

N-Subjettiness for boosted jets

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- Jet-Substructure techniques for boosted jets for signal/bkg disc
- Particularly important for hadronically decaying colourless particles
- Here focus on two-pronged decays, like a W

N-Subjettiness: $\tau_N^{(\beta)} = \frac{1}{p_{\perp} R^{\beta}} \sum_{i \in \text{jet}} p_{\perp i} \min(\theta_{ia_1}, \dots, \theta_{ia_N})$

N-subjettiness

$$d\mathcal{P}_W = dR_W(x) = \lambda^2 dz d\theta^2 P_W$$



- Uniform and flat in energy and angle

$$d\mathcal{P}_{\text{QCD}} = dR = \frac{\alpha_s C_i}{2\pi} d \log z d \log \theta^2 P_W$$



- more soft and collinear

$$\tau_2 = \sum_{i \in \text{jet}} \min(z_i \theta_{i1}^2, z_i \theta_{i2}^2)$$

$$\rightarrow \tau_2^W \sim \tau_2^{\text{QCD}}$$

However!

$$\tau_1 = \sum_{i \in \text{jet}} \min(z_i \theta_{i1}^2)$$

$$\rightarrow \tau_1^W > \tau_1^{\text{QCD}}$$

N-subjettiness

$$d\mathcal{P}_W = dR_W(x) = \lambda^2 dz d\theta^2 P_W$$



- Uniform and flat in energy and angle

$$d\mathcal{P}_{\text{QCD}} = dR = \frac{\alpha_s C_i}{2\pi} d \log z d \log \theta^2 P_W$$



- more soft and collinear

One then expects $\tau_{21} = \frac{\tau_2}{\tau_1}$ to be a good discriminant

N-Subjettines cut

- First efforts resumming double logs in the small tau region

$$\alpha_s^n \log^{2n} \rho, \quad \alpha_s^n \log^{2n} \tau, \quad \alpha_s^n \log^n \rho \log^n \tau$$

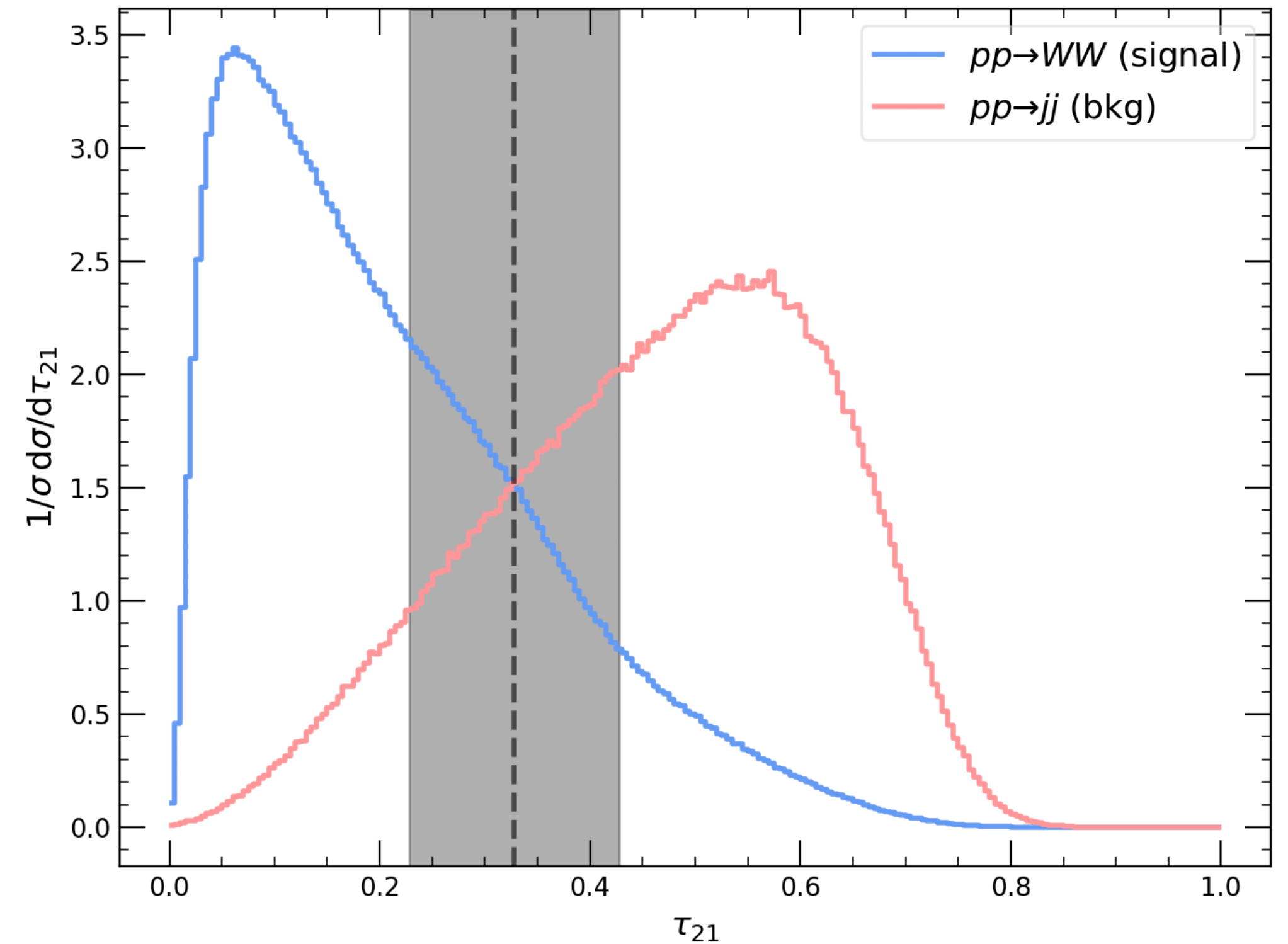
N-Subjettines cut

- First efforts resumming double logs in the small tau region

$$\alpha_s^n \log^{2n} \rho, \quad \alpha_s^n \log^{2n} \tau, \quad \alpha_s^n \log^n \rho \log^n \tau$$

- The interesting region's clearly at much larger tau!
- Structure now more complicated, as finite terms have to be included to have correct accuracy in rho

$$\alpha_s^n \log^{2n} \rho, \quad \alpha_s^n \log^{2n} \tau, \quad \alpha_s^n \log^n \rho f_n(\tau)$$



This is (in essence) the contribution we compute

N-Subjettiness cut

Small tau

$$\rho_1 \gg \rho_2 \gg \dots \gg \rho_n$$

$$\theta_1 \ll \theta_2 \ll \dots \ll \theta_n$$



Mass is set by first emission

$$\rho \sim \rho_1$$
$$\tau_{21} \sim \frac{\rho_2}{\rho_1}$$

N-Subjettiness cut

Small tau

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$$\theta_1 \ll \theta_2 \ll \dots \ll \theta_n$$

Mass is set by first emission

$$\rho \sim \rho_1$$
$$\tau_{21} \sim \frac{\rho_2}{\rho_1}$$

Finite tau

$$\rho_1 \sim \rho_2 \sim \dots \sim \rho_n$$

$$\theta_1 \ll \theta_2 \ll \dots \ll \theta_n$$

All emissions contribute

$$\rho = \sum_{i=1}^n \rho_i$$
$$\tau_{21} = 1 - \frac{\max_i \rho_i}{\rho}$$

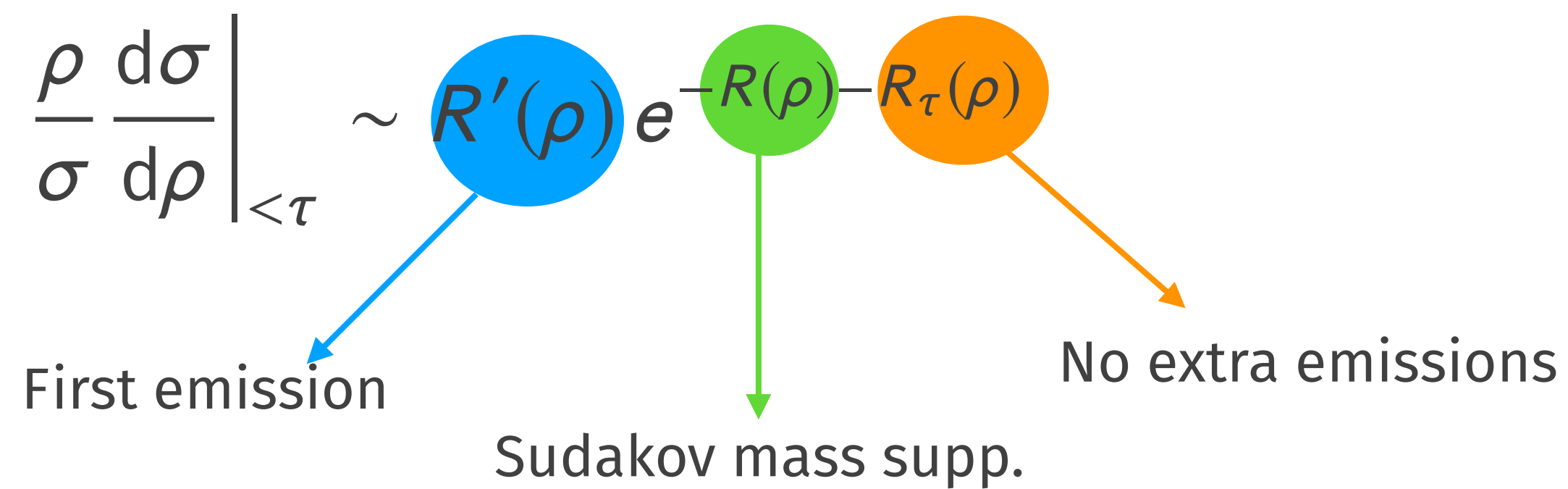
N-Subjettiness cut

Small tau

$$\rho_1 \gg \rho_2 \gg \dots \gg \rho_n$$

$$\theta_1 \ll \theta_2 \ll \dots \ll \theta_n$$

Mass is set by first emission



Finite tau

$$\rho_1 \sim \rho_2 \sim \dots \sim \rho_n$$

$$\theta_1 \ll \theta_2 \ll \dots \ll \theta_n$$

All emissions contribute

Next slides!!!

Finite tau

- Single out the emission with the largest rho: ρ_a

$$\frac{\rho}{\sigma} \frac{d\sigma}{d\rho} \Big|_{<\tau} = \int_0^1 \frac{d\rho_a}{\rho_a} R'(\rho_a) \lim_{\epsilon \rightarrow 0} e^{-\int_\epsilon^1 \frac{d\rho_v}{\rho_v} R'(\rho_v)}$$

$$\sum_{\rho=1}^{\infty} \frac{1}{\rho!} \int_\epsilon^{\rho_a} \prod_{i=1}^{\rho} \frac{d\rho_i}{\rho_i} R'(\rho_i) \rho \delta\left(\rho - \rho_a - \sum_{i=1}^{\rho} \rho_i\right) \Theta\left(\frac{\rho_a}{\rho} > 1 - \tau\right)$$

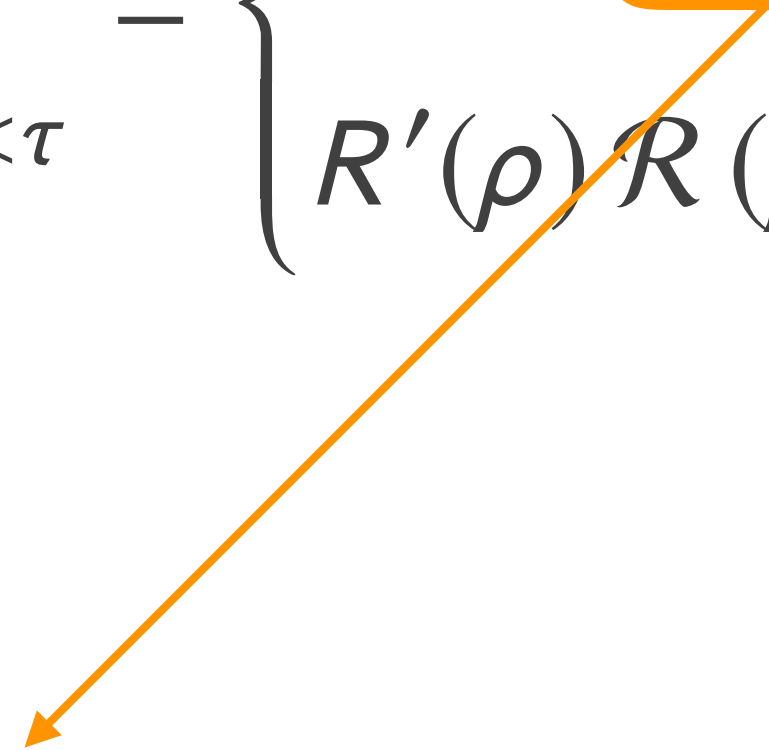
- Identify two regions $\tau \lesssim \frac{1}{2}$ and expand around $\frac{\tau}{1-\tau}$

$$R(\rho - \rho_a) \approx R\left(\rho \frac{\tau}{1-\tau}\right) + R'\left(\rho \frac{\tau}{1-\tau}\right) \log\left(\frac{\rho\tau}{(1-\tau)(\rho - \rho_a)}\right)$$

$$R'(\rho - \rho_a) \approx R'\left(\rho \frac{\tau}{1-\tau}\right).$$

Finite tau

$$\left. \frac{\rho}{\sigma} \frac{d\sigma}{d\rho} \right|_{<\tau} = \begin{cases} R'(\rho) \mathcal{R} \left(\frac{\rho \tau}{1 - \tau} \right) (1 - \tau)^{R'} {}_2F_1 & \tau < \frac{1}{2} \\ R'(\rho) \mathcal{R}(\rho) \left[2^{-R'(\rho)} {}_2F_1 + R'(\rho) \mathcal{I}_{ME} \right] & \tau > \frac{1}{2} \end{cases}$$


$$\mathcal{R}(x) = \frac{e^{-R(x) - \gamma_E R'(x)}}{\Gamma(1 + R'(x))}$$

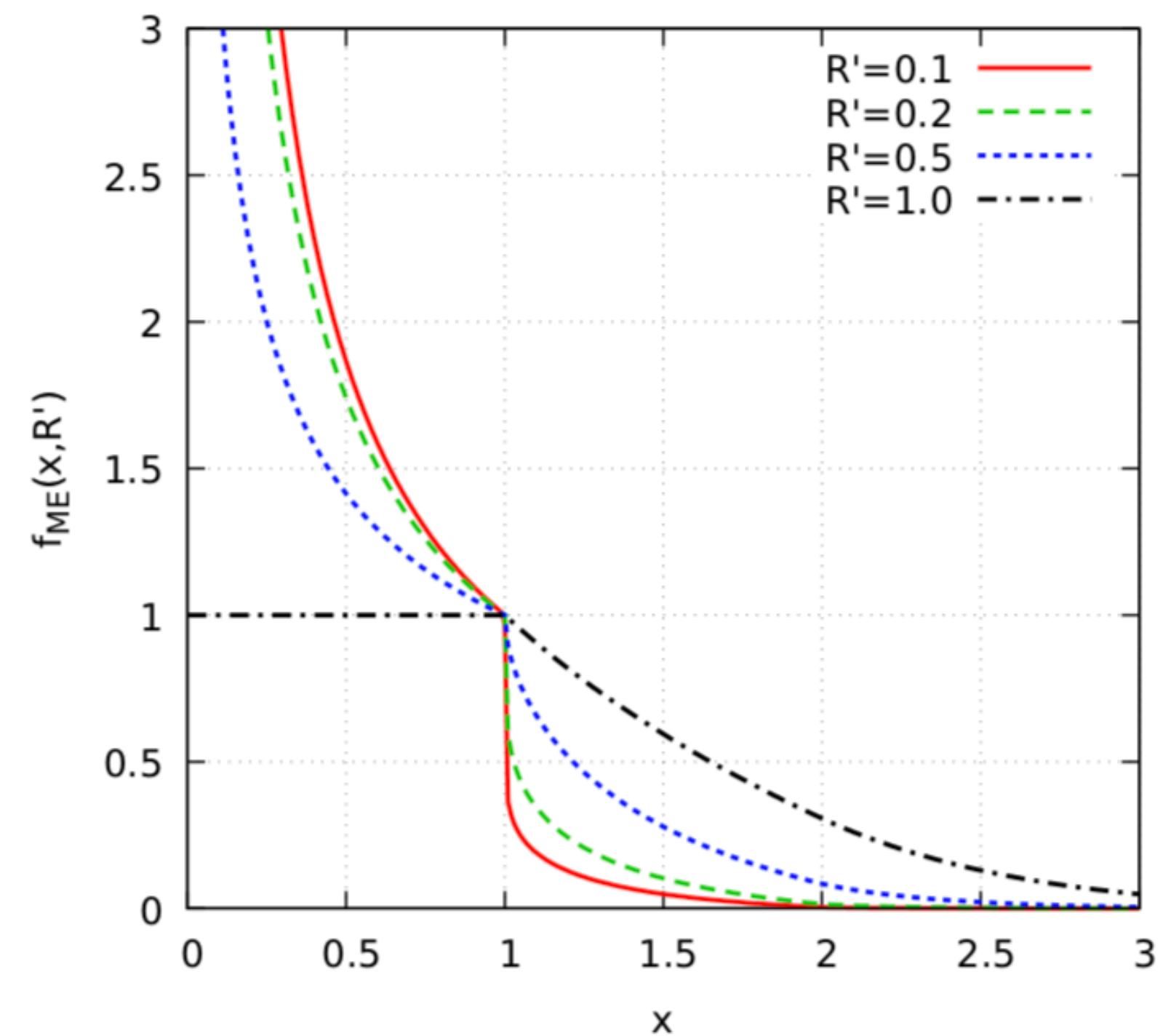
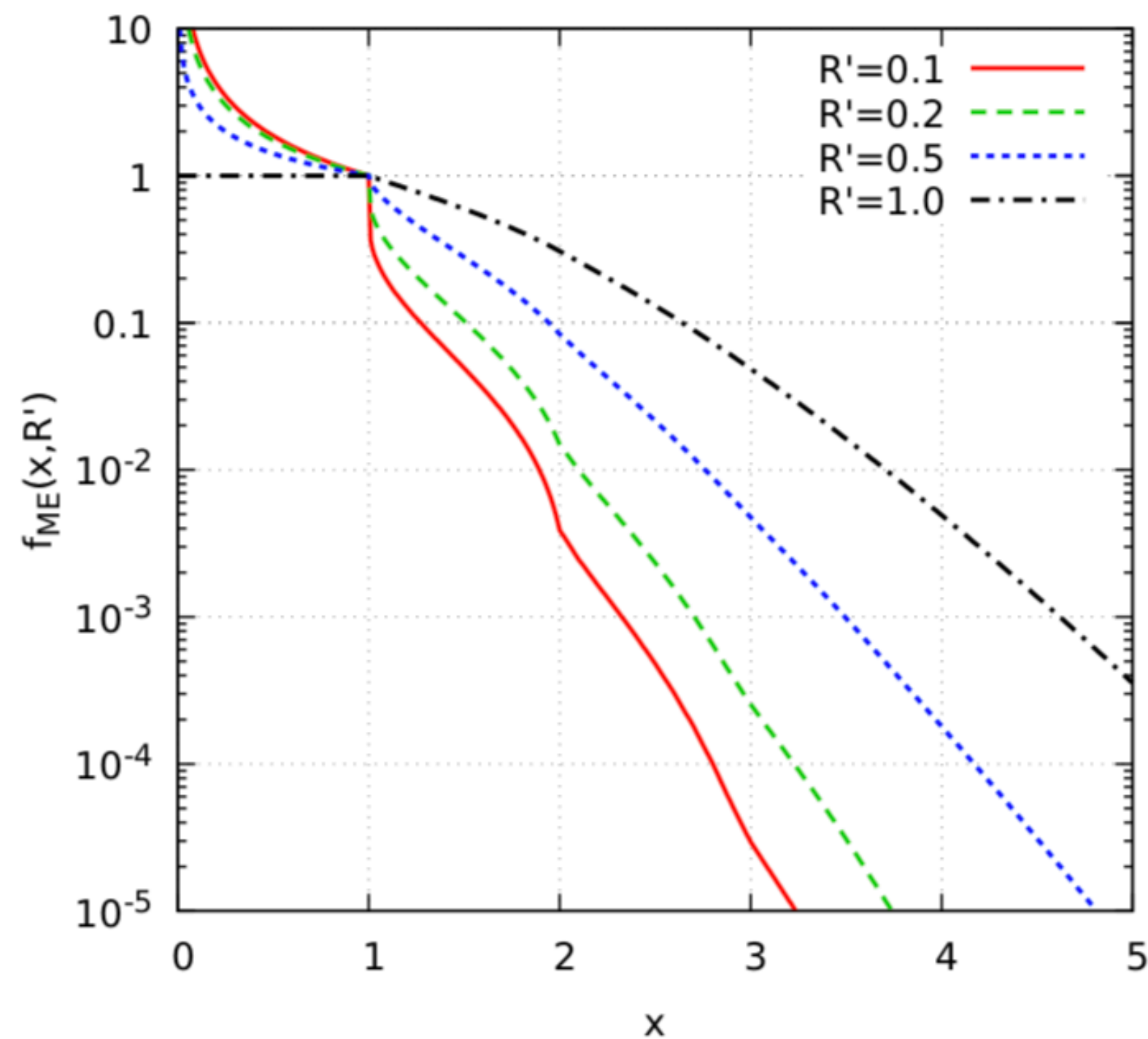
- Structure of the result is remarkably simple: scale depends on value of tau

Finite tau

- Multiple emission function, keeps track of transition points at $\tau = \frac{n-1}{n}$

$$\mathcal{I}_{ME} = \int_1^x \frac{du}{(1+u)^{R'}} \Gamma(R') \oint \frac{dv}{2i\pi} e^{v^x} \exp \left\{ \frac{R'}{2} \text{Ei}(-v) \left[\log(-v) - \log\left(\frac{1}{v}\right) \right] \right\}$$

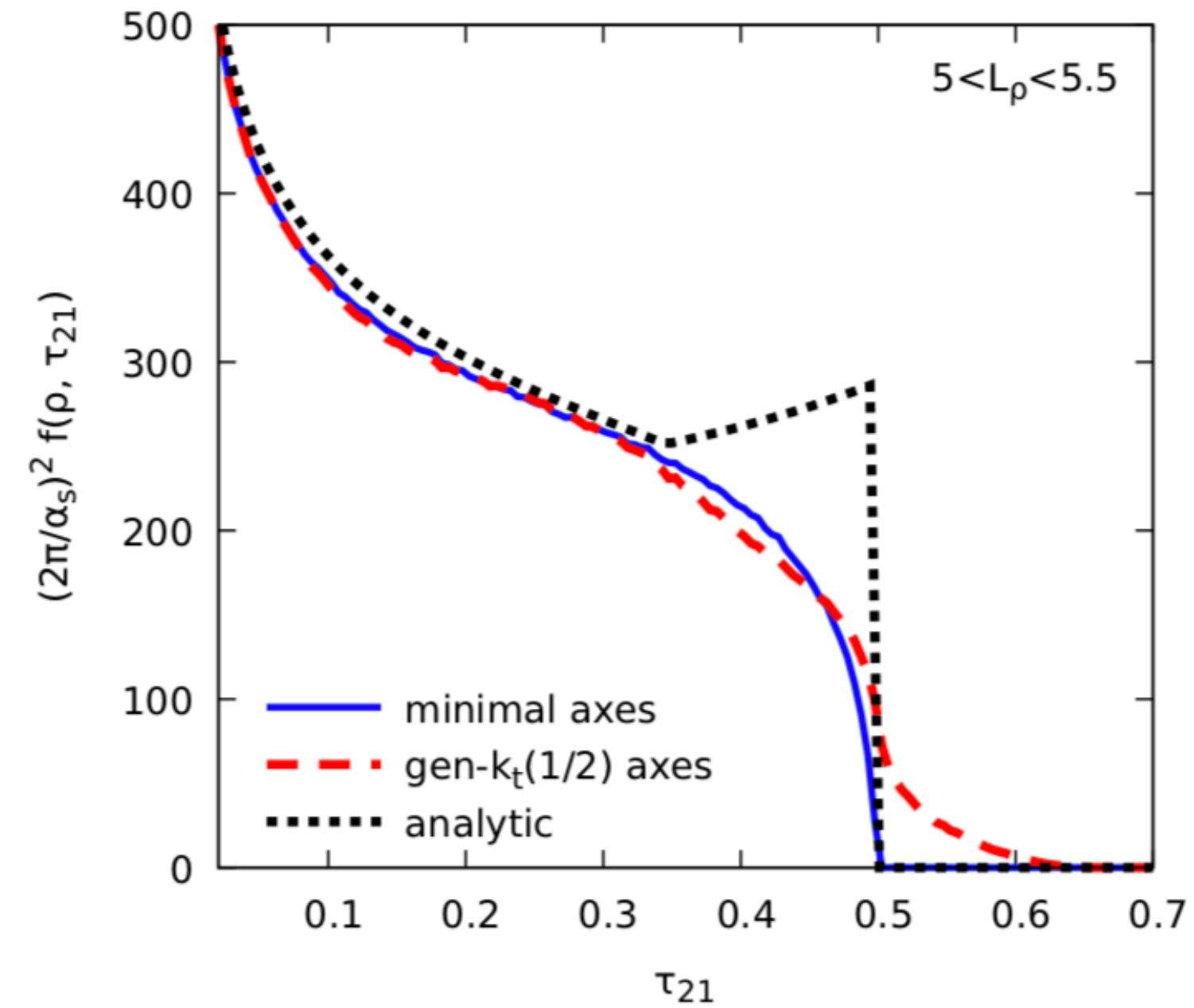
$f_{ME}(x, R')$



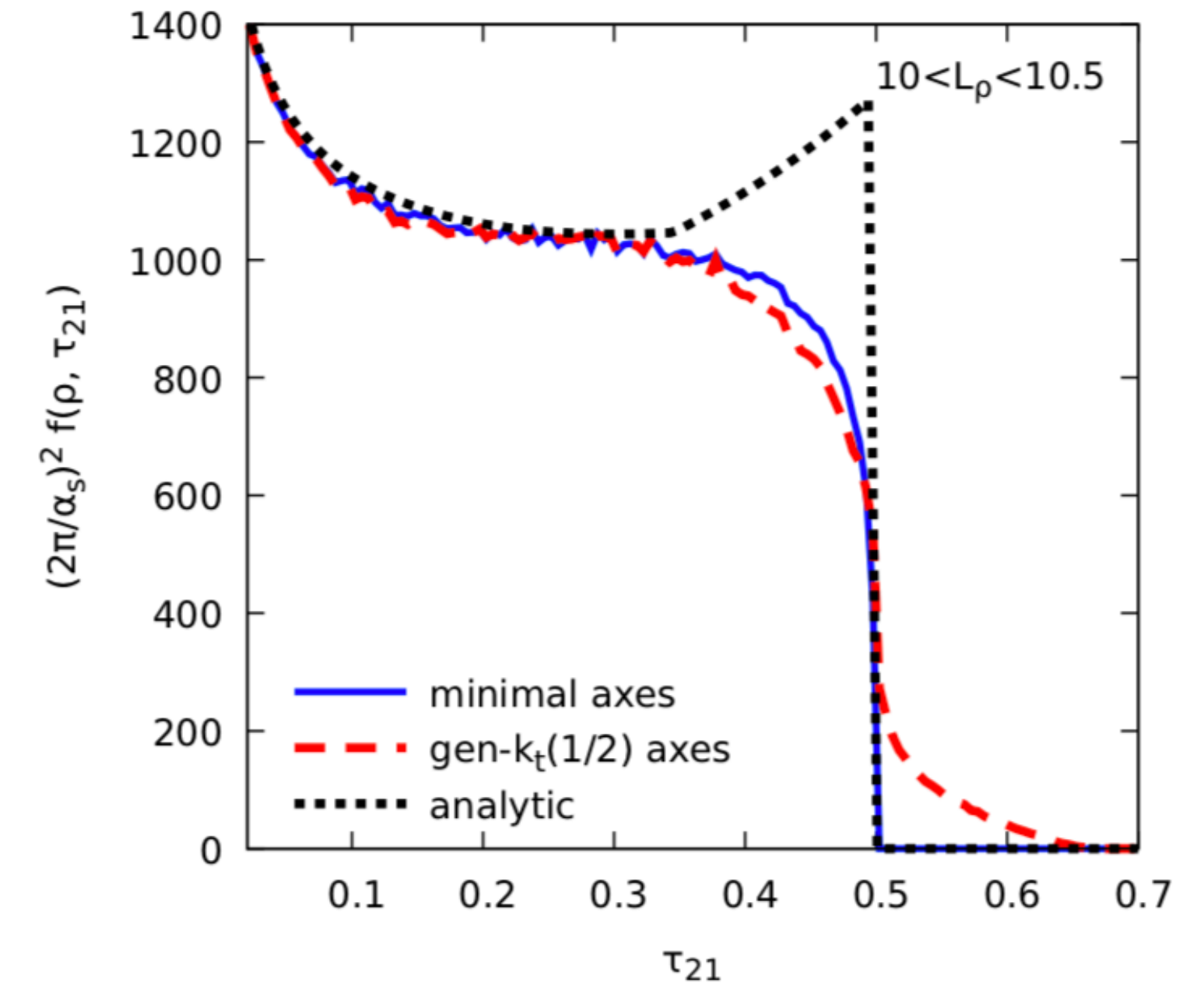
Results

- Checks of the minimisation procedure

- minimal axes quite similar to k_t



- Analytical vs exact: quite good at small/large tau

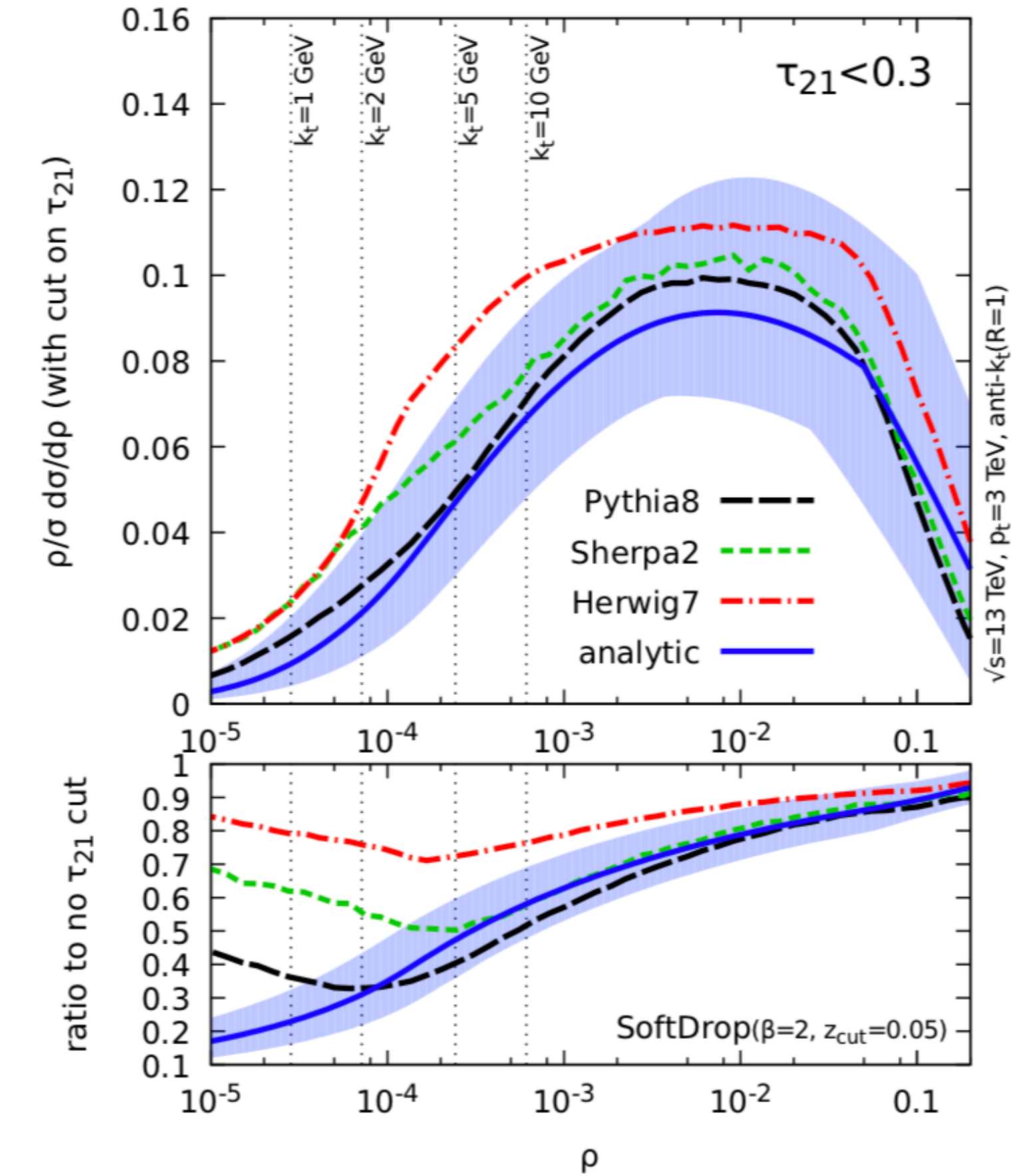
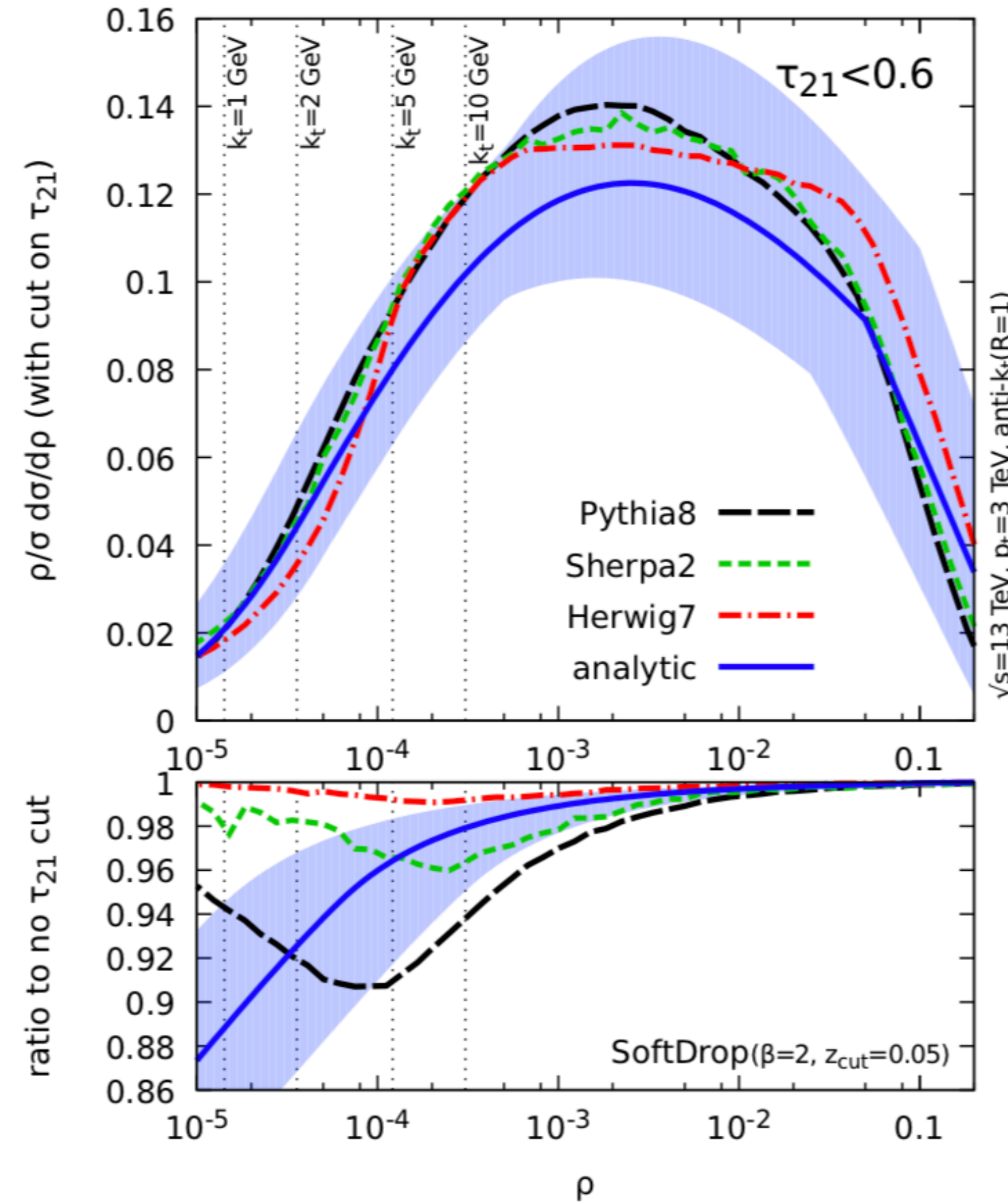


- **But**, it's integrated over, so beyond our accuracy (originates in neglecting large angle effects)

Results

- Good excuse to do a generators comparison!
- Analytical scale var is quite large (it's only LL)
- Overall quite good agreement (Herwig was before color fix)

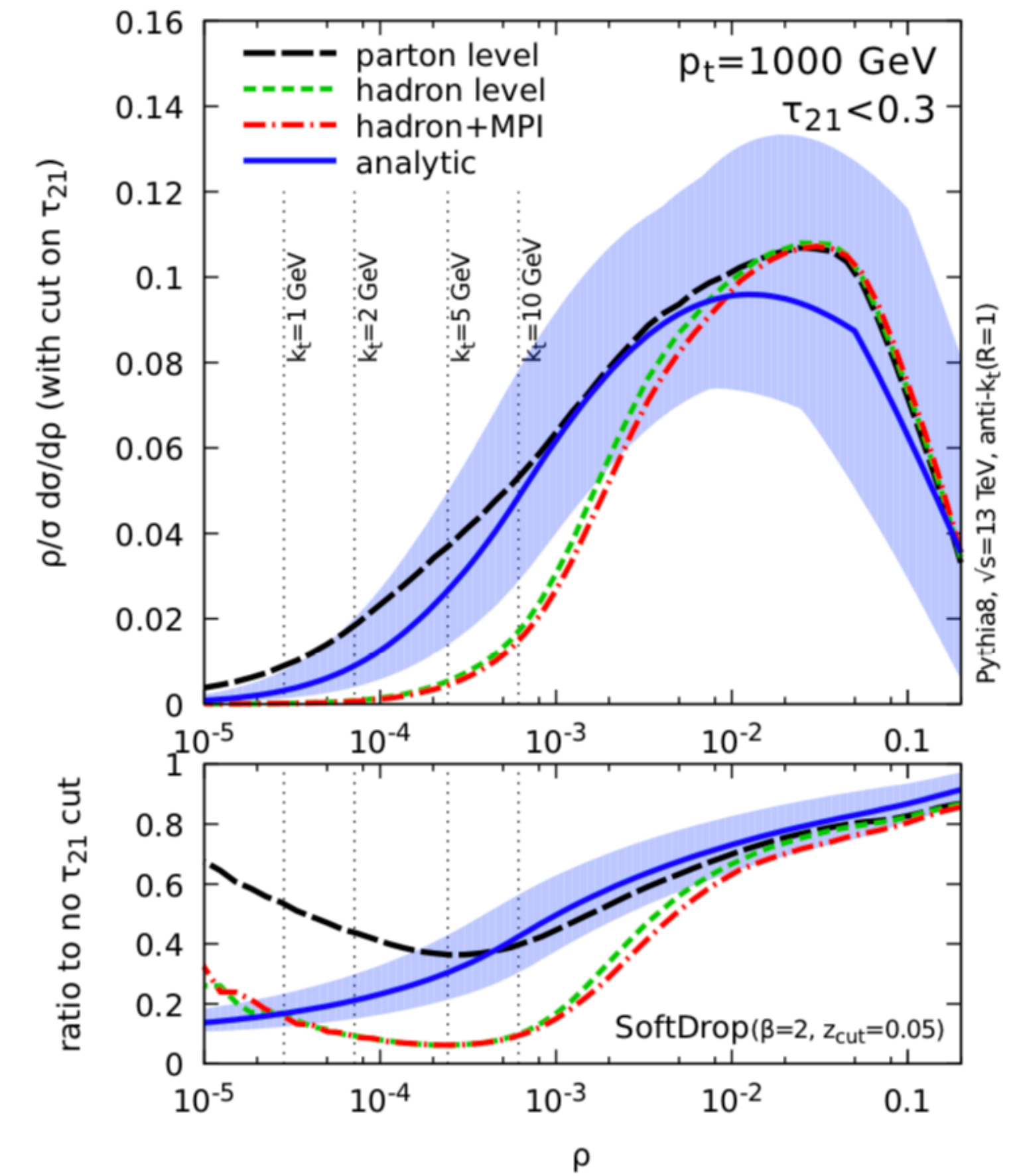
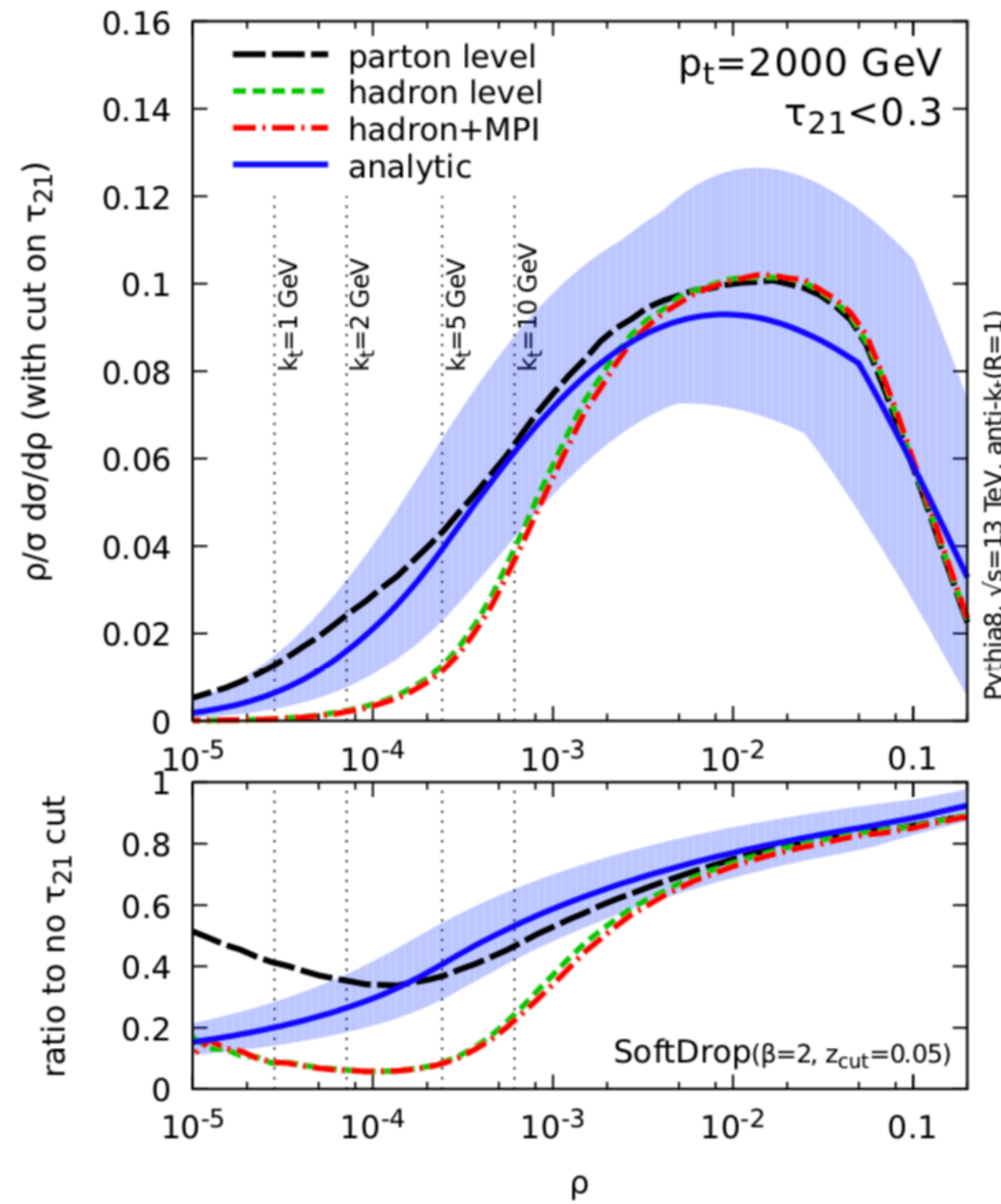
Parton Level



Results

- Clearly no control over NP-region
- still quite large effects in perturbative region
- rest is, over all, well under control

Hadron Level



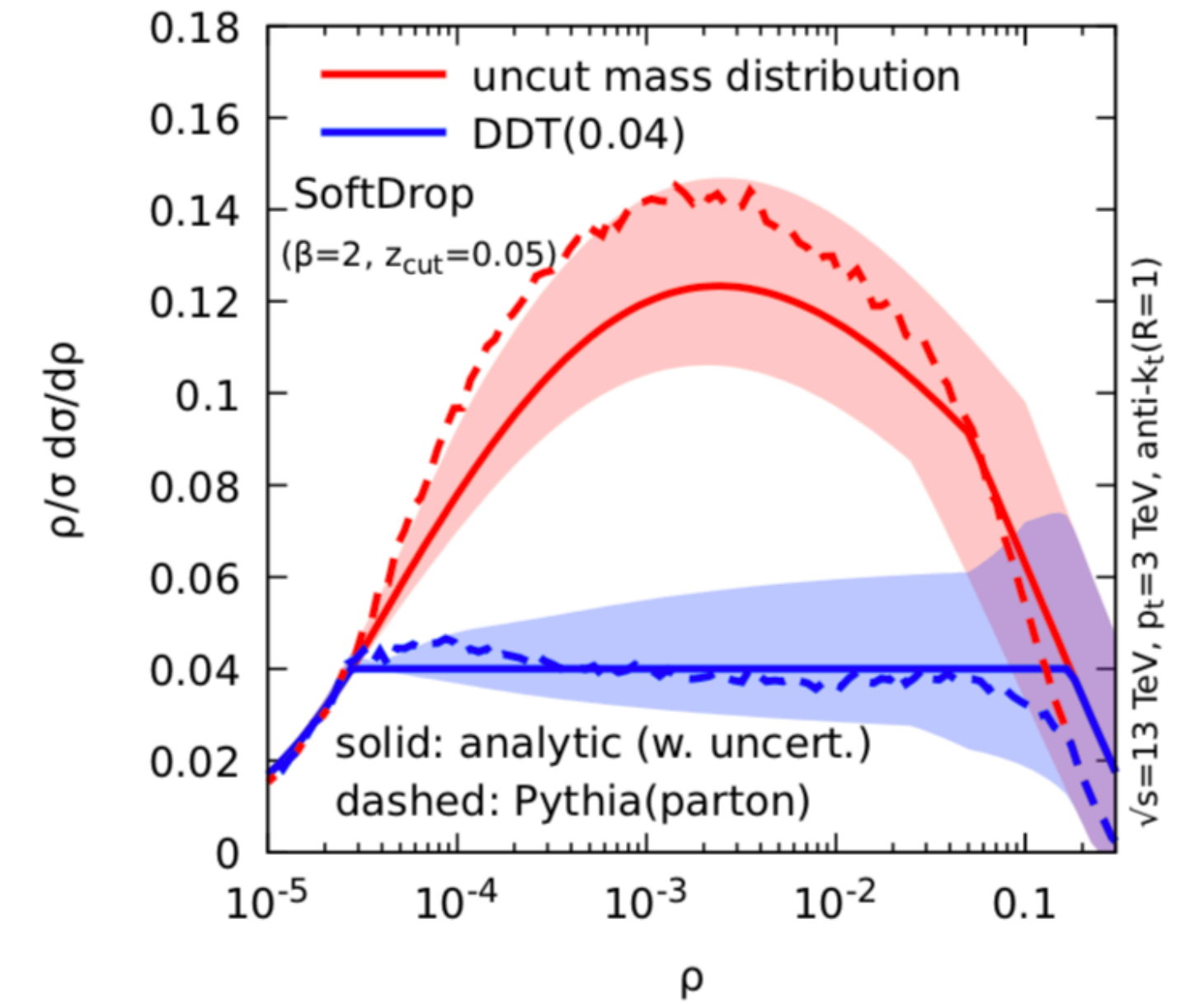
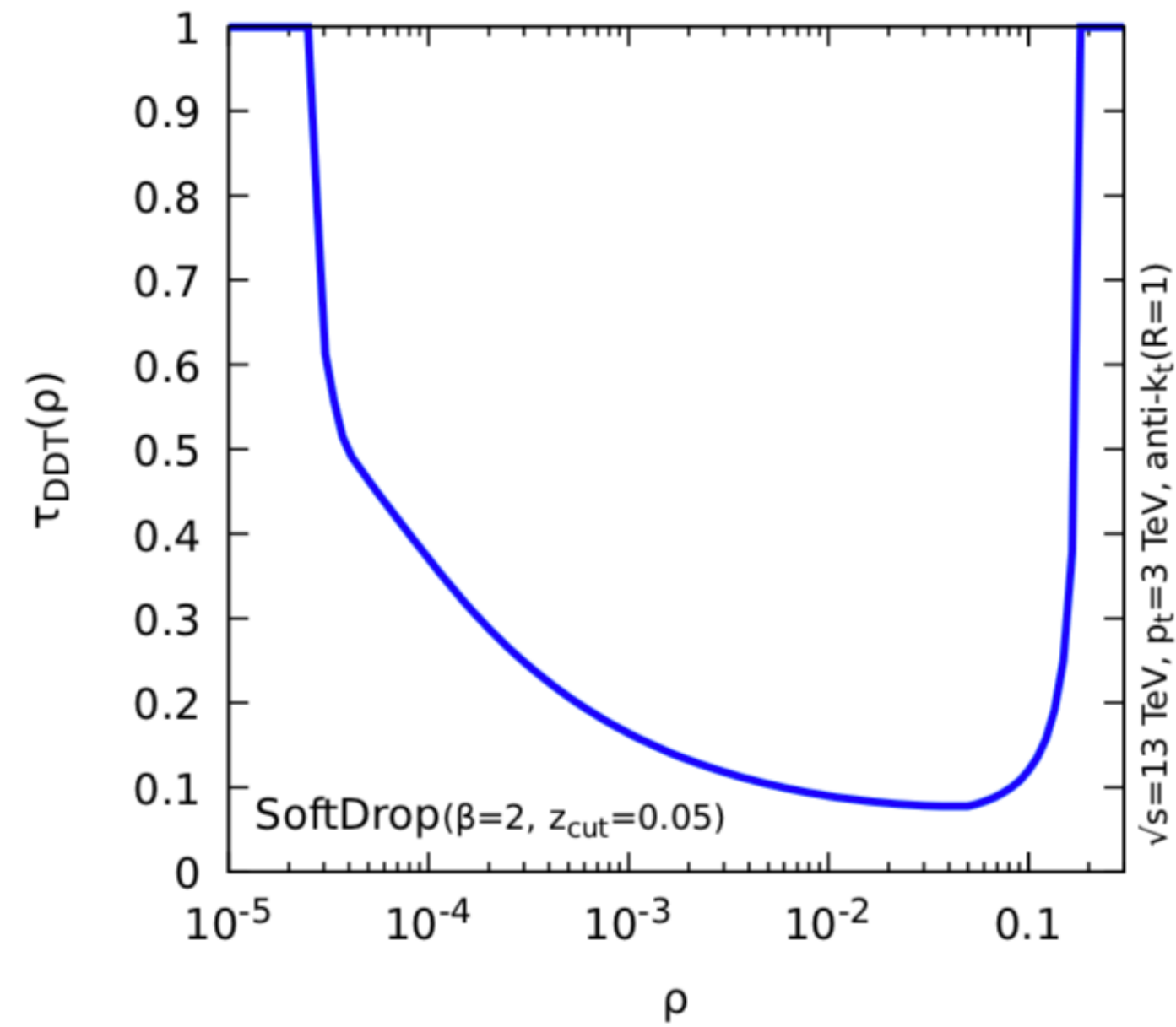
Results

- Find the value of tau that makes distribution flat

- Not necessarily an easy task in general, quite easy having analytical control

- To show-off we do it and seems to be working pretty reliably

Decorrelated taggers



Conclusions

- Jet-substructure has been successfully used for discrimination problems
- Analytical calculations in this field have also helped introducing new observables
- However often oversimplifications to overcome difficulties
- This calculation addresses some of them, and we produce some results with it!
 - The important thing is that with an analytical calculation one is able to extract some physics information on the problem at hand