Dispersive analysis of $\gamma^*\gamma^* \to \pi\pi, \pi\eta$

arXiv:1909.04158

Igor Danilkin

in coll. with O. Deineka & M. Vanderhaeghen

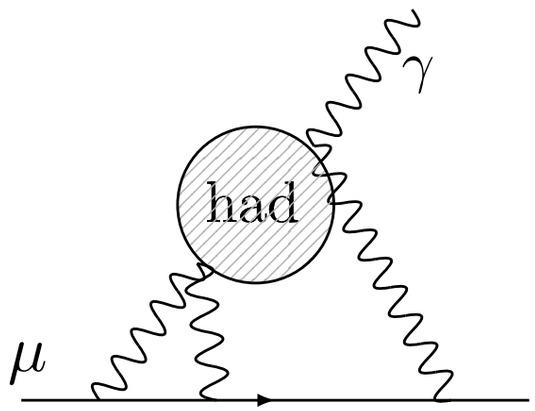
September 12, 2019

SFB 1044
THE LOW-ENERGY FRONTIER OF THE STANDARD MODEL

JOHANNES GUTENBERG UNIVERSITÄT MAINZ

JG|U

QCD contribution to (g-2)



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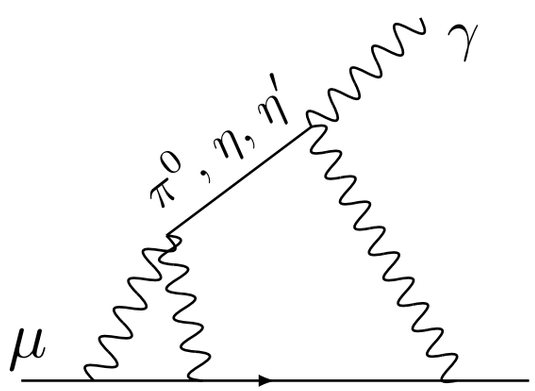

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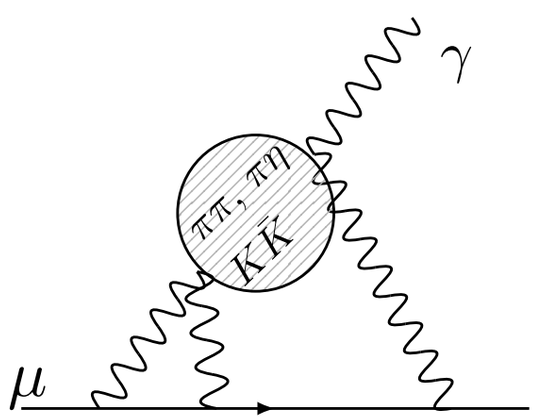
Relies on measurements of **TFF** $\pi^0\gamma^*\gamma^{(*)}, \eta\gamma^*\gamma^{(*)}, \ldots$ to reduce the model dependence



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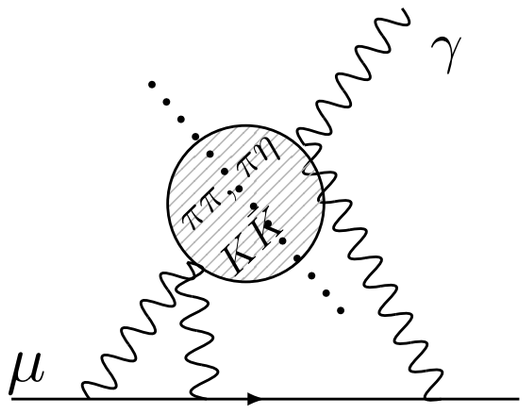
Dispersive analysis for ππ, πη, … loops is needed

…

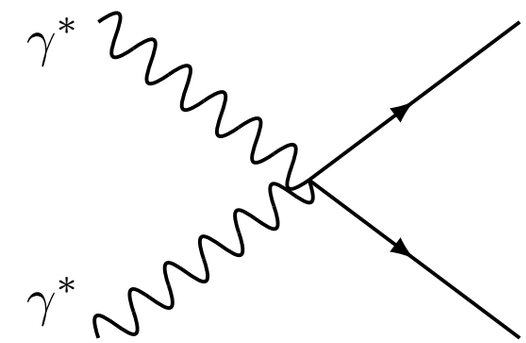
Multi-meson production



Important ingredient:

$\gamma^*\gamma^* \to \pi\pi, \pi\eta, \ldots$

$q^2 = -Q^2 < 0 \quad \text{space-like } \gamma^*$



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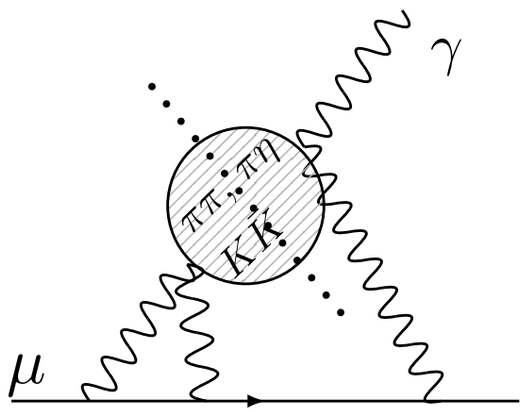

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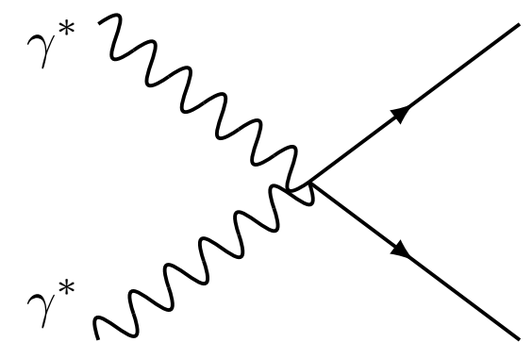
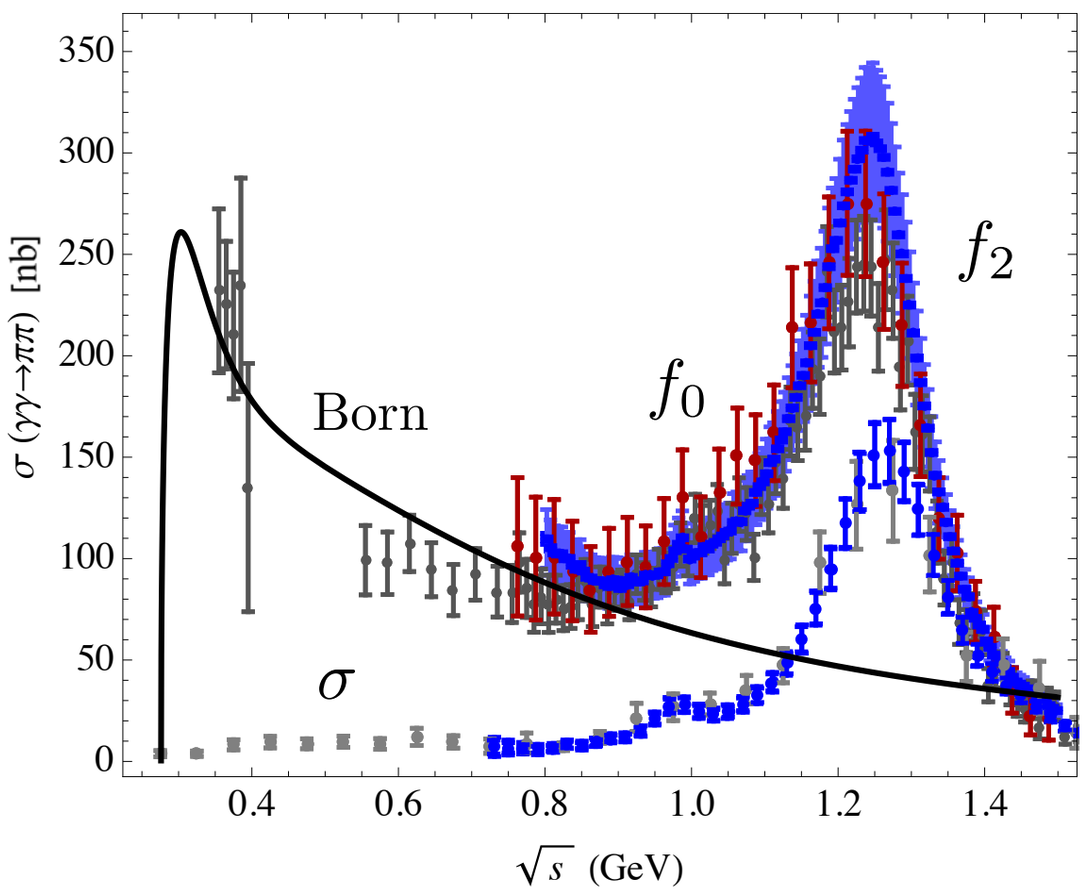
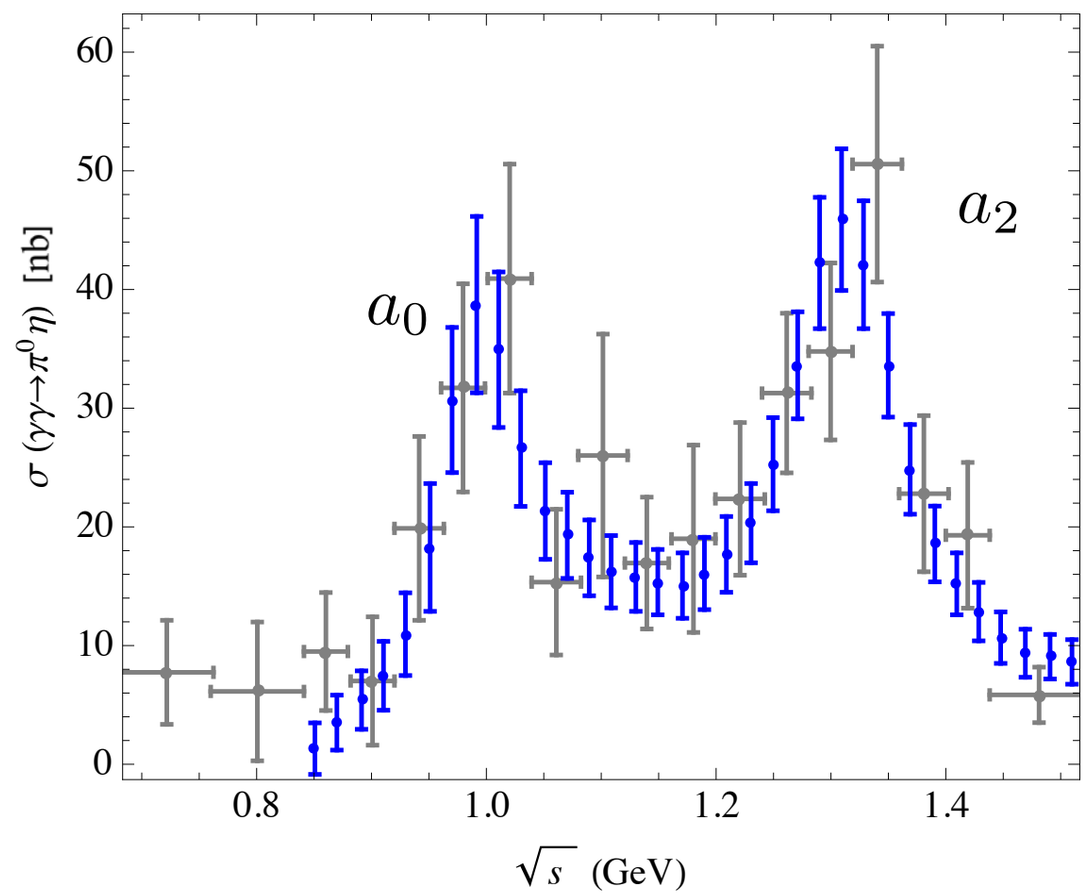


$\gamma\gamma \to \pi^+\pi^-, \pi^0\pi^0$

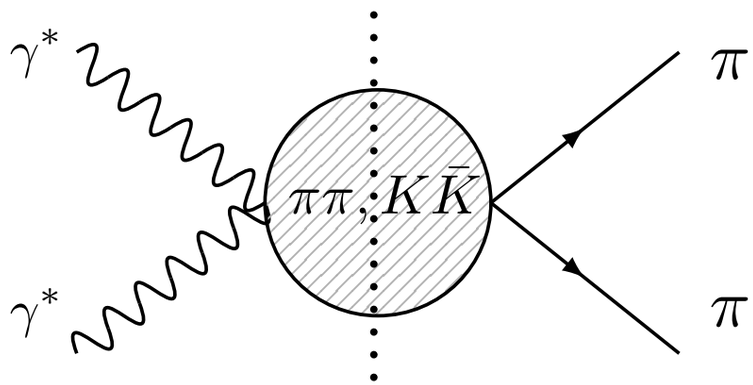


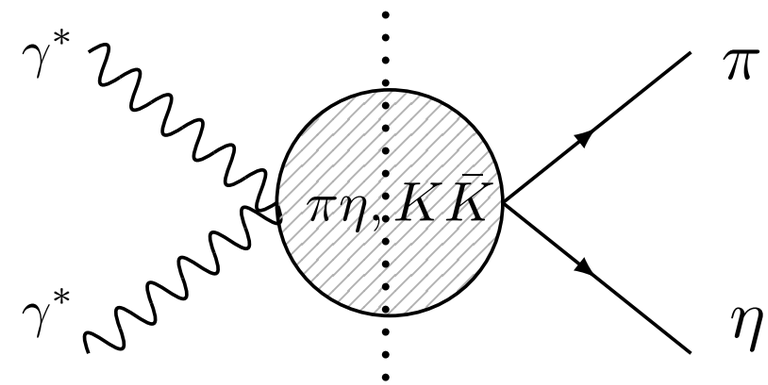
$\gamma\gamma \to \pi^0\eta$



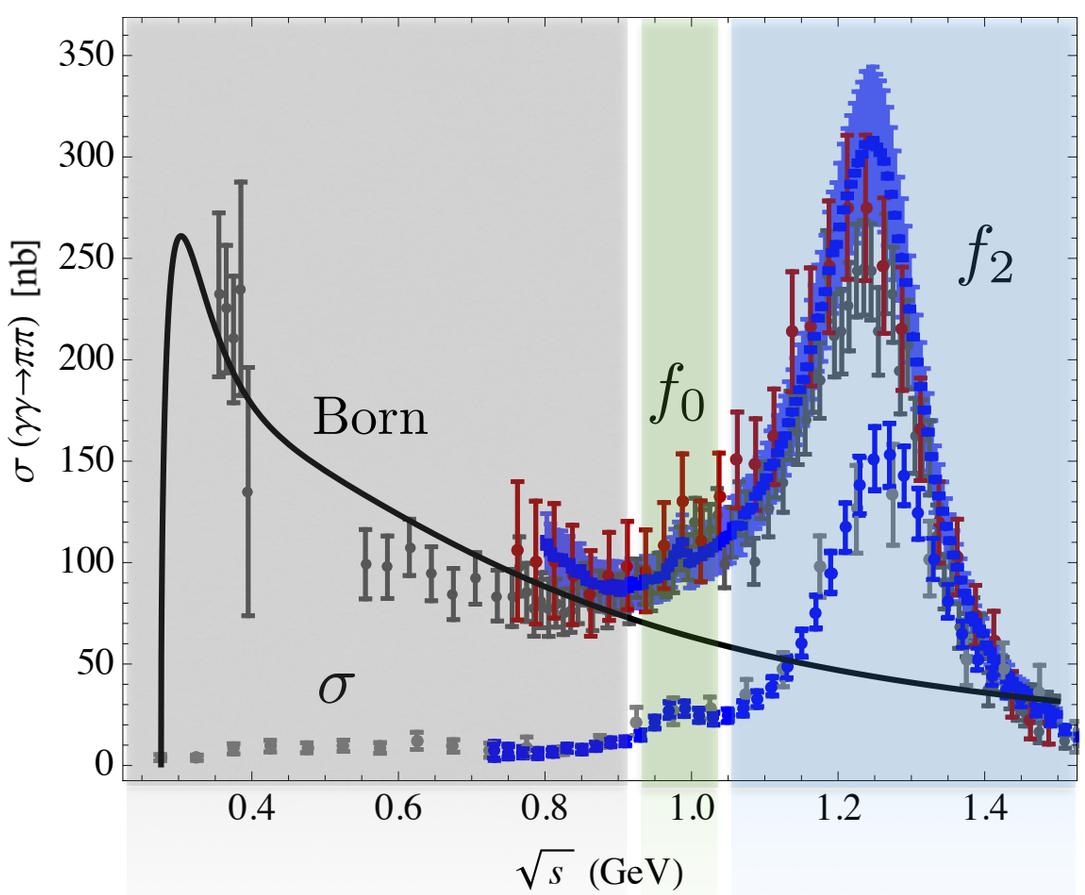
MarkII('90), CELLO ('92), Crystal Ball ('90), Belle ('07 '09)
Crystal Ball ('90), Belle ('09)

Inelastic contributions





$\gamma\gamma \to \pi^+\pi^-, \pi^0\pi^0$



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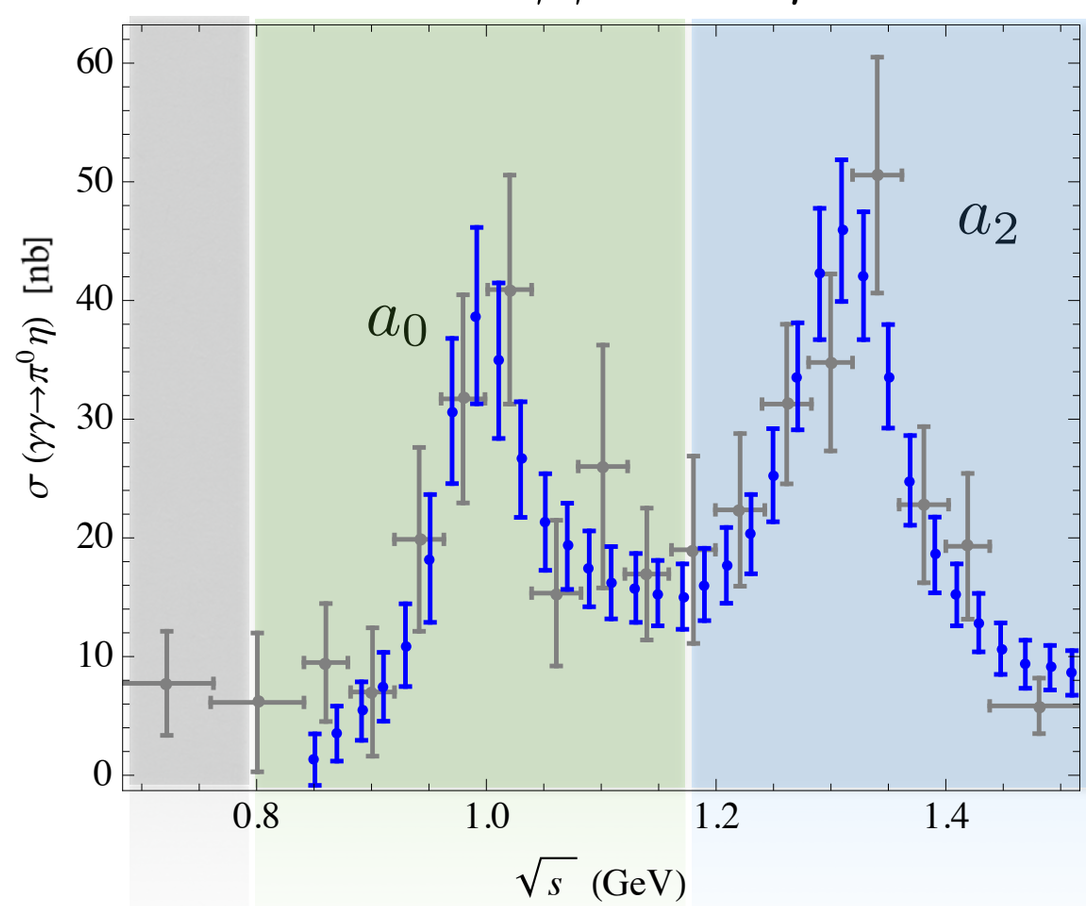

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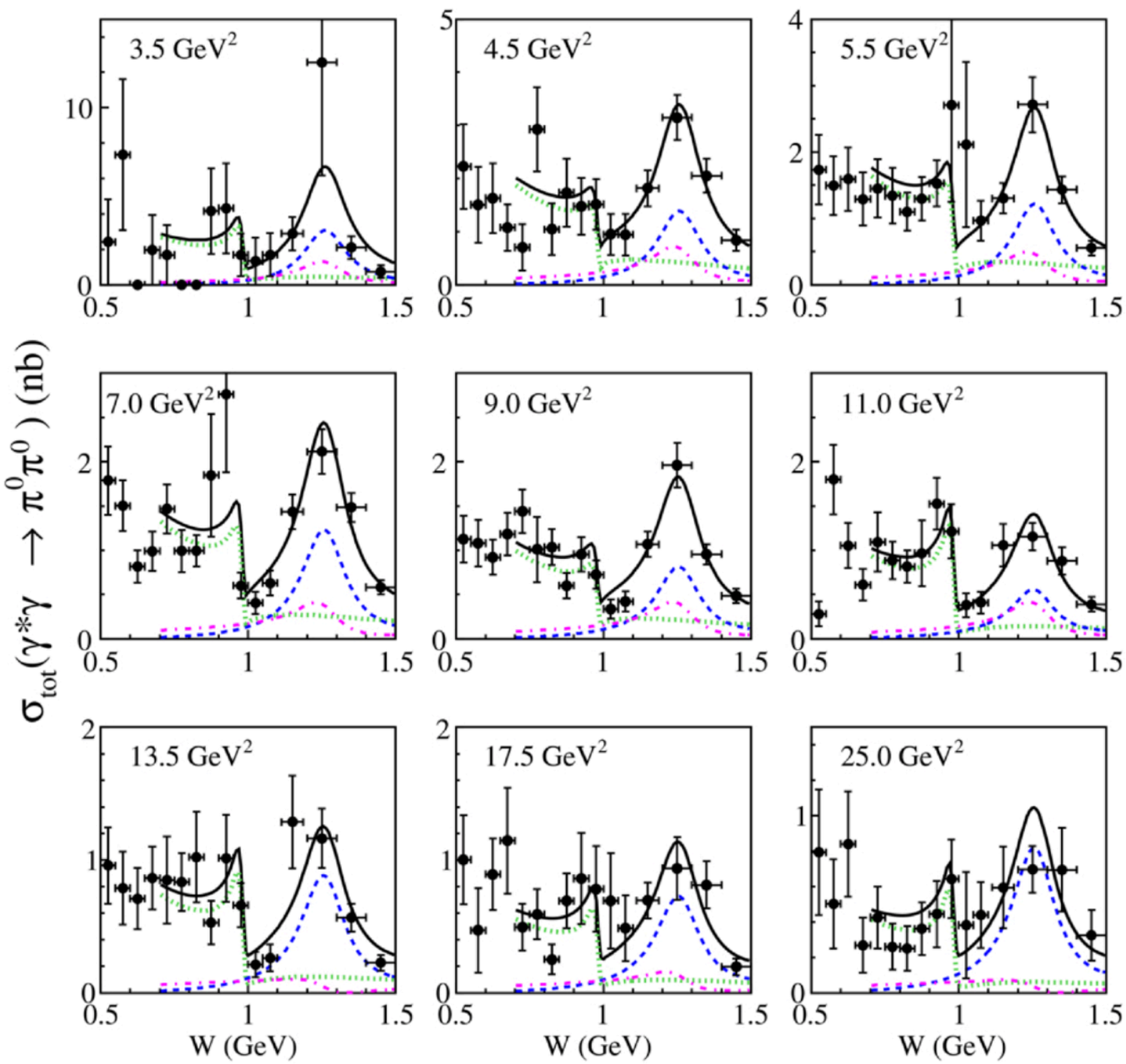
Single virtual data



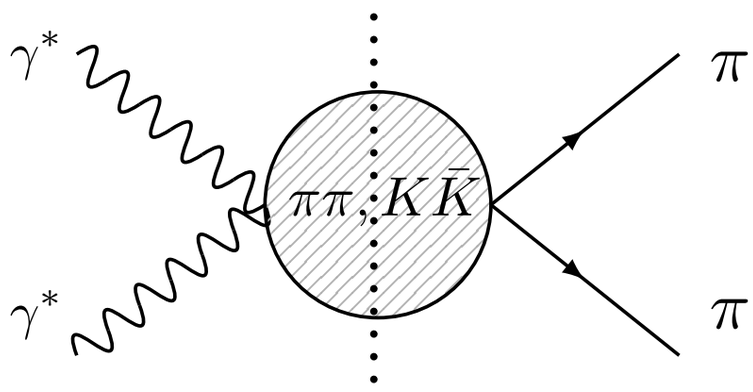
- Belle (2016)

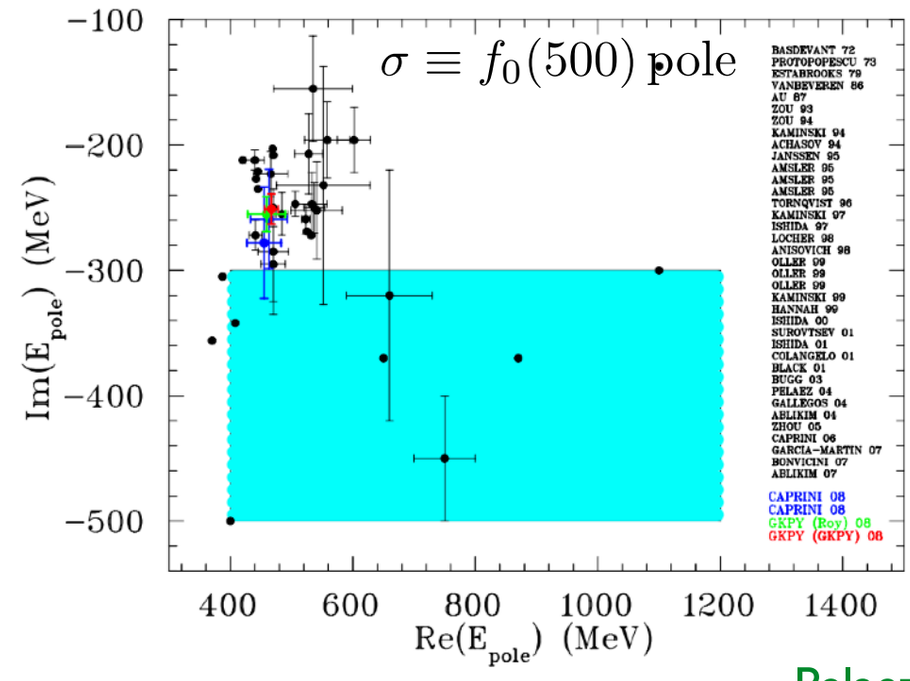
$$\gamma\gamma^* \to \pi^0\pi^0$$
$$Q^2 = 3.5 - 25\,\mathrm{GeV}^2$$

- BESIII (under analysis)

$$\gamma\gamma^* \to \pi^+\pi^-$$
$$\gamma\gamma^* \to \pi^0\pi^0, \pi^0\eta$$
$$Q^2 = 0.2 - 2.0\,\mathrm{GeV}^2$$

Right-hand cuts (hadronic input)





Pelaez (2016)

$$\sqrt{s^{\mathrm{II}}_{\mathrm{Roy}}} = \left(449^{+22}_{-16}\right) \pm i\,(275 \pm 12)\ \mathrm{MeV}$$

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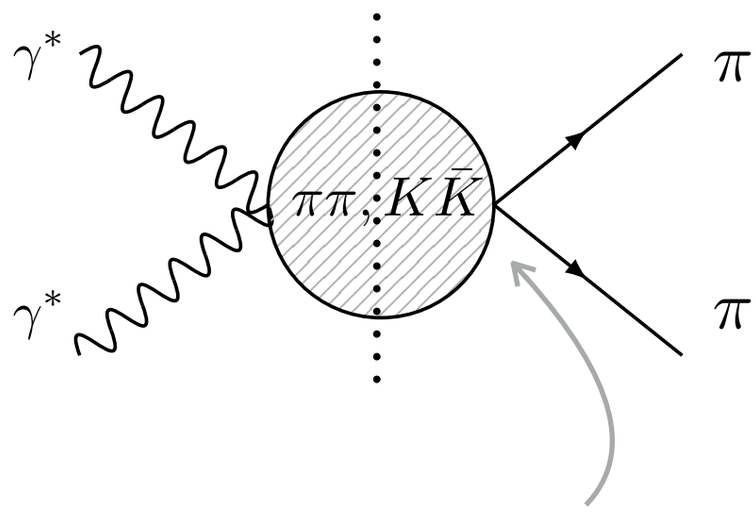

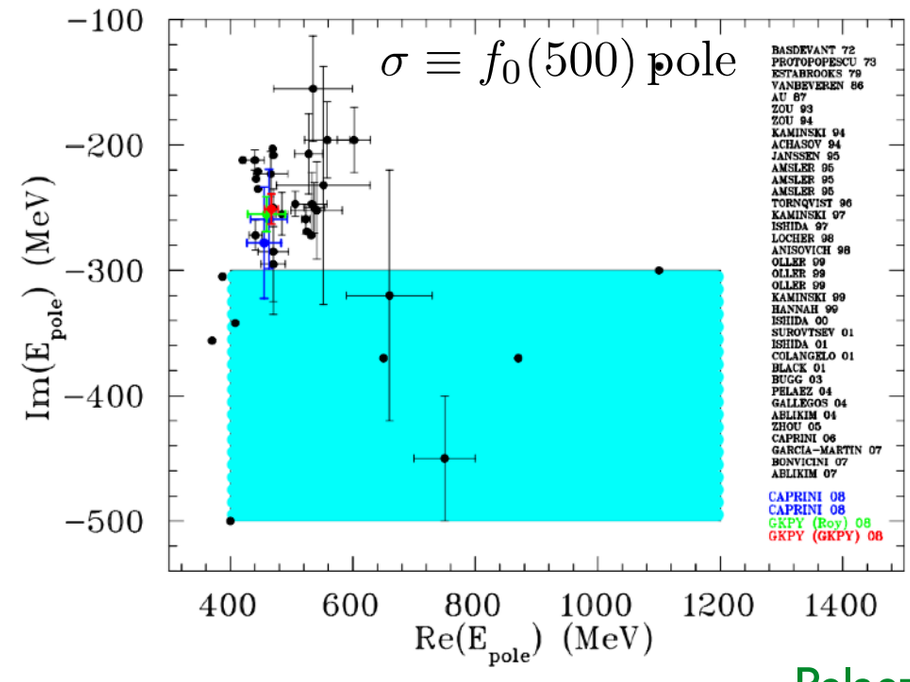

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- Coupled-channel Omnès formalism

$$t(s) = U(s) + \int_R \frac{ds'}{\pi} \frac{\rho(s')\,|t(s')|^2}{s' - s}$$

$$t(s) = \frac{N(s)}{D(s)} = \Omega(s)\,N(s)$$ Chew, Mandelstam (1960)

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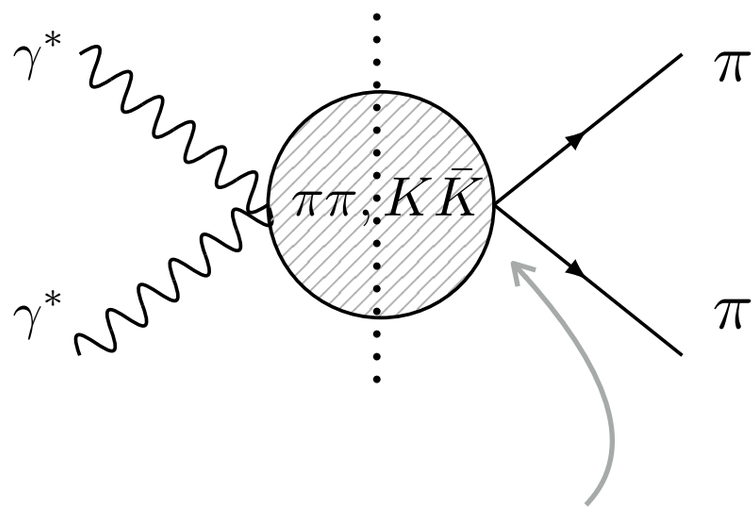

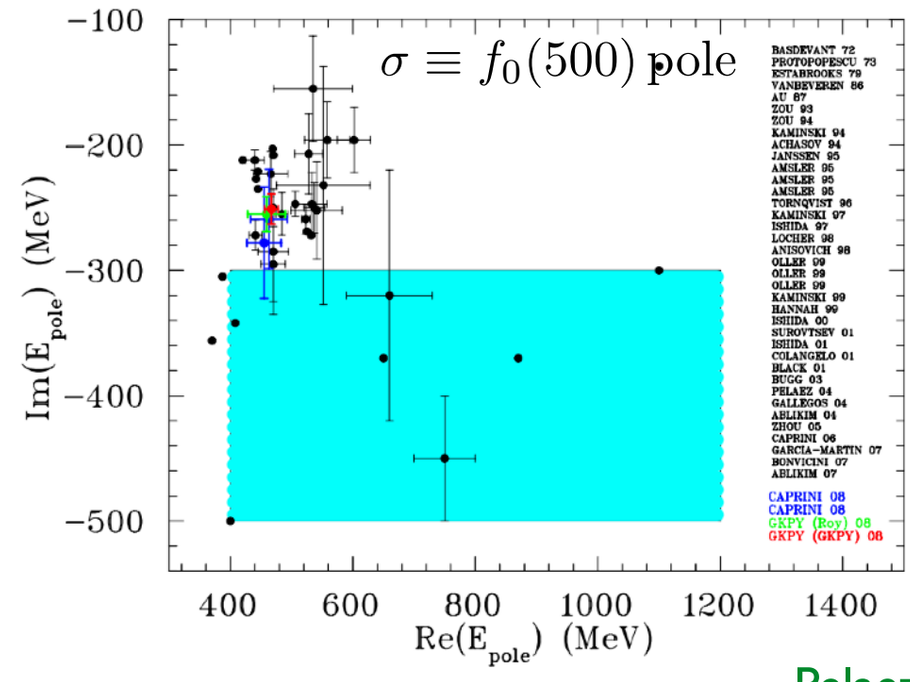

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$$U(s) = \sum_k C_k\,\xi(s)^k$$

I.D., Lutz, Gasparyan (2011)

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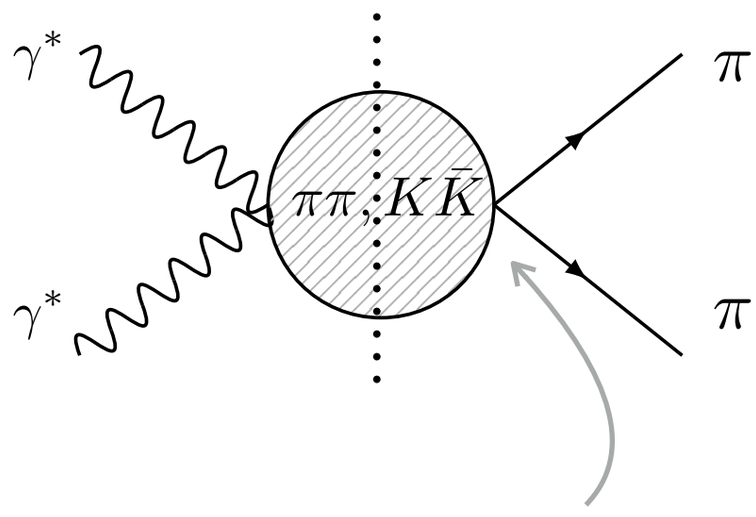

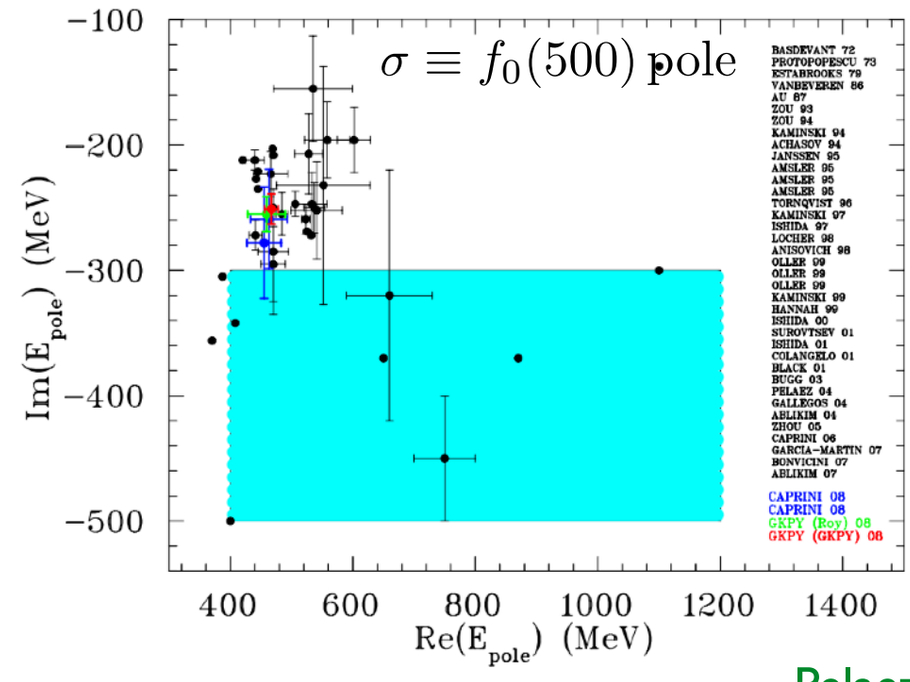

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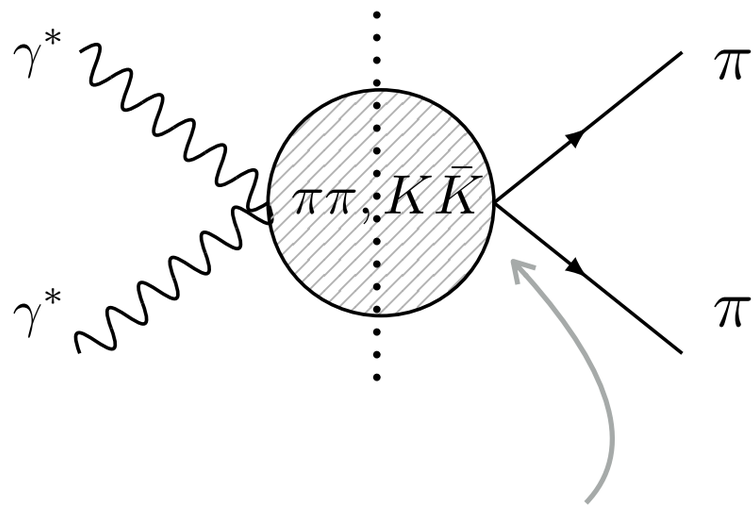

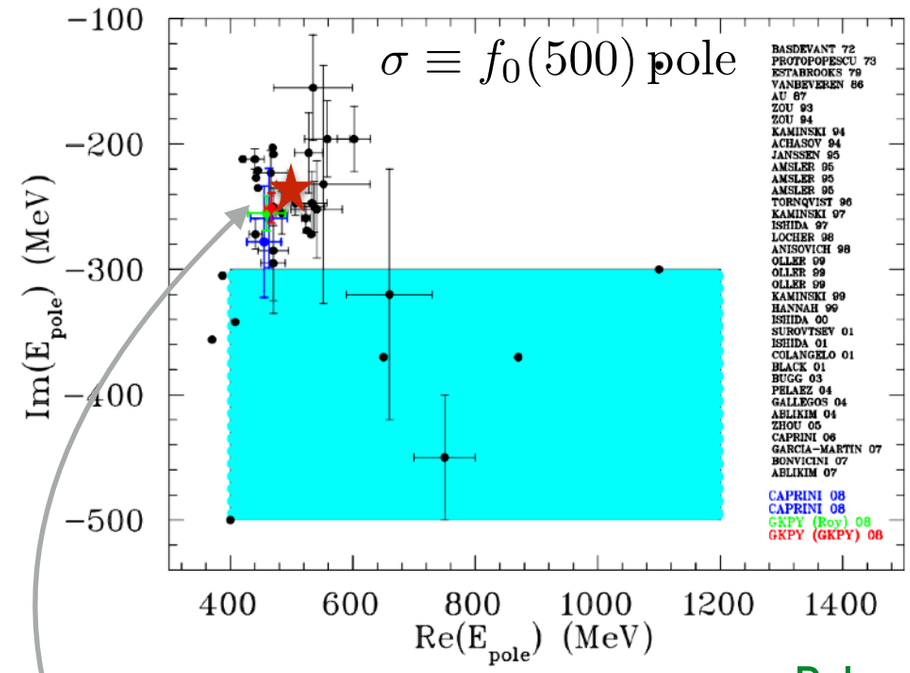

Pelaez (2016)

$$\sqrt{s^{\mathrm{II}}_{\mathrm{Roy}}} = \left(449^{+22}_{-16}\right) \pm i\,(275 \pm 12)\ \mathrm{MeV}$$

$$\sqrt{s^{\mathrm{II}}_{N/D}} = (496 \pm 48) \pm i\,(226 \pm 30)\ \mathrm{MeV}$$

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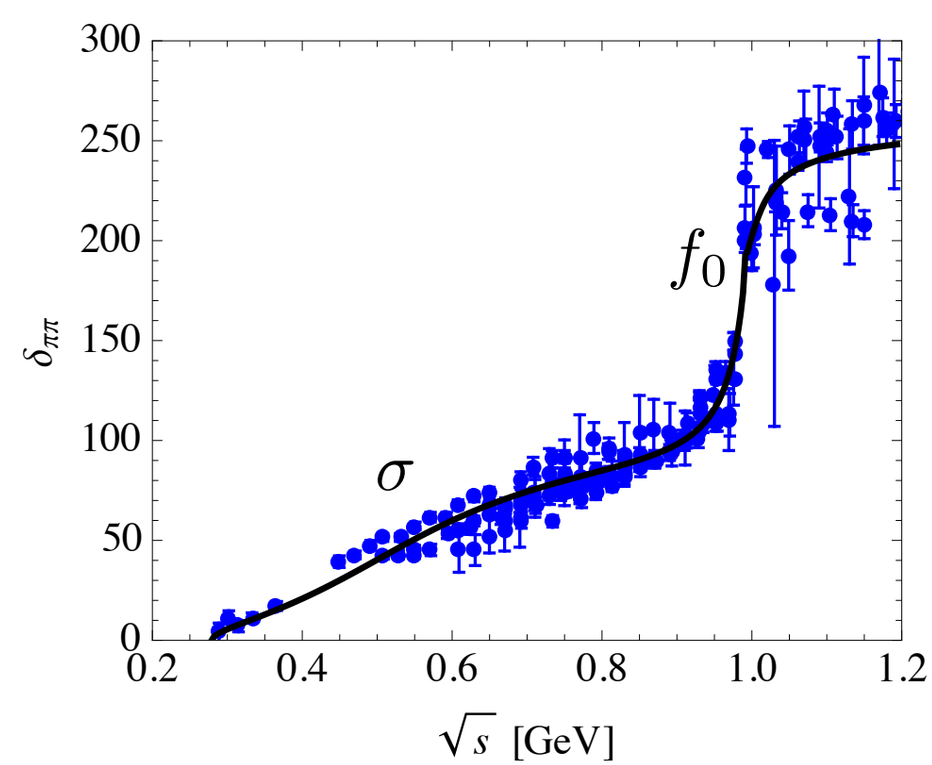

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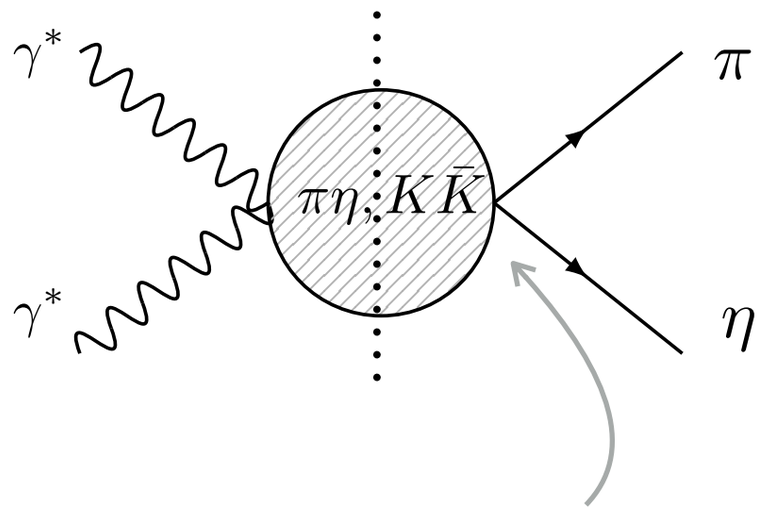

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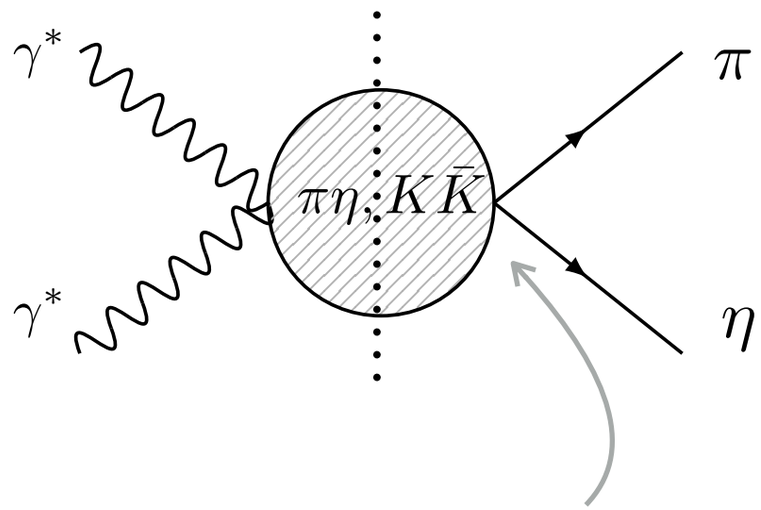

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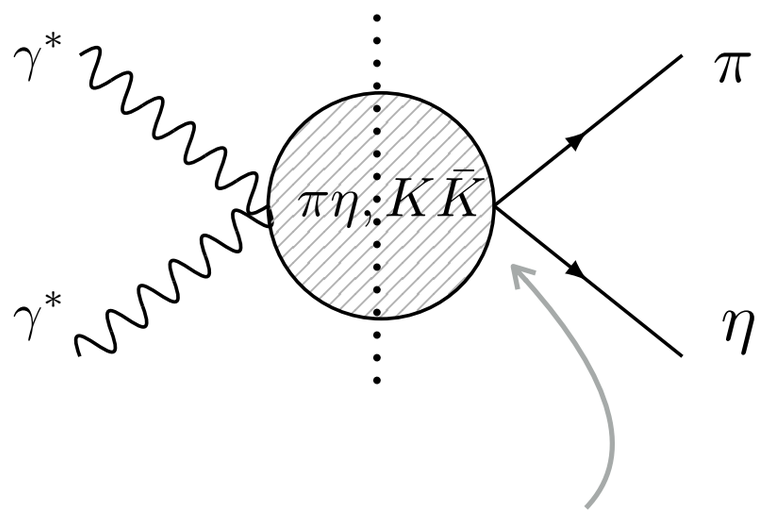

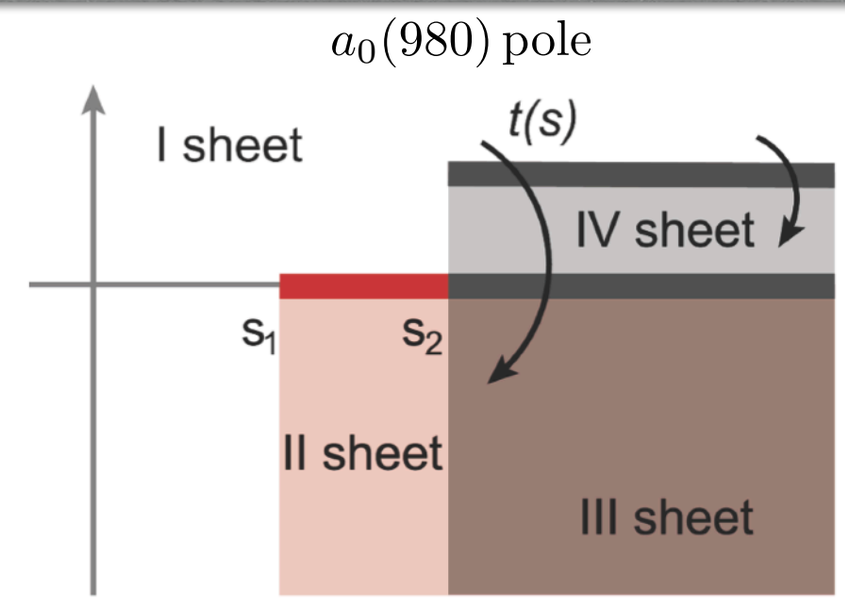

$$\sqrt{s_{a_0}^{\rm IV}} = \left(1.12^{-0.07}_{+0.02}\right) \pm \frac{i}{2}\left(0.28^{+0.08}_{-0.13}\right)\ \text{GeV}$$

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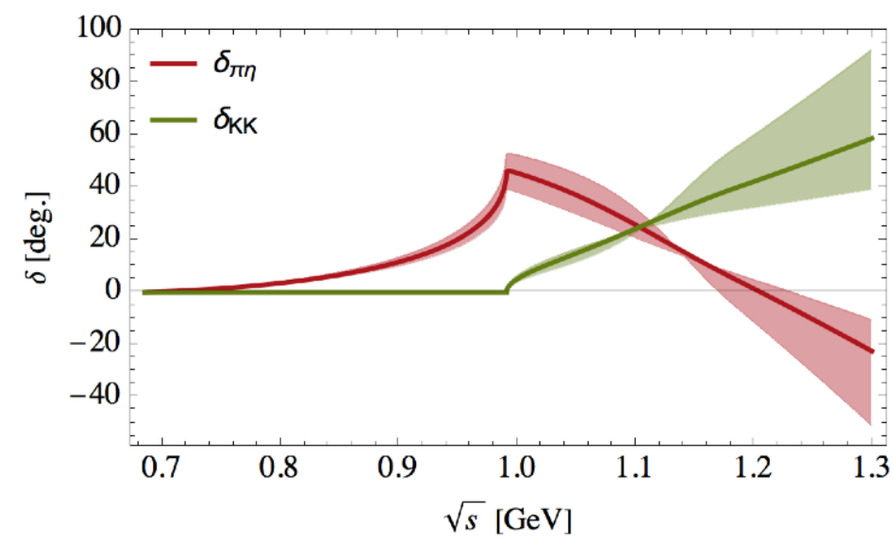
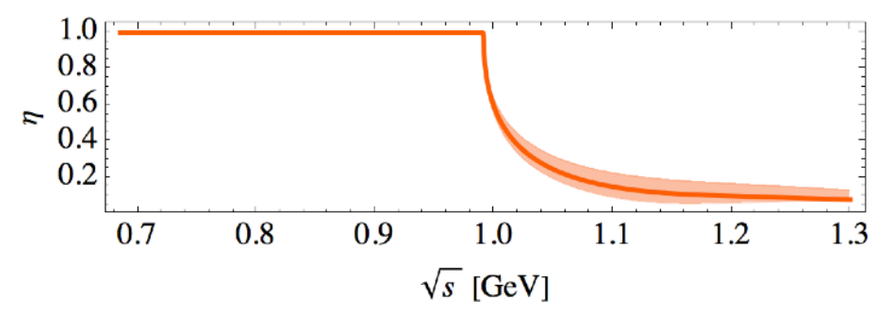
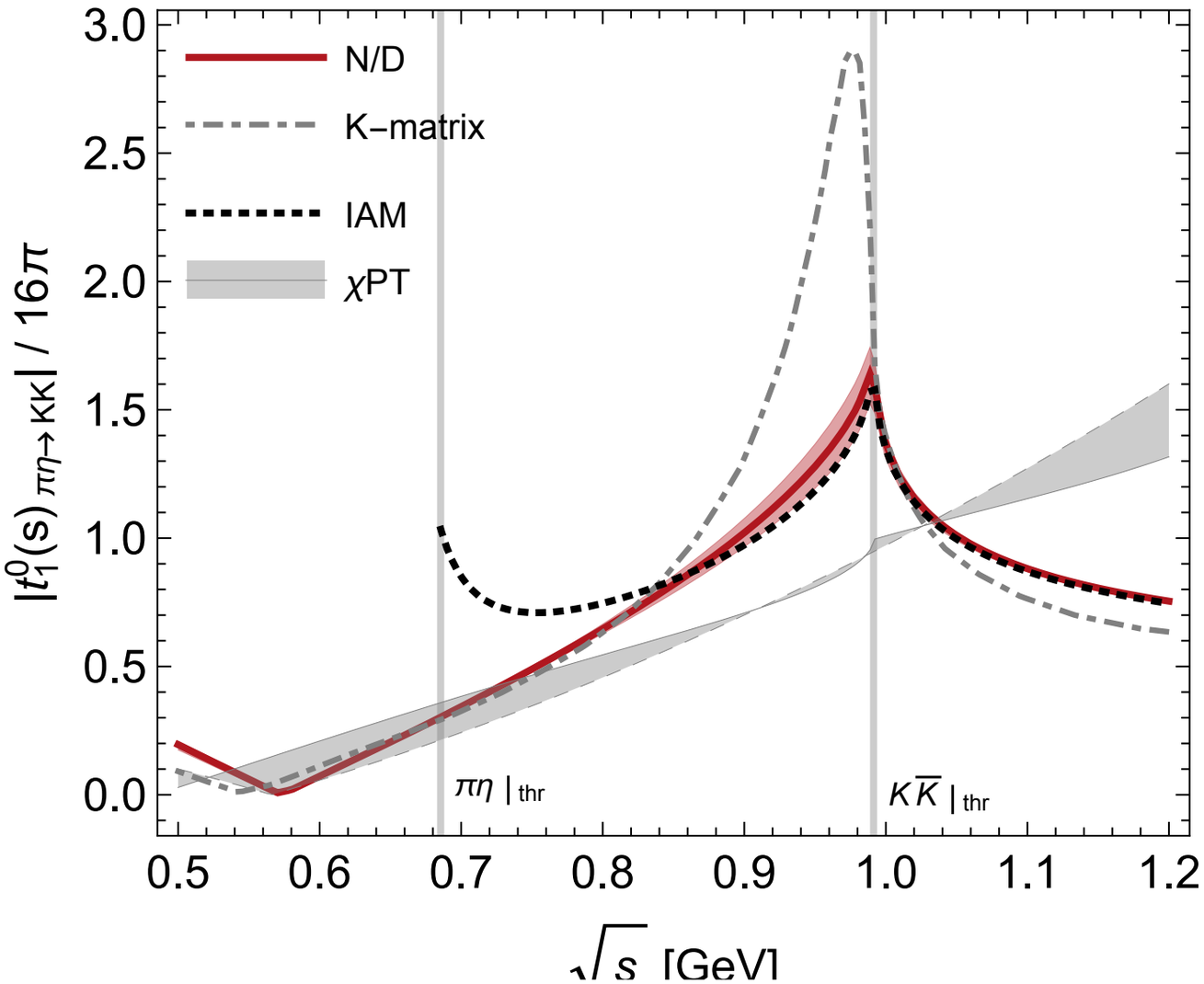




cf. also HadSpec Coll. (2016)

Scattering amplitude $\pi\eta \to KK$



N/D
I.D., Gil, Lutz (2013), I.D., Deineka, Vanderhaeghen (2017)

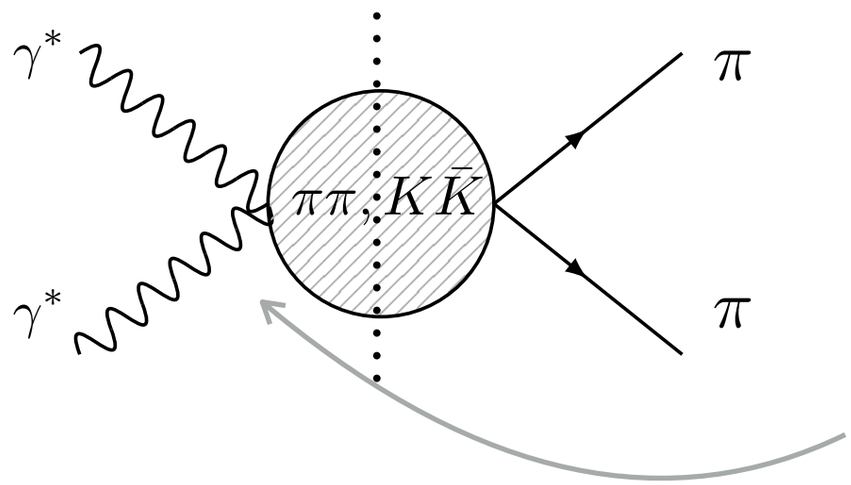
K-matrix
Albaladejo et. al. (2017)

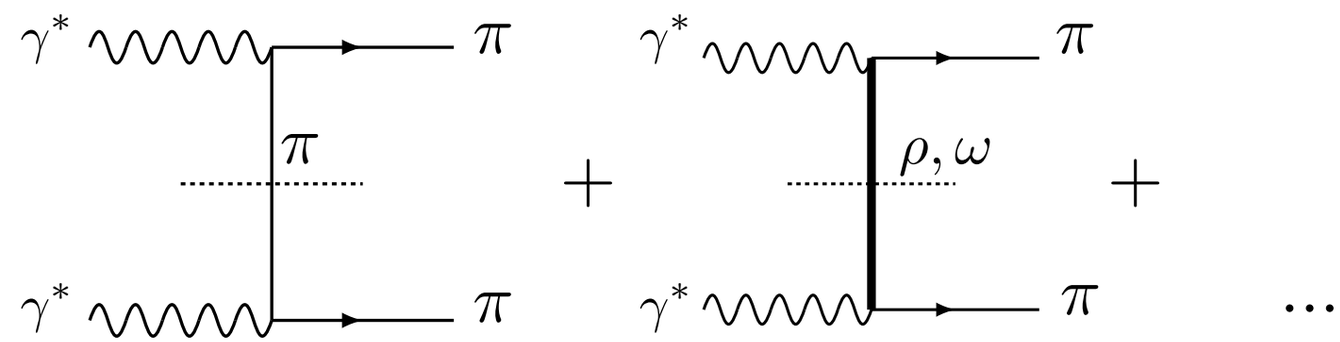
Inverse Amplitude Method (IAM)
Gomez Nicola et.al. (2002)

Chiral Perturbation Theory
Gasser et. al. (1985)

- First lattice analysis for m_π=391 MeV HadSpec Coll. (2016)
- Chiral extrapolation of the lattice results Zhi-Hui Guo et. al. (2017)

Left-hand cuts (pion pole)





- Left-hand cuts requires knowledge from $\gamma^*\pi\pi$, $\gamma^*\pi\omega$, $\gamma^*\pi\rho$ transition form factors

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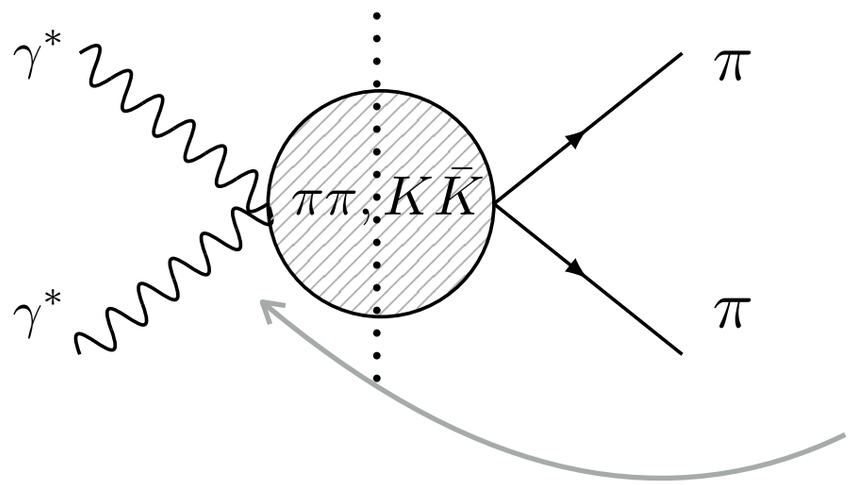

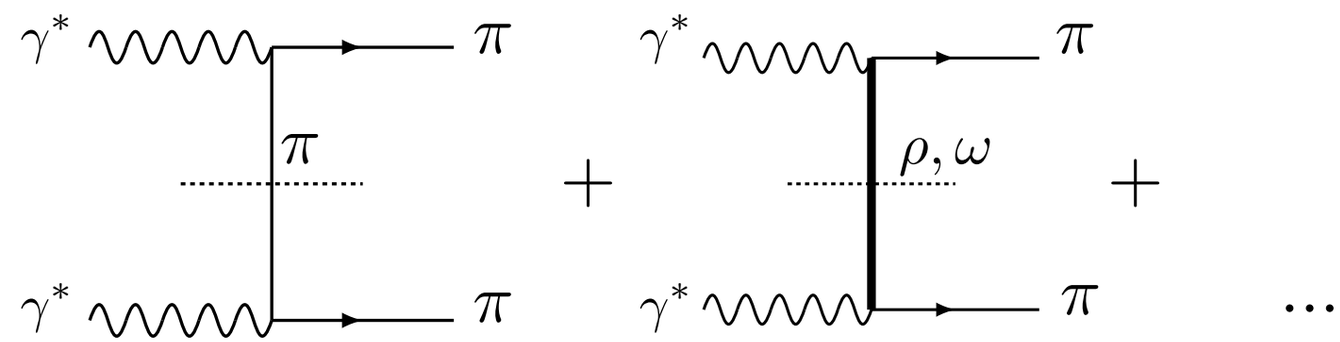

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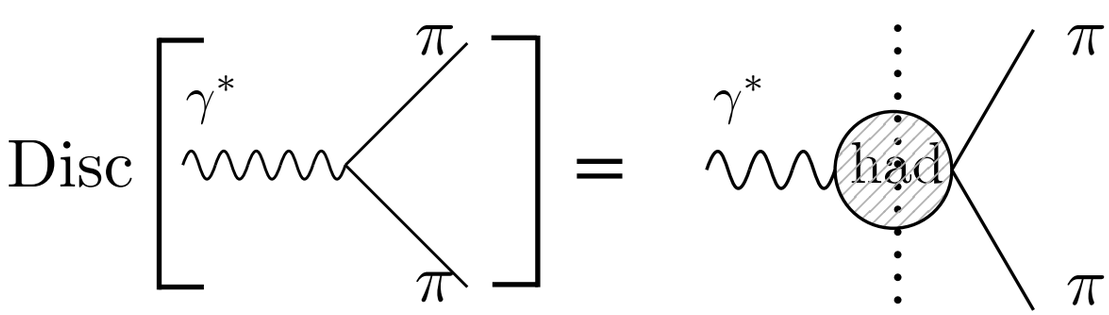
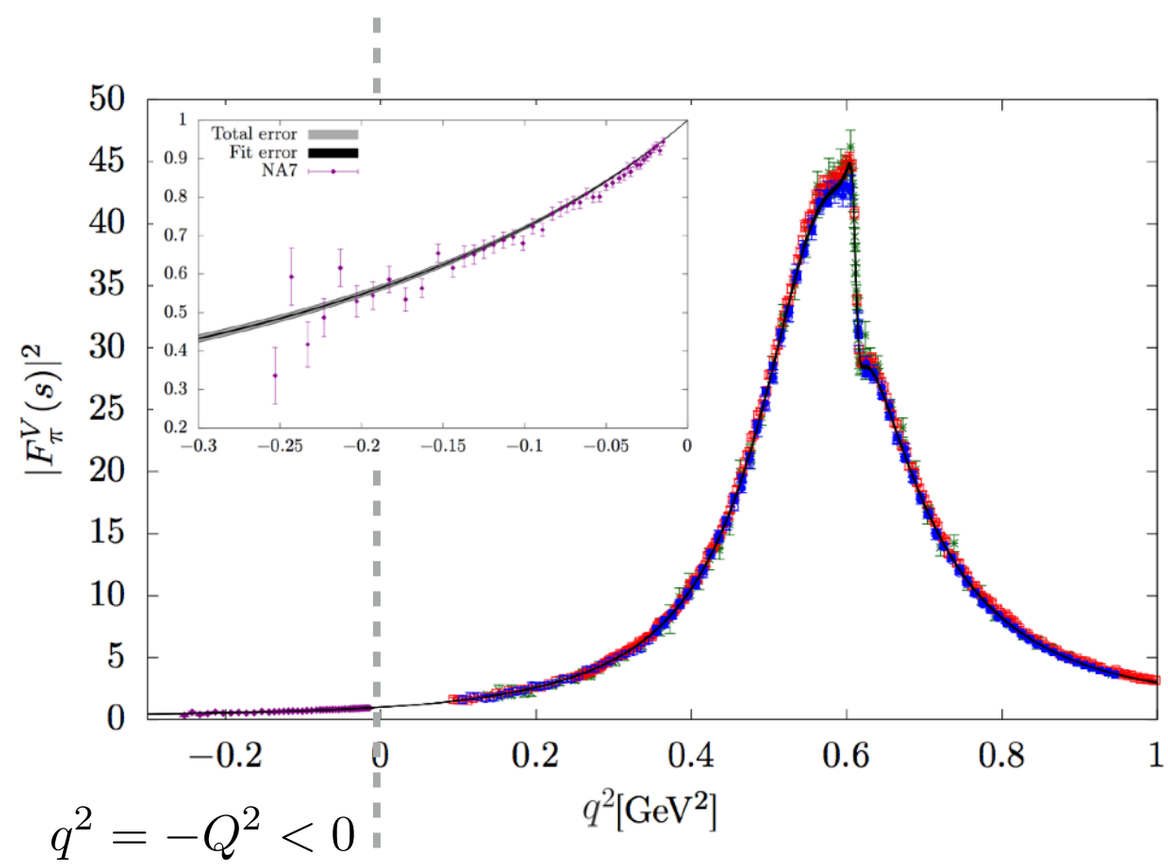
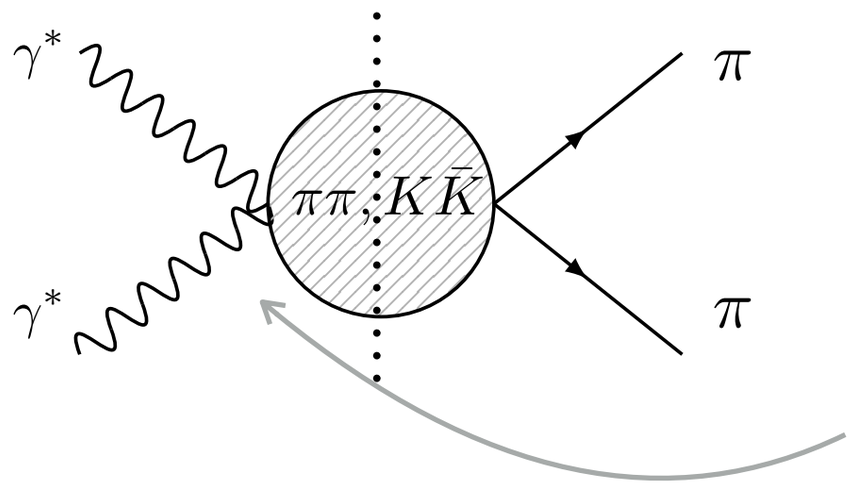


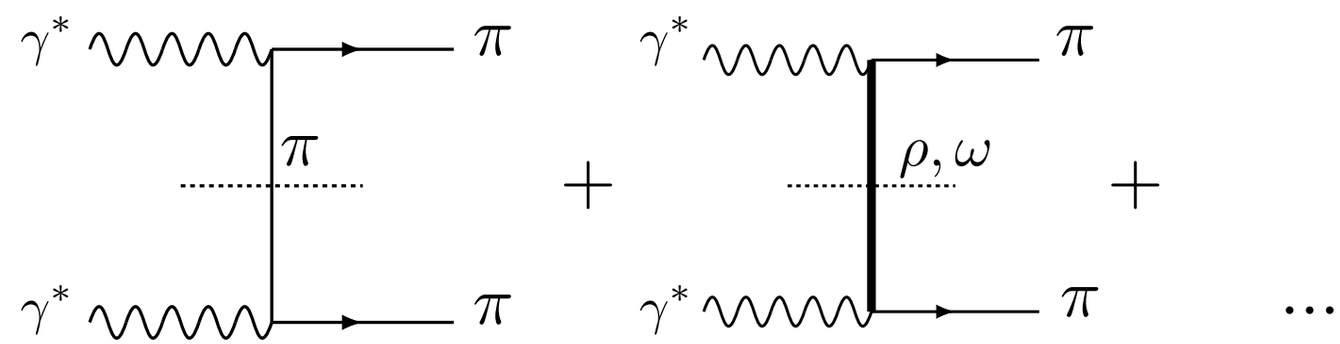
- Input from data: monopole TFF works well for space like region



Colangelo et al. (2019)

Left-hand cuts (vector poles)





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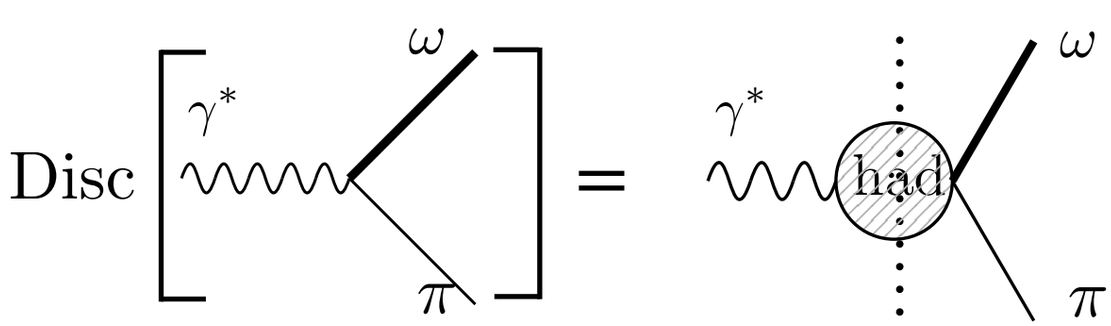
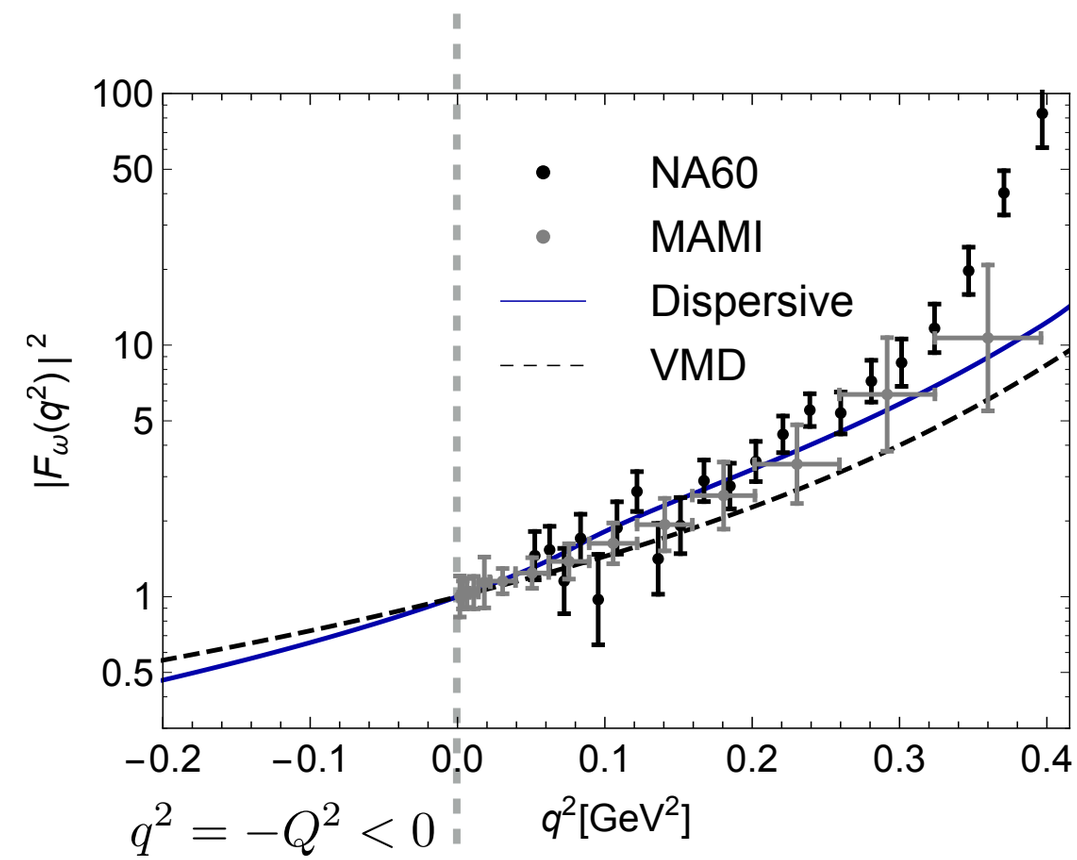
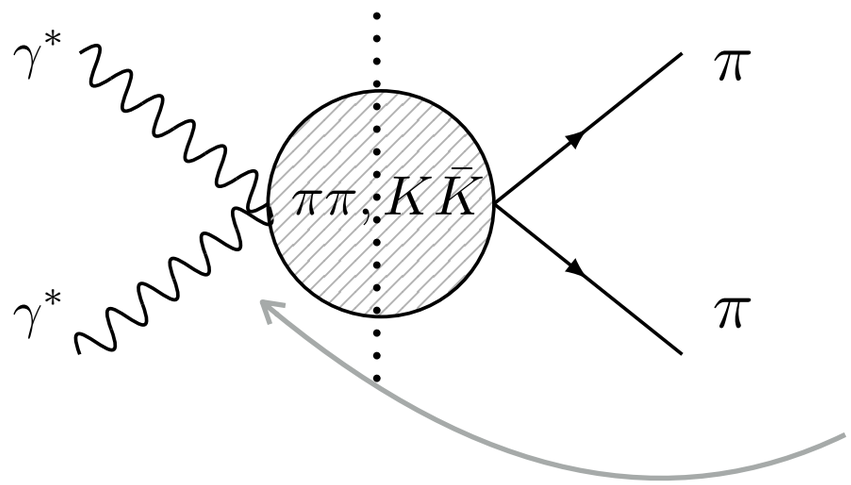


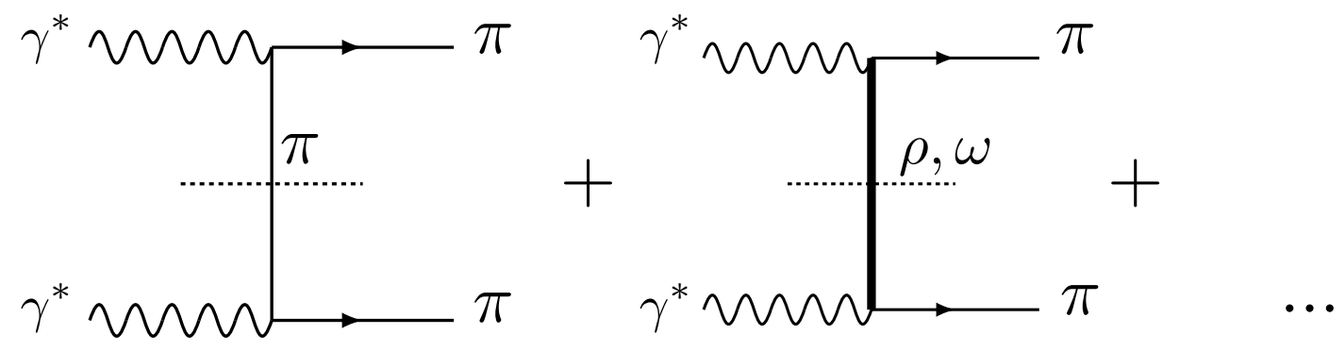


- Input from: $\omega \to 3\pi$

Khuri-Treiman (1960)
Schneider et al. (2012)
I.D. & JPAC (2015)

Left-hand cuts (vector poles)





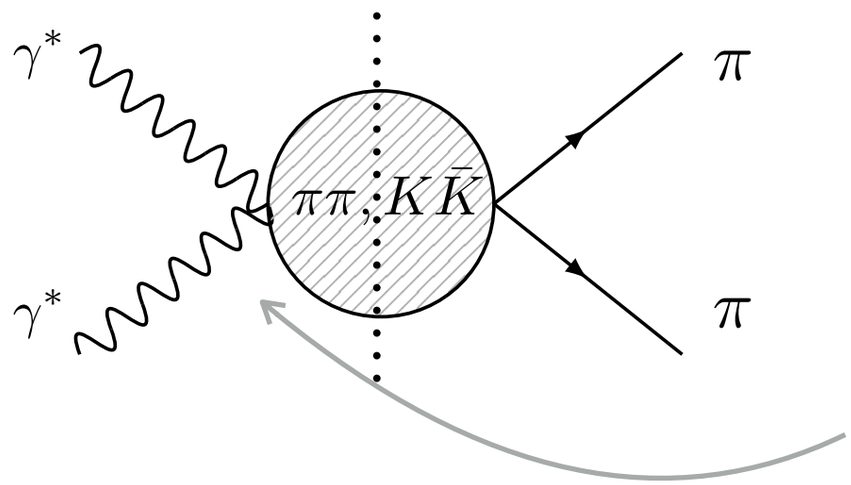
- Left-hand cuts requires knowledge from $\gamma^*\pi\pi$, $\gamma^*\pi\omega$, $\gamma^*\pi\rho$ transition form factors

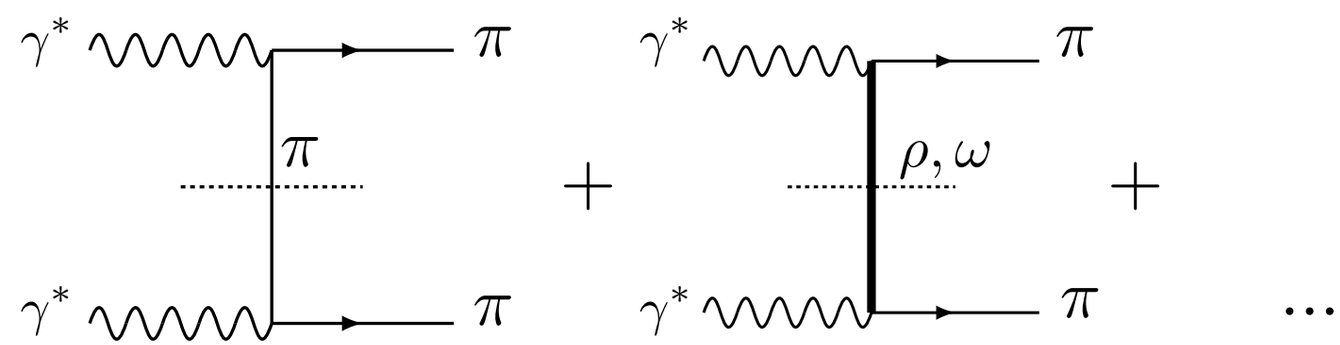
- Fitted parameter is the coupling: $g_{V\to\pi\gamma} \simeq C_{\rho^{\pm,0}\to\pi^{\pm,0}\gamma} \simeq \frac{1}{3} C_{\omega\to\pi^0\gamma} \overset{\text{PDG}}{=} 0.37(2)\ \text{GeV}^{-1}$

  $g_{V\to\pi\gamma} = 0.33\ \text{GeV}^{-1}$

I.D., Vanderhaeghen (2018)

Left-hand cuts (vector poles)





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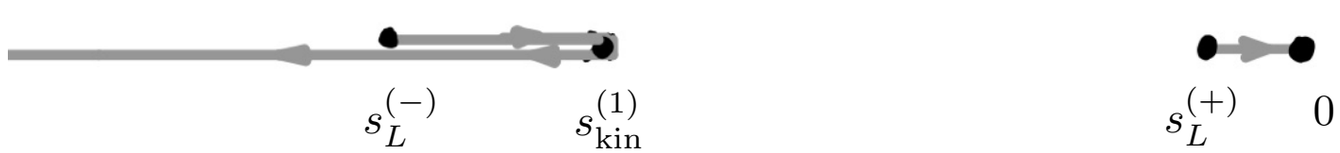
  $g_{V\to\pi\gamma} = 0.33\ \text{GeV}^{-1}$

  I.D., Vanderhaeghen (2018)

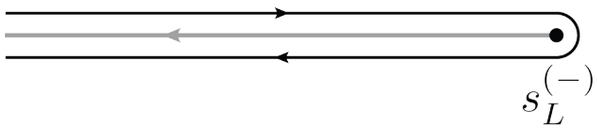
- Left-hand cuts: "anomalous thresholds" for large virtualities $Q_1^2 Q_2^2 > (M_V^2 - m_\pi^2)^2$

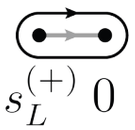
  Hoferichter, Stoffer (2019)
  Mandelstam (1960)



Left-hand cuts (vector poles)

- Normal case





Left-hand cuts (vector poles)

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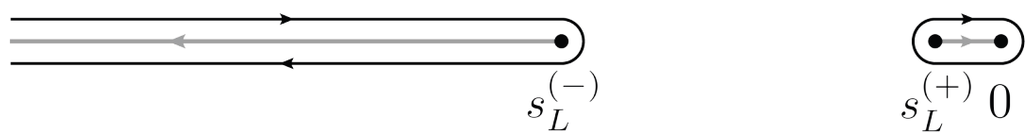
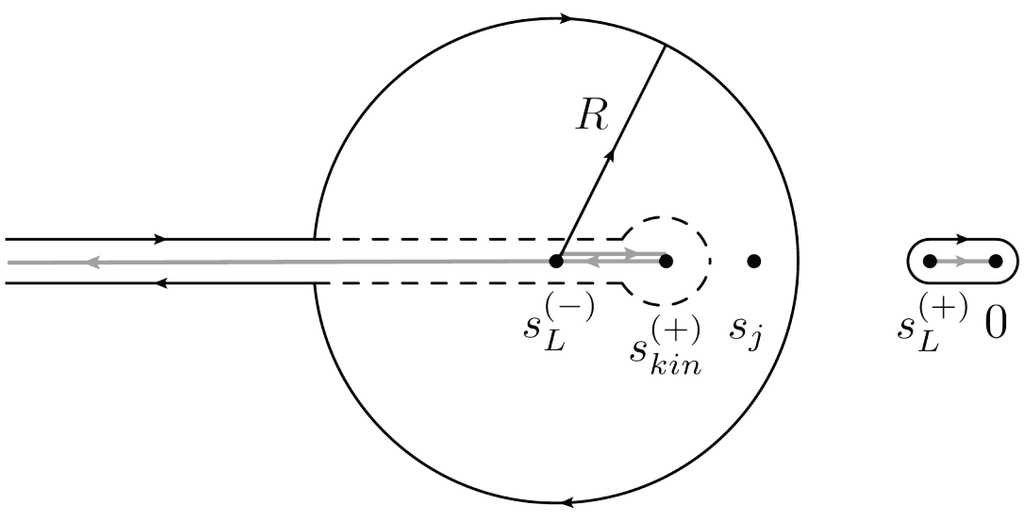


- Anomaly case requires contour deformation



$$h^V(s) \sim \frac{1}{(s - s_{\rm kin}^{(+)})^{9/2}} \log\left(\frac{X+1}{X-1}\right)$$ for J=2

- Dashed curve: cancellations of singular pieces
  It requires a careful numerical implementation

Hoferichter, Stoffer (2019)

- Solid curve: enlarged contour such that one stays away from possible numerical issues related to the anomaly piece

I.D., Deineka, Vanderhaeghen (2019)

  - Easy to implement
  - Independent on the degree of singularity
  - Generalisation to the physical case is straightforward

Kinematic constraints

- Helicity amplitudes

$$H_{\lambda_1,\lambda_2} = \epsilon_\mu(\lambda_1)\epsilon_\nu(\lambda_2) \sum_{n=1}^{5} F_n(s,t)\, L_n^{\mu\nu}$$

Bardeen et al. (1968), Tarrach (1975)
Metz et al. (1998), Colangelo et al. (2015)

where $\lambda_{1,2} = \pm 1, 0$ are photon helicities
(minimal basis for Born subtracted amplitudes)

Low et al. (1954)

- p.w. helicity amplitudes suffer from kinematic constraints

$$h^{(J)}_{\lambda_1\lambda_2} = \int \frac{d\cos\theta}{2}\, d^J_{\lambda_1-\lambda_2,0}(\theta)\, H_{\lambda_1\lambda_2}$$

$$A^{(J)}_n = \frac{1}{(p\,q)^J} \int \frac{d\cos\theta}{2}\, P_J(\theta)\, F_n(s,t)$$ ← object free of kinematic constraints

Lutz et al. (2010, 2014)

- Unconstrained basis for Born subtracted p.w. amplitudes

$$\bar{h}^{(J)}_i \equiv h^{(J)}_i - h^{(J),\mathrm{Born}}_i$$

$$\bar{h}^{(J)}_i = K_{ij}\, \bar{h}^{(J)}_j \qquad j \equiv \lambda_1\lambda_2 = \{++,\, +-,\, +0,\, 0+,\, 00\}$$

K_{ij} is 5×5 matrix

Kinematic constraints

- For s-wave

$$\bar{h}^{(0)}_{++} \pm \bar{h}^{(0)}_{00} \sim (s - s^{(\mp)}_{\rm kin}), \quad s^{(\pm)}_{\rm kin} \equiv -(Q_1 \pm Q_2)^2$$

Colangelo et al. (2017)
Pennington (1988), Moussallam (2013)

- For d-wave

$$(s + Q_1^2 + Q_2^2)\,\bar{h}^{(2)}_{+-} + \frac{2\sqrt{2\,s}\,Q_1^2 Q_2^2}{Q_1^2 - Q_2^2}\left(\frac{\bar{h}^{(2)}_{+0}}{Q_2} - \frac{\bar{h}^{(2)}_{0+}}{Q_1}\right) \sim (s - 4m_\pi^2)\left(s - s^{(+)}_{\rm kin}\right)\left(s - s^{(-)}_{\rm kin}\right)$$

$$\bar{h}^{(2)}_{+-} + \frac{\sqrt{2\,s}}{Q_1^2 - Q_2^2}\left(Q_2\,\bar{h}^{(2)}_{+0} - Q_1\,\bar{h}^{(2)}_{0+}\right) \sim (s - 4m_\pi^2)\left(s - s^{(+)}_{\rm kin}\right)\left(s - s^{(-)}_{\rm kin}\right)$$

$$\sqrt{2}\,\bar{h}^{(2)}_{+-} + \frac{\left(Q_1^2 + Q_2^2 + s\right)\sqrt{s}}{Q_1^2 - Q_2^2}\left(\frac{\bar{h}^{(2)}_{+0}}{Q_2} - \frac{\bar{h}^{(2)}_{0+}}{Q_1}\right) \sim (s - 4m_\pi^2)\left(s - s^{(+)}_{\rm kin}\right)\left(s - s^{(-)}_{\rm kin}\right)$$

+ 2 more

- Unconstrained basis for Born subtracted p.w. amplitudes

I.D., Deineka, Vanderhaeghen (2019)
cf. also Hoferichter, Stoffer (2019)

$$\bar{h}^{(J)}_i \equiv h^{(J)}_i - h^{(J),\rm Born}_i$$

$$\bar{h}^{(J)}_i = K_{ij}\,\bar{h}^{(J)}_j \qquad j \equiv \lambda_1\lambda_2 = \{++,\ +-,\ +0,\ 0+,\ 00\}$$

K_{ij} is 5×5 matrix

Dispersion relation

- <u>Unsubtracted</u> dispersion relation for kinematically unconstrained p.w. amplitudes

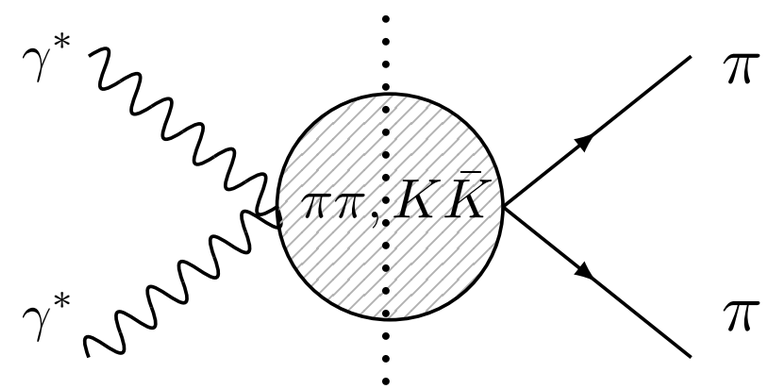
$$\bar{h}_i^{(J)} = \int_{-\infty}^{0} \frac{ds'}{\pi} \frac{\mathrm{Disc}\,\bar{h}_i^{(J)}(s')}{s'-s} + \int_{4m_\pi^2}^{\infty} \frac{ds'}{\pi} \frac{\mathrm{Disc}\,\bar{h}_i^{(J)}(s')}{s'-s}$$

$$\bar{h}_i^{(J)} \equiv h_i^{(J)} - h_i^{(J),\mathrm{Born}}$$

Garcia-Martin et. al (2010)
Hoferichter et. al. (2011,19)
Dai et al. (2014)
Moussallam (2013)

- Omnès solution of the unitarity relation

$$\mathrm{Disc}\, h_i^{(J)} = h_i^{(J)}\ \rho\ t_{\pi\pi}^{(J)*}$$
$$\mathrm{Disc}\, \Omega^{(J)} = \Omega^{(J)}\ \rho\ t_{\pi\pi}^{(J)*} \quad |_{s>4m_\pi^2}$$



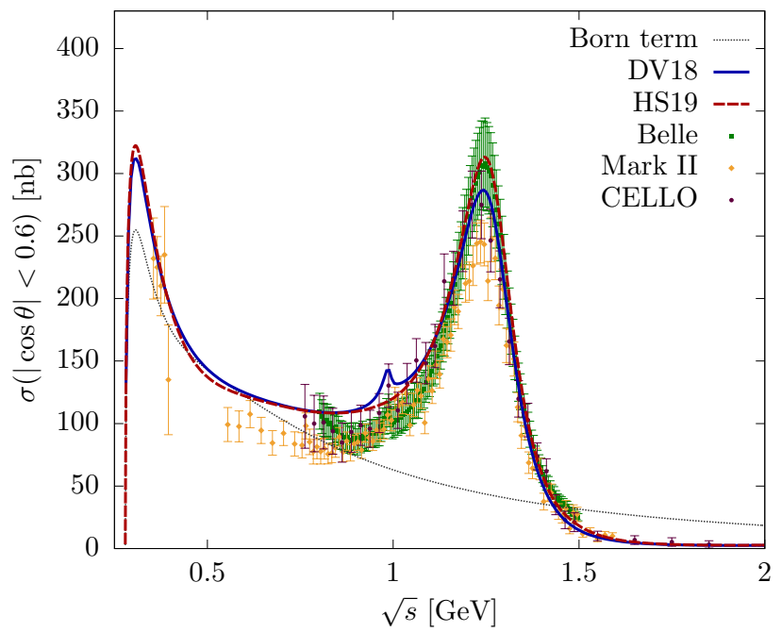
leads to

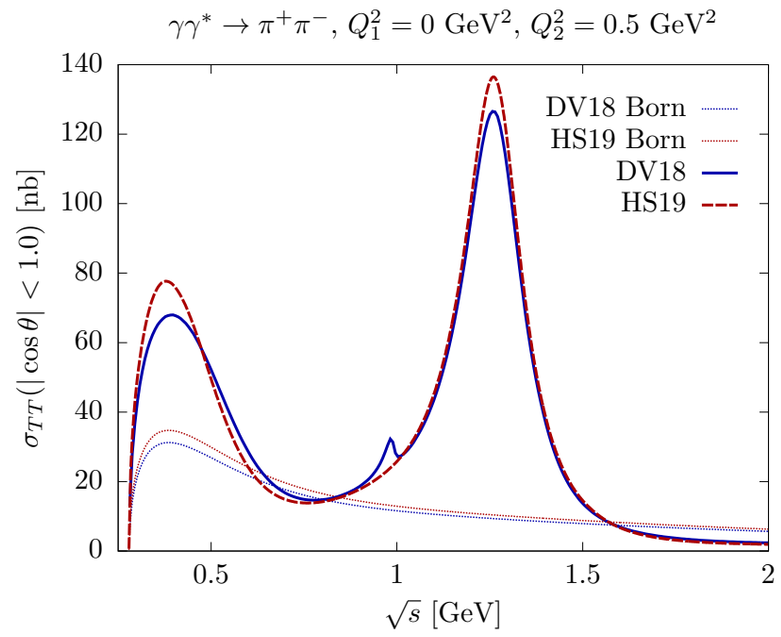
$$h_i^{(J)} = h_i^{(J),\mathrm{Born}} + \Omega^{(J)} \left(\int_{-\infty}^{0} \frac{ds'}{\pi} \frac{\mathrm{Disc}(\bar{h}_i^{(J)}(s'))\,\Omega^{(J)}(s')^{-1}}{s'-s} - \int_{4m_\pi^2}^{\infty} \frac{ds'}{\pi} \frac{h_i^{(J),\mathrm{Born}}(s')\,\mathrm{Im}\,\Omega^{(J)}(s')^{-1})}{s'-s} \right)$$

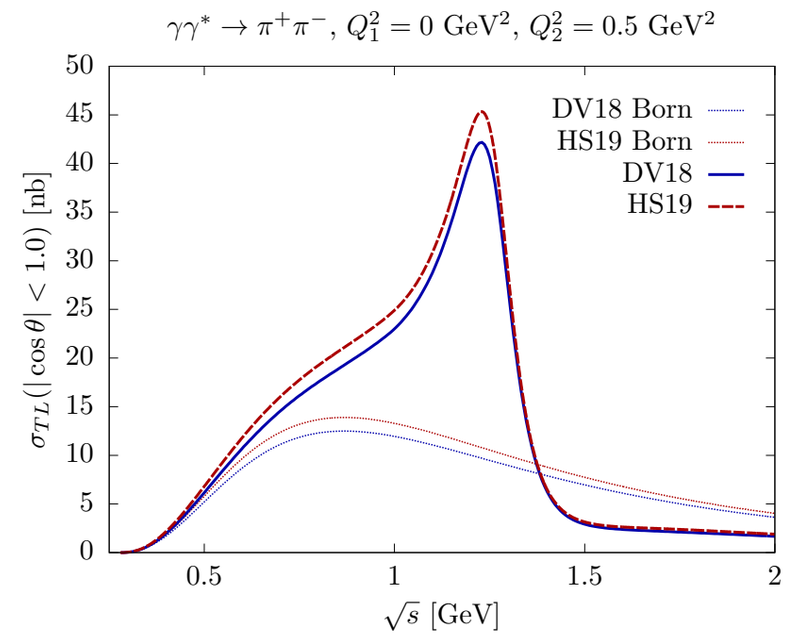
V-exch

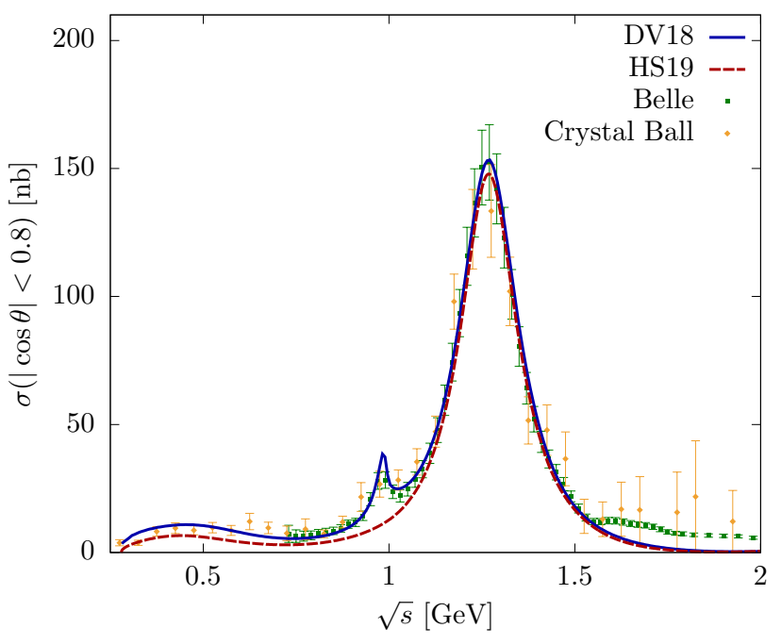
Omnès (1958)
Muskhelishvili (1953)

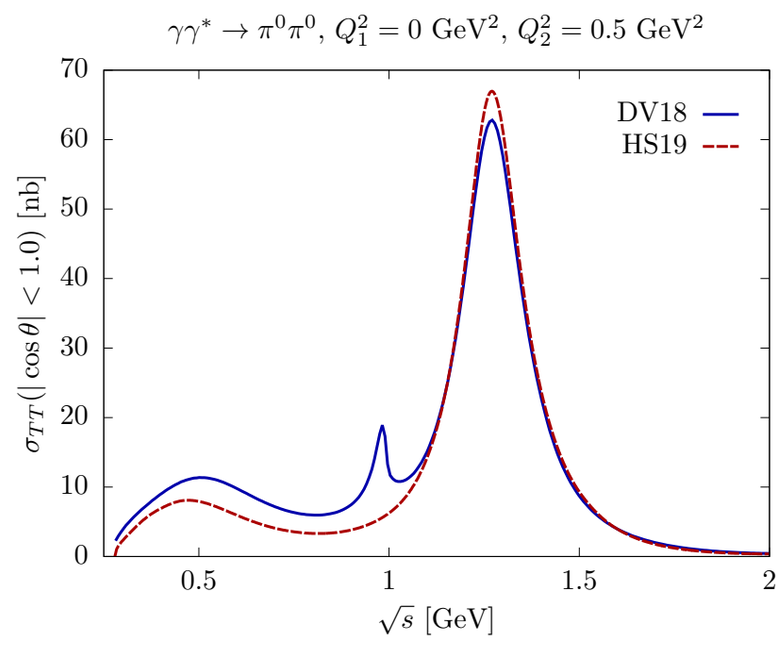
Results for ππ

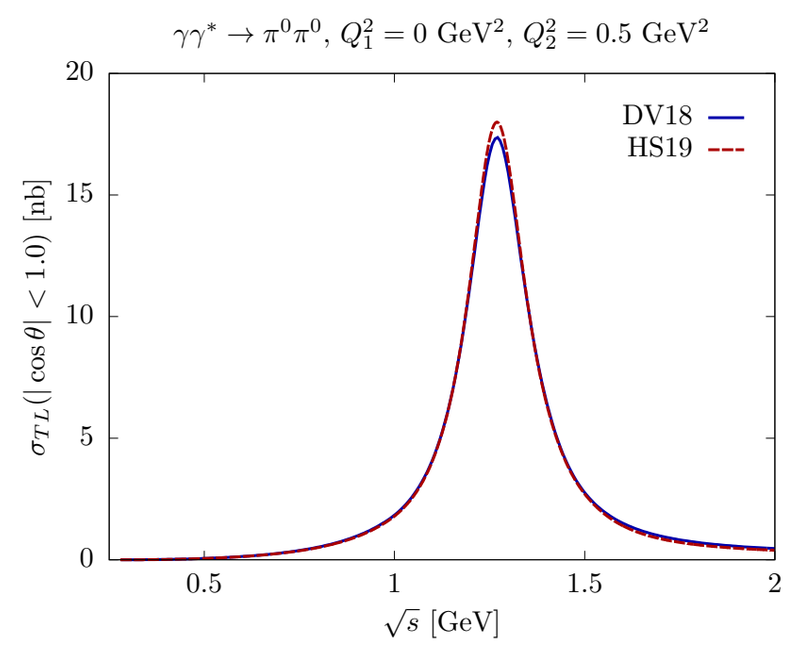










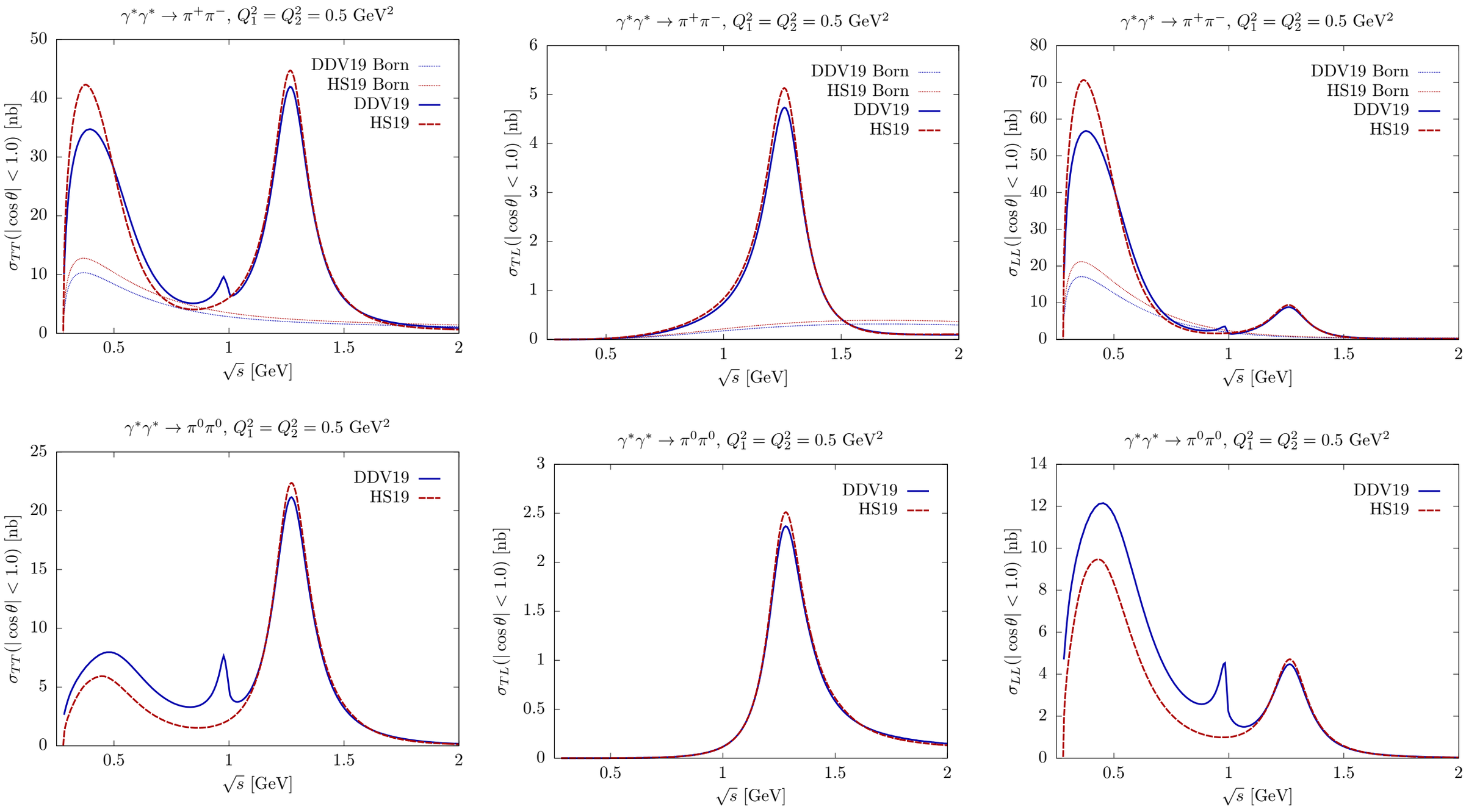


- Coupled-channel dispersive treatment of $f_0(980)$
- $f_2(1270)$ described dispersively through Omnès function

I.D., Vanderhaeghen (2018)
Hoferichter, Stoffer (2019)

cf. also Dai et al. (2014)
Garcia-Martin et al (2010)
Moussallam (2013)

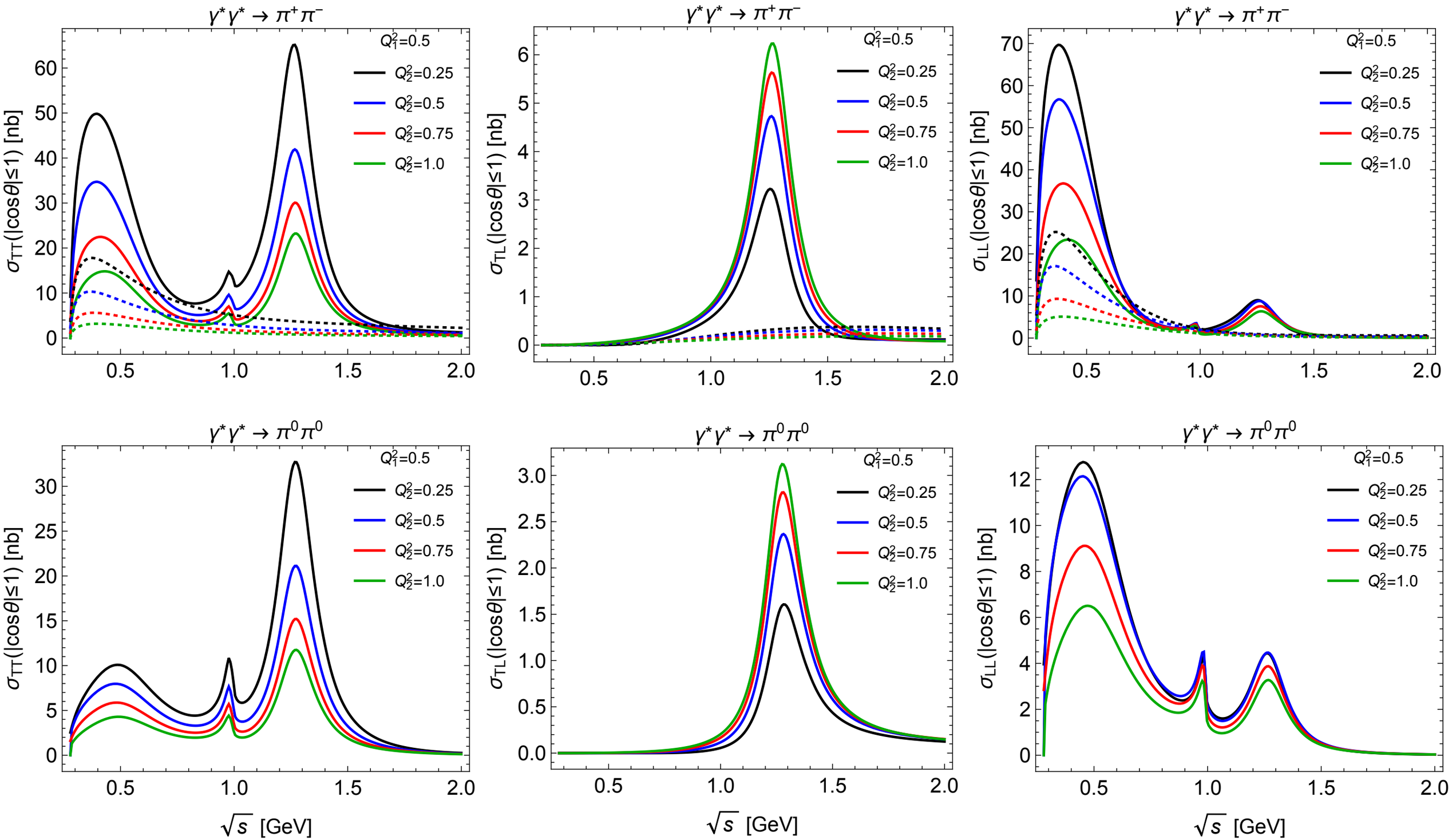
Results for ππ



$$\frac{d\sigma_{TT}}{d\cos\theta} \sim |H_{++}|^2 + |H_{+-}|^2, \quad \frac{d\sigma_{TL}}{d\cos\theta} \sim |H_{+0}|^2, \quad \frac{d\sigma_{LL}}{d\cos\theta} \sim |H_{00}|^2$$

I.D., Deineka, Vanderhaeghen (2019)
Hoferichter, Stoffer (2019)

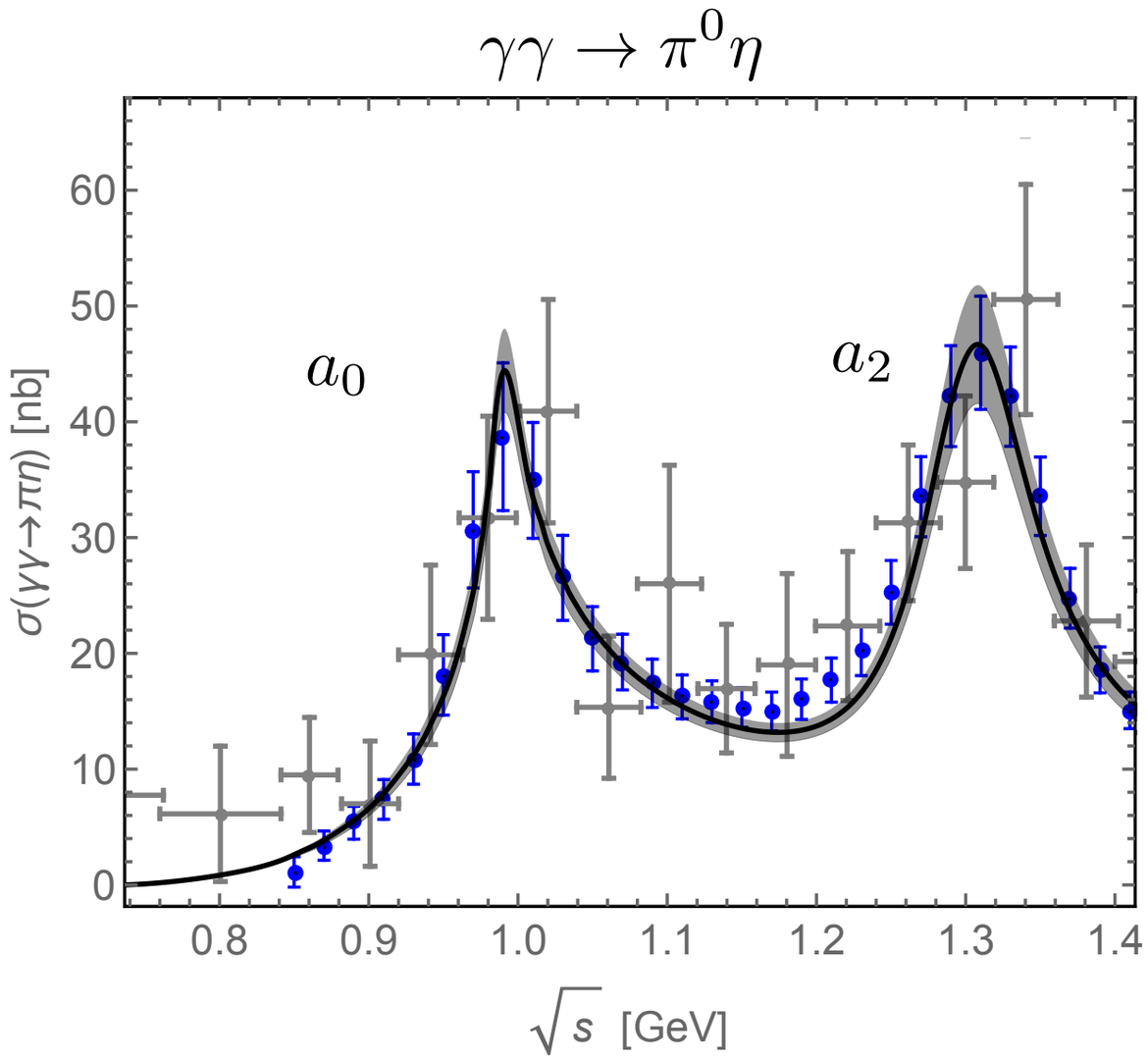
Results for ππ

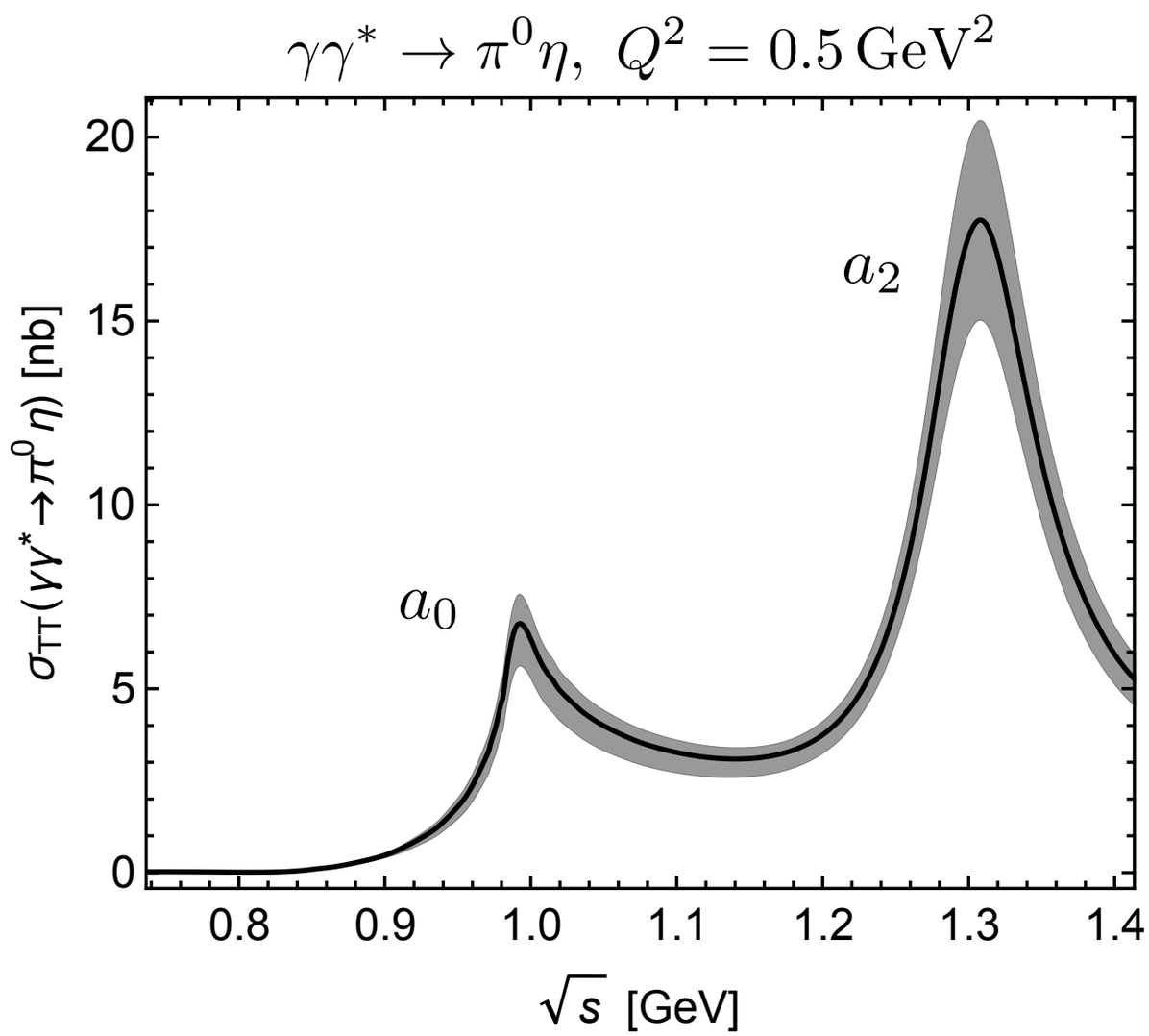


$$\frac{d\sigma_{TT}}{d\cos\theta} \sim |H_{++}|^2 + |H_{+-}|^2, \quad \frac{d\sigma_{TL}}{d\cos\theta} \sim |H_{+0}|^2, \quad \frac{d\sigma_{LL}}{d\cos\theta} \sim |H_{00}|^2$$

I.D., Deineka, Vanderhaeghen (2019)

Results for πη

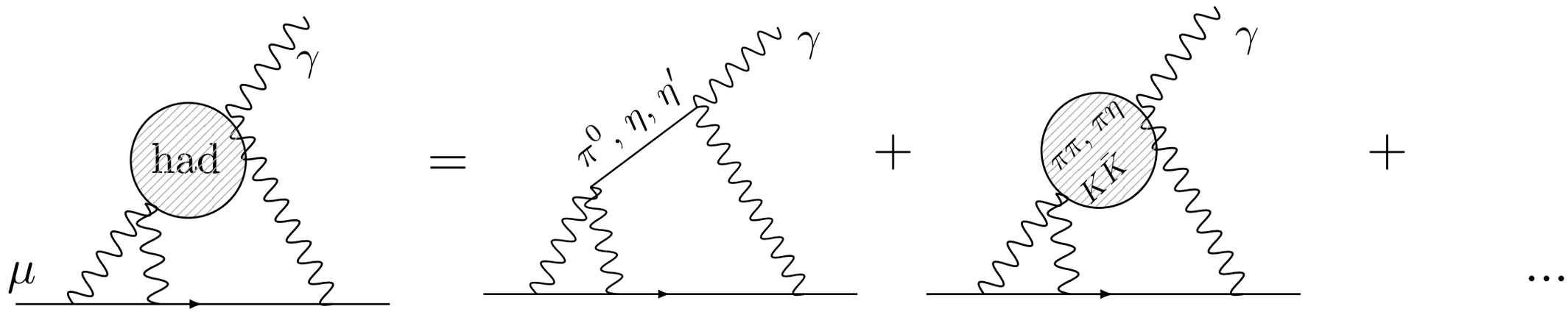




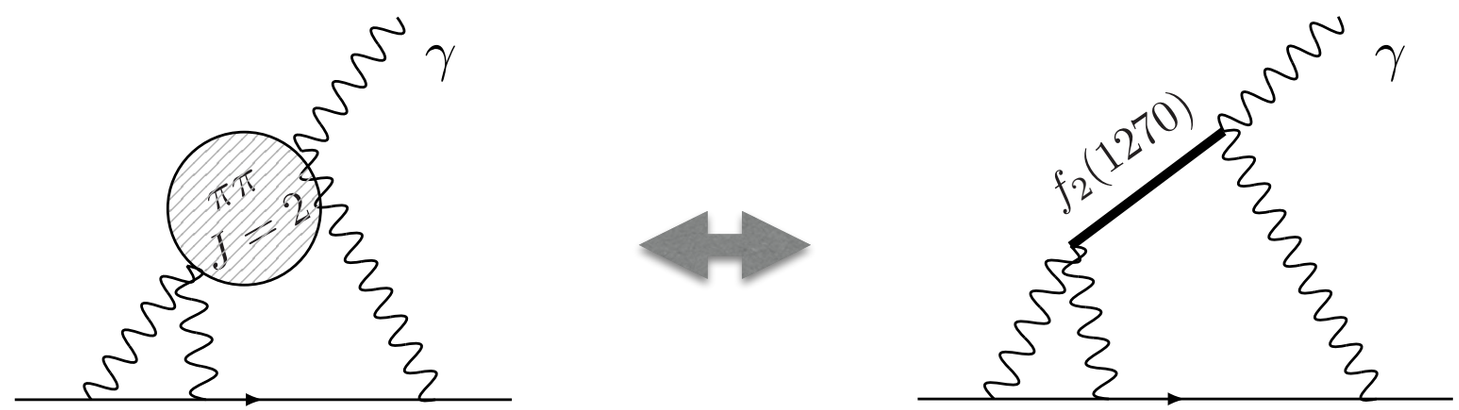
- Coupled-channel dispersive treatment of $a_0(980)$
- $a_2(1320)$ described as a Breit Wigner resonance
- $\gamma^*\gamma^*\to\pi\eta$ (work in progress)

I.D., Deineka, Vanderhaeghen (2017, 2018)

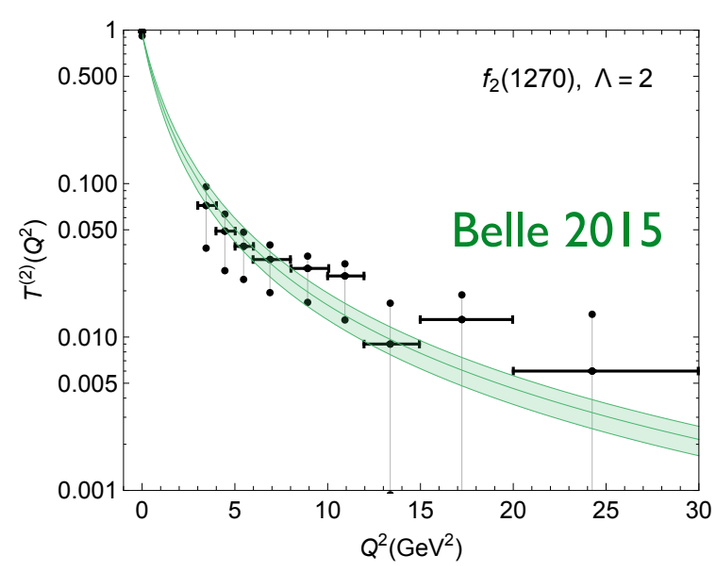
Effective resonance description



- One needs to compare



ongoing work



Pascalutsa, Pauk Vanderhaeghen (2012)

I.D., Vanderhaeghen (2016)

$$a_\mu = (0.05 \pm 0.01) \times 10^{-10}$$

Summary and Outlook

- new a_μ Fermilab and J-Parc experiments ongoing:
  aim: factor 4 improvement in experimental value

- Need to take into account $f_0(500)$, $f_0(980)$, $a_0(980)$, $f_2(1270)$ … and non resonant contributions in a dispersive approach to (g-2)

$$a_\mu^{\pi\text{-box}} + a_{\mu,J=0}^{\pi\pi,\pi\text{-pole LHC}} = -2.4(1) \times 10^{-10}$$

Colangelo et al. (2014-2017)

- Main ingredients: $\gamma^*\gamma^* \to \pi\pi$, $\pi\eta$ and (KK work in progress). Can be used in different (g-2) dispersive approaches.

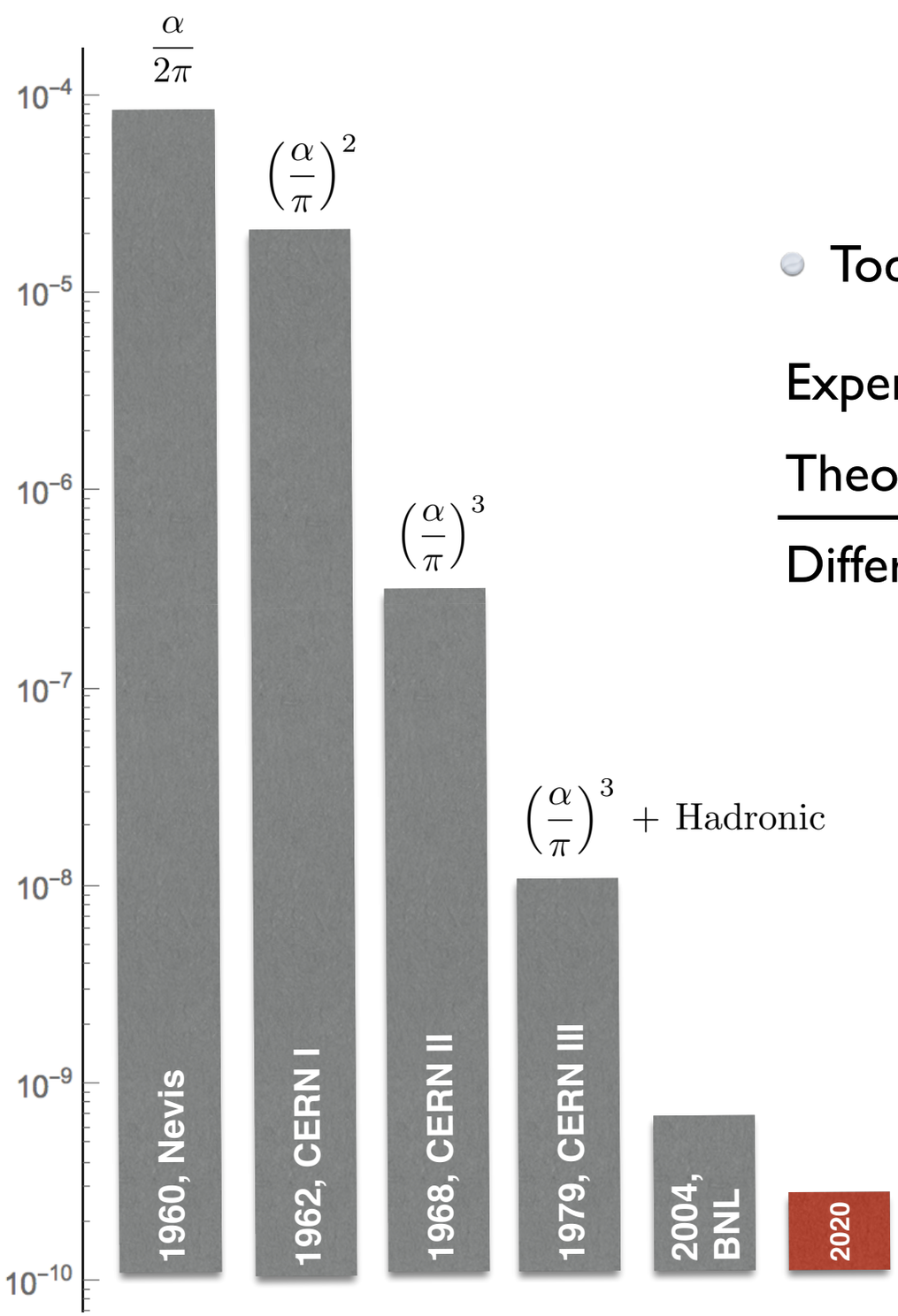
- Need to implement high energy asymptotic on TFF ($\gamma^*\pi\omega$, $\gamma^*\pi\rho$) used as input in dispersive formalism

- It is important to validate dispersive treatment of $\gamma\gamma^* \to \pi\pi$, $\pi\eta$, KK… with upcoming BESIII data

Thank you!

Motivation

- Accuracy



- Magnetic moment of the muon

$$\vec{\mu} = \frac{Q}{2m}\, g\, \vec{S}$$

- Anomalous part

$$a_\mu = \frac{(g-2)_\mu}{2}$$

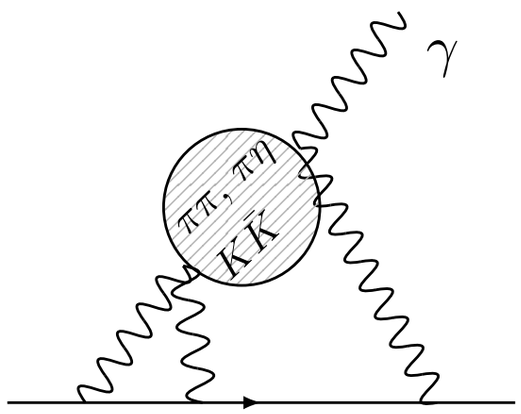
- Today

Experiment (BNL 2004)	$a_\mu^{\rm exp} = 0.0011659209(6)$
Theory (Standard model)	$a_\mu^{\rm th} = 0.0011659182(4)$
Difference	$\sim (3-4)\,\sigma$

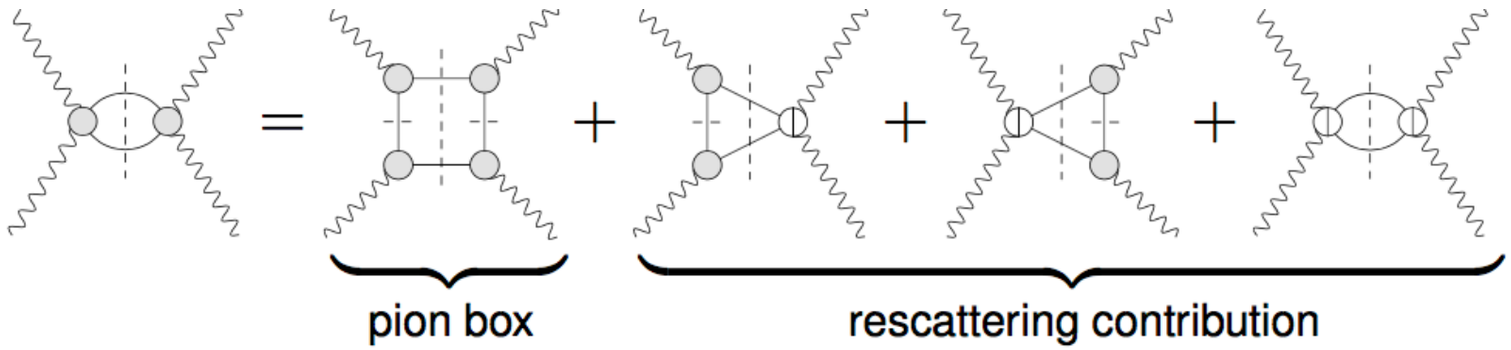
Keshavarzi et al. (2018)

$$\Delta a_\mu^{\rm exp} = 6.3 \times 10^{-10}$$
$$\rightarrow 1.6 \times 10^{-10}$$

Multi-meson contribution to (g-2)



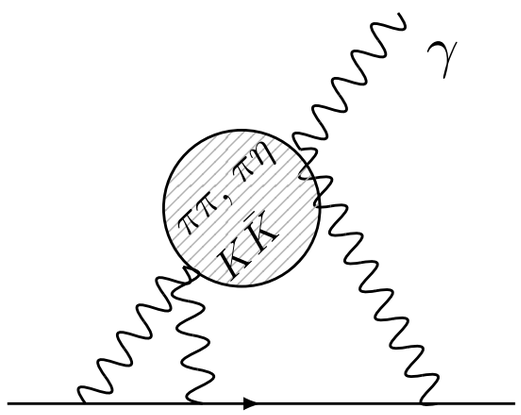
- Pioneering dispersive analyses for ππ loop contribution to a_μ



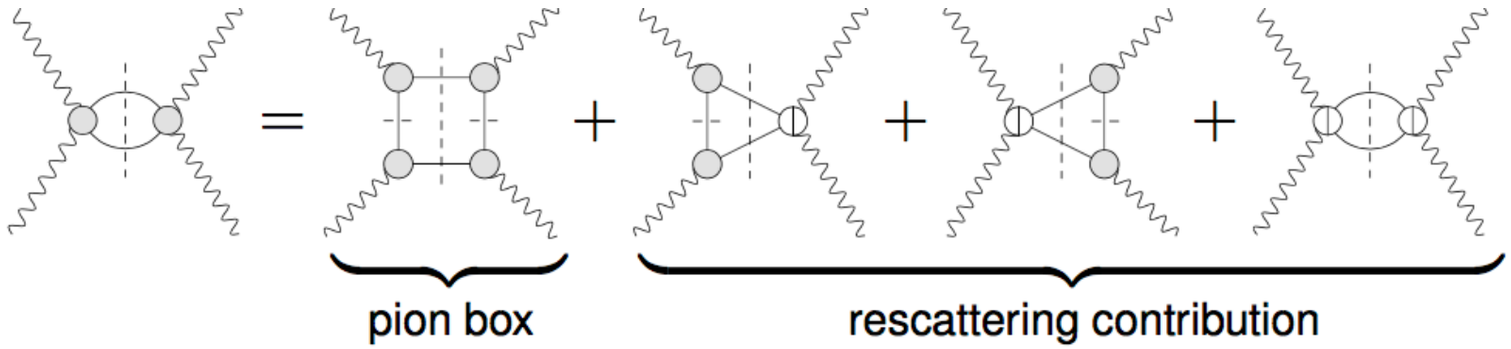
$$a_\mu^{\pi\text{-box}} + a_{\mu,J=0}^{\pi\pi,\pi\text{-pole LHC}} = -2.4(1) \times 10^{-10}$$

Colangelo et al. (2014-2017)

Multi-meson contribution to (g-2)



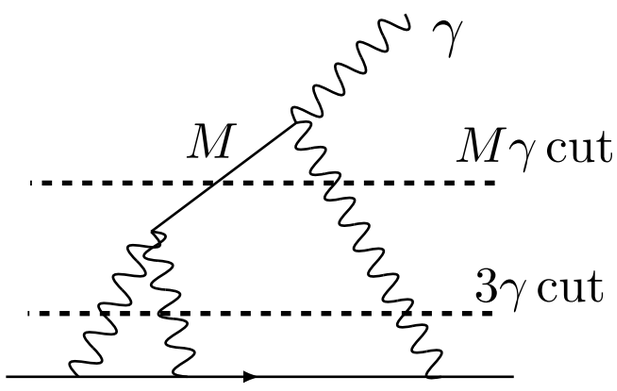
- Pioneering dispersive analyses for ππ loop contribution to a_μ

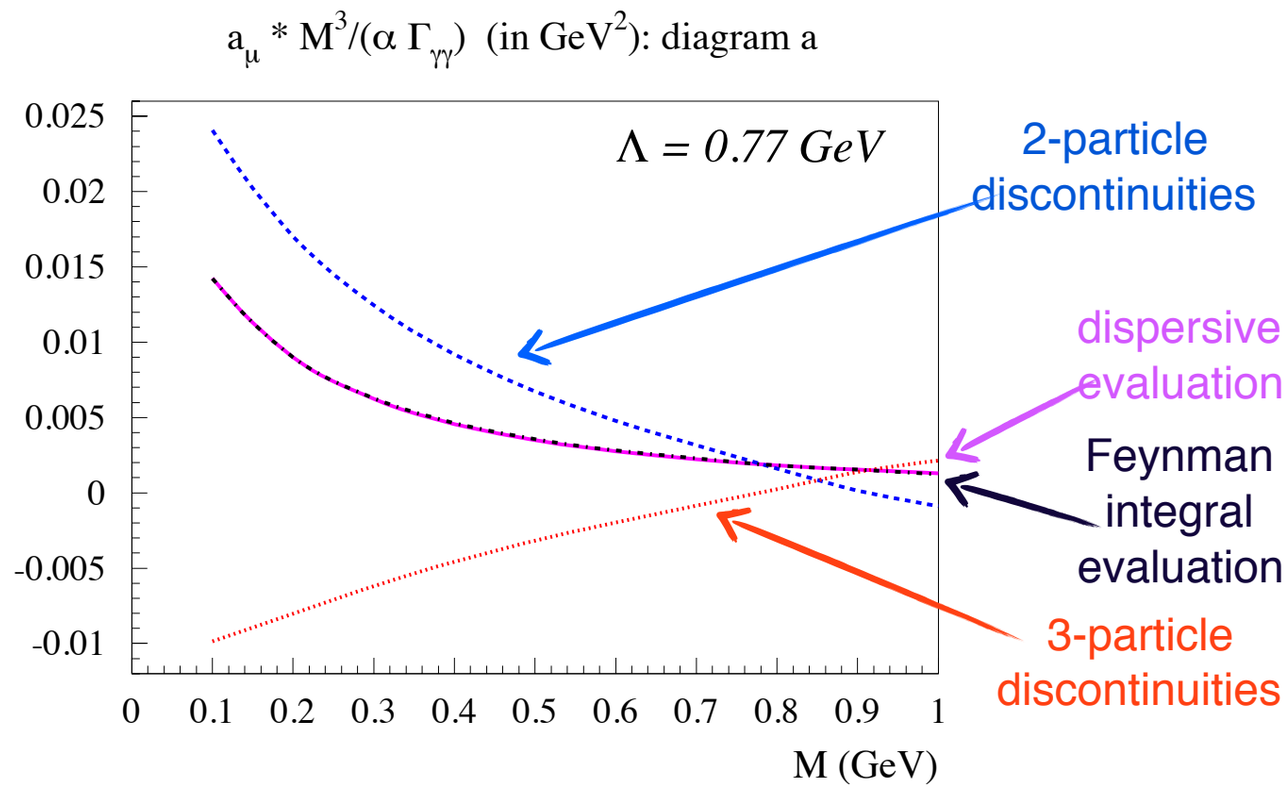


$$a_\mu^{\pi\text{-box}} + a_{\mu,J=0}^{\pi\pi,\pi\text{-pole LHC}} = -2.4(1) \times 10^{-10}$$

Colangelo et al. (2014-2017)

- Dispersive analysis for muon Pauli Form Factor

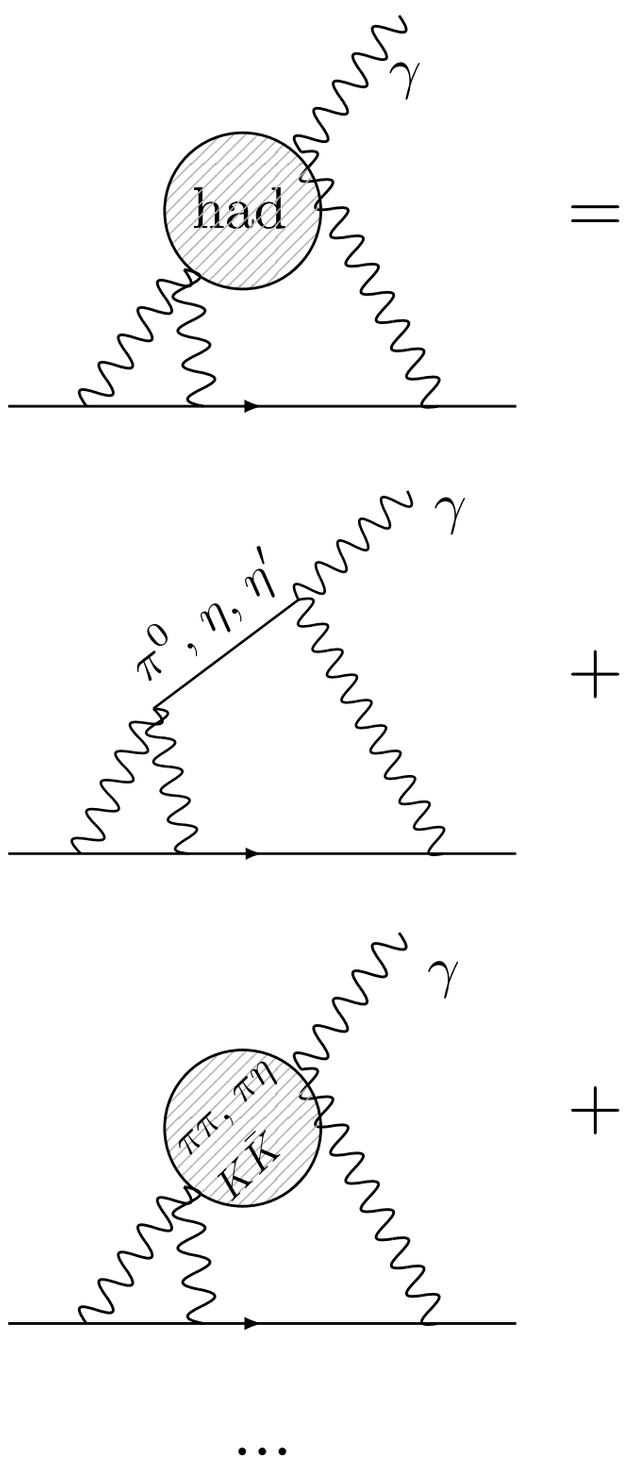


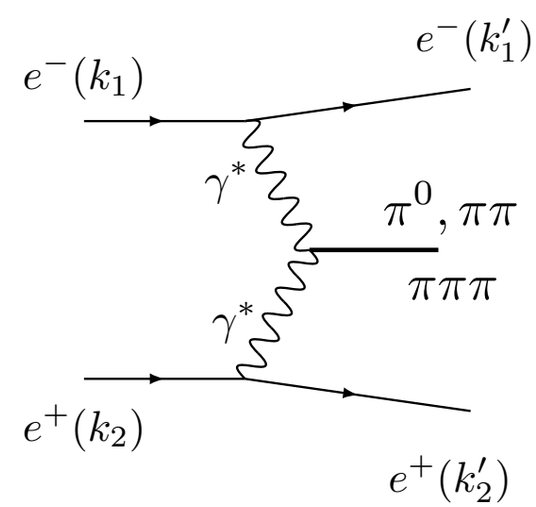


Pauk, Vanderhaeghen (2014)

Observables in experiment

$q^2 > 0$	time-like γ^*
$q^2 = -Q^2 < 0$	space-like γ^*





$\gamma\gamma^*\pi^0$, $\gamma\gamma^*\eta$, $\gamma\gamma^*\eta'$

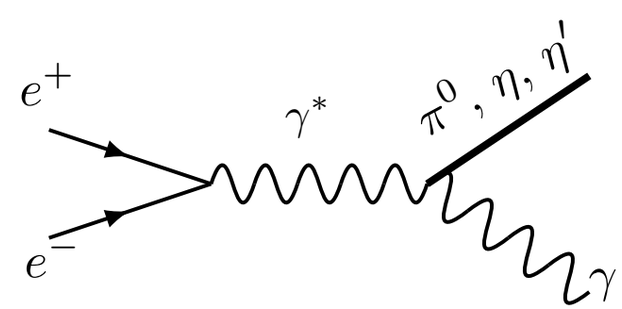
CLEO, BaBar, Belle, BESIII, ...

$\gamma\gamma \to \pi\pi, \pi\eta, K\bar{K}$

Mark II, CELLO, Crystal Ball, Belle

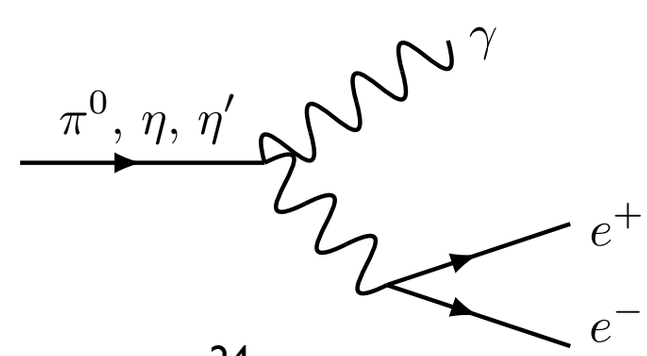
$\gamma\gamma^* \to \pi\pi$

BESIII



$\gamma\gamma^*\pi^0$, $\gamma\gamma^*\eta$

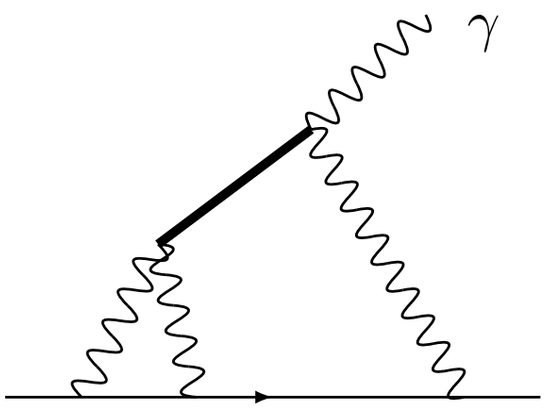
SND, CMD-2, BESIII, ...



$\gamma\gamma^*\pi^0$, $\gamma\gamma^*\eta$, $\gamma\gamma^*\eta'$

KLOE, MAMI/A2

Meson pole contributions to (g-2)



$$a_\mu^{LbL} = \lim_{k\to 0} ie^6 \int \frac{d^4q_1}{(2\pi)^4} \int \frac{d^4q_2}{(2\pi)^4} T^{\mu\nu\lambda\sigma}(q_1, k-q_1-q_2, q_2)\, \Pi_{\mu\nu\lambda\sigma}(q_1, k-q_1-q_2, q_2)$$

$$\frac{1}{q_1^2}\frac{1}{q_2^2}\frac{1}{(k-q_1-q_2)^2}\frac{1}{(p+q_1)^2-m^2}\frac{1}{(p'-q_2)^2-m^2}$$

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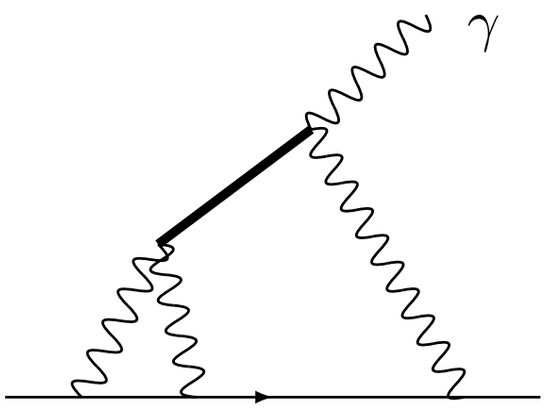

Lepton tensor: well known

$$a_\mu^{LbL} = \lim_{k\to 0} ie^6 \int \frac{d^4q_1}{(2\pi)^4} \int \frac{d^4q_2}{(2\pi)^4} T^{\mu\nu\lambda\sigma}(q_1, k-q_1-q_2, q_2)\, \Pi_{\mu\nu\lambda\sigma}(q_1, k-q_1-q_2, q_2)$$

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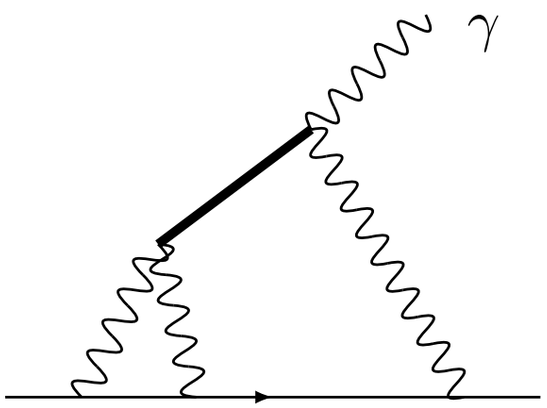

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Hadron tensor: requires input from **TFFs**

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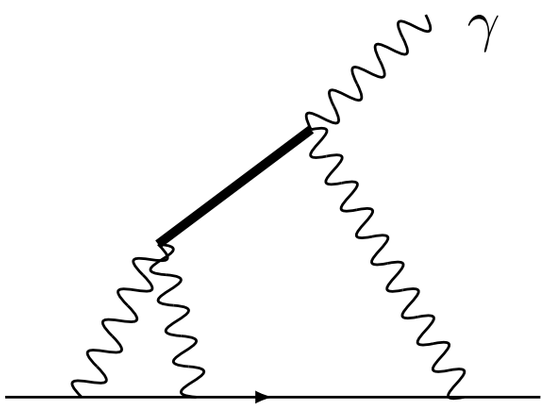

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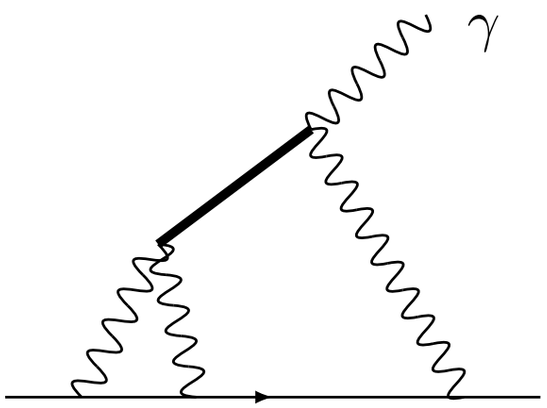
Results (excluding low energy region):

$$a_\mu[f_2(1270), f_2(1565)] = (0.1 \pm 0.01) \times 10^{-10}$$

Meson pole contributions to (g-2)



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New evaluation of axial vector contributions (satisfying Landau-Yang theorem)

$$a_\mu[f_1(1285), f_1(1420)] = (0.64 \pm 0.20) \times 10^{-10}$$
$$= (0.75 \pm 0.27) \times 10^{-10}$$

Pauk, Vdh (2013)
Jegerlehner (2015)

$$\delta a_\mu^{exp} = 1.6 \times 10^{-10}$$

FNAL, J-PARC experiments