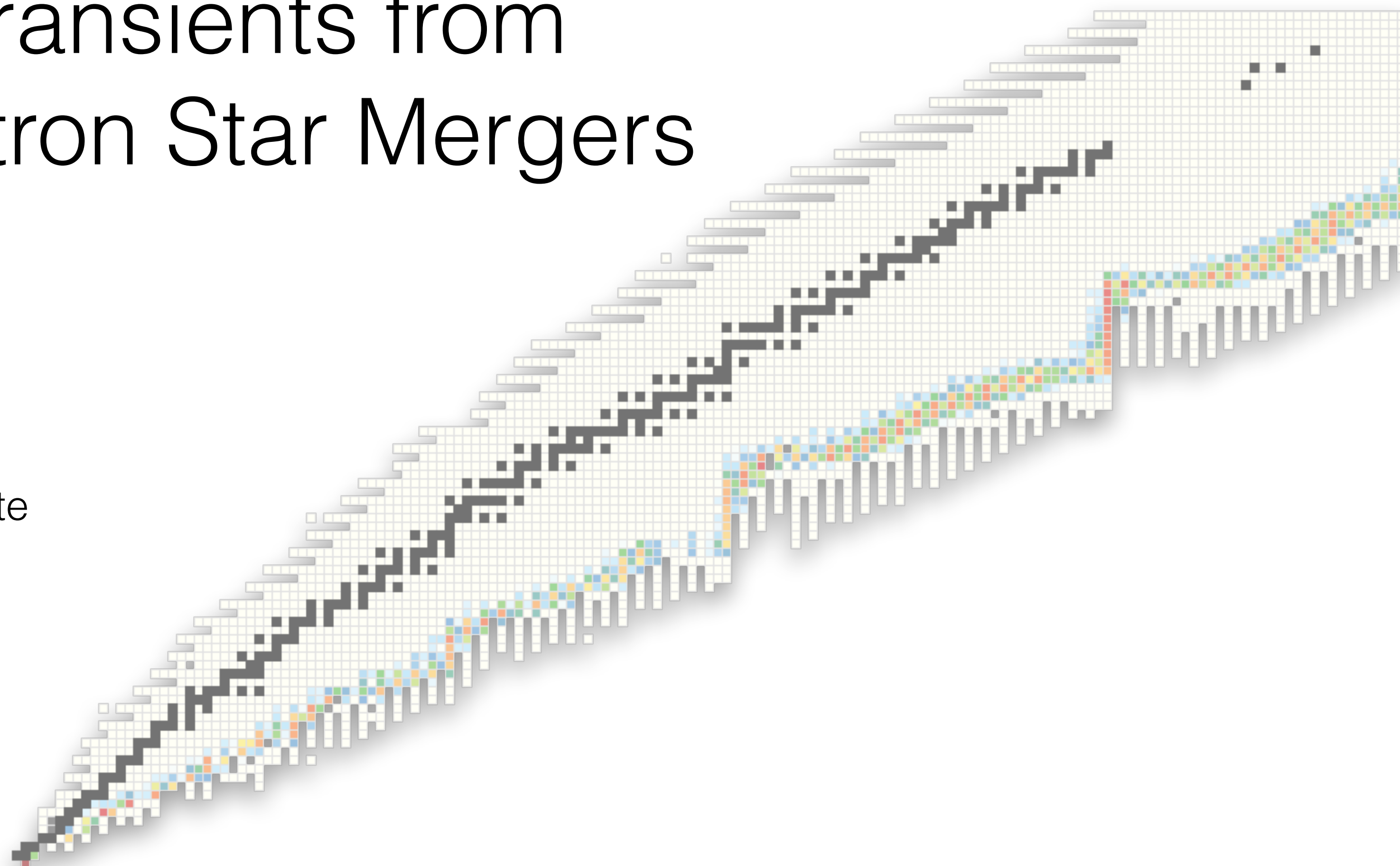


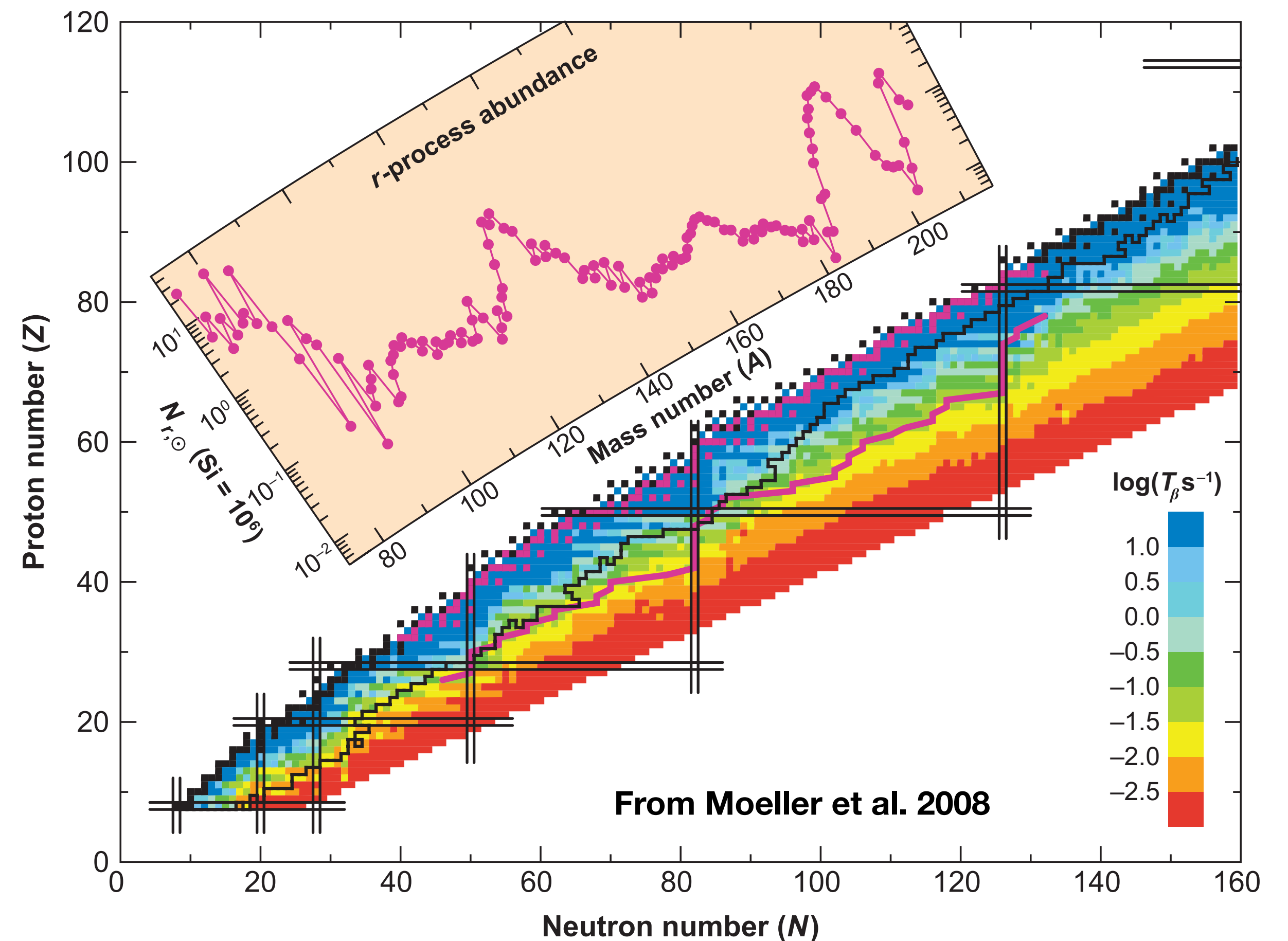
# Nucleosynthesis and Electromagnetic Transients from Neutron Star Mergers

Luke Roberts  
NSCL, Michigan State  
University



# What is the source of the *r*-process nuclei?

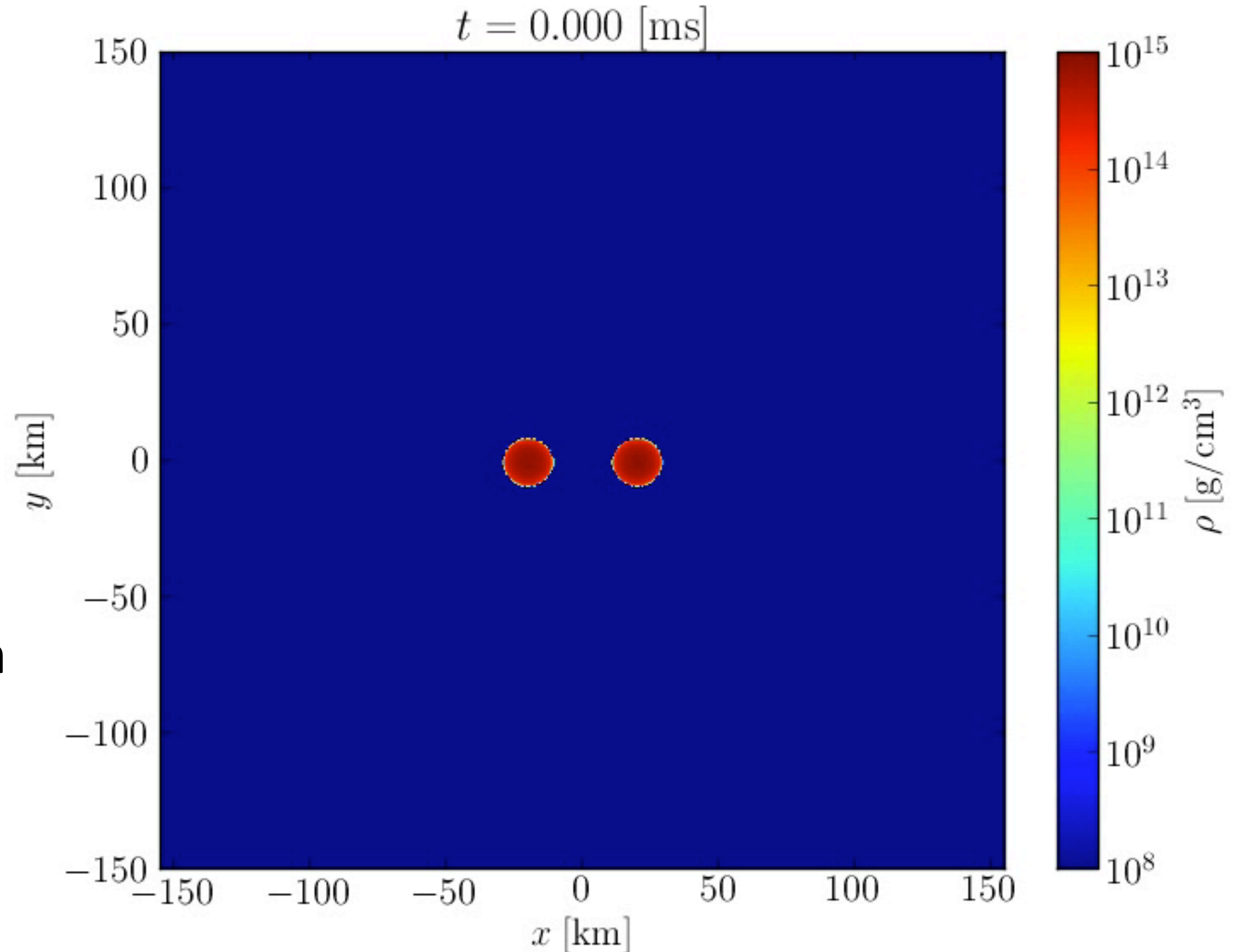
- *r*-process elements present in very low metallicity halo stars, suggesting it must be a primary process
- Abundance pattern of second and third peak *r*-process elements in low metallicity halo stars is remarkably similar to the pattern found in the sun
- Need lots of free neutrons
  - Site is one of the biggest questions of nuclear astrophysics
  - CCSNe have long been implicated as the site of the *r*-process
  - With GW170817, mounting evidence that NS mergers may be the site





# Merger Mass Ejection

- Dynamical Ejecta
  - Tidal Ejecta (BHNS)
  - GR  $\rightarrow$  matter ejected from collision region (NSNS)
- Disk winds  
(e.g. Surman et al. '08, Wanajo et al. '11)
- Disk outflows from viscous heating and alpha recombination  
(e.g. Fernandez & Metzger '13, Just '14)



# Nuclear Evolution of the Ejecta

Dynamical Timescale for the Ejected Material:

$$\tau_{ej} \approx 10 \text{ ms}$$

Ejected Material is neutron rich:

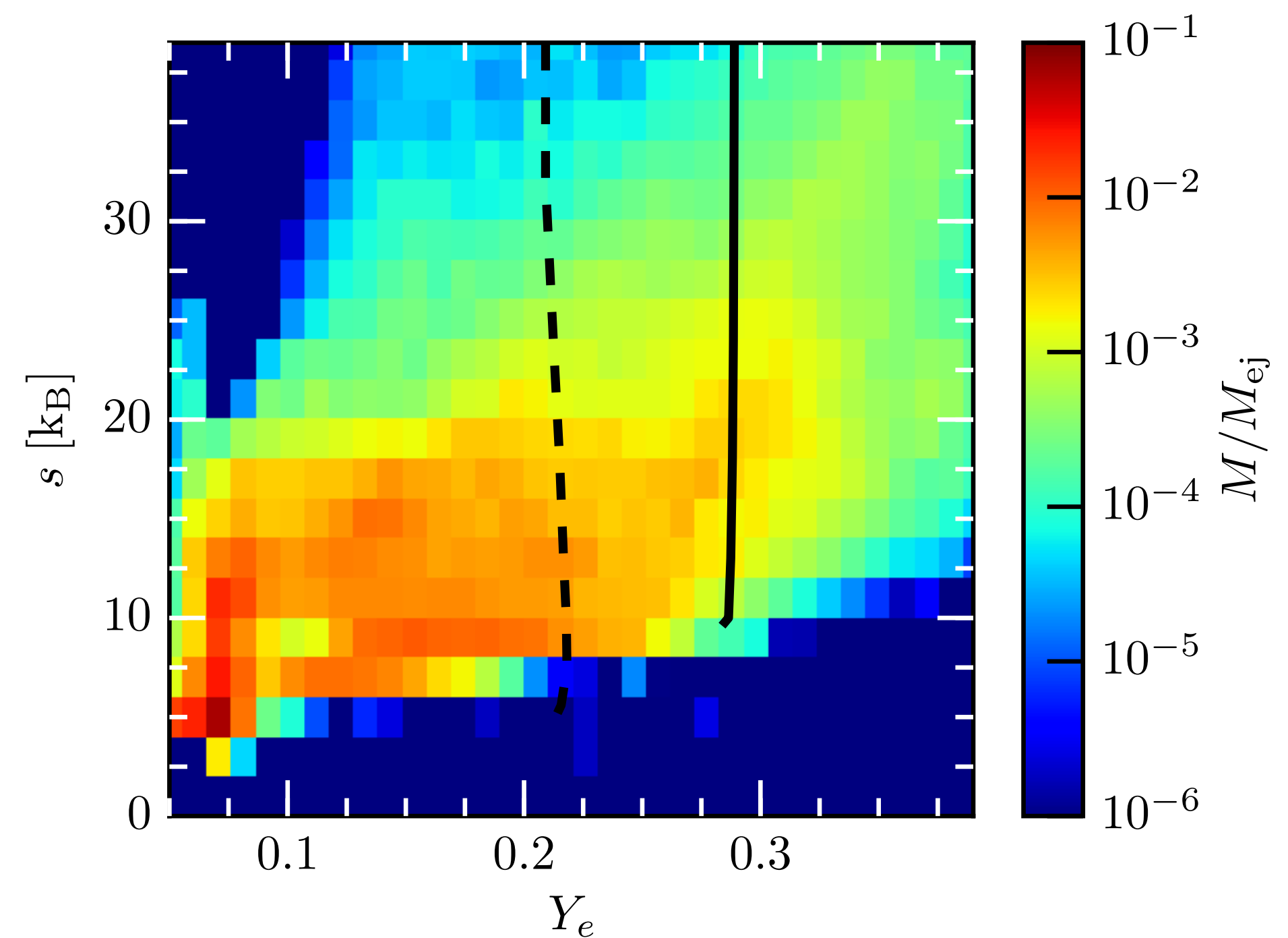
$$Y_e \sim 0.05 - 0.4$$

Low initial entropy:

$$S \sim 1 - 30$$

Which implies a neutron to seed ratio  
greater than 100

$$Y_e = 1 - \frac{n_{\text{neutrons,tot}}}{n_{\text{baryons}}}$$



Radice, et al. '16



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Initial distribution will be in NSE, clustered around doubly magic nuclei

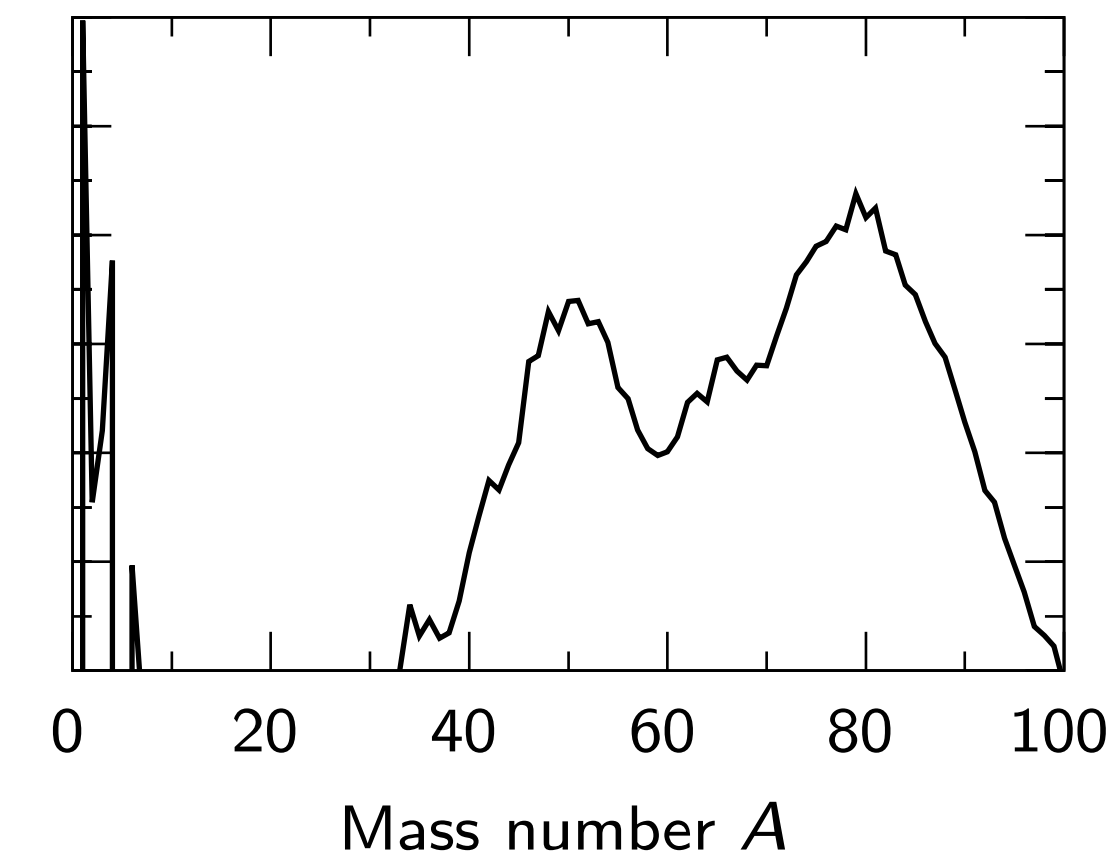
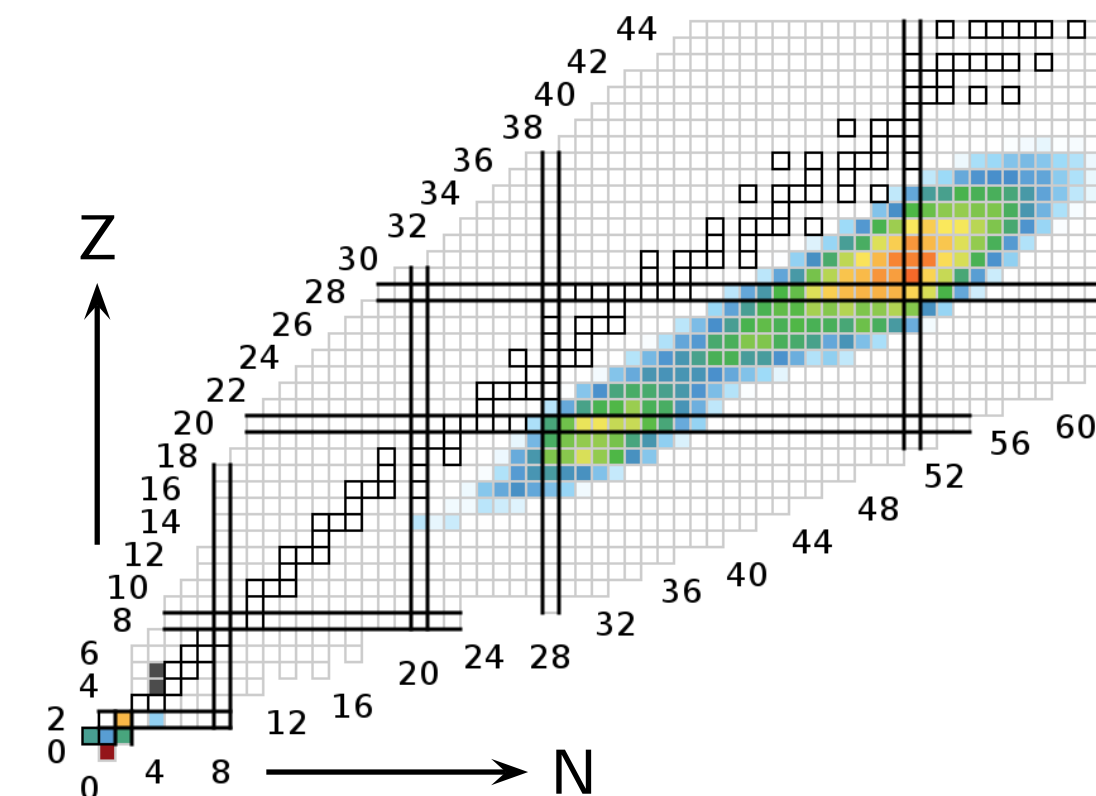
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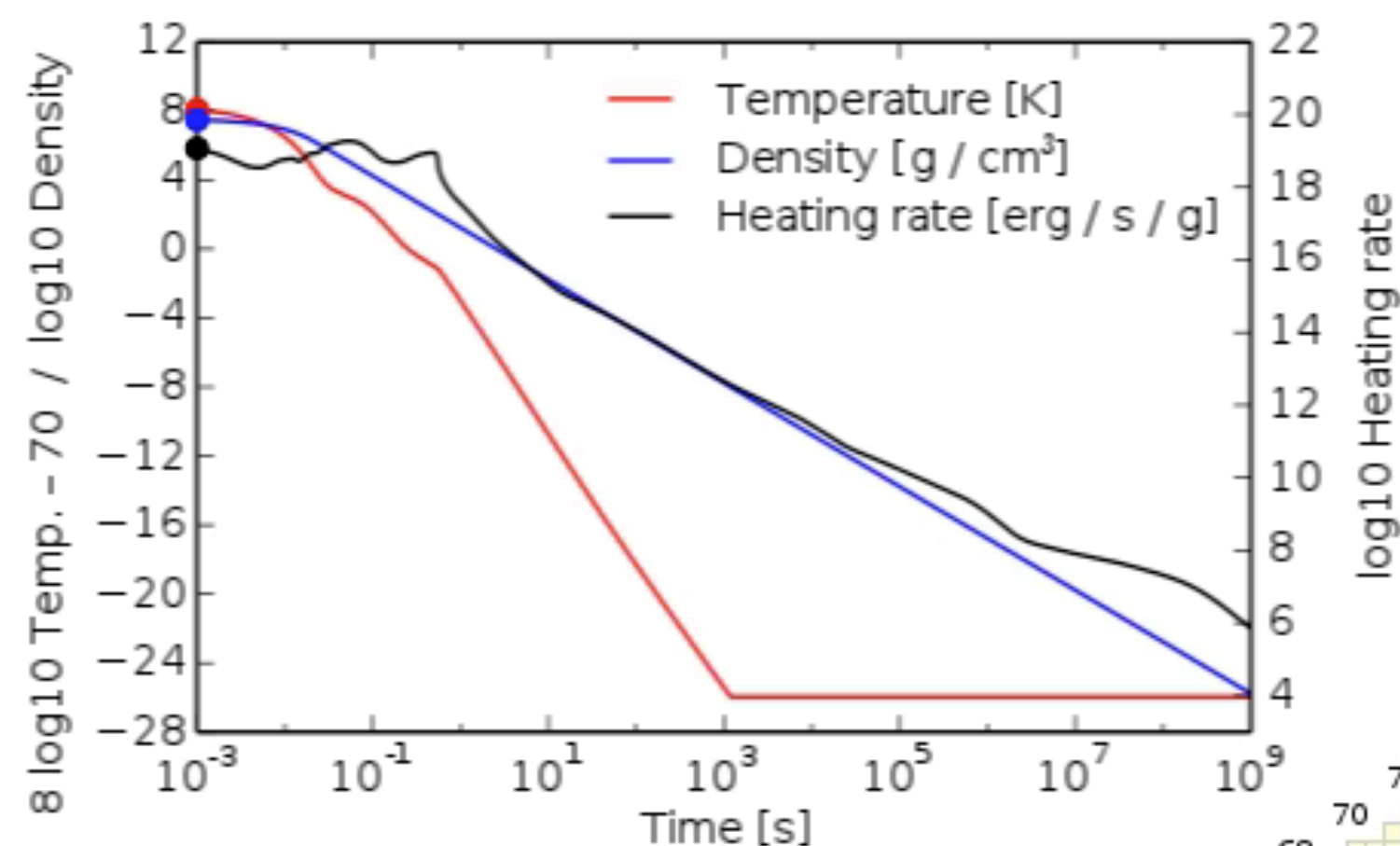
$$Y_e = 1 - \frac{n_{\text{neutrons,tot}}}{n_{\text{baryons}}}$$

$$T = 7.0 \text{ GK}$$

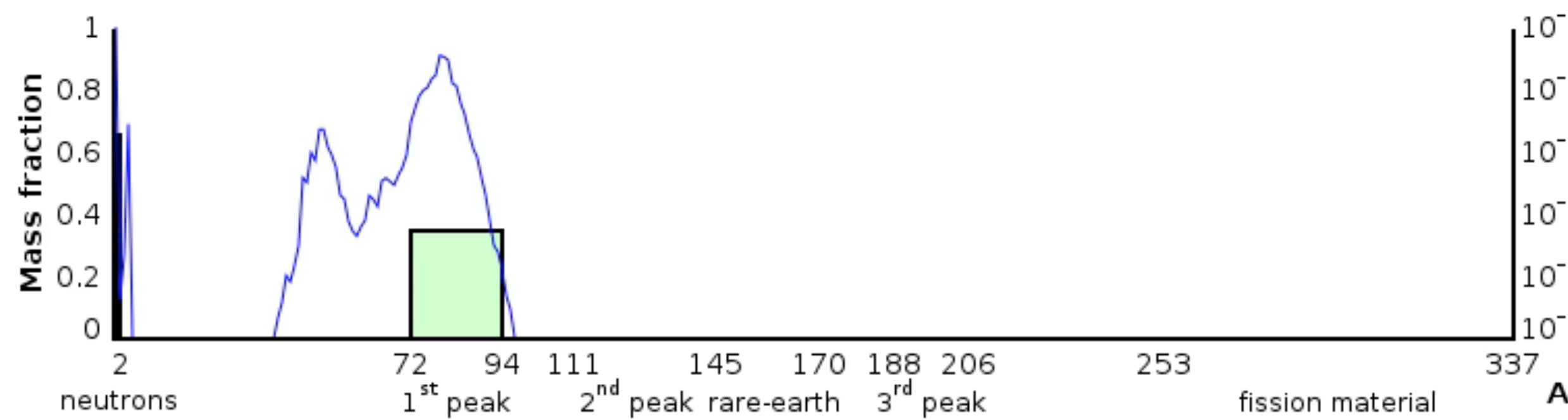
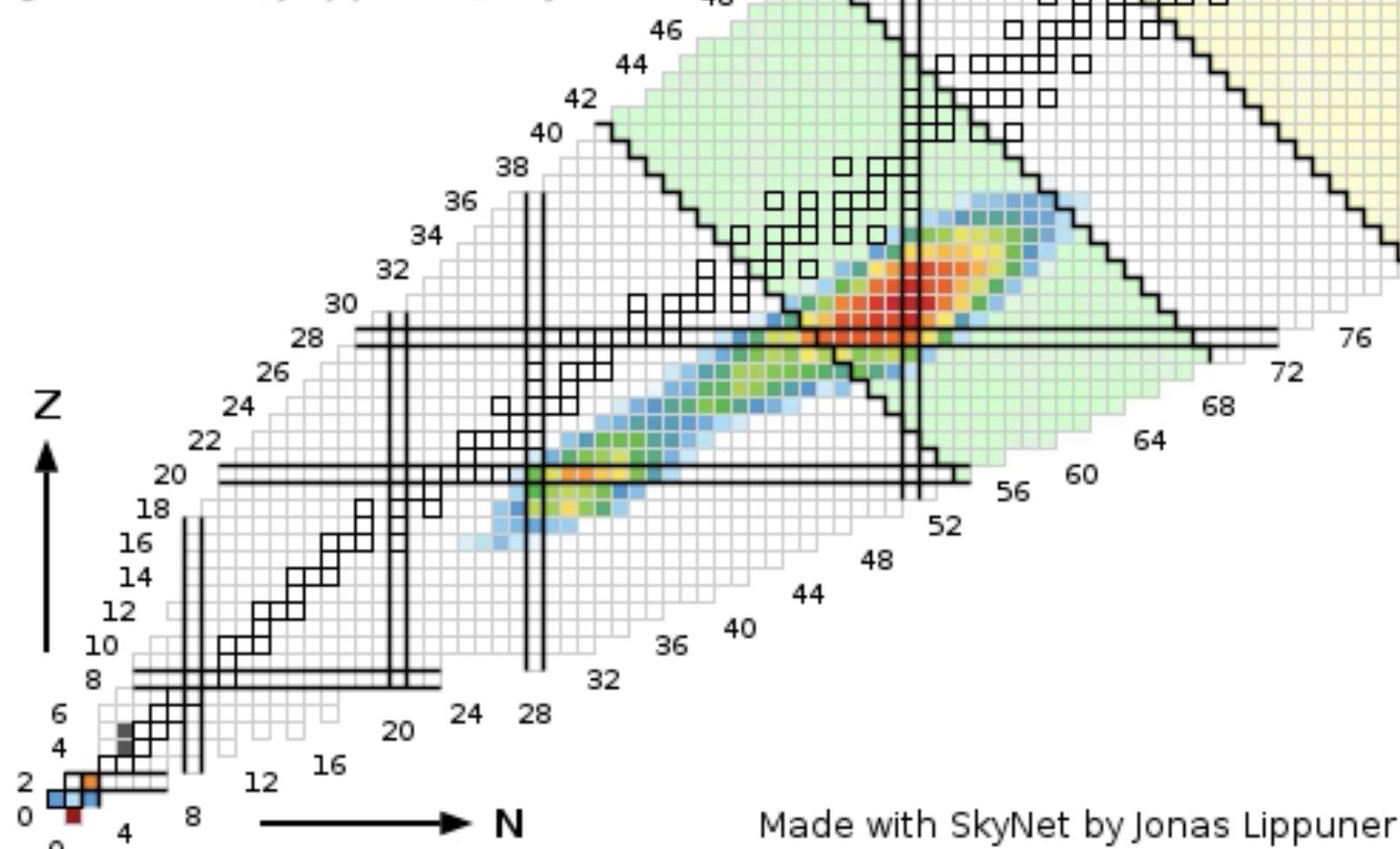
$$\rho = 2.2 \times 10^8 \text{ g cm}^{-3}$$

$$Y_e = 0.051$$



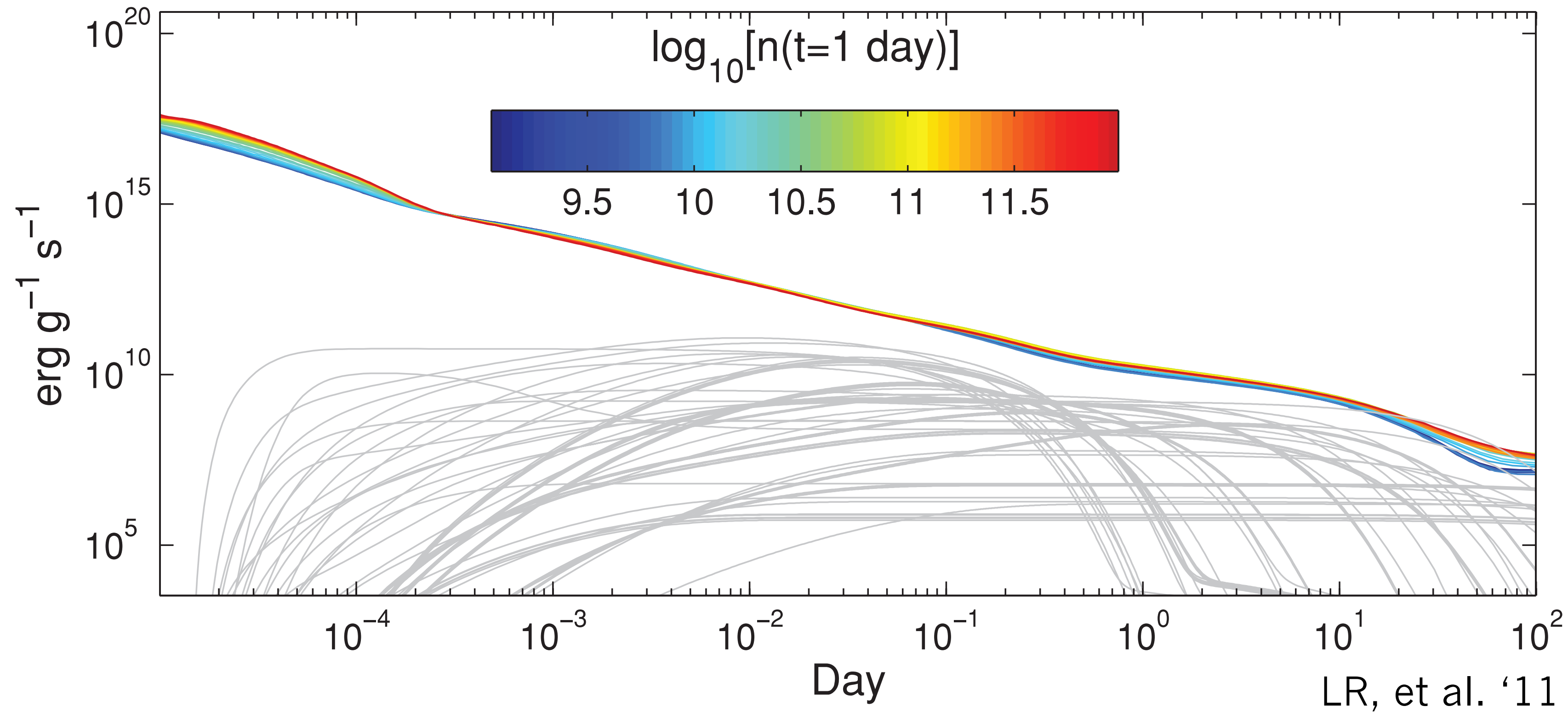


Temperature =  $5.71\text{E}+09$  K  
Density =  $3.04\text{E}+07$  g / cm<sup>3</sup>  
Heating rate =  $1.17\text{E}+19$  erg / s / g  
Entropy =  $9.99\text{E}+00$  kB / baryon  
Ye = 0.130



from Lippuner & LR, et al. '15

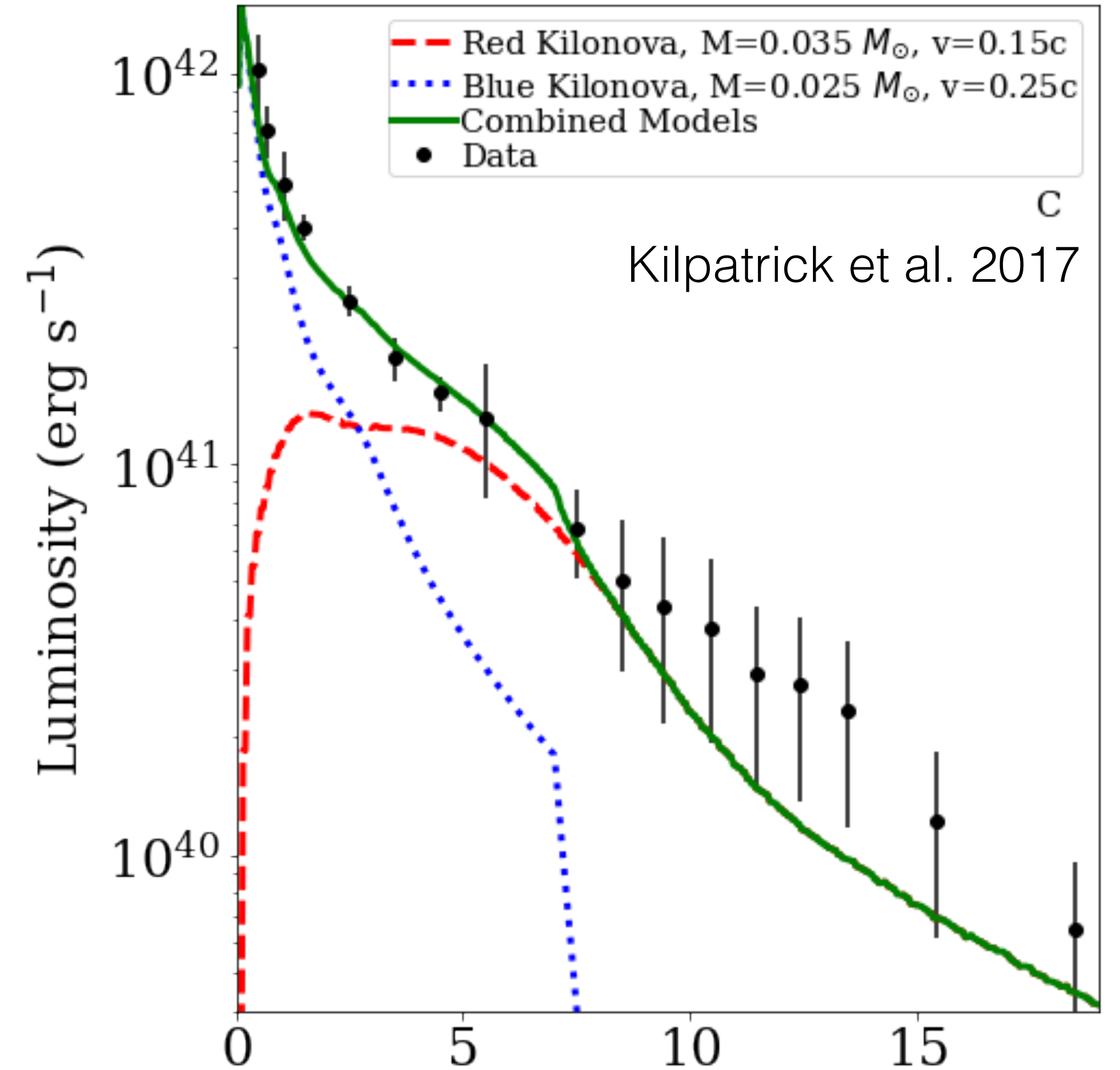
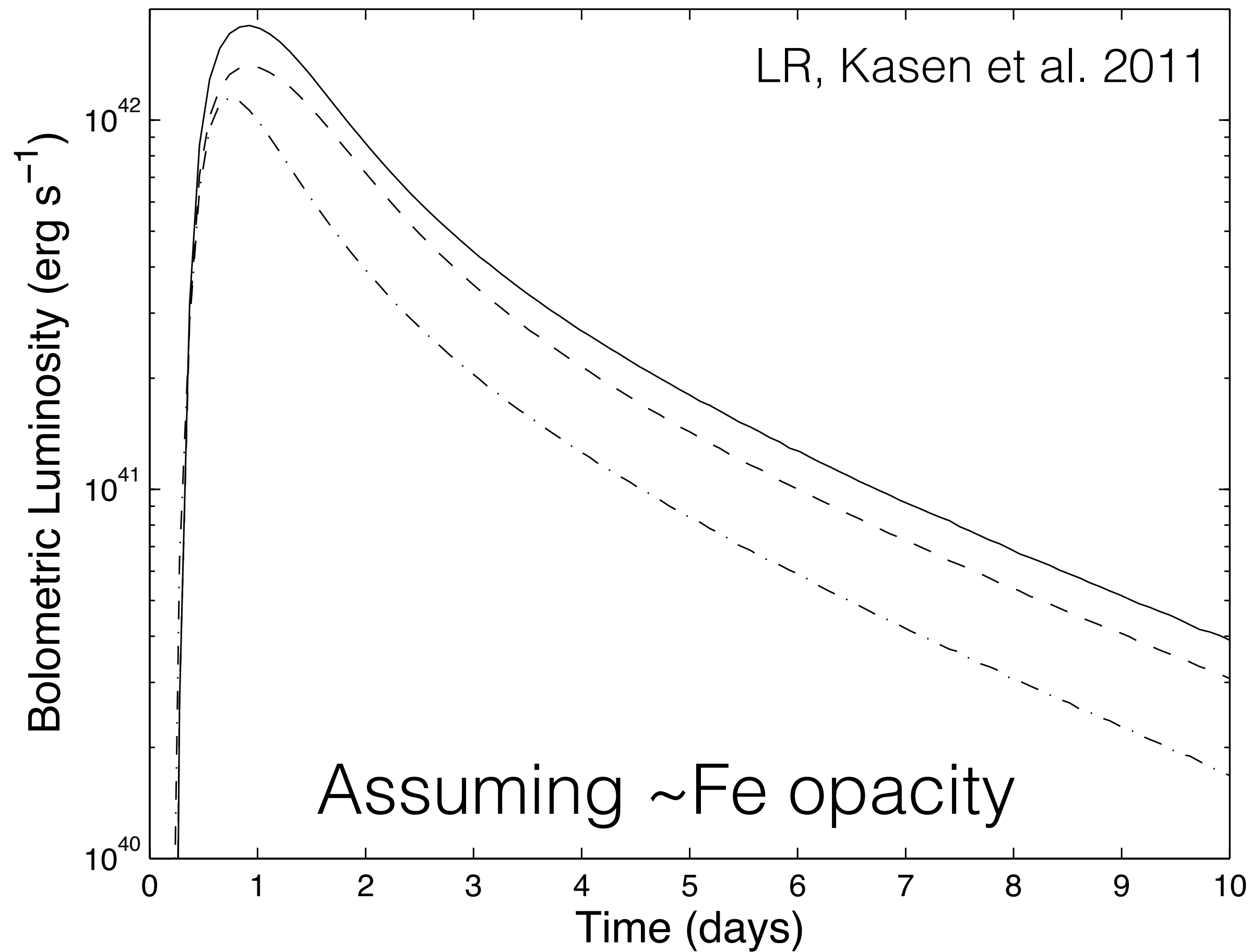
# Nuclear Heating Rate



- Power law heating rate (Metzger et al. '10, Roberts et al. '11, ...)
- Larger number of isotopes involved, sum of numerous individual decays
- Beta-decays, alpha decays and fission

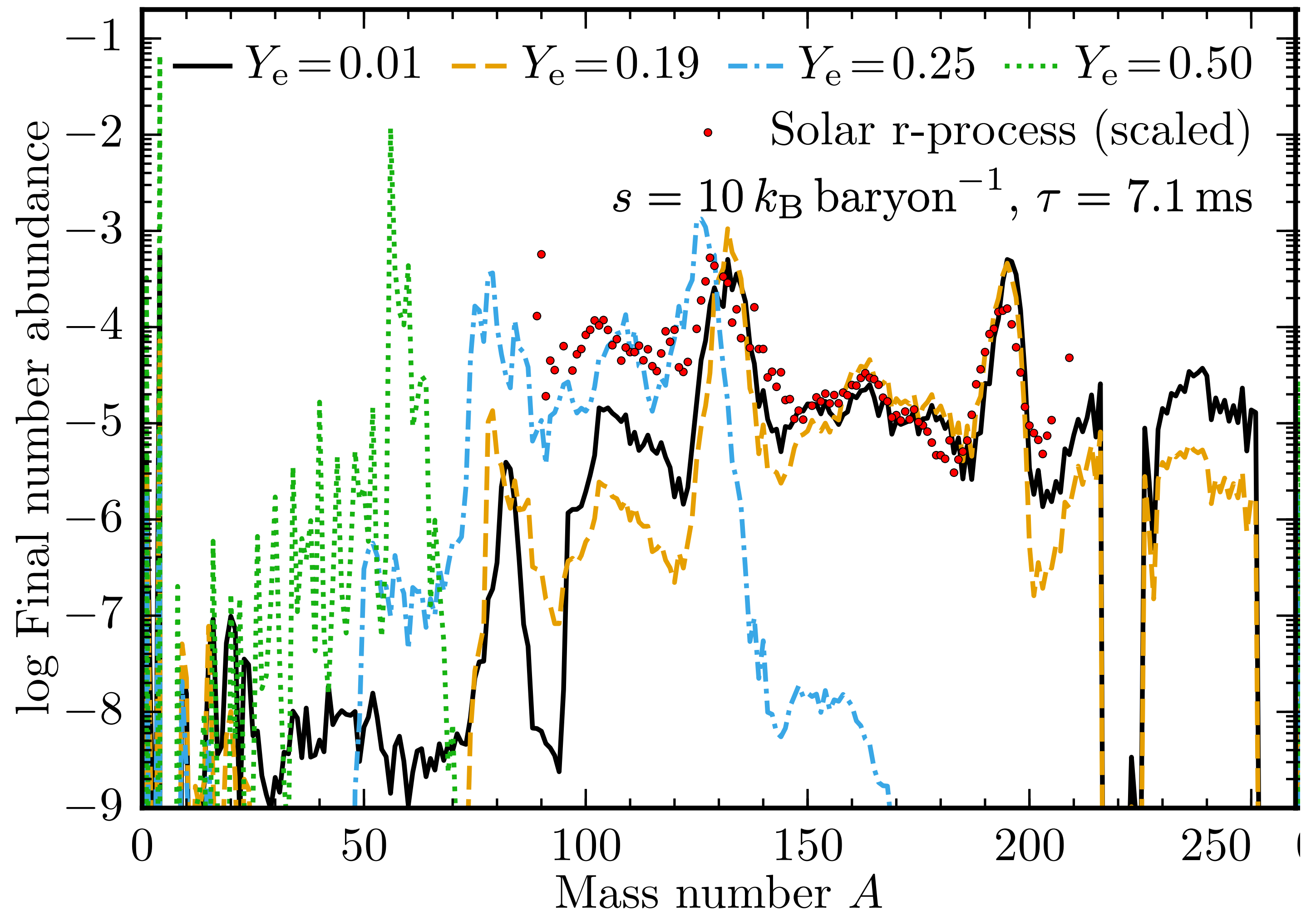


# Electromagnetic displays from nuclear decay

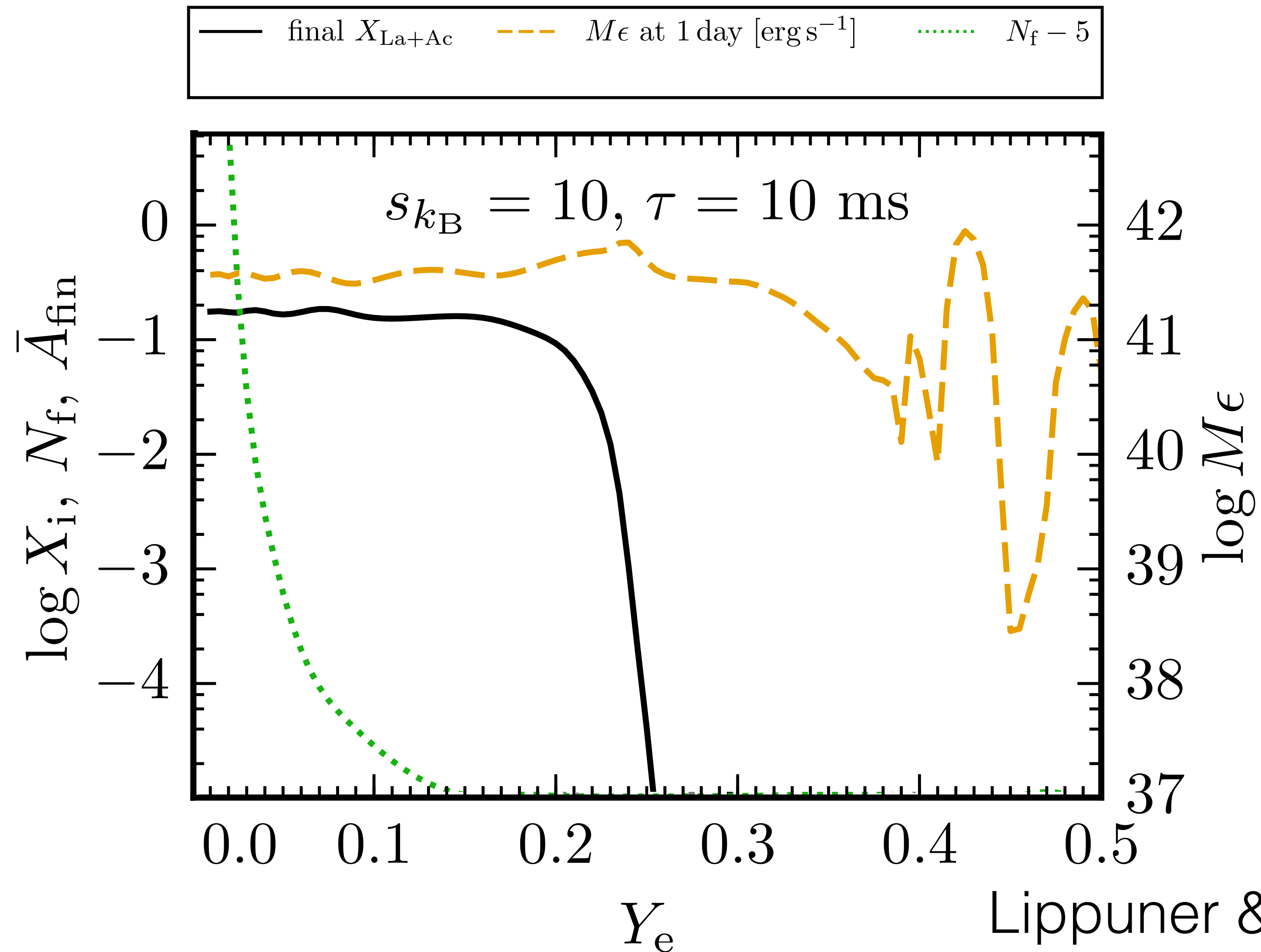


# Dependence of Nucleosynthesis on Initial Conditions

- Changing the electron fraction can substantially alter nucleosynthesis in neutron rich outflows
- Neutron rich nucleosynthesis is most sensitive to  $Y_e$
- How far does nucleosynthesis get before neutron exhaustion?

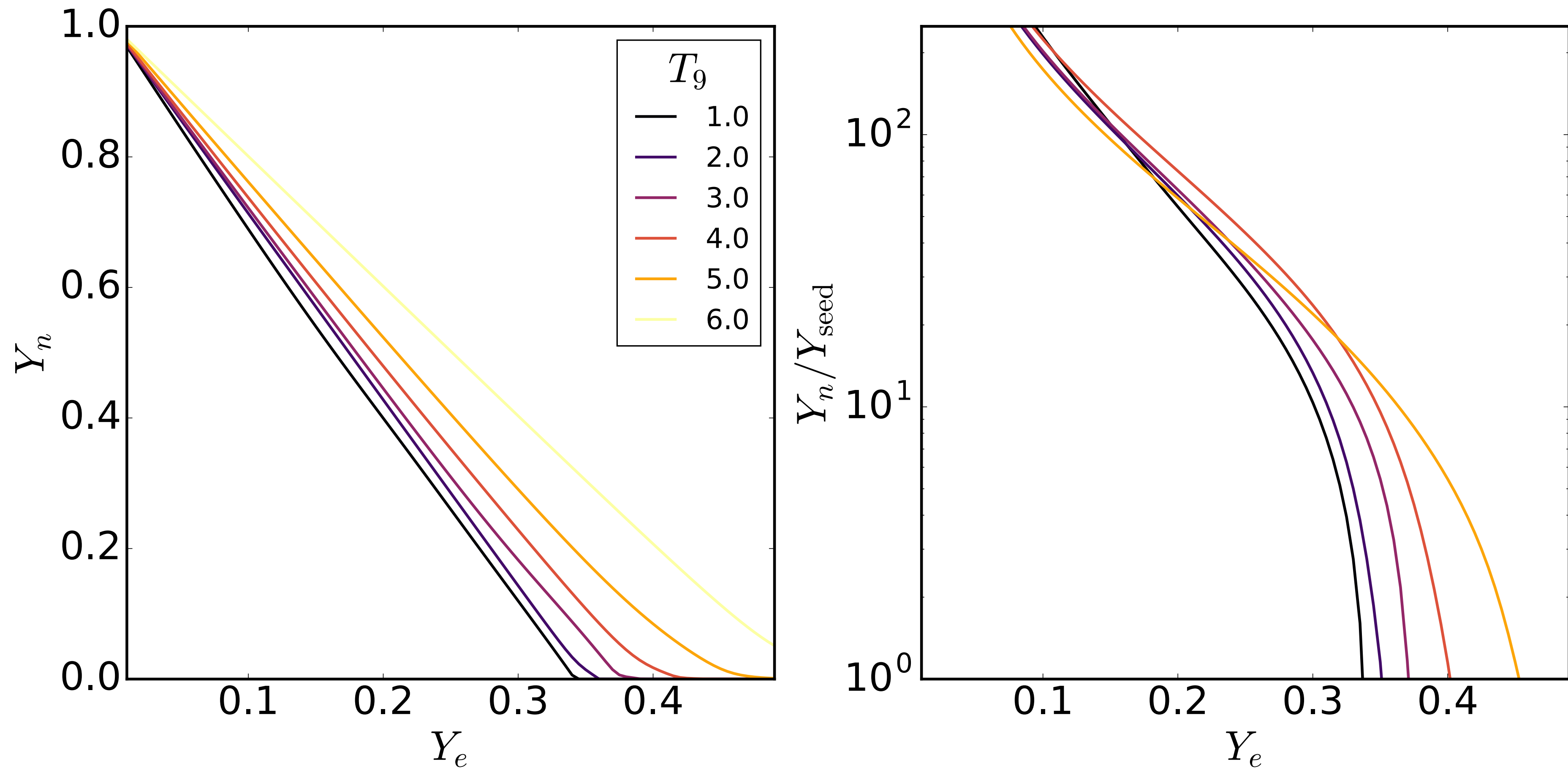


# Dependence of Nucleosynthesis on Initial Conditions





# Neutron-to-Seed Ratio of initial NSE distributions

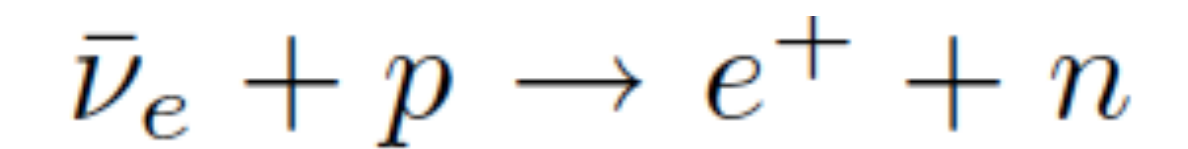
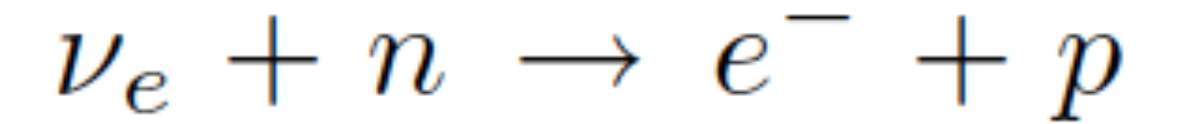


Can trace  $Y_e$  cutoff back to the initial conditions

# Setting $Y_e$ in the Ejecta

Evolution of the electron fraction is governed by

$$\frac{dY_e}{dt} = (\lambda_{\nu_e} + \lambda_{e+})Y_n - (\lambda_{\bar{\nu}_e} + \lambda_{e-})Y_p + \dots$$



$$Y_e(t) \approx Y_{e,0} \exp(-t/\tau_w) + [1 - \exp(-t/\tau_w)]Y_{e,eq}$$

Characteristic Rates:

where

$$\lambda_{e-p} \approx \lambda_{e+n} \approx 0.448 T_{\text{MeV}}^5 \text{ s}^{-1}$$

$$\tau_w = [\lambda_{e-} + \lambda_{e+} + \lambda_{\nu_e} + \lambda_{\bar{\nu}_e}]^{-1}$$

$$\lambda_{\nu_e n} \approx 4.83 L_{\nu_e, 51} \left( \epsilon_{\nu_e, \text{MeV}} + 2\Delta_{\text{MeV}} + 1.2 \frac{\Delta_{\text{MeV}}^2}{\epsilon_{\nu_e, \text{MeV}}} \right) r_6^{-2} \text{ s}^{-1}$$

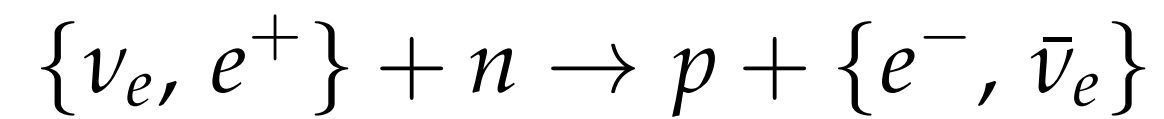
$$\lambda_{\bar{\nu}_e p} \approx 4.83 L_{\bar{\nu}_e, 51} \left( \epsilon_{\bar{\nu}_e, \text{MeV}} - 2\Delta_{\text{MeV}} + 1.2 \frac{\Delta_{\text{MeV}}^2}{\epsilon_{\bar{\nu}_e, \text{MeV}}} \right) r_6^{-2} \text{ s}^{-1}$$

$$Y_{e,eq} = \frac{\lambda_{\nu_e} + \lambda_{e+}}{\lambda_{\nu_e} + \lambda_{e+} + \lambda_{\bar{\nu}_e} + \lambda_{e-}}$$

# Weak Interactions in NS Mergers

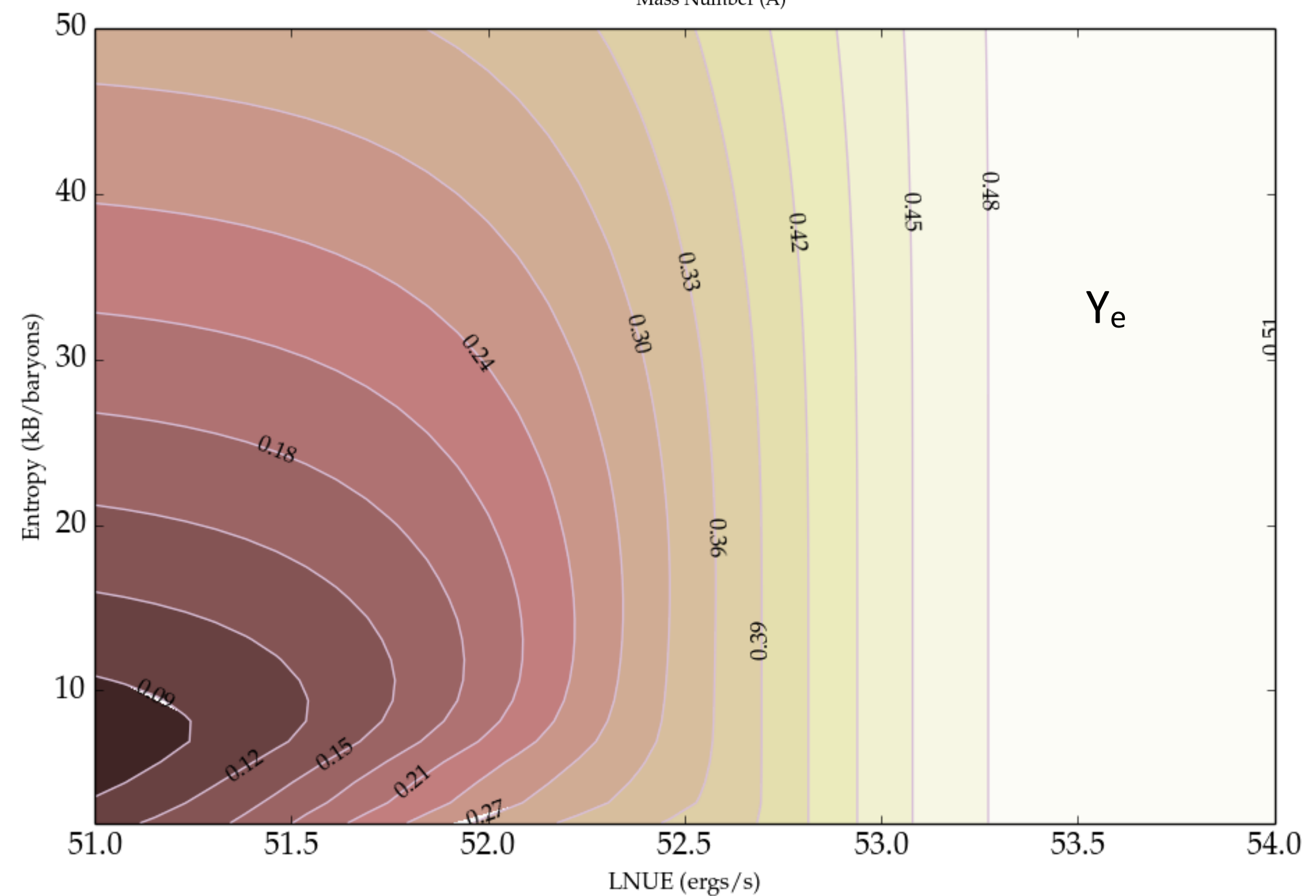
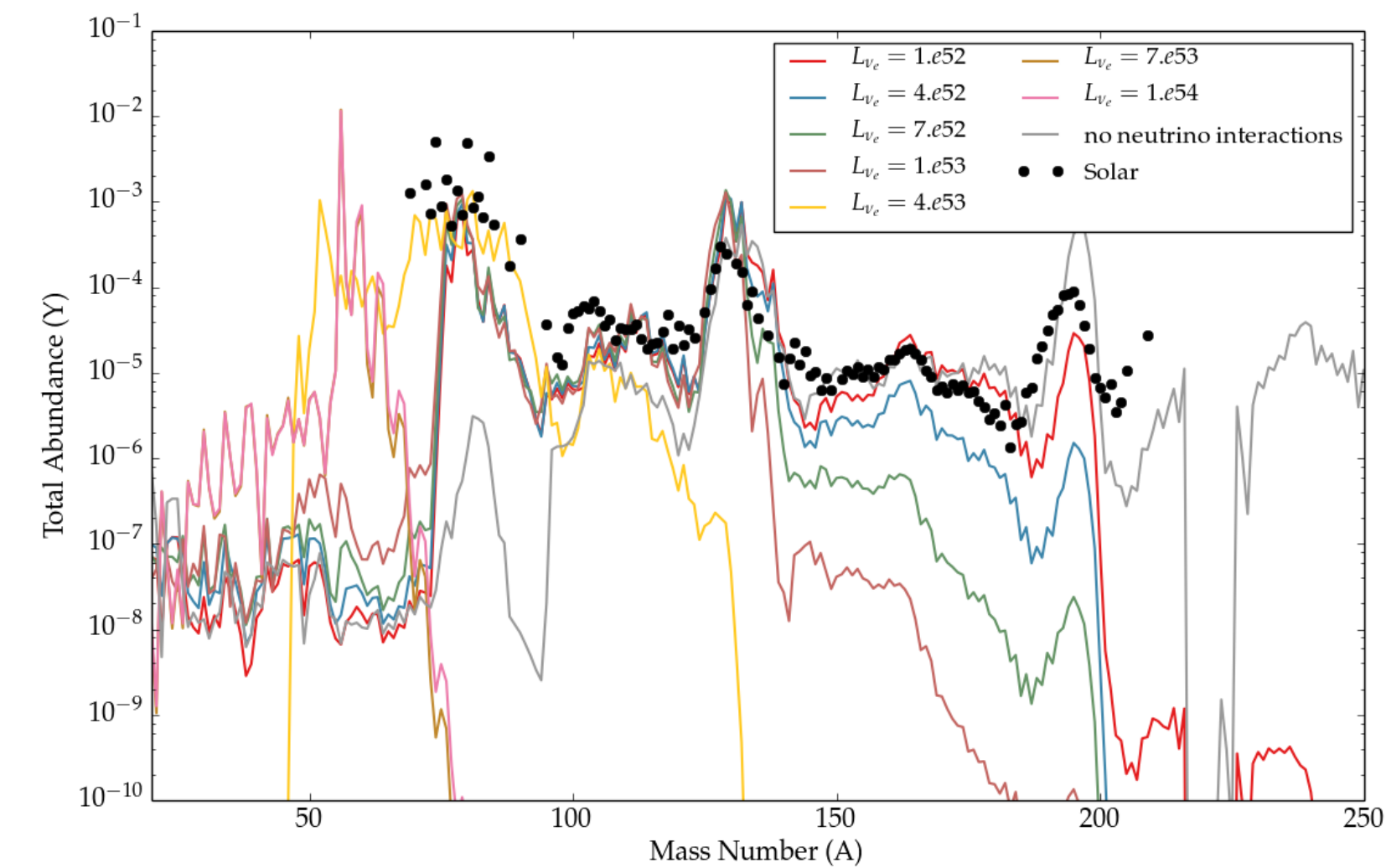
See Wanajo et al. (2014) and  
Goriely et al. (2015)

Destroy neutron at early times in hot,  
neutrino rich environment  
at early times via:



NSE favors more seed nuclei, fewer neutrons,  
thereby gives lower neutron to seed ratio

Incomplete r-process, material builds up at first  
peak



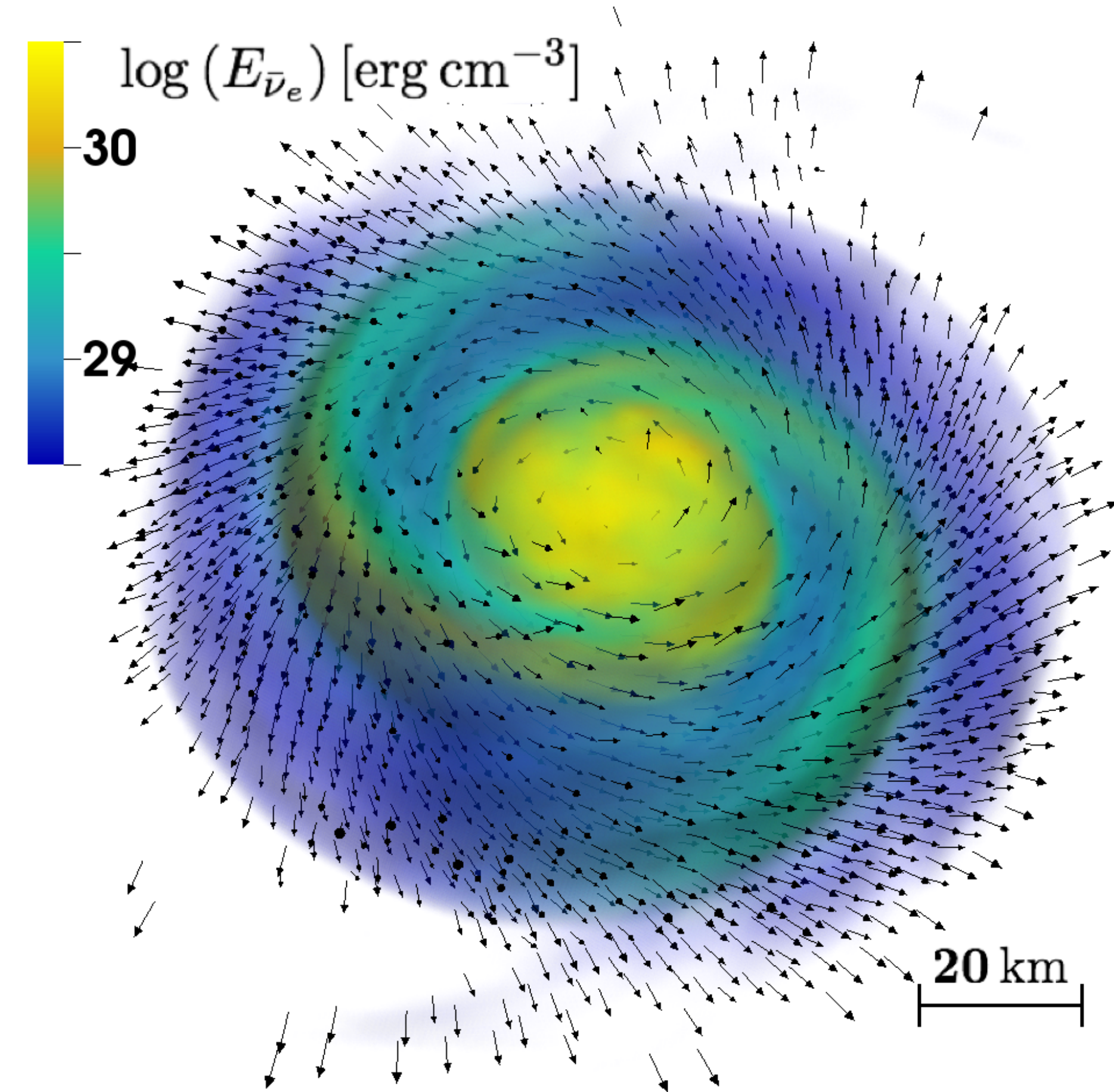
From summer student Sandra Ning



# Neutrino Transport in the Ejecta

see Wanajo, et al. '14, Radice et al. 16, Palenzuela et al. 16

- Large neutrino luminosities provided by central remnant of the NS merger
- Hierarchy of neutrino energies similar to proto-NS neutrino emission because neutrino decoupling physics is similar

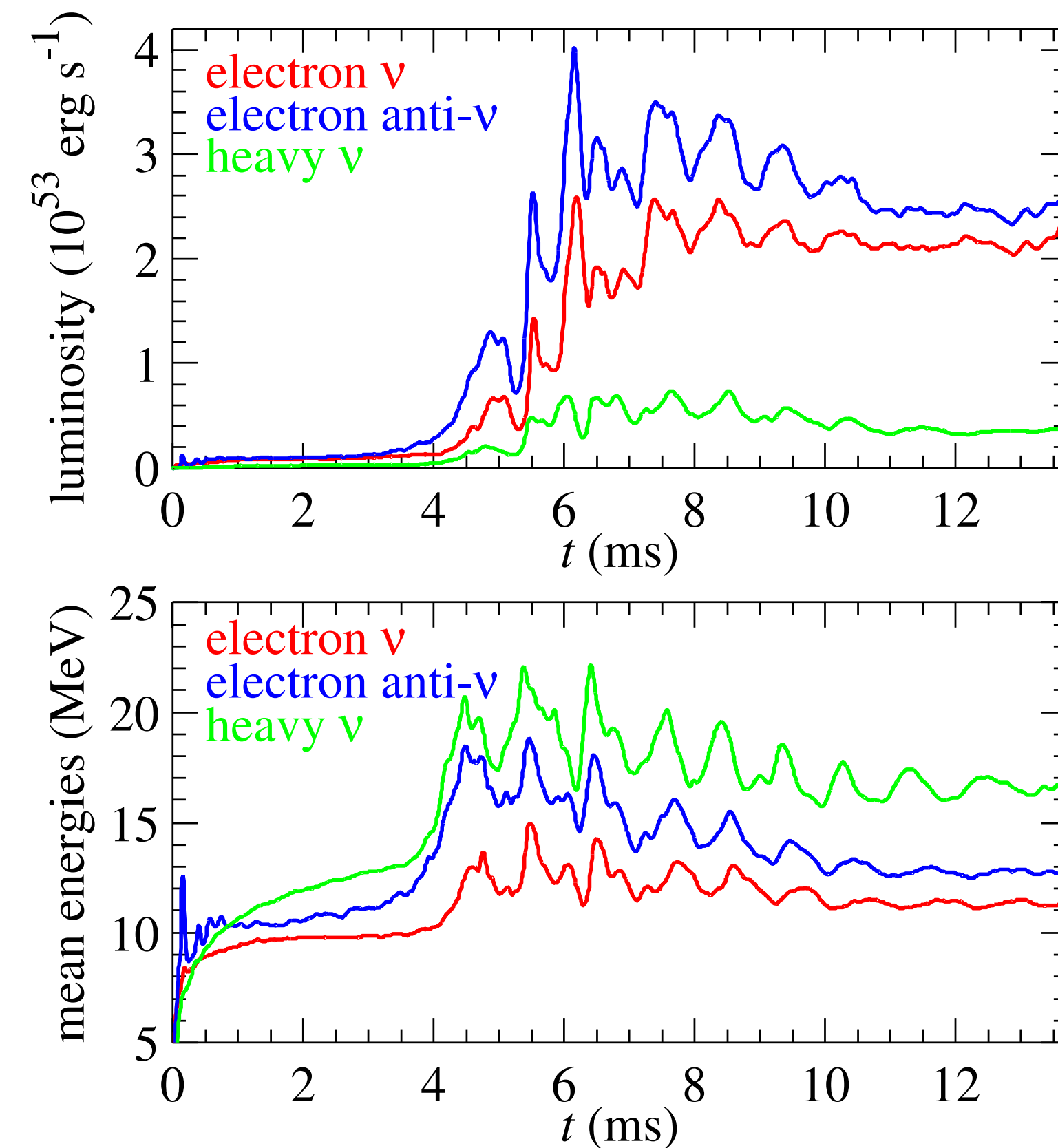


Foucart, O'Connor, LR et al. '16

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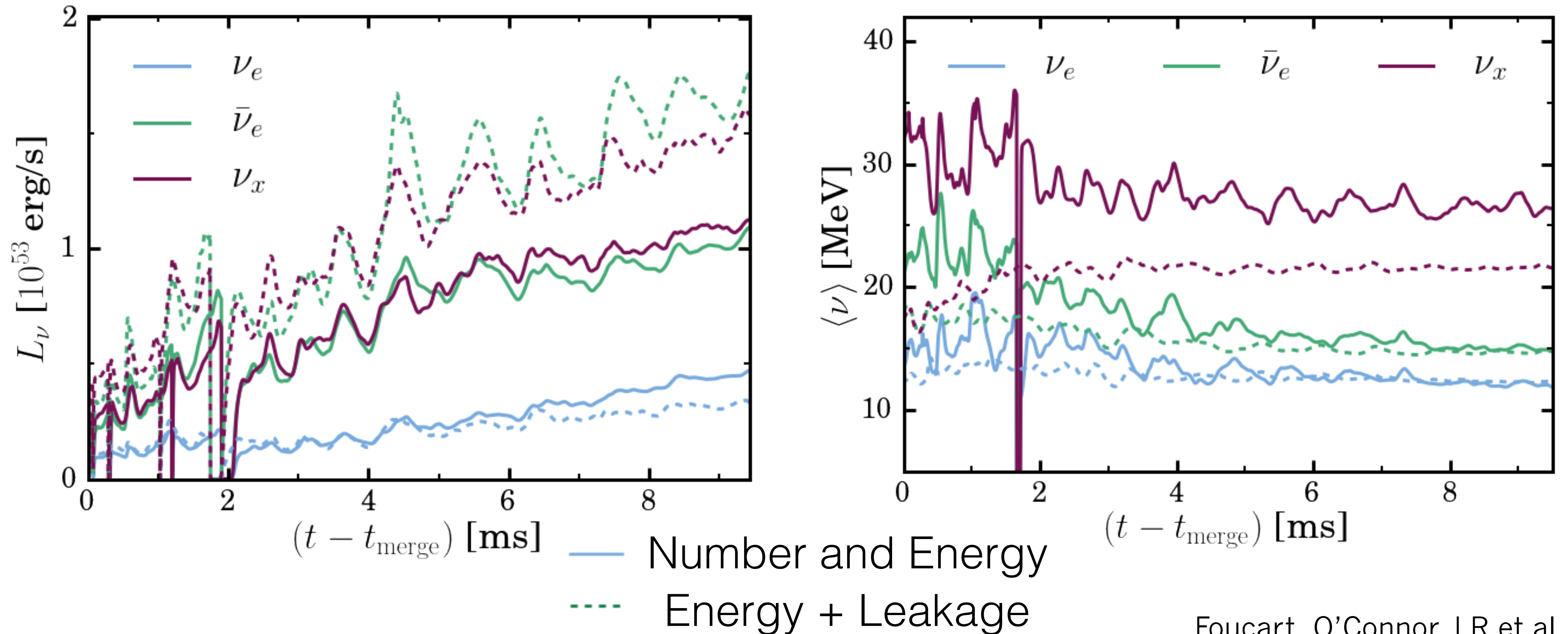


from Wanajo (2014)



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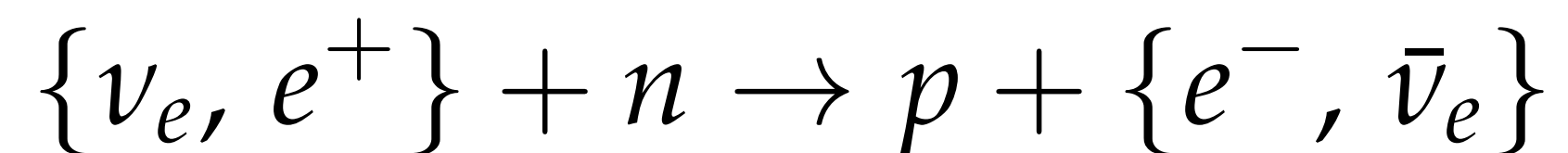
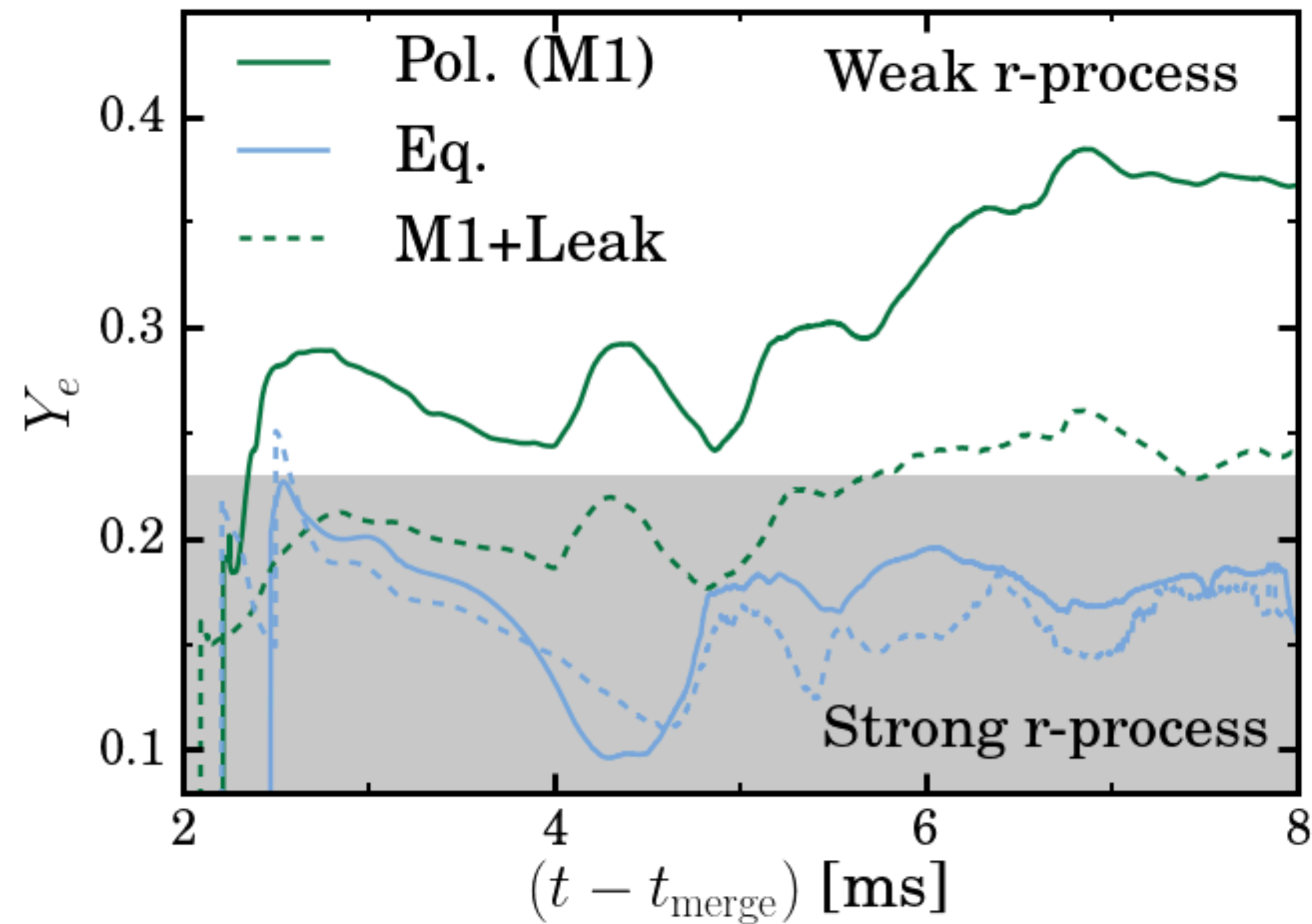
see Wanajo, et al. '14, Radice et al. 16, Palenzuela et al. 16





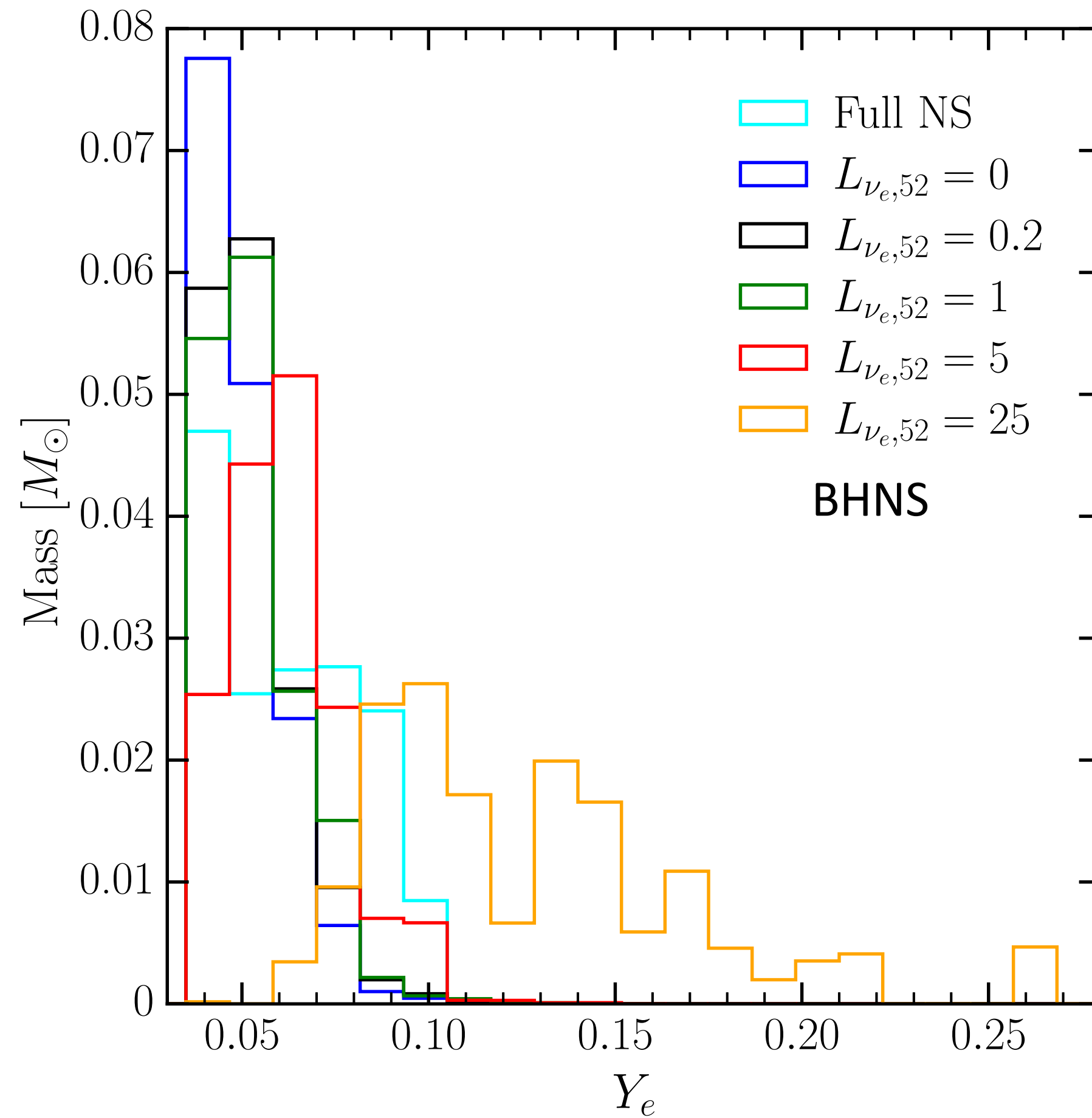
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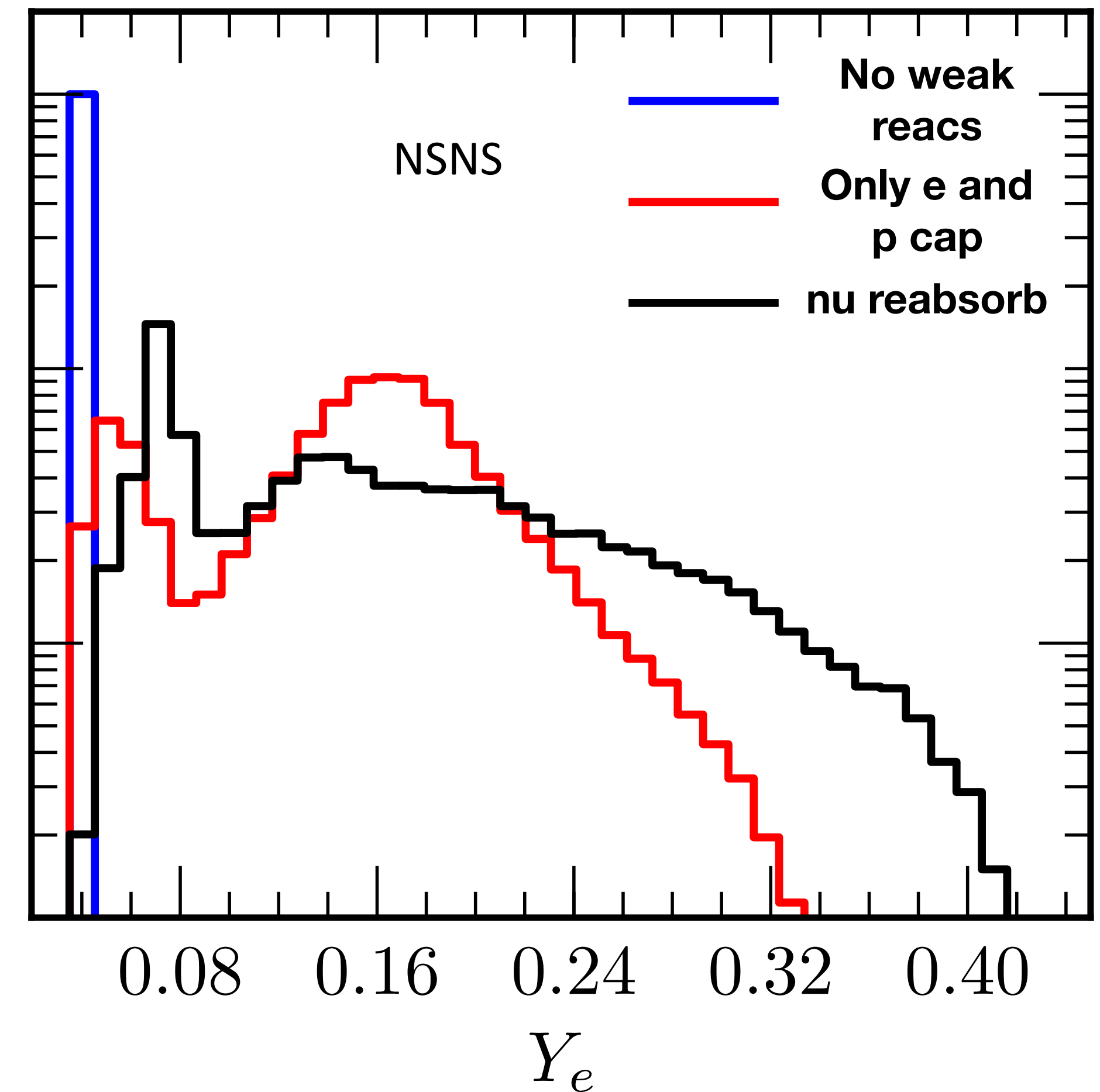


Foucart, O'Connor, LR et al. '16

# Dynamical Ejecta in BHNS mergers vs NSNS mergers

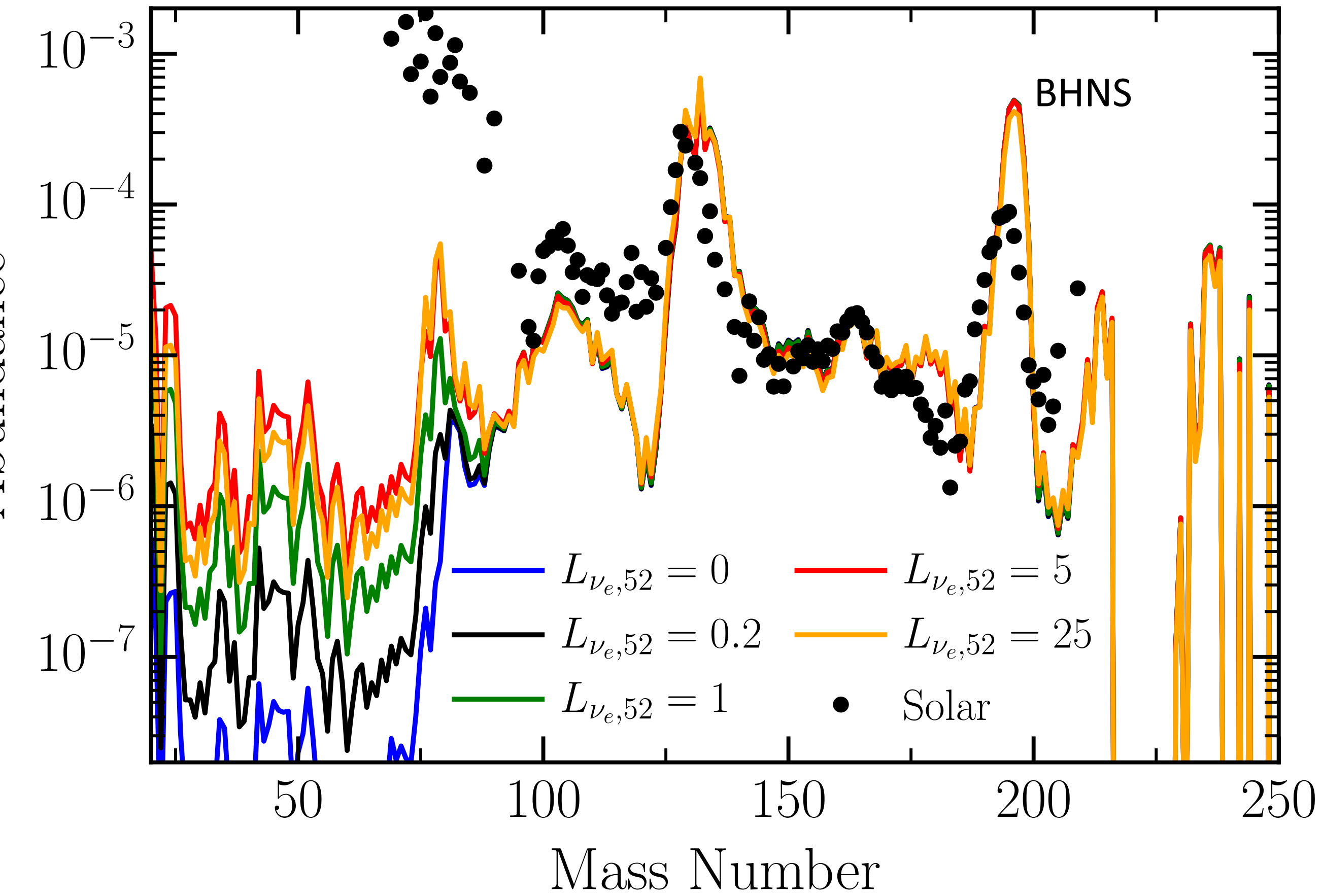


LR, et al. '16

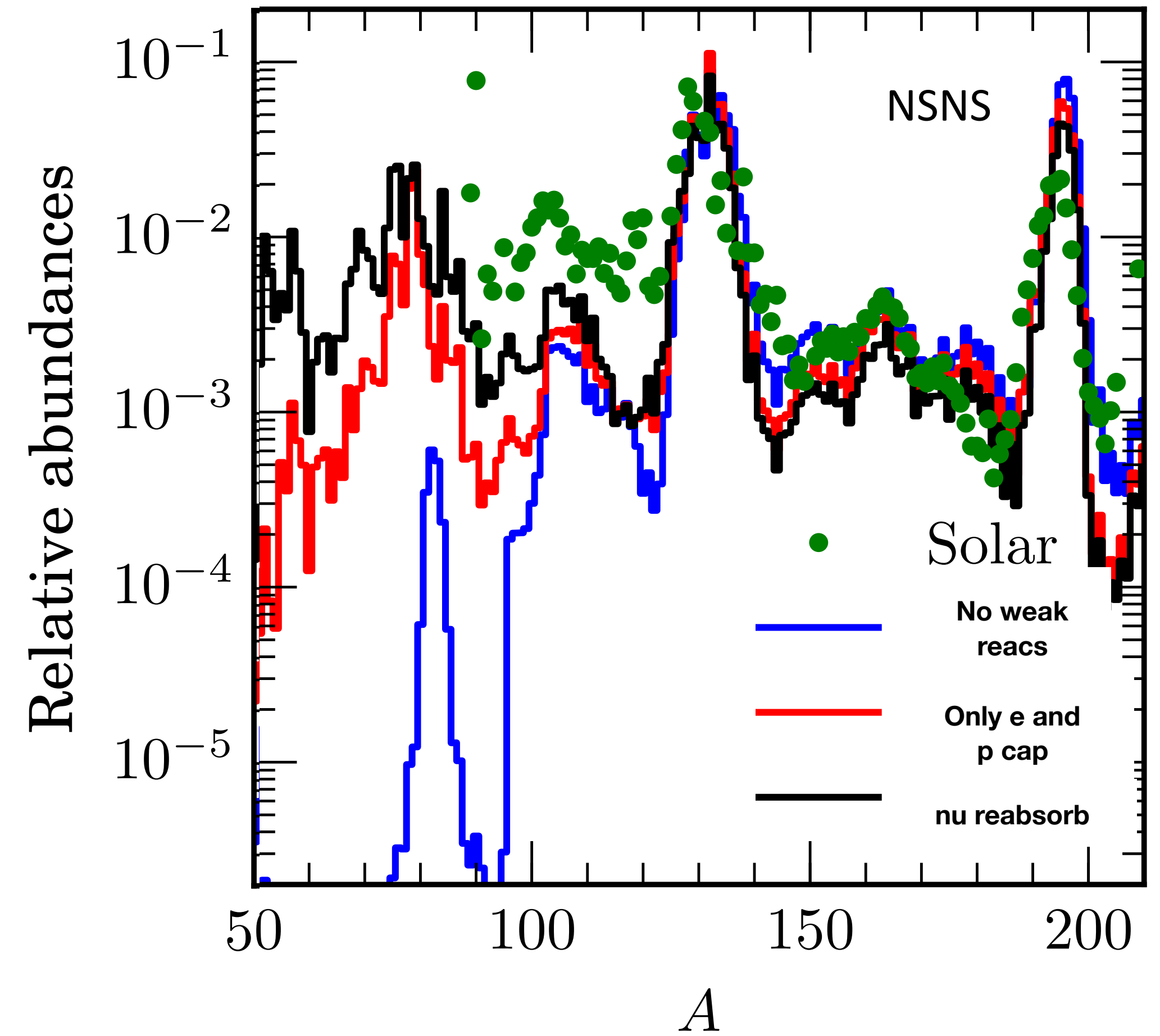


Radice, ..., LR et al. '16

# Ejecta Composition



LR, et al. '16



Radice, ..., LR et al. '16



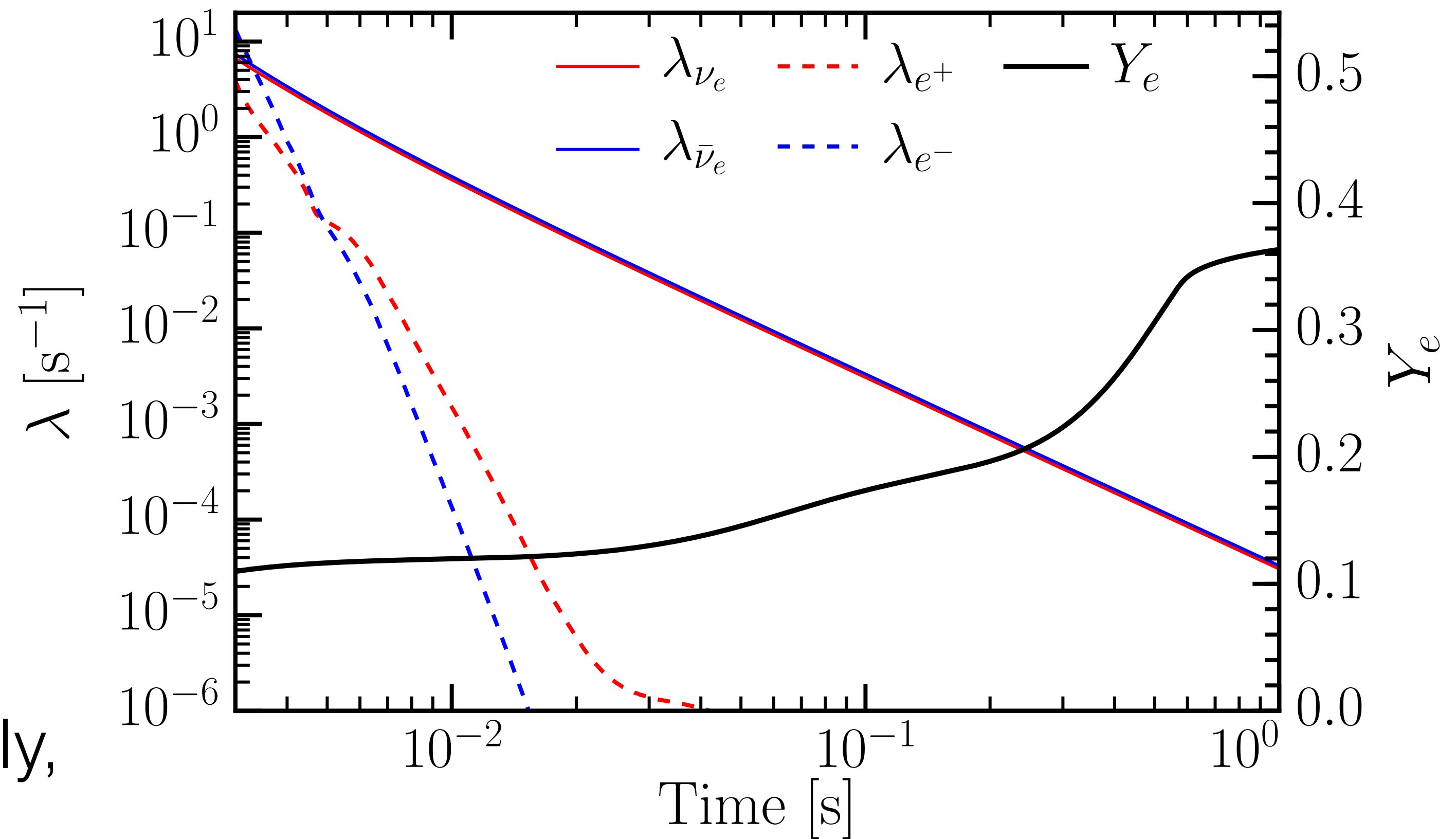
# Weak interactions in BHNS dynamical ejecta

Low entropy tidal ejecta -> small  
electron/positron capture rates

Neutrino reactions are  
somewhat faster

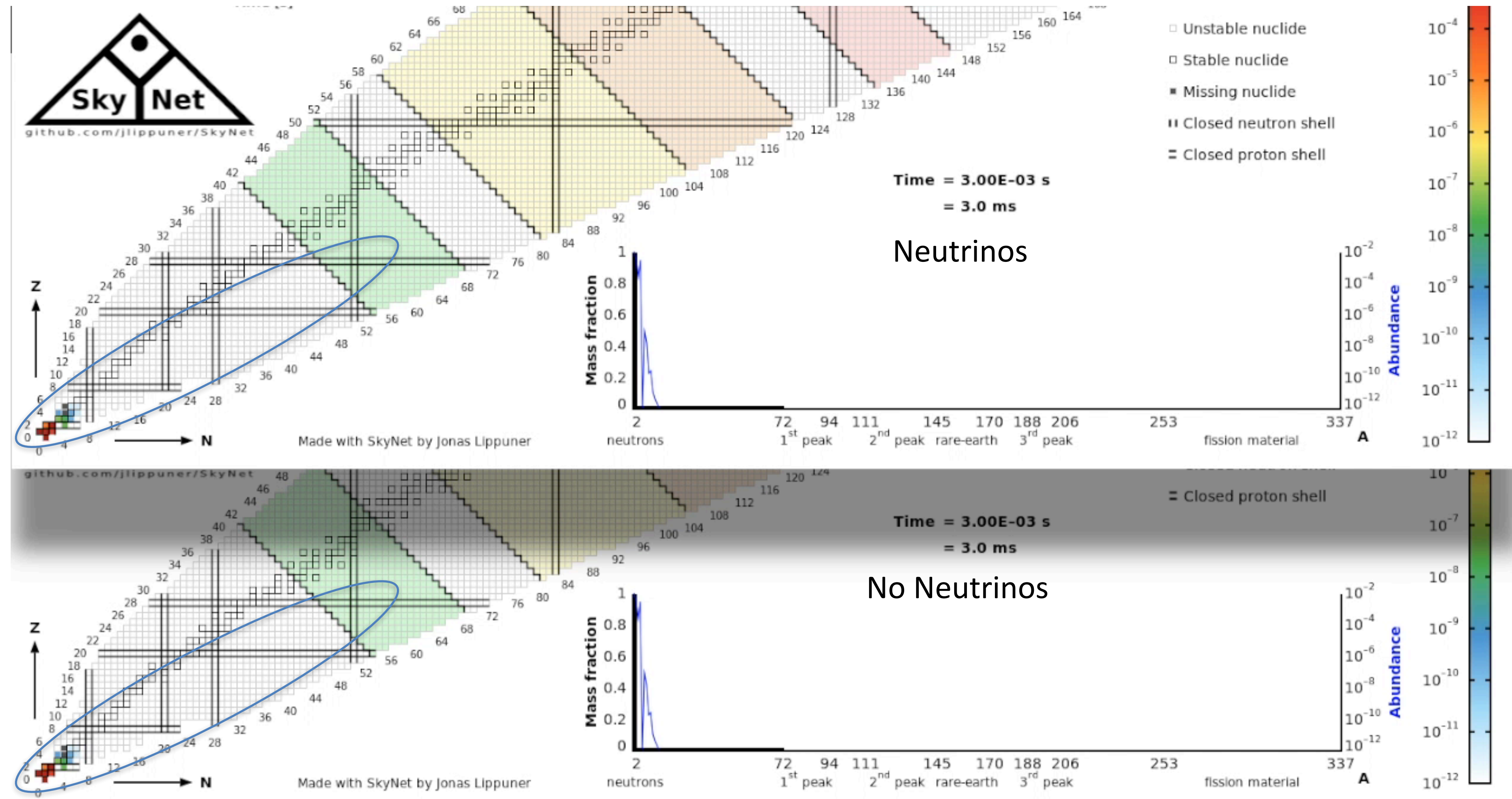
$$\tau_\nu(r) \approx 67.8 \text{ ms} \left( \frac{r}{250 \text{ km}} \right)^2 L_{\nu_e,53}^{-1} T_{\nu_e,5}^{-1}$$

Still too slow to impact  $Y_e$  significantly,  
but can impact the first peak  
nucleosynthesis in the dynamical ejecta



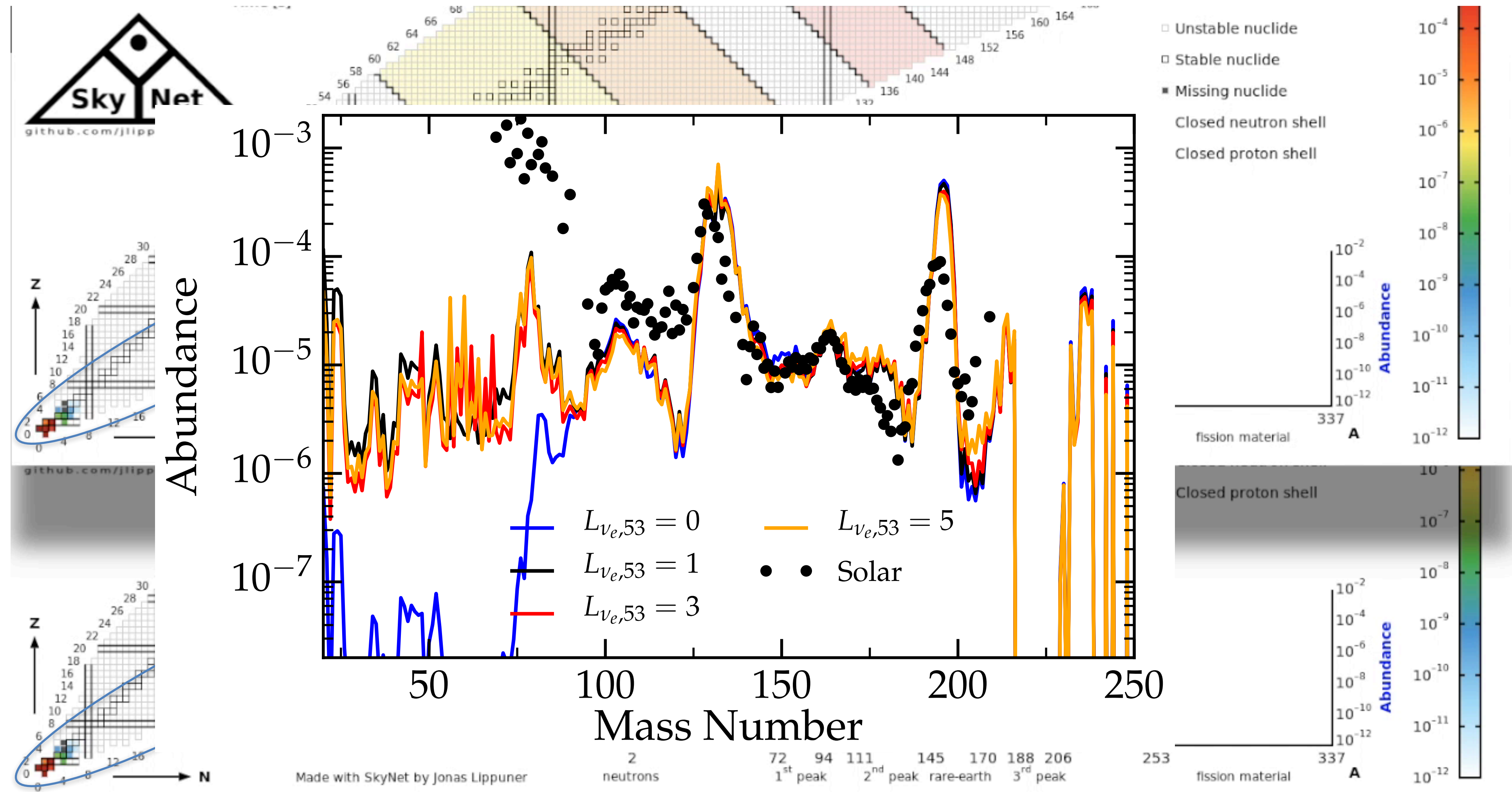
LR, et al. '16

# First r-Process Peak Production in BHNS Mergers





# First r-Process Peak Production Method 2

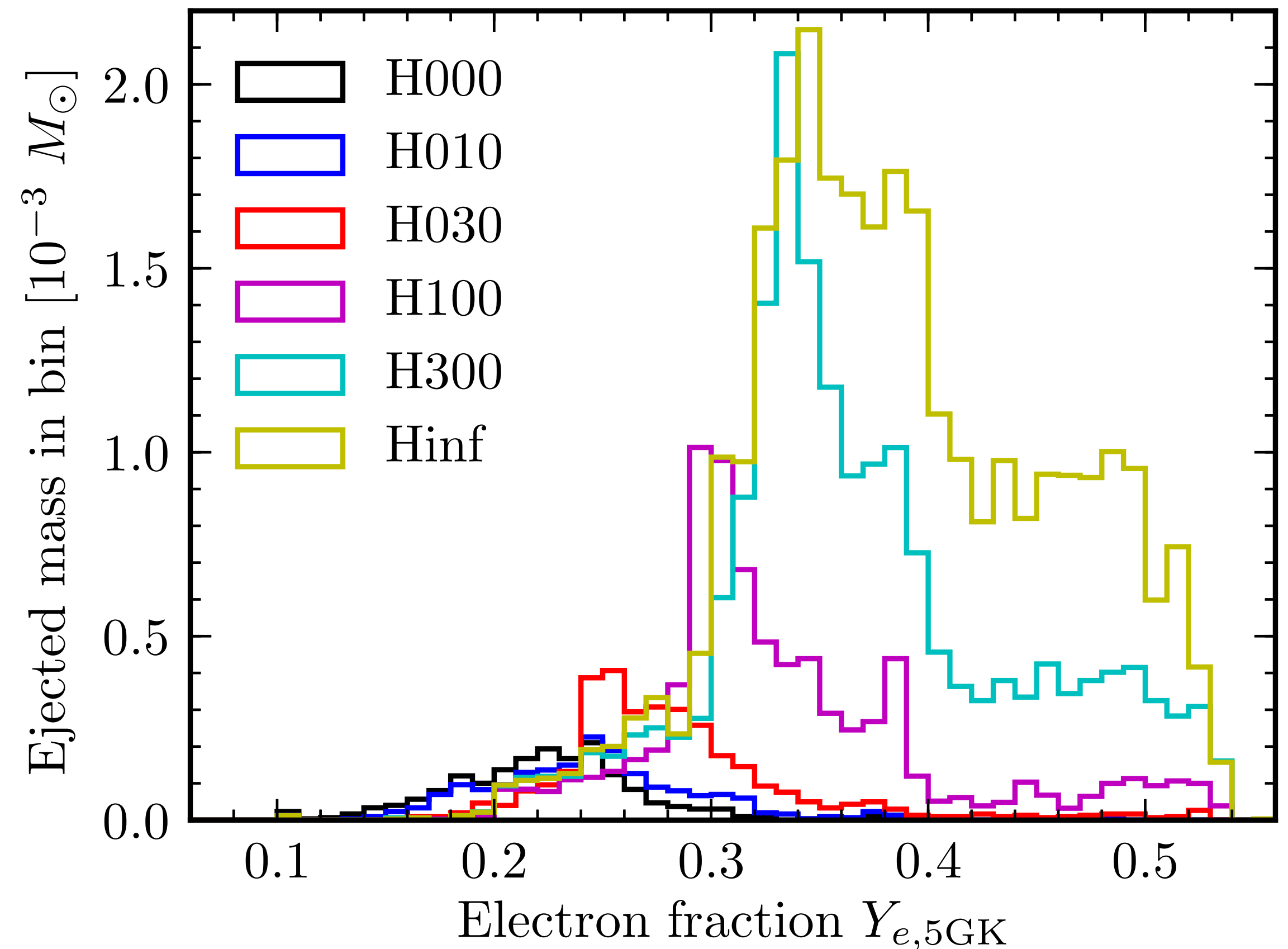




# Disk ejecta

see e.g. Metzger & Fernandez 14, Just et al. 15, Siegel & Metzger 2018

- Material in the remnant disk also experiences a large number of weak interactions, beta-equilibrates
- Broad range of  $Y_e$ , depending on the lifetime of the hyper-massive neutron star
- Ratio of weak to strong r-process sensitive to the lifetime of the central object

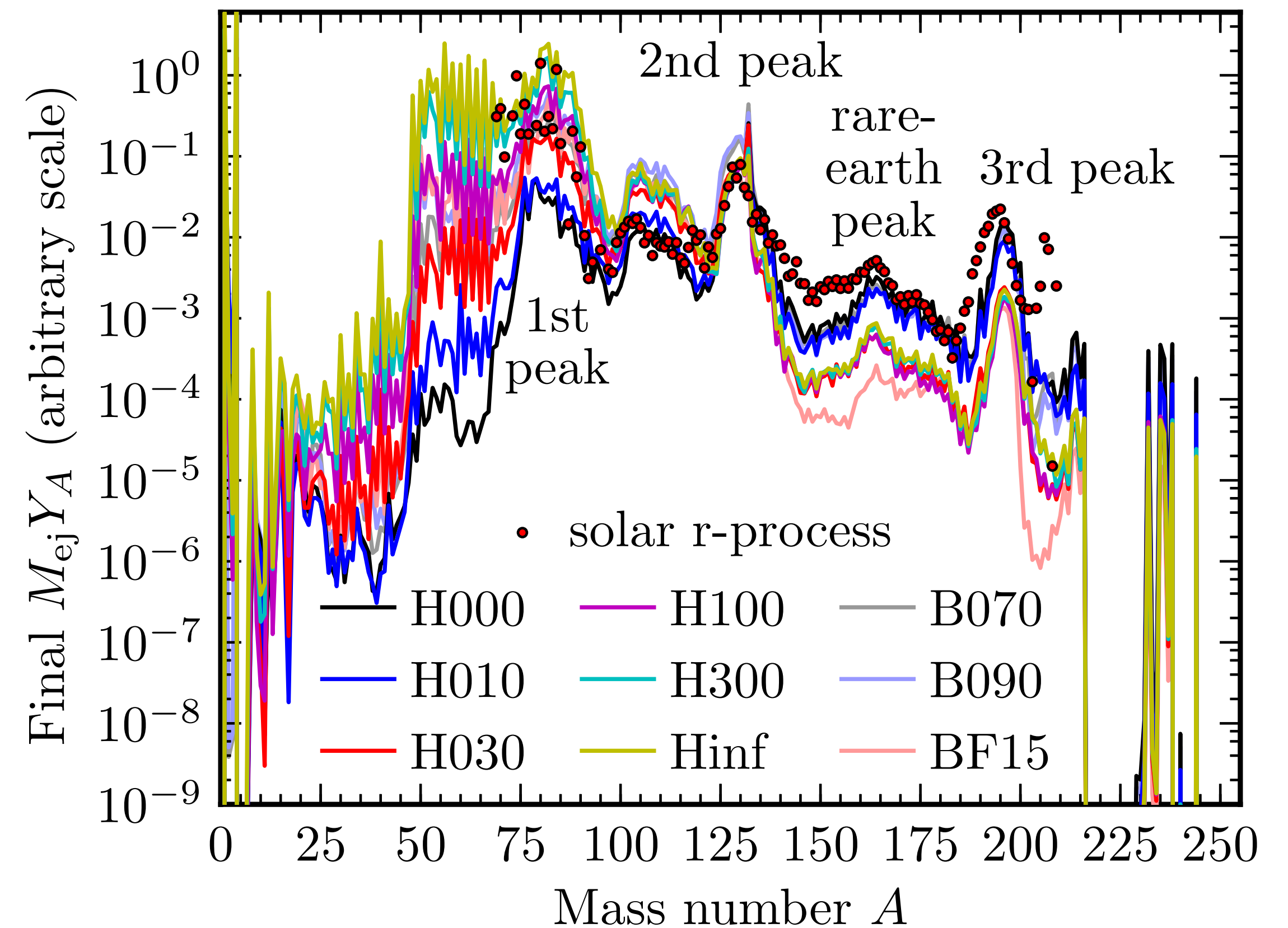


from Lippuner, Fernandez, LR, et al. (2017)

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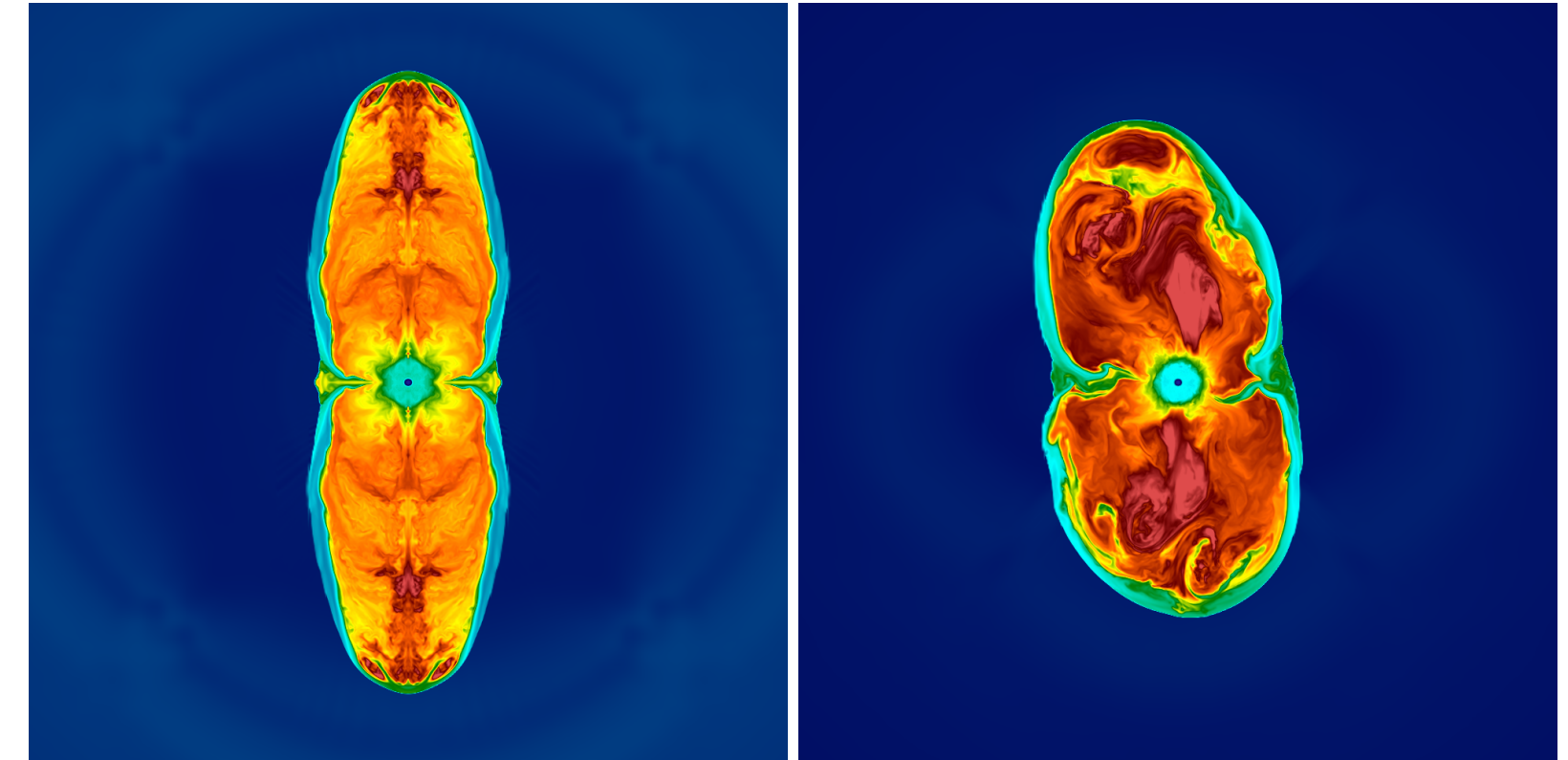
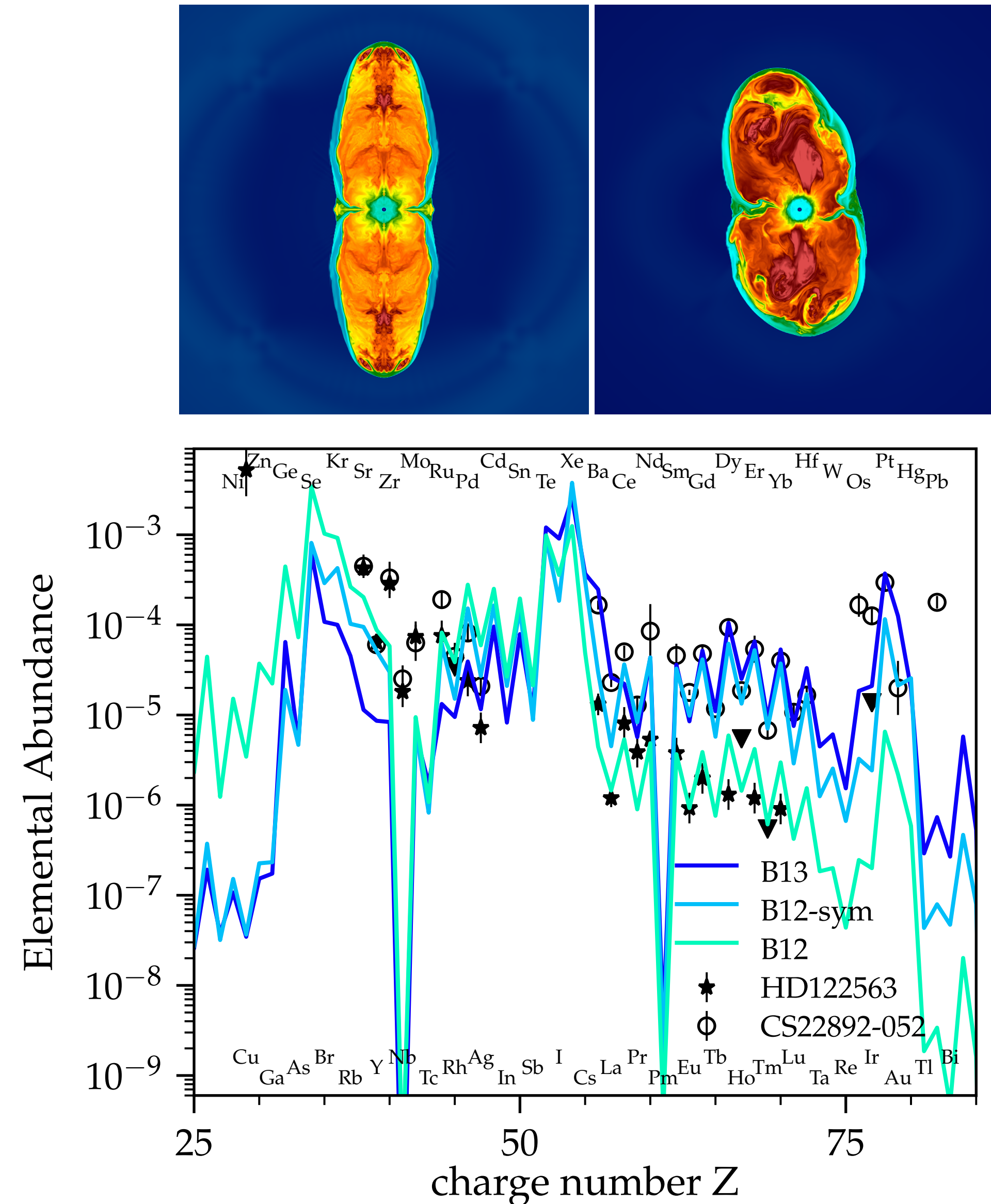
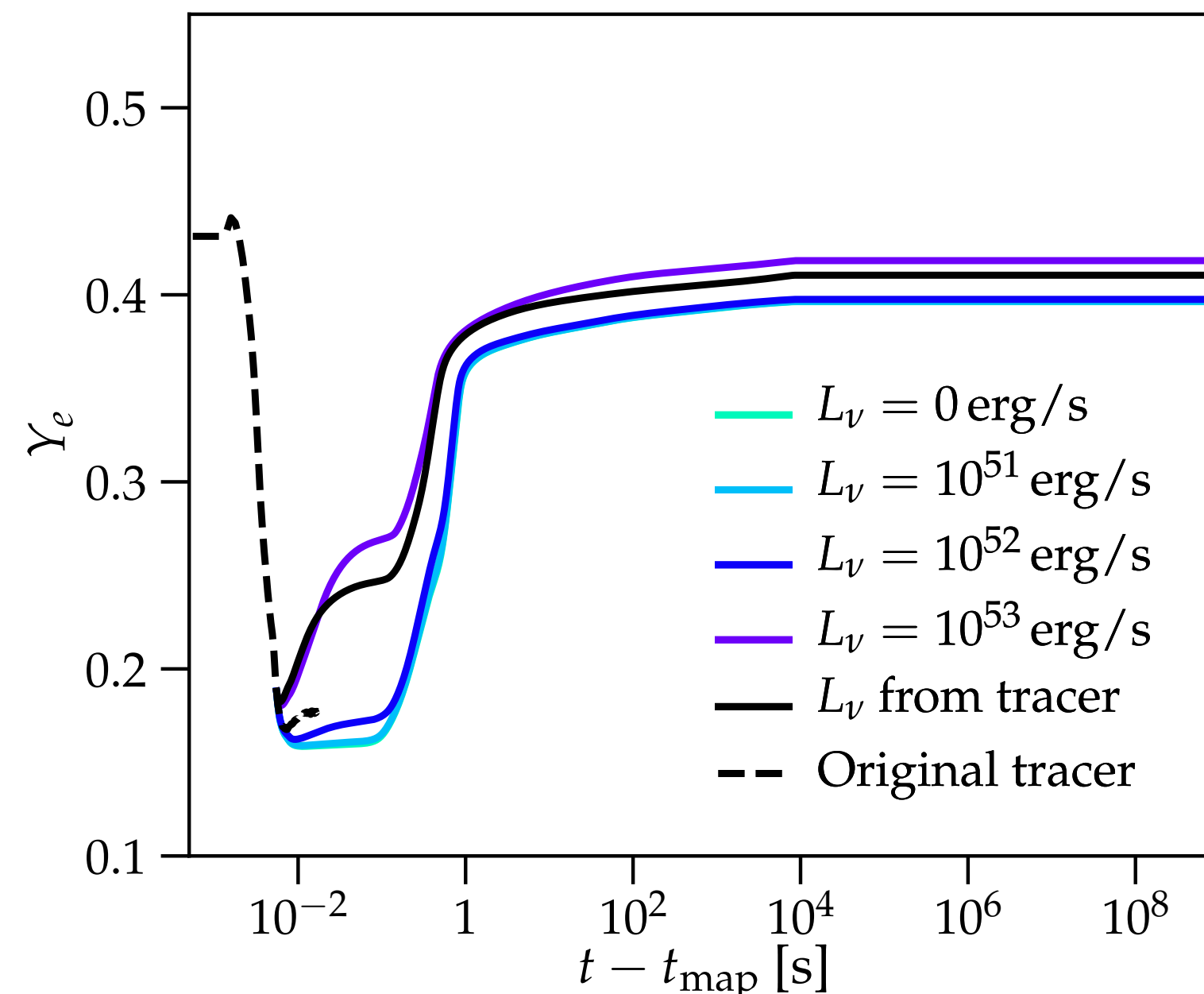
from Lippuner, Fernandez, LR, et al. (2017)

# Jet Driven Supernovae

see Winteler et al. 2012, Nishimura et al. 2015, Moesta et al. 2018

Moesta, LR, et al. (2018)

- Rapidly rotating, magnetized SNe
- Full 3D Dynamics also important here
- Kink instabilities in jet significantly change dynamics and impact nucleosynthesis





# Summary and Outlook

- $Y_e$  distribution of the ejecta determining factor in the final composition and properties of the transient
- Weak interactions play a substantial role in setting the initial conditions for nucleosynthesis
- Going forward need better treatment of neutrinos during the dynamical phase -> important to setting the electron fraction distribution via weak interactions
- Sensitivity of r-process nucleosynthesis to input nuclear data of nuclear reaction network calculations. How well is the lanthanide cutoff  $Y_e$  known?
- Still some possible SN sites of the r-process
- Hopefully observe a BHNS merger