

# Phase structure of three flavor QCD in external magnetic fields using HISQ



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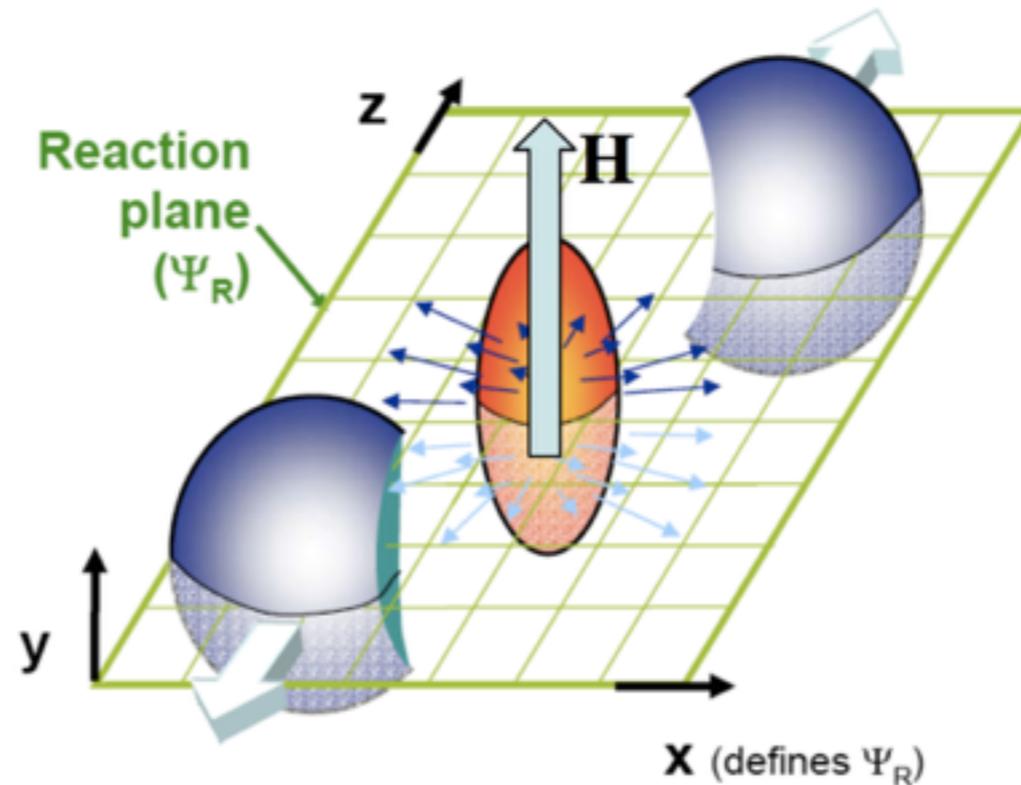
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Xiao-Dan Wang (CCNU)

Yu Zhang (CCNU)

# Our purpose 1/4

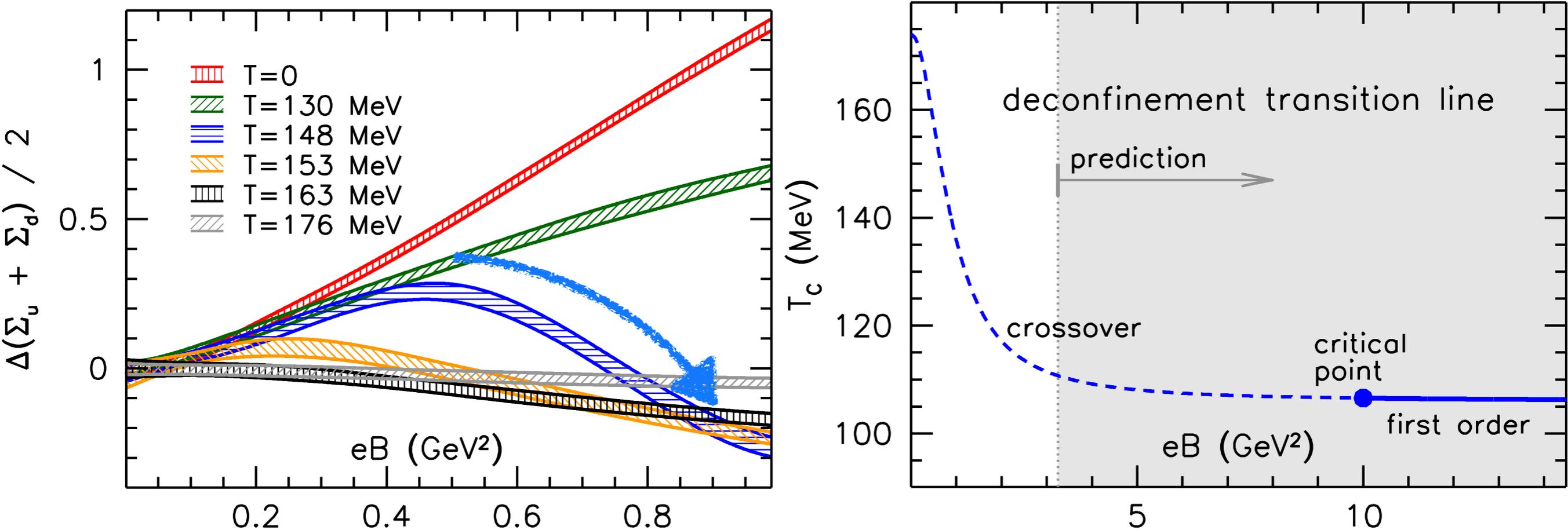
Relativistic Heavy Ion Collisions = QCD + External magnetic field



Non-central heavy-ion collisions produce huge magnetic field,  
 $eB \sim 10^{18}$  Gauss ( $\sqrt{eB} \sim 442$  MeV)  
 (c.f.  $eB^{\text{lab}} = 4.5 \times 10^5$  Gauss)  
 → It might affect to QCD phase structure

# Our purpose 2/4

Why  $m_\pi = 320$  &  $80$  MeV with HISQ in external magnetic fields?



## Stout staggered fermions at the physical point

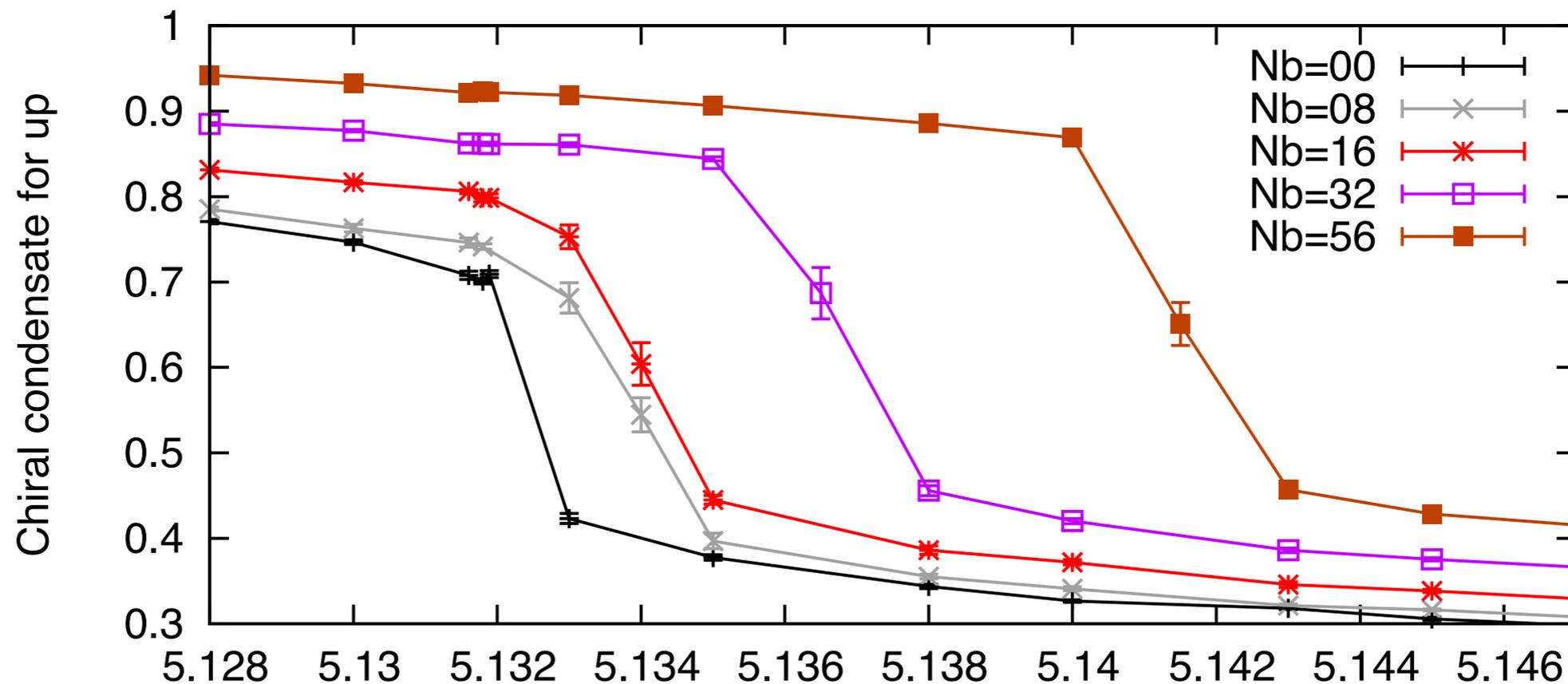
- **Inverse** magnetic catalysis at the critical temperature
- $T_c$  decreases along with external field  $eB$
- New critical endpoint for large  $eB$  is suggested by an effective model

# Our purpose 3/4

Why  $m_\pi = 320$  &  $80$  MeV with HISQ in external magnetic fields?

Unimproved staggered fermions  $N_f = 3$

A.T. *et al* 1711.02884



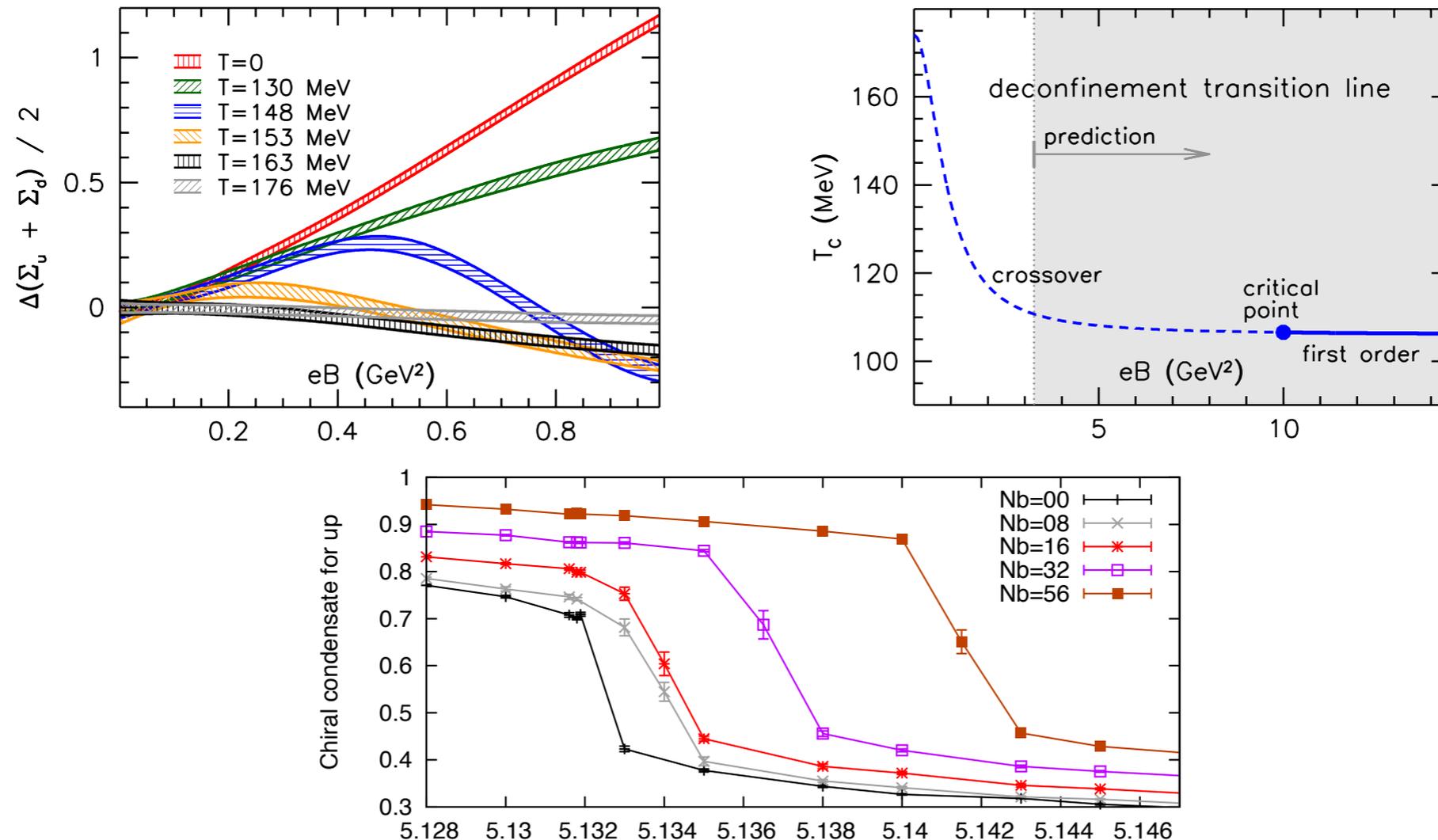
$\beta$

In our previous work, we found for  $m_\pi \sim 300$  MeV,

- **Normal** magnetic catalysis for whole temperature
- $T_c$  increases along with external field  $eB$  ( $\propto Nb$ )
- The Binder cumulant indicates 1st order phase transition

# Our purpose 4/4

Why  $m_\pi = 320$  &  $80$  MeV with HISQ in external magnetic fields?



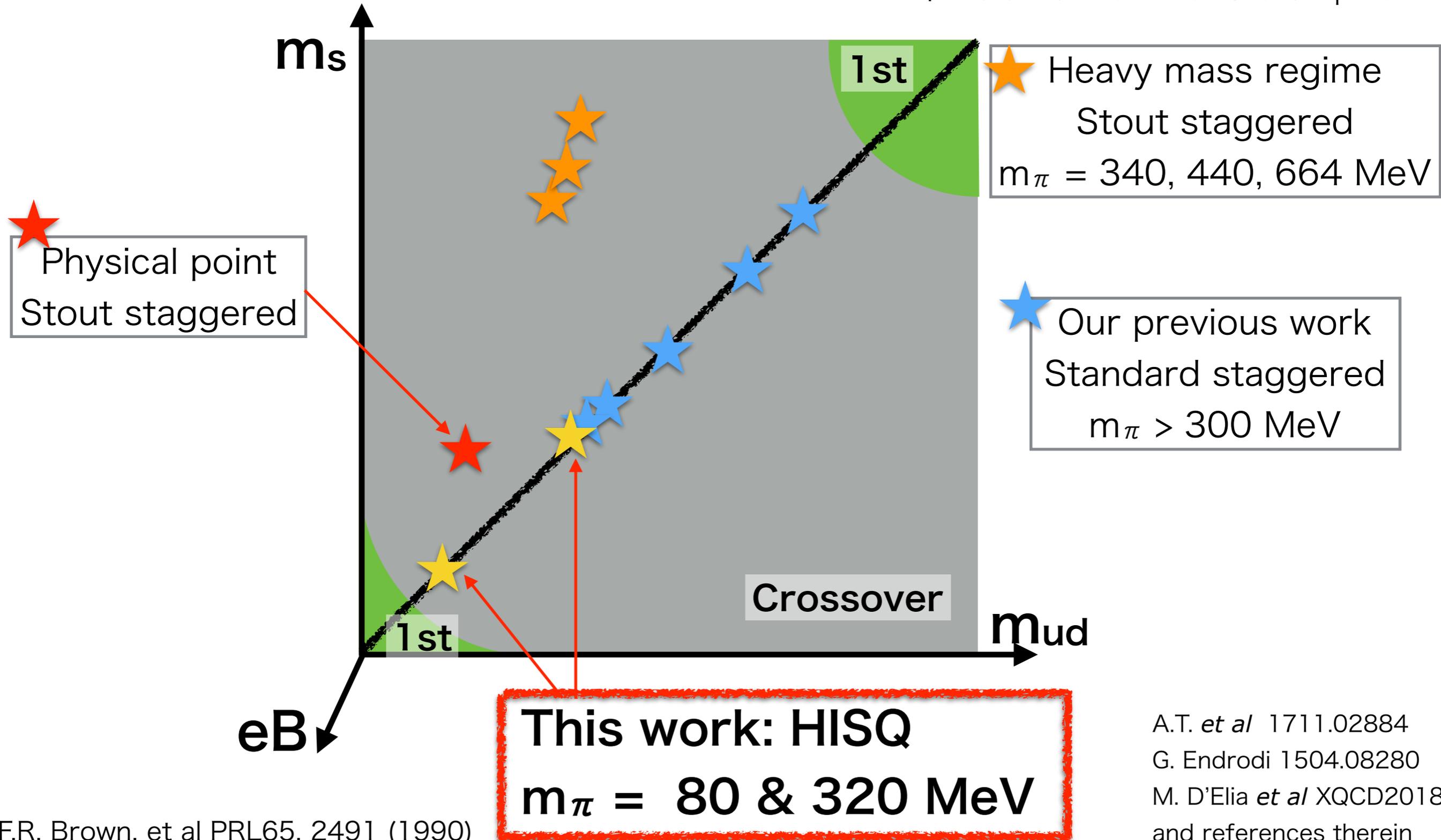
We employ **Highly Improved Staggered Quarks (HISQ):**

- To see cutoff effect and mass dependence on the catalysis
- To see  $T_c(eB)$  behavior. Does it go up or down?
- To find first order at large external field  $eB$

# Our target regime

Works with external magnetic field

† There are also Wilson & overlap works



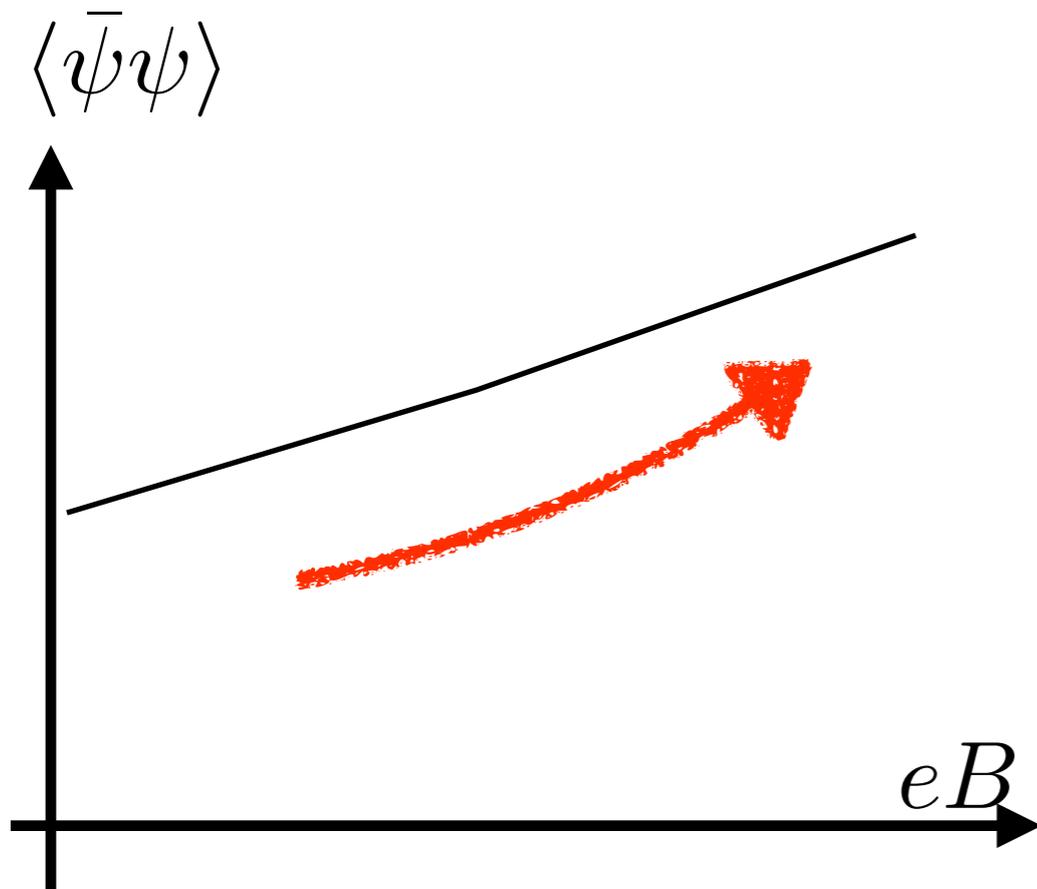
A.T. *et al* 1711.02884  
 G. Endrodi 1504.08280  
 M. D'Elia *et al* XQCD2018  
 and references therein

F.R. Brown, *et al* PRL65, 2491 (1990)

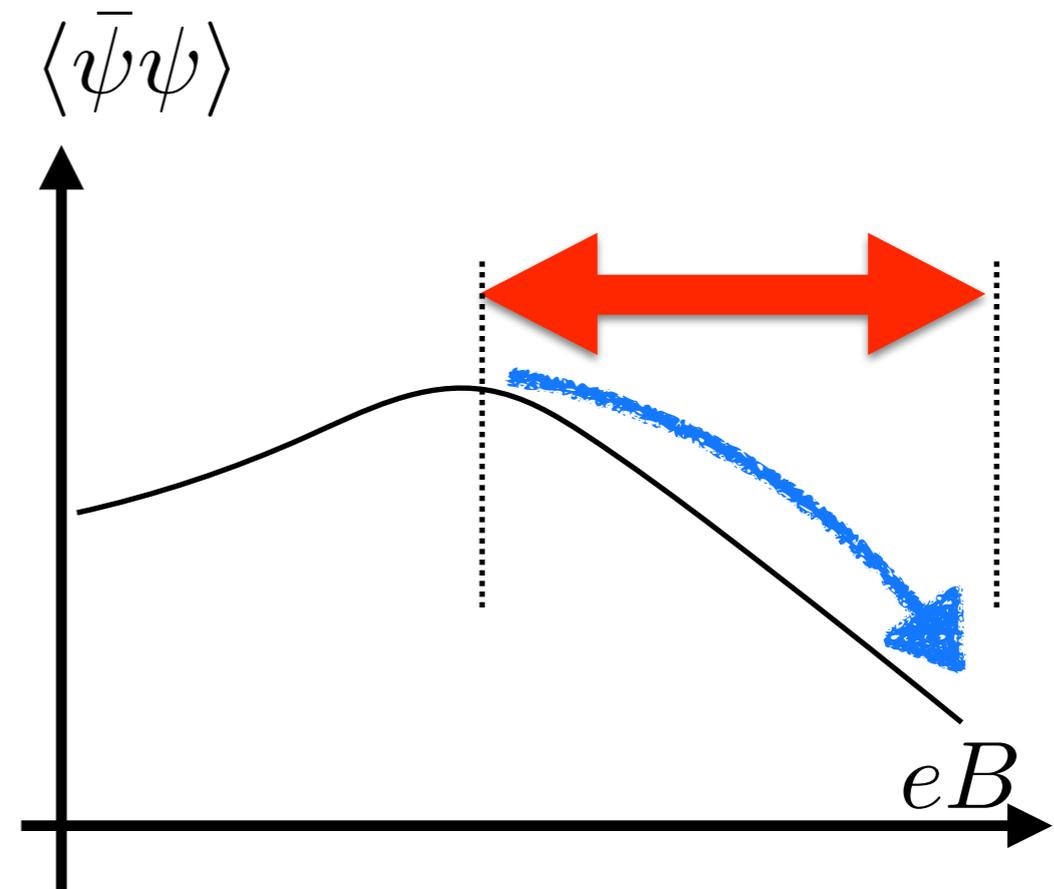
# Terminology

External magnetic field (inversely) catalyzes the chiral condensate

(Normal)  
Magnetic Catalysis



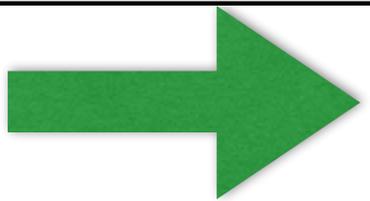
Inverse  
Magnetic Catalysis



# Our setup

$m_\pi = 320$  &  $80$  MeV with HISQ

	3 degenerate flavors
<u>Previous</u>	Unimproved staggered quarks with RHMC $m_\pi > 300$ MeV
<u>This work</u>	3 degenerate flavors with <b>HISQ</b> (RHMC) $m_\pi = \underline{320}$ MeV (~ lightest in unimproved case) & $\underline{80}$ MeV (near to SU(3) chiral limit)



The Dirac spectrum and topological susceptibility using the stochastic estimator.

[L. Giusti et al, 2009. G. Cossu et al 2016. P. d. Forcrand et al 2018.]

# Our setup

$m_\pi = 320$  &  $80$  MeV with HISQ

$N_f = 3$  QCD with **HISQ** ( $m_\pi = 320$  and  $80$  MeV)

Tree level Symanzik gauge action

$N_\sigma^3 \times N_\tau = 16^3 \times 6$  ( $a \sim 0.27$  fm)

The number of configurations  $\sim O(1000)$

Fermilab and CCNU GPU machines

$$N_b \propto a^2 eB \quad : \text{ \# of magnetic flux}$$

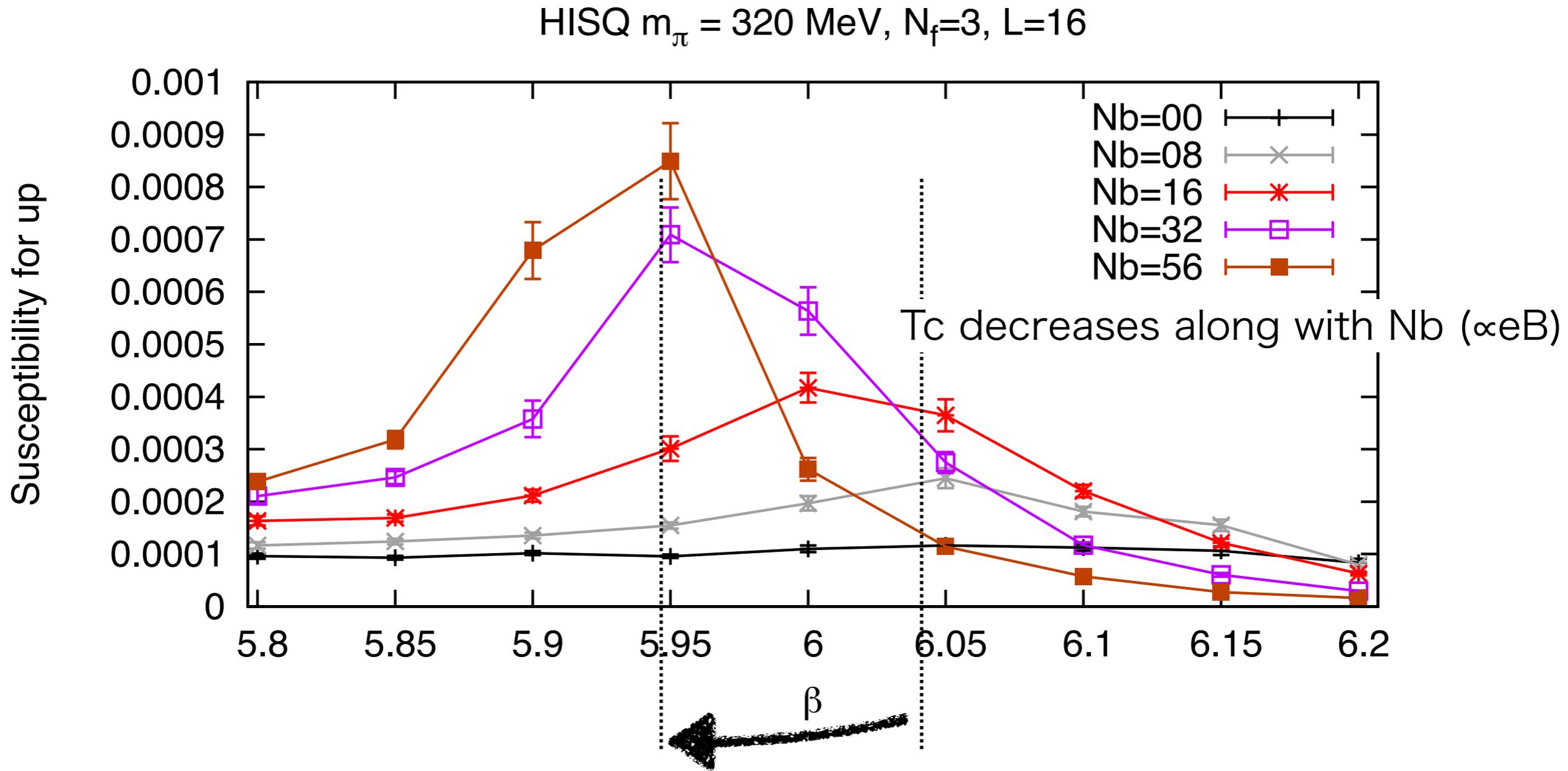
Nb(Nx=16)	0	8	16	32	56
$\sqrt{eB}$ (MeV)	0	330	460	660	870

m	320 MeV	80 MeV	80 MeV
Nx=Ny	16	16	24
$m_\pi L$	7	1.7	2.6 (Not showed)

$$m_{\pi} = 320 \text{ MeV}$$

# Results with $m_\pi = 320$ MeV

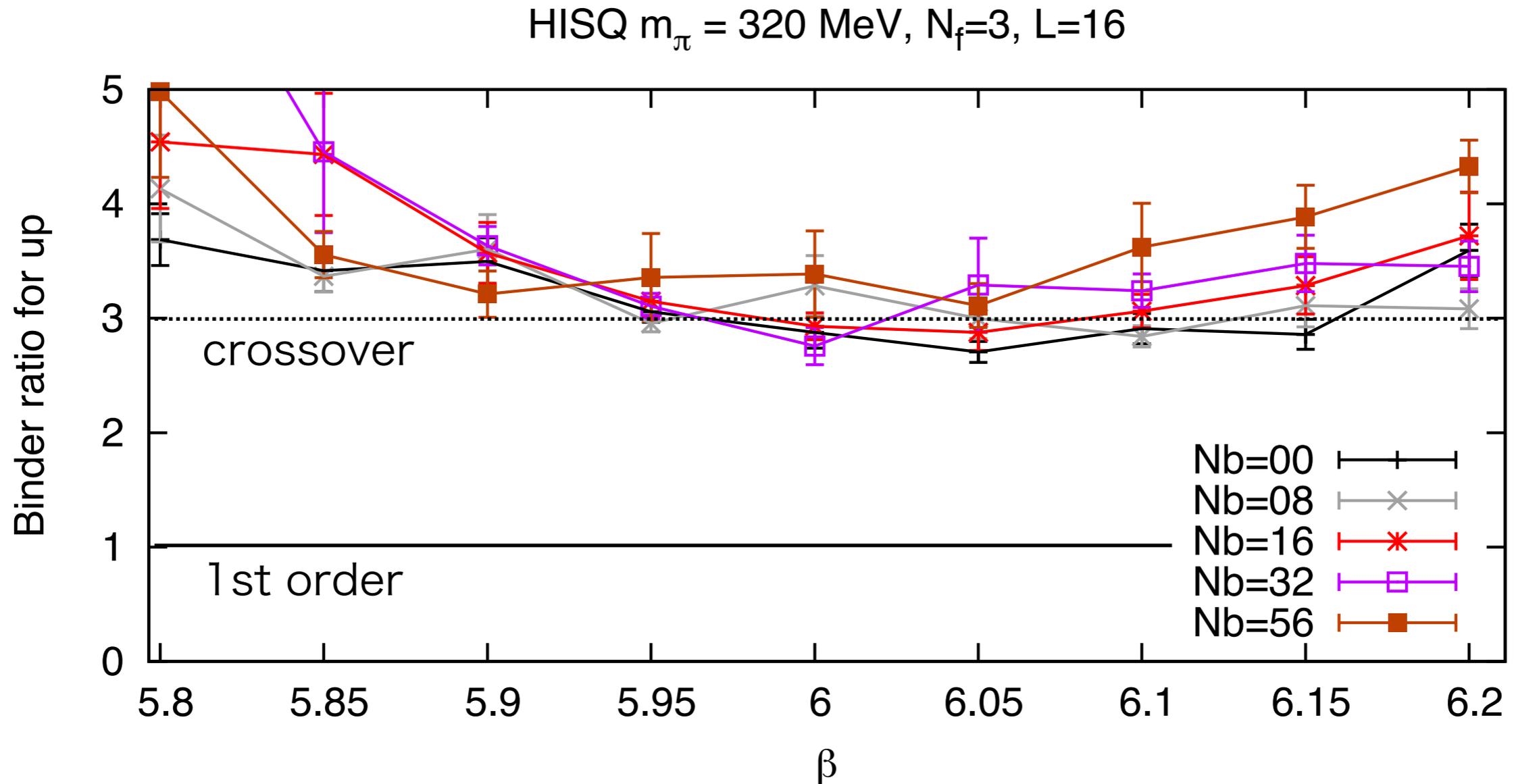
Susceptibility for heavier regime (~ lightest mass in unimproved staggered)



Consistent with results with physical point

# Results with $m_\pi = 320$ MeV

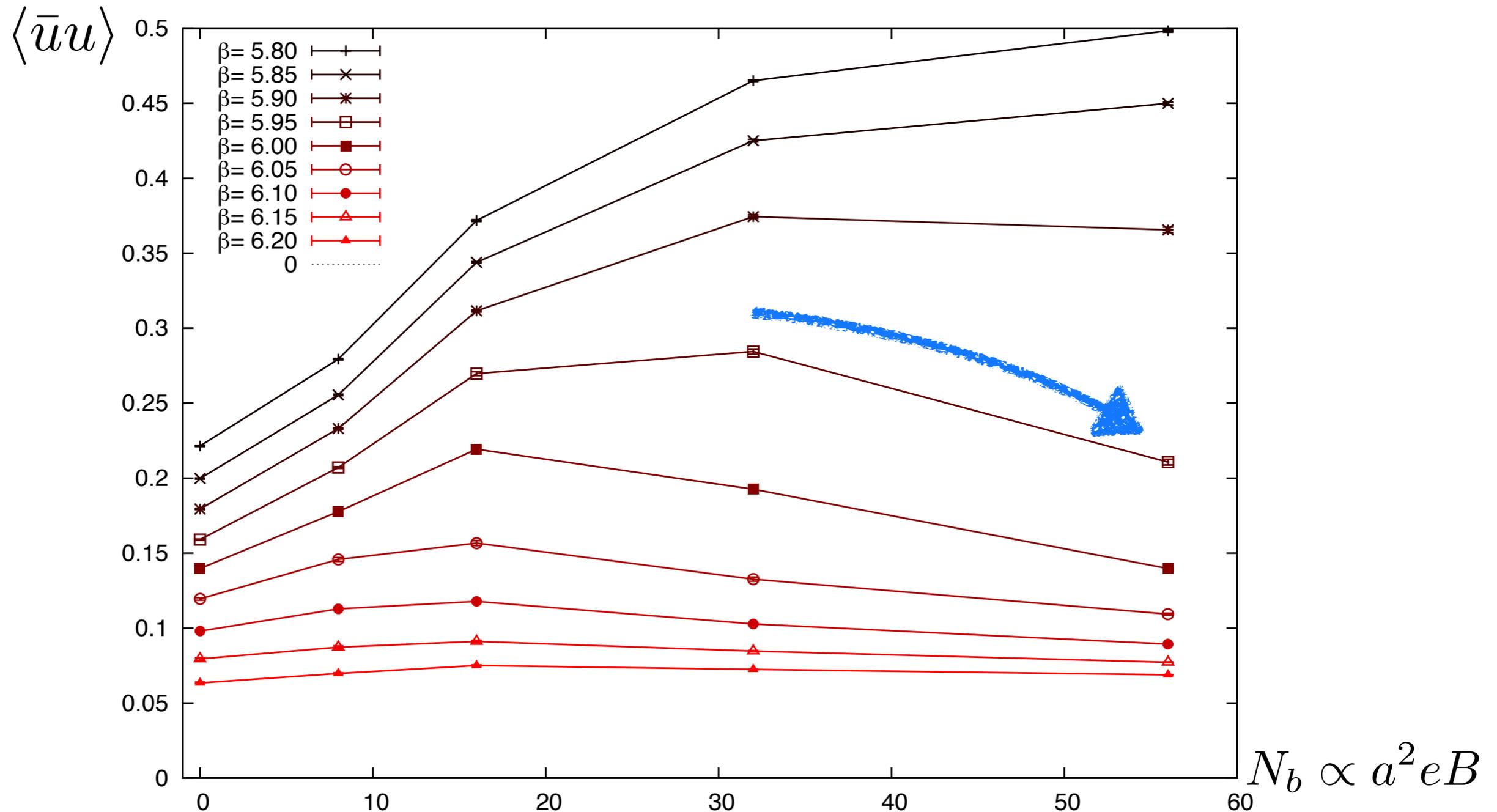
Binder cumulant for Heavier regime (~ lightest mass in unimproved staggered)



No indication of 1st order, different from unimproved staggered results

# Results with $m_\pi = 320$ MeV

Chiral condensate for Heavier regime (~ lightest mass in unimproved staggered)



$m_\pi = 320$  MeV data near  $T_c$  shows inverse magnetic catalysis  
 -> Unimproved staggered results are probably artifacts

# Results with $m_\pi = 320$ MeV

## Decomposition of relative increase of the chiral condensate

Relative increase of the chiral condensate

$$r(B) = \frac{\langle \bar{\psi}\psi \rangle(B) - \langle \bar{\psi}\psi \rangle(0)}{\langle \bar{\psi}\psi \rangle(0)}$$

$$r^{\text{full}}(B) = \underbrace{r^{\text{sea}}(B)}_{\substack{\text{eB only in} \\ \text{sea quarks}}} + \underbrace{r^{\text{val}}(B)}_{\substack{\text{eB only in} \\ \text{valence quarks}}} + O(B^4)$$

We can see sea and valence effects on normal/inverse magnetic catalysis

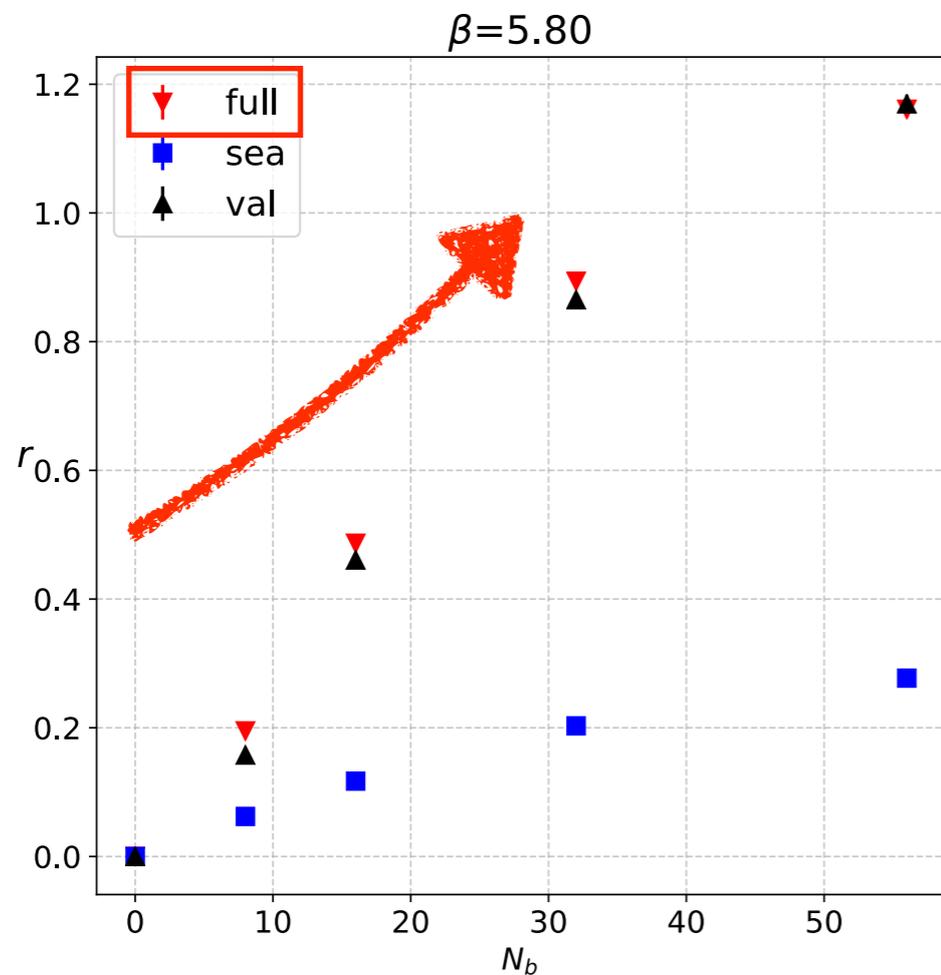
# Results with $m_\pi = 320$ MeV

Heavier regime: Consistent with previous results at physical points

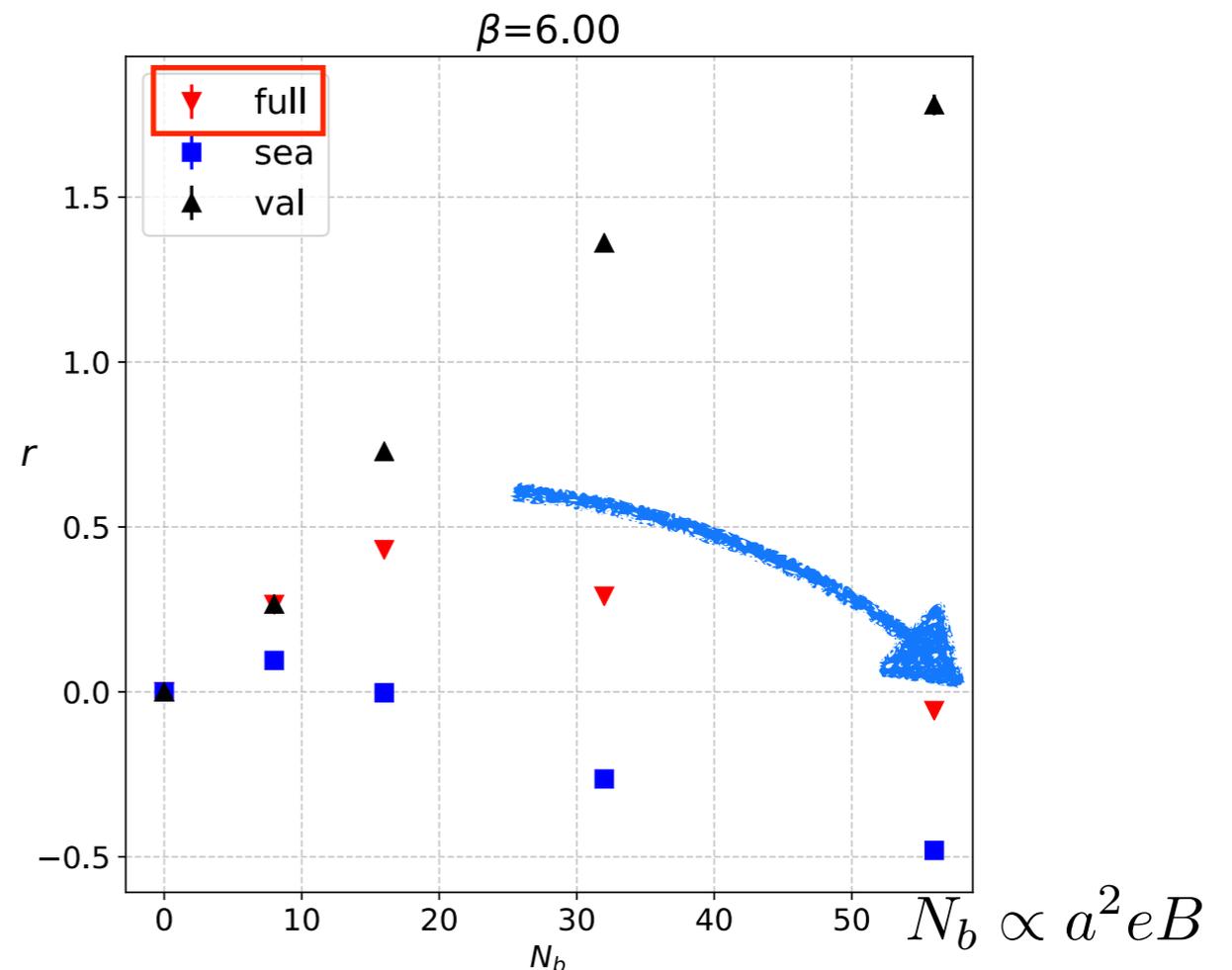
$$r(B) = \frac{\langle \bar{\psi}\psi \rangle(B) - \langle \bar{\psi}\psi \rangle(0)}{\langle \bar{\psi}\psi \rangle(0)}$$

$T < T_c$

$T > T_c$



(Normal) Magnetic Catalysis



**Inverse Magnetic Catalysis**

Sea quarks drive inverse catalysis  
Similar to previous results at the physical point

Using stochastic estimator [L. Giusti et al, 2009. G. Cossu et al 2016. P. d. Forcrand et al 2018.]

# Results with $m_\pi = 320$ MeV

Dirac spectrum: sea quarks drive inverse catalysis

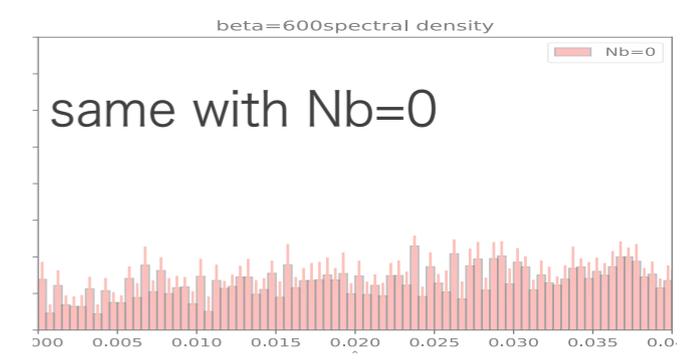
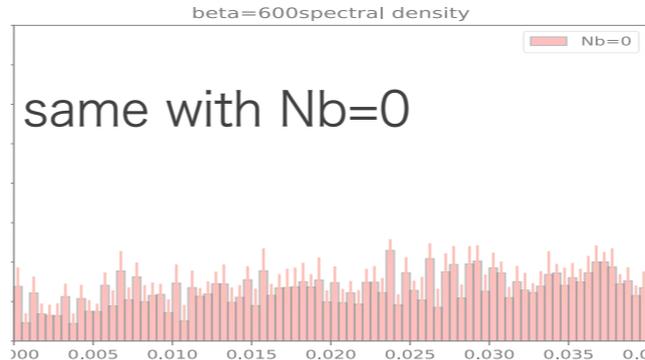
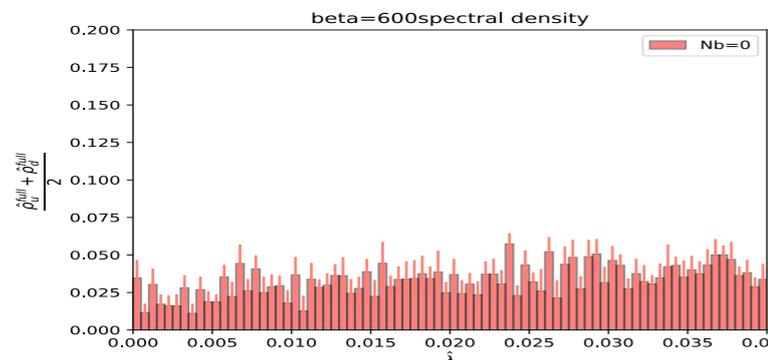
$$\beta = 6(T \gtrsim T_c)$$

Full

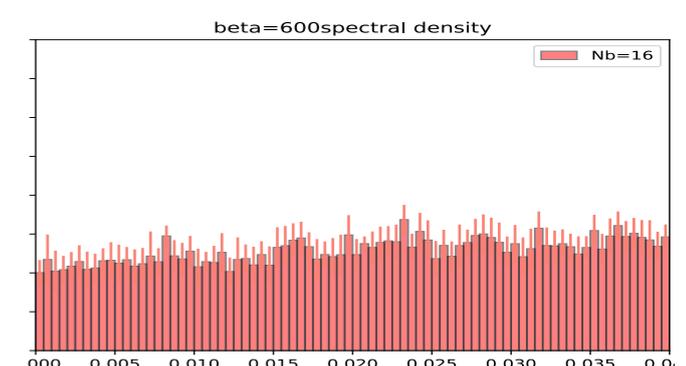
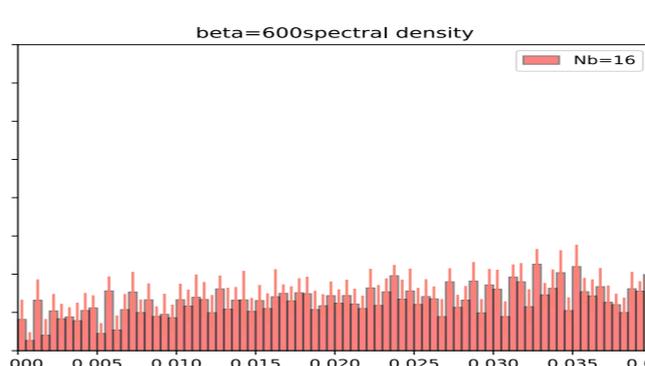
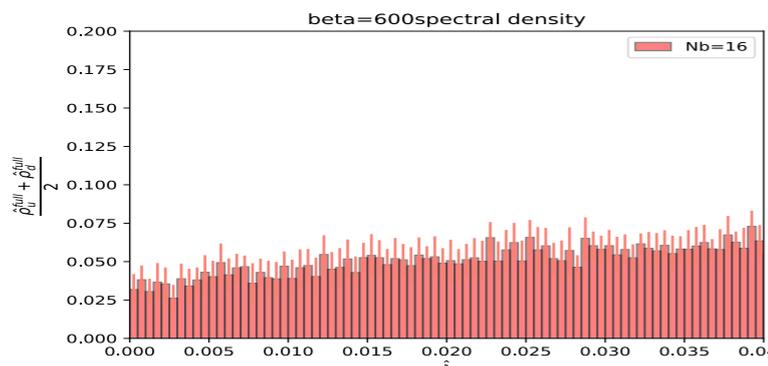
$eB$  on sea

$eB$  on valence

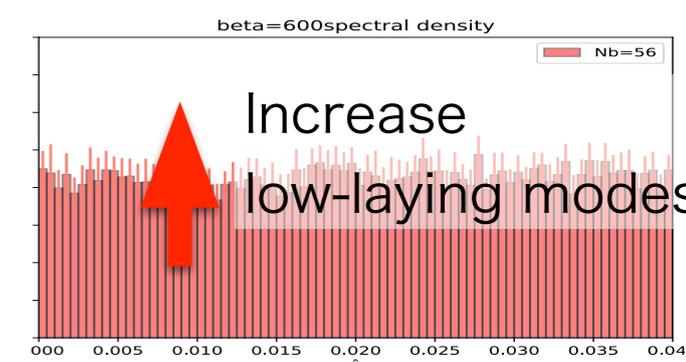
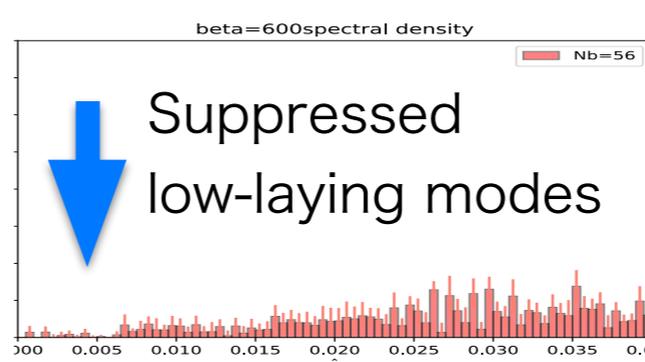
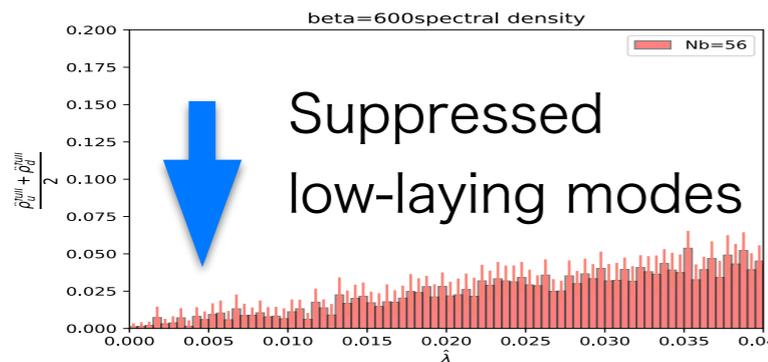
Nb=0



Nb=16



Nb=56

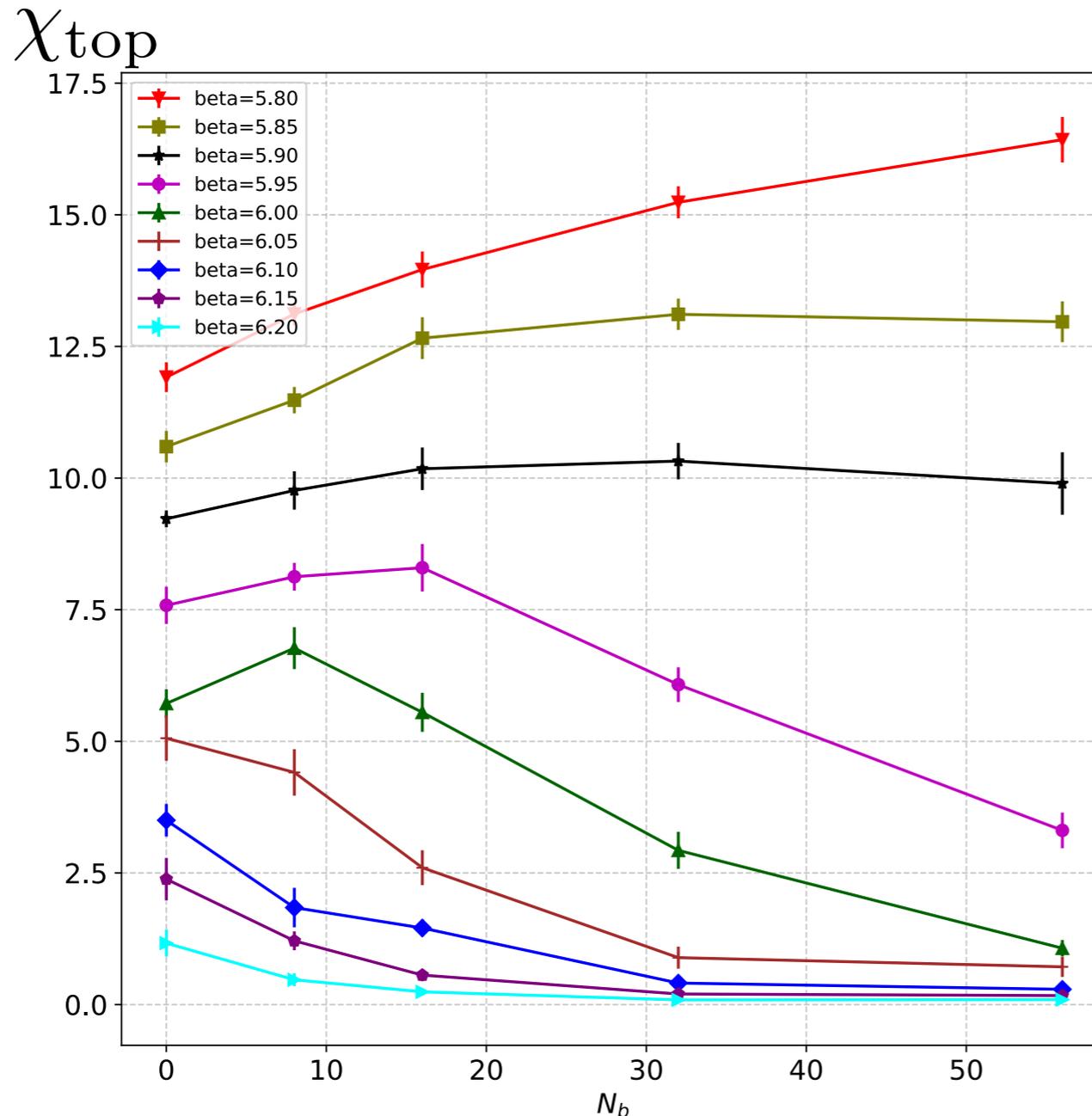


Valence quarks only contribute to normal catalysis

Using stochastic estimator [L. Giusti et al, 2009. G. Cossu et al 2016. P. d. Forcrand et al 2018.]

# Results with $m_\pi = 320$ MeV

Topological susceptibility shows similar tendency to the chiral condensate



$$\langle \bar{\psi}\psi \rangle = \int_0^\infty d\lambda \frac{2m\rho(\lambda)}{\lambda^2 + m^2}.$$

$$\chi_{\text{top}} = m^2 \int_0^\infty d\lambda \frac{4m^2\rho(\lambda)}{(\lambda^2 + m^2)^2}.$$

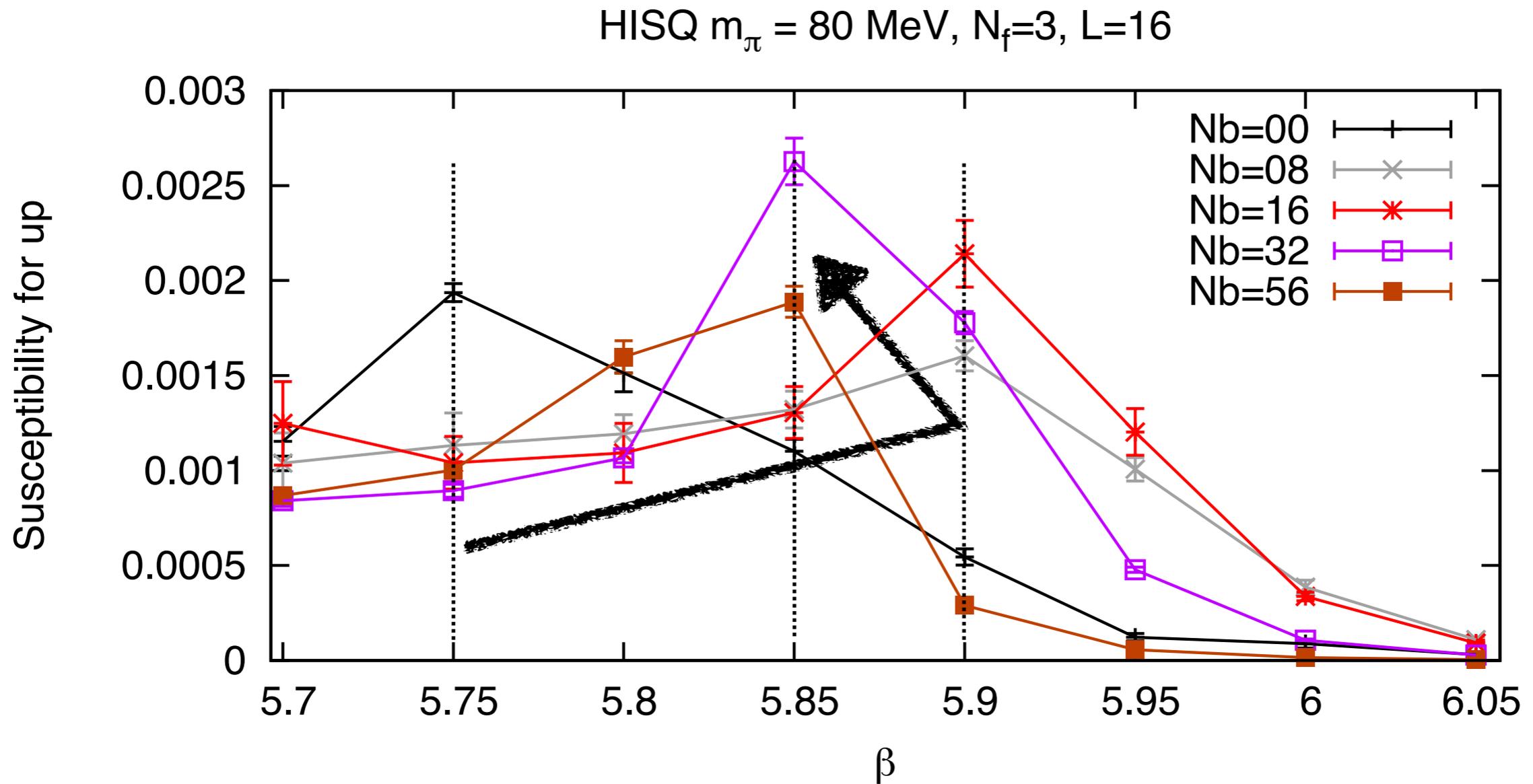
Similar tendency to the chiral condensate

Using stochastic estimator [L. Giusti et al, 2009. G. Cossu et al 2016. P. d. Forcrand et al 2018.]

$$m_{\pi} = 80 \text{ MeV}$$

# Results with $m_\pi = 80$ MeV

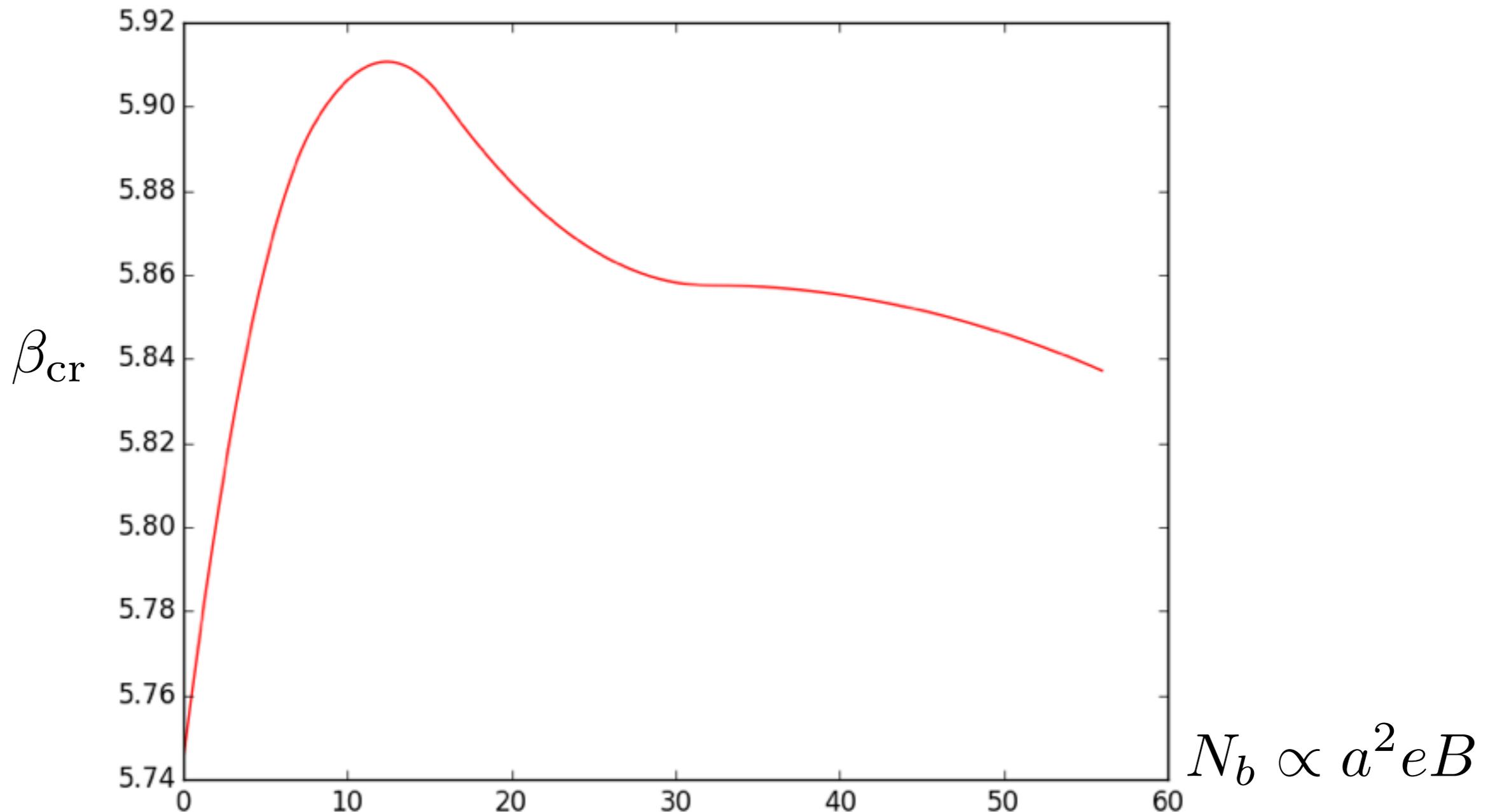
Susceptibility in light pion mass regime (near the SU(3) chiral limit)



$T_c$  goes up and down along with Nb??

# Results with $m_\pi = 80$ MeV

Susceptibility in light pion mass regime (near the SU(3) chiral limit)

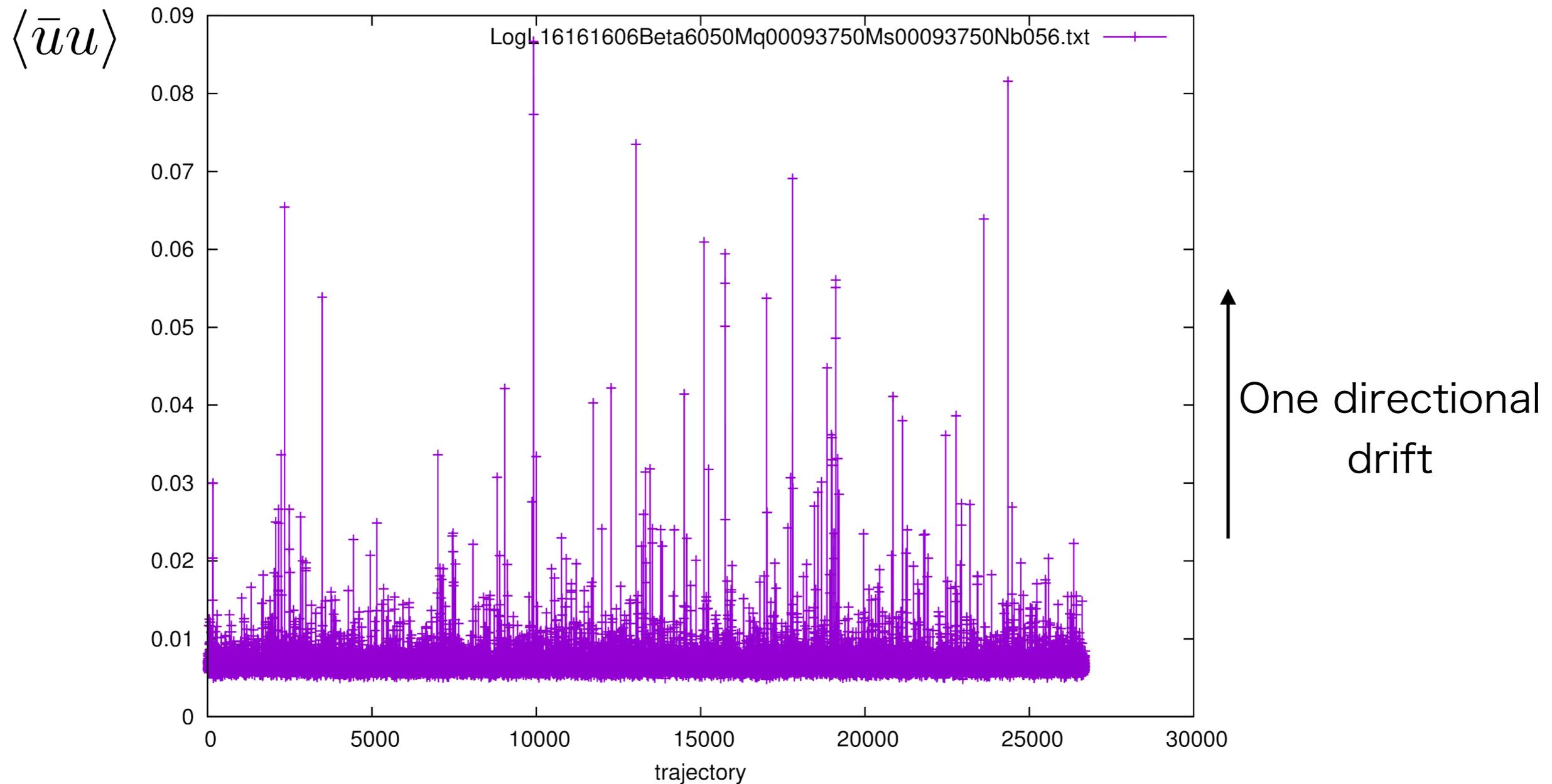


Tc goes up and down along with  $N_b$ ??

->  $M_{\pi} = 80$  MeV with  $N_x=16$  is affected by finite size effects.

# Results with $m_\pi = 80$ MeV

Lighter regime for  $N_x = 16$  is affected by finite volume effects



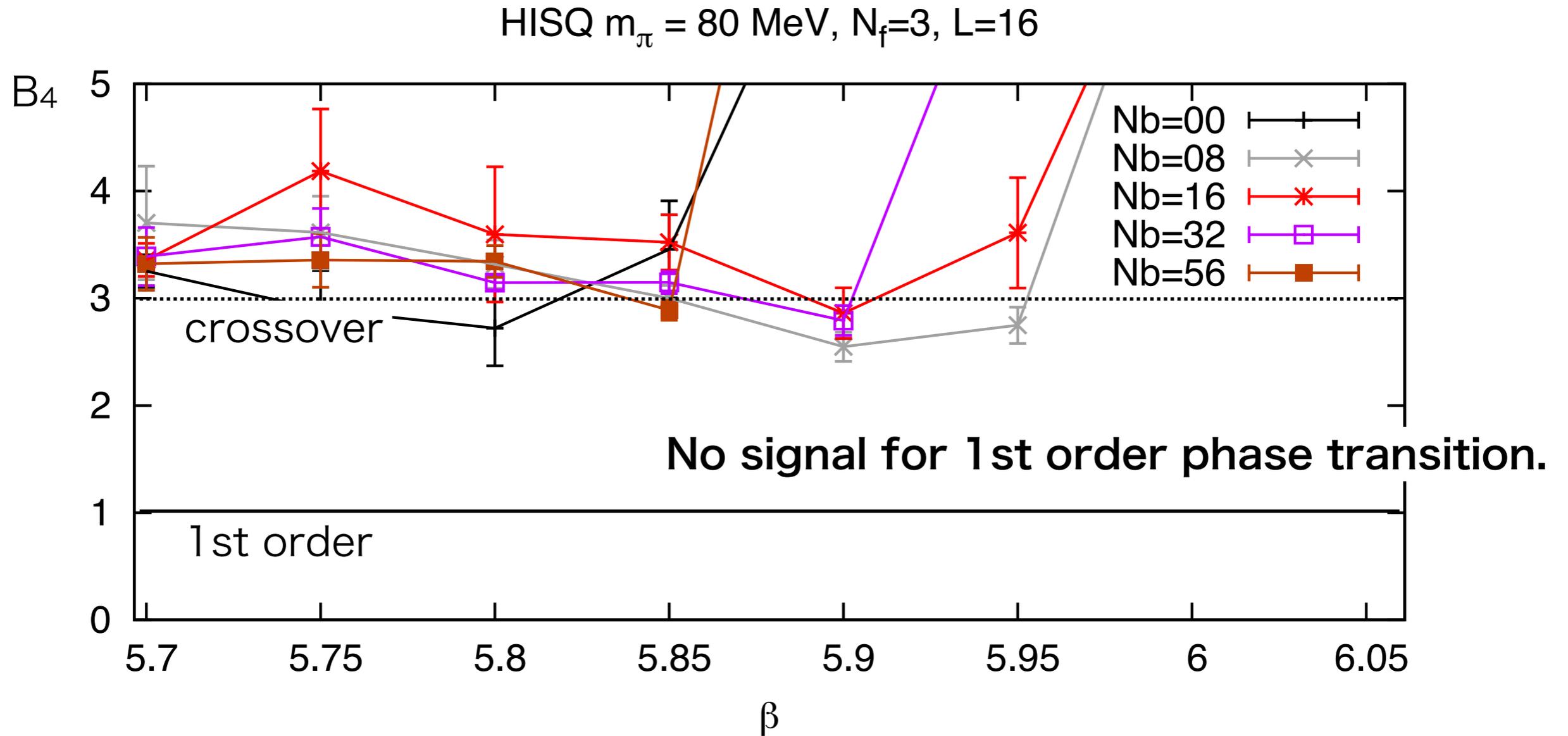
This mass is strongly affected by finite volume effects:

$$M_\pi L = 1.7 < 4$$

Especially, high temperature and large Nb is serious.

# Results with $m_\pi = 80$ MeV

Binder cumulant in light pion mass regime (near the SU(3) chiral limit)



# Summary: QCD phase trs. with magnetic field

No 1st order is found for  $m_\pi = 320$  and  $80$  MeV ( $\sqrt{eB} < 870$  MeV)

## Summary of HISQ with magnetic field

1.  $m_\pi = 320$  MeV shows crossover,  $T_c(eB)$  decreases.
  1. 1st order for standard staggered fermions are probably an artifact
  2. Dirac spectrum for sea part are consistent with inverse magnetic catalysis
  3. Qualitatively similar to the physical point results by stout staggered fermions
2.  $m_\pi = 80$  MeV shows crossover ( $\sqrt{eB} < 870$  MeV)
  1.  $T_c$  goes up and down along with  $N_b$ , but  $N_x=16$  lattice affected by finite volume effect, especially high temperature and large  $N_b$ . (highly suspected)
  2. Except for  $T_c(eB)$  behavior,  $m_\pi = 80$  MeV system shows qualitatively same behavior to  $320$  MeV (not showed in this talk)
  3. No indication of the first order
3. In both cases, topological susceptibility behave similar to the chiral condensate. ( $m_\pi = 80$  MeV is not showed)

## Outlook

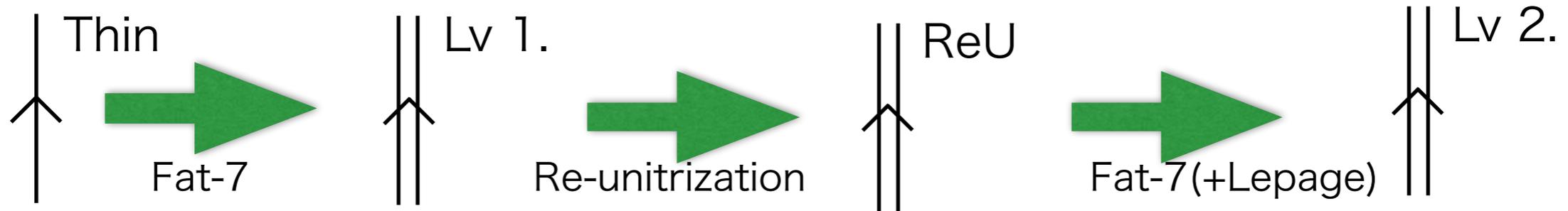
1. To improve the analyses for  $m_\pi = 80$  MeV, simulations with larger volumes are needed. Analyses on  $N_x=24$  is ongoing.



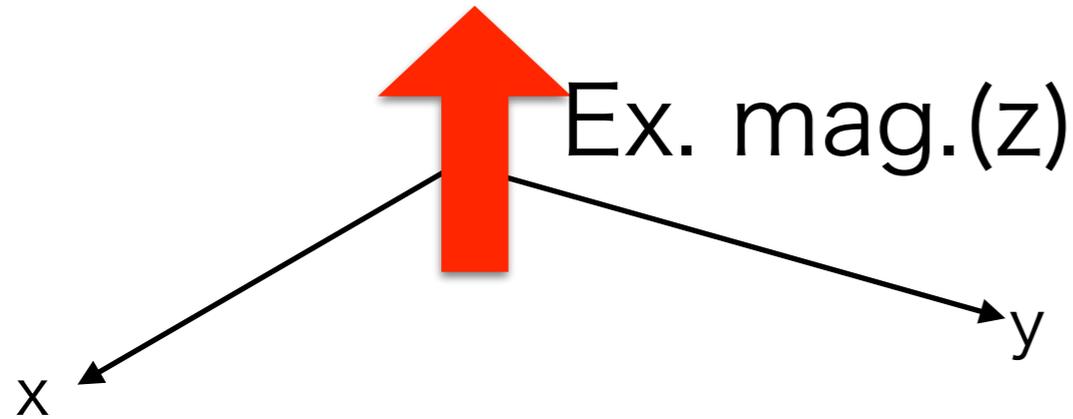
**Backup**

# HISQ with the magnetic field

HISQ = Highly improved staggered quarks



$$D_{\text{HISQ}}[U] = D_{\text{KS}} \left[ \begin{array}{c} \text{Lv 2.} \\ \uparrow \\ \parallel \end{array} \right] + D_{\text{Naik}} \left[ \begin{array}{c} \text{ReU} \\ \uparrow \\ \parallel \end{array} \right]$$



$$u_x(n_x, n_y, n_z, n_t) = \begin{cases} \exp[-iqBN_x n_y] & (n_x = N_x - 1) \\ 1 & (\text{Otherwise}) \end{cases}$$

$$u_y(n_x, n_y, n_z, n_t) = \exp[iqBn_x],$$

Multiplying  $u(1)$  links **after** the smearing

$$D_{\text{KS}} \left[ \mathbf{u} \begin{array}{c} \text{Lv 2.} \\ \uparrow \\ \parallel \end{array} \right] + D_{\text{Naik}} \left[ \mathbf{u} \begin{array}{c} \text{ReU} \\ \uparrow \\ \parallel \end{array} \right]$$

# Topological susceptibility

## Previous works by HISQ

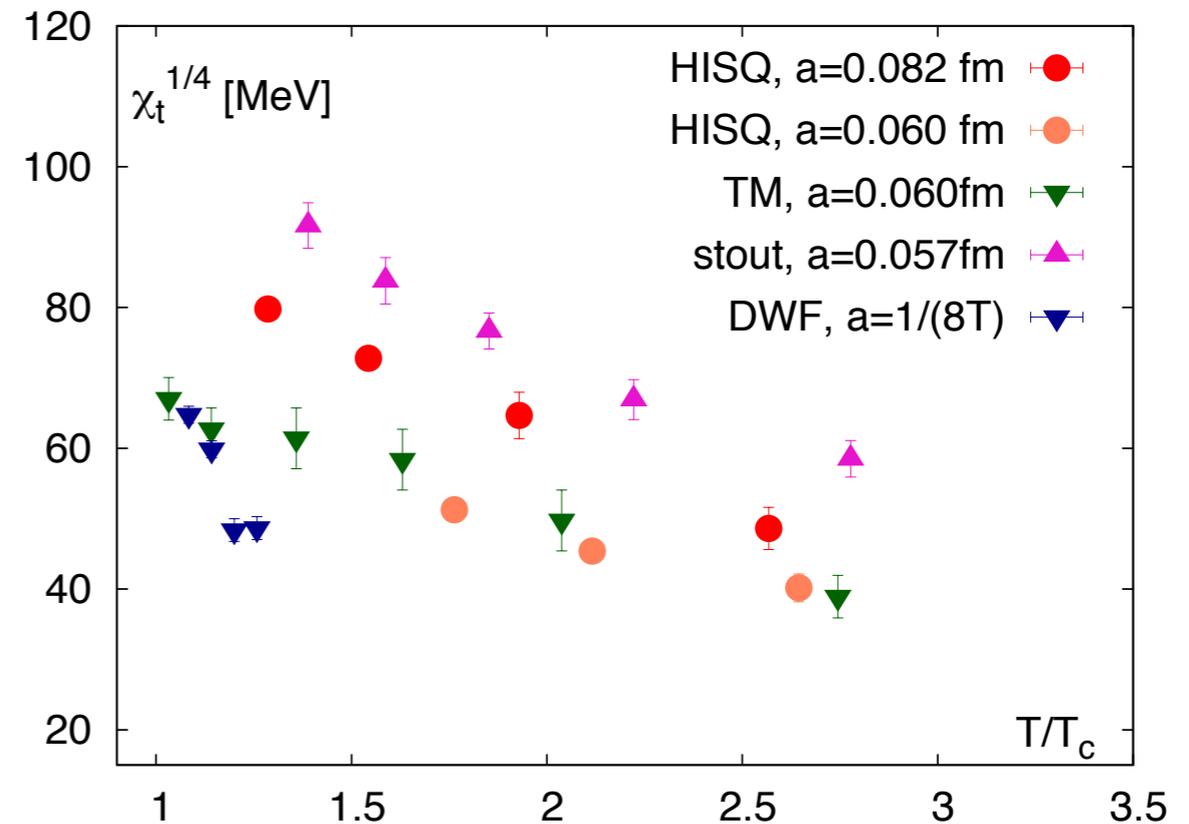
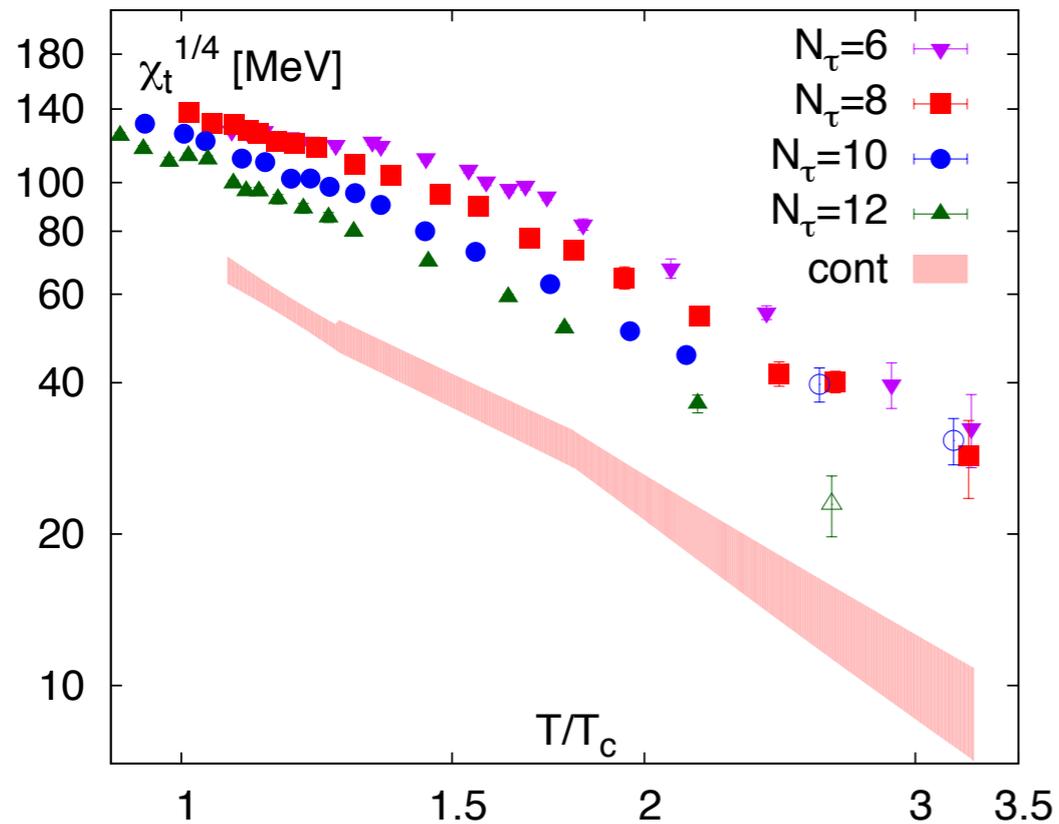


Figure 3: The temperature dependence of  $\chi_t$  in QCD for HISQ action on lattices with different  $N_\tau$  (left) and  $\chi_t$  for HISQ action at two lattice spacings compared with recent results using different fermion actions [19, 36, 49] (right). In the left panel, we also show the continuum result for  $\chi_t$  discussed in section 4 and open symbols represent the data points that have not been used in the continuum extrapolation.

Peter Petreczky, Hans-Peter Schadler and Sayantan Sharma

1606.03145

# Comparison to conventional method

The projector method gives consistent results to the chiral condensate

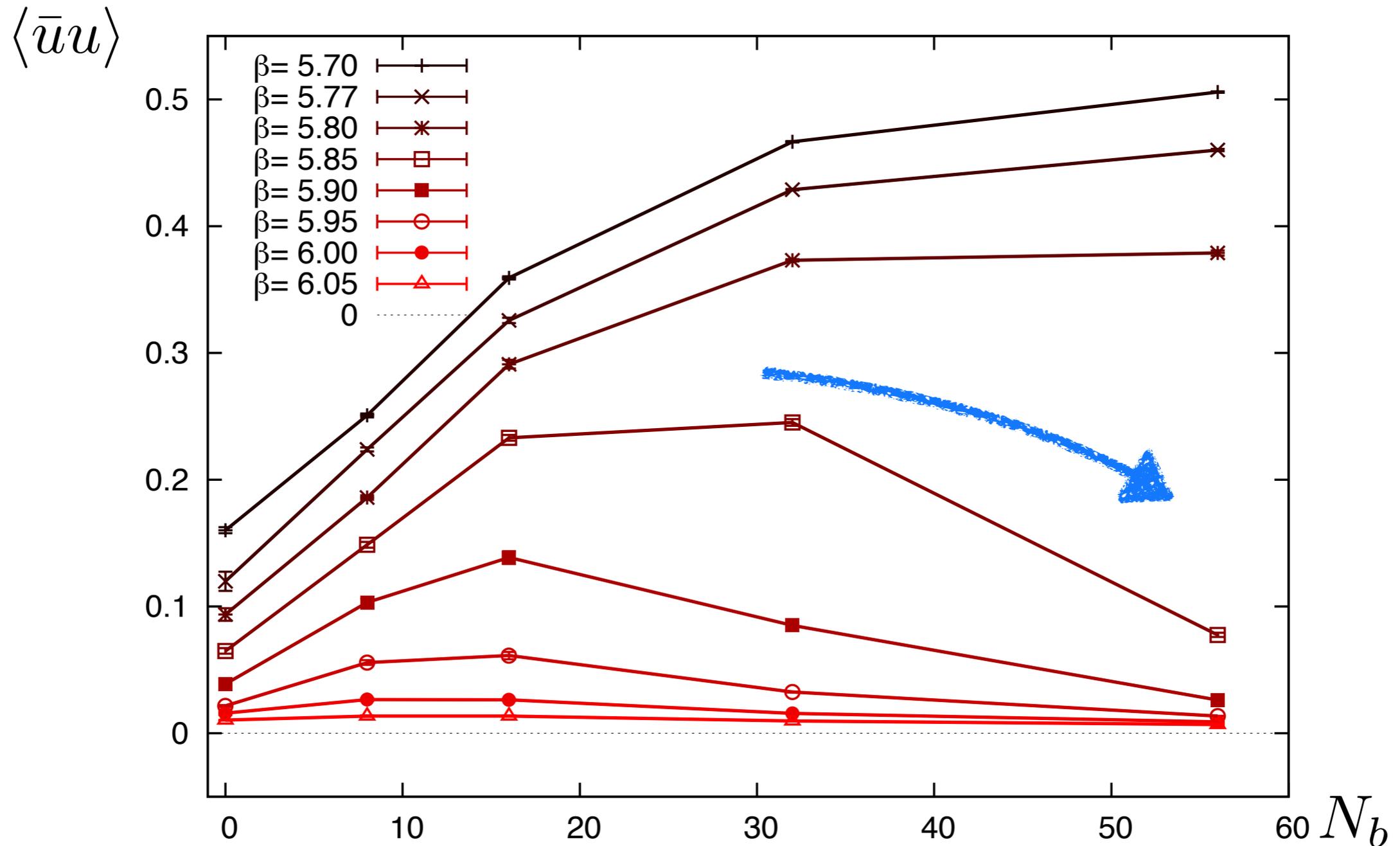
We have used stochastic estimator for the Dirac spectrum and calculate the chiral condensate

$$\begin{aligned}
 |\langle \bar{\psi}\psi \rangle| &= \int_0^\infty d\lambda \frac{2m\rho(\lambda)}{\lambda^2 + m^2} \\
 &= \frac{1}{V} \left\langle \text{Tr} \left[ \frac{1}{D + m} \right] \right\rangle \approx \frac{1}{N_r N_{\text{conf}}} \sum_c^{N_{\text{conf}}} \sum_r^{N_r} \xi_r^\dagger \frac{1}{D[U_c] + m} \xi_r
 \end{aligned}$$

Using stochastic estimator [L. Giusti et al, 2009. G. Cossu et al 2016. P. d. Forcrand et al 2018.]

# Results with $m_\pi = 80$ MeV

## Lighter regime



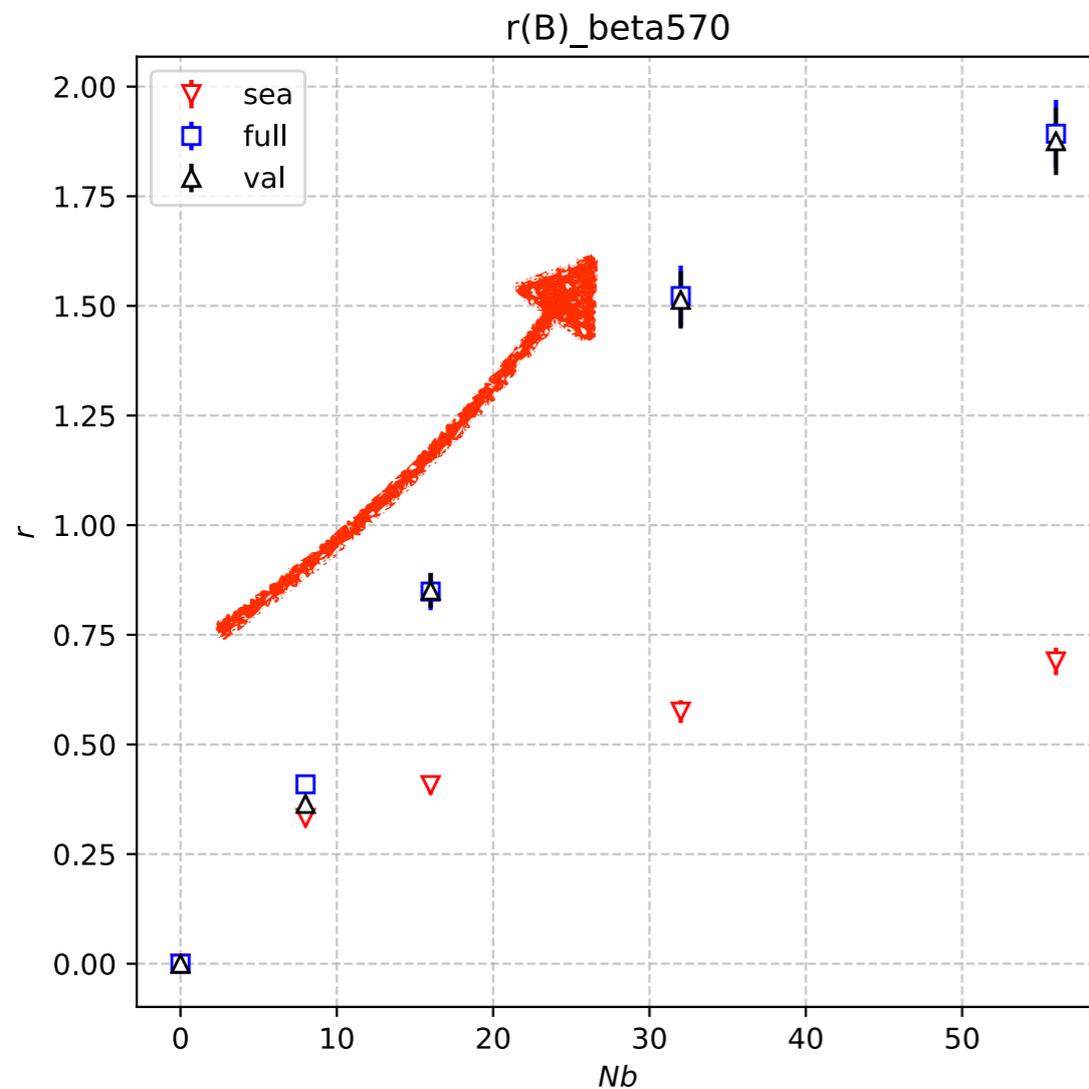
It shows inverse mag. catalysis

Using stochastic estimator [L. Giusti et al, 2009. G. Cossu et al 2016. P. d. Forcrand et al 2018.]

# Results with $m_\pi = 80$ MeV

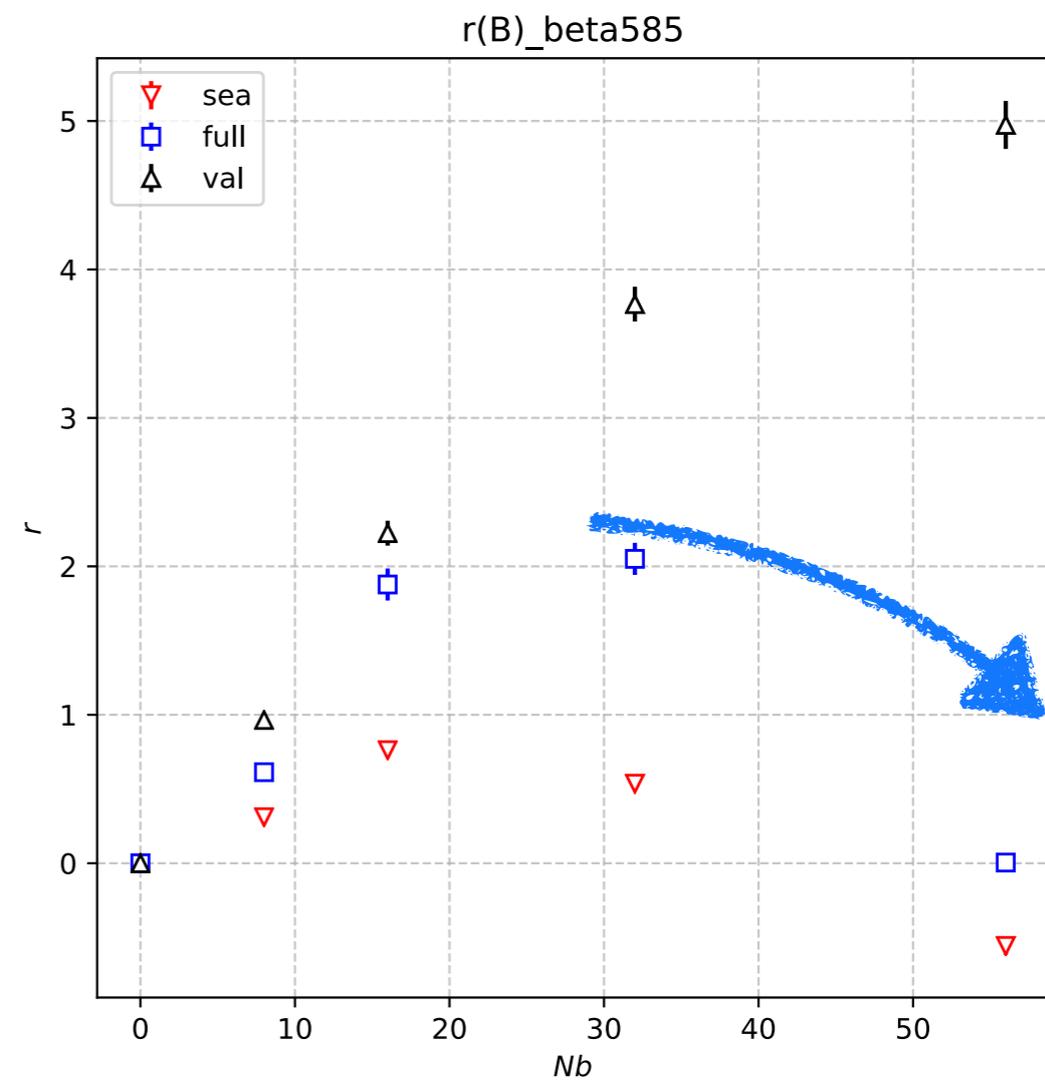
## Lighter regime

### $T < T_c$



(Normal) Magnetic Catalysis

### $T \sim T_c$



Inverse Magnetic Catalysis

**Sea quarks drive inverse catalysis**  
**Similar to previous results at the physical point**

Using stochastic estimator [L. Giusti et al, 2009. G. Cossu et al 2016. P. d. Forcrand et al 2018.]

# Results with $m_{\pi} = 320$ MeV

## Effects from valence and sea Dirac operators

External magnetic field come into two slots,

$$\begin{aligned} \langle \bar{\psi}\psi \rangle^{\text{full},f}(B) &= \int \mathcal{D}UP[U; B] \text{Tr} \left[ \frac{1}{D^f[U; B] + m} \right], \\ \langle \bar{\psi}\psi \rangle^{\text{val},f}(B) &= \int \mathcal{D}UP[U; 0] \text{Tr} \left[ \frac{1}{D^f[U; B] + m} \right], \\ \langle \bar{\psi}\psi \rangle^{\text{sea},f}(B) &= \int \mathcal{D}UP[U; B] \text{Tr} \left[ \frac{1}{D^f[U; 0] + m} \right], \end{aligned}$$

where

$$P[U; B] = \frac{1}{Z(B)} e^{-S_g[U]} \text{Det} [D^{\text{up}}[U; B] + m]^{1/4} \text{Det} [D^{\text{down}}[U; B] + m]^{1/2}$$

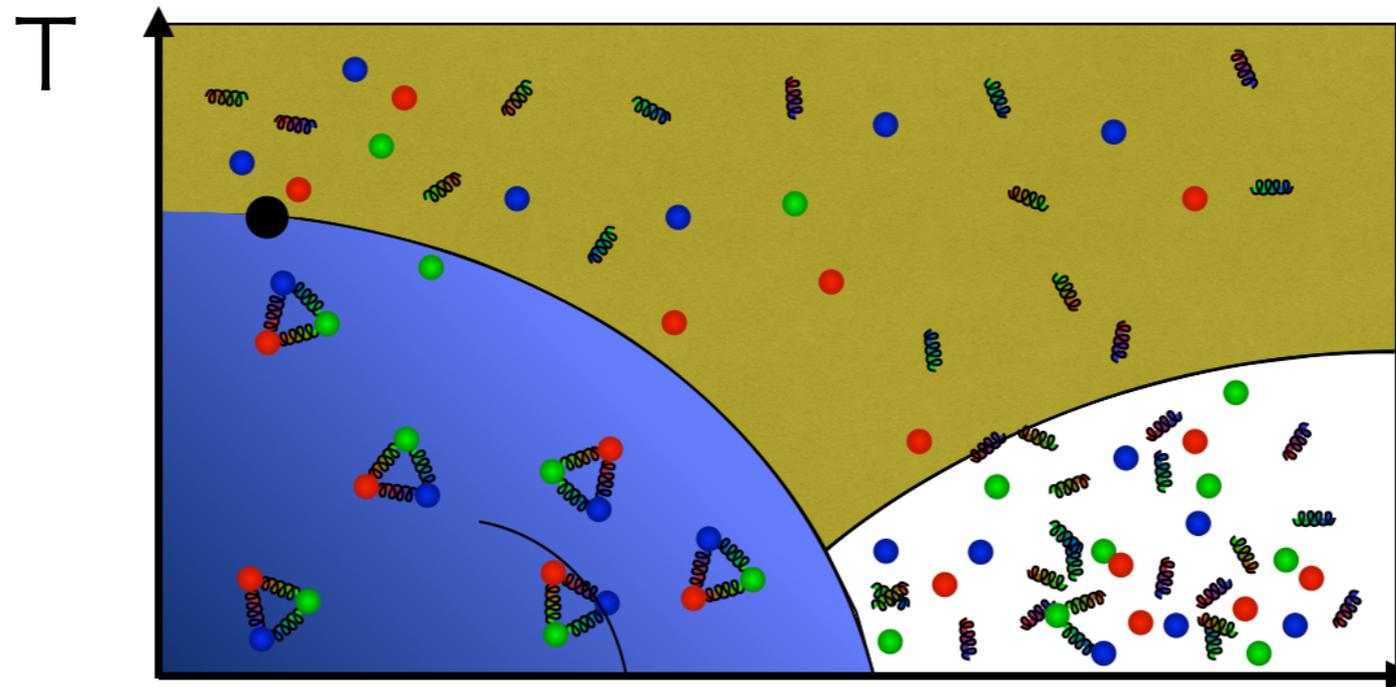
f = up, down

[M. D'Elia *et al* 1103.2080]

# QCD phase diagram with magnetic field

Works with external magnetic field (improved KS types)

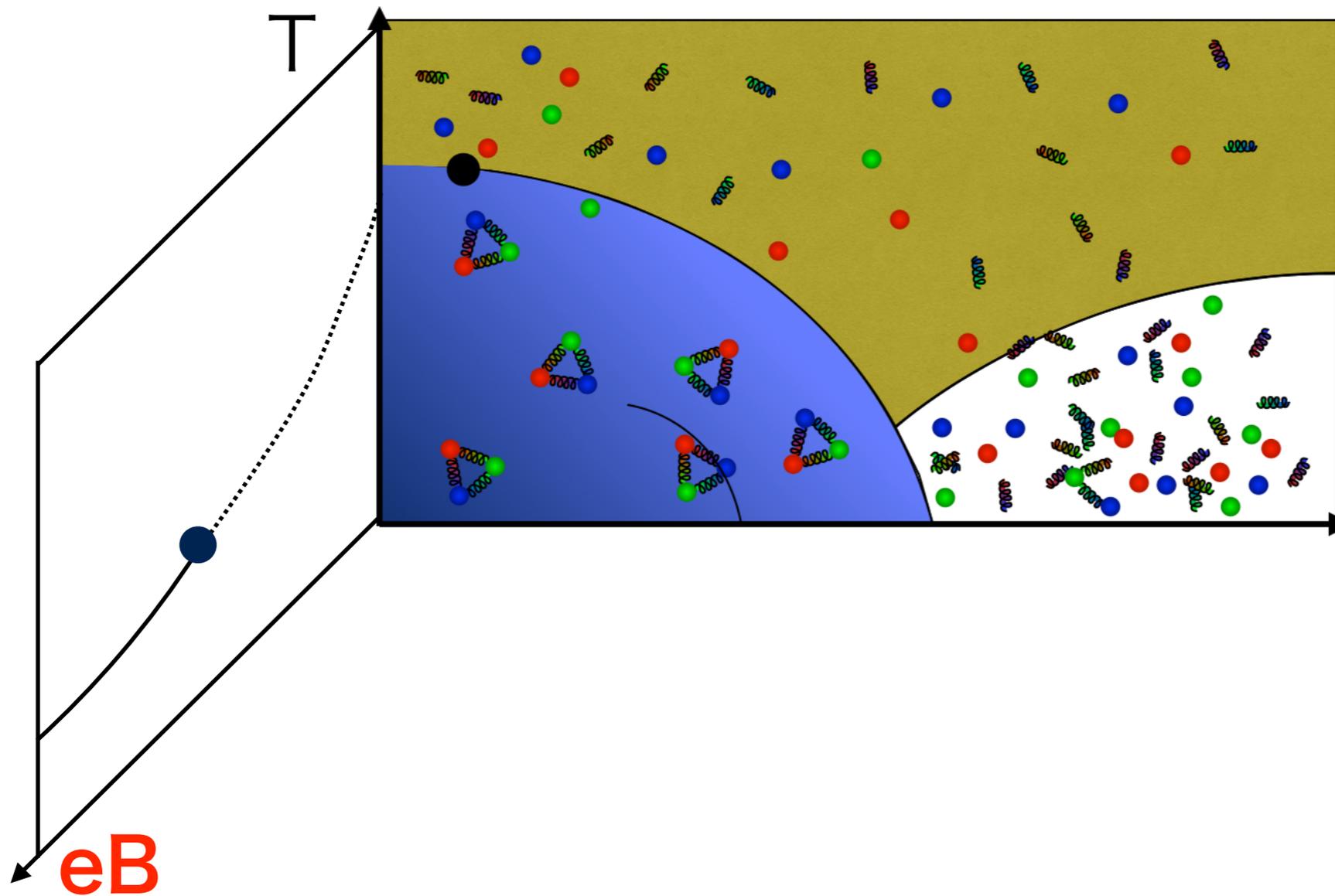
Phase diagram @  $m_\pi = \text{Physical mass}$



# QCD phase diagram with magnetic field

Works with external magnetic field (improved KS types)

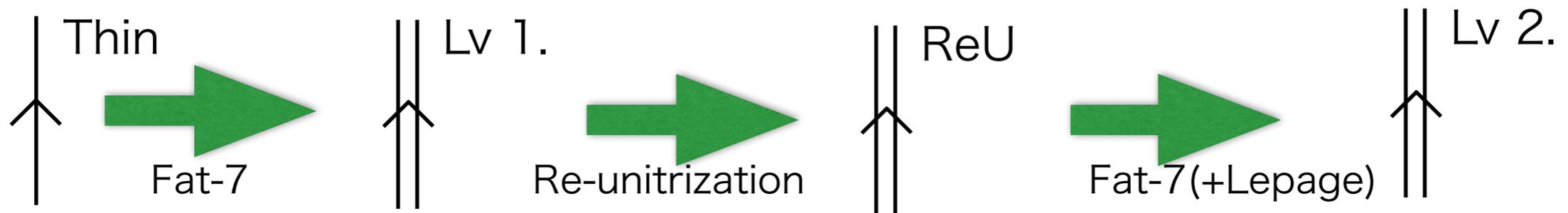
Phase diagram @  $m_\pi = \text{Physical mass}$



# HISQ with the magnetic field

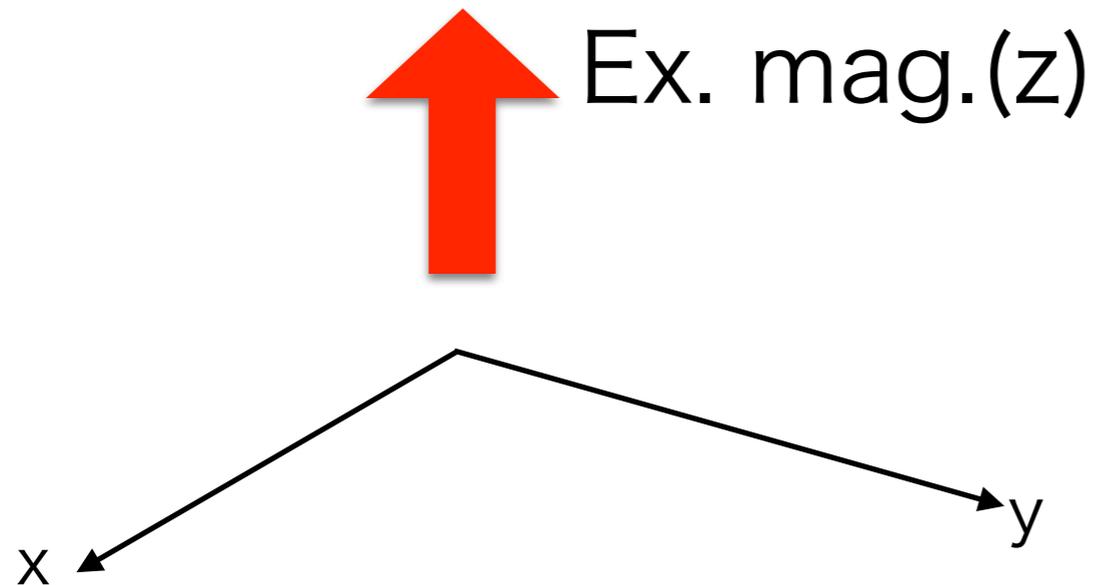
HISQ = Hight improved staggered quarks

HISQ = KS term with 2 level smearing + Naik term with reunitarized links



# Quantization of the magnetic field

Magnetic flux and upper limit.



External magnetic field for z-dir.

$$u_x(n_x, n_y, n_z, n_t) = \begin{cases} \exp[-iqBN_x n_y] & (n_x = N_x - 1) \\ 1 & (\text{Otherwise}) \end{cases}$$

$$u_y(n_x, n_y, n_z, n_t) = \exp[iqBn_x],$$

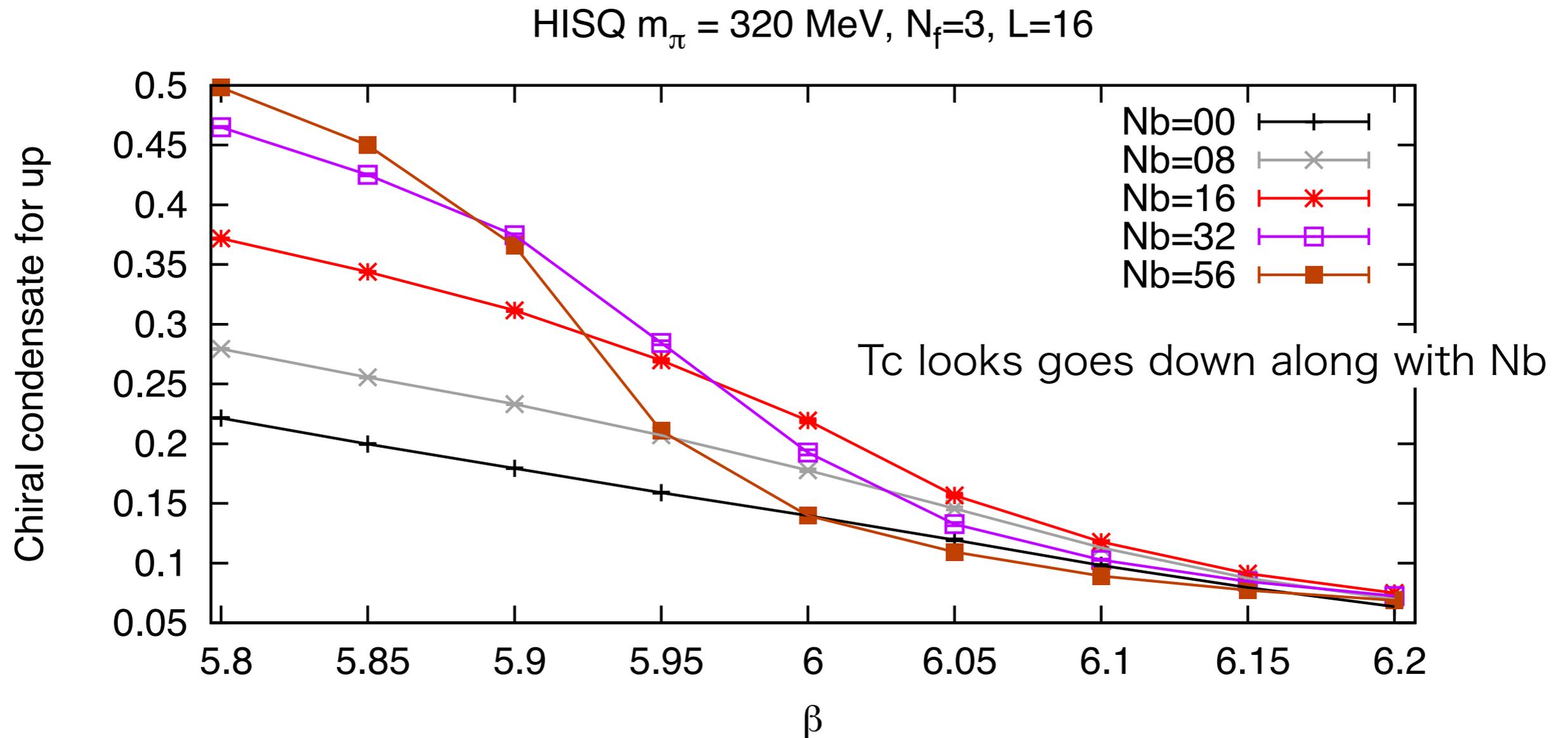
$$eB = \frac{2\pi N_b}{N_x N_y} a^{-2}$$

$$0 \leq N_b < \frac{N_x N_y}{4} \quad (\text{By a periodicity})$$

$N_x = N_y = 16$  case,  $N_b = 64$  is the maximum

# Results with $m_{\pi} = 320$ MeV

Chiral condensate for heavier regime ( $\sim$  standard setup)



# Results with $m_{pi} = 320$ MeV

Dirac spectrum sea quarks drive inverse catalysis

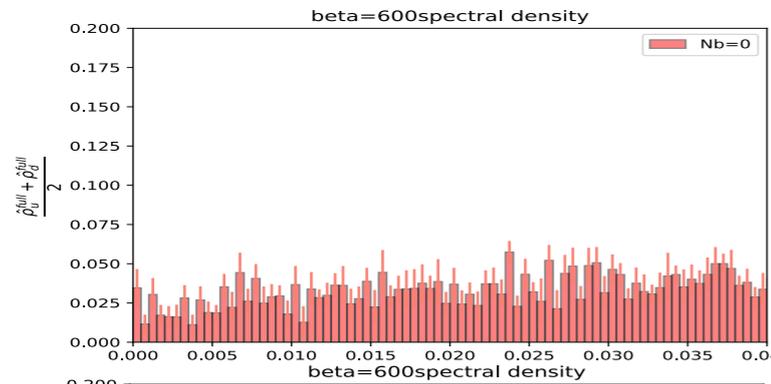
$\beta = 6$  ( $T > T_c$ )  
Nb

Full

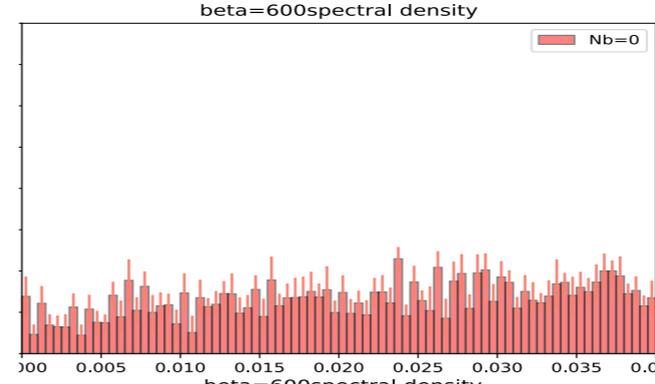
Sea

Valence

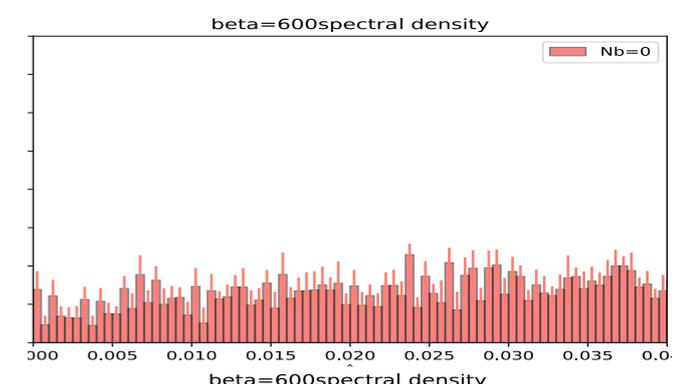
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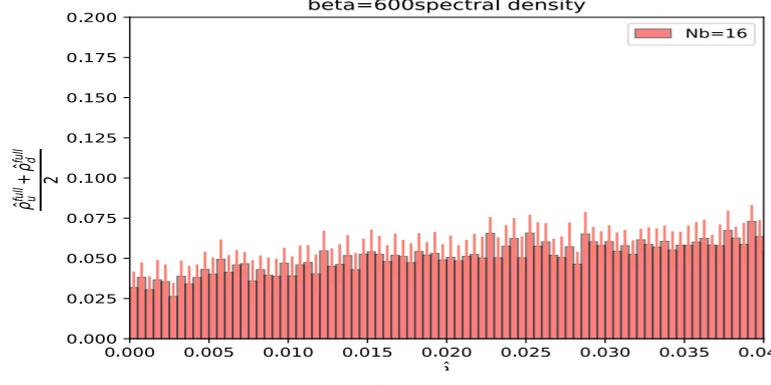
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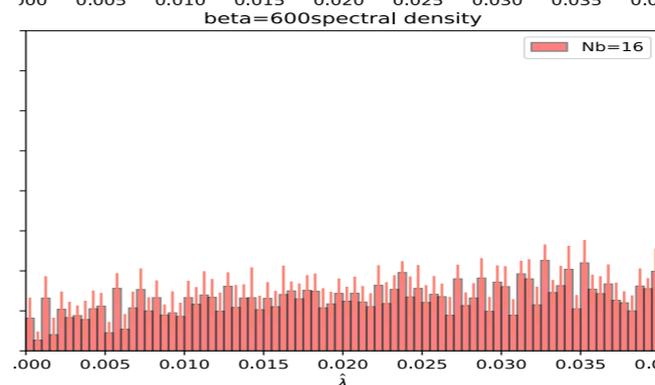
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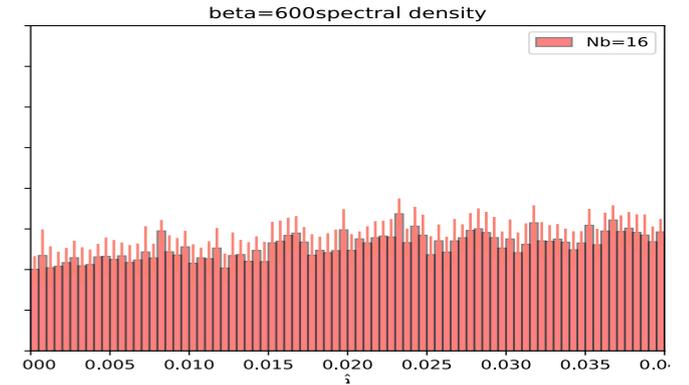
16



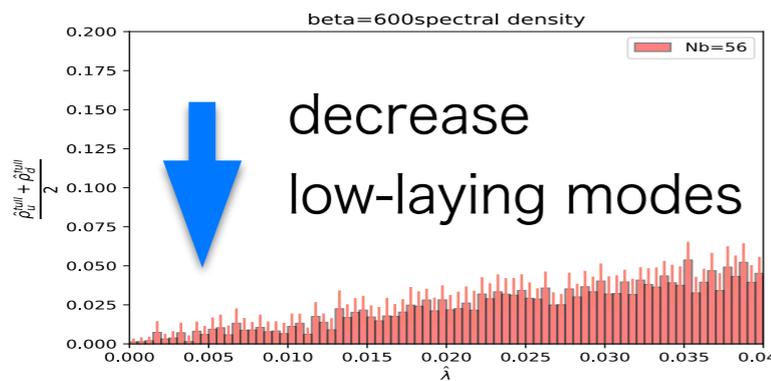
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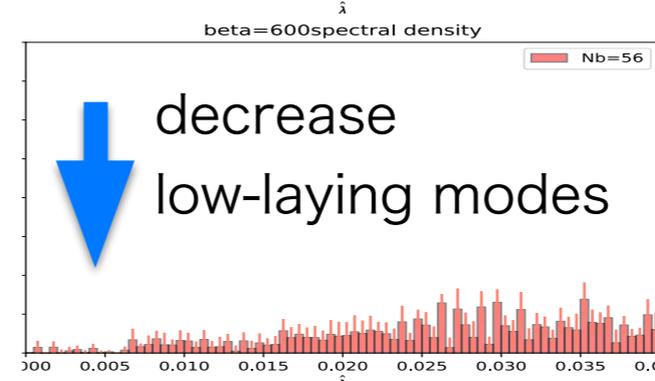
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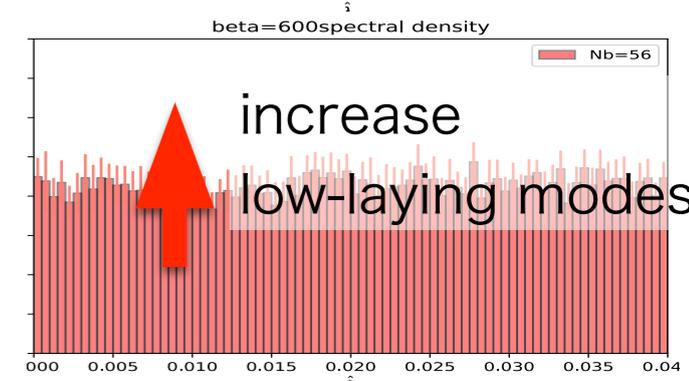
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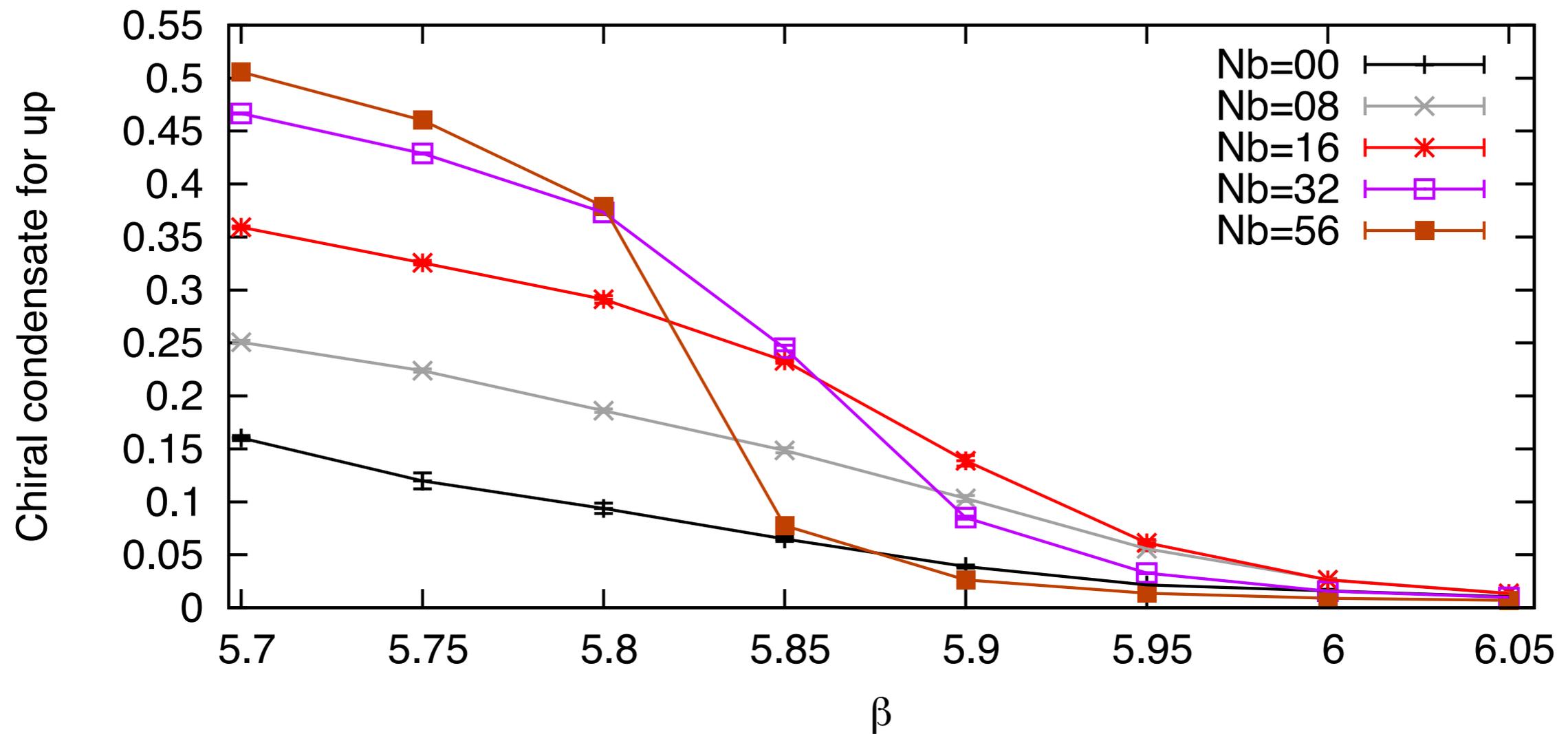
Valence quarks only contribute to normal catalysis

Using stochastic estimator [L. Giusti et al, 2009. G. Cossu et al]

# Results with $m_{\pi} = 80$ MeV

Chiral cond. in light pion mass regime (near the SU(3) chiral limit)

HISQ  $m_{\pi} = 80$  MeV,  $N_f=3$ ,  $L=16$



# Results with $m_{\pi} = 80 \text{ MeV}$

Dirac spectrum for lighter regime

$$\beta = X \quad (T > T_c)$$

Nb                      Full                      Sea                      Valence

0

16

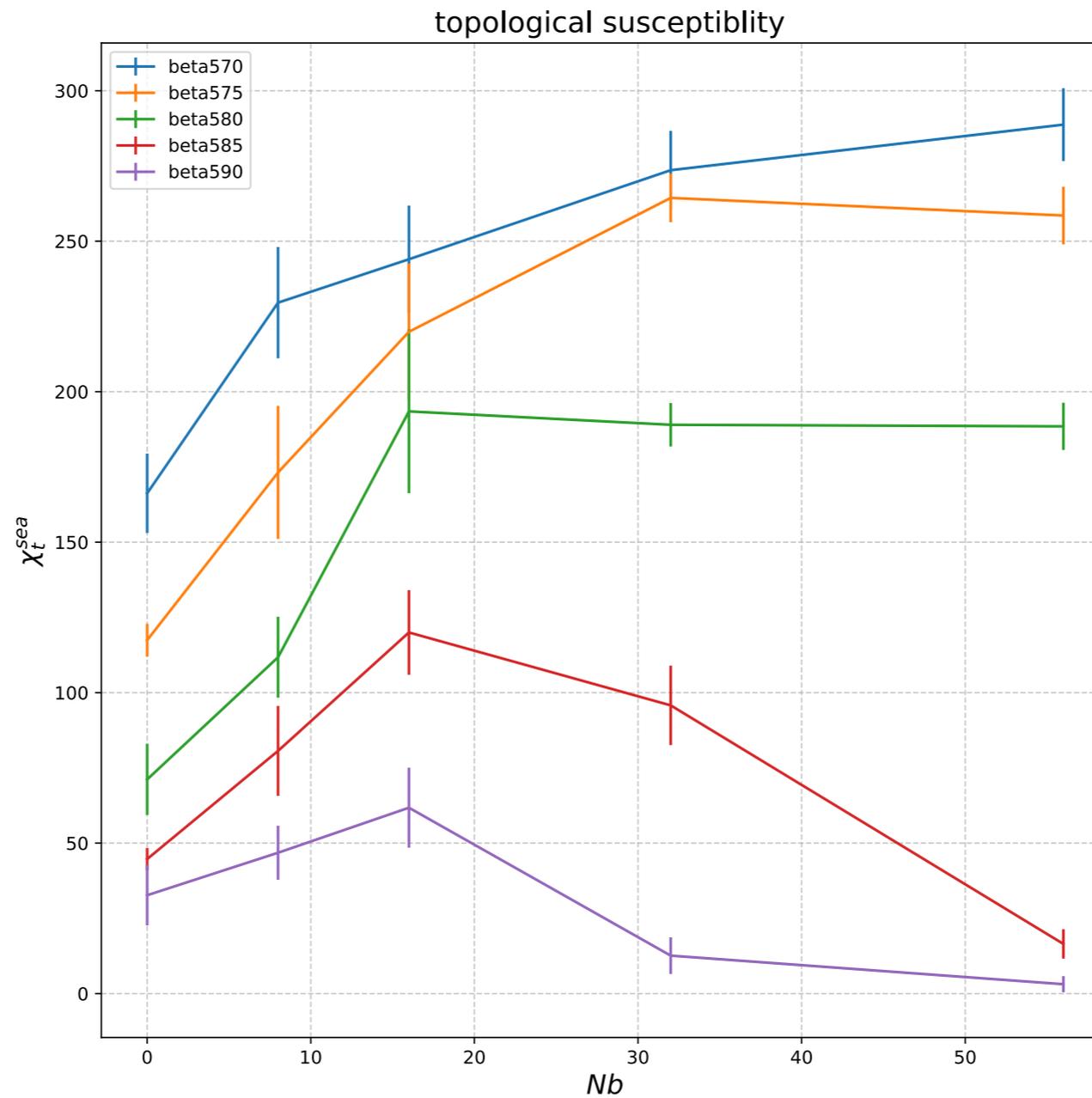
56

Valence only contribute to normal catalysis

Using stochastic estimator [L. Giusti et al, 2009. G. Cossu et al 2010]

# Results with $m_{\pi} = 80$ MeV

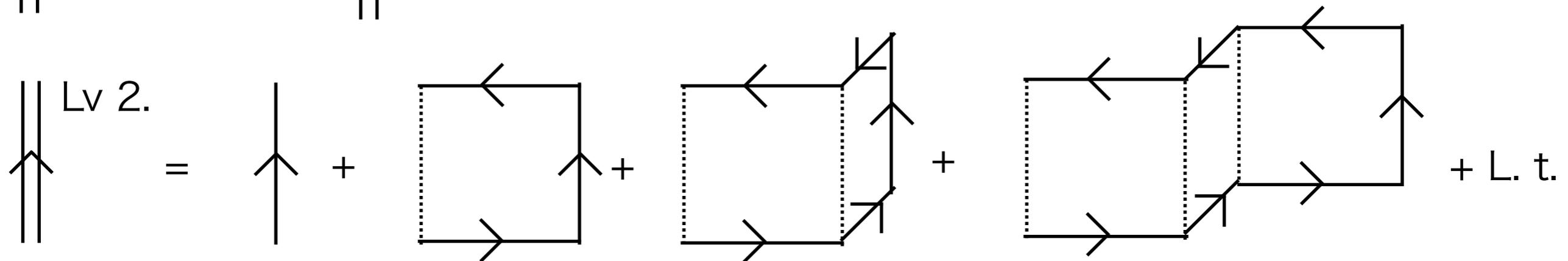
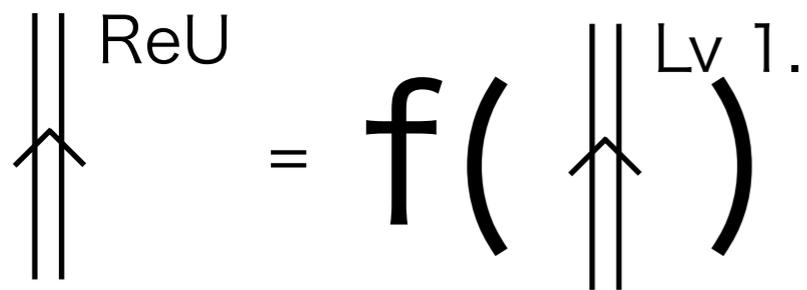
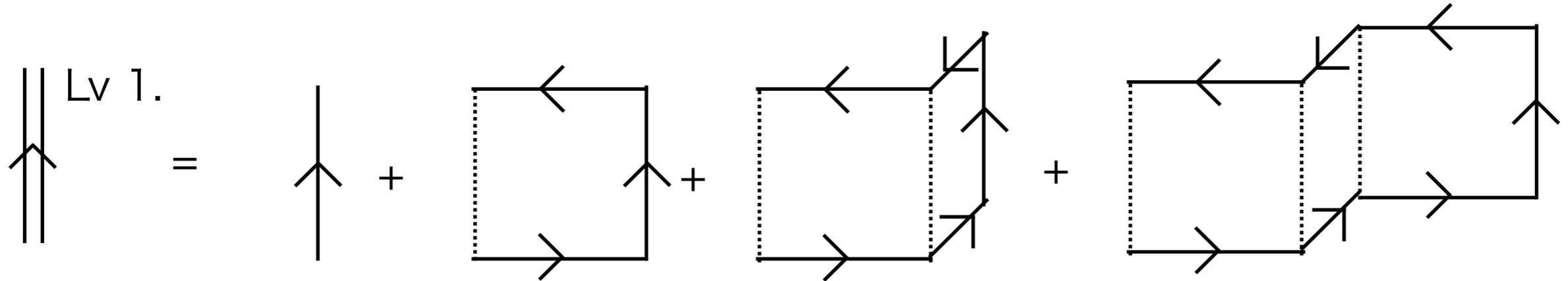
## Topological susceptibility for lighter regime



# HISQ with the magnetic field

HISQ = Hight improved staggered quarks

HISQ = KS term with 2 level smearing + Naik term with reunitarized links



$$D_{\text{HISQ}}[U] = D_{\text{KS}} \left[ \begin{array}{c} \text{Lv 2.} \\ \text{double line with staggered links} \end{array} \right] + D_{\text{Naik}} \left[ \begin{array}{c} \text{ReU} \\ \text{double line with staggered links} \end{array} \right]$$

# Related works

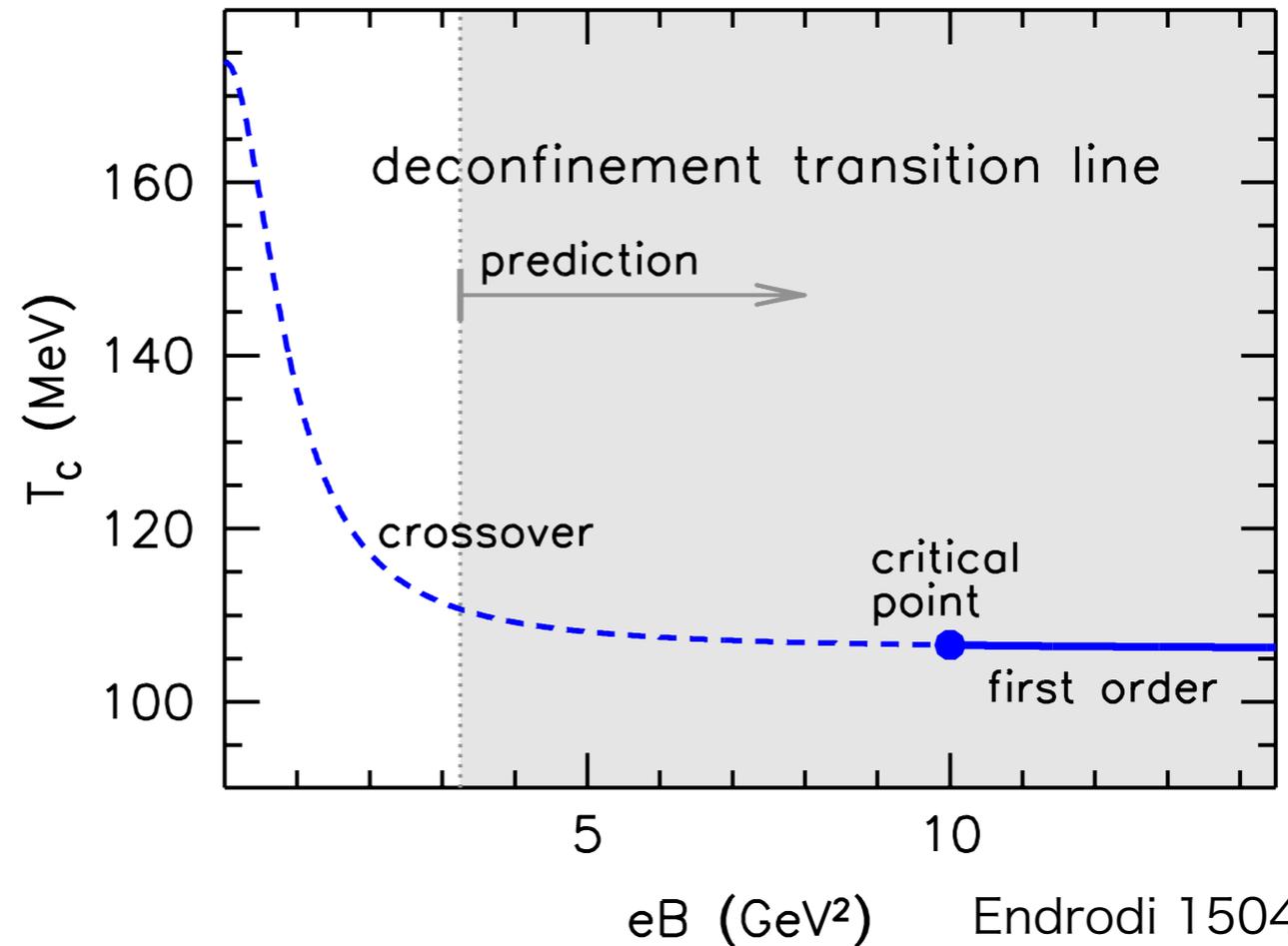
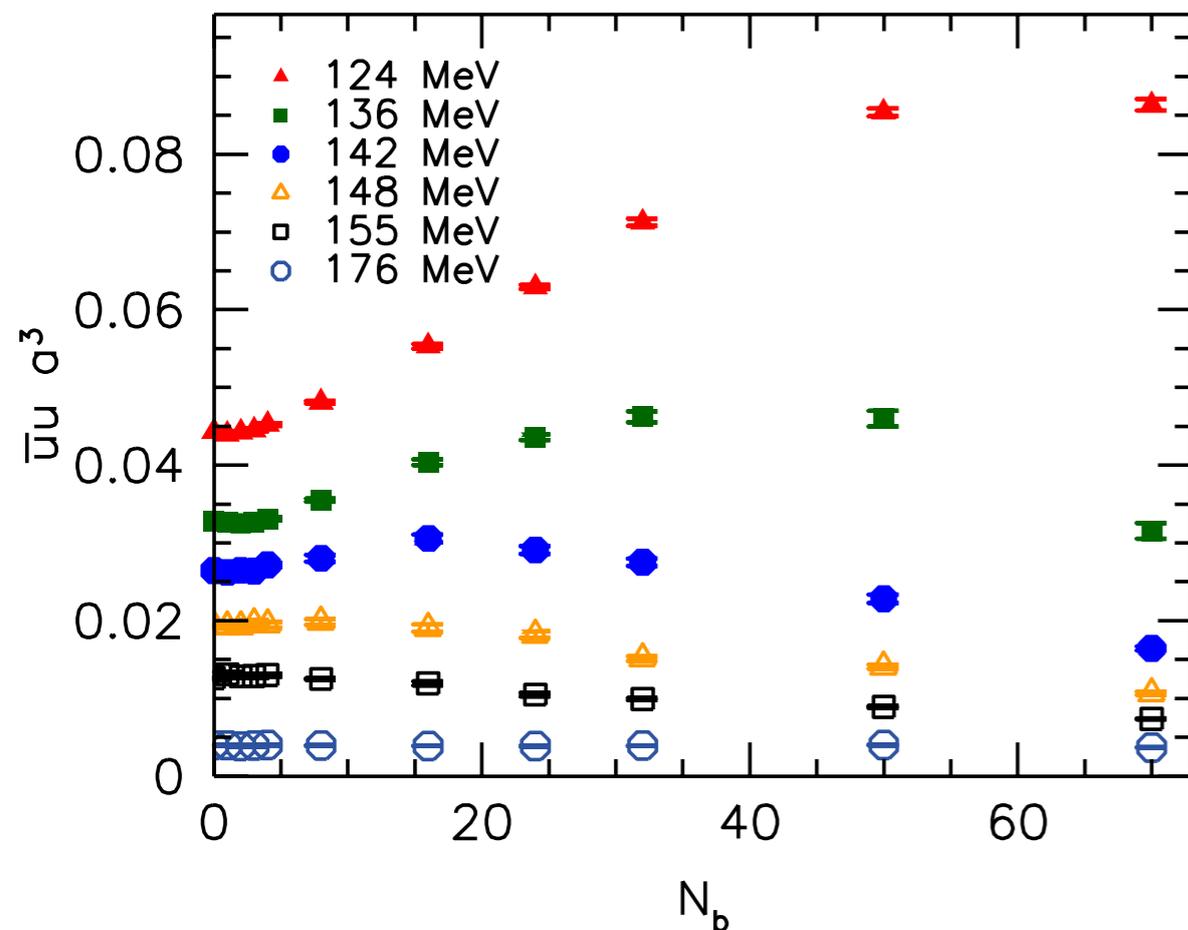
$M_{\pi} \sim 300 \text{ MeV}$

Our last year's one

- $T_c$  goes up
- normal catalysis
- First order??

# Related works

Mpi = physical, F. Bruckman et al, arXiv: 1111.4956, 1303.3972

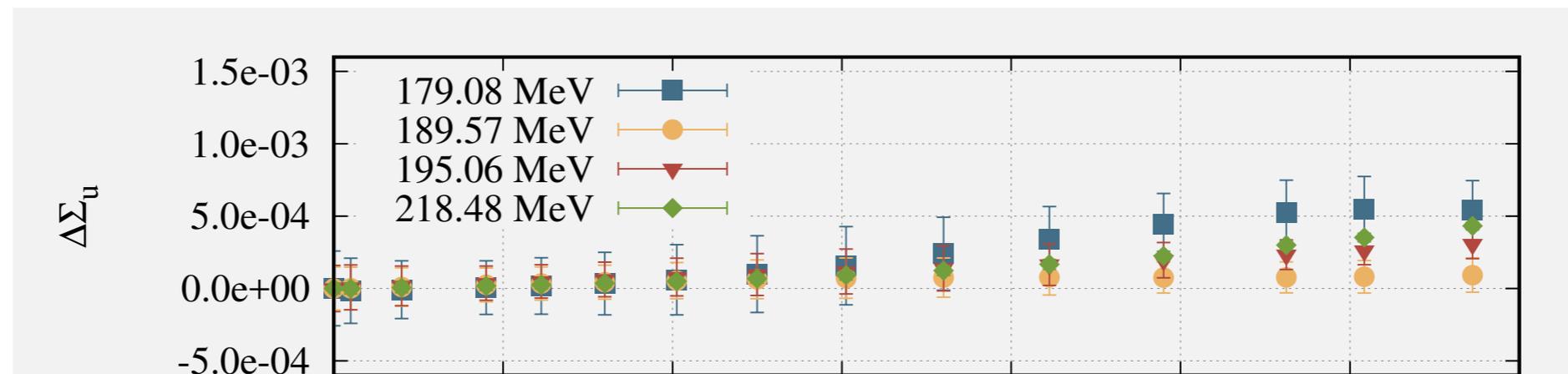
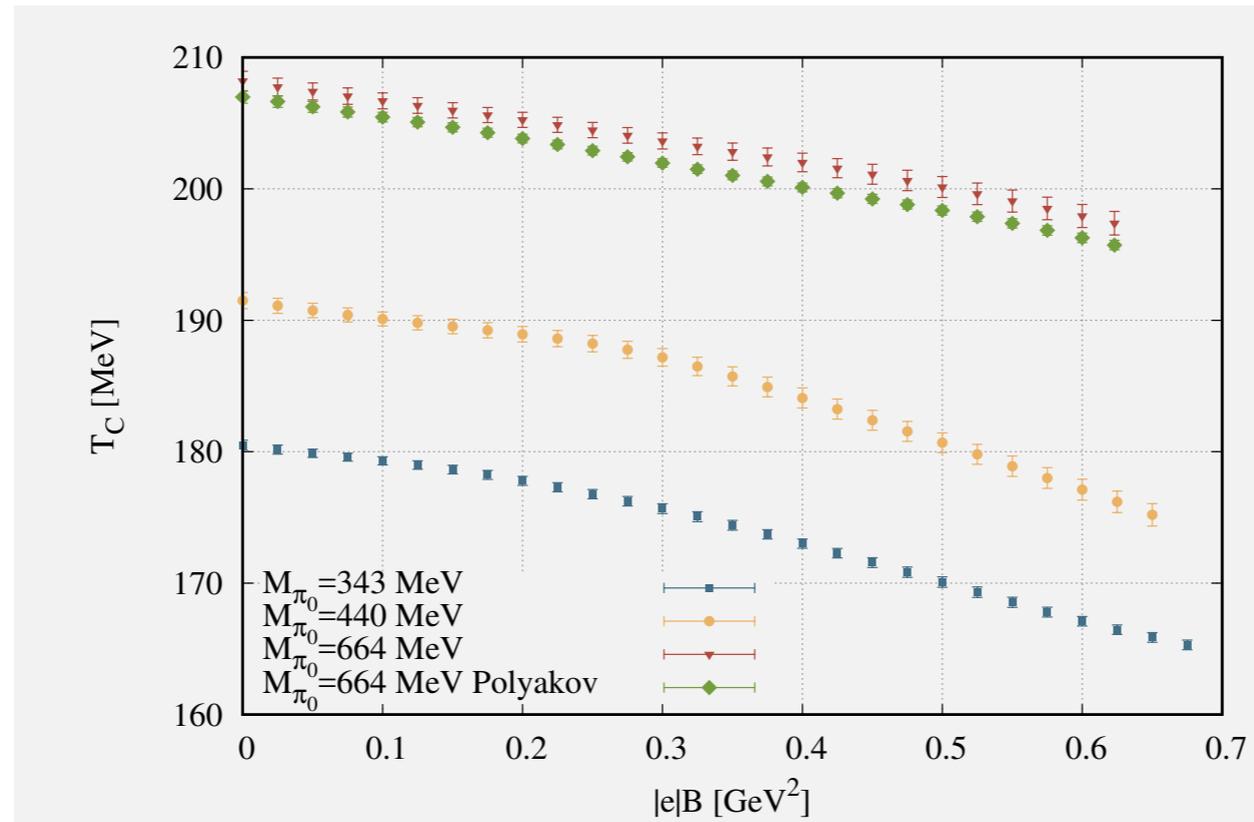


- $T_c$  goes down
- Inverse magnetic Catalysis around  $T_c$
- First order at extremely strong  $eB$  (Conjecture)

Endrodi 1504.08280

# Related works

$M_{\pi} = 600$  MeV, Massimo, XQCD2018



- $T_c$  goes down
- No Inverse magnetic Catalysis for the range

# Statistics with $m_{\pi} = 320$ MeV

## Heavier regime

L beta mass Nb Trajectories

16 5.80 0.015 00 5500  
 16 5.80 0.015 08 5000  
 16 5.80 0.015 16 5000  
 16 5.80 0.015 32 5000  
 16 5.80 0.015 56 5000  
 16 5.85 0.015 00 6500  
 16 5.85 0.015 08 6500  
 16 5.85 0.015 16 6500  
 16 5.85 0.015 32 6500  
 16 5.85 0.015 56 6500  
 16 5.90 0.015 00 7500  
 16 5.90 0.015 08 7500  
 16 5.90 0.015 16 6500  
 16 5.90 0.015 32 7500  
 16 5.90 0.015 56 5466  
 16 5.95 0.015 00 7500  
 16 5.95 0.015 08 7500  
 16 5.95 0.015 16 4500  
 16 5.95 0.015 32 7500  
 16 5.95 0.015 56 7500

L beta mass Nb Trajectories

16 6.00 0.015 00 4500  
 16 6.00 0.015 08 4500  
 16 6.00 0.015 16 4500  
 16 6.00 0.015 32 4500  
 16 6.00 0.015 56 4500  
 16 6.05 0.015 00 4000  
 16 6.05 0.015 08 4000  
 16 6.05 0.015 16 4000  
 16 6.05 0.015 32 4000  
 16 6.05 0.015 56 4000  
 16 6.10 0.015 00 6500  
 16 6.10 0.015 08 6500  
 16 6.10 0.015 16 6500  
 16 6.10 0.015 32 6500  
 16 6.10 0.015 56 6500  
 16 6.15 0.015 00 4170  
 16 6.15 0.015 08 6500  
 16 6.15 0.015 16 5500  
 16 6.15 0.015 32 5500  
 16 6.15 0.015 56 5500

L beta mass Nb Trajectories

16 6.20 0.015 00 4485  
 16 6.20 0.015 08 7500  
 16 6.20 0.015 16 7500  
 16 6.20 0.015 32 7500  
 16 6.20 0.015 56 7500

$m_a = 0.015$

# Statistics with $m_{\pi} = 80 \text{ MeV}$

## Lighter regime

L beta mass Nb Trajectories

16 5.700 0.0009375 00 630  
 16 5.700 0.0009375 08 1637  
 16 5.700 0.0009375 16 1640  
 16 5.700 0.0009375 32 10702  
 16 5.700 0.0009375 56 9638  
 16 5.750 0.0009375 00 700  
 16 5.750 0.0009375 08 1652  
 16 5.750 0.0009375 16 1644  
 16 5.750 0.0009375 32 11289  
 16 5.750 0.0009375 56 9363  
 16 5.800 0.0009375 00 700  
 16 5.800 0.0009375 08 1721  
 16 5.800 0.0009375 16 1467  
 16 5.800 0.0009375 32 12928  
 16 5.800 0.0009375 56 11486  
 16 5.850 0.0009375 00 700  
 16 5.850 0.0009375 08 1734  
 16 5.850 0.0009375 16 1623  
 16 5.850 0.0009375 32 11693  
 16 5.850 0.0009375 56 8512

L beta mass Nb Trajectories

16 5.900 0.0009375 00 700  
 16 5.900 0.0009375 08 2329  
 16 5.900 0.0009375 16 2327  
 16 5.900 0.0009375 32 8002  
 16 5.900 0.0009375 56 8276  
 16 5.950 0.0009375 00 700  
 16 5.950 0.0009375 08 1840  
 16 5.950 0.0009375 16 2337  
 16 5.950 0.0009375 32 13374  
 16 5.950 0.0009375 56 11990  
 16 6.000 0.0009375 00 700  
 16 6.000 0.0009375 08 3603  
 16 6.000 0.0009375 16 3700  
 16 6.000 0.0009375 32 13098  
 16 6.000 0.0009375 56 11690  
 16 6.050 0.0009375 00 700  
 16 6.050 0.0009375 08 3700  
 16 6.050 0.0009375 16 3700  
 16 6.050 0.0009375 32 14769  
 16 6.050 0.0009375 56 14010

$m_a = 0.0009375$

# $eB^4$ corrections

sea - valence decomposition can be justified below  $N_b=16$

$N_x = N_y = 16$  case

$$eB = \frac{2\pi N_b}{N_x N_y} a^{-2}$$

$$r^{\text{full}}(B) = r^{\text{sea}}(B) + r^{\text{val}}(B) + O(B^4)$$

$$\det(D + m)(U, B) = \det(D + m) + c(e\hat{B})^2 + O(\hat{B}^4)$$

$$(D + m)^{-1}(U, B) = (D + m)^{-1}(U) + c'(e\hat{B})^2 + O(\hat{B}^4)$$

Nb	8	16	32	56
$a^2 eB$	0.20	0.39	0.79	1.37
$(a^2 eB)^2$	0.039	0.154	0.617	1.889
$(a^2 eB)^4$	0.001	0.024	0.381	3.569
$(eB)^4/(eB)^2$	4%	15%	62%	189%

↑ NNLO contribution is large

# $eB^4$ corrections

sea - valence decomposition can be justified below  $N_b=16$

$N_x = N_y = 16$  case

