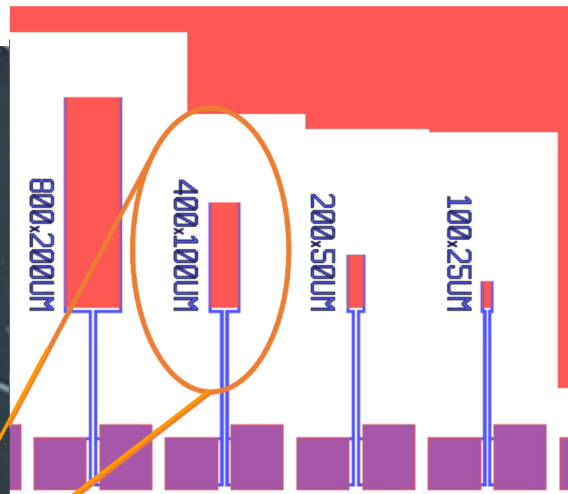
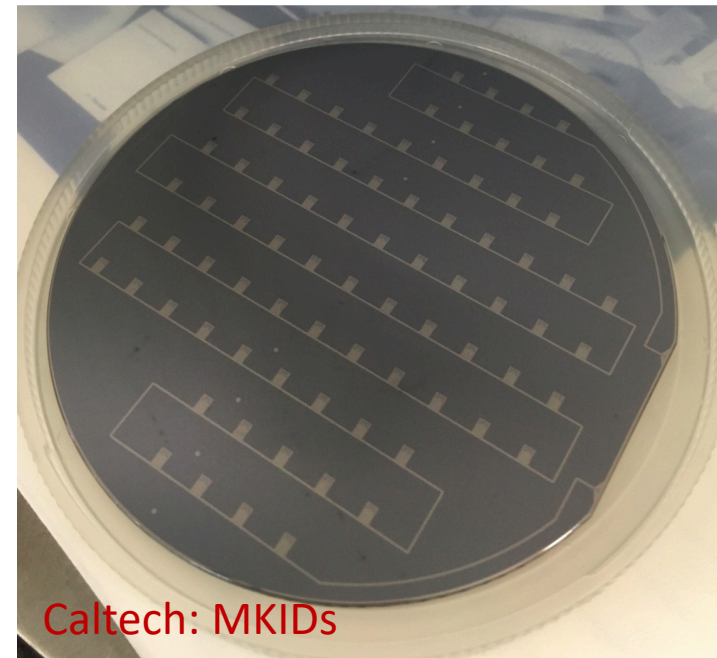


NTD, TES, MKID and SNSPD Sensors in Dark Matter



Matt Pyle

UC Berkeley

New Directions Fermilab

June 4 2019

Outline

- Light Mass Dark Matter Design Drivers
- TES
 - Athermal phonon collection and concentration schemes
- NTDs
- MKIDs
- SNSPDs
 - Energy Reconstruction
- Proposed/Running Experiments

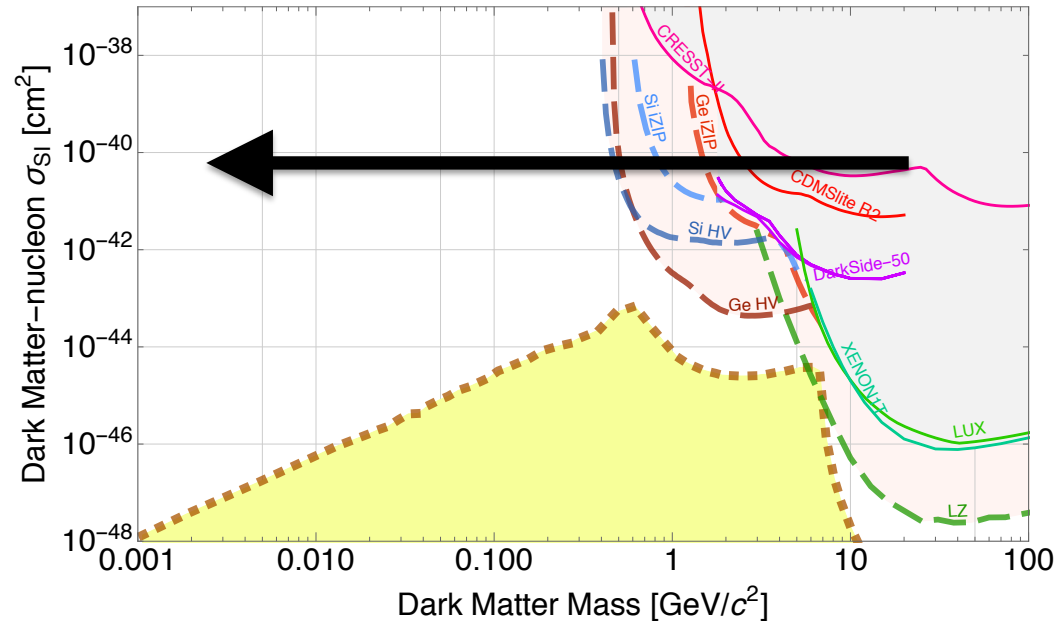
Light Mass Dark Matter Direct Detection Design Drivers

Light Mass DM Design Drivers: ~~Exposure~~

$$R = \sigma n_{DM} v_{DM} N_{exp}$$

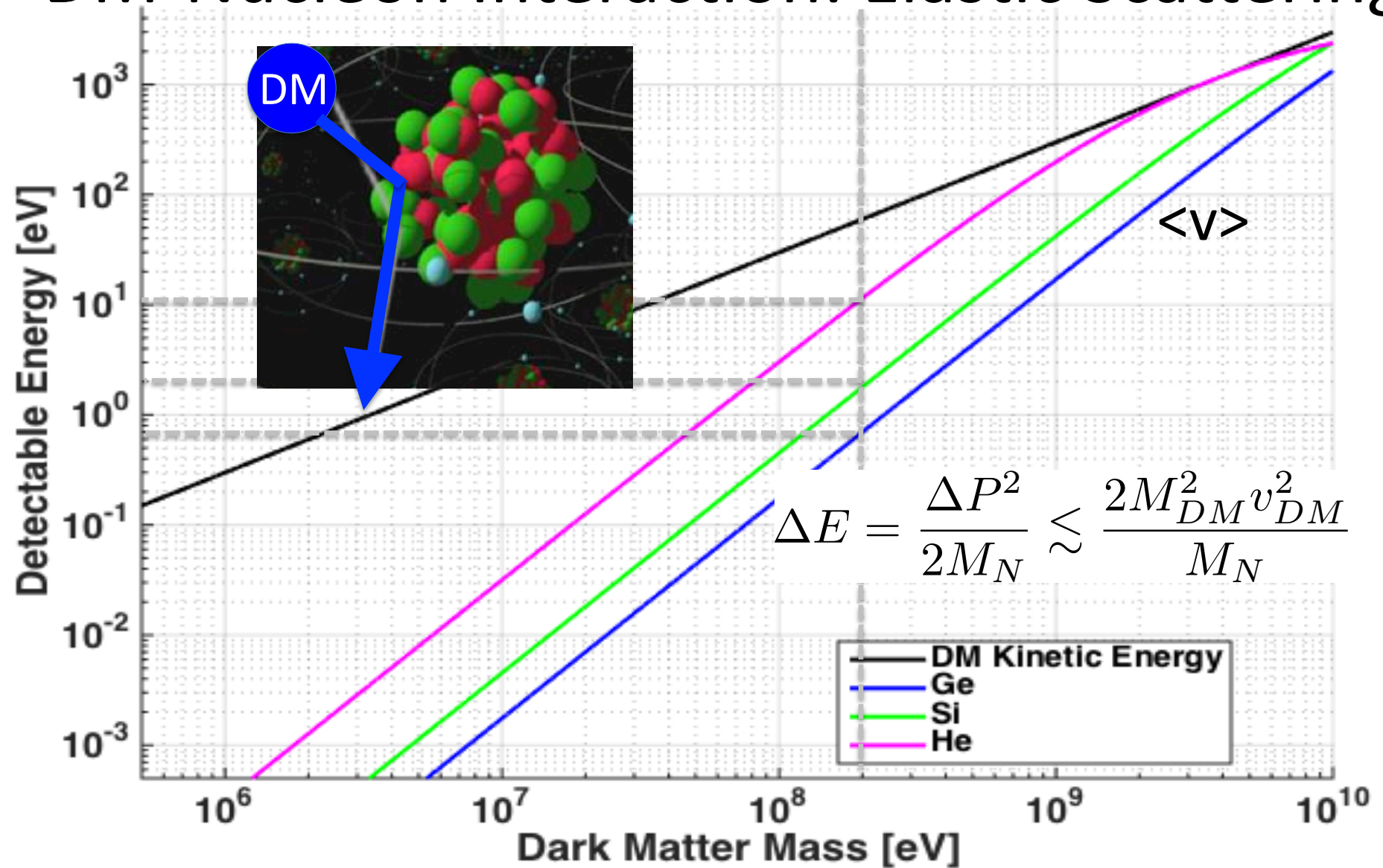
$$= \sigma \frac{\rho_{DM}}{M_{DM}} v_{esc} N_{exp}$$

Interaction
Rate scales
with $1/M_{DM}$



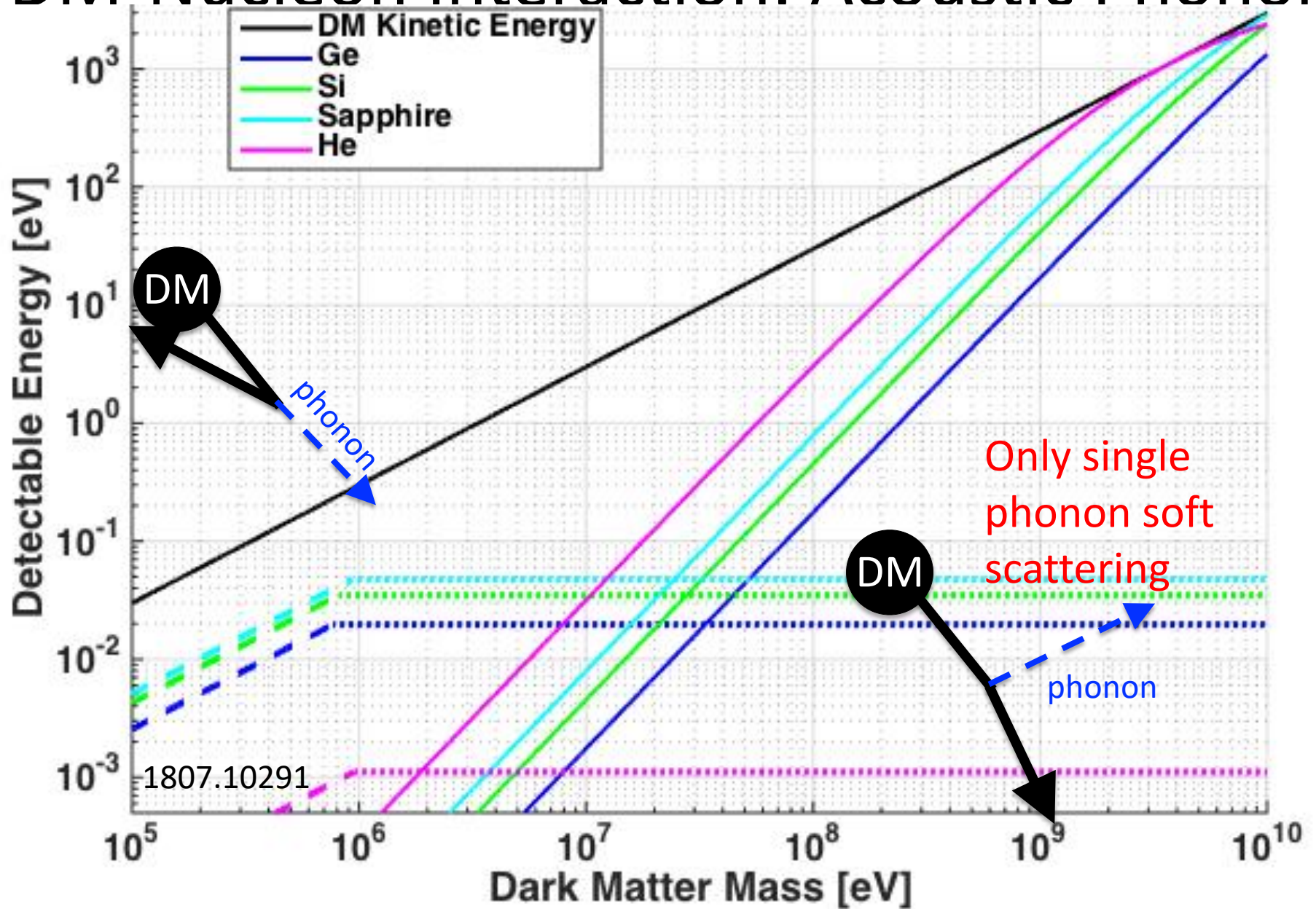
LZ needs 10 tons to get to 10^{-47} cm^2 at 100GeV, Light Mass DM searches only needs 10kg to reach the same level at 100MeV.

DM-Nucleon Interaction: Elastic Scattering



Primary Design Driver: Vibrational Energy Sensitivity

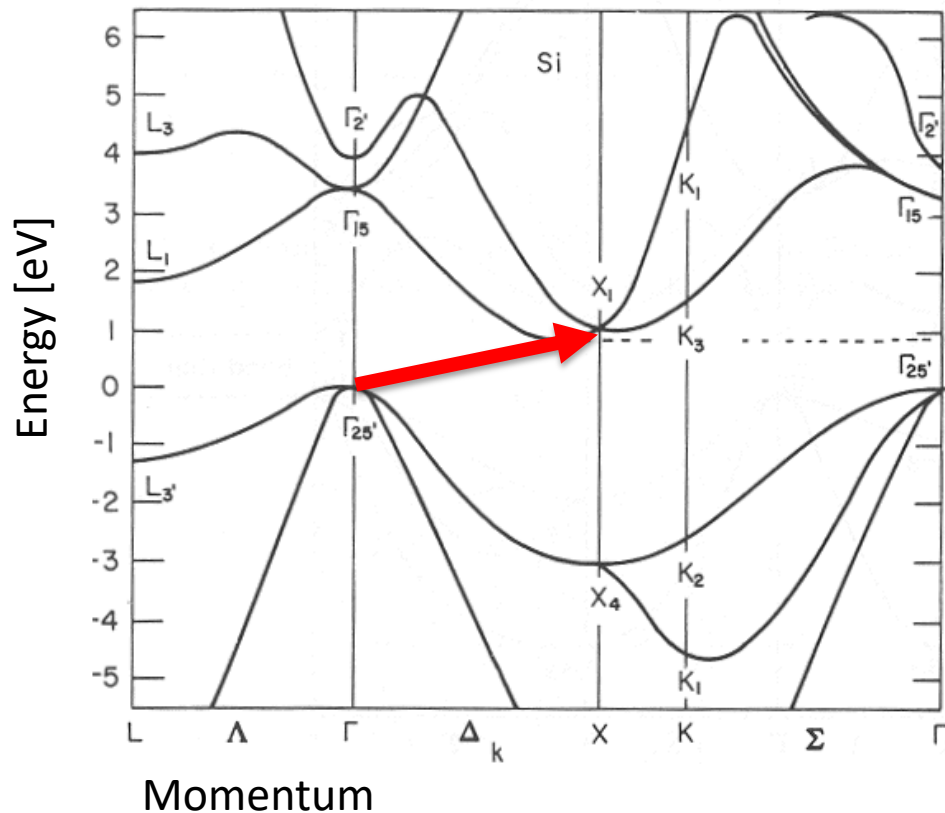
DM-Nucleon Interaction: Acoustic Phonon



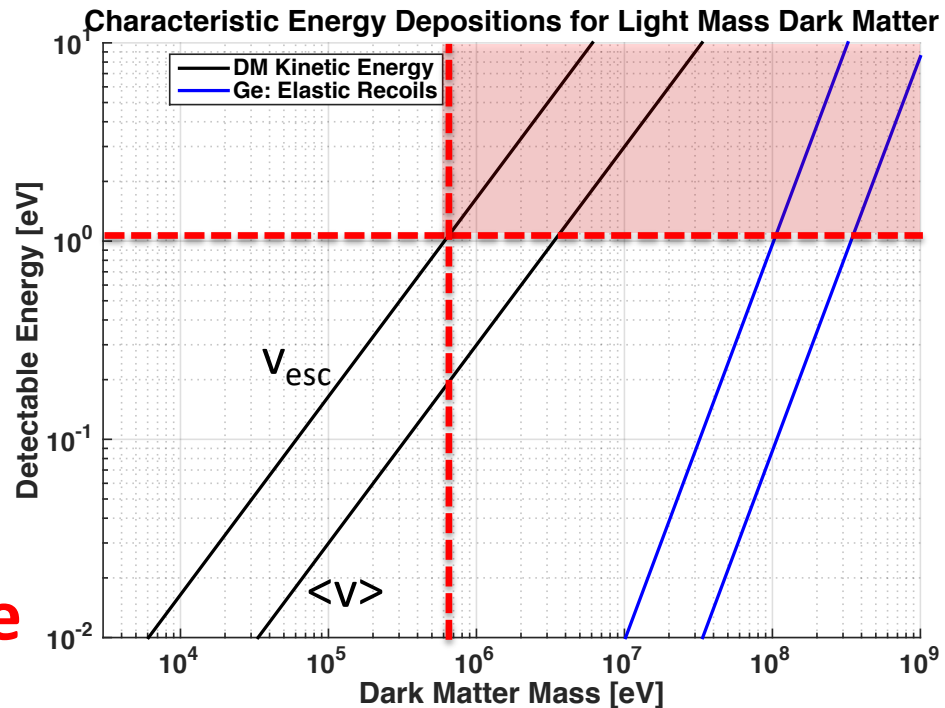
Primary design driver: Still Vibrational Energy Sensitivity

DM-Electronic Interactions: Inelastic e^- recoil

E [eV] Band Diagram for Si

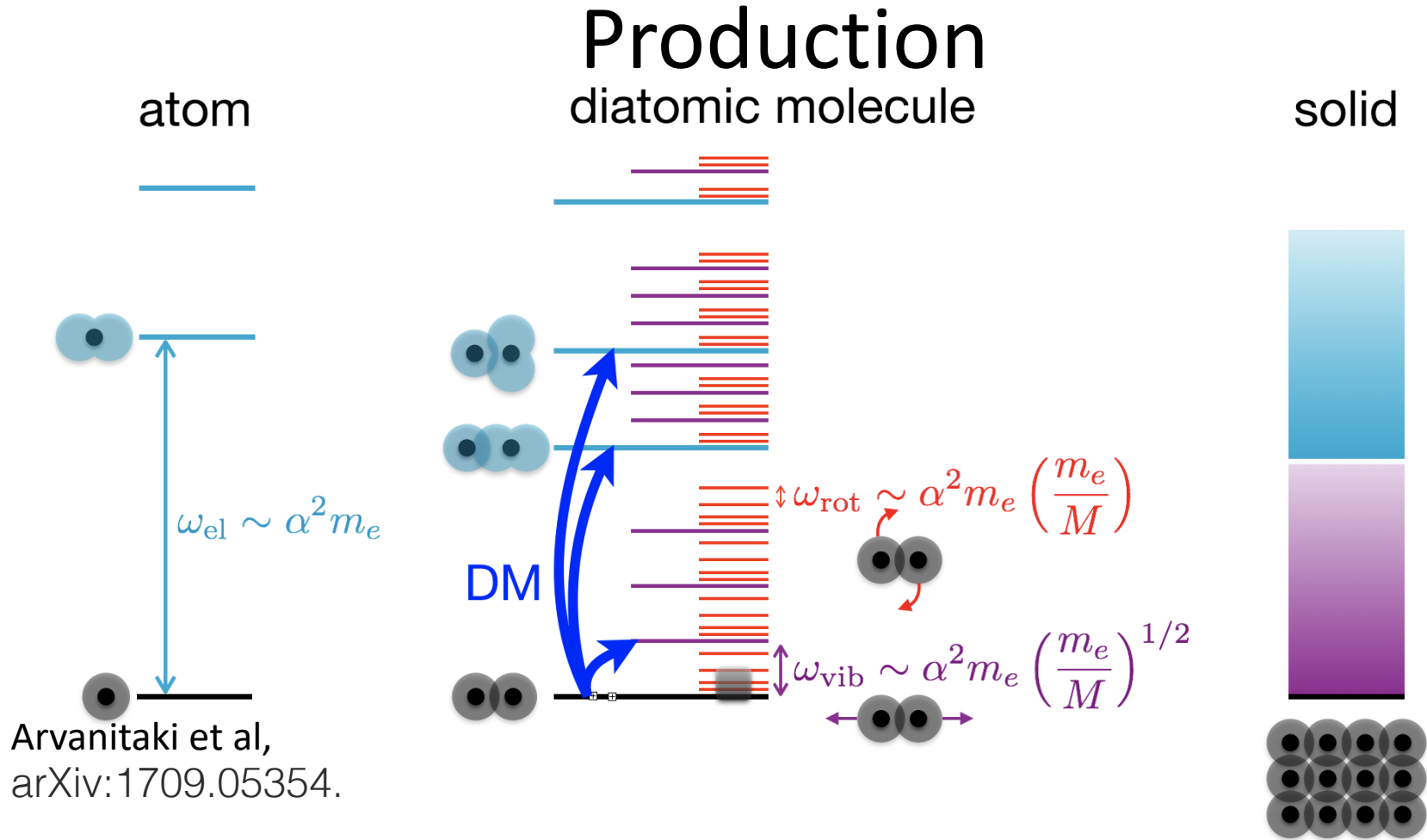


- e^- excitation momentum and energy scales in semiconductors well matched to 1 MeV-100MeV DM
- Essig et al: 1108.5383



Design Driver: Sensitivity to single e/h pairs (eV energy) with negligible dark count rate

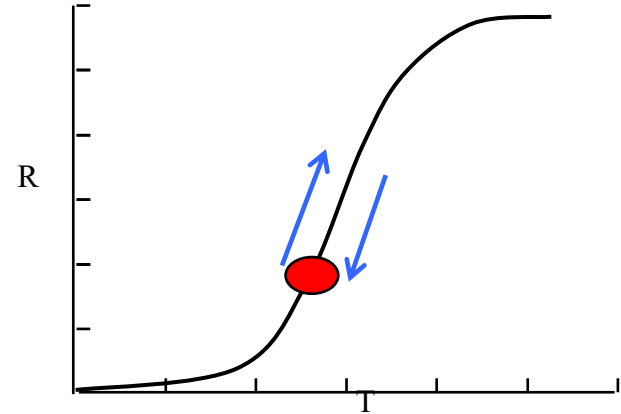
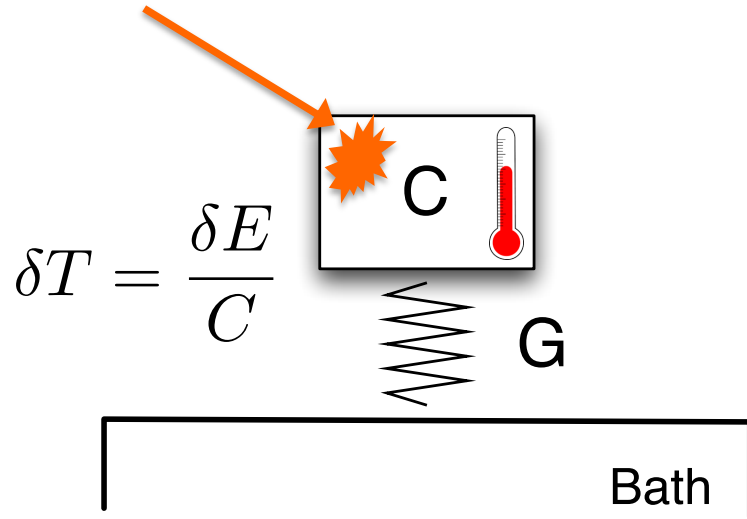
DM-Molecular/Atom Absorption: Photon



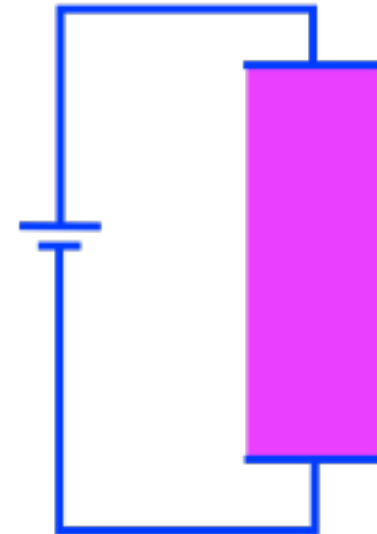
**Design Driver: 100meV –eV scale
single photon sensitivity**

Transition Edge Sensor

The Simplest Thermal Calorimeter

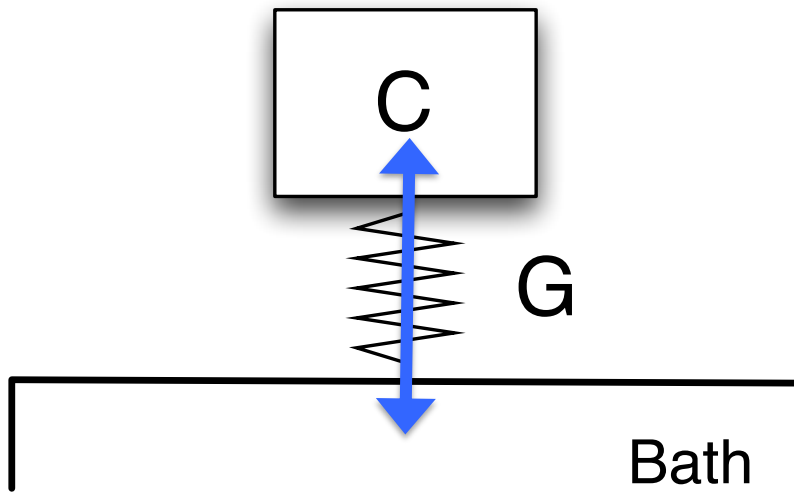


Transition Edge Sensor (TES): A superconducting metal film (W) that is externally biased so as to be within its superconducting/normal transition



TES Sensitivity

$$\begin{aligned}\sigma_{\langle E \rangle}^2 &= \sum_i (E_i - \langle E \rangle)^2 \frac{e^{-\beta E_i}}{\sum_j e^{-\beta E_j}} \\ &= \frac{\sum_i E_i^2 e^{-\beta E_i}}{\sum_j e^{-\beta E_j}} - \langle E \rangle^2 \\ &= -\frac{\partial \langle E \rangle}{\partial \beta} = \frac{\partial \langle E \rangle}{\partial T} k_b T^2 = C k_b T^2\end{aligned}$$

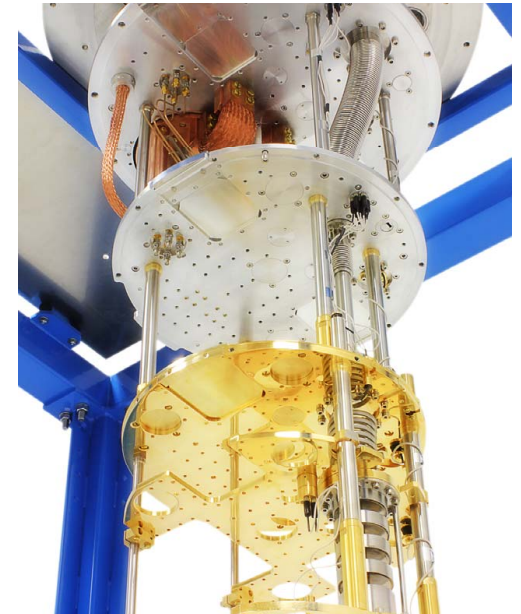


~ Intrinsic Thermal Noise
of Calorimeters

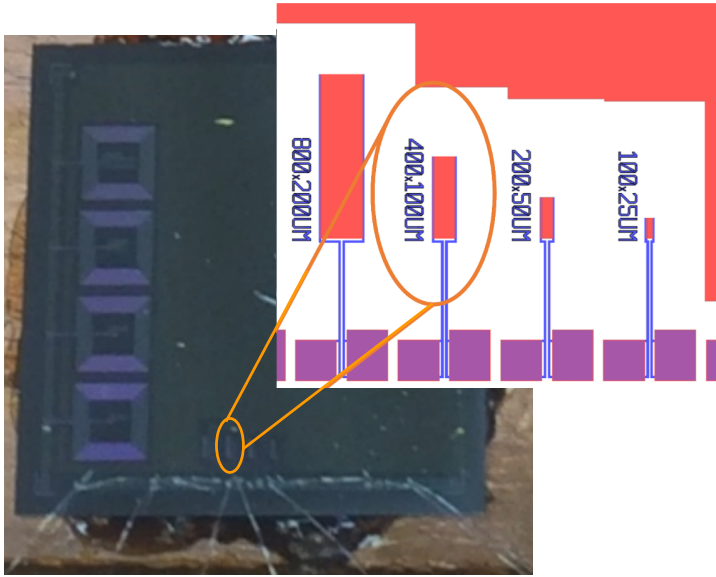
TES Scaling Laws & Optimization

$$\sigma_{\langle E \rangle}^2 = C k_b T_c^2$$
$$\propto V T_c^3$$

- Minimize Temperature
 - Dilution Refrigerators can cool detectors to 5mK
 - $T_c \rightarrow 1.5 \times T_{\text{bath}}$
- Minimize Volume



100meV Threshold TES are really small!



$$\sigma_{\text{theory}} \sim 20\text{meV}$$

$$T_c = 41\text{mK}$$

$$V = 400\mu\text{m} \times 100\mu\text{m} \times 40\text{nm}$$

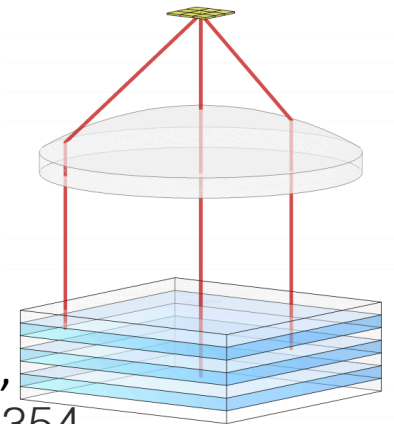
$$= 1.6 \times 10^{-9} \text{ cm}^3$$

$$= 30\text{ng}$$

1 gram of active target mass \rightarrow 30M channels

1 cm² of active area \rightarrow 2500 channels

TES can only be used by themselves when photon collecting over a very small area

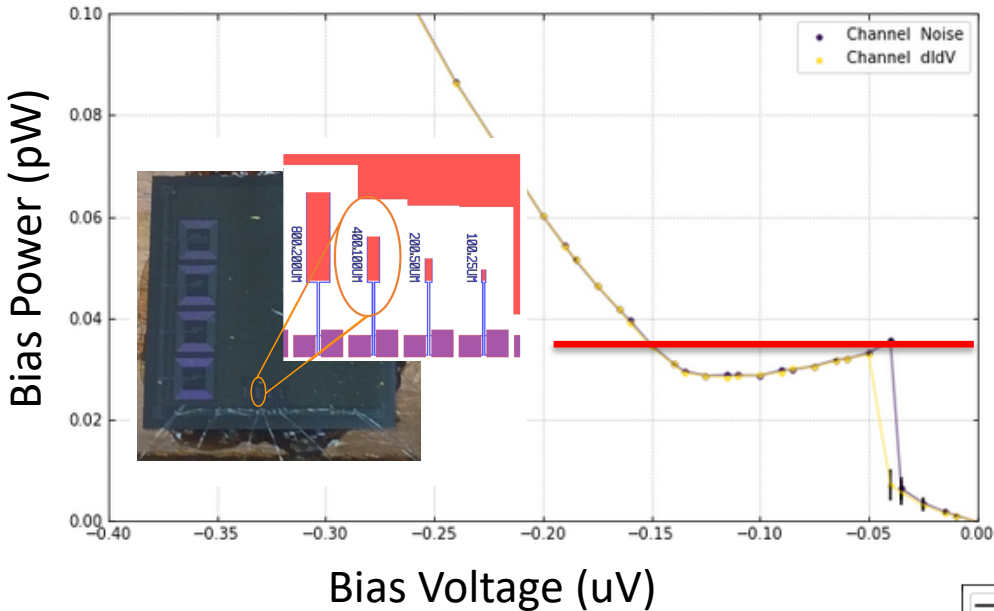


Arvanitaki et al,
arXiv:1709.05354.

Currently Achieved TES Energy Sensitivity:

$100\mu\text{m} \times 400\mu\text{m} \times 40\text{nm}$

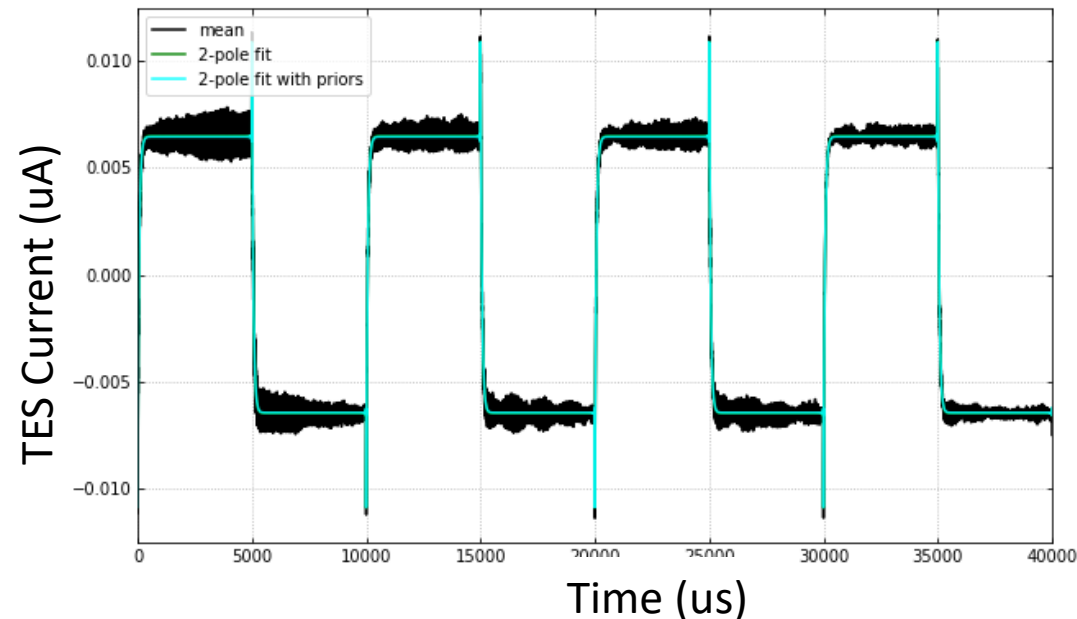
TES Power vs TES Bias Voltage



- T. Aramaki, P. Brink, C. Fink, R. Harris, Y. Kolomensky, R. Mahapatra, N. Mirabolfathi, R. Partridge, M. Platt, MP, B. Sadoulet, B. Serfass, S. Watkins
- $100\mu\text{m} \times 400\mu\text{m} \times 40\text{nm}$ W TES
- $T_c = 41\text{mK}$
- Normal Resistance: $630\text{m}\Omega$
- Bias Power: 35fW

TES Response to Square Wave Jitter

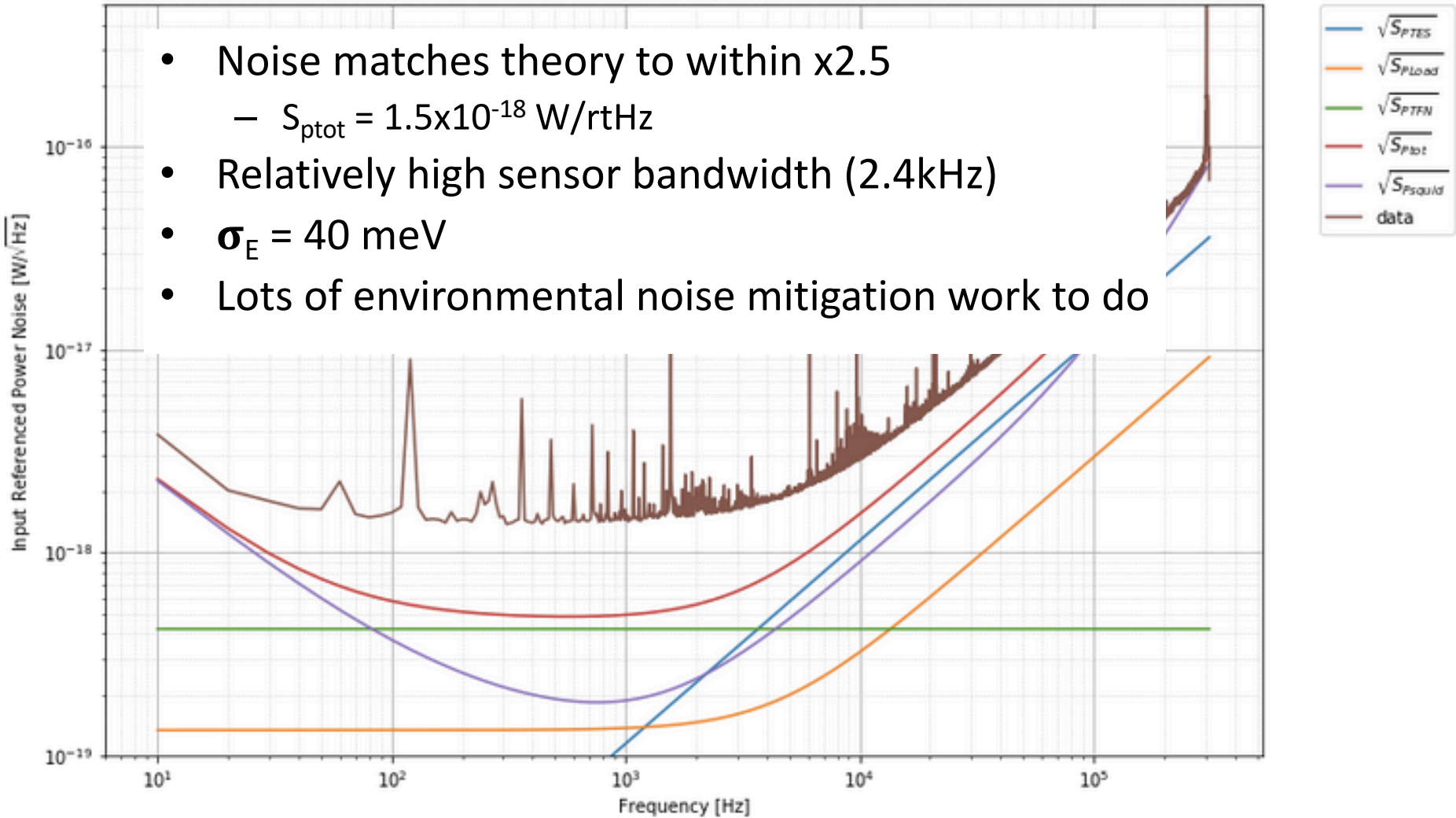
- Complex Impedance
 - Simple 2 pole TES dynamical model perfectly fits response
- TES falltime: $\sim 66\mu\text{s}$ (2.4kHz)
 - Relatively fast
 - Long term $\rightarrow 1\text{ms}$ and allow athermal phonons to ballistically bounce



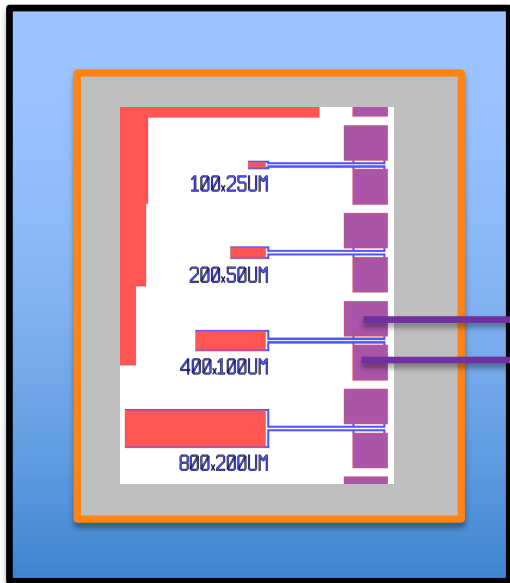
100 μ m x400 μ m TES Noise

Power Noise For $R_0 : 47.85 \text{ m}\Omega$

- Noise matches theory to within x2.5
 - $S_{\text{ptot}} = 1.5 \times 10^{-18} \text{ W/rHz}$
- Relatively high sensor bandwidth (2.4kHz)
- $\sigma_E = 40 \text{ meV}$
- Lots of environmental noise mitigation work to do

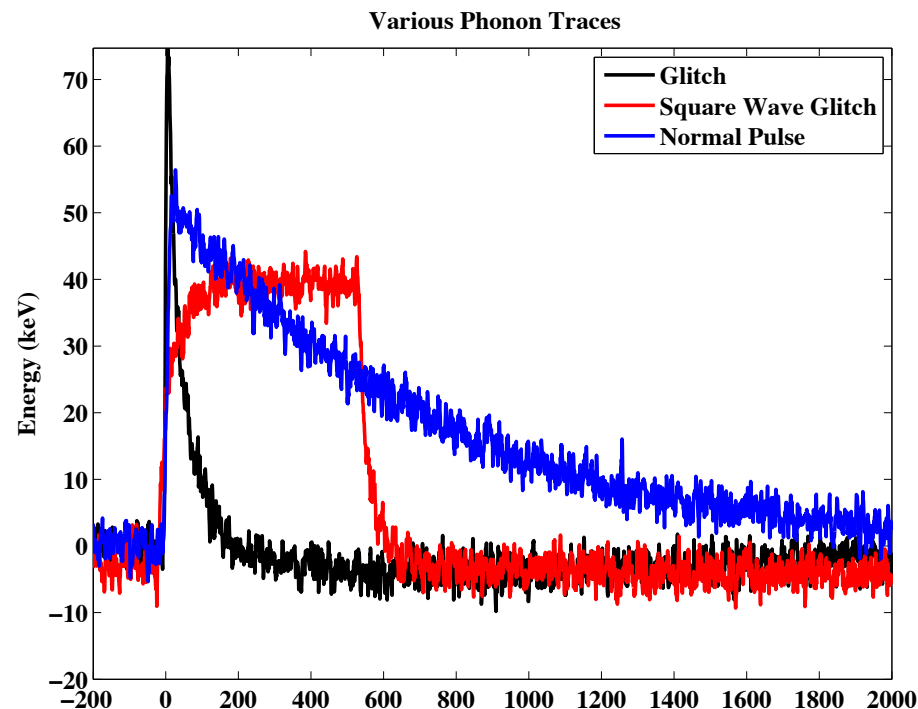


TES: Environmental Noise Susceptibility



- TES is a resistor ... you can heat a resistor with an E&M wave of any frequency
- 5fW of DC Environmental EMI coming down the TES bias lines
- Lots of AC power glitches seen too

Big Challenge: Need to continue to improve Environmental Isolation

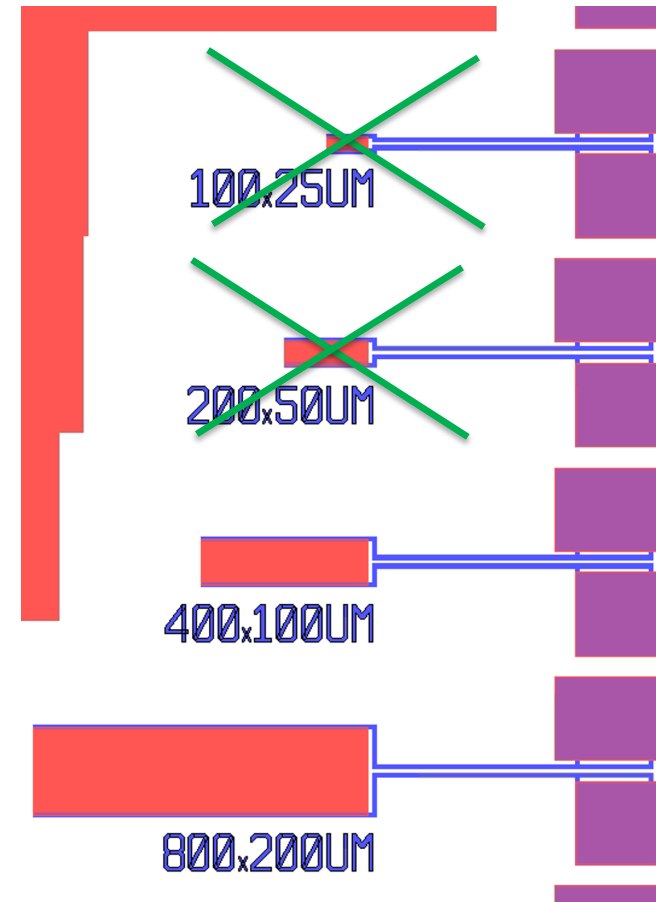


TES R&D

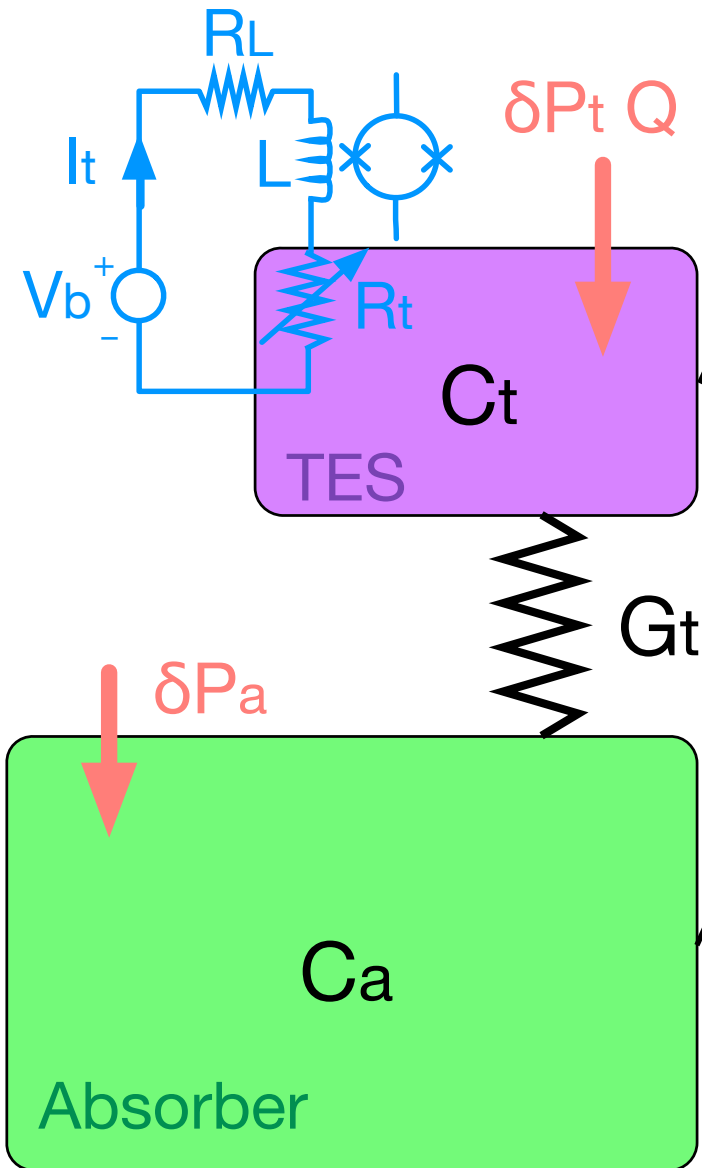
Energy Sensitivity: 40meV -> 1meV

R&D Work Plan

- Lower T_c from 40mK -> 10mK.
 - x8 sensitivity improvement
- Lower volume by x16
 - x4 sensitivity improvement
- Decrease environmental noise by 50dB ... there is a reason I'm not showing the performance of the 200umx50um TES

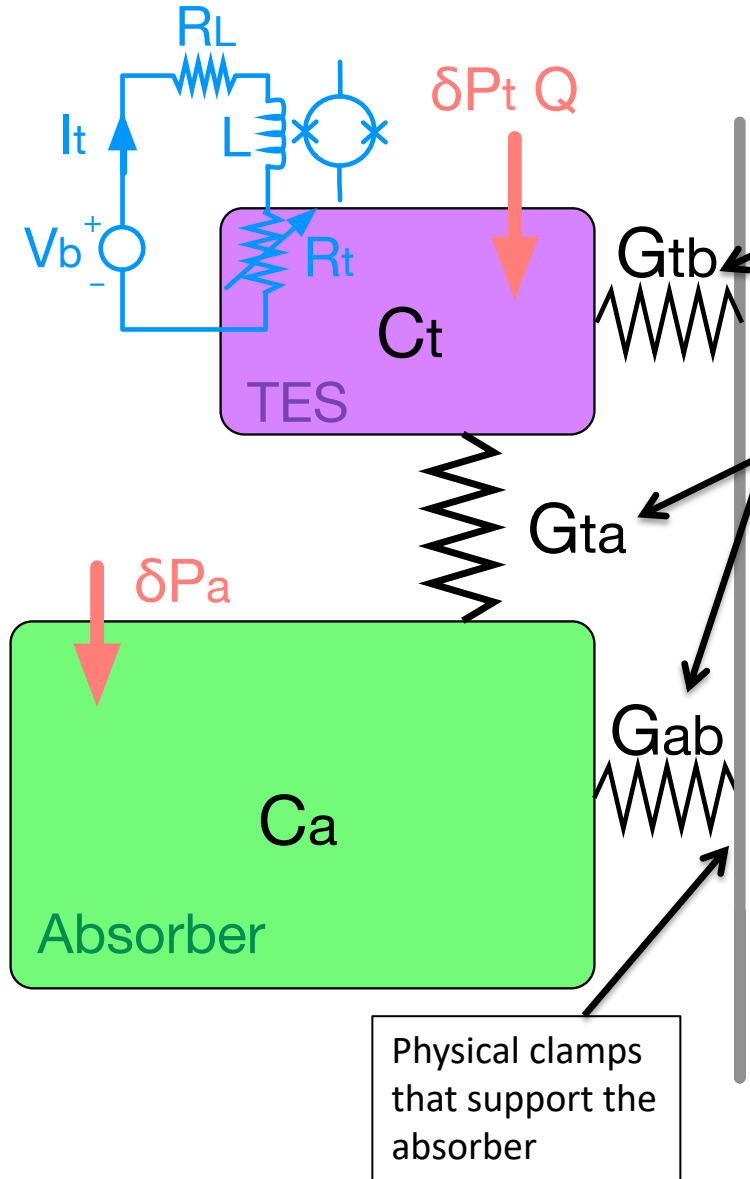


Increase Target Volume: Add Insulating Absorber



- Couple the sensor to a large volume insulator -> low heat capacity.
- Vibrational Energy in the Absorber is transported through G_t to the electronic system of the TES

Problem: Decoupling between the Sensor and Absorber



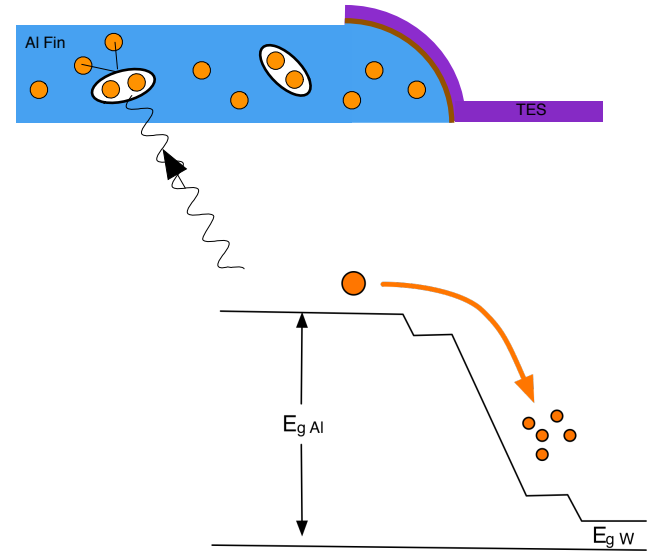
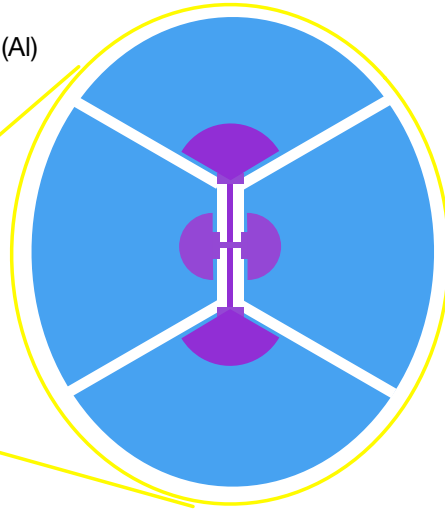
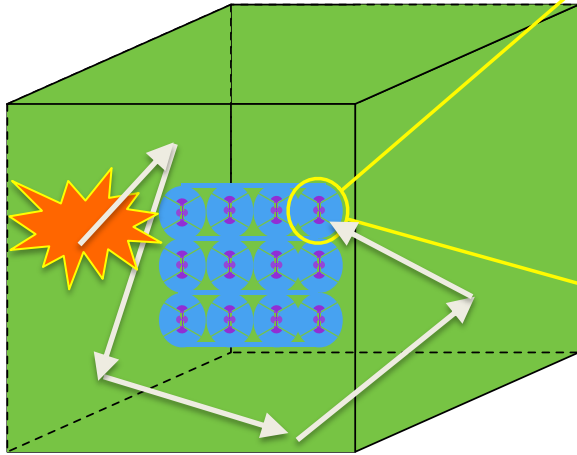
- Kapitza boundary conductance scale as T^3
- e-/phonon thermal conductance scales as T^4

As T is decreased, it's harder and harder to keep the sensor thermally coupled to the absorber

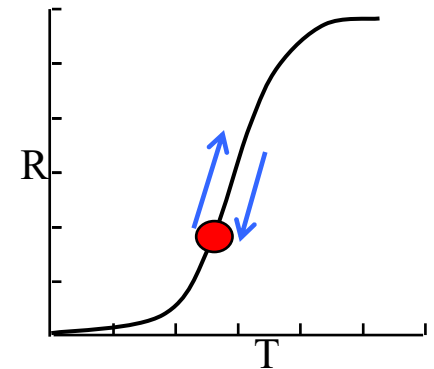
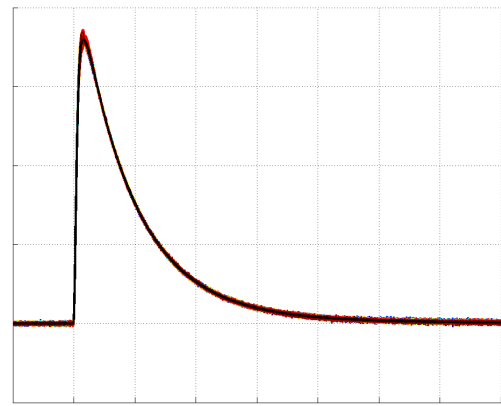
- Energy leaks out of the absorber through G_{ab} before its measured
- TES sensitive to power fluctuations through G_{tb}
- EDELWEISS has done this successfully

Athermal Phonon Sensor Technology

- TES and QP collection antennas (W)
- Athermal Phonon Collection Fins (Al)
- 1cm³Polar Crystal



Collect and Concentrate
Athermal Phonons (>4K) into
Sensor before they thermalize



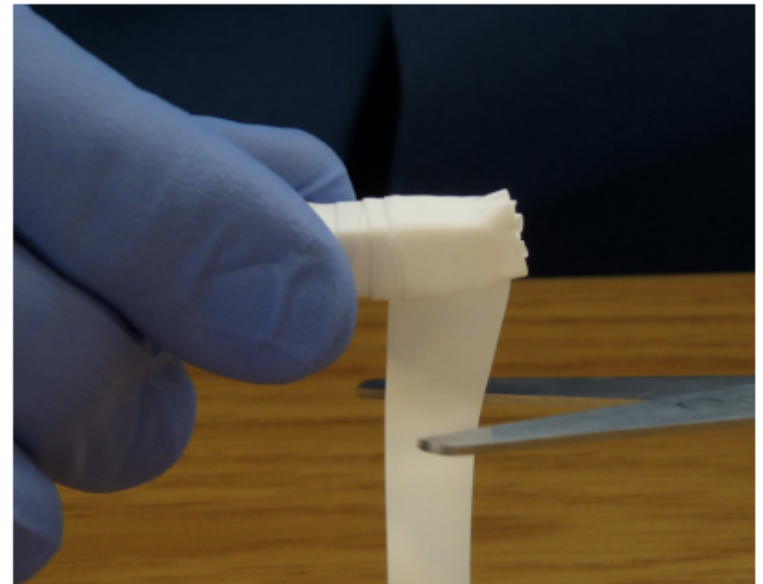
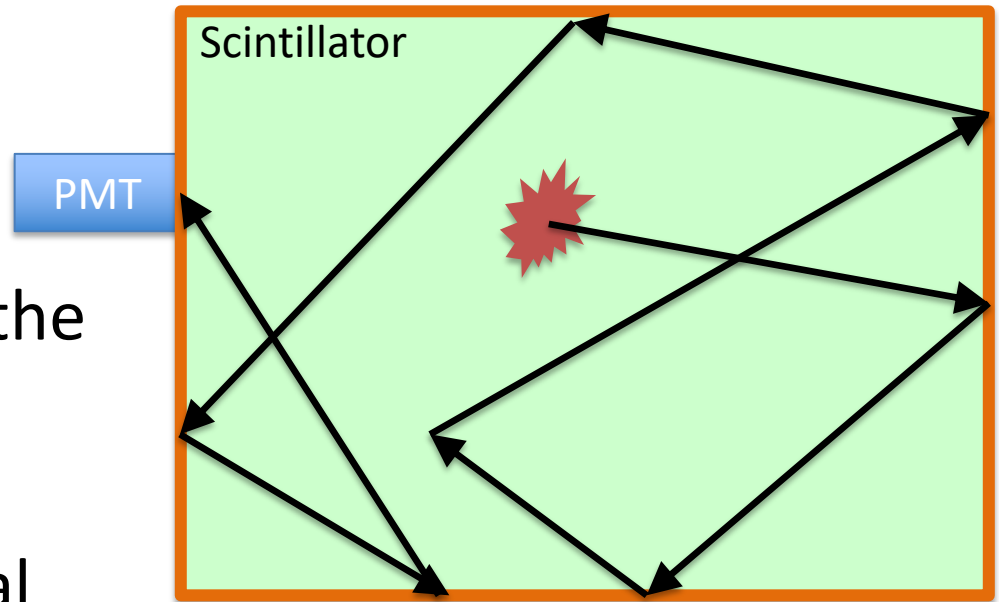
Excitation Detectors & Volume Scaling



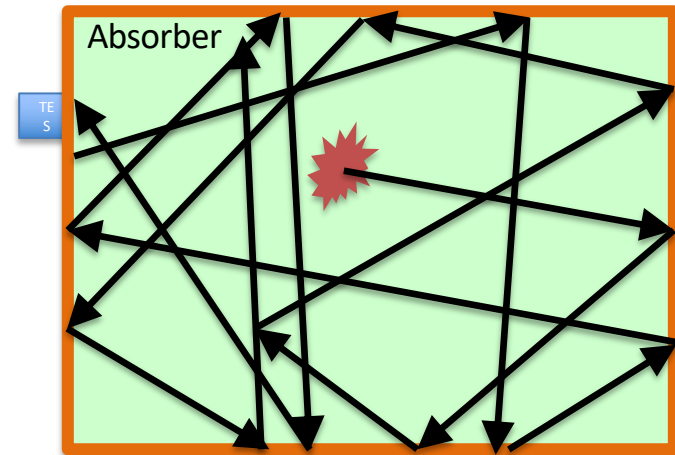
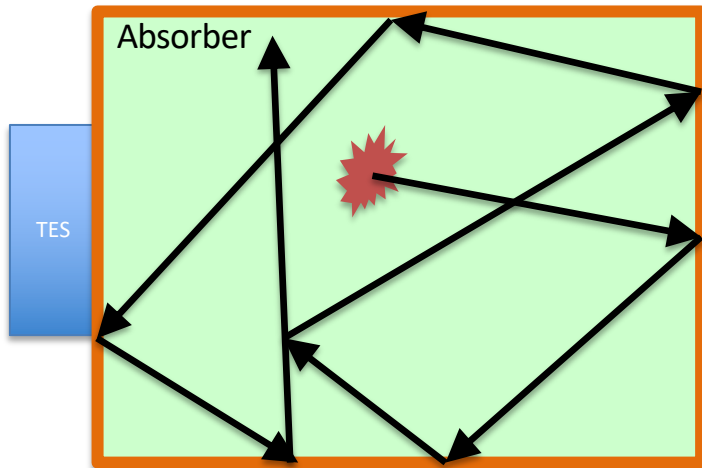
Will these detectors have the same energy sensitivity?

Yes, if:

- Lifetime of the athermal excitation (photon) is really long
- Excitation absorption dominated by sensor
- ~~Position Sensitivity~~

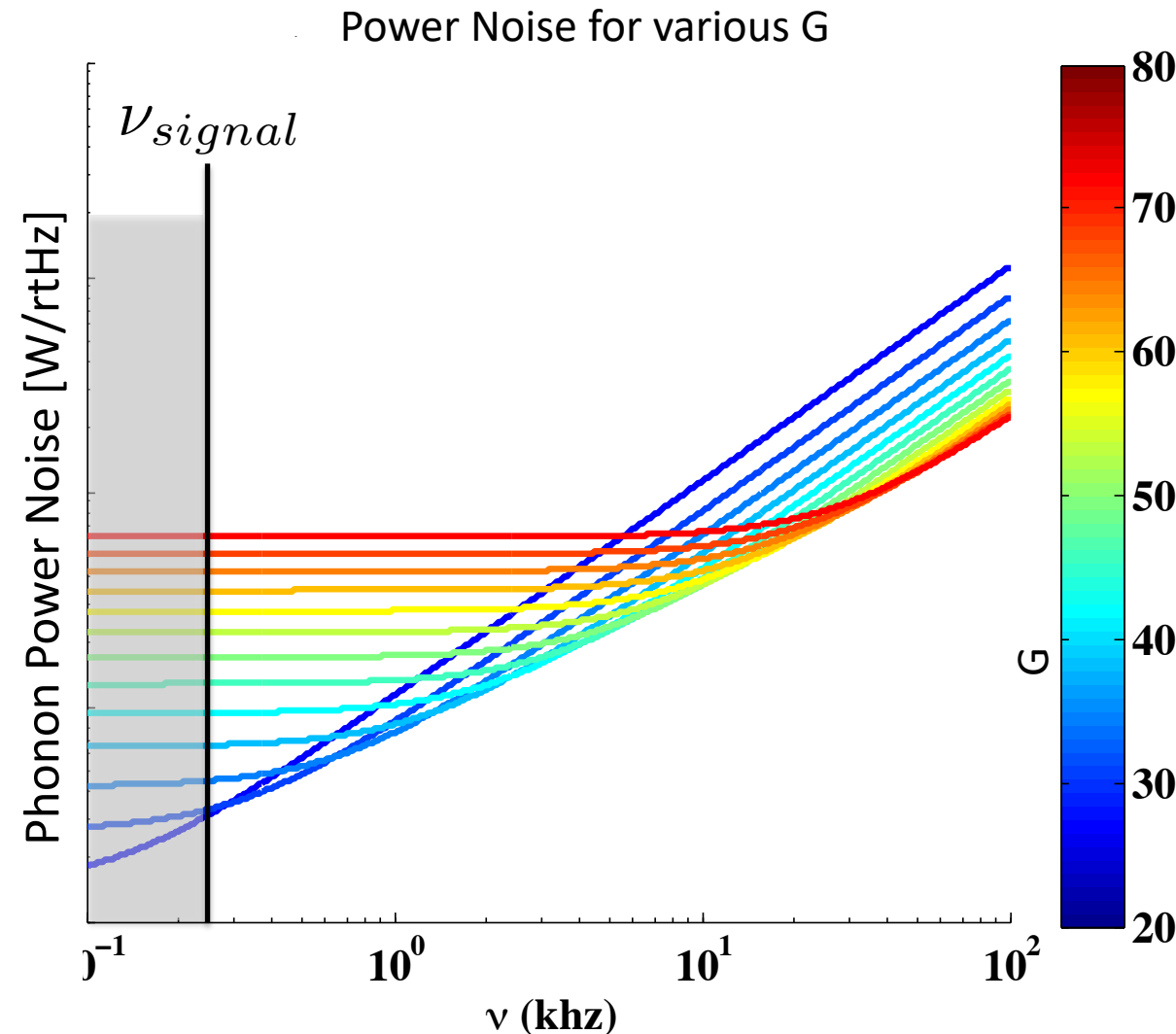


Optimizing the Athermal Phonon Excitation Detectors



Minimize the number/volume of the TES sensors instrumented on the surface to the point that you begin to see the bare surface thermalization rate

Athermal Phonon Sensor Sensitivity Scaling



$$G \propto T_c^4$$

$$S_{ptfn} = 4k_b T_c^2 G$$

$$\propto T_c^6$$

$$\sigma_E \propto T_c^3$$

- Lower ν_{sensor} (lower T_c) if $\nu_{signal} < \nu_{sensor}$
- Lower ν_{signal} (decrease Al coverage) if $\nu_{signal} > \nu_{sensor}$

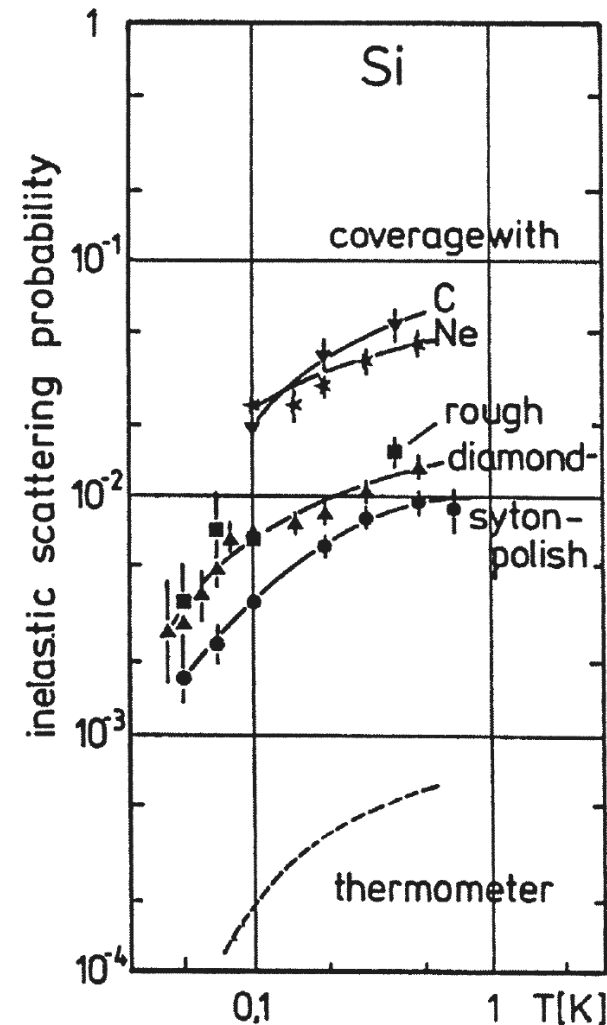
You can always say on νT_c^3 scaling (in principle)
 45mK \rightarrow 10mK: 2eV \rightarrow 20meV

Athermal Phonon Thermalization at Surfaces

- Athermal phonon surface thermalization probability found to depend upon
 - Crystal
 - Surface roughness
 - Surface cleanliness

(W. Knaak et al, Phonon Scattering in Condensed Matter V, 1986)

- 0.1%-1% of the crystal surface covered with athermal phonon sensors ... 1/1000-1/100 thermalization probability needed
- Si, Ge -> ok



Large Area Photon Detectors

Large Area Photon Detector



CRESST 2 Light Detector

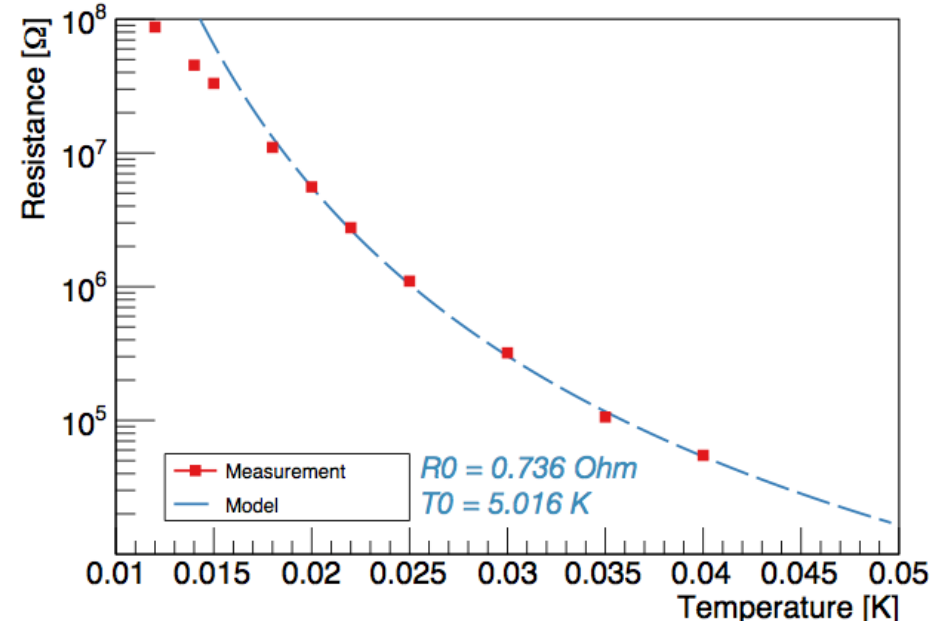
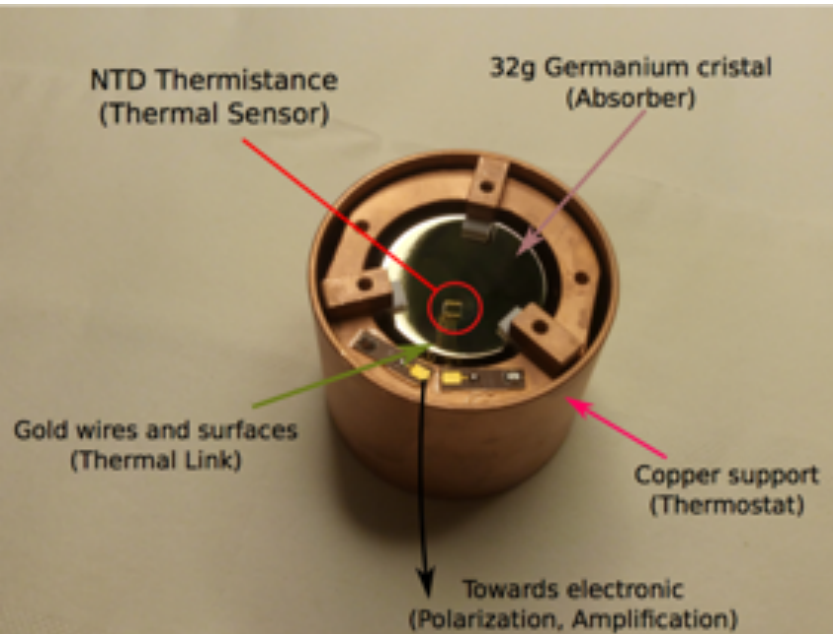


| | Sensor | Area (cm ²) | σ [eV] | σ/\sqrt{A} [eV] |
|---|--------|-------------------------|---------------|------------------------|
| CRESST 2 LD Rothe et al JLTP 193,1160 (2018) | W TES | 12.5 | 4-7 | 1.1-2.0 |
| LAPD (CDMS tech) | W TES | 45 | 3.9 | 0.58 |

NTD

[Neutron Transmutation Doped Ge
Thermistor]

NTD

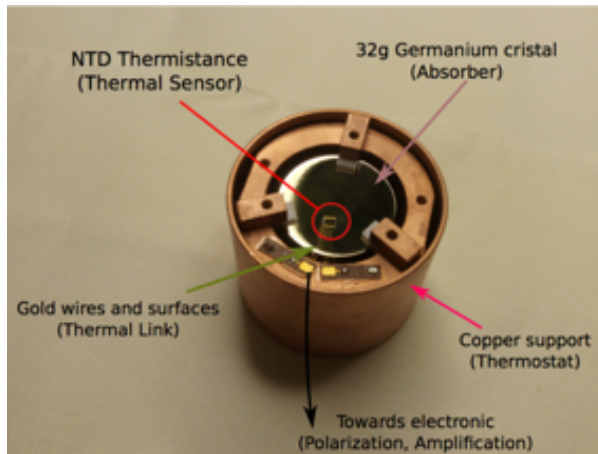


Thermal Calorimeter: $\sigma^2 \propto VT^3$

- Provided limited by TFN noise

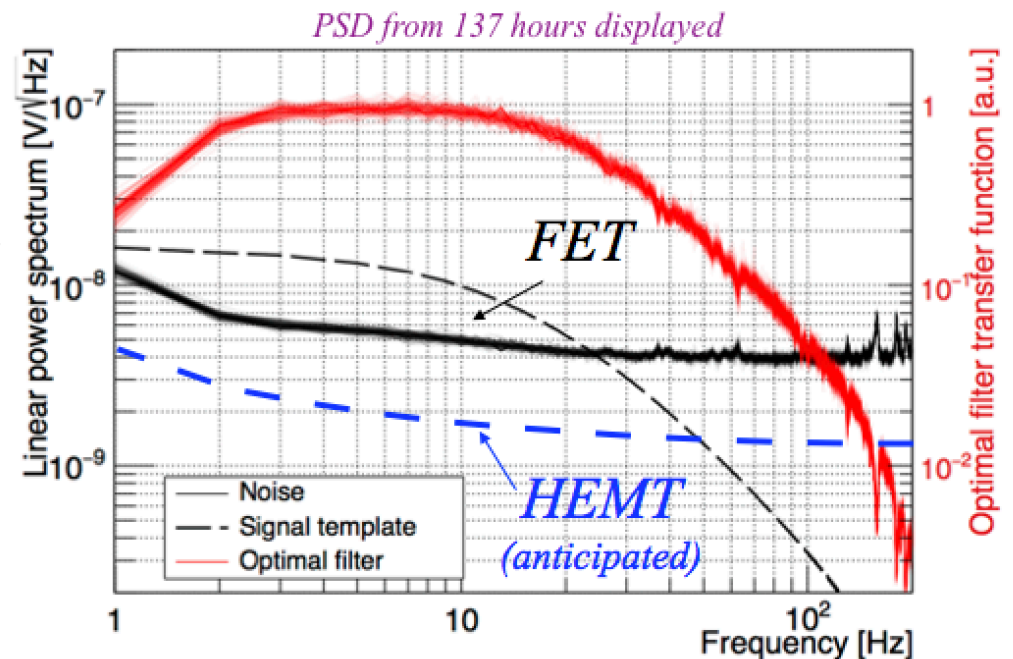
Thermal/Athermal Phonon Detector

Ge Athermal/Phonon Detector



- 17.7eV on 33.4g Ge detector
- 50eV on 200g Ge detector (x8 better than EDELWEISS 3 800g)

Hope to get to 10eV with improved electronics



Large Area Photon Detectors

Large Area Photon Detector



CRESST 2 Light Detector



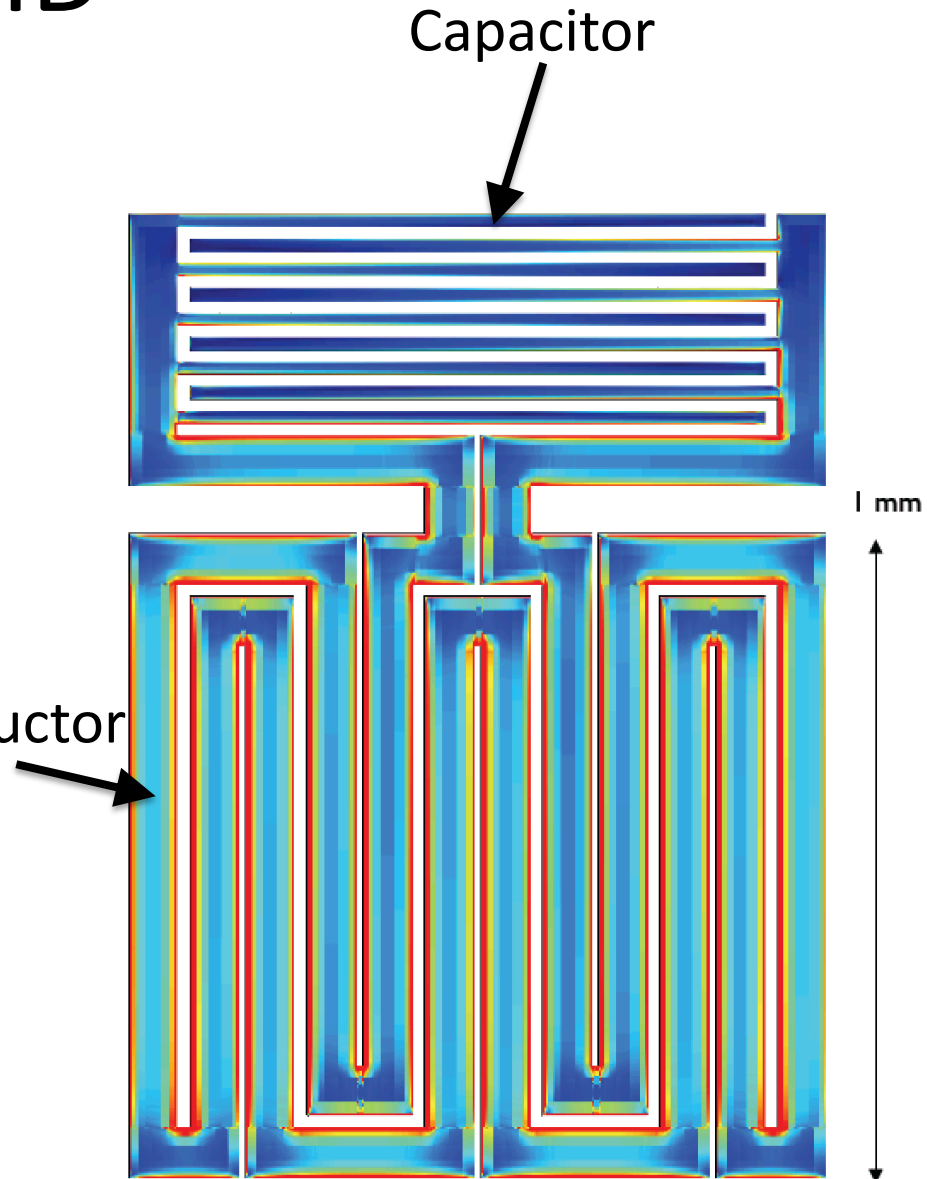
| | Sensor | Area (cm ²) | σ [eV] | σ/\sqrt{A} [eV] |
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| LAPD (CDMS tech) | W TES | 45 | 3.9 | 0.58 |
| LMO-3 LD E. Armengaud et al, Eur. Phys. J. C (2017) 77 :785 | NTD | 5 | 7.7 | 3.4 |

MKID

MKID

Superconductors have an AC inductance due to inertia of cooper pairs

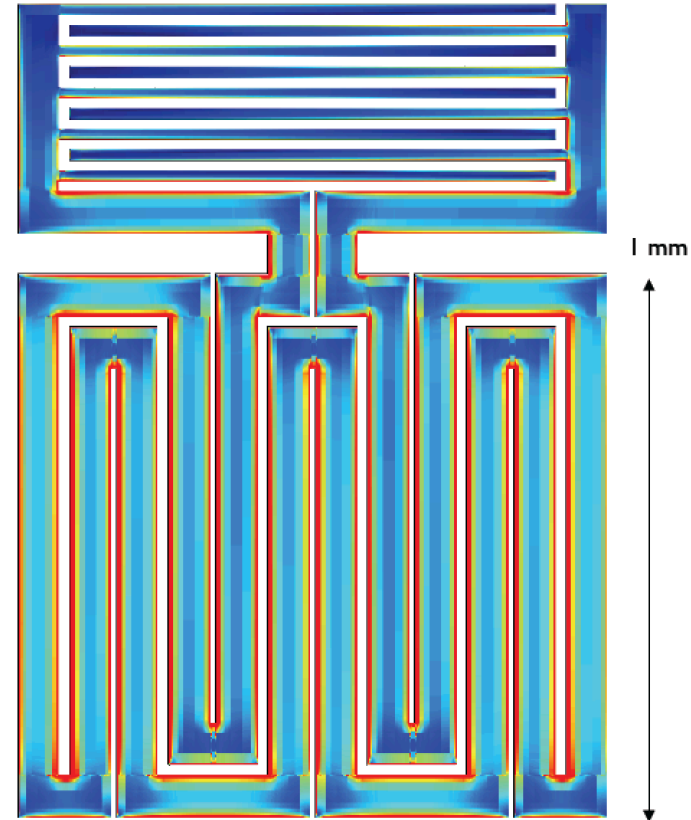
- fewer cooper pairs -> more inductance -> lower resonant frequency
- Easily Multiplexed
- 1mmx1mmx40nm active volume
- Frequency tuned by changing inductor length
- Capacitor: standard interdigitated capacitor to minimize TLS noise ... can be made out of Nb to avoid phonon absorption



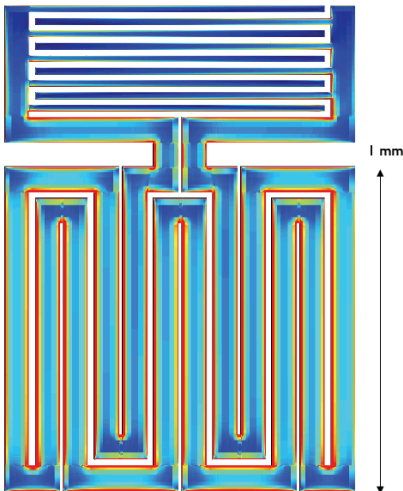
MKID Sensitivity

$$\sigma_E = 2\Delta \sqrt{\frac{\eta_a}{\chi_c \chi_{qp}}} \sqrt{\frac{k T_a N_0 V_r}{\left[2N_0 \Delta \frac{\partial(\sigma_1/|\sigma(0)|)}{\partial n_{qp}} \right] Q_s}}$$

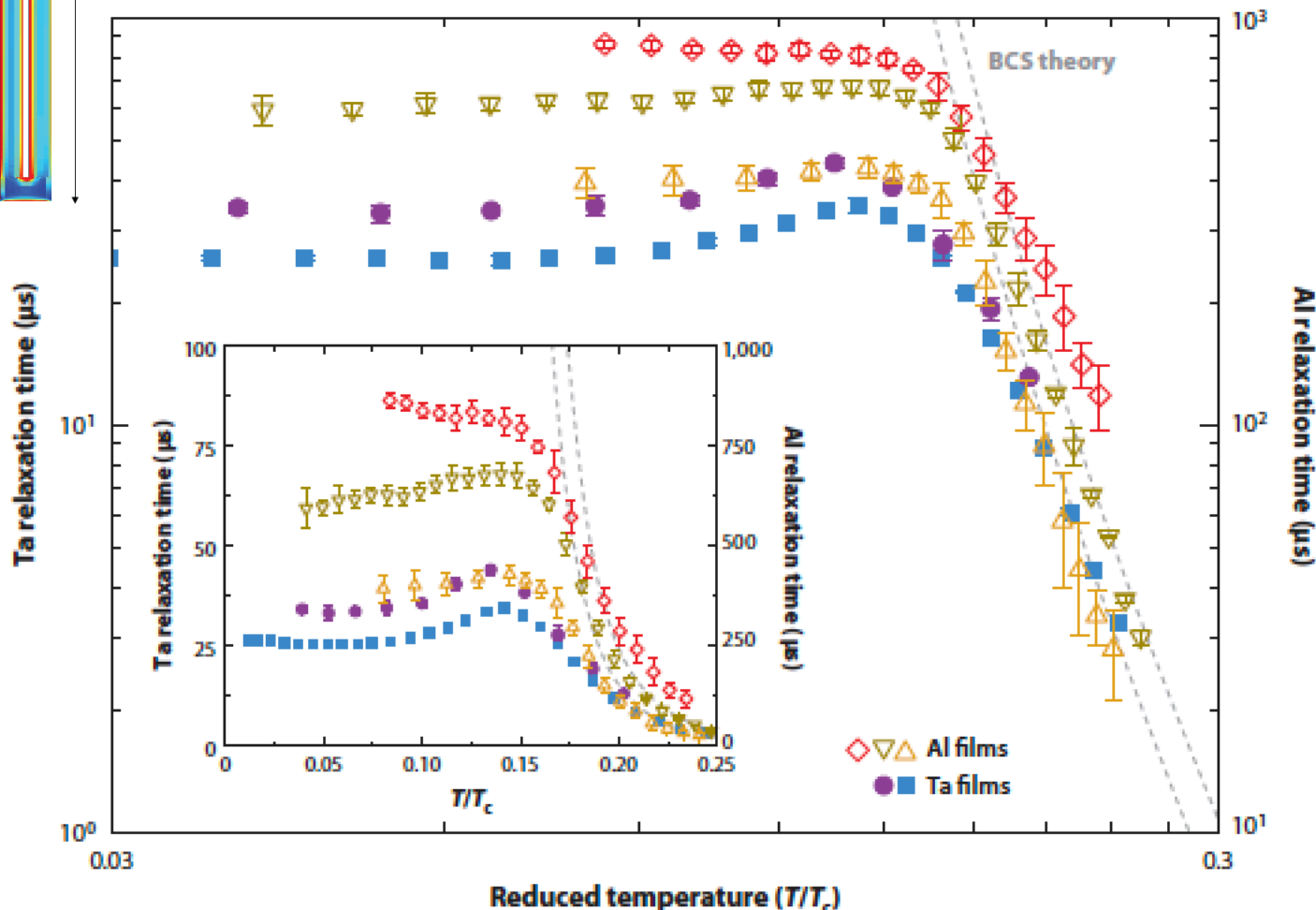
- Scales with sensor volume
- Scales with amplifier noise temperature
 - 4K -> 70mK with parametric amplifier
- Scales with T_c / Δ
- Assumes no excess quasi-particle density
- Not sensitive to all frequencies of EMI



Excess Quasi-particle Density

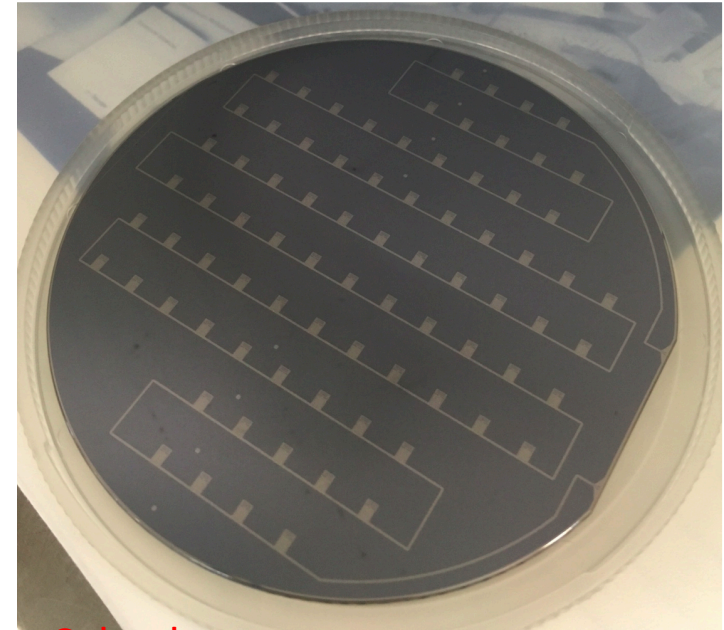
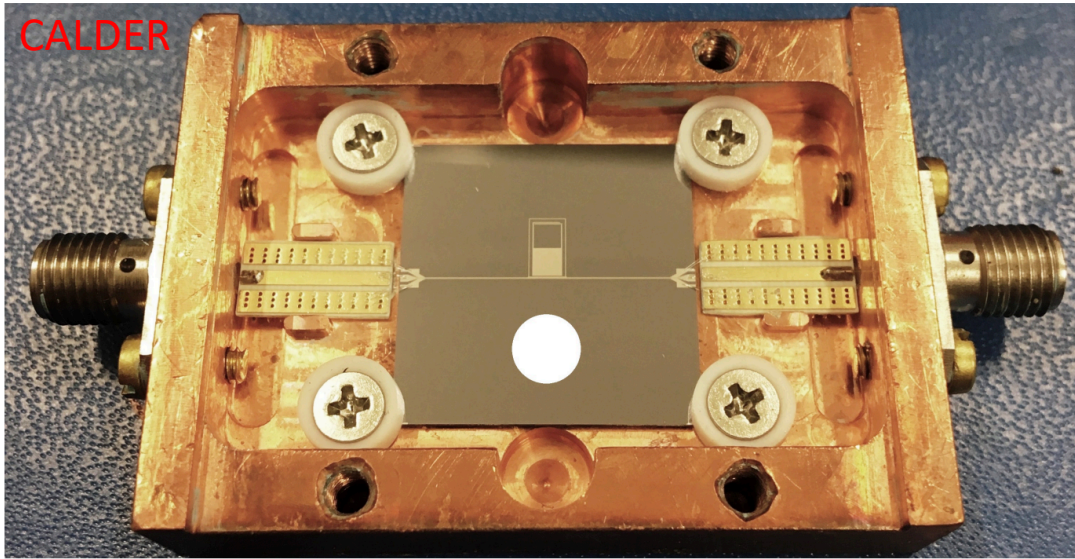


asymptotic regime; limiting excess qp density n_+ , or something else? related to disorder? (Barends et al implantation experiment)



Barends et al PRL (2008)
as reproduced in Zmuidzinas, ARCOMP (2012)

MKID: Athermal Phonon Sensor



| | Sensor | Area (cm ²) | σ [eV] | σ/\sqrt{A} [eV] |
|--|------------------|-------------------------|---------------|------------------------|
| CRESST 2 LD Rothe et al JLTP 193,1160 (2018) | W TES | 12.5 | 4-7 | 1.1-2.0 |
| LAPD (CDMS tech) | W TES | 45 | 3.9 | 0.58 |
| LMO-3 LD E. Armengaud et al, Eur. Phys. J. C (2017) 77 :785 | NTD | 5 | 7.7 | 3.4 |
| CALDER 1801.08403 | Al/Ti/Al MKID | 4 | 26 | 13 |

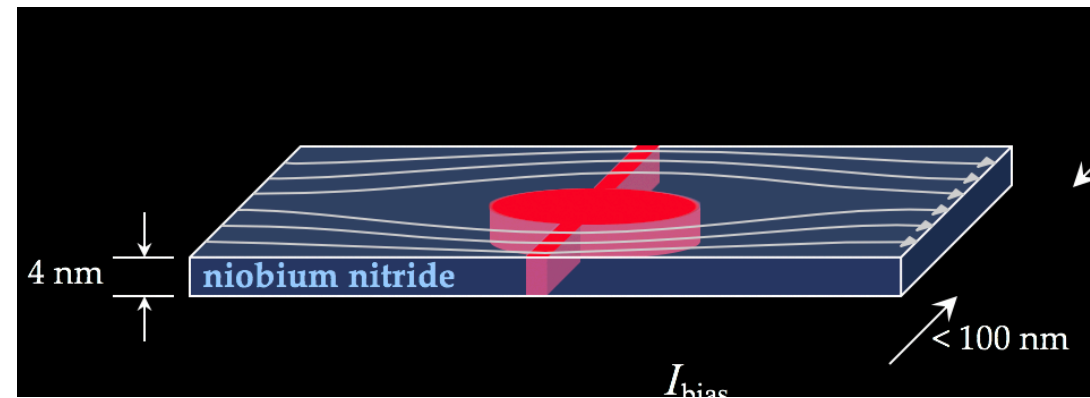
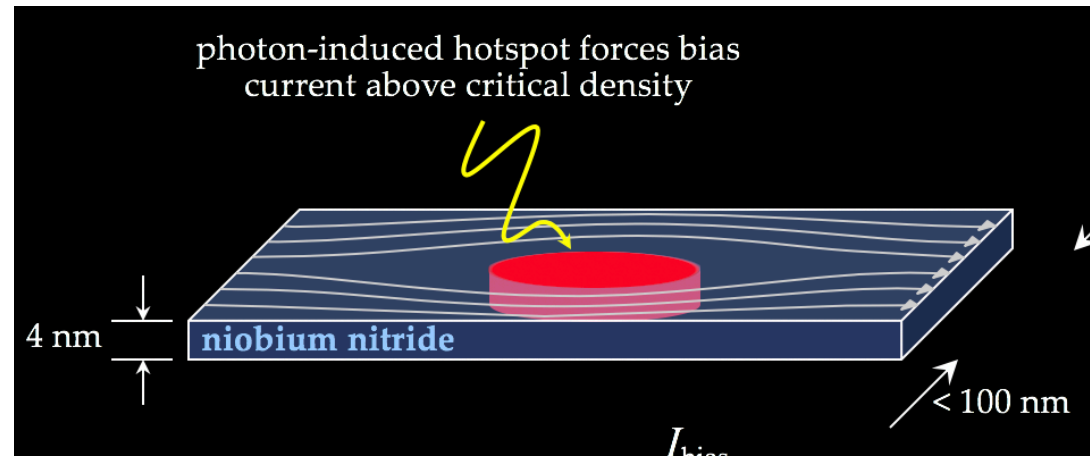
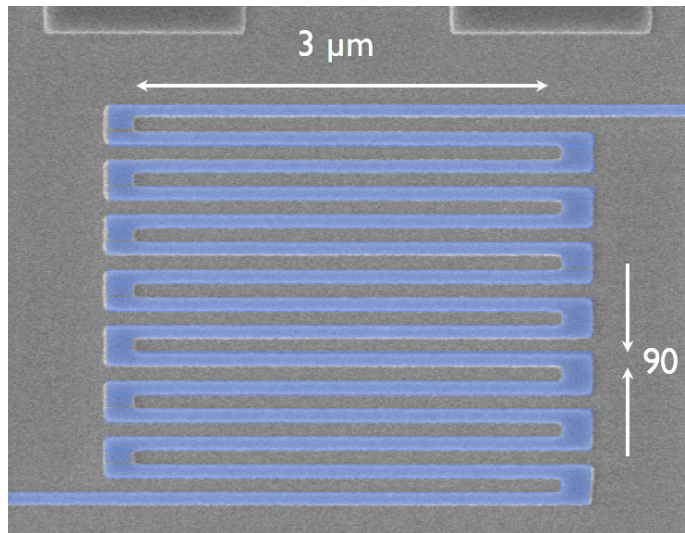
MKID R&D:

- Better Amplifiers (parametric amplifier at 70mK)
- Excess Quasi-particle density
- Two Level Systems

SNSPD

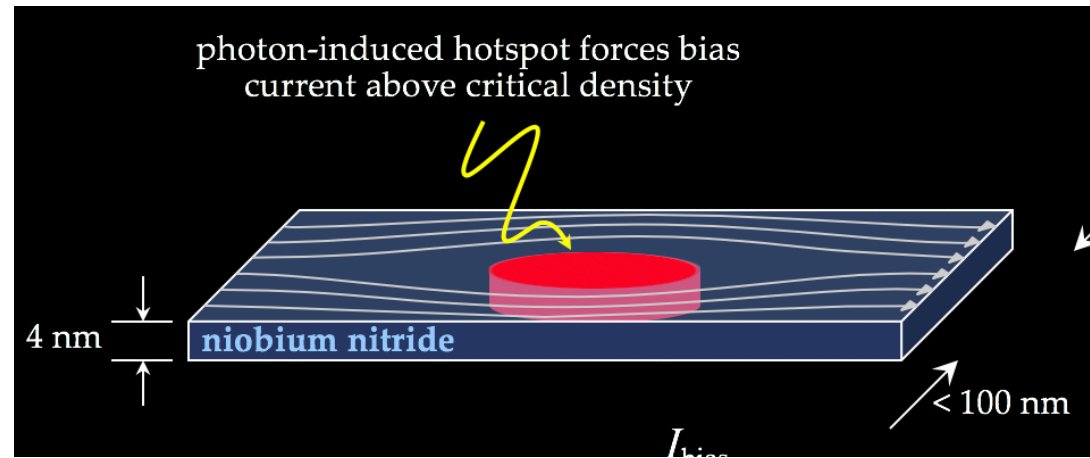
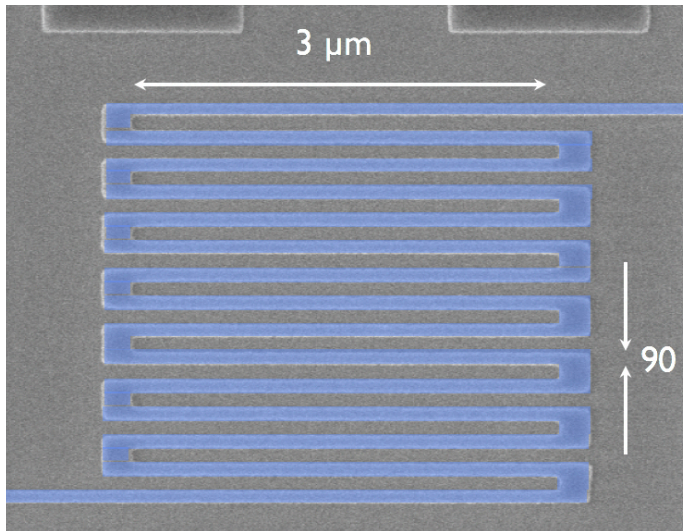
Superconducting Nanowire Single
Photon Detector

SNSPD: Operating principle



- SC wire biased very near critical current. Event causes nanowire to switch to normal state
- Natural Energy Amplification
- Event Measurement ... doesn't measure energy

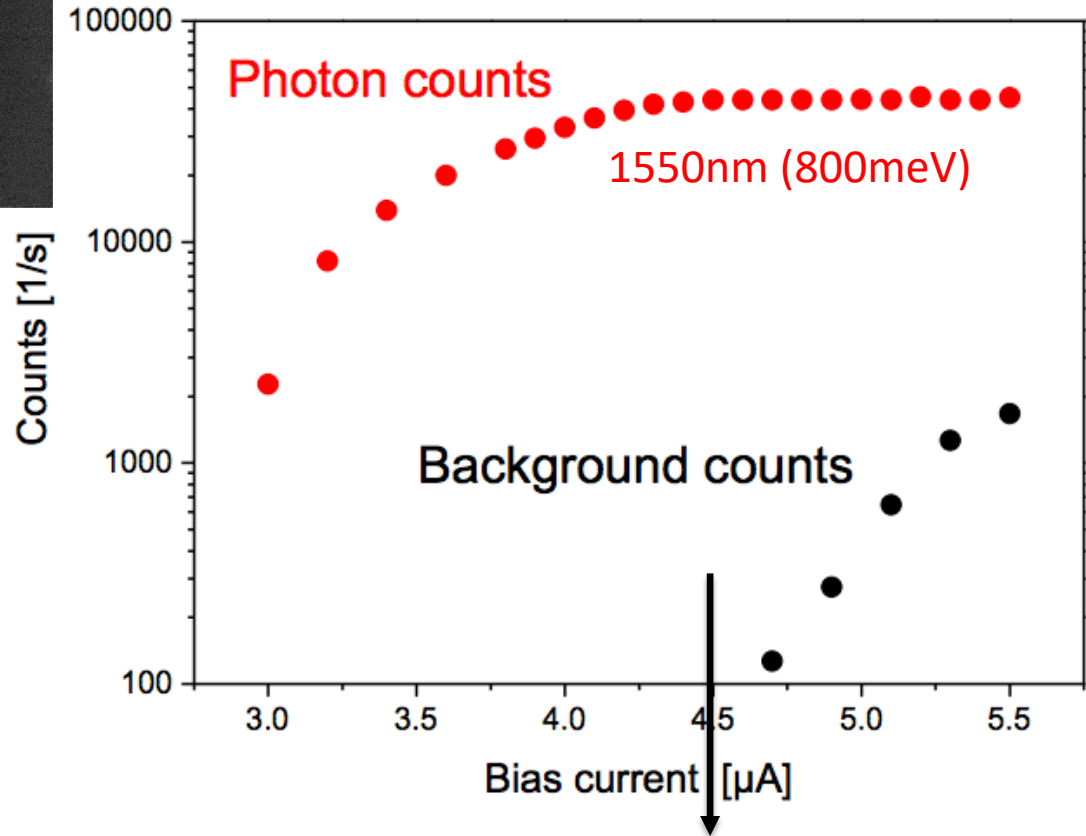
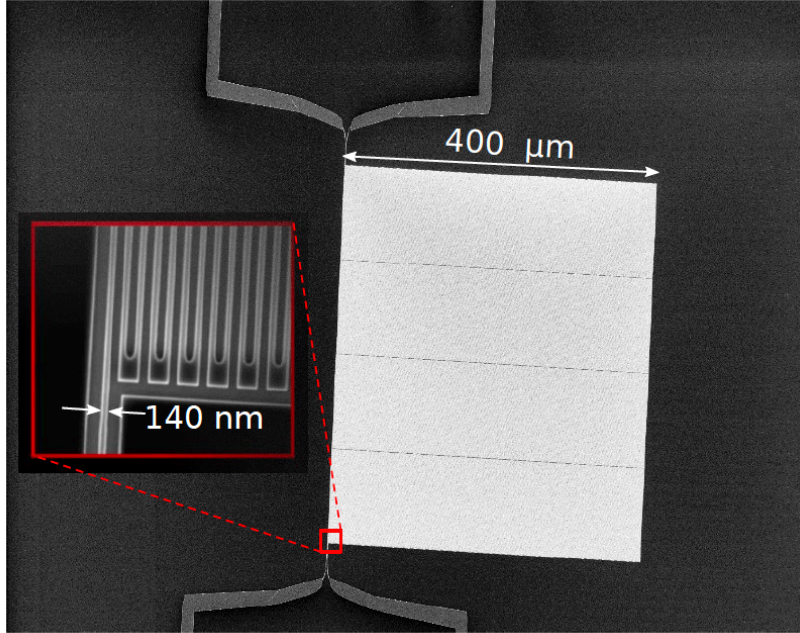
SNSPD: Energy Threshold Scalings



Scales with cross sectional area, not volume.
Can make the meander as long as you want

Current Performance

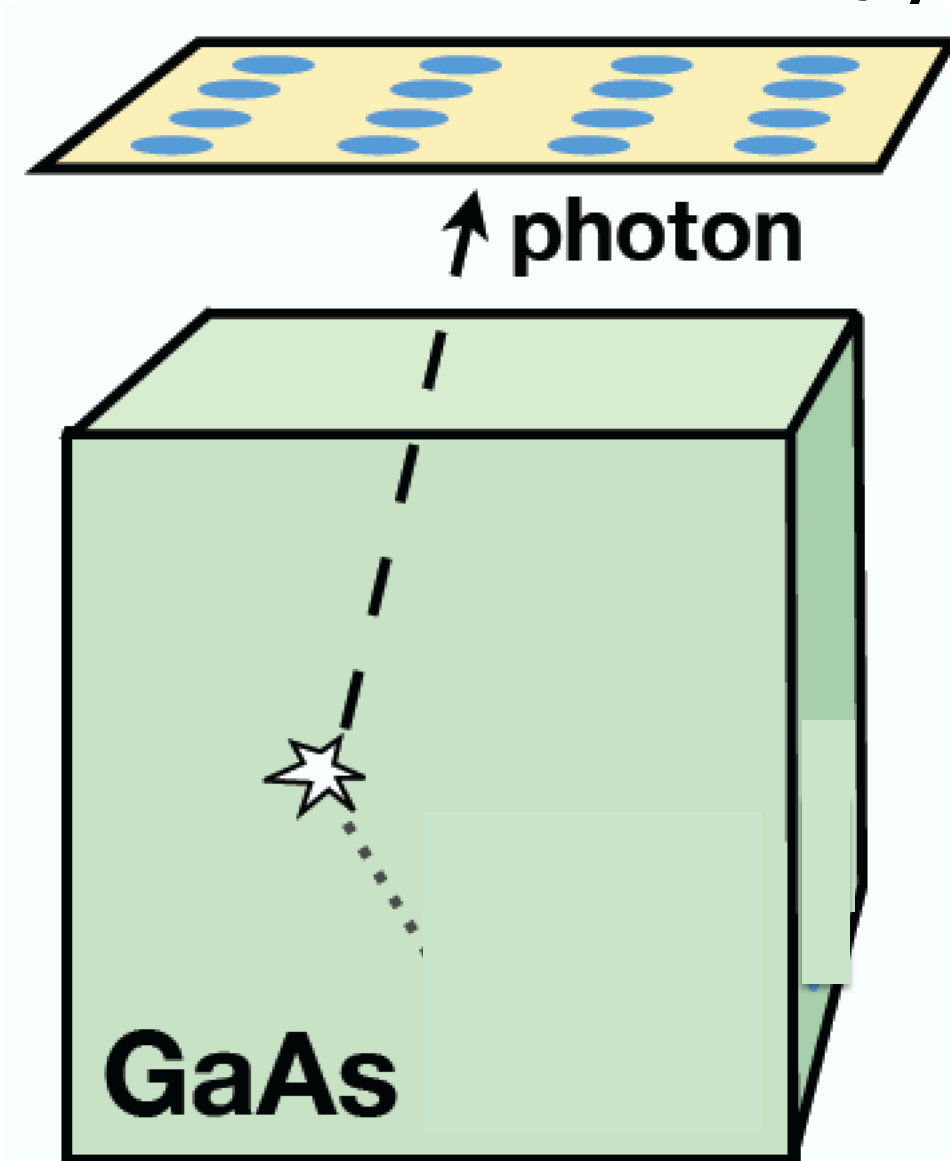
- 1903.05101
- WSi 400umx400umx7nm Meander



$< 10^{-4}$ Hz dark count rate
@ 4.5 μA

Measuring Energy With Multipixel SNSPD Arrays

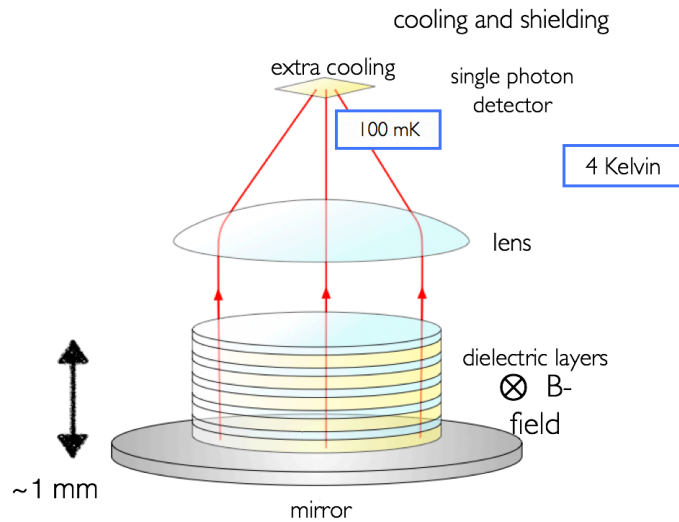
Arrays



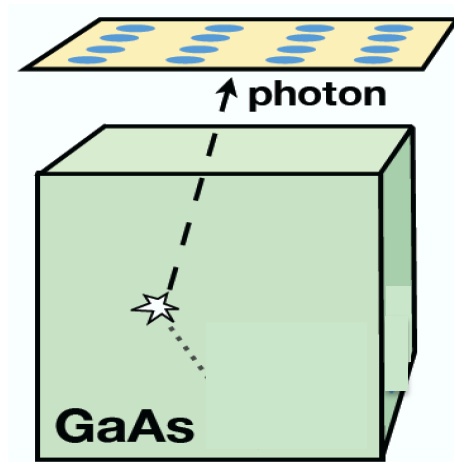
In the limit of large pixel number, the probability of 2 photons hitting the same pixel is small (and can be corrected) ... it's the same probability distribution as the shared birthday problem

SNSPD Applications

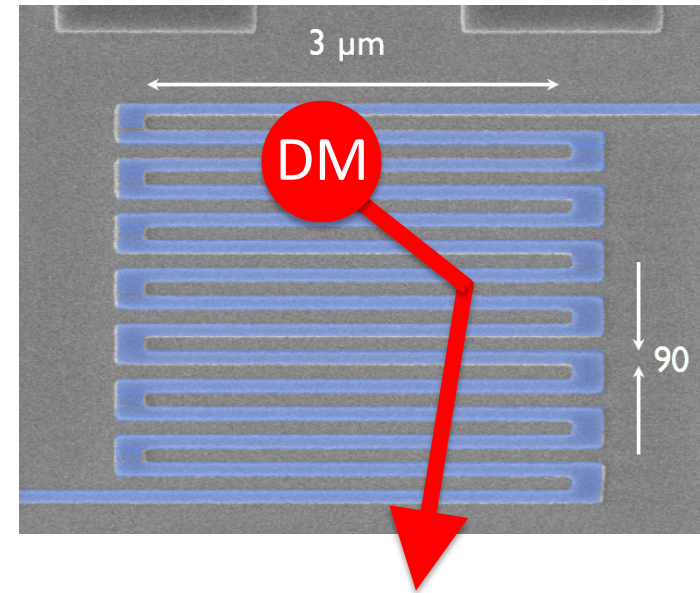
Optical/IR Haloscopes



Inelastic Electronic Recoil DM Search with Scintillator



Bulk DM search

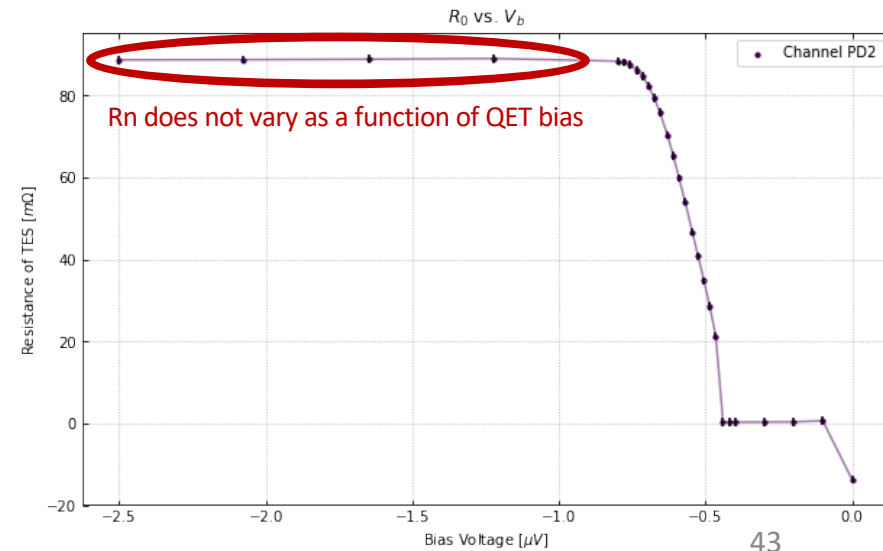
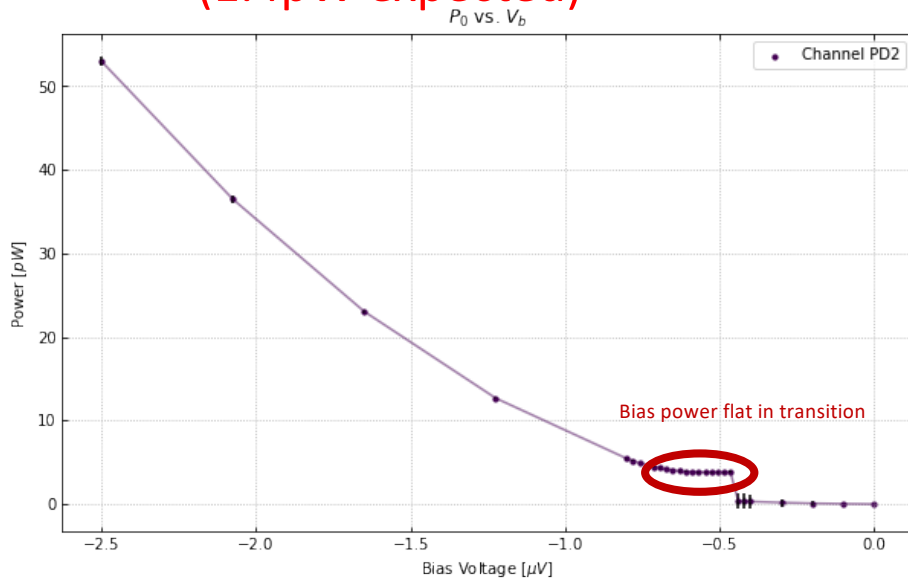
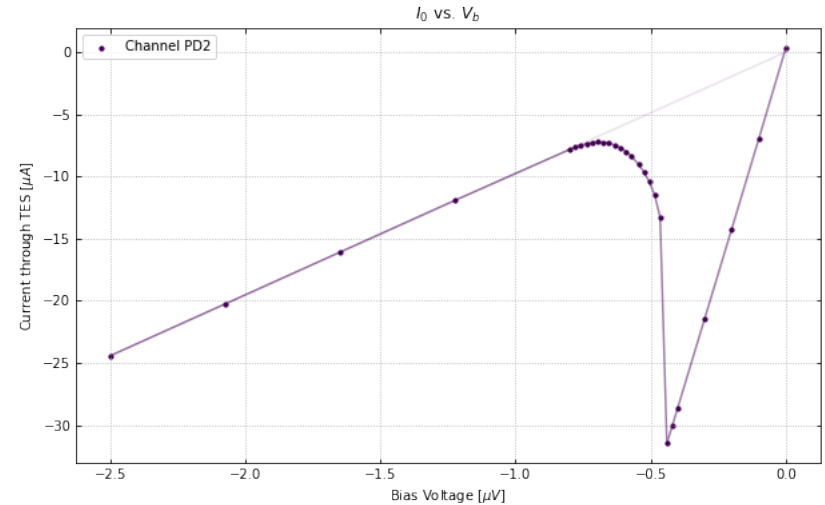


- 1903.05101
- 4.3ng per pixel
- 2.3×10^8 channels for 1g

Backup

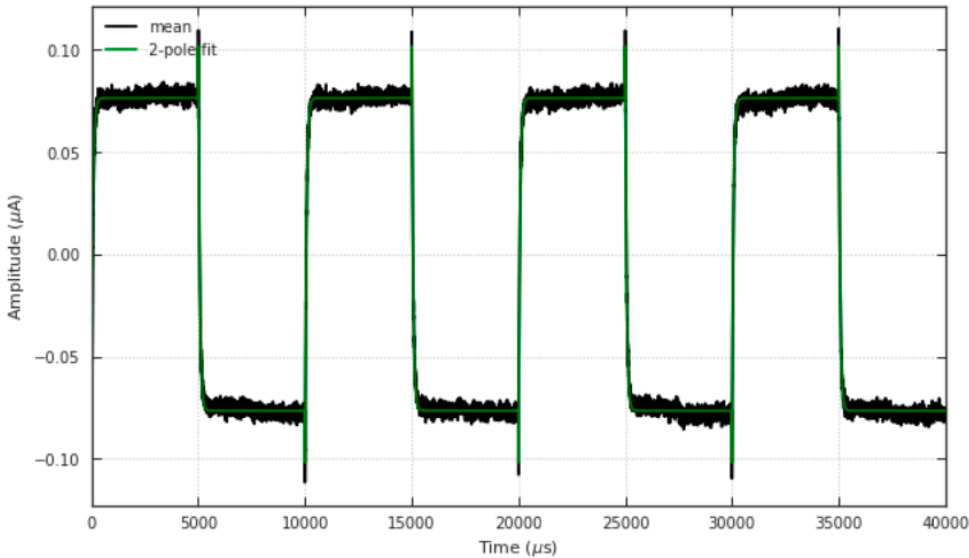
Measured Performance: Tc & IV

- $T_c = 41.5\text{mK}$
- IV curves show that the detector and electronics are behaving well
 - $R_n = 88\text{ mOhms}$ (300mOhm Expected ... TES too wide!)
 - $R_p = 8\text{ mOhms}$
 - Bias Power (P_0) = 3.9 pW (1.4pW expected)



Measured Performance: dIdV

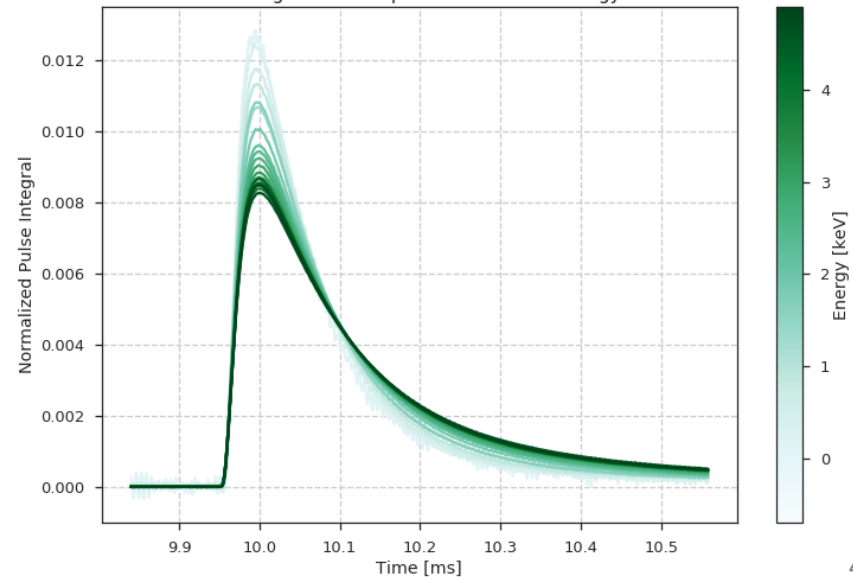
Full Trace of dIdV



- Therefore, we expect phonon signals to have a 20us rise time (athermal phonon collection) and a 60us fall time. **Seen for low energy comptons in average pulse shape!**
- Pulse shape varies with energy due to local TES saturation.

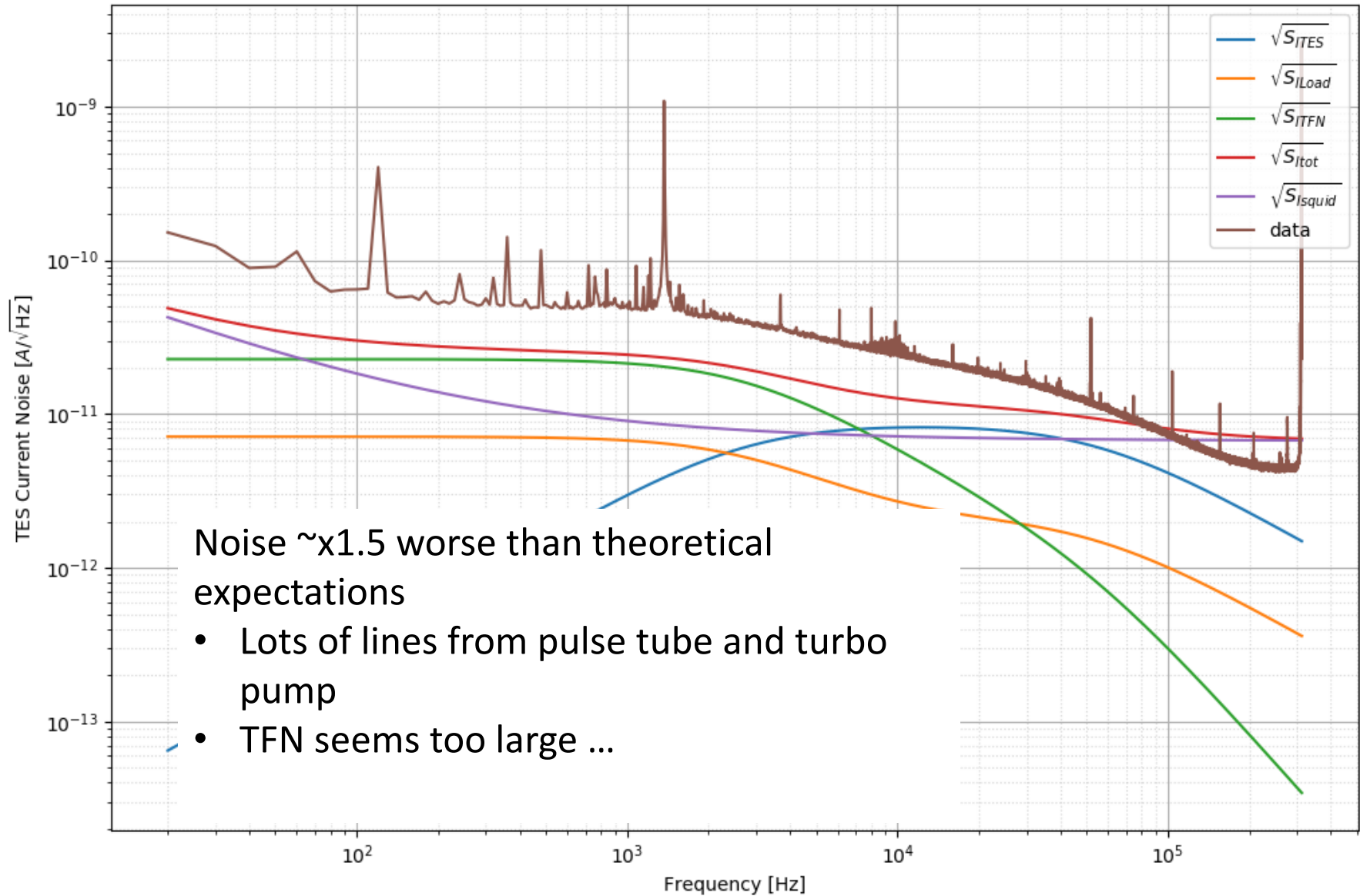
- TES sensor pretty fast @ 60us. However, it's not as fast as the estimated athermal phonon collection (20us)

Average Pulse Shape Variance With Energy



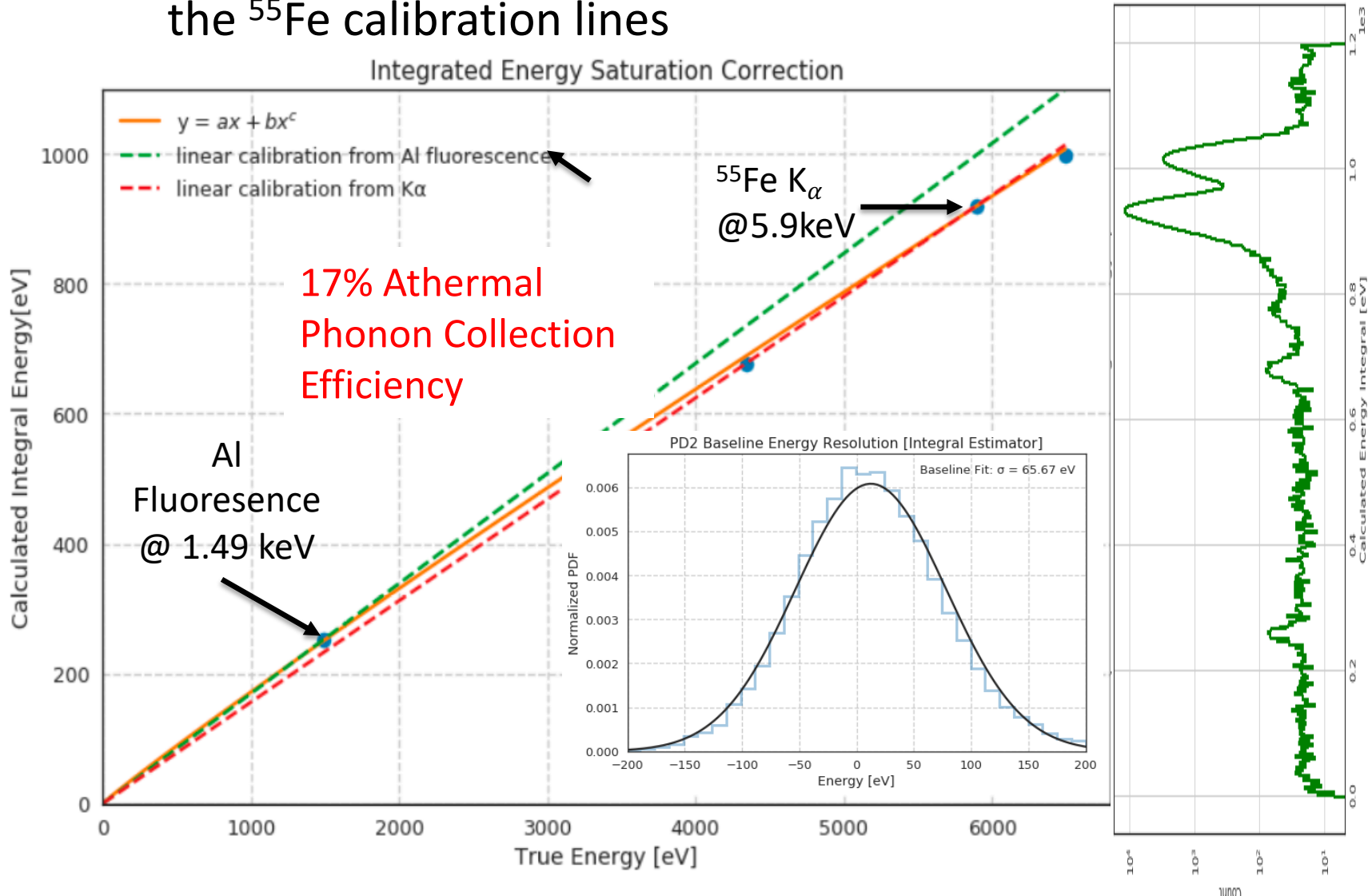
Measured/Theoretical Noise

Current Noise For $R_0 : 32.00 \text{ m}\Omega$



Integral Estimators for relative ^{55}Fe calibration

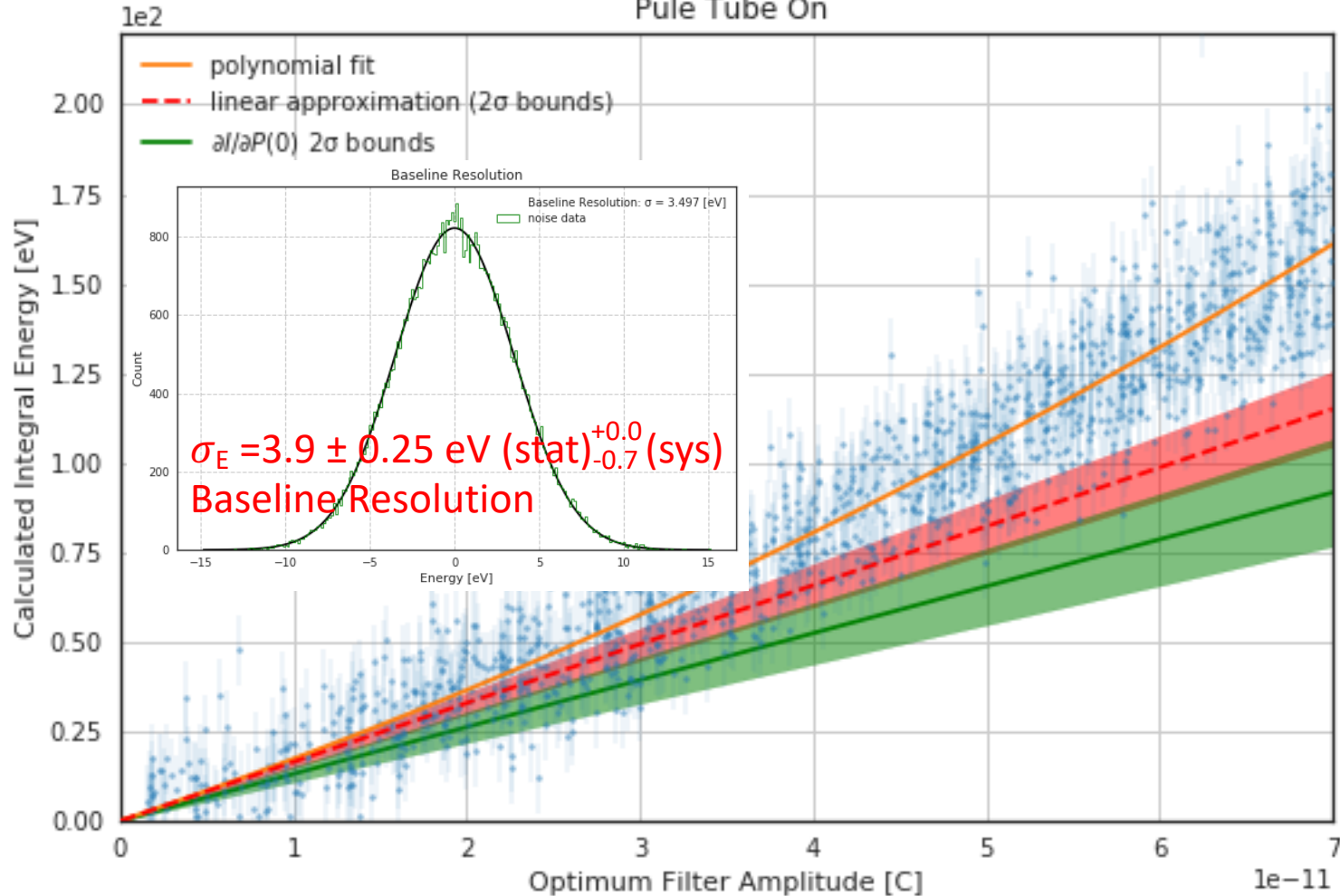
- Since pulse shape has significant variation with energy, we must use noisy but minimally biased DC estimators to fit the ^{55}Fe calibration lines



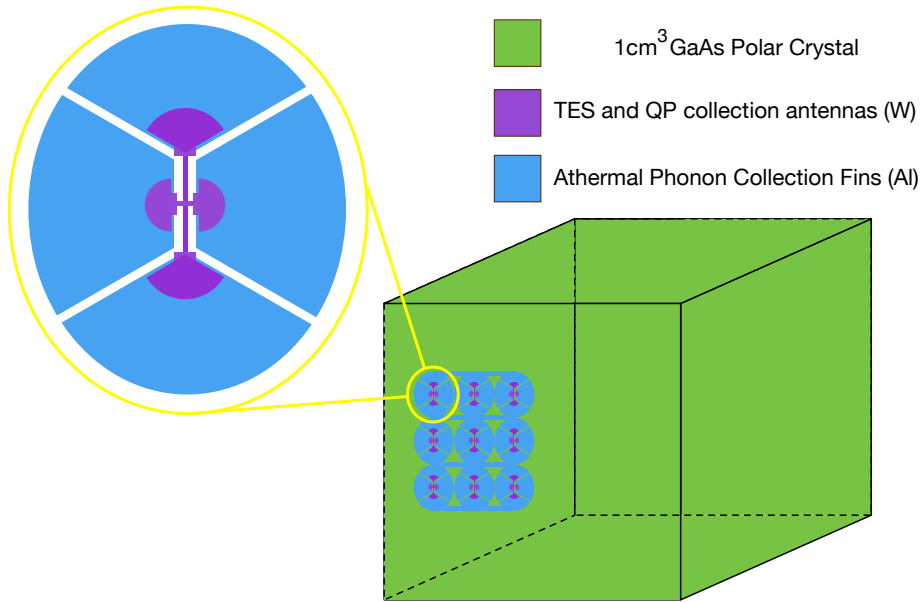
Calibrating Pulse Shape Dependent Energy Estimators to the DC estimator (Pulse Tube On)

Optimum Filter Amplitude vs Integral Energy

Pulse Tube On



Surface Dark Matter Search Completed at SLAC



Prototype Design

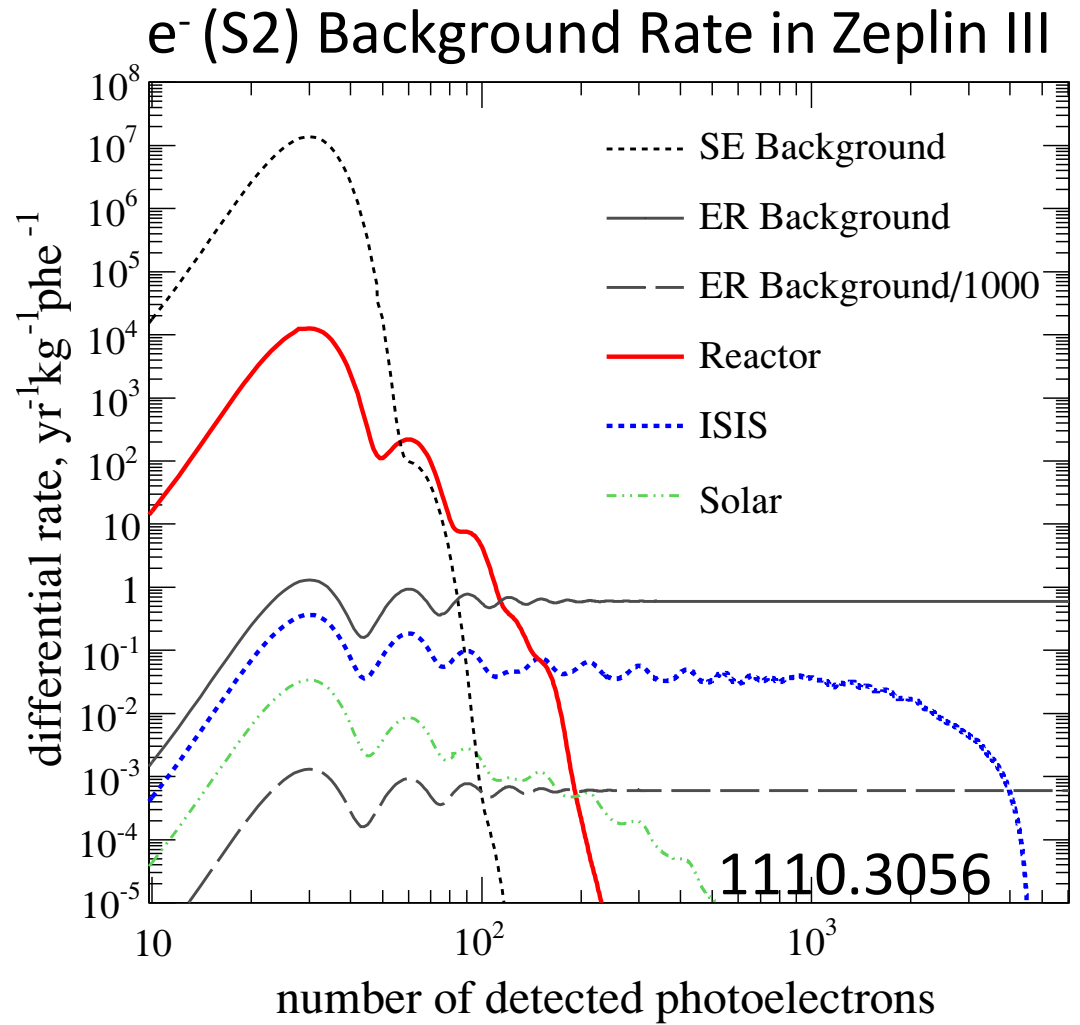
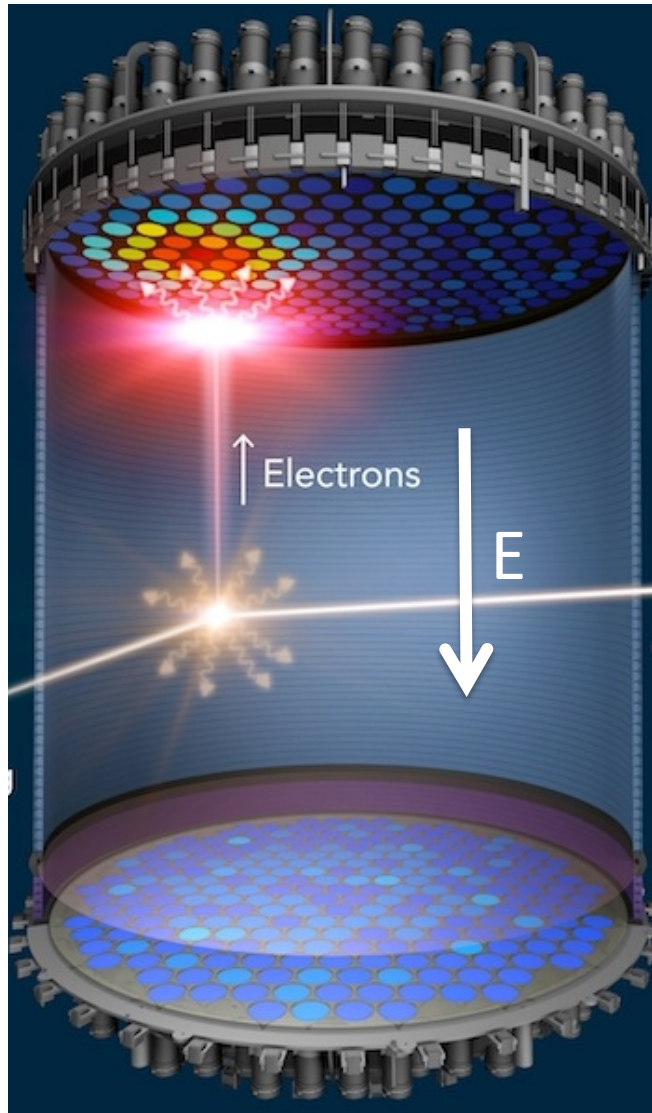
Estimated Sensitivities

New 1cm³ Prototype Test Design

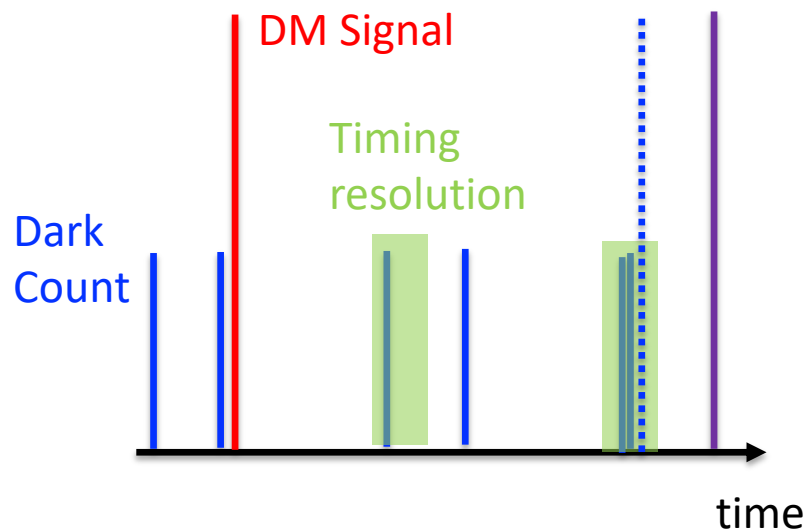
| | |
|------------------------|-------------------|
| # TES | 100 |
| TES Dimensions | 50um x 2um x40 nm |
| TES Rn | 320mOhm |
| Fin Length | 125um |
| W/Al Overlap | 15um |
| Fractional Al Coverage | 1% |
| Tc | 40mK |
| Bias Power | 48fW |
| Power Noise | 5.1e-19 W/rtHz |
| Phonon absorption time | 106us |
| Sensor fall time | 97us |
| Collection efficiency | 19% |
| σ_E | 219 meV |

- With a Si Absorber: single e/h sensitivity without Luke-Neganov gain. Can be used for inelastic electronic recoil DM
- World Leading Elastic Nuclear Recoil DM search potential

Problem: Detector Backgrounds in TPCs



Dark Leakage Needs to be Poissonian



$N e^-$ background

- $N 1e^-$ events occur within detector timing resolution (Poissonian Leakage)
- $N e^-$ leakage event (Non-Poissonian Leakage)

Xenon TPCs:

- $R_{1e^-} = 10\text{Hz}$
- $\Delta t = 100\text{ns}$
- $R_{2e^-}(\text{Poissonian}) = 10^{-5}\text{Hz}$

Due to fast timing Xe TPCs can handle a relatively high $1e^-$ rate and still have $2e^-$ bin free. Unfortunately, leakage is non-poissonian (R&D needed)