

Investigation of the 1+1 dimensional Thirring model using the method of matrix product states

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Outline

- Preliminaries
- Lattice simulations, the MPS and DMRG
- Phase structure of the Thirring model
- Remarks and outlook

Preliminaries

Motivation

- Tensor network for lattice field theory
- Topological phase transitions
- Real-time dynamics (long-term goal)

The 1+1 dimensional Thirring model and its duality to the sine-Gordon model

$$S_{\text{Th}} [\psi, \bar{\psi}] = \int d^2x \left[\bar{\psi} i \gamma^\mu \partial_\mu \psi - m_0 \bar{\psi} \psi - \frac{g}{2} (\bar{\psi} \gamma_\mu \psi)^2 \right]$$



strong-weak duality $g \leftrightarrow \kappa$

$$S_{\text{SG}} [\phi] = \int d^2x \left[\frac{1}{2} \partial_\mu \phi(x) \partial^\mu \phi(x) + \frac{\alpha_0}{\kappa^2} \cos(\kappa \phi(x)) \right]$$

$$\xrightarrow{\phi \rightarrow \phi/\kappa, \text{ and } \kappa^2 = t} \frac{1}{t} \int d^2x \left[\frac{1}{2} \partial_\mu \phi(x) \partial^\mu \phi(x) + \alpha_0 \cos(\phi(x)) \right]$$

Works in the zero-charge sector

Dualities and phase structure

Thirring	sine-Gordon	XY
g	$\frac{4\pi^2}{t} - \pi$	$\frac{T}{K} - \pi$

★ The K-T phase transition at $T \sim K\pi/2$ in the XY model.

$g \sim -\pi/2$, Coleman's instability point

★ The phase boundary at $t \sim 8\pi$ in the sine-Gordon theory.

➡ The cosine term becomes relevant or irrelevant.

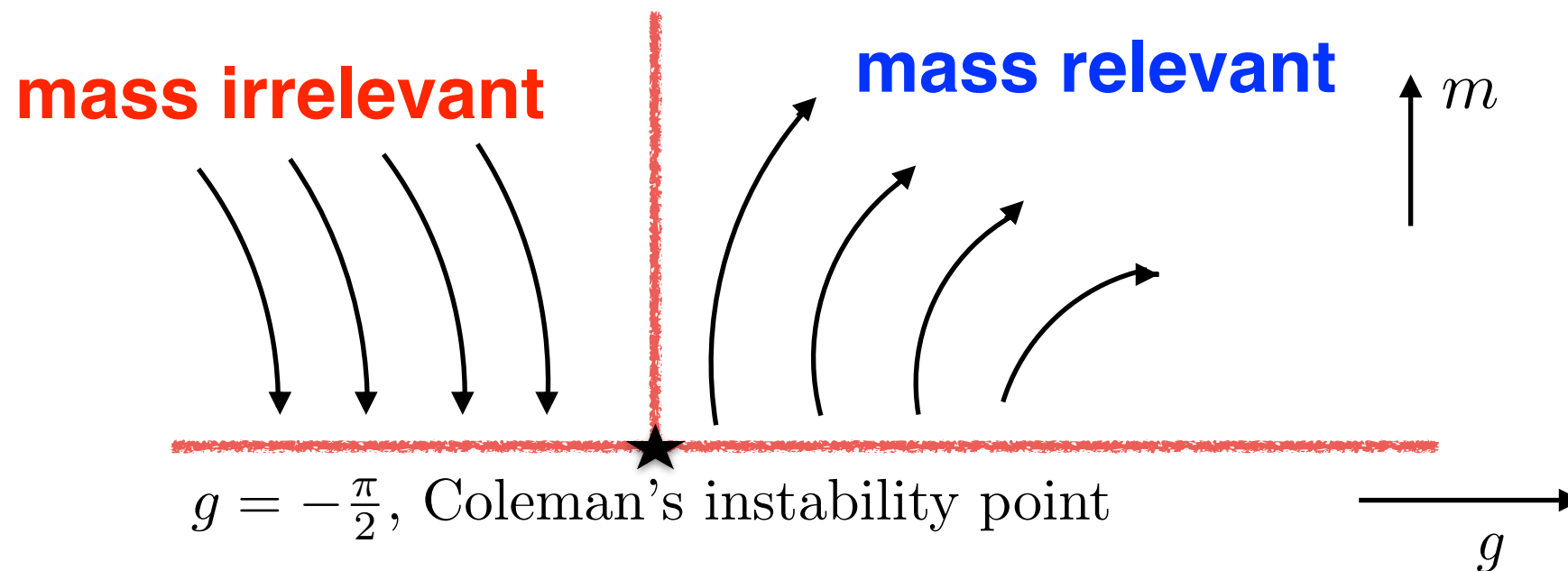
Thirring	sine-Gordon
$\bar{\psi}\gamma_\mu\psi$	$\frac{1}{2\pi}\epsilon_{\mu\nu}\partial_\nu\phi$
$\bar{\psi}\psi$	$\frac{\Lambda}{\pi}\cos\phi$

RG flows of the Thirring model

$$\beta_g \equiv \mu \frac{dg}{d\mu} = -64\pi \frac{m^2}{\Lambda^2},$$

$$\beta_m \equiv \mu \frac{dm}{d\mu} = \frac{-2(g + \frac{\pi}{2})}{g + \pi} m - \frac{256\pi^3}{(g + \pi)^2 \Lambda^2} m^3$$

★ Massless Thirring model is a conformal field theory



Lattice simulations, the MPS and DMRG

Operator formalism and the Hamiltonian

- Operator formalism of the Thirring model Hamiltonian

C.R. Hagen, 1967

$$H_{\text{Th}} = \int dx \left[-i\bar{\psi}\gamma^1\partial_1\psi + m_0\bar{\psi}\psi + \frac{g}{4}(\bar{\psi}\gamma^0\psi)^2 - \frac{g}{4}\left(1 + \frac{2g}{\pi}\right)^{-1}(\bar{\psi}\gamma^1\psi)^2 \right]$$

- Staggering, J-W transformation ($S_j^\pm = S_j^x \pm iS_j^y$):

J. Kogut and L. Susskind, 1975; A. Luther, 1976

$$\bar{H}_{X X Z} = \nu(g) \left[-\frac{1}{2} \sum_n^{N-2} (S_n^+ S_{n+1}^- + S_{n+1}^+ S_n^-) + a\tilde{m}_0 \sum_n^{N-1} (-1)^n \left(S_n^z + \frac{1}{2} \right) + \Delta(g) \sum_n^{N-1} \left(S_n^z + \frac{1}{2} \right) \left(S_{n+1}^z + \frac{1}{2} \right) \right]$$

$$\nu(g) = \frac{2\gamma}{\pi \sin(\gamma)}, \quad \tilde{m}_0 = \frac{m_0}{\nu(g)}, \quad \Delta(g) = \cos(\gamma), \quad \text{with } \gamma = \frac{\pi - g}{2}$$

$$\bar{H}_{X X Z}^{(\text{penalty})} = \bar{H}_{X X Z} + \lambda \left(\sum_{n=0}^{N-1} S_n^z - S_{\text{target}} \right)^2$$

projected to a sector of total spin

JW-trans of the total fermion number,
Bosonise to topological index in the SG theory.

Simulation details

- Matrix product operator for the Hamiltonian

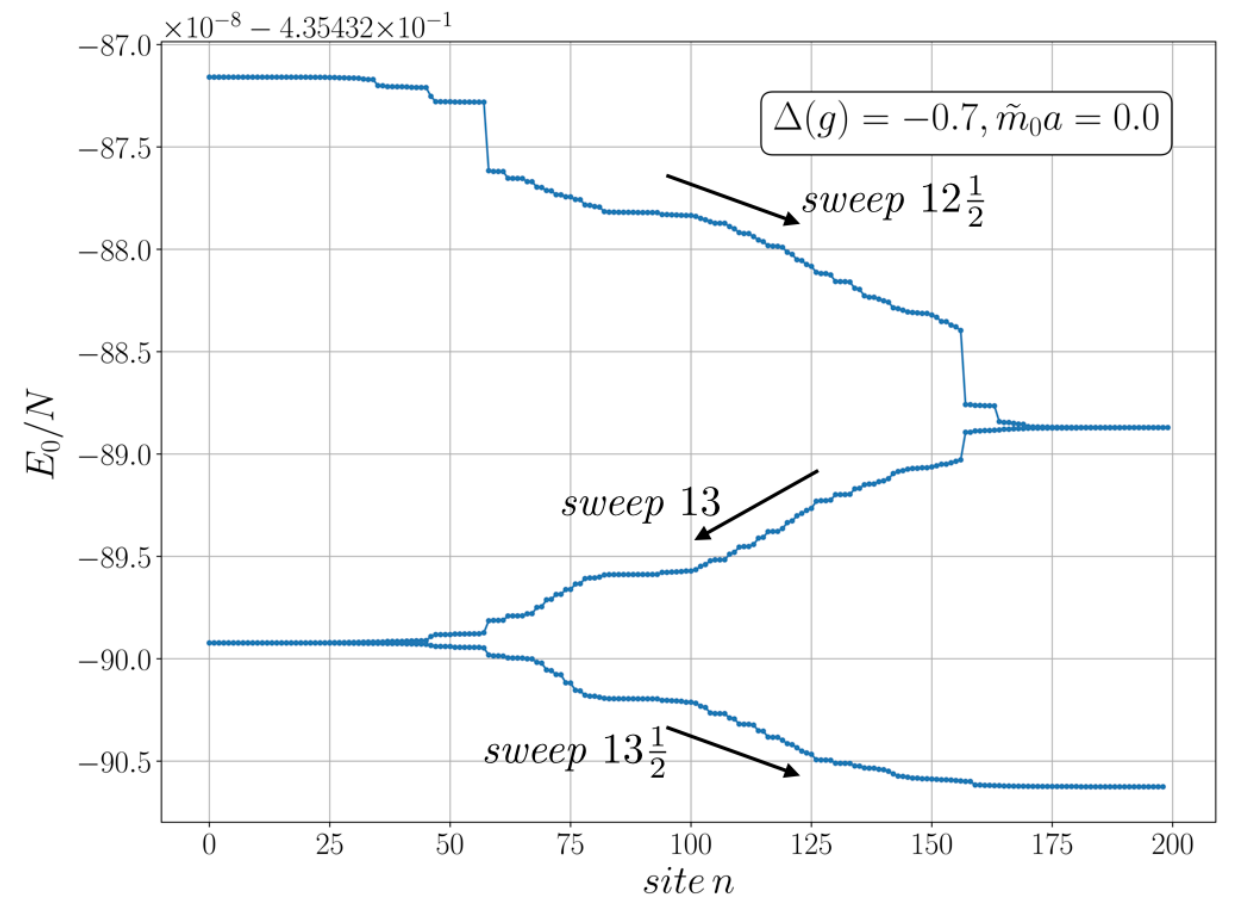
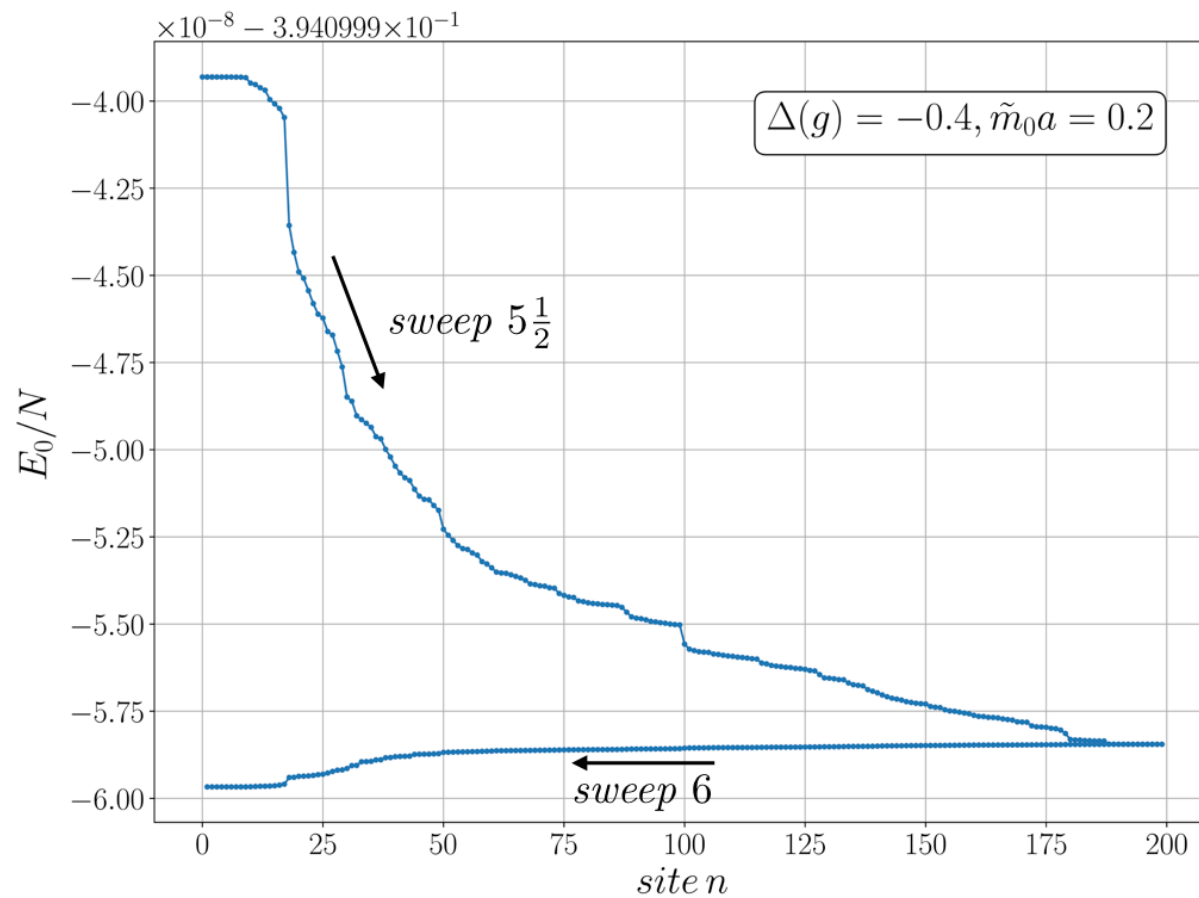
$$W^{[n]} = \begin{pmatrix} 1_{2 \times 2} & -\frac{1}{2}S^+ & -\frac{1}{2}S^- & 2\lambda S^z & \Delta S^z & \beta_n S^z + \alpha 1_{2 \times 2} \\ 0 & 0 & 0 & 0 & 0 & S^- \\ 0 & 0 & 0 & 0 & 0 & S^+ \\ 0 & 0 & 0 & 1 & 0 & S^z \\ 0 & 0 & 0 & 0 & 0 & S^z \\ 0 & 0 & 0 & 0 & 0 & 1_{2 \times 2} \end{pmatrix}$$

$$\beta_n = \Delta + (-1)^n \tilde{m}_0 a - 2\lambda S_{\text{target}}, \quad \alpha = \lambda \left(\frac{1}{4} + \frac{S_{\text{target}}^2}{N} \right) + \frac{\Delta}{4}$$

- Choices of parameters
 - ★ Twenty values of $\Delta(g)$, ranging from -0.9 to 1.0
 - ★ $\tilde{m}_0 a = 0.0, 0.1, 0.2, 0.3, 0.4$ (run 1)
 - ★ $\tilde{m}_0 a = 0.005, 0.01, 0.02, 0.03, 0.04, 0.06, 0.08, 0.13, 0.16$ (run 2)
 - ★ Bond dimension $D = 50, 100, 200, 300, 400, 500, 600$
 - ★ System size $N = 400, 600, 800, 1000$

Convergence of DMRG

- Start from random tensors at $D=50$, then go up in D
- DMRG converges fast at $\tilde{m}_0 a \neq 0$ and $\Delta(g) \gtrsim -0.7$

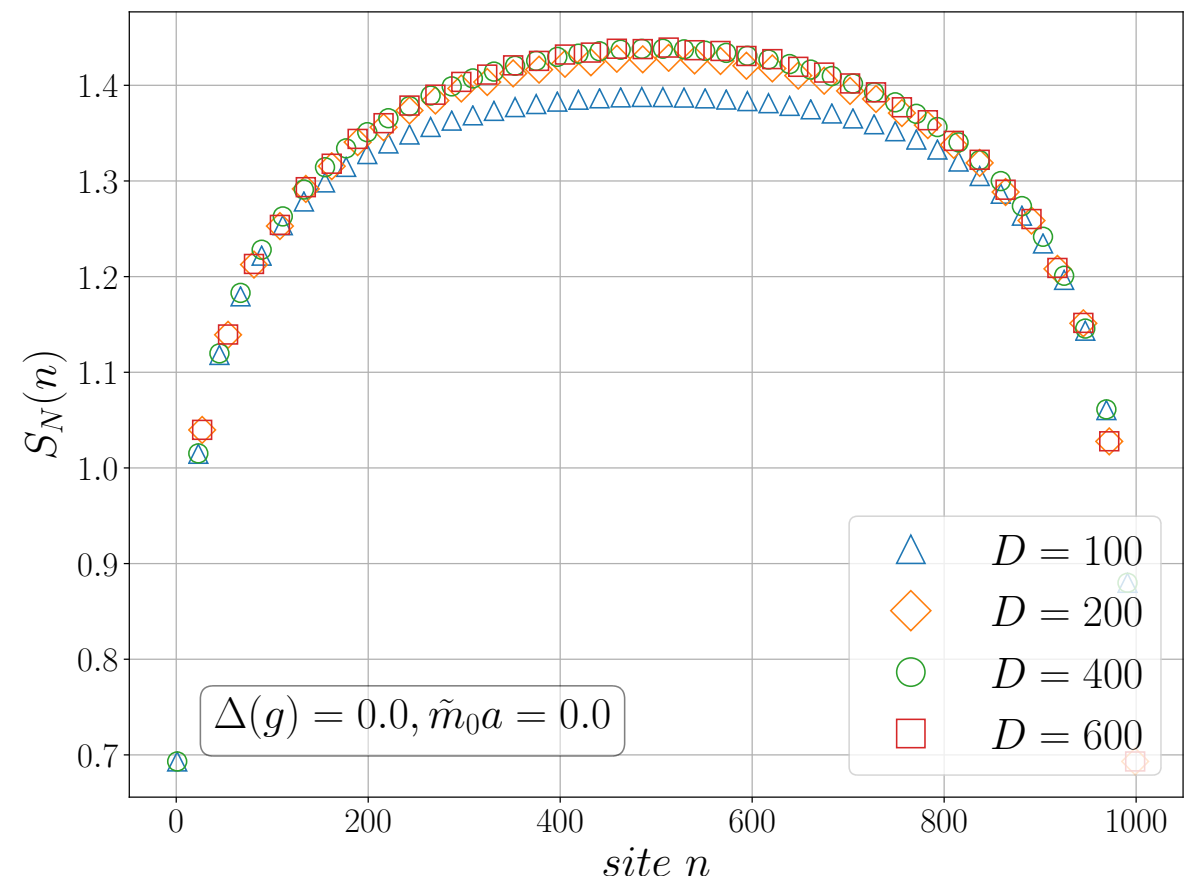
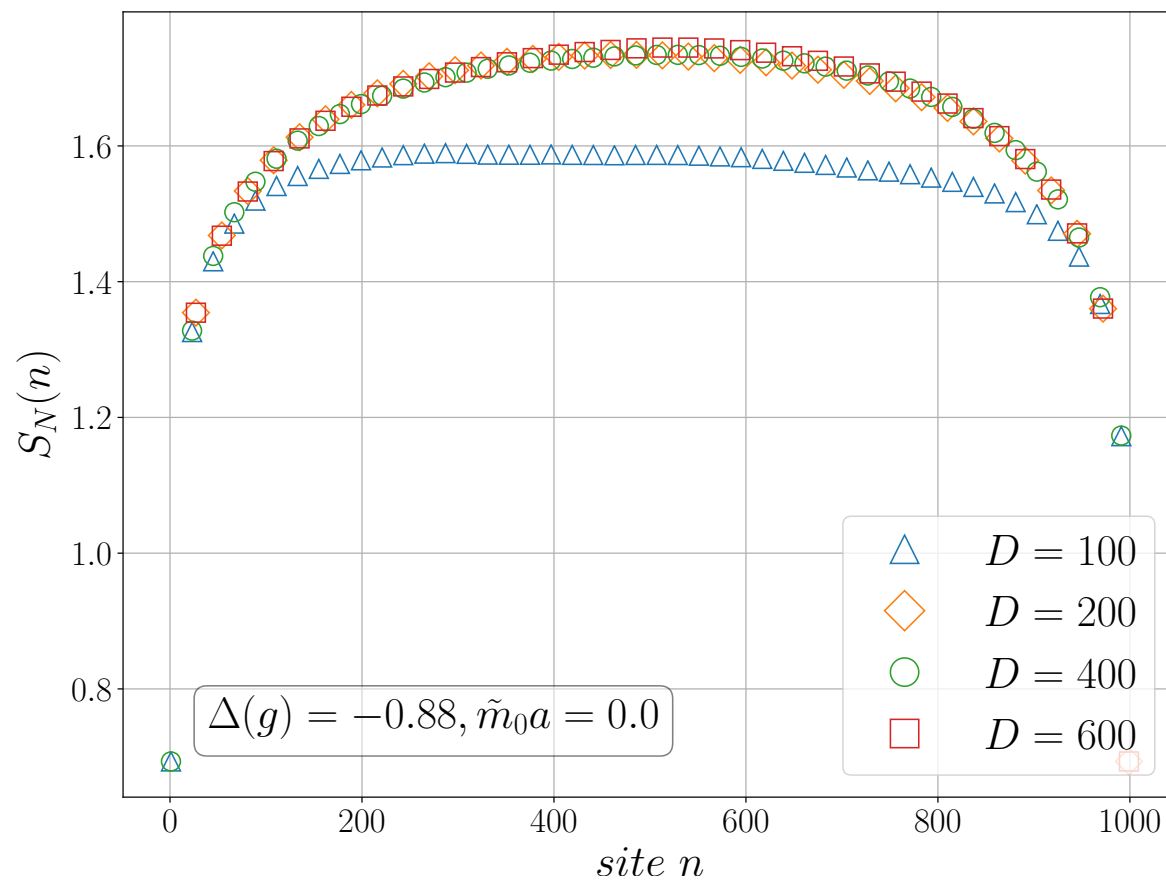


Results for the phase structure

Entanglement entropy

Calabrese-Cardy scaling and the central charge

$$S_N(n) = \frac{c}{6} \ln \left[\frac{N}{\pi} \sin \left(\frac{\pi n}{N} \right) \right] + k$$

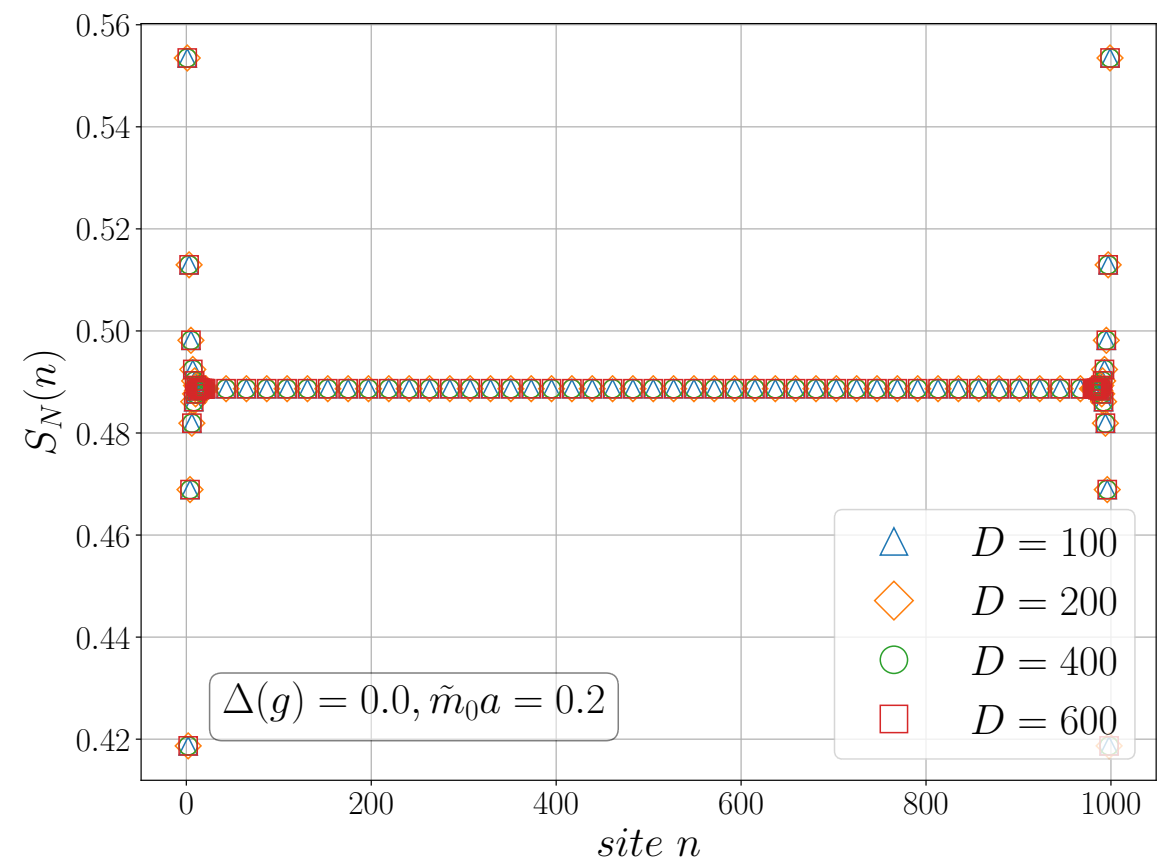
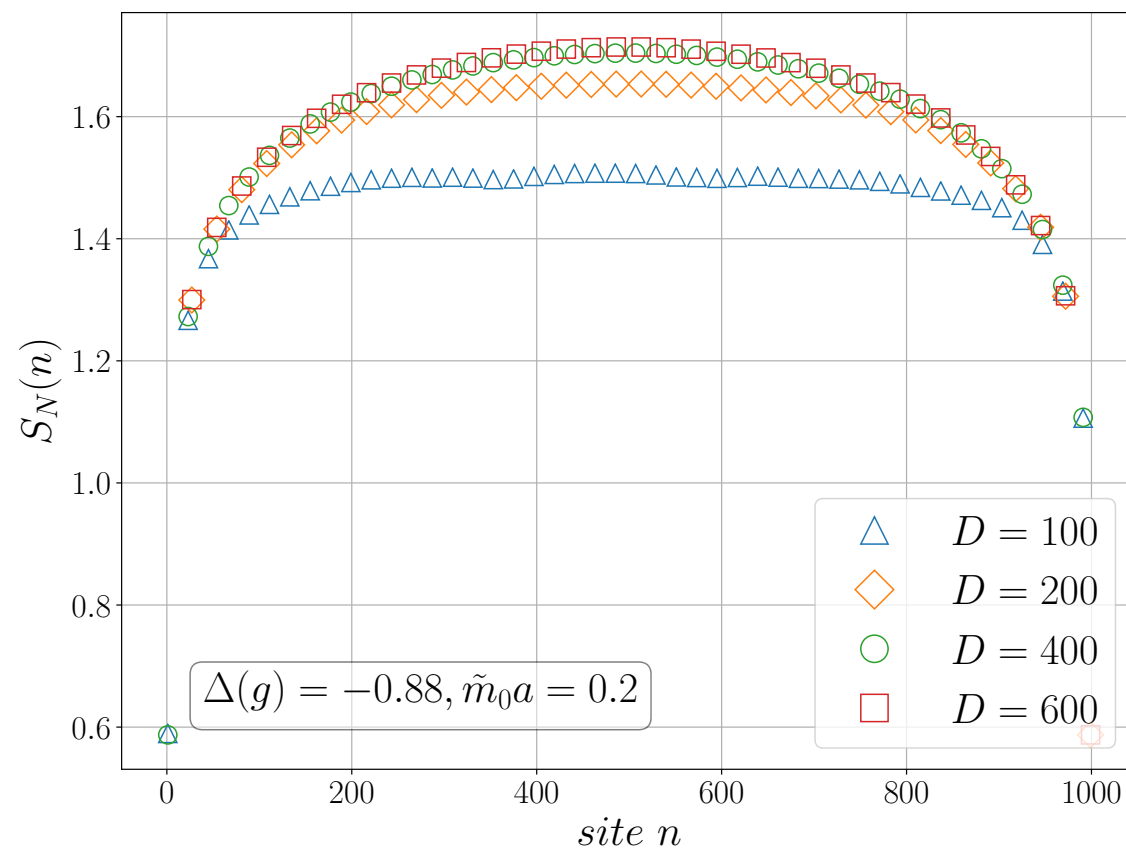


★ Calabrese-Cardy scaling observed at all values of $\Delta(g)$ for $\tilde{m}_0 a = 0$

Entanglement entropy

Calabrese-Cardy scaling and the central charge

$$S_N(n) = \frac{c}{6} \ln \left[\frac{N}{\pi} \sin \left(\frac{\pi n}{N} \right) \right] + k$$

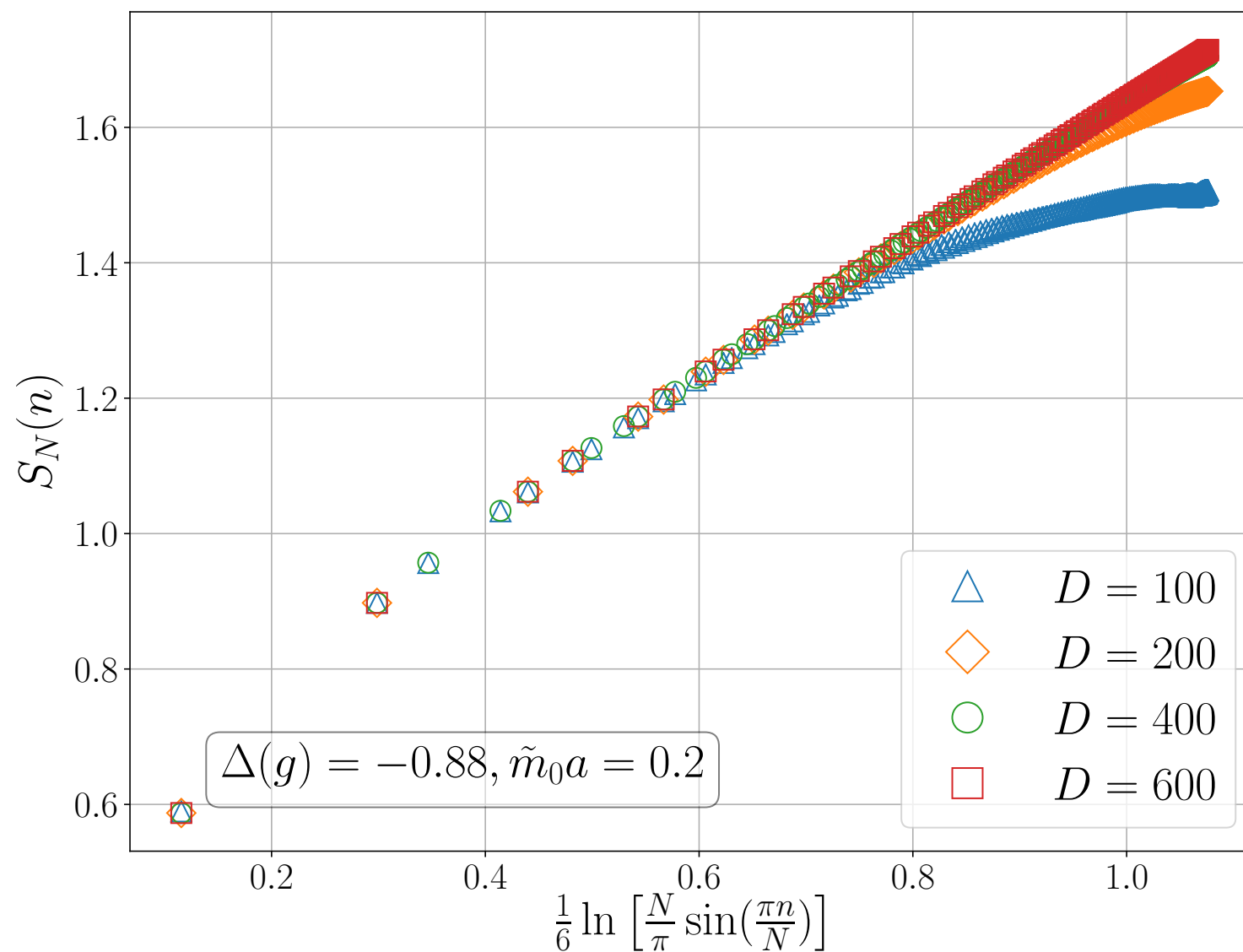


★ Calabrese-Cardy scaling observed at $\Delta(g) \lesssim -0.7$ for $\tilde{m}_0 a \neq 0$

Entanglement entropy

Calabrese-Cardy scaling and the central charge

$$S_N(n) = \frac{c}{6} \ln \left[\frac{N}{\pi} \sin \left(\frac{\pi n}{N} \right) \right] + k$$



★ Central charge is unity in the critical phase

Soliton correlators

S. Mandelstam, 1975

$$\psi_{\alpha}^{\dagger}(x)\psi_{\alpha}(y) = \mp i|2\pi(x-y)|^{-1}|c\mu(x-y)|^{-\beta^2 g^2/(2\pi)^3} \\ \times : \exp \left\{ \underbrace{-2\pi i\beta^{-1} \int_x^y d\xi \dot{\phi}(\xi)}_{\text{red}} \mp \underbrace{\frac{1}{2}i\beta [\phi(y) - \phi(x)]}_{\text{black}} + O(x-y)^2 \right\} :$$

Soliton operators

connecting vortex and anti-vortex

- ★ Power-law in the critical phase
- ★ Exponential-law in the gapped phase

Vertex operators

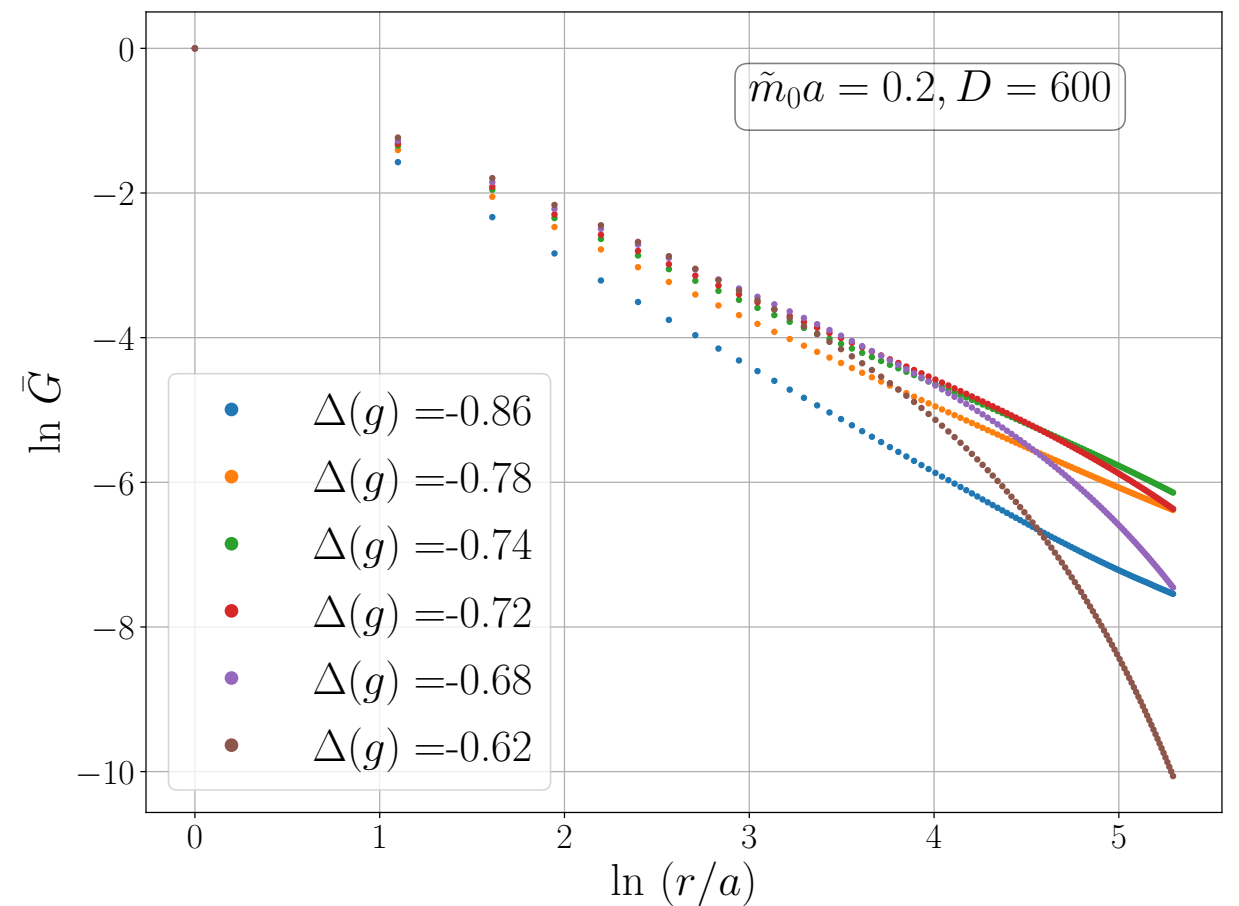
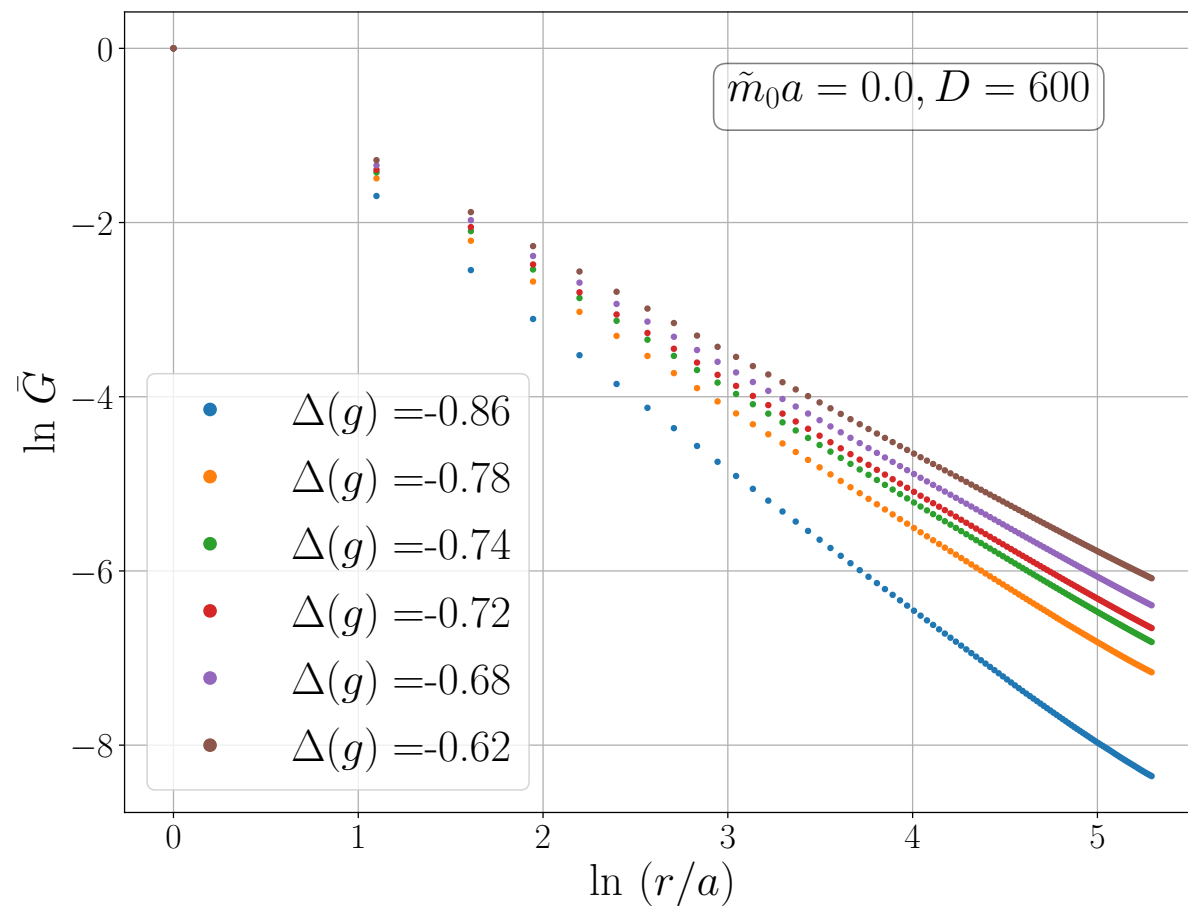
Power-law

Jordan-Wigner transformation

$$S_m^+ e^{i\pi \sum_{j=m+1}^{n-1} S_j^z} S_n^-$$

Soliton correlators

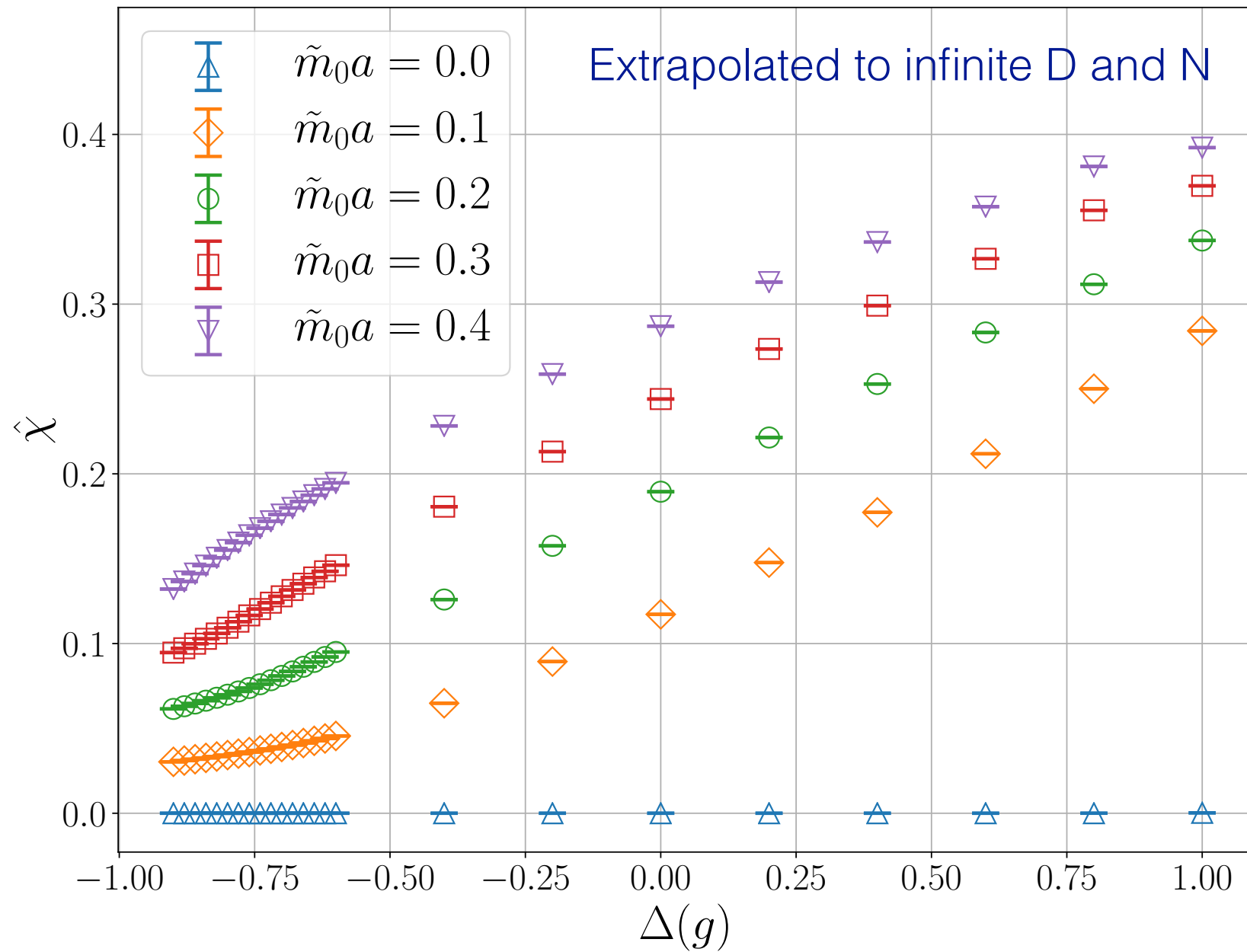
$$G(r) = \langle \psi_+^\dagger(r) \psi_+(0) \rangle, \quad \bar{G}(r) = G(r)/G(0)$$



★ Evidence for BKT phase transition

Chiral condensate

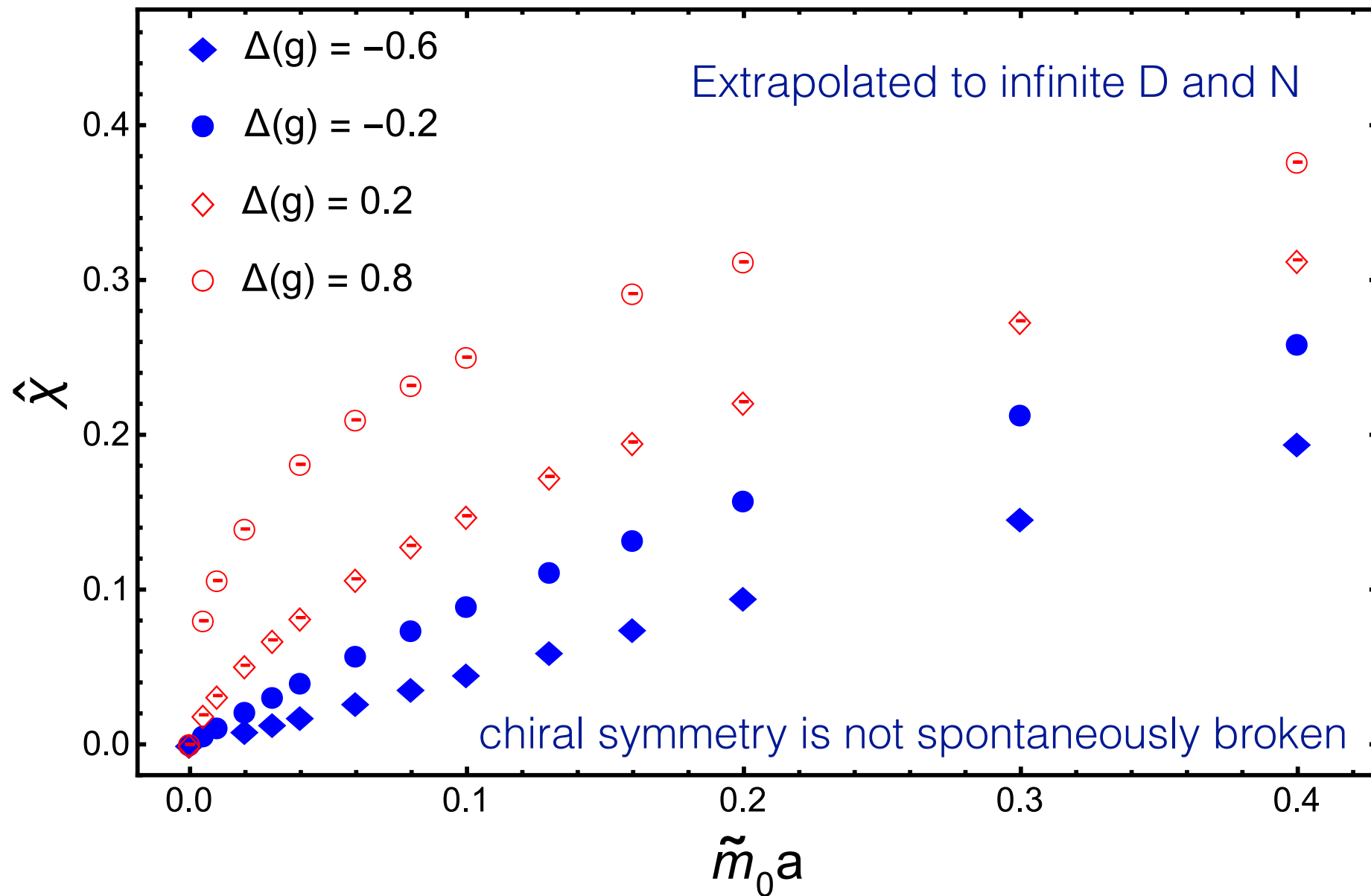
$$\hat{\chi} = a |\langle \bar{\psi} \psi \rangle| = \frac{1}{N} \left| \sum_n (-1)^n S_n^z \right|$$



★ Zero-mass results reproduced using uMPS

Chiral condensate

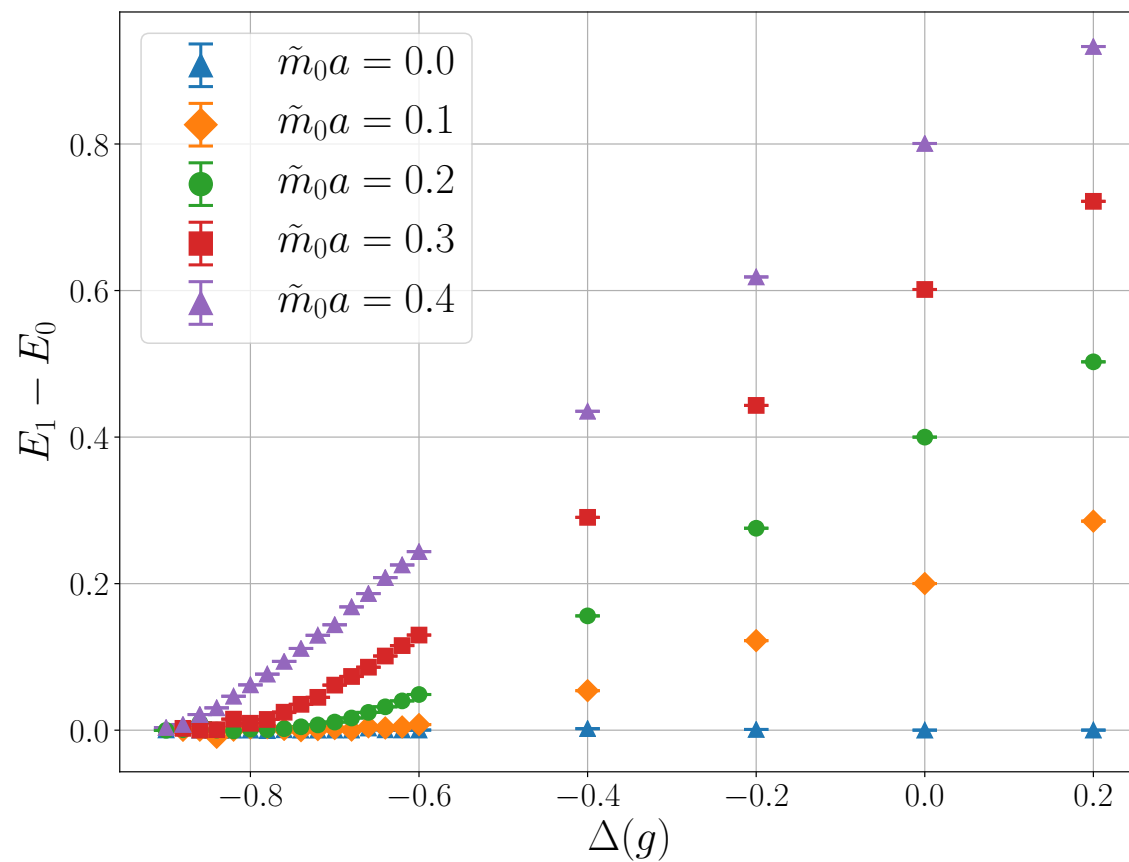
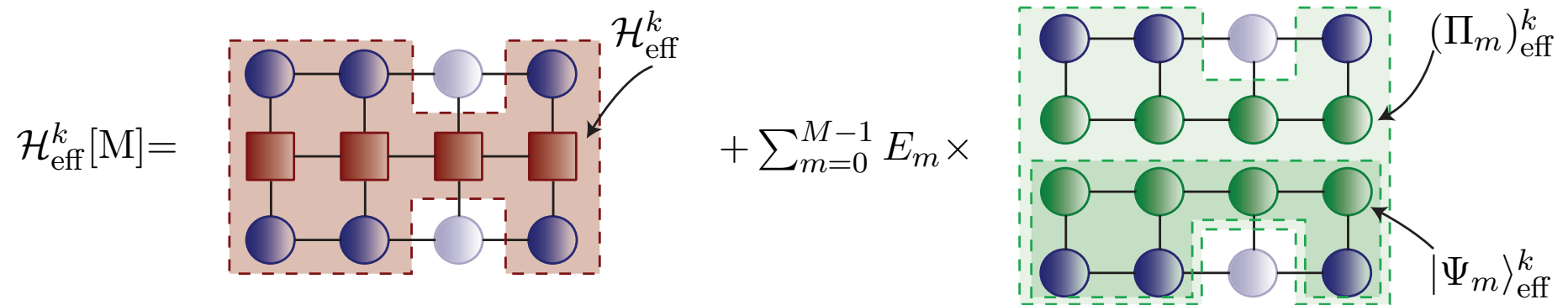
$$\hat{\chi} = a |\langle \bar{\psi} \psi \rangle| = \frac{1}{N} \left| \sum_n (-1)^n S_n^z \right|$$



★ Curvature at small mass in the gapped phase

Mass gap

$$H_{\text{eff}}[M] = \Pi_{M-1} \dots \Pi_0 H \Pi_0 \dots \Pi_{M-1} = H - \sum_{k=0}^{M-1} E_k |\Psi_k\rangle \langle \Psi_k|$$

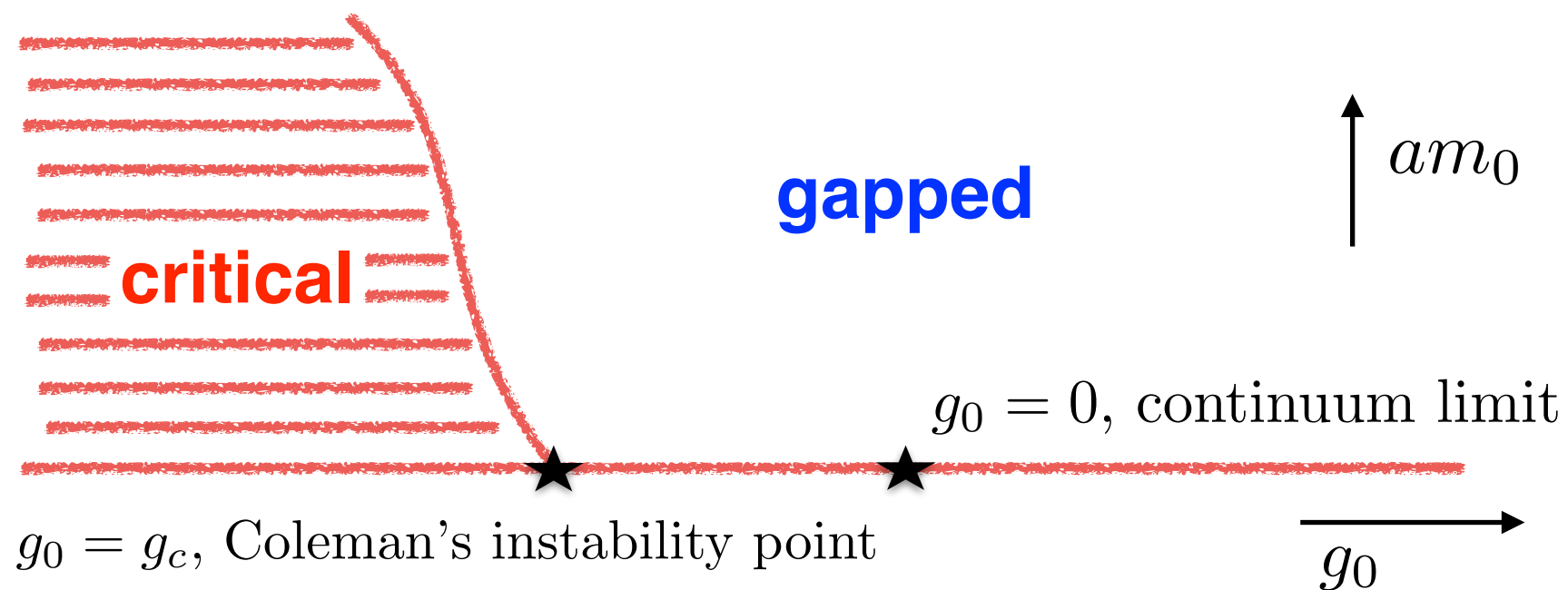


Phase structure of the Thirring model

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Massless Thirring model is a conformal field theory



Conclusion and outlook

- Evidence for BKT phase transition found using MPS
 - ★ Chiral symmetry is not spontaneously broken
- Current work for more detailed probe of the phase structure:
 - ★ More simulations at small fermion mass
 - ★ Eigenvalue spectrum of the transfer matrix
- Future projects:
 - ★ Chemical potential
 - ★ Real-time evolution with a quench