

High Precision Statistical Landau Gauge Lattice Gluon Propagator Computation

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Abstract

We report on results for the Landau gauge gluon propagator computed from large statistical ensembles and look at the compatibility of the results with the Gribov-Zwanziger tree level prediction for its refined and very refined versions. Our results show that the data is well described by the tree level estimate only up to momenta $p \lesssim 1$ GeV, while clearly favoring the so-called Refined Gribov-Zwanziger scenario. We also provide a global fit of the lattice data which interpolates between the above scenario at low momenta and the usual continuum one-loop renormalization improved perturbation theory after introducing an infrared log-regularizing term.

Lattice Landau gauge gluon propagator

$$\langle A_\mu^a(\hat{p}) A_\nu^b(\hat{p}') \rangle = L^4 \delta^{ab} \delta(\hat{p} + \hat{p}') P_{\mu\nu}(p) D(p^2)$$

► lattice momentum

$$\hat{p}_\mu = \frac{2\pi}{aL} n_\mu, \quad n_\mu = 0, 1, \dots, L-1$$

► continuum momentum

$$p_\mu = \frac{2}{a} \sin\left(\frac{\pi}{L} n_\mu\right)$$

► orthogonal projector

$$P_{\mu\nu}(p) = \delta_{\mu\nu} - \frac{p_\mu p_\nu}{p^2}$$

Lattice simulations

► two large physical volume lattice simulations

Duarte, Oliveira, Silva, PRD94(2016)014502

► Wilson gauge action, $\beta = 6.0$

$1/a = 1.943$ GeV, $a = 0.1016(25)$ fm,

► 64^4 and 80^4 lattices

► physical volumes: $(6.57 \text{ fm})^4$, $(8.21 \text{ fm})^4$

► number of configurations: 2000, 550

► rotated to the Landau gauge

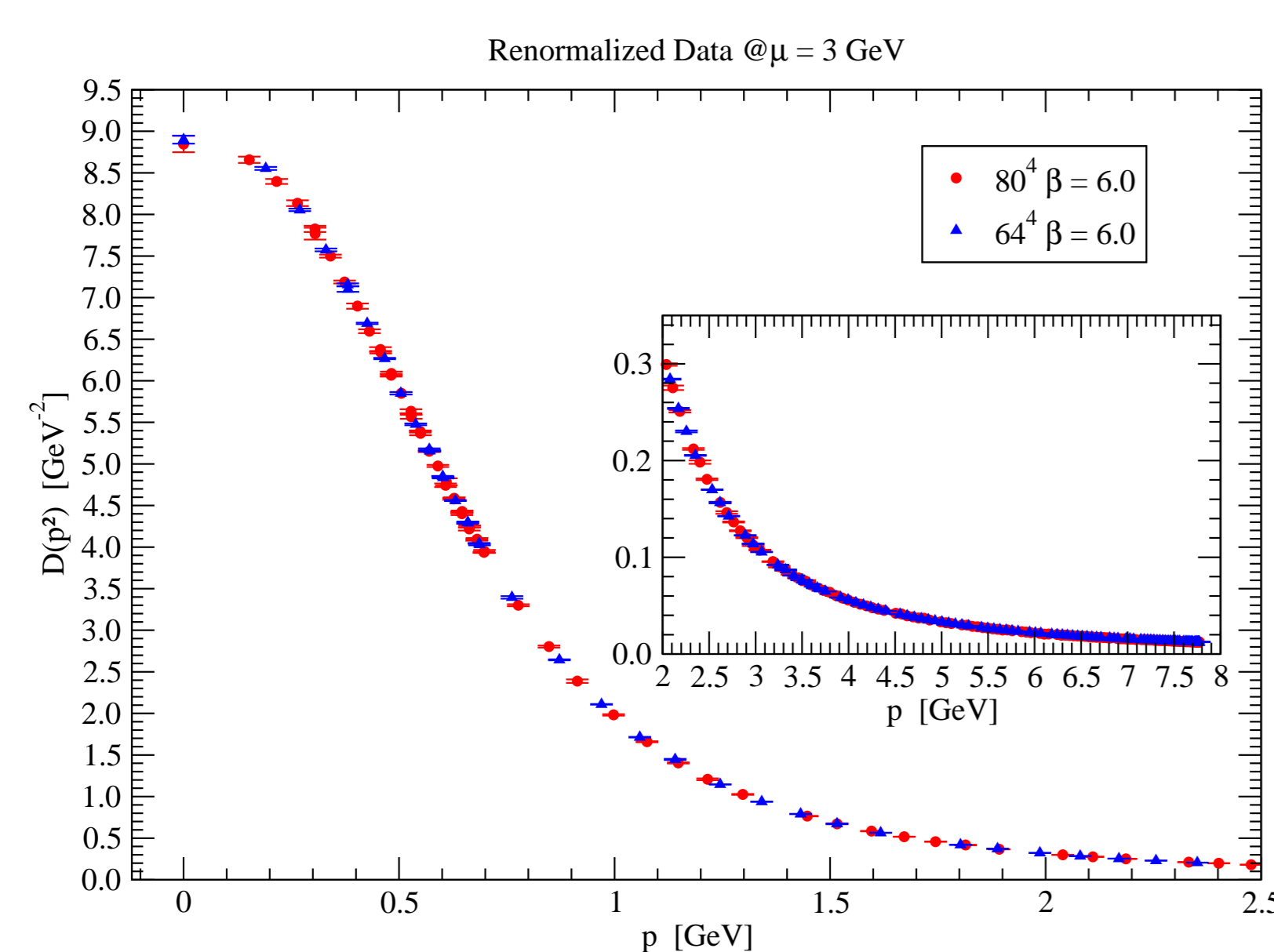
► $p_{min} = 191$ MeV, 153 MeV; $p_{max} = 7.7$ GeV

► renormalization: MOM scheme, scale $\mu = 3$ GeV

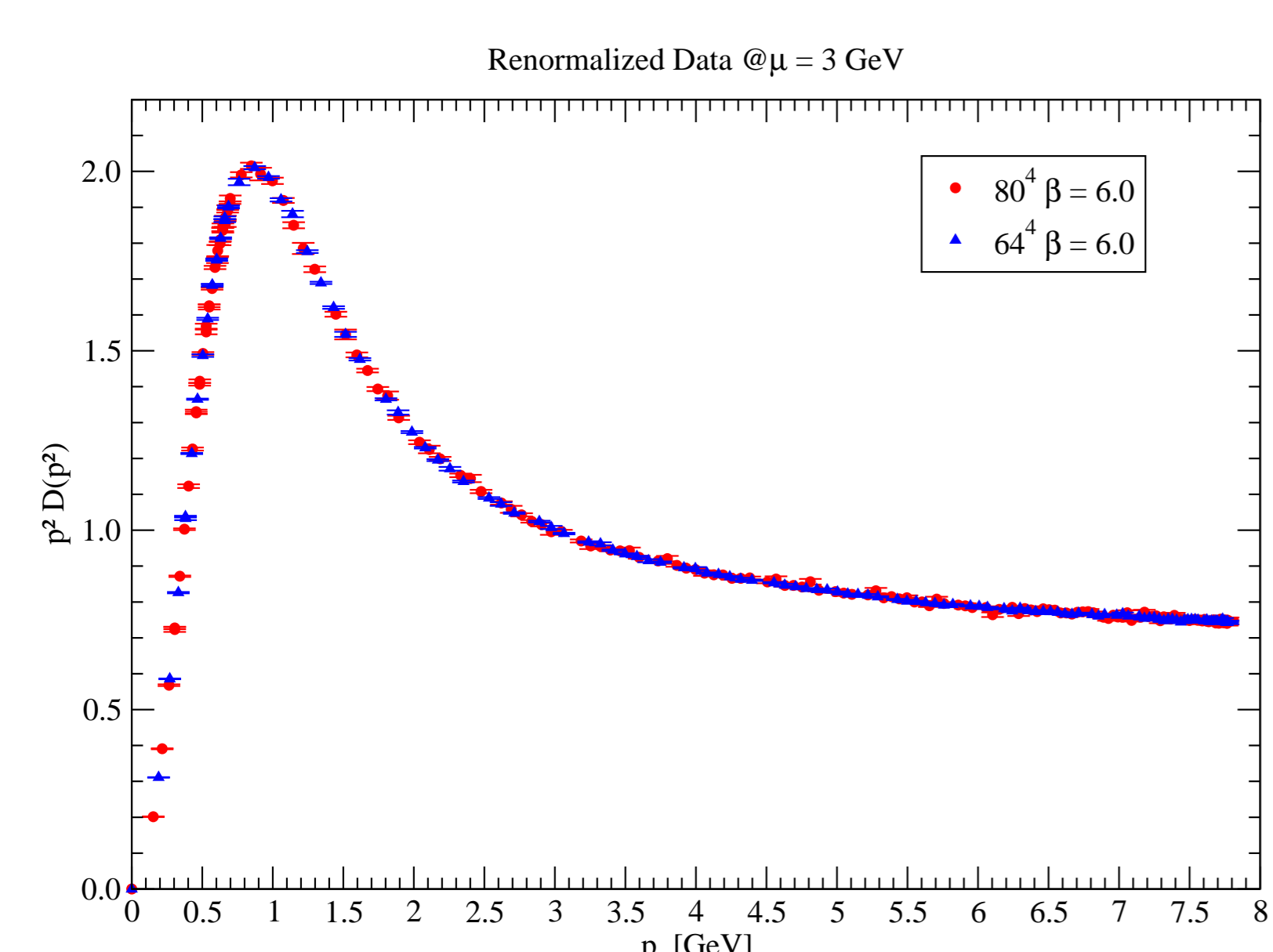
► conical and cylindrical cuts for $p > 0.7$ GeV

► all lattice data for $p < 0.7$ GeV

Gluon Propagator



Gluon Dressing Function



Refined Gribov-Zwanziger (RGZ) framework

► implements the functional restriction to the first Gribov horizon

► defines a local renormalizable quantum field theory

Dudal, Sorella, Vandersickel, Verschelde, PRD77(2008)071501
 Dudal, Gracey, Sorella, Vandersickel, Verschelde, PRD78(2008)065047
 Capri, Fiorentini, Pereira, Sorella, PRD96(2017)054022
 Gracey, PRD82(2010)085032
 Dudal, Sorella, Vandersickel, Verschelde, PRD84(2011)065039

► RGZ takes into account some $d = 2$ condensates

► we study the compatibility of tree level predictions for the gluon propagator with lattice data

► (Very) Refined Gribov-Zwanziger *ansätze*

$$D_{RGZ}(p^2) = Z \frac{p^2 + M_1^2}{p^4 + M_2^2 p^2 + M_3^4}$$

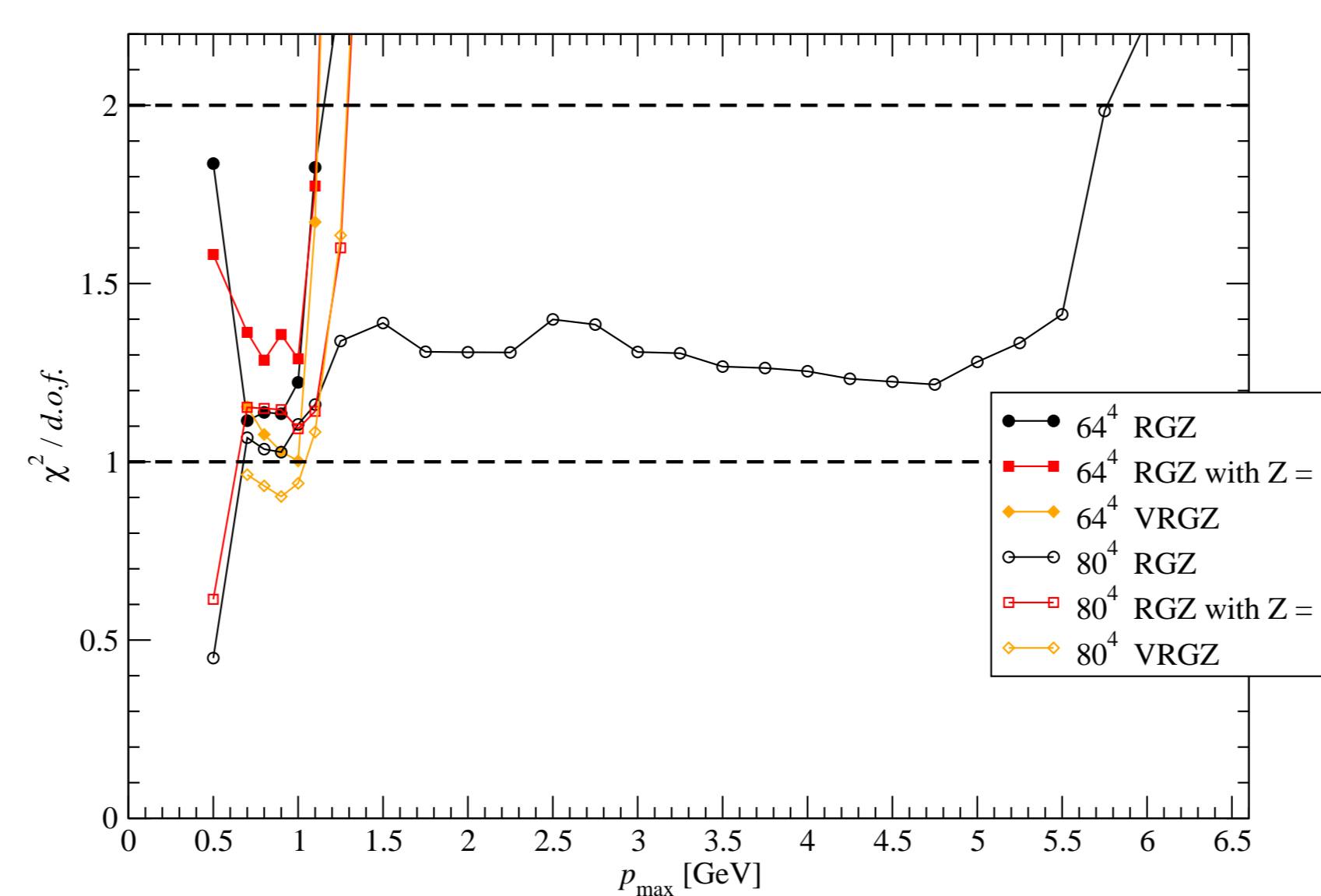
$$D_{VRGZ}(p^2) = \frac{p^4 + M_2^2 p^2 + M_1^4}{p^6 + M_5^2 p^4 + M_4^4 p^2 + M_3^6}$$

► M_i are fitting parameters with dimensions of mass

► for $D_{VRGZ}(p^2)$ we do not include an overall normalization factor Z

► fits to the lattice data from $p = 0$ up to p_{max}

Quality of the fits



► 80^4 lattice: $p_{max} \sim 6$ GeV (smaller statistics)

► 64^4 lattice: $p_{max} \sim 1$ GeV (larger statistics)

Fitting $D_{RGZ}(p^2)$

$$\nu = \chi^2 / \text{d.o.f.}$$

► $L = 64$

p_{max}	ν	Z	M_1^2	M_2^2	M_3^4
1.10	1.83	0.959(62)	2.67(32)	0.511(26)	0.285(15)
1.19	1.77	1	2.478(27)	0.496(9)	0.275(3)

► $L = 80$

p_{max}	ν	Z	M_1^2	M_2^2	M_3^4
1.10	1.16	0.957(66)	2.73(34)	0.527(29)	0.290(16)
3.00	1.31	0.730(4)	4.16(7)	0.592(10)	0.335(3)
1.25	1.60	1	2.454(35)	0.489(11)	0.273(3)

► lattice data compatible with RGZ up to $p \sim 1$ GeV

► agreement between ensembles for similar p_{max}

Fitting $D_{VRGZ}(p^2)$

► no significant difference between the RGZ and the VRGZ scenarios

► M_1^4, M_3^6 compatible with zero: VRGZ reduces to RGZ

$$D(p^2) = \frac{p^2 + M_2^2}{p^4 + M_5^2 p^2 + M_4^4}$$

$$M_2^2|_{VRGZ} \sim M_2^2|_{RGZ}, \quad M_5^2|_{VRGZ} \sim M_2^2|_{RGZ}, \quad M_4^4|_{VRGZ} \sim M_3^4|_{RGZ}$$

► IR lattice data ($p \lesssim 1$ GeV) well described by tree level Refined Gribov-Zwanziger prediction

► if the VRGZ reduces to the RGZ, the equality between the condensates $\langle \bar{\varphi} \varphi \rangle = \langle \varphi \varphi \rangle$ should hold.

Global Fits: from Infrared to Ultraviolet

► for large p^2 one expects to recover the usual perturbative behaviour

$$D(p^2) \propto \frac{1}{p^2} \left[\ln\left(\frac{p^2}{\Lambda_{QCD}^2}\right) \right]^{\gamma_{gl}}$$

► $\gamma_{gl} = -\frac{13}{22} \rightarrow$ 1-loop gluon anomalous dimension

► Interpolation between

► RGZ for low p

► 1-loop RG-improved expression for high p

$$D(p^2) = Z \frac{p^2 + M_1^2}{p^4 + M_2^2 p^2 + M_3^4} \left[\omega \ln\left(\frac{p^2 + m_g^2(p^2)}{\Lambda_{QCD}^2}\right) + 1 \right]^{\gamma_{gl}}$$

► $\omega = 11N \alpha_s(\mu) / (12\pi)$

► $\Lambda_{QCD} = 0.425$ GeV, $\alpha_s(3 \text{ GeV}) = 0.3837$

Dudal, Felix, Guimarães, Sorella, PRD96(2017)074036
 Aguilar, Binosi, Papavassiliou, JHEP07(2010)002

► regularisation of the leading log — introduction of a “log-regularisation mass” $m_g^2(p^2)$ which should become negligible at high momentum

► several regularizing mass definitions:

$$m_g^2(p^2) = \frac{M_3^4 + (M_2^2 - M_1^2)p^2}{M_1^2 + p^2} \quad (\text{RGZ inspired})$$

$$m_g^2(p^2) = m_0^2 \quad (\text{Constant})$$

$$m_g^2(p^2) = \frac{m_0^4}{p^2 + \lambda^2} \quad (\text{Cornwall})$$

$$m_g^2(p^2) = \lambda_0^2 + \frac{m_0^4}{p^2 + \lambda^2} \quad (\text{Corrected Cornwall})$$

► best 64^4 fit: corrected Cornwall regularisation mass ($\chi^2 / \text{d.o.f.} = 1.11$)

$$Z = 1.36992 \pm 0.00072$$

$$M_1^2 = 2.333 \pm 0.042 \text{ GeV}^2$$

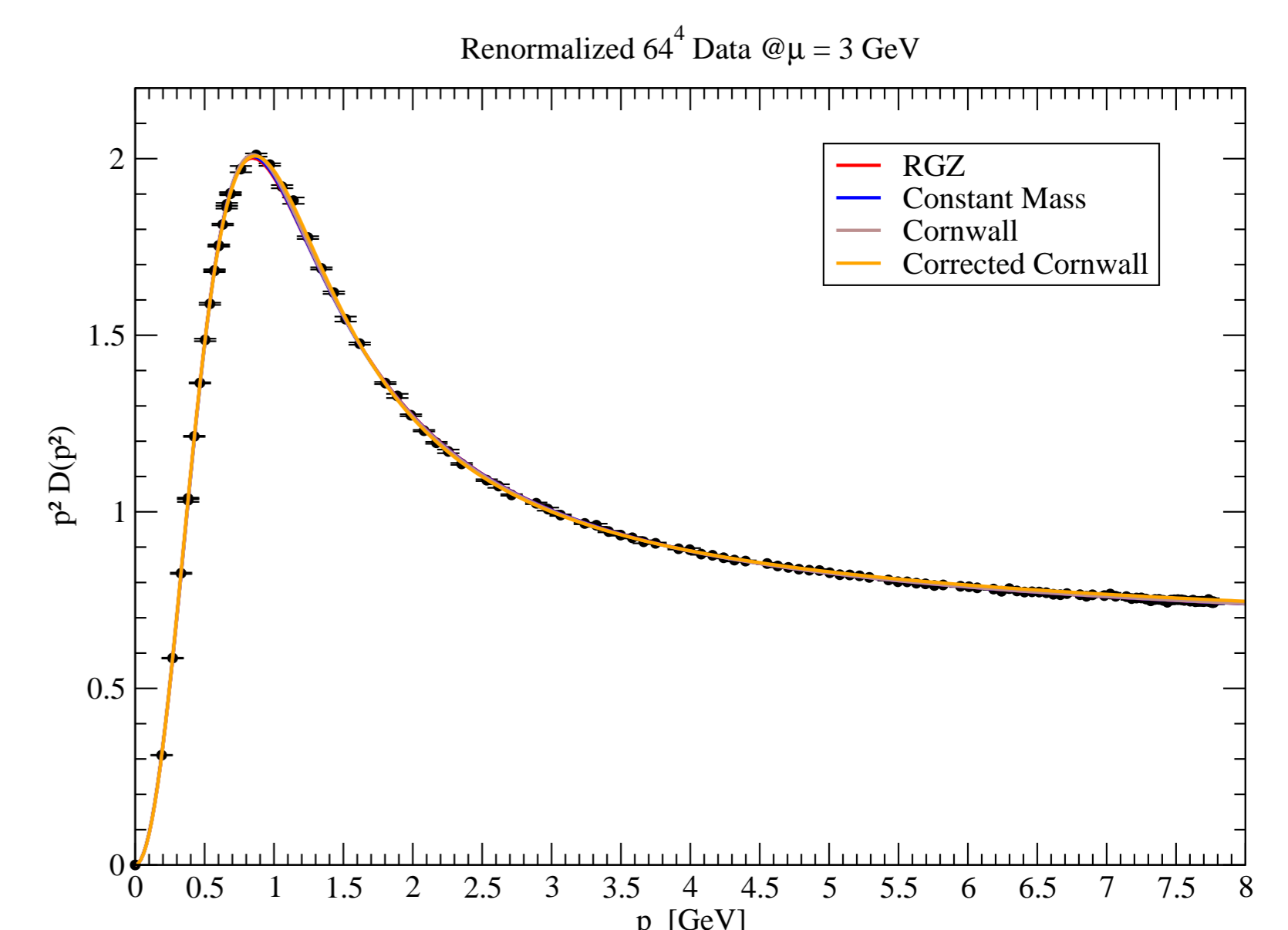
$$M_2^2 = 0.514 \pm 0.024 \text{ GeV}^2$$

$$M_3^4 = 0.2123 \pm 0.0032 \text{ GeV}^4$$

$$m_0^2 = 1.33 \pm 0.13 \text{ GeV}^4$$

$$\lambda^2 = 0.100 \pm 0.035 \text{ GeV}^2$$

$$\lambda_0^2 = -0.954 \pm 0.070 \text{ GeV}^2$$



For more information ...

... see arXiv:1803.02281 [hep-lat]

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