

Exploring the convergence radius of $\text{HB}\chi\text{PT}$

Karl Sallmen

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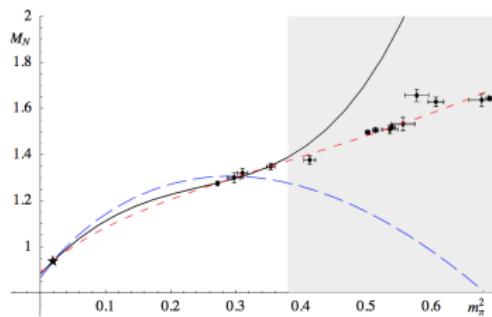


Introduction

Exploring the
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Due to the non-perturbative nature of QCD at low-energies and the frequent use of chiral extrapolations of lattice QCD data, there exists looming questions about the convergence of chiral expansions. In particular for the range of quark-masses being used on the lattice.



motivation

Show that including higher order terms in the correction of m_N from chiral EFT ruins the agreement with data for values of the pion mass larger than ($m_\pi = 600\text{MeV}$)

This is also the case for g_A !

With data for m_N and g_A in a range of "smaller" pion masses, perhaps we could attempt to address the question of convergence.

At higher orders in $\text{HB}\chi\text{PT}$, starting at N2LO ($\mathcal{O}(m_\pi^3)$), g_A couples to m_n

$$\delta m_n = -\frac{3g_A^2}{32\pi F_\pi} m_\pi^3$$

To capture the effect of the pion mass dependence of g_A , it is advantageous to perform a combined analysis (simultaneous fitting) of both m_N and g_A

HB χ PT "quickly"

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Incorporating heavy fermionic fields into χ PT, despite being chirally consistent, unfortunately has the undesired effect of introducing a mass that is of the same order as the chiral symmetry breaking scale $m_{hf}/\Lambda_\chi \sim 1$, and thus there is no longer a consistent derivative expansion for processes involving these heavy fermionic fields.

Additionally Weinberg's power counting fails since higher order loop graphs can produce amplitudes which are no longer suppressed by Λ_χ .

A solution to this problem was first provided by Elizabeth Jenkins and Aneesh V. Manohar.

HB χ PT "quickly"

Nearly on-shell Baryon with velocity v_μ has a momentum given by

$$p_\mu = m_B v_\mu + k_\mu$$

Where k_μ is a small off-shell contribution

The effective theory is then written in terms Baryon fields B_v

$$B_v = e^{im_B \not{v} \mu x^\mu} B(x)$$

Derivatives act on B_v by producing k , the higher derivative terms are thus suppressed by $1/\Lambda_\chi$. In addition the heavy Baryon Lagrangian has a $1/m_B$ expansion. For the lightest Baryon multiplets this can be combined with $1/\Lambda_\chi$ into a single expansion in $1/\Lambda_\chi$. Our theory has now a consistent power counting expansion.

Gathering expressions for m_N , g_A

Up to order m_π^5

$$m_N = m_0 + \delta m_N^{2+4} + \delta m_N^{3+5}$$

$$\begin{aligned}\delta m_N^{2+4} = & -4c_1 m_\pi^2 + m_\pi^4 \left[-e'(L_\chi) + \frac{3}{128\pi^2 f_\pi^2} \left(c_2 - \frac{2g_A^2}{M} \right) \right] \\ & - \frac{m_\pi^4}{32\pi^2 f_\pi^2} \ln \frac{m_\pi}{L_\chi} \left(\frac{3g_A^2}{M} - 32c_1 + 3c_2 + 12c_3 \right)\end{aligned}$$

and

$$\begin{aligned}\delta m_N^{3+5} = & -\frac{3g_A m_\pi^3}{32\pi f_\pi^2} + \frac{3g_A m_\pi^5}{32\pi f_\pi^2} \left[\frac{2l_4^r(L_\chi)}{f^2} - \frac{4(2d_{16}^r(L_\chi) - d_{18})}{g} \right. \\ & \left. + 16d_{28}^r(L_\chi) + \frac{g^2}{8\pi^2 f^2} + \frac{1}{8M^2} \right] + \frac{3g_A^4 m_\pi^5}{64\pi^3 f_\pi^4} \ln \frac{m_\pi}{L_\chi}\end{aligned}$$

Gathering expressions for m_N , g_A

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$$g_A = g_0 \left[1 + \left(\frac{\alpha_2}{(4\pi F_\pi)^2} \ln \frac{m_\pi}{L_\chi} + \beta_2 \right) m_\pi^2 + \alpha_3 m_\pi^3 \right. \\ \left. + \left(\frac{\alpha_4}{(4\pi F_\pi)^4} \ln^2 \frac{m_\pi}{L_\chi} + \frac{\gamma_4}{(4\pi F_\pi)^2} \ln \frac{m_\pi}{L_\chi} + \beta_4 \right) m_\pi^4 + \alpha_5 m_\pi^5 \right]$$

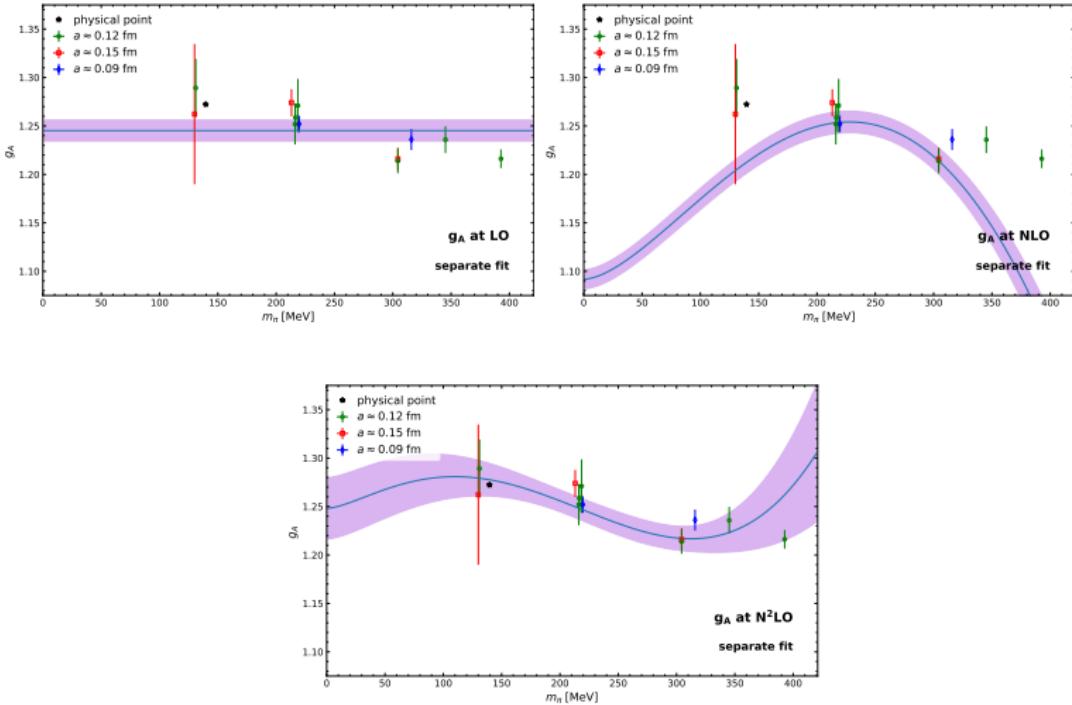
m_N	g_A	$\mathcal{O}(m_\pi)$
c_1	d_{16}^r, d_{28}^r	m_π^2
	c_3, c_4	m_π^3
e', c_2, c_1, c_3	$l_4^r, c_4, c_3, d_{16}^r$	m_π^4
$l_4^r, d_{16}^r, d_{18}, d_{28}^r$	c_3, c_4, l_4^r	m_π^5

table of LEC's appearing at different orders of m_π

plots A

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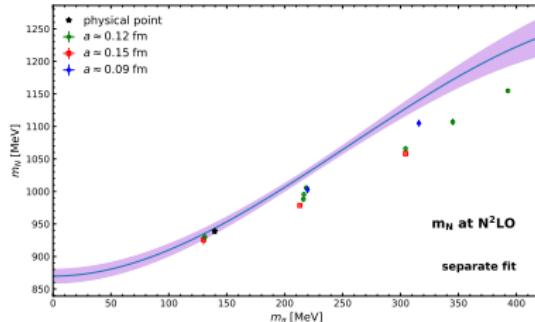
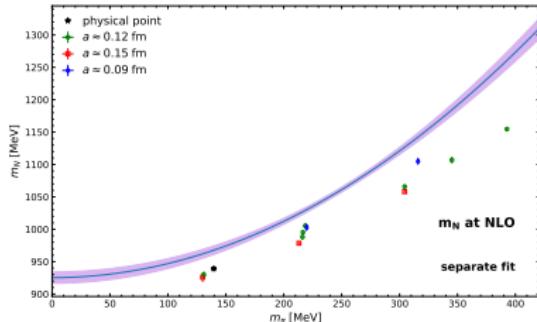
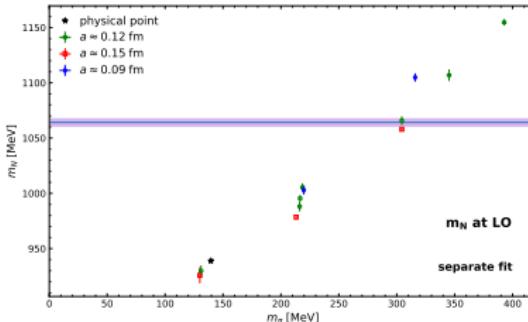
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plots B

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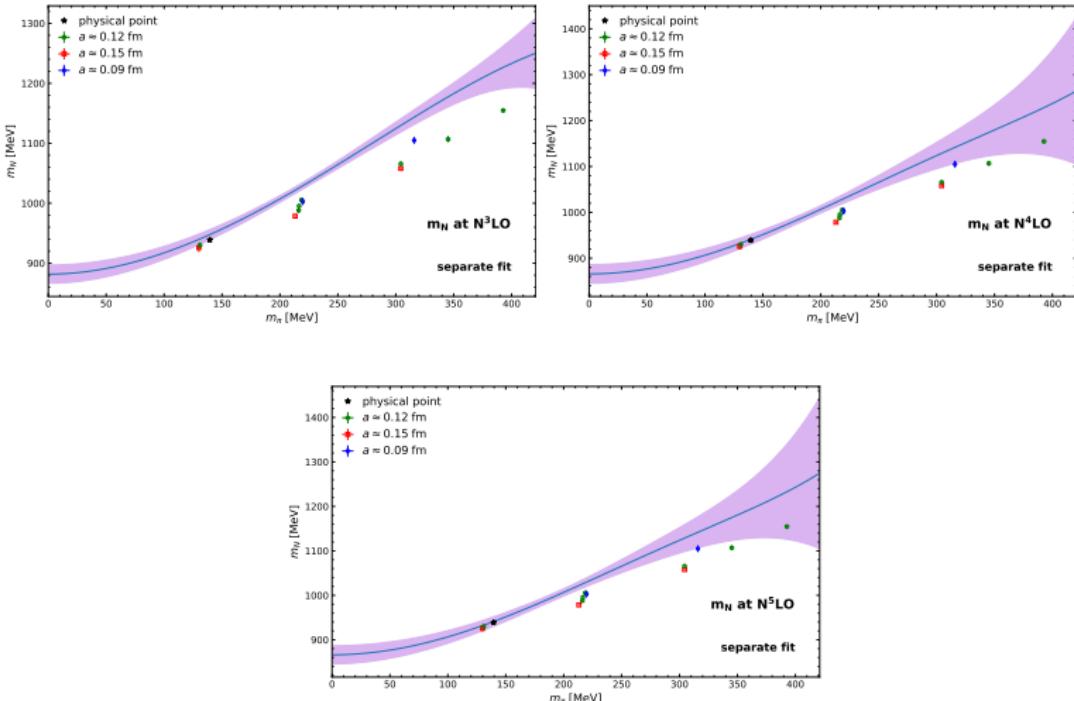
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plot C

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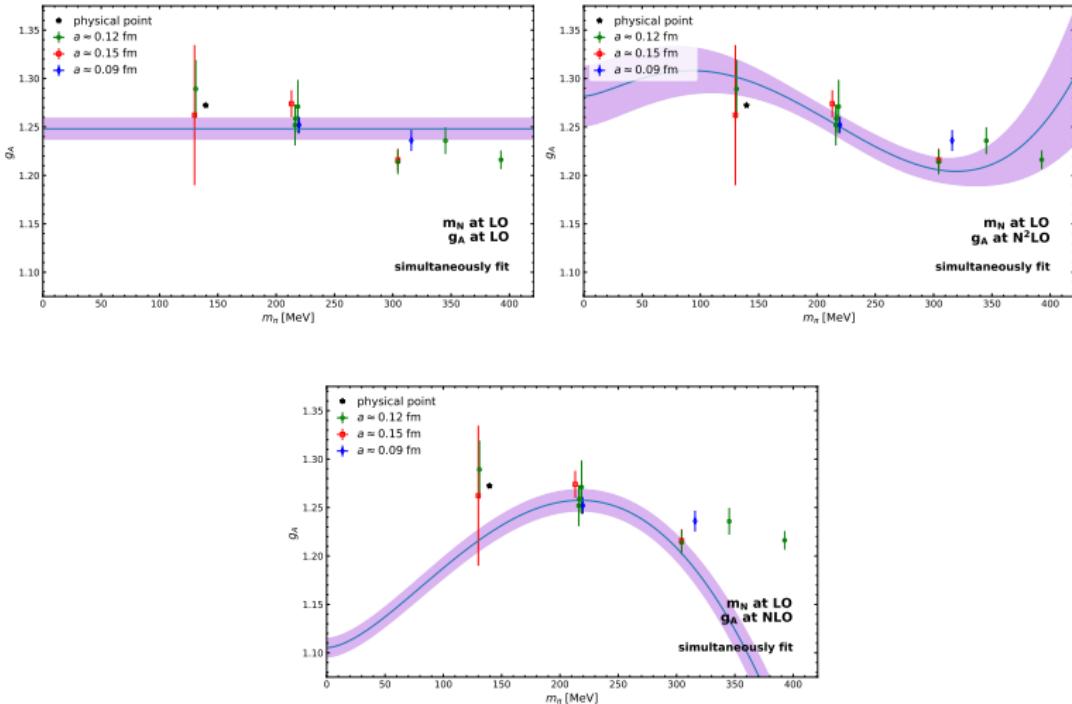
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plot D

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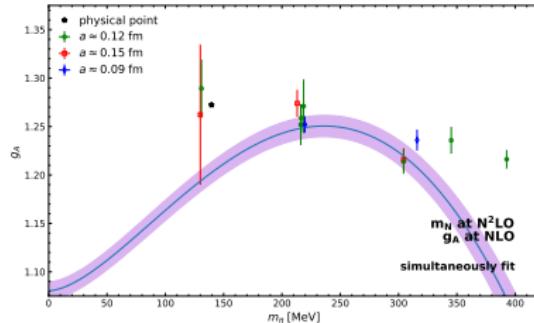
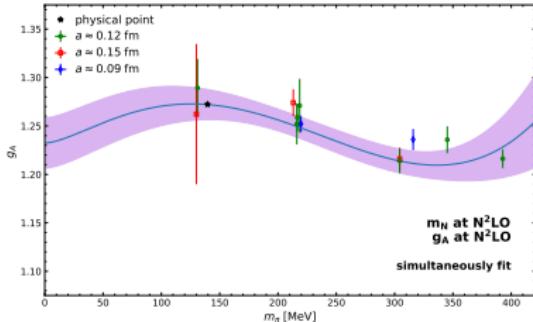
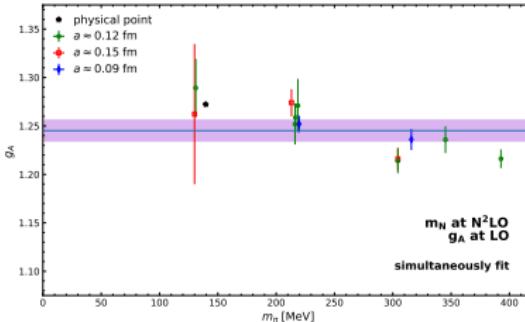
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plot E

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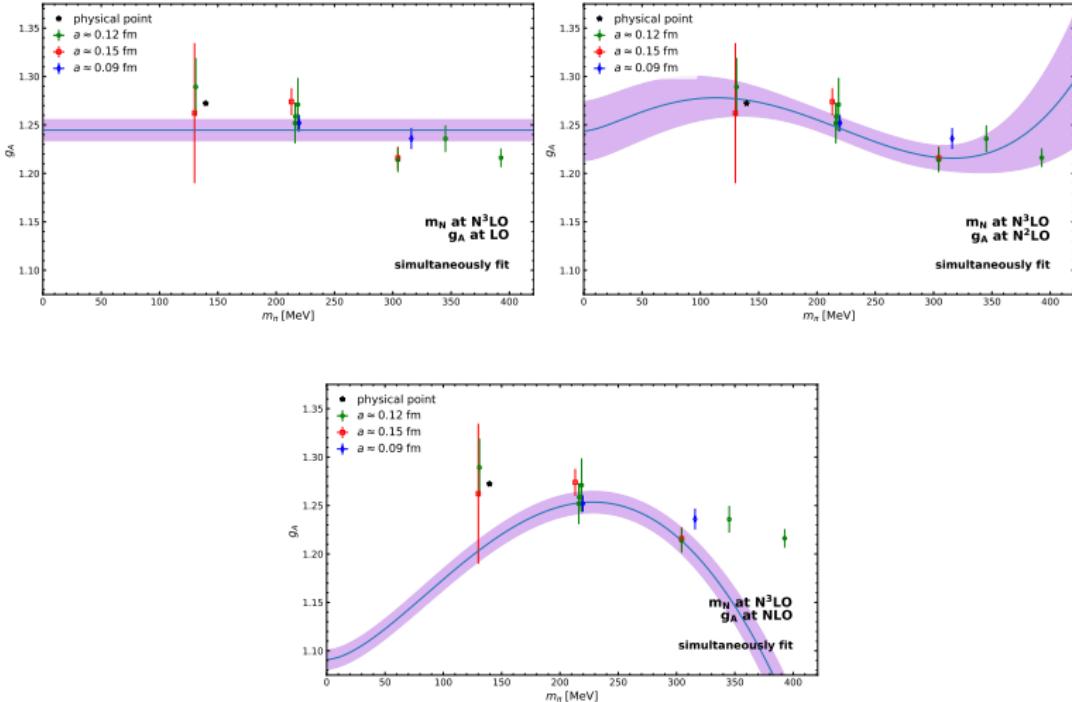
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plot F

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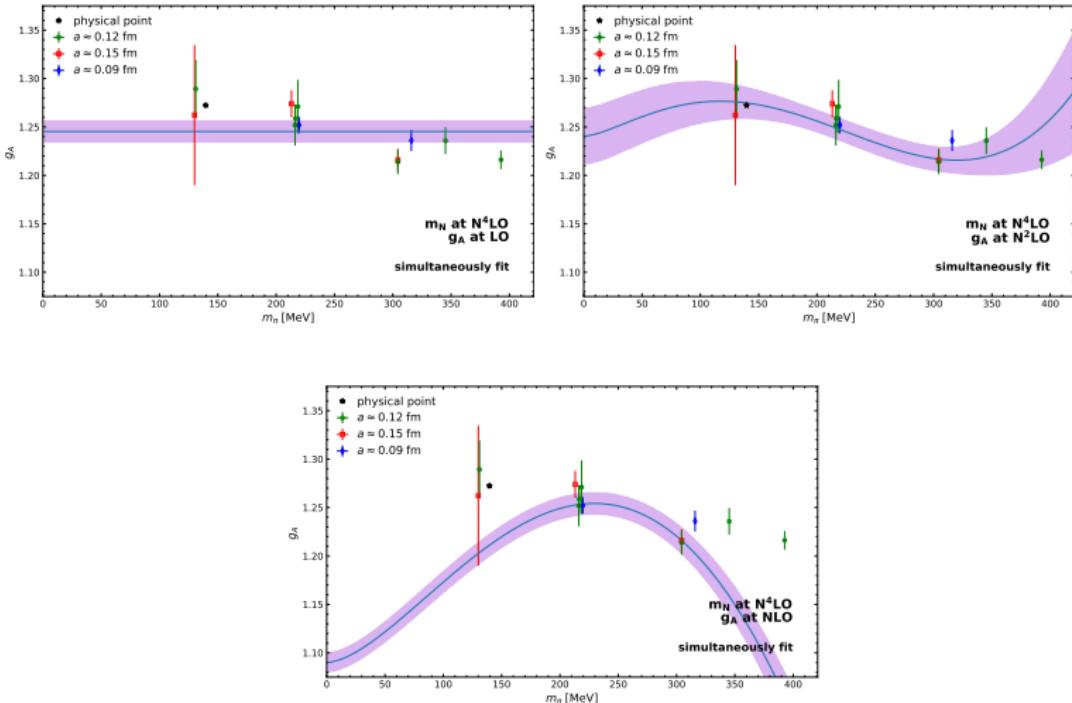
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plot G

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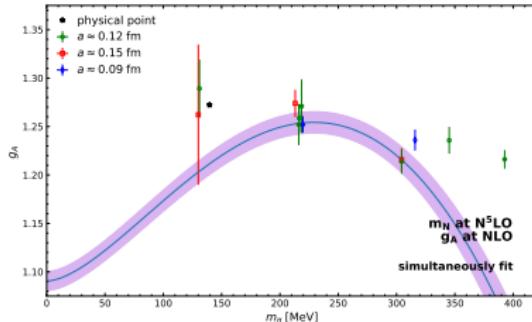
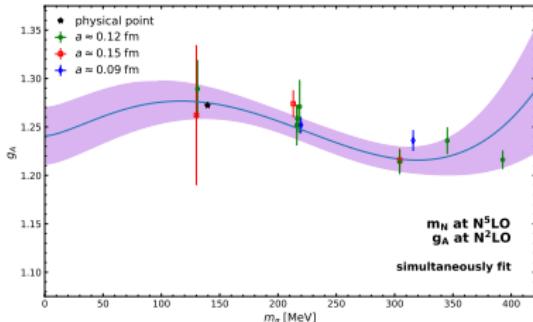
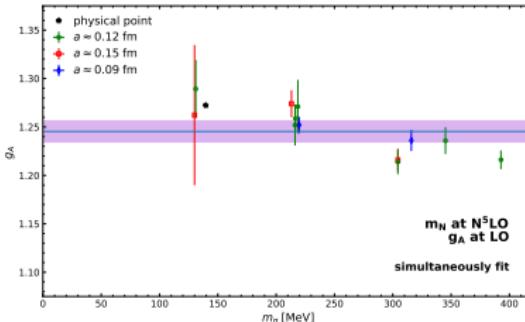
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plot H

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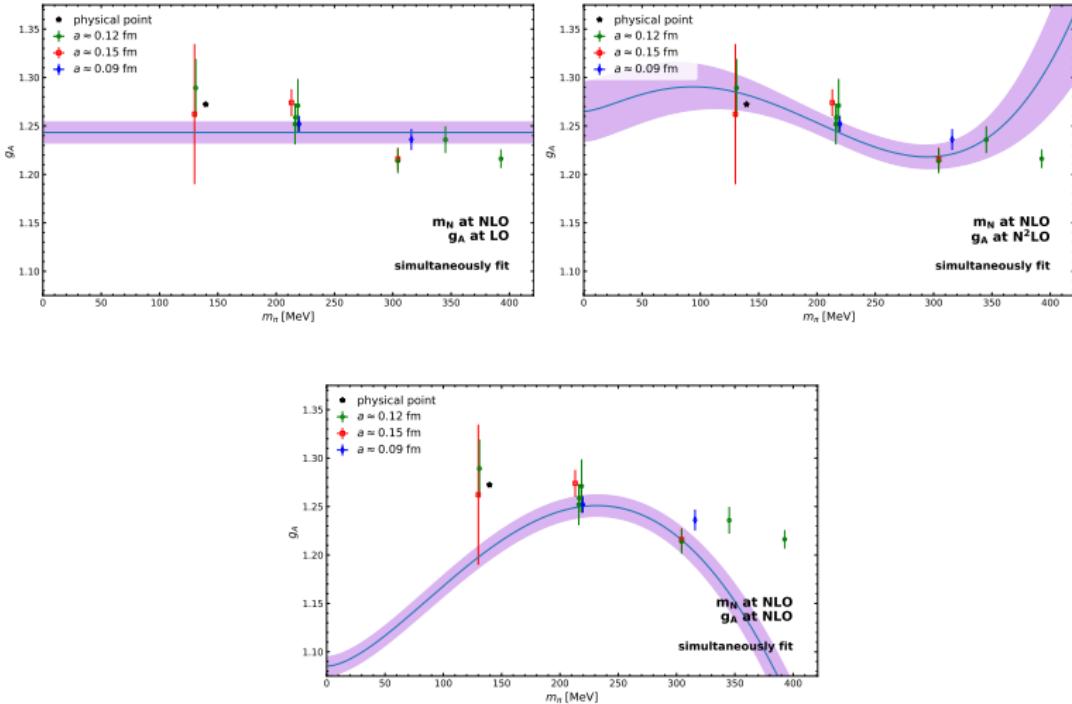
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plot I

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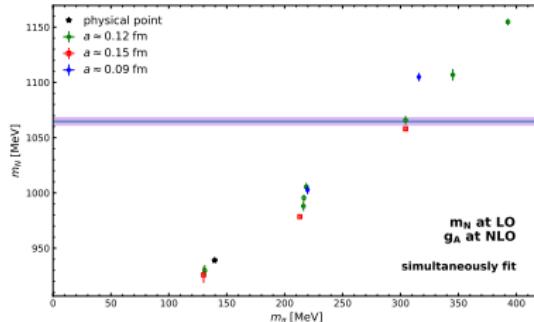
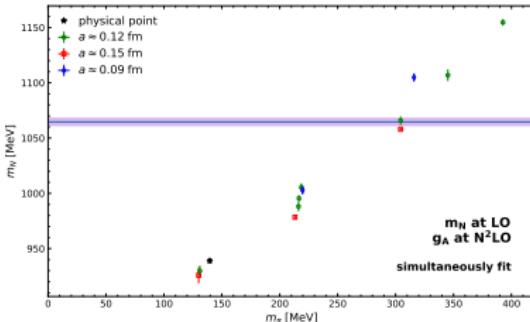
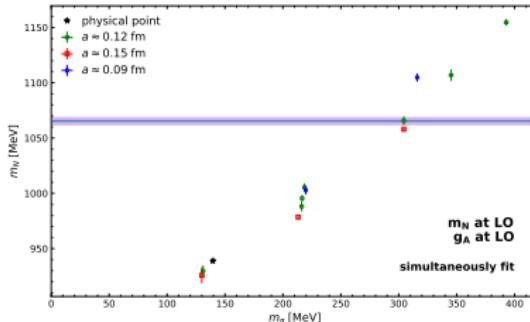
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plot J

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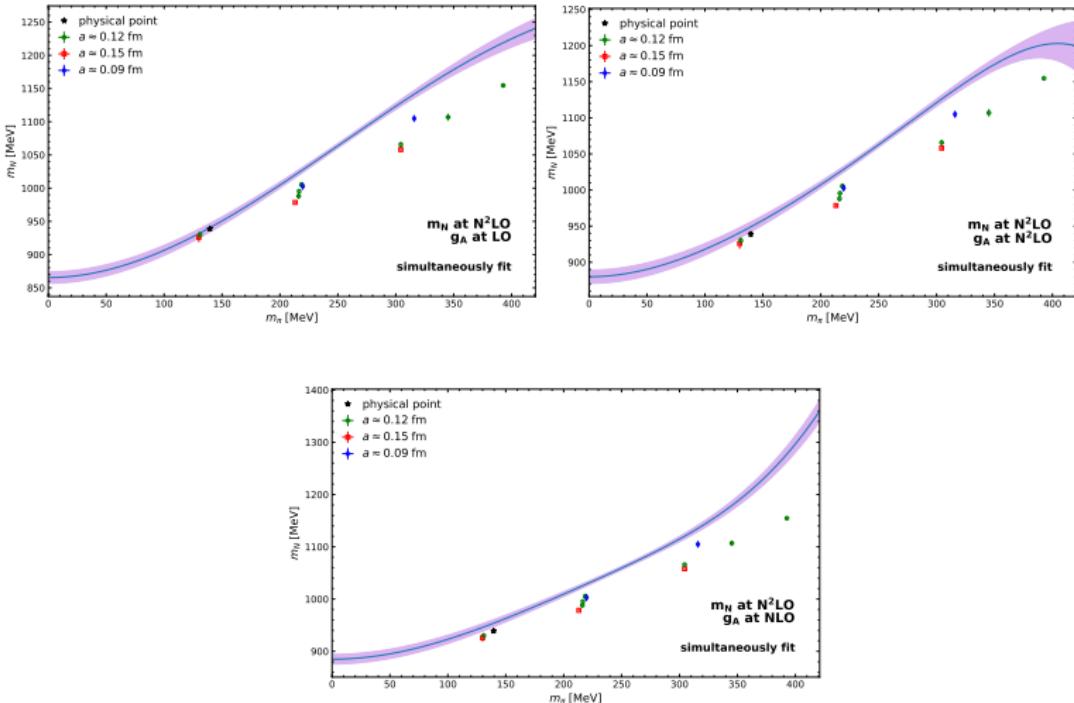
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plot K

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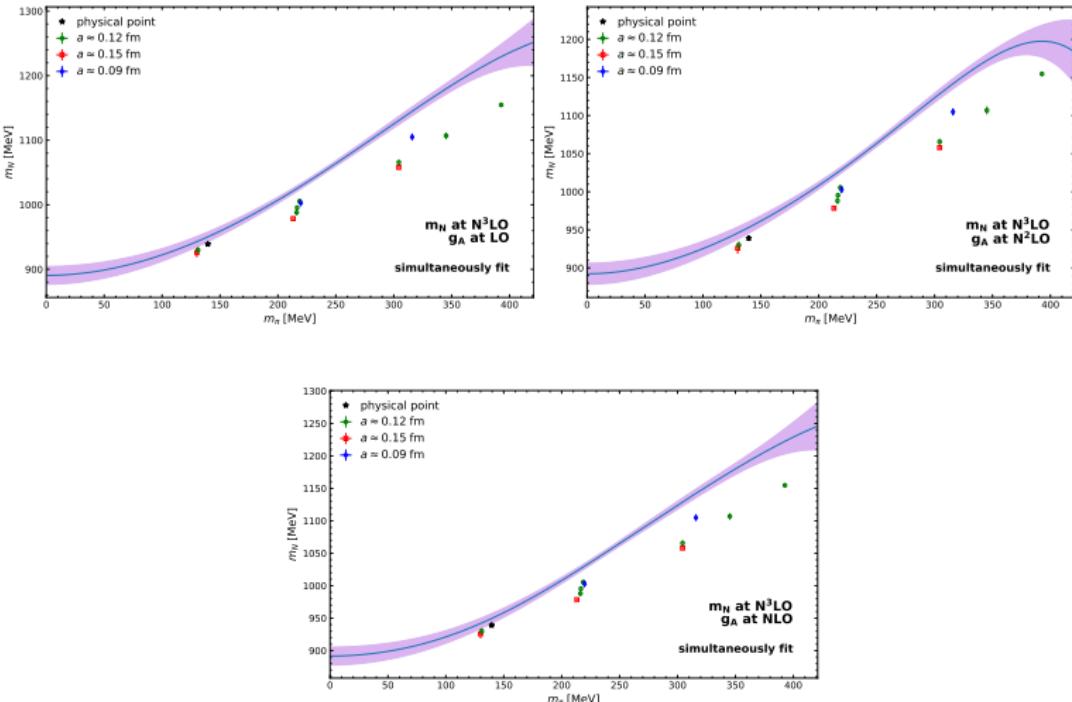
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plot L

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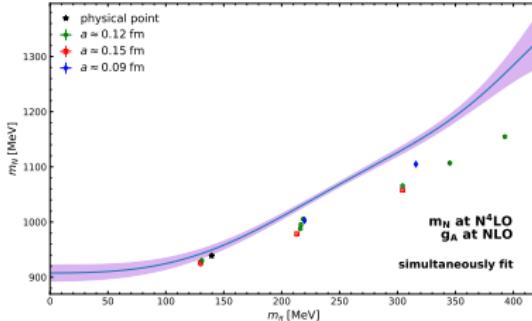
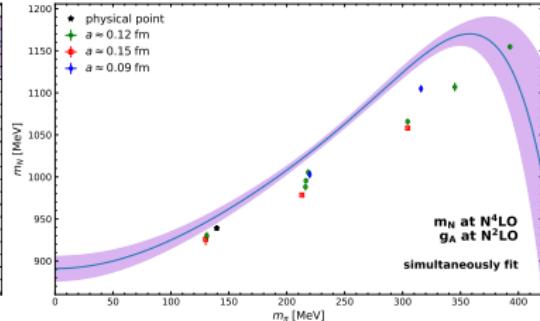
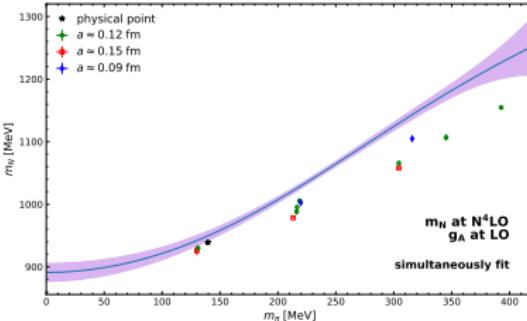
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plot M

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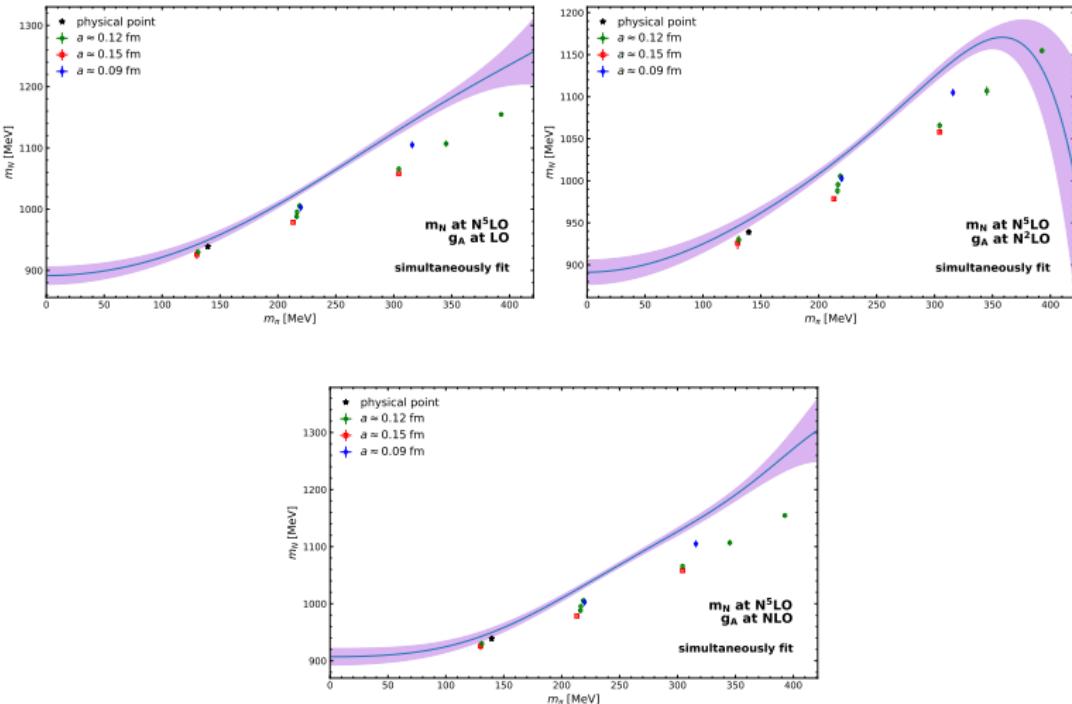
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plot N

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plot 0

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