QED Contribution to electron and muon $g-2$

T. Aoyama (KEK)

based on collaboration with
T. Kinoshita (Cornell and UMass Amherst),
M. Nio (RIKEN),
M. Hayakawa (Nagoya University)

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Hadronic Contribution to $(g-2)_\mu$
University of Washington, Seattle

Anomalous magnetic moment of leptons

- Electron $g-2$ is explained almost entirely by QED interaction between electron and photons. It has been the most stringent test of QED and the standard model.

in units of 10^{-12}

a_e(expr:HV08)	1 159 652 180.73 (28)
a_e(theory)	1 159 652 181.61 (23)
QED: e and γ	1 159 652 177.14 (23)
QED: μ, τ contributions	2.747 5720 (14)
Hadronic	1.693 (12)
Weak	0.03053 (23)

- Electron $g-2$ provides one of the most precise determination of the fine structure constant α.
- Muon $g-2$ is also dominated by QED contribution, which has been evaluated precisely for the on-going experiments.

Anomalous magnetic moment of electron

- The best measurement of the anomalous magnetic moment of electron obtained by Harvard group is:

$$a_e(\text{HV08}) = 1\ 159\ 652\ 180.73\ (28) \times 10^{-12} \qquad [0.24\text{ppb}]$$

Hanneke, Fogwell, Gabrielse, PRL100, 120801 (2008)
Hanneke, Fogwell Hoogerheide, Gabrielse, PRA83, 052122 (2011)

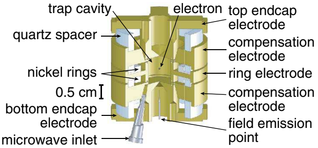


FIG. 2 (color). Cylindrical Penning trap cavity used to confine a single electron and inhibit spontaneous emission.

This result is 15-fold improvement over the previous measurement by the University of Washington group:

$$a_{e^-}(\text{UW87}) = 1\ 159\ 652\ 188.4\ (43) \times 10^{-12} \qquad [3.7\text{ppb}]$$

$$a_{e^+}(\text{UW87}) = 1\ 159\ 652\ 187.9\ (43) \times 10^{-12} \qquad [3.7\text{ppb}]$$

Van Dyck, Schwinberg, Dehmelt, PRL59, 26 (1987)

- Further improvement of electron anomaly as well as new measurement of positron is ongoing.

Standard Model prediction of a_e

- Contributions to electron $g-2$ within the context of the standard model consist of:

$$a_e = a_e(QED) + a_e(Hadronic) + a_e(Weak)$$

- QED contribution is further divided according to its lepton-mass dependence through mass-ratio:

$$a_e(\mathrm{QED}) = \underbrace{A_1}_{e,\gamma} + \underbrace{A_2(m_e/m_\mu)}_{e,\mu,\gamma} + \underbrace{A_2(m_e/m_\tau)}_{e,\tau,\gamma} + \underbrace{A_3(m_e/m_\mu, m_e/m_\tau)}_{e,\mu,\tau,\gamma}$$

- Each contribution is evaluated by perturbation theory:

$$A_i = A_i^{(2)}\left(\frac{\alpha}{\pi}\right) + A_i^{(4)}\left(\frac{\alpha}{\pi}\right)^2 + A_i^{(6)}\left(\frac{\alpha}{\pi}\right)^3 + A_i^{(8)}\left(\frac{\alpha}{\pi}\right)^4 + \cdots$$

These coefficients are calculated by using Feynman-diagram techniques.

	# diagrams	w/o fermion loop	w/ fermion loop
2nd	1	1	0
4th	7	6	1
6th	72	50	22
8th	891	518	373
10th	12,672	6536	6318

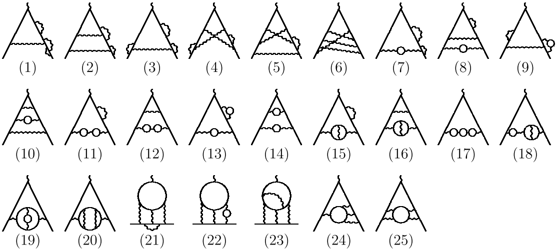
QED contribution: Summary

Coefficient $A_i^{(2n)}$	Value (Error)	References
$A_1^{(2)}$	0.5	Schwinger 1948
$A_1^{(4)}$	$-0.328\,478\,965\,579\,193\cdots$	Petermann 1957, Sommerfield 1958
$A_2^{(4)}(m_e/m_\mu)$	$0.519\,738\,676\,(24)\times10^{-6}$	Elend 1966
$A_2^{(4)}(m_e/m_\tau)$	$0.183\,790\,(25)\times10^{-8}$	Elend 1966
$A_1^{(6)}$	$1.181\,241\,456\,587\cdots$	Laporta-Remiddi 1996, Kinoshita 1995
$A_2^{(6)}(m_e/m_\mu)$	$-0.737\,394\,164\,(24)\times10^{-5}$	Samuel-Li, Laporta-Remiddi, Laporta
$A_2^{(6)}(m_e/m_\tau)$	$-0.658\,273\,(79)\times10^{-7}$	Samuel-Li, Laporta-Remiddi, Laporta
$A_3^{(6)}(m_e/m_\mu, m_e/m_\tau)$	$0.1909\,(1)\times10^{-12}$	Passera 2007
$A_1^{(8)}$	$-1.912\,245\,764\cdots$	Laporta 2017, AHKN 2015
$A_2^{(8)}(m_e/m_\mu)$	$0.916\,197\,070\,(37)\times10^{-3}$	Kurz et al 2014, AHKN 2012
$A_2^{(8)}(m_e/m_\tau)$	$0.742\,92\,(12)\times10^{-5}$	Kurz et al 2014, AHKN 2012
$A_3^{(8)}(m_e/m_\mu, m_e/m_\tau)$	$0.746\,87\,(28)\times10^{-6}$	Kurz et al 2014, AHKN 2012
$A_1^{(10)}$	$6.737\,(159)$	AKN 2018,2019
$A_2^{(10)}(m_e/m_\mu)$	$-0.003\,82\,(39)$	AHKN 2012,2015
$A_2^{(10)}(m_e/m_\tau)$	$\mathcal{O}(10^{-5})$	
$A_3^{(10)}(m_e/m_\mu, m_e/m_\tau)$	$\mathcal{O}(10^{-5})$	

All terms up to 8th order are well-known. 10th order term is obtained numerically.

QED contribution: 8th order term

- 891 Feynman diagrams contribute to 8th order $A_1^{(8)}$ term.



- Laporta obtained near-analytic precise value upto 1100 digits.

−1.91224576492644557415264716743983005406087339065872534517132984800 60
384439806517061427608927000036315837558415331473270056378514912854 5391
902804327050273822304345578957045562729309941296699760277782211578 4720
339064151908166527097970867438115012155147972274322164273431927975 9586
074050057837384960701874328314024838025192249460742298558930463506 1404
922526634310944240002356356881280620645494013224977594300429288836 7617
488992369151808780869897052635785337537769641170245361960134975744 9436
126848617516260683238718674730383150596274187801530551487940053697 7798
369464278684326918431175889581159743566950433048349073613426586499 5311
638781174347538542348836408558444188223721745670687104182330743051 7443
055739459611715508589611489952612660612469940731184039274723400234 6496
953173548258481799822409737371077365740464513521123091242528111137 2153
021544537210148111211598489708842232798797204842014451228284515165 8523
656178659459260099173303172130286546721234534050034910470072892448 7200
616044261325449069000431915198230047488181494311038495378299406296 7586
787538524978194698979313216219797575067670114290489796208505078559 2…

Laporta, PLB772, 232 (2017)

QED contribution: 8th order term

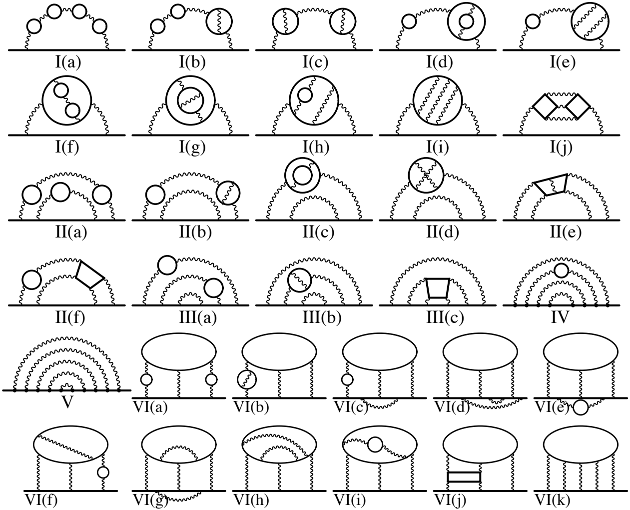
- Mass-independent term $A_1^{(8)}$
  - Near-analytic result
    $-1.9122457649264455741526471674\ldots$ — Laporta, PLB772, 232 (2017)
  - Alternative semi-analytic result
    $-1.87(12)$ — Marquad et al, arXiv:1708.07138
  - Numerical result
    $-1.91298(84)$ — AHKN, PRL109, 111809 (2012); PRD91, 033006 (2015)
- Mass-dependent terms $A_2^{(8)}$ and $A_3^{(8)}$
  - Numerical evaluation. — AHKN, PRL109, 111809 (2012)
  - Analytic calculation by the series expansion in mass-ratio $m_e/m_\ell \ll 1$.
    Kurz et al. PRD93, 053017 (2016)

	Analytic	Numerical
$A_2^{(8)}(m_e/m_\mu)$	$0.916\,197\,070\,(37) \times 10^{-3}$	$0.9222\,(66) \times 10^{-3}$
$A_2^{(8)}(m_e/m_\tau)$	$0.742\,92\,(12) \times 10^{-5}$	$0.738\,(12) \times 10^{-5}$
$A_3^{(8)}(m_e/m_\mu, m_e/m_\tau)$	$0.746\,87\,(28) \times 10^{-6}$	$0.7465\,(18) \times 10^{-6}$

- Now the 8th order term is well-known.

QED contribution: 10th order term

- 12 672 Feynman diagrams contribute to 10th order term.
  They are classified into 32 gauge invariant sets within 6 supersets.

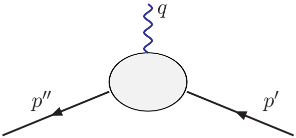


Most difficult is Set V that consists of 6354 diagrams w/o lepton loops.

Magnetic moment contribution

- Magnetic property of lepton can be studied through examining its scattering by a static magnetic field.
  The amplitude can be represented as:

$$e\bar{u}(p'')\left[\gamma^{\mu}\,F_1(q^2) + \frac{i}{2m}\sigma^{\mu\nu}\,q_{\nu}\,F_2(q^2)\right]u(p')\,A^{e}_{\mu}(\vec{q})$$



- The anomalous magnetic moment is the static limit of the magnetic form factor $F_2(q^2)$:

$$a_{\ell} = F_2(0) = Z_2 M, \qquad M = \lim_{q^2 \to 0} \mathrm{Tr}(P_{\nu}(p,q)\Gamma^{\nu})$$

  where Γ^{ν} is the proper vertex function with the external lepton on the mass shell, and $P_{\nu}(p,q)$ is the magnetic projection operator.

Numerical Approach

- A set of vertex diagrams Λ obtained by inserting an external vertex into each lepton line of self-energy diagram Σ can be related by Ward-Takahashi identity.

$$\Lambda^{\nu}(p,q) \simeq -q_{\mu} \left. \frac{\partial \Lambda^{\mu}(p,q)}{\partial q_{\nu}} \right|_{q\to 0} - \frac{\partial \Sigma(p)}{\partial p_{\nu}} .$$

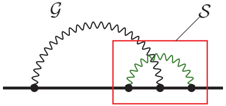
- Amplitude is given by an integral over loop momenta according to Feynman-Dyson rule.
  It is converted into Feynman parametric integral over $\{z_i\}$. Momentum integration is carried out analytically that yields

$$M_{\mathcal{G}}^{(2n)} = \left(-\frac{1}{4}\right)^{n} \Gamma(n-1) \int (dz)_{\mathcal{G}} \left[\frac{F_0}{U^2 V^{n-1}} + \frac{F_1}{U^3 V^{n-2}} + \cdots \right]$$

- Integrand is expressed by a rational function of terms called *building blocks*, U, V, B_{ij}, A_j, and C_{ij}.
  Building blocks are given by functions of $\{z_i\}$, reflecting the topology of diagram, flow of momenta, etc.

Subtraction of UV Divergences

- UV divergence occurs when loop momenta in a subdiagram go to infinity. It corresponds to the region of Feynman parameter space $z_i \sim \mathcal{O}(\epsilon)$ for $i \in \mathcal{S}$.



- In order to carry out subtraction numerically, the singularities are cancelled point-by-point on Feynman parameter space.

$$M_{\mathcal{G}} - L_{\mathcal{S}} M_{\mathcal{G}/\mathcal{S}} \quad \longrightarrow \quad \int (dz)_{\mathcal{G}} \left[m_{\mathcal{G}} - \mathbb{K}_{\mathcal{S}} m_{\mathcal{G}} \right]$$

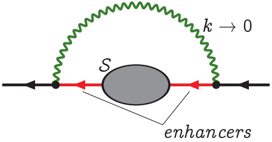
- The subtraction integrand $\mathbb{K}_{\mathcal{S}} m_{\mathcal{G}}$ is derived from $m_{\mathcal{G}}$ by simple power-counting rule called ***K*-operation**. Cvitanović and Kinoshita, 1974
- By construction, subtraction terms can be factorized into (UV-divergent part of) renormalization constant and lower-order magnetic part.

$$\int (dz)_{\mathcal{G}} \left[\mathbb{K}_{\mathcal{S}} m_{\mathcal{G}} \right] \quad = \quad L_{\mathcal{S}}^{\text{UV}} M_{\mathcal{G}/\mathcal{S}}$$

$L_{\mathcal{S}}^{\text{UV}}$ is the leading UV-divergent part of $L_{\mathcal{S}}$.

IR subtraction Scheme

- A diagram may have IR divergence when some momenta of photon go to zero. It is really divergent by "enhancer" leptons that are close to on-shell by kinematical constraint.



- We adopt subtraction approach for these divergences point-by-point on Feynman parameter space.

- There are two types of sources of IR divergence in $M_{\mathcal{G}}$ associated with a self-energy subdiagram. To handle these divergences, we introduce two subtraction operations:
  - ***R*-subtraction** to remove the residual self-mass term

$$\mathbb{R}_{\mathcal{S}} M_{\mathcal{G}} = \widetilde{\delta m}_{\mathcal{S}} M_{\mathcal{G}/\mathcal{S}(i^*)}$$

  - ***I*-subtraction** to subtract remaining logarithmic IR divergence

$$\mathbb{I}_{\mathcal{S}} M_{\mathcal{G}} = \widetilde{L}_{\mathcal{G}/\mathcal{S}(k)} M_{\mathcal{S}}$$

Amplitude as a finite integral

- Finite amplitude $\Delta M_{\mathcal{G}}$ free from both UV and IR divergences is obtained by Feynman-parameter integral as:

$$\Delta M_{\mathcal{G}} = \int (dz) \Big[F_{\mathcal{G}}$$

← unrenormalized amplitude

$$+ \sum_{f} \prod_{\mathcal{S} \in f} (-\mathbb{K}_{\mathcal{S}}) F_{\mathcal{G}}$$

← UV subtraction terms

f: **Zimmermann's forests**:
combinations of UV divergent subdiagrams.

$$+ \sum_{\tilde{f}} (-\mathbb{I}_{\mathcal{S}_i}) \cdots (-\mathbb{R}_{\mathcal{S}_j}) \cdots F_{\mathcal{G}} \Big]$$

← IR subtraction terms

$\widetilde{f}$: **annotated forests**:
combinations of self-energy subdiagrams
with distinction of I-/R-subtractions.

Residual renormalization

- We adopt the standard on-shell renormalization to ensure that the coupling constant α and the electron mass m_e are the ones measured by experiments.
- The sum of all these finite integrals defined by K-operation and I-/R-subtraction operations does not correspond to physical contribution to $g-2$.
- The difference is adjusted by the step called the residual renormalization.

$$
\begin{aligned}
a_e &= M(\text{bare}) - \text{on-shell renormalization} \\
&= \underbrace{\Big[M(\text{bare}) - \text{UV subtr.} - \text{IR subtr.}\Big]}_{\text{Finite integral } \Delta M} \\
&\quad + \underbrace{\Big[-\text{on-shell renorm.} + \text{UV subtr.} + \text{IR subtr.}\Big]}_{\text{finite residual renormalization}}
\end{aligned}
$$

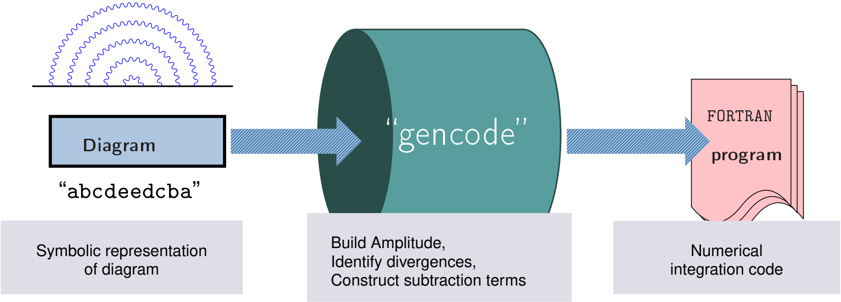
Deriving residual renormalization

- Sum up over 389 integrals of 10th order Set V, which requires analytic sum of $\sim$ 16,000 symbolic terms.
- The physical contribution from 10th order Set V is given as:

$$
\begin{aligned}
A_1^{(10)}[\text{Set V}] = {} & \Delta M_{10}[\text{Set V}] \\
& + \Delta M_8(-7\Delta LB_2) \\
& + \Delta M_6\{-5\Delta LB_4 + 20(\Delta LB_2)^2\} \\
& + \Delta M_4\{-3\Delta LB_6 + 24\Delta LB_4\Delta LB_2 - 28(\Delta LB_2)^3 + 2\Delta L_{2^*}\Delta dm_4\} \\
& + M_2\{-\Delta LB_8 + 8\Delta LB_6\Delta LB_2 - 28\Delta LB_4(\Delta LB_2)^2 \\
& \qquad + 4(\Delta LB_4)^2 + 14(\Delta LB_2)^4 + 2\Delta dm_6\Delta L_{2^*}\} \\
& + M_2\Delta dm_4(-16\Delta L_{2^*}\Delta LB_2 + \Delta L_{4^*} - 2\Delta L_{2^*}\Delta dm_{2^*}),
\end{aligned}
$$

- The terms with Δ are the finite nth order quantities.
  - ΔM_n, M_2: finite magnetic moment.
  - ΔLB_n: sum of vertex and wave-function renormalization constants.
  - Δdm_n: mass-renormalization constants.
  - ΔL_n^*, Δdm_n^*: $*$ denotes mass insertion.

Construction of numerical integration code



- We need to evaluate a large number of Feynman diagrams.
  It should be error-prone by writing numerical integration code for these huge integrals by hand. We developed an automated code-generating program.
- "gencode*N*" takes a single-line information that represents a diagram, and generates numerical integration code in FORTRAN.
- These integrals are evaluated on computers using numerical integration routines.

AHKN, NPB740, 138 (2006); NPB796, 184 (2008)

Numerical integration

- Multi-dimensional integral
  - The amplitude is expressed as a $14-1$ dimensinal integral for 10th order diagrams.
  - The integrands are huge. (approx. $\mathcal{O}(10^5)$ FORTRAN lines for each integral.)

- Digit-deficiency problem
  - The point-by-point subtraction suffers from severe digit-deficiency problem by rounding-off of floating-point numbers.

    We employ extended numerical precision arithmetic using *double-double* and *quadruple-double* of `qd` library.

    Bailey, Hida, Li. c.f. http://crd.lbl.gov/~dhbailey/mpdist/

- Sharp peaks
  - Integrands have sharp peaks due to divergences, and therefore requires robust integration method.

    We employ VEGAS, an adaptive-iterative Monte-Carlo integration algorithm.

    Lepage, J.Comput.Phys.27, 192 (1978)
    A new version of VEGAS: https://github.com/gplepage/vegas

QED contribution: 10th order term

- Numerical evaluation of the complete 10th order contribution was reported in 2012 and an updated result was published in 2015. Latest value is:

$$A_1^{(10)} = 6.737\ (159)$$

- Contribution to $A_1^{(10)}$ mainly comes from Set V that consists of 6354 vertex diagrams without closed lepton loops.

  Recently, Volkov announced their preliminary result by an independent numerical method.

$$A_1^{(10)}[\text{Set V}] = \begin{cases} 7.668\ (159) & \text{AKN, Atoms, 7, 28 (2019)} \\ 6.782\ (113) & \text{Volkov, ACAT2019, arXiv:1905.08007} \end{cases}$$

  Difference $-0.89\ (20)\ [4.5\sigma]$ does not affect seriously in the current precision.

- Mass-dependent term is also evaluated:

$$A_2^{(10)}(m_e/m_\mu) = -0.003\,82\ (39)$$

  tau-lepton contribution is negligibly small for the current experimental precision.

Numerical checks of Set V integrals

- 13 integration variables in $[0,1]^D$ are mapped to 14 Feynman parameters. Any mapping should yield the same result.
- As a cross check, we performed integrals with different mappings. They are regarded as independent evalations.
- Numerical results are in good agreement.

List of results that exhibit relatively large differences:

Diagram	Expression	Results in 2015	Results in 2017	Difference	Weighted average
$X141$	*abbcadedec*	-12.5567 (350)	-12.4879 (207)	-0.0688	-12.5057 (178)
$X113$	*abacddeebc*	-4.3847 (322)	-4.4412 (176)	0.0565	-4.4282 (155)
$X100$	*abacdcdeeb*	-15.2919 (331)	-15.2360 (203)	-0.0559	-15.2513 (173)
$X256$	*abccdeedba*	-14.0405 (342)	-13.9856 (194)	-0.0549	-13.9990 (169)
$X146$	*abbcdadeec*	-2.2990 (335)	-2.2458 (202)	-0.0532	-2.2600 (173)
$X075$	*abacbddeec*	-8.1138 (340)	-8.0608 (195)	-0.0531	-8.0739 (169)
$X144$	*abbccdedea*	23.7239 (368)	23.6713 (189)	0.0526	23.6823 (168)
$X252$	*abccdedeab*	-10.9091 (343)	-10.8565 (179)	-0.0526	-10.8677 (158)
$X236$	*abcbdedcea*	2.0560 (180)	2.1072 (205)	-0.0512	2.0782 (135)
$X325$	*abcdceedba*	11.5958 (343)	11.5456 (198)	0.0503	11.5582 (172)
$X158$	*abbcdeceda*	0.4607 (329)	0.4106 (206)	0.0502	0.4247 (174)

AKN, PRD97, 036001 (2018)

Fine Structure Constant α

- To obtain the theoretical prediction of a_e, we need a value of the fine-structure constant α determined independent of QED.

- Two high-precision values of α are obtained from the measurement of $h/m(X)$ of the Rb and Cs by the atom interferometer through the relation:

$$\alpha^{-1} = \left[\frac{2R_\infty}{c}\frac{A_r(X)}{A_r(e)}\frac{h}{m(X)}\right]^{-1/2}$$

where

- R_∞ the Rydberg constant

$$R_\infty(\mathrm{MPQ}) = 10\,973\,731.568\,076\,(096)\ m^{-1}$$ Beyer et al. Science, 358, 79 (2017)

$$R_\infty(\mathrm{Orsay}) = 10\,973\,731.568\,530\,(140)\ m^{-1}$$ Fleurbaey et al, PRL720, 183001 (2018)

- $A_r(X)$ relative atomic mass of an atom X
- $m(X)$ mass of an atom X

It leads to

$$\alpha^{-1}(\mathrm{Rb}) = 137.035\,998\,995\,(85)\ [0.62\mathrm{ppb}]$$ Bouchendira et al, PRL106, 080801 (2011)

$$\alpha^{-1}(\mathrm{Cs}) = 137.035\,999\,046\,(27)\ [0.20\mathrm{ppb}]$$ Parker et al, Science, 360, 191 (2018)

Theoretical Prediction of a_e

- Using α(Cs) and including the hadronic and weak contributions, the theoretical prediction of a_e becomes:

QED	mass-independent	mass-dependent	sum
2nd	1 161 409 733.21 (23)	0	1 161 409 733.21 (23)
4th	−1 772 305.063 85 (70)	2.814 1613 (13)	−1 772 302.249 69 (70)
6th	14 804.203 6740 (88)	−0.093 240 76 (10)	14 804.110 4333 (88)
8th	−55.667 989 379 (44)	0.026 909 719 (35)	−55.641 079 660 (56)
10th	0.456 (11)	−0.000 258 (26)	0.455 (11)
a_e(QED)	1 159 652 177.14 (23)	2.747 5720 (14)	1 159 652 179.88 (23)
Weak			
a_e(weak)			0.030 53 (23)
Hadron			
VP LO			1.849 (10)
VP NLO			−0.2213 (11)
VP NNLO			0.027 99 (17)
LbyL			0.037 (5)
a_e(hadron)			1.693 (12)
a_e(theory)			1 159 652 181.61 (23)

Theoretical Prediction of a_e

- We obtain the theoretical prediction of a_e as

$$a_e(\text{theory: } \alpha(\text{Rb})) = 1\ 159\ 652\ 182.037\ (720)(11)(12) \times 10^{-12}$$
$$a_e(\text{theory: } \alpha(\text{Cs})) = 1\ 159\ 652\ 181.606\ (229)(11)(12) \times 10^{-12}$$

  where uncertainties are due to fine-structure constant α, QED 10th order, and hadronic contribution.

- The measurement of a_e is

$$a_e(\text{expt.}) = 1\ 159\ 652\ 180.73\ (28) \times 10^{-12}$$

- The differences between theory and measurement are

$$a_e(\text{theory: } \alpha(\text{Rb})) - a_e(\text{expt.}) = 1.31\ (77) \times 10^{-12}\ [1.7\sigma]$$
$$a_e(\text{theory: } \alpha(\text{Cs})) - a_e(\text{expt.}) = 0.88\ (36) \times 10^{-12}\ [2.4\sigma]$$

Fine Structure Constant α from a_e

- From the measurement and the theory of electron $g-2$, the value of fine-structure constant can be determined.

Theoretical calculations

Experimental value

$$a_e = A^{(2)}\left(\frac{\alpha}{\pi}\right) + A^{(4)}\left(\frac{\alpha}{\pi}\right)^2 + A^{(6)}\left(\frac{\alpha}{\pi}\right)^3 + A^{(8)}\left(\frac{\alpha}{\pi}\right)^4 + A^{(10)}\left(\frac{\alpha}{\pi}\right)^5 + \cdots$$

$$+(\text{small contributions})$$

- Newly obtained value of fine-structure constant is:

(α^5) (had) (exp)

$$\alpha^{-1}(a_e) = 137.035\,999\,1496\,(13)(14)(330) \quad [0.24\text{ppb}]$$

AKN, Atoms, 7, 28 (2019)

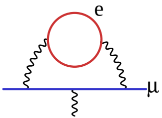
- The differences in α from the atomic recoil determinations are

$$\alpha^{-1}(a_e) - \alpha^{-1}(\text{Rb}) = 0.155\,(91) \times 10^{-6}\ [1.7\sigma],$$
$$\alpha^{-1}(a_e) - \alpha^{-1}(\text{Cs}) = 0.104\,(43) \times 10^{-6}\ [2.4\sigma].$$

Muon $g-2$: QED contribution

- What distinguishes a_e(QED) and a_μ(QED) is the mass-dependent component.
- Light lepton loop contribution yields large logarithmic enhancement involving a factor $\ln(m_e/m_\mu)$.
  - Vacuum polarization loop:

$$\frac{2}{3}\ln(m_\mu/m_e) - \frac{5}{9} \simeq 3.$$

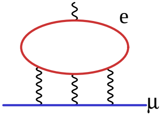


  - Light-by-light scattering loop:

$$\frac{2}{3}\pi^2 \ln(m_\mu/m_e) \simeq 35.$$

6th-order l-by-l effect is important.

c.f. Aldins, Kinoshita, Brodsky, Dufner, PRL8, 441 (1969)



- Therefore, the sets of diagrams giving the leading contribution can be identified and were evaluated in the earlier stage.
  The entire contribution including non-leading diagrams have been evaluated.

Muon $g-2$: QED contribution

- a_μ(QED) is known up to 10th order. Their values contributing to mass-dependent terms are:

	$A_2(m_\mu/m_e)$	$A_2(m_\mu/m_\tau)$	$A_3(m_\mu/m_e, m_\mu/m_\tau)$
4th	1.094 258 3093 (76)	0.000 078 076 (11)	—
6th	22.868 379 98 (20)	0.000 360 671 (94)	0.000 527 738 (75)
8th	132.685 2 (60)	0.042 4941 (53)	0.062 722 (10)
10th	742.32 (86)	$-$0.0656 (45)	2.011 (10)

Elend, PL20, 682 (1966); Samuel and Li, PRD44, 3935 (1991); Li, Mendel and Samuel, PRD47, 1723 (1993)
Laporta, Nuovo Cim. A106, 675 (1993); Laporta and Remiddi, PLB301, 440 (1993); Czarnecki and Skrzypek, PLB449, 354 (1999)
Laporta, PLB312, 495 (1993); Kinoshita and Nio, PRD70, 113001 (2004); Kurz, Liu, Marquard, Steinhauser, NPB879, 1 (2014)
Laporta, PLB328, 522 (1994); Kinoshita and Nio, PRD73, 053007 (2006)
TA, Hayakawa, Kinoshita, Nio, Watanabe, PRD78, 053005 (2008)
TA, Asano, Hayakawa, Kinoshita, Nio, Watanabe, PRD81, 053009 (2010)
TA, Hayakawa, Kinoshita, Nio, PRD78, 113006 (2008); 82, 113004 (2010); 83, 053002 (2011)
83, 053003 (2011); 84, 053003 (2011); 85, 033007 (2012); 85, 093013 (2012)

- Together with the mass-independent term A_1, we obtain:

$$a_\mu(\mathrm{QED} : \alpha(\mathrm{Cs})) = 116\,584\,718.931\,(7)\,(17)\,(6)\,(100)\,(23)\,[104] \times 10^{-11}$$

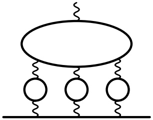
$$a_\mu(\mathrm{QED} : \alpha(a_e)) = 116\,584\,718.842\,(7)\,(17)\,(6)\,(100)\,(28)\,[106] \times 10^{-11}$$

(mass ratio)(8th)(10th)(12th)(α)[combined]

12th order contribution

- In view of rather large values of $A_2(m_\mu/m_e)$, one might wonder how much the twelfth order contribution.

- The leading contribution will come from three insertions of 2nd-order vacuum-polarization loop into the 6th-order light-by-light diagram. It is estimated as:



$$\sim \text{(6th light-by-light)} \times \text{(2nd VP)}^3 \times 10 \times \left(\frac{\alpha}{\pi}\right)^6$$

$$\sim 0.08 \times 10^{-11}.$$

It is larger than the uncertainty of 10th order term. A crude evaluation may be desirable.

Summary

- QED contribution to electron $g-2$ up to 8th order has been firmly established.
- QED contribution of 10th order has been evaluated by extensive numerical calculation.
- QED contributions are now ready for the on-going new measurements of electron and position $g-2$, and muon $g-2$.
- Electron $g-2$ provides one of most precise determination of fine structure constant α.
  It serves for new SI as a significant factor of the uncertainty of many physical constants.

Backup

Numerical Approach

- Procedure:

  Step 1. Find distinct set of Feynman diagrams.

  Step 2. Construct **amplitude** in terms of Feynman parametric integral.

  Step 3. Construct subtraction terms of **UV divergence**.
  • K-operation

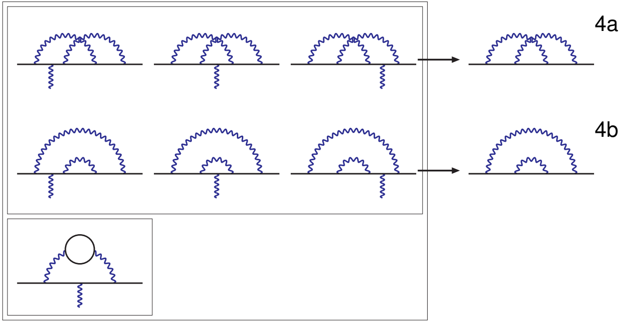
  Step 4. Construct subtraction terms of **IR divergence**.
  • R-subtraction of residual mass-renormalization.
  • I-subtraction of logarithmic IR divergences.

  Step 5. Carry out **residual renormalization**
  to achieve the standard on-shell renormalization.

  Step 6. Evaluate the finite amplitude by numerical integration.

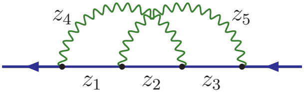
Step-by-step example with 4th-order diagrams : Step 1

- Let us illustrate the steps by simpler case,
  e.g. 4th-order diagrams.
- There are 7 diagrams of 4th order;
  6 of them have no closed lepton loop (q-type).
- They are *WT-sum*ed into 2 self-energy-like diagrams, 4a and 4b.



Step 2: Amplitude

- Introduce Feynman parameters $z_1, \ldots, z_5$ to propagators:



- Anomalous magnetic moment M_{4a} is converted analytically into the form:

$$M_{4a} = \int (dz)\, \mathcal{F}_{4a} = \int (dz) \left[\frac{E_0 + C_0}{U^2 V} + \frac{N_0 + Z_0}{U^2 V^2} + \frac{N_1 + Z_1}{U^3 V} \right]$$

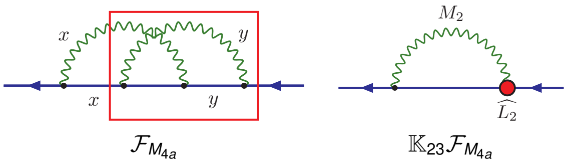
where integrand and building blocks are given as follows:

$$(dz) = dz_1 dz_2 dz_3 dz_4 dz_5 \delta(1 - z_{12345})$$
$$B_{11} = z_{235},\, B_{12} = z_{35},\, B_{13} = -z_2,$$
$$B_{23} = z_{14},\, B_{22} = z_{1345},\, B_{33} = z_{124},$$
$$U = z_2 B_{12} + z_{14} B_{11},$$
$$A_i = 1 - (z_1 B_{1i} + z_2 B_{2i} + z_3 B_{3i})/U,$$
$$G = z_1 A_1 + z_2 A_2 + z_3 A_3,\, V = z_{123} - G,$$
$$z_{ijk\cdots} = z_i + z_j + z_k + \cdots .$$

$$\begin{aligned}
E_0 &= 8(2A_1A_2A_3 - A_1A_2 - A_1A_3 - A_2A_3) \\
C_0 &= -24 Z_4 Z_5 / U \\
N_0 &= G(E_0 - 8(2A_2 - 1)) \\
Z_0 &= 8z_1(-A_1 + A_2 + A_3 + A_1A_2 + A_1A_3 - A_2A_3) \\
&\quad + 8z_2(1 - A_1A_2 + A_1A_3 - A_2A_3 + 2A_1A_2A_3) \\
&\quad + 8z_3(A_1 + A_2 - A_3 - A_1A_2 + A_1A_3 + A_2A_3) \\
N_1 &= 8G(B_{12}(2 - A_3) + 2B_{13}(1 - 2A_2) + B_{23}(2 - A_1)) \\
Z_1 &= -8z_1(B_{12}(1 - A_3) + B_{13} + B_{23}A_1) \\
&\quad + 8z_2(B_{12}(1 - A_3) - 4B_{13}A_2 + B_{23}(1 - A_1)) \\
&\quad - 8z_3(B_{12}A_3 + B_{13} + B_{23}(1 - A_1))
\end{aligned}$$

Step 3: UV subtraction

- M_{4a} is not well-defined — it has UV divergences when the loop momenta goes to infinity.
- This corresponds to a region of z_i's when all z_i on the loop vanish simultaneously.
- We prepare an integral which has the same UV divergent profile by K-operation, and perform subtraction point-by-point on the integrand.



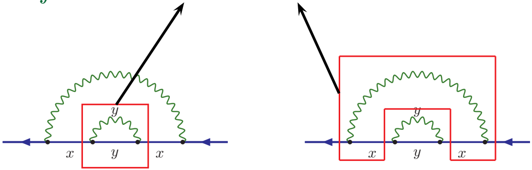
- Then the finite part of the anomalous magnetic moment ΔM_{4a} is obtained by the integral:

$$\Delta M_{4a} = \int (dz) \left[\mathcal{F}_{4a} - \mathbb{K}_{12}\mathcal{F}_{4a} - \mathbb{K}_{23}\mathcal{F}_{4a} \right]$$

Step 4: IR subtraction

- M_{4b} has IR divergence as well, from vanishing of virtual photon momentum.
- This logarithmic IR divergence is handled by an integral which is constructed by I-subtraction.
- Then the finite part of the anomalous magnetic moment ΔM_{4b} is obtained by the integral:

$$\Delta M_{4b} = \int (dz) \left[\mathcal{F}_{4b} - \mathbb{K}_{22} \mathcal{F}_{4b} - \mathbb{I}_{13} \mathcal{F}_{4b} \right]$$



Step 5: Residual renormalization

- Finite part of amplitude is given in terms of integral with appropriate UV and/or IR subtraction terms.

$$\Delta M_{4a} = \int (dz) \left[\mathcal{F}_{4a} - \mathbb{K}_{12}\mathcal{F}_{4a} - \mathbb{K}_{23}\mathcal{F}_{4a} \right]$$
$$= M_{4a} - \widehat{L}_2 M_2 - \widehat{L}_2 M_2$$
$$\Delta M_{4b} = \int (dz) \left[\mathcal{F}_{4b} - \mathbb{K}_{22}\mathcal{F}_{4b} - \mathbb{I}_{13}\mathcal{F}_{4b} \right]$$
$$= M_{4b} - (\delta m_2 M_{2^\star} + \widehat{B}_2 M_2) - \widetilde{L}_2 M_2$$

- Subtraction terms are analytically factorized into products of lower-order quantities.

- Standard on-shell renormalization is denoted by

$$a^{(4)}[\text{q-type}] = M_{4a} - 2L_2 M_2$$
$$+ M_{4b} - (\delta m_2 M_{2^\star} + B_2 M_2)$$

- By substitution, magnetic moment is given

$$a^{(4)}[\text{q-type}] = (\Delta M_{4a} + \Delta M_{4b}) - \Delta LB_2\, M_2$$

where ΔLB_2 is finite part of $L_2 + B_2$.