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Review on Composite Higgs Models

Oliver Witzel



Lattice 2018 East Lansing, MI, USA, July 24, 2018



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Experimental observations

- ► Discovery of the Higgs boson in 2012 [Atlas PLB716(2012)1] [CMS PLB716(2012)30]
- Higgs boson
 - $\rightarrow M_{H^0} = 125.18(16) \text{ GeV}$ [PDG 2018]
 - \rightarrow Spin 0 preferred over spin 2; spin 1 excluded ($H^0 \rightarrow \gamma \gamma)$
 - \rightarrow CP difficult to determine (mixing of e/o eigenstates)
 - \rightarrow SM decay width too small for LHC measurement
 - \rightarrow Improving precision on coupling to SM particles
- ▶ So far no other states found
 - \Rightarrow No supersymmetric particles
 - ⇒ No heavier resonances
 - \rightarrow What is the origin of the electro-weak sector?
 - → Maybe new resonances of a few TeV?



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General idea: composite Higgs models

▶ Extend the Standard Model by a new, strongly coupled gauge-fermion system

- ▶ The Higgs boson arises as bound state of this new sector
 - \rightarrow Mass and quantum numbers match experimental values when accounting for SM interactions/corrections
- ▶ System exhibits a large separation of scales
 - \rightarrow Explaining why a 125 GeV Higgs boson but no other states have been found
 - \rightarrow Indications that such a system cannot be QCD-like (e.g. quark mass generation)
 - \rightsquigarrow near-conformal gauge theories
- Exhibits mechanism to generate masses for SM fermions and gauge bosons
- ▶ In agreement with electro-weak precision constraints (e.g. S-parameter)?

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Composite Higgs models

Aim: describe states of the SM as well as particles originating from new physics

$$\mathcal{L}_{\textit{UV}}
ightarrow \mathcal{L}_{\textit{SD}} + \mathcal{L}_{\textit{SM}_0} + \mathcal{L}_{\textit{int}}
ightarrow \mathcal{L}_{\textit{SM}} + \dots$$

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Composite Higgs models

▶ Aim: describe states of the SM as well as particles originating from new physics

 \blacktriangleright Start with a Higgs-less, massless SM

$$\mathcal{L}_{UV} \rightarrow \mathcal{L}_{SD} + \mathcal{L}_{SM_0} + \mathcal{L}_{int} \rightarrow \mathcal{L}_{SM} + \dots$$

Composite Higgs models

- Aim: describe states of the SM as well as particles originating from new physics
- ► Start with a Higgs-less, massless SM
- ► Add new strong dynamics coupled to SM

$$\mathcal{L}_{UV} \to \mathcal{L}_{SD} + \mathcal{L}_{SM_0} + \mathcal{L}_{int} \to \mathcal{L}_{SM} + \dots$$

full SM + states from \mathcal{L}_{SD}

- ▶ Leads to an effective theory giving mass to
 - \rightarrow the SM gauge fields
 - \rightarrow the SM fermions fields: 4-fermion interaction or partial compositeness

Composite Higgs models

- Aim: describe states of the SM as well as particles originating from new physics
- ▶ Start with a Higgs-less, massless SM
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$$\mathcal{L}_{UV} \to \mathcal{L}_{SD} + \mathcal{L}_{SM_0} + \mathcal{L}_{int} \to \mathcal{L}_{SM} + \dots$$
full SM + states from \mathcal{L}_{SD}

► Leads to an effective theory giving mass to

 \rightarrow the SM gauge fields

 \rightarrow the SM fermions fields: 4-fermion interaction or partial compositeness

▶ Does not explain mass of \mathcal{L}_{SD} fermions and 4-fermion interactions: \mathcal{L}_{UV}

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Two scenarios for a composite Higgs

- Light iso-singlet scalar (0^{++})
 - \rightarrow "Dilaton-like"
 - ightarrow Scale: ${\it F}_{\pi}=$ SM vev \sim 246 GeV
 - \rightarrow ideal 2 massless flavors
 - \Rightarrow giving rise to 3 Goldstone bosons
 - \Rightarrow longitudinal components of W^{\pm} and Z^0
- 2-flavor sextet [LatHC, CP3] (Kuti Wed 2:00 PM, Wong Wed 2:20 PM)
- 8-flavor fundamental [LatKMI, LSD] (Rebbi Thu 11:00 AM, Neil Thu 12:00 PM)
- 2-flavor fundamental [Drach et al.] see appendix

- ▶ pseudo Nambu Goldstone Boson (pNGB)
 - → Spontaneous breaking of flavor symmetry $\Rightarrow N_f > 3$
 - \rightarrow Mass emerges from its interactions
 - \rightarrow Non-trivial vacuum alignment

 $\mathit{F}_{\pi} = (\mathsf{SM} \; \mathsf{vev}) / \sin(\chi) > 246 \; \mathsf{GeV}$

- Two-representation model by Ferretti [TACoS] (Jay Thu 12:20 PM)
- ▶ Mass-split models [4+8, LSD]
- ► SU(4)/Sp(4) composite Higgs [Bennett et al.] (Lee Tue 2:20 PM) see appendix

near-conformal gauge theories

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Near-conformal gauge theories

- ▶ Gauge-fermion system with $N_c \ge 2$ colors and N_f flavors
 - in some representation



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Near-conformal gauge theories

- ▶ Gauge-fermion system with $N_c \ge 2$ colors and N_f flavors
 - in some representation



summarv

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Conformal window

- ► Indications of the conformal window for different representations, N_c, and N_f [Dietrich, Sannino PRD75(2007)085018]
- Derived from perturbative and Schwinger-Dyson arguments
- Lower bonds of conformal window typically require nonperturbative calculations



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SU(2) gauge theory with fermions in the adjoint representation

- Nonperturbative investigations include e.g.
 - \rightarrow Scaling of hadron masses
 - \rightarrow Mode number of Dirac operator
 - \Rightarrow Determination of the anomalous dimension
- Conclusions
 - $\rightarrow N_f = 2$ is conformal [Bergner et al. PRD96(2017)034504]
 - \rightarrow $N_f=1$ likely conformal [Athenodorou et al. PRD91(2015)114508]
 - \rightarrow \textit{N}_{f} = 1/2 (1 Majorana fermion) is QCD-like

[Bergner et al. JHEP03(2016)080]

→ $N_f = 3/2$ (3 Majorana fermions) is conformal Mode number: $\gamma^* \approx 0.38(2)$; fit spectrum $\gamma^* \approx 0.37(2)$

► Mixed fundamental-adjoint action (Bergner Fri 5:10 PM) → Investigations of supersymmetric QCD

(Scior poster)



[Bergner et al. JHEP01(2018)119]



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The step-scaling β -function

- \blacktriangleright IRFP: β function has zero for $g^2>0$
- For large g^2 nonperturbative methods are required
- **•** Calculate discretized β function (step scaling)
 - \rightarrow Requires calculations on a set of different volumes
 - \rightarrow Well established in QCD [Lüscher et al. NPB359(1991)221]
- ► Gradient flow step scaling [Lüscher JHEP08(2010)071] [Fodor et al. JHEP11(2012)007][Fodor et al. JHEP09(2014)018]



[Dalla Brida et al. PRD95(2017)014507]

$$g_c^2(L) = rac{128\pi^2}{3(N_c^2 - 1)} \; rac{1}{C(c,L)} t^2 \langle E(t)
angle \quad ext{with } \sqrt{8t} = c \cdot L; \qquad eta_s^c(g_c^2;L) = rac{g_c^2(sL) - g^2(L)}{\log(s^2)}$$

• Extrapolate $L \rightarrow \infty$ to remove discretization effects and take the continuum limit • Expect to find agreement for results based on different actions, flows, operators ...



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The challenge of establishing an IRFP







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Active research

- \blacktriangleright Larger volumes might be required for $L \rightarrow \infty$ extrapolation
 - \rightarrow Small c-values certainly require larger volumes than larger c-values
 - ---- Larger c-values have larger statistical uncertainties
 - \rightarrow Different actions have different discretization errors
- **•** DWF with Symanzik gauge action feature a fully $O(a^2)$ improved set-up à la Symanzik
 - \rightarrow Zeuthen flow [Ramos, Sint EPJC76(2016)15]
 - \rightarrow Symanzik operator
 - \rightarrow Perturbative tree-level normalization [Fodor et al. JHEP09(2014)018] works for N_f = 12 and 10
- \rightarrow Perturbative improvement breaks down for staggered with $N_f = 8$ [Fodor et al JHEP06(2015)19]
- ▶ [Rooted] Staggered Fermions: Good, Bad or Ugly? [Sharpe Plenary Lattice 2006] → Are staggered and DW/Wilson fermions in conformal systems in the same universality class?

(Hasenfratz poster, Kuti Wed 2:00 PM, Holland Wed 2:40 PM, Nogradi Wed 3:00 PM)

Higgs as a light 0^{++} scalar



SU(3) with $N_f = 2$ sextet flavors (two-index symmetric representation)

- Minimal flavor content to describe EW symmetry breaking
- Likely very close to the onset of the conformal window

- ▶ LatHC
 - \rightarrow Chirally broken spectrum [Fodor et al. PLB718(2012)657]
 - \rightarrow 0⁺⁺ (f_0) is light
 - → No IRFP in the explored range of the β -function [Fodor et al. JHEP09(2015)039]

(Kuti Wed 2:00 PM, Wong Wed 2:20 PM)





SU(3) with $N_f = 2$ sextet flavors (two-index symmetric representation)

- Minimal flavor content to describe EW symmetry breaking
- Likely very close to the onset of the conformal window

- ▶ Hansen, Drach, Pica
 - → Two different phases one chirally broken, one looking IR conformal [Hansen, Drach, Pica PRD96(2017)0345]
- ▶ Hasenfratz, Liu, Yu-Han Huang
 - \rightarrow Indications for a possible IRFP [Hasenfratz et al. 1507.08260]





SU(3) with $N_f = 8$ fundamental flavors

- ▶ Theory considered to be chirally broken but close to the onset of the conformal window
 - \rightarrow Step scaling analysis of the discrete β function

[Hasenfratz et al. JHEP06(2015)143][Fodor et al. JHEP06(2015)019]

 \rightarrow Finite temperature phase diagram

[Deuzeman et al. PLB670(2008)41][Jin, Mawhinney PoS LATTICE2010 055][Schaich et al. PoS LATTICE2012 028]

\rightarrow Studies of the low-lying meson spectrum

[Aoki et al. PRD89(2014)111502][Appelquist et al. PRD93(2016)114514][Aoki et al. PRD96(2017)014508] [Appelquist et al. 1807.08411]

- ► Theory has 63 Goldstone boson not an ideal candidate to explain EW symmetry breaking
 - \rightarrow Allows to investigate qualitative features of near-conformal gauge theories
 - \rightarrow Reduce number of light Goldstones by assigning e.g. mass or charge to some flavors
- $ightarrow 0^{++}$ is light, degenerate with the pion!

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Determining the iso-singlet scalar 0^{++}

- ▶ Receives quark line connected and disconnected contributions
 - \rightarrow Stochastic estimator (noisy)
- Its quantum numbers are the same vacuum
 - \rightarrow Large vacuum subtraction
- ▶ Lighter in near-conformal systems than in QCD
 - \rightarrow Easier to determine, stable particle
 - (energetically protected from decaying)
- Nevertheless most expensive state in the spectrum
- Idea: take advantage of correlators at non-zero momenta to avoid the vacuum subtraction (Rebbi Thu 11:00 PM)





[Appelquist et al. PRD93(2016)114514]

▶ 0⁺⁺ is light, degenerate with the pion $\Rightarrow \chi PT$ not applicable



Comparison of SU(3) with $N_f = 4$ and 8 fundamental flavors



- Wilson flow scale $\sqrt{8t_0}$ vs fermion mass am_f
- Strong mass-dependence for $N_f = 8$
 - ⇒ Show quantities in units of $\sqrt{8t_0}$ or dimensionless ratios



Spectrum in units of F_π
 N_f = 8: pion and σ (0⁺⁺) degenerate rather different than in N_f = 4

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Effective field theories (a selection of recent work)

Holdom, Koniuk: A bound state model for a light scalar [Holdom, Koniuk JHEP12(2017)102]

- \rightarrow Existence of light scalar, well separated from heavier states related to existence of near conformal gauge dynamics extending over a wide range
- \rightarrow Scalar mass and form factor close to parity doubled limit
- \rightarrow Light scalar has characteristics different from light dilaton

Appelquist, Ingoldby, Piai: Dilaton EFT [Appelquist, Ingoldby, Piai JHEP07(2017)035][JHEP03(2018)039]

- \rightarrow Light singlet scalar interpreted as dilaton (spontaneous breaking of conformal symmetry)
- \rightarrow Treat dilaton together with pions (spontaneous breaking of chiral symmetry)
- \rightarrow Add general form for the dilaton potential to be determined from lattice data
- \rightarrow EFT "fits" lattice data ($N_f = 8$ fundamental and $N_f = 2$ sextet)

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Effective field theories (a selection of recent work)

Golterman, Shamir: The large-mass regime of the dilaton-pion low-energy effective theory [Golterman, Shamir 1805.00198]

- ▶ Investigate dilaton-pion EFT in the Veneziano limit ($N_f \rightarrow \infty$ for N_f/N_c fixed)
- Expand around $n_f n_f^*$, with $n_f = N_f/N_c$ and n_f^* onset of conformal window in Veneziano limit
 - \rightarrow Small mass region: dilaton decouples from pions, typical chiral behavior
 - \rightarrow Large mass region: hadron masses, decay constants scale with $m_f^{1/(1+\gamma^*)}$ (hyperscaling)
 - \Rightarrow LSD $N_f = 8$ data is in the large mass region
 - \rightarrow Explains characteristics of LSD data
 - \rightarrow To reach small mass region: reduce $m_f
 ightarrow m_f/100$

 \rightarrow Small mass region may show that $N_f = 8$ is indeed confining (Golterman Thu 11:20 PM)



C 0.2

0.0

-0.2

- \Rightarrow Effect on spectrum and onset of conformal window
- \rightarrow Tree-level spectrum leads to dimensionless ratios:

$$R_{\sigma} = (M_{\sigma}^2 - M_{\pi}^2)/M_{\eta'}^2 + (1 - 2/N_f)(1 - M_{\pi}^2/M_{\eta'}^2)$$

$$R_{a_0} = (M_{a_0}^2 - M_{\pi}^2)/M_{\eta'}^2 - (2/N_f)(1 - M_{\pi}^2/M_{\eta'}^2)$$

 \rightarrow LatKMI data: almost flat, no N_f dependence \rightsquigarrow e.g. bound on N_{fc}



- Consider flavors with two masses m_1 and $m_2 = m_1 + \delta_m$ (δ_m small)
 - \rightarrow Spectrum exhibits light-light, heavy-light, and heavy-heavy mesons

→ If $M_{\pi II}^2 < M_{\pi hI}^2 < M_{\pi hh}^2$, then inverse ordering for scalars $M_{a_0II}^2 > M_{a_0hI}^2 > M_{a_0hI}^2 > M_{a_0hI}^2$ [Floor, Gustafson, Meurice, 1807.05047] (Gustafson poster)

0.4

0.5

0.6

0.3

 M_{π}/M_{π}

0.2

0.1

Higgs as a pseudo Nambu-Goldstone boson

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Ferretti's Model [Ferretti JHEP06(2014)142]

- \blacktriangleright SU(4) gauge theory with fermions in two representations
 - \rightarrow $N_6^{\cal W}=5$ Weyl massless flavors of sextet (Q) (two-index antisymmetric) fermions with EW charge
 - \rightarrow \textit{N}_{4} = 3 fundamental Dirac flavors (q) with color charge
- Mesons
 - \rightarrow sextet QQ, $Q\overline{Q}$, $\overline{Q}\overline{Q}$ pNGBs, vectors
 - \rightarrow fundamental $q\bar{q}$ pNGBs, vectors

- Baryons
 - ▶ sextet *QQQQQQ* bosons
 - ▶ fundamental *qqqq* bosons
 - chimera Qqq fermions
- Ferretti limit ($m_6 \rightarrow 0$) Higgs is a massless sextet NGB, its potential arises from SM interactions
- Fermion acquire mass from quartic mixing $u\bar{u}H \rightarrow u\bar{u}QQ$
- ▶ Non-anomalous superposition of $U_{A(4)}(1)$ and $U_{A(6)}(1) \rightarrow \text{axial singlet pNGB}$ (ζ meson)
- ▶ top quark mixes linearly with chimera \Rightarrow large m_t



Adaption of Ferretti's model on the lattice

[Ayyar et al. PRD97(2018)074505][PRD97(2018)114502][PRD97(2018)114505]

- ▶ SU(4) gauge theory
 - $\rightarrow N_6 = 2$ Dirac flavors ($N_6^4 = 4$ Weyl) sextet flavors
 - $\rightarrow N_4 = 2$ fundamental Dirac flavors
- Finite temperature phase diagram: only two phases
 Low-temperature

both fermion species confined and chirally broken High-temperature

both fermion species deconfined and chirally restored

 Single phase transition appears to be first order as theoretically predicted



[Ayyar et al. PRD97(2018)114502]



Results in the Ferretti limit $(m_6 \rightarrow 0)$

$\triangleright \zeta$ meson

- $\rightarrow M_{P56} = 0 \text{ (sextet pNGB exactly massless)}$ $\rightarrow M_{C} < M_{P54}$
- $\rightarrow \zeta$ meson lightest, massive state
- → Reconstruct M_{ζ} from chiral fit (function of m_4 and m_6)





Results in the Ferretti limit $(m_6 \rightarrow 0)$

- Spectrum in units of F_6
 - \rightarrow (1/2,0) chimera (Qqq) is top partner and lightest baryon
 - $_{\rightarrow}$ Experimental constraint: $\mathit{F}_{6} \gtrsim 1.1 \text{ TeV}$
 - \Rightarrow Mass of top partner chimera $M\gtrsim 6.5~{\rm TeV}$

▶ Further details (Jay Thu 12:20 PM)



[[]Ayyar et al. PRD97(2018)114505]

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Mass-split models

▶ Promising candidates are chirally broken in the IR but conformal in the UV

[Luty, Okui JHEP09(2006)070], [Dietrich, Sannino PRD75(2007)085018], [Vecchi 1506.00623], [Ferretti, Karateev JHEP03(2014)077]



Mass-split models e.g. SU(3) gauge theory with "heavy" and "light" (massless) fundamental flavors
 Add N_h heavy flavors to push the system ► N_ℓ = 4 light flavors are chirally broken in the IR near an IRFP of a conformal theory

heavy flavors could be invisible to SM

fundamental composite 2HDM with 4 flavors in SU(3) gauge [Ma, Cacciapaglia JHEP03(2016)211]

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The mass-split paradigm

▶ In QCD: $g^2 \rightarrow 0$ (continuum limit); fermion mass $m_f \rightarrow 0$ (chiral limit)

- ▶ Theory with degenerate $N_f = N_h + N_\ell$ is (mass-deformed) conformal and exhibits an IRFP
 - All ratios of hadron masses scale with the anomalous dimension (hyperscaling)
 - \rightarrow Continuum limit is taken by sending fermion mass $m_f \rightarrow 0$

▶ Mass-split models live in the basin of attraction of the IRFP of N_f degenerate flavors

- \rightarrow Inherit hyperscaling of ratios of hadron masses but are chirally broken
- $_{
 m
 m \rightarrow}$ Continuum limit: $m_h
 m
 m
 m
 m 0$ keeping m_ℓ/m_h fixed
- ightarrow Chiral limit: $m_\ell
 ightarrow$ 0 i.e. $m_\ell/m_h
 ightarrow$ 0
- \rightarrow Gauge coupling is irrelevant
- \rightarrow No free parameters after taking the chiral and continuum limit, but light-light, heavy-light, and heavy-heavy bound states

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Results for four light and eight heavy flavors

▶ Hyperscaling in the light-light sector



► $M_n/F_\pi \approx 11$ ► $M_\varrho/F_\pi \approx 8$ ► $M_{0^{++}}/F_\pi \approx 4-5$ → taking the chiral limit is difficult but 0⁺⁺ well separated from the ϱ and degenerate with the pion

Statistical errors only

 "Scatter" indicates corrections to scaling

▶ Gauge coupling is irrelevant

[Brower et al. PRD 93 (2016) 075028]



Results for four light and eight heavy flavors



- ▶ 4+8 heavy-heavy spectrum is not QCD-like; QCD is not hyperscaling
- M^{hh}/F_{π} increases but F_{π} is finite in the chiral limit
- ▶ $M_{\rho}^{hh} \sim 3M_{\varrho} \Rightarrow$ could be accessible at the LHC
- ▶ Data at $\beta = 4.0$ and 4.4: gauge coupling is irrelevant

[Hasenfratz, Rebbi, Witzel PLB773(2017)86]

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Results for four light and eight heavy flavors

▶ The system is chirally broken



- All data points in a_★ units
 a_★F_π is finite
- ► Linearity in M_{π}^2 for small m_{ℓ} ► $N_f = 4$ (QCD-like): ratio diverges ► QCD: $m_d/m_s = 4.7/96 \approx 0.05$ ► $N_f = 12$: almost constant ratio [Cheng at al. PRD90(2014)014509]

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Outlook

- ▶ Mass-split models using 4 light and 6 heavy flavors of MDWF
 - \rightarrow If degenerate $\mathit{N_f}=10$ is conformal, expect to see hyperscaling
 - \rightarrow First data with eventually large systematics look promising
 - \rightarrow \textit{N}_{f} = 10 would have larger anomalous dimension
- Simpler to calculate phenomenologically interesting quantities
 - \rightarrow Generation of mass for SM fermions (partial compositeness, four-fermion interaction
 - → Baryon anomalous dimension e.g. via new gradient flow method [Carosso et al. 1806.01385] (Hasenfratz Fri 4:50 PM)
 - \rightarrow S-parameter [Appelquist et al. PRL106(2011)231601], Higgs-potential, ...



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Summary

- ▶ The experiments will tell us whether the Higgs is a composite particle
 - \rightarrow Performing nonperturbative simulations we can guide experimentalists and model builders
 - \rightarrow Even (old) QCD calculations can be useful (DeGrand Thu 11:40 AM)
- ► Proposal of a new, alternative Higgs mechanism based on dynamical mass generation (Garofalo Mon 3:00 PM and Frezzotti Mon 3:20 PM), appendix
- ► Simulating near-conformal systems is more costly than QCD but can be as controversial → Particular challenge: identifying an IRFP at strong coupling
- ▶ Simulations of near-conformal systems revealed a light 0⁺⁺ with mass $M_{0^{++}} \sim M_{\pi}$ → Different effective field theories are required/explored
- ► Models based on two representations or mass-split systems exhibit novel features → E.g. chimera baryons, hyperscaling in a chirally broken system

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Acknowledgments

Step-scaling

A. Hasenfratz, K. Holland, J. Kuti, D. Nogradi $N_f = 8$

A. Hasenfratz, G.T. Fleming, E.T. Neil, E. Rinaldi, C. Rebbi, D. Schaich EFTs Fundar

T. Appelquist, M. Golterman, R. Koniuk, Y. Meurice Ferretti model

T. DeGrand, D.C. Hackett, W.I. Jay, E.T. Neil

SU(2) adjoint

G. Bergner

Dynamical mass generation

M. Garofalo, R. Frezzotti

Fundamental Higgs A. Maas Mass-split models A. Hasenfratz, C. Rebbi SU(2) fundamental V. Drach SU(4)/Sp(4) J.W. Lee



$N_f = 12$ step-scaling



- Hasenfratz, Schaich nHYP-smeared staggered,
 Wilson gauge w/ adjoint term
 Wilson flow, clover operator
- Lin, Ogawa, Ramos unimproved staggered, Wilson gauge, twisted BC Wilson flow, clover operator

Fodor et al.

3× stout smeared staggered Symanzik gauge Symanzik flow, clover operator Newer results discussed in a moment!

$N_f = 12$ step-scaling



(Hasenfratz poster)

- ▶ Hasenfratz, Rebbi, Witzel
 - Möbius domain-wall fermions, Symanzik gauge Zeuthen flow, Symanzik operator
 - → Perturbative tree-level normalization [Fodor et al. JHEP09(2014)018]

(working for $N_f = 12$ and 10)

- Result robust
 - \rightarrow Alternative flow/operators
 - \rightarrow Without tree-level normalization
 - \rightarrow Alternative $L \rightarrow \infty$ extrapolation
 - \rightarrow Changing scheme, e.g., c = 0.3

$N_f = 12$ step-scaling

▶ Page provided by Julius Kuti



 $(g^{2}(sL) - g^{2}(L))/log(s^{2}) = c_{-} + c_{-} \cdot a^{2}/L^{2} + a^{4}/L^{4}$ terr $c_{o} = 0.143 \pm 0.038$ nf=12 new c = -98.1 + 39v²/dof= 0.39 cu0.2 target F continuum limit s₁= 0.118 + 8-01 Autor 3 3.5 4.5 ×10⁻³ 202

LatHC PLB B779 (2018) 230-236 arXiv:1710.09262 confirmed with new updated results:

L=32 -> L=64 step at several targets adds evidence against nf=12 IRFP

Talk: J. Kuti Wed, 14:00 BSM room 104

Dynamical generation of elementary particle masses — an alternative to the Higgs mechanism

[Dimopoulos, Frezzotti, Garofalo, Kostrzewa, Pittler, Rossi, Urbach]

- Non-abelian strongly interacting fermions coupled via Yukawa couplings to a scalar field and a Wilson-like term
 - \rightarrow Exact symmetry acting on fermions and scalars prevents power divergent fermion mass terms
 - \rightarrow Fermionic chiral invariance broken by Yukawa and Wilson-like term,
 - but restored at critical Yukawa coupling
- ► Scalar field with double-well potential
 - \rightarrow Left-over breaking of chiral symmetry at cutoff scale polarizes vacuum
 - \Rightarrow spontaneous chiral symmetry breaking dynamically generates PCAC fermion mass
- Dynamical fermion mass can be naturally "small" and fermion masses exhibit natural hierarchy
 Higgs boson is a composite state in WW + ZZ channel bound by new strongly interacting particles

Dynamical generation of elementary particle masses — an alternative to the Higgs mechanism

 Dynamical generation of fermion mass demonstrated by numerical simulations (Garofalo Mon 3:00 PM)

▶ Electro-weak interactions and how electro-weak boson acquire mass by this mechanism

- \rightarrow Dynamical EW symmetry breaking due to a super-strongly sector
- \rightarrow No "unnatural" fine tuning of effective four-fermion coupling

(Frezzotti Mon 3:20 PM)

Effects of a fundamental Higgs

[Maas, Törek]

- ▶ Physical spectrum must be gauge invariant (in QCD guaranteed by confinement)
- ▶ Weak sector: perturbative description BRST-invariant, but gauge dependent
 - → Experimental results match predictions due to the Fröhlich-Morchio-Strocchi (FMS) mechanism [Fröhlich, Morchio, Strocchi PLB97 (1981)249][NPB160(1981)553]
 - \rightarrow SM: weak gauge group matches global custodial symmetry
 - \rightarrow Not guaranteed for BSM models
 - \Rightarrow discrepancy between physical and elementary spectrum
- ► Investigate SU(3) gauge theory with a fundamental Higgs field [Maas, Törek 1804.04453]
 - \rightarrow Blue gauge invariant spectrum
 - \rightarrow Red predictions from gauge-invariant PT $_{[Maas 1712.04721]}$
 - → Standard PT fails



SU(2) with $N_f = 2$ fundamental flavors

[Drach, Janowski, Pica, Prelovsek]

▶ Starting new investigations using Wilson-clover fermion with Symanzik gauge action

▶ Improving on earlier work with (unimproved) Wilson fermions and plaquette gauge action



Little changes w.r.t. unimproved setup



- Approaching the chiral limit
- \blacktriangleright Investigate scattering and $\rho\pi\pi$ coupling

Fundamental composite 2HDM with four flavors

[Ma, Cacciapaglia JHEP03(2016)211]

▶ Global symmetry at low energies:

 $SU(4) \times SU(4)$ broken to $SU(4)_{diag}$

▶ 15 pNGB transform under custodial symmetry

 $SU(2)_L \times SU(2)_R$

$$\Rightarrow \mathbf{15}_{SU(4)_{\text{diag}}} = (2,2) + (2,2) + (3,1) + (1,3) + (1,1)$$

 \rightarrow One doublet plays the role of the Higgs doublet field

 \rightarrow Other doublet and triplets are stable; could play role of dark matter

► Vecchi: "choose the right couplings to RH top" [Edinburgh talk] $\Rightarrow (2,2) + (2,2) + (3,1) + (1,3) + (1,1)$ $\rightarrow \text{ effectively SU(4)/Sp(4)}$

SU(4)/Sp(4) composite Higgs

[Bennett, Hong, Lee, Lin, Lucini, Piai, Vadacchino]

Systematic program to investigate Sp(2N) gauge theories for $N_f = 2$ fund. flavors and N > 1

 \rightarrow Quenched results for Sp(4) published [Bennett et al. JHEP 1803(2018)185]

- \rightarrow First dynamical results for masses and decay constants (Lee Thu 2:20 PM)
- \rightarrow Qualitative agreement between quenched and dynamical results

 \rightarrow Comparison between $N_f = 2$ fundamental and anti-symmetric

