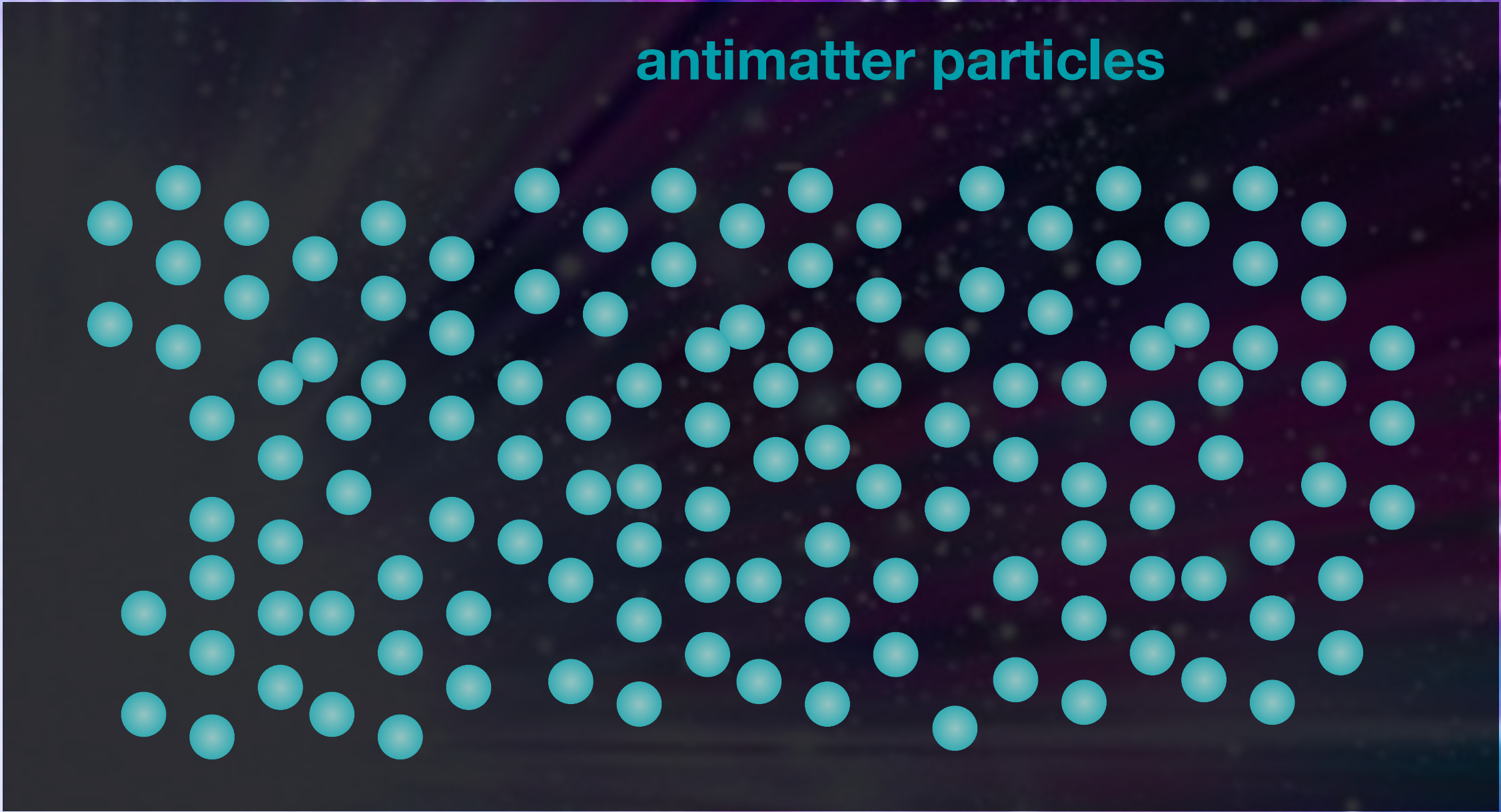
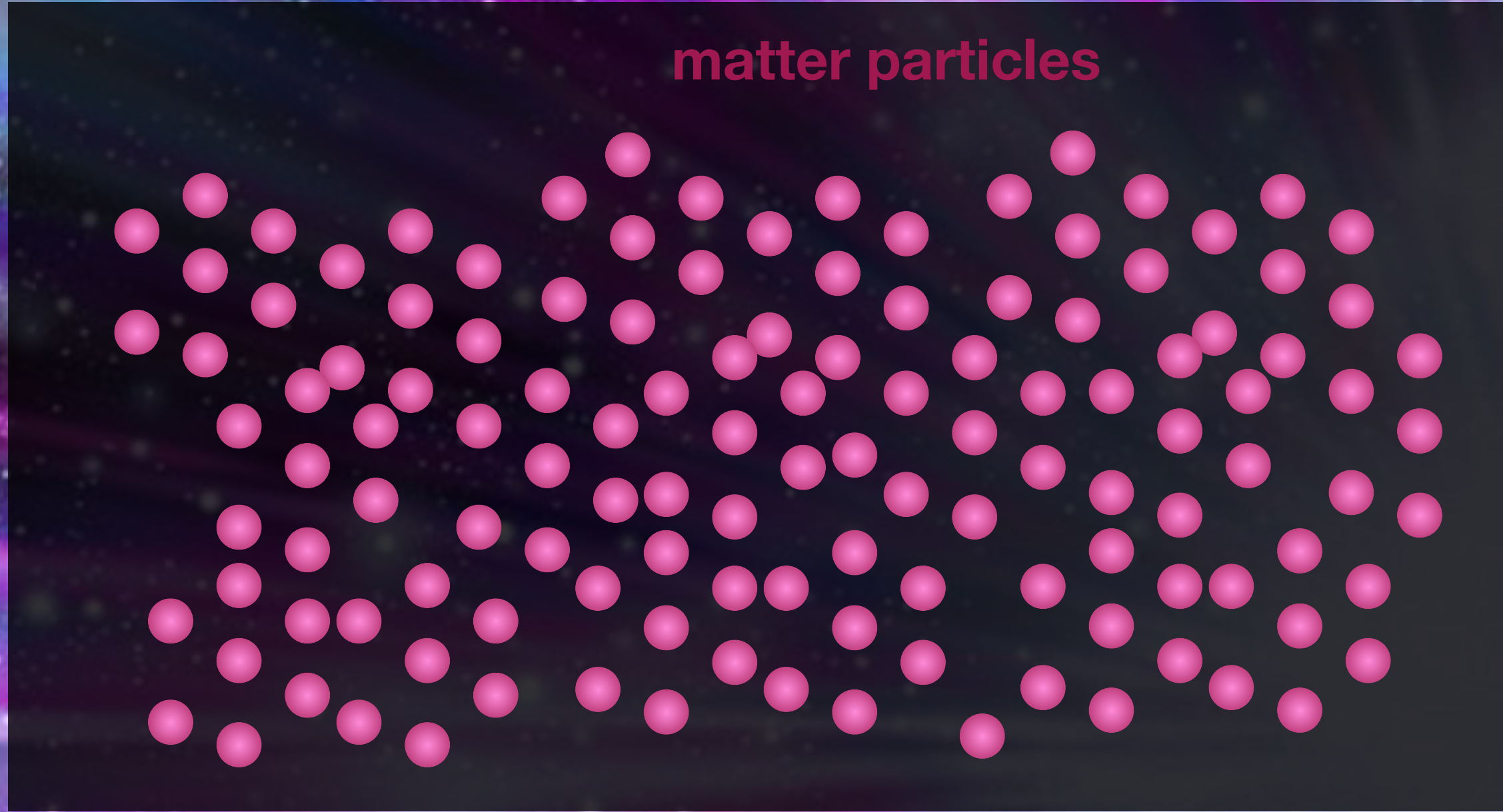


Neutrinoless $\beta\beta$ decay: physics that matters (... but doesn't antimatter)

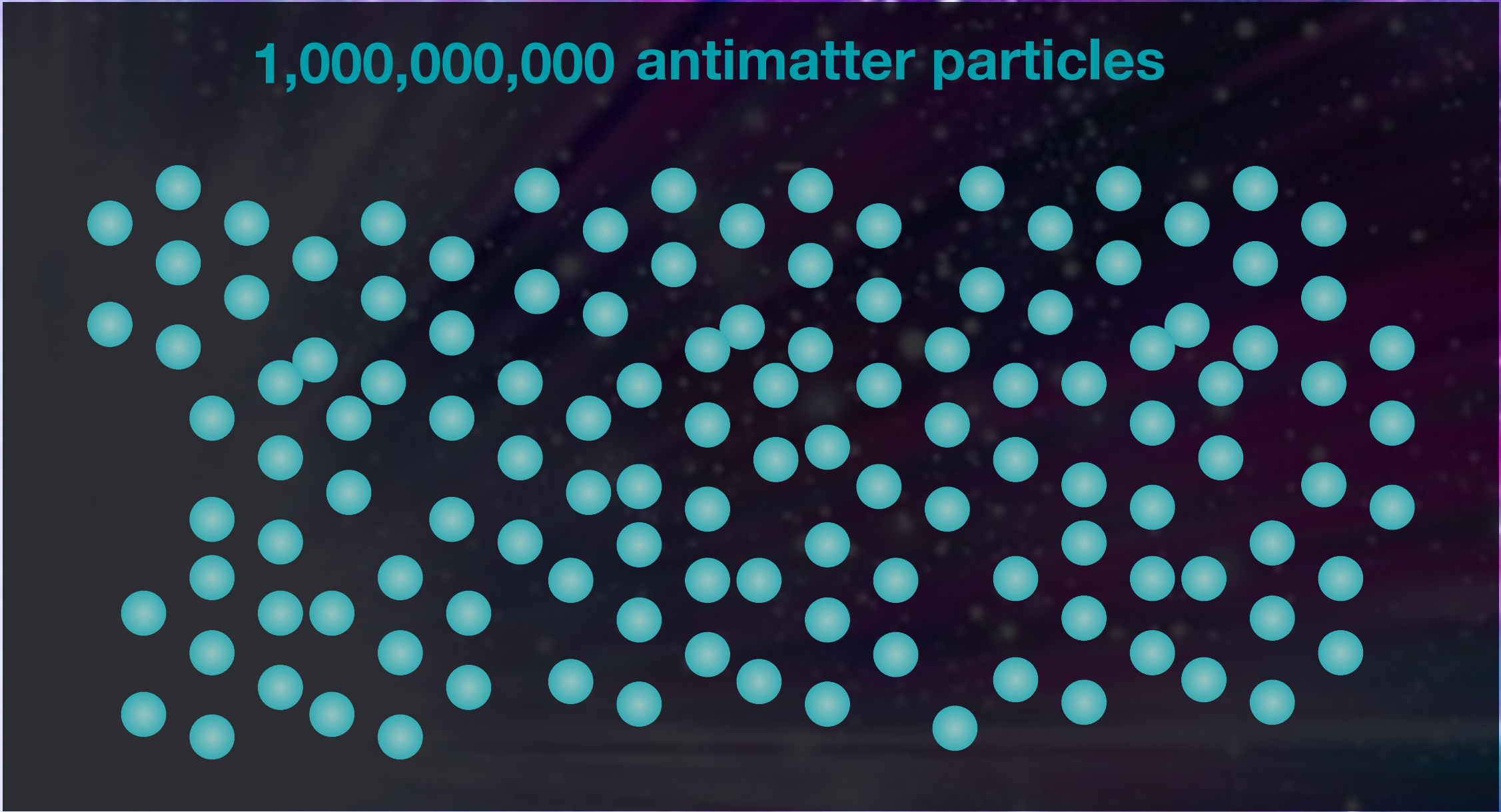
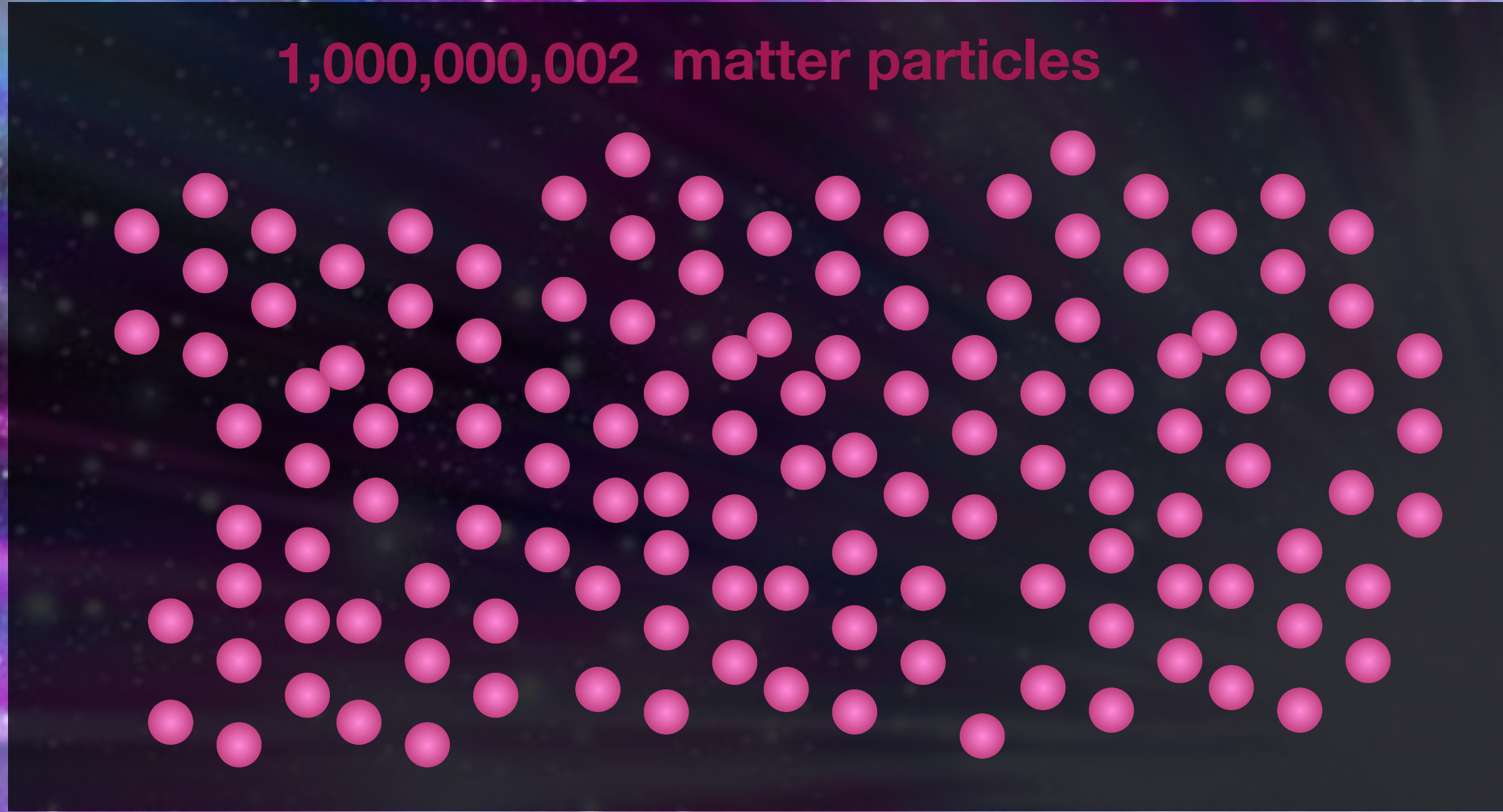
UCL

Cheryl Patrick, University College London
International Neutrino Summer School 2019, Fermilab

In the beginning was...



In the beginning was...



In the beginning was...

2 matter particles



No antimatter particles

2 matter particles



No antimatter particles

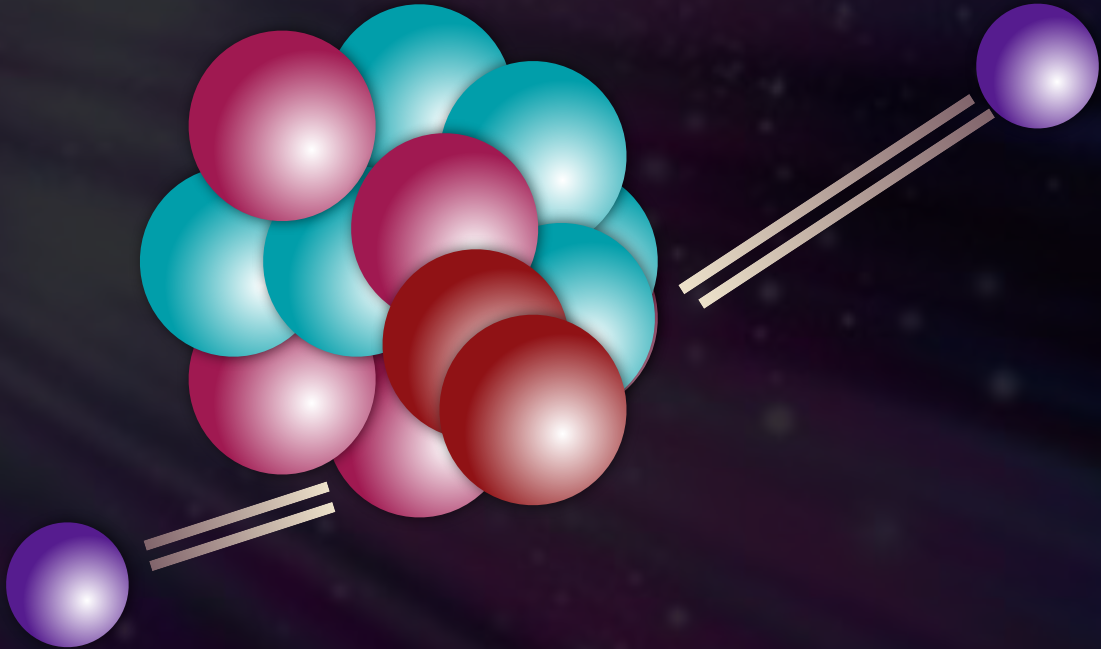
- How did we get this imbalance?

2 matter particles

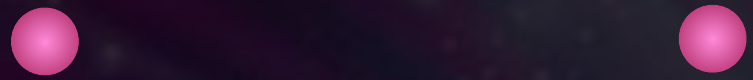


No antimatter particles

- How did we get this imbalance?
- **Neutrinoless double-beta decay** would CREATE matter...



2 matter particles



No antimatter particles

- How did we get this imbalance?
- **Neutrinoless double-beta decay** would CREATE matter...
- ... tells us about the **nature of the neutrino** (that isn't there)...



2 matter particles

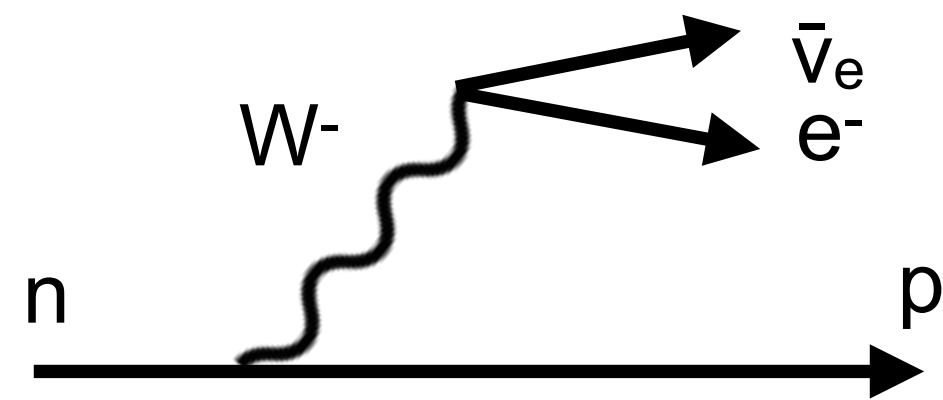
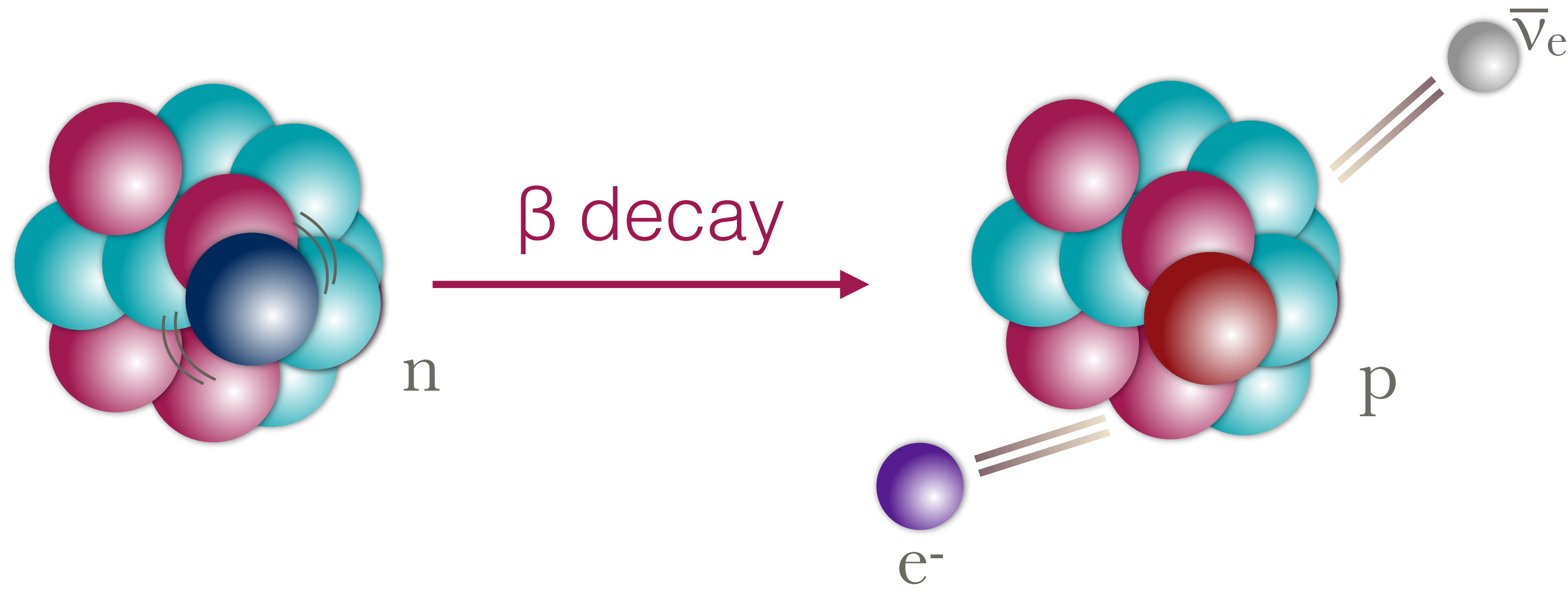


No antimatter particles

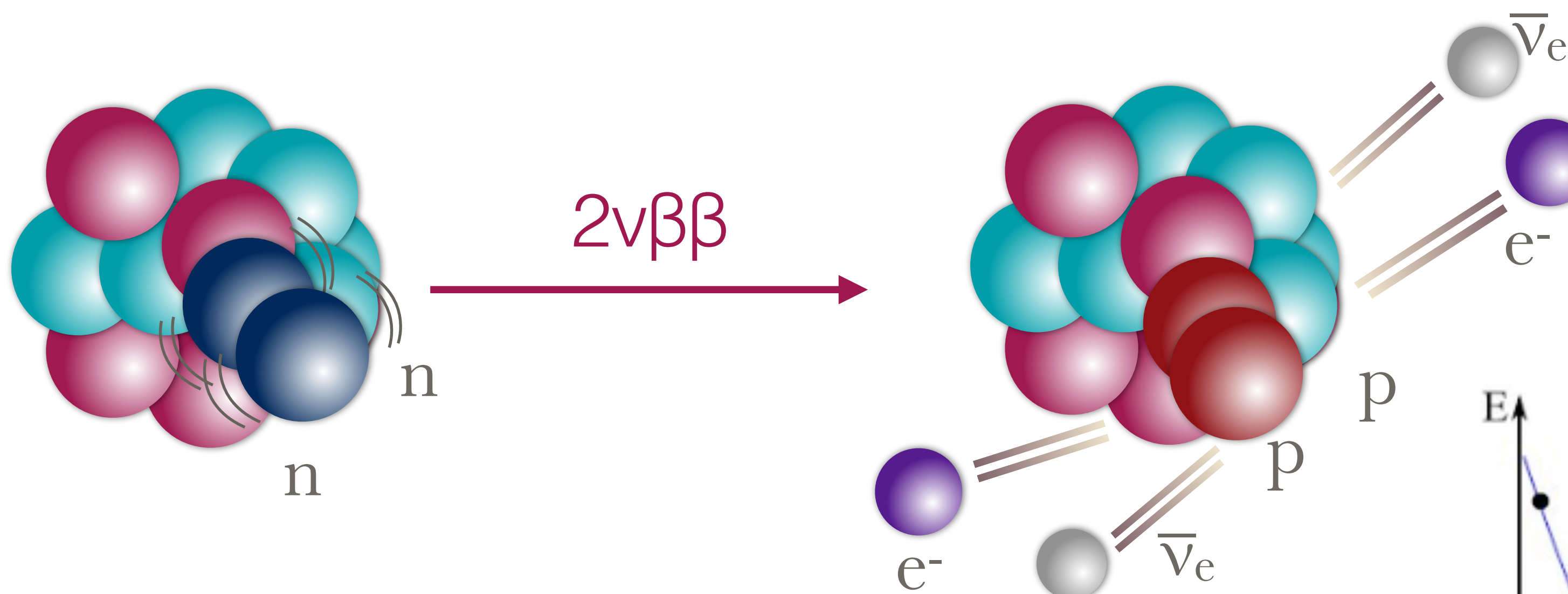
- How did we get this imbalance?
- **Neutrinoless double-beta decay** would **CREATE** matter...
- ... tells us about the **nature of the neutrino** (that isn't there)...
- ... and could help us measure its **mass**



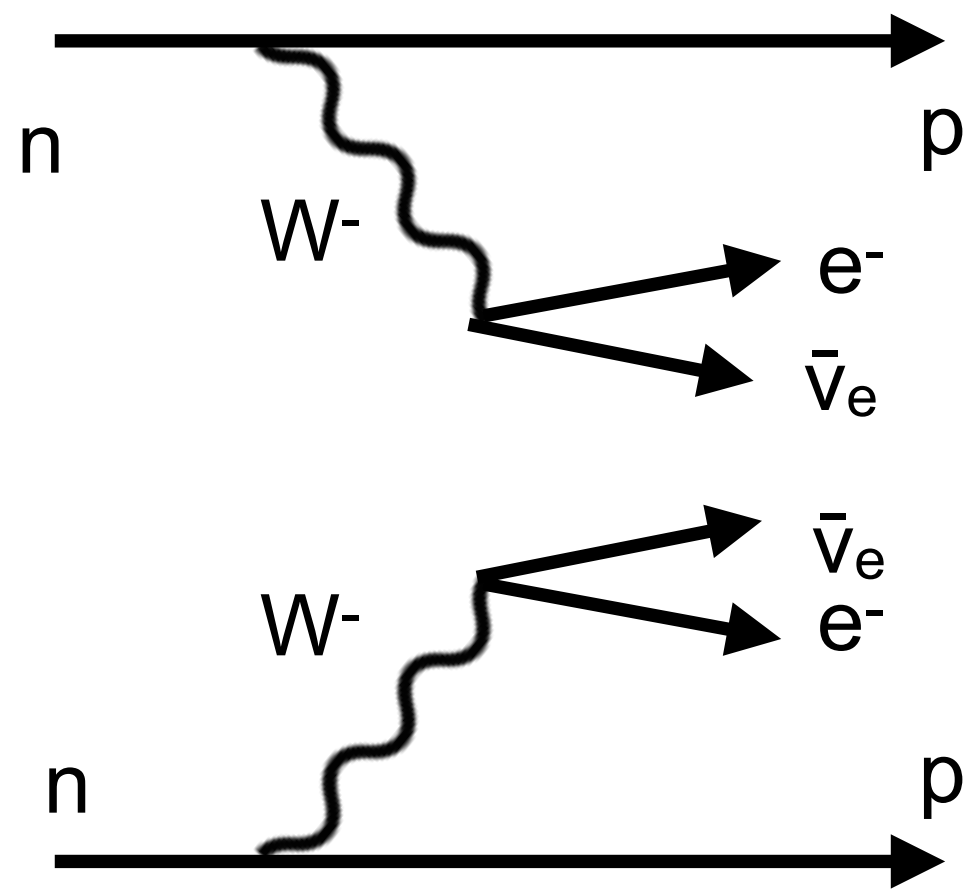
beta decay



double-beta decay



1 H 1.008	2 He 4.0026											13 Al 26.982	14 Si 28.085	15 P 30.974	16 S 32.06	17 Cl 35.45	18 Ar 39.948	
3 Li 6.94	4 Be 9.0122											5 B 10.81	6 C 12.011	7 N 14.007	8 O 15.999	9 F 18.998	10 Ne 20.180	
11 Na 22.990	12 Mg 24.305	3 Sc 44.956	4 Ti 47.867	5 V 50.942	6 Cr 51.996	7 Mn 54.938	8 Fe 55.845	9 Co 58.933	10 Ni 58.693	11 Cu 63.546	12 Zn 65.38	31 Ga 69.723	32 Ge 72.63	33 As 74.922	34 Se 78.971	35 Br 79.904	36 Kr 83.798	
19 K 39.098	48 Ca 40.078	21 Sc 44.956	22 Ti 47.867	23 V 50.942	24 Cr 51.996	25 Mn 54.938	26 Fe 55.845	27 Co 58.933	28 Ni 58.693	29 Cu 63.546	30 Zn 65.38	49 In 114.82	50 Sn 118.710	51 Sb 121.76	52 Te 127.60	53 I 126.90	54 Xe 131.29	
37 Rb 85.468	38 Sr 87.62	39 Y 88.906	96 Zr 91.224	41 Nb 92.906	42 Mo 95.94	43 Tc (98)	44 Ru 101.07	45 Rh 102.91	46 Pd 106.32	47 Ag 107.87	48 Cd 112.411	80 Hg 200.59	81 Tl 204.38	82 Pb 207.2	83 Bi 208.98	84 Po (209)	85 At (210)	86 Rn (222)
55 Cs 132.91	56 Ba 137.33	57-71 * Lanthanide series	72 Hf 178.49	73 Ta 180.95	74 W 183.84	75 Re 186.21	76 Os 190.23	77 Ir 192.22	78 Pt 195.08	79 Au 196.97	80 Hg 200.59	81 Tl 204.38	82 Pb 207.2	83 Bi 208.98	84 Po (209)	85 At (210)	86 Rn (222)	
87 Fr (223)	88 Ra (226)	89-103 * Actinide series	104 Rf (261)	105 Db (262)	106 Sg (263)	107 Bh (264)	108 Hs (265)	109 Mt (266)	110 Ds (271)	111 Rg (272)	112 Cn (285)	113 Nh (286)	114 Fl (289)	115 Mc (290)	116 Lv (293)	117 Ts (294)	118 Og (294)	
			57 La 138.91	58 Ce 140.12	59 Pr 140.91	60 Nd 144.24	61 Pm (145)	62 Sm 150.36	63 Eu 151.96	64 Gd 157.25	65 Tb 158.93	66 Dy 162.50	67 Ho 164.93	68 Er 167.26	69 Tm 168.93	70 Yb 173.05	71 Lu 174.97	
			89 Ac (227)	90 Th 232.04	91 Pa 231.04	92 U 238.03	93 Np (237)	94 Pu (244)	95 Am (243)	96 Cm (247)	97 Bk (247)	98 Cf (251)	99 Es (252)	100 Fm (257)	101 Md (258)	102 No (259)	103 Lr (262)	



Lepton number conserved
 $\Delta L = 0$
 $T_{1/2} \sim 10^{20}$ years

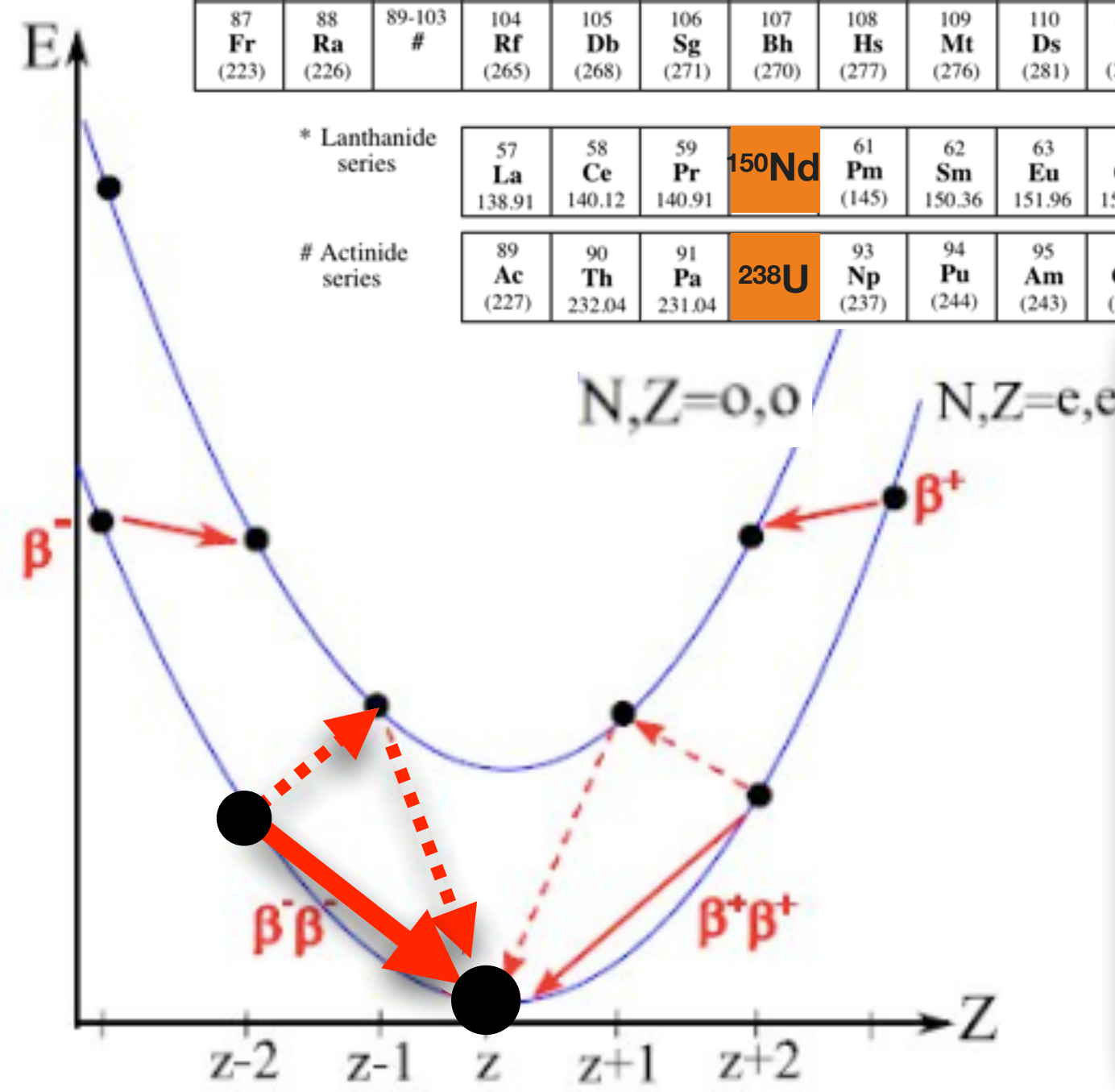
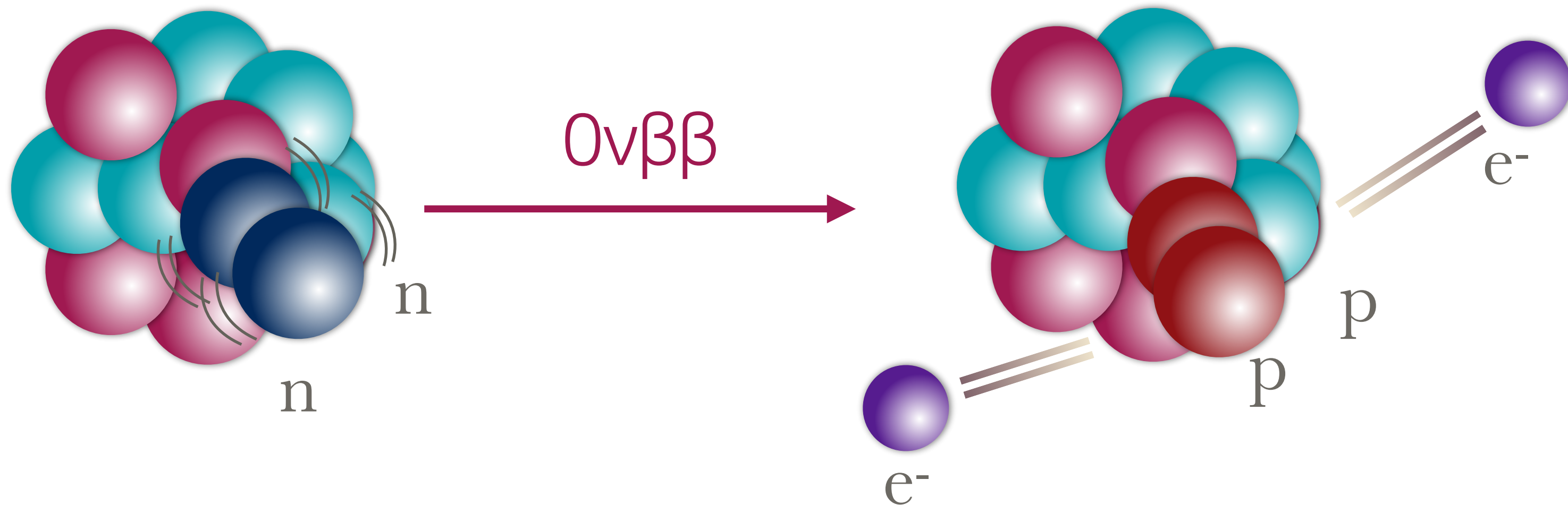


Diagram - E. Falk

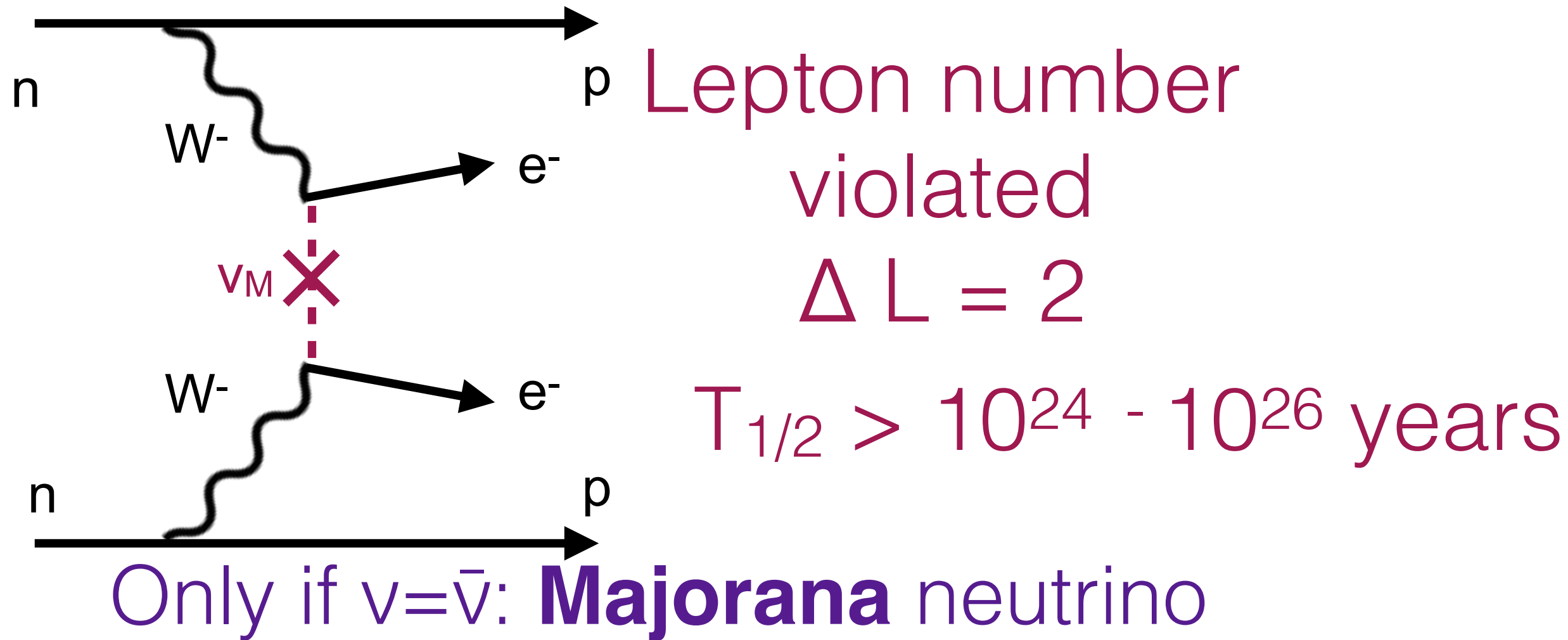


Maria Göppert Mayer

Neutrinoless double-beta decay : the smoking gun for Majorana



Wendell Furry



$0\nu\beta\beta$ could tell us about neutrino mass

$0\nu\beta\beta$ rate

$$\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z) g_A^4 |M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2} \longleftarrow \text{Neutrino mass}$$

$0\nu\beta\beta$ could tell us about neutrino mass

$0\nu\beta\beta$ rate $\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z)g_A^4|M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$ ← Neutrino mass

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

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PMNS mixing angles

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

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$0\nu\beta\beta$ rate $\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z)g_A^4|M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$ ← **Neutrino mass**

PMNS mixing angles

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Individual neutrino masses

$0\nu\beta\beta$ could tell us about neutrino mass

$0\nu\beta\beta$ rate $\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z)g_A^4|M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$ ← **Neutrino mass**

PMNS mixing angles

Majorana phases

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$$U = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha_1/2} & 0 \\ 0 & 0 & e^{i\alpha_2/2} \end{bmatrix}$$

PMNS mixing angles

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PMNS mixing angles

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$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

Individual neutrino masses

$$m_{\beta\beta} = \underbrace{c_{12}^2 c_{13}^2}_{0.68} m_{\nu_1} + \underbrace{s_{12}^2 c_{13}^2}_{0.30} m_{\nu_2} e^{i\alpha_1} + \underbrace{s_{13}^2}_{0.02} m_{\nu_3} e^{i\alpha_3}$$

$0\nu\beta\beta$ could tell us about neutrino mass

$0\nu\beta\beta$ rate $\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z)g_A^4|M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$ ← **Neutrino mass**

Beta decay endpoint

$$m_{\bar{\nu}_e} =$$

Neutrinoless double beta decay

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

$$m_{\beta\beta} = \underbrace{c_{12}^2 c_{13}^2}_{0.68} m_{\nu_1} + \underbrace{s_{12}^2 c_{13}^2}_{0.30} m_{\nu_2} e^{i\alpha_1} + \underbrace{s_{13}^2}_{0.02} m_{\nu_3} e^{i\alpha_3}$$

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Beta decay endpoint

$$m_{\bar{\nu}_e} = \sqrt{\sum_{i=1}^3 |U_{ei}|^2 m_i^2}$$

Neutrinoless double beta decay

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Beta decay endpoint

Incoherent sum

$$m_{\bar{\nu}_e} = \sqrt{\sum_{i=1}^3 |U_{ei}|^2 m_i^2}$$

Neutrinoless double beta decay

Coherent sum

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

$$m_{\beta\beta} = \underbrace{c_{12}^2 c_{13}^2}_{0.68} m_{\nu_1} + \underbrace{s_{12}^2 c_{13}^2}_{0.30} m_{\nu_2} e^{i\alpha_1} + \underbrace{s_{13}^2}_{0.02} m_{\nu_3} e^{i\alpha_3}$$

$0\nu\beta\beta$ could tell us about neutrino mass

$0\nu\beta\beta$ rate $\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z)g_A^4|M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$ ← **Neutrino mass**

Beta decay endpoint

$$m_{\bar{\nu}_e} = \sqrt{\sum_{i=1}^3 |U_{ei}|^2 m_i^2}$$

Is there a lower bound?

Neutrinoless double beta decay

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

$$m_{\beta\beta} = \underbrace{c_{12}^2 c_{13}^2}_{0.68} m_{\nu_1} + \underbrace{s_{12}^2 c_{13}^2}_{0.30} m_{\nu_2} e^{i\alpha_1} + \underbrace{s_{13}^2}_{0.02} m_{\nu_3} e^{i\alpha_3}$$

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Beta decay endpoint

$$m_{\bar{\nu}_e} = \sqrt{\sum_{i=1}^3 |U_{ei}|^2 m_i^2}$$

Neutrinoless double beta decay

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$0\nu\beta\beta$ rate $\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z)g_A^4|M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$ ← **Neutrino mass**

Beta decay endpoint

$$m_{\bar{\nu}_e} = \sqrt{\sum_{i=1}^3 |U_{ei}|^2 m_i^2}$$

Neutrinoless double beta decay

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

Is there a lower bound?

It's complicated...

$$m_{\beta\beta} = \underbrace{c_{12}^2 c_{13}^2}_{0.68} m_{\nu_1} + \underbrace{s_{12}^2 c_{13}^2}_{0.30} m_{\nu_2} e^{i\alpha_1} + \underbrace{s_{13}^2}_{0.02} m_{\nu_3} e^{i\alpha_3}$$

$0\nu\beta\beta$ could tell us about neutrino mass

$0\nu\beta\beta$ rate $\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z)g_A^4|M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$ ← **Neutrino mass**

Beta decay endpoint

$$m_{\bar{\nu}_e} = \sqrt{\sum_{i=1}^3 |U_{ei}|^2 m_i^2}$$

Neutrinoless double beta decay

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

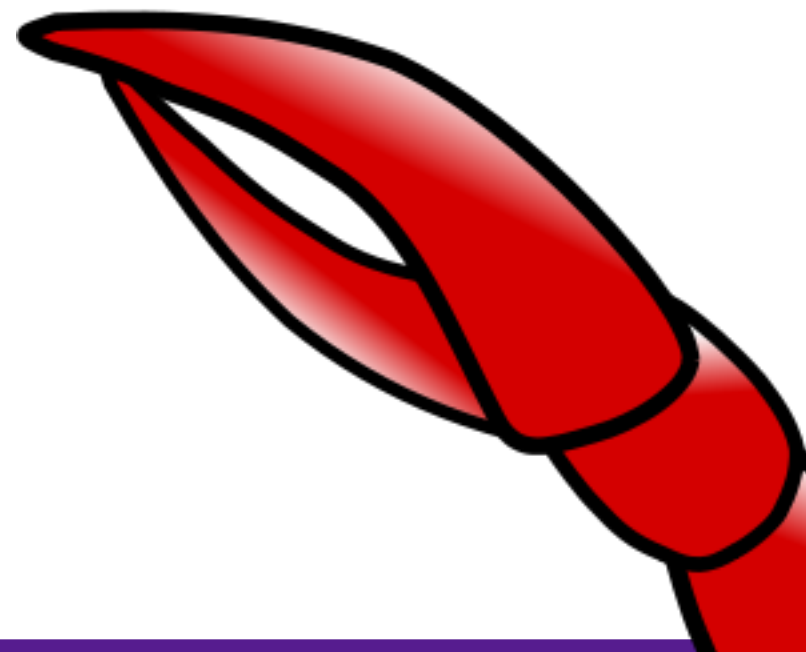
Is there a lower bound?

It's ~~complicated...~~ complex

$$m_{\beta\beta} = \underbrace{c_{12}^2 c_{13}^2}_{0.68} m_{\nu_1} + \underbrace{s_{12}^2 c_{13}^2}_{0.30} m_{\nu_2} e^{i\alpha_1} + \underbrace{s_{13}^2}_{0.02} m_{\nu_3} e^{i\alpha_3}$$

All about $m_{\beta\beta}$ featuring: the double-beta decay lobster!

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$



All about $m_{\beta\beta}$ featuring: the double-beta decay lobster!

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

Real number between 0 and 1

All about $m_{\beta\beta}$ featuring: the double-beta decay lobster!

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

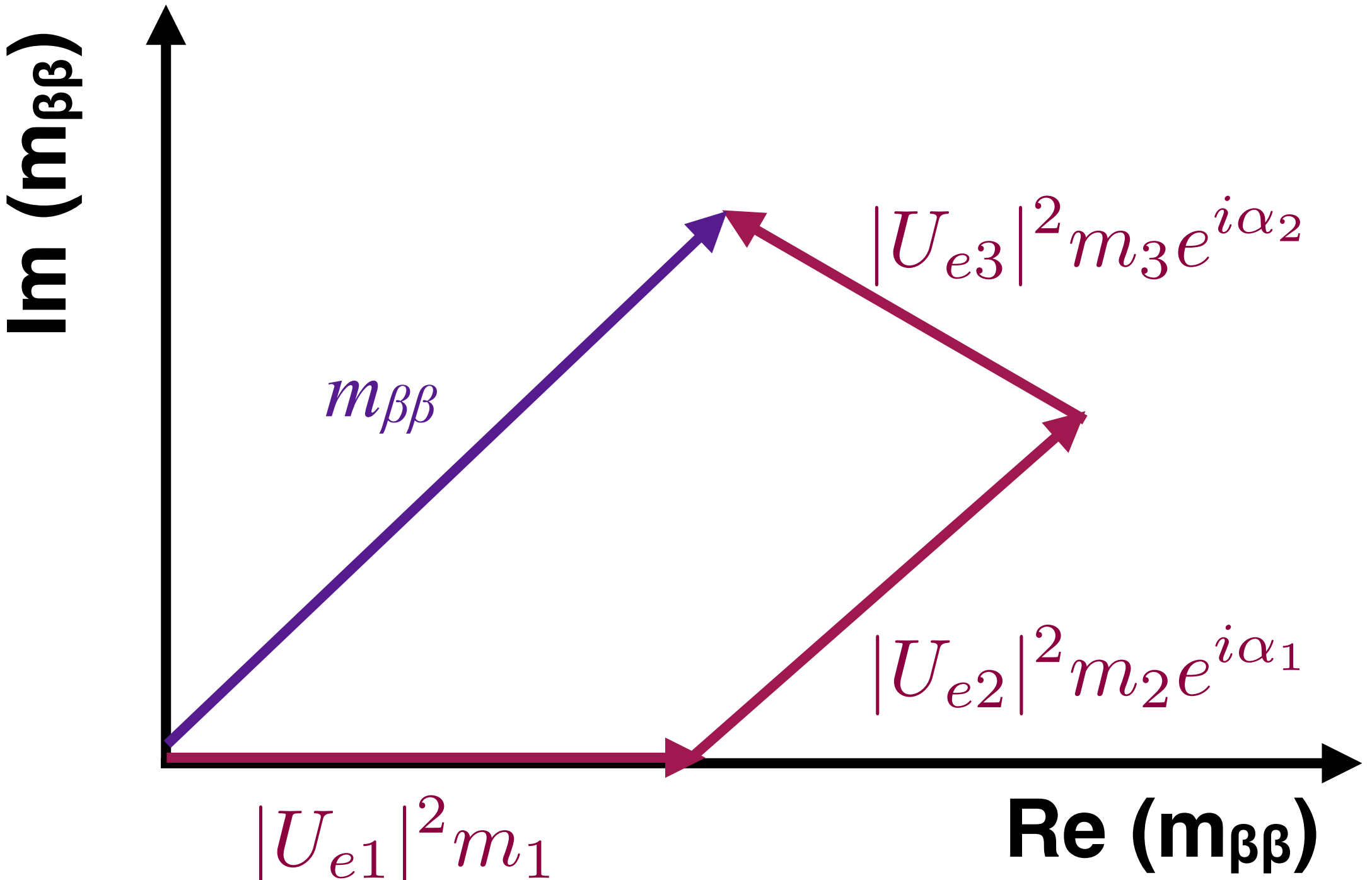
Complex phases



All about $m_{\beta\beta}$ featuring: the double-beta decay lobster!

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

Complex phases



No Majorana phase for m_1 component
(choice in PMNS matrix)

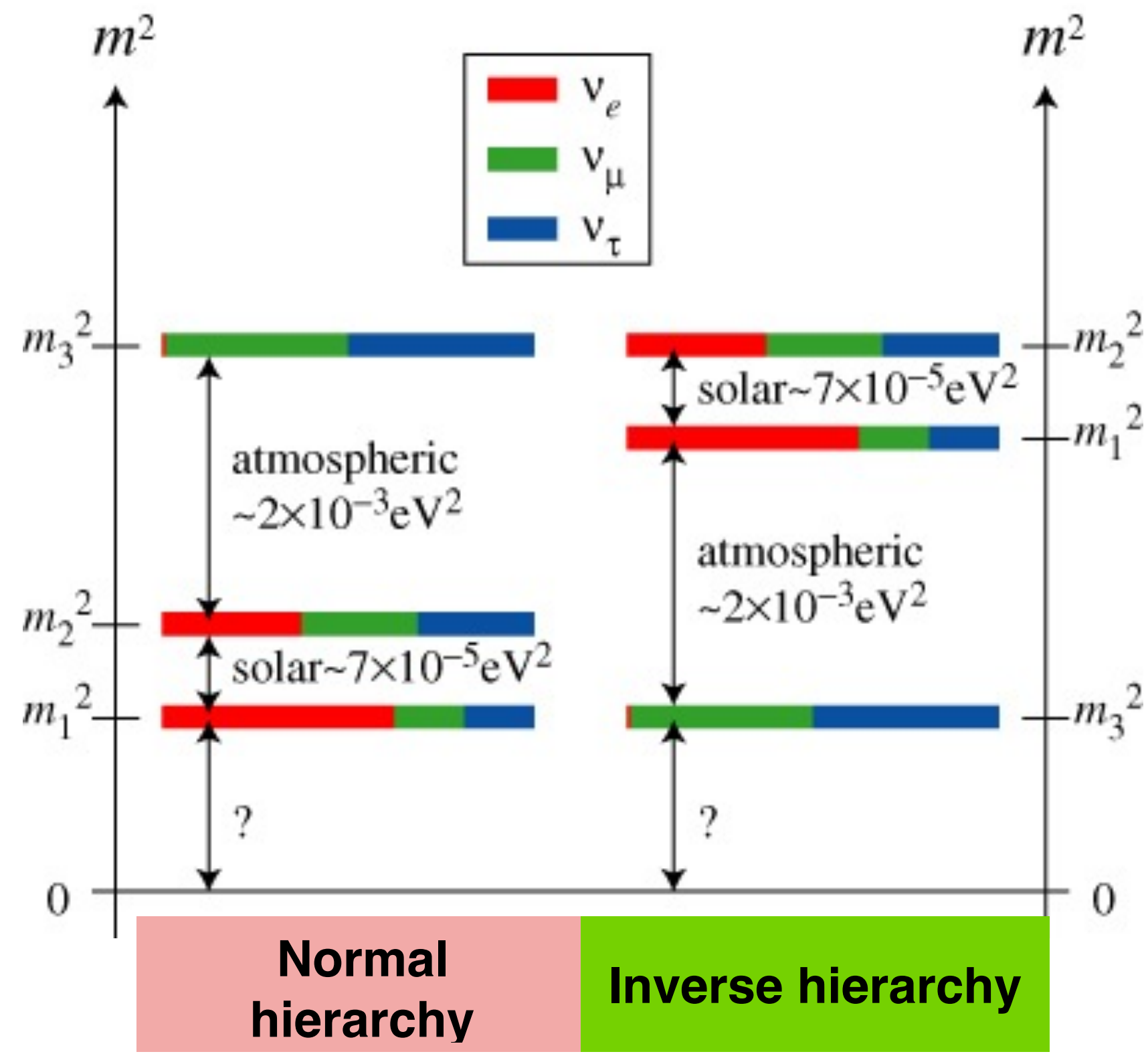
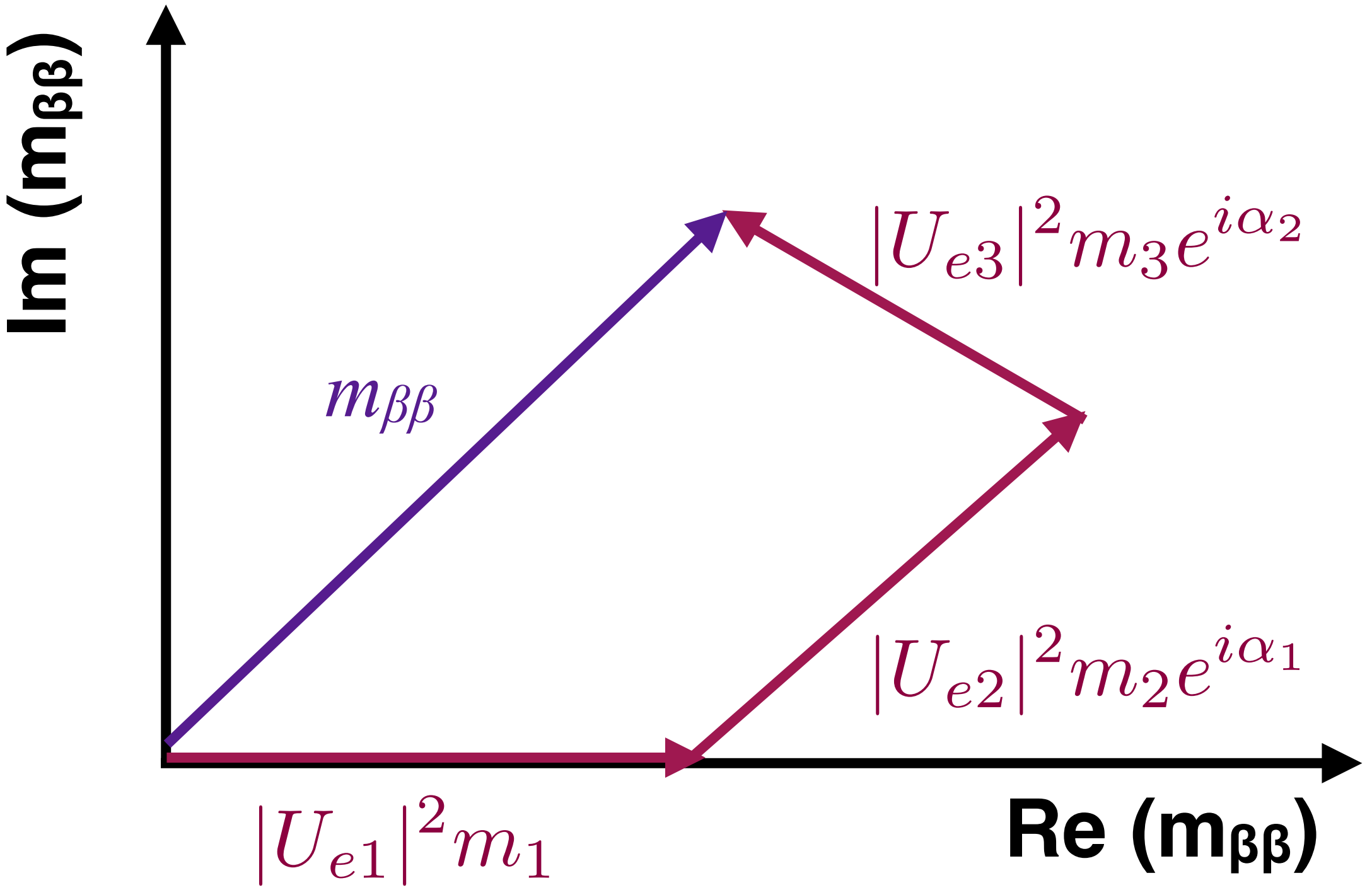
$$U = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha_1/2} & 0 \\ 0 & 0 & e^{i\alpha_2/2} \end{bmatrix}$$

$$\begin{bmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha_1/2} & 0 \\ 0 & 0 & e^{i\alpha_2/2} \end{bmatrix}$$

All about $m_{\beta\beta}$ featuring: the double-beta decay lobster!

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

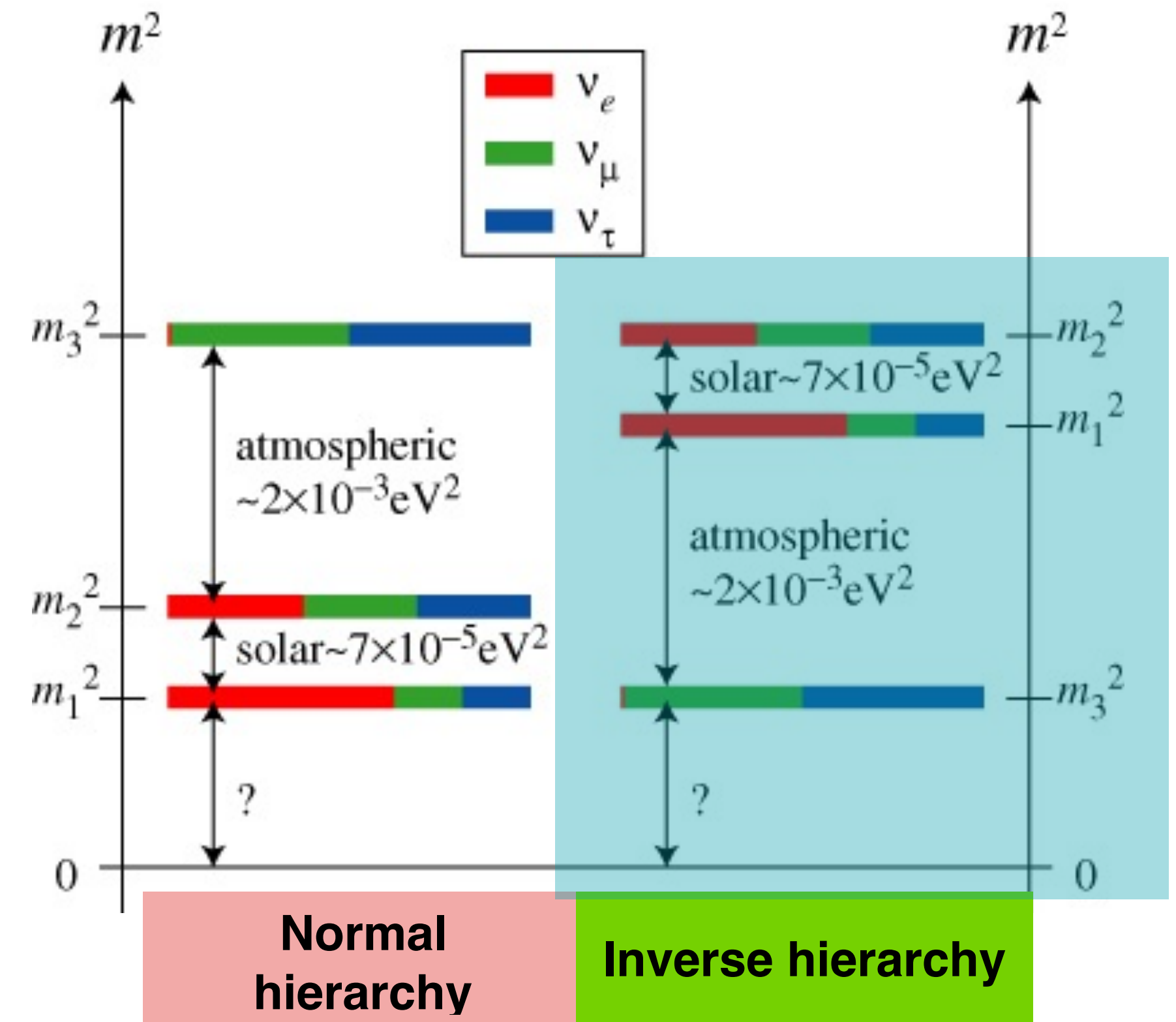
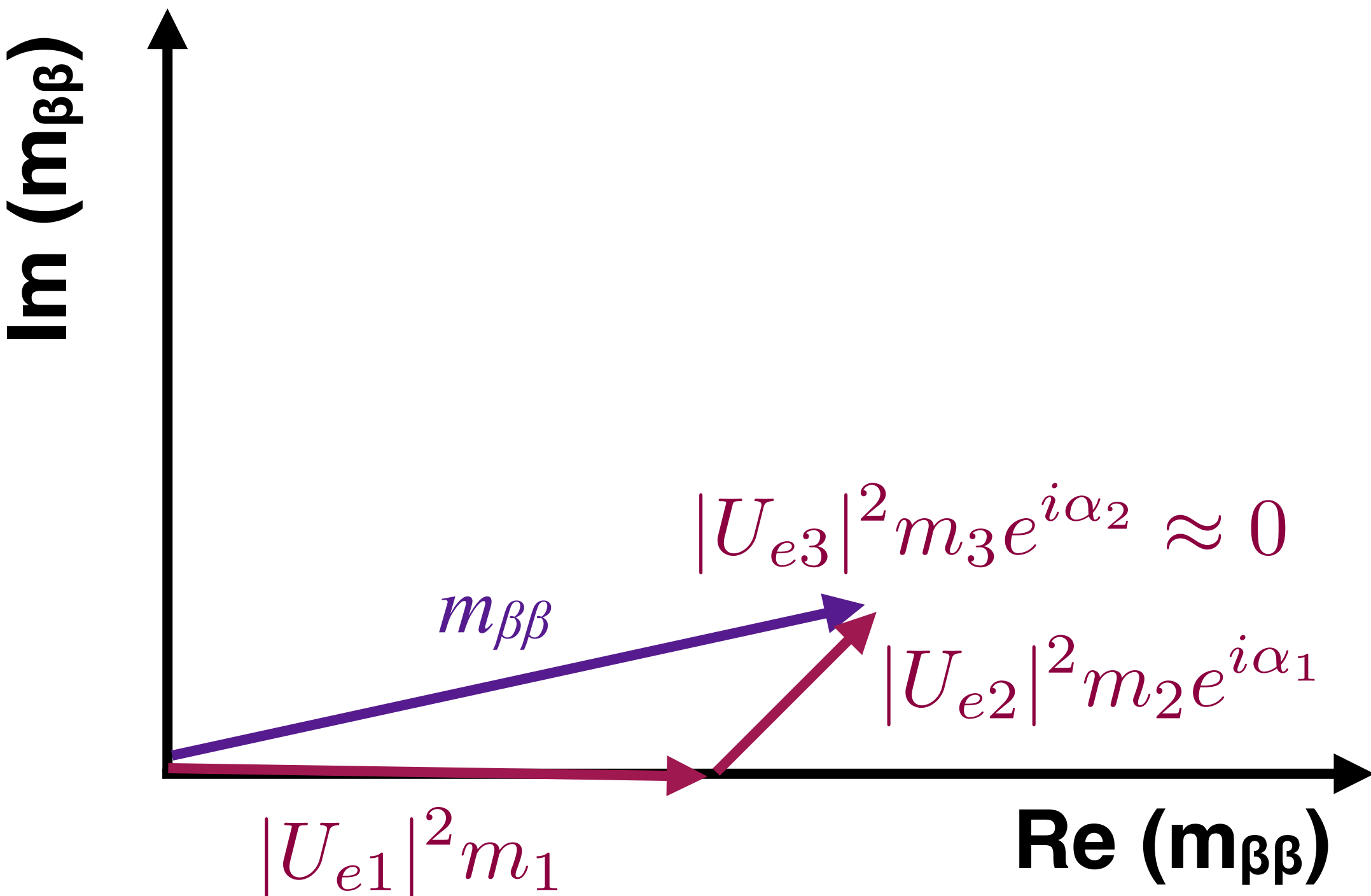
Complex phases



All about $m_{\beta\beta}$ featuring: the double-beta decay lobster!

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

Complex phases



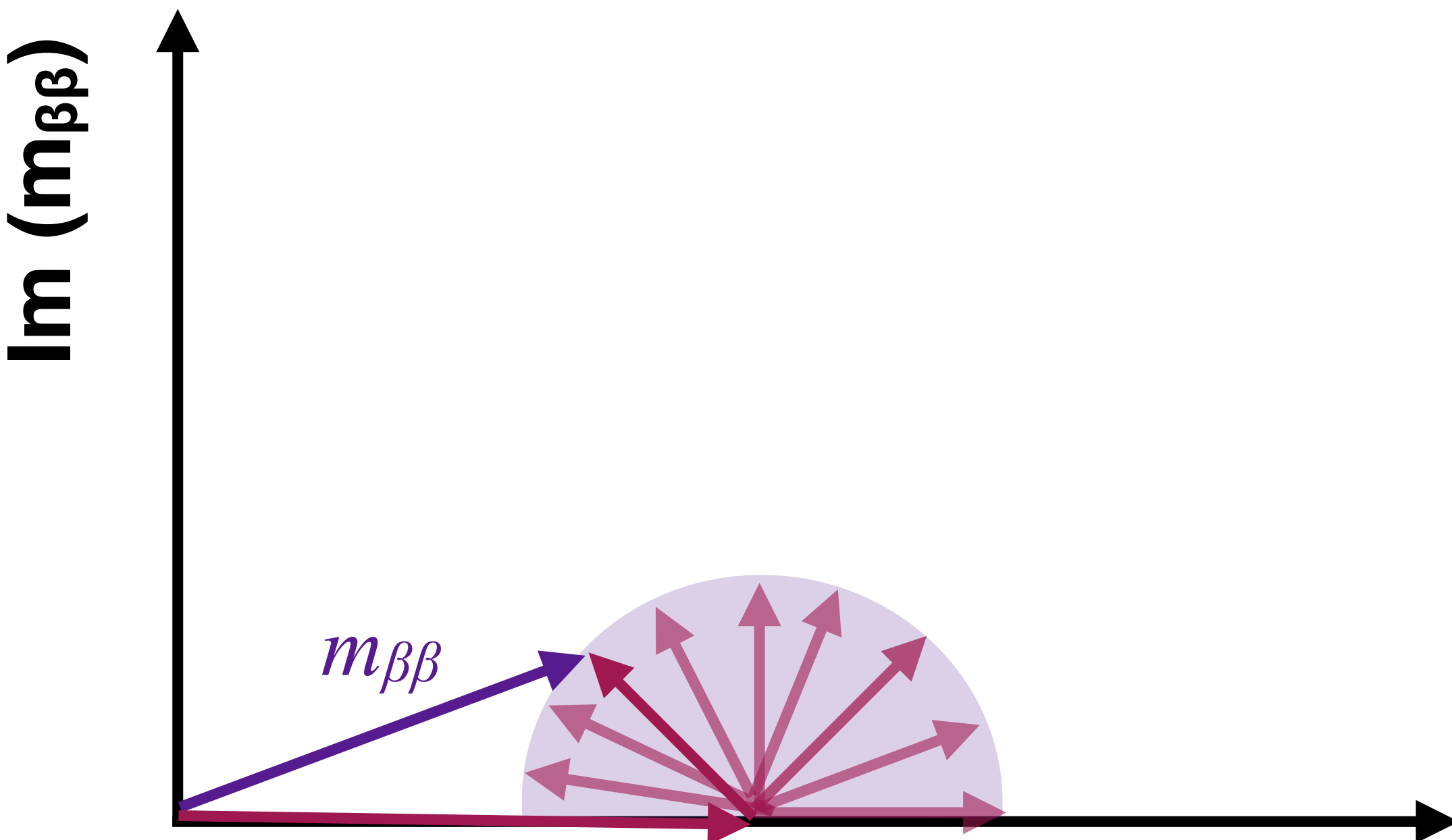
Scenario 1: Inverse hierarchy: $m_2 \sim m_1 \gg m_3$

$$m_{\beta\beta} = \underbrace{c_{12}^2 c_{13}^2}_{0.68} m_{\nu_1} + \underbrace{s_{12}^2 c_{13}^2}_{0.30} m_{\nu_2} e^{i\alpha_1} + \underbrace{s_{13}^2}_{0.02} m_{\nu_3} e^{i\alpha_3}$$

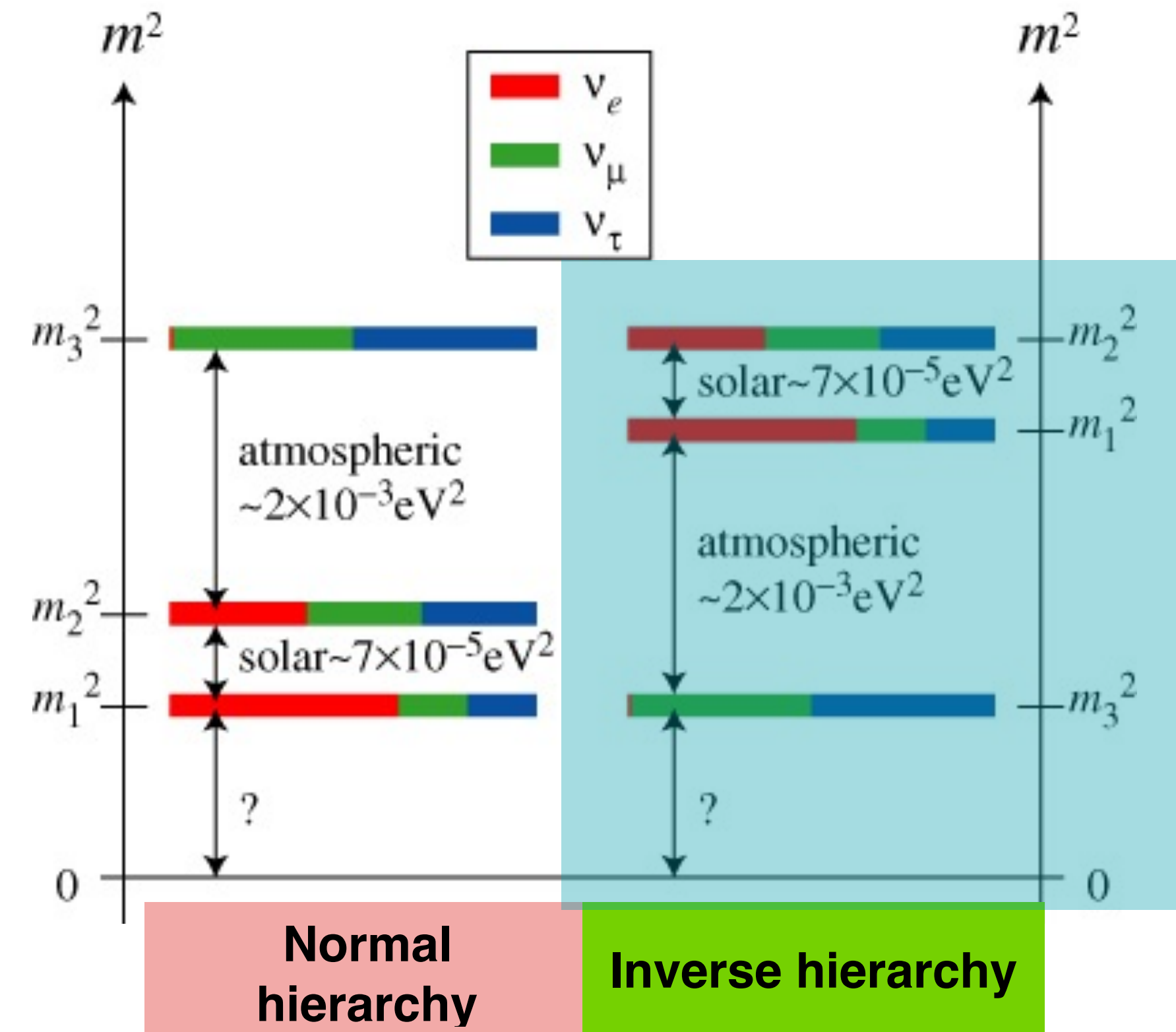
All about $m_{\beta\beta}$ featuring: the double-beta decay lobster!

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

Complex phases



We know length (mixing angles) **Re ($m_{\beta\beta}$)** but not direction (Majorana phase) of m_2 vector

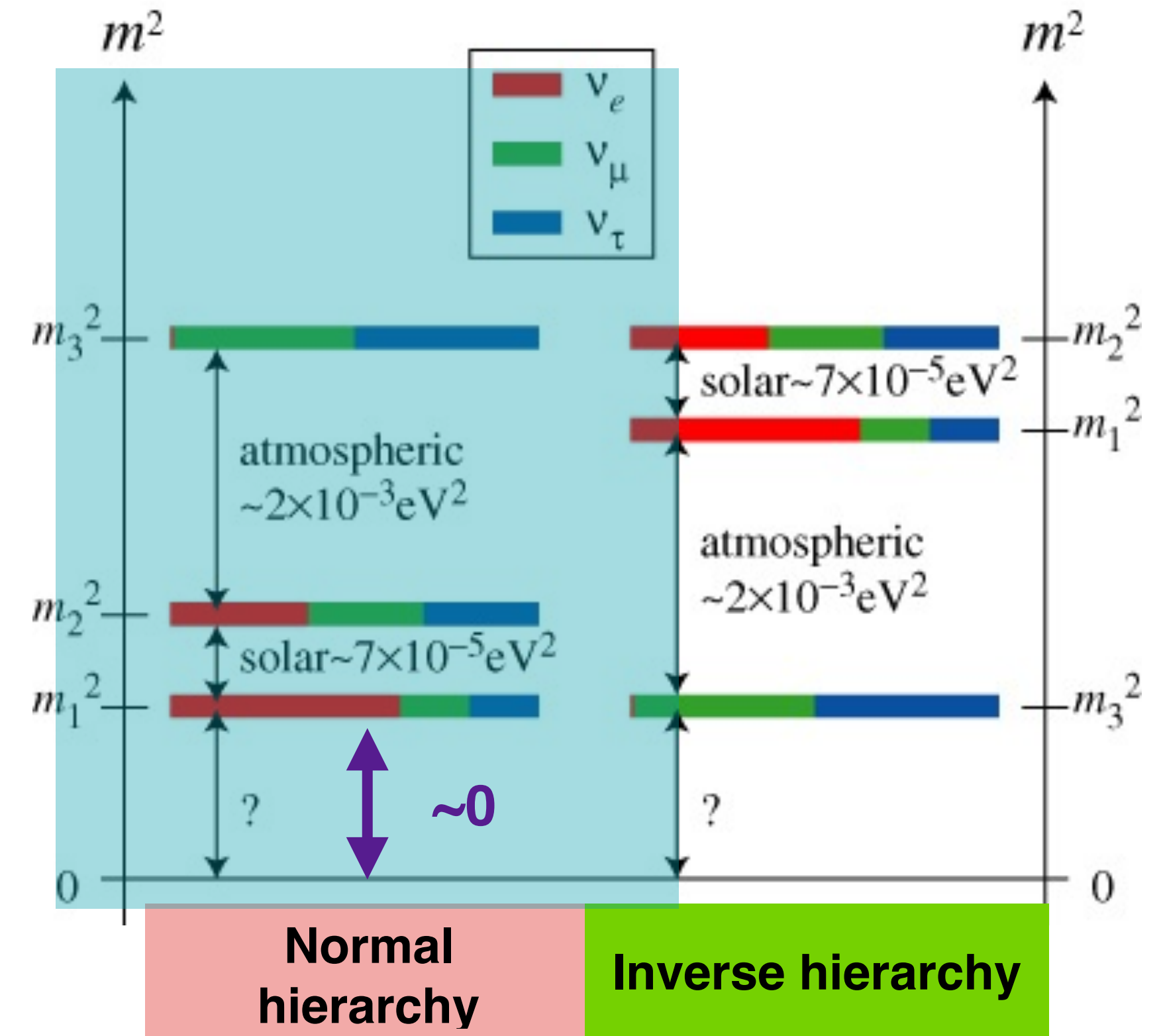
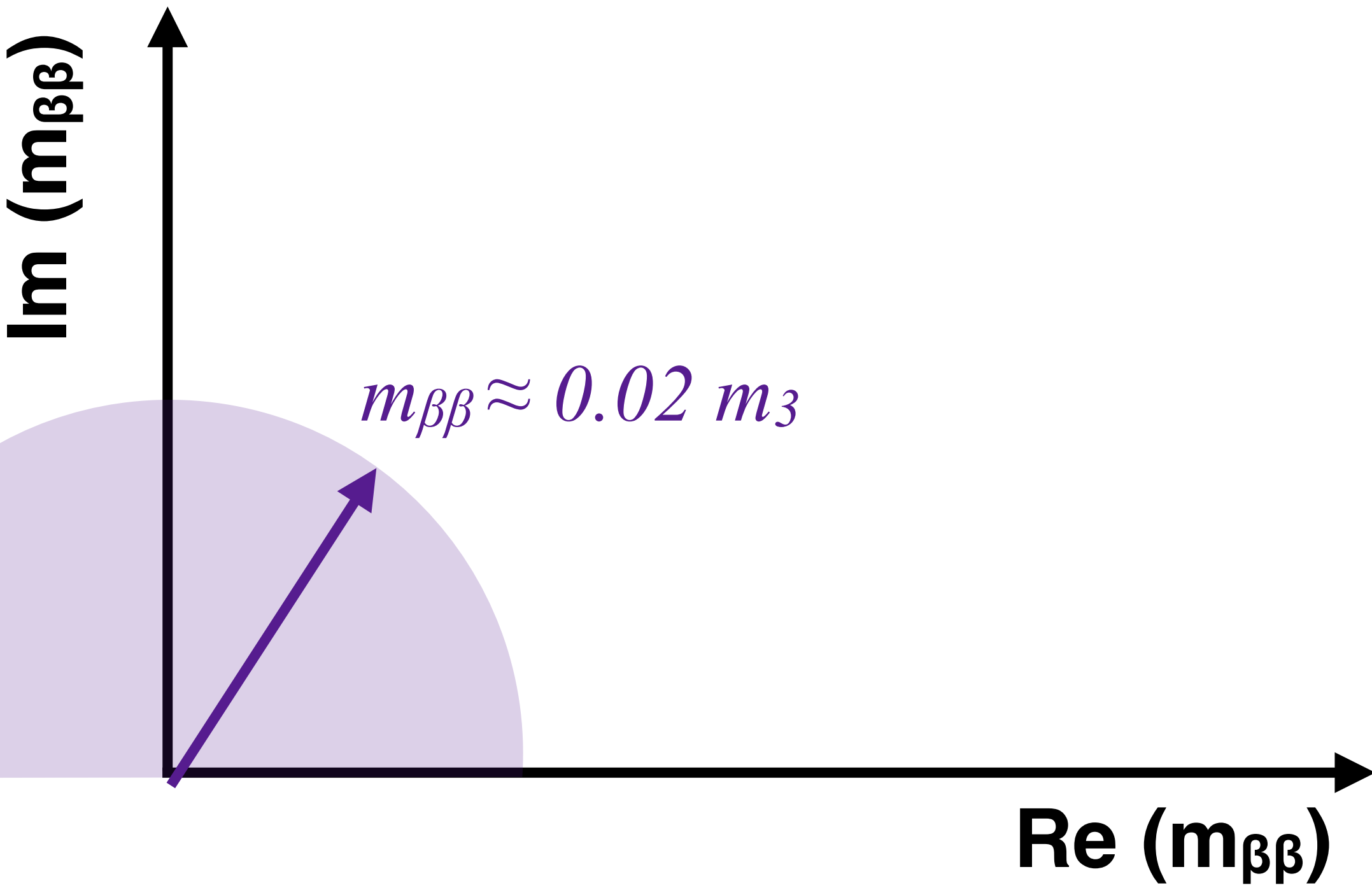


$$m_{\beta\beta} = \underbrace{c_{12}^2 c_{13}^2}_{0.68} m_{\nu_1} + \underbrace{s_{12}^2 c_{13}^2}_{0.30} m_{\nu_2} e^{i\alpha_1} + \underbrace{s_{13}^2}_{0.02} m_{\nu_3} e^{i\alpha_3}$$

All about $m_{\beta\beta}$ featuring: the double-beta decay lobster!

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Complex phases



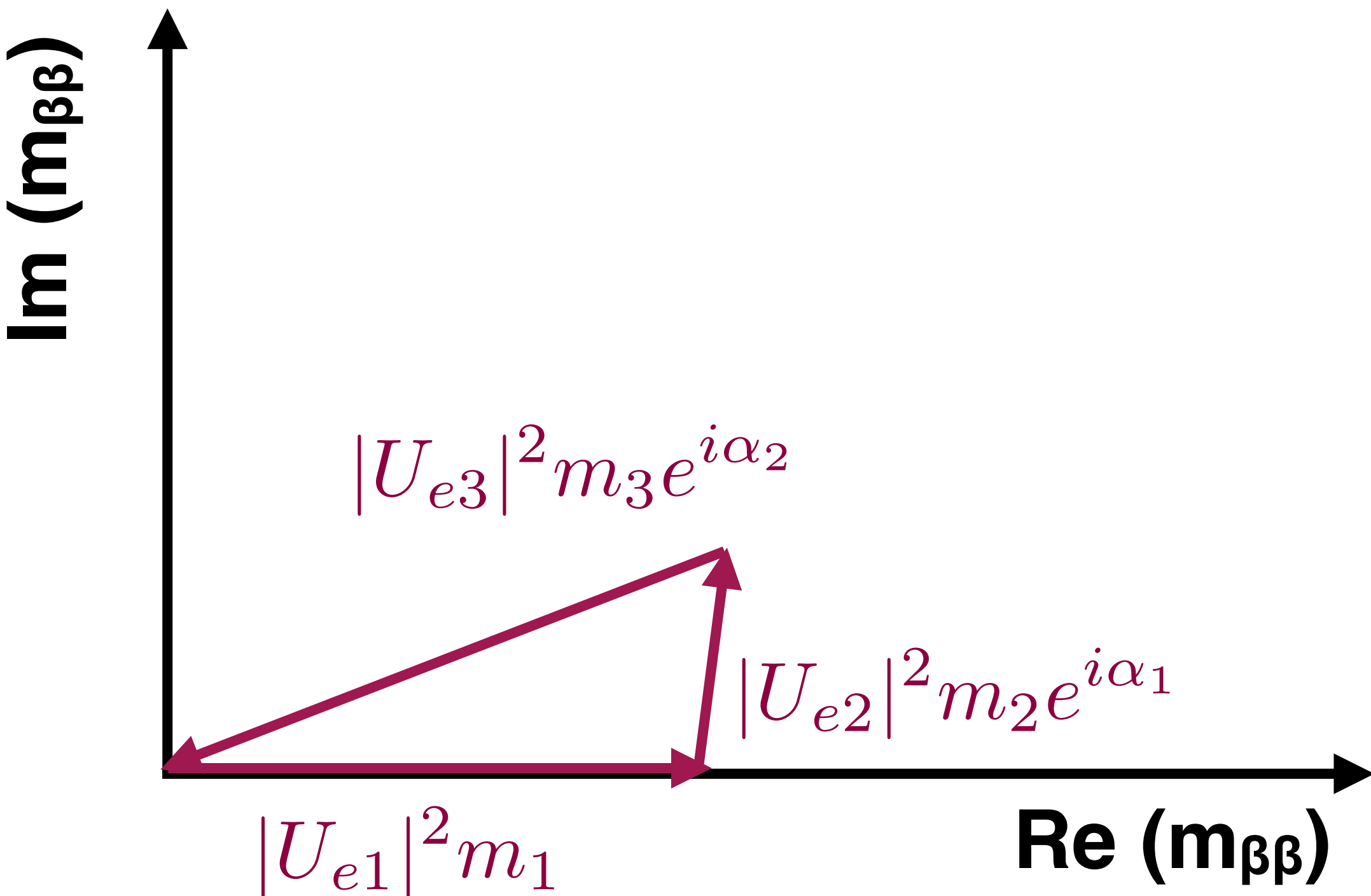
Scenario 2: Normal hierarchy: $m_3 \gg m_2 > m_1$

$$m_{\beta\beta} = \underbrace{c_{12}^2 c_{13}^2}_{0.68} m_{\nu_1} + \underbrace{s_{12}^2 c_{13}^2}_{0.30} m_{\nu_2} e^{i\alpha_1} + \underbrace{s_{13}^2}_{0.02} m_{\nu_3} e^{i\alpha_3}$$

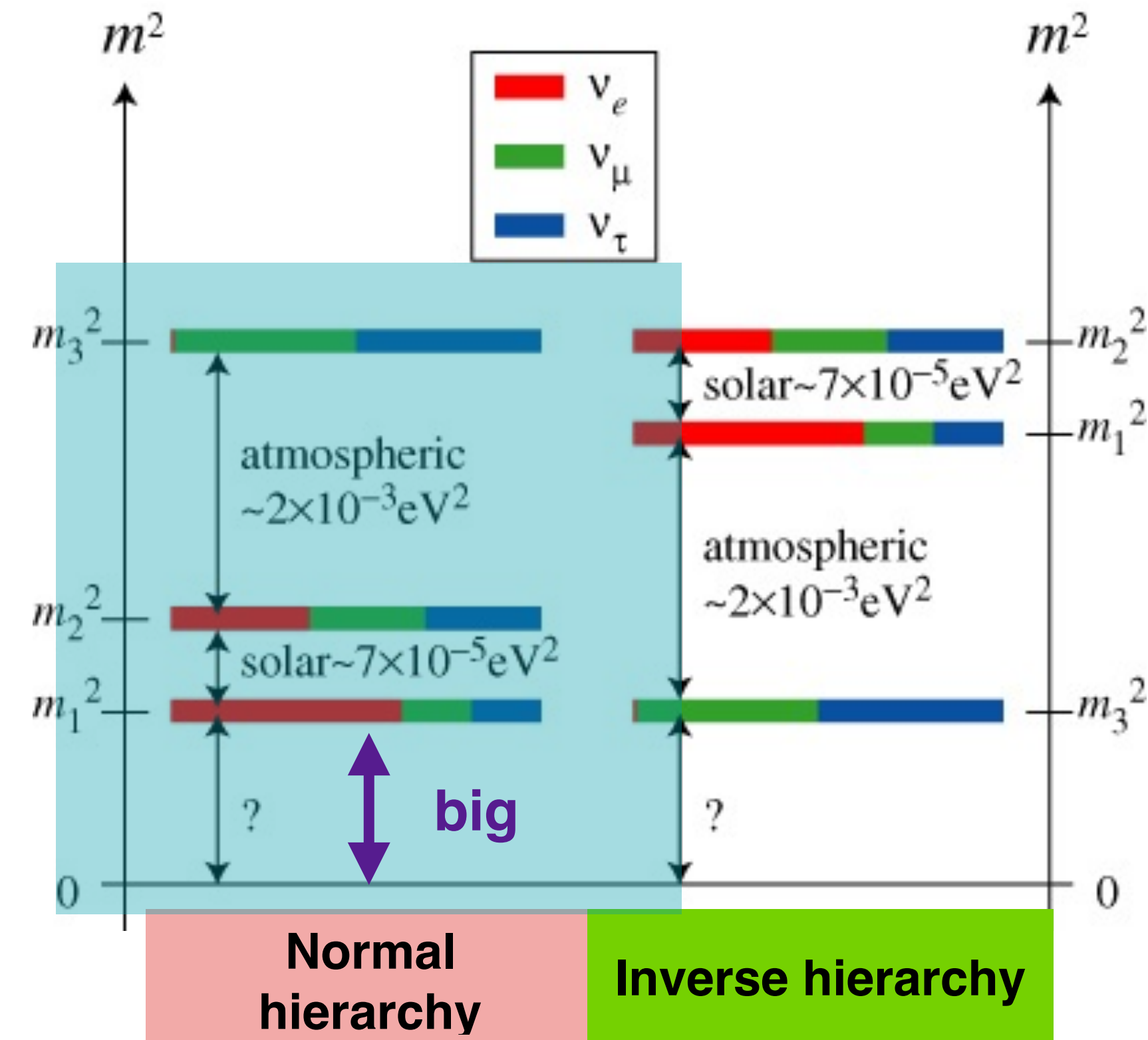
All about $m_{\beta\beta}$ featuring: the double-beta decay lobster!

$$m_{\beta\beta} = \left| \sum_{i=1}^3 |U_{ei}|^2 m_i e^{i\alpha_i} \right|$$

Complex phases



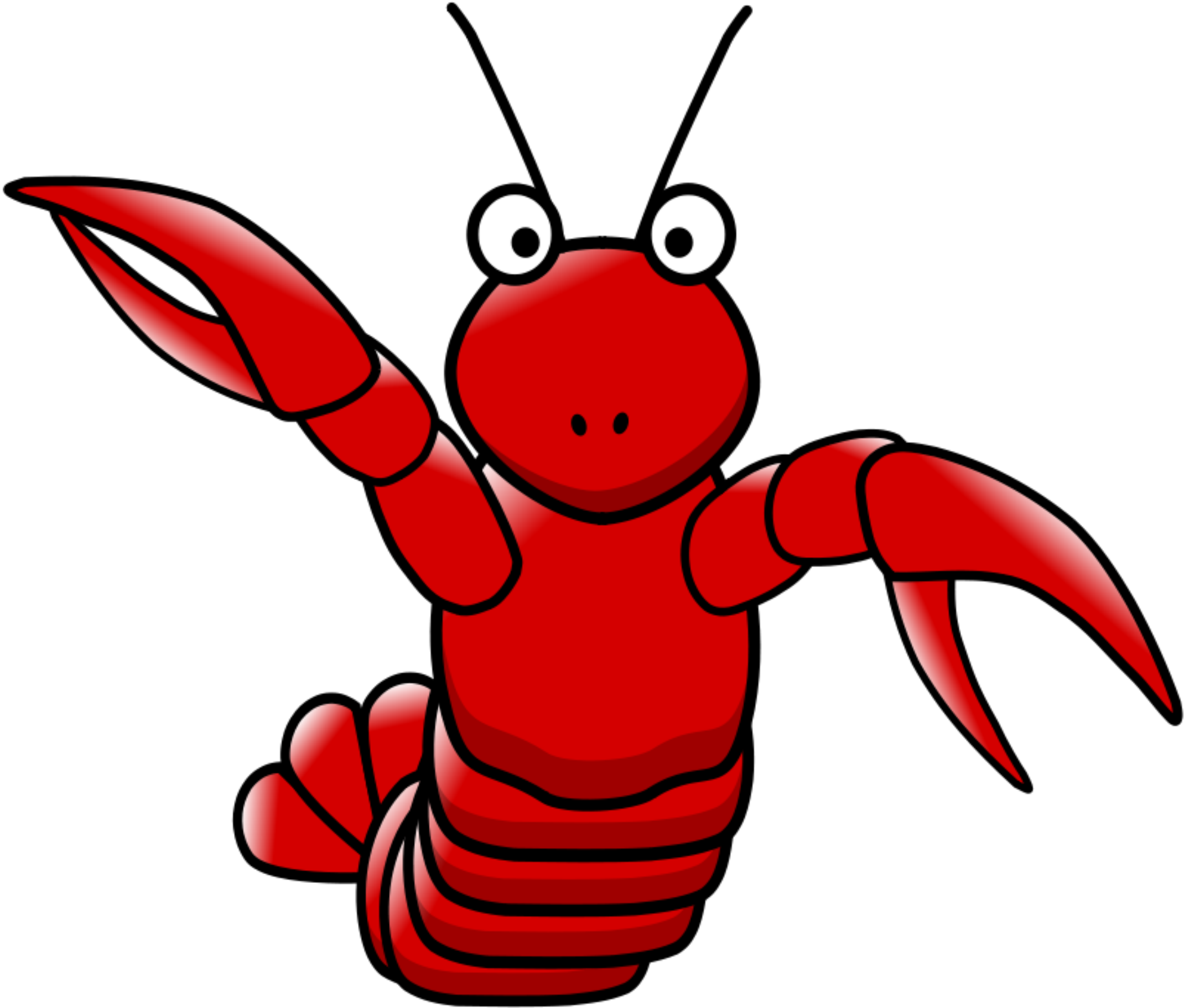
$m_{\beta\beta} = 0$ is possible for the some masses / phases!



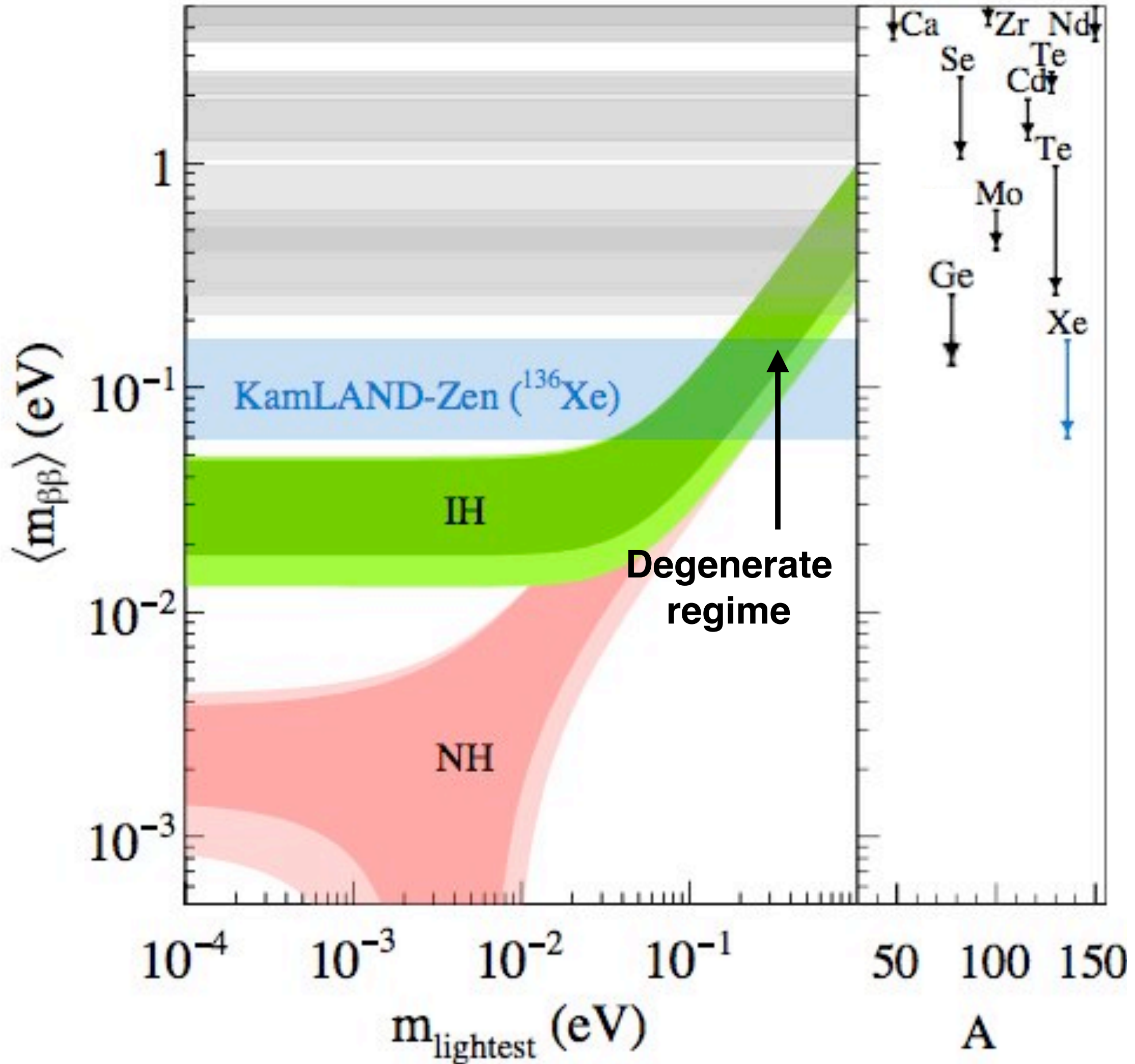
Scenario 3: Normal hierarchy: $m_3 > m_2 > m_1$

$$m_{\beta\beta} = \underbrace{c_{12}^2 c_{13}^2}_{0.68} m_{\nu_1} + \underbrace{s_{12}^2 c_{13}^2}_{0.30} m_{\nu_2} e^{i\alpha_1} + \underbrace{s_{13}^2}_{0.02} m_{\nu_3} e^{i\alpha_3}$$

Introducing the lobster!



Adapted from PRL 117, 082503 (2016)



Introducing the lobster!

Nuclear effects

$$0\nu\beta\beta \text{ rate } \frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z) g_A^4 |M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$$

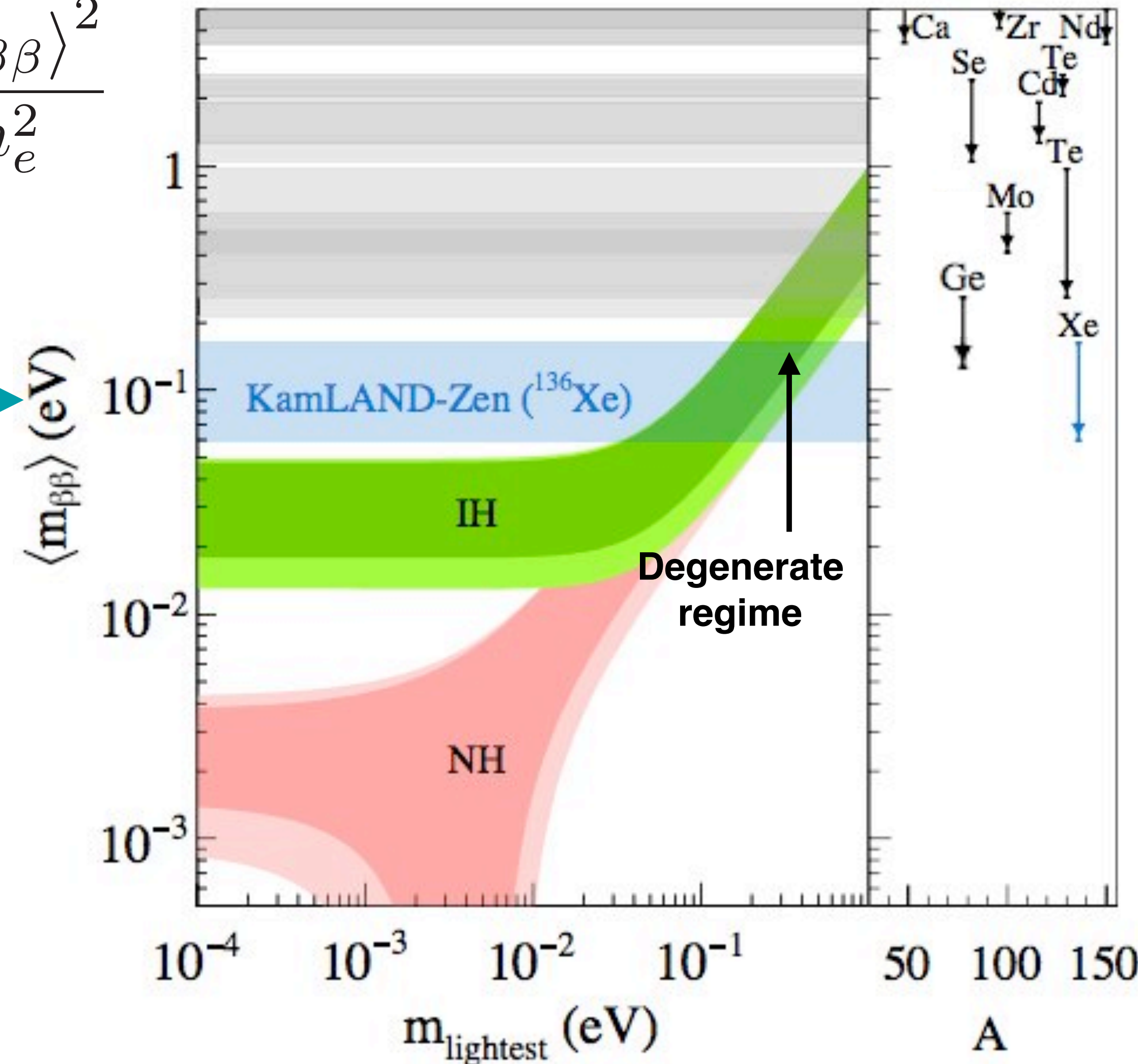
- **Best** current limit on the **$0\nu\beta\beta$ rate**

$$T_{1/2} > 1.07 \times 10^{26} \text{ years}$$

- Corresponds to **lowest $m_{\beta\beta}$ limit**

$$\langle m_{\beta\beta} \rangle < 61\text{-}165 \text{ meV}$$

- **Range** due to model uncertainty on **nuclear effects**

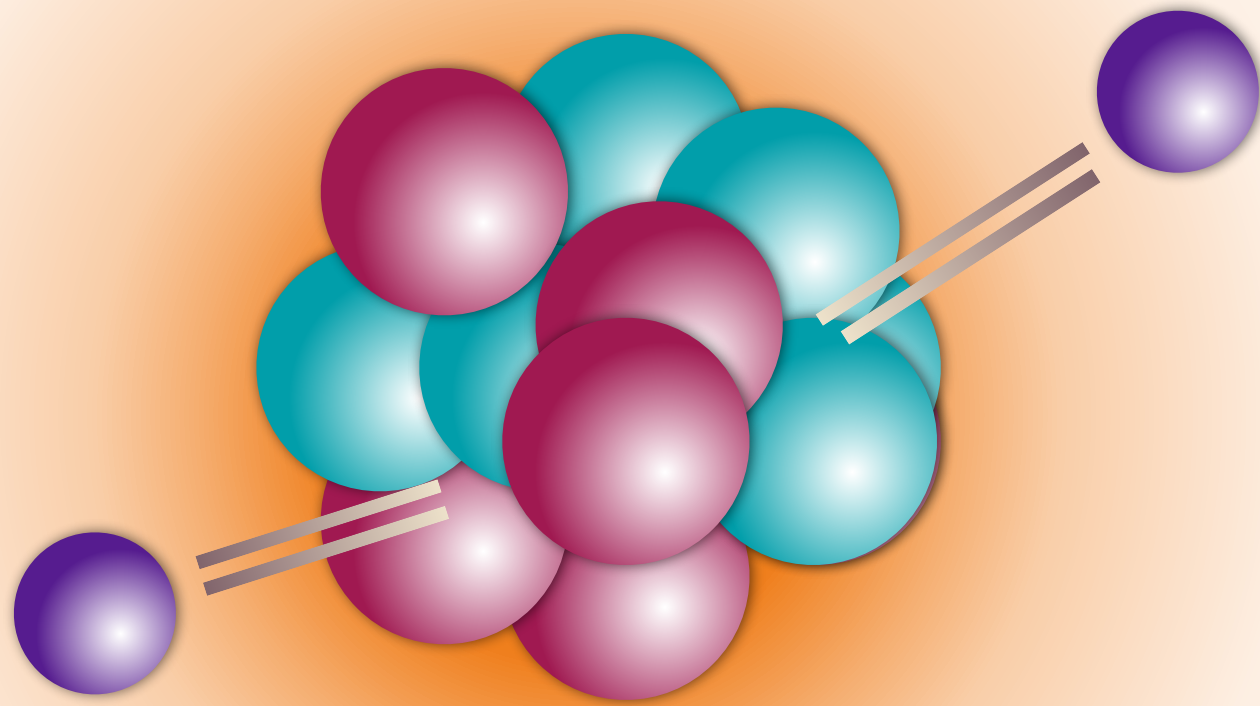


$0\nu\beta\beta$ rate

$$\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z) g_A^4 |M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$$

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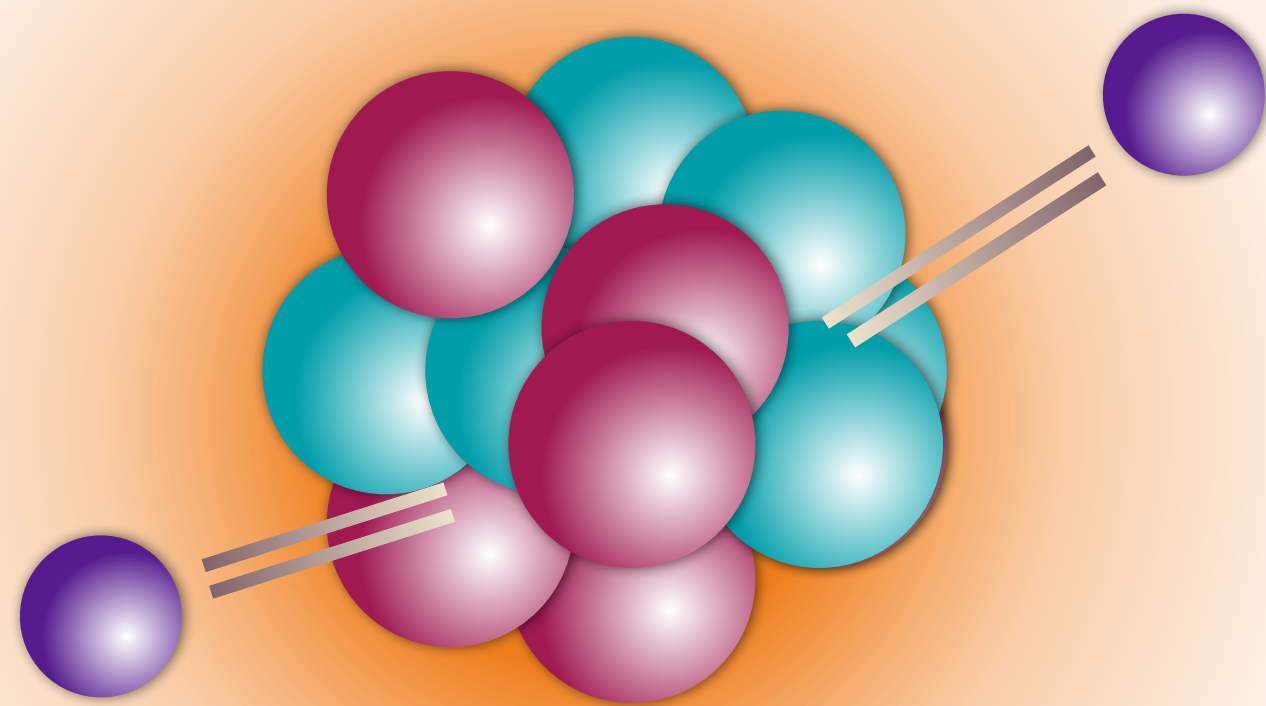
Phase space factor



Coulomb effects of the nucleus on the emitted electrons

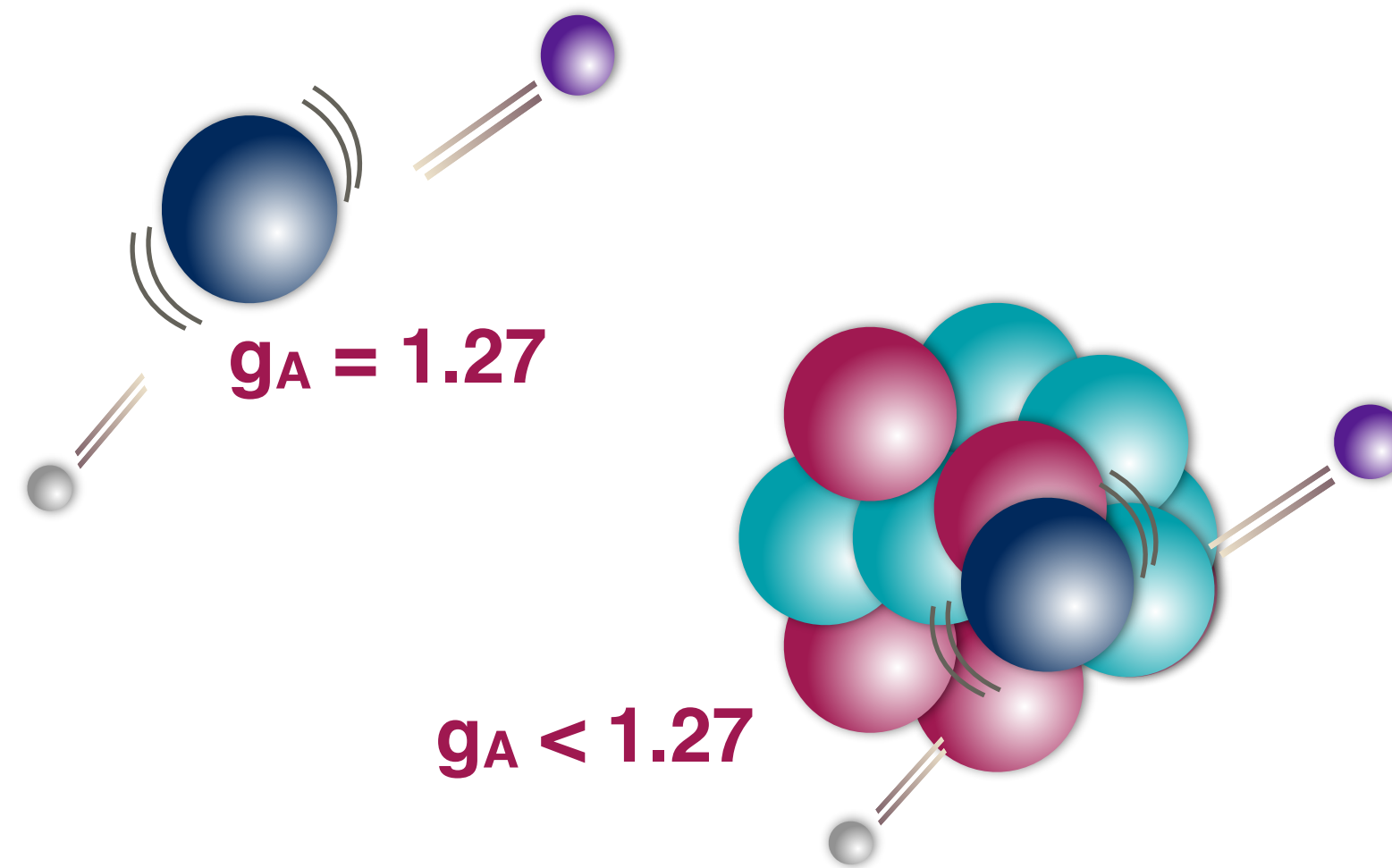
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Phase space factor



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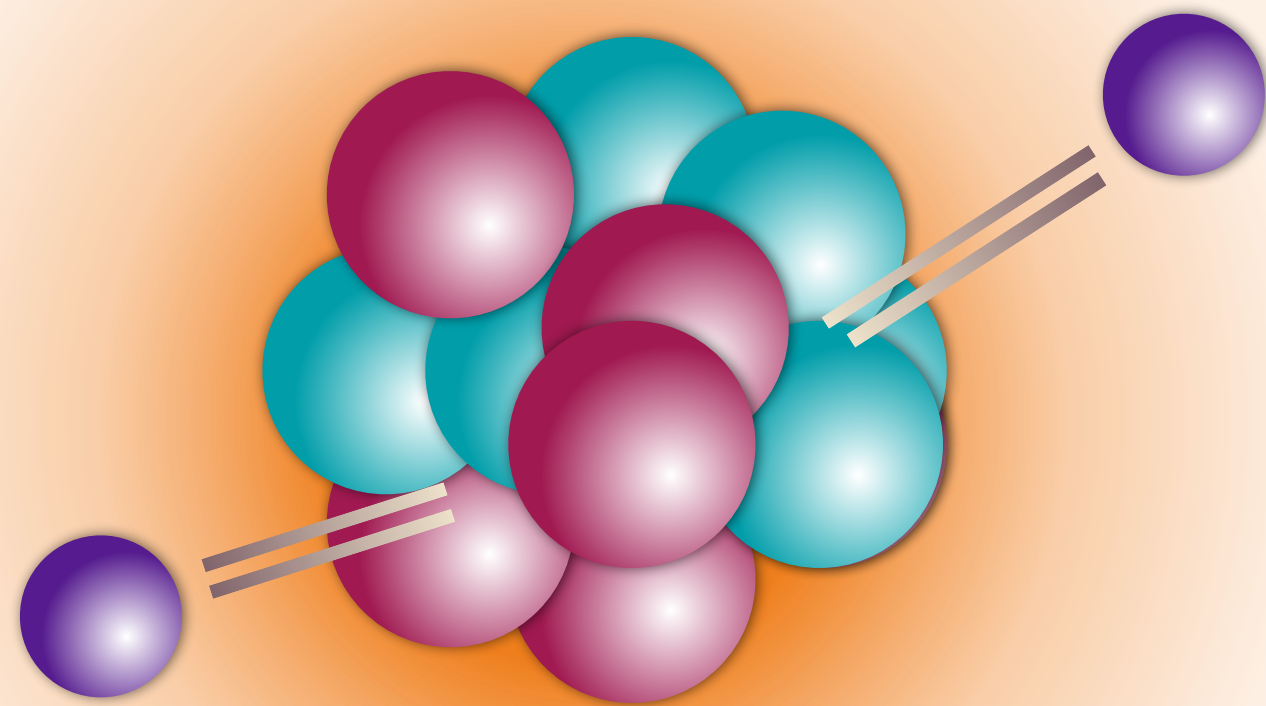
Axial coupling constant



Axial vector coupling constant appears quenched in heavy nuclei

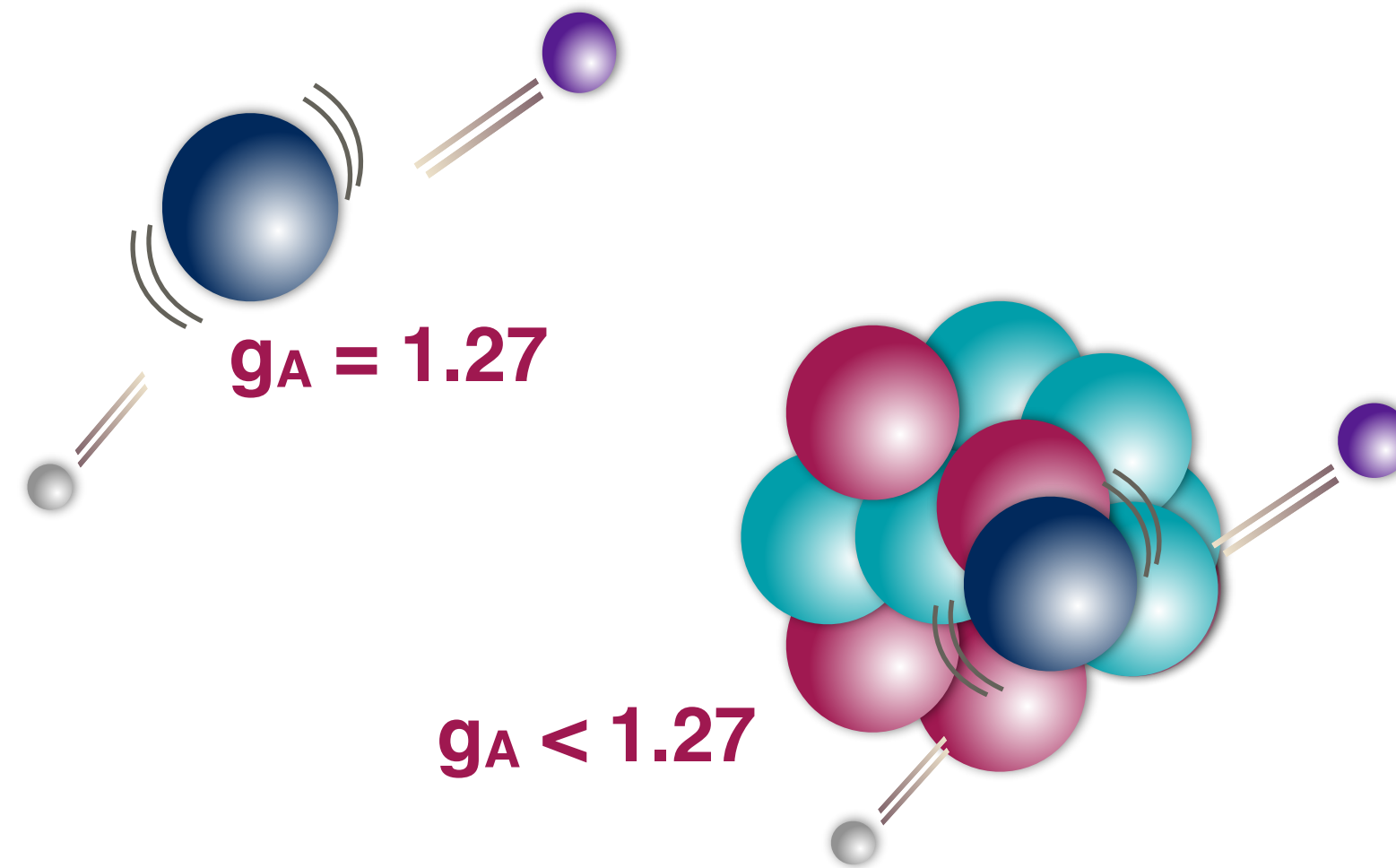
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Phase space factor



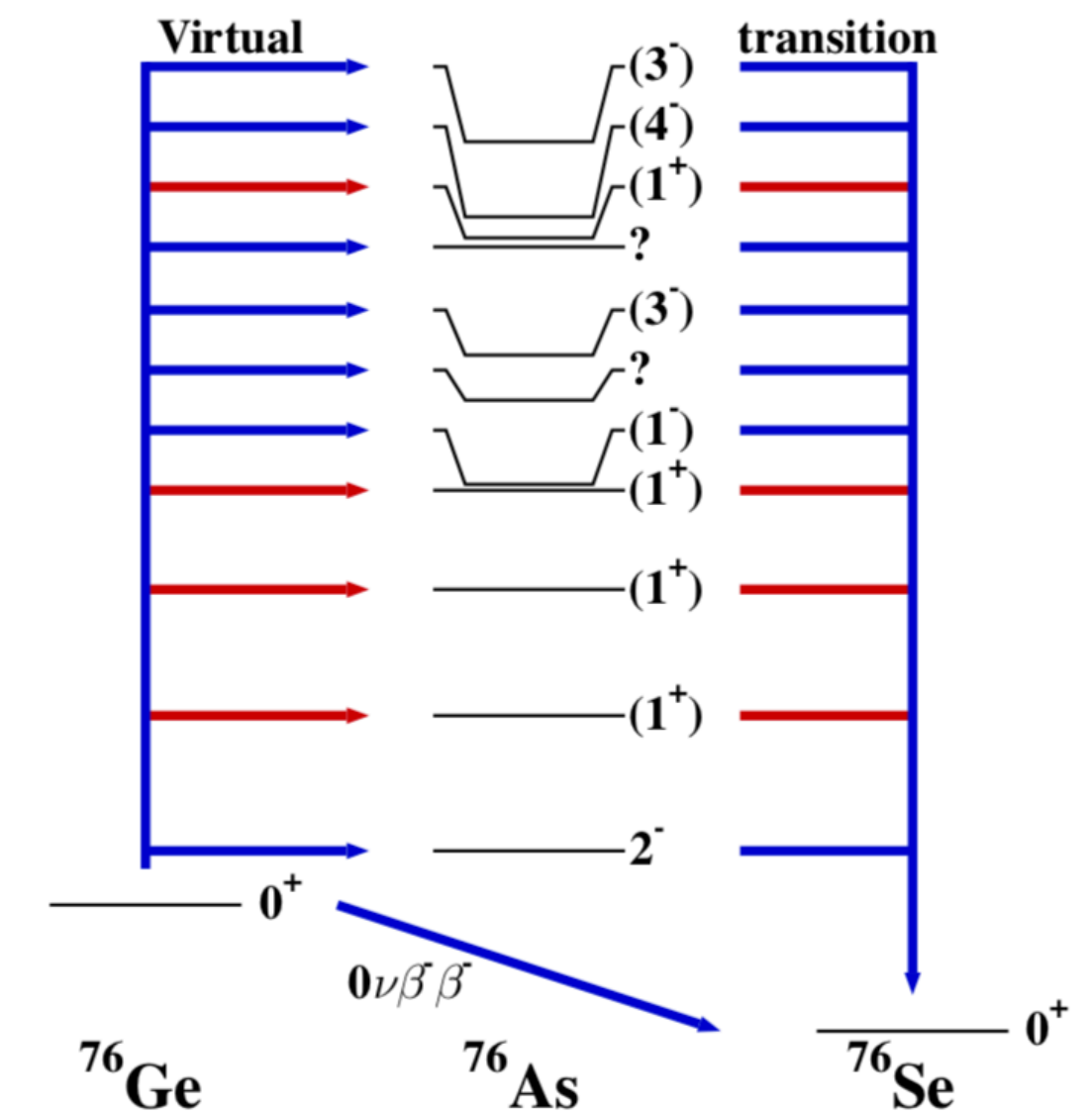
Coulomb effects of the nucleus on the emitted electrons

Axial coupling constant



Axial vector coupling constant appears quenched in heavy nuclei

Nuclear matrix element



Nuclear structure effects of the parent, daughter, and intermediate nuclei

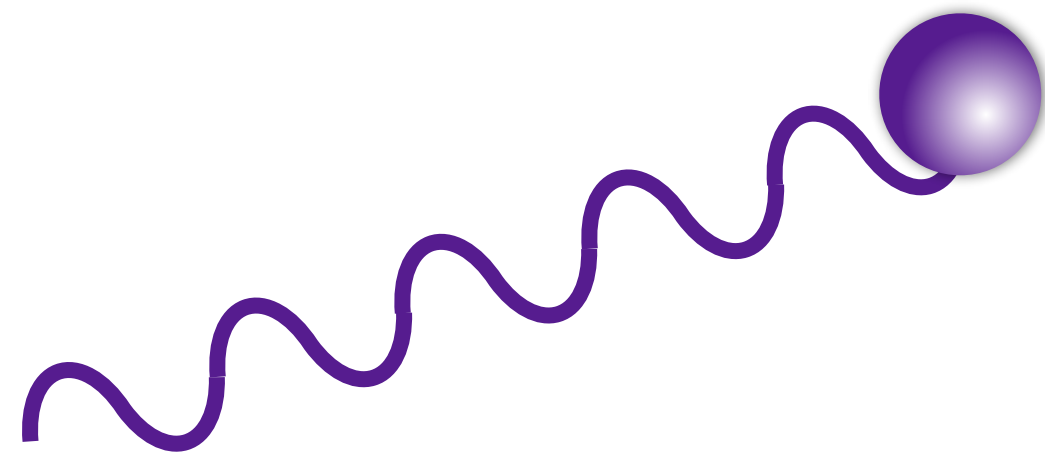
IPPP/05/56, DCPT/05/114

$$0\nu\beta\beta \text{ rate } \frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z) g_A^4 |M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$$



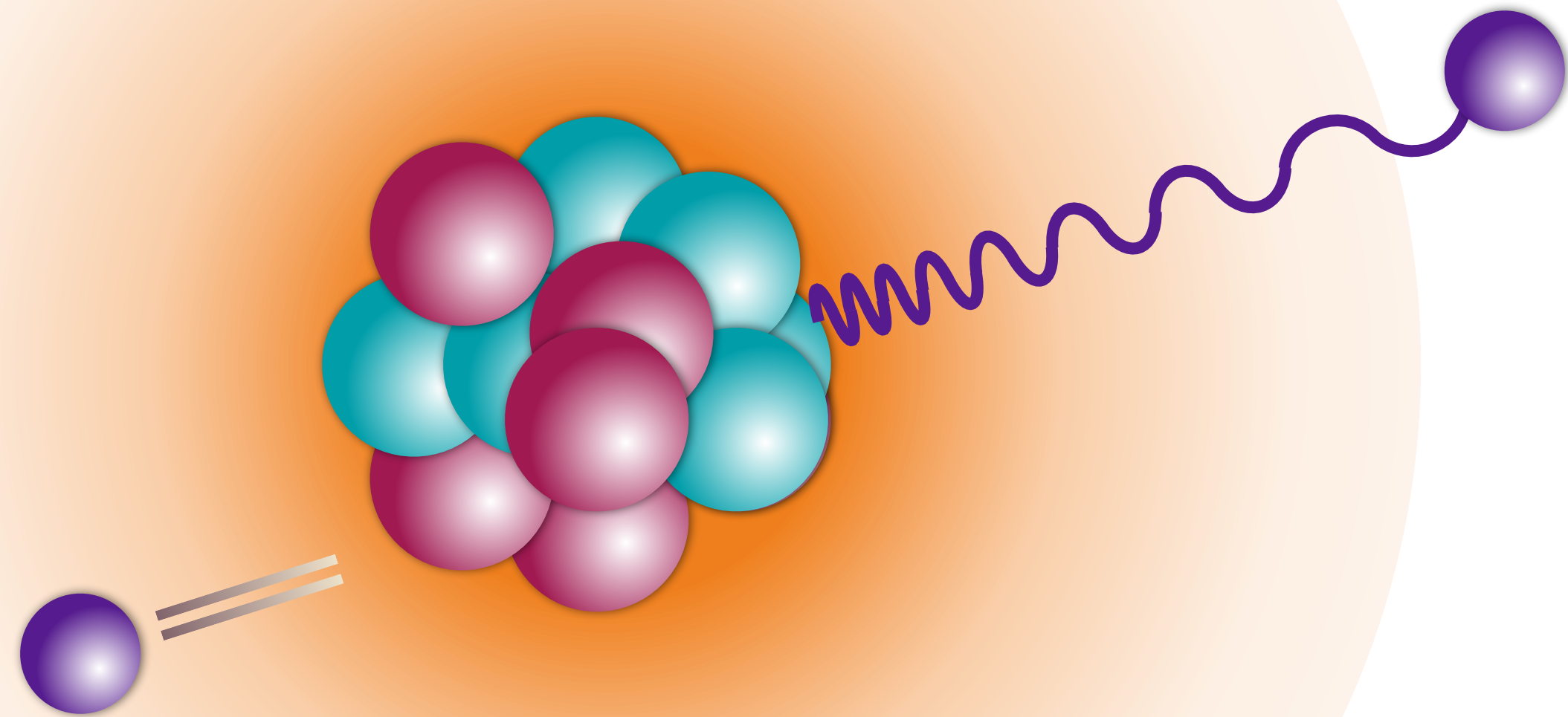
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- The wave function for a free electron: a **plane wave**



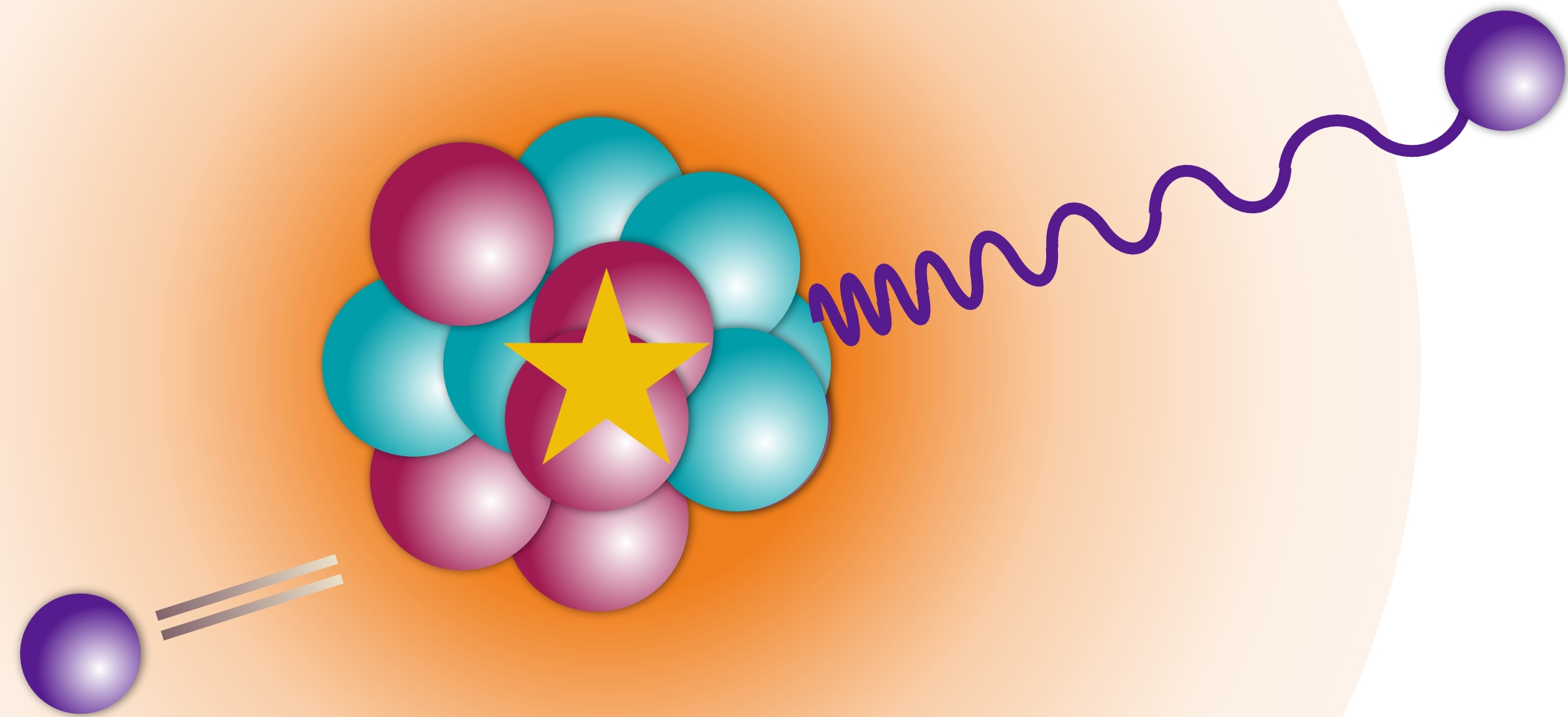
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- The wave function for a free electron: a **plane wave**
- ... is distorted by the **Coulomb field** of the nucleus



DOI: 10.3389/fphy.2019.00012

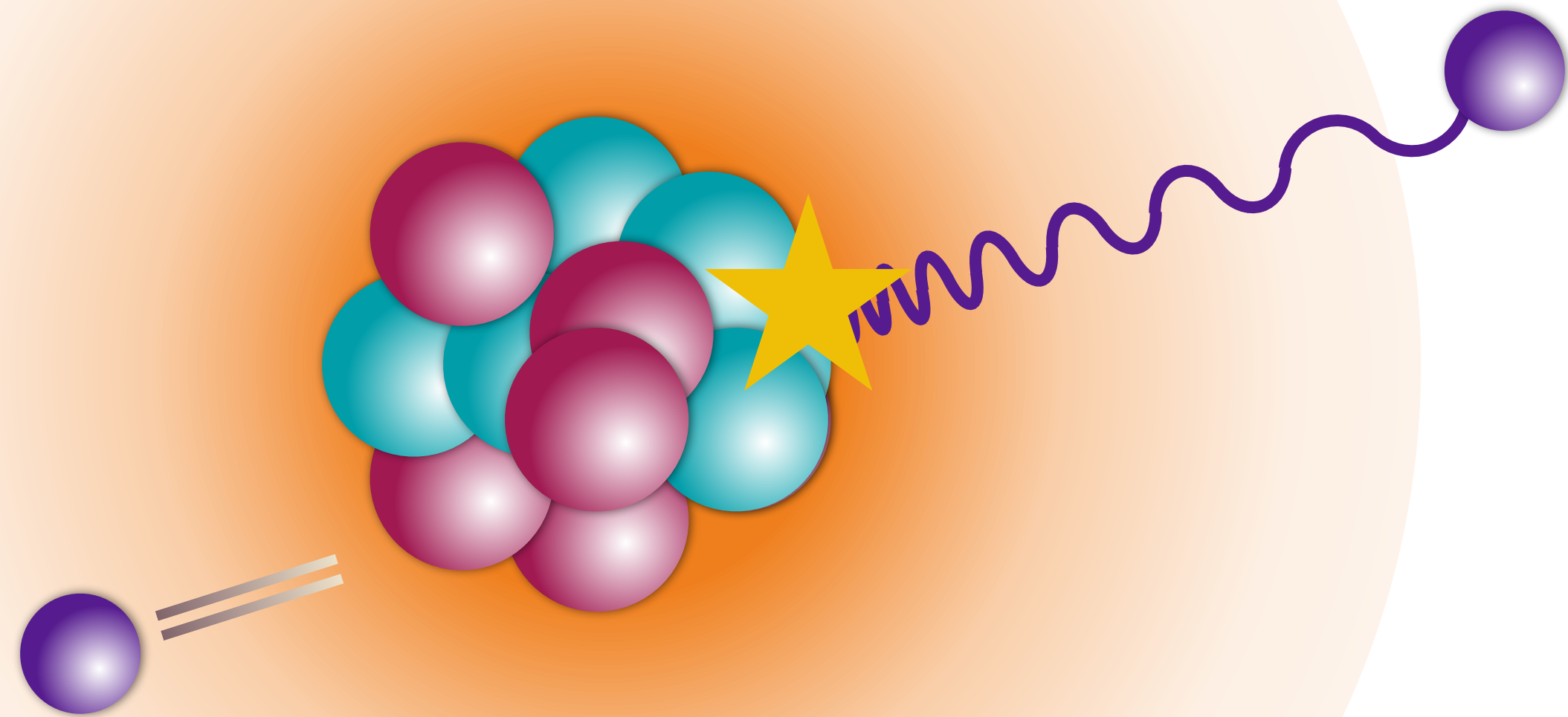
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- The wave function for a free electron: a **plane wave**
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- **Nonrelativistic** model - multiply electron wave function by a **Fermi factor**
 - Solve Schrödinger equation for a Coulomb potential of **point charge** equal to nuclear charge Z
 - Divide by **plane wave** solution
 - Evaluate at the **origin**

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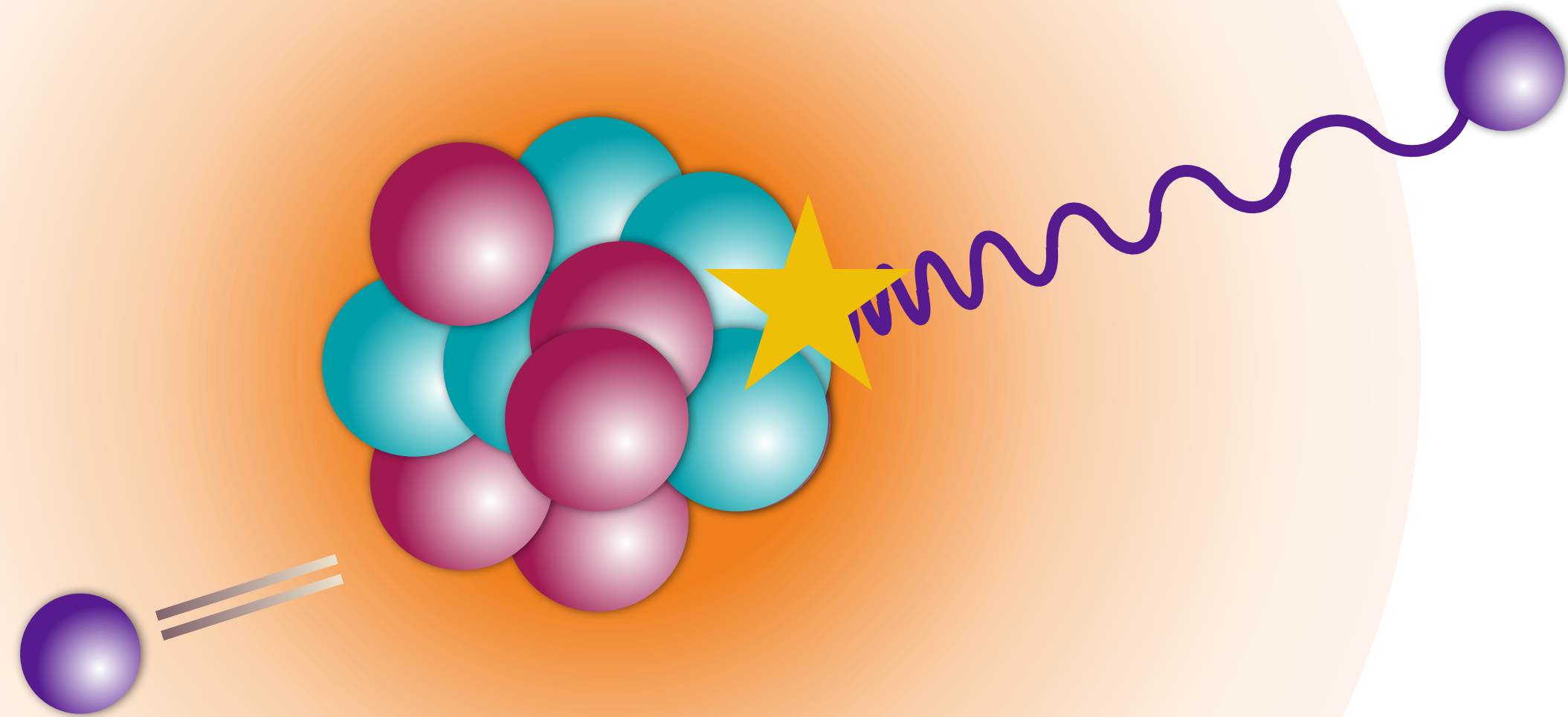
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 - Solve **Dirac equation** in point charge potential
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 - Solve **Dirac equation** in point charge potential
 - Divide by plane wave at **nuclear surface**
- Recent **enhancements** include
 - **Screening**
 - Realistic **proton density** functions
 - Including transitions to **excited nuclear states**

DOI: 10.3389/fphy.2019.00012

Phase space factor cheat sheet

Nucleus	Z	$Q_{\beta\beta}$	$G_{0\nu}$ (10^{-15} yr $^{-1}$): various models					
^{48}Ca	20	4.267	24.65	24.81	26.1	26	24.83	24.55
^{150}Nd	60	3.371	61.94	63.03	85.9	78.4	63.16	66
^{100}Mo	42	3.034	15.84	15.92	18.7	45.6	15.95	15.74
^{82}Se	34	2.996	10.14	10.16	11.4	11.1	10.18	9.96
^{130}Te	52	2.528	14.24	14.22	19.4	16.7	14.25	14.1
^{136}Xe	54	2.458	14.54	14.58	19.4	17.7	14.62	14.49
^{76}Ge	32	2.039	2.372	2.363	2.62	2.55	2.37	2.28

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$G_{0\nu}$ increases strongly with $Q_{\beta\beta}$

$$T_{1/2}^{0\nu} \propto Q_{\beta\beta}^{-5}$$

Adapted from DOI: 10.3389/fphy.2019.00012

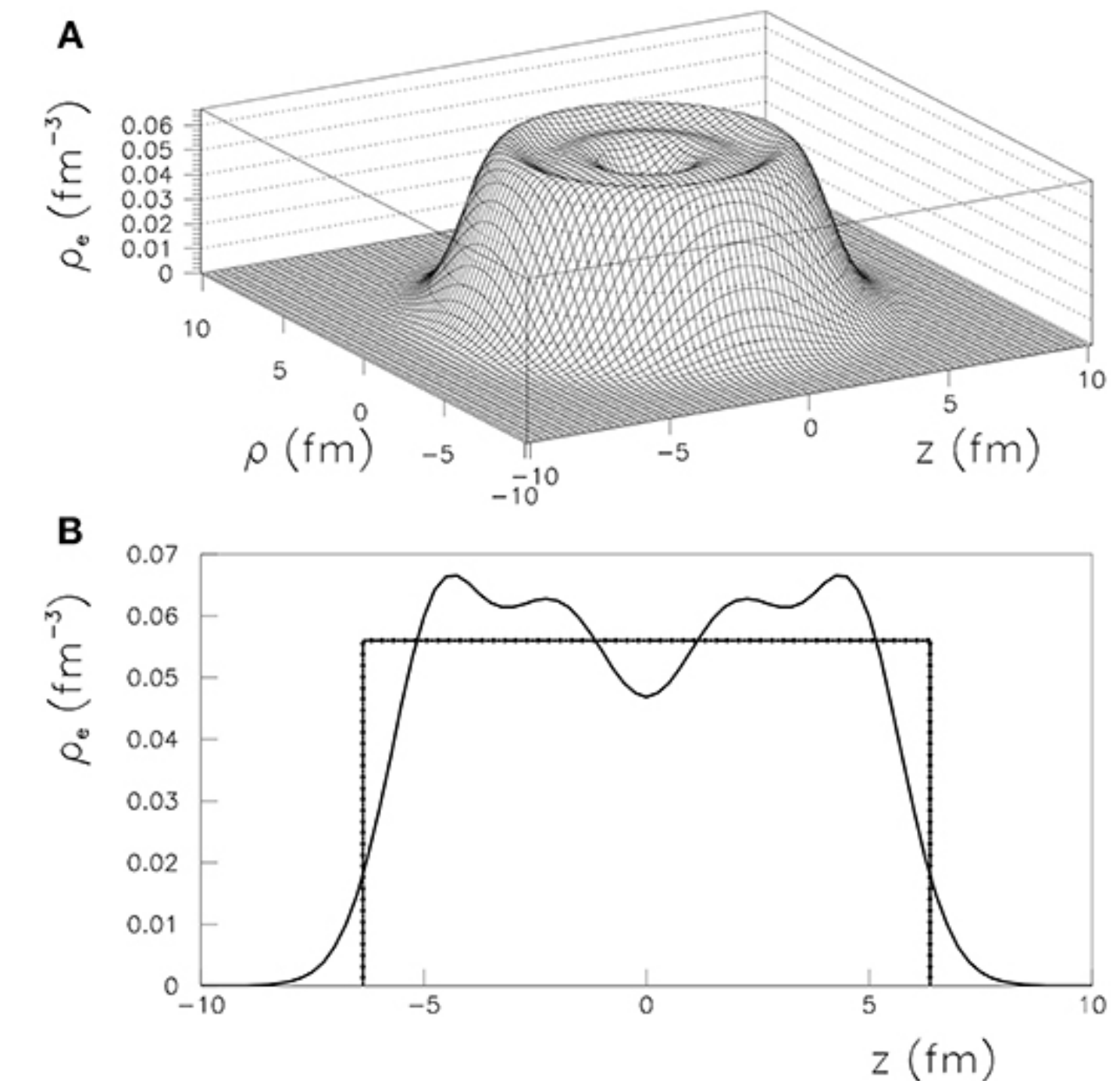
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Also depends on Z, nuclear radius ($\sim A^{1/3}$), charge density



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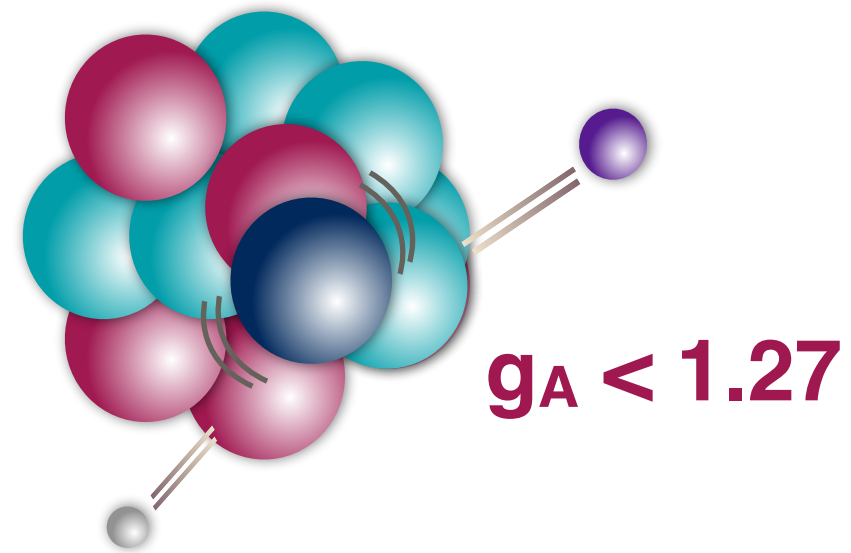
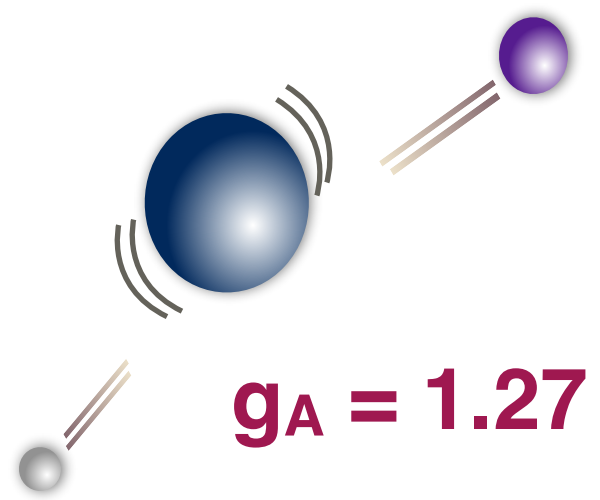
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Also depends on Z, nuclear radius ($\sim A^{1/3}$), charge density

Relatively little variation between models

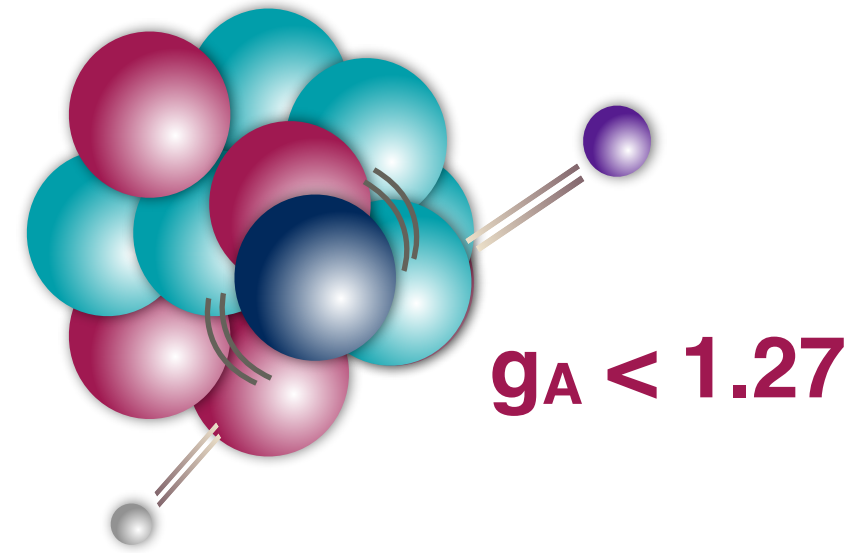
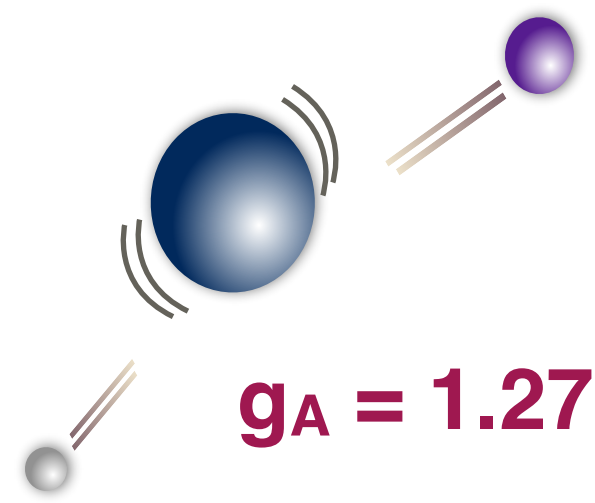
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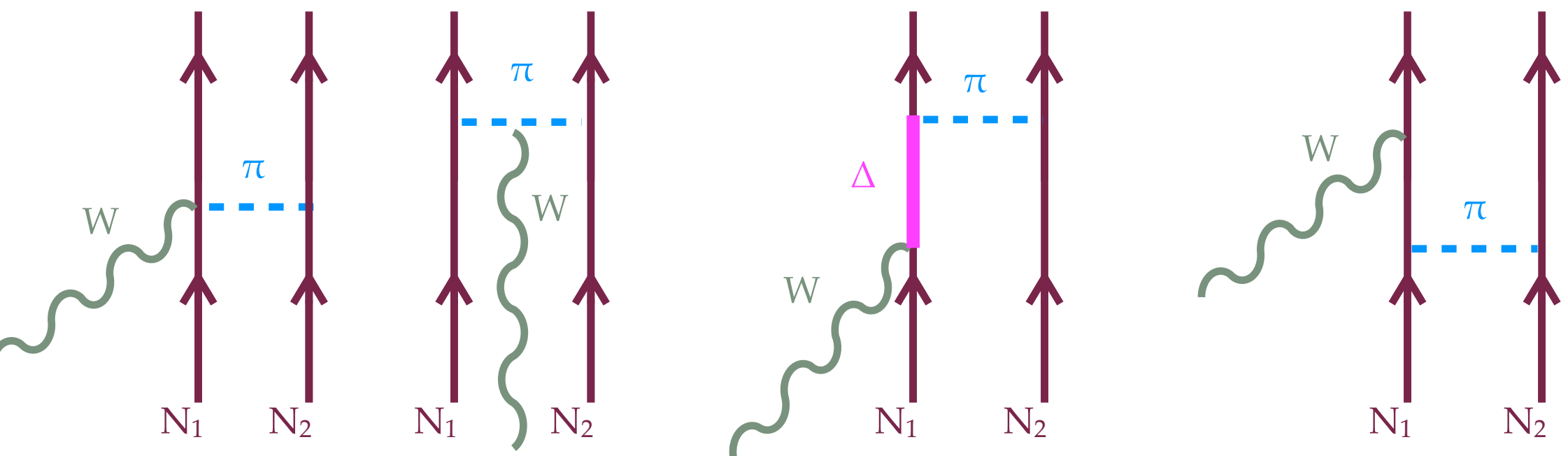


Quenching effect of the **axial vector coupling** for heavy nuclei has been observed for $2\nu\beta\beta$ and β -decay...

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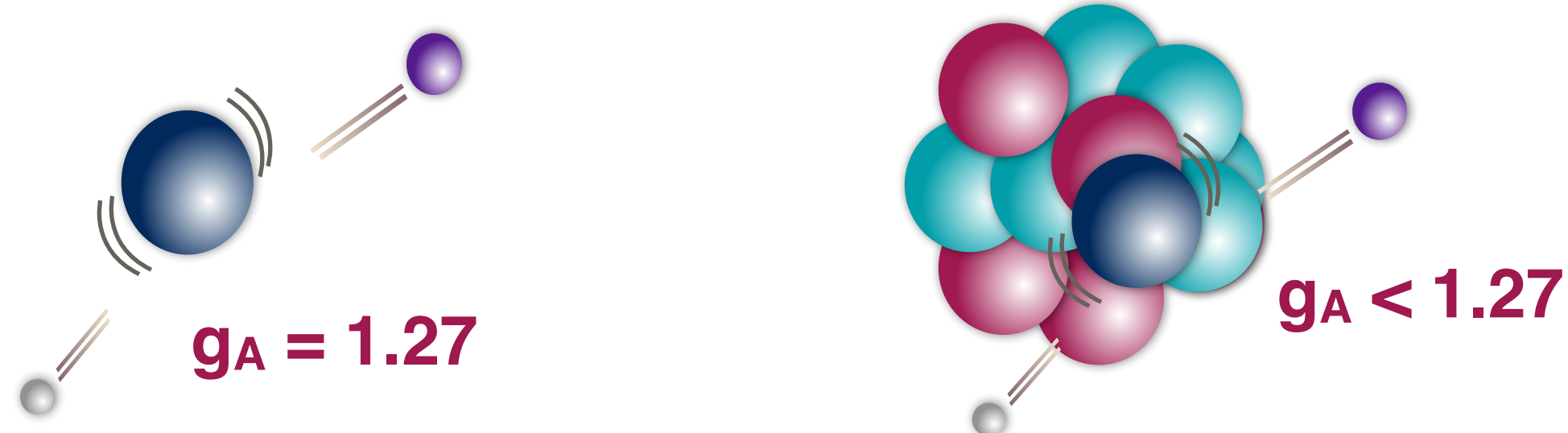
Contact/pion-in-flight

Δ -meson exchange current

Correlation

...due to **multi-nucleon effects** like meson-exchange currents

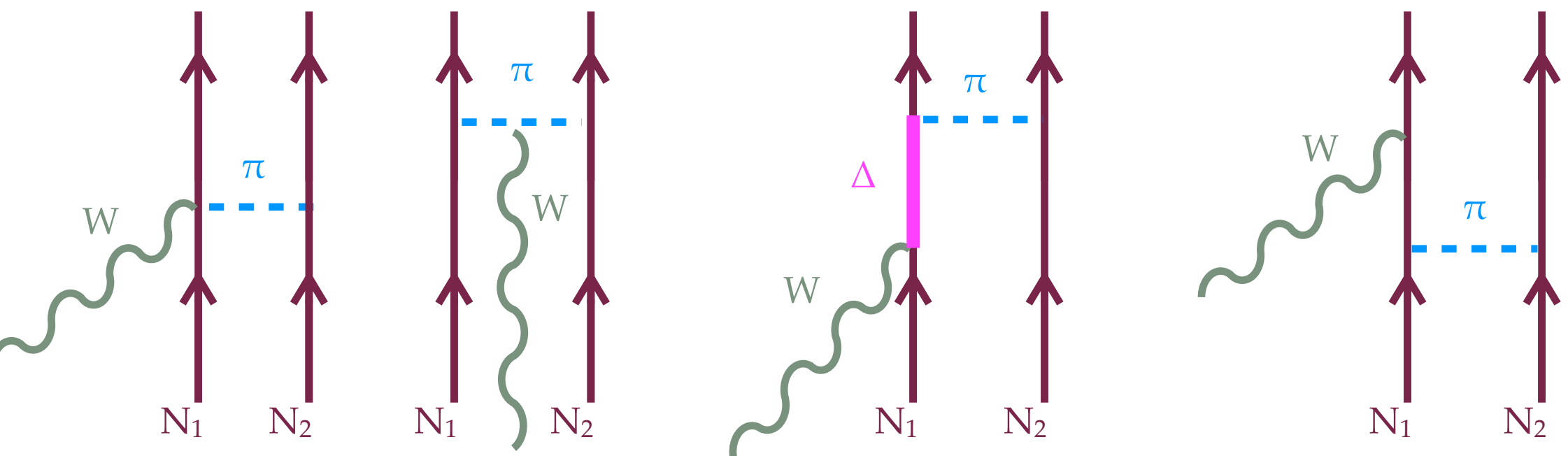
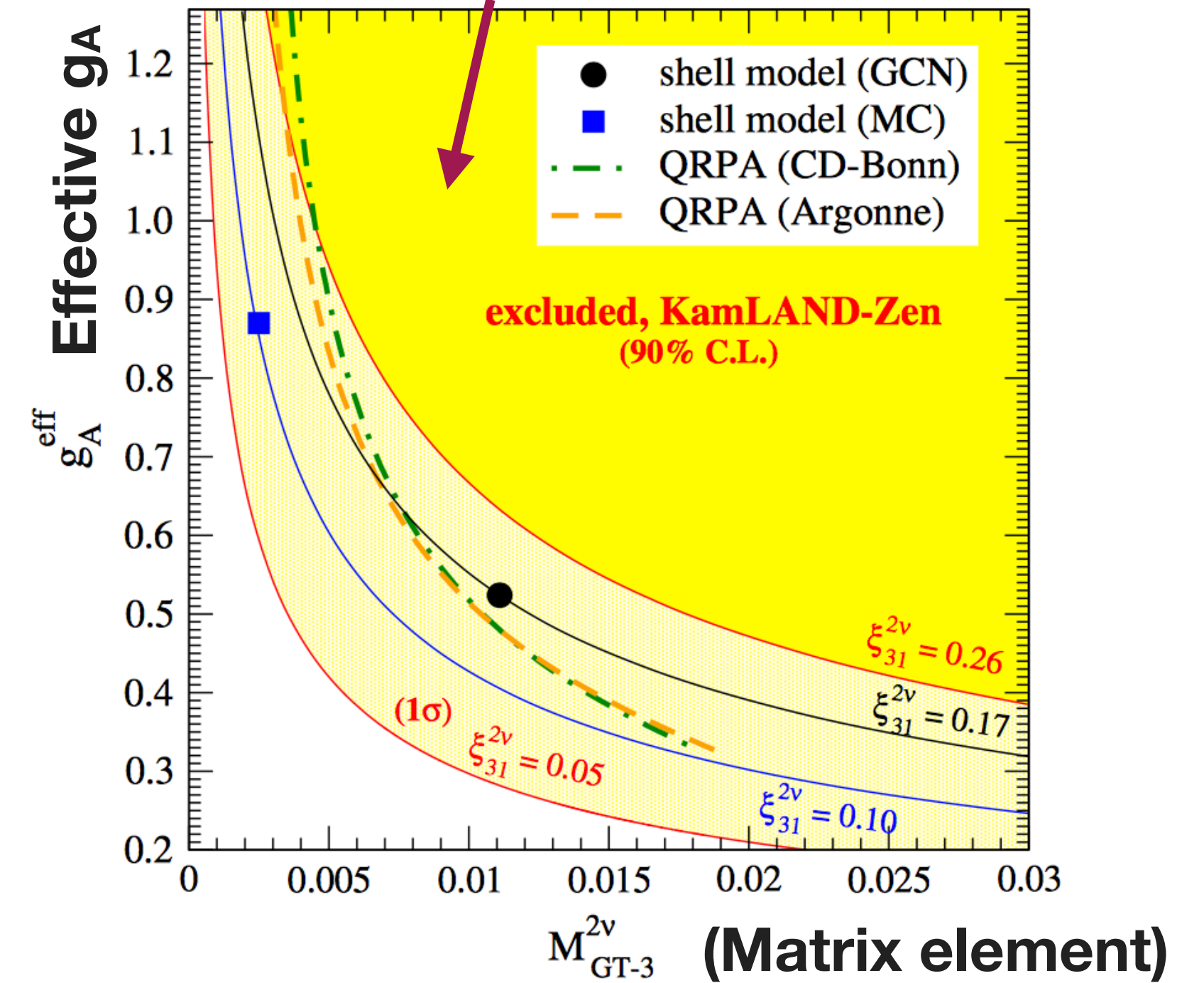
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$2\nu\beta\beta$ experiments can help set **limits**

KamLAND-Zen excludes this region



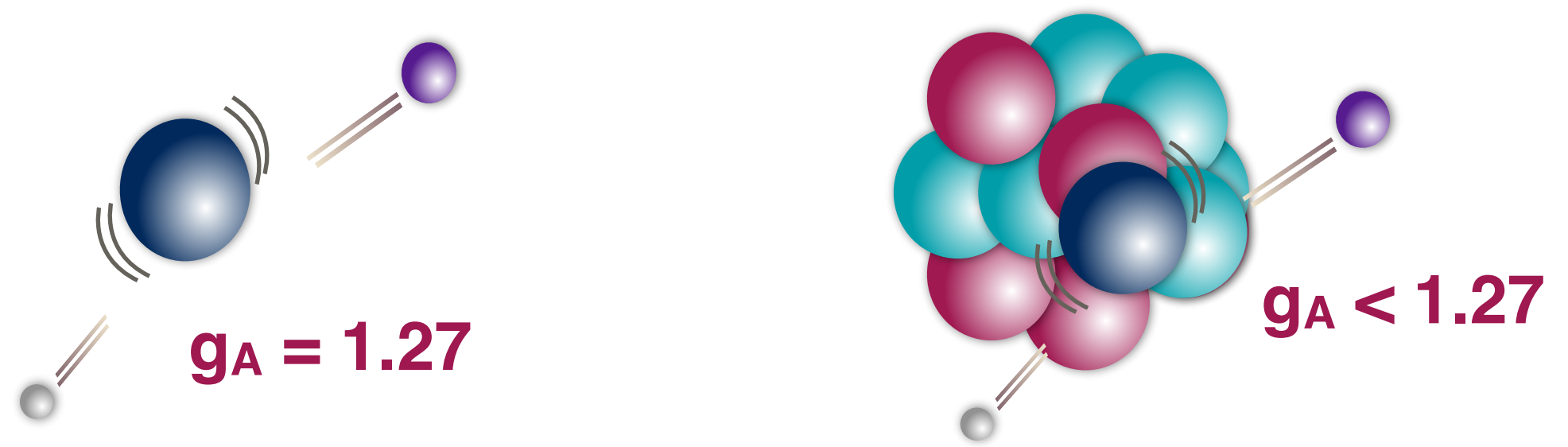
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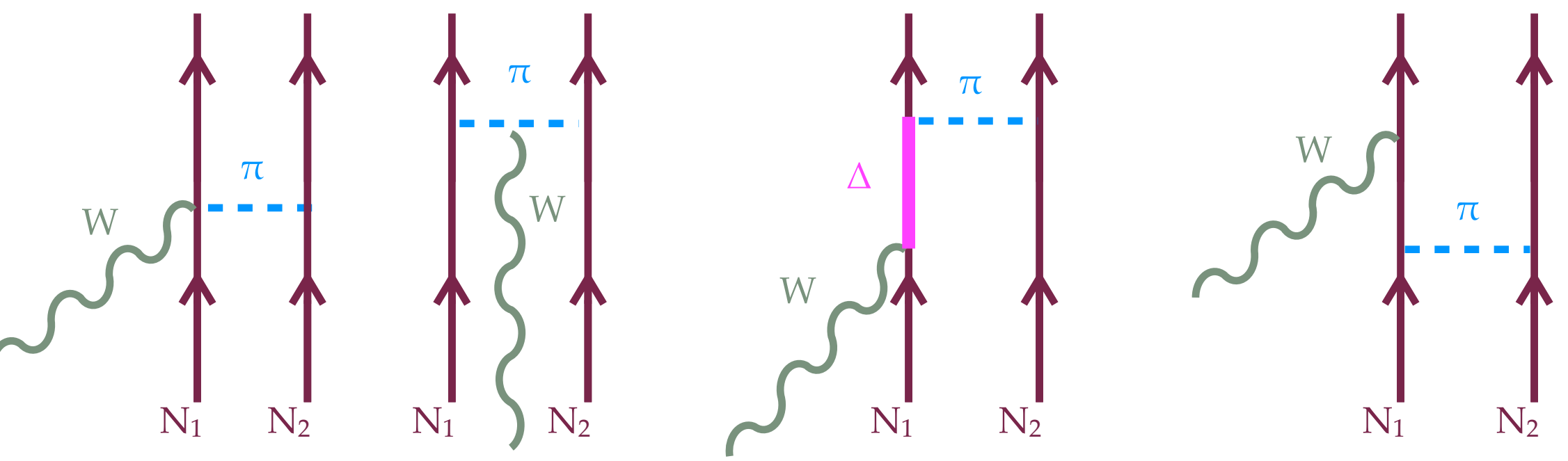
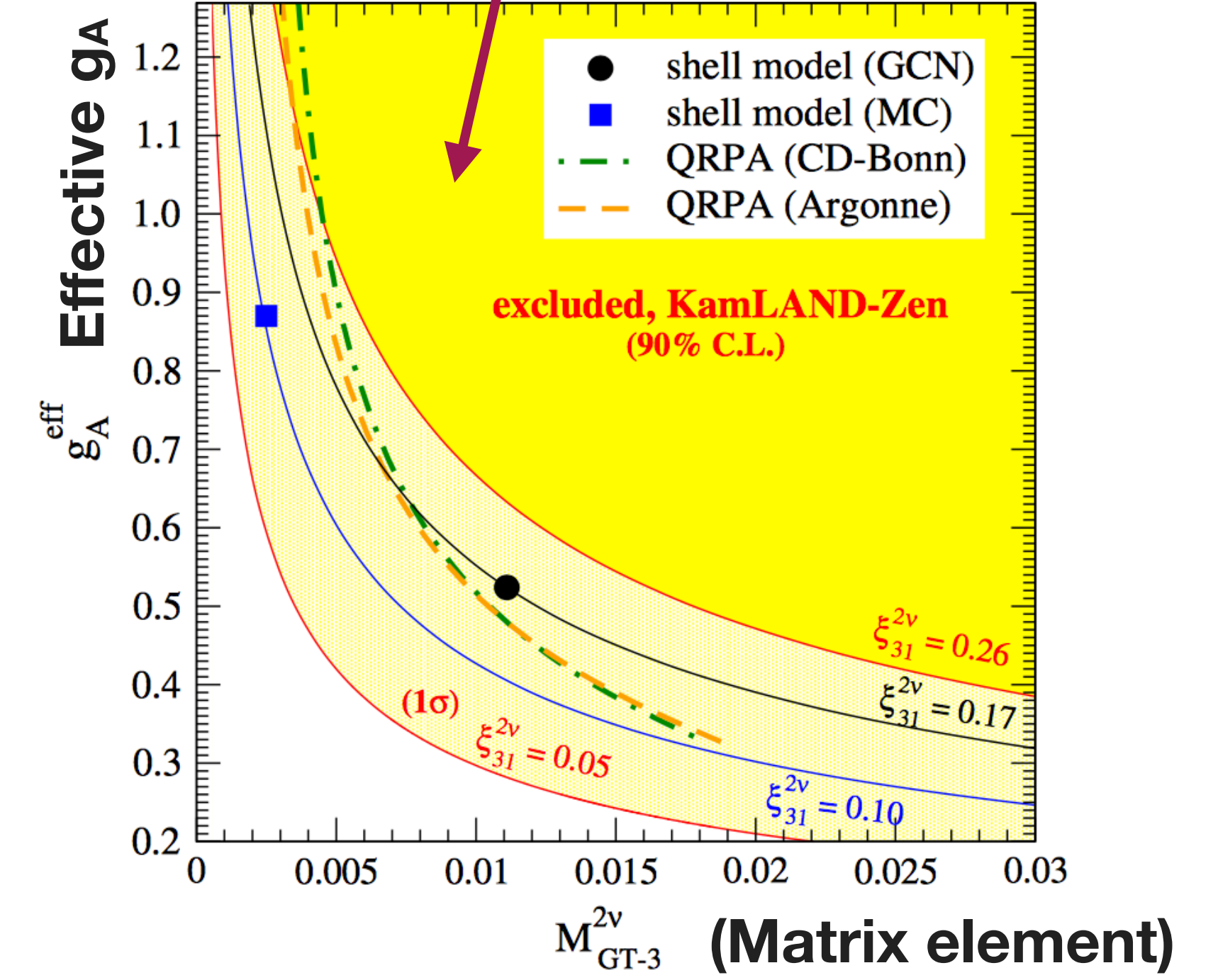
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nature physics

Discrepancy between experimental and theoretical β -decay rates resolved from first principles

New *ab initio* calculations can **model** these effects in the **nuclear matrix element**

Nature Physics volume 15, pages 428–431 (2019)

...due to **multi-nucleon effects** like meson-exchange currents

$$0\nu\beta\beta \text{ rate } \frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z) g_A^4 |M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$$

- Describe the **nuclear structure effects** of the parent, daughter and intermediate nuclei
- Different NMEs correspond to **2νββ (testable)** and each **0νββ mechanism**
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- **Biggest uncertainty in $m_{\beta\beta}$ calculation**

Rept.Prog. Phys. 80, 046301 (2017)


Nuclear matrix elements in very little detail

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Gamow-Teller
Fermi
Tensor



electron spins
parallel

g_A again!



electron spins
anti-parallel

Rept. Prog. Phys. 80, 046301 (2017)

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↑↑
electron spins parallel

g_A again!
↑↓
electron spins anti-parallel

Spherical Bessel functions

Inter-nucleon position

Neutrino potential

Operators act on nuclear states

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$$M_{0\nu}^F = \frac{2R}{\pi g_A^2} \int_0^\infty |q| d|q| \langle f | \sum_{a,b} \frac{j_0(|q|r_{ab}) h_F(|q|)}{|q| + \bar{E} - (E_i + E_f)/2} \tau_a^+ \tau_b^+ | i \rangle ,$$

$$M_{0\nu}^T = \frac{2R}{\pi g_A^2} \int_0^\infty |q| d|q| \langle f | \sum_{a,b} \frac{j_2(|q|r_{ab}) h_T(|q|) [3\sigma_j \cdot \hat{r}_{ab} \sigma_k \cdot \hat{r}_{ab} - \sigma_a \cdot \sigma_b]}{|q| + \bar{E} - (E_i + E_f)/2} \tau_a^+ \tau_b^+ | i \rangle$$

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electron spins parallel

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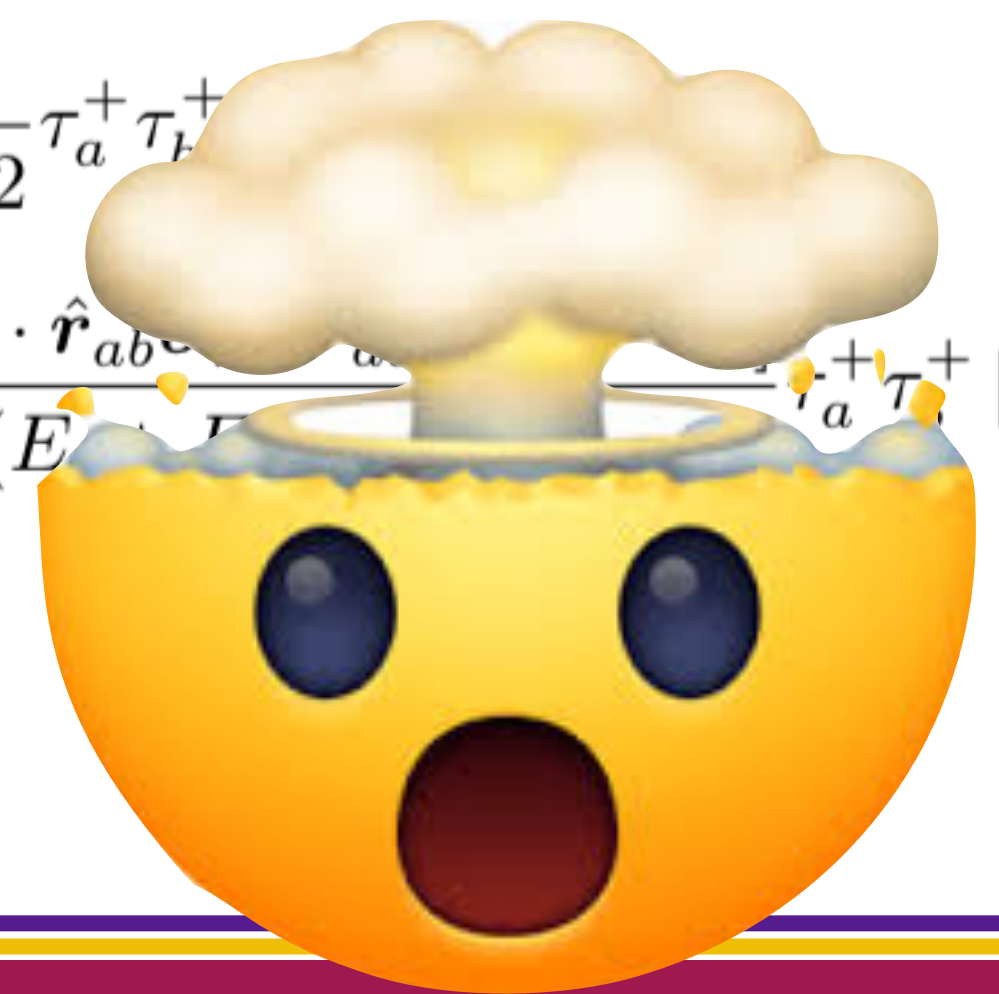
Neutrino potential

Operators act on nuclear states

$$M_{0\nu}^{GT} = \frac{2R}{\pi g_A^2} \int_0^\infty |\mathbf{q}| d|\mathbf{q}| \langle f | \sum_{a,b} \frac{j_0(|\mathbf{q}|r_{ab}) h_{GT}(|\mathbf{q}|) \sigma_a \cdot \sigma_b}{|\mathbf{q}| + \bar{E} - (E_i + E_f)/2} \tau_a^+ \tau_b^+ | i \rangle ,$$

$$M_{0\nu}^F = \frac{2R}{\pi g_A^2} \int_0^\infty |\mathbf{q}| d|\mathbf{q}| \langle f | \sum_{a,b} \frac{j_0(|\mathbf{q}|r_{ab}) h_F(|\mathbf{q}|)}{|\mathbf{q}| + \bar{E} - (E_i + E_f)/2} \tau_a^+ \tau_b^+ | i \rangle ,$$

$$M_{0\nu}^T = \frac{2R}{\pi g_A^2} \int_0^\infty |\mathbf{q}| d|\mathbf{q}| \langle f | \sum_{a,b} \frac{j_2(|\mathbf{q}|r_{ab}) h_T(|\mathbf{q}|) [3\sigma_j \cdot \hat{r}_{ab}]}{|\mathbf{q}| + \bar{E} - (E_i + E_f)/2} \tau_a^+ \tau_b^+ | i \rangle ,$$



Rept.Prog. Phys. 80, 046301 (2017)

Two approaches: both difficult

QRPA: Quasi-particle Random Phase Approximation

NSM: Nuclear Shell Model

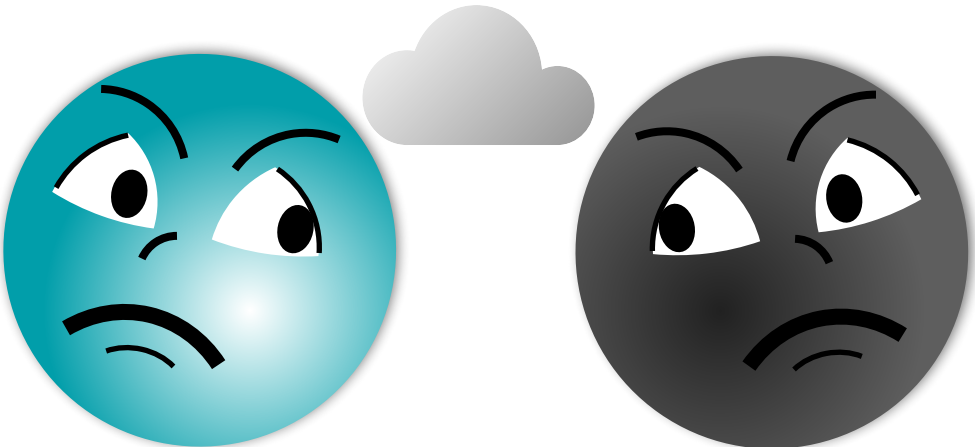
Ann.Rev.Nucl.Part.Sci.52:115-151,2002

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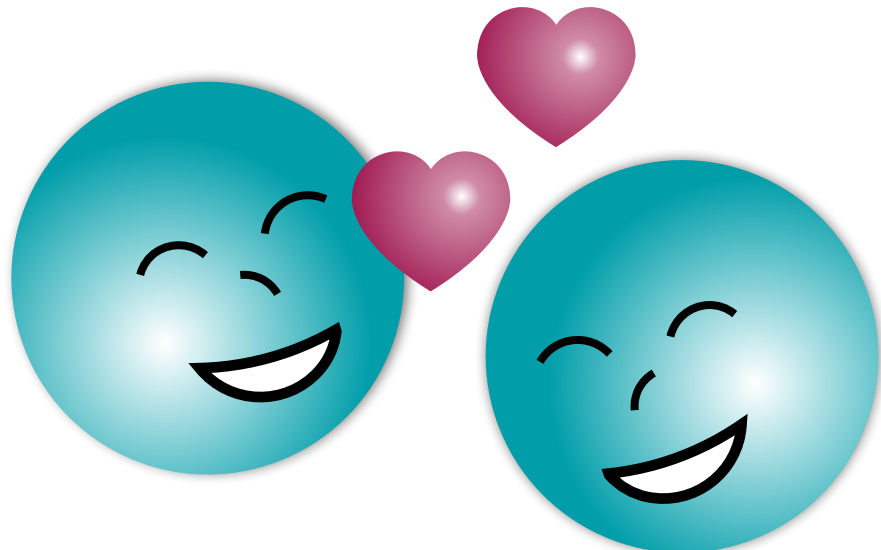
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Most popular method, relying on **small deviations** from ground state, and modeling two main ingredients:



Repulsive particle-hole spin-isospin interaction



Attractive particle-particle interaction

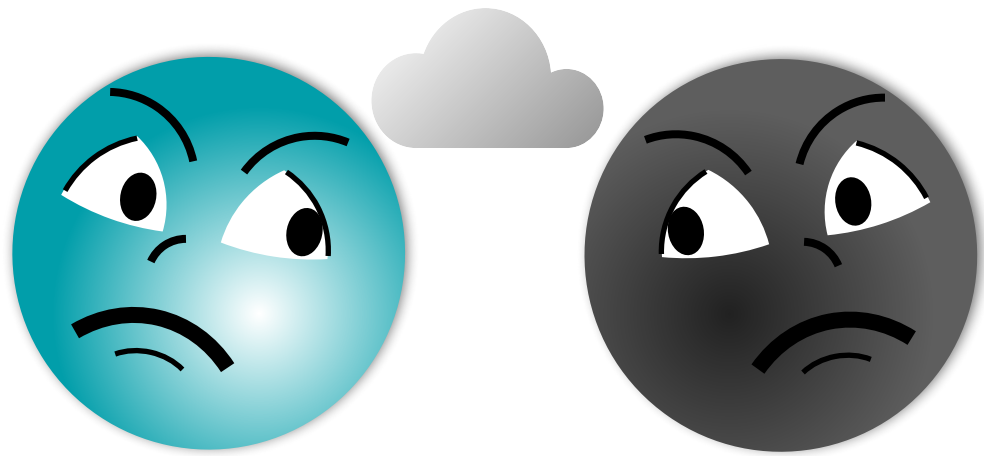
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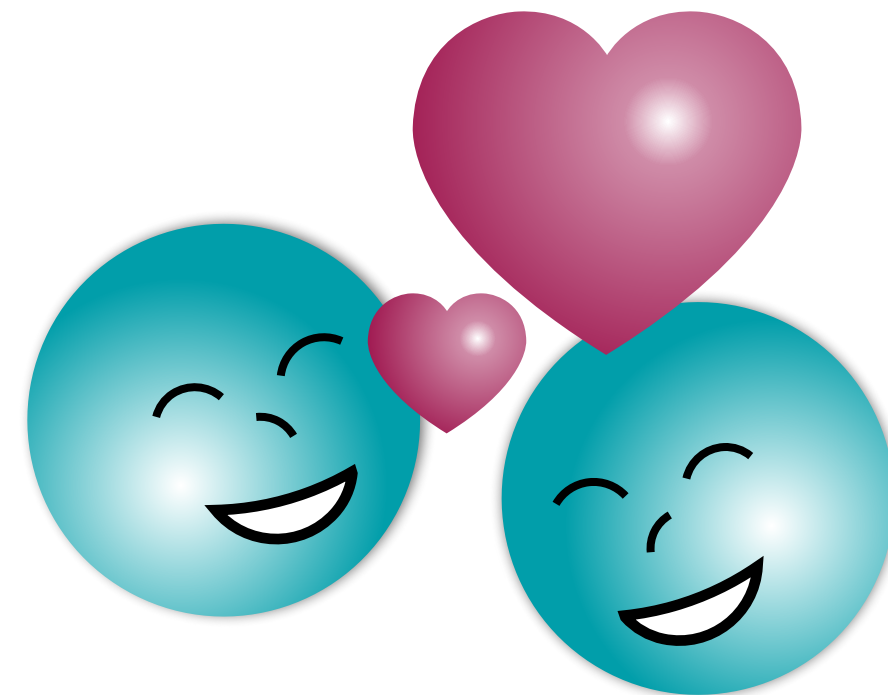
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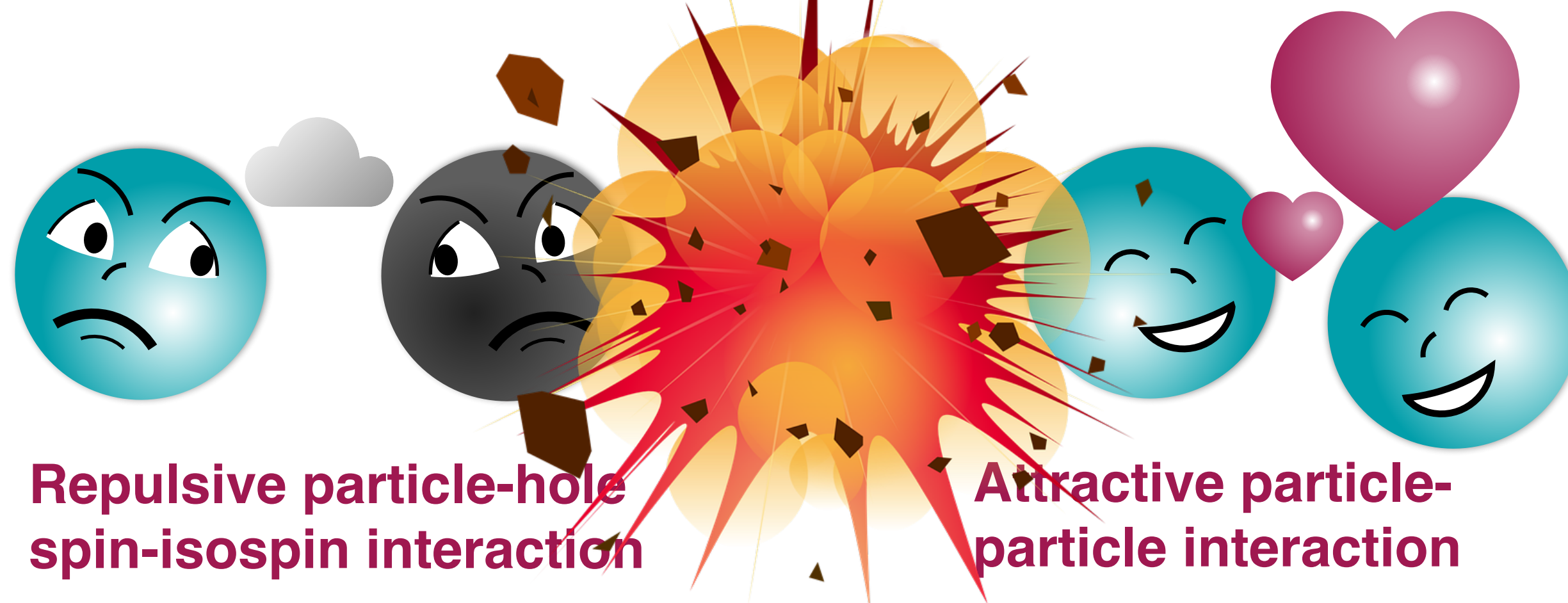
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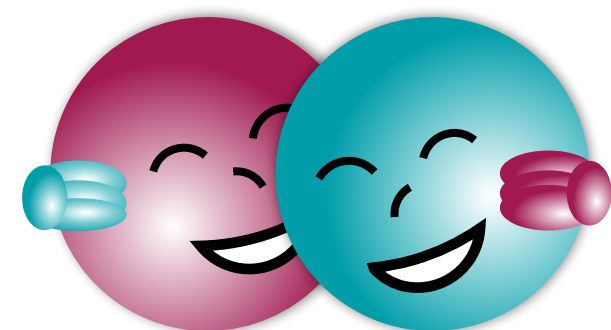
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Basic versions don't include **multi-particle** states



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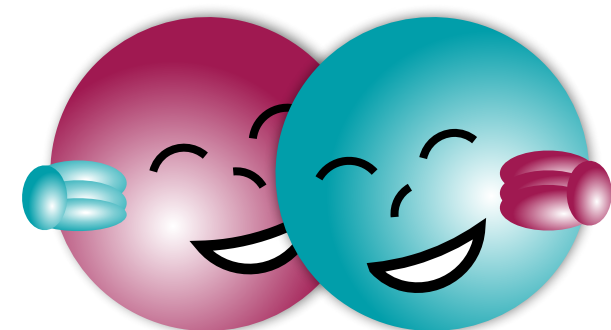
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NSM: Nuclear Shell Model

Theoretically **better** method but limited by computing power for all but **lighter isotopes** (up to ^{48}Ca)

- Start from set of all **single-particle states**
- Find **effective Hamiltonian** based on modified nucleon-nucleon interaction to describe them
- Use all configurations, including **correlations** to diagonalize Hamiltonian and **extract NME**

Ann.Rev.Nucl.Part.Sci.52:115-151,2002

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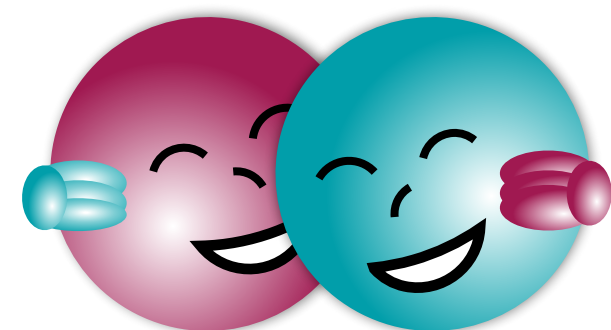


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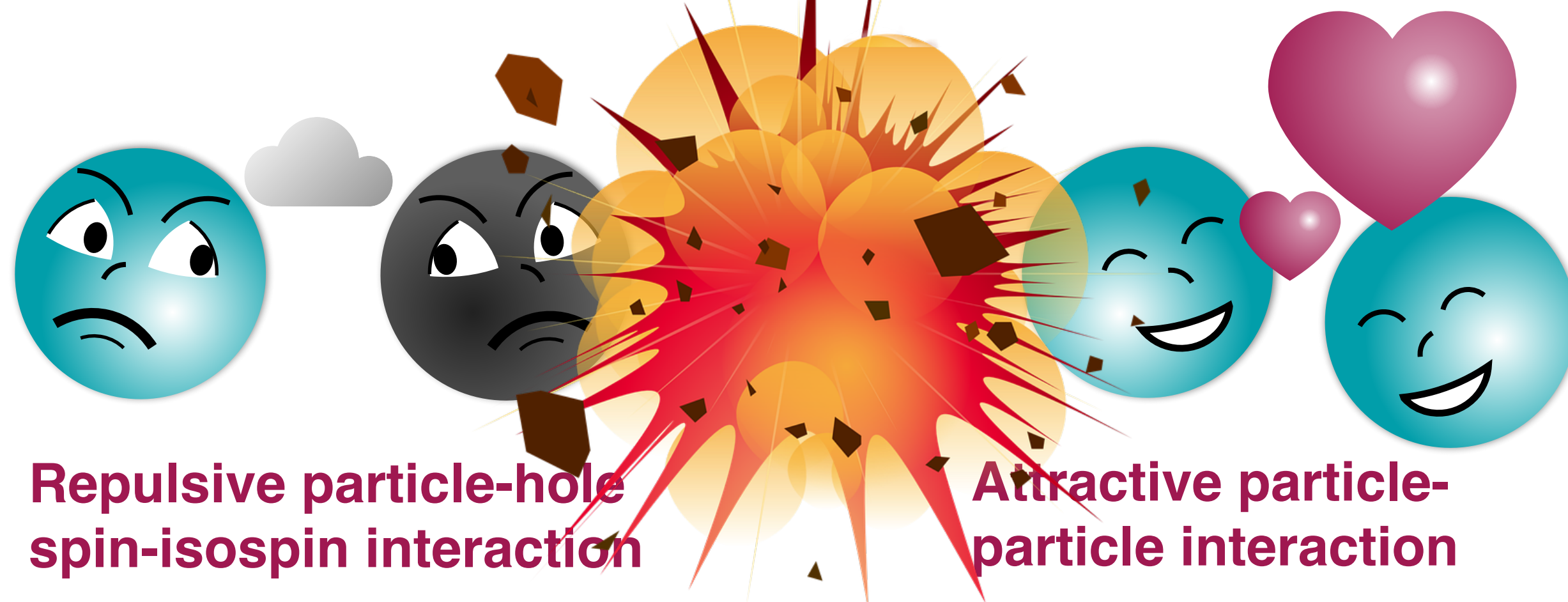


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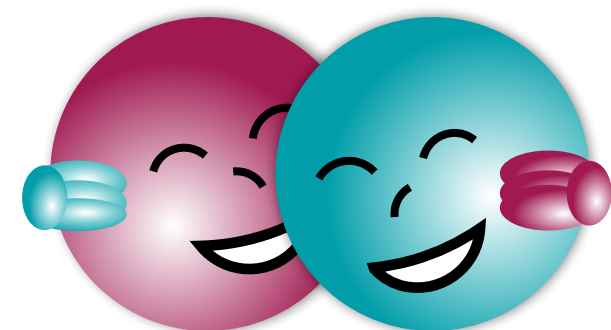


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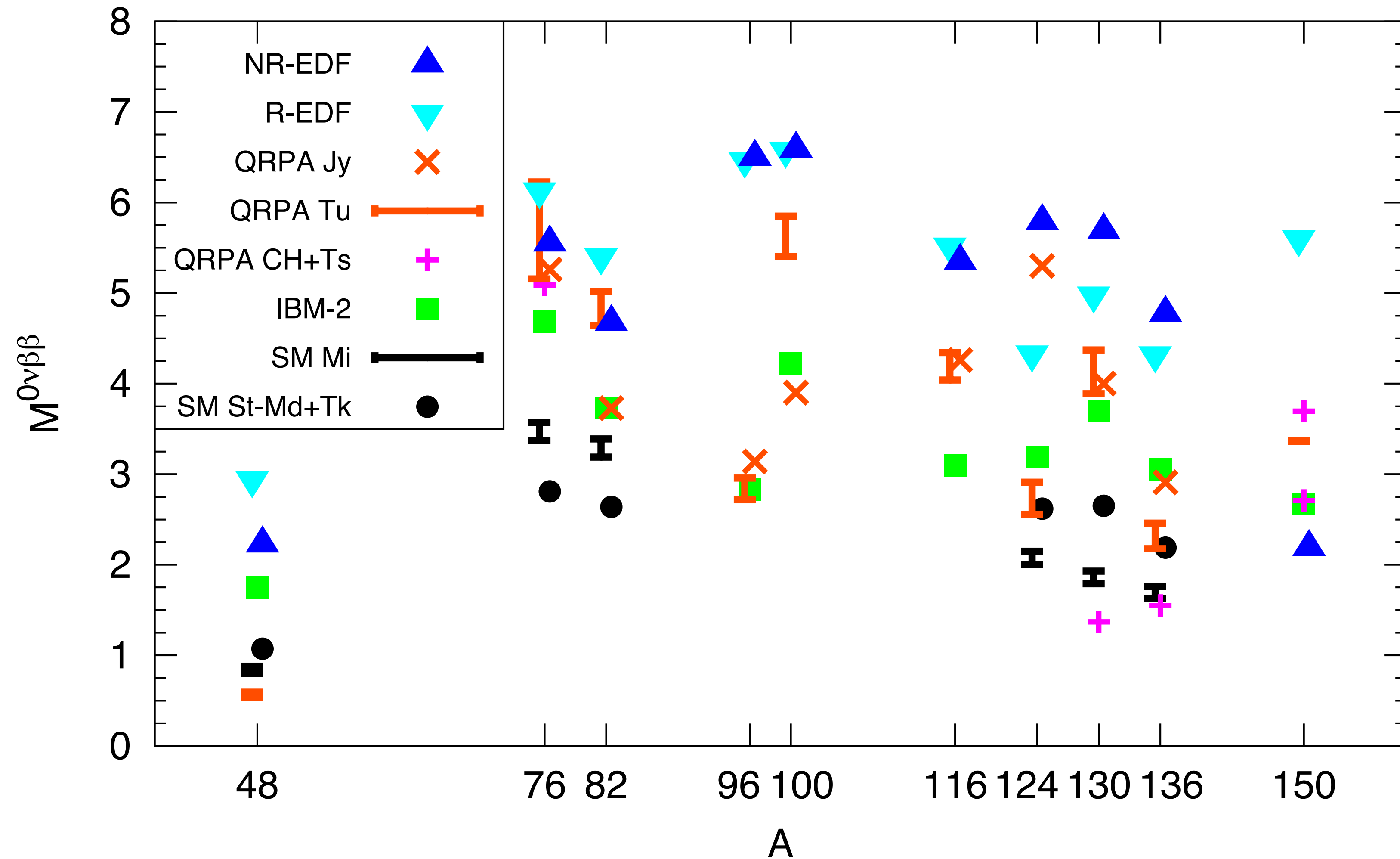
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Computationally impossible - have to use a subset

Not clear what effective operators will simulate the excluded states

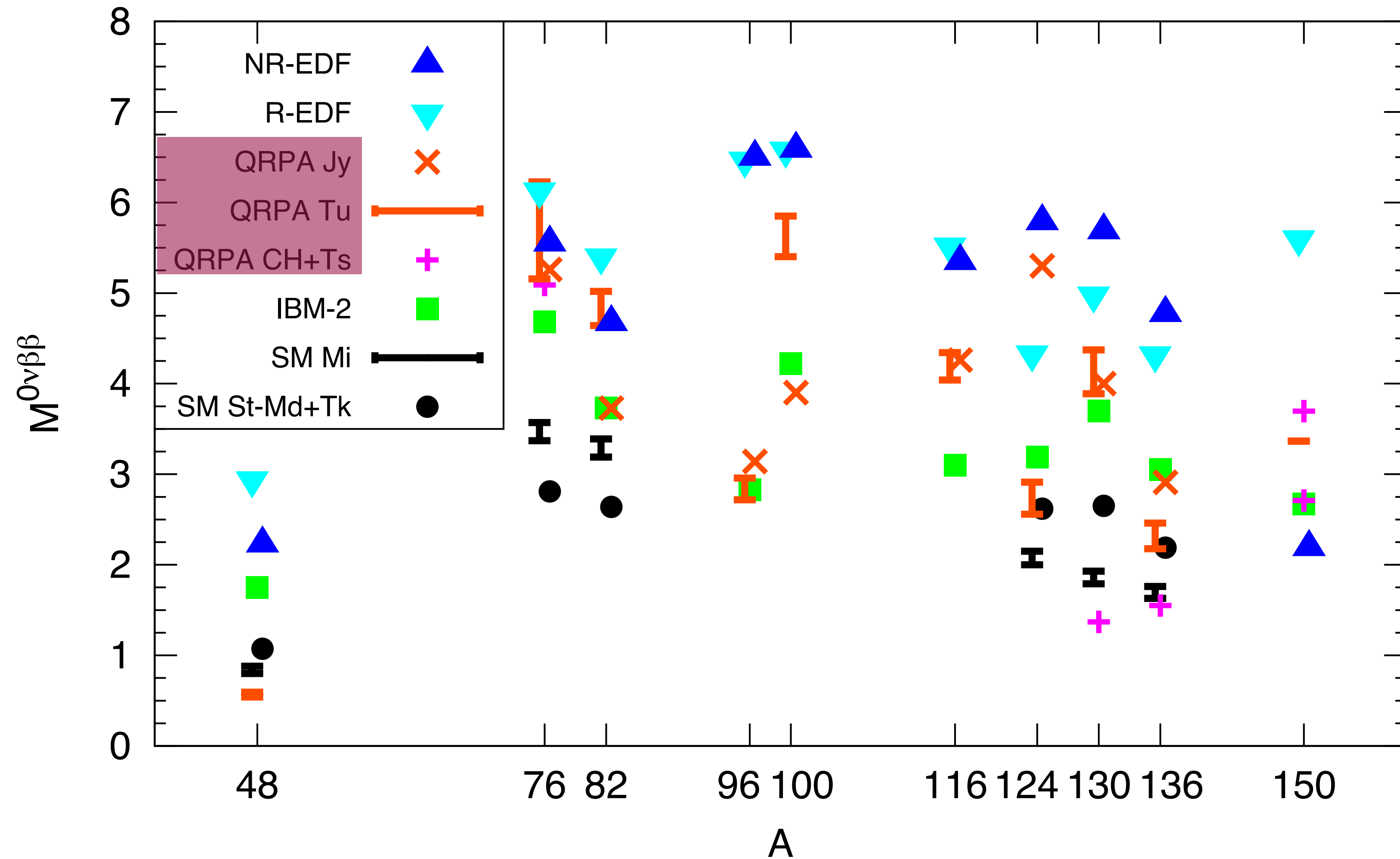
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- Use all configurations, including **correlations** to diagonalize Hamiltonian and **extract NME**

Different calculations give very different results



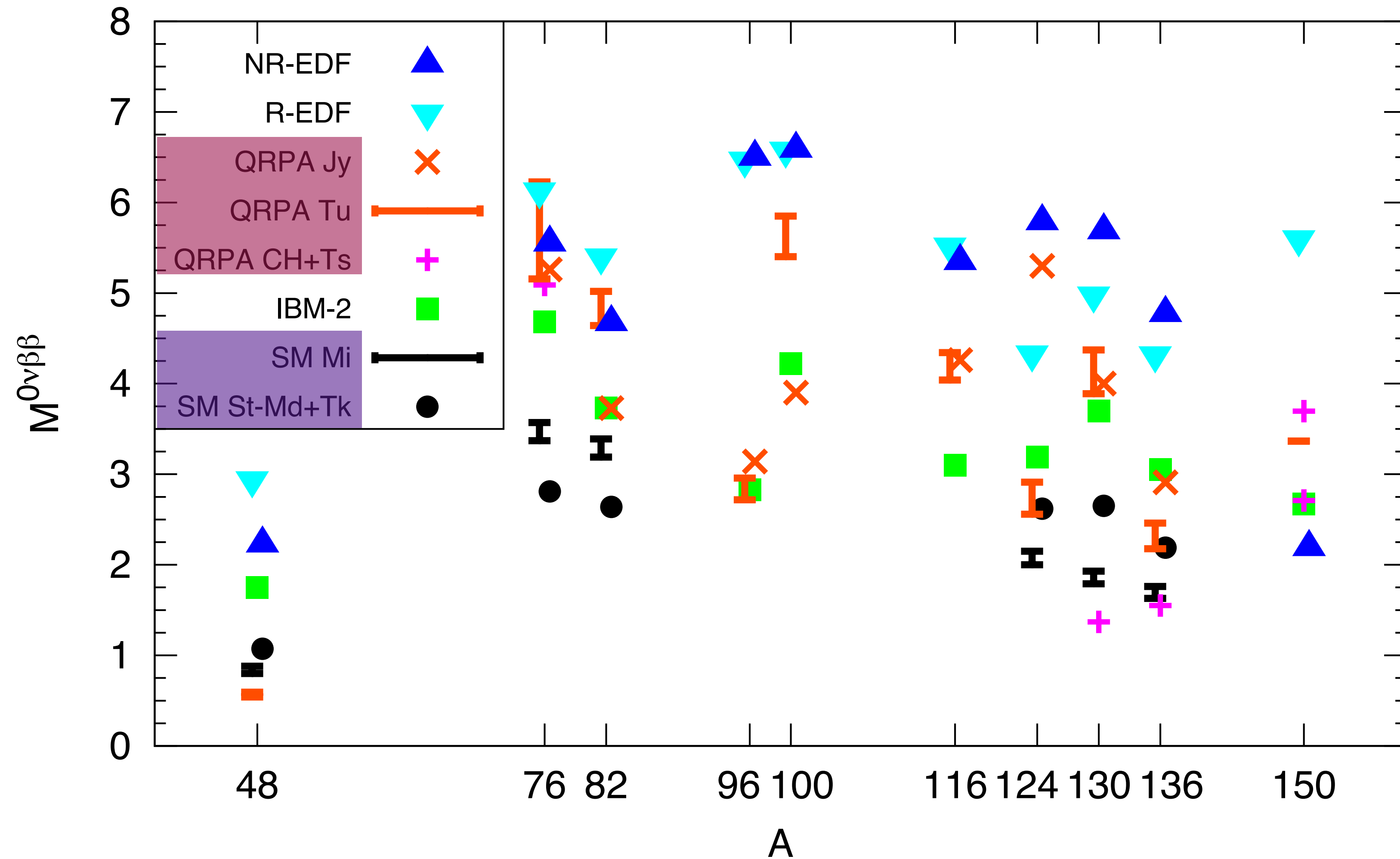
NME's calculated by various groups using:

Different calculations give very different results



NME's calculated by various groups using:
QRPA

Different calculations give very different results

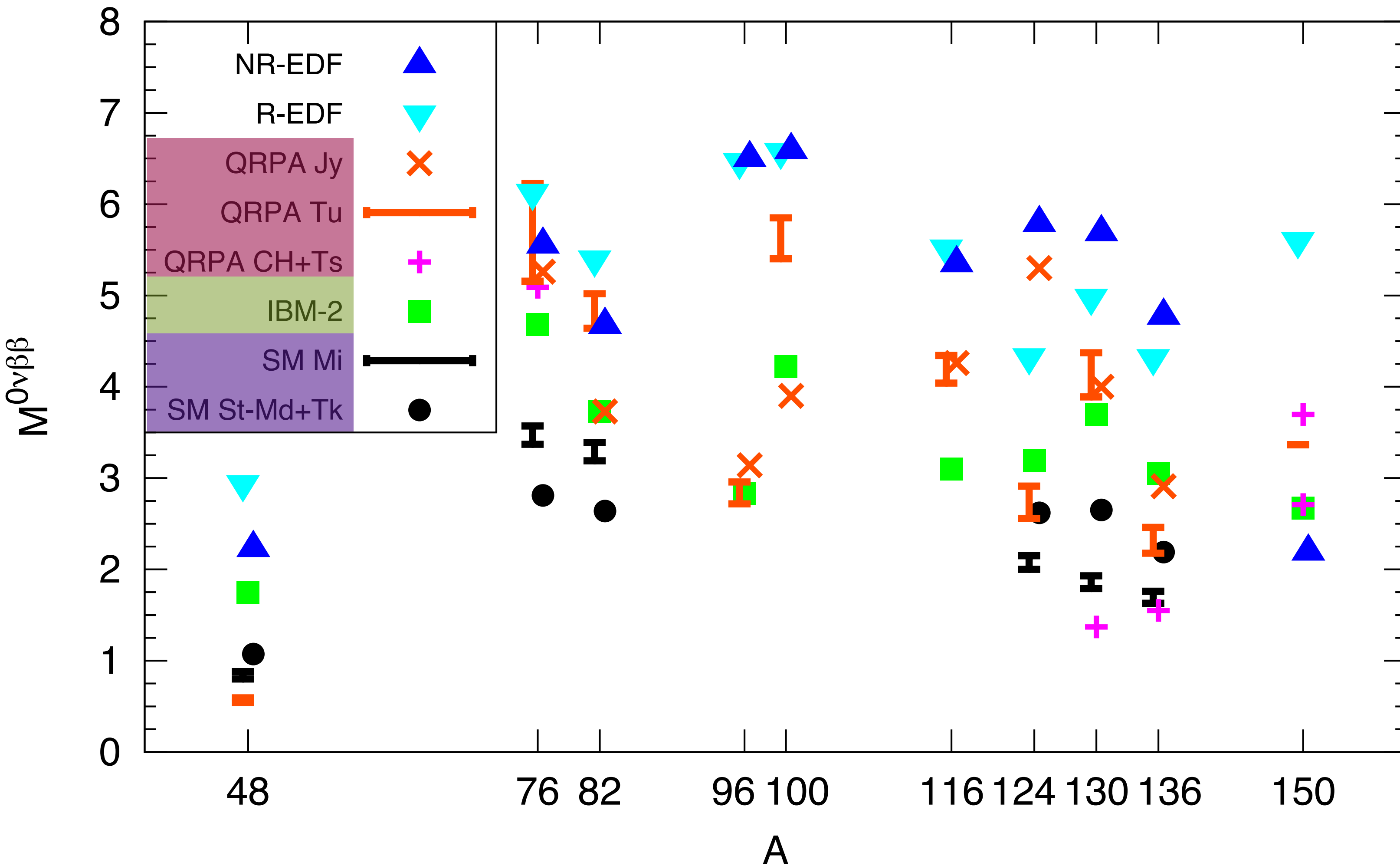


NME's calculated by various groups using:

- QRPA**
- Shell Model**

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Different calculations give very different results



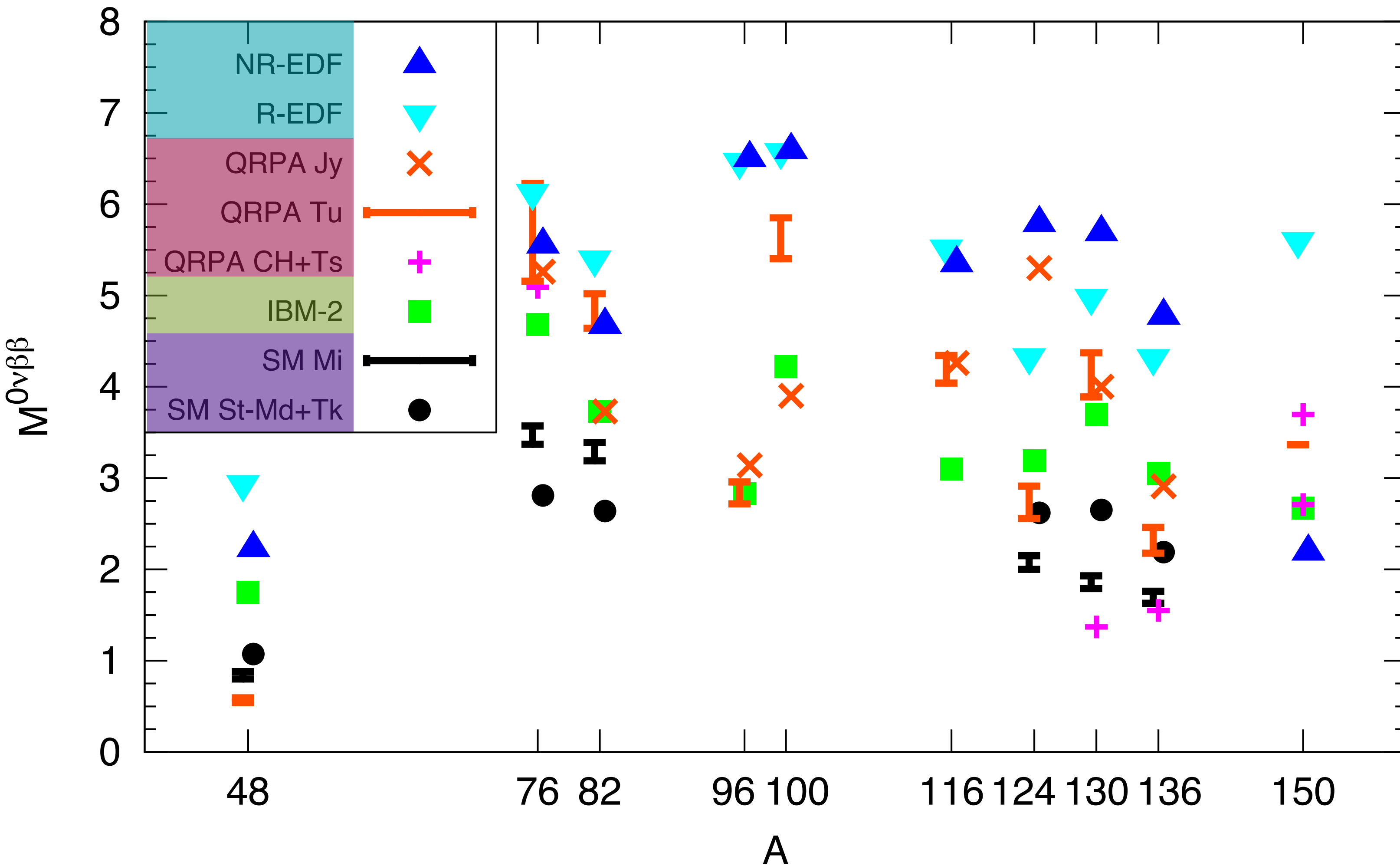
NME's calculated by various groups using:

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- Interacting boson model**

Variation on shell model that aims to describe collective excitations up to heavy nuclei, modelling low-energy states as bosons

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Different calculations give very different results



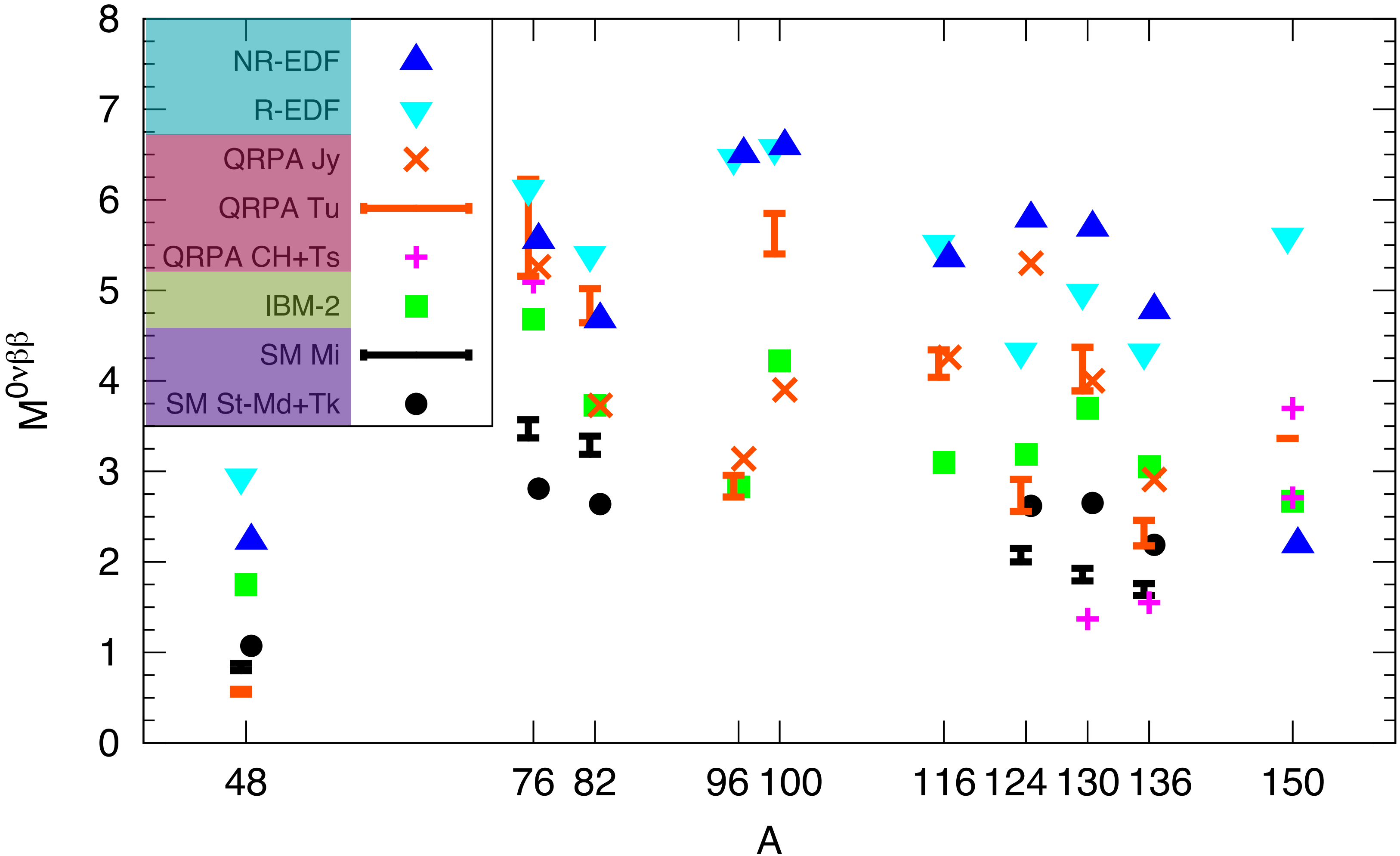
NME's calculated by various groups using:

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- Interacting boson model**
- Energy density functional theory**

Can be used when QRPA breaks down

Rept.Prog. Phys. 80, 046301 (2017)

Different calculations give very different results

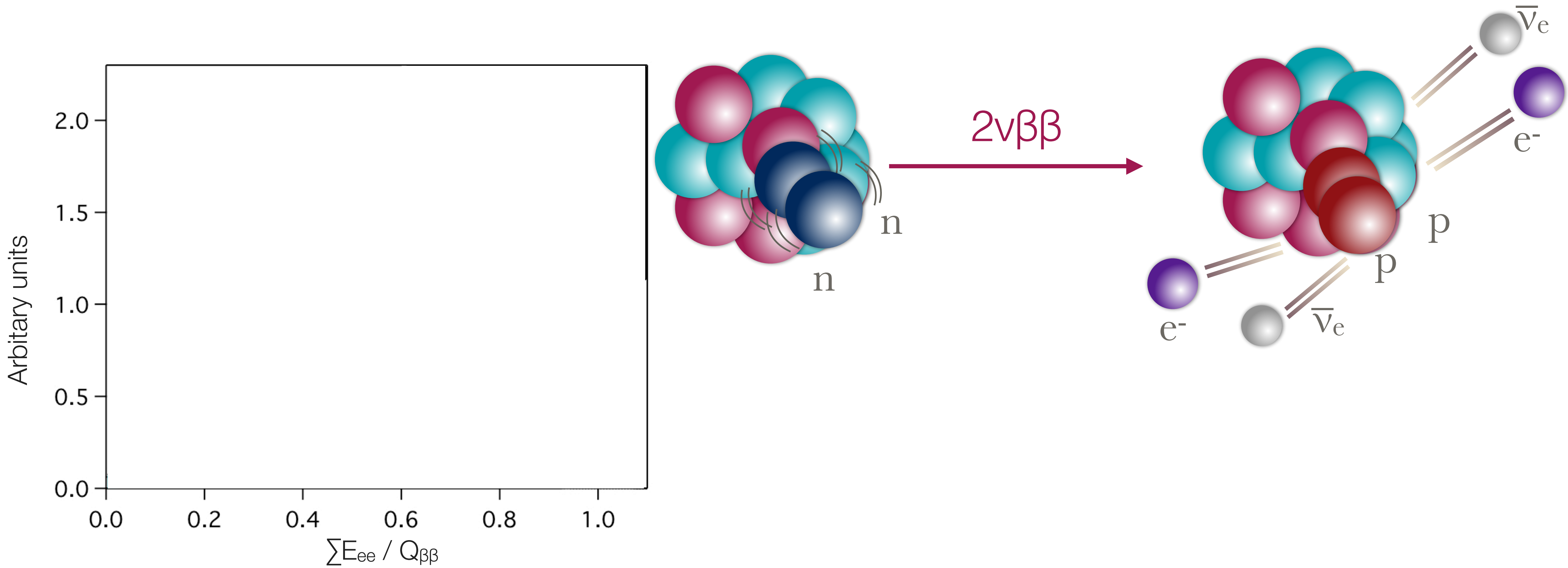


NME's calculated by various groups using:

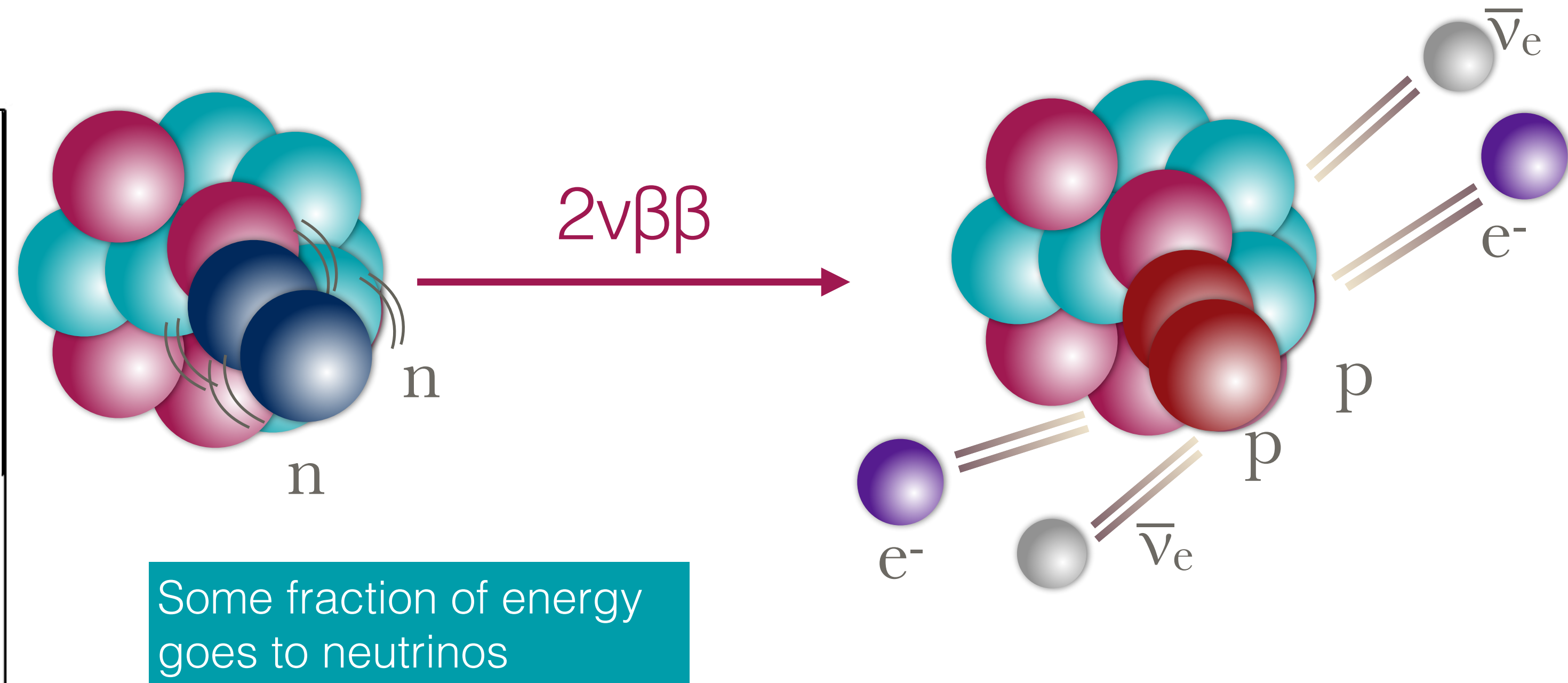
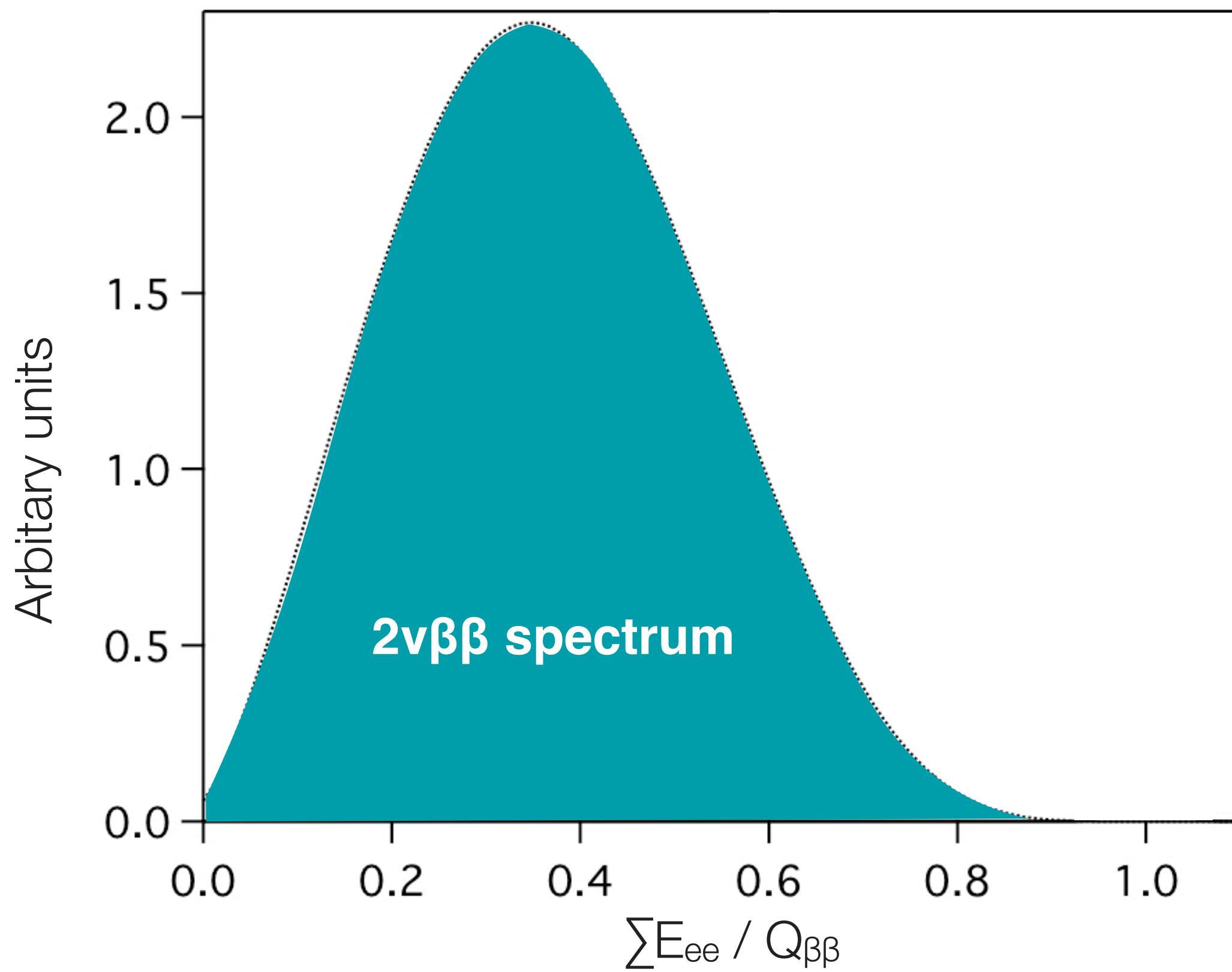
- QRPA
- Shell Model
- Interacting boson model
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Calculated results can vary by a factor of 3 or more

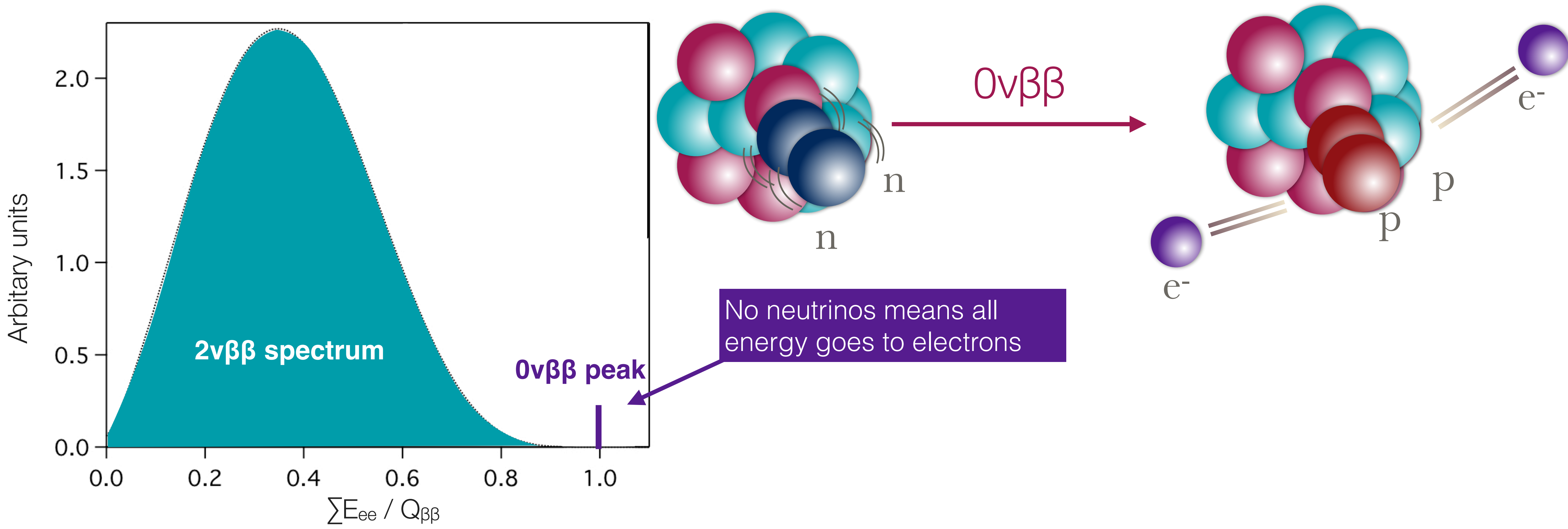
Rept.Prog. Phys. 80, 046301 (2017)



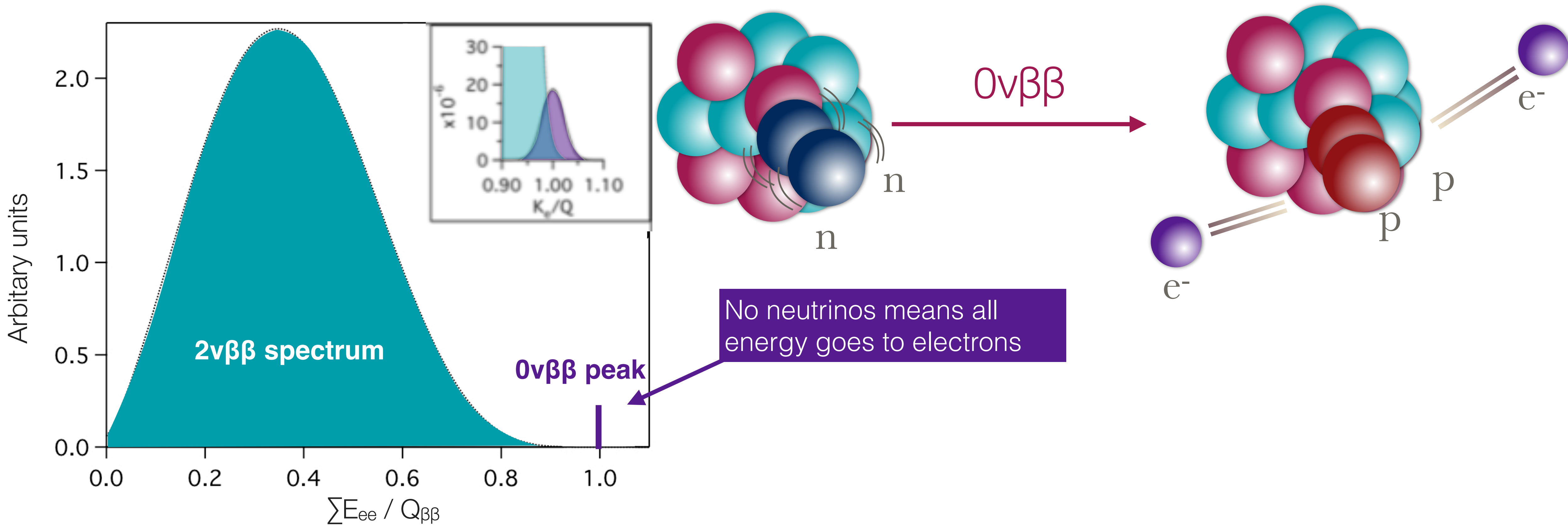
Sum of the **2 electron energies**, as fraction of $\beta\beta$ decay energy



Sum of the **2 electron energies**, as fraction of $\beta\beta$ decay energy



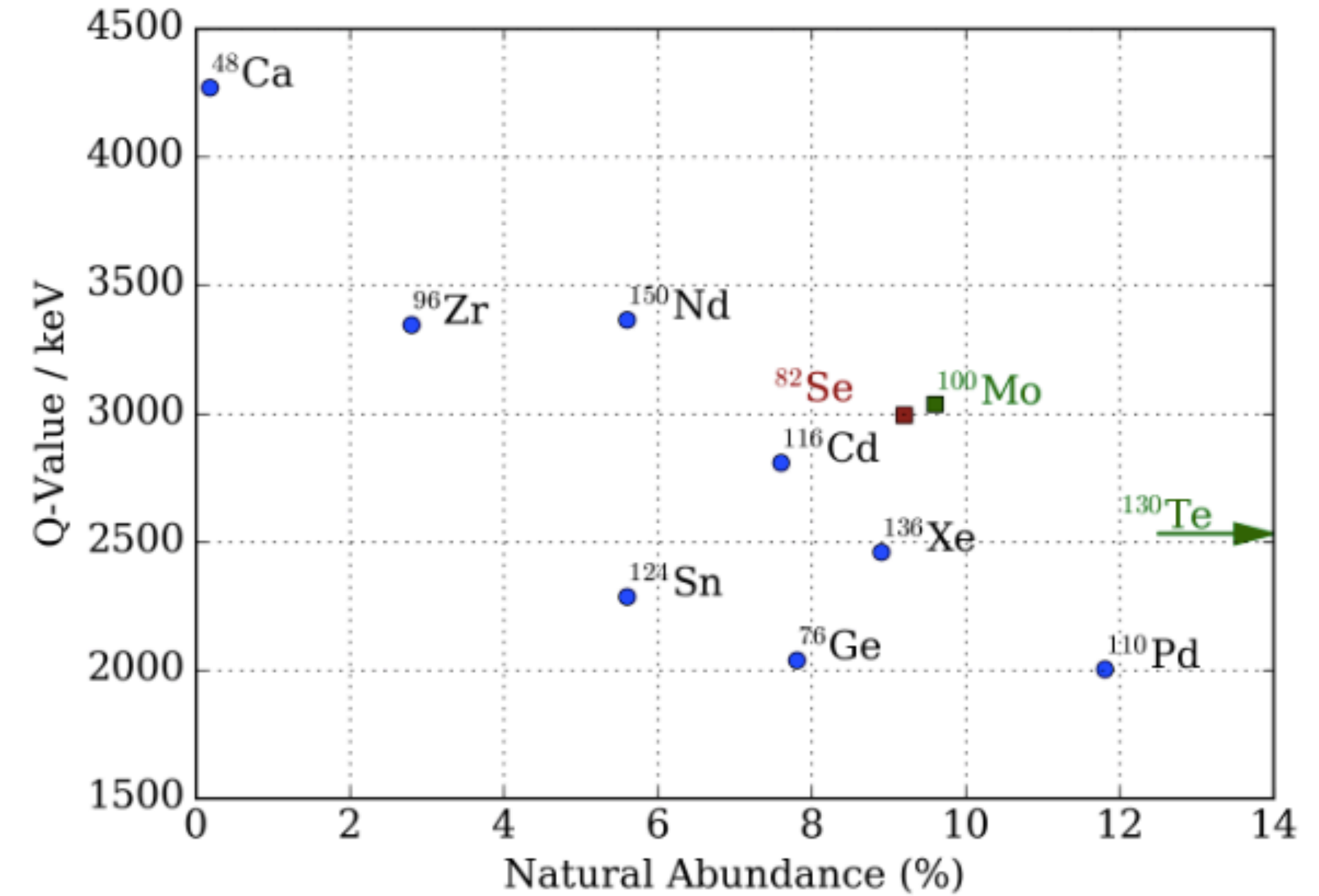
Sum of the **2 electron energies**, as fraction of $\beta\beta$ decay energy



Sum of the **2 electron energies**, as fraction of $\beta\beta$ decay energy

Choosing an isotope

1 H 1.008																	2 He 4.0026
3 Li 6.94	4 Be 9.0122											5 B 10.81	6 C 12.011	7 N 14.007	8 O 15.999	9 F 18.998	10 Ne 20.180
11 Na 22.990	12 Mg 24.305											13 Al 26.982	14 Si 28.085	15 P 30.974	16 S 32.06	17 Cl 35.45	18 Ar 39.948
19 K 39.098	20 Ca 40.08	21 Sc 44.956	22 Ti 47.867	23 V 50.942	24 Cr 51.996	25 Mn 54.938	26 Fe 55.845	27 Co 58.933	28 Ni 58.693	29 Cu 63.546	30 Zn 65.38	31 Ga 69.723	32 Ge 72.64	33 As 74.922	34 Se 78.96	35 Br 79.904	36 Kr 83.798
37 Rb 85.468	38 Sr 87.62	39 Y 88.906	40 Zr 91.224	41 Nb 92.906	42 Mo 95.94	43 Tc (98)	44 Ru 101.07	45 Rh 102.91	46 Pd 106.92	47 Ag 107.87	48 Cd 112.41	49 In 114.82	50 Sn 118.71	51 Sb 121.76	52 Te 127.6	53 I 126.90	54 Xe 131.29
55 Cs 132.91	56 Ba 137.33	57-71 * Lanthanide series	72 Hf 178.49	73 Ta 180.95	74 W 183.84	75 Re 186.21	76 Os 190.23	77 Ir 192.22	78 Pt 195.08	79 Au 196.97	80 Hg 200.59	81 Tl 204.38	82 Pb 207.2	83 Bi 208.98	84 Po (209)	85 At (210)	86 Rn (222)
87 Fr (223)	88 Ra (226)	89-103 * Actinide series	104 Rf (261)	105 Db (262)	106 Sg (263)	107 Bh (264)	108 Hs (265)	109 Mt (266)	110 Ds (267)	111 Rg (268)	112 Cn (269)	113 Nh (270)	114 Fl (271)	115 Mc (272)	116 Lv (273)	117 Ts (274)	118 Og (284)
			57 La 138.91	58 Ce 140.12	59 Pr 140.91	60 Nd 144.24	61 Pm (145)	62 Sm 150.36	63 Eu 151.96	64 Gd 157.25	65 Tb 158.93	66 Dy 162.50	67 Ho 164.93	68 Er 167.26	69 Tm 168.93	70 Yb 173.05	71 Lu 174.97
			89 Ac (227)	90 Th 232.04	91 Pa 231.04	92 U 238.03	93 Np (237)	94 Pu (244)	95 Am (243)	96 Cm (247)	97 Bk (247)	98 Cf (251)	99 Es (252)	100 Fm (257)	101 Md (258)	102 No (259)	103 Lr (262)

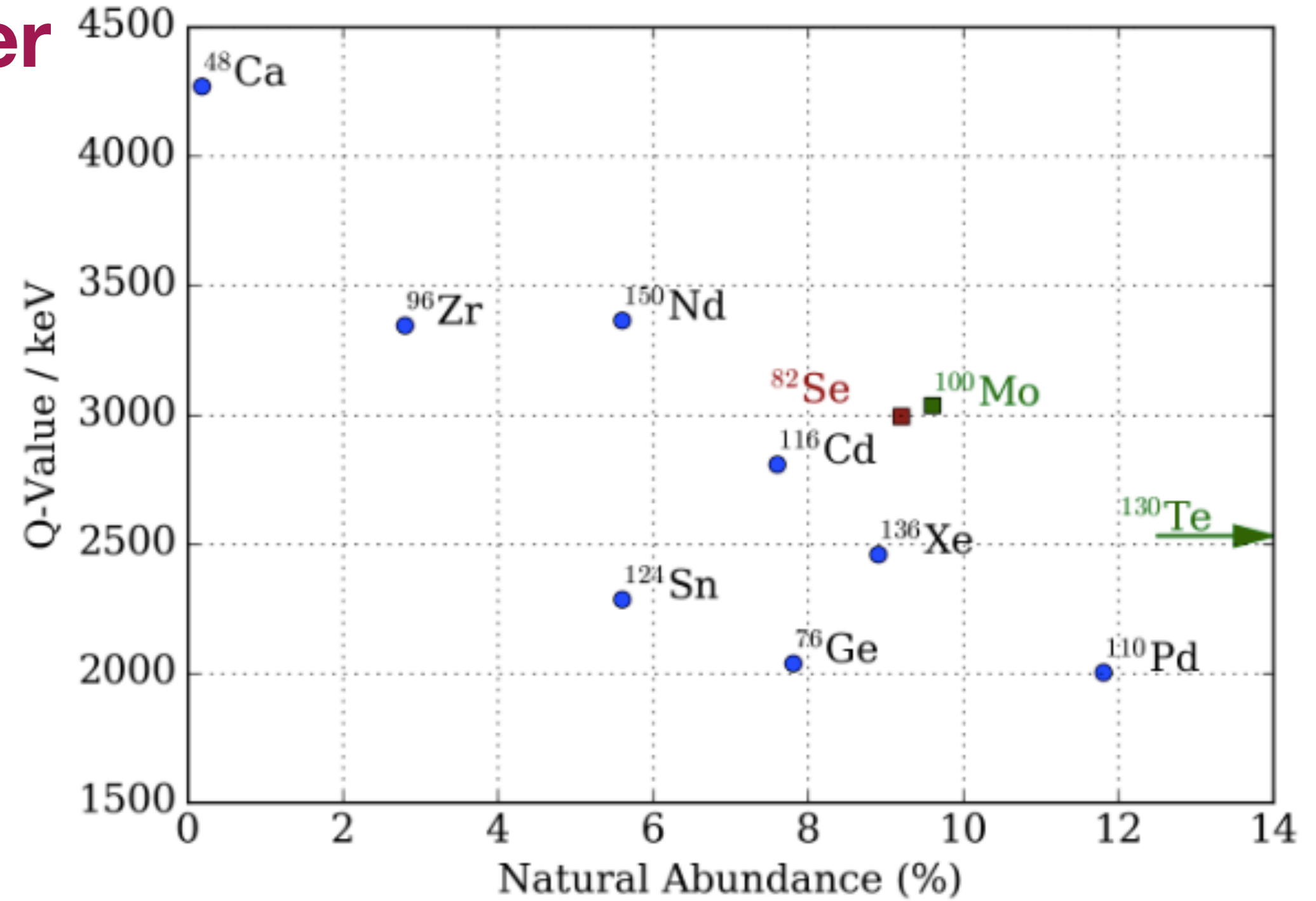


- “Short” half-life

$$\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z) |M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$$

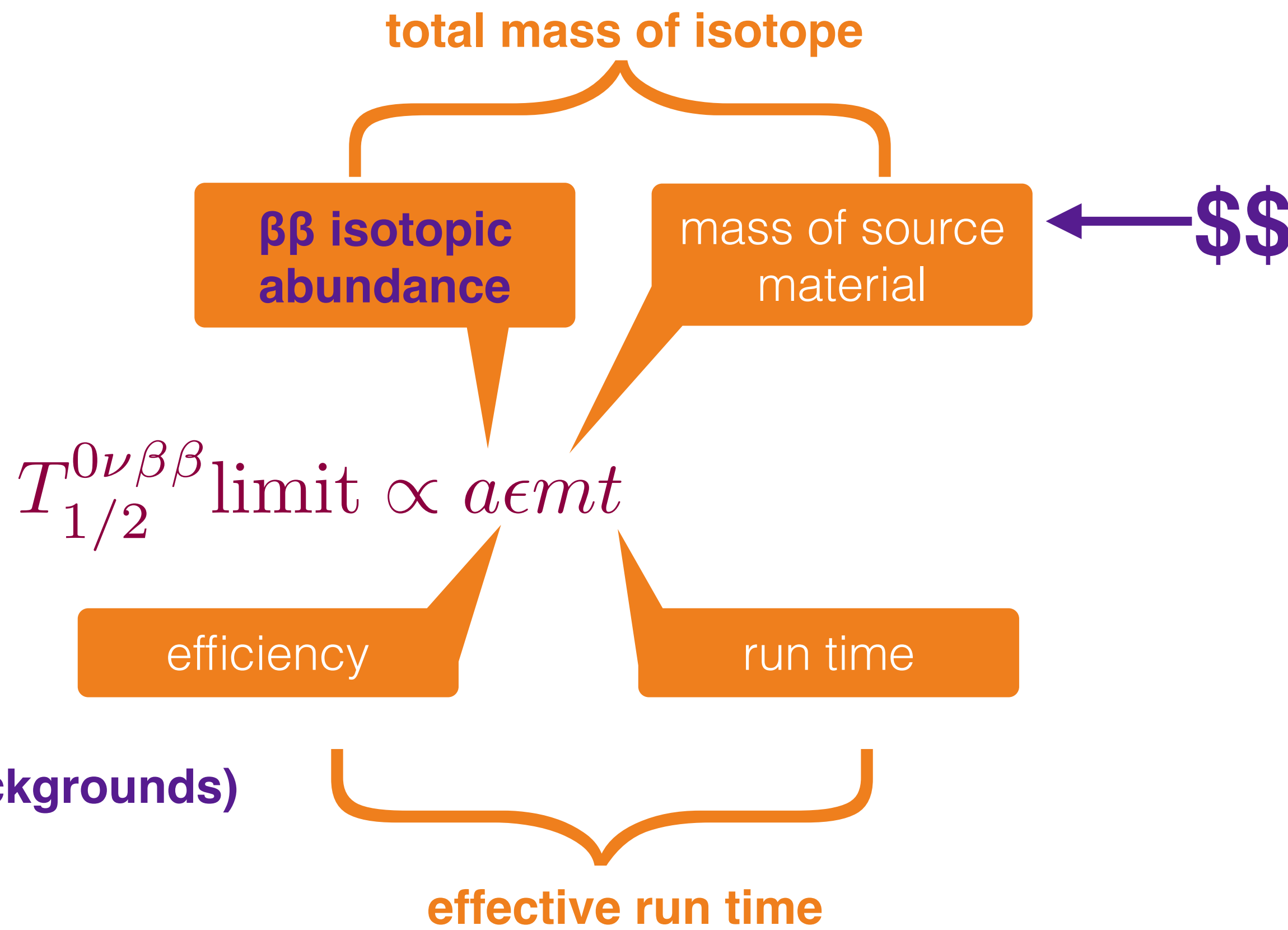
$$T_{1/2}^{0\nu} \propto Q^{-5}$$

Better

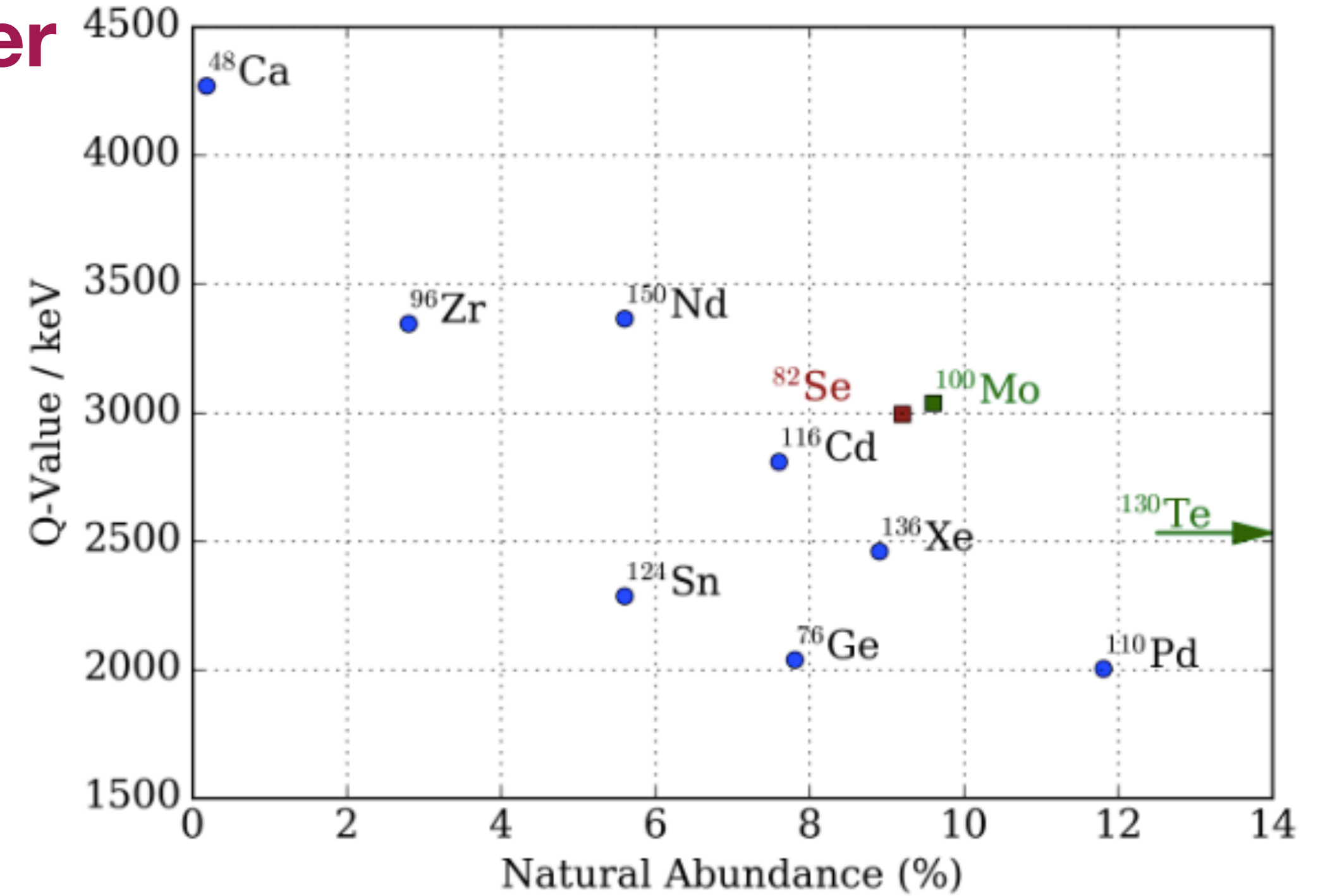


Choosing an isotope

- “Short” half-life
- Lots of isotope



Better



Better

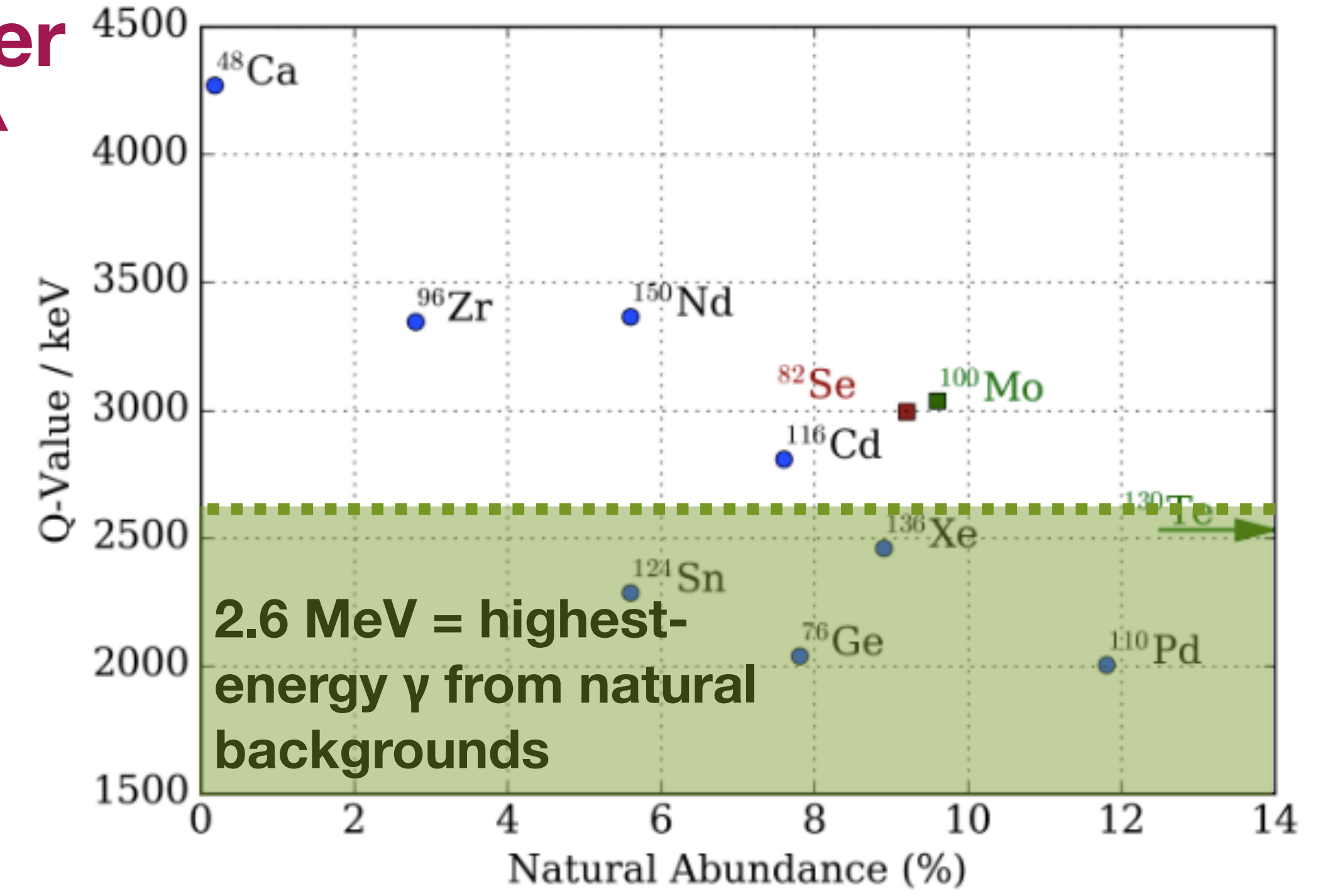
Choosing an isotope

- “Short” half-life
- Lots of isotope
- (Ultra-) Low backgrounds

$$T_{1/2}^{0\nu} \text{ limit} \propto a\epsilon \sqrt{\frac{mt}{b\Delta E}}$$

background events
energy resolution
less “value” for increasing exposure

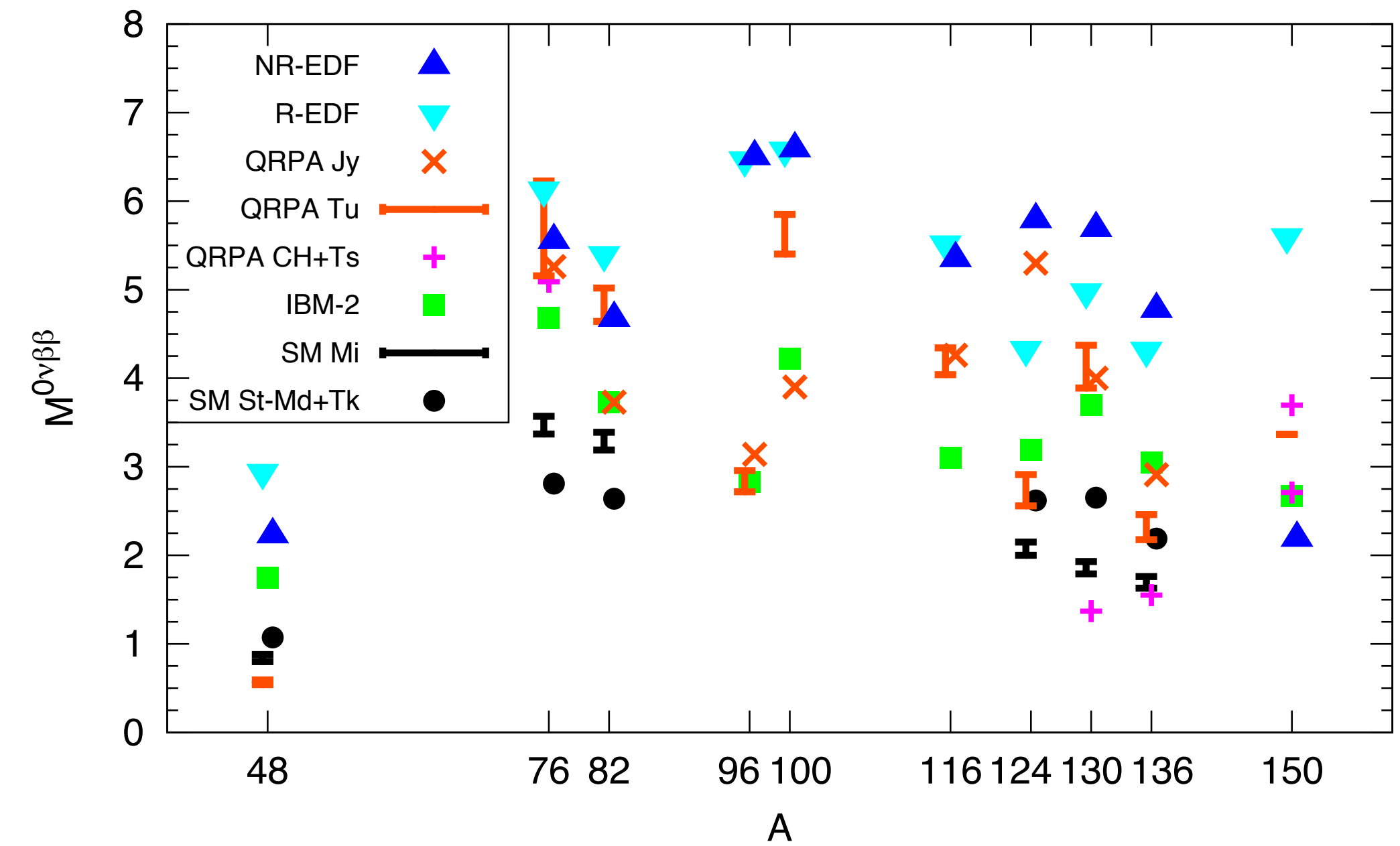
Better



Better

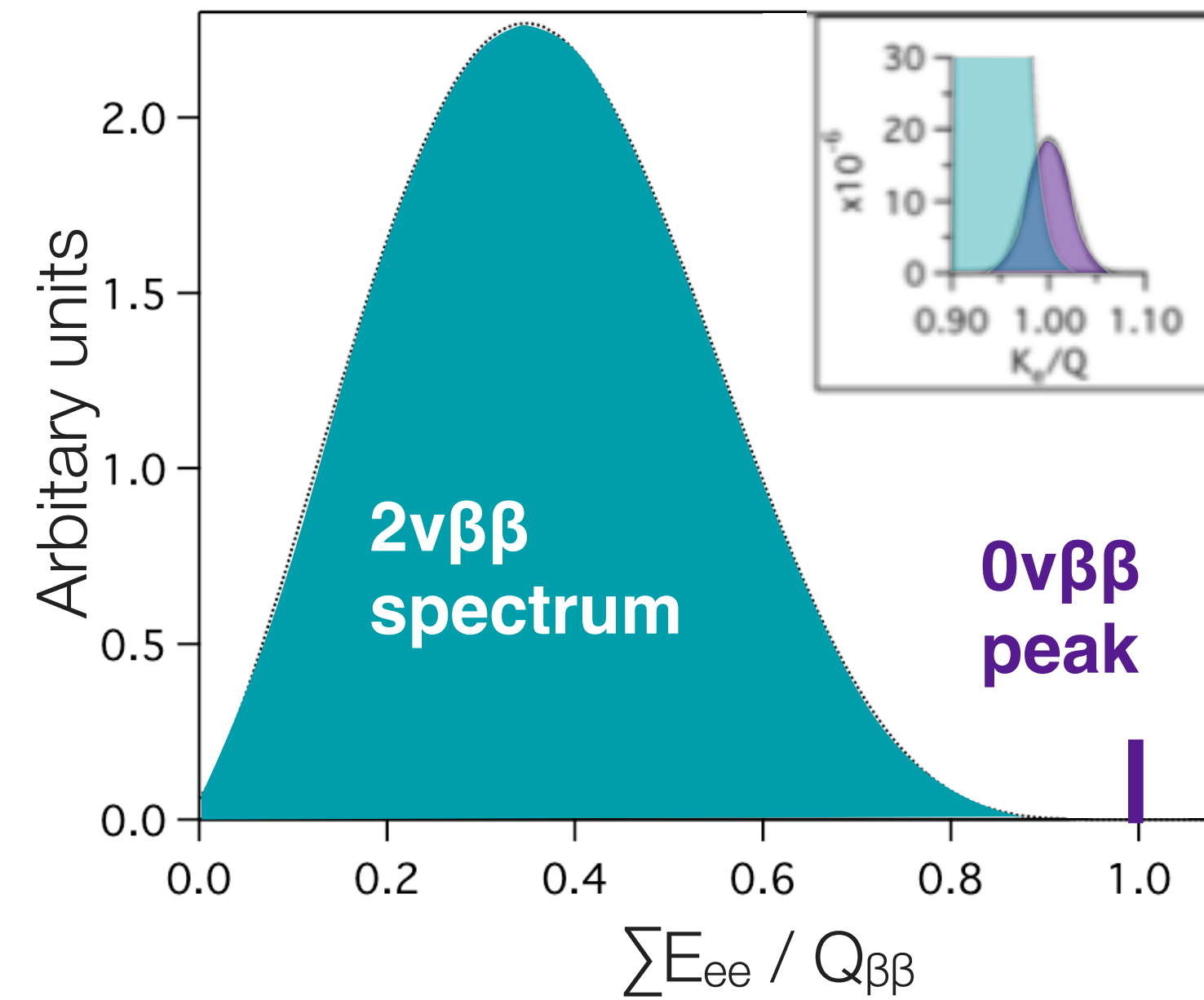
- “Short” half-life
- Lots of isotope
- (Ultra-) Low backgrounds
- Nuclear matrix element

$$\frac{1}{T_{1/2}^{0\nu\beta\beta}} = G_{0\nu}(Q_{\beta\beta}, Z) g_A^4 |M_{0\nu}|^2 \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$$





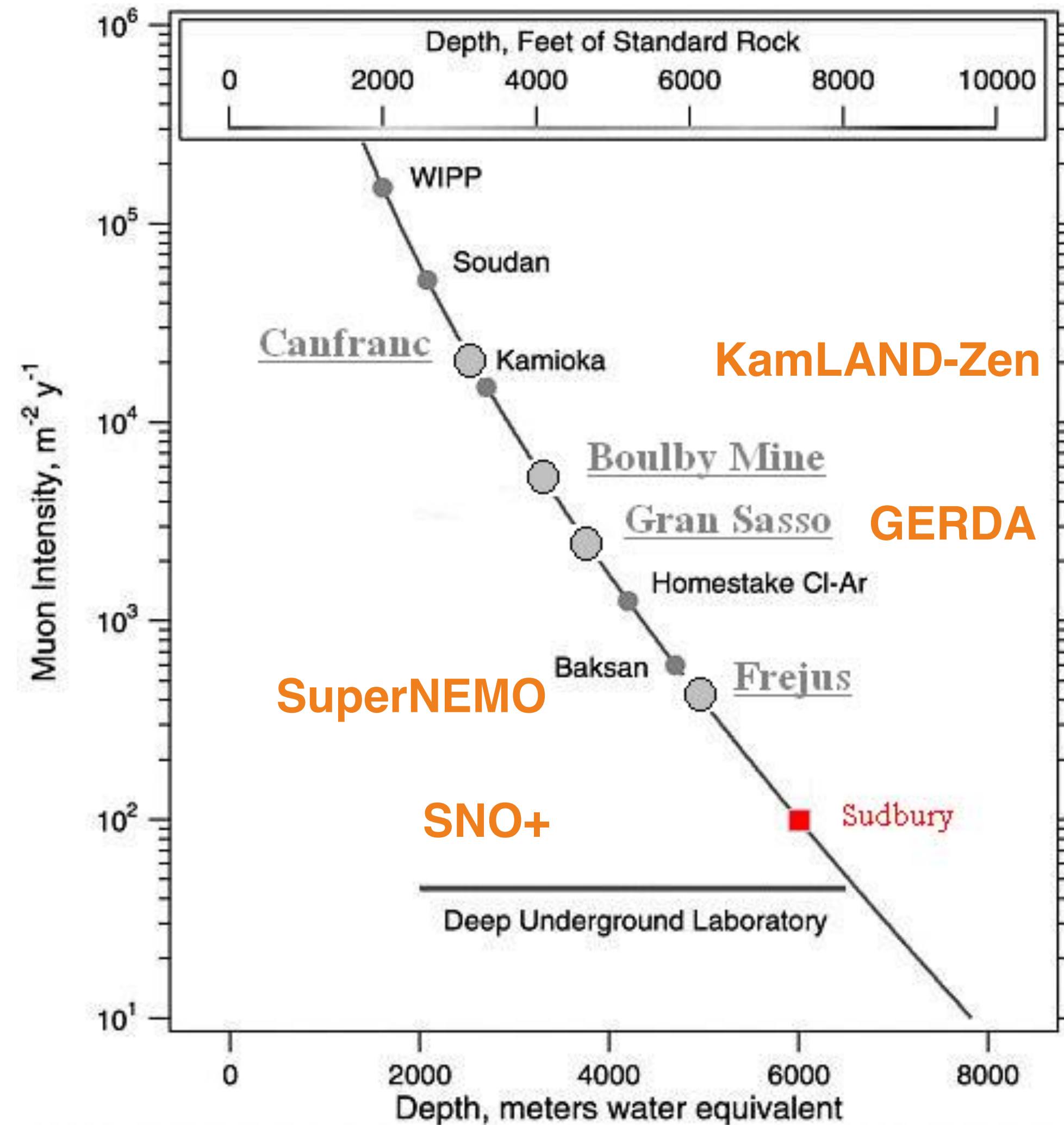
$2\nu\beta\beta$ high-energy tail



- **Irreducible** - all $2\nu\beta\beta$ isotopes also exhibit $2\nu\beta\beta$
- Exactly the **same signature** apart from energy difference
- Excellent **energy resolution** is the only way to mitigate

$2\nu\beta\beta$ high-energy tail

Cosmic rays



Alps at Modane (SuperNEMO)

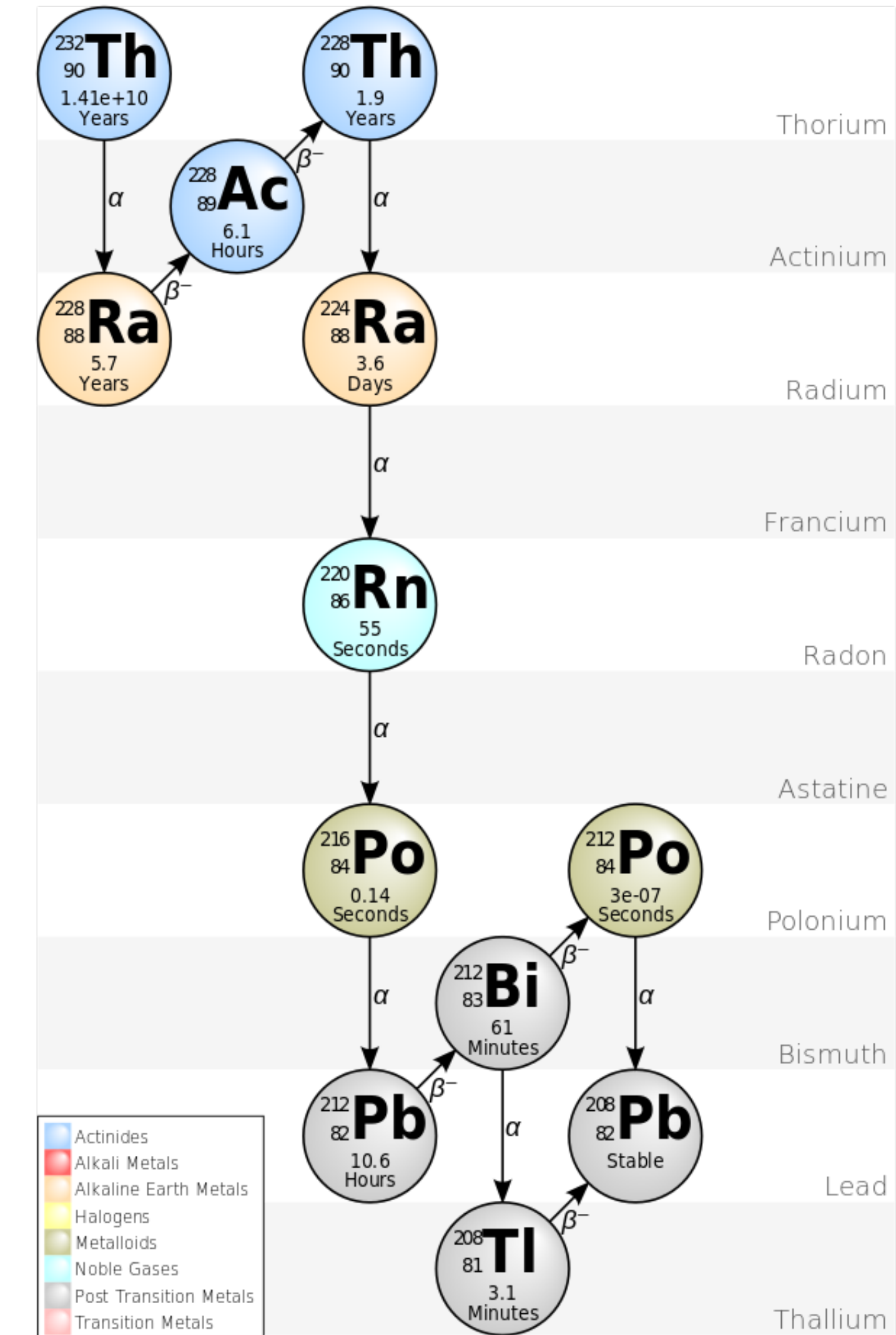
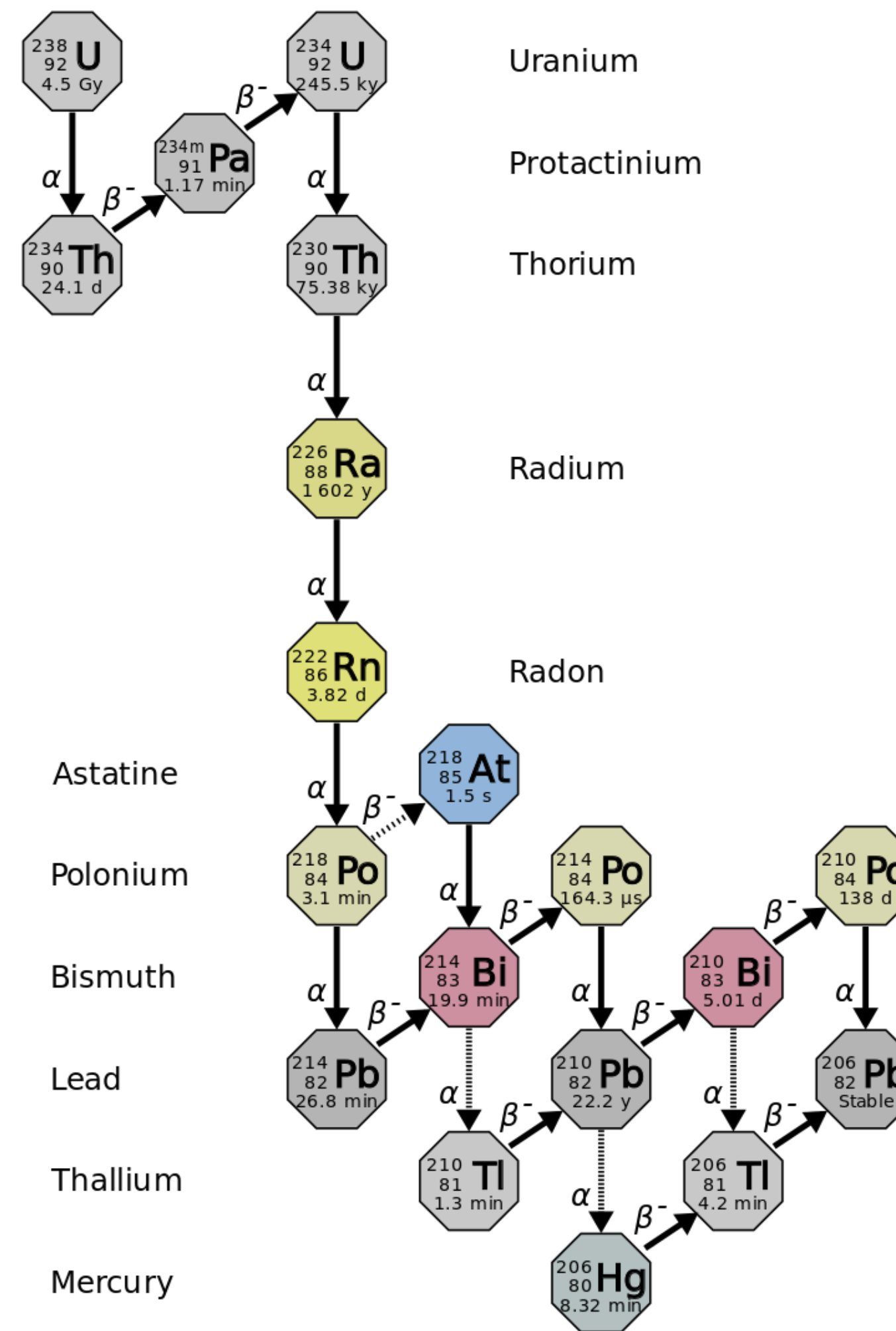


Sudbury nickel mine (SNO+)

$2\nu\beta\beta$ high-energy tail

Cosmic rays

β -decaying isotopes

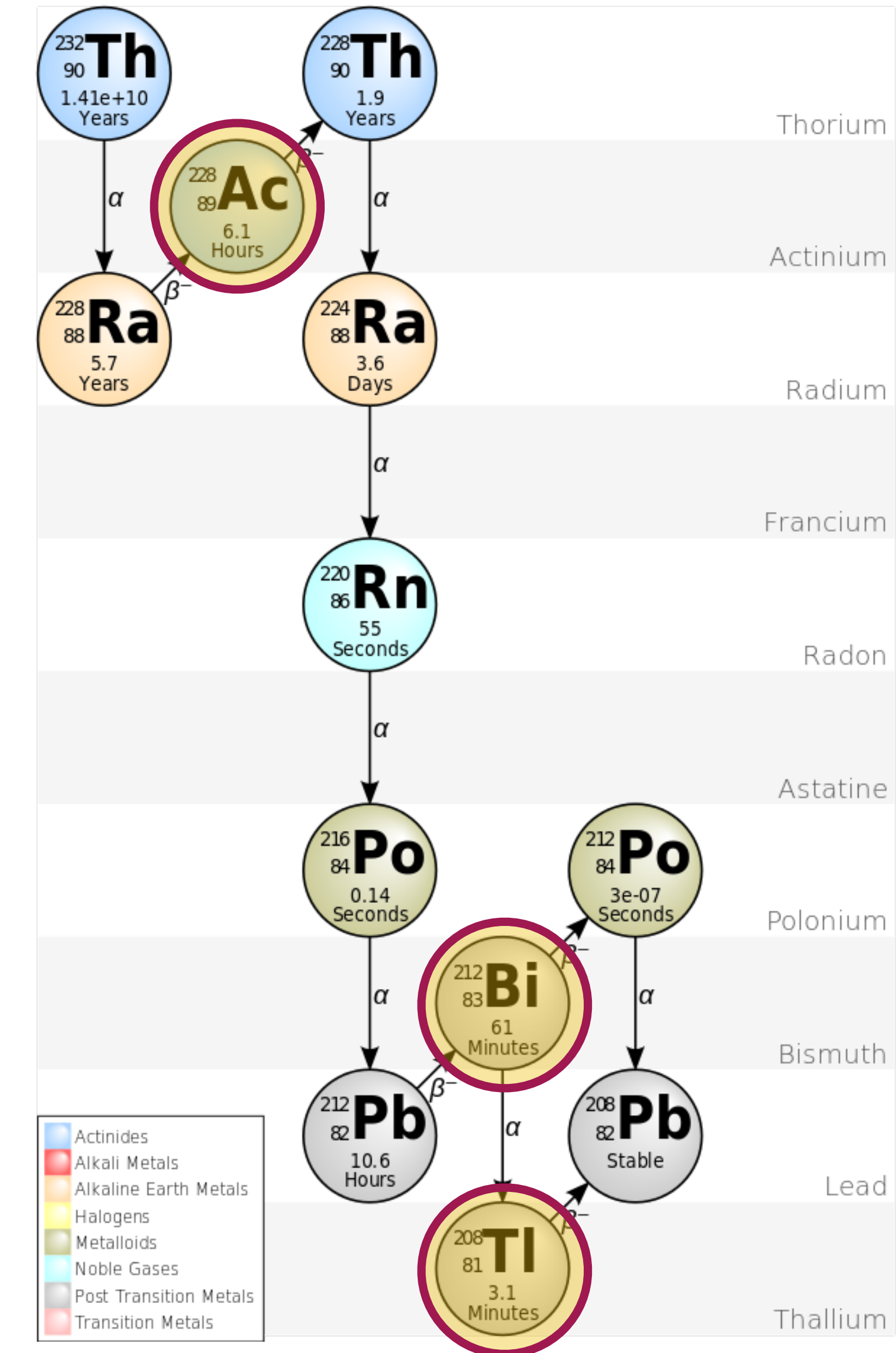
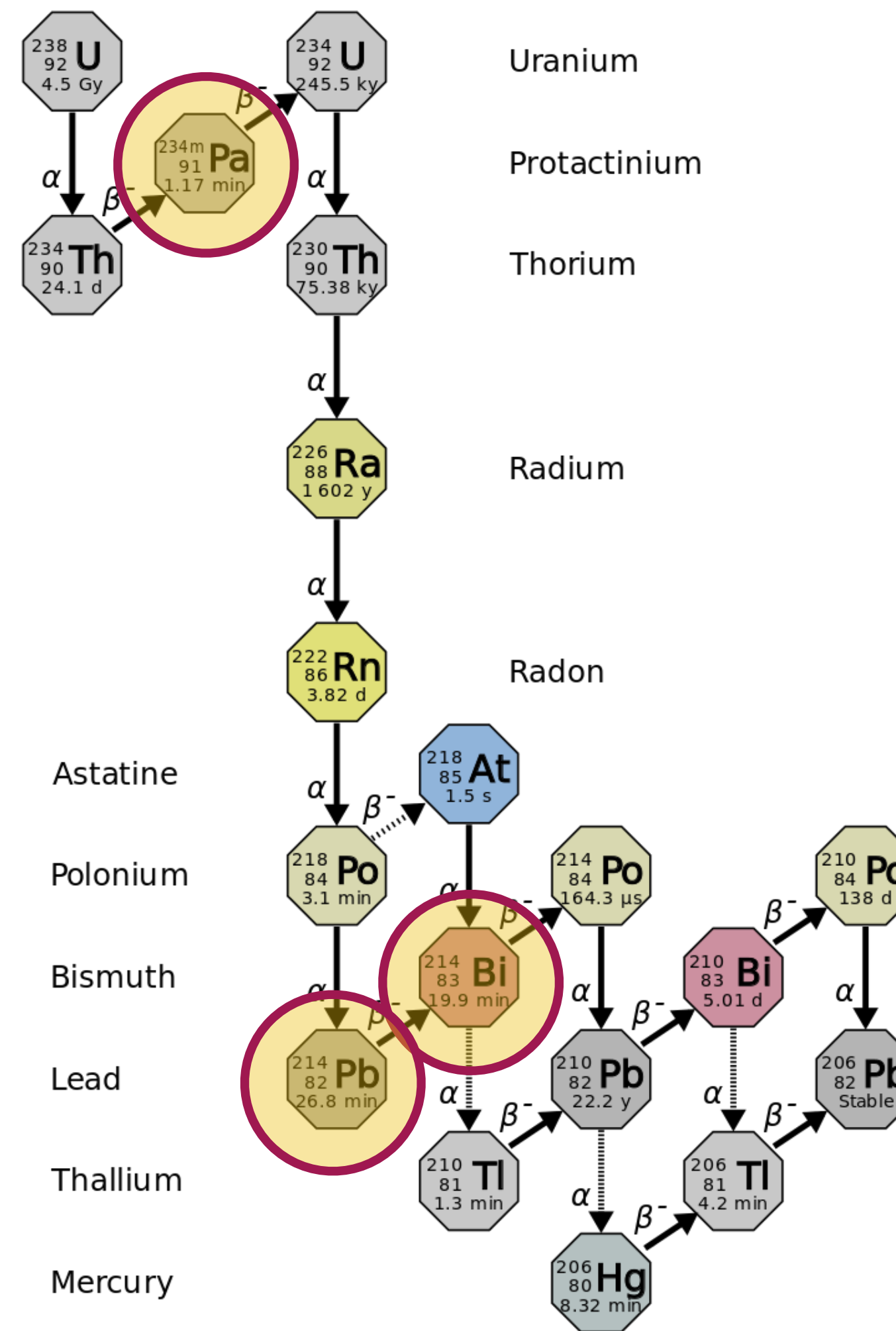


Backgrounds to double-beta experiments

$2\nu\beta\beta$ high-energy tail

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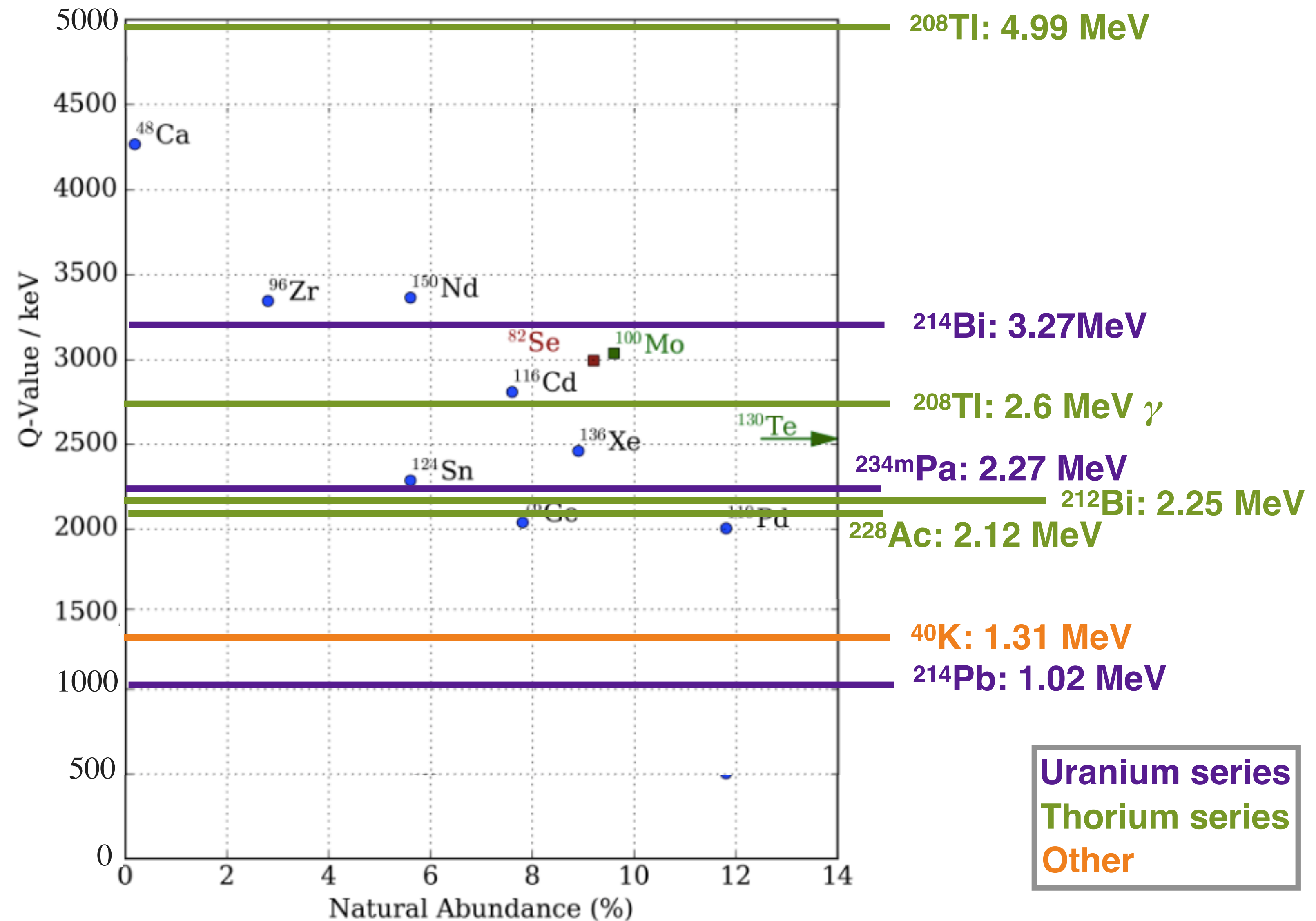


- Actinides
- Alkali Metals
- Alkaline Earth Metals
- Halogens
- Metalloids
- Noble Gases
- Post Transition Metals
- Transition Metals

$2\nu\beta\beta$ high-energy tail

Cosmic rays

β -decaying isotopes



$2\nu\beta\beta$ high-energy tail

Cosmic rays

β -decaying isotopes

External contamination



Building the GERDA detector array in the glovebox © J. Suvorov, GERDA collaboration

Working in the SuperNEMO cleanroom



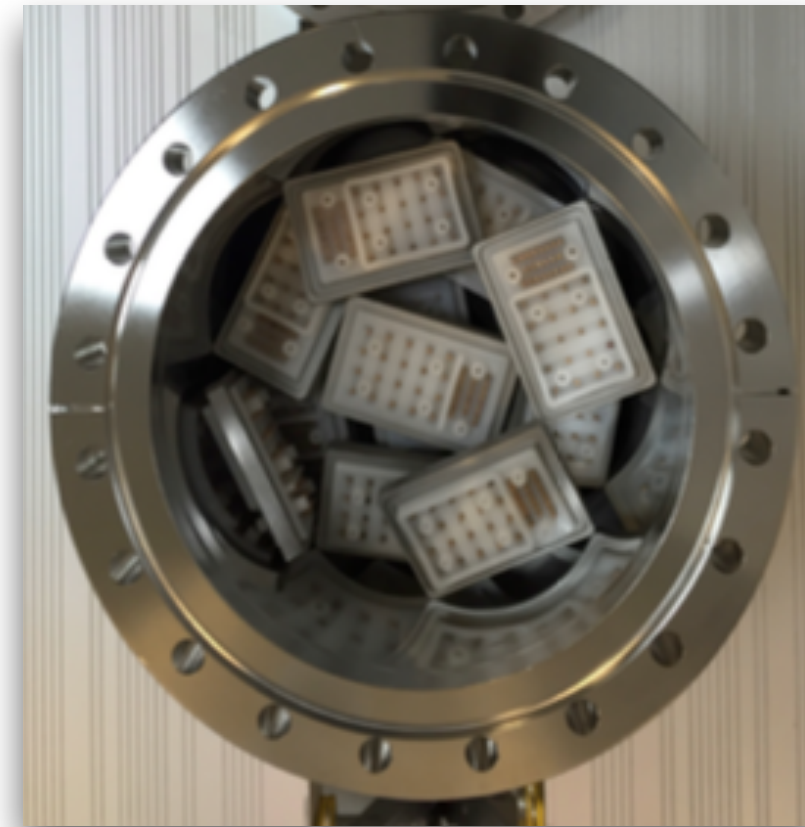
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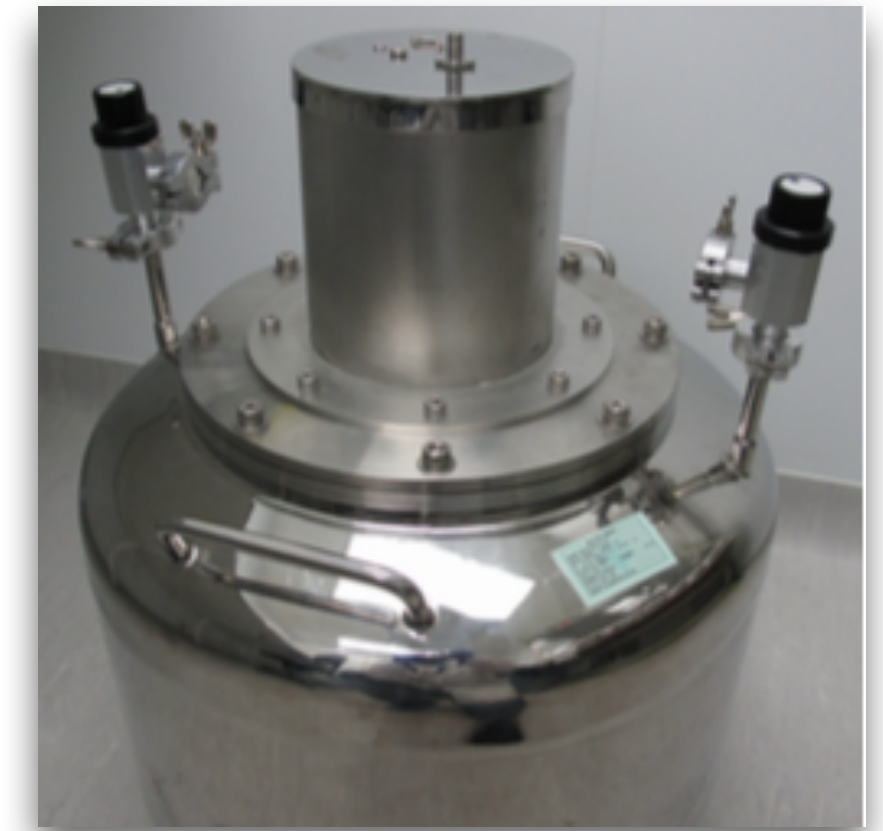
β -decaying isotopes

External contamination

Detector components



Radiopure components selected & allowed to **emanate radon**. Activity is measured with an **electrostatic detector** (SuperNEMO)



Components machined and cleaned in underground cleanroom (MAJORANA)



2 $\nu\beta\beta$ high-energy tail

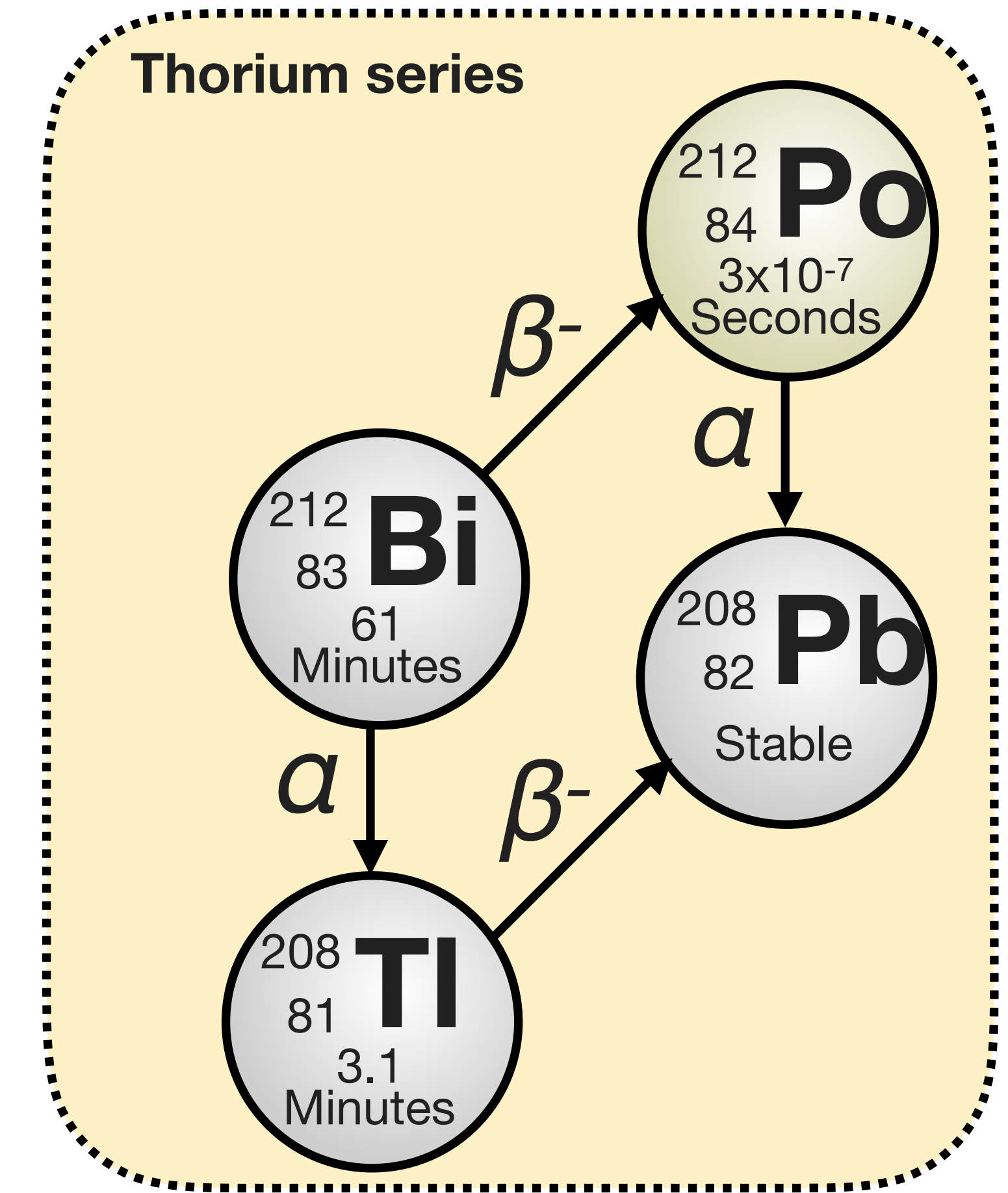
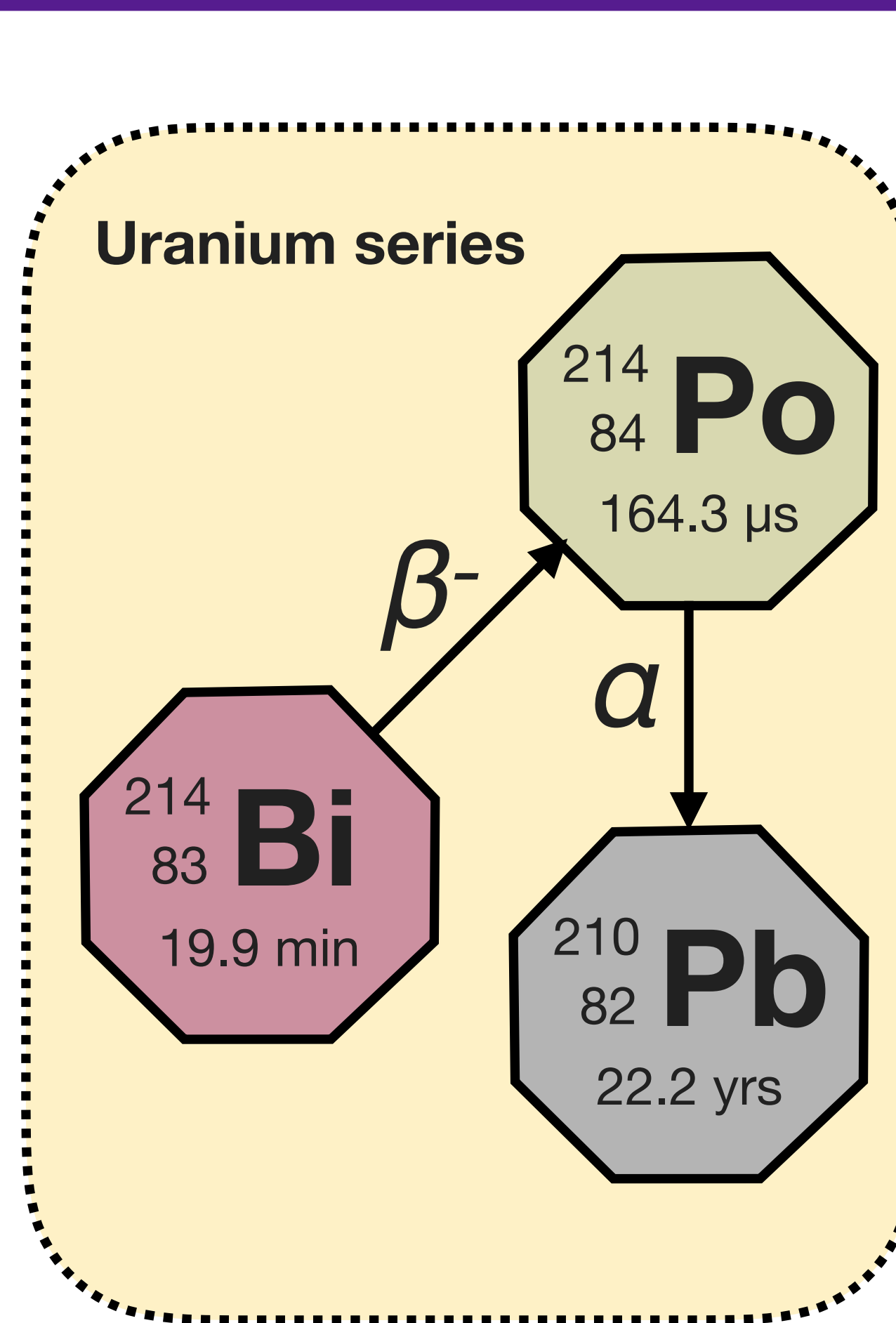
Cosmic rays

β -decaying isotopes

External contamination

Detector components

BiPo events



$2\nu\beta\beta$ high-energy tail

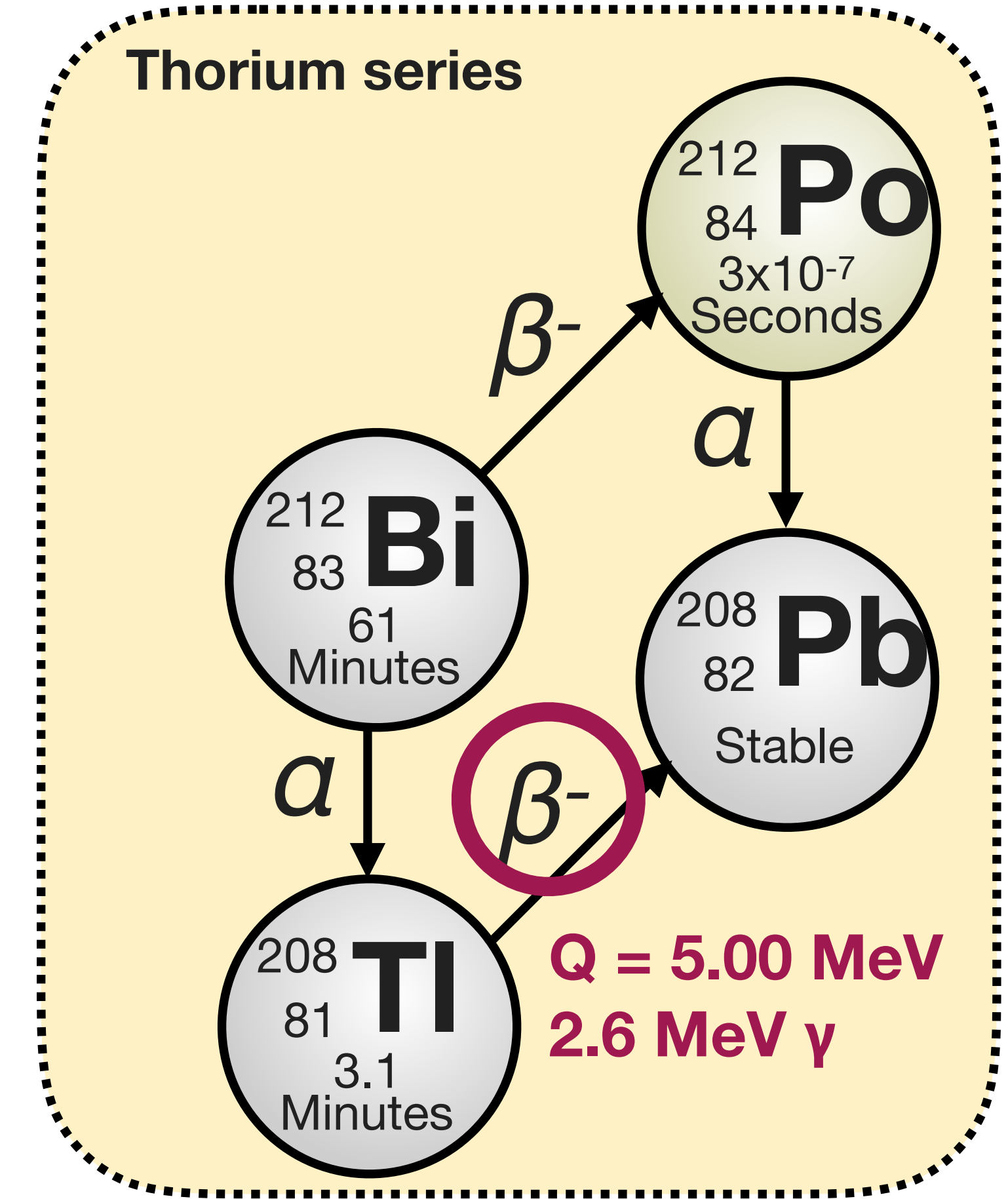
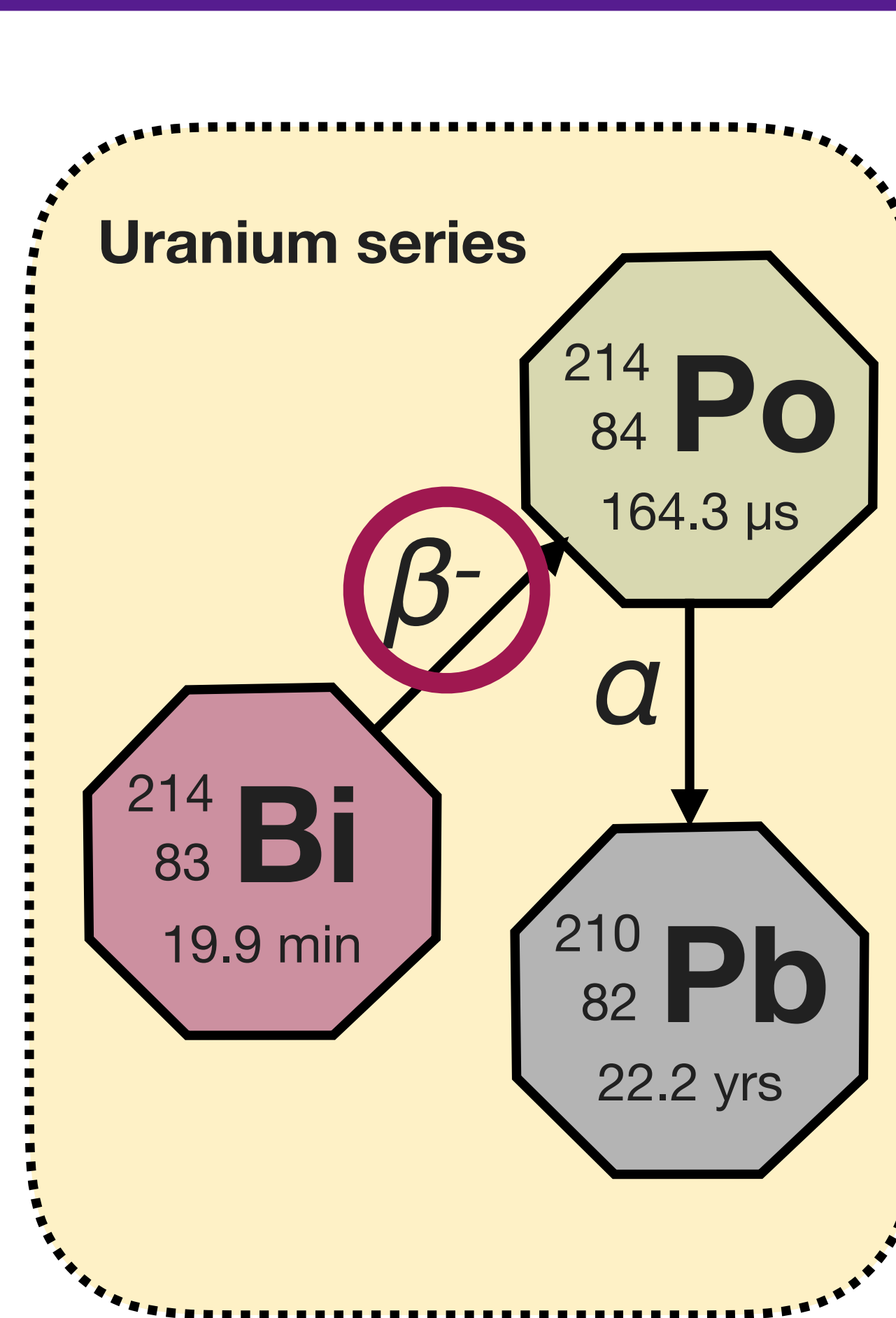
Cosmic rays

β -decaying isotopes

External contamination

Detector components

BiPo events



2νββ high-energy tail

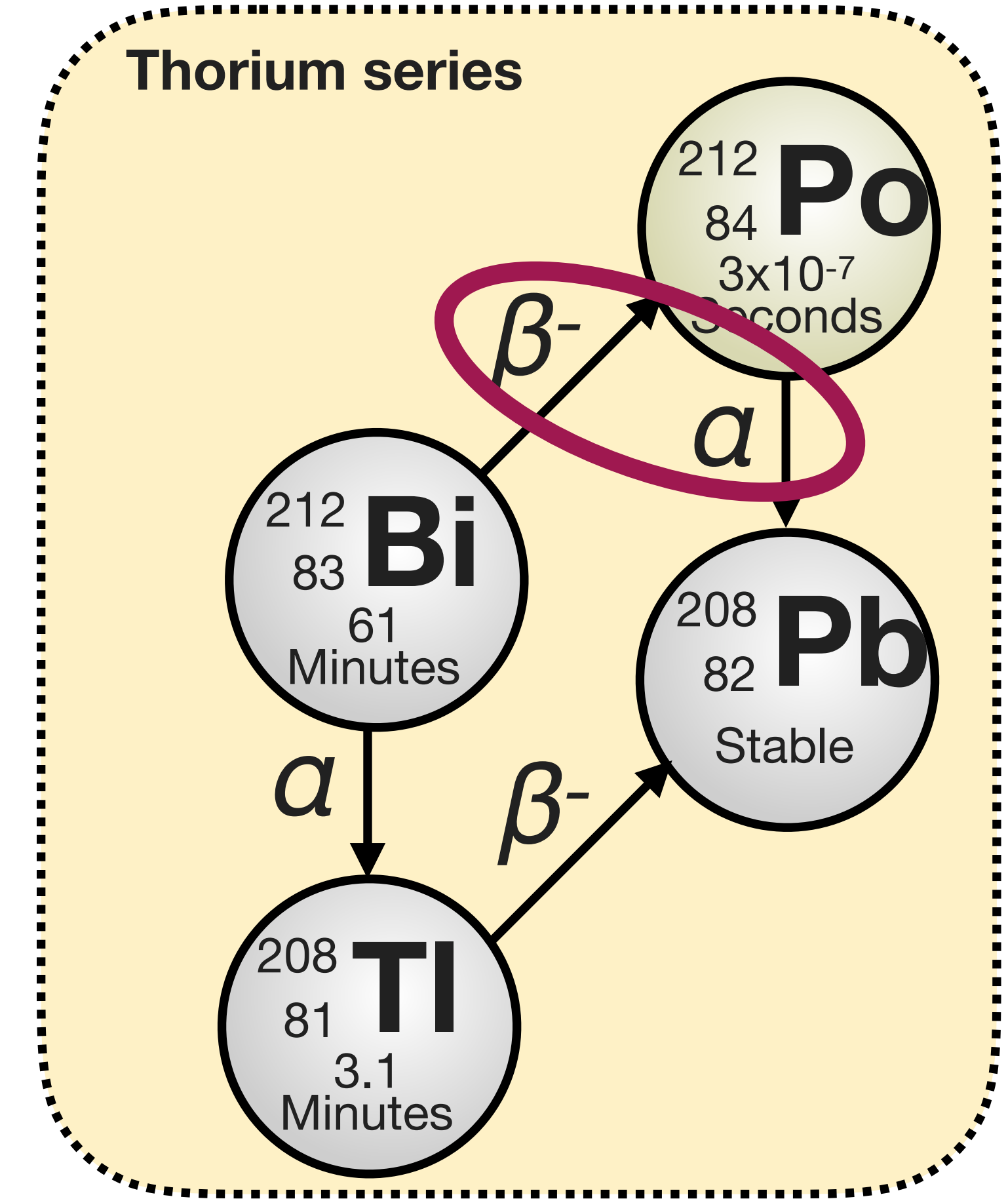
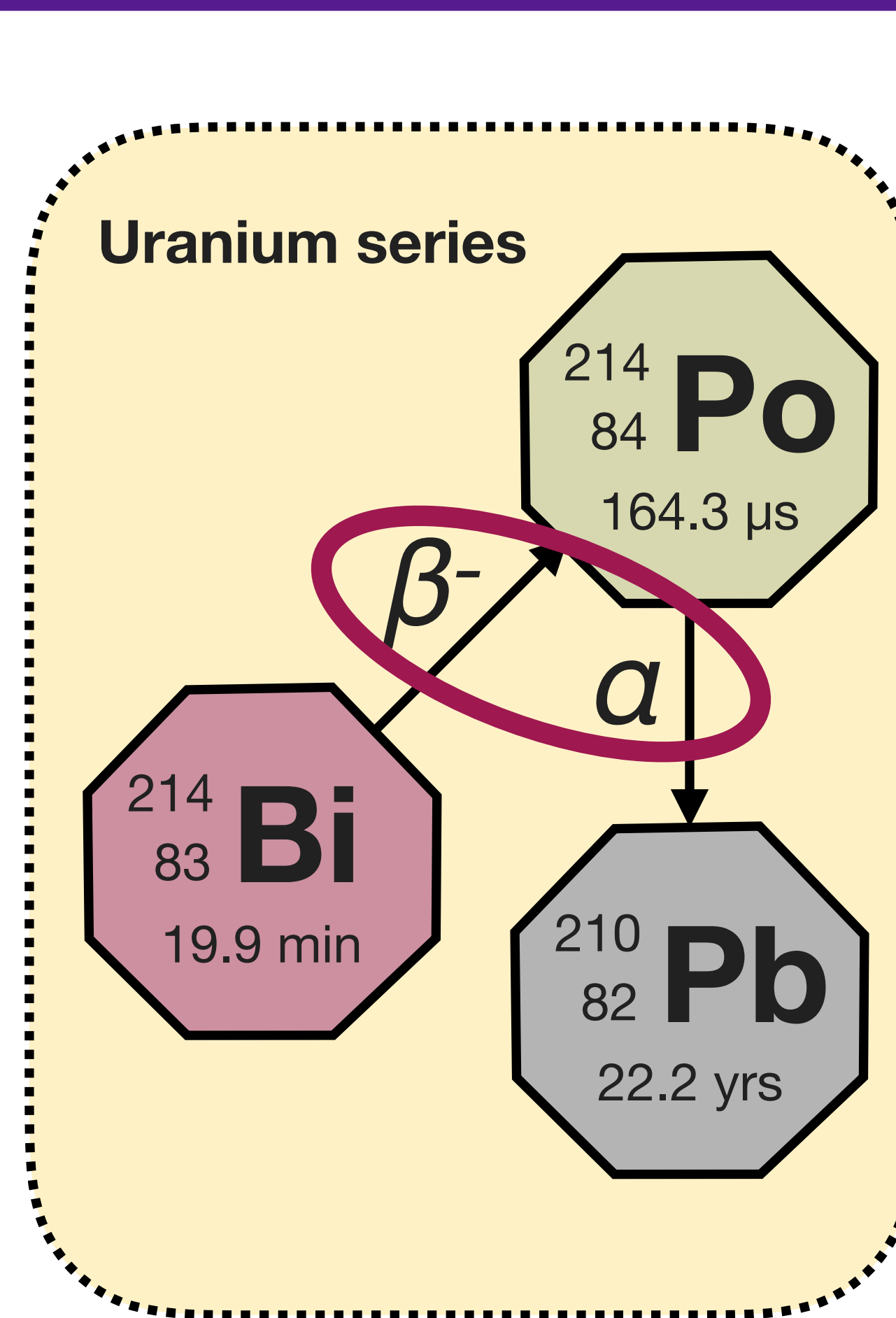
Cosmic rays

β-decaying isotopes

External contamination

Detector components

BiPo events



$2\nu\beta\beta$ high-energy tail

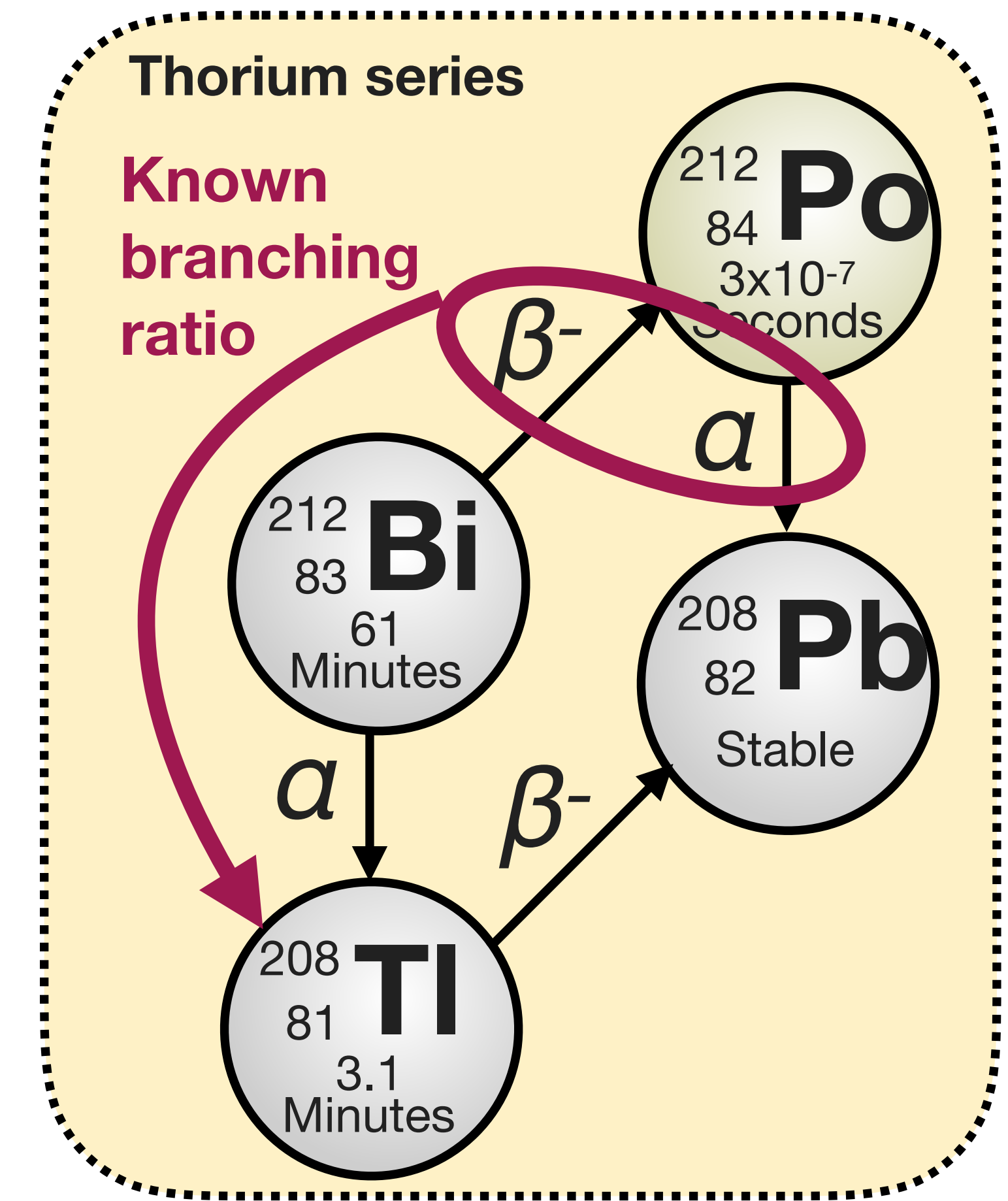
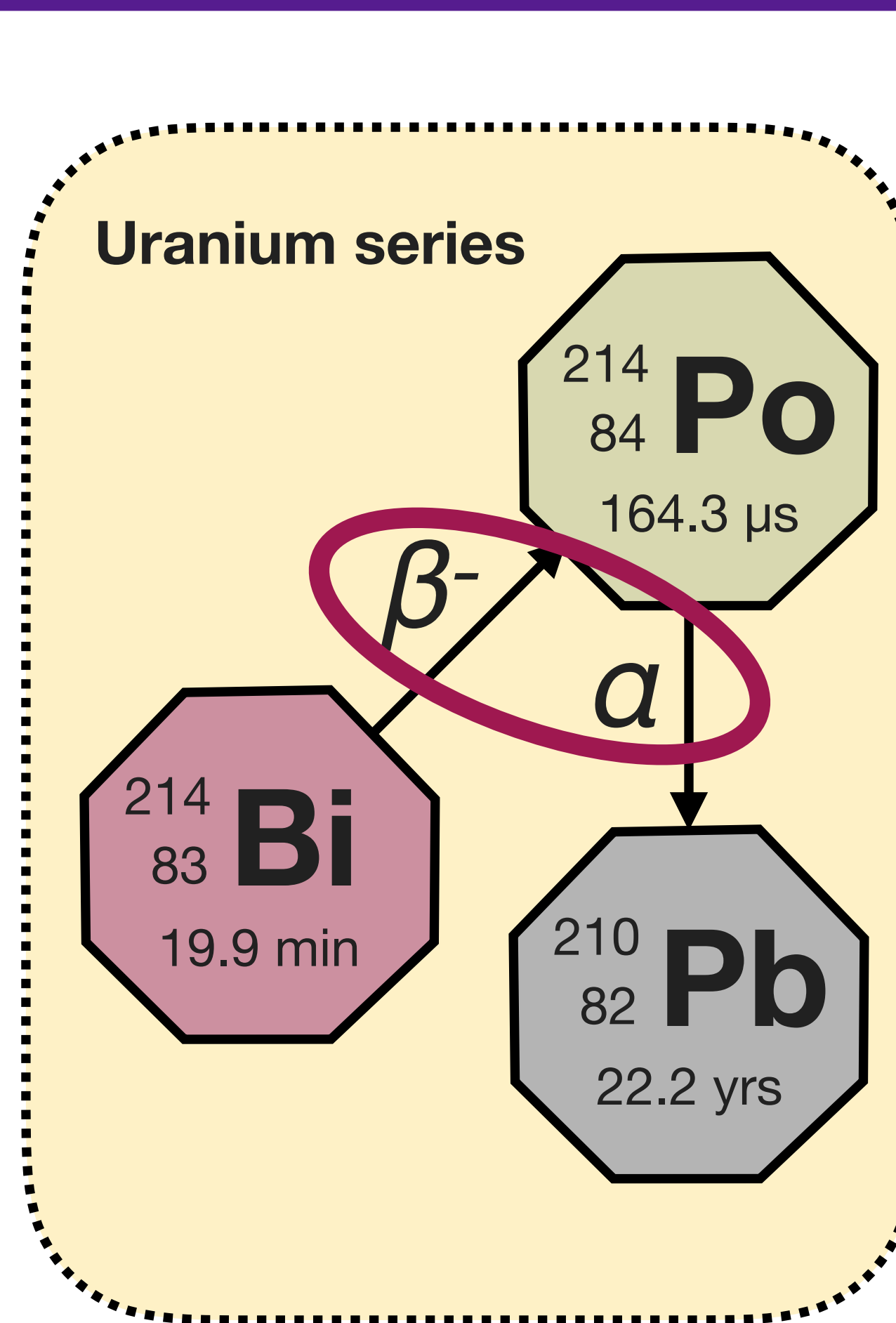
Cosmic rays

β -decaying isotopes

External contamination

Detector components

BiPo events



2νββ high-energy tail

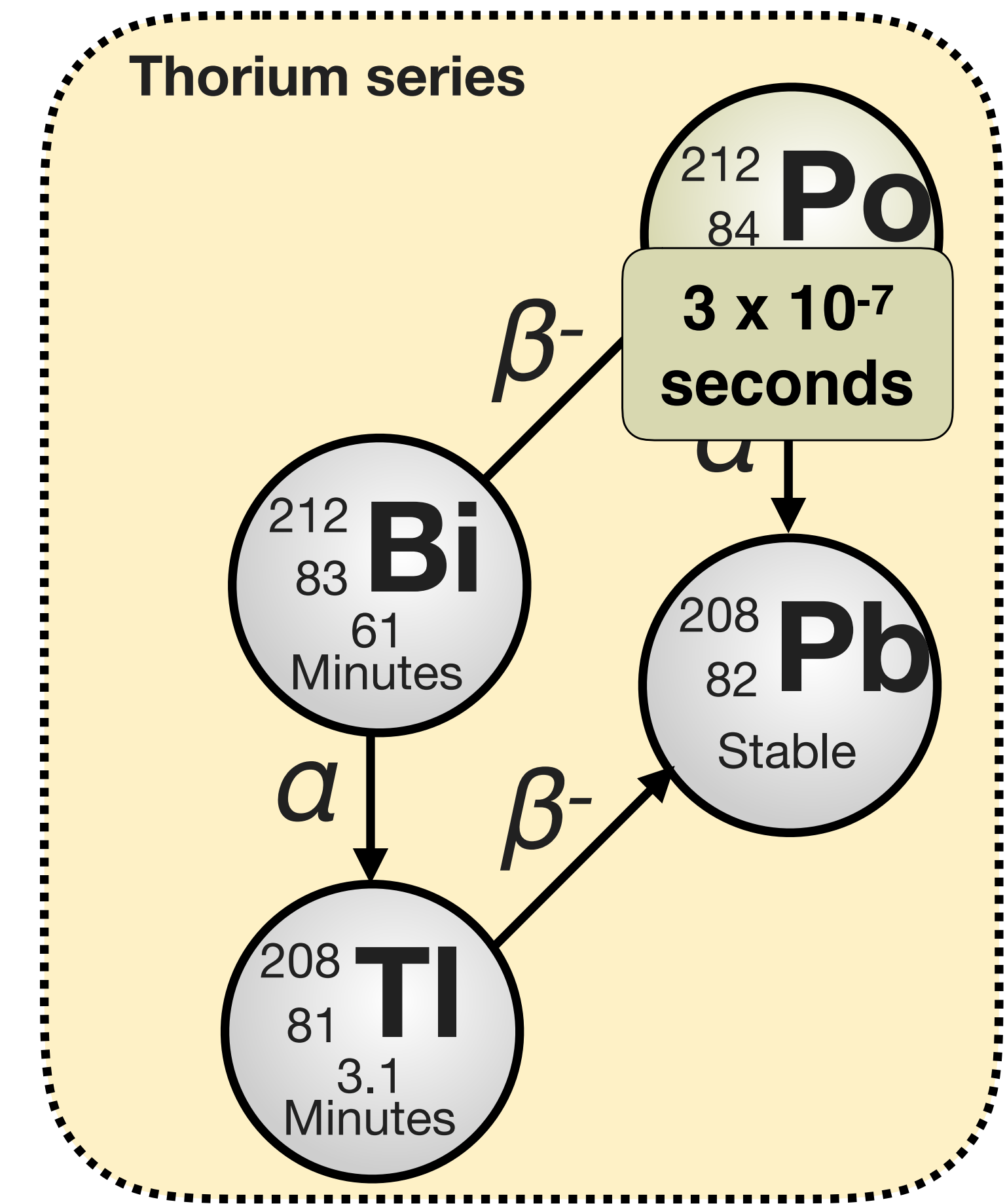
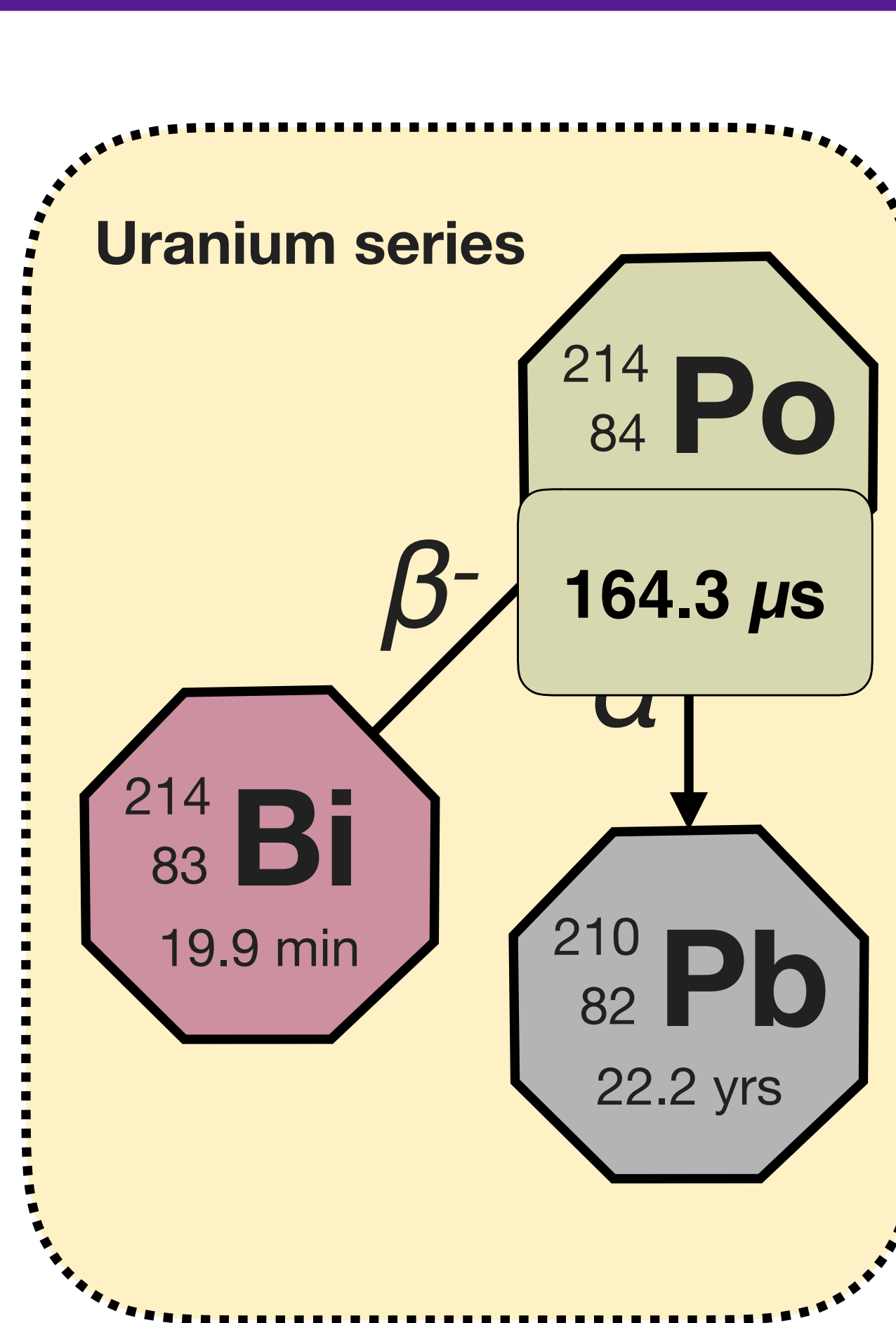
Cosmic rays

β-decaying isotopes

External contamination

Detector components

BiPo events



$2\nu\beta\beta$ high-energy tail

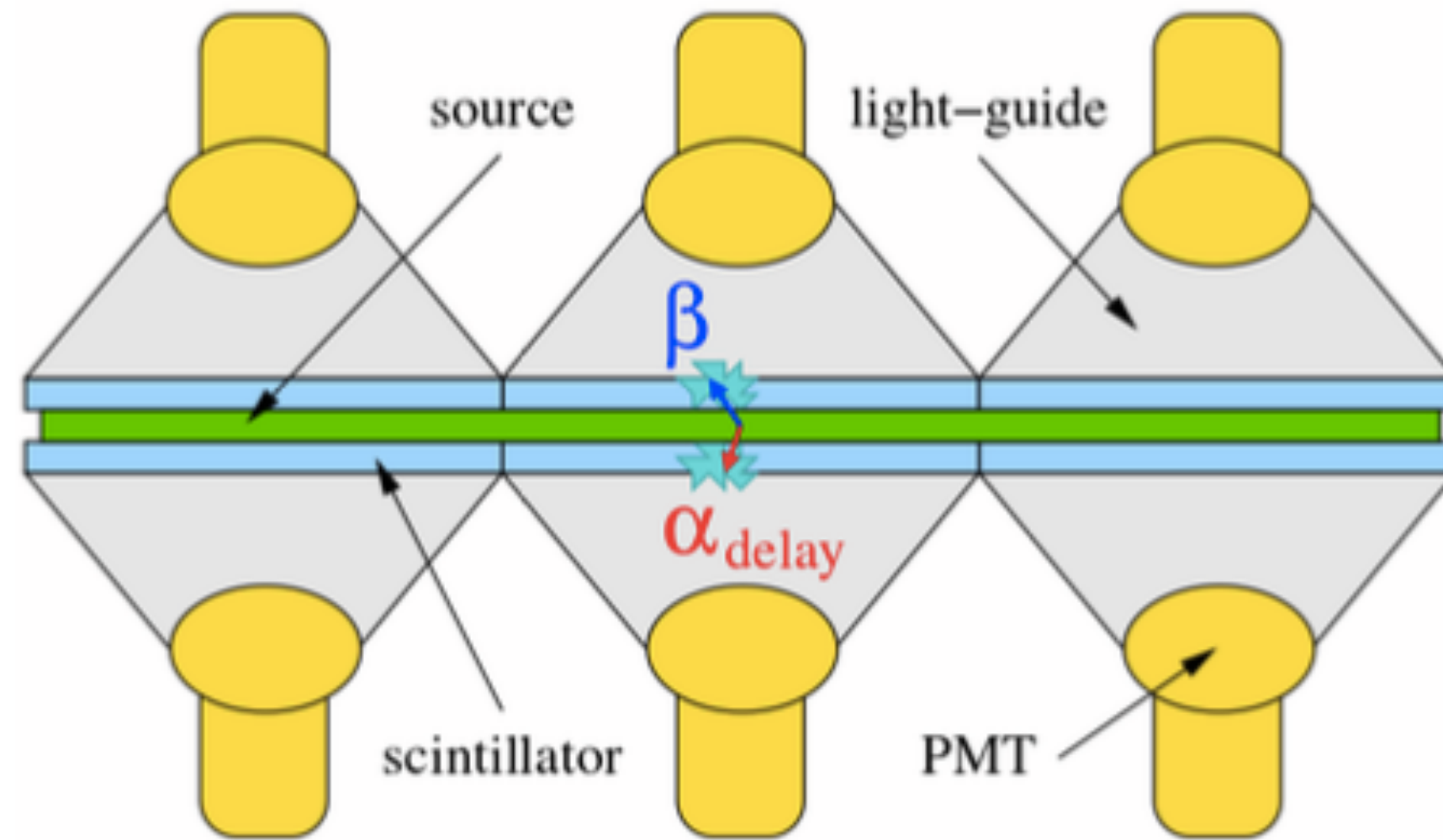
Cosmic rays

β -decaying isotopes

External contamination

Detector components

BiPo events

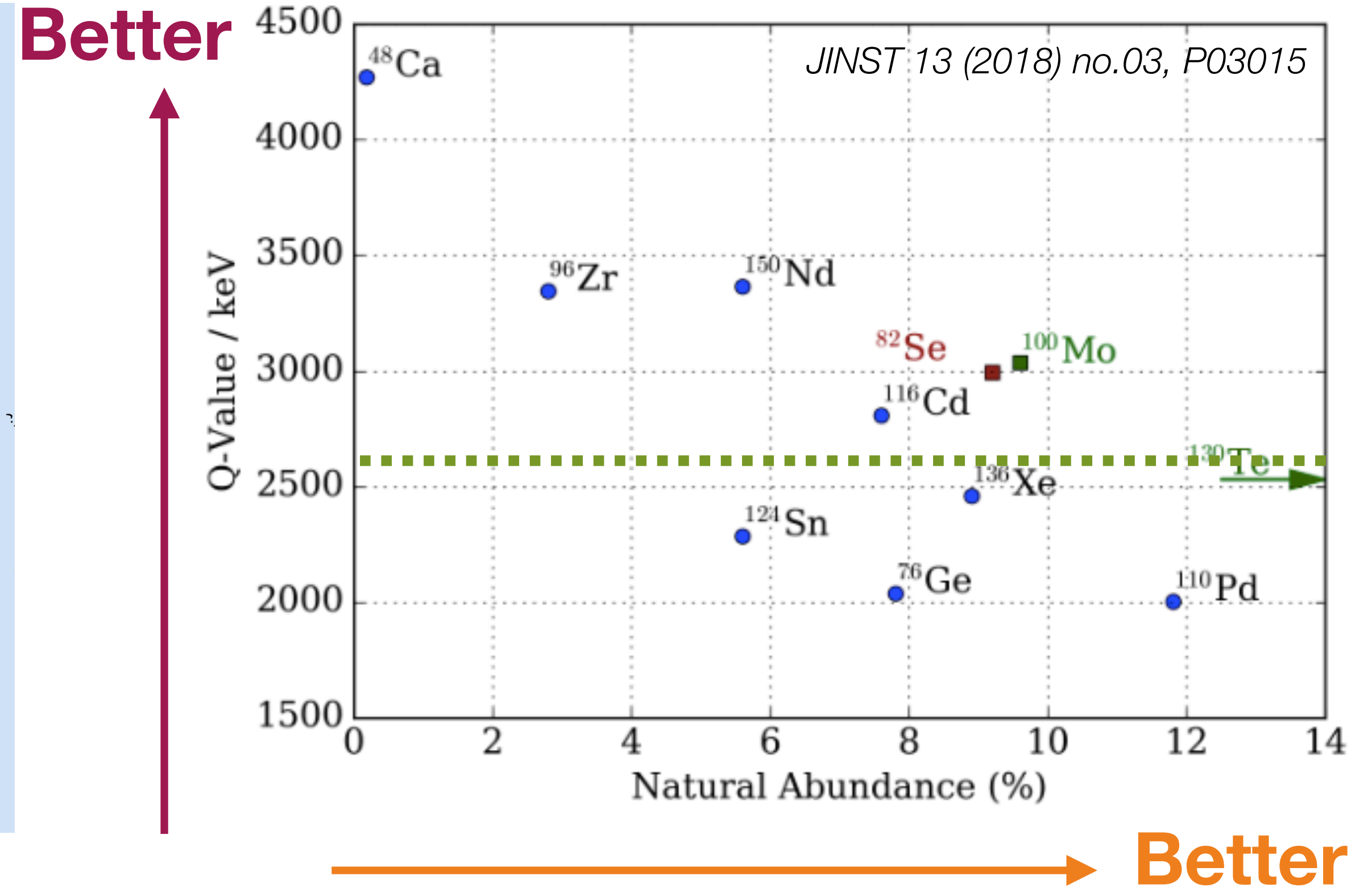


BiPo-3 detector looks for this topology

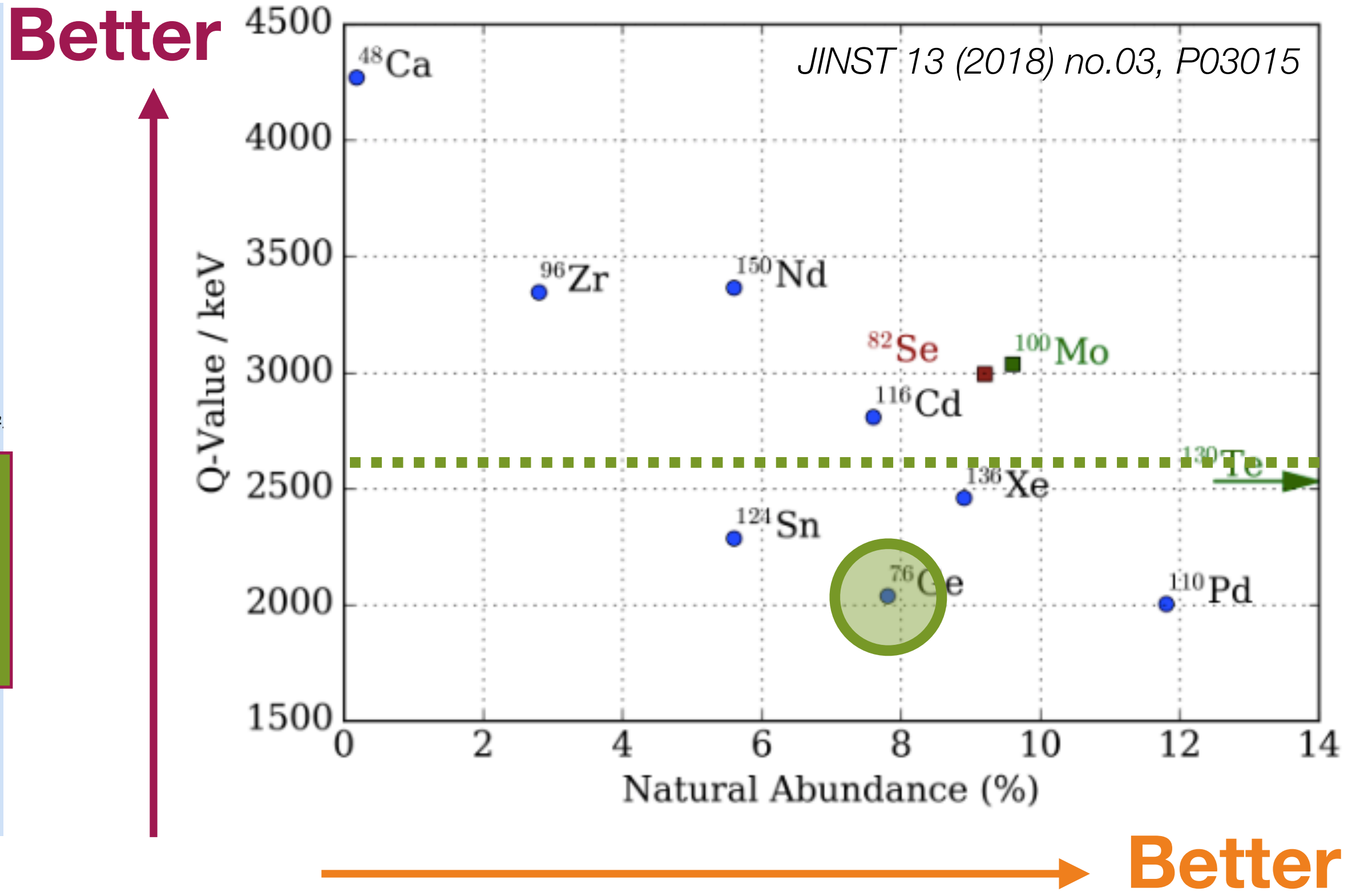
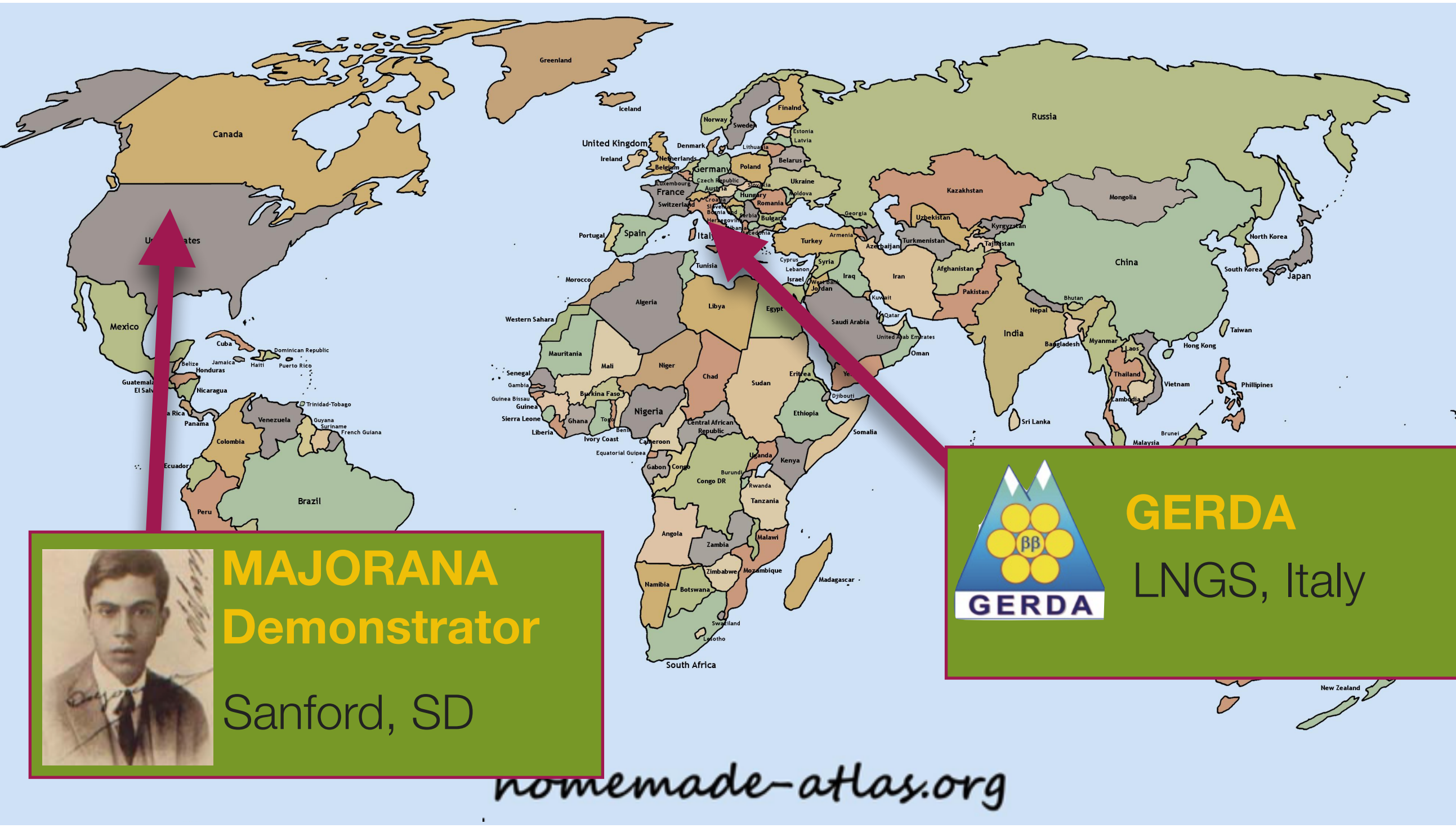
Canfranc underground lab was formerly Europe's biggest train station



Double-beta decay experiments



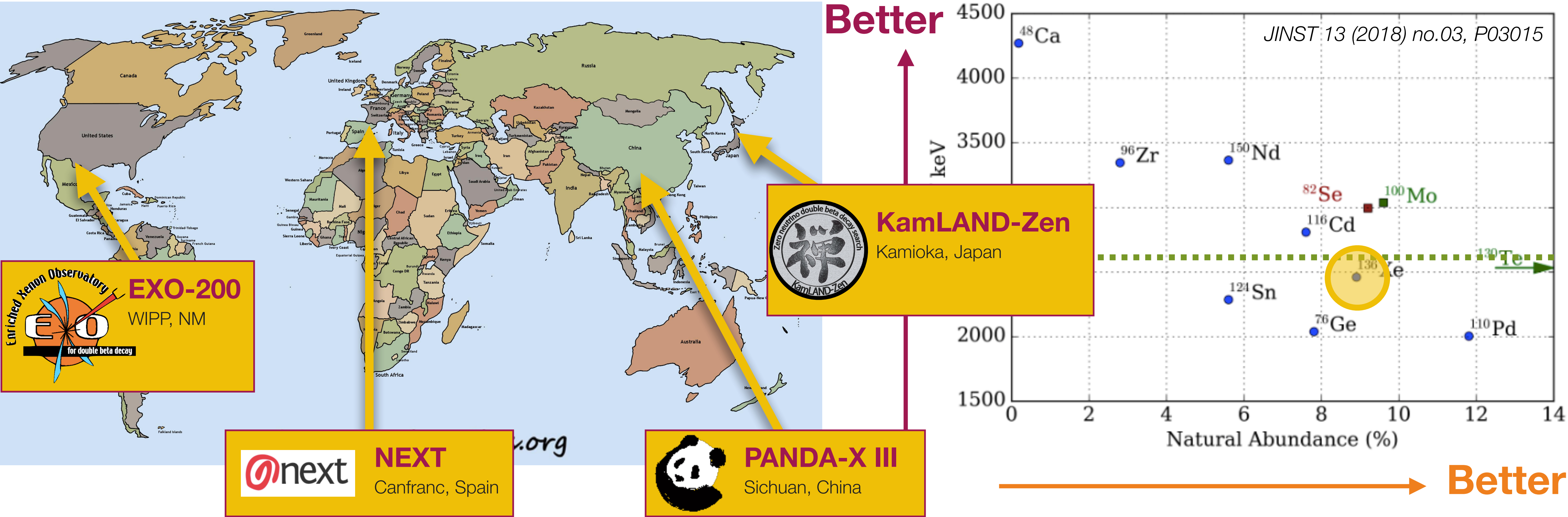
Double-beta decay experiments



- $\beta\beta$ source = detector (^{76}Ge)
- Excellent **efficiency** and **resolution**
- Future : **LEGEND**

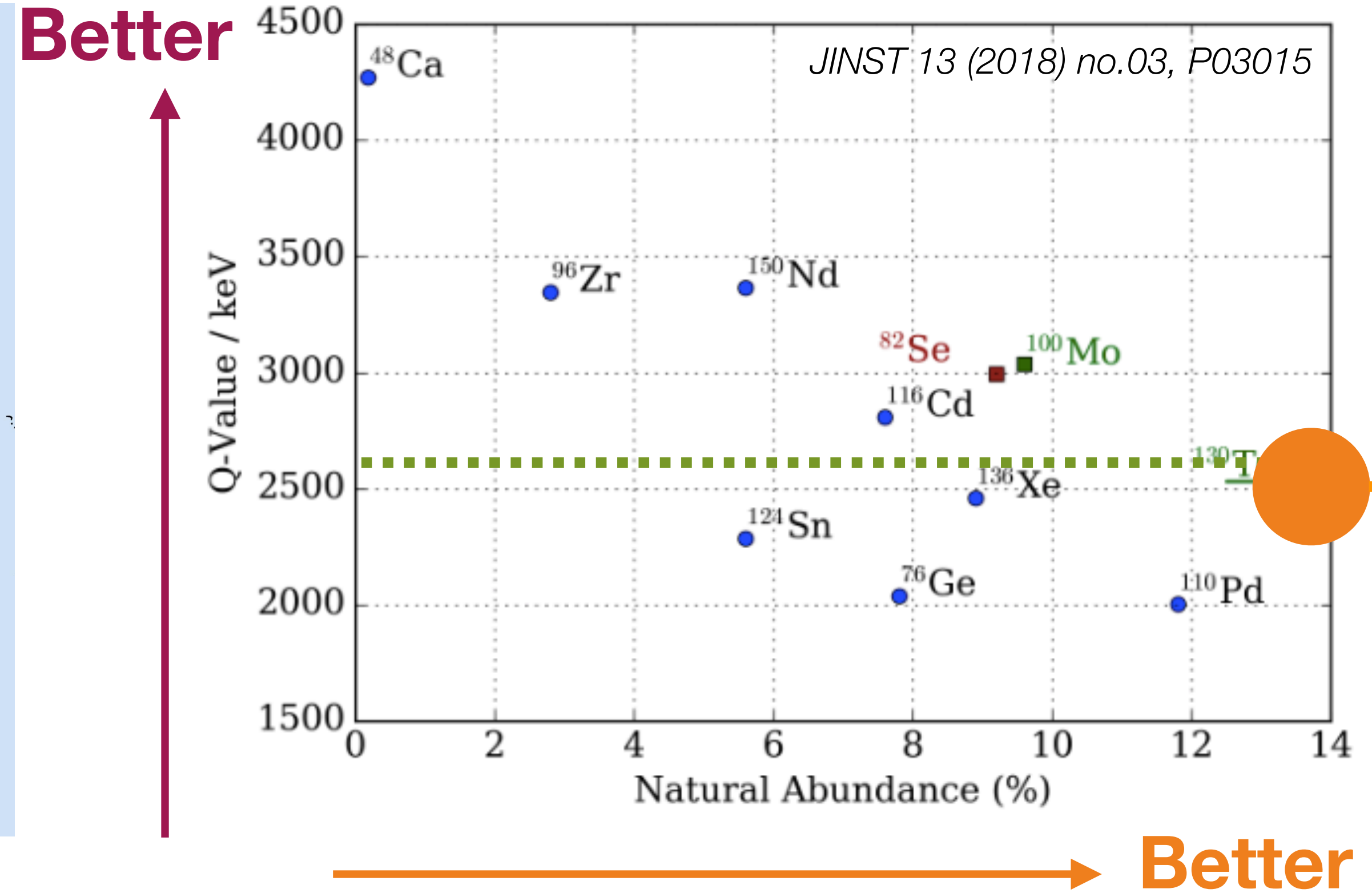
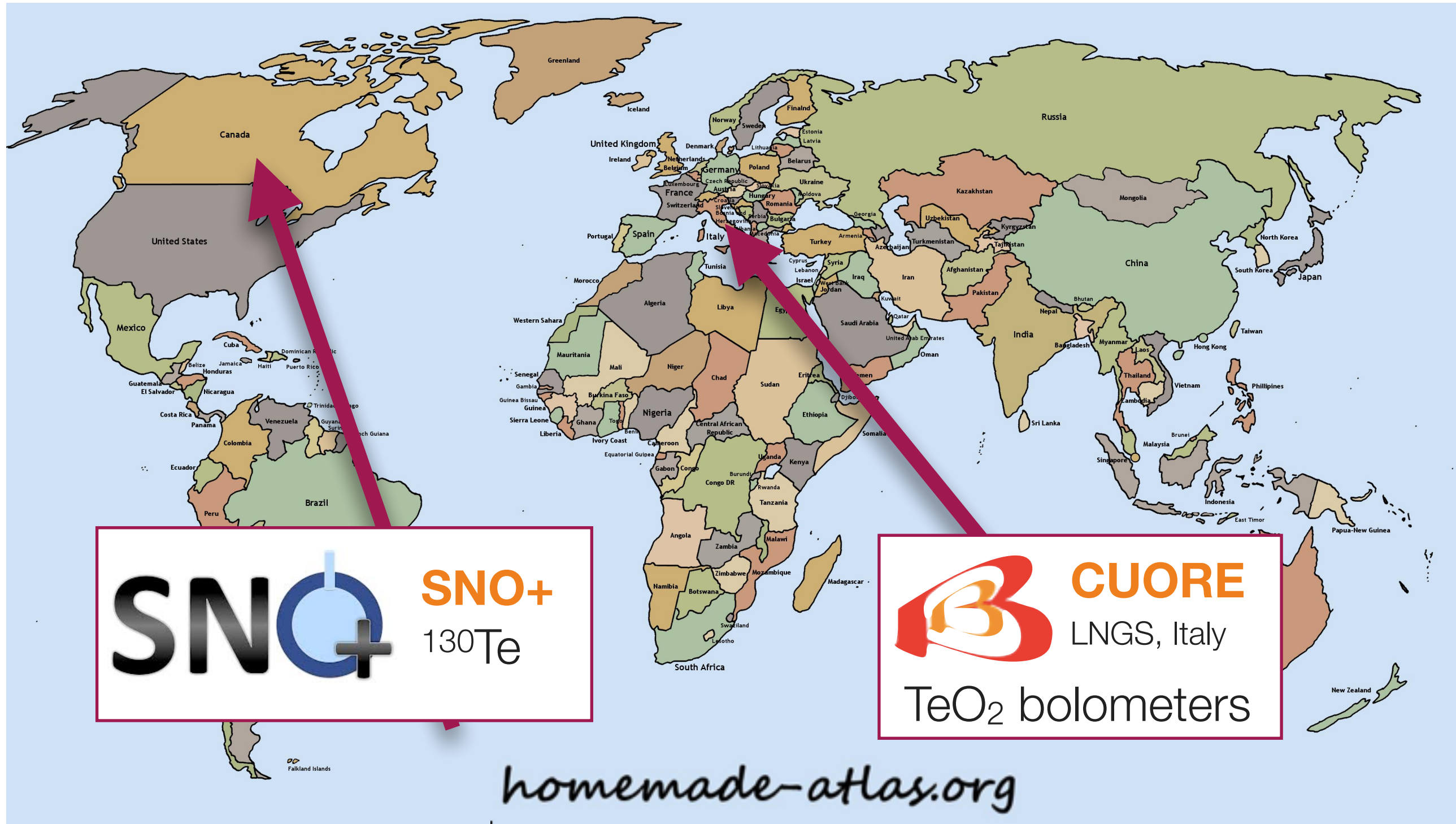


Double-beta decay experiments



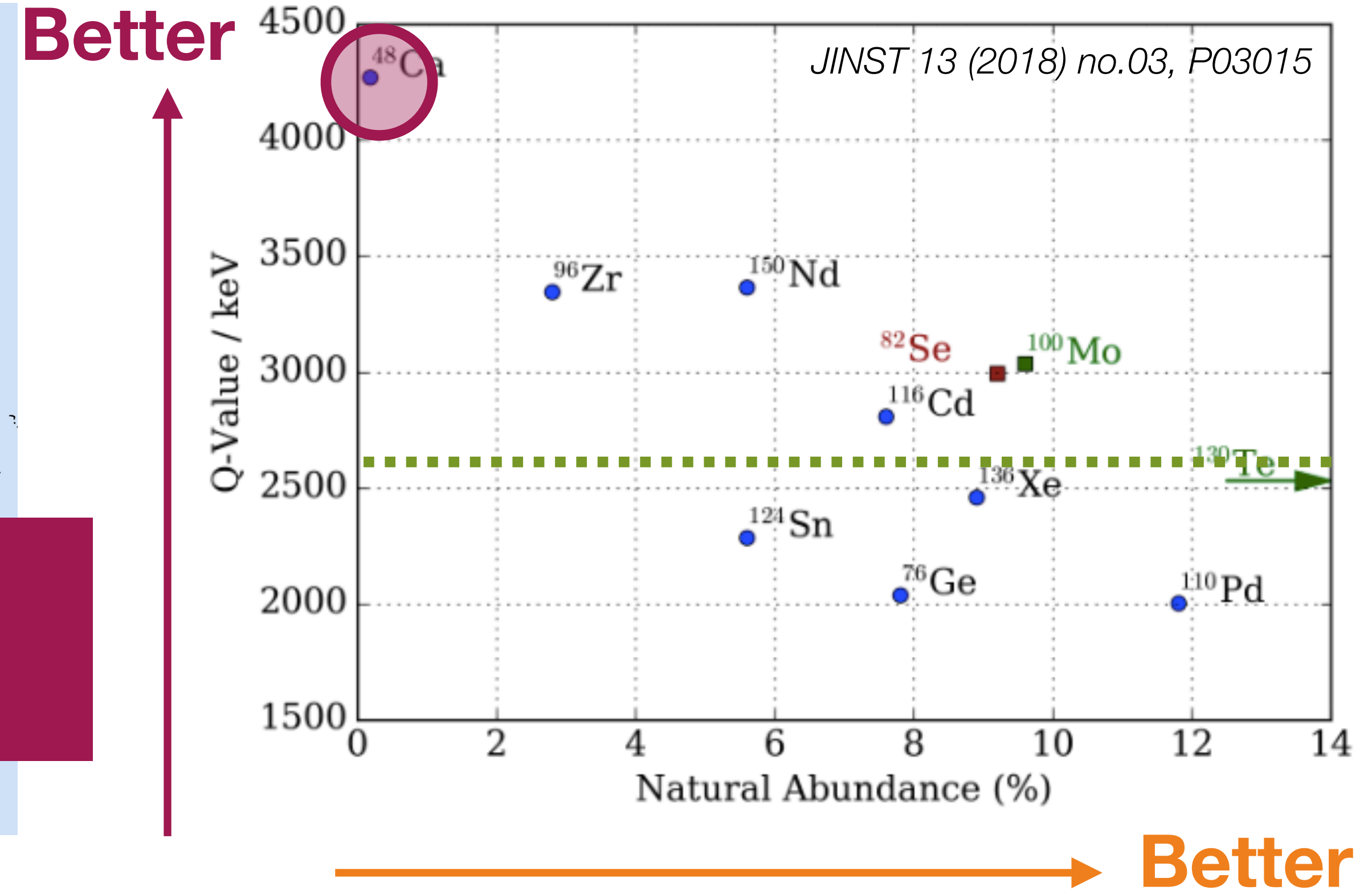
- Large detectors : **hundreds of kg** of isotope
- KamLAND-Zen has current best $0\nu\beta\beta$ half-life / $m_{\beta\beta}$ mass limit ($\langle m_{\beta\beta} \rangle < 61-165 \text{ meV}$)
- Future detectors - **nEXO, KamLAND2 Zen**

Double-beta decay experiments



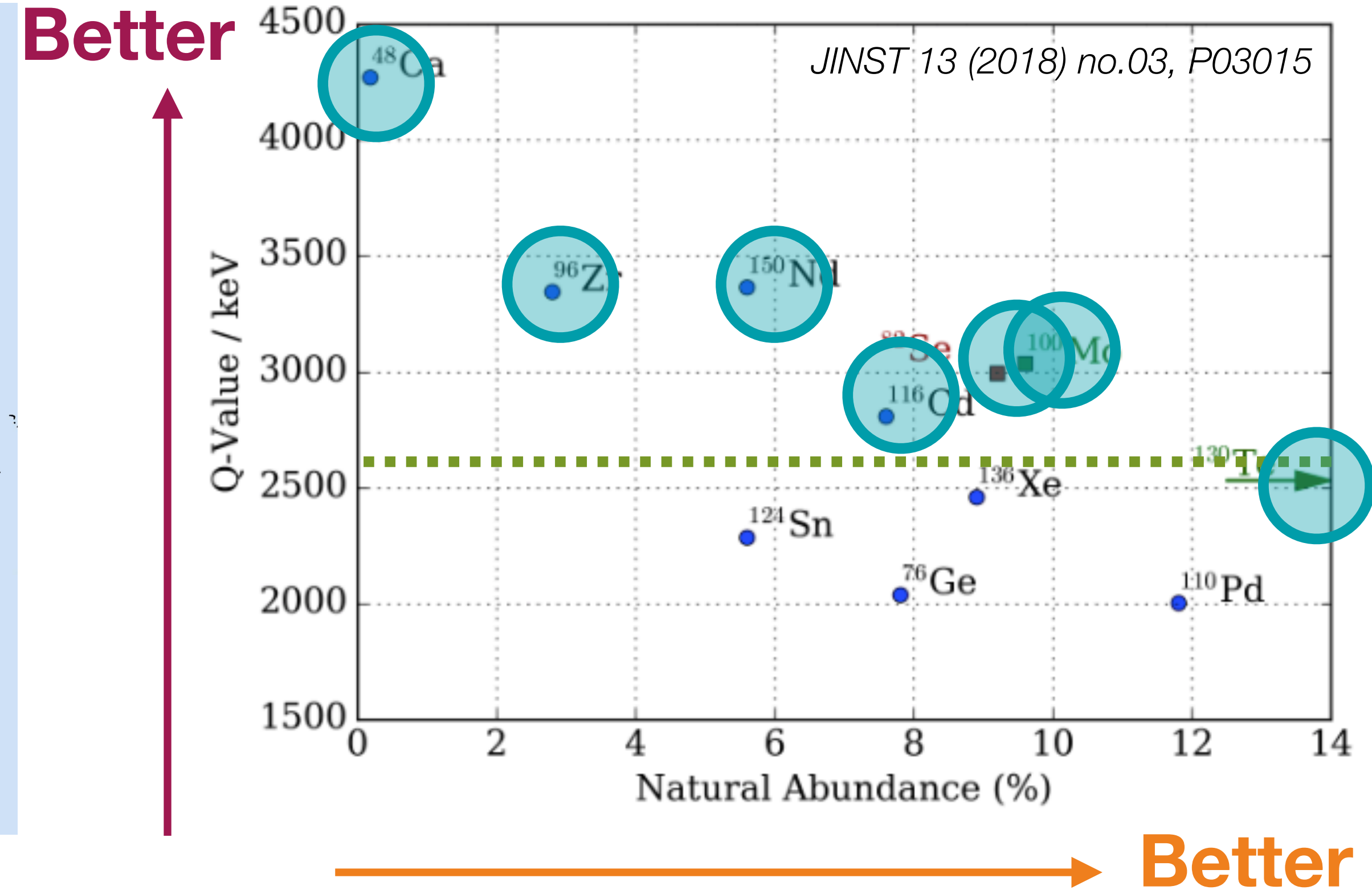
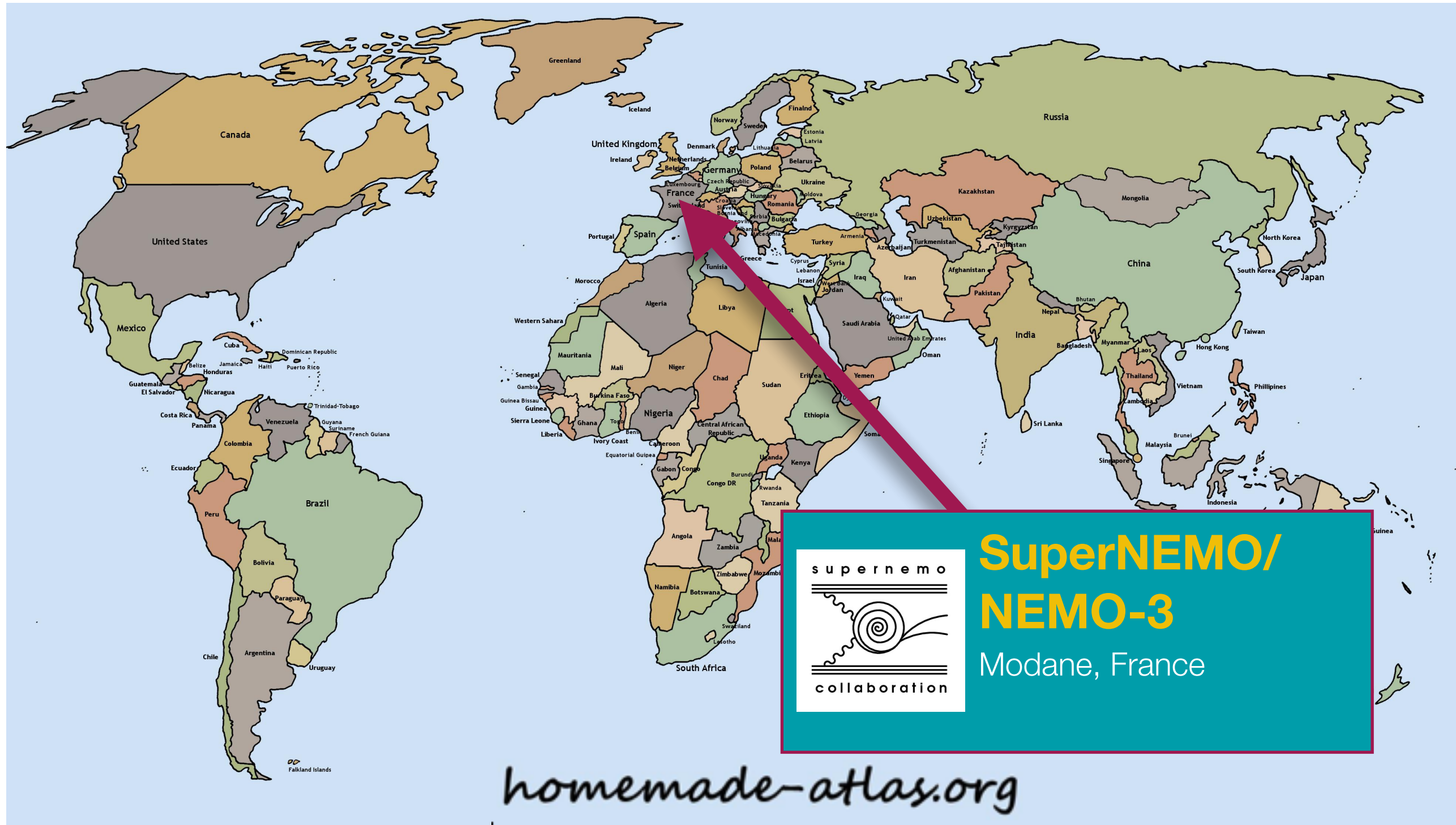
- ^{130}Te has 34% natural **abundance**

Double-beta decay experiments

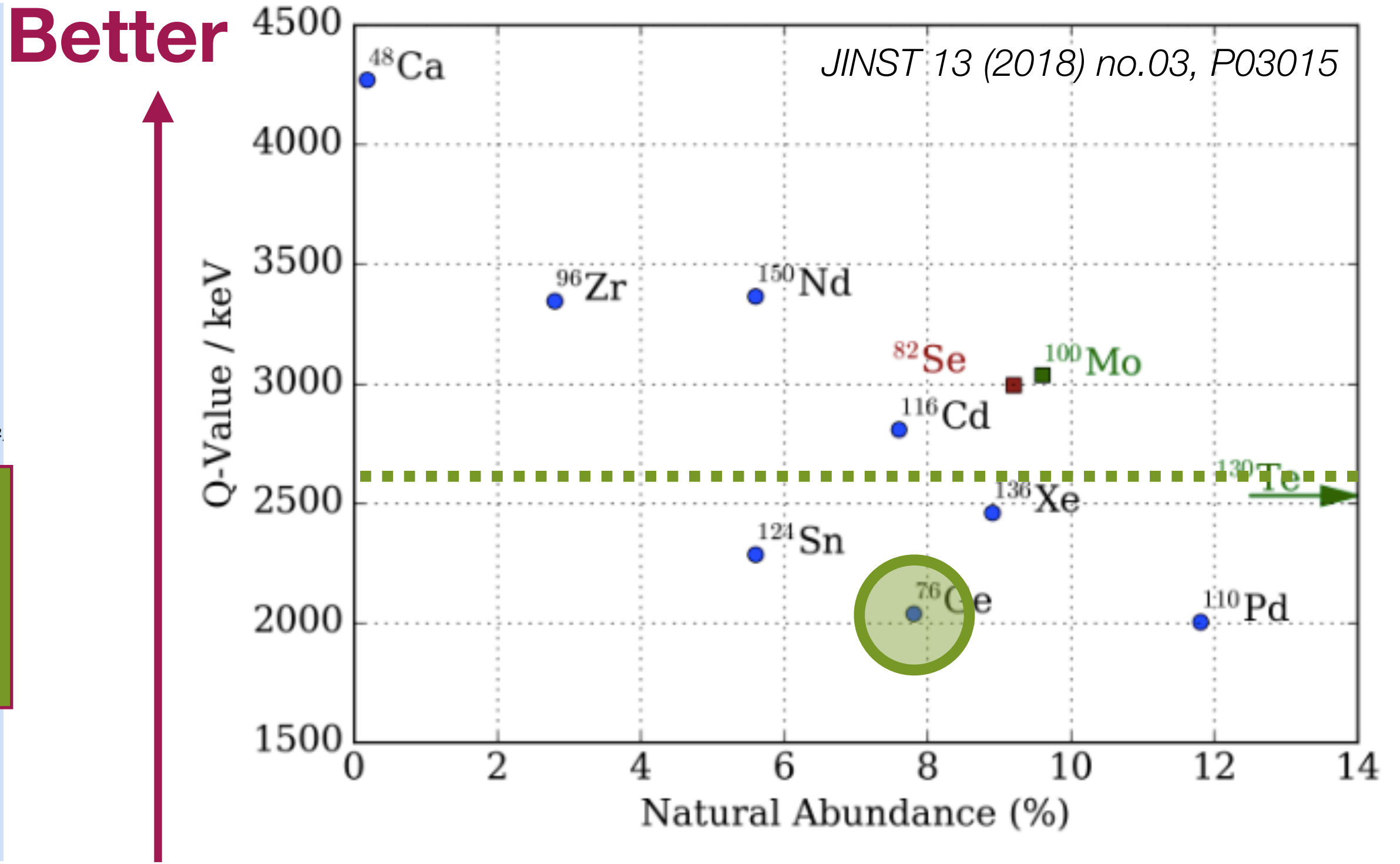
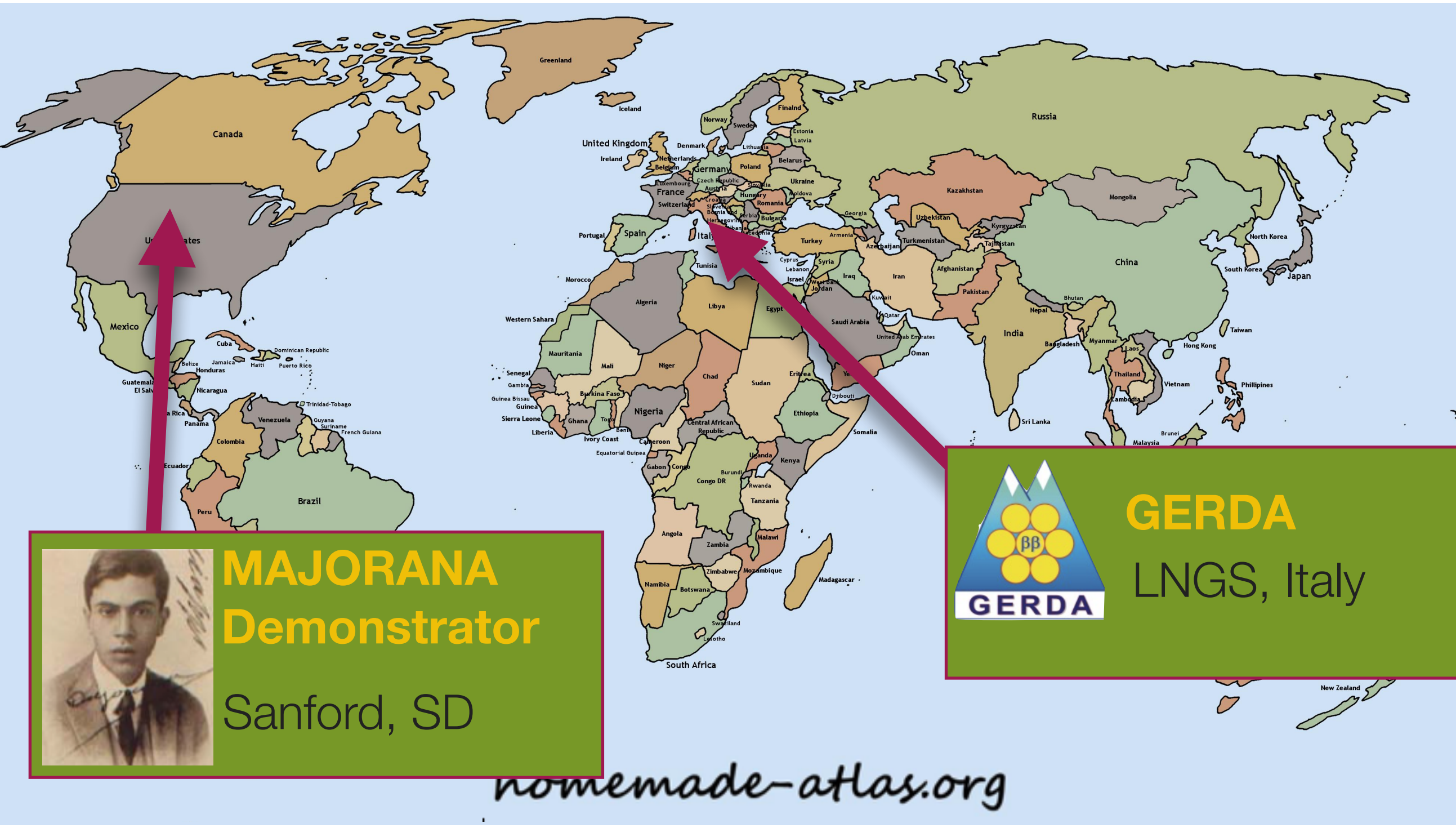


- High **Q-value** makes this an attractive isotope
- Low abundance makes **enrichment** a challenge

Double-beta decay experiments

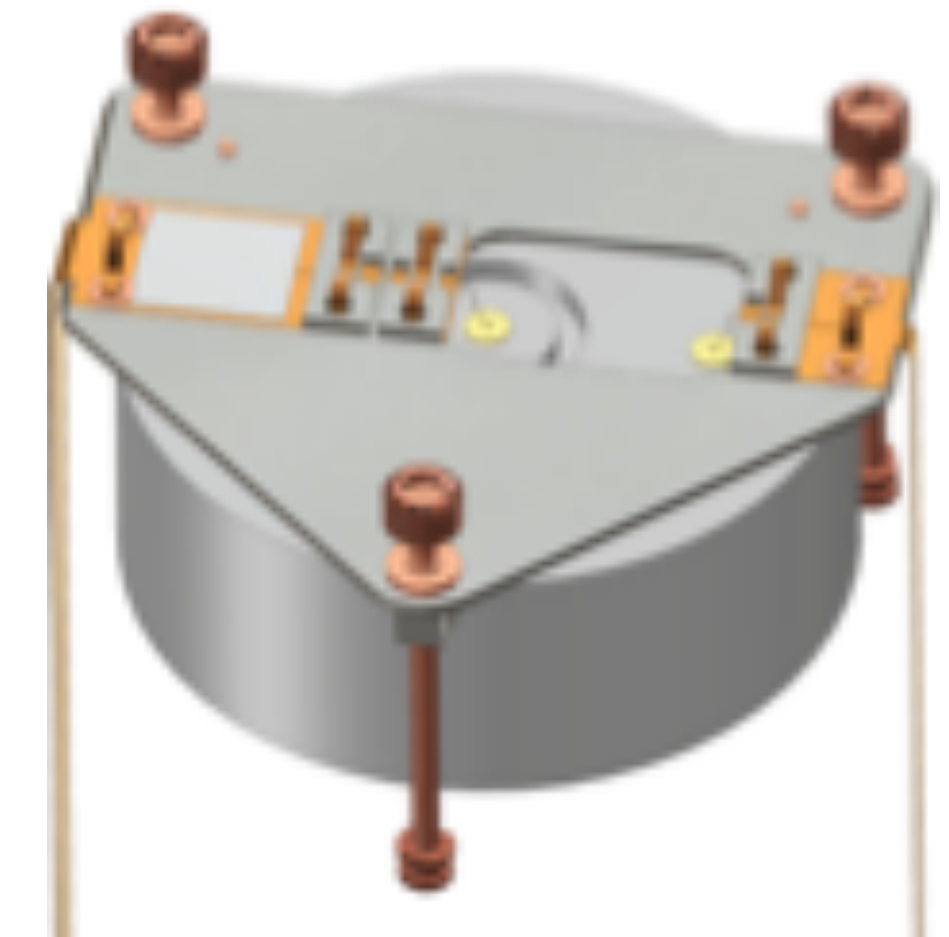
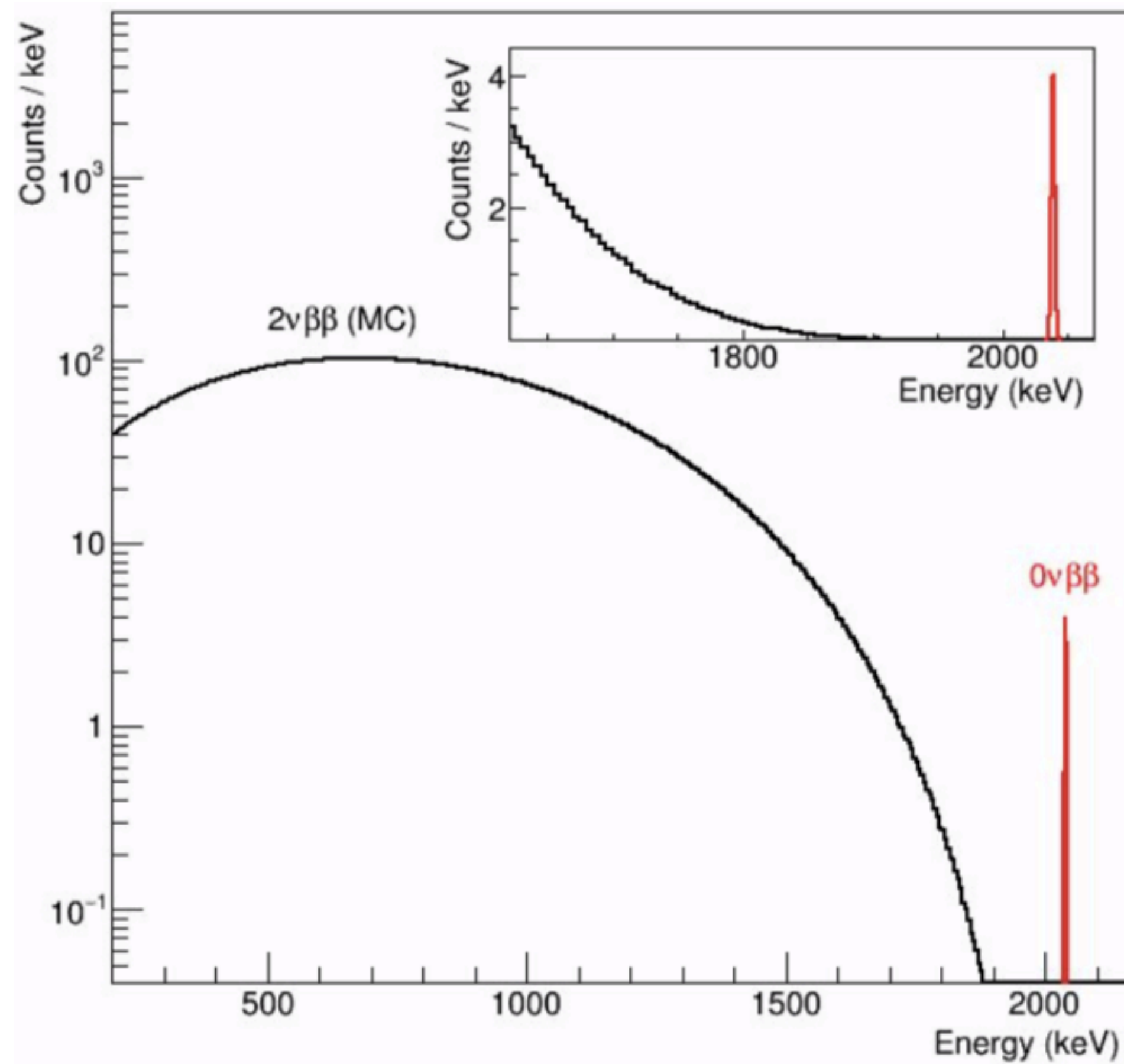


- Source separated from detector - use **any solid isotope**
- SuperNEMO - ^{82}Se , NEMO-3 mostly ^{100}Mo
- High-granularity tracker-calorimeter design gives **individual electron** energies and trajectories



- $\beta\beta$ source = detector (^{76}Ge)
- Excellent **efficiency** and **resolution**
- Future : **LEGEND**



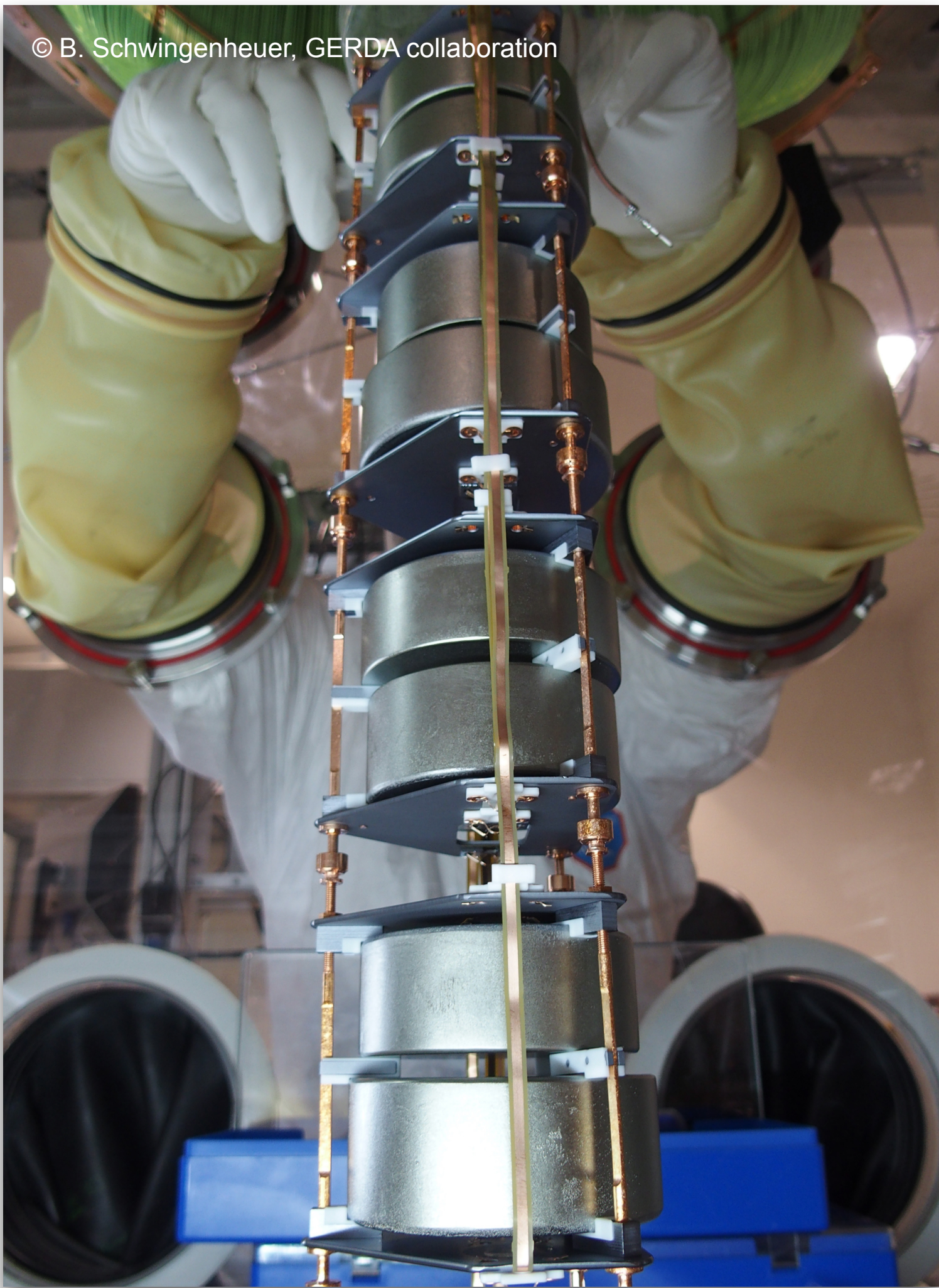


- Germanium detector, enriched to **88% ^{76}Ge** ($Q_{\beta\beta} = 2039 \text{ keV}$)
- Intrinsically **pure material** reduces **backgrounds**
- Source = detector: high detection **efficiency**
- Excellent **energy resolution** ($\sigma_E/E \approx 0.2\%$) **zero 2νββ background** in region of interest (3-4 keV)
- Electron range in Ge ~ 1 mm: ββ events are **point-like** (vs. multisite gamma backgrounds)

M Agostini, Neutrino 2016

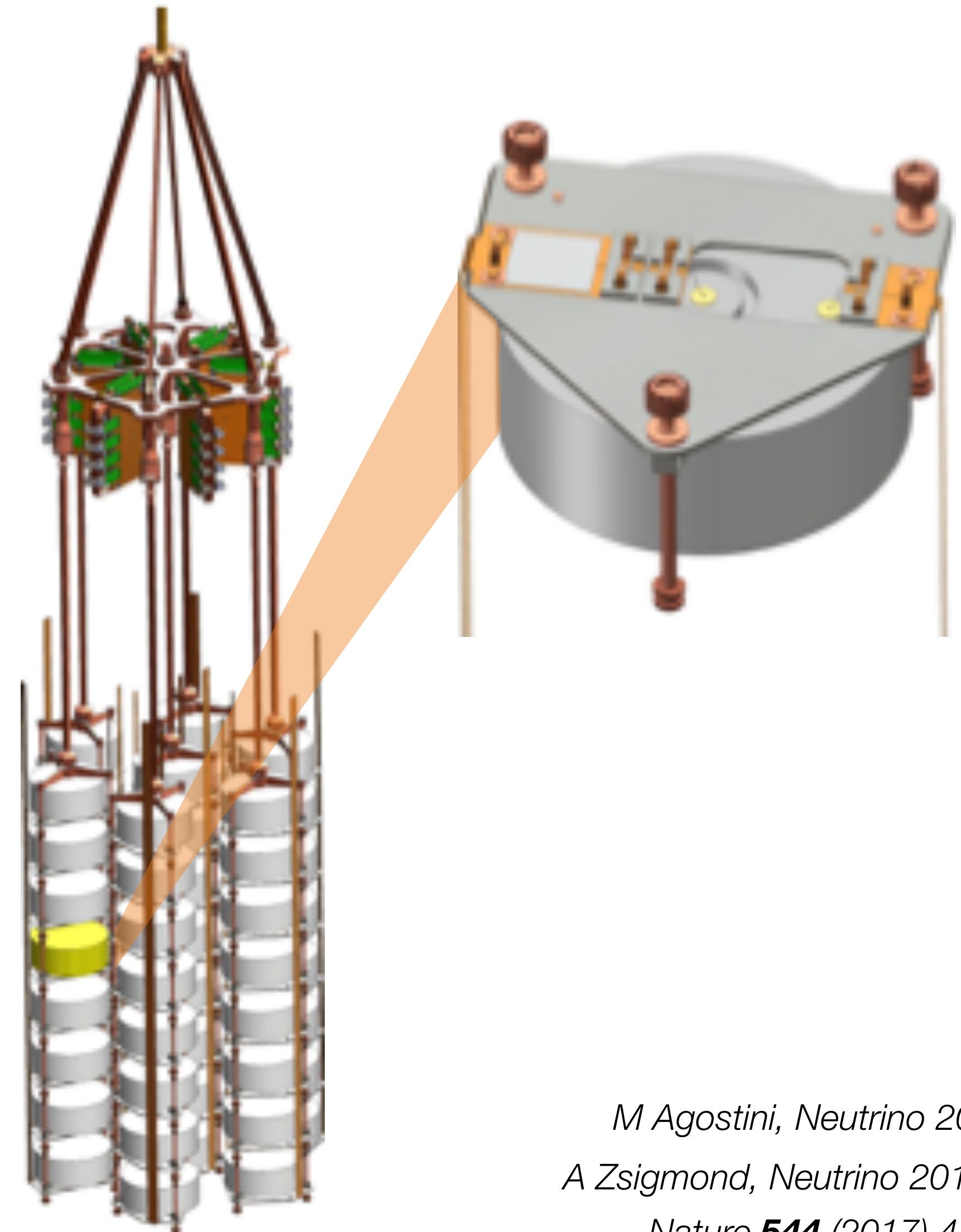
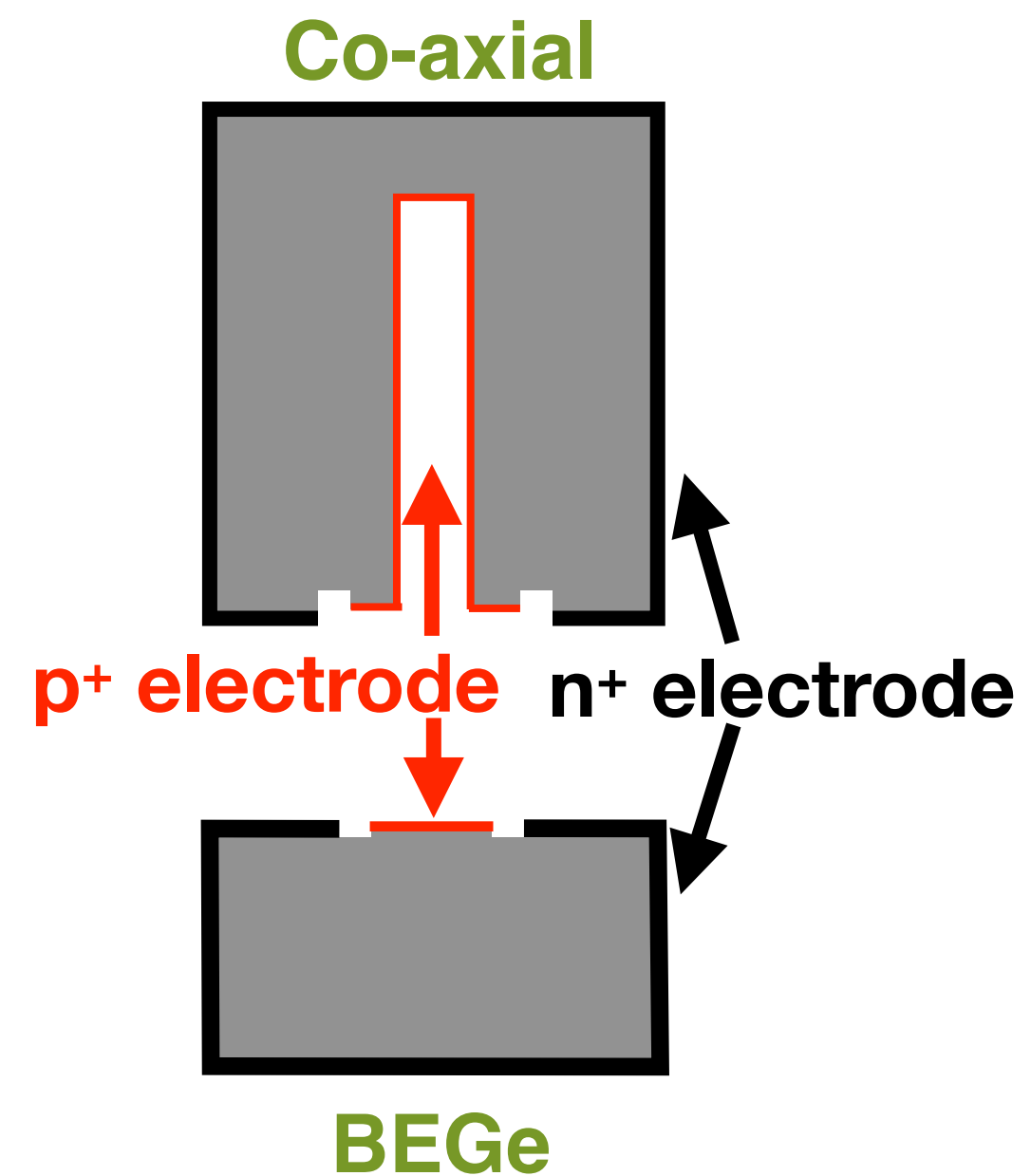
A Zsigmond, Neutrino 2018

Nature **544** (2017) 47



7 strings with **40 detectors** in total
(**43.2kg**)

- 7 enriched semi-coaxial (15.6 kg)
- 30 enriched thick window Broad Energy germanium detectors (BEGe) (20.0 kg)
- 3 natural semi-coaxial (7.6 kg)



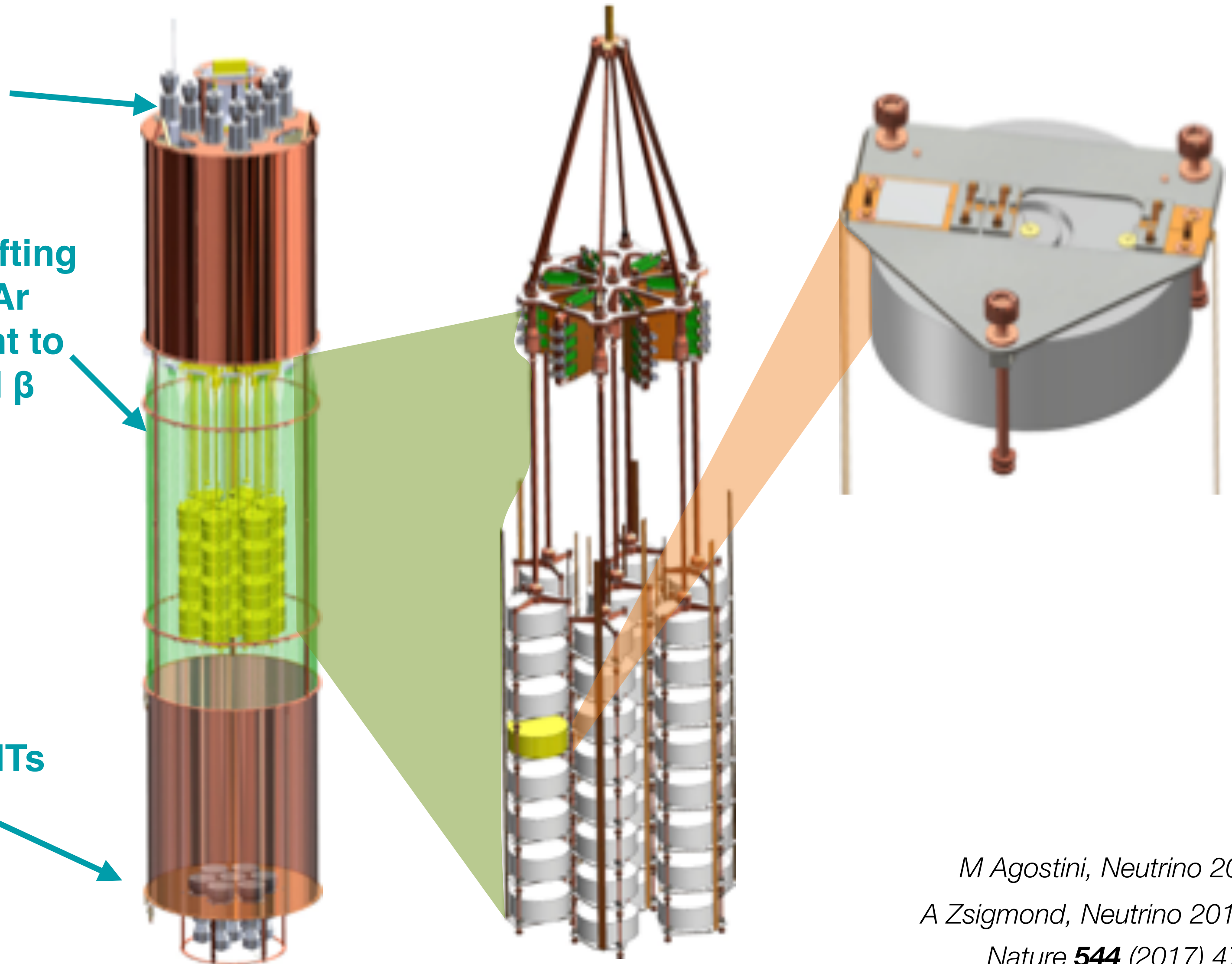
M Agostini, Neutrino 2016
A Zsigmond, Neutrino 2018
Nature **544** (2017) 47



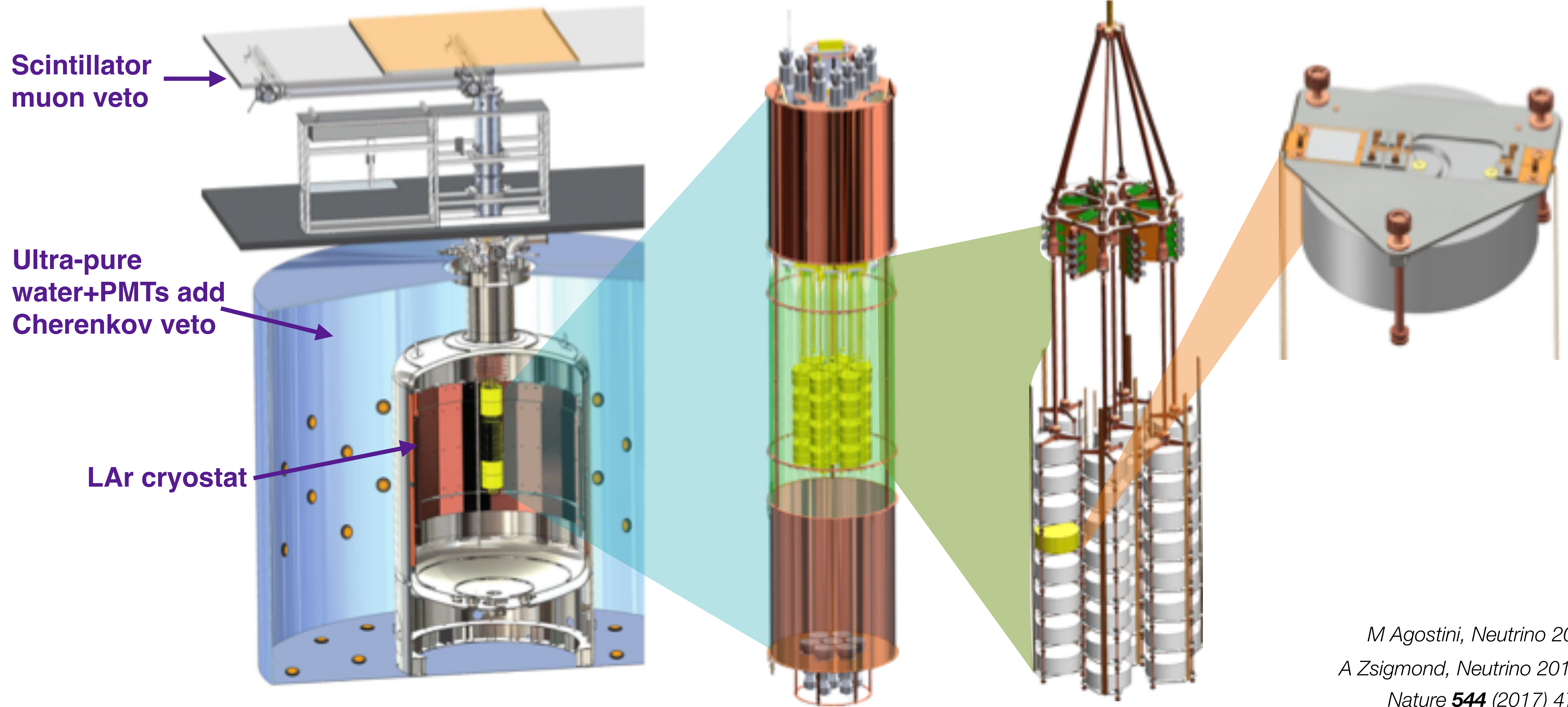
LAr electronics

Wavelength-shifting fibres deliver LAr scintillation light to SiPMs for γ and β veto

Low-activity PMTs



M Agostini, Neutrino 2016
A Zsigmond, Neutrino 2018
Nature **544** (2017) 47

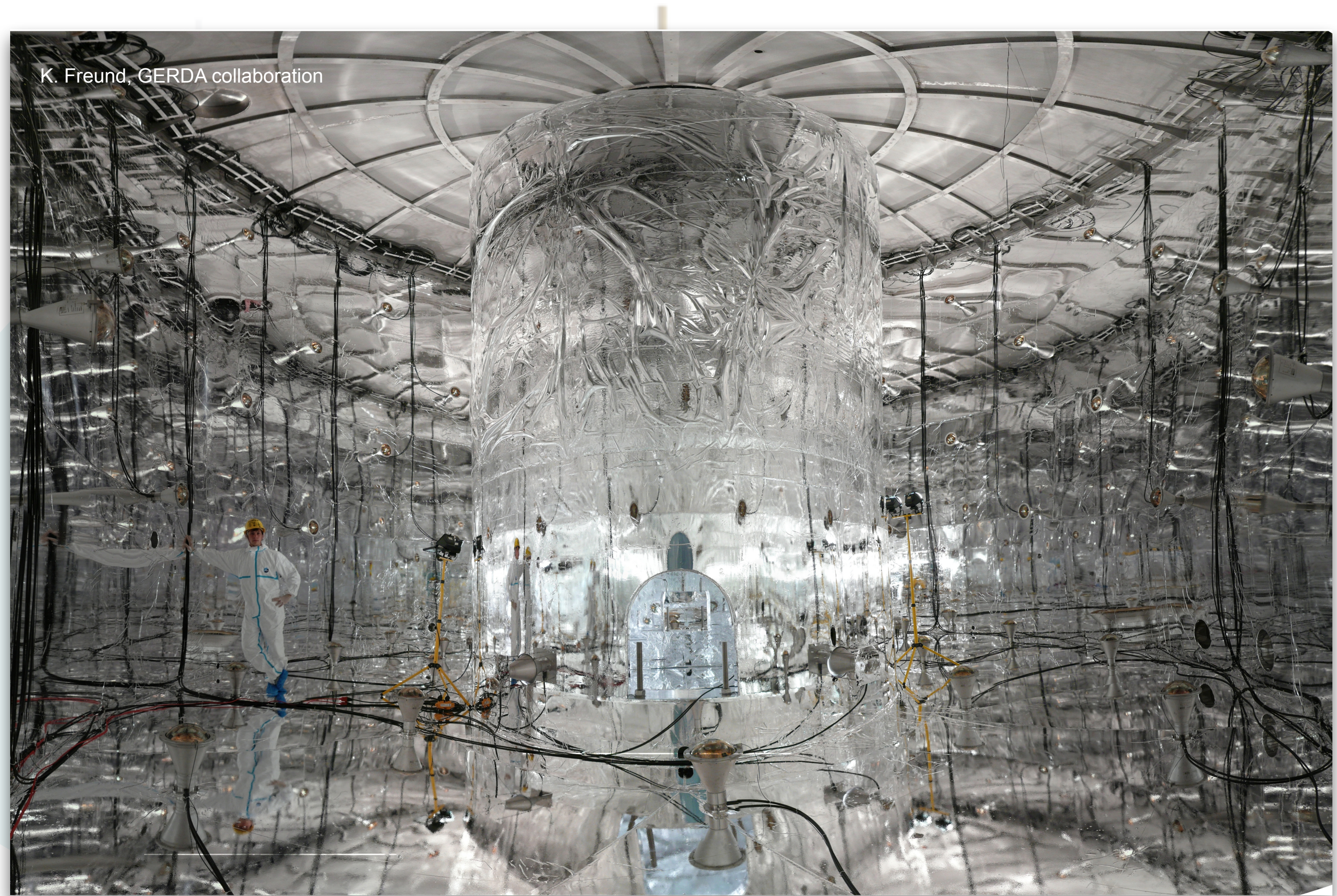
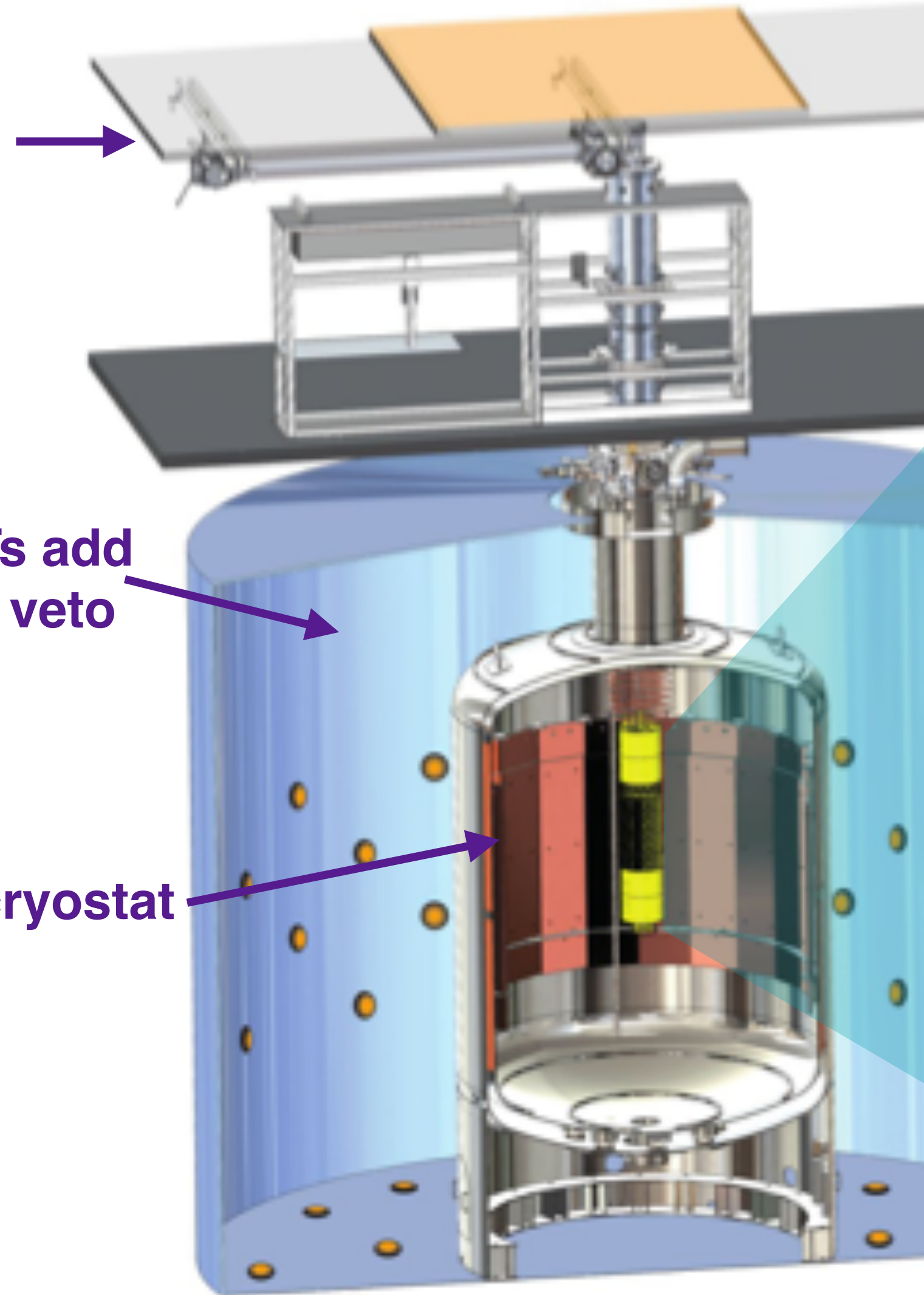


M Agostini, Neutrino 2016
A Zsigmond, Neutrino 2018
Nature **544** (2017) 47

Scintillator
muon veto

Ultra-pure
water+PMTs add
Cherenkov veto

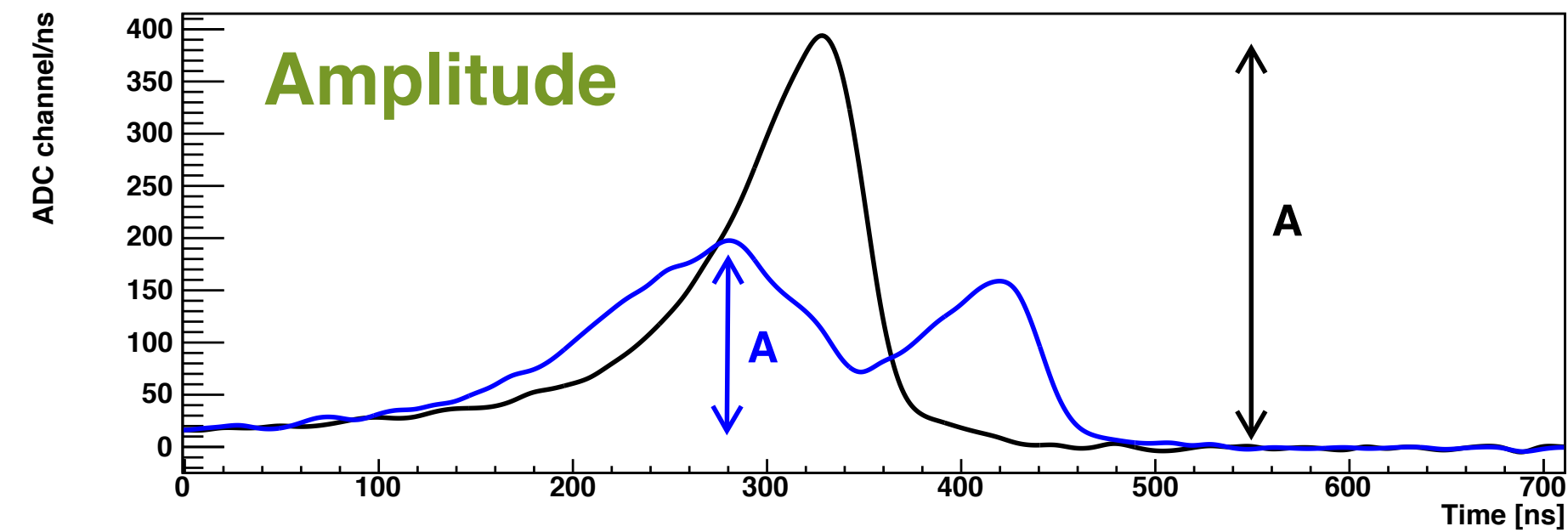
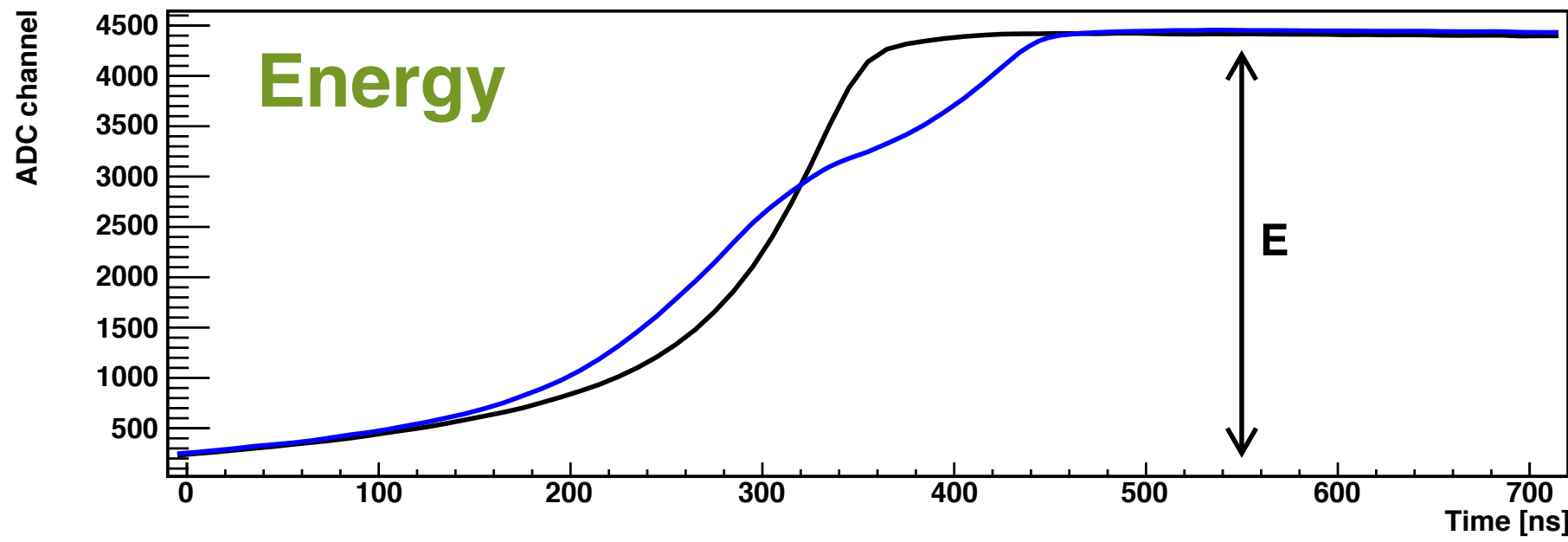
LAr cryostat



M Agostini, Neutrino 2016

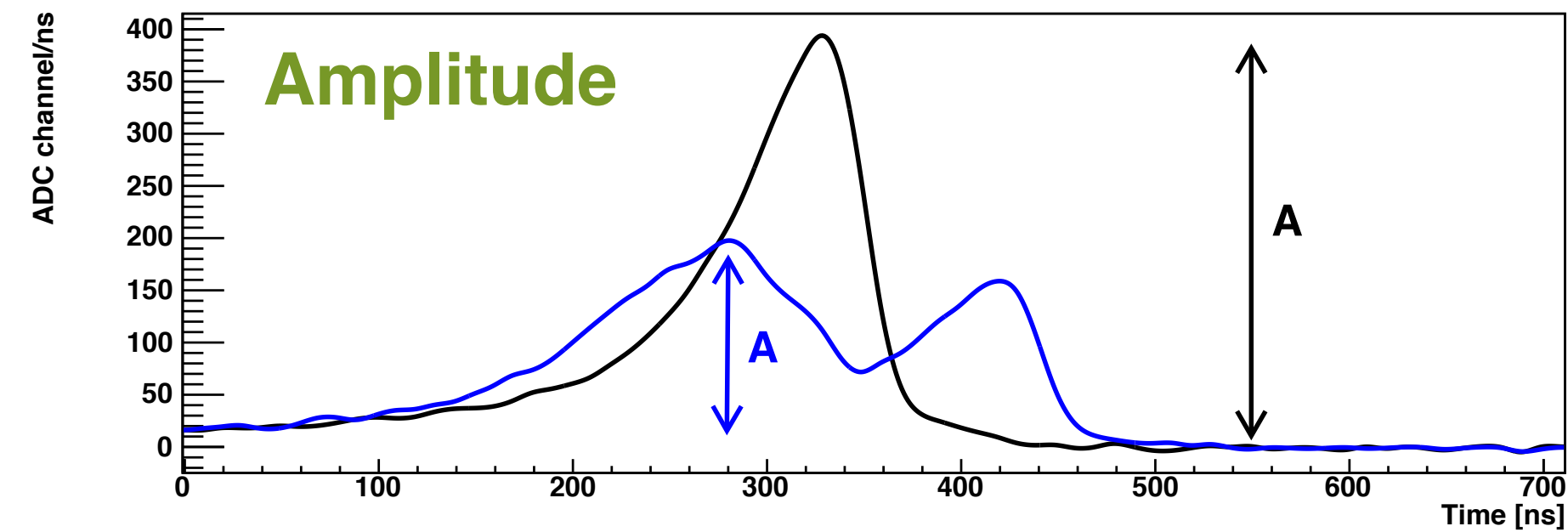
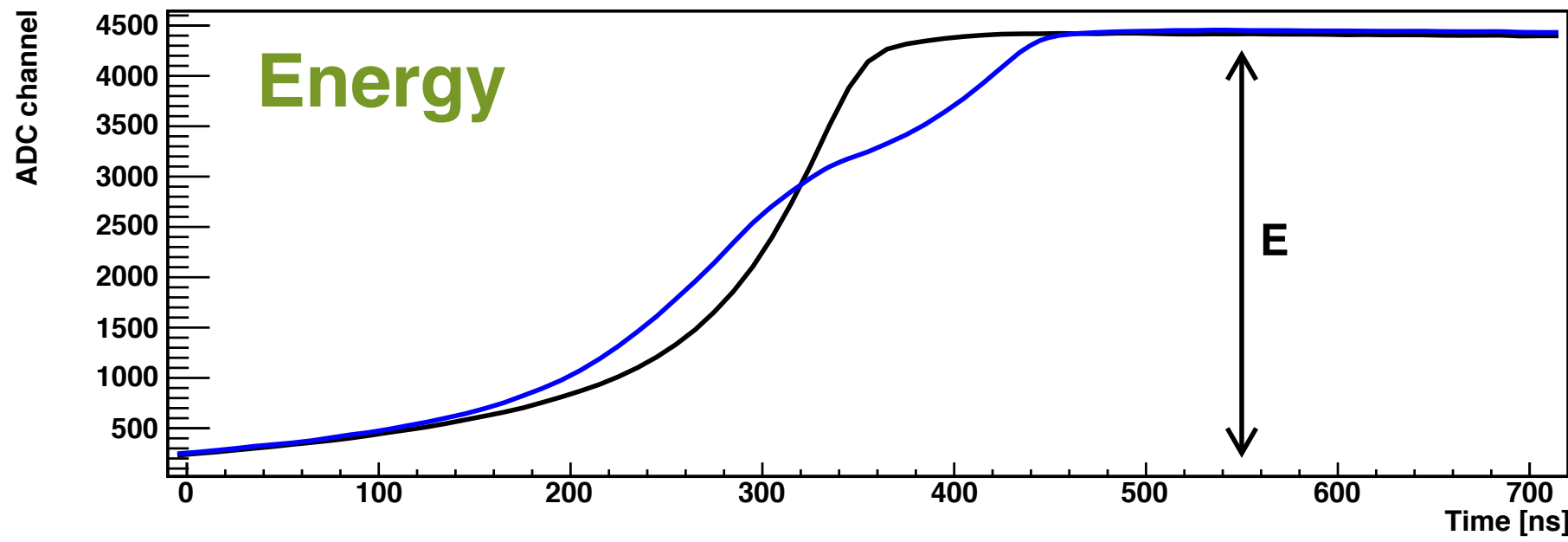
A Zsigmond, Neutrino 2018

Nature **544** (2017) 47



Pulse shape (energy and amplitude) distinguishes **single ($\beta\beta$)** vs **multi-site (background)** events. Plots for BEGe detector.

2011 JINST 6 P04005



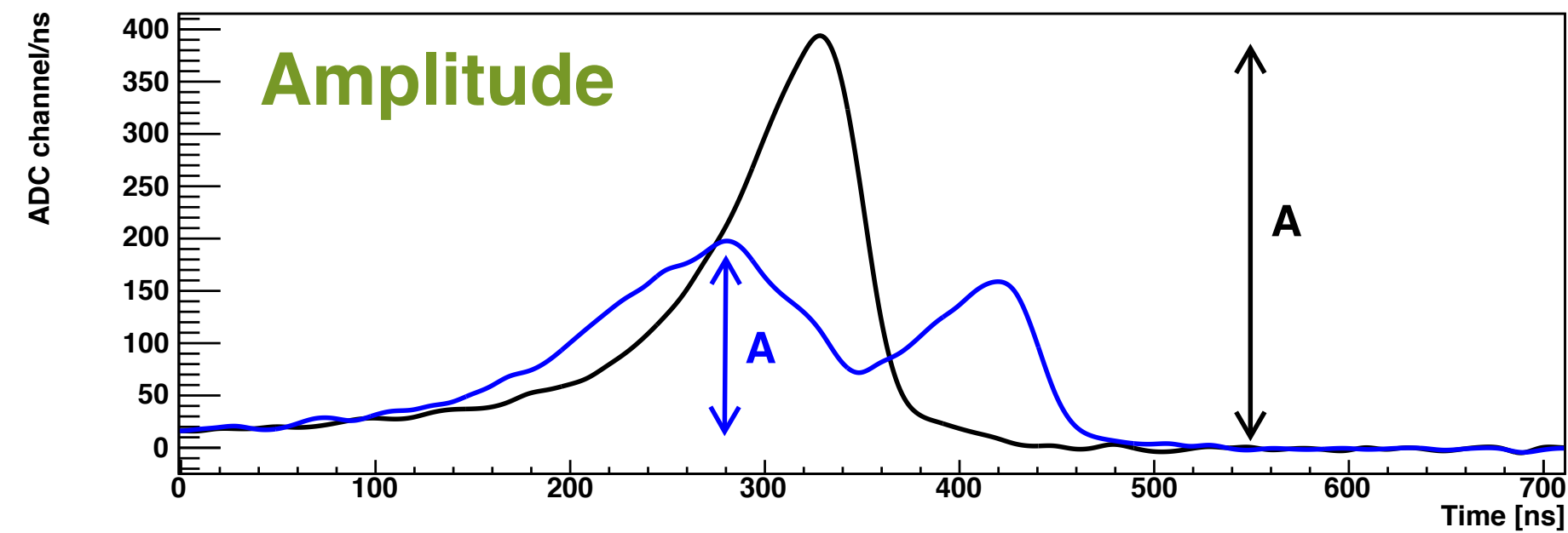
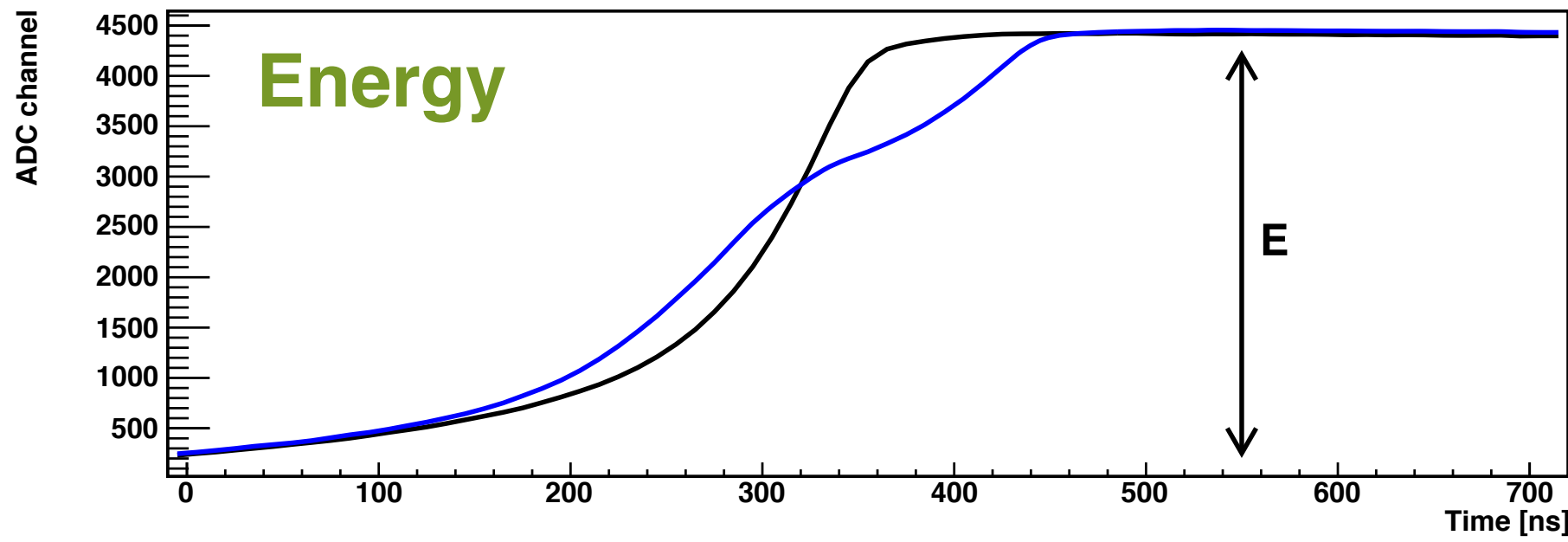
Pulse shape (energy and amplitude) distinguishes **single ($\beta\beta$)** vs **multi-site (background)** events. Plots for BEGe detector.

2011 JINST 6 P04005

Muon veto from plastic scintillator and Cherenkov light in water: **> 99% rejection efficiency**



EPJ C, 2016, Vol 76, Number 5, Page 1



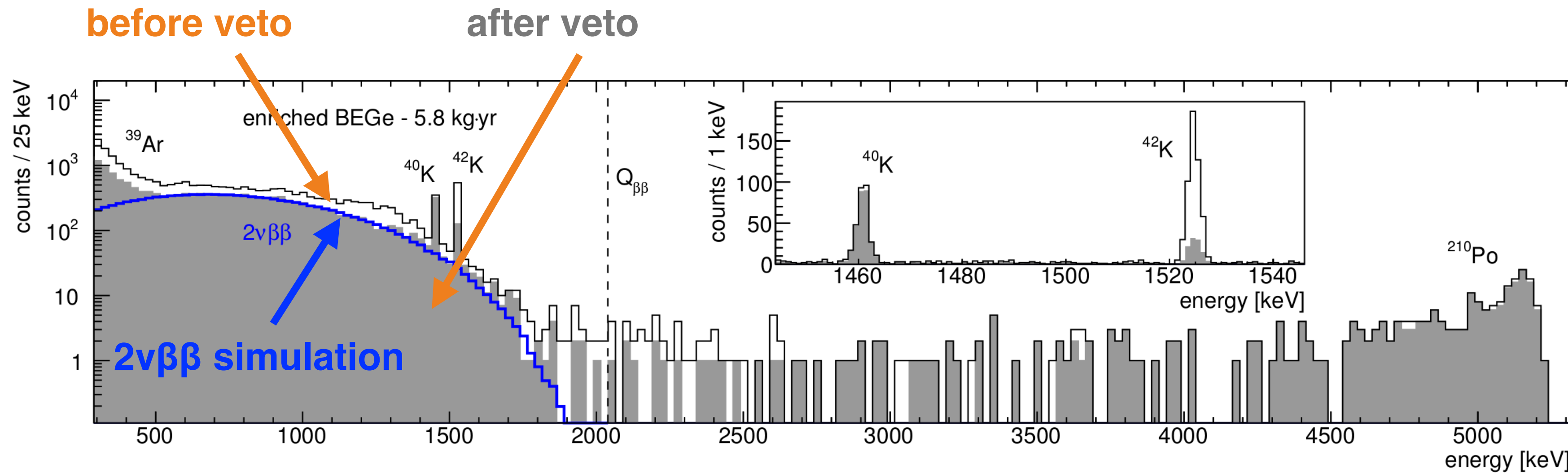
Pulse shape (energy and amplitude) distinguishes **single ($\beta\beta$)** vs **multi-site (background)** events. Plots for BEGe detector.

2011 JINST 6 P04005

Muon veto from plastic scintillator and Cherenkov light in water: **> 99% rejection efficiency**

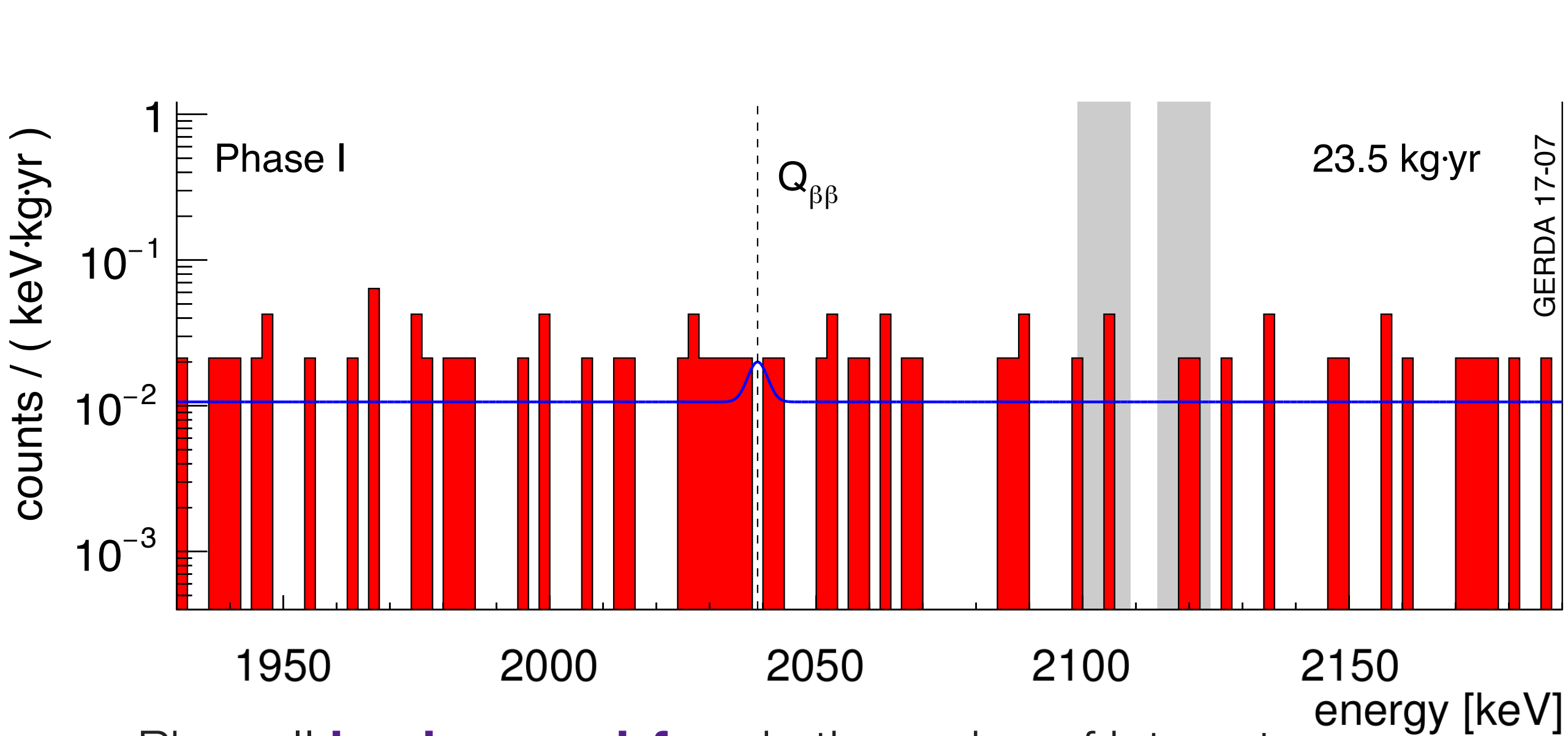


EPJ C, 2016, Vol 76, Number 5, Page 1



Liquid argon veto removes external backgrounds to give clean **$2\nu\beta\beta$ spectrum** (98% efficiency). Plot for BEGe detectors.

Nature **544** (2017) 47-52



- Phase II **background-free** in the region of interest

$$1.0^{+0.6}_{-0.4} \times 10^{-3} \text{cts} / (\text{keV} \cdot \text{kg} \cdot \text{yr})$$

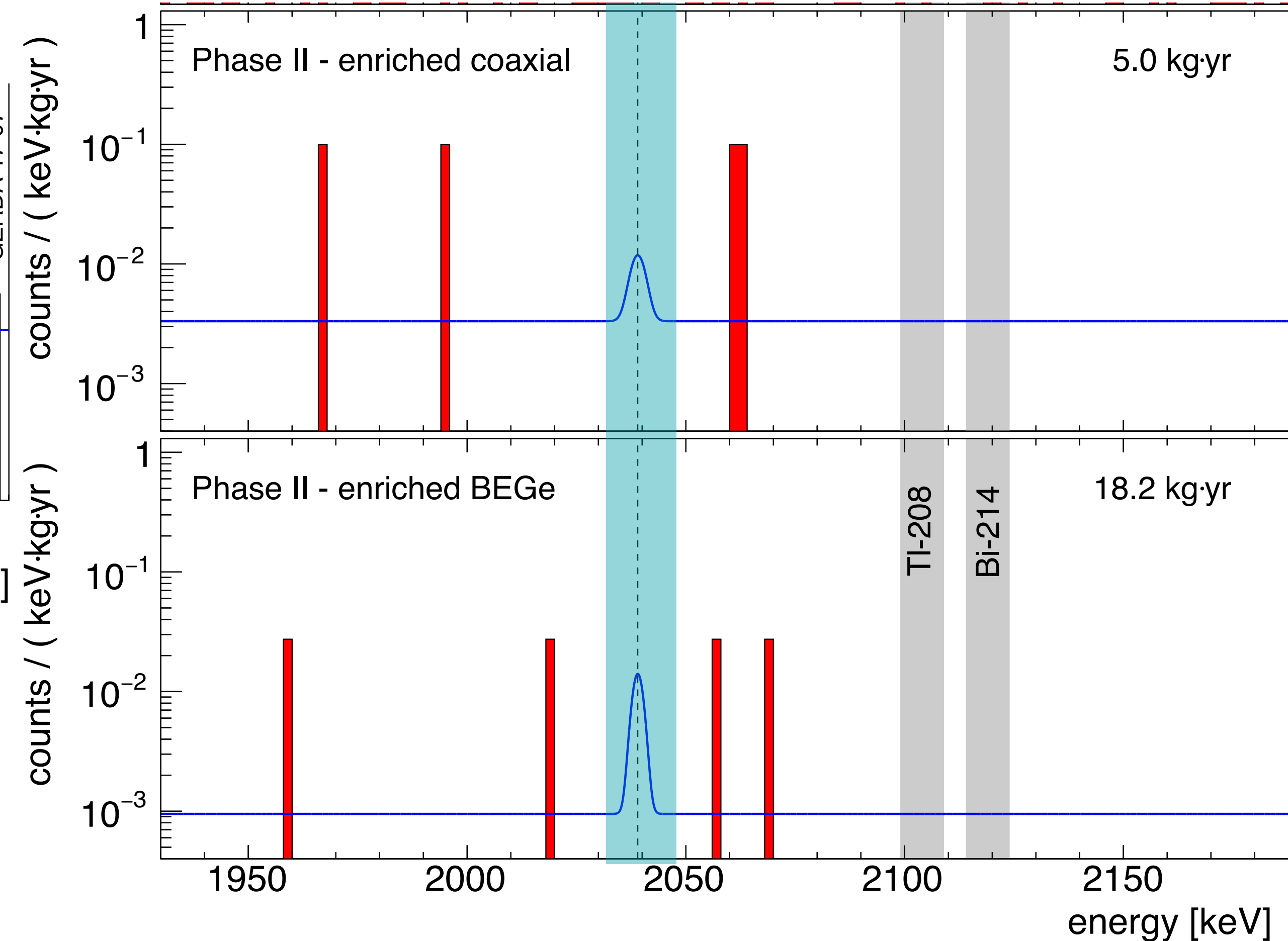
- Limit** set at 90% confidence level: 2 events in 34.4 kg·yr, or

$$T_{1/2}^{0\nu} > 8.0 \times 10^{25} \text{yr}$$

for light Majorana neutrino exchange:

$$m_{\beta\beta} < 0.12 - 0.26 \text{eV}$$

- At **design exposure** of 100 kg·yr, **sensitivity** $> 10^{26}$ yr



GERDA cryostat can hold 200kg of detectors

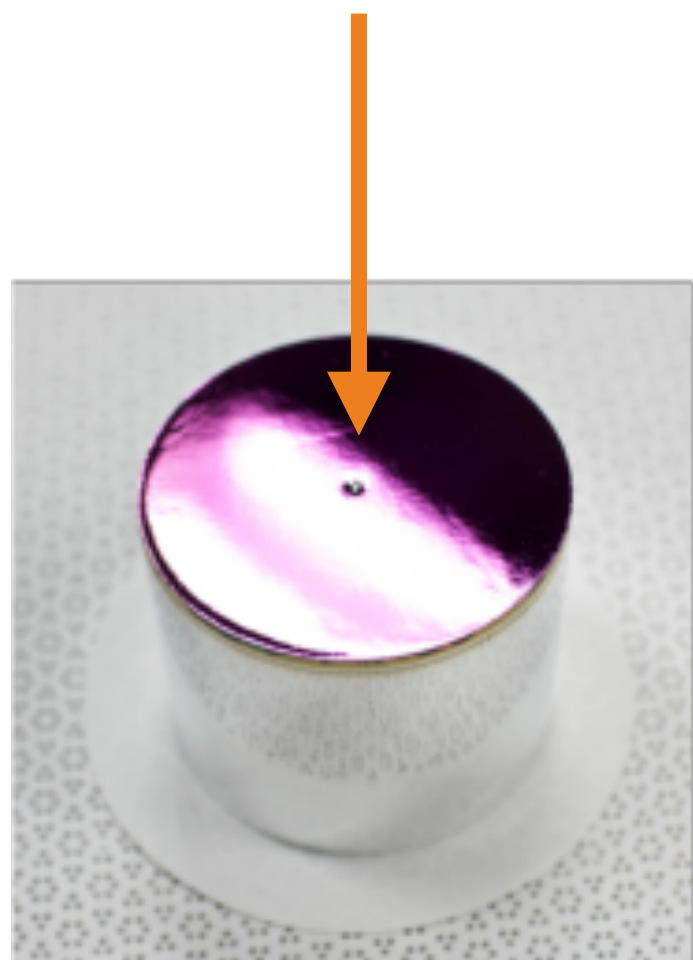
Phys. Rev. Lett. 120 (2018) 132503



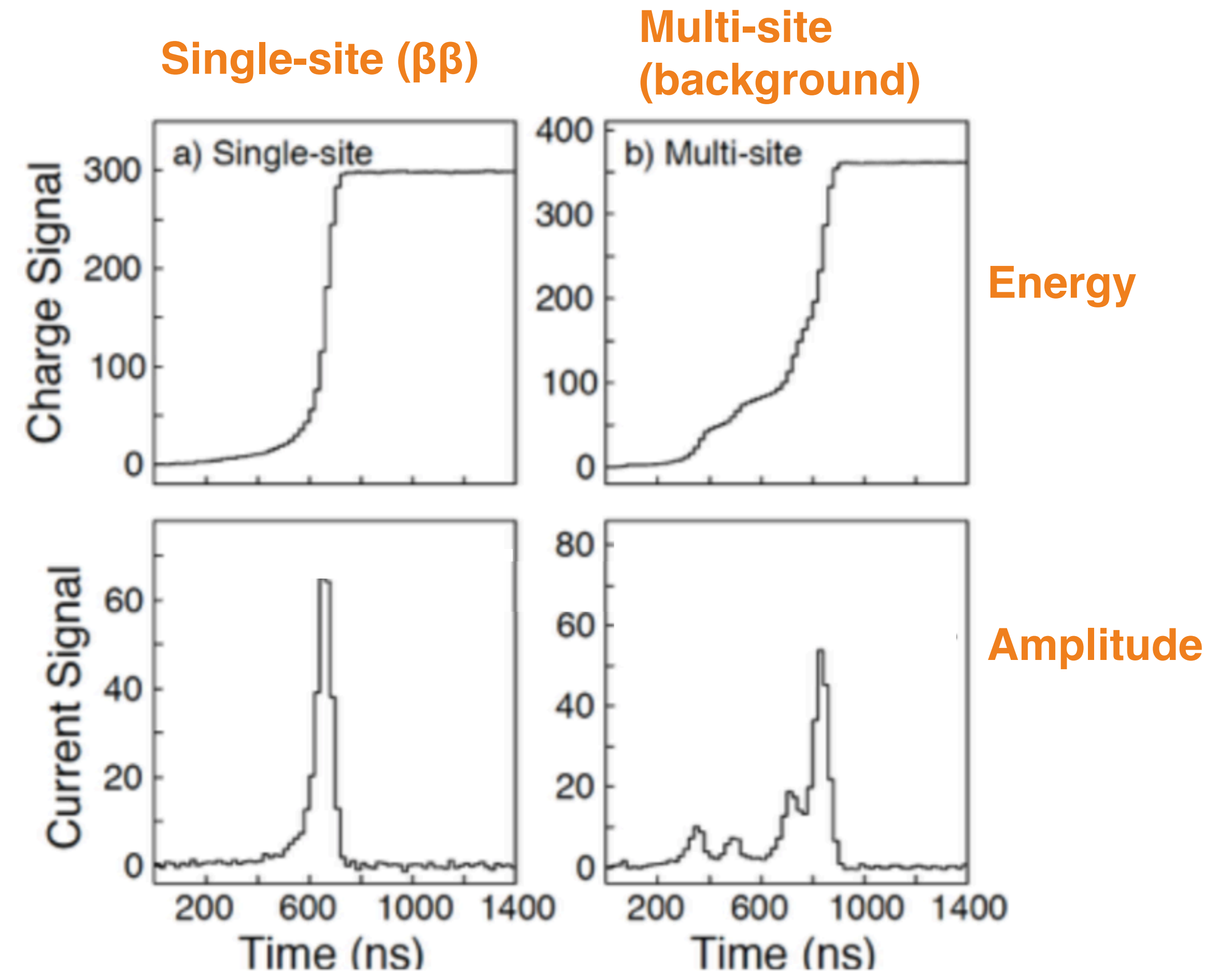
More ^{76}Ge : MAJORANA Demonstrator at Sanford, SD

- Enriched **p-type point-contact HPGe detectors** have
 - Excellent **pulse-shape discrimination**
 - Benefits of standard **coaxial HPGe** detectors
 - Low capacitance ($\sim 1\text{pF}$): excellent low-energy resolution

Small, dimple-like central contact

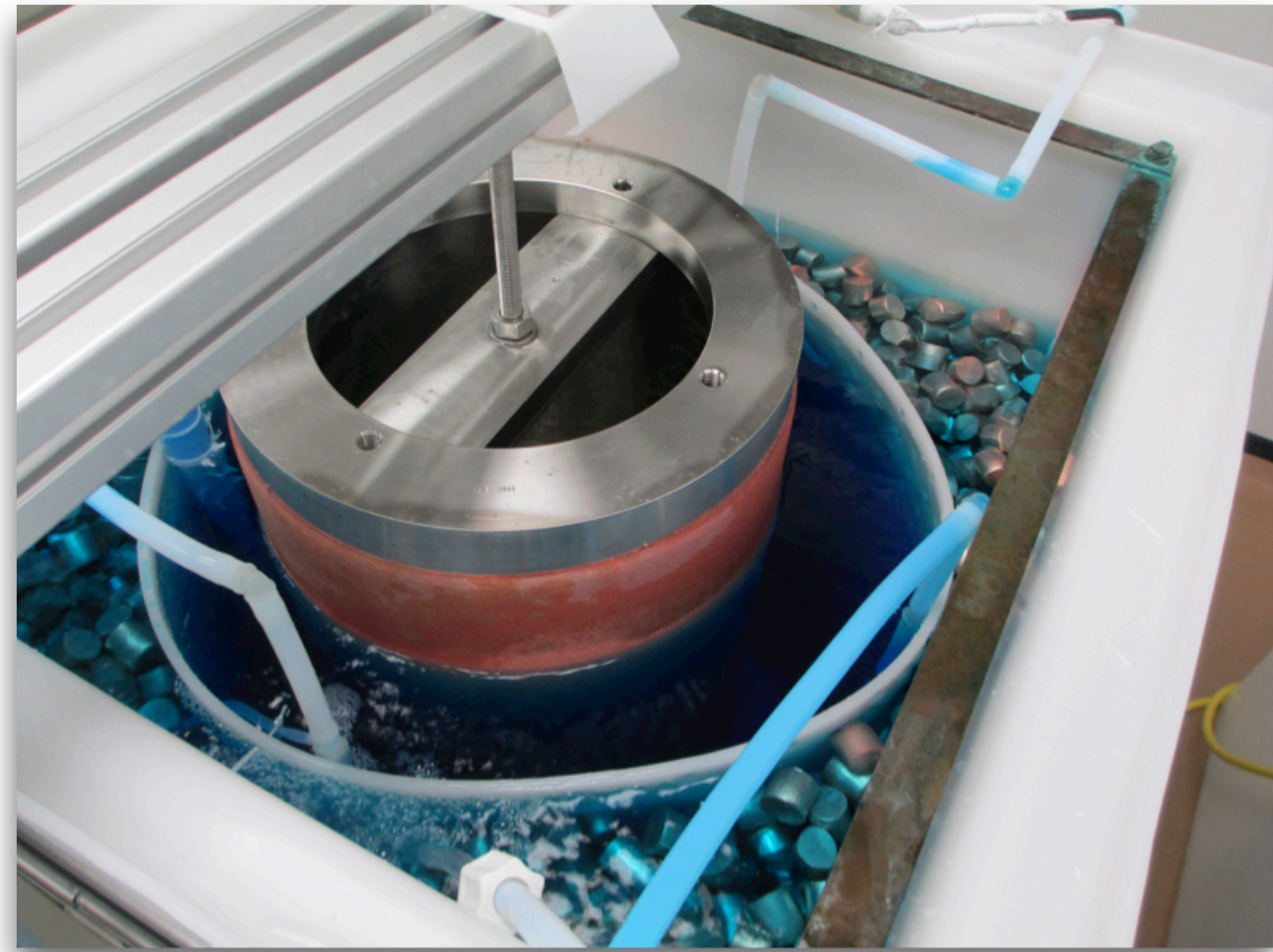


PPC HPGe
Detector



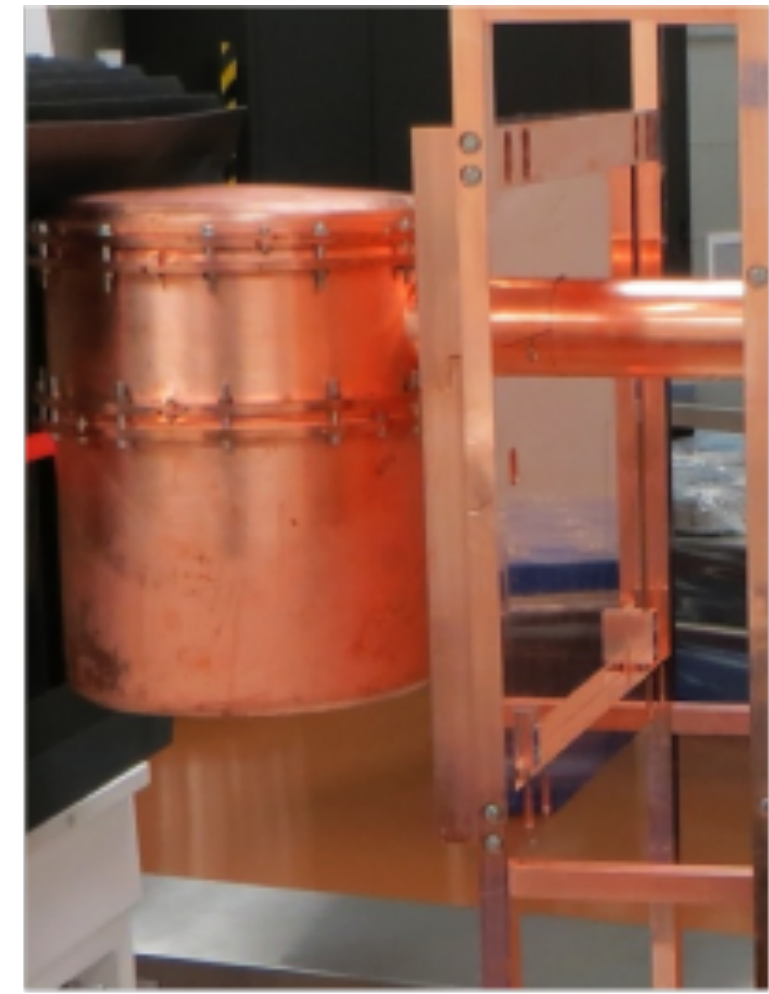
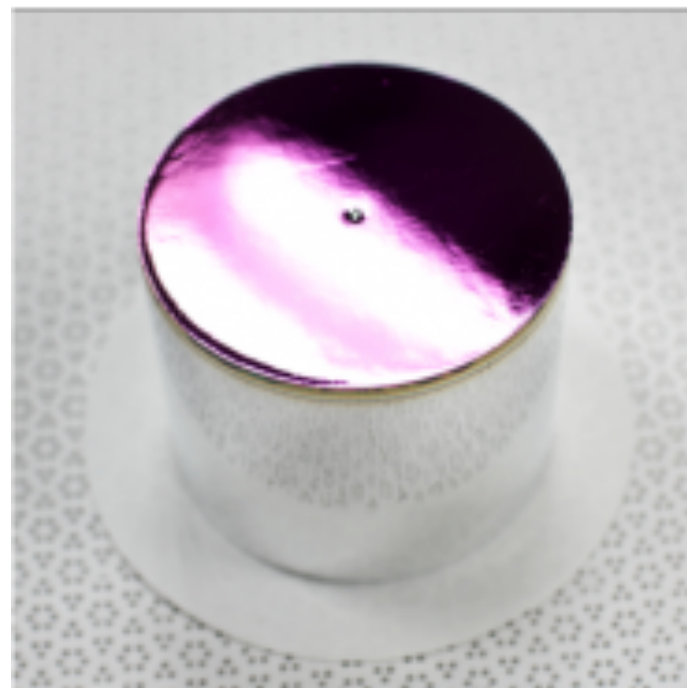


More ^{76}Ge : MAJORANA Demonstrator at Sanford, SD



- All components made from highly **radiopure materials**
 - **Underground electroformed copper**
 - **NXT-85** cleanroom-manufactured Teflon

10x less radioactive than commercial copper



PPC HPGe Detector

Low-Mass Mount

String

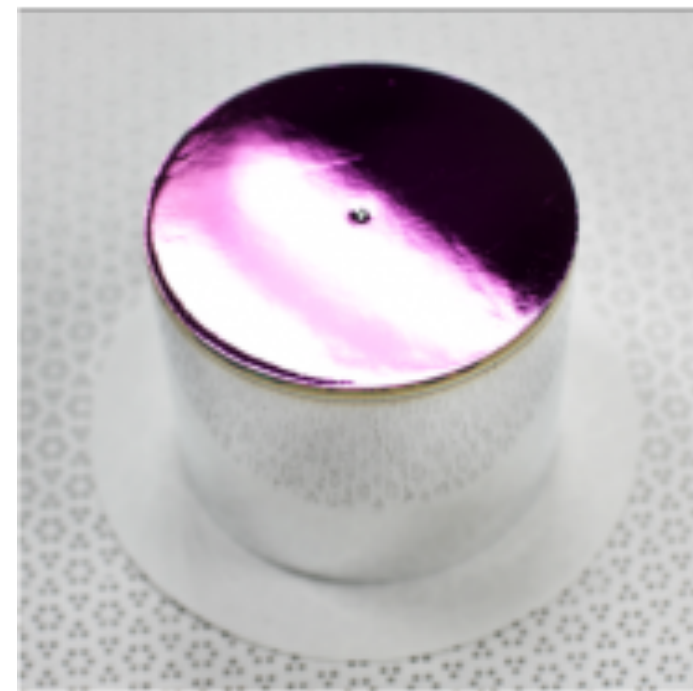
7-String Array

Cryostat



More ^{76}Ge : MAJORANA Demonstrator at Sanford, SD

- **Modular** design:
 - 2 **independent cryostats** with their own vacuum, electronics, and cooling
 - Each with **7 strings of 4-5 detectors**
- Total **44.1 kg** of germanium detectors
- Assembled in glove box with **purged N_2 environment**



PPC HPGe Detector

Low-Mass Mount

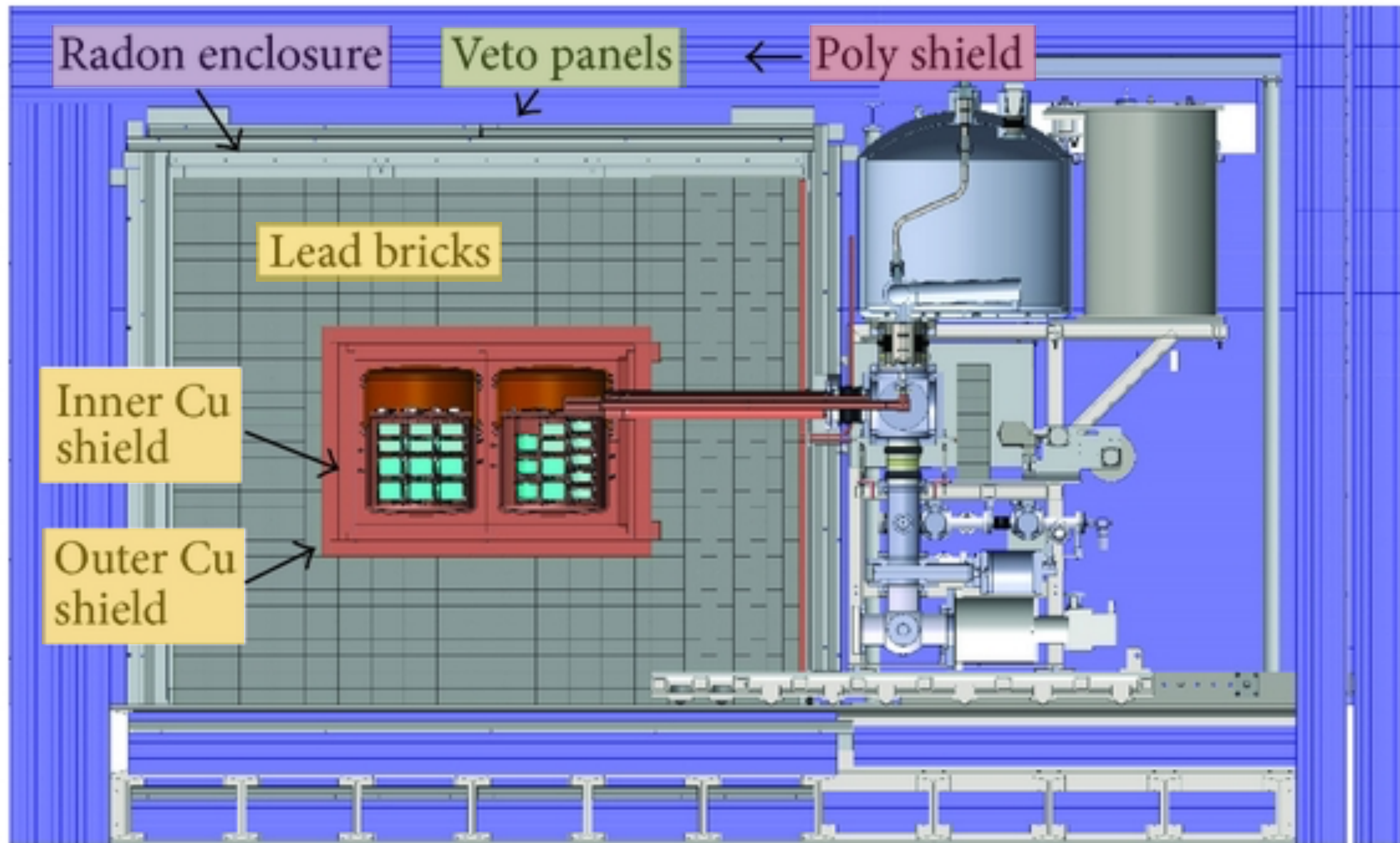
String

7-String Array

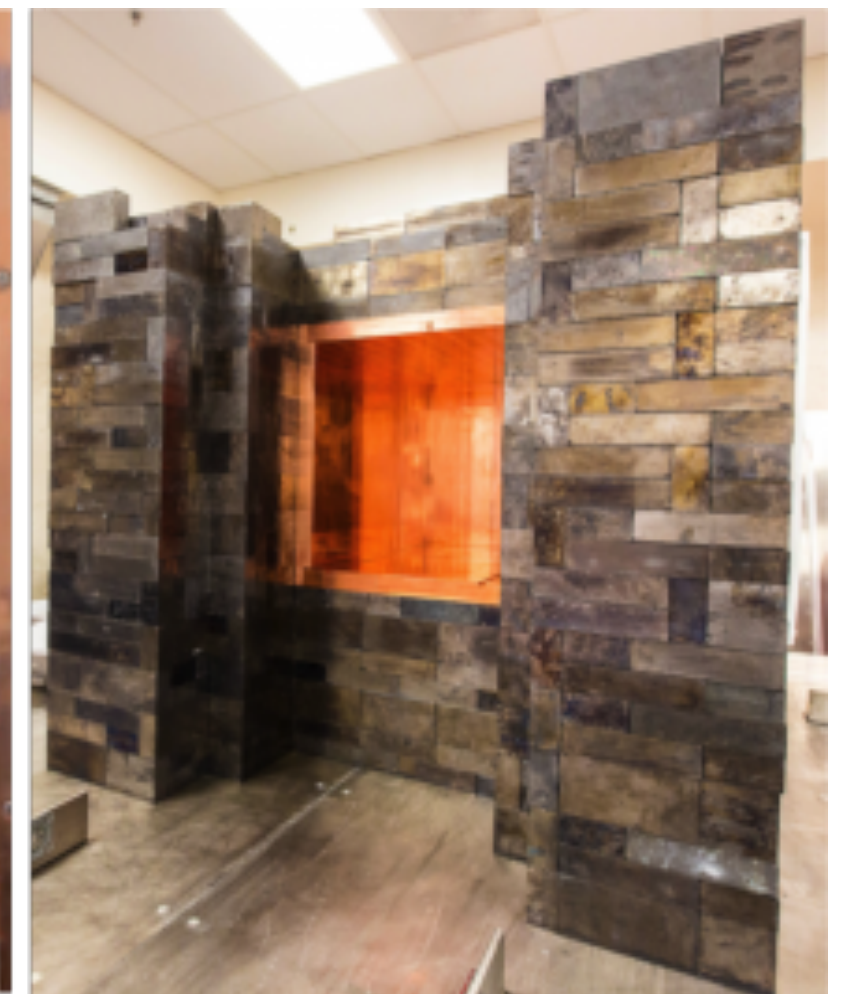
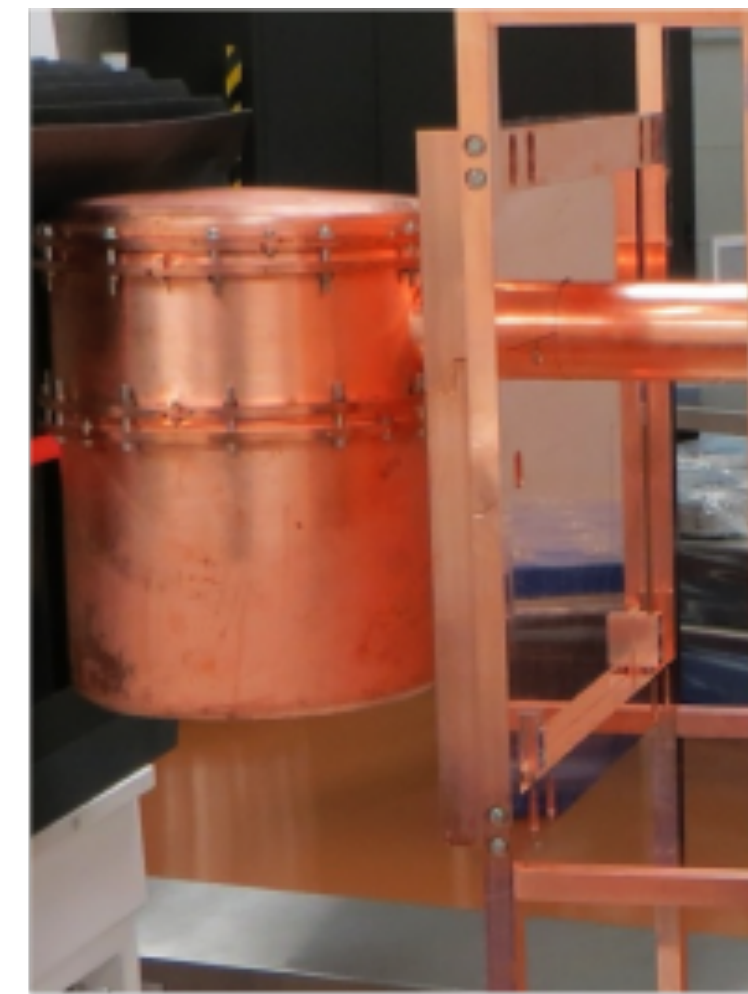
Cryostat



More ^{76}Ge : MAJORANA Demonstrator at Sanford, SD



- External **gammas** blocked by **copper and lead** shielding
- **Neutron** background reduced by borated **poly shield**
- Active **muon** veto panels made of plastic **scintillator**
- Ultra-low **radon** N_2 **purge** gas



PPC HPGe
Detector

Low-Mass
Mount

String

7-String Array

Cryostat

Shield

GERDA + MAJORANA = LEGEND



- Liquid argon veto
- Low-A shield



New inverted coaxial point-contact HPGe detectors



- Low-noise electronics
- Radio-pure parts (mounts, cables etc)

Underground argon

Reduce passive materials

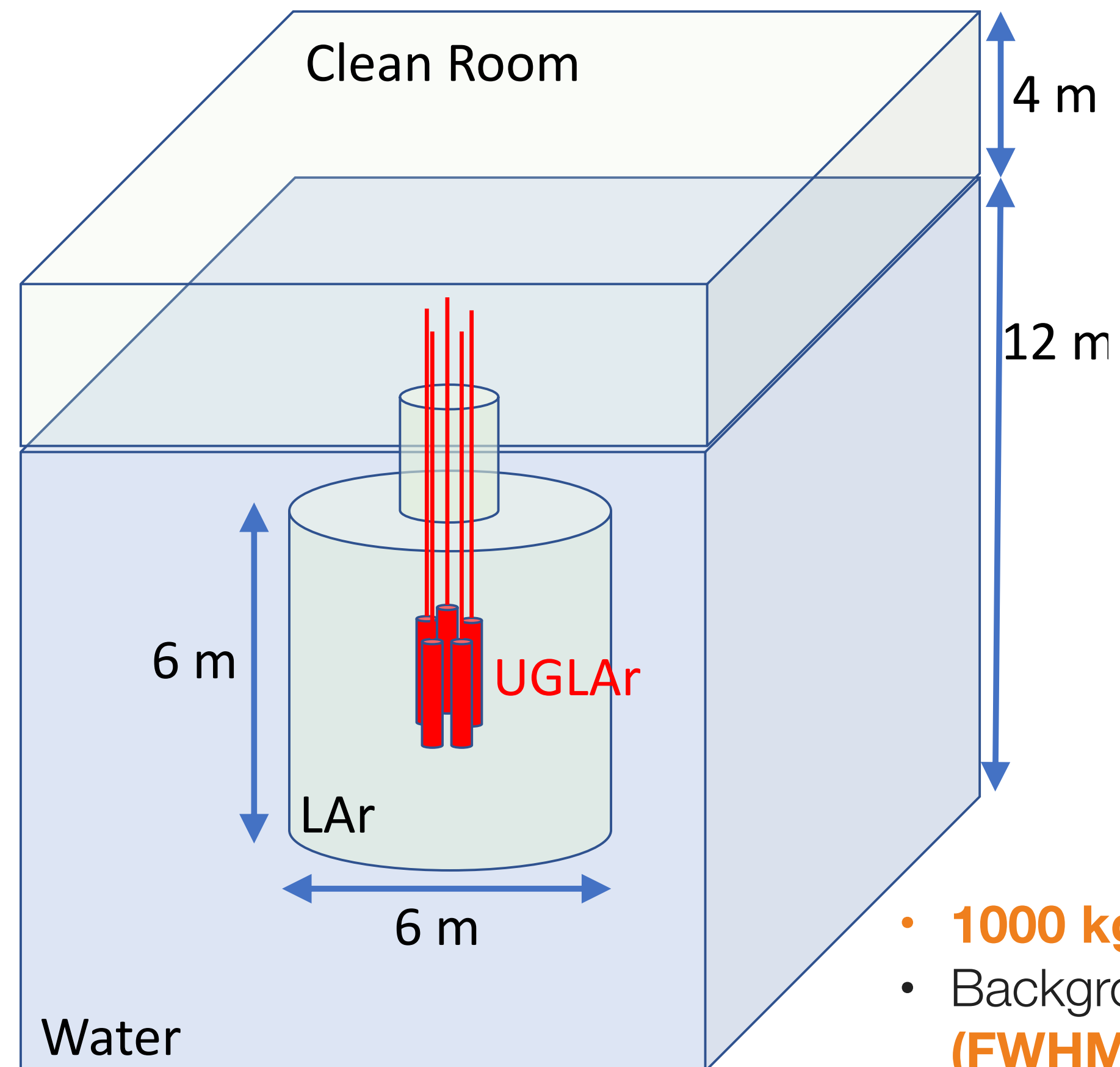


LEGEND-200

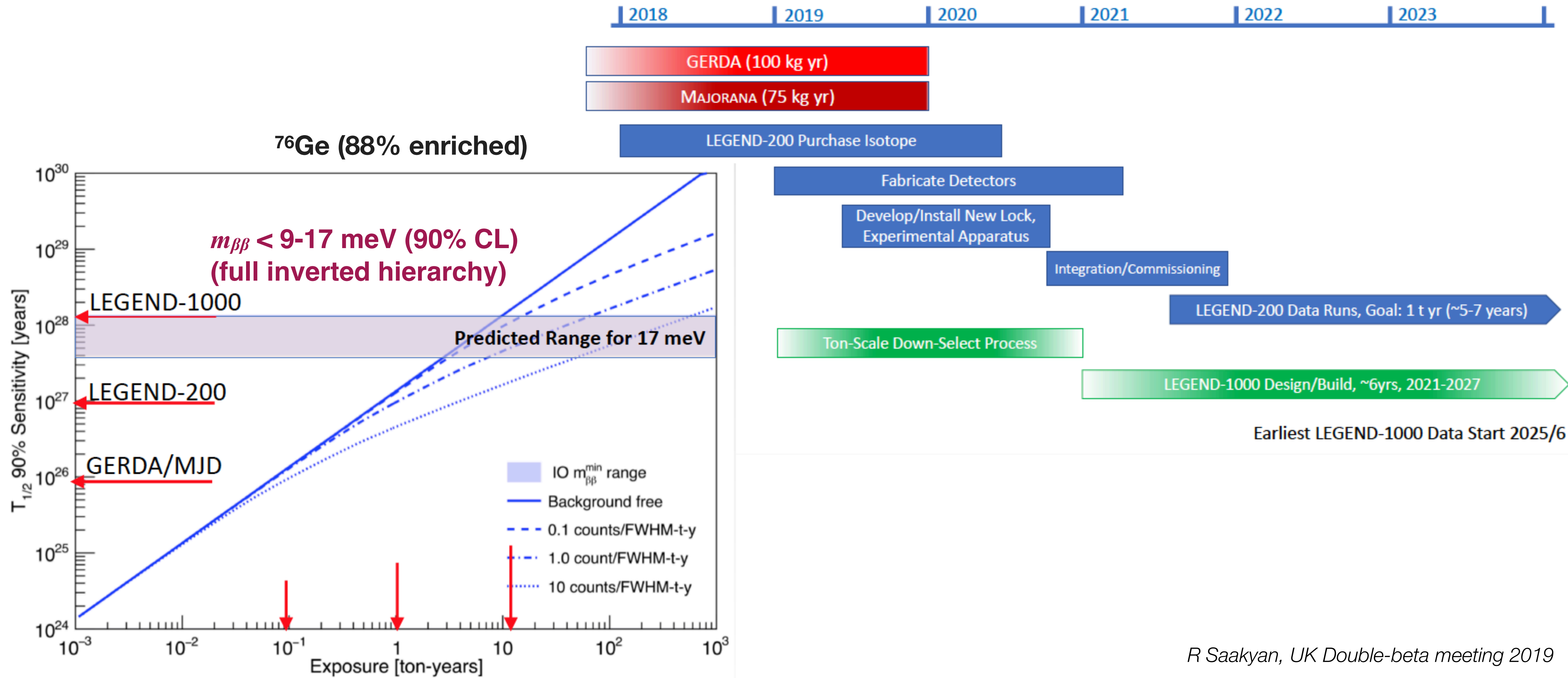


- **200 kg** upgrade of GERDA at LNGS
- Background goal **0.6 cts / (FWHM t yr)**
- Data start ~2021

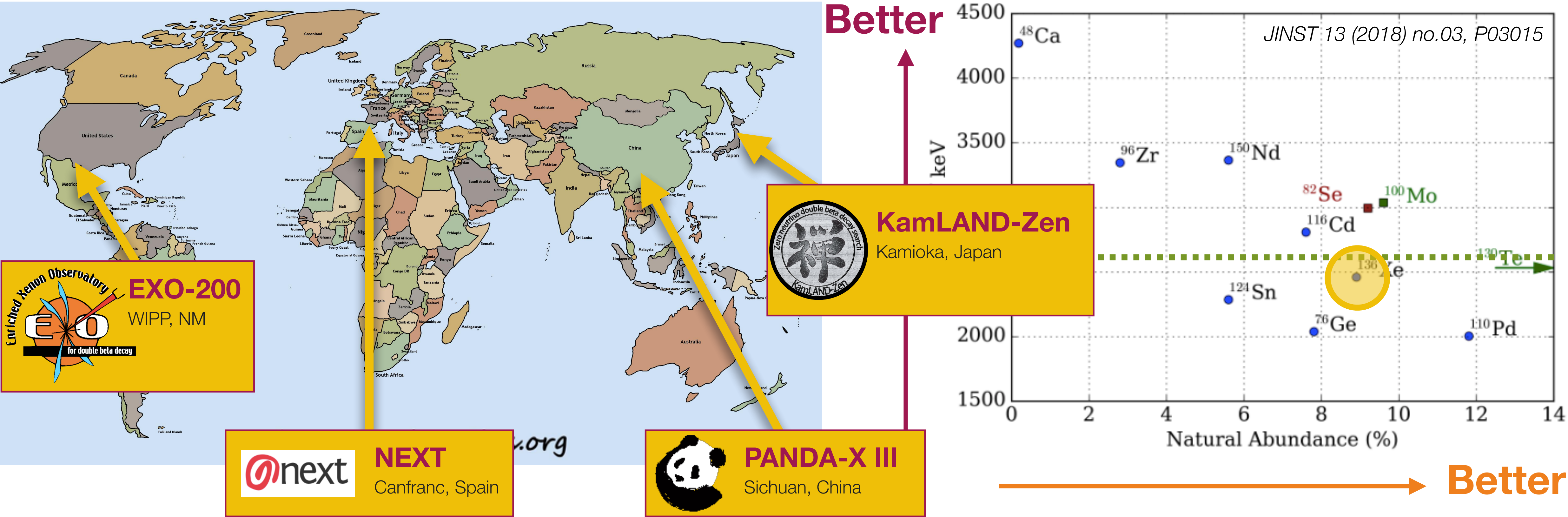
LEGEND-1000



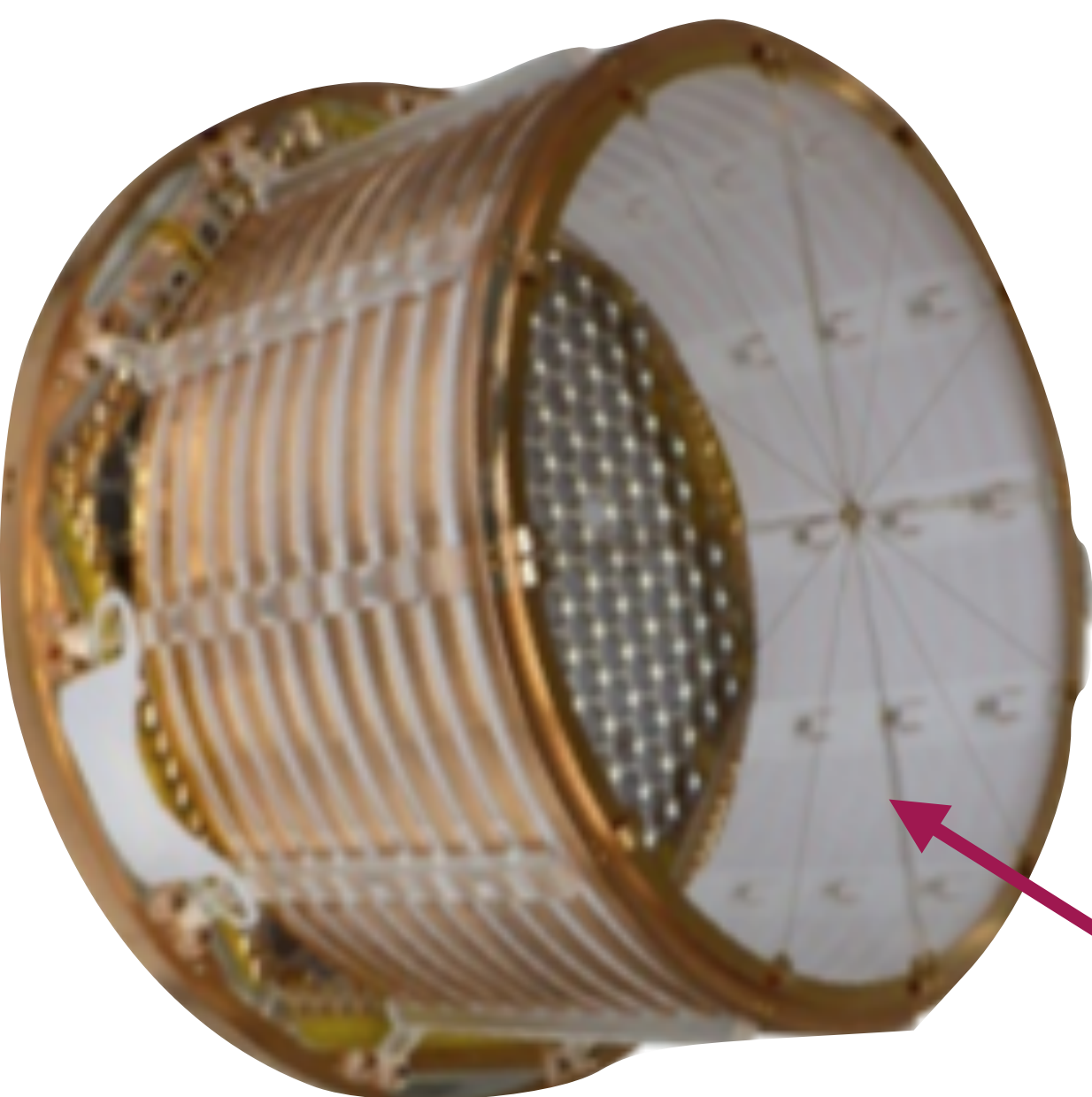
- **1000 kg**
- Background target **0.1 / (FWHM t yr)**
- Location not yet chosen



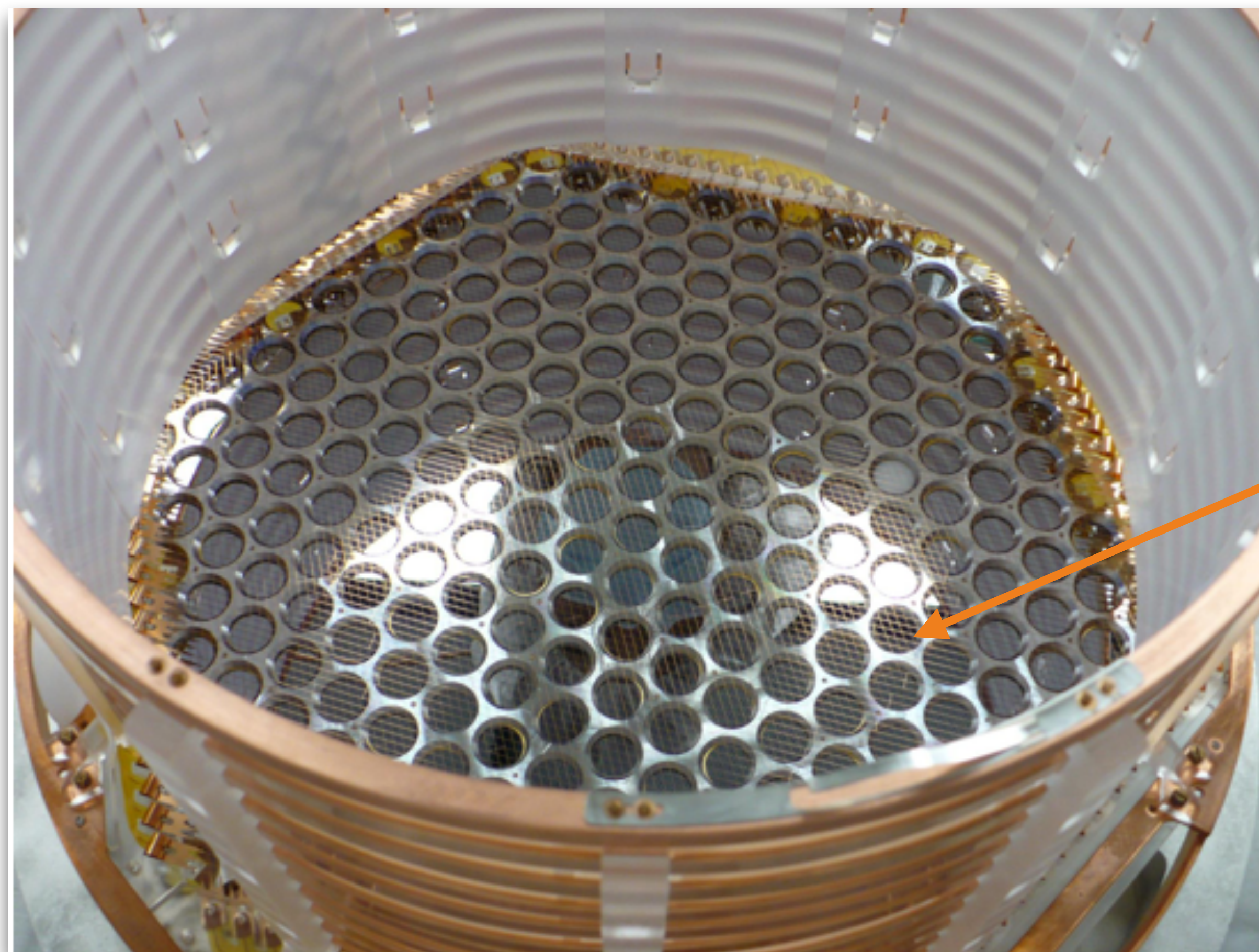
R Saakyan, UK Double-beta meeting 2019



- Large detectors : **hundreds of kg** of isotope
- KamLAND-Zen has current best $0\nu\beta\beta$ half-life / $m_{\beta\beta}$ mass limit ($\langle m_{\beta\beta} \rangle < 61-165 \text{ meV}$)
- Future detectors - **nEXO, KamLAND2 Zen**



Central cathode held at high negative voltage (-8 to-12 kV)

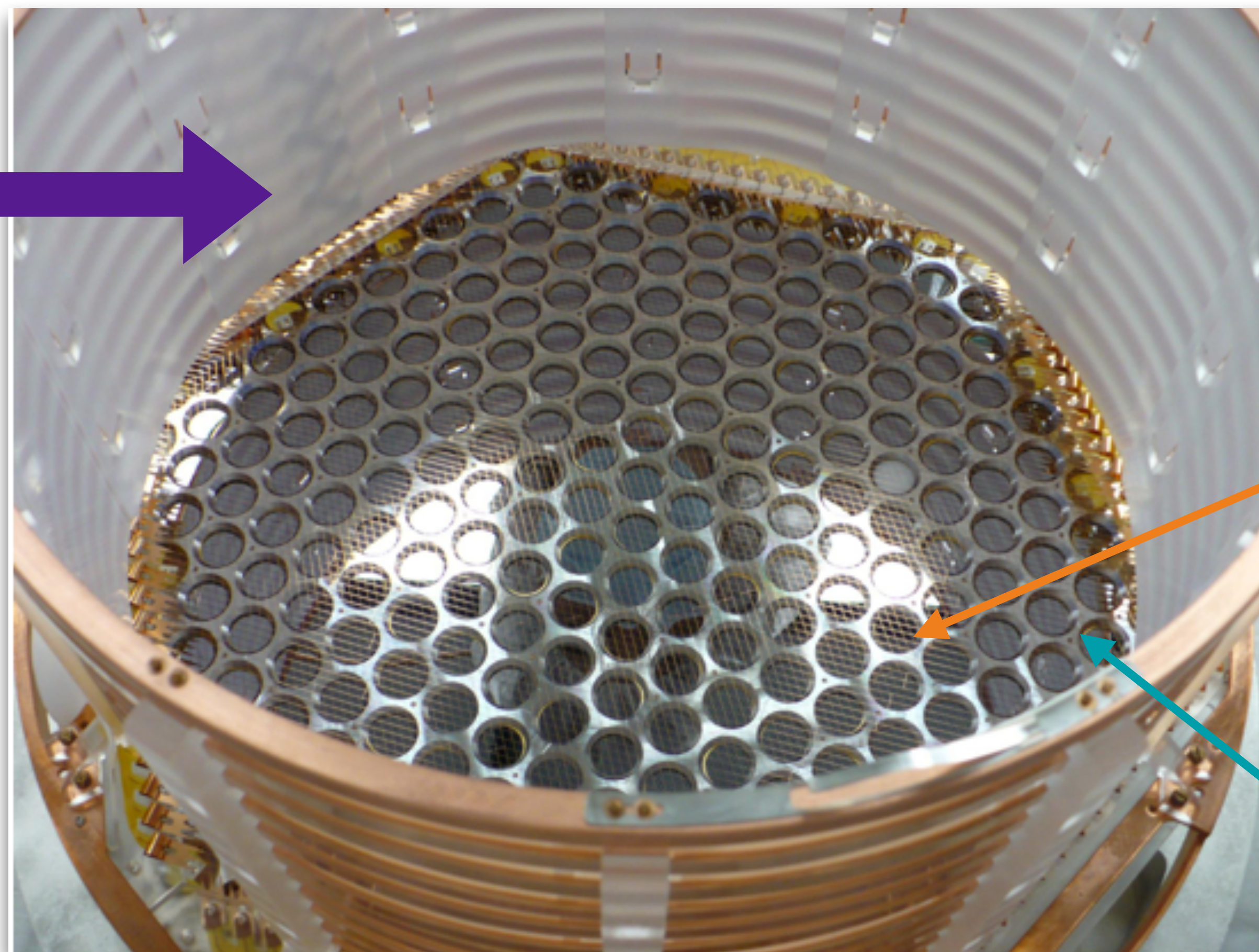


Charge collection grid

← 40cm →

Liquid xenon (80.6% ^{136}Xe) has 3 roles

- TPC ionizer
- $\beta\beta$ emitter
- Scintillator

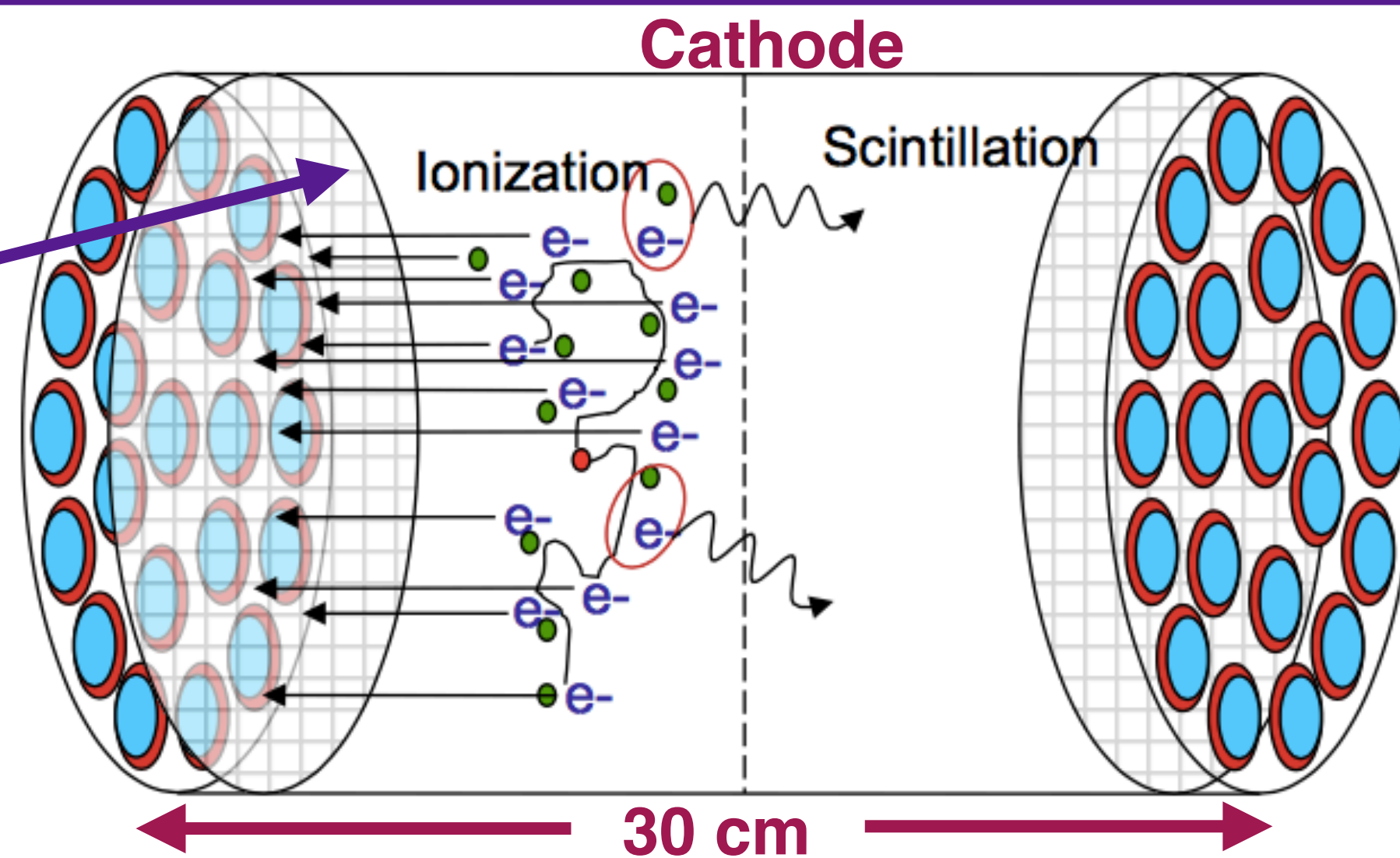


Charge collection grid

Avalanche photodiodes collect 175nm scintillation light

40cm

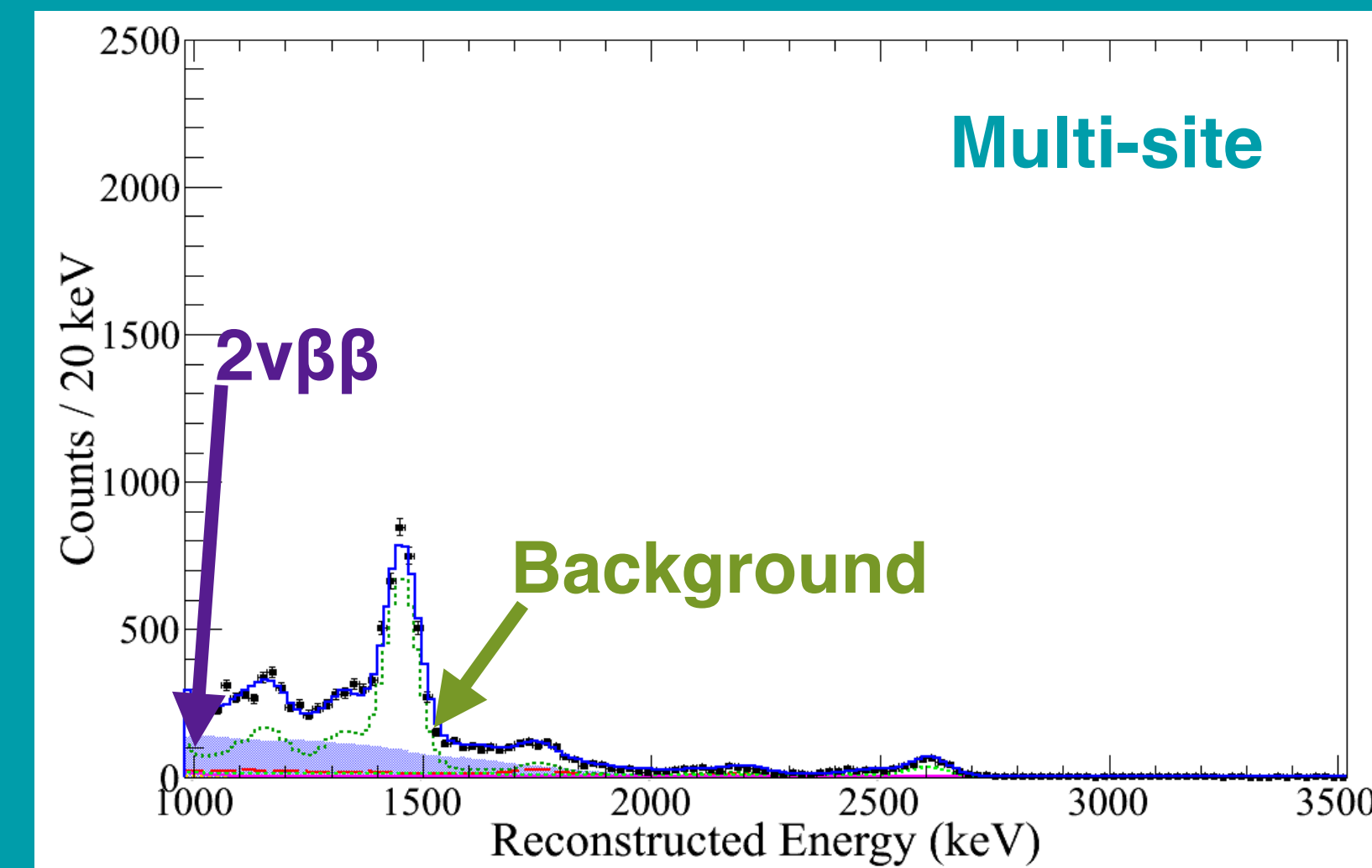
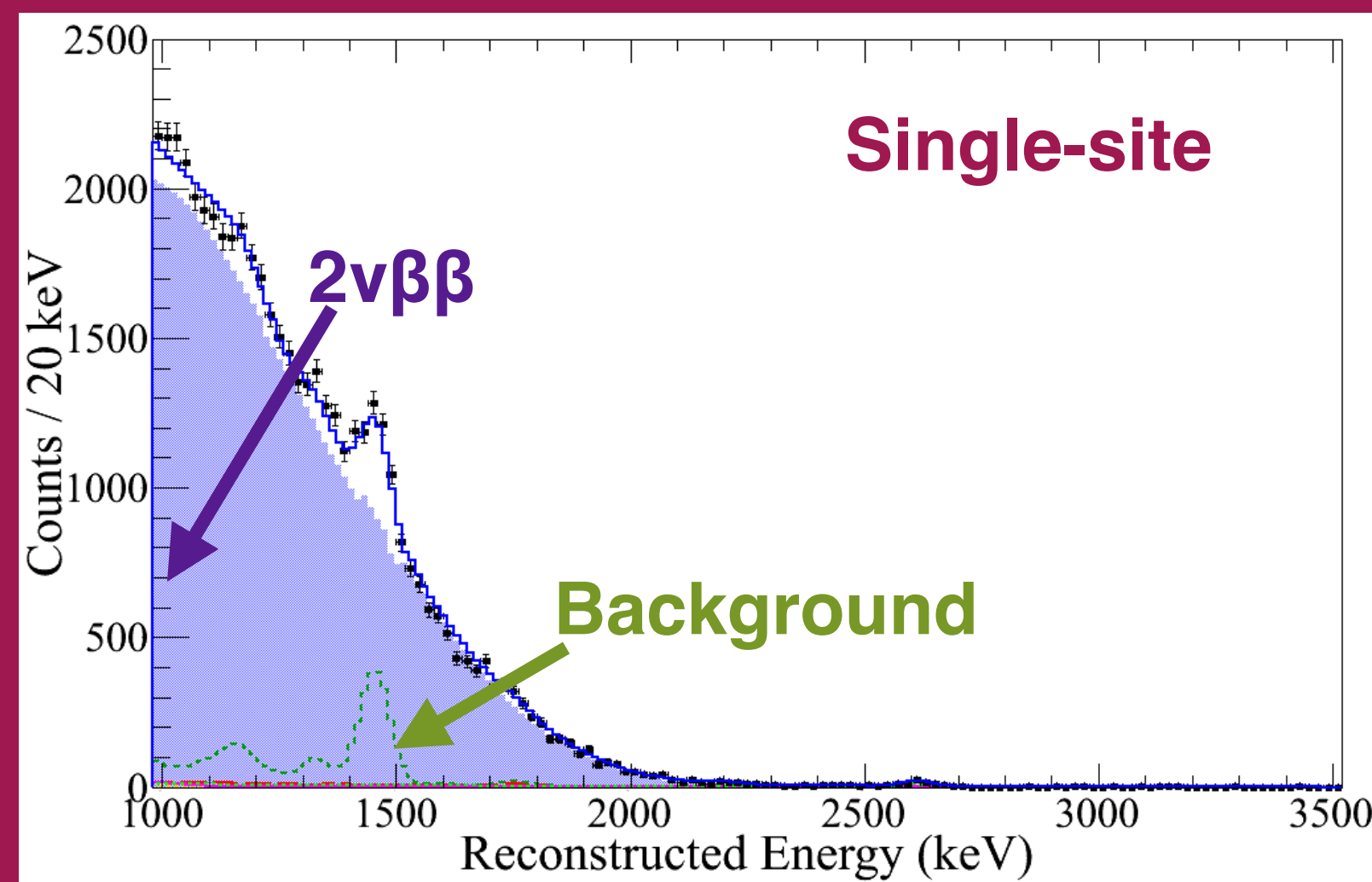
L Yang, Neutrino 2016



TPC wire signals give x-y position

Time delay between scintillation light and wire signals gives longitudinal position

3-D reconstruction distinguishes between **single-site ($\beta\beta$)** and **multi-site (γ background)** events



Single-site cut

Energy cut
($\sigma_E / E = 1.23\%$)

Boosted decision tree

Phys. Rev. Lett. 120, 072701 (2018)

Single-site cut

Energy cut
($\sigma_E / E = 1.23\%$)

Boosted decision tree

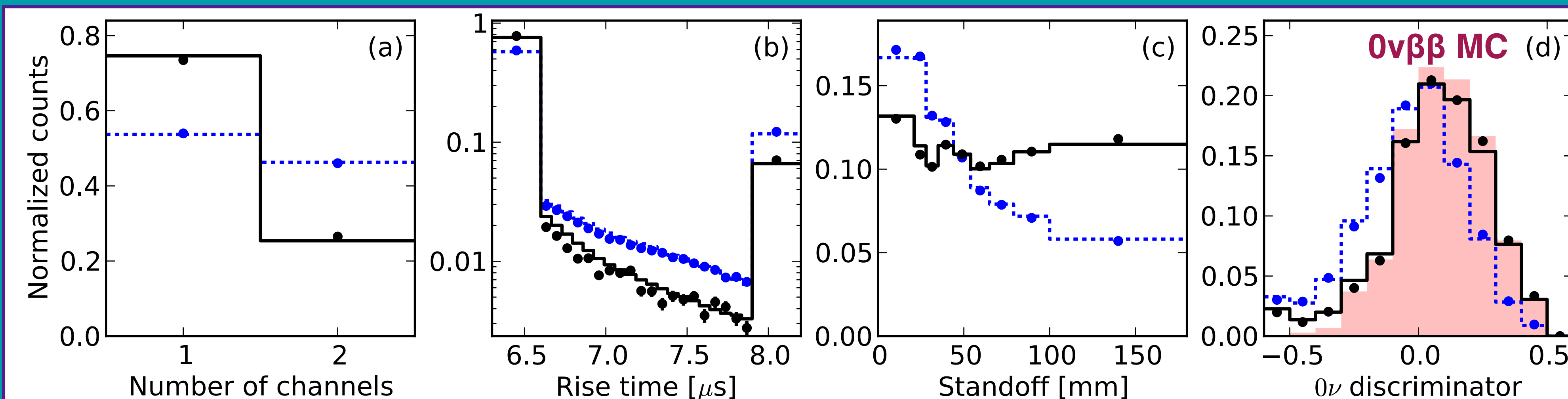
Time between 5 and 95% of charge being collected

$2\nu\beta\beta$

— simulation
● background-subtracted data

Calibration source
(^{226}Ra)

⋯ simulation
● data



Backgrounds tend to fire neighbouring channels

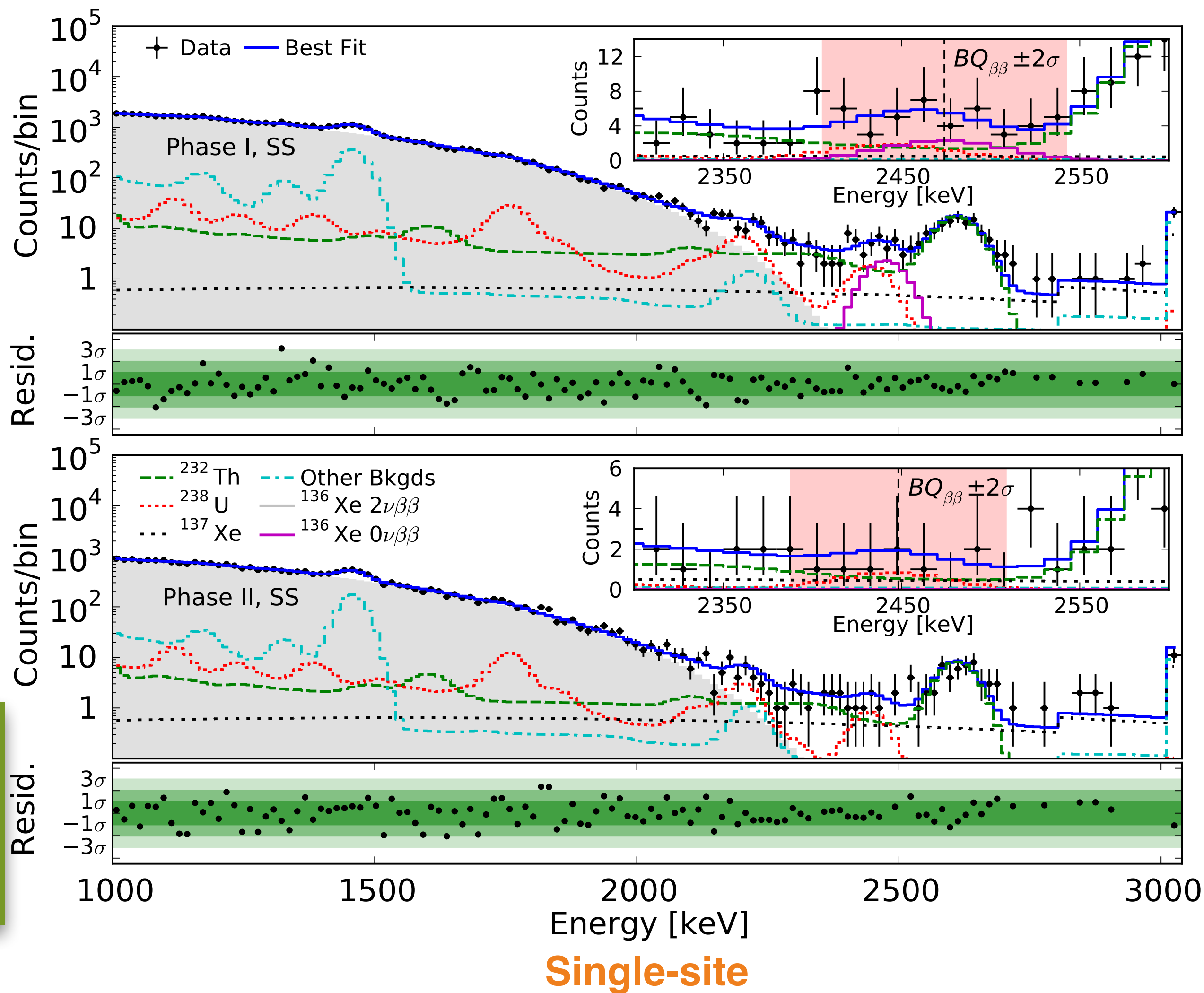
External backgrounds generate clusters near TPC surface

15% sensitivity improvement

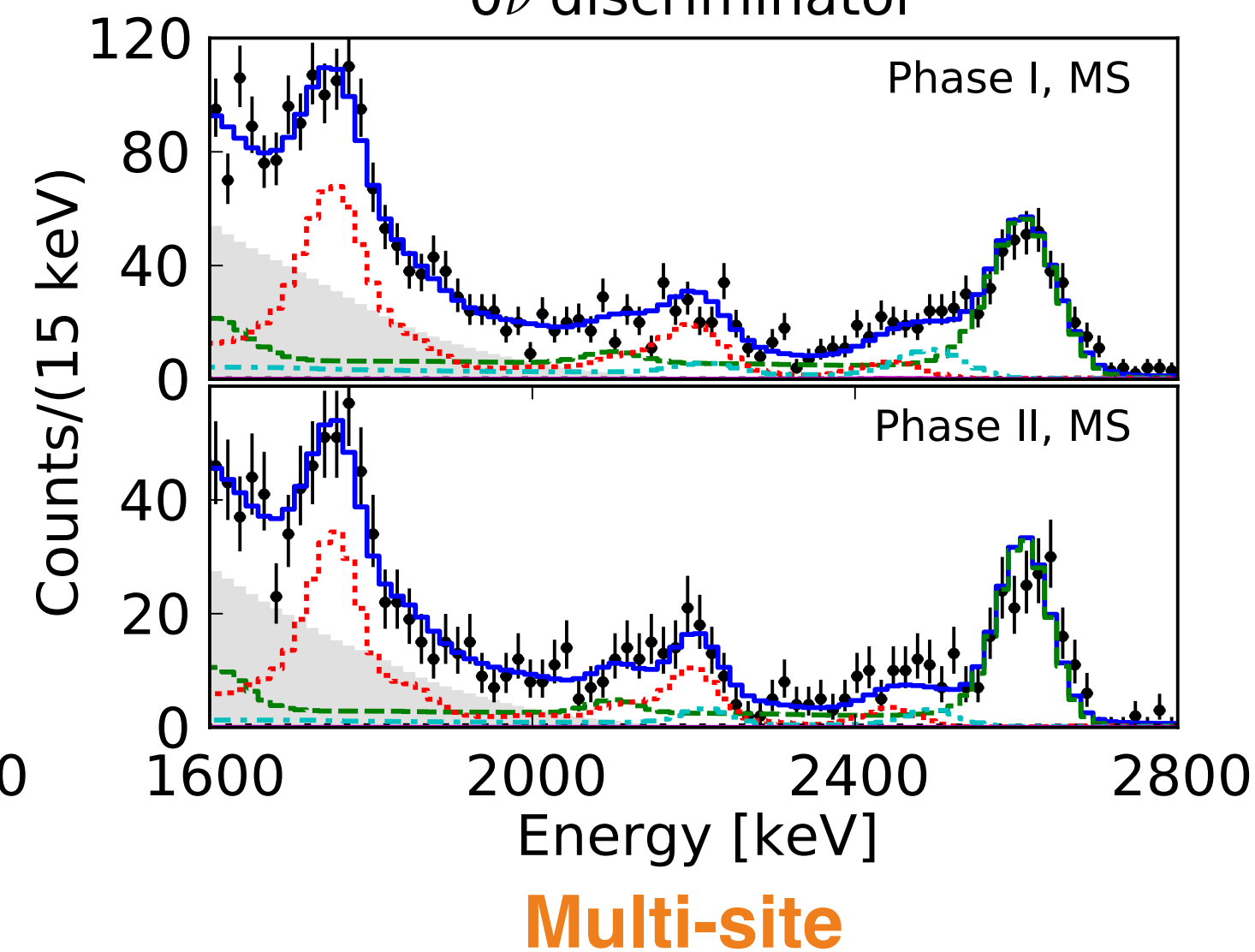
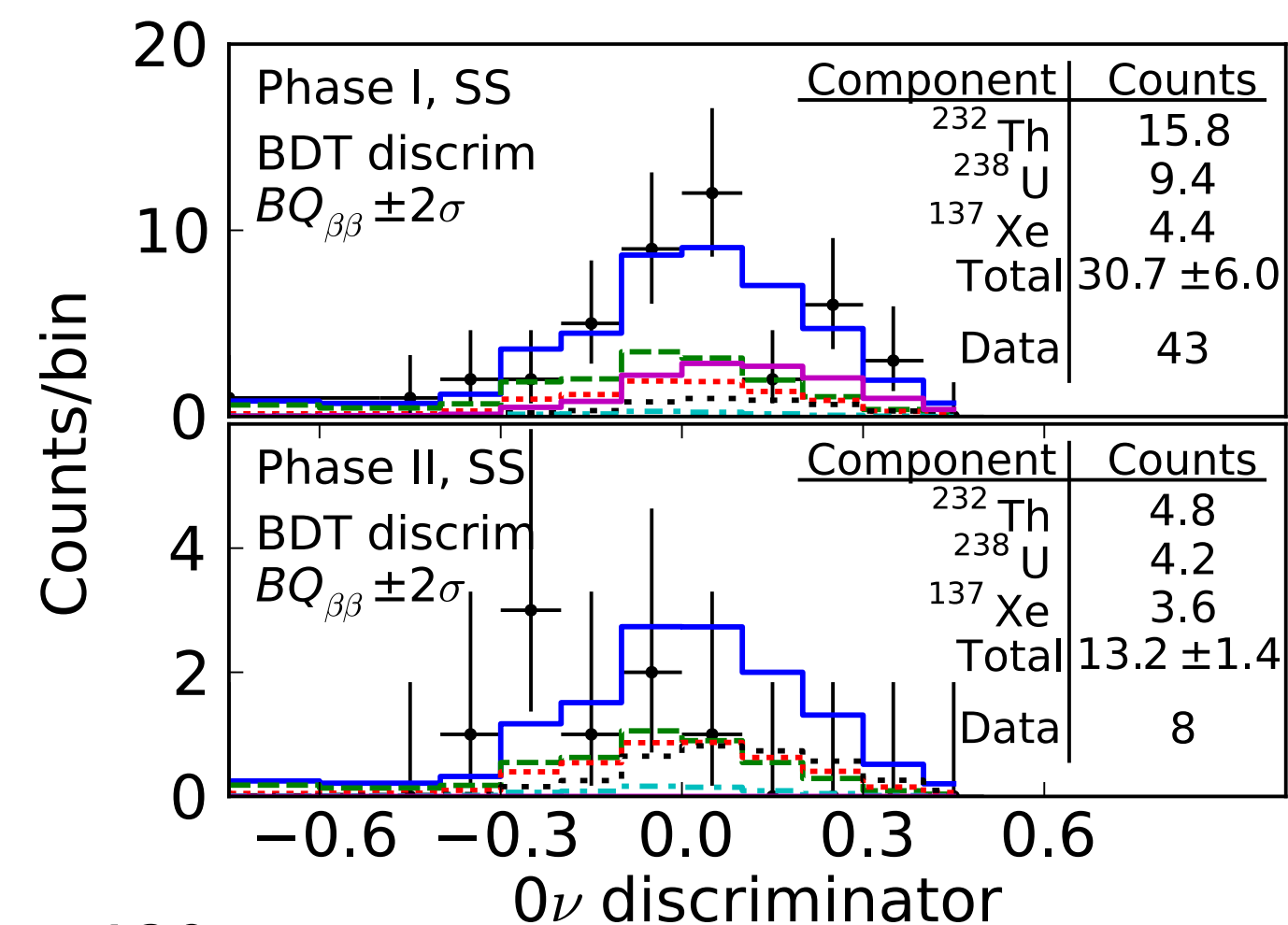
Phys. Rev. Lett. 120, 072701 (2018)

Phys. Rev. Lett. 120, 072701 (2018)

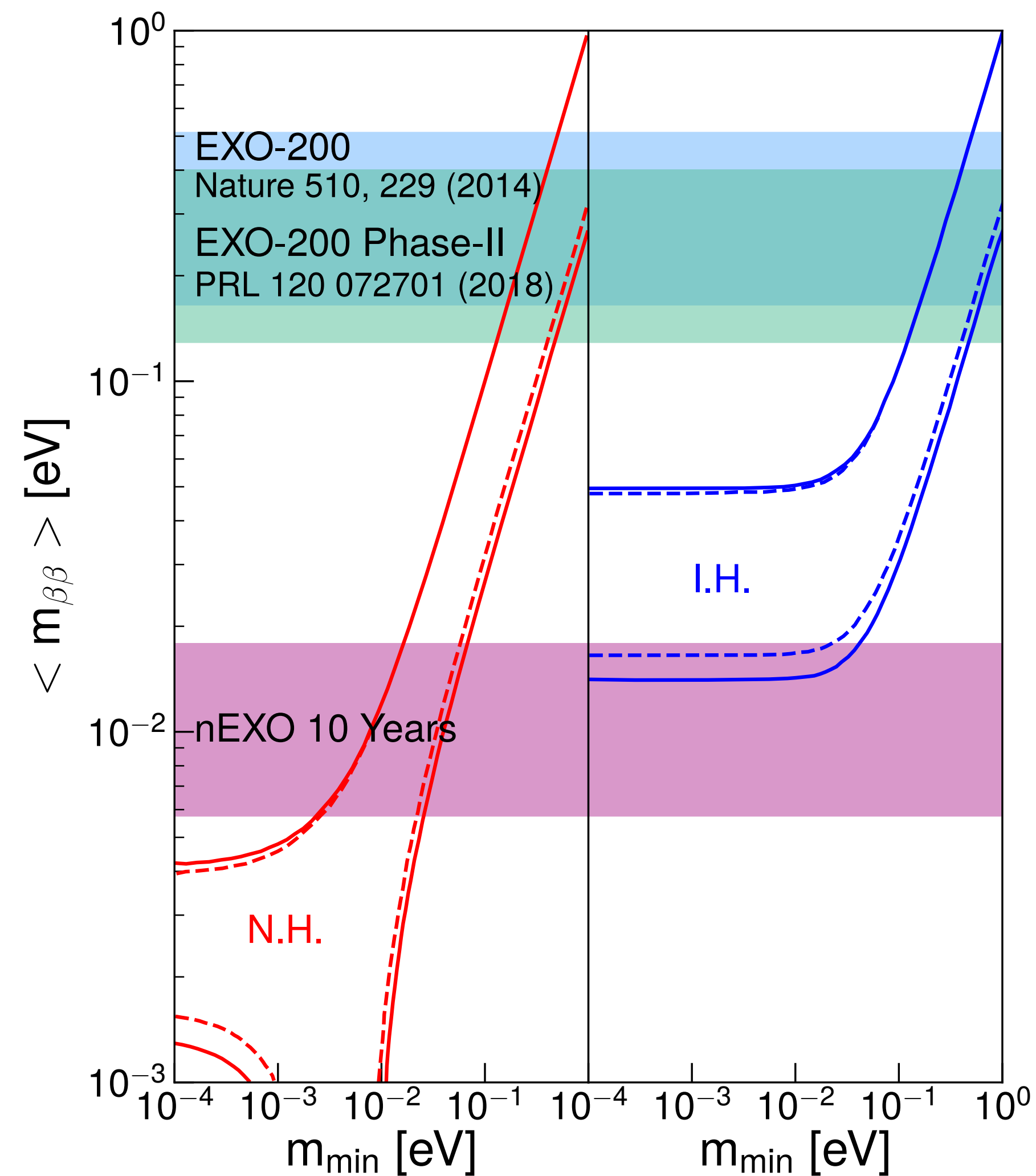
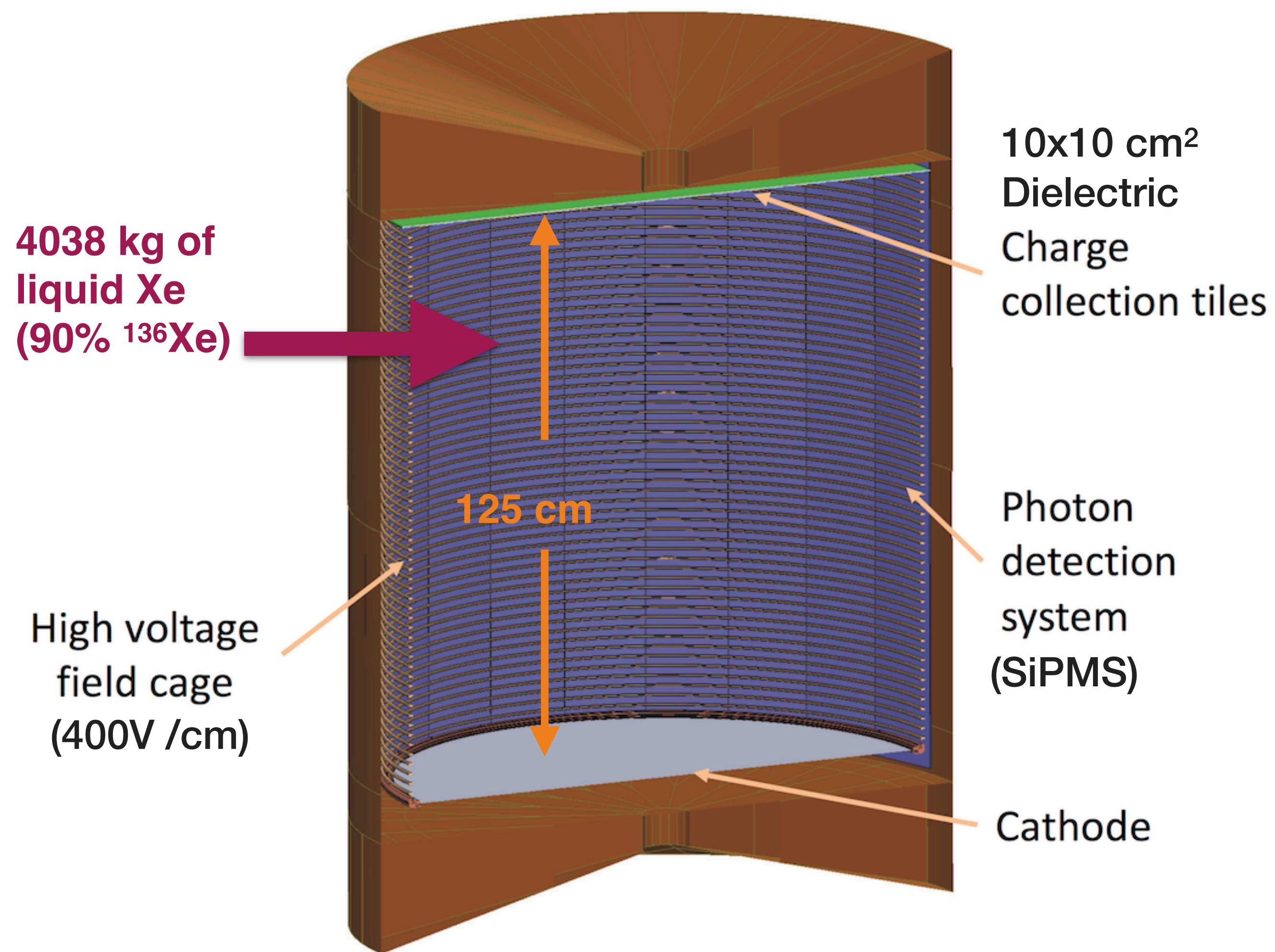
**Phase I vs Phase II
- 2014 accident at WIPP; 2 year
pause and
upgrades**



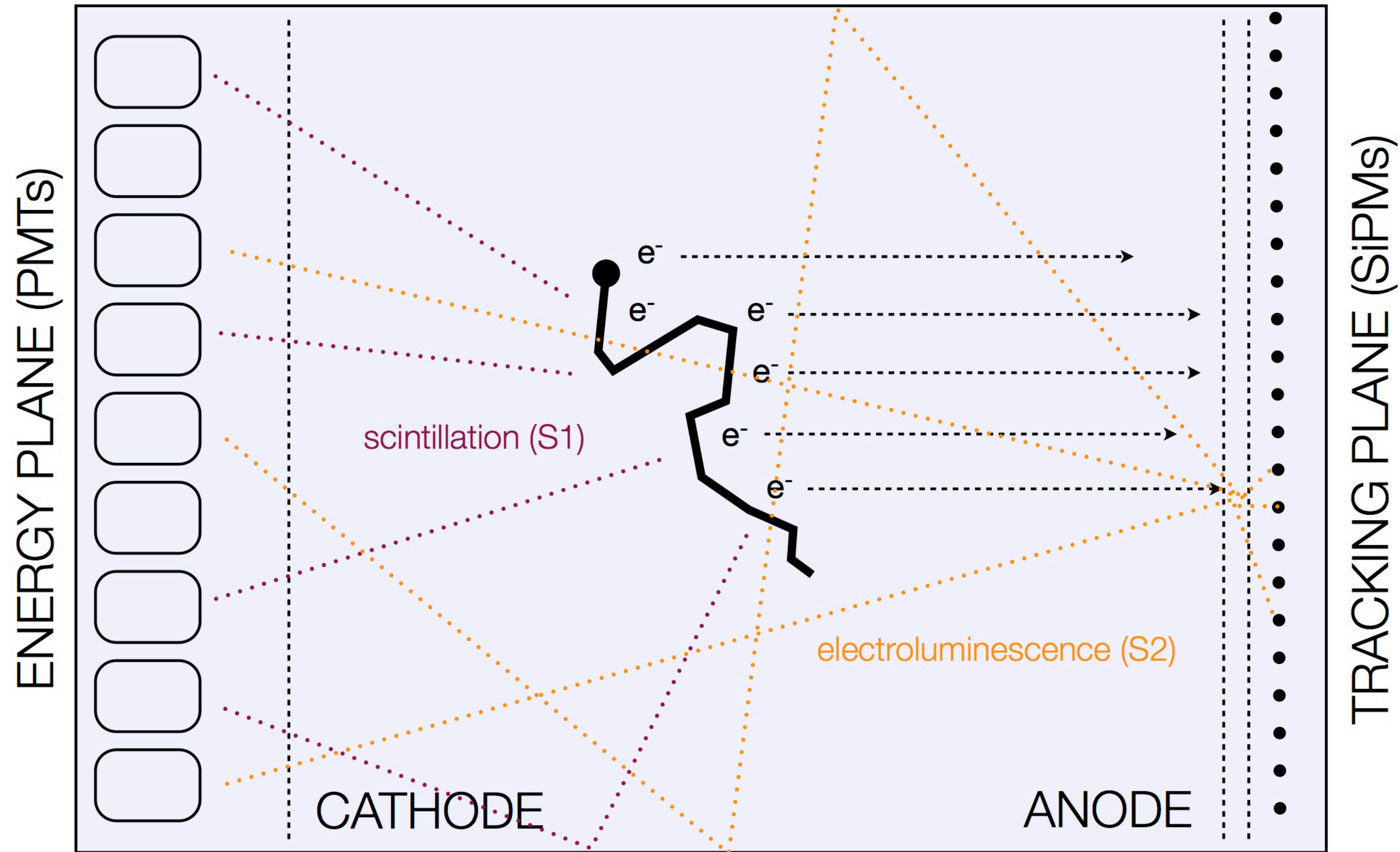
BDT discriminator



$T_{1/2} > 1.8 \times 10^{25}$ years
(90% C.L.)
($\langle m_{\beta\beta} \rangle < 147 - 398$ meV)
177.6 kg.years

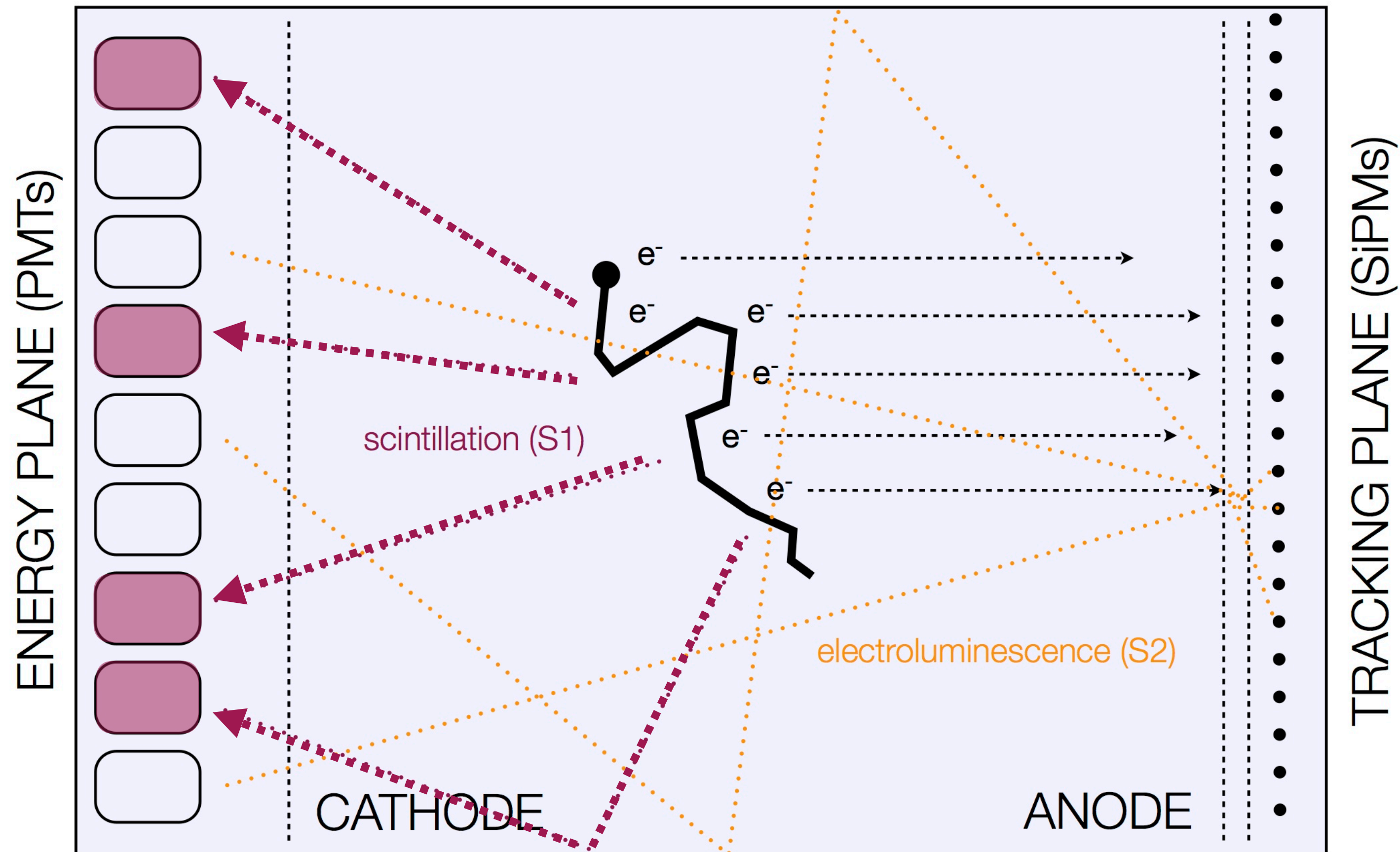


Phys. Rev. C 97, 065503 (2018)



SOFT TPC Separated Optimized Function

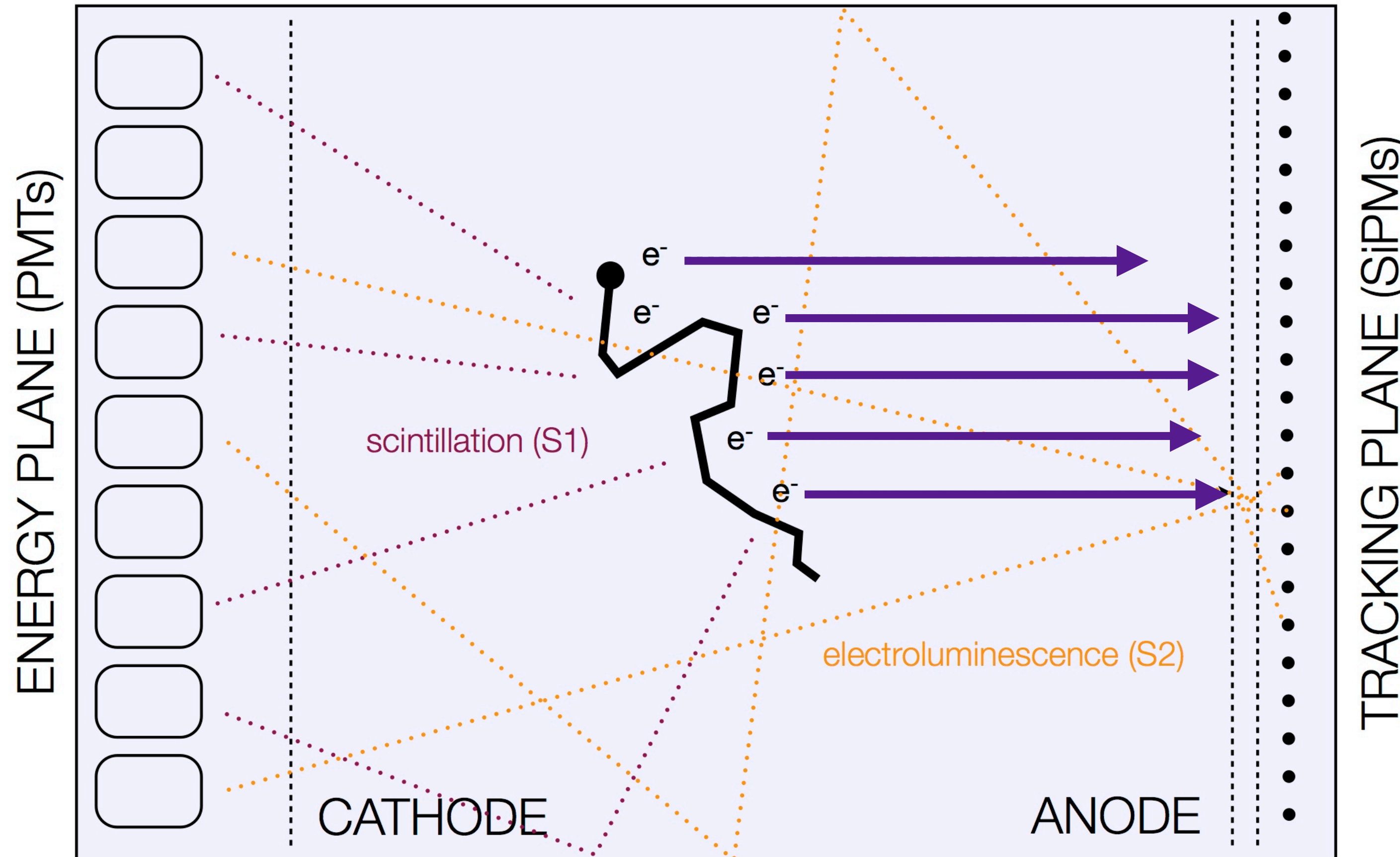
Charged particle travels through TPC



SOFT TPC Separated Optimized Function

Charged particle travels through TPC

Primary scintillation light gives event trigger time

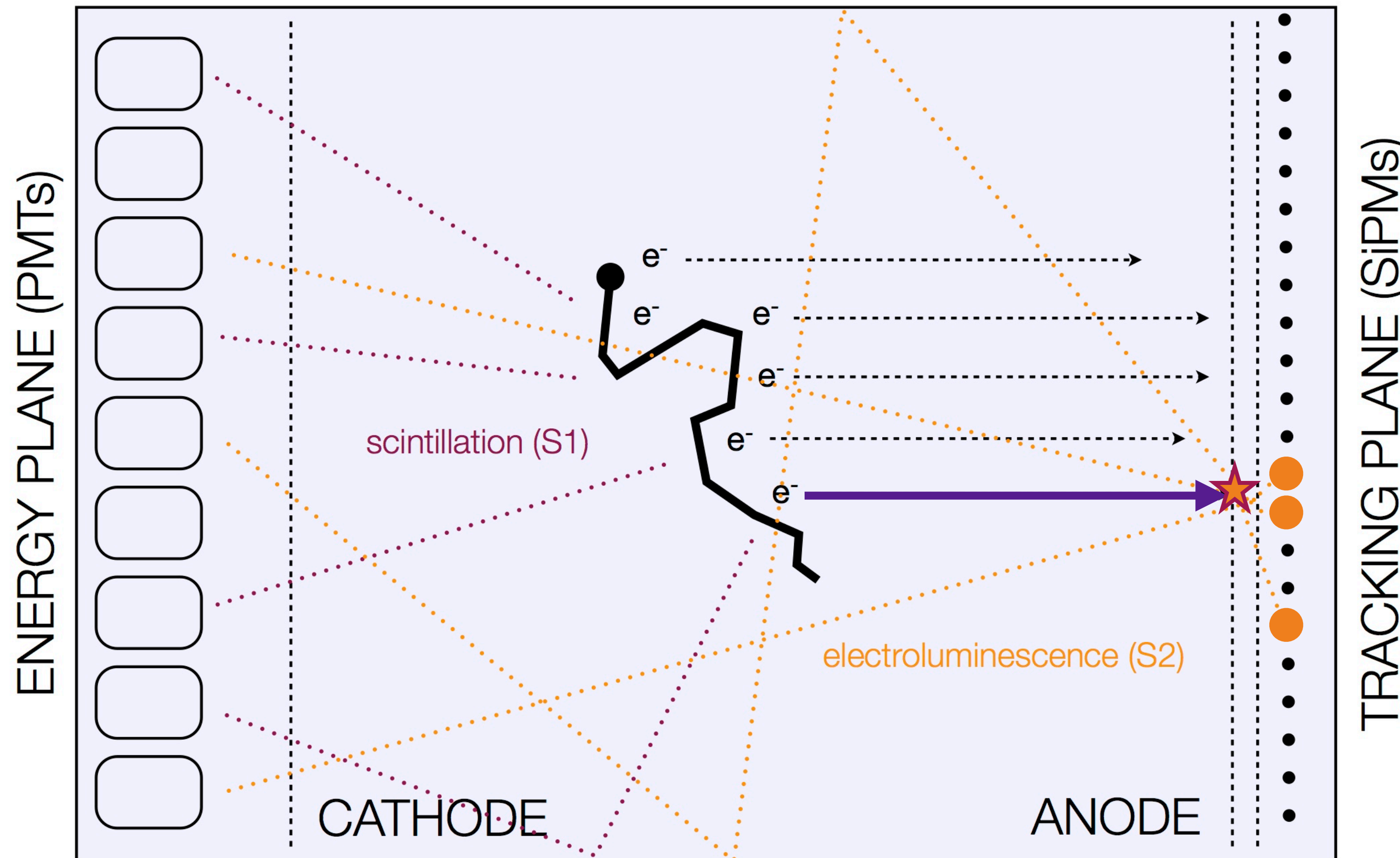


SOFT TPC Separated Optimized Function

Charged particle travels through TPC

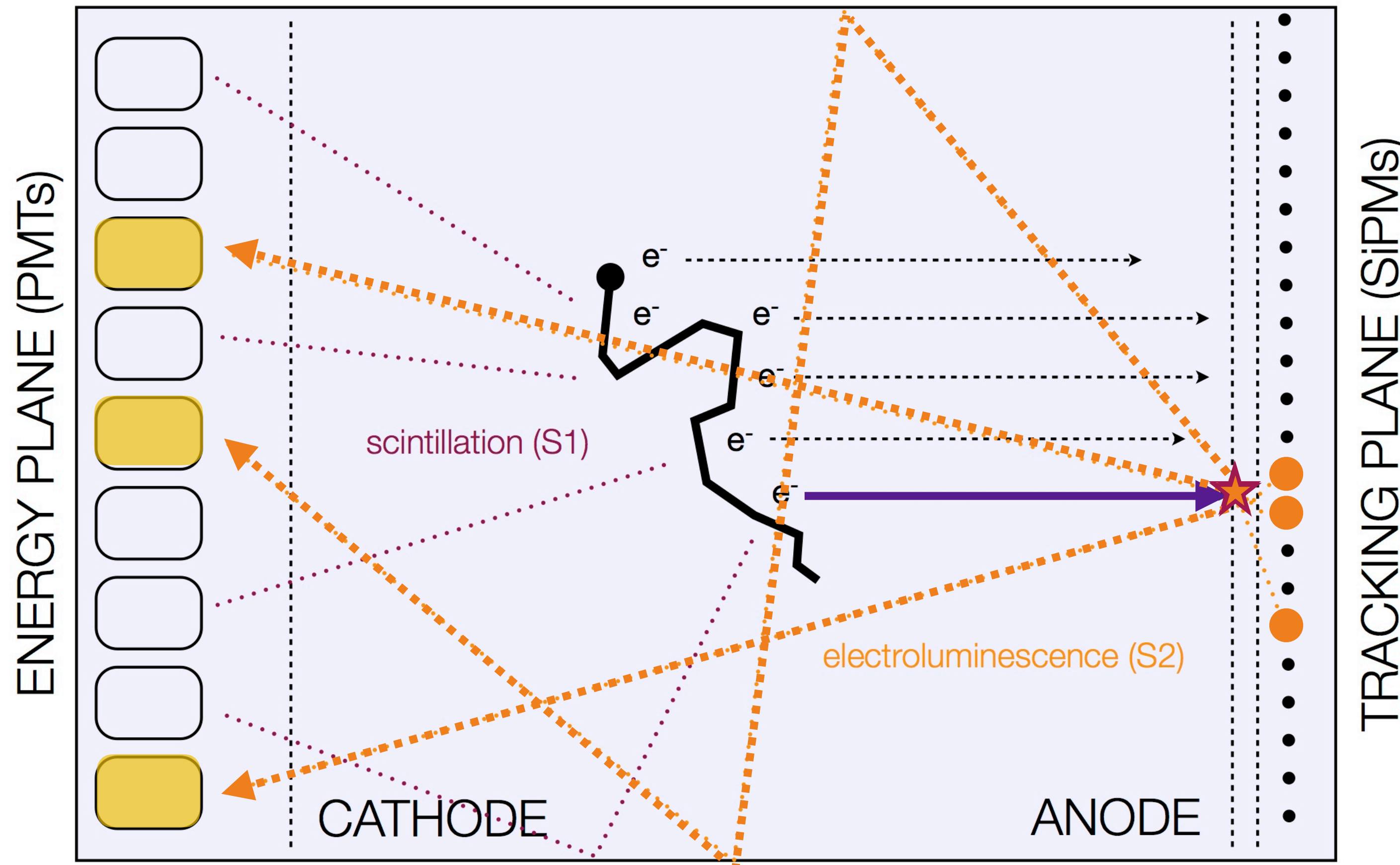
Primary scintillation light gives event trigger time

Ionization electrons drift to anode



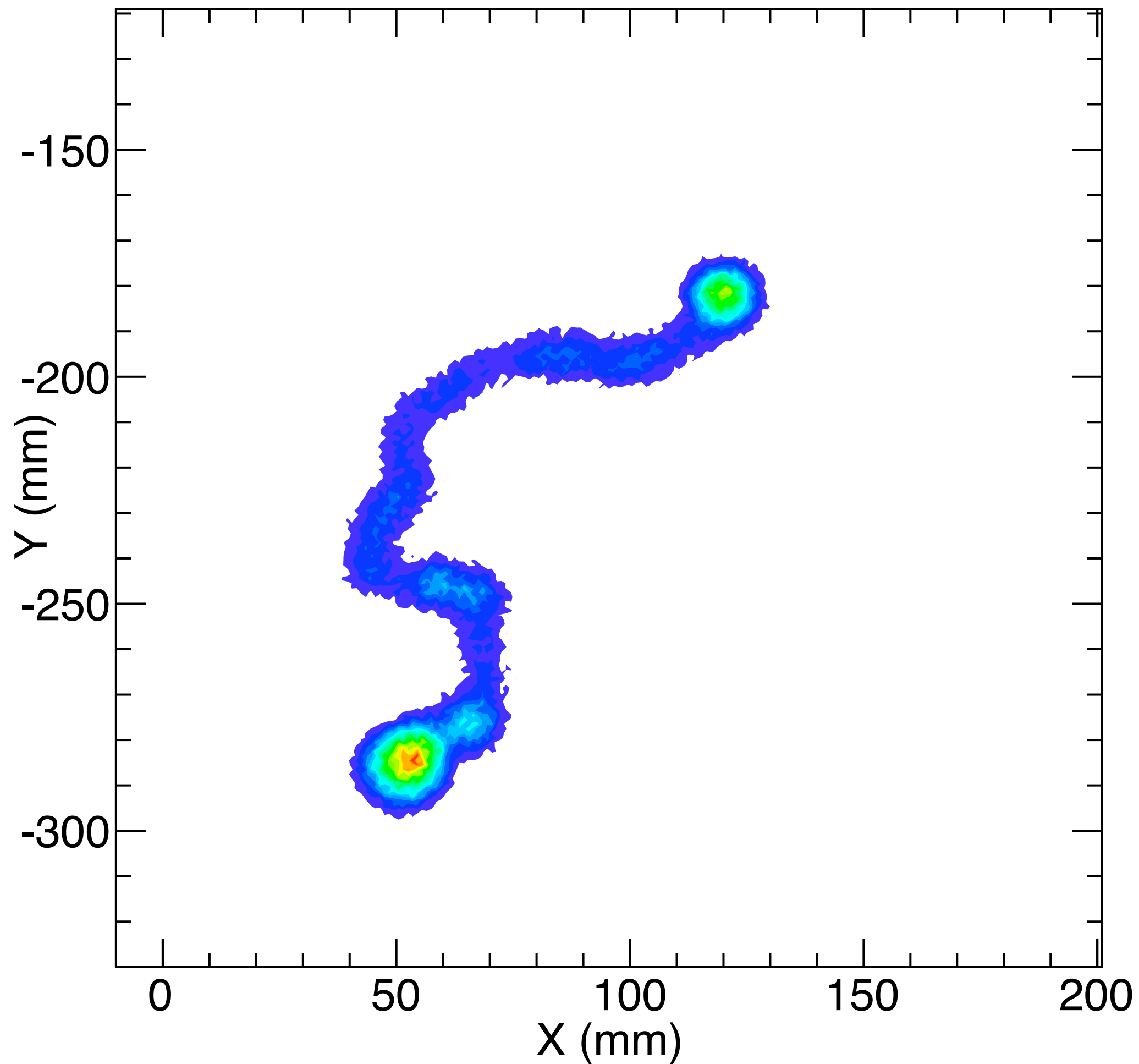
SOFT TPC Separated Optimized Function

- Charged particle travels through TPC**
- Primary scintillation light gives event trigger time**
- Ionization electrons drift to anode**
- Electroluminescent light generated at anode is recorded in SiPMs and used for tracking**



SOFT TPC Separated Optimized Function

- Charged particle travels through TPC**
- Primary scintillation light gives event trigger time**
- Ionization electrons drift to anode**
- Electroluminescent light generated at anode is recorded in SiPMs and used for tracking**
- Electroluminescent light also recorded in PMTs and used for energy measurement**



SOFT TPC Separated Optimized Function

Charged particle travels through TPC

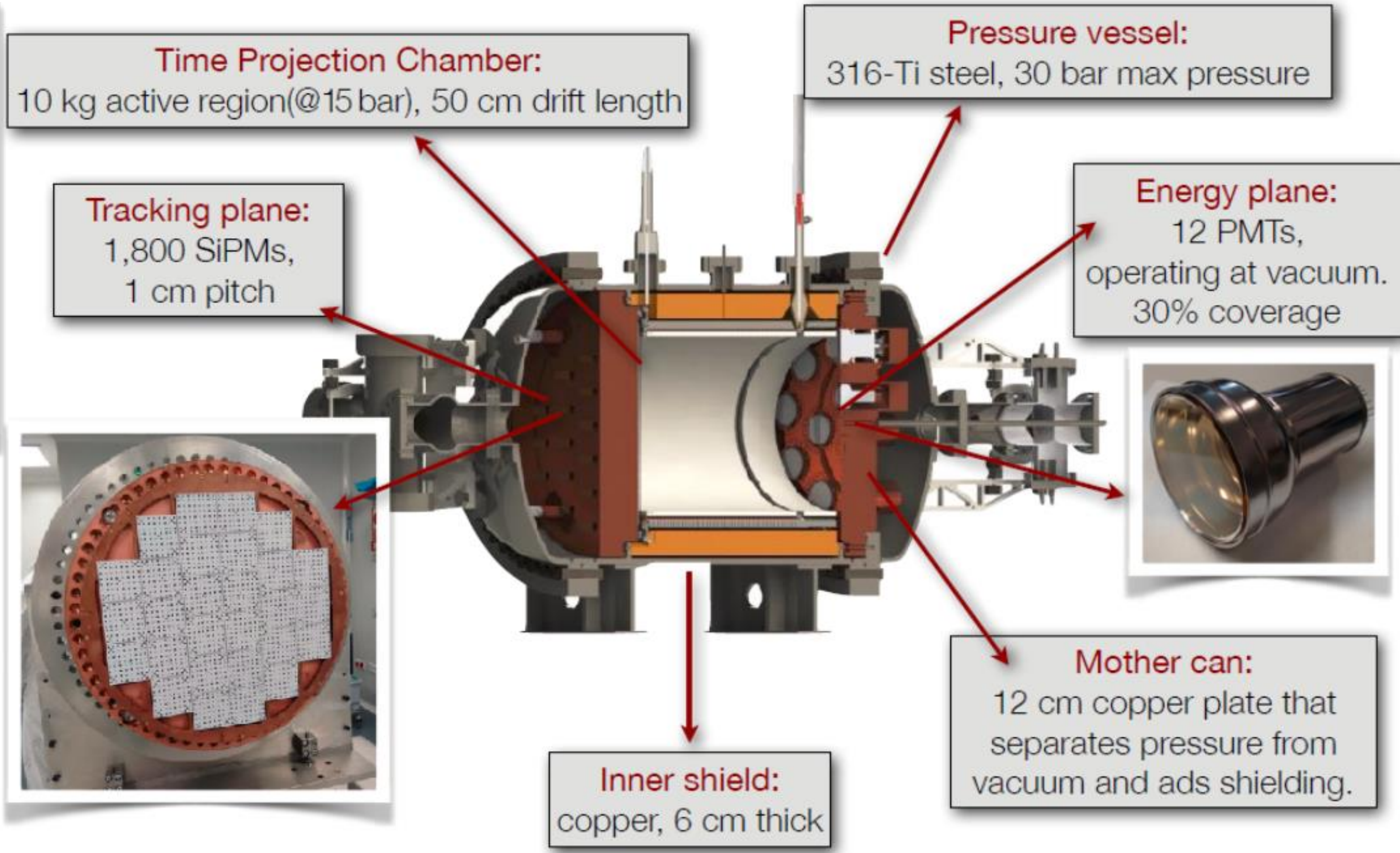
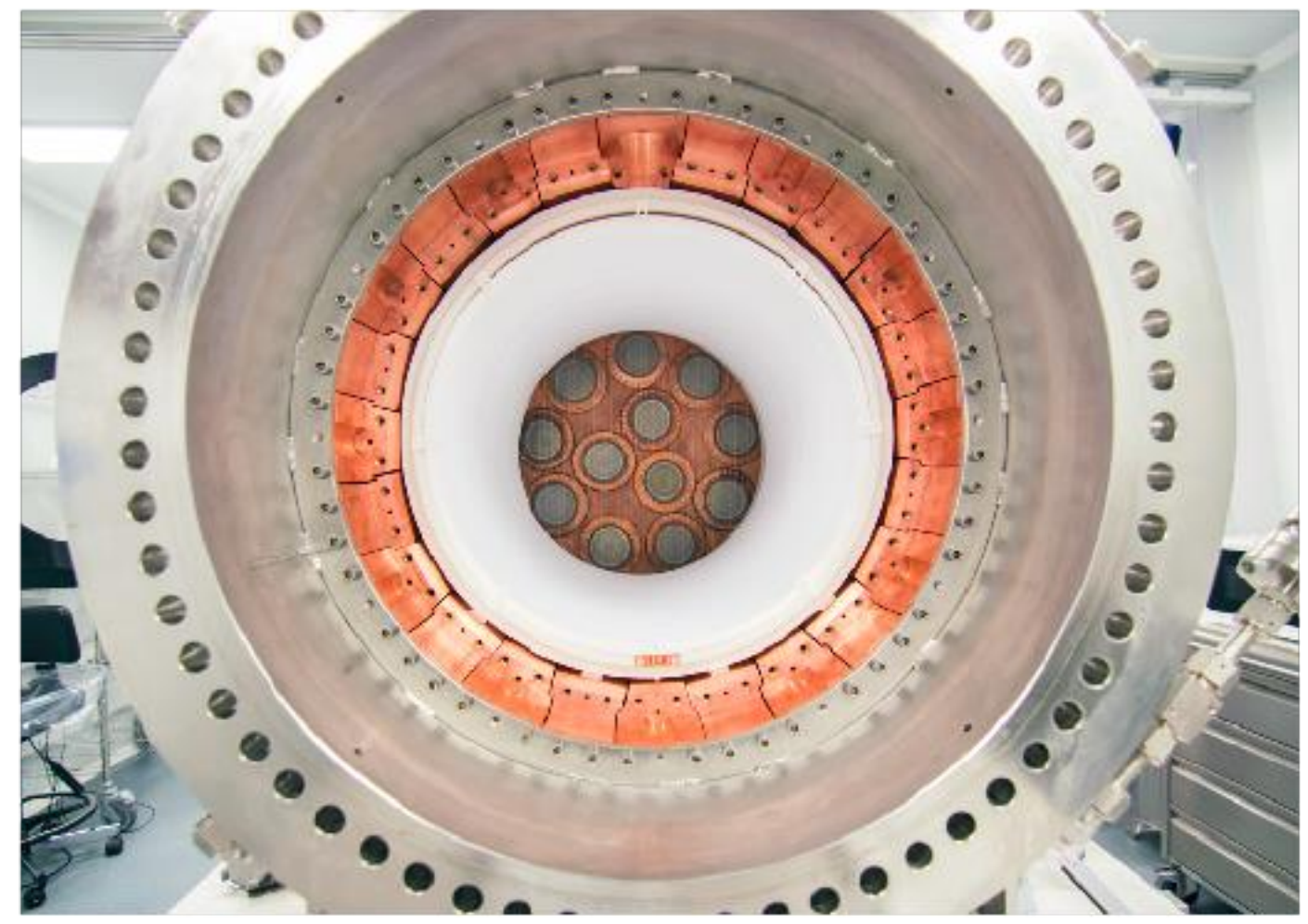
Primary scintillation light gives event trigger time

Ionization electrons drift to anode

Electroluminescent light generated at anode is recorded in SiPMs and used for tracking

Electroluminescent light also recorded in PMTs and used for energy measurement

Each electron generates a constant deposition, followed by a “blob” of energy at the end - count blobs to separate signal from background

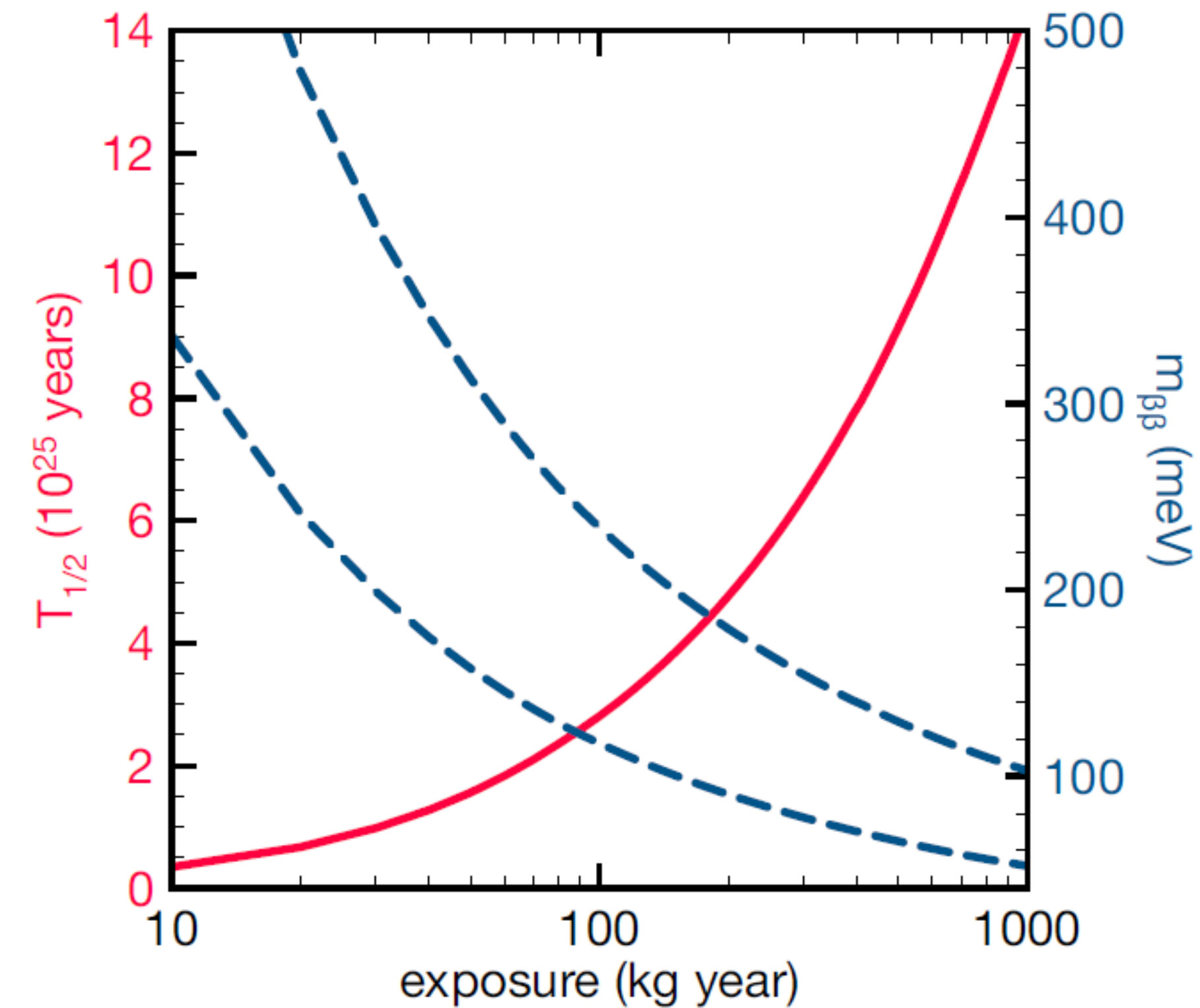
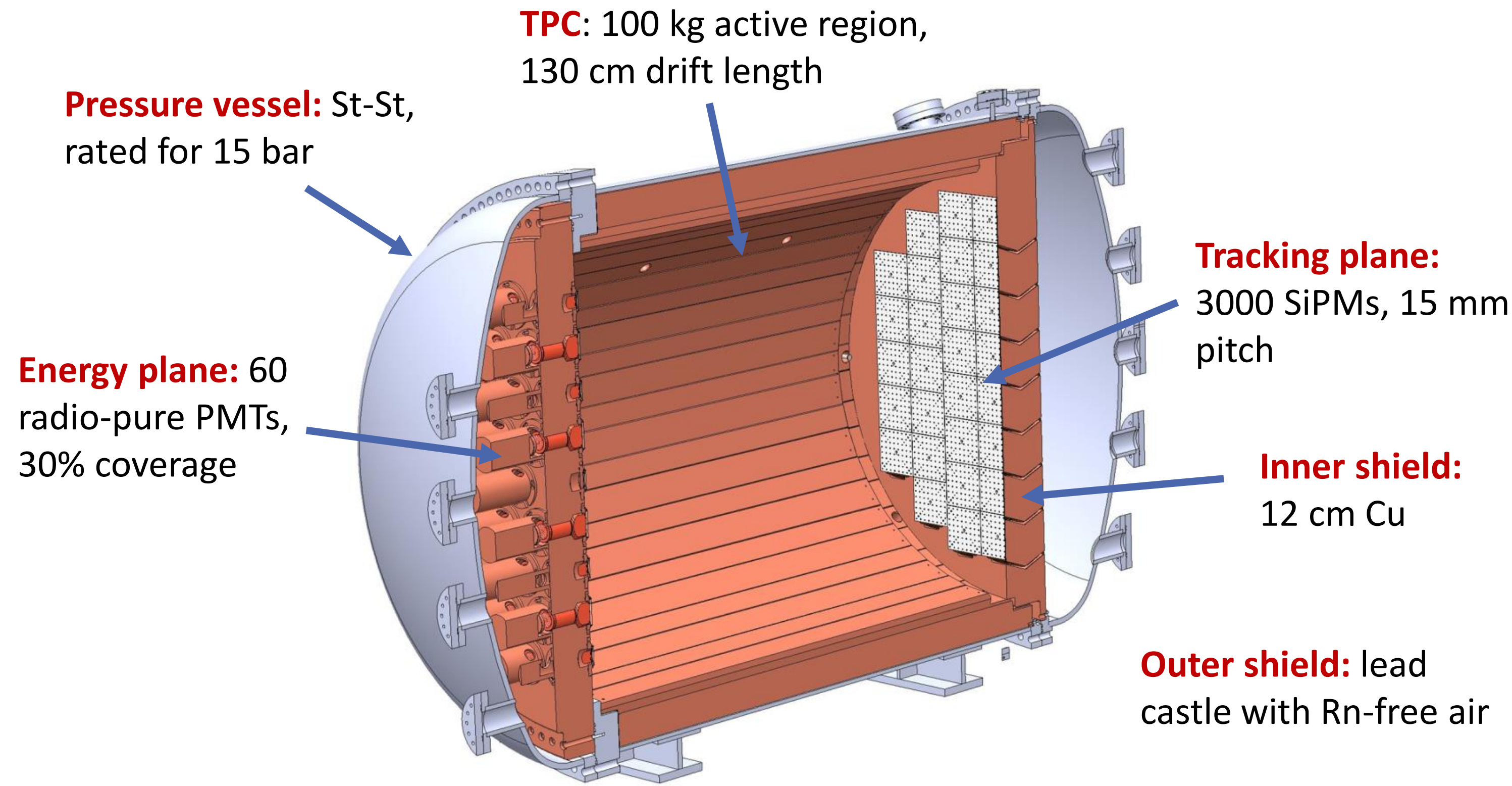


Ran with depleted Xe in 2018 -
now running with enriched Xe

arXiv:1804.02409

L Arazi

JHEP 2016 159

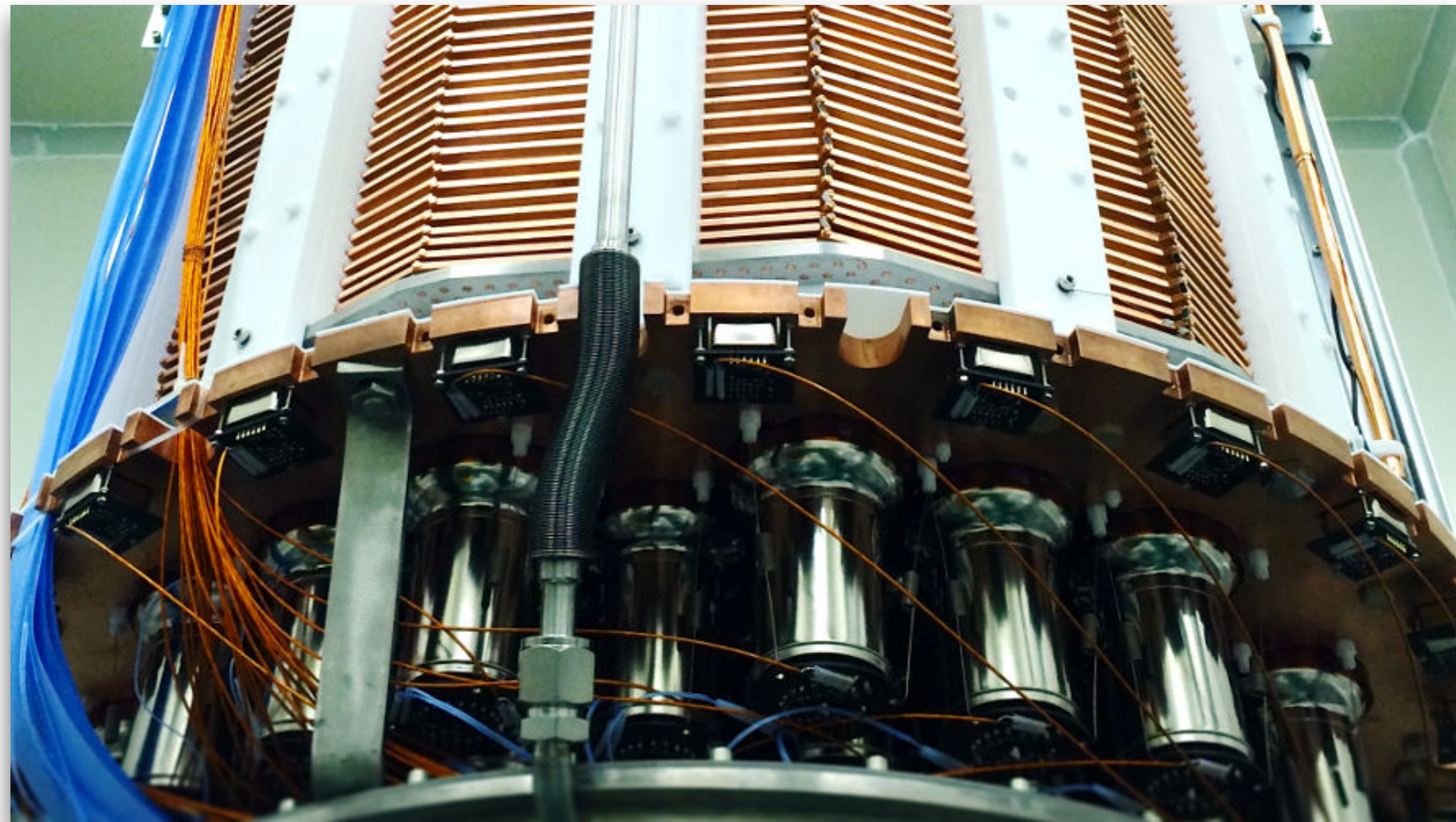


Resolution 1% FWHM
1 background count / 100 kg / year
Demonstrator for ton-scale experiment

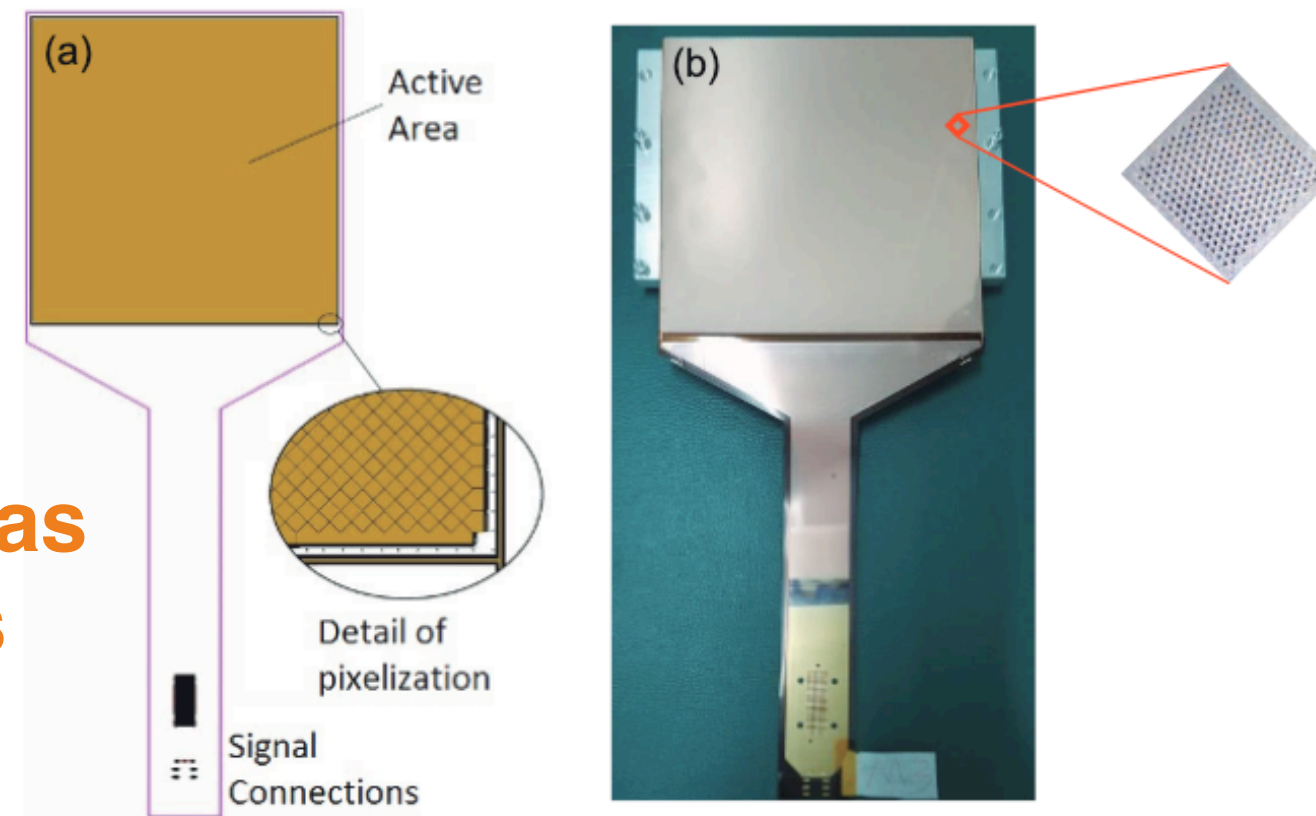
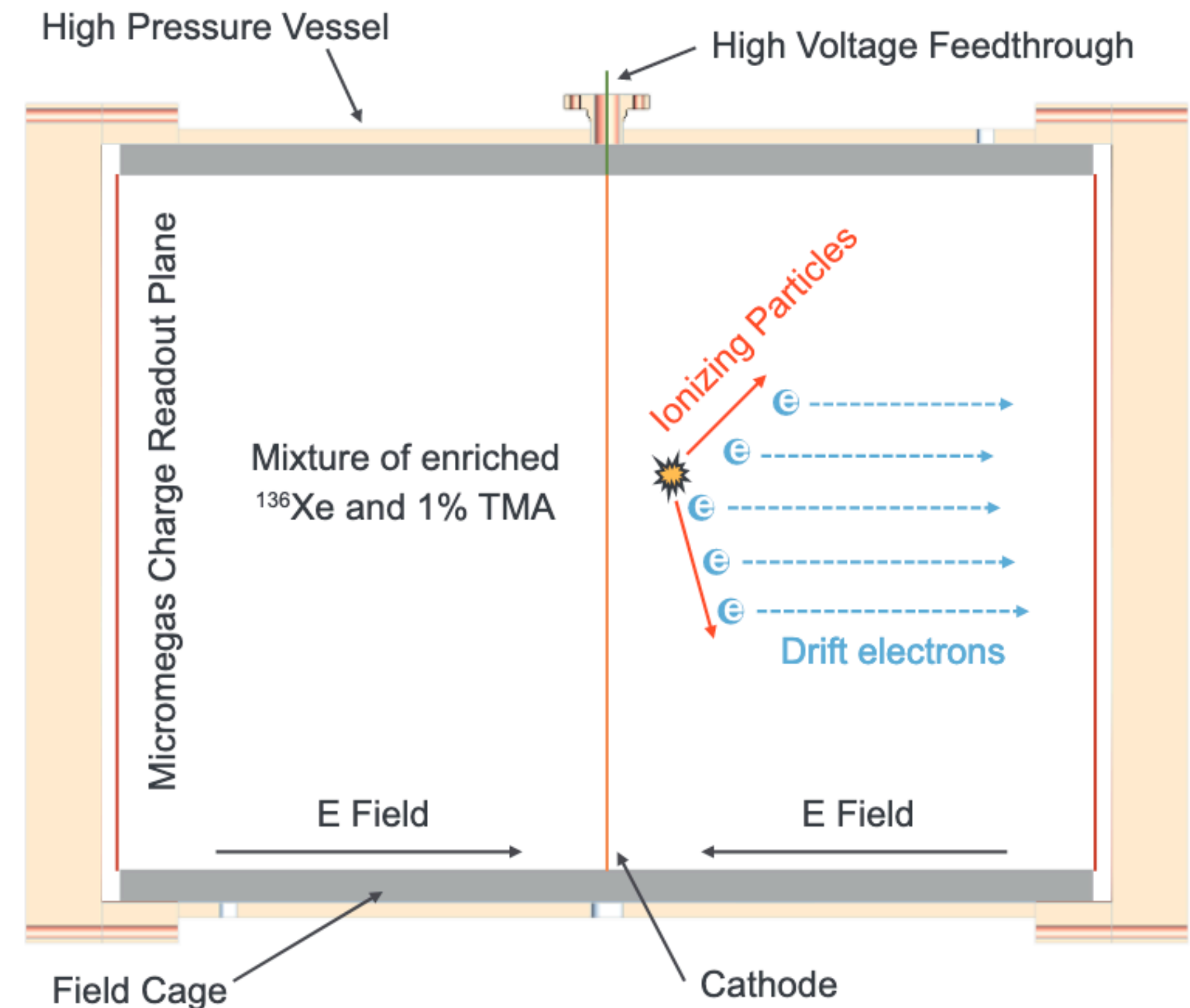


PandaX-III at China Jinping Underground Laboratory

PandaX-II

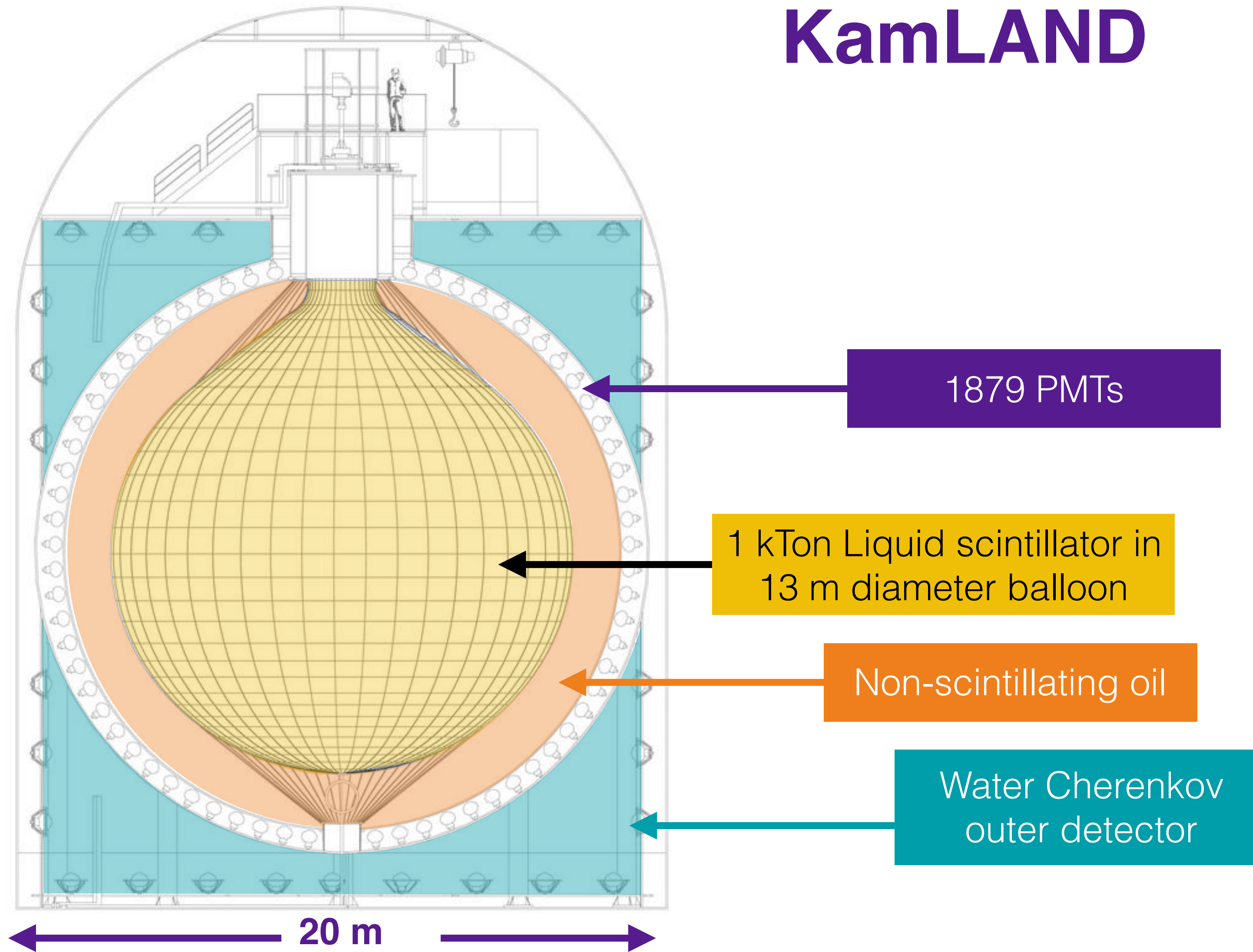


- **PandaX-II** is a **dark matter** direct detection experiment equipped with a half-ton scale **dual-phase TPC**
- **PandaX-III** will enrich with 200kg- 1 ton of 90% ^{136}Xe

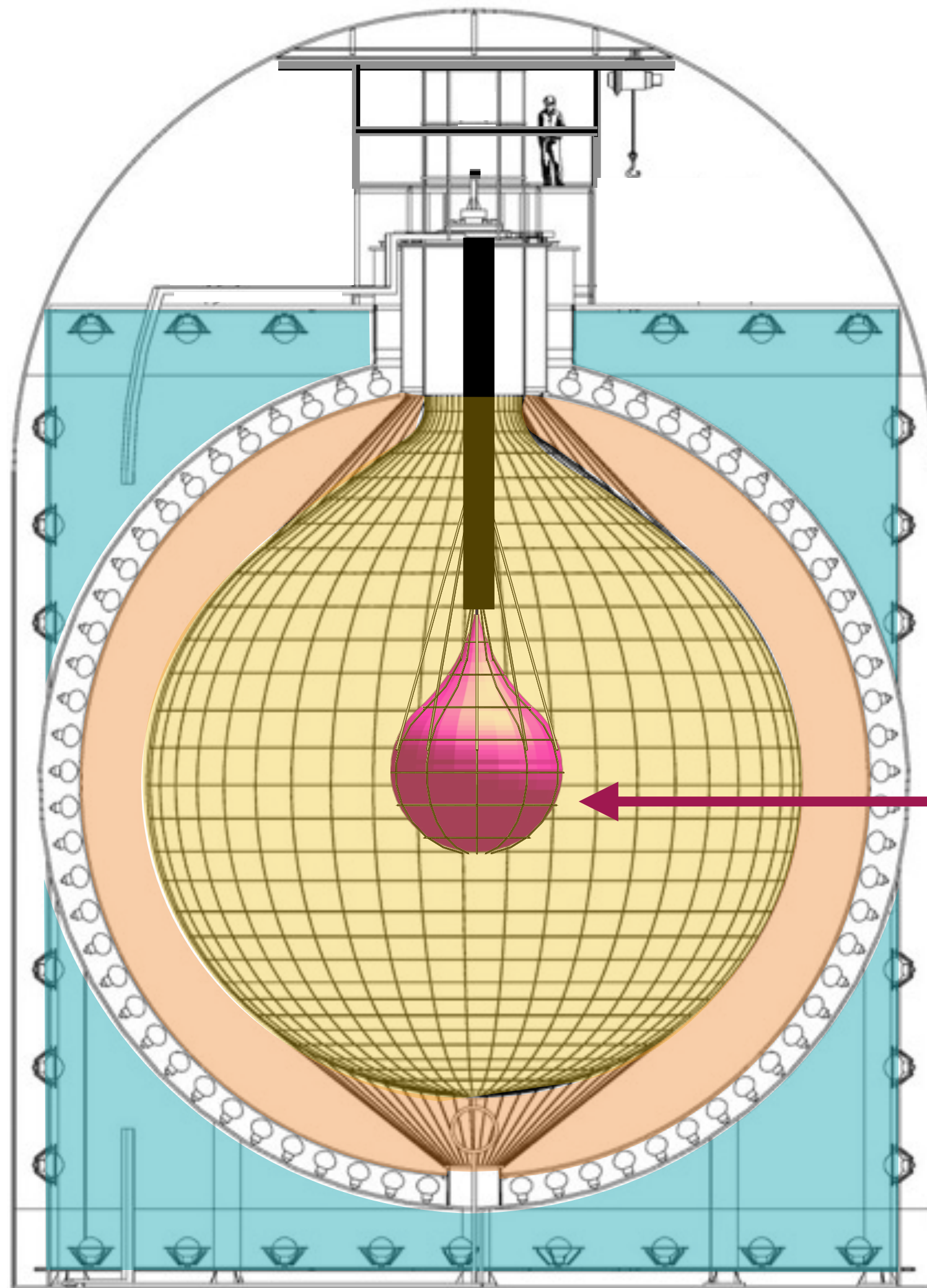


Microbulk Micromegas readout at both ends (no PMTs)

KamLAND

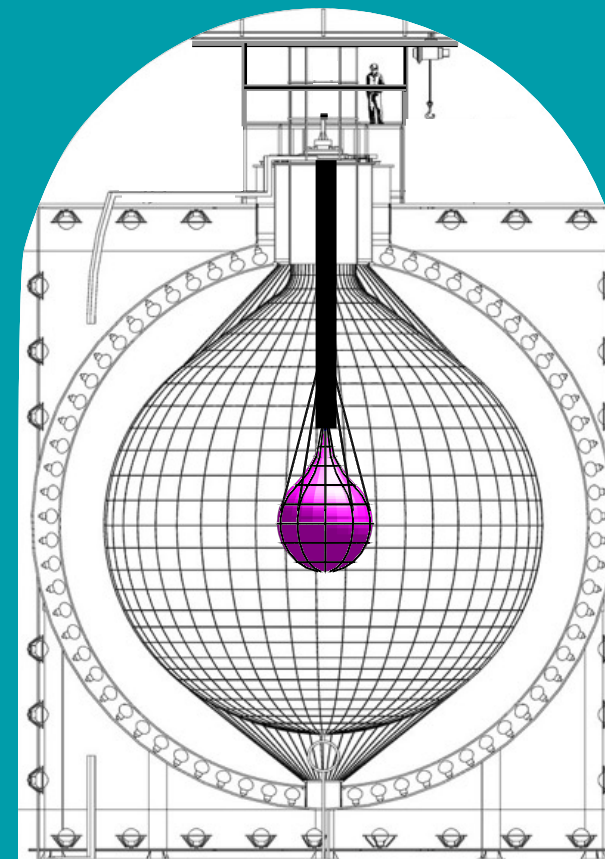
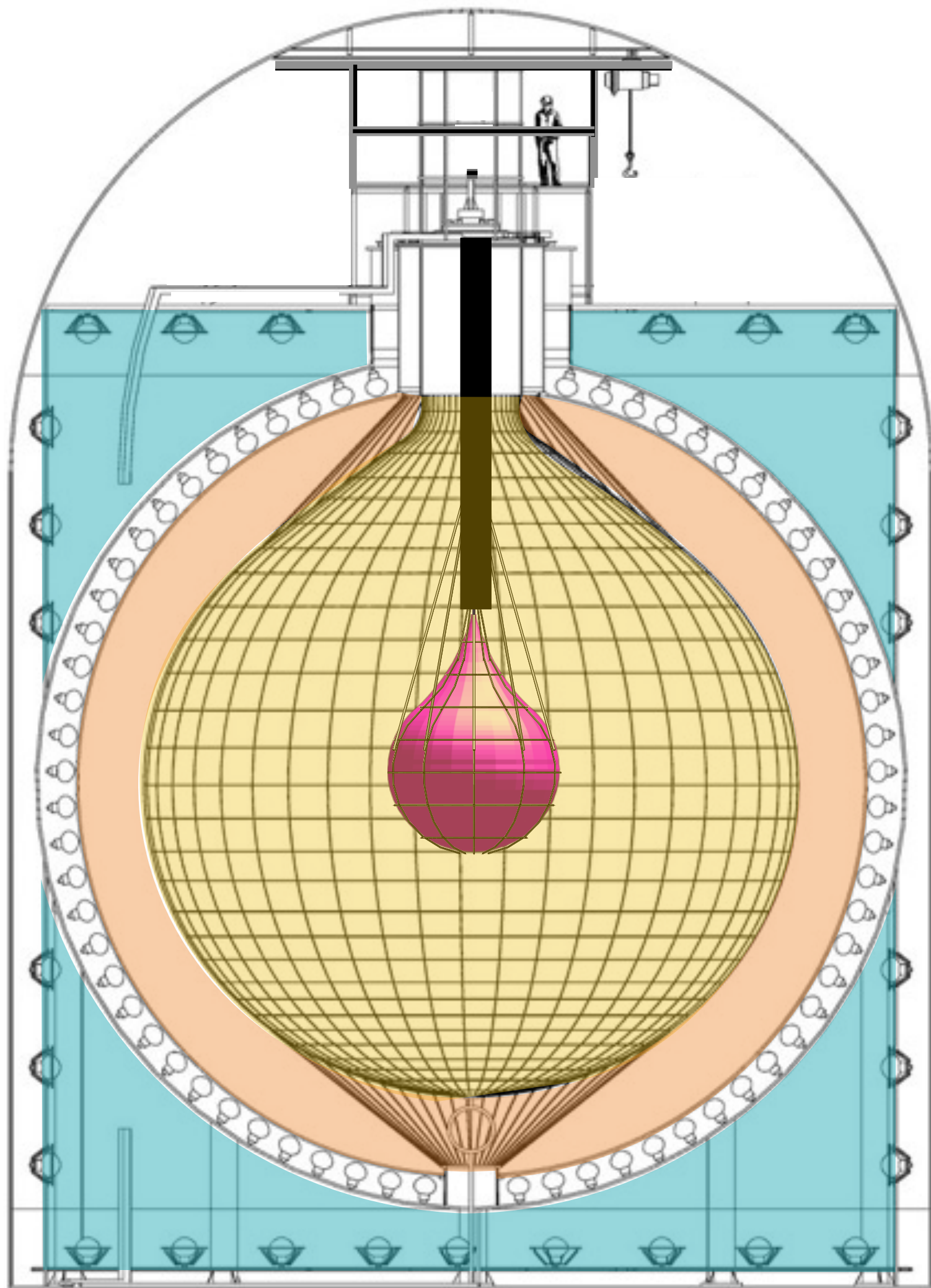


KamLAND - Zen



^{136}Xe -loaded liquid scintillator in inner balloon

KamLAND - Zen

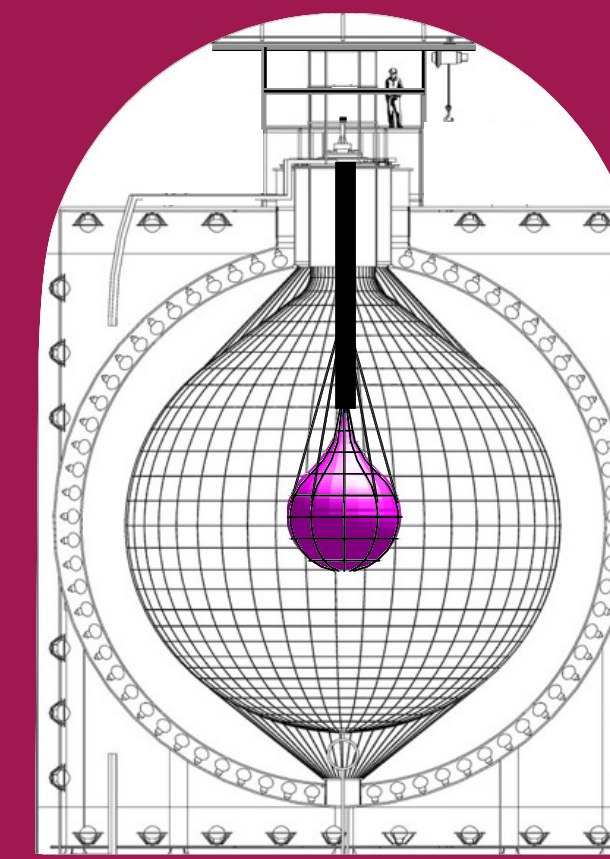


2011-2015

KamLAND-Zen 400

320-380 kg of Xe

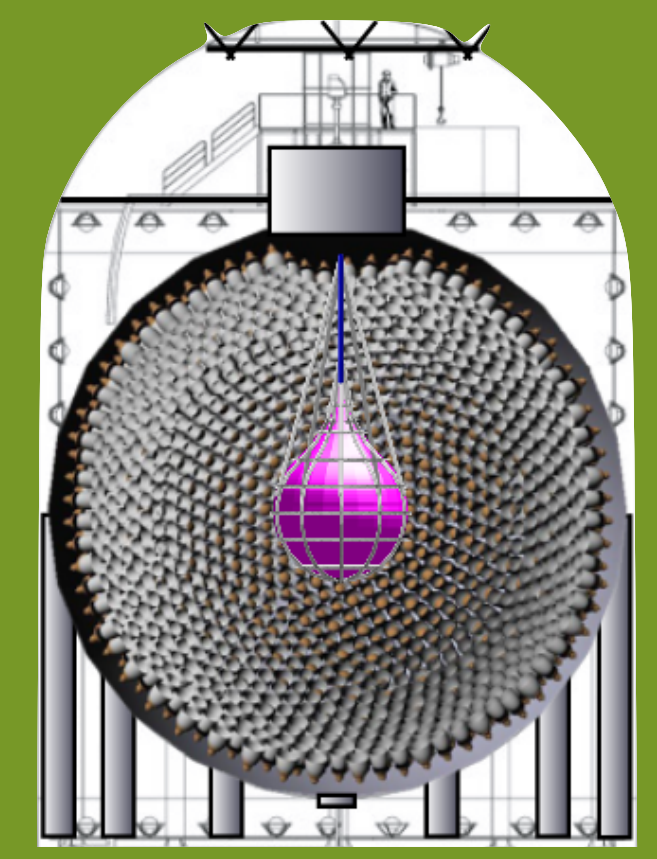
91% enriched



2018-

KamLAND-Zen 800

750 kg of Xe



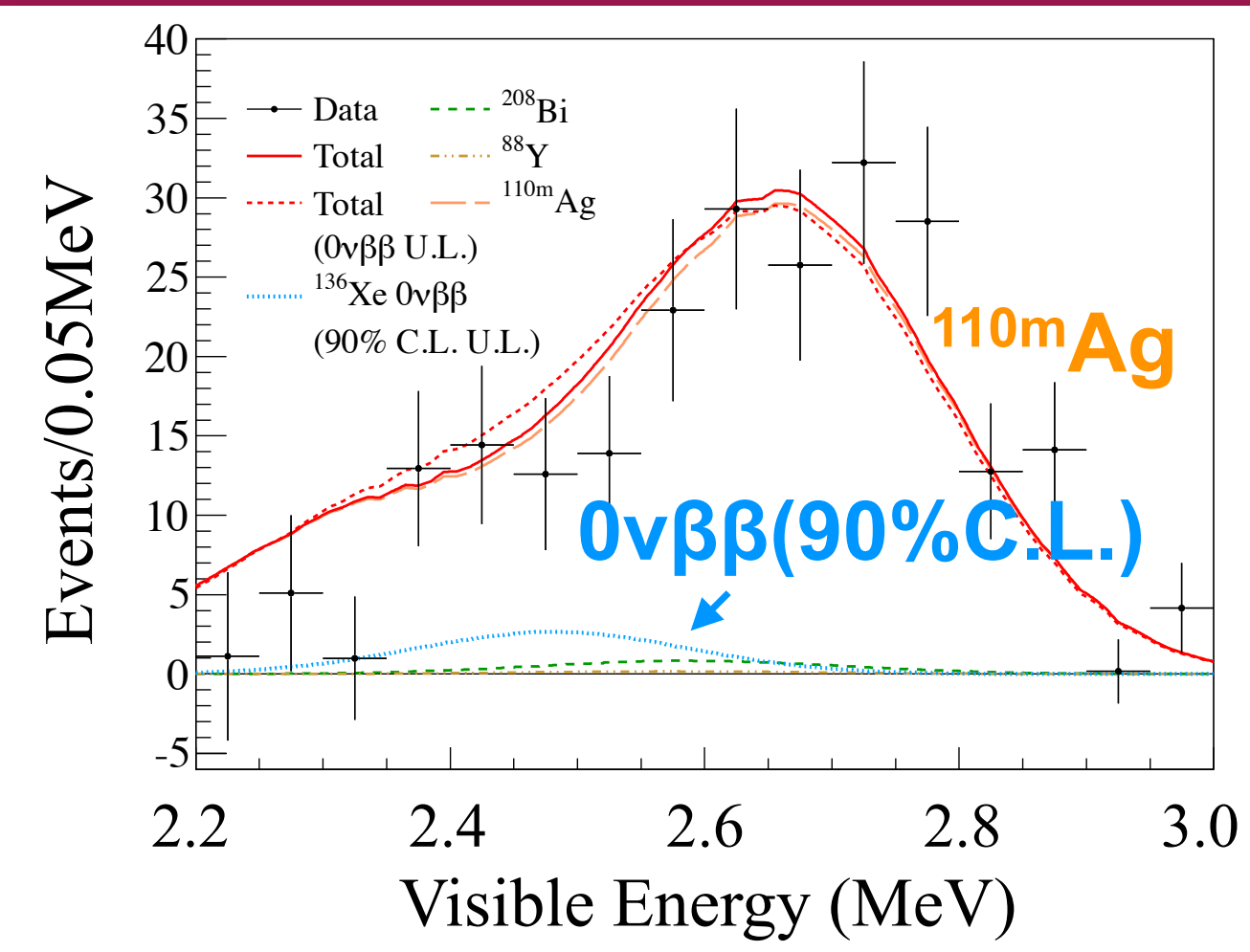
Future

KamLAND2-Zen

1 ton of Xe

Better energy resolution

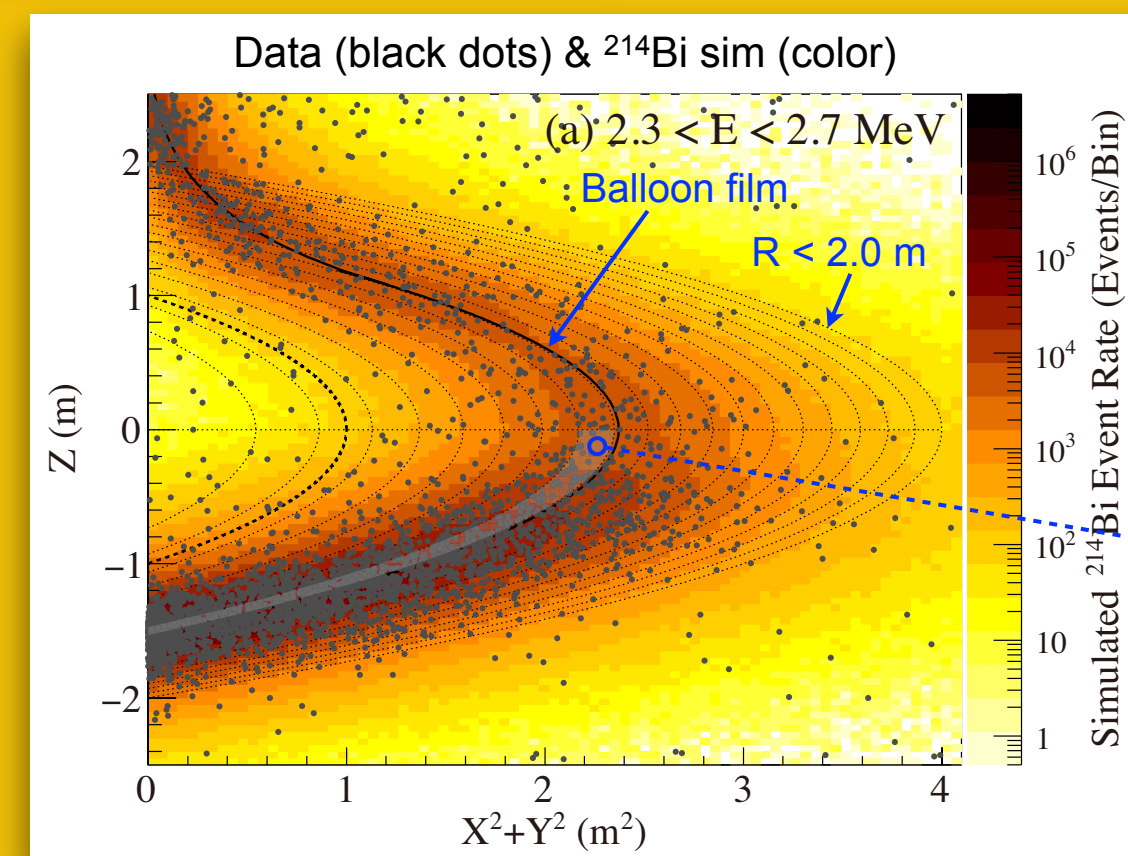
^{110m}Ag (β emitter) in XeLS from Fukushima fallout: scintillator cleaned for Phase II



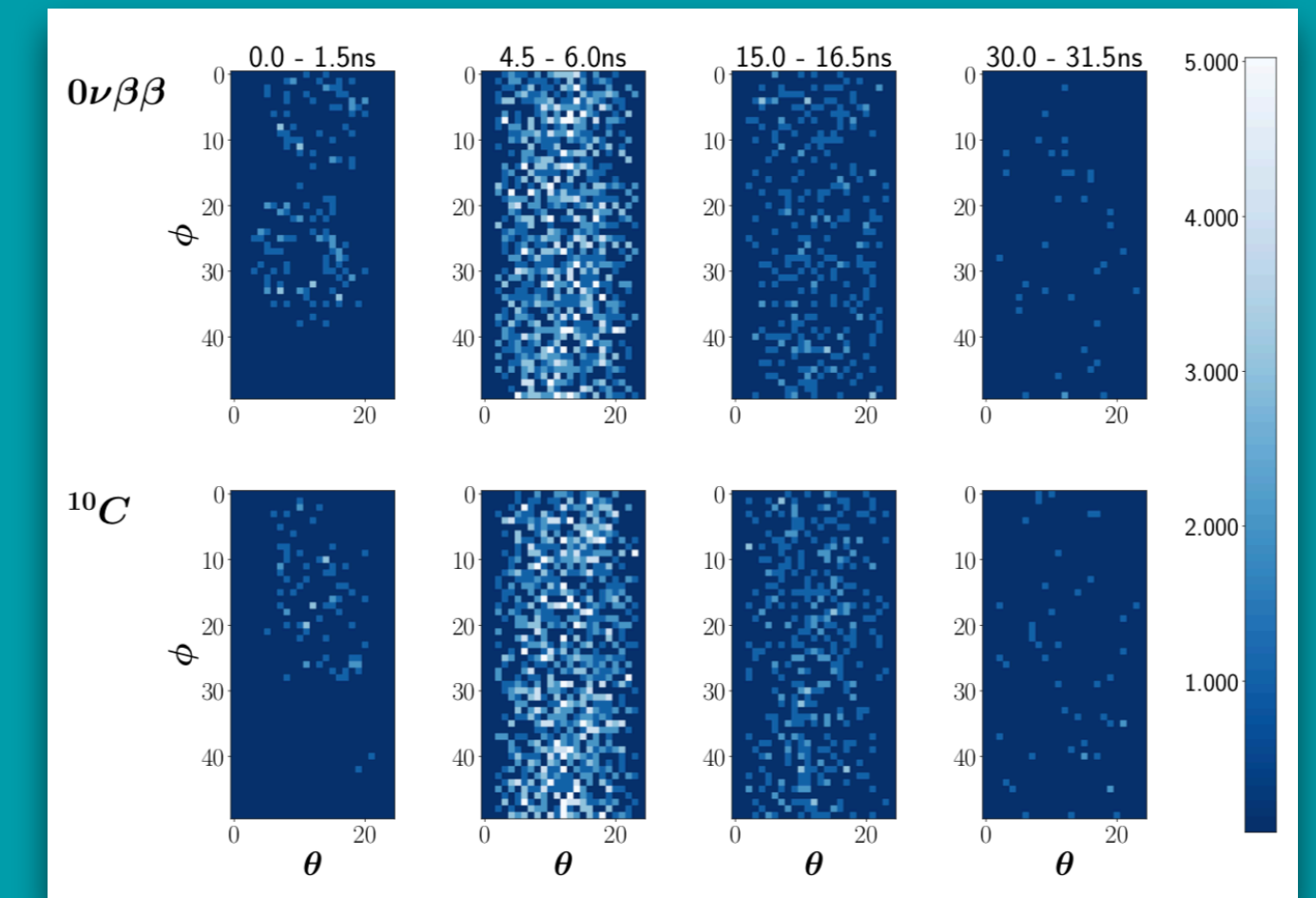
Inner balloon contamination from mine air



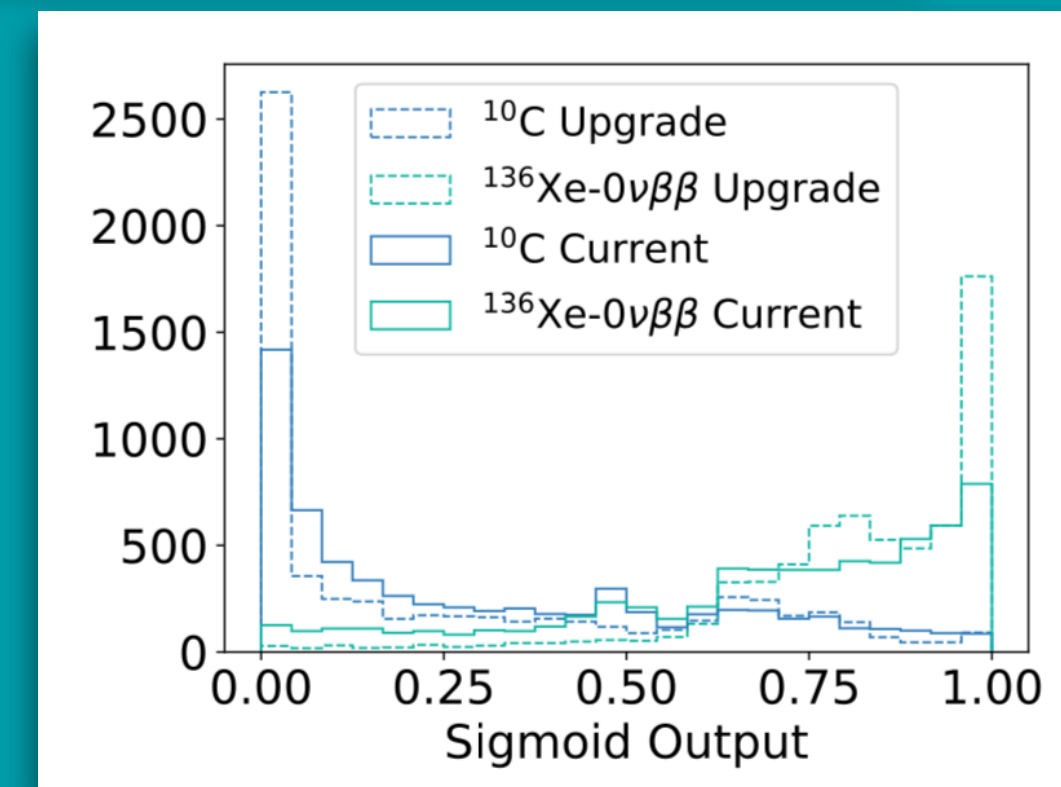
Large, monolithic detector lets you limit fiducial volume



Muon-induced spallation product ^{10}C : triple coincidence muon + neutron + ^{10}C β^+ decay



Separation with computer-vision algorithm



[arXiv:1812.02906](https://arxiv.org/abs/1812.02906) [physics.ins-det]

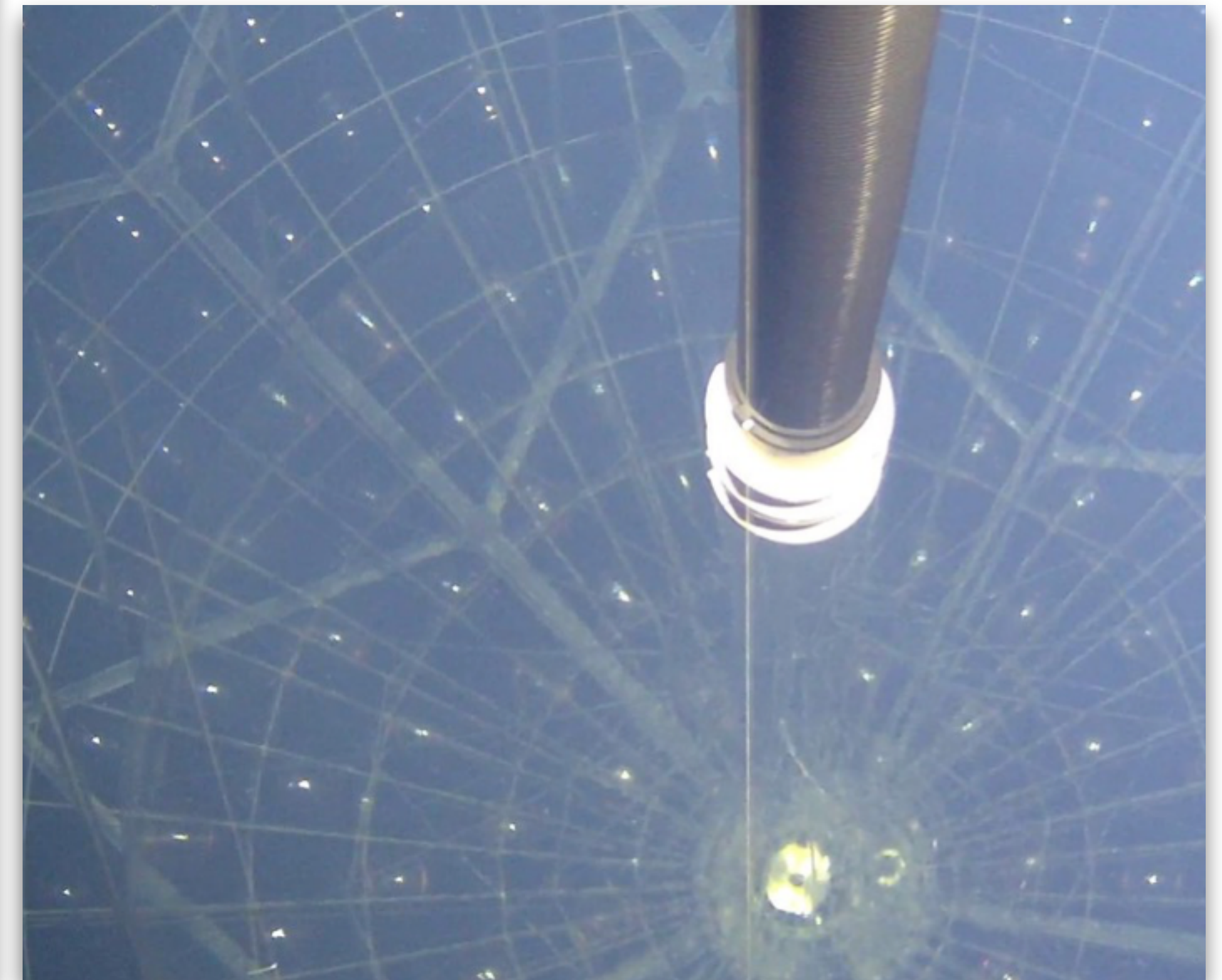


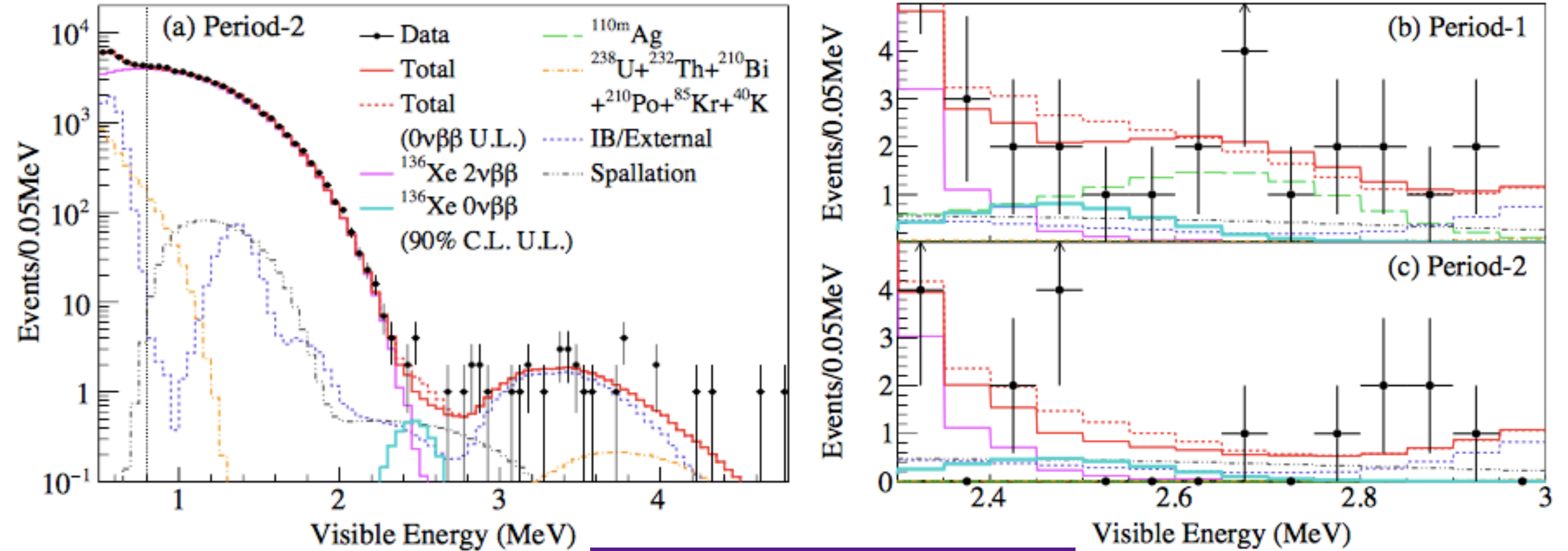
Inner balloon being prepared...



... and installed into the detector...

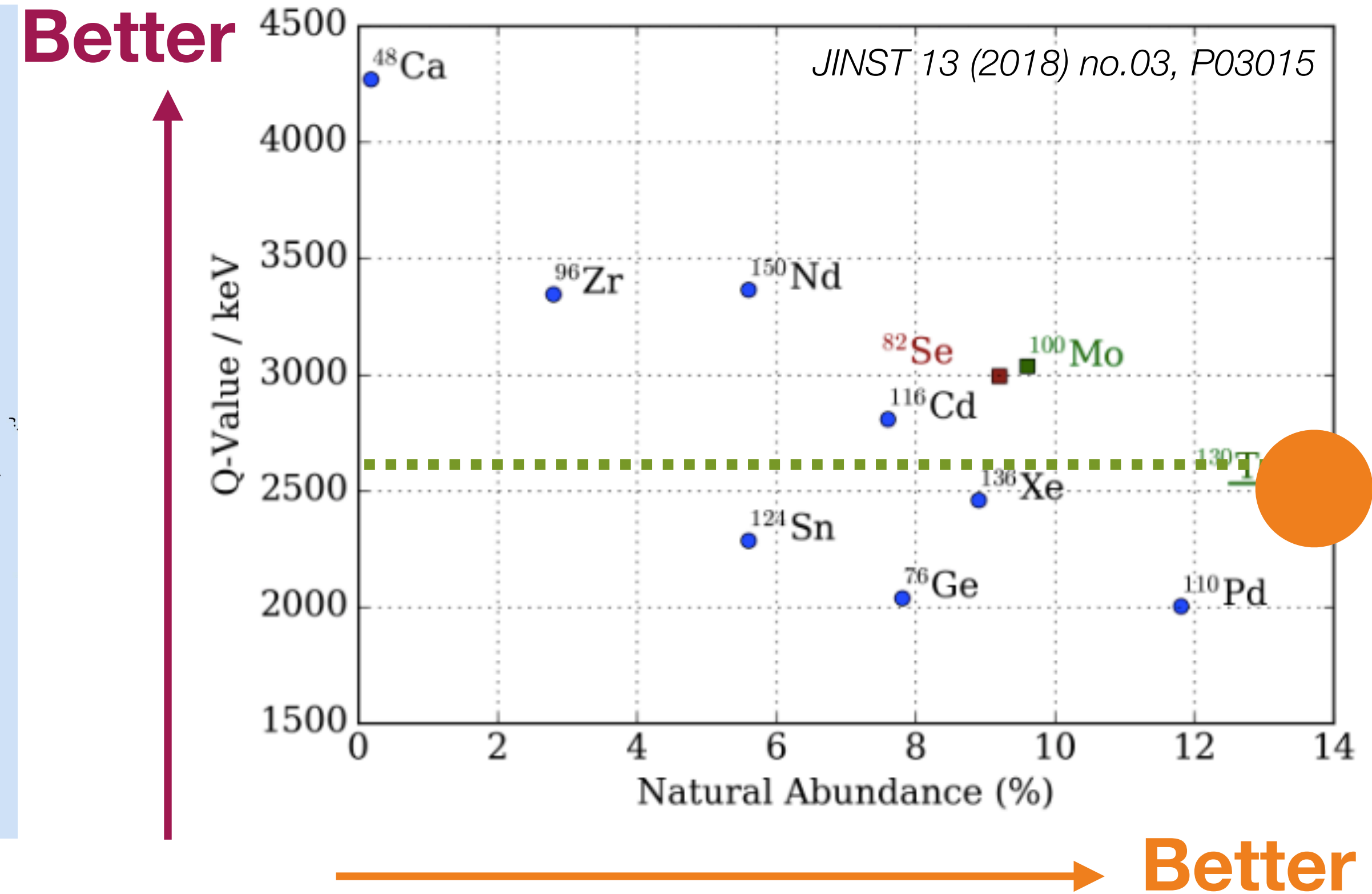
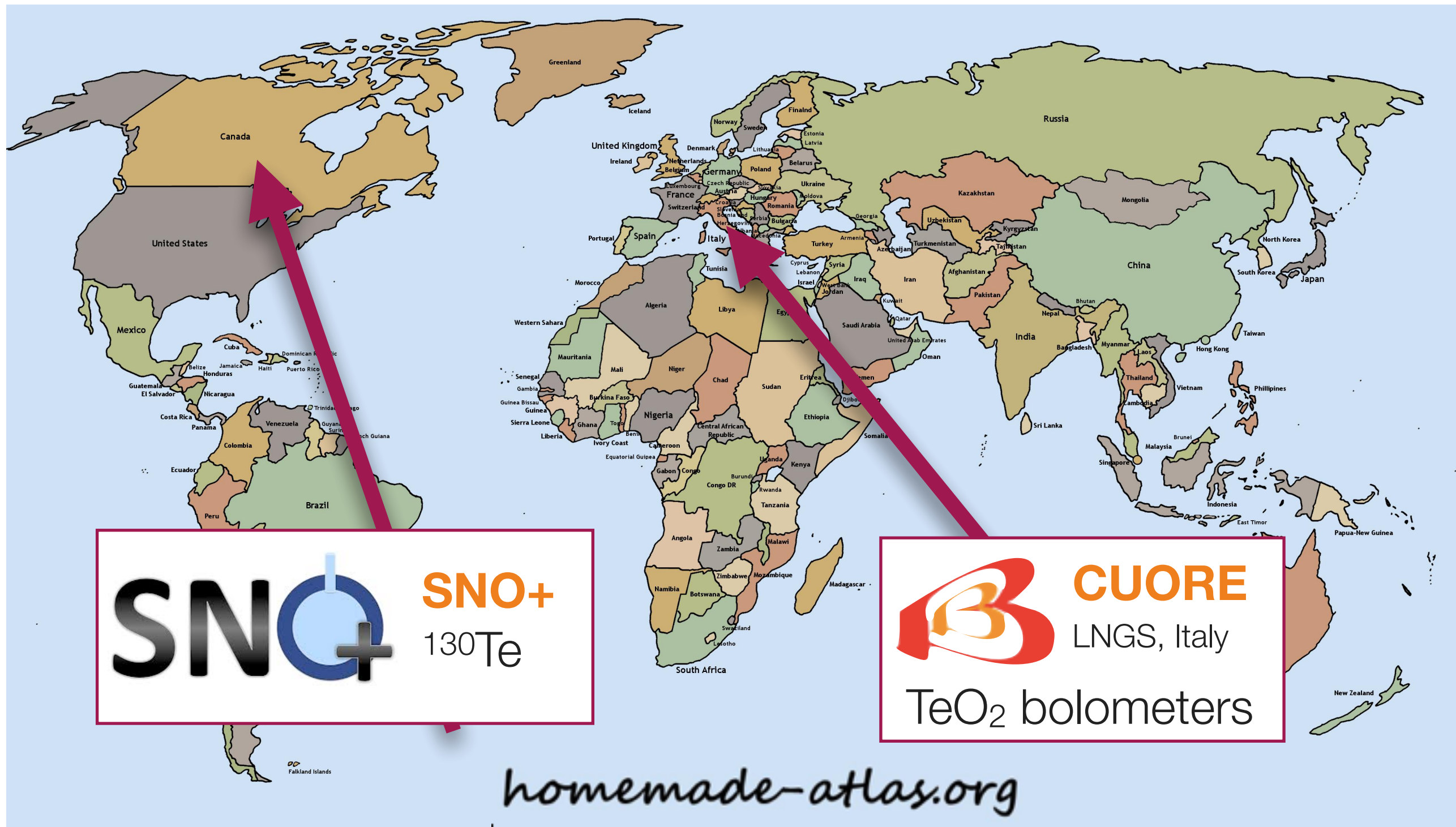
...ready to be filled with scintillator and ^{136}Xe



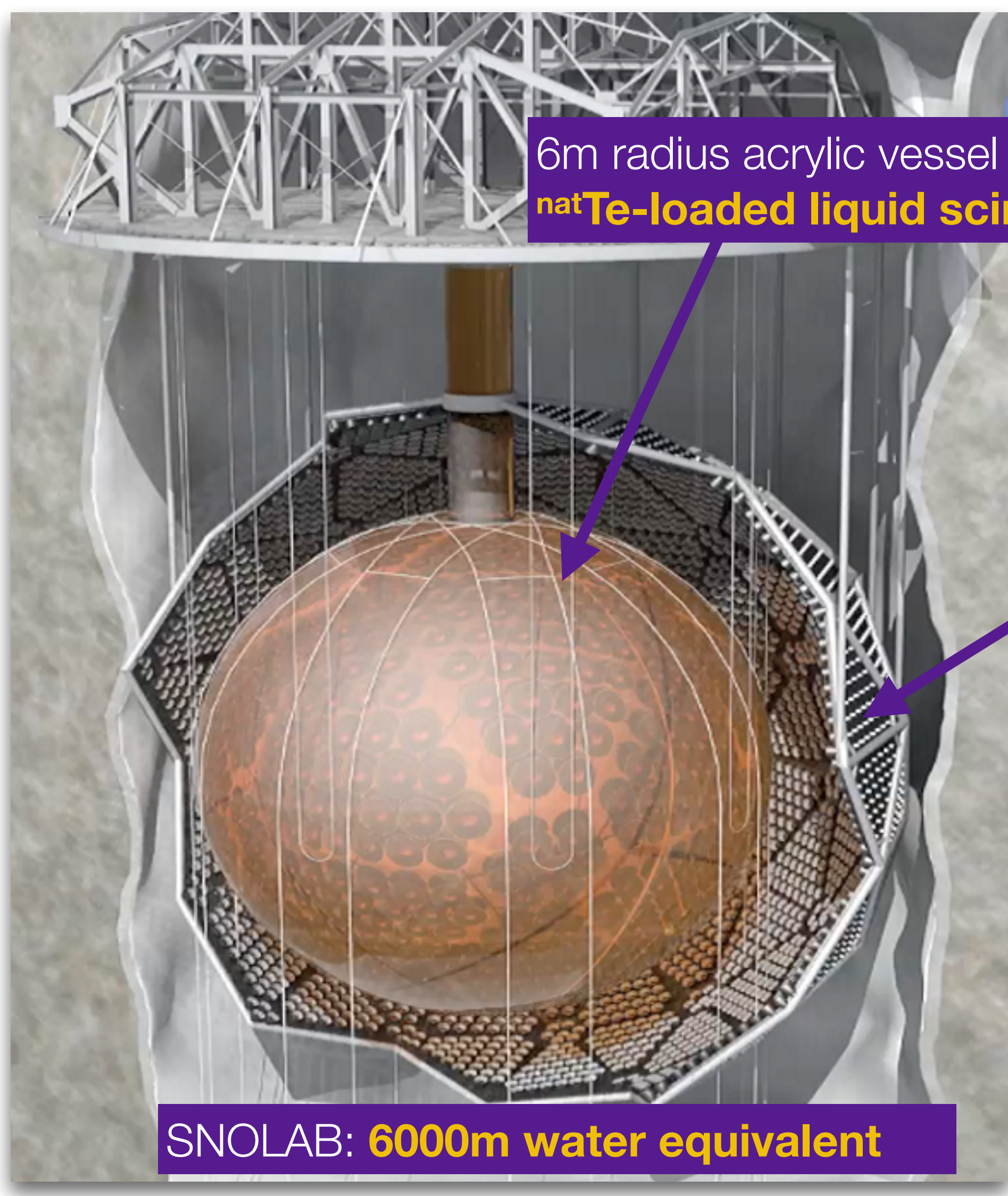


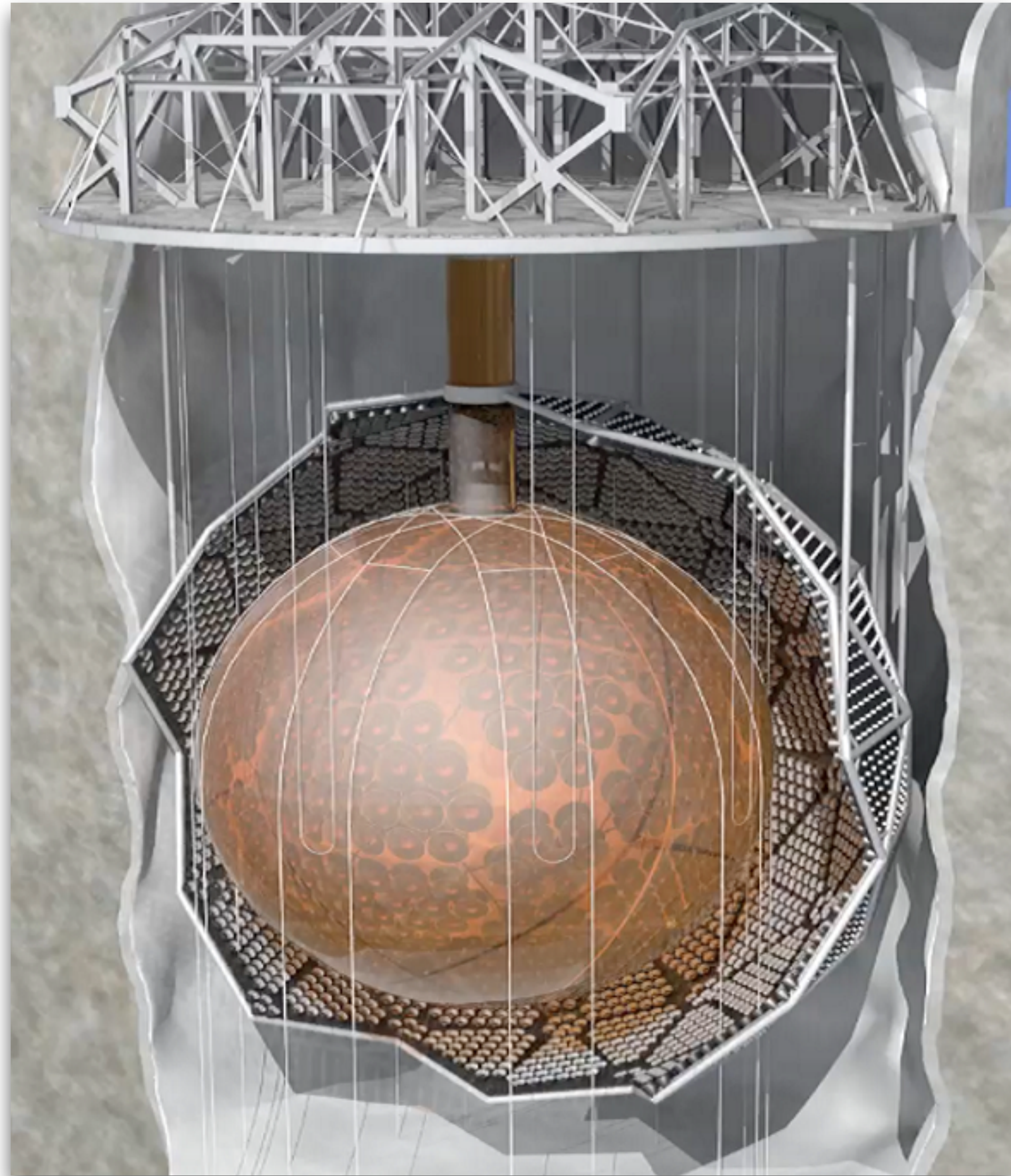
$T_{1/2} > 1.07 \times 10^{26}$ years
 ($\langle m_{\beta\beta} \rangle < 61-165$ meV)
 504 kg.years

Phys Rev Lett 117, 082503 (2016)



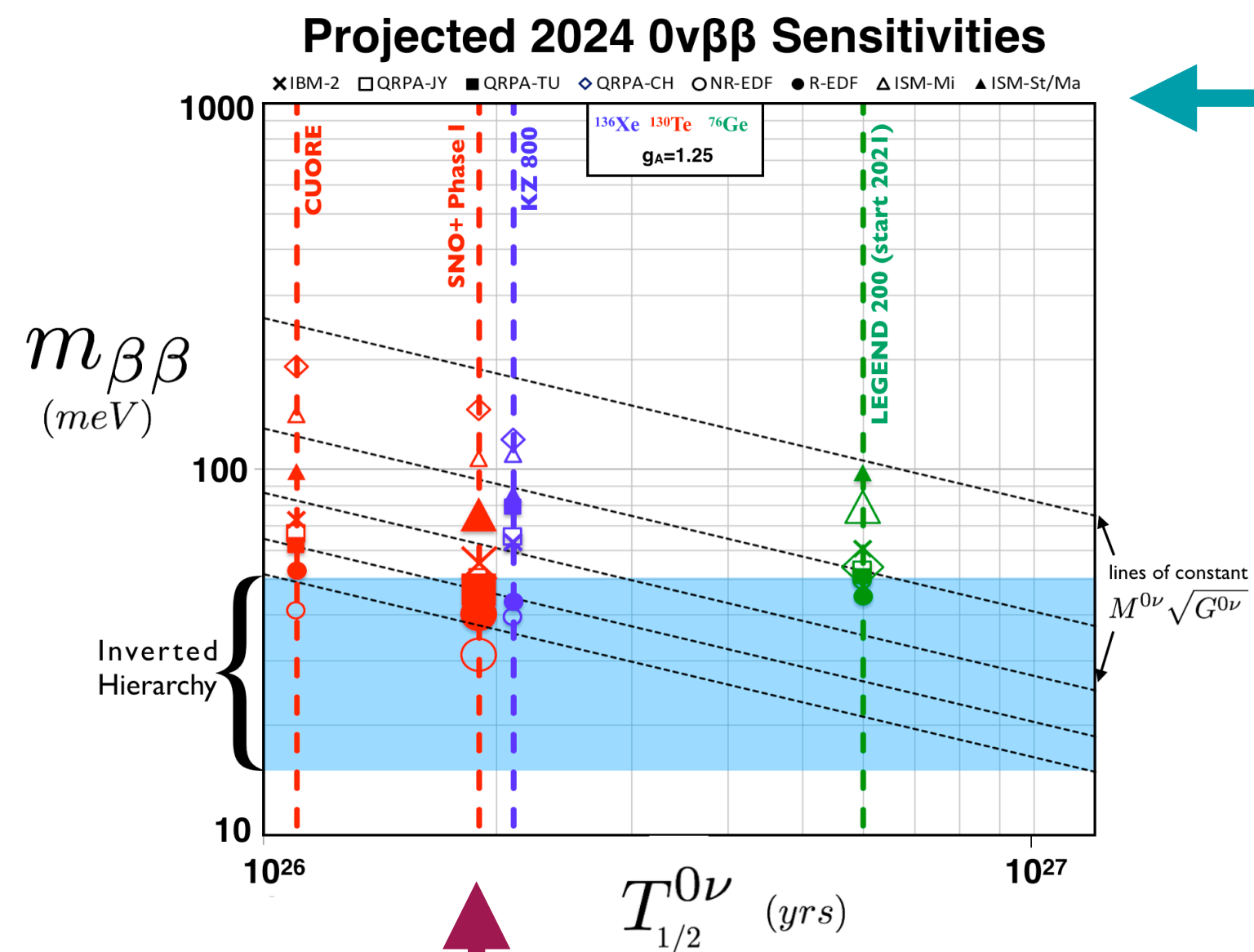
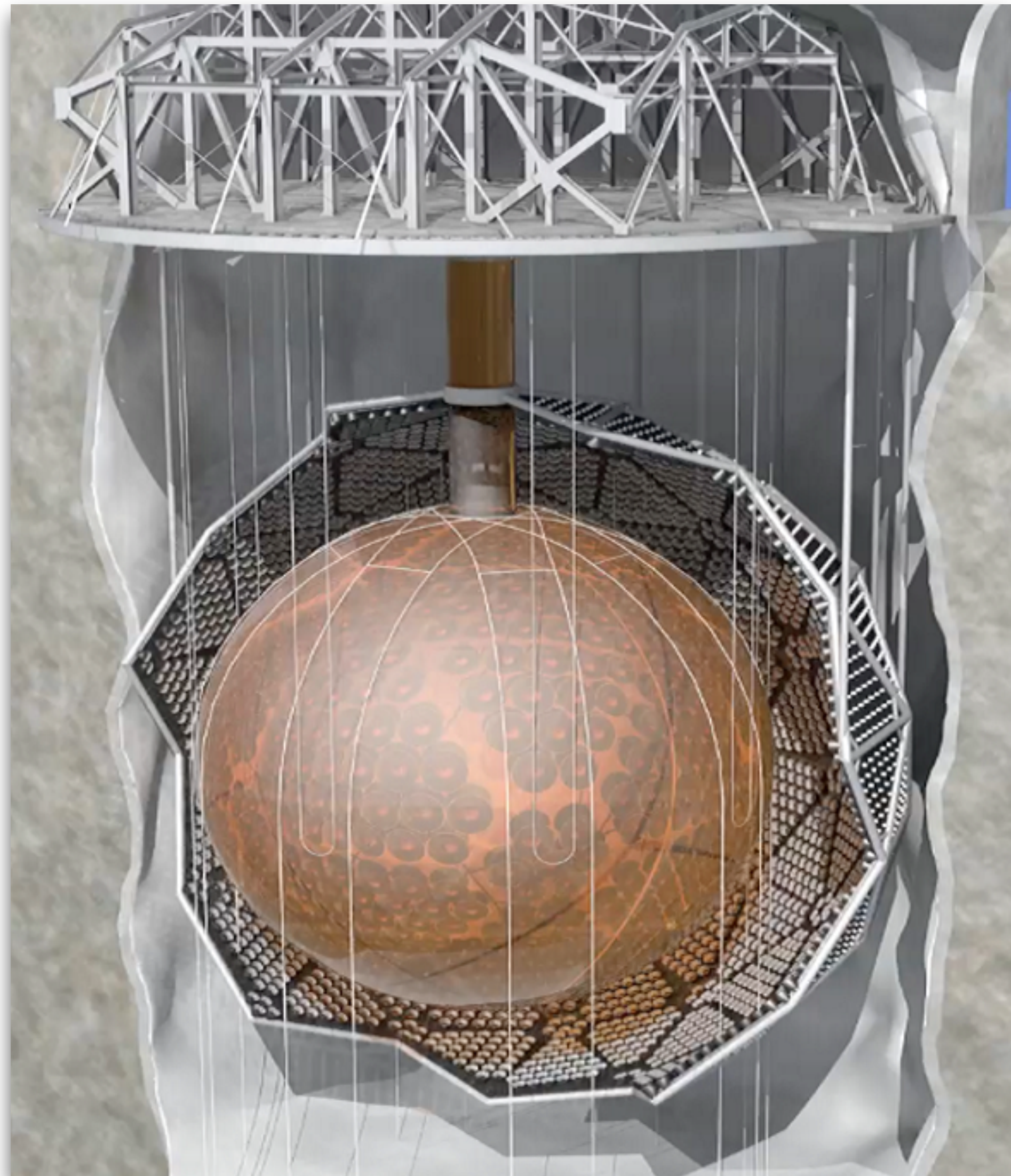
- ^{130}Te has 34% natural **abundance**





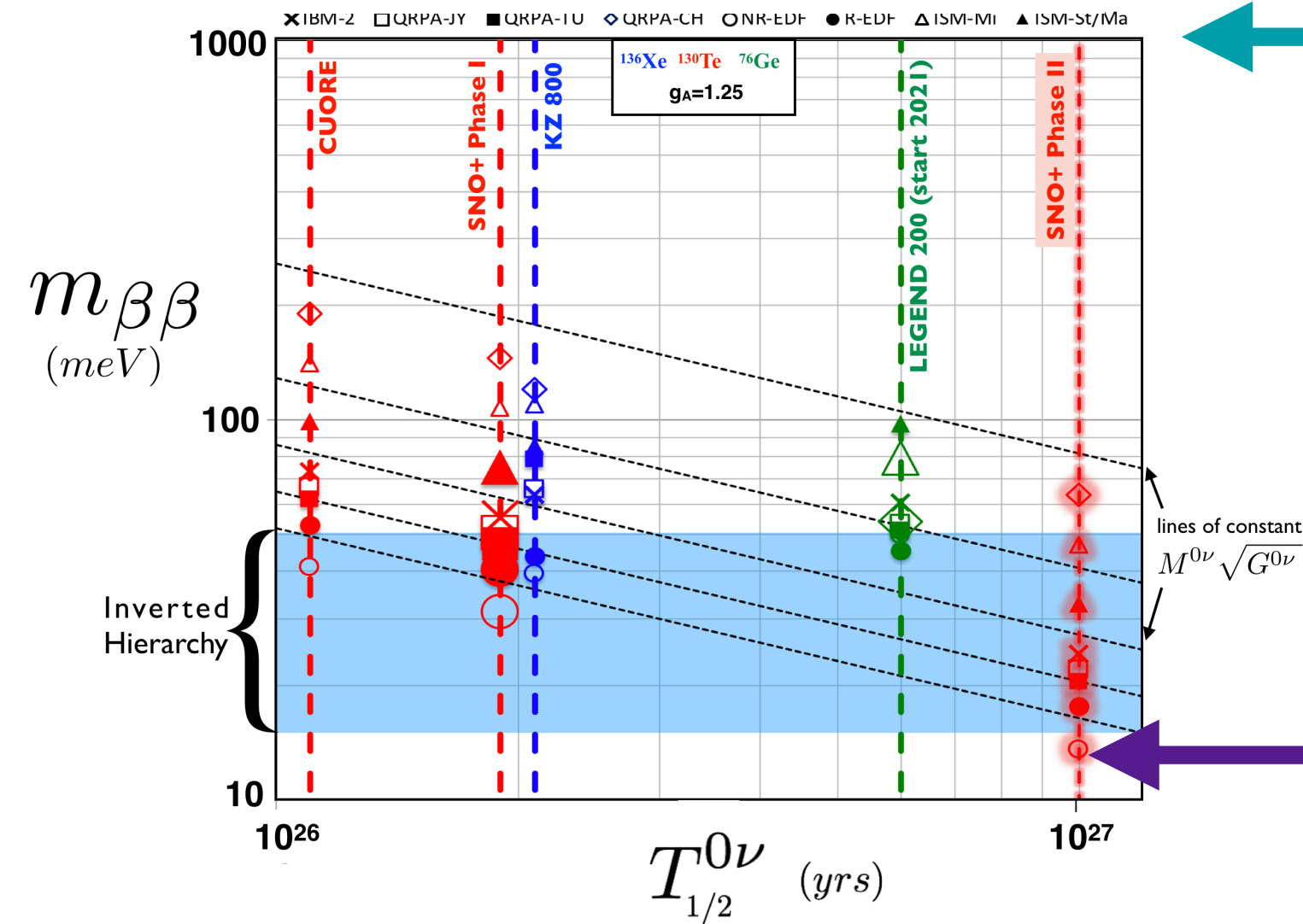
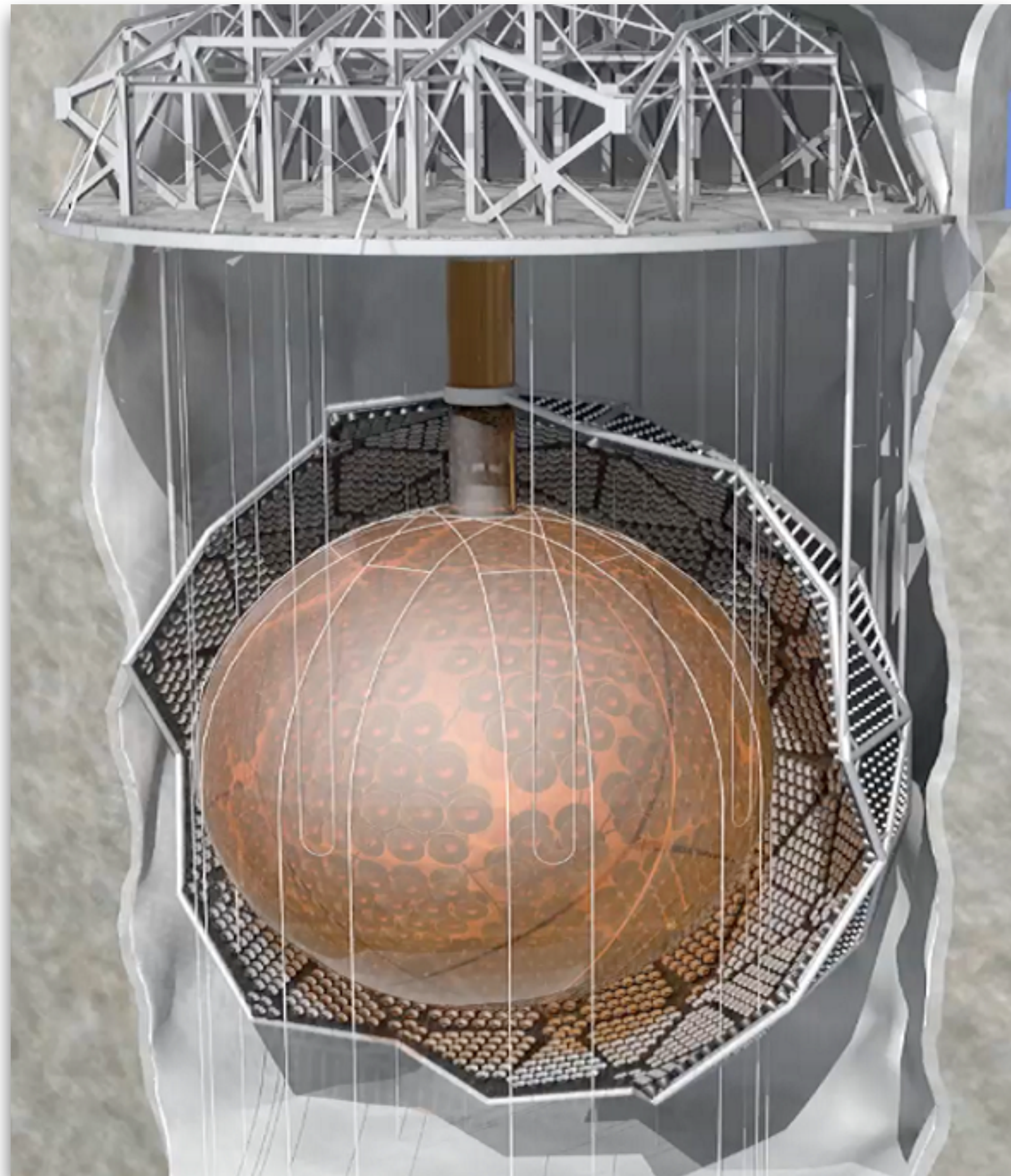
Thanks to Steve Biller & Esther Turner for slide content

- Highly **economical**
 - ^{130}Te is the most economically scalable **isotope** (high natural abundance);
 - Liquid scintillator also very economically scalable **detector technology!**
- Potential for dramatic **scale-up**



Different matrix element calculations give different masses for the same half-life

- Highly **economical**
 - ^{130}Te is the most economically scalable **isotope** (high natural abundance);
 - Liquid scintillator also very economically scalable **detector technology!**
- Potential for dramatic **scale-up**
- Allows **sensitivity** above current leading measurement:
 - $T_{1/2}^{0\nu\beta\beta} > 2.1e^{26}$ years (**$m_{\beta\beta} < 37- 89$ meV**) after 5 years of running



Different matrix element calculations give different masses for the same half-life

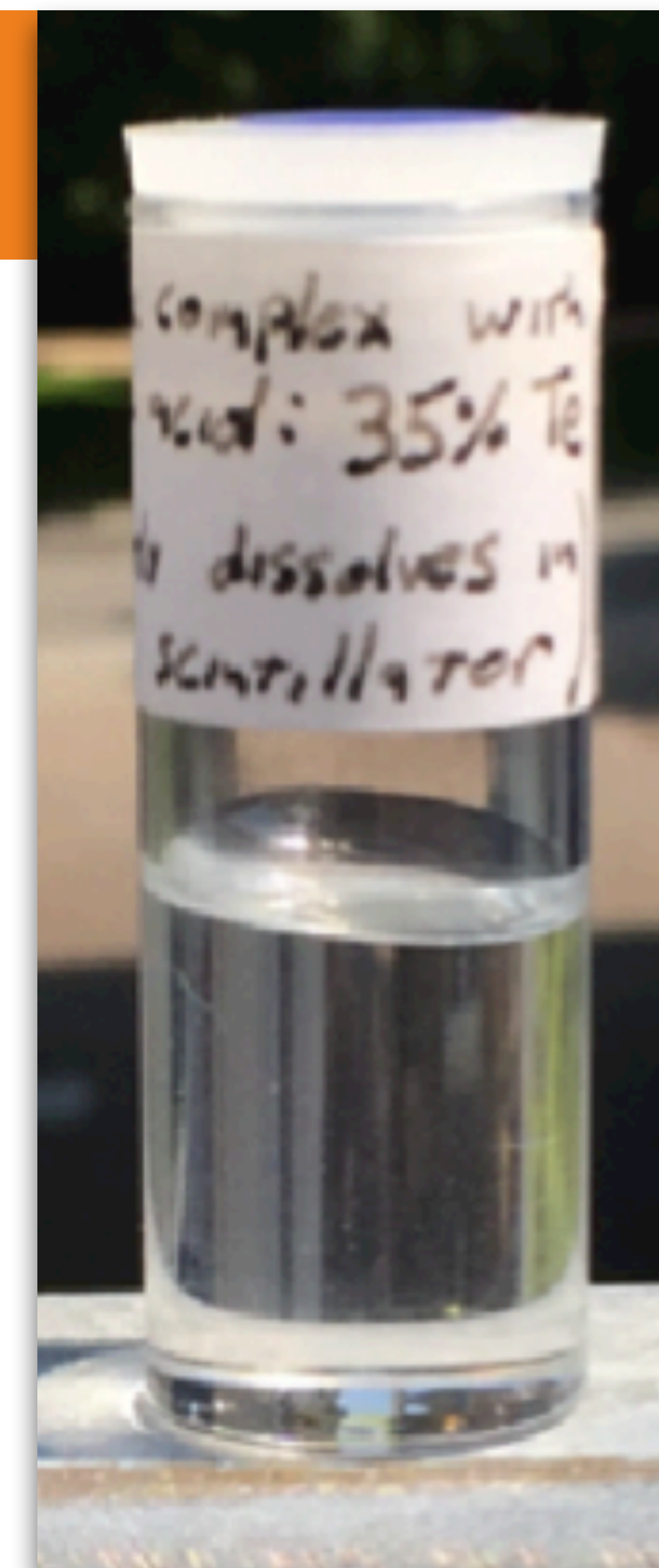
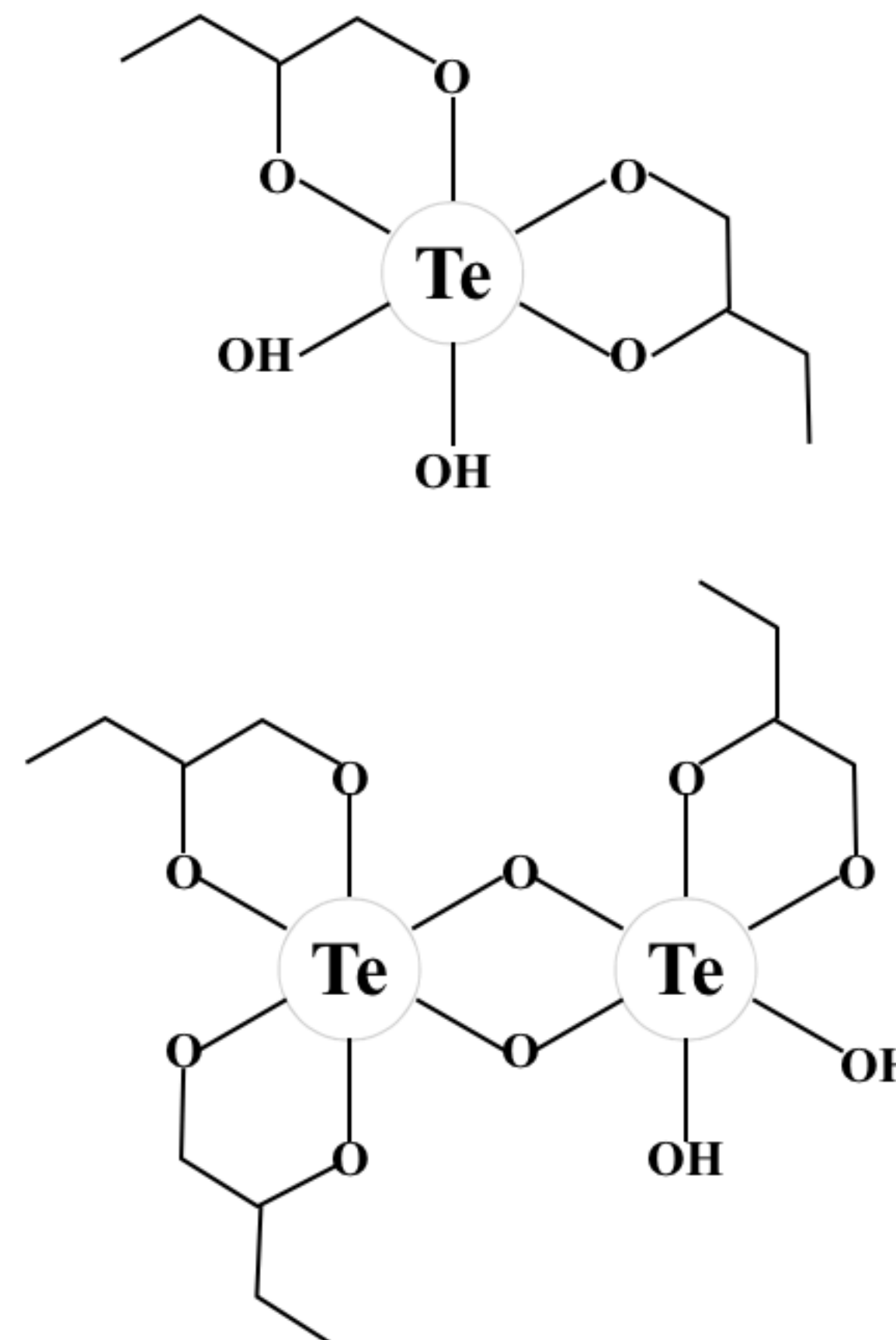
New chemistry developments may make it possible to simply increase the loading in the current instrument to achieve this

- Highly **economical**
 - ^{130}Te is the most economically scalable **isotope** (high natural abundance);
 - Liquid scintillator also very economically scalable **detector technology!**
- Potential for dramatic **scale-up**
- Allows **sensitivity** above current leading measurement:
 - $T_{1/2}^{0\nu\beta\beta} > 2.1 \times 10^{26}$ years ($m_{\beta\beta} < 37-89$ meV) after 5 years of running
- **Phase II** could reach 10^{27} years with the same detector but **higher loading**

New loading method: Te-butenediol complex dissolved in liquid scintillator

- Simple **synthesis**
- Single **safe**, distillable chemical
- Low **radioactivity** levels
- Minimal optical **absorption**
- High **light levels** at 0.5% ^{nat}Te loading

Natural tellurium is 34% ¹³⁰Te



- Operating with **water** from 2017

- Invisible nucleon decay
- Solar neutrinos
- Supernova neutrinos

PHYSICAL REVIEW D **99**, 012012 (2019)

Measurement of the ^8B solar neutrino flux in SNO+ with very low backgrounds

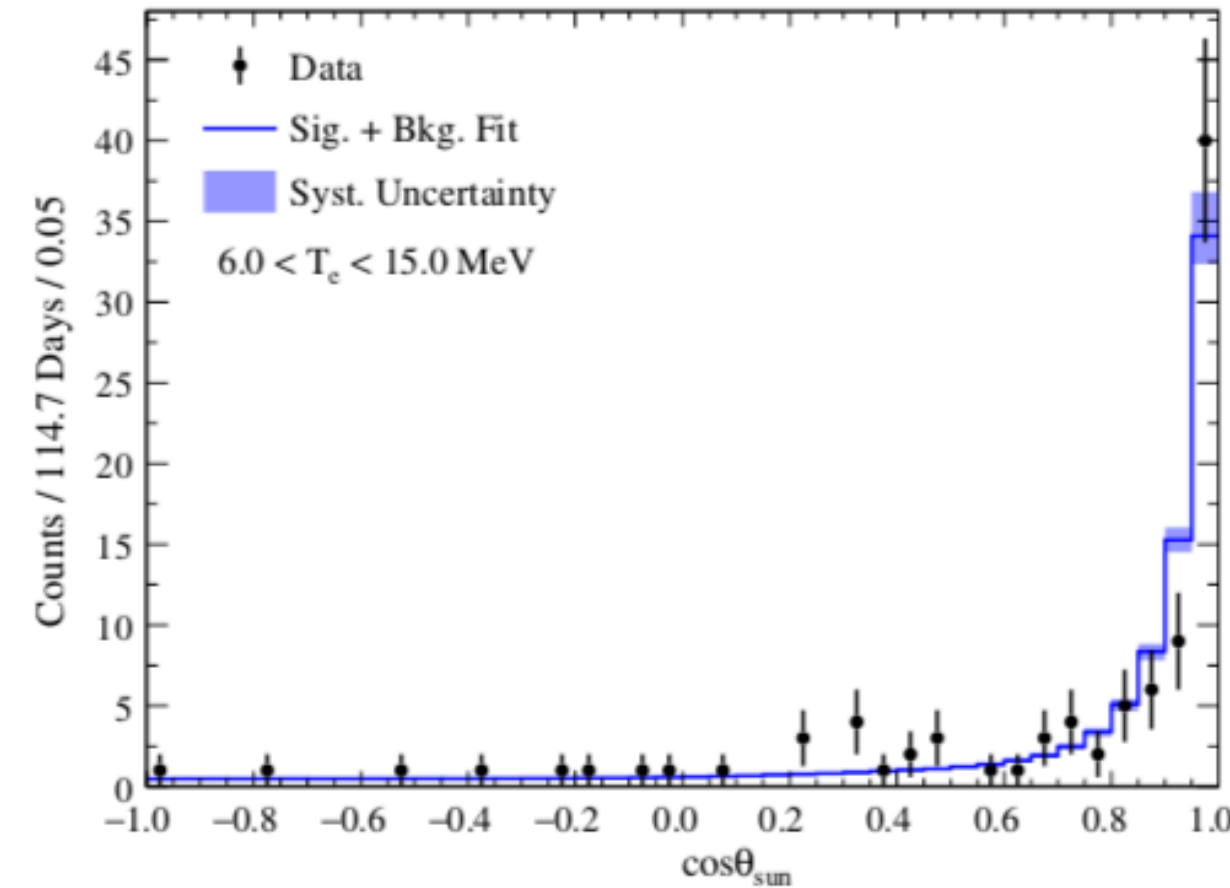
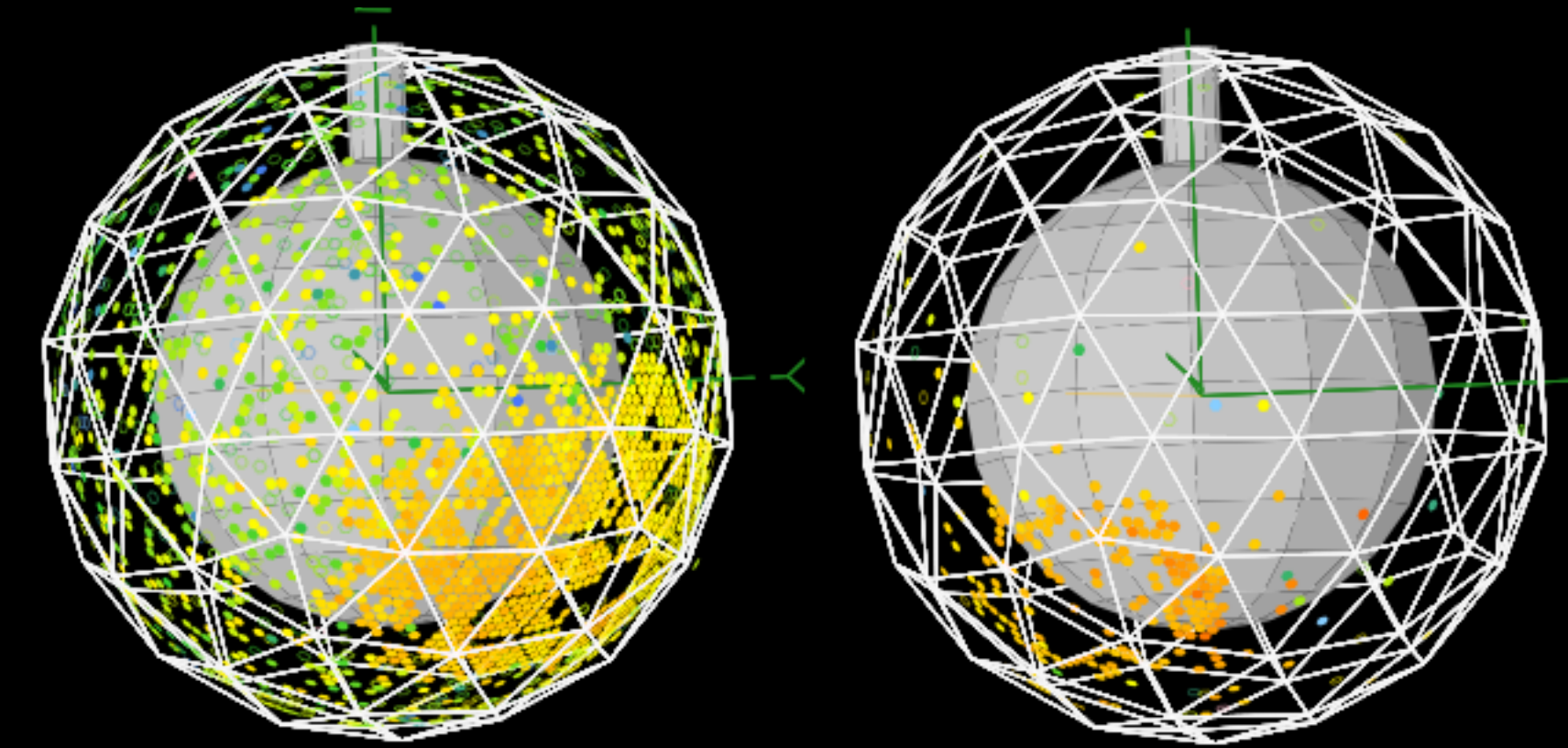


FIG. 4. Distribution of event directions with respect to solar direction for events with energy in the range 6.0–15.0 MeV.



Largest background to $0\nu\beta\beta$

Plus neutron capture, muon-induced neutrons (in progress)



PHYSICAL REVIEW D **99**, 032008 (2019)

Search for invisible modes of nucleon decay in water with the SNO+ detector

TABLE VI. Lifetime limits at 90% C.I. for the spectral and counting analysis, including statistical and systematic uncertainties alongside the existing limits.

	Spectral analysis	Counting analysis	Existing limits
n	2.5×10^{29} y	2.6×10^{29} y	5.8×10^{29} y [9]
p	3.6×10^{29} y	3.4×10^{29} y	2.1×10^{29} y [10]
pp	4.7×10^{28} y	4.1×10^{28} y	5.0×10^{25} y [11]
pn	2.6×10^{28} y	2.3×10^{28} y	2.1×10^{25} y [13]
nn	1.3×10^{28} y	0.6×10^{28} y	1.4×10^{30} y [9]

- Operating with **water** from 2017
- Transition to **scintillator** happening now

- Scintillator purification plant commissioned
- **LAB solvent** successfully distilled underground
- **PPO fluor** prep underway
- N₂/steam stripping tested



- Invisible nucleon decay
- Solar neutrinos
- Supernova neutrinos
- **Reactor antineutrinos (Δm^2_{12})**
- **Geo-neutrinos**



- Operating with **water** from 2017
- Transition to **scintillator** happening now
- **Tellurium** loading for $\beta\beta$ due in 2019-20 (1330 kg ^{130}Te)

- Invisible nucleon decay
- Solar neutrinos
- Supernova neutrinos
- Reactor neutrinos (Δm^2_{12})
- Geo-neutrinos
- **Neutrinoless double-beta decay**

Te needed for Phase I **all underground**

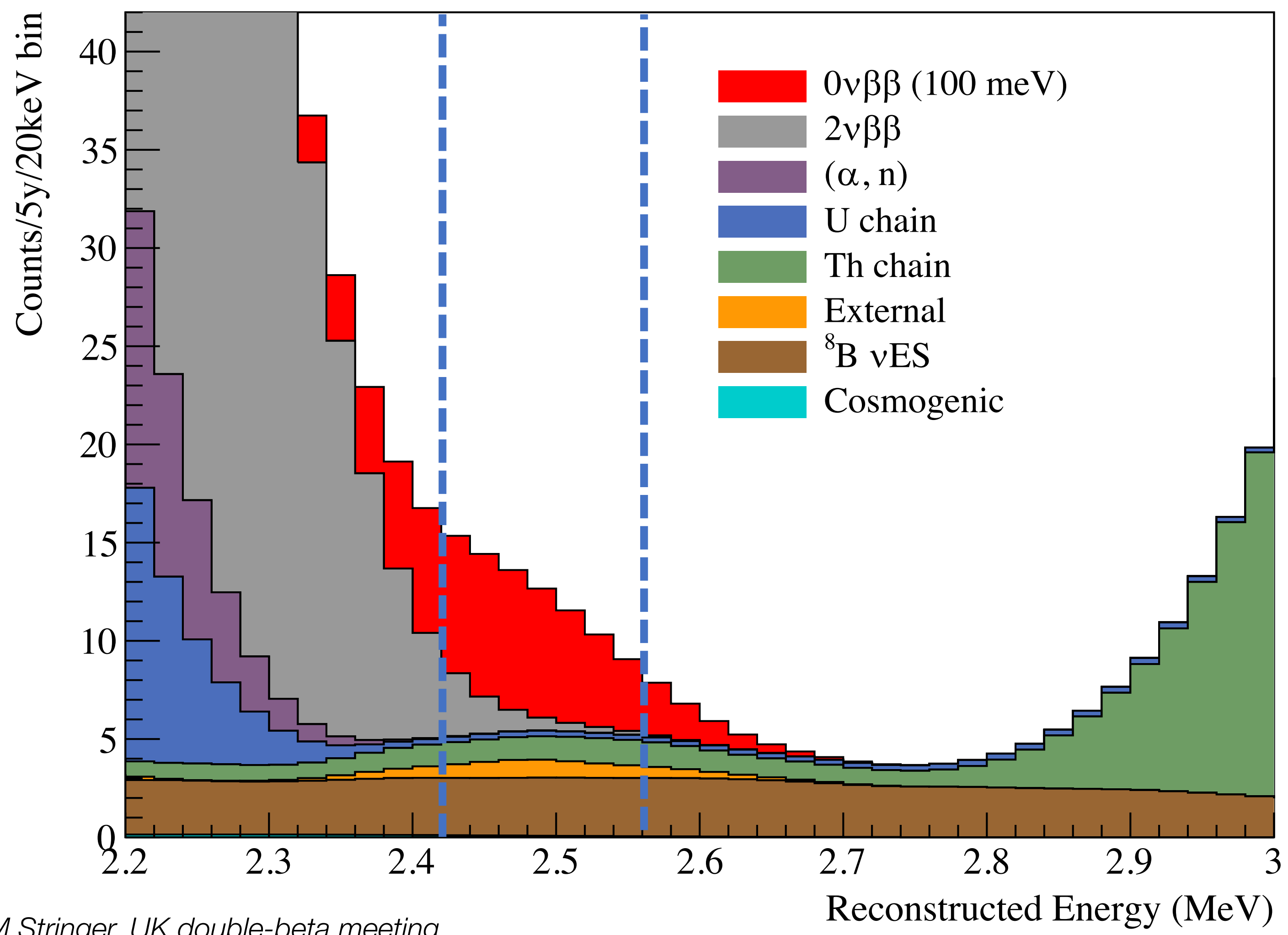


Te **purification system** almost complete



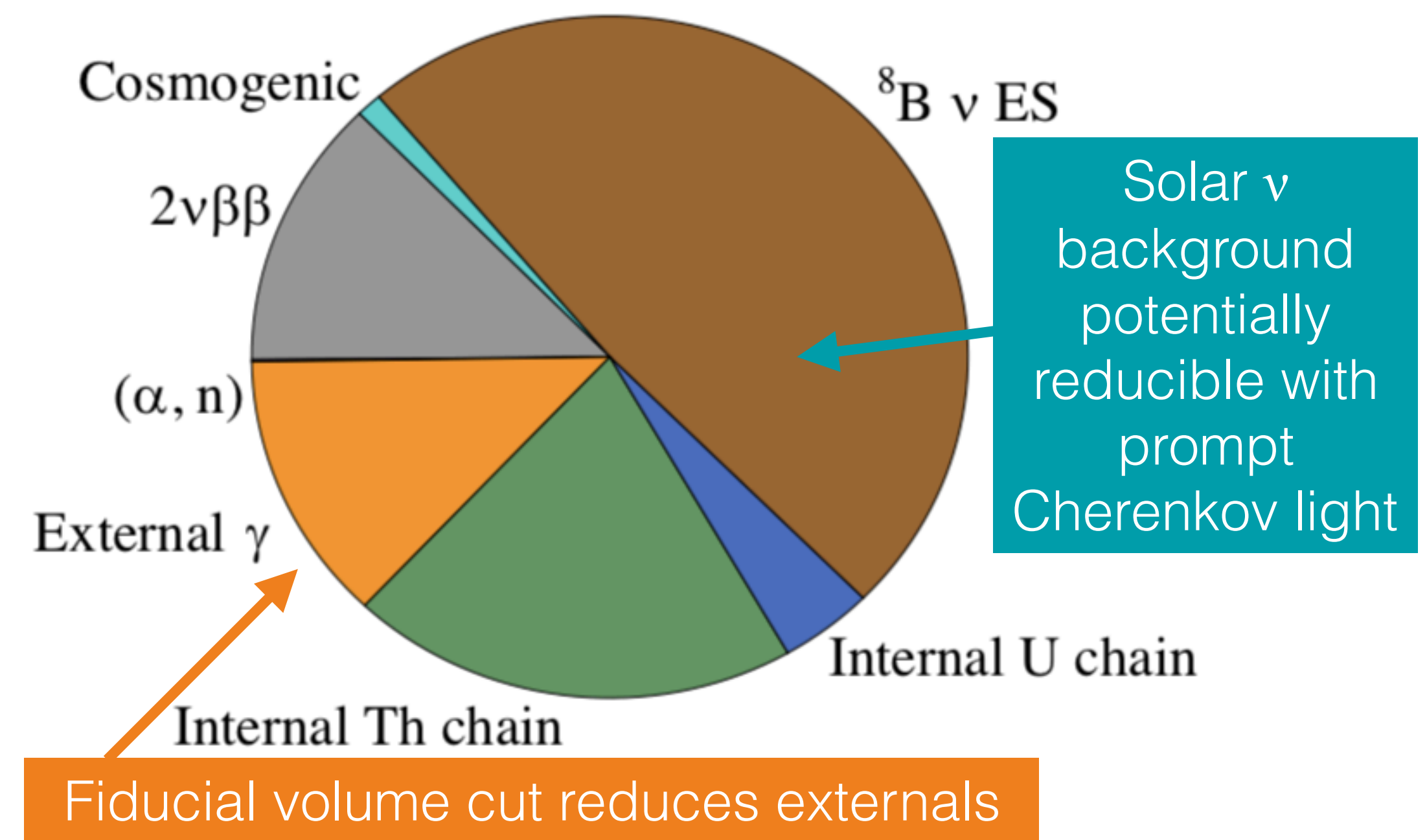
Te-diol synthesis plant construction is well advanced (synthesised from telluric acid)





M Stringer, UK double-beta meeting

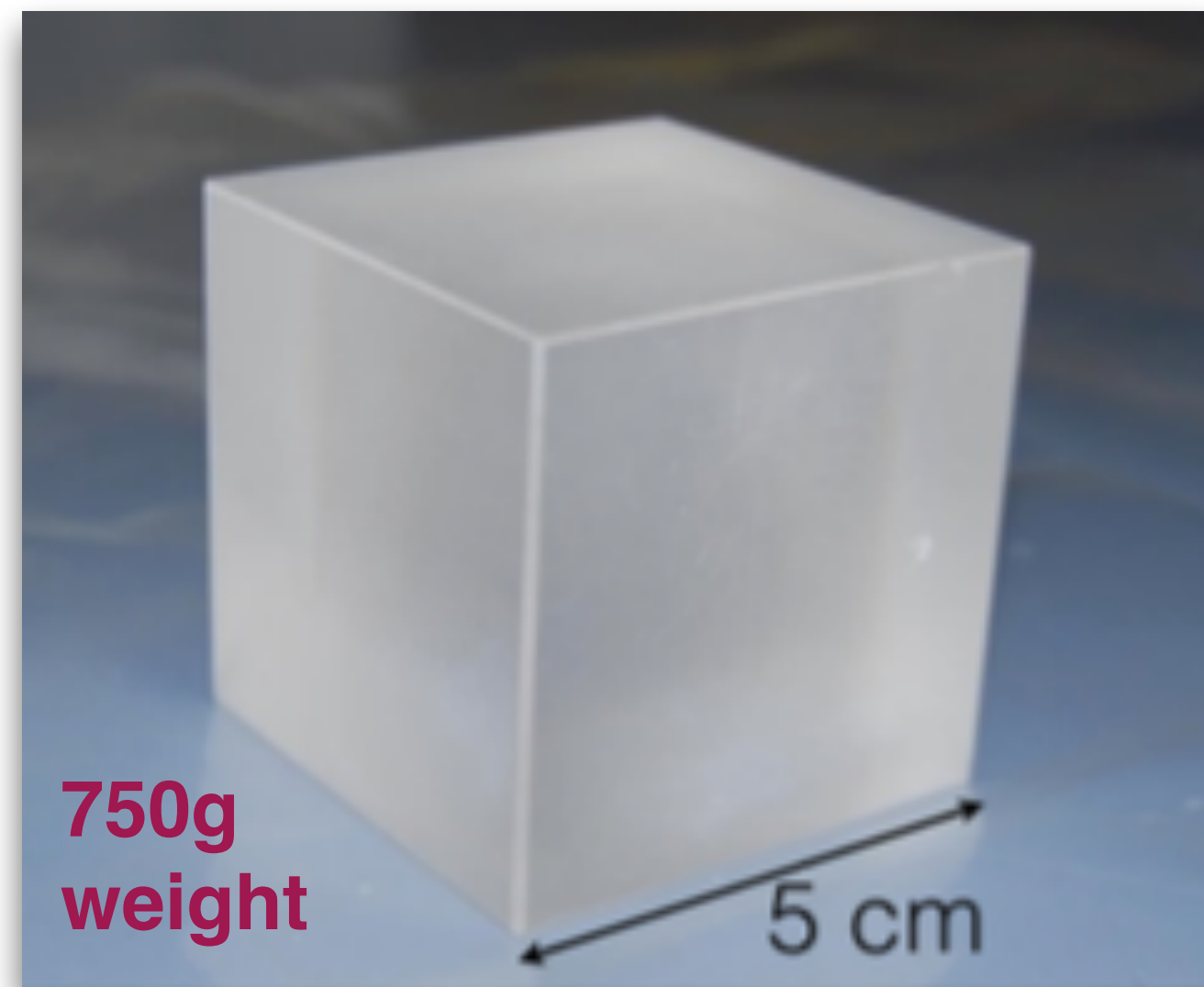
ROI: 2.42 - 2.56 MeV [-0.5σ - 1.5σ]
 Counts/Year: 9.47



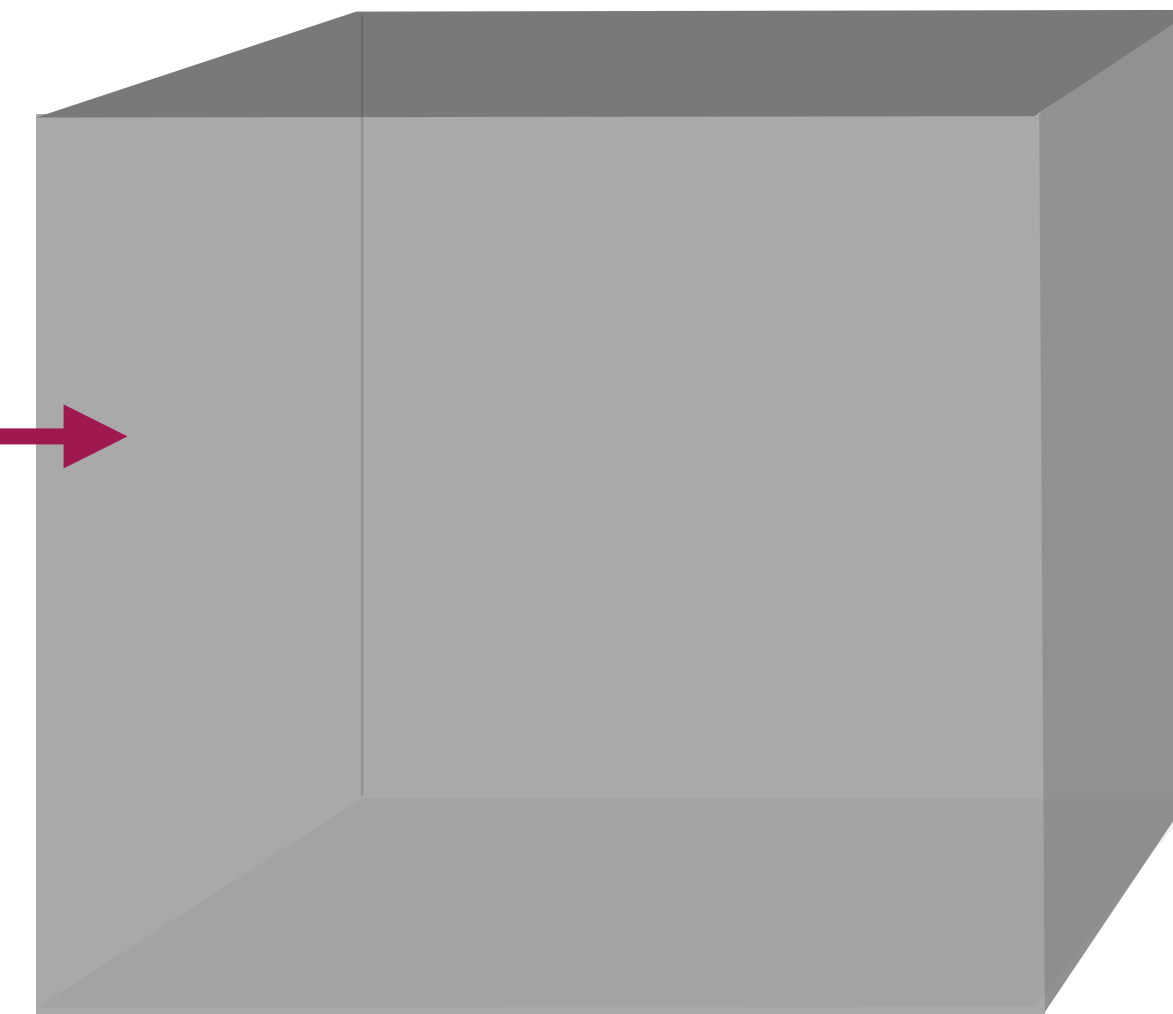
$T_{1/2} > 2.1 \times 10^{26}$ years
 Predicted for 5 years of running



988 natural TeO₂ crystals (206 kg of ¹³⁰Te)



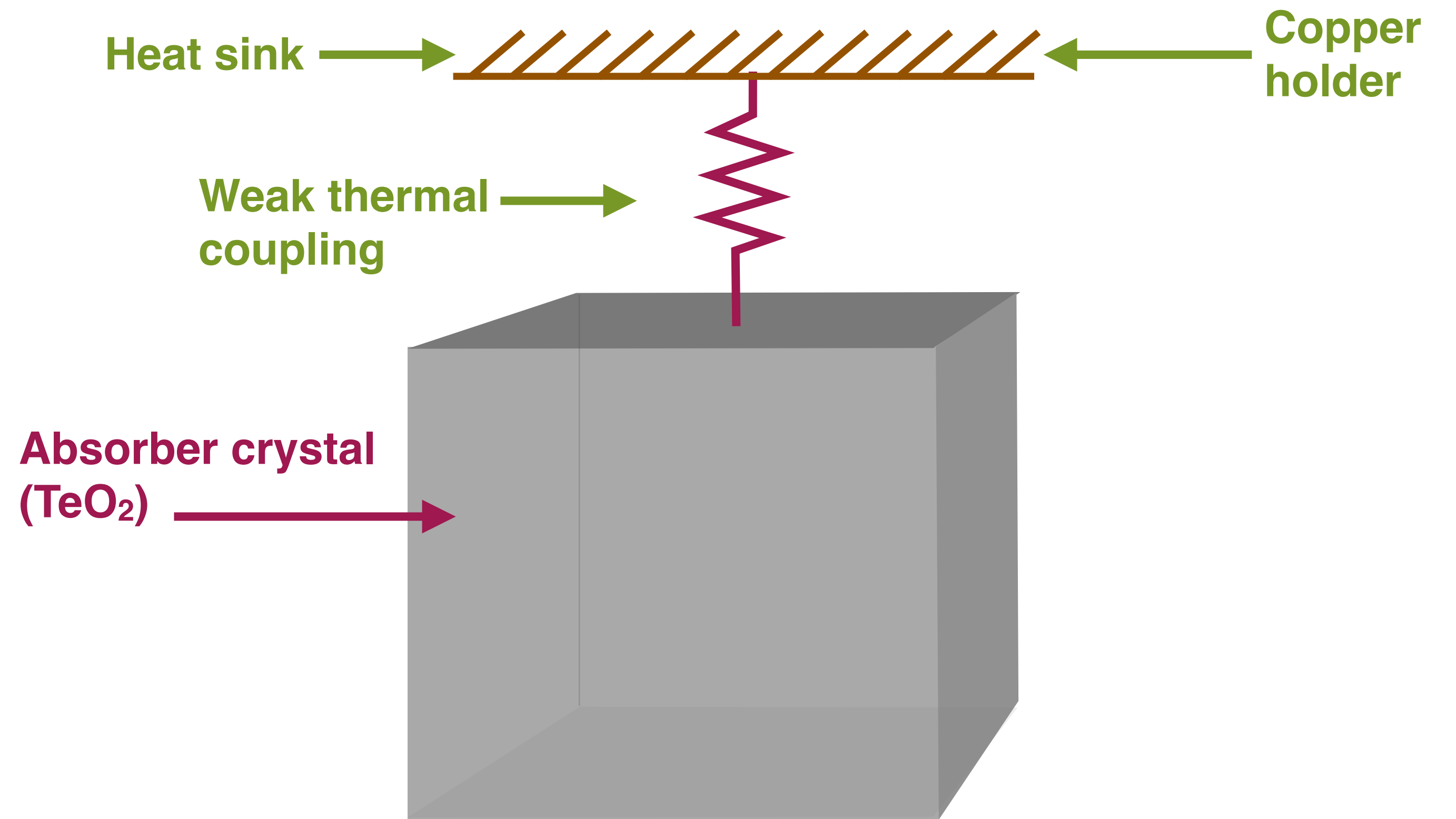
Absorber crystal
(TeO₂)



- Crystal is **source and detector** (solid-state bolometer)
- Future plans to use **other isotopes** (⁸²Se)

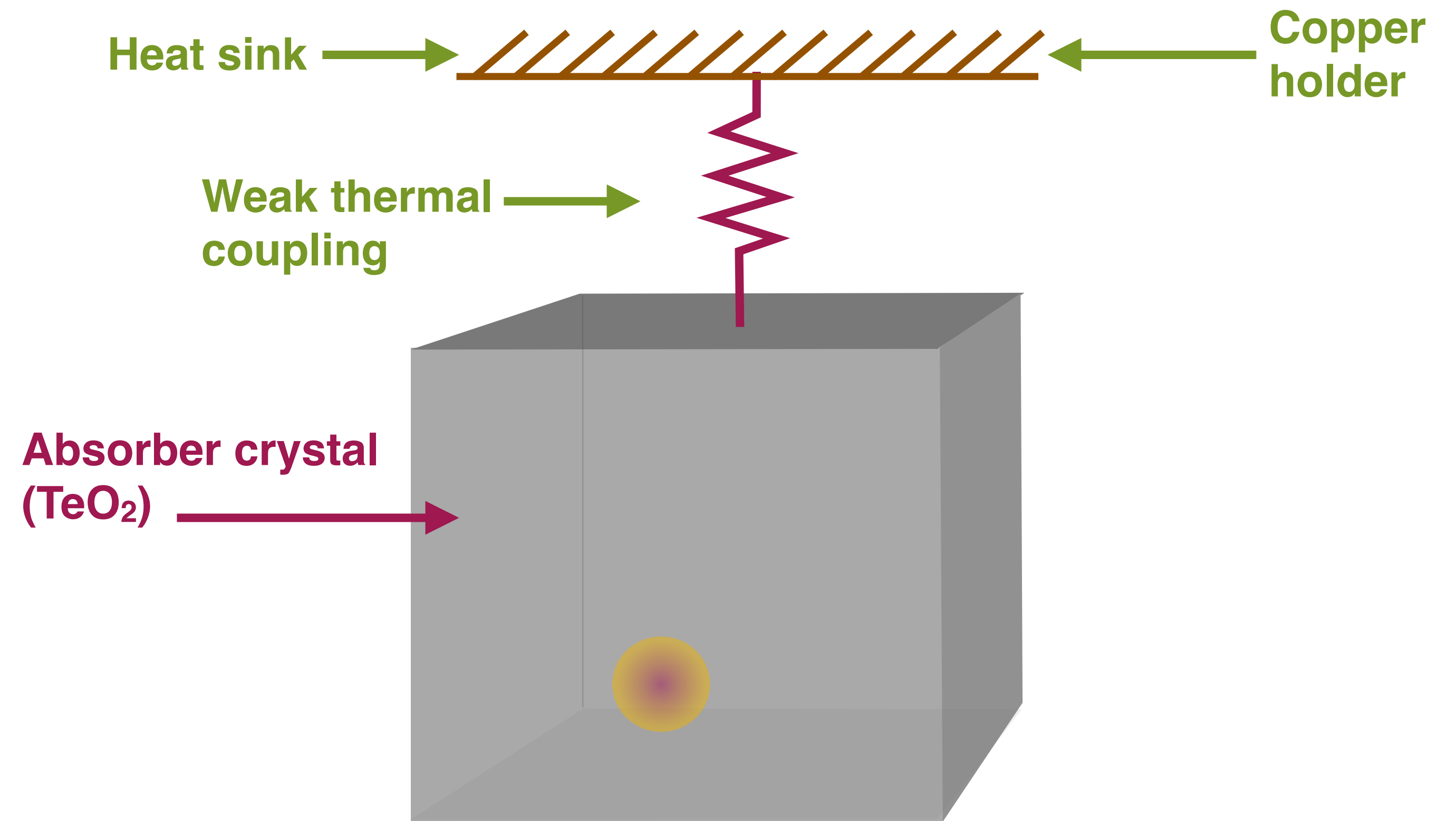


- Crystals at low temperature: **10 mK**



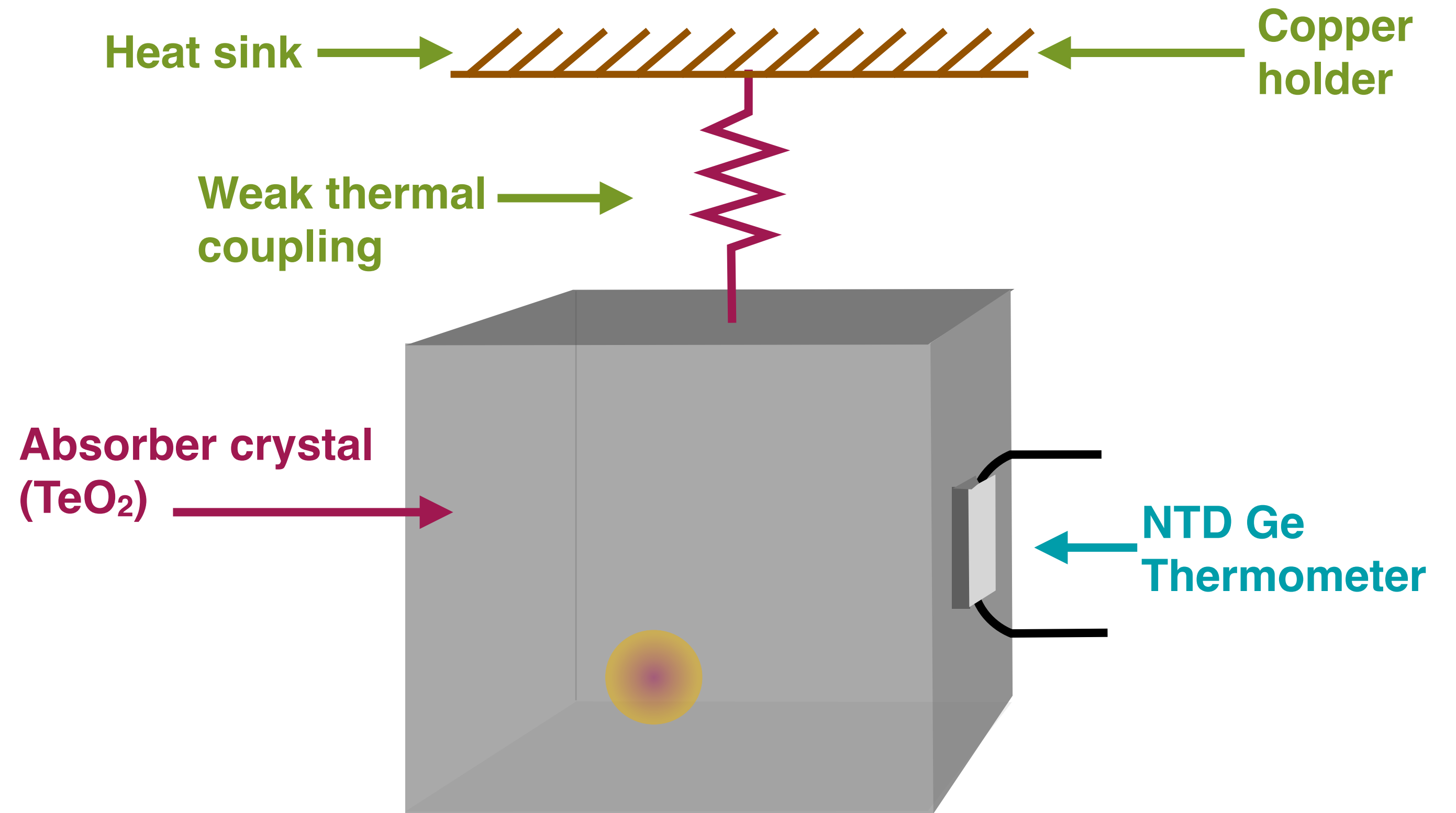
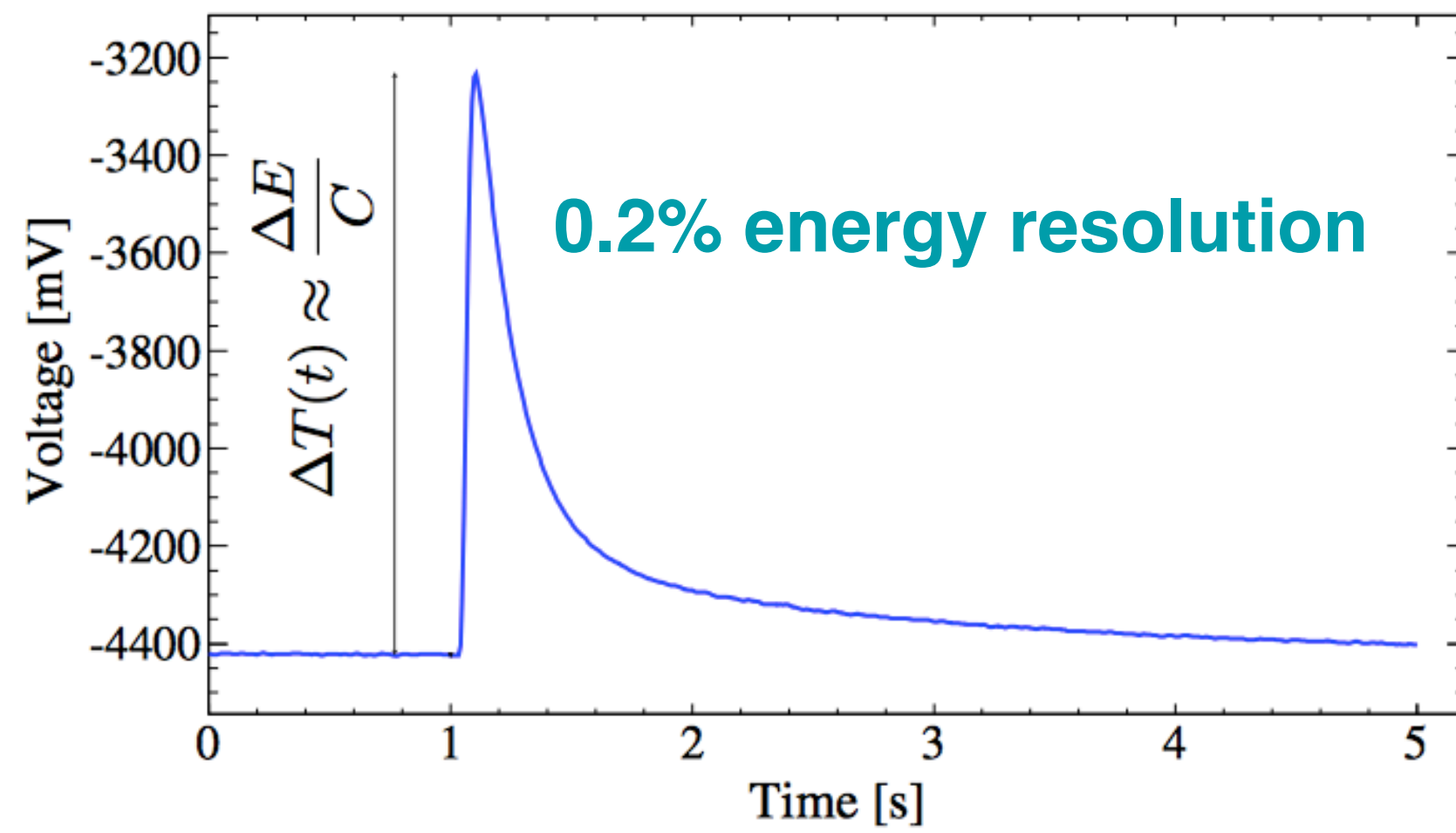


- Crystals at low temperature: **10 mK**
- Energy deposition (e.g. $\beta\beta$ decay) **raises temperature** by \sim **100 μ K / MeV**



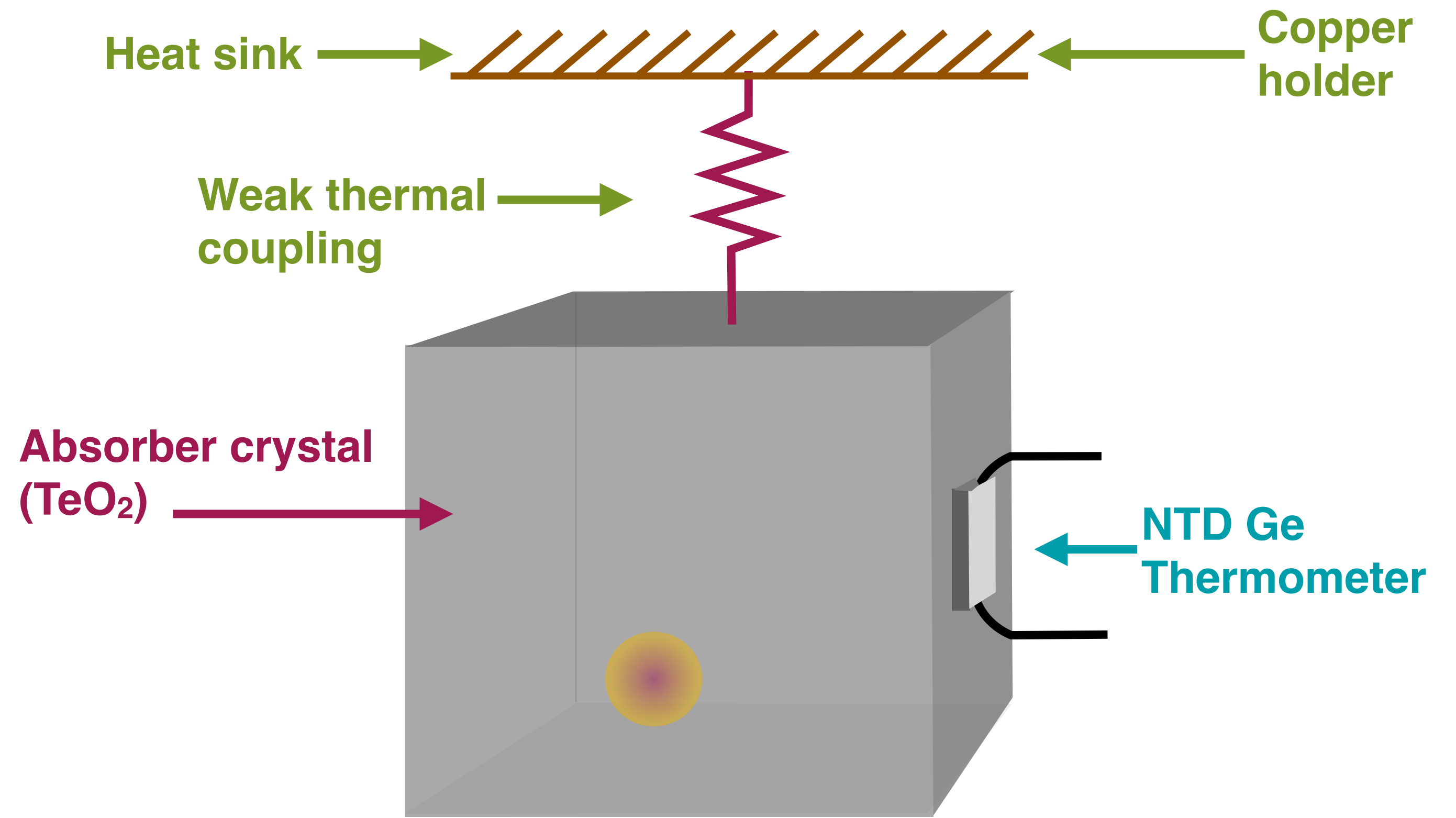
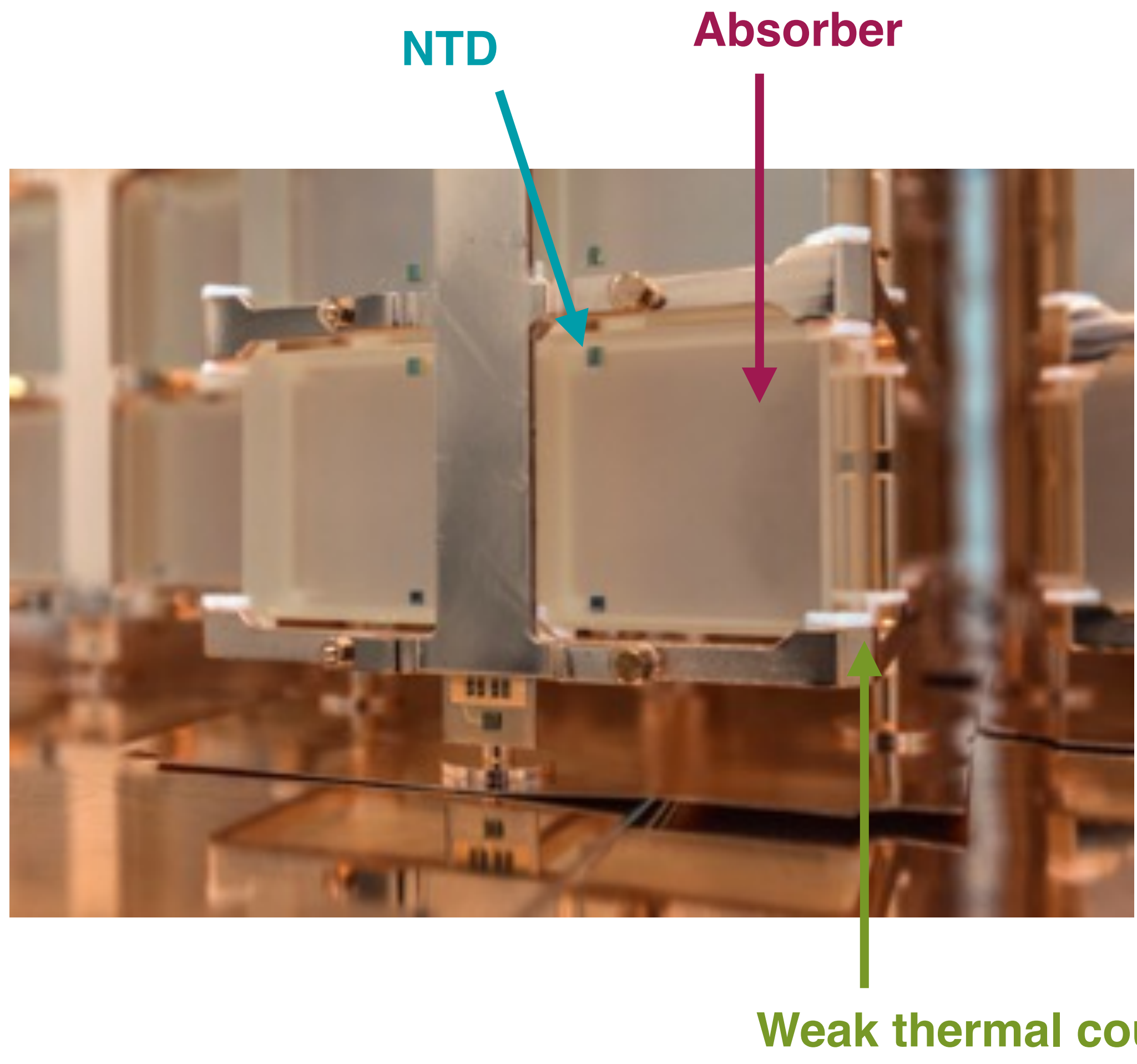


- Crystals at low temperature: **10 mK**
- Energy deposition (e.g. $\beta\beta$ decay) **raises temperature** by $\sim 100 \mu\text{K} / \text{MeV}$
- Read out by a germanium neutron transmutation doped **thermometer**





CUORE: TeO₂ bolometers at LNGS, Italy

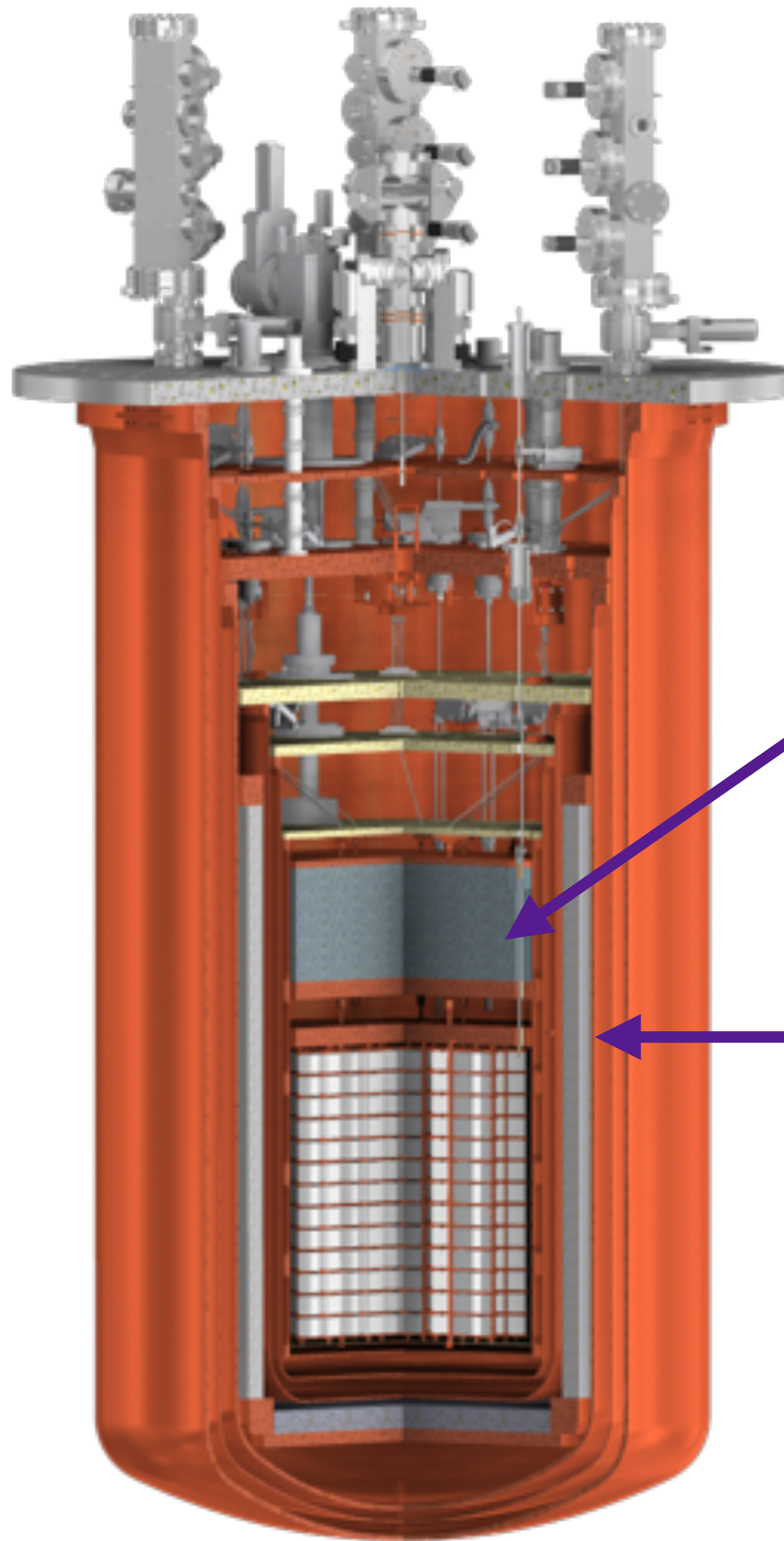


Cooling

- Detectors cooled to 10 mK
- Long-term **stable** base temperature **6.3 mK**



Coldest cubic metre in the known universe



Shielding

- Radiopure copper
- 70 tons of lead including 4 of **Roman lead**



Cooling

Shielding

Shielding

- Detectors
- Long-term
mK



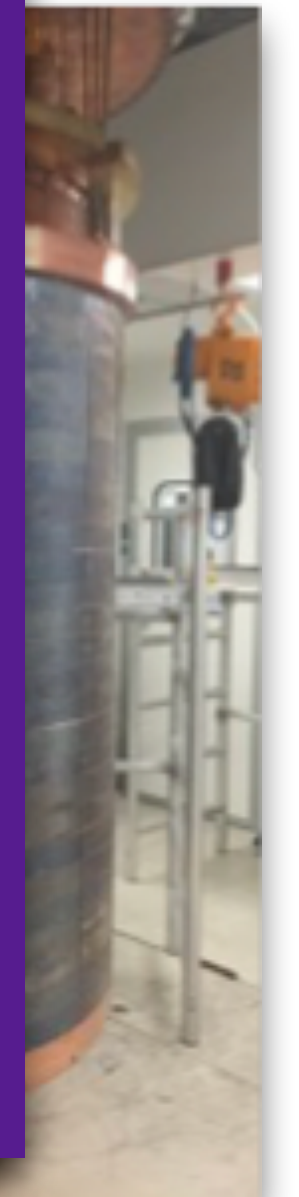
- Found off Sardinia coast in 1988
- Shipwreck from 80-50 B.C.E
- 120 lead ingots have lost their natural radioactivity

in lead



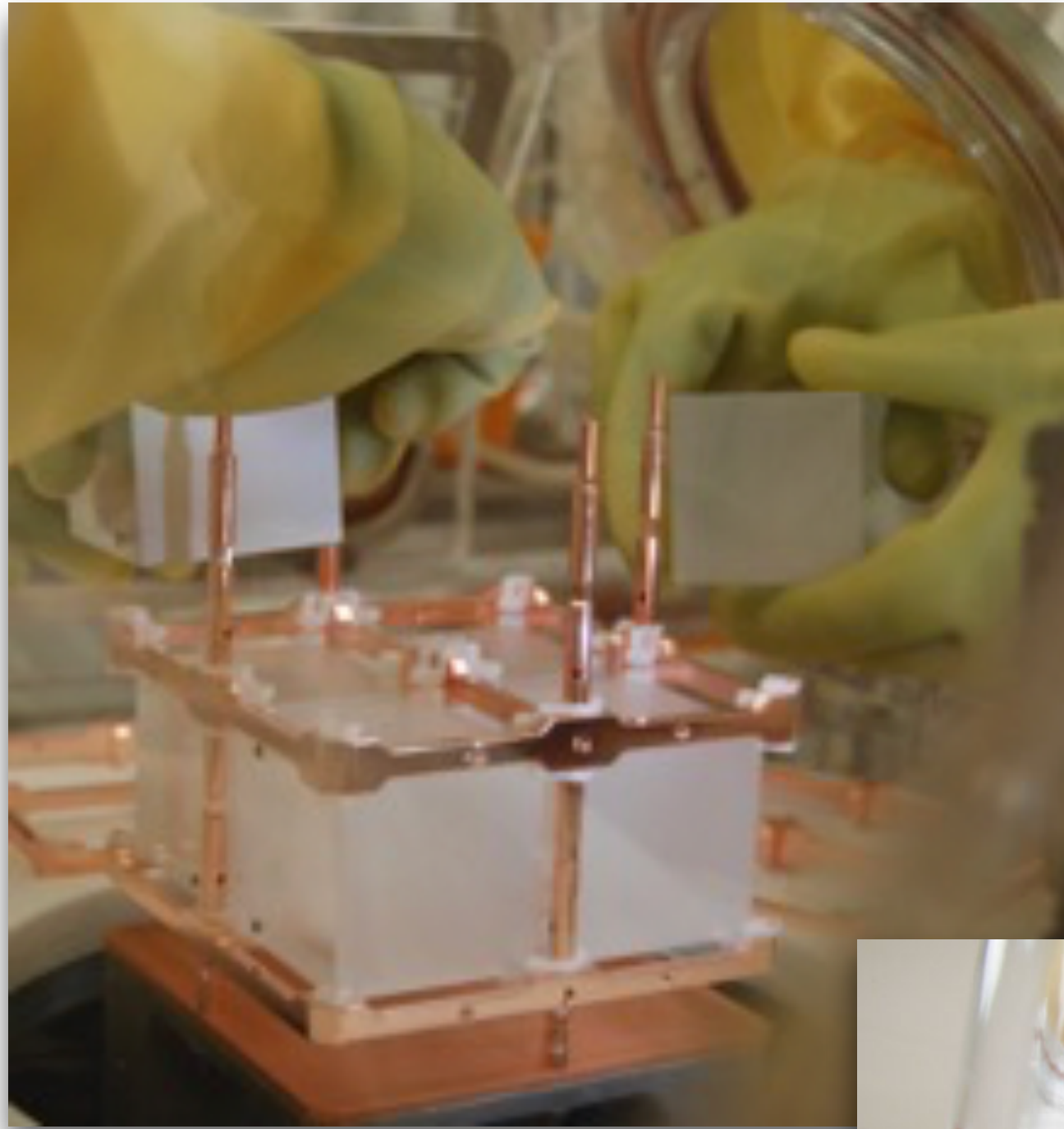
- Ingots originated in Spain
- Used for ammunition
- Inscriptions preserved for archaeology

<https://www.nature.com/news/2010/100415/full/news.2010.186.html>



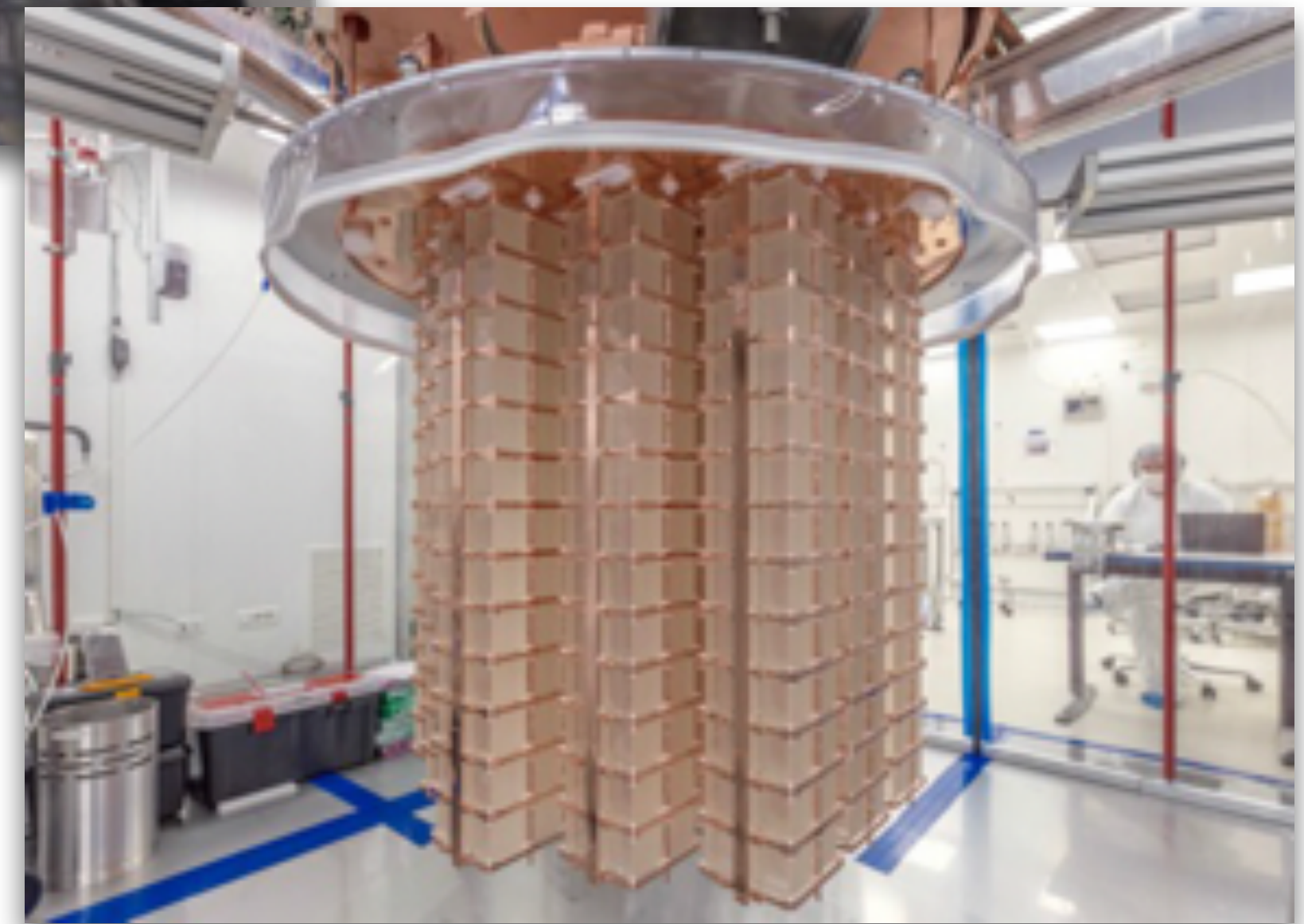
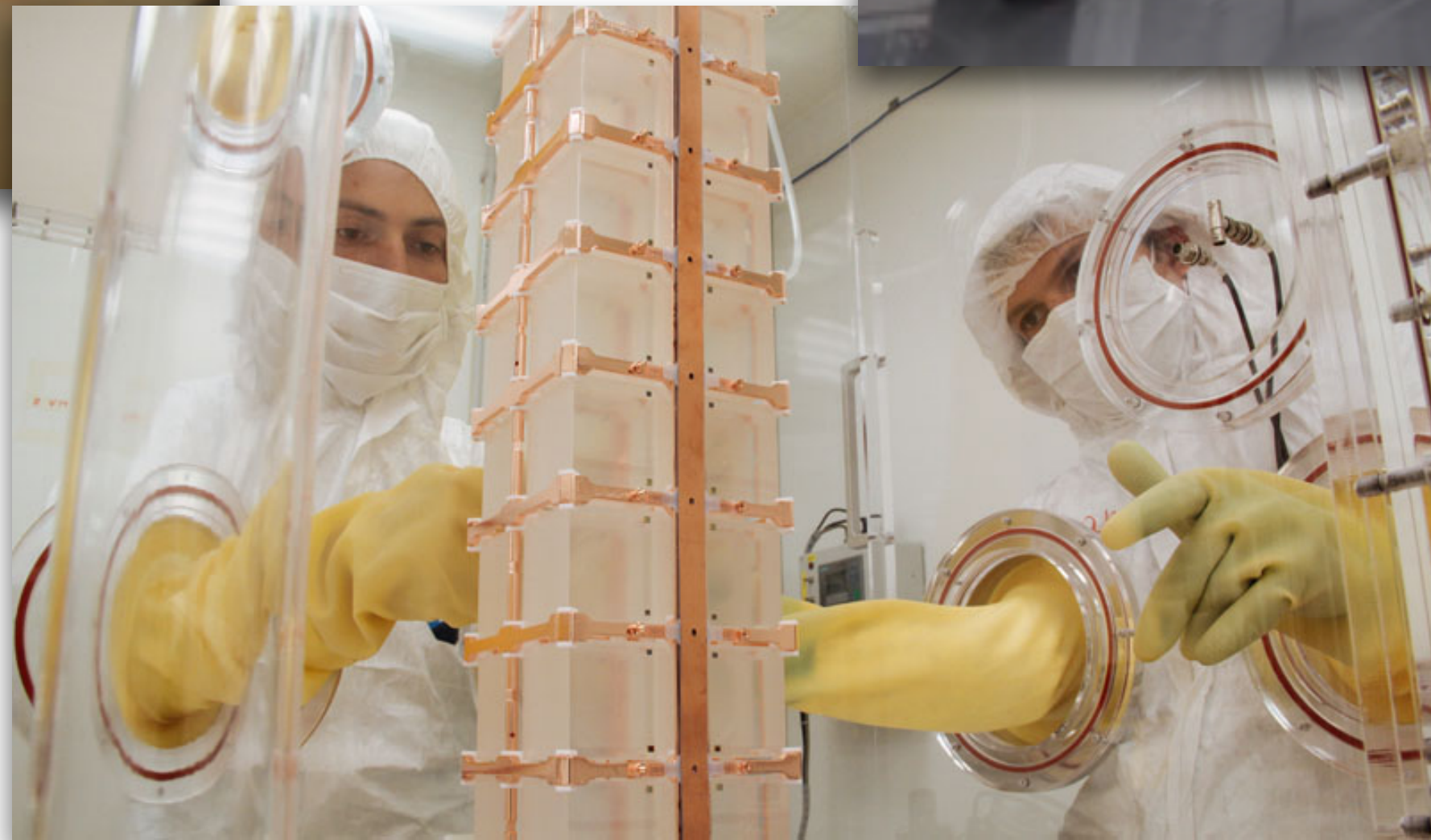


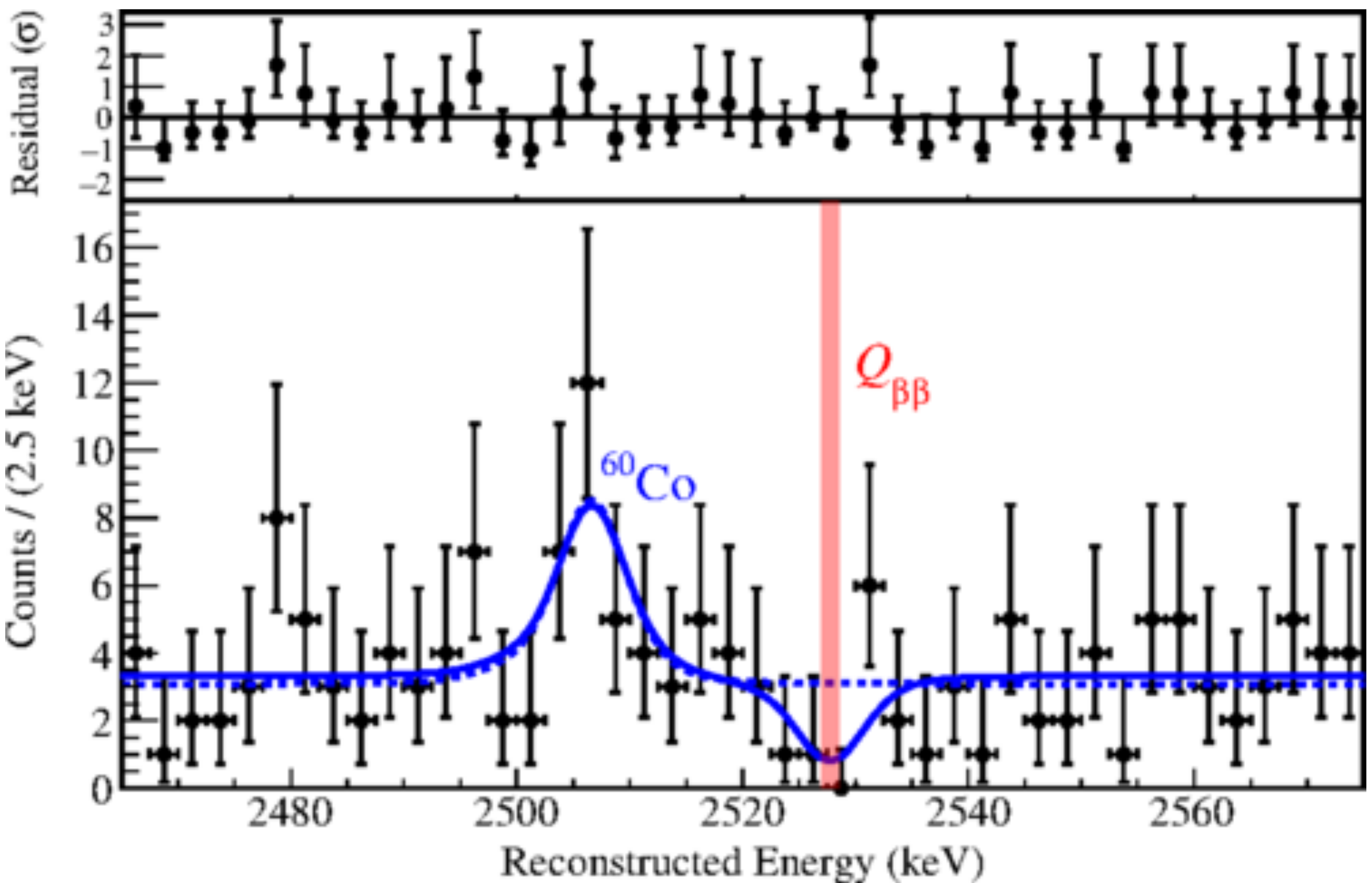
J Ouellet, Neutrino 2018



Installing the detector tower

Assembling in N₂ environment

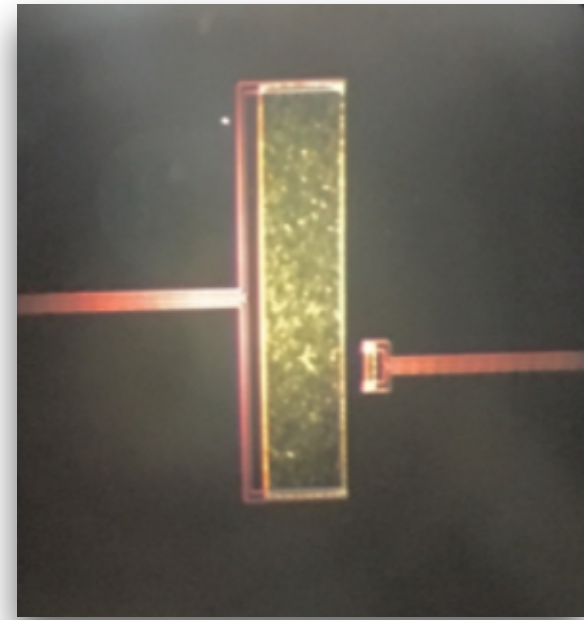
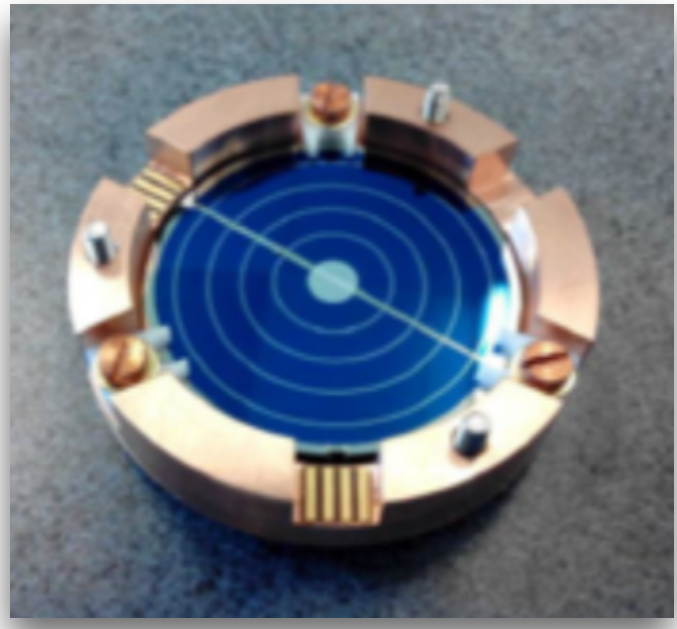




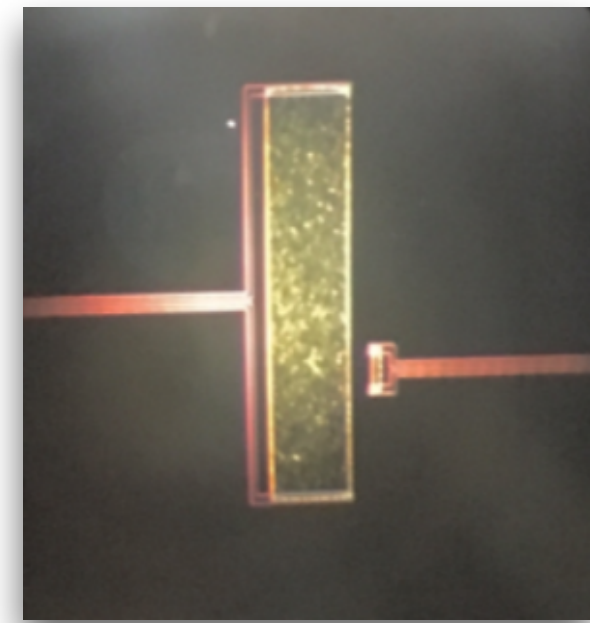
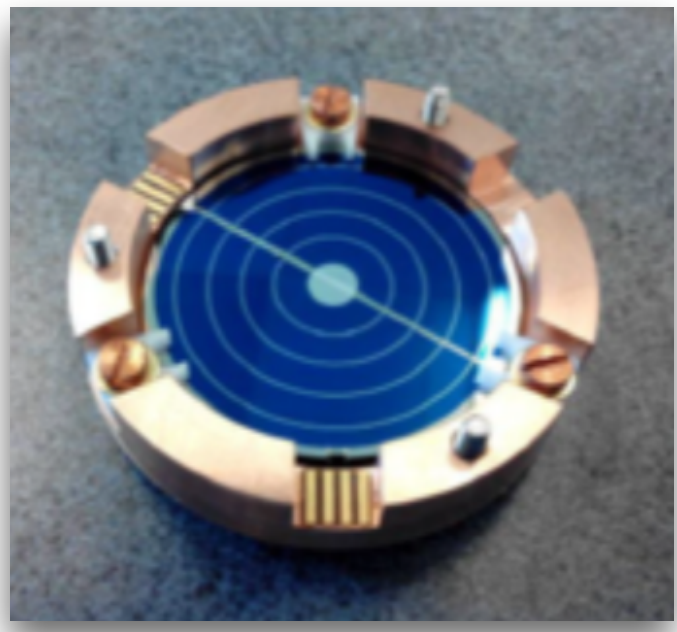
Phys. Rev. Lett. **120**,
132501 (2018)

$T_{1/2} > 1.5 \times 10^{25}$ years
 ($\langle m_{\beta\beta} \rangle < 110\text{-}520$ meV)
 86.3 kg.years
 5 years: 9.5×10^{25} years

Eur. Phys. J. C **77**, 13 (2017)

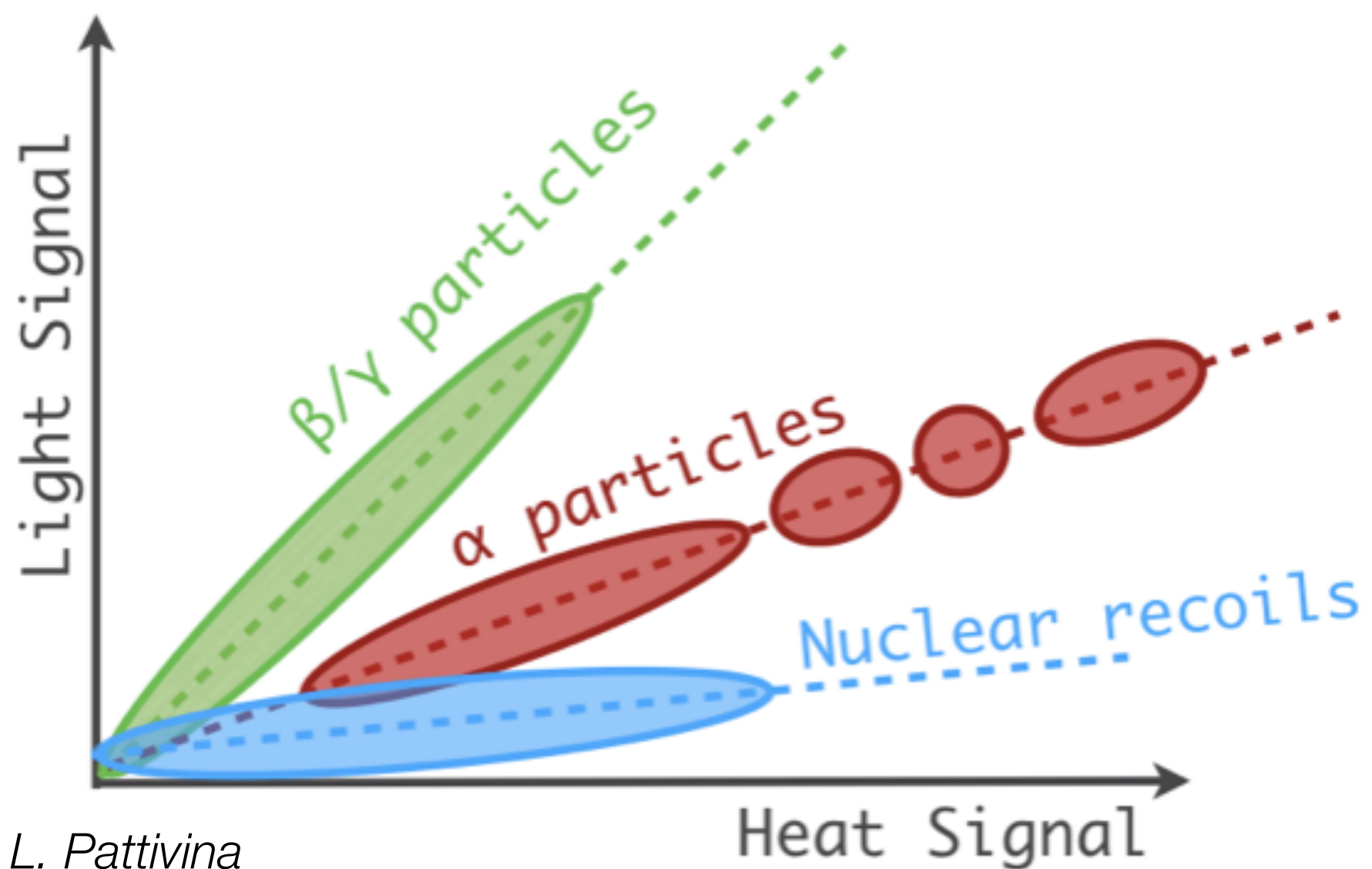


Light-emitting bolometers allow particle ID

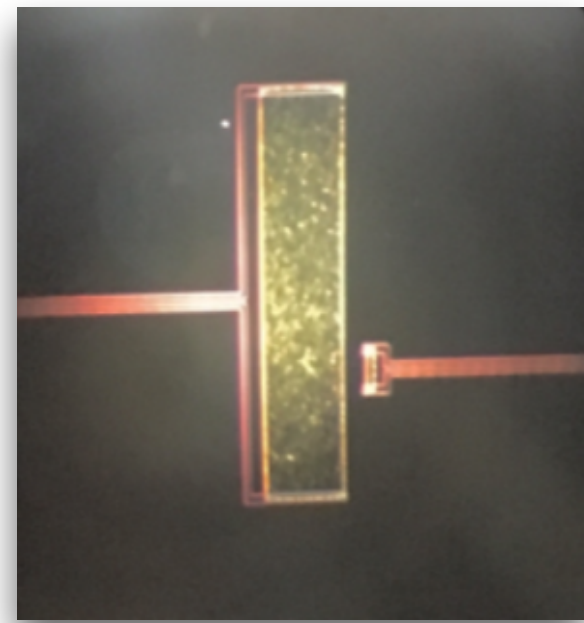
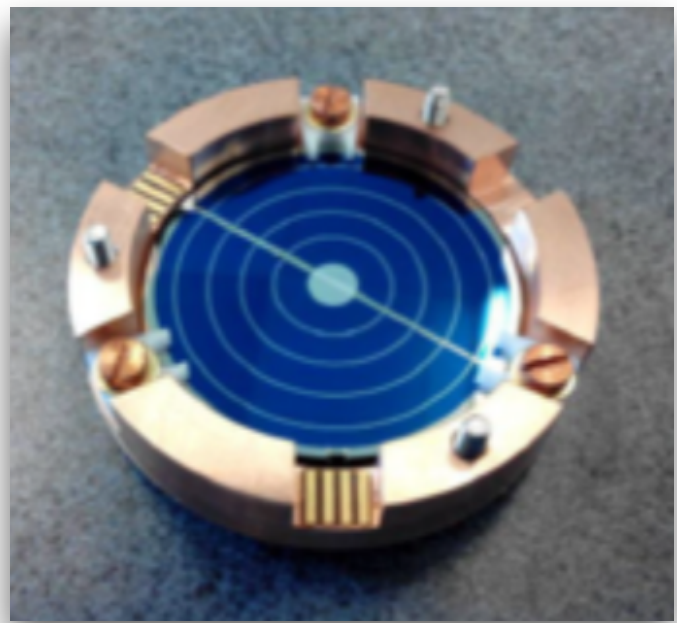


Light-emitting bolometers allow particle ID

Light shape parameter separates **α events** from β and γ decays

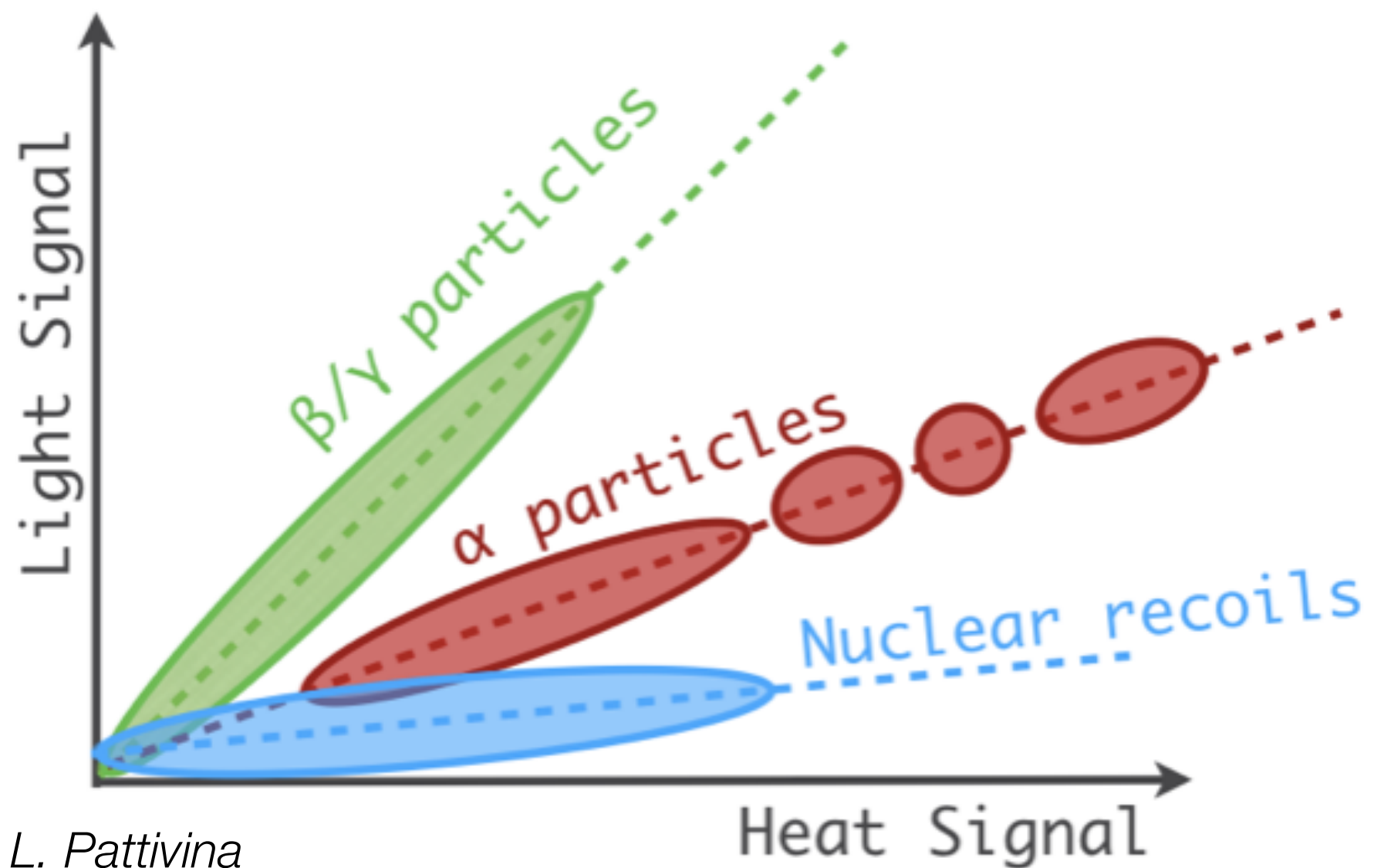


L. Pattivina

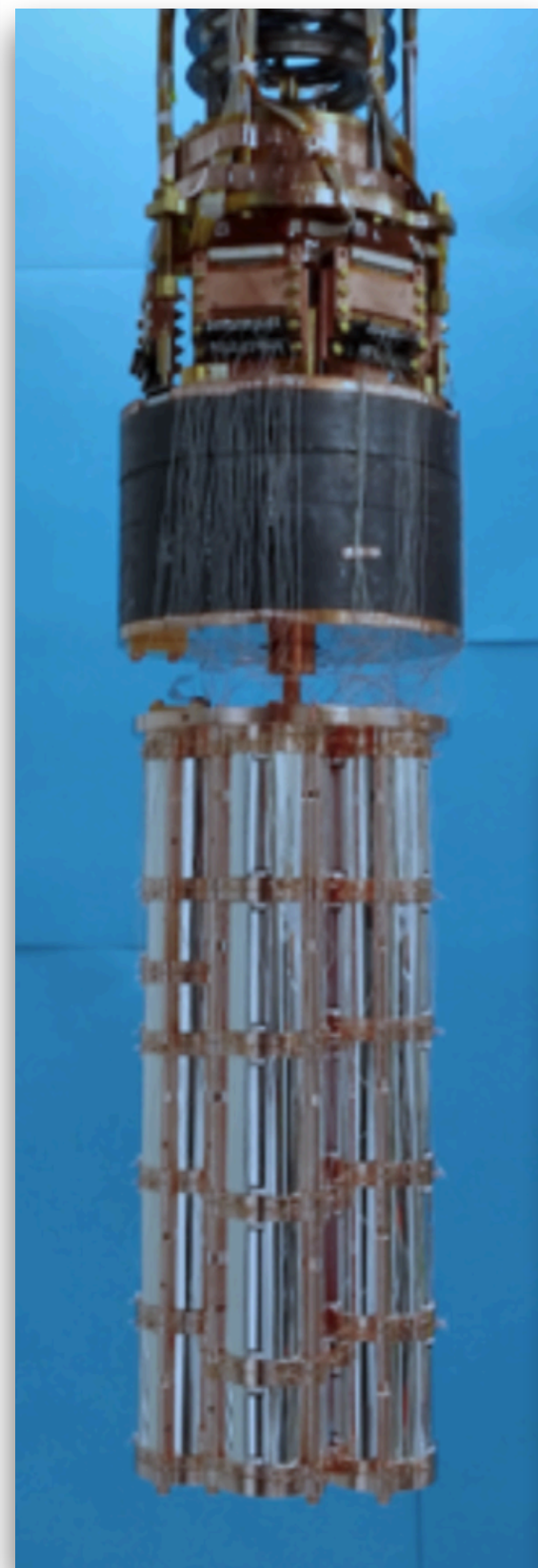


Light-emitting bolometers allow particle ID

Light shape parameter separates **α events** from β and γ decays

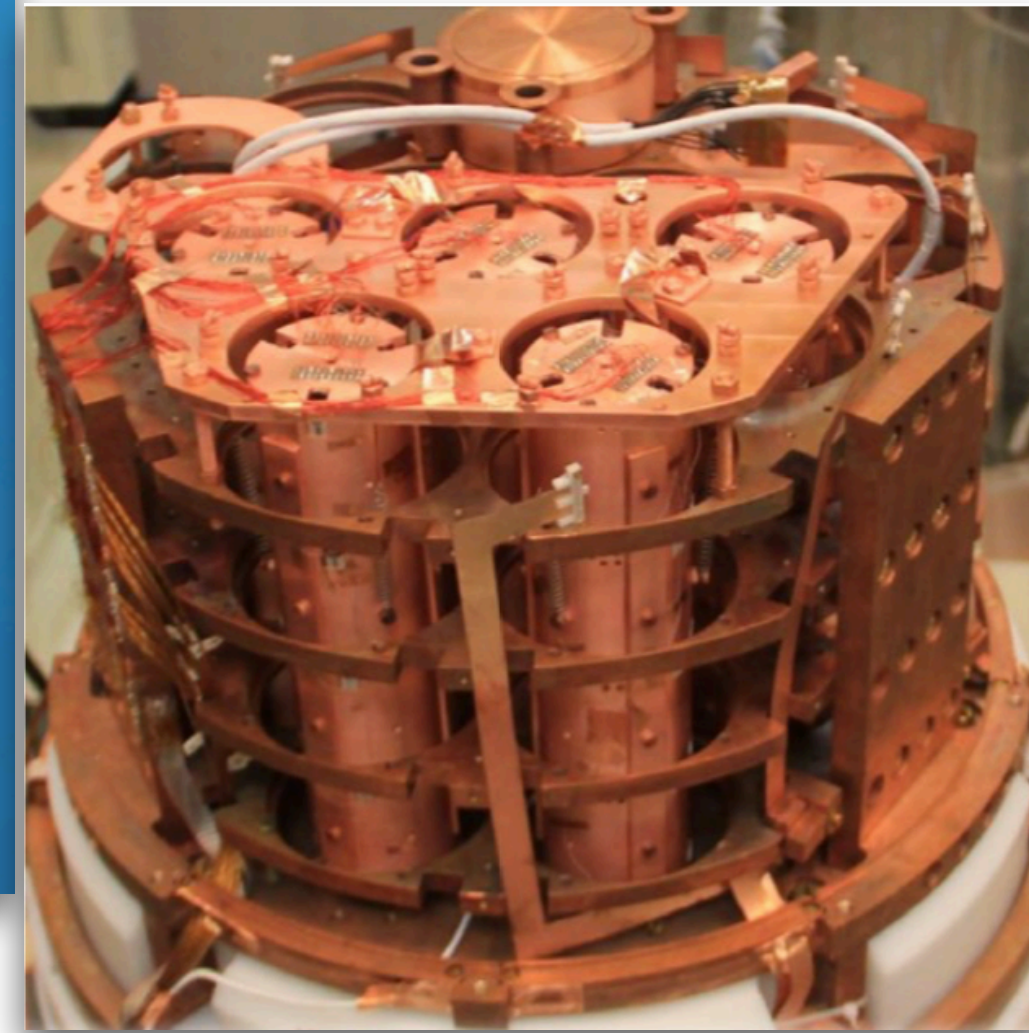


L. Pattivina

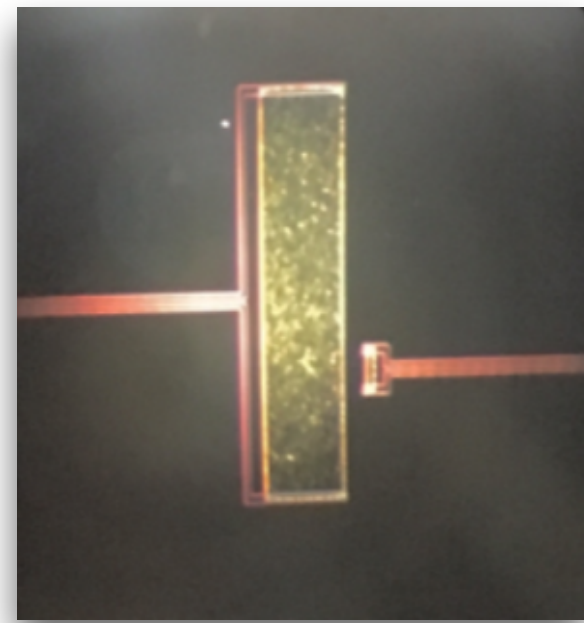
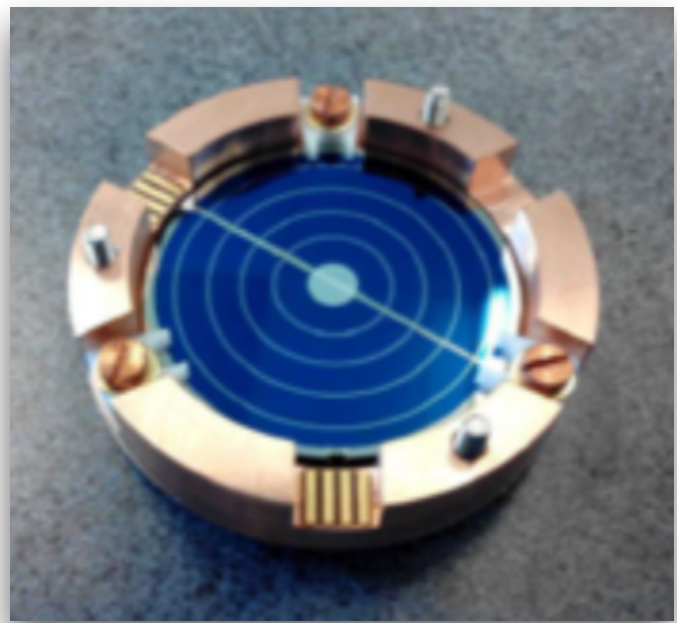


CUPID-0:
Zn⁸²Se
bolometers (5.28 kg ⁸²Se)

CUPID-Mo uses
Li₂¹⁰⁰MoO₄ crystals

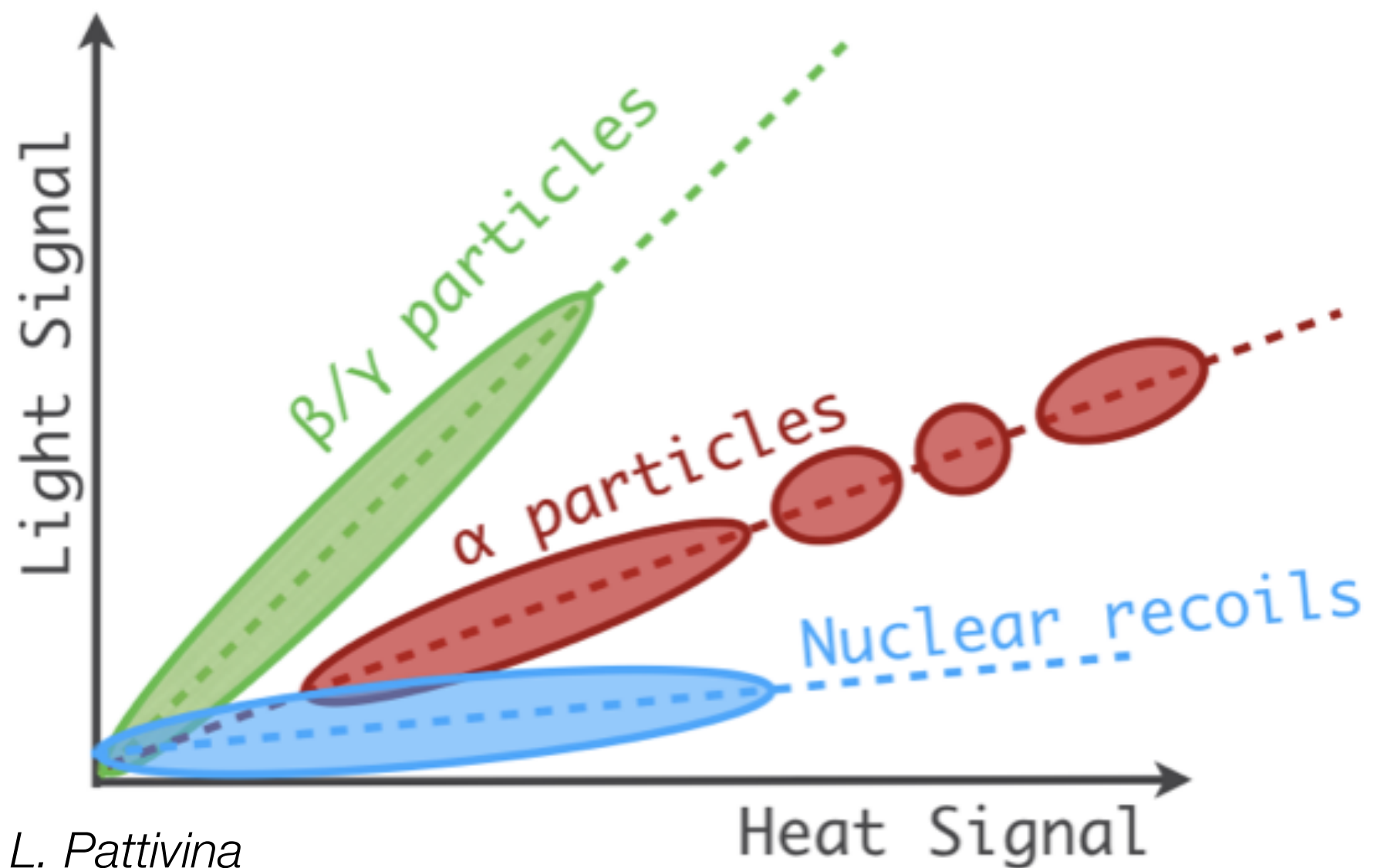


TeO₂ scintillating bolometers under R&D



Light-emitting bolometers allow particle ID

Light shape parameter separates α events from β and γ decays

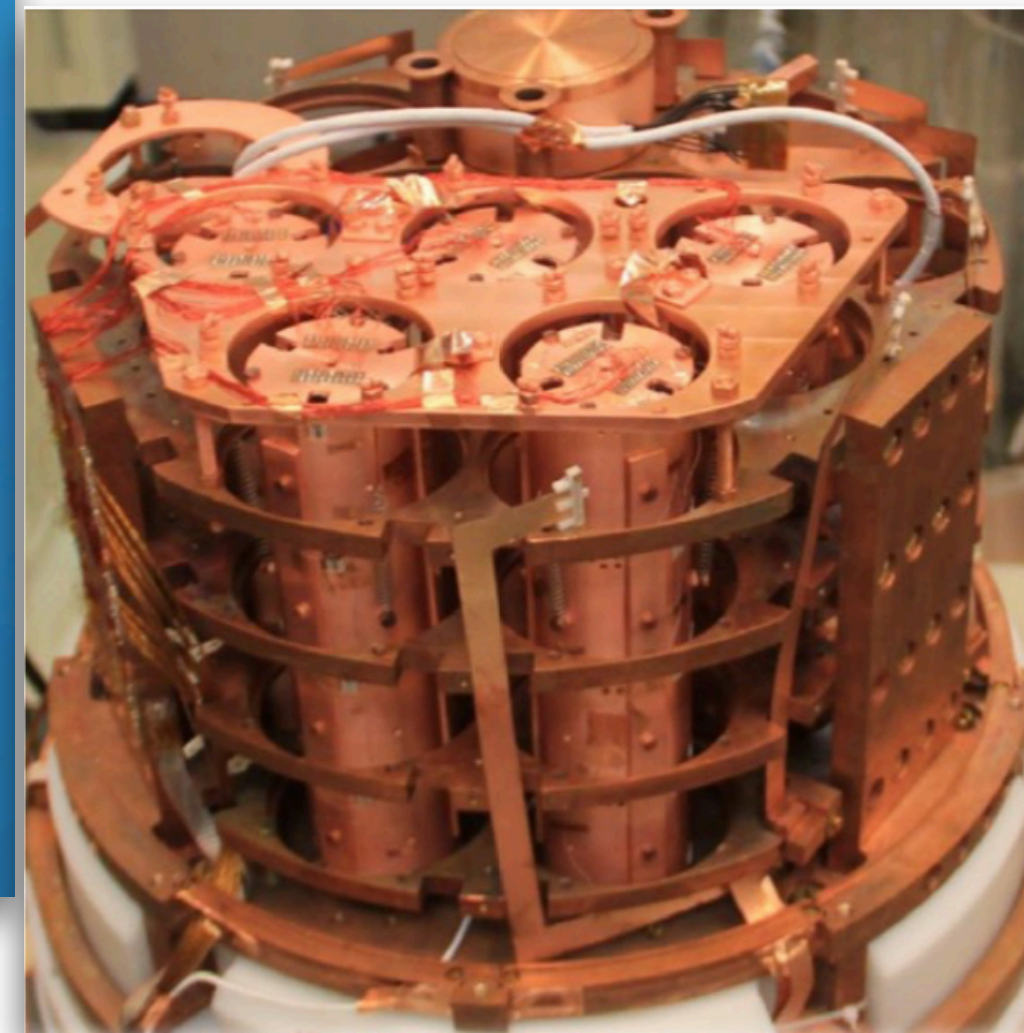


L. Pattivina

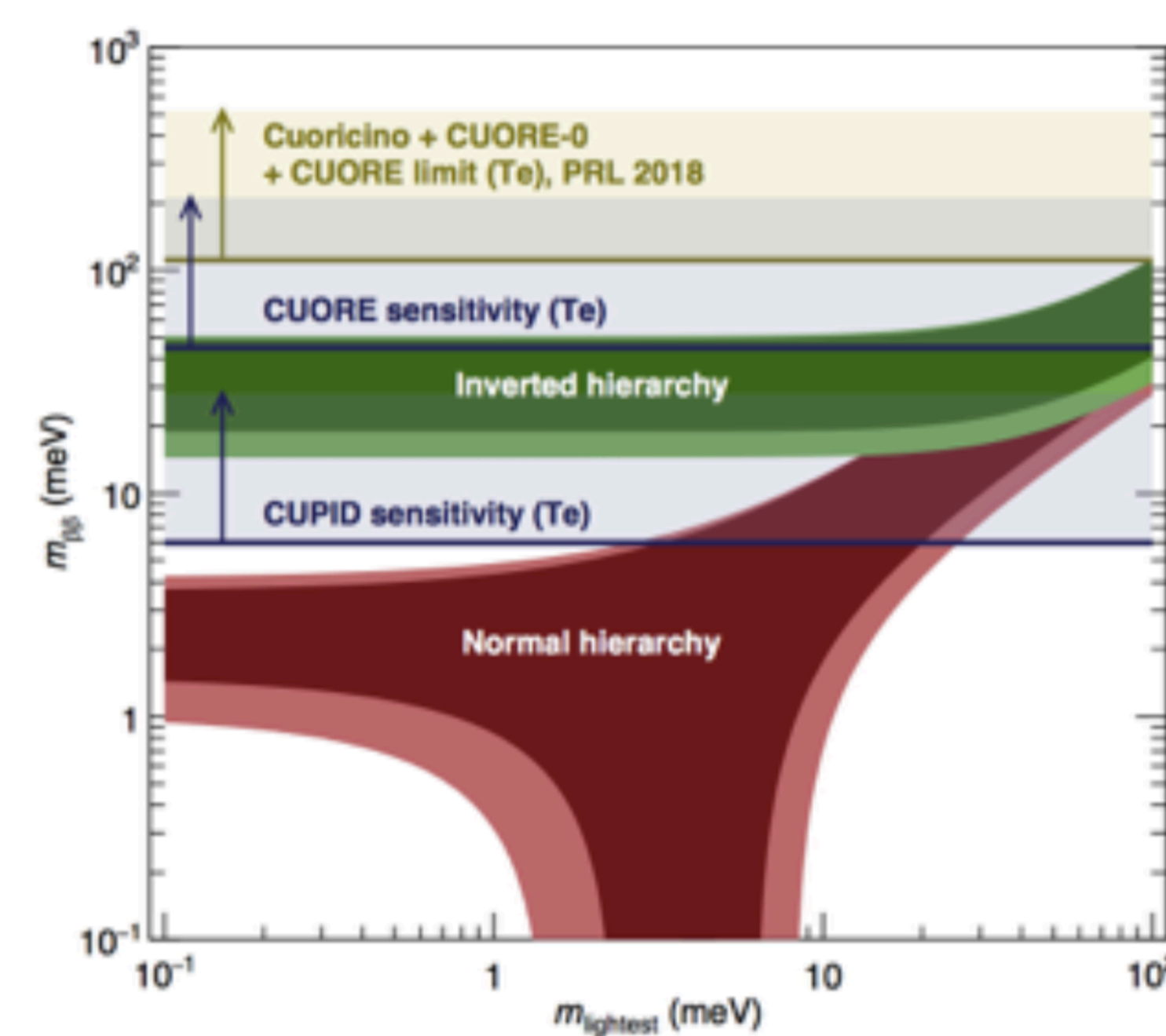


CUPID-0:
 $Zn^{82}Se$
bolometers (5.28 kg ^{82}Se)

CUPID-Mo uses $Li_2^{100}MoO_4$ crystals



TeO_2 scintillating bolometers under R&D

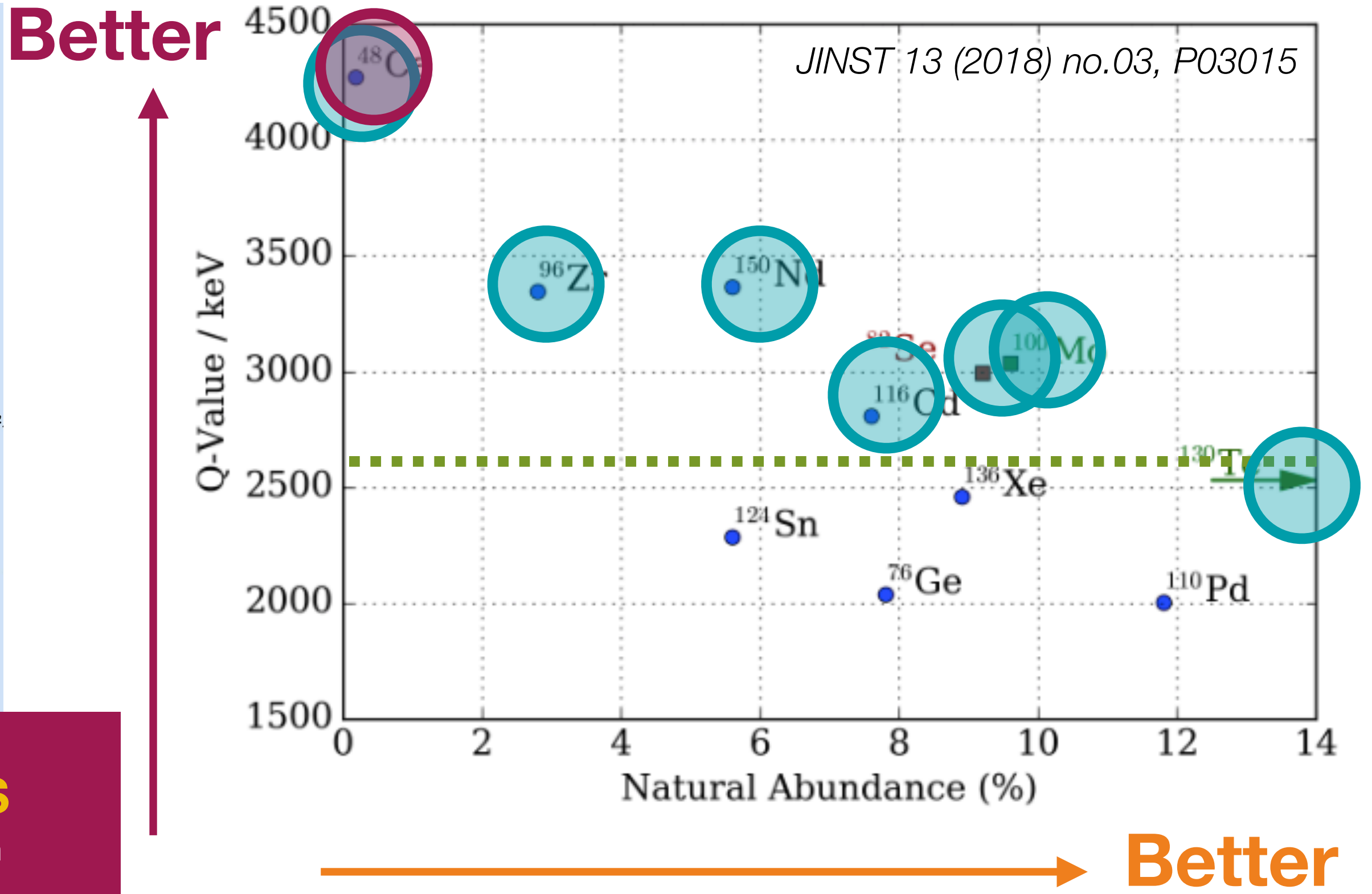


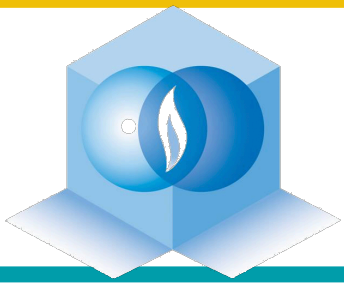
Target CUPID sensitivity covers whole **inverted hierarchy** ($m_{\beta\beta} \sim 6 - 20$ meV, $T_{1/2} \sim 10^{27}$ yr)



**SuperNEMO/
NEMO-3**
Modane, France

CANDLES
Kamioka, Japan





Highest Q-value

- low backgrounds
- big phase space factor = high rate

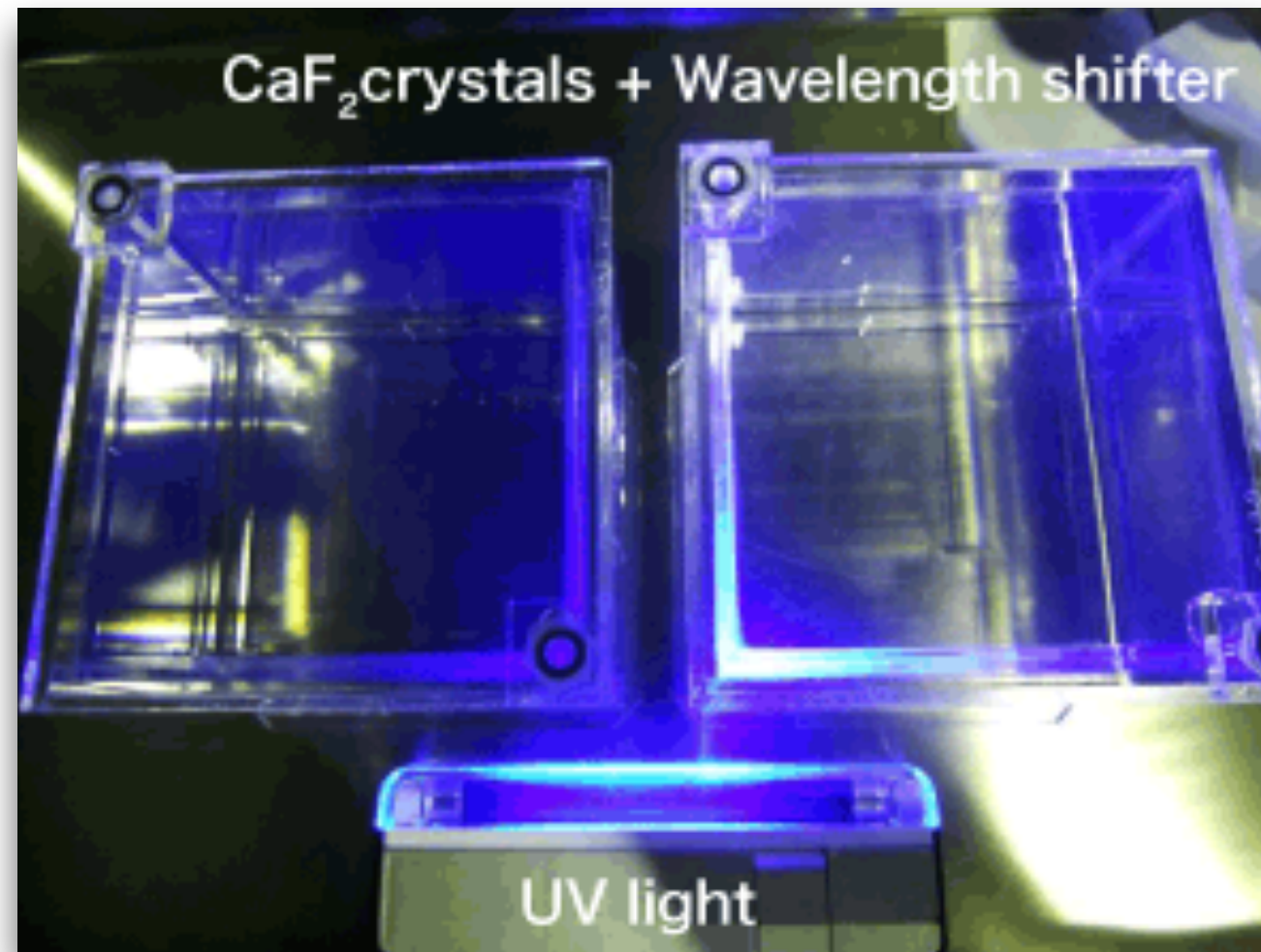
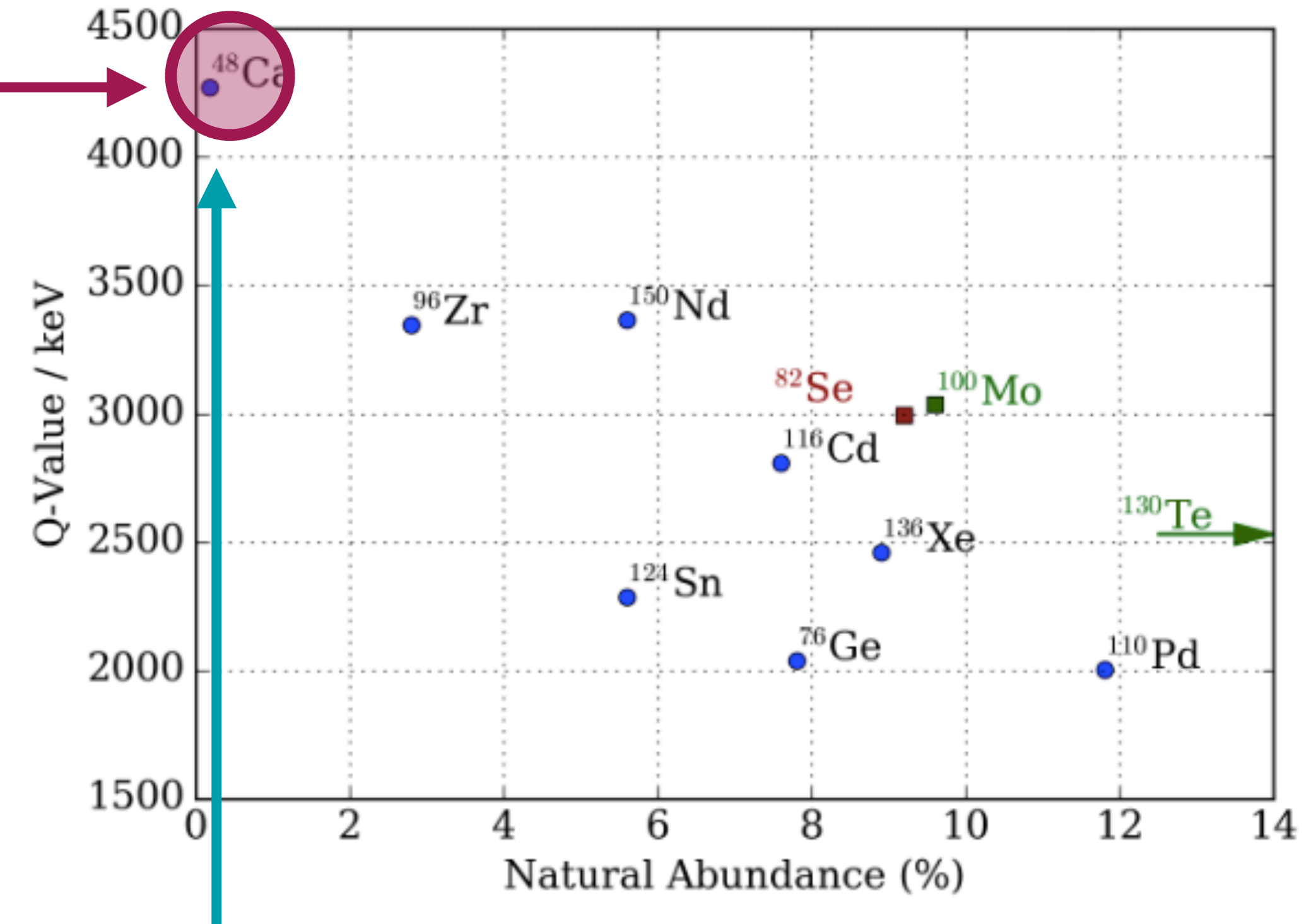


Image: Osaka University



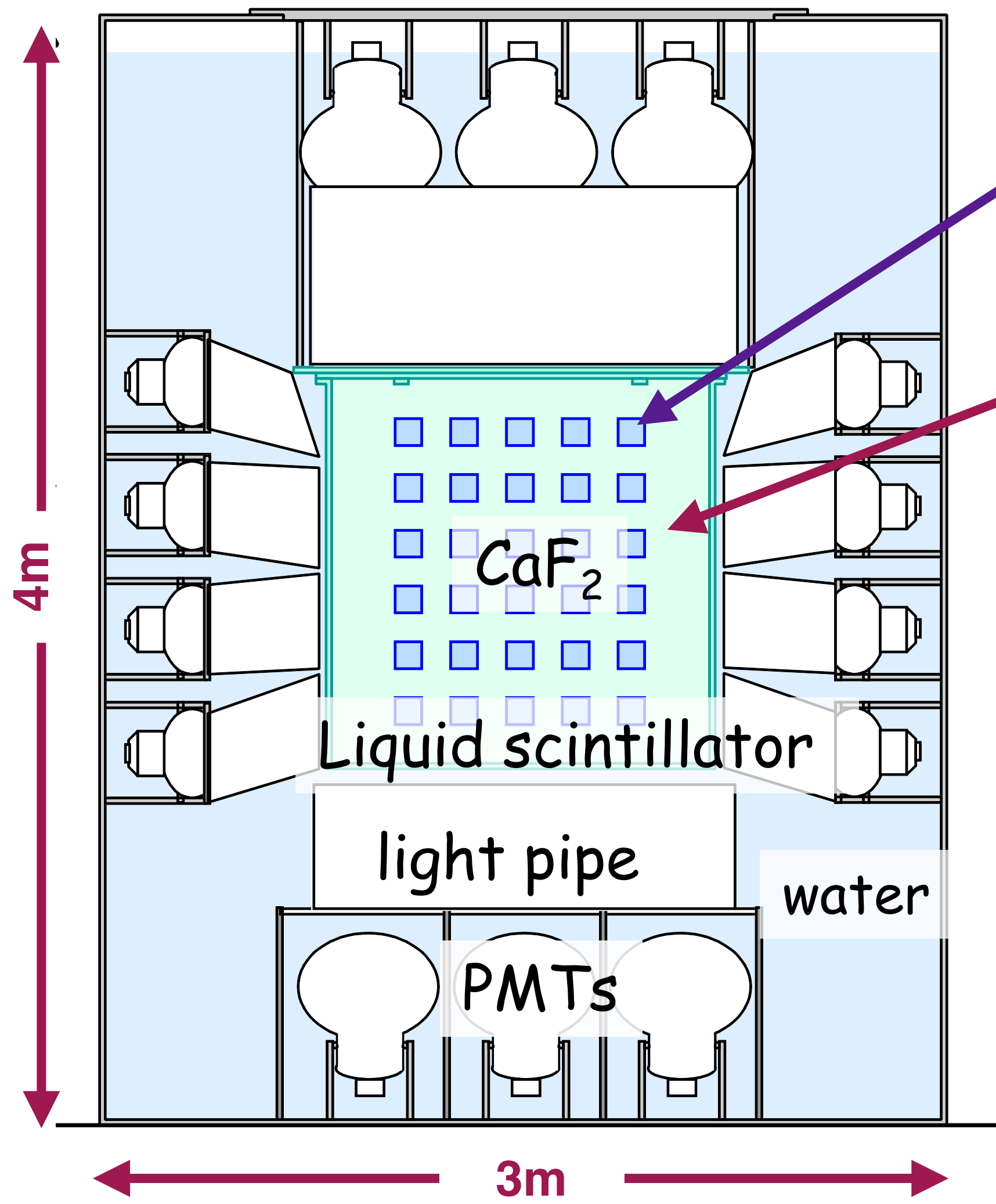
lowest abundance (< 0.2%)

- **zero background** (i.e. high resolution) vital
- CANDLES working on **enrichment** techniques & scintillating bolometers



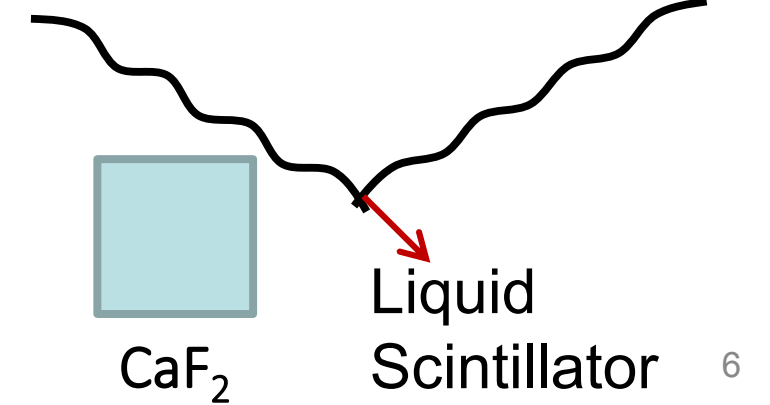
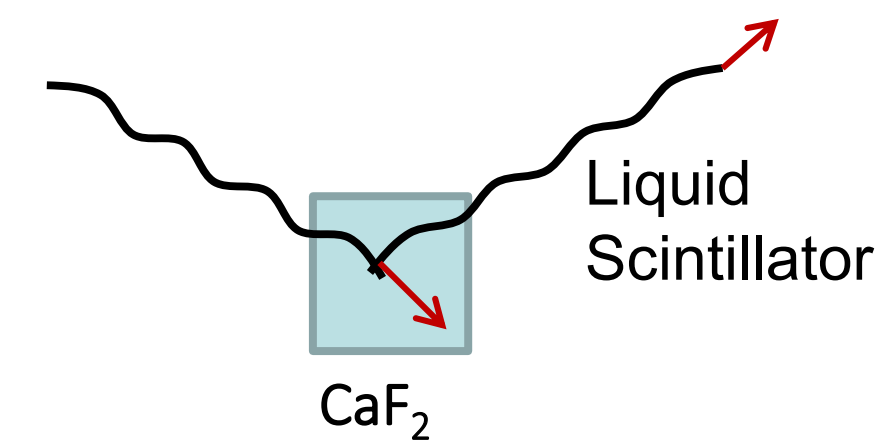
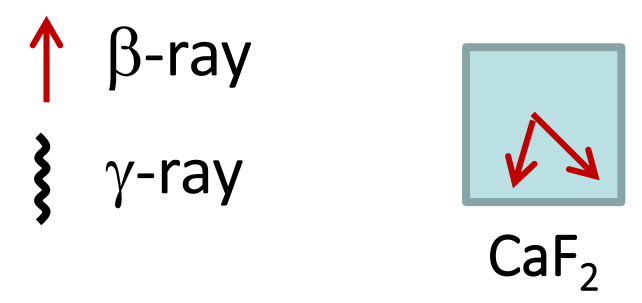
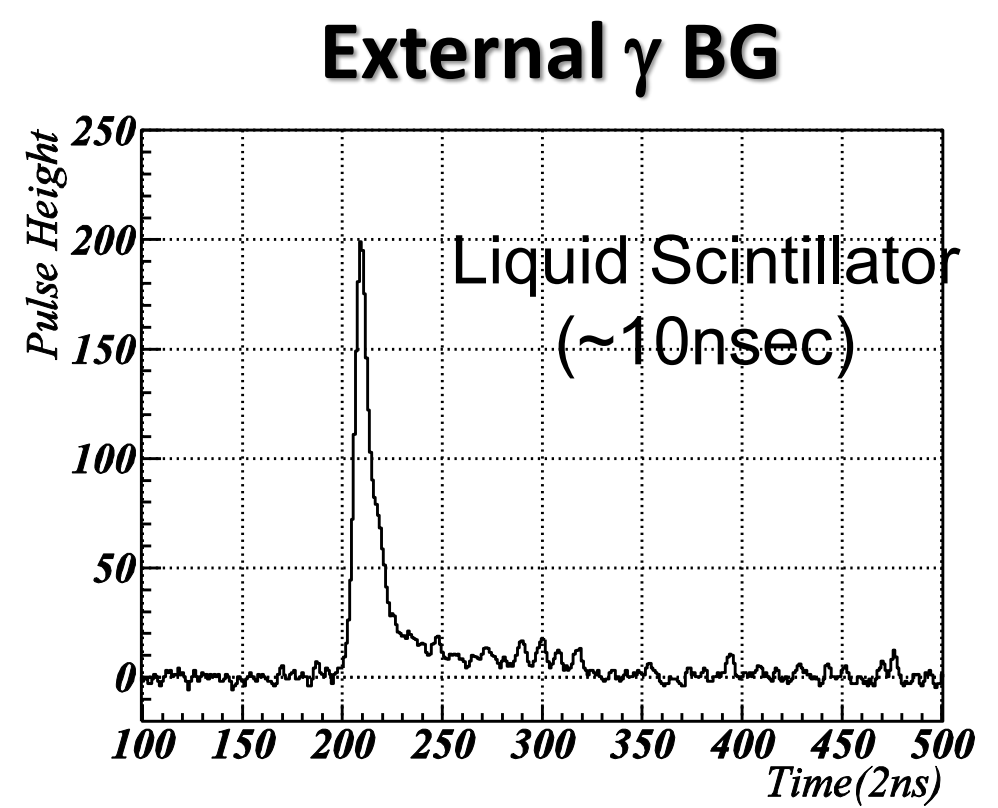
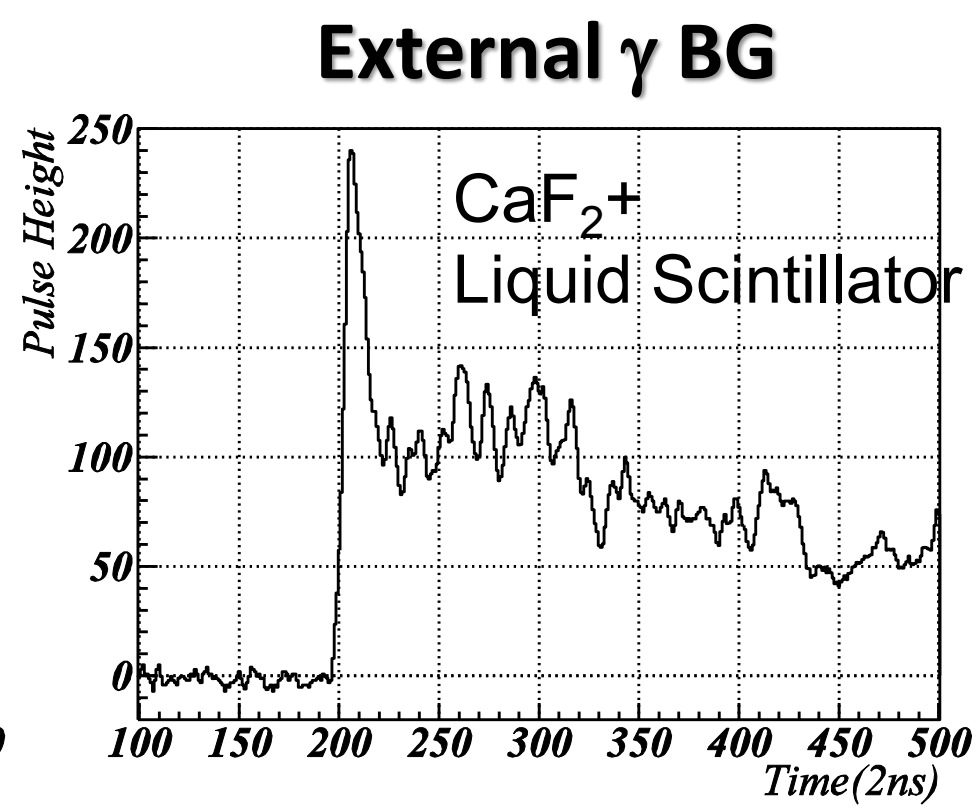
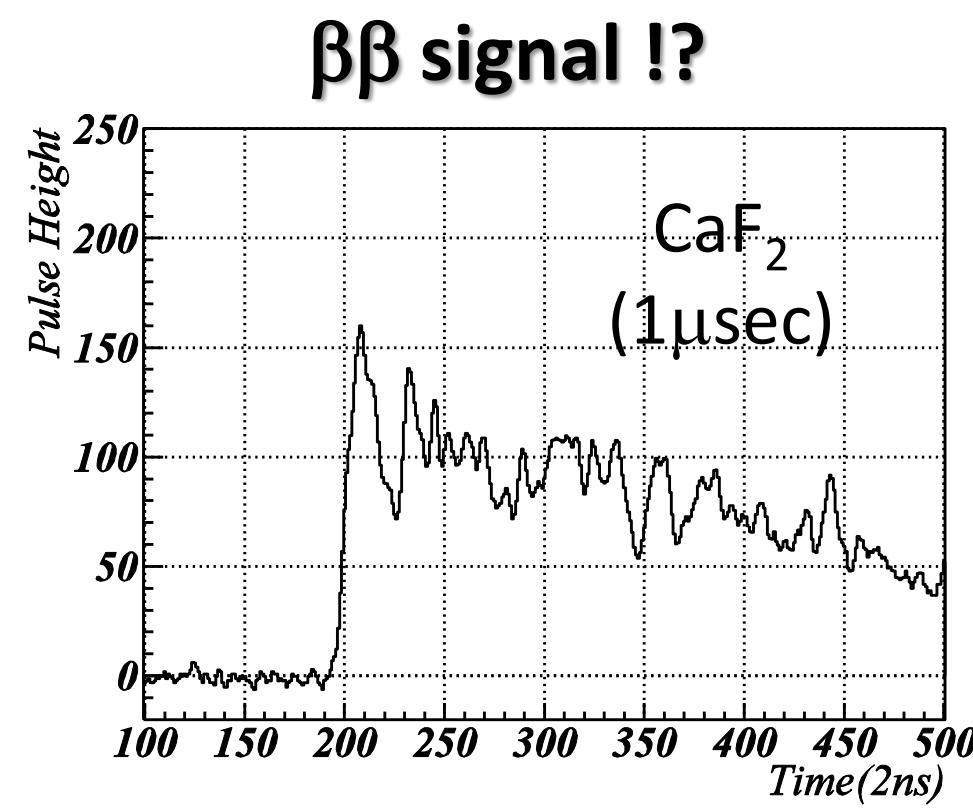
CANDLES III detector

T Kishimoto, DBD18



96 x 3.2 kg of unenriched CaF₂ scintillator crystals

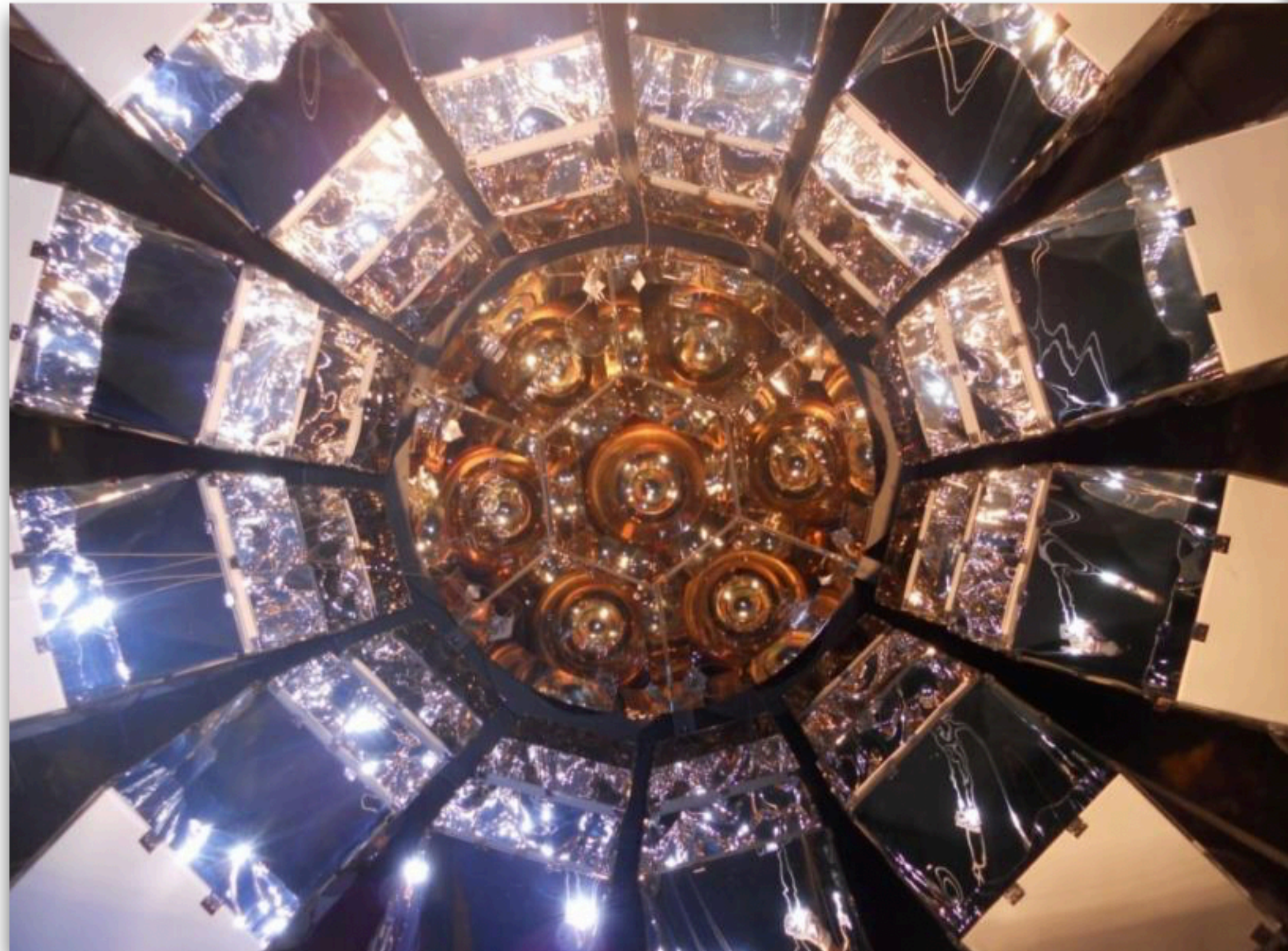
LS response time ~ 100x faster than CaF₂: background rejection





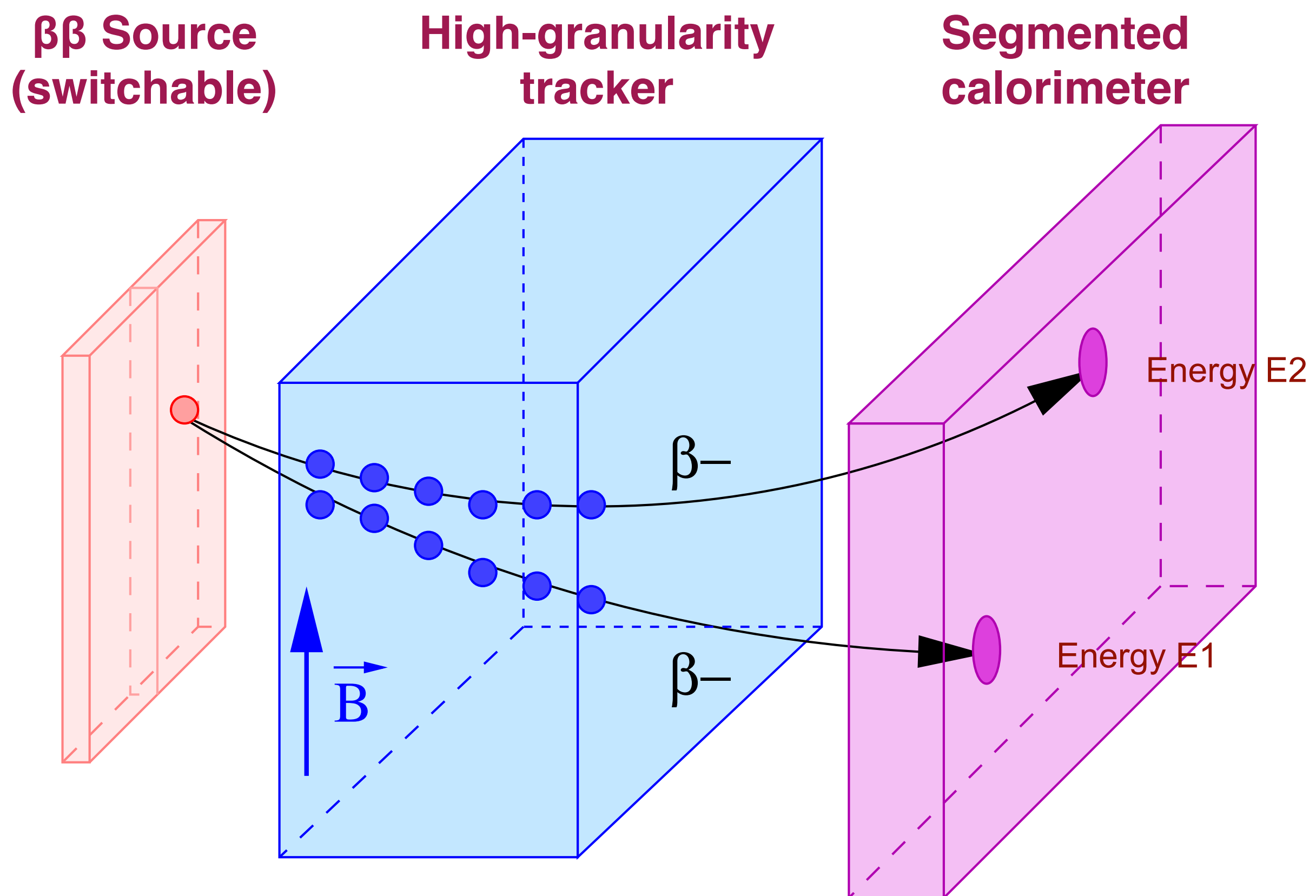
CANDLES III detector

T Kishimoto, DBD18



$T_{1/2}^{0\nu} > 6.2 \times 10^{22}$ years
(world's best)

But enrichment needed...



Strengths



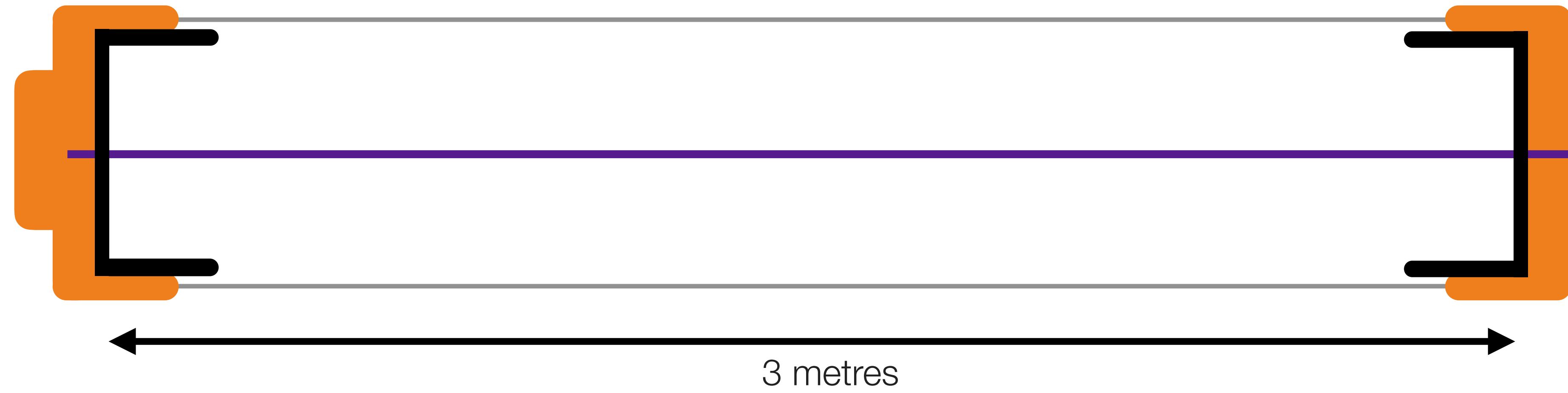
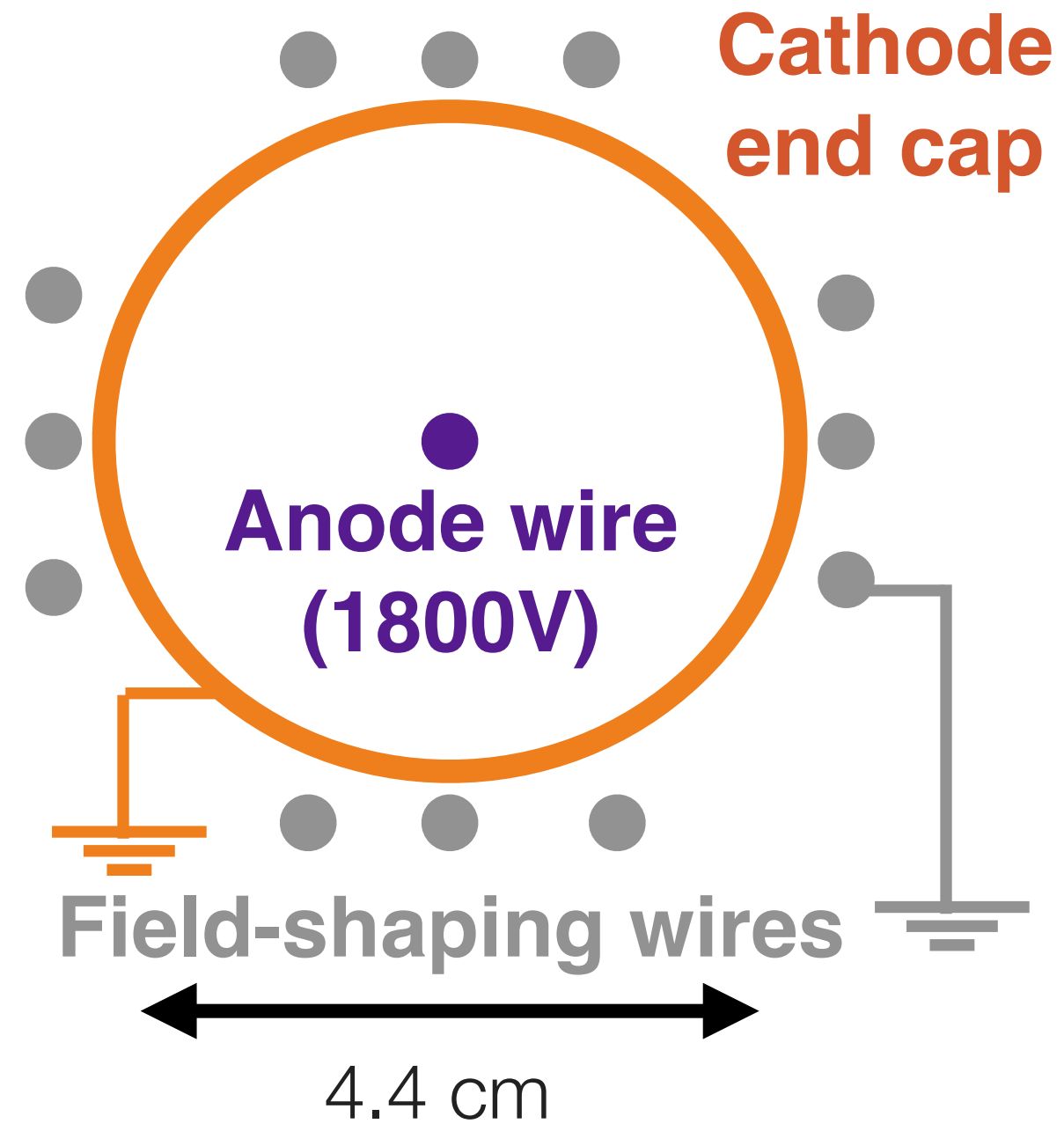
- Source decoupled from detector - use **any solid $\beta\beta$ source** isotope
- Track reconstruction gives **particle identification**
- Combine with timings to identify topologies for ultra-high **background rejection**
- Tracking info (angle between tracks) & individual energy distributions can distinguish between **$\beta\beta$ mechanisms**

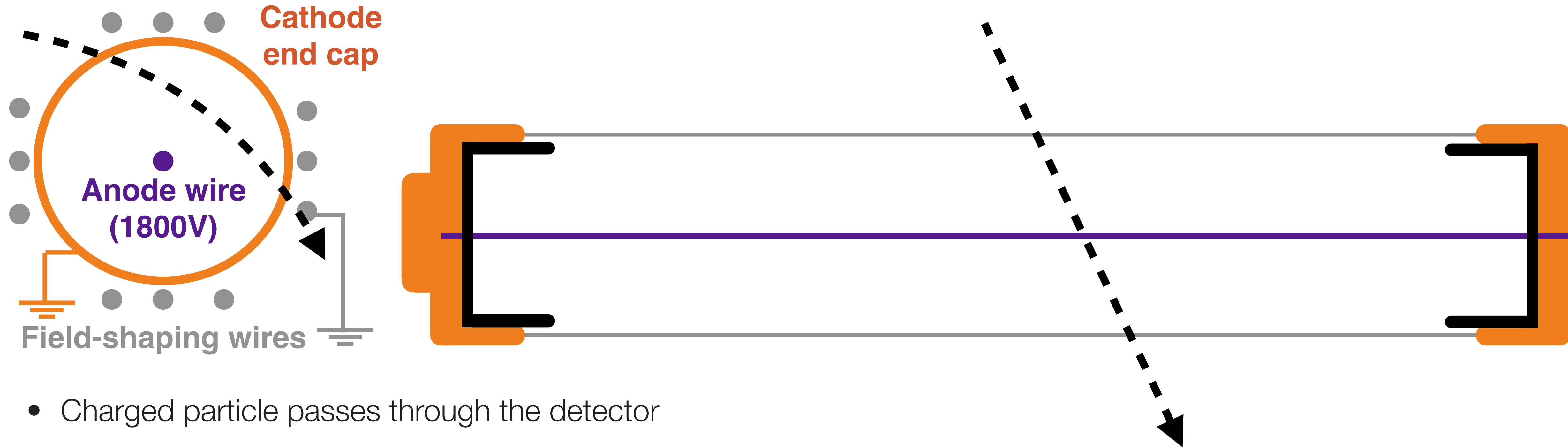
Weaknesses



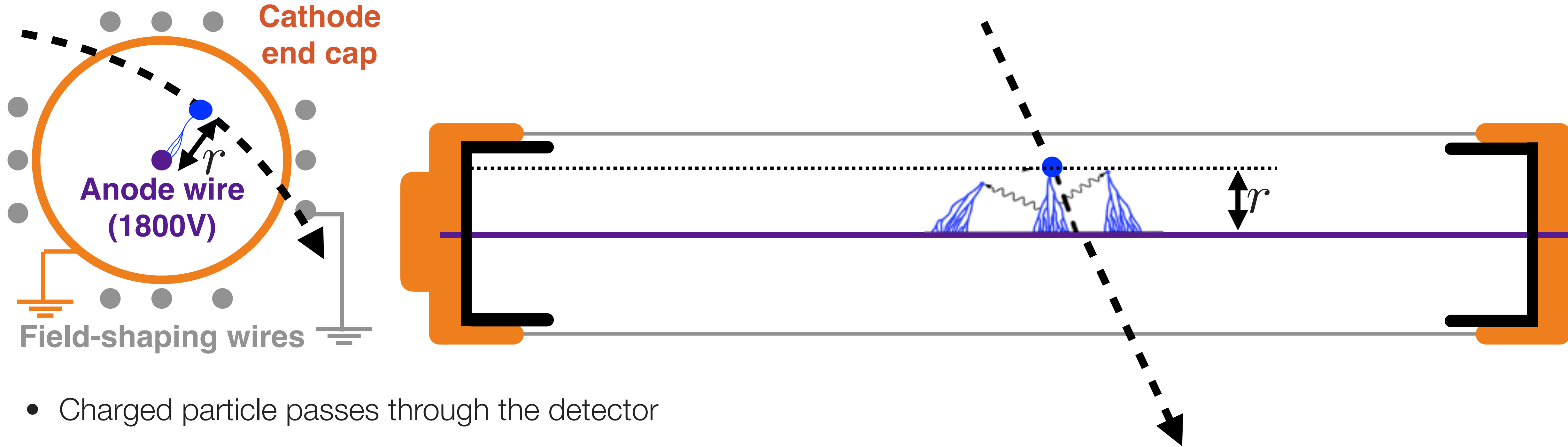
- **Energy resolution** poorer than for most homogenous detectors
- Doesn't scale as well as some other designs

The NEMO tracker

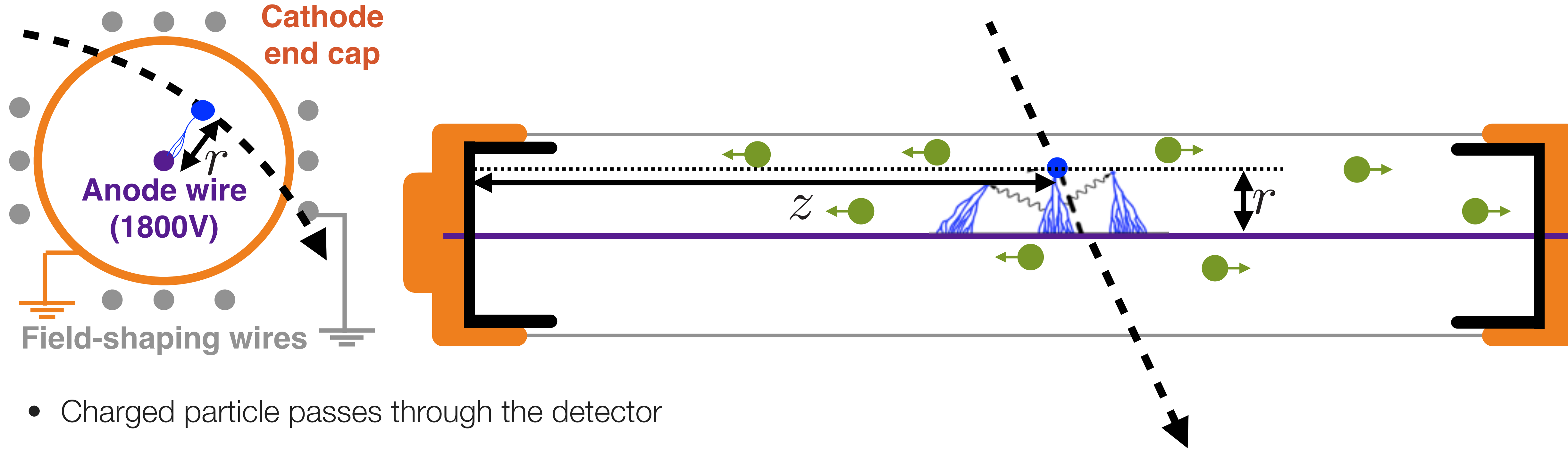




- Charged particle passes through the detector



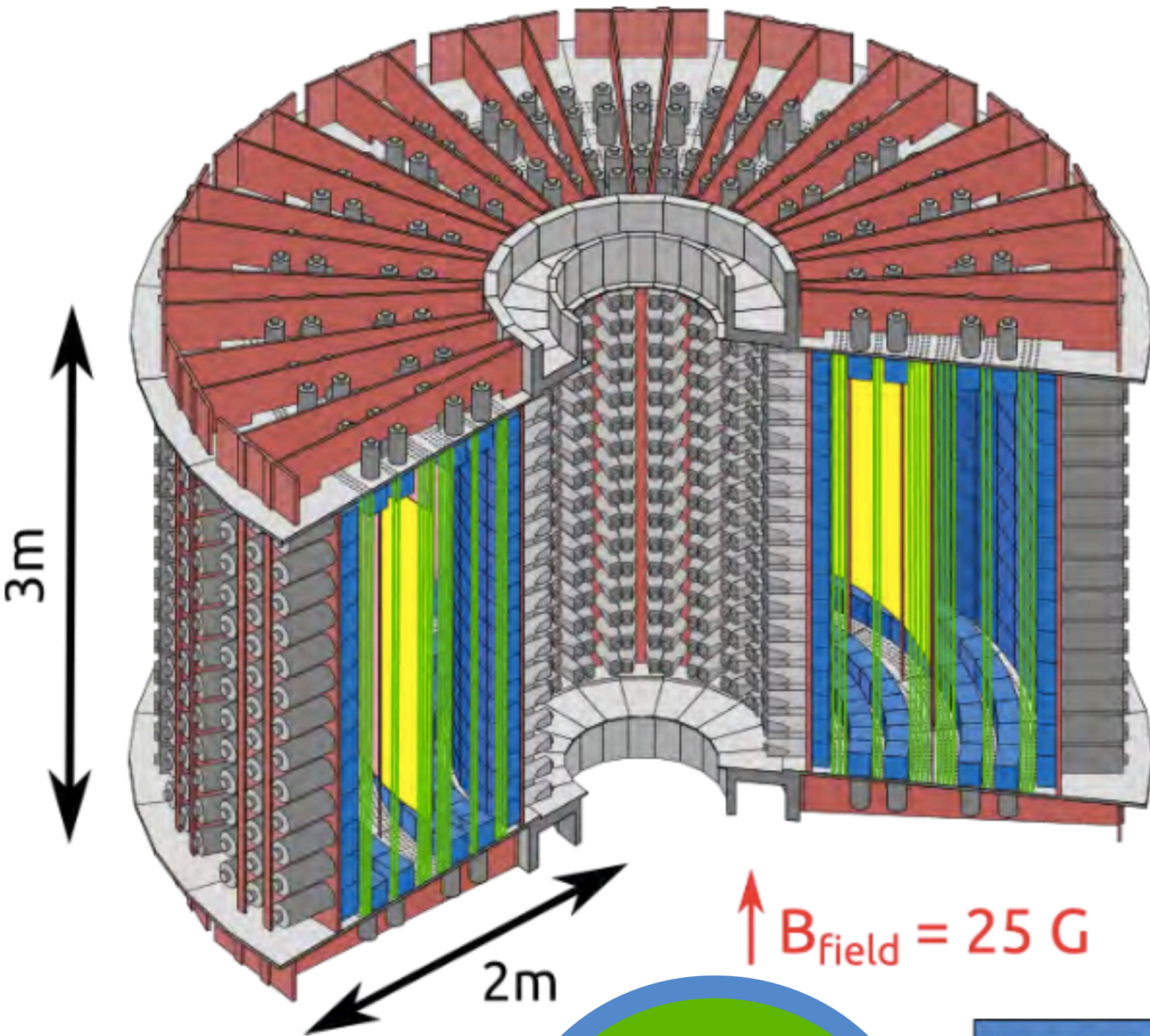
- Charged particle passes through the detector
- Electron avalanche drifts to anode (Geiger mode)
- Drift time gives radius of closest approach r



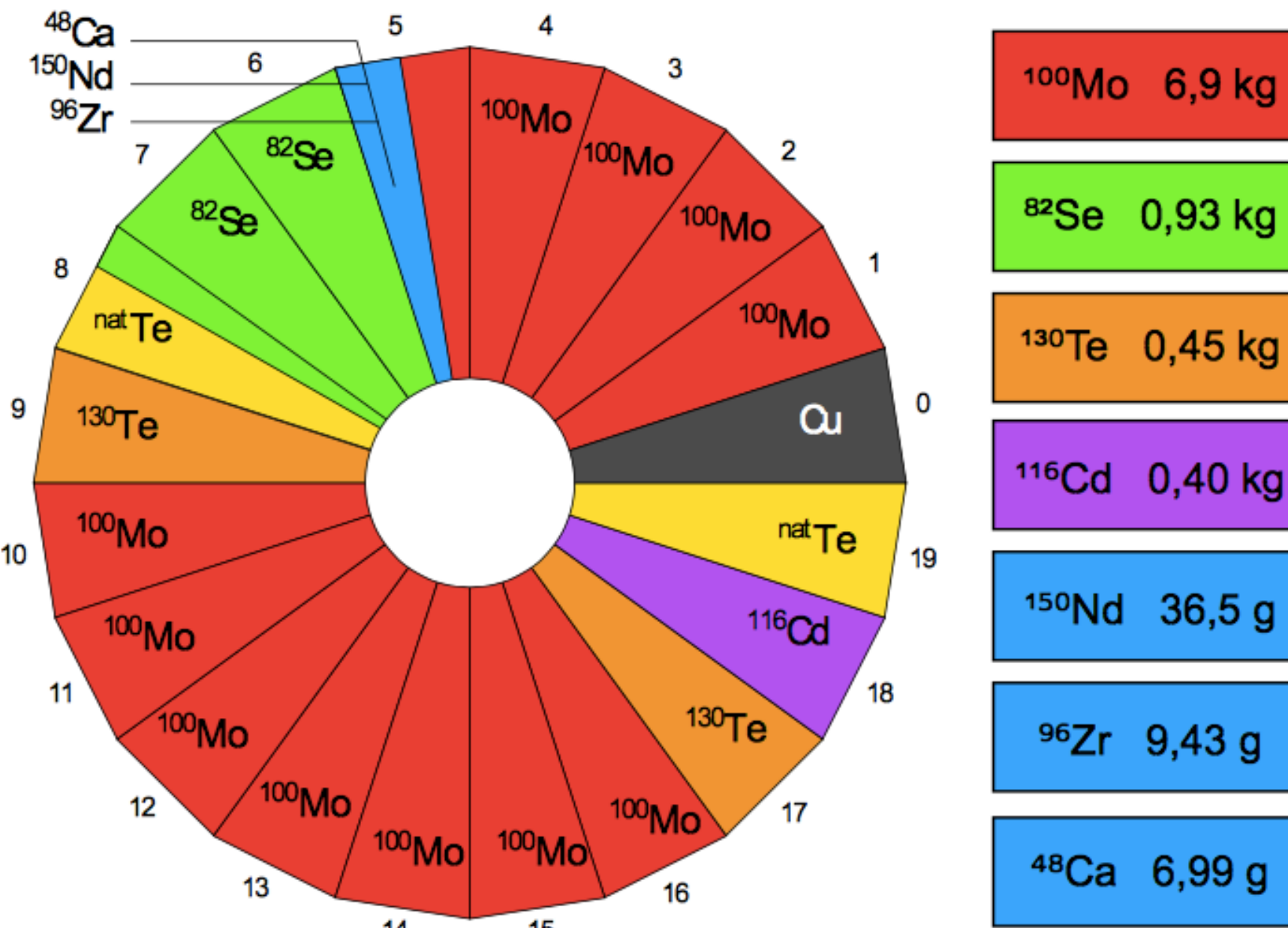
- Charged particle passes through the detector
- Electron avalanche drifts to anode (Geiger mode)
 - Drift time gives radius of closest approach r
- Plasma propagates towards the two cathode end caps
 - Difference in drift times gives distance along wire z

Allows 3-d track reconstruction

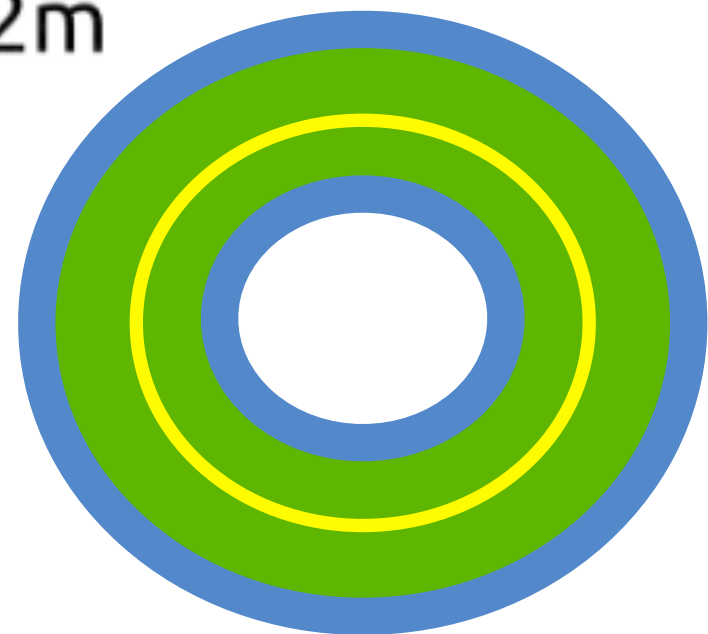
NEMO-3 (2003-11)



NEMO-3 "camembert" (source top view)



$\uparrow B_{\text{field}} = 25 \text{ G}$



calorimeter
1940 optical modules :
polystyren scintillators
+ 3" and 5" PMTs
 $\text{FWHM}_E \sim 15\% / \sqrt{E_{\text{MeV}}}$
 $\sigma_t \sim 250 \text{ ps}$

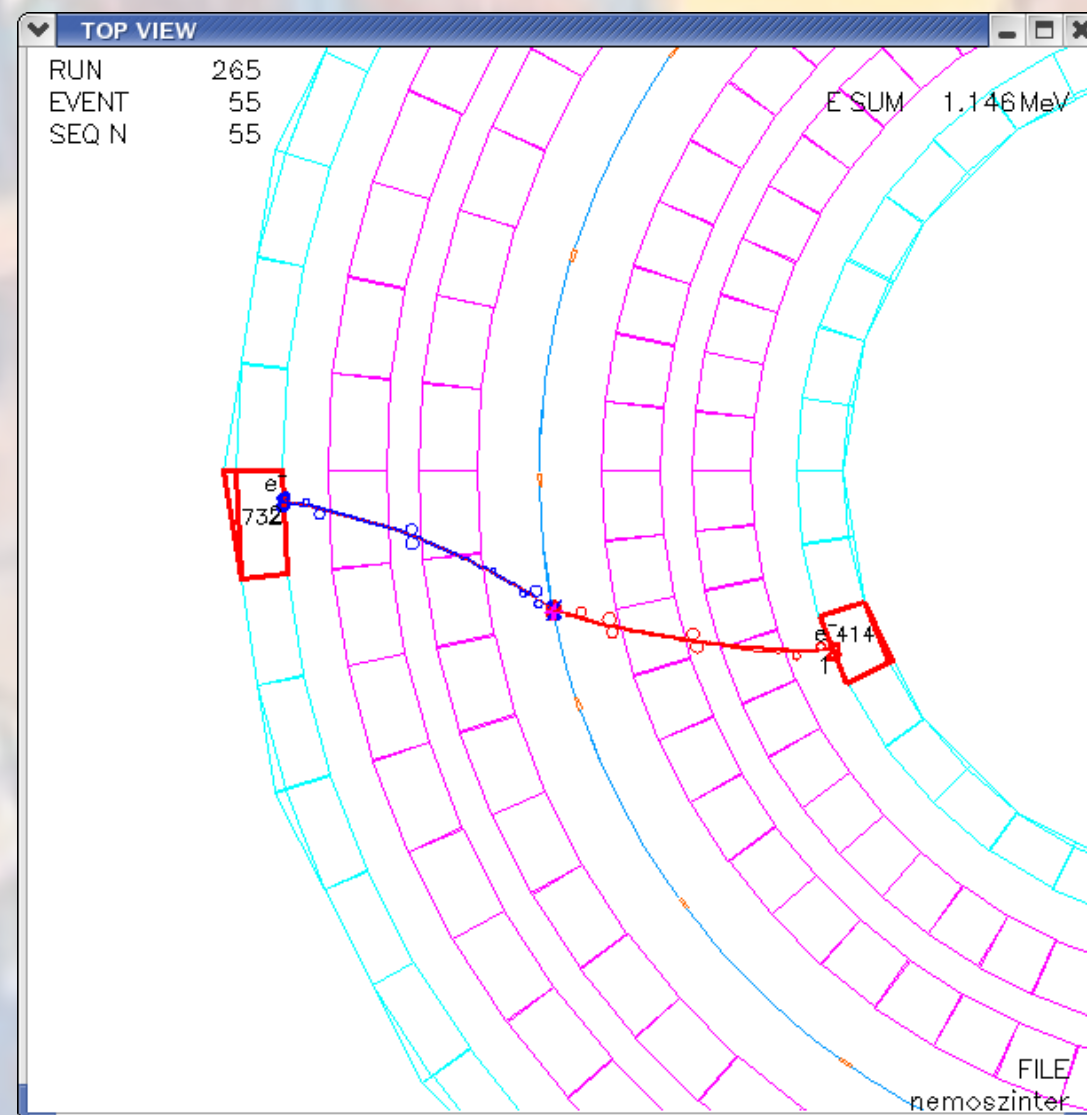
tracker
6180 Geiger cells
vertex resolution :
 $\sigma_{xy} \sim 3 \text{ mm}$ $\sigma_z \sim 10 \text{ mm}$

sources
60 mg/cm² foils
10 kg of $\beta\beta$ isotopes

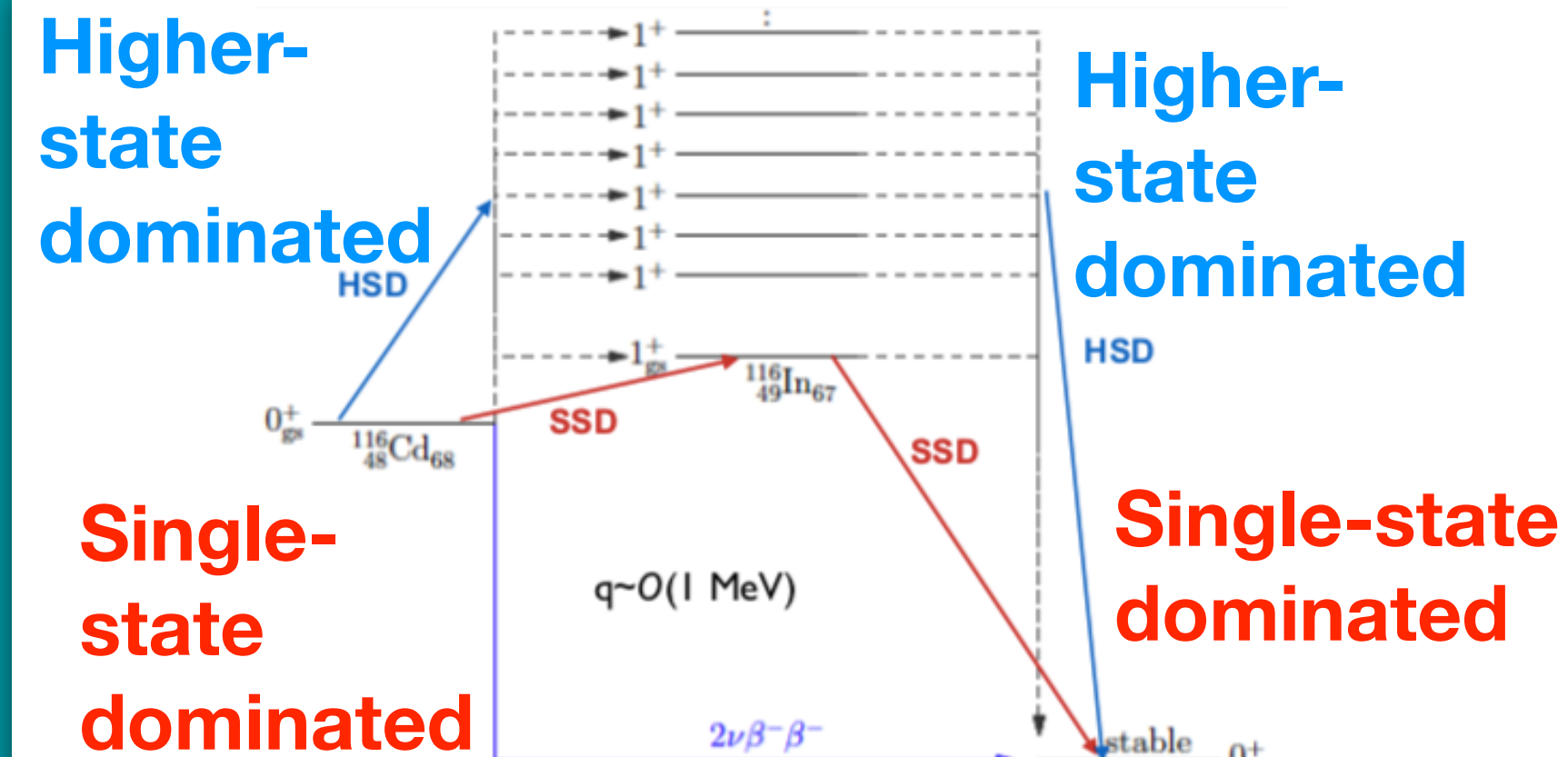


$2\nu\beta\beta$ measurements and $0\nu\beta\beta$ limit

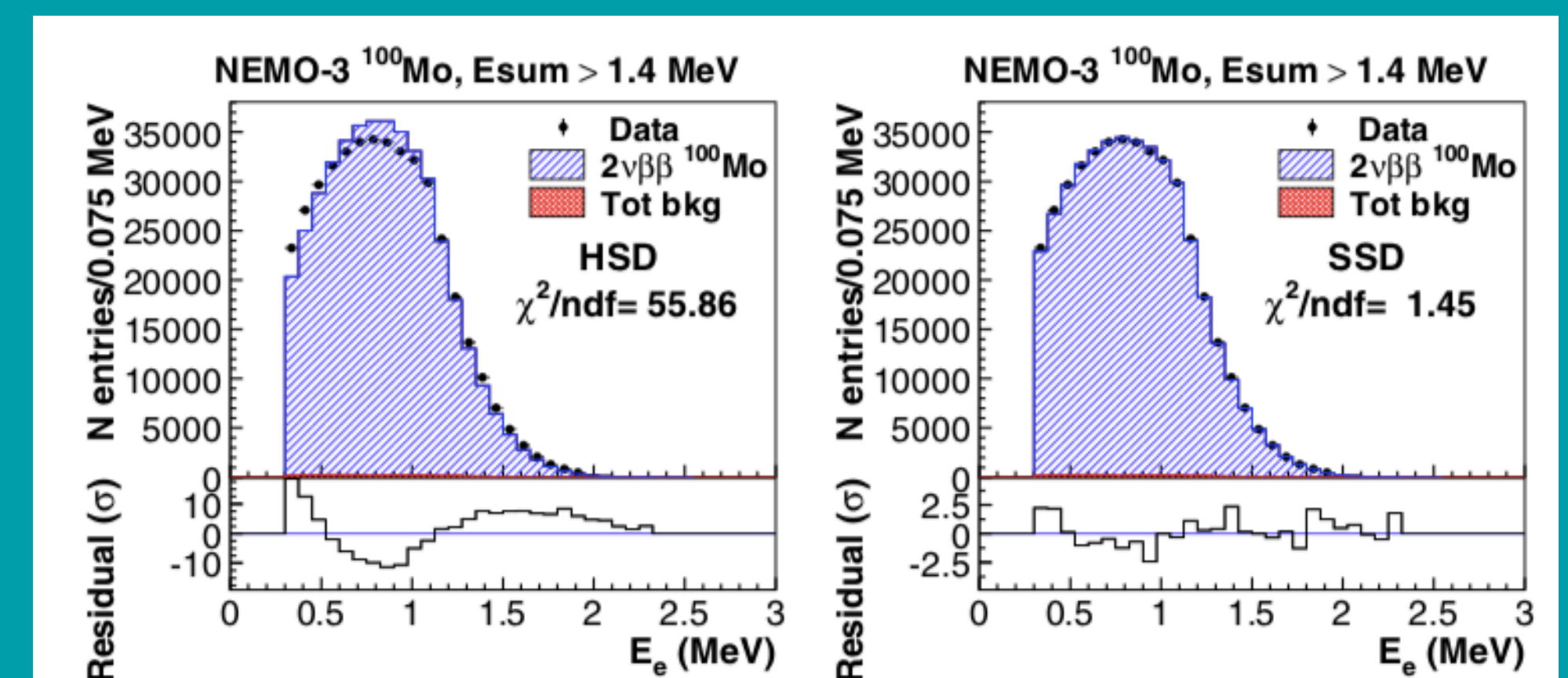
- **^{100}Mo** (Eur. Phys. J. C (2019) 79: 440)



Individual electron track angles and energies give extra information

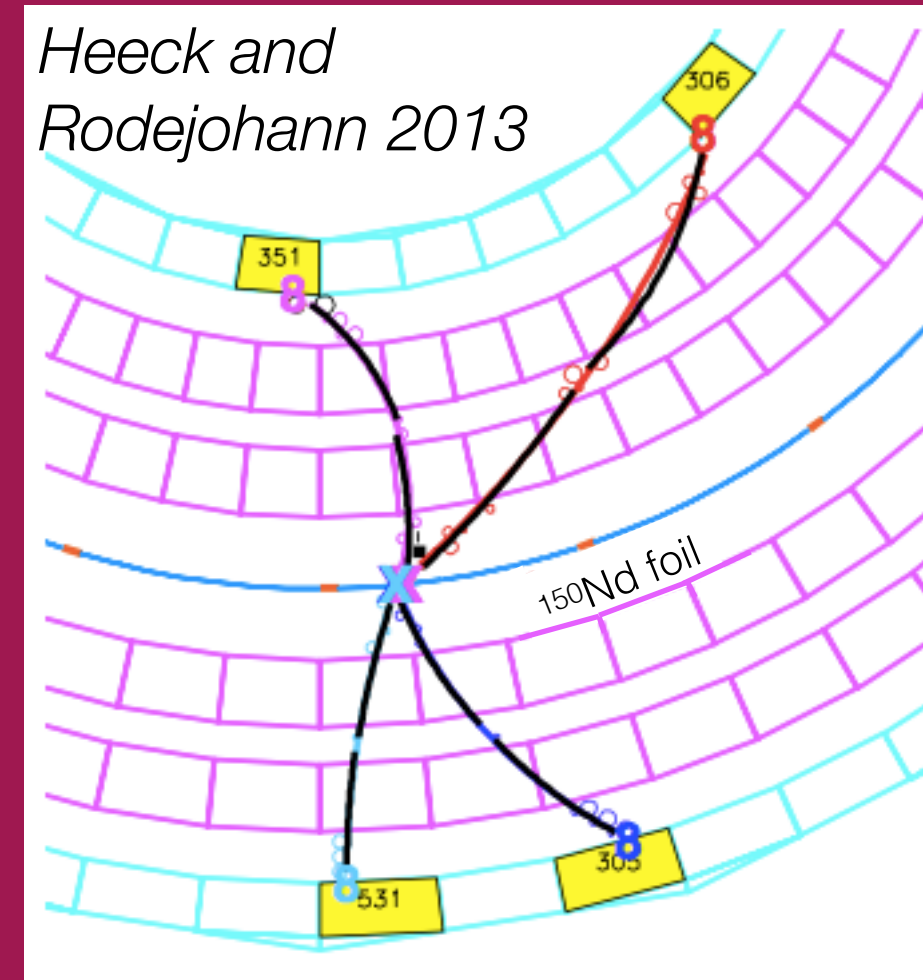


5 x 10⁵ events
Signal / bkgd ~ 80



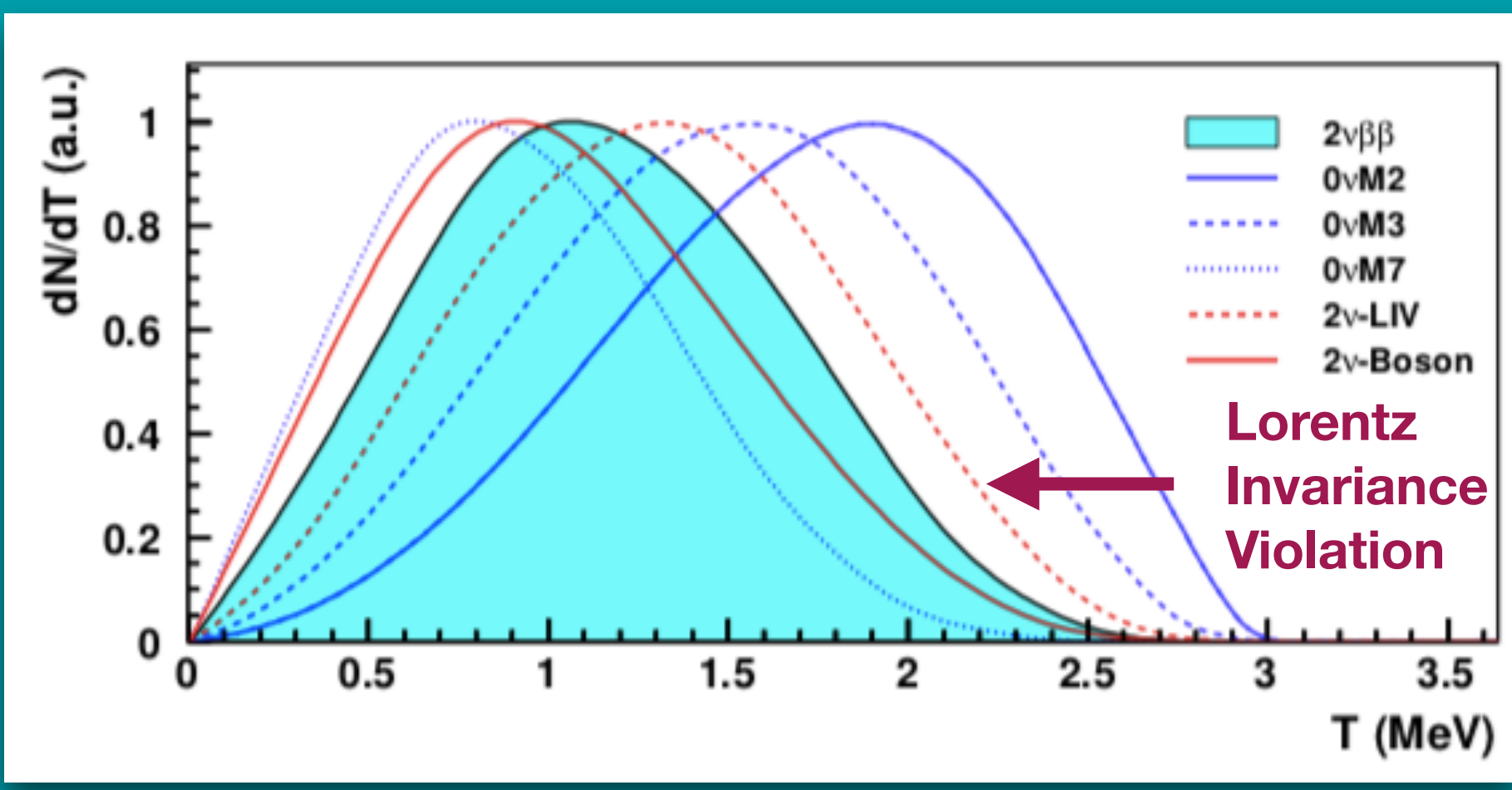
2νββ measurements and 0νββ limit

- **¹⁰⁰Mo** (Eur. Phys. J. C (2019) 79: 440)
- **⁸²Se** (Eur. Phys. J. C (2018) 78: 821)
- **⁴⁸Ca** (Phys. Rev. D 93, 112008)
- **¹⁵⁰Nd** (Phys. Rev. D 94, 072003)
- **¹¹⁶Cd** (Phys. Rev. D 95, 012007)
- **¹³⁰Te** (Phys. Rev. Lett. 107, 062504)
- **⁹⁶Zr** (Nucl.Phys.A847:168-179)



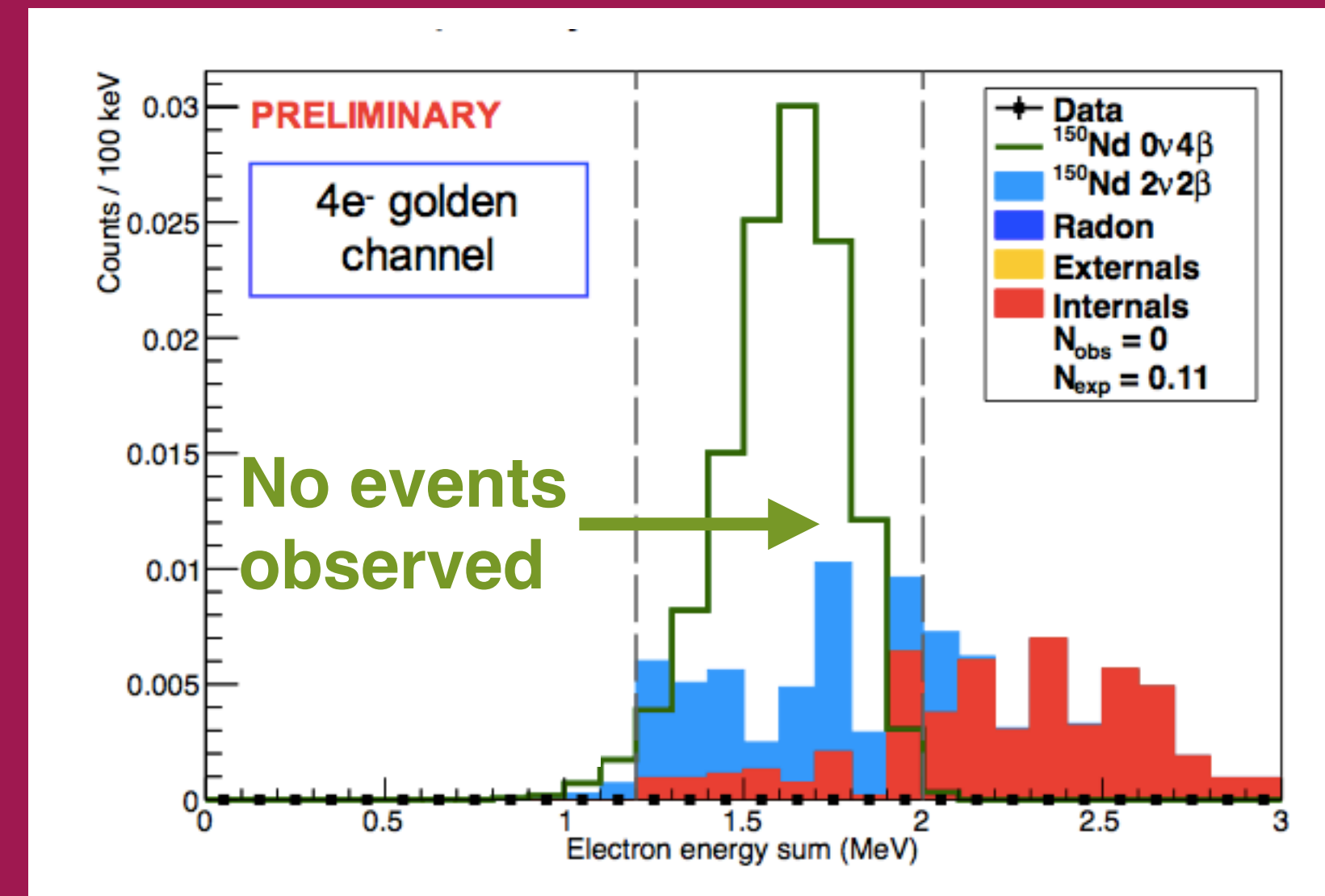
0ν4β decays would violate lepton number, but could occur even if neutrinos are Dirac

Allowed for **3 isotopes**, including ¹⁵⁰Nd



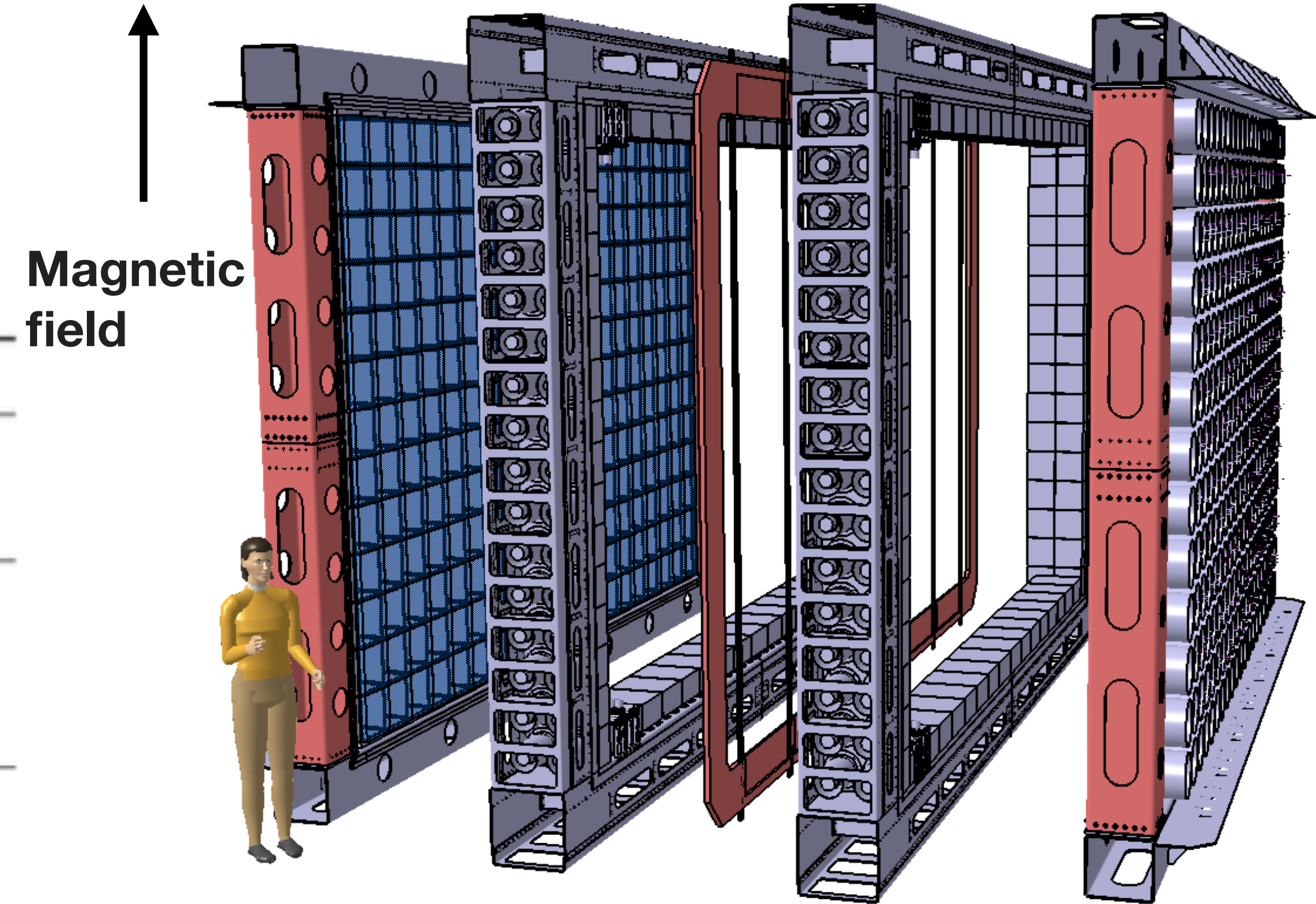
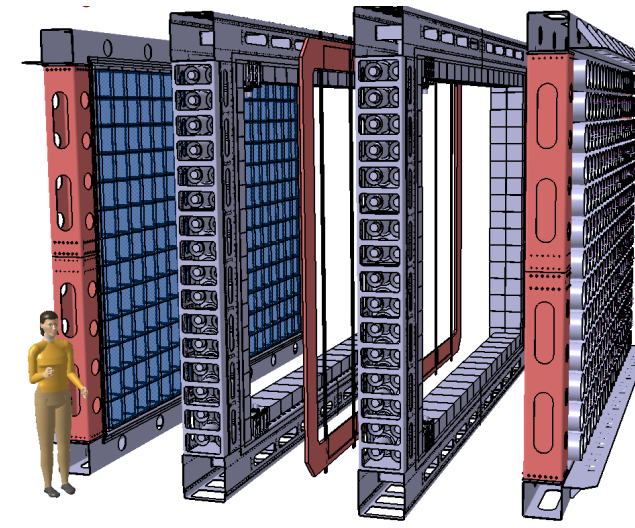
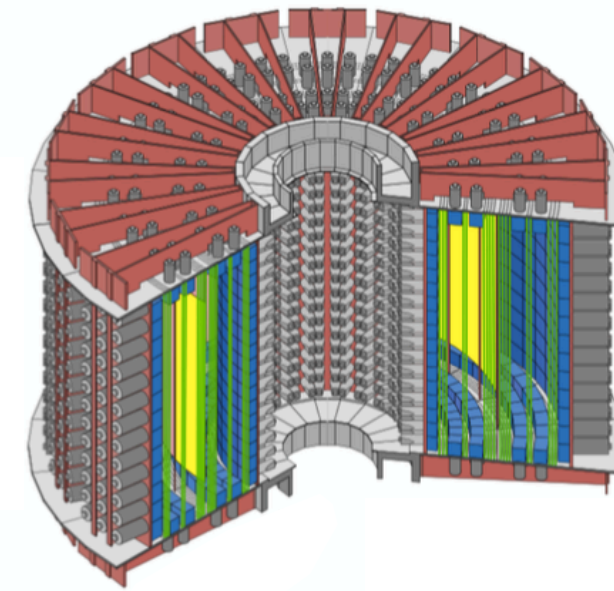
Lorentz Invariance Violation and exotic 0νββ mechanisms would modify energy spectrum
Limit set on contribution from Lorentz-Invariance violating events

$$-4.2 \times 10^{-7} \text{ GeV} < \dot{a}_{of}^{(3)} < 3.5 \times 10^{-7} \text{ GeV} \quad (90\% \text{ C.L.}).$$



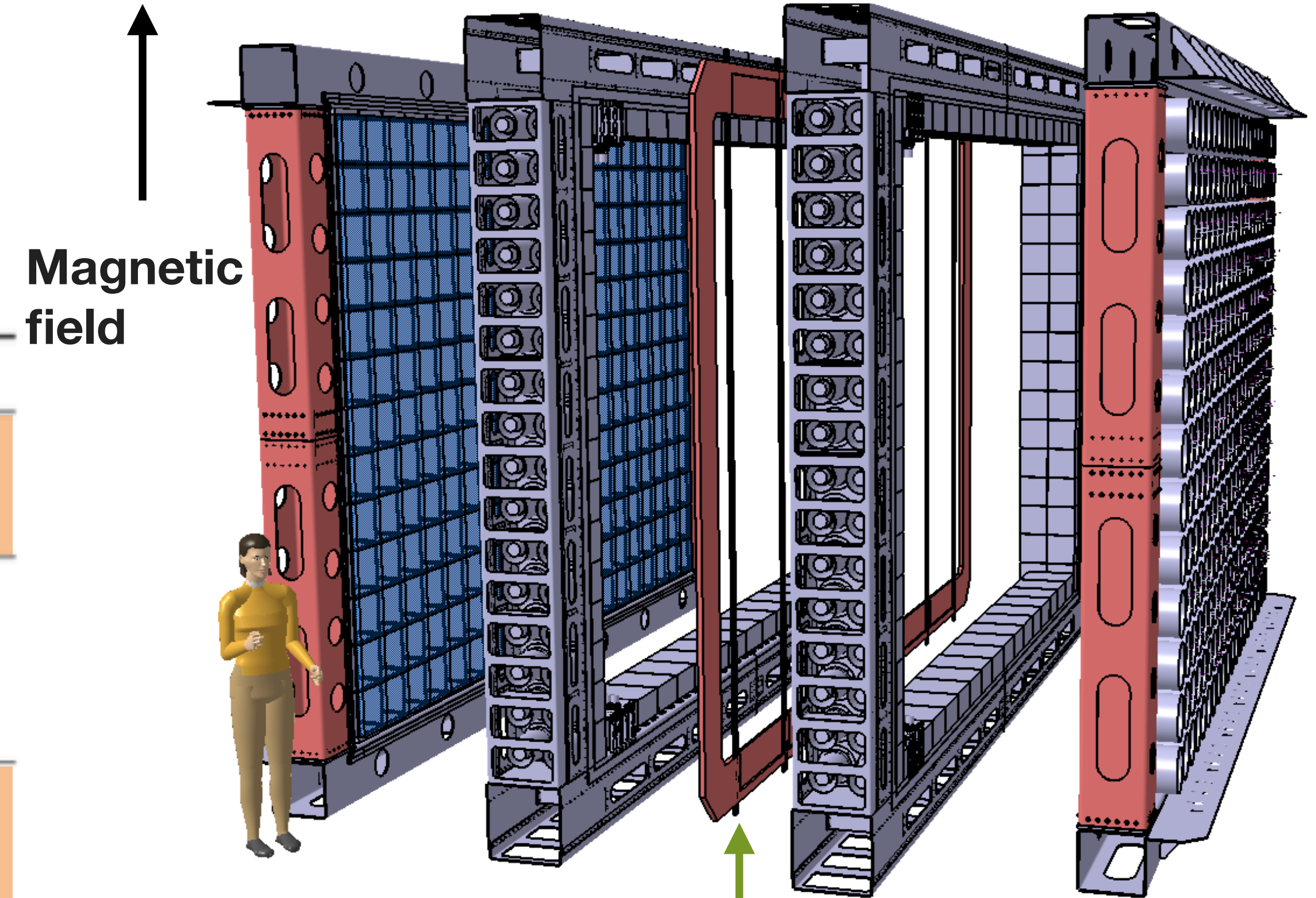
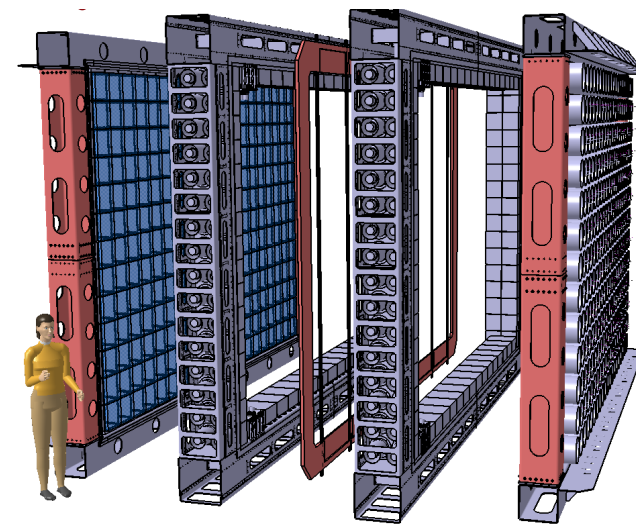
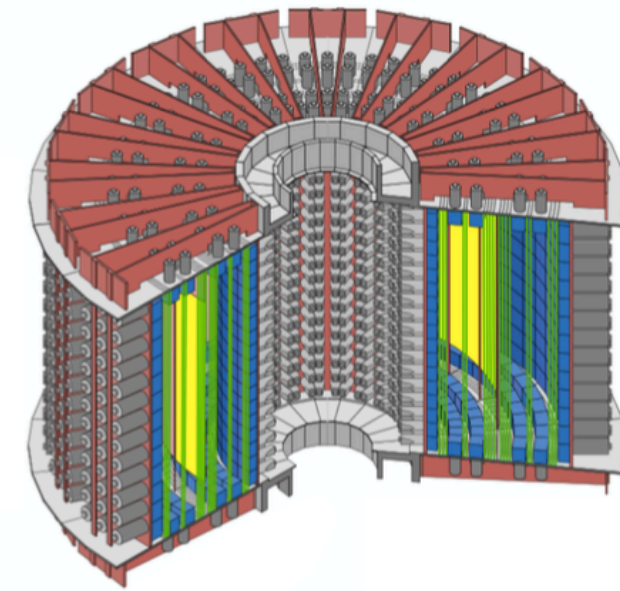
$$T_{1/2} > 2.6 \times 10^{21} \text{ yr} \quad (90\% \text{ CL})$$

Now at LSM: SuperNEMO Demonstrator



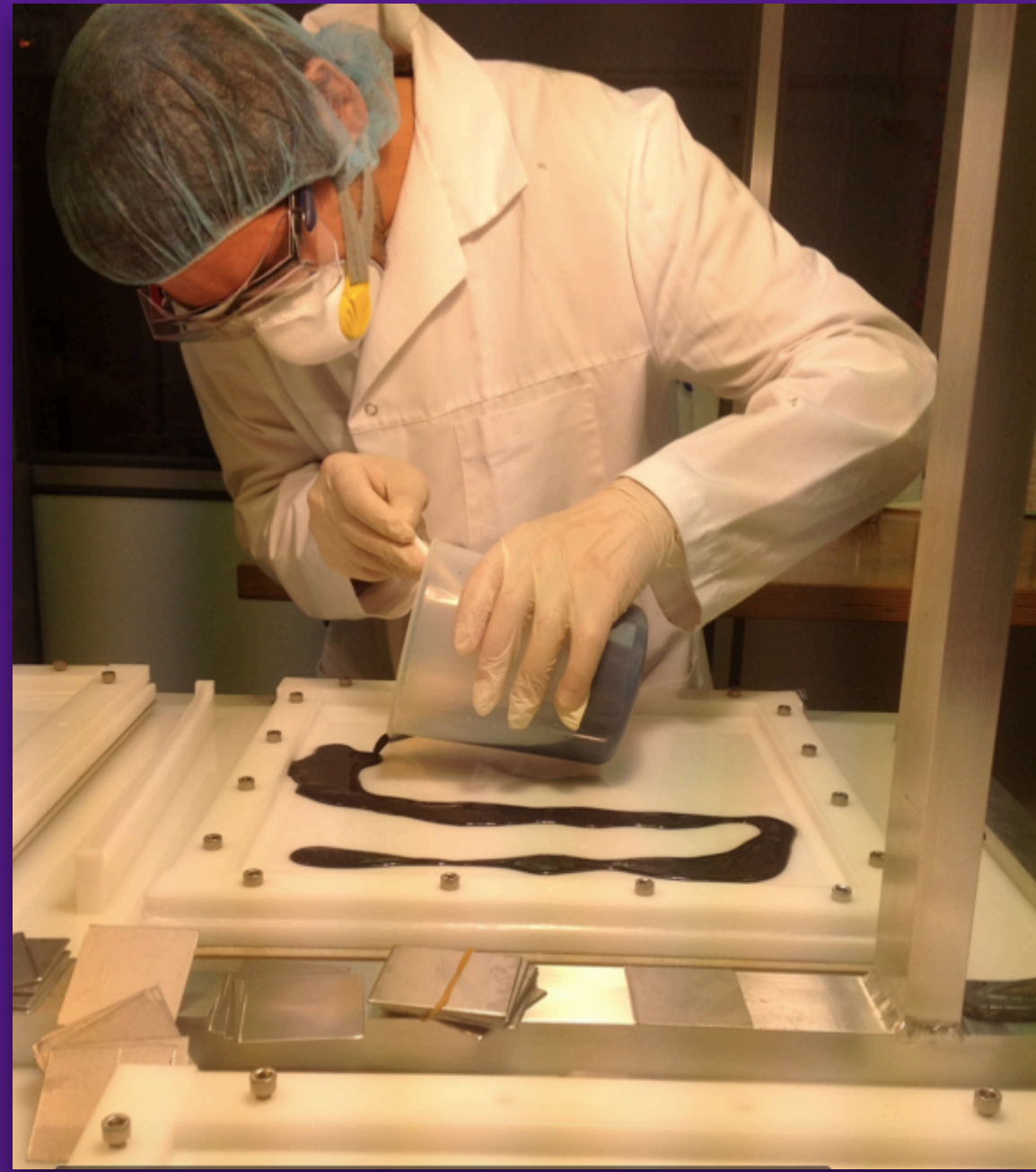
	NEMO-3	SuperNEMO demonstrator
Mass [kg] (main isotopes)	7 (^{100}Mo)	6.3 (^{82}Se)
$T_{1/2}^{2\nu}$ [y]	6.8×10^{18}	9.4×10^{19}
Energy resolution FWHM at 1 MeV	15 %	8 %
FWHM at 3 MeV	8 %	4 %
Source radiopurity $A(^{208}\text{Tl})$	$\sim 100 \mu\text{Bq/kg}$	$< 2 \mu\text{Bq/kg}$
$A(^{214}\text{Bi})$	$< 300 \mu\text{Bq/kg}$	$< 10 \mu\text{Bq/kg}$
Level of radon $A(^{222}\text{Rn})$	$\sim 5.0 \text{ mBq/m}^3$	$< 0.15 \text{ mBq/m}^3$
Sensitivity after 5 (2.5) y data taking	$T_{1/2}^{0\nu} > 10^{24} \text{ y}$	$T_{1/2}^{0\nu} > 6 \times 10^{24} \text{ y}$

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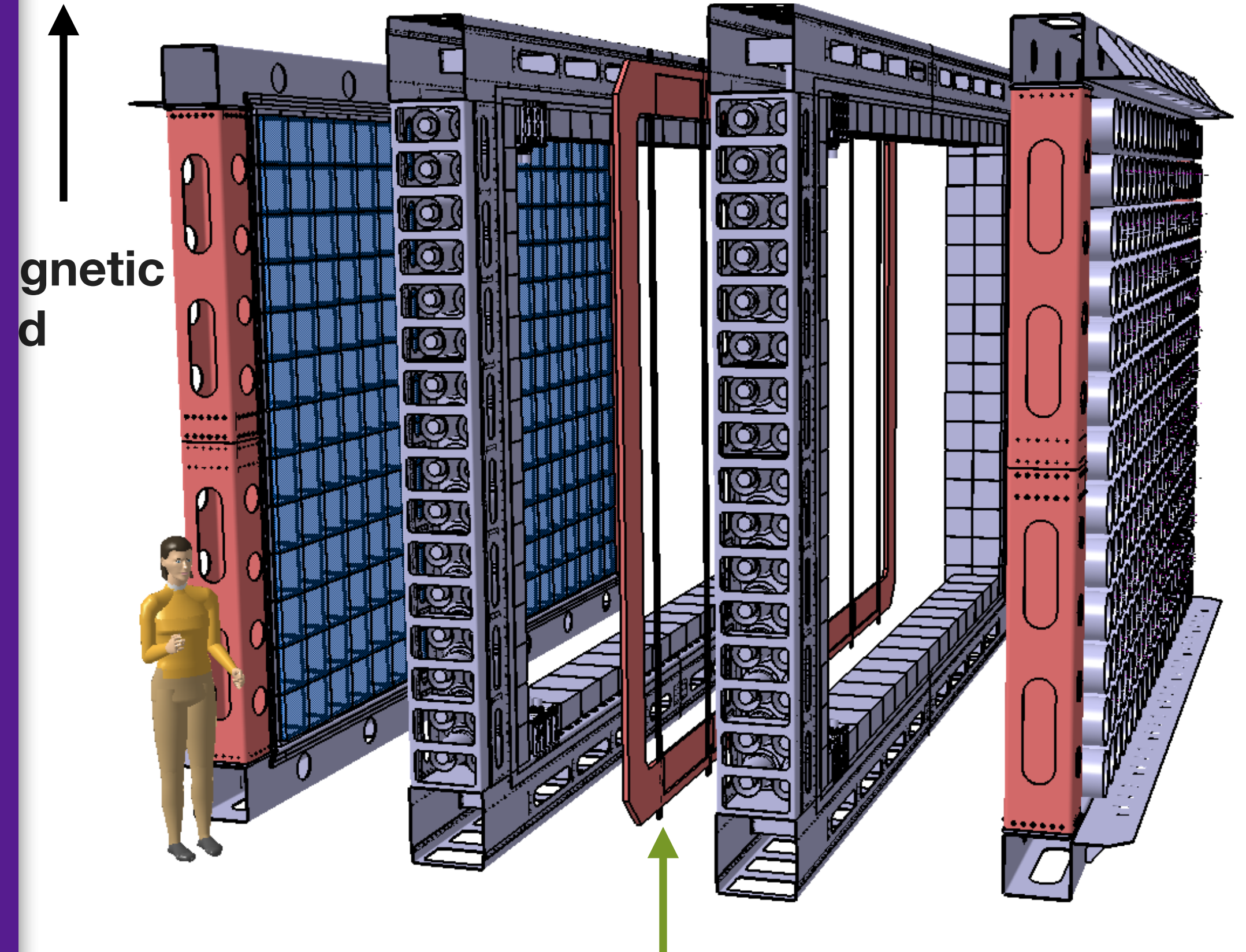


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Source frame holding 6.3kg of $\beta\beta$ emitter (^{82}Se)

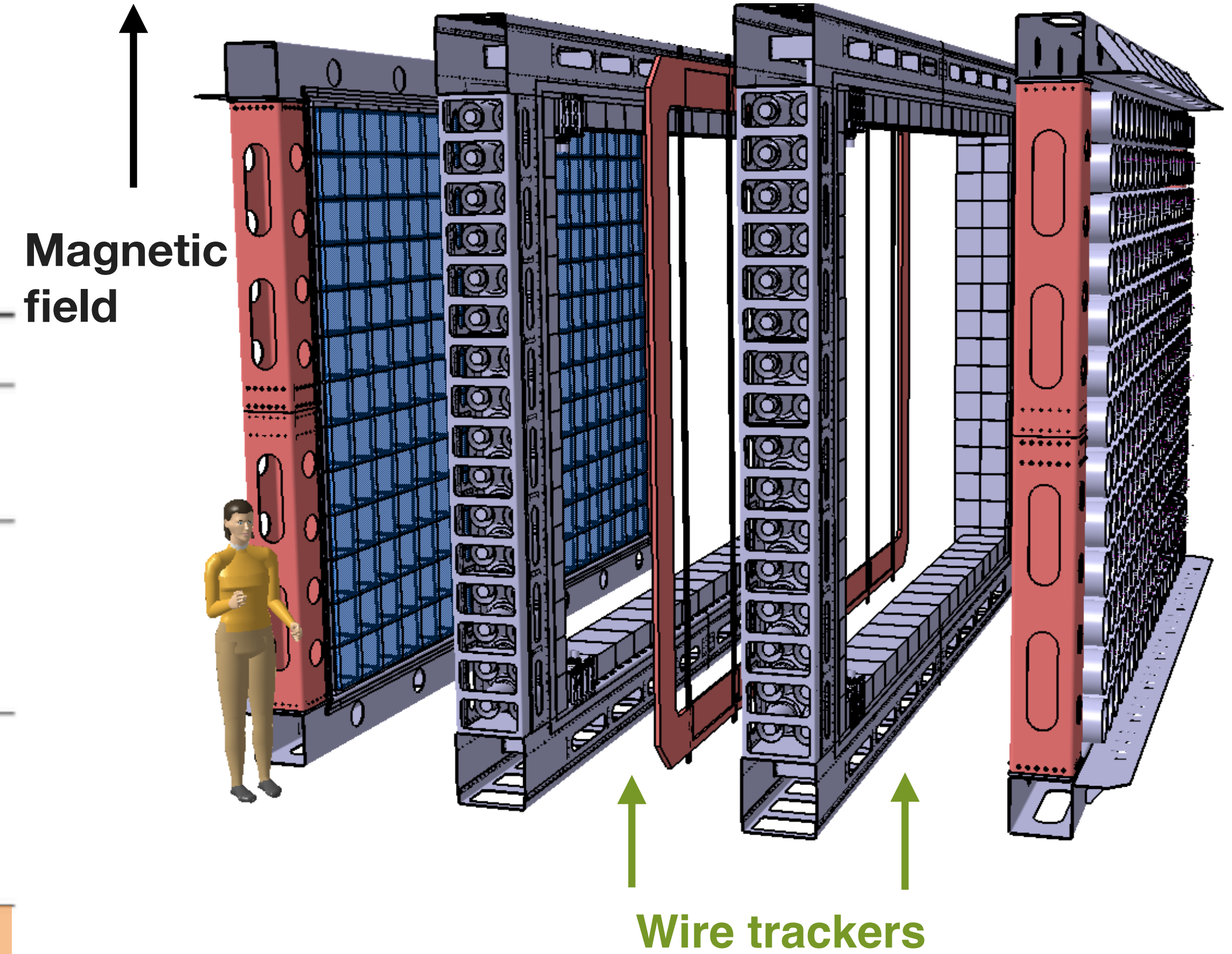
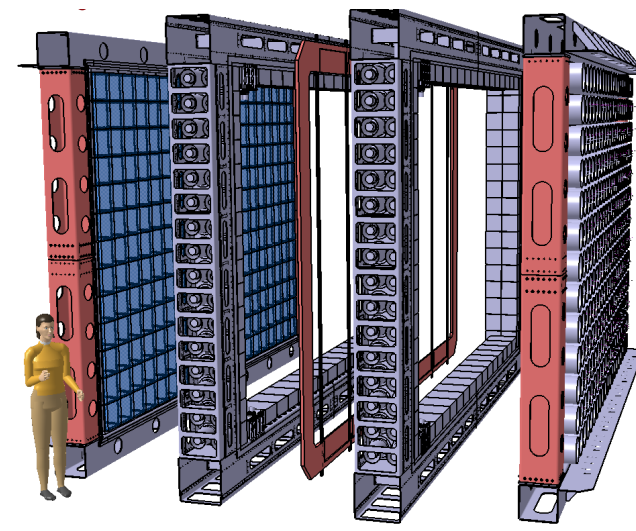
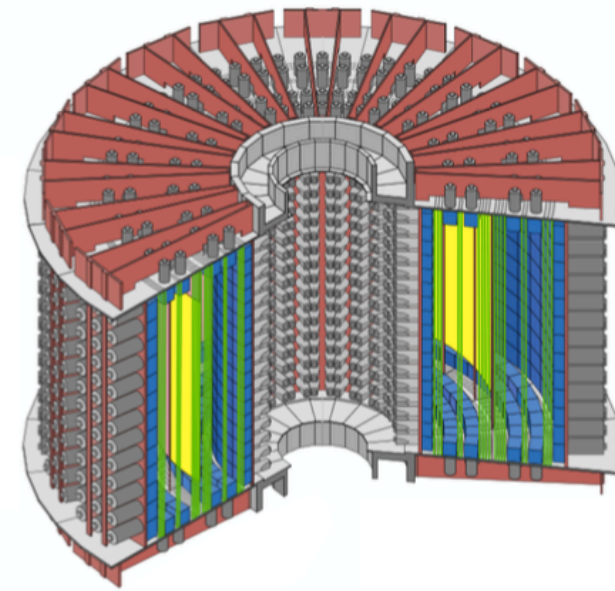


- 34 foils
- Enriched Se powder mixed with PVA
- Increased radio purity through distillation / chromatography / chemical precipitation



Source frame holding 6.3kg of $\beta\beta$ emitter (^{82}Se)

Now at LSM: SuperNEMO Demonstrator

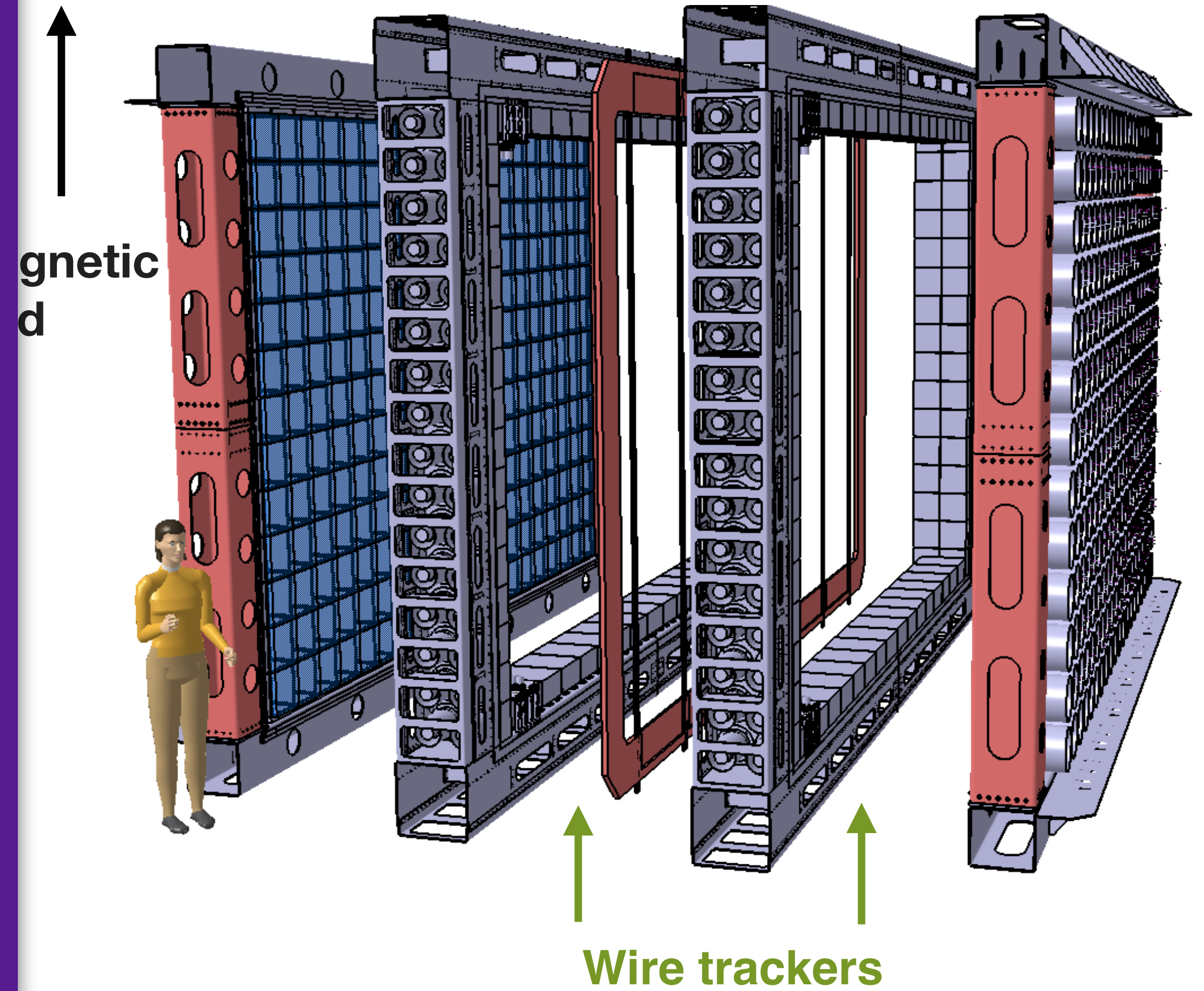


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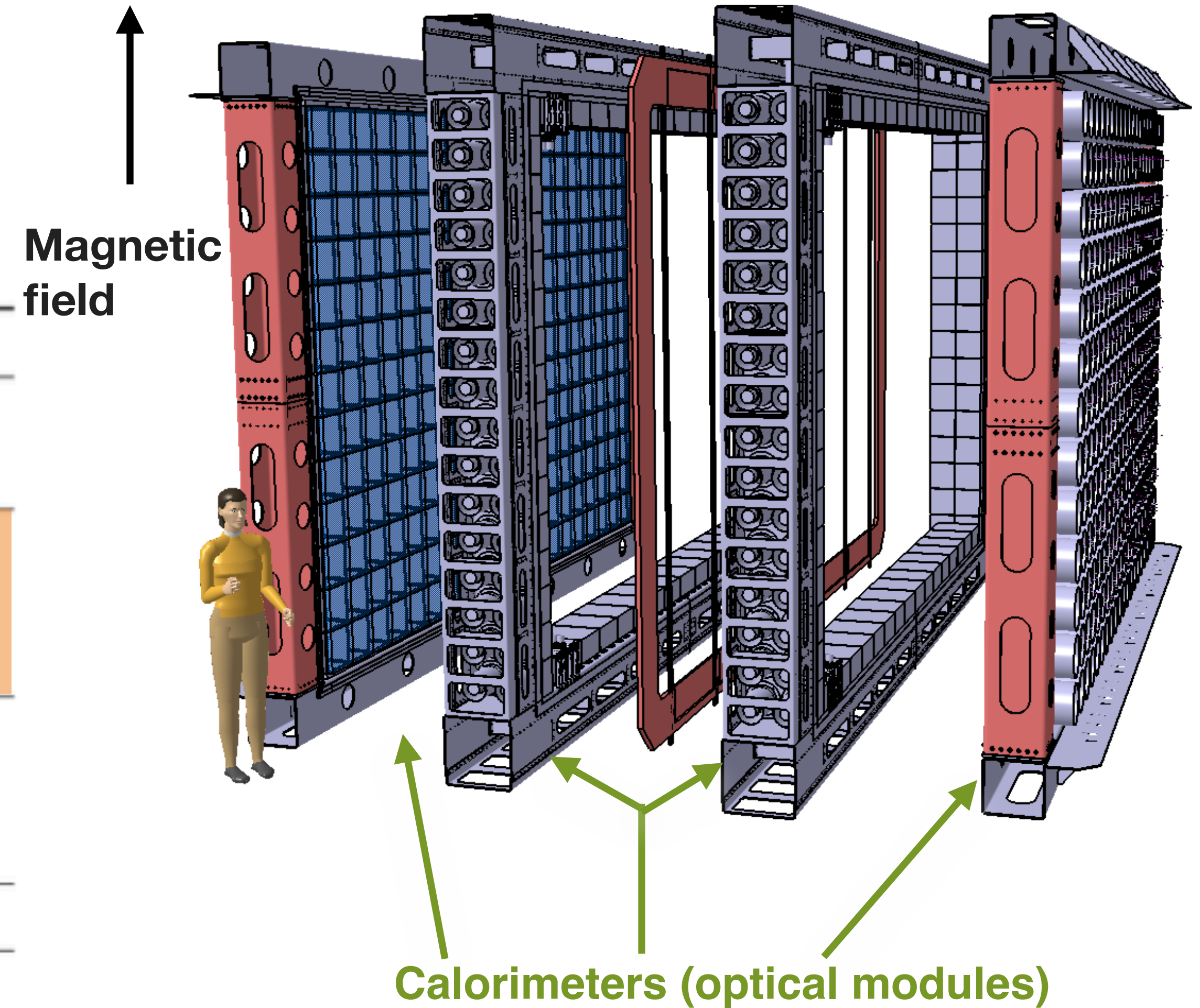
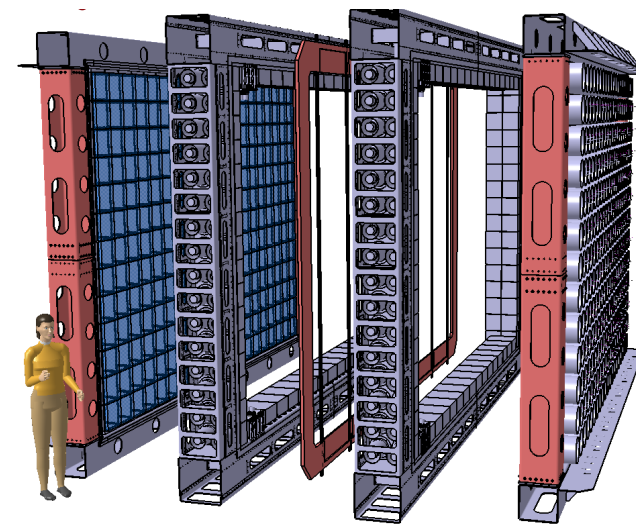
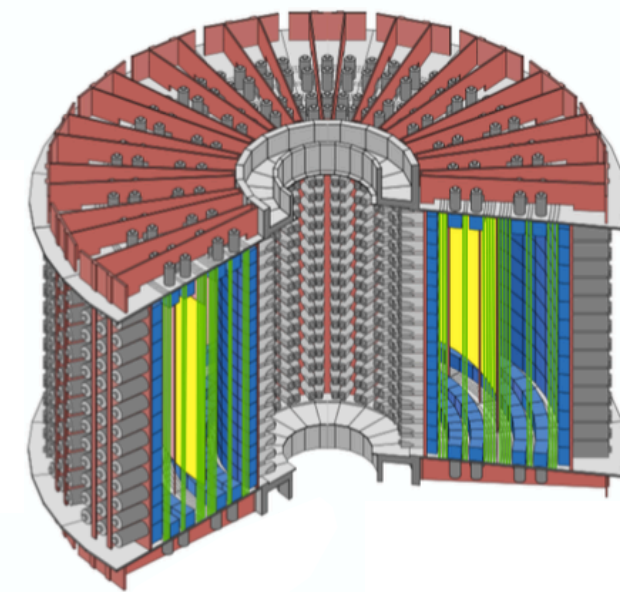
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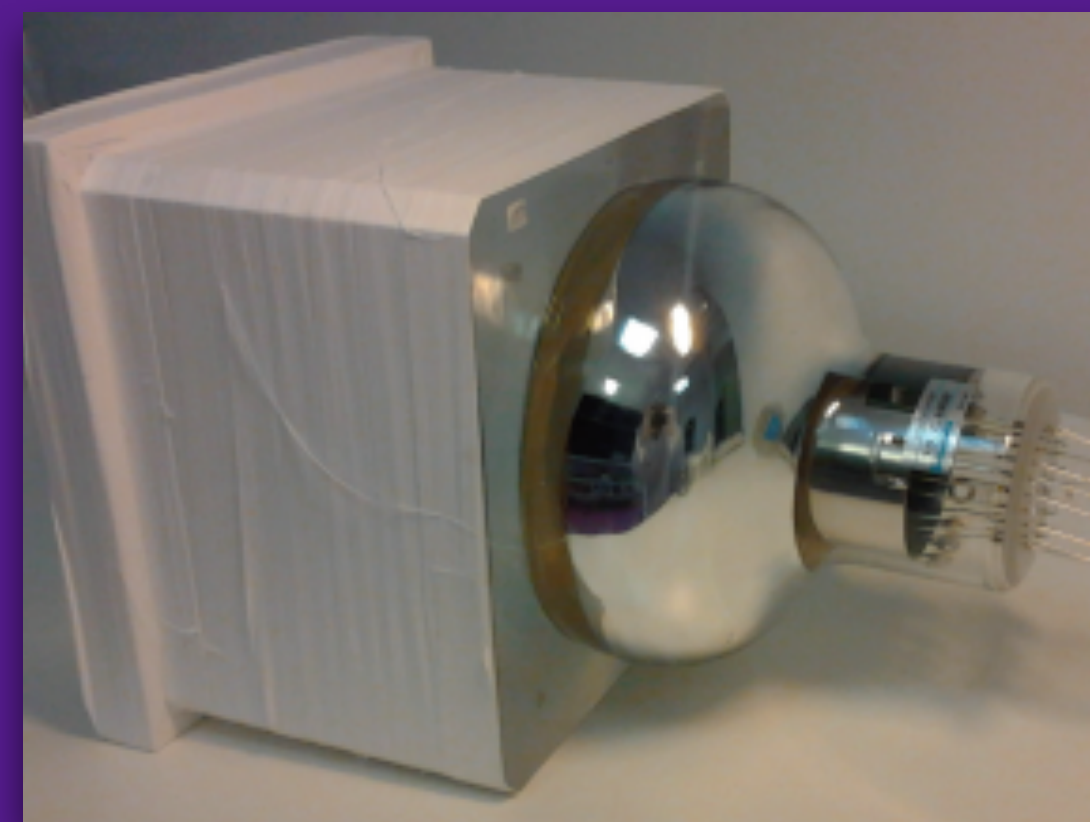
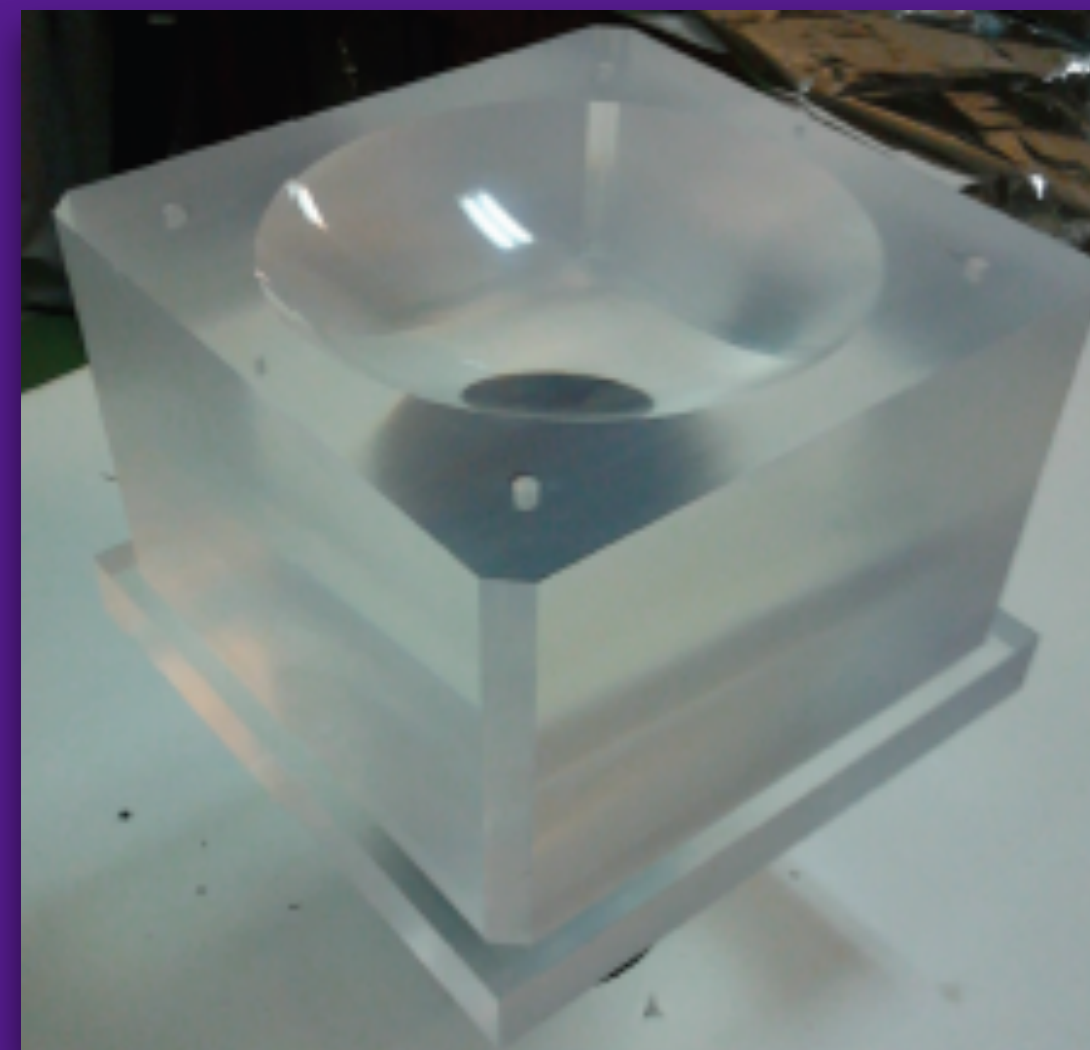
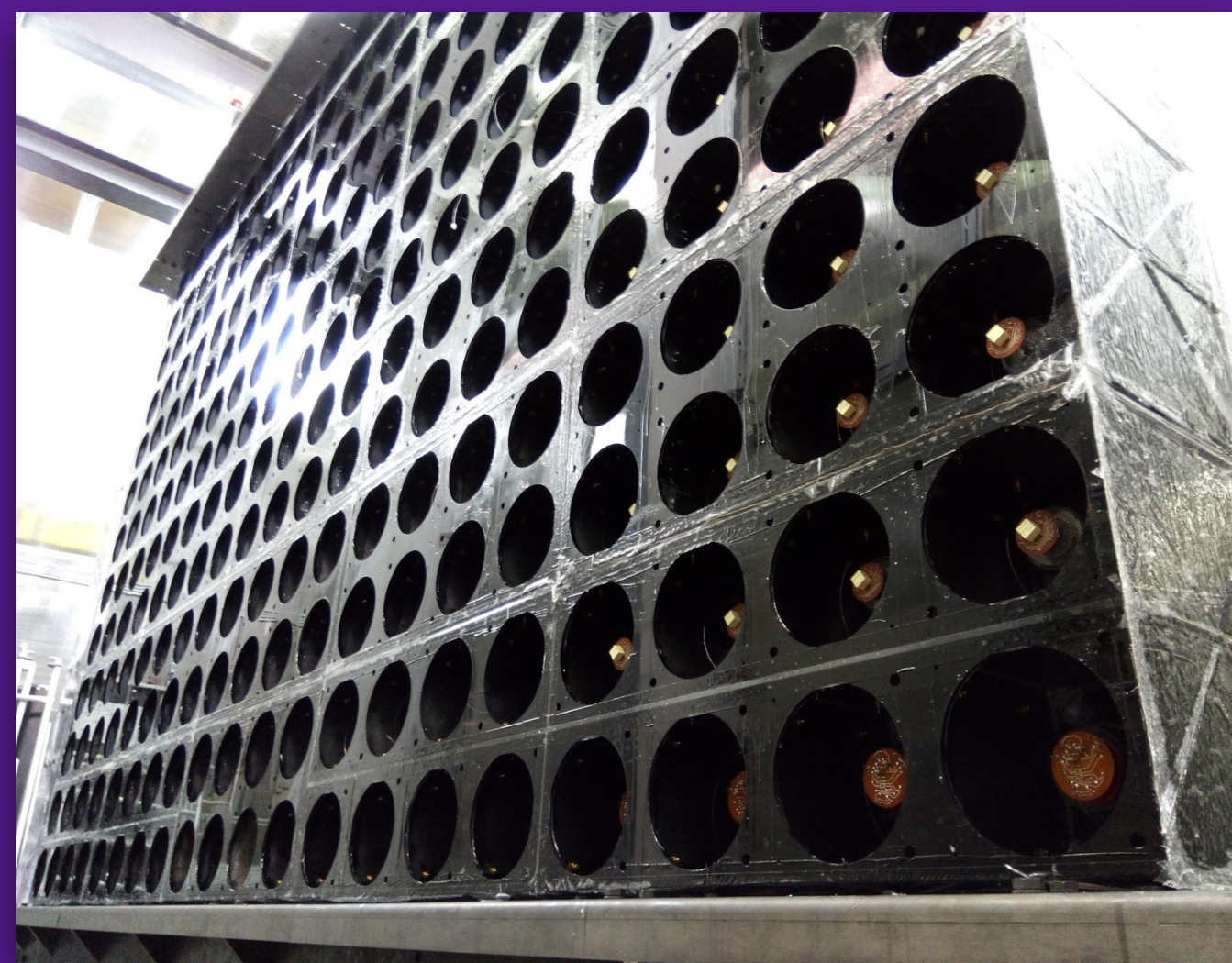
- 2034 drift cells (13,000 wires!)
- He/Ar/Ethanol gas mixture with advanced radon reduction program



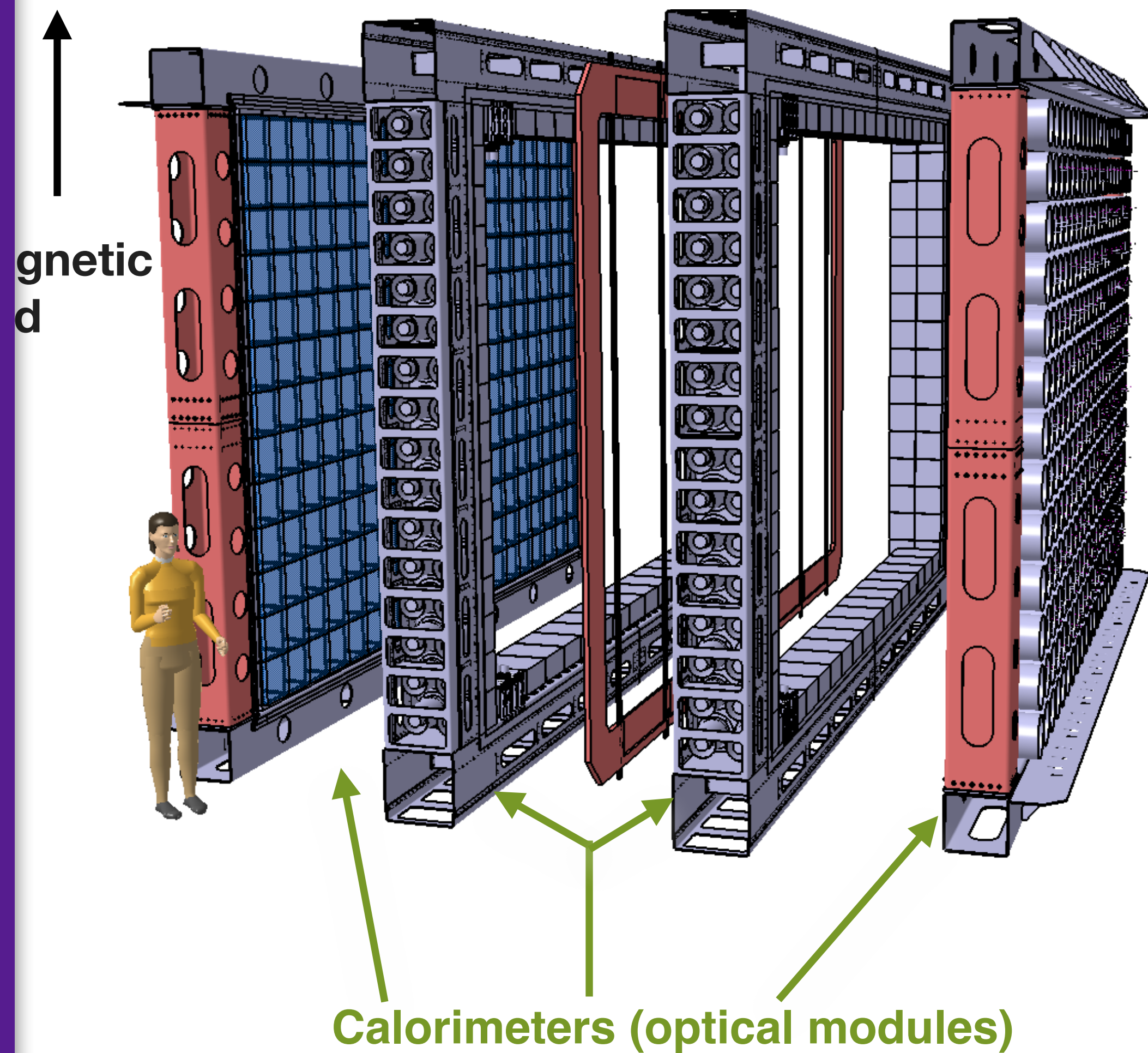
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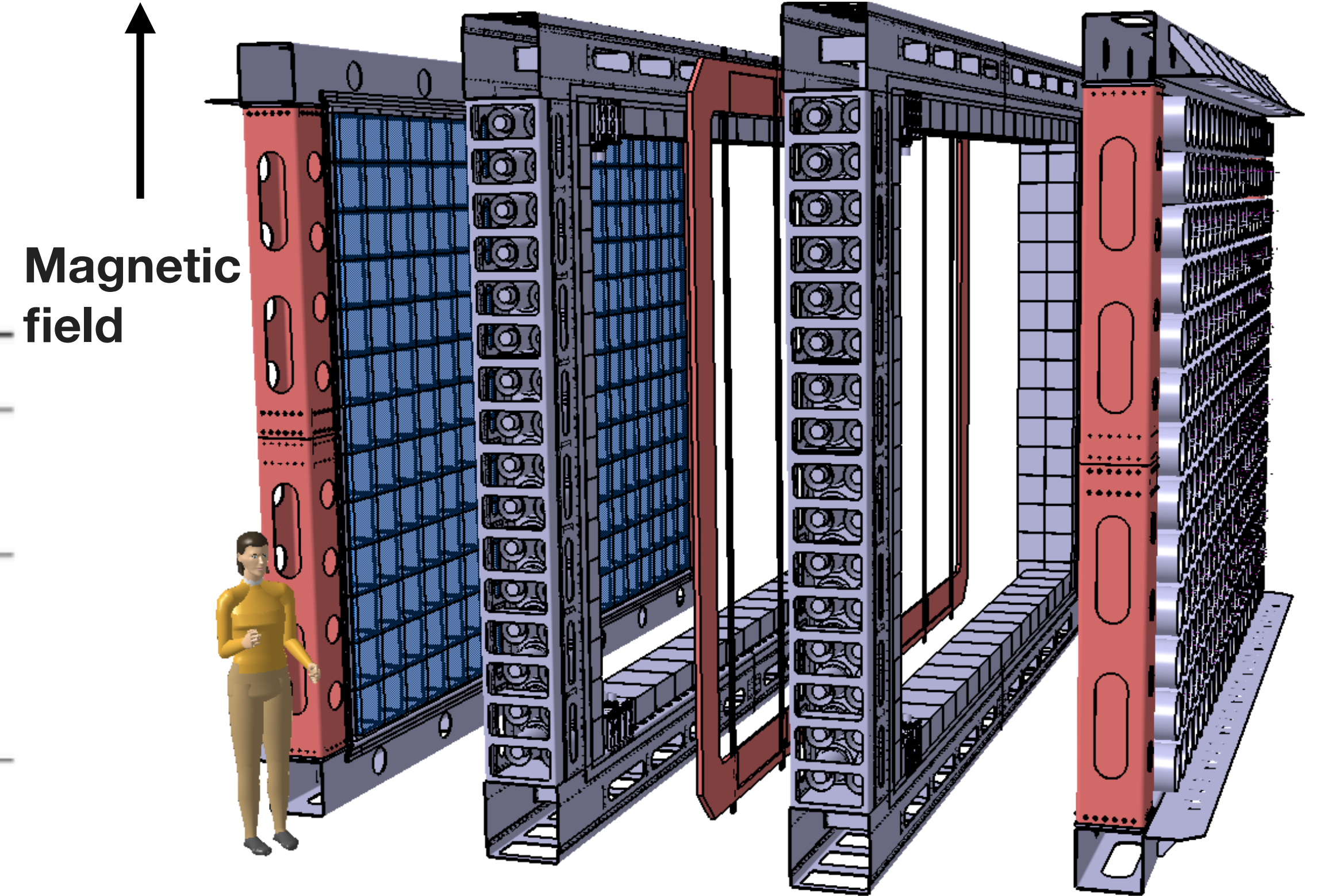
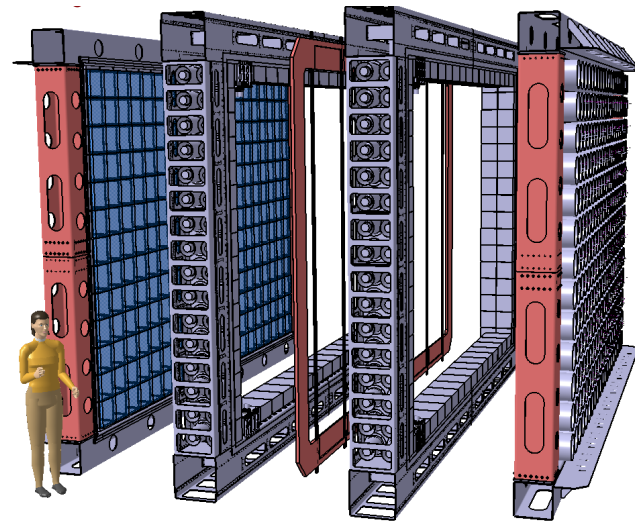
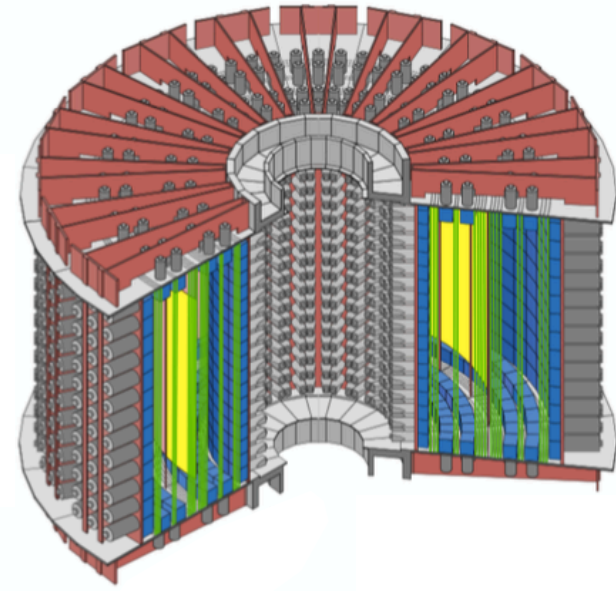
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- 440 8" radiopure PMTs (plus 5" NEMO-3 PMTs)
- Improved photocathode quantum efficiency
- Directly coupled to polystyrene scintillator (no light guide)
(*Nucl.Inst.Meth. A 868 98-108*)

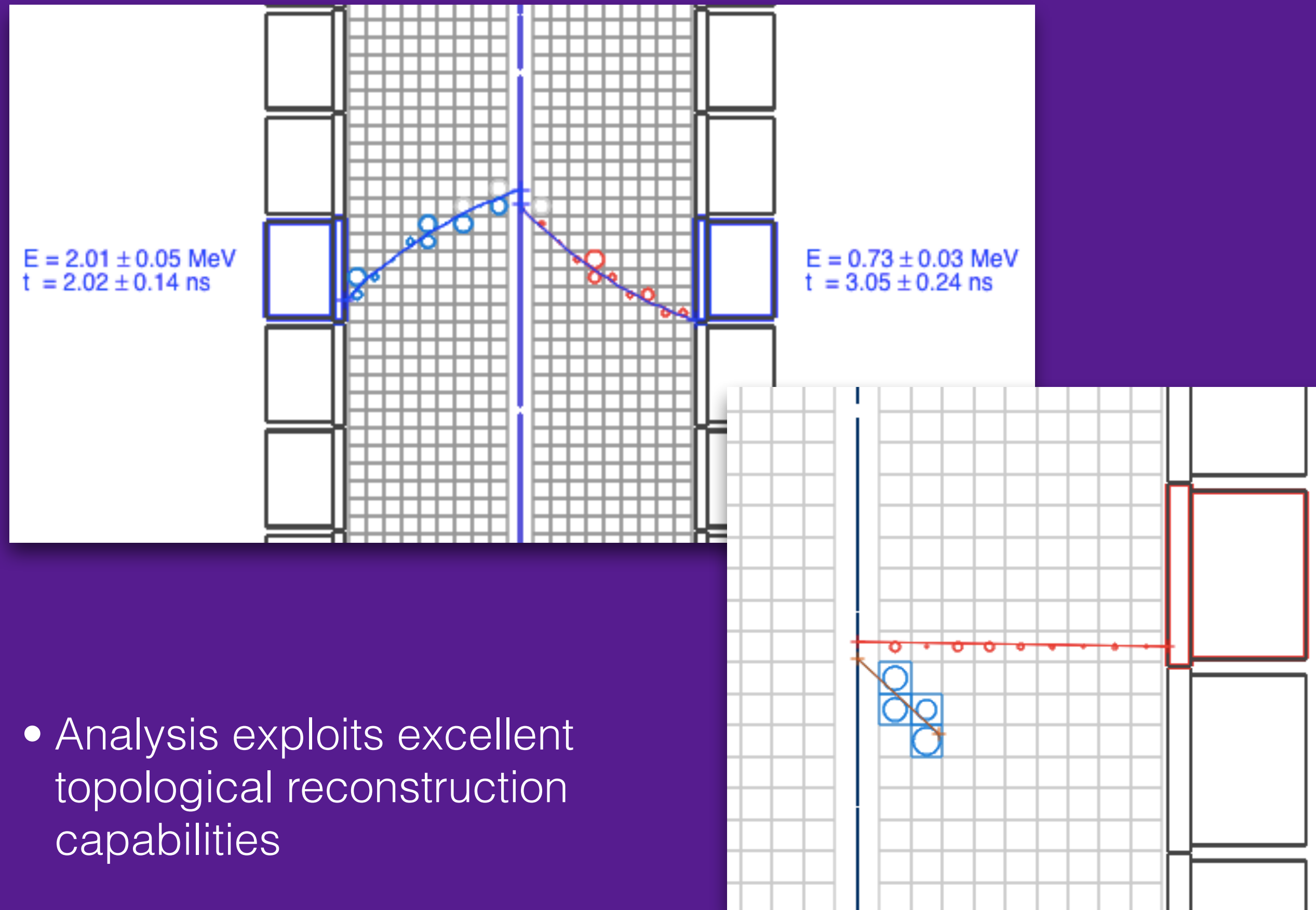


Now at LSM: SuperNEMO Demonstrator

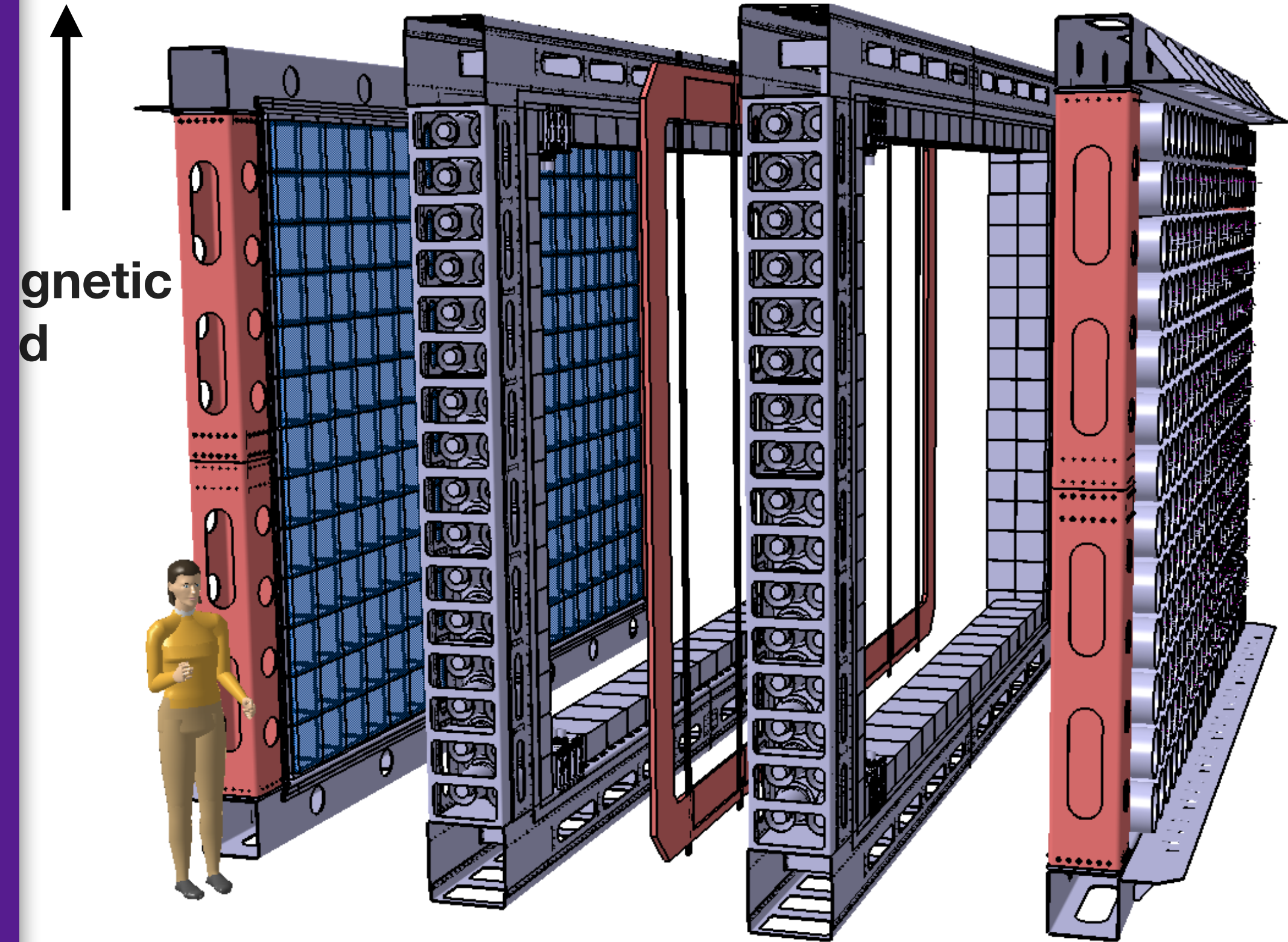


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$T_{1/2}^{0\nu} > 6 \times 10^{24} \text{ years}$

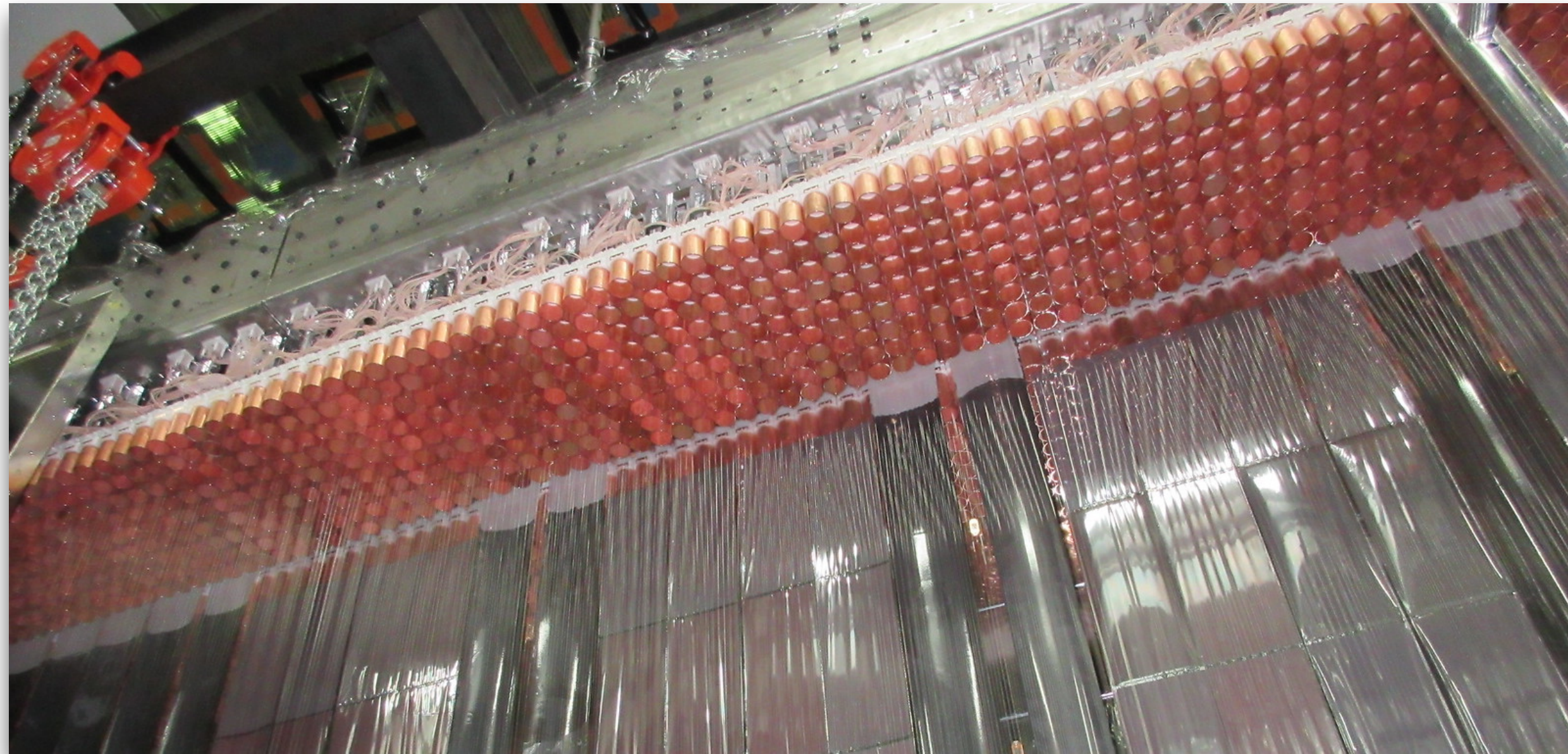


- Analysis exploits excellent topological reconstruction capabilities

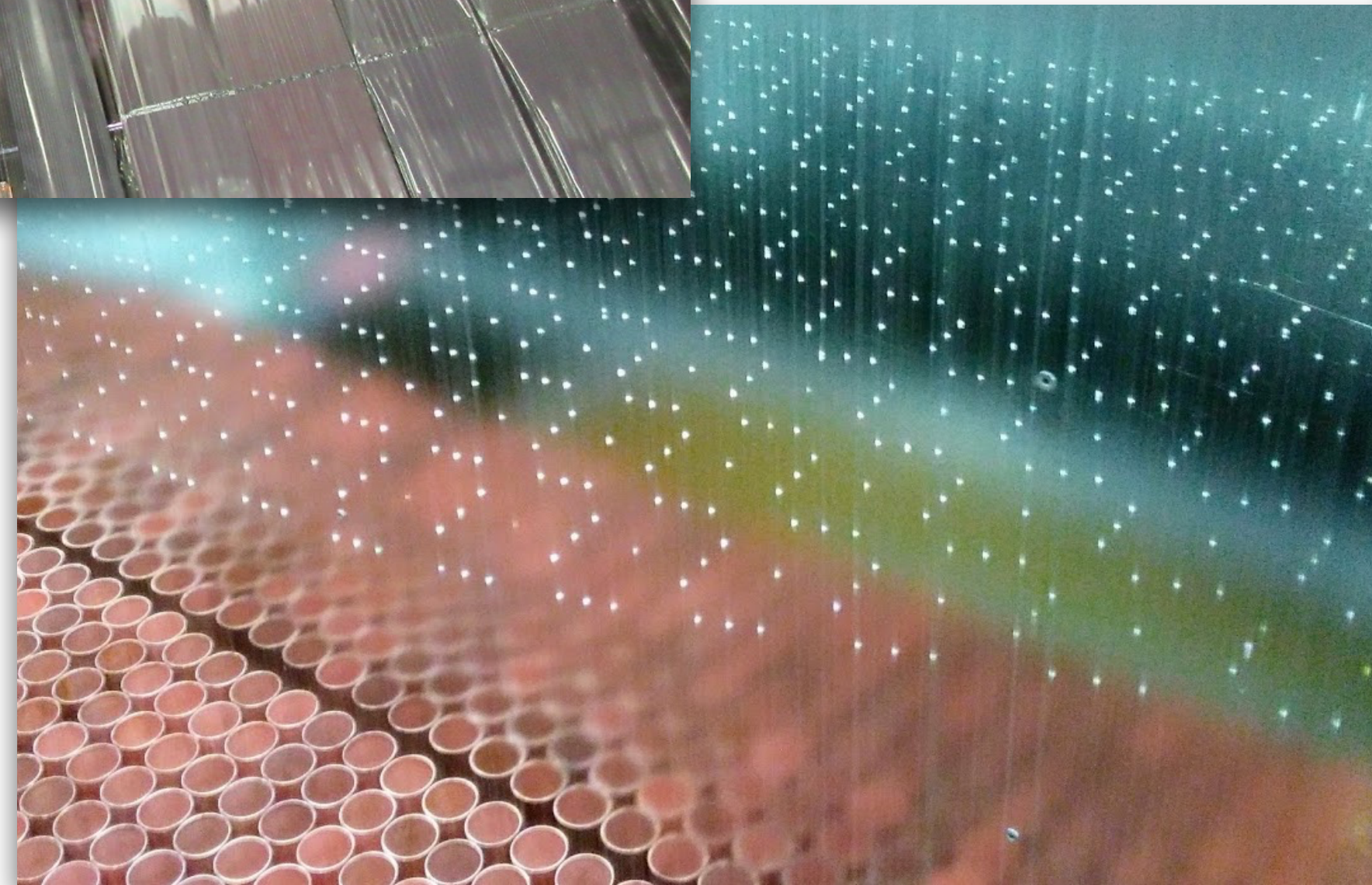


$T_{1/2}^{0\nu} > 6 \times 10^{24} \text{ years}$

Se

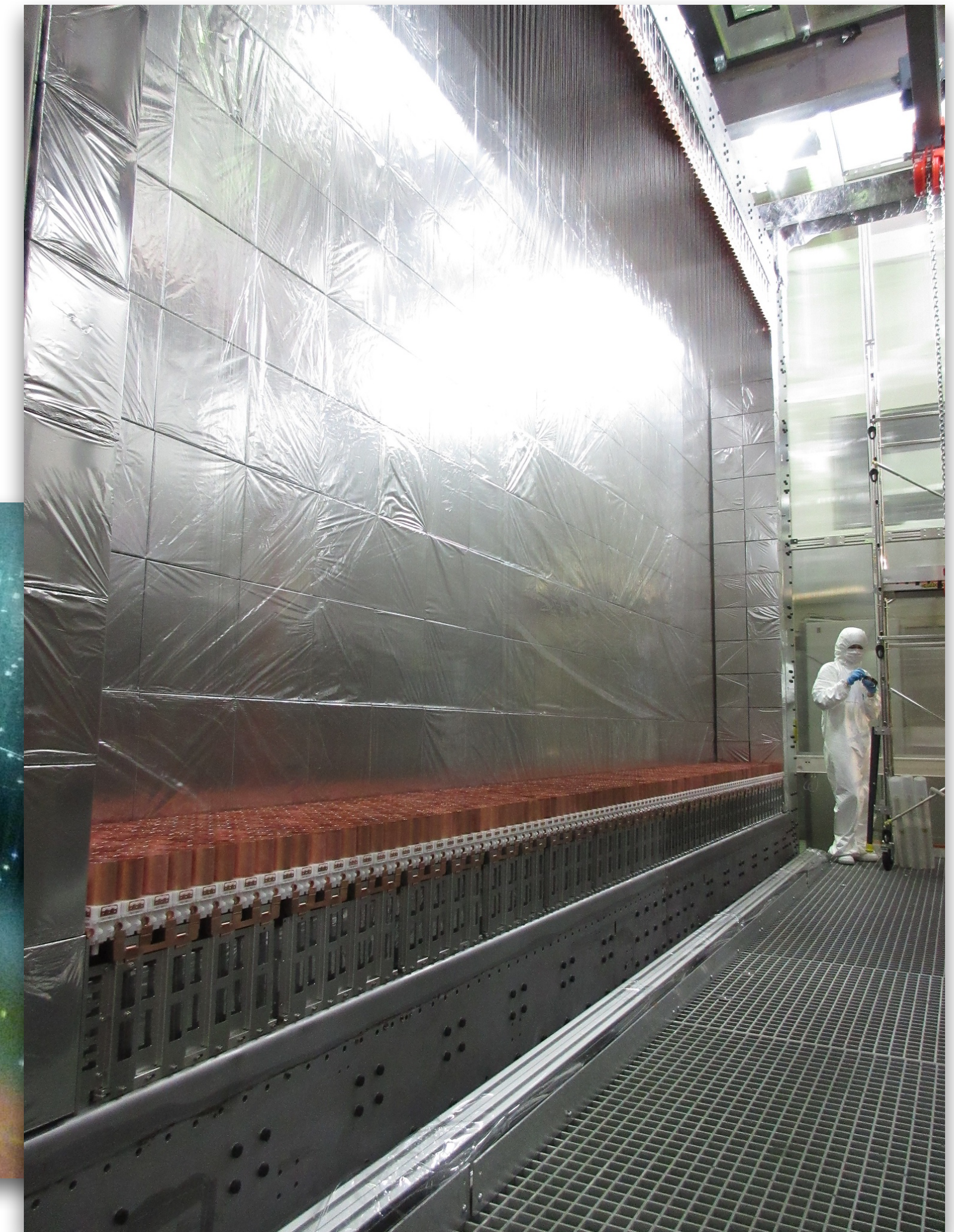


Source foils through the tracker



Tracker wire pattern

Calorimeter through the tracker

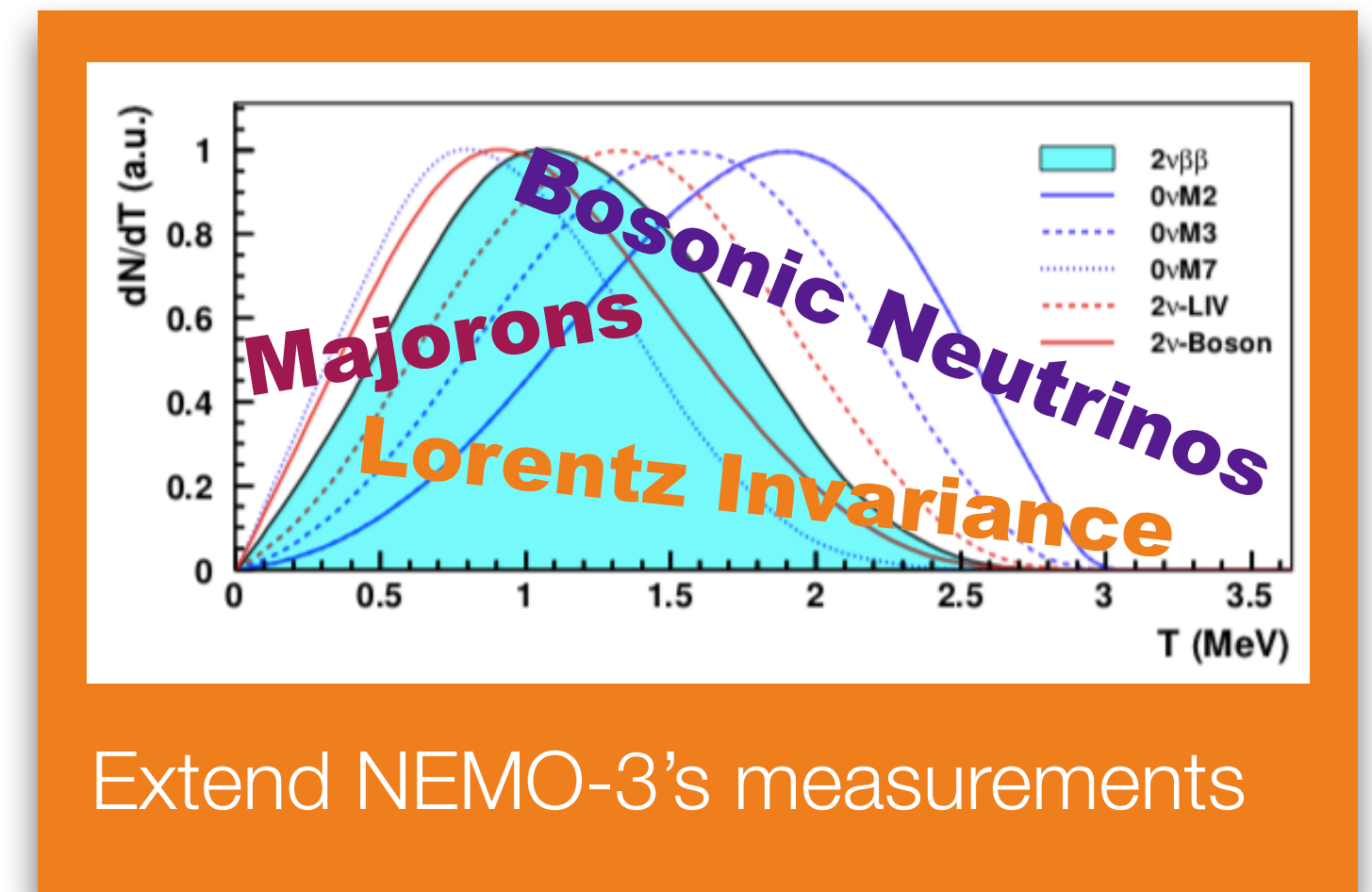


$0\nu\beta\beta$: $T_{1/2} > 6 \times 10^{24}$ years; $\langle m_\nu \rangle < 160\text{-}400$ meV

$0\nu\beta\beta$: $T_{1/2} > 6 \times 10^{24}$ years; $\langle m_\nu \rangle < 160-400$ meV

Exotic $0\nu\beta\beta$ mechanisms

Lorentz invariance violation test

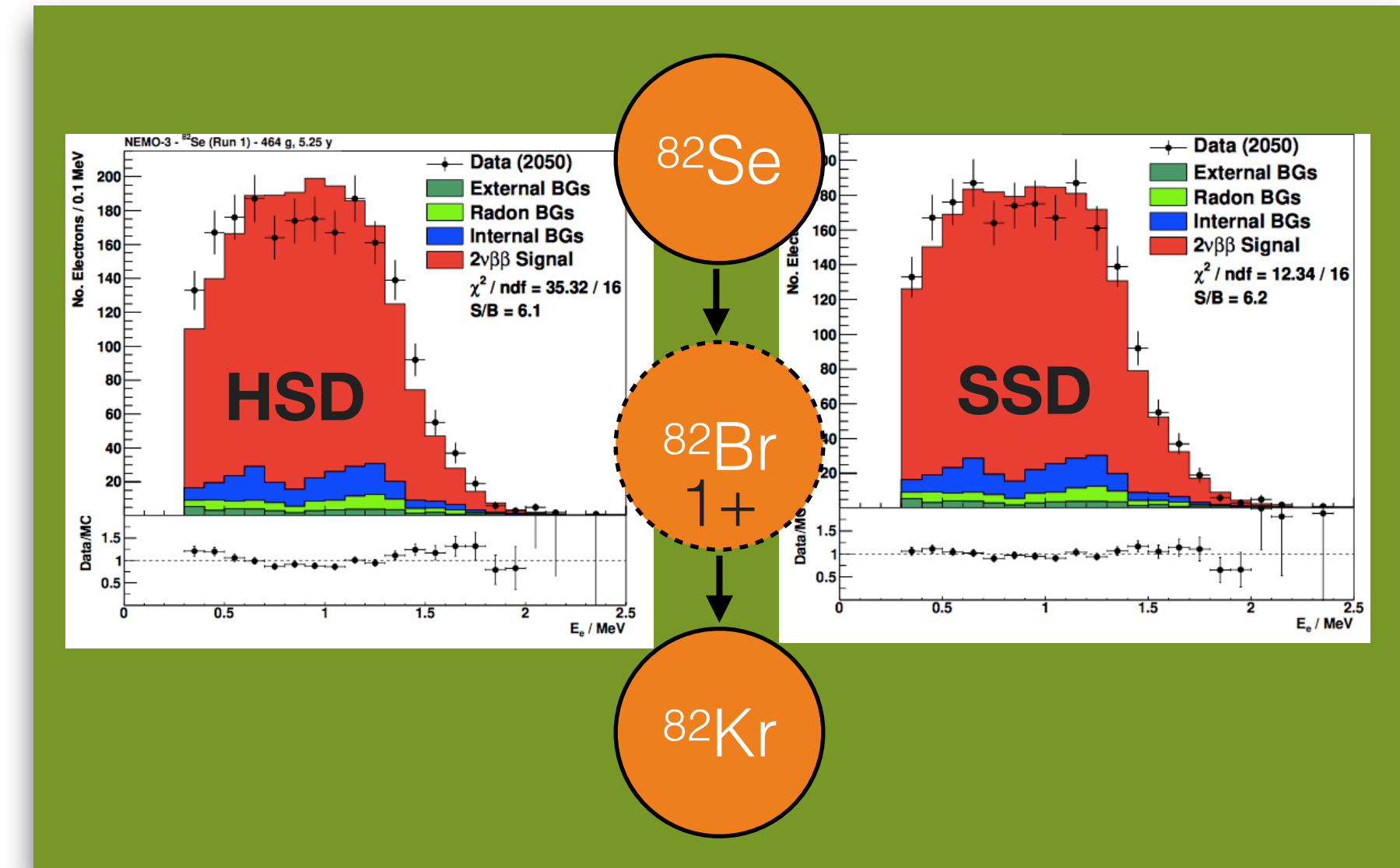


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$2\nu\beta\beta$: SSD/HSD discrimination at 5σ level



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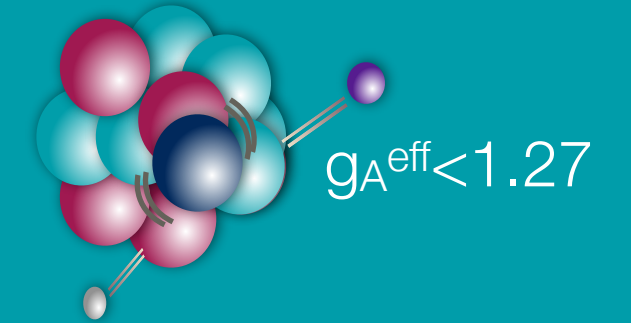
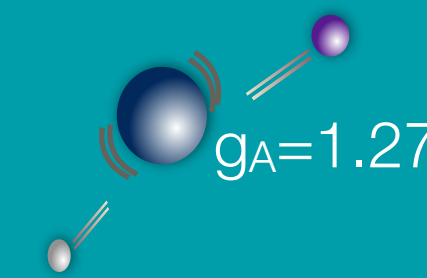
Exotic $0\nu\beta\beta$ mechanisms

Lorentz invariance violation test

$2\nu\beta\beta$: SSD/HSD discrimination at 5σ level

Probe nuclear physics by investigating g_A

- Axial-vector coupling constant g_A is **quenched** in heavy nuclei



- $2\nu\beta\beta$ rate proportional to g_A^4
- $$\left(T_{1/2}^{2\nu}\right)^{-1} = \left(g_A^{\text{eff}}\right)^4 \left|M_{GT}^{2\nu}\right|^2 G^{2\nu}$$
- New KamLAND-Zen paper investigates this quenching <https://arxiv.org/pdf/1901.03871.pdf>

Precision measurement of the ^{136}Xe two-neutrino $\beta\beta$ spectrum in KamLAND-Zen and its impact on the quenching of nuclear matrix elements

Gando,¹ Y. Gando,¹ T. Hachiya,¹ M. Ha Minh,¹ S. Hayashida,¹ Y. Honda,¹ K. Hosokawa,¹ H. Ikeda,¹ K. Inoue,¹ C. Ishidoshiro,¹ Y. Kamei,¹ K. Kamizawa,¹ T. Kinoshita,¹ M. Koga,^{1,2} S. Matsuda,¹ T. Mitsui,¹ K. Nakamura,^{1,2} Ono,¹ N. Ota,¹ S. Otsuka,¹ H. Ozaki,¹ Y. Shibukawa,¹ I. Shimizu,¹ Y. Shirahata,¹ J. Shirai,¹ T. Sato,¹ K. Soma,

- NEMO's topological capabilities mean it could do even **better!**

$0\nu\beta\beta$: $T_{1/2} > 6 \times 10^{24}$ years; $\langle m_\nu \rangle < 160\text{-}400$ meV

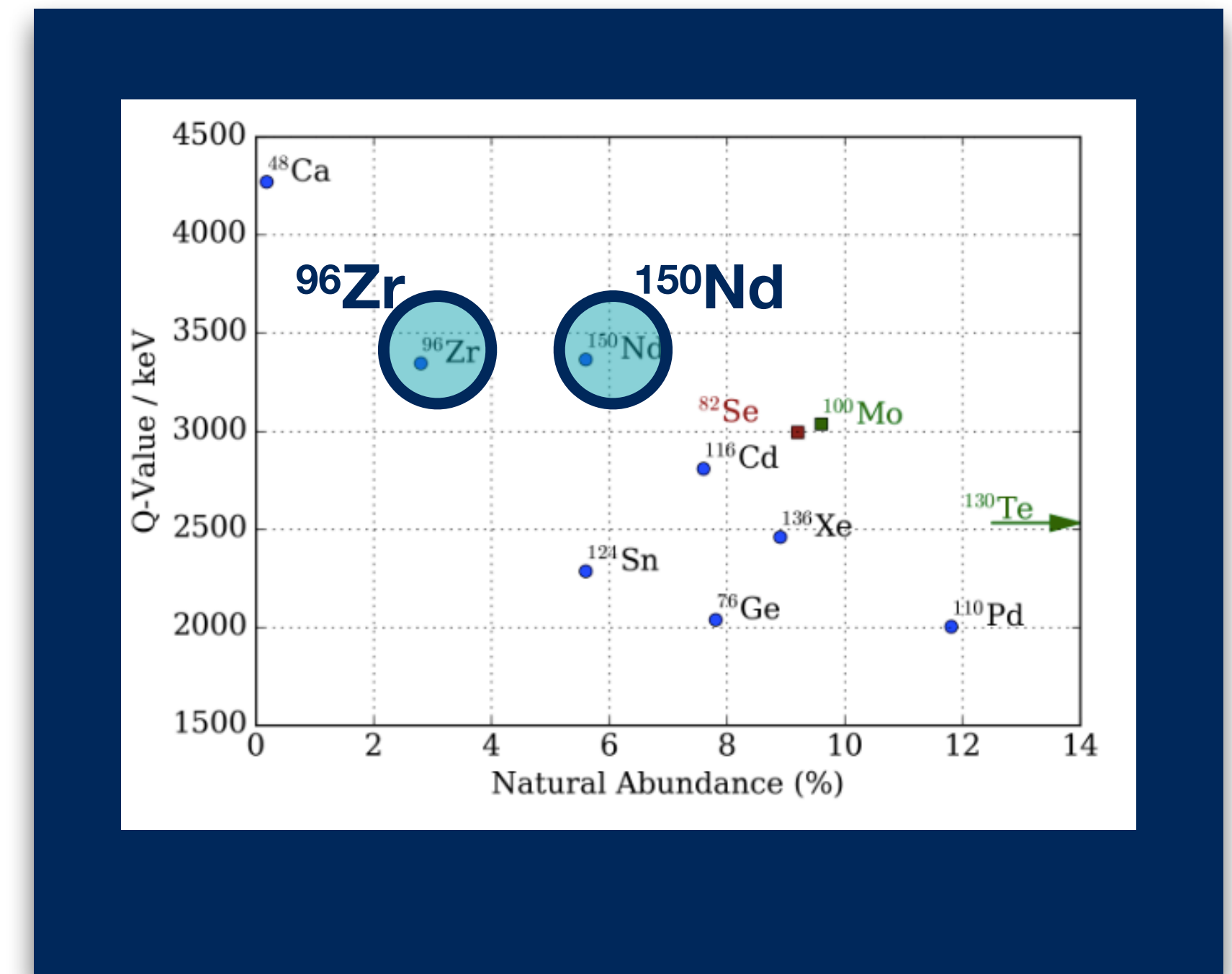
Exotic $0\nu\beta\beta$ mechanisms

Lorentz invariance violation test

$2\nu\beta\beta$: SSD/HSD discrimination at 5σ level

Probe nuclear physics by investigating g_A

Alternative isotopes: ^{150}Nd and ^{96}Zr



$0\nu\beta\beta$: $T_{1/2} > 6 \times 10^{24}$ years; $\langle m_\nu \rangle < 160\text{-}400$ meV

Exotic $0\nu\beta\beta$ mechanisms

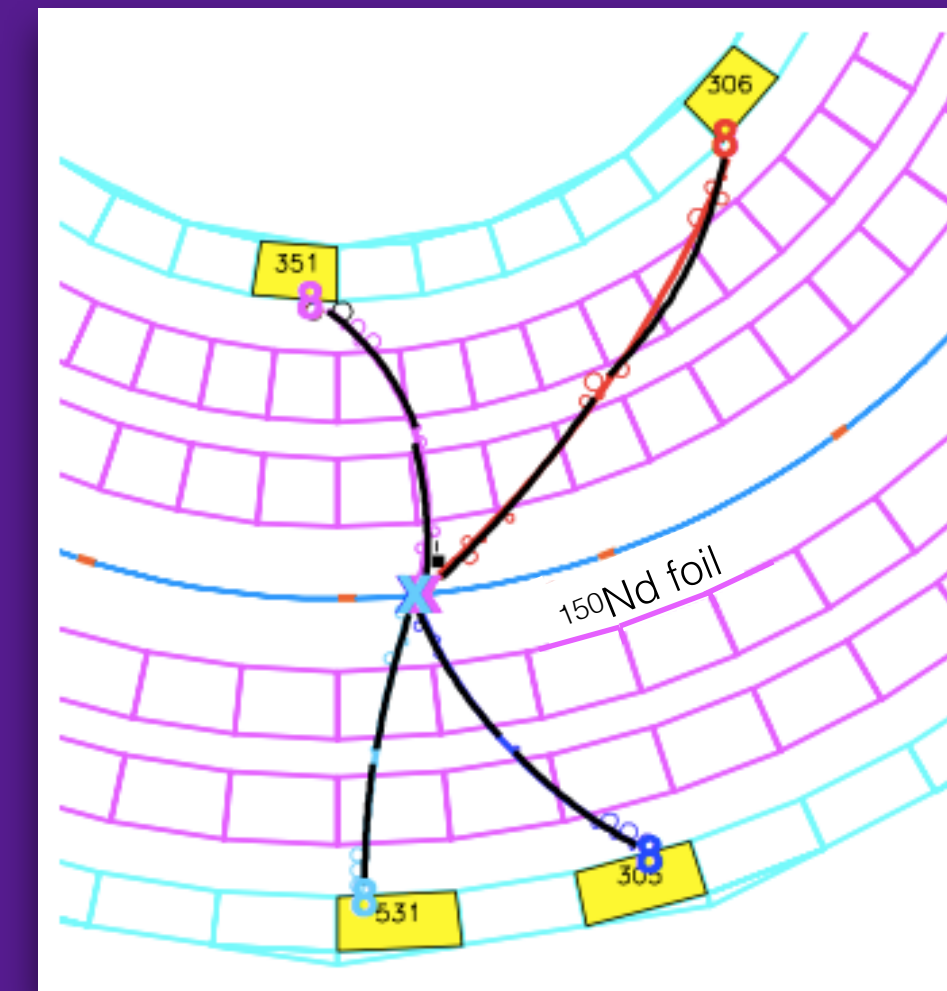
Lorentz invariance violation test

$2\nu\beta\beta$: SSD/HSD discrimination at 5σ level

Probe nuclear physics by investigating g_A

Alternative isotopes: ^{150}Nd and ^{96}Zr

$0\nu 4\beta$: for ^{150}Nd



NEMO-3 placed limit on lepton number-violating process, which could affect even Dirac neutrinos
Phys. Rev. Lett. 119, 041801

$0\nu\beta\beta$: $T_{1/2} > 6 \times 10^{24}$ years; $\langle m_\nu \rangle < 160-400$ meV

Exotic $0\nu\beta\beta$ mechanisms

Lorentz invariance violation test

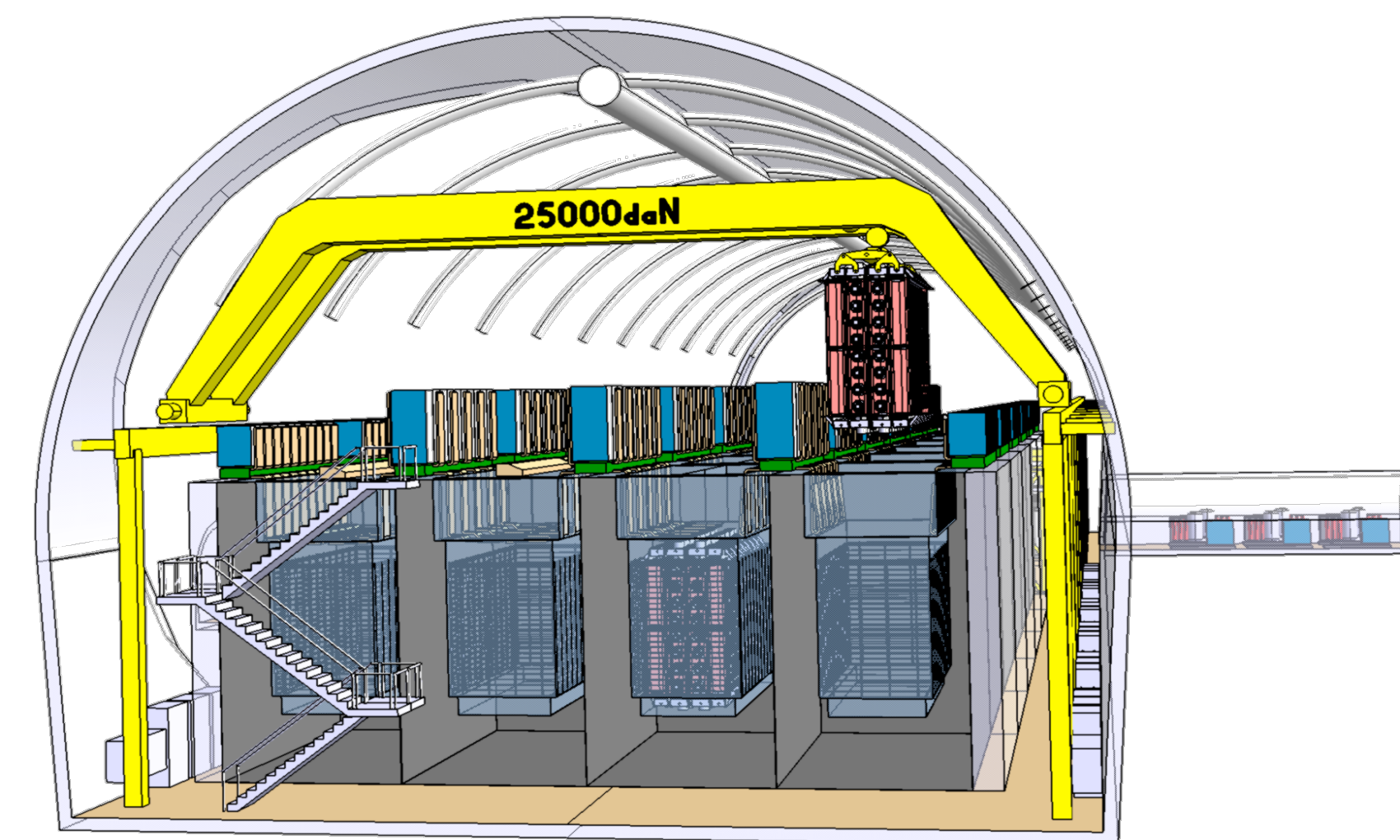
$2\nu\beta\beta$: SSD/HSD discrimination at 5σ level

Probe nuclear physics by investigating g_A

Alternative isotopes: ^{150}Nd and ^{96}Zr

$0\nu4\beta$: for ^{150}Nd

plus proof of concept for...

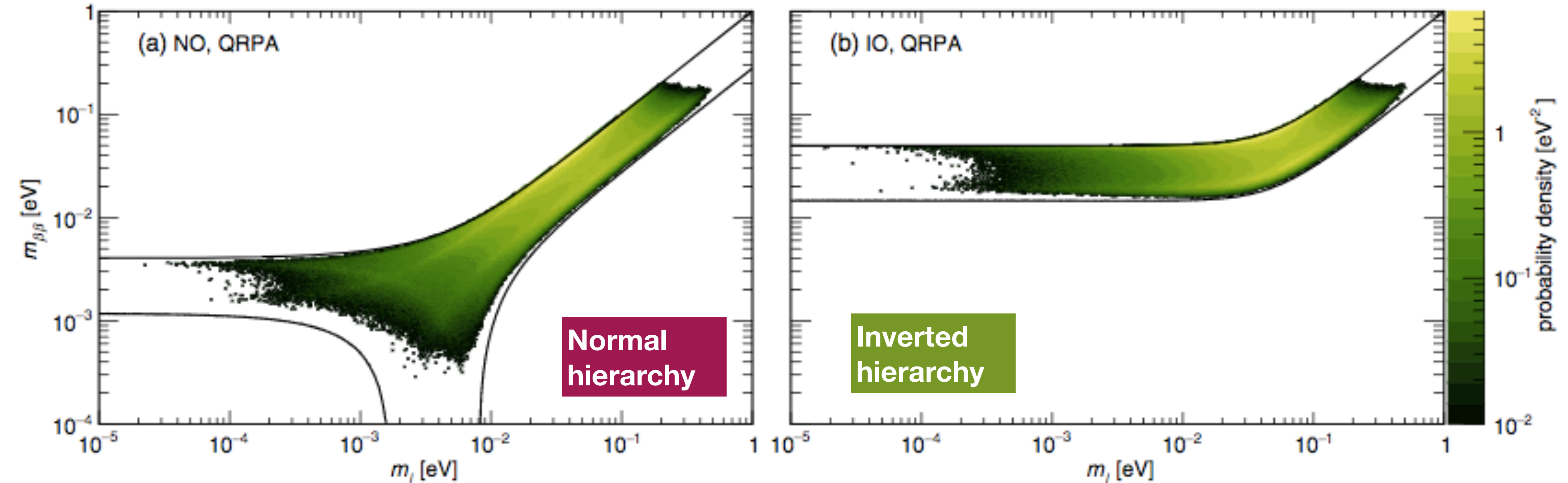


Full SuperNEMO

- **Modular** design allows easy scaling up
- 20 modules x 5 years (500 kg year) gives sensitivity comparable or better than current **leading experiments**
- Best technique to understand more about **$0\nu\beta\beta$ mechanism** in the event of discovery

Look to the future...

Bayesian probability density fit by D'Agostini, Benato & Detwiler: *Phys. Rev. D* 96, 053001 (2017)



Known parameters:

- Mass splittings $\Delta m^2_{12}, \Delta m^2_{13}$ (NO) or Δm^2_{13} (IO)
- Mixing angles θ_{12}, θ_{13}

Thrown from a flat prior:

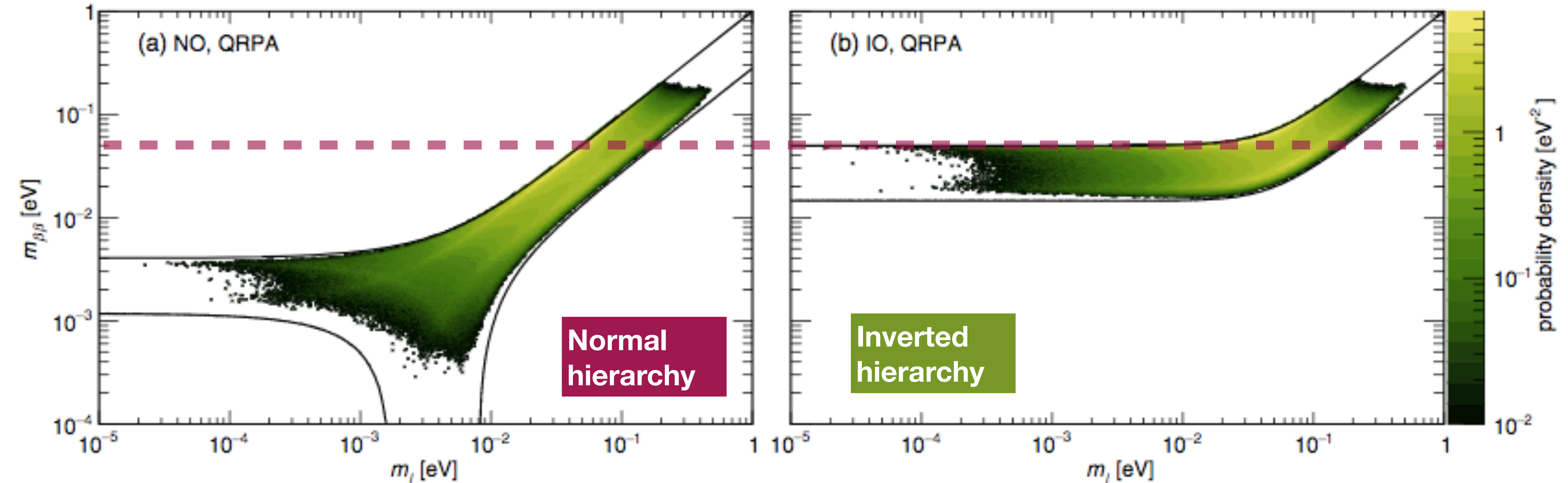
- Majorana phases α_1, α_2
- CP violating phase δ

Thrown from a logarithmic prior:

- Sum of neutrino masses Σ (upper limit from cosmology)

Next-generation $0\nu\beta\beta$ searches

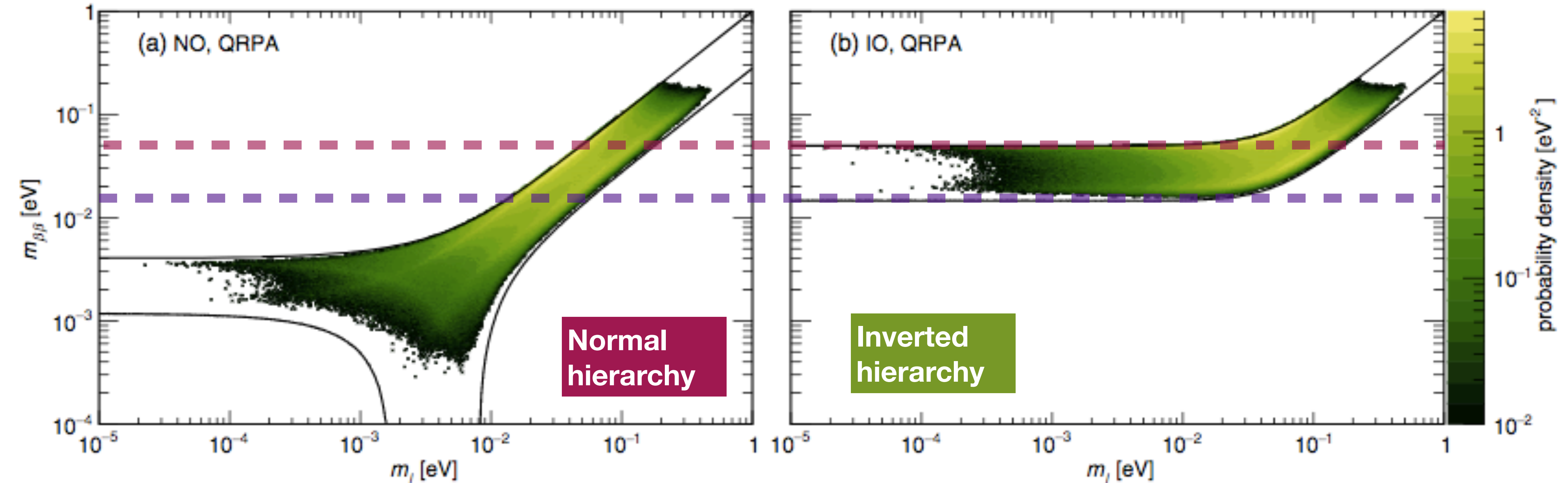
Bayesian probability density fit by D'Agostini, Benato & Detwiler: *Phys. Rev. D* 96, 053001 (2017)



- Current experiments probe the **degenerate** regime

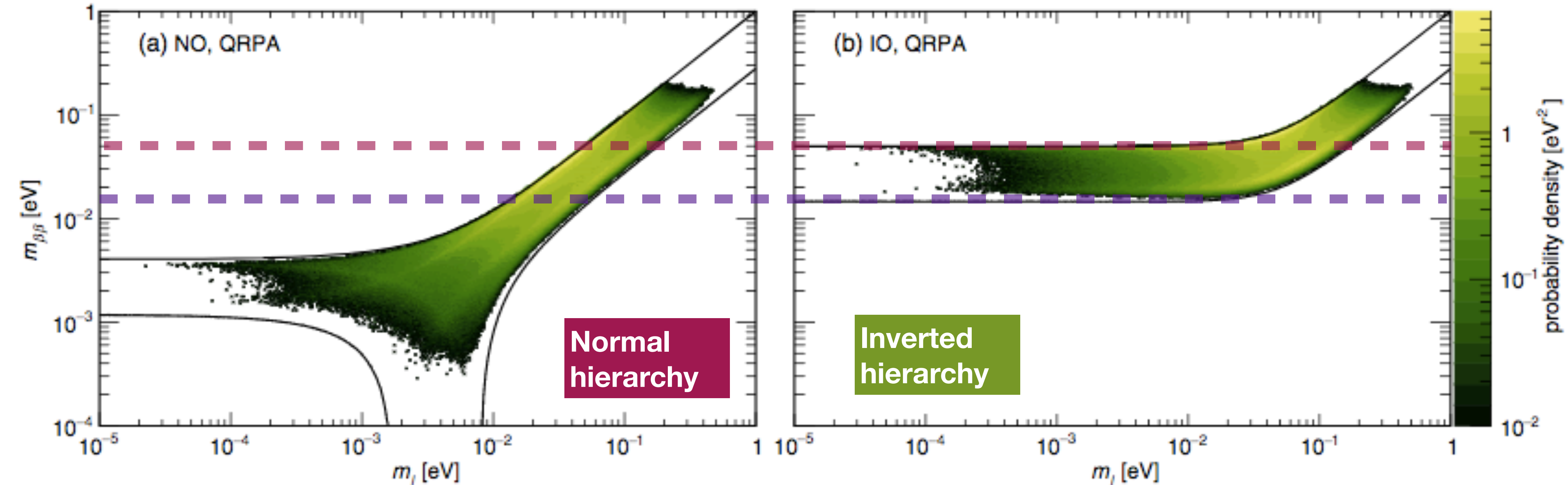
Next-generation $0\nu\beta\beta$ searches

Bayesian probability density fit by D'Agostini, Benato & Detwiler: *Phys. Rev. D* 96, 053001 (2017)



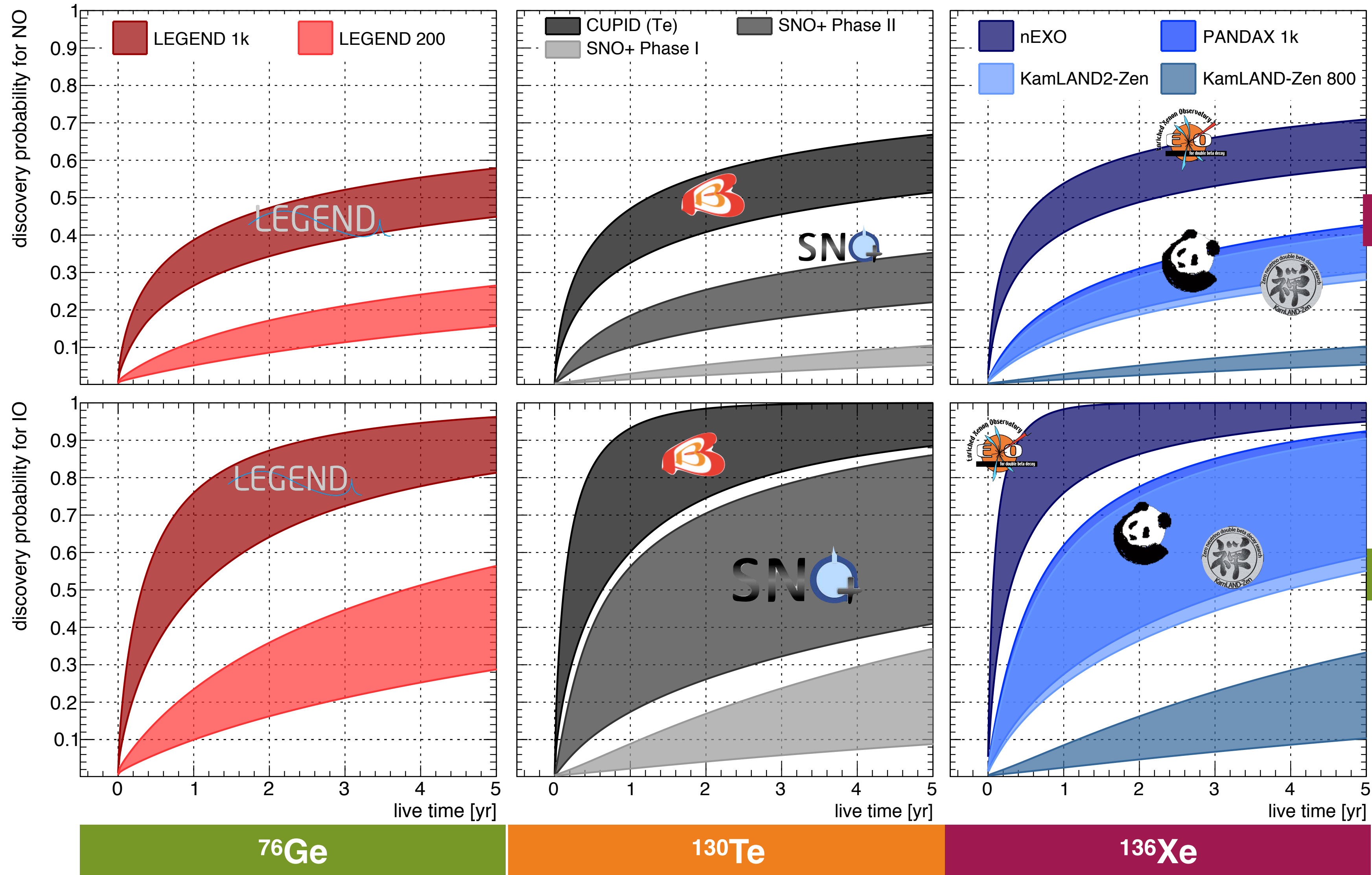
- Current experiments probe the **degenerate** regime
- Next-generation will cover **full inverted hierarchy** region

Bayesian probability density fit by D'Agostini, Benato & Detwiler: *Phys. Rev. D* 96, 053001 (2017)



- Current experiments probe the **degenerate** regime
- Next-generation will cover **full inverted hierarchy** region
- When likelihood density is considered, this mass range also covers more than **50% of normal hierarchy** probability

Will the next generation find it?

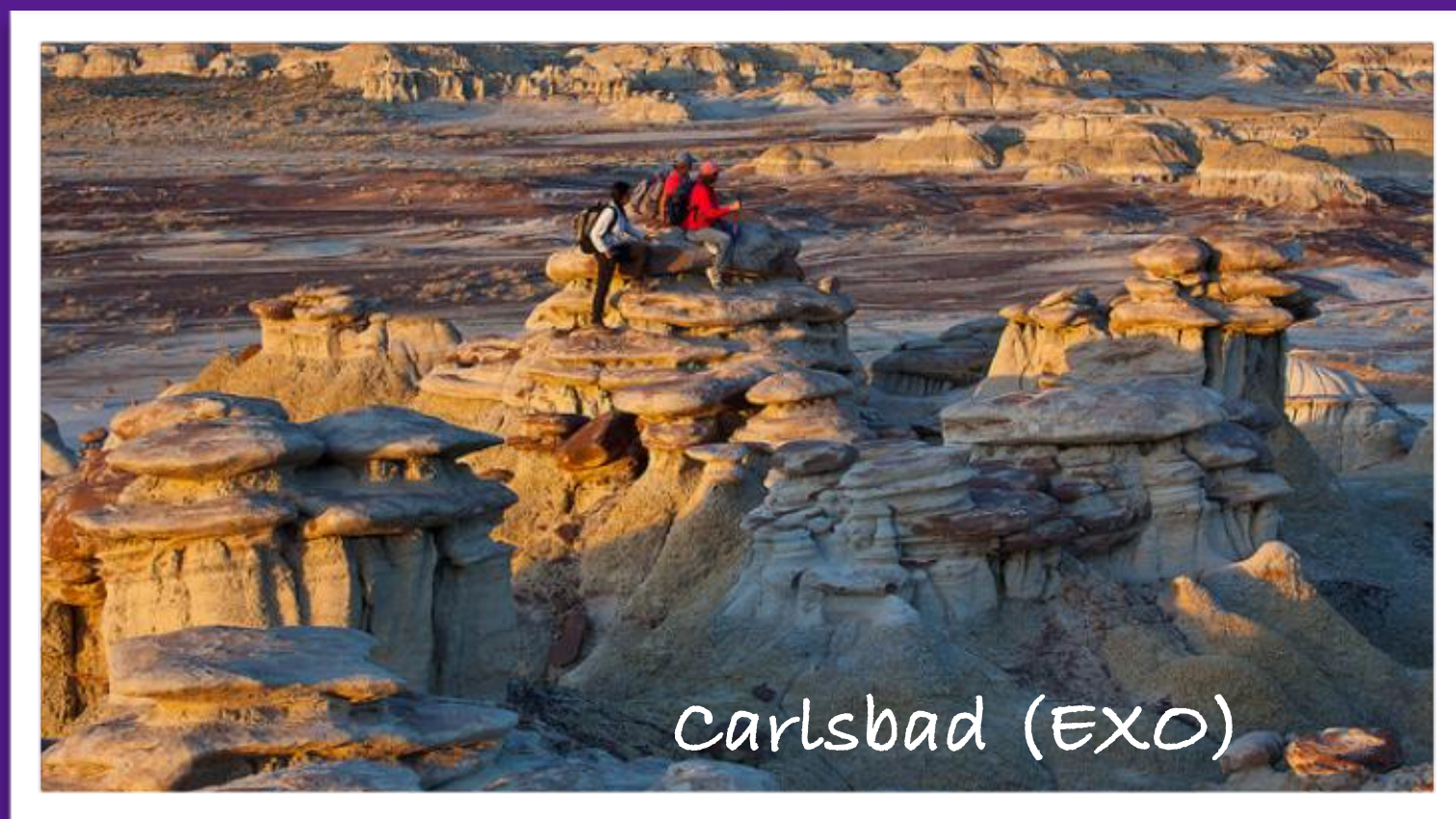


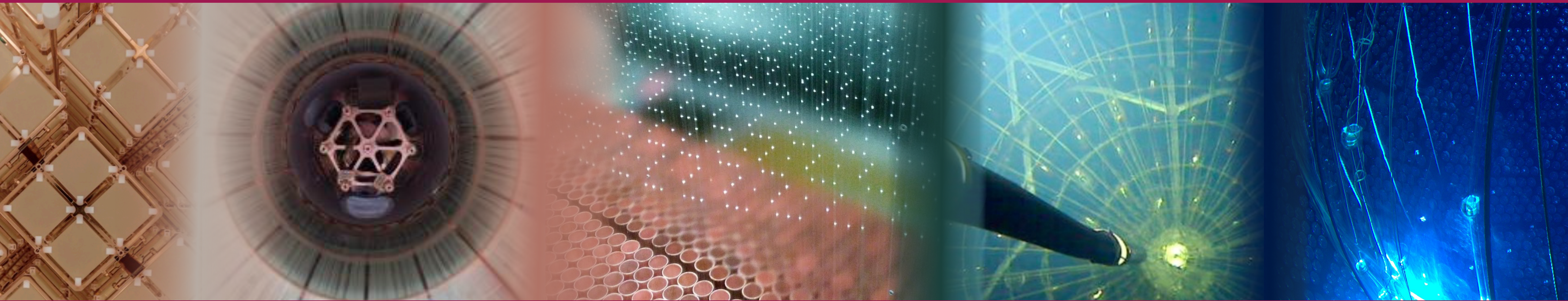
Normal hierarchy

Inverted hierarchy

Phys. Rev. D 96, 053001 (2017)

Postcards from double-beta land





Did you spot the pictures?



CUORE

Did you spot the pictures?



CUORE

GERDA

Did you spot the pictures?



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