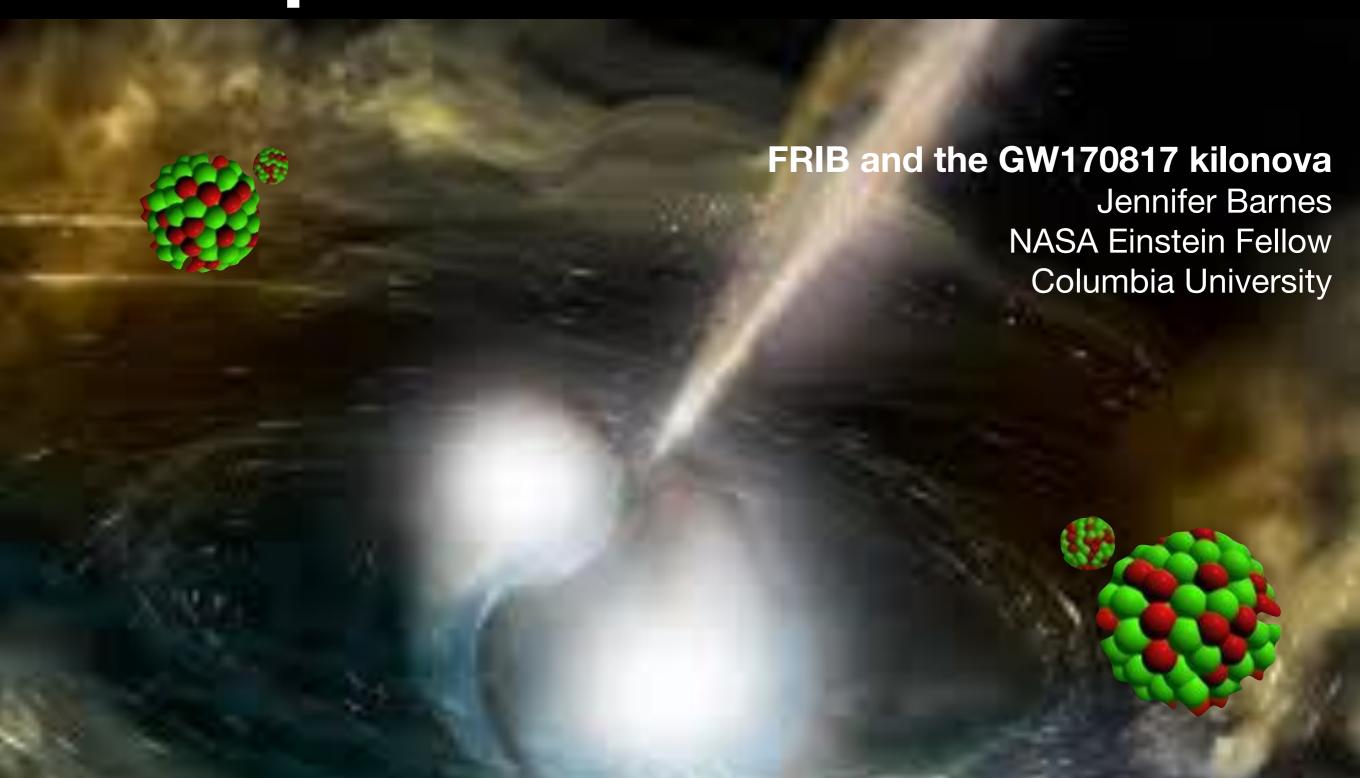
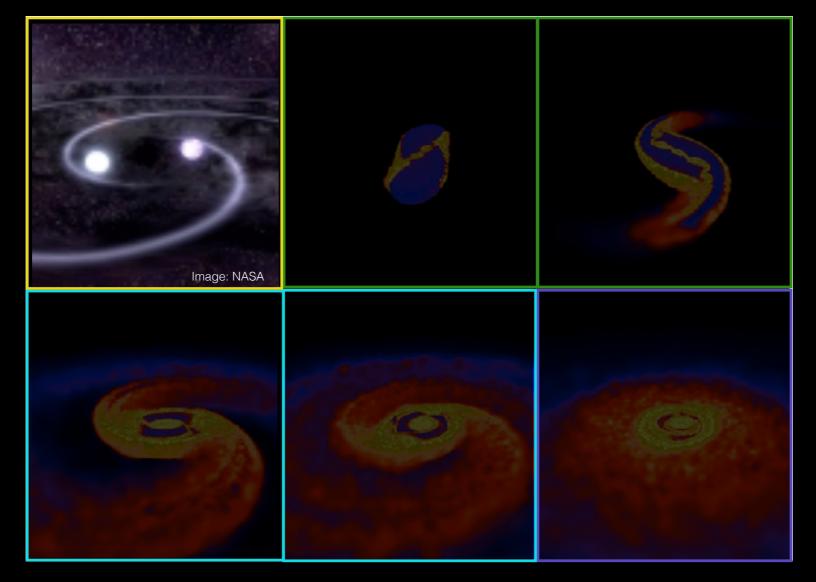
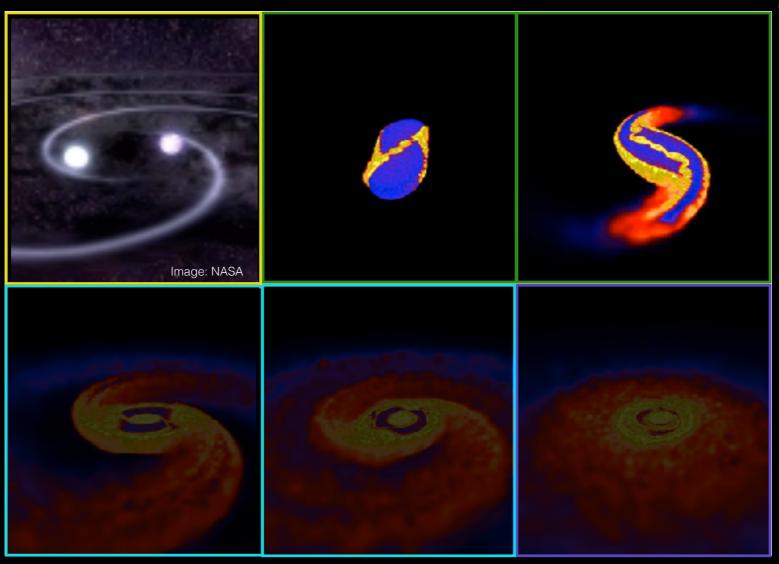
# Kilonova signatures and the *r*-process



### final few orbits: strong GW source



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merger: neutron star is partially disrupted, central remnant forms

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ejecta: some material is escapes; some is bound

e.g. Lattimer & Schramm 1974, 1976 Li & Paczynski 1998

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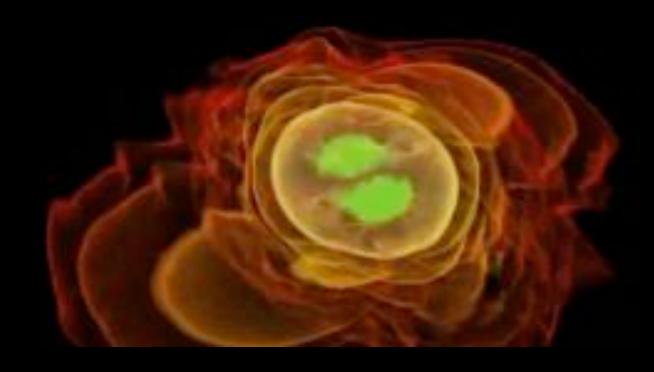
ejecta: some material is escapes; some is bound

final: a central NS or BH, an accretion disk, unbound ejecta

e.g. Lattimer & Schramm 1974, 1976 Li & Paczynski 1998

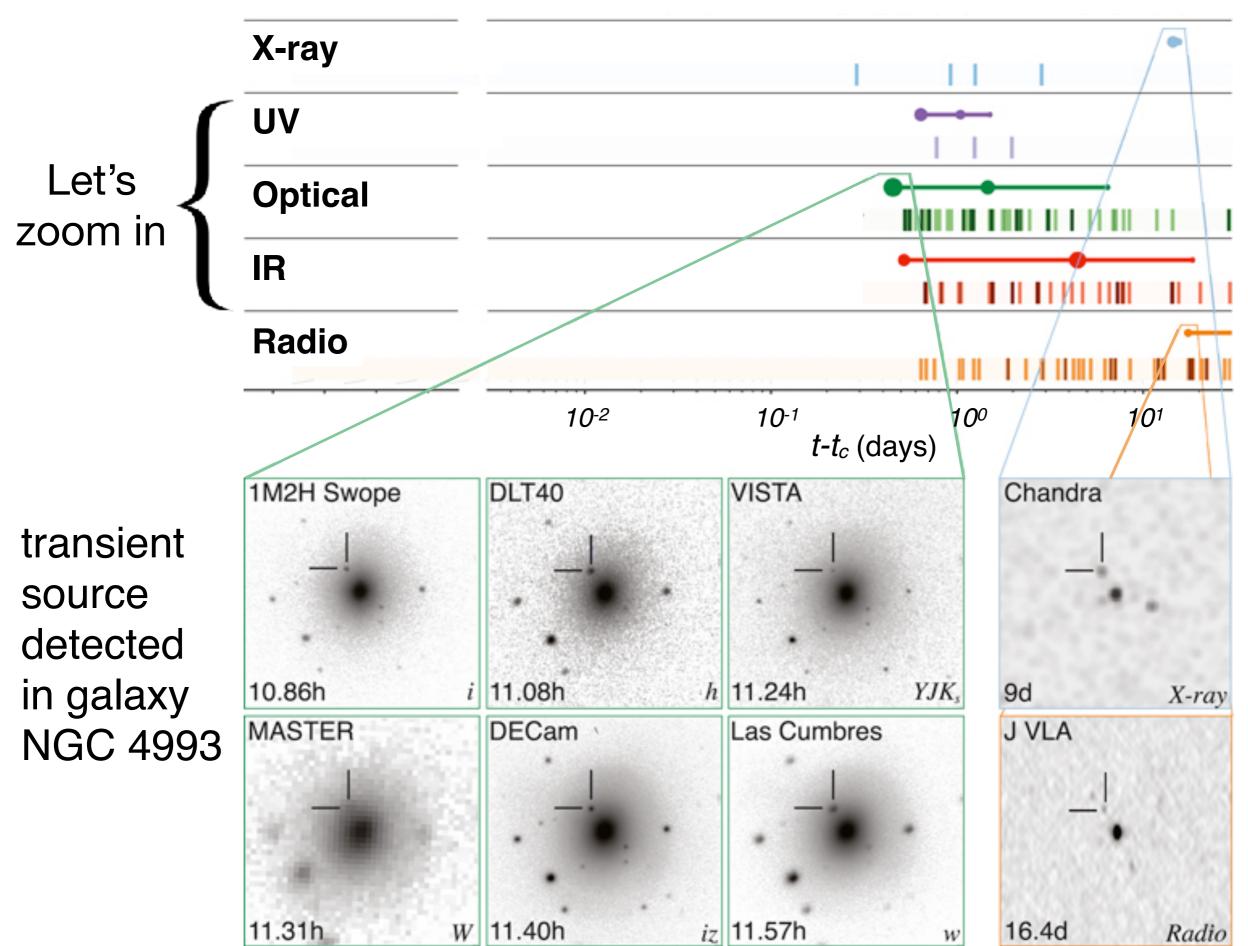
# radioactive transients are probes of the *r*-process

#### "kilonova"

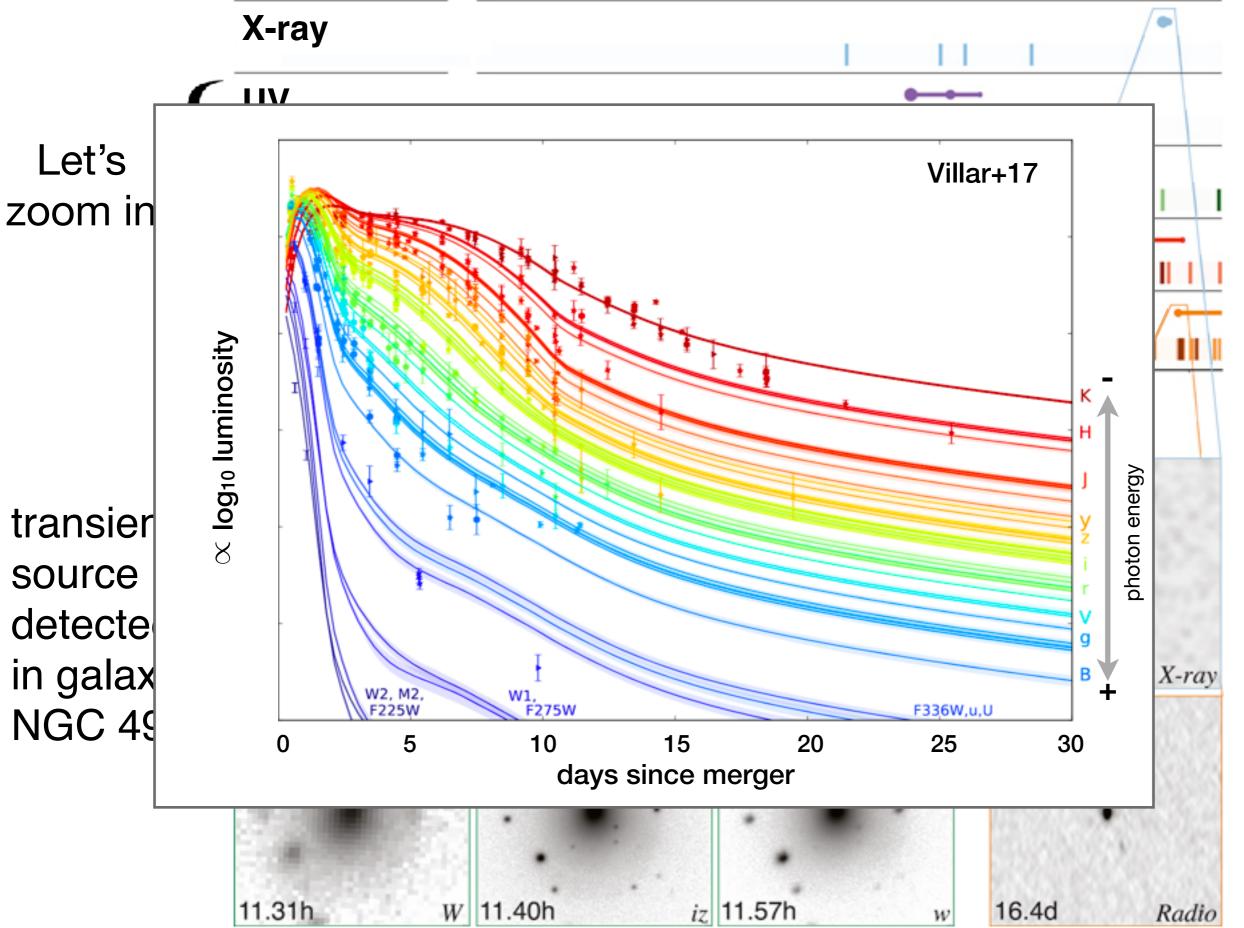


- Mildly relativistic unbound material
- Heavy elements are synthesized

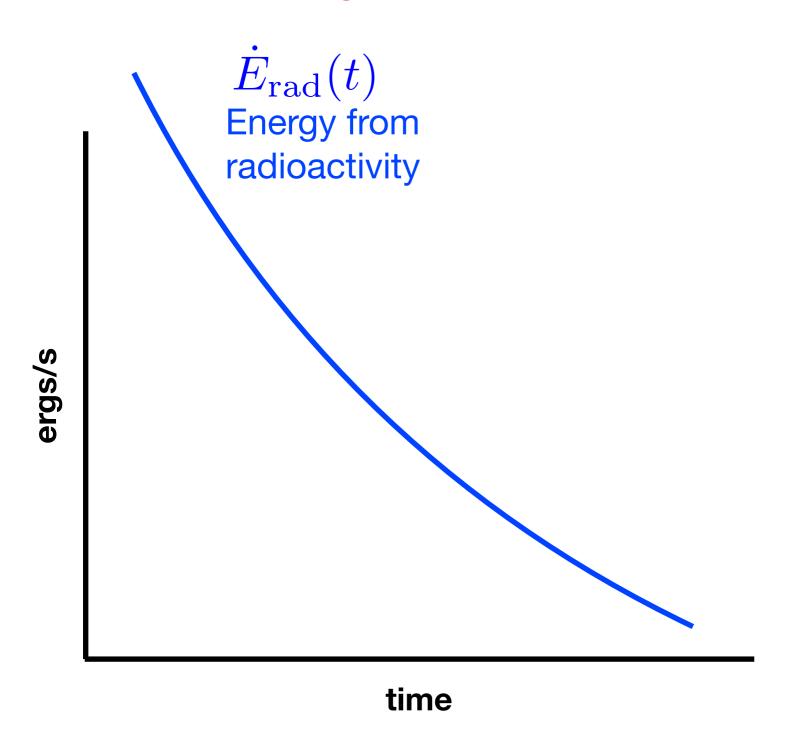
An expanding cloudheated by radioactive decays

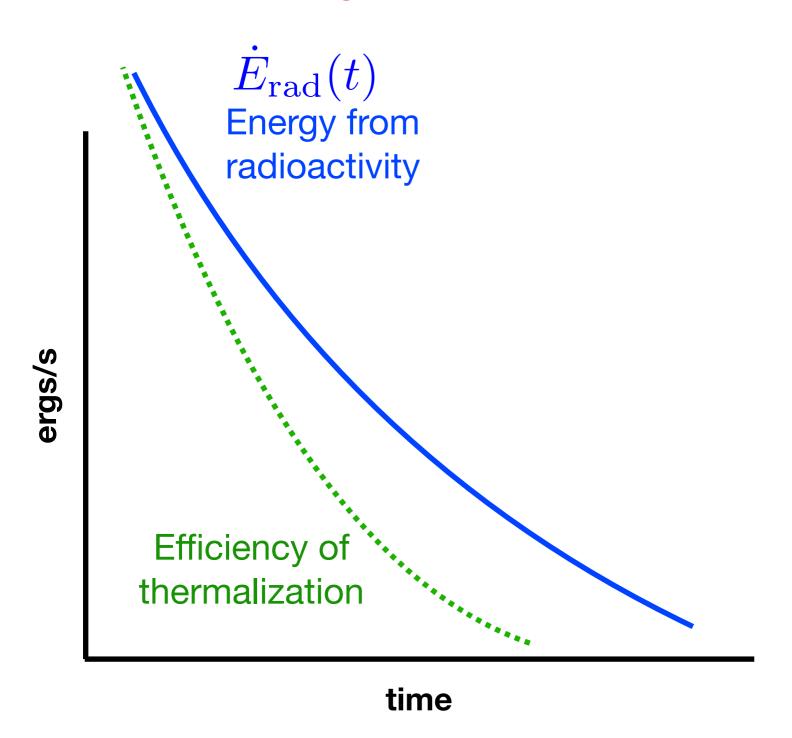


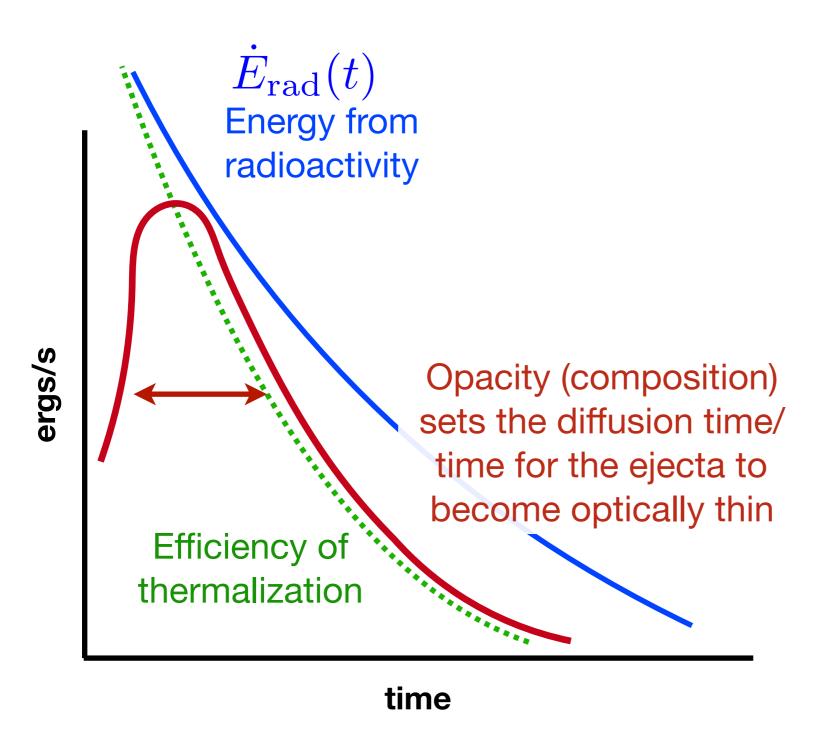
ad. from ALV + EM Partners 17



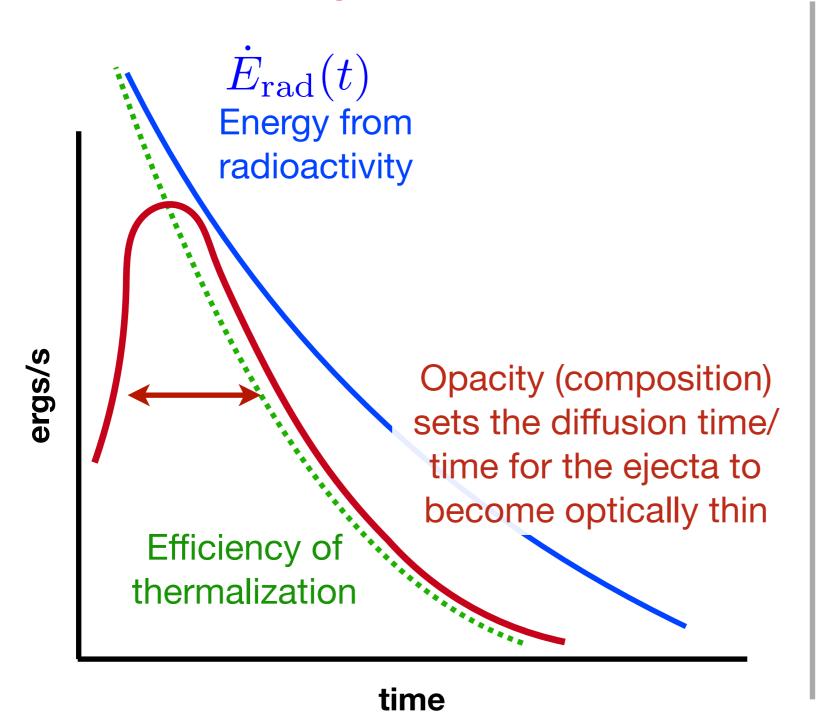
ad. from ALV + EM Partners 17







#### (bolometric) light curves



#### colors & spectra

- Quasi-blackbody with temperature set by the net effect of radioactivity, thermalization, photon absorption/ emission, and cooling
- Line-blanketing can affect the spectrum
- Individual features correspond to particular atoms or ions

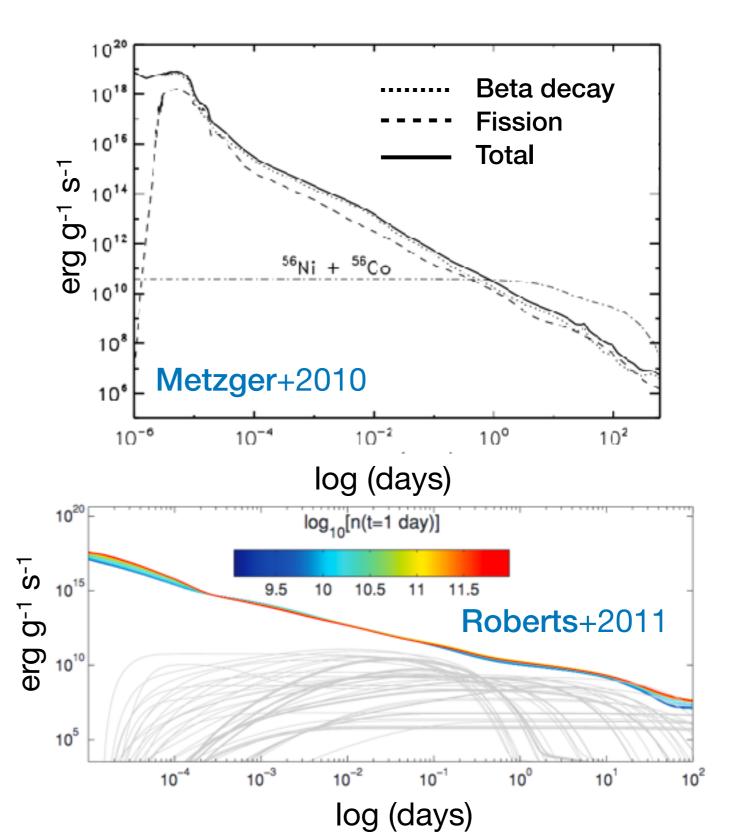
# Because there are many contributing decays, $\dot{E}_{\rm rad}(t)$ follows a ~power law

The expression 
$$\dot{E}_{\rm rad}(t) = \frac{fc^2}{t}$$

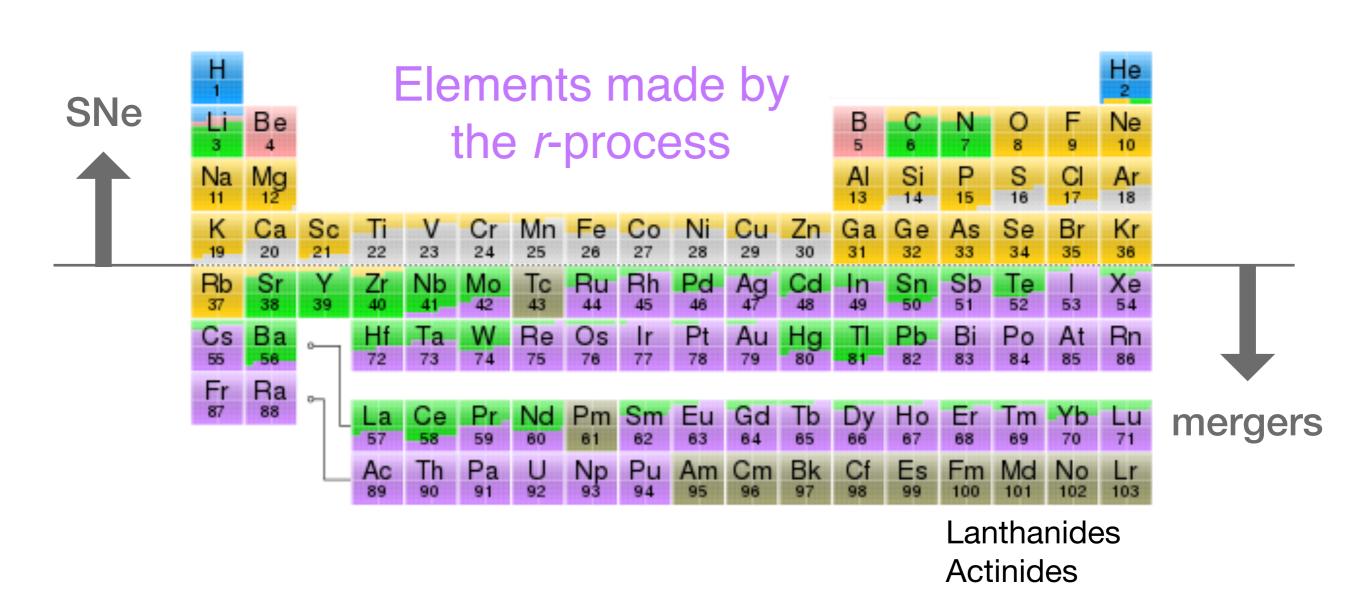
was first derived **analytically** (Li & Paczyński 1998; see also Hotokezaka+17)

The basic behavior has since been borne out by nuclear network calculations

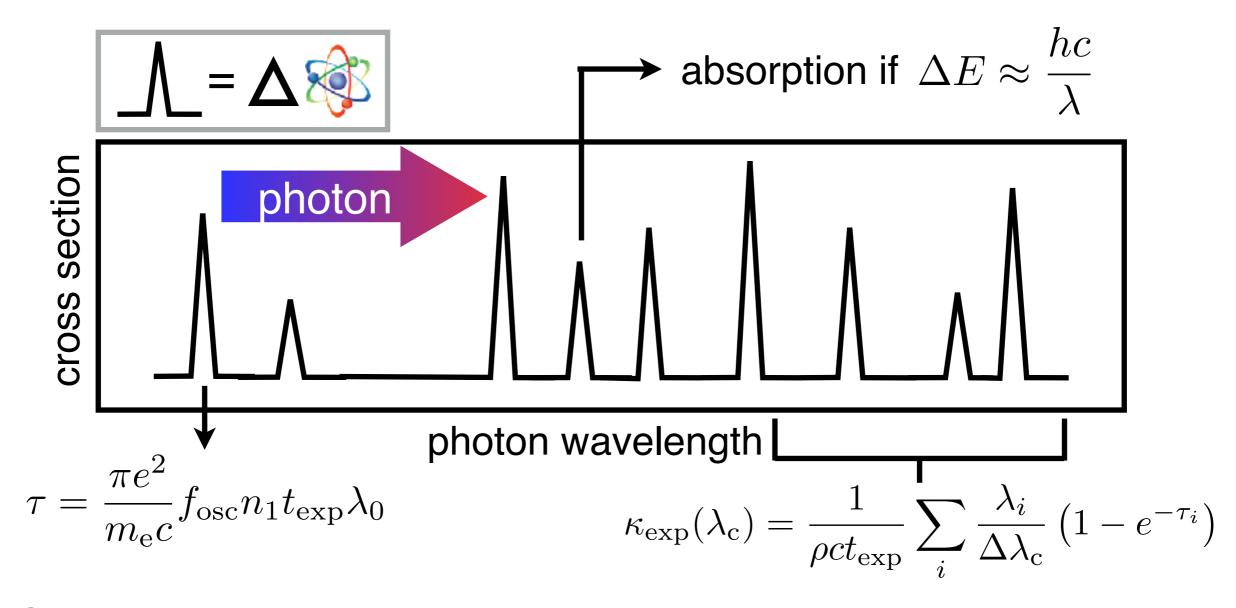
 however, the power-law behavior may break down at late times.



The *r*-process produces elements with atomic structures that are unique among explosively-synthesized compositions.



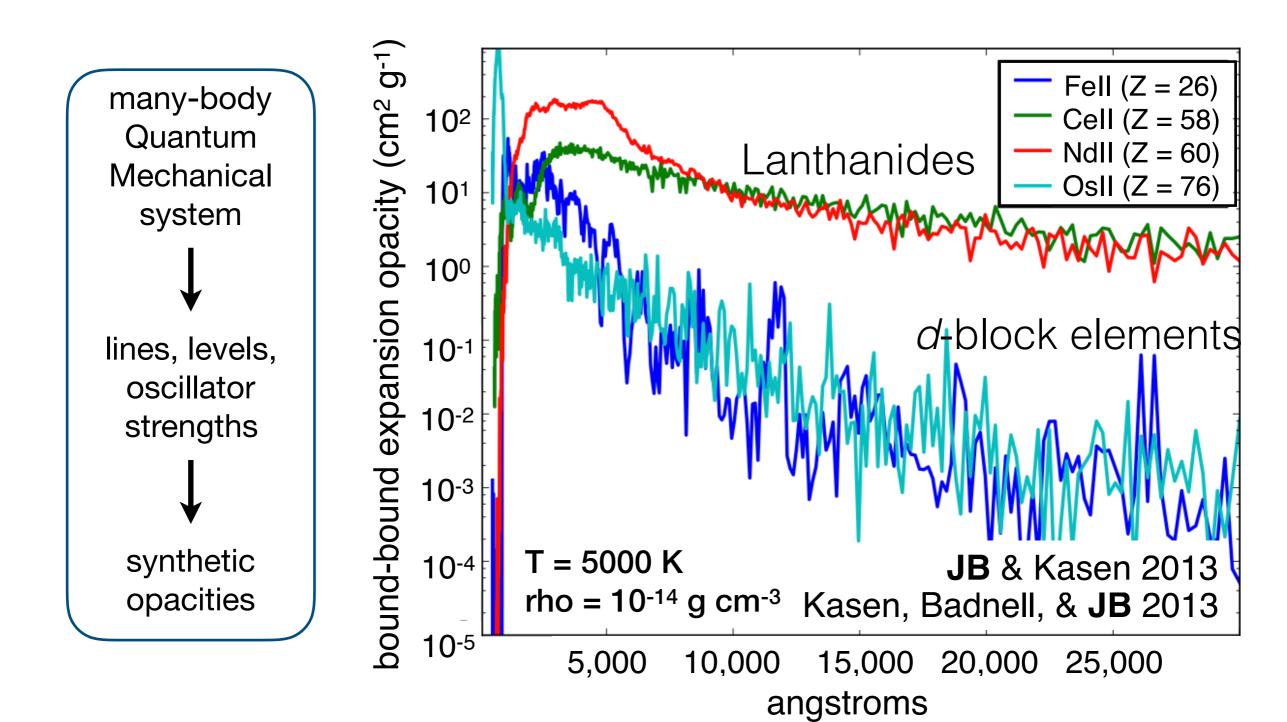
• **Bound-bound** opacity (cm<sup>2</sup> g<sup>-1</sup>) sets the photon mean free path.



Sobolev optical depth sets interaction probability with a particular line

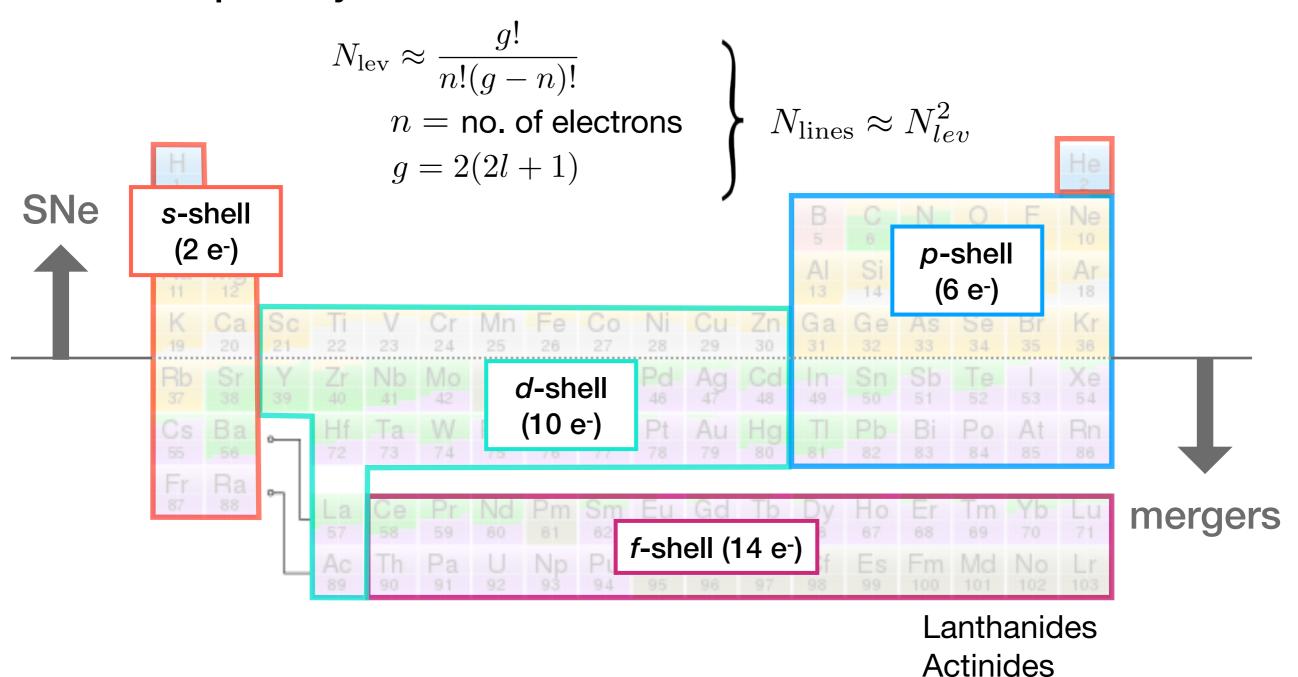
The expansion opacity determines the effective continuum opacity

- Atomic structure modeling compensates for missing data
- Lanthanides/actinides increase the opacity



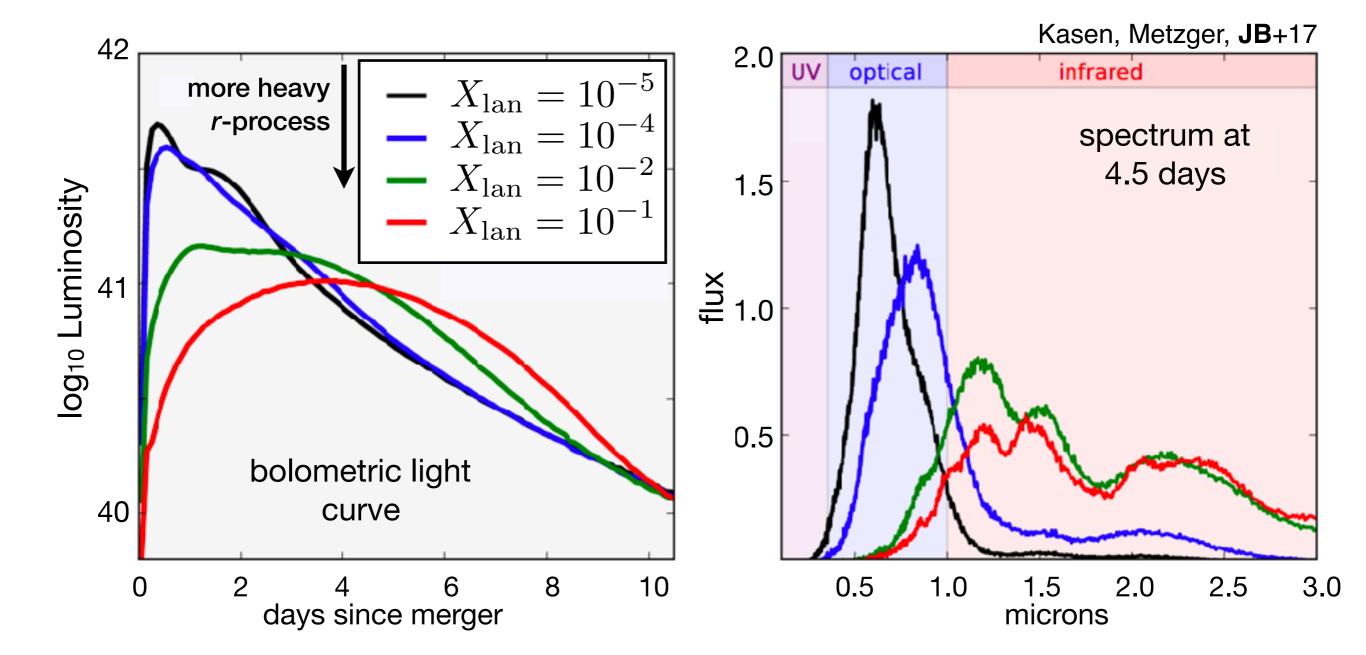
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#### Simple analytic estimates:



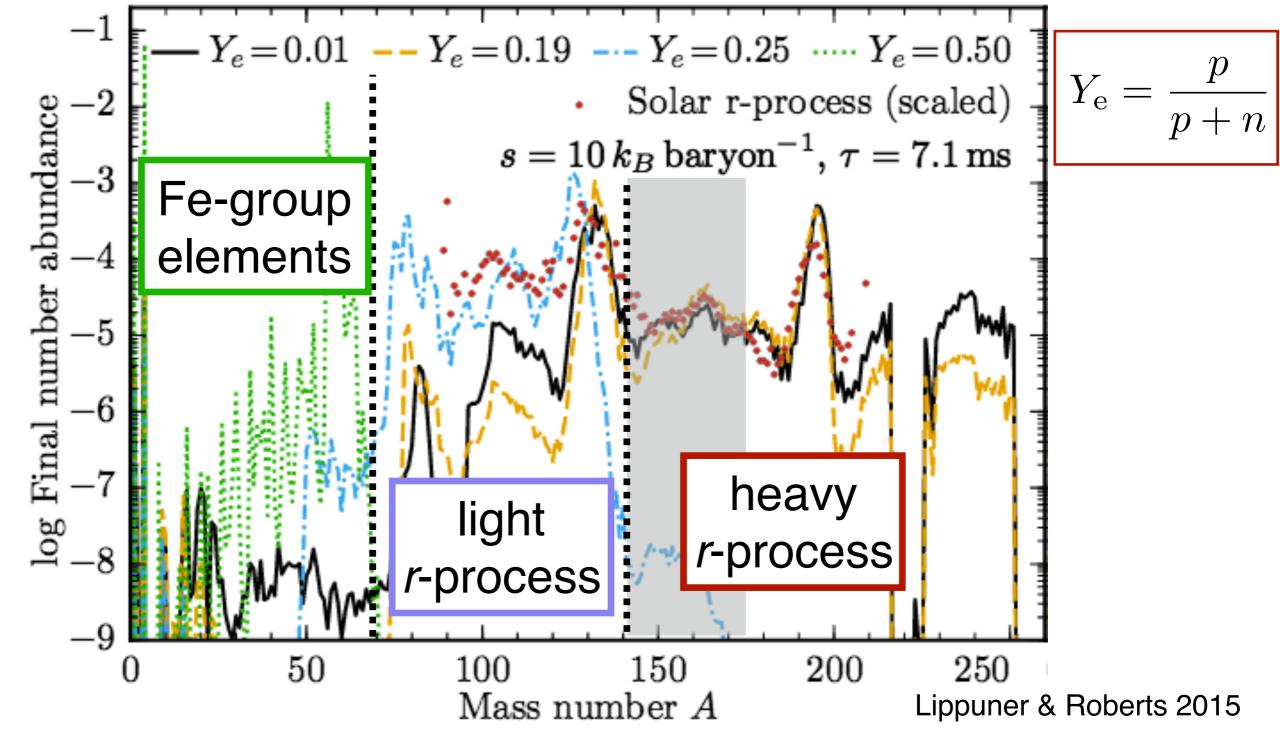
## higher opacities lead to longer, dimmer, redder light curves

diffusion time:  $t_{\rm diff} pprox \left(\frac{M\kappa}{vc}\right)^{1/2}$  adiabatic losses:  $E_{\rm phot} \sim t^{-1}$  line blanketing at optical wavelengths



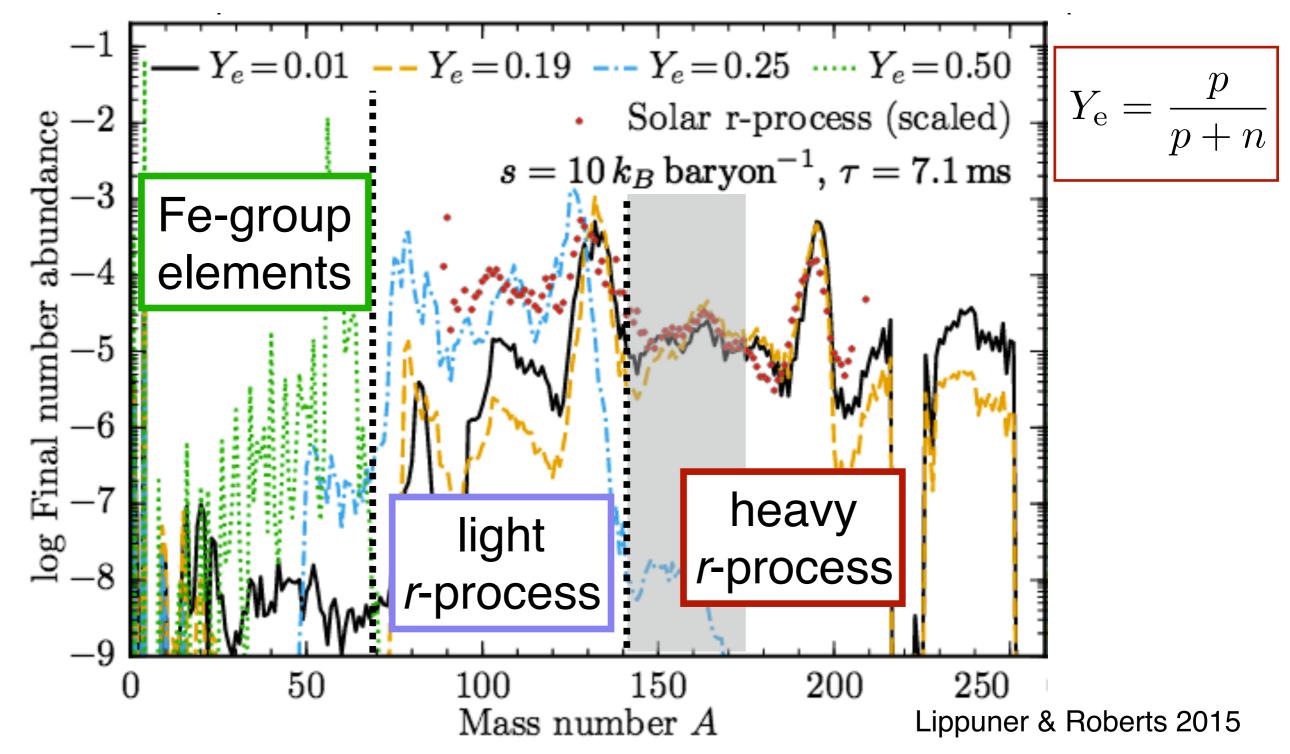
# kilonova emission is tied to the strength of the *r*-process!

fewer free n per seed ← \_\_\_\_ more free n per seed

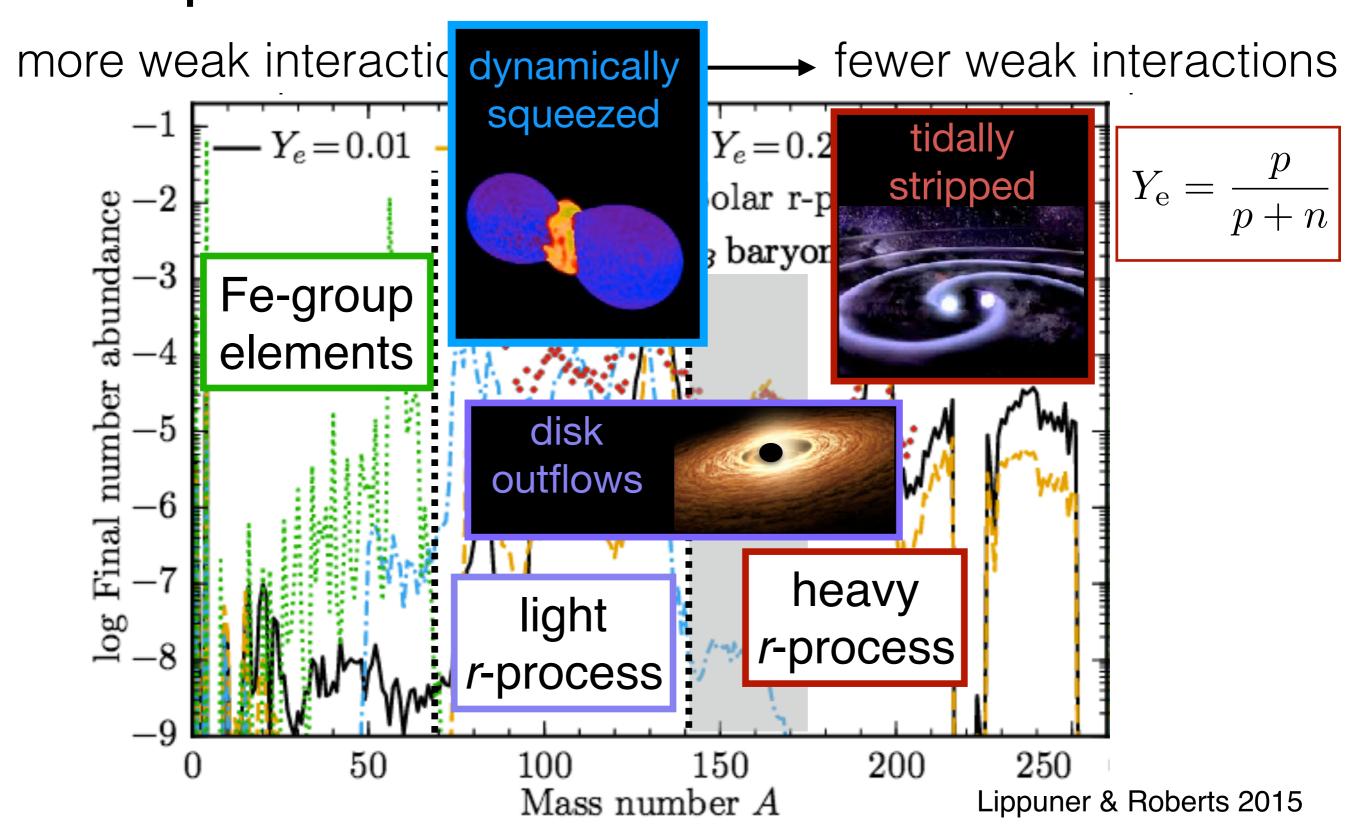


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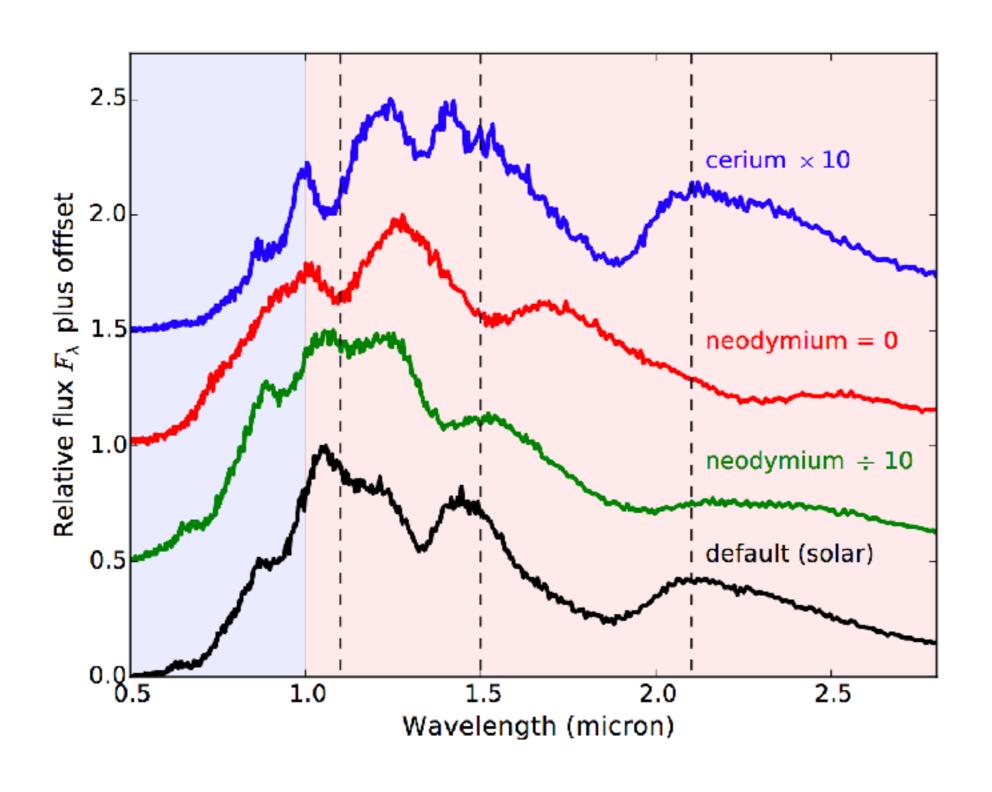
more weak interactions ————— fewer weak interactions



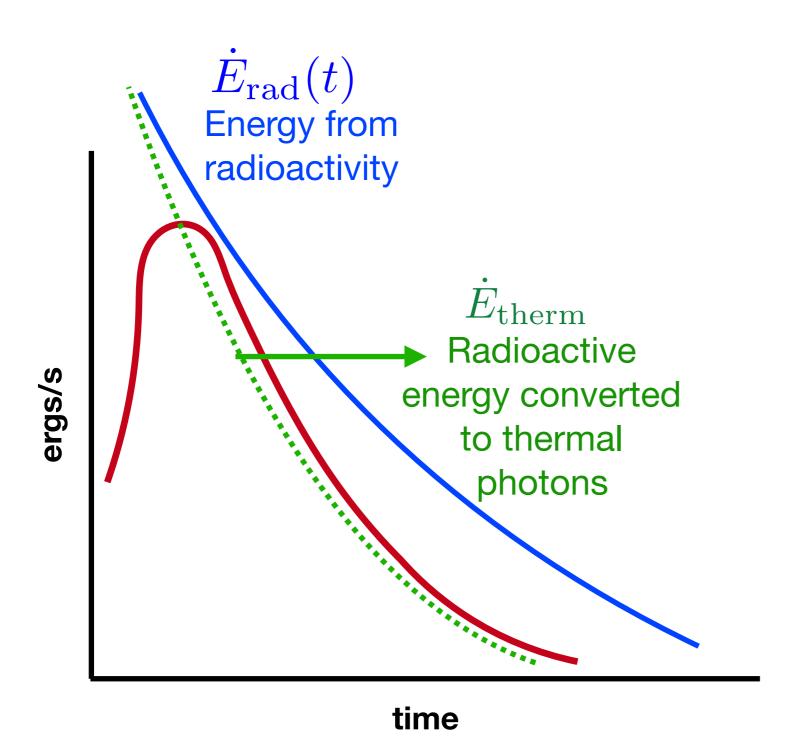
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#### spectral identification: the next frontier!

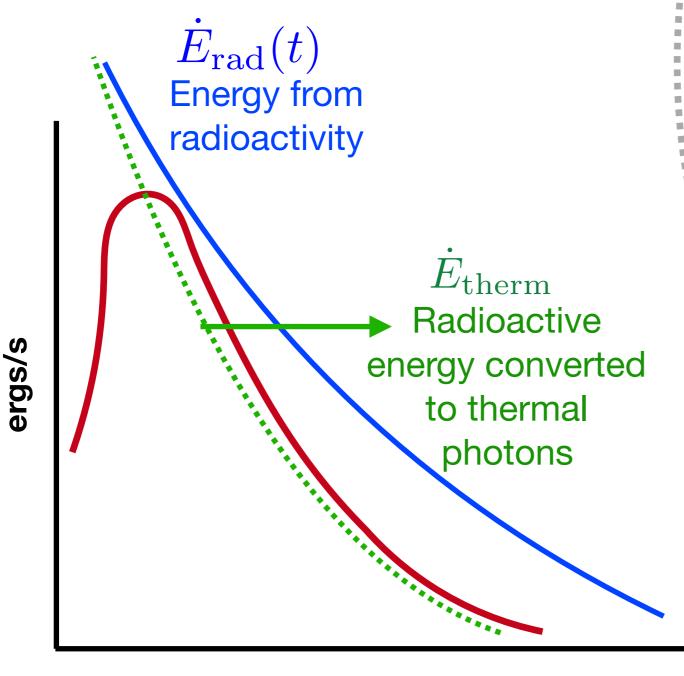


### the r-process and kilonova thermalization



#### the r-process and kilonova thermalization

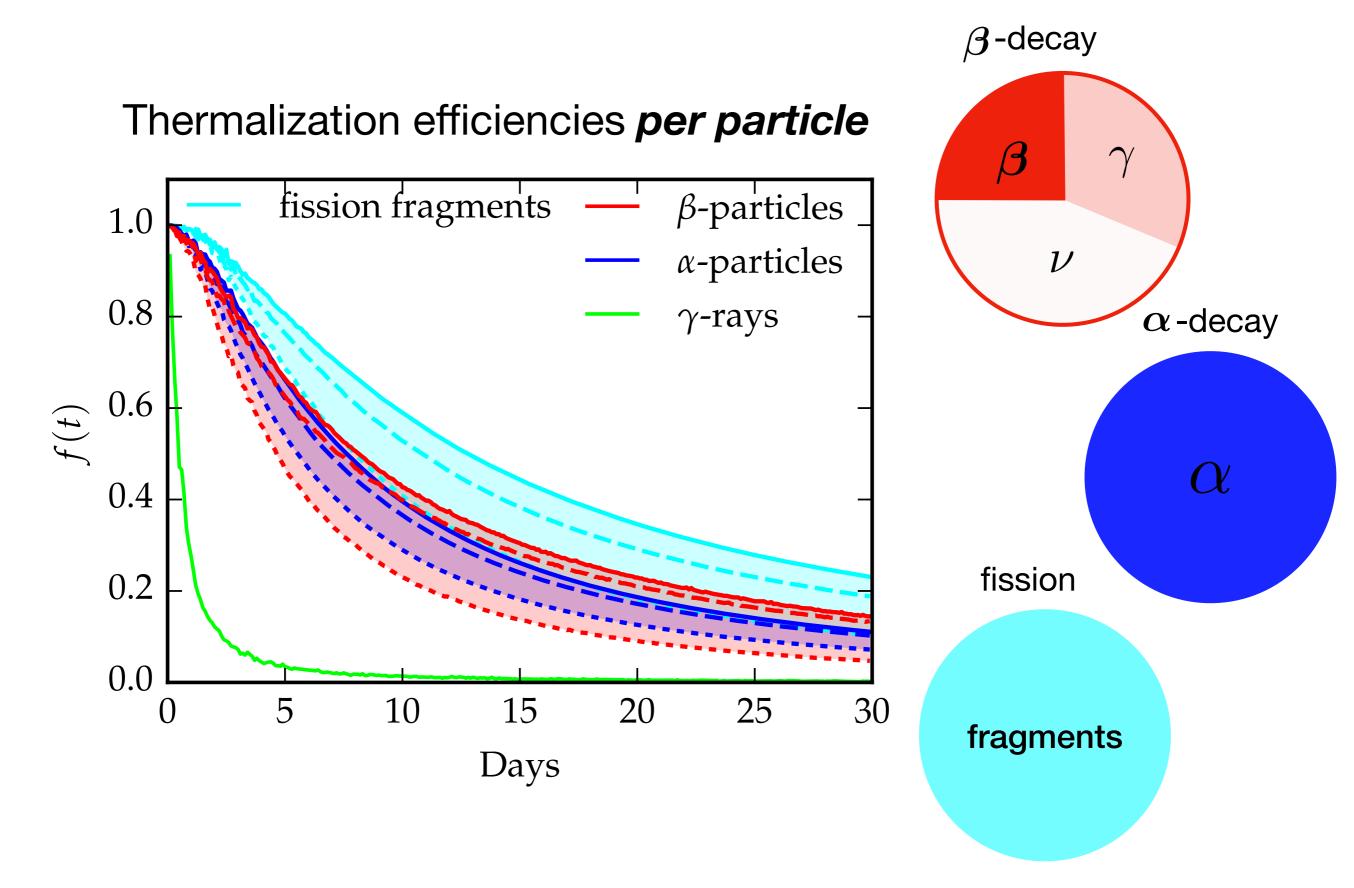




### thermalization efficiency depends on:

- decay mode
- decay spectra
- composition (crosssections)
- ejecta mass, velocity

### thermalization depends on decay mode



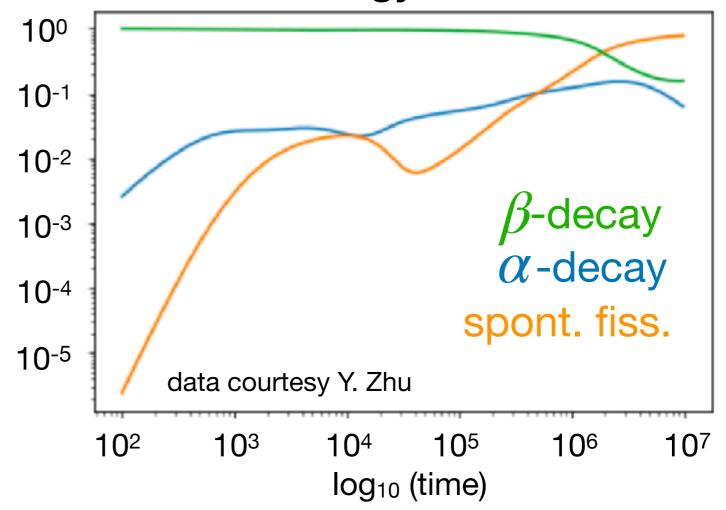
### At late times, a few decays may dominate the heating

Fissioning or  $\alpha$ -decaying nuclei with  $\tau \sim$  weeks or months could substantially affect the luminosity

- High Q-values (compared to  $\beta$ -decay)
- Efficient thermalization

#### Fraction of energy in each channel

# Californium-254 $Q_{\rm SF} \approx 200 \, {\rm days}$ $au_{1/2} = 60.5 \, {\rm MeV}$

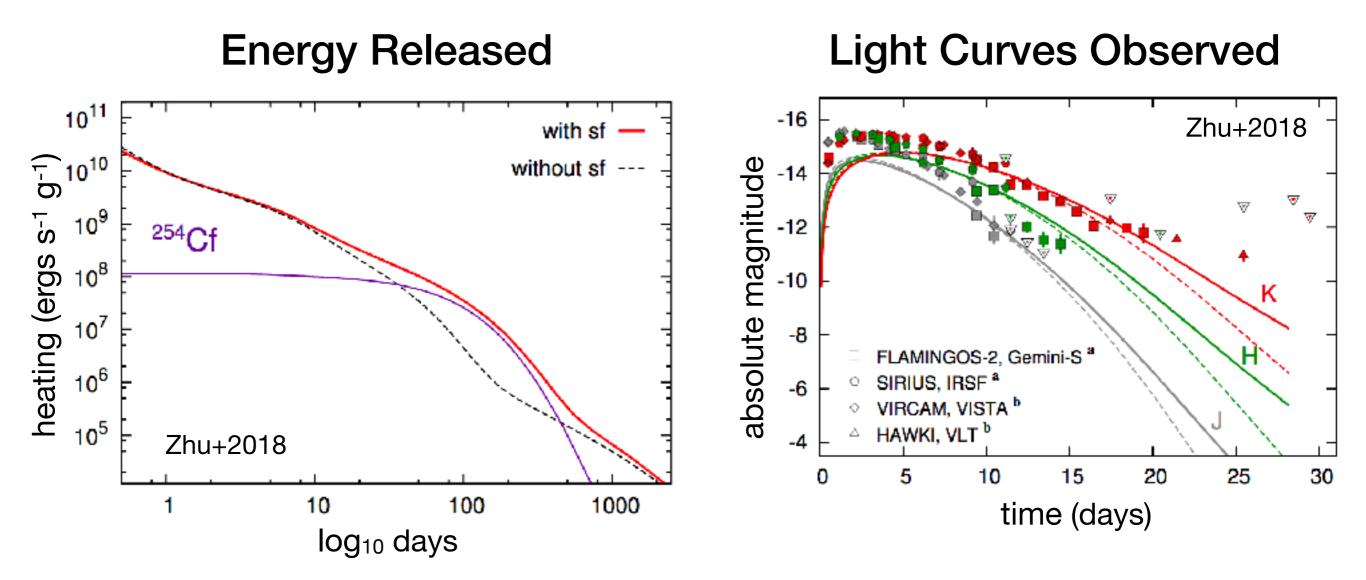


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# Heating from spontaneous fission of Cf-254 impacts kilonova light curves



Late-time light curves can probe the production of the heaviest nuclei and give more detailed information about the composition