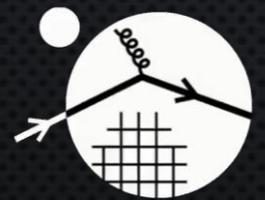

Digitization of Scalar Fields for NISQ-era Quantum Computing

arXiv 1808.10378: NK and Martin J. Savage

Next Steps in Quantum Science for HEP
Fermilab, September 12-14, 2018

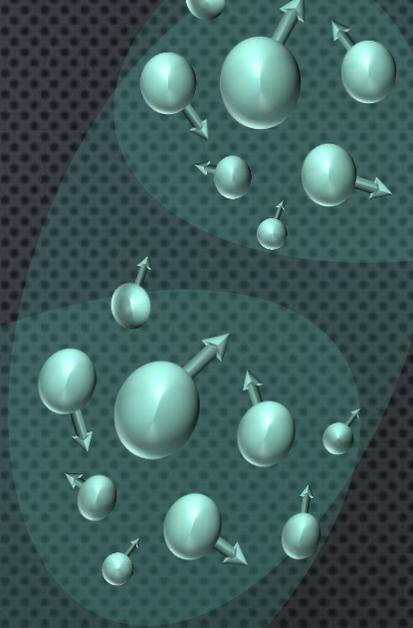
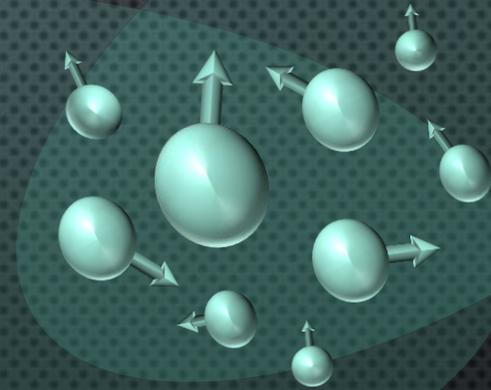
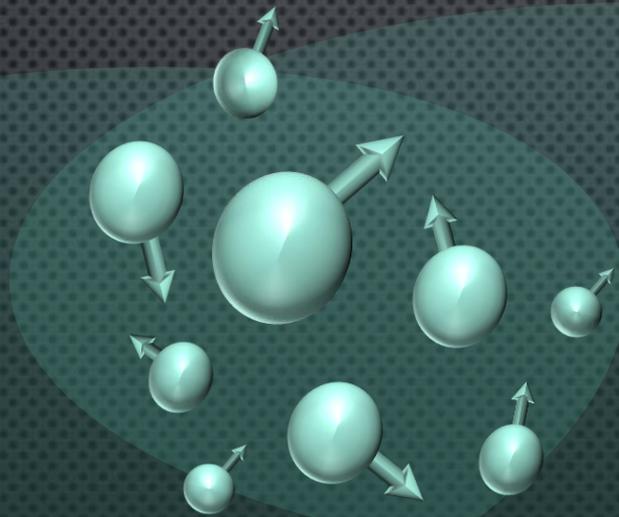
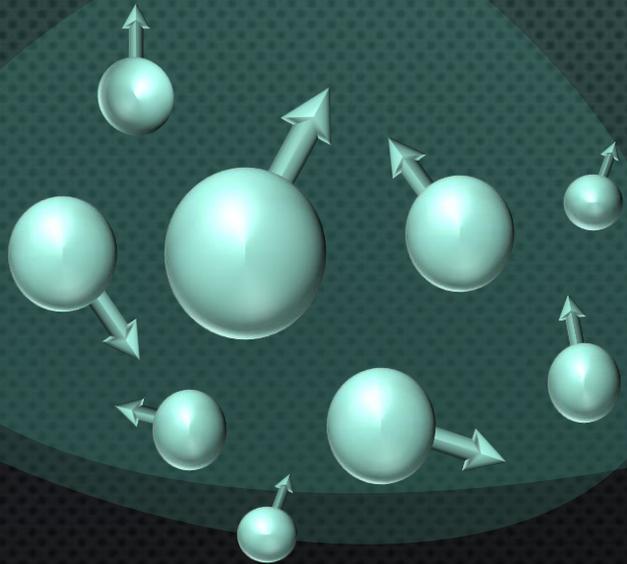


Natalie Klco



Structure of Site Qubit Registers

- Entanglement
- Interactions
- Communications



$$\bar{H} = \sum_x \frac{1}{2} \bar{\Pi}^2 + \frac{1}{2} \bar{\phi}^2 + \frac{\bar{\lambda}_0}{4!} \bar{\phi}^4 - \frac{1}{2} \bar{\phi} \bar{\nabla}_a \bar{\phi}$$

Jordan-Lee-Preskill

arXiv: 1111.3633

arXiv: 1112.4833

Digitization Bases

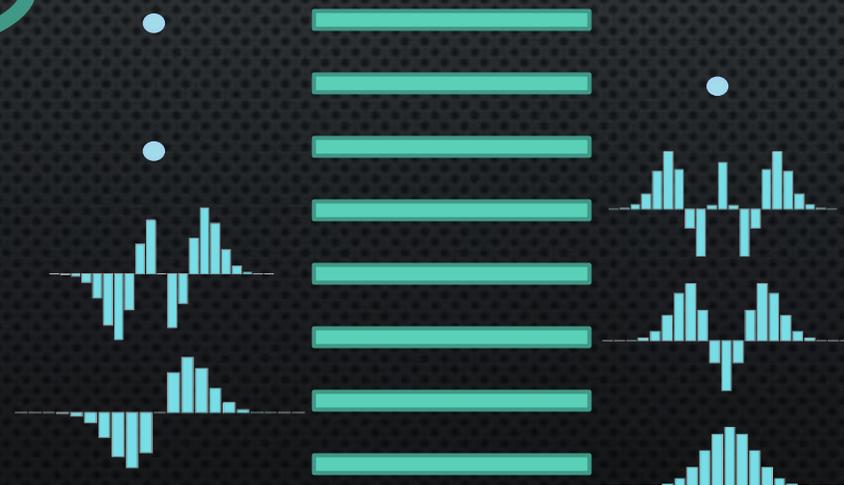
$$\bar{\phi} = \frac{\bar{\phi}_{max}}{2^{n_Q} - 1} \sum_{j=0}^{n_Q-1} 2^j \sigma_j^Z$$

Somma arXiv: 1503.06319

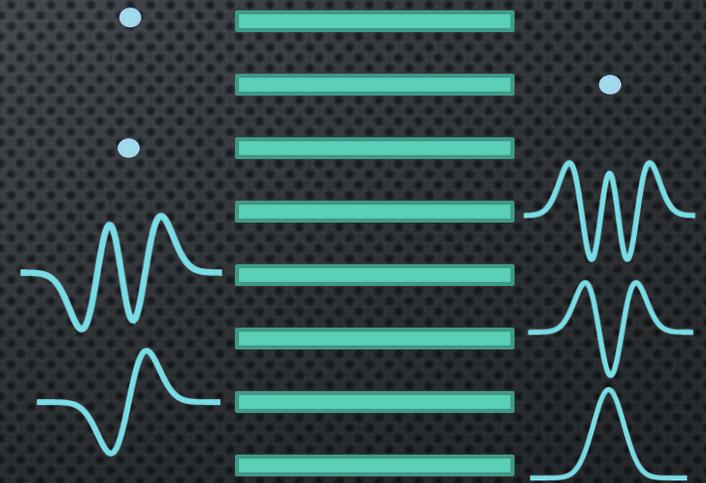
Macridin-Spentzouris-
Amundson-Harnik

arXiv: 1802.07347

arXiv: 1805.09928



Harmonic Oscillator



\vec{k} -space mode expansion
 Yeter-Aydeniz and Siopsis
 arXiv: 1709.02355

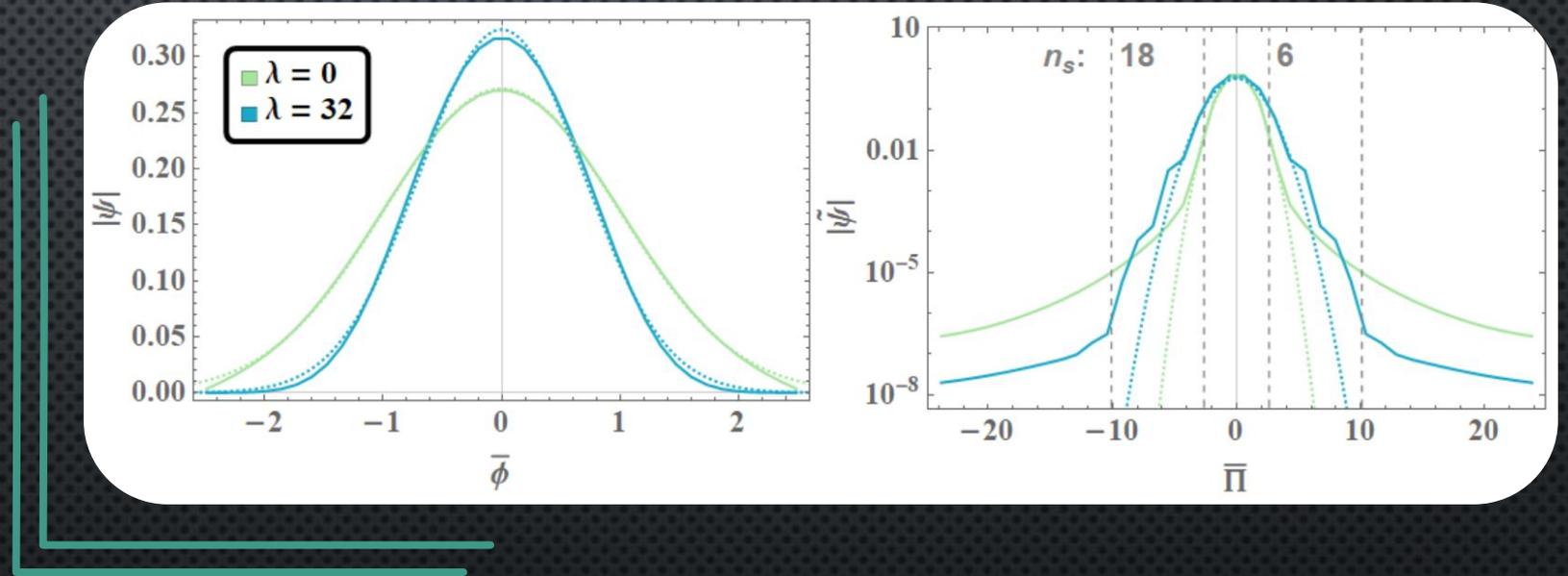
$$\delta\phi = \frac{2\bar{\phi}_{max}}{2^{n_Q} - 1}$$

Nyquist-Shannon Sampling Theorem

continuous (field)
signal with
finite bandwidth
(conjugate-momentum)

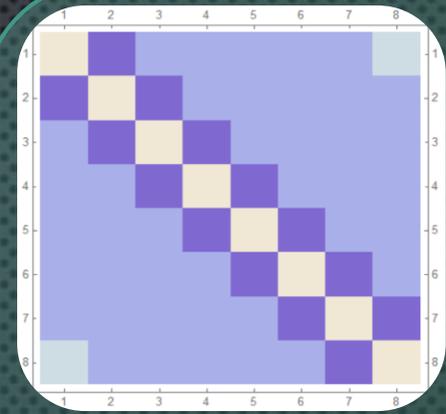
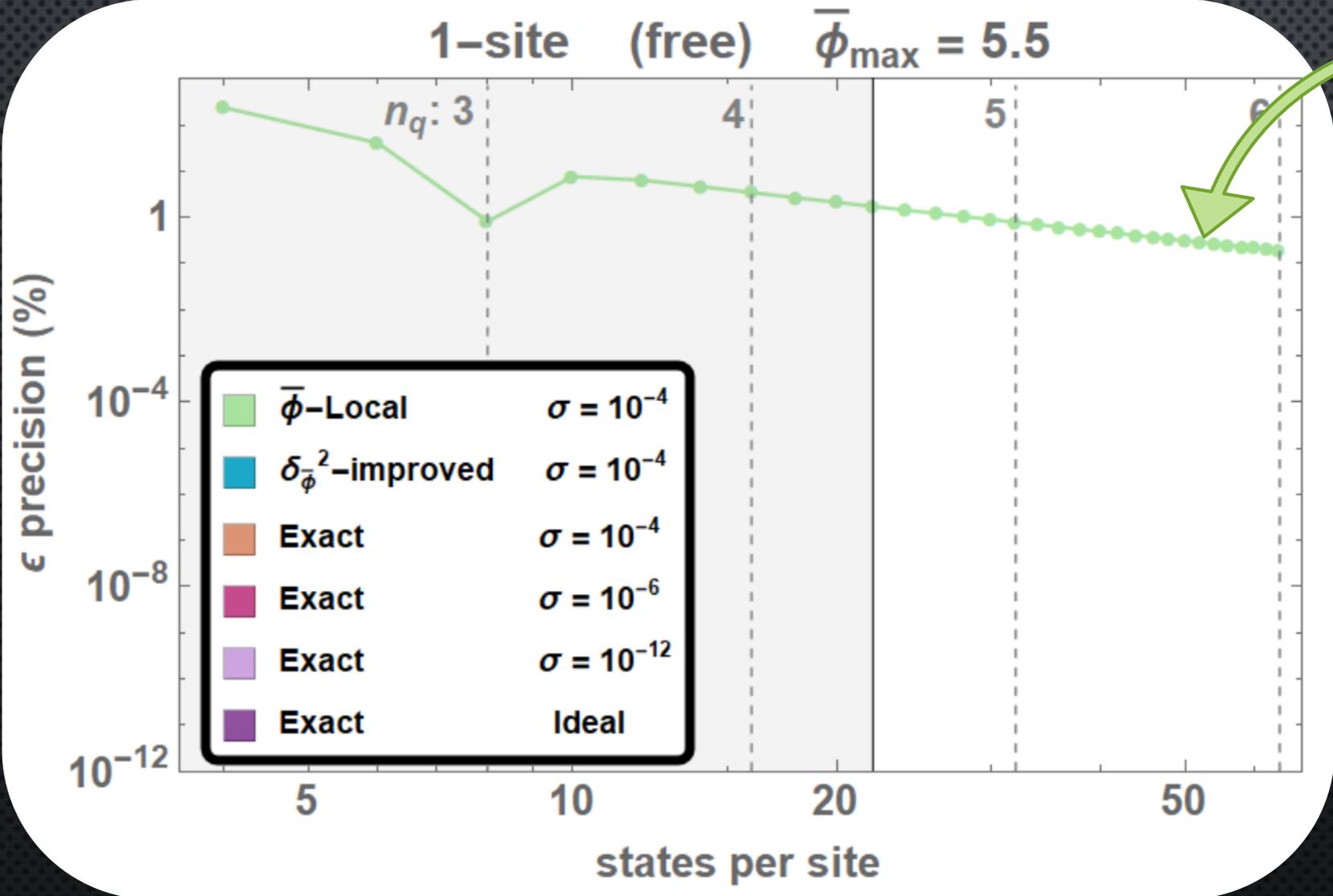


finite temporal sample rate
sufficient

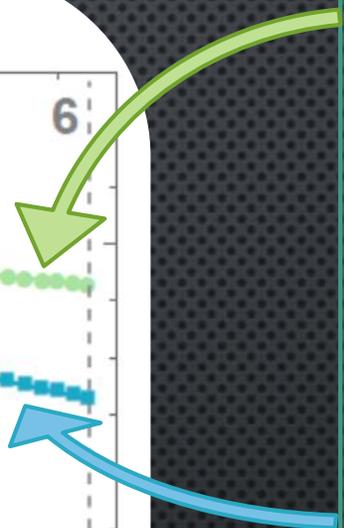
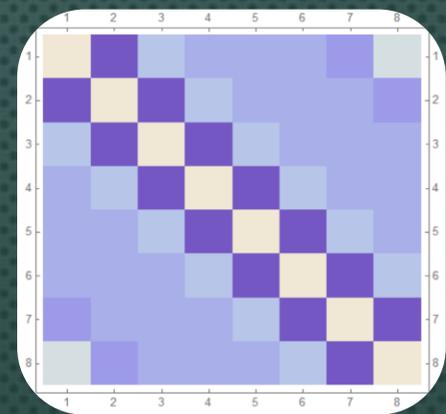
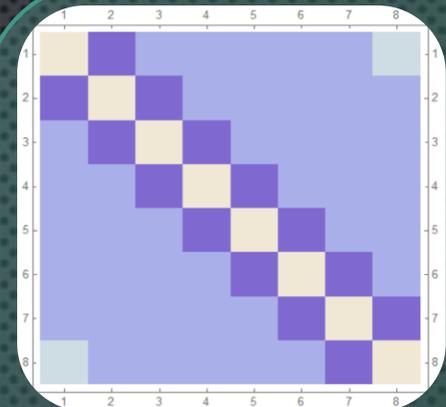
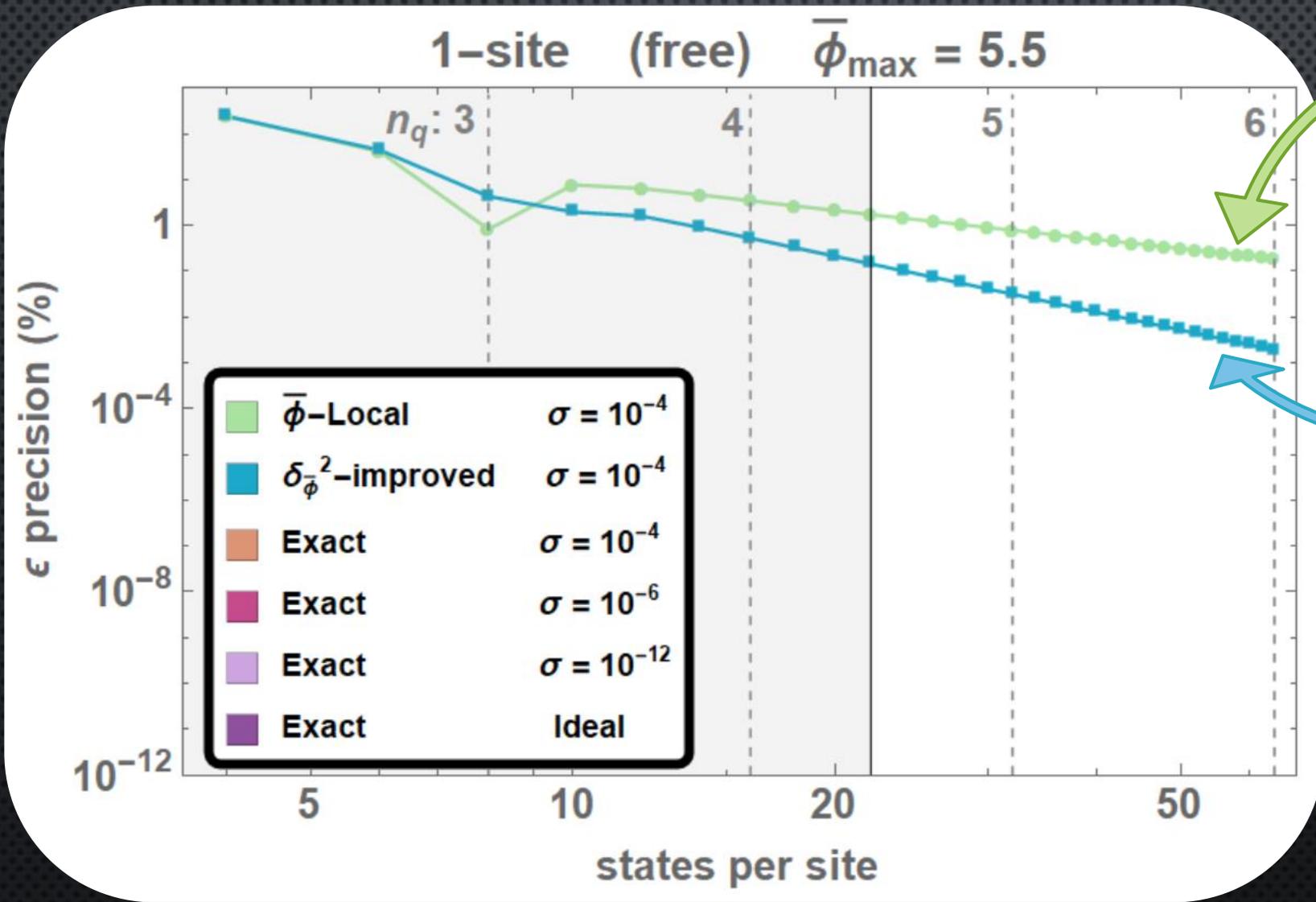


Informs optimal use of Quantum Resources

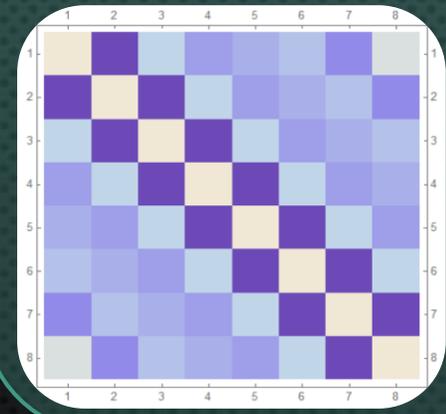
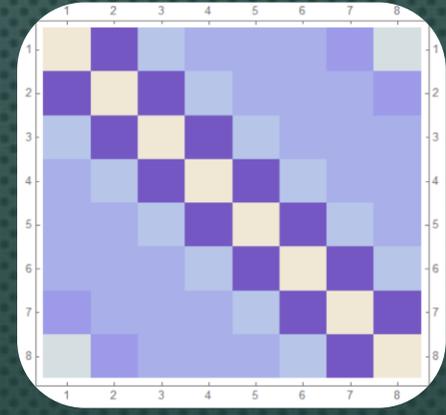
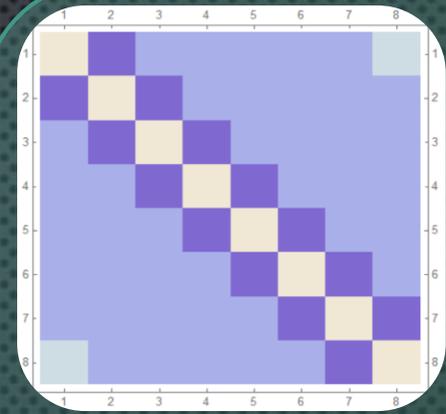
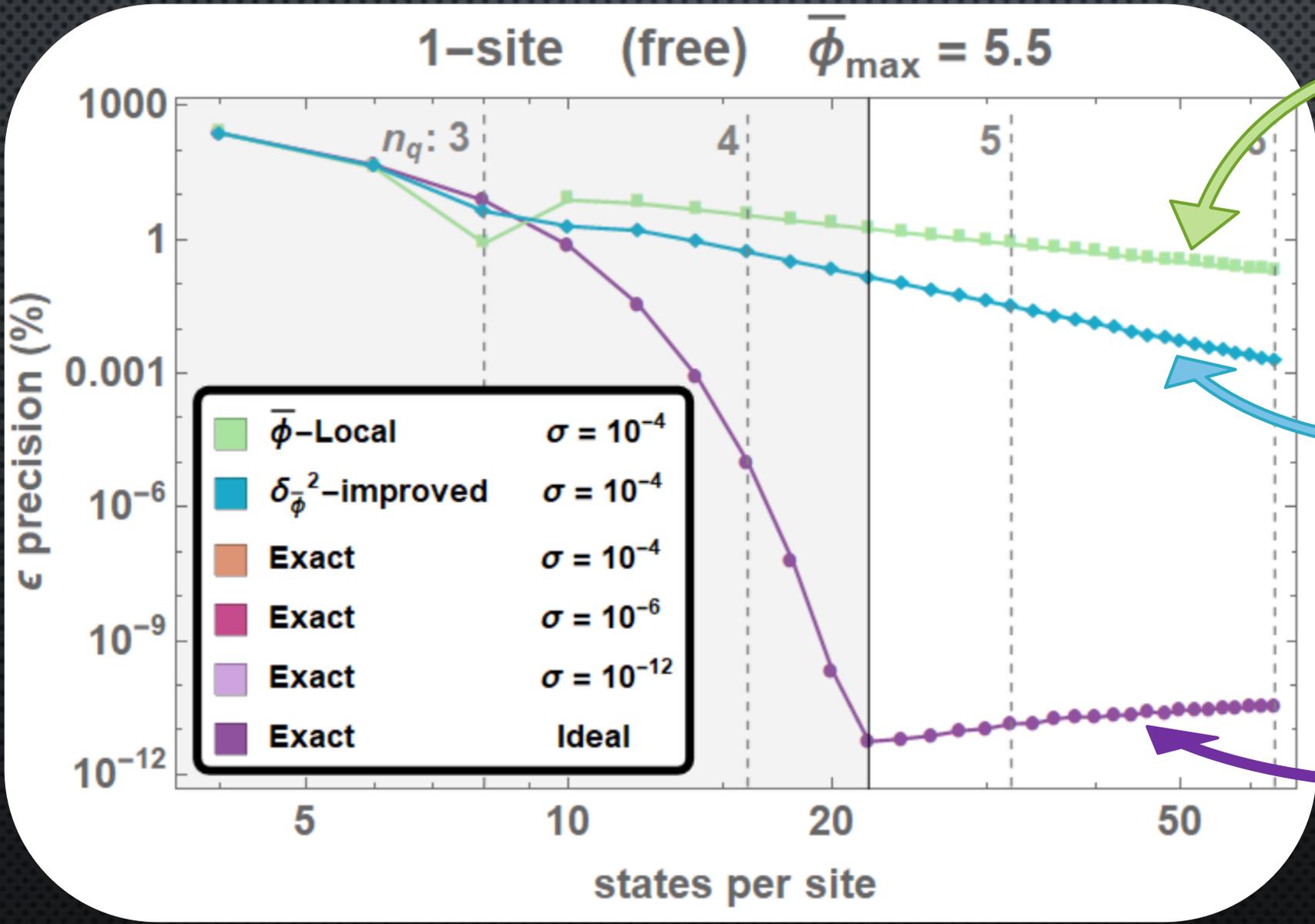
QFT and digitization Improvement



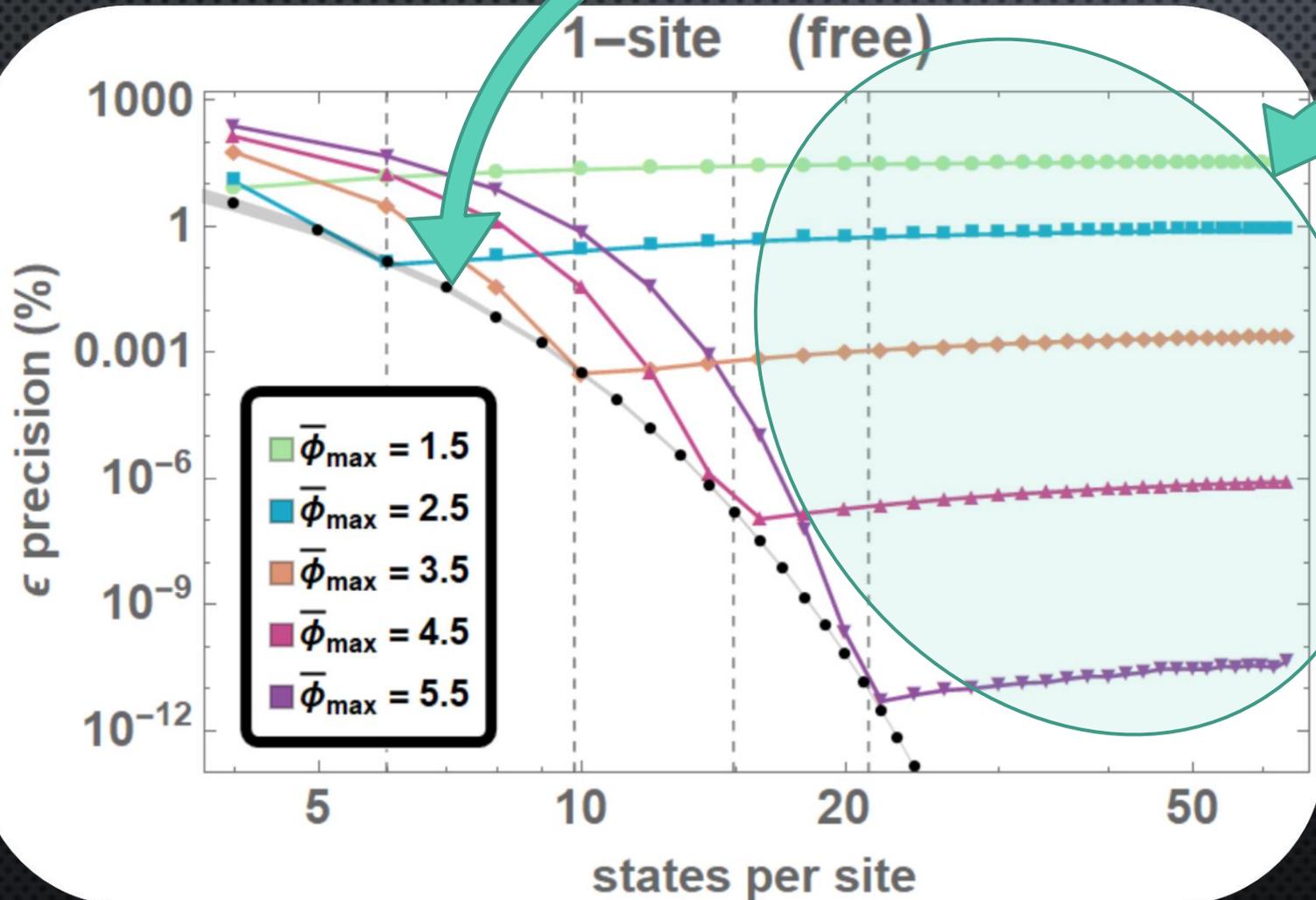
QFT and digitization Improvement



QFT and digitization Improvement



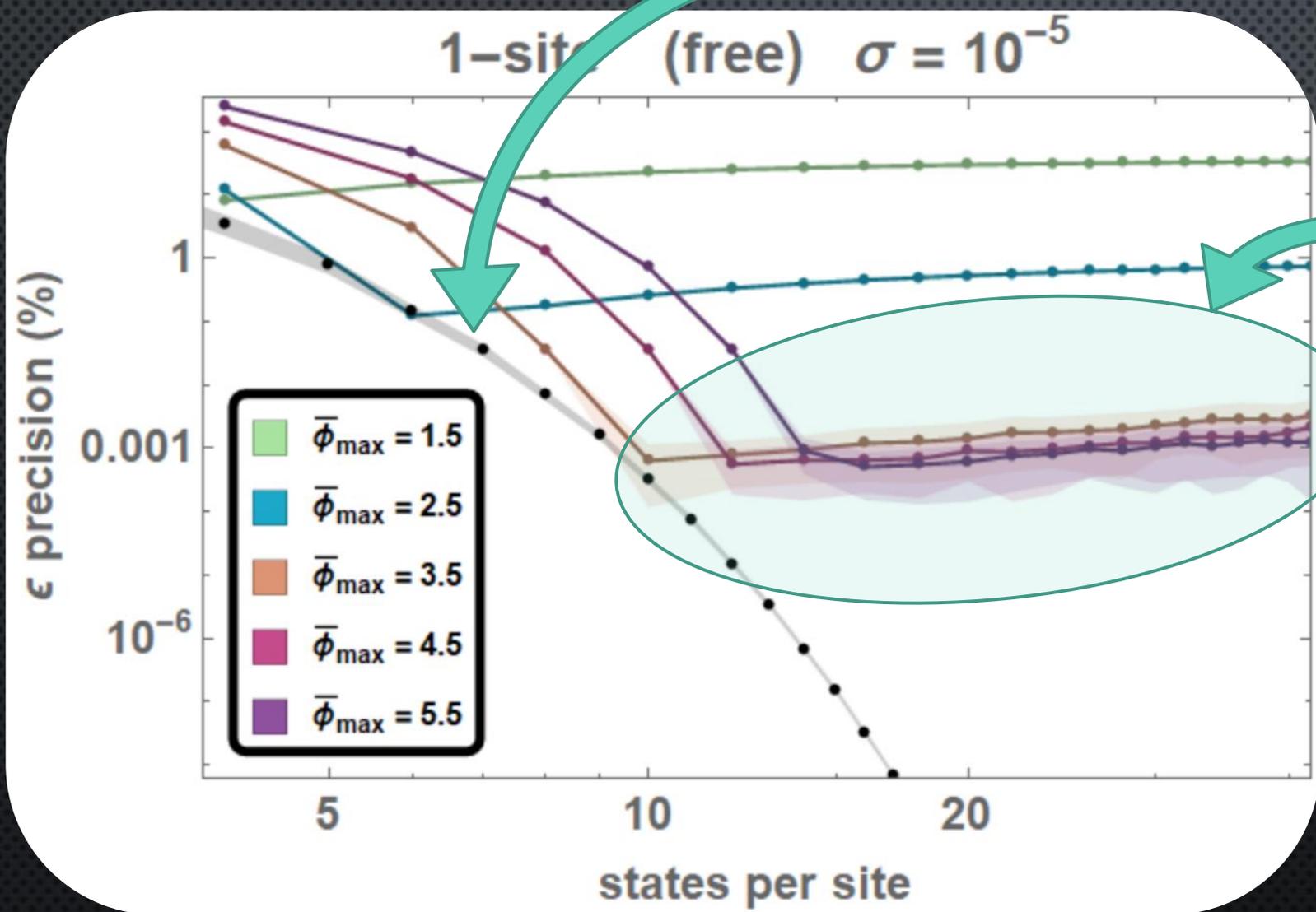
$$\epsilon \sim (1.8(2) \times 10^3) 2^{-2.234(4)2^{n_Q}}$$



Improvements
after NS
saturation are
exponentially
small



$$\epsilon \sim (1.8(2) \times 10^3) 2^{-2.234(4)2^{n_Q}}$$

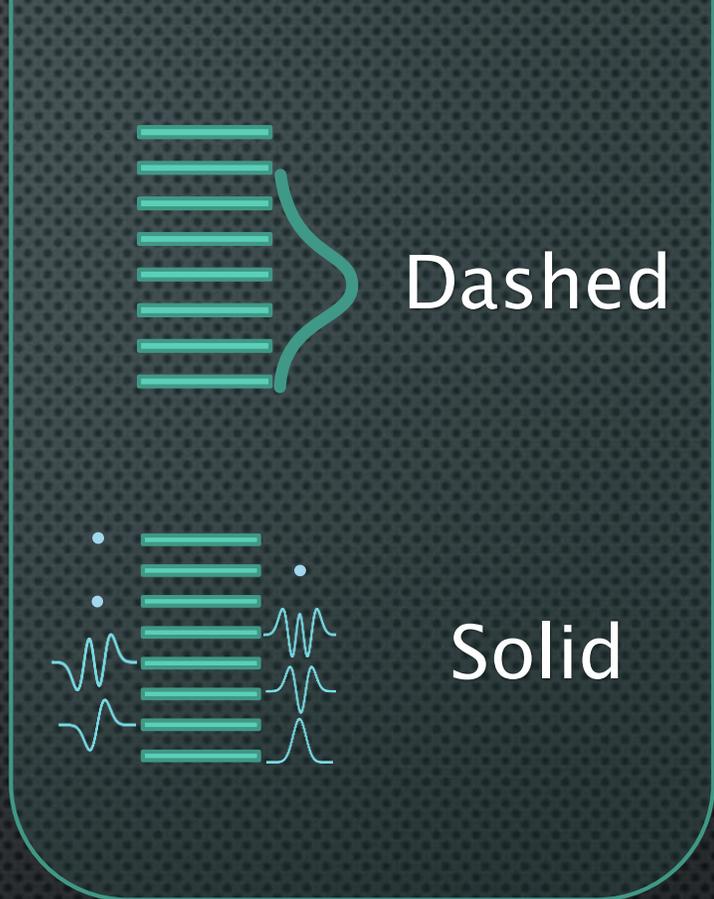
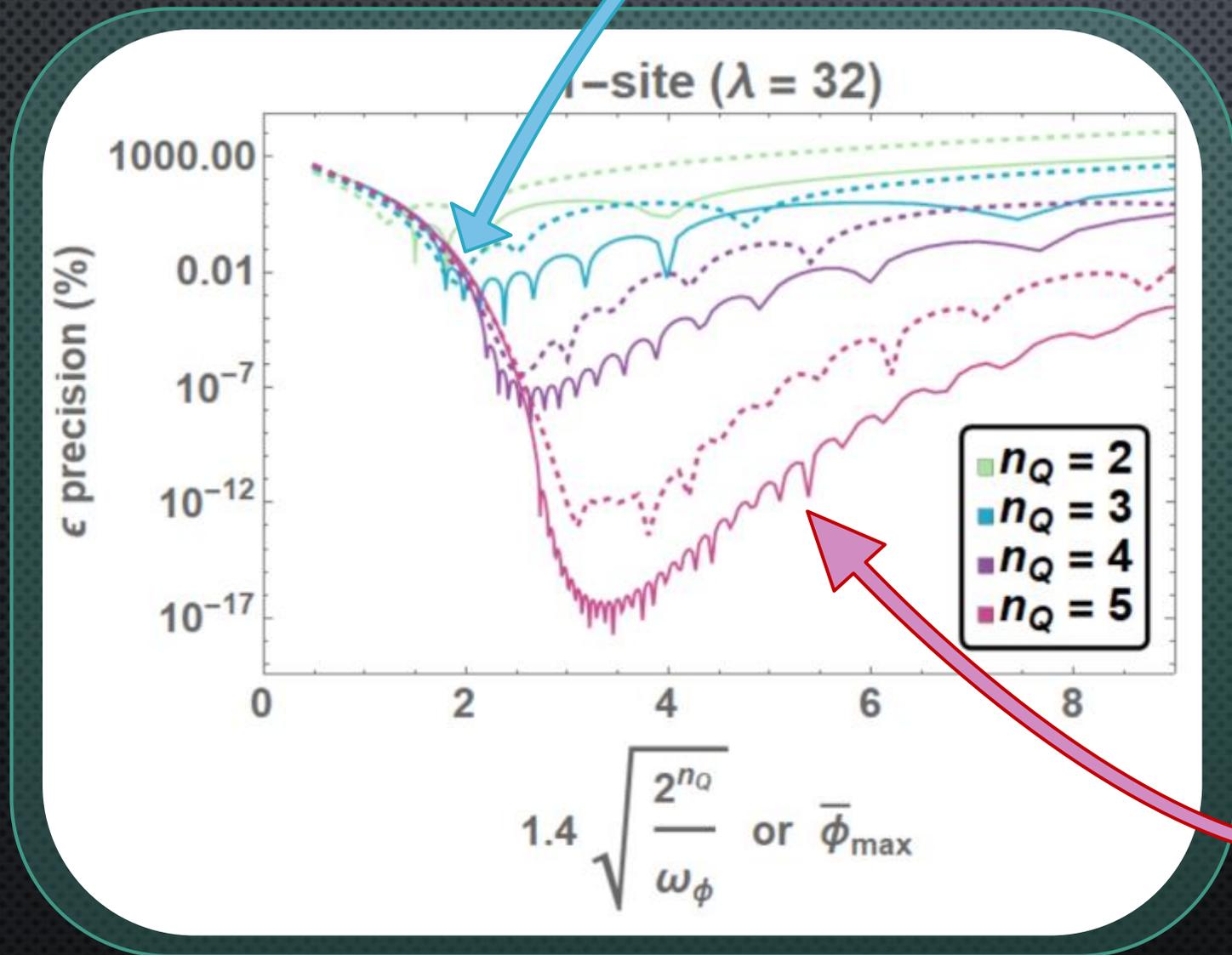


Quantum noise limits achievable precision

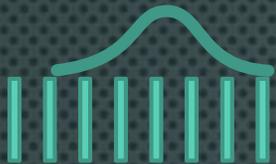


Tuning the basis

Exponentially
converged in n_Q

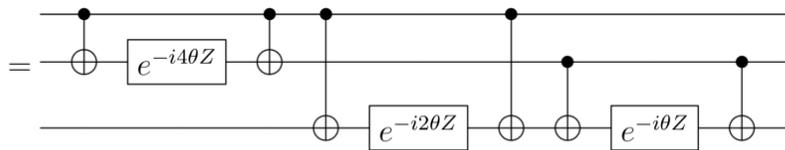


Under-sampling
conjugate-
momentum
space



$$\epsilon_{\text{tuned}}^\lambda \sim 10^{-3}$$

$$\Phi_3(\theta) = e^{-i\theta I \otimes Z \otimes Z} e^{-i2\theta Z \otimes I \otimes Z} e^{-i4\theta Z \otimes Z \otimes I}$$



$$e^{-i\tilde{H}_3 t} = \lim_{M \rightarrow \infty} \left(\Phi_3 \left(\frac{t}{2M}, \frac{4}{49} \bar{\phi}_{max}^2 \right) QFT_{sym} \left(\frac{xvuz\phi_{19} Mz}{z^{26} t} \right) \epsilon_\Phi QFT_{sym}^{-1} \right)^M$$

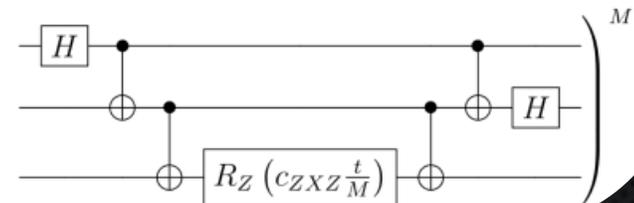
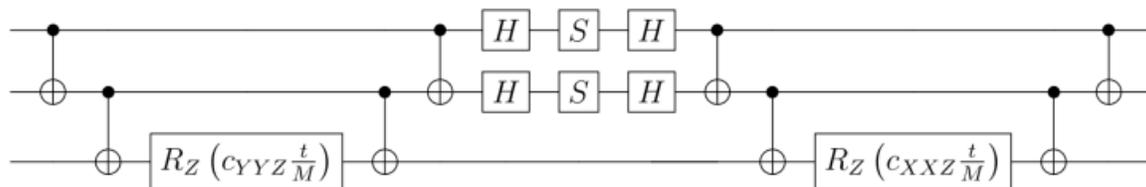
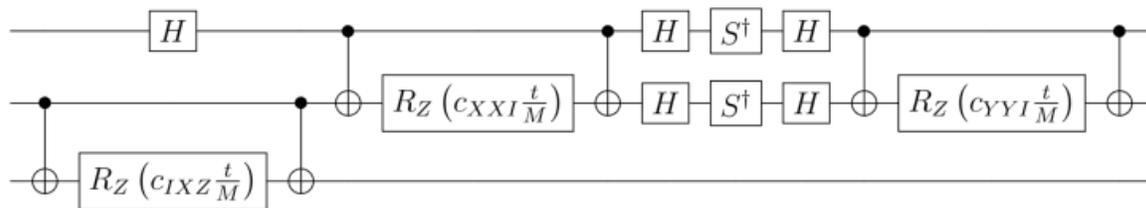
3-qubit real-time evolution



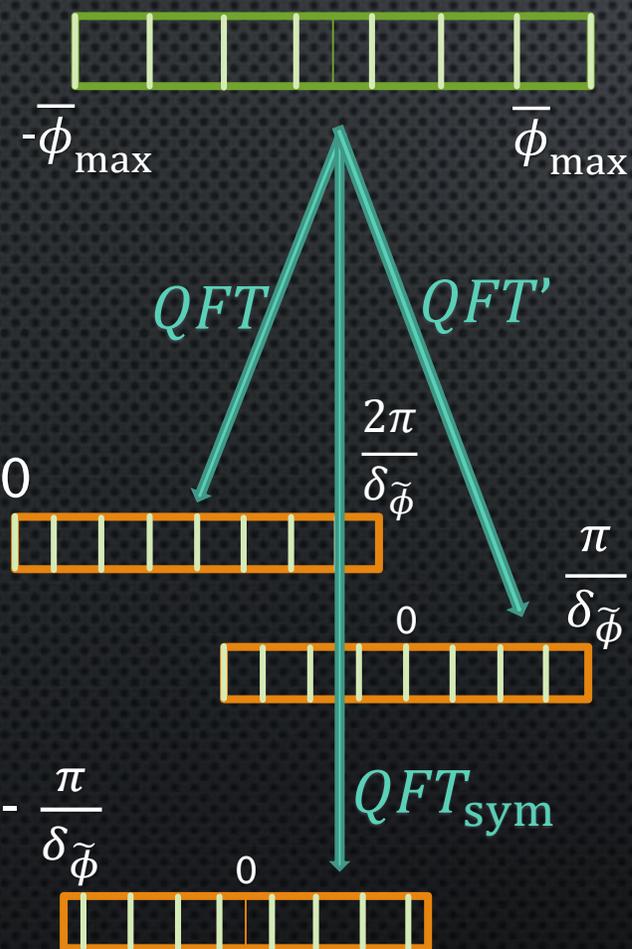
$$e^{-iHt} = \lim_{N \rightarrow \infty} \left(\prod_j e^{-iH_j \frac{t}{N}} \right)^N$$

$$\epsilon_{\text{tuned}}^\lambda \sim 10^{-4}$$

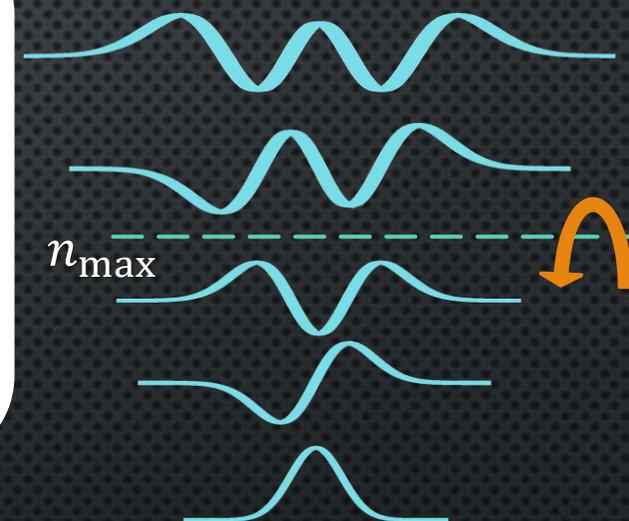
$$e^{-i(H_{\text{basis}} + \delta H_{\omega_\phi})t} = \lim_{M \rightarrow \infty} \left(\begin{array}{l} R_Z(c_{ZII} \frac{t}{M}) \\ R_Z(c_{IZI} \frac{t}{M}) \text{---} H \text{---} R_Z(c_{IXI} \frac{t}{M}) \oplus R_Z(c_{ZXI} \frac{t}{M}) \\ R_Z(c_{IIZ} \frac{t}{M}) \end{array} \right)$$



Circuits, Operators and Gates



Basis	n_Q	0-body	1-body	2-body	3-body	4-body	5-body	6-body	QFT	CNOTs
	2	1	8	2					✓	8
	3	1	14	6					✓	24
JLP	4	1	20	12		1			✓	54
	5	1	26	20		5			✓	110
	6	1	32	30		15			✓	210
JLP	n_Q	1	$4n_Q - 6$	$2 * \binom{n_Q}{2}$		$\binom{n_Q}{4}$			✓	$8 \binom{n_Q}{2} + 6 \binom{n_Q}{4}$
	2	1	3	2						4
	3	1	5	9	4					34
HO	4	1	6	16	18	10				164
	5	1	7	22	32	44	22			612
	6	1	8	29	44	84	98	46		1982



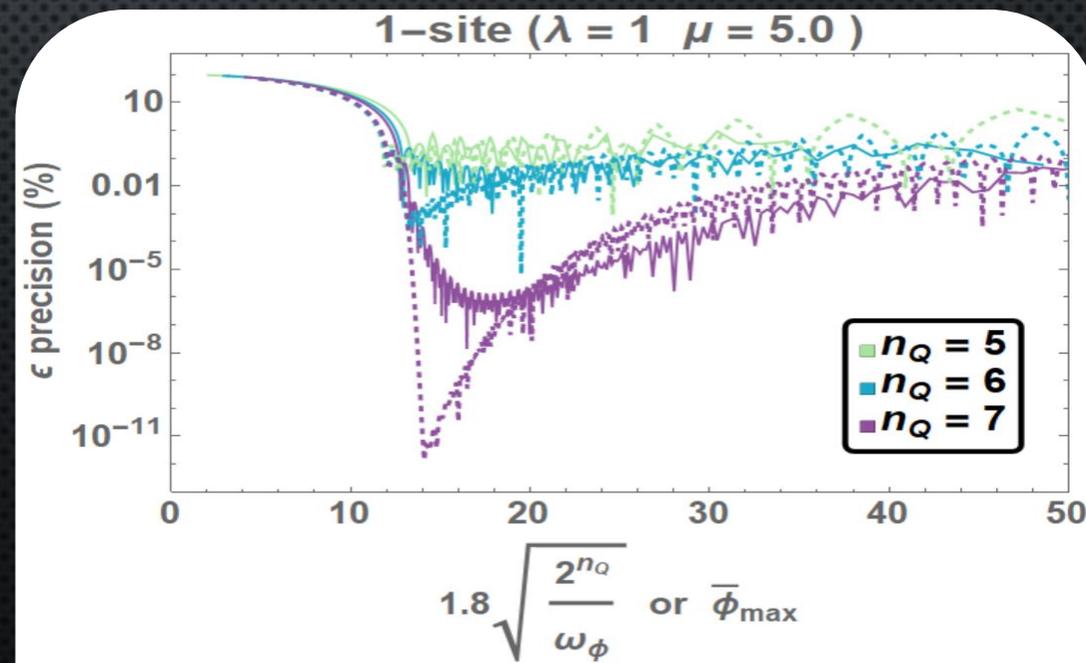
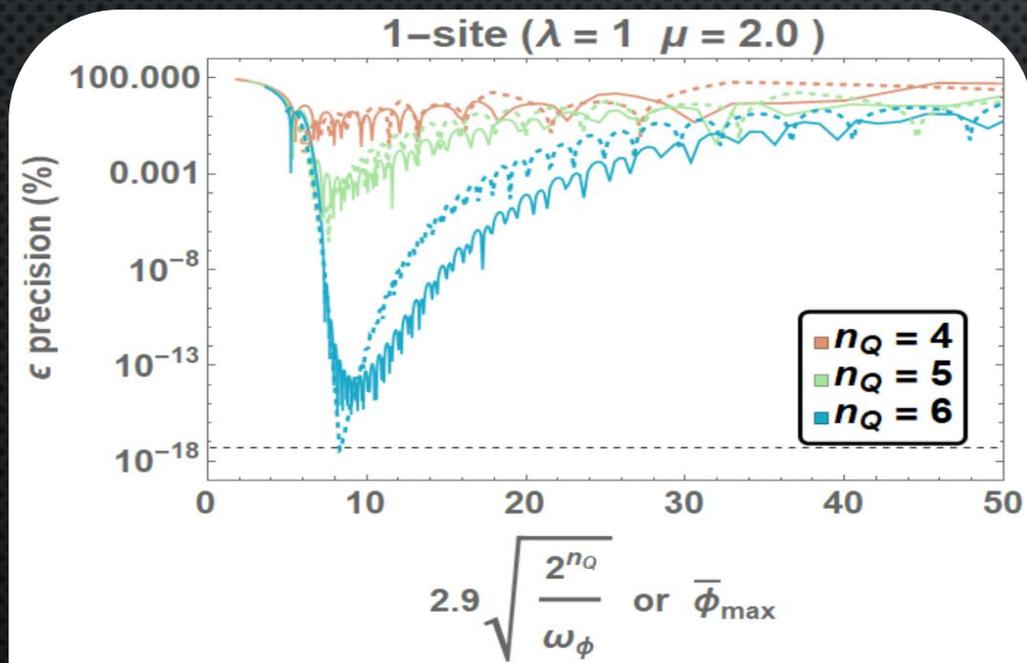
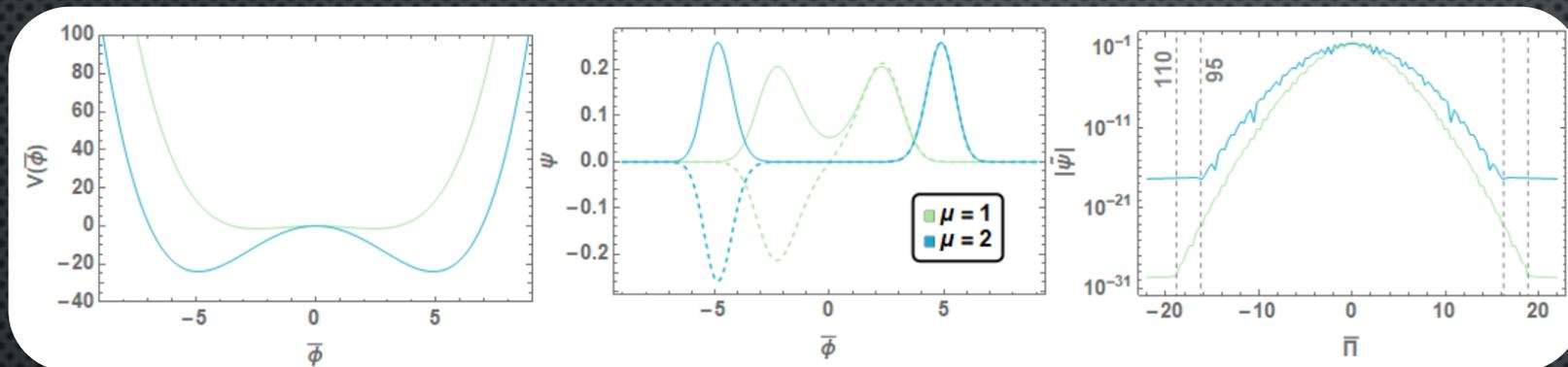
Quantum resources sensitive to truncation details

Delocalized Wavefunctions

$$-\mu^2 = m^2 < 0$$

$$\text{Minima: } \varphi = \pm \frac{\sqrt{3!}\mu}{\sqrt{\lambda}}$$

$$\lim_{\mu \rightarrow \infty} E_1 = E_0$$

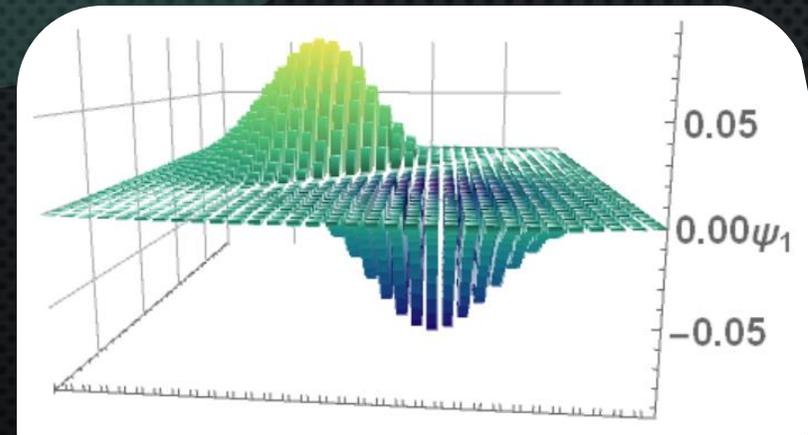
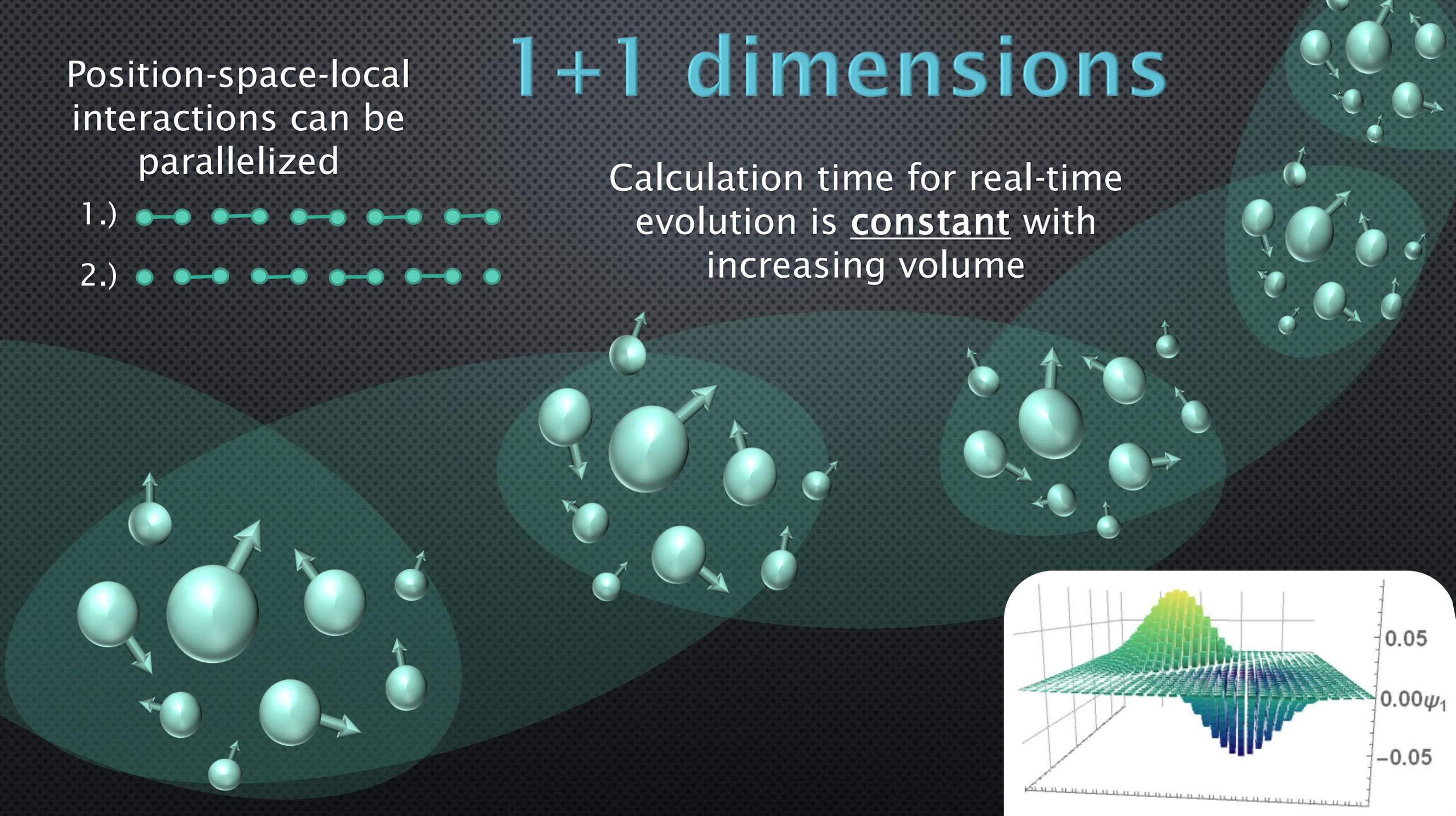


1+1 dimensions

Position-space-local interactions can be parallelized



Calculation time for real-time evolution is constant with increasing volume



Summary

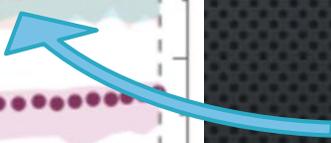
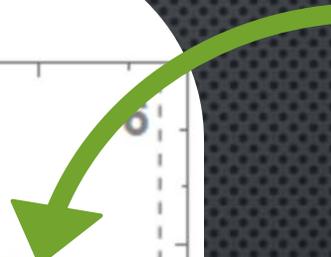
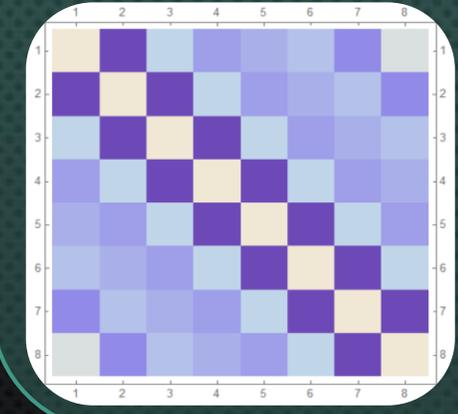
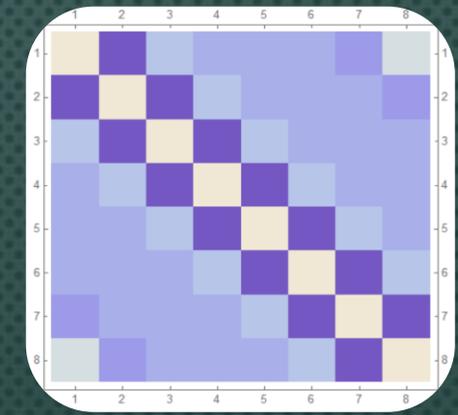
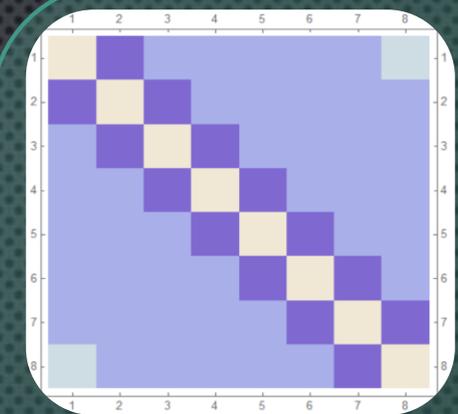
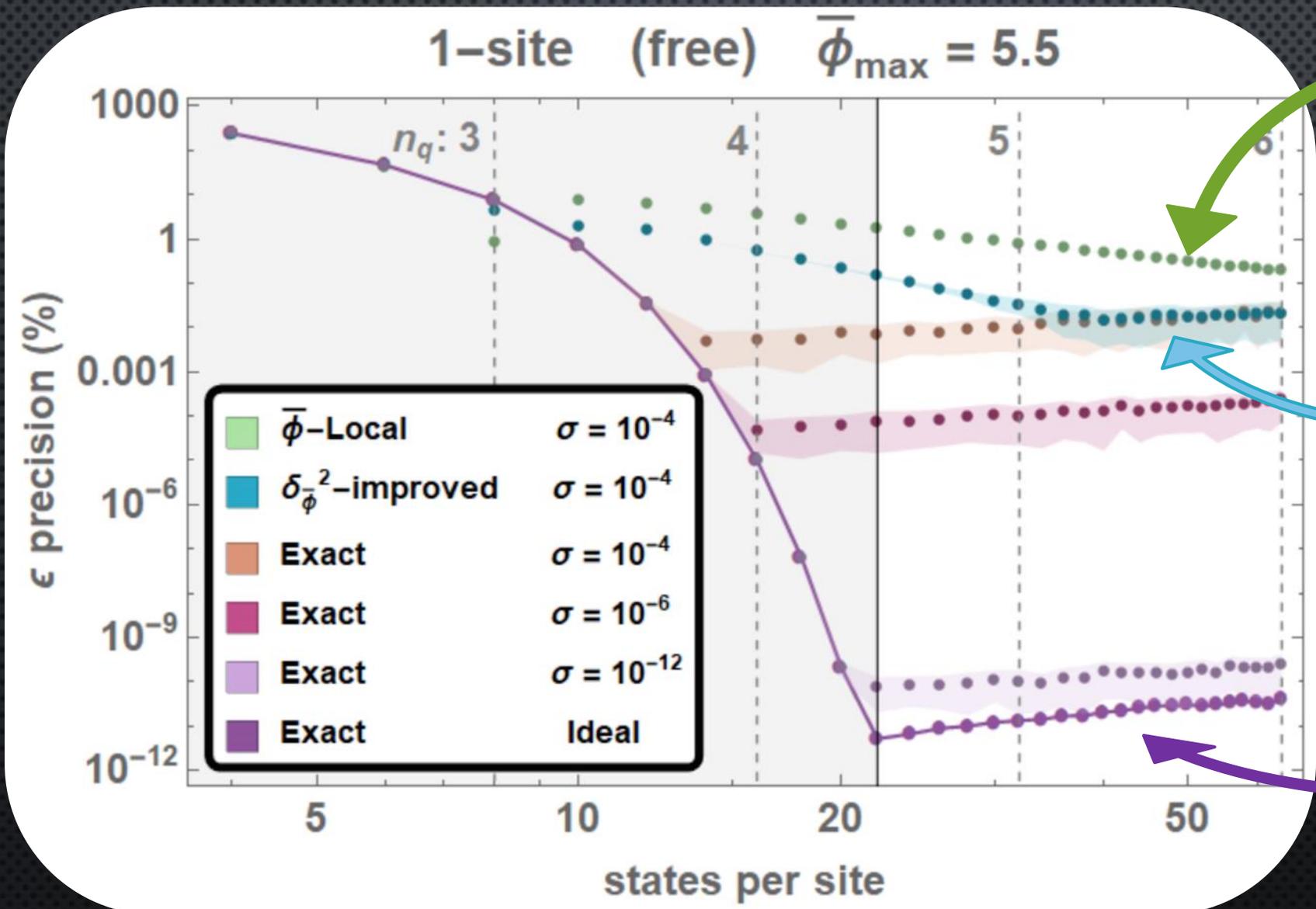
- ❑ Parallel exploration of quantum implementations is vital --- optimal basis may change dramatically with hardware properties.
- ❑ Hardware-specific, informed decisions in formulating calculations for quantum device requires analysis of the multi-dimensional, correlated space of necessary qubits, operators, and gates.
- ❑ Small scalar lattices have communication and digitization requirements amenable to precise calculation with near-term quantum devices.
- ❑ Lessons learned in scalar field implementation will be carried forward for gauge field theories.

UW collaborators:

Silas Beane
David Kaplan
Kenneth Roche
Martin Savage
Jesse Stryker



Introduction of Quantum Noise

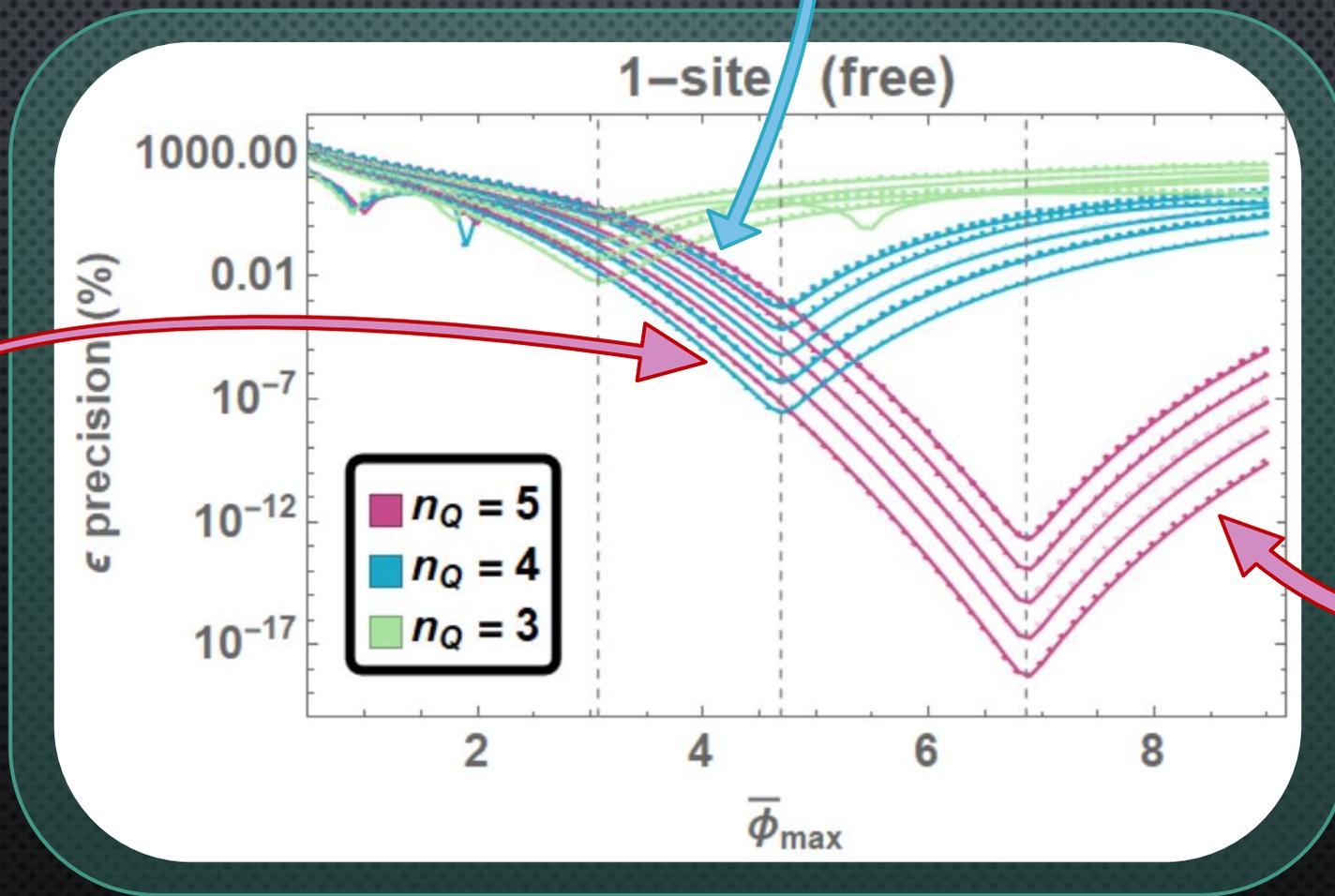


Tuning Basis for Fixed Qubits



Exponentially converged in n_Q

$\bar{\phi}_{\max}$ truncation of field-space wavefunction



Under-sampling conjugate-momentum space

Gradient Operator

Basis	n_Q	2-body	3-body	4-body	5-body	6-body	7-body	8-body	9-body	10-body	11-body	12-body	CNOT
JLP	2	4											8
	3	9											18
	4	16											32
	5	25											50
	6	36											72
	n_Q	n_Q^2											$2n_Q^2$
HO	2	1	6	9									80
	3	1	8	30	56	49							1,152
	4	1	10	47	140	271	330	225					11,264
	5	1	12	68	244	630	1204	1668	1612	961			89,600
	6	1	14	93	392	1186	2772	5154	7560	8541	7182	3969	626,688

Free

Basis	n_Q	0-body	1-body	2-body	3-body	4-body	5-body	6-body	<i>QFT</i>	CNOTs
JLP	2	1	8	2					✓	8
	3	1	14	6					✓	24
	4	1	20	12					✓	48
	5	1	26	20					✓	80
	6	1	32	30					✓	120
JLP	n_Q	1	$6n_Q - 4$	$2 * \binom{n_Q}{2}$					✓	$8 \binom{n_Q}{2}$
$\text{HO}_{\omega \equiv 1}$	2	1	2							0
	3	1	3							0
	4	1	4							0
	5	1	5							0
	6	1	6							0
$\text{HO}_{\omega \equiv 1}$	n_Q	1	n_Q							0
$\text{HO}_{\omega \neq 1}$	2	1	3	1						2
	3	1	4	4	3					20
	4	1	5	5	11	7				96
	5	1	6	6	16	26	15			352
	6	1	7	7	22	42	57	31		1120