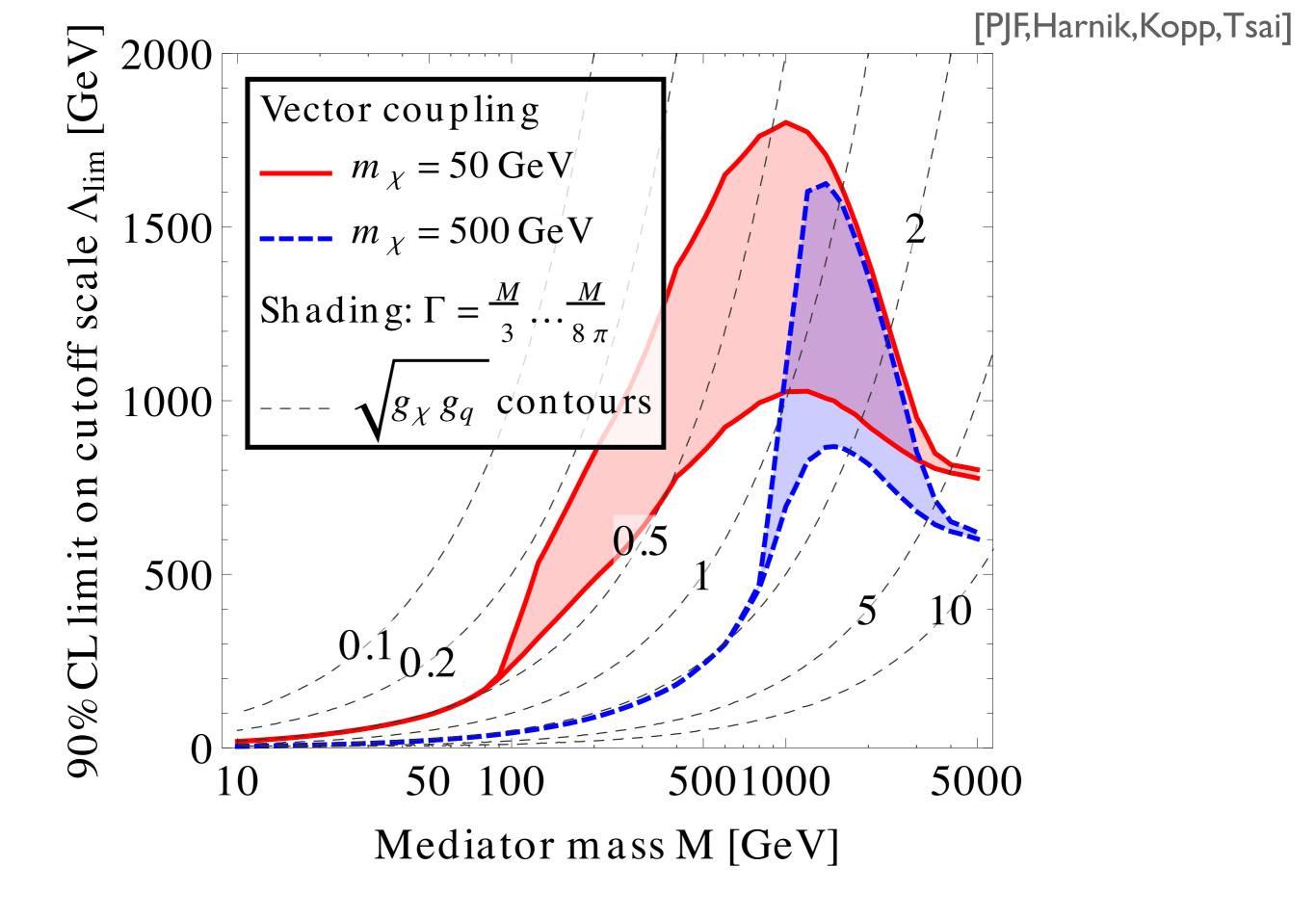


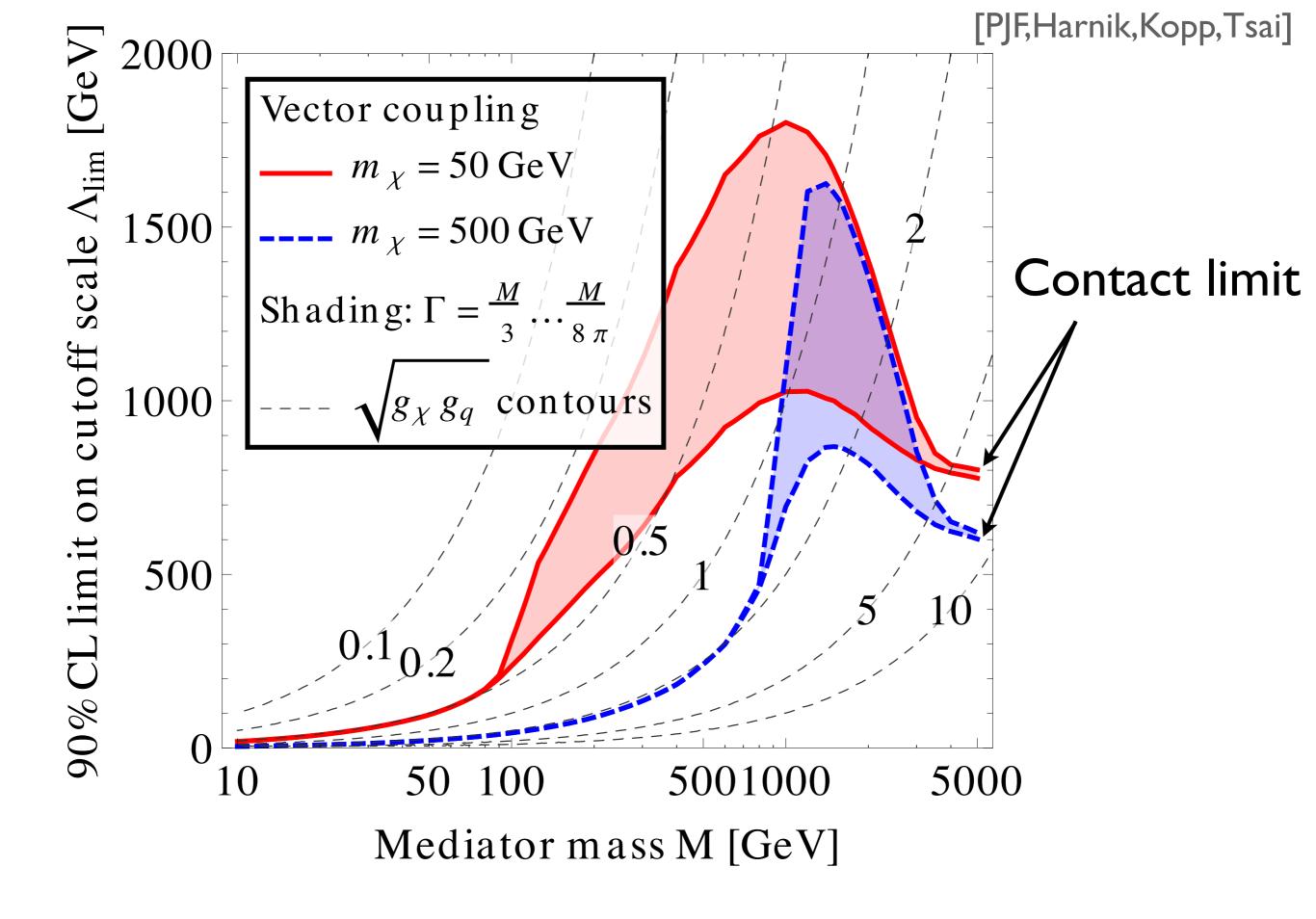
Simplified Models

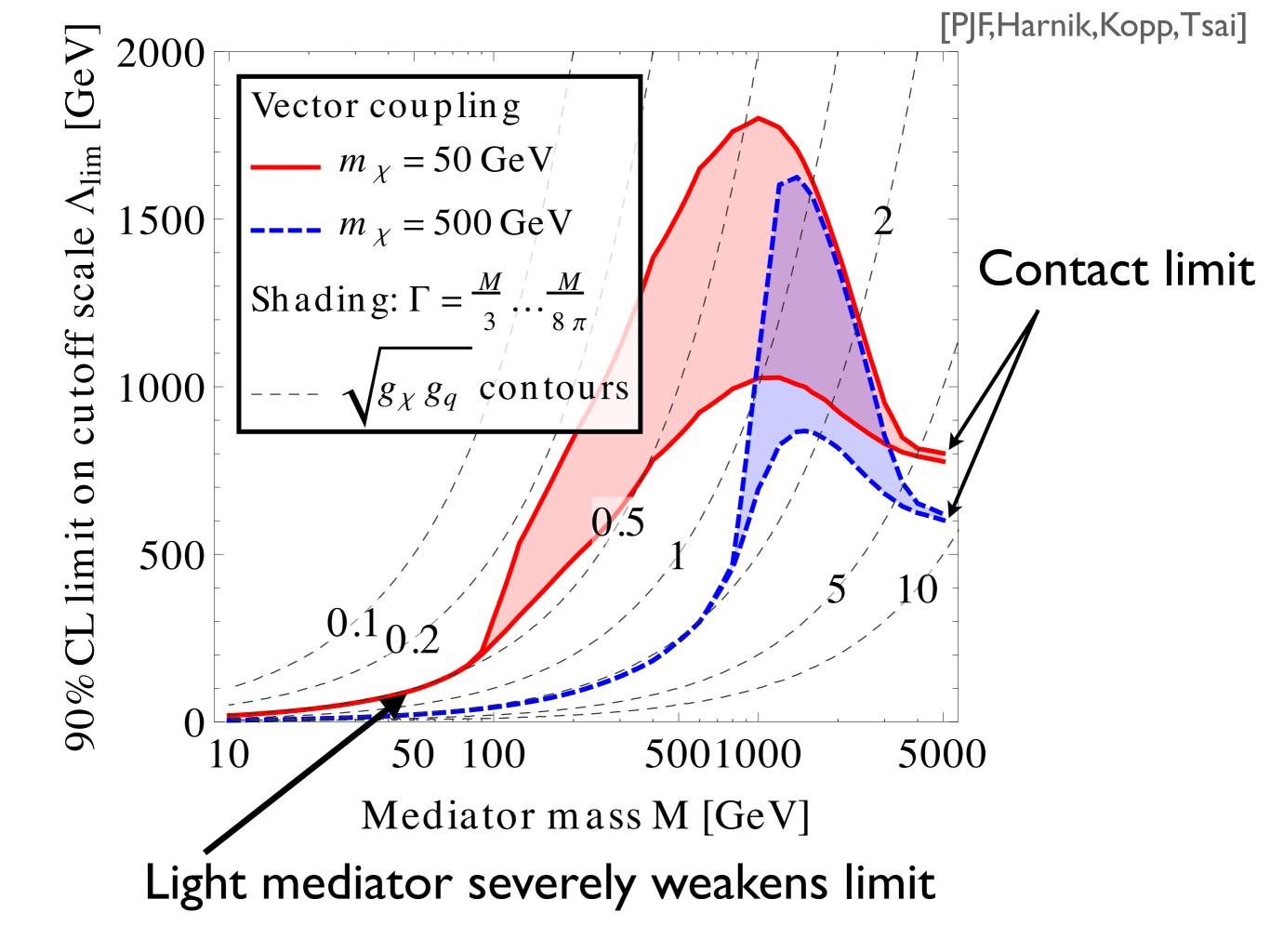


Collider only sensitive to all 4 parameters over a narrow range

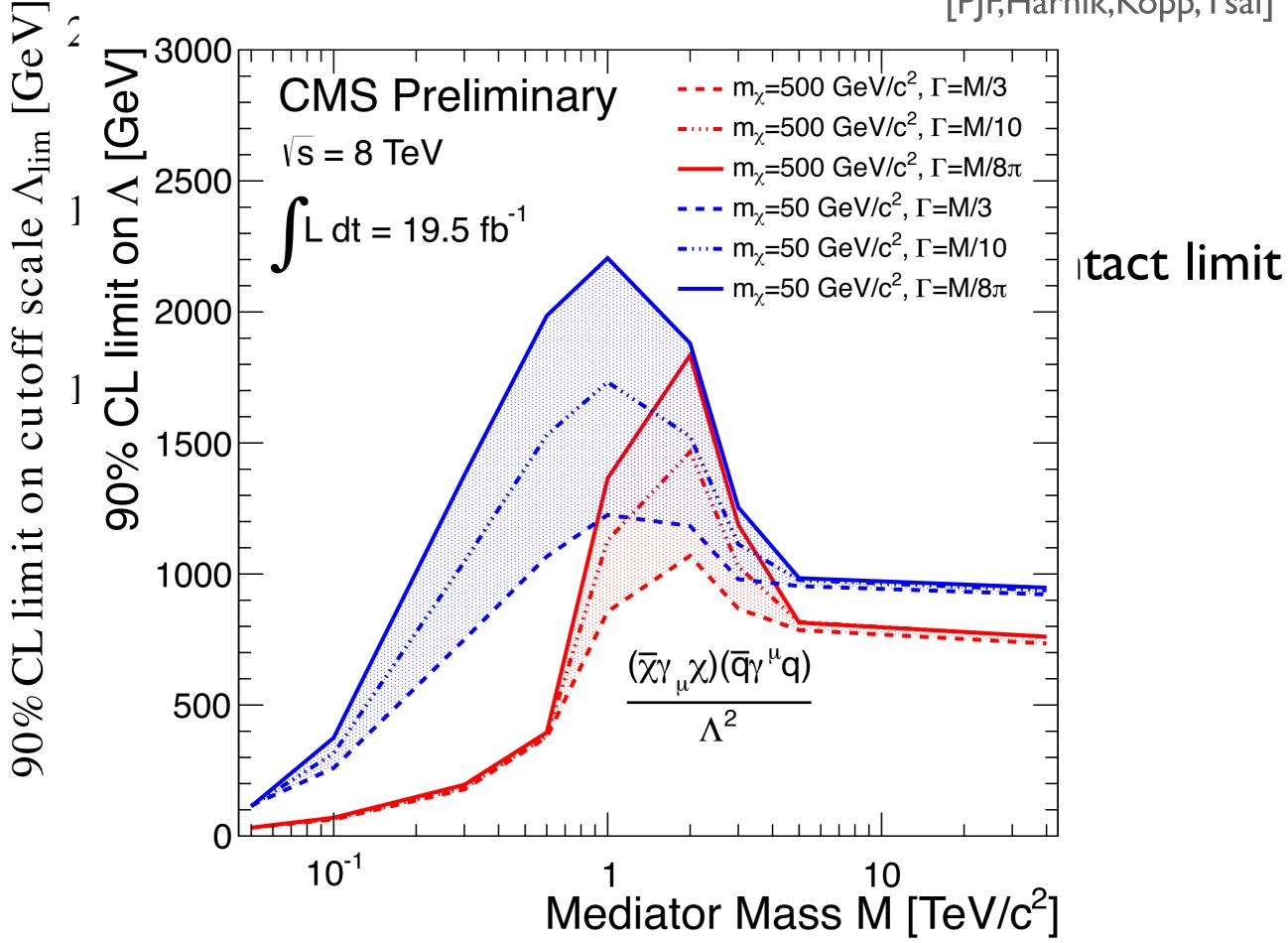
But mapping collider constraints to direct/indirect detection now requires assumptions







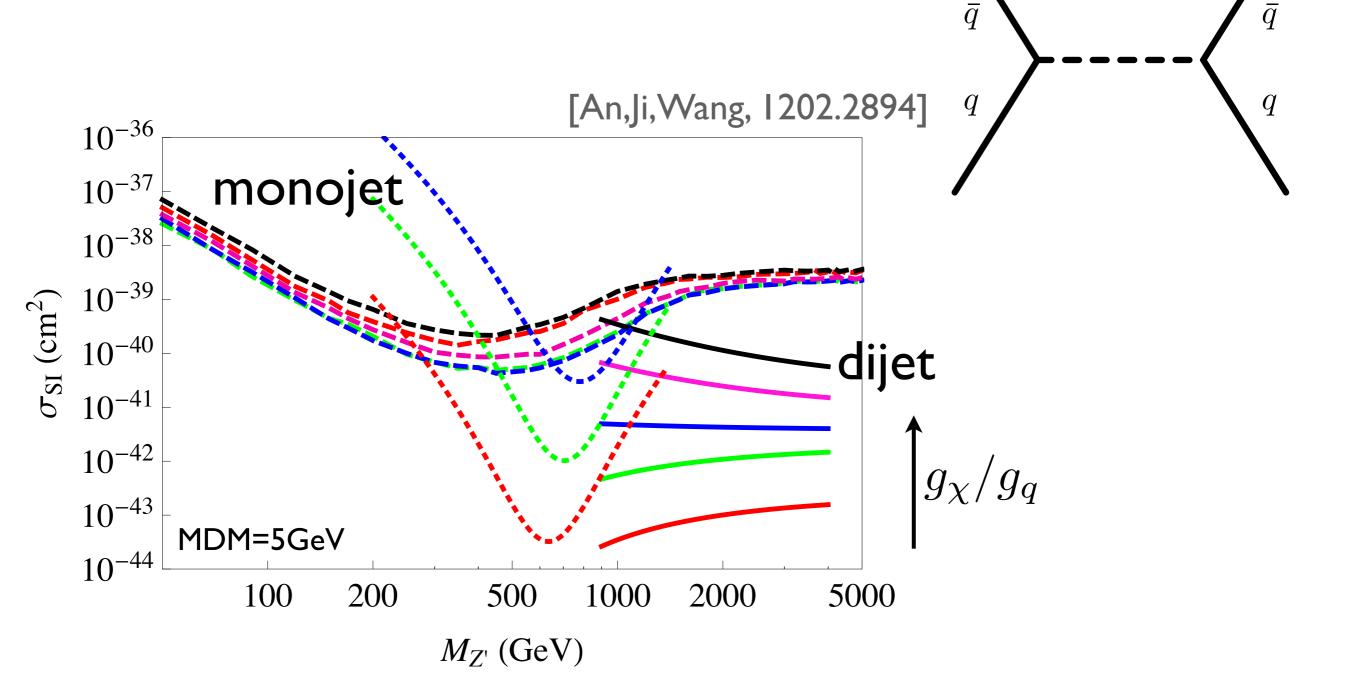
[PJF, Harnik, Kopp, Tsai]



Light Mediators

[An,Ji,Wang:1202.2894;March-Russell, Unwin,West: 1203.4854]

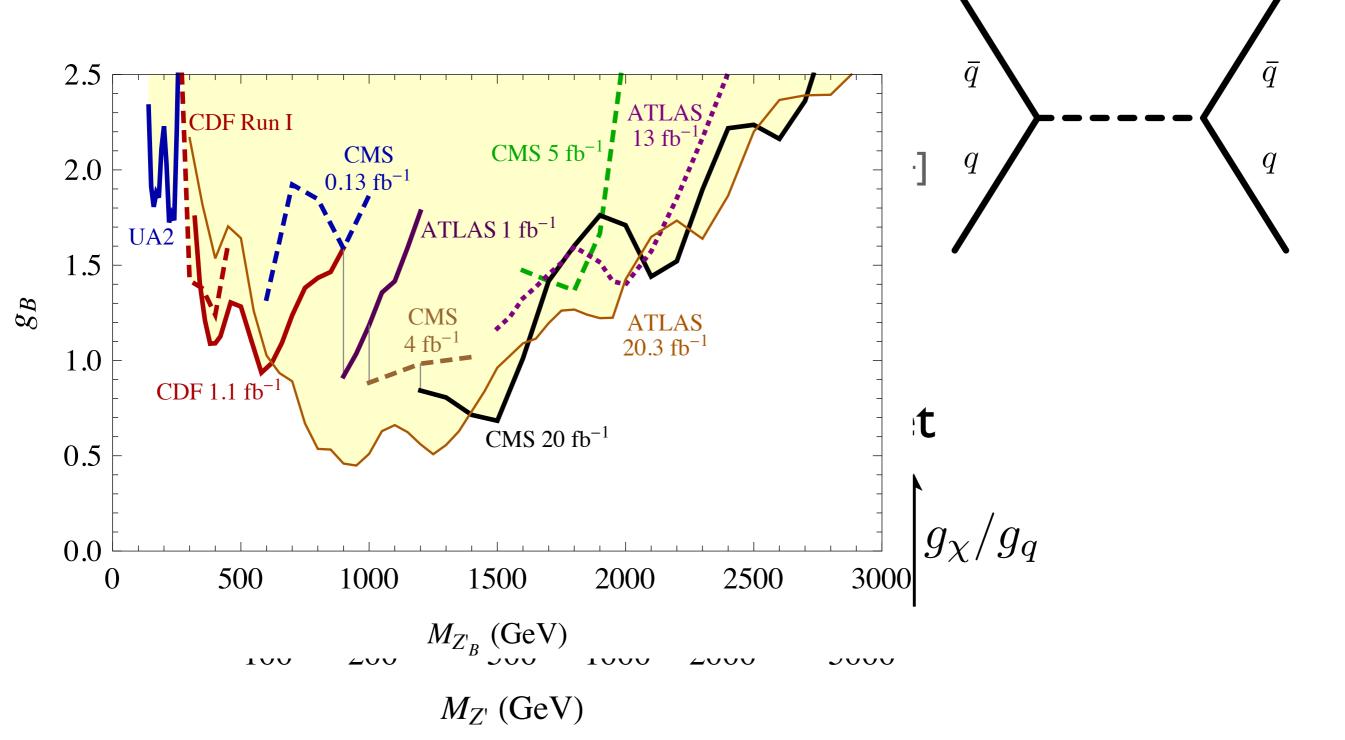
Look for the light mediator directly-dijet resonance/angular distributions



Light Mediators

[An,Ji,Wang:1202.2894;March-Russell, Unwin,West: 1203.4854]

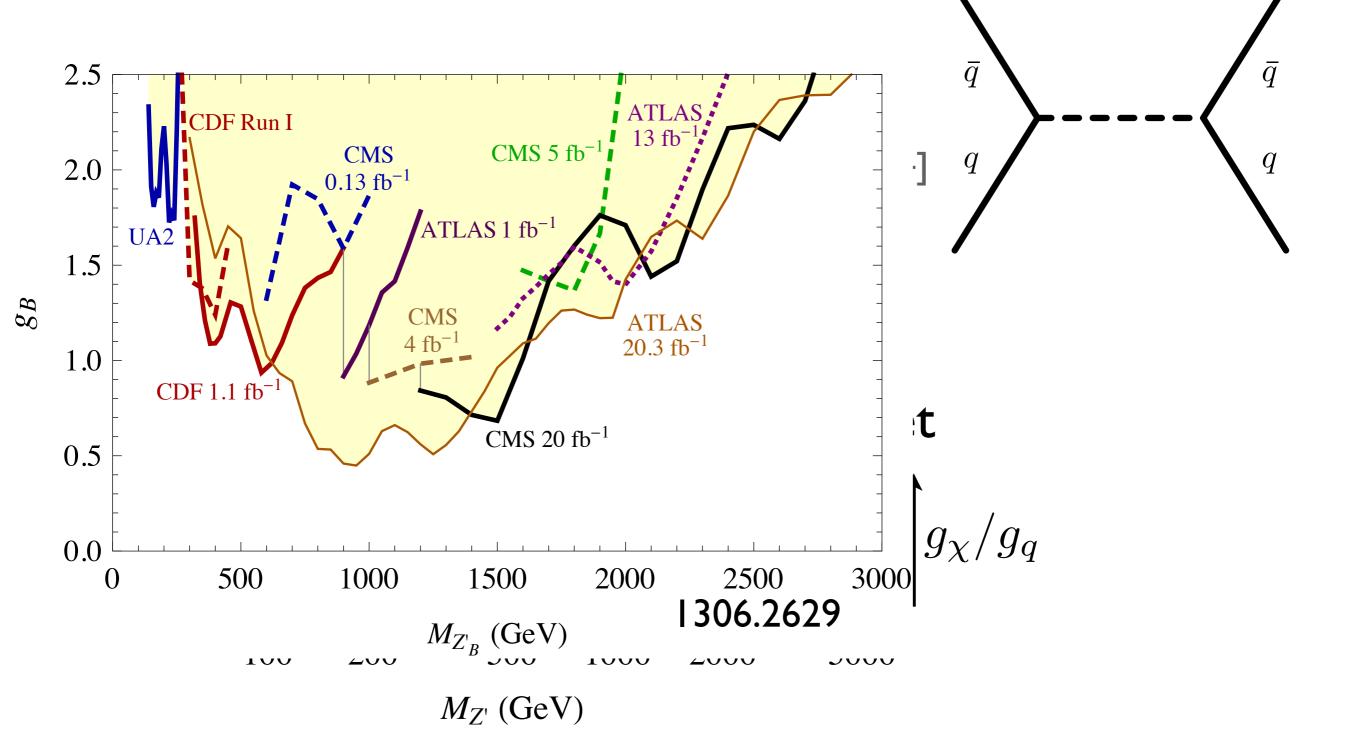
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Light Mediators

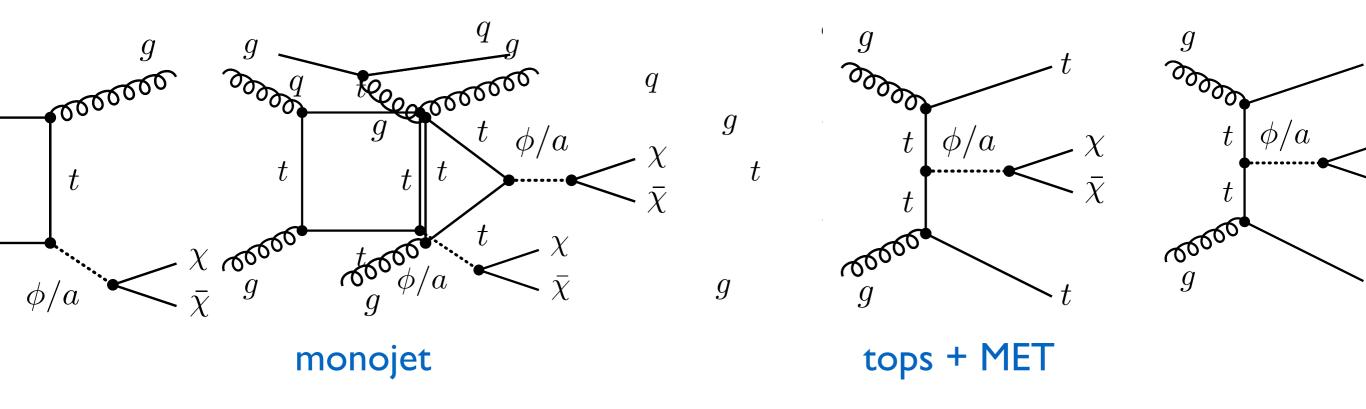
[An,Ji,Wang:1202.2894;March-Russell, Unwin,West: 1203.4854]

Look for the light mediator directly-dijet resonance/angular distributions



-channel scalar/psuedo-scalar

MFV: $\lambda_{\chi}\phi\bar{\chi}\chi + \lambda_U\phi\left(Y_U^{ij}Q_iHU_j^c\right)$ **Physics dominated by top**

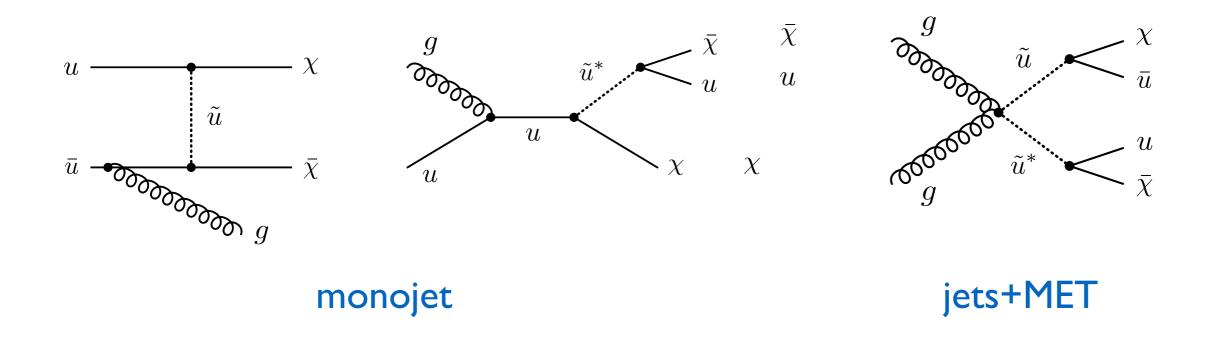


- Scalars have helicity suppressed annihilation, and SI DD
- Pseudo scalars do not, and have SD momentum suppressed DD

-channel scalar/psuedo-scalar

MFV requires DM or mediator to carry flavour $\lambda \phi_i \bar{\chi} q_i$

(Like in SUSY MFV allows for separation of 1,2 from 3 gen.)

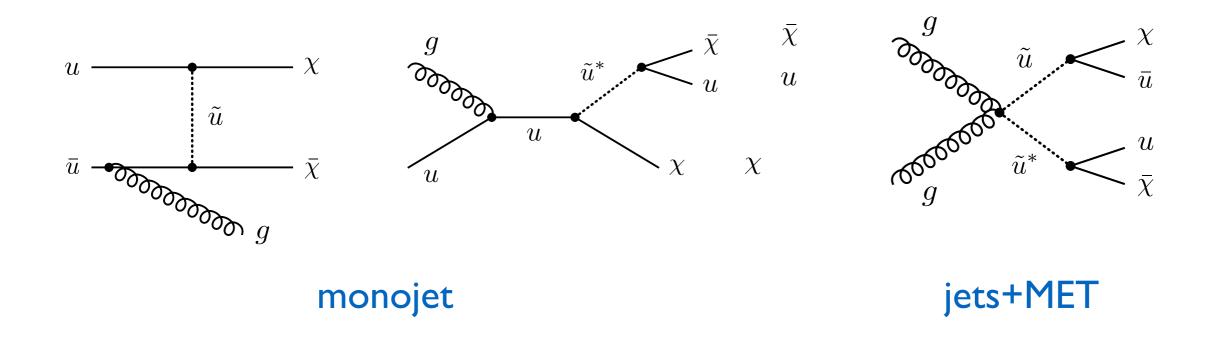


Majorana has only SD, Dirac has both Dirac cannot be a thermal relic, Majorana can if > 100 GeV

-channel scalar/psuedo-scalar

"squarks" who SUSY prior "squarks" $\lambda \phi_i \overline{\chi} \alpha$: MFV requires DM or mediator to c

(Like in SUSY MFV allows for separation of 1,2 from 3 gen.)



Majorana has only SD, Dirac has both Dirac cannot be a thermal relic, Majorana can if > 100 GeV

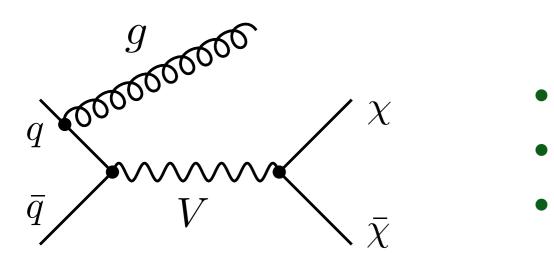
s-channel vector/axial-scalar

Spontaneously broken U(1)' accessible, can alter physics)

Consistency of model? How does DM get mass, anomalies...

$m_{\chi} \lesssim \frac{\sqrt{4\pi}}{g_{\chi}^{A}} M_{V}$

Bounds on dileptons, leptophobic Z'



$$\begin{array}{ccc} g & g \\ \bullet \operatorname{Vectors} \operatorname{are} \operatorname{SI} & \chi & q \\ \bullet \operatorname{Axial} \operatorname{vectors} \operatorname{SD} & q \\ \bullet \operatorname{If} \operatorname{thermal} \operatorname{often} \operatorname{underproduc}_{\chi} \operatorname{cd}_{\chi} \\ \overline{q} \end{array}$$

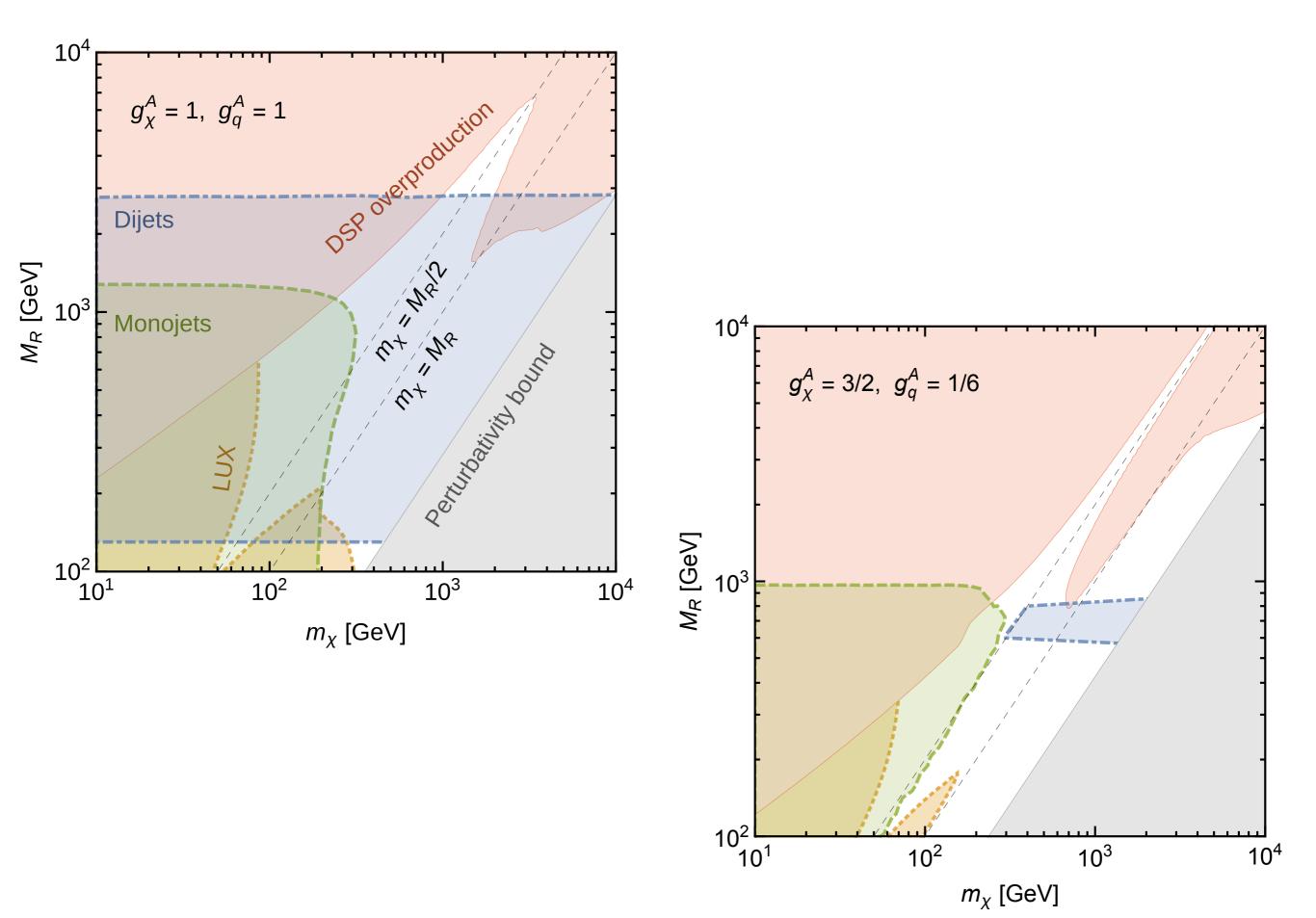
(Higgs mode may be

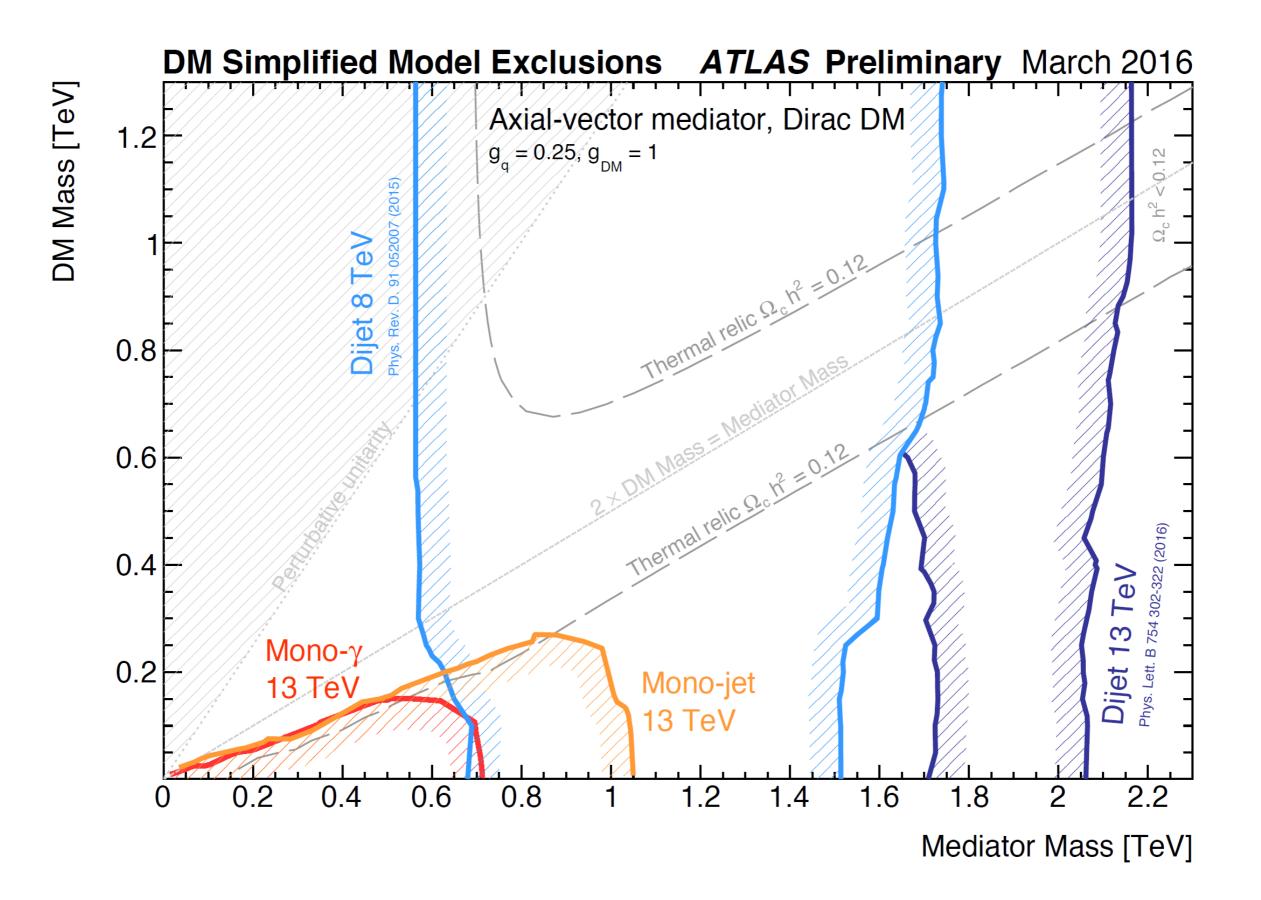
monojet

V

- Landscape of simplified models is broad and varied
- Spin/parity of DM and mediator
- MFV
- Kinetic mixing
- Higgs portal
- Vector DM
- •Other dark sector states alter thermal history & BRs
- •Electroweak-inos, singlet-doublet DM, etc

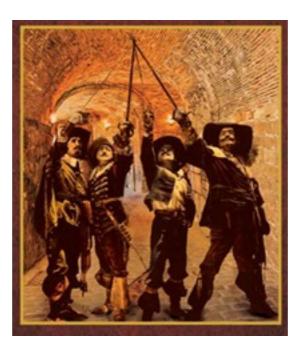
[Chala, Kahlhoefer, McCullough, Nardini, Schmidt-Hoberg]





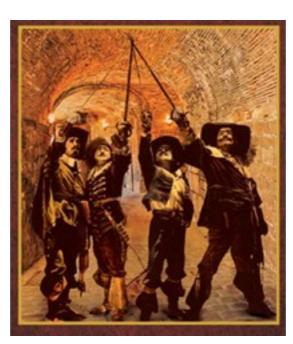
Complementarity

- Direct detection limited to DM above GeV, needs DM nearby moving in the right way
- No upper limit on mass probed, learn about DM in cosmos
- Indirect detection very sensitive to astrophysics
- Halo shapes can probe DM-DM interactions
- Collider searches have kinematic upper limit, no astrophysics systematics, but many others
- Complementary taken together provide complete picture



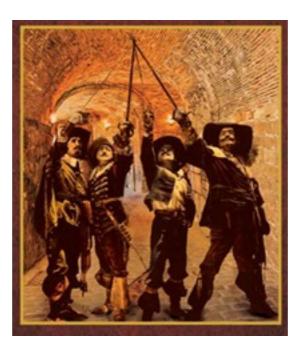
Complementarity

- Direct detection line
 Many exciting new ideas for probing light PM e.g. scattering off electrons in semi/super conductors
 probed, learn about DM in cosmos
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Complementarity

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"If you like laws and sausages, you should never watch either one being made"

Otto von Bismark





Clever field theory idea, cute new symmetry, deep new underlying principle

Clever field theory idea, cute new symmetry, deep new underlying principle

or

Clever field theory idea, cute new symmetry, deep new underlying principle

or

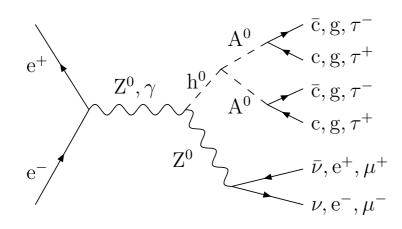
New data needs explaining, signal not being searched for

"Top down"

Identify "grand problem" e.g. weak hierarchy, cosmological constant, flavour
Introduce "grand principle" e.g. extra dimensions, supersymmetry, new strong dynamics
Define new theory obeying principle that has SM as long energy limit

Outcome: theoretically very appealing model, often highly correlated signals, complicated parameter space

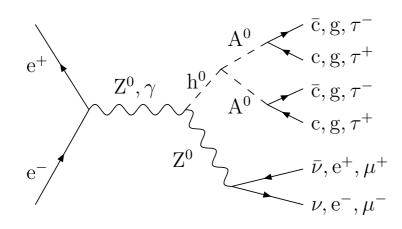
A cautionary tale



OPAL Higgs search $m_A [GeV/c^2]$ theoretically inaccessible theoretically inaccessible excluded by excluded LEP1 searches by OPAL 80 90 100 $m_{h} [GeV/c^{2}]$

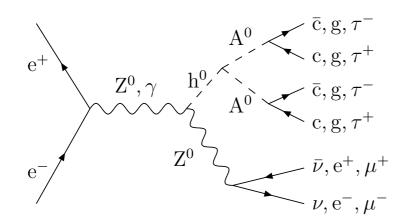
Eur.Phys.J.C27:483-495,2003.

A cautionary tale



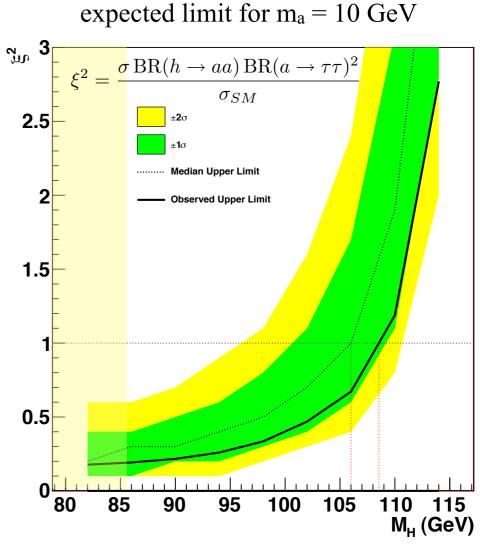
OPAL Higgs search $m_A [GeV/c^2]$ theoretically inaccessible theoretically inaccessible excluded by excluded LEP1 searches by OPAL 80 90 100 $m_{h} [GeV/c^{2}]$ Eur.Phys.J.C27:483-495,2003.

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OPAL Higgs search $m_A [GeV/c^2]$ theoretically inaccessible theoretically inaccessible excluded by excluded LEP1 searches by OPAL $m_h [GeV/c^2]$ Eur.Phys.J.C27:483-495,2003.

New ALEPH search



Cranmer, Yavin, Beacham, Spagnolo

"Bottom up"

- •Data disagrees with SM in some channel(s)
- Add new states and couplings to SM to explain deviations
- •Must have some concept of minimality: degrees of freedom, parameters
- Outcome: build up the new physics piece by piece, correlations may not be apparent initially, simple parameter space
 - Easy for us to talk...exchange MadGraph/SHERPA model files that contain a few dials

"Bottom up".... without anomaly

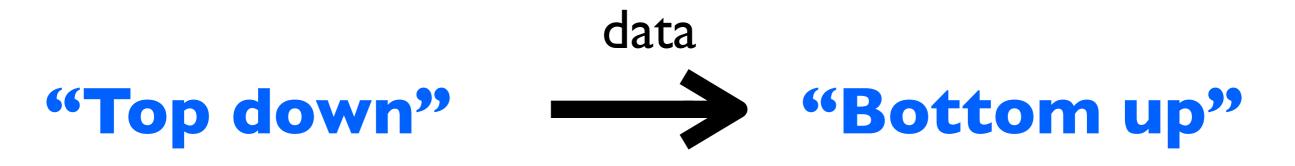
Bottom up without excess = "signal building"

- Build simple modules that contain interesting new signatures not necessarily contained in other models
- Motivate new analyses
- Again allows simple communication

"Bottom up".... without anomaly

Bottom up without excess = "signal building"

- Build simple modules that contain interesting new signatures not necessarily contained in other models
- Motivate new analyses
- Again allows simple communication



Rules of model building

- "First do no harm"
 - •FCNC's, PEWT, LEP, B-physics, proton decay, existing searches,.. (often reason for new parity...DM)
- •Describe physics with a local, Lorentz invariant, unitary field theory, causal
- •Preserve gauge invariance, anomaly free
- •Prefer renormalizable field theories
- •Occam's razor? cf. Hickam's Dictum
- Perturbativity
- •Running of gauge couplings, unification