

Connecting microscopic NN interactions with heavy nuclei through DFT

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MSU. June 20th, 2018

FRIB-TA Workshop. From bound states to the continuum



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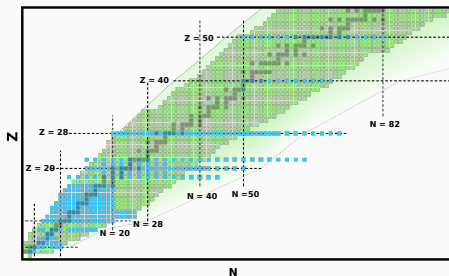
- Density Functional Theory:
 - Motivation
 - Overview of self-consistent mean field
 - Implementing a Microscopic Functional
 - Validation calculations

- Elastic Nucleus-Nucleus scattering
 - Coulomb threshold “anomaly”
 - ${}^6\text{Li} + {}^{209}\text{Bi}$ with uncertainties

- Summary and Outlook

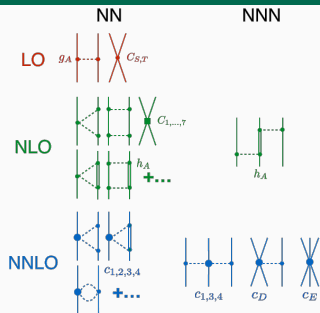
Microscopic DFT

Ab-initio methods with χ -EFT

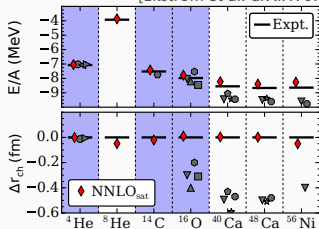


[Hergert et al. Phys. Rept. 621 (2016) 165]

- Systematic, order by order
- Light and medium nuclei ✓
- Heavy nuclei out of reach
 - A problem of scaling

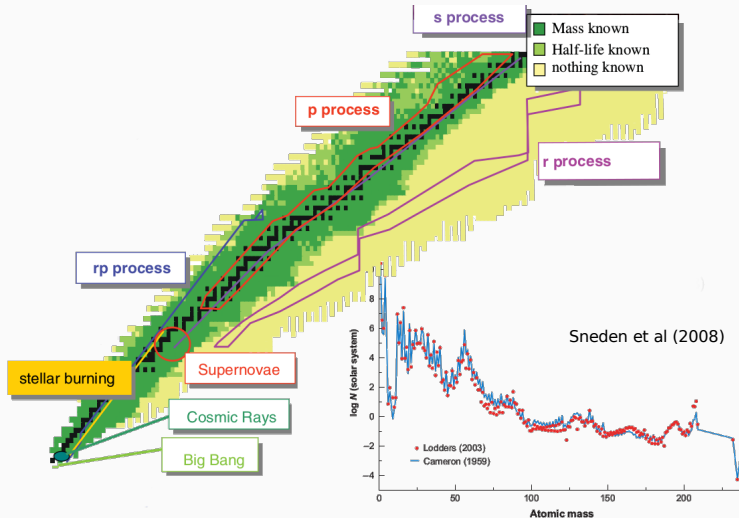


[Ekström et al. arXi:1707.09028]



[Hagen et al. Phys. Scr. 91 (2016) 063006]

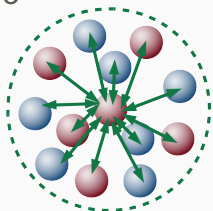
FRIB and astrophysical processes



[Grawe et al. Rept. Prog. Phys. 70 (2007) 1525]

Phenomenological DFT

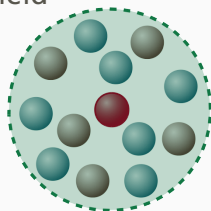
Ab-initio



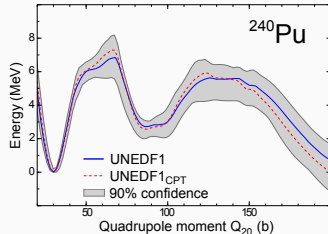
$$(T + \sum_{ij}^A V_{ij})\Psi = E\Psi$$

- Good computational scaling
- Static and dynamic properties of nuclei
- Can go “beyond mean field”
- No systematic improvement

Mean field



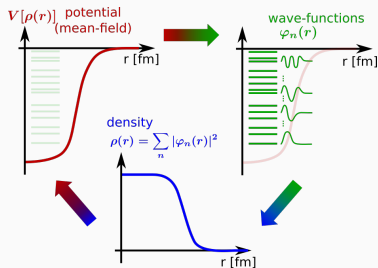
$$(T + \sum_i^A \tilde{V}_i)\Psi = E\Psi$$



[McDonnell et al. PRL114 (2015) 122501]

Self-Consistent Mean Field

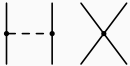



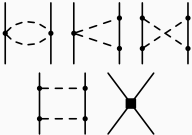
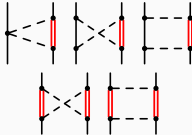



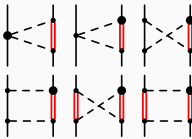
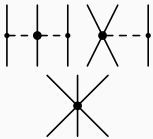

- DFT, based on HFB theory
- Non-linear eigenvalue problem
- Iterative solutions
- Phenomenological interactions
 - Fitted to hundreds of nuclear data
 - Contact interactions: Skyrme
 - Finite range: Gogny



Self-consistent mean-field theory

[<https://commons.wikimedia.org/>]

Microscopically constrained EDF: Chiral expansion

	NN force		$3N$ force	
	Δ -less EFT	Δ contributions	Δ -less EFT	Δ contributions
LO				
NLO				
N^2 LO				

[A. Dyhdalo, et al. PRC95 (2017) 054314]

Structure of the potential and the EDF

- (Extremely) schematic 2 body chiral potential

$$\hat{V}_{NN}(r_1, r_2) = \underbrace{\delta(r_1 - r_2)}_{\text{yields Skyrme-like}} + \underbrace{V_F(r_1 - r_2)}_{\text{sum of Yukawas}}$$

- Yukawas expanded as a sum of Gaussians
 - Treat Hartree term of the EDF like a Gogny force
- 3N force has similar structure but:
 - Hartree term is zero in time-reversal invariant systems
 - Fock term induced by finite range has non-local densities
- Use density matrix expansion for all exchange terms
 - Recast V_{NN} and V_{3N} into Skyrme-like terms
 - Reduces computational cost significantly
 - Also reduces numerical stability issues

The best of both worlds

DFT Component

- Contact terms for short range physics
- Adjusted to nuclear properties
- Encodes many body correlations

+

Microscopic Component

- Derived from χ -EFT
- Long range physics, pions
- Order by order
- Non-Local density!

A scalable framework with systematic improvements

Microscopically constrained DFT

- Non-local densities for finite range potentials

$$E_H^{NN} \sim \int dRdr \langle r | V^{NN} | r \rangle \rho_1 \left(R + \frac{r}{2} \right) \rho_2 \left(R - \frac{r}{2} \right)$$

$$E_F^{NN} \sim \int dRdr \langle r | V^{NN} | r \rangle \rho_1 \left(R - \frac{r}{2}, R + \frac{r}{2} \right) \rho_2 \left(R + \frac{r}{2}, R - \frac{r}{2} \right) \hat{P}_{12}$$

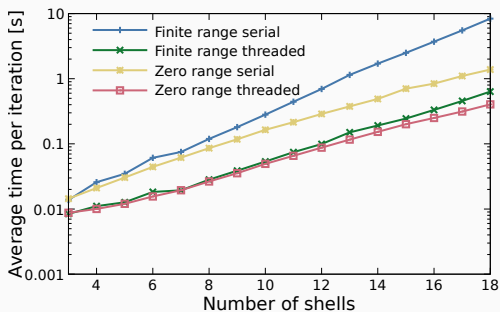
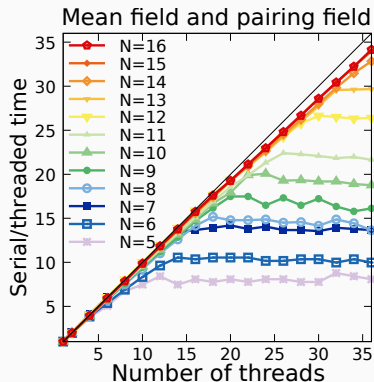
- Density Matrix Expansion (DME)

$$\begin{aligned} \rho \left(R + \frac{r}{2}, R - \frac{r}{2} \right) &\approx \Pi_0^\rho(k_F r) \rho(R) \\ &+ \frac{r^2}{6} \Pi_2^\rho(k_F r) \left[\frac{1}{4} \Delta \rho(R) - \tau(R) + \frac{3}{5} k_F^2 \rho(R) \right] \end{aligned}$$

Like a Taylor expansion for the non-local density

DFT Numerical Solver

HFBTHO v3 takes advantage of current HPC architectures

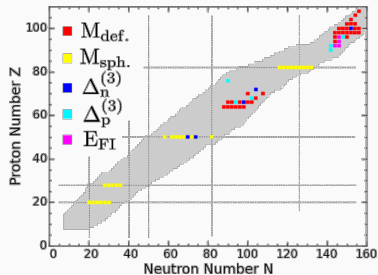


- Finite-range calculation not much more expensive than zero-range when threading

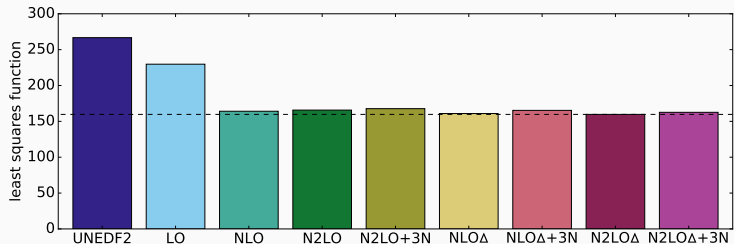
[RNP, Schunck, Lasserri, Zhang CPC220 (2017) 363]

Optimizing DFT component

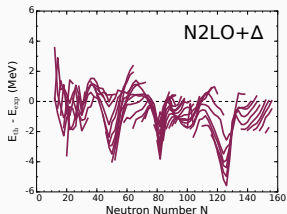
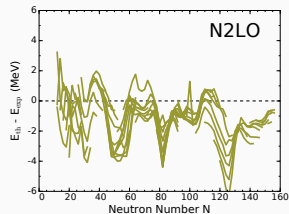
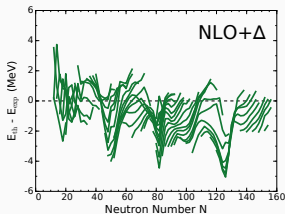
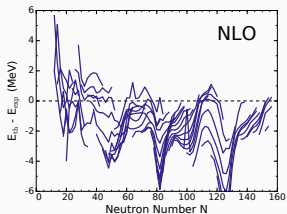
- UNEDF2 optimization protocol
- 130 data, 14 parameters
- Masses, radii, fission isomers, spin-orbit splittings, nuclear matter



[Schunck et al. EPJA51 (2015) 169]



Mass Tables

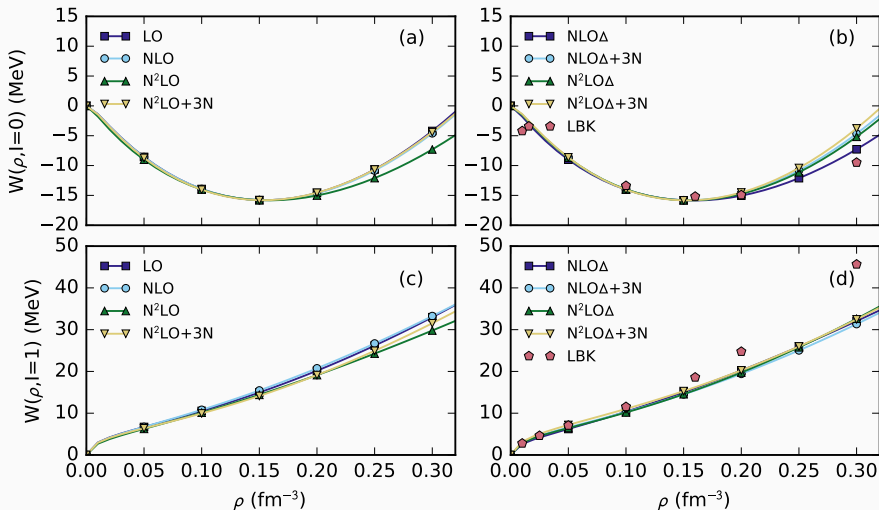


Order by order improvement

EDF	r.m.s. deviation
UNEDF2	1.98
LO	1.99
NLO Δ	1.41
N2LO Δ	1.26

[RNP, Schunck, Dyhdalo, Furnstahl, Bogner. PRC97 (2018) 05430]

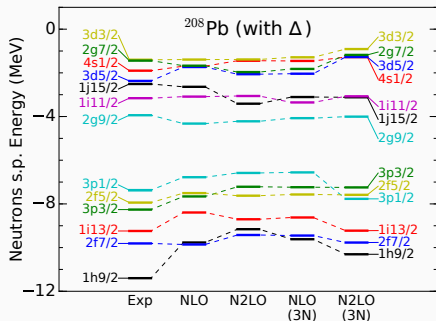
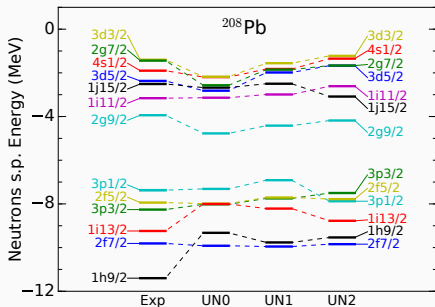
Nuclear Matter and Neutron Matter



[RNP, Schunck, Dyhdalo, Furnstahl, Bogner. PRC97 (2018) 05430]

Single-Particle Spectra

Quantitatively comparable to UNEDF results

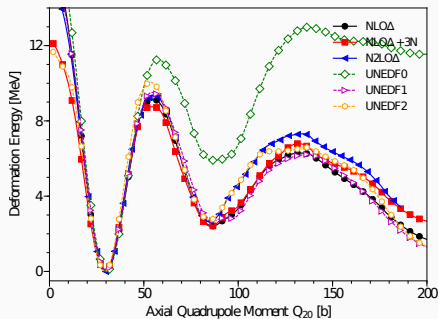
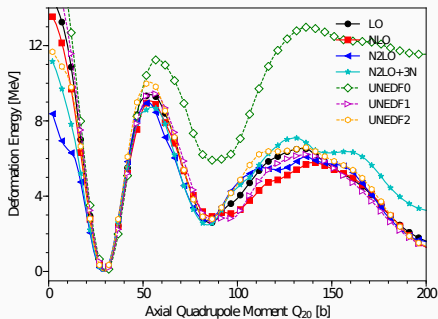


[RNP, Schunck, Dyhdalo, Furnstahl, Bogner. PRC97 (2018) 05430]

- Single-particle energies from blocking calculations
- Exactly the same conditions for all EDFs

Deformation Properties

Quality of fission barriers is comparable to other EDFs



[RNP, Schunck, Dyhdalo, Furnstahl, Bogner. PRC97 (2018) 05430]

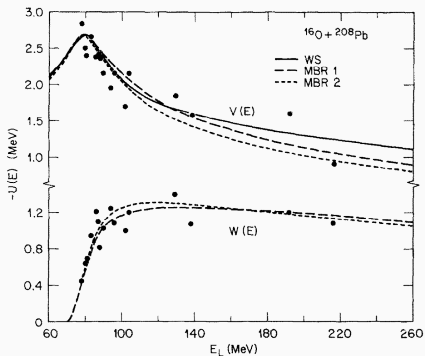
- Inclusion of fission isomers in fitting protocols constrains fission barriers
- Variations up to 2 MeV in height of fission barriers

Elastic Nucleus-Nucleus scattering

Elastic Nucleus-Nucleus scattering

- Near the Coulomb barrier
- Energy dependence of the effective (optical) potential
- Imaginary part decreases with energy
- Real part shows an "anomaly"
- A consequence of causality
 - Dispersion relation

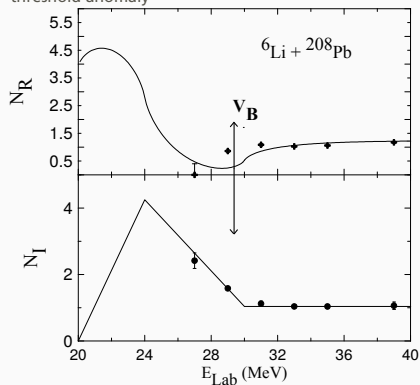
$$V(r; E) = V_0(r) + \frac{P}{\pi} \int_0^{\infty} \frac{W(r; E')}{E' - E} dE'$$



[Mahaux et al. NPA449 (1986) 354]

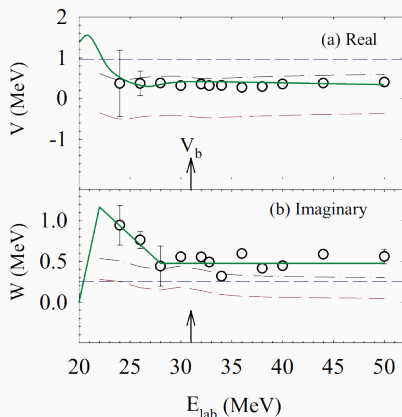
Looking at weakly bound Nuclei

New manifestation of the dispersion relation: Breakup threshold anomaly



[Hussein et al. PRC 73 (2006) 044610]

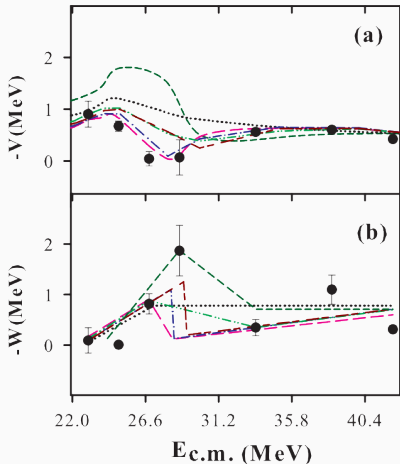
Reaction mechanisms involving weakly bound ${}^6\text{Li}$ and ${}^{209}\text{Bi}$ at energies near the Coulomb barrier



[Santra et al. PRC 83 (2011) 034616]

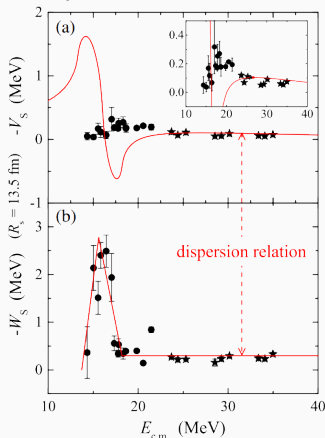
Looking at weakly bound Nuclei

Investigation of the threshold anomaly for the ${}^7\text{Li} + {}^{159}\text{Tb}$ system



[Patel et al. PRC 91 (2015) 054614]

Is the Dispersion Relation Applicable for Exotic Nuclear Systems? The Abnormal Threshold Anomaly in the ${}^6\text{He} + {}^{209}\text{Bi}$ System



[L. Yang et al. PRL 119 (2017) 042503]

Looking at weakly bound Nuclei

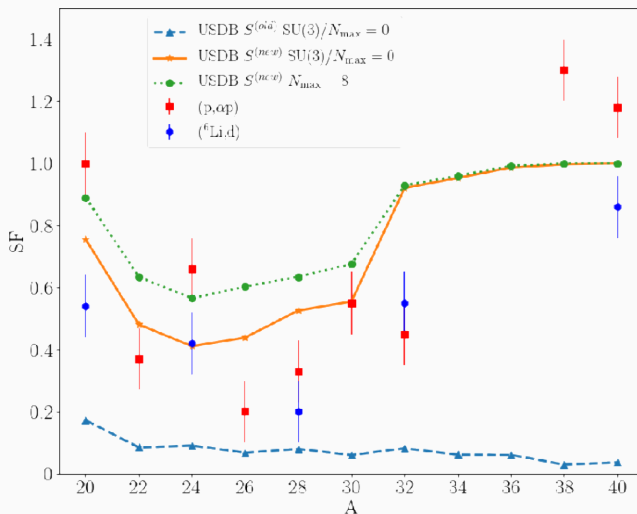
- This is a cautionary tale in:
 - Underestimated experimental uncertainties
 - Misunderstood theoretical uncertainties
 - Good software used as a black box
 - Ignoring the underlying physics
- Let's focus on one of these four examples
 - Elastic ${}^6\text{Li} + {}^{209}\text{Bi}$ cross sections
 - Measured at energies from 24 to 50 MeV
 - Use an effective potential with Woods Saxon shape

$$V(r; E) = -\frac{V_E}{1 + e^{(r-R_v)/a_v}} - i\frac{W_E}{1 + e^{(r-R_w)/a_w}}$$

- Integrate with Numerov and match to Coulomb WF

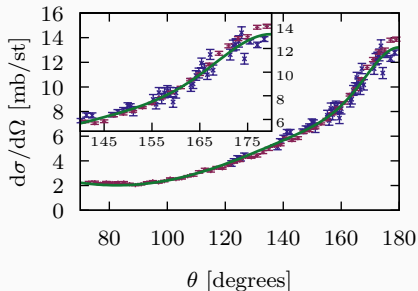
Incompatible experimental data

During Heiko's talk ...



Incompatible experimental data

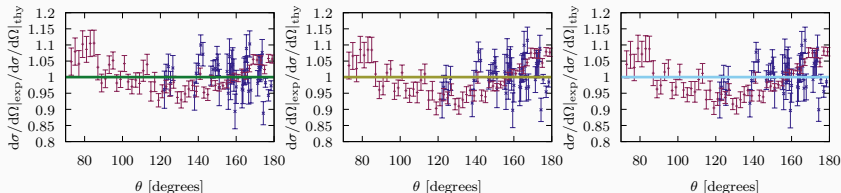
- Fits **NEVER** give $\chi^2/\text{d.o.f.} \lesssim 1$
 - Restrictive model ? \rightarrow Improve model
 - Mutually incompatible data ? \rightarrow Reject data
- np $d\sigma/d\Omega$ at 162 MeV
- Statistical and systematic errors over or underestimated
- 3σ criterion
 - Fit to all data ($\chi^2/\text{d.o.f.} > 1$)
 - Remove data sets with improbably high or low χ^2
 - Refit parameters



[RNP, Amaro, Arriola PRC88 (2013) 064002]

Incompatible experimental data

- Mutually incompatible data
 - Which experiment is correct?
 - Is any of the two correct?
- Maximization of experimental consensus
 - Fit to all data ($\chi^2/\text{d.o.f.} > 1$)
 - Apply 3σ criterion
 - Refit parameters
 - Re-apply 3σ criterion to all data
 - Repeat until no more data is excluded or recovered

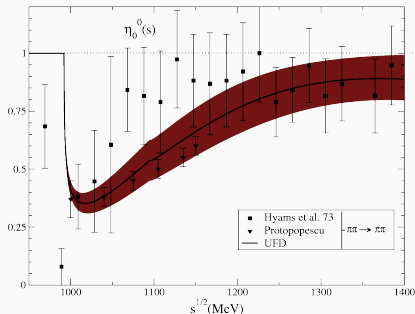
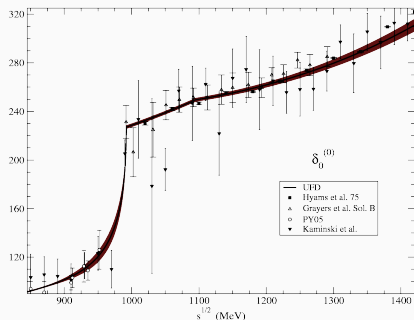


[RNP, Amaro, Arriola PRC88 (2013) 064002]

Incompatible experimental data

$\pi - \pi$ scattering data analysis

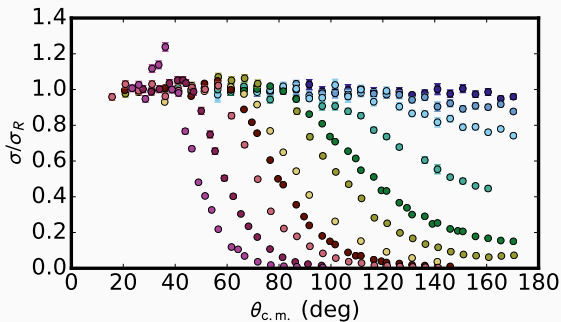
Elastic and inelastic data



[RNP, Arriola, Ruiz de Elvira, PRD91 (2015) 074014]

Incompatible experimental data

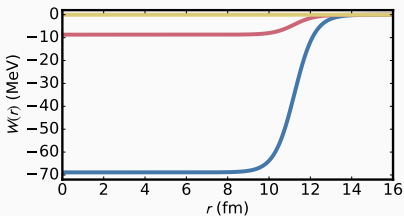
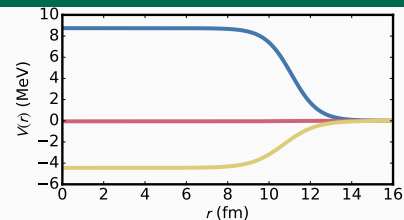
${}^6\text{Li} + {}^{209}\text{Bi}$ data



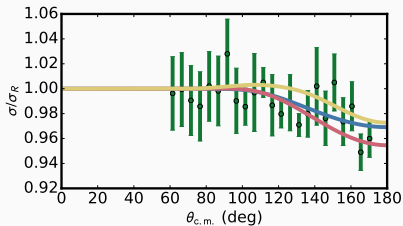
One single experiment at different energies

3σ criterion still applicable

Large uncertainties at low energies (24 MeV)



Fit simultaneously at different energies and “Bootstrap” data



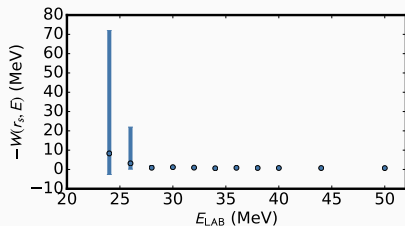
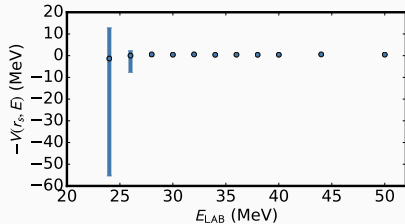
3 different effective potentials

$$V(r, E) = -V_E / (1 + e^{(r-r_v)/a_v}) - iW_E / (1 + e^{(r-r_w)/a_w})$$

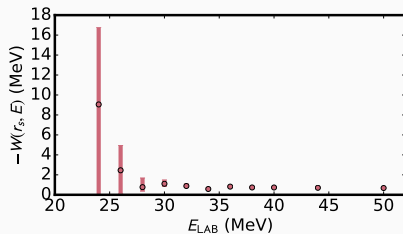
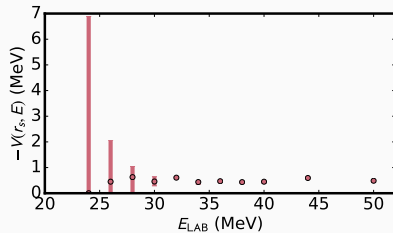
- No constraints
- Attractive potentials
- Small imaginary part

3 Types of analysis

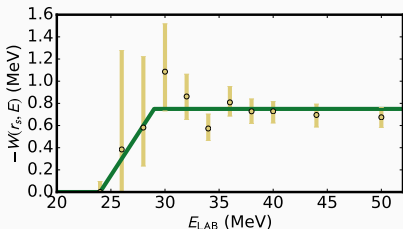
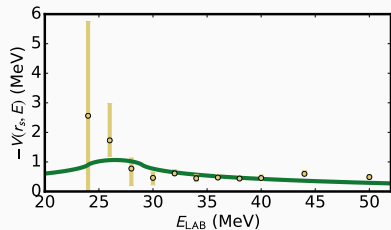
Unconstrained fits



Attractive potentials



3 Types of analysis



- Constraining some parameters
- Penalty function
- $\chi^2 \rightarrow \chi^2 + (W_{24})^2 + (W_{26})^2 + (W_{28})^2$
- This is a prior from the Bayesian point of view
- Results are compatible with dispersion relation

Microscopic DFT

- New family of EDFs constrained by χ -EFT
 - Up to N2LO with/o 3N and Δ -excitations
- Quality EDFs with global predictive power and microscopic underpinning
- Surprising improvement in mass calculations

Elastic Nucleus-Nucleus scattering

- UQ has meaningful consequences
- Physics needs to be included in any analysis

Microscopic DFT

- Effect of 3N forces
 - Δ 's improve performance, 3N terms don't
- Quantification and propagation of uncertainties
 - Truncation, fitting bias, sensitivity to LEC's
- Use the same EDF in pairing channel

Elastic Nucleus-Nucleus scattering

- Propagate uncertainty to dispersion relation
- Include dispersion relation by construction
- Analyze other reactions with weakly bound nuclei



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