Jet Pull

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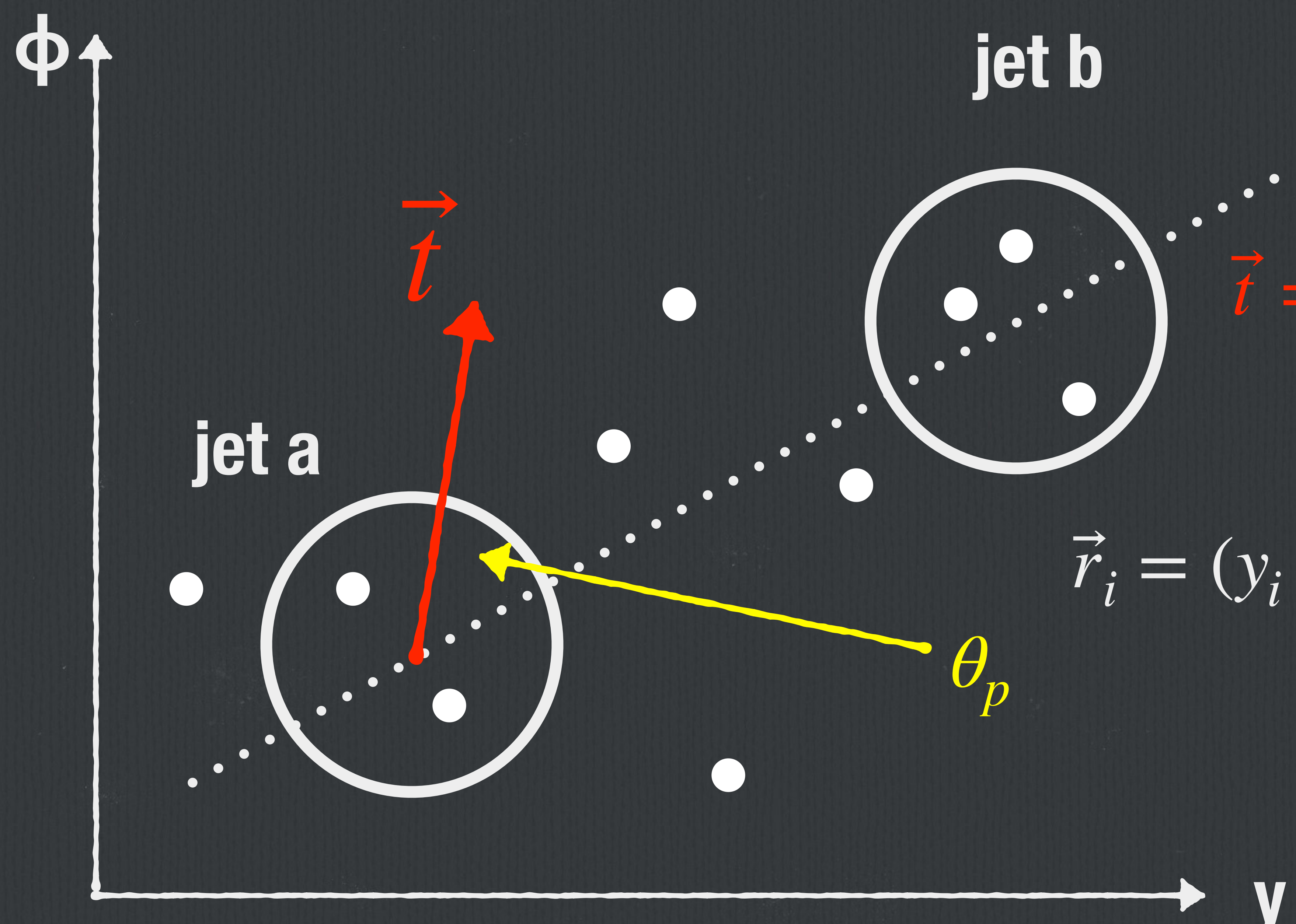
Outline

- Probing colour flow with jet pull
- Theory predictions for the pull angle
- IRC safe projections of jet pull
- Conclusion and Outlook

Probing colour properties

- one key aspect of the LHC physics program is the characterisation of (new) particle properties
- these include spin, CP properties and gauge charges
- in particular we would like to understand whether new particles carry colour charge
- a powerful way to extract this information is to study QCD radiation that accompanies the hard process
- this is often done with jet vetoes (central jet vetoes are a way of enhancing VBF against ggF Higgs production)
- jet shapes, which measure energy flow within a jet, are also sensitive to the colour flow of the jet environment

Jet Pull



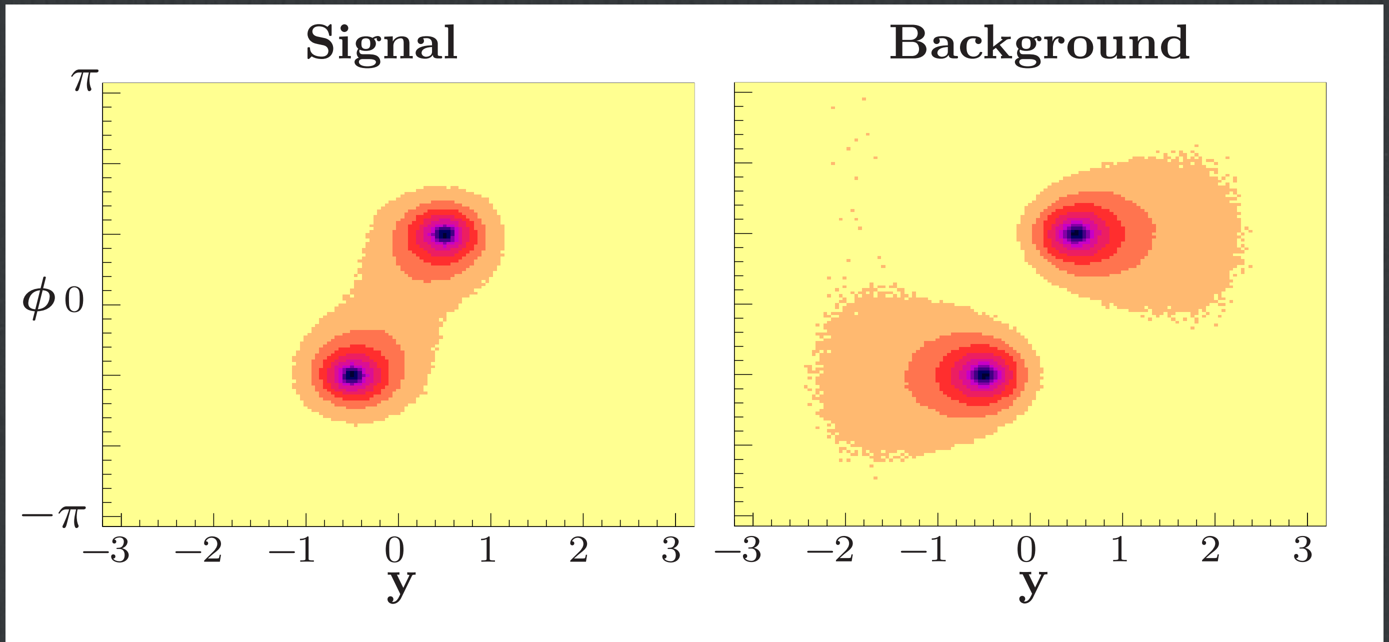
$$\vec{t} = \sum_{i \in jet\, a} \frac{p_{ti}}{p_t} |\vec{r}_i| \vec{r}_i$$

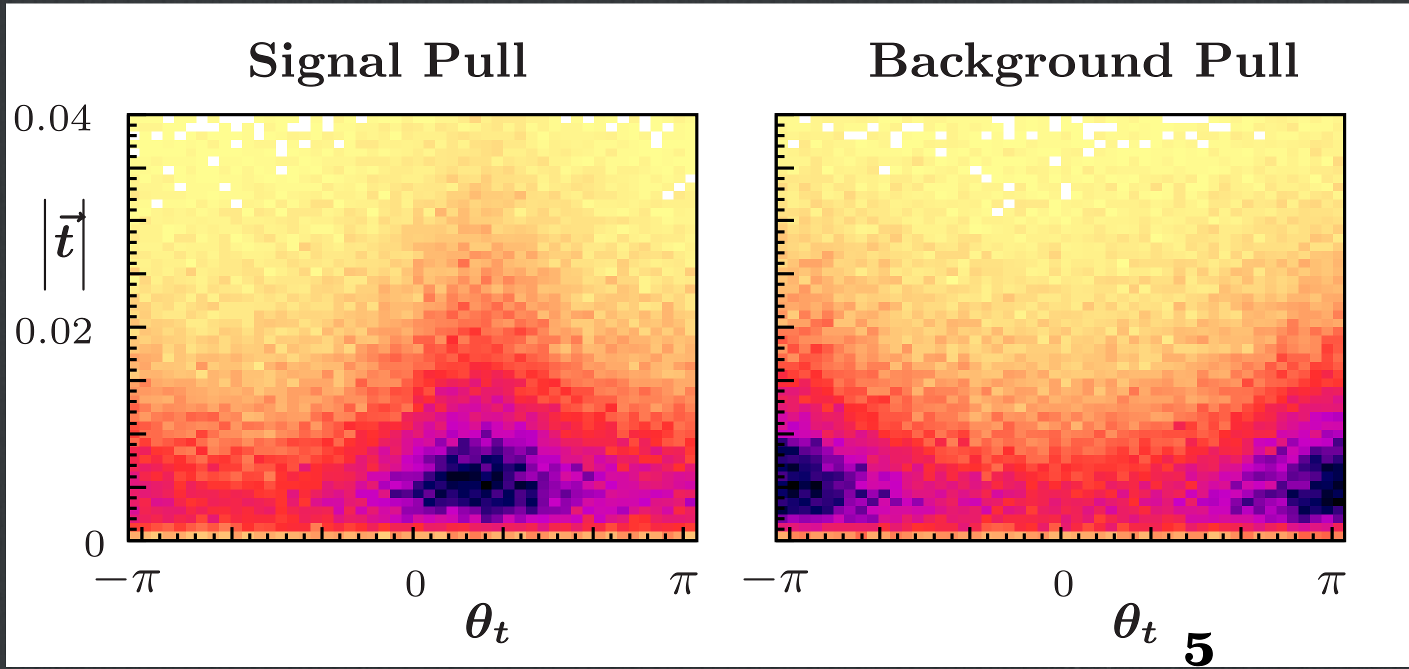
$$\vec{r}_i = (y_i - y_a, \phi_i - \phi_a)$$

Gallichio and Schwartz (2010)

H→bb vs g→bb

- consider radiation pattern of a colour singlet (signal) vs colour octet (background)
- simulation shows dominant colour-connections:
- between the two b's for the singlet
- between each b and the initial-state for the background

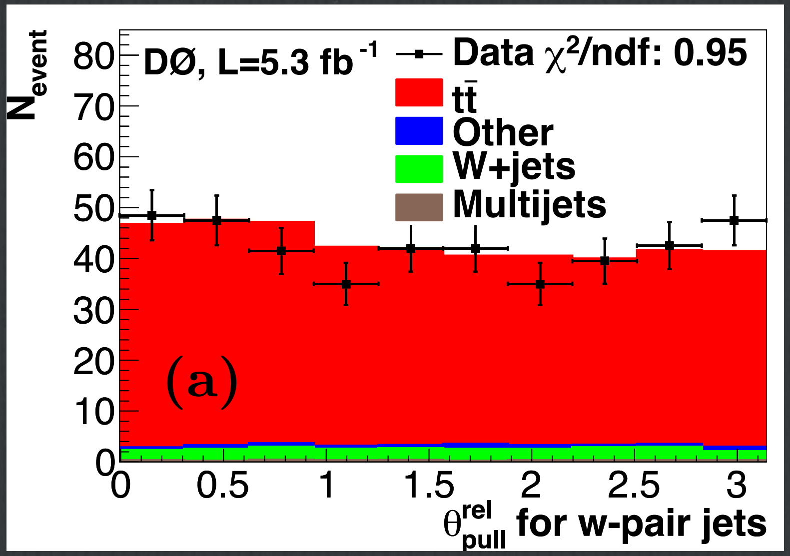




- pull angle shows much more sensitivity to colour flow than the pull magnitude

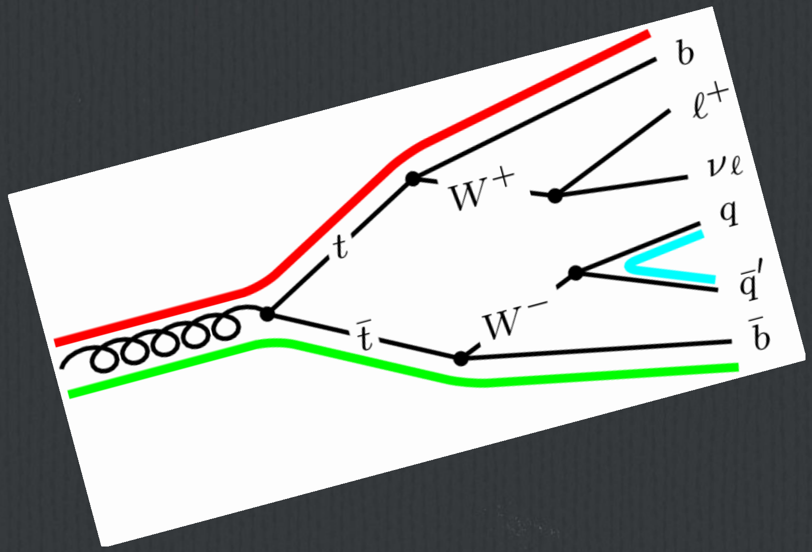
Gallichio and Schwartz (2010)

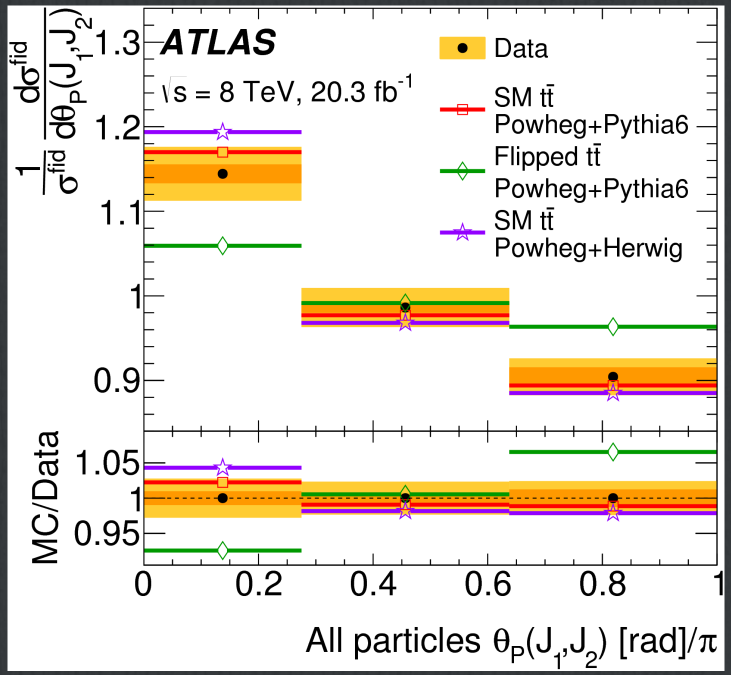
Experimental measurements



Phys.Rev. D83 (2011) 092002

- abundant production of top quarks offer nice lab for these studies
- pull angle can be measured on different types of colour connections





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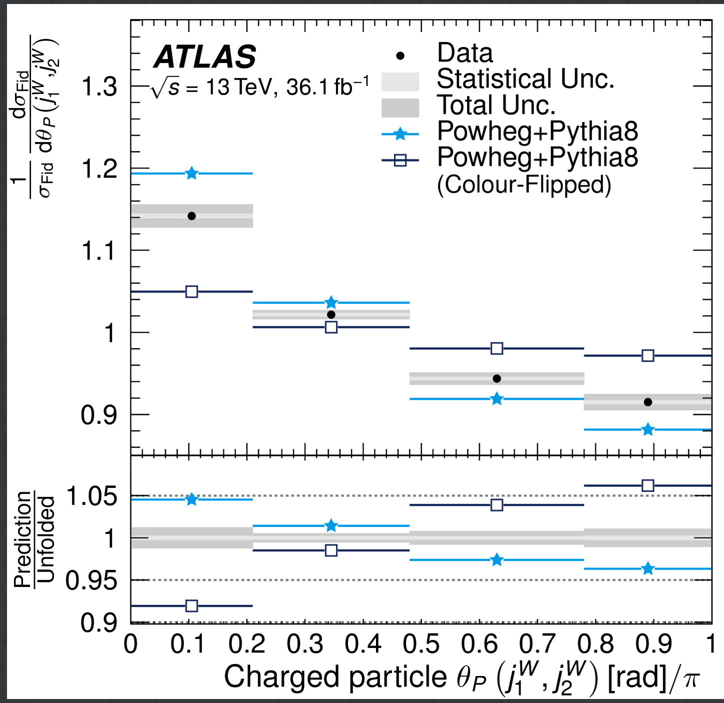
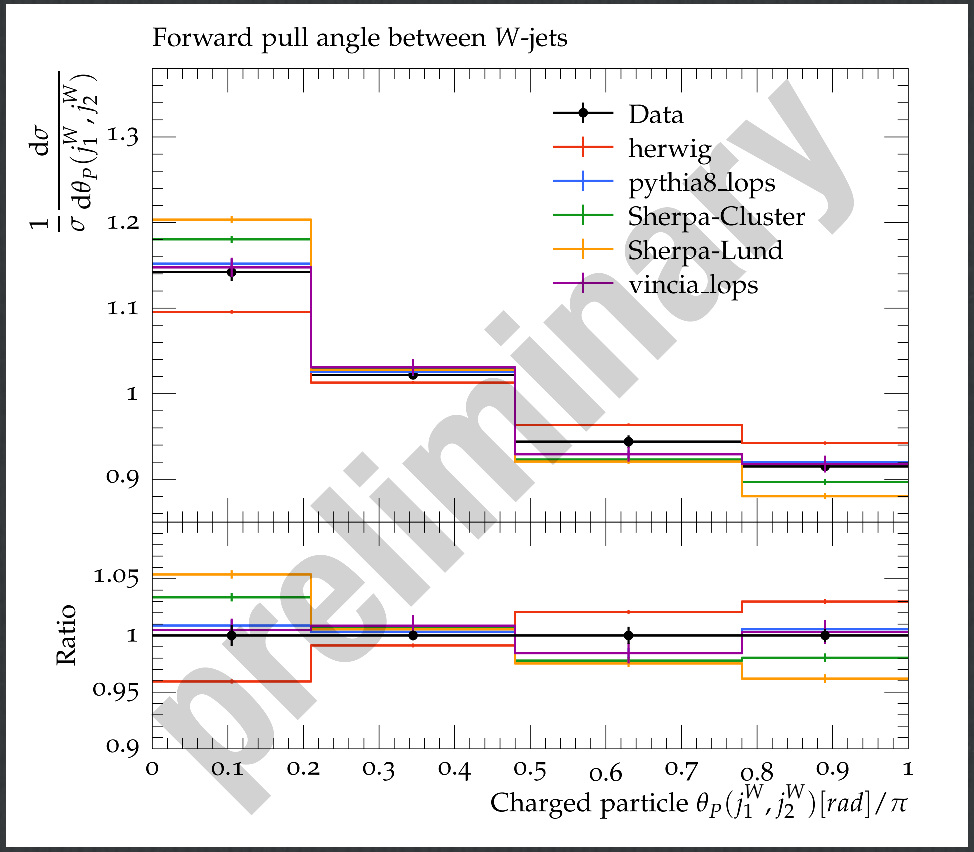


CMS ???

Eur. Phys. J. C 78 (2018)

Les Houches studies



courtesy of Helen Brooks

- astonishing precision of 13 TeV data allows for stringent test of different MC tools
- on-going study started at LH 19 aimed at understanding this observable better
- intricate interplay of different ingredients:
    - spread in parton shower modelling is comparable (if not bigger) than spread due to non-perturbative contributions

- we hope to achieve a clearer picture for the proceedings!

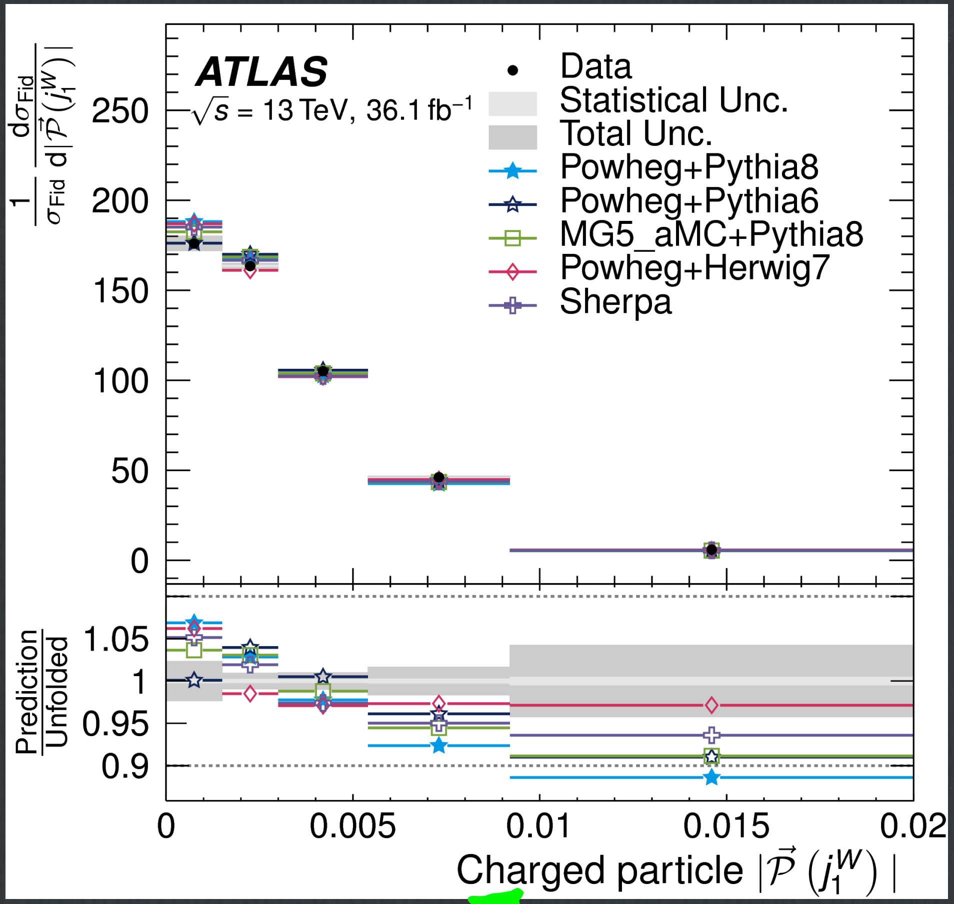
Can we make firmer theory predictions?

- besides MC study we can try and understand jet pull with analytic calculations
- we aim at a description that matches together fixed-order and resummed prediction
- let’s look at next-to-leading log (NLL)

- pull magnitude (or any component of the pull vector) is IRC safe:

- if two emissions p_1, p_2 become collinear, we are only sensitive to $p_1 + p_2$
- if emission p_1 becomes soft $p_{t1} \to 0$ and it does not contribute to the magnitude

- we can calculate this distribution in perturbation theory!



not quite IRC safe!

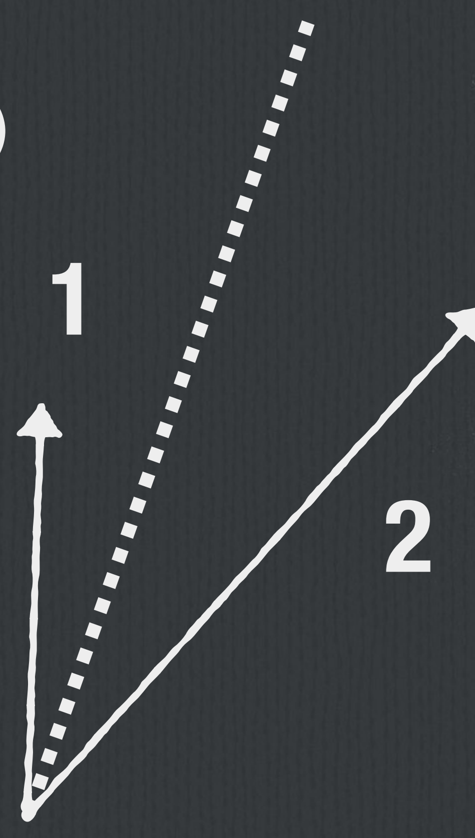
Resummation of the pull magnitude

- we consider the resummation of the magnitude of a 2-D vector:

$$t = |\vec{t}\,| = \left| \sum_{i \in jet\, a} \frac{p_{ti}}{p_t} |\vec{r}_i| \vec{r}_i \right|$$

- situation very similar to well-known Q_T resummation (but in the final state)
- we we work in a conjugate Fourier space
- on the other hand, its scaling properties are the similar to the jet mass

$$t = |\vec{t}_1 + \vec{t}_2| \simeq |z(1-z)^2 - (1-z)z^2|\,\theta_{12}^2$$
$$= |z(1-z)(1-2z)|\,\theta_{12}^2$$



NLL resummation

- in the collinear limit, i.e. up to term of O(R^2) the all-order behaviour is

$$\frac{1}{\sigma_0}\frac{d\sigma}{dt} = \int_0^\infty dbbJ_0(bt)\exp\left[Lf_1(\lambda) + f_{2c}(\lambda)\right] S_{ng}(\lambda)$$

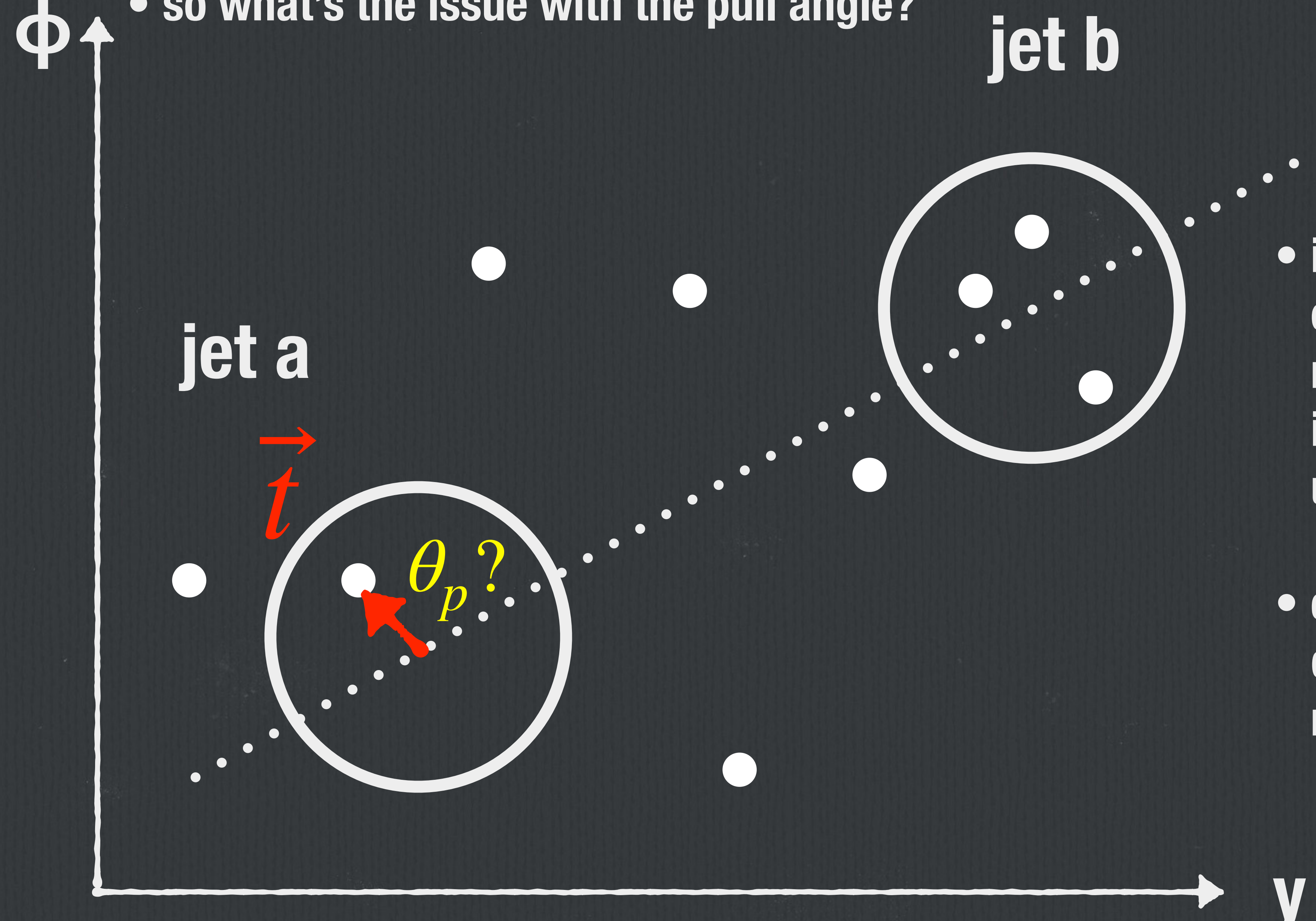
$$\lambda = \alpha_s \beta_0 \log\left(b\frac{e^{\gamma_E}}{2}\right)$$

- remarkably, the f_1 and f_{2c} have the same function have the same functional form as for the jet mass distribution

- S_{ng} accounts for non-global logarithms, not investigated yet

Pull angle: IRC un-safety

- so what's the issue with the pull angle?



- if particle i is soft, its contribution to the magnitude vanishes but its pull angle is undefined
- equivalently: in order to define the pull angle we need two emissions

aside: Sudakov safety

- we know of other observables that suffers from similar problems (ratio of angularities, soft-drop momentum balance)
- we can make sense of these observables if we are able to resolve the singularities with the help of a (safe) companion variable
- we say that the IRC unsafe variable u is Sudakov safe if there exists an IRC safe observable s such that

$$p(u) = \int ds\, p(u \mid s)\, p(s)$$

- p(s) must be calculated to all-orders in order to (Sudakov) suppress the s=0 singularity

Larkoski, Thaler (2013)
Larkoski, SM, Thaler (2015)

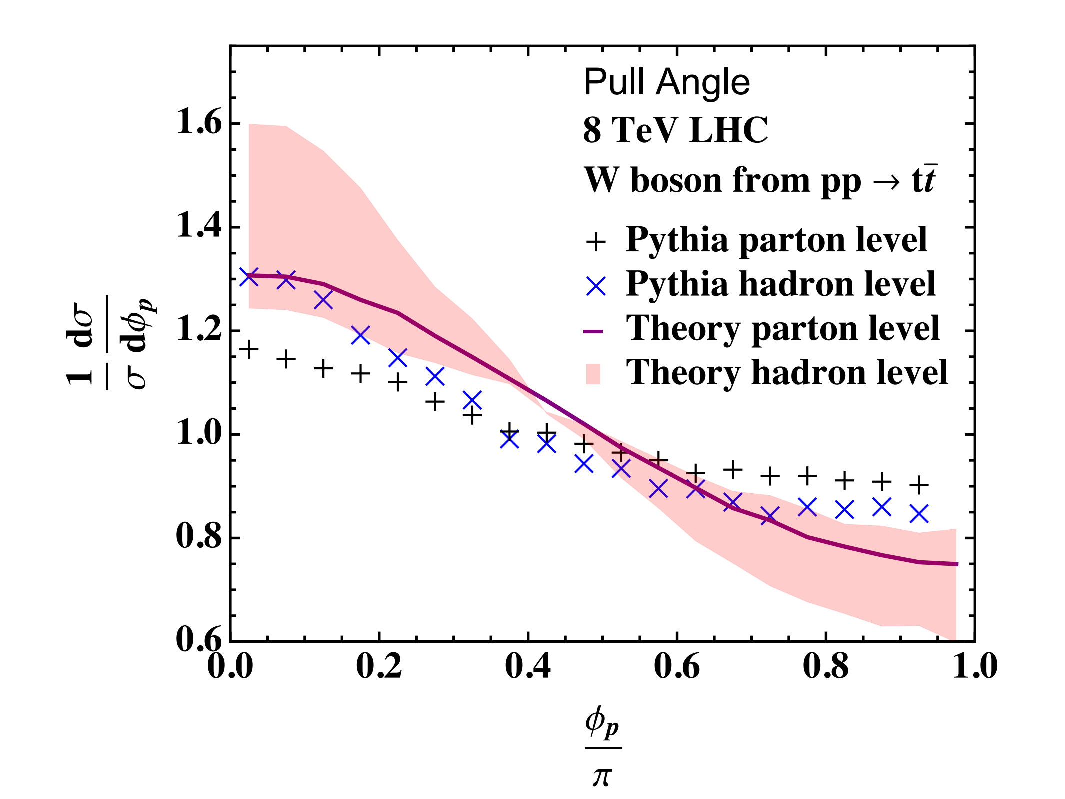
Perturbative calculation

- our natural candidate for the safe companion is the pull magnitude itself
- thus we can write (apologies but in our paper we call the pull angle ϕ_p rather than θ_p)

$$\frac{1}{\sigma}\frac{d\sigma}{d\phi_p} = \int dt\, p(\phi_p\,|\,t)\, p^{res}(t)$$

computed at fixed-order

resummed at NLL in the collinear limit (but with no non-global logs a the moment)



Non-perturbative effects

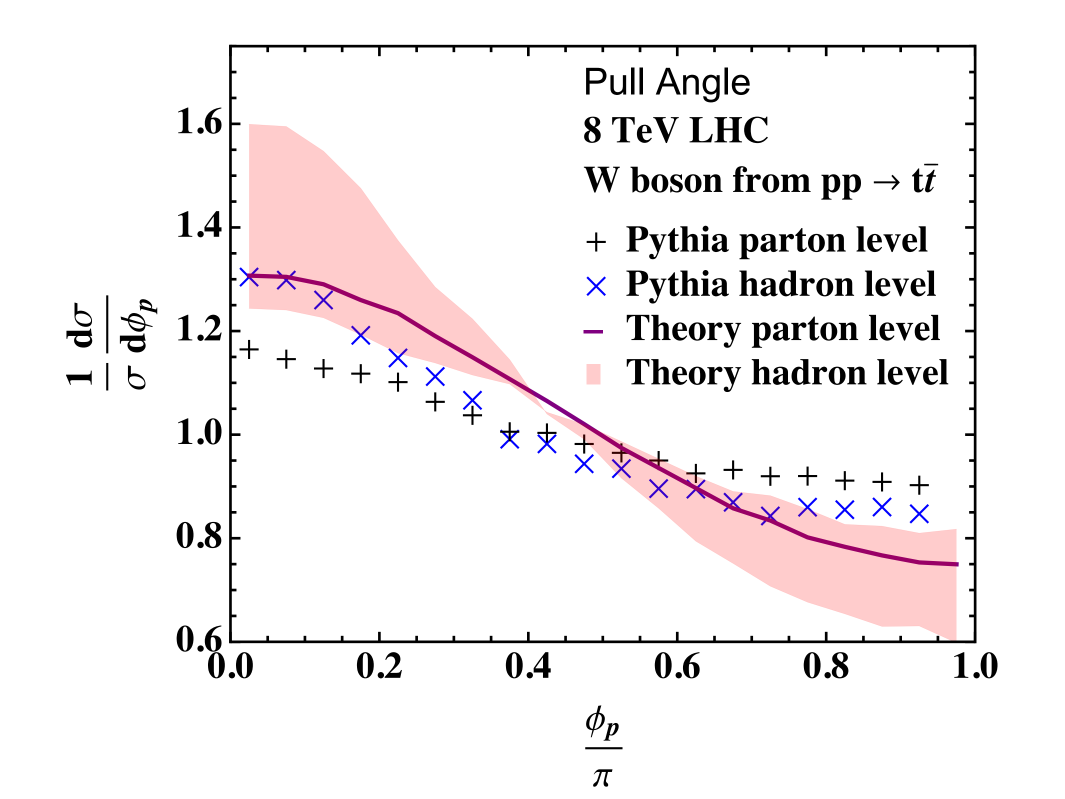
- for IRC safe observable we can make relatively strong statements about the scaling of non-perturbative corrections
- for Sudakov safe observable we don't have such luxury
- we make some rough estimate

$$p_{np}(t,\phi_p) \propto \tanh\left(\frac{1}{a\phi_p(2\pi-\phi_p)}\right)\delta\left(t-\frac{\Omega}{E_J}\right)$$

$$\Omega \simeq \Lambda_{QCD}\,, \quad 0 \le a \le 0.25$$

uniform to peaked in ϕ_p=0

$$p_{tot} = p_{np} \overrightarrow{\otimes} p_{pert}$$

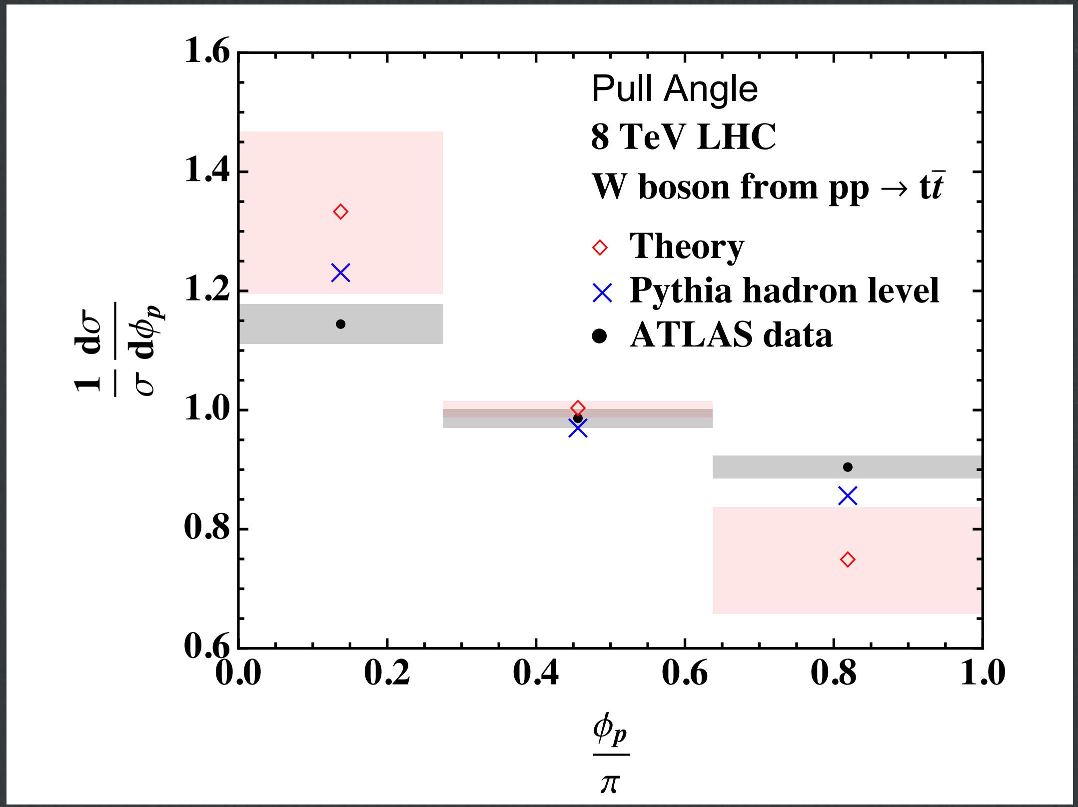


Comparison to the data

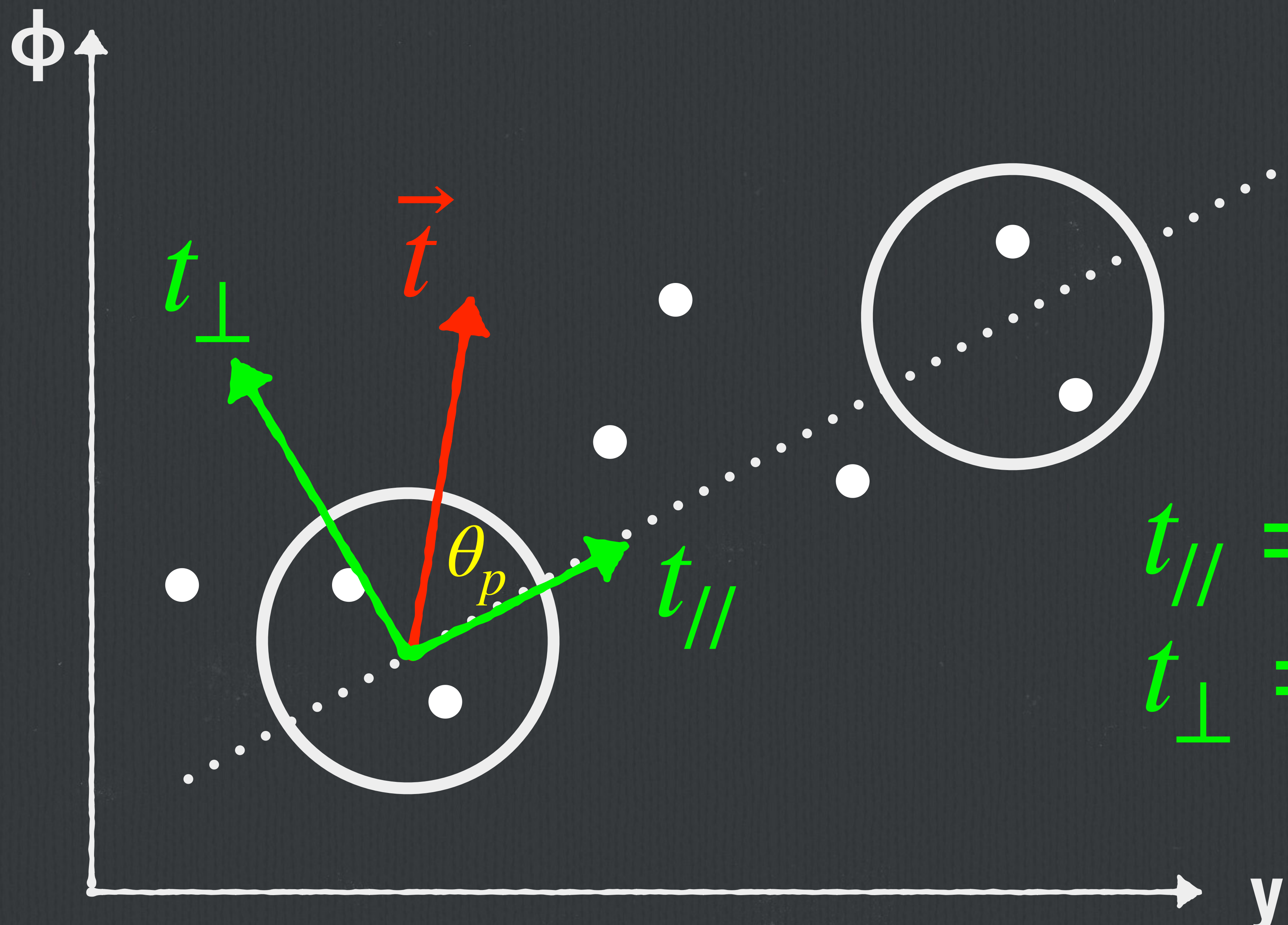
- we can now compare to the Run I measurement by the ATLAS collaboration
- our calculation is in fair agreement with the data (similar to the Monte Carlo prediction)
- however it suffers from large theory uncertainties

- perturbative uncertainties: Sudakov safe observables do not admit standard expansion in terms of Feynman diagrams: we have to combine fixed-order and resummed ingredients (and several questions remain open)

- non-pert. uncertainties: lack of IRC safety prevents us from clearly separating perturbative and non-perturbative regions



Safe uses of jet pull



$$t_{//} = |\vec{t} \cdot \hat{n}_{//}|$$
$$t_\perp = |\vec{t} \cdot \hat{n}_\perp|$$

NLL resummation

- the new variables $t_{//}$ and $t_{\perp}$ are IRC safe: we can combine fixed-order and resummed prediction in the standard way
- their resummation is relatively straightforward
- it closely follows the formalism develop for analogous projections of Q_T i.e. a_T (and ϕ^*)

Banfi, Dasgupta
and Duran Delgado (2010)

$$\frac{1}{\sigma_0}\frac{d\sigma^{collinear}}{dt_{\perp}} = \frac{1}{\pi}\int_0^{\infty} db \cos(bt_{\perp}) \exp\left[L f_1(\lambda) + f_{2c}(\lambda)\right]$$

Towards phenomenology

- in order achieve full NLL resummation we will have to include
- soft radiation at wide angle (through an expansion in powers of the jet radius)
- non-global logarithms
- notice that soft radiation is crucial (this is what we want to probe) but does depend on the jet environment
- we have to specify the process and the event selection, e.g.

$$pp \to W(\to qq') + jet$$

- and measure the pull between the W subjets

Conclusions & Outlook

- jet pull is an interesting observable to probe colour flow in LHC events
- pull angle distributions measured by D0 and ATLAS: fair agreement with MC but small experimental uncertainties expose shortcomings of simulations
- theory analysis possible but conclusions are not very firm because of IRC unsafely
- this can be alleviated by looking at safe projections of the jet pull vector (work in progress)

THANK YOU VERY MUCH!