

# Prediction for Neutrino Masses, CP Violation and Neutrinoless Double Beta Decay from

## $\mathcal{T}_{13}$ Family Symmetry

Moinul Hossain Rahat

Institute for Fundamental Theory, Department of Physics, University of Florida

Collaborators: M. Jay Pérez, **Pierre Ramond**, Alexander Stuart, Bin Xu

### The Asymmetric Texture

The **minimal Yukawa texture** based on  $SU(5)$  GUT that explains GUT-scale mass ratios and mixings of quarks and leptons:

$$Y^{(\frac{2}{3})} \sim \text{diag}(\lambda^8, \lambda^4, 1),$$

$$Y^{(-\frac{1}{3})} \sim \begin{pmatrix} bd\lambda^4 & a\lambda^3 & b\lambda^3 \\ a\lambda^3 & c\lambda^2 & g\lambda^2 \\ d\lambda & g\lambda^2 & 1 \end{pmatrix},$$

$$Y^{(-1)} \sim \begin{pmatrix} bd\lambda^4 & a\lambda^3 & d\lambda \\ a\lambda^3 & -3c\lambda^2 & g\lambda^2 \\ b\lambda^3 & g\lambda^2 & 1 \end{pmatrix},$$

$$\mathcal{U}_{\text{Seesaw}} = \text{diag}(1, 1, e^{i\delta}) \mathcal{U}_{\text{TBM}},$$

where  $a = c = \frac{1}{3}$ ,  $g = A$ ,  $b = A\sqrt{\rho^2 + \eta^2}$ ,  $d = \frac{2a}{g} = \frac{2}{3A}$  and  $\cos \delta = 0.2$ . The origin of the texture is traced back to an  $SU(5) \times \mathcal{T}_{13}$  model.

### Model Building Features

- **diagonal**  $Y^{(\frac{2}{3})}$ : should be explained from group theory
- **asymmetry** along (13) – (31) of  $Y^{(-\frac{1}{3})}$  and  $Y^{(-1)}$ : need a discrete subgroup of  $SU(3)$  with at least two inequivalent triplet representations
- **zero subdeterminant** with respect to the (22) element: GUT-scale mass ratios suggest  $\det Y^{(-\frac{1}{3})} = \det Y^{(-1)}$
- **TBM** seesaw mixing with a phase: should originate from breaking of family symmetry

The minimum order group that fits the bill is  $\mathcal{T}_{13} = \mathcal{Z}_{13} \rtimes \mathcal{Z}_3$  at order 39.

### Key Predictions

- **Dirac CP violation**:  $\delta_{CP} = 1.32\pi$ ,  $\mathcal{J}_{CP} = -0.028$   
PDG 2020 fit:  $\delta_{CP} = 1.36 \pm 0.17\pi$  at  $1\sigma$   
 $5\sigma$  measurement by DUNE and Hyper-K in  $\sim 5 - 10$  yrs
- **Normal ordering of neutrino masses**:  $m_{\nu_1} = 27.6$  meV,  $m_{\nu_2} = 28.9$  meV,  $m_{\nu_3} = 57.8$  meV  
Normal ordering over inverted ordering is preferred over  $3.6\sigma$  from T2K and  $\text{NO}\nu\text{A}$
- **Sum of neutrino masses**:  $\sum_i m_{\nu_i} = 114.3$  meV  
Planck 2018:  $\sum_i m_{\nu_i} < 120$  meV, error bar less than 10 meV expected from combining Euclid, LSST data with DESI, WFIRST
- **Majorana invariants**:  $\mathcal{I}_1 = -0.106$ ,  $\mathcal{I}_2 = -0.011$
- **Invariant mass parameter for  $0\nu\beta\beta$** :  $|m_{\beta\beta}| = 13.02$  or  $25.21$  meV  
KamLAND-Zen:  $|m_{\beta\beta}| < 61 - 165$  meV, prediction is sensitive to LEGEND (11 – 28 meV), nEXO (8 – 22 meV), CUPID (6 – 17 meV)
- **Baryon asymmetry of the universe** can be explained through leptogenesis with **right handed neutrino masses**  $\sim 10^8 - 10^{12}$  GeV.

### Seesaw Sector

- Introduce four right handed neutrinos  $\bar{N} \sim (\mathbf{1}, \mathbf{3}_2)$  and  $\bar{N}_4 \sim (\mathbf{1}, \mathbf{1})$
- **Dirac Yukawa matrix**,  $Y^{(0)}$ :  $F\bar{N}\bar{H}_5\varphi_A$
- Majorana matrix,  $\mathcal{M}$ :  $\bar{N}\bar{N}\varphi_B$  and  $\bar{N}\bar{N}_4\varphi_z$
- **Seesaw matrix**,  $\mathcal{S} = Y^{(0)} \mathcal{M}^{-1} Y^{(0)T}$  is diagonalized by the complex TBM matrix:  $\mathcal{S} \propto \mathcal{U}_{\text{TBM}}(\delta) \text{diag}(m_{\nu_1}, m_{\nu_3}/2, m_{\nu_3}) \mathcal{U}_{\text{TBM}}^T(\delta)$ .
- Neutrino masses are determined from  $m_{\nu_2} = \frac{1}{2}m_{\nu_3}$  using oscillation data
- TBM phase  $\delta = \cos^{-1} 0.2$  yields the Dirac and Majorana CP phases and  $|m_{\beta\beta}|$
- Successful leptogenesis can occur for right handed neutrino masses  $\sim 10^8 - 10^{12}$  GeV

### References

- **M.H.R.**, P. Ramond, and B. Xu, *Phys. Rev. D*, 98(5), 055030, 2018, [arXiv:1805.10684](https://arxiv.org/abs/1805.10684) [hep-ph]
- M. J. Pérez, **M.H.R.**, P. Ramond, A. J. Stuart, and B. Xu, *Phys. Rev. D*, 100(7), 075008, 2019, [arXiv:1907.10698](https://arxiv.org/abs/1907.10698) [hep-ph]
- M. J. Pérez, **M.H.R.**, P. Ramond, A. J. Stuart, and B. Xu, *Phys. Rev. D*, 101(7), 075008, 2020, [arXiv:2001.04019](https://arxiv.org/abs/2001.04019) [hep-ph]

### Acknowledgements

This work was supported in part by the U.S. Department of Energy under Grant No. DE-SC0010296.

### Contact Information

- Email: [mrahat@ufl.edu](mailto:mrahat@ufl.edu)

### Building the Yukawas

$$Y^{(-\frac{1}{3})} \leftarrow FTH_5\varphi^{(-\frac{1}{3})}: \begin{pmatrix} F_1 \\ F_2 \\ F_3 \end{pmatrix}_{\mathbf{3}_1} \otimes \begin{pmatrix} T_1 \\ T_3 \\ T_2 \end{pmatrix}_{\mathbf{3}_2} = \begin{pmatrix} F_3T_2 \\ F_1T_1 \\ F_2T_3 \end{pmatrix}_{\mathbf{3}_1} \oplus \begin{pmatrix} F_3T_1 \\ F_1T_3 \\ F_2T_2 \end{pmatrix}_{\mathbf{3}_2} \oplus \begin{pmatrix} F_3T_3 \\ F_1T_2 \\ F_2T_1 \end{pmatrix}_{\mathbf{3}_2}$$

$\mathcal{T}_{13}$  can dial *individual* matrix elements!

$$Y^{(\frac{2}{3})} \leftarrow TTH_5\varphi^{(\frac{2}{3})}: \begin{pmatrix} T_1 \\ T_3 \\ T_2 \end{pmatrix}_{\mathbf{3}_2} \otimes \begin{pmatrix} T_1 \\ T_3 \\ T_2 \end{pmatrix}_{\mathbf{3}_2} = \begin{pmatrix} T_3T_3 \\ T_2T_2 \\ T_1T_1 \end{pmatrix}_{\mathbf{3}_1} \oplus \begin{pmatrix} T_3T_2 \\ T_2T_1 \\ T_1T_3 \end{pmatrix}_{\mathbf{3}_2} \oplus \begin{pmatrix} T_2T_3 \\ T_1T_2 \\ T_3T_1 \end{pmatrix}_{\mathbf{3}_2}$$

Diagonals are distinguished from off-diagonals!

### Zero Subdeterminant

$$\begin{pmatrix} bd\lambda^4 & \times & \times \\ \times & \times & \times \\ d\lambda & \times & 1 \end{pmatrix} = \begin{pmatrix} b\lambda^3 \cdot d\lambda & \times & b\lambda^3 \cdot 1 \\ \times & \times & \times \\ d\lambda & \times & 1 \end{pmatrix}$$

$$(33) \text{ term} \rightarrow FTH_5\varphi^{(1)}$$

$$(31) \text{ term} \rightarrow FTH_5\varphi^{(2)}$$

$$(13) \text{ term} \rightarrow FTH_5\varphi^{(1)}\varphi^{(3)}$$

$$(11) \text{ term} \rightarrow FTH_5\varphi^{(2)}\varphi^{(3)}$$

Higgs  $H_5$  is a family singlet and the familons  $\varphi^{(i)}$  are gauge singlets and family (anti)triplets.