# Excited state analysis in the quasi-PDF matrix elements

#### Yi-Bo Yang Michiaan State University



**L**attice **P**arton **P**hysics **P**roject

yangyibo@pa.msu.edu

# Large momentum effective theory

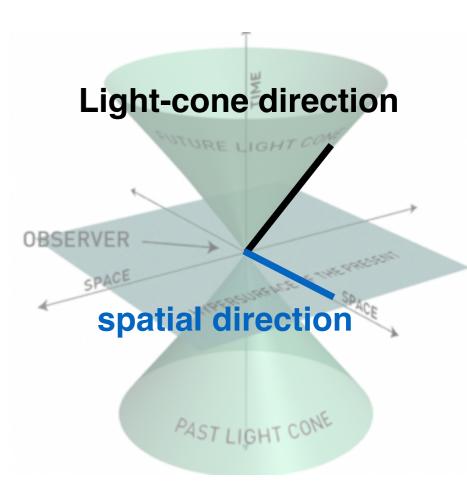
The light-cone PDF is defined by

$$\begin{split} q(x,\mu^2) \; &= \; \int \frac{d\xi^-}{4\pi} e^{-ix\xi^- P^+} \langle P | \overline{\psi}(\xi^-) \gamma^+ \\ & \times \exp\left(-ig \int_0^{\xi^-} d\eta^- A^+(\eta^-)\right) \psi(0) | P \rangle \end{split}$$

#### When nucleon is boosted:

 The axial gauge conditions become the light-cone one,

## quasi-PDF



and the FT of the RI/MOM renormalized quasi-PDF,

$$\langle P|\bar{\psi}(z)\gamma_t U_z(z,0)\psi(0)|P\rangle_{p^2=\mu_R^2,p_z=p_z^R}^R$$

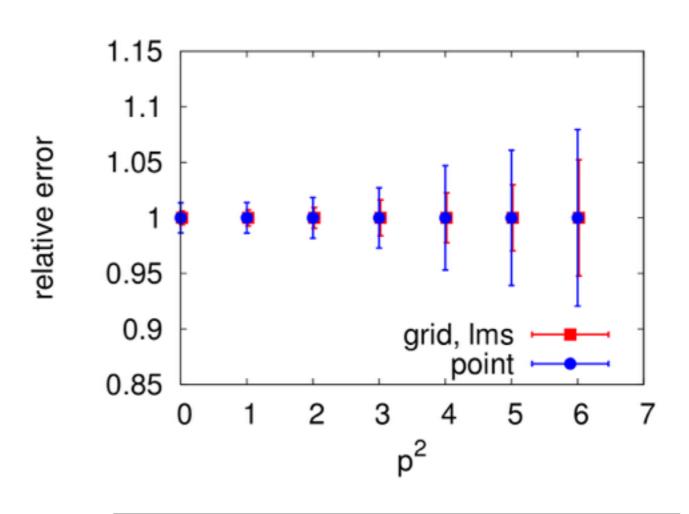
X. Ji, PRL 110 (2013) 262002, 1305.1539 X. Ji. SCPMA 57 (2014) 1407, 1404.6680

becomes the light-cone PDF up to perturbative matching.

# Large momentum...

## Really possible?

- The expectation value of a moving hadron decays as  $\sim e^{-Et}$ , where  $E=\sqrt{(m^2+p^2)}$ .
- Its statistical uncertainty decays as  $\sim e^{-m_0 t}$ , where  $m_0 \propto m_\pi$  and  $m_0 \leq m$ .
- Seems to be hopeless to reach large momentum as required by LaMET.

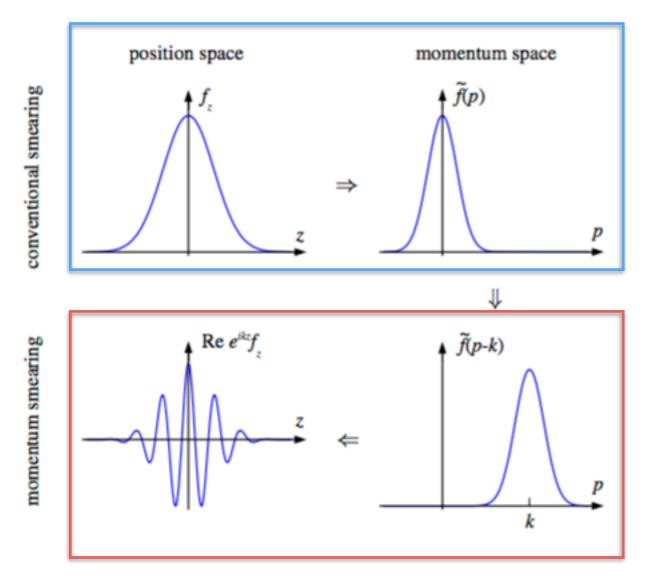


An example from: YBY, et al., χQCD collaboration, PRD93 (2016) 034503, 1509.04616

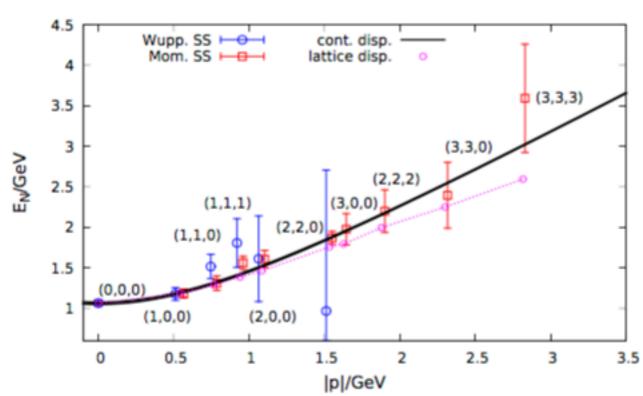
# Momentum smearing

### in 2pt

#### Wuppertal (Gaussian) Smearing



Momentum Smearing

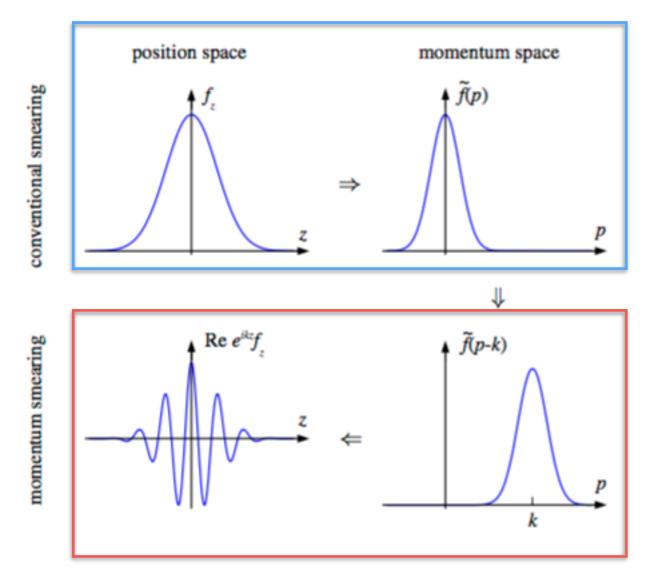


Momentum smearing can provide much better signal in 2pt!

# Momentum smearing

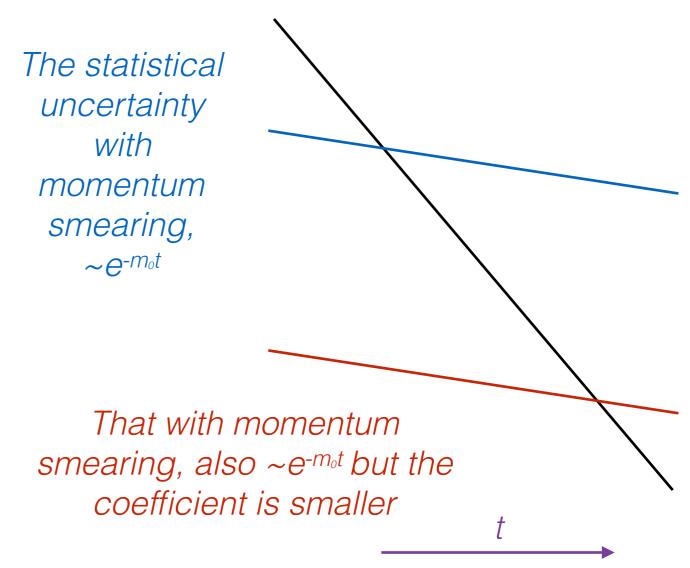






Momentum Smearing

The expectation value of the 2pt with momentum, ~e<sup>-Et</sup>

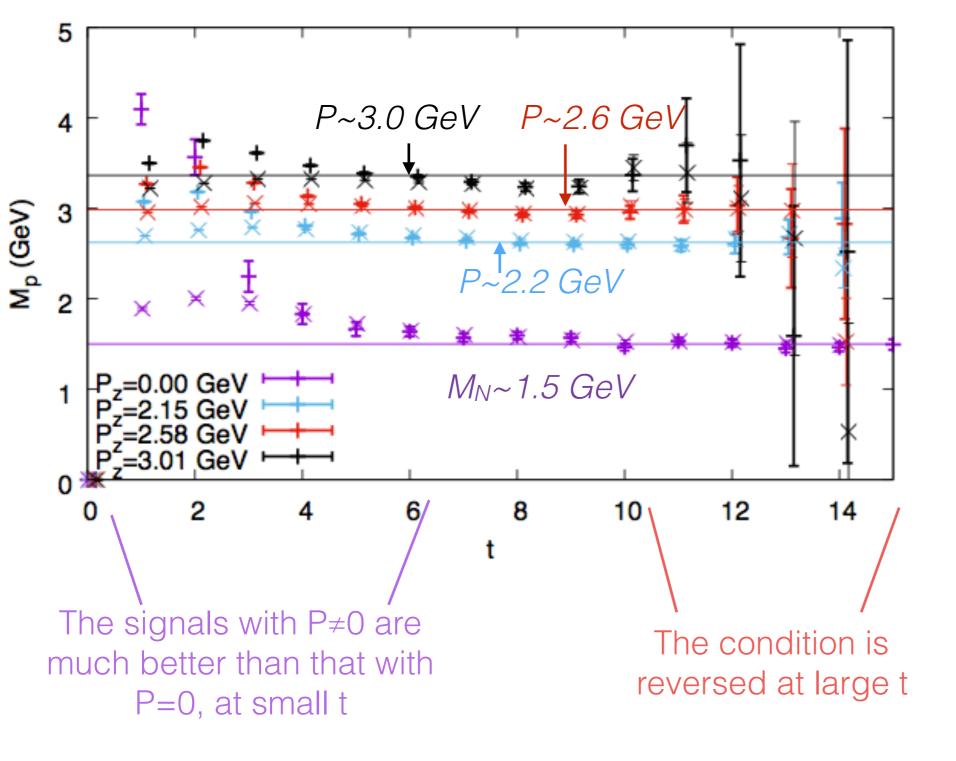


The momentum smearing is not a magic. Just effectively gain some statistics.

# 2pt

#### a09m310

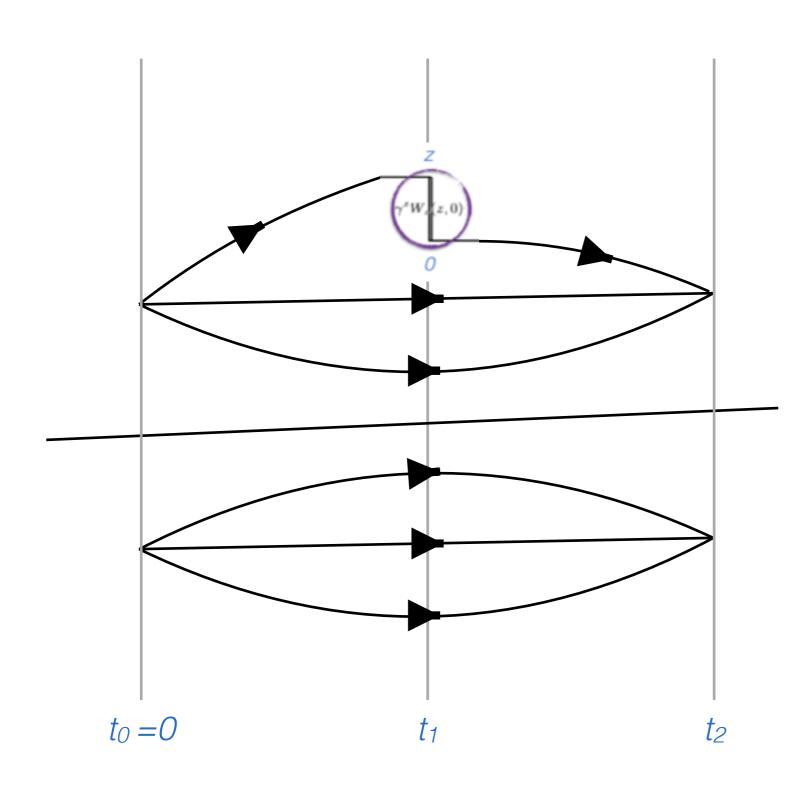
- $m_{\pi}$ ~670 MeV,  $m_{\pi}$ sea~310 MeV, a=0.09 fm;
- $L^3xT = 32^3x96$ ;
- 1 step HYP smearing on everything.
- 1152 measurements=(288 configurations) x (4 sources/configuration);



+: smeared src. to smeared sink

**x**: smeared src. to point sink

## Quasi-PDF matrix element

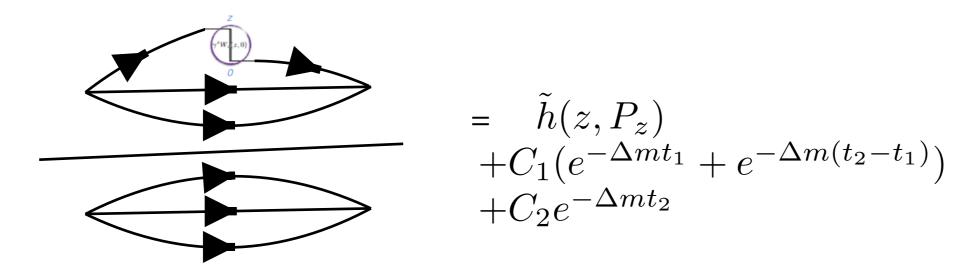


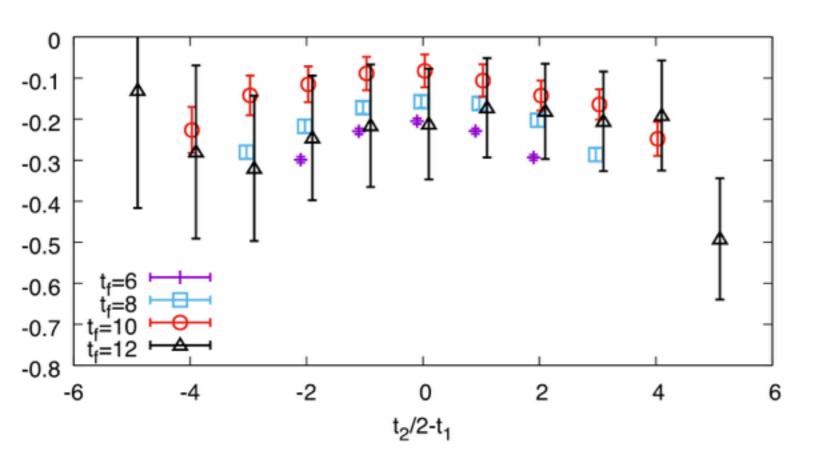
In practical, the following ratio is calculated on the lattice:

$$= \tilde{h}(z, P_z) + C_1(e^{-\Delta m t_1} + e^{-\Delta m (t_2 - t_1)}) + C_2 e^{-\Delta m t_2}$$

where the  $C_{1,2}$  terms correspond to the excited state contaminations vanishing in the  $t_2 \gg t_1 \gg 0$  limit.

## Quasi-PDF matrix element



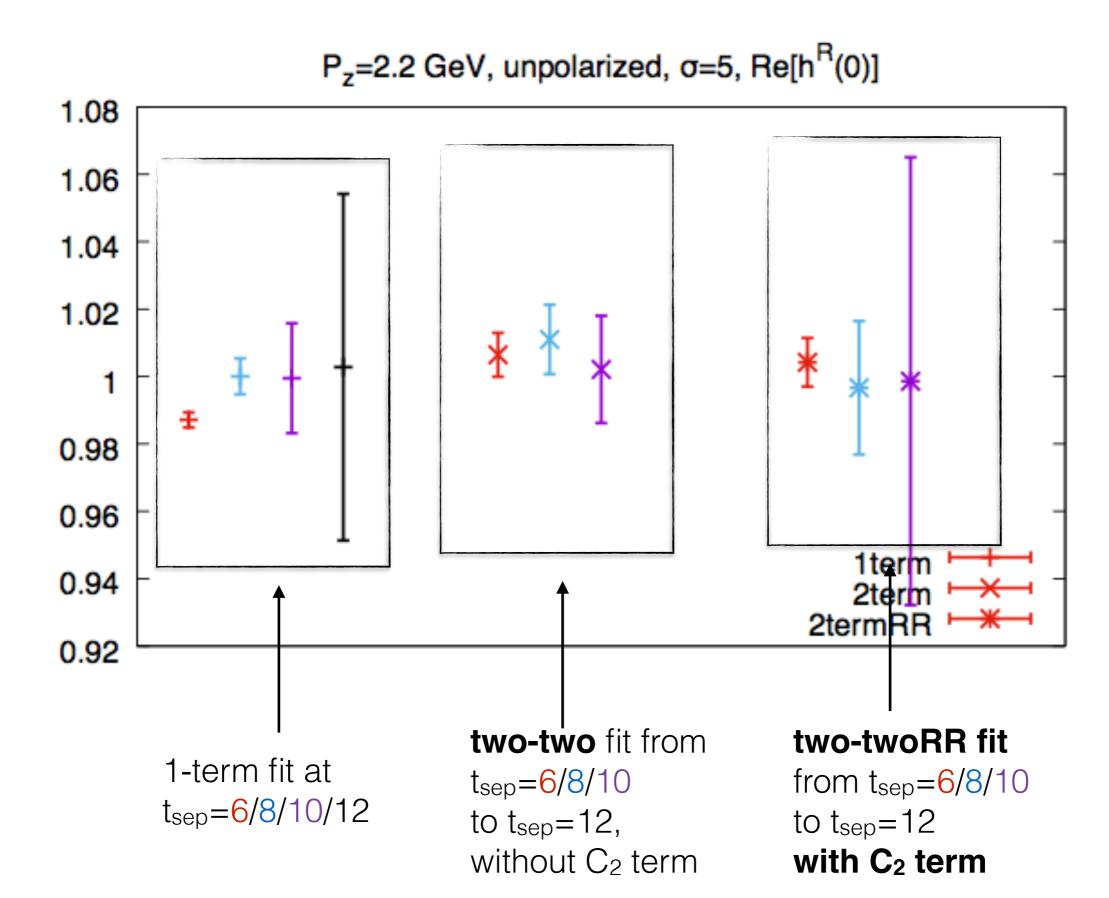


1-term fit: A constant fit around t<sub>1</sub>~t<sub>2</sub>/2

2-term fit: drop the contribution from the C<sub>2</sub> term

2-termRR fit: including all the terms

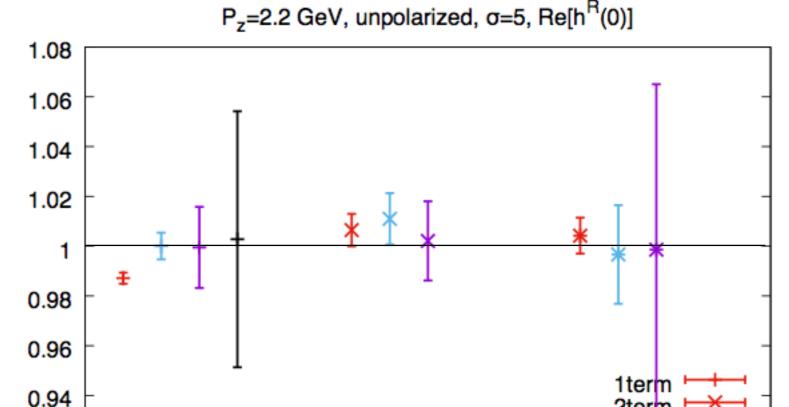
#### The fit results of the ME



# z=0, unpolarized

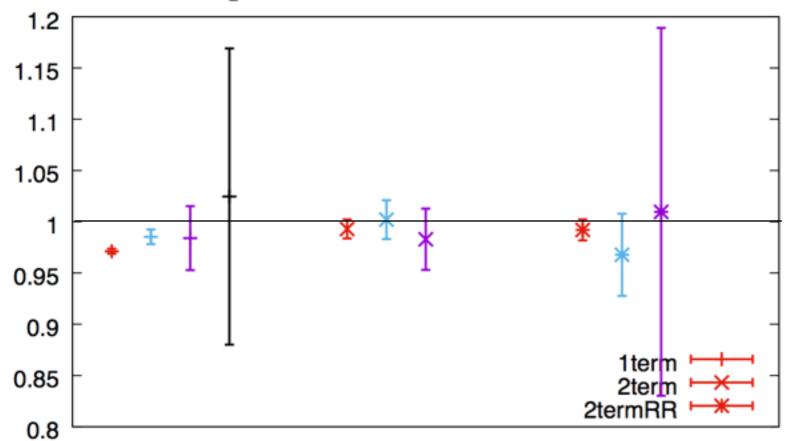
Three groups for three types of the fits

Different colors for different minimum separations.



 $P_z$ =2.6 GeV, unpolarized,  $\sigma$ =5, Re[h<sup>R</sup>(0)]

0.92



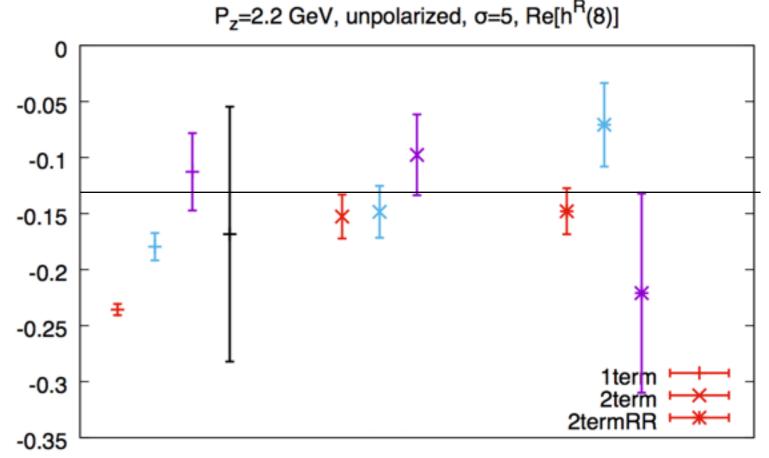
The results with different P\_z are consistent with the same normalization at z=0.

The horizontal line with the same value are placed to guide your eyes.

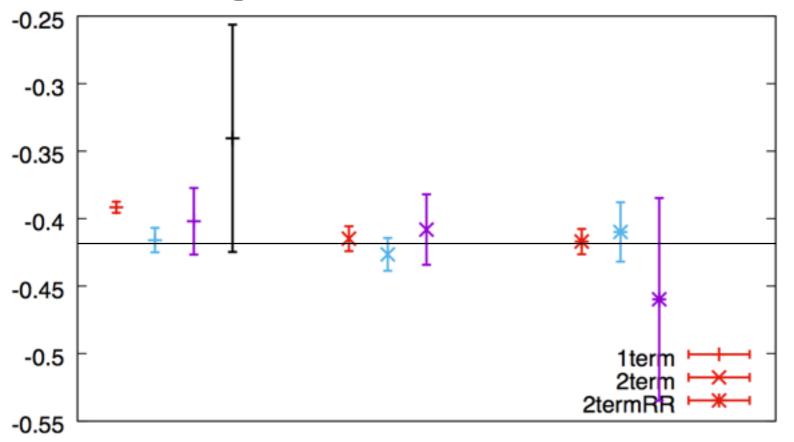
# Pz=2.2 GeV, z=8, unpolarized

Three groups for three types of the fits

Different colors for different minimum separations.

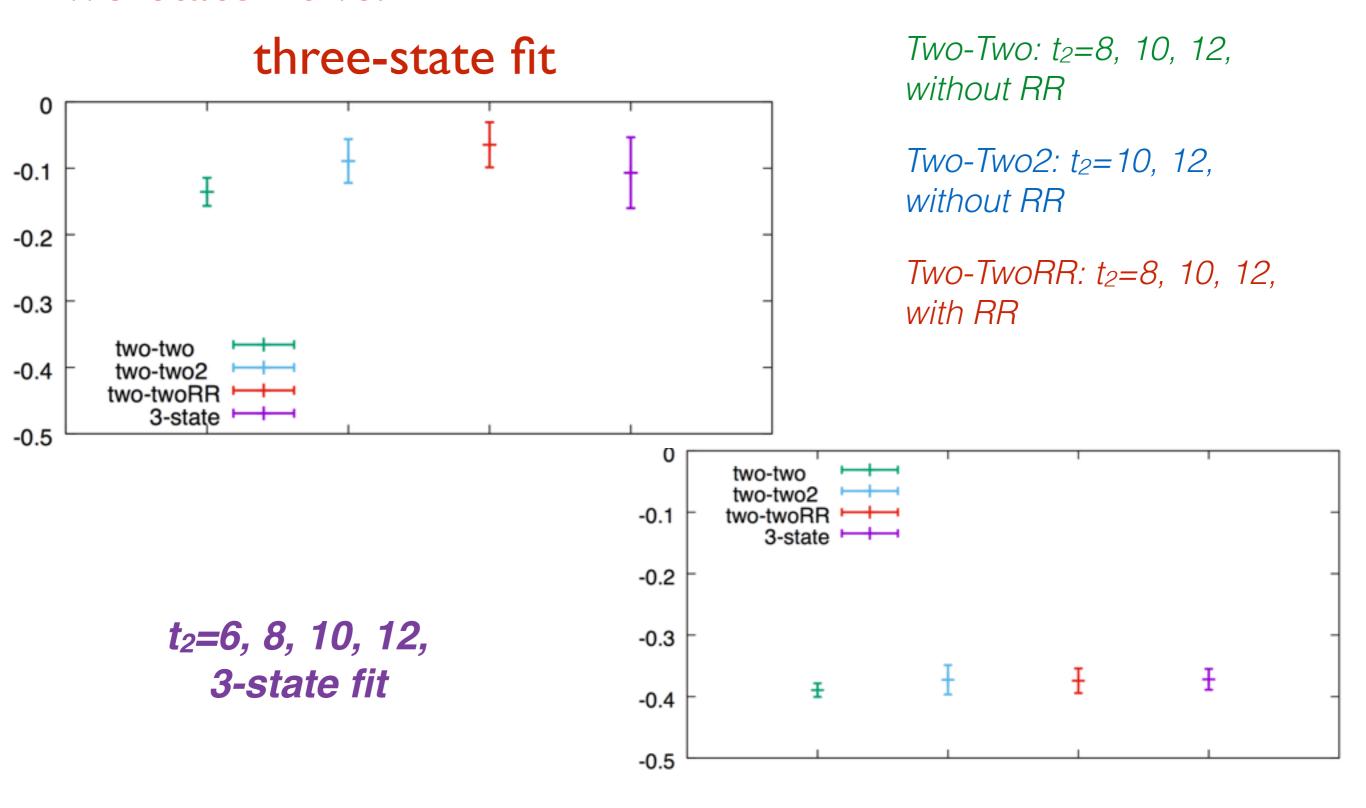


 $P_z$ =2.2 GeV, unpolarized,  $\sigma$ =5, Im[ $h^R$ (8)]

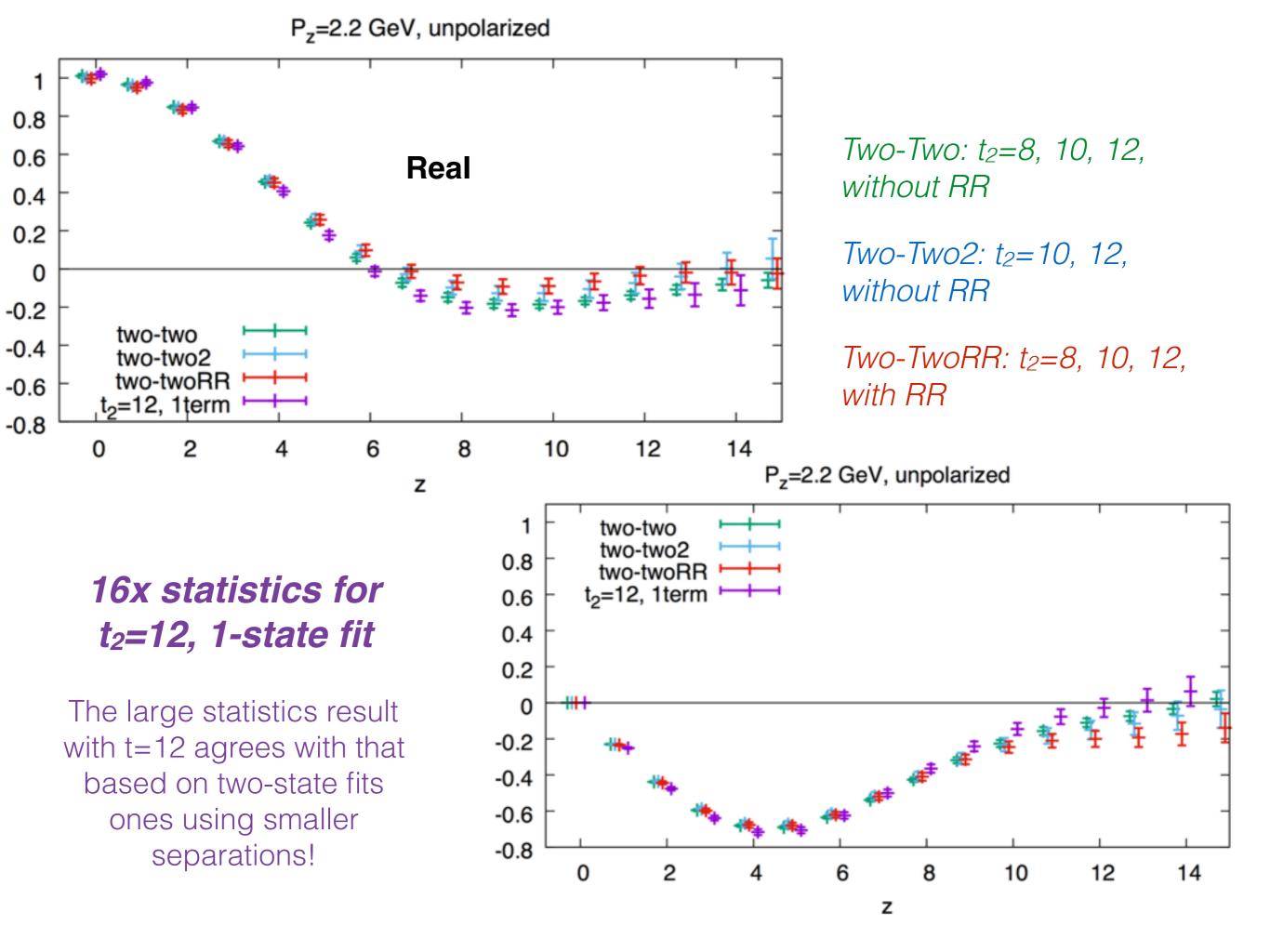


At z=8, the fits for the imaginal part have good agreement but the real part require more statistics at large t<sub>sep</sub> to get a solid conclusion.

#### Two-state fit vs.



Ruizi. Li, et al., LP<sup>3</sup> collaboration, in preparation



# Summary

- The momentum smearing allow us to achieve good signal for the matrix elements with large hadron momentum, at small source-sink separation.
- The multi-state fit can provide a good subtraction on the excited state contamination with smaller source-sink separations.
- The production with another smearing size are ongoing to confirm the multi-state fit results.