
Reconstruction

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Hadron Collider Physics Summer School 2018

You've heard excellent lectures on theory, experimental measurements and searches, and detector technologies

Reconstruction: algorithms to select/combine detector signals into representative physics observables for experimental analysis

Drawing off of other lectures:

Silicon Detectors, Calorimetry, Machine Learning, Heavy Ions, Precision Measurements, Fast Timing

Caveat 1: my experience is in ATLAS/CMS style reconstruction, so I will focus on that, with a few special topics for heavy ions and b-physics.

Caveat 2: More CMS results mostly because I know where to find those plots more easily — but most everything I will say will be generic

Part 1: Building blocks

- a. Charged particle tracking, vertexing
- b. Precision Timing
- c. Calorimetry

Part 2: Particle reconstruction

- a. Muons
- b. Photons/Electrons
- c. Taus, Hadrons
- d. special topic: LHCb RICH detector
- e. Particle Flow

Part 3: Composite objects and beyond

- a. Jets, MET
- b. Jet substructure
- c. Pileup Mitigation
 - c.ii. special topic: Underlying event in heavy ions
- d. Displaced/Exotic objects

A LOT OF GROUND TO COVER!
MY STRATEGY: GIVE YOU AN IDEA
OF MANY THINGS RATHER THAN
FOCUS ON A FEW

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Lecture 1

Lecture 2

Part 3: Composite objects and beyond

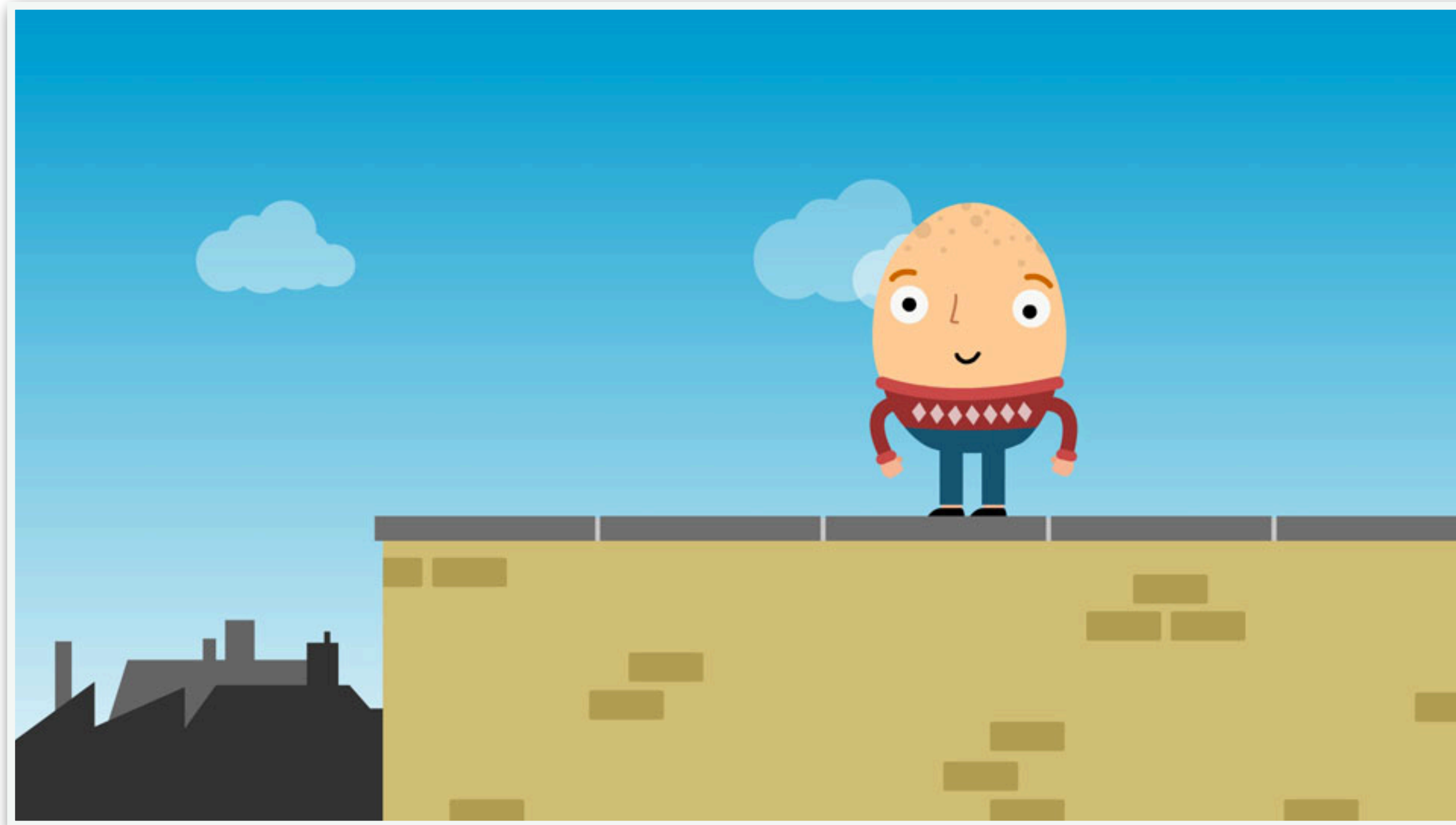
- a. Jets, MET
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I'm drawing a lot from
different sources, but
great references are
lectures from previous
years of HCPSS.

Special credit to my
predecessors at the
FNAL school: Phil
Harris and Rick
Cavanaugh

INTRODUCTION



Collision

Physics process

Partons

Stable particles

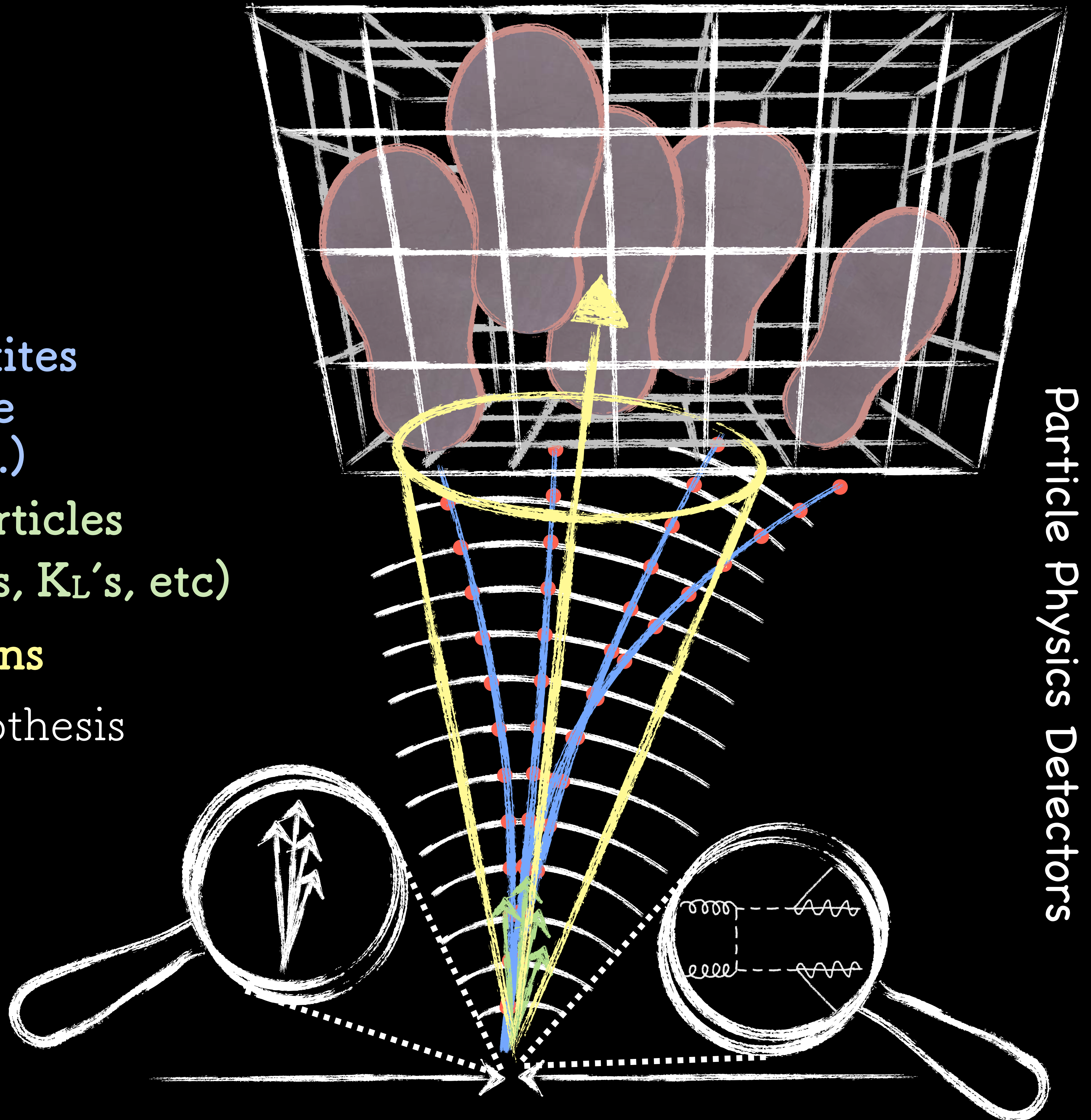
Detector hits

Reconstructed quantities
(momenta, charge
energy, angles, ...)

List of ID'd reco. particles
(e 's, μ 's, γ 's, π 's, K_L 's, etc)

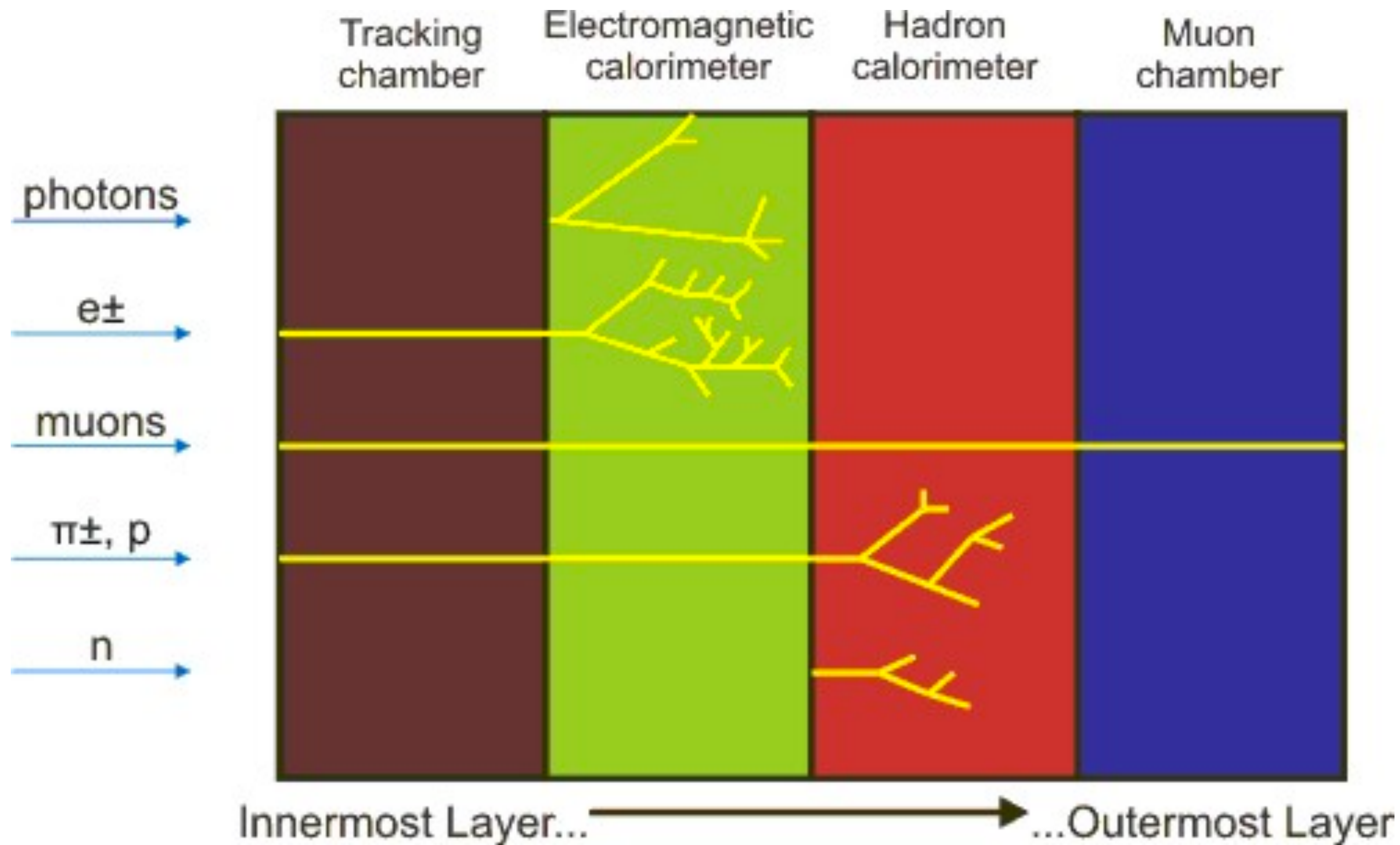
Reconstructed partons

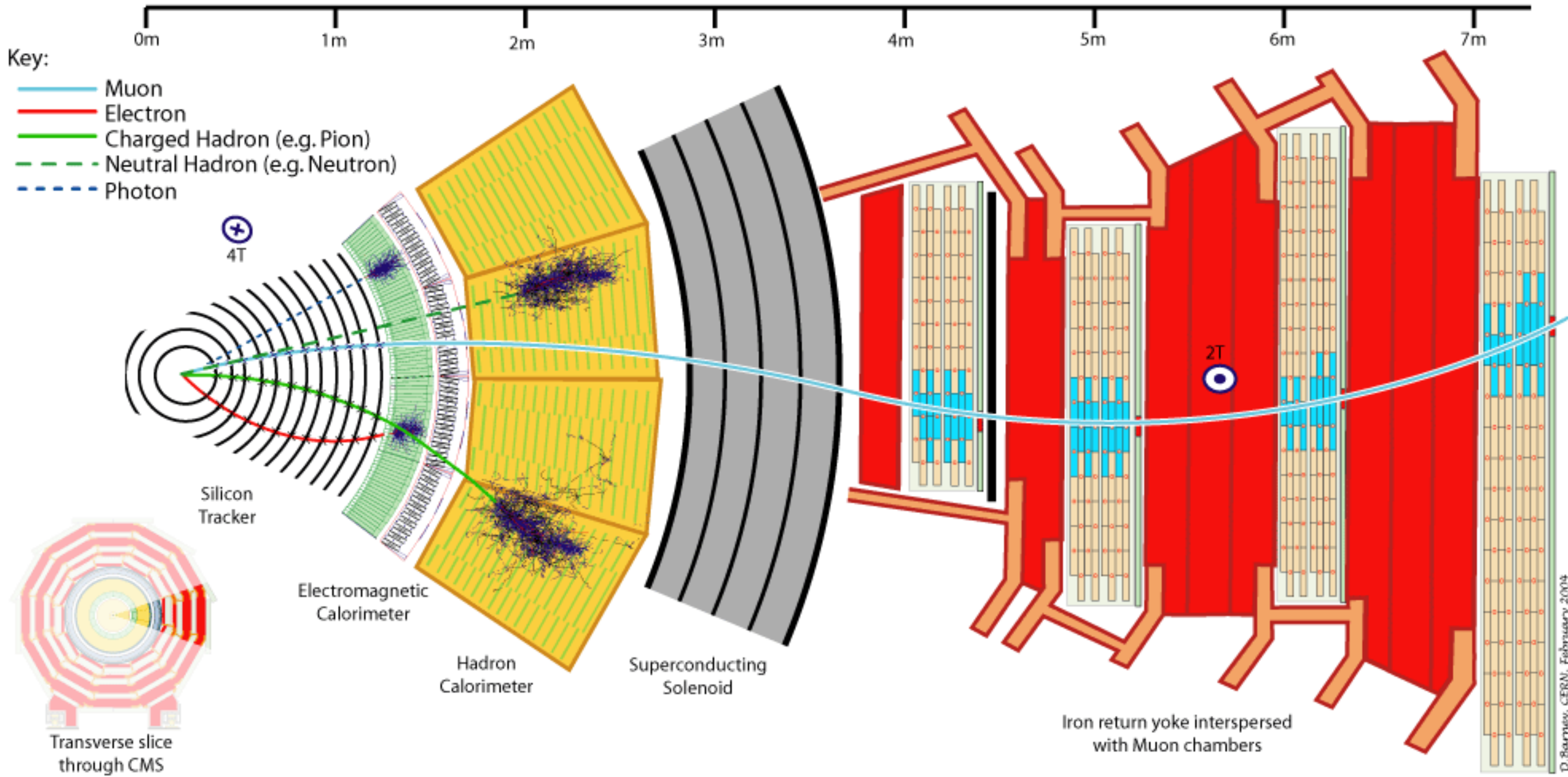
Physics process hypothesis

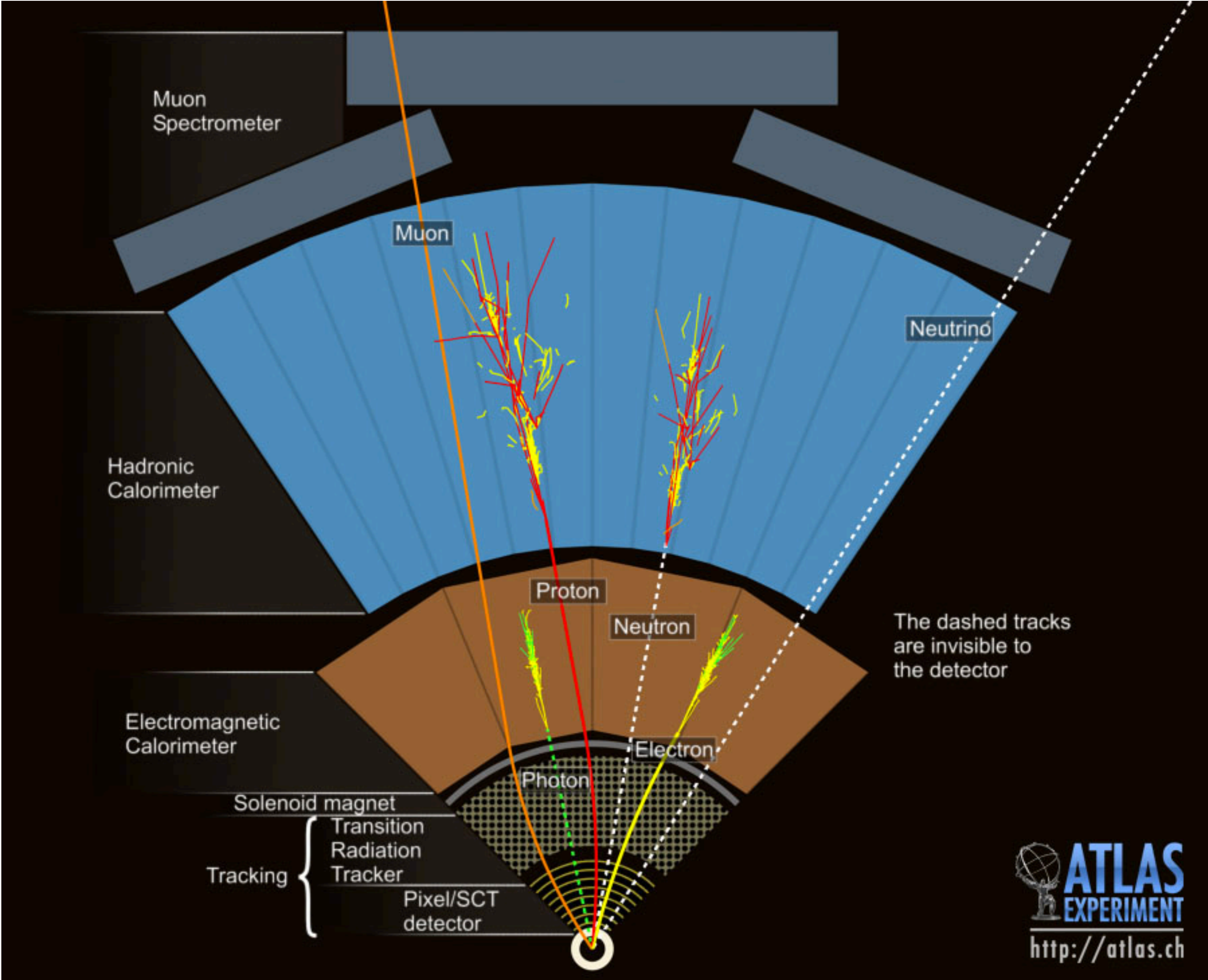


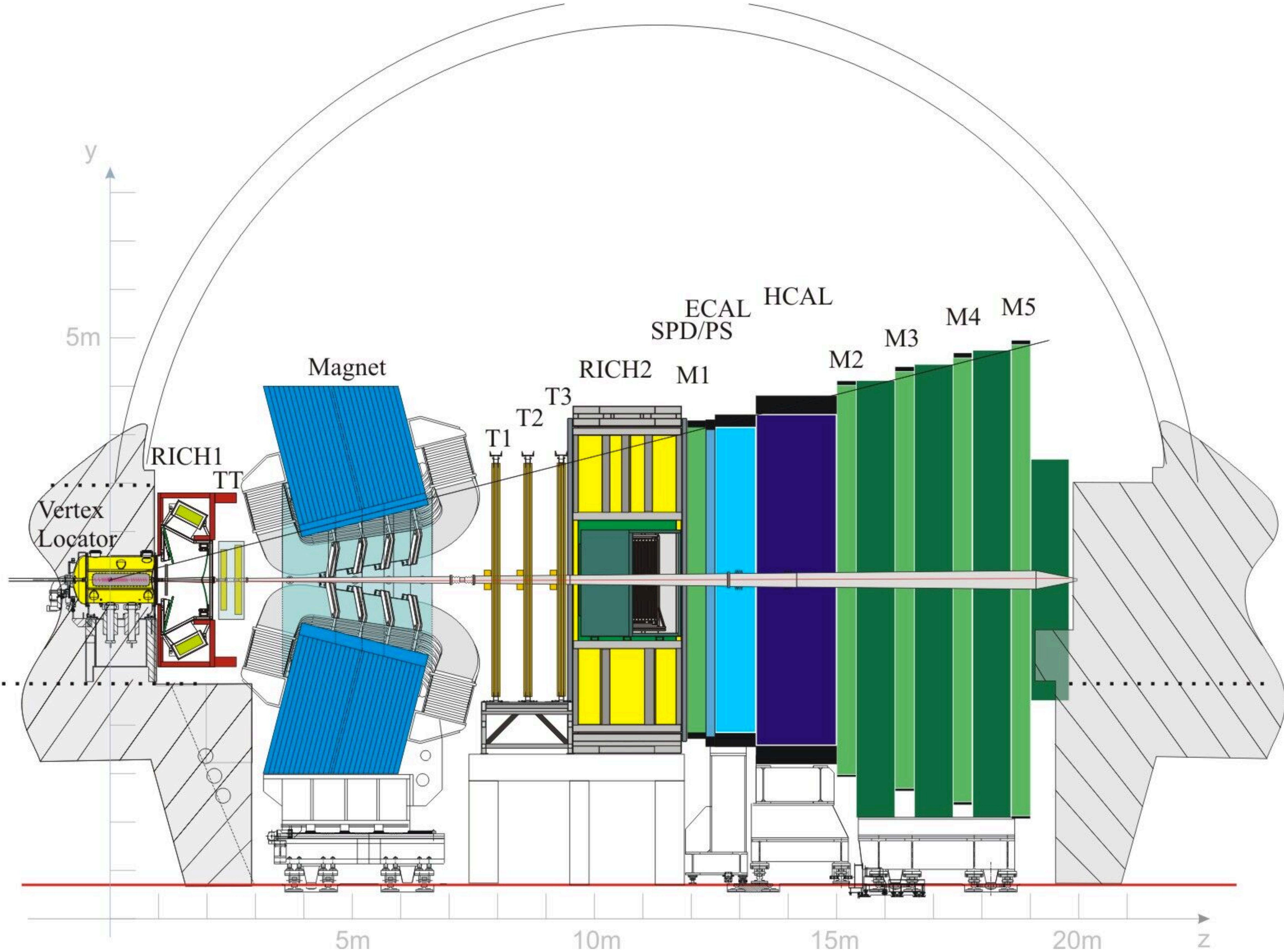
Particle Physics Detectors

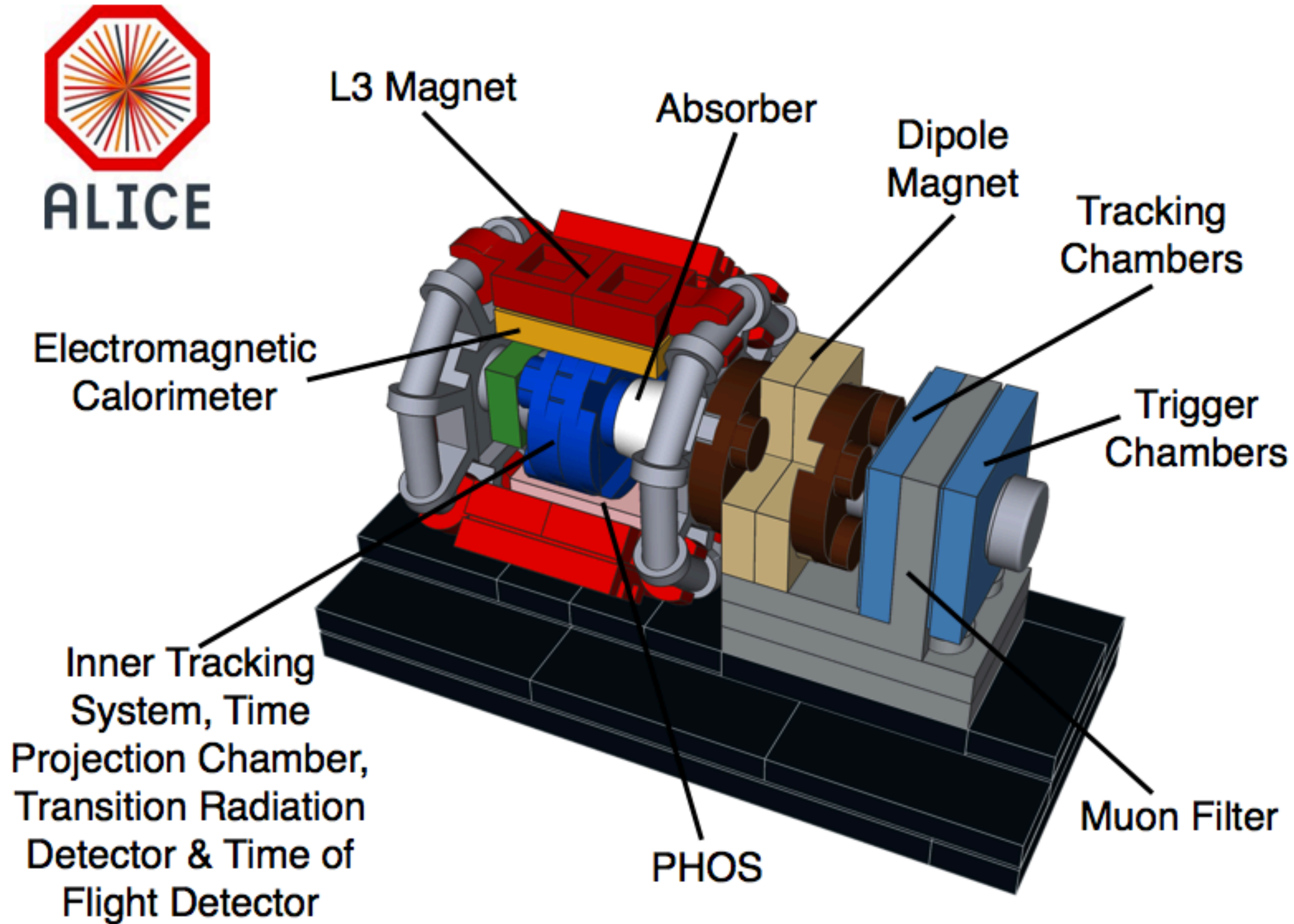
PARTICLE IDENTIFICATION











Detectors are built in layers to detect different species of (semi-) stable particles

Goal: determine momentum, energy, charge, mass

Techniques:

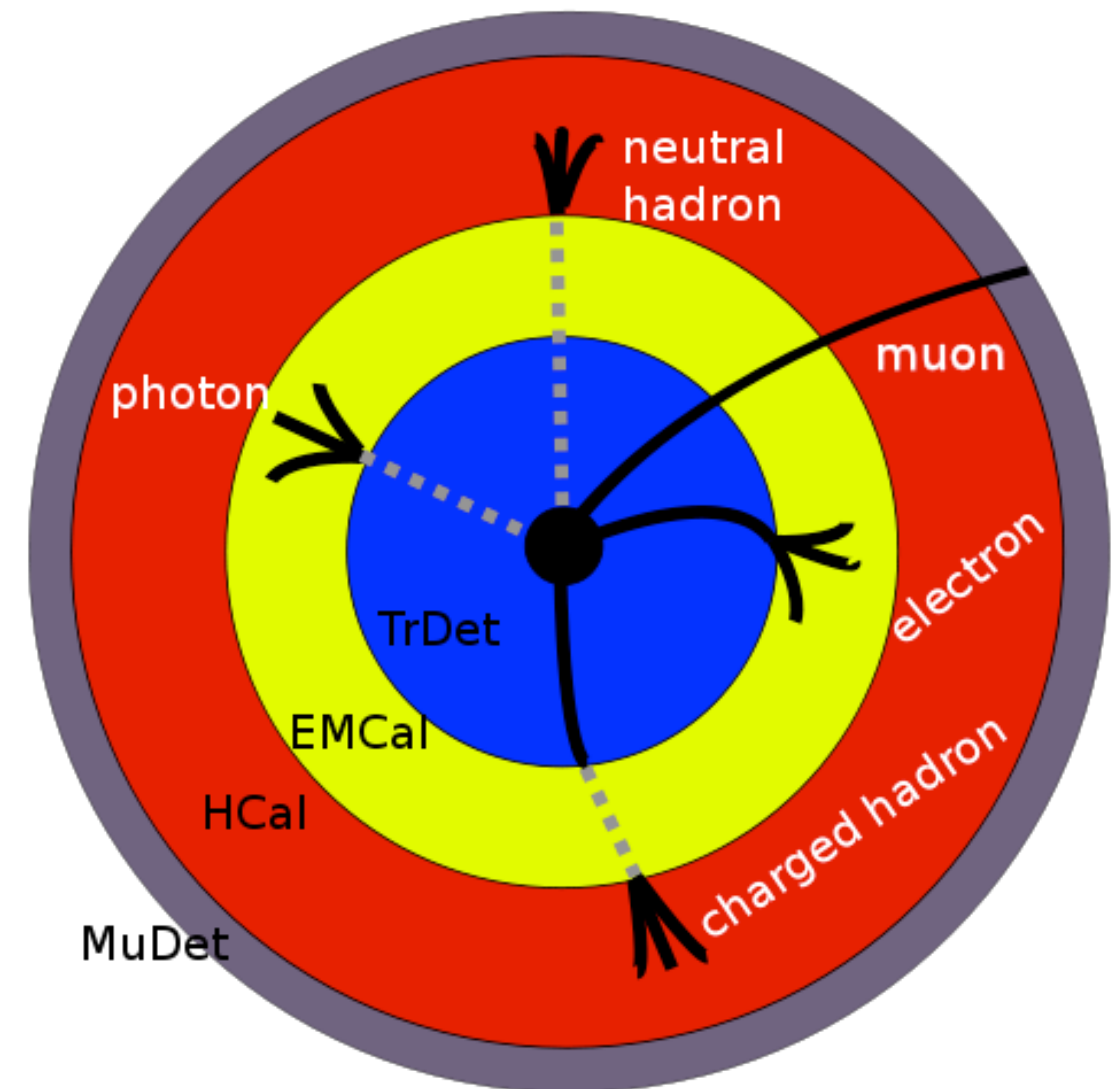
Energy loss (dE/dx)

Total Energy (E_{dep})

Velocity (β)

Curvature ($1/\rho$)

DETECTORS ARE LIKE OGRES



MuDet: muon detectors
TrDet: trace detector + vertex detector
EMCal: electromagnetic calorimeter
HCal: hadron calorimeter

Introduce the basic way we identify particle types and measure particle properties

Important: the resolution effects associated with performance of that reconstruction

Next:

Explore the **complementarity** of those measurements

Build up those objects to get to more complex objects

Goals:

Understand why we have all these different layers of detector and how they complement each other!

Understand reconstruction strategies, from the simplest to the most complex objects, and the physics concepts behind them

1. BUILDING BLOCKS



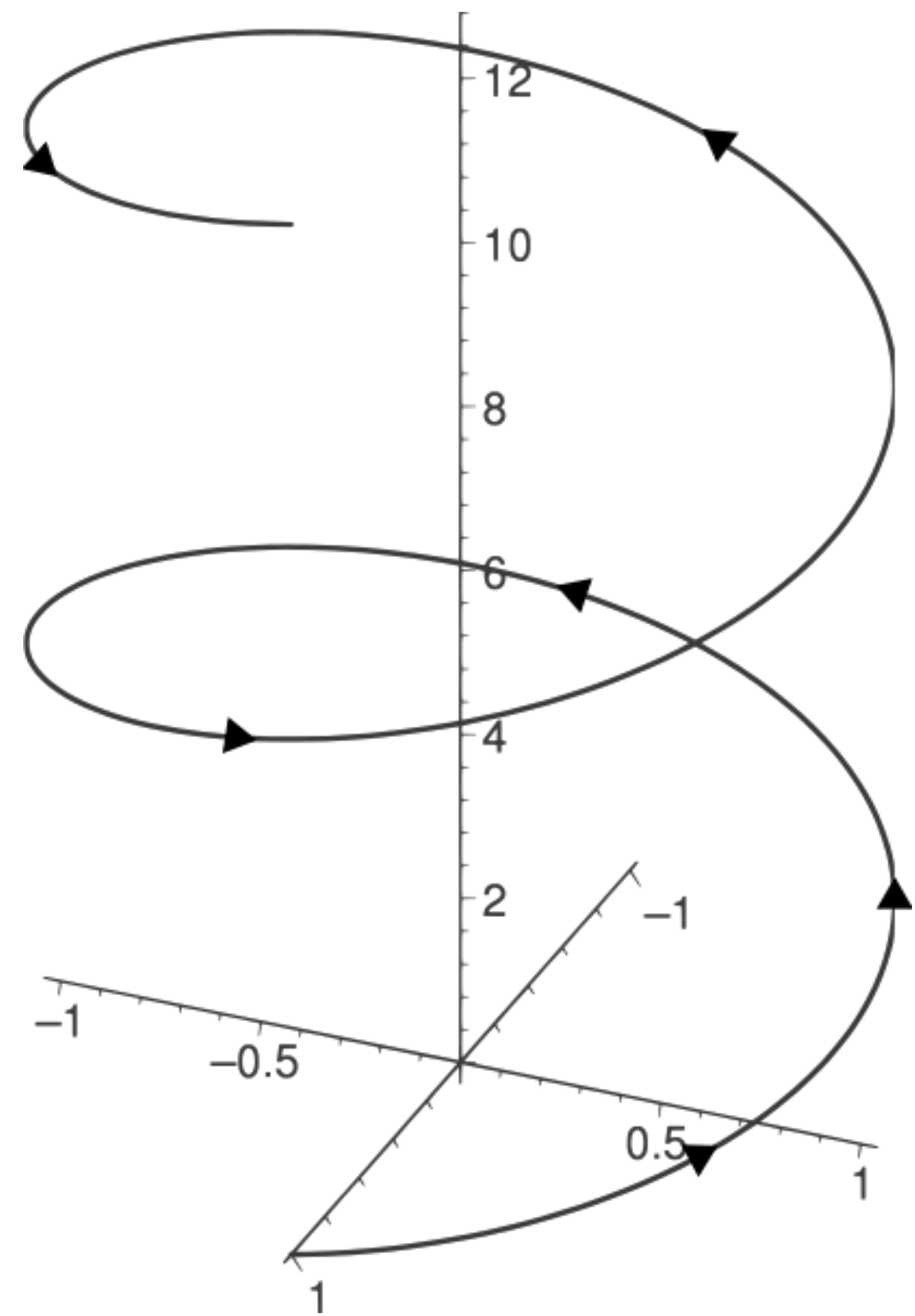
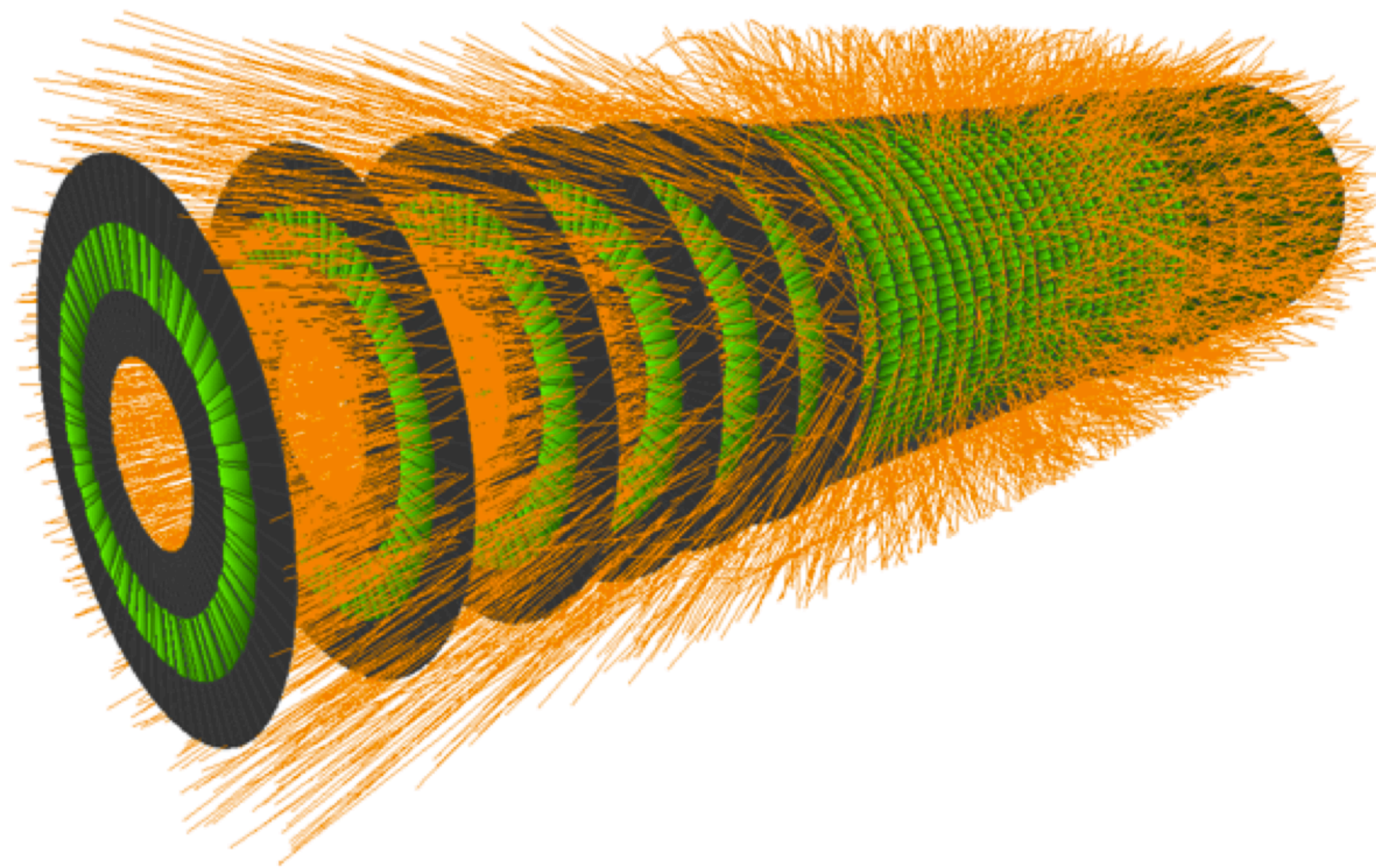
Tracking, Timing, Calorimetry

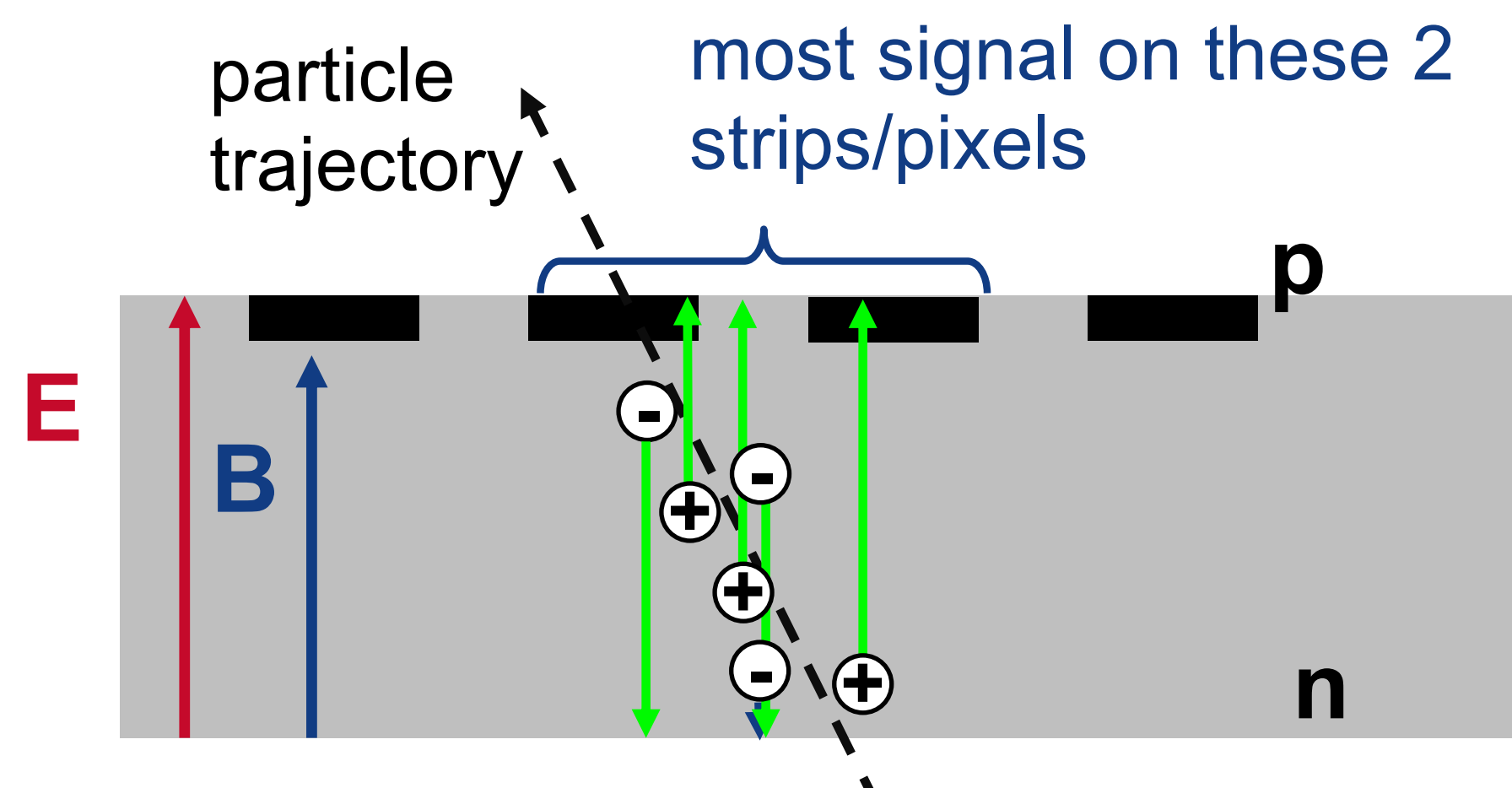
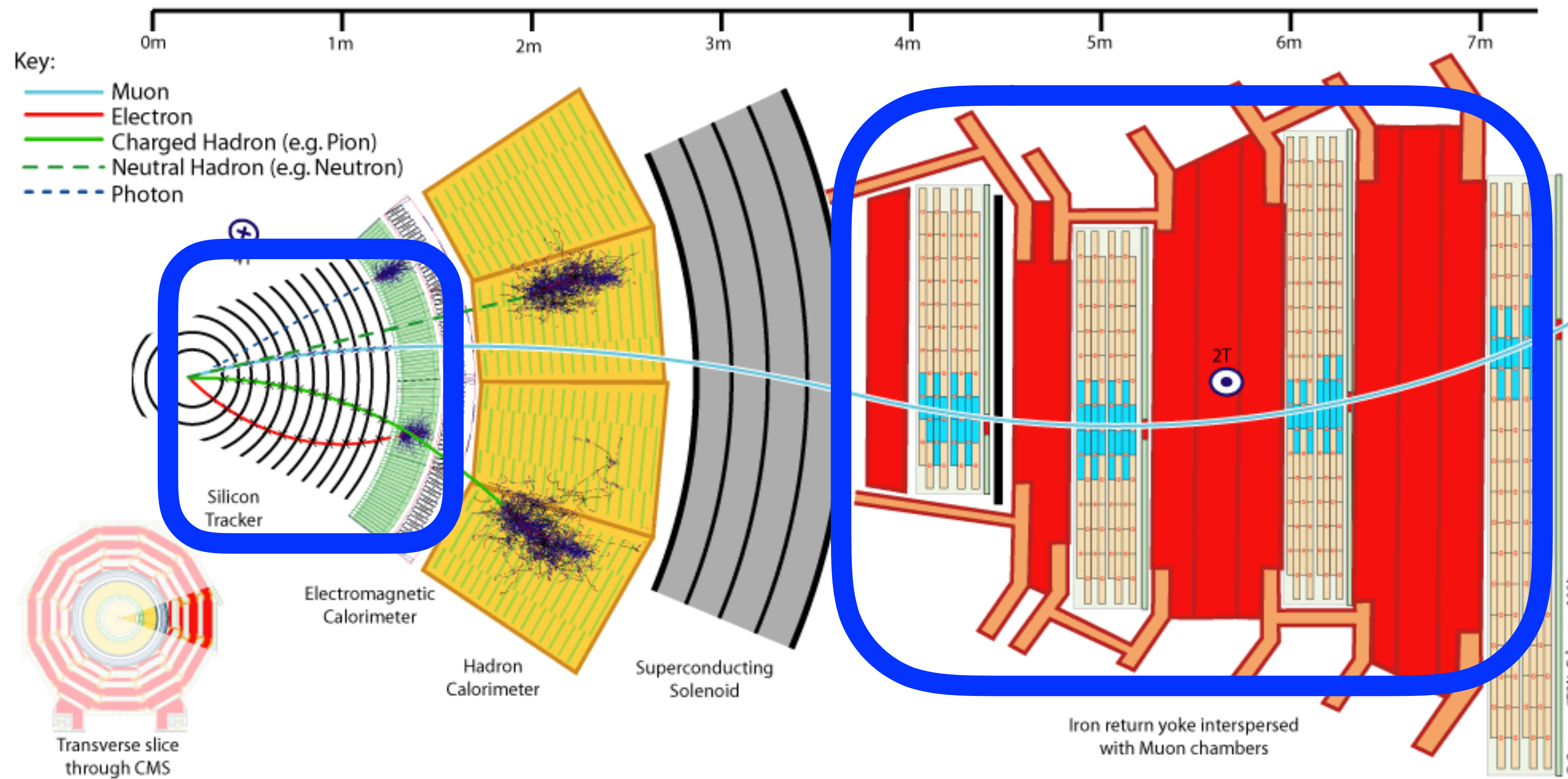
Some overlap with previous lectures, but I'll pull out the most relevant parts for reconstruction

Charged particles in a strong magnetic field follow a helical trajectory with curvature proportional to momentum

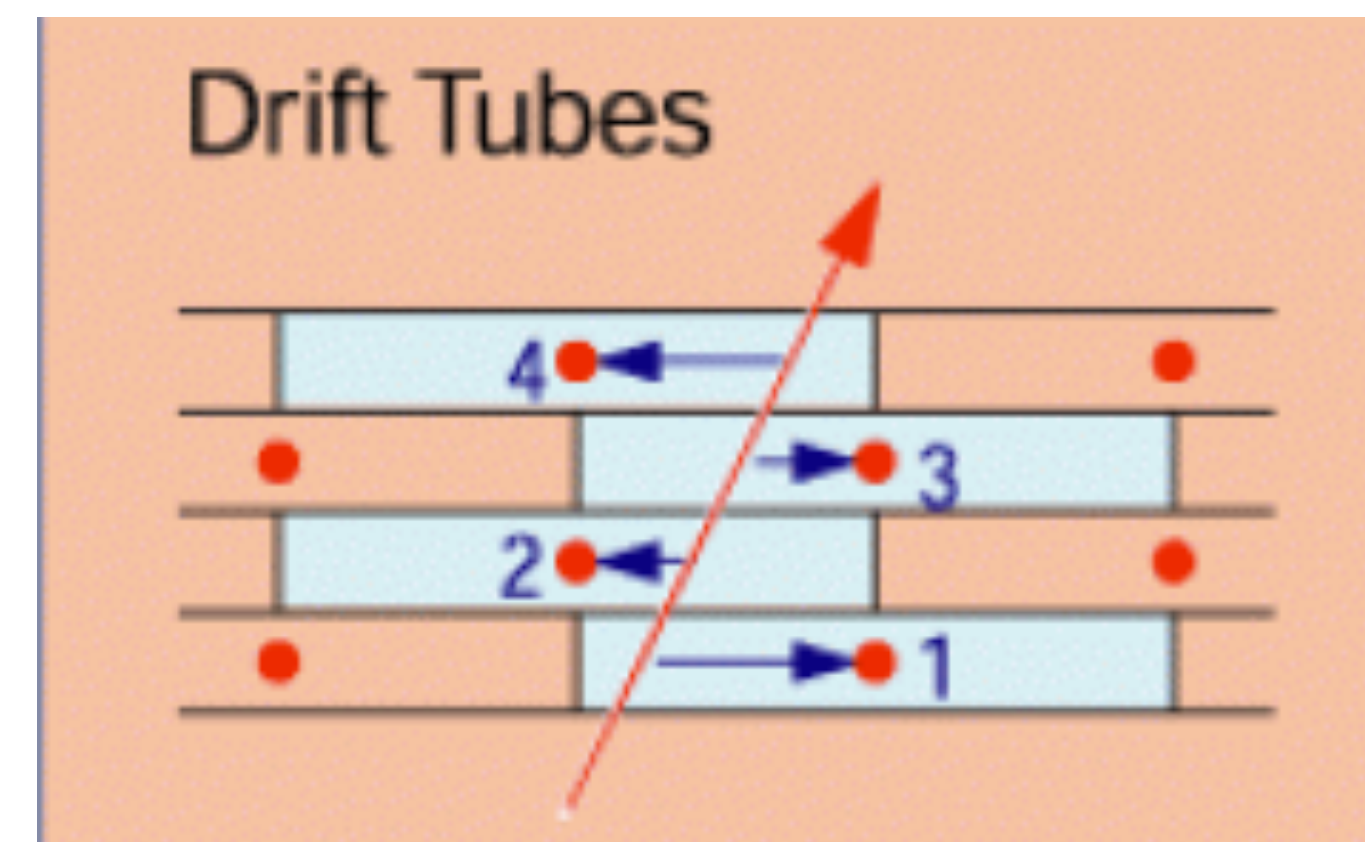
Determine track parameters:

- p_T
- theta, phi
- impacts parameters: d_0 , d_z





Precise, high-granularity silicon pixel and strip detectors are the workhorse



Muon trackers have to economically cover a lot of ground! Example are gaseous drift tube detectors

Tracking in the inner tracking volume is an important and compute intensive task

A constant challenge and one of the big bottlenecks in the reconstruction chain

Combinatorics are huge!

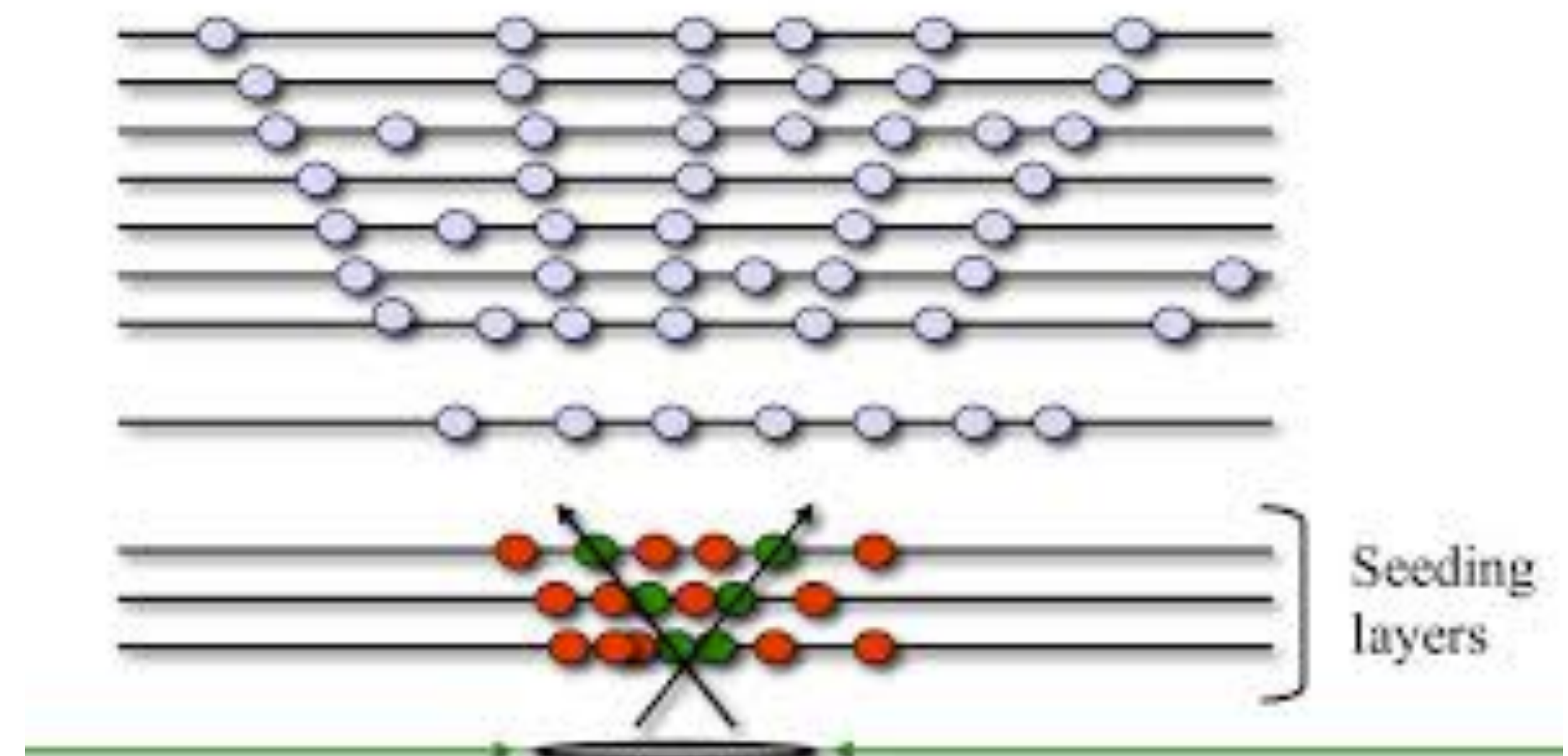
4 Basic Steps:

Seeding: initial candidate from a few hits

Finding: extrapolating from seeds with Kalman filter

Fitting: smooth trajectory and fit params

Selection: apply quality cuts



FITTING FOR MOMENTUM

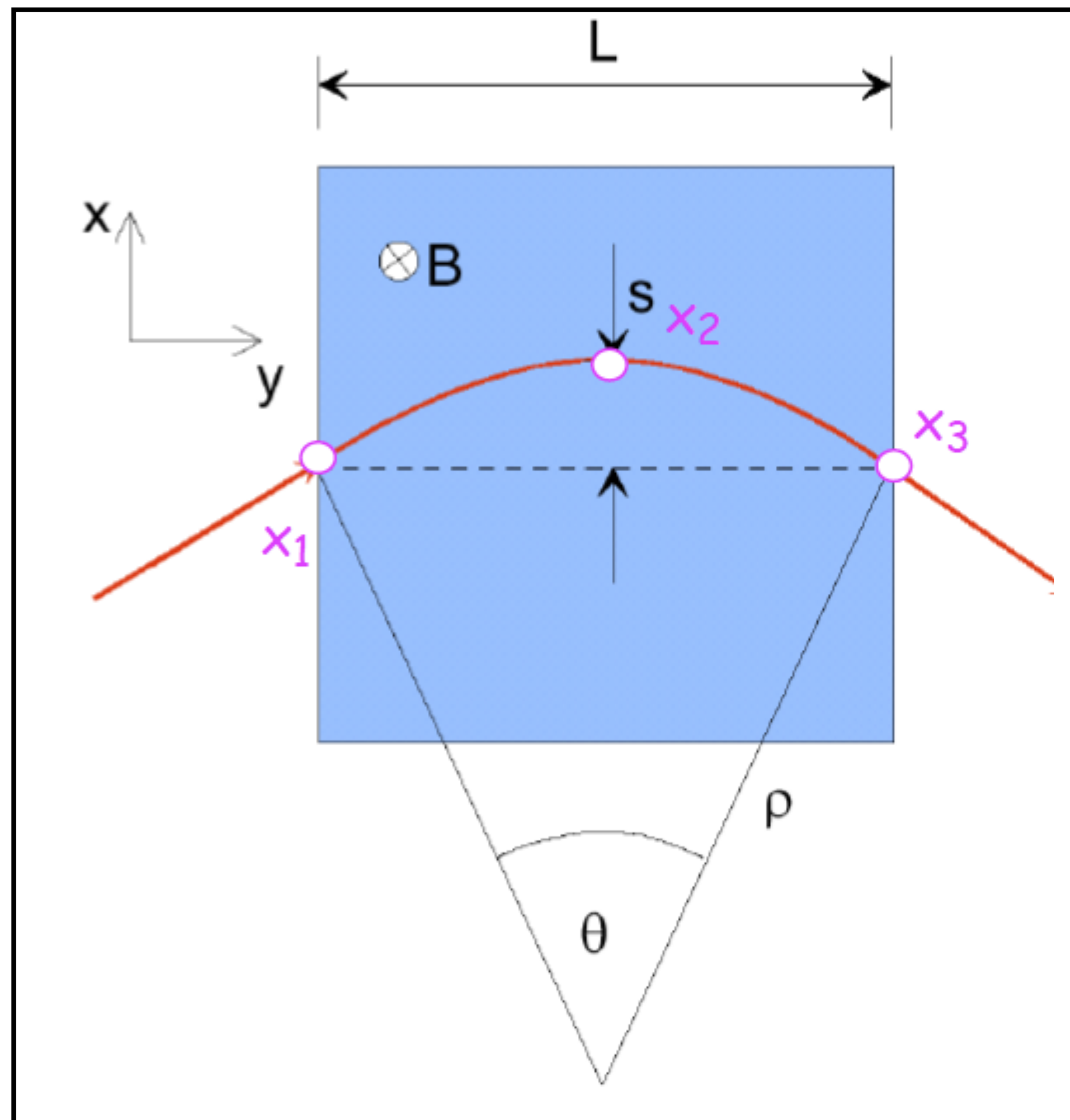
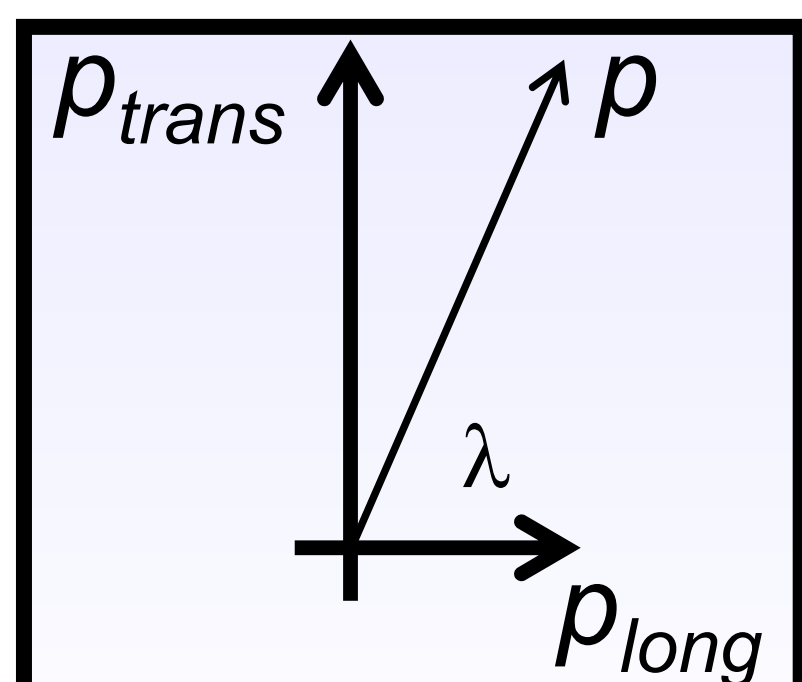
To get the pT of the track, we fit for its curvature

Useful formula:

$$p_T[\text{GeV}/c] = 0.3 \times B[T] \times r[m]$$

The full momentum is related by the polar angle

$$p = \frac{p_T}{\sin\lambda}$$



Lorentz Force

$$\vec{F}_L = q \cdot \vec{v} \times \vec{B}$$

Centripetal Force

$$F_c = m \cdot v^2 / r$$

$$p = q \cdot B \cdot r$$

MOMENTUM RESOLUTION

The transverse momentum resolution is driven by:

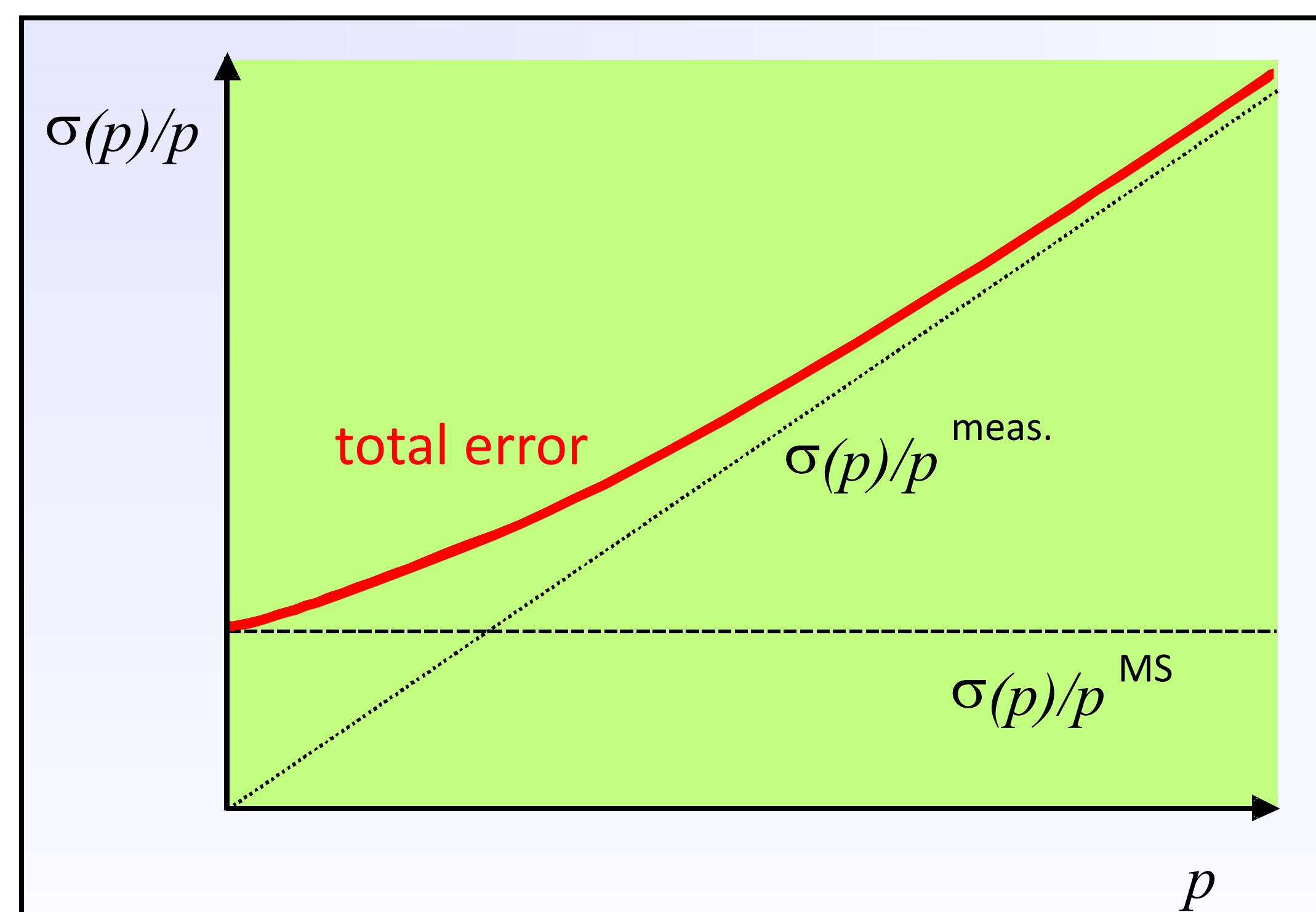
Curvature measurement and hit resolution

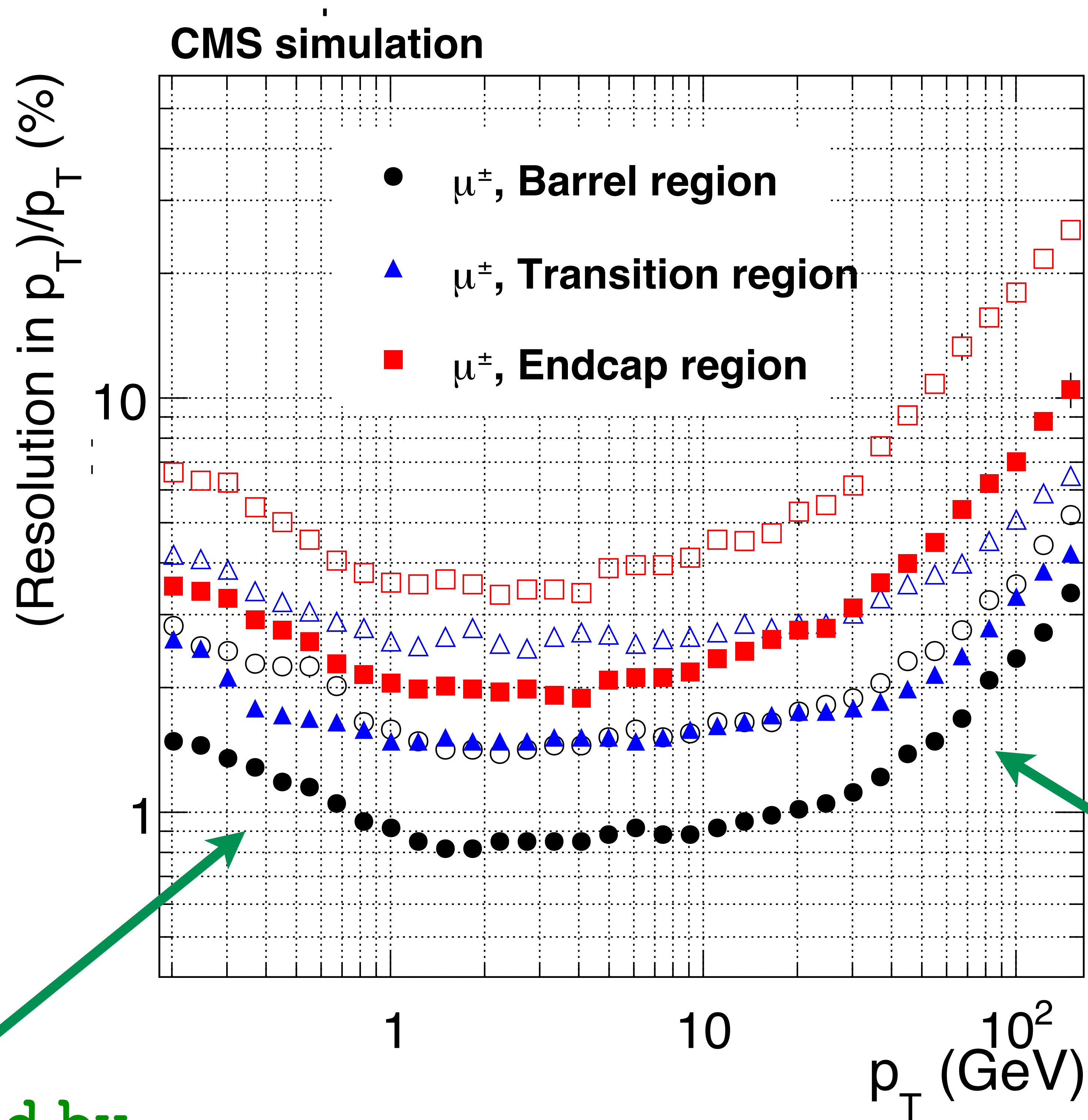
Multiple scattering

$$\frac{\delta p}{p} \sim \frac{0.0136}{\beta} \sqrt{\frac{X}{X_0}} \frac{1}{0.3BL} \frac{\sqrt{4A_N}}{N}$$

$$\left(\frac{\sigma_{p_T}}{p_T}\right)^2 \propto c_1 \cdot \left(\frac{p_T}{BL^2} \sqrt{\frac{720}{N+4}}\right)^2 + c_2 \cdot \left(\frac{1}{B\sqrt{LX_0}}\right)^2$$

curvature **multiple scattering**

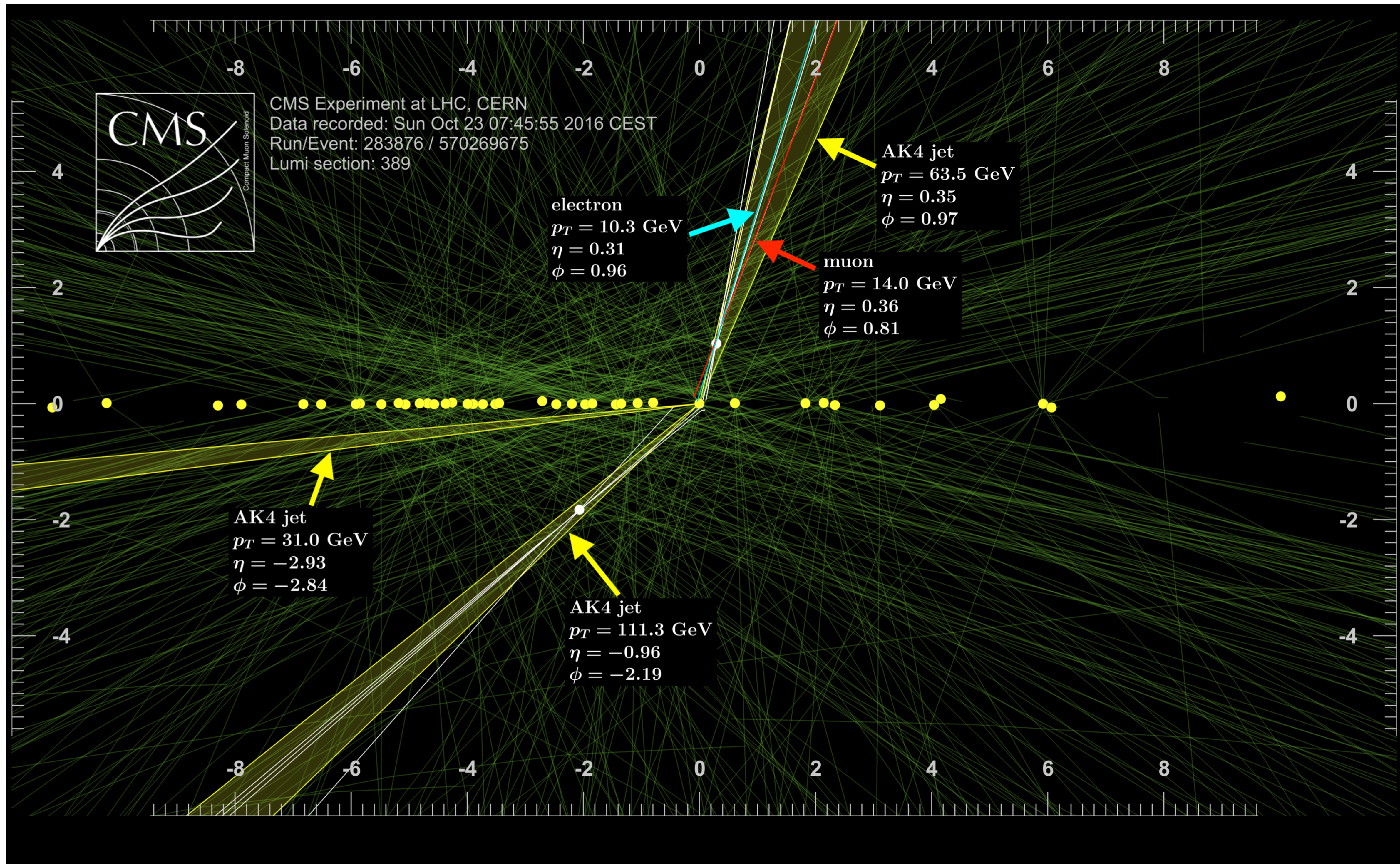




WHY IS RESOLUTION WORSE IN MORE FORWARD REGIONS?

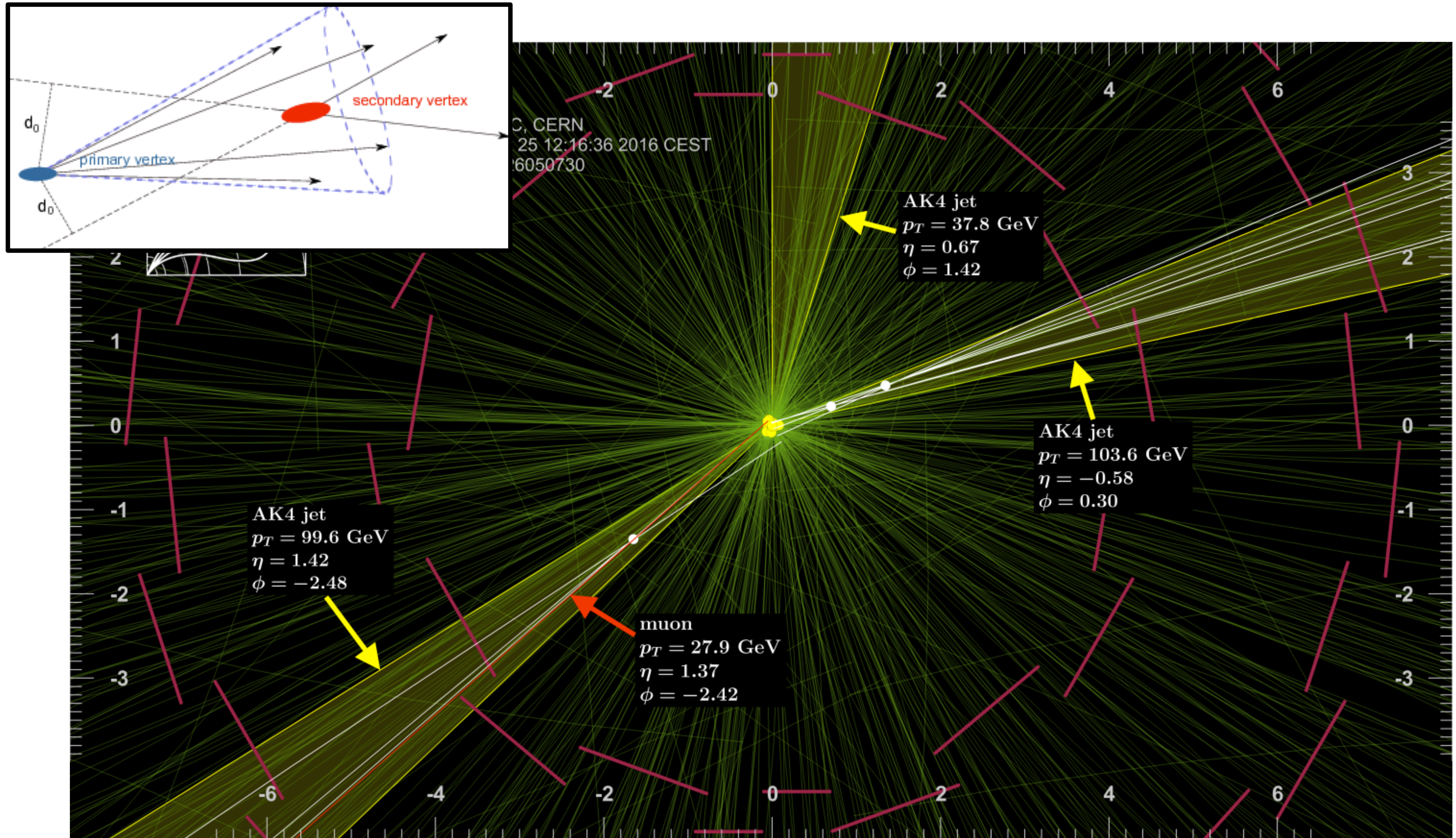
Dominated by multiple scattering

Dominated by curvature meas.



Use Z position of the primary vertex to separate pileup
(much more on this later)

VERTEX RECONSTRUCTION [XY]

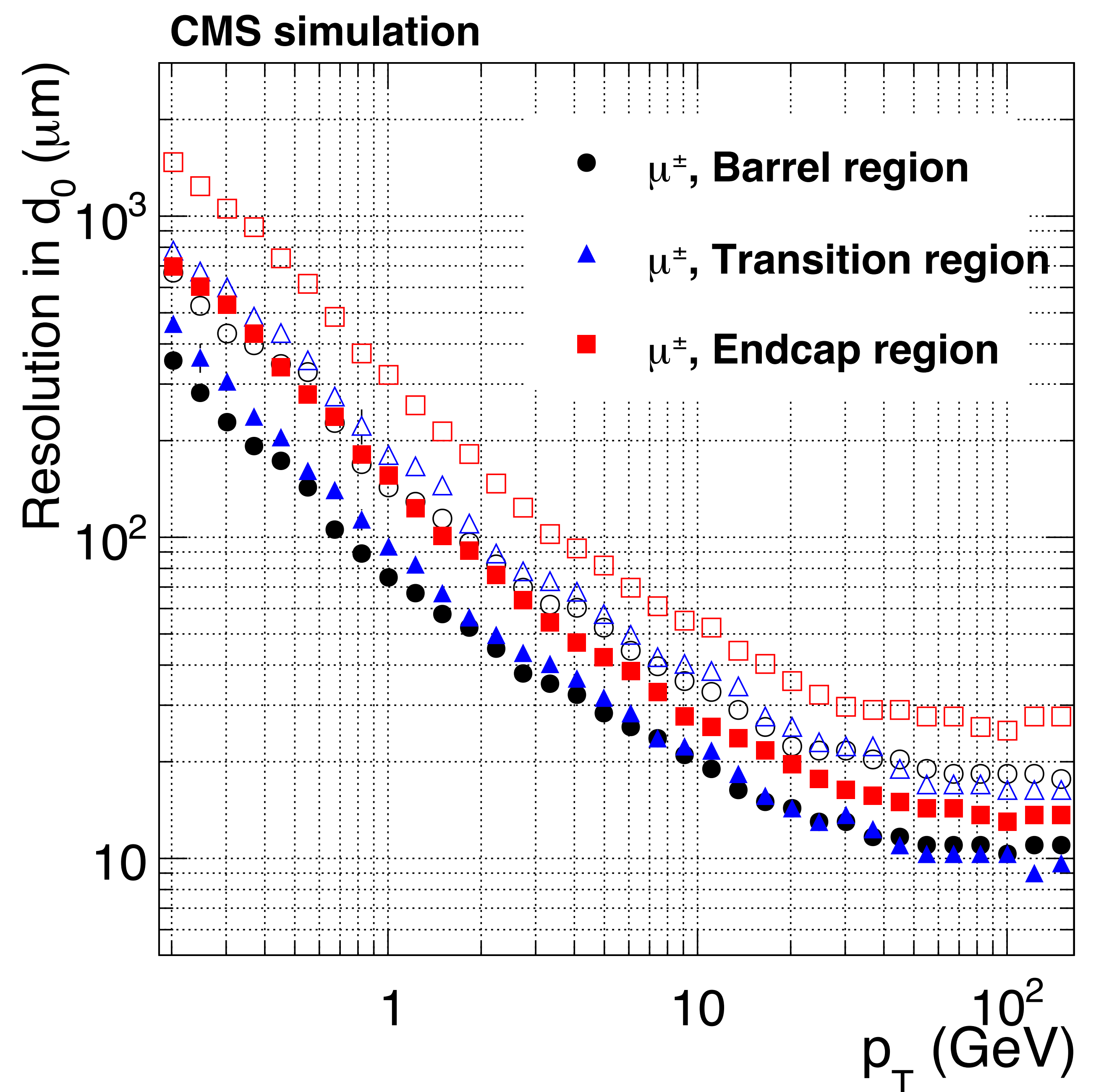


Use impact parameter of the secondary vertex to identify displaced vertices

The main drivers of the vertex resolution are the position measurement and the lever arm of the measurement (how far are you away from the vertex)

For example:

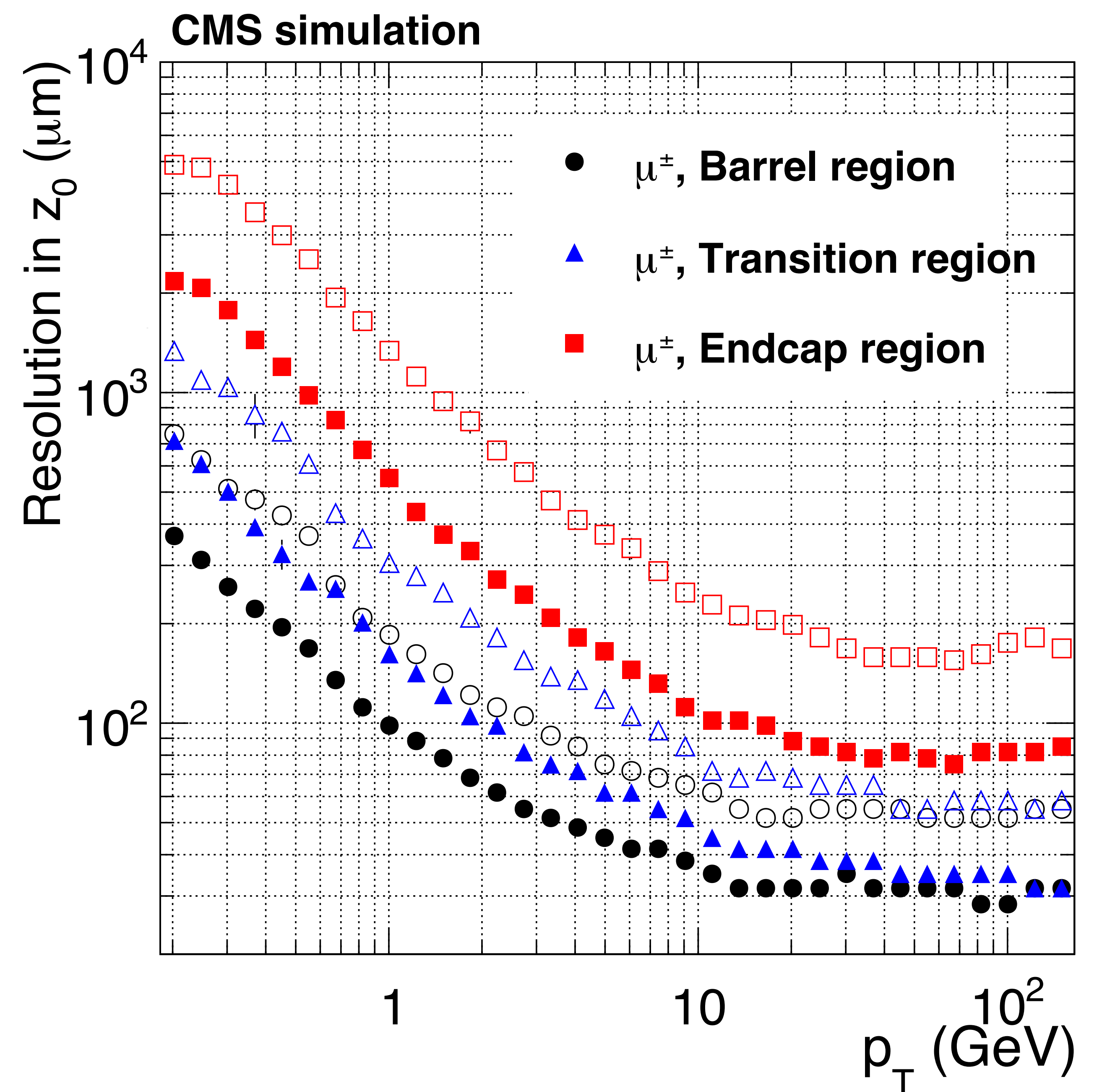
$$\sigma_{d_0}^2 = \frac{r_2^2 \sigma_1^2 + r_1^2 \sigma_2^2}{(r^2 - r^1)^2} + \sigma_{MS}^2$$

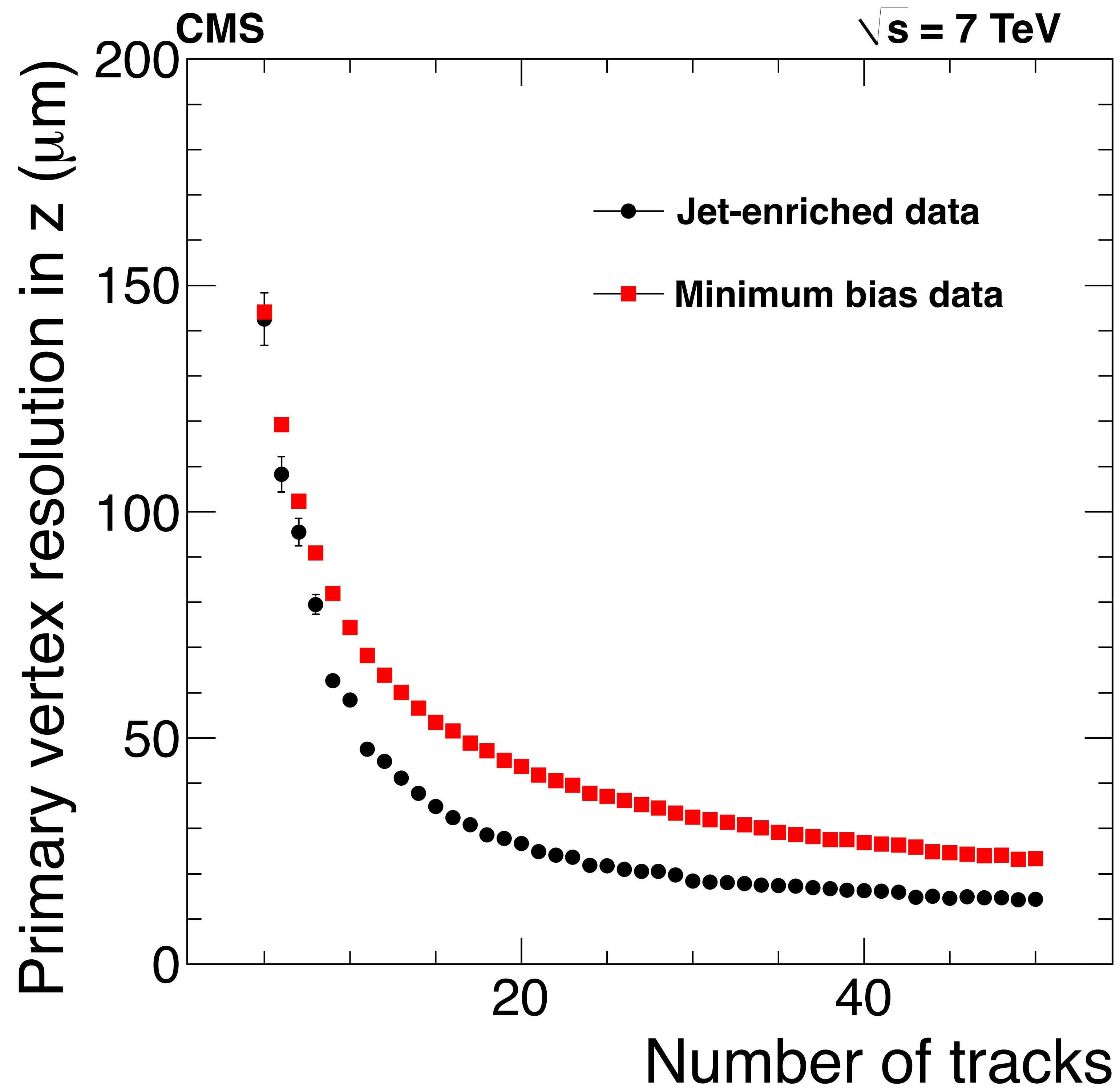


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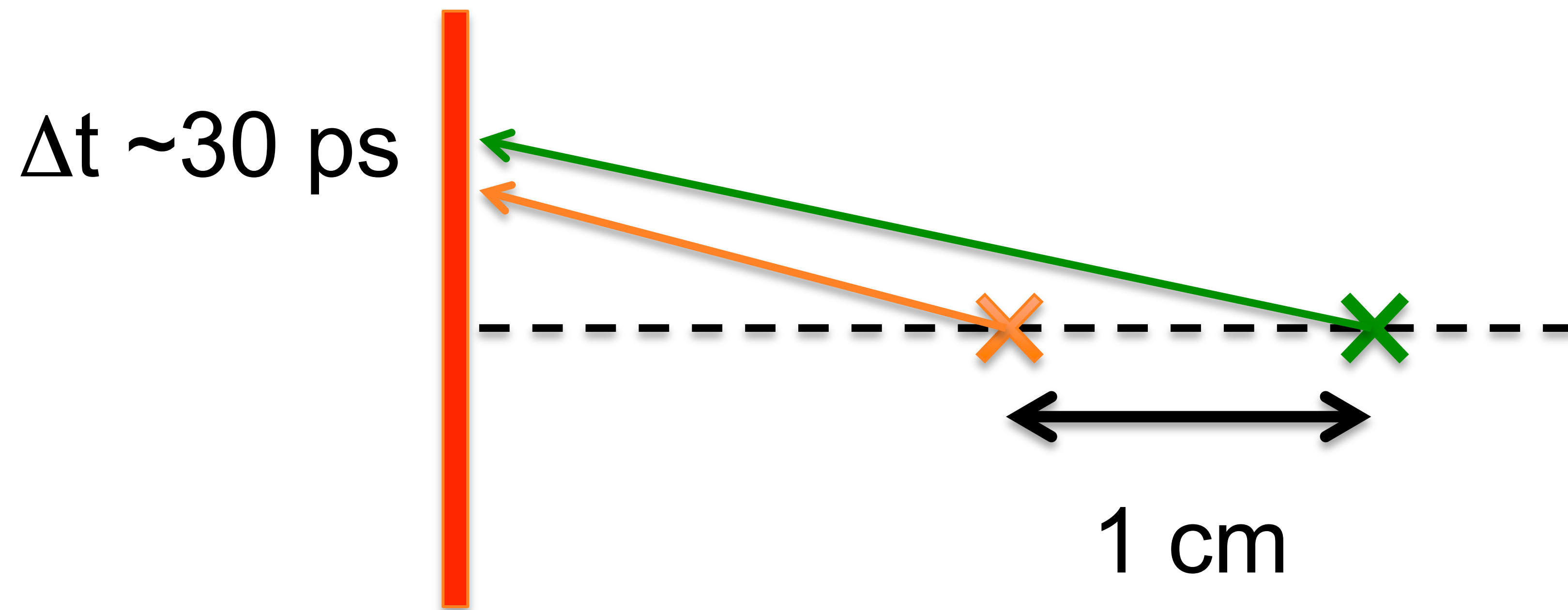




~similar
performance of
the vertex
resolution in
 x, y too

Precision fast timing has promise to be a powerful additional piece of information for reconstruction

There are plans by ATLAS and CMS to include precision timing detectors for HL-LHC upgrades

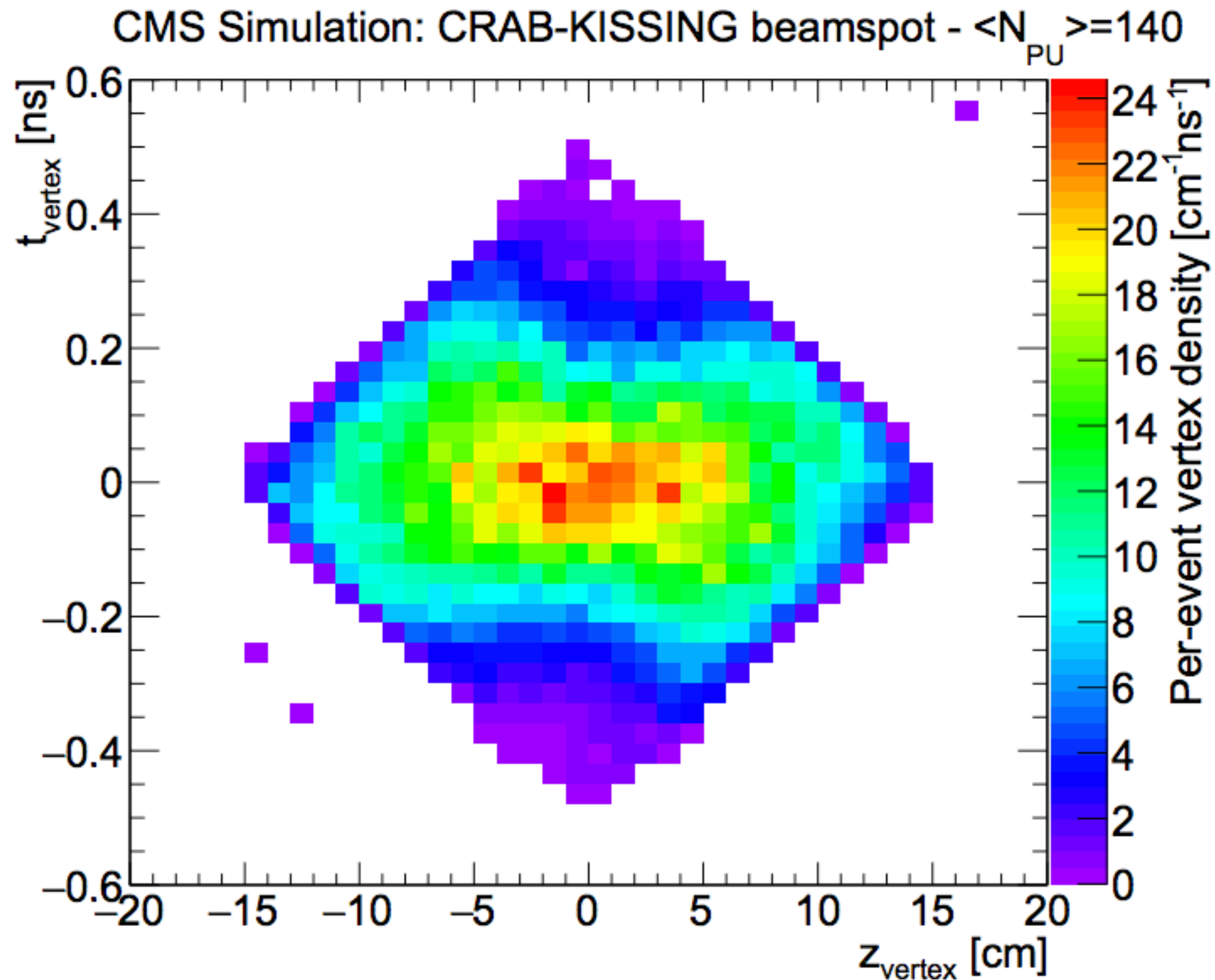


Preliminary!

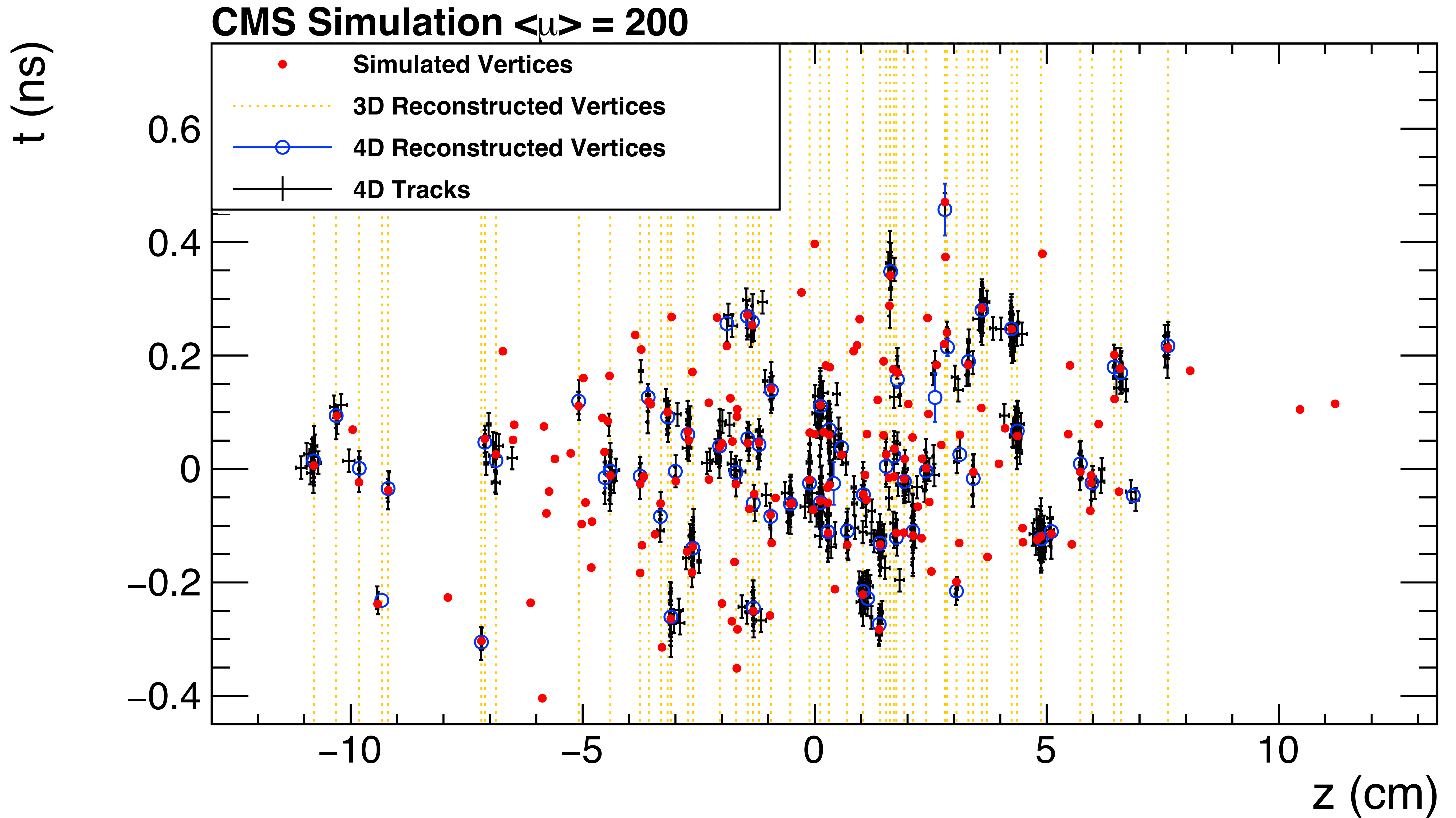
Resolution for charged particles is around $\sim 30 \text{ ps}$.

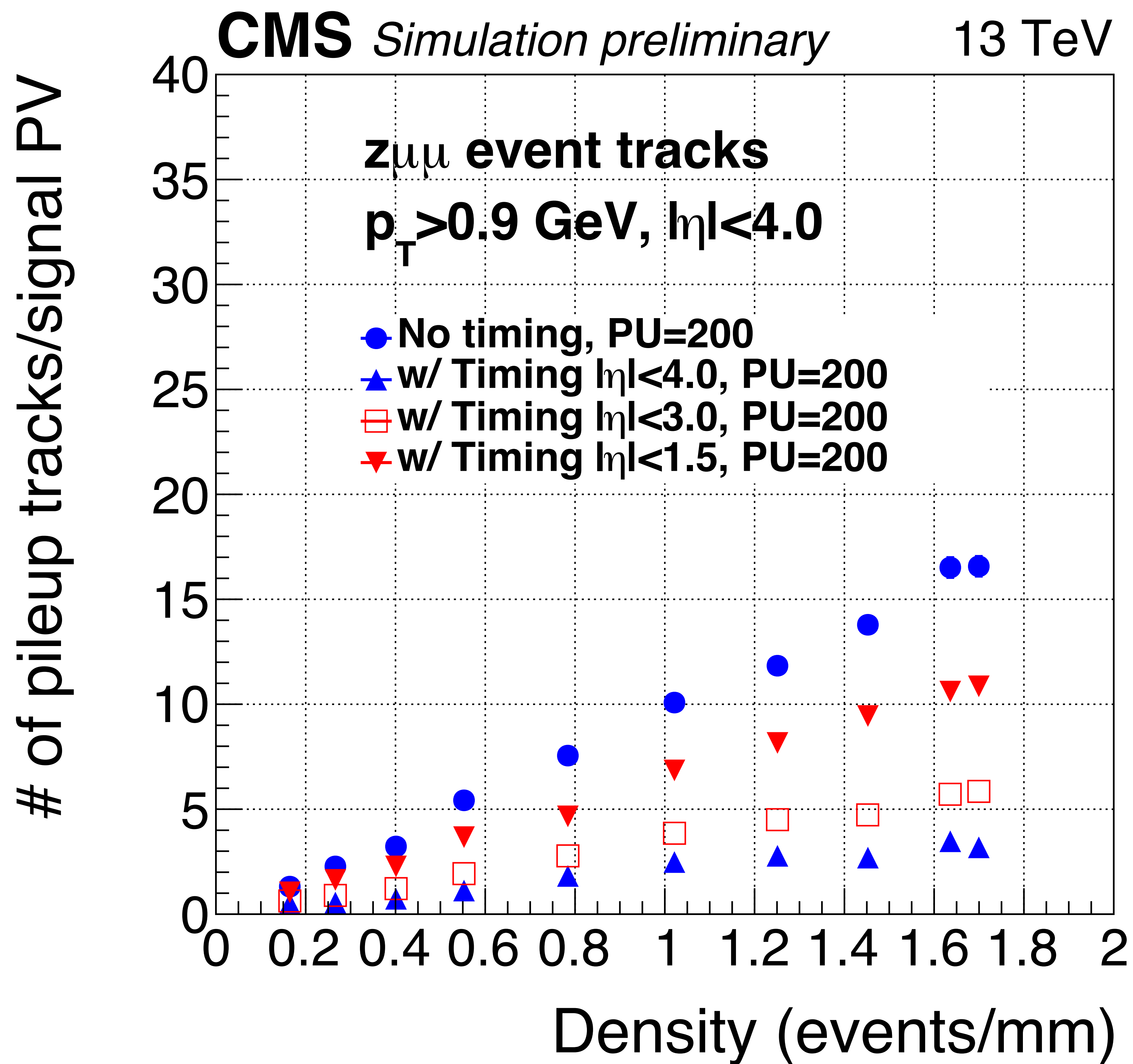
Neutral resolution is energy dependent: $\sim 30\text{--}300 \text{ ps}$ for 100–few GeV

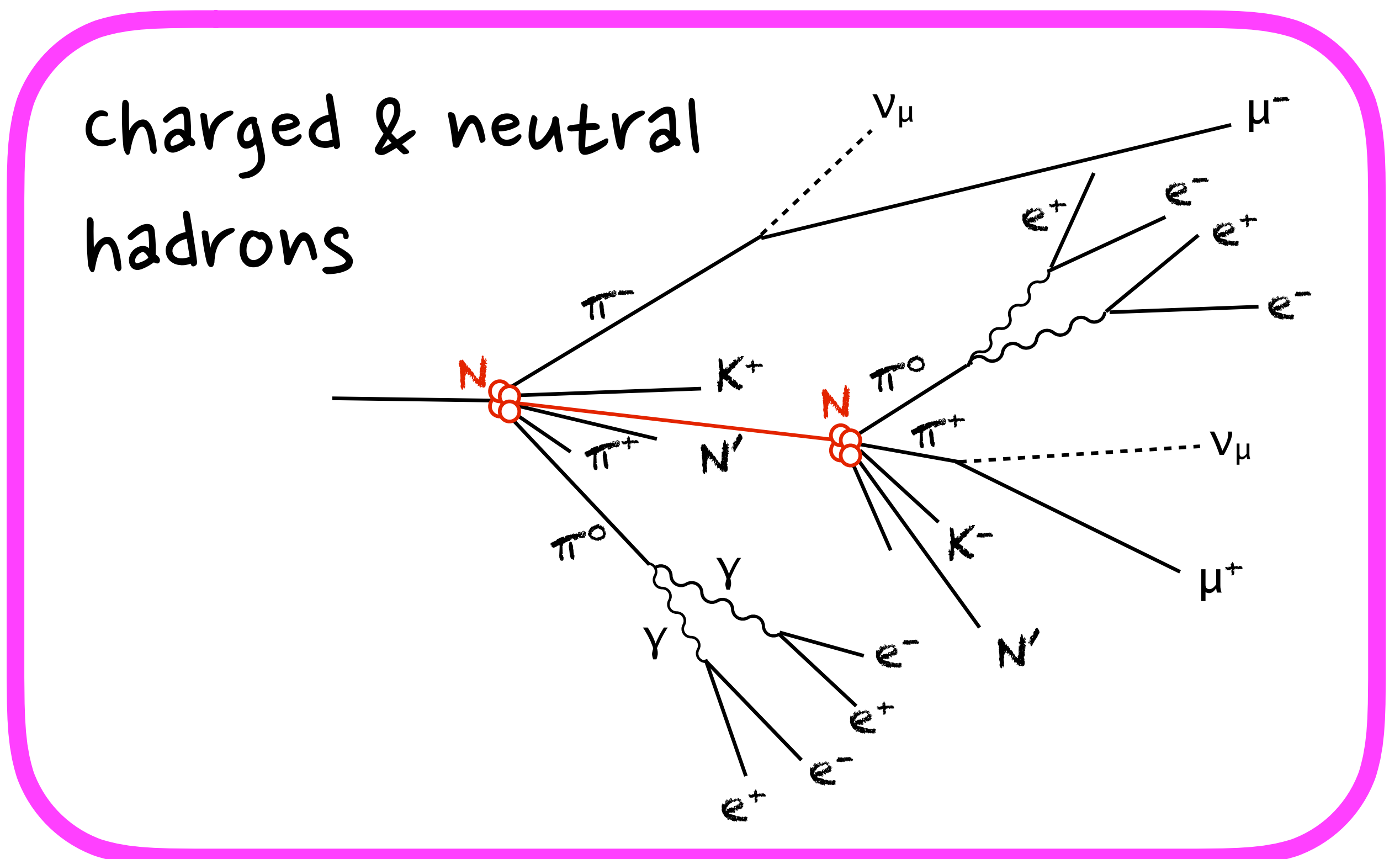
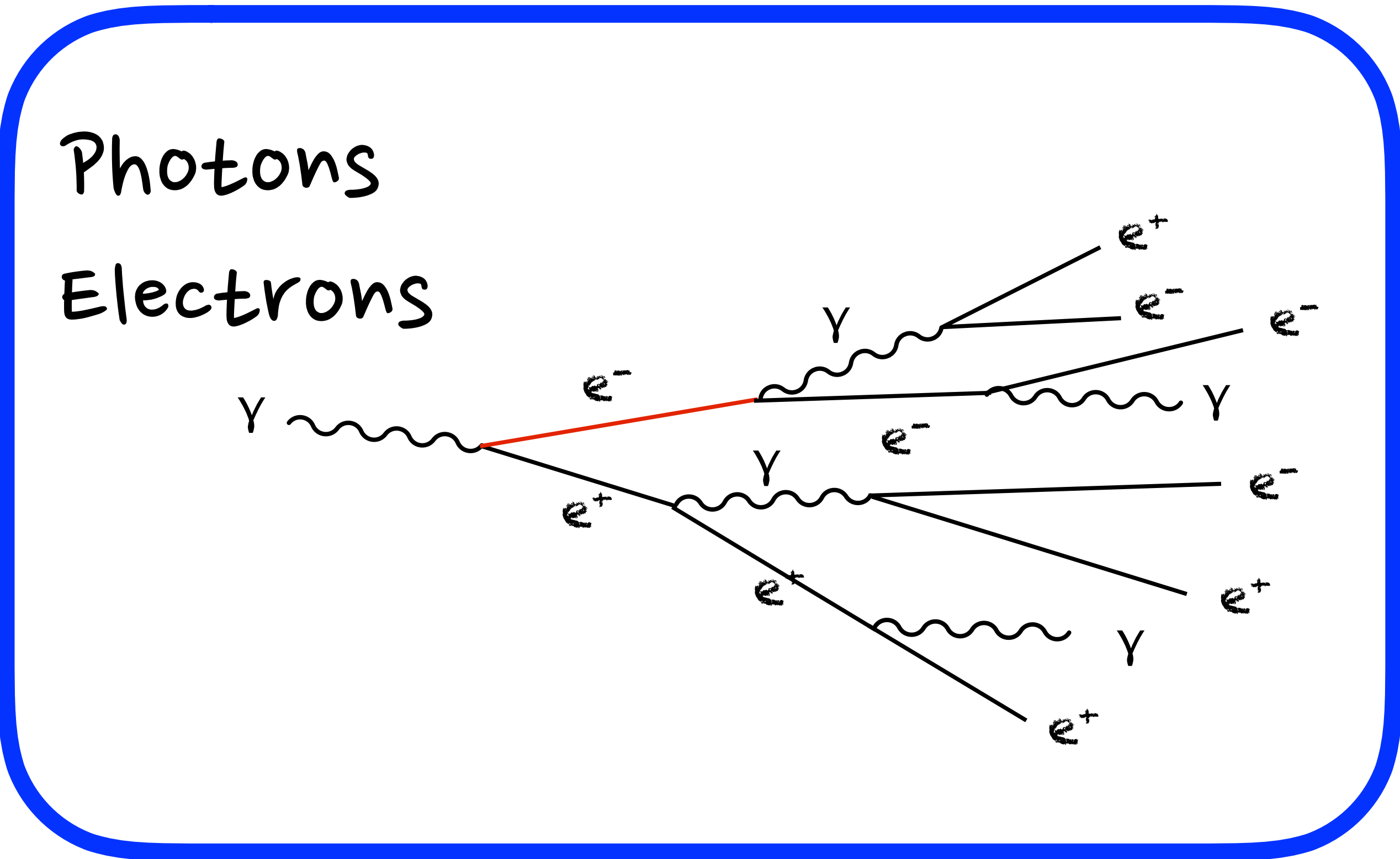
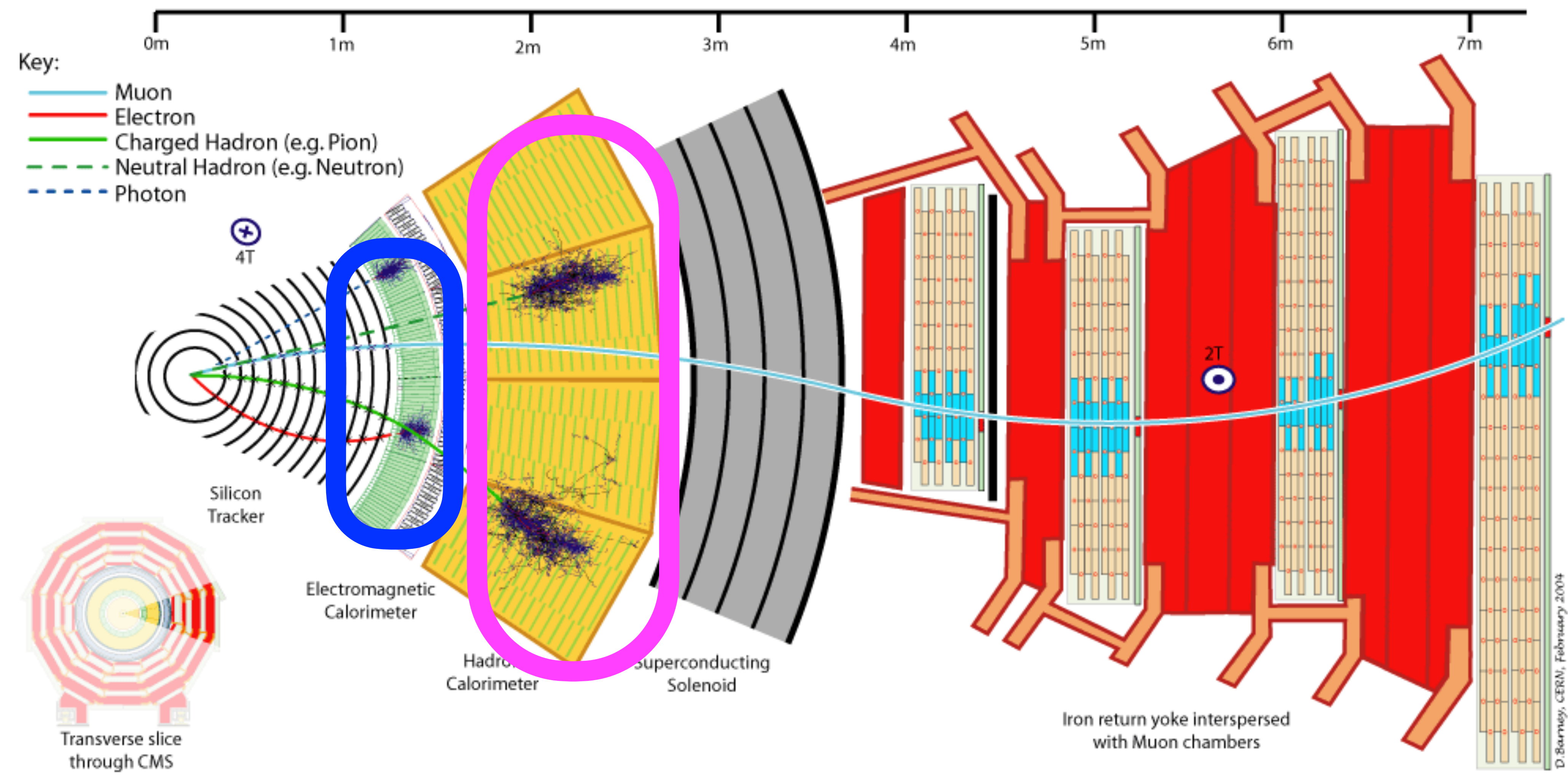
Time of flight can be used to disentangle the origin of particles as well — particularly useful for neutral particles



In future conditions of ~ 200 pileup, timing can be used to disentangle pileup vertices. Proton beam crossing spread out of z and time!







Cf. Calorimetry lectures from R. Wigmans

A reminder of the basics: energy resolution and characteristic size of electromagnetic and hadronic showers

Resolution:

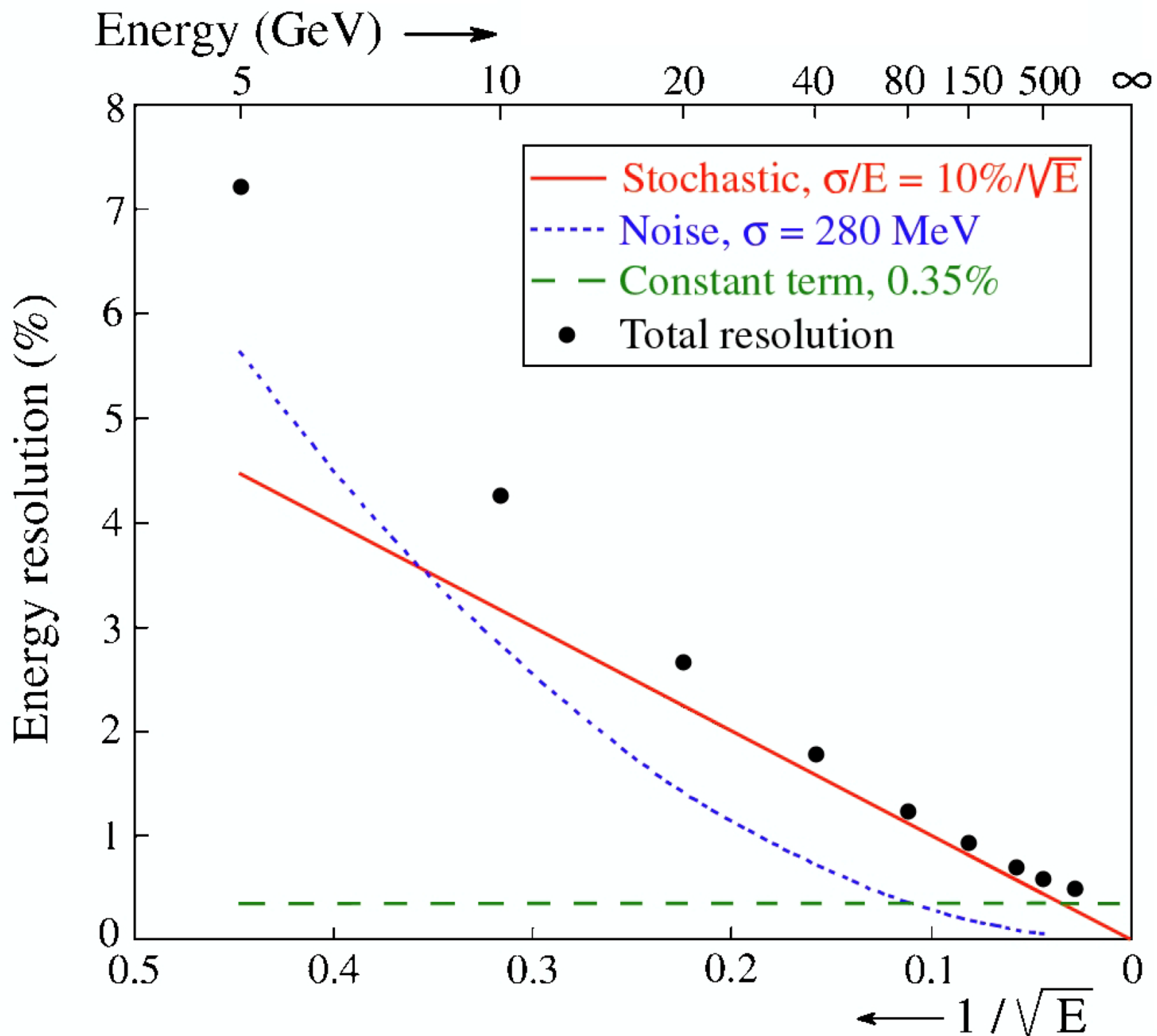
$$\frac{\sigma_E}{E} = \frac{a}{E} \oplus \frac{b}{\sqrt{E}} \oplus c$$

Noise term:
fixed vs. energy
Typically important at
low energies

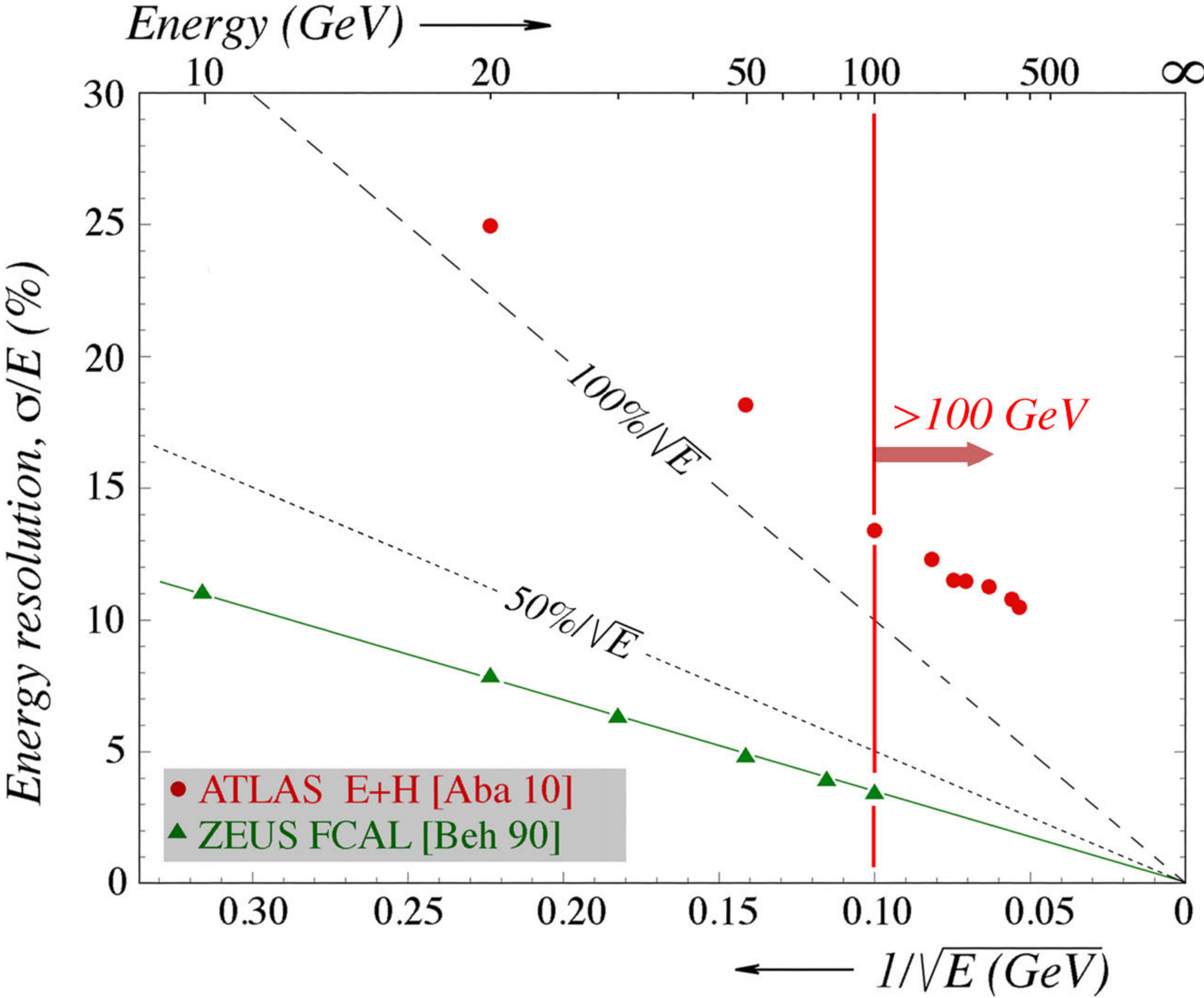
Stochastic term:
Error is $\sim E^{-1/2}$, as a
counting error

Constant term:
instrumental effect,
shower leakage, etc.

EXAMPLE: ATLAS EM CALORIMETER



Note the change in scale!!



Another important consideration in reconstruction are the size of the showers

EM showers are much smaller, uniform

Hadronic showers are larger, less-uniform

Important concept

$$X_0 = \frac{716.4 \text{ g cm}^{-2} A}{Z(Z+1) \ln(287/\sqrt{Z})}$$

X_0 , radiation length: characteristic length of a energy loss of particles interacting electromagnetically

Moliere radius: transverse size of the shower is related to X_0

$$R_M = 0.0265 X_0 (Z + 1.2)$$

λ , interaction length: characteristic length of particles interacting with nuclei

<http://pdg.lbl.gov/2017/AtomicNuclearProperties/>

The screenshot shows a periodic table with elements color-coded by group. Below the table is a list of material categories with corresponding dropdown menus:

- Inorganic compounds (Al through Fe) - Aluminum oxide through ferrous oxide
- Inorganic compounds (Freon through Pu) - Freon through plutonium oxide
- Inorganic compounds (Potassium thru yttrium) - Potassium iodide through water
- Inorganic scintillators (BaF2 through Y2SiO5) - Barium fluoride through Y2SiO5
- Simple organic compounds - Acetone through Xylene
- Polymers - Polymers
- Mixtures - Aerogel through standard rock
- Biological materials - A-mn-dimethyl_formamide through tissue-equivalent gas

Atomic and nuclear properties of iron (Fe)

| | | | | |
|----------------------------|-------|--------------------|-------|----|
| Nuclear collision length | 81.7 | g cm ⁻² | 10.37 | cm |
| Nuclear interaction length | 132.1 | g cm ⁻² | 16.77 | cm |
| Pion collision length | 107.0 | g cm ⁻² | 13.59 | cm |
| Pion interaction length | 160.8 | g cm ⁻² | 20.42 | cm |
| Radiation length | 13.84 | g cm ⁻² | 1.757 | cm |

For high Z materials
 $X_0 \ll \lambda$

The different detector technologies and their intrinsic resolution are **complementary**!

As one starts to get worse, the other starts to get better
Tracking has the best intrinsic spatial resolution

Representative numbers for the CMS case

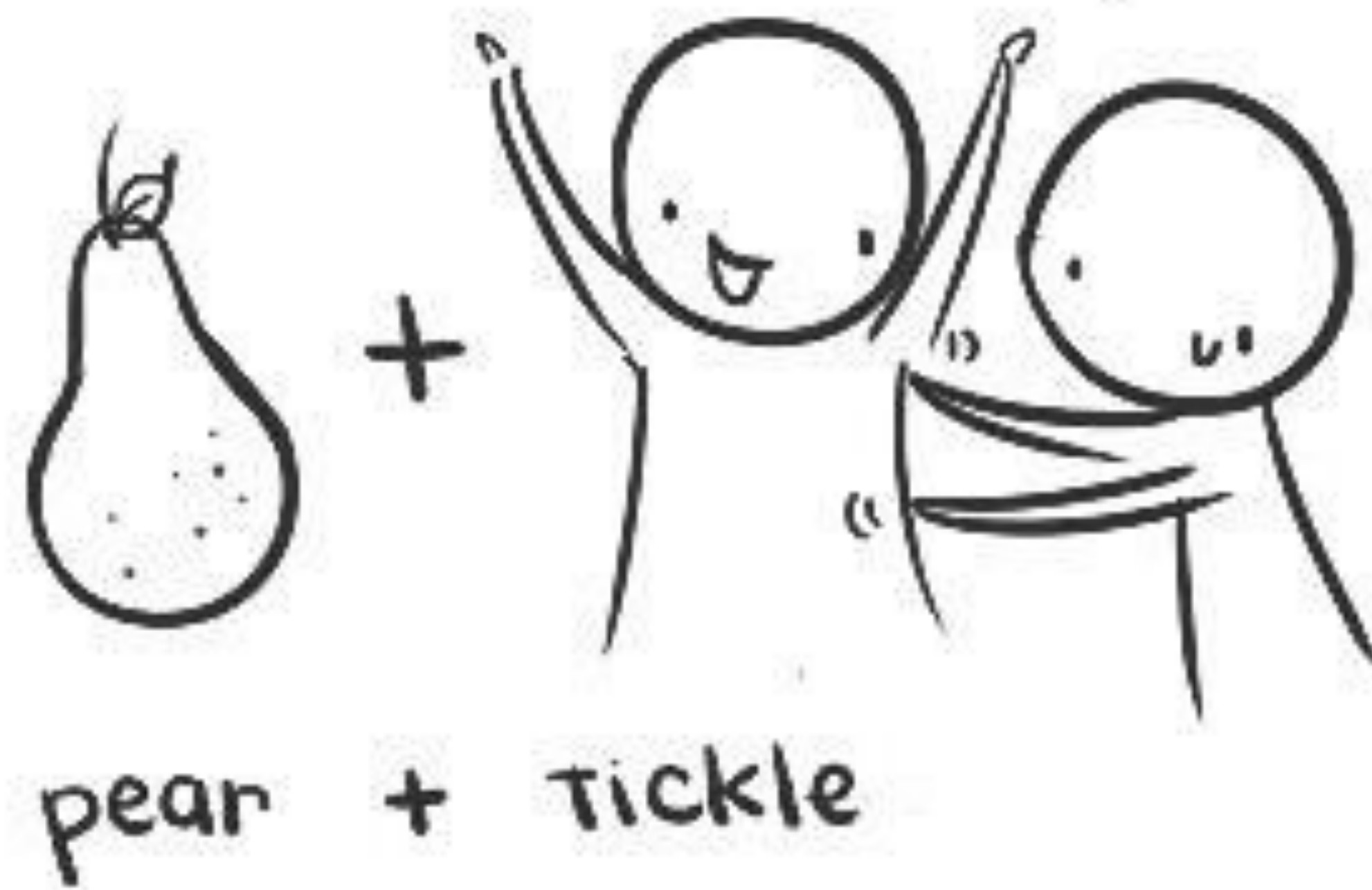
| Detector | p_T -resolution | η/Φ -segmentation |
|----------|-------------------------------|-----------------------------------|
| Tracker | 0.6% (0.2 GeV) – 5% (500 GeV) | 0.002 x 0.003 (first pixel layer) |
| ECAL | 1% (20 GeV) – 0.4% (500 GeV) | 0.017 x 0.017 (barrel) |
| HCAL | 30% (30 GeV) – 5% (500 GeV) | 0.087 x 0.087 (barrel) |

Vertexing numbers:

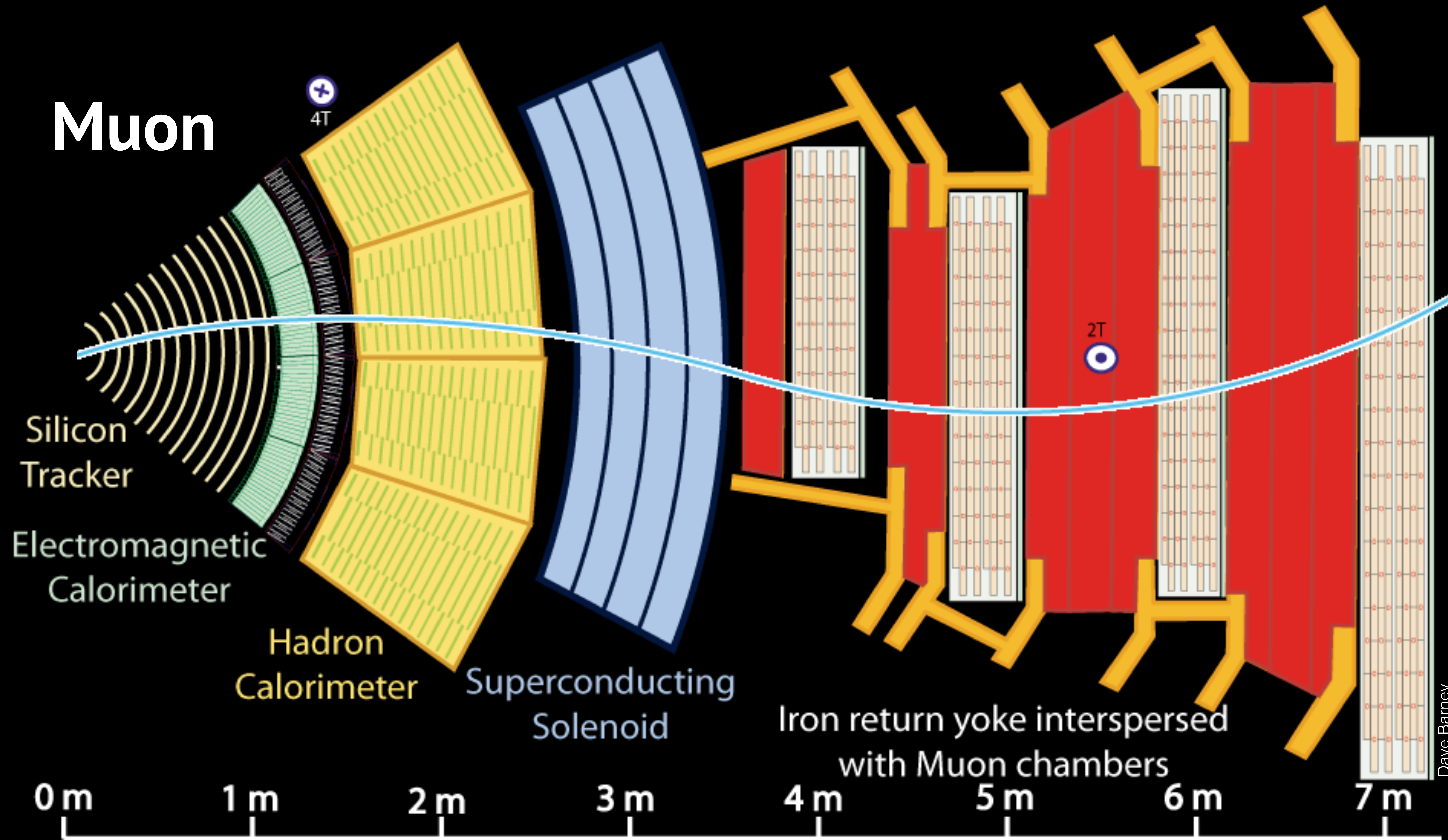
Primary vertex resolution: ~25-100 mm

Timing detector resolution: ~30-300 ps

2. PARTICLE RECONSTRUCTION

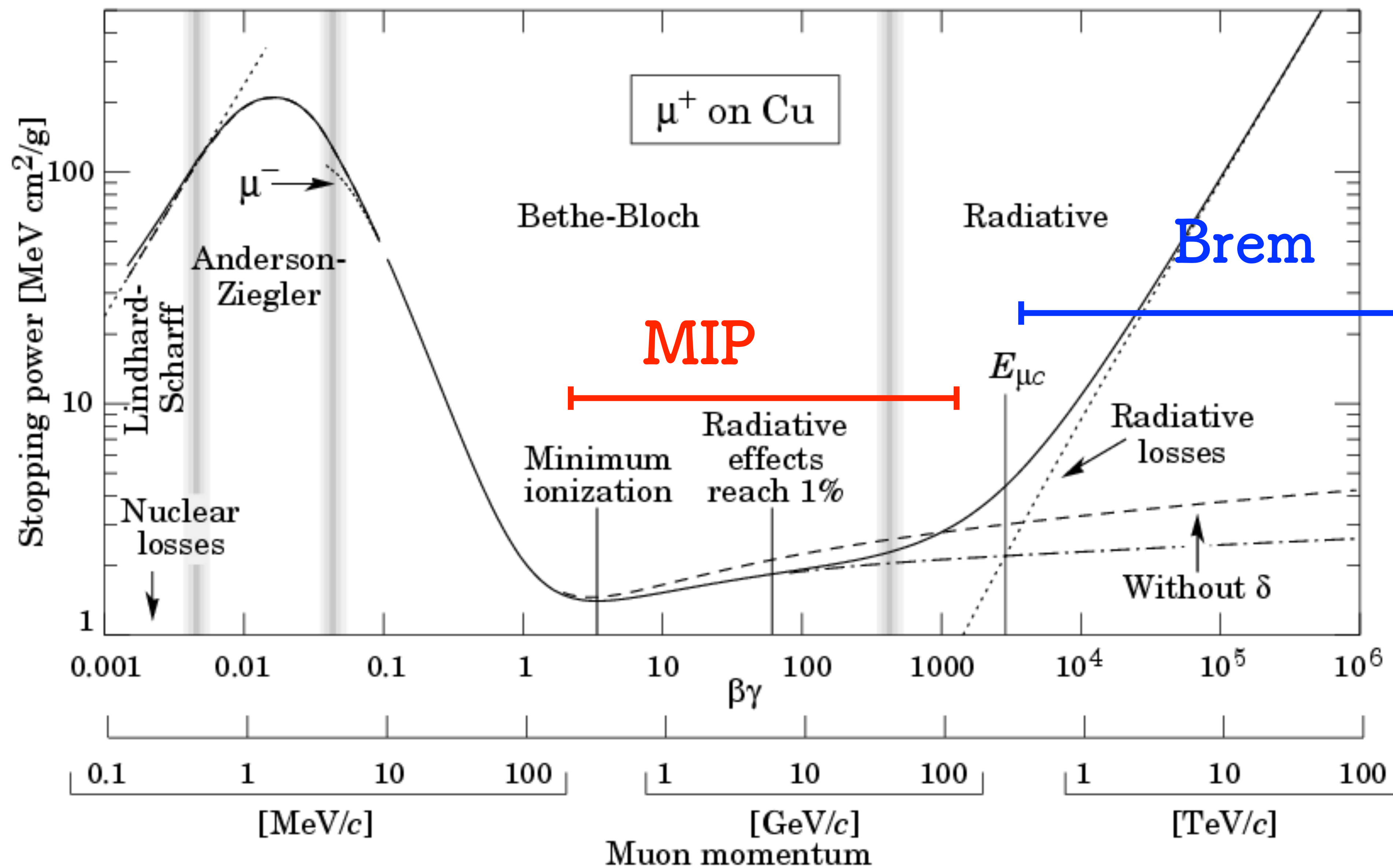


Muon

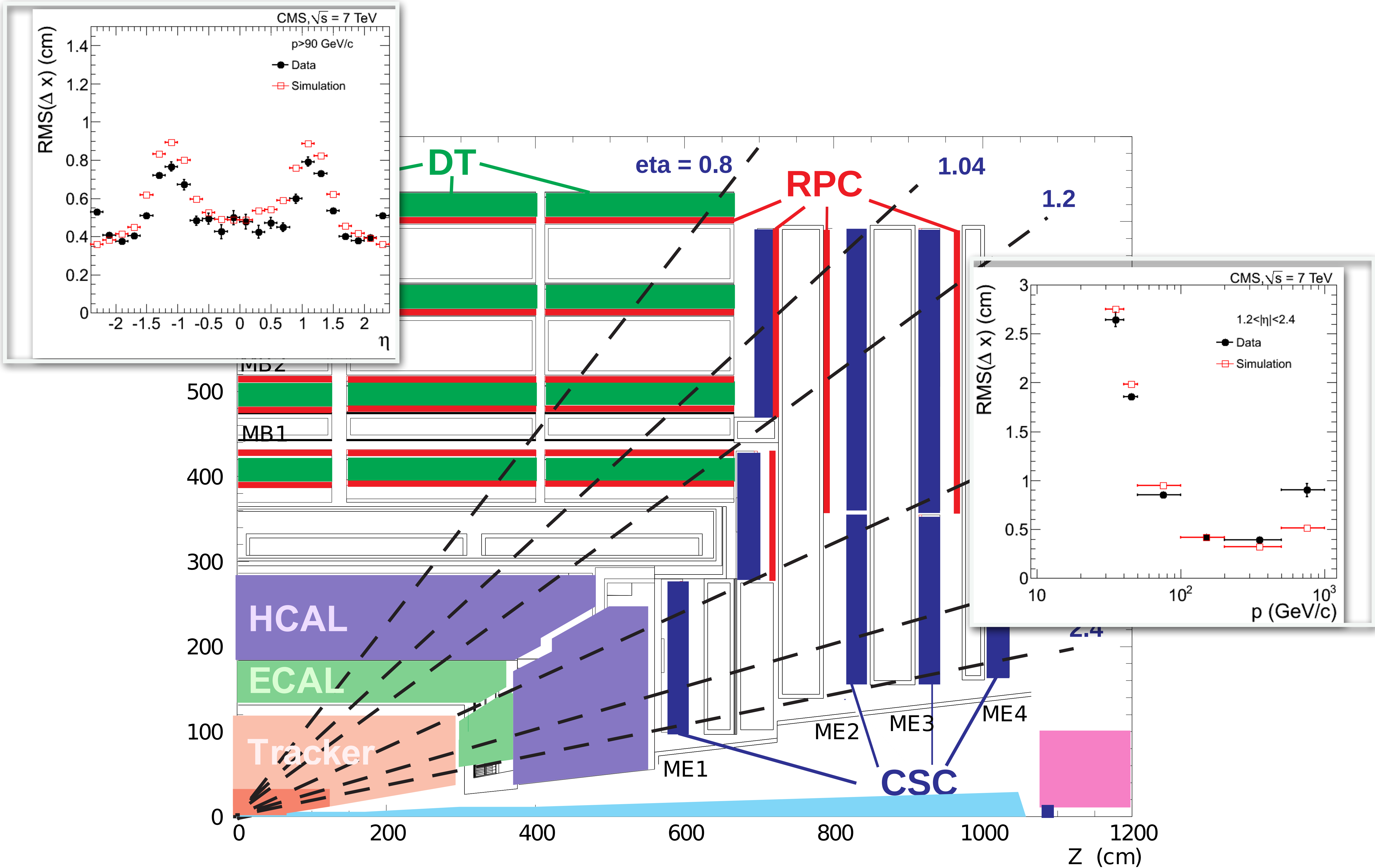


Because of its long lifetime — the muon is a stable particle for our purposes ($c \tau = 700\text{m}$)

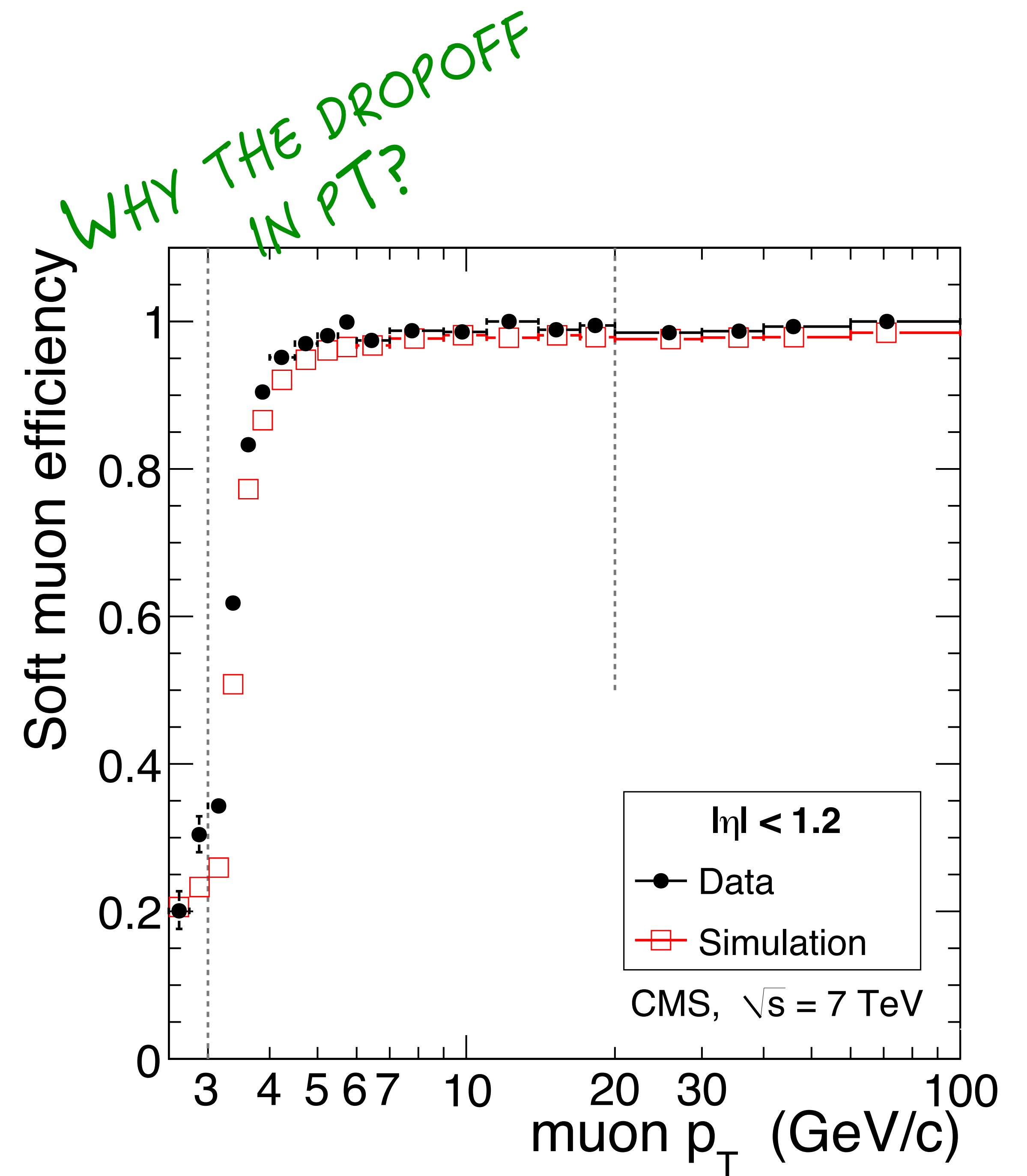
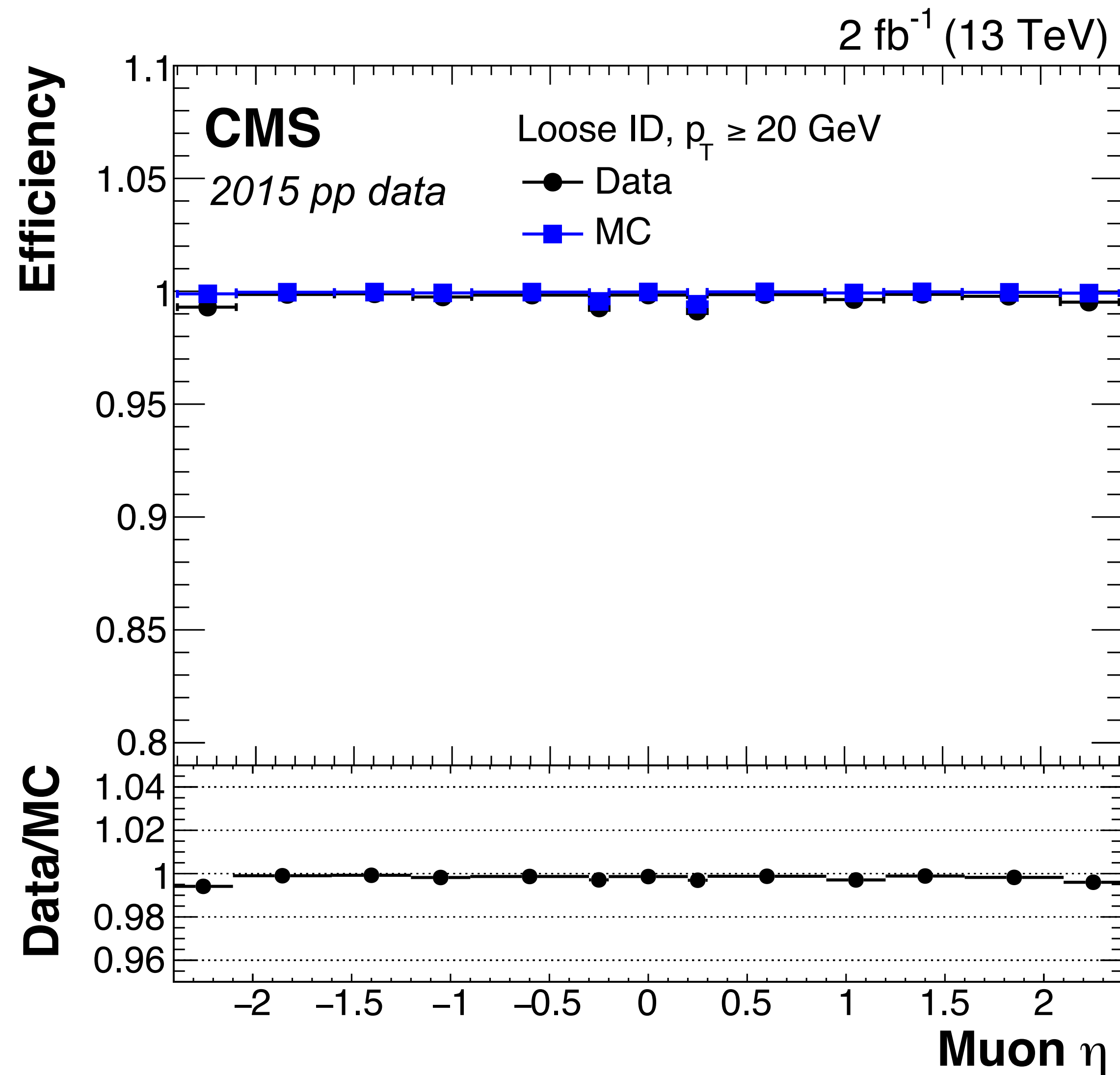
It does not feel the strong interaction, so it's only minimum ionizing particle ... except at high energies where it acts like an electron ($> 1 \text{ TeV}$)



MUON DETECTORS



Muons are very penetrating and primarily interact as a MIP
Very high ID efficiency!

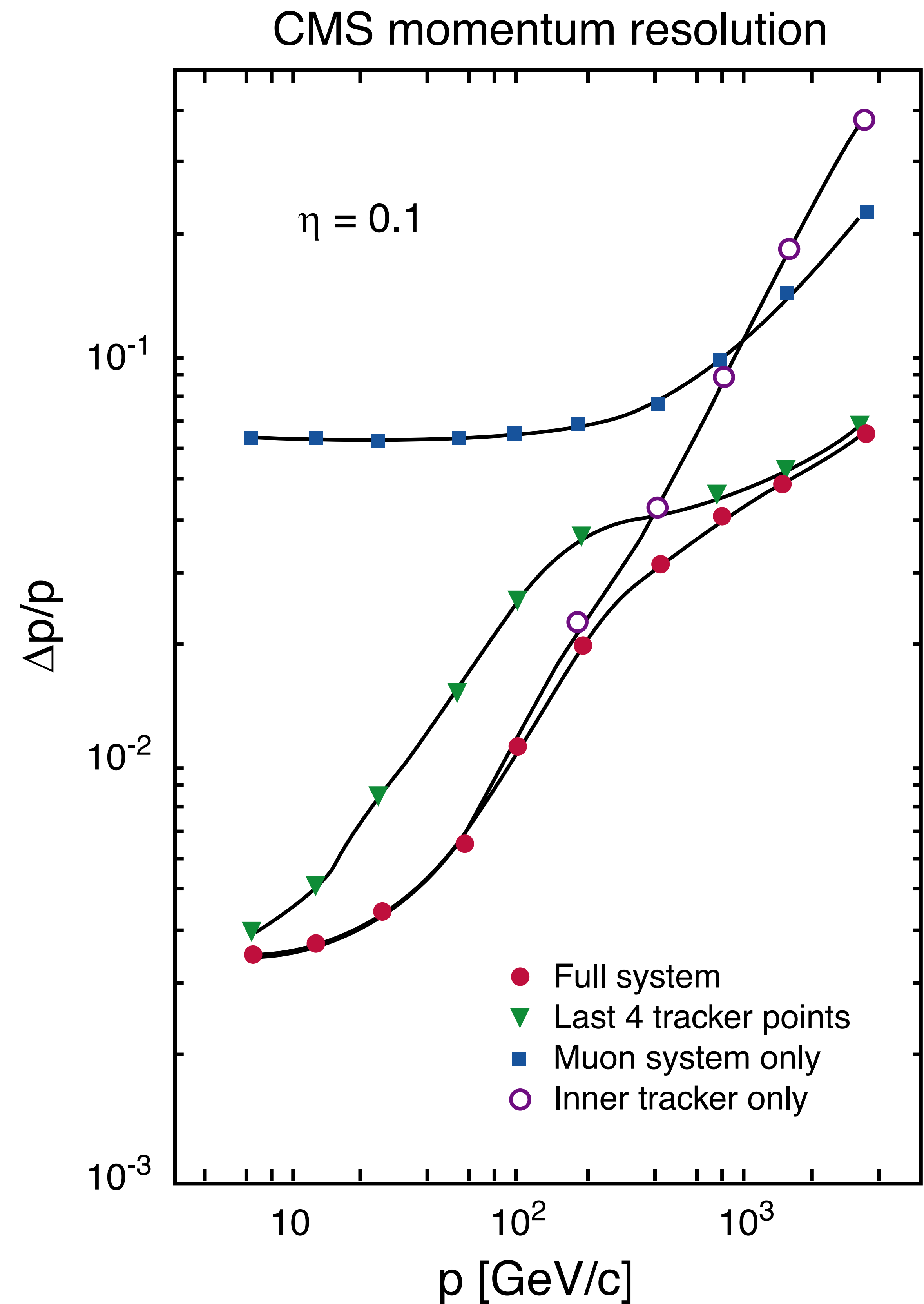
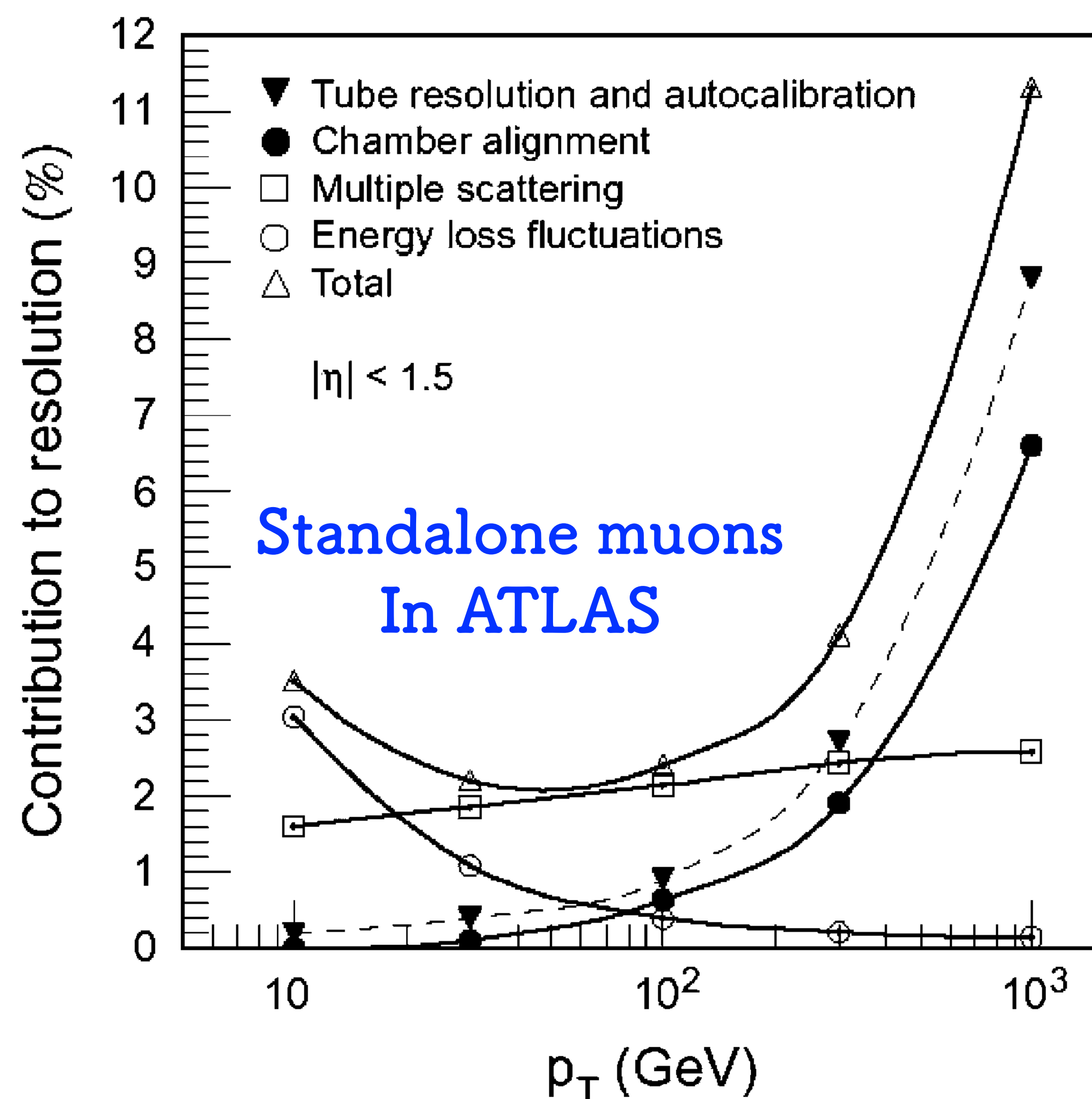


MUON MOMENTUM RESOLUTION

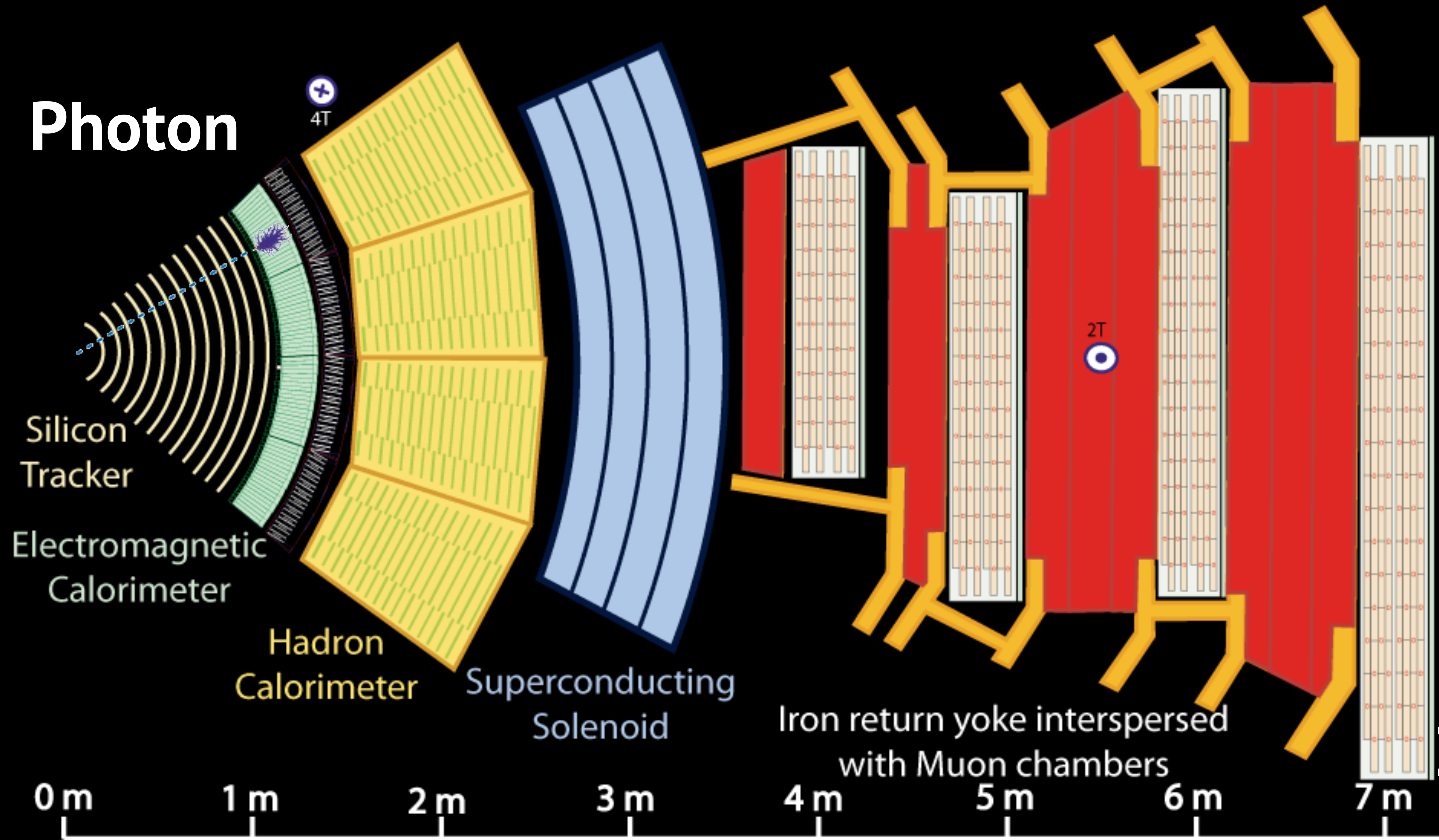
The muon system should be very efficient for identifying muons

Momentum measurements important at the **trigger level**

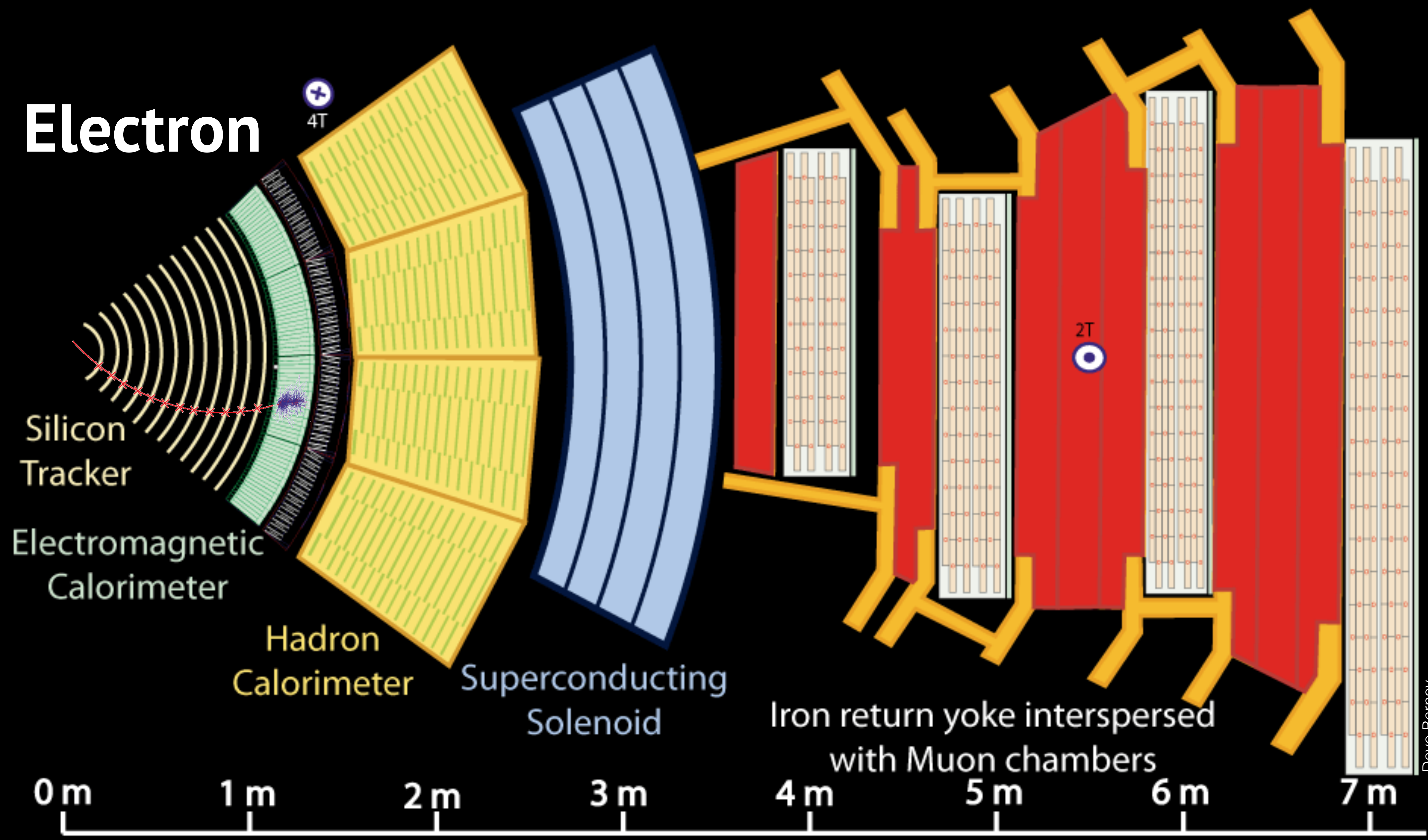
And also for **high pT muons**



Photon



Electron



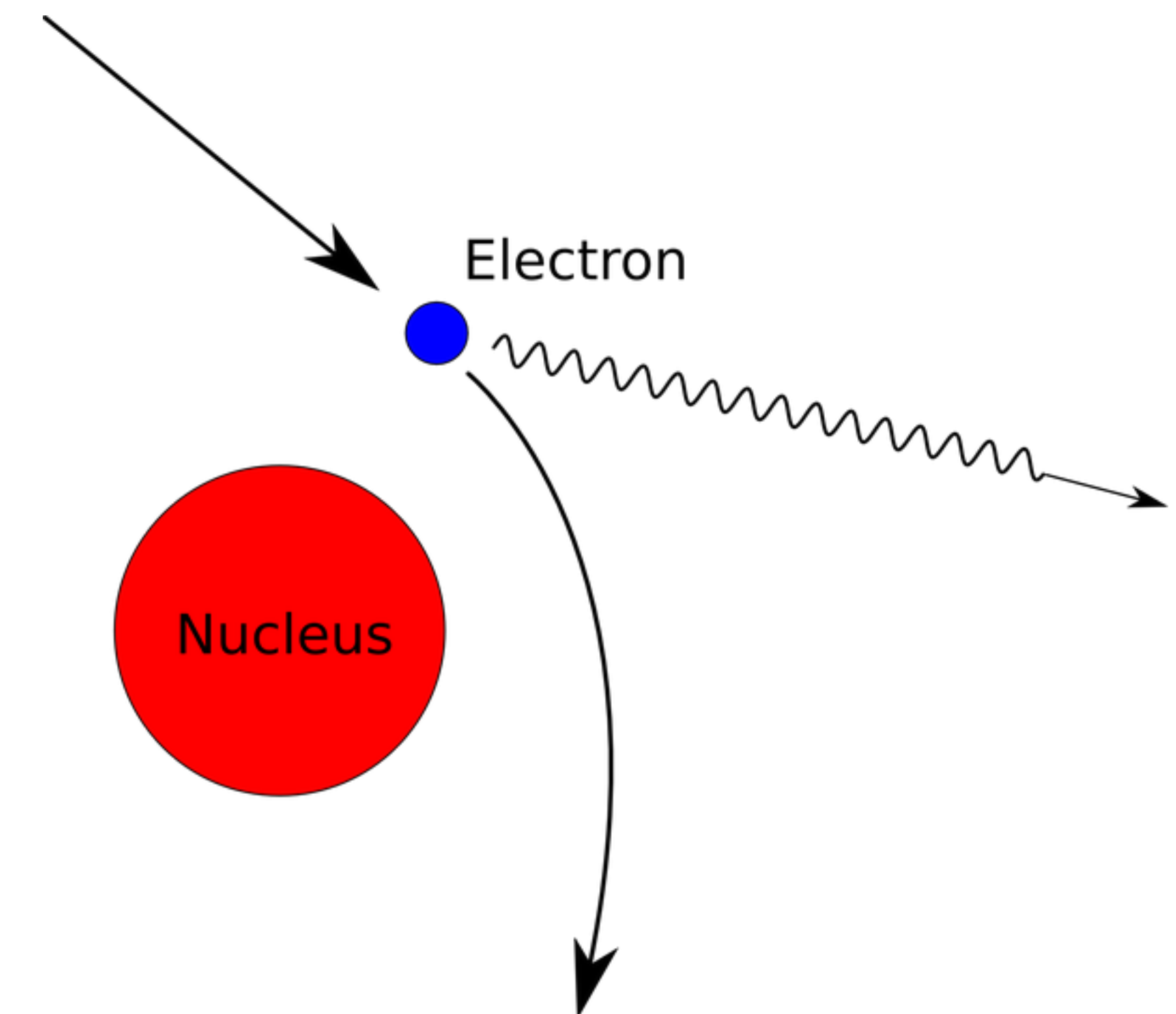
The problem with electrons...

They interact a lot more! Primarily through **bremsstrahlung**

Energy loss from bremsstrahlung:

(energy loss is proportional to energy)

$$-\frac{dE}{dx} = \frac{E}{X_0}$$



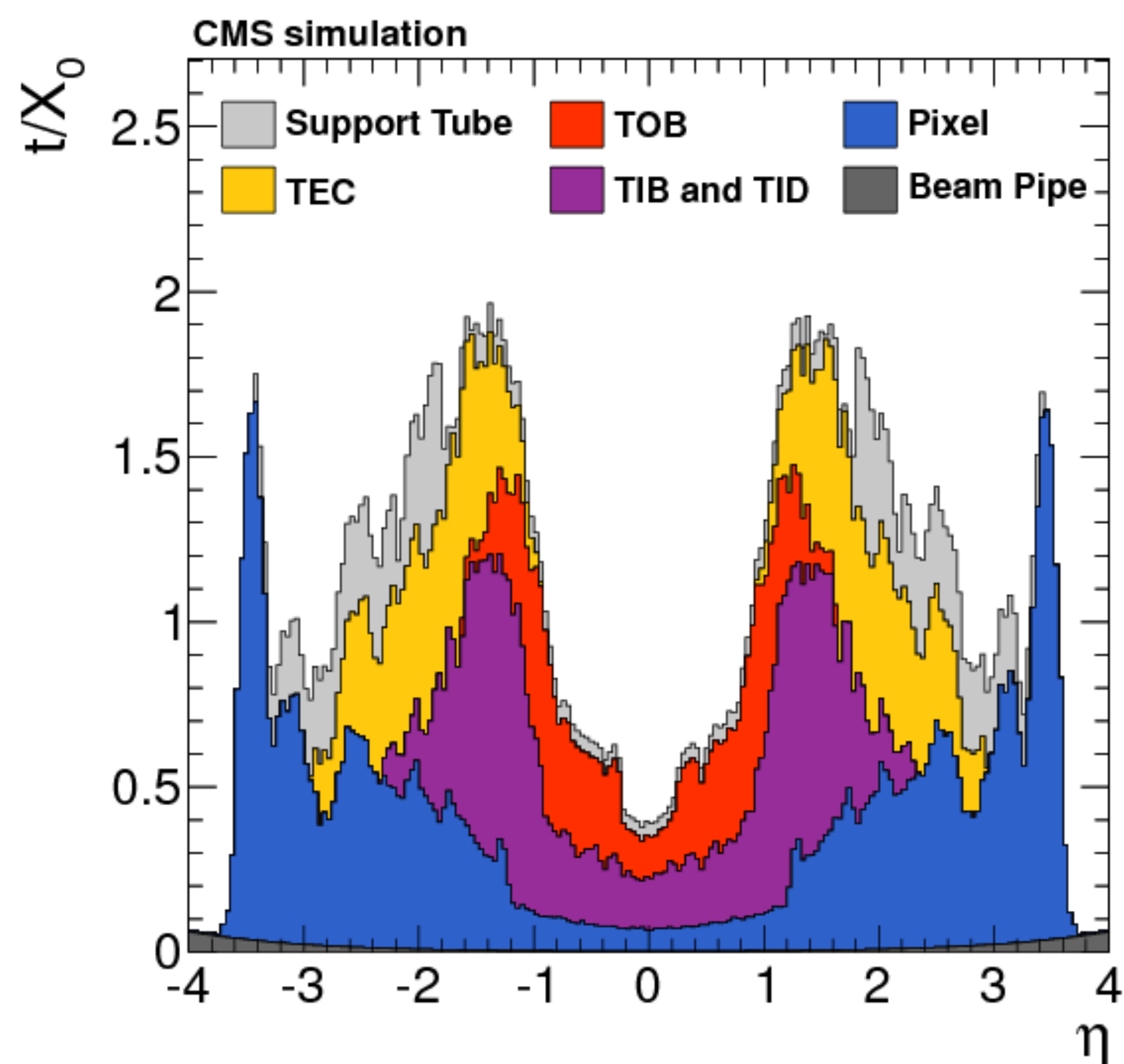
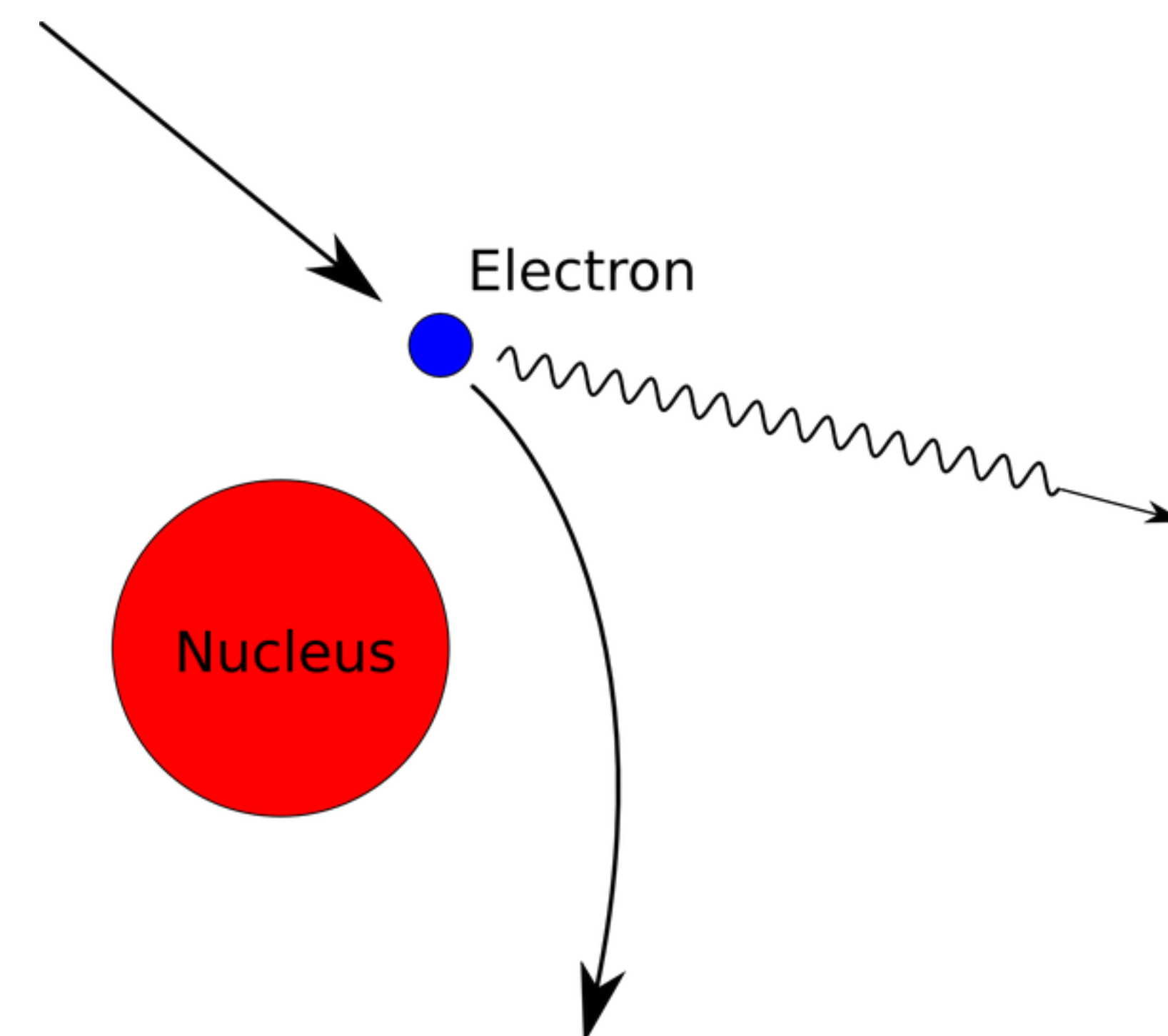
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Mind your material!

Important to consider the material budget in the tracker detector design

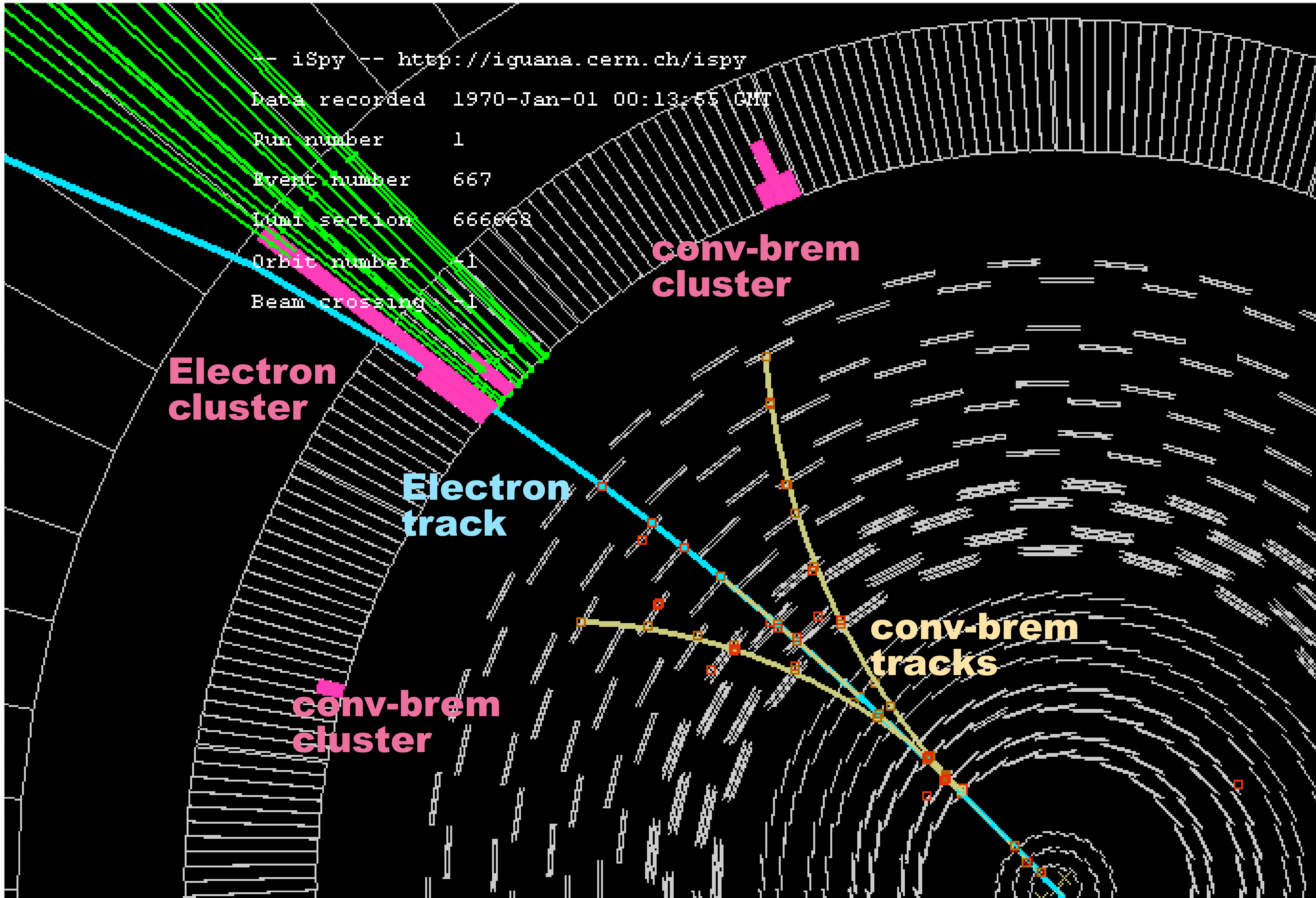
COMPLICATIONS WITH ELECTRONS

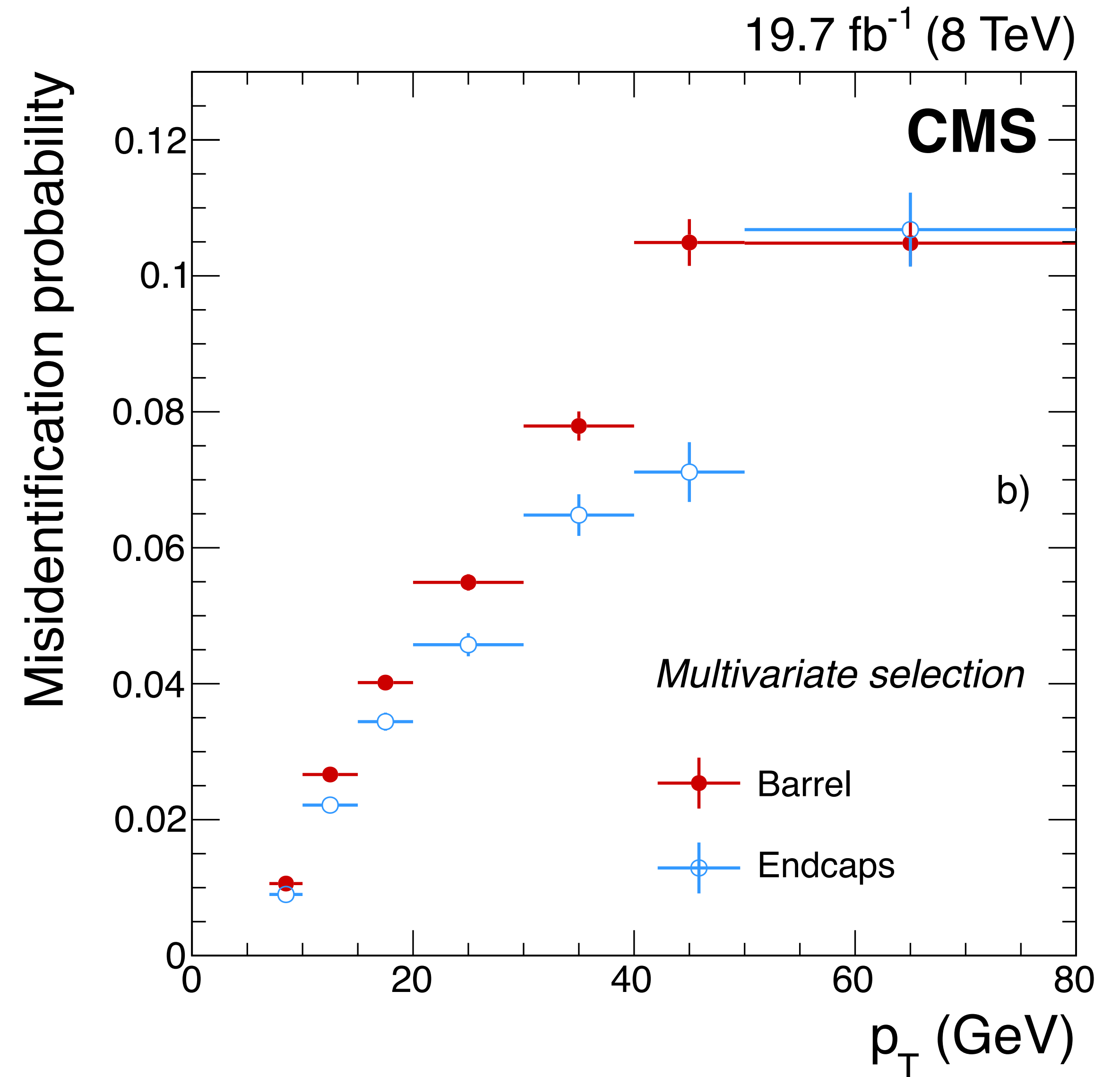
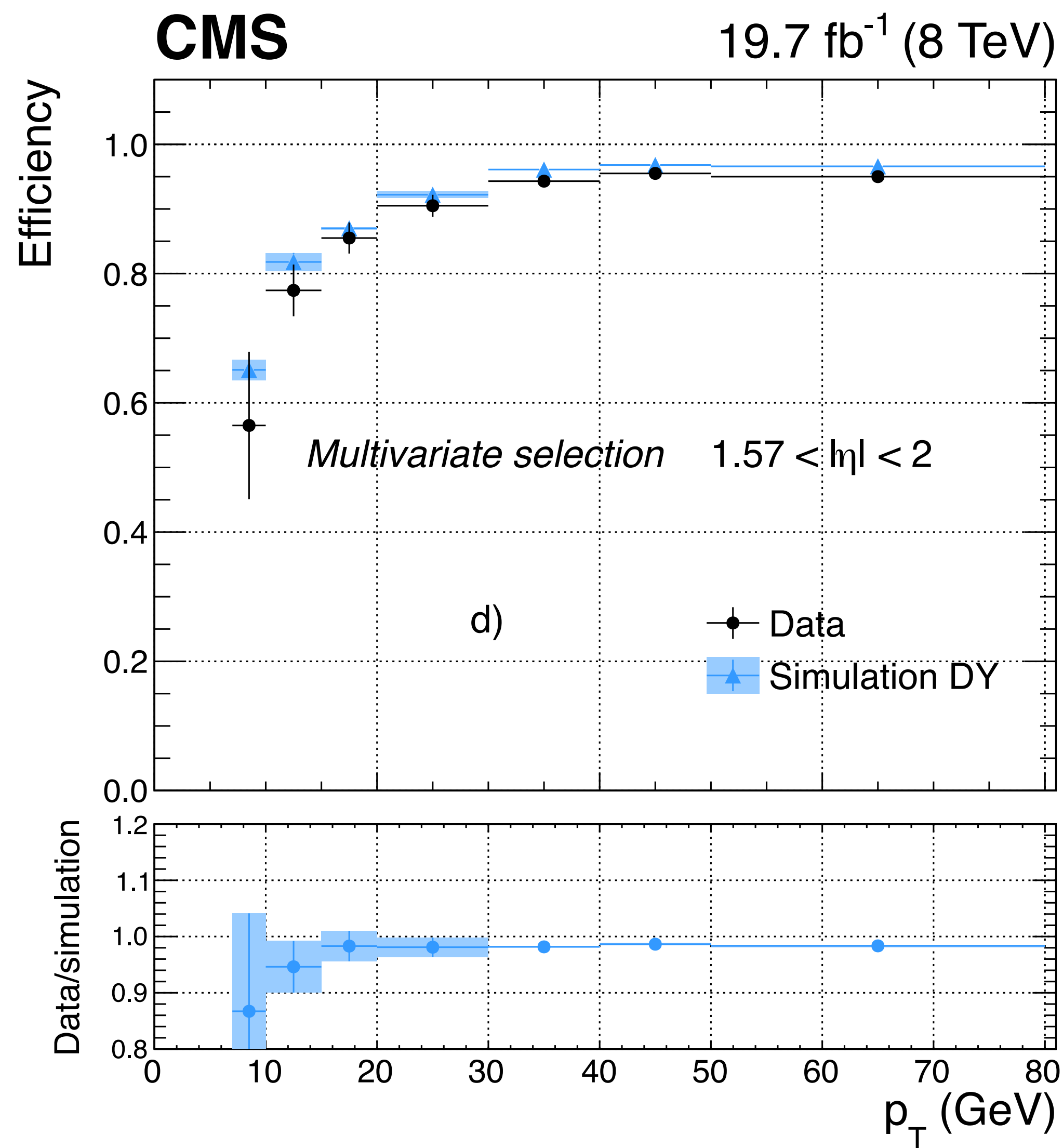
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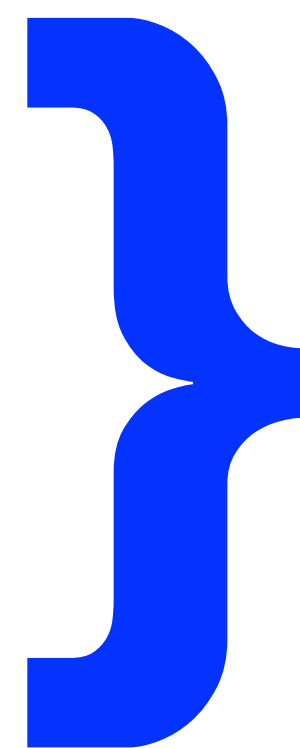
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$|\Delta\eta|$
 $|\Delta\phi|$
 H/E_{SC}
 $\sigma_{\eta\eta}$
 $|1/E_{SC} - 1/p|$
 $ISO_{PF} (\Delta R=0.3) / p_T$
 $|d_0|$
 $|d_z|$
 Missing hits
 Conversion-fit probability



What variables go into the selection?

Identifying **prompt** and **isolated** photons important

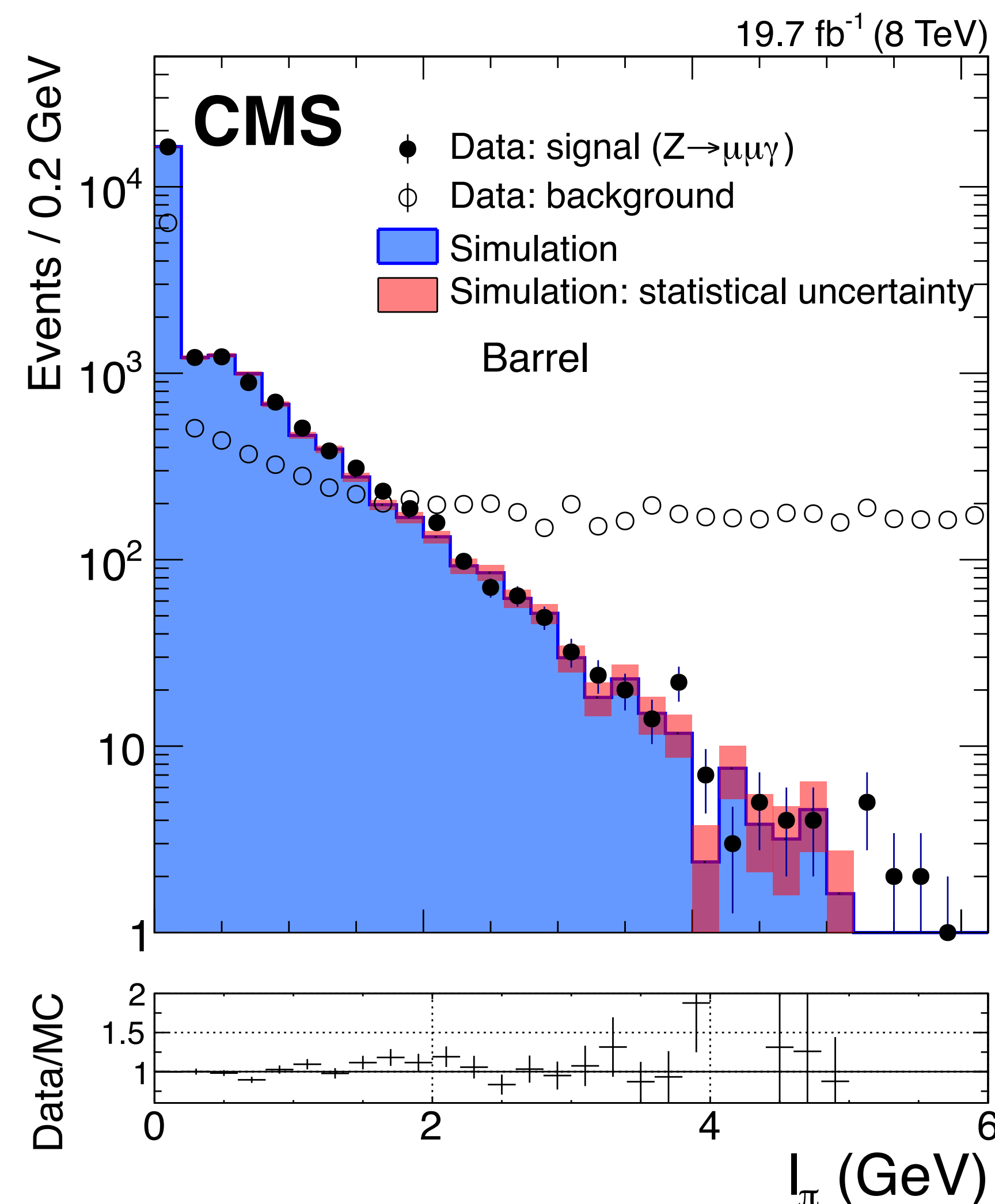
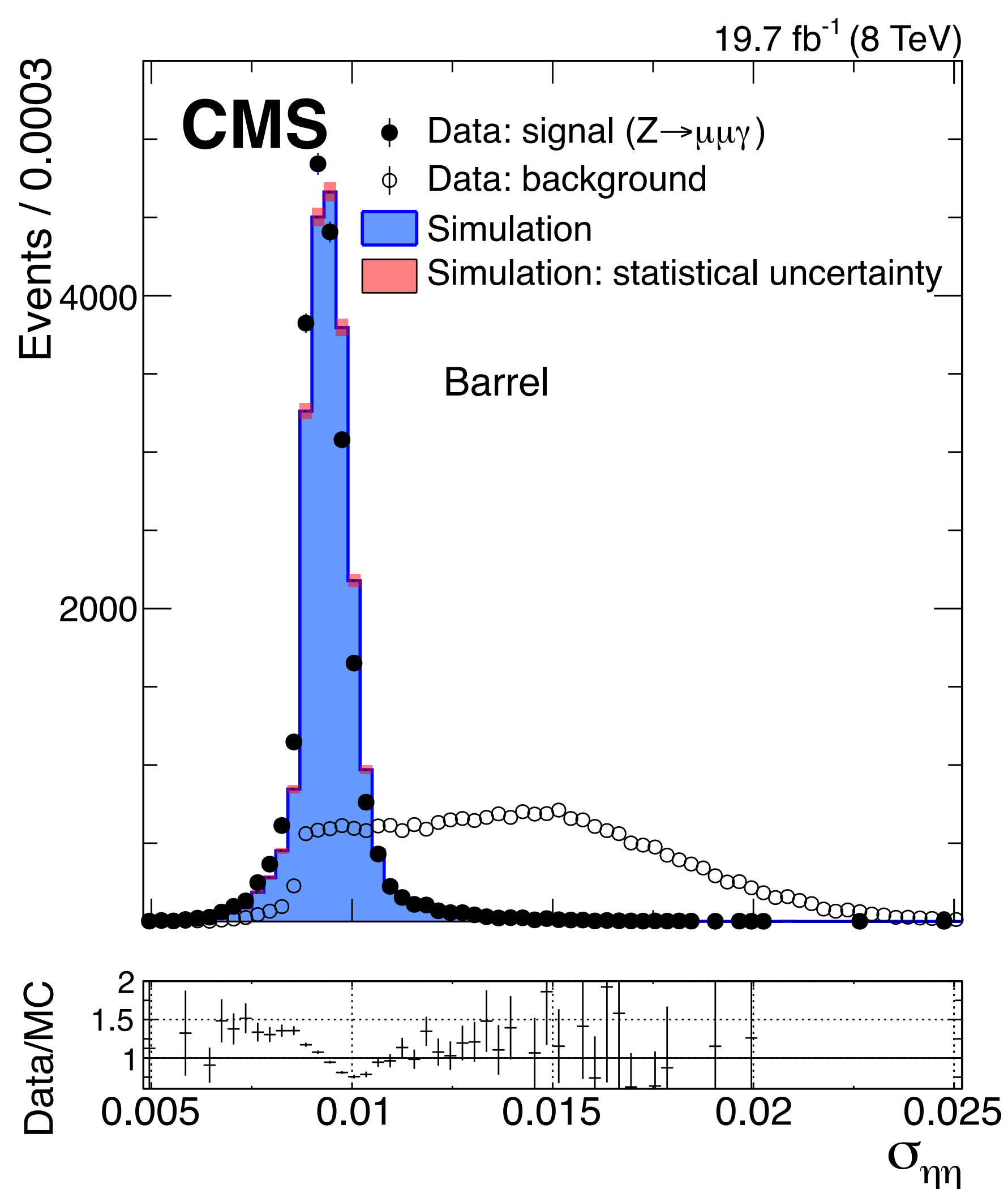
Particularly for analyses like $H(\gamma\gamma)$

Primary variables for photon identification are shower-shape and isolation (more on this later) variables

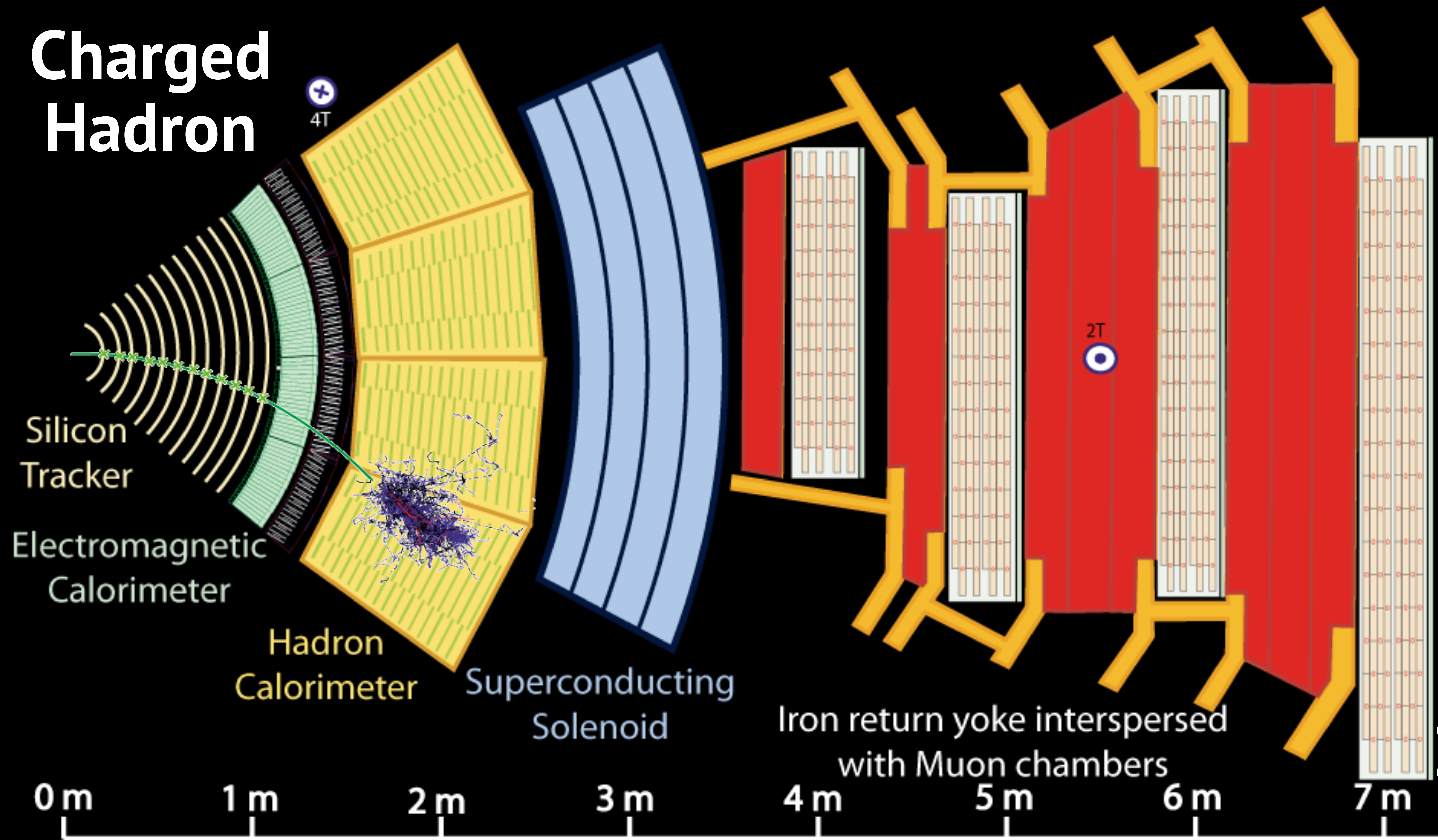
No matched track to separate from electrons

signal
Isolated FSR
photons from
 $Z\mu\mu$

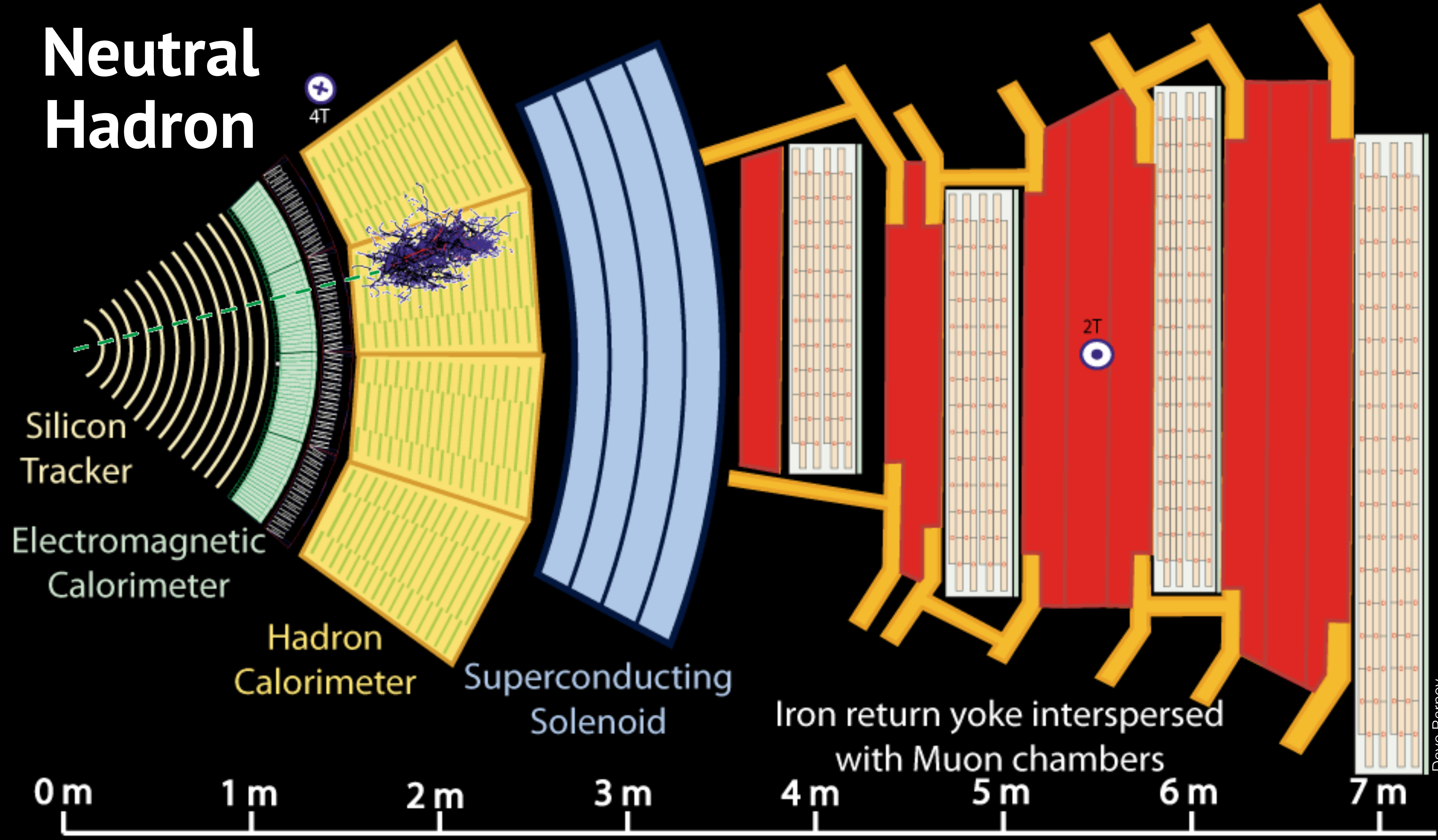
background
Photons from
jets



Charged Hadron



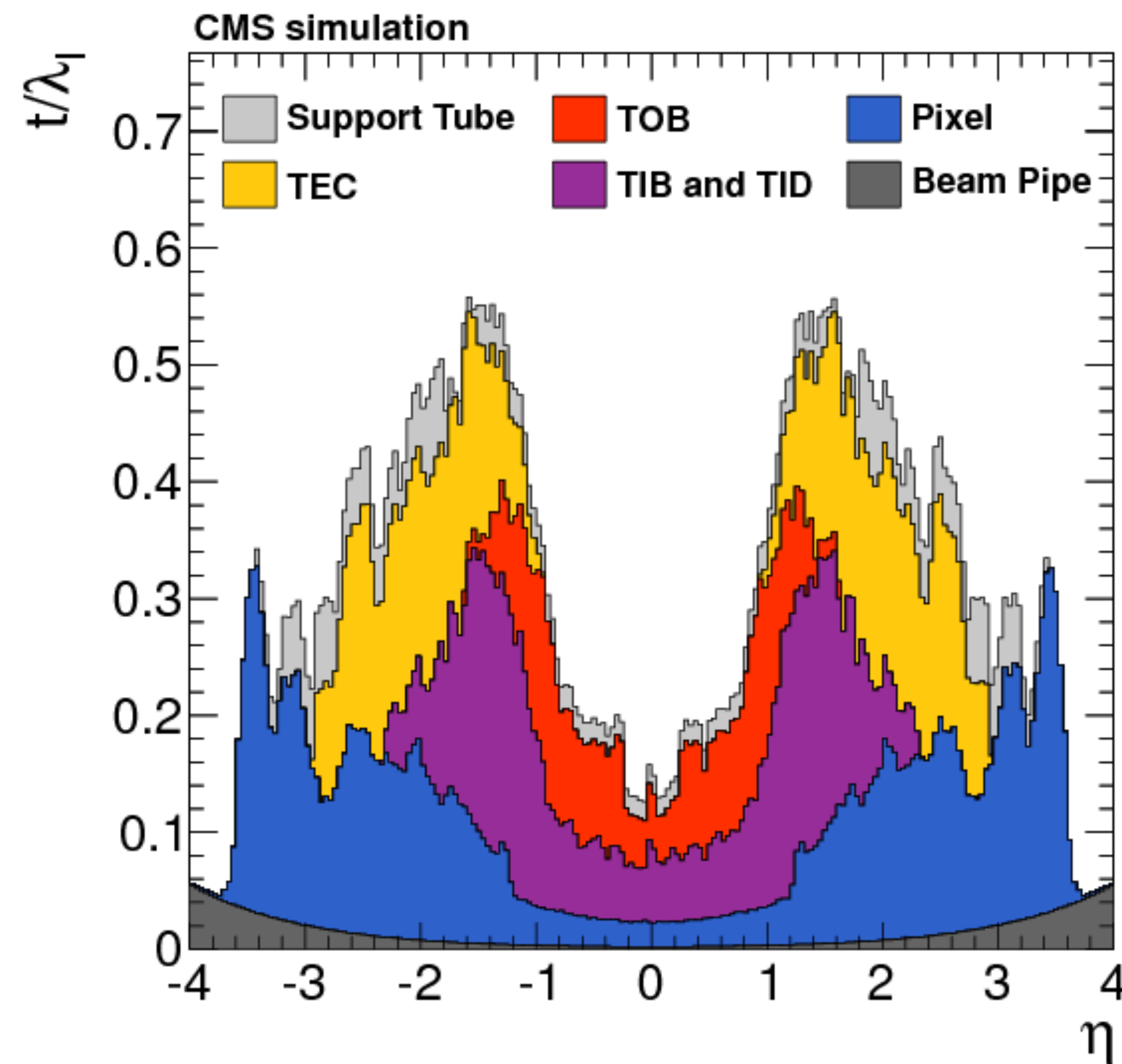
Neutral Hadron



Match tracks to hadronic clusters to form charged hadrons

Again, mind your materials!

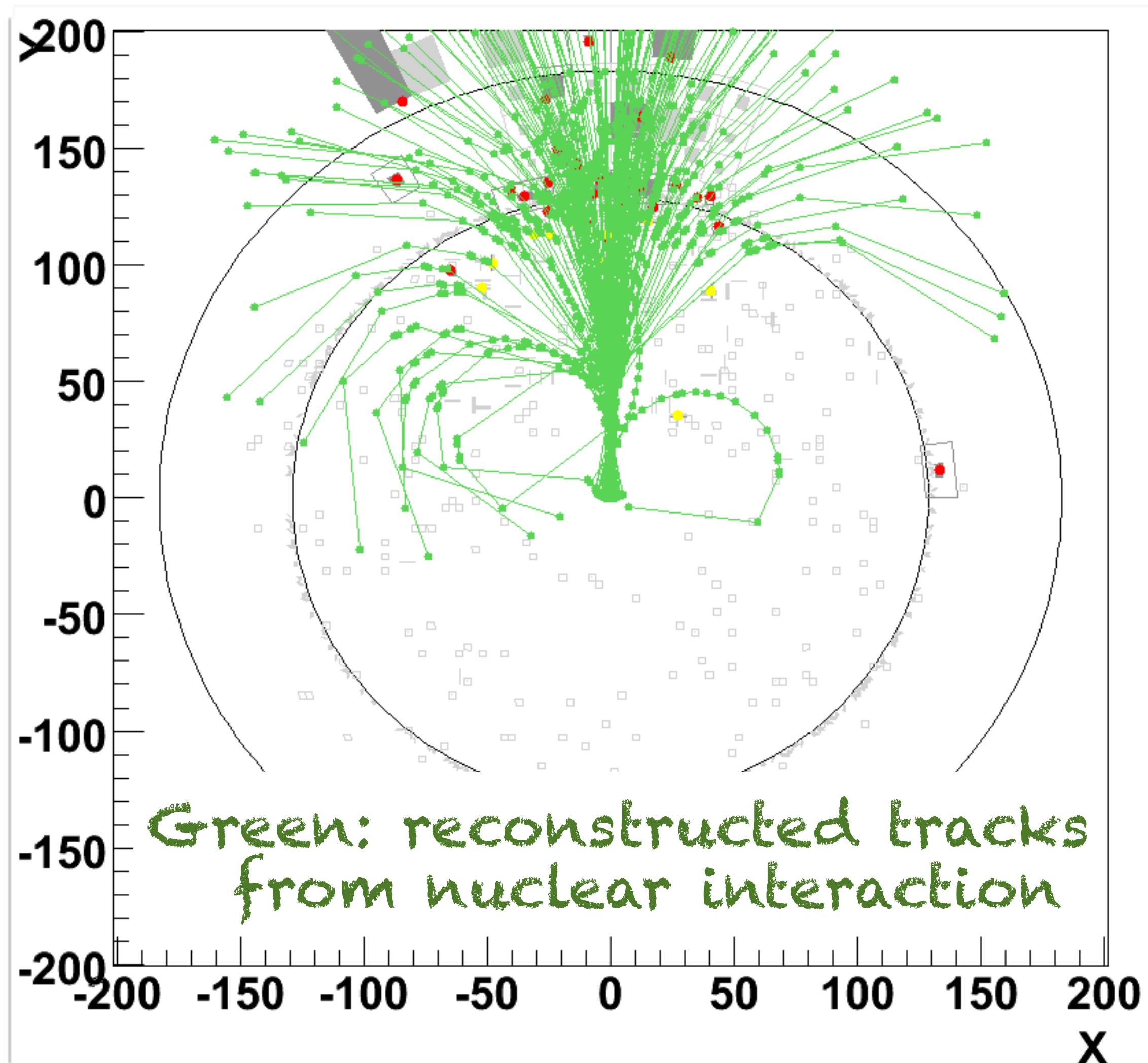
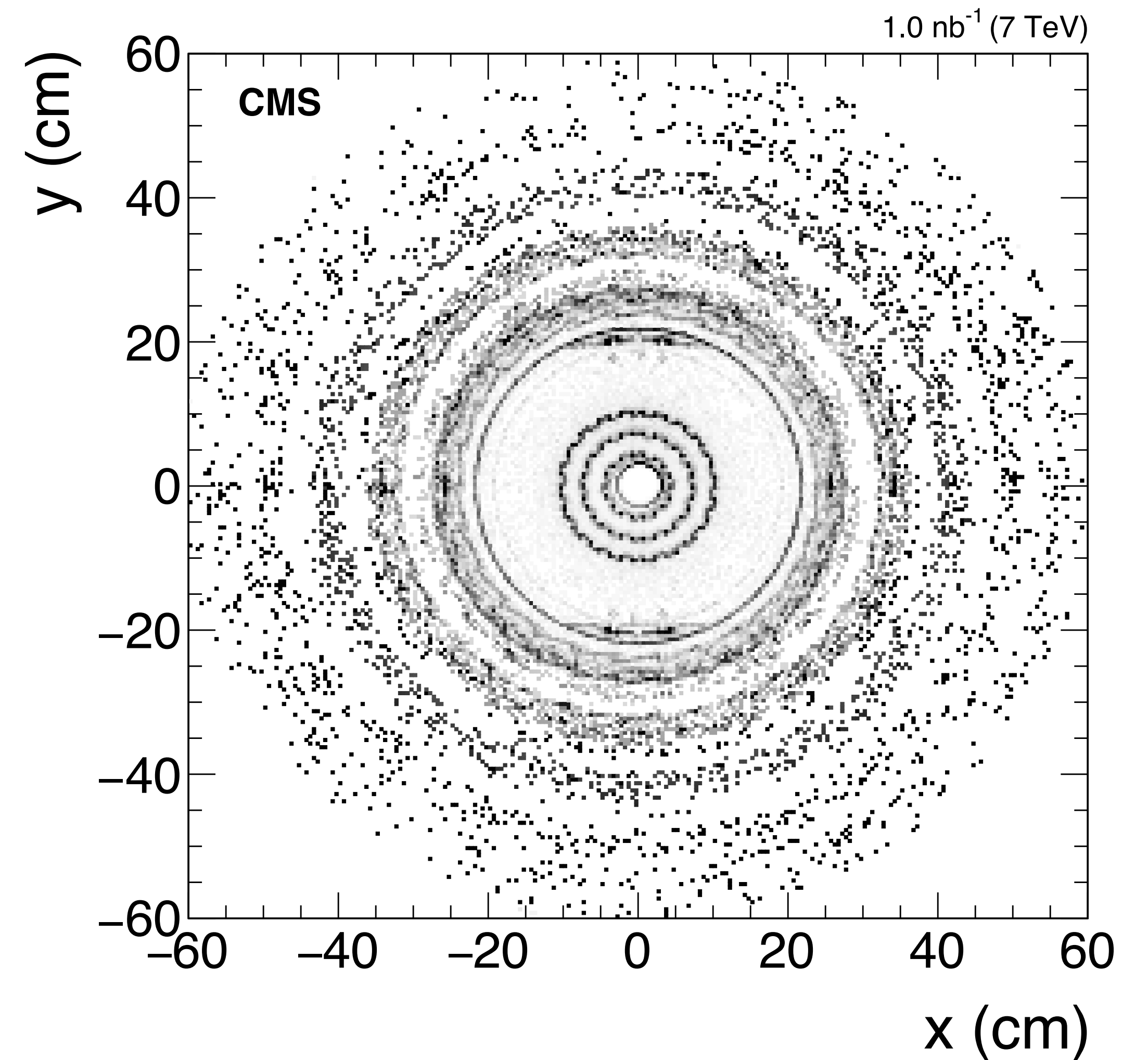
The tracker material acts as a hadronic preshower
(for both charged and neutral hadrons)



Nuclear interactions often result in kinks in the track or a production of secondary particles

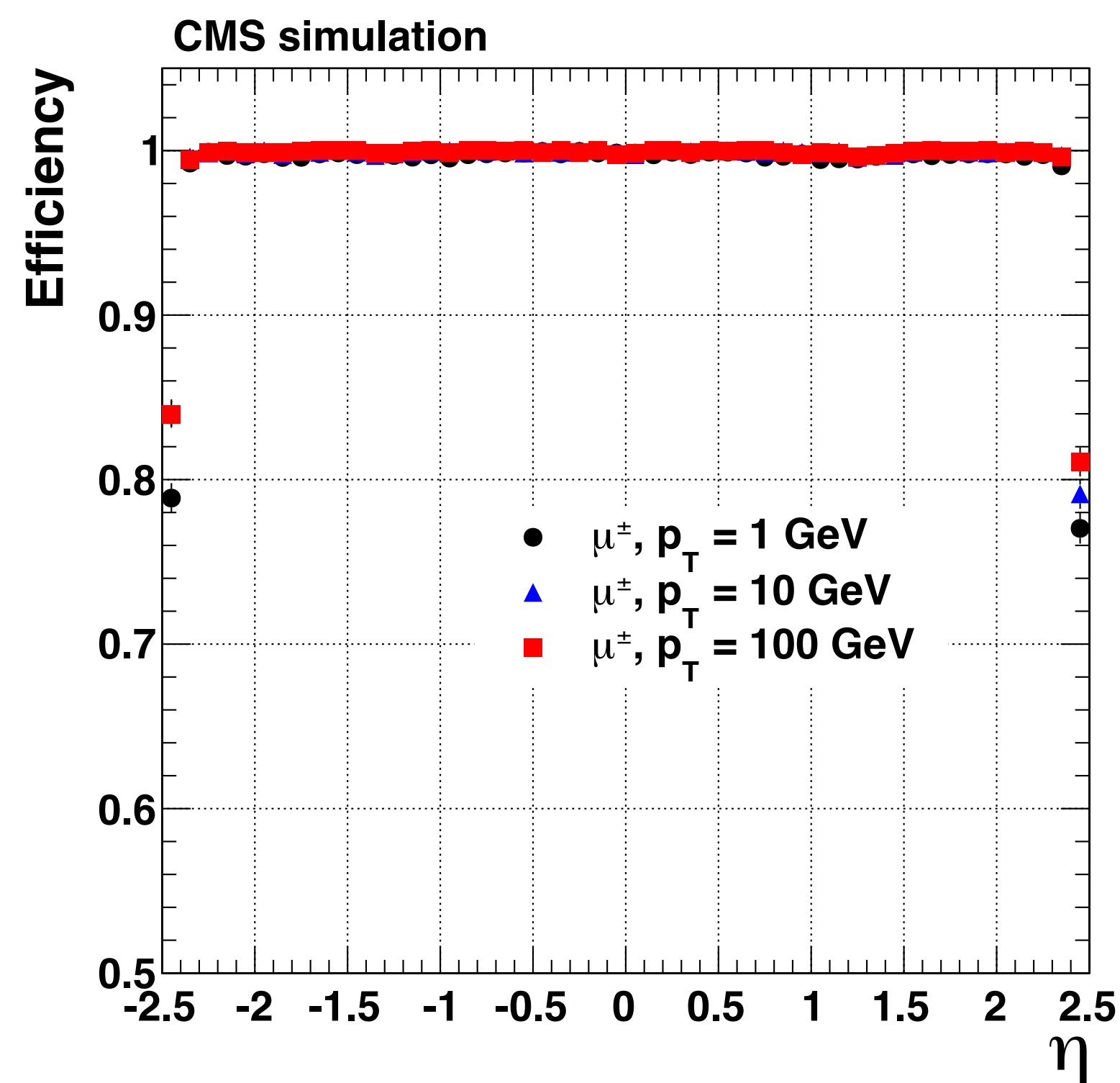
Can be recovered with displaced track reconstruction

Map of nuclear interactions

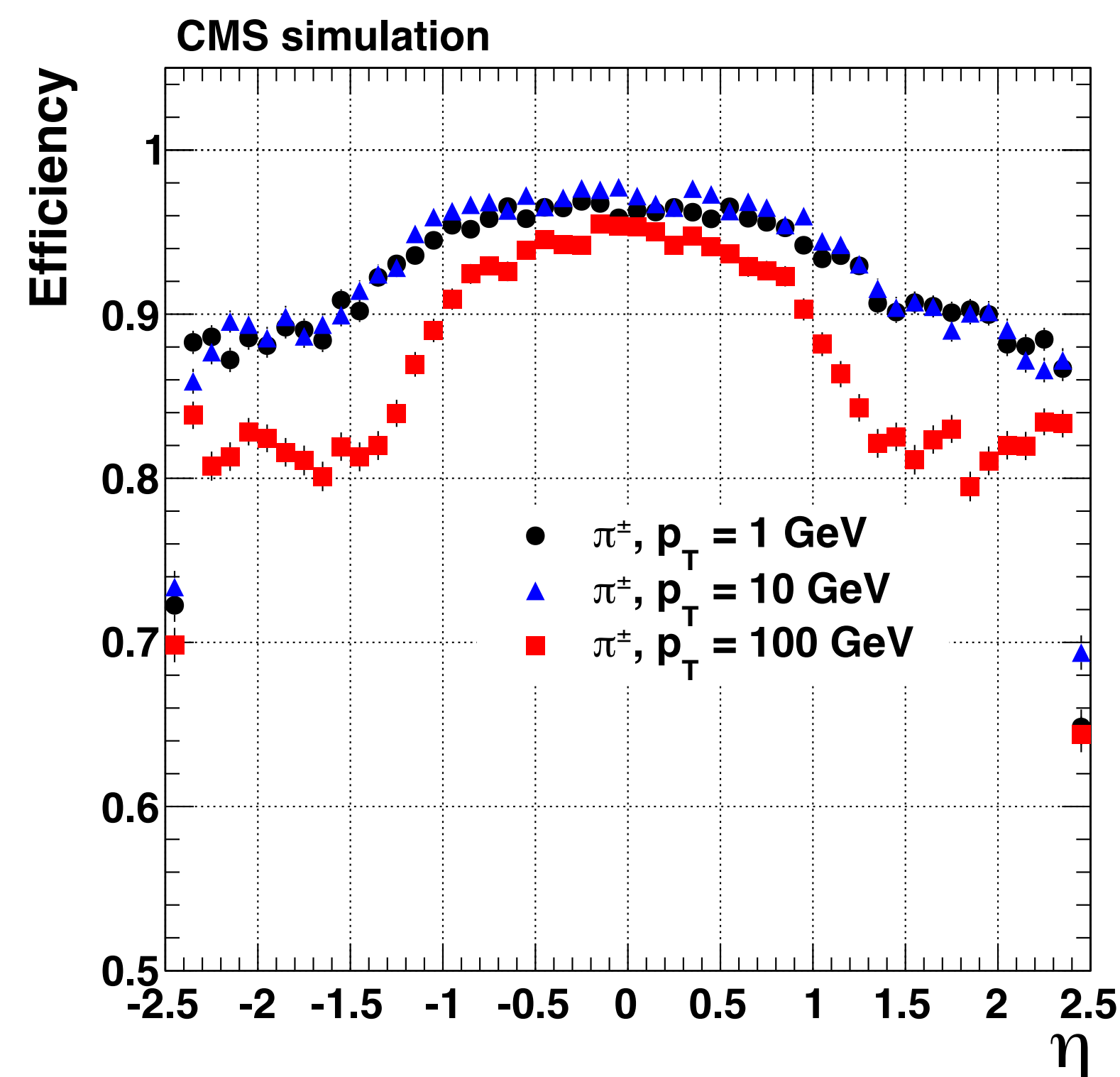


To avoid double counting, nuclear interactions need to be identified and combined into primary particles (part of particle flow, see later)

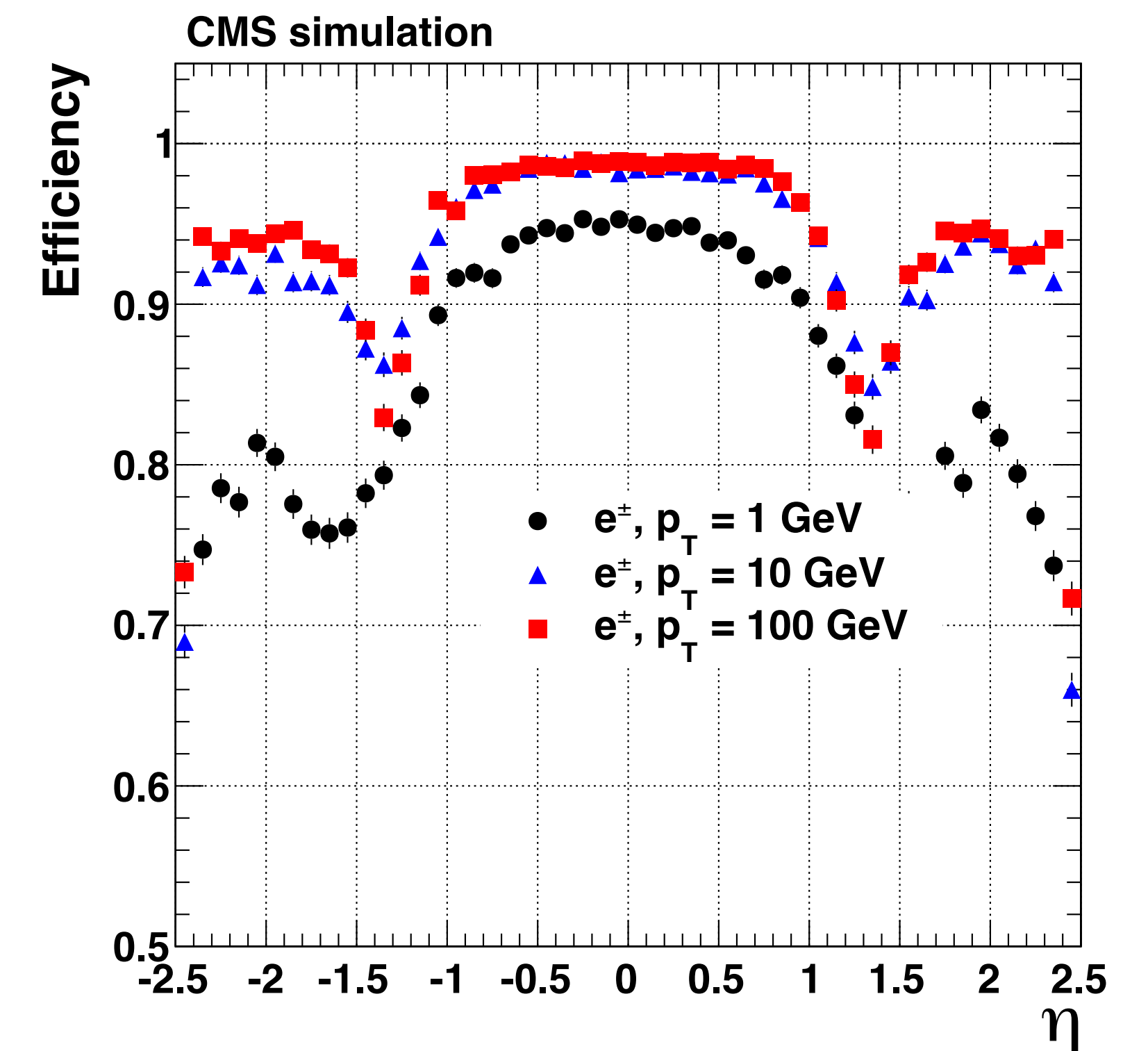
Muons



Pions



Electrons

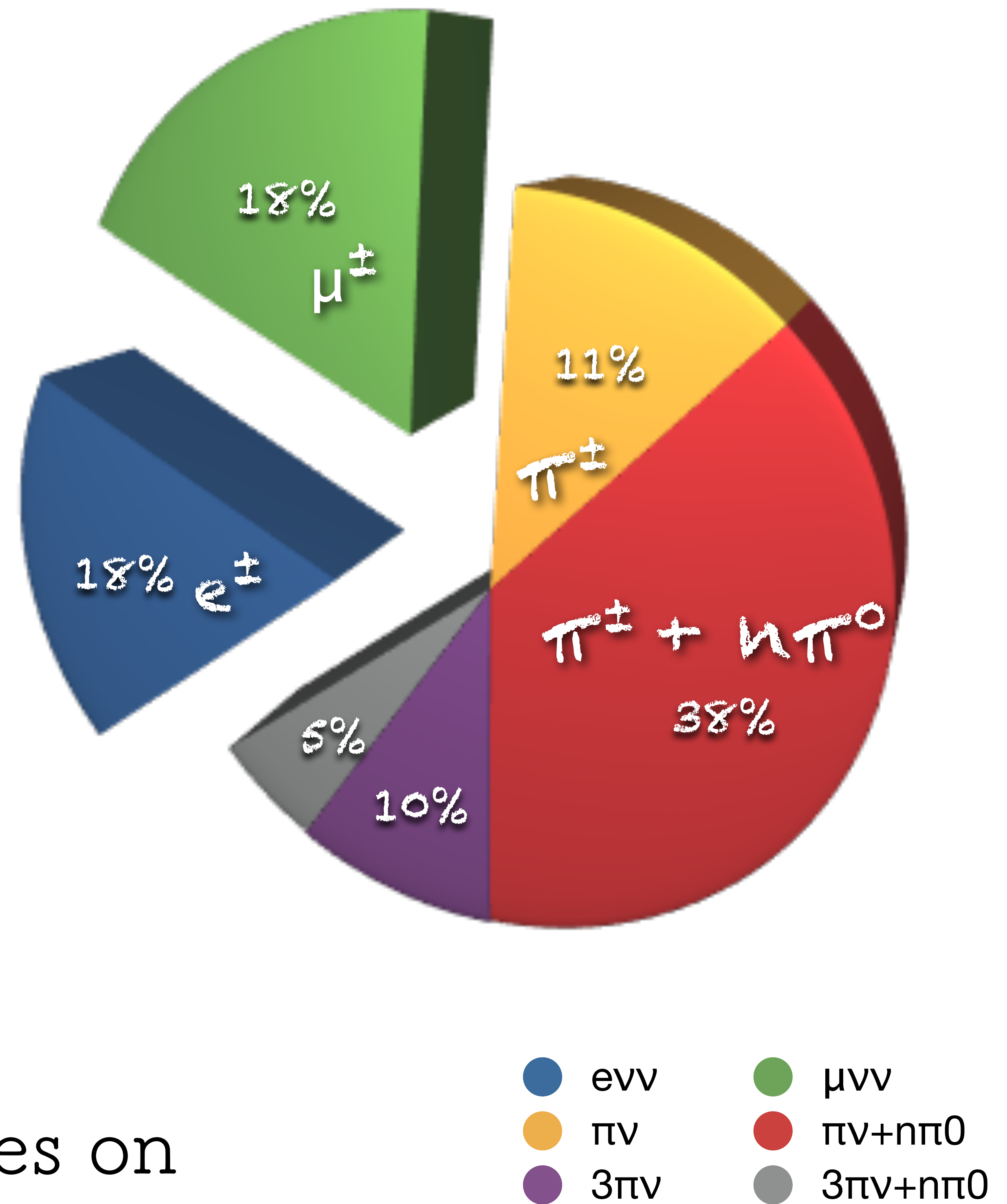
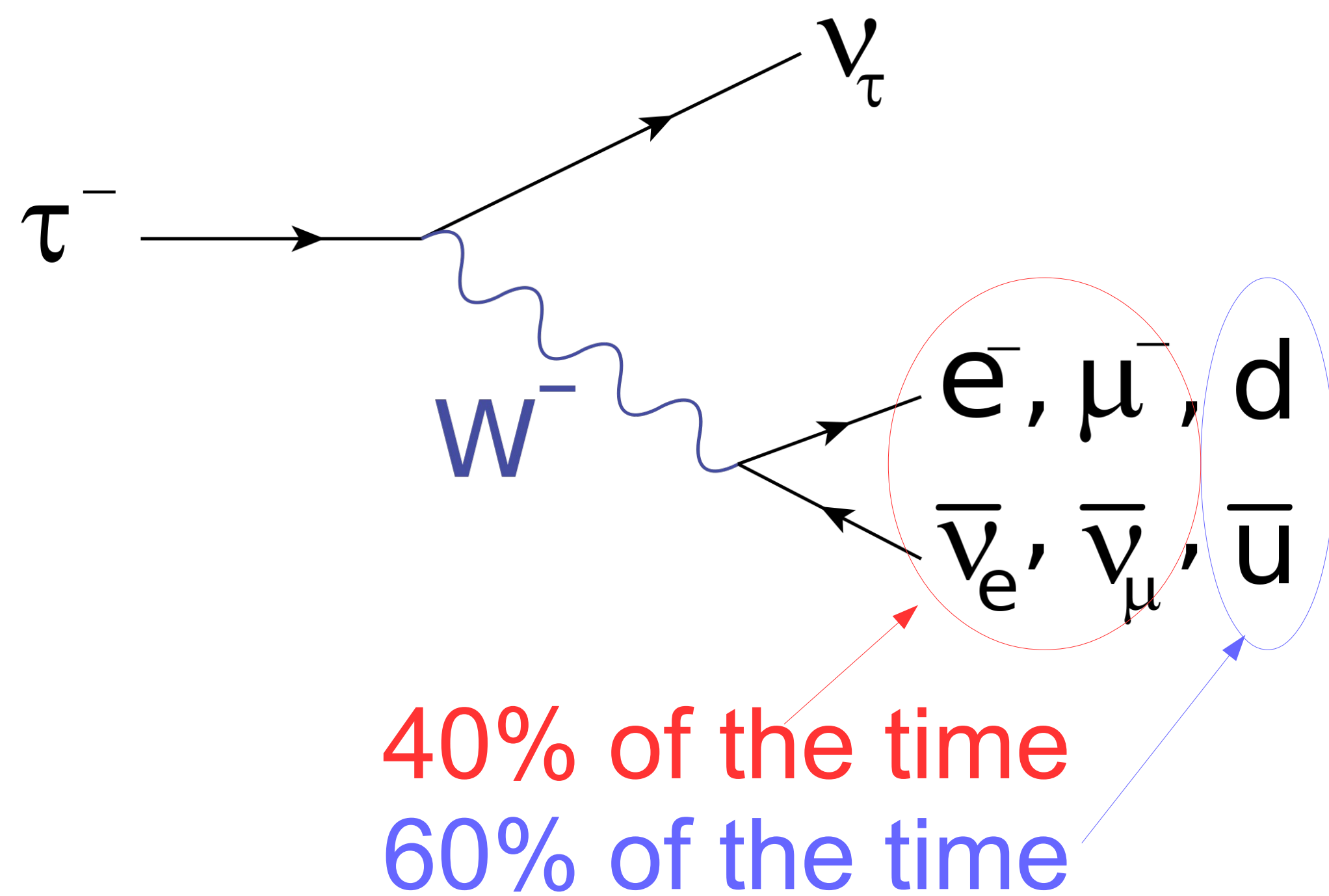


Side-by-side comparison of muon, pion, electron tracking efficiency — this illustrates the challenge of tracker material for charged hadrons and electrons

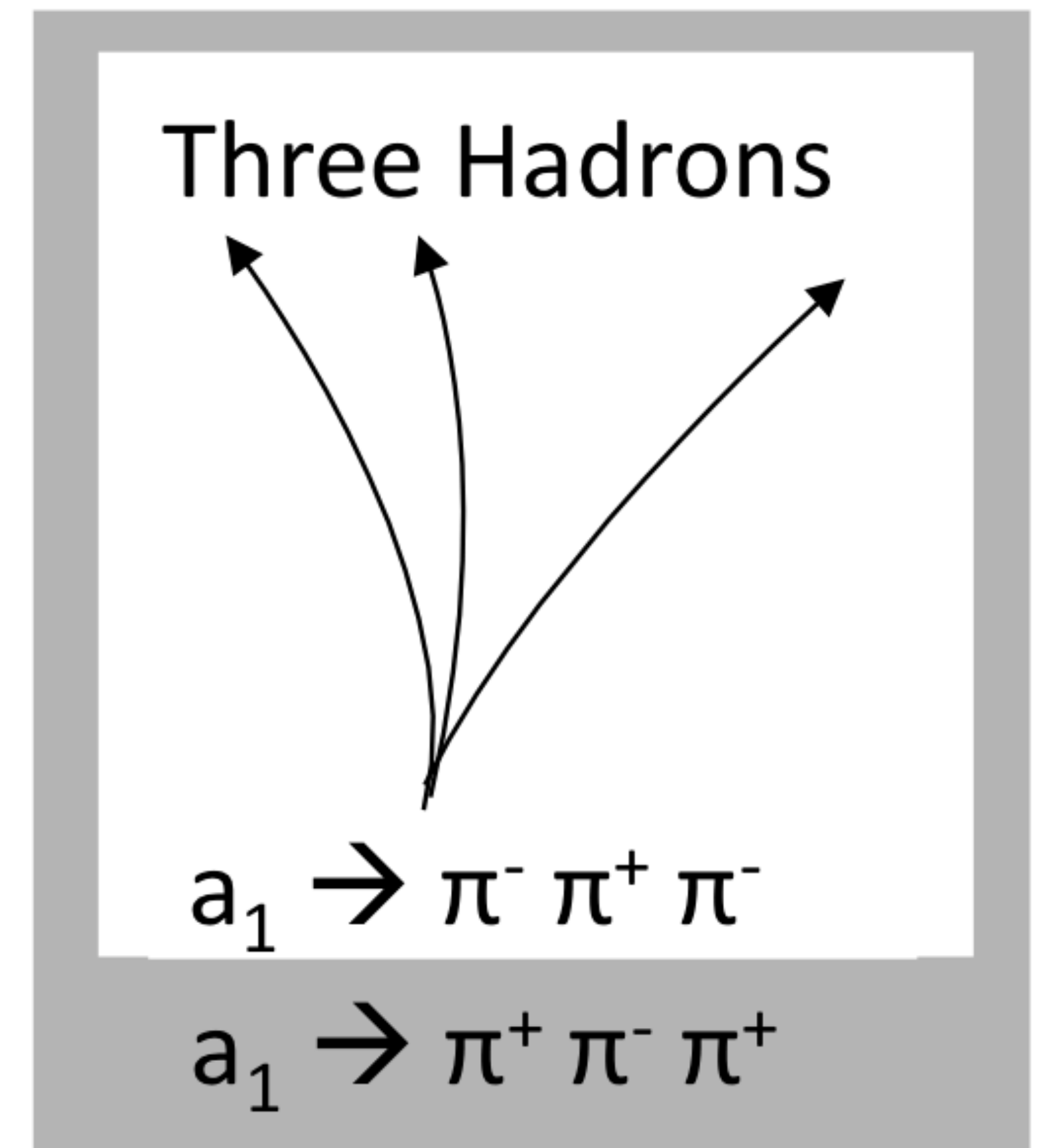
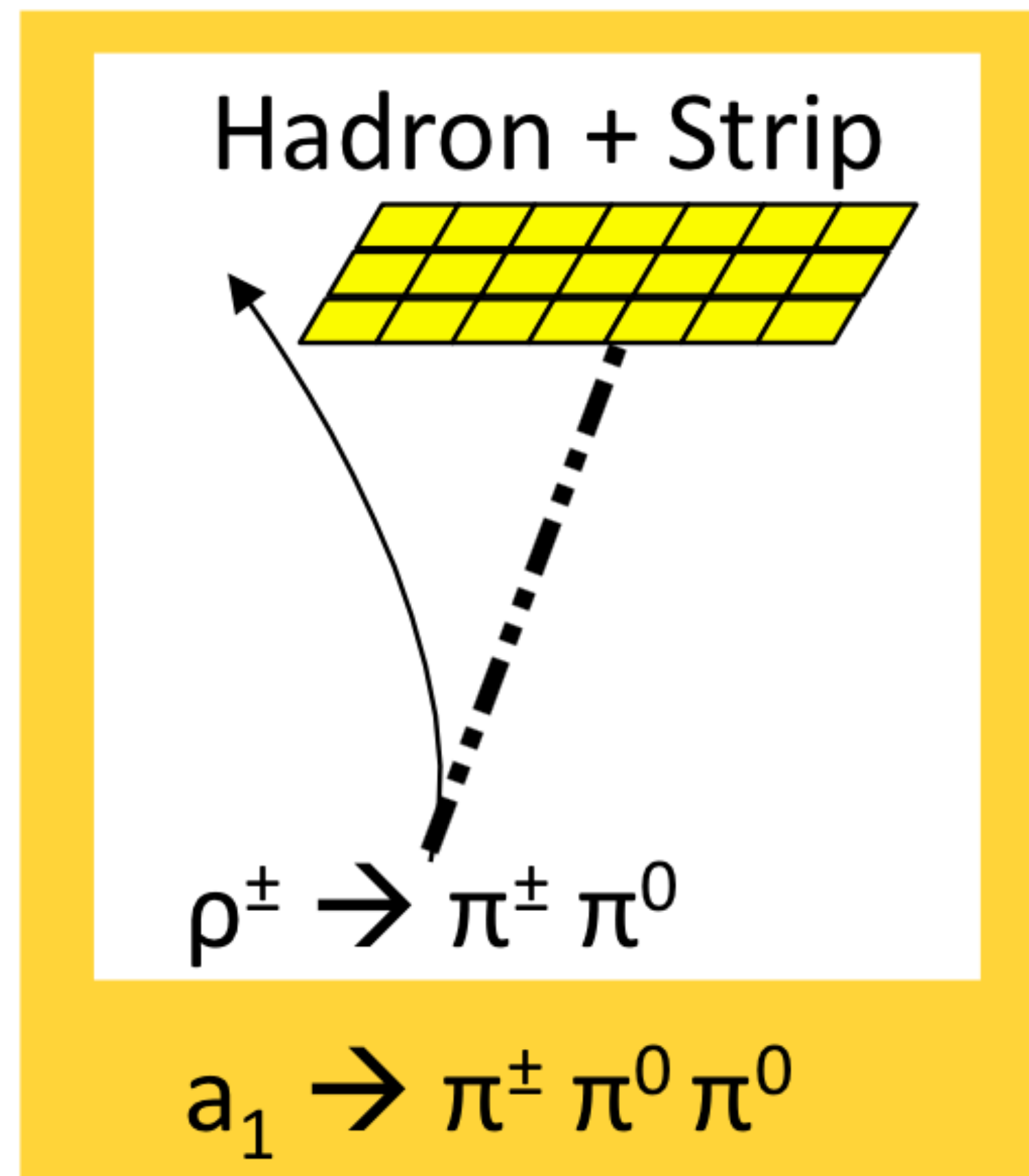
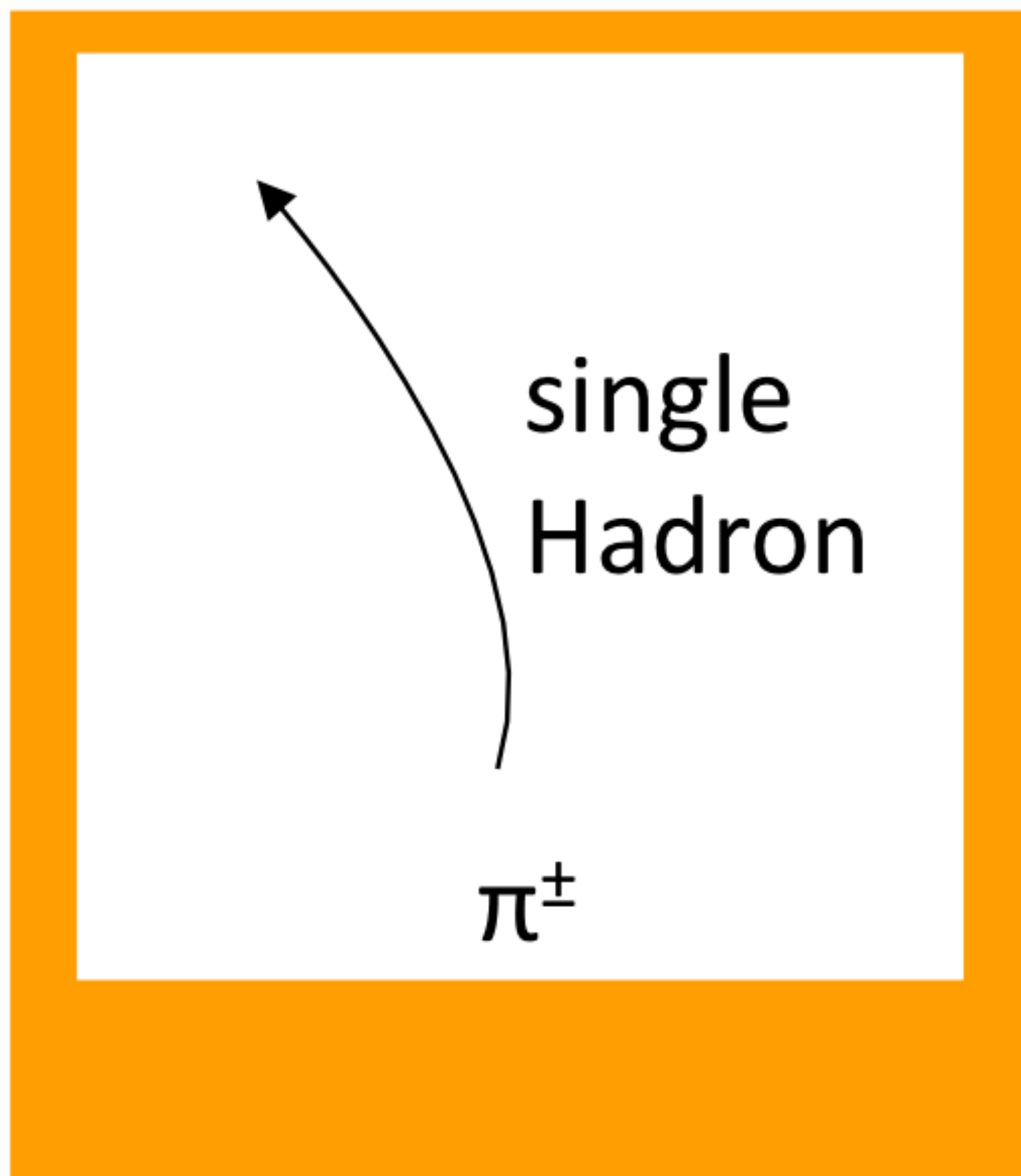
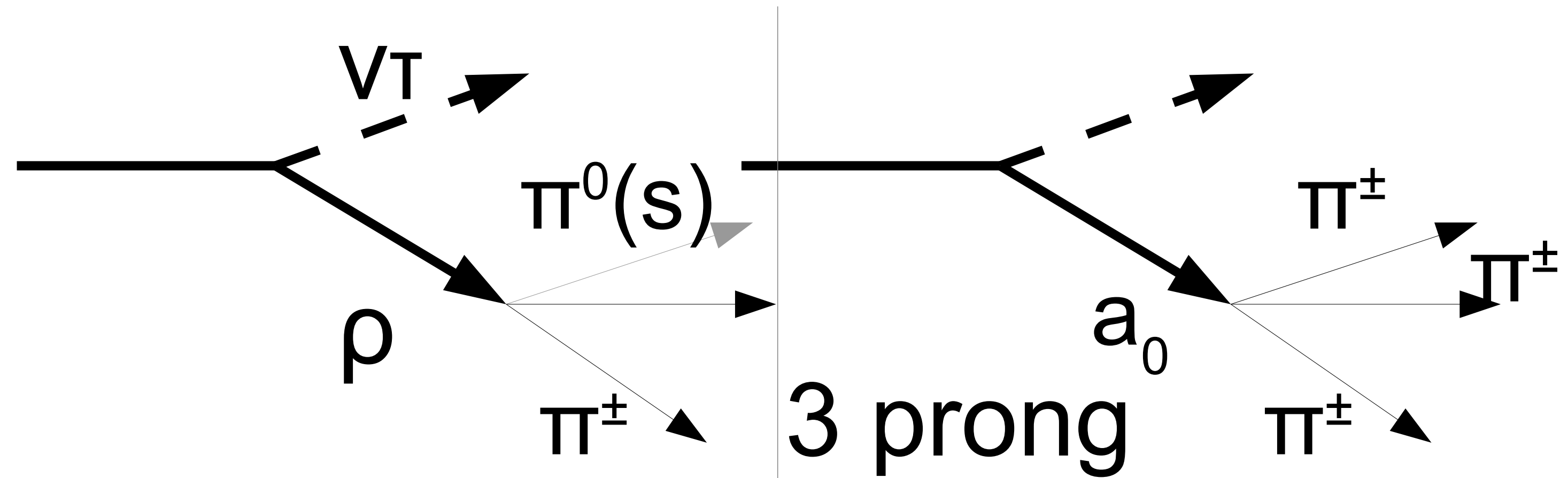
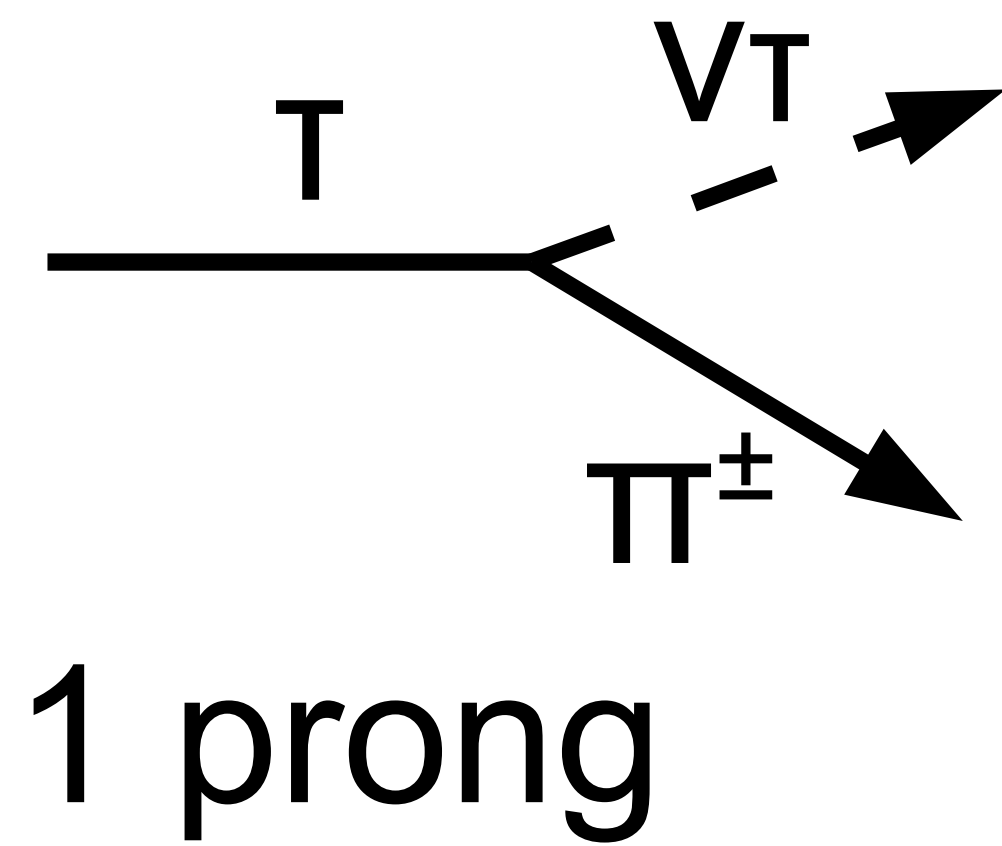
Massive and relatively long lived

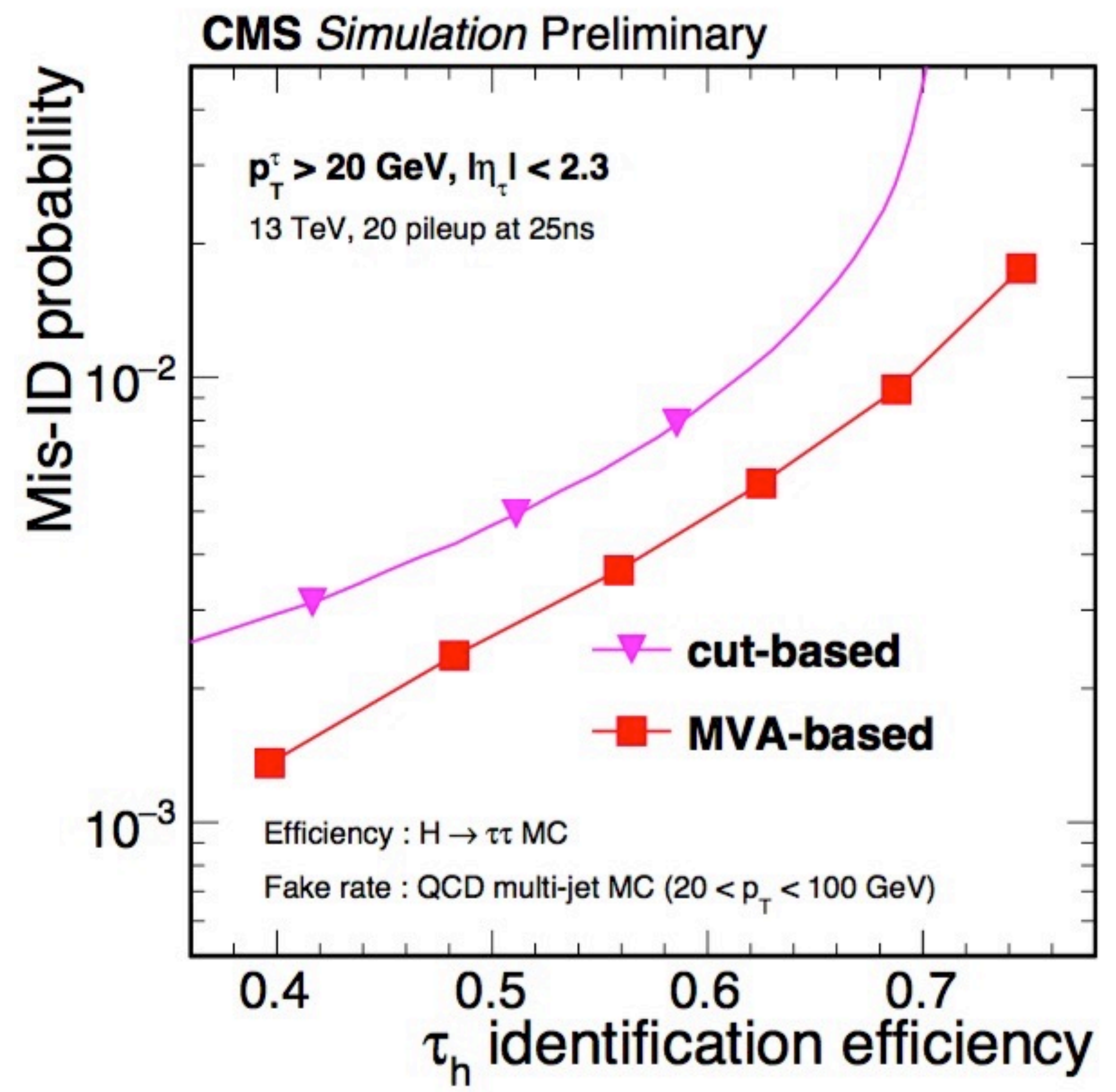
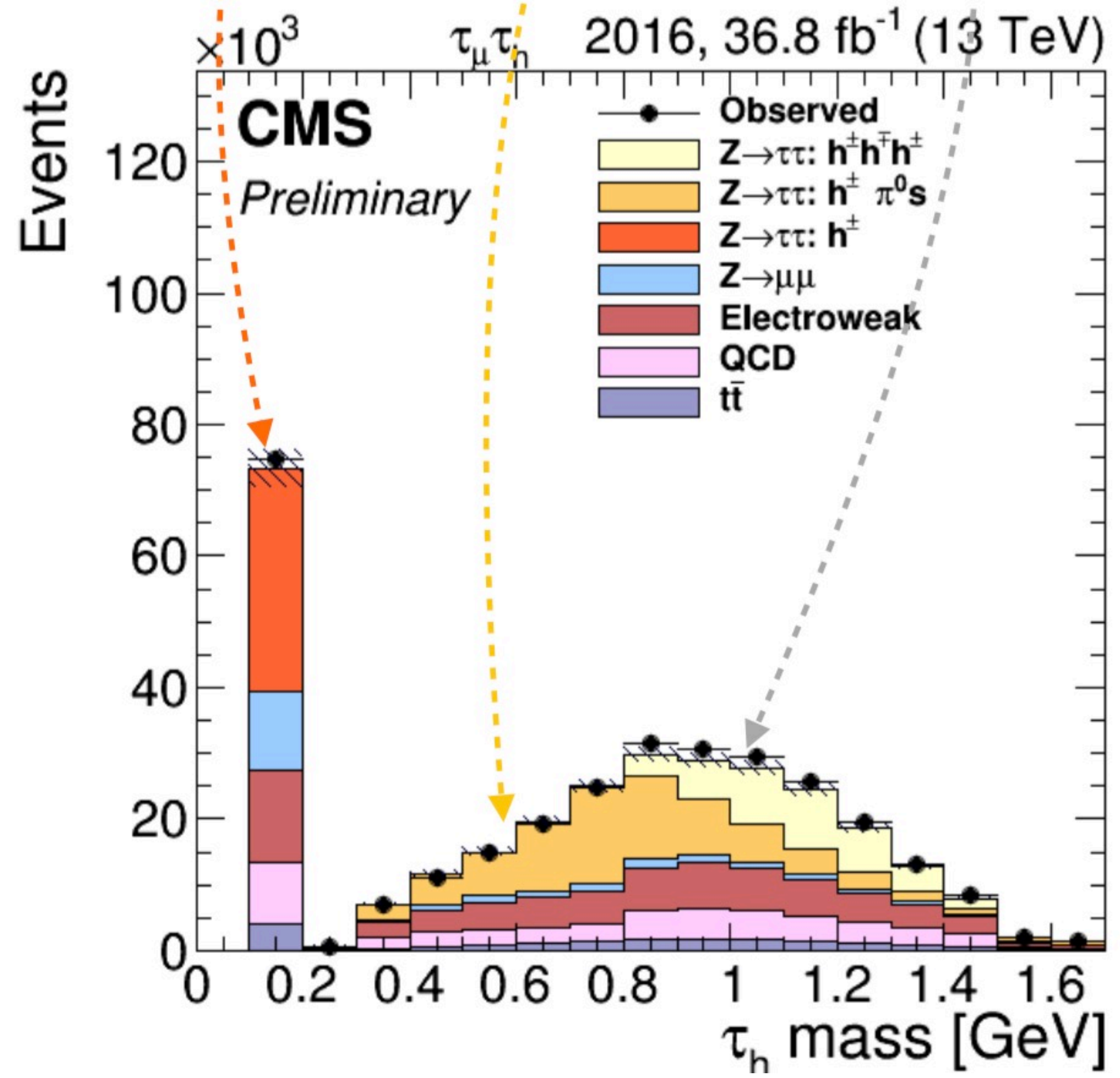
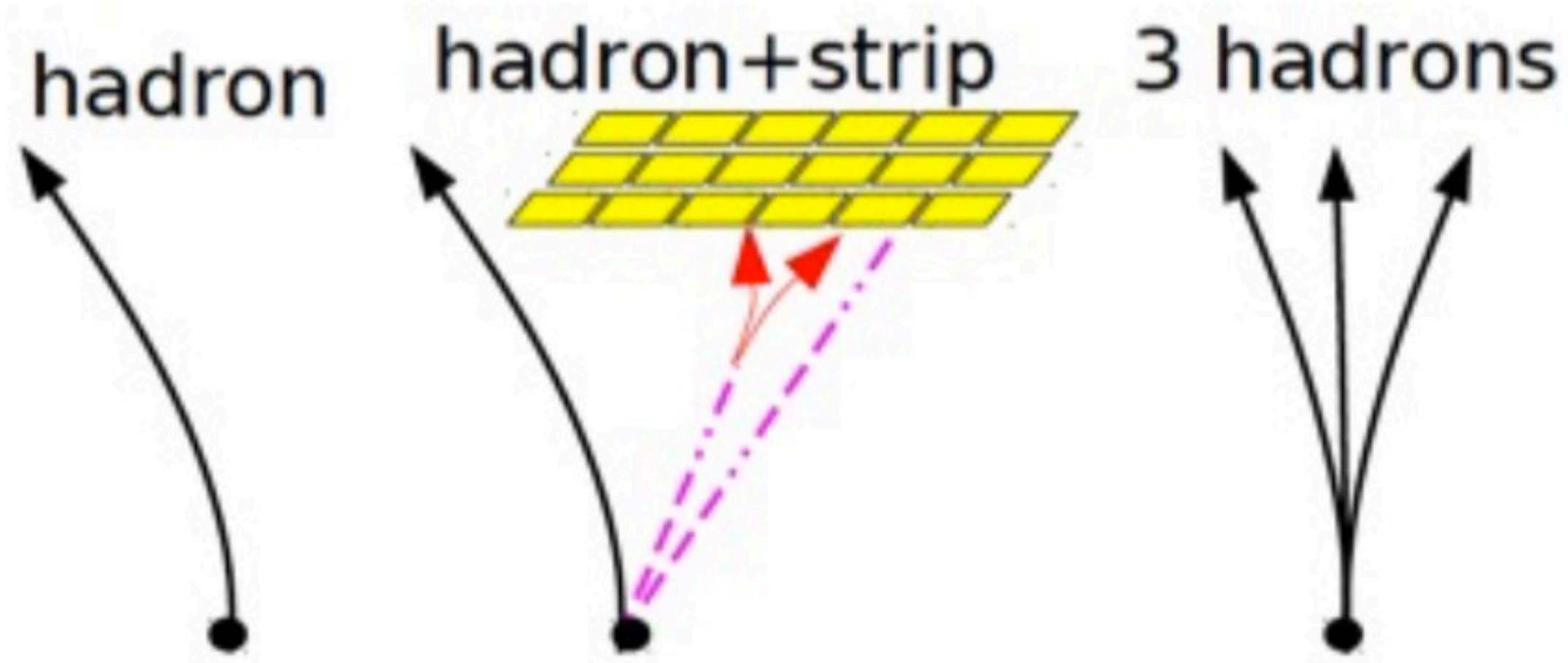
$$m(\tau) = 1.7 \text{ GeV}$$

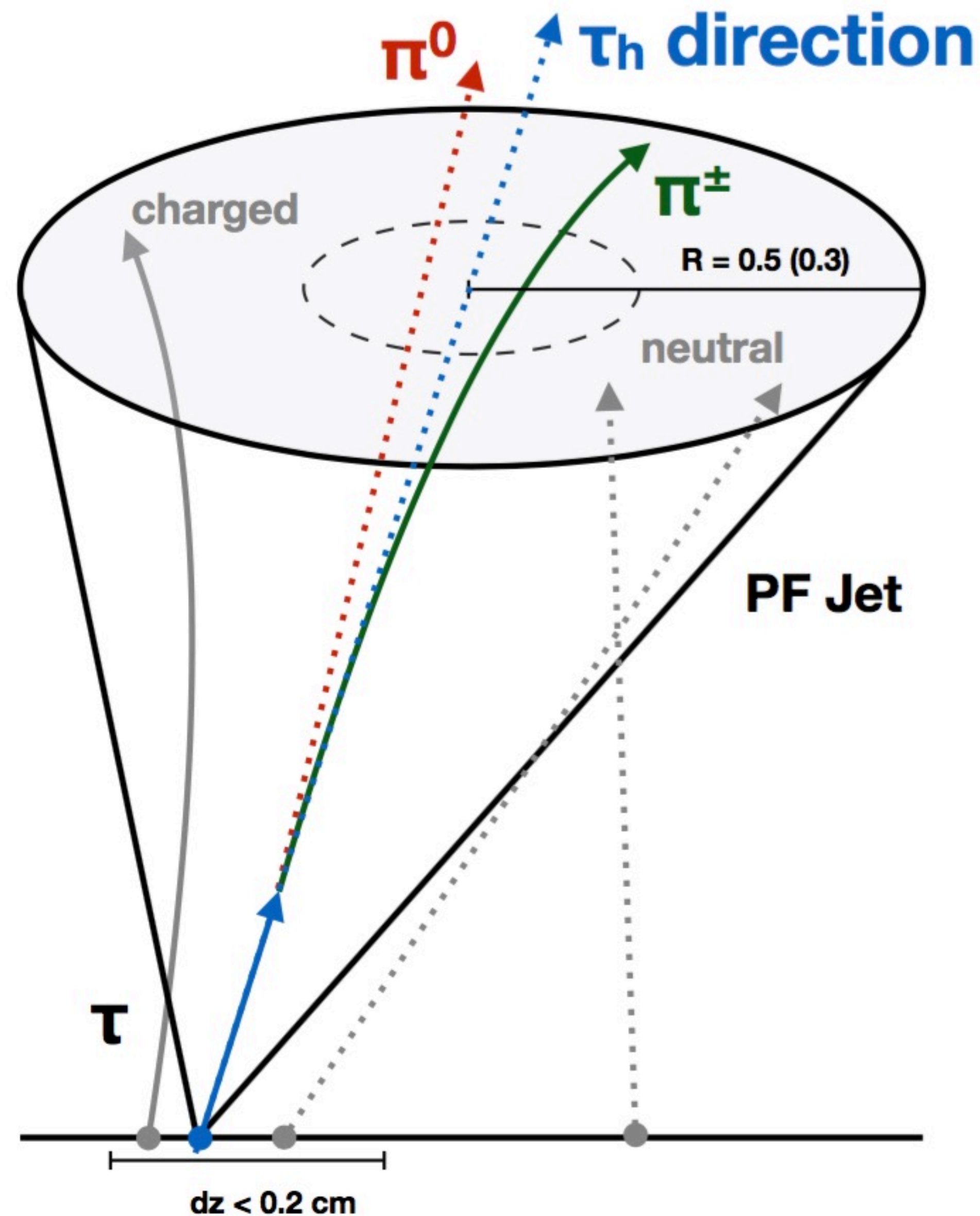
$$c\tau = 87 \mu\text{m}$$



Leptonic tau reconstruction relies on missing energy from the neutrinos







So far isolation has been mentioned in many contexts

Isolation very important to identify prompt muon, electron, photon, tau signals

For example:

Prompt:

Hadronic Tau vs. jet

Photon vs. jet

Muon vs. b jet

Isolation: the extra amount of energy around the object of interest

Often relative isolation is the quantity of interest

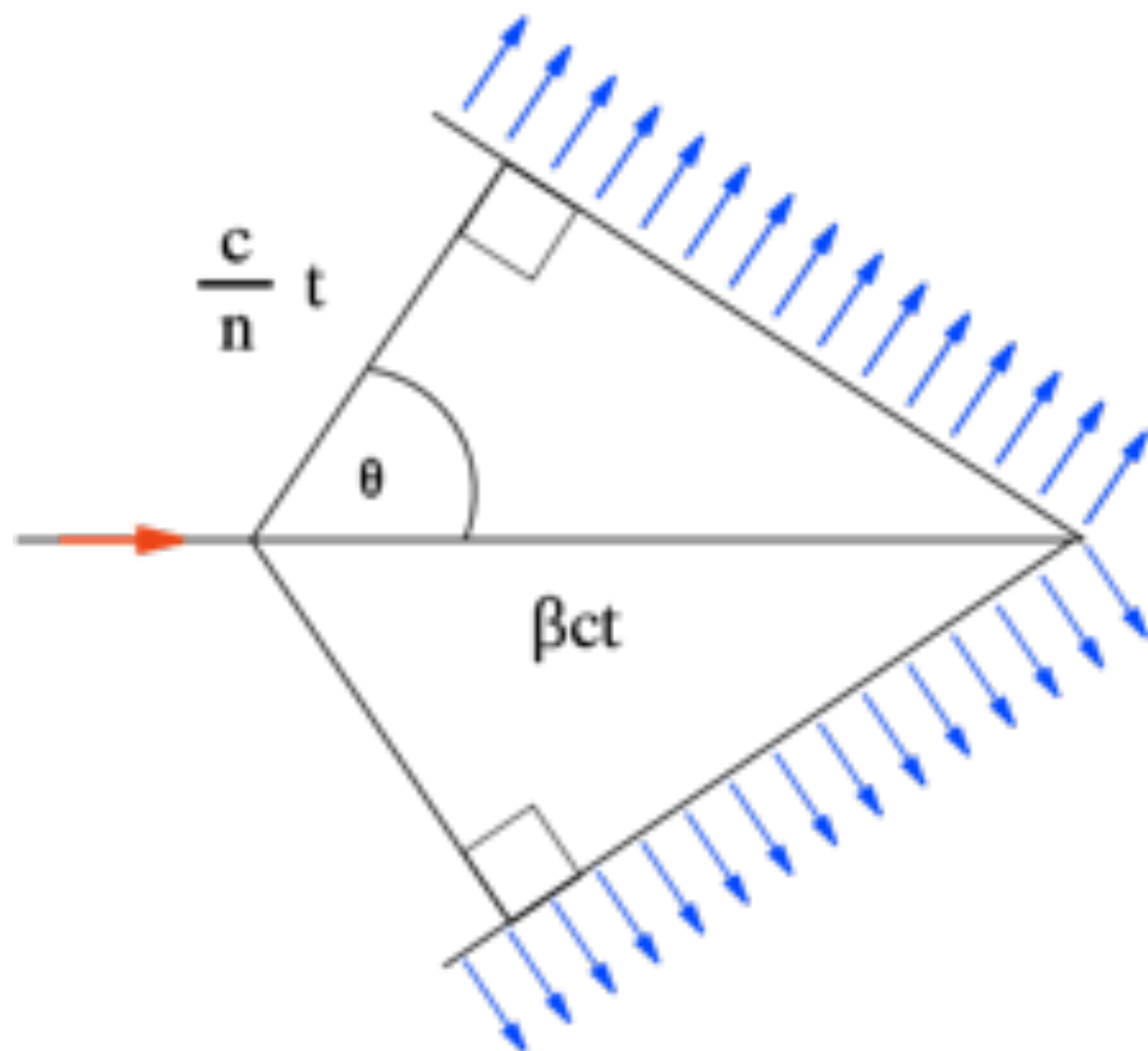
Will come back to this later with pileup discussion

Hadron ID is very important, particular in b physics

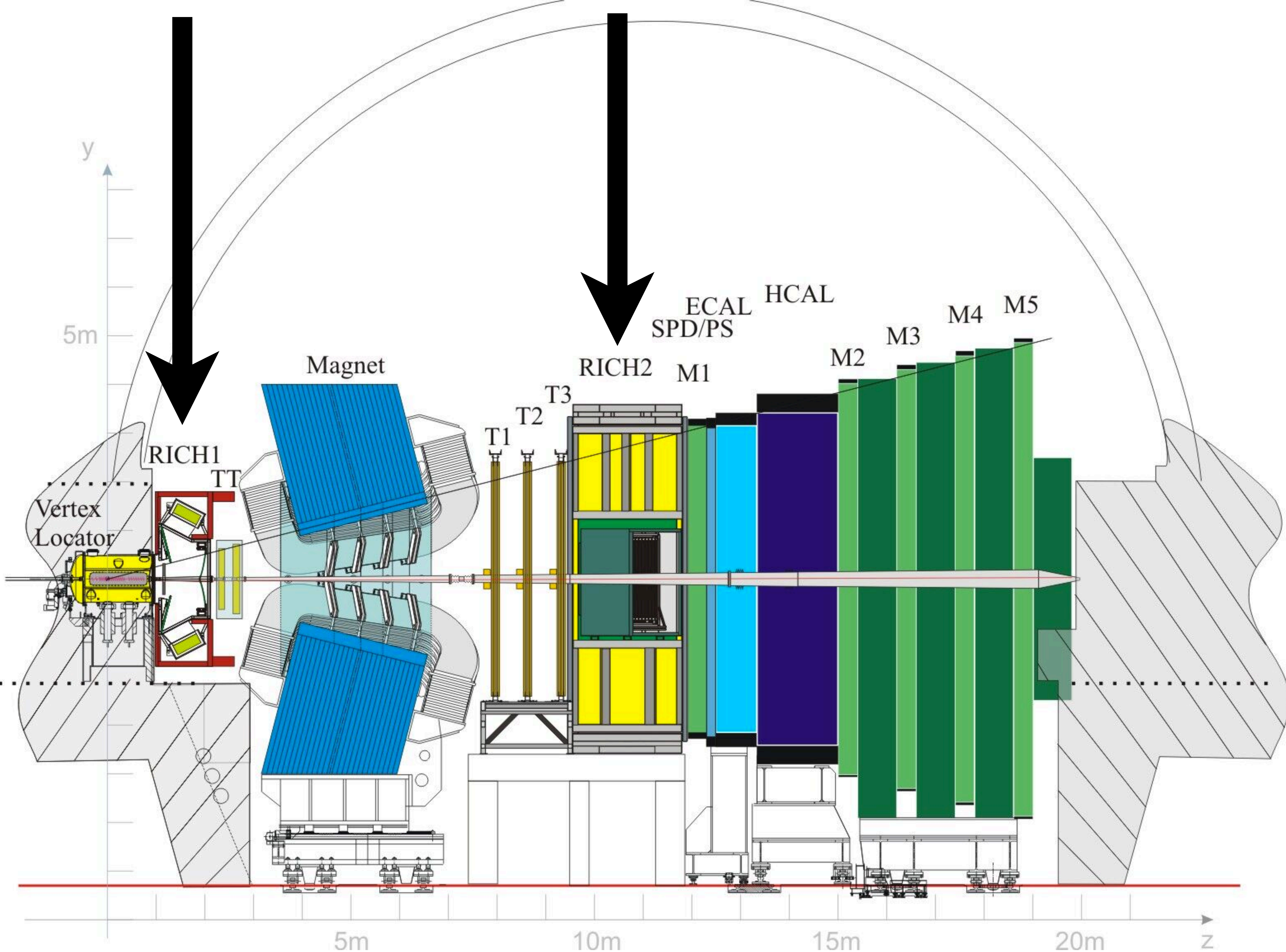
LHCb has a dedicated detector, RICH, for particle ID

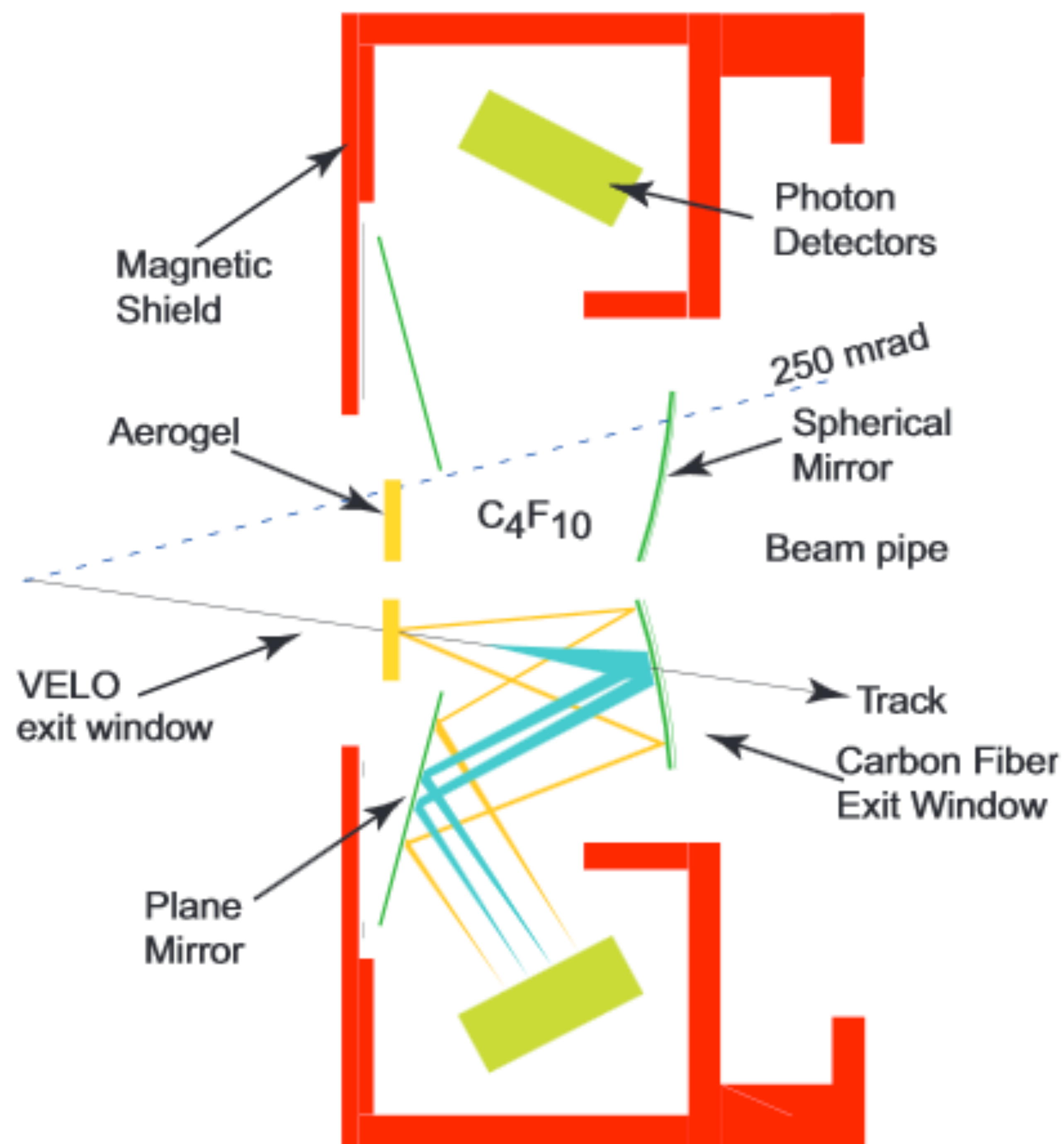
RICH: Ring Imaging Cherenkov Detector

Cherenkov radiation: Particles moving in material with index of refraction greater than 1 travel faster than the speed of light and emit radiation at an angle θ_c

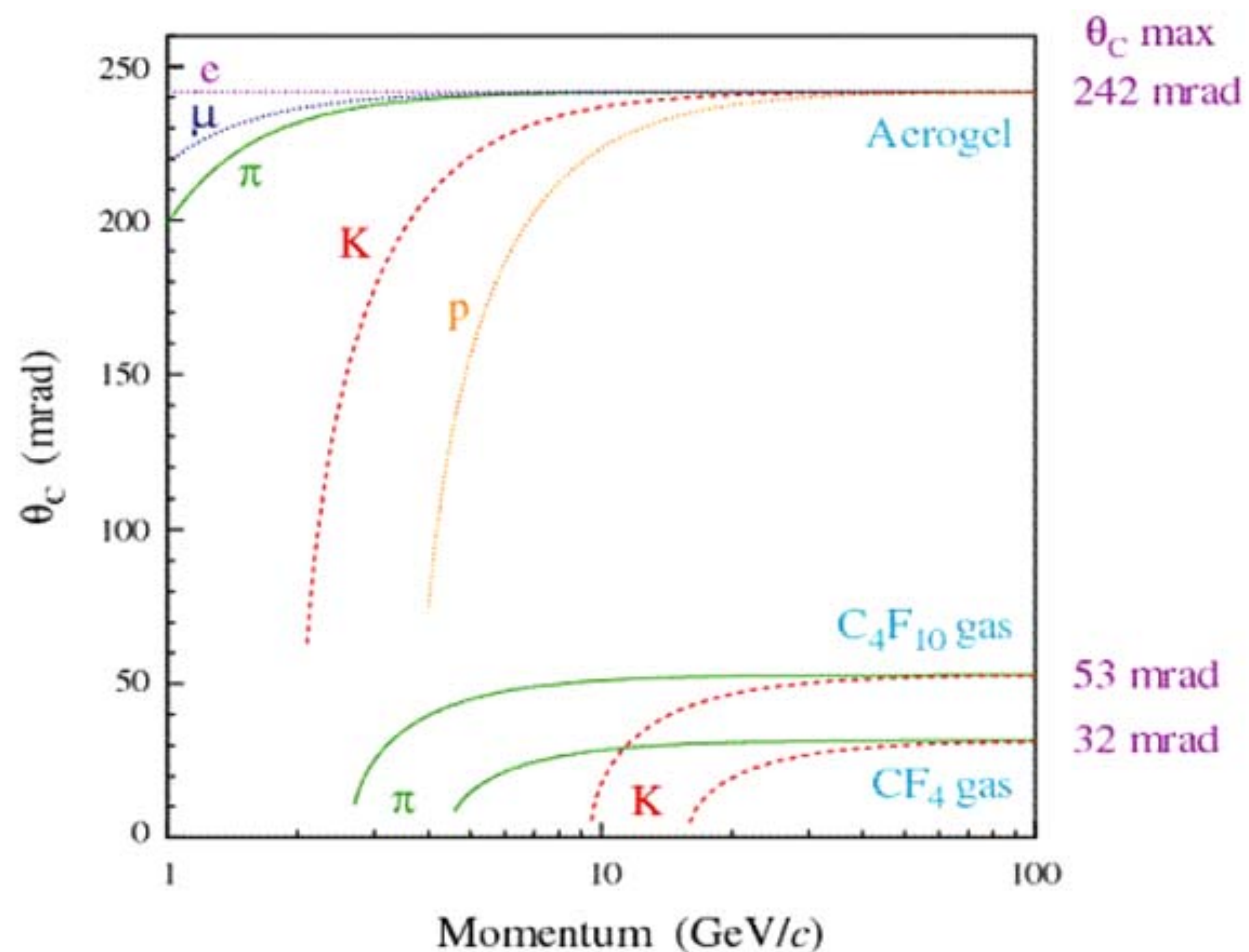


$$\cos \theta_c = \frac{1}{\beta n}$$

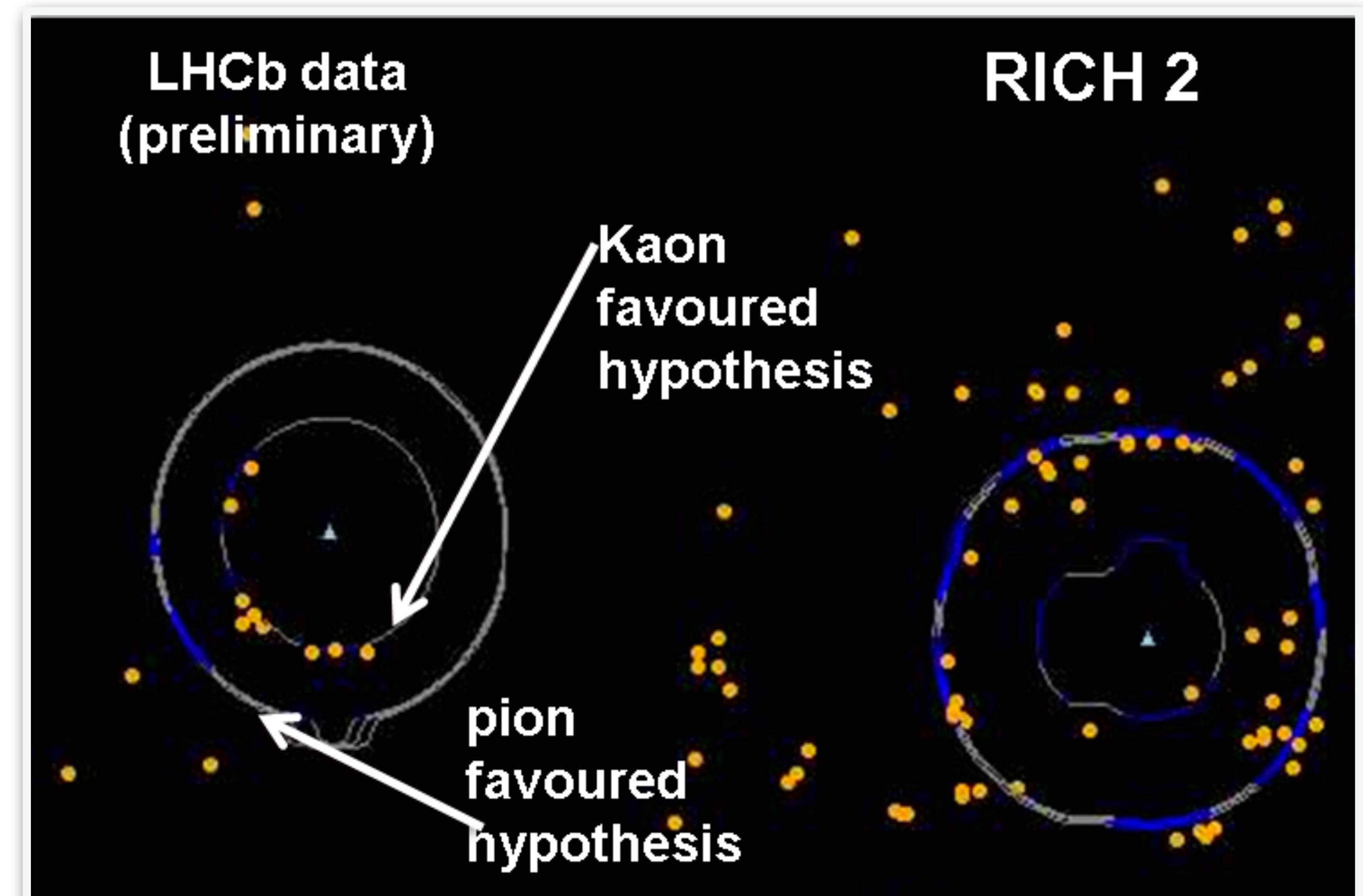
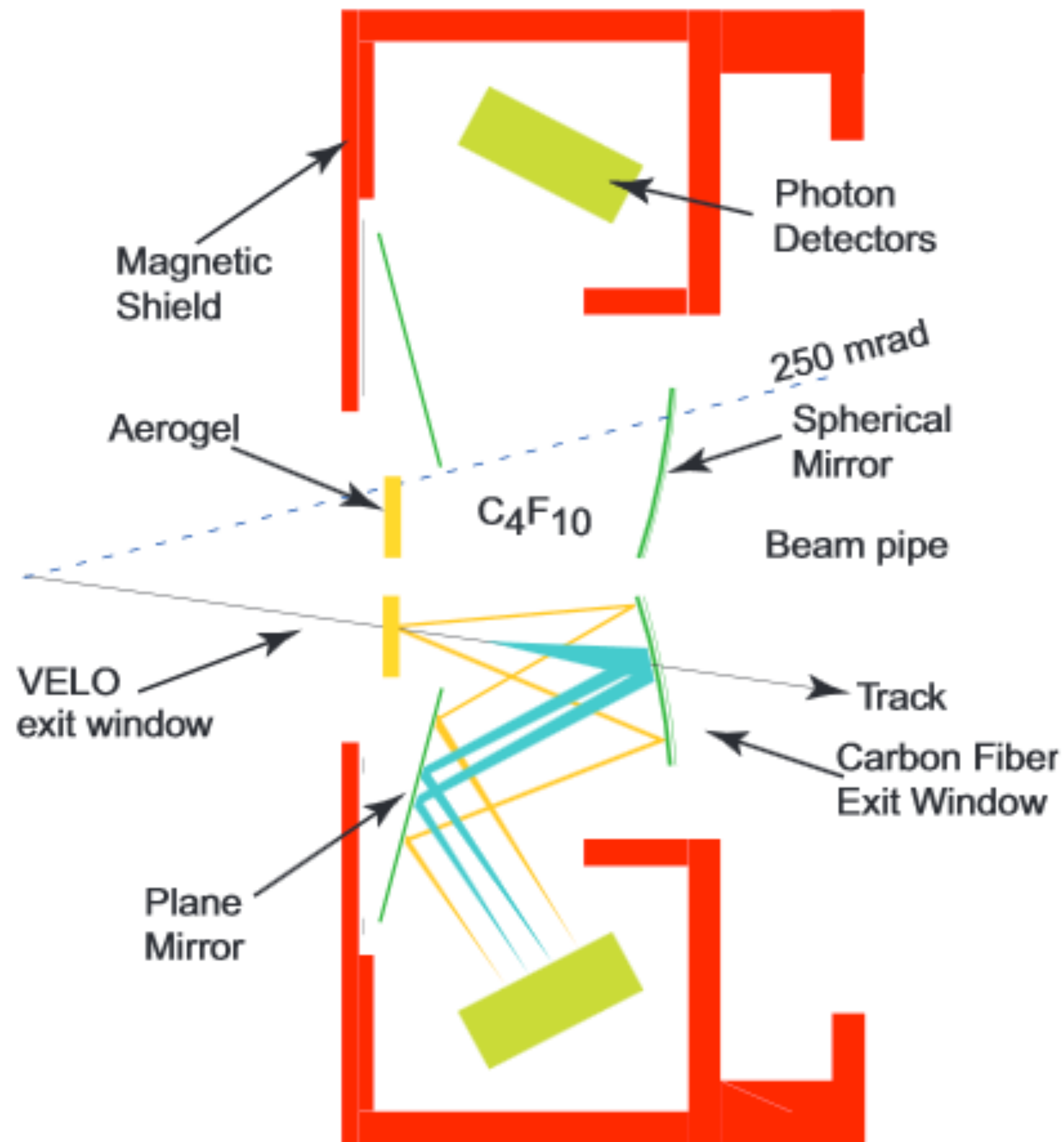


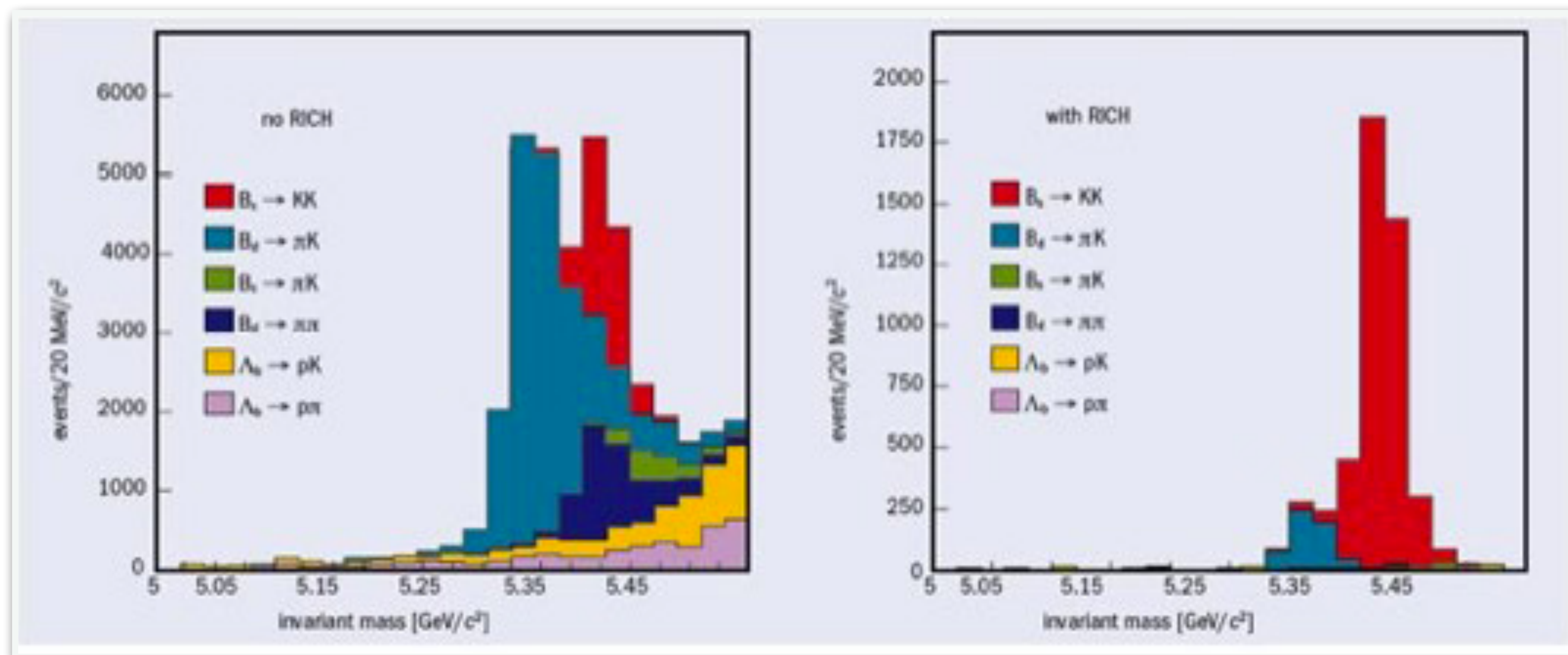
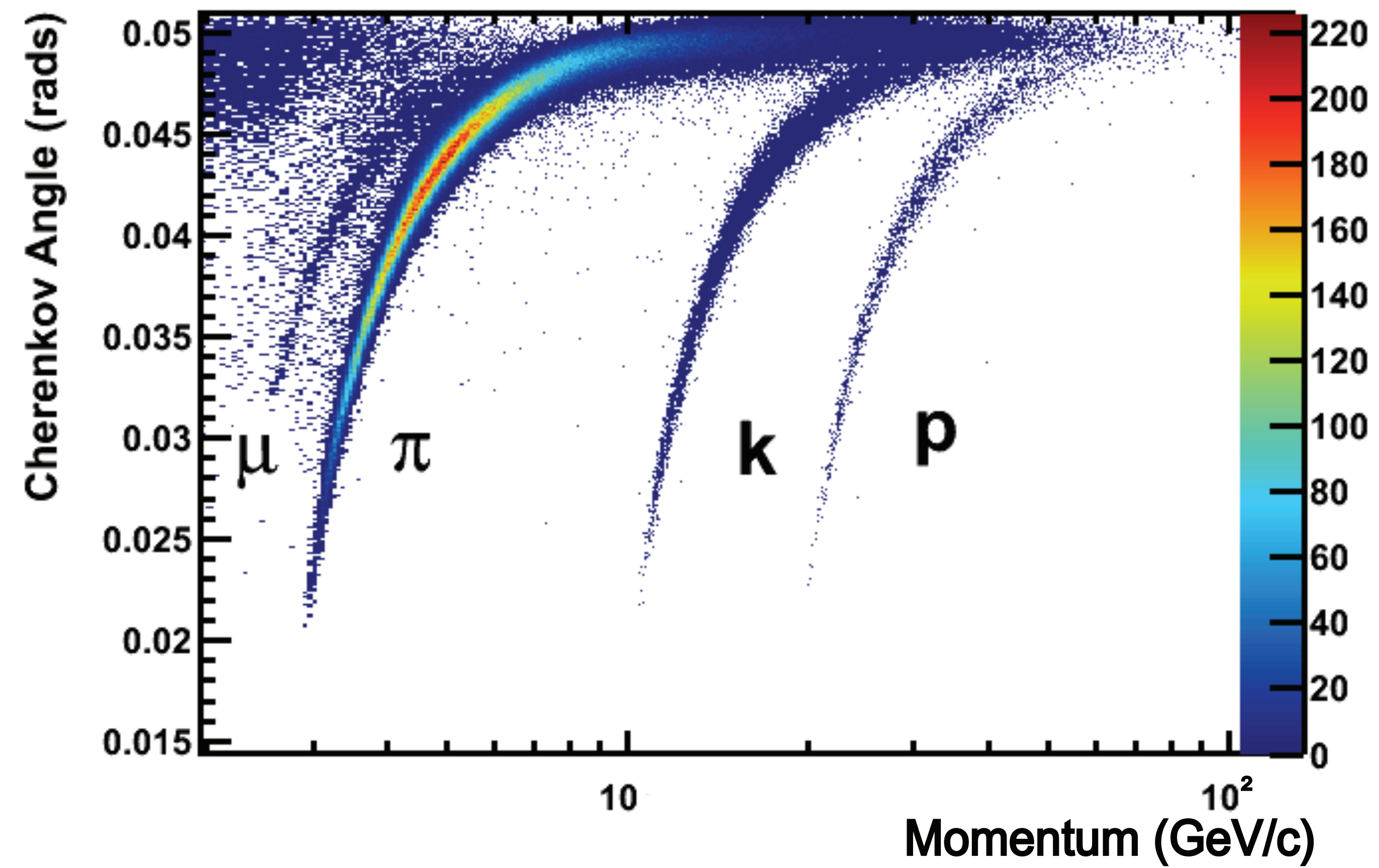


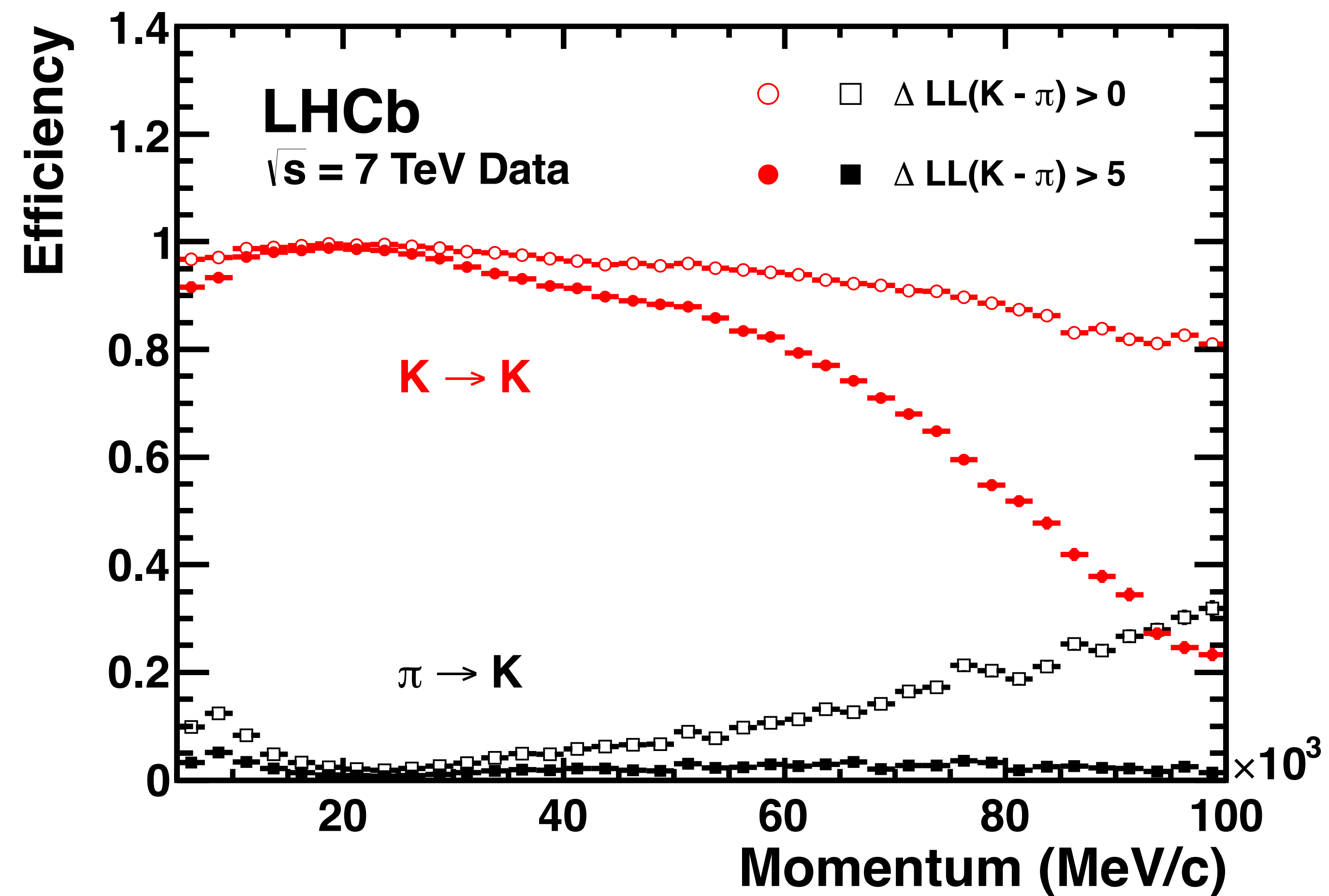
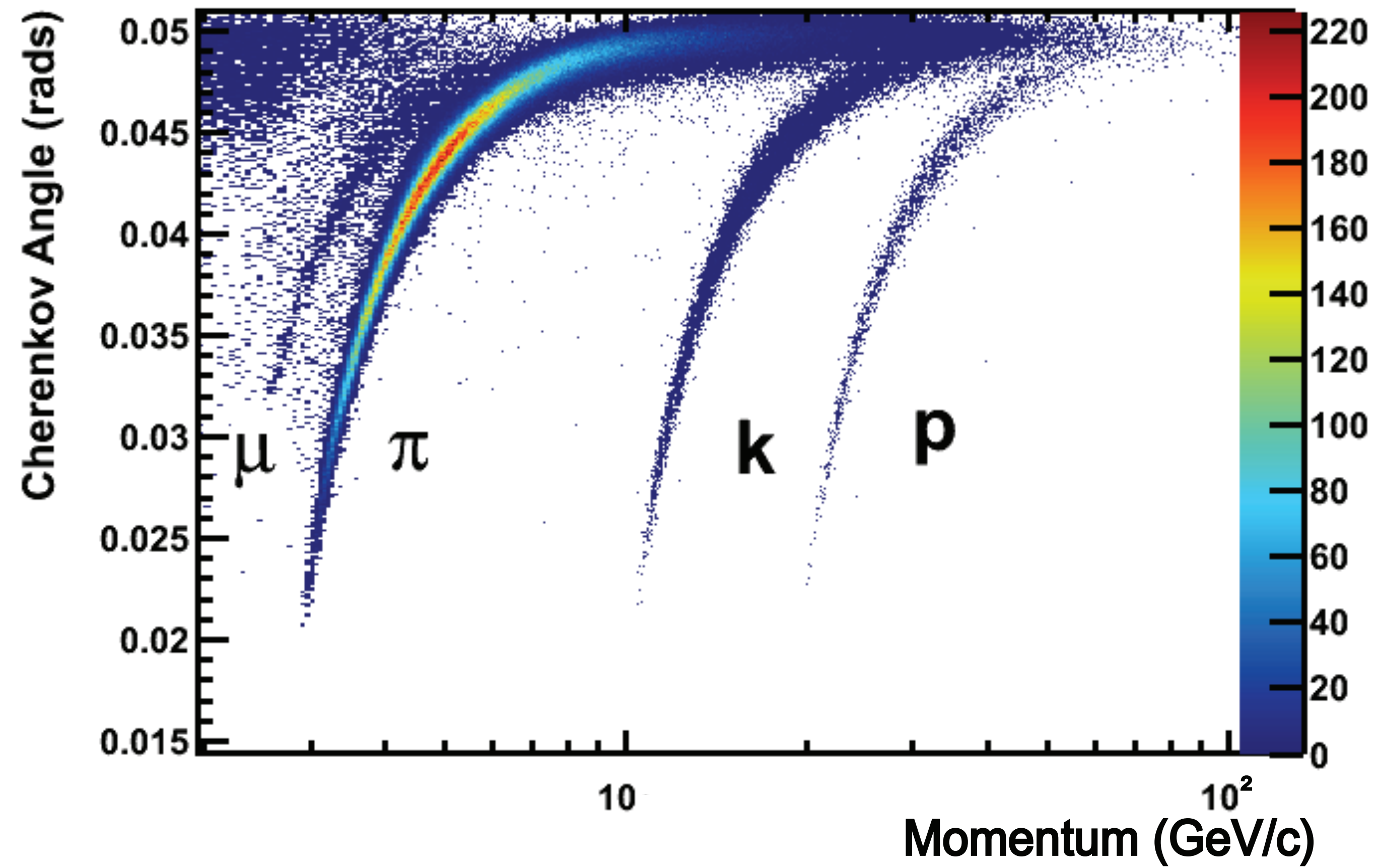
By measuring the track momentum and θ_c , one can identify the particle type



RICH1 = aerogel and C_4F_{10} gas
 RICH2 = CF_4 gas







End of lecture 1

Tomorrow: Let's get ADVANCED

Particle flow

Jets and MET

Jet substructure

Pileup and underlying event in HI

Exotic and beyond