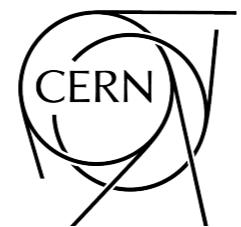


Solar and Reactor Neutrinos

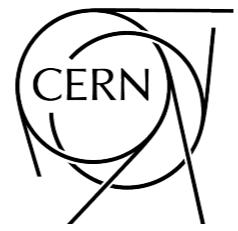
Joachim Kopp | CERN & JGU Mainz | Lectures at INSS 2019, Fermilab



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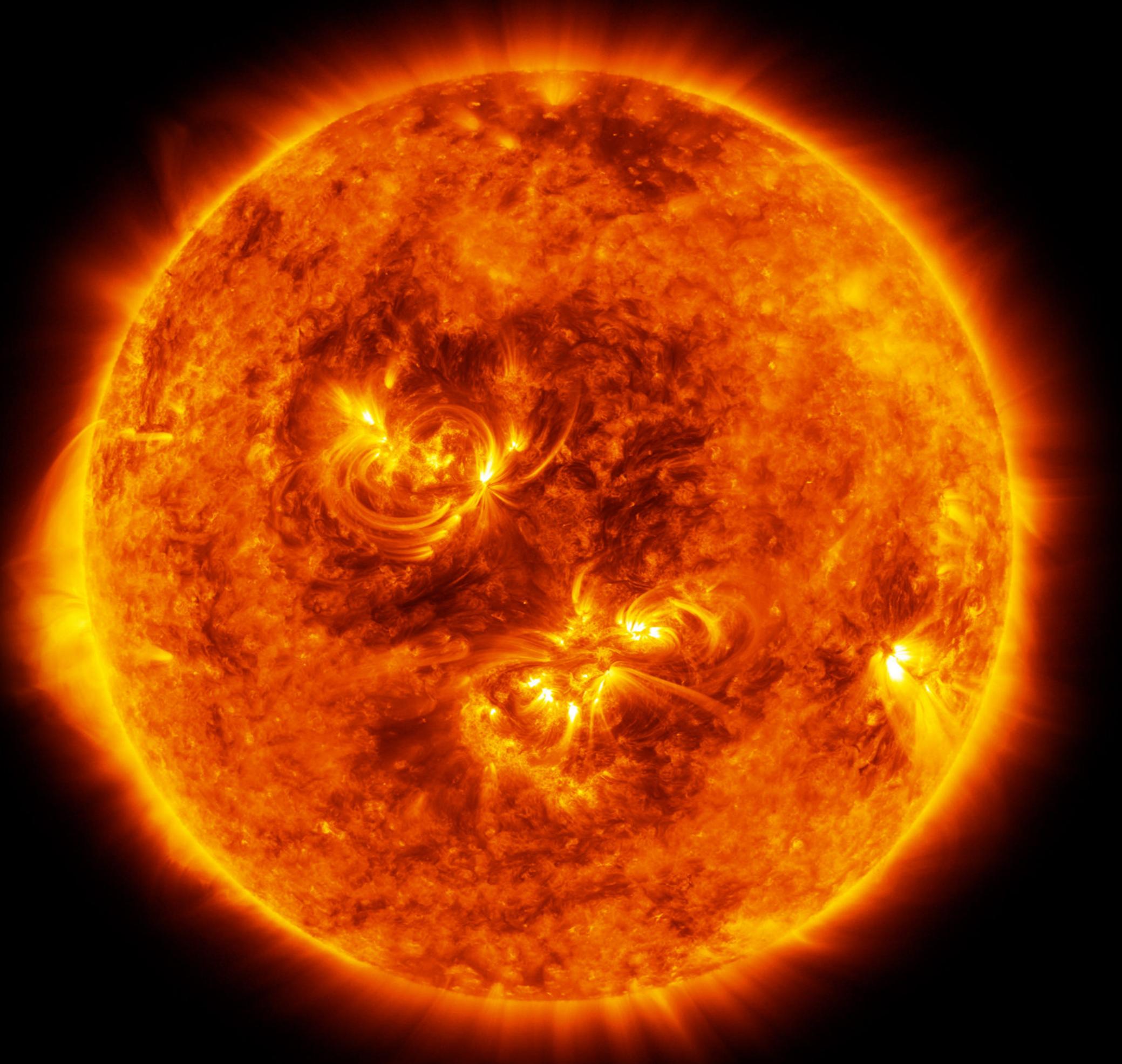


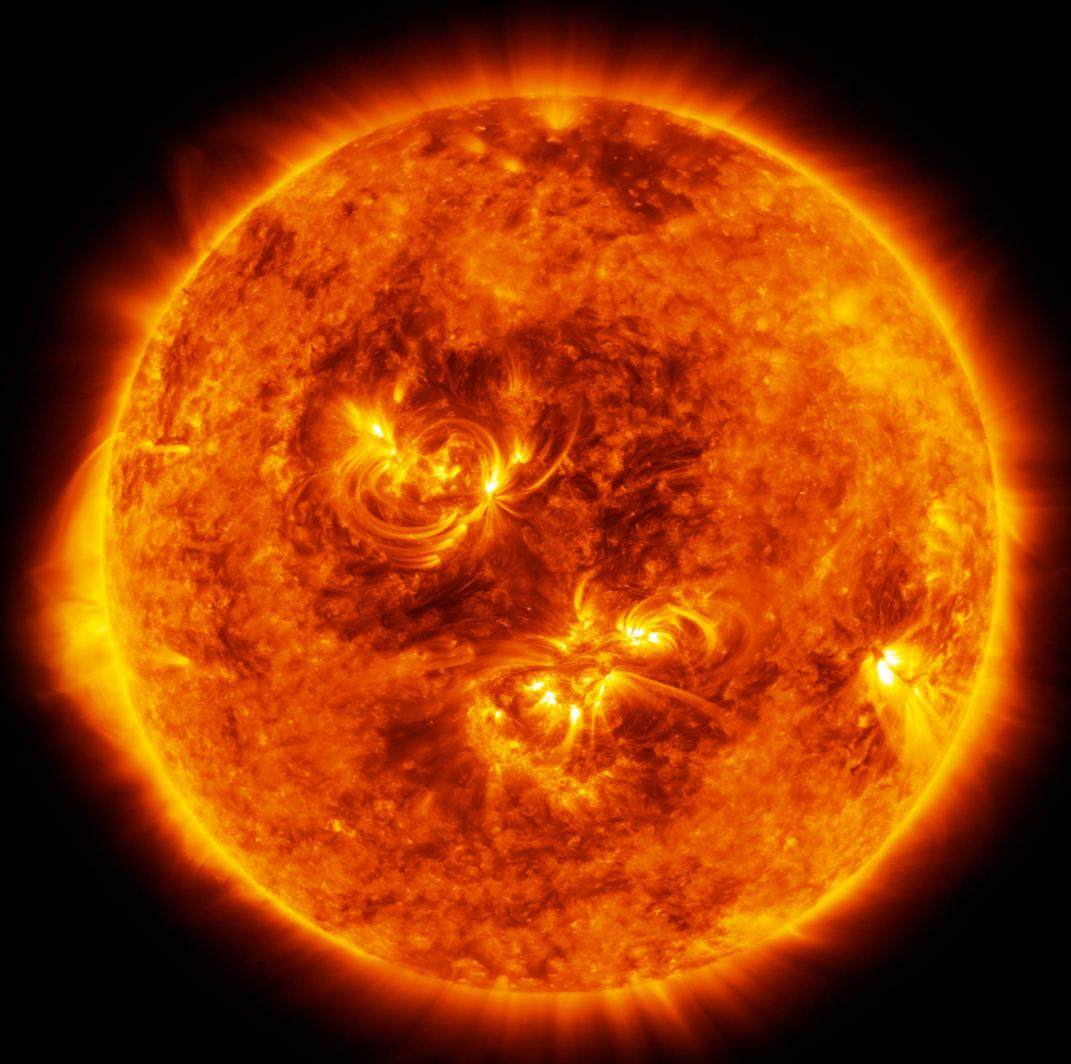
Solar Neutrinos

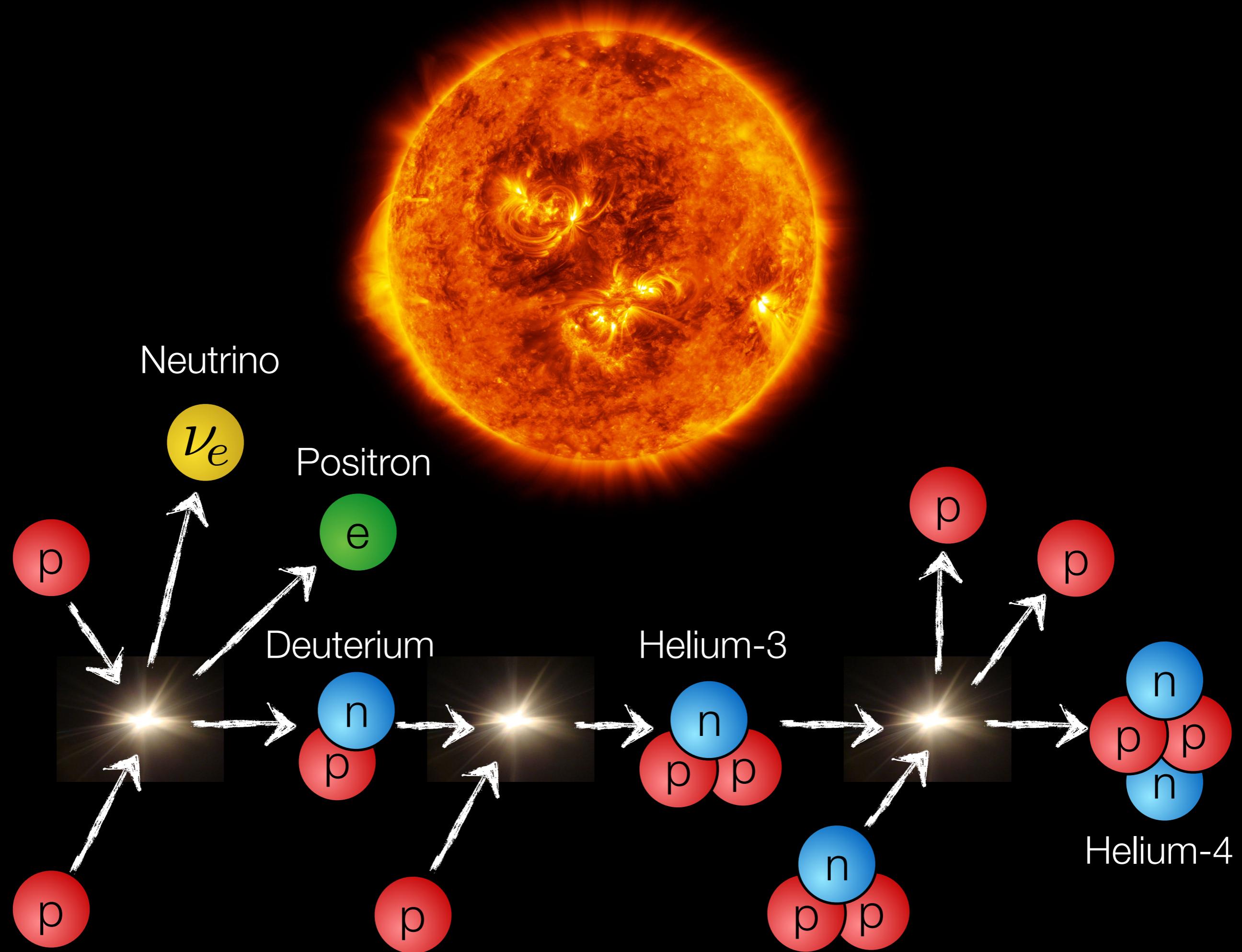


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Stellar Evolution

Stellar Evolution Equations

$$\frac{\partial r}{\partial m} = \frac{1}{4\pi r^2 \varrho} ,$$

$$\frac{\partial P}{\partial m} = -\frac{Gm}{4\pi r^4} ,$$

$$\frac{\partial l}{\partial m} = \varepsilon_n - \varepsilon_v - c_P \frac{\partial T}{\partial t} + \frac{\delta}{\varrho} \frac{\partial P}{\partial t} ,$$

$$\frac{\partial T}{\partial m} = -\frac{GmT}{4\pi r^4 P} \nabla ,$$

$$\frac{\partial X_i}{\partial t} = \frac{m_i}{\varrho} \left(\sum_j r_{ji} - \sum_k r_{ik} \right) , \quad i = 1, \dots, I .$$

from Kippenhahn Weigert Weiss “Stellar Structure and Evolution”

Stellar Evolution

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Conservation of mass
(Relation between radius, enclosed mass, and density)

Stellar Evolution

Stellar Evolution Equations

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Conservation of momentum

(Pressure difference across a mass shell
is due to mass of the shell)

Stellar Evolution

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Conservation of energy

(change in **luminosity** across a mass shell due to energy generation in **nuclear reactions**, energy loss due to **neutrino emission**, and energy production or loss by **expansion/contraction**)



Stellar Evolution

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Thermal Structure

(change in temperature across a mass shell is related to equation of state $\Delta = d \ln T / d \ln P$)

Stellar Evolution

Stellar Evolution Equations

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**Evolution of Mass Fraction
of different species**

$$\frac{\partial T}{\partial m} = -\frac{GmT}{4\pi r^4 P} \nabla,$$

$$\frac{\partial X_i}{\partial t} = \frac{m_i}{\varrho} \left(\sum_j r_{ji} - \sum_k r_{ik} \right), \quad i = 1, \dots, I.$$

from Kippenhahn Weigert Weiss “Stellar Structure and Evolution”

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from Kippenhahn Weigert Weiss “Stellar Structure and Evolution”

Stellar Evolution

Stellar Evolution Equations

- Solve numerically on *discretized grid*
- Additional complication: *convection*

Requires input parameters / boundary conditions

- total mass
- surface temperature
- initial chemical composition

Most relevant output for neutrino physics

- Core temperature T_{core}
- Neutrino flux depends on the 25^{th} power of T_{core} !

Helioseismology

- Need to determine the Sun's chemical composition
- Need tool to verify solar models
- Helioseismology
 - Study oscillation modes of the Sun
 - Generated in the convective zone
 - Observed via Doppler shift of spectral lines
 - Oscillation modes depend on the Sun's internal structure, so they allow us to learn about the latter

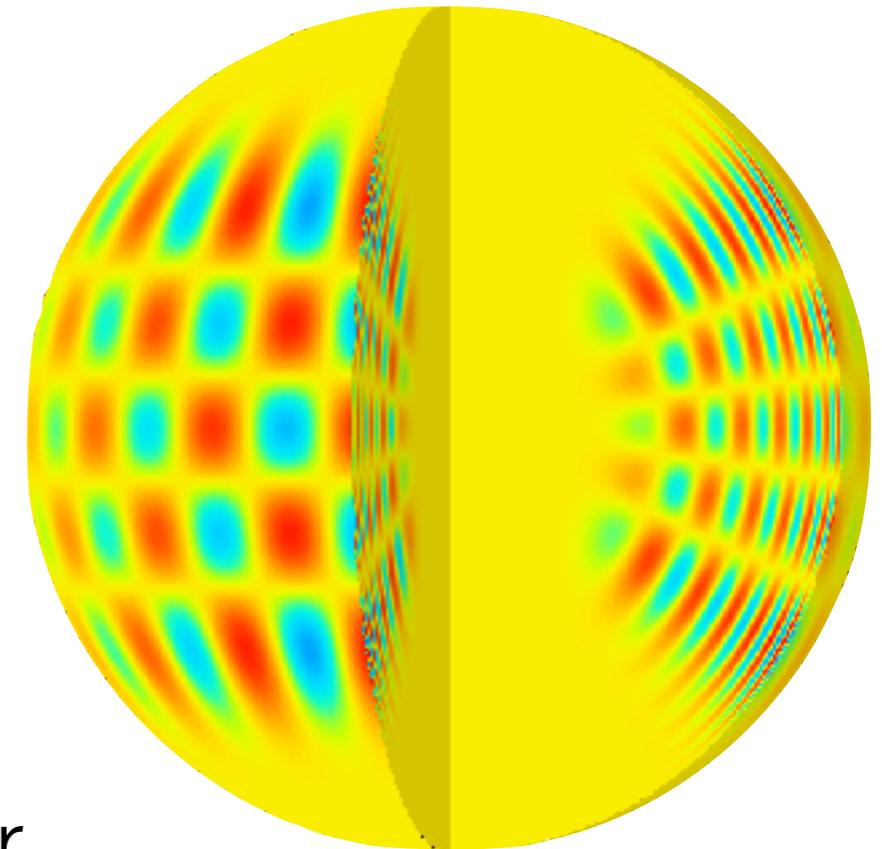
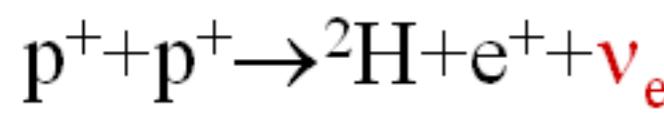
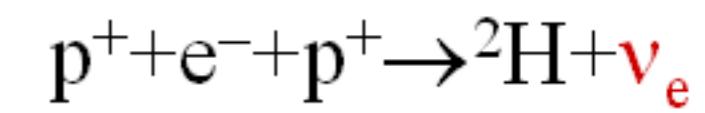


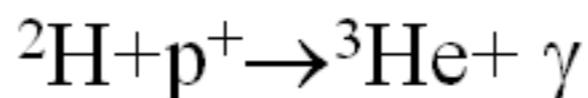
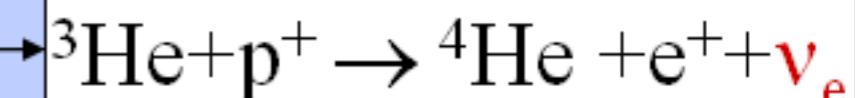
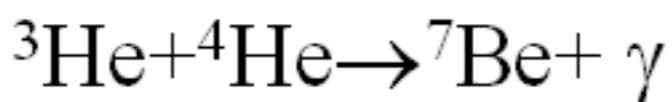
Image from <http://soi.stanford.edu/results/heliowhat.html>

pp**pep**

99,77 %

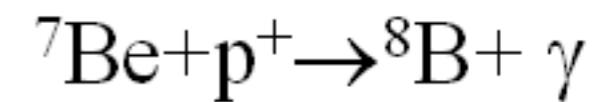
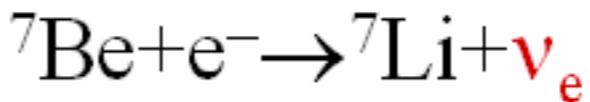
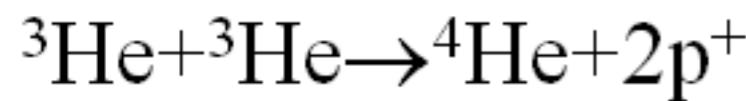


0,23 %

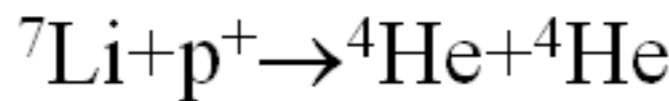
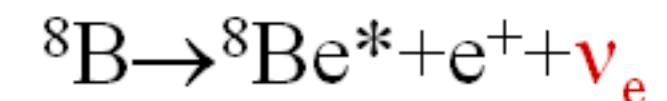
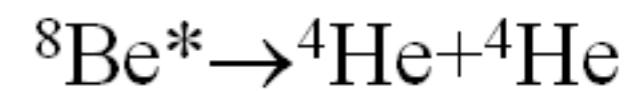
10⁻⁵ %**hep****⁷Be**

99,9 %

0,1 %

**BB****ppI**

84,92 %

**ppII****ppIII**

Predicted Solar Neutrino Flux

- Pure ν_e
- Flux on Earth:
 $\sim 10^{11} \text{ cm}^2/\text{sec}$
- Energy:
 $\lesssim 10 \text{ MeV}$

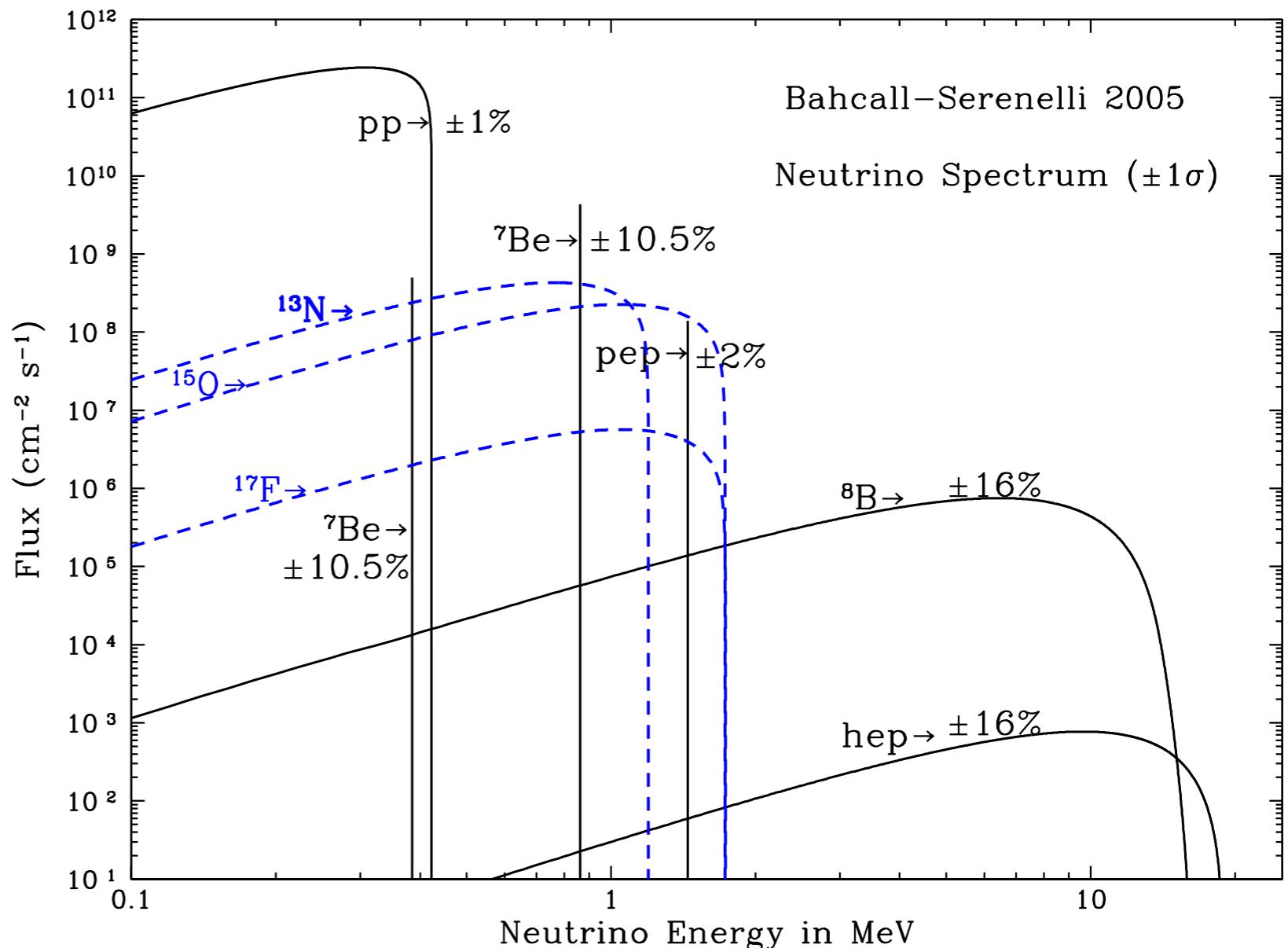
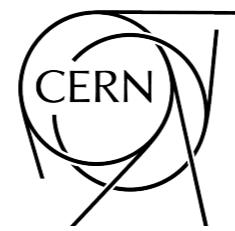


Image by John Bahcall

Reactor Neutrinos



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Predicting the Reactor Neutrino Spectrum

Method:

Mueller *et al.* [1101.2663](#), Huber [1106.0687](#)

- Use measured β spectra from ^{235}U , ^{238}U , ^{239}Pu , ^{241}Pu fission
- Convert to $\bar{\nu}_e$ spectrum
- For single β decay: $E_\nu = Q - E_e$

$$\frac{dN_\nu(E_\nu)}{dE_\nu} \equiv \frac{dN_e(Q - E_e)}{dE_e}$$

- For energy-independent nuclear matrix elements:
simple phase space argument

$$\begin{aligned} dN_\nu &\propto d^3 p_e d^3 p_\nu \delta(E_e + E_\nu - Q) \\ &\propto p_e^2 dp_e p_\nu^2 dp_\nu \delta(E_e + E_\nu - Q) \\ &= p_e E_e p_\nu E_\nu dE_\nu \\ &= \sqrt{(Q - E_\nu)^2 - m_e^2} (Q - E_\nu) E_\nu^2 dE_\nu. \end{aligned}$$

Corrections to the Reactor ν Spectrum

- Fermi function $F(A, Z, E_\nu)$
 - describes interactions of final state electron with Coulomb field of the nucleus
- Screening of the nuclear charge by bound electrons
- Non-zero nuclear radius
- Final state radiation: $(A, Z) \rightarrow (AA, Z+1) + e^- + \bar{\nu}_e + \gamma$
- Approximation of energy-independent nuclear matrix elements valid only for allowed beta decays

Corrections to the Reactor v Spectrum (contd.)

 Weak Magnetism: impact of finite nuclear size on weak interactions

- Weak interaction vertex:

$$\mathcal{L}_{\text{weak}} \supset \frac{g}{\sqrt{2}} J_W^\mu W_\mu + h.c.$$

with the weak current

$$J_{W,\text{point-like}}^\mu = \bar{u} \gamma^\mu \frac{1 - \gamma^5}{2} d$$

- For non-pointlike objects, extra terms and form factors appear

$$J_{W,\text{extended}}^\mu = \bar{u} \left[c_V(q^2) \gamma^\mu + c_A(q^2) \gamma^\mu \gamma^5 + F_2(q^2) \frac{i \sigma^{\mu\nu} q_\nu}{2M} \right] d,$$

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weak vector charge

Fermi form factor

Corrections to the Reactor v Spectrum (contd.)

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weak axial charge
Gamow-Teller form factor

Corrections to the Reactor v Spectrum (contd.)

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weak magnetism
(analog of **magnetic moment** in QED)

Corrections to the Reactor v Spectrum (contd.)

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Predicting the Reactor Neutrino Spectrum

- This was the story for a single beta decay ...
- Reality: ~6000 decays contribute
 - would need to know Q-value, relative importance, and all correction factors for each of them
 - this information is available only for some decays
 - many isotopes are too short-lived to be studied in the lab
- Method:
 - Use information from nuclear data tables where available ...
 - ... complemented by a fit to “effective decay branches”
(a set of beta decays with parameters fitted in order to match the observed electron spectrum)

Mueller *et al.* [1101.2663](#), Huber [1106.0687](#)

Verification

- Simulate mock electron spectrum
- Convert to neutrino spectrum using aforementioned method
- Compare to MC truth

The Reactor Neutrino Anomaly

Result: predicted flux is $\sim 3.5\%$ ($\sim 3\sigma$) higher than observation

The Reactor Neutrino Anomaly

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