

Proton decay matrix element on the lattice at the physical point

Jun-Sik Yoo

Stony Brook University

jun-sik.yoo@stonybrook.edu

July 24, 2018

co-author: Yasumichi Aoki, Sergey Syritsyn, Taku Izubuchi

Acknowledgement

Our project is supported by USQCD and the computation was done using 2017-2018 allocation at JLab.

I appreciate Sergey Syritsyn, Yasumichi Aoki, Taku Izubuchi, Amarjit Soni, Daniel Hoving and many of the BNL lattice group members having productive discussion with me.

Introduction

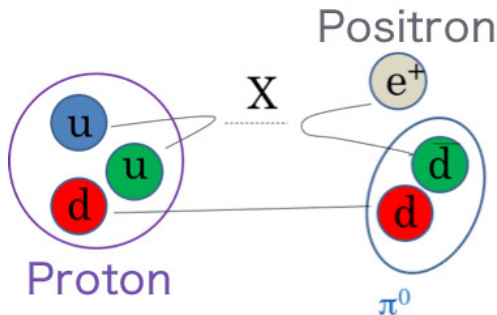


Figure: Proton decay schematic diagram (Hyper-k.org)

- GUT, SUSY-GUT: new interactions between quarks and leptons
 - Nucleon stability
 - Proton decay : Baryon number violation
- Is one of Sakharov's necessary conditions for Baryogenesis

Introduction

Experimental Effort

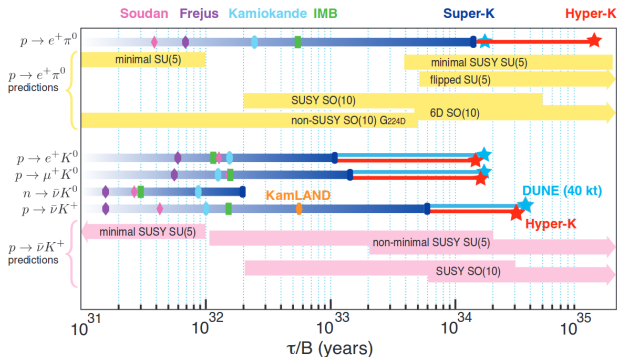


Figure: Proton decay schematic diagram [DUNE arXiv:1512.06148]

- Super-Kamiokande
- DUNE (Deep Underground Neutrino Experiment)

Introduction

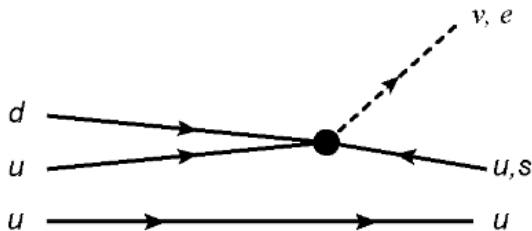


Figure: Proton decay effective diagram

$$C^I \langle \ell, PS | [\bar{\ell} O_{\Gamma\Gamma'}] | N \rangle = C^I \bar{v}_\ell \langle PS | O_{\Gamma\Gamma'} | N \rangle$$

$$\rightarrow C^I P_{\Gamma'} \left[W_0^{\Gamma\Gamma'}(q^2) - \frac{i \not{q}}{m_N} W_1^{\Gamma\Gamma'}(q^2) \right] u_N(p, s)$$

, where C^I being the Wilson coefficient of I-th kind of operator.

Introduction

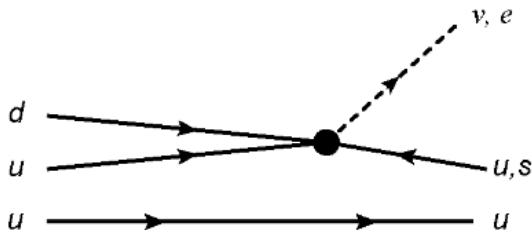


Figure: Proton decay schematic diagram

$$\Gamma(N \rightarrow P + \bar{\ell}) = \frac{m_N}{32\pi} [1 - (\frac{m_P}{m_N})^2]^2 |\sum_l C^l W_0^l(N \rightarrow P)|^2$$

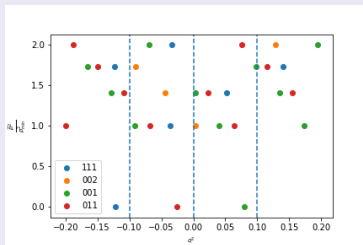
, where C^l being the Wilson coefficient of l-th kind of operator.

Lattice settings

- RBC/UKQCD generated $N_f = 2 + 1$ dynamic Domain wall Fermion, gauge action Iwasaki-DSDR
- Lattice size $24^3 \times 64 (L \sim 4.8 fm)$, $L_5 = 24$,
 $\beta = 1.633$, $m_\ell a = 0.00107$, $m_h a = 0.0850$, $m_{res} = 0.00228$
- $a^{-1} = 1.015$ GeV, $m_\pi = 139$ MeV, $m_K = 505$ MeV, $m_\pi L \sim 3.4$
- Deflated CG with 2000 Eigenvectors (basis 1000)
- Generated 32+1 AMA samples on 52 gauge configurations with 3 source-sink separation, i.e., $t_{sep} \in \{8, 9, 10\}$
- To meet the kinematic condition, chose the most suitable two sets of \vec{p} for each meson.

Lattice settings

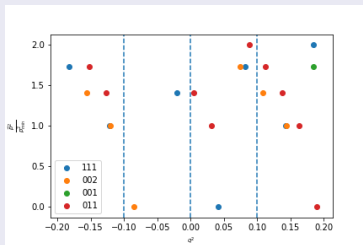
$$p = p' + q$$



- Physical $q^2 \sim 10^{-8} \text{ GeV}^2$
- Chose [001] [011] for Kaons
- Chose [111] [002] for pions

Lattice settings

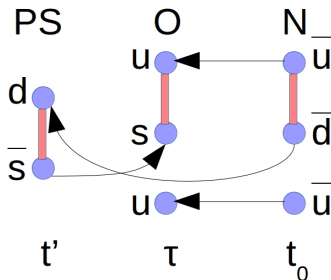
$$p = p' + q$$



- Physical $q^2 \sim 10^{-8} \text{ GeV}^2$
- Chose [001] [011] for Kaons
- Chose [111] [002] for pions

Lattice settings

$$\langle PS(t') | O_{\Gamma\Gamma'}(\tau) | N(t_0) \rangle$$



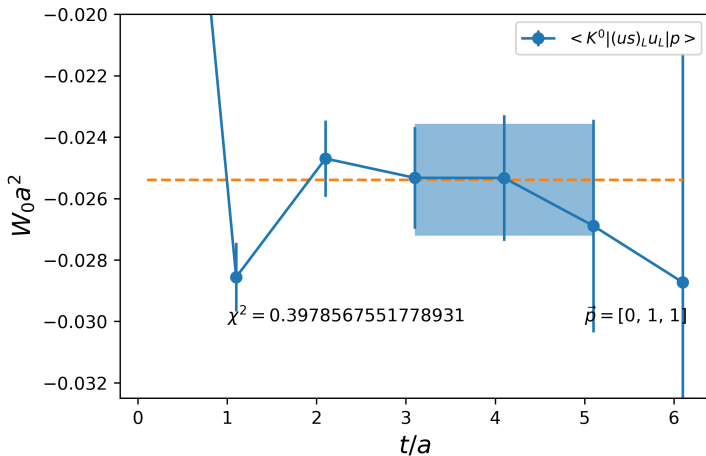
$$R_3^{\Gamma\Gamma'}(t', \tau, t_0; \vec{p}', q; P) = \frac{\text{tr}[PC_3^{\Gamma\Gamma'}(t', \tau, t_0; q)]}{C_{PS}(t', \tau; \vec{p}') \text{tr}[P_4 C_N(\tau, t_0)]} \sqrt{Z_{PS} Z_N}$$

As $t' - \tau \rightarrow \infty$, $\tau - t_0 \rightarrow \infty$,

$$R_3 \longrightarrow P_{\Gamma'} \left[W_0^{\Gamma\Gamma'}(q^2) - \frac{i \not{q}}{m_N} W_1^{\Gamma\Gamma'}(q^2) \right] u_N(p, s)$$

Form Factors

$$W_0^{\Gamma\Gamma'} = R_3^{\Gamma\Gamma'}(t', t, t_0; \vec{p}', q; \textcolor{red}{P}_4) - \frac{m_N - E_\pi}{q_j} R_3^{\Gamma\Gamma'}(t', t, t_0; \vec{p}', q; i\textcolor{red}{P}_4\gamma_j)$$



Form Factors

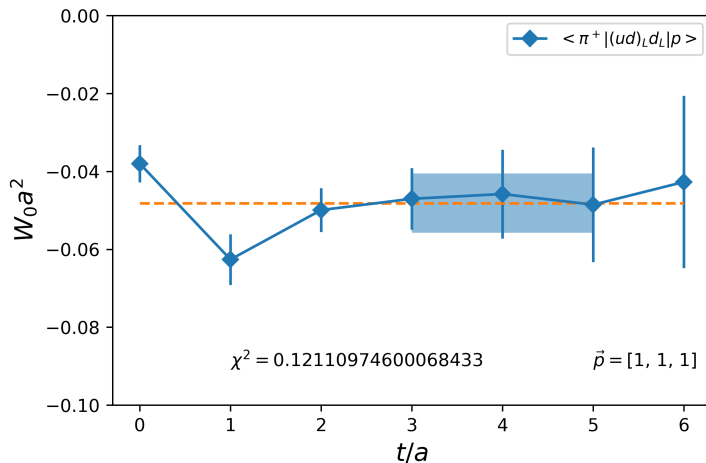


Figure: Pion Channel form factor with $\vec{p}_\pi = [111]$

Form Factors

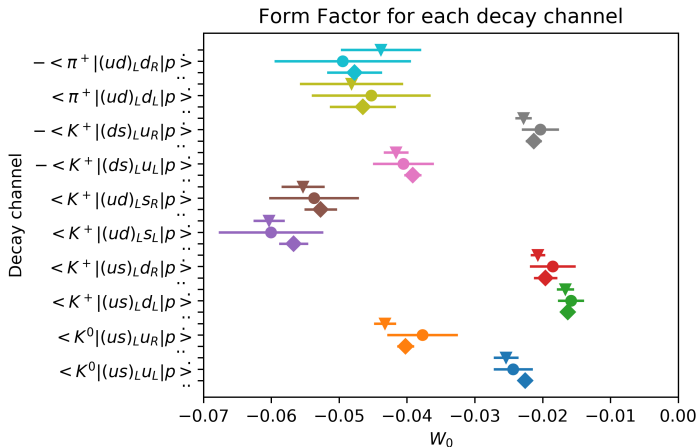


Figure: Form factors Channel by Channel with $\vec{p}_\pi = [111]$ $\vec{p}_K = [011]$

Interpolation to physical point

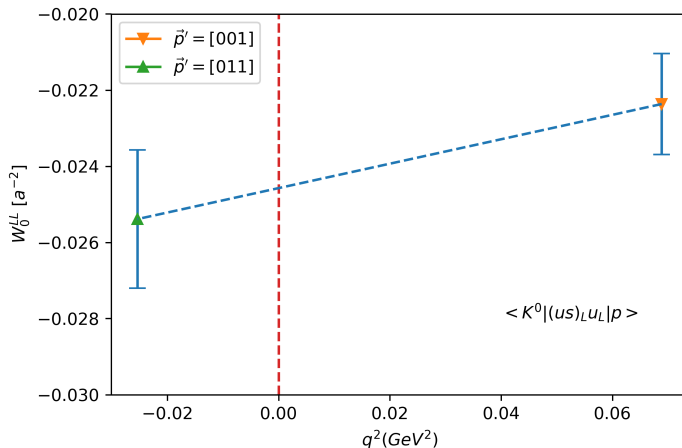


Figure: Form factors Across different \vec{p}_K

Consistency check with earlier study

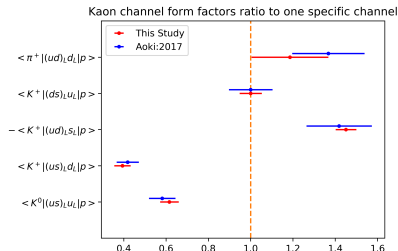


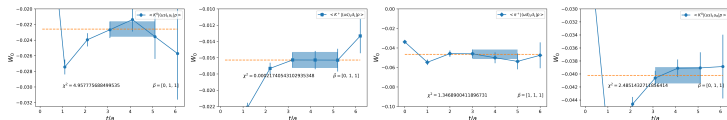
Figure: Normalized Form factors Comparison with Earlier study and our data at $t_{sep} = 10$

	Stat. [%] (This study)	Stat. [%] (Aoki:2017)	Chiral Extrapol. [%]	a^2 [%]	Δ_Z [%]
$\langle K^0 (us)_L u_L p \rangle$	5.1	3.5	3.1	5.0	8.1
$\langle K^+ (us)_L d_L p \rangle$	17	4.4	7.5	5.0	8.1
$\langle K^+ (ud)_L s_L p \rangle$	4.8	3.0	3.9	5.0	8.1
$-\langle K^+ (ds)_L u_L p \rangle$	4.1	2.8	2.8	5.0	8.1
$\langle \pi^+ (ud)_L d_L p \rangle$	15.8	3.4	5.7	5.0	8.1

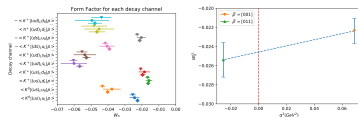
Summary

What is done

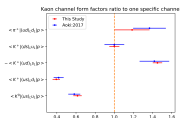
Preliminary bare form factor is extracted



Preliminary Analysis



Comparison with earlier study



Tasks Ahead

Multi-state fit(Excited states)

Renormalization

Increase samples(Esp. pion channel)

Different Lattice (Different volume, Different cutoff)

Indirect method checkup (using α, β)

Exploring New channels

References



Y. Aoki, E. Shintani, A. Soni

Proton decay matrix element on lattice, 2013



Hyper Kamiokande Collaboration



Y. Aoki, T. Izubuchi, E. Shintani, A. Soni

Improved lattice computation of proton decay matrix element, 2017



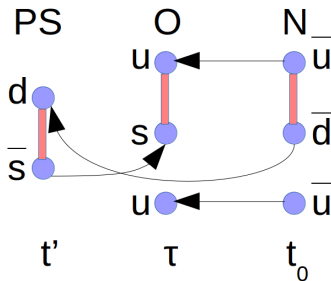
Deep Underground Neutrino Experiment

Conceptual Design Report

The End

Lattice settings

$$\langle PS(t') | O_{\Gamma\Gamma'}(\tau) | N(t_0) \rangle$$



$$\begin{aligned}
 C_{PS}(t) &= \frac{Z_{PS}}{2m_{PS}} e^{-m_{PS}t} \\
 &\quad \langle 0 | J_{PS} | PS \rangle \\
 &= \sqrt{Z_{PS}}
 \end{aligned}$$

Effective mass plots

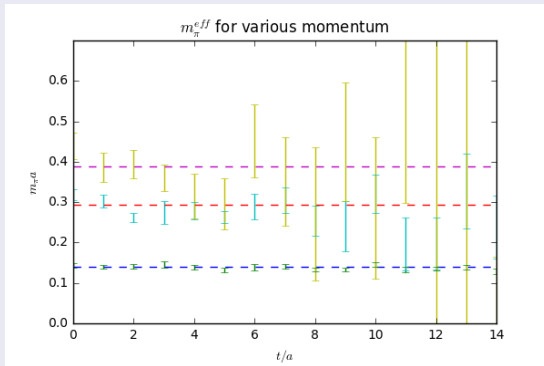


Figure: Pion mass dispersion with $\vec{p} = [000], [001], [011]$

Effective mass plots

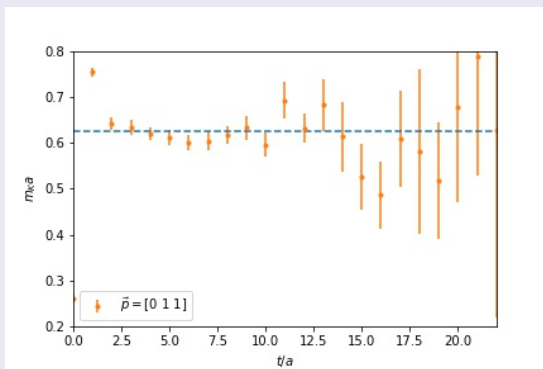
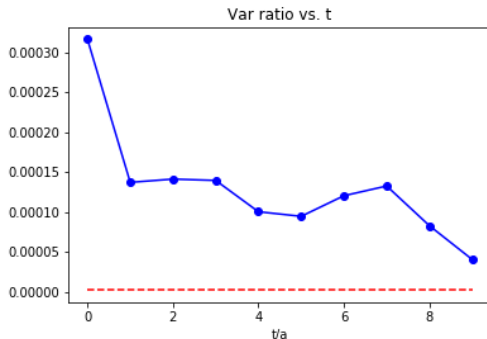


Figure: Kaon effective mass with $\vec{p} = [011]$

AMA method



- possible distribution of computation time to exact, approximate samples enhance statistics
- Should be easy to devise approximate sample for Domain Wall Fermion :zmobius
- the most cost efficient variance around

$$N_d/N_{ex} = \sqrt{\frac{1}{1-10^{-4}}} \cdot 11 \approx 270 \text{ for both protons and pions.}$$

Using two-state fit to exclude excited states contribution

$$C^{2pt}(t_f, t_i) = \text{Tr}[\mathcal{P} \sum_{\vec{x}} \langle 0 | J_N(t, \vec{x}) \bar{J}_N(0, \vec{0}) | 0 \rangle]$$

$$C^{3pt}(t_f, \tau, t_i) = \text{Tr}[\mathcal{P} \sum_{\vec{x}, \vec{x}'} \langle 0 | J_N(t, \vec{x}) \mathcal{O}(\tau, \vec{x}') \bar{J}_N(0, \vec{0}) | 0 \rangle]$$

, where J_N , nucleon interpolator on the lattice, \mathcal{P} spin projection operator.
Spectral decomposition:

$$C^{2pt}(t_f, t_i) = |\mathcal{A}_0|^2 e^{-M_0(t_f - t_i)} + |\mathcal{A}_1|^2 e^{-M_1(t_f - t_i)} + \dots$$

$$C^{3pt}(t_f, \tau, t_i) = |\mathcal{A}_0|^2 \langle \mathcal{O}_{00} \rangle e^{-M_0(t_f - t_i)} + \mathcal{A}_0 \mathcal{A}_1^* \langle \mathcal{O}_{01} \rangle e^{-M_0(t_f - t_i)} + \\ \mathcal{A}_1 \mathcal{A}_0^* \langle \mathcal{O}_{10} \rangle e^{-M_0(t_f - t_i)} + |\mathcal{A}_1|^2 \langle \mathcal{O}_{11} \rangle e^{-M_1(t_f - t_i)} + \dots$$

Using two-state fit to exclude excited states contribution

- Two-state fit to two point correlation function

\mathcal{A}_0	$M_0 a$	\mathcal{A}_1	$M_1 a$
0.00027891	-0.961095	0.640617	-2.278580756339

Chiral Perturbation theory

- $\langle PS(p') | O^{\Gamma\Gamma'}(q) | N(p) \rangle$ to be approximated to chiral perturbation to $\langle 0 | O^{\Gamma\Gamma'}(q) | N(p) \rangle \rightarrow$ Low Energy Constant (LEC)
- LEC to be calculated :
 $\langle 0 | O^{LL}(q) | N(p) \rangle = \alpha P_L u_s$
 $\langle 0 | O^{LR}(q) | N(p) \rangle = \beta P_L u_s$
, where O being specifically $(ud)_{\Gamma} P_{\Gamma'} u$

!sic! Pending slides!