

Is the muon just a heavy electron?

11/08/19

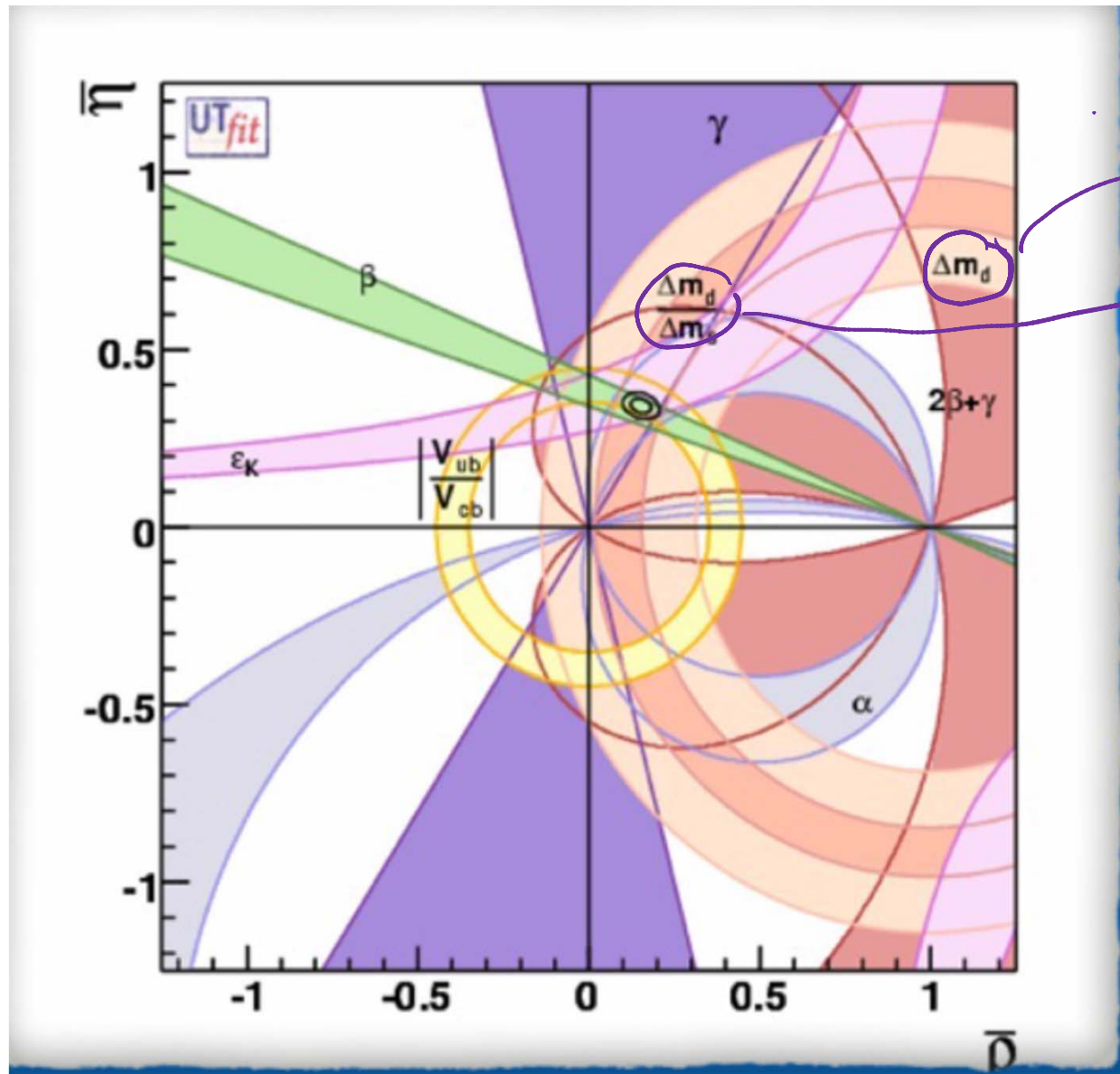
Paul Mackenzie Symposium, FermiL

Amarjit Soni
HET@BNL
[adlersoni@gmail.com]

outline

- **Core my talk will be on LFUV with some extras**

ODE TO LGT!



f_B, B_B
 ϕ SU(3) breaking
 $B \rightarrow l \nu \pi$
 $\rightarrow l \nu D$

MILESTONE in our understanding of CPV: SM-CKM

Courtesy: Tom Browder

Critical Role of the B factories in the verification of the KM hypothesis was recognized and cited by the Nobel Foundation

A single irreducible phase in the weak interaction matrix accounts for most of the CPV observed in kaons and B's.

CP violating effects in the B sector are $O(1)$ rather than $O(10^{-3})$ as in the kaon system.

5

小林益川理論が正解だった！ Bファクトリーが放った決定打



Bファクトリー実験に参加している研究教育機関

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ギーゼン大学 キョンスン大学 ハワイ大学
広島工業大学 北京 高能所
モスクワ 高エネルギー理論実験物理研
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高エネルギー加速器研究機構

Poster Designed by T. Iijima, Y. Iwasaki,
S. Kataoka, N. Katayama, K. Miyabayashi

Seminal paper providing basis for dealing with HQs on the Lattice

PHYSICAL REVIEW D

VOLUME 55, NUMBER 7

1 APRIL 1997

Massive fermions in lattice gauge theory

Aida X. El-Khadra

Department of Physics, University of Illinois, 1110 W. Green Street, Urbana, Illinois 61801

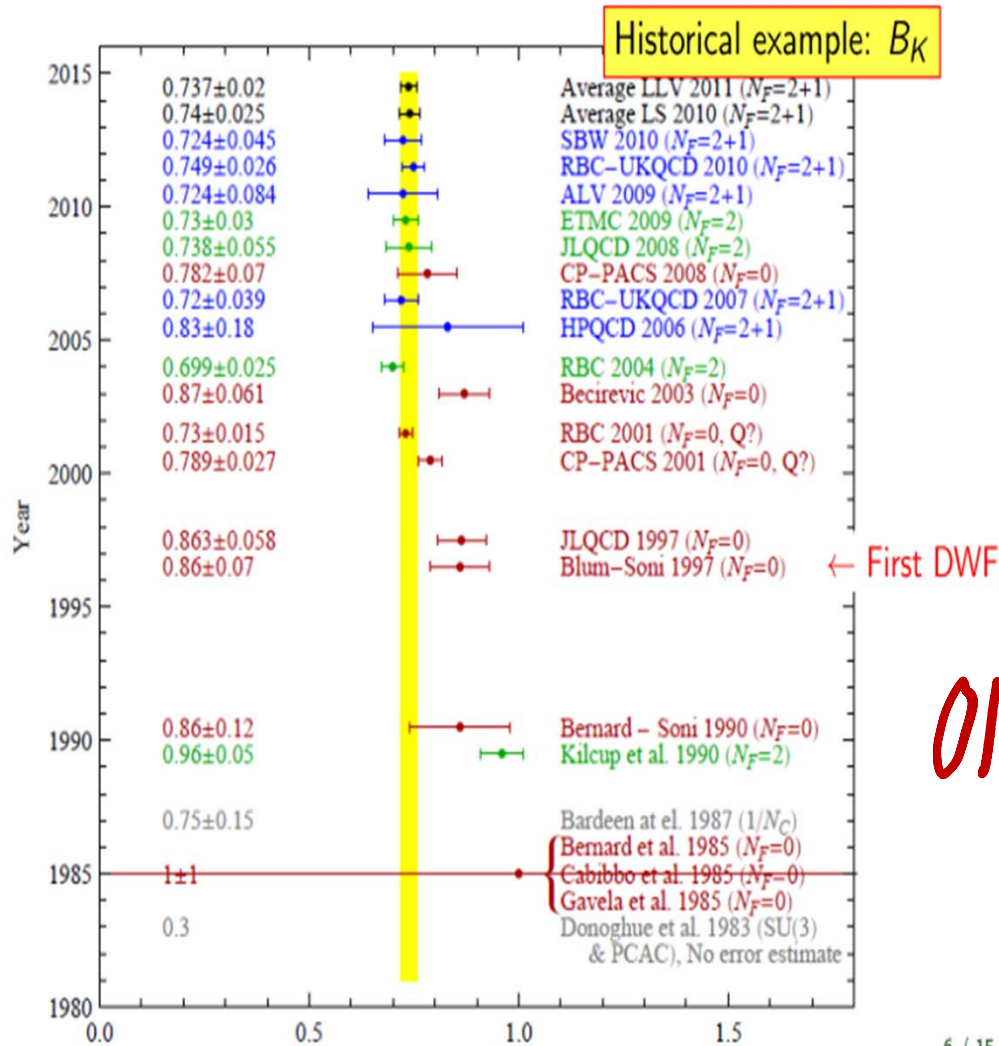
Andreas S. Kronfeld and Paul B. Mackenzie

Theoretical Physics Group, Fermi National Accelerator Laboratory, P.O. Box 500, Batavia, Illinois 60510

(Received 4 April 1996)

This paper presents a formulation of lattice fermions applicable to all quark masses, large and small. We incorporate interactions from previous light-fermion and heavy-fermion methods, and thus ensure a smooth connection to these limiting cases. The couplings in improved actions are obtained for arbitrary fermion mass m_q , without expansions around small- or large-mass limits. We treat both the action and external currents. By interpreting on-shell improvement criteria through the lattice theory's Hamiltonian, one finds that cutoff artifacts factorize into the form $b_n(m_q a)[\mathbf{p}a]^{s_n}$ where \mathbf{p} is a momentum characteristic of the system under study, s_n is related to the dimension of the n th interaction, and $b_n(m_q a)$ is a bounded function, numerically always of order 1 or less. In heavy-quark systems \mathbf{p} is typically rather smaller than the fermion mass m_q . Therefore, artifacts of order $(m_q a)^s$ do not arise, even when $m_q a \gtrsim 1$. An important by-product of our analysis is an interpretation of the Wilson and Sheikholeslami-Wohlert actions applied to nonrelativistic fermions. [S0556-2821(97)03607-2]

Power of the lattice: Only method to systematically reduce the NP error!



AB-initio Calculations

$$B_K = \frac{\langle K | (\bar{s} \gamma_\mu \gamma_5 d)^2 | K \rangle}{8/3 g^2 m_K^2}$$

ONE ILLUSTRATION

6 / 15

Flag 2019: sample [Nf=2 +1]

- $BK^{\text{hat}} = 0.7625(97) \dots 1.5\%$

- $f_B = 192.0 (4.3) \text{ MeV} \dots 2.2 \%$

- $\xi = 1.206 (17) \dots 1.5\%$

- $B \Rightarrow D, \delta[\text{sl FF } (q^2)] \sim 5-10\%$

- $RD \sim 4\%$

)) Need lot more
wakesp due
expts on the
horizon

Couplings between CWB+ AS & PBM et al

- **With Claude, we identified rather early on important observables in K, B...weak decays for lattice studies**
- **George Hockney, UCLA ~85=> FermiL**
- **Aida E-K, UCLA/BNL ~'91 => FermiL**
- **Jim S, UCLA/BNL/Edinburgh ~'94=> FermiL**
- **Jack L, Princeton/BNL ~'04 => FermiL**
- **Ruth VdW FermiL=>BNL => FermiL**

And of course also in Physics

- **For long CWB + AS used suggestions of Lepage and Mackenzie in renormalization of operators via “boosted coupling”, see L&M, PRD’93.**
- **And in heavy quark treatments we dealt with large am difficulties as suggested originally by Kronfeld and Mackenzie [inspired Jim Labrenz PhD work]. All that eventually evolved into a systematic treatment as in El-K, K & M, PRD’97.**

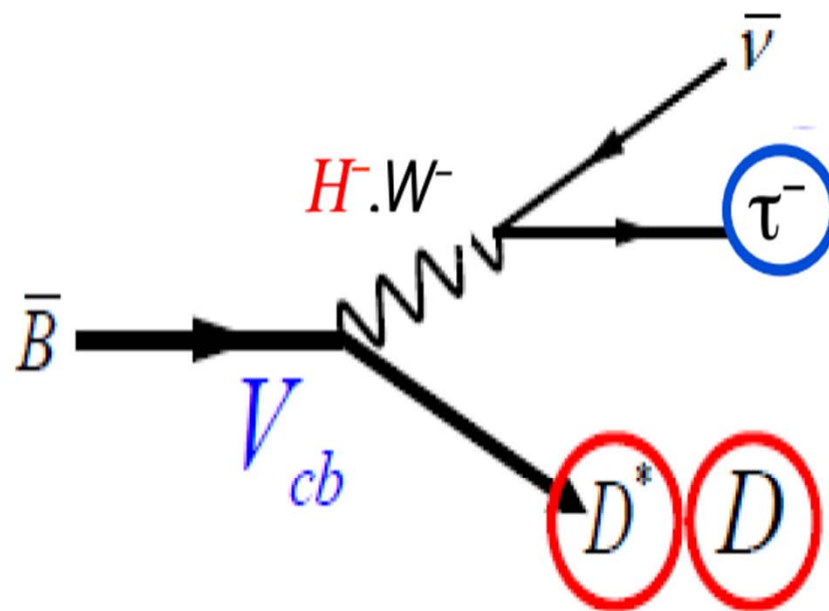
Anomalies galore!

- $RD(*) \sim 46(?)$; ALSO $R_4 \sim 26$ LHC
- $RK(*)$: $2.66(R_K)$; 2.2 & 2.56 R_{K^*}
- $g-2 \dots BNL \Rightarrow FNAL$ expt... ~ 3.66 *main lattice progress*
- ϵ' : a personal obsession....for a long^{^3} time= \Rightarrow 'cause of the strong belief that it is super-sensitive to NP
 216[PRL 2015] $\Rightarrow \sim 1400$ of which ~ 740 g c analyzed } *see NHC talk*
 $[2.1\sigma \Rightarrow ??]$ few more months to new results
- **Notice in each case, because of the omnipresence of non-perturbative effects, lattice methods provide crucial info for experiments to be able to use the data in the most economical manner**

BEST CHANCE IN A VERY LONG TIME OF POSSIBLE SIGHTINGS OF BSM



Exclusive $B \rightarrow D^{(*)}\tau\nu$



RA LUTH (BABAR)

'CP May 2012
(HEFEL, China)

MANUEL FRANCO
SEVILLA
PhD Thesis

Independent of
 V_{cb} !

- To test the SM Prediction, we measure

$$R(D) = \frac{\Gamma(\bar{B} \rightarrow D\tau\nu)}{\Gamma(\bar{B} \rightarrow D\ell\nu)} \quad R(D^*) = \frac{\Gamma(\bar{B} \rightarrow D^*\tau\nu)}{\Gamma(\bar{B} \rightarrow D^*\ell\nu)}$$

Leptonic τ
decays only

Several experimental and theoretical uncertainties cancel in the ratio!

- DD^* events are fully reconstructed.

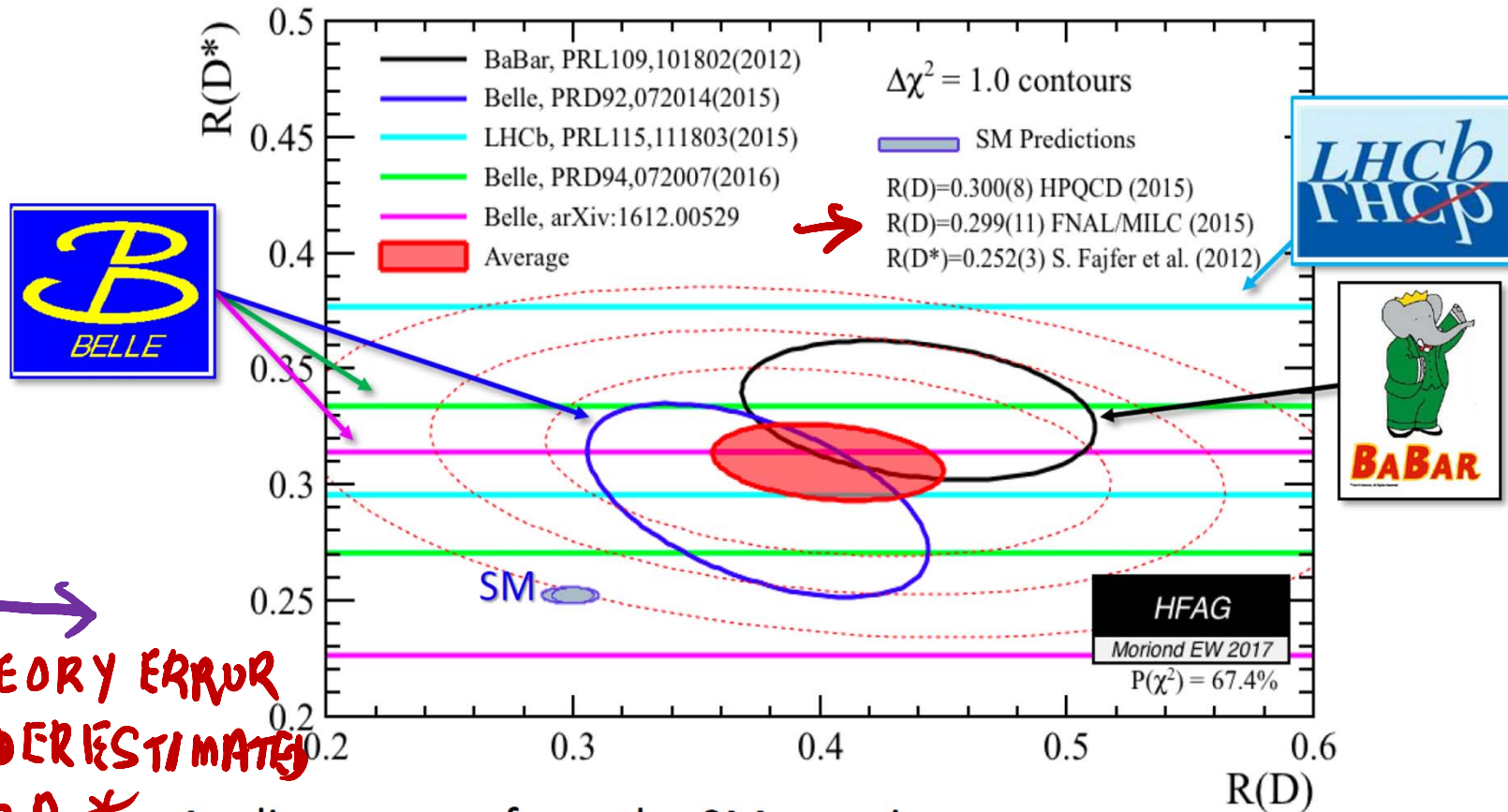
some expt + also some
The only errors tend to cancel

$l = \mu \text{ or } e$

■ $R(D^{(*)})$ by HFAG

Hirose [BELLE]@EW
MORIOND Mar. 2017

11/15



THEORY ERROR
UNDERESTIMATED
ESP R_D^*

$\sim 4\sigma$ discrepancy from the SM remains

- All the experiments show the larger $R(D^{(*)})$ than the SM
- More precise measurements at Belle II and LHCb are essential

Belle deviations quite mild

Rencontres de Moriond EW 2017

P Mackenzie Symp Nov 2019; soni-HET-BNL

Semileptonic B decays

BaBar measured an excess of $B^0 \rightarrow D^{(*)} \tau^- \nu_\tau$ (**3 σ away from SM!**) [PRD 88 (2013) 072012]
[Nature 546 (2017) 227]

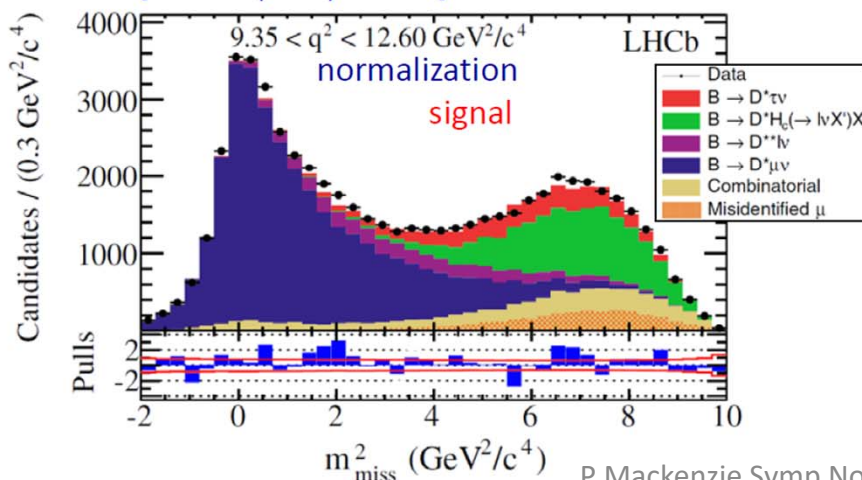
LHCb:

- $R(D^*)$
 - $B^0 \rightarrow D^{*+} \tau^- \bar{\nu}_\tau$, with $\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$ [PRL 115 (2015) 111803]
 - $B^0 \rightarrow D^{*-} \tau^+ \nu_\tau$, with $\tau^+ \rightarrow \pi^+ \pi^- \pi^+ (\pi^0) \bar{\nu}_\tau$ [PRL 120 (2018) 171802]
- $R(J/\psi)$
 - $B_c^+ \rightarrow J/\psi \tau^+ \nu_\tau$, with $\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$ [PRL 120 (2018) 121801]

■ Using $\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$

Information from the missing mass squared $m_{\text{miss}}^2 = (P_B - P_{D^*} - P_\mu)^2$ and muon energy

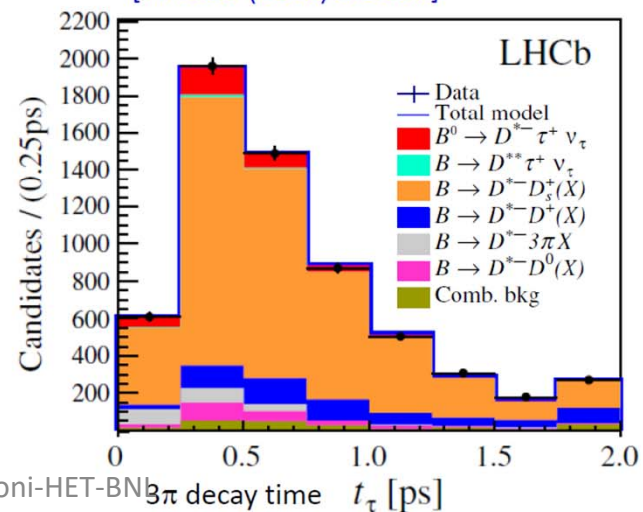
[PRL 115 (2015) 111803]



■ Using $\tau^+ \rightarrow \pi^+ \pi^- \pi^+ \bar{\nu}_\tau$

Information from the position of the pions. Normalized to $B^0 \rightarrow D^{*-} \pi^+ \pi^- \pi^+$

[PRL 120 (2018) 121801]

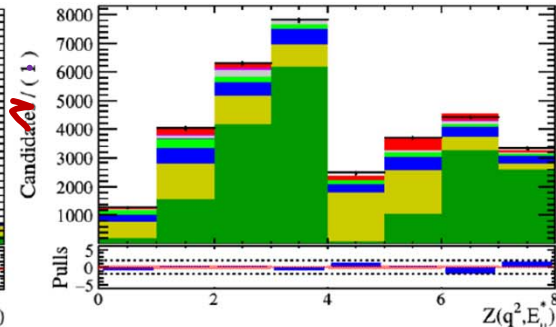
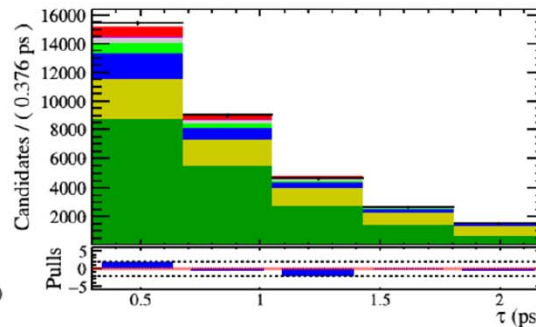
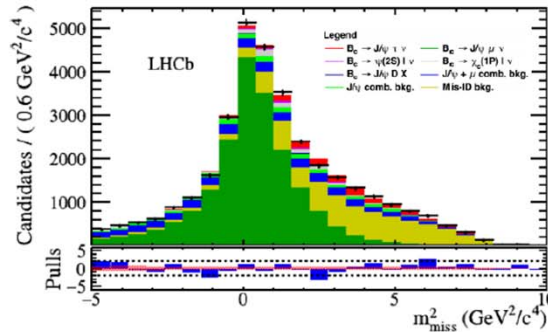


$$B_c \rightarrow J/\psi \tau \nu$$

2 PM Jan 2018

Greg Ciezarek,
on behalf of the LHCb collaboration

$$B_c \rightarrow \frac{b \rightarrow W \rightarrow \tau \nu}{\tau} \psi$$



- $R_{J/\psi} \equiv B_c \rightarrow J/\psi \tau \nu / B_c \rightarrow J/\psi \mu \nu$
- Measured using very similar techniques to $\mathcal{R}(D^*)$, on run 1 data
- $R_{J/\psi} = 0.71 \pm 0.17 \pm 0.18$

- $\sim 2\sigma$ from SM

- But nearly as far from consistency with $\mathcal{R}(D^*)$

- LHCb-PAPER-2017-035 (Run 1 data)

SM $R_{J/\psi} \sim 0.265 \pm 0.015$

QUITE ROBUST! ESSENTIALLY A NR Bound State

ALSO MARK SMITH PAPER 2018

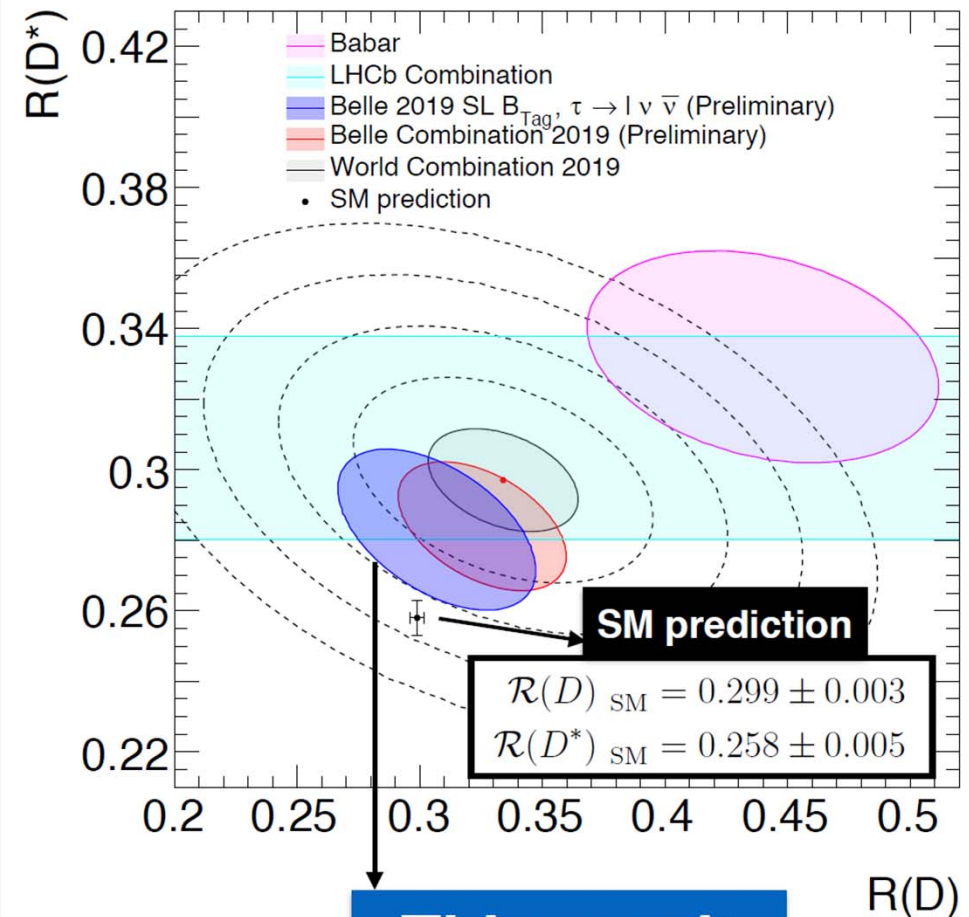
REMAINING ISSUES

PRIMARYLY EXPTAL

1. Stat 2. D^{**} 3. $\tau \rightarrow h_c + 2\ell$

Conclusion / Preliminary $R(D^{(*)})$ averages

- **Most precise measurement** of $R(D)$ and $R(D^*)$ to date
- First **$R(D)$** measurement performed with a **semileptonic tag**
- Results **compatible with SM** expectation within **1.2σ**
- **$R(D) - R(D^*)$ Belle average** is now within **2σ** of the SM prediction
- **$R(D) - R(D^*)$ exp. world average** tension with SM expectation **decreases from 3.8σ to 3.1σ**



This result

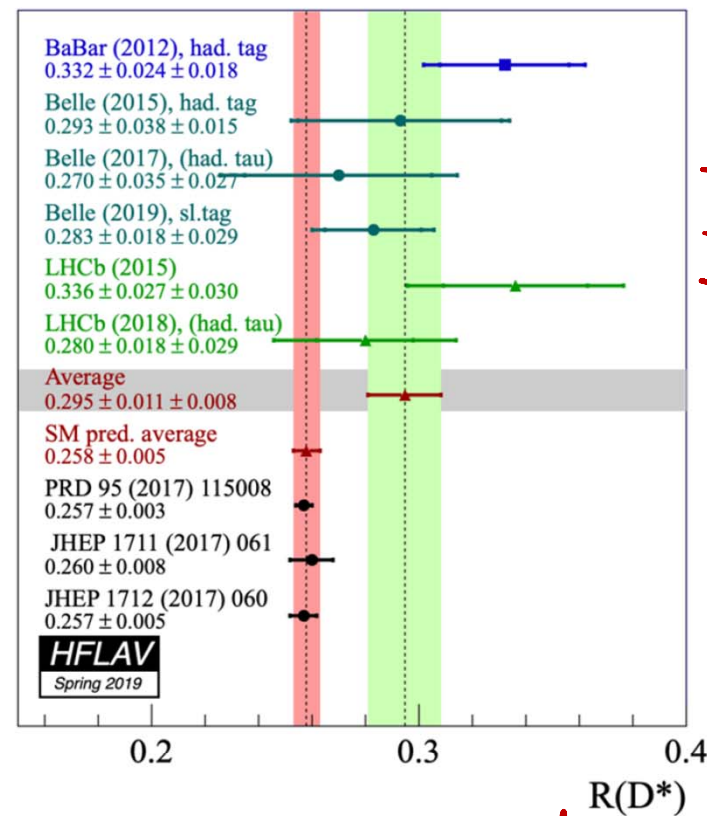
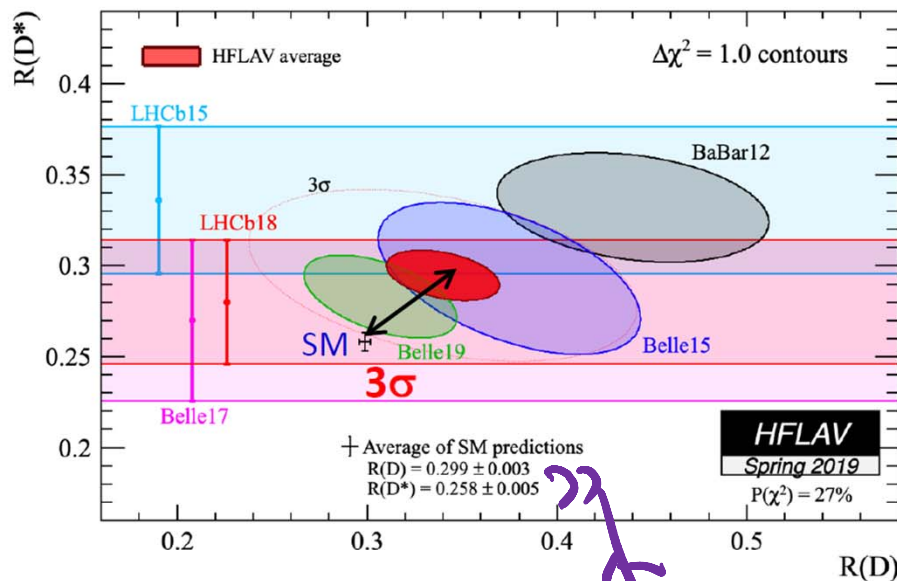
$$\begin{aligned}\mathcal{R}(D) &= 0.307 \pm 0.037 \pm 0.016 \\ \mathcal{R}(D^*) &= 0.283 \pm 0.018 \pm 0.014\end{aligned}$$

Semileptonic B decays

6 R_{D^*} measurements

New results (Moriond 2019) from Belle:

- Global picture of R_D and R_{D^*}



→ New results from Belle: $4\sigma \rightarrow 3\sigma$ deviation from SM

R_D : 3 BaBar, + Belle

↑↑

$B \rightarrow D^*$ Form factors from Lattice still cooking

FACT OR FARCE?

1) Exptal results [not all independent], AhL central values above Theory

4

experiment	tag method	τ decay mode	R_D	R_D^*	R_ψ
Babar (2012)[1]	hadronic	$1 \nu \nu$	$0.440 \pm 0.058 \pm 0.042$	$0.332 \pm 0.024 \pm 0.018$	
Belle (2015)[2]	hadronic	$1 \nu \nu$	$0.375 \pm 0.064 \pm 0.026$	$0.293 \pm 0.038 \pm 0.015$	
LHCb (2015)[5]	hadronic	$1 \nu \nu$	-	$0.336 \pm 0.027 \pm 0.030$	
Belle (2016)[2]	semileptonic	$1 \nu \nu$	-	$0.302 \pm 0.030 \pm 0.011$	
Belle (2017)[4]	hadronic	$\pi(\rho)\nu$	-	$0.270 \pm 0.035 \pm 0.027$	
LHCb (2017)[6]	hadronic	$3\pi\nu$	-	$0.291 \pm 0.019 \pm 0.029$	
Belle (2019)[7]	semileptonic	$1 \nu \nu$	$0.307 \pm 0.037 \pm 0.016$	$0.283 \pm 0.018 \pm 0.014$	
LHCb(2016) [9]	hadronic	$1 \nu \nu$	-	-	$0.71 \pm 0.17 \pm 0.18$
SM	-	-	0.299 ± 0.011	0.260 ± 0.008	0.26 ± 0.02

TABLE I: All experimental results announced to date on R_D , R_{D^*} and on R_ψ versus the predictions of those for the SM

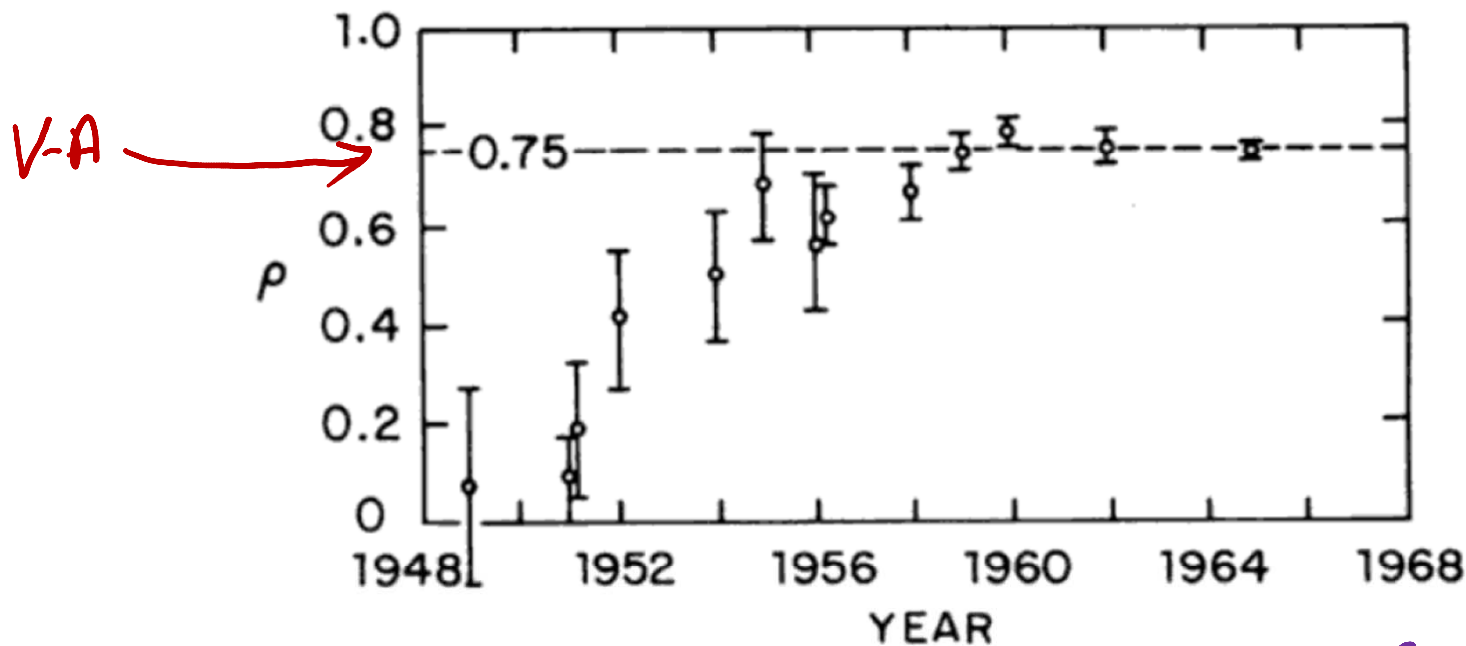
ALTMANNSTOFER, DeV+AS, Yicong Suo (in prep)

Excellent 9.5

RECAP

- 3 different major B-experiments
 - 3 with $B \Rightarrow D$
 - 7 with $B \Rightarrow D^*$
 - 1 with $B_c \Rightarrow \psi$
 - 9 with $\tau \Rightarrow l$ ($l=\mu$ or e) $\nu \nu'$
 - 2 with $\tau \Rightarrow \text{hadron} + \nu$
-
- Each and everyone of the 11 experimental results seem to imply tau is NOT just a heavy muon(electron) as dictated by SM.

Imp. Historical ~~Aside~~ CAUTION



FROM M PUROHIT

Figure 16. The change of the Michel parameter ρ from year to year.

From T. D. Lee's text

Lepton universality tests

- In the SM, ratios

$\bar{b} \rightarrow \bar{u} \mu^+ \mu^-$
 $\bar{b} \rightarrow \bar{u} e^+ e^-$

$$R_K = \frac{\int d\Gamma[B^+ \rightarrow K^+ \mu^+ \mu^-]/dq^2 \cdot dq^2}{\int d\Gamma[B^+ \rightarrow K^+ e^+ e^-]/dq^2 \cdot dq^2}$$

LHCb introduced such & well defined ratios

only differ from unity by phase space — the dominant SM processes couple equally to the different lepton flavours.

- Theoretically clean since hadronic uncertainties cancel in the ratio.
- Experimentally challenging due to differences in muon/electron reconstruction (in particular Bremsstrahlung from the electrons).
 - Take double ratios with $B \rightarrow J/\psi X$ decays to cancel possible sources of systematic uncertainty.
 - Correct for migration of events in q^2 due to FSR/Bremsstrahlung using MC (with PHOTOS).

LHCb

Lepton Flavour Universality

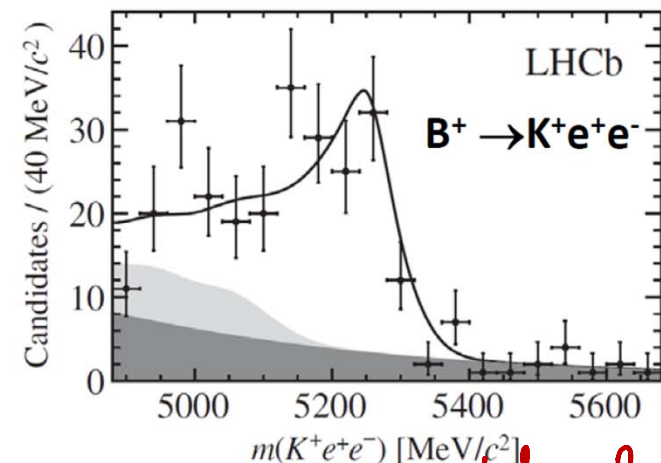
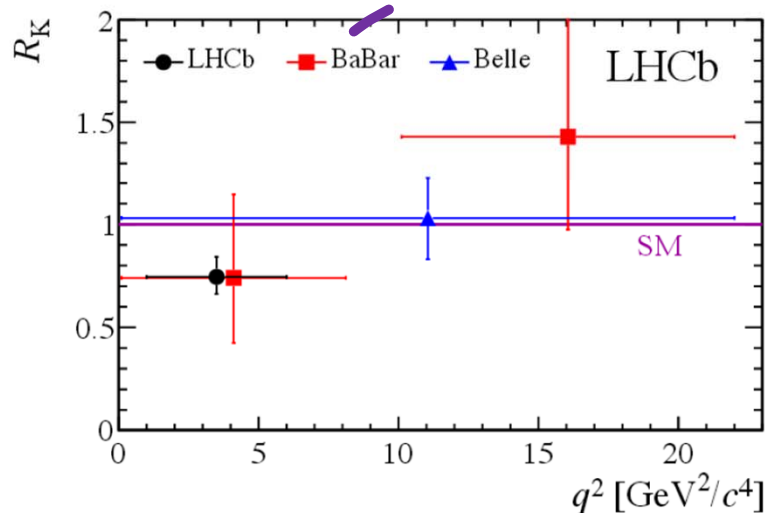
Arantza Oyanguren

- In the SM all leptons are expected to behave in the sa

$$R_K = \frac{\mathcal{B}(B^+ \rightarrow K^+ \mu^+ \mu^-)}{\mathcal{B}(B^+ \rightarrow K^+ e^+ e^-)} = 1.000 + \mathcal{O}(m_\mu^2/m_b^2) \text{ (SM)}$$

[PRL 113 (2014) 151601]

- Experimentally, use the $B^+ \rightarrow K^+ J/\psi (\rightarrow e^+ e^-)$ and $B^+ \rightarrow K^+ J/\psi (\rightarrow \mu^+ \mu^-)$ to perform a double ratio
- Precise theory prediction due to **cancellation of hadronic form factor uncertainties**



1 GeV < q² < 6 GeV *← should be safe from Rad Corr*

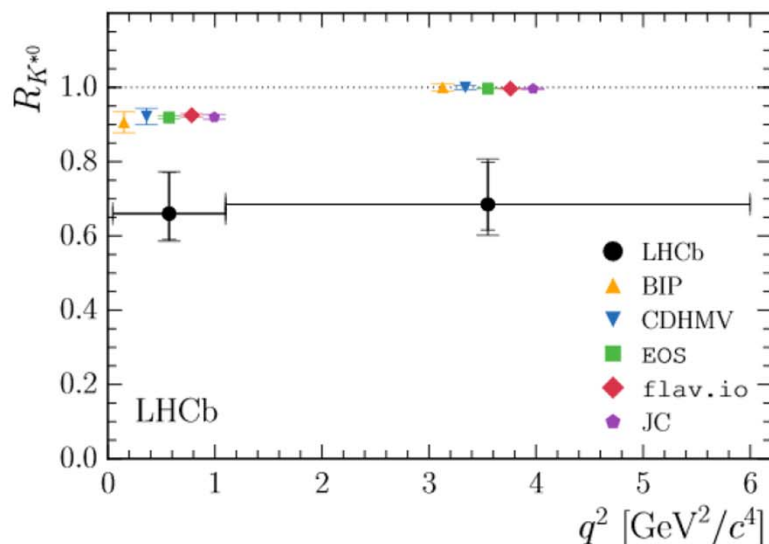
$$R_K = 0.745^{+0.090}_{-0.074} \text{ (stat)} \pm 0.036 \text{ (syst)}$$

→ Consistent, but lower, than the SM at **2.6σ**

Lepton Flavour Universality

• Results:

LHCb, JHEP08(2017)055



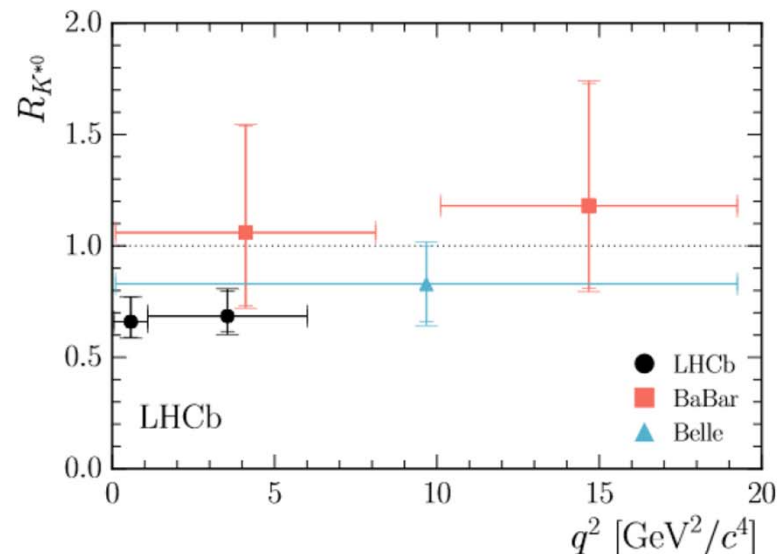
▲ BIP [EPJC 76 (2016) 440]
 ▼ CDH MV [JHEP 04 (2017) 016]
 ■ EOS [PRD 95 (2017) 035029]
 ◆ flav.io [EPJC 77 (2017) 377]
 ★ JC [PRD 93 (2016) 014028]

Low q^2 [0.045-1.1 GeV²]: $SM_{\nabla} = 0.922(22)$

$$R_{K^{*0}} = 0.66 \pm 0.11 \text{ (stat)} \pm 0.03 \text{ (syst)}$$

Central q^2 : [1.1-6 GeV²]: $SM_{\nabla} = 1.000(6)$

$$R_{K^{*0}} = 0.69 \pm 0.11 \text{ (stat)} \pm 0.05 \text{ (syst)}$$



● LHCb [PRL 113 (2014) 151601]
 ▲ Belle [PRL 103 (2009) 171801]
 ■ BaBar [PRD 86 (2012) 032012]

$\approx 3.4\sigma$

→ Consistent, but lower than the SM at **2.1-2.3 σ** (low q^2) and **2.4-2.5 σ** (central q^2)

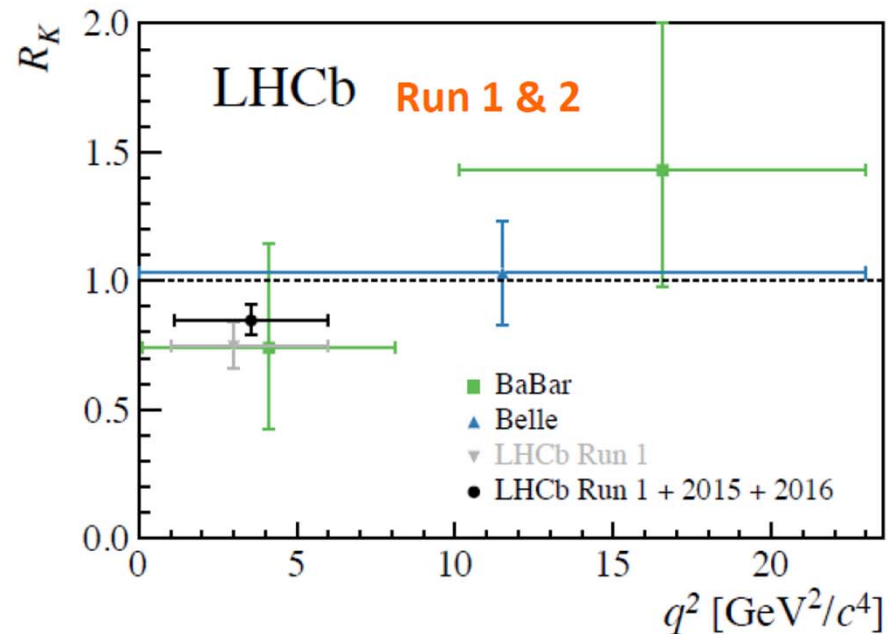
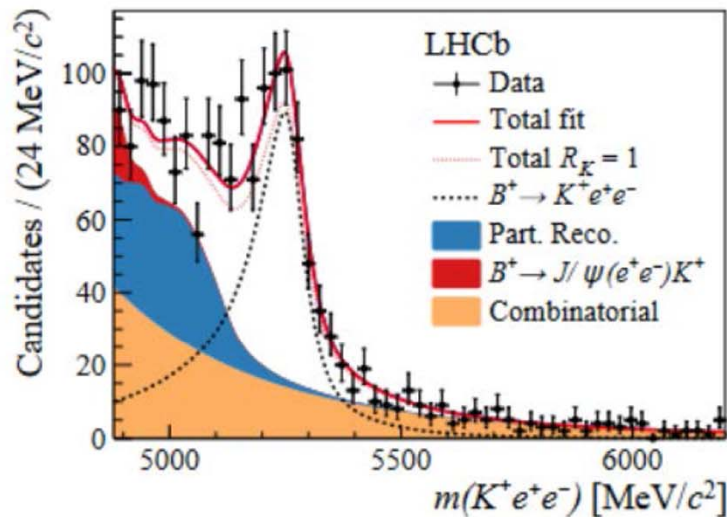
Rare B decays: R_K

New results (Moriond 2019):

Including partial sample of
Run2 (2fb^{-1})

[LHCb, *PRL* 122 (2019) 191801]

With improved reconstruction and
re-optimized analysed strategy



$1.1 \text{ GeV} < q^2 < 6 \text{ GeV}$

$$R_K = 0.846^{+0.060}_{-0.054}(\text{stat.})^{+0.016}_{-0.014}(\text{syst.})$$

→ Still consistent, lower, than the SM at 2.5σ

Not confirmed, not ruled out...

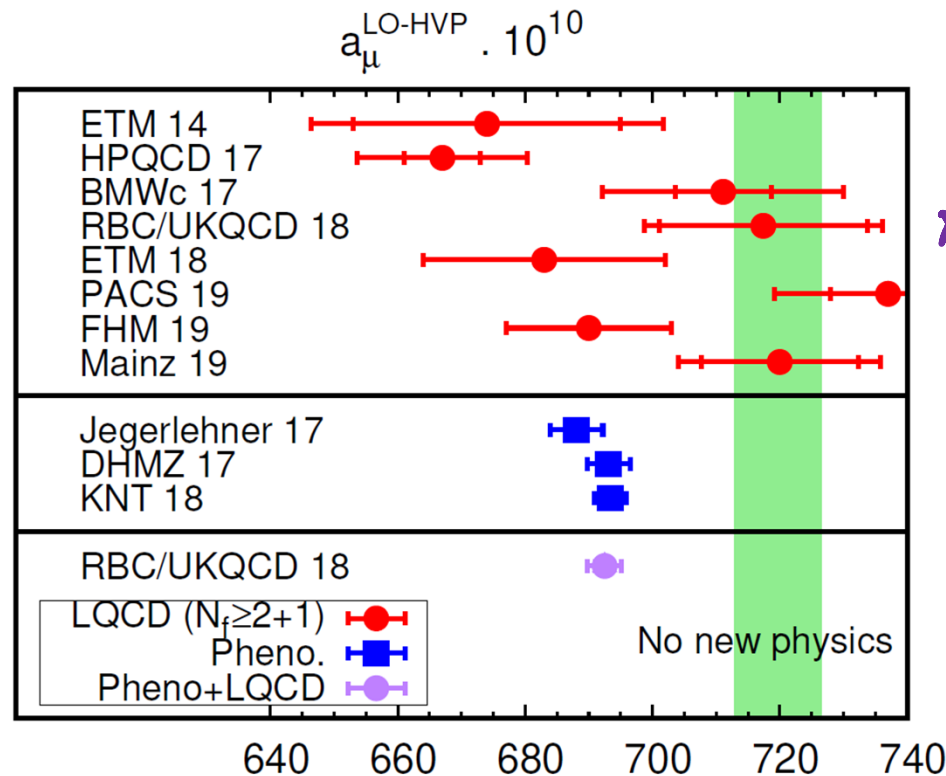
A. El-Khadra's talk at Seattle INT workshop.

Sept 2019

$(g-2)_\mu$

[prepared by K. Miura for WP]

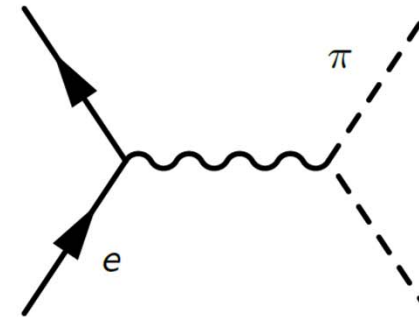
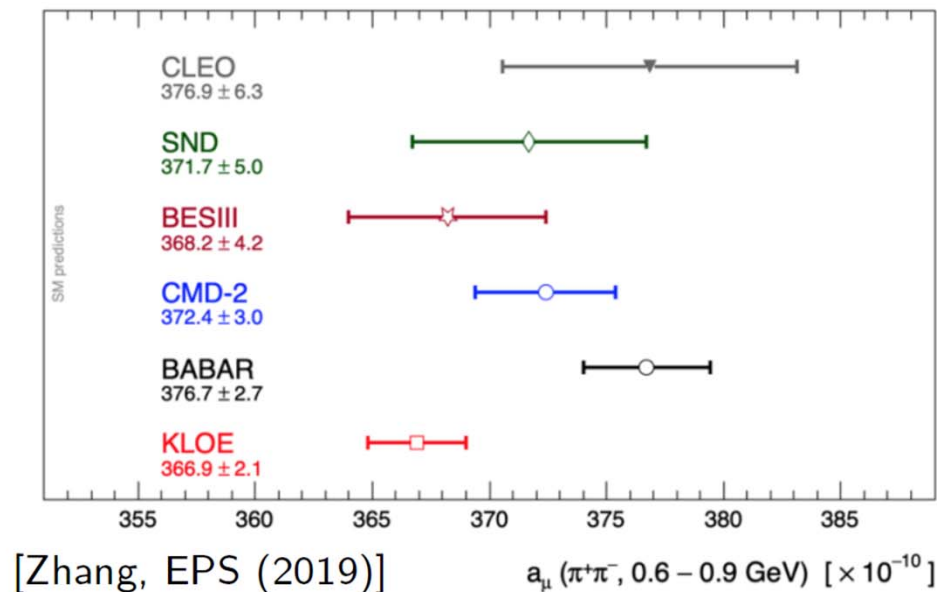
[V.]



x 1/3 reduction
in errors
very soon

Lattice work started by Tom Blum
EBL ~ '04

Tensions in Experiment



R-ratio data for $ee \rightarrow \pi\pi$ exclusive channel, $\sqrt{s} = 0.6 - 0.9 \text{ GeV}$ region

Tension between most precise measurements (BABAR/KLOE)

R-ratio a_μ^{HVP} uncertainty $<$ difference in this channel

Avoid tension by computing precise lattice-only estimate of a_μ^{HVP}

Use lattice QCD to inform experiment, resolve discrepancy

KILC, KOBACH
+ AS
PRD 2015

Table 1

Constraints on lepton-flavor violating and conserving processes. For the last four observables, the experimental null results are given in terms of a dimension-6 operator, suppressed by two orders of Λ , which can be interpreted as the nominal scale of new physics.

Observable	Limit
$\text{Br}(\mu \rightarrow 3e)$	$< 1.0 \times 10^{-12}$ [1]
$\text{Br}(\mu \rightarrow e\gamma)$	$< 5.7 \times 10^{-13}$ [1]
$\text{Br}(\tau \rightarrow 3e)$	$< 2.7 \times 10^{-8}$ [1]
$\text{Br}(\tau \rightarrow e^- \mu^+ \mu^-)$	$< 2.7 \times 10^{-8}$ [1]
$\text{Br}(\tau \rightarrow e^+ \mu^- \mu^-)$	$< 1.7 \times 10^{-8}$ [1]
$\text{Br}(\tau \rightarrow \mu^- e^+ e^-)$	$< 1.8 \times 10^{-8}$ [1]
$\text{Br}(\tau \rightarrow \mu^+ e^- e^-)$	$< 1.5 \times 10^{-8}$ [1]
$\text{Br}(\tau \rightarrow 3\mu)$	$< 2.1 \times 10^{-8}$ [1]
$\text{Br}(\tau \rightarrow \mu\gamma)$	$< 4.4 \times 10^{-8}$ [1]
$\text{Br}(\tau \rightarrow e\gamma)$	$< 3.3 \times 10^{-8}$ [1]
μ - e conversion	$\Lambda \gtrsim 10^3 \text{ TeV}$ [5]
$e^+ e^- \rightarrow e^+ e^-$	$\Lambda \gtrsim 5 \text{ TeV}$ [3]
$e^+ e^- \rightarrow \mu^+ \mu^-$	$\Lambda \gtrsim 5 \text{ TeV}$ [3]
$e^+ e^- \rightarrow \tau^+ \tau^-$	$\Lambda \gtrsim 4 \text{ TeV}$ [3]

Ist gem not
sensitive to
NP
+
(g-2)_μ
R_K(*)
+
R_D(*)

-

Possibly interesting inter-related story reg. LUV evolving over the past ~15 years

- μ on ($g-2$).....BNL ~2004; FermiL ~2017=>?
- $RD(*)$ BaBar, Belle, LHCb~2012-----→
- $RK(*)$ LHCb ~2014-----→
- Intriguing rather long tell- tale signs of LUV or few sigma flukes?
- Fortunately, wont have to wait too long < ~ 2 years due to FermiL, LHCb & Belle-II AND LATTICE

IF
exptl hints stay then



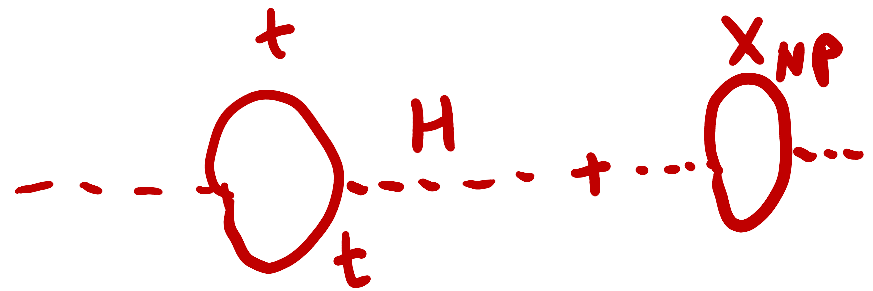
Altmanmshofer, Dev, A.S. 2017
+WIP

**ANOMALIES: POSSIBLY A HINT FOR
(NATURAL) SUSY-WITH RPV3**

then RPV - Susy extremely well
motivated - -

- ASSUMING the anomaly is REAL & HERE TO STAY [BIG ASSUMPTION due to caveats mentioned]
- Anomaly involves simple tree-level semi-leptonic decays
- Also $b \Rightarrow \tau$ (3rd family)
- **Speculate: May be related to Higgs naturalness**
- Seek minimal solution: perhaps 3rd family super-partners(a lot) lighter than other 2 gens > proton decay concerns may not be relevant=> RPV [“natural” SUSY]
- **RPV natural setting for LUV ...can accommodate g-2, RK(*) if needs be**
- Collider signals tend to get a lot harder than (usual-RPC) SUSY
- RPV makes leptoquarks natural [and respectable]
- Moreover, RPV should be viewed as an umbrella i.e. under appropriate limits other models are incorporated

$$m_H \approx 126 \text{ GeV}$$



RPV₃ preserves gauge coupling unification irrespective of # of effective gens. 1, 2 or 3.

ADS-PD'17

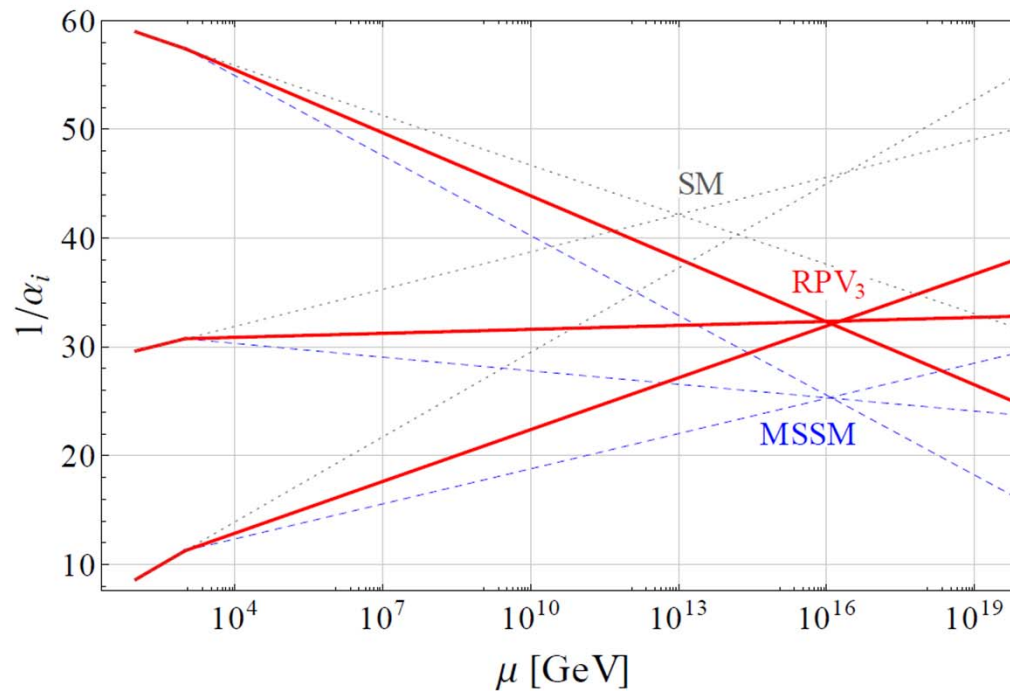


FIG. 2. RG evolution of the gauge couplings in the SM, MSSM and with partial supersymmetrization.

Unification scale ~ stays same, only value of couplings shifts

For pheno relevant terms:

ADS'PRD 2017

$$\mathcal{L} = \lambda'_{ijk} [\tilde{\nu}_{iL} \bar{d}_{kR} d_{jL} + \tilde{d}_{jL} \bar{d}_{kR} \nu_{iL} + \tilde{d}_{kR}^* \bar{\nu}_{iL}^c d_{jL} \\ - \tilde{e}_{iL} \bar{d}_{kR} u_{jL} - \tilde{u}_{jL} \bar{d}_{kR} e_{iL} - \tilde{d}_{kR}^* \bar{e}_{iL}^c u_{jL}] + \text{H.c.}$$

Analytic
of
 G_F
 $\sqrt{2}$

$$\mathcal{L}_{\text{eff}} \supset \frac{\lambda'_{ijk} \lambda'^*_{mnk}}{2m_{\tilde{d}_{kR}}^2} \left[\bar{\nu}_{mL} \gamma^\mu \nu_{iL} \bar{d}_{nL} \gamma_\mu d_{jL} \right. \\ \left. - \nu_{mL} \gamma^\mu e_{iL} \bar{d}_{nL} \gamma_\mu \left(V_{\text{CKM}}^\dagger u_L \right)_j + \text{h.c.} \right] \\ - \frac{\lambda'_{ijk} \lambda'^*_{mjn}}{2m_{\tilde{u}_{jL}}^2} \bar{e}_{mL} \gamma^\mu e_{iL} \bar{d}_{kR} \gamma_\mu d_{nR} ,$$

) RPV₃ interaction

← Dim-6

→ $f_n \Gamma_D(x)$

NOTE:

ITS
SM-like!

2 Theoretical framework

2.1 Effective Hamiltonian

We adopt the most general $SU(3)_C \otimes U(1)_Q$ -invariant effective Hamiltonian describing $b \rightarrow c \ell \bar{\nu}_\ell$ transitions at the bottom quark scale, not considering the possibility of light right-handed neutrinos:

$$\mathcal{H}_{\text{eff}}^{b \rightarrow c \ell \bar{\nu}} = \frac{4G_F}{\sqrt{2}} V_{cb} \left[(1 + C_{V_L}) \mathcal{O}_{V_L} + C_{V_R} \mathcal{O}_{V_R} + C_{S_R} \mathcal{O}_{S_R} + C_{S_L} \mathcal{O}_{S_L} + C_T \mathcal{O}_T \right] + \text{h.c.} \quad (2.1)$$

The above fermionic operators are given by⁴

$$\mathcal{O}_{V_{L,R}} = (\bar{c} \gamma^\mu b_{L,R}) (\bar{\ell}_L \gamma_\mu \nu_{\ell L}), \quad \mathcal{O}_{S_{L,R}} = (\bar{c} b_{L,R}) (\bar{\ell}_R \nu_{\ell L}), \quad \mathcal{O}_T = (\bar{c} \sigma^{\mu\nu} b_L) (\bar{\ell}_R \sigma_{\mu\nu} \nu_{\ell L}), \quad (2.2)$$

Very useful work includes constraints from $B_c(x)$, q^2 distribution, $B_c \rightarrow \tau \nu$

not only

$C_{V_L} \neq 0 \Rightarrow$ corresponds to our RPV3
MURGULETOL
1964.09311
MARTIN J.

MURCU et al

- On the other hand, considering scenarios with only a single Wilson coefficient present, there is a clear preference for C_{V_L} : removing the other three Wilson coefficients increases χ^2 only by 1.4, corresponding to 0.14σ . Hence, Min 1 is well compatible with a global modification of the SM, that is, C_{V_L} being the only non-zero coefficient.

Grac. Belle EWM χ^2 χ^2 χ^2

central values do change. Again all individual coefficients are roughly compatible with zero at 1σ . C_{V_L} alone also still provides an excellent fit to all the data, now with a smaller central value of ~ 0.08 . Interestingly, the fit with only C_T is improved by the new results, which,

A striking endorsement for RPV_3 . IFNP is needed!

But CVL only solution also has a bad news

[esp small non-vanishing CVL]

- **It may just be SM + (expt + theory) errors**
- **Assuming** NP is needed then RPV3 is a very good candidate

Sequel to our work on
RPV₃

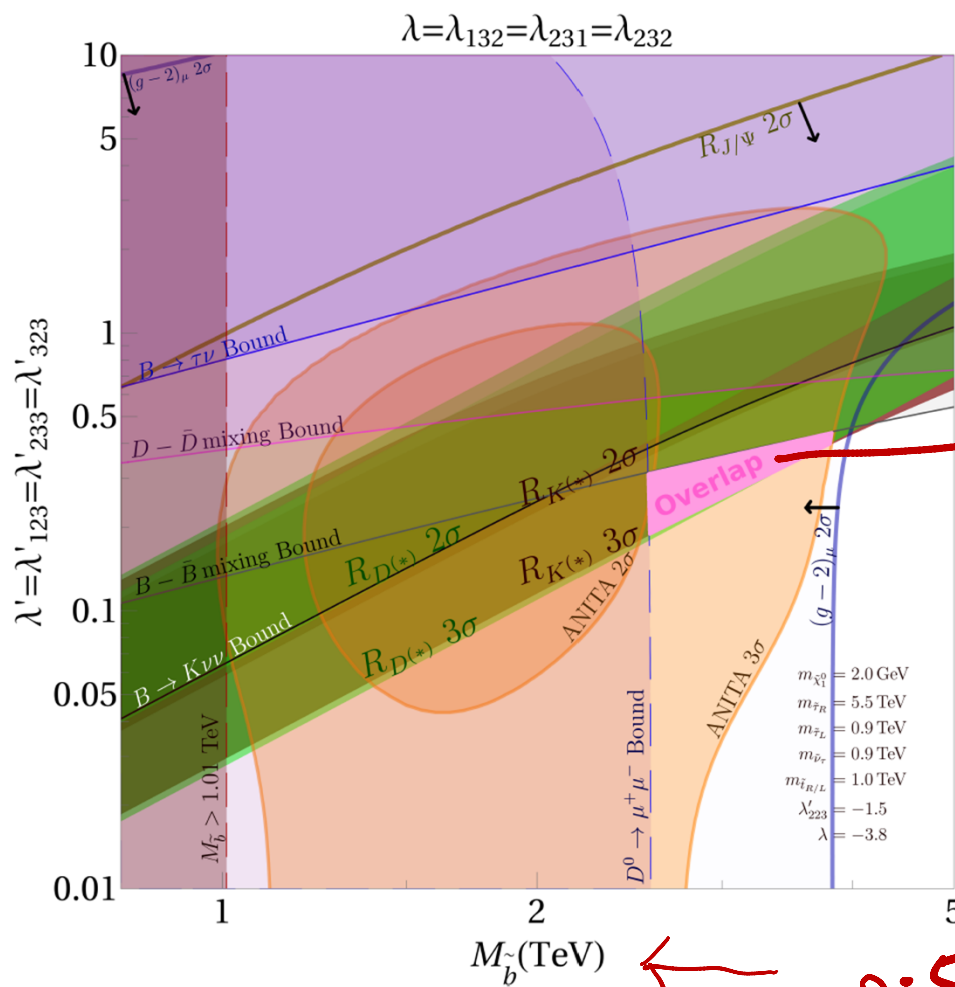
Addressing B -anomalies, muon $g - 2$ and ANITA anomaly in a Minimal R -parity violating supersymmetric framework

Wolfgang Altmannshofer¹, P. S. Bhupal Dev², Amarjit Soni³, Yicong Sui²

¹Max-Planck-Institut für Physik, Werner-Heisenberg-Institut, D-80805 München, Germany
²Department of Physics, University of Wisconsin at Madison, 480 Lincoln Drive, Madison, WI 53706, USA
³Department of Physics, University of Illinois at Chicago, 55 S. Chicago Ave., Chicago, IL 60607, USA

Excellent grad student
@ WUSTL

10 +
ALL known
exptal
constraints



→ overlap

$m_{\text{sup}} \rightarrow 5 \text{ TeV}$

$5 \lesssim m_{\tilde{t}_2} \leq 3.5 \text{ TeV}$

$$2.5 \lesssim m_{T_4} \leq 3.5 \text{ TeV}$$

IN CLOSING: A REMINDER

Importance of the “IF”: score card

- Beta decay $\Rightarrow G_f \Rightarrow W \dots$
 - Huge suppression of $KL \Rightarrow \mu \mu$; miniscule $\Delta m_K \Rightarrow$ charm
 - $KL \Rightarrow 2 \pi$ but very rarely; mostly to $3 \pi \Rightarrow$ CP violation \Rightarrow 3 families
 - Largish B_d –mixing \Rightarrow large top mass
 - etc.....
 - \Rightarrow extremely unwise to put all eggs in HEF
 - info from IF complementary to HEF can be a crucial guide
for pointing to new thresholds as well as to provide important clues
to the nature of the signals there from
- History may repeat yet again!*

Summary + Outlook + . . .

- Hints of LUV are extremely interesting, intriguing and important. *There is nothing we know of that tells us that these hints cannot be true.*
- While these indications are rather serious, **they are not yet compelling. They ask for too radical a departure from conventional understanding so we must exercise extreme caution and care before accepting them. Moreover, in each of the 3 cases there are features that cause concern.**
- Fortunately significant experimental/theoretical progress should occur in $< \sim 2$ years and is eagerly awaited.
- **Given all the above hints, may be just may be with some luck the IF will lead us to the gem of NP and once again, as many times in the past, guide collider physics et al**

Thank you, Paul

- For many contributions to physics via leading the FermiLab lattice group for long. In much of the topics I touched on, Paul + FermiLab Lattice group indeed played the central role.
- And also for the major impact on USQCD via his leadership in its Executive Committee.
- **We all wish you the very best in your retirement years!**

LQ EFT SBS + JW et al

LHCb also finds deficit in
observed $B_s \rightarrow \phi \mu \mu$ compared to "SM"
 $\sim 2.5 \sigma$
Less reliable
due to LD contamination

URGENTLY Needed! $\Rightarrow R_\phi \equiv \frac{B_s \rightarrow \phi \mu \mu}{B_s \rightarrow \phi e e}$
FROM LHCb

Revisiting R-parity violating interactions as an explanation of the B-physics anomalies

Sokratis Trifinopoulos^{*1}

"S-T" sym

¹*Physik-Institut, Universität Zürich, CH-8057 Zürich, Switzerland*

introduces a nice formal sym argument to implement
 Ok RPV₃ [3rd gen \bar{s} lightest]
 ↑ → Abstract

In the last few years, the ratios $R_{D^{(*)}}$ and of $R_{K^{(*)}}$ have reportedly exhibited significant deviations from the relevant Standard Model predictions, hinting towards a possible violation of Lepton Flavor Universality and a window to New Physics. We investigate to what extent the inclusion of R-parity violating couplings in the Minimal Supersymmetric Standard Model can provide a better fit to the anomalies simultaneously. We perform this analysis employing an approximate, non-abelian $\mathcal{G}_f = U(2)_q \times U(2)_\ell$ flavour symmetry, which features a natural explanation of the appropriate hierarchy of the R-parity violating couplings. We show that, under the requirement of a supersymmetric spectrum with much heavier left-handed doublet superpartners, our assumption favors a considerable enhancement in the tree-level charged-current $B \rightarrow D^{(*)} \tau \bar{\nu}$, while the anomalies induced by $b \rightarrow s \ell^+ \ell^-$ receive up to an approximate 30% improvement. The consistency with all relevant low-energy constraints is assessed.

Implementing S-T sym simplifies things considerably by reducing # of parameters.

BUT
can not a/c
for $(g-2)_\mu$

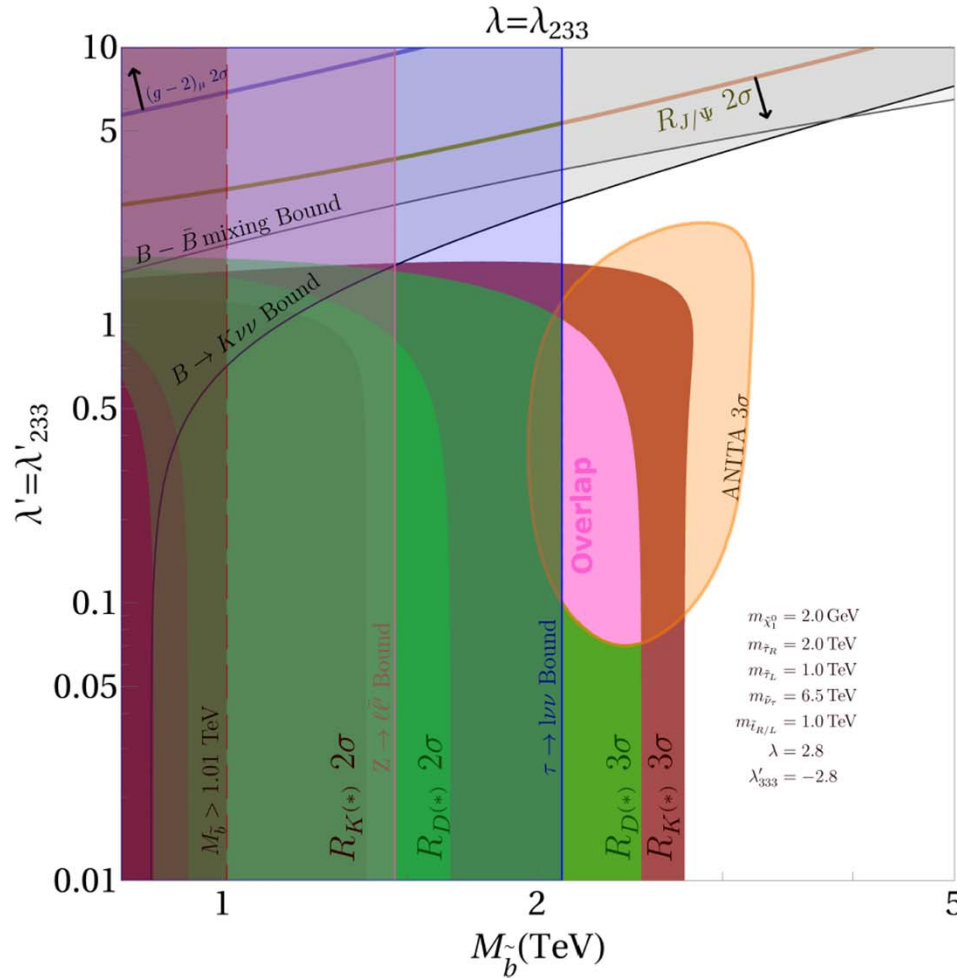


FIG. 5: Benchmark scenario for ST symmetry with overlapping $R_D^{(*)}$, $R_{J/\Psi}$, $R_K^{(*)}$ and ANITA regions. The total overlap is shown as pink area. $R_D^{(*)}$ 2, 3 σ flavored regions are denoted as green regions; $R_K^{(*)}$ 2, 3 σ flavored regions are shown in red regions; $(g-2)_\mu$ and $R_{J/\Psi}$ 2 σ flavored region is marked by thick blue and dark yellow edges, respectively, with arrows pointing inwards to the allowed regions; ANITA anomaly 3 σ flavored region is shown in orange regions. $B \rightarrow K\nu\nu$ bound is shown as black curve with forbidden region indicated as dark gray region while $B - \bar{B}$ mixing bound shown as gray curve with forbidden region in gray; $Z \rightarrow \ell\bar{\ell}'$ bound is shown as pink vertical line with forbidden region in light pink; $\tau \rightarrow \ell\nu\nu$ bound is shown as blue vertical line with forbidden region in light blue.

BACK TO LOW PT: NOTABLE LFV DECAYS OF TAU AND B [IN RPV3]

LFV of τ & q of B 's
 \rightarrow 3-g centric RPV_3 : Altmannshofer
 Dev, AS, YS

* Impreg

Mode	Model dependent BR	Current bound
$\tau \rightarrow \mu \phi$	2×10^{-10}	8×10^{-8}
$\tau \rightarrow \mu K K$	3×10^{-11}	4×10^{-8}
$\tau \rightarrow \mu K_s^0$	6×10^{-11}	2×10^{-8}
$\tau \rightarrow 3\mu$	1.5×10^{-10}	2×10^{-8}
$\tau \rightarrow \mu \gamma$	1.1×10^{-11}	4×10^{-8}
$\tau \rightarrow \mu l^+ l^-$	6×10^{-12}	2×10^{-8}
$b \rightarrow s \mu \tau$	7×10^{-7}	4.5×10^{-5}
$B_s \rightarrow \tau \mu$	1.3×10^{-8}	N/A

TABLE I: Few examples of lepton flavor violating decay modes of τ and of B - mesons. Shown are also loop decays $\tau \rightarrow \mu \gamma$ and $\tau \rightarrow \mu l^+ l^-$; see text

$B_u[B \rightarrow K^* \mu \tau \sim 10^{-7}]!$ \Leftarrow Recall $B \rightarrow X_c \ell \nu$
 $\propto V_{cb} \sim \alpha^2$
 $\& B \rightarrow K \tau \tau$ also deserves attention

In passing, a side remark, please

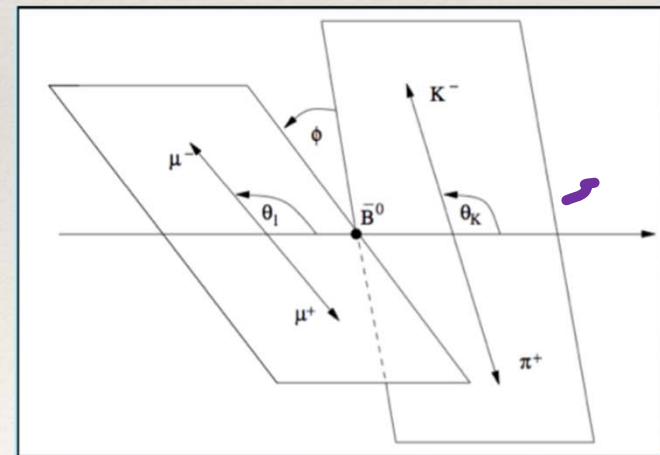
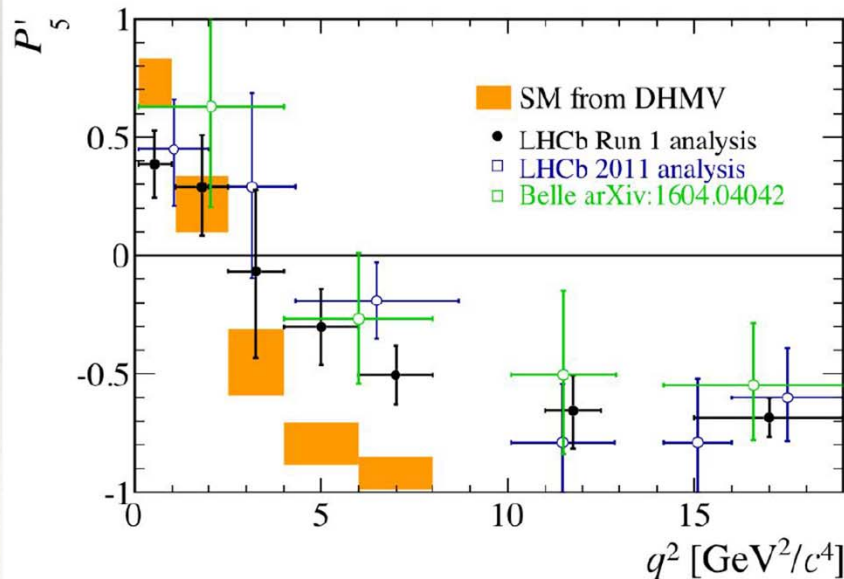
- QCD and therefore non-perturbative dynamics critically effects SM and or BSM
- In almost all of these “IF” experiments, quantitative understanding of non-perturbative [non-P] effects is of crucial importance to make most economical use of experimental data, often obtained at huge cost. **The non-P methods do not just need humungous computing hardware, (wo)man power needs are also very large. Given their vital use, they deserve greater support from the (experimental) community.**

B-flavor anomalies: P_5' $B \rightarrow K \mu \mu$

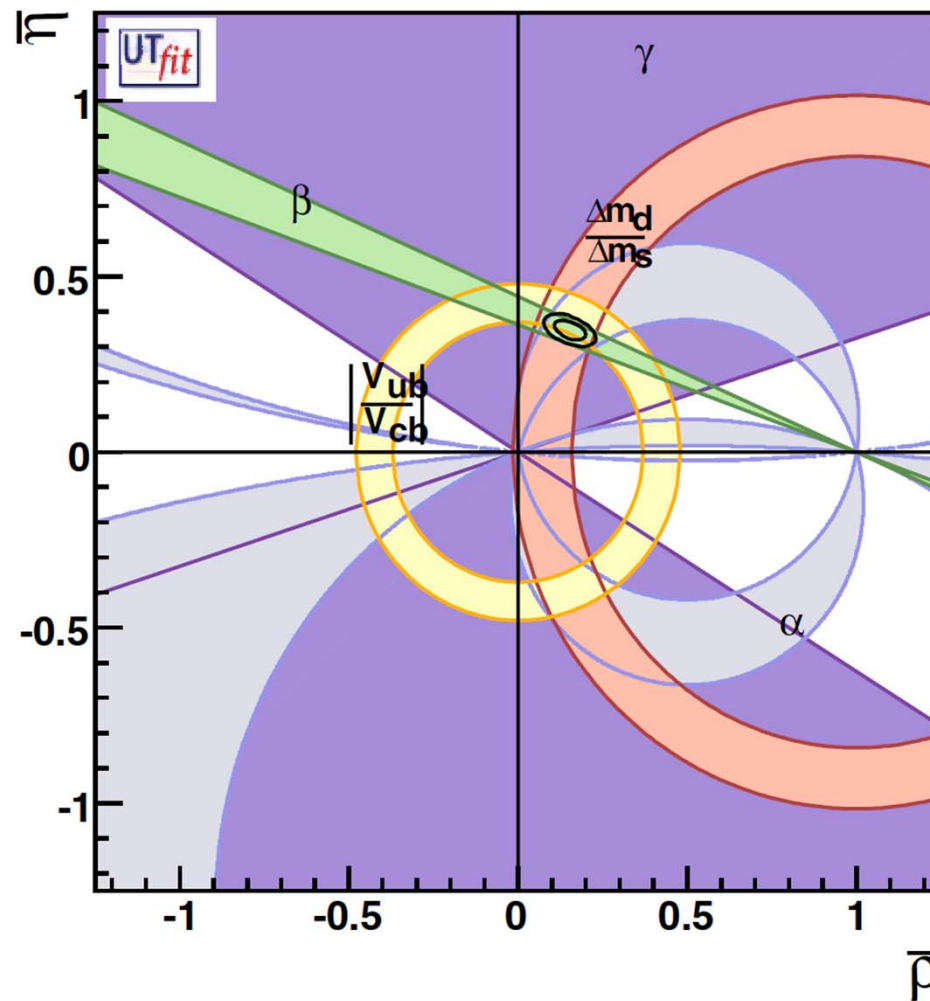
REMAIN CONCERNED
ABOUT NON-local
contributions

Much LESS clean esp
in low q^2

- ❖ Several angular observables measured as functions of q^2
- ❖ Some, like P_5' , are optimized to be insensitive to hadronic uncertainties: [\[Descotes-Genon, Matias, Ramon, Virto: 1207.2753\]](#)



PRL **97**, 151803 (2006)



MARCELLA BONA ~2007

COUPLINGS BETWEEN CWB+ AS & PBM ET AL

Improving constraints on $\tan\beta/m_H$ using $B \rightarrow D \tau \bar{\nu}$

Ken Kiers* and Amarjit Soni†

Department of Physics, Brookhaven National Laboratory, Upton, New York 11973-5000

(Received 12 June 1997)

We study the q^2 dependence of the exclusive decay mode $B \rightarrow D \tau \bar{\nu}$ in type-II two Higgs doublet models (2HDM's) and show that this mode may be used to put stringent bounds on $\tan\beta/m_H$. There are currently rather large theoretical uncertainties in the q^2 distribution, but these may be significantly reduced by future measurements of the analogous distribution for $B \rightarrow D(e, \mu) \bar{\nu}$. We estimate that this reduction in the theoretical uncertainties would eventually (i.e., with sufficient data) allow one to push the upper bound on $\tan\beta/m_H$ down to about 0.06 GeV^{-1} . This would represent an improvement on the current bound by about a factor of 7. We

FF
f₁ f₀

used HQS

⇒ followed up by Nierste et al; Fajfer et al '12
/08

**ALL 11 EXPERIMENTAL CENTRAL
VALUES ARE ABOVE THEORY!**