

## TOWARDS THE P-WAVE $N\pi$ SCATTERING AMPLITUDE IN THE $\Delta(1232)$

Interpolating fields and spectra

July 27, 2018 | Giorgio Silvi | Forschungszentrum Jülich

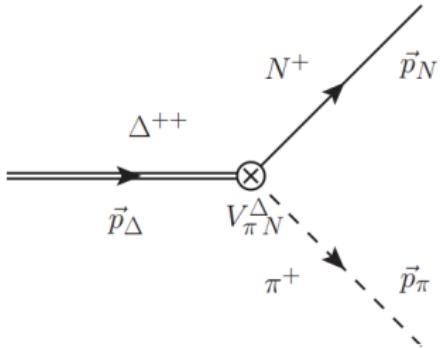
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# THE DELTA(1232)

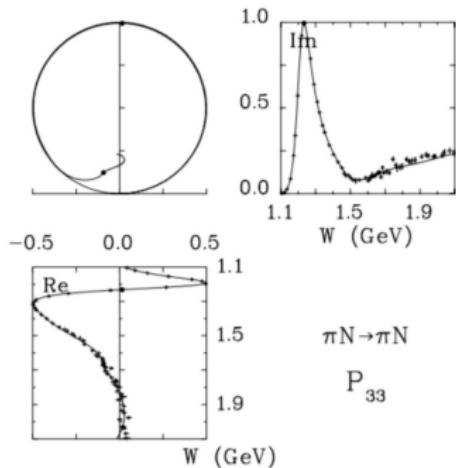
the first baryon resonance

- In nature:  $\Delta^- \Delta^0 \Delta^+ \Delta^{++}$  (u,d quarks) - mass  $\sim 1232$  MeV
  - On the lattice: isospin symmetry
- The unstable  $\Delta(1232)$  decay predominantly to stable  $N\pi$
- Study: Pion-Nucleon scattering  $J = 3/2$ ,  $I = 3/2$ ,  $I_3 = +3/2$ 
  - Orbital angular momentum:  $L = 1$



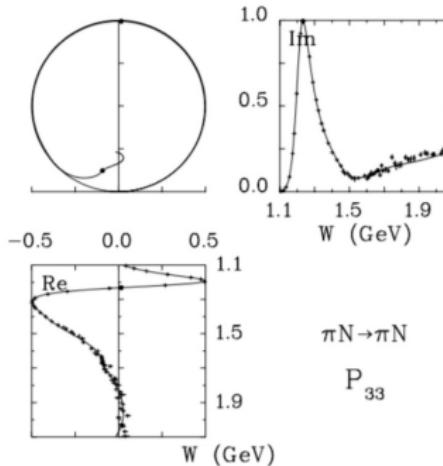
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$N\pi(\rightarrow \Delta(1232)) \rightarrow N\pi$   
completely elastic...



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Particle	$J^P$	$\Gamma_{N\pi} [MeV]$
$\Delta(1232)$	$3/2^+$	$112.4(5)$
$\Delta(1600)$	$3/2^+$	$18(4)$
$\Delta(1620)$	$1/2^-$	$37(2)$
$\Delta(1700)$	$3/2^-$	$36(2)$
...	...	...

# LÜSCHER METHOD

## Lüscher quantization condition for baryons

$$\det[M_{Jlm, J'l'm'} - \delta_{JJ'}\delta_{ll'}\delta_{mm'} \cot \delta_{JI}] = 0 \quad [\text{Gockeler et al. (2012)}]$$

This relation connects the energy  $E$  from a lattice simulation in a finite volume to the unknown phases  $\delta_{JI}$  in the infinite volume via the calculable non-diagonal matrix  $M_{Jlm, J'l'm'}$  (depends on symmetry)

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## Simplify!

With a proper transformation, the matrix  $M_{Jlm, J'l'm'}$  can be block diagonalized in the basis of the irreps  $\Lambda$  of the lattice.

# MOVING FRAMES

## Problem

Due to quantized momenta  $p = 2\pi n/L$  we have energy levels spaced from each other. Chances of hitting the energy region of interest are low.

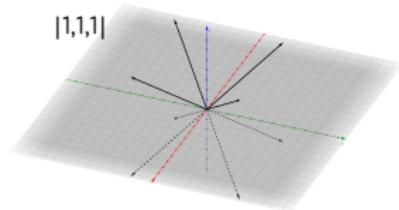
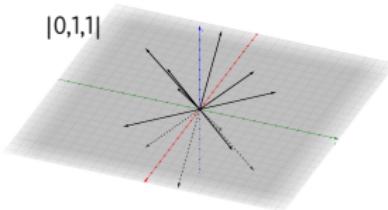
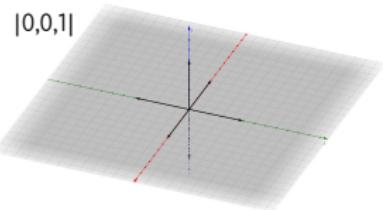
# MOVING FRAMES

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## Solution: Moving frames!

The Lorentz boost contracts the box giving a different effective value of the size L. Allow access to phase shift at different energies!



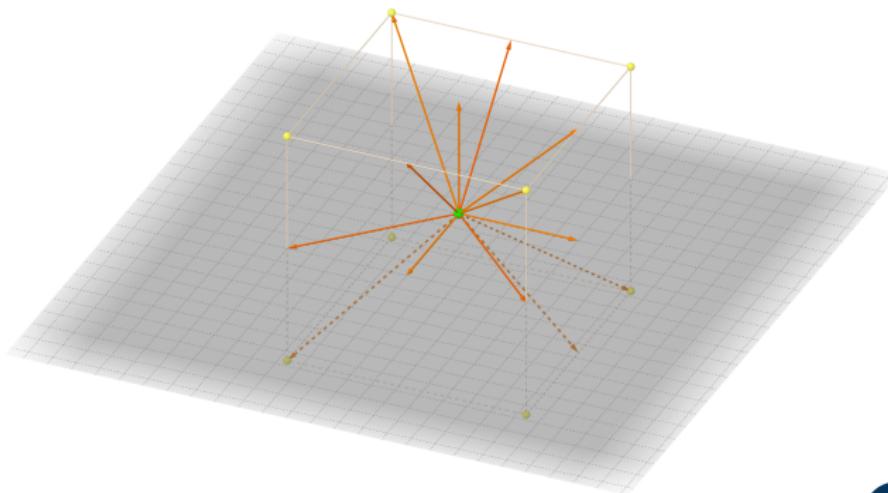
momentum directions

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- In the continuum, states are classified according to angular momentum  $J$  and parity  $P$ 
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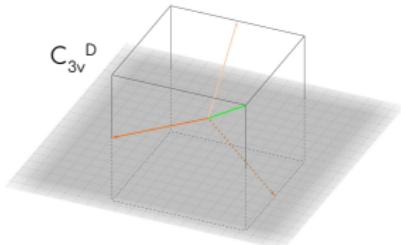
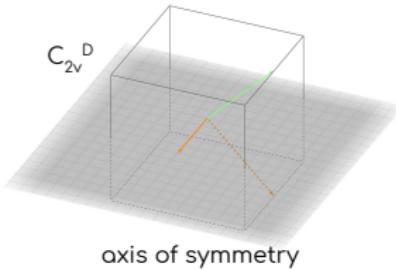
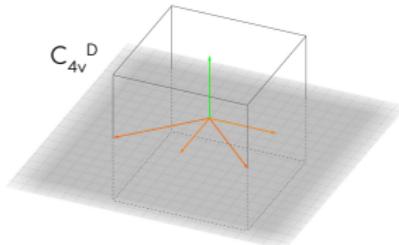
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  - For half-integer  $J$  we need the double cover  $O_h^D$  (96 elements) which include the negative identity ( $2\pi$  rotation)
- Each of the infinite irreps  $J^P$  in the continuum get mapped to one of the finite irreps  $\Lambda$  of the group  $O_h^D$  on the lattice.

# GROUND PLAN

## Frames, Groups & Irreps $\Lambda$ (with ang. mom. content)

$P_{ref}[N_{dir}]$	Group	$N_{elem}$	$\Lambda(J) : \pi(0^-)$	$\Lambda(J) : N(\frac{1}{2}^+)$	$\Lambda(J) : \Delta(\frac{3}{2}^+)$
(0, 0, 0) [1]	$O_h^D$	96	$A_{1u}(0, 4, \dots)$	$G_{1g}(\frac{1}{2}, \frac{7}{2}, \dots) \oplus G_{1u}(\frac{1}{2}, \frac{7}{2}, \dots)$	$H_g(\frac{3}{2}, \frac{5}{2}, \dots) \oplus H_u(\frac{3}{2}, \frac{5}{2}, \dots)$
(0, 0, 1) [6]	$C_{4v}^D$	16	$A_2(0, 1, \dots)$	$G_1(\frac{1}{2}, \frac{3}{2}, \dots)$	$G_1(\frac{1}{2}, \frac{3}{2}, \dots) \oplus G_2(\frac{3}{2}, \frac{5}{2}, \dots)$
(0, 1, 1) [12]	$C_{2v}^D$	8	$A_2(0, 1, \dots)$	$G(\frac{1}{2}, \frac{3}{2}, \dots)$	$G(\frac{1}{2}, \frac{3}{2}, \dots)$
(1, 1, 1) [8]	$C_{3v}^D$	12	$A_2(0, 1, \dots)$	$G(\frac{1}{2}, \frac{3}{2}, \dots)$	$G(\frac{1}{2}, \frac{3}{2}, \dots) \oplus F_1(\frac{3}{2}, \frac{5}{2}, \dots) \oplus F_2(\frac{3}{2}, \frac{5}{2}, \dots)$



# SINGLE HADRON OPERATORS

Delta interpolators:

$$\Delta_{i\mu}^{(1)} = \epsilon_{abc} u_\mu^a (u^{bT} C \gamma_i u^c) \quad (1)$$

$$\Delta_{i\mu}^{(2)} = \epsilon_{abc} u_\mu^a (u^{bT} C \gamma_i \gamma_0 u^c) \quad (2)$$

Nucleon interpolators:

$$\mathcal{N}_\mu^{(1)} = \epsilon_{abc} u_\mu^a (u^{bT} C \gamma_5 d^c) \quad (3)$$

$$\mathcal{N}_\mu^{(2)} = \epsilon_{abc} u_\mu^a (u^{bT} C \gamma_0 \gamma_5 d^c) \quad (4)$$

Pion interpolator:

$$\pi = \bar{d} \gamma_5 u \quad (5)$$

# PROJECTION METHOD

how it works...

$$O^{G^D, \Lambda, r, m}(p) = \frac{d_\Lambda}{g_{G^D}} \sum_{\tilde{R} \in G^D} \Gamma_{r,r}^\Lambda(\tilde{R}) U_{\tilde{R}} \phi(p) U_{\tilde{R}}^{-1} \quad [\text{C. Morningstar et al. (2013)}]$$

- to get an operators  $O^{G^D, \Lambda, r, m}(p)$  for a specific:

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  - 4 proper transformation matrices  $U_{\tilde{R}}$

# OCCURENCES OF IRREPS

It is possible to find the occurrence ( $\sim$ multiplicity)  $m$  of the irrep  $\Gamma^\Lambda$  in the transformation matrices  $U_{\tilde{R}}$  using the character  $\chi$ : [Moore,Fleming (2006)]

$$m = \frac{1}{g_G^D} \sum_{\tilde{R} \in G^D} \chi^{\Gamma^\Lambda}(\tilde{R}) \chi^U(\tilde{R})$$

NUCLEON

$O_h^D$

[0,0,0]

$C_{4v}^D$

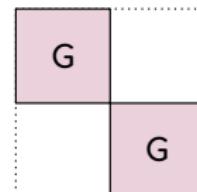
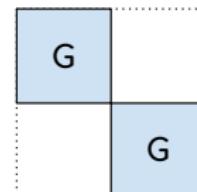
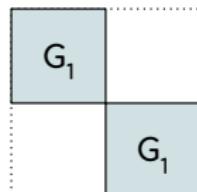
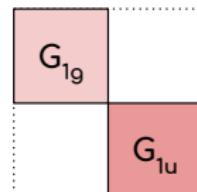
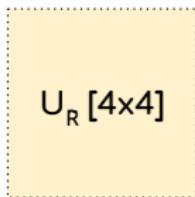
[0,0,1]

$C_{2v}^D$

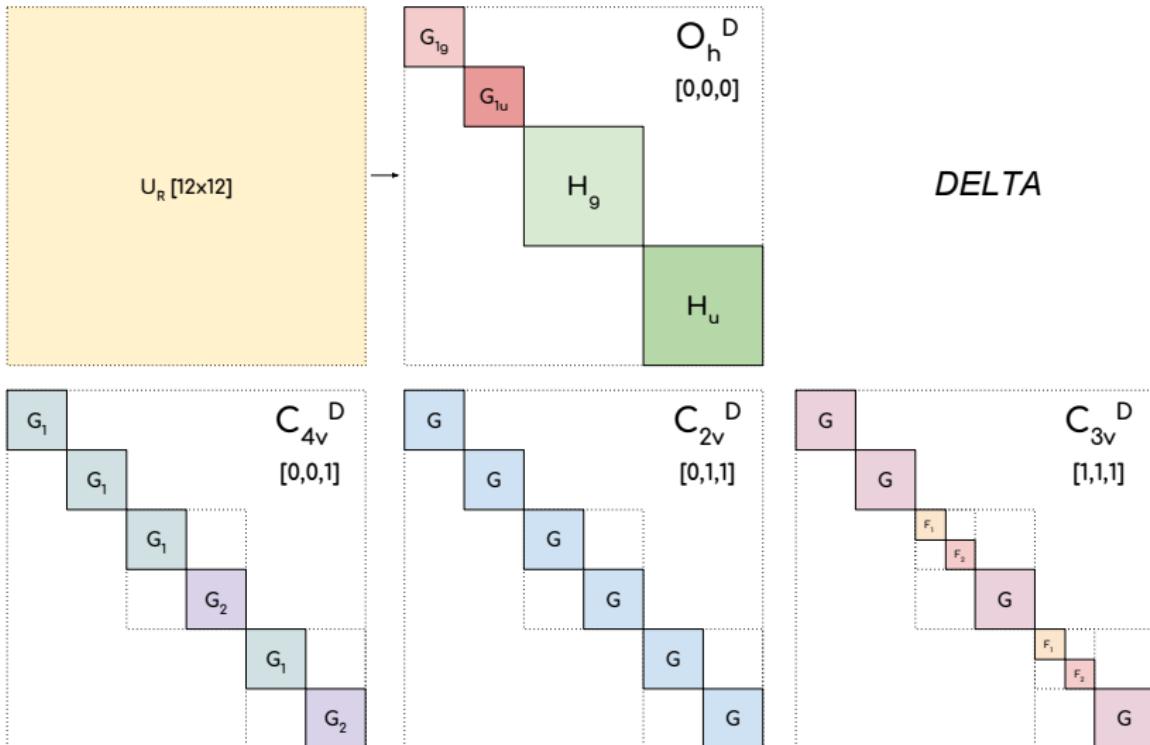
[0,1,1]

$C_{3v}^D$

[1,1,1]



# OCCURENCES OF IRREPS



# EXAMPLE OF PROJECTION: NUCLEON

C4v - G1 - first instance (row 1 , row 2)

$$\left( \begin{array}{cc} \{Nc(0, 0, 1)(3)\} & \{Nc(0, 0, 1)(4)\} \\ \{Nc(0, 0, -1)(3)\} & \{Nc(0, 0, -1)(4)\} \\ \{Nc(1, 0, 0)(4) - Nc(1, 0, 0)(2)\} & \{Nc(1, 0, 0)(2) + Nc(1, 0, 0)(4)\} \\ \{Nc(-1, 0, 0)(4) - Nc(-1, 0, 0)(2)\} & \{Nc(-1, 0, 0)(2) + Nc(-1, 0, 0)(4)\} \\ \{Nc(0, -1, 0)(4) - Nc(0, -1, 0)(2)\} & \{Nc(0, -1, 0)(2) + Nc(0, -1, 0)(4)\} \\ \{Nc(0, 1, 0)(4) - Nc(0, 1, 0)(2)\} & \{Nc(0, 1, 0)(2) + Nc(0, 1, 0)(4)\} \end{array} \right)$$

C4v - G1 - second instance (row 1 , row 2)

$$\left( \begin{array}{cc} \{Nc(0, 0, 1)(1)\} & \{Nc(0, 0, 1)(2)\} \\ \{Nc(0, 0, -1)(1)\} & \{Nc(0, 0, -1)(2)\} \\ \{Nc(1, 0, 0)(1) + Nc(1, 0, 0)(3)\} & \{Nc(1, 0, 0)(3) - Nc(1, 0, 0)(1)\} \\ \{Nc(-1, 0, 0)(1) + Nc(-1, 0, 0)(3)\} & \{Nc(-1, 0, 0)(3) - Nc(-1, 0, 0)(1)\} \\ \{Nc(0, -1, 0)(1) + Nc(0, -1, 0)(3)\} & \{Nc(0, -1, 0)(3) - Nc(0, -1, 0)(1)\} \\ \{Nc(0, 1, 0)(1) + Nc(0, 1, 0)(3)\} & \{Nc(0, 1, 0)(3) - Nc(0, 1, 0)(1)\} \end{array} \right)$$

# EXAMPLE OF PROJECTION: DELTA

## C4v - G1 - row 1 - instance 1

$$\frac{\frac{\Delta(0, 0, 1)(3, 1)}{\Delta(1, 0, 0)(1, 1)}}{\sqrt{2}} - \frac{\frac{\Delta(1, 0, 0)(1, 3)}{\Delta(1, 0, 0)(1, 1)}}{\sqrt{2}}$$
$$\frac{1}{2}\Delta(0, -1, 0)(1, 1) - \frac{1}{2}\Delta(0, -1, 0)(1, 3) - \frac{1}{2}\Delta(0, -1, 0)(3, 2) + \frac{1}{2}\Delta(0, -1, 0)(3, 4)$$

## C4v - G1 - row 1 - instance 2

$$\frac{\frac{\Delta(0, 0, 1)(1, 2)}{\Delta(1, 0, 0)(2, 1)}}{\sqrt{2}} + \frac{i\Delta(0, 0, 1)(2, 2)}{\sqrt{2}}$$
$$\frac{1}{2}\Delta(1, 0, 0)(2, 1) - \frac{1}{2}\Delta(1, 0, 0)(2, 3) - \frac{1}{2}i\Delta(1, 0, 0)(3, 2) + \frac{1}{2}i\Delta(1, 0, 0)(3, 4)$$
$$\frac{\frac{\Delta(0, -1, 0)(2, 1)}{\Delta(0, -1, 0)(2, 2)}}{\sqrt{2}} - \frac{\frac{\Delta(0, -1, 0)(2, 3)}{\Delta(0, -1, 0)(2, 2)}}{\sqrt{2}}$$

## C4v - G1 - row 1 - instance 3

$$\frac{\Delta(0, 0, 1)(3, 3)}{\frac{1}{2}i\Delta(1, 0, 0)(2, 2) + \frac{1}{2}i\Delta(1, 0, 0)(2, 4) + \frac{1}{2}\Delta(1, 0, 0)(3, 1) + \frac{1}{2}\Delta(1, 0, 0)(3, 3)}$$
$$\frac{1}{2}\Delta(0, -1, 0)(1, 2) + \frac{1}{2}\Delta(0, -1, 0)(1, 4) + \frac{1}{2}\Delta(0, -1, 0)(3, 1) + \frac{1}{2}\Delta(0, -1, 0)(3, 3)$$

## C4v - G1 - row 1 - instance 4

$$\frac{\frac{\Delta(0, 0, 1)(1, 4)}{\Delta(1, 0, 0)(1, 2)}}{\sqrt{2}} + \frac{i\Delta(0, 0, 1)(2, 4)}{\sqrt{2}}$$
$$\frac{\frac{\Delta(1, 0, 0)(1, 4)}{\Delta(0, -1, 0)(2, 2)}}{\sqrt{2}} + \frac{\Delta(1, 0, 0)(1, 4)}{\sqrt{2}}$$
$$\frac{\frac{\Delta(0, -1, 0)(2, 4)}{\Delta(0, -1, 0)(2, 2)}}{\sqrt{2}} + \frac{\Delta(0, -1, 0)(2, 4)}{\sqrt{2}}$$

# COMPLETE SET OF OPERATOR: $\Delta, N\pi$

Group	$P_{ref}$ [dir. used]	Irrep[rows]	Op. type	n	Op. per irrep
$O_h^D$	(0, 0, 0)[1]	$H_g[4]$	$\Delta(\gamma_i)$ $\Delta(\gamma_i\gamma_0)$ $N\pi(\gamma_5)$ $N\pi(\gamma_0\gamma_5)$	1 1 8 8	
				18	$\times 4\text{rows}$
		$H_u[4]$		18	$\times 4\text{rows}$
$C_{4v}^D$	(0, 0, 1)[3]	$G_1[2]$	$\Delta$ $N\pi$	$4 \times 2$ $20 \times 2$	
				48	$\times 2\text{rows} \times 3\text{dir}$
		$G_2[2]$	$\Delta$ $N\pi$	$2 \times 2$ $16 \times 2$	
				36	$\times 2\text{rows} \times 3\text{dir}$
$C_{2v}^D$	(0, 1, 1)[6]	$G[2]$	$\Delta$ $N\pi$	$6 \times 2$ $24 \times 2$	
				60	$\times 2\text{rows} \times 6\text{dir}$
					720
$C_{3v}^D$	(1, 1, 1)[4]	$G[2]$	$\Delta$ $N\pi$	$4 \times 2$ $12 \times 2$	
				32	$\times 2\text{rows} \times 4\text{dir}$
		$F_1[1]$	$\Delta$ $N\pi$	$2 \times 2$ $4 \times 2$	
				12	$\times 4\text{dir}$
		$F_2[1]$		12	$\times 4\text{dir}$
Total					1720

# PLANNING

## Ensemble from BMW collaboration

$N_s$	$N_t$	$\beta$	$am_{u,d}$	$am_s$	$c_{sw}$	$a(fm)$	$L(fm)$	$m_\pi(MeV)$	$m_\pi L$
24	48	3.31	-0.0953	-0.040	1.0	0.116	2.8	254	3.6

[ensemble description in S. Durr et al.(2011)]

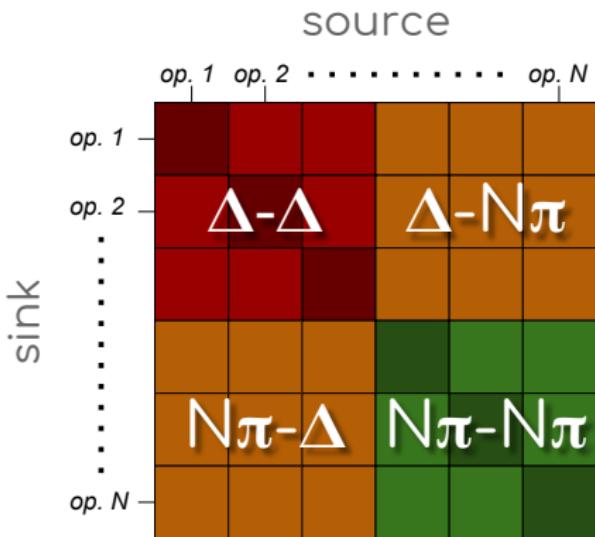
- Lattice action: Wilson-Clover with  $N_f = 2 + 1$  dynamical fermions
- Two-point correlators built from a combination of smeared forward, sequential and stochastic propagators
- This talk: 192 configurations, 16 source location per conf.

## Beginning

- Project the operators on all relevant irreps (Wolfram Mathematica)
- Compute contraction for all diagrams and momentum direction
- Apply the projected operators on correlators

# CORRELATION MATRICES

After projectioning the correlators...



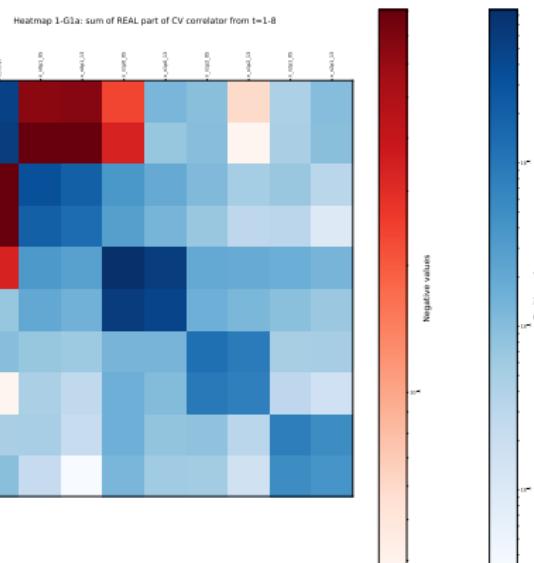
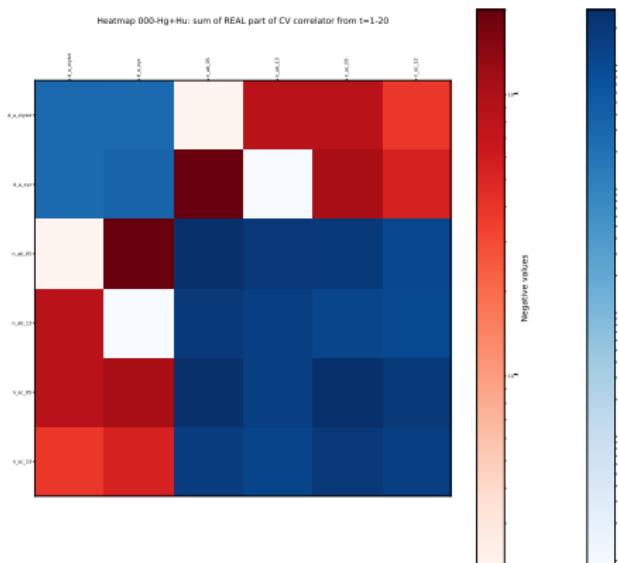
✖ N irreps

(H,G1,G2,G,F1,F2,G)

- Use GEVP to determine the spectra
- Search the best basis of operators

# HEATMAP

## Rest frame / Moving frame

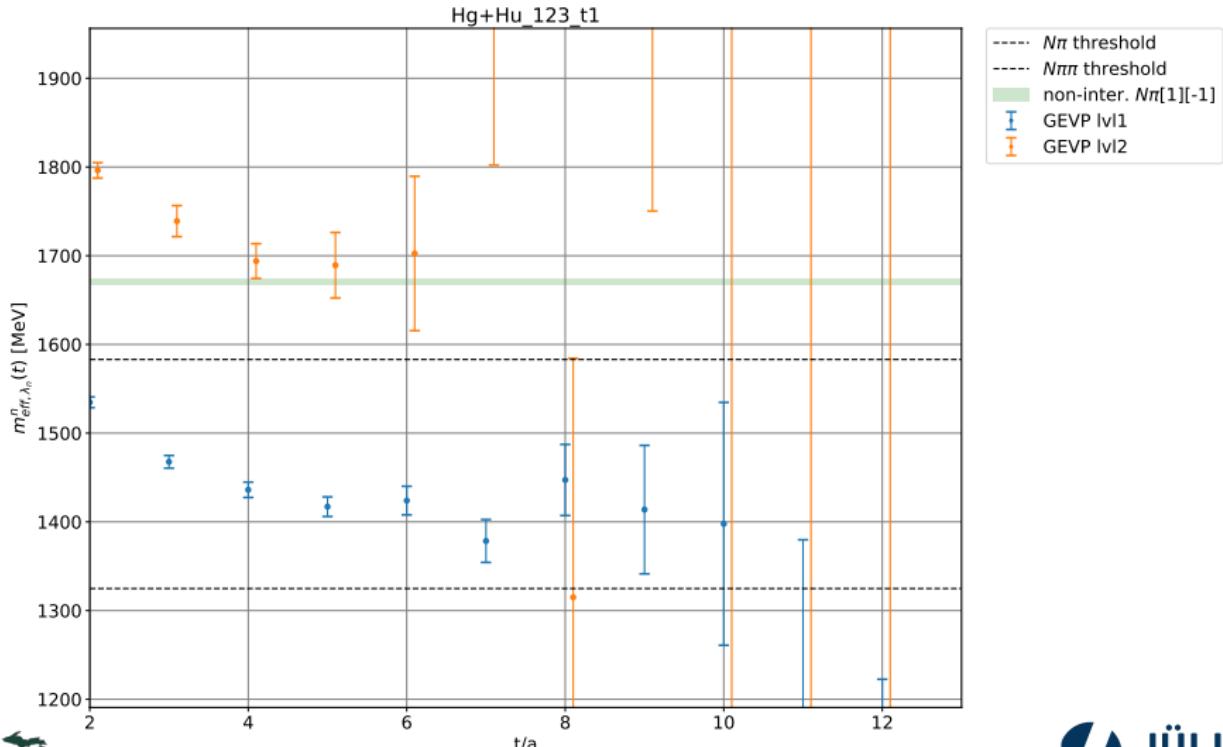


$$O_h^D - irrepHg$$

$$C_{4v}^D - irrepG1$$

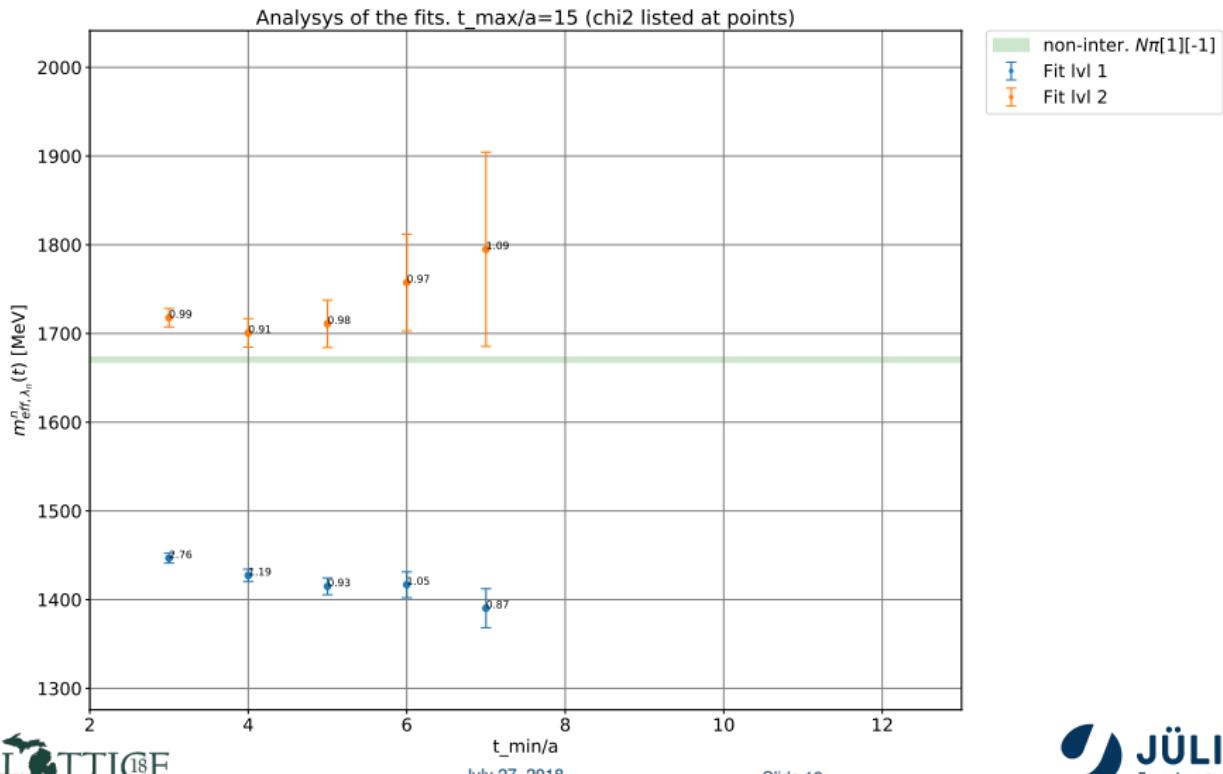
# SPECTRA - REST FRAME

$O_h^D - irrepHg(+Hu)$



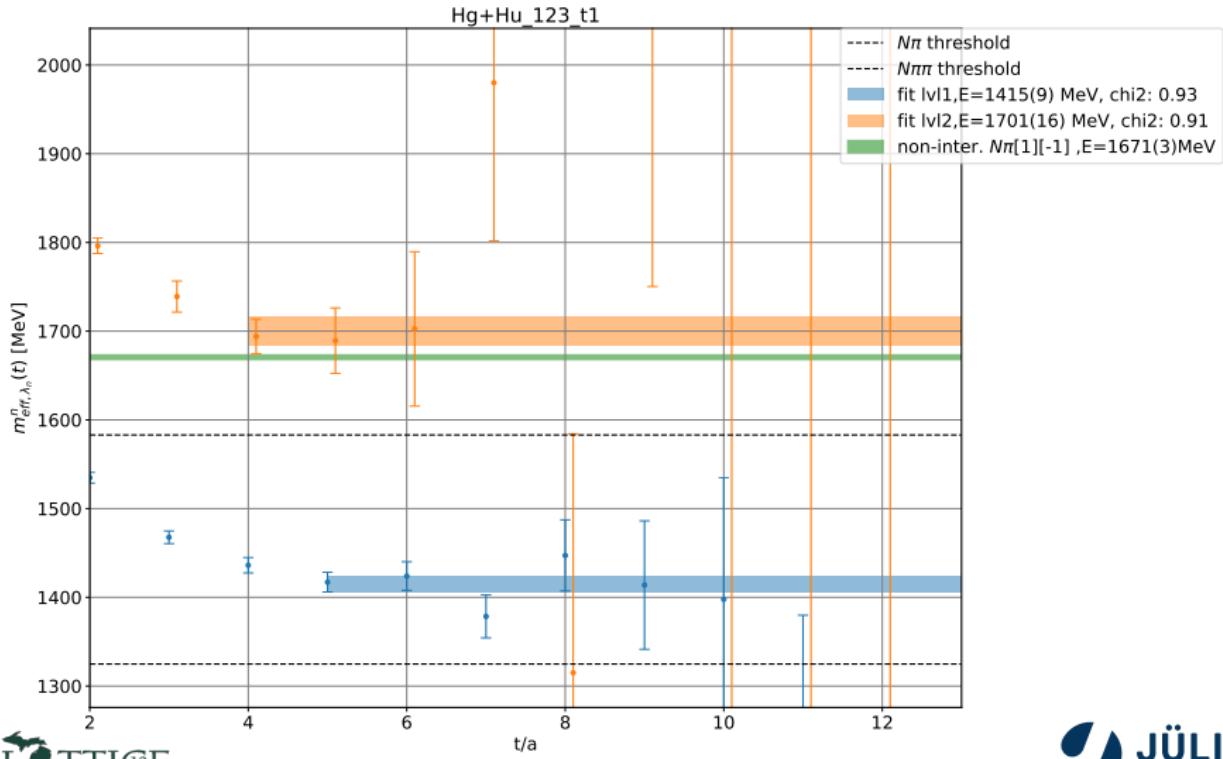
# ANALYSIS OF THE FITS - REST FRAME

$$O_h^D - irrepHg(+Hu)$$



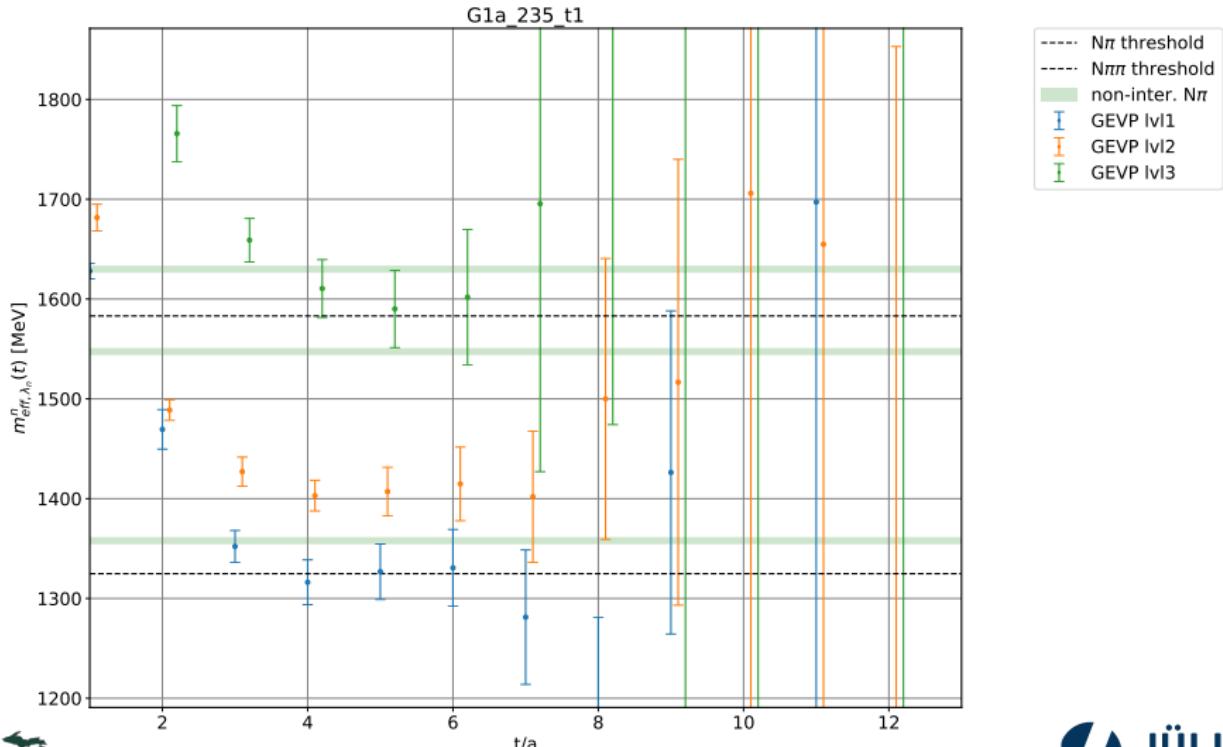
# SPECTRA - REST FRAME

$O_h^D - \text{irrep} : Hg(+Hu)$



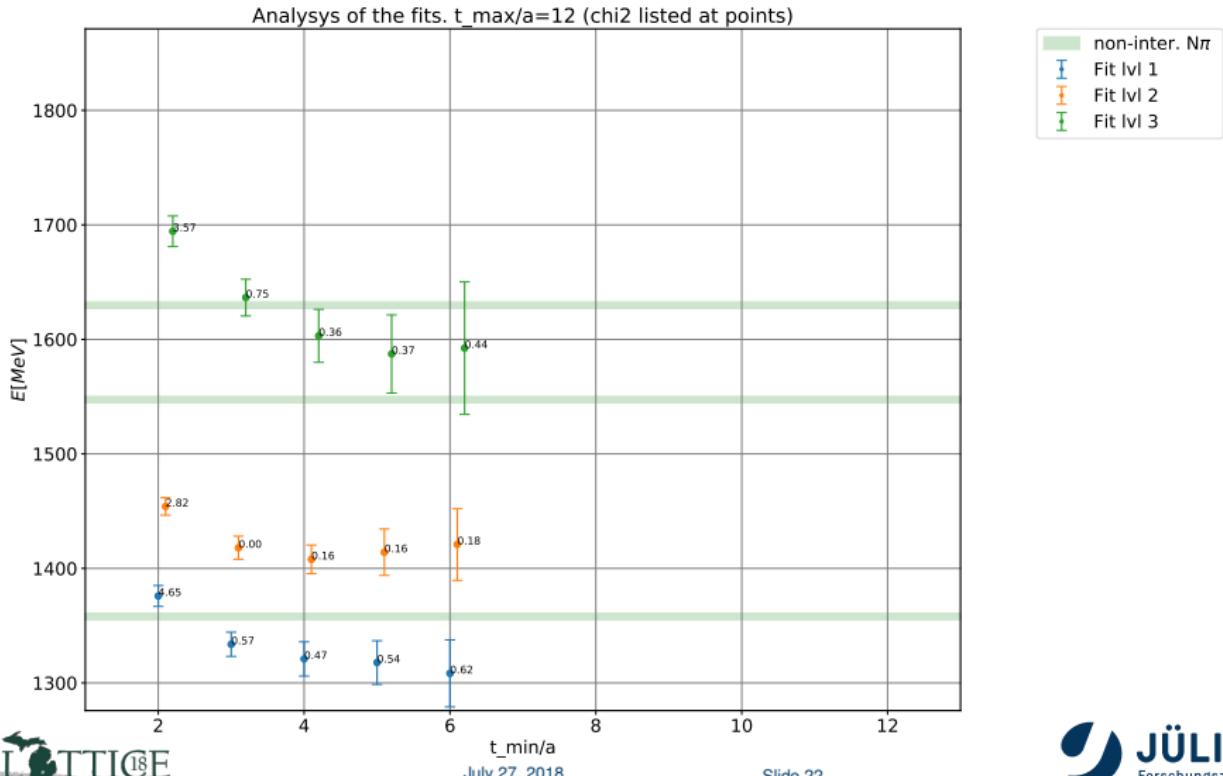
# SPECTRA - MOVING FRAME

$C_{4v}^D$  – irrep : G1



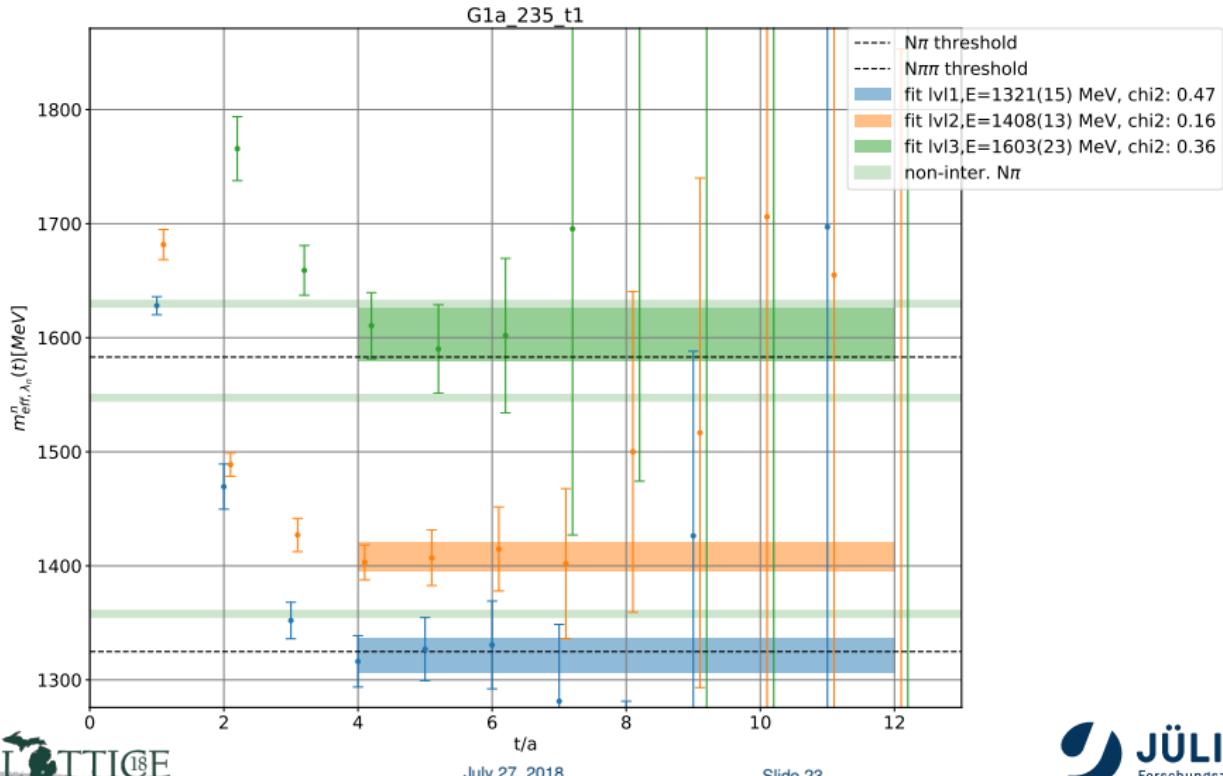
# ANALYSIS OF THE FITS - MOVING FRAME

$C_{4v}^D - irrepG1$



# SPECTRA - MOVING FRAME

$C_{4v}^D - irrepG1$



# CONCLUSION

- Complete projection of operators for  $N$ ,  $\Delta$  and  $N\pi$  in relevant irreps
- Room for improvement of spectra with a thorough search for best basis
- Additional configurations coming(3 $\times$ )
- In the future aim to add a bigger box, more  $m_\pi$  and smaller lattice spacing

# CONCLUSION

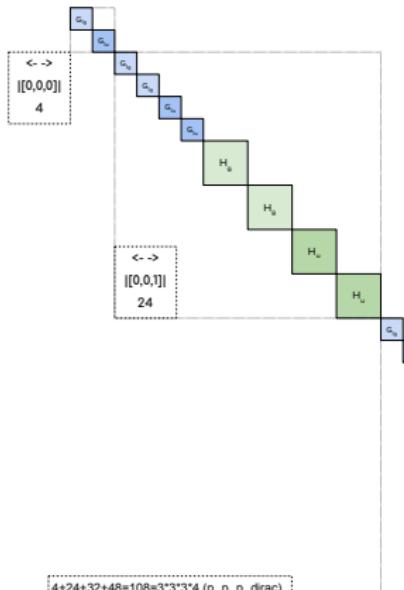
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*Thank you for the attention!*

(second part in the next talk by Srijit)

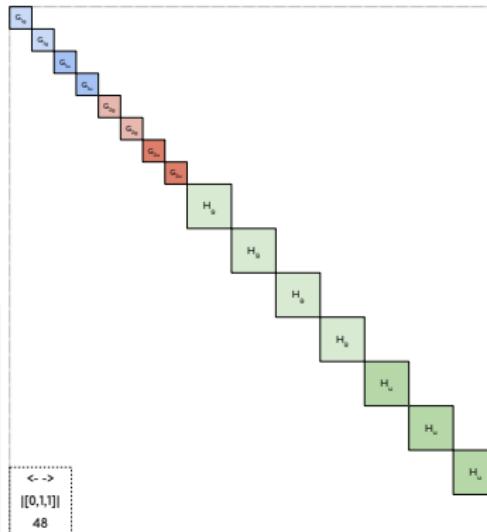


# BACK-UP SLIDE



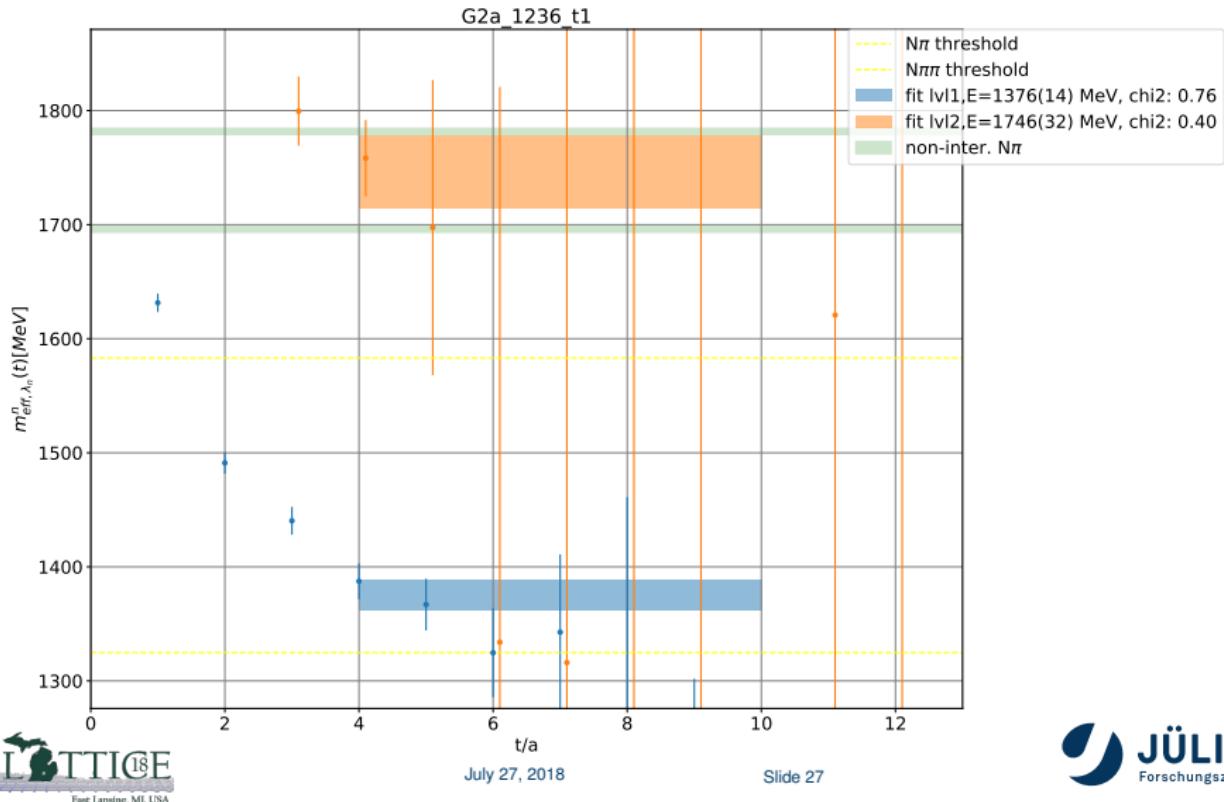
$N-\pi$   
 $O_h^D$   
 $P_{tot}=[0,0,0]$

NUCLEON- PION TENSOR



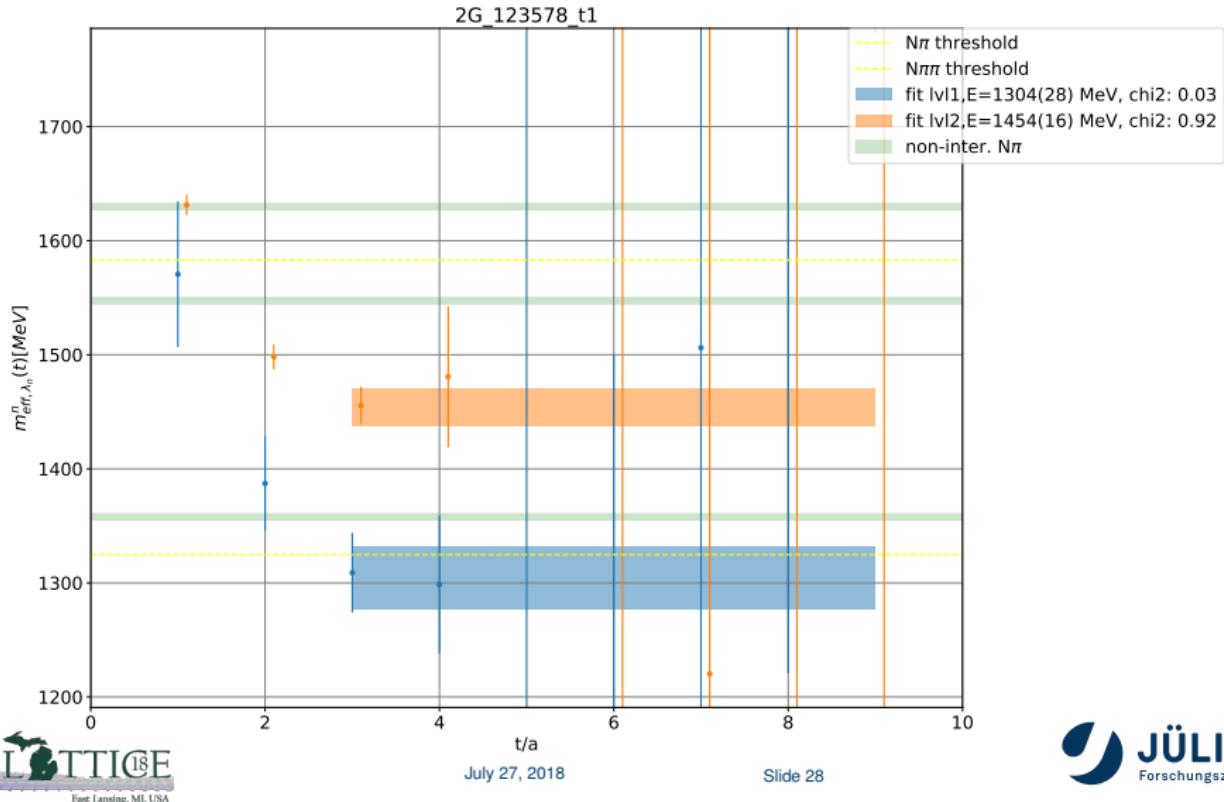
# BACK-UP SLIDE

$C_{4v}^D - irrepG2$



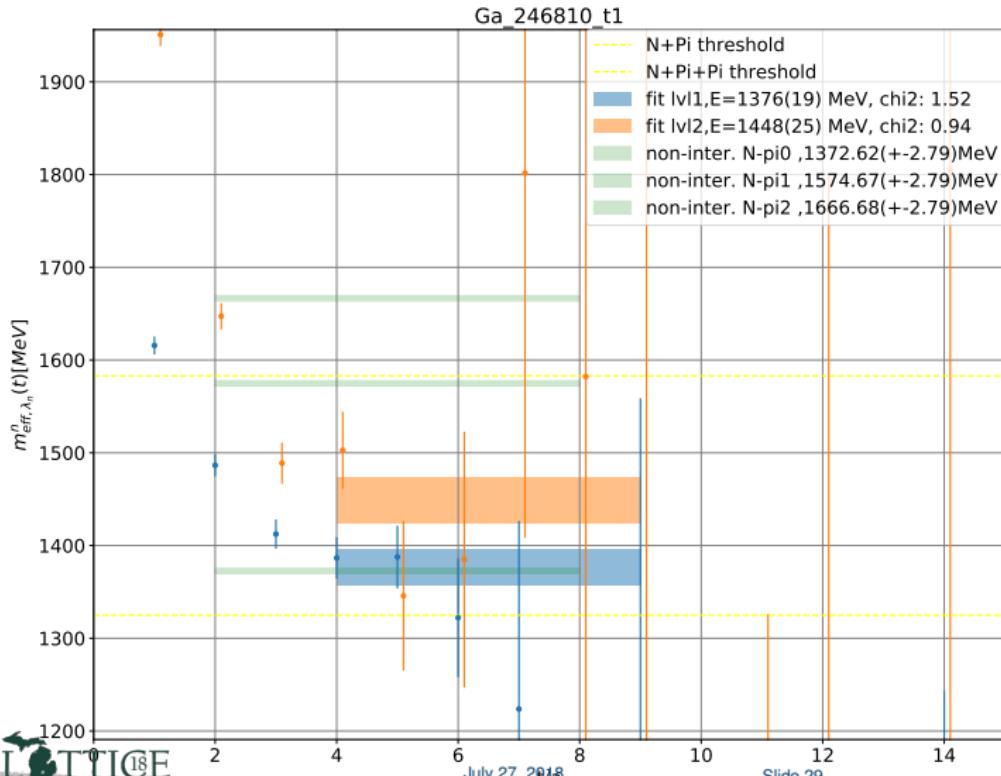
# BACK-UP SLIDE

$C_{2v}^D - irrepG$



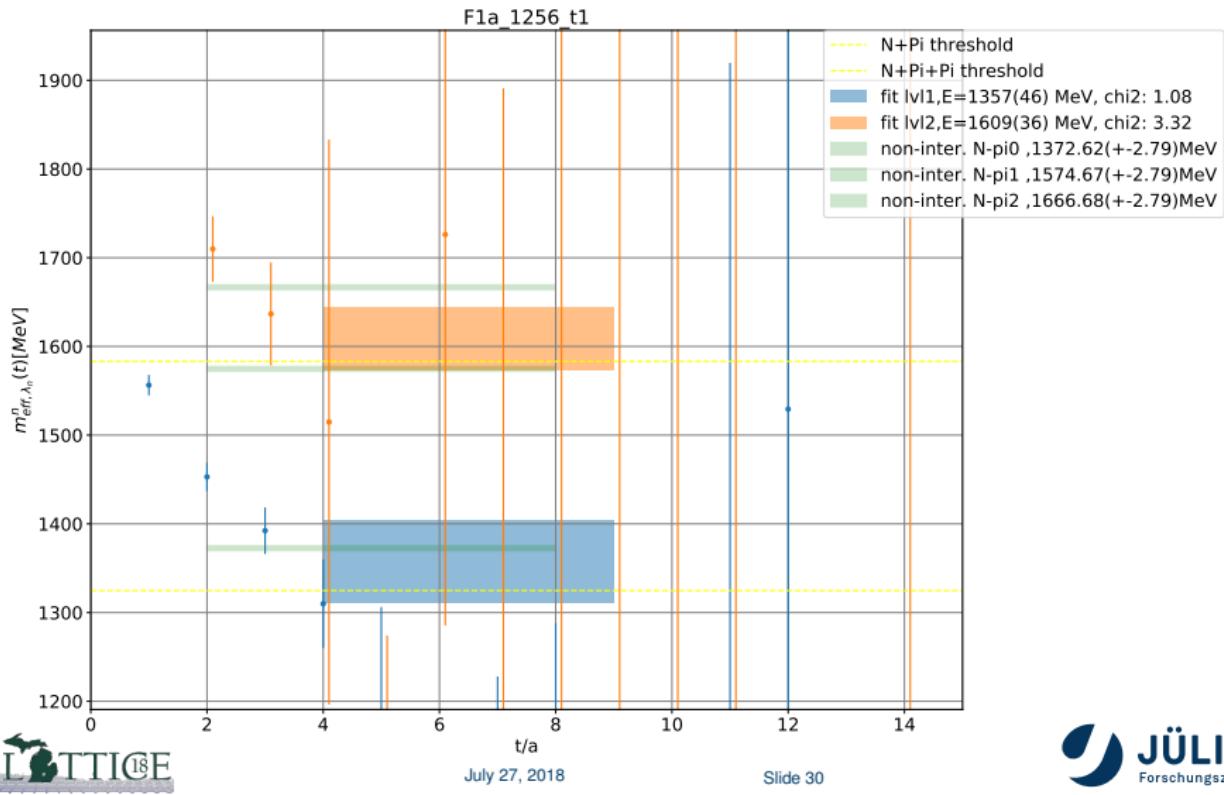
# BACK-UP SLIDE

$C_{3v}^D$  – irrepG preliminary



# BACK-UP SLIDE

$C_{3v}^D - irrepF1$  preliminary



# BACK-UP SLIDE

$C_{3v}^D$  – irrepF2 preliminary

