

# Is $SU(3)$ gauge theory with 13 massless flavors conformal?

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# Context

LatHC studies of many SU(3) gauge theories shows trend of continuing decrease in beta function with increasing  $N_f$  and a steadily lighter scalar

Extensive simulations with a range of lattice spacings, especially for  $N_f = 12$

Consistent picture that  $N_f = 12$  is close to being conformal, paradigm of walking model

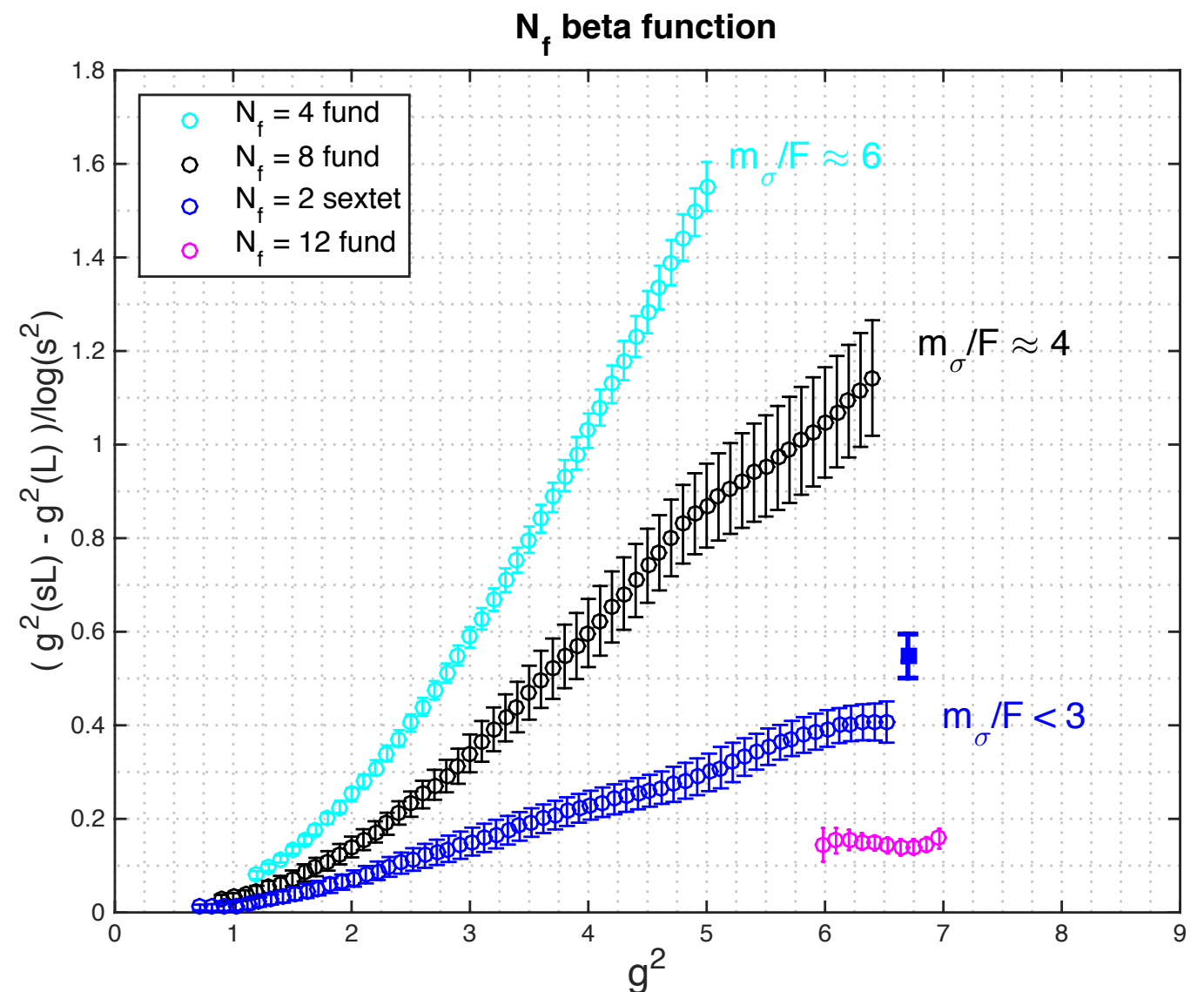
Recent critiques:

1. “staggered fermions are in the wrong universality class” - Hasenfratz, Rebbi, Witzel
2. “ $N_f = 10$  with domain wall fermions has an infrared fixed point” - TW Chiu

Our response:

1. you can't add relevant operators to staggered fermions: spin models with extra fixed points not applicable
2. holding the renormalized gauge coupling fixed in some physical volume controls taste symmetry recovery
3. the Dirac operator eigenvalue spectrum from MC simulations shows the correct structure
4. we have simulated  $N_f = 10$  with staggered fermions with a range of lattice spacings and find no IRFP

One role of an  $N_f = 13$  study: can staggered fermions find an IRFP and a conformal theory?



# Why $N_f = 13$ ?

5-loop  $\overline{\text{MS}}$ -bar calculation (20 year project)  
of gauge coupling beta function

$\text{SU}(3)$   $N_f = 12$  has an infrared fixed point  
(IRFP) at 2, 3, and 4-loop in  $\overline{\text{MS}}$ -bar, which  
disappears at 5-loop

$\text{SU}(3)$   $N_f = 13$  has an intriguing structure at 5-loop  
of an IRFP and a non-trivial UV fixed point

$N_f = 13$  IRFP also present at 2, 3 and 4-loop

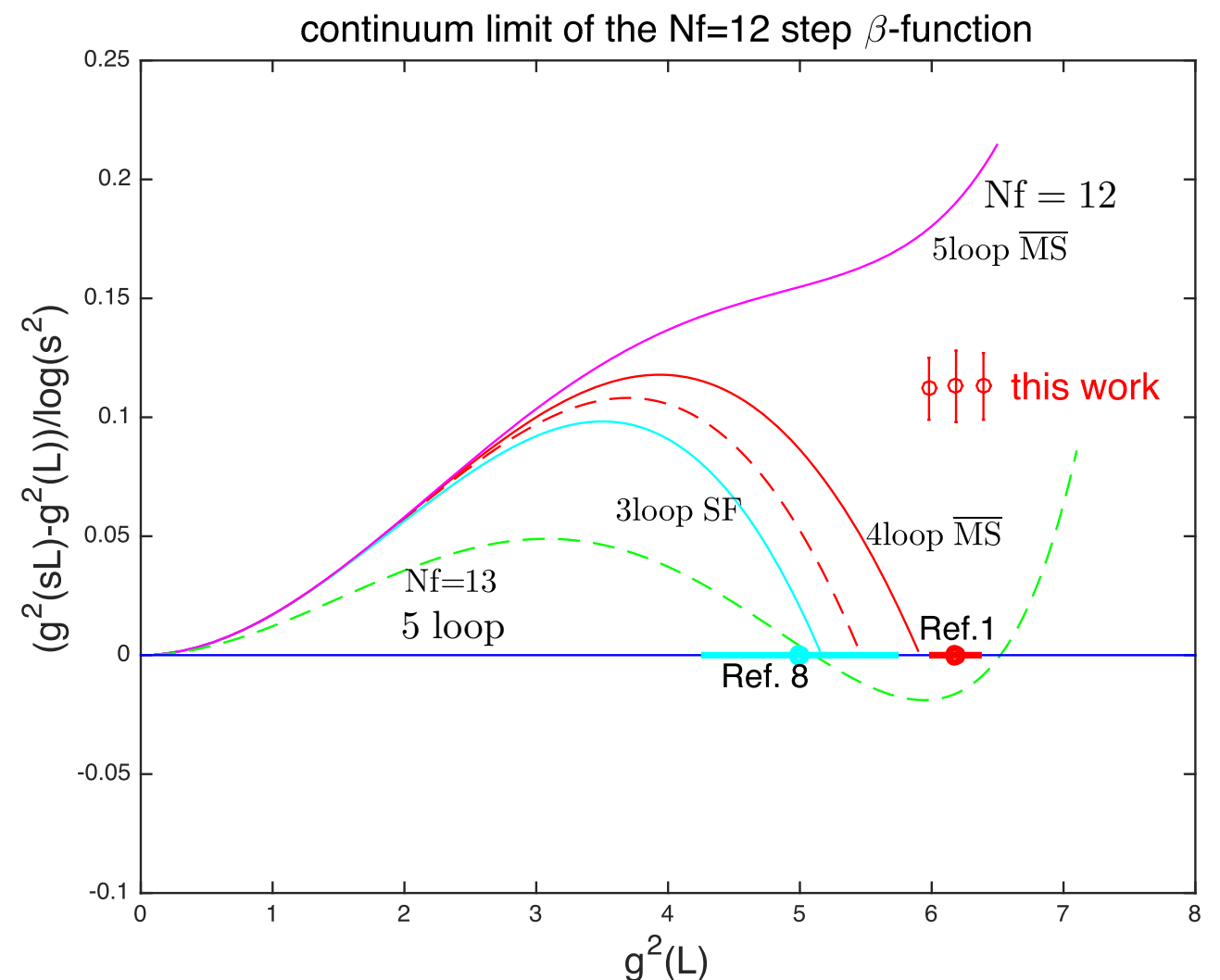
suggestive of merger of IRFP and UVFP at lower edge  
of conformal window (Kaplan-Son-Stephanov, Vecchi)

Can we test with staggered fermions if  $N_f = 13$  is conformal?

Interesting setting to test new predictions for conformal behavior (5-loop  $\overline{\text{MS}}$ -bar, Rytov-Shrock delta)

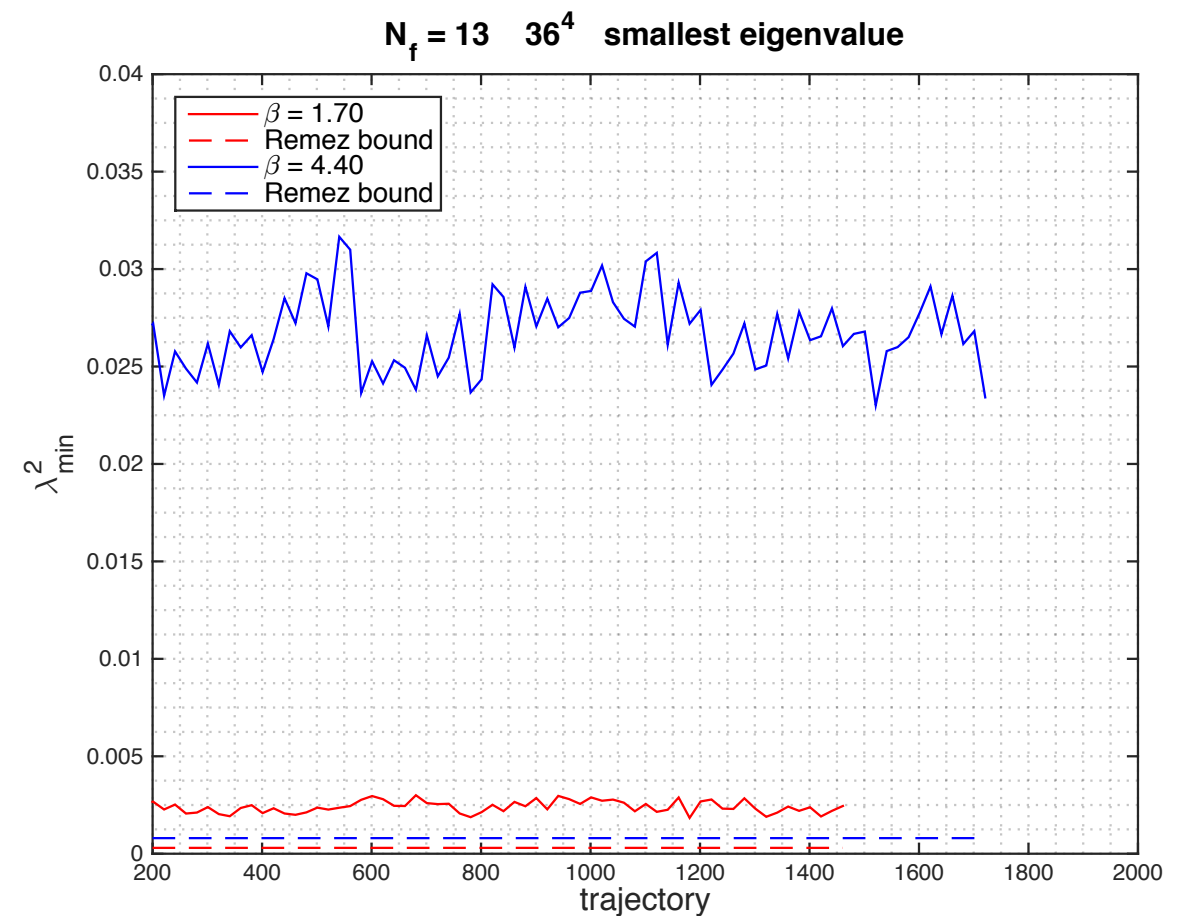
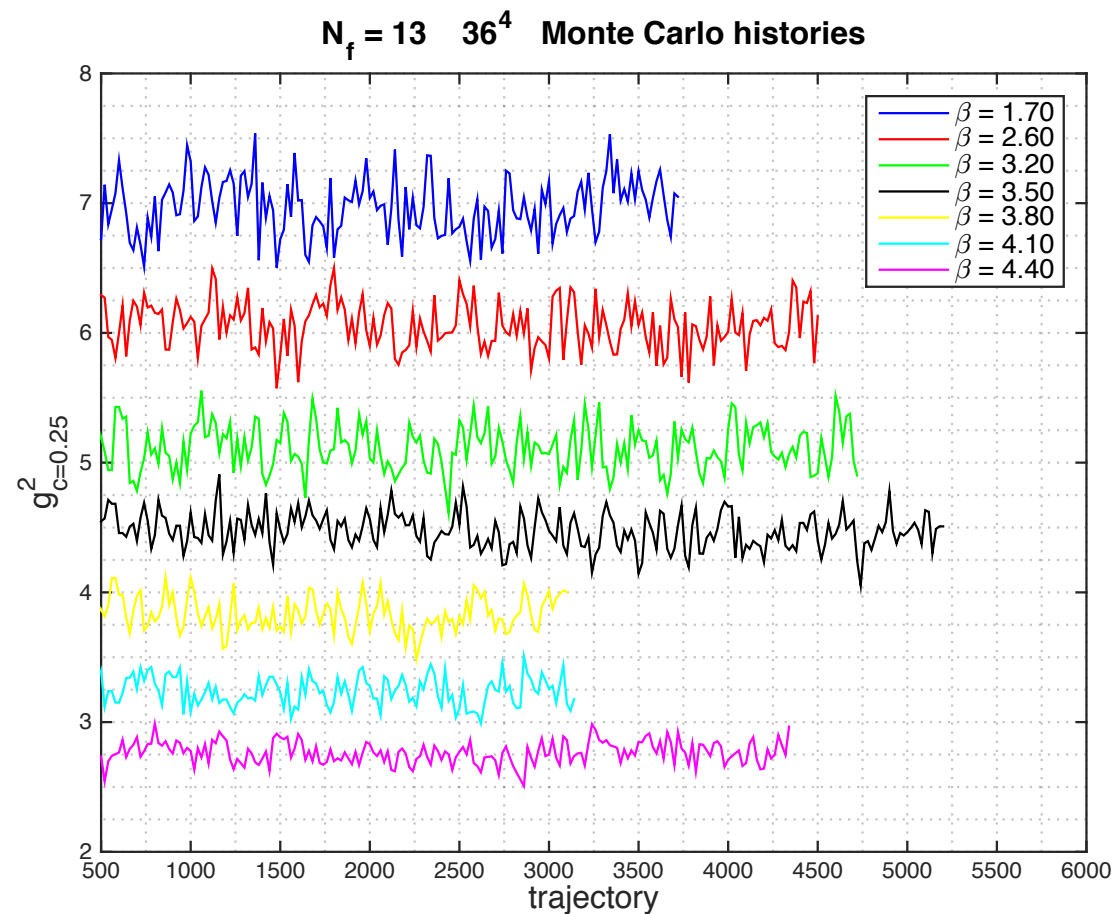
Lattice 2017: Daniel Nogradi presented results for  $\text{SU}(3)$   $N_f = 14$  fund rep and  $N_f = 3$  sextet rep  
perturbative beta function much smaller:  $< 0.014$  ( $N_f = 14$ ),  $< 0.001$  ( $N_f = 3$  sextet)  
continuum extrapolation difficult

$N_f = 13$  might be less difficult for continuum extrapolation



Fodor et al LatHC PRD94 091501R (2016)

EPJ Web Conf. 175 (2018) 08028



stout-smearred staggered fermions, Symanzik gauge action, RHMC algorithm: standard technology

for beta function with gradient flow and step-scaling, need a range of volumes and couplings:

$L = 12, 16, 18, 20, 24, 30, 32, 36, 48$  (with some 54)

bare coupling beta from 1.7 to 4.4 in steps of 0.3 (gives renormalized coupling in range 2 — 9)

large set of ensembles (9 volumes x 10 couplings) which allows various tests e.g.

step-scaling  $s = 3/2$ :  $12 \rightarrow 18$ ,  $16 \rightarrow 24$ ,  $20 \rightarrow 30$ ,  $24 \rightarrow 36$ ,  $32 \rightarrow 48$

## 5 lattice spacings for continuum extrapolation — crucial

several thousand trajectories — goal few per mille accuracy in renormalized couplings

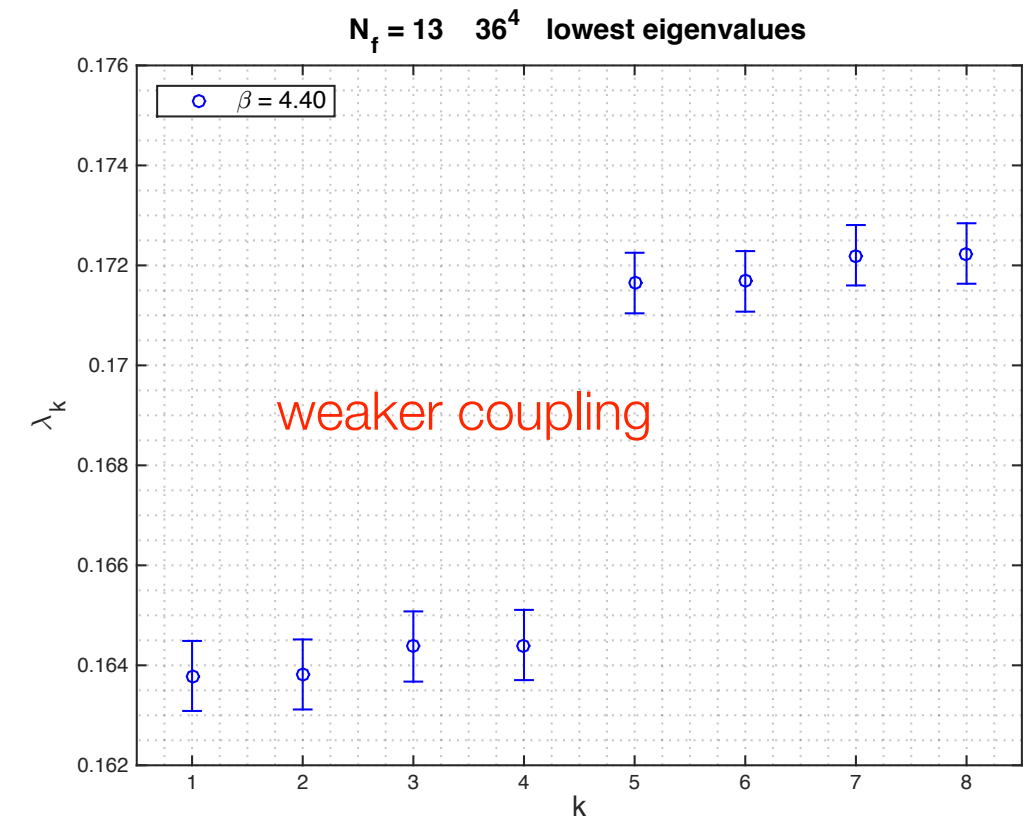
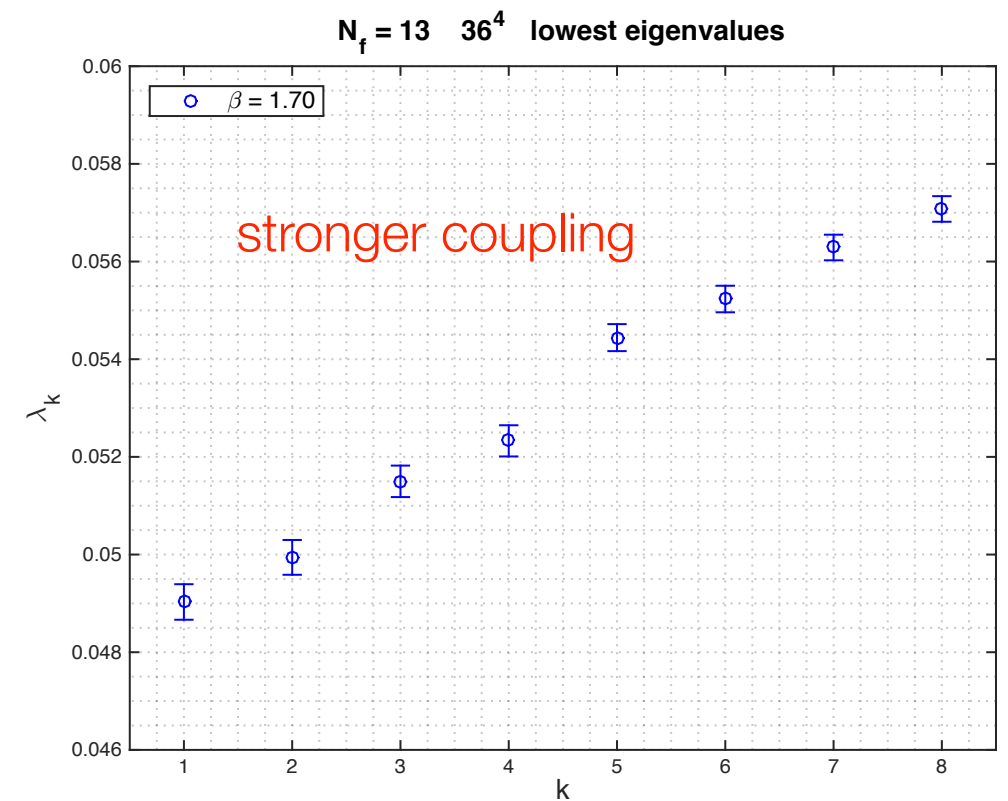
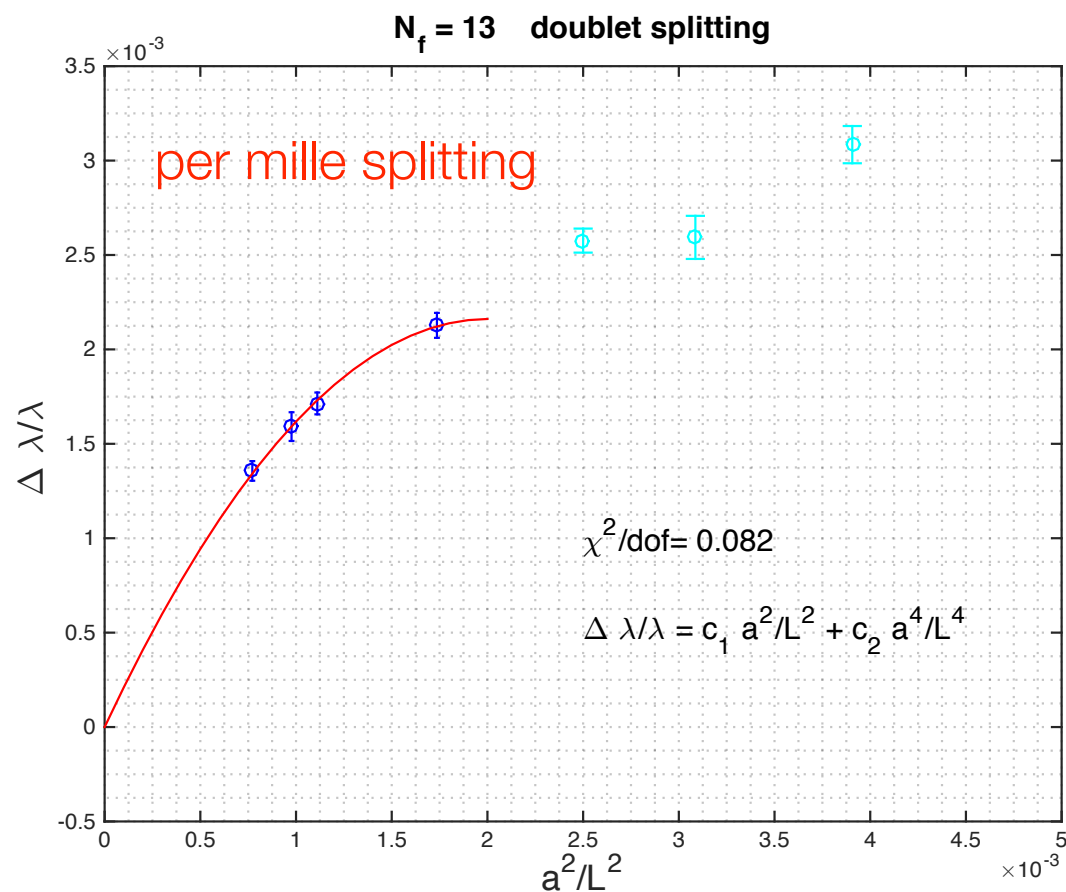
monitor RHMC algorithm throughout: lowest Dirac eigenvalue well above Remez bounds

is taste symmetry of staggered fermions recovered?

evident in Dirac operator eigenvalues

from strong to weak coupling: split eigenvalues first gather in pairs, then in quartets

eigenvalue splitting vanishes in the continuum limit, reaching the correct continuum theory: 1 of 3 arguments in response to claim staggered fermions are in wrong universality class



gradient flow: many possible discretizations e.g **SSC**

Symanzik gauge action for Monte Carlo generation and gradient flow

Clover operator for  $E = 1/2 \text{ tr } G_{\mu\nu}^2$

first look: step function at fixed bare coupling

polynomial interpolation of step in  $g^2$

data in groups of 5 for the possible steps  
 $L \rightarrow sL$  with  $s = 3/2$

interesting structure:

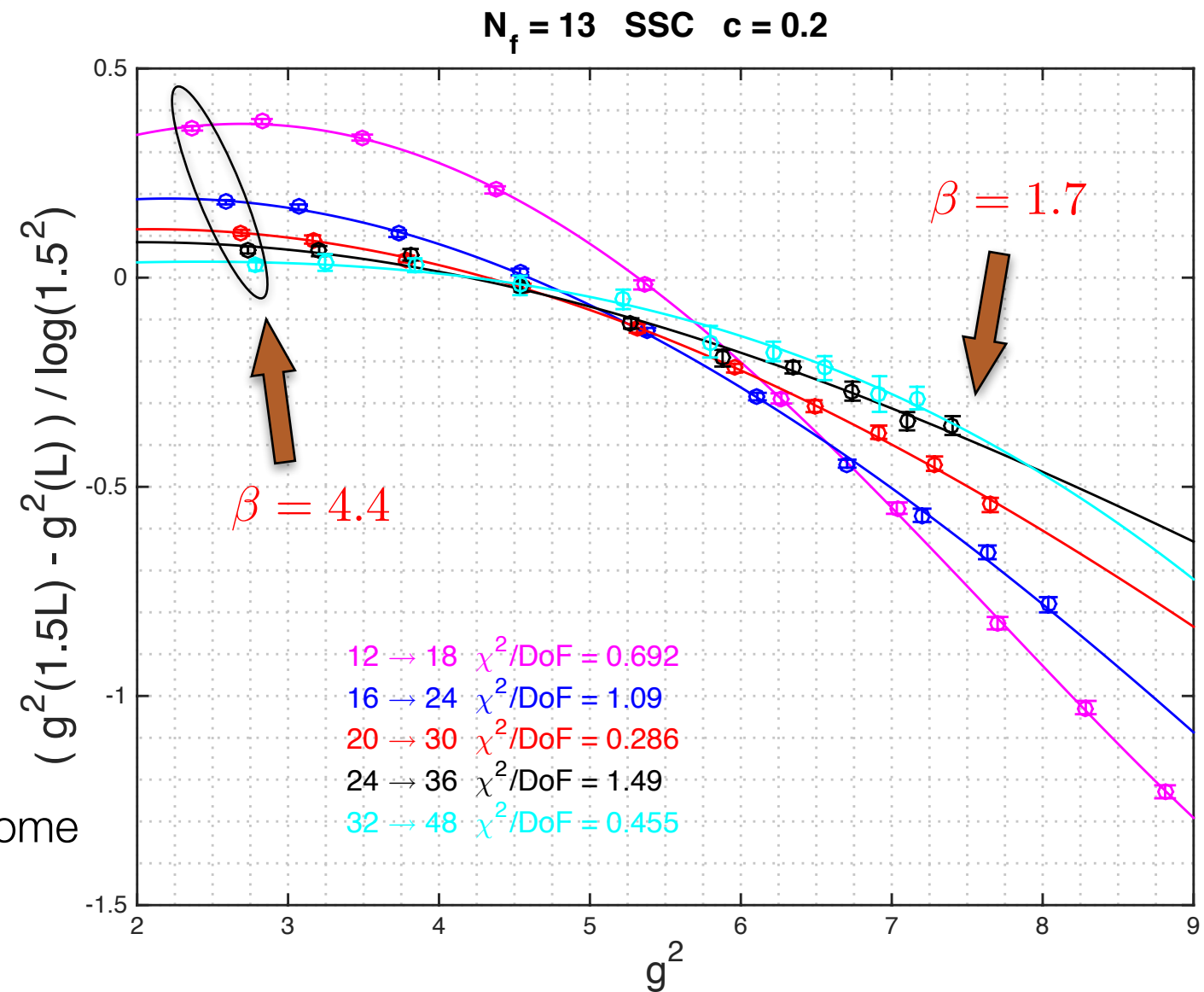
for each  $L \rightarrow sL$  pair, step function crosses zero at some renormalized coupling

cutoff effects appear to change sign as renormalized coupling increases i.e. approach continuum step function from above at weak coupling and from below at strong coupling

indication that continuum step function has a zero at an IRFP

gradient flow finite-volume scheme:  $c = \sqrt{8t}/L = 0.20$

balance between good statistical accuracy (smaller  $c$ ) and small cutoff effects (larger  $c$ )

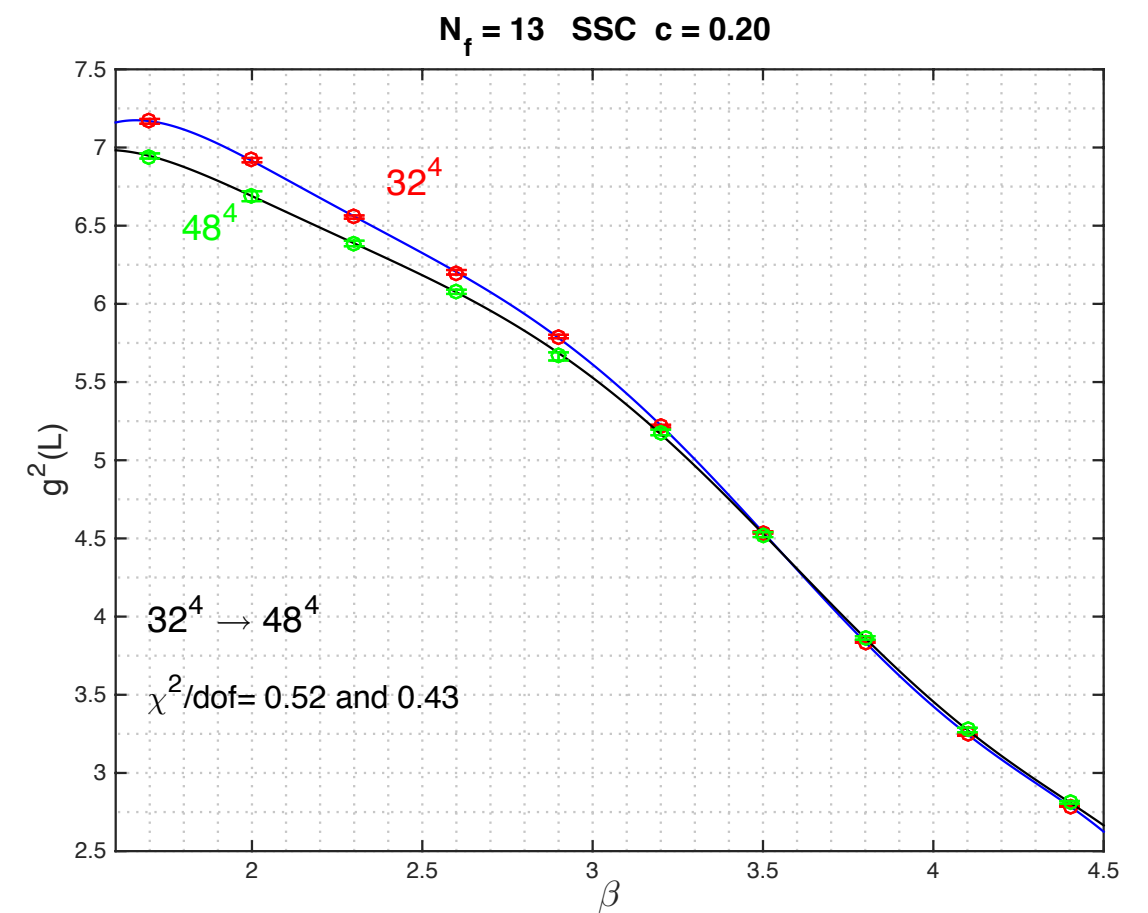
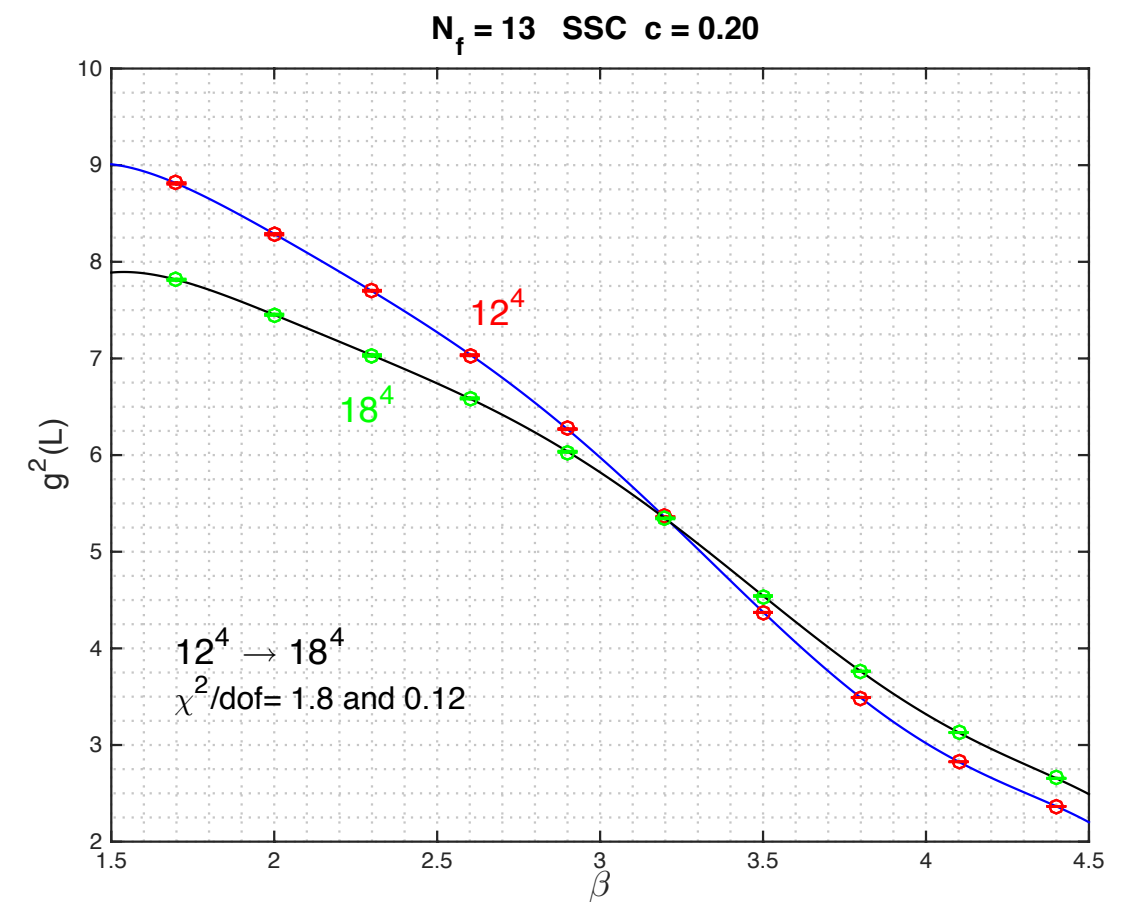




further analysis: polynomial interpolation of  $g^2(L)$  in  $\beta$

zero in step-function where curves cross

use interpolation to tune to chosen  $g^2(L)$



continuum extrapolation at tuned coupling:  $g^2(L) = 4.5$

leading cutoff effects expected to be  $\mathcal{O}(a^2)$

with next order terms  $\mathcal{O}(a^4)$

**blue curve**: quadratic fit in  $a^2$

few per mille accuracy for individual data points  
and continuum result

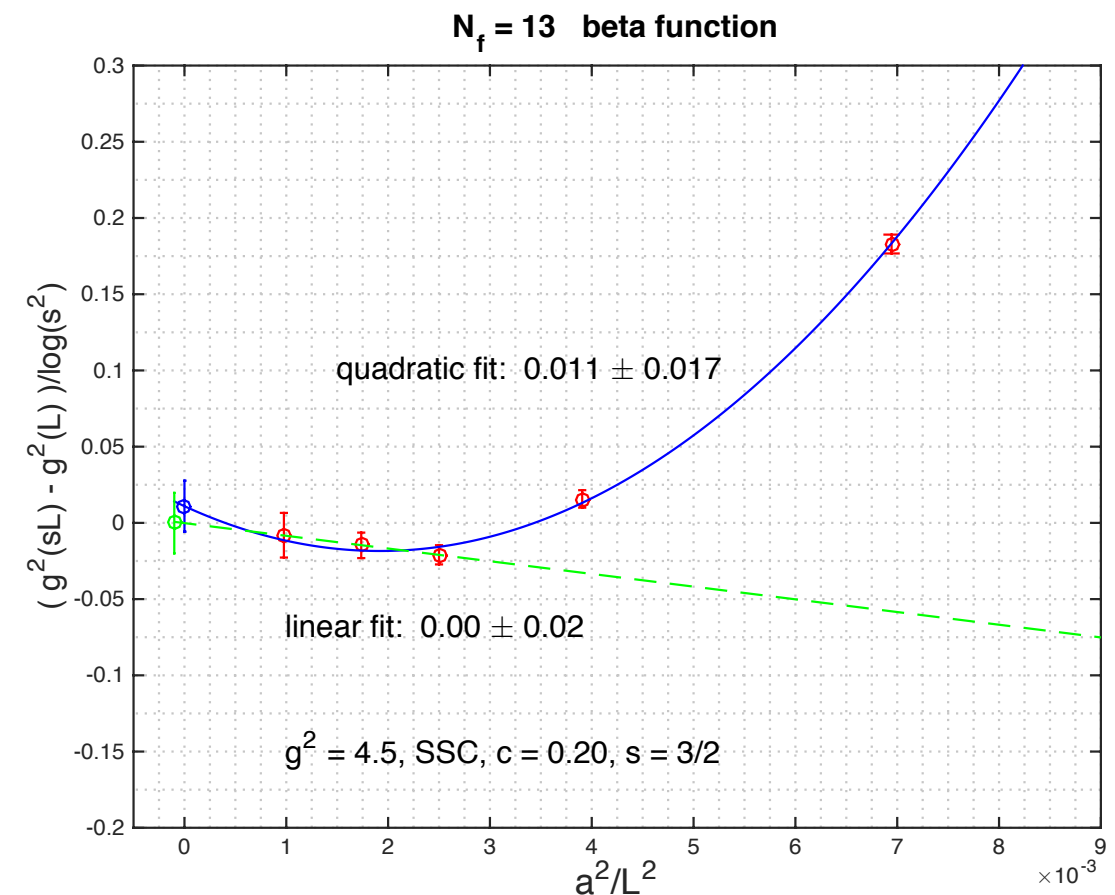
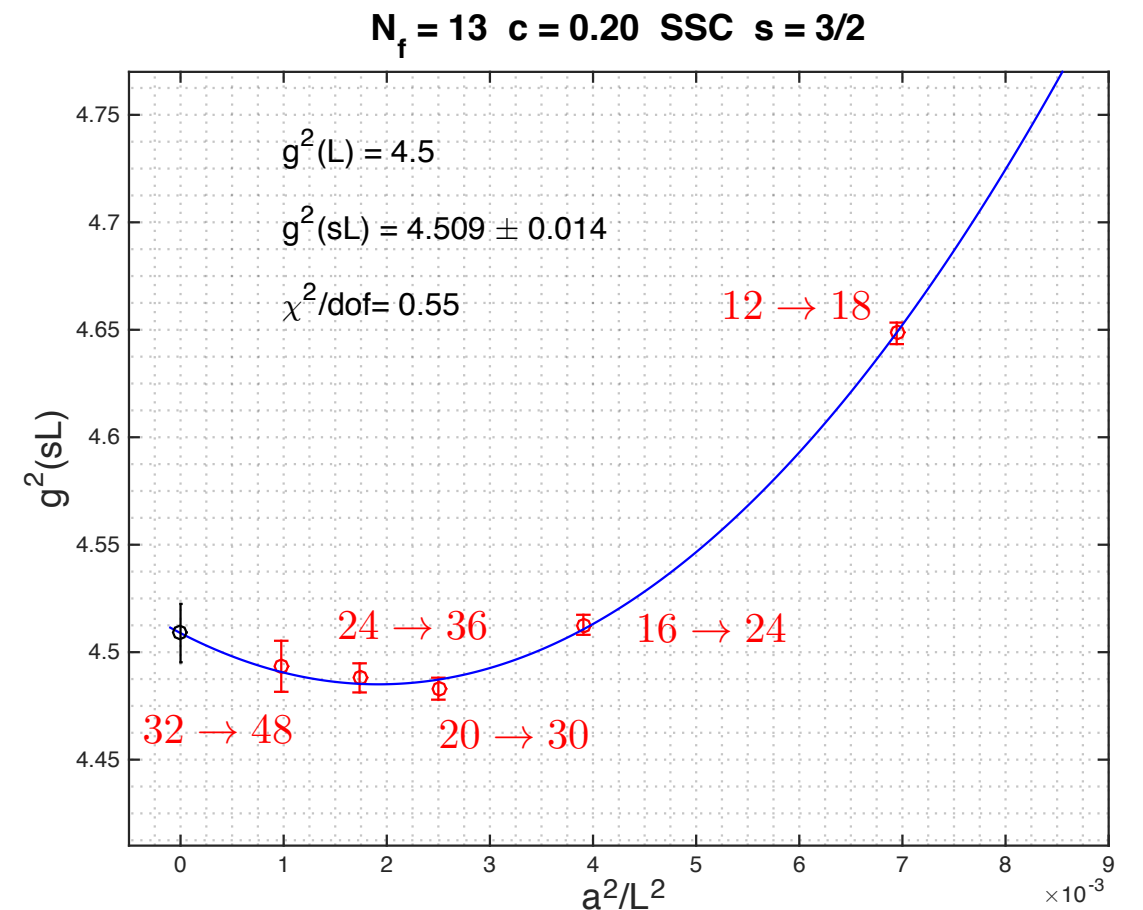
based on fit, coarser lattice spacing  $8 \rightarrow 12$

would be significantly removed from continuum value

discrete step function  $\frac{g^2(sL) - g^2(L)}{\log(s^2)}$ ,  $s = 3/2$

step function at this coupling is zero within error

linear fit in  $a^2$  using 3 finest lattice spacings only  
gives consistent result (**green curve**)





repeat procedure for range of tuned renormalized couplings from  $\sim 3$  to 4.5

shown here: result of quadratic continuum extrapolation

small but non-zero discrete step function

with an IRFP at  $g^2 \sim 4.5$

for comparison: 5-loop MS-bar beta function

has local max  $\sim 0.05$  and IRFP at  $g^2 \sim 5$

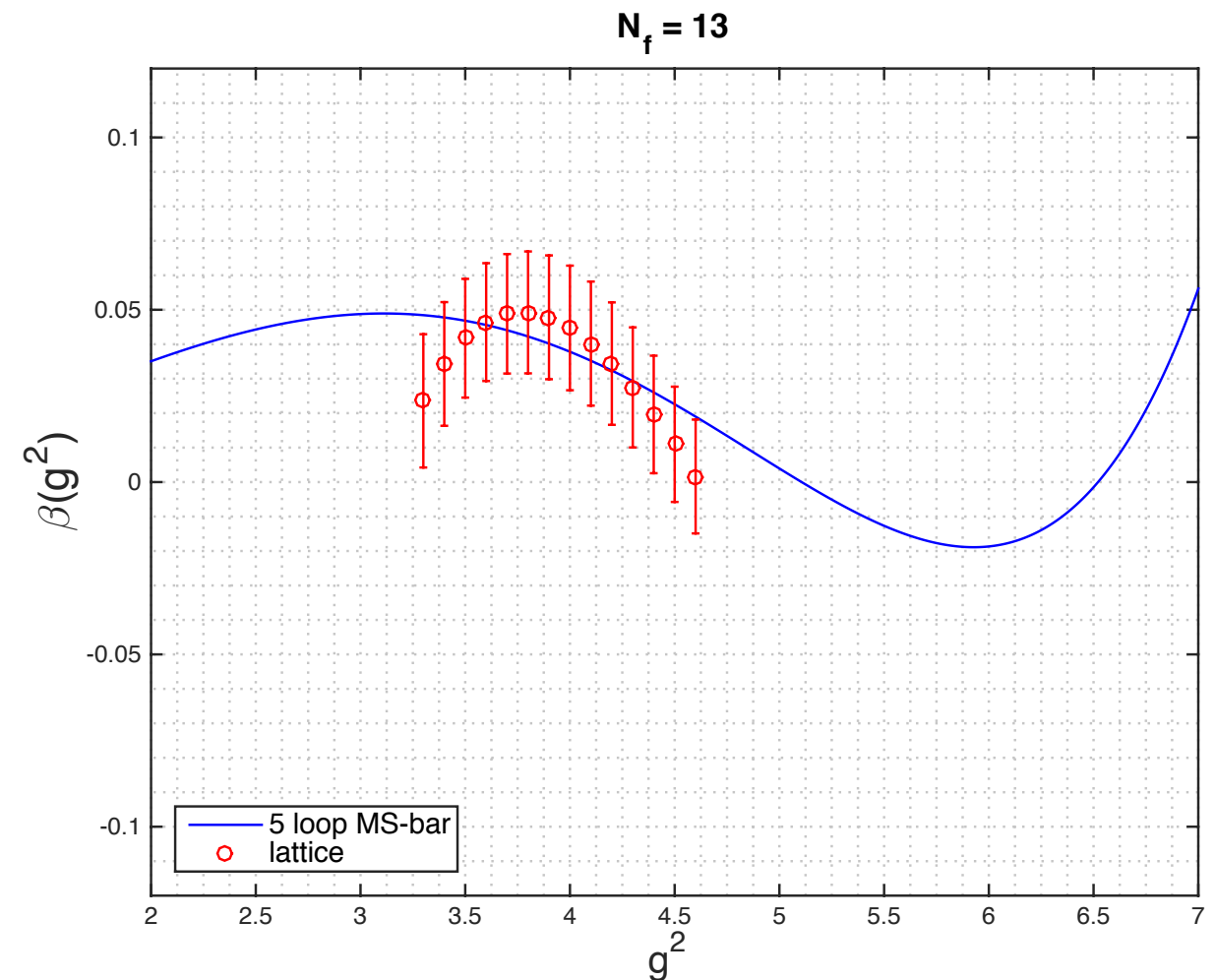
non-perturbative results have similar qualitative and even quantitative behavior

at lower end of coupling expect beta function is too small to distinguish from zero at this level of accuracy

**if we think beta function has an IRFP, can we demonstrate if it has conformal behavior?**

additional properties: anomalous mass dimension  $\gamma^*$

derivative of beta function at IRFP  $\beta' = d\beta/dg^2|_{g_*^2}$



new technique presented by LatHC group at Lattice 2015:  
Fodor et al PoS Lattice 2015 (2016) 310

full reconstruction of Dirac operator eigenvalue spectrum  
using Chebyshev polynomials to high order

$$\rho(t) = \frac{1}{\sqrt{1-t^2}} \sum_{k=0}^{\infty} c_k T_k(t)$$

uses recursive properties of Chebyshev polynomials and  
stochastic measurement of traces to extract  $c_k$

example:

20 stochastic noises

maximum polynomial order 8000

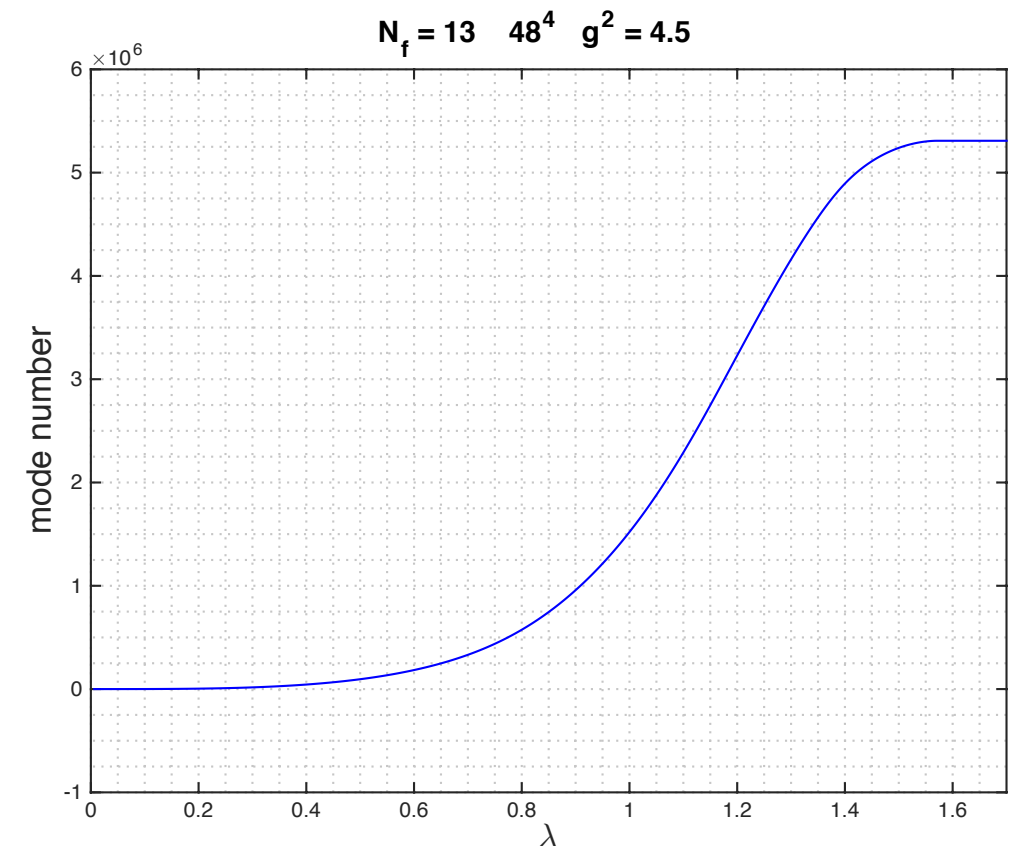
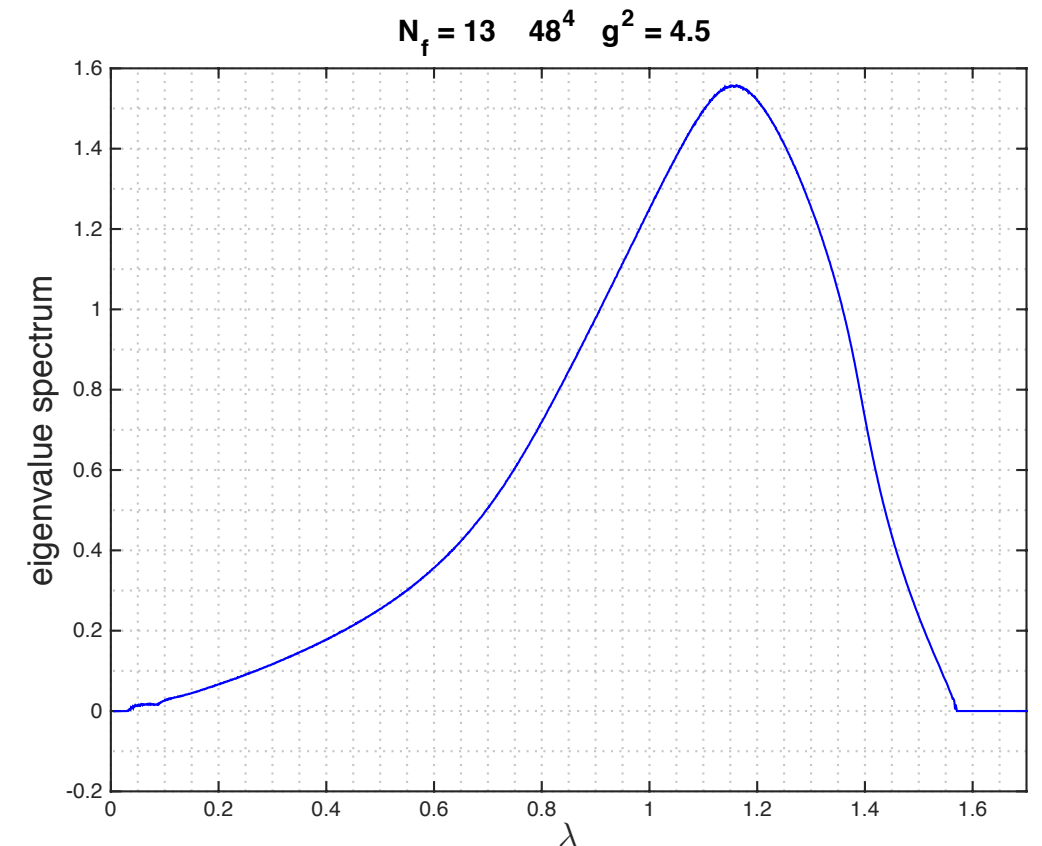
(bonus: all lower orders automatically generated)

statistical error is small, not visible on this scale

can reconstruct not only eigenvalue density but also e.g.  
mode number

$$\nu(\lambda) = \int_0^{\lambda} \rho(\lambda') d\lambda'$$

i.e. number of eigenvalues below some cut



new method presented by J Kuti for LatHC group at Simons Center Conformal Gauge Theories Workshop January 2018

anomalous dimension  $\gamma^*$  via step-scaling of mode number  
continuum limit similar to step-function

mode number is renormalized:

$$\nu_R(\lambda_R) = \nu(\lambda), \quad \lambda_R = Z_p^{-1} \cdot \lambda$$

measure mode number for eigenvalue cut set via  
finite volume

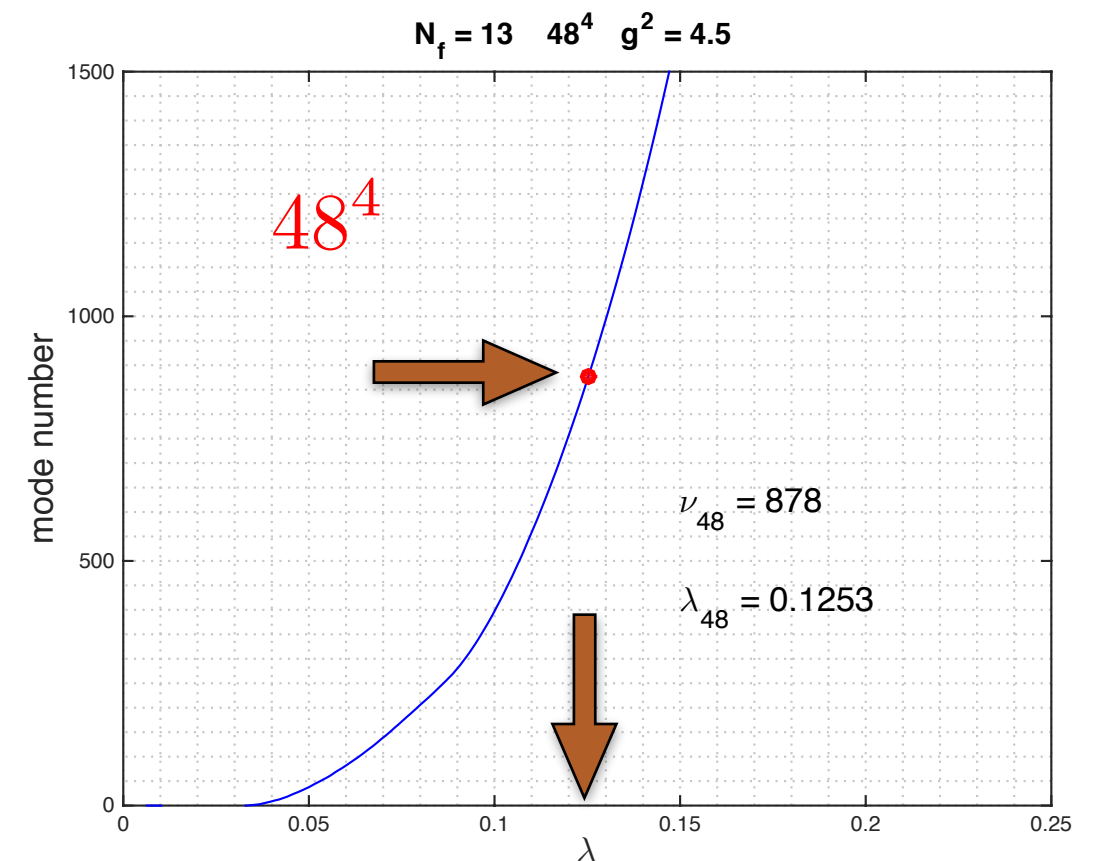
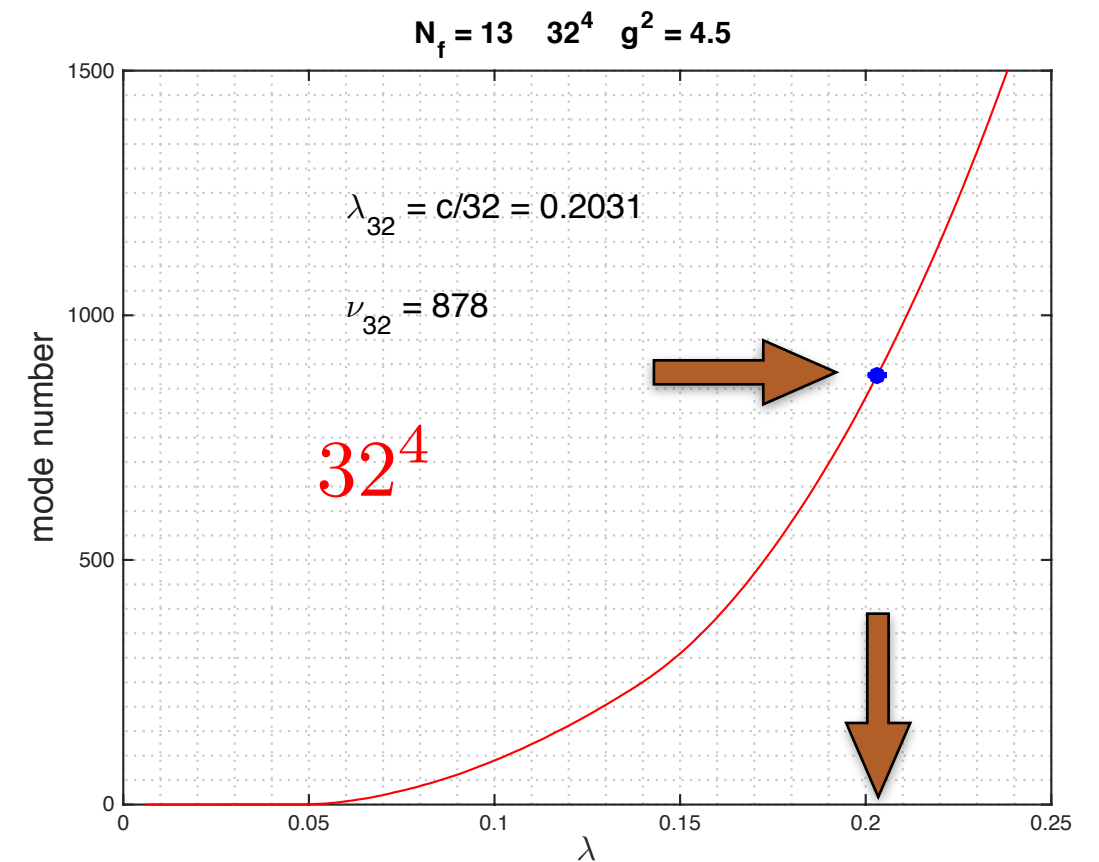
$$\lambda_L = \frac{c}{L}$$

example upper figure:  $L = 32, \quad c = 6.5$

match mode number with eigenvalue cut  $\lambda_{sL}$   
for larger volume  $sL$  e.g.  $s = 3/2, \quad sL = 48$

renormalization factors via eigenvalue ratio

$$\frac{\lambda_L}{\lambda_{sL}} = \frac{sZ_p(g_0, L/a)}{Z_p(g_0, sL/a)}$$



repeat for sequence of bare couplings and paired volumes  $L \rightarrow sL$  at tuned renormalized coupling

cutoff effects removed as  $a^2/L^2 \rightarrow 0$

**blue curve:** linear fit in  $a^2$

anomalous dimension

$$\gamma = \frac{\log Z_p(L)/Z_p(sL)}{\log(s)}$$

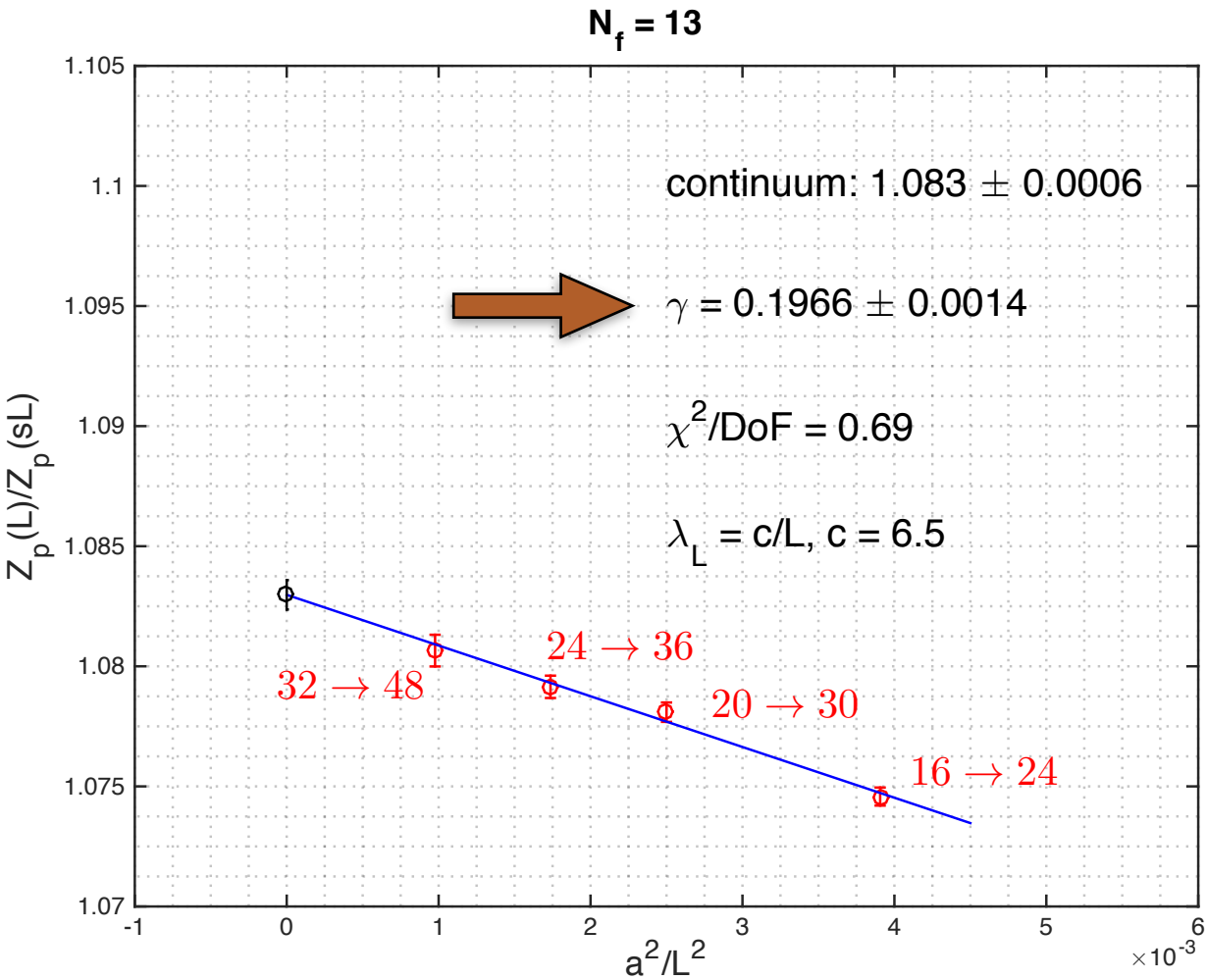
at renormalized coupling  $\sim 4.5$ , anomalous dimension 0.1966(14)

comparison with perturbation theory

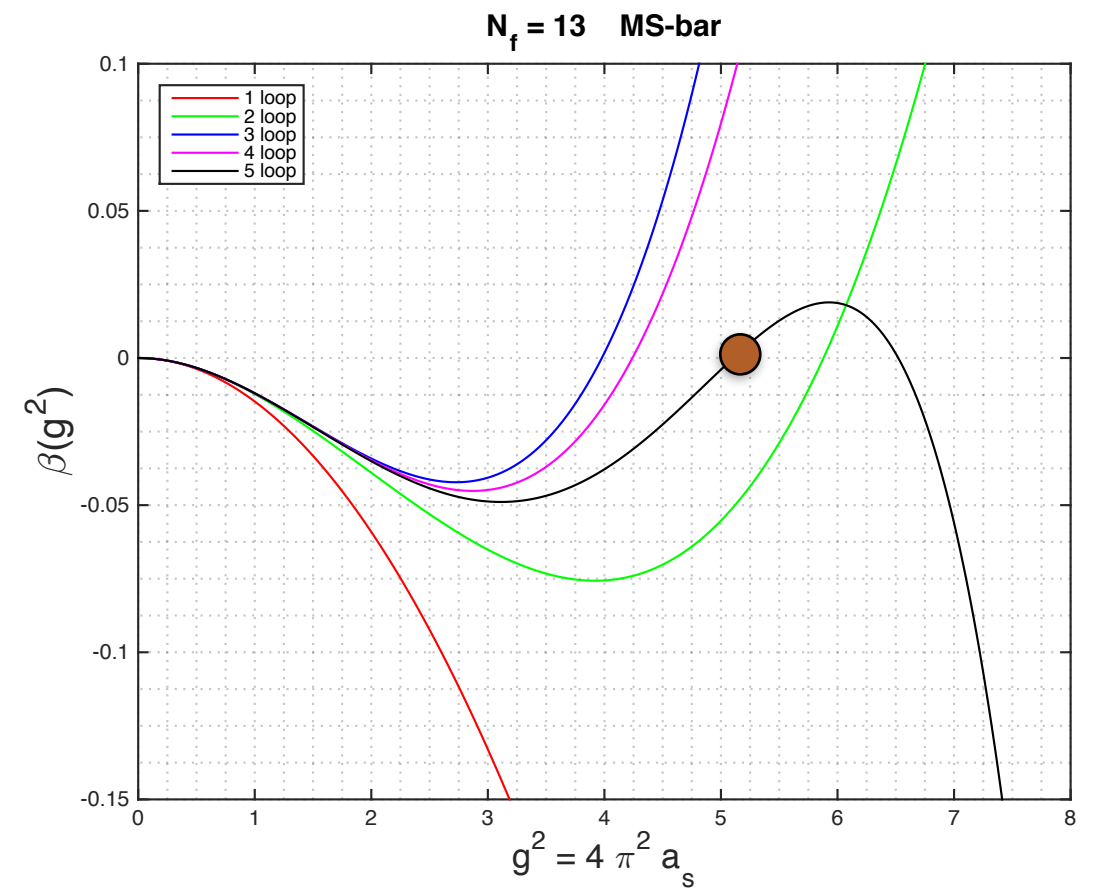
MS-bar: Baikov, Chetyrkin, Kühn  
JHEP 1410 (2014) 076

delta scheme: Rytlov, Shrock  
PRD95 (2017) 105004

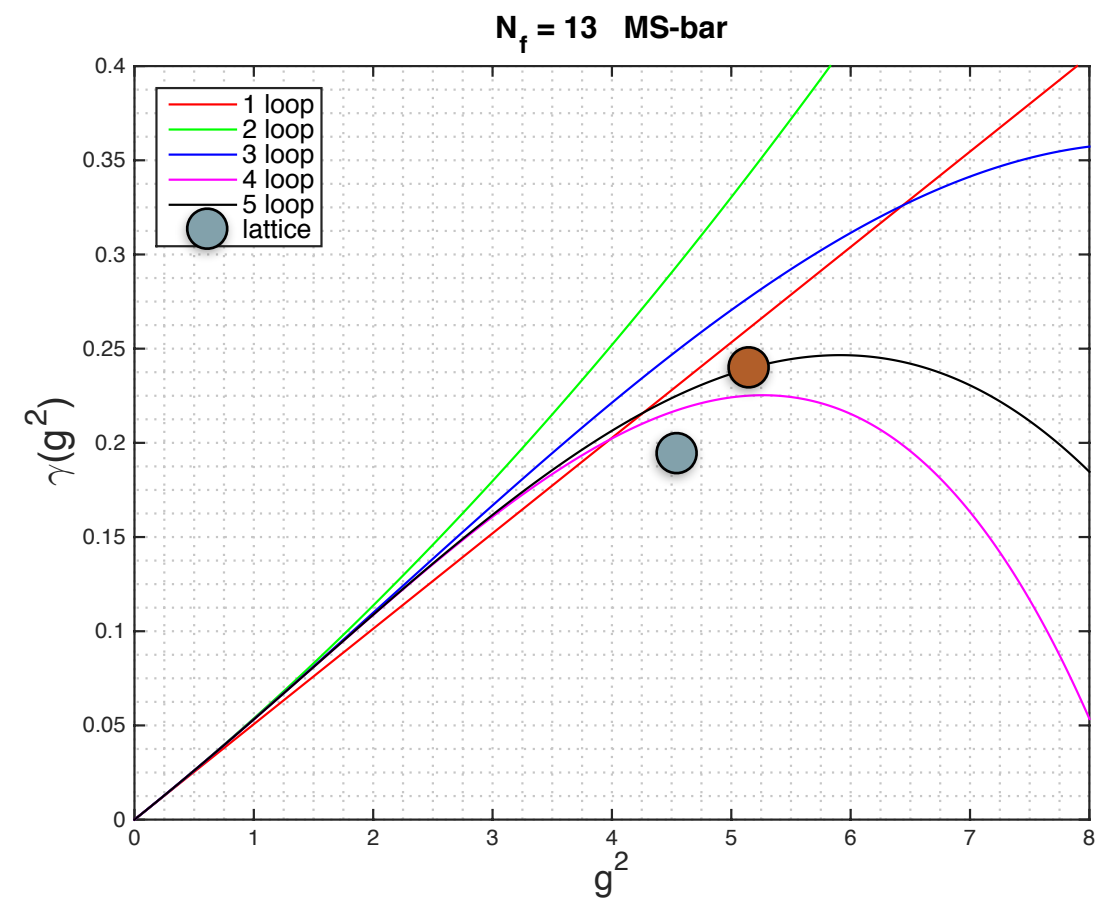
MS-bar loop order	$\gamma^*$	Ryttov-Shrock order	$\gamma^*$
2	0.404	1	0.174
3	0.220	2	0.221
4	0.210	3	0.231
5	0.239	4	0.237



perturbative predictions for anomalous dimension  
rather stable with increasing order



non-perturbative lattice result quite close to high loop  
order prediction



what happens at stronger coupling?

**blue curve:** quadratic fit in  $a^2$

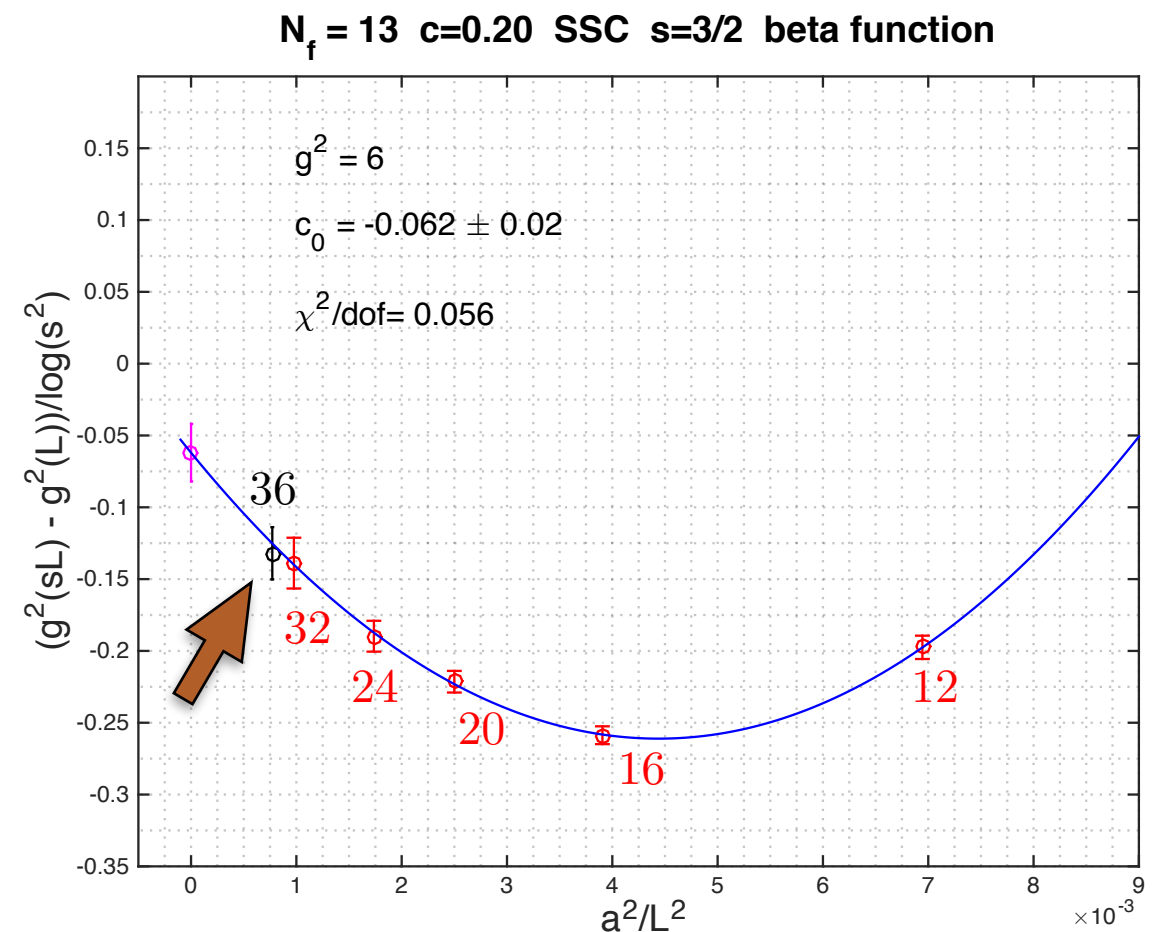
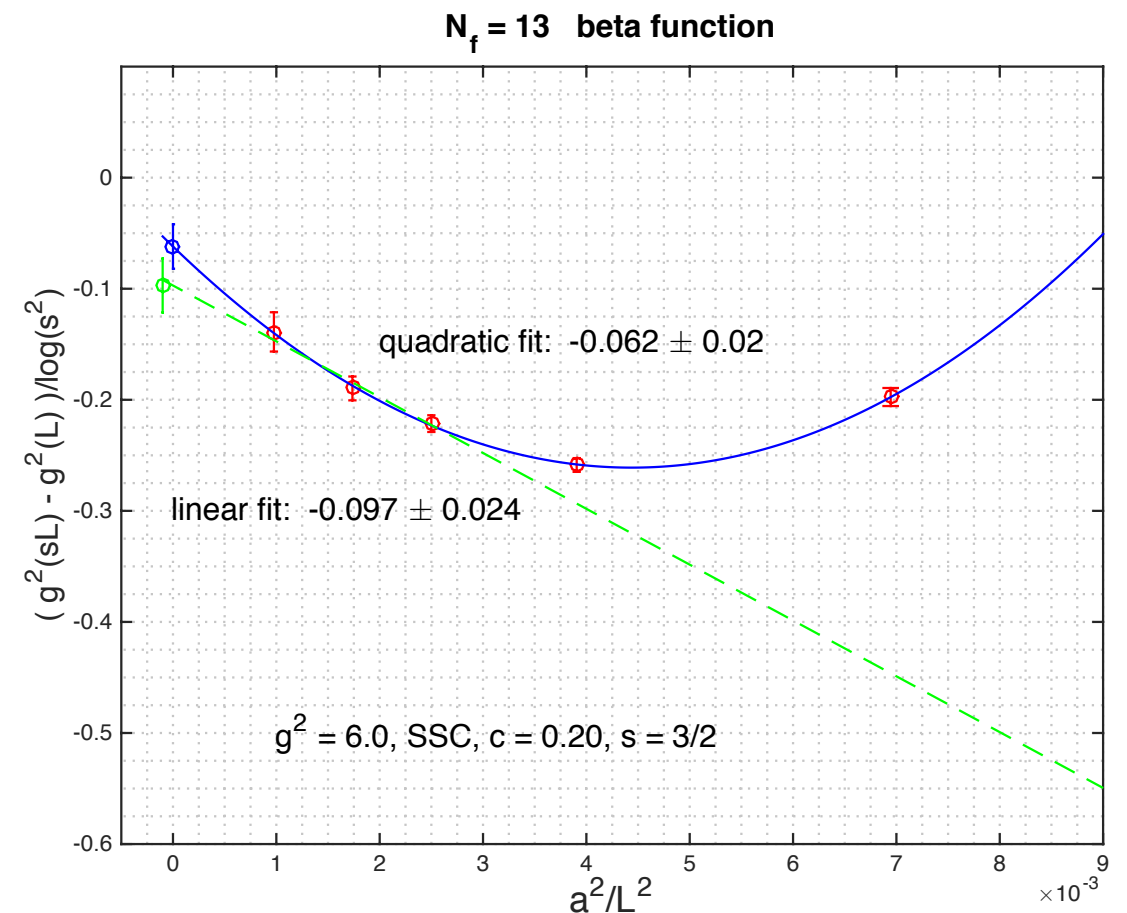
**green curve:** linear fit in  $a^2$   
of 3 smallest lattice spacings

continuum step-function appears to change sign  
with increasing renormalized coupling

to check extrapolation: additional ensemble  $54^4$

one more step  $36 \rightarrow 54$

**black data point** lies on curve predicted by fit of  
other lattice spacing data





quadratic continuum extrapolation across range of renormalized couplings

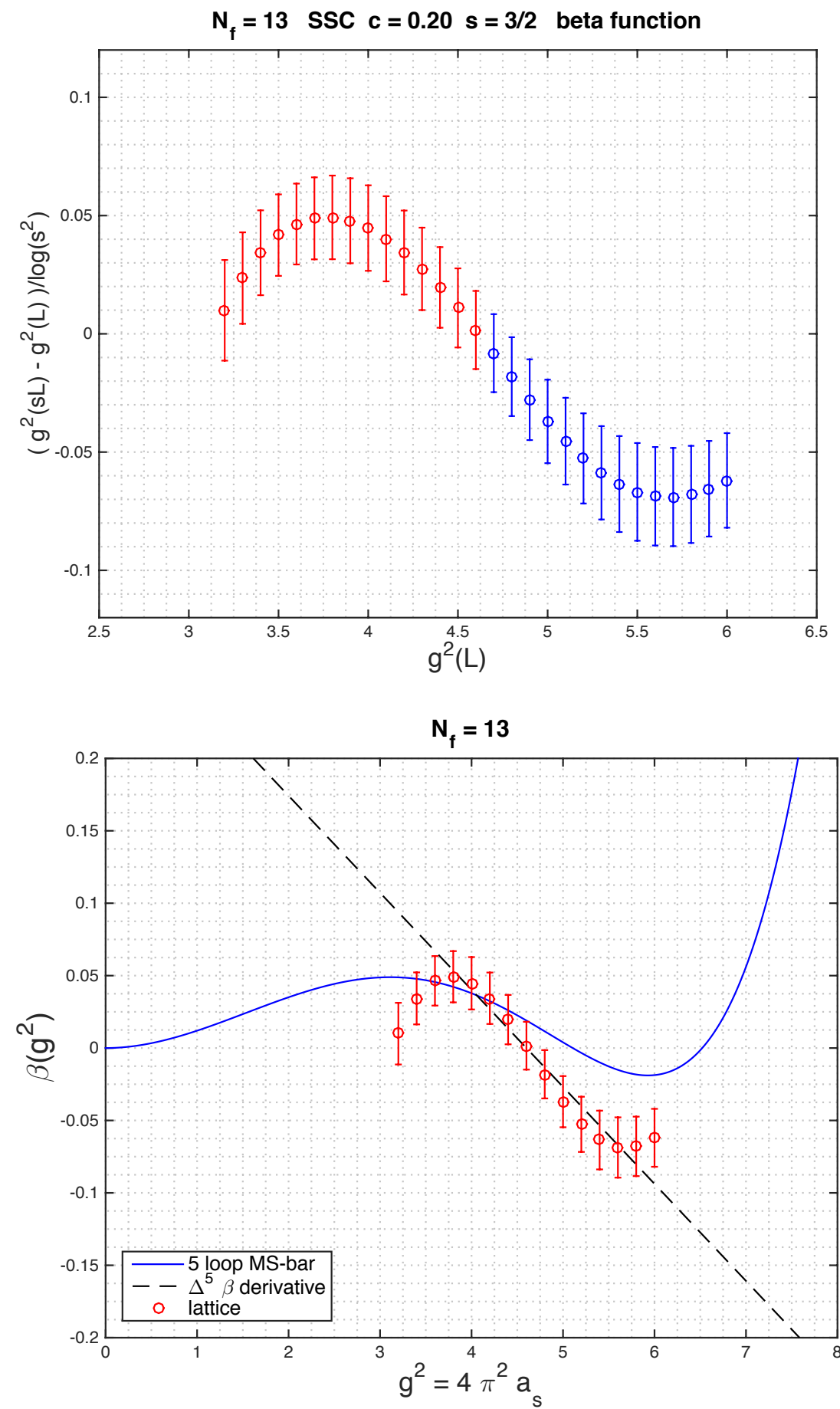
physical property: slope of beta function at IRFP

$$\beta' = d\beta/dg^2|_{g_*^2}$$

MS-bar loop order	$\beta'$	Rytov-Shrock delta order	$\beta'$
2	0.087	2	0.051
3	0.078	3	0.072
4	0.075	4	0.069
5	0.037	5	0.067

perturbative results quite close to lattice data

5-loop MS-bar beta function less steep than lower orders



compare discretization choice

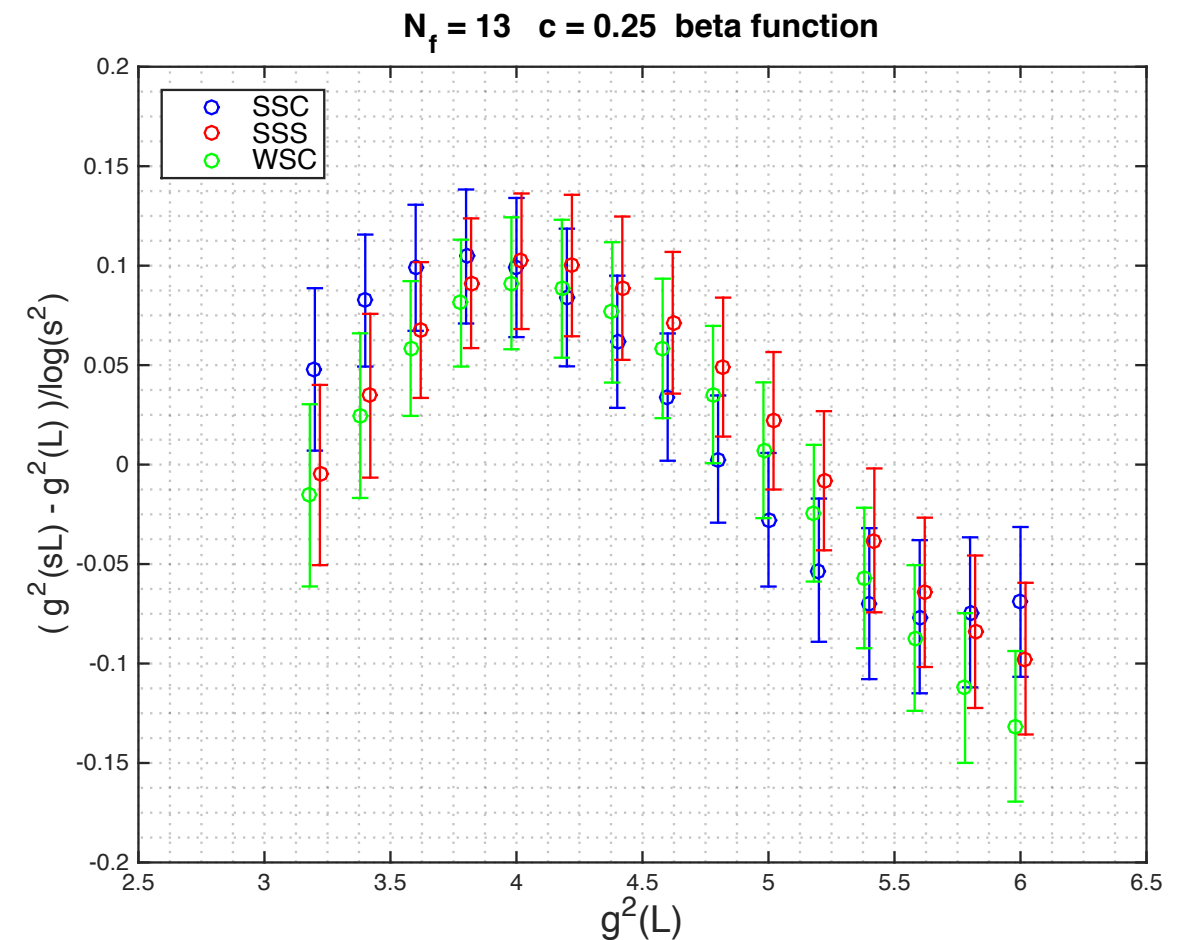
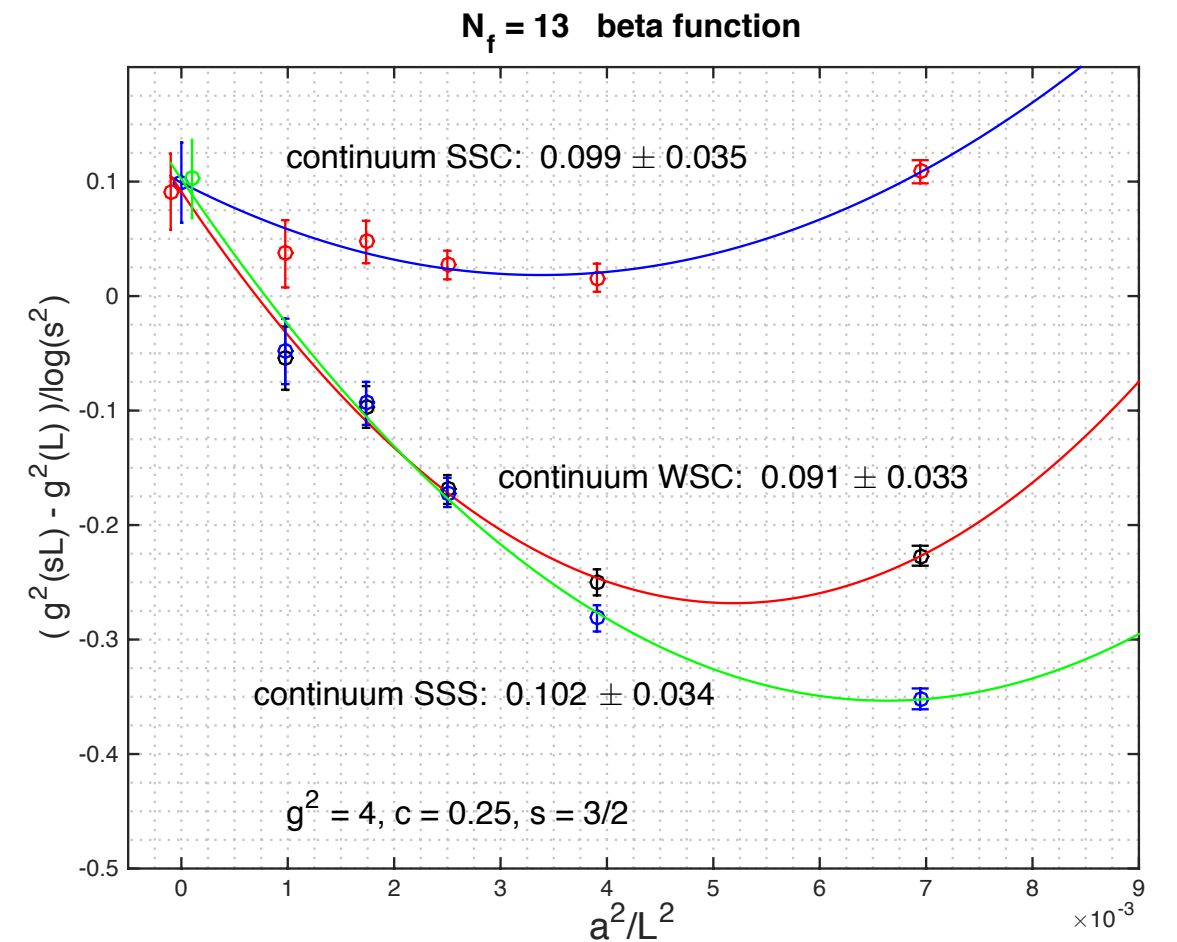
**SSC**: Symanzik for MC and Flow, Clover for E

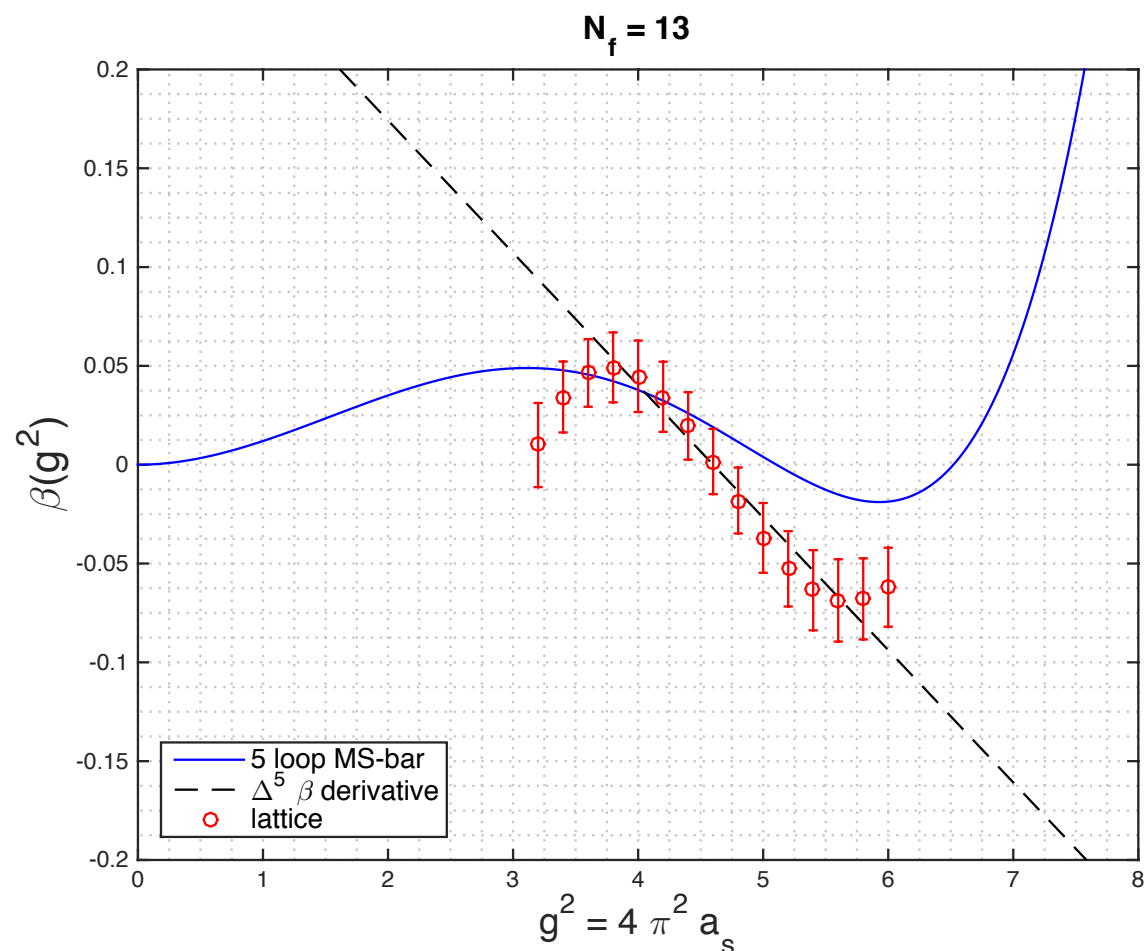
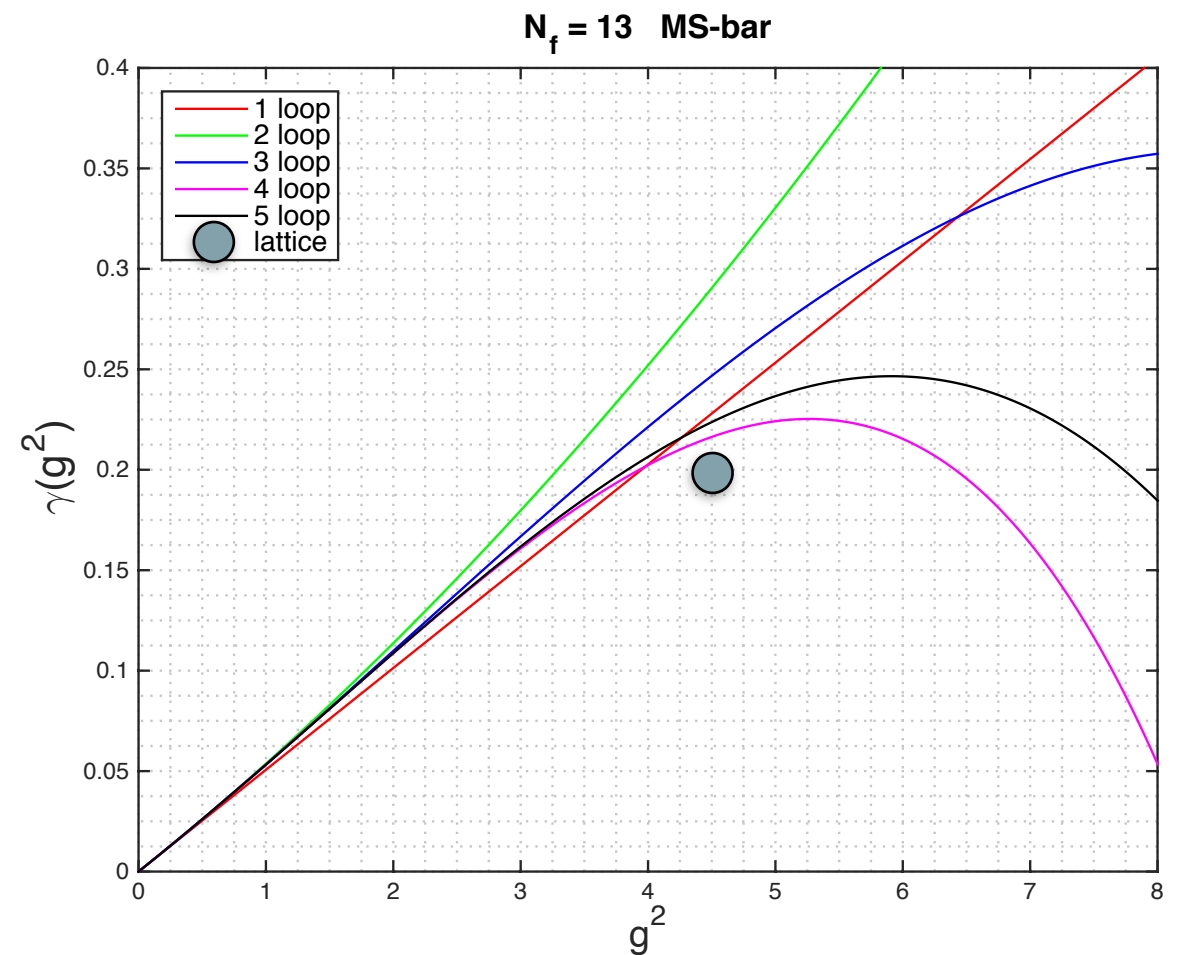
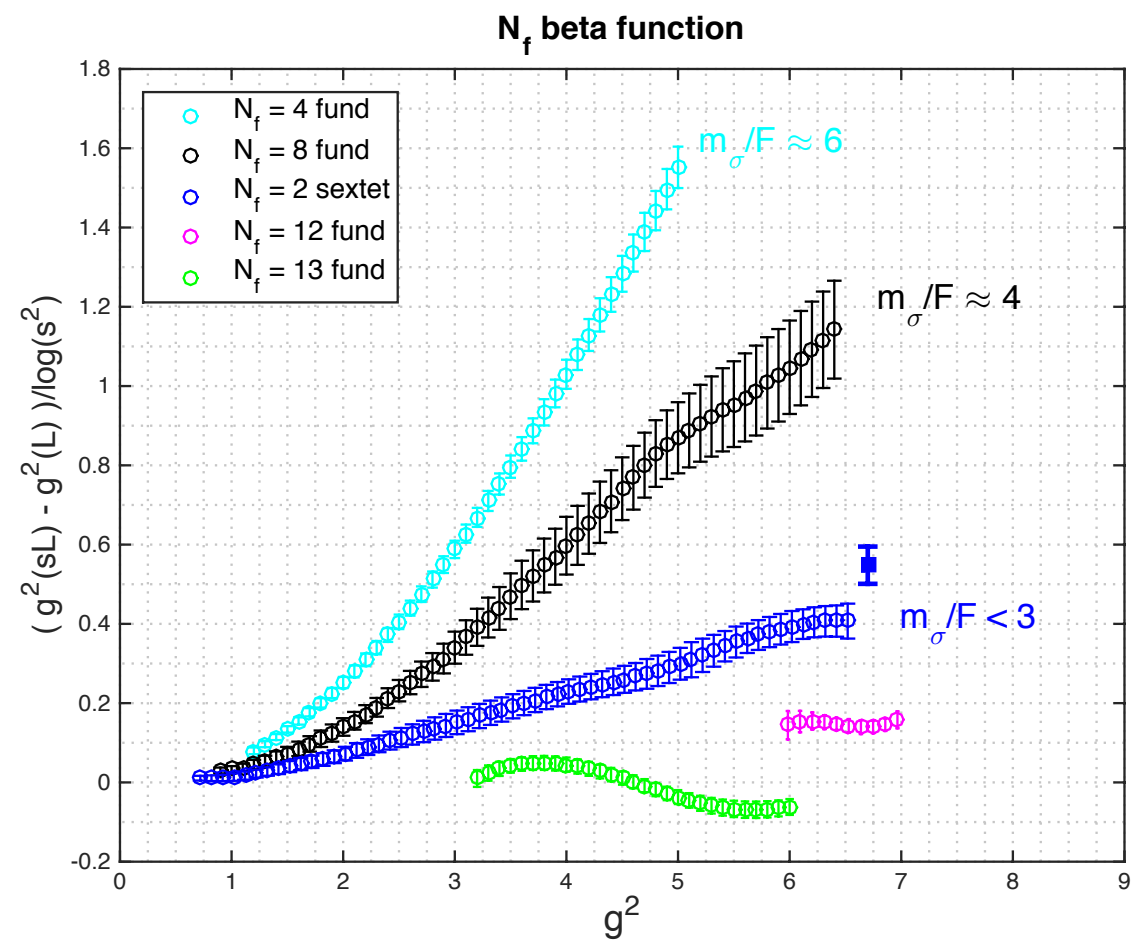
**SSS**: Symanzik for all 3

**WSC**: Wilson for Flow, Symanzik for MC, Clover for E

larger cutoff effects for WSC and SSS — similar to our experience in other  $N_f$  beta function studies

consistent results in the continuum limit from independent fits — a very useful crosscheck





**$N_f = 13$  does appear conformal:** not just a zero in the beta function, but also conformal like in physical properties, can test new perturbative results

to date, looks like  $N_f = 12$  walking, but not inside the conformal window

where does  $N_f = 10$  appear in this picture?

wait for more: Daniel Nogradi on  $N_f = 10$  (next talk)