Nucleon EDMs on a Lattice at the Physical Point

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LATTICE 2018 East Lansing, MI, July 22-28, 2018

Outline

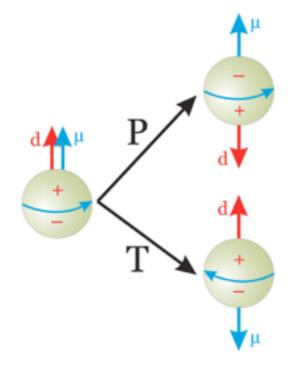
Nucleon Electric Dipole Moments: Introduction

- Motivation
- Experimental status & outlook
- Lattice methodology
- Physical point calculations with chiral quarks
 - Form Factors → [T.Izubuchi's talk, July 27 5:30pm @106 (Hadron Structure)]
 - Electric dipole moments induced by quark chromo-EDM

Studies of θ_{QCD} -induced nucleon EDM

- Noise reduction with subvolume top.charge sampling
- Results from mπ≥330 MeV lattices
- Outlook for physical point calculations

Nucleon Electric Dipole Moments



$$\vec{d}_N = d_N \frac{\vec{S}}{S} \qquad \mathcal{H} = -\vec{d}_N \cdot \vec{E}$$

OR $\mathcal{L}_{int} = eA_{\mu}^{em} \mathcal{V}^{\mu}$ (P,T-even)
 $+ eA_{\mu}^{em} \mathcal{A}^{\mu}$ (P,T-odd)

EDMs are the most sensitive probes of CPv:

- Signals for beyond SM physics (SM = 10^{-5} of the current exp.bound)
- Prerequisite for Baryogenesis
- θ_{QCD} -induced EDM : Strong CP problem

$$\langle N_{p'} | J^{\mu} | \bar{N}_{p} \rangle_{CP} = \bar{u}_{p'} \begin{bmatrix} F_{1} \gamma^{\mu} + (F_{2} + iF_{3} \gamma_{5}) \frac{\sigma^{\mu\nu} (p' - p)_{\nu}}{2m_{N}} \end{bmatrix} u_{p}$$

$$\text{Dirac} \qquad \text{Pauli} \qquad \text{Electric dipole}$$

$$\text{(anom.magnetic)}$$

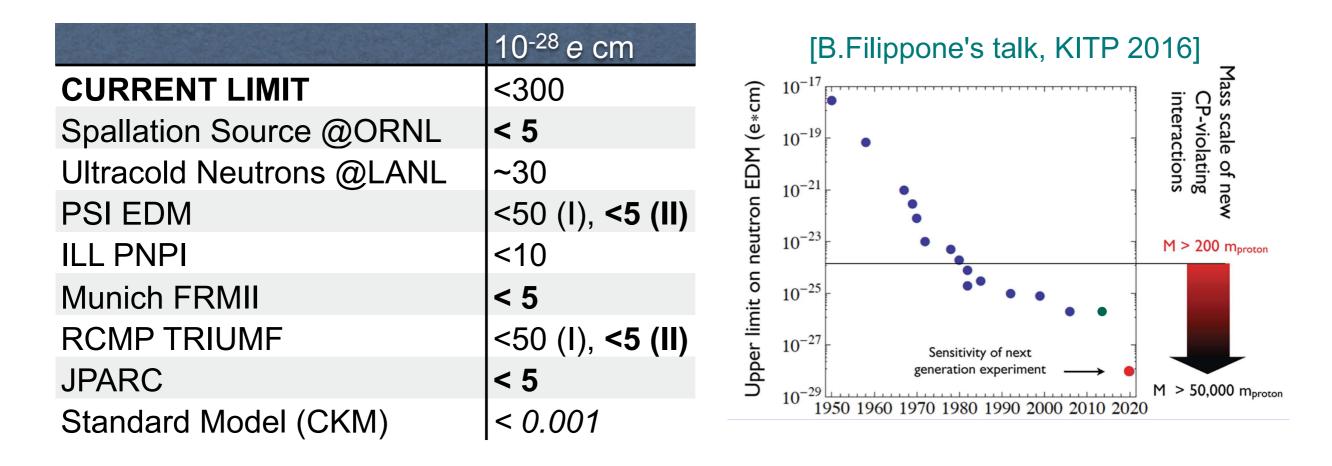
Experimental Outlook

Current nEDM limits:

• $|d_n| < 2.9 \times 10^{-26} e \cdot cm$ [Baker et al, PRL97: 131801(2006)] • $|d_n| < 1.6 \times 10^{-26} e \cdot cm$ [Graner et al, PRL116:161601(2016)]

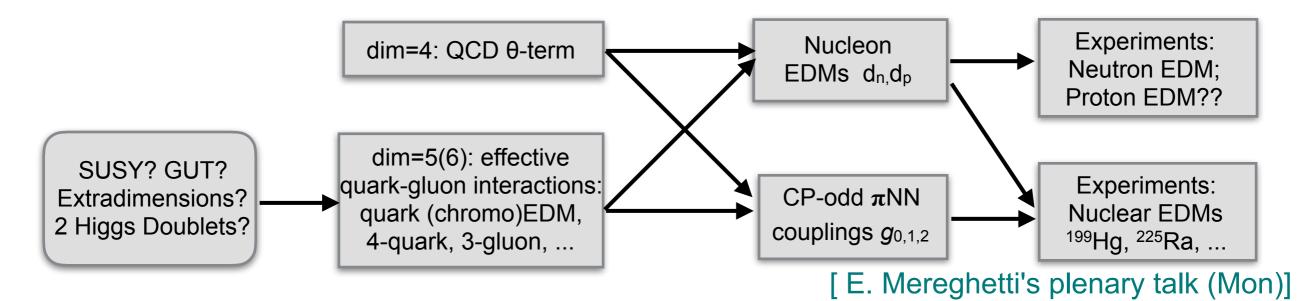
Future nEDM sensitivity :

- 1–2 years : next best limit?
- 3–4 years : x10 improvement
- 7-10 years : x100 improvement



Other experiments: light nuclei in storage rings, octupole-deformed ²²⁵Ra, etc

Nucleon EDMs: a Window into New Physics



Effective quark-gluon CPv interactions organized by dimension

$$\mathcal{L}_{eff} = \sum_{i} \frac{c_i}{[\Lambda_{(i)}]^{d_i - 4}} \mathcal{O}_i^{[d_i]}$$
$$d=4: \theta_{QCD}$$

[J.Engel, M. Ramsey-Musolf, U. van Kolck, Prog.Part.Nucl.Phys. 71 (2013), pp. 21-74]

$$d=5(6)$$
: quark EDM, quark-gluon chromo EDM
 $d=6$: 4-fermion CPv, 3-gluon (Weinberg)

$$egin{aligned} & d_{n,p} \ & F_3^{n,p}(Q^2) \end{aligned}$$

 $c_i \iff d_{n,p}$?

•
$$d_{n,p} = d_{n,p}^{\theta} \theta_{\text{QCD}} + d_{n,p}^{cEDM} c_{cEDM} + \dots$$

lattice QCD calculations are needed to constrain θ_{QCD} , C_{CEDM} , ...

CP-odd Nucleon Structure on a Lattice

CPv interaction induces a chiral phase in fermion fields:

To determine $F_{2,3}$ correctly, one has to use positive-parity spinors [M.Abramczyk, S.Aoki, S.N.S, *et al* (2017) arXiv:1701.07792]

$$\langle N_{p'} | \bar{q} \gamma^{\mu} q | N_p \rangle_{\mathcal{CP}} = \bar{u}_{p'} \Big[F_1 \gamma^{\mu} + (F_2 + i F_3 \gamma_5) \frac{i \sigma^{\mu\nu} (p' - p)_{\nu}}{2m_N} \Big] u_p \text{, with } \frac{\gamma_4 u = +u}{\bar{u}\gamma_4 = +\bar{u}}$$

Prior to 2017, lattice determinations of EDM were subject to large bias from $F_{2,3}$ mixing

$${}^{``}F_3{}^{"} \approx [F_3]_{\text{true}} - 2\alpha [F_2]_{\text{true}}$$
$${}^{`'}d_{n,p}{}^{"} \approx [d_{n,p}]_{\text{true}} - 2\alpha \frac{\kappa_{n,p}}{2m_N}$$

Quark Chromo-EDM on a Lattice

$$\mathcal{L}_{\text{cEDM}} = \sum_{q=u,d} \frac{\tilde{\delta}_q}{2} \,\bar{q} \left[G_{\mu\nu} \sigma^{\mu\nu} \gamma_5 \right] q$$

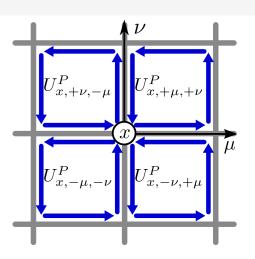
dim-5 operator : O(a⁻²) mixing with dim-3 pseudoscalar density \Rightarrow evaluate&subtract p,nEDM induced by PS density $P = \bar{q}\gamma_5 q$ [T.Bhattacharya et al, 1502.07325]

Chiral symmetry is important:

O(a) clover term in, e.g., Wilson fermion action \equiv chromo-magnetic DM $\mathcal{L}^{clover} = a \frac{c}{4} \bar{q} \left[G_{\mu\nu} \sigma^{\mu\nu} \right] q$

In presense of CPv, condensate is realigned $q \to e^{i\gamma_5\Omega}q$ so that $\langle \operatorname{vac} | \mathcal{L}_m + \mathcal{L}_{CP} | \pi^a \rangle = 0$

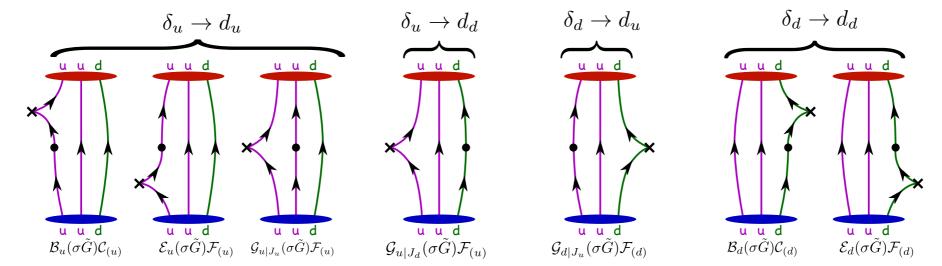
leading to mixing (chromo)EDM \iff (chromo)MDM: $\delta \mathcal{L}_{cEDM} = \delta(\bar{q} [\tilde{D}_q G_{\mu\nu} \sigma^{\mu\nu} \gamma_5] q) = \bar{q} [\{\Omega, \tilde{D}_q\} G_{\mu\nu} \sigma^{\mu\nu}] q) \sim \delta \mathcal{L}_{cMDM}$



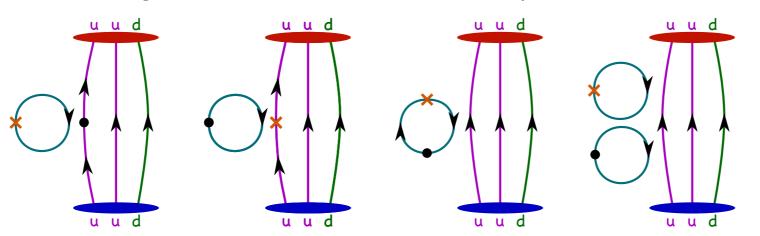
Quark-Gluon EDM: Insertions of dim-5 Operators

$$\mathcal{L}^{(5)} = \sum_{q} \tilde{d}_{q} \,\bar{q}(G \cdot \sigma) \gamma_{5} q \qquad \longleftrightarrow \qquad \langle N(y) \,\bar{N}(0) \int d^{4}x \,\bar{q}(G \cdot \sigma) \gamma_{5} q \rangle \\ \langle N(y) \,[\bar{\psi}\gamma^{\mu}\psi]_{z} \,\bar{N}(0) \,\int d^{4}x \,\bar{q}(G \cdot \sigma) \gamma_{5} q \rangle$$

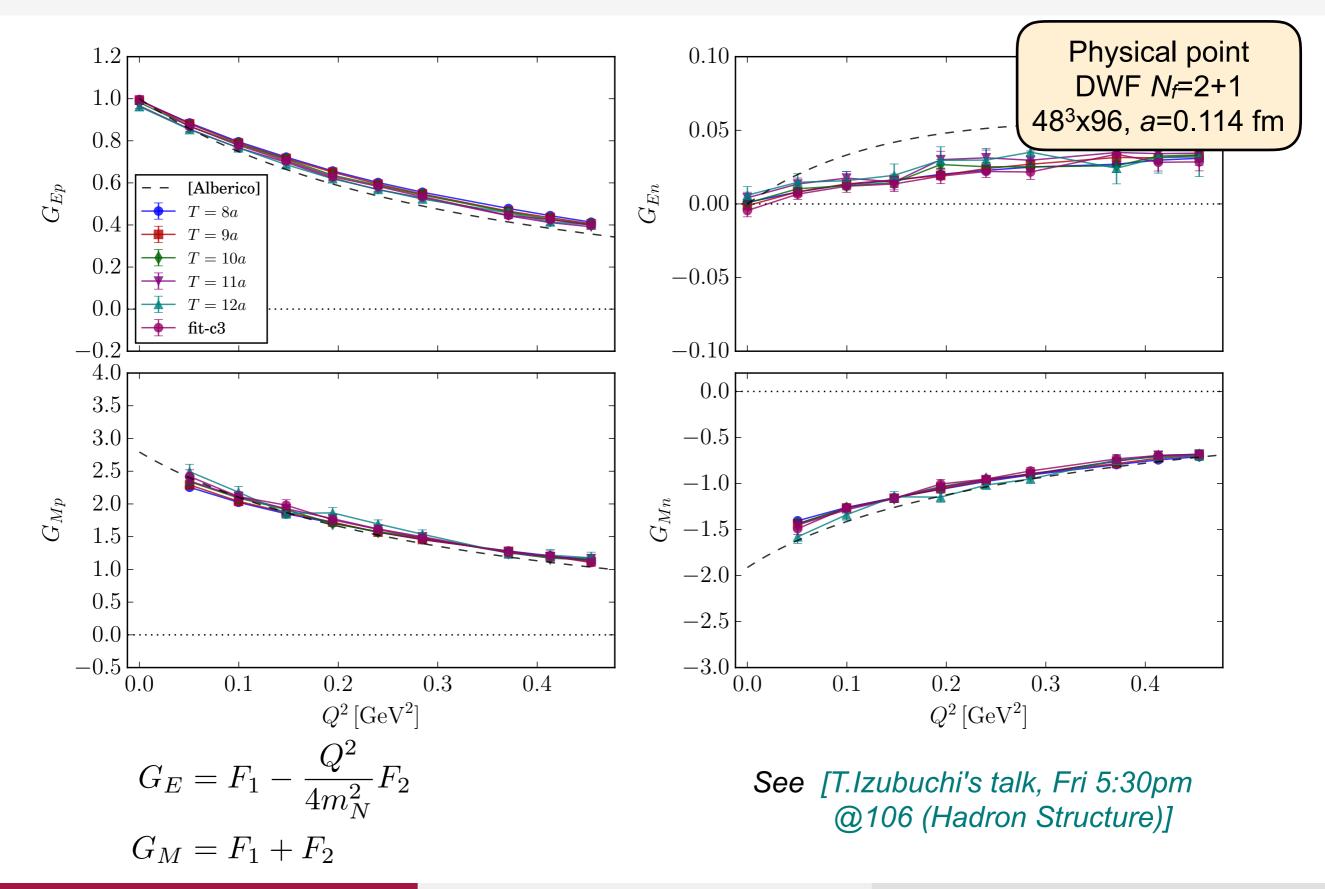
This work: Only quark-connected insertions



In future: Single- and double-disconnected diagrams (contribute to isosinglet cEDM, mix with θ-term)



Nucleon Vector (Sachs) Form Factors



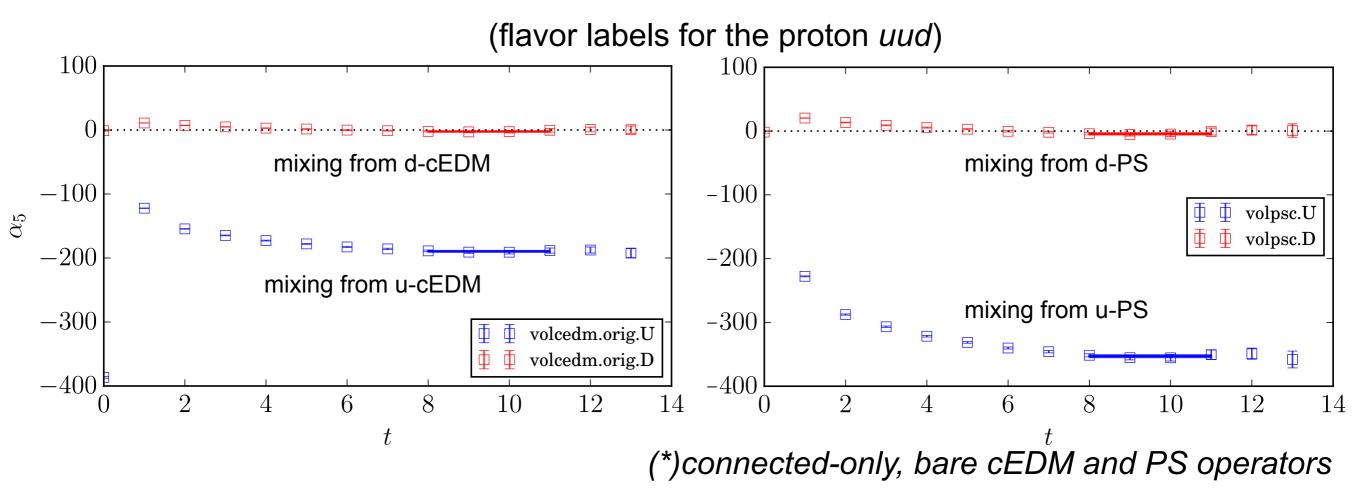
Parity Mixing : cEDM and pseudoscalar(*)

$$N_{\delta} = \epsilon^{abc} \, u^a_{\delta} \left(u^{aT} \mathcal{C} \gamma_5 d^c \right)$$

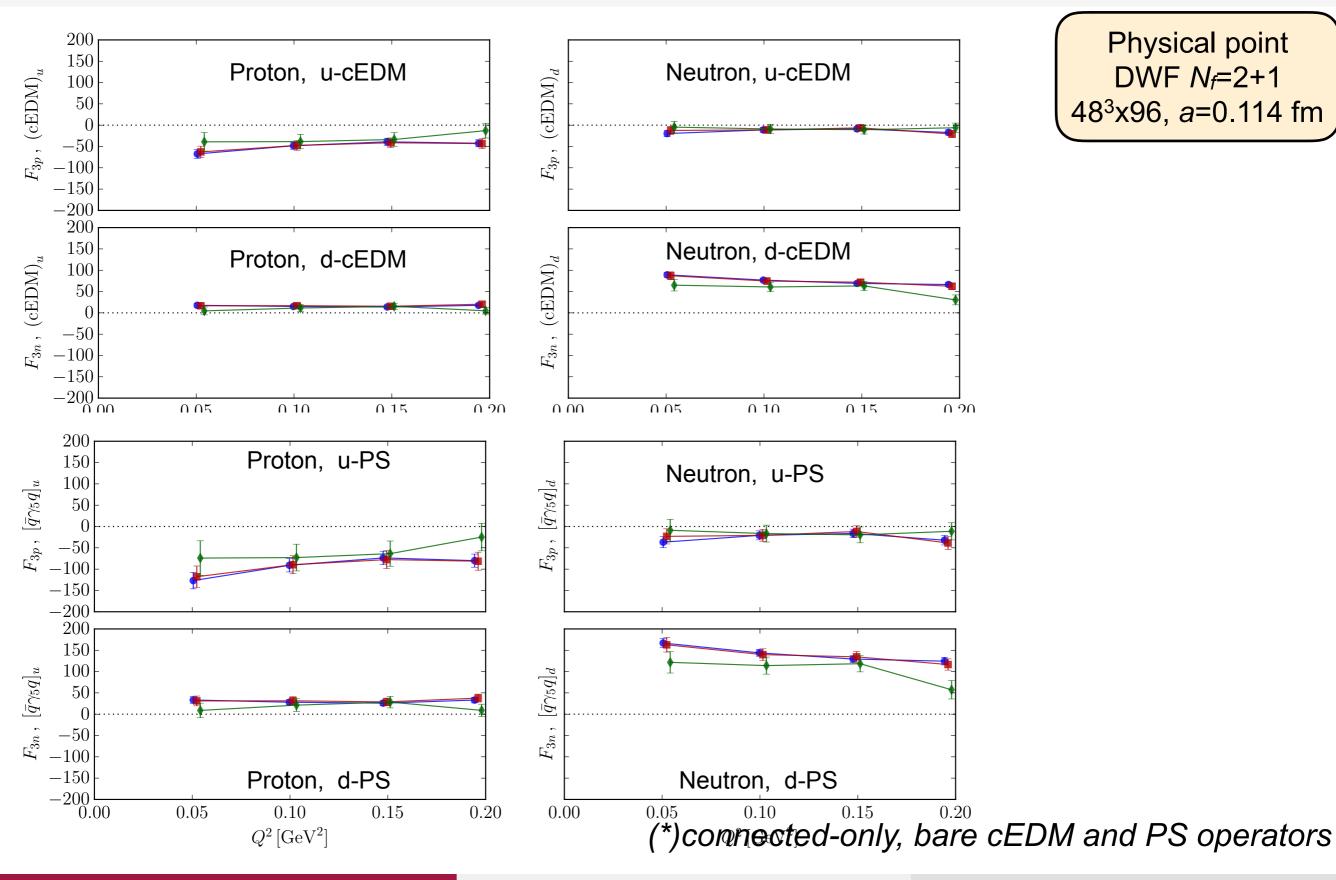
$$\langle N(t)\bar{N}(0)\rangle_{\mathcal{CP}} = \frac{-i\not\!\!\!/ + m_N e^{2i\alpha_5\gamma_5}}{2m_N} e^{-E_N t}$$

$$\hat{\alpha}_5 = \frac{\alpha_5}{\tilde{d}} = -\frac{\operatorname{ReTr}\left[T^+\gamma_5 \cdot C_{2pt}^{\overline{CP}}(t)\right]}{\operatorname{ReTr}\left[T^+ \cdot C_{2pt}^{CP}(t)\right]}, \quad t \to \infty$$

Physical point DWF *N_f*=2+1 48³x96, *a*=0.114 fm



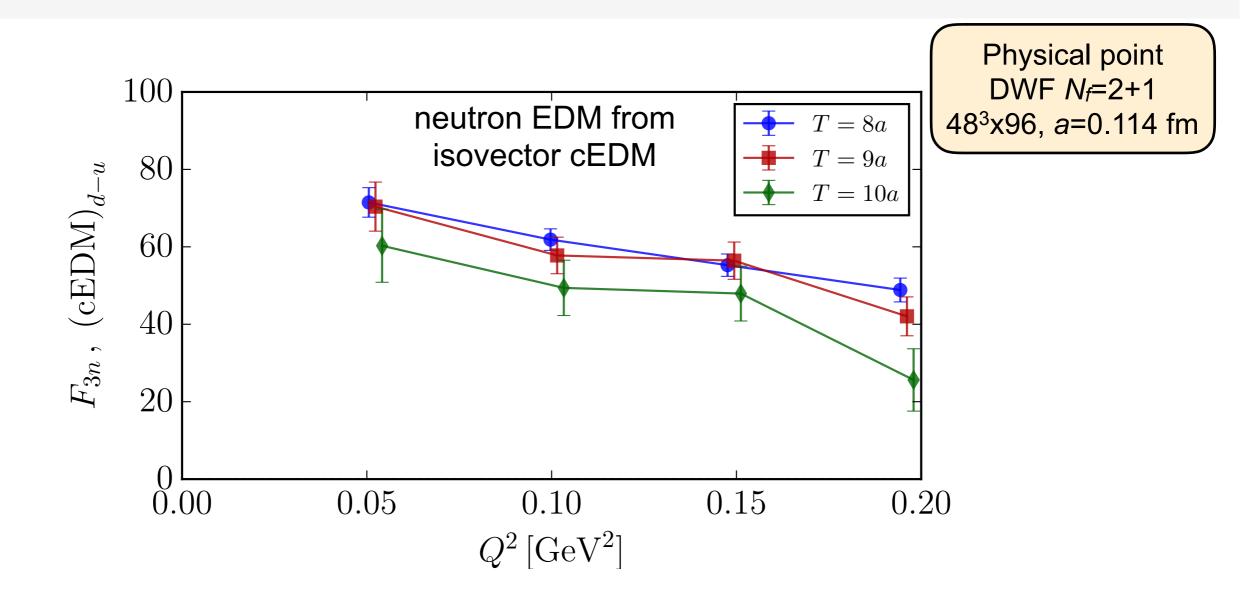
Proton & Neutron EDM Form Factors (*)



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Nucleon EDMs on a Lattice at the Physical Point

Neutron EDM from Isovector Quark cEDM



Outlook for cEDM-induced p,nEDM

- Renormalization & mixing subtractions : work underway using position-space scheme
- Flavor-dependent CPv from cEDM : disconnected diagrams are required, will be challenging due to noise and mixing with θ_{QCD} term

θ_{QCD}-induced nEDM : Status

Correction to previous results:

 $[F_3]_{\text{true}} = "F_3" + 2\alpha F_2$

0.10.0 -0.1 $(0)^{\circ}H^{\circ} = -0.2$ -0.3-0.4 $m_{\pi} = 465 MeV$ -0.5-0.62.02.53.0 1.0 1.50.5

[M.Abramczyk, S.A $^{\theta}$ oki, S.N.S., *et al*, (2017)]

		$m_{\pi} [{ m MeV}]$	$m_N [{ m GeV}]$	F_2	α	$ ilde{F}_3$	F_3
[ETMC 2016]	\overline{n}	373	1.216(4)	$-1.50(16)^{a}$	-0.217(18)	-0.555(74)	0.094(74)
[Shintani et al 2005]	$\int n$	530	1.334(8)	-0.560(40)	$-0.247(17)^{b}$	-0.325(68)	-0.048(68)
	p	530	1.334(8)	0.399(37)	$-0.247(17)^{b}$	0.284(81)	0.087(81)
[Berruto et al 2006]	$\hat{\mathbf{f}}$ n	690	1.575(9)	-1.715(46)	-0.070(20)	-1.39(1.52)	-1.15(1.52)
	n	605	1.470(9)	-1.698(68)	-0.160(20)	0.60(2.98)	1.14(2.98)
[Guo et al 2015]	$\int n$	465	1.246(7)	$ -1.491(22)^{c}$	$-0.079(27)^d$	-0.375(48)	$-0.130(76)^d$
	n	360	1.138(13)	$ -1.473(37)^c$	$-0.092(14)^d$	-0.248(29)	$0.020(58)^d$

After removing the spurious contribution,

• no lattice signal for θ_{QCD} -induced nEDM $rightarrow d_N$ is very small

• no more conflict with phenomenology values or m_q scaling

θ-Term Noise Reduction for EDM

Variance of lattice θ -induced nEDM signal ~ (Volume)_{4d} :

$$\begin{split} d_N &\sim \langle Q \cdot (N J_\mu \bar{N}) \rangle \\ \text{Top. charge} \ \ Q &\sim \int_{V_4} (G \tilde{G}) \ , \ \ \text{with} \ \langle |Q|^2 \rangle \sim V_4 \end{split}$$

Constrain Q sum to the fiducial volume

 \bigcirc in time around current, $|t_Q - t_J| < \Delta t$ [E.Shintani et al (2015)] \bigcirc in time around source, $|t_Q - t_{source}| < \Delta t$ [J. Dragos, talk on Tue] \bigcirc 4-d sphere around sink, $|x_Q - x_{sink}| < R$ [K.-F.Liu et al, (2017)]:

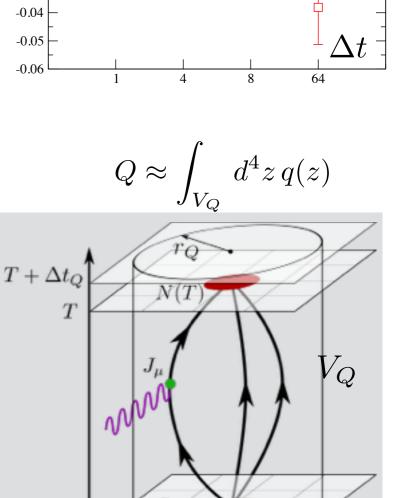


Proper treatment of nucleon parity mixing is critical for correct determination of F₃ \implies nucleon must "settle" in the new $\theta \neq 0$ vacuum

 $N^{(+)} \rightarrow \tilde{N}^{(+)} \approx N^{(+)} + i\alpha N^{(-)}$ $N^{(-)} \rightarrow \tilde{N}^{(-)} \approx N^{(-)} - i\alpha N^{(+)}$

 \implies constrain time and space differently :

4d "cylinder" V_Q : $|\vec{z}| < r_Q$, $-\Delta t_Q < z_0 < T + \Delta t_Q$



[E.Shintani et al (2015)]

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 $F_3(Q^2_{min})$

Ο

0.05

0.04

0.03

0.02

0.01

-0.01

-0.02

-0.03

-0.04

-0.05

-0.06

 $F_3(2\pi/L)/2m$

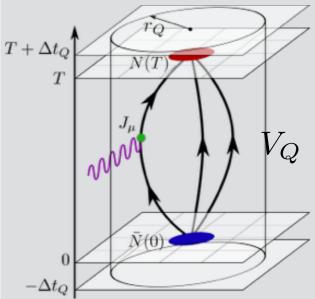
NO

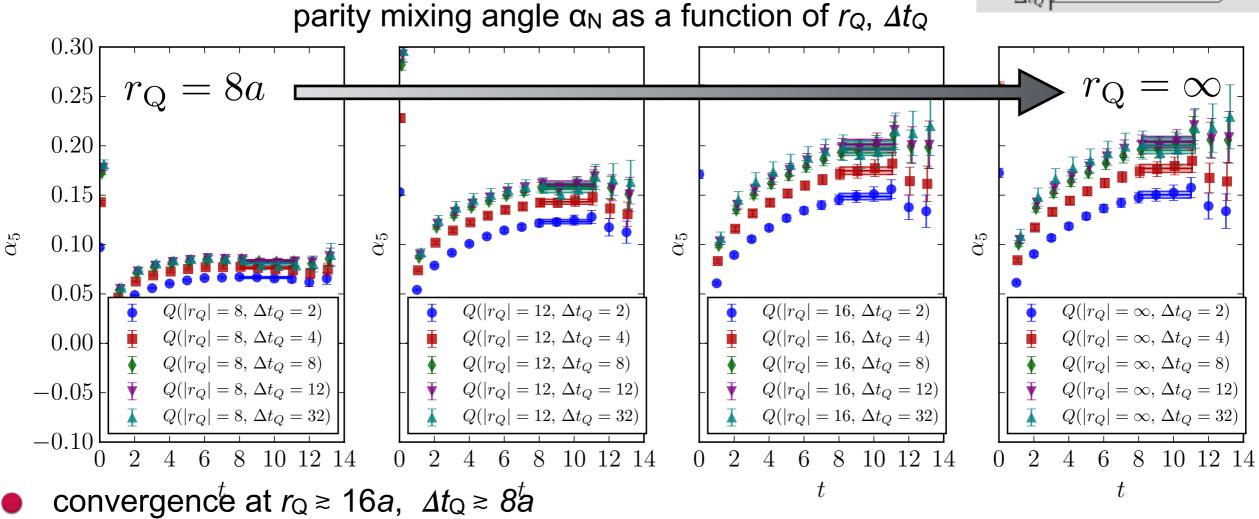
0

 $-\Delta t_O$

Tests on m_{π} =330 MeV Lattices: Parity Mixing

- $N_f=2+1$ Domain Wall (RBC/UKQCD) 24³x64 a = 0.114 fm
- 1400 confiigs * (64sloppy+1exact) samples \implies 89.6k stat.
- Top.charge with 5-loop improved GG [P. de Forcrand et al '97] on Wilson-flowed (t=8a²) gauge links [M.Luscher, 1006.4518]

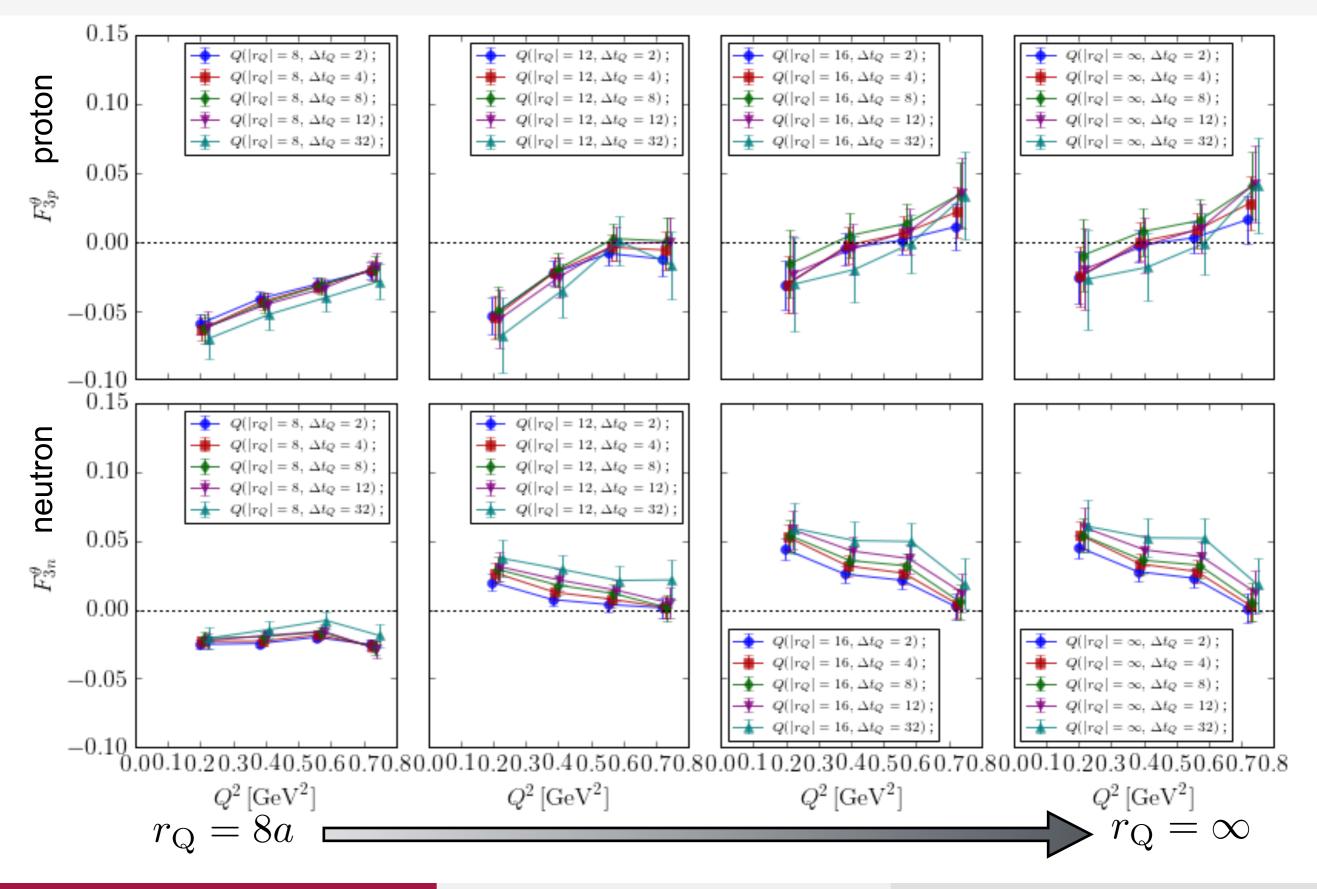




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Tests on m_π=330 MeV Lattices: EDM(Form Factor)



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How Hard is θ -nEDM at the Physical Point?

• chiral fermions, m_{π} = 330 MeV [this work]

$$2m_n d_n | = |F_{3n}(0)| \approx 0.05 \cdot \theta$$

• Wilson fermions, $m\pi$ =360 MeV [Guo et al 2015] after correction

 $|2m_n d_n| = |F_{3n}(0)| \leq 0.06 \cdot \theta$

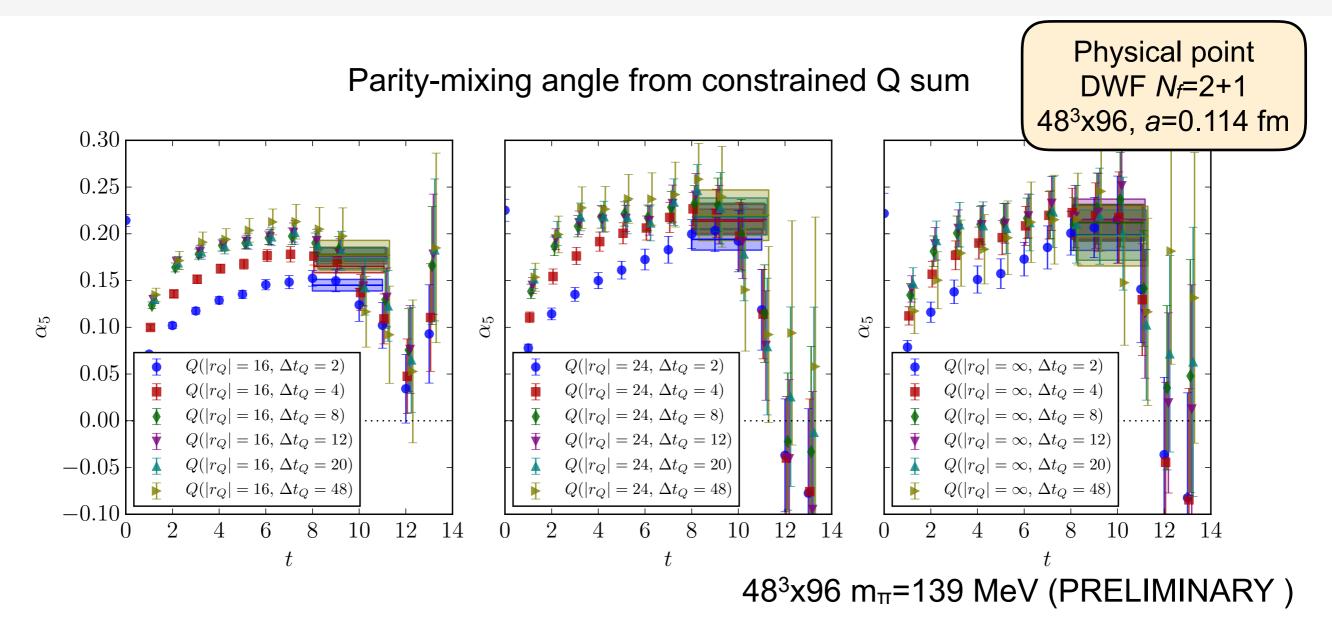
• best guess for the physical point with $|d_n| \sim m_q \sim (m_\pi)^2$

 \implies phys.point

 $|\mathsf{F}_{3n}(0)| \approx 0.01 \cdot \theta, |d_n| \approx 0.001 \cdot \theta e \text{ fm}$

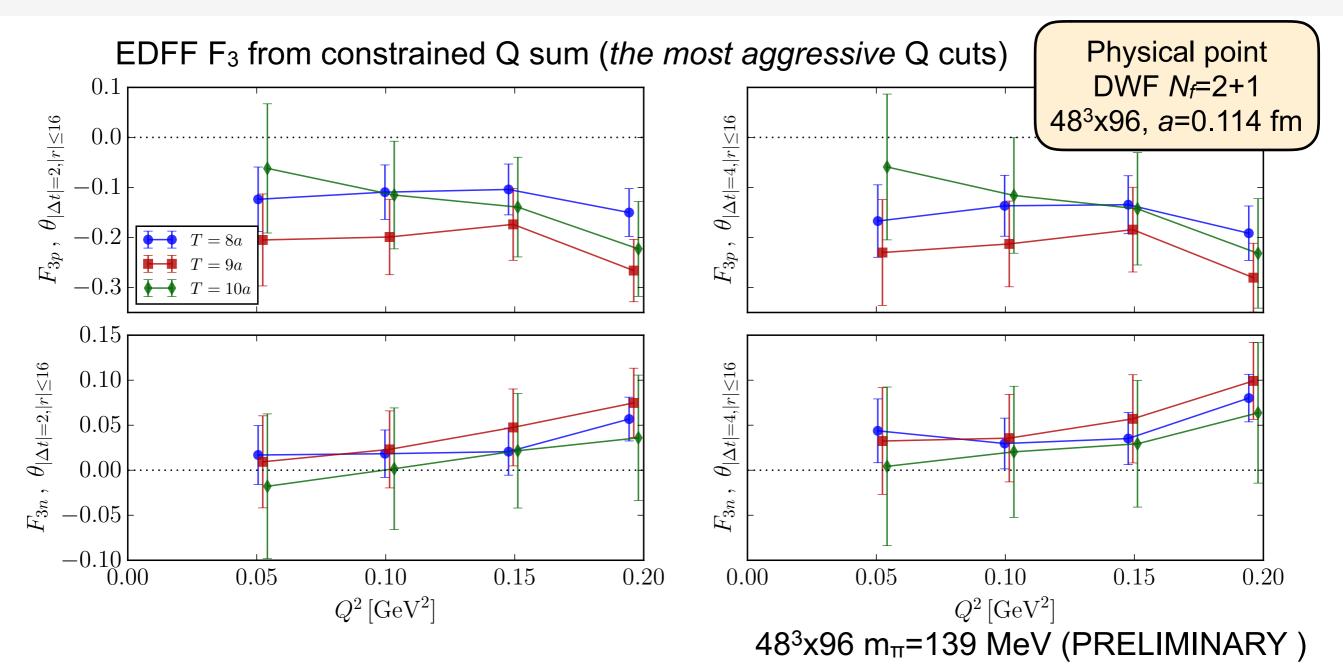
 $|F_{3n}^{\text{phys}}(0)| \sim O(10^{-2}) \,\theta, \quad |d_n| \sim O(10^{-3}) \,e\,\text{fm}\,\theta$

Physical point : θ_{QCD} **-induced Parity Mixing** α_N



Reassuring results for noise reduction at the physical point

Physical point : θ_{QCD} **-induced EDFF F**₃



- 33k lattice samples, ~ 30 M core-hours on Argonne BlueGene/Q
- connected diagrams only
- result compatible with zero, $|F_{3n}| \le 0.05$ constraint

Need x30..100 more statistics to constrain $|F_{3n}| \approx 0.01$: θ -nEDM remains difficult at the physical point

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Nucleon EDM : Summary

Encouraging physical-point results for nucleon EDM induced by quark chromo-EDM

~20% stochastic uncertainty for quark cEDM-induced EDM Renormalization & mixing subtractions are underway Full flavor dependence will require disconnected diagrams & θ_{QCD}-term

Clear signal for θ_{QCD}-induced nEDM at m_π = 330 MeV Variance-reduction for Q sampling is essential Physical |d_{n,p}|≈10⁻³ e fm values are in agreement with phenomenology

Constraining θ_{QCD}-induced nEDM at the physical point will be challenging

O(300-1000) M core*hours may be required even with variance reduction Shall look for alternative methods: dynamical θ -therm?