

# Exclusive nucleus- $\nu$ cross sections from quantum computers

Alessandro Roggero

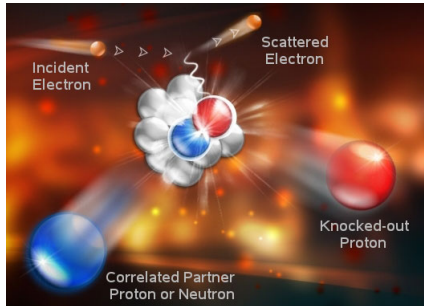


figure credit: JLAB collab.

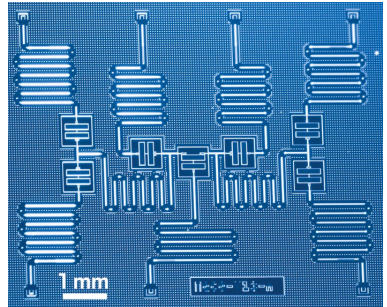


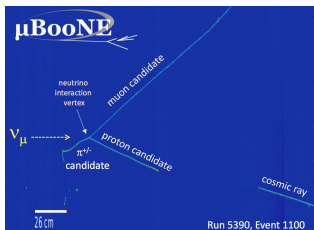
figure credit: IBM



FNAL – 09 May, 2019



# Exclusive cross sections in neutrino oscillation experiments



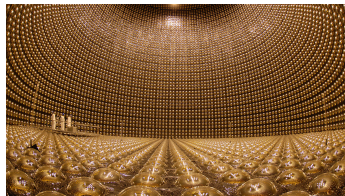
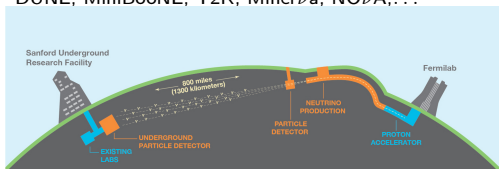
## Goals for $\nu$ oscillation exp.

- neutrino masses
- accurate mixing angles
- CP violating phase

$$P(\nu_\alpha \rightarrow \nu_\alpha) = 1 - \sin^2(2\theta) \sin^2 \left( \frac{\Delta m^2 L}{4E_\nu} \right)$$

- need to use measured reaction products to constrain  $E_\nu$  of the event

DUNE, MiniBooNE, T2K, Minerva, NO $\nu$ A, ...



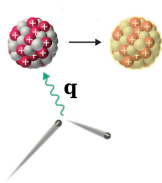
# Idealized algorithm for exclusive processes at fixed $q$

- prepare the target ground state



## Idealized algorithm for exclusive processes at fixed $q$

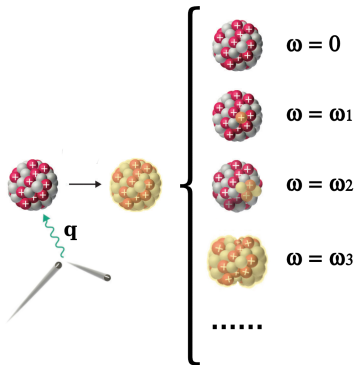
- prepare the target ground state
- right after scattering vertex the target is left in excited state





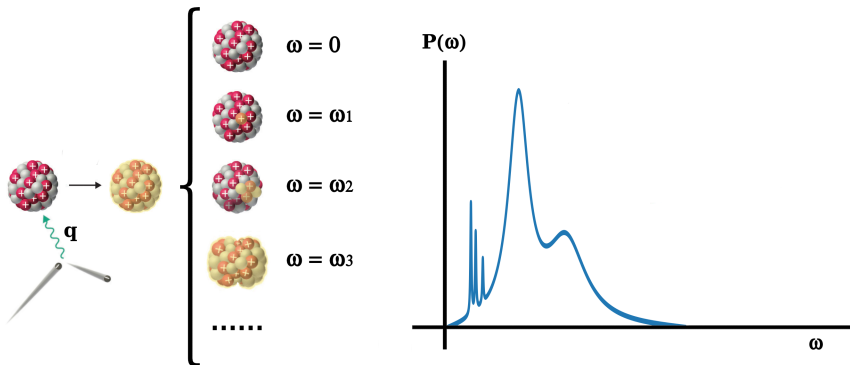
# Idealized algorithm for exclusive processes at fixed $q$

- prepare the target ground state
- right after scattering vertex the target is left in excited state



# Idealized algorithm for exclusive processes at fixed $q$

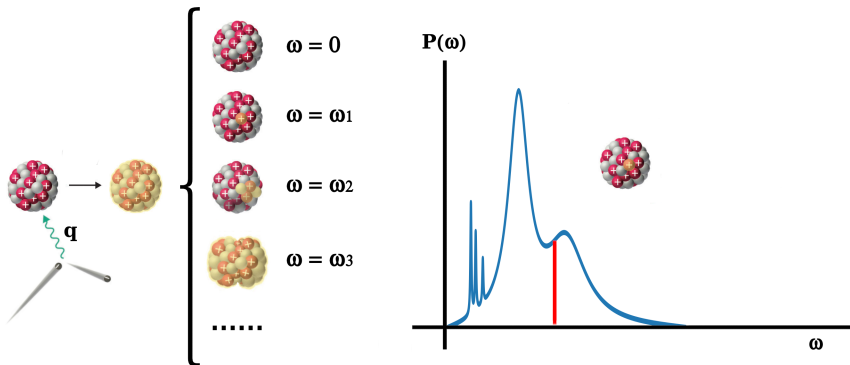
- prepare the target ground state
- right after scattering vertex the target is left in excited state



Roggero & Carlson (2018)

# Idealized algorithm for exclusive processes at fixed $q$

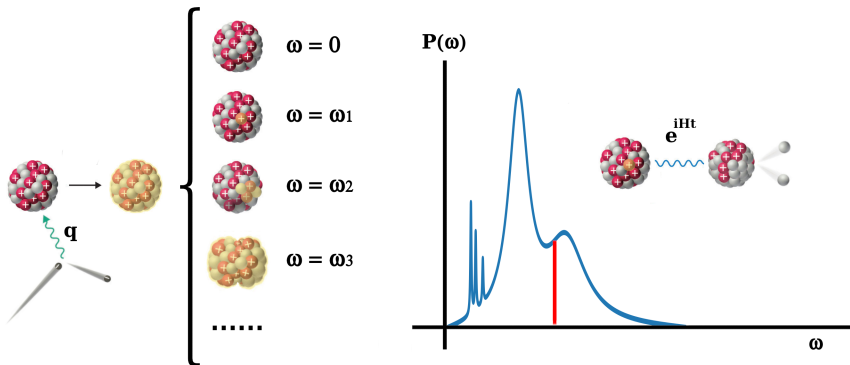
- prepare the target ground state
- right after scattering vertex the target is left in excited state
- energy measurement selects subset of final nuclear states



Roggero & Carlson (2018)

# Idealized algorithm for exclusive processes at fixed $q$

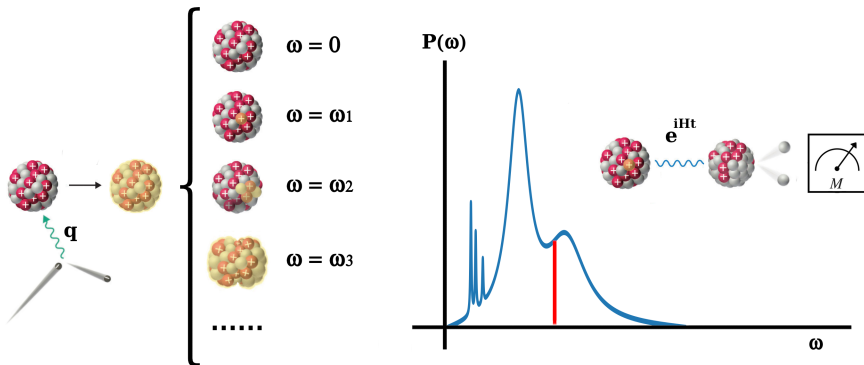
- prepare the target ground state
- right after scattering vertex the target is left in excited state
- energy measurement selects subset of final nuclear states
- further time evolution to let system decay



Roggero & Carlson (2018)

# Idealized algorithm for exclusive processes at fixed $q$

- prepare the target ground state
- right after scattering vertex the target is left in excited state
- energy measurement selects subset of final nuclear states
- further time evolution to let system decay
- measure asymptotic state in detector



Roggero & Carlson (2018)

# The nuclear many-body problem

## A controllable theory for nuclear systems

$$H = \sum_i \frac{p^2}{2m} + \frac{1}{2} \sum_{i,j} V_{ij} + \frac{1}{6} \sum_{i,j,k} W_{ijk} + \dots$$

- much easier to deal with than not the QCD lagragian
- being non-perturbative it is still extremely challenging
  - nuclear states live in huge Hilbert spaces:  $\dim(\mathcal{H}) > 4^A$

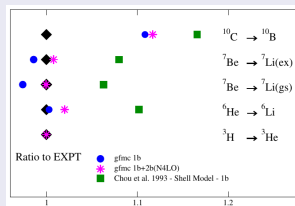
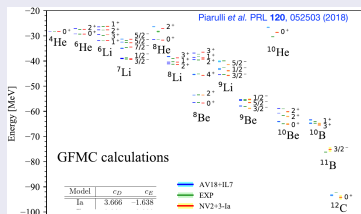
# The nuclear many-body problem

## A controllable theory for nuclear systems

$$H = \sum_i \frac{p^2}{2m} + \frac{1}{2} \sum_{i,j} V_{ij} + \frac{1}{6} \sum_{i,j,k} W_{ijk} + \dots$$

- much easier to deal with than not the QCD lagragian
- being non-perturbative it is still extremely challenging
  - nuclear states live in huge Hilbert spaces:  $\dim(\mathcal{H}) > 4^A$

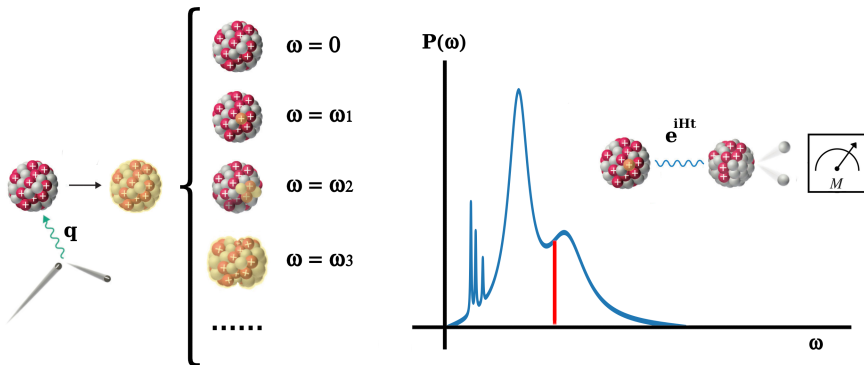
## Great success for light systems with regular (super) computers



Pastore, Baroni et al. (2018)

# Classical algorithm for exclusive processes at fixed $q$

- prepare the target ground state (closed-shell and/or small)
- right after scattering vertex the target is left in excited state
- energy measurement selects subset of final nuclear states (Lovato's talk)
- further time evolution to let system decay
- measure asymptotic state in detector

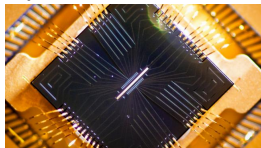


Roggero & Carlson (2018)

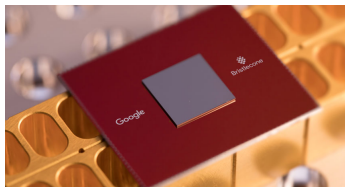


# What is a Quantum computer?

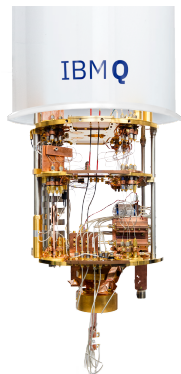
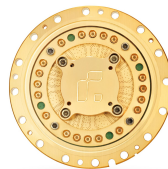
JQI@Univ. of MD



Google

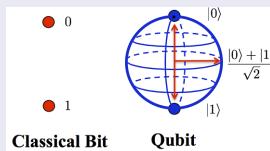


Righetti



- Microsoft?

## Bits vs Qubits



- N bits: an integer number  $< 2^N$
- N qubits: a vector  $|\psi\rangle$  in  $2^N$ -dim Hilbert-space  
 $\Rightarrow$  exponentially more information available

# Quantum Simulations with qubits

*“Nature isn’t classical, dammit, and if you want to make a simulation of nature, you’d better make it quantum mechanical.”*

— R.Feynman (1982)

- in 1996 S.Lloyd shows the conjecture is correct for local interactions

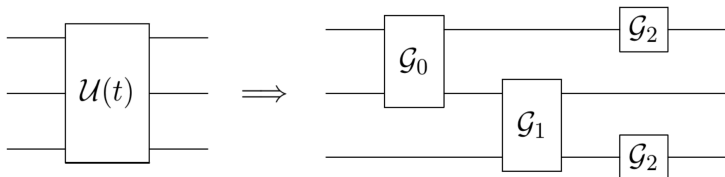
# Quantum Simulations with qubits

*“Nature isn’t classical, dammit, and if you want to make a simulation of nature, you’d better make it quantum mechanical.”*

— R.Feynman (1982)

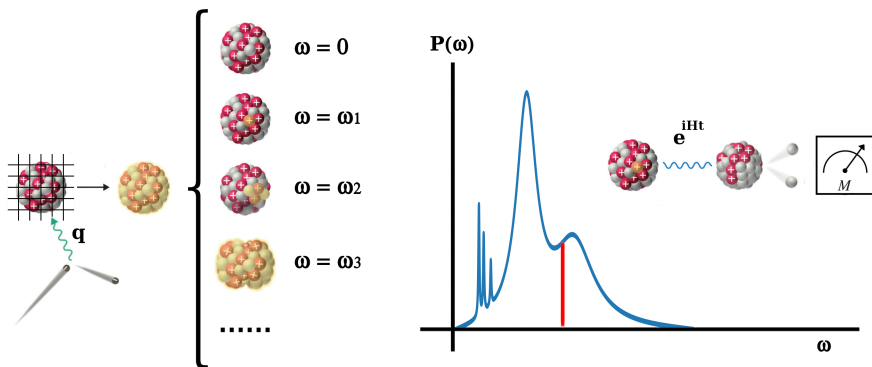
- in 1996 S.Lloyd shows the conjecture is correct for local interactions
- choose a finite basis to discretize system  $\rightarrow \dim(\mathcal{H}) = \Omega \propto e^A$
- physical states can be mapped in states of  $\sim \log_2(\Omega)$  qubits

$$|\Psi(t)\rangle = U(t) |\Psi(0)\rangle$$



# Quantum algorithm for exclusive processes at fixed $q$

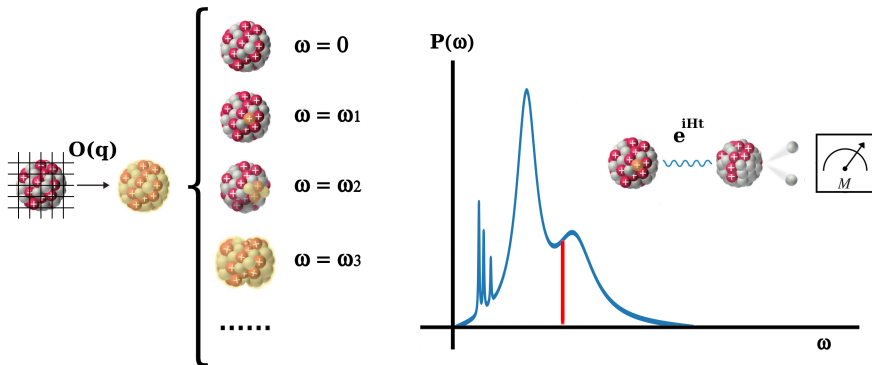
- prepare the target ground state **on a finite qubit basis**
- right after scattering vertex the target is left in excited state
- energy measurements selects subset of final nuclear states
- further time evolution to let system decay
- measure asymptotic state in detector



Roggero & Carlson (2018)

# Quantum algorithm for exclusive processes at fixed $q$

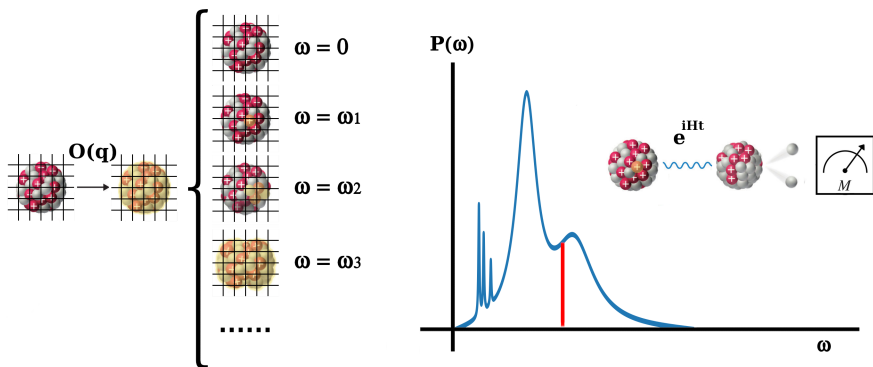
- prepare the target ground state **on a finite qubit basis**
- right after scattering vertex the target is left in excited state
- energy measurements selects subset of final nuclear states
- further time evolution to let system decay
- measure asymptotic state in detector



Roggero & Carlson (2018)

# Quantum algorithm for exclusive processes at fixed $q$

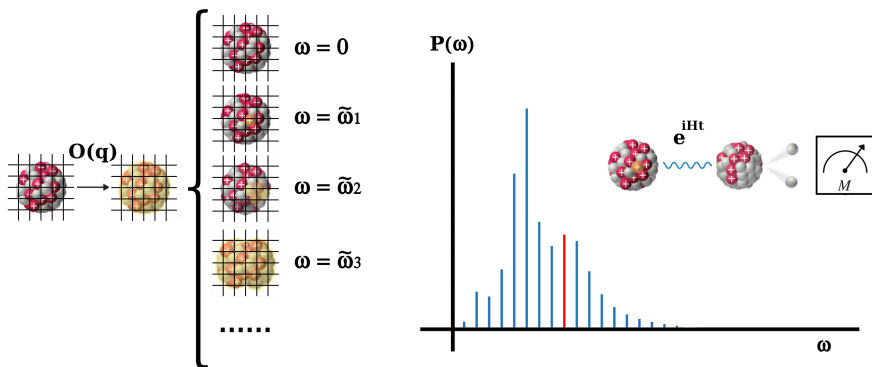
- prepare the target ground state **on a finite qubit basis**
- right after scattering vertex the target is left in excited state
- energy measurements selects subset of final nuclear states
- further time evolution to let system decay
- measure asymptotic state in detector



Roggero & Carlson (2018)

# Quantum algorithm for exclusive processes at fixed $q$

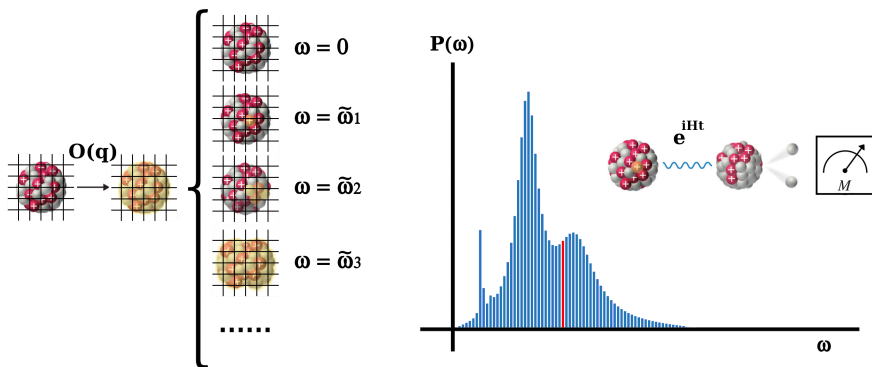
- prepare the target ground state **on a finite qubit basis**
- right after scattering vertex the target is left in excited state
- energy measurements selects subset of final nuclear states (**finite  $\Delta\omega$** )
- further time evolution to let system decay
- measure asymptotic state in detector



Roggero & Carlson (2018)

# Quantum algorithm for exclusive processes at fixed $q$

- prepare the target ground state **on a finite qubit basis**
- right after scattering vertex the target is left in excited state
- energy measurements selects subset of final nuclear states (**finite  $\Delta\omega$** )
- further time evolution to let system decay
- measure asymptotic state in detector

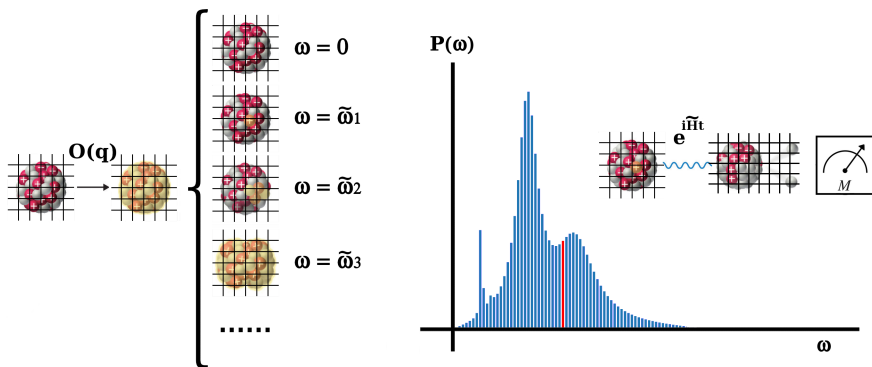


Roggero & Carlson (2018)



# Quantum algorithm for exclusive processes at fixed $q$

- prepare the target ground state **on a finite qubit basis**
- right after scattering vertex the target is left in excited state
- energy measurements selects subset of final nuclear states (**finite  $\Delta\omega$** )
- further **approximate** time evolution to let system decay
- measure asymptotic state in detector



Roggero & Carlson (2018)

## How practical is all this?

- pionless EFT on a  $10^3$  lattice of size 20 fm [ $a = 2.0$  fm]
- 10x faster gates and negligible error correction cost (very optimistic)
- want  $R(q, \omega)$  with 20 MeV energy resolution

## How practical is all this?

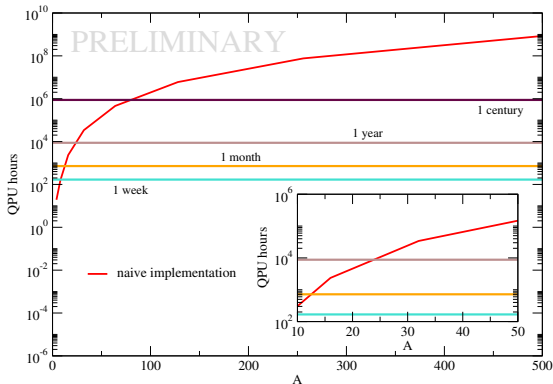
- pionless EFT on a  $10^3$  lattice of size 20 fm [ $a = 2.0$  fm]
- 10x faster gates and negligible error correction cost (very optimistic)
- want  $R(q, \omega)$  with 20 MeV energy resolution

we need a quantum device with  $\approx 4000$  qubits (current record is 72)

# How practical is all this?

- pionless EFT on a  $10^3$  lattice of size 20 fm [ $a = 2.0$  fm]
- 10x faster gates and negligible error correction cost (very optimistic)
- want  $R(q, \omega)$  with 20 MeV energy resolution

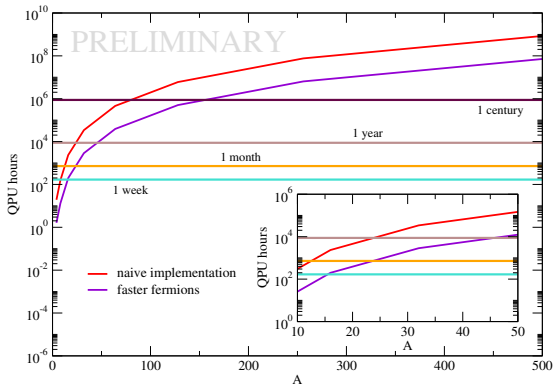
we need a quantum device with  $\approx 4000$  qubits (current record is 72)



# How practical is all this?

- pionless EFT on a  $10^3$  lattice of size 20 fm [ $a = 2.0$  fm]
- 10x faster gates and negligible error correction cost (very optimistic)
- want  $R(q, \omega)$  with 20 MeV energy resolution

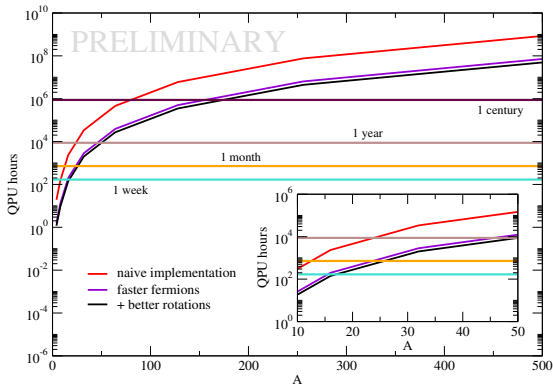
we need a quantum device with  $\approx 4000$  qubits (current record is 72)



# How practical is all this?

- pionless EFT on a  $10^3$  lattice of size 20 fm [ $a = 2.0$  fm]
- 10x faster gates and negligible error correction cost (very optimistic)
- want  $R(q, \omega)$  with 20 MeV energy resolution

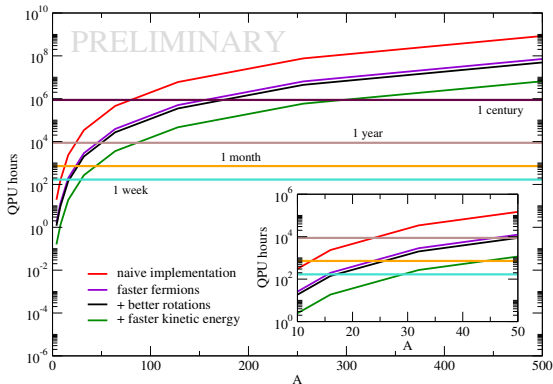
we need a quantum device with  $\approx 4000$  qubits (current record is 72)



# How practical is all this?

- pionless EFT on a  $10^3$  lattice of size 20 fm [ $a = 2.0$  fm]
- 10x faster gates and negligible error correction cost (very optimistic)
- want  $R(q, \omega)$  with 20 MeV energy resolution

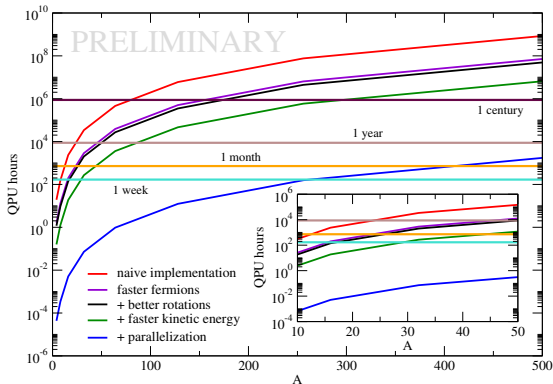
we need a quantum device with  $\approx 4000$  qubits (current record is 72)



# How practical is all this?

- pionless EFT on a  $10^3$  lattice of size 20 fm [ $a = 2.0$  fm]
- 10x faster gates and negligible error correction cost (very optimistic)
- want  $R(q, \omega)$  with 20 MeV energy resolution

we need a quantum device with  $\approx 4000$  qubits (current record is 72)

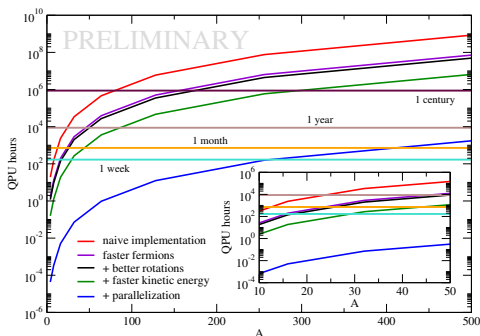




# How practical is all this?

- pionless EFT on a  $10^3$  lattice of size 20 fm [ $a = 2.0$  fm]
- 10x faster gates and negligible error correction cost (very optimistic)
- want  $R(q, \omega)$  with 20 MeV energy resolution

we need a quantum device with  $\approx 4000$  qubits (current record is 72)



coherence time for  $^{40}\text{Ar}$

naive  $\approx 9$  years

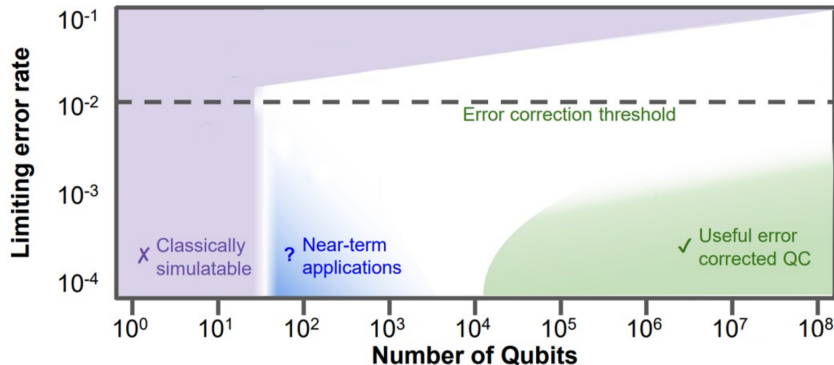
optimized  $\approx 3$  minutes

- algorithm efficiency is critical
- there is still a long way to go
- find new algorithms and/or approximations for near term

# Where are we right now?

figure adapted from Google AI

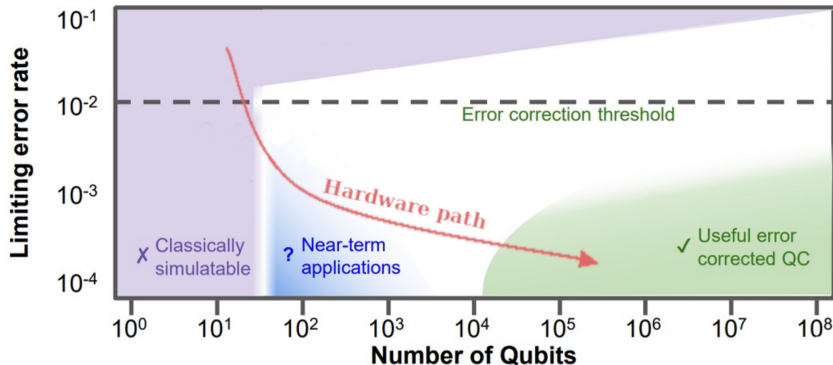
## Need Both Quality and Quantity



# Where are we right now?

figure adapted from Google AI

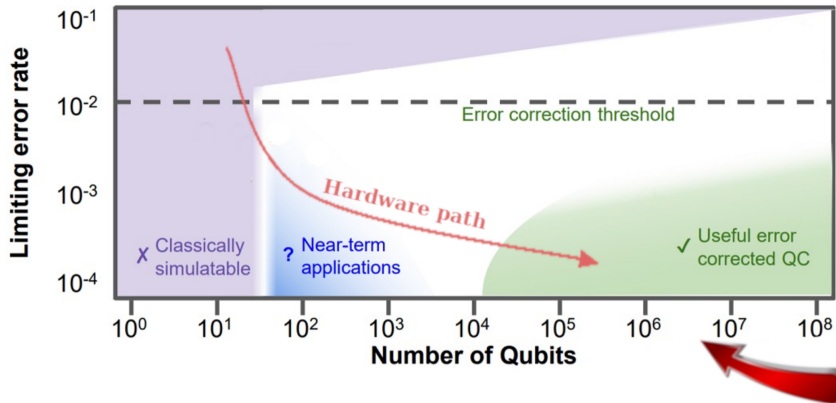
## Need Both Quality and Quantity



# Where are we right now?

figure adapted from Google AI

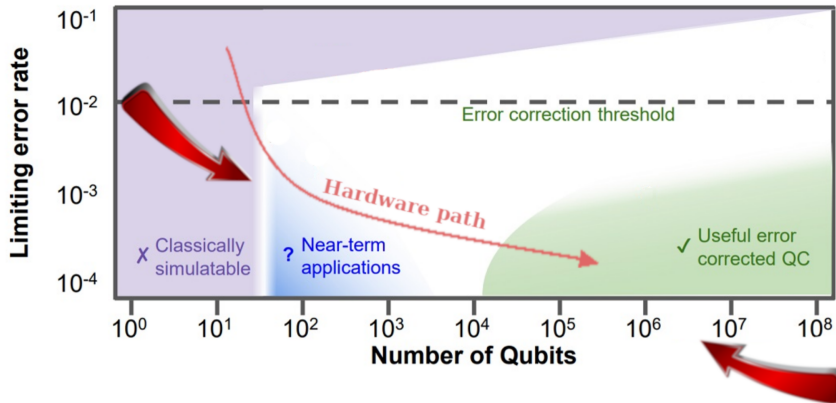
## Need Both Quality and Quantity



# Where are we right now?

figure adapted from Google AI

## Need Both Quality and Quantity



FNAL - INT - LANL effort

A.R. (INT), J. Carlson & R. Gupta (LANL), G. Perdue, A. Li & A. Macridin (FNAL)

# Summary

- understanding low-energy dynamics of nuclear many-body systems is important for current and planned neutrino oscillation experiments
- remarkable progress for inclusive x-sec of light nuclei ([See Lovato's talk](#))
  - still not enough for exclusive scattering off  $^{40}\text{Ar}$ , need new ideas: short time approximation([See Carlson's talk](#))? quantum computing?
- QC is an emerging technology with the potential of revolutionize the way theory calculations are done
- we already know how to simulate efficiently the time-evolution of non relativistic systems and how to study exclusive scattering
- more work has to be done to make all this viable in the near term

# Summary

- understanding low-energy dynamics of nuclear many-body systems is important for current and planned neutrino oscillation experiments
- remarkable progress for inclusive x-sec of light nuclei ([See Lovato's talk](#))
  - still not enough for exclusive scattering off  $^{40}\text{Ar}$ , need new ideas: short time approximation([See Carlson's talk](#))? quantum computing?
- QC is an emerging technology with the potential of revolutionize the way theory calculations are done
- we already know how to simulate efficiently the time-evolution of non relativistic systems and how to study exclusive scattering
- more work has to be done to make all this viable in the near term

## Collaborators:

- J. Carlson (LANL)

