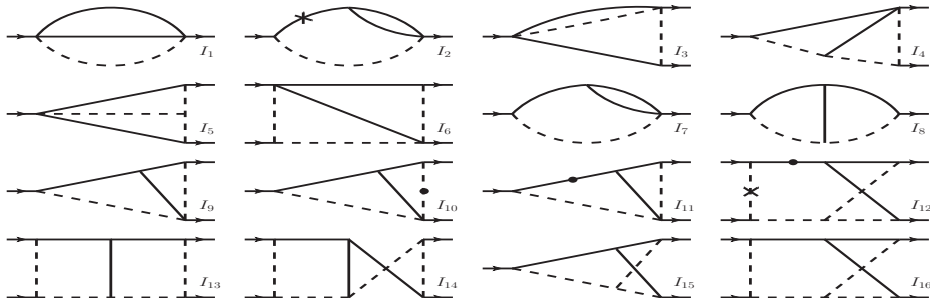


NNLO real corrections to Higgs boson pair production in the large top mass limit

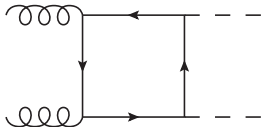
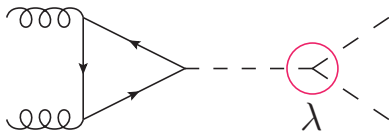
Florian Herren | 14.08.2019

based on [Davies, FH, Mishima, Steinhauser JHEP 1905 (2019) 157, [WIP]]

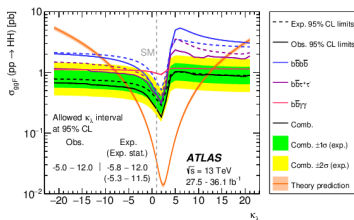
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Why Higgs boson pair production?



- direct probe of cubic term of the scalar potential
- completely fixed in the SM

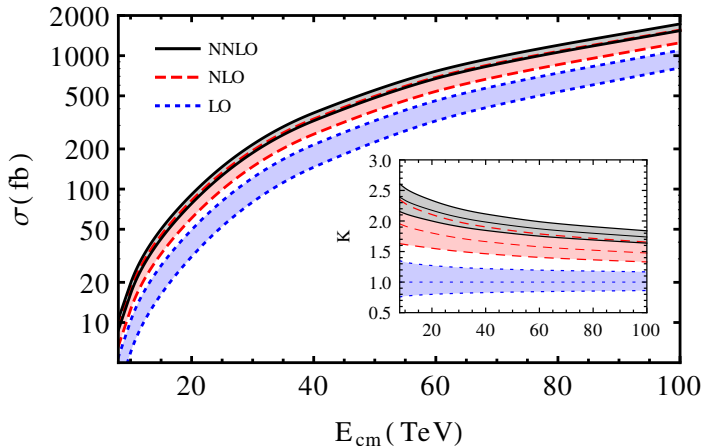


taken from arXiv:1906.02025 [hep-ex]

- current limit: $-5.0 < \kappa_\lambda < 12.0$
- HL-LHC: $0.5 < \kappa_\lambda < 2.4$ [WG2 Report '19]
- FCC-hh: 5% accuracy [FCC CDR '19]

Why higher orders?

$$pp \rightarrow HH + X$$



Hadronic c.o.m energy, taken from [de Florian, Mazzitelli '13]

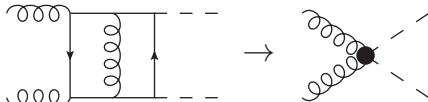
Status at NLO

exact LO results known for 3 decades:

[Eboli, Marques, Novaes, Natale, '87],[Glover, van der Bij '88],[Plehn, Spira, Zerwas '96]

Various approximations used at NLO, mostly in the limit of (infinitely) heavy top quark:

- Born-improved HEFT [Dawson, Dittmaier, Spira '98]
- LME [Grigo, Hoff, Melnikov, Steinhauser '13], [Degrassi, Giardino, Gröber '16]
- $FT_{\text{approx}}, FT'_{\text{approx}}$ [Maltoni, Vryonidou, Zaro '14]
- LME + Threshold expansion [Gröber, Maier, Rauh '17]



Furthermore there exist approximations addressing other parts of the phase-space:

- High-energy expansion [Davies, Mishima, Steinhauser, Wellmann '18 + '19],[Mishima '18]
- Small- p_T expansion [Bonciani, Degrassi, Giardino, Gröber '18]

Exact numerical calculations available, covering the whole phase-space

[Borowka, Greiner, Herinrich, Jones, Kerner, Schlenk, Zirke '16].

[Baglio, Campanario, Glaus, Mühlleitner, Spira, Streicher '18]

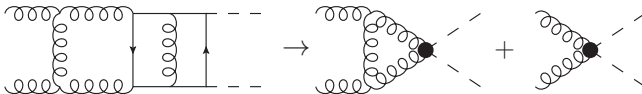
Recently have been combined with high-energy expansion

[Davies, Heinrich, Jones, Kerner, Mishima, Steinhauser, Wellmann '19]

Status beyond NLO

The next step is to improve the situation at NNLO, here only few results exist:

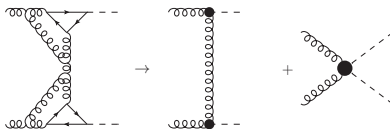
- HEFT [de Florian, Mazzitelli '13], [Grigo, Melnikov, Steinhauser '14]
- LME in SV-approximation [Grigo, Hoff, Steinhauser '14]
- FT_{approx} [Grazzini, Heinrich, Jones, Kallweit, Kerner, Lindert, Mazzitelli '18]



→ our goal is to improve upon the LME by also considering the real radiation

At N^3 LO the basic building differing from single Higgs production recently have been computed for HEFT:

- Wilson coefficient C_{HH} [Spira '16],[Gerlach, FH, Steinhauser '18]
- 2 loop box-type diagrams [Banerjee, Borowka, Dhani, Gehrmann, Ravindran '18]



- inclusive, partonic cross-section at NNLO

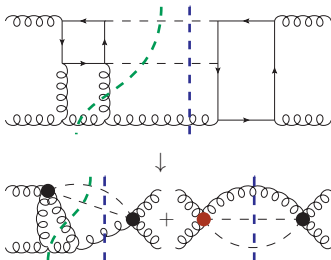
$$\sigma_{ij}(s, m_H, m_t) = \delta_{ig}\delta_{jg}\sigma_{gg}^{(0)} + \frac{\alpha_s}{\pi}\sigma_{ij}^{(1)} + \left(\frac{\alpha_s}{\pi}\right)^2\sigma_{ij}^{(2)} + \dots$$
$$\sigma_{ij}^{(2)} = \underbrace{\delta_{ig}\delta_{jg}\sigma_{gg,\text{virt}}^{(2)}}_{\text{[Grigo,Hoff,Steinhauser '14]}} + \underbrace{\sigma_{ij,\text{real}}^{(k)}}_{\text{this work}}$$

- so far only known in $m_t \rightarrow \infty$ limit, full computation not feasible
→ asymptotic expansion in large top quark mass

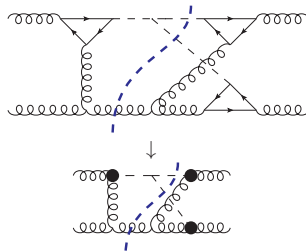
$$\sigma_{ij}^{(l)}(s, m_H, m_t) \approx \sum_k \rho^k \sigma_{ij}^{(l),k}(x)$$
$$\rho = m_H^2/m_t^2, \quad x = m_H^2/s$$

What are we computing?

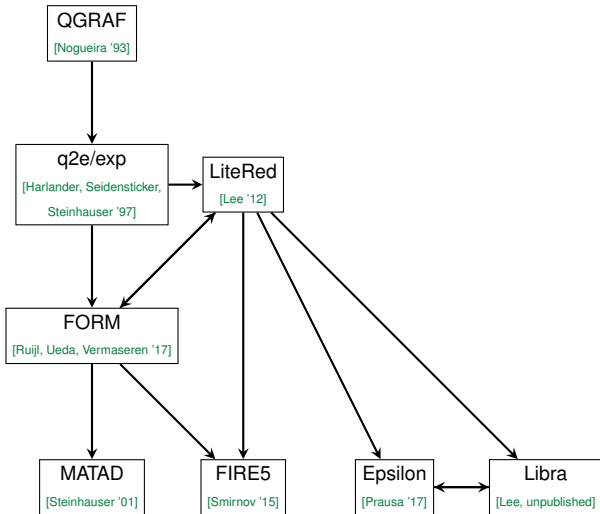
We use reverse unitarity to express the real corrections through cuts of 5-loop diagrams with 2 masses in forward scattering kinematics.



- 2 1-loop tadpoles + 3-loop phase space integral
- 1-loop tadpole + 2-loop tadpole + 2-loop phase space integral



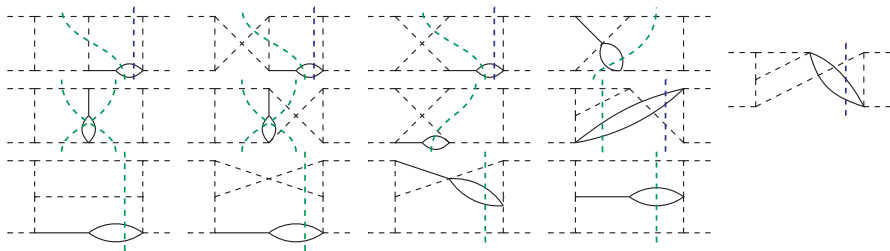
- 3 1-loop tadpoles + 2-loop phase space integral



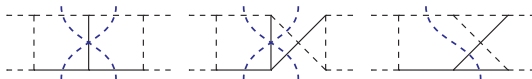
Families

We use LiteRed [Lee '12] for minimizing families, identifying common sectors and performing the reduction

We end up with 13 three-loop,



and 3 two-loop phase space integral families:



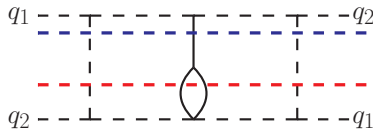
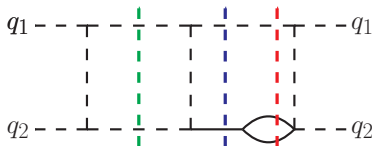
Differential equations

- 3 systems of differential equations
- We bring them to Fuchsian form using Epsilon [Prausa '17]:

$$\partial_x \vec{T} = \sum_i \frac{M^{(i)}(\epsilon)}{x - x_i} \vec{T} \quad x = m_h^2/s$$

- They contain the letters:

$x_i =$	0	1	$\frac{1}{4}$	-1	$-\frac{1}{4}$	$e^{\pm i\frac{\pi}{3}}$	$-\frac{1}{3}$
2 loop, 3 particle	✓	✓	✓			✓	✓
3 loop, 3 particle	✓	✓	✓				
3 loop, 4 particle	✓	✓	✓	✓	✓		



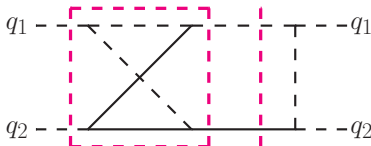
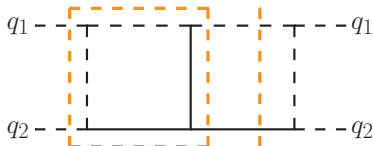
Differential equations

- 3 systems of differential equations
- We bring them to Fuchsian form using Epsilon [Prausa '17]:

$$\partial_x \vec{l} = \sum_i \frac{M^{(i)}(\epsilon)}{x - x_i} \vec{l} \quad x = m_h^2/s$$

- They contain the letters:

$x_i =$	0	1	$\frac{1}{4}$	-1	$-\frac{1}{4}$	$e^{\pm i \frac{\pi}{3}}$	$-\frac{1}{3}$
2 loop, 3 particle	✓	✓	✓			✓	✓
3 loop, 3 particle	✓	✓	✓				
3 loop, 4 particle	✓	✓	✓	✓	✓		



- We proceed to transform the differential equations to ϵ -form using Libra [\[Lee, unpublished\]](#):

$$\partial_x \vec{T} = \sum_i \frac{M^{(i)}(\epsilon)}{x - x_i} \vec{T} \quad \rightarrow \quad \partial_x \vec{J} = \epsilon \sum_i \frac{\tilde{M}^{(i)}}{x - x_i} \vec{J}$$

- The solution to DEs in this form can be written as Goncharov Polylogarithms
- Some of the $M^{(i)}(\epsilon)$ contain Eigenvalues of the form $\epsilon \pm 1/2$:

	$\sqrt{1 - 4x}$	$\sqrt{1 + 4x}$	$\sqrt{(1 + 3x)(1 - x)}$
2 loop, 3 particle	✓		✓
3 loop, 3 particle	✓		
3 loop, 4 particle	✓	✓	

→ Need to find change of variable to transform square roots into rational functions

- We proceed to transform the differential equations to ϵ -form using Libra [Lee, unpublished]:

$$\partial_x \vec{I} = \sum_i \frac{M^{(i)}(\epsilon)}{x - x_i} \vec{I} \quad \rightarrow \quad \partial_x \vec{J} = \epsilon \sum_i \frac{\tilde{M}^{(i)}}{x - x_i} \vec{J}$$

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- Some of the $M^{(i)}(\epsilon)$ contain Eigenvalues of the form $\epsilon \pm 1/2$:

	$\sqrt{1-4x}$	$\sqrt{1+4x}$	$\sqrt{(1+3x)(1-x)}$
2 loop, 3 particle	✓		✓
3 loop, 3 particle	✓		
3 loop, 4 particle	✓	✓	

- In the DE for the three-loop, three-particle cuts only $\sqrt{1-4x}$ appears

$$x = \frac{y}{(1+y)^2} \quad x \in (0, 1/4] \rightarrow y \in (0, 1]$$

- We proceed to transform the differential equations to ϵ -form using Libra [Lee, unpublished]:

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- Some of the $M^{(i)}(\epsilon)$ contain Eigenvalues of the form $\epsilon \pm 1/2$:

	$\sqrt{1-4x}$	$\sqrt{1+4x}$	$\sqrt{(1+3x)(1-x)}$
2 loop, 3 particle	✓		✓
3 loop, 3 particle	✓		
3 loop, 4 particle	✓	✓	

- In the DE for the three-loop, four-particle cuts both $\sqrt{1-4x}$ and $\sqrt{1+4x}$ appear

$$x = \frac{t^4 + 1}{8t^2} \quad x \in (0, 1/4] \rightarrow t = e^{i\phi} : \phi \in (1/4, 0]$$

- We proceed to transform the differential equations to ϵ -form using Libra [Lee, unpublished]:

$$\partial_x \vec{T} = \sum_i \frac{M^{(i)}(\epsilon)}{x - x_i} \vec{T} \quad \rightarrow \quad \partial_x \vec{J} = \epsilon \sum_i \frac{\tilde{M}^{(i)}}{x - x_i} \vec{J}$$

- The solution to DEs in this form can be written as Goncharov Polylogarithms
- Some of the $M^{(i)}(\epsilon)$ contain Eigenvalues of the form $\epsilon \pm 1/2$:

	$\sqrt{1-4x}$	$\sqrt{1+4x}$	$\sqrt{(1+3x)(1-x)}$
2 loop, 3 particle	✓		✓
3 loop, 3 particle	✓		
3 loop, 4 particle	✓	✓	

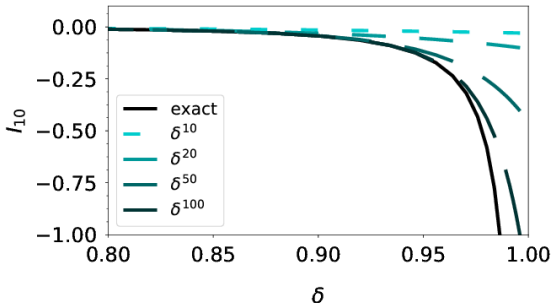
- In the DE for the two-loop, three-particle cuts $\sqrt{1-4x}$ and $\sqrt{(1+3x)(1-x)}$ appear
- Can not be rationalized together, however integrations over $\sqrt{(1+3x)(1-x)}$ do not contribute to cross-section

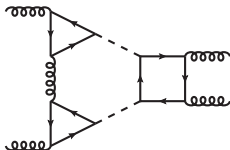
$$x = \frac{y}{(1+y)^2} \quad x \in (0, 1/4] \rightarrow y \in (0, 1]$$

Boundary conditions and δ -expansion

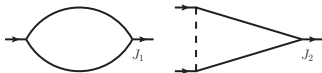
- Parametrize necessary phase-space integrals through $\delta = 1 - 4x$
- Expand around threshold $\delta \approx 0$
- Minimization and other manipulation of Goncharov Polylogarithms with PolyLogTools [Duhr, Dulat '19]
- We also obtained solutions to the DE with the ansatz

$$\vec{T} = \sum_{i,j,k} \vec{c}_{i,j,k} \epsilon^i \delta^{(2j+1)/2} \ln^k \delta$$





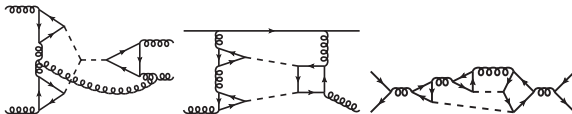
- only gg channel contributes
- already known in expansion around threshold [Grigo, Hoff, Melnikov, Steinhauser '13]
- recomputed exactly in $m_h^2/s \rightarrow 2$ MIs



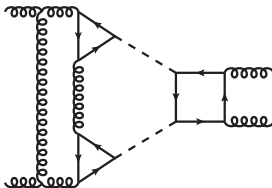
- collinear counterterm

$$\sigma_{gg, \text{coll}}^{(2), n_h^3} = 2 \int_{1-\delta}^1 dz P_{gg}(z) \sigma_{gg}^{(1), n_h^3}(x/z)$$

- three contributing channels gg , gq and $q\bar{q}$:

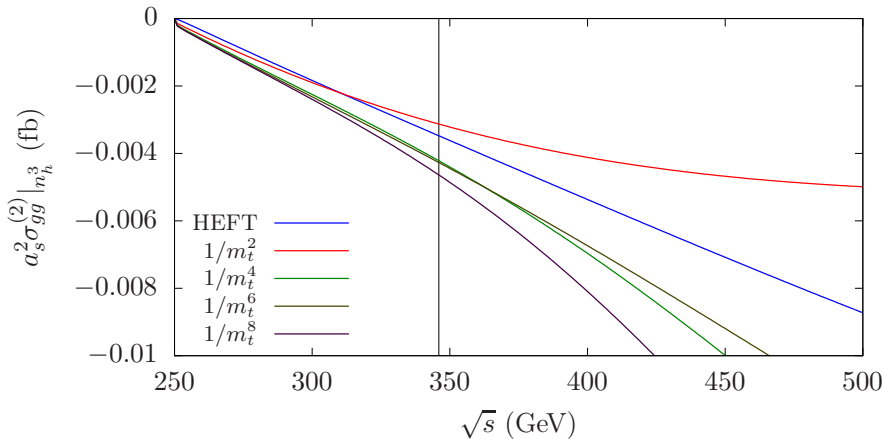


- expanded to $1/m_t^8$
- combine them with virtual corrections



Results

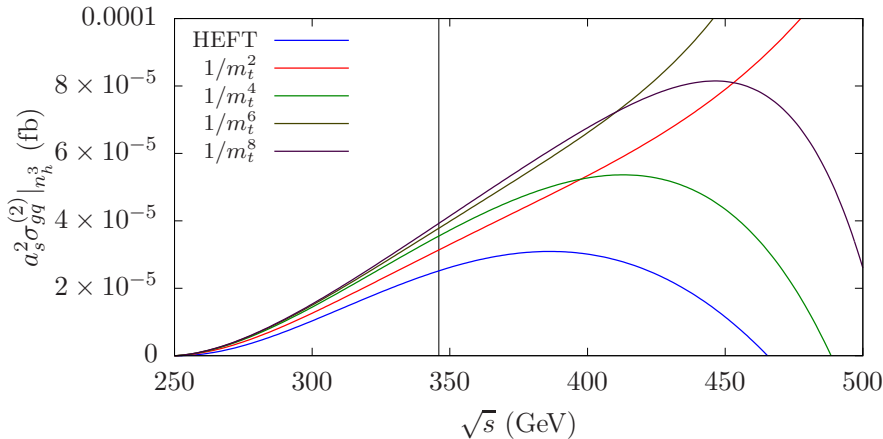
[Davies, FH, Mishima, Steinhauser JHEP 1905 (2019) 157]



■ Convergence only below $2m_t$

Results

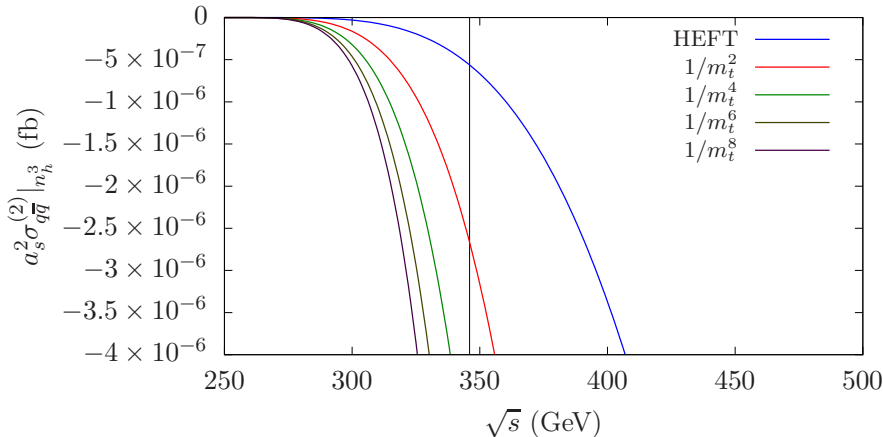
[Davies, FH, Mishima, Steinhauser JHEP 1905 (2019) 157]



■ Convergence only below $2m_t$

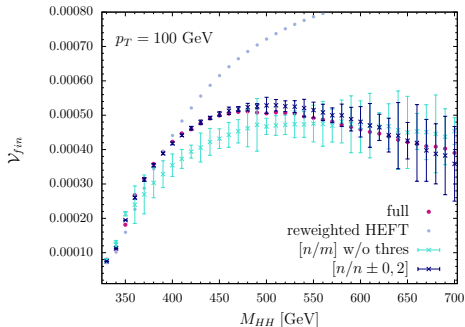
Results

[Davies, FH, Mishima, Steinhauser JHEP 1905 (2019) 157]



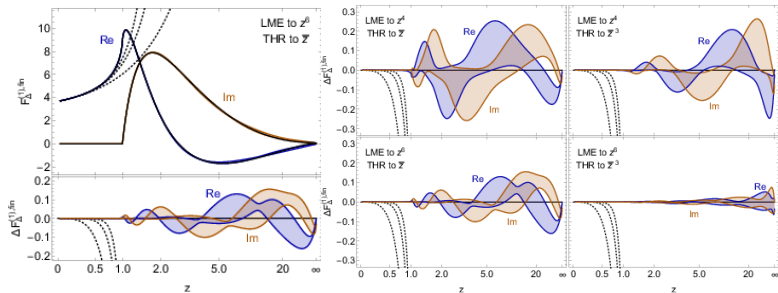
■ No convergence, even below $2m_t$

- LME is only useful in a limited part of phase space
- Can be combined with threshold expansion
- Works well for virtual part:



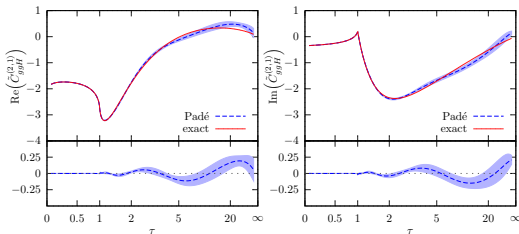
taken from [Gröber, Maier, Rauh '17]

- LME is only useful in a limited part of phase space
- Can be combined with threshold expansion
- Works well for virtual part
- Higher orders in LME crucial, e.g. for NLO ggH formfactor:



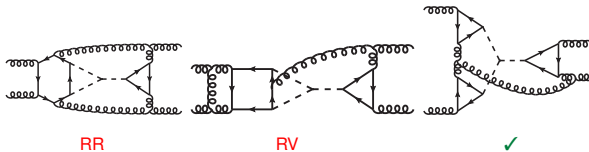
taken from [Davies, Gröber, Maier, Rauh, Steinhauser '19]

- LME is only useful in a limited part of phase space
- Can be combined with threshold expansion
- Works well for virtual part
- Higher orders in LME crucial
- Recently contributions with light fermion loops at NNLO confirmed by analytic computation:



taken from [Harlander, Prausa, Usovitch '19]

- Exact computation already very challenging for $gg \rightarrow H$
- Expansions needed, in particular for $gg \rightarrow HH$



- Obtain top-mass suppressed terms at NNLO to Higgs boson pair production
- Computed phase-space MIs, both, in an expansion around threshold and exactly
→ rather generic to any pair-production process of massive particles
- Byproduct: NLO real corrections without expanding around threshold
- Work in progress: cross-sections for double real and remaining real-virtual contributions