

# High efficiency free electron lasers

C. Emma

IOTA/FAST Collaboration meeting

May 10, 2018

Fermilab

# Presentation Outline

## 1. Physics of tapered FELs

1.1. Motivation

1.2. Review of theory: 1-D, 3-D, and time dependent effects

## 2. Experimental studies

2.1. Past

2.1.1. Fresh-bunch self-seeding experiment (SLAC)

2.1.2. NOCIBUR experiment (BNL-ATF)

2.2. Present

2.2.1. TESSA-266 (Argonne)

2.3. Future

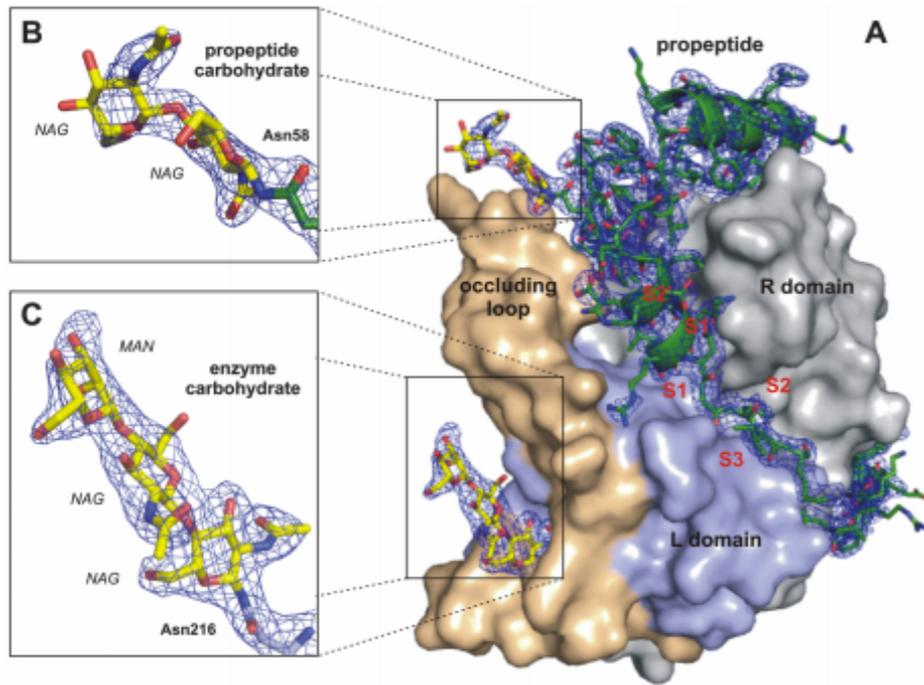
2.3.1. Opportunities at FAST

## 3. Conclusions

# Motivation for high efficiency FEL

## Imaging single molecules via “diffraction before destruction”

Redecke et al., Science 339, 6116, (2012)



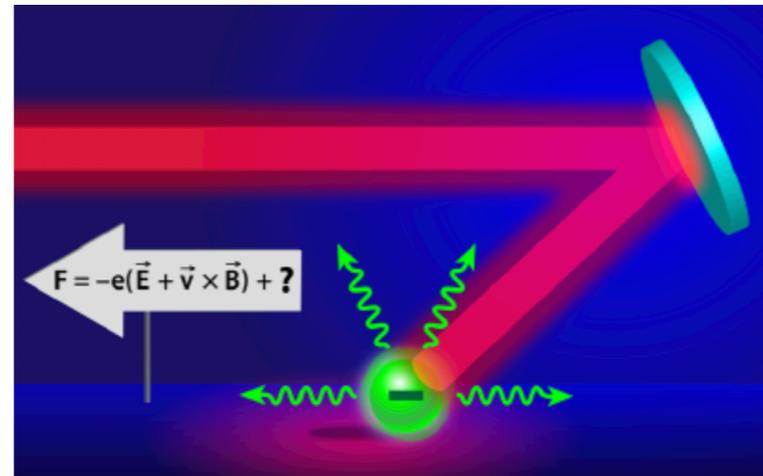
2.1 Å resolution

*Trypanosoma brucei* cysteine protease cathepsine B

Single Molecule Imaging Goal  
10 fs - 10 mJ - 2020

Aquila et al., “The LCLS single particle imaging roadmap” Stuct. Dynam. 2, 041701

## High field electrodynamics



From Phys. viewpoints, A. Macchi, Physics 11, 13 (2018)

Unexplored physics phenomena e.g. radiation reaction occur when electrons interact with an E-field close to the Schwinger critical field

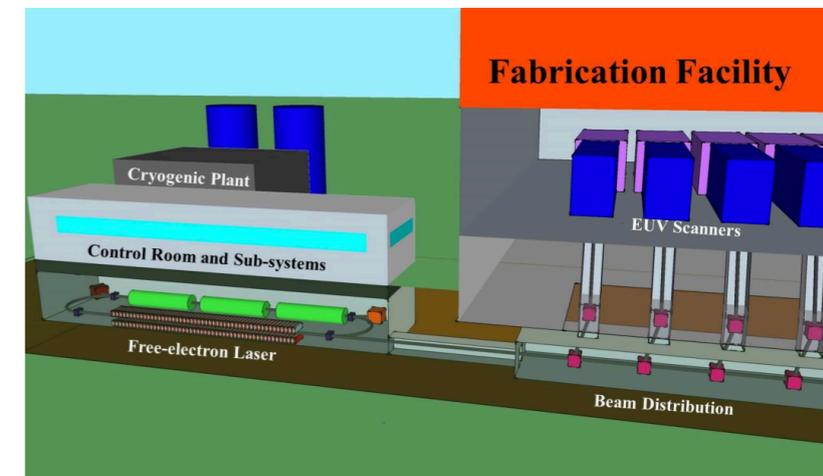
$$E_{Schwinger} = \frac{m_e^2 c^3}{\hbar^3 e} = 1.3 \times 10^{18} \text{ V/m}$$

Is being pursued actively at the moment using PW lasers and LWFA e-beams with *interesting* results e.g.

Cole *et al*, Experimental evidence of radiation reaction in the collision of a high-intensity laser pulse with a LWFA electron beam PRX 8, 011020 (2018)

Poder *et al* “Evidence of strong radiation reaction in the field of an ultra-intense laser”, arXiv:1709.01861

## Industrial FEL for EUV lithography



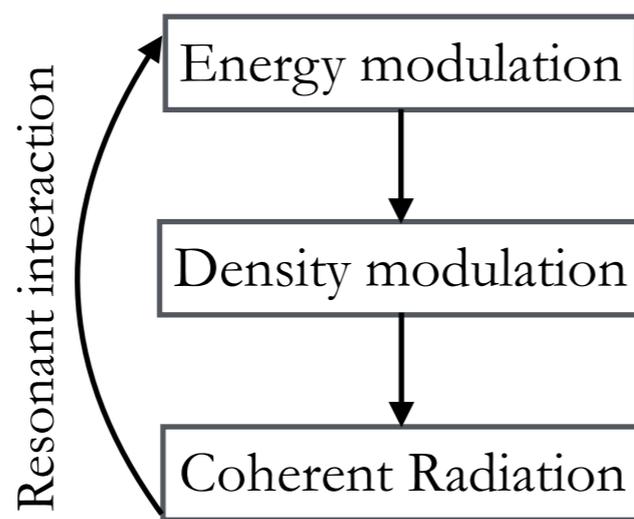
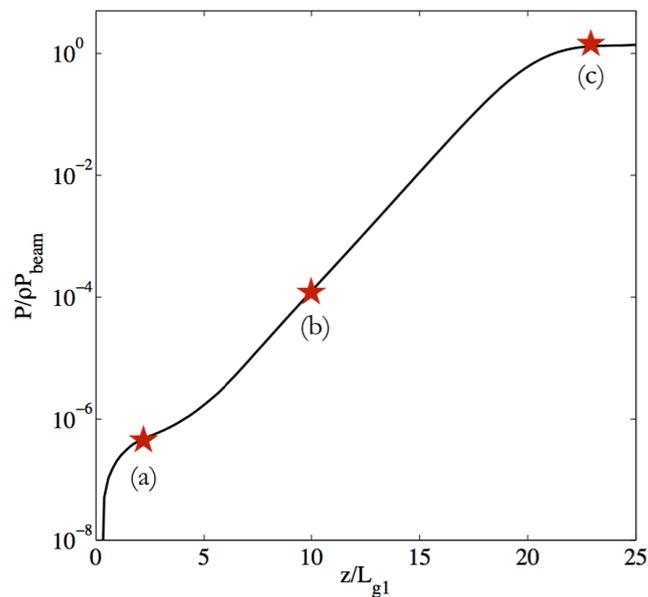
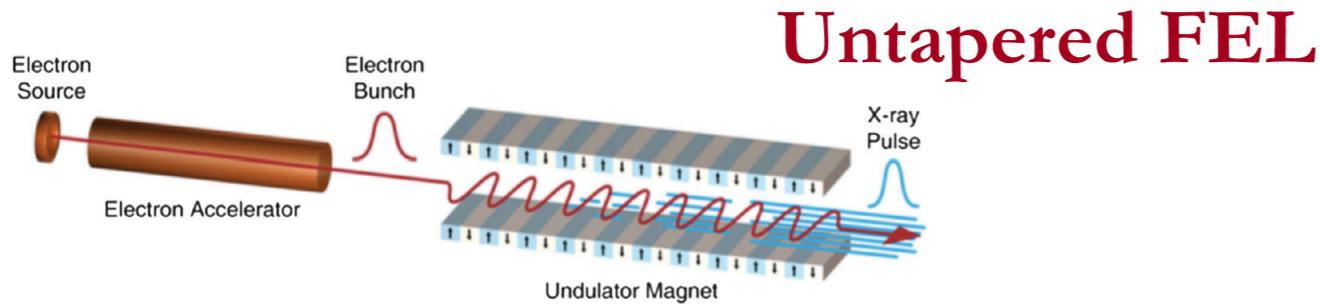
From Hosler, Wood “Free electron lasers: beyond EUV lithography insertions”, Global Foundries

Semiconductor industry is a very large economic driver \$2.0T global electronics maker

Current technology uses EUV from laser produced plasma sources. Expected cost is \$15-20B for leading-edge fab.

FELs being considered as an alternative but need to fulfill many requirements, one of which is *high efficiency* ~ 10 %

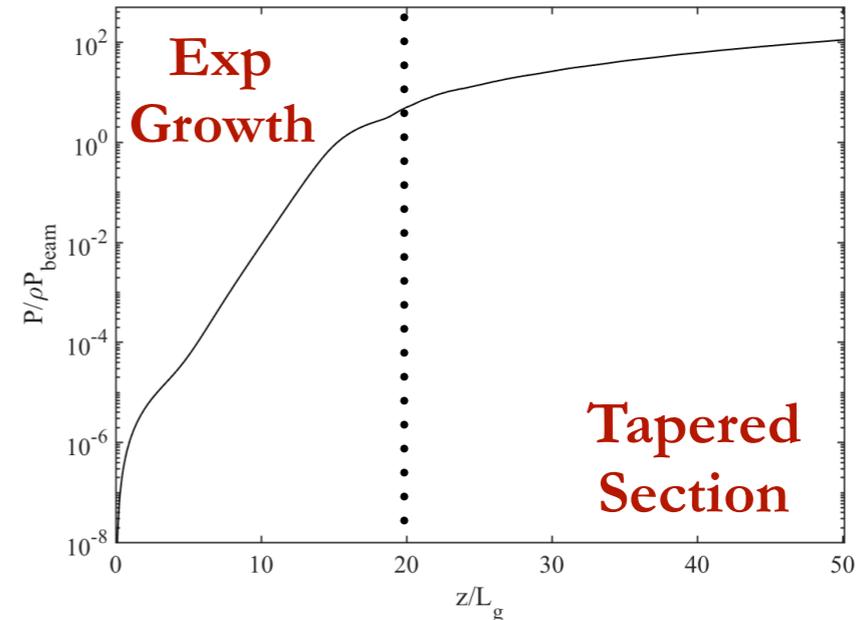
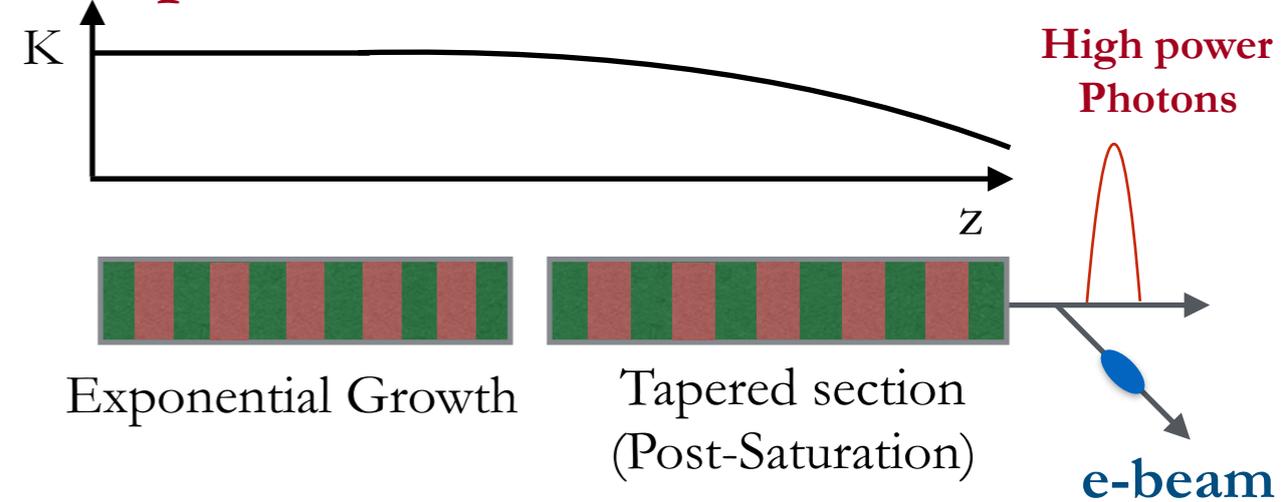
# Undulator tapering for high efficiency FEL



Resonance condition

$$\lambda = \frac{\lambda_u}{2\gamma_r^2} (1 + K^2)$$

## Tapered FEL



- Resonant interaction can continue *past saturation* by tapering the magnetic field  $K(z)$  to match the e-beam energy loss  $\gamma(z)$
- Questions are:
  - What is the max achievable efficiency?
  - What limits the max achievable efficiency?
  - How do you optimize the taper to achieve the max efficiency?

- 1-D physics (BPN Opt Comm. 1984) described by single parameter  $\rho$ .
 
$$P_{sat} = \rho P_{beam} \quad \rho = \frac{1}{\gamma} \left( \frac{I}{I_A} \left( \frac{K}{4k_u \sigma_x} \right)^2 \right)^{1/3}$$
- Some numbers:  
 $E_e = 250 \text{ MeV}, I = 0.2 \text{ kA}, \lambda = 250 \text{ nm}, \rho \sim 0.5\%$
- For high efficiency applications we want  $> 20x$  in efficiency to  $> 10\%$

# 1-D effects: How to choose the taper for max. power

Power scaling in  
post-sat regime

$$P_{rad} = P_0 + P_1 \bar{z} + P_2 \bar{z}^2$$



$$K = K_0 - bz - \frac{cz^2}{2}$$

Dominant for short  
undulators or large seed

Dominant for  
long undulators

# 1-D effects: How to choose the taper for max. power

Power scaling in post-sat regime

$$P_{rad} = P_0 + \cancel{P_1 \bar{z}} + P_2 \bar{z}^2$$

$$\sin \psi_r \propto \frac{|K'|}{E}$$

Dominant for short undulators or large seed

Dominant for long undulators

$$P_2 = \frac{Z_0}{8\pi} \left( \frac{K}{\gamma} \frac{\lambda_u}{\sigma_e} I \right)^2 (f_t \sin \psi_r)^2$$

Initial Condition contribution

Tapering contribution

# 1-D effects: How to choose the taper for max. power

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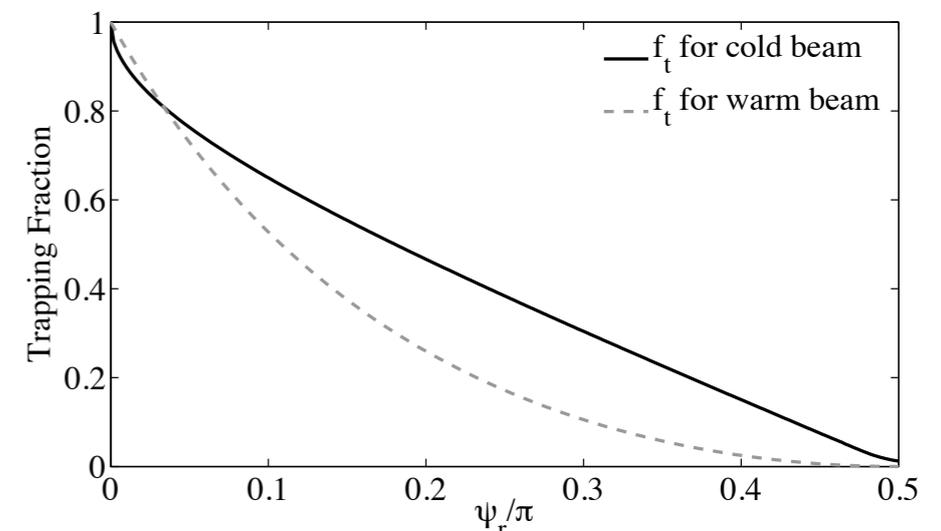
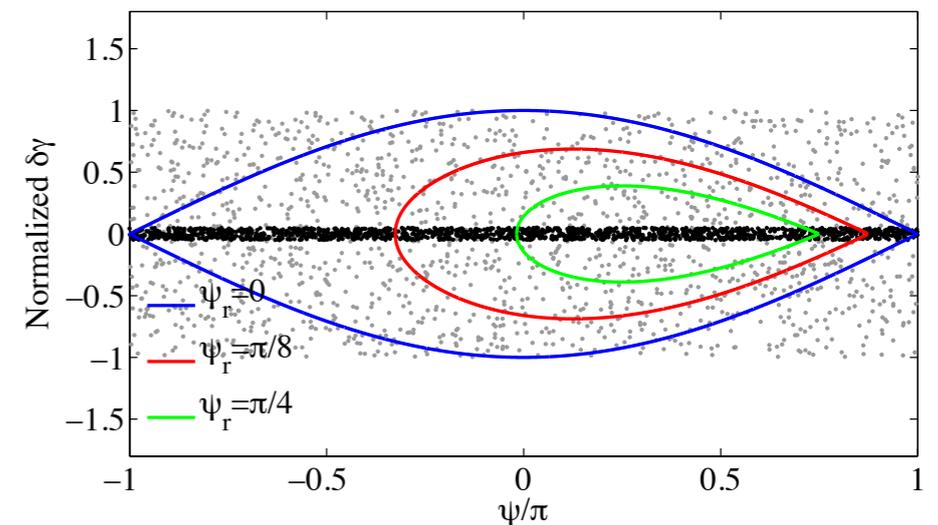
Dominant for short undulators or large seed

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Initial Condition contribution

Tapering contribution

## Take home messages from 1-D theory

(1) **Resonant phase  $\psi_r$  sets the speed of the taper and the size of the bucket**

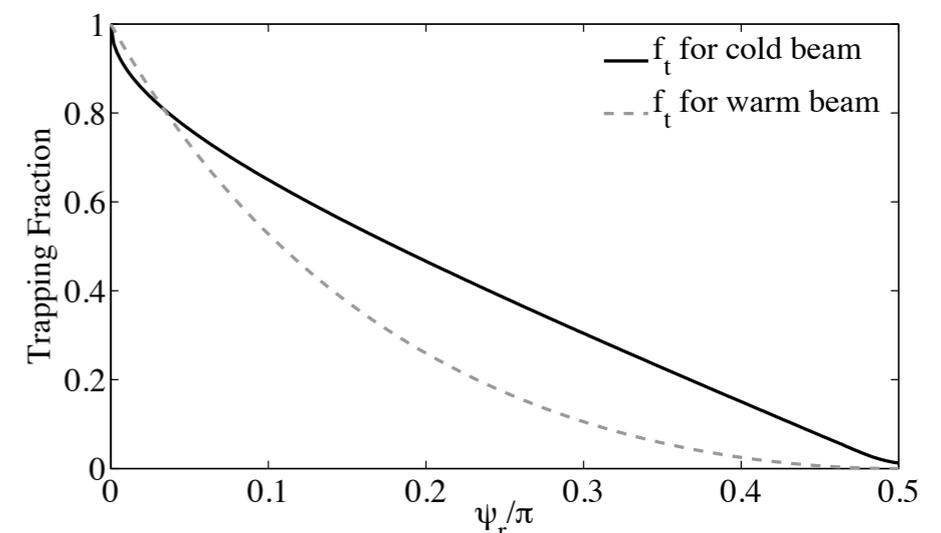
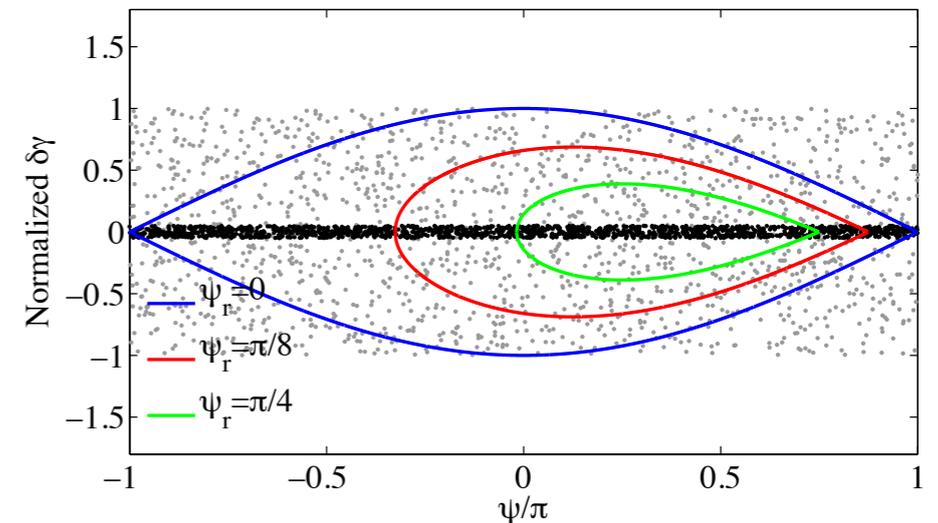
$\therefore$  Trade-off between number of electron trapped and how quickly the electrons are decelerated

(2) **Power scales like  $(f_t \sin \psi_r)^2$**

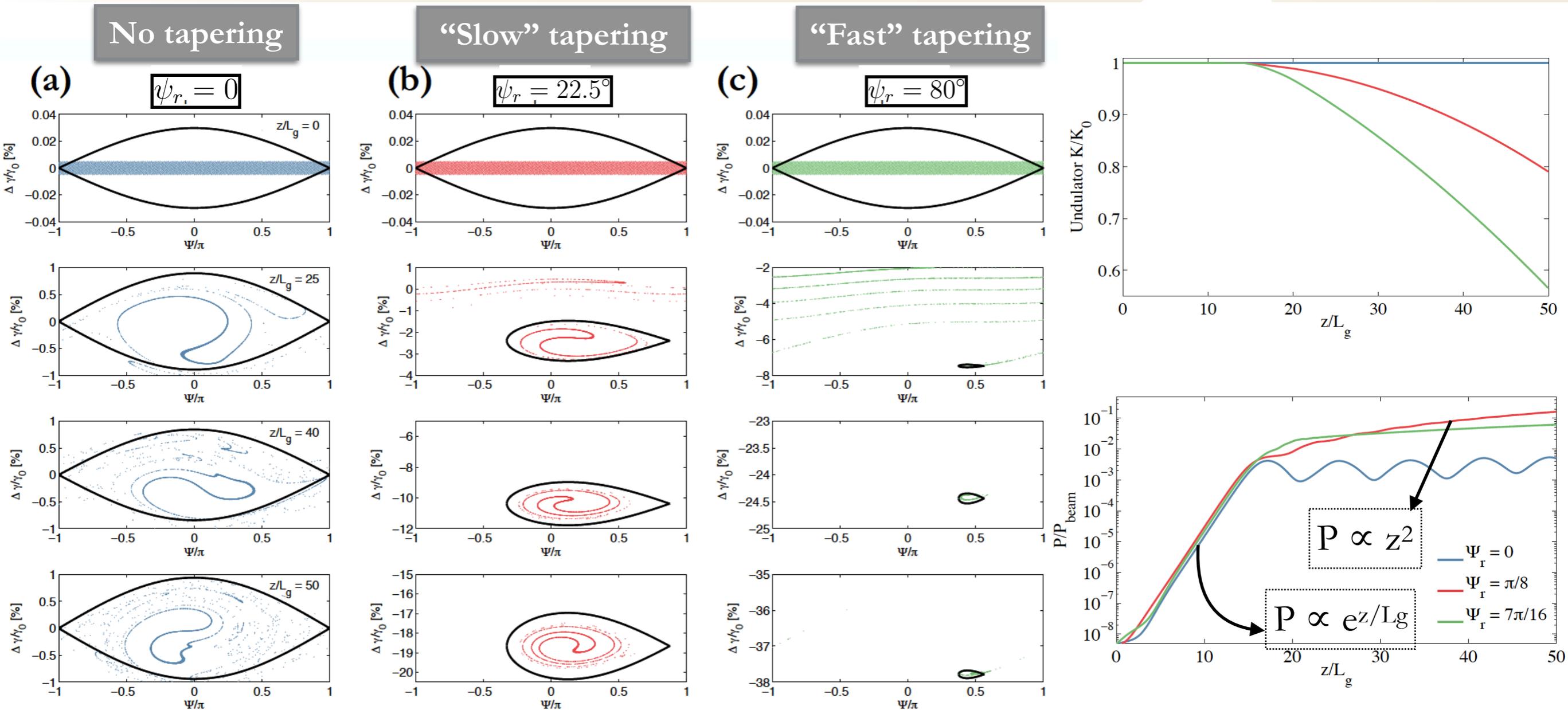
$\therefore$  Increasing the trapping by e.g. pre-bunching can increase P

(3) **Power scales like  $I^2 / \sigma_e^2 = I^2 / \beta \epsilon_n$**

$\therefore$  Brighter beam/smaller beta conducive to high efficiency



# 1-D effects: trade-offs and design considerations



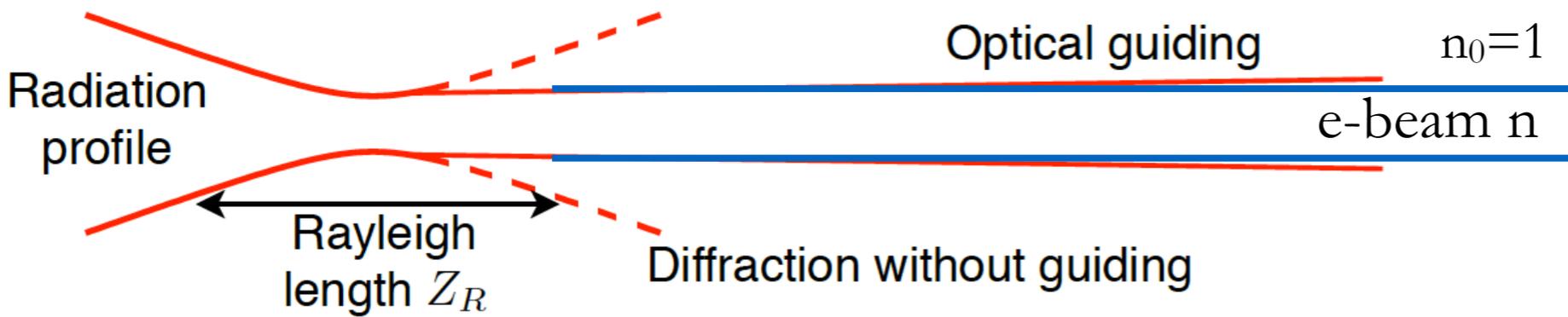
No tapering  
efficiency is the  
same as saturation

“Slow” taper  
strikes the balance  
between total  
energy loss and  
trapping fraction

“Fast” taper has  
larger net energy  
loss but smallest  
fraction captured

In 1-D theory, with a  
judiciously chosen taper  
you can continue to  
increase power by  
adding undulators

# 3-D effects: diffraction limits to the 1-D model



E-beam refractive index

$$n - 1 = \frac{\chi_2 K}{k \gamma} \langle e^{i\psi} \rangle E$$

Growth of field reduces guiding sets limit on max. E field

Microbunching and trapping must be kept high to maintain good guiding

$$E_{max} \approx \frac{Z_0 I K}{\lambda \gamma} \cos \psi_r$$

$$P_{rad} = \frac{2\pi}{Z_0} E^2 \sigma_r^2$$

D. Prosnitz, A. et al, Phys. Rev. A 24, 1436 (1981)

Scharlemann, T. et al, Phys. Rev. Lett. 54, 17 (1981)

Fawley W., NIMA 375 (1996)

Yao, J., PRSTAB. 15, 050704 (2012)

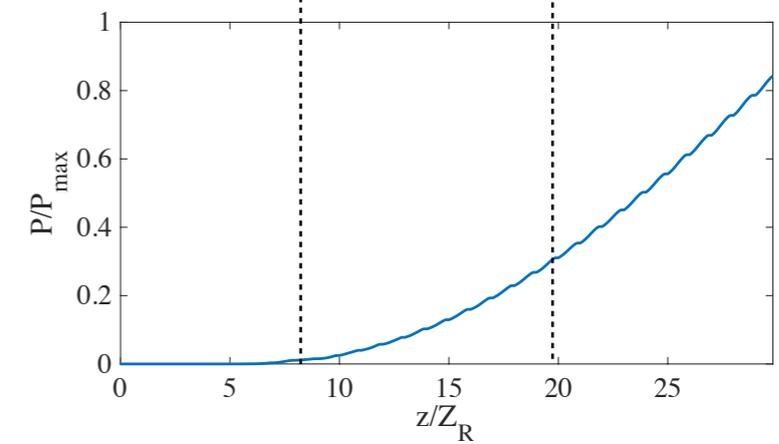
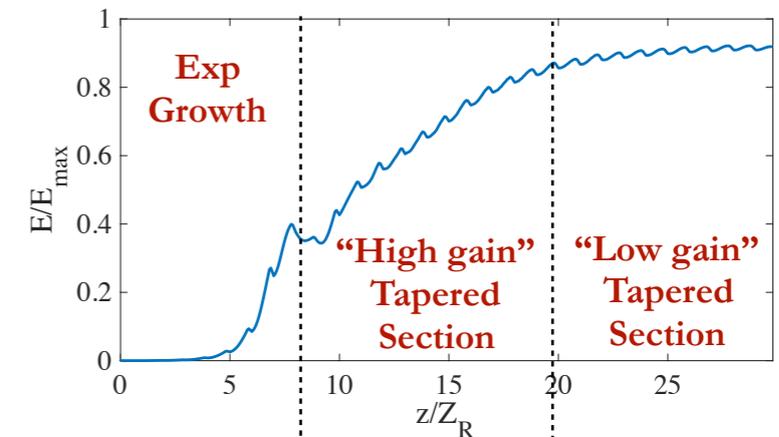
Schneidmiller, et al., PRSTAB. 18, 030705 (2015)

## Take home messages from 3-D theory

Limit on field and radiation growth region **in contrast with 1-D theory**

Needs to be considered for long undulators  $L_u \gg Z_R$

**Want to extract energy (taper) as fast as possible to outrun diffraction limit**

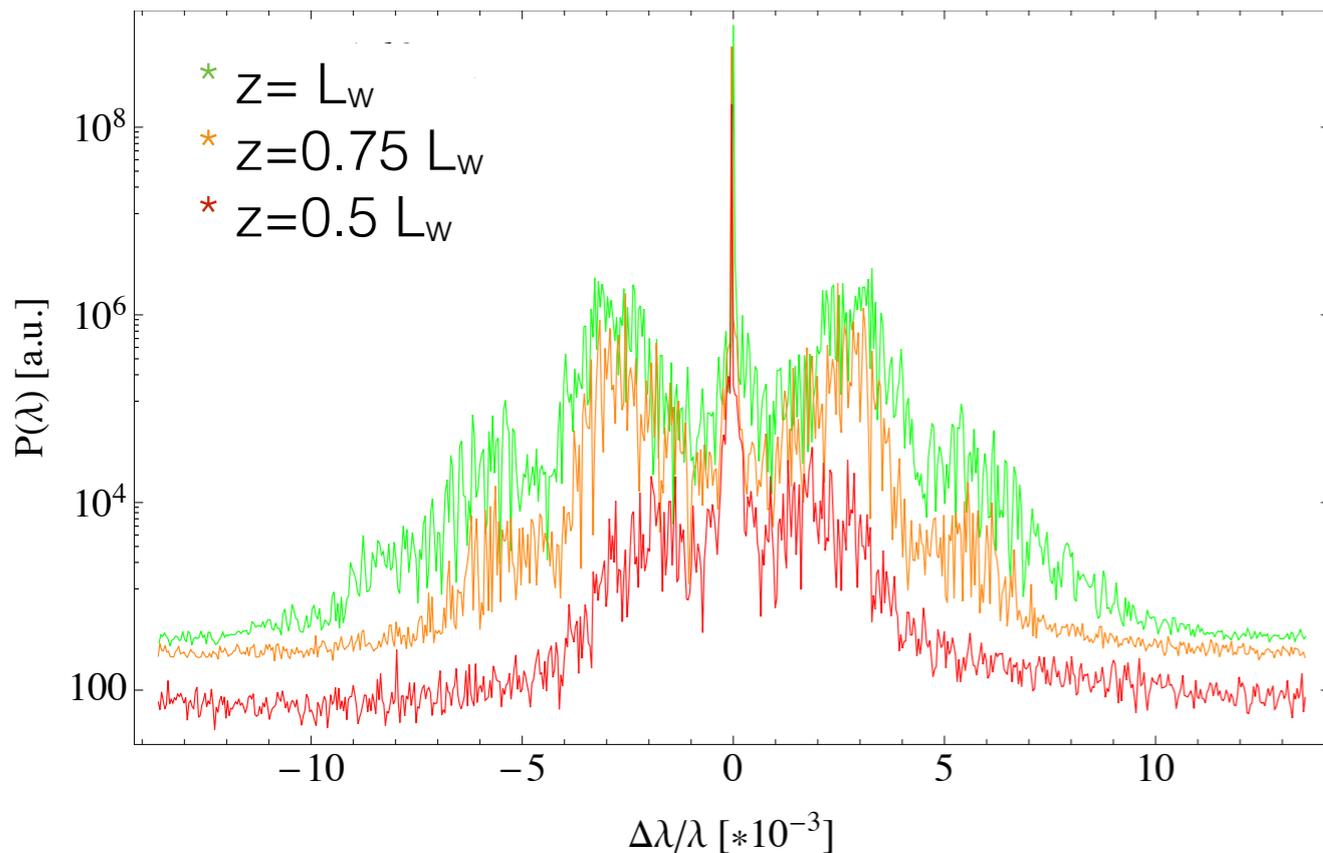


# Time Dependent effects: limits to the 1 frequency model

Electron beam shot noise and synchrotron motion

$$E'(z, t) \propto I(t) \left\langle \frac{\sin \psi(z, t)}{\gamma(z, t)} \right\rangle$$

$$\phi'(z, t) \propto \frac{I(t)}{E(z, t)} \left\langle \frac{\cos \psi(z, t)}{\gamma(z, t)} \right\rangle$$



Amplitude and phase modulations of the radiation field

Resonance between sideband radiation and synchrotron motion

Radiation field saturation from reduced optical guiding gives  $\sim$  constant  $L_{\text{synch}}$

**Sideband Instability**

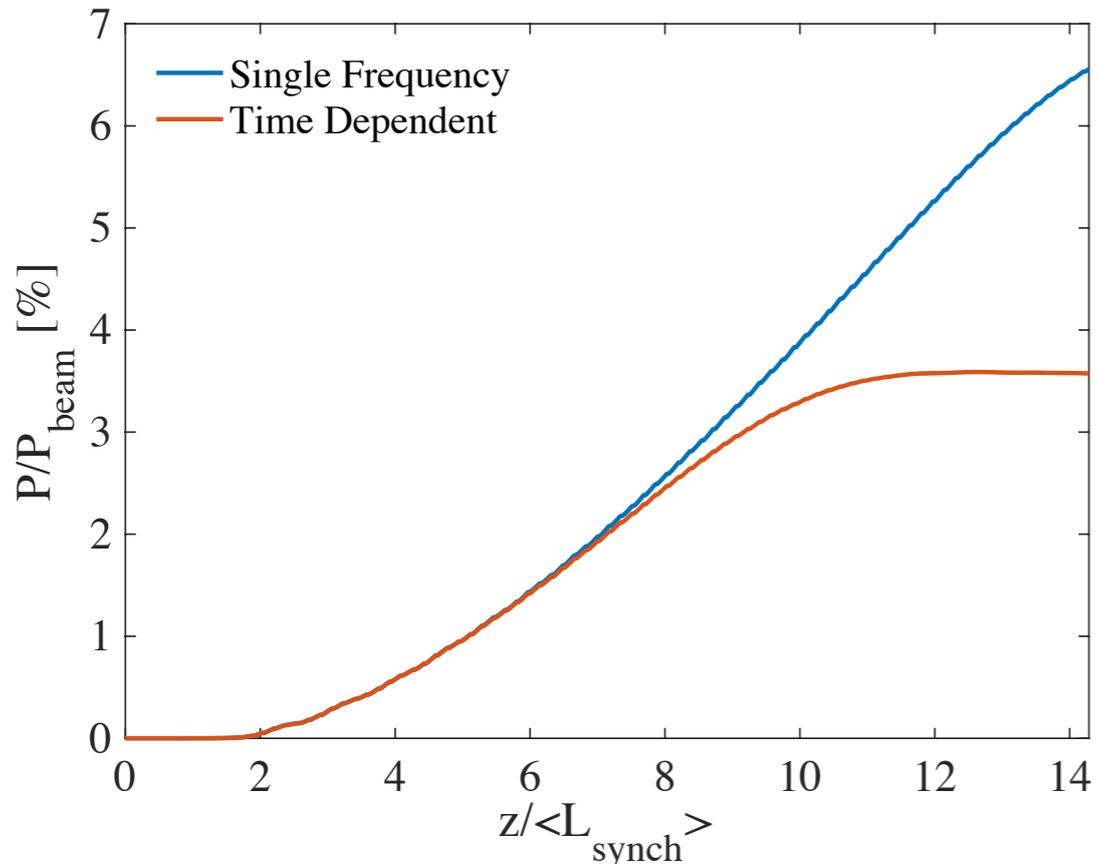
“...the electron motion in a FEL will become chaotic when the sideband amplitude exceeds a certain threshold. This, in turn, will result in significant electron detrapping. Since it is the deceleration of the trapped electron bucket that provides the energy for the radiation in the case of tapered wigglers, detrapping will cause loss of amplification for the FEL signal”

# Time Dependent effects: limits to the 1 frequency model

Electron beam shot noise and synchrotron motion

$$E'(z, t) \propto I(t) \left\langle \frac{\sin \psi(z, t)}{\gamma(z, t)} \right\rangle$$

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Amplitude and phase modulations of the radiation field

Resonance between sideband radiation and synchrotron motion

Radiation field saturation from reduced optical guiding gives ~ constant L<sub>synch</sub>

**Sideband Instability**

**Take home messages from TDP theory**

Sideband instability can cause second saturation of radiation power in tapered FEL

Want to reduce the sideband growth along tapered undulator to continue extracting power

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2.1.2. NOCIBUR experiment (BNL-ATF)

2.2. Present

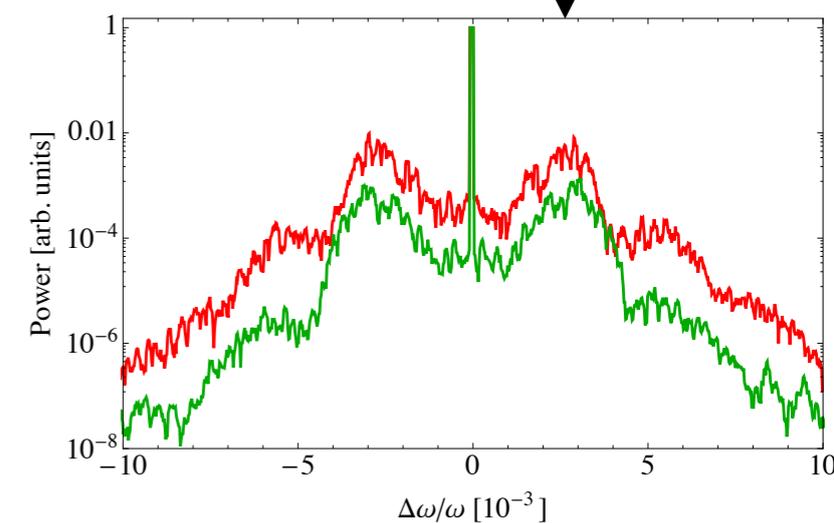
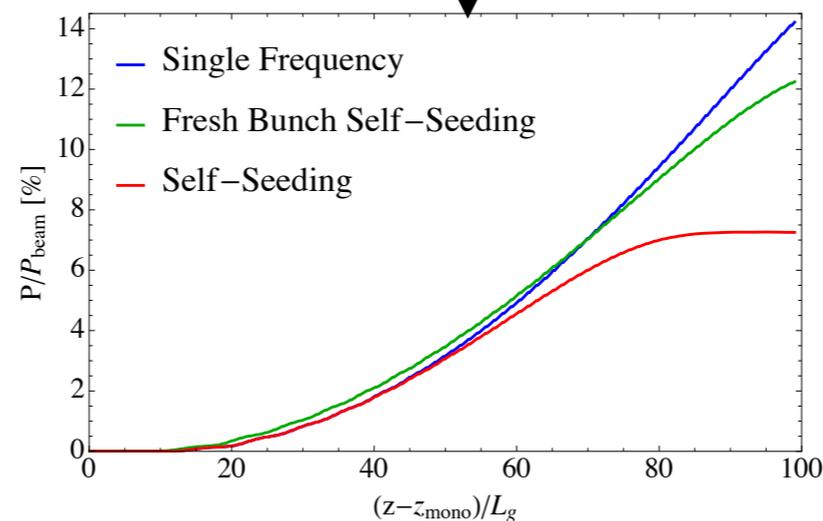
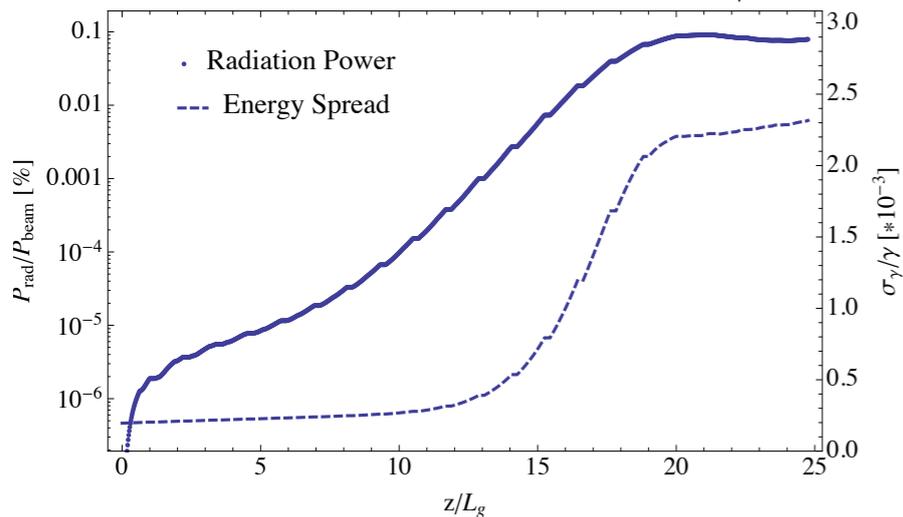
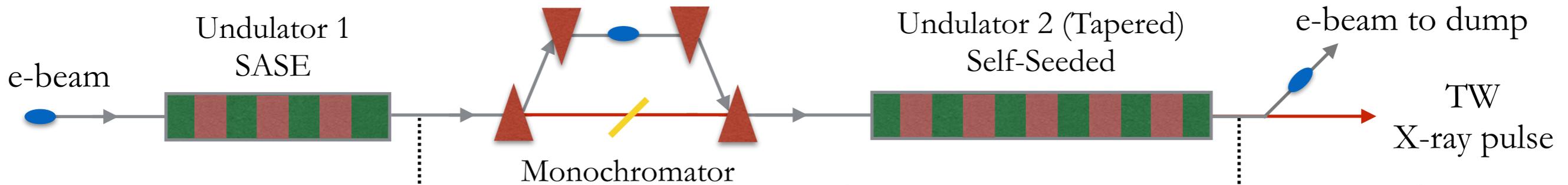
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## 3. Conclusions

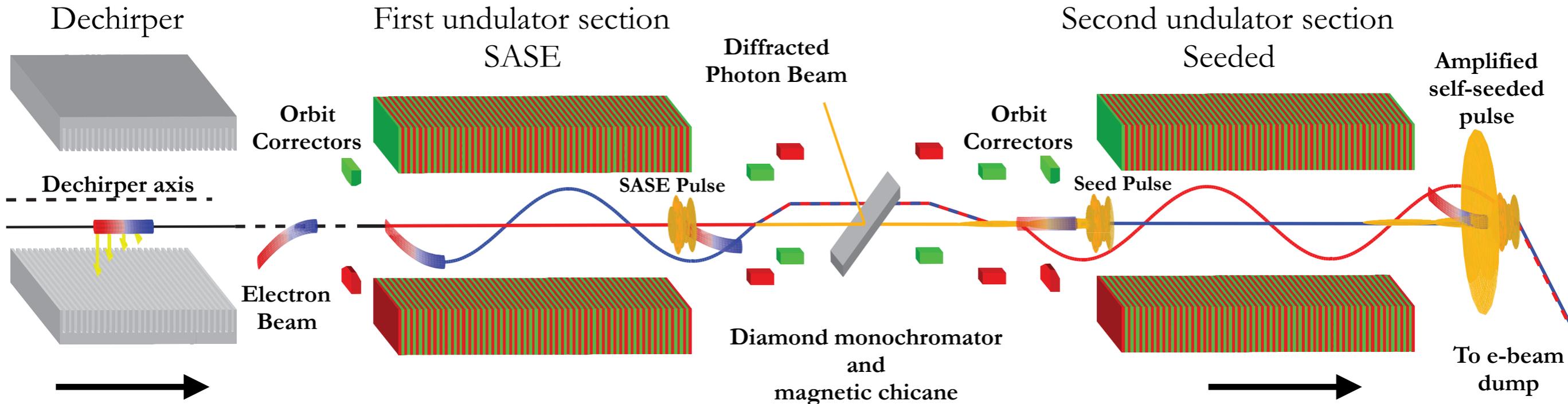
# Fresh bunch self-seeding experiment - simulations



C. Emma et al., PRAB 19, 020705 (2016)

- GENESIS simulations show time dependent losses from sideband instability can be overcome using a large seed ( $P_{\text{seed}}/P_{\text{noise}} \sim 10^3$ ).
- In a self-seeded FEL having a large seed comes at the expense of a large energy spread at the start of the seeded section.
- Escaping the trade-off between seed power and energy spread requires **fresh bunch self-seeding**.

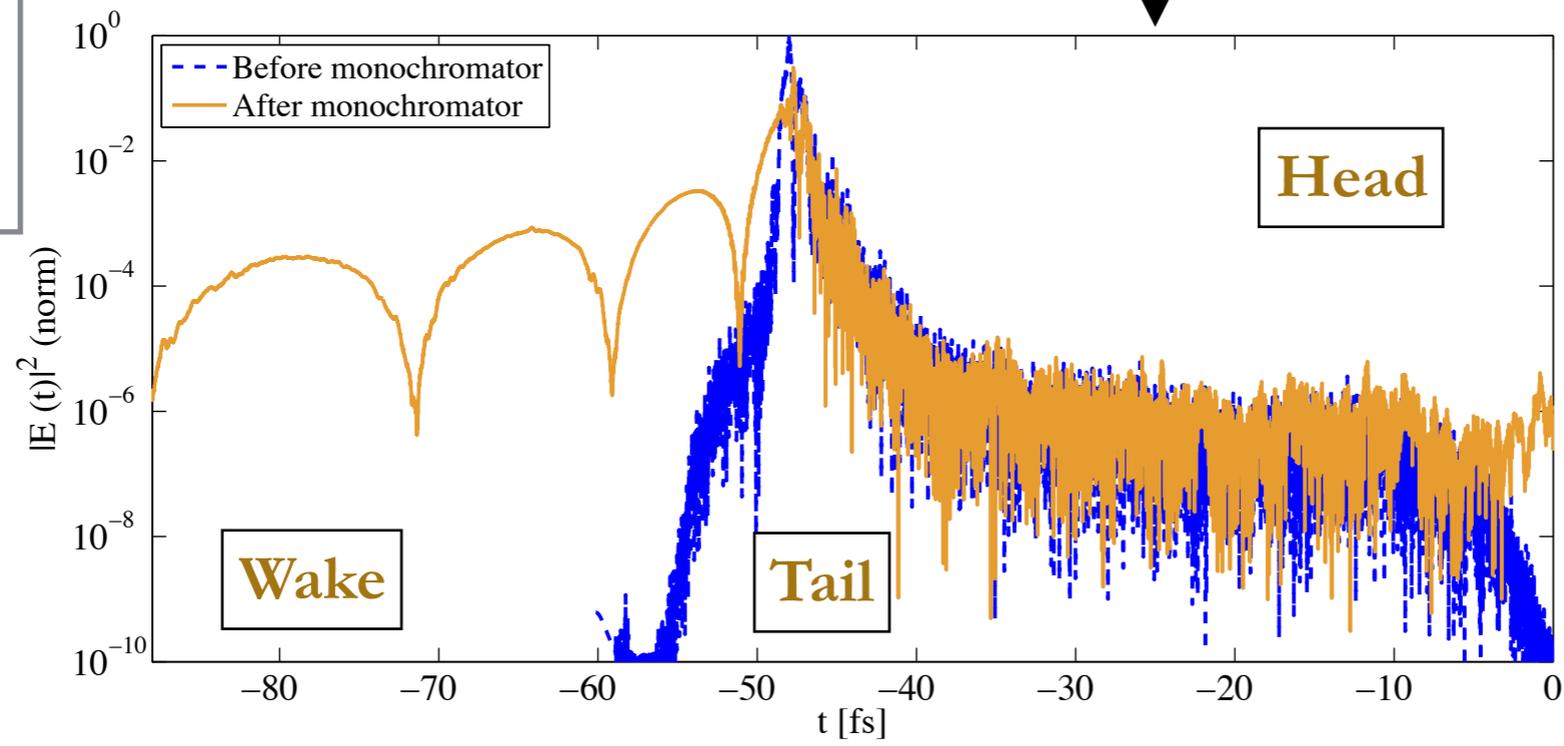
# Fresh slice self seeding experiment at LCLS



$$E_{\text{X-Ray}} = 5.5 \text{ keV}$$

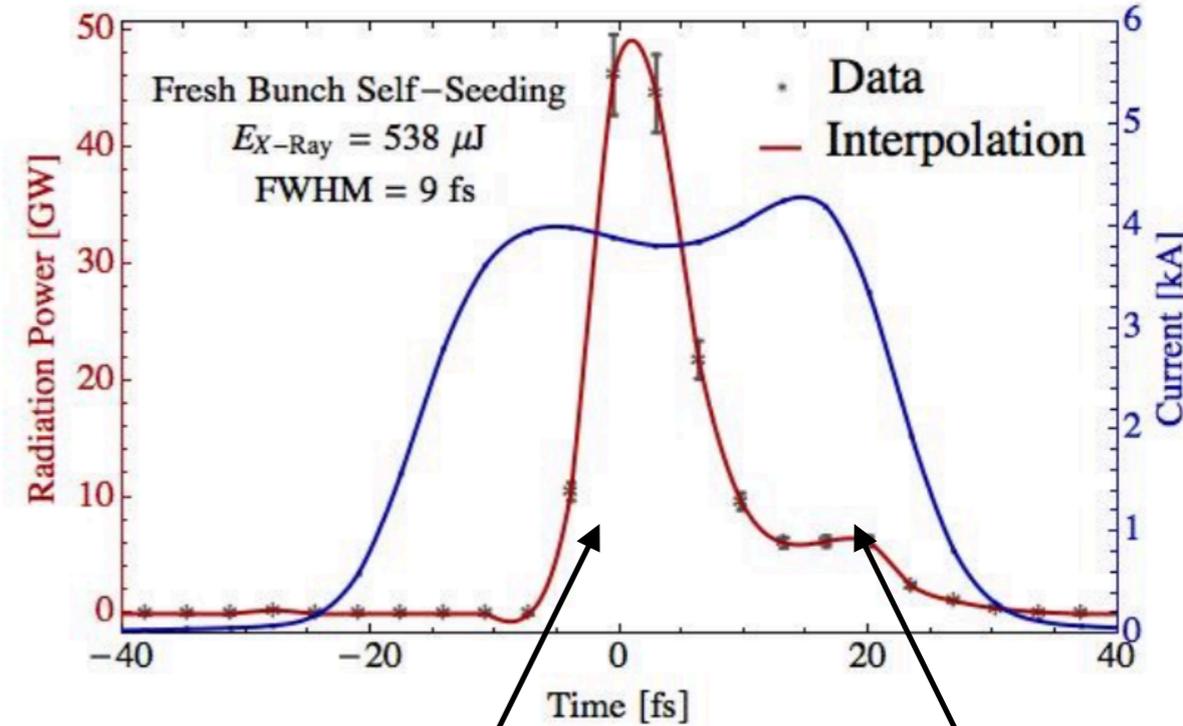
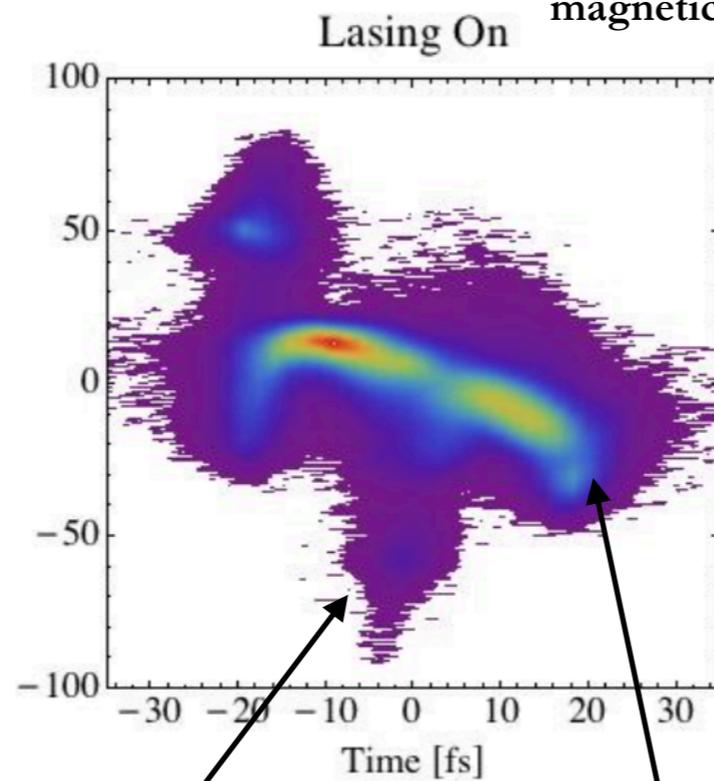
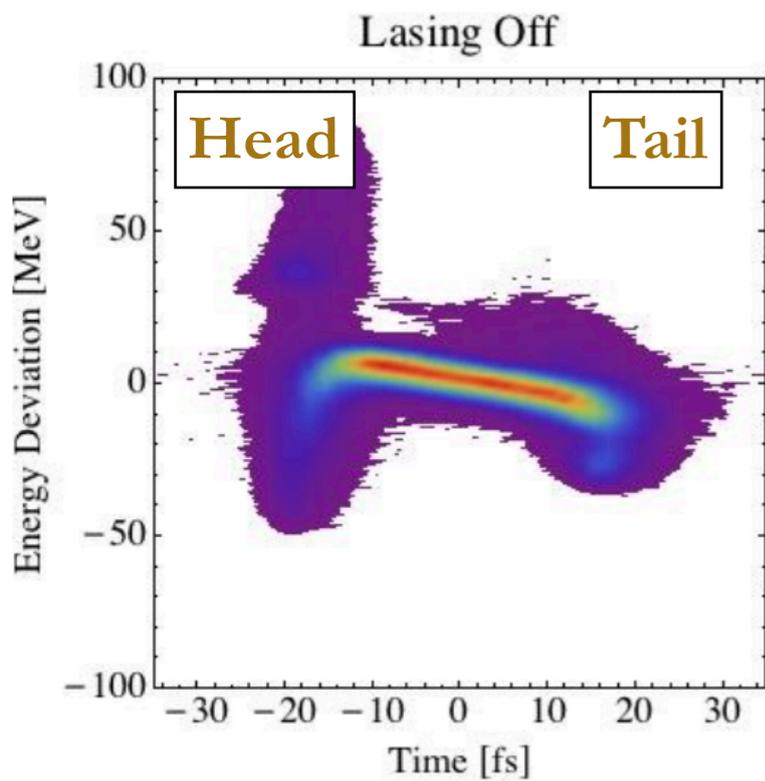
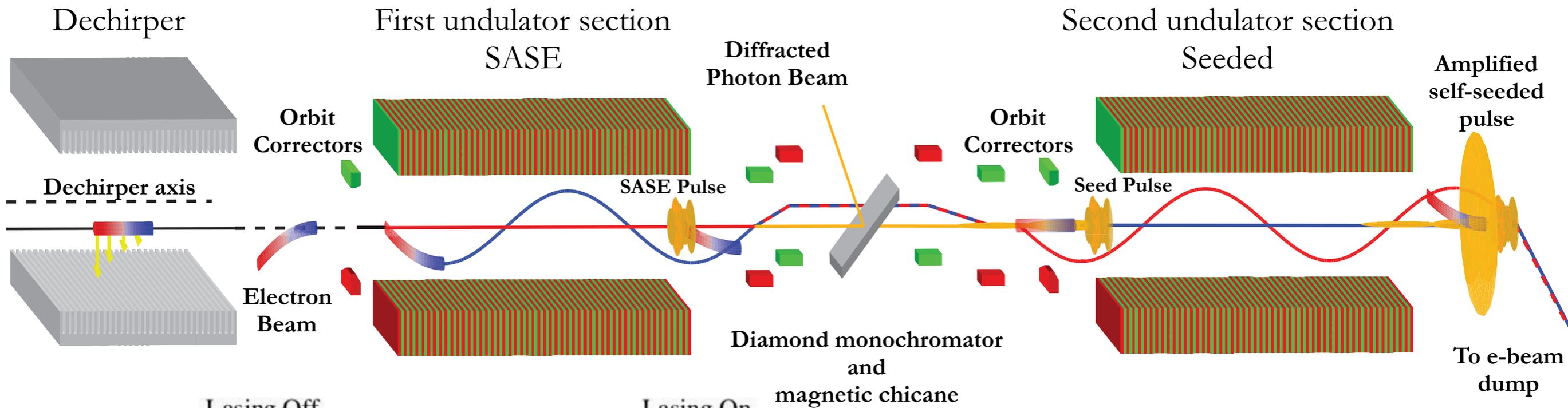
E-beam

$I_{\text{pk}} = 4 \text{ kA}$   
 $E = 11 \text{ GeV}$   
 $Q = 180 \text{ pC}$



- Diagnostics
- 1) **Transverse deflecting cavity**  
Electron beam energy loss (time resolved)
  - 2) **Gas detector**  
X-ray intensity
  - 3) **X-ray spectrometer**

# Fresh slice self seeding experiment at LCLS



Seeded core

SASE lasing slice

Seeded core

SASE lasing slice

# Fresh slice self seeding experiment at LCLS

| Brightness ratio            | Average | Average filtered on e-beam energy | Peak |
|-----------------------------|---------|-----------------------------------|------|
| $B_{FBSS}/B_{SASE}$         | 12.5    | 15.5                              | 35.4 |
| $B_{FBSS}/B_{self-seeding}$ | 2.4     | 2.1                               | 2.3  |

Scientific Achievements

Short  $\sim 10$ fs pulses with 50 GW power and  $<10^{-4}$  b.w.

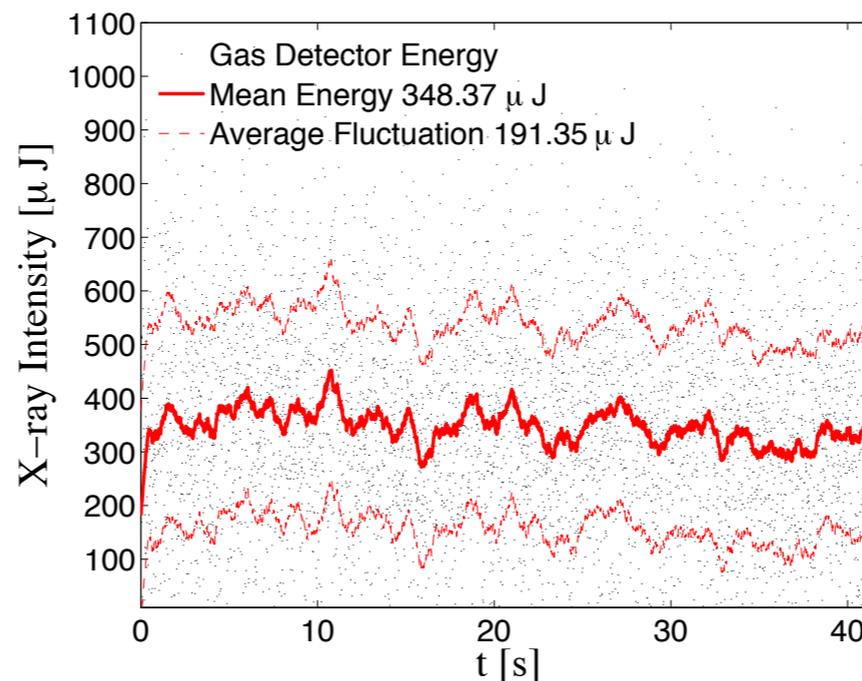
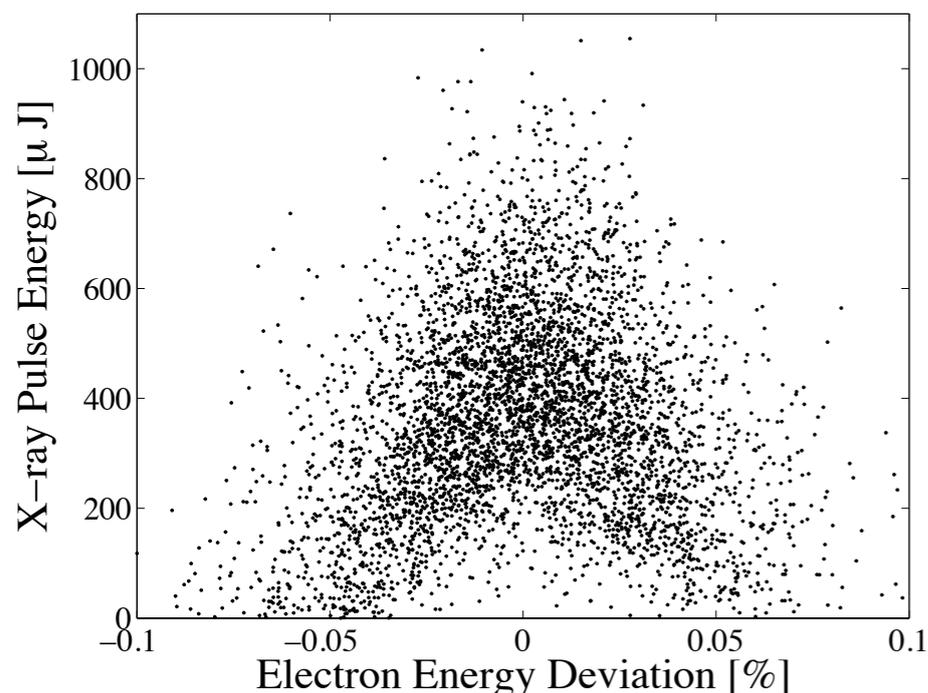
$\sim 2^*$  increase in X-ray power / brightness compared to self-seeding

**Table 6.1:** Comparison between the average and peak brightness of the FBSS scheme with SASE and self-seeding at the same photon energy. The left column is an average without filtering the data based on the incoming electron beam energy. The middle column is an average of the data within the energy jitter window  $\Delta E/E_0 = 0.5\%$ . The right column is calculating using the best shot for each of the three schemes.

Issues to work on

Large shot-to-shot intensity fluctuations due to:

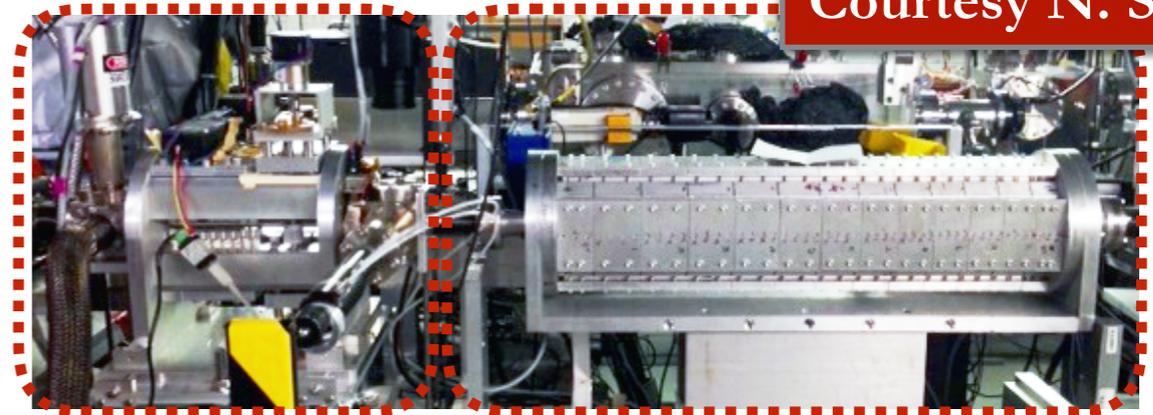
- i) Energy jitter
- ii) Seed power fluctuations from self-seeding monochromator



# TESSA demonstration NOCIBUR experiment

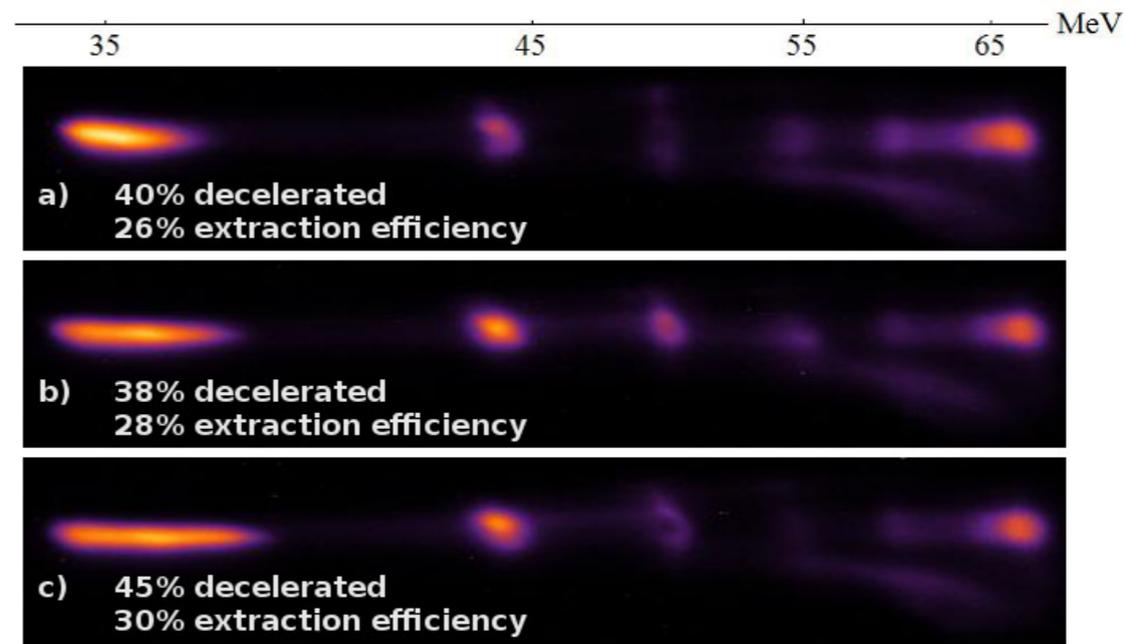
Courtesy N. Sudar

- strongly tapered helical undulator (decelerating)
  - $\lambda_w$ : 6 - 4 cm,  $K$ : 2 - 1.2,  $\Psi_r = \pi/4$
- pre-bunched electron beam
- 200 GW, 10.3  $\mu\text{m}$  CO<sub>2</sub> laser seed
- 45% of particles decelerated from 65  $\rightarrow$  35 MeV ( $L_u=0.54\text{m}$ )
- **30% conversion efficiency**
- good agreement with GPT simulations
- Genesis simulations show expected radiation growth for electron beam energy loss

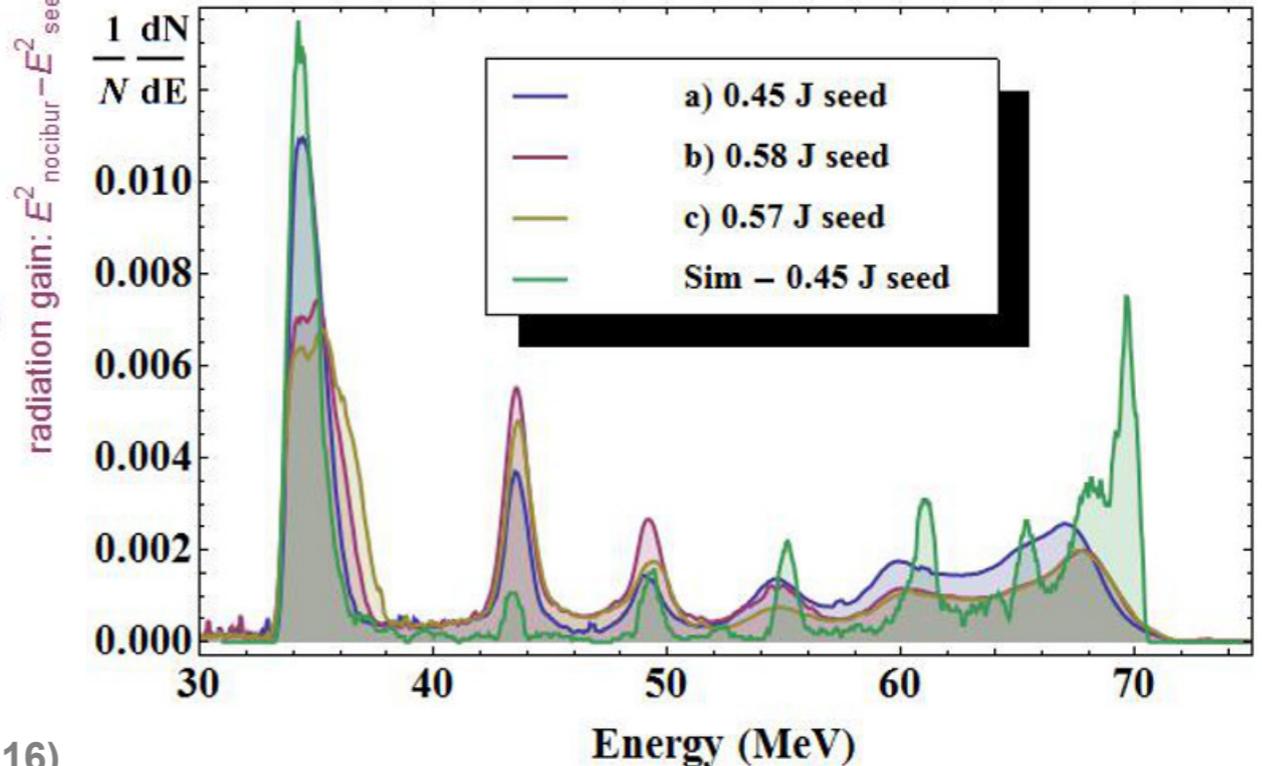
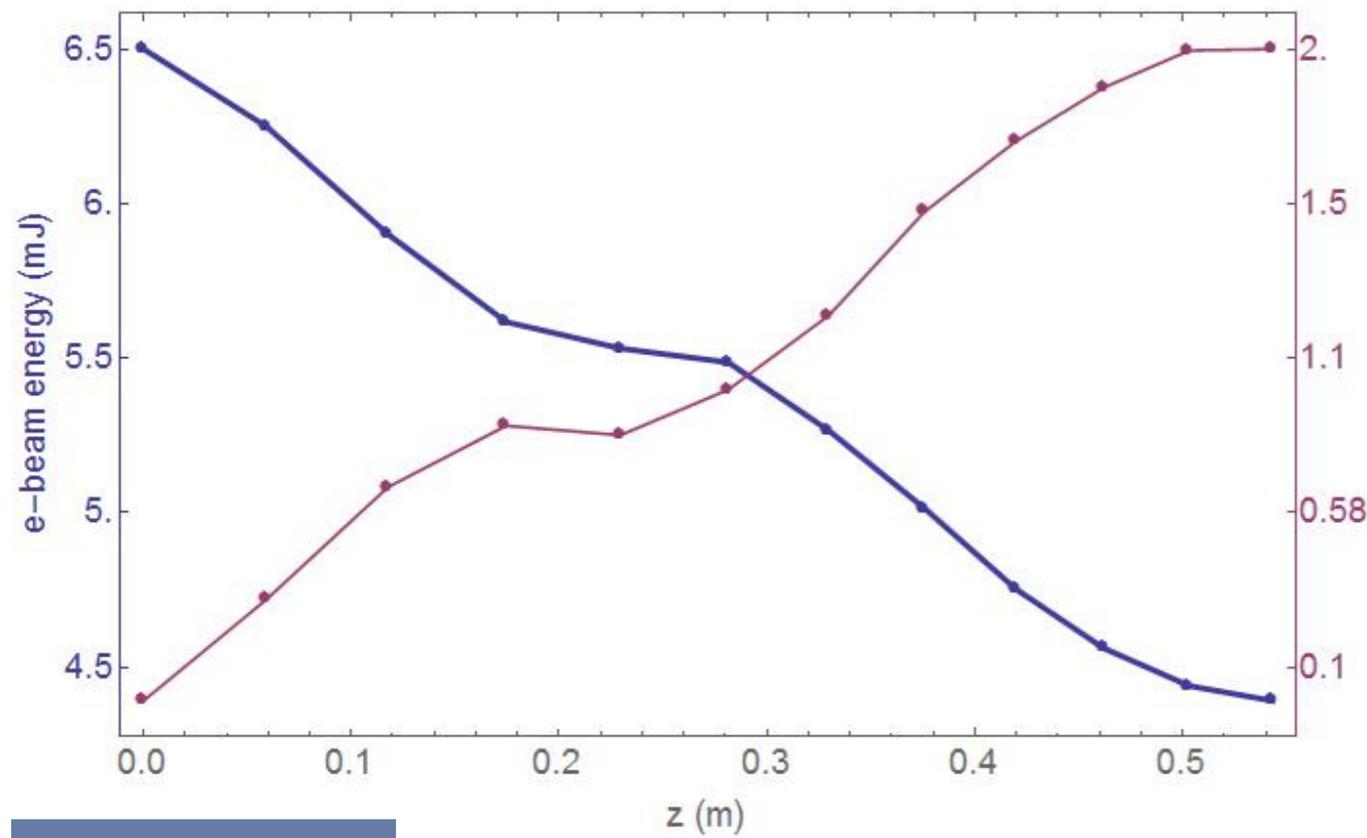


Pre-buncher

Nocibur undulator

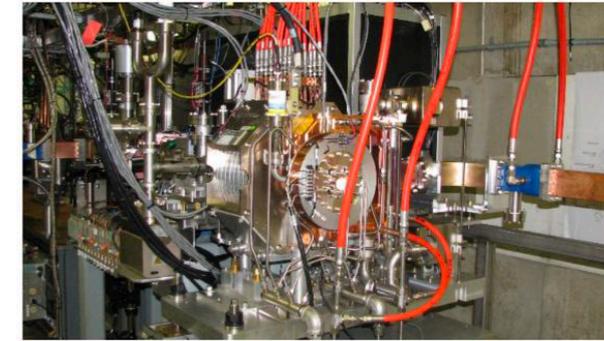


**Issue to work on**  
Very large external seed laser necessary. Would like a self-contained system



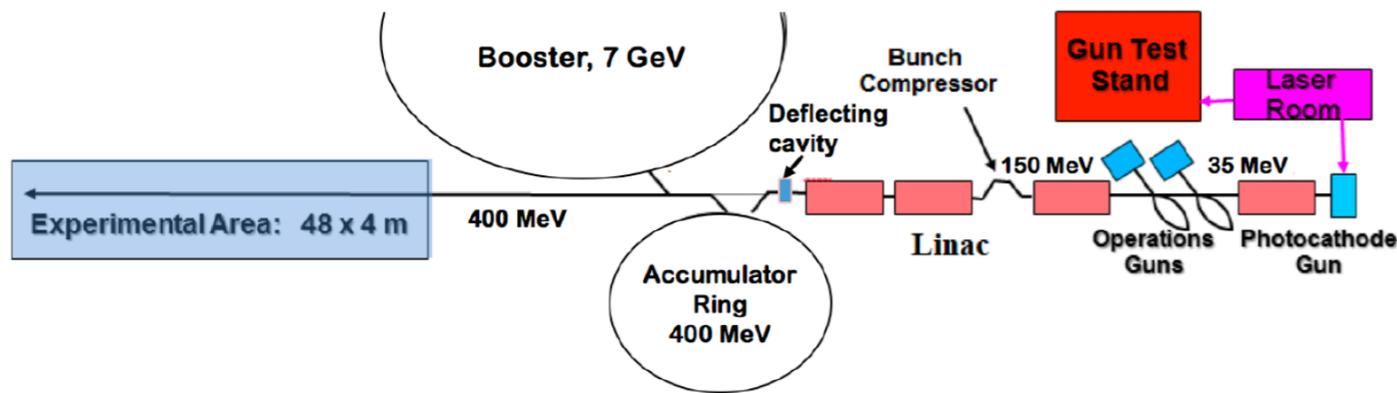
Presently ongoing  
experimental work

Courtesy  
P. Musumeci, Y. Park



## TESSA-266

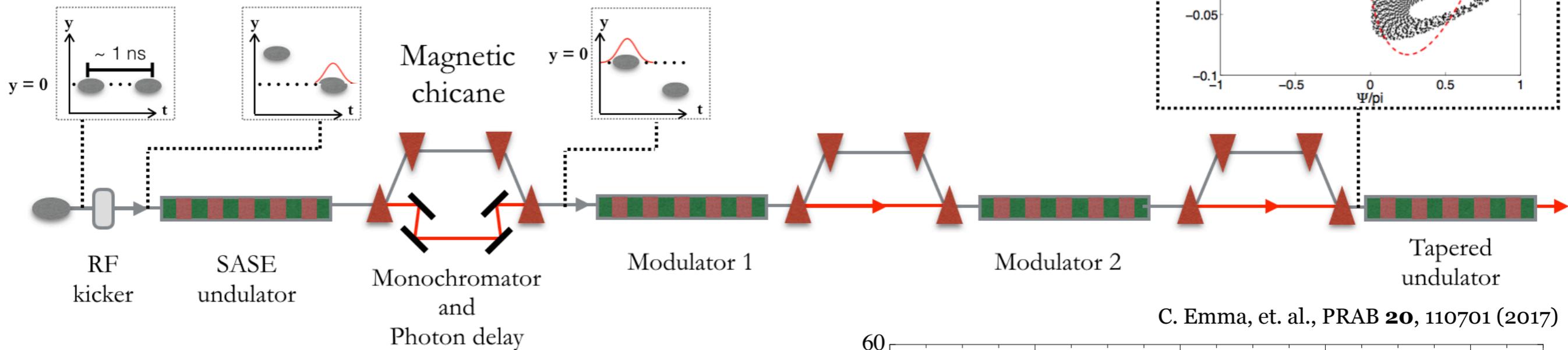
- Higher gain and higher current for radiation measurement
- Injector Linac at Argonne National Laboratory will operate at 375 MeV, 1kA and provide injection for 1.5 minutes in every 2 minutes
- decelerate ebeam for ~10% efficiency in 4 m undulator.
- Significant improvement from <1% efficiency in previous short wavelength FEL.



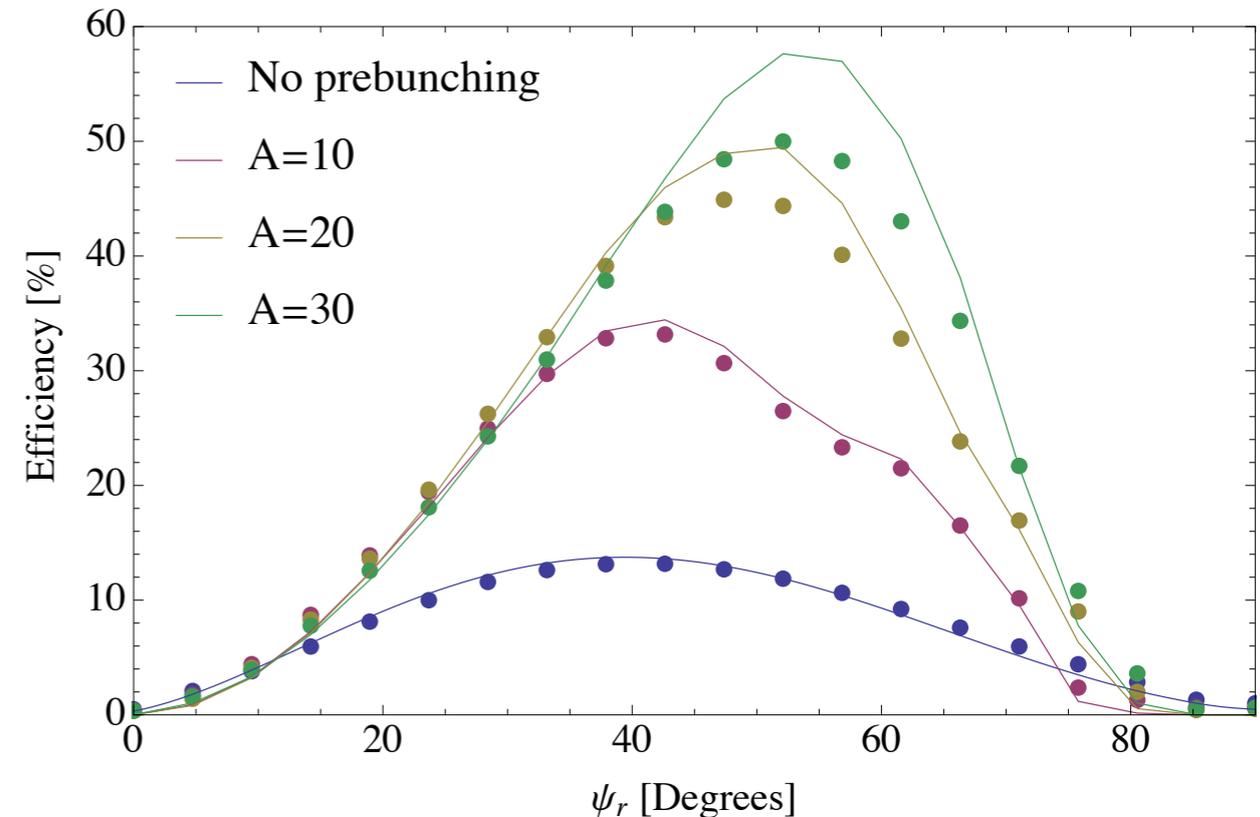
|                      |        |
|----------------------|--------|
| E-beam Energy        | 375MeV |
| Radiation Wavelength | 266nm  |
| Emittance            | 2 um   |
| Seed Power           | < 1GW  |
| Energy spread        | ~.1%?  |
| Peak Current         | 1 kA   |
| Bunch length (RMS)   | 50um?  |

# Future advanced schemes to be tested #1

Working title: Double-bunch, pre-bunched, fresh-bunch, self-seeded XFEL

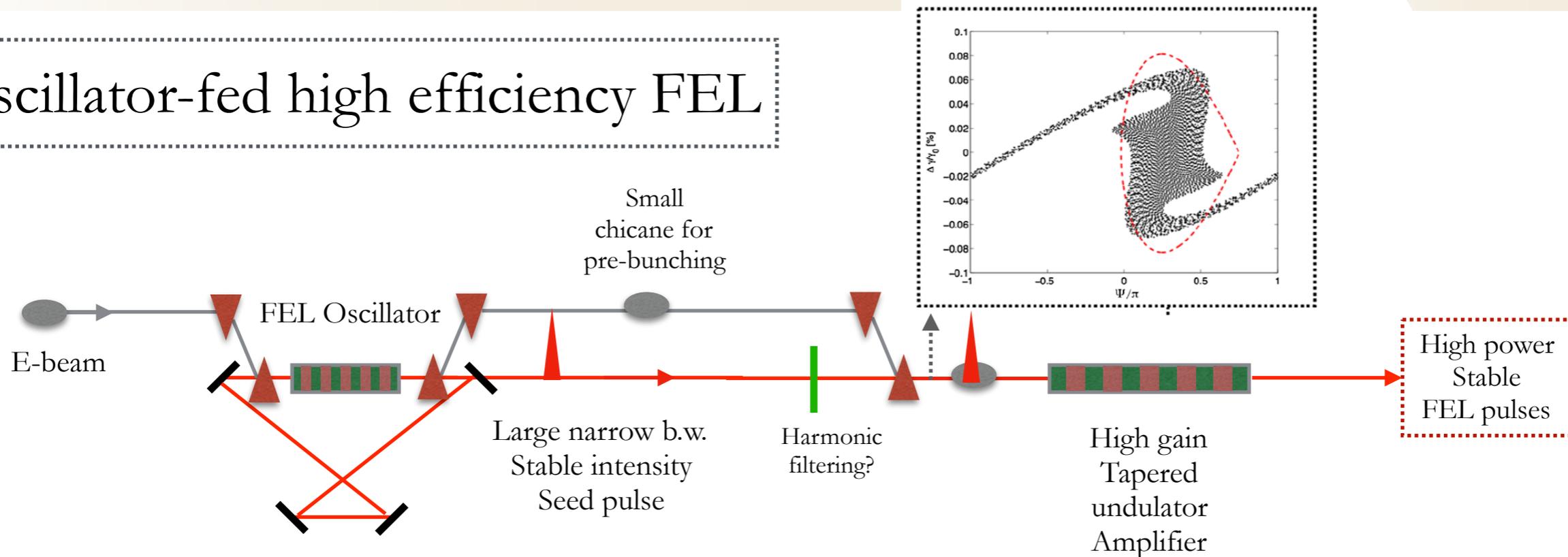


- Pre-bunching with a strong seed (10-30 \* the electron energy spread) can increase efficiency to 30-50 % (1-D sims)
- The second advantage is the peak efficiency occurs at larger resonant phase. This allows faster energy extraction, countering the effects of diffraction and sideband instability. Still suffers from intensity fluctuations.
- Has not been tested: would be nice to do the experiment!



# Future advanced schemes to be tested #2

## Oscillator-fed high efficiency FEL



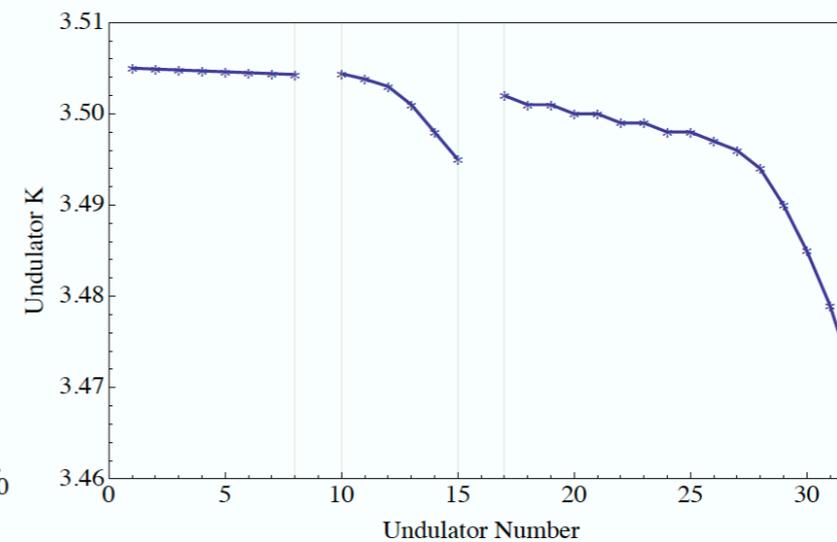
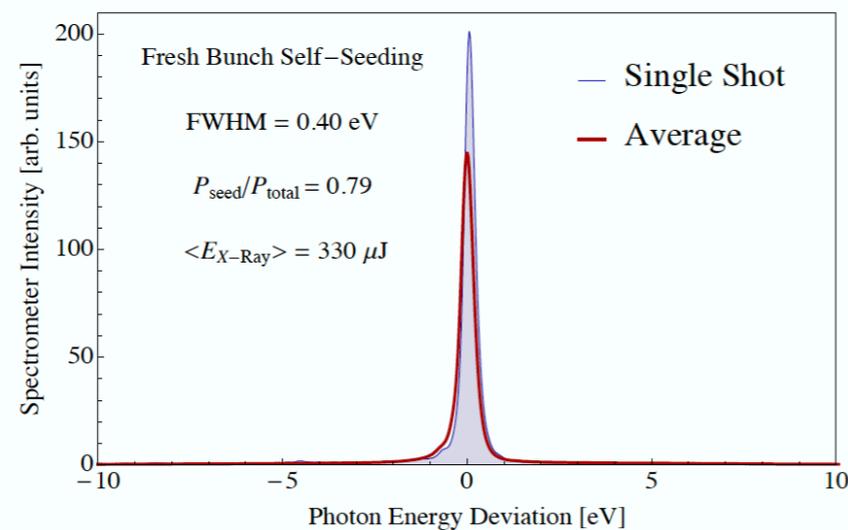
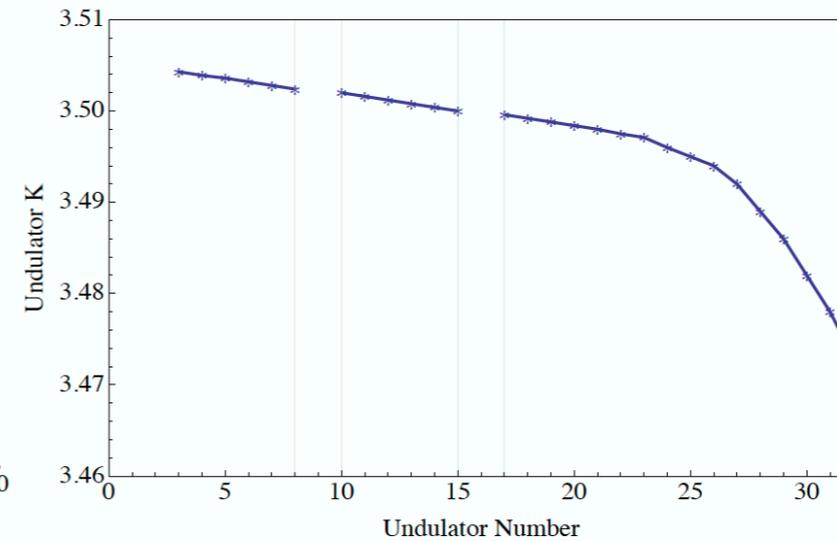
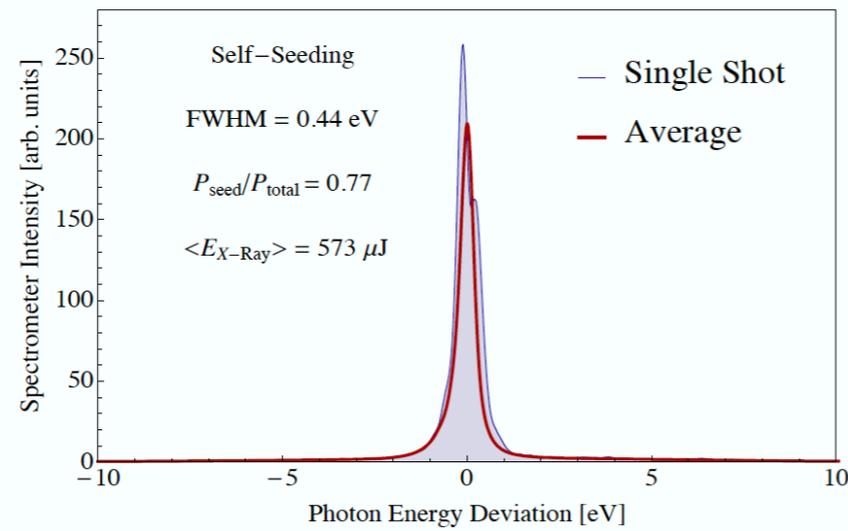
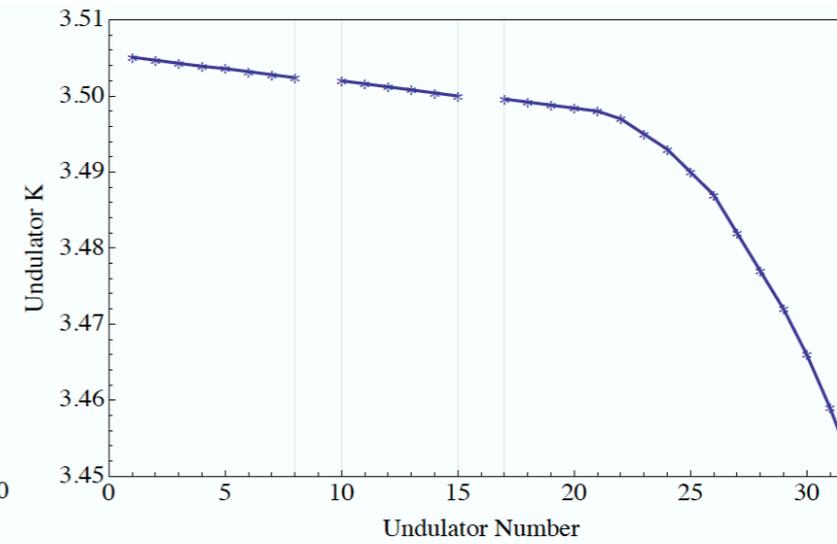
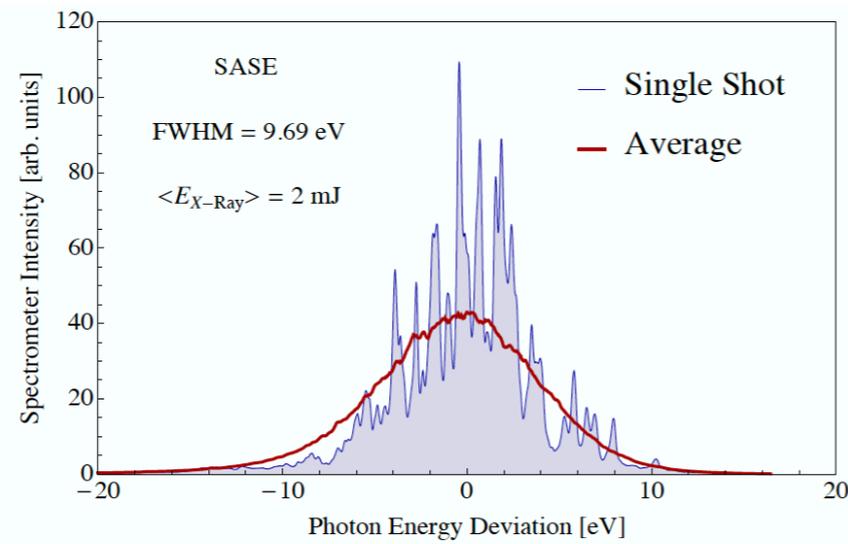
- Oscillator-fed tapered FELs have two advantages over fresh-bunch seeding:
  - Narrow bandwidth seed
  - Stable intensity, no shot-to-shot fluctuations
- Pre-bunching with a small chicane can yield the ideal set of initial conditions
- Harmonic filters and appropriate tuning of the high gain tapered section can extend the scheme to shorter wavelengths.
- Oscillator-fed high efficiency/harmonic FEL has not been tested experimentally but has received recent theoretical interest e.g. J. Duris *et al.*, <https://arxiv.org/pdf/1704.05030.pdf>, K. Kim, "A Harmonic X-ray FEL Oscillator," in High-Brightness Sources and Light-Driven Interactions

# Conclusion

- (1) We studied **undulator tapering strategies** to increase the efficiency of FELs and reach TW peak power levels.
- (2) **Diffraction** and the **sideband instability** were identified as the fundamental processes which limit the efficiency of tapered FELs.
- (3) We determine that a **large, stable seed** and a **pre-bunched** electron beam are the **ideal initial conditions** for high efficiency FEL.
- (4) We have studied this combination of pre-bunching and large seed as a in simulation. Results from 1-D sims are encouraging, 3-D sims to come.
- (5) The above situations have been investigated separately in experiment **but not yet together...**
- (6) Given recent interest in the FEL/science community/industry and limited FEL R&D time at large user facilities, a facility for testing advanced concepts for high efficiency FELs would be a welcome development for the field.

# Backup slides

# Fresh slice self seeding experiment at LCLS



# Sideband suppression via gain modulation

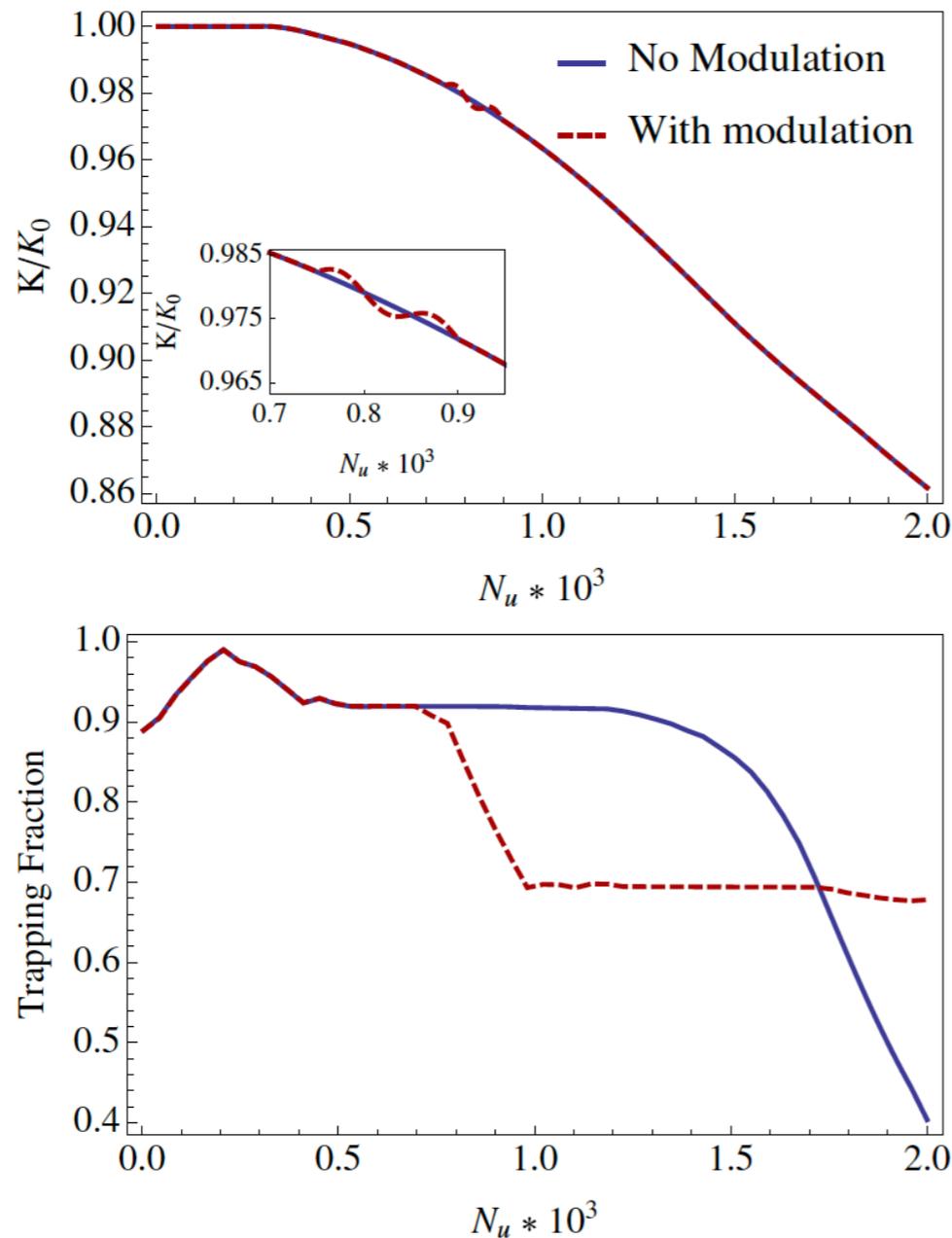


FIG. 5. (Top) Undulator taper profile for a gain modulated tapered FEL. The modulation section at  $N_u = 750$  changes the synchrotron frequency and damps the sideband growth (see Fig. 6). (Bottom) The trapping fraction drops after the modulation section but remains constant compared to the unmodulated case which suffers from severe sideband-induced detraping after  $N_u = 1500$ .

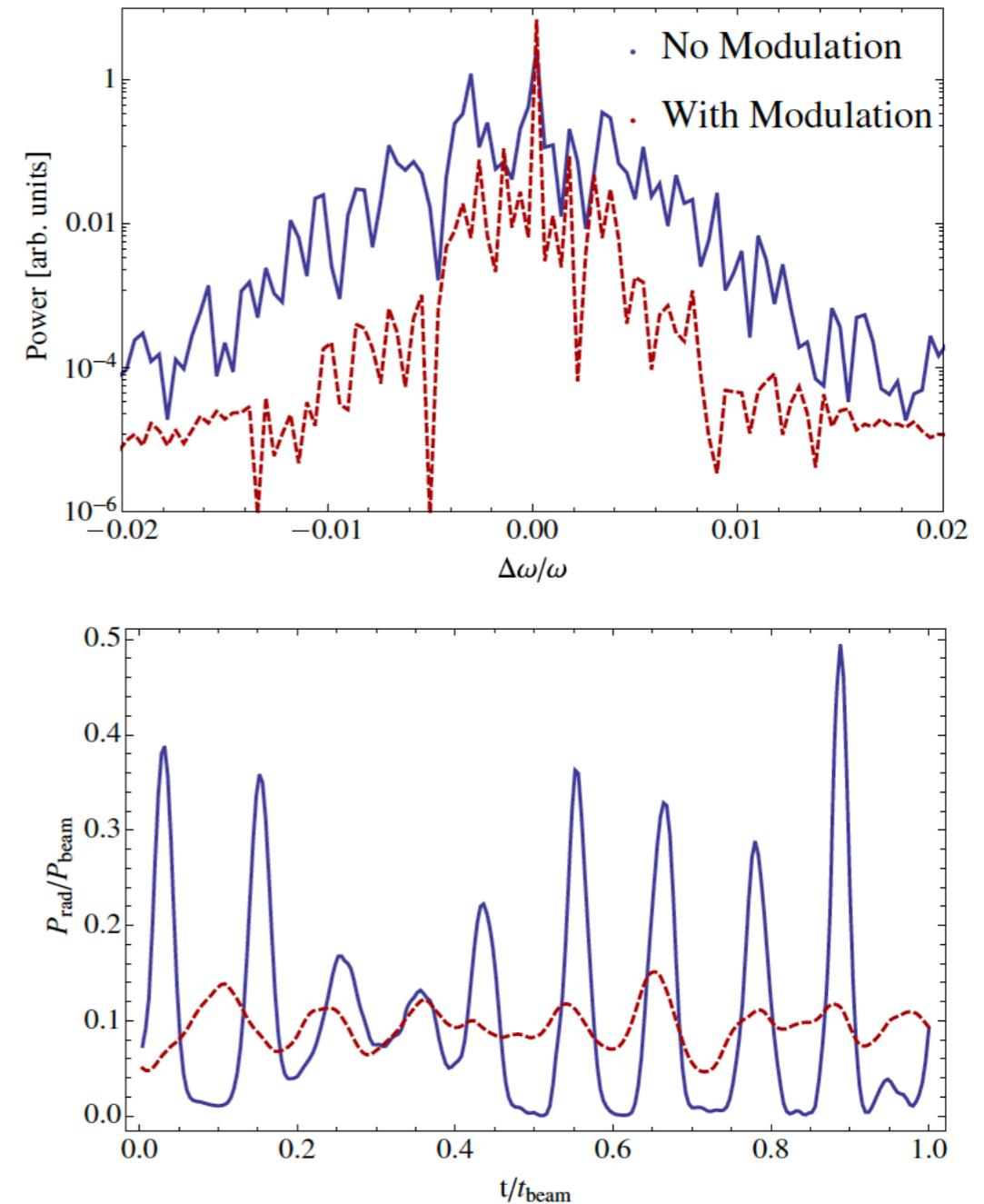


FIG. 6. Radiation spectrum (top) and temporal profile (bottom) with and without gain modulation showing sideband reduction for a gain modulated high efficiency FEL. The ratio of sideband to total power is 55% in the unmodulated case and 4% in the modulated case.