Nuclear Structure Physics in Coherent Elastic Neutrino-Nucleus Scattering



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I. Introduction

- The primary source of uncertainty in the evaluation of the coherent elastic neutrino-nucleus scattering (CEvNS) cross section comes from the underlying nuclear dynamics embedded in the nucleon form factors.
- Proton density distributions are relatively well constrained through elastic electron scattering experiments while the neutron density distributions are poorly known.
- An accurate estimation of form factors is vital to the CEvNS program, since any experimentally measured deviation from the expected CEvNS event rate can either be attributed to new physics or to unconstrained nuclear physics.

II. Formalism

• A neutrino scatters off a nucleus, ν_{α} (E_f, initially at rest, exchanging a single Z^0 boson. The nucleus remains in its ground state and receives a small recoil energy T. The differential cross section is ν_{α} (E expressed as:

$$\begin{array}{c} A \langle f | (M_A + T, p'_A) \rangle \\ Z^0 (T, \vec{q}) \\ F_i, \vec{k}_i \rangle \\ \end{array}$$

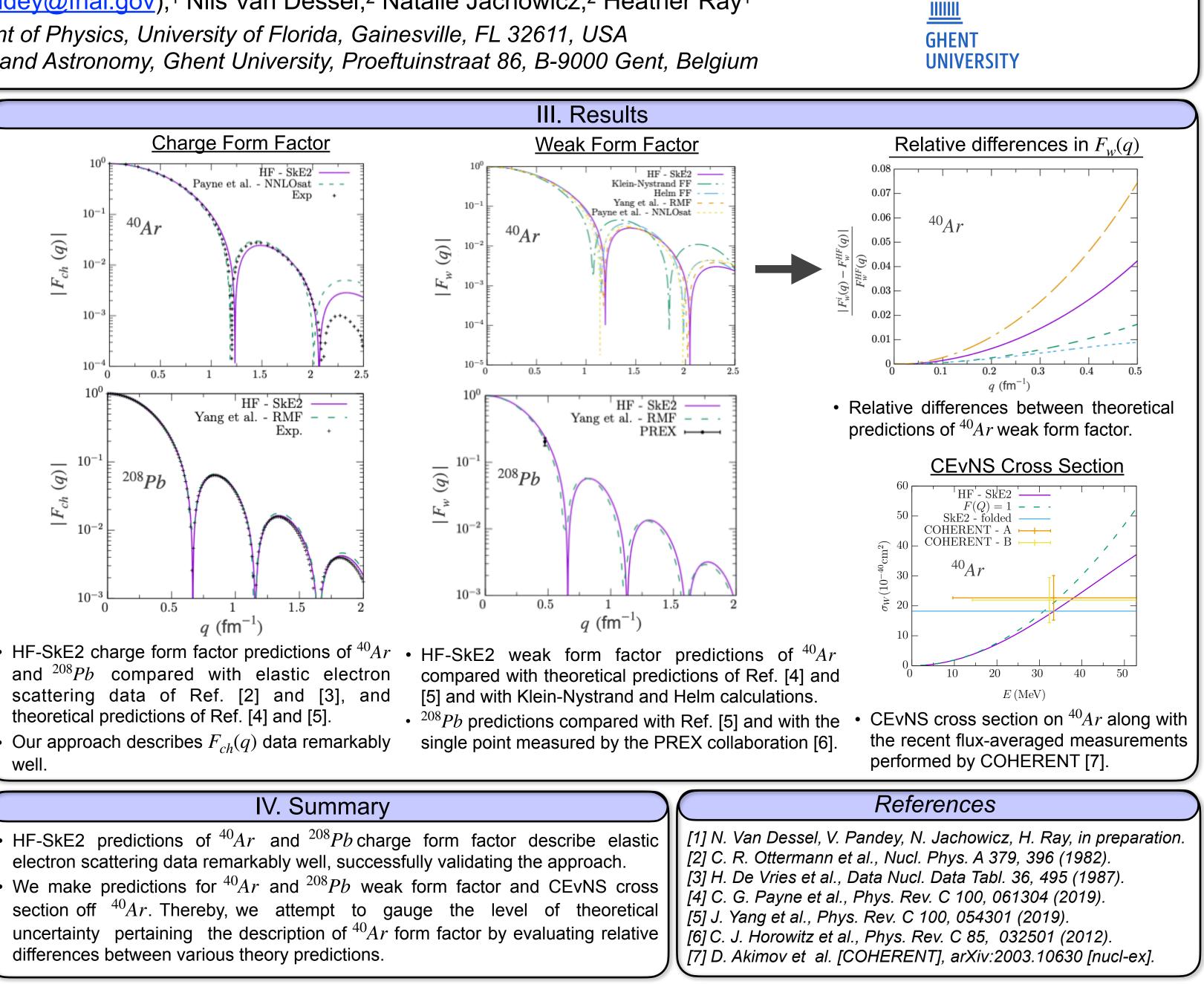
$$\frac{d\sigma}{dT} = \frac{G_F^2}{\pi} M_A \left(1 - \frac{T}{E_i} - \frac{M_A T}{2E_i^2} \right) \frac{Q_W^2}{4} F_W^2(Q^2)$$

where $Q_{\rm W}$ is weak nuclear charge, and $F_{\rm W}(Q^2)$ is weak form factor written as:

$$F_W(q^2) = \frac{1}{Q_W} \left[(1 - 4\sin^2\theta_W) f_p(q) F_p(q^2) - f_n(q) F_n(q^2) \right]$$
$$f_n(q) = \frac{4\pi}{N} \int dr \ r^2 \ \frac{\sin(qr)}{qr} \ \rho_n(r), \qquad f_p(q) = \frac{4\pi}{Z} \int dr \ r^2 \ \frac{\sin(qr)}{qr} \ \rho_p(r)$$

with $f_n(f_p)$ as neutron (proton) form factors, and $\rho_n(\rho_p)$ as density distributions.

• We obtain proton and neutron densities, and charge and weak form factors, in a microscopic many-body nuclear theory model where the nuclear ground state is described in a Hartree-Fock (HF) approach with a Skyrme (SkE2) nuclear potential [1].



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