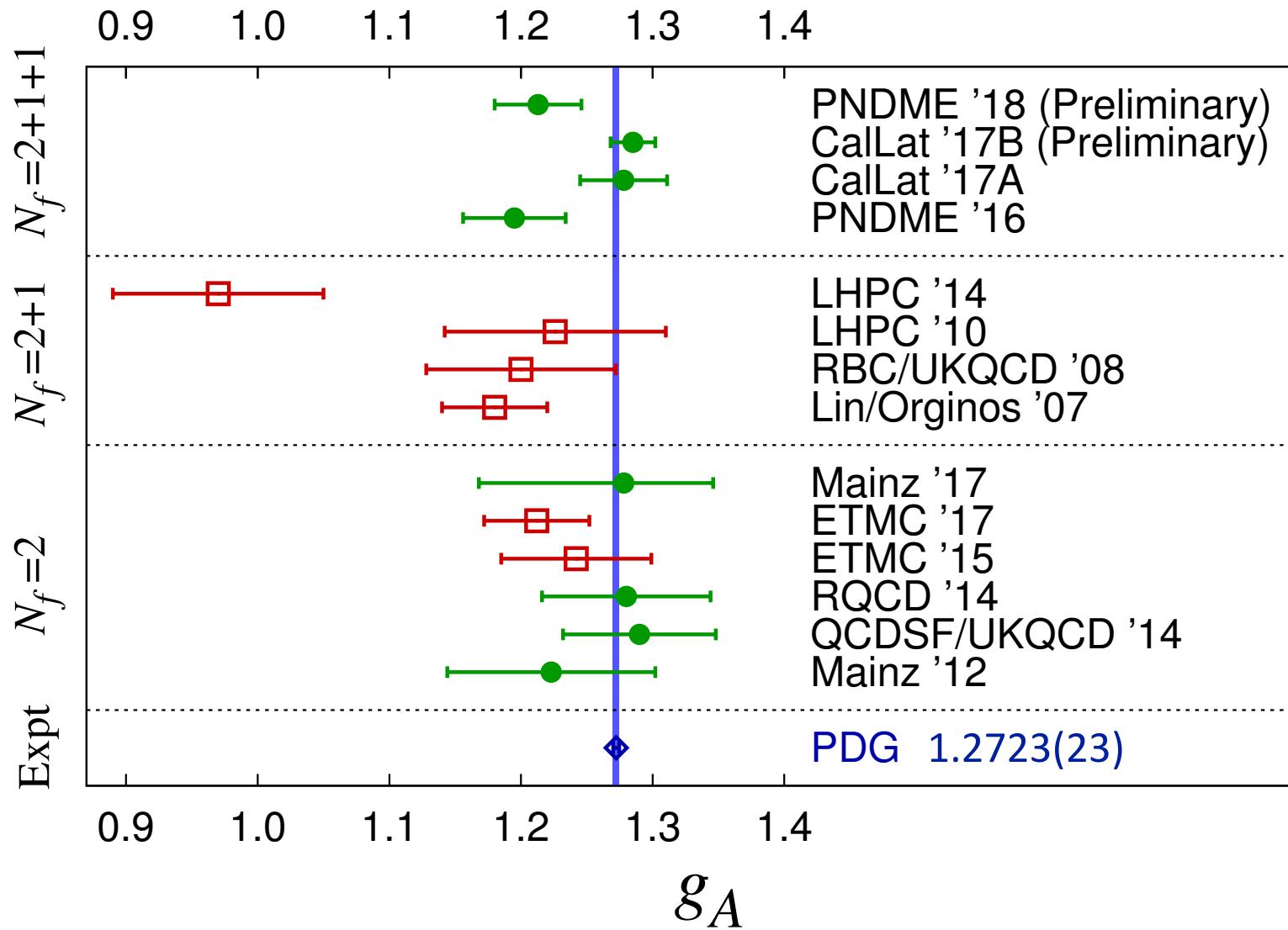


$g_A$

Boram Yoon  
Los Alamos National Laboratory

USQCD AHM, Fermilab, April 20-21, 2018

# Current Status - $g_A$



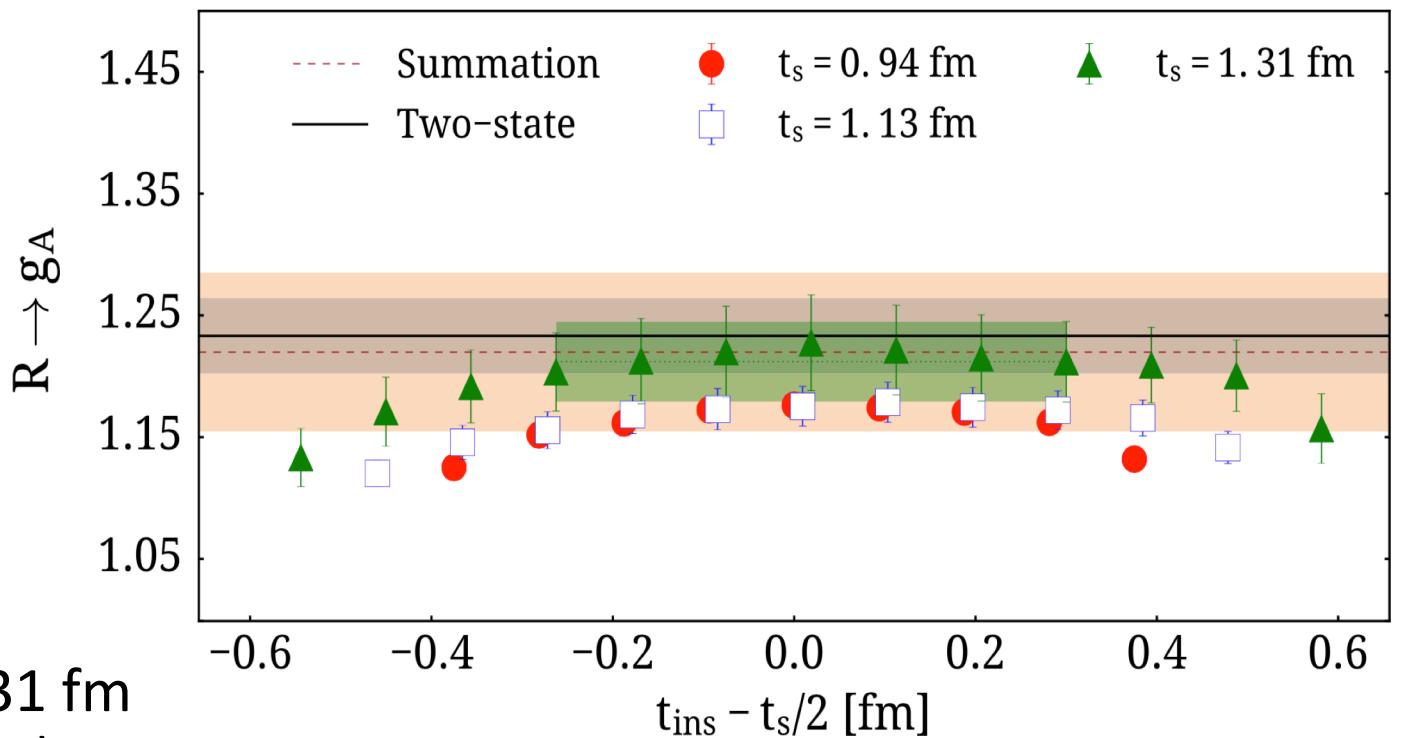
# ETMC '17 (PRD 96, 054507, 2017)

- **Lattice**

- TMF,  $N_f=2$
- $a = 0.094$  fm
- $m_\pi = 130$  MeV
- $m_\pi L = 3.0$

- **Measurement/Analysis**

- APE-smeared links
- Fixed-sink method
- *Plateau method* at  $t_{sep} = 1.31$  fm  
(compared with Two-state fit)
- NPR
- Single lattice spacing at physical pion mass



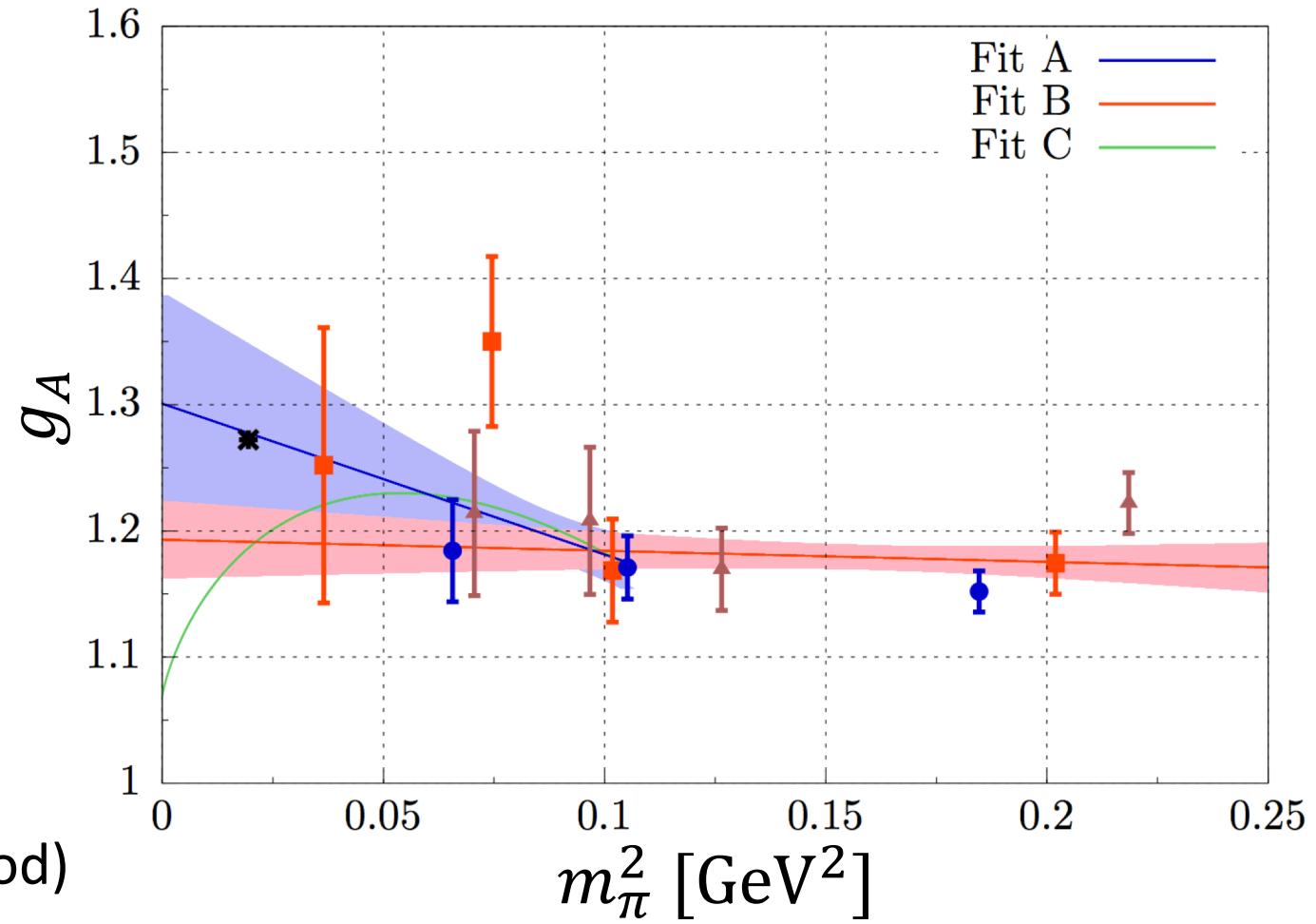
# Mainz '17 (arXiv:1705.06186)

- **Lattices**

- CLS,  $N_f=2$
- $a = 0.050, 0.063$  and  $0.079$  fm
- $m_\pi = 190 - 473$  MeV,  $m_\pi L \gtrsim 4$

- **Measurement/Analysis**

- APE-smeared links
- Fixed-sink method
- $O(a)$ -improved axial current
- *Two-state fit* with assumption of the  $N\pi\pi$  dominant excited-state (compared with summation method)
- NPR
- No  $a$ -dependence, central value from linear fit in  $m_\pi^2$  for  $m_\pi < 330$  MeV



# CallLat '17 (EPJ WoC 175, 01008, 2018)

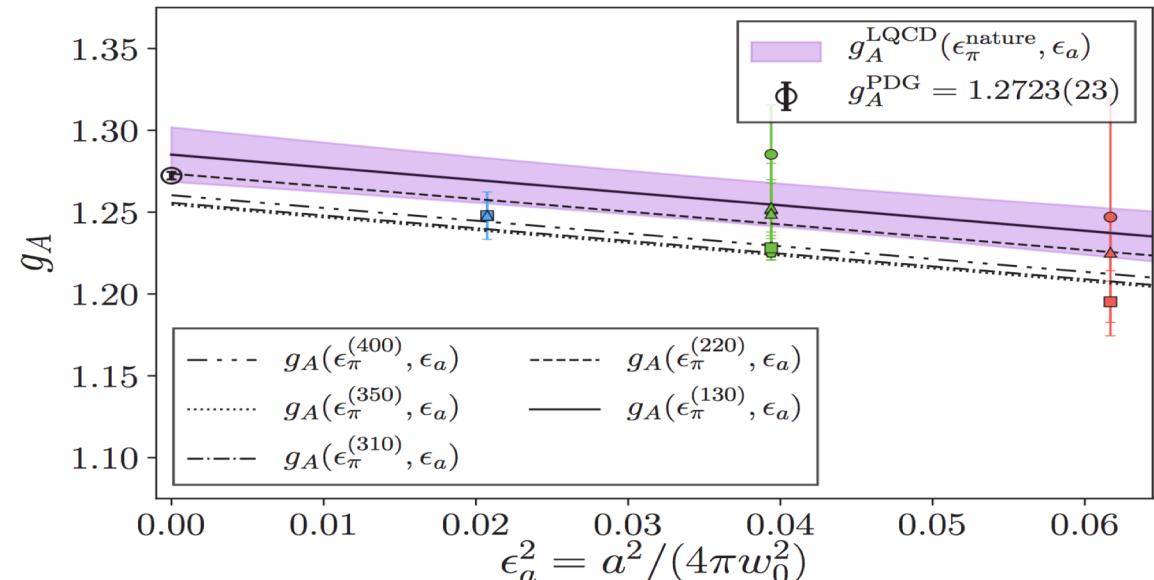
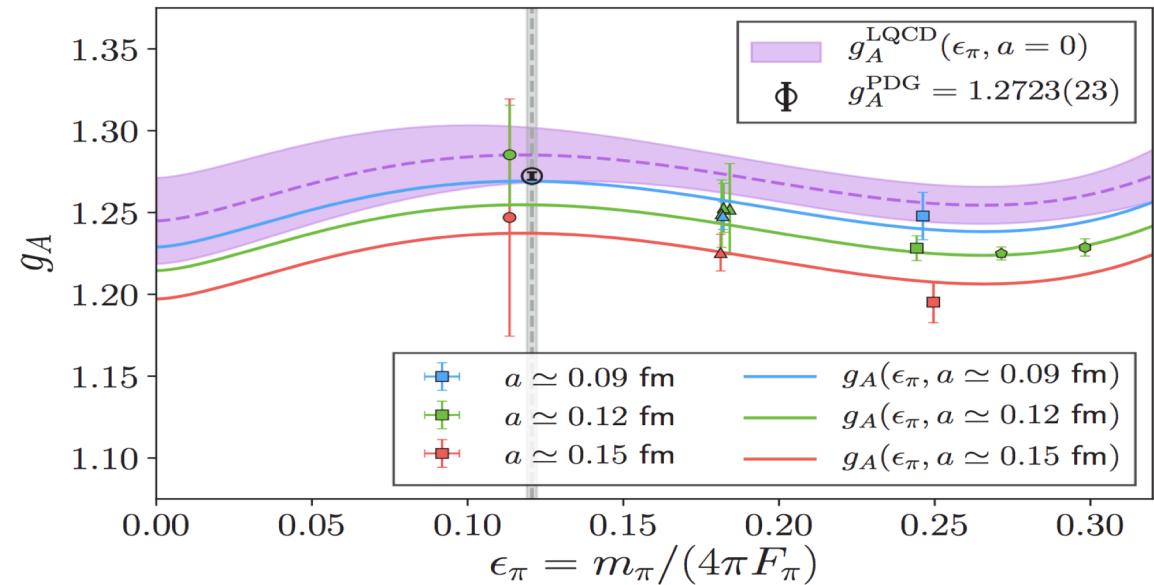
- **Lattices**

- DWF-on-HISQ,  $N_f = 2+1+1$
- $a = 0.09, 0.12$  and  $0.15$  fm
- $m_\pi = 130 - 400$  MeV

- **Measurement/Analysis**

- Gradient-flowed links
- Feynman-Hellmann-inspired current-at-all-timeslices method
- Summation method + two-state fit
- $g_A/g_V$  with  $Z_A/Z_V=1$ ,  $Z_V g_V=1$
- $(a, m_\pi, L)$  extrapolation:

$$g_A = g_0 - \epsilon_\pi^2 [(g_0 + 2g_0^3) \ln \epsilon_\pi^2 - c_2] + g_0 c_3 \epsilon_\pi^3 \\ + a_2 \epsilon_a^2 + c_4 \epsilon_\pi^4 + b_4 \epsilon_a^2 \epsilon_\pi^2 + a_4 \epsilon_a^4 \\ + (8/3) \epsilon_\pi^2 [g_0^3 F_1(m_\pi L) + g_0 F_3(m_\pi L)]$$



# PNDME '18 (Preliminary)

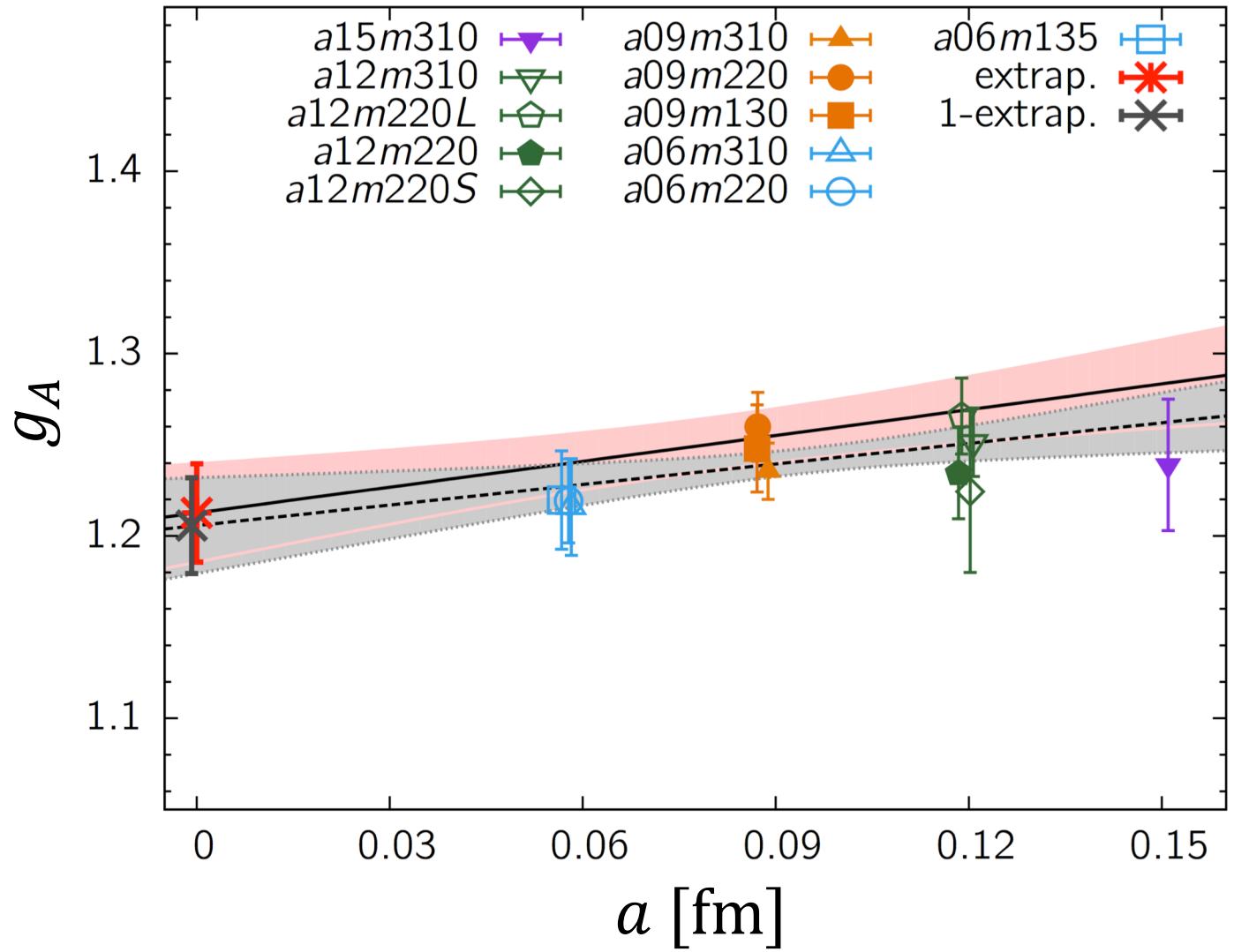
- **Lattices**

- Clover-on-HISQ,  $N_f = 2+1+1$
- $a = 0.06, 0.09, 0.12$  and  $0.15$  fm
- $m_\pi = 130 - 320$  MeV

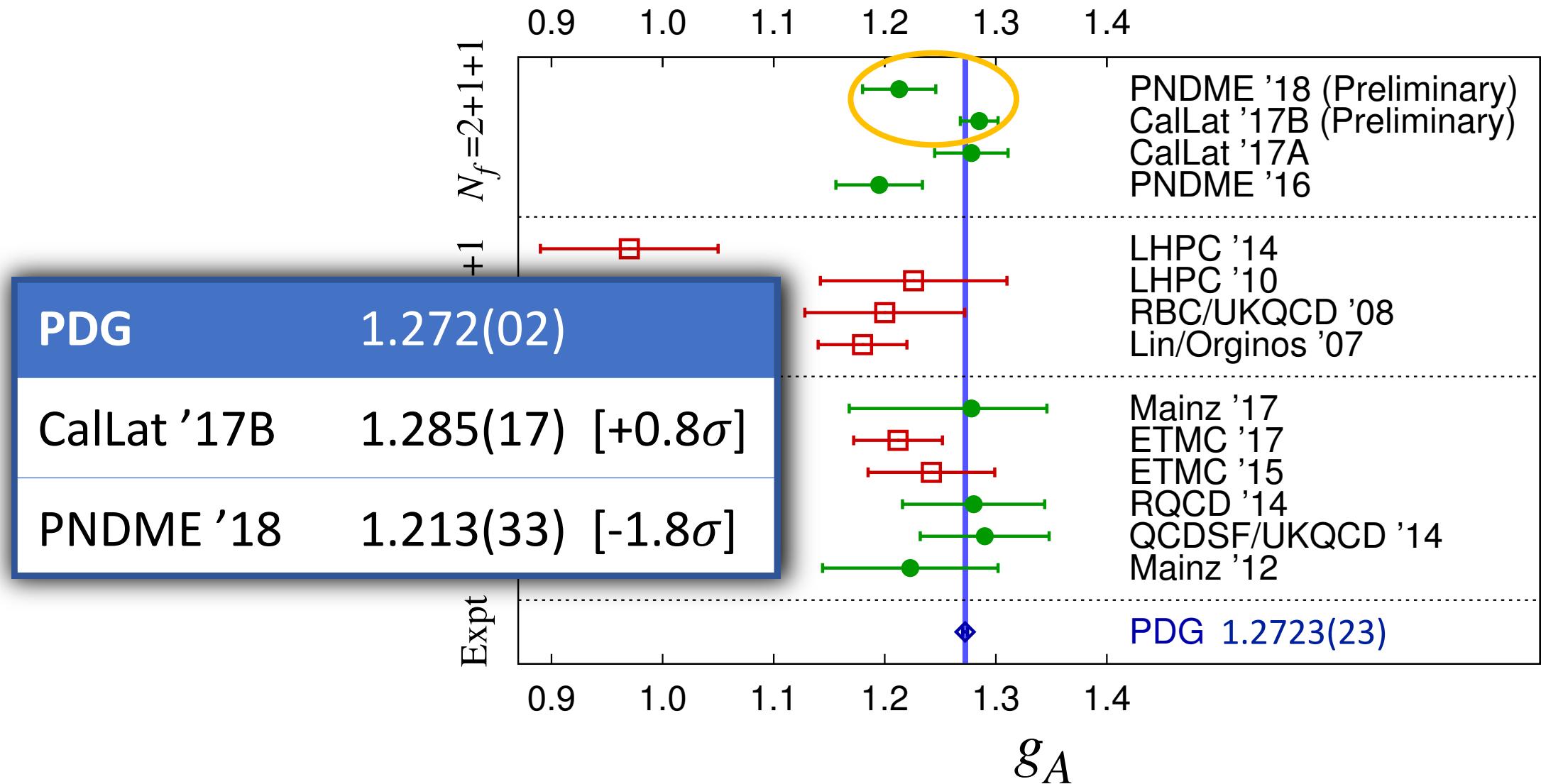
- **Measurement/Analysis**

- HYP-smeared links
- Fixed-sink-method
- 3-state fit for excited state
- NPR
- $(a, m_\pi, L)$  extrapolation:

$$g_A = c_1 + c_2 a + c_3 m_\pi^2 + c'_3 m_\pi^2 \ln(m_\pi/m_\rho)^2 + c_4 m_\pi^2 e^{-m_\pi L}$$



# $g_A$ Comparison between CalLat and PNDME

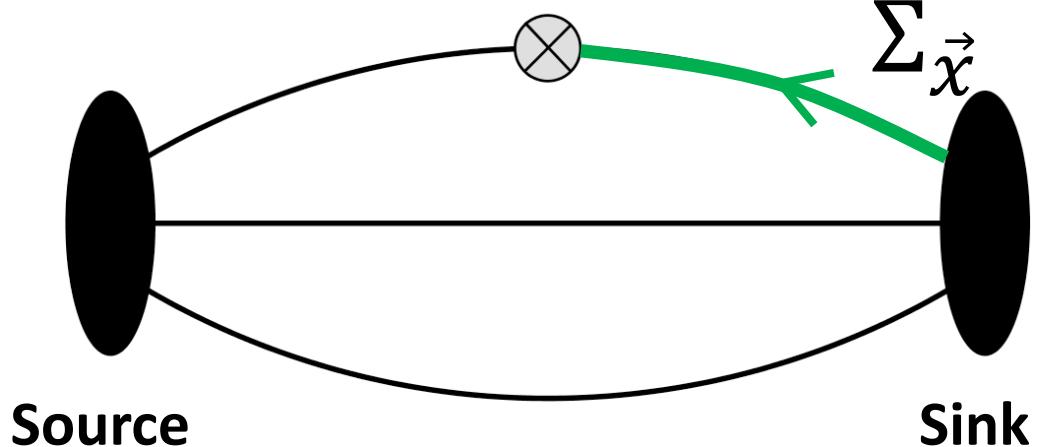


# CaLLat vs PNDME

	PNDME	CaLLat
Sequential source position in $C_{3pt}$	At sink All $\vec{p}$ & $\gamma$ , <i>single</i> $t_{sep}$	At current (at all $t$ ) Single $\vec{p}$ & $\gamma$ , <i>all</i> $t_{sep}$
Excited State	three-state fit $t_{sep} = 0.9 - 1.5$ fm	Summation method + two-state fit $t_{sep} = 0.5 - 1.5$ fm
Renormalization	NPR	$Z_A/Z_V = 1$
Leading ( $a, m_\pi$ )	$a, m_\pi^2$	$(a/w_0)^2, (m_\pi/F_\pi)^2$

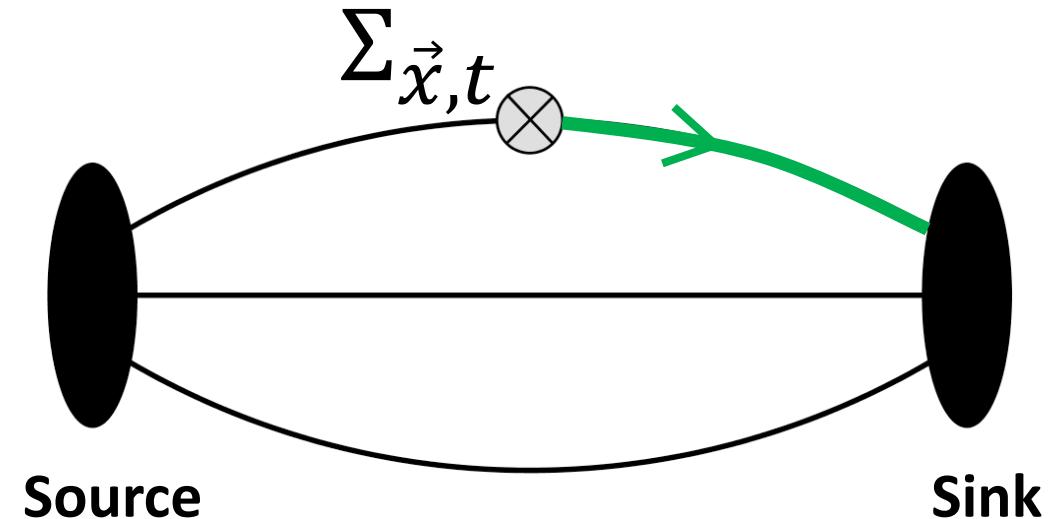
## Fixed-sink method

- Fixed sink timeslice, momentum, projection
- Sequential prop from sink
- Any operator insertions with any momenta can be obtained by contraction
- Need sequential prop for each  $t_{sep}$



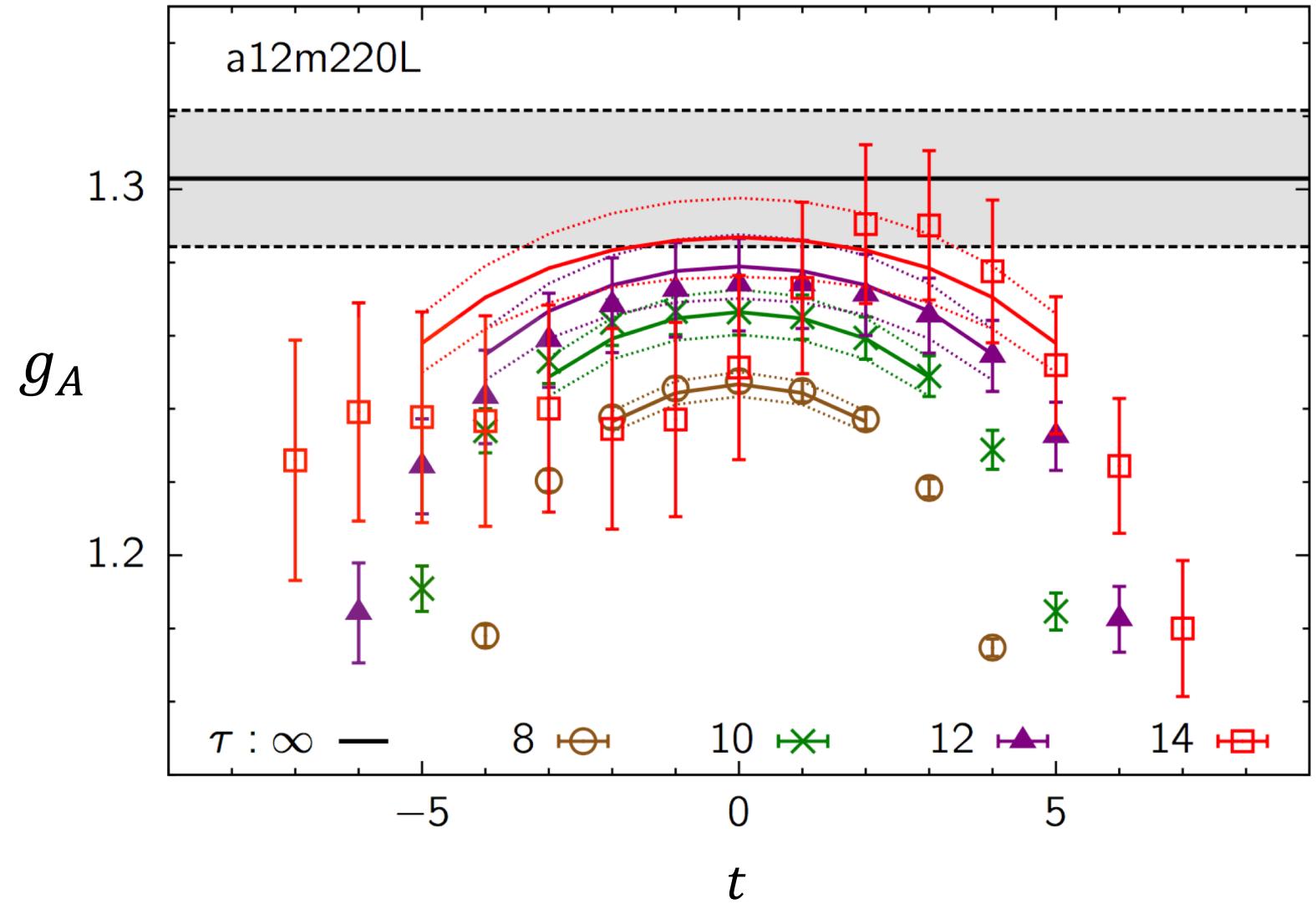
## Current-at-all-timeslices

- Current inserted at all lattice space-time
- Sequential prop from current insertion
- Any sink timeslices obtained by contraction
- Need sequential propagator calculation for each current (A, V) and momentum



# Excited States - PNDME

- 3-state fit with  
 $t_{sep} \approx 0.9 - 1.5$  fm
- Data at  $t_{sep} = 14a$   
are noisy



# Excited States - CalLat

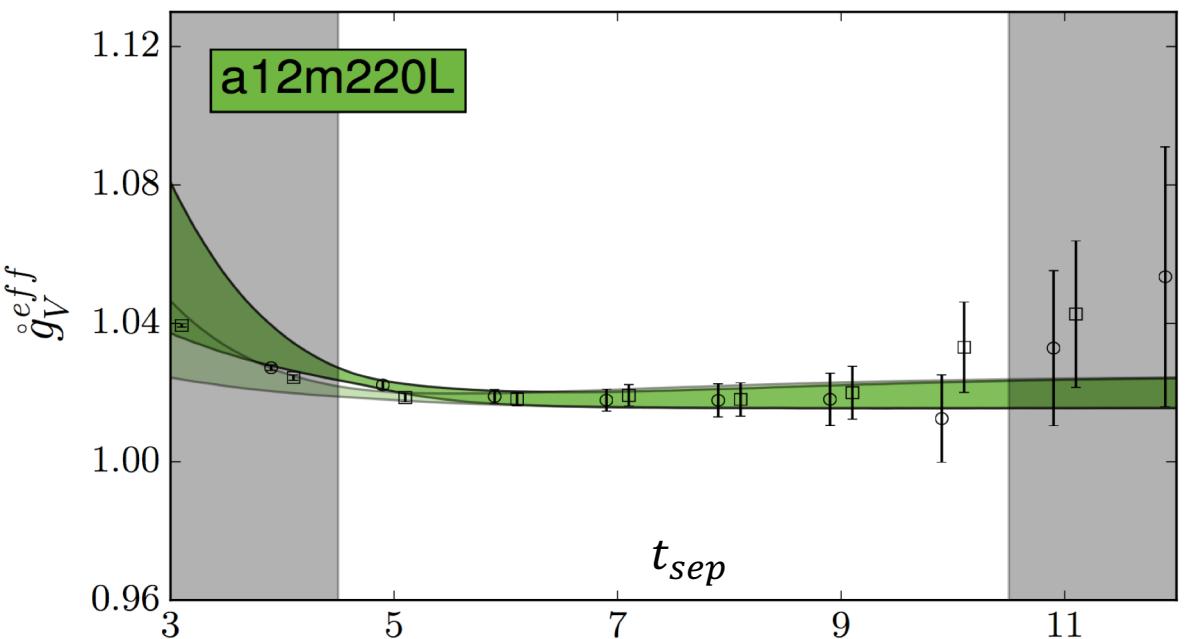
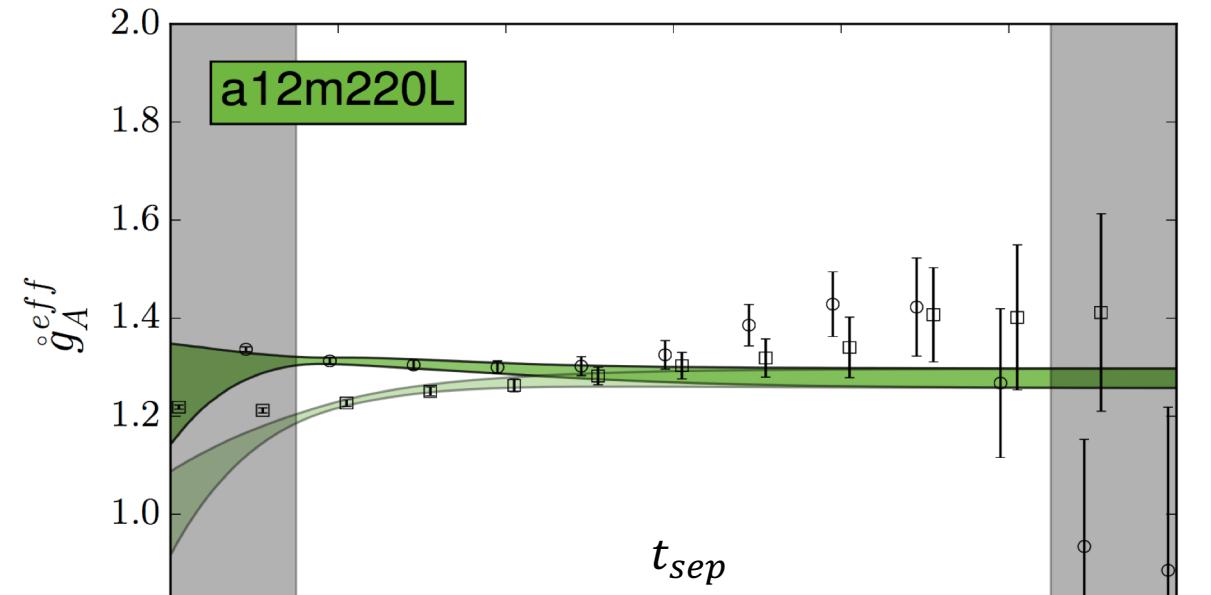
- Remove excited states by 2-point difference in **summation method**:

$$g_{A,V}^{\text{eff}}(t_{\text{sep}}) = \frac{N_{\text{3pt}}(t_{\text{sep}} + 1)}{C_{\text{2pt}}(t_{\text{sep}} + 1)} - \frac{N_{\text{3pt}}(t_{\text{sep}})}{C_{\text{2pt}}(t_{\text{sep}})}$$

Usual summation method:

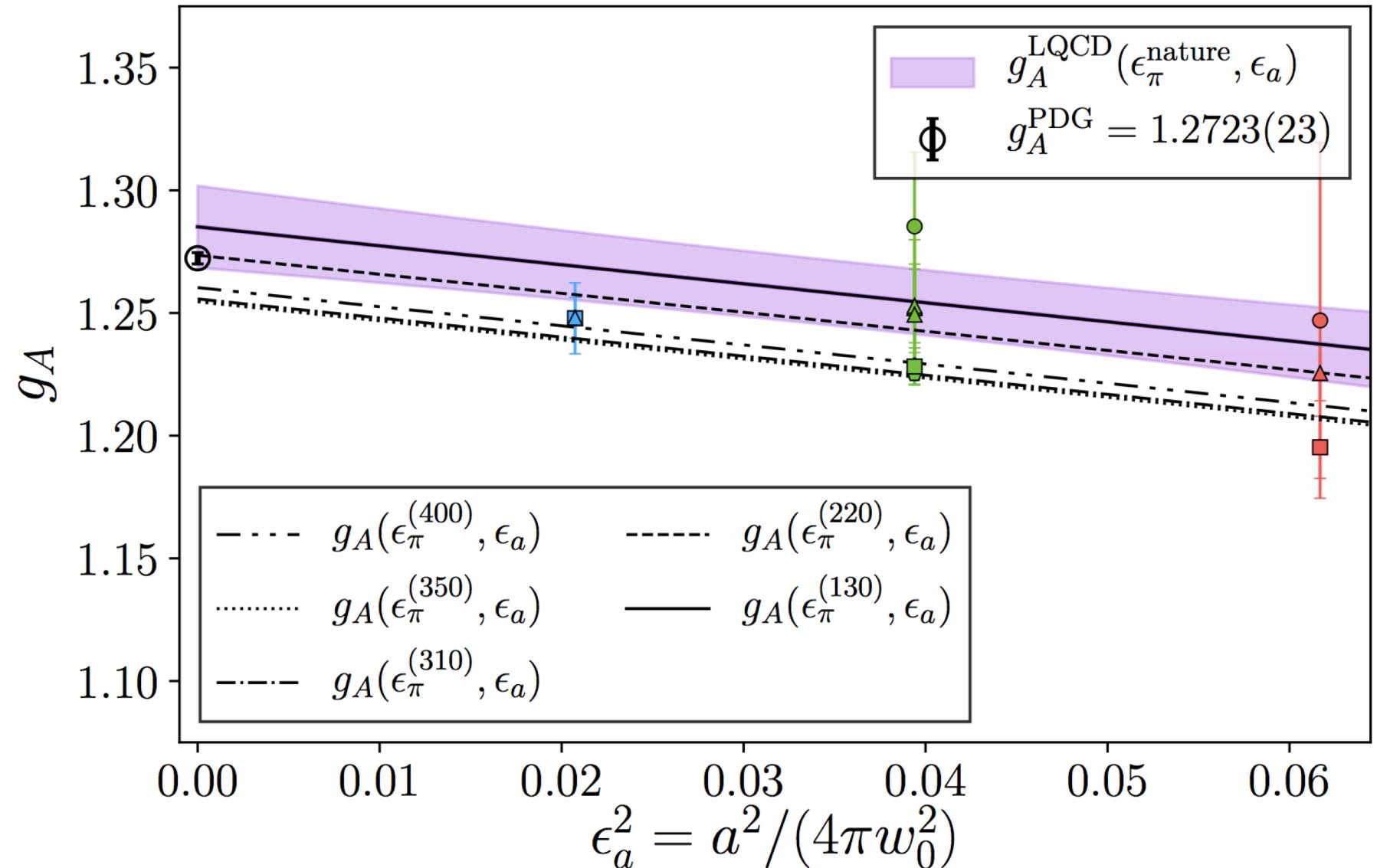
$$\sum_{t=1}^{t_{\text{sep}}-1} \frac{C_{\text{3pt}}(t_{\text{sep}}, t)}{C_{\text{2pt}}(t_{\text{sep}})} = \text{Const} + t_{\text{sep}} g_A + \dots$$

- Remove remaining excited state by two-state fit to  $g_{A,V}^{\text{eff}}$  with  $t_{\text{sep}} \approx 0.5 - 1.5 \text{ fm}$



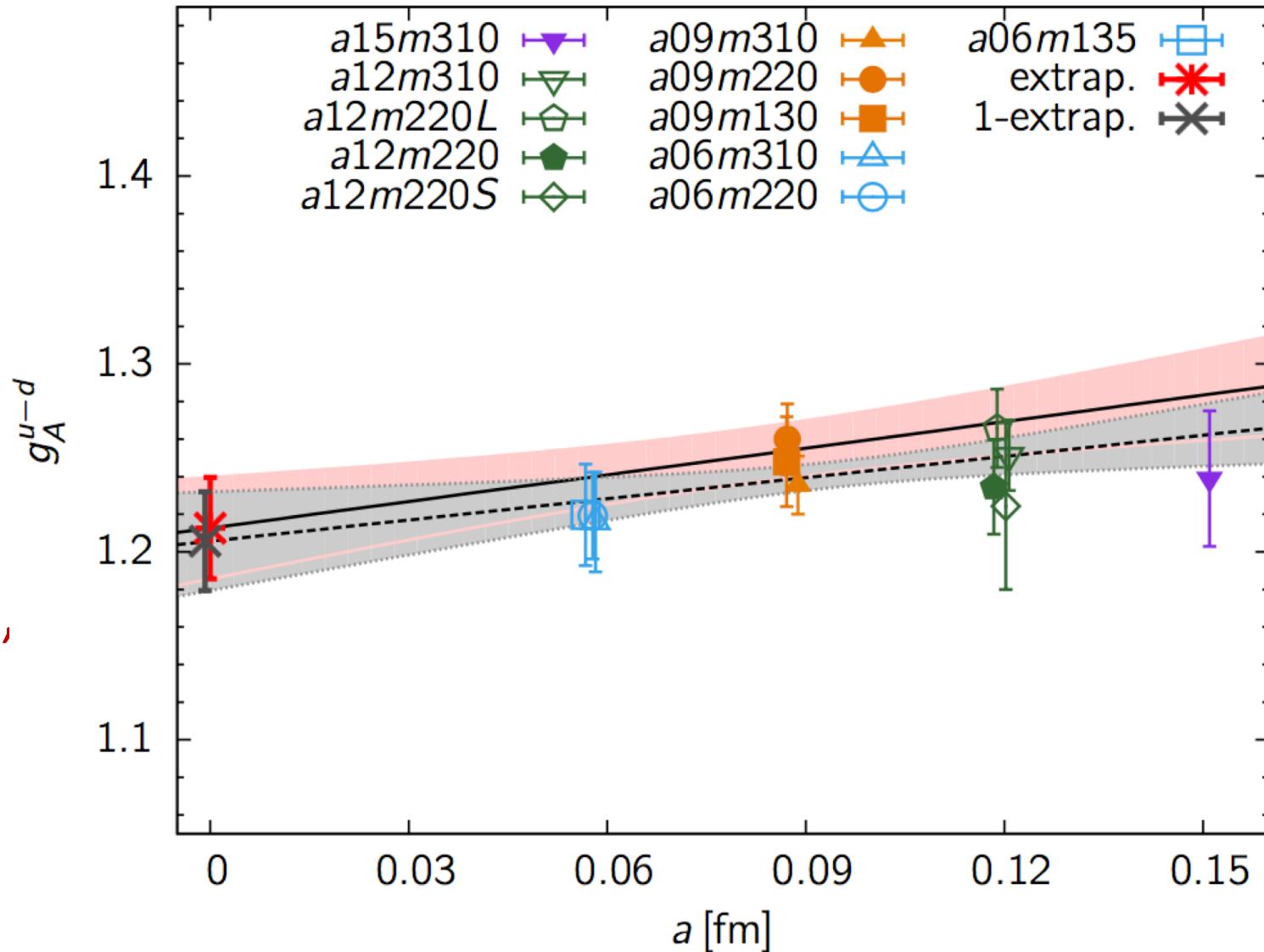
# Continuum extrapolation - CalLat

- Upward trend in  $a^2$
- $g_A = 1.285(17)$



# Continuum extrapolation - PNDME

- Downward trend in  $a$
- $g_A = 1.213(33)$
- Largely driven by  $a \sim 0.06$  fm points
- Without  $a \sim 0.06$  fm,  $g_A = 1.238(49)$



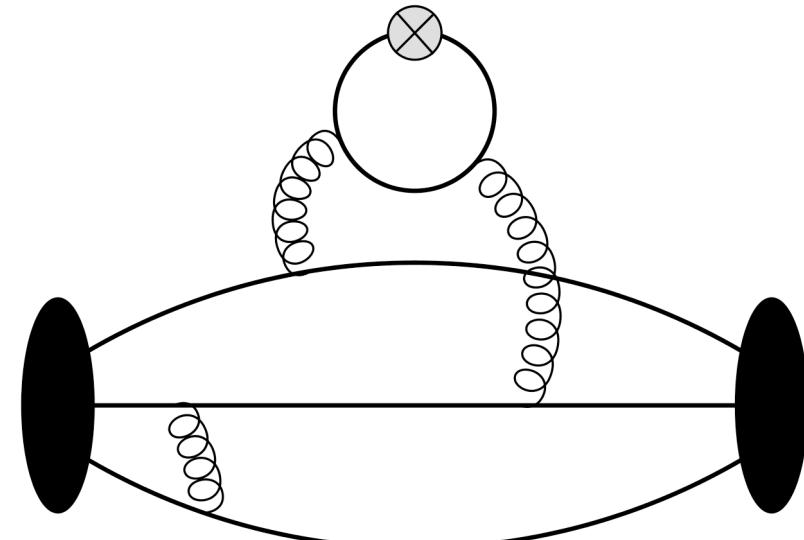
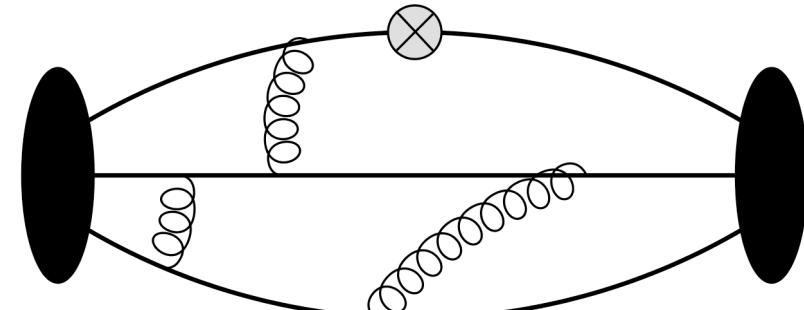
# Flavor-diagonal $g_A^u$ , $g_A^d$ , and $g_A^s$

- Proton spin decomposition

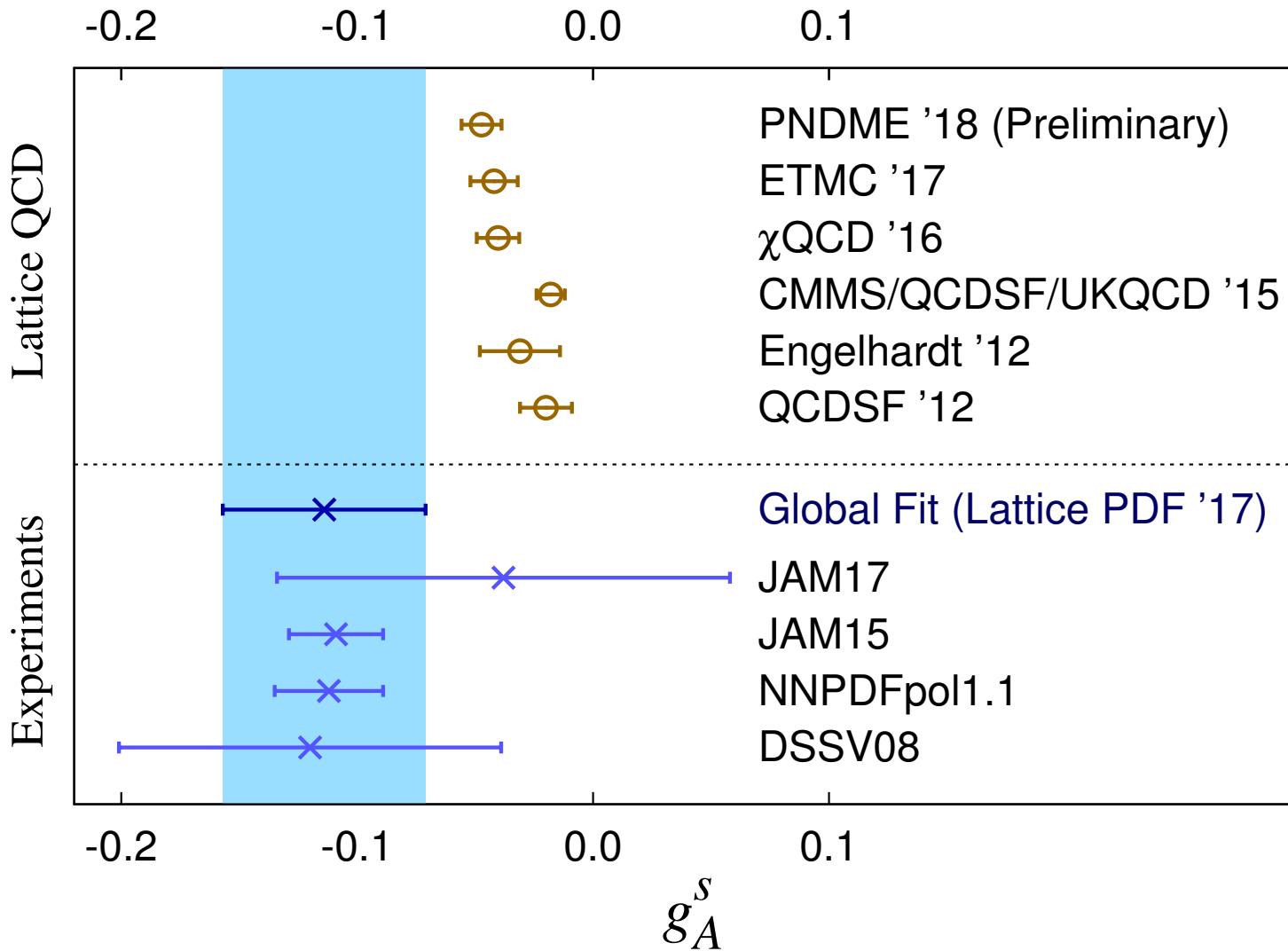
$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + L_q + \Delta G$$

$$\Delta\Sigma = g_A^u + g_A^d + g_A^s + \dots$$

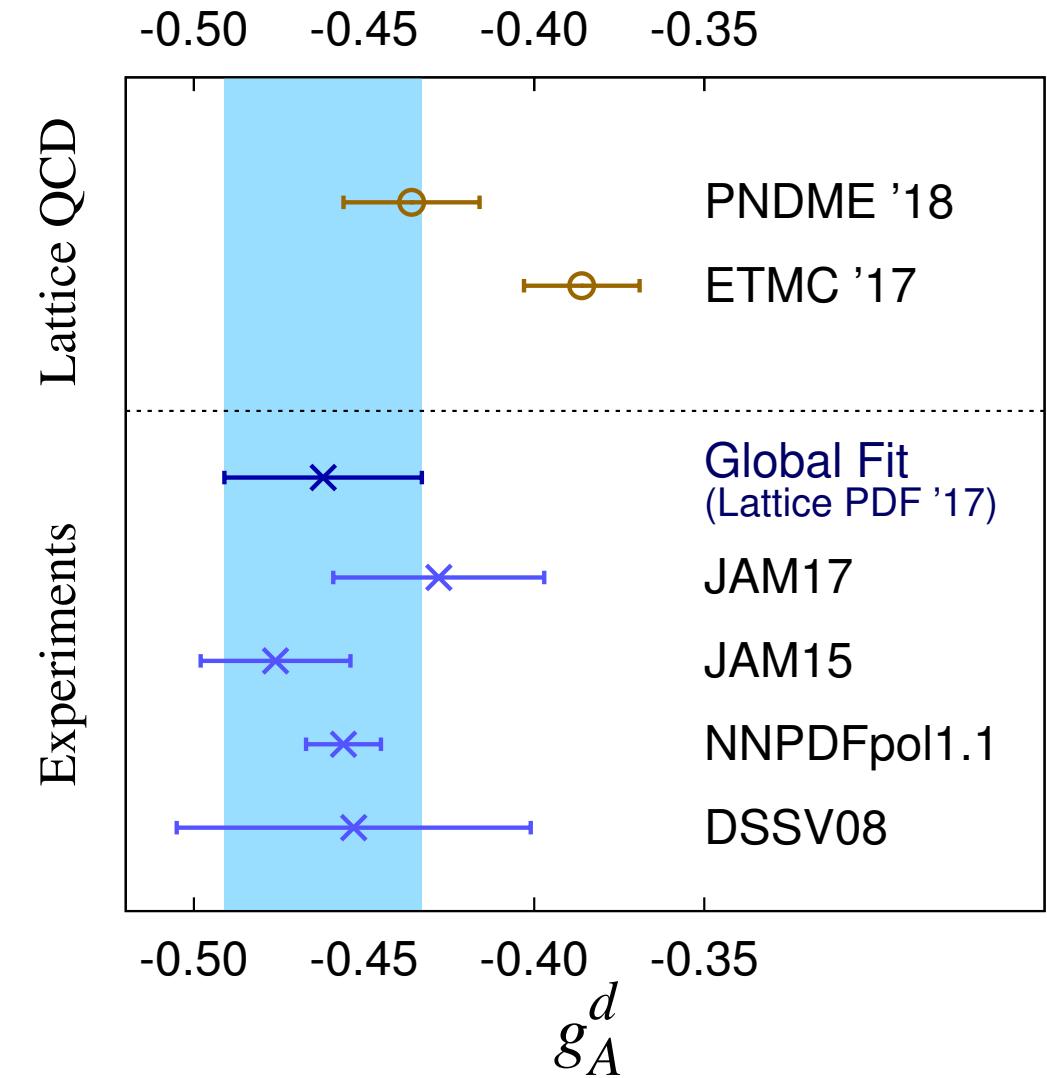
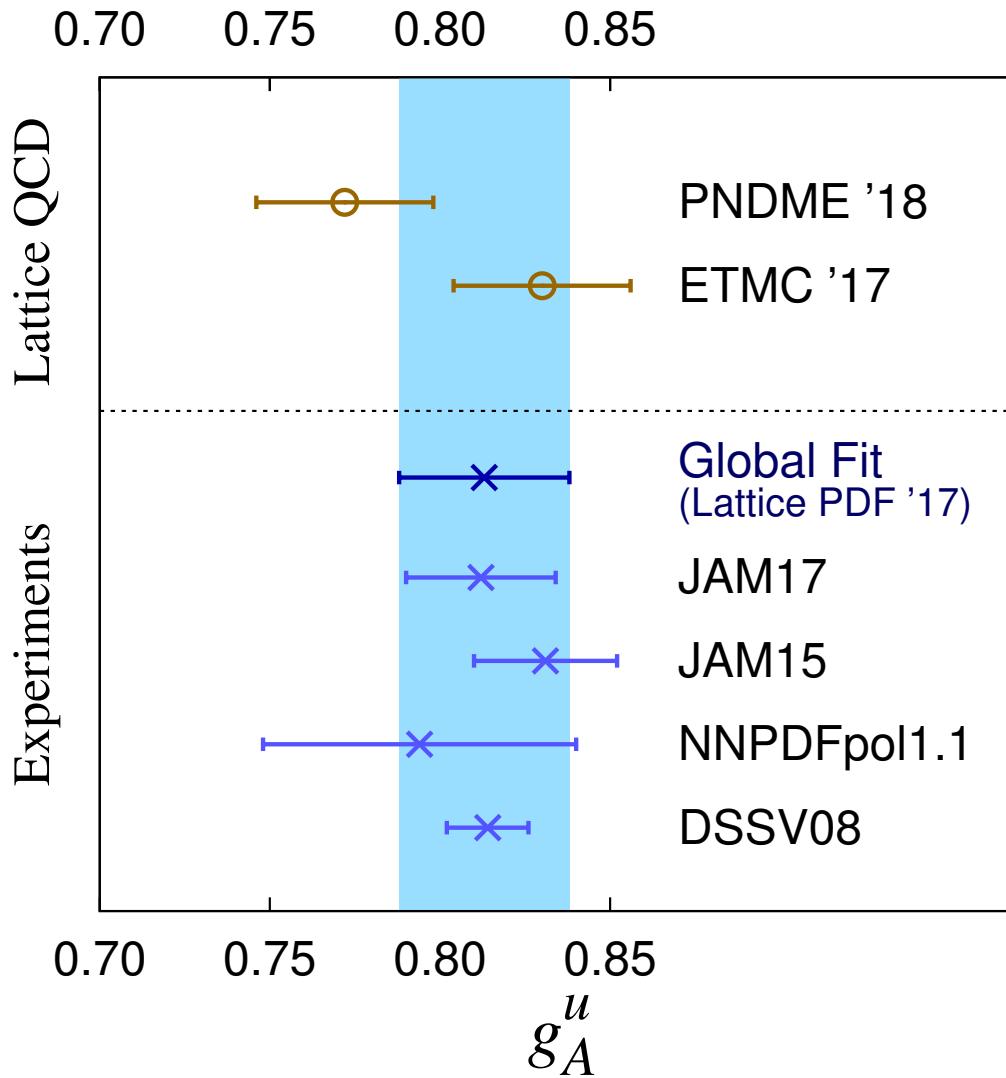
- Need evaluation of quark-line  
**connected & disconnected** diagrams



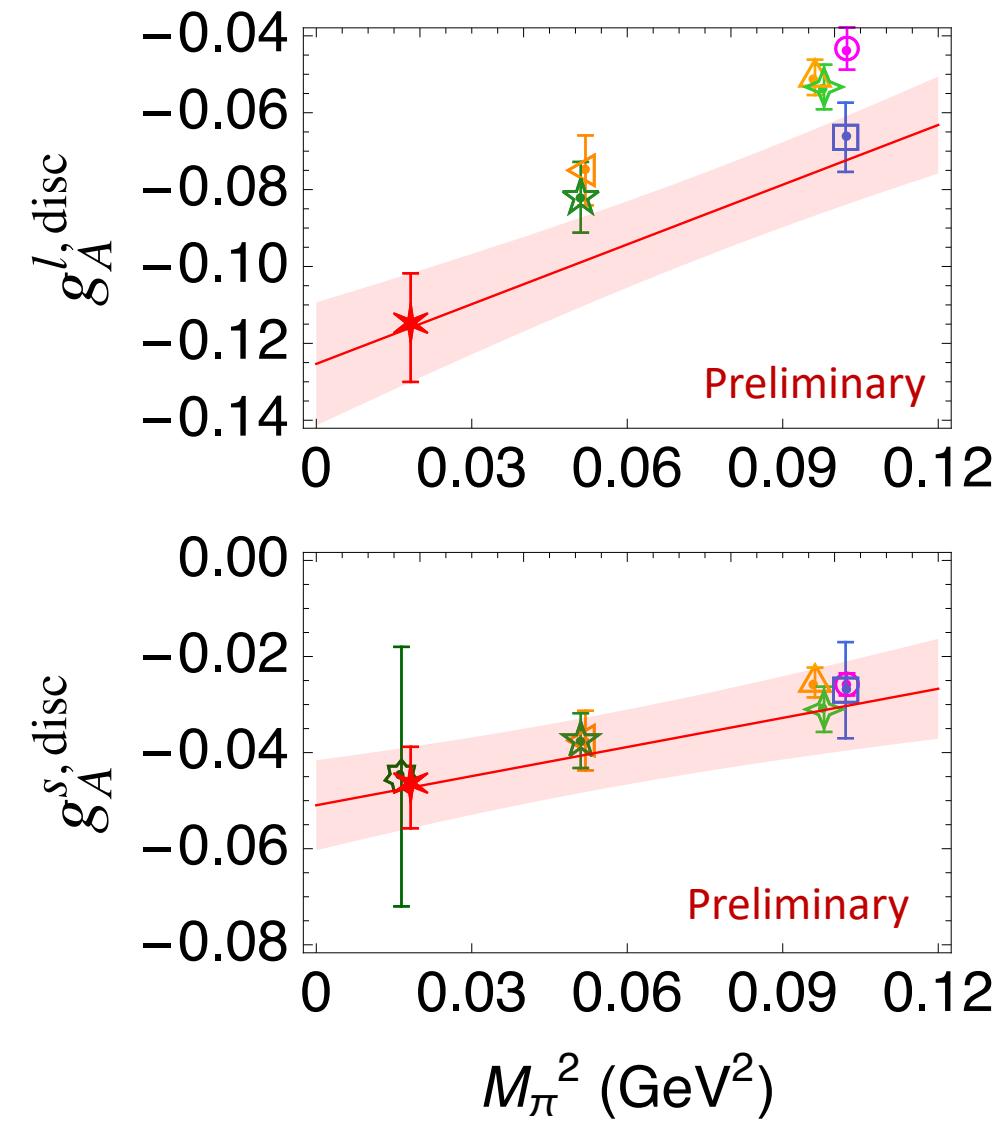
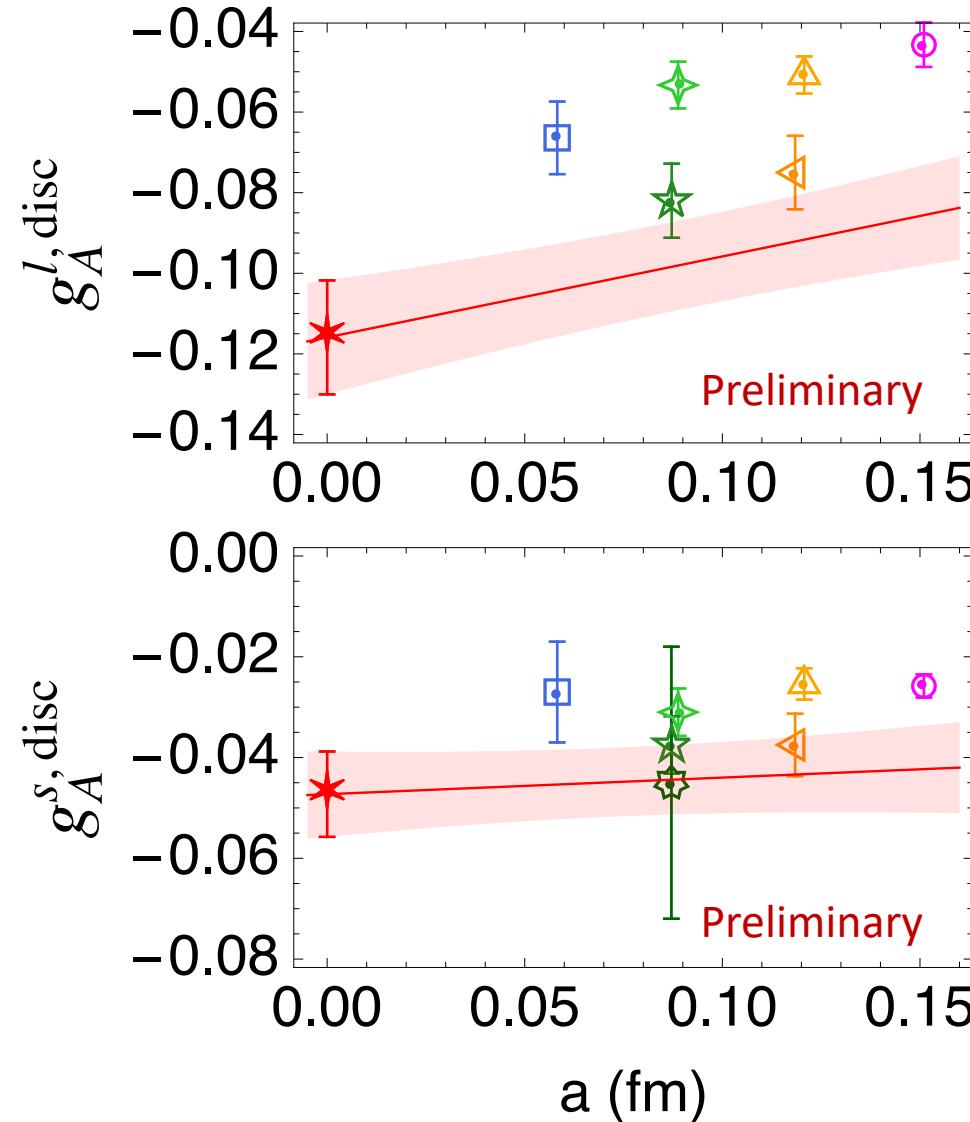
# Current Status - $g_A^S$



# Current Status - $g_A^u$ and $g_A^d$



# Disconnected contribution: PNDME



# Conclusion

- $g_A$ : CalLat provides most precise lattice QCD estimate based on fits with small  $t_{sep}$  and exact  $Z_A/Z_V = 1$
- Difference between PNDME and CalLat comes mostly from difference in  $a$ -extrapolation with results at  $a \sim 0.06$  fm
- $g_A^{u,d}$ : PNDME show a significant dependence on  $a$  and  $m_\pi$ . The larger negative disconnected contribution reduces the contribution of quarks to the proton spin:

$$\frac{1}{2}\Sigma = 0.201(18) \text{ [ETMC '17]}$$

$$\frac{1}{2}\Sigma = 0.144(17) \text{ [PNDME '18 (preliminary)]}$$